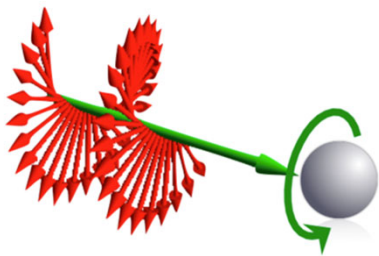
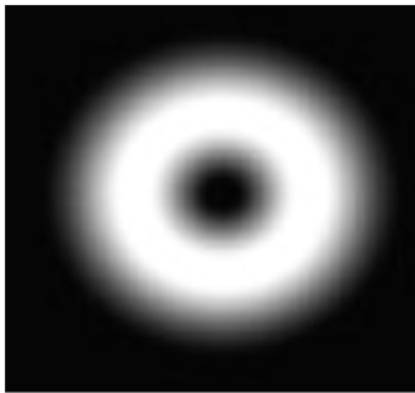
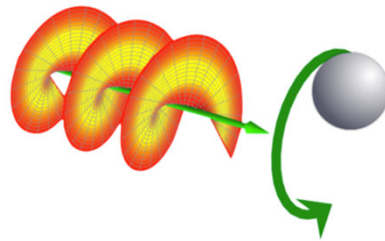


Structured light manipulation in strongly anisotropic metamaterials

Natalia M. Litchinitser, Jingbo Sun
Duke University

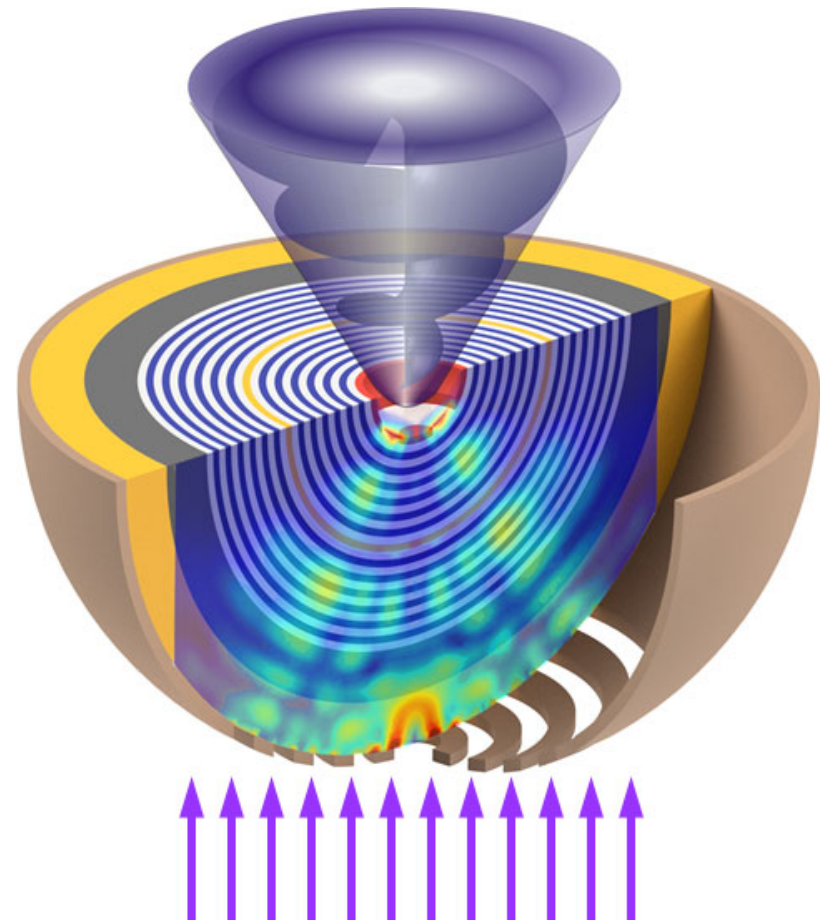


SAM



OAM

SAM – spin angular momentum
OAM – orbital angular momentum

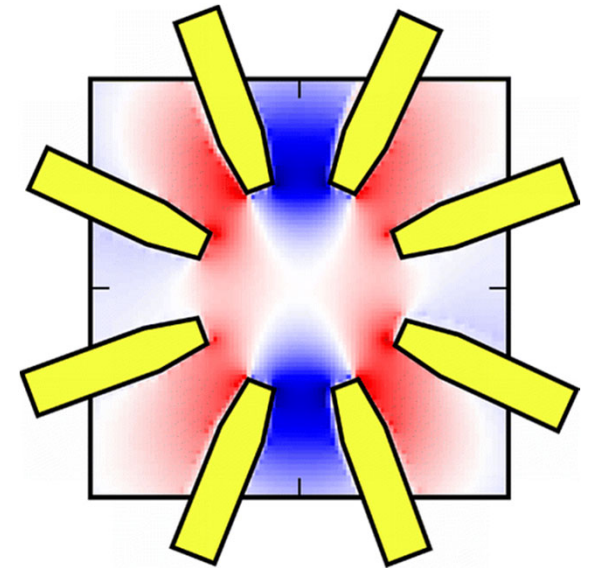
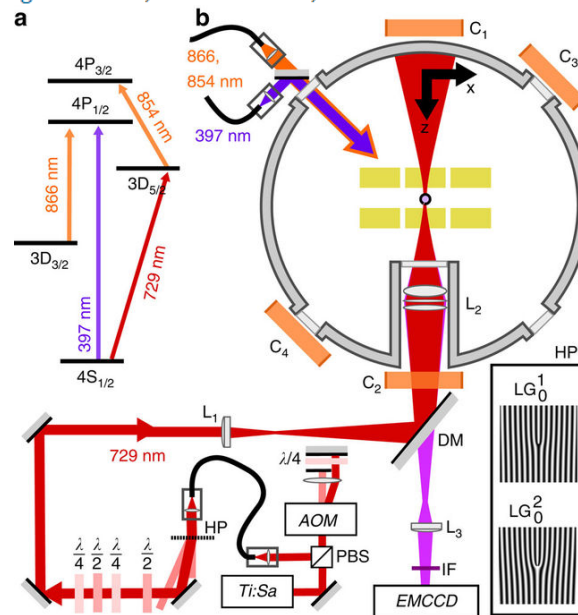


Hyperlens

Transfer of optical orbital angular momentum to a bound electron

Christian T. Schmiegelow , Jonas Schulz, Henning Kaufmann, Thomas Ruster, Ulrich G. Poschinger & Ferdinand Schmidt-Kaler

Nature Communications **7**,
Article number: 12998 (2016)
doi:10.1038/ncomms12998



Subwavelength Focusing of Light with Orbital Angular Momentum

Reinier W. Heeres* and Valery Zwiller

Kavli Institute of Nanoscience, Delft University of Technology, P.O. Box 5046, 2600 GA Delft, The Netherlands

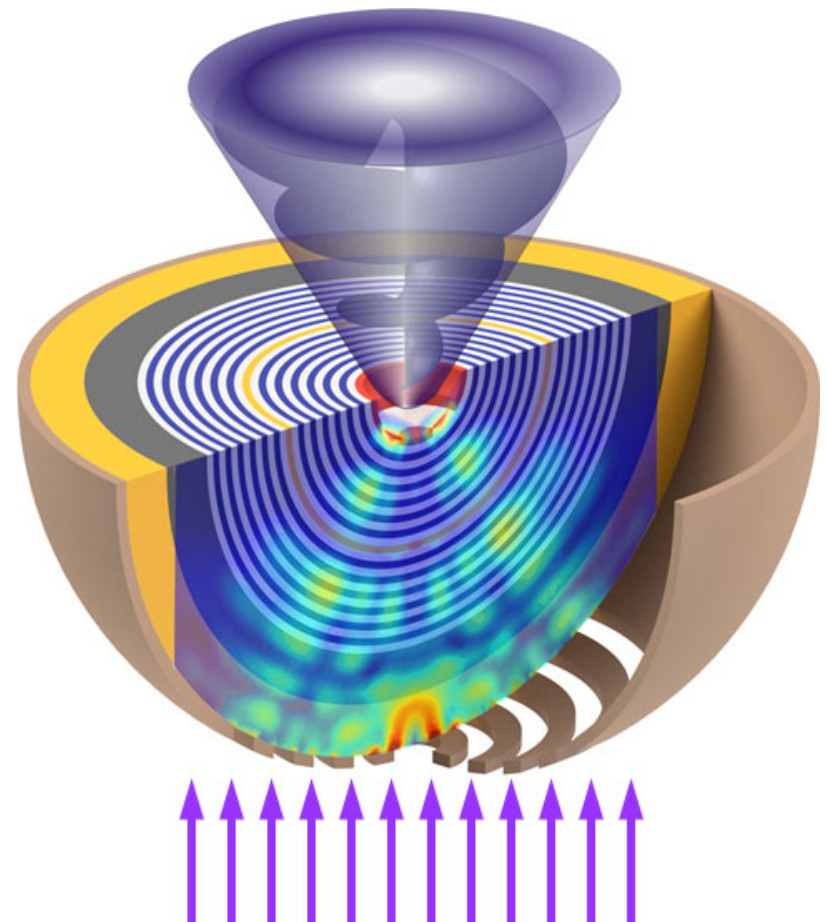
Nano Lett., **2014**, 14 (8), pp 4598–4601

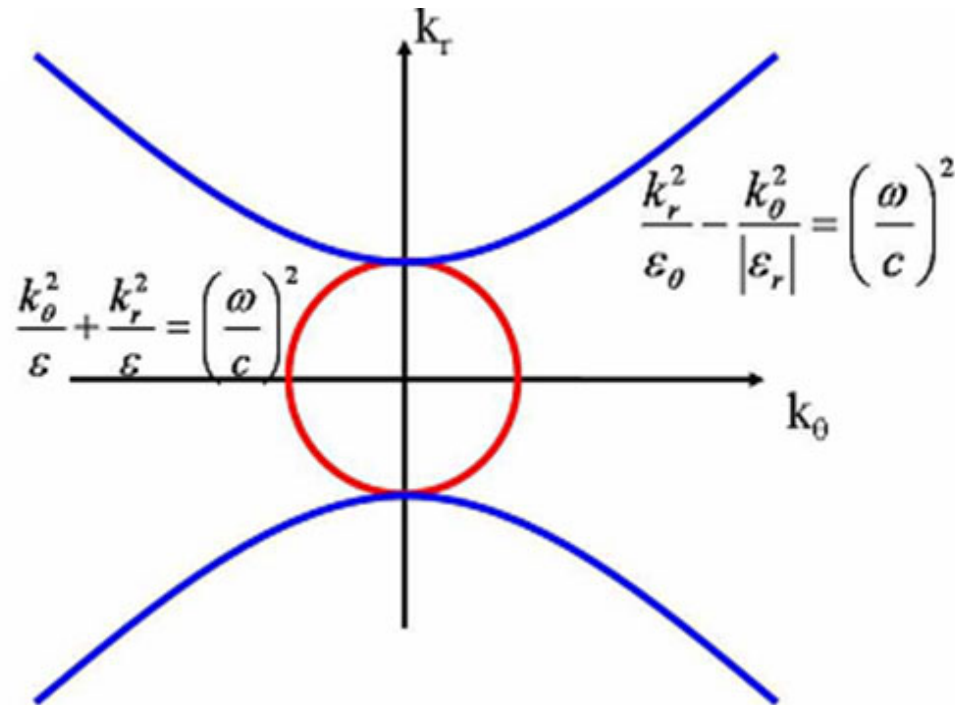
The spatial structure of light with OAM closely resembles the shape of atomic wave functions. It could therefore make symmetry-forbidden transitions possible in quantum dots, or “artificial atoms” [Eur. Phys. J. D 66, 157–165 (2012)]

However, the symmetry of the transition probability integral can be modified by choosing the spatial properties of the driving field, making symmetry-forbidden transitions possible.

Yet, the vanishing intensity in the center of an OAM beam usually makes this effect weak.

In order to strengthen the interaction an OAM beam size has to be reduced to subwavelength dimensions.

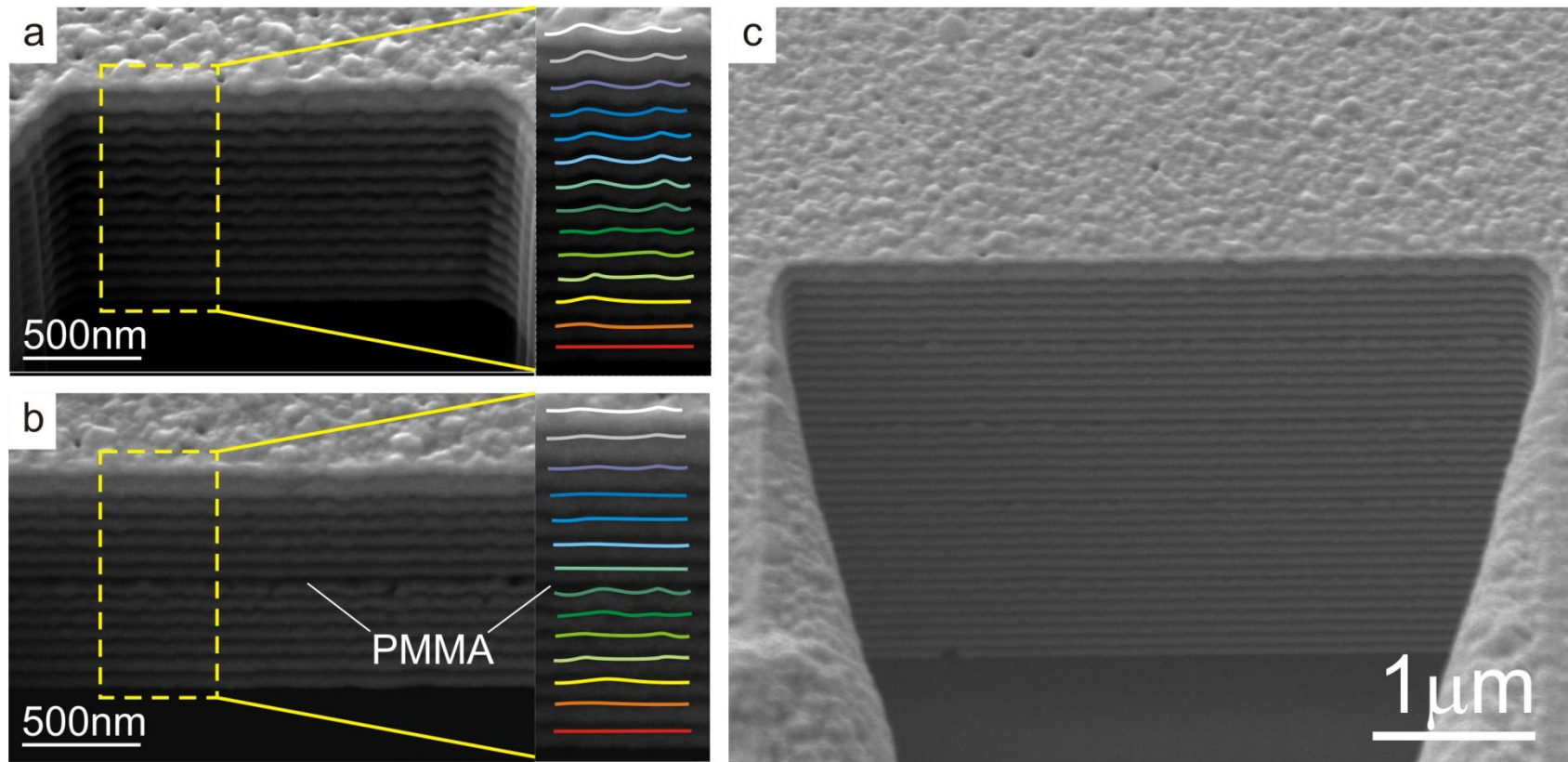




$$\tilde{\epsilon} = \begin{bmatrix} \epsilon_r & 0 \\ 0 & \epsilon_\theta \end{bmatrix}$$

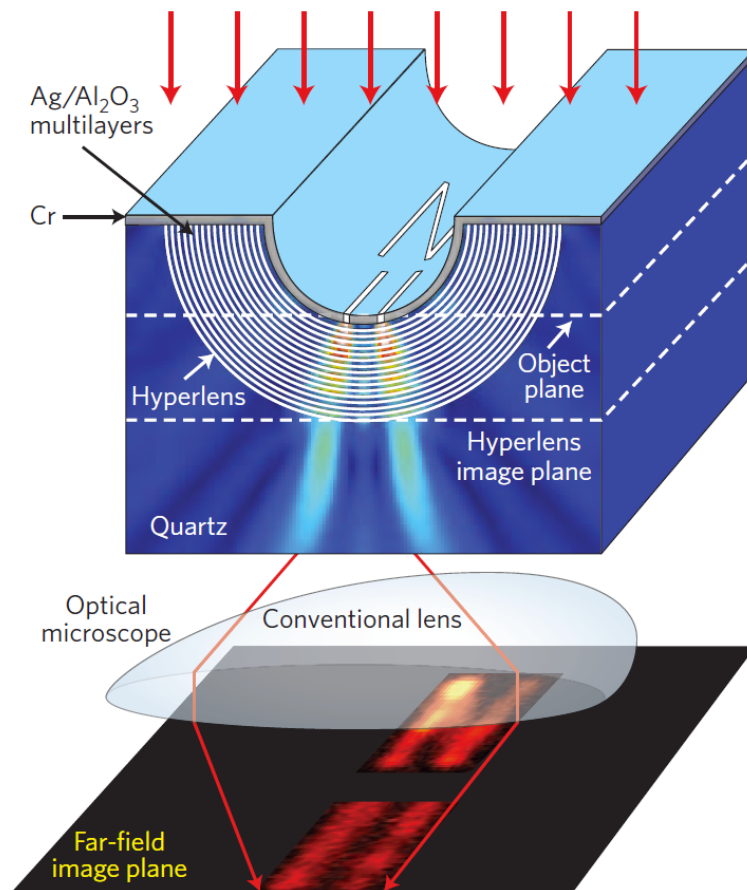
$$\epsilon_r < 0, \quad \epsilon_\theta > 0$$

- The most important property of HMMs is related to the behavior of waves with large magnitude wavevectors.
- In vacuum, such large wavevector waves are evanescent and decay exponentially.
- In HMMs waves with infinitely large wavevectors are allowed to propagate.

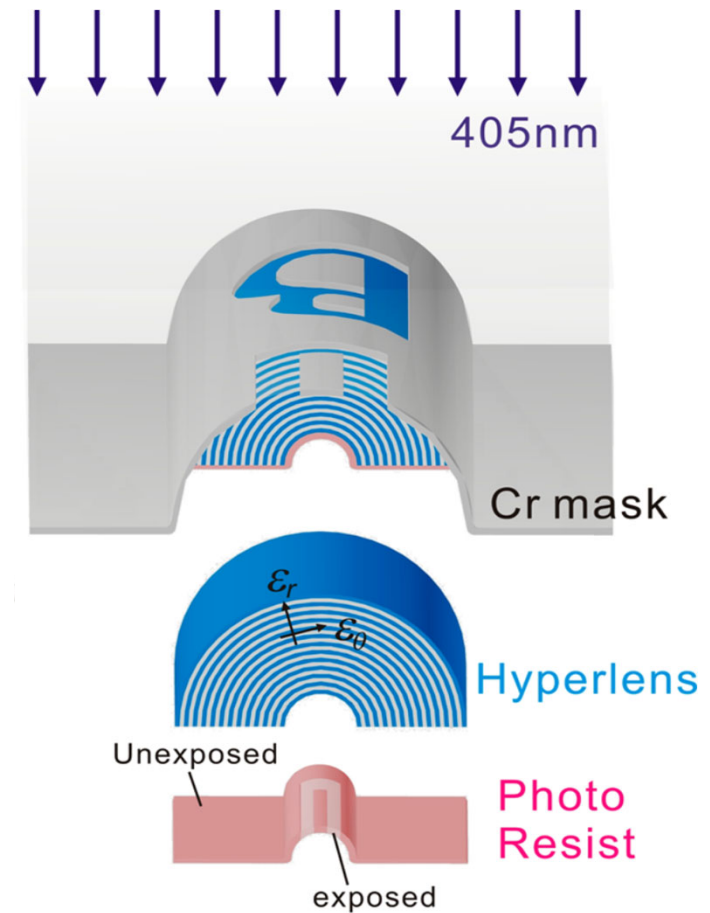


- (a) 27 layers Ag/Ti₃O₅ with no intermediate PMMA layer;
- (b) Two-step deposition of 26 layers Ag/Ti₃O₅ with a PMMA layer separating two 13-layered structures;
- (c) Five-step deposition of 65 layers of Ag/Ti₃O₅ with 4 intermediate PMMA layers separating every 13 layers.

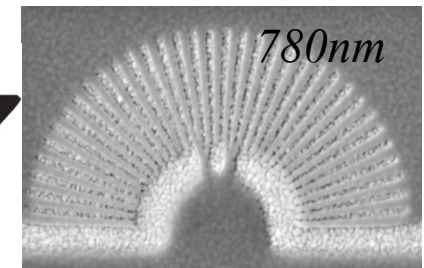
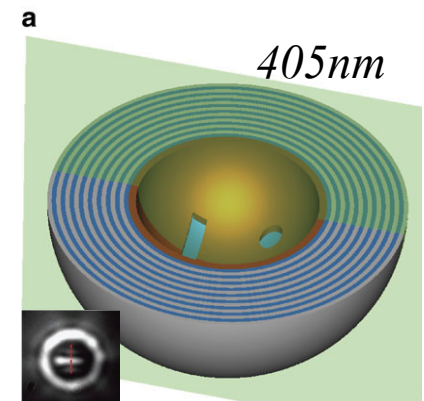
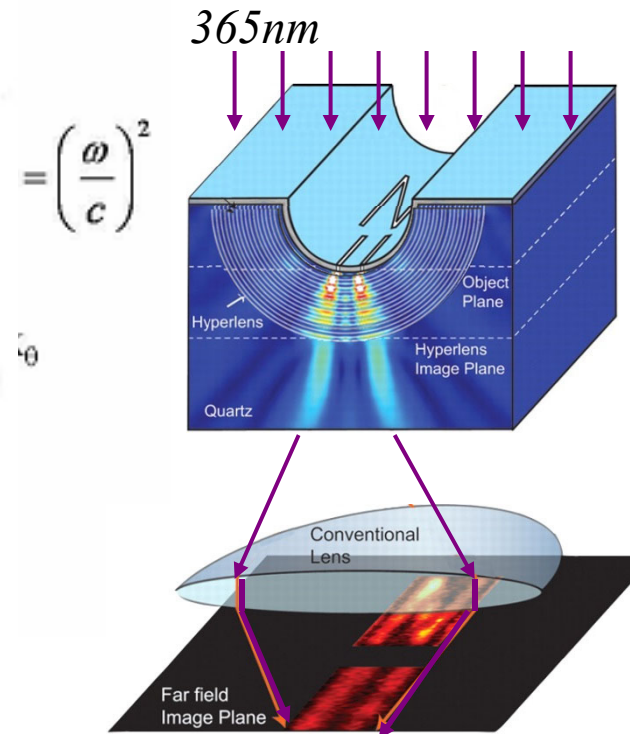
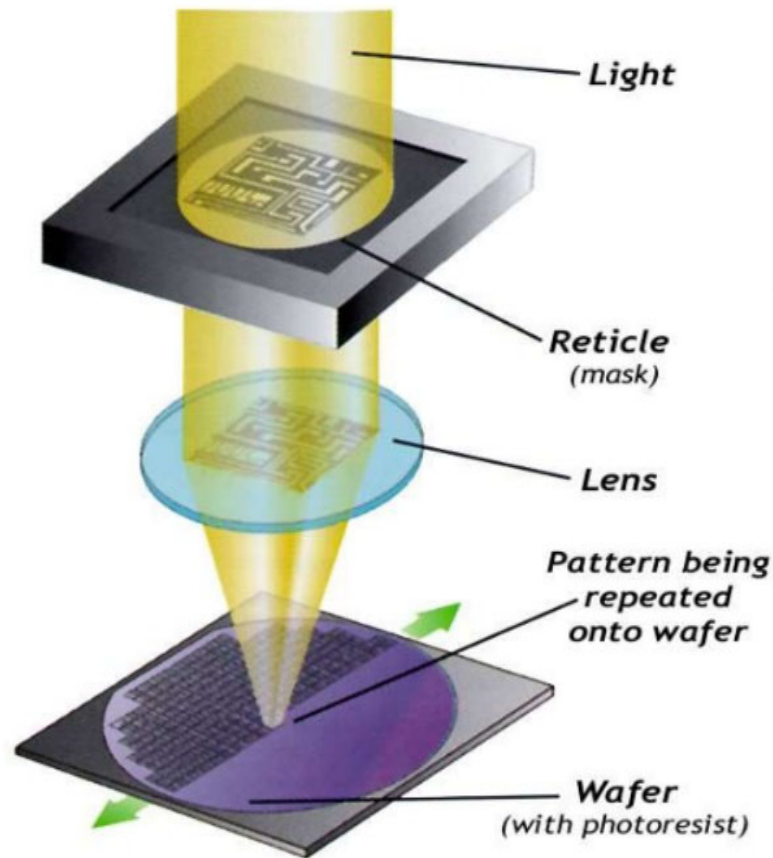
Magnifying hyperlens



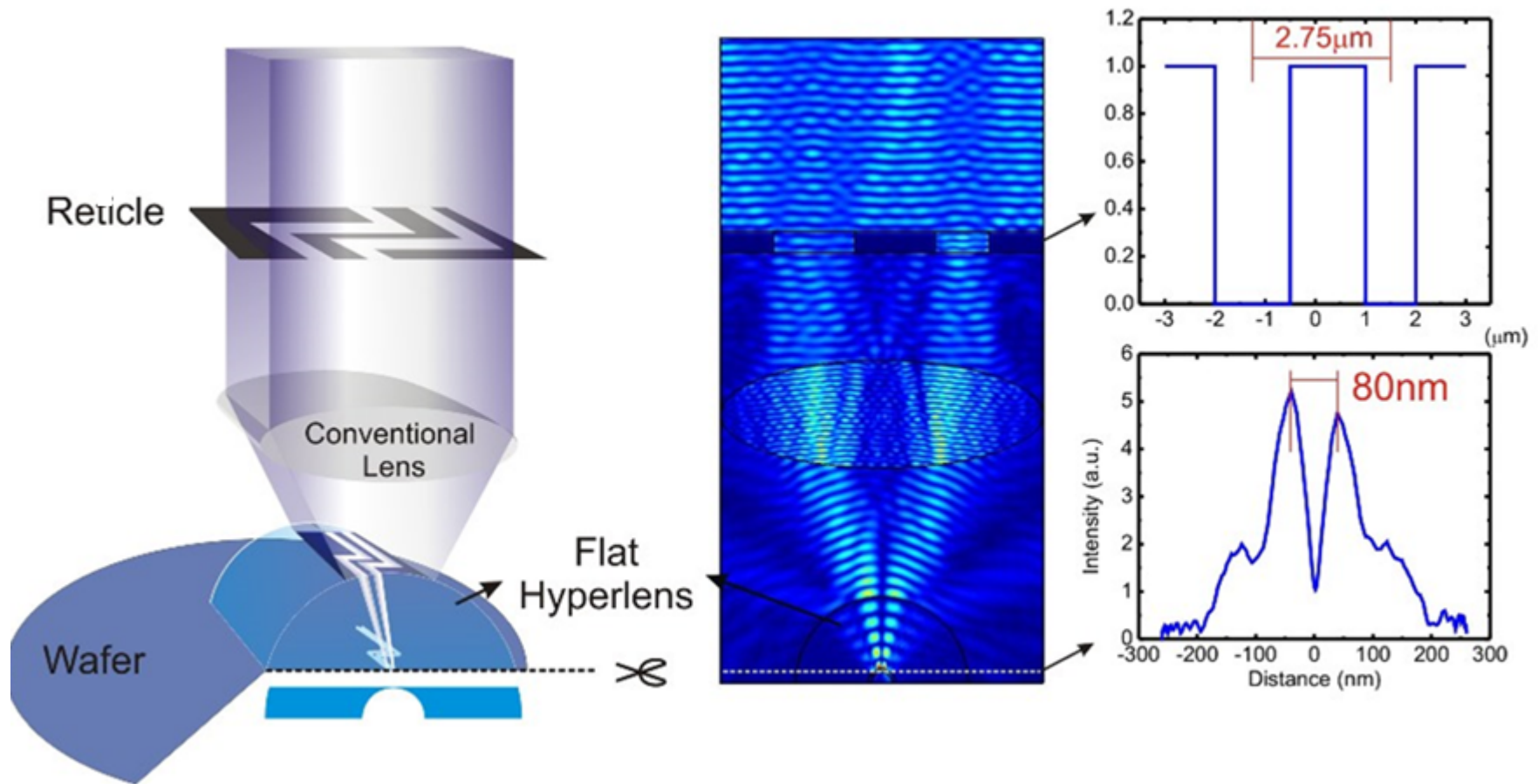
De-magnifying hyperlens



A promising way to achieve subwavelength imaging by supporting the transmission of the evanescent wave



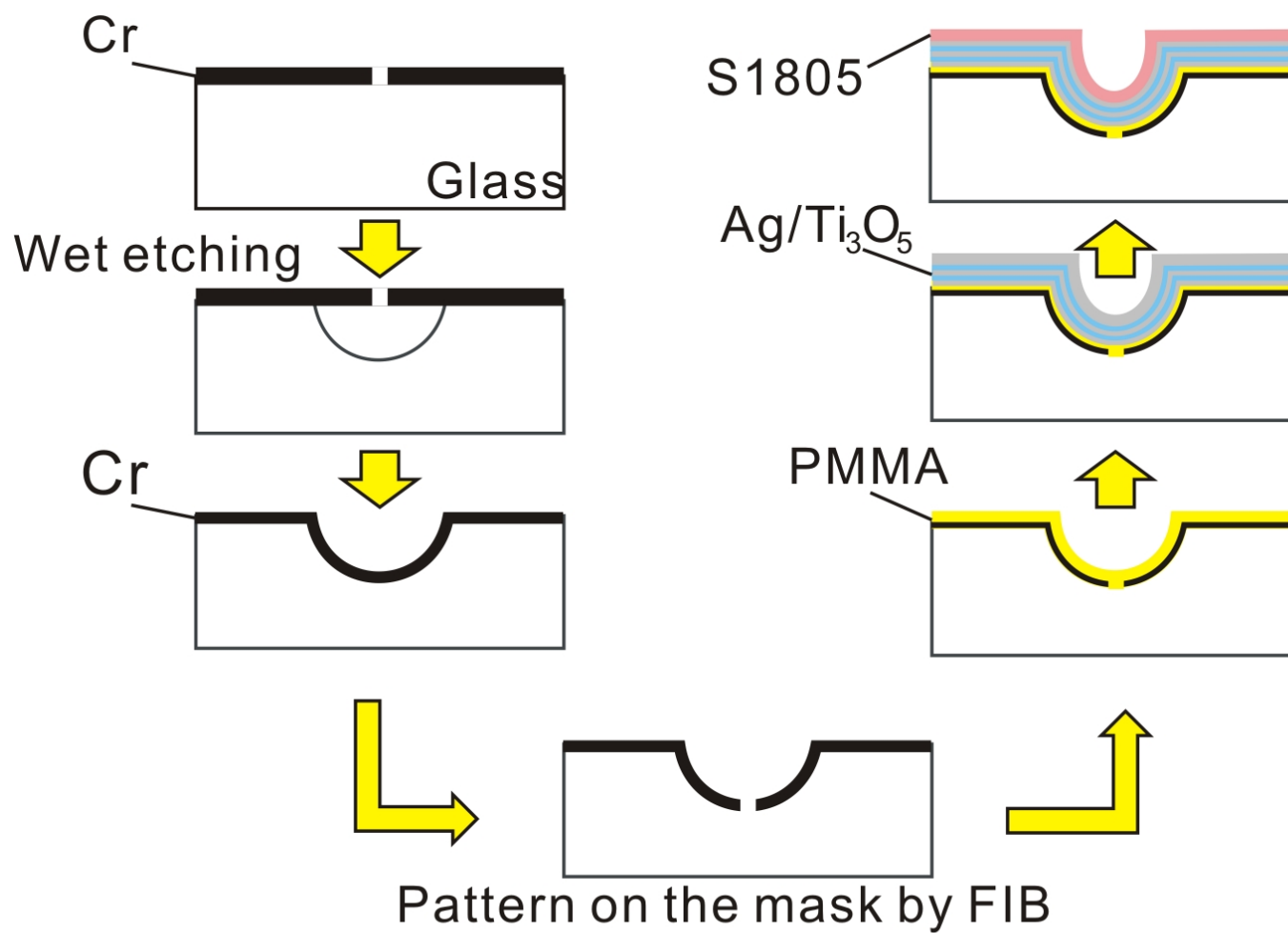
- Jacob, Alekseyev, Narimanov, Opt. Exp. **14**, 8247–8256 (2006).
 Salandrino, Engheta Phys. Rev. B **74**, 075103 (2006).
 Liu, Y. Xiong, Science **315**, 1686 (2007).
 Liu, Y. Xiong, Science **315**, 1686 (2007).
 Rho, Xiong, Nat. Comm. **1**, 143 (2010).
 J. Sun, W. Dai, Opt. Exp. **23**, 18510 (2015).



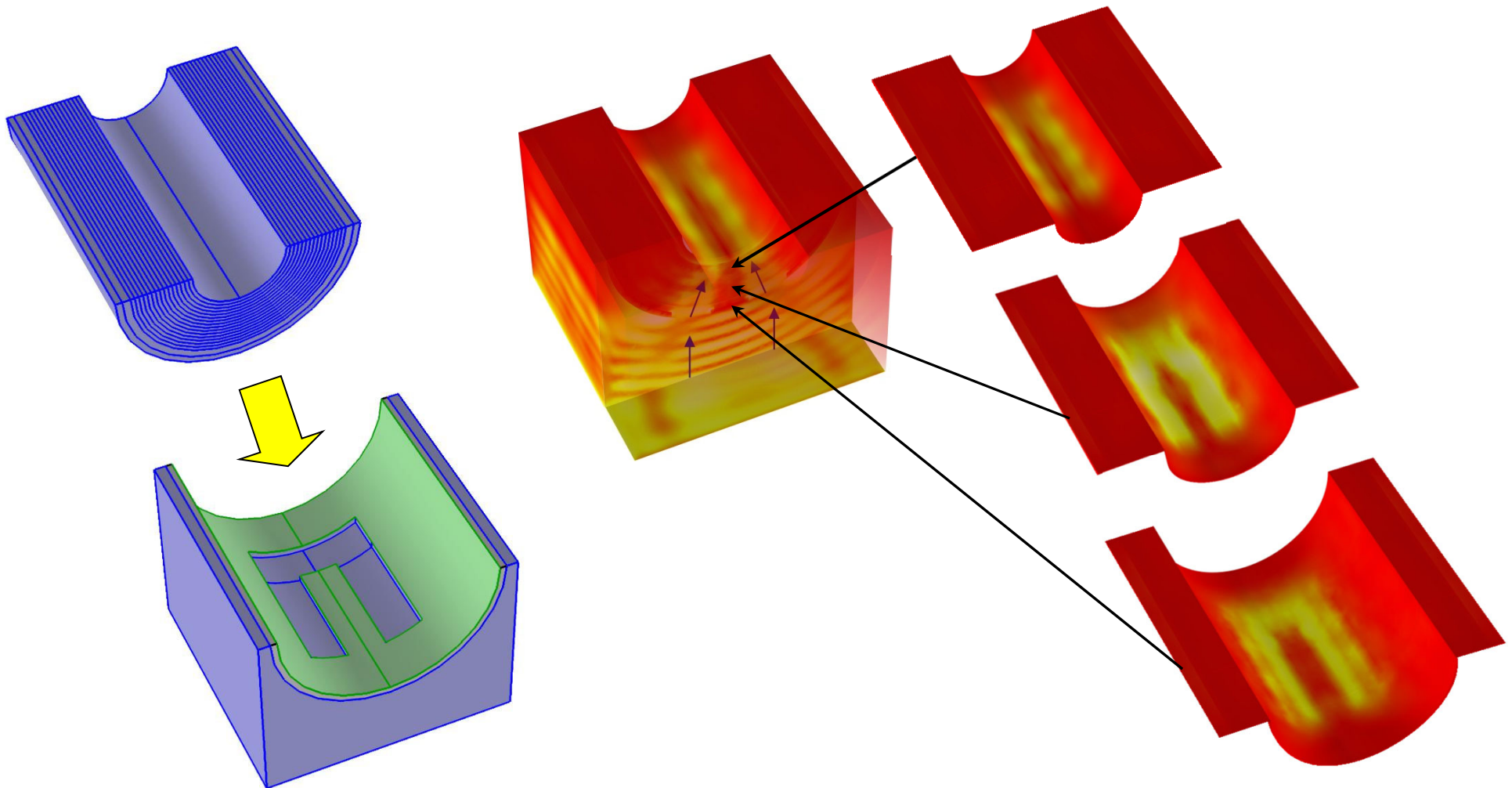
J. Sun, NML, Nano Lett. 2016

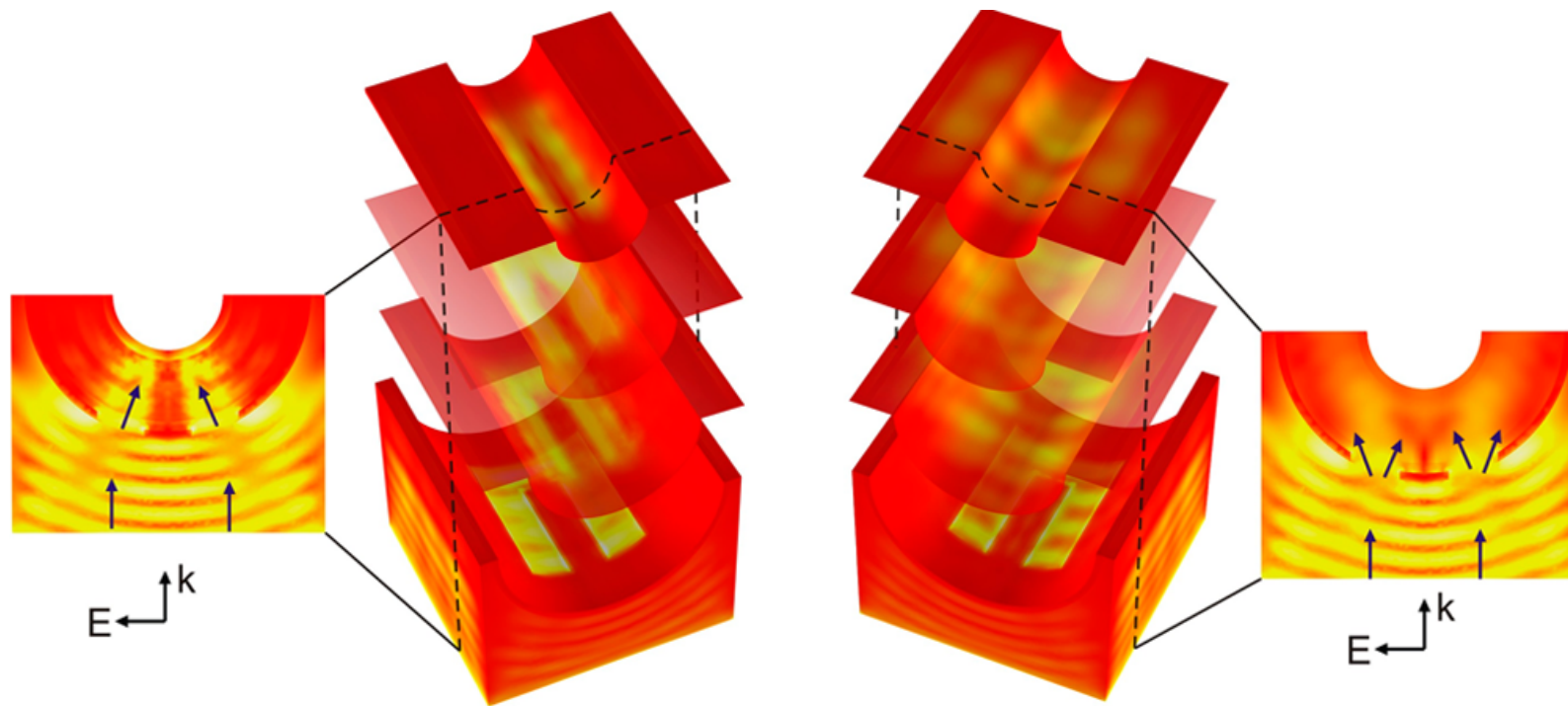
J. Sun, NML, ACS Nano. 2018

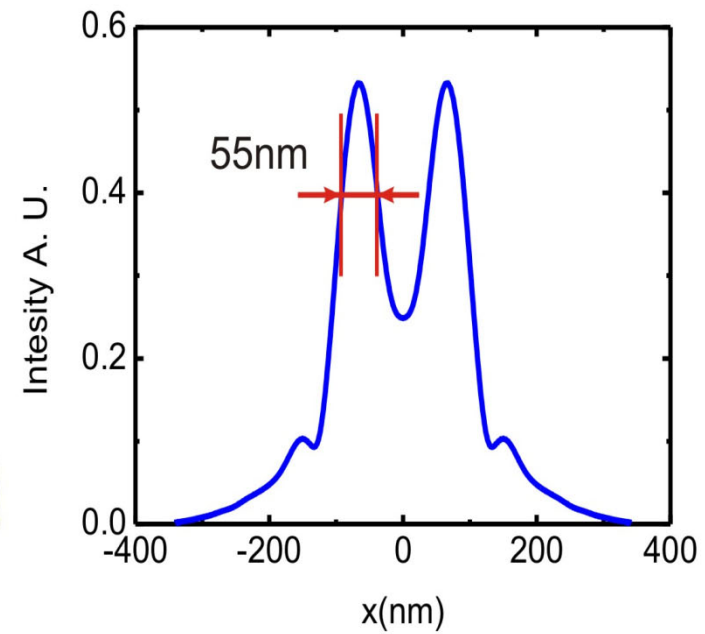
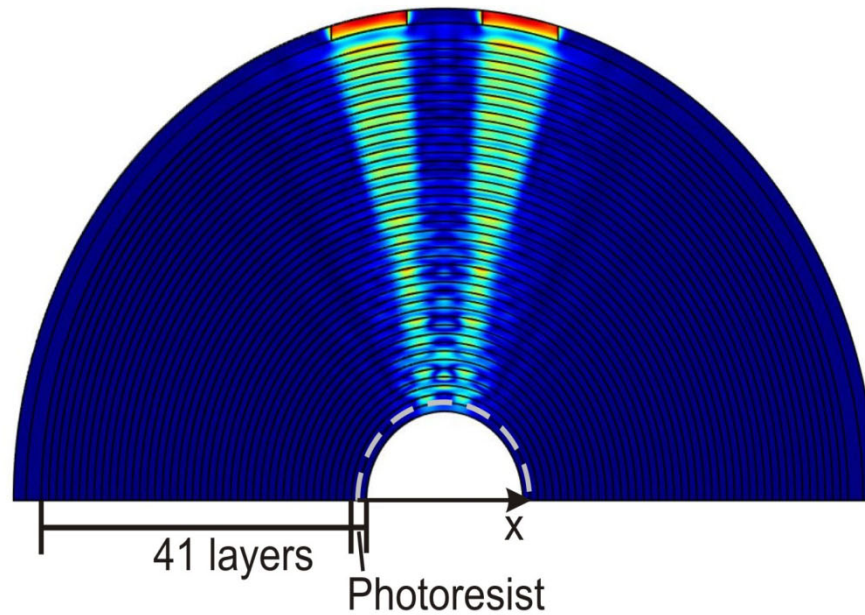
J. Sun, NML, J. Opt. 2018

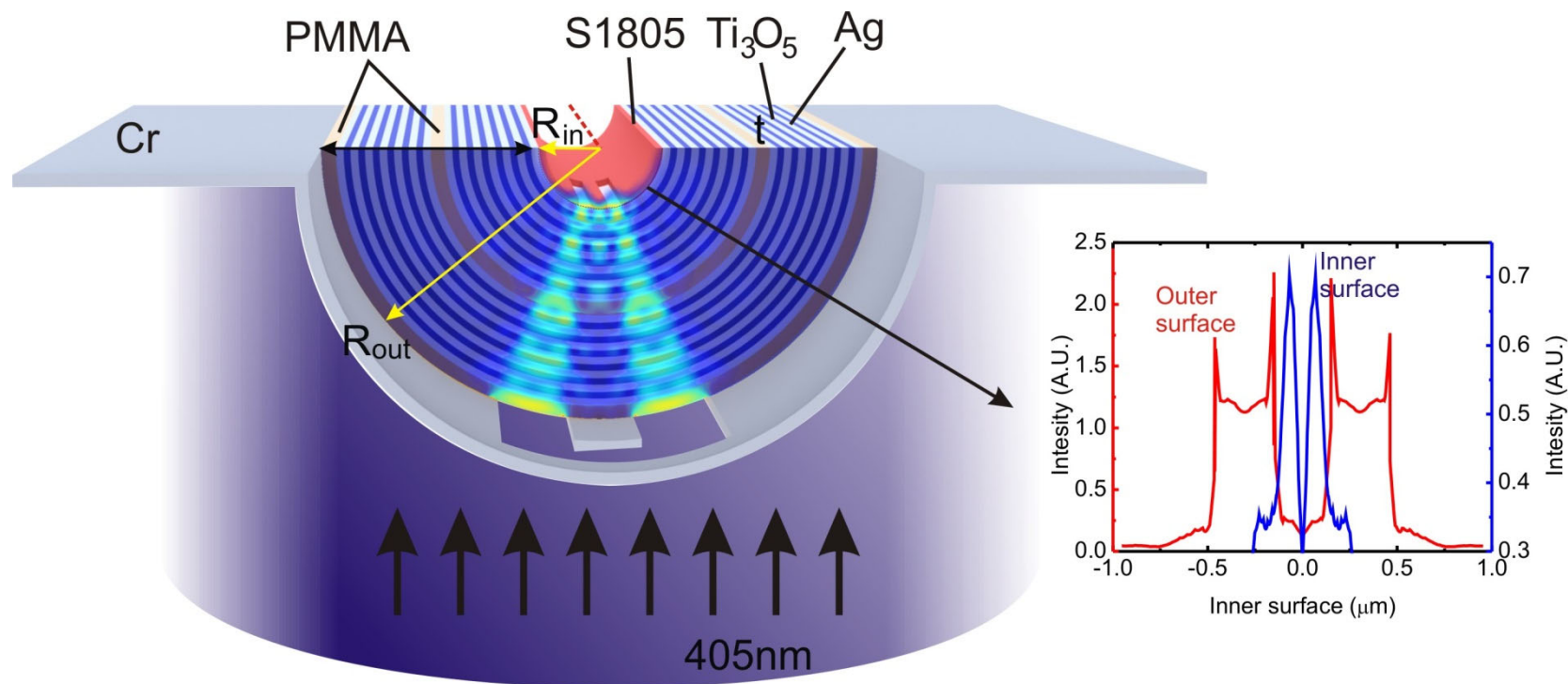


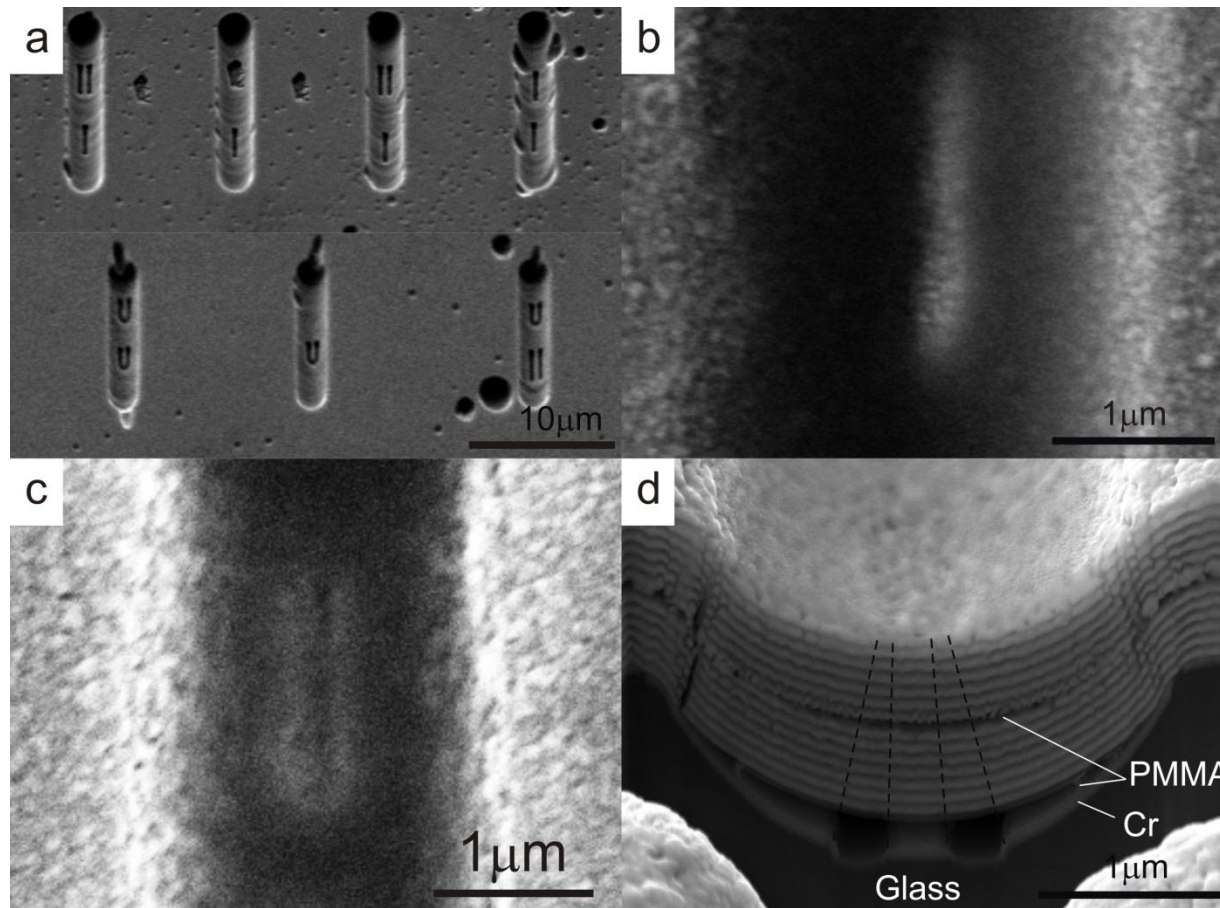
Simulations of the wave propagation in a hyperlens - de-magnification











Smallest feature size is 80 nm, corresponding to 3.75 de-magnification

An experimentally realized de-magnifying hyperlens consisted of 13 layers of alternating layers of Ag and Ti_3O_5 . The permittivities of the Ti_3O_5 and Ag at 405nm are 5.85 and $-4.84+0.22i$, respectively.

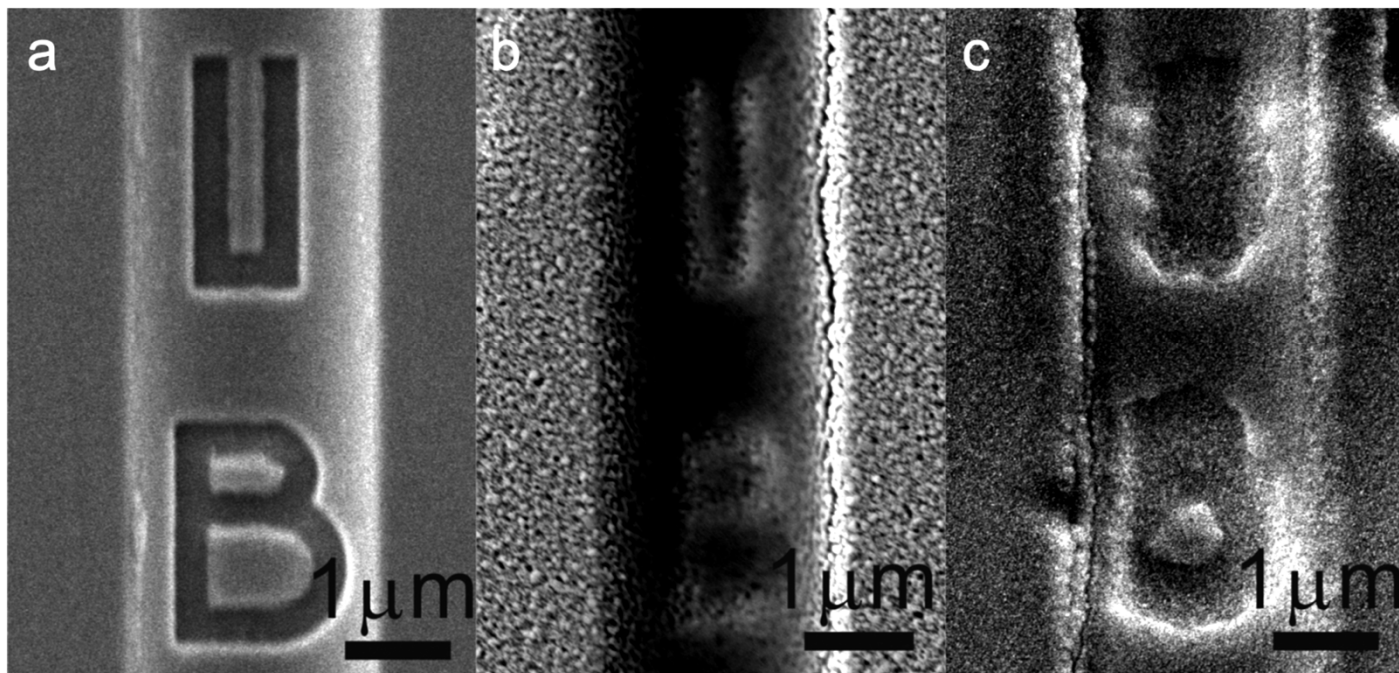
$$\varepsilon_\theta = \varepsilon_d f + \varepsilon_m (1-f) = 0.51+0.11i$$

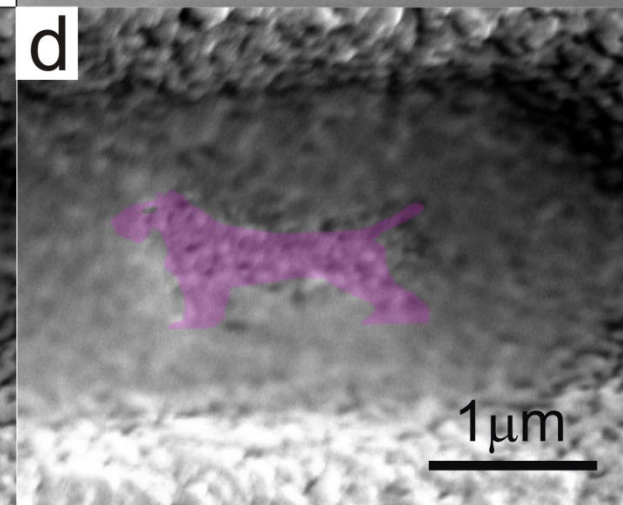
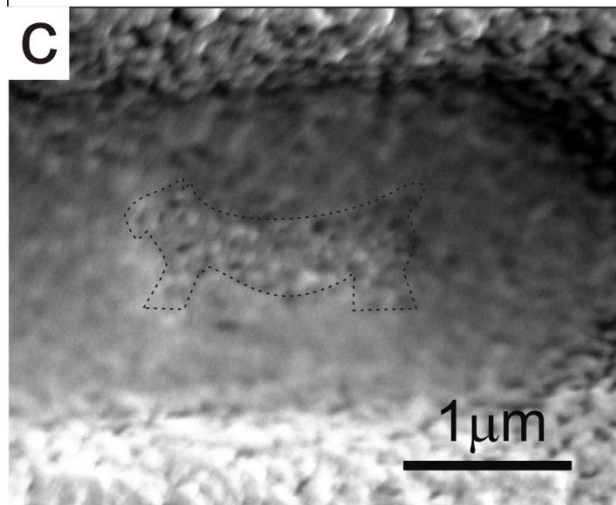
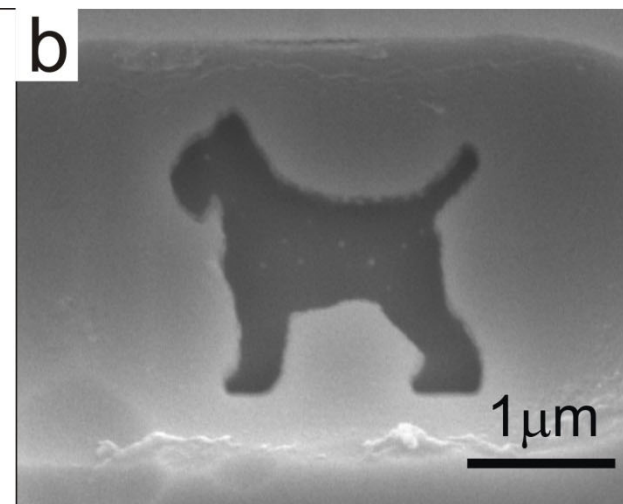
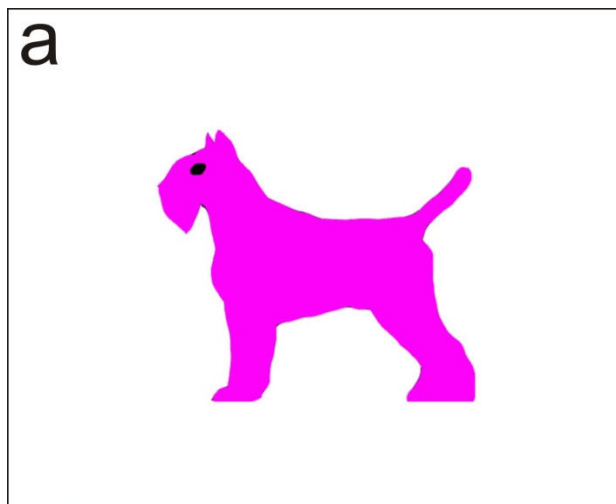
$$\varepsilon_r = \frac{\varepsilon_d \varepsilon_m}{\varepsilon_m f + \varepsilon_d (1-f)} = -53+14.1i$$

Original Pattern

With hyperlens

Reference (MgF_2) sample





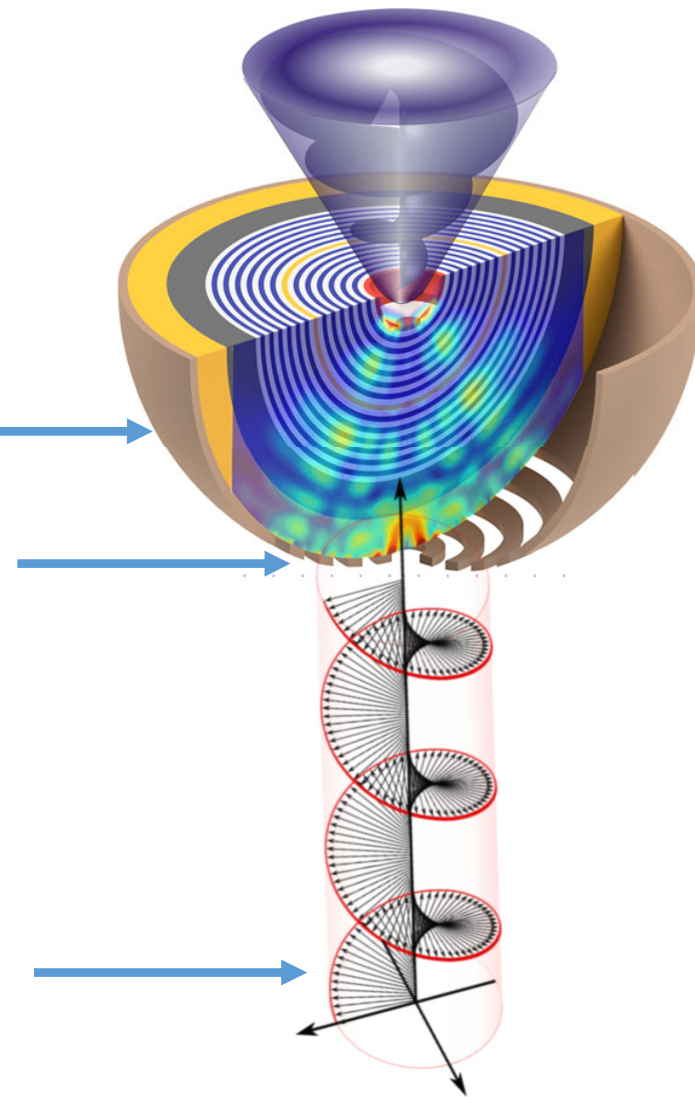
Demagnifying hyperlens

$\rightarrow$ Focusing below diffraction limit
OAM-to-SAM conversion

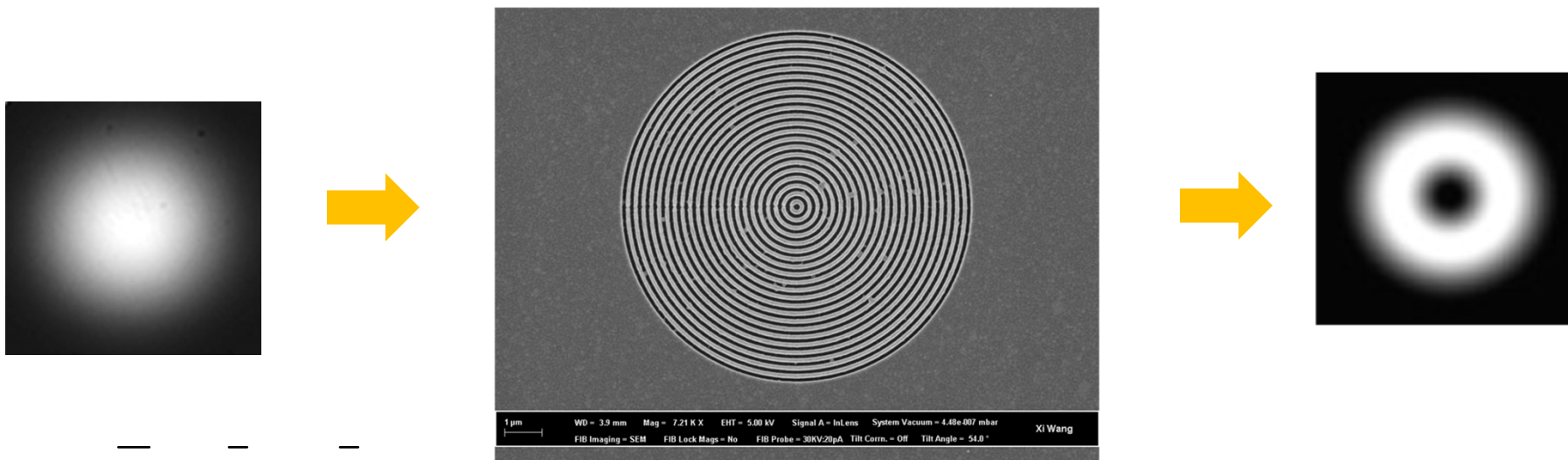
SAM-to-OAM converter

$\rightarrow$ Radially polarized OAM beam
with topological charge 1

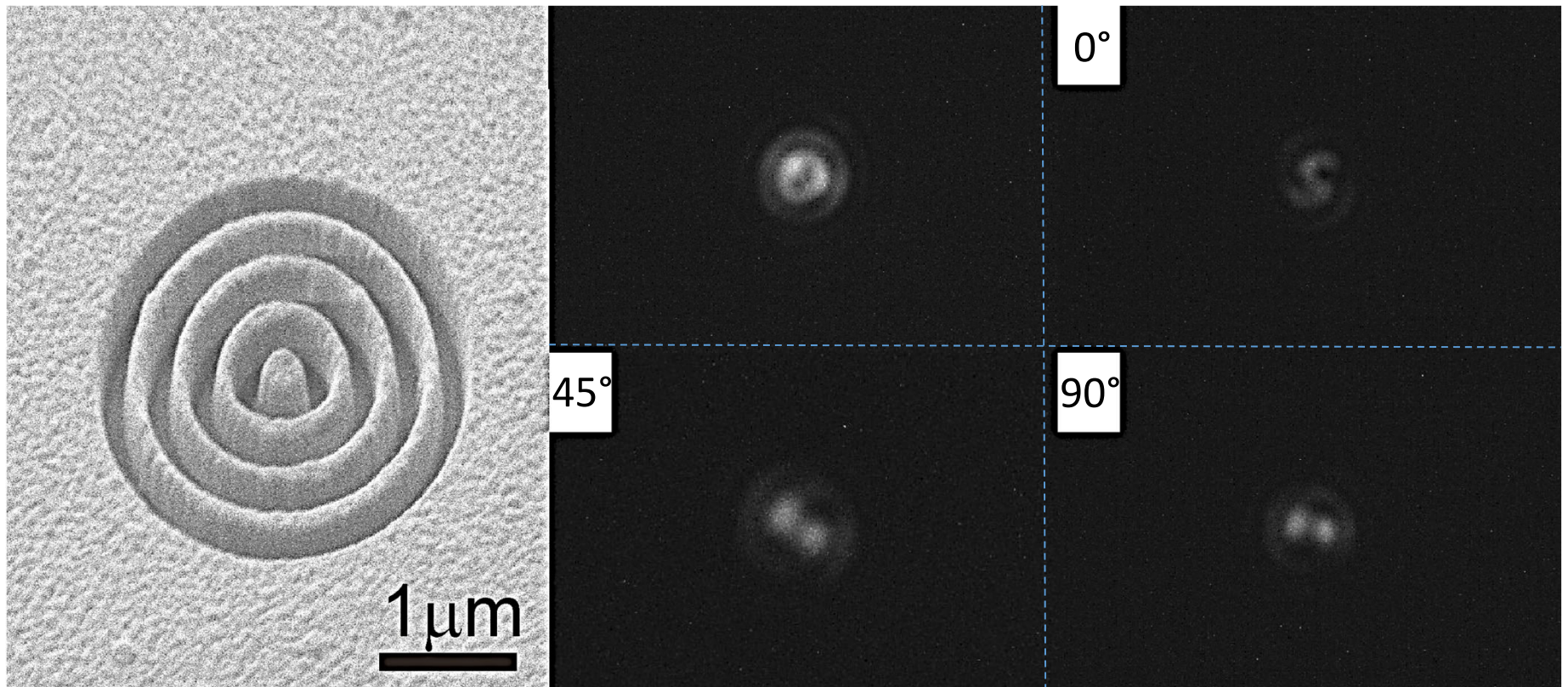
Circularly polarized Gaussian beam

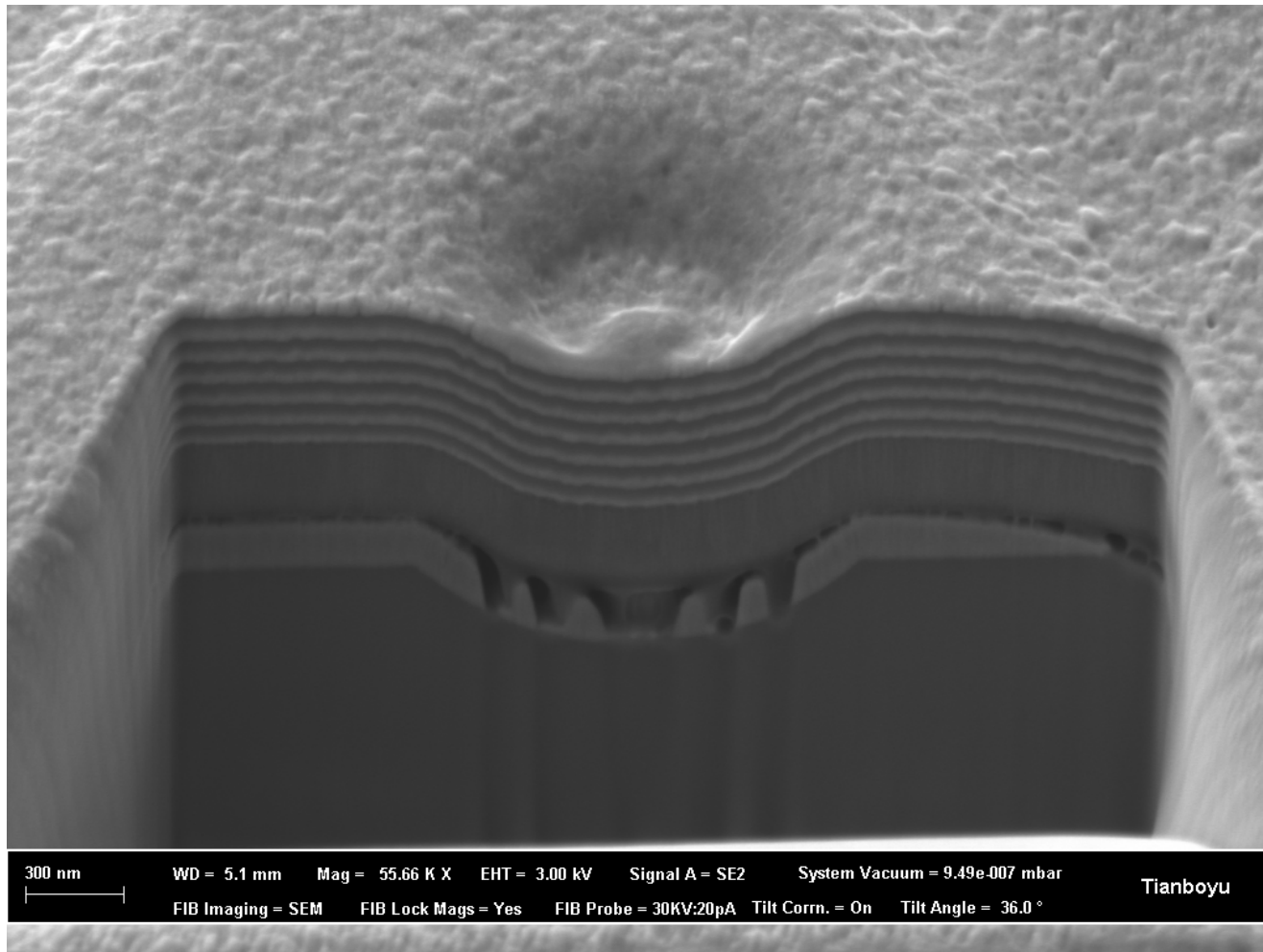


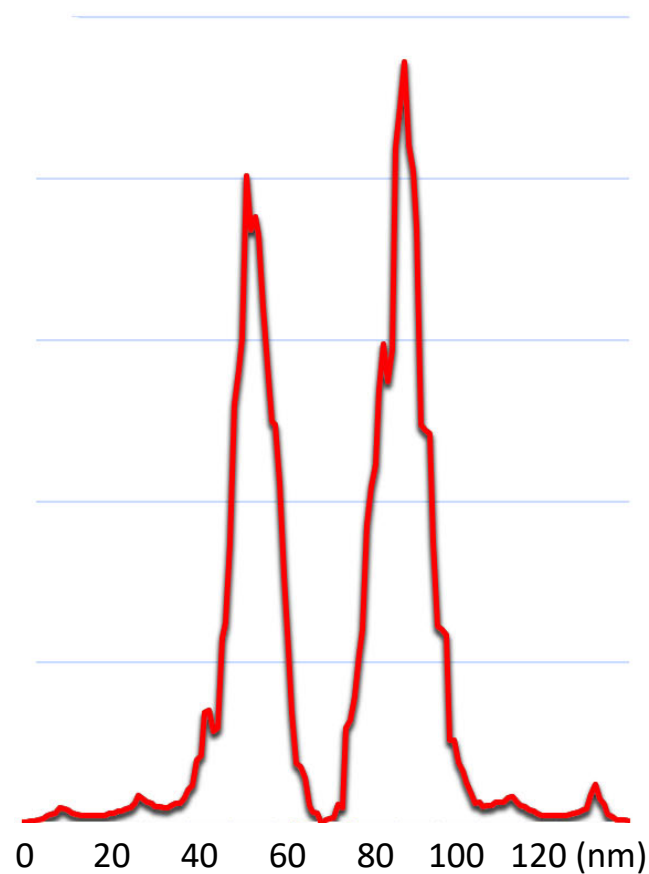
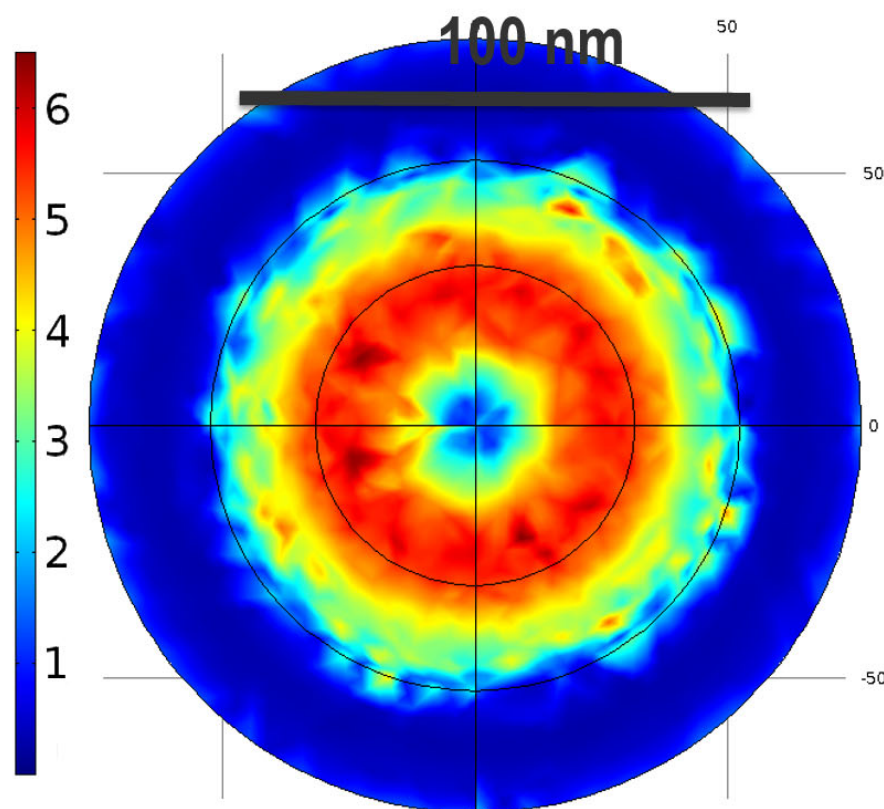
The circular polarization can be decomposed into the combination of radial polarization and azimuthal polarization components with a spiral phase wavefront. Concentric subwavelength metal grating reflects most of azimuthal polarization components while allow radial polarization vortex to pass through.

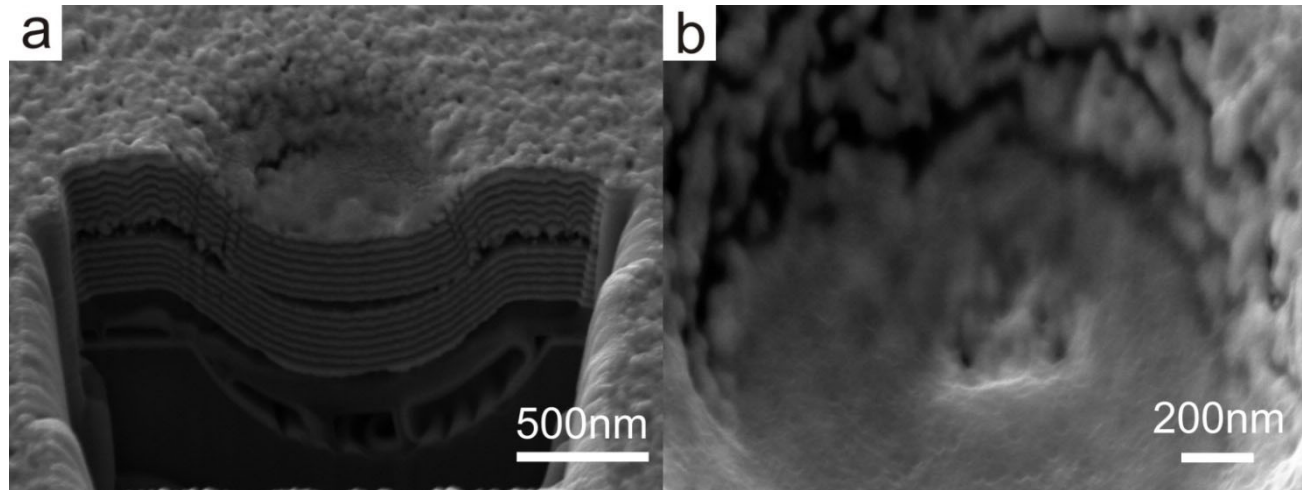


$$\begin{aligned}
 \bar{E}_{in} &= \bar{e}_x + j\bar{e}_y \\
 &= (\cos \phi \bar{e}_r - \sin \phi \bar{e}_\phi) + j(\sin \phi \bar{e}_r + \cos \phi \bar{e}_\phi) \\
 &= \bar{e}_r (\cos \phi + j \sin \phi) + j\bar{e}_\phi (\cos \phi + j \sin \phi) \\
 &= e^{j\phi} (\bar{e}_r + j\bar{e}_\phi)
 \end{aligned}$$

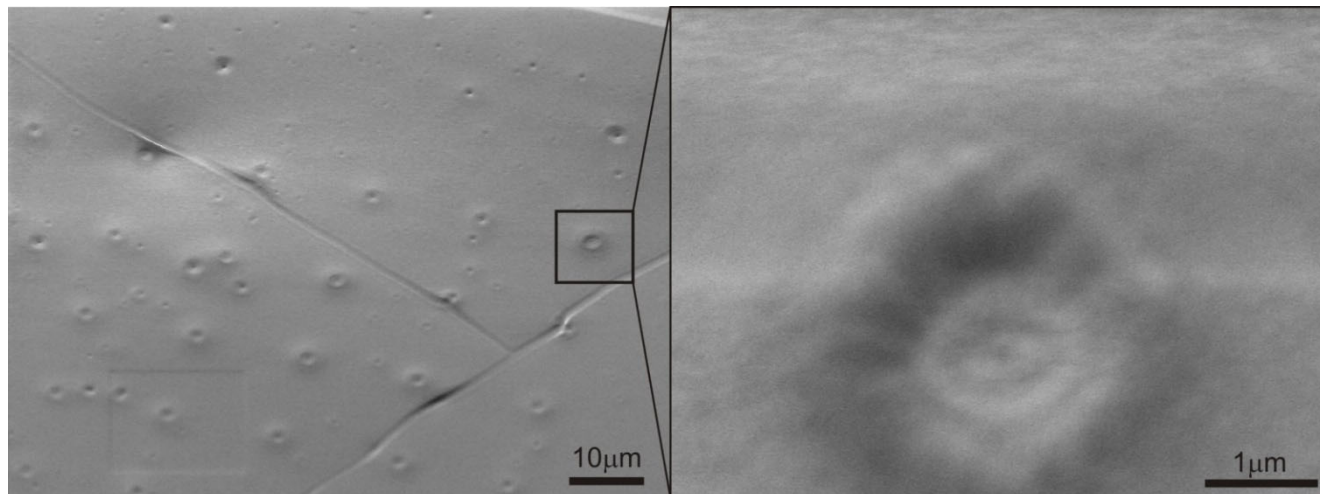




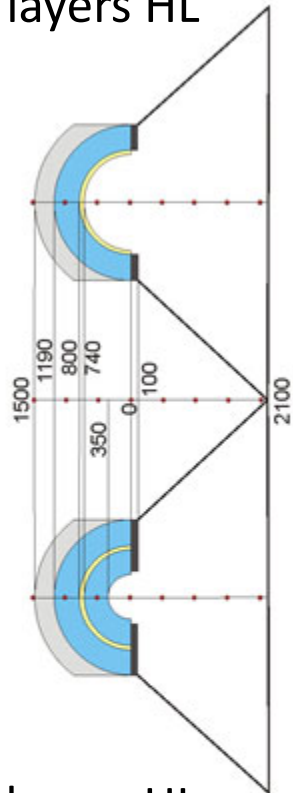




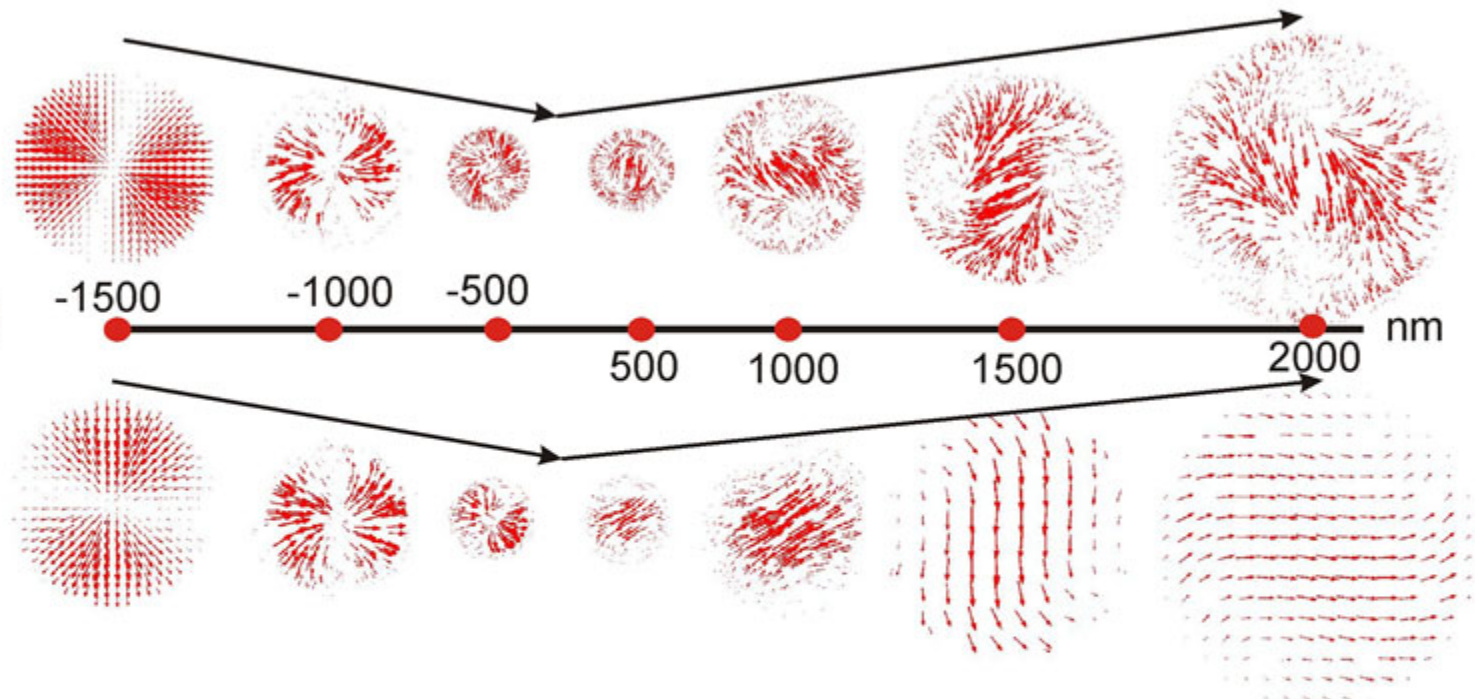
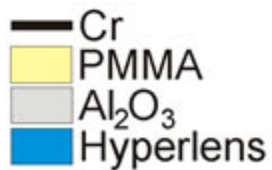
300 nm doughnut-shaped developed pattern on the photoresist

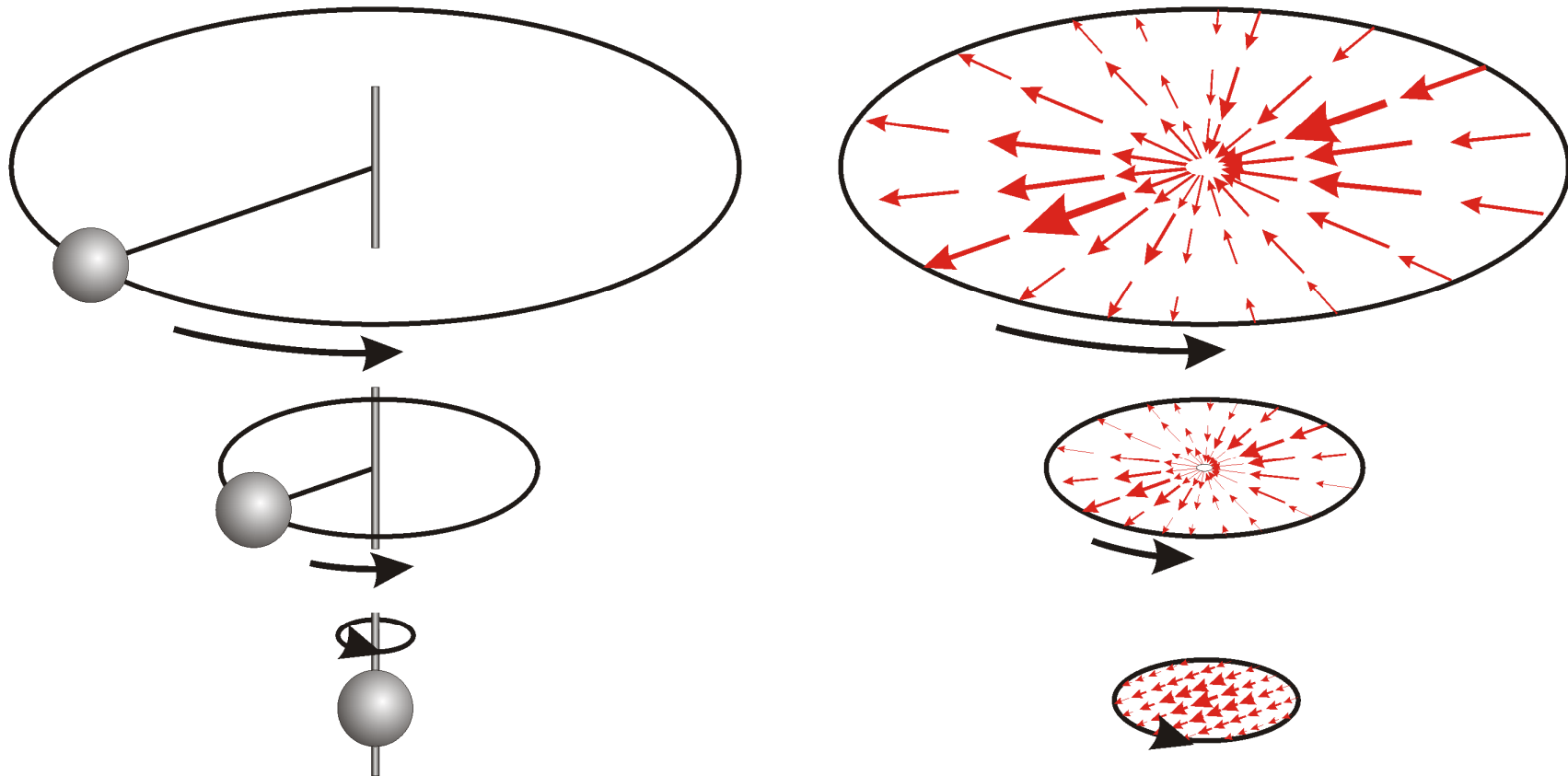


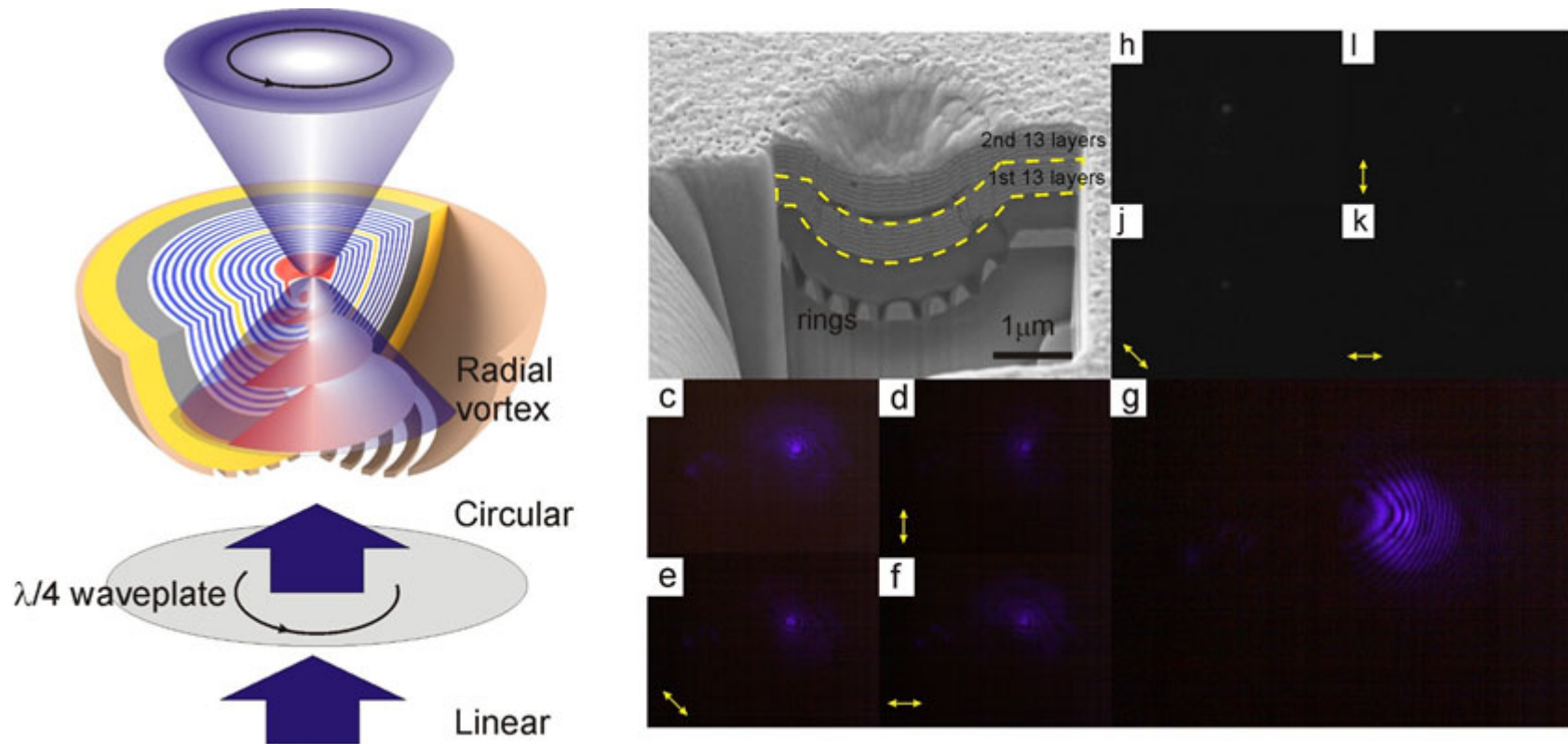
13 layers HL



26 layers HL



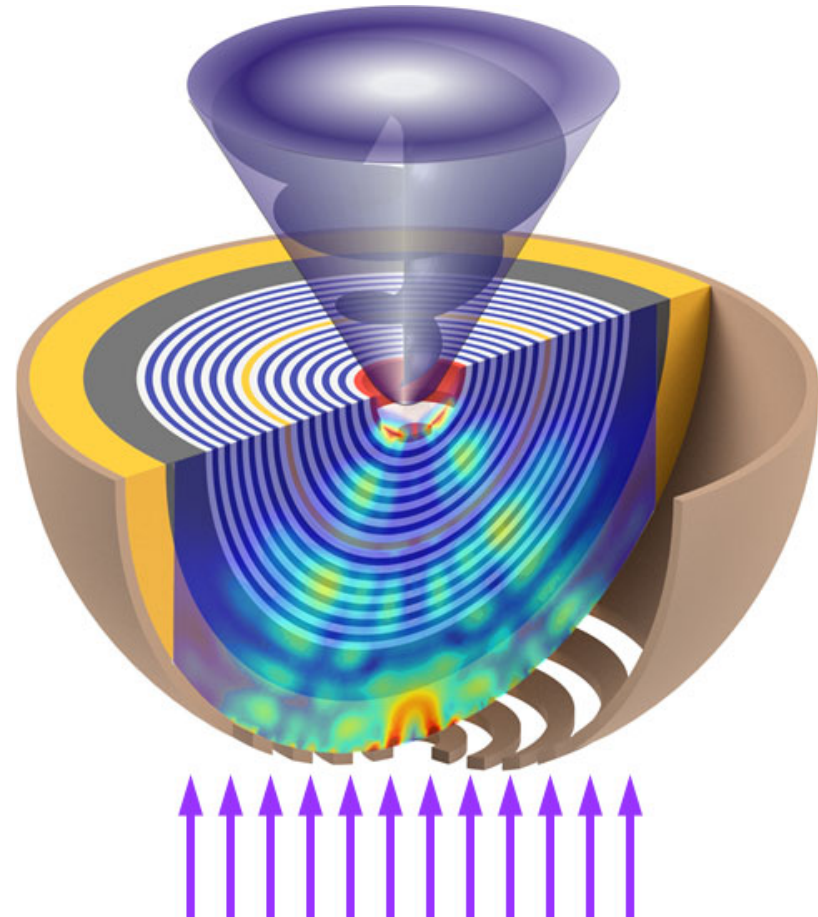
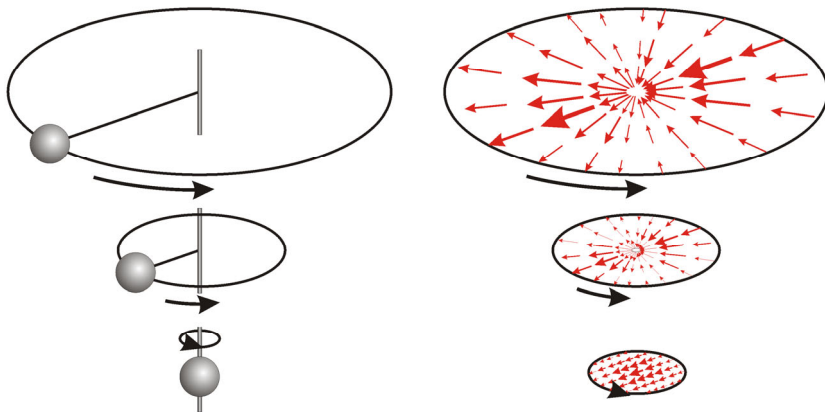




- Incident beam of circular polarization is transformed into a radially polarized vortex beam at the input and then focused by the hyperlens
- The output beam is circularly polarized with no OAM

We predicted and demonstrated the OAM-to-SAM conversion that can be

- achieved by focusing a radially polarized OAM beam down to the subwavelength scale using a hyperlens
- understood using a simple mechanical “throwing hammer” analogy



This work was supported by the NSF award Award 1809518