

Proton Network Flexibility Enables Robustness and Large Electric Fields in the Ketosteroid Isomerase Active Site

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Abstract

Hydrogen bond networks play vital roles in biological functions ranging from protein folding to enzyme catalysis. Here we combine electronic structure calculations and *ab initio* path integral molecular dynamics simulations, which incorporate both nuclear and electronic quantum effects, to show why the network of short hydrogen bonds in the active site of ketosteroid isomerase is remarkably robust to mutations along the network and how this gives rise to large local electric fields. We demonstrate that these properties arise from the network’s ability to respond to a perturbation by shifting proton positions and redistributing electronic charge density. This flexibility leads to small changes in properties such as the partial ionization of residues and pK_a isotope effects upon mutation of the residues, consistent with recent experiments. This proton flexibility is further enhanced when an extended hydrogen bond network forms in the presence of an intermediate analog, which allows us to explain the chemical origins of the large electric fields in the enzyme’s active site observed in recent experiments.

Introduction

Hydrogen bonds play essential roles in biological systems, carrying out diverse functions ranging from stabilizing protein structures, to modulating molecular recognition and catalyzing enzymatic reactions.¹⁻⁷ An important class of these are low-barrier hydrogen bonds (LBHBs), which can occur when the donor-acceptor distance between heavy atoms, R , is below 2.6 Å and the proton affinities of the hydrogen bonding partners are closely matched.⁸⁻¹⁹ LBHBs have been implicated in a wide variety of biological processes such as catalyzing chemical reactions,^{12,20} promoting structural stability²¹ and facilitating proton transfer.^{22,23} LBHBs are more commonly observed in proteins than in liquids as the protein’s three-dimensional fold can help position the residues closer together. While the values for R in LBHBs are only slightly shorter than typical hydrogen bonds in liquids, such as water where $R \sim 2.8$ Å, this shortening significantly alters the potential energy surface for transferring the proton between the donor and acceptor. As an illustration, consider an O–H bond with a typical length of 1 Å. When $R = 2.8$ Å, the distance between the two minima, arising from the H being bound to the donor or to the acceptor, in the potential curve is ~ 0.8 Å. Decreasing R to 2.6 Å shortens this distance to ~ 0.6 Å. This $\sim 25\%$ decrease brings the distance between the potential energy minima close to the proton’s thermal wavelength. In these cases the proton transfer barrier can become comparable to the zero point energy (ZPE) for the Donor–H bond (typically ~ 5 kcal/mol), allowing the proton to delocalize between the donor and acceptor. Upon further compression of the donor-acceptor distance ($R < 2.4$ Å), the potential energy surface for proton transfer becomes increasingly flat such that even the classical thermal energy is sufficient to allow proton delocalization.

While many biological systems exploit single hydrogen bonds in their functions, others invoke networks of hydrogen bonds.^{14,22-26} The enzyme ketosteroid isomerase (KSI) provides one such intriguing system, which contains a network of short hydrogen bonds in its active site.²⁷⁻²⁹ KSI catalyzes the isomerization of steroid substrates with a 10^{11} -fold rate enhancement compared to the analogous reaction in solution.³⁰⁻³⁶ Its catalytic mechanism

involves proton transfer from the substrate to the enzyme to form a negatively charged intermediate, which is stabilized by the enzyme’s hydrogen bond network. This is followed by proton donation from the enzyme back to the substrate at a different position, leading to a net isomerization of the substrate (Fig. S1).^{33,37} KSI’s ability to stabilize its reaction’s intermediate through the active-site hydrogen bond network can be investigated using its mutant, KSI^{D40N}, which preserves the structure of the wild-type enzyme and mimics its ionization state in the enzyme-intermediate complex.^{27–29,38,39} In the absence of an intermediate analog, the side chain of residues Tyr16, Tyr32 and Tyr57 in KSI^{D40N} form a hydrogen bonded triad with $R \sim 2.6$ Å, as shown in Fig. 1a.^{27,29} This network of LBHBs promotes quantum proton delocalization and facilitates the deprotonation of Tyr57, leading to a very high acidity of this residue ($pK_a=5.8$;⁴⁰ tyrosine in aqueous solution has a pK_a of 10.1) and a prominent H/D isotope effect ($\Delta pK_a=0.9$;⁴⁰ typical organic acids have a ΔpK_a of 0.4-0.6⁴¹). Upon binding of an inhibitor, an extended network of short hydrogen bonds is formed to incorporate the inhibitor and residue Asp103 (Fig. 1b).^{28,29,36,42,43} KSI^{D40N} thus provides an ideal system to assess the behavior of networks of LBHBs in a biological system. The donor-acceptor distances, R , are short and the pK_a ’s of the hydrogen bond participants in solution are closely matched, suggesting that they share similar proton affinities in the network. These are generally considered as two key prerequisites for forming LBHBs.^{8,11,12,17}

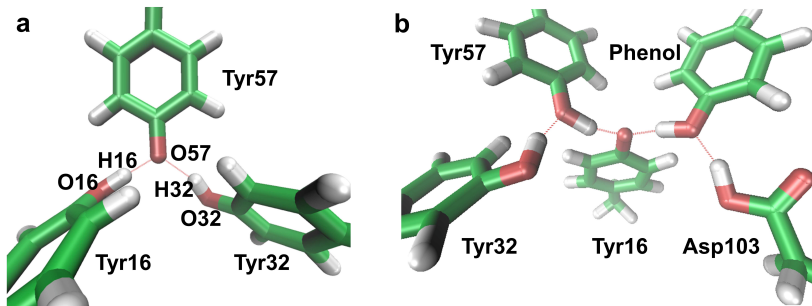


Figure 1: Hydrogen bond network in the active site of KSI^{D40N}. (a) Local network formed from residues Tyr16, Tyr32 and Tyr57 with a triad structural motif, with the hydrogen bonded oxygen and hydrogen atoms labeled. (b) Extended network formed in the presence of an intermediate analog phenol. Red, green and white represent oxygen, carbon and hydrogen atoms, respectively.

Recent experiments have revealed a number of intriguing observations about how KSI responds to chemical perturbations in its active site.^{40,44–49} In particular, mutagenesis experiments have been conducted that employ unnatural amino acids to perturb R and the proton affinities of the active-site residues while keeping the overall hydrogen bond network intact. These studies have shown that KSI’s catalytic rate and the properties of the network are remarkably robust with respect to such perturbations. For example, a recent experiment substituted Tyr16 with a series of fluorotyrosines and found that, despite the fact that fluorination results in a 40-fold increase in tyrosine’s acidity in aqueous solution, only a 1.2-fold decrease in KSI’s catalytic rate constant was observed.⁴⁵ In another study, the tyrosine residues in the triad were substituted with their chlorinated counterparts (Fig. 2a and b), which led to only minor changes in properties such as the partial ionization and H/D isotope effects.⁴⁰ In addition, recent vibrational Stark effect spectroscopy experiments used a product-like inhibitor that binds to wild-type KSI and identified a very large electric field of -144 ± 6 MV/cm in the enzyme’s active site (compared to -80 ± 2 MV/cm in water).⁴⁶ The strength of the fields was demonstrated to strongly correlate with the enzyme’s catalytic rate enhancement and hence it was suggested that they may play a role in the enzyme’s function.^{46,49} However, despite the apparent importance of these fields, previous studies invoking both point charge and polarizable empirical potentials have observed average electric fields of less than half the experimental value.^{46,50}

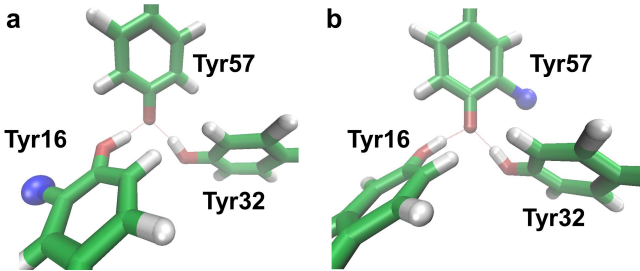


Figure 2: Perturbations to the active-site hydrogen bond network of KSI^{D40N}: (a) Tyr16 replaced by 3-chlorotyrosine and (b) Tyr57 replaced by 3-chlorotyrosine. Red, green, blue and white represent oxygen, carbon, chlorine and hydrogen atoms, respectively.

Here, we use QM/MM simulations to analyze the quantum fluctuations of the electrons

and nuclei in the active-site hydrogen bond network of KSI^{D40N} and uncover the origins of the robustness of the network in response to chemical perturbations. We then show that extending the size of the hydrogen bond network by binding an intermediate analog (Fig. 1b)²⁹ further enhances the movement of protons and electrons in the active site, resulting in the large electric fields acting on the intermediate analog. Finally, we discuss the possible biological implications of the ability of the extended hydrogen bond to respond to perturbations arising from mutations in the network.

Simulation methods

To examine the role of structural fluctuations, we carried out *ab initio* molecular dynamics (AIMD) simulations via an MPI interface⁵¹ to the TeraChem program.⁵² Since nuclear quantum effects (NQEs) are important in hydrogen bonded systems,^{23,43,53–66} we also examined the effect of NQEs using *ab initio* path integral molecular dynamics (AI-PIMD) simulations, which treat both the nuclei and electrons quantum mechanically.^{67–69}

AIMD and AI-PIMD simulations were performed using a QM/MM setup for the following systems: KSI^{D40N} (initial structure from PDB 1OGX),²⁷ KSI^{D40N} with Tyr16 or Tyr57 replaced by 3-chlorotyrosine (initial structures from PDB 5D82 and 5D81,⁴⁰ respectively), KSI^{D40N} with a bound phenol (initial structure modified from PDB 3VGN²⁹) and tyrosine in aqueous solution. Each simulation was performed in the NVT ensemble at 300 K, using a time step of 0.5 fs. The AI-PIMD simulations of KSI^{D40N} with a bound phenol were run for 30 ps. All the other simulations have the lengths of 20 ps. Electronic structure in the QM region was evaluated at the B3LYP level⁷⁰ with the D3 dispersion corrections,⁷¹ combined with the 6-31G* basis set. In the AI-PIMD simulations, each particle was represented by six path integral beads to account for nuclear quantum effects, using the path integral generalized Langevin equation approach.⁷² The MM region was modeled using classical dynamics with the AMBER03 force field⁷³ and the TIP3P water model.⁷⁴ Additional simulation details are

given in the Supporting Information. We have also summarized these simulations methods in a recent review.⁷⁵

Partial ionizations of residues were calculated using a linear interpolation scheme, which is consistent with the method to extract ionization values from ^{13}C NMR experiments.^{40,59} To justify the scheme, we calculated ^{13}C NMR chemical shifts of Tyr16’s C_ζ with the proton H16 moving along the O16–O57 vector. As shown in Fig. S2, the calculated ^{13}C chemical shifts have a strong linear dependence on the O–H bond length. For each hydrogen bonded pair O–H–O’, we define the atom O to be completely neutral if the distance between O and H, $d_{OH} \leq 0.96 \text{ \AA}$ and it to be completely ionized if $d_{O'H} \leq 0.96 \text{ \AA}$. The amount of partial ionization is obtained by linearly interpolating the two limits. We note that this definition is different from the one we exploited previously, where we used the proton transfer coordinate $\nu = d_{OX,HX} - d_{O57,HX}$ (X=16 or 32) and defined a tyrosine to be ionized when $\nu \geq 0$ for H16 or H32.⁴³

The excess isotope effect, $\Delta\Delta pK_a$ (the difference between the active site tyrosines H/D pKa isotope effect and that of tyrosine in water), was calculated from the free energy differences using the thermodynamic cycles shown in Fig. S3. $\Delta\Delta pK_a$ was further decomposed into three orthogonal coordinates, one along the O–H bond direction, one in the plane of the O–H–O’ hydrogen bond but perpendicular to the O–H bond vector, and one perpendicular to the O–H–O’ hydrogen bond plane. The decomposition was calculated by projecting the free energies along the three directions.

The active-site electric fields in KSI^{D40N}, E_{CO} , were calculated as the electric fields exerted by the active-site atoms on the C–O chemical bond of the phenol molecule, projected along the C–O bond unit vector. We extracted 4200 and 6142 complex configurations from AIMD and AI-PIMD simulations, respectively, of KSI^{D40N} with phenol bound. Each complex configuration contains Tyr16, Tyr32, Tyr57, Asp103 and phenol, and we evaluated the total electric fields on the C and O atoms in the phenol molecule, which were then projected along the C–O bond direction. The total electric field on the C–O bond in the enzyme

complex, $E_{complex}$, was calculated by averaging over the fields on the C and O atoms. To remove the contribution from the phenol molecule itself, E_{phenol} (i.e., the self-field), we repeated the calculation on the deprotonated phenol molecule as isolated from each complex configuration. E_{CO} was obtained by $E_{complex} - E_{phenol}$. The calculations were performed using the Gaussian 09 program⁷⁶ using the B3LYP functional⁷⁰ and 6-31G* basis set. The expectation values of the electrostatic properties were calculated using the PRISM integral algorithm.⁷⁷ We did not include the protein environment in these calculations. To justify this, we extracted 422 configurations from AI-PIMD simulations, and calculated E_{CO} with and without a point charge environment from the rest of the protein, water and counter-ions (partial charges according to the AMBER03 force field⁷³ and TIP3P water model⁷⁴). The resulting difference in electric fields was less than 2%.

Results and discussion

By incorporating unnatural amino acids to perturb the active-site hydrogen bond network in KSI^{D40N}, recent experiments have shown that this network is surprisingly robust in that the ionization states of the residues and the pK_a isotope effects are almost unchanged upon making these perturbations.⁴⁰ After binding an inhibitor, Stark effect spectroscopy measurements were used to show that the active site of the enzyme exerts a large electric field on a bond of the inhibitor, which undergoes a charge redistribution during the catalytic cycle.^{46,49} Here, we discuss these properties in turn and elucidate how the delocalization of the protons in the hydrogen bond network gives rise to these observations.

Change in partial ionizations in response to mutations

We begin by considering the effect of chlorination on the hydrogen-bonded triad of tyrosine residues in the active site of KSI^{D40N}. In the absence of mutations (Fig. 1a), the triad motif formed in KSI^{D40N} allows the protons H16 and H32 to be quantum mechanically delocalized

in the network.⁴³ This delocalization causes the side-chain phenol groups of the three tyrosine residues to be partially ionized. Equilibrium partial ionizations of residues Tyr16, Tyr32 and Tyr57 were obtained from ^{13}C NMR experiments assuming a linear change in the NMR chemical shift of the carbon adjacent to the ionizable oxygen ($^{13}\text{C}_\zeta$) between the ionized and unionized state, and yielded values of 20%, 13%, and 67%, respectively.^{40,43} Our electronic structure calculations exhibit a close-to-linear relationship between the calculated ^{13}C NMR chemical shift and the position of the proton when it is moved between the hydrogen bond donor and acceptor (Fig. S2), giving *post hoc* justification to this experimental assumption. Using this relation, our AI-PIMD simulations predict that residues Tyr16, Tyr32 and Tyr57 are 14%, 13% and 73% ionized, respectively, in very good agreement with experiment.

We now consider two mutants, which have been investigated in recent experiments,⁴⁰ where either Tyr16 or Tyr57 is replaced by 3-chlorotyrosine. The resulting proteins are referred to as Cl-Tyr16 and Cl-Tyr57 KSI^{D40N}, which are depicted in Fig. 2a and 2b, respectively. In aqueous solution, chlorination lowers the side-chain pK_a of tyrosine by 1.8 units,⁴⁰ which is equivalent to stabilizing the deprotonated state by 2.5 kcal/mol. If one were to assume a simple two-state model to describe the proton transfer between Tyr16 and Tyr57, i.e., H16 is attached to either one tyrosine or the other (Fig. 1a), then one would expect chlorination of Tyr16 ($\delta_{pK_a} = 1.8$) to increase the relative ionization of Tyr16 to Tyr57, denoted as f_{16}/f_{57} , by $e^{2.303\delta_{pK_a}} = 63$ fold compared to the unlabeled protein. In addition, a key assumption of LBHBs is that the proton affinities of the hydrogen bonding partners are closely matched. As chlorination breaks the matching of the solution pK_a 's of the residues in the triad, one might expect a large change in the properties of the triad.^{8,9,11,12}

In contrast to this prediction, our AI-PIMD simulations show that chlorinating the tyrosine residues only slightly affects the ionization states in the triad. As shown in Table 1, in Cl-Tyr16 KSI^{D40N} f_{16}/f_{57} is increased by 1.34-fold compared to the unlabeled protein, which is in excellent agreement with the experimentally measured value of 1.45-fold.⁴⁰ In contrast, chlorinating Tyr57 reduces proton sharing in the triad, which decreases f_{16}/f_{57} by

1.2 fold compared to the unlabeled case.

To assess how much of the change can be explained by simply considering the change in the proton transfer potential energy surface, Fig. S4 shows the energy of moving H16 between Tyr16 and Tyr57 for the unlabeled, Cl-Tyr16 and Cl-Tyr57 KSI^{D40N} in the crystal structure.²⁷ Chlorinating Tyr16 lowers the potential energy difference between the two minima by ~ 2 kcal/mol relative to the unlabeled protein, while chlorinating Tyr57 increases it by ~ 1.5 kcal/mol. This is smaller than the 2.5 kcal/mol change predicted by using the pK_a change in solution, but still would be expected to lead to significant changes in the ionization that are not observed in the simulation. In both mutants the barrier to proton delocalization remains significantly higher than the thermal energy at 300 K (~ 0.6 kcal/mol). However, as the ZPE of the O–H bond is comparable to the barrier, it can flood the potential energy well and allow quantum delocalization of the proton between the residues.

Table 1: Partial ionizations of residues Tyr16, Tyr32 and Tyr57 in the unlabeled KSI^{D40N} and their changes in Cl-Tyr16 and Cl-Tyr57 KSI^{D40N} compared to the unlabeled protein ($\Delta(\text{Cl-Tyr16})$ and $\Delta(\text{Cl-Tyr57})$ respectively). Also included in brackets are the values obtained from experiments.⁴⁰

	Unlabeled	$\Delta(\text{Cl-Tyr16})$	$\Delta(\text{Cl-Tyr57})$
Tyr16	14% (20%)	+4% (+7%)	-1%
Tyr32	13% (13%)	-1% (-2%)	-5%
Tyr57	73% (67%)	-3% (-5%)	+6%

To provide insight into the origins of the small changes to the partial ionization of Tyr16 (I_{tot}) upon replacement of tyrosine residues by 3-chlorotyrosines, we decompose I_{tot} as

$$I_{tot} = \int P(R)I(R)dR, \quad (1)$$

where $P(R)$ represents the probability distribution of the donor-acceptor distance, R , and $I(R)$ is the partial ionization of Tyr16 at a given R . This allows us to separate the effects on partial ionization arising directly from the change in the distance between the hydrogen donor and acceptor, R , upon perturbation from those arising from the change in the ion-

ization at a given distance. As shown in Fig. 3a, chlorinating Tyr16 shortens the average distance between Tyr16 and Tyr57, $R_{16/57}$, by about 0.03 Å compared to the unlabeled protein. Chlorinating Tyr57 has almost no impact and both are consistent with the previous experimental observation that $R_{16/57}$ is unaffected by chlorination within the instrument resolution of ~ 0.1 Å.⁴⁰ Due to the minor change in $R_{16/57}$ in Cl-Tyr57 KSI^{D40N}, the ionization transferred from Tyr57 to Tyr16 is similarly small (1%) compared to the unlabeled protein.

Fig. 3b shows that the ionization depends strongly on the donor-acceptor distance ($\sim 4\%$ per 0.1 Å). Chlorination, which alters the relative proton affinities, is seen to make the largest change at the shortest $R_{16/57}$ values (making a change of $\sim 2\%$ at 2.35 Å) while making almost no difference at 2.8 Å. From this one can calculate that of the total ionization change of 4% upon chlorination of Tyr16, 1.5% comes from the change in $P(R_{16/57})$ and 2.5% from the change in the slope of $I(R_{16/57})$. Therefore, both the short R and the matching of the proton affinity of the residues in the active-site hydrogen bond network of KSI^{D40N} are almost equally crucial factors in determining the residues' change in partial ionization in the presence of perturbations. The invariance of I_{tot} comes from the fact that chlorination leads to a small change (less than 20% at any R) in both $P(R_{16/57})$ and $I(R_{16/57})$.

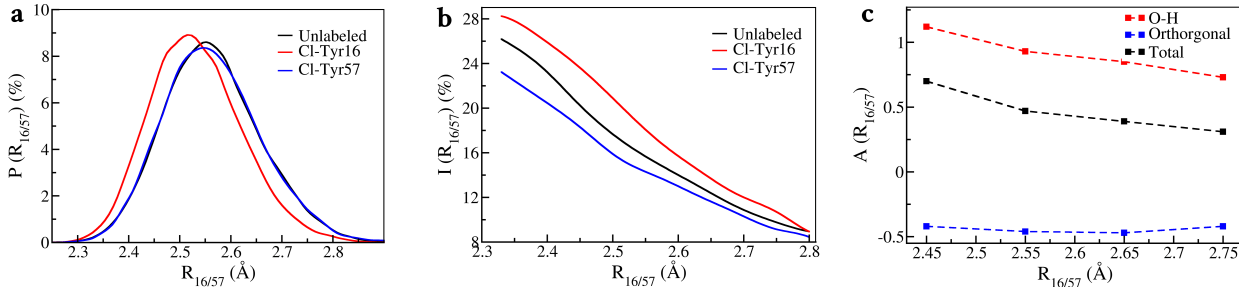


Figure 3: (a) The distribution of $R_{16/57}$ and (b) the ionization propensity $I(R_{16/57})$ for unlabeled, Cl-Tyr16 and Cl-Tyr57 KSI^{D40N}. (c) The pK_a isotope effect $A(R_{16/57})$ of the unlabeled protein.

Change in pK_a isotope effect in response to mutations

The excess isotope effect, $\Delta\Delta pK_a$, defined as the difference between the H/D pK_a isotope effect in KSI^{D40N} and that in tyrosine solution, provides another experimental measure of

the change in the network upon chlorination. We have previously shown that the KSI^{D40N} hydrogen bond network results in a large excess isotope effect for the triad ($\Delta\Delta pK_a = 0.50 \pm 0.03$) which arises from the close match between the zero-point energy in the O–H bond and the energetic barrier needed to move a proton across a hydrogen bond from a donor to the acceptor.⁵⁹ If the nuclei were to behave classically, the pK_a would have no dependence on the mass of the protons. Hence $\Delta\Delta pK_a$ arises purely from the quantum mechanical nature of the nuclei, which can be treated using path integral simulations.

In Cl-Tyr16 KSI^{D40N}, one might expect residue 16 to be more likely to donate its proton due to its lower pK_a , which would decrease the proton transfer barrier and enhance the quantum delocalization of protons in the biological hydrogen bond network. This would lead to an increase in $\Delta\Delta pK_a$ as compared to the unlabeled protein. Similarly, one might expect less proton sharing and a much smaller $\Delta\Delta pK_a$ in Cl-Tyr57 KSI^{D40N}, since the deprotonated residue 57 can be stabilized by the electron-withdrawing chlorine rather than by hydrogen bonding. However, as shown in Fig. 4a, our AI-PIMD simulations predict that the $\Delta\Delta pK_a$ values are very similar for the unlabeled, Cl-Tyr16 and Cl-Tyr57 KSI^{D40N}. Compared to the unlabeled protein, chlorinating Tyr16 has no impact on $\Delta\Delta pK_a$ within the error bar (± 0.03), while chlorinating Tyr57 only slightly decreases the isotope effect.

To uncover the origins of the small changes in $\Delta\Delta pK_a$, we again evaluate the impact of the donor acceptor distance R and matching of proton affinity in the triad by decomposing $\Delta\Delta pK_a$ as

$$\Delta\Delta pK_a = \int P(R)A(R)dR, \quad (2)$$

where $A(R)$ is the $\Delta\Delta pK_a$ value obtained at donor-acceptor distance R . KSI^{D40N} and its two chlorine-labeled mutants share the same $A(R_{16/57})$ profile within the error bar (± 0.05), which suggests that it is the O–O distances, rather than the difference in the residues’ proton affinity, that dominate $\Delta\Delta pK_a$ in the hydrogen bonded triad. $A(R_{16/57})$ for the unlabeled protein is shown in Fig. 3c. Thus, unlike the partial ionization, the small change in $\Delta\Delta pK_a$ is nearly entirely due to the largely unchanged distance distribution $P(R_{16/57})$

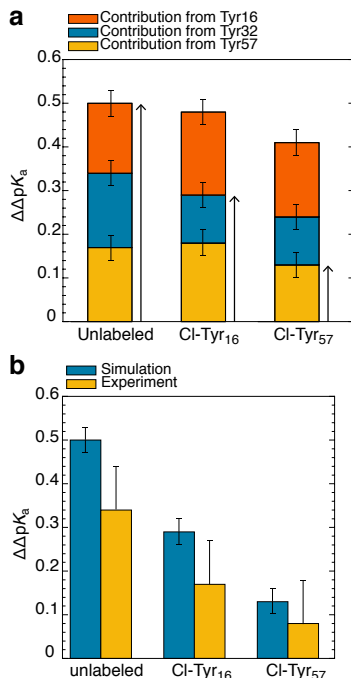


Figure 4: $\Delta\Delta pK_a$ for unlabeled, Cl-Tyr₁₆, and Cl-Tyr₅₇ KSI^{D40N}. (a) Contributions of Tyr16, Tyr32 and Tyr57 to the overall $\Delta\Delta pK_a$. The arrows show the residues measured in experiment. (b) Comparison of $\Delta\Delta pK_a$ obtained from AI-PIMD simulations and UV-Vis experiments.⁴⁰ The theoretical and experimental error bars are 0.03 and 0.10, respectively.

upon chlorination.

As shown in Fig. 3c, the isotope effect $A(R_{16/57})$ drops by 55% when $R_{16/57}$ moves from 2.45 Å to 2.75 Å. To understand this trend, we decompose $A(R_{16/57})$ into three orthogonal internal directions – one along the O–H covalent bond (O–H), one in the plane of the O–H–O hydrogen bond and one perpendicular to the plane of the hydrogen bond. The O–H direction gives a large positive contribution to $\Delta\Delta pK_a$ since the larger ZPE of the O–H stretch motion allows the H atom to be more quantum mechanically delocalized compared to D. This strengthens the hydrogen bond in KSI^{D40N} compared to tyrosine in solution. To illustrate this point, as shown in Fig. 3c, if only the O–H contribution were present, $\Delta\Delta pK_a$ would be 0.85 for $R = 2.65$ Å. However, ZPE in the two directions orthogonal to the O–H direction distorts the hydrogen bond and hence has the opposite effect. These orthogonal contributions reduce $\Delta\Delta pK_a$ by 0.46, resulting in the observed value of 0.39 at that distance. Fig. 3c shows that for every $R_{16/57}$ sampled in AI-PIMD simulations, a large

amount of these two effects cancel. This observation demonstrates the established principle of competing quantum effects in hydrogen bonds.^{19,56,57,59,64,78–80} The amount of cancellation increases from 55% to 73% as $R_{16/57}$ moves from 2.45 Å to 2.75 Å, giving rise to the large drop in $A(R_{16/57})$. As shown in Fig. 3c, this increase in the cancellation all arise from the fall in the O–H contribution, with the orthogonal contribution staying roughly constant as R increases. This reflects the fact that small changes in R can dramatically modulate the barrier to delocalize the proton along the O–H direction. The observed $A(R_{16/57})$ is positive for all relevant $R_{16/57}$ and thus contributions along the O–H direction dominate. This is in contrast to the situation observed in liquid water where $R \approx 2.8$ Å, leading to the two competing quantum effects almost perfectly cancelling each other.^{56,57} In KSI^{D40N}, NQEs act to strength the active-site hydrogen bond network because the overall protein fold keeps the tyrosine residues in such close proximity that the strengthening effect arising from the O–H contribution dominates.

These results allow us to suggest an explanation for two seemingly contradictory observations made in recent experiments: the unlabeled and Cl-Tyr16 KSI^{D40N} were observed to have a large difference in $\Delta\Delta pK_a$, in contradiction to the small changes observed in their partial ionization values.⁴⁰ Our results suggest that the UV-Vis experiments used to extract $\Delta\Delta pK_a$ were complicated by the overlapping spectral signatures of tyrosine and chlorotyrosine. Specifically, in the unlabeled protein, Tyr16, Tyr32 and Tyr57 contribute to a single absorption band in the UV-Vis spectrum and thus the experimentally measured $\Delta\Delta pK_a$ probes the triad as a whole, not the individual residues. However, when a chlorine-substituted tyrosine is incorporated, the band is split into two absorption peaks arising from tyrosinate and 3-chlorotyrosinate, which have significant spectral overlap.⁴⁰ As a result, in UV-Vis measurements of ionization of Cl-Tyr16 KSI^{D40N}, only the contributions of ionized Tyr32 and Tyr57 are detected, while in UV-Vis measurements of Cl-Tyr57 KSI^{D40N}, only the contribution of ionized Cl-Tyr57 is detected, as shown with arrows in Fig. 4a. In Fig. 4b we show that when we isolate these contributions from our simulations, the resulting

$\Delta\Delta pK_a$ are in agreement with experiments within the error bars.

Electric fields in the extended hydrogen bond network

Upon docking an intermediate analog, phenol, an extended hydrogen bond network is formed which incorporates phenol and residue Asp103 (Fig. 1b). As shown in Fig. 5b, the average distances between the five oxygen atoms in the network from our AI-PIMD simulations are all below 2.7 Å. In particular, the network is composed of three triads where the central triad is composed of Tyr57, Tyr16 and phenol. It is interesting to note that upon formation of the extended network, the distance between Tyr57 and Tyr16 ($R_{57/16}$) is reduced by 0.06 Å compared to that in the tyrosine triad, while Tyr32, which is on the outside of the network, increases its distance from Tyr57 ($R_{32/57}$) by 0.13 Å. This suggests that forming the hydrogen bond between Tyr16 and phenol enhances the Tyr16–Tyr57 interaction at the expense of the Tyr57–Tyr32 one. The central triad in the extended hydrogen bond network composed of Tyr57, Tyr16 and phenol thus has shorter O–O distances compared to the tyrosine triad in the apo protein, which allow the protons, in particular H57 and HP, to sample a wide range of distances along the ν coordinates (~ 1.8 Å), as shown in Fig. 5c.

Due to the delocalization of the protons, the hydrogen bonding partners in the extended network are partially ionized. As Tyr57, Tyr16 and phenol are all at the center of their corresponding triad structures, they are 22%, 37% and 21% ionized, respectively, compared with the experimental values of $\sim 39\%$, $\sim 38\%$ and $\sim 23\%$.²⁹ Residues Tyr57 and Tyr16 thus share the ionization of phenol, stabilizing the negatively charged intermediate analog.⁵⁹

The ability of KSI to stabilize the charged dienolate intermediate more than its uncharged ground state has been suggested to be connected to the electric fields in the enzyme’s active site.⁴⁶ Recent experiments used Stark effect spectroscopy to show that wild-type KSI exerts a large electric field of -144 ± 6 MV/cm on the C=O bond of a bound inhibitor 19-nortestosterone (19-NT), while KSI^{D40N} generates a field of -135 ± 4 MV/cm.⁴⁶ In addition, KSI mutants with the active-site tyrosine residues replaced by 3-chlorotyrosine were observed

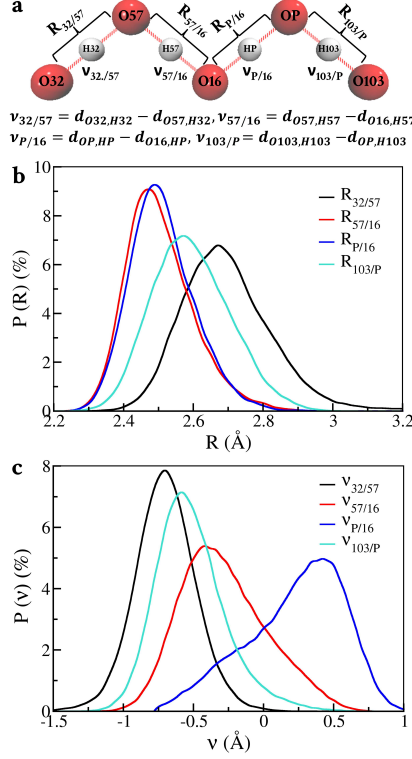


Figure 5: Schematic depiction of the extended hydrogen bond network (a) and the distributions of the O–O distances R (b) and the proton transfer coordinate ν (c).

to retain the large electric field (between -119 and -131 MV/cm depending on the mutation site),⁴⁷ which is consistent with the relatively unperturbed structural arrangement of the active-site residues and the modest change in tyrosine’s O–H dipole following chlorination. It further suggests that the large field an inhibitor experiences in this environment arises from the geometric arrangement of the active site. Our QM/MM simulations give a total electric field experienced by the phenol’s C–O bond, E_{CO} , of -152 ± 4 MV/cm when the nuclei are treated quantum mechanically. This is within the error bars of that obtained from our simulations with classical nuclei (-159 ± 6 MV/cm). These fields are in much closer agreement with experiment (-135 MV/cm) than previous simulations (-57.6 MV/cm and -60.4 MV/cm).^{46,50} In both our AI-PIMD simulations and the previous simulations using polarizable force fields,⁵⁰ the contribution from the protein scaffold is a small part of the total electric field (-2 MV/cm and -16 MV/cm respectively), but in our case the contribution from the active site is 3.7 fold larger, bringing it into closer agreement with the observed

electric field.

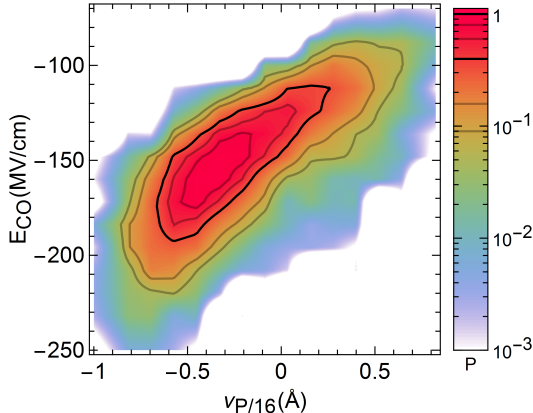


Figure 6: Joint probability of $\nu_{P/16}$ and E_{CO} in KSI^{40N} with a phenol bound, as obtained from AI-PIMD simulations and normalized by the maximum probability.

Where does the large electric field come from? By comparing E_{CO} in the presence and absence of the protein environment, we find that the fields exerted by residues in the active-site hydrogen bond network (Fig. 1b) account for 98% of the overall electric field. Considering the large degree of proton sharing between Tyr16 and phenol in this network, we plot E_{CO} as a function of the position of H16, as captured by $\nu_{P/16} = d_{OP,H16} - d_{O16,H16}$. As shown in Fig. 6, E_{CO} fluctuates as the proton moves along the hydrogen bond, which are accompanied by electronic rearrangements. E_{CO} increases dramatically in magnitude when the proton moves towards the intermediate analog (i.e, when $\nu_{P/16}$ becomes more negative). The large average E_{CO} (-152 MV/cm) thus results from a short average $\nu_{P/16}$ of 0.1 Å in the AI-PIMD simulations (Fig. 5c), which is close to the value of zero expected for a proton which sits on average exactly equidistant from both residues. Our calculated E_{CO} is slightly larger than the experimentally measured fields.⁴⁶ One likely reason for this is that phenol in the simulations is an intermediate analog, which should allow for greater proton sharing with Tyr16 than 19-NT in the experiment, which is a substrate analog that possesses a C=O group (in place of C–O). Therefore, the calculated average E_{CO} from the simulations can be viewed as an upper limit to the electric field that the C=O in 19-NT would experience. Importantly, this calculation is consistent with the suggestion that KSI’s substrate undergoes

small geometric changes during the catalytic cycle, and hence the electric field experienced by the C=O bond in the ground state is similar to that experienced by the C–O bond in the transition state/intermediate state.⁴⁹

Since the active-site hydrogen bond network gives rise to most of E_{CO} , we can now further decompose the electric field into contributions from different residues. We performed calculations on KSI^{D40N} in the absence and presence of individual active-site residues, and found that Tyr16 and Asp103 dominate E_{CO} as they form direct hydrogen bonds to the intermediate analog. As shown in Fig. 7, the values of E_{CO} from Tyr16 and Asp103 are -97 ± 2 MV/cm and -47 ± 1 MV/cm, respectively. These decomposition results are in good agreement with the experimental values of -84 ± 7 MV/cm and -52 ± 7 MV/cm, which were obtained by making the Tyr16Phe and Asp103Asn mutations.⁴⁶¹

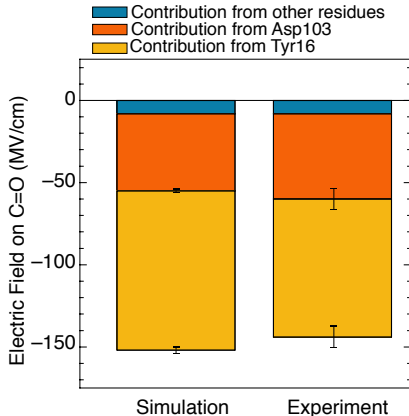


Figure 7: Decomposition of the average E_{CO} on the C–O bond of phenol (from AI-PIMD simulations) and on the C=O bond of 19-NT (from experiments⁴⁶).

However, it is essential to note that, while the protein environment is not a major contributor to the net electric fields, it plays a crucial role of maintaining the short O–O distances and orientations in the extended hydrogen bond network, which would otherwise not be stable. This positioning is essential: for example, in AI-PIMD simulations of KSI^{D40N} with the bound phenol, the average $R_{P/16}$ is 2.5 Å, whereas in previous classical molecular dy-

¹We note that these mutation experiments do not simply isolate the contribution from the removed residue but rather replace it with another amino acid, and thus the measured electric fields might also have contributions from the residues they are mutated to.

namics simulations, the steep repulsive walls in the force field push the O–O distances in the active-site hydrogen bond network to more than 4 Å, which lead to them to form weaker hydrogen bonds and hence to under-predict the E_{CO} value by a factor of ~ 2.7 .⁴⁶

Hence, the rigid protein environment which stabilizes the heavy atoms geometry permits the protons to flexibly move within the network (Fig. 5). The flexibility of the protons leads simultaneously to large but also quite heterogeneous electric fields (Fig. 6). Experiments have shown that the carbonyl stretch of 19-NT exhibits a very narrow IR linewidth;⁴⁶ however, these findings are consistent with the field fluctuations shown in Fig. 6 being associated primarily with movement of the protons, not the heavy atoms. Although AI-PIMD simulations do not provide dynamical information, we expect the proton movements occur on an ultrafast timescale, and hence would be averaged out on the IR timescale.

Conclusions

Here we have shown that electronic and nuclear quantum fluctuations in the active-site tyrosine triad in KSI^{D40N} give rise to small changes in the residues’ partial ionizations and $\Delta\Delta pK_a$ in response to perturbations such as unnatural amino acid mutagenesis, by shifting the proton positions and redistributing the electron density to buffer the perturbation effects. In the extended hydrogen bond network upon binding an intermediate analog, we showed that the quantum mechanical treatment of the electrons results in a large electric field in the enzyme’s active site, in excellent agreement with recent experiments. We have demonstrated how this flexibility and the large electric fields arise from the fact that the short O–O distances lower the potential energy barrier for transferring protons along the hydrogen bond network and that NQEs further flood the potential energy well, thus allowing the electrons and protons to move over a wide range of distances. This work further highlights how AI-PIMD simulations of biological systems can be powerful tools that can reproduce and explain experimental observables such as isotope effects and electric fields.^{23,43,61,75,80,81}

Finally, it is worth commenting on how the flexibility of the hydrogen bond network might impact the enzyme’s function. For example, our observation of the small change in the ionization of the residues upon chlorination is consistent with the recent experimental observation that replacement of Tyr16 with a chloro- or fluoro-tyrosine leads to very little change to KSI’s catalytic activity.⁴⁵ In addition, the ability to form an extended hydrogen bond network with enhanced proton flexibility, rather than a single hydrogen bond, may also be of importance. In particular, a single Tyr57Phe or Asp103Leu mutation, both of which retain an extended network where the substrate can dock, lead to modest 6 or 100 fold decrease in KSI’s catalytic activity, respectively.^{82,83} However, the combined Tyr57Phe/Asp103Leu mutation, which leaves Tyr16 as the only hydrogen bond donor to the substrate with no extended network and might be expected to yield a 600 fold decrease if the effects of these mutations were independent, reduces the activity by over 15,800 fold.⁸² This suggests that Tyr57 can partially compensate for the impact of mutating Asp103 since the substrate can still be incorporated in an extended hydrogen bond network, which possesses electronic and nuclear flexibility. This robustness of the network thus might protect the enzyme from losing its function unless multiple mutations occur concurrently.

Supporting Information

Simulation details, methods for calculating partial ionization, $\Delta\Delta pK_a$ and the potential energy surface of proton transfer in the tyrosine triad systems

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TOC Graphic

