

Significant Excess of Electronlike Events in the MiniBooNE Short-Baseline Neutrino Experiment

A. A. Aguilar-Arevalo,¹³ B. C. Brown,⁶ L. Bugel,¹² G. Cheng,⁵ J. M. Conrad,¹² R. L. Cooper,^{10,15} R. Dharmapalan,^{1,2} A. Diaz,¹² Z. Djurcic,² D. A. Finley,⁶ R. Ford,⁶ F. G. Garcia,⁶ G. T. Garvey,¹⁰ J. Grange,⁷ E.-C. Huang,¹⁰ W. Huelsnitz,¹⁰ C. Ignarra,¹² R. A. Johnson,³ G. Karagiorgi,⁵ T. Katori,^{12,16} T. Kobilarcik,⁶ W. C. Louis,¹⁰ C. Mariani,¹⁹ W. Marsh,⁶ G. B. Mills,^{10,*} J. Mirabal,¹⁰ J. Monroe,¹⁸ C. D. Moore,⁶ J. Mousseau,¹⁴ P. Nienaber,¹⁷ J. Nowak,⁹ B. Osmanov,⁷ Z. Pavlovic,⁶ D. Perevalov,⁶ H. Ray,⁷ B. P. Roe,¹⁴ A. D. Russell,⁶ M. H. Shaevitz,⁵ J. Spitz,¹⁴ I. Stancu,¹ R. Tayloe,⁸ R. T. Thornton,¹⁰ M. Tzanov,^{4,11} R. G. Van de Water,¹⁰ D. H. White,^{10,*} D. A. Wickremasinghe,³ and E. D. Zimmerman⁴

(MiniBooNE Collaboration)

¹University of Alabama, Tuscaloosa, Alabama 35487, USA

²Argonne National Laboratory, Argonne, Illinois 60439, USA

³University of Cincinnati, Cincinnati, Ohio 45221, USA

⁴University of Colorado, Boulder, Colorado 80309, USA

⁵Columbia University, New York, New York 10027, USA

⁶Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA

⁷University of Florida, Gainesville, Florida 32611, USA

⁸Indiana University, Bloomington, Indiana 47405, USA

⁹Lancaster University, Lancaster LA1 4YB, United Kingdom

¹⁰Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA

¹¹Louisiana State University, Baton Rouge, Louisiana 70803, USA

¹²Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

¹³Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, CDMX 04510, Mexico

¹⁴University of Michigan, Ann Arbor, Michigan 48109, USA

¹⁵New Mexico State University, Las Cruces, New Mexico 88003, USA

¹⁶Queen Mary University of London, London E1 4NS, United Kingdom

¹⁷Saint Mary's University of Minnesota, Winona, Minnesota 55987, USA

¹⁸Royal Holloway, University of London, Egham TW20 0EX, United Kingdom

¹⁹Center for Neutrino Physics, Virginia Tech, Blacksburg, Virginia 24061, USA

 (Received 30 May 2018; revised manuscript received 28 September 2018; published 26 November 2018)

The MiniBooNE experiment at Fermilab reports results from an analysis of ν_e appearance data from 12.84×10^{20} protons on target in neutrino mode, an increase of approximately a factor of 2 over previously reported results. A ν_e charged-current quasielastic event excess of 381.2 ± 85.2 events (4.5σ) is observed in the energy range $200 < E_\nu^{\text{QE}} < 1250$ MeV. Combining these data with the $\bar{\nu}_e$ appearance data from 11.27×10^{20} protons on target in antineutrino mode, a total ν_e plus $\bar{\nu}_e$ charged-current quasielastic event excess of 460.5 ± 99.0 events (4.7σ) is observed. If interpreted in a two-neutrino oscillation model, $\nu_\mu \rightarrow \nu_e$, the best oscillation fit to the excess has a probability of 21.1%, while the background-only fit has a χ^2 probability of 6×10^{-7} relative to the best fit. The MiniBooNE data are consistent in energy and magnitude with the excess of events reported by the Liquid Scintillator Neutrino Detector (LSND), and the significance of the combined LSND and MiniBooNE excesses is 6.0σ . A two-neutrino oscillation interpretation of the data would require at least four neutrino types and indicate physics beyond the three neutrino paradigm. Although the data are fit with a two-neutrino oscillation model, other models may provide better fits to the data.

DOI: [10.1103/PhysRevLett.121.221801](https://doi.org/10.1103/PhysRevLett.121.221801)

Evidence for short-baseline neutrino anomalies at an $L/E_\nu \sim 1$ m/MeV, where E_ν is the neutrino energy and L is the distance that the neutrino traveled before detection, comes from both neutrino appearance and disappearance experiments. The appearance anomalies include the excess of ν_e and $\bar{\nu}_e$ charge-current quasielastic (CCQE) events observed by the Liquid Scintillator Neutrino Detector (LSND) [1] and MiniBooNE [2,3] experiments, while the disappearance anomalies, although not completely consistent, include the deficit of ν_e and $\bar{\nu}_e$ events observed by reactor [4] and radioactive-source experiments [5]. As the masses and mixings within the three-generation neutrino matrix have been attached to solar and long-baseline neutrino experiments, more exotic models are typically used to explain these anomalies, including, for example, $3 + N$ neutrino oscillation models involving three active neutrinos and N additional sterile neutrinos [6–14], resonant neutrino oscillations [15], Lorentz violation [16], sterile neutrino decay [17], sterile neutrino nonstandard interactions [18], and sterile neutrino extra dimensions [19]. This Letter presents improved MiniBooNE ν_e and $\bar{\nu}_e$ appearance results, assuming two-neutrino oscillations with probability $P = \sin^2(2\theta) \sin^2(1.27\Delta m^2 L/E)$, where θ is the mixing angle, Δm^2 (eV²/c⁴) is the difference in neutrino mass eigenstates squared, L (m) is the distance traveled by the neutrino, and E (MeV) is the neutrino energy.

The booster neutrino beam (BNB) at Fermilab delivers to the MiniBooNE experiment a flux of neutrinos and antineutrinos that is simulated using information from external measurements [20]. The BNB is produced by 8 GeV protons from the Fermilab booster interacting on a beryllium target inside a magnetic focusing horn. Depending on the polarity of the horn, either π^+ are focused and π^- are defocused to produce a fairly pure beam of ν_μ , or π^- are focused and π^+ are defocused to produce a somewhat pure beam of $\bar{\nu}_\mu$. In neutrino mode, the ν_μ , $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ flux contributions at the detector are 93.5%, 5.9%, 0.5%, and 0.1%, respectively, while in antineutrino mode, the flux contributions are 15.7%, 83.7%, 0.2%, and 0.4%, respectively. The ν_μ and $\bar{\nu}_\mu$ fluxes peak at approximately 600 and 400 MeV, respectively.

The MiniBooNE detector is described in detail in Ref. [21]. The detector consists of a 12.2 m diameter sphere filled with 818 tonnes of pure mineral oil (CH₂) and is located 541 m from the beryllium target. The detector is covered by 1520 8-inch photomultiplier tubes (PMTs), where 1280 PMTs are in the interior detector region and 240 PMTs are located in the optically isolated outer veto region. Charged particles produced by neutrino interactions in the mineral oil emit both directed Cherenkov light and isotropic scintillation light that is detected by the PMTs. Event reconstruction [22] and particle identification make use of the hit PMT charge and time information, and the reconstructed neutrino energy E_ν^{QE} is estimated from the

measured energy and angle of the outgoing muon or electron, assuming the kinematics of CCQE scattering [23].

From 2002–2017, the MiniBooNE experiment has collected a total of 11.27×10^{20} protons on target (POT) in antineutrino mode, 12.84×10^{20} POT in neutrino mode, and a further 1.86×10^{20} POT in a special beam-off target mode to search for sub-GeV dark matter [24]. The neutrino sample has approximately doubled in size since the previous publication [3]. The published neutrino-mode data correspond to 6.46×10^{20} POT, while 6.38×10^{20} POT were obtained in 2016 and 2017. During the 15 years of running, the BNB and MiniBooNE detector have been stable to within 2% in neutrino energy.

The analysis is optimized to measure ν_e and $\bar{\nu}_e$ induced CCQE events, and the event reconstruction [22] and selection are identical to the previous analysis [3]. The average selection efficiency is $\sim 20\%$ ($\sim 0.1\%$) for ν_e -induced CCQE events (ν_μ -induced background events) generated over the fiducial volume. The fraction of CCQE events in antineutrino mode that are from wrong-sign neutrino events was determined from the angular distributions of muons created in CCQE interactions and by measuring CC single π^+ events [25].

The predicted but unconstrained ν_e and $\bar{\nu}_e$ CCQE background events for the neutrino energy range $200 < E_\nu^{\text{QE}} < 1250$ MeV are shown in Table I for both neutrino

TABLE I. The expected (unconstrained) number of events for the $200 < E_\nu^{\text{QE}} < 1250$ MeV neutrino energy range from all of the backgrounds in the ν_e and $\bar{\nu}_e$ appearance analysis before using the constraint from the CC ν_μ events. Also shown are the constrained background, as well as the expected number of events corresponding to the LSND best fit oscillation probability of 0.26%, assuming oscillations at large Δm^2 . The table shows the diagonal-element systematic plus statistical uncertainties, which become substantially reduced in the oscillation fits when correlations between energy bins and between the electron and muon neutrino events are included. The antineutrino numbers are from a previous analysis [3].

Process	Neutrino mode	Antineutrino mode
ν_μ & $\bar{\nu}_\mu$ CCQE	73.7 ± 19.3	12.9 ± 4.3
NC π^0	501.5 ± 65.4	112.3 ± 11.5
NC $\Delta \rightarrow N\gamma$	172.5 ± 24.1	34.7 ± 5.4
External events	75.2 ± 10.9	15.3 ± 2.8
Other ν_μ & $\bar{\nu}_\mu$	89.6 ± 22.9	22.3 ± 3.5
ν_e & $\bar{\nu}_e$ from $\mu^\pm$ decay	425.3 ± 100.2	91.4 ± 27.6
ν_e & $\bar{\nu}_e$ from $K^\pm$ decay	192.2 ± 41.9	51.2 ± 11.0
ν_e & $\bar{\nu}_e$ from K_L^0 decay	54.5 ± 20.5	51.4 ± 18.0
Other ν_e & $\bar{\nu}_e$	6.0 ± 3.2	6.7 ± 6.0
Unconstrained bkgd.	1590.6 ± 176.9	398.2 ± 49.7
Constrained bkgd.	1577.8 ± 85.2	398.7 ± 28.6
Total data	1959	478
Excess	381.2 ± 85.2	79.3 ± 28.6
0.26% (LSND) $\nu_\mu \rightarrow \nu_e$	463.1	100.0

mode and antineutrino mode [26]. See Supplemental Material [27] for more information on backgrounds. The upper limit of 1250 MeV corresponded to a small value of L/E and was chosen by the collaboration before unblinding the data in 2007. The lower limit of 200 MeV is chosen because we constrain the ν_e events with the CCQE ν_μ events and our CCQE ν_μ event sample only goes down to 200 MeV, as we require a visible Cherenkov ring from the muon. The estimated sizes of the intrinsic ν_e and gamma backgrounds are based on MiniBooNE event measurements and uncertainties from these constraints are included in the analysis. The intrinsic $\nu_e/\bar{\nu}_e$ background from muon decay is directly related to the large sample of observed $\nu_\mu/\bar{\nu}_\mu$ events, as these events constrain the muons that decay in the 50 m decay region. This constraint uses a joint fit of the observed $\nu_\mu/\bar{\nu}_\mu$ and $\nu_e/\bar{\nu}_e$ events, assuming that there are no substantial $\nu_\mu/\bar{\nu}_\mu$ disappearance oscillations. The other intrinsic ν_e background component, from kaon decay, is constrained by fits to kaon production data and SciBooNE measurements [28]. The intrinsic ν_e background from pion decay (1.2×10^{-4} branching ratio) and hyperon decay are very small. Other backgrounds from misidentified ν_μ or $\bar{\nu}_\mu$ [29,30] events are also constrained by the observed CCQE sample.

The gamma background from neutral-current (NC) π^0 production and $\Delta \rightarrow N\gamma$ radiative decay [31,32] are constrained by the associated large two-gamma sample (mainly from Δ production) observed in the MiniBooNE data, where π^0 measurements [33] are used to constrain the π^0 background. The π^0 background measured in the first and second neutrino data sets were found to be consistent, resulting in a lower statistical background uncertainty for the combined data. Other neutrino-induced single gamma production processes are included in the theoretical predictions, which agree well with the MiniBooNE estimates [31,34]. Single-gamma backgrounds from external neutrino interactions (“dirt” backgrounds) are estimated using topological and spatial cuts to isolate the events whose vertices are near the edge of the detector and point towards the detector center [35]. With the larger data set, the background from external neutrino interactions is now better determined to be approximately 7% larger, but with smaller uncertainty than in the previous publication [3]. A new technique to measure or constrain the gamma and dirt backgrounds based on event timing relative to the beam is in development.

Systematic uncertainties are determined by considering the predicted effects on the ν_μ , $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ CCQE rates from variations of uncertainty parameters. The parameters include uncertainties in the neutrino and antineutrino flux estimates, uncertainties in neutrino cross sections, most of which are determined by *in situ* cross-section measurements at MiniBooNE [29,33], uncertainties from nuclear effects, and uncertainties in detector modeling and reconstruction. A covariance matrix in bins of E_ν^{QE} is

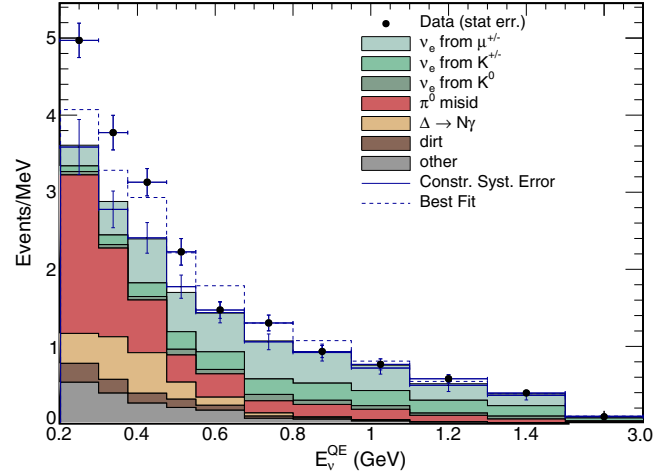


FIG. 1. The MiniBooNE neutrino mode E_ν^{QE} distributions, corresponding to the total 12.84×10^{20} POT data, for ν_e CCQE data (points with statistical errors) and background (histogram with systematic errors). The dashed curve shows the best fit to the neutrino-mode data assuming two-neutrino oscillations. The last bin is for the energy interval from 1500–3000 MeV.

constructed by considering the variation from each source of systematic uncertainty on the ν_e and $\bar{\nu}_e$ CCQE signal and background, and the ν_μ and $\bar{\nu}_\mu$ CCQE prediction as a function of E_ν^{QE} . This matrix includes correlations between any of the ν_e and $\bar{\nu}_e$ CCQE signal and background and ν_μ and $\bar{\nu}_\mu$ CCQE samples, and is used in the χ^2 calculation of the oscillation fits.

Table I also shows the expected number of events corresponding to the LSND best fit oscillation probability of 0.26%, assuming oscillations at large Δm^2 . LSND and MiniBooNE have the same average value of L/E , but MiniBooNE has a larger range of L/E . Therefore, the appearance probabilities for LSND and MiniBooNE should not be exactly the same at lower L/E values.

Figure 1 shows the E_ν^{QE} distribution for ν_e CCQE data and background in neutrino mode for the total 12.84×10^{20} POT data. Each bin of reconstructed E_ν^{QE} corresponds to a distribution of “true” generated neutrino energies, which can overlap adjacent bins. In neutrino mode, a total of 1959 data events pass the ν_e CCQE event selection requirements with $200 < E_\nu^{\text{QE}} < 1250$ MeV, compared to a background expectation of $1577.8 \pm 39.7(\text{stat}) \pm 75.4(\text{syst})$ events. The excess is then 381.2 ± 85.2 events or a 4.5σ effect. Note that the 162.0 event excess in the first 6.46×10^{20} POT data is approximately 1σ lower than the average excess, while the 219.2 event excess in the second 6.38×10^{20} POT data is approximately 1σ higher than the average excess. Figure 2 shows the excess events in neutrino mode from the first 6.46×10^{20} POT data and the second 6.38×10^{20} POT data (top plot). Combining the MiniBooNE neutrino and anti-neutrino data, there are a total of 2437 events in the $200 < E_\nu^{\text{QE}} < 1250$ MeV energy region, compared to a background

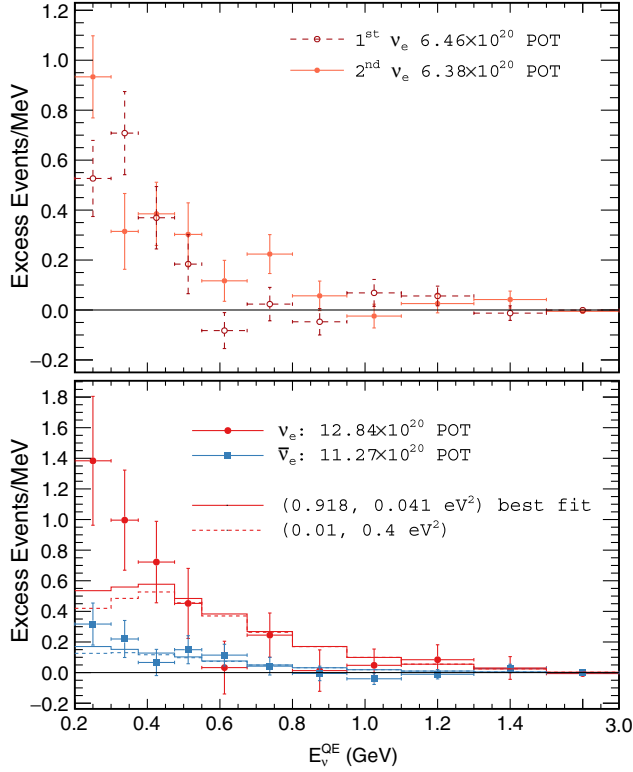


FIG. 2. The top plot shows the MiniBooNE event excesses in neutrino mode as a function of E_ν^{QE} from the first 6.46×10^{20} POT data and the second 6.38×10^{20} POT data. The bottom plot shows the total event excesses in both neutrino mode and antineutrino mode, corresponding to 12.84×10^{20} POT and 11.27×10^{20} POT, respectively. The solid (dashed) curve is the best fit (1σ fit point) to the neutrino-mode and antineutrino-mode data assuming two-neutrino oscillations. The last bin is for the energy interval from 1500–3000 MeV. Error bars include only statistical uncertainties for the top plot and both statistical and correlated systematic uncertainties for the bottom plot.

expectation of $1976.5 \pm 44.5(\text{stat}) \pm 88.5(\text{syst})$ events. This corresponds to a total ν_e plus $\bar{\nu}_e$ CCQE excess of 460.5 ± 99.0 events with respect to expectation or a 4.7σ excess. Figure 2 (bottom plot) shows the total event excesses as a function of E_ν^{QE} in both neutrino mode and antineutrino mode. The dashed curves show the two-neutrino oscillation predictions at the best-fit point ($\Delta m^2 = 0.041 \text{ eV}^2$, $\sin^2 2\theta = 0.92$), as well as at a point within 1σ of the best-fit point ($\Delta m^2 = 0.4 \text{ eV}^2$, $\sin^2 2\theta = 0.01$).

A two-neutrino model is assumed for the MiniBooNE oscillation fits in order to compare with the LSND data. However, the appearance neutrino experiments appear to be incompatible with the disappearance neutrino experiments in a $3 + 1$ model [10,12], and other models [15–19] may provide better fits to the data. The oscillation parameters are extracted from a combined fit of the observed E_ν^{QE} event distributions for muonlike and electronlike events using the full covariance matrix described previously in the full

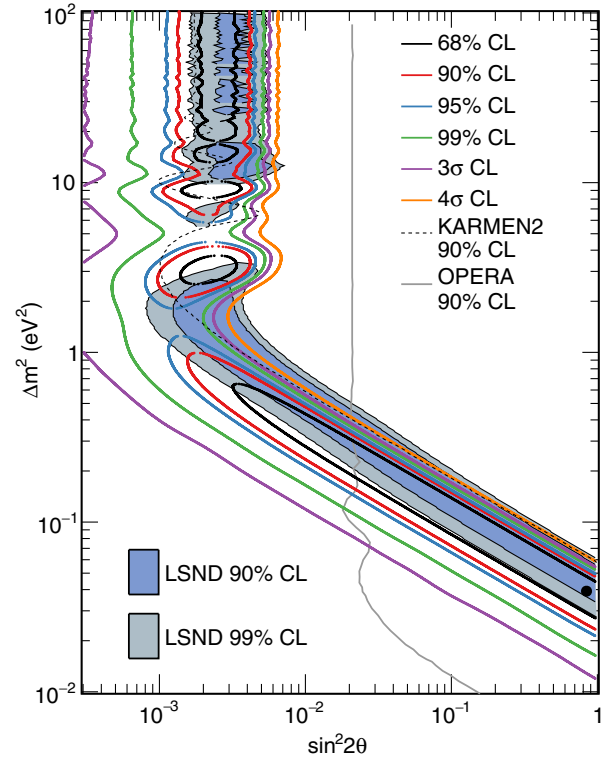


FIG. 3. MiniBooNE allowed regions in neutrino mode (12.84×10^{20} POT) for events with $200 < E_\nu^{\text{QE}} < 3000$ MeV within a two-neutrino oscillation model. The shaded areas show the 90% and 99% C.L. LSND $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ allowed regions. The black point shows the MiniBooNE best fit point. Also shown are 90% C.L. limits from the KARMEN [36] and OPERA [37] experiments.

energy range $200 < E_\nu^{\text{QE}} < 3000$ MeV. The fit assumes the same oscillation probability for both the right-sign ν_e and wrong-sign $\bar{\nu}_e$, and no ν_μ , $\bar{\nu}_\mu$, ν_e , or $\bar{\nu}_e$ disappearance. Using a likelihood-ratio technique [3], the confidence level values for the fitting statistic, $\Delta\chi^2 = \chi^2(\text{point}) - \chi^2(\text{best})$, as a function of oscillation parameters, Δm^2 and $\sin^2 2\theta$, is determined from frequentist, fake data studies. The fake data studies also determine the effective number of degrees of freedom and probabilities. With this technique, the best neutrino oscillation fit in neutrino mode occurs at $(\Delta m^2, \sin^2 2\theta) = (0.039 \text{ eV}^2, 0.84)$, as shown in Fig. 3. The χ^2/ndf for the best-fit point in the energy range $200 < E_\nu^{\text{QE}} < 1250$ MeV is $9.9/6.7$ with a probability of 15.5%. The background-only fit has a χ^2 probability of 0.06% relative to the best oscillation fit and a $\chi^2/\text{ndf} = 24.9/8.7$ with a probability of 0.21%. Figure 3 shows the MiniBooNE closed confidence level (C.L.) contours for ν_e appearance oscillations in neutrino mode in the $200 < E_\nu^{\text{QE}} < 3000$ MeV energy range.

Nuclear effects associated with neutrino interactions on carbon can affect the reconstruction of the neutrino energy, E_ν^{QE} , and the determination of the neutrino oscillation

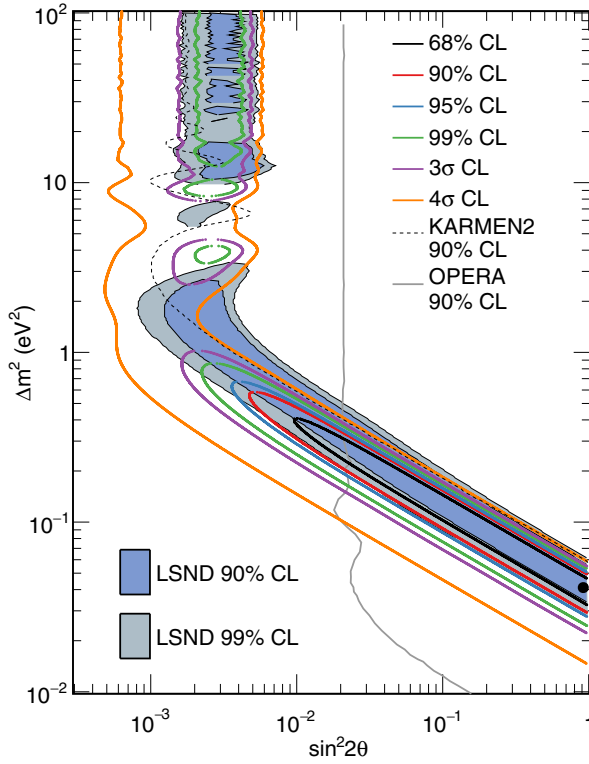


FIG. 4. MiniBooNE allowed regions for a combined neutrino mode (12.84×10^{20} POT) and antineutrino mode (11.27×10^{20} POT) data sets for events with $200 < E_\nu^{\text{QE}} < 3000$ MeV within a two-neutrino oscillation model. The shaded areas show the 90% and 99% C.L. LSND $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ allowed regions. The black point shows the MiniBooNE best fit point. Also shown are 90% C.L. limits from the KARMEN [36] and OPERA [37] experiments.

parameters [38]. These effects were studied previously [3,39] and were found to not affect substantially the oscillation fit. In addition, they do not affect the gamma background, which is determined from direct measurements of NC π^0 and dirt backgrounds.

Figure 4 shows the MiniBooNE allowed regions in both neutrino mode and antineutrino mode [3] for events with $200 < E_\nu^{\text{QE}} < 3000$ MeV within a two-neutrino oscillation model. For this oscillation fit the entire data set is used and includes the 12.84×10^{20} POT data in neutrino mode and the 11.27×10^{20} POT data in antineutrino mode. As shown in the figure, the MiniBooNE 1σ allowed region lies mostly within the LSND 90% C.L. band, which demonstrates good agreement between the LSND and MiniBooNE signals. Also shown are 90% C.L. limits from the KARMEN [36] and OPERA [37] experiments. The KARMEN2 90% C.L. limits are outside the MiniBooNE 95% C.L. allowed region, while the OPERA 90% C.L. limits disfavor the MiniBooNE allowed region below approximately 0.3 eV². The best combined neutrino oscillation fit occurs at $(\Delta m^2, \sin^2 2\theta) = (0.041 \text{ eV}^2, 0.92)$. The χ^2/ndf for the best-fit point in the energy range $200 < E_\nu^{\text{QE}} < 1250$ MeV

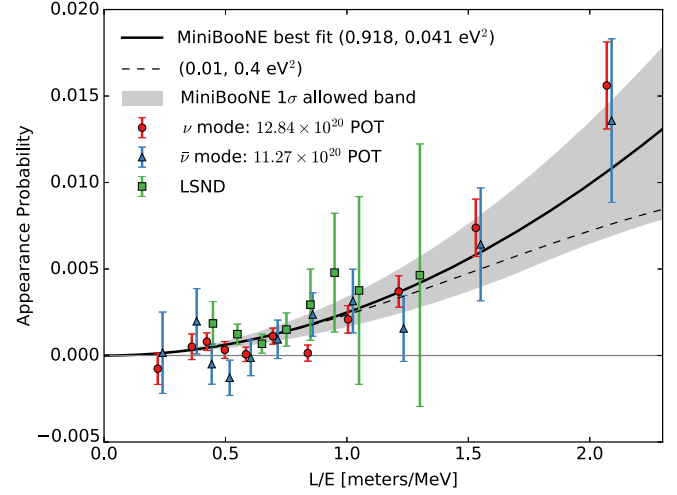


FIG. 5. A comparison between the L/E_ν^{QE} distributions for the MiniBooNE data excesses in neutrino mode (12.84×10^{20} POT) and antineutrino mode (11.27×10^{20} POT) to the L/E distribution from LSND [1]. The error bars show statistical uncertainties only. The curves show fits to the MiniBooNE data, assuming two-neutrino oscillations, while the shaded area is the MiniBooNE 1σ allowed band. The best-fit curve uses the reconstructed neutrino energy E_ν^{QE} for the MiniBooNE data. The dashed curve shows the example 1σ fit point.

is 19.4/15.6 with a probability of 21.1%, and the background-only fit has a χ^2 probability of 6×10^{-7} relative to the best oscillation fit and a $\chi^2/\text{ndf} = 47.1/17.3$ with a probability of 0.02%.

Figure 5 compares the L/E_ν^{QE} distributions for the MiniBooNE data excesses in neutrino mode and antineutrino mode to the L/E distribution from LSND [1]. The error bars show statistical uncertainties only. As shown in the figure, there is agreement among all three data sets. Assuming two-neutrino oscillations, the curves show fits to the MiniBooNE data described above. Fitting both MiniBooNE and LSND data, by adding LSND L/E data as additional terms, the best fit occurs at $(\Delta m^2, \sin^2 2\theta) = (0.041 \text{ eV}^2, 0.96)$ with a $\chi^2/\text{ndf} = 22.4/22.4$, corresponding to a probability of 42.5%. The MiniBooNE excess of events in both oscillation probability and L/E spectrum is, therefore, consistent with the LSND excess of events. The significance of the combined LSND (3.8σ) [1] and MiniBooNE (4.7σ) excesses is 6.0σ , which is obtained by adding the significances in quadrature, as the two experiments have completely different neutrino energies, neutrino fluxes, reconstructions, backgrounds, and systematic uncertainties.

In summary, the MiniBooNE experiment observes a total ν_e CCQE event excess in both neutrino and antineutrino running modes of 460.5 ± 99.0 events (4.7σ) in the energy range $200 < E_\nu^{\text{QE}} < 1250$ MeV. The MiniBooNE allowed region from a two-neutrino oscillation fit to the data, shown in Fig. 4, is consistent with the allowed region reported by

the LSND experiment [1]. On the other hand, a two-neutrino oscillation interpretation of the data would require at least four neutrino types and indicate physics beyond the three neutrino paradigm. The significance of the combined LSND and MiniBooNE excesses is 6.0σ . All of the major backgrounds are constrained by *in situ* event measurements, so nonoscillation explanations would need to invoke new anomalous background processes. Although the data are fit with a two-neutrino oscillation model, other models may provide better fits to the data. The MiniBooNE event excess will be further studied by the Fermilab short-baseline neutrino (SBN) program [40].

We acknowledge the support of Fermilab, the Department of Energy, and the National Science Foundation, and we acknowledge Los Alamos National Laboratory for LDRD funding.

*Deceased.

- [1] C. Athanassopoulos *et al.*, *Phys. Rev. Lett.* **75**, 2650 (1995); **77**, 3082 (1996); **81**, 1774 (1998); *Phys. Rev. C* **54**, 2685 (1996); **58**, 2489 (1998); A. Aguilar *et al.*, *Phys. Rev. D* **64**, 112007 (2001).
- [2] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. Lett.* **98**, 231801 (2007); **102**, 101802 (2009); **105**, 181801 (2010).
- [3] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. Lett.* **110**, 161801 (2013).
- [4] G. Mention, M. Fechner, T. Lasserre, T. A. Mueller, D. Lhuillier, M. Cribier, and A. Letourneau, *Phys. Rev. D* **83**, 073006 (2011).
- [5] C. Giunti and M. Laveder, *Phys. Rev. C* **83**, 065504 (2011).
- [6] M. Sorel, J. M. Conrad, and M. H. Shaevitz, *Phys. Rev. D* **70**, 073004 (2004).
- [7] G. Karagiorgi, Z. Djurcic, J. M. Conrad, M. H. Shaevitz, and M. Sorel, *Phys. Rev. D* **80**, 073001 (2009); **81**, 039902(E) (2010).
- [8] G. H. Collin, C. A. Argüelles, J. M. Conrad, and M. H. Shaevitz, *Phys. Rev. Lett.* **117**, 221801 (2016).
- [9] C. Giunti and M. Laveder, *Phys. Lett. B* **706**, 200 (2011); *Phys. Rev. D* **84**, 073008 (2011).
- [10] S. Gariazzo, C. Giunti, M. Laveder, and Y. F. Li, *J. High Energy Phys.* **06** (2017) 135.
- [11] J. Kopp, M. Maltoni, and T. Schwetz, *Phys. Rev. Lett.* **107**, 091801 (2011); J. Kopp and P. A. N. Machado, *J. High Energy Phys.* **05** (2013) 050.
- [12] M. Dentler, A. Hernandez-Cabezudo, J. Kopp, P. Machado, M. Maltoni, I. Martinez-Soler, and T. Schwetz, *J. High Energy Phys.* **08** (2018) 010.
- [13] K. N. Abazajian *et al.*, [arXiv:1204.5379](https://arxiv.org/abs/1204.5379).
- [14] J. M. Conrad, C. M. Ignarra, G. Karagiorgi, M. H. Shaevitz, and J. Spitz, *Adv. High Energy Phys.* **2013**, 163897 (2013).
- [15] J. Asaadi, E. Church, R. Guenette, B. J. P. Jones, and A. M. Szelc, *Phys. Rev. D* **97**, 075021 (2018); G. Karagiorgi, M. H. Shaevitz, and J. M. Conrad, [arXiv:1202.1024](https://arxiv.org/abs/1202.1024); H. Pas, S. Pakvasa, and T. J. Weiler, *Phys. Rev. D* **72**, 095017 (2005); D. Doring, H. Paes, P. Sicking, and T. J. Weiler, [arXiv:1808.07460](https://arxiv.org/abs/1808.07460).
- [16] V. A. Kostelecky and M. Mewes, *Phys. Rev. D* **69**, 016005 (2004); T. Katori, V. A. Kostelecky, and R. Tayloe, *Phys. Rev. D* **74**, 105009 (2006); J. S. Diaz and V. A. Kostelecky, *Phys. Lett. B* **700**, 25 (2011); *Phys. Rev. D* **85**, 016013 (2012).
- [17] S. N. Gninenko, *Phys. Rev. Lett.* **103**, 241802 (2009); S. N. Gninenko and D. S. Gorbunov, *Phys. Rev. D* **81**, 075013 (2010); Y. Bai, R. Lu, S. Lu, J. Salvado, and B. A. Stefanek, *Phys. Rev. D* **93**, 073004 (2016); Z. Moss, M. H. Moulai, C. A. Argüelles, and J. M. Conrad, *Phys. Rev. D* **97**, 055017 (2018); E. Baertuzzo, S. Jana, P. A. N. Machado, and R. Z. Funchal, [arXiv:1807.09877](https://arxiv.org/abs/1807.09877); P. Ballett, S. Pascoli, and M. Ross-Lonergan, [arXiv:1808.02915](https://arxiv.org/abs/1808.02915).
- [18] J. Liao and D. Marfatia, *Phys. Rev. Lett.* **117**, 071802 (2016).
- [19] M. Carena, Y.-Y. Li, C. S. Machado, and P. A. N. Machado, and C. E. M. Wagner, *Phys. Rev. D* **96**, 095014 (2017).
- [20] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. D* **79**, 072002 (2009).
- [21] A. A. Aguilar-Arevalo *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. A* **599**, 28 (2009).
- [22] R. B. Patterson, E. M. Laird, Y. Liu, P. D. Meyers, I. Stancu, and H. A. Tanaka, *Nucl. Instrum. Methods Phys. Res., Sect. A* **608**, 206 (2009).
- [23] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. D* **81**, 092005 (2010); *Phys. Rev. D* **88**, 032001 (2013).
- [24] A. A. Aguilar-Arevalo *et al.* (MiniBooNE Collaboration), *Phys. Rev. Lett.* **118**, 221803 (2017); [arXiv:1807.06137](https://arxiv.org/abs/1807.06137).
- [25] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. D* **84**, 072005 (2011).
- [26] R. B. Patterson, Ph. D. thesis, Princeton University, 2007, http://www-boone.fnal.gov/publications/Papers/rbpatter_thesis.pdf.
- [27] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevLett.121.221801> for more information on backgrounds and plots showing agreement between data and the Monte Carlo simulation.
- [28] G. Cheng *et al.*, *Phys. Rev. D* **84**, 012009 (2011); C. Mariani, G. Cheng, J. M. Conrad, and M. H. Shaevitz, *Phys. Rev. D* **84**, 114021 (2011).
- [29] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. D* **81**, 092005 (2010); *Phys. Rev. Lett.* **100**, 032301 (2008).
- [30] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. D* **83**, 052007 (2011); *Phys. Rev. Lett.* **103**, 081801 (2009).
- [31] V. P. Efrosinin, Yu. G. Kudenko, and A. N. Khotjantsev, *Phys. At. Nucl.* **72**, 459 (2009); R. J. Hill, *Phys. Rev. D* **81**, 013008 (2010); *Phys. Rev. D* **84**, 017501 (2011); X. Zhang and B. D. Serot, *Phys. Lett. B* **719**, 409 (2013); *Phys. Rev. C* **86**, 035502 (2012); **86**, 035504 (2012); B. D. Serot and X. Zhang, *Phys. Rev. C* **86**, 015501 (2012).
- [32] E. Wang, L. Alvarez-Ruso, and J. Nieves, *Phys. Rev. C* **89**, 015503 (2014); *Phys. Lett. B* **740**, 16 (2015).
- [33] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. D* **81**, 013005 (2010); *Phys. Lett. B* **664**, 41 (2008).
- [34] D. Rein and L. M. Sehgal, *Phys. Lett.* **104B**, 394 (1981); S. S. Gershtein, Yu. Ya. Komachenko, and M. Yu. Khlopov, *Sov. J. Nucl. Phys.* **33**, 860 (1981); J. A. Harvey, C. T. Hill, and R. J. Hill, *Phys. Rev. Lett.* **99**, 261601 (2007); J. P. Jenkins and T. Goldman, *Phys. Rev. D* **80**, 053005 (2009); A. M. Ankowski, O. Benhar, T. Mori, R. Yamaguchi, and M. Sakuda, *Phys. Rev.*

- Lett. **108**, 052505 (2012); K. M. Graczyk, D. Kielczewska, P. Przewlocki, and J. T. Sobczyk, *Phys. Rev. D* **80**, 093001 (2009).
- [35] A. A. Aguilar-Arevalo *et al.*, *Phys. Rev. Lett.* **102**, 101802 (2009).
- [36] B. Armbruster *et al.*, *Phys. Rev. D* **65**, 112001 (2002).
- [37] N. Agafonova *et al.*, [arXiv:1803.11400](#).
- [38] M. Martini, M. Ericson, G. Chanfray, and J. Marteau, *Phys. Rev. C* **80**, 065501 (2009); M. Martini, M. Ericson, and G. Chanfray, *Phys. Rev. D* **85**, 093012 (2012); **87**, 013009 (2013); D. Meloni and M. Martini, *Phys. Lett. B* **716**, 186 (2012); J. Nieves, I. R. Simo, and M. J. Vicente Vacas, *Phys. Rev. C* **83**, 045501 (2011); J. Nieves, F. Sanchez, I. R. Simo, and M. J. Vicente Vacas, *Phys. Rev. D* **85**, 113008 (2012); O. Lalakulich, K. Gallmeister, and U. Mosel, *Phys. Rev. C* **86**, 014614 (2012); U. Mosel, O. Lalakulich, and K. Gallmeister, *Phys. Rev. Lett.* **112**, 151802 (2014); A. Meucci and C. Giusti, *Phys. Rev. D* **85**, 093002 (2012); G. D. Megias *et al.*, *Phys. Rev. D* **91**, 073004 (2015); P. Coloma and P. Huber, *Phys. Rev. Lett.* **111**, 221802 (2013); J. T. Sobczyk, *Phys. Rev. C* **86**, 015504 (2012).
- [39] M. Ericson, M. V. Garzelli, C. Giunti, and M. Martini, *Phys. Rev. D* **93**, 073008 (2016).
- [40] M. Antonello *et al.*, [arXiv:1503.01520](#).