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Secondary twin variant selection in four types of double twins in titanium

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ABSTRACT

We experimentally characterized four types of double twins in pure titanium and discussed the selection of secondary twin variants based on the apparent Schmid factor (a-SF), the displacement gradient accommodation (DGA), the modified DGA (m-DGA), and the nucleation via dislocation dissociation (NDD). The four types of double twins are classified into two families according to their zone axes: co-family double twins and non-family double twins. Co-family double twins share the zone axis $\langle 10\bar{1}0 \rangle$. Non-family double twins have different zone axes with one along $\langle 11\bar{2}0 \rangle$ and the other along $\langle 10\bar{1}0 \rangle$. Experimental observations reveal that the co-family double twins have one secondary twin variant prevailed over others and the non-family double twins have two secondary twin variants preferred than others. The a-SFs associated with secondary twins are always positive once the primary twin is activated, thus predicting no obvious preference for the selection of secondary twin variants. The m-DGA and NDD correctly predict the selection of secondary twin variants, implying that the preferred secondary twin variant would, to the greatest extent, relax plastic deformation associated with the primary twinning and the nucleation of the secondary twin variant is facilitated by the accumulation of available gliding dislocations at the primary twin boundary.

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1. Introduction

α -Titanium (Ti) is attractive in aerospace engineering, chemical industry, and medical implants due to light weight, high strength, excellent corrosion resistance, and good biocompatibility [1]. Research work has been focused on an understanding of the mechanisms and mechanics of plastic deformation associated with dislocation slips and twins in hexagonal close packed (HCP) crystals in the context of temperature and strain rate [2,3], cyclic loading [4–8], strain path change [9–11], texture [12,13], grain size [14,15], and sample size [16]. At room temperature, the main slip systems in Ti include $\langle a \rangle$ slips on basal {0002}, prismatic {1100}, and pyramidal {1101} planes [17,18]. Dislocations associated with these $\langle a \rangle$ slips, however, cannot accommodate deformation along the c -axis. $\langle c+a \rangle$ slips on the {1011} planes and twinning are possible shear

mechanisms accommodating deformation in the c -axis direction. The dislocations associated with $\langle c+a \rangle$ slips have a large Burgers vector often with a non-planar core, making their activation energetically unfavourable. Twinning, on the other hand, takes place by the combined effect of non-lattice Burgers vector glide and atomic shuffling, and its activation results in the formation of a sheared domain where the crystal experiences a large reorientation with respect to the parent matrix.

Twinning competes favourably with $\langle c+a \rangle$ slips at low temperature or at high deformation rate when dislocations are difficult to overcome the threshold barriers. Tension twinning introduces a tensile strain and compression twinning generates a compressive strain along the c -axis of the crystal. {1012} $\langle$ 1011 $\rangle$ tension twinning and {1122} $\langle$ 1123 $\rangle$ compression twinning are commonly observed in α -titanium at room temperature [19–21] and are referred to as T_1^1 and C_1^1 , respectively, where the subscript represents variants from 1 to 6. Other twinning modes, {1121} $\langle$ 1126 $\rangle$ tension twinning [22–25] and {1124} $\langle$ 2243 $\rangle$ compression twinning

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[26–28], can occur depending on temperature and the loading condition, and are referred to as T_i^I and C_i^I , respectively. For example, $\{11\bar{2}1\}$ and $\{11\bar{2}4\}$ twins are often observed in coarse-grained α -titanium at a high strain rate [26,28–30]. Unlike dislocation slips, twinning causes the reorientation of the twinned domain and generates a localized shear strain along the twinning direction. Once a twin forms, another twin could be activated inside the twin. According to the sequence in nucleating twin variants, deformation twins can be classified into primary twin, secondary twin, and tertiary twin [31,32].

Twin nucleation must precede twin propagation. Twin nucleation is driven by the local stress state and the local atomistic configurations while twin propagation is driven by long-range stress state across the grain. Several criteria have been proposed to address the selection of twin variants observed in experiments [10,30,33–37] and to determine twin variants in polycrystalline plastic models [9,10,38–40]. The standard Schmid factor criterion (referred to as SF, i.e., the highest resolved shear stress selection rule) is extensively applied to account for variant selection of primary twins with the assumption that the local stress is identical to the nominal stress in polycrystalline aggregates [41]. Without considering specific nucleation sites and mechanisms, the SF offers a simple argument that a twin variant with the highest resolved shear stress is favourably activated. However, the local stresses inside a grain indeed differ from the externally applied stress [34,37,42]. To consider the local stress influence on twin nucleation at the meso and micro scales, Beyerlein et al. [43] introduced stress fluctuation criterion that was inspired from the atomistic simulations of dislocation and grain boundary (GB) interactions. The nucleation of primary twins is treated as a stochastic event in terms of the selection of twin mode and twin variant. Such a theoretical speculation is consistent with the experimental observations that some twins are nucleated at GBs with low or negative SFs [37,44,45]. The criterion provides a geometric measure of how well a twin system is oriented with respect to an externally applied stress.

Nucleation of primary twins to some extent does show a preferred selection of twin variants [46]. To explain such a preference observed in experiments, Jonas et al. [34] proposed a displacement gradient accommodation (DGA) criterion for the selection of a primary twin variant. To predict twin variant, the displacement gradients associated with the primary twin are transformed into the crystal reference frame in a neighbouring grain. A twin variant is selected when basal slips in the neighbouring grain can accommodate the transformed deformation. Such a notion provides a measurement of plastic deformation compatibility across a GB. Later, Shi et al. [37] extended the original DGA operation to all possible slip/twinning reference frames in the neighbouring grains, and found that the accommodation could be also realized through twinning often being referred to as a pair of cross-boundary twins [20,47–49]. Inspired by atomistic simulations [50–54], Beyerlein et al. [33] proposed a nucleation criterion for secondary twinning according to dislocation dissociation (NDD) without using the highest resolved shear stress concept and the plastic deformation compatibility. The criterion is based on an assumption that twin nucleation is facilitated by the accumulation, reaction, and dissociation of gliding dislocations at primary twin boundary [33].

The aforementioned three criteria have advantages and disadvantages. Efforts have been made to improve the criteria, such as apparent-SF (a-SF) and modified DGA [29,37,55]. Xu et al. [29] modified the DGA criterion (referred to as m-DGA) to predict secondary twin variant with a focus on minimizing the resultant plastic deformation in the matrix associated with double twinning. In other words, a selected secondary twin variant should, to the

greatest extent, diminish the shear deformation resulting from the primary twin [29]. On the other hand, there has been no systematic assessment of these existing criteria in terms of their capability to predict primary and secondary twinning. In this work, coarse-grained polycrystalline α -titanium was subjected to uniaxial compression at room temperature along the normal direction (ND) of the rolled sheet. We detected four twin modes, $\{11\bar{2}2\}$ and $\{11\bar{2}4\}$ compression twins and $\{10\bar{1}2\}$ and $\{11\bar{2}1\}$ tension twins. More interestingly, four types of double twins were observed. By comparing with the experimental observations, we systematically evaluate the existing criteria for their capabilities to predict secondary twin variant selection in the double twins. Advantages of each criterion are discussed based on the experimental and theoretical results.

2. Secondary twin variants in four types of double twins

2.1. Crystallographic characteristics of double twins

$\{10\bar{1}m\}$ ($m = 1, 2, 3$) type and $\{11\bar{2}n\}$ ($n = 1, 2, 4$) type twins are often observed in hexagonal metals, as depicted in Fig. 1a and Fig. 1b, respectively. The former type has the zone axis parallel to $\langle 11\bar{2}0 \rangle$ and the latter type has the zone axis parallel to $\langle 10\bar{1}0 \rangle$. Due to 12 rotational symmetry matrices in an HCP structure, there are six equivalent variants for each twinning system. Under mechanical deformation at room temperature, $\{10\bar{1}2\}$ tension twins and $\{11\bar{2}2\}$ compression twins are commonly observed in α -Titanium [20,30,56]. $\{11\bar{2}1\}$ tension twins and $\{11\bar{2}4\}$ compression twins can also be produced but have a low volume fraction. When primary twins form, the twinned domain may be twinned further, forming double twins. To facilitate discussions, T_i^I and T_i^{II} ($i = 1, 2, \dots, 6$) are used to denote the six $\{10\bar{1}2\}$ and the six $\{11\bar{2}1\}$ tension twin variants, respectively, and C_i^I and C_i^{II} ($i = 1, 2, \dots, 6$) represent respectively the six $\{11\bar{2}2\}$ and the six $\{11\bar{2}4\}$ compression twin variants. For example, T_1^I is the $(10\bar{1}2)[1011]$ twin variant, T_1^{II} is the $(11\bar{2}1)[1\bar{1}26]$ twin variant, C_1^I is the $(11\bar{2}2)[11\bar{2}3]$ twin variant, and C_1^{II} is the $(11\bar{2}4)[22\bar{4}3]$ variant, as depicted in Fig. 1. According to the crystallography of an HCP structure, T_i^I twin variant can be obtained by rotating the T_1^I variant ($i - 1$) $\times 60$ degrees about the $\langle 0001 \rangle$ axis. The same operation is applicable to other twins.

Four types of double twins that were observed in deformed Ti specimens can be classified into two families: co-family and non-family double twins. The co-family double twins include $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ ($C_i^I \rightarrow T_j^{II}$) and $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ ($T_i^{II} \rightarrow C_j^{II}$) where both twins share the zone axis $\langle 10\bar{1}0 \rangle$ and the arrow in the notation refers to the process from the primary twinning to the secondary twinning. The non-family double twins include $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ ($C_i^I \rightarrow T_j^I$) and $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ ($C_i^{II} \rightarrow T_j^I$) where both twins have different zone axes with one along $\langle 11\bar{2}0 \rangle$ and the other along $\langle 10\bar{1}0 \rangle$.

Each type of double twins can be further classified into a few groups according to the misorientation between the secondary twin variant and the parent grain. Taking $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ ($C_i^I \rightarrow T_j^{II}$) co-family double twins as an example, $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ double twins can be categorized into four groups according to the misorientation angle between a secondary twin variant $\{11\bar{2}1\}$ and the parent grain. Group I includes $C_i^I \rightarrow T_i^{II}$ double twins that hold a misorientation angle of 29.5° about the axis $\langle 1\bar{1}00 \rangle$. Group II includes $C_i^I \rightarrow T_{i+1}^{II}$ and $C_i^I \rightarrow T_{i+5}^{II}$ double twins that have a misorientation angle of 55° about the axis $\langle \bar{1}0 \ 5 \ 5 \ \bar{3} \rangle$. Group III includes $C_i^I \rightarrow T_{i+3}^{II}$ double twins that hold a misorientation angle of 80.6° around the axis $\langle 1\bar{1}00 \rangle$. Group IV includes $C_i^I \rightarrow T_{i+2}^{II}$ and $C_i^I \rightarrow T_{i+4}^{II}$ double twins that hold a misorientation angle of 86.8° around the axis $\langle \bar{1}5 \ 10 \ 5 \ \bar{3} \rangle$. Table 1 summaries the geometric characteristics of the six $\{11\bar{2}1\}$ secondary twin variants inside the $(11\bar{2}2)[11\bar{2}3]$

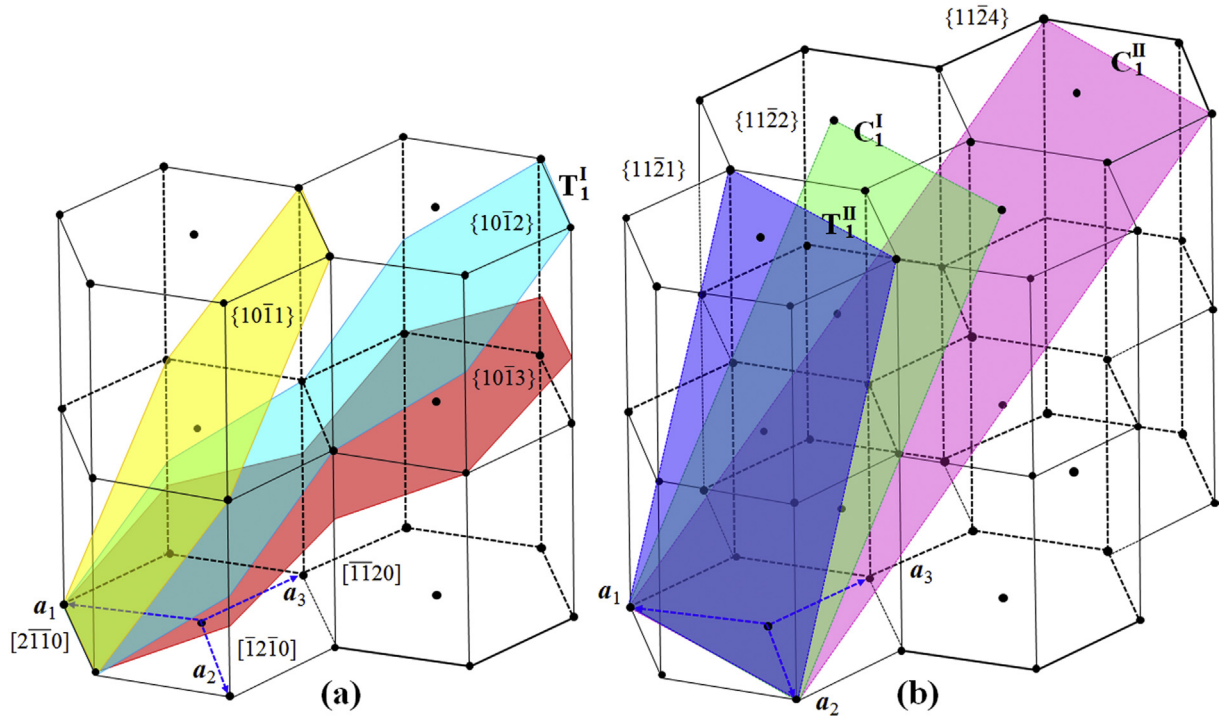


Fig. 1. Main twinning systems in α -titanium: (a) $\{10\bar{1}1\}$ twinning plane in yellow, $\{10\bar{1}2\}$ twinning plane in light blue and $\{10\bar{1}3\}$ twinning plane in red (b) $\{11\bar{2}1\}$ twinning plane in purple, $\{11\bar{2}2\}$ twinning plane in green and $\{11\bar{2}4\}$ twinning plane in pink. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

Table 1
Geometric characters between the primary twin and secondary twins.

Secondary twin	Misorientation	Group	Nomination	$l_{(S-P)}$	$l_{(B-P)}$	$\theta_{(SB/P)}$
$(11\bar{2}2)[11\bar{2}3](C_1^I) \rightarrow \{11\bar{2}1\}$ double twins (Co-family)						
$(11\bar{2}1)[1\bar{1}26]$	$\langle 1\bar{1}00 \rangle 29.5^\circ$	I	$C_1^I \rightarrow T_{i+1}^{II}$	$[1\bar{1}00]$	$[1\bar{1}00]$	0°
$(\bar{1}2\bar{1}1)[1\bar{2}16]$	$\langle \bar{1}0\ 5\ 5\ \bar{3} \rangle 55^\circ$	II	$C_1^I \rightarrow T_{i+1}^{II}$	$[\bar{1}\bar{1}23]$	$[1\bar{1}00]$	90°
$(2\bar{1}11)[2\bar{1}16]$	$\langle \bar{1}5\ 10\ 5\ \bar{3} \rangle 86.8^\circ$	IV	$C_1^I \rightarrow T_{i+2}^{II}$	$[1323]$	$[1\bar{1}00]$	58.4°
$(\bar{1}\bar{1}21)[1\bar{1}26]$	$\langle 1\bar{1}00 \rangle 80.6^\circ$	III	$C_1^I \rightarrow T_{i+3}^{II}$	$[1\bar{1}00]$	$[1\bar{1}00]$	0°
$(1\bar{2}11)[\bar{1}2\bar{1}6]$	$\langle \bar{1}5\ 10\ 5\ \bar{3} \rangle 86.8^\circ$	IV	$C_1^I \rightarrow T_{i+4}^{II}$	$[\bar{3}123]$	$[1\bar{1}00]$	58.4°
$(2\bar{1}\bar{1}1)[\bar{2}116]$	$\langle \bar{1}0\ 5\ 5\ \bar{3} \rangle 55^\circ$	II	$C_1^I \rightarrow T_{i+5}^{II}$	$[\bar{1}\bar{1}23]$	$[1\bar{1}00]$	90°
$(11\bar{2}1)[\bar{1}\bar{1}26](T_1^{II}) \rightarrow \{11\bar{2}4\}$ double twins (Co-family)						
$(11\bar{2}4)[2\bar{2}4\bar{3}]$	$\langle 1\bar{1}00 \rangle 68.2^\circ$	III	$T_1^{II} \rightarrow C_{i+1}^{II}$	$[1\bar{1}00]$	$[1\bar{1}00]$	0°
$(\bar{1}2\bar{1}4)[2\bar{4}2\bar{3}]$	$\langle \bar{3}5\ 10\ 25\ 9 \rangle 88.3^\circ$	IV	$T_1^{II} \rightarrow C_{i+1}^{II}$	$[3\bar{4}13]$	$[1\bar{1}00]$	39.5°
$(2\bar{1}14)[\bar{4}22\bar{3}]$	$\langle 6\ \bar{1}0\ 4\ 3 \rangle 65.6^\circ$	II	$T_1^{II} \rightarrow C_{i+2}^{II}$	$[4\bar{5}13]$	$[1\bar{1}00]$	32.6°
$(\bar{1}\bar{1}24)[2\bar{2}4\bar{3}]$	$\langle 1\bar{1}00 \rangle 41.9^\circ$	I	$T_1^{II} \rightarrow C_{i+3}^{II}$	$[1\bar{1}00]$	$[1\bar{1}00]$	0°
$(1\bar{2}14)[2\bar{4}2\bar{3}]$	$\langle 6\ \bar{1}0\ 4\ 3 \rangle 65.6^\circ$	II	$T_1^{II} \rightarrow C_{i+4}^{II}$	$[5\bar{4}1\bar{3}]$	$[1\bar{1}00]$	32.6°
$(2\bar{1}\bar{1}4)[\bar{4}22\bar{3}]$	$\langle \bar{3}5\ 10\ 25\ 9 \rangle 88.3^\circ$	IV	$T_1^{II} \rightarrow C_{i+5}^{II}$	$[4\bar{3}1\bar{3}]$	$[1\bar{1}00]$	39.5°
$(11\bar{2}2)[11\bar{2}3](C_1^I) \rightarrow \{10\bar{1}2\}$ double twins (Non-family)						
$(10\bar{1}2)[\bar{1}011]$	$\langle \bar{5}503 \rangle 48.4^\circ$	II	$C_1^I \rightarrow T_i^I$	$[\bar{4}223]$	$[1\bar{1}00]$	47.3°
$(01\bar{1}2)[0\bar{1}11]$	$\langle \bar{5}503 \rangle 48.4^\circ$	II	$C_1^I \rightarrow T_{i+1}^I$	$[2\bar{4}23]$	$[1\bar{1}00]$	47.3°
$(\bar{1}102)[1\bar{1}01]$	$\langle 7\bar{4}30 \rangle 87.9^\circ$	III	$C_1^I \rightarrow T_{i+2}^I$	$[2\bar{4}23]$	$[1\bar{1}00]$	47.3°
$(1012)[10\bar{1}1]$	$\langle 5\bar{1}43 \rangle 41.3^\circ$	I	$C_1^I \rightarrow T_{i+3}^I$	$[\bar{1}0\ 8\ 2\ 3]$	$[1\bar{1}00]$	19.8°
$(0\bar{1}12)[01\bar{1}1]$	$\langle 5\bar{1}43 \rangle 41.3^\circ$	I	$C_1^I \rightarrow T_{i+4}^I$	$[8\ \bar{1}0\ 2\ 3]$	$[1\bar{1}00]$	19.8°
$(1\bar{1}02)[\bar{1}101]$	$\langle 7\bar{4}30 \rangle 87.9^\circ$	III	$C_1^I \rightarrow T_{i+5}^I$	$[\bar{4}223]$	$[1\bar{1}00]$	47.3°
$(11\bar{2}4)[2\bar{2}4\bar{3}](C_1^{II}) \rightarrow \{10\bar{1}2\}$ double twins (Non-family)						
$(10\bar{1}2)[\bar{1}011]$	$\langle 11\bar{2}3 \rangle 39.5^\circ$	I	$C_1^{II} \rightarrow T_i^I$	$[2\bar{2}4\bar{3}]$	$[1\bar{1}00]$	90°
$(01\bar{1}2)[0\bar{1}11]$	$\langle 11\bar{2}3 \rangle 39.5^\circ$	I	$C_1^{II} \rightarrow T_{i+1}^I$	$[2\bar{2}4\bar{3}]$	$[1\bar{1}00]$	90°
$(\bar{1}102)[1\bar{1}01]$	$\langle 7\ 20\ \bar{2}7\ 20 \rangle 44.3^\circ$	II	$C_1^{II} \rightarrow T_{i+2}^I$	$[1543]$	$[1\bar{1}00]$	55.8°
$(1012)[10\bar{1}1]$	$\langle \bar{1}3\ \bar{2}6\ 39\ 4 \rangle 89.5^\circ$	III	$C_1^{II} \rightarrow T_{i+3}^I$	$[10\ \bar{1}4\ 4\ 3]$	$[1\bar{1}00]$	20.2°
$(0\bar{1}12)[01\bar{1}1]$	$\langle \bar{1}3\ \bar{2}6\ 39\ 4 \rangle 89.5^\circ$	III	$C_1^{II} \rightarrow T_{i+4}^I$	$[\bar{1}4\ 10\ 4\ 3]$	$[1\bar{1}00]$	20.2°
$(1\bar{1}02)[\bar{1}101]$	$\langle 7\ 20\ \bar{2}7\ 20 \rangle 44.3^\circ$	II	$C_1^{II} \rightarrow T_{i+5}^I$	$[\bar{5}143]$	$[1\bar{1}00]$	55.8°

(c_1^1) primary twin. For convenience in performing the NDD analysis, we also characterize the intersection lines between twin planes, slip planes and the primary twin boundary. The intersection line between a secondary twin plane and a primary twin plane is denoted by $l_{(S-P)}$ and that between a basal plane and a primary twin plane is represented by $l_{(B-P)}$. The angle between $l_{(S-P)}$ and $l_{(B-P)}$ is denoted by $\theta_{(S-B/P)}$. Similarly, $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins are divided into four groups (refer to Table 1). Group I and Group III have one variant while Group II and Group IV have two variants. Unlike co-family double twins, $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ and $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ non-family double twins are divided into three groups (Table 1) with each group containing two twin variants.

2.2. Experimental observations

A commercial rolled pure titanium T40 sheet (ASTM grade 2) with a thickness of 1.5 mm was used in this work. Compressive samples with a dimension of 10 mm (RD) $\times$ 8 mm (TD) $\times$ 1.5 mm (ND) were processed from the as-received sheet. Here, RD, TD and ND represent rolling direction, transverse direction and normal direction of the rolled sheet, respectively. Then, the prepared samples were annealed in a vacuum furnace at 800 °C for 2 h. After that, one-step compression was conducted at room temperature at

a strain rate $1 \times 10^{-3} \text{ s}^{-1}$ on a Zwick 120T machine. The reduction of the thickness along the ND is 8.7%. After compression, the RD-TD plane (rolling surface) of the deformed sample was ground with SiC papers of grits from 1200# to 4000# to remove the surface layer. Electrolytic polishing was then performed in a solution of 10% perchloric acid and 90% methanol at 35 V for 5 s at 5 °C and electron backscatter diffraction (EBSD) measurements were made in a JEOL JSM-6500F field emission gun scanning electron microscopy (SEM) equipped with an EBSD camera and the AZtec acquisition software package (Oxford Instruments).

In Ti, $\{11\bar{2}1\}$ tension twinning rotates the crystal by $\sim 35^\circ$ around a $\langle 1\bar{1}00 \rangle$ axis. $\{10\bar{1}2\}$ tension twinning results in a rotation of the crystal by $\sim 87^\circ$ around a $\langle 1\bar{2}10 \rangle$ axis. $\{11\bar{2}2\}$ compression twinning rotates the twinned domain by $\sim 64^\circ$ around a $\langle 1\bar{1}00 \rangle$ axis. $\{11\bar{2}4\}$ compression twinning rotates the twinned domain by $\sim 77^\circ$ around a $\langle 1\bar{1}00 \rangle$ axis. According to these crystallographic characteristics, we further characterized four types of double twins in deformed specimens.

Fig. 2a shows an EBSD pattern where a $\{11\bar{2}2\}$ compression twin contains secondary $\{11\bar{2}1\}$ and $\{10\bar{1}2\}$ tension twins. The boundaries formed by the $\{11\bar{2}1\}$ tension twin, $\{10\bar{1}2\}$ tension twin, and $\{11\bar{2}2\}$ compression twin with the parent crystal are indicated in Fig. 2b in yellow, red, and green, respectively, with a tolerance of

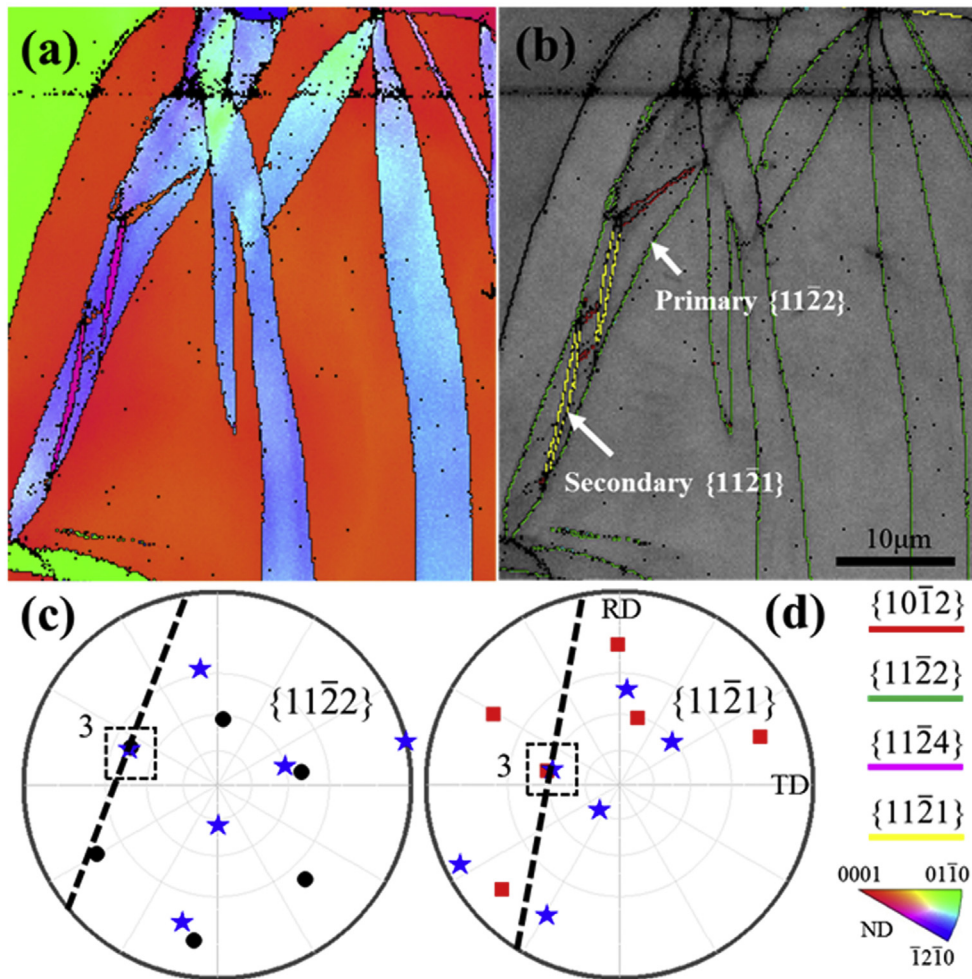


Fig. 2. (a) An EBSD map of $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ double twins with twin boundaries indicated by color lines based on the misorientation axis/angle between the twin and grain in (b). In (c), the black dots represent six $\{11\bar{2}2\}$ twin planes in the matrix and the blue stars indicate six $\{11\bar{2}2\}$ twin planes in the primary twin. In (d), the blue stars represent six $\{11\bar{2}1\}$ twin planes in the primary twin and the red squares indicate six $\{11\bar{2}1\}$ twin planes in the secondary twin. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

$\pm 5^\circ$ deviation from the ideal crystallographic axis and angle. Such a tolerance is used for the other double twin observations made throughout the current work. With the help of pole figures in Fig. 2c, the misorientation angle between $(\bar{2}112)$ planes in the primary twin and in the matrix is 0.89° , which is the smallest among all possibilities. The dashed black line in Fig. 2c indicates the trace of the twinning plane, which is in agreement with the detected trace of the twinning plane shown in Fig. 2b. Therefore, C_3^I twin variant is identified as the primary twin. In the pole figure of $\{11\bar{2}1\}$ planes (Fig. 2d), the misorientation angle between $(\bar{2}111)$ planes in the C_3^I primary twin and the secondary twin is 3.12° . The trace of this twinning plane in Fig. 2d is practically parallel to that in Fig. 2b. T_3^I is, therefore, identified as a secondary twin variant. The double twin configuration can be regarded as $C_3^I \rightarrow T_3^I$, which belongs to Group I as indicated in Table 1. In the EBSD results, 43 $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ double twins are detected including 41 Group I and two Group II double twins. Therefore, Group I, i.e., $C_3^I \rightarrow T_3^I$ is the prevailed double twin over the others.

Fig. 3a shows an EBSD map where the $\{11\bar{2}1\}$ tension twin contains secondary $\{11\bar{2}4\}$ compression twins and the twin boundaries are indicated in Fig. 3b. With the help of pole figures in Fig. 3c, the misorientation angle between $(\bar{2}1\bar{1}1)$ planes in the primary twin and that in the matrix is 1.46° , which is the closest

among all possibilities. In addition, the dashed black line in Fig. 3c indicates the trace of the twinning plane, which is in agreement with the detected trace of the twinning plane as shown in Fig. 3a. Therefore, T_2^I twin variant is identified as the primary twin. In the pole figure of $\{11\bar{2}4\}$ planes (Fig. 3d), the misorientation angle between $(\bar{2}1\bar{1}4)$ planes in the T_2^I primary twin and the secondary twin is 2.37° . The trace of this twinning plane in Fig. 3d is almost parallel to that in Fig. 3b. C_2^I is thus the secondary twin variant. The double twin configuration can be described as $T_2^I \rightarrow C_2^I$, which belongs to Group III as indicated in Table 1. It is noted that such kind of double twins are rarely observed in our experiments. More experiments are required for a solid conclusion.

Fig. 4a shows an EBSD map where the $\{11\bar{2}2\}$ compression twin contains secondary $\{10\bar{1}2\}$ tension twins and the twin boundaries are indicated in Fig. 4b. From the pole figures shown in Fig. 4c, C_2^I twin variant can be identified as the primary twin. In the pole figure of $\{10\bar{1}2\}$ planes (Fig. 4d), T_2^I is the secondary twin variant. The double twin can be described as $C_2^I \rightarrow T_2^I$ which belongs to Group II. In the C_2^I primary twin, there are two variants T_2^I and T_3^I belonging to Group I. A similar analysis was applied to $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ double twins in Fig. 4e, and the results indicate a detection of $C_4^I \rightarrow T_5^I$ double twin (Group II). In the EBSD results, 425 $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins were detected and 85.6% of them belong to Group II

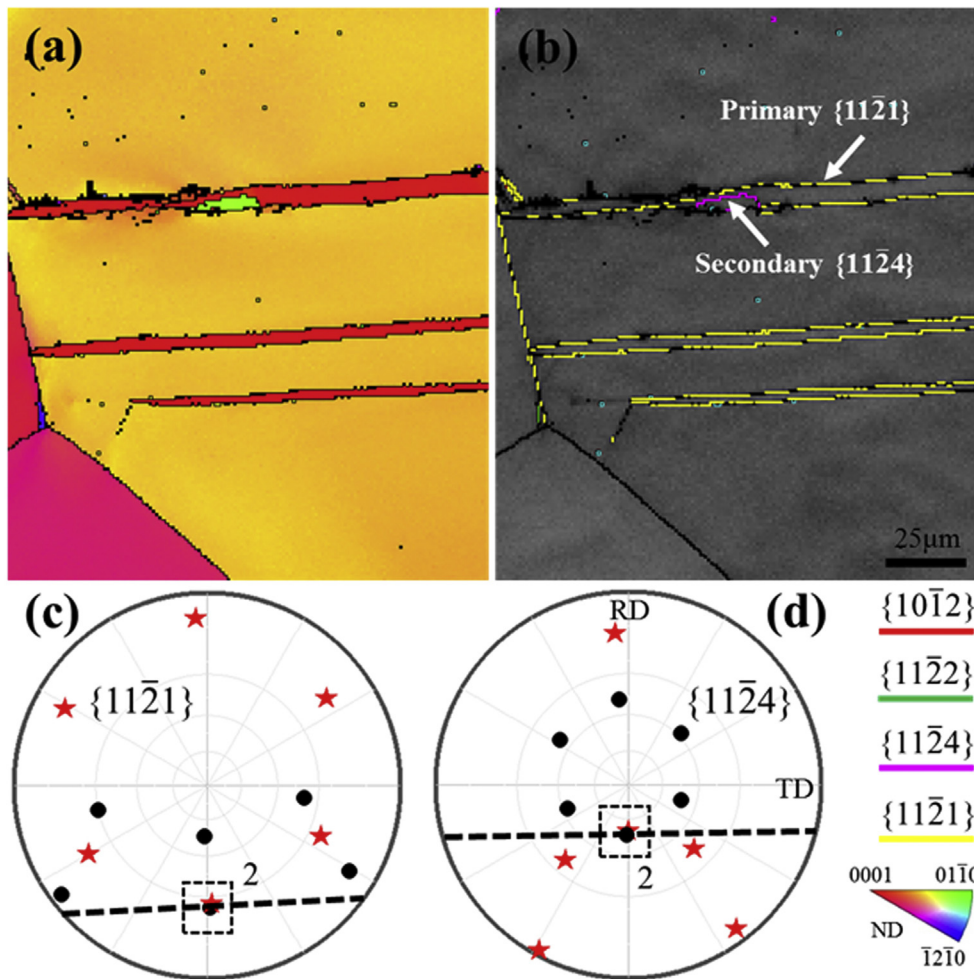


Fig. 3. (a) An EBSD map of $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins with twin boundaries indicated by different color according to the misorientation axis/angle between the twin and grain in (b). In (c), the black dots represent six $\{11\bar{2}1\}$ twin planes in the matrix and the blue stars indicate six $\{11\bar{2}1\}$ twin planes in the primary twin. In (d), the blue stars represent six $\{11\bar{2}4\}$ twin planes in the primary twin and the red squares indicate six $\{11\bar{2}4\}$ twin planes in the secondary twin. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

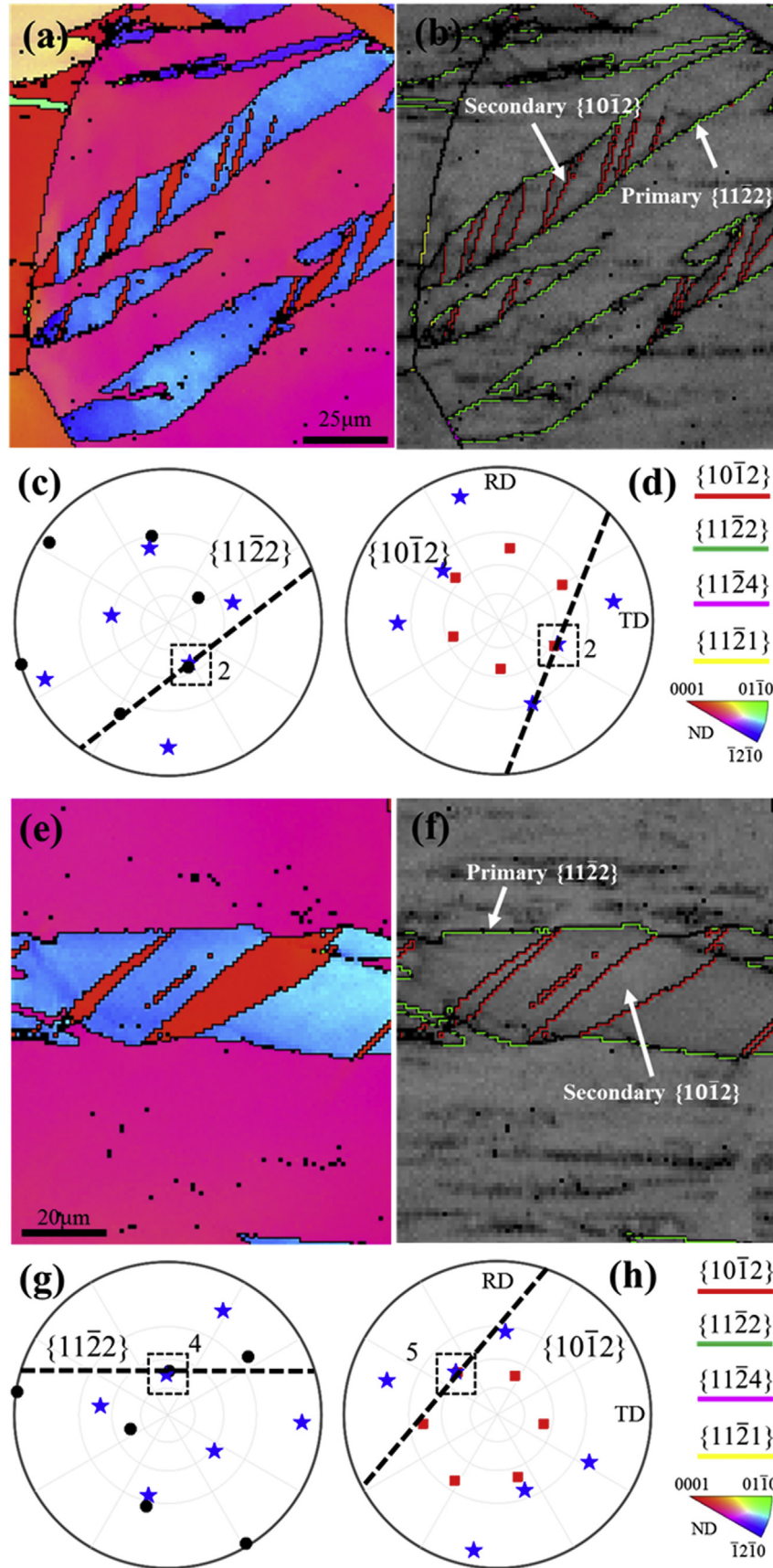


Fig. 4. Examples of $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ double twins: (a) $C_2^I \rightarrow T_1^I$ double twins with twin boundaries indicated by color lines according to the misorientation axis/angle between the twin and grain in (b). (e) $C_4^I \rightarrow T_5^I$ double twins with indicating the twin boundaries by colorful lines according to the misorientation axis/angle between the twin and grain in (f). In (c) and (g), the black dots represent six $\{11\bar{2}2\}$ twin planes in the matrix and the blue stars indicate six $\{11\bar{2}2\}$ twin planes in the primary twin. In (d) and (h), the blue stars represent six $\{10\bar{1}2\}$ twin planes in the primary twin and the red squares indicate six $\{10\bar{1}2\}$ twin planes in the secondary twin. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

double twins. Therefore, the preferred double twins can be described as $C_i^I \rightarrow T_i^I$ or T_{i+1}^I . The a-SF analysis will be used to further identify the variant selection in the preferred group.

Fig. 5a shows an EBSD map where the $\{11\bar{2}4\}$ compression twin contains secondary $\{10\bar{1}2\}$ tension twins. The twin boundaries are indicated in Fig. 5b. The pole figures in Fig. 5c indicates that C_6^I twin variant is the primary twin. In the pole figure of $\{10\bar{1}2\}$ planes (Fig. 5d), T_6^I is the active secondary twin variant. The double twin configuration can be described as $C_6^I \rightarrow T_6^I$, which belongs to Group I. A similar analysis was applied to $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins in Fig. 5e, $C_4^I \rightarrow T_5^I$ double twin (Group I) can be detected. In the EBSD results, 41 $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins were identified and all of them belong to Group I double twins. Therefore, Group I double twin, $C_i^I \rightarrow T_i^I$ or T_{i+1}^I , is the prevailed double twins over other two groups.

In summary, co-family double twins can be categorized into four groups and only one double twin variant is prevailed in each primary twin. For $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ double twins, the dominant double twins are described as $C_i^I \rightarrow T_i^I$ (Group I). For $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins, the dominant double twins are described as $T_i^I \rightarrow C_i^I$ (Group III). Non-family double twins can be divided into three groups, and two variants are possibly selected in the prevailed group. For $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ double twins, the preferred double twins can be described as $C_i^I \rightarrow T_i^I$ or T_{i+1}^I (Group II). For $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins, the preferred double twins can be described as $C_i^I \rightarrow T_i^I$ or T_{i+1}^I (Group I).

3. Selection of secondary twin variants

The mode and variant of secondary twins are inevitably related to the local stresses. An analysis of the stress field could be simplified regarding the secondary twinning knowing that the resolved shear stress associated with primary twinning is positive. The resolved shear stress associated with a secondary twin can be directly related to the resolved shear stress activating the primary twinning. With this notion in mind, Xu et al. [29] recently proposed a concept of apparent SF (a-SF). First, a primary loading domain in which the active primary twin is the most favorable variant among all possibilities is identified in an inverse pole figure. The SFs associated with the six secondary twin variants are calculated under an applied loading condition in the primary loading domain, and the results are plotted in the primary loading domain. The prevailed secondary twin variant is corresponding to the variant with relative high SFs in most of the primary loading domain. By comparing the SFs with the experiments [21], Xu et al. found that a twin variant can be activated with a relatively large SF but not necessary at the highest SF value among all the variants.

Experimental observations reveal that a secondary twin variant associated with sequential twinning or transmission twinning [27,48,49,57,58] has a strong correlation with the incoming twin or slip [59–61], and a secondary twin variant associated with double twinning is related to the primary twin [10,33,62,63]. The DGA and the m-DGA criteria have been applied to predict the selection of these secondary twin variants [29,55,64]. Here we will compare the predictions according to the DGA and m-DGA with experimental observations. Furthermore, we perform a nucleation analysis on the primary twin boundary based on the gliding dislocation dissociation (NDD). In what follows, we will examine these criteria for selecting secondary twin variant in four types of double twins experimentally observed in Ti.

3.1. $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ co-family double twins

3.1.1. Apparent SF (a-SF) analysis

$\{11\bar{2}2\}$ twin variant C_1^I is chosen as the primary twin for the

analysis without losing generality. The primary loading domain is determined according to the SF associated with C_1^I twin when a grain is subjected to uniaxial compression. SFs of the six $\{11\bar{2}2\}$ twin variants are determined for a given loading direction relative to the grain orientation. The loading directions with SFs being larger than 0.3 are grouped into the primary loading domain (associated with primary twin C_1^I) in an inverse pole figure (IPF) as shown in Fig. 6a. Afterward, the SFs associated with the six $\{11\bar{2}1\}$ secondary twin variants T_j^I ($j = 1 \dots 6$) inside the primary twin C_1^I are determined for a loading case in the primary loading domain and are plotted into IPFs (Fig. 6b–g) where the red color indicates positive SFs and the blue color denotes negative SFs. The results show that twin variants T_1^I and T_4^I hold high SFs under most loading directions, while other four secondary twin variants have moderate SFs only at specific loading cases and low SFs at most loading cases. Results in Fig. 6b–g suggest that activation of secondary twin variants T_1^I and T_4^I has a higher probability than the other four primary compression twin variants. When the SFs of the secondary twin variants larger than 0.2 are selected for an IPF, the area of the loading domain is virtually zero for T_2^I , T_3^I , T_5^I , and T_6^I , while the area for T_1^I and T_4^I are non-zero and identical. Therefore, an a-SF analysis seems to be able to explain the high probability for nucleating secondary twin variants in Group I and Group III, and the low frequency for Group II and Group IV twin variants. More importantly, the a-SF analysis clearly shows that the SFs for six twin variants are positive in most loading domain as long as the primary twin is subjected to a positive resolved shear stress.

3.1.2. Nucleation via dislocation dissociation (NDD)

Assuming that a secondary twin is nucleated via twinning dislocations from the primary twin boundary, we first analyze the geometric characteristics of the intersection between secondary twin planes and the primary twin plane, as shown in Fig. 7. The intersection line between secondary twin variants (belonging to Group I and Group III) and the primary twin plane is parallel to $[1\bar{1}00]$, which also lies on the basal plane (Fig. 7a). Therefore, the emission of twinning dislocations can be considered as the dissociation of basal $\langle a \rangle$ dislocations on the secondary twin plane. For Group II, the intersection line between the primary twin plane and secondary twin plane is parallel to $[\bar{1}123]$, which lies on prismatic plane $\{1\bar{1}00\}$ and is perpendicular to the intersection line between the basal and the primary twin plane (Fig. 7b). As a result, emitting twinning dislocations can be considered as the dissociation of prismatic $\langle a \rangle$ dislocation on the secondary twin plane. If $\langle c+a \rangle$ dislocations on the prismatic plane can be activated, emitting twinning can be also dissociated into emissary twin dislocations (TDs) on the secondary twin plane. For Group IV, the intersection line between the primary and secondary twin planes is along $\langle \bar{3}123 \rangle$ which does not lie on any usual slip planes (Fig. 7c). Therefore, dislocation dissociation mechanisms would be difficult to activate secondary twins unless dislocations can easily climb on the primary twin plane and are parallel to the intersection line along $\langle \bar{3}123 \rangle$.

Fig. 7d illustrates secondary twinning mechanisms associated with the dissociation of $\langle a \rangle$ dislocations into TDs. When $(0002) \langle 11\bar{2}0 \rangle$ dislocations approach the primary twin boundary, Fig. 7e illustrates the dissociation of the dislocation into TDs associated with Group I (T_1^I) and Group III (T_4^I) secondary twins. For Groups I to III, an $\langle a \rangle$ dislocation b_a can be dissociated into x secondary twinning dislocations $xb_t^{\{11\bar{2}1\}}$ and a residual dislocation b_r^m , where the subscript m represents the m th ($m = 1, 2, 4, 6$) secondary twin,

$$b_i \Rightarrow b_r^m + xb_t^{\{11\bar{2}1\}} \quad (1)$$

In Eq. (1) $b_i = a_i$ with a magnitude of 0.295 nm for α -Ti. The

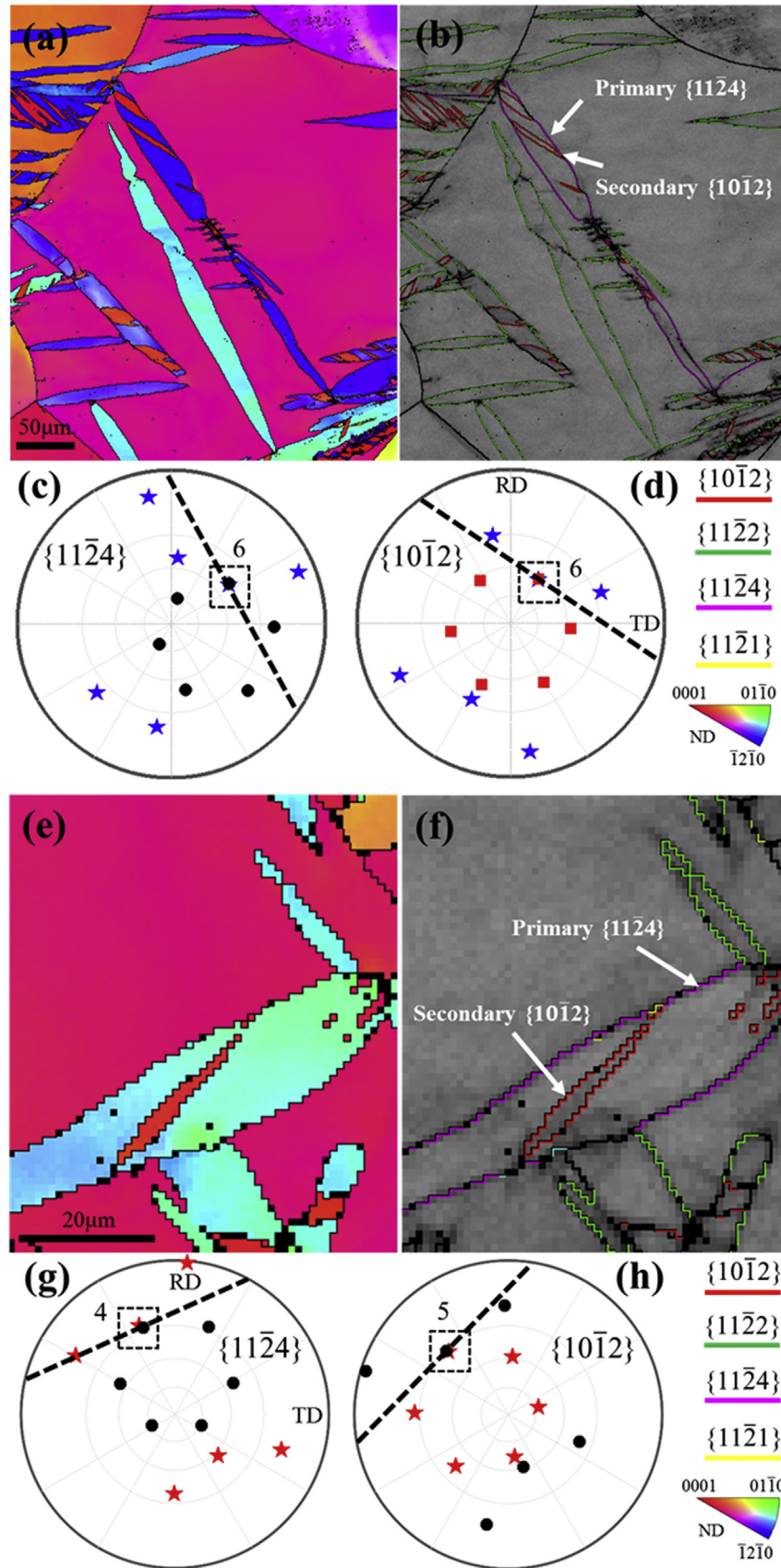


Fig. 5. Examples of $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins: (a) $C_6^{\text{II}} \rightarrow T_6^{\text{I}}$ double twins with twin boundaries indicated by color lines according to the misorientation axis/angle between the twin and grain in (b). (e) $C_4^{\text{II}} \rightarrow T_5^{\text{I}}$ double twins with indicating the twin boundaries by colorful lines according to the misorientation axis/angle between the twin and grain in (f). In (c) and (g), the black dots represent six $\{11\bar{2}4\}$ twin planes in the matrix and the blue stars indicate six $\{11\bar{2}4\}$ twin planes in the primary twin. In (d) and (h), the blue stars represent six $\{10\bar{1}2\}$ twin planes in the primary twin and the red squares indicate six $\{10\bar{1}2\}$ twin planes in the secondary twin. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

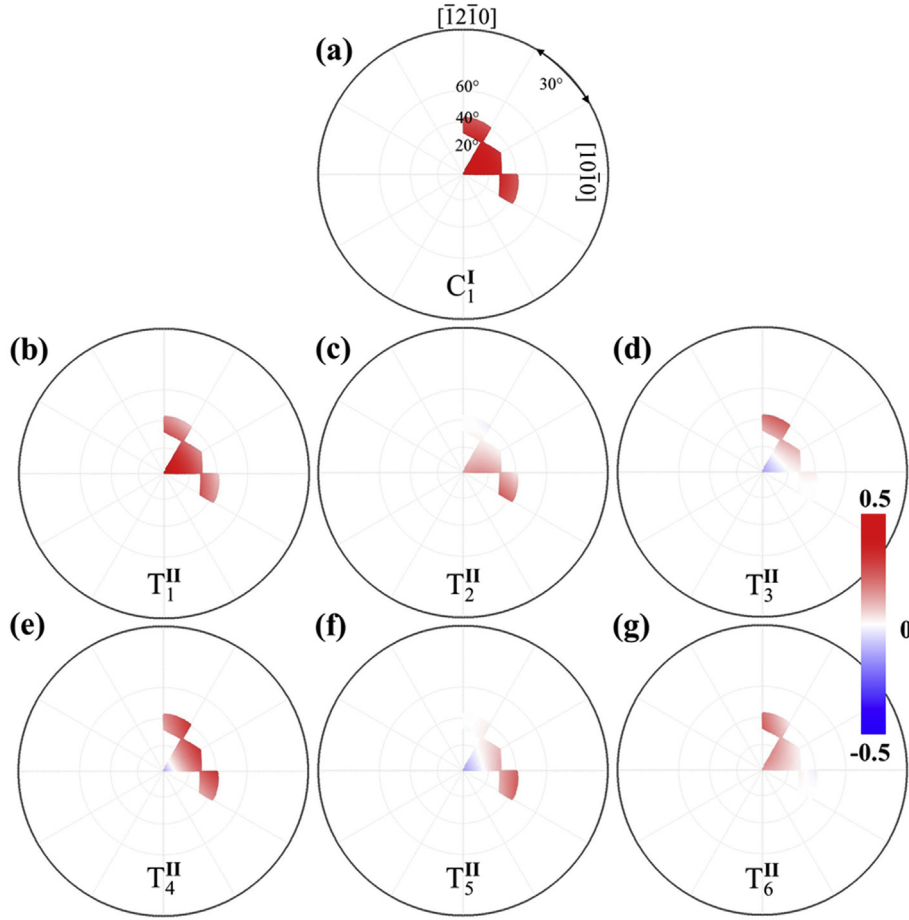


Fig. 6. (a) Stress domain where the SF of a C_1^I primary twin variant is positive and the largest among six possibilities. Under this stress domain, the SF of: (b) T_1^{II} , (c) T_2^{II} , (d) T_3^{II} , (e) T_4^{II} , (f) T_5^{II} , (g) T_6^{II} secondary twin variants.

three base vectors a_1 , a_2 , and a_3 are equal to $\frac{1}{3}[2\bar{1}10]$, $\frac{1}{3}[\bar{1}2\bar{1}0]$, $\frac{1}{3}[\bar{1}\bar{1}20]$, respectively. The Burgers vector of $\{11\bar{2}1\}$ TD $b_t^{(11\bar{2}1)} = \lambda <\bar{1}126>$, $\lambda = \frac{1}{3+12k^2} = 0.0301$ for Ti ($a = 0.295$ nm and $c = 0.4683$ nm, $k = c/a$). For T_1^{II} and T_4^{II} secondary twins, three $<a>$ dislocations on basal plane can act as sources for the dissociation. However, there is only one b_3 dislocation on prismatic plane that acts as the source for the dissociation into secondary twin dislocations belonging to T_2^{II} and T_6^{II} secondary twins. According to Frank's law, the change in the elastic energy associated with these dissociations is described as follows for T_1^{II} :

$$|b_3|^2 = |b_r^1|^2 + |xb_t^{(11\bar{2}1)}|^2 \quad (2)$$

In the process of the dislocation dissociation, we use the maximum allowable number of twinning dislocations (x value) that satisfy Frank's law. When $x = 1$, the elastic energy before and after the dissociation is identical. Results summarized in Table 2 reveal that the dissociation of $<a>$ dislocations only favourably produces twinning dislocation associated with T_1^{II} secondary twin variant. The other dissociations corresponding to T_2^{II} , T_4^{II} and T_6^{II} secondary twins are accompanied with an increase in the elastic energy after dissociation. It is noted that only the dissociation of b_3 dislocation is energetically favorable for producing twinning dislocation of T_1^{II} secondary twin variant. In other words, $C_1^I \rightarrow T_1^{II}$ is predominant

among all double twinning. Such a speculation is consistent with the a-SF analysis and agrees with experimental observation.

3.1.3. Displacement gradient accommodation analysis (DGA and m-DGA)

The original DGA criterion predicts a preference of nucleation of a twin variant whose deformation can be accommodated by easy slips in the vicinity of the twin domain. For a $(11\bar{2}2)[1123]$ primary twin (C_1^I), the displacement gradients induced by six possible secondary $\{11\bar{2}1\}$ twin (D_{ST}) variants in its twinning frame can be resolved into three directions, $x' \parallel [10\bar{1}0]$, $y' \parallel [\bar{1}2\bar{1}0]$, and $z' \parallel [0001]$, in the crystal frame of the parent grain. $e_{x'z'}$ and $e_{y'z'}$ represent the accommodation by double and single basal slips, respectively; $e_{z'x'}$ and $e_{z'y'}$ indicate the accommodation by twinning, and $e_{x'y'}$ and $e_{y'x'}$ are associated with the accommodation by double and single prismatic slips, respectively [34]. The transformed displacement gradient tensors, D_{ST} for the six potential secondary twin variants are listed in the first column in Table 2, where the accommodation associated with basal slip, prismatic slip and twinning is indicated in red, blue and green font, respectively. The accommodation components required in the parent grain are defined as $[d_1, d_2, d_3]$, where $d_1 = |e_{x'z'}| + |e_{y'z'}|$, $d_2 = |e_{x'y'}| + |e_{y'x'}|$, and $d_3 = |e_{z'x'}| + |e_{z'y'}|$ represent the accommodation in the parent grain through basal slip, prismatic slip, and twinning, respectively. It can be seen that the induced deformation gradient by the secondary twins belonging to Group I can be easily accommodated by basal slip and

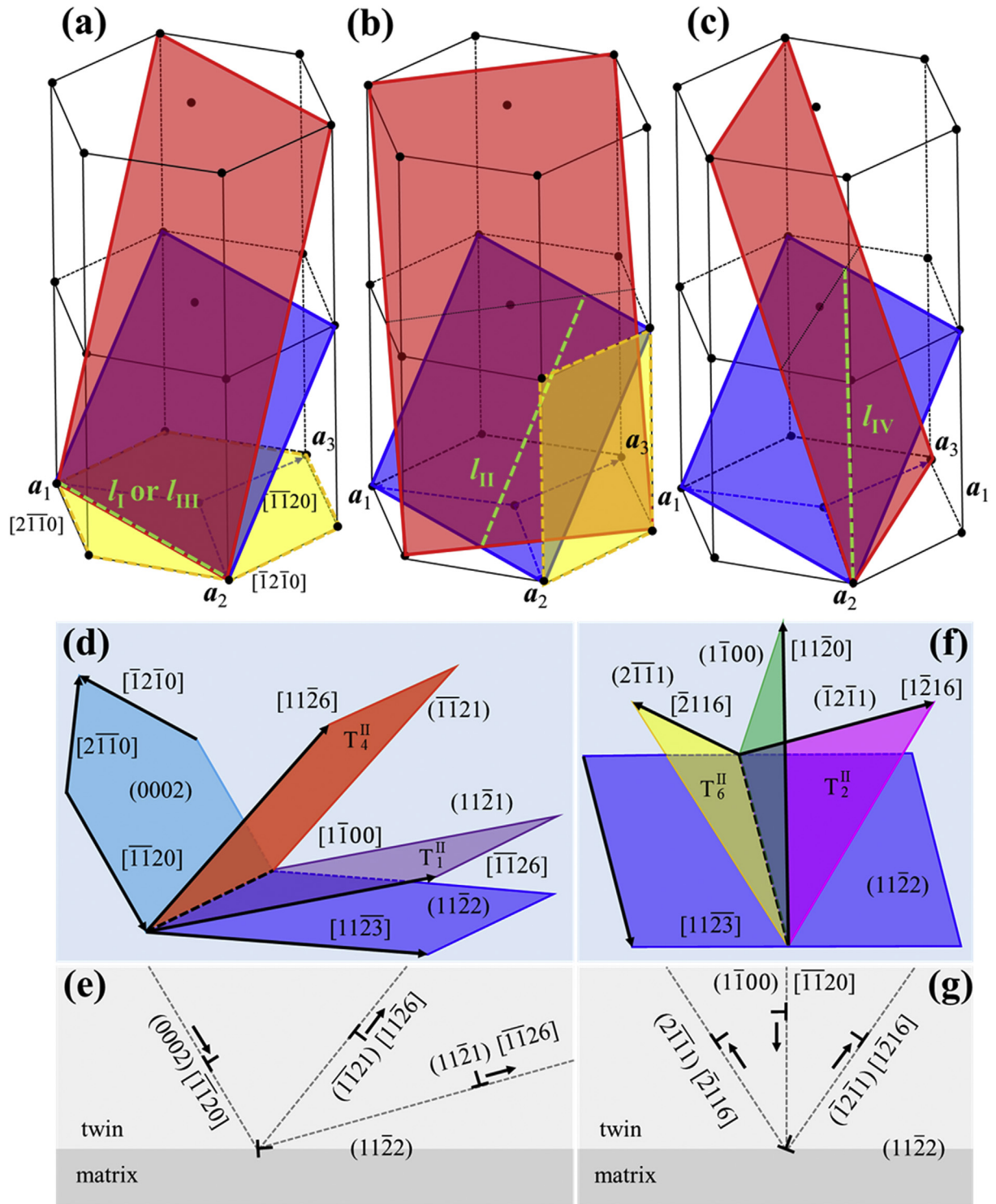


Fig. 7. Crystallography of four groups of $\{11\bar{2}\} \rightarrow \{11\bar{1}\}$ double twins: (a) Group I and Group III, (b) Group II, (c) Group IV. The blue plane denotes the primary twin C_I^I , and the red plane represents the secondary twin T_j^{II} . The green dashed lines are the intersection between the primary twin C_I^I and the secondary twin T_j^{II} . The yellow planes outline the common slip planes where the intersection lines of the primary and secondary twin planes lie. (d) Illustration of the common line among a basal plane, a T_1^{II} twinning plane, a T_4^{II} twinning plane and a primary twin plane along $[1\bar{1}00]$. (e) Schematic of reaction of a basal $\langle a \rangle$ dislocation at a primary twin boundary into a T_1^{II} or a T_4^{II} twin dislocation plus a residual dislocation. (f) Illustration of the common line between a prismatic plane, a T_2^{II} twinning plane, a T_6^{II} twinning plane and a primary twin plane along $[11\bar{2}3]$. (g) Schematic of reaction of a prismatic $\langle a \rangle$ dislocation at a primary twin boundary into a T_2^{II} or a T_6^{II} twin dislocation plus a residual dislocation. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

Table 2
Nucleation of $\{11\bar{2}1\}$ secondary twin variants in the $\{11\bar{2}2\}[\bar{1}1\bar{2}3]$ primary twin.

NDD Analysis							
Secondary twin	b_t	b_l	b_r^m	$ b_l ^2$	x	$ xb_l ^2$	$ b_r^m ^2$
$(11\bar{2}1)[\bar{1}1\bar{2}6] T_1^{\text{II}}$	$\lambda[\bar{1}1\bar{2}6]$	b_1	$[0.63 \ 0.36 \ 0.27 \ 0.18]$	8.70	1	0.78	8.70
		b_2	$[0.36 \ 0.64 \ 0.27 \ 0.18]$	8.70	1	0.78	8.70
		b_3	$[-0.31 \ 1 \ 2 \ 0.6]$	8.70	1	0.78	7.92
$(\bar{1}2\bar{1}1)[1\bar{2}16] T_2^{\text{II}}$	$\lambda[1\bar{2}16]$	b_3	$-0.09 [4 \ 3 \ 7 \ 2]$	8.70	1	0.78	8.70
$(\bar{1}\bar{1}21)[11\bar{2}6] T_4^{\text{II}}$	$\lambda[11\bar{2}6]$	b_1	$-0.09 [\bar{7} \ 4 \ 3 \ 2]$	8.70	1	0.78	8.70
		b_2	$-0.09 [4 \ 7 \ 3 \ 2]$	8.70	1	0.78	8.70
		b_3	$-0.18 [2 \ 2 \ 4 \ 1]$	8.70	1	0.78	11.06
$(2\bar{1}\bar{1}1)[\bar{2}116] T_6^{\text{II}}$	$\lambda[\bar{2}116]$	b_3	$-0.09 [3 \ 4 \ 7 \ 2]$	8.70	1	0.78	8.70
DGA and m-DGA Analysis							
Secondary twin	Group	D_{ST} in the crystal reference frame of the parent grain			D_{ST} in the twinning reference frame of the primary twin		
$(11\bar{2}1)[\bar{1}1\bar{2}6] T_1^{\text{II}}$	I	$\begin{bmatrix} 0.0655 & 0.0378 & 0.0809 \\ 0.0378 & 0.0218 & 0.0467 \\ -0.0707 & -0.0408 & -0.0873 \end{bmatrix}$			$\begin{bmatrix} -0.1549 & 0 & -0.5893 \\ 0 & 0 & 0 \\ 0.0407 & 0 & 0.1549 \end{bmatrix}$		
$(\bar{1}2\bar{1}1)[1\bar{2}16] T_2^{\text{II}}$	II	$\begin{bmatrix} 0.0808 & -0.0466 & 0.0294 \\ 0.0901 & -0.052 & 0.0328 \\ -0.0791 & 0.0457 & -0.0288 \end{bmatrix}$			$\begin{bmatrix} 0 & -0.4614 & -0.3150 \\ 0 & -0.1354 & -0.0924 \\ 0 & 0.1984 & 0.1354 \end{bmatrix}$		
$(\bar{2}111)[2\bar{1}\bar{1}6] T_3^{\text{II}}$	IV	$\begin{bmatrix} 0.0551 & -0.0802 & -0.0685 \\ 0.0565 & -0.0822 & -0.0674 \\ -0.0227 & 0.033 & 0.0271 \end{bmatrix}$			$\begin{bmatrix} 0.2328 & -0.3783 & 0.1115 \\ 0.0833 & -0.1354 & 0.0399 \\ -0.2035 & 0.3307 & -0.0974 \end{bmatrix}$		
$(\bar{1}\bar{1}21)[11\bar{2}6] T_4^{\text{II}}$	III	$\begin{bmatrix} -0.0183 & -0.0106 & -0.1485 \\ -0.0106 & -0.0061 & -0.0858 \\ 0.003 & 0.0017 & 0.0244 \end{bmatrix}$			$\begin{bmatrix} 0.3108 & 0 & 0.2636 \\ 0 & 0 & 0 \\ -0.3664 & 0 & -0.3108 \end{bmatrix}$		
$(1\bar{2}11)[\bar{1}2\bar{1}6] T_5^{\text{II}}$	IV	$\begin{bmatrix} -0.0581 & 0.1219 & -0.0912 \\ -0.0148 & 0.0311 & -0.0232 \\ 0.0173 & -0.0362 & 0.0271 \end{bmatrix}$			$\begin{bmatrix} 0.2328 & 0.3783 & 0.1115 \\ -0.0833 & -0.1354 & -0.0399 \\ -0.2035 & -0.3307 & -0.0974 \end{bmatrix}$		
$(2\bar{1}\bar{1}1)[\bar{2}116] T_6^{\text{II}}$	II	$\begin{bmatrix} 0 & 0.1367 & 0.0431 \\ 0 & 0.0288 & 0.0091 \\ 0 & -0.0914 & -0.0288 \end{bmatrix}$			$\begin{bmatrix} 0 & 0.4616 & -0.3150 \\ 0 & -0.1354 & 0.0924 \\ 0 & -0.1984 & 0.1354 \end{bmatrix}$		

twinning with accommodation components of **[0.1276, 0.0756, 0.1115]**. For Group II, the accommodation is realized through prismatic slip and twinning corresponding to **[0.0622, 0.1367, 0.1248]** for $(\bar{1}2\bar{1}1)[1\bar{2}16]$ and **[0.0522, 0.1367, 0.0914]** for $(2\bar{1}\bar{1}1)[\bar{2}116]$. For Group III, basal slip is the main accommodation system according to **[0.2343, 0.0212, 0.0043]**. For Group IV, much accommodation is required through basal and prismatic slip according to **[0.1359, 0.1367, 0.0557]** for $(\bar{2}111)[2\bar{1}\bar{1}6]$ and **[0.1144, 0.1367, 0.0535]** for $(1\bar{2}11)[\bar{1}2\bar{1}6]$. Therefore, the original DGA cannot distinguish the prevailed secondary twin variant over other variants.

When the DGA criterion is modified to focus on minimizing the resultant plastic deformation associated with the double twinning (m-DGA), the displacement gradient created by the six potential secondary twin variants is transformed into the twinning reference frame of the primary twin. The coordinate is: 1-direction || the shear direction of the primary twin, 3-direction || the normal of the twinning plane, and 2-direction is the cross product of 3-direction and 1-direction. Therefore, e_{13} component of the rotated tensor indicates the displacement gradient that happens along the shear direction in the twinning plane of the primary twin, which reflects the ability of a secondary twin variant to diminish the strain created by the primary twin. A larger magnitude of the e_{13} component (light blue) will more effectively diminish the resultant strain. The results are listed in Table 2. It can be seen that Group I can maximally diminish the strain induced by the $\{11\bar{2}2\}[\bar{1}1\bar{2}3]$ primary twin. The m-DGA predicts a clear order in terms of activating the secondary twin variants, Group I ($C_i^I \rightarrow T_i^{\text{II}}$) > Group II ($C_i^I \rightarrow T_{i+1}^{\text{II}}$, $C_i^I \rightarrow T_{i+5}^{\text{II}}$) > Group IV ($C_i^I \rightarrow T_{i+2}^{\text{II}}$, $C_i^I \rightarrow T_{i+4}^{\text{II}}$) > Group III ($C_i^I \rightarrow T_{i+3}^{\text{II}}$). Such predictions are consistent with the a-SF and NDD analysis, more importantly, agree with the experimental observation.

3.2. $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins

3.2.1. Apparent SF (a-SF) analysis

Fig. 8 shows IPFs of the primary loading domain associated with the primary twin T_1^{II} . The SFs of the six $\{11\bar{2}4\}$ twin variants C_j^{II} ($j = 1 \dots 6$) inside the primary twin T_1^{II} are plotted in Fig. 8b–g. We integrated the area of the domains where the SFs are larger than 0.2. The ratios of the domain areas for four Groups are 1.79:2.27:1:1.66. The average SFs for four Groups are 0.40 for Group I, 0.38 for Group II, 0.34 for Group III, and 0.36 for Group IV. Secondary twin variants in Group II are predicted to be activated with a slightly higher possibility. However, experimental observations do not support such a prediction.

3.2.2. Nucleation via dislocation dissociation (NDD)

Fig. 9 shows the geometric characteristics of four groups of $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins. The intersection line between the primary twin plane and the secondary twin variants belonging to Group III and Group I is parallel to $[\bar{1}100]$, which is parallel to the intersection line between the basal plane and the primary twin plane, as depicted in Fig. 9a. Correspondingly, basal $\langle a \rangle$ dislocation can be dissociated to produce Group I and Group III secondary twins. For Group II, the common axis between the primary and secondary twin plane aligns parallel to $\langle 4513 \rangle$, which lies in pyramidal $\{\bar{1}011\}$ planes and has an angle of 32.6° with the intersection line between the primary twin plane and the basal plane. Dissociation of dislocations on $\{\bar{1}011\}$ pyramidal planes is required to activate Group II secondary twins, which is relatively difficult. For Group IV, the common lines between the primary and secondary twin planes are along $\langle 3413 \rangle$, which do not lie on any usual slip planes, and have an angle of 39.5° with the intersection line

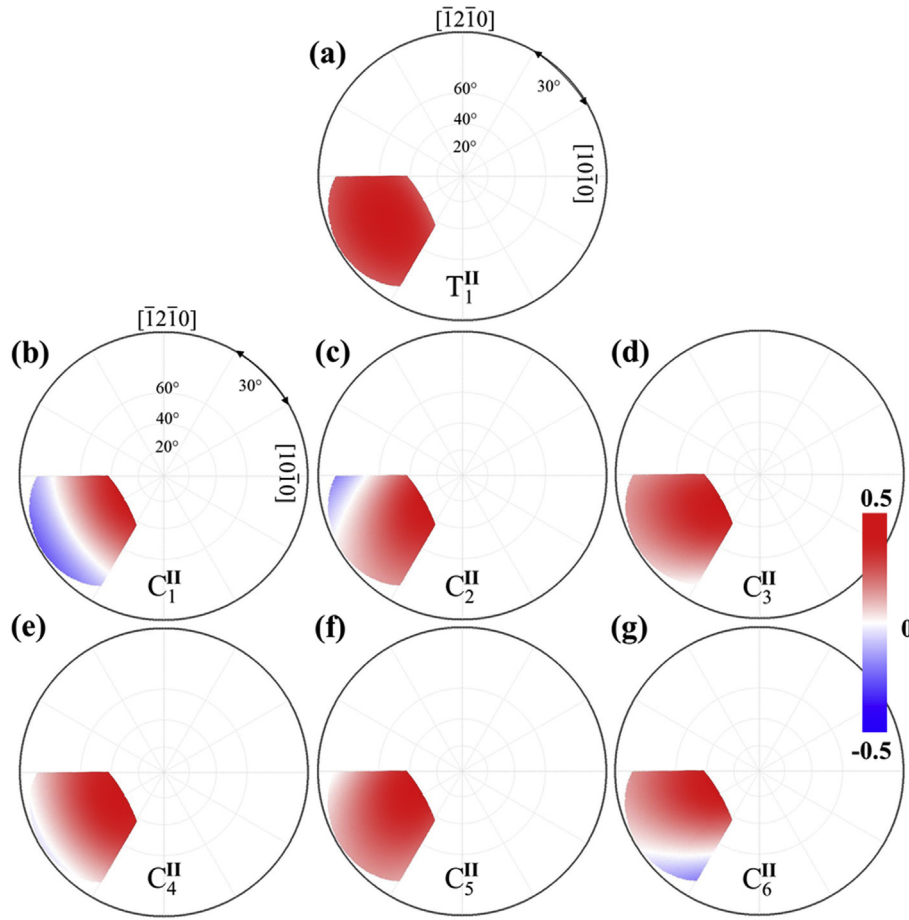


Fig. 8. (a) Stress domain where the SF of a T_1^{II} primary twin variant is positive and the largest among six possibilities. Under this stress domain, the SF of: (b) C_1^{II} , (c) C_2^{II} , (d) C_3^{II} , (e) C_4^{II} , (f) C_5^{II} , (g) C_6^{II} secondary twin variants.

between the primary twin plane and the basal plane. As a result, dislocation dissociation is difficult unless dislocations can easily climb on the primary twin plane. Nucleation based on dissociation of basal $\langle a \rangle$ dislocation is more likely for the secondary twin variants belonging to Group I and Group III. When a $(0002)\langle 11\bar{2}0 \rangle$ dislocation approaches the primary twin boundary, it can be dissociated into x twinning dislocations $xb_t^{\{11\bar{2}4\}}$ associated with Group III (C_1^{II}) and Group I (C_4^{II}) secondary twins and residual dislocations b_r^m , where the subscript m represents the m th ($m = 1, 4$) secondary twin.

$$b_i \Rightarrow b_r^m + xb_t^{\{11\bar{2}4\}} \quad (3)$$

Table 3 summarizes the dissociation of a basal $\langle a \rangle$ dislocation into secondary twins belonging to Groups I, II, and III. For C_1^{II} (Group III), a $(0002)\langle 11\bar{2}0 \rangle$ dislocation, a b_3 can be maximally dissociated into three secondary twinning dislocations $3b_t^{\{11\bar{2}4\}}$ ($x = 3$), and a residual dislocation b_r^m ($m = 1$). The dissociation process is energetically favorable according to Frank's law ($|b_3|^2 > |b_r^1|^2 + |3b_t^{\{11\bar{2}4\}}|^2$). When a b_1 or b_2 dislocation acts as the source for dissociation, the production of only one C_1^{II} secondary twinning dislocation $b_t^{\{11\bar{2}4\}}$ ($x = 1$) is energetically favored as listed in Table 3. However, for C_4^{II} (Group III), a $(0002)\langle 11\bar{2}0 \rangle$ dislocation, b_i dissociates into one secondary twinning dislocation $b_t^{\{11\bar{2}4\}}$ ($x = 1$) and a residual one b_r^m ($m = 4$). This dissociation process is energetically unfavorable due to the increase in the line energy ($|b_i|^2 < |b_r^4|^2 + |b_t^{\{11\bar{2}4\}}|^2$). The magnitude of Burgers vector

of $\{11\bar{2}4\}$ twinning dislocation is equal to $\frac{c^2 - 2a^2}{\sqrt{4a^2 + c^2}}$, 0.6 nm for α -Ti. According to the theory, the energetically favorable dissociation is associated with the formation of secondary twin variant C_1^{II} (Group III), and the other twin variants are difficult to be activated according to dislocation dissociation. The NDD prediction agrees with the experimental observations.

3.2.3. Displacement gradient accommodation analysis (DGA and m-DGA)

For a DGA analysis, the displacement gradients induced by six possible secondary $\{11\bar{2}2\}$ twins (D_{ST}) in $(11\bar{2}1)[\bar{1}1\bar{2}6]$ primary twin (T_1^{II}) can be expressed in the crystal frame of the parent grain as listed in the first column of Table 3. Much accommodation for Group III is made through twinning according to the accommodation components [0.0195, 0.0414, **0.2196**]. For Groups I and II, the accommodation is mainly realized by basal slip according to the accommodation components [0.2382, 0.0092, 0.0009] for Group I and [0.1684, 0.0869, 0.0402] for Group II ($\bar{2}114$)[$\bar{4}22\bar{3}$] and [0.1685, 0.0957, 0.0331] for Group II ($\bar{1}214$)[$\bar{2}42\bar{3}$]. For Group IV, the accommodation is conducted through twinning according to the accommodation components [0.0741, 0.0992, **0.1275**] for $(\bar{1}2\bar{1}4)$ [$\bar{2}42\bar{3}$] and [0.0762, 0.0868, **0.1488**] for $(\bar{2}1\bar{1}4)$ [$\bar{4}22\bar{3}$]. The results suggest that Group I and II are prevailed over Groups III and IV because plastic relaxation via twinning is more difficult than that via basal slip. Such a prediction is different from that based on the NDD and the experimental results.

According the m-DGA analysis, we transformed the

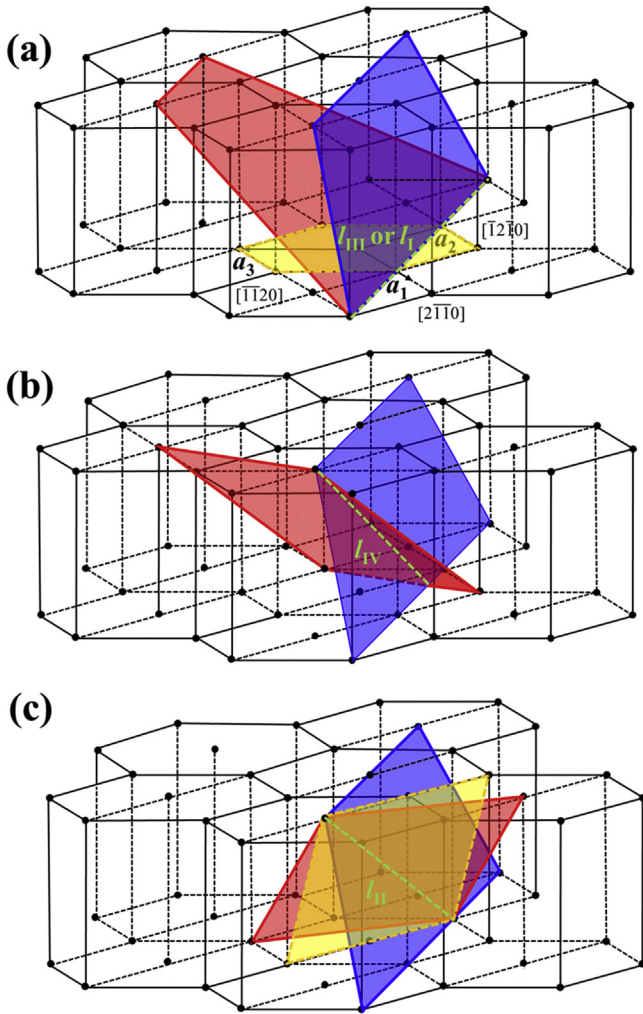


Fig. 9. Crystallography of four groups of $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins. The blue plane denotes the primary twin T_1^I , and the red plane represents the secondary twin C_1^I . The green dashed line indicates the intersection between the primary twin T_1^I and the secondary twin C_1^I . The yellow plane outlines the common slip plane where the intersection line of the primary and secondary twin plane lies. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

displacement gradient induced by six potential secondary twins into the twinning reference frame of the primary twin. The results listed in the second column in Table 3 indicate that Group III can, to a great extent, accommodate the strain induced by the $(11\bar{2}1)[11\bar{2}6]$ primary twin. Group IV secondary twins can also partially accommodate the strain associated with the primary twin. Other secondary twins do not reduce the resultant deformation produced by the double twins. The prediction according to the m-DGA is consistent with that based on the NDD analysis and experimental observations.

3.3. $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ non-family double twins

3.3.1. Apparent SF (α -SF) analysis

Fig. 10a shows IPFs of the primary loading domain associated with the primary twin C_1^I . Under this primary loading, the SFs associated with six potential $\{10\bar{1}2\}$ twin variants T_j^I ($j = 1 \dots 6$) are calculated inside the primary twin C_1^I and plotted in Fig. 10b–g. The results show that secondary twin variants (T_1^I and T_6^I) exhibit very similar SFs due to the approximate 90° angle between their

twin planes. The SFs of T_3^I and T_6^I secondary twin variants are very low at most loading cases, while other four secondary twin variants hold high SFs. When the SFs of secondary twins larger than 0.2 are selected, the area of the loading domain is almost zero for T_3^I and T_6^I , while the area of this loading domain and the average SF in this loading area are almost identical for other four secondary twin variants. Therefore, α -SF analysis is unable to distinguish the four secondary twin variants.

3.3.2. Nucleation via dislocation dissociation (NDD)

As listed in Table 4, the intersection line between the primary twinning plane and the secondary twinning planes in Group II and Group III is along $\langle 4223 \rangle$, which lies in prismatic $\{10\bar{1}0\}$ plane and has an angle of 47.3° with the intersection between the primary twin plane and the basal plane. This provides a geometric favorite for the dissociation of prismatic $\langle a \rangle$ dislocation into the secondary twin dislocations belonging to the two groups. However, the intersection line between Group I secondary twinning plane and the primary twinning plane does not lie on any usual slip planes. The dissociation mechanism for nucleating secondary twins in Group I requires dislocation climb on the primary twin plane. A prismatic $\langle a \rangle$ dislocation can dissociate into x secondary $\{10\bar{1}2\}$ twinning dislocation $xb_t^{(10\bar{1}2)}$ and a residual dislocation b_r^m ($m = 1, 2, 3, 6$).

$$b_i \Rightarrow b_r^m + xb_t^{(10\bar{1}2)} \quad (4)$$

As listed in Table 4, a b_1 or b_2 dislocation can dissociate into three secondary $\{10\bar{1}2\}$ twinning dislocations ($x = 3$) associated with T_1^I and T_2^I variants (Group II), respectively, which is energetically favorable according to Frank's law. However, a b_2 or b_1 dislocation dissociates into only one secondary $\{10\bar{1}2\}$ twinning dislocation ($x = 1$) associated with T_3^I and T_6^I variant (Group III), respectively, and the dissociation is energetically unfavourable due to the increase in the elastic energy after dissociation. Therefore, the dislocation dissociation analysis shows a preference of $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ double twins, i.e., the $C_1^I \rightarrow T_1^I$ or T_{i+1}^I double twins (Group II).

3.3.3. Displacement gradient accommodation analysis (DGA and m-DGA)

According to the DGA analysis, the displacement gradients induced by the six possible secondary $\{10\bar{1}2\}$ twins in $(11\bar{2}2)[11\bar{2}3]$ primary twin (C_1^I) can be expressed in the crystal frame of the parent grain and are listed in the first column of Table 4 for a DGA analysis. It can be seen that for Group I, the displacement gradient induced by the secondary twin can be mainly accommodated through basal slip according to the accommodation components $[0.199, 0.0789, 0.0304]$ for $(\bar{1}0\bar{1}2)[10\bar{1}1]$ and $[0.1591, 0.0871, 0.0342]$ for $(0\bar{1}12)[01\bar{1}1]$. For Group II, twinning is the main accommodation carrier according to the accommodation components $[0.0229, 0.0789, 0.2091]$ for $(10\bar{1}2)[\bar{1}011]$ and $[0.0216, 0.085, 0.1715]$ for $(01\bar{1}2)[0\bar{1}11]$. For Group III, the accommodation can be achieved through prismatic slip in the parent according to the accommodation components $[0.0614, 0.1578, 0.0596]$ for $(\bar{1}102)[\bar{1}101]$ and $[0.0685, 0.1578, 0.0591]$ for $(1\bar{1}02)[\bar{1}101]$. According to the DGA analysis, Group III twins are the preferred secondary twins since prismatic slip is the easiest slip to activate in α -titanium. Such a prediction is inconsistent with the experimental observation and is different from that obtained from the NDD.

Displacement gradients created by the six potential double twins were transformed into the twinning reference frames of the primary twin for a m-DGA analysis. The results listed in the secondary column in Table 4 indicate that Group II can, to a great extent, accommodate the strain induced by the $(11\bar{2}2)[11\bar{2}3]$

Table 3Nucleation of $\{11\bar{2}4\}$ secondary twin variants in the $(11\bar{2}1)[\bar{1}\bar{1}26]$ primary twin.

NDD Analysis							
Secondary twin	b_t	b_i	b_r^m	$ b_i ^2$	x	$ xb_t ^2$	$ b_r^m ^2$
$(11\bar{2}4)[22\bar{4}3] C_1^{\text{II}}$	$\lambda[22\bar{4}3]$	b_1	[0.61 0.39 $\overline{0.22}$ 0.08]	8.70	1	0.36	7.68
		b_2	[$\overline{0.39}$ 0.61 $\overline{0.22}$ 0.08]	8.70	1	0.36	7.68
		b_3	[0.17 0.17 $\overline{0.34}$ 0.24]	8.70	3	3.24	3.62
$(\bar{1}\bar{1}24)[\bar{2}\bar{2}4\bar{3}] C_4^{\text{II}}$	$\lambda[\bar{2}\bar{2}4\bar{3}]$	b_1	[0.72 0.28 0.44 0.08]	8.70	1	0.36	10.45
		b_2	[$\overline{0.28}$ 0.72 $\overline{0.44}$ 0.08]	8.70	1	0.36	10.45
		b_3	[0.39 0.39 $\overline{0.78}$ 0.08]	8.70	1	0.36	11.84
DGA and m-DGA Analysis							
Secondary twin	Group	D _{ST} in the crystal reference frame of the parent grain			D _{ST} in the twinning reference frame of the primary twin		
$(11\bar{2}4)[22\bar{4}3] C_1^{\text{II}}$	III	$\begin{bmatrix} 0.0359 & 0.0207 & -0.0124 \\ 0.0207 & 0.012 & -0.0071 \\ 0.1392 & 0.0804 & -0.0479 \end{bmatrix}$			$\begin{bmatrix} -0.1021 & 0 & -0.1509 \\ 0 & 0 & 0 \\ -0.0691 & 0 & 0.1021 \end{bmatrix}$		
$(\bar{1}\bar{2}14)[\bar{2}4\bar{2}\bar{3}] C_2^{\text{II}}$	IV	$\begin{bmatrix} 0.0474 & -0.0062 & -0.025 \\ -0.093 & 0.0121 & 0.0491 \\ 0.1128 & -0.0147 & -0.0596 \end{bmatrix}$			$\begin{bmatrix} -0.1022 & 0.0842 & -0.0831 \\ -0.0976 & 0.0803 & -0.0794 \\ 0.0269 & -0.0221 & 0.0219 \end{bmatrix}$		
$(\bar{2}114)[\bar{4}2\bar{2}\bar{3}] C_3^{\text{II}}$	II	$\begin{bmatrix} -0.0189 & 0.016 & 0.0355 \\ -0.0709 & 0.0598 & 0.1329 \\ 0.0218 & -0.0184 & -0.0409 \end{bmatrix}$			$\begin{bmatrix} -0.0879 & 0.0563 & 0.0064 \\ -0.1254 & 0.0803 & 0.0091 \\ -0.1036 & 0.0664 & 0.0075 \end{bmatrix}$		
$(\bar{1}\bar{1}24)[\bar{2}\bar{2}4\bar{3}] C_4^{\text{II}}$	I	$\begin{bmatrix} 0.0079 & 0.0046 & 0.151 \\ 0.0046 & 0.0026 & 0.0872 \\ -0.0006 & -0.0003 & -0.0106 \end{bmatrix}$			$\begin{bmatrix} -0.0735 & 0 & 0.0281 \\ 0 & 0 & 0 \\ -0.1919 & 0 & 0.0735 \end{bmatrix}$		
$(1\bar{1}24)[\bar{2}4\bar{2}\bar{3}] C_5^{\text{II}}$	II	$\begin{bmatrix} 0.0164 & -0.0913 & 0.1328 \\ -0.0044 & 0.0245 & -0.0357 \\ -0.005 & 0.0281 & -0.0409 \end{bmatrix}$			$\begin{bmatrix} -0.0879 & -0.0563 & 0.0064 \\ 0.1254 & 0.0803 & -0.0091 \\ -0.1036 & -0.0664 & 0.0075 \end{bmatrix}$		
$(2\bar{1}\bar{1}4)[\bar{4}2\bar{2}\bar{3}] C_6^{\text{II}}$	IV	$\begin{bmatrix} -0.022 & -0.0529 & 0.03 \\ 0.0339 & 0.0816 & -0.0462 \\ 0.0437 & 0.1051 & -0.0596 \end{bmatrix}$			$\begin{bmatrix} -0.1022 & -0.0842 & -0.0831 \\ 0.0976 & 0.0803 & 0.0794 \\ 0.0269 & 0.0221 & 0.0219 \end{bmatrix}$		

primary twin. In contrast, Group III double twins can accommodate little strain induced by the primary twin, and Group I cannot accommodate the strain. The m-DGA analysis is consistent with experimental observations and the prediction by the NDD.

3.4. $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ non-family double twins

3.4.1. Apparent SF (a-SF) analysis

Fig. 11 shows the primary loading domain associated with the primary twin C_1^{II} and six potential $\{10\bar{1}2\}$ twin variants T_j^{I} ($j = 1 \dots 6$) into an IPF. Except for the difference in the primary loading domain from that in Fig. 10, the conclusion is the same as that in the case of $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ double twins. In other words, a-SF cannot distinguish the preference of secondary twin variants.

3.4.2. Nucleation via dislocation dissociation (NDD)

The crystallographic feature of three groups of $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins is illustrated in Fig. 12 when $(11\bar{2}4)[22\bar{4}3] (C_1^{\text{II}})$ is the primary twin. The intersection line between two secondary twin variants belonging to Group I (T_1^{I} and T_2^{I}) is parallel to $[22\bar{4}3]$, which lies in prismatic $(\bar{1}\bar{1}00)$ plane and is perpendicular to the intersection line between the primary twin plane and basal plane as indicated in Fig. 12a. Thus, basal $\langle a \rangle$ dislocations cannot dissociate into twinning dislocations associated with the two secondary twins. Instead, $\langle a \rangle$ dislocation on the prismatic plane can dissociate to produce secondary twins. For T_3^{I} and T_6^{I} in Group II, the common axis between the primary and secondary twin plane aligns parallel to $[1\bar{5}43]$ and $[5143]$, which lies on pyramidal planes $(10\bar{1}1)$ and $(01\bar{1}1)$ as indicated in Fig. 12b. However, dislocations on pyramidal planes are rarely activated. In addition, the angle between the common axis and the intersection line between the basal plane and the primary twin plane is 55.8° . The dissociation based on basal $\langle a \rangle$ dislocations requires the rotation of basal $\langle a \rangle$ dislocation line for 55.8° via the climb on the twin plane. Thus,

nucleating secondary twins in Group II lacks dislocation sources. For T_4^{I} and T_5^{I} in Group III, the common axis between the primary and secondary twin plane is parallel to $[10\bar{1}4\ 4\ 3]$ and $[14\ 10\ 4\ 3]$ (refer to Fig. 12c), which do not lie on any usual slip planes. However, the angle between the common axis and the intersection line between the basal plane and the primary twin plane is 20.2° . Compared with Group II, there is a less dependence on dislocation climb. Regarding secondary twins T_1^{I} and T_2^{I} in Group I, there is only one $\langle a \rangle$ dislocation with the Burgers vector b_3 that can be dissociated into x secondary $\{10\bar{1}2\}$ twinning dislocation $xb_t^{(10\bar{1}2)}$ and a residual one b_r^m ($m = 1, 2$),

$$b_i = b_r^m + xb_t^{(10\bar{1}2)} \quad (5)$$

As listed in Table 5, b_3 is dissociated into three secondary twinning dislocation $3b_t^{(10\bar{1}2)}$ ($x = 3$). The dissociation is energetically favorable due to the decrease in the elastic energy after dissociation. Therefore, secondary twins in Group I should be prevailed over others, and secondary twins in Group III are likely activated than those in Group II.

3.4.3. Displacement gradient accommodation analysis (DGA and m-DGA)

For a DGA analysis, the displacement gradients induced by six possible secondary $\{10\bar{1}2\}$ twin in its twinning frame can be expressed in the crystal frame of the parent grain as listed in the first column of Table 5. For Group I, the main accommodation system is basal slip according to the accommodation components **[0.1241, 0.095, 0.0757]** for $(10\bar{1}2)[\bar{1}011]$ and **[0.1589, 0.0852, 0.0554]** for $(01\bar{1}2)[0\bar{1}11]$. For Group II, the accommodation can be mainly realized by prismatic slip according to the accommodation components **[0.0342, 0.1705, 0.0319]** for $(\bar{1}\bar{1}02)[\bar{1}\bar{1}01]$ and **[0.0366, 0.1705, 0.0329]** for $(1\bar{1}02)[\bar{1}\bar{1}01]$. For Group III, the accommodation is carried out by twinning according to the accommodation

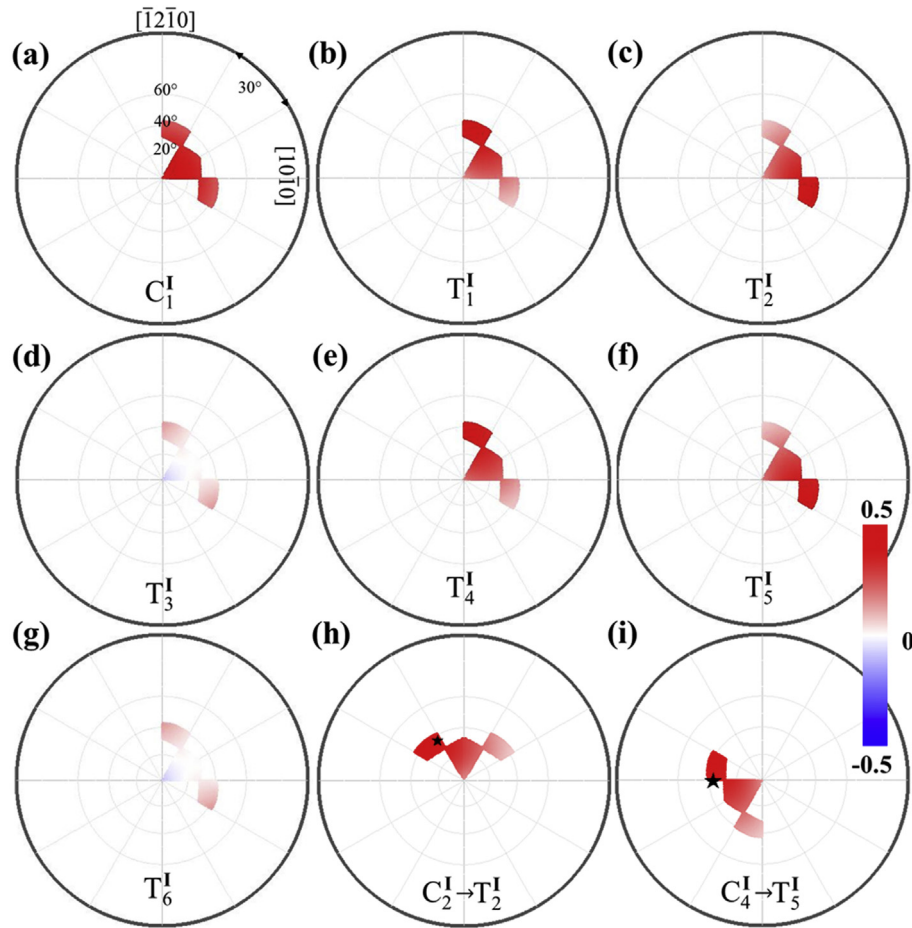


Fig. 10. (a) Stress domain where the SF of a C_1^I primary twin variant is positive and the largest among six possibilities. Under this stress domain, the SF of: (b) T_1^I , (c) T_2^I , (d) T_3^I , (e) T_4^I , (f) T_5^I and (g) T_6^I secondary twin variants. (h) The SFs for secondary twin T_2^I in primary twin C_2^I . The black star indicates the loading direction associated with the grain in Fig. 4a. (i) The SFs for secondary twin T_5^I in primary twin C_4^I . The black star indicates the loading direction associated with the grain in Fig. 4e.

Table 4Nucleation of $\{10\bar{1}2\}$ secondary twin variants in the $(11\bar{2}2)[11\bar{2}3]$ primary twin.

NDD Analysis							
Secondary twin	b_t	b_i	b_r^m	$ b_t ^2$	x	$ xb_t ^2$	$ b_r^m ^2$
$(10\bar{1}2)[\bar{1}011] T_1^I$	$\lambda[\bar{1}011]$	b_1	$[0.41 \ 0.33 \ 0.08 \ 0.26]$	8.70	3	3.24	5.16
$(01\bar{1}2)[0\bar{1}11] T_2^I$	$\lambda[0\bar{1}11]$	b_2	$[0.33 \ 0.41 \ 0.08 \ 0.26]$	8.70	3	3.24	5.16
$(\bar{1}102)[1\bar{1}01] T_3^I$	$\lambda[1\bar{1}01]$	b_2	$[0.42 \ 0.75 \ 0.33 \ 0.09]$	8.70	1	0.36	11.34
$(1\bar{1}02)[\bar{1}101] T_6^I$	$\lambda[\bar{1}101]$	b_1	$[0.75 \ 0.42 \ 0.33 \ 0.09]$	8.70	1	0.36	11.34
DGA and m-DGA Analysis							
Secondary twin	Group	D_{ST} in the crystal reference frame of the parent grain			D_{ST} in the twinning reference frame of the primary twin		
$(10\bar{1}2)[\bar{1}011] T_1^I$	II	$\begin{bmatrix} 0.0518 & 0.0622 & 0.0173 \\ -0.0167 & -0.02 & -0.0056 \\ -0.095 & -0.1141 & -0.0318 \end{bmatrix}$			$\begin{bmatrix} 0.0498 & 0.0539 & -0.1417 \\ -0.0201 & -0.0218 & 0.0573 \\ 0.0098 & 0.0106 & -0.0280 \end{bmatrix}$		
$(01\bar{1}2)[0\bar{1}11] T_2^I$	II	$\begin{bmatrix} 0.0177 & 0.003 & 0.0038 \\ 0.082 & 0.0141 & 0.0178 \\ -0.1463 & -0.0252 & -0.0318 \end{bmatrix}$			$\begin{bmatrix} 0.0498 & -0.0539 & -0.1417 \\ 0.0201 & -0.0218 & -0.0573 \\ 0.0098 & -0.0106 & -0.0280 \end{bmatrix}$		
$(\bar{1}102)[1\bar{1}01] T_3^I$	III	$\begin{bmatrix} 0.0255 & -0.007 & -0.0089 \\ 0.1508 & -0.0417 & -0.0525 \\ -0.0467 & 0.0129 & 0.0163 \end{bmatrix}$			$\begin{bmatrix} 0.0624 & -0.0676 & -0.0393 \\ 0.0805 & -0.0872 & -0.0507 \\ -0.0393 & 0.0426 & 0.0248 \end{bmatrix}$		
$(\bar{1}012)[10\bar{1}1] T_4^I$	I	$\begin{bmatrix} 0.0536 & -0.0088 & -0.0862 \\ 0.0701 & -0.0115 & -0.1128 \\ 0.0261 & -0.0043 & -0.0421 \end{bmatrix}$			$\begin{bmatrix} 0.0379 & -0.0137 & 0.0041 \\ 0.0604 & -0.0218 & 0.0066 \\ -0.1474 & 0.0532 & -0.0161 \end{bmatrix}$		
$(0\bar{1}12)[01\bar{1}1] T_5^I$	I	$\begin{bmatrix} 0.0313 & 0.083 & -0.1408 \\ 0.0041 & 0.0108 & -0.0183 \\ 0.0094 & 0.0248 & -0.0421 \end{bmatrix}$			$\begin{bmatrix} 0.0379 & 0.0137 & 0.0041 \\ -0.0604 & -0.0218 & -0.0066 \\ -0.1474 & -0.0532 & -0.0161 \end{bmatrix}$		
$(1\bar{1}02)[\bar{1}101] T_6^I$	III	$\begin{bmatrix} 0.0373 & 0.1439 & -0.0499 \\ -0.0139 & -0.0539 & 0.0186 \\ -0.0122 & -0.0469 & 0.0163 \end{bmatrix}$			$\begin{bmatrix} 0.0624 & 0.0676 & -0.0393 \\ -0.0805 & -0.0872 & 0.0507 \\ -0.0393 & -0.0426 & 0.0248 \end{bmatrix}$		

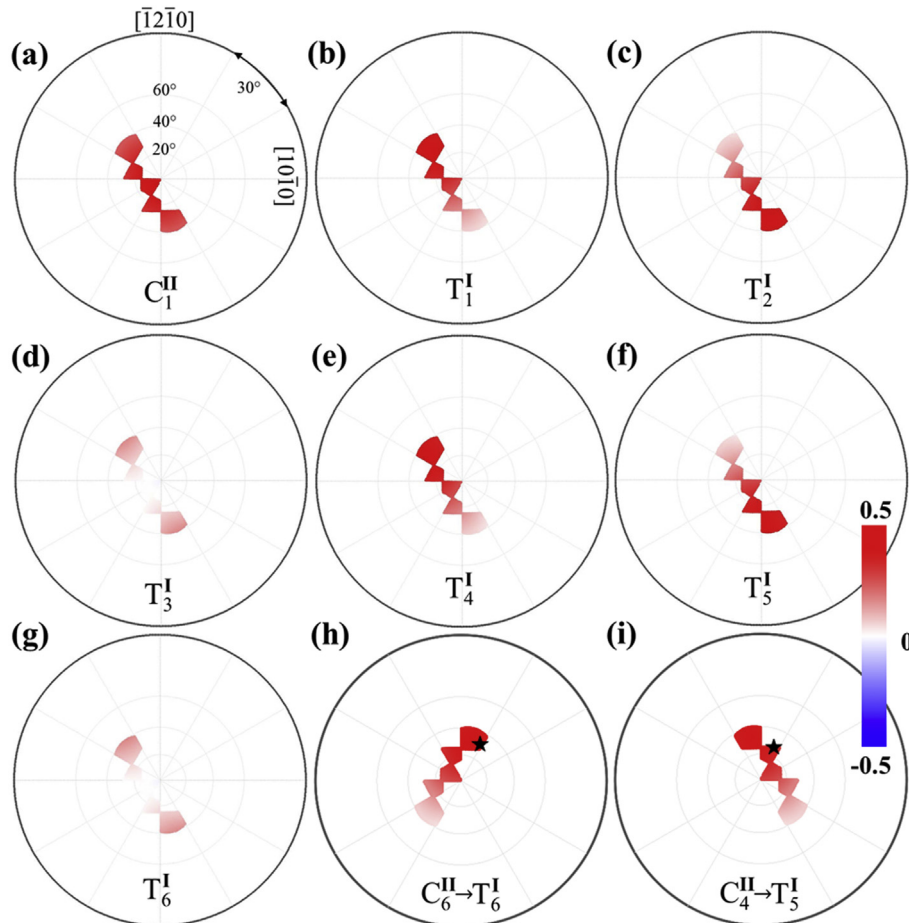


Fig. 11. (a) Stress domain where the SF of a C_1^{II} primary twin variant is positive and the largest among six possibilities. Under this stress domain, the SF of: (b) T_1^I , (c) T_2^I , (d) T_3^I , (e) T_4^I , (f) T_5^I and (g) T_6^I secondary twin variants. (h) The SFs for secondary twin T_6^I in primary twin C_6^{II} . The black star indicates the loading direction associated with the grain in Fig. 5a. (i) The SFs for secondary twin T_5^I in primary twin C_4^{II} . The black star indicates the loading direction associated with the grain in Fig. 5e.

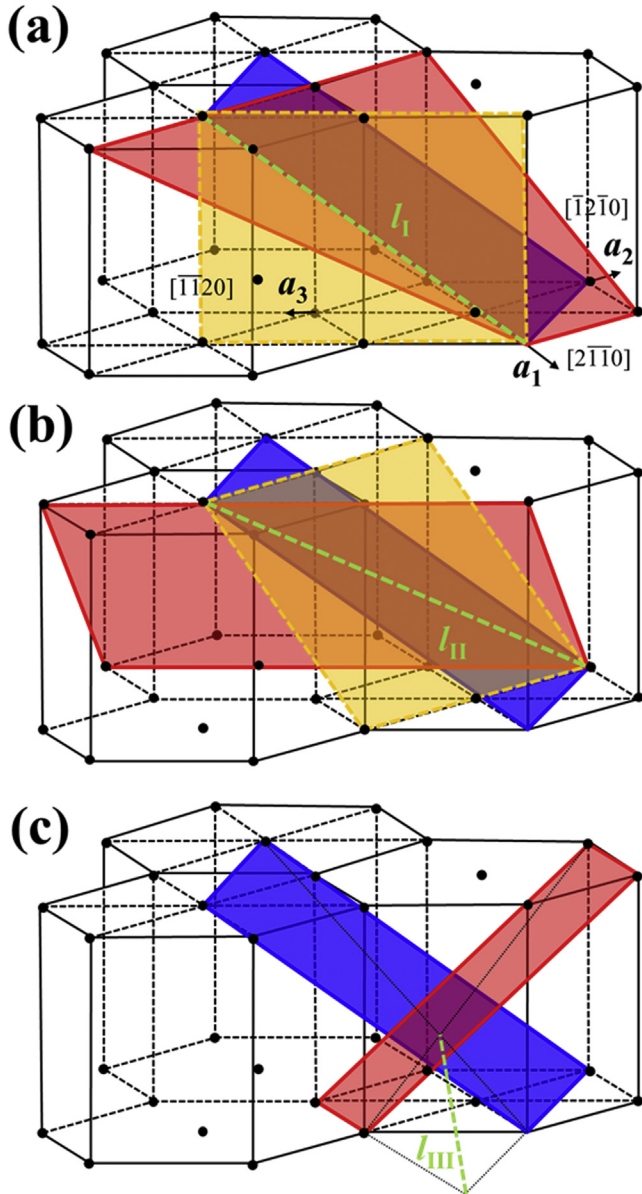


Fig. 12. Crystallography of four groups of $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins. The blue plane denotes the primary twin C_1^I , and the red plane represents the secondary twin T_1^I . The green dashed line indicates the intersection between the primary twin C_1^I and the secondary twin T_1^I . The yellow plane outlines the common slip plane where the intersection line of the primary and secondary twin plane lies. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

components [0.0589, 0.0973, **0.1409**] for $(\bar{1}012)[10\bar{1}1]$ and [0.0487, 0.0852, **0.1744**] for $(0\bar{1}12)[01\bar{1}1]$. Accordingly, Group II variants should be prevailed among three groups because prismatic slip is the easiest slip to activate in α -titanium. Group III variants are unlikely activated because of the accommodation by twinning based on the DGA analysis.

According to the m-DGA analysis, the displacement gradient created by six potential double twins was transformed into the twinning reference frame of the primary twin, and the results are listed in the secondary column in Table 5. Group I can, to the greatest extent, accommodate the strain induced by the $(11\bar{2}4)[22\bar{4}3]$ primary twin. In contrast, Group II double twins can accommodate little strain induced by the primary twin while

Group III cannot accommodate the strain. Therefore, the m-DGA analysis gives a different prediction as does the DGA analysis, but is consistent with the NDD analysis and the experimental observations.

4. Discussion

For co-family double twins, only one double twin variant is selected in the preferred group of the double twins, which can be accurately predicted by using the m-DGA and NDD criteria. However, two double twin variants are possibly activated in the predominant group of the double twins belonging to non-family type, which fails to be distinguished by using the m-DGA and NDD criteria. The a-SF criterion is helpful for determining the competition between two double twin variants in the prevailed group of non-family double twins. In Fig. 4a, $C_2^I \rightarrow T_2^I$ is identified to be a $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ non-family double twin. In C_2^I primary twin, two secondary twins belong to predominant Group (Group II), i.e., T_4^I and T_5^I . According to the a-SF analysis, the loading direction identified for the grain in Fig. 4a, where C_2^I forms is located at the predicted loading domain as shown in Fig. 10h, where the SF of T_2^I secondary twin variant is greater than T_3^I . In Fig. 4e, a detection of $C_4^I \rightarrow T_5^I$ double twin (Group II) is indicated. In C_4^I primary twin, both T_4^I and T_5^I belong to Group II. The a-SF analysis in Fig. 10i indicates that the loading direction identified for the grain is corresponding to a greater SF of T_5^I secondary twin variant than T_4^I .

As to $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ non-family double twins, the criteria based on m-DGA and NDD can fully account for the preference of Group I double twin variants. To determine the selection of two variants in Group I, the a-SF criterion is applied. In Fig. 5a, $C_6^{II} \rightarrow T_1^I$ double twin belonging to Group I is detected. In C_6^{II} primary twin, T_6^I and T_1^I are both Group I double twin variants. According to the a-SF analysis in Fig. 11h, the loading direction identified for the grain where C_6^{II} forms is located at the predicted loading domain, where the SF of T_6^I secondary twin variant is greater than T_1^I . In Fig. 5e, the detection of $C_4^{II} \rightarrow T_5^I$ double twin (Group I) is determined. In C_4^{II} primary twin, there are two secondary twins T_4^I and T_5^I belonging to Group I. The a-SF analysis in Fig. 11i indicates that the loading direction identified for the grain where C_4^{II} forms is located at the loading domain, where the SF of T_5^I secondary twin variant is greater than T_4^I .

Regarding the selection criteria, the a-SF clearly shows that the SFs for six twin variants are positive in most of the primary loading domain as long as the primary twin is subjected to a positive resolved shear stress, although it cannot predict the preferred secondary twin variants. The DGA focusing on the maximum relaxation of the strain created by the primary twins in the vicinity of the primary twin is not effective to predict the preferred secondary twin variant, while the m-DGA with the focus on minimizing the resultant plastic deformation due to the double twinning successfully predicts the preferred secondary twin variant. The success of the m-DGA prediction implies that the formation of double twin may be related to the minimization of the local shear associated with the primary twin in the matrix. Both m-DGA and NDD predict results that agree with the experimental observations. These two theories speculate that nucleation of a secondary twin requires nucleation sources at the primary twin boundary and the accumulation and reaction of gliding dislocations at the primary twin boundary provides such nucleation sources. Such speculations reflect in general the physics of twinning.

5. Conclusions

Four types of double twins observed in deformed Ti specimens are classified into two families according to their zone axes: co-

Table 5Nucleation of $\{10\bar{1}2\}$ secondary twin variants in the $(11\bar{2}4)[22\bar{4}3]$ primary twin.

NDD Analysis							
Secondary twin	b_t	b_i	b_r^m	$ b_i ^2$	x	$ xb_r ^2$	$ b_r^m ^2$
$(10\bar{1}2)[\bar{1}011] T_1^I$	$\lambda[\bar{1}011]$	b_3	$[0.07 \ 0.33 \ 0.40 \ 0.26]$	8.70	3	3.24	5.16
$(01\bar{1}2)[0\bar{1}11] T_2^I$	$\lambda[0\bar{1}11]$	b_3	$[0.33 \ 0.07 \ 0.40 \ 0.26]$	8.70	3	3.24	5.16
DGA and m-DGA Analysis							
Secondary twin	Group	D_{ST} in the crystal reference frame of the parent grain			D_{ST} in the twinning reference frame of the primary twin		
$(10\bar{1}2)[\bar{1}011] T_1^I$	I	$\begin{bmatrix} 0.052 & 0.0901 & 0.1135 \\ 0.0049 & 0.0084 & 0.0106 \\ -0.0277 & -0.048 & -0.0604 \end{bmatrix}$			$\begin{bmatrix} 0 & 0.0544 & -0.1516 \\ 0 & -0.0218 & 0.0607 \\ 0 & -0.0078 & 0.0218 \end{bmatrix}$		
$(01\bar{1}2)[0\bar{1}11] T_2^I$	I	$\begin{bmatrix} 0.0604 & 0 & 0.0659 \\ 0.0852 & 0 & 0.093 \\ -0.0554 & 0 & -0.0604 \end{bmatrix}$			$\begin{bmatrix} 0 & -0.0544 & -0.1516 \\ 0 & -0.0218 & -0.0607 \\ 0 & 0.0078 & 0.0218 \end{bmatrix}$		
$(\bar{1}102)[1\bar{1}01] T_3^I$	II	$\begin{bmatrix} 0.0338 & -0.008 & 0.0059 \\ 0.1625 & -0.0383 & 0.0283 \\ 0.0258 & -0.0061 & 0.0045 \end{bmatrix}$			$\begin{bmatrix} 0.0337 & -0.0497 & -0.0424 \\ 0.0591 & -0.0872 & -0.0745 \\ -0.0424 & 0.0626 & 0.0535 \end{bmatrix}$		
$(\bar{1}012)[10\bar{1}1] T_4^I$	III	$\begin{bmatrix} 0.0412 & 0.006 & -0.0183 \\ 0.0913 & 0.0134 & -0.0406 \\ 0.1229 & 0.0180 & -0.0546 \end{bmatrix}$			$\begin{bmatrix} -0.0128 & 0.0047 & 0.0030 \\ 0.0591 & -0.0218 & -0.0138 \\ -0.1486 & 0.0548 & 0.0346 \end{bmatrix}$		
$(0\bar{1}12)[01\bar{1}1] T_5^I$	III	$\begin{bmatrix} 0.0625 & 0.079 & -0.0443 \\ -0.0062 & -0.0079 & 0.0044 \\ 0.077 & 0.0974 & -0.0546 \end{bmatrix}$			$\begin{bmatrix} -0.0128 & -0.0047 & 0.0030 \\ -0.0591 & -0.0218 & 0.0138 \\ -0.1486 & -0.0548 & 0.0346 \end{bmatrix}$		
$(1\bar{1}02)[\bar{1}101] T_6^I$	II	$\begin{bmatrix} 0.0466 & 0.1551 & 0.0275 \\ -0.0154 & -0.0511 & -0.0091 \\ 0.0076 & 0.0253 & 0.0045 \end{bmatrix}$			$\begin{bmatrix} 0.0337 & 0.0497 & -0.0424 \\ -0.0591 & -0.0872 & 0.0745 \\ -0.0424 & -0.0626 & 0.0535 \end{bmatrix}$		

family double twins and non-family double twins. The co-family double twins include $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ ($C_i^I \rightarrow T_{ij}^{II}$) and $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ ($T_{ij}^{II} \rightarrow C_j^{II}$), where both twins share the zone axis $\langle 10\bar{1}0 \rangle$. The non-family double twins include $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ ($C_i^I \rightarrow T_j^I$) and $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ ($C_i^{II} \rightarrow T_j^I$) that have different zone axes, one along $\langle 11\bar{2}0 \rangle$ and the other along $\langle 10\bar{1}0 \rangle$. Secondary twin variant selection is examined according to the apparent Schmid factor (a-SF) analysis, the displacement gradient accommodation (DGA) and the modified DGA (m-DGA) analysis, and the nucleation via dislocation dissociation (NDD). The theories were assessed by comparing the predictions with the experimental observations.

Co-family double twins can be categorized into four groups. For $\{11\bar{2}2\} \rightarrow \{11\bar{2}1\}$ double twins, the dominant double twins are described as $C_i^I \rightarrow T_{ij}^{II}$. For $\{11\bar{2}1\} \rightarrow \{11\bar{2}4\}$ double twins, the dominant double twins are described as $T_{ij}^{II} \rightarrow C_j^{II}$. In each primary twin, only one secondary twin variant is prevailed over other secondary twin variants according to the NDD and m-DGA analysis. The a-SF and DGA analyses reveal the relative possibility among six variants but cannot predict the preferred one.

Non-family double twins can be divided into three groups. For $\{11\bar{2}2\} \rightarrow \{10\bar{1}2\}$ double twins, the preferred double twins can be described as $C_i^I \rightarrow T_j^I$ or T_{i+1}^I . For $\{11\bar{2}4\} \rightarrow \{10\bar{1}2\}$ double twins, the preferred double twins can be described as $C_i^{II} \rightarrow T_j^I$ or T_{i+1}^I . The selection between the two variants obeys the SF rule. More importantly, the NDD and m-DGA can properly predict the prevailed group of secondary twin variants. The a-SF is helpful for determining the competition of two double twin variant belonging to the prevailed group. The DGA fails to predict correctly the preferred secondary twin variants.

A correct prediction based on m-DGA and NDD also implies that selection of secondary twin variant in a double twin is deterministic because of two facts, both m-DGA and NDD are geometric models and the resolved shear stress associated with secondary twinning is positive. The current study can help develop micro- and macro-scale predictive models of the deformation behavior of hexagonal materials [38–40].

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