

# *Dynamics and Stability of Surface Waves with Bulk-Soluble Surfactants*

**Ian Tice & Lei Wu**

**Acta Applicandae Mathematicae**

ISSN 0167-8019

Volume 161

Number 1

Acta Appl Math (2019) 161:35–70  
DOI 10.1007/s10440-018-0203-0

**ACTA APPLICANDAE MATHEMATICA**

ISSN 0167-8019

Vol. 161, No. 1, June 2019

*Corresponding Editors*

J. Berestycki, University of Oxford, UK  
A. DeSimone, The BioRobotics Institute, Scuola Superiore Sant'Anna, Pisa, Italy  
A. Figalli, ETH Zurich, Switzerland  
G. Kutyniok, Technische Universität Berlin, Germany  
G. Naldi, University of Milan, Italy  
P. Zhang, Chinese Academy of Sciences, Beijing, China  
E. Zuazua, DeustoTech Bilbao, Basque Country - Spain

*Associate Editors*

|                             |                          |                        |
|-----------------------------|--------------------------|------------------------|
| F. Alouges, France          | A. Garoni, Italy         | H. Rauhut, Germany     |
| K. Aoki, Japan              | A. Grothey, UK           | S. Reich, Germany      |
| S. Boyarchenko, USA         | Y. Guo, USA              | H. Schaeffer, USA      |
| E. Carlen, USA              | B. Hamby, UK             | A. Scheel, USA         |
| J.A. Carrillo, UK           | P. Jorgensen, USA        | V. Schulz, Germany     |
| Y. Chitour, France          | A. Kohatsu-Higa, Japan   | Y. Shkolniksky, Israel |
| M. Chybá, USA               | I. Kontoyannidis, Greece | S. Sotin, France       |
| M. Cicalese, Germany        | K. Kunisch, Austria      | G. Steidl, Germany     |
| P.A. Clarkson, UK           | U. Ledzewicz, USA        | S. Terracini, Italy    |
| R. Dalang, Switzerland      | D. Levy, USA             | X.J. Wang, Australia   |
| R. Danchin, France          | B. Maury, France         | M. Winkler, Germany    |
| S. Dereich, Germany         | S. Mîră, Romania         | H. Yin, China          |
| R. Donat, Spain             | C. Mouhot, UK            | X. Zhang, China        |
| M. Escobedo Martínez, Spain | R.G. Plaza, Mexico       | Z. Zhang, China        |
| M. Feldman, USA             | A. Porretta, Italy       |                        |

 Springer

**Your article is protected by copyright and all rights are held exclusively by Springer Nature B.V.. This e-offprint is for personal use only and shall not be self-archived in electronic repositories. If you wish to self-archive your article, please use the accepted manuscript version for posting on your own website. You may further deposit the accepted manuscript version in any repository, provided it is only made publicly available 12 months after official publication or later and provided acknowledgement is given to the original source of publication and a link is inserted to the published article on Springer's website. The link must be accompanied by the following text: "The final publication is available at [link.springer.com](http://link.springer.com)".**

# Dynamics and Stability of Surface Waves with Bulk-Soluble Surfactants

Ian Tice<sup>1</sup>  · Lei Wu<sup>1</sup>

Received: 26 September 2017 / Accepted: 16 July 2018 / Published online: 23 July 2018  
© Springer Nature B.V. 2018

**Abstract** In this paper we study the dynamics of a layer of incompressible viscous fluid bounded below by a rigid boundary and above by a free boundary, in the presence of a uniform gravitational field. We assume that a mass of surfactant is present both at the free surface and in the bulk of the fluid, and that conversion from one species to the other is possible. The surfactants couple to the fluid dynamics through the coefficient of surface tension, which depends on the surface density of surfactants. Gradients in this concentration give rise to Marangoni stress on the free surface. In turn, the fluid advects the surfactants and distorts their concentration through geometric distortions of the free surface. We model the surfactants in a way that allows absorption and desorption of surfactant between the surface and bulk. We prove that small perturbations of the equilibrium solutions give rise to global-in-time solutions that decay to equilibrium at an exponential rate. This establishes the asymptotic stability of the equilibrium solutions.

**Keywords** Surfactants · Navier-Stokes equations · Interfacial stability

**Mathematics Subject Classification** Primary 35Q30 · 35R35 · 76D45 · Secondary 35B40 · 76E17 · 76E99

---

I. Tice was supported by a Simons Foundation Grant (#401468) and an NSF CAREER Grant (DMS #1653161).

L. Wu was supported by NSF grant 0967140.

---

✉ I. Tice  
iantice@andrew.cmu.edu

L. Wu  
lwu2@andrew.cmu.edu

<sup>1</sup> Department of Mathematical Sciences, Carnegie Mellon University, Pittsburgh, PA 15213, USA

## 1 Introduction

### 1.1 Model Presentation: Soluble Surfactants

Surfactants (a portmanteau for “surface active agents”) are chemicals that can change the strength of surface tension when they collect at fluid free interfaces. Variations in the surfactant concentration on the surface also give rise to tangential surface forces called Marangoni forces. The surfactant dynamics are driven by several effects: absorption and desorption from the free surface, fluid transport along the surface and in the bulk, and both bulk and surface diffusion. We refer to the books [9, 17] and the review [21] for a more thorough discussion of surfactant physics. In manufacturing and industrial applications surfactants are a fundamental tool for stabilizing bubble formation in processes such as foaming, emulsifying, and coating (see the books [19, 20] for an exhaustive list of surfactant applications). Surfactants also play a critical role in preventing the collapse of the lungs during breathing (see [14] and the references therein) and are currently being developed as tools to aid in drug delivery in the lungs (see for example [4, 18]).

We consider a model of surfactants in which a single layer of fluid occupies the three-dimensional domain  $\Omega(t)$  with free boundary surface  $\Gamma(t)$ . In order to allow for surfactants concentrated on the free surface and in the bulk, we define the surface concentration  $\hat{c}(\cdot, t) : \Gamma(t) \rightarrow [0, \infty)$  as well as the bulk concentration  $\hat{b}(\cdot, t) : \Omega(t) \rightarrow [0, \infty)$ . The coefficient of surface tension on  $\Gamma(t)$  depends on the surfactant concentration  $\hat{c}$  via a relation  $\sigma = \sigma(\hat{c})$ , where we assume that the surface tension function satisfies:

$$\begin{cases} \sigma \in C^3([0, \infty)) \\ \sigma \text{ is positive and strictly decreasing.} \end{cases} \quad (1.1)$$

The latter assumption comes from the fact that surfactants decrease the surface tension in higher concentration, and the former assumption is merely a technical assumption needed for our PDE analysis.

We will assume that the fluid is incompressible and viscous (with viscosity  $\mu > 0$ ) and that a uniform gravitational field  $-ge_3 = (0, 0, -g) \in \mathbb{R}^3$  is applied to the fluid (here  $g > 0$ ). The fluid and surfactant dynamics then couple through the following system of equations (see [9, 17] for derivations and precise definitions of the operators: we will soon reformulate these equations so do not fully define the operators here)

$$\begin{cases} \partial_t u + u \cdot \nabla u + \nabla \hat{p} = \mu \Delta u - ge_3 & \text{in } \Omega(t) \\ \nabla \cdot u = 0 & \text{in } \Omega(t) \\ \hat{p} v - \mu(\nabla u + \nabla u^T)v = -\sigma(\hat{c})\mathcal{H}_{\Gamma(t)}v - \nabla_{\Gamma(t)}(\sigma(\hat{c})) & \text{on } \Gamma(t) \\ D_t \hat{c} + \hat{c} \nabla_{\Gamma(t)} \cdot u = \gamma \Delta_{\Gamma(t)} \hat{c} - \beta \nabla \hat{b} \cdot v & \text{on } \Gamma(t) \\ \partial_t \hat{b} + u \cdot \nabla \hat{b} = \beta \Delta \hat{b} & \text{in } \Omega(t) \\ \beta \nabla \hat{b} \cdot v = \omega(\hat{c}, \hat{b}) & \text{on } \Gamma(t). \end{cases} \quad (1.2)$$

Here  $\nabla_{\Gamma(t)}$  denotes the surface gradient on  $\Gamma(t)$ ,  $D_t$  is a temporal derivative along the flowing surface,  $\nabla_{\Gamma(t)} \cdot$  is the surface divergence,  $\Delta_{\Gamma(t)}$  is the surface Laplacian,  $\gamma > 0$  is the surface surfactant diffusion constant,  $\beta > 0$  is the bulk surfactant diffusion constant, and  $\mathcal{H}_{\Gamma(t)}$  is twice the mean-curvature operator on  $\Gamma(t)$ . The first two equations in (1.2) are the usual incompressible Navier-Stokes equations. The third equation is the balance of stress on the free surface, and the right-hand side shows that two stresses are generated by the

surfactants. The first term is a normal stress caused by surface curvature, and the second is a tangential stress, known as the Marangoni stress, caused by gradients in the surfactant concentration on the surface. The fourth equation in (1.2) shows that the surface surfactant concentration changes due to flow on the surface as well as diffusion and the normal derivative of  $\hat{b}$ , and the fifth equation is an advection-diffusion equation for the bulk surfactant concentration.

The sixth equation in (1.2) is of fundamental importance in surfactant modeling, as it gives the law for the conversion of bulk surfactant into surface surfactant at the free interface. We refer, for instance, to the book [9] for a derivation of such a law from physical kinetic arguments. More recent work [7, 8, 11] has derived laws of a similar form for two-phase flows using free-energy arguments. Here the function  $\omega : [0, \infty) \times [0, \infty) \rightarrow \mathbb{R}$  describes the surfactant flow from the surface to bulk phase. We will assume that  $\omega$  obeys the following:

$$\begin{cases} \omega \in C^4([0, \infty) \times [0, \infty)) \\ \text{there exists } f \in C^2([0, \infty); [0, \infty)) \text{ such that } \omega(c, b) > 0 \Leftrightarrow c > f(b) \\ f'(b) > 0 \text{ for all } b \geq 0. \end{cases} \quad (1.3)$$

In the parlance of surface chemistry, the set  $\{\omega(c, b) = 0\}$  is called an isotherm. It represents the values of surface and bulk surfactant that do not result in conversion between species, and consequently any equilibrium configuration must lie on the isotherm. Our first assumption on  $\omega$  is purely technical, but the second and third indicate that the isotherm is given as the graph of the function  $f$ . The sign condition on  $f'$  guarantees that more surface surfactant is needed to form equilibrium when more bulk surfactant is present in the fluid. As an example, in the Langmuir absorption model (see for instance [9]) we have

$$\omega(c, b) = -b(k_1 - c) + k_2 c \quad \text{for physical constants } k_1, k_2 > 0. \quad (1.4)$$

In this case  $f(b) = bk_1/(k_2 + b)$ .

Surfactants have been extensively studied in the physics literature, and we will not attempt to survey that literature here. In contrast, they have not received extensive attention in the mathematics literature. Kwan-Park-Shen [16], Xu-Li-Lowengrub-Zhao [25], and Barrett-Garcke-Nürnberg [2] developed numerical studies of surfactant dynamics. The local well-posedness of a two-phase bubble model without gravity was proved by Bothe-Prüss-Simonett [6] in the context of diffusion-limited absorption, which means that the last condition in (1.2) is replaced by the condition  $c = f(b)$  so that the surfactant phases are in constant equilibrium. The linear stability of the same model was studied by Bothe-Prüss [5]. In [15] Kim-Tice studied the dynamics of surfactants without absorption and proved the existence of global-in-time solutions near equilibrium as well as their asymptotic stability. Diffuse interface models with absorption laws similar to (1.2) have also recently been studied by Garcke-Lam-Stinner [11] and Abels-Garcke-Lam-Weber [1].

The principal goal of the present paper is to extend the techniques developed in [15] to handle the absorption model presented in (1.2). We will construct global solutions near equilibrium and prove their long-time decay to equilibrium.

## 1.2 Problem Presentation: Formulation of Equations

We now give a precise description of the equations of motion. To begin we assume that the fluid occupies a three-dimensional domain  $\mathcal{Q}(t)$  that is horizontally periodic with a flat rigid bottom and free upper surface:

$$\mathcal{Q}(t) = \{y \in \Sigma \times \mathbb{R} : -L_3 < y_3 < \eta(t, y_1, y_2)\}. \quad (1.5)$$

Here the horizontal cross-section is given by

$$\Sigma = (L_1 \mathbb{T}) \times (L_2 \mathbb{T}) = (\mathbb{R}/L_1 \mathbb{Z}) \times (\mathbb{R}/L_2 \mathbb{Z}), \quad (1.6)$$

where  $L_1, L_2 > 0$  are fixed periodicity lengths and  $L_3 > 0$  is the fixed depth. The upper free surface is the graph of the free surface function  $\eta : [0, \infty) \times \Sigma \rightarrow (-L_3, \infty)$ . We will denote the fixed lower boundary by

$$\Sigma_b = \{y \in \Sigma \times \mathbb{R} : y_3 = -L_3\}, \quad (1.7)$$

and the moving upper boundary by

$$\Gamma(t) = \{y \in \Sigma \times \mathbb{R} : y_3 = \eta(t, y_1, y_2)\}, \quad (1.8)$$

which is the graph of the unknown function  $\eta$ .

For each  $t \geq 0$ , the dynamics of the fluid and surfactants are described by the following unknowns. The fluid velocity and pressure are  $u(\cdot, t) : \Omega(t) \rightarrow \mathbb{R}^3$  and  $\hat{p}(\cdot, t) : \Omega(t) \rightarrow \mathbb{R}$ , and the free surface function is  $\eta(\cdot, t) : \Sigma \rightarrow \mathbb{R}$ . It will be convenient to study the projection onto  $\Sigma$  of the surface surfactant concentration, so we define  $\tilde{c}(\cdot, t) : \Sigma \rightarrow \mathbb{R}$  via  $\tilde{c}(y_*, t) = \hat{c}(y_*, \eta(y_*, t), t)$  where  $\hat{c}(\cdot, t)$  is defined on the moving surface  $\Gamma(t)$  and  $y_* = (y_1, y_2) \in \Sigma$ . To be consistent with the new notation we also write  $\tilde{b}(\cdot, t) = \hat{b}(\cdot, t) : \Omega(t) \rightarrow [0, \infty)$  for the bulk surfactant concentration. It will also be convenient to redefine the pressure by subtracting off the hydrostatic pressure, i.e. to set  $p(\cdot, t) : \Omega(t) \rightarrow \mathbb{R}$  via  $p(x, t) = \hat{p}(x, t) + g x_3$ .

In the following, we will employ the horizontal differential operators

$$\nabla_* F = \sum_{i=1}^2 (\partial_i F e_i) \quad \text{and} \quad \nabla_* \cdot G = \sum_{i=1}^2 (\partial_i G_i) \quad (1.9)$$

acting on scalar and vector fields, respectively. Similarly, we define  $v_* = (v_1, v_2)$  to be the horizontal components of outward unit normal vector on the moving surface:

$$v = (v_1, v_2, v_3) = \frac{(-\nabla_* \eta, 1)}{\sqrt{1 + |\nabla_* \eta|^2}}. \quad (1.10)$$

The surface differential operators are defined as

$$\partial_{\Gamma,1} = \frac{1 + |\partial_2 \eta|^2}{1 + |\nabla \eta|^2} \partial_1 - \frac{\partial_1 \eta \partial_2 \eta}{1 + |\nabla \eta|^2} \partial_2, \quad (1.11)$$

$$\partial_{\Gamma,2} = \frac{1 + |\partial_1 \eta|^2}{1 + |\nabla \eta|^2} \partial_2 - \frac{\partial_1 \eta \partial_2 \eta}{1 + |\nabla \eta|^2} \partial_1, \quad (1.12)$$

$$\partial_{\Gamma,3} = \frac{\partial_1 \eta}{1 + |\nabla \eta|^2} \partial_1 + \frac{\partial_2 \eta}{1 + |\nabla \eta|^2} \partial_2. \quad (1.13)$$

Then we have for  $f$  a scalar and  $g$  a vector field defined in  $\Sigma$ ,

$$\nabla_{\Gamma} f = \sum_{i=1}^3 (\partial_{\Gamma,i} f e_i) \quad \text{and} \quad \nabla_{\Gamma} \cdot g = \sum_{i=1}^3 (\partial_{\Gamma,i} g_i). \quad (1.14)$$

Also, for  $F$  and  $G$  defined in  $\Gamma(t)$ ,

$$\nabla_\Gamma F = \nabla_\Gamma(F \circ \eta) \quad \text{and} \quad \nabla_\Gamma \cdot G = \nabla_\Gamma \cdot (G \circ \eta). \quad (1.15)$$

Thus, we define the surface Laplacian  $\Delta_\Gamma = \nabla_\Gamma \cdot \nabla_\Gamma$ .

The fluid, surface, and surfactant unknowns  $(u, p, \eta, \tilde{c}, \tilde{b})$  must satisfy the following system of equations:

$$\left\{ \begin{array}{ll} \partial_t u + u \cdot \nabla u + \nabla p = \mu \Delta u & \text{in } \Omega(t), \\ \nabla \cdot u = 0 & \text{in } \Omega(t), \\ (pI - \mu \mathbb{D}u)v = g\eta v - \sigma(\tilde{c})H(\eta)v - \nabla_\Gamma \sigma(\tilde{c}) & \text{on } \Gamma(t), \\ u = 0 & \text{on } \Sigma_b, \\ \partial_t \eta = u_3 - u_1 \partial_{y_1} \eta - u_2 \partial_{y_2} \eta & \text{on } \Gamma(t), \\ \partial_t \tilde{c} + u \cdot \nabla_* \tilde{c} + \tilde{c} \nabla_\Gamma \cdot u = \gamma \Delta_\Gamma \tilde{c} - \beta \nabla \tilde{b} \cdot v & \text{on } \Gamma(t), \\ \partial_t \tilde{b} + u \cdot \nabla \tilde{b} = \beta \Delta \tilde{b} & \text{in } \Omega(t), \\ \beta \nabla \tilde{b} \cdot v = \omega(\tilde{c}, \tilde{b}) & \text{on } \Gamma(t), \\ \nabla \tilde{b} \cdot v = 0 & \text{on } \Sigma_b, \\ u(t=0) = u_0 & \text{in } \Omega_0, \\ \eta(t=0) = \eta_0 & \text{on } \Sigma, \\ \tilde{b}(t=0) = \tilde{b}_0 & \text{in } \Omega_0, \\ \tilde{c}(t=0) = \tilde{c}_0 & \text{on } \Sigma. \end{array} \right. \quad (1.16)$$

The first set of equations describes the motion of fluid under the influence of surfactant and the second set describes the absorption and desorption of surfactant with convection and diffusion both on the surface and in the bulk. Here viscosity  $\mu$ , gravity  $g$ , convection strength  $\beta$ , and diffusion strength  $\gamma$  are positive constants,  $I$  the  $3 \times 3$  identity matrix,  $(\mathbb{D}u)_{ij} = \partial_j u_i + \partial_i u_j$  the symmetric gradient, and the mean-curvature operator is

$$H(\eta) = \nabla_* \cdot \left( \frac{\nabla_* \eta}{\sqrt{1 + |\nabla_* \eta|^2}} \right). \quad (1.17)$$

Note that the gravitational term has shifted from the first to the third equation due to the switch in pressure unknowns  $\hat{p} \mapsto p$  described above.

Employing a standard scaling argument in space and time, we can eliminate two of the physical constants at the expense of possibly renaming the rest. We will do this in order to set  $\mu = g = 1$ . We will employ this convention throughout the rest of this paper. Also, we assume the initial surface  $\eta_0$  satisfies the zero-average condition

$$\int_\Sigma \eta_0 = 0, \quad (1.18)$$

and the total mass of surfactant

$$\int_\Sigma \tilde{c}_0 \sqrt{1 + |\nabla_* \eta_0|^2} + \int_{\Omega(t)} \tilde{b}_0 := M > 0. \quad (1.19)$$

Standard calculations, which we omit here for the sake of brevity, show that these conditions persist in time, i.e.

$$\int_\Sigma \eta(\cdot, t) = 0 \quad \text{and} \quad \int_\Sigma \tilde{c}(\cdot, t) \sqrt{1 + |\nabla_* \eta(\cdot, t)|^2} + \int_{\Omega(t)} \tilde{b}_0(\cdot, t) = M \quad (1.20)$$

for  $t > 0$ . Note, though, that the latter condition only guarantees that the total surfactant mass is preserved in time. The portion of mass in each phase (bulk and surface) does not have to be preserved.

### 1.3 Energy-Dissipation Structure

We may follow the computations in [15], with the help of Lemma A.2, to derive the following energy-dissipation structure from the fluid equations:

$$\frac{d}{dt} \left( \frac{1}{2} \int_{\Omega(t)} |u|^2 + \frac{1}{2} \int_{\Sigma} |\eta|^2 \right) + \frac{1}{2} \int_{\Omega(t)} |\mathbb{D}u|^2 = - \int_{\Sigma} \sigma(\tilde{c})(\nabla_{\Gamma} \cdot u) \sqrt{1 + |\nabla_* \eta|^2}. \quad (1.21)$$

The first term in parentheses is the kinetic energy of the fluid (the density is 1), and the second term is the total gravitational potential energy (we have set  $g = 1$  at this point) stored in the fluid. The third term on the left is the viscous dissipation (with  $\mu = 1$ ). We might hope to use similar computations with  $\tilde{c}$  and  $\tilde{b}$  to produce further energy-dissipation equations that sum with the above to produce an equation with no exchange terms on the right. Unfortunately, this does not work if we directly work with the  $\tilde{c}$  and  $\tilde{b}$ . Instead we extend an idea used in [5, 15] for the case without absorption and in [10] for a related thin-film model, which shifts from an  $L^2$ -based estimate to a more complicated estimate related to the free energy associated to  $\sigma$  (we refer to [22] for a discussion of the relevant thermodynamics). To this end, for any  $r \in (0, \infty)$  we define the auxiliary functions  $\zeta_r, \phi_r : [0, \infty) \rightarrow \mathbb{R}$  via

$$\zeta_r(s) = s \left( \frac{\sigma(r)}{r} - \int_r^s \frac{\sigma(z)}{z^2} dz \right) = \sigma(s) - s \int_r^s \frac{\sigma'(z)}{z} dz, \quad (1.22)$$

$$\phi_r(s) = C_0 + \int_0^s \frac{z\sigma'(f(z))f'(z)}{f(z)} dz - s \int_r^{f(s)} \frac{\sigma'(z)}{z} dz \quad (1.23)$$

for a constant  $C_0 > 0$  chosen such that  $\phi_r(s) > 0$  for all  $s \geq 0$ . This is possible since  $f'(0) > 0$  implies that

$$\lim_{z \rightarrow 0} \frac{zf'(z)}{f(z)} = \begin{cases} 0 & \text{if } f(0) > 0 \\ 1 & \text{if } f(0) = 0. \end{cases} \quad (1.24)$$

Note also that

$$0 \leq s \leq r \Rightarrow \left| s \int_r^s \frac{\sigma'(z)}{z} dz \right| \leq \|\sigma'\|_{L^\infty([0,r])} s \log\left(\frac{r}{s}\right), \quad (1.25)$$

so the  $\sigma'$  integral in the definition of  $\zeta_r$  is well-defined for all  $s \geq 0$ . A similar argument shows that the second integral in the definition of  $\phi_r$  is well-defined even when  $f(0) = 0$ .

By construction, we have

$$\zeta'_r(s) = - \int_r^s \frac{\sigma'(z)}{z} dz \quad \text{and} \quad \phi'_r(s) = - \int_r^{f(s)} \frac{\sigma'(z)}{z} dz, \quad (1.26)$$

as well as

$$\zeta_r(s), \phi_r(s), \zeta_r''(s), \phi_r''(s) > 0, \quad (1.27)$$

$$\zeta_r(s) - s\zeta_r'(s) = \sigma(s). \quad (1.28)$$

Moreover,  $\zeta_r$  attains its minimum at  $s = r$  and  $\phi_r$  attains its minimum if there exists  $s \geq 0$  such that  $f(s) = r$ . With these auxiliary functions in hand, we may then follow [15], using the identities in Lemma A.3, to see that

$$\begin{aligned} \frac{d}{dt} \int_{\Sigma} \zeta_r(\tilde{c}) \sqrt{1 + |\nabla_* \eta|^2} + \gamma \int_{\Sigma} \zeta_r''(\tilde{c}) |\nabla_{\Gamma} \tilde{c}|^2 \sqrt{1 + |\nabla_* \eta|^2} + \int_{\Sigma} \zeta_r'(\tilde{c}) \omega(\tilde{c}, \tilde{b}) \sqrt{1 + |\nabla_* \eta|^2} \\ = \int_{\Sigma} \sigma(\tilde{c}) (\nabla_{\Gamma} \cdot u) \sqrt{1 + |\nabla_* \eta|^2}, \end{aligned} \quad (1.29)$$

and

$$\frac{d}{dt} \int_{\Omega(t)} \phi_r(\tilde{b}) + \beta \int_{\Omega(t)} \phi_r''(\tilde{b}) |\nabla \tilde{b}|^2 - \int_{\Sigma} \phi_r'(\tilde{b}) \omega(\tilde{c}, \tilde{b}) \sqrt{1 + |\nabla_* \eta|^2} = 0. \quad (1.30)$$

Summing up (1.21), (1.29), and (1.30), we obtain the full energy-dissipation structure:

$$\begin{aligned} \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega(t)} |u|^2 + \frac{1}{2} \int_{\Sigma} |\eta|^2 + \int_{\Omega(t)} \phi_r(\tilde{b}) + \int_{\Sigma} \zeta_r(\tilde{c}) \sqrt{1 + |\nabla_* \eta|^2} \right) \\ + \left( \frac{1}{2} \int_{\Omega(t)} |\mathbb{D}u|^2 + \beta \int_{\Omega(t)} \phi_r''(\tilde{b}) |\nabla \tilde{b}|^2 + \gamma \int_{\Sigma} \zeta_r''(\tilde{c}) |\nabla_{\Gamma} \tilde{c}|^2 \sqrt{1 + |\nabla_* \eta|^2} \right) \\ + \int_{\Sigma} (\zeta_r'(\tilde{c}) - \phi_r'(\tilde{b})) \omega(\tilde{c}, \tilde{b}) \sqrt{1 + |\nabla_* \eta|^2} = 0. \end{aligned} \quad (1.31)$$

The terms appearing on the first and second lines of the energy-dissipation equality are clearly positive semi-definite, but at first glance it is not clear that the term on the third line possesses a good sign. To show that this is in fact true we first note that the definitions of  $\zeta_r$  and  $\phi_r$  allow us to compute

$$(\zeta_r'(\tilde{c}) - \phi_r'(\tilde{b})) \omega(\tilde{c}, \tilde{b}) = \left( \int_{\tilde{c}}^{f(\tilde{b})} \frac{\sigma'(z)}{z} dz \right) \omega(\tilde{c}, \tilde{b}). \quad (1.32)$$

Employing the assumptions on  $\sigma$  (1.1) and  $\omega$  (1.3), we then find that

$$\begin{cases} \tilde{c} > f(\tilde{b}) \Rightarrow \int_{\tilde{c}}^{f(\tilde{b})} \frac{\sigma'(z)}{z} dz > 0 \quad \text{and} \quad \omega(\tilde{c}, \tilde{b}) > 0, \\ \tilde{c} < f(\tilde{b}) \Rightarrow \int_{\tilde{c}}^{f(\tilde{b})} \frac{\sigma'(z)}{z} dz < 0 \quad \text{and} \quad \omega(\tilde{c}, \tilde{b}) < 0. \end{cases} \quad (1.33)$$

Thus, in both cases we have

$$\int_{\Sigma} (\zeta_r'(\tilde{c}) - \phi_r'(\tilde{b})) \omega(\tilde{c}, \tilde{b}) \sqrt{1 + |\nabla \eta|^2} \geq 0. \quad (1.34)$$

Therefore, if we define

$$E = \frac{1}{2} \int_{\Omega(t)} |u|^2 + \frac{1}{2} \int_{\Sigma} |\eta|^2 + \int_{\Omega(t)} \phi_r(\tilde{b}) + \int_{\Sigma} \zeta_r(\tilde{c}) \sqrt{1 + |\nabla \eta|^2}, \quad (1.35)$$

$$\begin{aligned} D = & \frac{1}{2} \int_{\Omega(t)} |\mathbb{D}u|^2 + \beta \int_{\Omega(t)} \phi_r''(\tilde{b}) |\nabla \tilde{b}|^2 + \gamma \int_{\Sigma} \zeta_r''(\tilde{c}) |\nabla_{\Gamma} \tilde{c}|^2 \sqrt{1 + |\nabla \eta|^2} \\ & + \int_{\Sigma} (\zeta_r'(\tilde{c}) - \phi_r'(\tilde{b})) \omega(\tilde{c}, \tilde{b}) \sqrt{1 + |\nabla \eta|^2}, \end{aligned} \quad (1.36)$$

then we know  $E$  and  $D$  are positive semi-definite and satisfy the energy-dissipation equation

$$\frac{dE}{dt} + D = 0. \quad (1.37)$$

## 1.4 Equilibria

We now turn our attention to a discussion of the equilibrium solutions. Assume we have a time-independent solution. This assumption and (1.37) imply that  $D = 0$ , and hence

$$|\mathbb{D}u|^2 = \phi_r''(\tilde{b}) |\nabla \tilde{b}|^2 = \zeta_r''(\tilde{c}) |\nabla_{\Gamma} \tilde{c}|^2 = (\zeta_r'(\tilde{c}) - \phi_r'(\tilde{b})) \omega(\tilde{c}, \tilde{b}) = 0. \quad (1.38)$$

Since  $u = 0$  on  $\Sigma_b$ , we then know from Korn's inequality that  $u = 0$ . This and the first equation in (1.16) then imply that  $p = p_0$  is a constant. Considering that  $\phi_r$  and  $\zeta_r$  are strictly convex, we know  $\nabla \tilde{b} = \nabla_{\Gamma} \tilde{c} = 0$ . Hence, we have  $\tilde{b} = b_0 \in (0, \infty)$ ,  $\tilde{c} = c_0 \in (0, \infty)$ . Plugging this information into the third equation of (1.16) then shows that  $p_0 = g\eta - \sigma(c_0)H(\eta)$  on  $\Sigma$ , and upon integrating and using the zero-average condition for  $\eta$  given in (1.20), we find that

$$p_0 L_1 L_2 = \int_{\Sigma} p_0 = \int_{\Sigma} g\eta - \sigma(c_0) \nabla_* \cdot \left( \frac{\nabla_* \eta}{\sqrt{1 + |\nabla_* \eta|^2}} \right) = 0. \quad (1.39)$$

Thus  $0 = p_0 = g\eta - \sigma(c_0)H(\eta)$ , and upon multiplying by  $\eta$  and integrating by parts, we see that

$$0 = \int_{\Sigma} g|\eta|^2 + \sigma(c_0) \frac{|\nabla_* \eta|^2}{\sqrt{1 + |\nabla_* \eta|^2}}, \quad (1.40)$$

from which we deduce that  $\eta = 0$ . Thus  $p = \eta = 0$ , and the equilibrium domain is the flat slab  $\Sigma \times (-L_3, 0)$ . Furthermore,  $b_0$  and  $c_0$  must preserve the total mass as in (1.20),

$$|\Sigma|_2 c_0 + |\Omega|_3 b_0 = \int_{\Sigma} c_0 + \int_{\Omega} b_0 = M, \quad (1.41)$$

where  $|\cdot|_k$  denotes  $k$ -dimensional Lebesgue measure, as well as satisfy the condition  $\omega(c_0, b_0) = 0$ . The latter requires that  $f(b_0) = c_0$ . Since  $f' > 0$ , for given total mass  $M > |\Sigma|_2 f(0)$ , we have a unique solution  $b_0, c_0 = f(b_0) > 0$ .

In summary, we have the equilibrium configuration

$$\begin{cases} u = 0 \\ p = 0 \\ \eta = 0 \\ \omega(c_0, b_0) = 0, \\ |\Sigma|_2 c_0 + |\Omega|_3 b_0 = M, \end{cases} \quad (1.42)$$

where  $c_0, b_0 \in (0, \infty)$  are uniquely determined by the choice of mass  $0 \leq |\Sigma|_2 f(0) < M < \infty$  and the form of the surfactant flux function  $\omega$ .

It is worth noting that in the definition of  $\zeta_r$  and  $\phi_r$  above we have that  $\zeta_r$  obtains its minimal value at  $s = r$  and  $\phi_r$  obtains its minimum at  $s$  such that  $f(s) = r$ . If we choose  $r = c_0$  then the minimal value of  $\zeta_r$  is obtained at  $c_0$  and the minimal value of  $\phi_r$  is obtained at  $f(c_0) = b_0$ . Thus (1.37) suggests that the equilibrium configuration will be dynamically stable.

We conclude our discussion of the equilibria by introducing some notation that will be useful later. We linearize  $\omega$  around  $(c_0, b_0)$  as

$$\begin{aligned} \omega(c_0 + h, b_0 + d) &= \omega(c_0, b_0) + \omega_{c0}h + \omega_{b0}d \\ &\quad + 2hd \int_0^1 (1-s) \partial_{cb} \omega(c_0 + sh, b_0 + sd) ds \\ &\quad + h^2 \int_0^1 (1-s) \partial_{cc} \omega(c_0 + sh, b_0 + sd) ds \\ &\quad + d^2 \int_0^1 (1-s) \partial_{bb} \omega(c_0 + sh, b_0 + sd) ds \\ &= \omega(c_0, b_0) + \omega_{c0}h + \omega_{b0}d + O(|h|^2 + |d|^2) \\ &= \omega_{c0}h + \omega_{b0}d + O(|h|^2 + |d|^2), \end{aligned} \quad (1.43)$$

where we have written

$$\omega_{c0} := \partial_c \omega(c_0, b_0) \text{ and } \omega_{b0} := \partial_b \omega(c_0, b_0). \quad (1.44)$$

Since  $\omega(f(b), b) = 0$  for all  $b \geq 0$ , we may differentiate to see that

$$\omega_{c0} f'(b_0) + \omega_{b0} = 0, \quad (1.45)$$

which implies

$$\frac{\omega_{b0}}{\omega_{c0}} = -f'(b_0). \quad (1.46)$$

Also, since  $\omega(c, b) > 0$  leads to  $c > f(b)$ , we have  $\omega_{c0} > 0$  and  $\omega_{b0} < 0$ .

## 1.5 Reformulation

In order to work in a fixed domain, we employ a frequently used transformation: see [3, 12, 13, 15, 23]. We define the equilibrium fluid domain  $\Omega$  via

$$\Omega = \{y \in \Sigma \times \mathbb{R} : -L_3 < y_3 < 0\}, \quad (1.47)$$

which possesses  $\Sigma$  as upper boundary and  $\Sigma_b$  as lower boundary. We can then view  $\eta$  as defined on  $\Sigma$ . Assume  $\bar{\eta}$  is the harmonic extension of  $\eta$  into  $\Sigma \times (-\infty, 0)$  defined as

in (B.1). Define mapping

$$\Phi : \Omega \ni x = (x_1, x_2, x_3) \mapsto \left( x_1, x_2, x_3 + \bar{\eta}(x, t) \left( 1 + \frac{x_3}{L_3} \right) \right) \in \Omega(t). \quad (1.48)$$

We may easily check that  $\Phi$  maps  $\Sigma$  to  $\Gamma(t)$  and maps  $\Sigma_b$  to  $\Sigma_b$ . The Jacobian and transform matrices are

$$\nabla \Phi = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ A & B & J \end{pmatrix} \quad \text{and} \quad \mathcal{A} = (\nabla \Phi)^{-T} = \begin{pmatrix} 1 & 0 & -AK \\ 0 & 1 & -BK \\ 0 & 0 & K \end{pmatrix}, \quad (1.49)$$

where

$$\begin{aligned} A &= \tilde{L} \frac{\partial \bar{\eta}}{\partial x_1}, & B &= \tilde{L} \frac{\partial \bar{\eta}}{\partial x_2}, \\ J &= 1 + \frac{\bar{\eta}}{L_3} + \tilde{L} \frac{\partial \bar{\eta}}{\partial x_3}, & K &= \frac{1}{J}, & \tilde{L} &= 1 + \frac{x_3}{L_3}. \end{aligned} \quad (1.50)$$

Note that  $J = \det \nabla \Phi$ , and that in principle  $J$  could vanish. We prevent this in our analysis by working in a small energy regime (see Lemma 4.1), which imposes a bound on  $\bar{\eta}$  sufficient to guarantee that  $J > 0$  and that  $\Phi$  is a diffeomorphism.

In this new coordinate system, (1.16) is transformed into

$$\begin{cases} \partial_t u - \partial_t \bar{\eta} \tilde{L} K \partial_3 u + u \cdot \nabla_{\mathcal{A}} u - \Delta_{\mathcal{A}} u + \nabla_{\mathcal{A}} p = 0 & \text{in } \Omega, \\ \nabla_{\mathcal{A}} \cdot u = 0 & \text{in } \Omega, \\ (pI - \mathbb{D}_{\mathcal{A}} u) \mathcal{N} = \eta \mathcal{N} - \sigma(\tilde{c}) H(\eta) \mathcal{N} - \sqrt{1 + |\nabla_* \eta|^2} \sigma'(\tilde{c}) \nabla_{\Gamma} \tilde{c} & \text{on } \Sigma, \\ u = 0 & \text{on } \Sigma_b, \\ \partial_t \eta = u \cdot \mathcal{N} & \text{on } \Sigma, \\ \partial_t \tilde{c} + u \cdot \nabla_* \tilde{c} + \tilde{c} \nabla_{\Gamma} \cdot u = \gamma \Delta_{\Gamma} \tilde{c} - \beta \nabla_{\mathcal{A}} \tilde{b} \cdot \frac{\mathcal{N}}{\sqrt{1 + |\nabla_* \eta|^2}} & \text{on } \Sigma, \\ \partial_t \tilde{b} - \partial_t \bar{\eta} \tilde{L} K \partial_3 \tilde{b} + u \cdot \nabla_{\mathcal{A}} \tilde{b} = \beta \Delta_{\mathcal{A}} \tilde{b} & \text{in } \Omega, \\ \nabla_{\mathcal{A}} \tilde{b} \cdot \frac{\mathcal{N}}{\sqrt{1 + |\nabla_* \eta|^2}} = \omega(\tilde{c}, \tilde{b}) & \text{on } \Sigma, \\ \partial_3 \tilde{b} = 0 & \text{on } \Sigma_b. \end{cases} \quad (1.51)$$

Here we have defined the transformed operators as follows:

$$\begin{aligned} (\nabla_{\mathcal{A}} f)_i &= \mathcal{A}_{ij} \partial_j f, \\ \nabla_{\mathcal{A}} \cdot g &= \mathcal{A}_{ij} \partial_j g_i, \\ \Delta_{\mathcal{A}} f &= \nabla_{\mathcal{A}} \cdot \nabla_{\mathcal{A}} f, \\ \mathcal{N} &= (-\partial_1 \eta, -\partial_2 \eta, 1), \\ (\mathbb{D}_{\mathcal{A}} u)_{ij} &= \mathcal{A}_{ik} \partial_k u_j + \mathcal{A}_{jk} \partial_k u_i, \\ \mathcal{S}_{\mathcal{A}}(p, u) &= pI - \mathbb{D}_{\mathcal{A}} u, \end{aligned} \quad (1.52)$$

where the summation should be understood in the Einstein convention. If we extend the divergence  $\nabla_{\mathcal{A}} \cdot$  to act on symmetric tensors in the natural way, then a straightforward computation reveals  $\nabla_{\mathcal{A}} \cdot \mathcal{S}_{\mathcal{A}}(p, u) = \nabla_{\mathcal{A}} p - \Delta_{\mathcal{A}} u$  for vector fields satisfying  $\nabla_{\mathcal{A}} \cdot u = 0$ .

## 1.6 Perturbation

We will study the case when the initial data is a perturbation of the equilibrium. Define the perturbation as  $c = \tilde{c} - c_0$  and  $b = \tilde{b} - b_0$ , then  $(u, p, \eta, c, b)$  satisfies

$$\left\{ \begin{array}{ll} \partial_t u - \partial_t \bar{\eta} \tilde{L} K \partial_3 u + u \cdot \nabla_{\mathcal{A}} u - \Delta_{\mathcal{A}} u + \nabla_{\mathcal{A}} p = 0 & \text{in } \Omega, \\ \nabla_{\mathcal{A}} \cdot u = 0 & \text{in } \Omega, \\ (pI - \mathbb{D}_{\mathcal{A}} u) \mathcal{N} = \eta \mathcal{N} - \sigma(c + c_0) H(\eta) \mathcal{N} - \sqrt{1 + |\nabla_* \eta|^2} \sigma'(c + c_0) \nabla_{\Gamma} c & \text{on } \Sigma, \\ u = 0 & \text{on } \Sigma_b, \\ \partial_t \eta = u \cdot \mathcal{N} & \text{on } \Sigma, \\ \partial_t c + u \cdot \nabla_* c + (c + c_0) \nabla_{\Gamma} \cdot u = \gamma \Delta_{\Gamma} c - \beta \nabla_{\mathcal{A}} b \cdot \frac{\mathcal{N}}{\sqrt{1 + |\nabla_* \eta|^2}} & \text{on } \Sigma, \\ \partial_t b - \partial_t \bar{\eta} \tilde{L} K \partial_3 b + u \cdot \nabla_{\mathcal{A}} b = \beta \Delta_{\mathcal{A}} b & \text{in } \Omega, \\ \nabla_{\mathcal{A}} b \cdot \frac{\mathcal{N}}{\sqrt{1 + |\nabla_* \eta|^2}} = \omega(c + c_0, b + b_0) & \text{on } \Sigma, \\ \partial_3 b = 0 & \text{on } \Sigma_b. \end{array} \right. \quad (1.53)$$

In the following, we will write  $\sigma_0 = \sigma(c_0)$  and  $\sigma'_0 = \sigma'(c_0)$ .

## 2 Main Results and Discussion

### 2.1 Main Result

To properly state the main results we must first define energy and dissipation functionals that will be used throughout the paper. We define the full energy to be

$$\begin{aligned} \mathcal{E} = & \|u\|_{H^2(\Omega)}^2 + \|\partial_t u\|_{H^0(\Omega)}^2 + \|p\|_{H^1(\Omega)}^2 + \|\eta\|_{H^3(\Sigma)}^2 + \|\partial_t \eta\|_{H^{\frac{3}{2}}(\Sigma)}^2 + \|\partial_t^2 \eta\|_{H^{-\frac{1}{2}}(\Sigma)}^2 \\ & + \|b\|_{H^2(\Omega)}^2 + \|\partial_t b\|_{H^0(\Omega)}^2 + \|c\|_{H^2(\Sigma)}^2 + \|\partial_t c\|_{H^0(\Sigma)}^2, \end{aligned} \quad (2.1)$$

and full dissipation to be

$$\begin{aligned} \mathcal{D} = & \|u\|_{H^3(\Omega)}^2 + \|\partial_t u\|_{H^1(\Omega)}^2 + \|p\|_{H^2(\Omega)}^2 + \|\eta\|_{H^{\frac{7}{2}}(\Sigma)}^2 + \|\partial_t \eta\|_{H^{\frac{5}{2}}(\Sigma)}^2 + \|\partial_t^2 \eta\|_{H^{\frac{1}{2}}(\Sigma)}^2 \\ & + \|b\|_{H^3(\Omega)}^2 + \|\partial_t b\|_{H^1(\Omega)}^2 + \|c\|_{H^3(\Sigma)}^2 + \|\partial_t c\|_{H^1(\Sigma)}^2. \end{aligned} \quad (2.2)$$

Here the spaces  $H^s$  denote the usual  $L^2$ -based Sobolev spaces of order  $s$ .

Now we present the main theorem on a priori estimates for solutions to (1.53).

**Theorem 2.1** (Proved later in Theorem 7.1) *Suppose that  $(u, p, \eta, c, b)$  solves (1.53) on the temporal interval  $[0, T]$ . Let  $\mathcal{E}$  and  $\mathcal{D}$  be as defined in (2.1) and (2.2). Then there exists a universal constant  $0 < \delta$  (independent of  $T$ ) such that if*

$$\sup_{0 \leq t \leq T} \mathcal{E}(t) \leq \delta \quad \text{and} \quad \int_0^T \mathcal{D}(t) dt < \infty, \quad (2.3)$$

then

$$\sup_{0 \leq t \leq T} e^{\lambda t} \mathcal{E}(t) + \int_0^T \mathcal{D}(t) dt \lesssim \mathcal{E}(0) \quad (2.4)$$

for all  $t \in [0, T]$ , where  $\lambda > 0$  is a universal constant.

This theorem tells us that if solutions exist for which the energy functional  $\mathcal{E}$  remains small and the dissipation functional is integrable in time, then in fact we have much stronger information: the energy decays exponentially and the integral of the dissipation is controlled by the initial energy. In order for this result to be useful we must couple it with a local existence result. By now it is well-understood how to construct local-in-time solutions for solutions to problems of the form (1.53) once the corresponding a priori estimates are understood: we refer for instance to [12, 23, 24] for local existence results in spaces determined by energies and dissipations of the form (2.1) and (2.2), and to [6] for results that employ  $L^p$ -maximal regularity techniques. Consequently, in the interest of brevity, we will not attempt to prove a local existence result in the present paper. Instead we will simply state the result that one can prove by modifying the known methods in straightforward ways.

Given the initial data  $u_0, \eta_0, \tilde{c}_0, \tilde{b}_0$ , we need to construct the initial data  $\partial_t u(\cdot, 0), \partial_t \eta(\cdot, 0), \partial_t c(\cdot, 0), \partial_t b(\cdot, 0)$ , and  $p(\cdot, 0)$ . To construct these we require a compatibility condition for the data. To state this properly we define the orthogonal projection onto the tangent space of the surface  $\Gamma(0) = \{x_3 = \eta_0(x_*)\}$  according to

$$\Pi_0 v = v - (v \cdot \mathcal{N}_0) \mathcal{N}_0 |\mathcal{N}_0|^{-2} \quad (2.5)$$

for  $\mathcal{N}_0 = (-\partial_1 \eta_0, -\partial_2 \eta_0, 1)$ . Then the compatibility conditions for the data read

$$\begin{cases} \Pi_0(\mathbb{D}_{\mathcal{A}_0} u_0 \mathcal{N}_0) - \sqrt{1 + |\nabla_* \eta_0|^2} \sigma'(\tilde{c}_0) \nabla_{\Gamma_0} \tilde{c}_0 = 0 & \text{on } \Sigma \\ \nabla_{\mathcal{A}_0} \cdot u_0 = 0 & \text{in } \Omega \\ u_0 = 0 & \text{on } \Sigma_b, \end{cases} \quad (2.6)$$

where here  $\mathcal{A}_0$  and  $\Gamma_0$  are determined by  $\eta_0$ . To state the local result we will also need to define  $\mathcal{H}^1 := \{u \in H^1(\Omega) \mid u|_{\Sigma_b} = 0\}$  and

$$\mathcal{X}_T = \{u \in L^2([0, T]; \mathcal{H}^1) \mid \nabla_{\mathcal{A}(t)} \cdot u(t) = 0 \text{ for a.e. } t\}. \quad (2.7)$$

Having stated the compatibility conditions, we can now state the local existence result.

**Theorem 2.2** *Let  $u_0 \in H^2(\Omega)$ ,  $\eta_0 \in H^3(\Sigma)$ ,  $\tilde{c}_0 \in H^2(\Sigma)$ ,  $\tilde{b}_0 \in H^2(\Omega)$ , and assume that  $\eta_0$  satisfies (1.18) and  $(\tilde{c}_0, \tilde{b}_0)$  satisfy (1.19), where  $M > |\Sigma|_2 f(b_0)$  is the equilibrium mass associated to the equilibrium surfactant concentration  $(c_0, b_0)$  given in (1.42). Further assume that the initial data satisfy the compatibility conditions of (2.6). Let  $T > 0$ . Then there exists a universal constant  $\kappa > 0$  such that if*

$$\|u_0\|_{H^2(\Omega)}^2 + \|\eta_0\|_{H^3(\Sigma)}^2 + \|\tilde{c}_0 - c_0\|_{H^2(\Sigma)}^2 + \|\tilde{b}_0 - b_0\|_{H^2(\Omega)}^2 \leq \kappa, \quad (2.8)$$

*then there exists a unique (strong) solution  $(u, p, \eta, c, b)$  to (1.53) on the temporal interval  $[0, T]$  satisfying the estimate*

$$\sup_{0 \leq t \leq T} \mathcal{E}(t) + \int_0^T \mathcal{D}(t) dt + \int_0^T [\|\partial_t^2 c(t)\|_{H^{-1}(\Sigma)}^2 + \|\partial_t^2 b(t)\|_{(H^1(\Omega))^*}^2] dt + \|\partial_t^2 u\|_{(\mathcal{X}_T)^*}^2 \lesssim \mathcal{E}(0). \quad (2.9)$$

Moreover,  $\eta$  is such that the mapping  $\Phi(\cdot, t)$ , defined by (1.48), is a  $C^1$  diffeomorphism for each  $t \in [0, T]$ .

**Remark 2.3** All of the computations involved in the a priori estimates that we develop in this paper are justified by Theorem 2.2.

The local existence and a priori estimate then combine to yield the following global-in-time existence and decay result.

**Theorem 2.4** Let  $u_0 \in H^2(\Omega)$ ,  $\eta_0 \in H^3(\Sigma)$ ,  $\tilde{c}_0 \in H^2(\Sigma)$ ,  $\tilde{b}_0 \in H^2(\Omega)$ , and assume that  $\eta_0$  satisfies (1.18) and  $(\tilde{c}_0, \tilde{b}_0)$  satisfy (1.19), where  $M > |\Sigma|_2 f(b_0)$  is the equilibrium mass associated to the equilibrium surfactant concentration  $(c_0, b_0)$  given in (1.42). Further assume that the initial data satisfy the compatibility conditions of (2.6). Then there exists a universal  $\kappa > 0$  such that if

$$\|u_0\|_{H^2(\Omega)} + \|\eta_0\|_{H^3(\Sigma)} + \|\tilde{b}_0 - b_0\|_{H^2(\Omega)} + \|\tilde{c}_0 - c_0\|_{H^2(\Sigma)} \leq \kappa, \quad (2.10)$$

then there exists a unique solution  $(u, p, \eta, c, b)$  to (1.53) in the global temporal interval  $[0, \infty)$  and a universal constant  $\lambda > 0$  such that

$$\sup_{t \in [0, \infty)} e^{\lambda t} \mathcal{E}(t) + \int_0^\infty \mathcal{D}(t) dt \lesssim \mathcal{E}(0). \quad (2.11)$$

Moreover,  $\eta$  is such that the mapping  $\Phi$  defined by (1.48) is a  $C^1$  diffeomorphism.

*Proof* This follows from Theorems 2.1 and 2.2 in a standard way. See the proof of [15, Theorem 3] for details.  $\square$

**Remark 2.5** Theorem 2.4 can be interpreted as an asymptotic stability result: the equilibria  $u = 0$ ,  $p = 0$ ,  $\eta = 0$ ,  $\tilde{c} = c_0$  and  $\tilde{b} = b_0$  are asymptotically stable, and solutions return to equilibrium exponentially fast.

**Remark 2.6** The surface function  $\eta$  is sufficiently small to guarantee that the mapping  $\Phi(\cdot, t)$ , defined in (1.48), is a diffeomorphism for each  $t \geq 0$ . As such, we may change coordinates to  $y \in \Omega(t)$  to produce a global-in-time, decaying solution to (1.2).

## 2.2 Plan of Paper and Summary of Arguments

Define the horizontal energy

$$\begin{aligned} \bar{\mathcal{E}} = & \sum_{|\alpha| \leq 2} \left( \frac{1}{2} \int_{\Omega} |\partial^\alpha u|^2 + \frac{1}{2} \int_{\Sigma} |\partial^\alpha \eta|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \partial^\alpha \eta|^2 \right. \\ & \left. + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |\partial^\alpha b|^2 + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |\partial^\alpha c|^2 \right), \end{aligned} \quad (2.12)$$

and the horizontal dissipation

$$\begin{aligned} \bar{\mathcal{D}} = & \sum_{|\alpha| \leq 2} \left( \frac{1}{2} \int_{\Omega} |\mathbb{D} \partial^\alpha u|^2 + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} |\nabla \partial^\alpha b|^2 \right. \\ & \left. + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* \partial^\alpha c|^2 + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial^\alpha c - f'(b_0) \partial^\alpha b)^2 \right), \end{aligned} \quad (2.13)$$

where  $\alpha = (\alpha_0; \alpha_1, \alpha_2, \alpha_3)$  with  $\alpha_3 = 0$  is the temporal and spatial derivatives index with  $|\alpha| = 2\alpha_0 + \alpha_1 + \alpha_2 + \alpha_3$ .

The global well-posedness and exponential decay follow from a delicate nonlinear energy method. As in [13, 15, 23, 24], we utilize energy and dissipation with higher-order derivatives. This includes the following three steps:

- (1) Horizontal energy-dissipation estimates: We apply horizontal derivatives (i.e.  $\partial_1$ ,  $\partial_2$ , and  $\partial_t$ ) to the equations and employ a variant of the basic energy-dissipation equality (1.37) to get estimates. The horizontal derivatives are needed because they are the only ones compatible with the boundary conditions. Since the derivatives do not commute with the operators in (1.53) we arrive at an equality of the form

$$\frac{d\bar{\mathcal{E}}}{dt} + \bar{\mathcal{D}} = \mathcal{I} \quad (2.14)$$

for a nonlinear interaction term  $\mathcal{I}$ .

- (2) Nonlinear estimates: Here we seek to bound  $\mathcal{I}$  by the full energy and dissipation  $\mathcal{E}$  and  $\mathcal{D}$  in the form  $|\mathcal{I}| \lesssim \sqrt{\mathcal{E}\mathcal{D}}$ , which means  $\mathcal{I}$  can be absorbed into  $\mathcal{D}$  on the left side of the energy-dissipation equality for small  $\mathcal{E}$ . This particular structure is essential for the absorbing argument.
- (3) Comparison estimates: We utilize the horizontal energy and dissipation to control their “full” counterparts defined in (2.1) and (2.2). Up to error terms we show that

$$\mathcal{E} \lesssim \bar{\mathcal{E}} \text{ and } \mathcal{D} \lesssim \bar{\mathcal{D}}, \quad (2.15)$$

which can then be used to get bounds and decay estimates for the full energy and dissipation.

We follow this general outline here as well, but there are a couple interesting new features in the analysis.

*Averaging estimate of  $c$  and  $b$ .* In order to prove exponential decay we must first prove the coercivity estimate  $\mathcal{E} \lesssim \mathcal{D}$ . The techniques for estimating the non-surfactant terms here can be found in [15], so it is the pair of surfactant terms that are of primary interest. Roughly speaking, we need to bound  $\|c\|_{H^0(\Sigma)}$  and  $\|b\|_{H^0(\Omega)}$  with  $\|\nabla_* c\|_{H^0(\Sigma)}$  and  $\|\nabla b\|_{H^0(\Omega)}$ . When the masses of  $\tilde{c}$  and  $\tilde{b}$  are conserved separately in the dynamics, this estimate may be proved via the Poincaré-Wirtinger inequality and estimates of the averages. However, the absorption and desorption of surfactant between surface and bulk yield non-constant averages of  $c$  and  $b$ . To get around this problem we introduce an enhanced Poincaré-type inequality, which combines the bound of  $c$  and  $b$  in one step, i.e. for  $C \in H^1(\Sigma)$  and  $B \in H^1(\Omega)$  satisfying

$$\int_{\Sigma} C + \int_{\Omega} B = 0, \quad (2.16)$$

we prove

$$\|C\|_{H^0(\Sigma)} + \|B\|_{H^0(\Omega)} \lesssim \|\nabla_* C\|_{H^0(\Sigma)} + \|\nabla B\|_{H^0(\Omega)} + \|C - f'(b_0)B\|_{H^0(\Sigma)}, \quad (2.17)$$

where  $C$  and  $B$  can be variants of  $c$  and  $b$ . We thus see an interesting feature of the dynamics: the conservation of the sum of two quantities combines with the linearization of the flux function  $\omega$  to provide a coercivity estimate for both surfactant phases at once.

*Nonlinear estimate of  $\omega$ .* The linearization (1.43) of  $\omega$  plays a key role in the derivation of the full energy-dissipation structure and the estimate of higher-order derivatives. As

Lemmas 3.1 and 3.2 reveal, the quantity  $c - f'(b_0)b$  creates a non-negative bound for the interaction between  $c$  and  $b$  and helps also in the above averaging estimate. Furthermore, the integral form of  $\omega(c_0 + c, b_0 + b) - \omega(c_0, b_0) - \omega_{c0}c - \omega_{b0}b$  leads to small constants which assist in the absorbing of nonlinear term  $\mathcal{I}$ .

## 2.3 Definitions and Terminology

We now mention some of the definitions, bits of notation, and conventions that we will use throughout the paper.

*Einstein summation and constants:* We will employ the Einstein convention of summing over repeated indices for vector and tensor operations. Throughout the paper  $C > 0$  will denote a generic constant that can depend on the parameters of the problem, and  $\Omega$ , but does not depend on the data, etc. We refer to such constants as “universal.” They are allowed to change from one inequality to the next. We will employ the notation  $a \lesssim b$  to mean that  $a \leq Cb$  for a universal constant  $C > 0$ .

*Norms:* We write  $H^k(\Omega)$  with  $k \geq 0$  and  $H^s(\Sigma)$  with  $s \in \mathbb{R}$  for the usual Sobolev spaces. We will typically write  $H^0 = L^2$ . To avoid notational clutter, we will avoid writing  $H^k(\Omega)$  or  $H^k(\Sigma)$  in our norms and typically write only  $\|\cdot\|_k$  for  $H^k(\Omega)$  norms and  $\|\cdot\|_{\Sigma,s}$  for  $H^s(\Sigma)$  norms.

## 3 Perturbation Form

In this section, we introduce two types of formulation for system (1.53) which will assist the estimate of temporal and horizontal derivatives.

### 3.1 Geometric Perturbed Form

Here we consider the first linearized formulation of (1.53), which we can think of as a geometric perturbed form due to the appearance of the geometric coefficients  $\mathcal{A}$ ,  $K$ , etc. We will state the formulation for the unknowns  $(v, q, \xi, h, d)$ , with given forcing terms  $F^1, \dots, F^7$  defined on appropriate sets. In later applications the unknowns will correspond to the solution or its derivatives, and the forcing terms will have specific structures related to the solution and its derivatives (see below for the specific forms). The geometric perturbed linear equations are then:

$$\left\{ \begin{array}{ll} \partial_t v - \partial_t \bar{\eta} \tilde{L} K \partial_3 v + u \cdot \nabla_{\mathcal{A}} v + \nabla_{\mathcal{A}} \cdot \mathcal{S}_{\mathcal{A}}(v, q) = F^1 & \text{in } \Omega, \\ \nabla_{\mathcal{A}} \cdot v = F^2 & \text{in } \Omega, \\ \mathcal{S}(v, q) \mathcal{N} = \xi \mathcal{N} - \sigma_0 \Delta_* \xi \mathcal{N} - \sigma'_0 \nabla_* h + F^3 & \text{on } \Sigma, \\ v = 0 & \text{on } \Sigma_b, \\ \partial_t \xi = v \cdot \mathcal{N} + F^4 & \text{on } \Sigma, \\ \partial_t h + c_0 \nabla_* \cdot v = \gamma \Delta_* h - \beta \nabla_{\mathcal{A}} d \cdot \mathcal{N} + F^5 & \text{on } \Sigma, \\ \partial_t d - \partial_t \bar{\eta} \tilde{L} K \partial_3 d + u \cdot \nabla_{\mathcal{A}} d = \beta \Delta_{\mathcal{A}} d + F^6 & \text{in } \Omega, \\ \nabla_{\mathcal{A}} d \cdot \mathcal{N} = \omega_c h + \omega_{b0} d + F^7 & \text{on } \Sigma, \\ \partial_3 d = 0 & \text{on } \Sigma_b. \end{array} \right. \quad (3.1)$$

The next result records the energy-dissipation equation associated to this problem.

**Lemma 3.1** *We have*

$$\begin{aligned}
& \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} J|v|^2 + \frac{1}{2} \int_{\Sigma} |\xi|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \xi|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} J|d|^2 + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |h|^2 \right) \\
& + \left( \frac{1}{2} \int_{\Omega} J|\mathbb{D}_{\mathcal{A}} v|^2 + \frac{-\beta\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} J|\nabla_{\mathcal{A}} d|^2 + \frac{-\gamma\sigma'_0}{c_0} \int_{\Sigma} |\nabla_* h|^2 \right) \\
& + \frac{-\beta\sigma'_0}{c_0} \int_{\Sigma} (h - f'(b_0)d)(\omega_{c0}h + \omega_{b0}d) \\
& = \int_{\Omega} Jv \cdot F^1 + \int_{\Omega} JqF^2 - \int_{\Sigma} v \cdot F^3 + \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) F^4 \\
& + \frac{-\sigma'_0}{c_0} \int_{\Sigma} hF^5 + \frac{-\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} JdF^6 + \frac{\beta\sigma'_0}{c_0} \int_{\Sigma} (h - f'(b_0)d) F^7. \tag{3.2}
\end{aligned}$$

*Proof* Multiplying  $Jv$  on both sides of the Navier-Stokes equations and integrating over  $\Omega$ , we have

$$\int_{\Omega} Jv \cdot \partial_t v + \int_{\Omega} Jv \cdot (-\partial_t \bar{\eta} \tilde{L} K \partial_3 v + u \cdot \nabla_{\mathcal{A}} v) + \int_{\Omega} Jv \cdot (\nabla_{\mathcal{A}} \cdot \mathcal{S}_{\mathcal{A}}(v, q)) = \int_{\Omega} Jv \cdot F^1. \tag{3.3}$$

We may directly compute

$$\begin{aligned}
& \frac{d}{dt} \left( \int_{\Omega} Jv \cdot \partial_t v \right) + \int_{\Omega} Jv \cdot (-\partial_t \bar{\eta} \tilde{L} K \partial_3 v + u \cdot \nabla_{\mathcal{A}} v) \\
& = \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} J|v|^2 \right) - \frac{1}{2} \int_{\Omega} \partial_t J|v|^2 + \frac{1}{2} \int_{\Omega} \partial_3 (\partial_t \bar{\eta} \tilde{L}) |v|^2 - \int_{\Sigma} \partial_t \eta |v^2| \\
& \quad - \int_{\Omega} (\nabla_{\mathcal{A}} \cdot u) |v^2| + \int_{\Sigma} (u \cdot \mathcal{N}) |v^2| \\
& = \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} J|v|^2 \right). \tag{3.4}
\end{aligned}$$

Also, we know

$$\begin{aligned}
& \int_{\Omega} Jv \cdot (\nabla_{\mathcal{A}} \cdot \mathcal{S}_{\mathcal{A}}(v, q)) \\
& = - \int_{\Omega} Jq(\nabla_{\mathcal{A}} \cdot v) + \frac{1}{2} \int_{\Omega} J|\mathbb{D}_{\mathcal{A}} v|^2 + \int_{\Sigma} (\mathcal{S}(v, q) \mathcal{N}) \cdot v \\
& = - \int_{\Omega} JqF^2 + \frac{1}{2} \int_{\Omega} J|\mathbb{D}_{\mathcal{A}} v|^2 + \int_{\Sigma} (\xi \mathcal{N} - \sigma_0 \Delta_* \xi \mathcal{N}) \cdot v - \sigma'_0 \int_{\Sigma} \nabla_* h \cdot v + \int_{\Sigma} F^3 \cdot v. \tag{3.5}
\end{aligned}$$

We may continue computing

$$\begin{aligned}
\int_{\Sigma} (\xi \mathcal{N} - \sigma_0 \Delta_* \xi \mathcal{N}) \cdot v & = \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) (\partial_t \xi - F^4) \\
& = \frac{d}{dt} \left( \frac{1}{2} \int_{\Sigma} |\xi|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \xi|^2 \right) + \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) F^4. \tag{3.6}
\end{aligned}$$

In surface surfactant equation, multiplying  $\frac{-\sigma'_0}{c_0}h$  on both sides and integrating over  $\Sigma$ , we have

$$\begin{aligned} \frac{d}{dt} \left( \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |h|^2 \right) + \sigma'_0 \int_{\Sigma} \nabla_* h \cdot v + \frac{-\gamma\sigma'_0}{c_0} \int_{\Sigma} |\nabla_* h|^2 + \frac{-\beta\sigma'_0}{c_0} \int_{\Sigma} h (\nabla_{\mathcal{A}} d \cdot \mathcal{N}) \\ = \frac{-\sigma'_0}{c_0} \int_{\Sigma} h F^5. \end{aligned} \quad (3.7)$$

In surface surfactant equation, multiplying  $\frac{-\beta\sigma'_0 f'(b_0)}{c_0} J d$  on both sides and integrating over  $\Omega$ , we have

$$\begin{aligned} \frac{d}{dt} \left( \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} J |d|^2 \right) + \frac{-\beta\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} J |\nabla_{\mathcal{A}} d|^2 - \frac{-\beta\sigma'_0}{c_0} \int_{\Sigma} f'(b_0) d (\nabla_{\mathcal{A}} d \cdot \mathcal{N}) \\ = \int_{\Omega} J d F^6. \end{aligned} \quad (3.8)$$

Note that

$$\int_{\Sigma} (h - f'(b_0) d) (\nabla_{\mathcal{A}} d \cdot \mathcal{N}) = \int_{\Sigma} (h - f'(b_0) d) (\omega_c h + \omega_{b0} d) + \int_{\Sigma} (h - f'(b_0) d) F^7. \quad (3.9)$$

Summarizing all above, we have the desired result.  $\square$

We will mainly use this form to study the estimate of temporal derivative. We apply  $\partial_t$  operator on the unknown variables and let  $(v, q, \xi, h, d) = (\partial_t u, \partial_t p, \partial_t \eta, \partial_t c, \partial_t b)$ . In this situation, we have

$$F^1 = F^{1,1} + F^{1,2} + F^{1,3} + F^{1,4}, \quad (3.10)$$

$$F_i^{1,1} = \partial_t (\partial_t \eta \tilde{L} K) \partial_3 u_i, \quad (3.11)$$

$$F_i^{1,2} = -\partial_t (u_j \mathcal{A}_{jk}) \partial_k u_i + \partial_t \mathcal{A}_{ik} \partial_k p, \quad (3.12)$$

$$F_i^{1,3} = \partial_t \mathcal{A}_{jl} \partial_l (\mathcal{A}_{im} \partial_m u_j + \mathcal{A}_{jm} \partial_m u_i), \quad (3.13)$$

$$F_i^{1,4} = \mathcal{A}_{jk} \partial_k (\partial_t \mathcal{A}_{il} \partial_l u_j + \partial_t \mathcal{A}_{jl} \partial_l u_i), \quad (3.14)$$

$$F^2 = -\partial_t \mathcal{A}_{ij} \partial_j u_i, \quad (3.15)$$

$$F^3 = F^{3,1} + F^{3,2} + F^{3,3}, \quad (3.16)$$

$$F_i^{3,1} = (\eta - p) \partial_t \mathcal{N}_i + (\mathcal{A}_{ik} \partial_k u_j + \mathcal{A}_{jk} \partial_k u_i) \partial_t \mathcal{N}_j + (\partial_t \mathcal{A}_{ik} \partial_k u_j + \partial_t \mathcal{A}_{jk} \partial_k u_i) \mathcal{N}_j, \quad (3.17)$$

$$\begin{aligned} F_i^{3,2} = -\sigma'(\tilde{c}) \partial_t c H \mathcal{N}_i - (\sigma(\tilde{c}) - \sigma_0) \partial_t H \mathcal{N}_i - (\sigma_0 \partial_t H - \sigma_0 \partial_t \Delta_* \eta) \mathcal{N}_i \\ - \sigma(\tilde{c}) H \partial_t \mathcal{N}_i, \end{aligned} \quad (3.18)$$

$$\begin{aligned} F_i^{3,3} = -\frac{\nabla_* \eta \cdot \nabla_* \partial_t \eta}{\sqrt{1 + |\nabla_* \eta|^2}} \sigma'(\tilde{c}) - \sqrt{1 + |\nabla_* \eta|^2} \sigma''(\tilde{c}) \partial_t \tilde{c} (\nabla_{\Gamma} \tilde{c})_i \\ + \sqrt{1 + |\nabla_* \eta|^2} \sigma'(\tilde{c}) v_i (\nu_* \cdot \nabla_*) \partial_t c \\ - \sqrt{1 + |\nabla_* \eta|^2} \sigma'(\tilde{c}) (\partial_t v_i (\nu_* \cdot \nabla_*) \tilde{c} + v_i (\partial_t \nu_* \cdot \nabla_*) \tilde{c}), \end{aligned} \quad (3.19)$$

$$F^4 = \partial_t \nabla_* \eta \cdot u, \quad (3.20)$$

$$F^5 = F^{5,1} + F^{5,2}, \quad (3.21)$$

$$F^{5,1} = \partial_t(-u \cdot \nabla_* c - c \nabla_\Gamma \cdot u + \gamma(\Delta_\Gamma c - \Delta_* c) - c_0(\nabla_\Gamma \cdot u - \nabla_* \cdot u)), \quad (3.22)$$

$$F^{5,2} = -\beta(\partial_t \mathcal{A}_{ik}) \partial_k b \mathcal{N}_i - \beta \mathcal{A}_{ik} \partial_k b (\partial_t \mathcal{N}_i), \quad (3.23)$$

$$F^6 = F^{6,1} + F^{6,2} + F^{6,3}, \quad (3.24)$$

$$F_i^{6,1} = \partial_t(\partial_t \tilde{\eta} \tilde{L} K) \partial_3 b, \quad (3.25)$$

$$F_i^{6,2} = -\partial_t(u_j \mathcal{A}_{jk}) \partial_k b, \quad (3.26)$$

$$F_i^{6,3} = \beta \partial_t \mathcal{A}_{kl} \partial_k \mathcal{A}_{jm} \partial_m b + \beta \mathcal{A}_{jk} \partial_k \partial_t \mathcal{A}_{ml} \partial_m b, \quad (3.27)$$

$$F^7 = F^{7,1} + F^{7,2}, \quad (3.28)$$

$$F^{7,1} = -(\partial_t \mathcal{A}_{ik}) \partial_k b \mathcal{N}_i - \mathcal{A}_{ik} \partial_k b (\partial_t \mathcal{N}_i), \quad (3.29)$$

$$F^{7,2} = \partial_t(\omega(\tilde{c}, \tilde{b}) - \omega_{c0}c - \omega_{b0}b). \quad (3.30)$$

### 3.2 Linear Perturbed Form

Now we consider the second linearized formulation of (1.53), which we call the linear perturbed form due to the lack of geometric terms. Again we phrase the problem for unknowns  $(v, q, \xi, h, d)$  with given forcing terms now called  $Z^1, \dots, Z^7$ , though in later use the unknowns will be related to the solution and its derivative and the forcing terms will have specific forms. The systems is:

$$\begin{cases} \partial_t v - \Delta v + \nabla q = Z^1 & \text{in } \Omega, \\ \nabla \cdot v = Z^2 & \text{in } \Omega, \\ (q I - \mathbb{D}v)e_3 = \xi e_3 - \sigma_0 \Delta_* \xi e_3 - \sigma'_0 \nabla_* h + Z^3 & \text{on } \Sigma, \\ v = 0 & \text{on } \Sigma_b, \\ \partial_t \xi = v_3 + Z^4 & \text{on } \Sigma, \\ \partial_t h + c_0 \nabla_* \cdot v = \gamma \Delta_* h - \beta \partial_3 d + Z^5 & \text{on } \Sigma, \\ \partial_t d = \beta \Delta d + Z^6 & \text{in } \Omega, \\ \partial_3 d = \omega_{c0} h + \omega_{b0} d + Z^7 & \text{on } \Sigma, \\ \partial_3 d = 0 & \text{on } \Sigma_b. \end{cases} \quad (3.31)$$

The next result records the energy-dissipation equation associated to this problem.

**Lemma 3.2** *We have*

$$\begin{aligned} \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} |v|^2 + \frac{1}{2} \int_{\Sigma} |\xi|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \xi|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |d|^2 + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |h|^2 \right) \\ + \left( \frac{1}{2} \int_{\Omega} \|\mathbb{D}v\|^2 + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} |\nabla d|^2 + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* h|^2 \right) \\ + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} (h - f'(b_0)d)(\omega_{c0}h + \omega_{b0}d) \\ = \int_{\Omega} v \cdot Z^1 + \int_{\Omega} (q Z^2 - v \cdot \nabla Z^2) - \int_{\Sigma} v \cdot Z^3 + \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) Z^4 \\ + \frac{-\sigma'_0}{c_0} \int_{\Sigma} h Z^5 + \frac{-\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} d Z^6 + \frac{\beta \sigma'_0}{c_0} \int_{\Sigma} (h - f'(b_0)d) Z^7. \end{aligned} \quad (3.32)$$

*Proof* From the Navier-Stokes equations, we have

$$\partial_t v + \nabla \cdot \mathcal{S}(v, q) = \partial_t v - \Delta v + \nabla q - \nabla Z^2 = Z^1 - \nabla Z^2. \quad (3.33)$$

Multiplying  $v$  on both sides of Navier-Stokes equations and integrating over  $\Omega$ , we have

$$\begin{aligned} \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} |v|^2 \right) + \frac{1}{2} \int_{\Omega} |\mathbb{D}v|^2 + \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) v_3 - \sigma'_0 \int_{\Sigma} \nabla_* h \cdot v \\ = \int_{\Omega} v \cdot Z^1 + \int_{\Omega} (q Z^2 - v \cdot \nabla Z^2) - \int_{\Sigma} v \cdot Z^3. \end{aligned} \quad (3.34)$$

Also, we have

$$\begin{aligned} \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) v_3 &= \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) (\partial_t \xi - Z^4) \\ &= \frac{d}{dt} \left( \frac{1}{2} \int_{\Sigma} |\xi|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \xi|^2 \right) - \int_{\Sigma} (\xi - \sigma_0 \Delta_* \xi) Z^4. \end{aligned} \quad (3.35)$$

In surface surfactant equation, multiplying  $\frac{-\sigma'_0}{c_0} h$  on both sides and integrating over  $\Sigma$ , we have

$$\frac{d}{dt} \left( \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |h|^2 \right) + \sigma'_0 \int_{\Sigma} \nabla_* h \cdot v + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* h|^2 + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} h \partial_3 d = \frac{-\sigma'_0}{c_0} \int_{\Sigma} h Z^5. \quad (3.36)$$

In surface surfactant equation, multiplying  $\frac{-\beta \sigma'_0 f'(b_0)}{c_0} d$  on both sides and integrating over  $\Omega$ , we have

$$\frac{d}{dt} \left( \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |d|^2 \right) + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} |\nabla d|^2 - \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} f'(b_0) d \partial_3 d = \int_{\Omega} d Z^6. \quad (3.37)$$

Note that

$$\int_{\Sigma} (h - f'(b_0) d) \partial_3 d = \int_{\Sigma} (h - f'(b_0) d) (\omega_{c0} h + \omega_{b0} d) + \int_{\Sigma} (h - f'(b_0) d) Z^7. \quad (3.38)$$

Summarizing all above, we have the desired result.  $\square$

We will mainly use this form to study the estimate of spatial derivative. In the system (1.53), we have

$$\begin{cases} \partial_t u - \Delta u + \nabla p = G^1 & \text{in } \Omega, \\ \nabla \cdot u = G^2 & \text{in } \Omega, \\ (pI - \mathbb{D}u)e_3 = \eta e_3 - \sigma_0 \Delta_* \eta e_3 - \sigma'_0 \nabla_* c + G^3 & \text{on } \Sigma, \\ u = 0 & \text{on } \Sigma_b, \\ \partial_t \xi = u_3 + G^4 & \text{on } \Sigma, \\ \partial_t c + c_0 \nabla_* \cdot u = \Delta_* c - \partial_3 b + G^5 & \text{on } \Sigma, \\ \partial_t b = \Delta b + G^6 & \text{in } \Omega, \\ \partial_3 b = \omega_{c0} c + \omega_{b0} b + G^7 & \text{on } \Sigma, \\ \partial_3 b = 0 & \text{on } \Sigma_b, \end{cases} \quad (3.39)$$

where

$$G^1 = G^{1,1} + G^{1,2} + G^{1,3} + G^{1,4} + G^{1,5}, \quad (3.40)$$

$$G_i^{1,1} = (\delta_{ij} - \mathcal{A}_{ij})\partial_j p, \quad (3.41)$$

$$G_i^{1,2} = u_j \mathcal{A}_{jk} \partial_k u_i, \quad (3.42)$$

$$G_i^{1,3} = (K^2(1 + A^2 + B^2) - 1)\partial_{33}u_i - 2AK\partial_{13}u_i - 2BK\partial_{23}u_i, \quad (3.43)$$

$$G_i^{1,4} = (-K^3(1 + A^2 + B^2)\partial_3 J + AK^2(\partial_1 J + \partial_3 A) + BK^2(\partial_2 J + \partial_3 B) - K(\partial_1 A + \partial_2 B))\partial_3 u_i, \quad (3.44)$$

$$G_i^{1,5} = \partial_t \tilde{\eta} \tilde{L} K \partial_3 u_i, \quad (3.45)$$

$$G^2 = AK\partial_3 u_1 + BK\partial_3 u_2 + (1 - K)\partial_3 u_3, \quad (3.46)$$

$$G^3 = G^{3,1} + G^{3,2} + G^{3,3} + G^{3,4}, \quad (3.47)$$

$$G^{3,1} = \partial_1 \eta \begin{pmatrix} p - \eta - 2(\partial_1 u_1 - AK\partial_3 u_1) \\ -\partial_2 u_1 - \partial_1 u_2 + BK\partial_3 u_1 + AK\partial_3 u_2 \\ -\partial_1 u_3 - K\partial_3 u_1 + AK\partial_3 u_3 \end{pmatrix} \quad (3.48)$$

$$+ \partial_2 \eta \begin{pmatrix} -\partial_2 u_1 - \partial_1 u_2 + BK\partial_3 u_1 + AK\partial_3 u_2 \\ p - \eta - 2(\partial_2 u_2 - BK\partial_3 u_2) \\ -\partial_2 u_3 - K\partial_3 u_2 + BK\partial_3 u_3 \end{pmatrix} + \begin{pmatrix} (K - 1)\partial_3 u_1 + AK\partial_3 u_3 \\ (K - 1)\partial_3 u_2 + BK\partial_3 u_3 \\ 2(K - 1)\partial_3 u_3 \end{pmatrix},$$

$$G^{3,2} = (\sigma(\tilde{c}) - \sigma_0)\Delta_* \eta e_3 + \sigma(\tilde{c})(H - \Delta_* \eta)\mathcal{N} + \sigma(\tilde{c})\Delta_* \eta(\mathcal{N} - e_3), \quad (3.49)$$

$$G^{3,3} = (\sqrt{1 + |\nabla_* \eta|^2} - 1)\sigma'(\tilde{c})\nabla_* c + (\sigma'(\tilde{c}) - \sigma'_0)\nabla_* c \\ + \sqrt{1 + |\nabla_* \eta|^2}\sigma'(\tilde{c})(\nabla_\Gamma c - \nabla_* c), \quad (3.50)$$

$$G^{3,4} = \sigma'(\tilde{c})v_* \cdot \nabla_* c e_3, \quad (3.51)$$

$$G^4 = -\partial_1 \eta u_1 - \partial_2 \eta u_2, \quad (3.52)$$

$$G^5 = G^{5,1} + G^{5,2}, \quad (3.53)$$

$$G^{5,1} = -u \cdot \nabla_* c - c \nabla_\Gamma \cdot u + \gamma(\Delta_\Gamma c - \Delta_* c) - c_0(\nabla_\Gamma \cdot u - \nabla_* \cdot u), \quad (3.54)$$

$$G^{5,2} = \beta(\partial_1 b - AK\partial_3 b)\partial_1 \eta + \beta(\partial_2 b - BK\partial_3 b)\partial_2 \eta + \beta(K - 1)\partial_3 b, \quad (3.55)$$

$$G^6 = G^{6,1} + G^{6,2} + G^{6,3} + G^{6,4}, \quad (3.56)$$

$$G_i^{6,1} = u_j \mathcal{A}_{jk} \partial_k b, \quad (3.57)$$

$$G_i^{6,2} = \beta(K^2(1 + A^2 + B^2) - 1)\partial_{33}b - 2AK\partial_{13}b - 2BK\partial_{23}b, \quad (3.58)$$

$$G_i^{6,3} = \beta(-K^3(1 + A^2 + B^2)\partial_3 J + AK^2(\partial_1 J + \partial_3 A) + BK^2(\partial_2 J + \partial_3 B) - K(\partial_1 A + \partial_2 B))\partial_3 b, \quad (3.59)$$

$$G_i^{6,4} = \partial_t \tilde{\eta} \tilde{L} K \partial_3 b, \quad (3.60)$$

$$G^7 = G^{7,1} + G^{7,2}, \quad (3.61)$$

$$G^{7,1} = (\partial_1 b - AK\partial_3 b)\partial_1 \eta + (\partial_2 b - BK\partial_3 b)\partial_2 \eta + (K - 1)\partial_3 b, \quad (3.62)$$

$$G^{7,1} = \omega(\tilde{c}, \tilde{b}) - \omega_{c0}c - \omega_{b0}b. \quad (3.63)$$

## 4 Preliminary Estimates

In this section, we mainly cover some preliminary estimates in energy-dissipation estimates, nonlinear estimates and comparison estimates.

### 4.1 A Basic Lemma

The following lemma provides basic estimates for terms that appear frequently in our analysis.

**Lemma 4.1** *There exists a universal  $0 < \delta \leq 1$  such that if  $\|\eta\|_{H^{\frac{5}{2}}(\Sigma)} \leq \delta$ , then the following holds*

(1) *We have*

$$\|J - 1\|_{L^\infty(\Omega)} + \|A\|_{L^\infty(\Omega)} + \|B\|_{L^\infty(\Omega)} \leq \frac{1}{4}, \quad (4.1)$$

$$\|K\|_{L^\infty(\Omega)} + \|\mathcal{A}\|_{L^\infty(\Omega)} \leq 1. \quad (4.2)$$

(2) *The mapping  $\Phi$  is a diffeomorphism from  $\Omega$  to  $\Omega(t)$ , and  $\|\eta\|_{L^\infty(\Sigma)} < L_3/2$ .*

(3) *For all  $v \in H^1(\Omega)$  with  $v = 0$  on  $\Sigma_b$ , we have*

$$\int_{\Omega} |\mathbb{D}v|^2 \leq \int_{\Omega} J|\mathbb{D}_{\mathcal{A}}v|^2 + C_0\sqrt{\mathcal{E}} \int_{\Omega} |\mathbb{D}v|^2. \quad (4.3)$$

(4) *For all  $\psi \in H^1(\Omega)$  we have*

$$\int_{\Omega} |\nabla \psi|^2 \leq \int_{\Omega} J|\nabla_{\mathcal{A}}\psi|^2 + C_0\sqrt{\mathcal{E}} \int_{\Omega} |\nabla \psi|^2. \quad (4.4)$$

(5) *We have the pointwise bounds*

$$-\frac{c_0}{2} \leq c \leq \frac{c_0}{2}, \quad \frac{c_0}{2} \leq \tilde{c} \leq \frac{3c_0}{2}, \quad (4.5)$$

$$-\frac{b_0}{2} \leq b \leq \frac{b_0}{2}, \quad \frac{b_0}{2} \leq \tilde{b} \leq \frac{3b_0}{2}. \quad (4.6)$$

*Proof* See the proof of [15, Lemma 4.1] with obvious modifications. The last inequality involving  $b_0$  can be proved as for  $c_0$ .  $\square$

### 4.2 Temporal Estimates

Now we turn our attention to estimating various nonlinearities that appear when we consider the temporally-differentiated problem.

**Lemma 4.2** *Suppose  $\mathcal{E} \leq \delta \ll 1$ . Let  $F^1, \dots, F^7$  be as defined in Sect. 3.1. We have*

$$\|JF^1\|_{H^0(\Omega)} + \|F^3\|_{H^0(\Sigma)} + \|F^4\|_{H^0(\Sigma)} + \|JF^6\|_{H^0(\Omega)} \leq \sqrt{\mathcal{E}}\sqrt{\mathcal{D}}, \quad (4.7)$$

$$\left| \int_{\Omega} JF^2 \partial_t p - \partial_t \int_{\Omega} JF^2 p \right| \lesssim \sqrt{\mathcal{E}}\mathcal{D}, \quad \left| \int_{\Omega} JF^2 p \right| \lesssim \mathcal{E}^{\frac{3}{2}}, \quad (4.8)$$

$$\left| \int_{\Sigma} F^5 \partial_t c \right| \lesssim \sqrt{\mathcal{E}} \mathcal{D}, \quad (4.9)$$

$$\| F^7 \|_{H^0(\Sigma)} \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (4.10)$$

*Proof* Equations (4.7) and (4.8) can be proved as [15, Theorem 4.2]. For (4.9), we already know the estimate of  $F^{5,1}$ , so we only need to bound

$$F^{5,2} = -\beta (\partial_t \mathcal{A}_{ik}) \partial_k b \mathcal{N}_i - \beta \mathcal{A}_{ik} \partial_k b (\partial_t \mathcal{N}_i). \quad (4.11)$$

We directly estimate

$$\begin{aligned} \| F^{5,2} \|_{H^0(\Sigma)} &\lesssim \| \partial_t \nabla \bar{\eta} \|_{H^0(\Sigma)} \| \nabla b \|_{L^\infty(\Sigma)} (1 + \| \nabla_* \eta \|_{L^\infty(\Sigma)}) \\ &\quad + \| \partial_t \nabla_* \eta \|_{H^0(\Sigma)} \| \nabla b \|_{L^\infty(\Sigma)} (1 + \| \nabla \bar{\eta} \|_{L^\infty(\Sigma)}) \\ &\lesssim \| \partial_t \bar{\eta} \|_{H^{\frac{3}{2}}(\Omega)} \| \nabla b \|_{H^{\frac{3}{2}}(\Sigma)} (1 + \| \nabla_* \eta \|_{H^{\frac{3}{2}}(\Sigma)}) \\ &\quad + \| \partial_t \eta \|_{H^1(\Sigma)} \| \nabla b \|_{H^{\frac{3}{2}}(\Sigma)} (1 + \| \nabla \bar{\eta} \|_{H^{\frac{3}{2}}(\Sigma)}) \\ &\lesssim \| \partial_t \eta \|_{H^1(\Sigma)} \| b \|_{H^3(\Omega)} (1 + \| \eta \|_{H^{\frac{5}{2}}(\Sigma)}) \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (4.12)$$

Thus,

$$\left| \int_{\Sigma} F^{5,2} \partial_t c \right| \lesssim \| F^{5,2} \|_{H^0(\Sigma)} \| \partial_t c \|_{H^0(\Sigma)} \lesssim \sqrt{\mathcal{E}} \mathcal{D}. \quad (4.13)$$

For (4.10), we need to estimate

$$F^7 = F^{7,1} + F^{7,2}, \quad (4.14)$$

$$F^{7,1} = -(\partial_t \mathcal{A}_{ik}) \partial_k b \mathcal{N}_i - \mathcal{A}_{ik} \partial_k b (\partial_t \mathcal{N}_i), \quad (4.15)$$

$$F^{7,2} = \partial_t (\omega(\tilde{c}, \tilde{b}) - \omega_{c0} c - \omega_{b0} b). \quad (4.16)$$

It is easy to see

$$\| F^{7,1} \|_{H^0(\Sigma)} = \beta \| F^{5,2} \|_{H^0(\Sigma)} \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (4.17)$$

On the other hand, using Taylor expansion as in (1.43), we obtain

$$\begin{aligned} \| F^{7,2} \|_{H^0(\Sigma)} &\lesssim (\| \partial_t c \|_{H^0(\Sigma)} + \| \partial_t b \|_{H^0(\Sigma)}) (\| c \|_{L^\infty(\Sigma)} + \| b \|_{L^\infty(\Sigma)} + \| c \|_{L^\infty(\Sigma)}^2 + \| b \|_{L^\infty(\Sigma)}^2) \\ &\lesssim (\| \partial_t c \|_{H^0(\Sigma)} + \| \partial_t b \|_{H^1(\Omega)}) (\| c \|_{H^{\frac{3}{2}}(\Sigma)} + \| b \|_{H^2(\Omega)} + \| c \|_{H^{\frac{3}{2}}(\Sigma)}^2 + \| b \|_{H^2(\Omega)}^2) \\ &\lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (4.18)$$

Therefore, we know

$$\| F^7 \|_{H^0(\Sigma)} \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (4.19)$$

□

### 4.3 Spatial Estimates

Now we control various nonlinearities that appear in the spatially-differentiated problem.

**Lemma 4.3** Suppose  $\mathcal{E} \leq \delta \ll 1$ . Let  $G^1, \dots, G^7$  be as defined in Sect. 3.2. We have

$$\begin{aligned} & \|G^1\|_{H^1(\Omega)} + \|G^2\|_{H^2(\Omega)} + \|G^3\|_{H^{\frac{3}{2}}(\Sigma)} + \|G^4\|_{H^{\frac{5}{2}}(\Sigma)} + \|G^5\|_{H^1(\Sigma)} \\ & + \|G^6\|_{H^1(\Omega)} + \|G^7\|_{H^{\frac{3}{2}}(\Sigma)} \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}, \end{aligned} \quad (4.20)$$

$$\begin{aligned} & \|G^1\|_{H^0(\Omega)} + \|G^2\|_{H^1(\Omega)} + \|G^3\|_{H^{\frac{1}{2}}(\Sigma)} + \|G^4\|_{H^{\frac{3}{2}}(\Sigma)} + \|G^5\|_{H^0(\Sigma)} \\ & + \|G^6\|_{H^0(\Omega)} + \|G^7\|_{H^{\frac{1}{2}}(\Sigma)} \lesssim \mathcal{E}, \end{aligned} \quad (4.21)$$

$$\|\partial_t G^4\|_{H^{\frac{1}{2}}(\Sigma)} \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (4.22)$$

*Proof* In (4.20), the estimates of  $G^1, G^2, G^3, G^4, \partial_t G^4, G^{5,1}$  and  $G^6$  can be proved as [15, Theorem 4.3]. It suffices to estimate  $G^7 = G^{7,1} + G^{7,2}$  since  $G^{5,2} = \beta G^{7,1}$ . By (B.5) with  $r = s_1 = s_2 = \frac{3}{2}$ , we have

$$\begin{aligned} \|G^{7,1}\|_{H^{\frac{3}{2}}(\Sigma)} &= \|(\partial_1 b - AK\partial_3 b)\partial_1 \eta + (\partial_2 b - BK\partial_3 b)\partial_2 \eta + (K-1)\partial_3 b\|_{H^{\frac{3}{2}}(\Sigma)} \\ &\lesssim (1 + \|\nabla \bar{\eta}\|_{H^{\frac{3}{2}}(\Sigma)}) \|\nabla b\|_{H^{\frac{3}{2}}(\Sigma)} \|\nabla \eta\|_{H^{\frac{3}{2}}(\Sigma)} + \|\nabla \bar{\eta}\|_{H^{\frac{3}{2}}(\Sigma)} \|\nabla b\|_{H^{\frac{3}{2}}(\Sigma)} \\ &\lesssim \|\eta\|_{H^{\frac{5}{2}}(\Sigma)} \|b\|_{H^3(\Omega)} \lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (4.23)$$

Using Taylor expansion as (1.43), we obtain

$$\begin{aligned} \|G^{7,2}\|_{H^{\frac{3}{2}}(\Sigma)} &= \|\omega(\tilde{c}, \tilde{b}) - \omega_{c0}c - \omega_{b0}b\|_{H^{\frac{3}{2}}(\Sigma)} \\ &\lesssim \|\omega\|_{C^4} (\|c\|_{H^3(\Sigma)} + \|b\|_{H^3(\Omega)}) (\|c\|_{H^2(\Sigma)} + \|b\|_{H^2(\Omega)})^2 \left| \int_0^1 s(1-s) ds \right| \\ &\lesssim \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (4.24)$$

Similarly, in (4.21), the estimates of  $G^1, G^2, G^3, G^4, G^{5,1}$  and  $G^6$  can be proved as before. By (B.6) with  $r = s_1 = \frac{1}{2}$  and  $s_2 = 2$ , we estimate  $G^7$  as

$$\begin{aligned} \|G^{7,1}\|_{H^{\frac{1}{2}}(\Sigma)} &= \|(\partial_1 b - AK\partial_3 b)\partial_1 \eta + (\partial_2 b - BK\partial_3 b)\partial_2 \eta + (K-1)\partial_3 b\|_{H^{\frac{1}{2}}(\Sigma)} \\ &\lesssim (1 + \|\nabla \bar{\eta}\|_{H^2(\Sigma)}) \|\nabla b\|_{H^{\frac{1}{2}}(\Sigma)} \|\nabla \eta\|_{H^2(\Sigma)} + \|\nabla \bar{\eta}\|_{H^2(\Sigma)} \|\nabla b\|_{H^{\frac{1}{2}}(\Sigma)} \\ &\lesssim \|\eta\|_{H^3(\Sigma)} \|b\|_{H^2(\Omega)} \lesssim \mathcal{E}. \end{aligned} \quad (4.25)$$

Also, using Taylor expansion as (1.43), we get

$$\begin{aligned} \|G^{7,2}\|_{H^{\frac{1}{2}}(\Sigma)} &= \|\omega(\tilde{c}, \tilde{b}) - \omega_{c0}c - \omega_{b0}b\|_{H^{\frac{1}{2}}(\Sigma)} \\ &\lesssim \|\omega\|_{C^3} (\|c\|_{H^2(\Sigma)} + \|b\|_{H^2(\Omega)})^3 \left| \int_0^1 s(1-s) ds \right| \lesssim \mathcal{E}. \end{aligned} \quad (4.26)$$

□

## 5 Energy-Dissipation Estimates

In this section, we begin to investigate the energy-dissipation structure for temporal and horizontal derivatives. Define

$$\mathcal{F} = \int_{\Omega} J F^2 p + \frac{1}{2} \int_{\Omega} (J-1) |\partial_t u|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} (J-1) |\partial_t b|^2. \quad (5.1)$$

**Theorem 5.1** *Suppose*

$$\sup_{t \in [0, T]} \mathcal{E}(t) \leq \delta \ll 1, \quad \int_0^T \mathcal{D}(t) dt < \infty. \quad (5.2)$$

We have

$$\frac{d}{dt} (\bar{\mathcal{E}} - \mathcal{F}) + \bar{\mathcal{D}} \lesssim \sqrt{\mathcal{E}} \mathcal{D}. \quad (5.3)$$

*Proof* We apply operator  $\partial^\alpha$  for  $\alpha = (\alpha_0, \alpha_1, \alpha_2, 0)$  to the system (1.53) and estimate each term. We divide it into several steps:

Step 1: Temporal derivatives.

We apply  $\partial_t$  to system (1.53), which yields  $F^1, \dots, F^7$  as defined in Sect. 3.1. Then we have from Lemma 3.1 that

$$\begin{aligned} \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} J |\partial_t u|^2 + \frac{1}{2} \int_{\Sigma} |\partial_t \eta|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \partial_t \eta|^2 \right. \\ \left. + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} J |\partial_t b|^2 + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |\partial_t c|^2 \right) \\ + \left( \frac{1}{2} \int_{\Omega} J |\mathbb{D}_A \partial_t u|^2 + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} J |\nabla_A \partial_t b|^2 + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* \partial_t c|^2 \right) \\ + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial_t c - f'(b_0) \partial_t b) (\omega_{c0} \partial_t c + \omega_{b0} \partial_t b) \\ = \int_{\Omega} J \partial_t u \cdot F^1 + \int_{\Omega} J \partial_t p F^2 - \int_{\Sigma} \partial_t u \cdot F^3 + \int_{\Sigma} (\partial_t \eta - \sigma_0 \Delta_* \partial_t \eta) F^4 \\ + \frac{-\sigma'_0}{c_0} \int_{\Sigma} \partial_t c F^5 + \frac{-\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} J \partial_t b F^6 + \frac{\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial_t c + f'(b_0) \partial_t b) F^7. \quad (5.4) \end{aligned}$$

For the first and fourth terms inside the time derivative, we write

$$\begin{aligned} \frac{1}{2} \int_{\Omega} J |\partial_t u|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} J |\partial_t b|^2 = \frac{1}{2} \int_{\Omega} |\partial_t u|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |\partial_t b|^2 \\ + \frac{1}{2} \int_{\Omega} (J-1) |\partial_t u|^2 \\ + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} (J-1) |\partial_t b|^2 \quad (5.5) \end{aligned}$$

in order to eliminate the  $J$  factors and absorb the error terms with  $(J - 1)$  into  $\mathcal{F}$ . We may estimate the right-hand side of (5.4) as

$$\begin{aligned}
 RHS &\lesssim \|\partial_t u\|_{H^0(\Omega)} \|JF^1\|_{H^0(\Omega)} + \frac{d}{dt} \int_{\Omega} JpF^2 + \left| \int_{\Sigma} p(\partial_t JF^2 + J\partial_t F^2) \right| \\
 &\quad + \|\partial_t u\|_{H^0(\Sigma)} \|F^3\|_{H^0(\Sigma)} \\
 &\quad + \|\partial_t \eta - \sigma_0 \Delta_* \partial_t \eta\|_{H^0(\Sigma)} \|F^4\|_{H^0(\Sigma)} + \left| \int_{\Sigma} \partial_t c F^5 \right| + \|\partial_t b\|_{H^0(\Omega)} \|JF^6\|_{H^0(\Omega)} \\
 &\quad + \|\partial_t c + f'(b_0) \partial_t b\|_{H^0(\Sigma)} \|F^7\|_{H^0(\Sigma)}. \tag{5.6}
 \end{aligned}$$

Combining this estimate with Lemmas 4.1 and 4.2 then shows that

$$\begin{aligned}
 &\frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} |\partial_t u|^2 + \frac{1}{2} \int_{\Sigma} |\partial_t \eta|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \partial_t \eta|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |\partial_t b|^2 \right. \\
 &\quad \left. + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |\partial_t c|^2 + \mathcal{F} \right) \\
 &\quad + \left( \frac{1}{2} \int_{\Omega} |\mathbb{D} \partial_t u|^2 + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} |\nabla \partial_t b|^2 + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* \partial_t c|^2 \right) \\
 &\quad + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial_t c - f'(b_0) \partial_t b) (\omega_{c0} \partial_t c + \omega_{b0} \partial_t b) \lesssim \sqrt{\mathcal{E}} \mathcal{D} \tag{5.7}
 \end{aligned}$$

Step 2: Spatial derivative—second order.

We apply  $\partial^\alpha$  with  $\alpha_0 = 0$  and  $|\alpha| = 2$  to system (1.53), which produces forcing terms  $G^1, \dots, G^7$  as defined in Sect. 3.2. Then from Lemma 3.2 we have

$$\begin{aligned}
 &\frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} |\partial^\alpha u|^2 + \frac{1}{2} \int_{\Sigma} |\partial^\alpha \eta|^2 + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \partial^\alpha \eta|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |\partial^\alpha b|^2 \right. \\
 &\quad \left. + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |\partial^\alpha c|^2 \right) \\
 &\quad + \left( \frac{1}{2} \int_{\Omega} |\mathbb{D} \partial^\alpha u|^2 + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} |\nabla \partial^\alpha b|^2 + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* \partial^\alpha c|^2 \right) \\
 &\quad + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial^\alpha c - f'(b_0) \partial^\alpha b) (\omega_{c0} \partial^\alpha c + \omega_{b0} \partial^\alpha b) \\
 &= \int_{\Omega} \partial^\alpha u \cdot \partial^\alpha G^1 + \int_{\Omega} (\partial^\alpha p \partial^\alpha G^2 - \partial^\alpha u \cdot \nabla \partial^\alpha G^2) - \int_{\Sigma} \partial^\alpha u \cdot \partial^\alpha G^3 \\
 &\quad + \int_{\Sigma} (\partial^\alpha \eta - \sigma_0 \Delta_* \partial^\alpha \eta) \partial^\alpha G^4 + \frac{-\sigma'_0}{c_0} \int_{\Sigma} \partial^\alpha c \partial^\alpha G^5 \\
 &\quad + \frac{-\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} \partial^\alpha b \partial^\alpha G^6 + \frac{\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial^\alpha c + f'(b_0) \partial^\alpha b) \partial^\alpha G^7. \tag{5.8}
 \end{aligned}$$

Assume  $\partial^\alpha = \partial^\vartheta \partial^\varphi$  for  $|\vartheta| = |\varphi| = 1$ . We may integrate by parts on the right-hand side to obtain

$$\begin{aligned}
 RHS = & - \int_{\Omega} \partial^{\alpha+\vartheta} u \cdot \partial^\varphi G^1 + \int_{\Omega} (\partial^\alpha p \partial^\alpha G^2 - \partial^{\alpha+\vartheta} u \cdot \nabla \partial^\varphi G^2) - \int_{\Sigma} \partial^\alpha u \cdot \partial^\alpha G^3 \\
 & + \int_{\Sigma} (\partial^\varphi \eta - \sigma_0 \Delta_* \partial^\varphi \eta) \partial^{\alpha+\vartheta} G^4 + \frac{-\sigma'_0}{c_0} \int_{\Sigma} \partial^{\alpha+\vartheta} c \partial^\varphi G^5 \\
 & + \frac{-\sigma'_0 f'(b_0)}{c_0} \int_{\Omega} \partial^{\alpha+\vartheta} b \partial^\varphi G^6 + \frac{\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial^\alpha c + f'(b_0) \partial^\alpha b) \partial^\alpha G^7 \\
 \lesssim & \|u\|_{H^3(\Omega)} \|G^1\|_{H^1(\Omega)} + \|p\|_{H^2(\Omega)} \|G^2\|_{H^2(\Omega)} \\
 & + \|u\|_{H^3(\Omega)} \|G^2\|_{H^2(\Omega)} + \|u\|_{H^{\frac{5}{2}}(\Sigma)} \|G^3\|_{H^{\frac{3}{2}}(\Sigma)} \\
 & + \|\eta\|_{H^{\frac{7}{2}}(\Sigma)} \|G^4\|_{H^{\frac{5}{2}}(\Sigma)} + \|c\|_{H^3(\Sigma)} \|G^5\|_{H^1(\Sigma)} + \|b\|_{H^3(\Omega)} \|G^6\|_{H^1(\Omega)} \\
 & + (\|c\|_{H^{\frac{5}{2}}(\Sigma)} + \|b\|_{H^3(\Omega)}) \|G^7\|_{H^{\frac{3}{2}}(\Sigma)}. \tag{5.9}
 \end{aligned}$$

We then employ Lemma 4.3 to deduce that

$$\begin{aligned}
 & \frac{d}{dt} \left( \frac{1}{2} \int_{\Omega} |\partial^\alpha u|^2 + \frac{1}{2} \int_{\Sigma} |\partial^\alpha \eta|^2 \right. \\
 & \quad \left. + \frac{\sigma_0}{2} \int_{\Sigma} |\nabla_* \partial^\alpha \eta|^2 + \frac{-\sigma'_0 f'(b_0)}{2c_0} \int_{\Omega} |\partial^\alpha b|^2 + \frac{-\sigma'_0}{2c_0} \int_{\Sigma} |\partial^\alpha c|^2 \right) \\
 & \quad + \left( \frac{1}{2} \int_{\Omega} |\mathbb{D} \partial^\alpha u|^2 + \frac{-\beta \sigma'_0 f'(b_0)}{c_0} \int_{\Omega} |\nabla \partial^\alpha b|^2 + \frac{-\gamma \sigma'_0}{c_0} \int_{\Sigma} |\nabla_* \partial^\alpha c|^2 \right) \\
 & \quad + \frac{-\beta \sigma'_0}{c_0} \int_{\Sigma} (\partial^\alpha c - f'(b_0) \partial^\alpha b) (\omega_{c0} \partial^\alpha c + \omega_{b0} \partial^\alpha b) \lesssim \sqrt{\mathcal{E}} \mathcal{D} \tag{5.10}
 \end{aligned}$$

whenever  $|\alpha| = 2$  and  $\alpha_0 = 0$ .

Step 3: Spatial derivative—first order and zeroth order.

We apply  $\partial^\alpha$  with  $\alpha_0 = 0$  and  $|\alpha| \leq 1$  to system (1.53). An argument simpler than the one used in Step 2 shows that (5.10) holds also for this range of  $\alpha$ .

Step 4: Conclusion.

To complete the estimate we sum (5.7) with (5.10), with the latter applied for all  $|\alpha| \leq 2$  and  $\alpha_0 = 0$ .  $\square$

## 6 Comparison Estimates

In this section, we turn to the crucial comparison estimate between horizontal quantities and full quantities.

### 6.1 Preliminaries

The next result is the key to proving coercivity of the dissipation over the energy space. It is a Poincaré-type inequality for pairs of functions.

**Lemma 6.1** For  $C \in H^1(\Sigma)$  and  $B \in H^1(\Omega)$  satisfying

$$\int_{\Sigma} C + \int_{\Omega} B = 0, \quad (6.1)$$

we have

$$\|C\|_{H^0(\Sigma)} + \|B\|_{H^0(\Omega)} \lesssim \|\nabla_* C\|_{H^0(\Sigma)} + \|\nabla B\|_{H^0(\Omega)} + \|C - f'(b_0)B\|_{H^0(\Sigma)}. \quad (6.2)$$

*Proof* Assume this claim is not true. Then we can find  $\{C_n, B_n\}_{n=1}^{\infty}$  satisfying

$$\int_{\Sigma} C_n + \int_{\Omega} B_n = 0, \quad (6.3)$$

such that

$$\|C_n\|_{H^0(\Sigma)} + \|B_n\|_{H^0(\Omega)} = 1, \quad (6.4)$$

and

$$\|\nabla_* C_n\|_{H^0(\Sigma)} + \|\nabla B_n\|_{H^0(\Omega)} + \|C_n - f'(b_0)B_n\|_{H^0(\Sigma)} \leq \frac{1}{n}. \quad (6.5)$$

This implies

$$\|C_n\|_{H^1(\Sigma)} \leq 2, \quad \|B_n\|_{H^1(\Omega)} \leq 2. \quad (6.6)$$

Thus, we can extract weakly convergent subsequence

$$C_n \rightharpoonup C_0 \quad \text{in } H^1(\Sigma), \quad (6.7)$$

$$B_n \rightharpoonup B_0 \quad \text{in } H^1(\Omega). \quad (6.8)$$

Also, by compactly embedding theorem and trace theorem, we know

$$C_n \rightarrow C_0 \quad \text{in } H^0(\Sigma), \quad (6.9)$$

$$B_n \rightarrow B_0 \quad \text{in } H^0(\Omega), \quad (6.10)$$

$$B_n \rightarrow B_0 \quad \text{in } H^0(\Sigma). \quad (6.11)$$

This implies

$$\int_{\Sigma} C_0 + \int_{\Omega} B_0 = 0, \quad (6.12)$$

$$\|C_0\|_{H^0(\Sigma)} + \|B_0\|_{H^0(\Omega)} = 1, \quad (6.13)$$

and by weak lower semi-continuity,

$$\|\nabla_* C_0\|_{H^0(\Sigma)} + \|\nabla B_0\|_{H^0(\Omega)} + \|C_0 - f'(b_0)B_0\|_{H^0(\Sigma)} \leq 0. \quad (6.14)$$

Naturally, (6.14) yields that  $C_0$  and  $B_0$  are constants satisfying

$$C_0 - f'(b_0)B_0 = 0, \quad |\Sigma|_2 C_0 + |\Omega|_3 B_0 = 0 \quad (6.15)$$

Since  $f'(b_0) > 0$ , the only solution to this linear system is  $C_0 = B_0 = 0$ . This contradicts the normalization condition (6.13). Therefore, the claim is verified.  $\square$

We next use Lemma 6.1 to derive a useful estimate.

**Lemma 6.2** *Assume  $\mathcal{E} \leq \delta \ll 1$ . We have*

$$\left| \int_{\Sigma} c \right| + \left| \int_{\Omega} b \right| \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.16)$$

*Proof* We decompose

$$\begin{aligned} \int_{\Sigma} c &= \int_{\Sigma} ((c_0 + c) - c_0) \\ &= \left( \int_{\Sigma} (c_0 + c) \sqrt{1 + |\nabla_* \eta|^2} - |\Sigma| c_0 \right) + \int_{\Sigma} (c_0 + c) (1 - \sqrt{1 + |\nabla_* \eta|^2}), \end{aligned} \quad (6.17)$$

$$\begin{aligned} \int_{\Omega} b &= \int_{\Omega} ((b_0 + b) - b_0) \\ &= \left( \int_{\Omega} J(b_0 + b) - |\Omega|_3 b_0 \right) + \int_{\Omega} (b_0 + b) (1 - J), \end{aligned} \quad (6.18)$$

which, using the conservation of surfactant mass (1.20), implies

$$\begin{aligned} \int_{\Sigma} c + \int_{\Omega} b &= \left( \int_{\Sigma} (c_0 + c) \sqrt{1 + |\nabla_* \eta|^2} + \int_{\Omega} J(b_0 + b) - |\Sigma|_2 c_0 - |\Omega|_3 b_0 \right) \\ &\quad + \int_{\Sigma} (c_0 + c) (1 - \sqrt{1 + |\nabla_* \eta|^2}) + \int_{\Omega} (b_0 + b) (1 - J) \\ &= \int_{\Sigma} (c_0 + c) (1 - \sqrt{1 + |\nabla_* \eta|^2}) + \int_{\Omega} (b_0 + b) (1 - J). \end{aligned} \quad (6.19)$$

Let  $K_c[\eta] = \int_{\Sigma} (c_0 + c) (1 - \sqrt{1 + |\nabla_* \eta|^2})$  and  $K_b[\eta] = \int_{\Omega} (b_0 + b) (1 - J)$ . Since for  $\mathcal{E} \leq \delta \ll 1$ , Lemma 4.1 guarantees that

$$\frac{c_0}{2} \leq c_0 + c \leq \frac{3c_0}{2}, \quad (6.20)$$

$$\frac{b_0}{2} \leq b_0 + b \leq \frac{3b_0}{2} \quad (6.21)$$

for  $c_0$  and  $b_0$  the positive equilibrium constants, we have

$$|K_c[\eta]| \lesssim \|\nabla_* \eta\|_{H^0(\Sigma)}^2 \lesssim \|\eta\|_{H^1(\Sigma)}^2 \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}, \quad (6.22)$$

$$|K_b[\eta]| \lesssim \|\bar{\eta}\|_{H^1(\Omega)} \lesssim \|\eta\|_{H^{\frac{1}{2}}(\Sigma)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.23)$$

Define

$$C = c - \frac{K_c[\eta]}{|\Sigma|_2}, \quad B = b - \frac{K_b[\eta]}{|\Omega|_3}, \quad (6.24)$$

which yields

$$\int_{\Sigma} C + \int_{\Omega} B = 0. \quad (6.25)$$

Based on Lemma 6.1, we have

$$\|C\|_{H^0(\Sigma)} + \|B\|_{H^0(\Omega)} \lesssim \|\nabla_* C\|_{H^0(\Sigma)} + \|\nabla B\|_{H^0(\Omega)} + \|C - f'(b_0)B\|_{H^0(\Sigma)}. \quad (6.26)$$

Then by the triangle inequality we know

$$\begin{aligned} \|c\|_{H^0(\Sigma)} + \|b\|_{H^0(\Omega)} &\lesssim \|\nabla_* c\|_{H^0(\Sigma)} + \|\nabla b\|_{H^0(\Omega)} + \|c - f'(b_0)b\|_{H^0(\Sigma)} \\ &+ \left\| \frac{K_c[\eta]}{|\Sigma|_2} \right\|_{H^0(\Sigma)} + \left\| \frac{K_b[\eta]}{|\Omega|_3} \right\|_{H^0(\Omega)} \\ &+ \left\| \frac{K_c[\eta]}{|\Sigma|_2} - \frac{f'(b_0)K_b[\eta]}{|\Omega|_3} \right\|_{H^0(\Sigma)} \\ &\lesssim \sqrt{\bar{\mathcal{D}}} + |K_c[\eta]| + |K_b[\eta]|, \\ &\lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}}\sqrt{\mathcal{D}}. \end{aligned} \quad (6.27)$$

By Cauchy's inequality, our desired result naturally follows.  $\square$

## 6.2 Energy Comparison Estimates

Now we use various auxiliary estimates to control the full energy functional in terms of the horizontal energy.

**Theorem 6.3** Suppose  $\mathcal{E} \leq \delta \ll 1$ . We have

$$\mathcal{E} \lesssim \bar{\mathcal{E}} + \mathcal{E}^2. \quad (6.28)$$

*Proof* It suffices to prove

$$\|u\|_{H^2(\Omega)}^2 + \|p\|_{H^1(\Omega)}^2 + \|\partial_t \eta\|_{H^{\frac{3}{2}}(\Sigma)}^2 + \|\partial_t^2 \eta\|_{H^{-\frac{1}{2}}(\Sigma)}^2 + \|b\|_{H^2(\Omega)}^2 \lesssim \bar{\mathcal{E}} + \mathcal{E}^2. \quad (6.29)$$

The previous four terms can be estimated as [15, Theorem 5.2], so we focus on the last term. Now, according to (1.53), we have

$$\begin{cases} -\beta \Delta b = -\partial_t b + G^6 & \text{in } \Omega, \\ \partial_3 b = \omega_{c0} c + \omega_{b0} b + G^7 & \text{on } \Sigma, \\ \partial_3 b = 0 & \text{on } \Sigma_b. \end{cases} \quad (6.30)$$

Based on elliptic estimate in Lemma B.4, we have

$$\begin{aligned} \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^1(\Omega)} &\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|\partial_t c + c_0 \nabla_* \cdot u - \Delta_* c - G^5\|_{H^0(\Sigma)} \\ &+ \|G^7\|_{H^0(\Sigma)} \end{aligned}$$

$$\begin{aligned}
&\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|\partial_t c\|_{H^0(\Sigma)} + \|u\|_{H^1(\Sigma)} + \|c\|_{H^2(\Sigma)} \\
&\quad + \|G^5\|_{H^0(\Sigma)} + \|G^7\|_{H^0(\Sigma)} \\
&\lesssim \sqrt{\bar{\mathcal{E}}} + \mathcal{E}.
\end{aligned} \tag{6.31}$$

Cauchy's inequality implies

$$\left| \int_{\Omega} b \right| \lesssim \|b\|_{H^0(\Omega)} \lesssim \sqrt{\bar{\mathcal{E}}}, \tag{6.32}$$

which further yields

$$\|b\|_{H^1(\Omega)} \lesssim \left| \int_{\Omega} b \right| + \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^1(\Omega)} \lesssim \sqrt{\bar{\mathcal{E}}} + \mathcal{E} \lesssim \sqrt{\bar{\mathcal{E}}} + \mathcal{E}. \tag{6.33}$$

Further, we have

$$\begin{aligned}
\left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^2(\Omega)} &\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|\omega_{c0}c + \omega_{b0}b + G^7\|_{H^{\frac{1}{2}}(\Sigma)} \\
&\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|c\|_{H^{\frac{1}{2}}(\Sigma)} + \|b\|_{H^1(\Omega)} + \|G^7\|_{H^{\frac{1}{2}}(\Sigma)} \\
&\lesssim \sqrt{\bar{\mathcal{E}}} + \mathcal{E}.
\end{aligned} \tag{6.34}$$

Hence, we know

$$\|b\|_{H^2(\Omega)} \lesssim \left| \int_{\Omega} b \right| + \sqrt{\bar{\mathcal{E}}} + \mathcal{E} \lesssim \sqrt{\bar{\mathcal{E}}} + \mathcal{E}. \tag{6.35}$$

□

### 6.3 Dissipation Comparison Estimates

Next we control the full dissipation in terms of the horizontal dissipation.

**Theorem 6.4** *Suppose  $\mathcal{E} \leq \delta \ll 1$ . We have*

$$\mathcal{D} \lesssim \bar{\mathcal{D}} + \mathcal{E}\mathcal{D}. \tag{6.36}$$

*Proof* In the full dissipation

$$\begin{aligned}
\mathcal{D} &= \|u\|_{H^3(\Omega)}^2 + \|\partial_t u\|_{H^1(\Omega)}^2 + \|p\|_{H^2(\Omega)}^2 + \|\eta\|_{H^{\frac{7}{2}}(\Sigma)}^2 + \|\partial_t \eta\|_{H^{\frac{5}{2}}(\Sigma)}^2 + \|\partial_t^2 \eta\|_{H^{\frac{1}{2}}(\Sigma)}^2 \\
&\quad + \|b\|_{H^3(\Omega)}^2 + \|\partial_t b\|_{H^1(\Omega)}^2 + \|c\|_{H^3(\Sigma)}^2 + \|\partial_t c\|_{H^1(\Sigma)}^2,
\end{aligned} \tag{6.37}$$

the estimates of  $(u, p, \eta)$  can be proved as [15, Theorem 5.3], so it suffices to show

$$\|c\|_{H^3(\Sigma)} + \|\partial_t c\|_{H^1(\Sigma)} + \|b\|_{H^3(\Omega)}^2 + \|\partial_t b\|_{H^1(\Omega)}^2 \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}}\sqrt{\bar{\mathcal{D}}}. \tag{6.38}$$

Based on Lemma 6.2 and Poincaré-Wirtinger inequality, we have

$$\|c\|_{H^3(\Sigma)} \lesssim \|\nabla_* c\|_{H^2(\Sigma)} + \left| \int_{\Sigma} c \right| \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.39)$$

Also, we know

$$\left| \int_{\Sigma} \partial_t c \right| \lesssim \|G^5\|_{H^0(\Sigma)} + \|\omega_{c0}c + \omega_{b0}b\|_{H^0(\Sigma)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}, \quad (6.40)$$

since  $\omega_{c0}c + \omega_{b0}b = C_0(c - f'(b_0)b)$  for some constant  $C_0 > 0$ . Hence, by Poincaré-Wirtinger inequality,

$$\|\partial_t c\|_{H^1(\Sigma)} \lesssim \|\nabla_* \partial_t c\|_{H^0(\Sigma)} + \left| \int_{\Sigma} \partial_t c \right| \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.41)$$

Then we need to reconsider the elliptic system (6.30) and try to bound by dissipation. Based on elliptic estimate in Lemma B.4, we have

$$\begin{aligned} \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^1(\Omega)} &\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|\omega_{c0}c + \omega_{b0}b\|_{H^0(\Sigma)} + \|G^7\|_{H^0(\Sigma)} \\ &\lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (6.42)$$

Combining this with Lemma 6.2, we have

$$\|b\|_{H^1(\Omega)} \lesssim \left| \int_{\Omega} b \right| + \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^1(\Omega)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.43)$$

Further, we have

$$\begin{aligned} \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^2(\Omega)} &\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|\omega_{c0}c + \omega_{b0}b + G^7\|_{H^{\frac{1}{2}}(\Sigma)} \\ &\lesssim \|\partial_t b\|_{H^0(\Omega)} + \|G^6\|_{H^0(\Omega)} + \|c\|_{H^{\frac{1}{2}}(\Sigma)} + \|b\|_{H^1(\Omega)} + \|G^7\|_{H^{\frac{1}{2}}(\Sigma)} \\ &\lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (6.44)$$

Similarly, we obtain

$$\|b\|_{H^2(\Omega)} \lesssim \left| \int_{\Omega} b \right| + \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^2(\Omega)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.45)$$

On the other hand, by Cauchy's inequality, we know

$$\begin{aligned} \left| \int_{\Omega} \partial_t b \right| &\lesssim \|G^6\|_{H^0(\Omega)} + \|\omega_{c0}c + \omega_{b0}b + G^7\|_{H^0(\Sigma)} \\ &\lesssim \|G^6\|_{H^0(\Omega)} + \|\omega_{c0}c + \omega_{b0}b\|_{H^0(\Sigma)} + \|G^7\|_{H^0(\Sigma)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \end{aligned} \quad (6.46)$$

By the Poincaré-Wirtinger inequality, this implies

$$\|\partial_t b\|_{H^1(\Omega)} \lesssim \|\nabla \partial_t b\|_{H^0(\Omega)} + \left| \int_{\Omega} \partial_t b \right| \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}} \sqrt{\mathcal{D}}. \quad (6.47)$$

Therefore, we continue using elliptic estimate to obtain

$$\begin{aligned}
 \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^3(\Omega)} &\lesssim \|\partial_t b\|_{H^1(\Omega)} + \|G^6\|_{H^1(\Omega)} + \|\omega_{c0}c + \omega_{b0}b + G^7\|_{H^{\frac{3}{2}}(\Sigma)} \\
 &\lesssim \|\partial_t b\|_{H^1(\Omega)} + \|G^6\|_{H^1(\Omega)} + \|c\|_{H^{\frac{3}{2}}(\Sigma)} + \|b\|_{H^2(\Omega)} + \|G^7\|_{H^{\frac{3}{2}}(\Sigma)} \\
 &\lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}}\sqrt{\mathcal{D}}.
 \end{aligned} \tag{6.48}$$

Naturally, we have

$$\|b\|_{H^3(\Omega)} \lesssim \left| \int_{\Omega} b \right| + \left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^3(\Omega)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}}\sqrt{\mathcal{D}}. \tag{6.49}$$

In summary, we have shown the desired estimate

$$\|\partial_t b\|_{H^1(\Omega)} + \|b\|_{H^3(\Omega)} \lesssim \sqrt{\bar{\mathcal{D}}} + \sqrt{\mathcal{E}}\sqrt{\mathcal{D}}. \tag{6.50}$$

□

## 7 Proof of Main Results

We now synthesize the previous results to prove the main a priori estimates.

**Theorem 7.1** *Suppose  $(u, p, \eta, c, b)$  are solution to (1.53) in  $t \in [0, T]$ . Then there exists a universal  $0 < \delta \ll 1$  such that for any*

$$\sup_{t \in [0, T]} \mathcal{E}(t) \leq \delta, \quad \int_0^T \mathcal{D}(t) dt < \infty, \tag{7.1}$$

we have

$$\sup_{t \in [0, T]} e^{\lambda t} \mathcal{E}(t) + \int_0^T \mathcal{D}(t) dt \lesssim \mathcal{E}(0), \tag{7.2}$$

for some universal  $\lambda > 0$ .

*Proof* Theorems 6.3 and 6.4 justify

$$\bar{\mathcal{E}} \lesssim \mathcal{E} \lesssim \bar{\mathcal{E}}, \quad \bar{\mathcal{D}} \lesssim \mathcal{D} \lesssim \bar{\mathcal{D}}. \tag{7.3}$$

Theorem 5.1 justifies

$$\frac{d}{dt}(\bar{\mathcal{E}} - \mathcal{F}) + \bar{\mathcal{D}} \lesssim 0. \tag{7.4}$$

Also, Lemmas 4.1 and 4.2 together with (7.3) show that

$$|\mathcal{F}| \leq \frac{\bar{\mathcal{E}}}{4} + C\mathcal{E}^{\frac{3}{2}} = \left( \frac{1}{4} + C\sqrt{\mathcal{E}} \right) \bar{\mathcal{E}}, \tag{7.5}$$

which implies that for  $\delta$  smaller than a universal constant

$$\frac{\bar{\mathcal{E}}}{2} \leq \bar{\mathcal{E}} - \mathcal{F} \leq \frac{3\bar{\mathcal{E}}}{2}. \quad (7.6)$$

Hence, integrating over  $[0, T]$  in (7.4), we have

$$\int_0^T \mathcal{D}(t) dt \lesssim (\bar{\mathcal{E}}(T) - \mathcal{F}(T)) + \int_0^T \mathcal{D}(t) dt \lesssim \bar{\mathcal{E}}(0) + \mathcal{F}(0), \quad (7.7)$$

which further implies

$$\int_0^T \mathcal{D}(t) dt \lesssim \mathcal{E}(0). \quad (7.8)$$

Also, it holds naturally that

$$\bar{\mathcal{E}} - \mathcal{F} \lesssim \bar{\mathcal{E}} \lesssim \mathcal{D}, \quad (7.9)$$

which means there exists  $\lambda > 0$  such that

$$\frac{d}{dt}(\bar{\mathcal{E}} - \mathcal{F}) + \lambda(\bar{\mathcal{E}} - \mathcal{F}) \leq 0. \quad (7.10)$$

Therefore, this implies

$$\bar{\mathcal{E}}(t) \lesssim \bar{\mathcal{E}}(0) - \mathcal{F}(t) \lesssim e^{-\lambda t} (\bar{\mathcal{E}}(0) - \mathcal{F}(0)) \lesssim e^{-\lambda t} \mathcal{E}(0), \quad (7.11)$$

which yields our desired decay estimate.  $\square$

## Appendix A: Surface Differential Operators

The proof of following lemmas can be found in [15, Appendix A].

**Lemma A.1** *We have the following identities:*

$$\nabla_\Gamma \cdot v = \nabla_* v_* = -H, \quad (A.1)$$

$$\partial_i \sqrt{1 + |\nabla_* \eta|^2} = -v_* \cdot \nabla_* \partial_i \eta, \quad (A.2)$$

$$\partial_t \sqrt{1 + |\nabla_* \eta|^2} = \nabla_* \left( \frac{\partial_t \eta \nabla_* \eta}{\sqrt{1 + |\nabla_* \eta|^2}} \right) - \partial_t \eta H, \quad (A.3)$$

$$\nabla_\Gamma f \cdot v = 0. \quad (A.4)$$

**Lemma A.2** *We have the following identities:*

$$\int_\Sigma \partial_{\Gamma,i} f g \sqrt{1 + |\nabla_* \eta|^2} = - \int_\Sigma (f \partial_{\Gamma,i} g + f g v_i H) \sqrt{1 + |\nabla_* \eta|^2}, \quad (A.5)$$

$$\int_\Sigma \nabla_\Gamma \cdot X \sqrt{1 + |\nabla_* \eta|^2} = - \int_\Sigma (X \cdot v) H \sqrt{1 + |\nabla_* \eta|^2}. \quad (A.6)$$

**Lemma A.3** Assume  $f \in C^2(\mathbb{R})$ . Suppose  $\tilde{c}$  and  $\eta$  satisfies

$$\begin{cases} \partial_t \tilde{c} + u \cdot \nabla_* \tilde{c} + \tilde{c} \nabla_\Gamma \cdot u = \gamma \Delta_\Gamma \tilde{c} - X, \\ \partial_t \eta - (u \cdot \nu) \sqrt{1 + |\nabla_* \eta|^2} = 0. \end{cases} \quad (\text{A.7})$$

Then we have

$$\begin{aligned} & \frac{d}{dt} \int_{\Sigma} f(\tilde{c}) \sqrt{1 + |\nabla_* \eta|^2} \\ &= \int_{\Sigma} ((f(\tilde{c}) - f'(\tilde{c})\tilde{c}) \nabla_\Gamma \cdot u - f''(\tilde{c}) |\nabla_\Gamma \tilde{c}|^2 - f'(\tilde{c}) X) \sqrt{1 + |\nabla_* \eta|^2}. \end{aligned} \quad (\text{A.8})$$

## Appendix B: Analytic Tools

We define the Poisson integral from  $\Sigma$  to  $\Sigma \times (-\infty, 0]$  as

$$\mathcal{P}[f](x) = \sum_{n \in (L_1^{-1}\mathbb{Z}) \times (L_2^{-1}\mathbb{Z})} e^{2\pi i n x_*} e^{2\pi |n| x_3} \hat{f}(n), \quad (\text{B.1})$$

where

$$\hat{f}(n) = \int_{\Sigma} f(x_*) \frac{e^{2\pi i n x_*}}{L_1 L_2} dx_* . \quad (\text{B.2})$$

$\mathcal{P}$  is a bounded linear operator from  $H^s(\Sigma)$  to  $H^{s+\frac{1}{2}}(\Sigma \times (-\infty, 0])$  for  $s > 0$ .

In what follows we recall that  $\Omega = \Sigma \times (-L_3, 0)$ . The proof of following lemmas can be found in [15, Appendix B].

**Lemma B.1** We have for  $q \geq 1$ ,

$$\|\nabla^q \mathcal{P}[f]\|_{H^0(\Omega)} \lesssim \|f\|_{\dot{H}^{q-\frac{1}{2}}(\Sigma)}. \quad (\text{B.3})$$

**Lemma B.2** We have for  $q \geq 1$  and  $s > 1$ ,

$$\|\nabla^q \mathcal{P}[f]\|_{L^\infty(\Omega)} \lesssim \|f\|_{\dot{H}^{q+s}(\Sigma)}. \quad (\text{B.4})$$

The same estimate holds for  $q = 0$  if  $\int_{\Sigma} f = 0$ .

**Lemma B.3** Let  $U$  denote either  $\Sigma$  or  $\Omega$ .

(1) Let  $0 \leq r \leq s_1 \leq s_2$  be such that  $s_1 > n/2$ . Let  $f \in H^{s_1}(U)$ ,  $g \in H^{s_2}(U)$ . Then  $fg \in H^r(U)$  and

$$\|fg\|_{H^r(U)} \lesssim \|f\|_{H^{s_1}(U)} \|g\|_{H^{s_2}(U)}. \quad (\text{B.5})$$

(2) Let  $0 \leq r \leq s_1 \leq s_2$  be such that  $s_2 > r + n/2$ . Let  $f \in H^{s_1}(U)$ ,  $g \in H^{s_2}(U)$ . Then  $fg \in H^r(U)$  and

$$\|fg\|_{H^r(U)} \lesssim \|f\|_{H^{s_1}(U)} \|g\|_{H^{s_2}(U)}. \quad (\text{B.6})$$

(3) Let  $0 \leq r \leq s_1 \leq s_2$  be such that  $s_2 > r + n/2$ . Let  $f \in H^{-s}(U)$ ,  $g \in H^{s_2}(U)$ . Then  $fg \in H^{-s_1}(U)$  and

$$\|fg\|_{H^{-s_1}(U)} \lesssim \|f\|_{H^{-r}(U)} \|g\|_{H^{s_2}(U)}. \quad (\text{B.7})$$

In the comparison estimate, we need the following classical elliptic estimate.

**Lemma B.4** Assume  $\phi \in H^{r-2}(\Omega)$  and  $\psi \in H^{r-3/2}(\Sigma)$  for  $r \geq 1$ . Then solutions to the elliptic problem

$$\begin{cases} -\Delta b = \phi & \text{in } \Omega, \\ \partial_3 b = \psi & \text{on } \Sigma, \\ \partial_3 b = 0 & \text{on } \Sigma_b, \end{cases} \quad (\text{B.8})$$

satisfy the estimate

$$\left\| b - \frac{1}{|\Omega|_3} \int_{\Omega} b \right\|_{H^r(\Omega)} \lesssim \|\phi\|_{H^{r-2}(\Omega)} + \|\psi\|_{H^{r-3/2}(\Sigma)}. \quad (\text{B.9})$$

## References

1. Abels, H., Garcke, H., Lam, K., Weber, J.: Two-phase flow with surfactants: diffuse interface models and their analysis. Preprint (2016). [arXiv:1610.08221](https://arxiv.org/abs/1610.08221)
2. Barrett, J., Garcke, H., Nürnberg, R.: Stable finite element approximations of two-phase flow with soluble surfactant. *J. Comput. Phys.* **297**, 530–564 (2015)
3. Beale, J.: Large-time regularity of viscous surface waves. *Arch. Ration. Mech. Anal.* **84**(4), 307–352 (1983/84)
4. Beppler, B.K., Varanasi, K.S., Garoff, S., Eymenenko, G., Woods, K.N.: Influence of fluid flow on the deposition of soluble surfactants through receding contact lines of volatile solvents. *Langmuir* **24**, 6705–6711 (2008)
5. Bothe, D., Prüss, J.: Stability of equilibria for two-phase flows with soluble surfactant. *Q. J. Mech. Appl. Math.* **63**(2), 177–199 (2010)
6. Bothe, D., Prüss, J., Simonett, G.: Well-posedness of a two-phase flow with soluble surfactant. In: Non-linear Elliptic and Parabolic Problems. *Progr. Nonlinear Differential Equations Appl.*, vol. 64, pp. 37–61. Birkhäuser, Basel (2005)
7. Diamant, H., Andelman, D.: Kinetics of surfactant absorption at fluid-fluid interfaces. *J. Phys. Chem.* **100**, 13732–13742 (1996)
8. Diamant, H., Ariel, G., Andelman, D.: Kinetics of surfactant adsorption: the free energy approach. *Colloids Surf. A, Physicochem. Eng. Asp.* **183**, 259–276 (2001)
9. Edwards, D., Brenner, H., Wasan, D.: *Interfacial Transport Processes and Rheology*. Butterworth-Heinemann, Boston (1991)
10. Garcke, H., Wieland, S.: Surfactant spreading on thin viscous films: nonnegative solutions of a coupled degenerate system. *SIAM J. Math. Anal.* **37**(6), 2025–2048 (2006)
11. Garcke, H., Lam, K., Stinner, B.: Diffuse interface modelling of soluble surfactants in two-phase flow. *Commun. Math. Sci.* **12**(8), 1475–1522 (2014)
12. Guo, Y., Tice, I.: Local well-posedness of the viscous surface wave problem without surface tension. *Anal. PDE* **6**(2), 287–369 (2013)
13. Guo, Y., Tice, I.: Almost exponential decay of periodic viscous surface waves without surface tension. *Arch. Ration. Mech. Anal.* **207**(2), 459–531 (2013)
14. Hills, B.: An alternative view of the role(s) of surfactant and the alveolar model. *J. Appl. Physiol.* **87**(5), 1567–1583 (1999)
15. Kim, C., Tice, I.: Dynamics and stability of surfactant-driven surface wave. *SIAM J. Math. Anal.* **49**(2), 1295–1332 (2017)
16. Kwan, Y.-Y., Park, J., Shen, J.: A mathematical and numerical study of incompressible flows with a surfactant monolayer. *Discrete Contin. Dyn. Syst.* **28**(1), 181–197 (2010)

17. Levich, V.: *Physiochemical Hydrodynamics*. Prentice Hall International, Englewood Cliffs (1962)
18. Marcinkowski, A.L., Garoff, S., Tilton, R., Pilewski, J.M., Corcoran, T.E.: Post-deposition dispersion of aerosol medications using surfactant carriers. *J. Aerosol Med. Pulm. Drug Deliv.* **21**, 361–369 (2008)
19. Myers, D.: *Surfactants Science and Technology*. Wiley, Hoboken (2006)
20. Rosen, M.: *Surfactants and Interfacial Phenomena*. Wiley, Hoboken (2004)
21. Sarpkaya, T.: Vorticity, free surface, and surfactants. *Annu. Rev. Fluid Mech.* **28**, 83–128 (1996), Annual Reviews, Palo Alto, CA
22. Venables, J.: *Introduction to Surface and Thin Film Processes*. Cambridge University Press, London (2000)
23. Wang, Y., Tice, I., Kim, C.: The viscous surface-internal wave problem: global well-posedness and decay. *Arch. Ration. Mech. Anal.* **212**(1), 1–92 (2014)
24. Wu, L.: Well-posedness and decay of the viscous surface wave. *SIAM J. Math. Anal.* **46**(3), 2084–2135 (2014)
25. Xu, J.-J., Li, Z., Lowengrub, J., Zhao, H.: A level-set method for interfacial flows with surfactant. *J. Comput. Phys.* **212**(2), 590–616 (2006)