

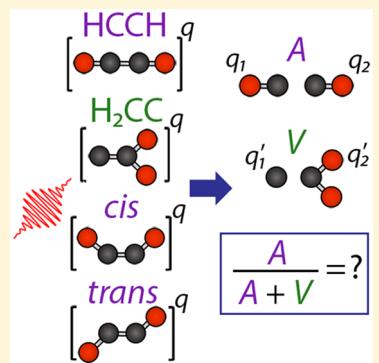
Dependence on the Initial Configuration of Strong Field-Driven Isomerization of C_2H_2 Cations and Anions

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ABSTRACT: We have investigated the femtosecond laser-induced fragmentation of $C_2H_2^q$ ion beam targets in various initial configurations, including acetylene (linear HCCH), vinylidene (H_2CC), and *cis/trans*. The initial configuration is shown to have a tremendous impact on the branching ratio of acetylene-like ($CH^{q_1} + CH^{q_2}$) and vinylidene-like ($C^{q_1} + CH_2^{q_2}$) dissociation of a specific $C_2H_2^q$ molecular ion. In particular, whereas $C_2H_2^+$ generated from C_2H_2 , a linear HCCH target, exhibits comparable levels of acetylene-like and vinylidene-like fragmentation, vinylidene or *cis/trans* configuration ion beams preferably undergo vinylidene-like fragmentation, with an acetylene branching ratio ranging from 13.9% to zero.



Photochemistry studies often involve probing molecular dynamics starting from stable configurations. The initial configuration, however, can have a great influence on the dynamics. Examining various geometries beyond those commonly studied can therefore offer a different perspective. One well-studied topic that may benefit from such an alternative approach is isomerization.

Widespread in nature^{1–4} and key to many industrial applications,^{5–7} isomerization reactions have been extensively studied in chemistry and physics, both experimentally and theoretically.^{8,9} The C_2H_2 molecule in particular has proven to be a useful testing ground for examining isomerization. Conversion between its acetylene (HCCH) and vinylidene (H_2CC) isomers, an example of 1,2 hydrogen migration,¹⁰ is one of the simplest bond rearrangement isomerization reactions.

A wide array of experimental techniques have been applied to examine isomerization of C_2H_2 molecules. Photodetachment spectroscopy measurements on $C_2H_2^-$ ions have collectively provided a fuller picture of the energetics and $H_2CC \rightarrow HCCH$ isomerization dynamics of the resulting transient neutral vinylidene.^{11–16} With an alternative approach, stimulated emission spectroscopy (SEP)^{17–19} and dispersed fluorescence (DF)^{17,19–21} measurements have examined the reverse reaction, preparing acetylene (HCCH) in highly excited vibrational states close to the isomerization threshold and looking for signatures of $HCCH \rightarrow H_2CC$ transitions.

In addition to being heavily studied in traditional linear spectroscopy,^{13–19,21} C_2H_2 isomerization has also garnered significant interest in the photofragmentation community. Here, a far less “gentle” approach is taken, breaking the molecule and extracting information from the measured

fragments. The signatures of isomerization are thus fundamentally different from those in the aforementioned types of spectroscopy. For example, common signatures in fragmentation studies are measurement of $C^{q_1} + CH_2^{q_2}$ final products for acetylene \rightarrow vinylidene isomerization and $CH^{q_1} + CH^{q_2}$ for vinylidene \rightarrow acetylene isomerization.

These photofragmentation studies have involved an assortment of excitation sources. Extreme ultraviolet (XUV) and X-ray sources have been fruitfully employed to study C_2H_2 isomerization through dissociative ionization.^{22–34} Following early suggestions and demonstrations of the potential utility of intense, ultrafast laser pulses in probing C_2H_2 isomerization dynamics, particularly the relevant time scale(s), a plethora of studies have used these sources.^{35–47}

Despite the abundance of studies of C_2H_2 isomerization, it is nevertheless an intriguing and important topic to explore further, as much remains unknown. For instance, in addition to the ongoing debate about topics like the lifetime of the neutral vinylidene,^{12,13,48–51} open questions about the ions of C_2H_2 are even greater in number.⁵²

In this Letter, we utilize our unique capability to prepare $C_2H_2^q$ ion beams in a variety of initial configurations to investigate the role these configurations play in ultrafast laser-induced isomerization and dissociation. To do so, we examine target configurations ranging from vinylidene to acetylene, including *cis/trans* configurations. These would not be easily accessible as neutral gas phase targets.

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In our experiment, C_2H_2^q ions are produced via fast electron impact in an ion source. Specifically, an electron cyclotron resonance (ECR) ion source is used to create the C_2H_2^+ beams, while the C_2H_2^- beam is generated in a duoplasmatron source. The ions are accelerated to an energy of ~ 8 keV, and the ion beam of interest is selected with a magnet according to its momentum-to-charge ratio. At the laser interaction region, the ion beam has a cross section of ~ 0.9 mm \times ~ 0.9 mm. The typical current varies widely depending on the specific ion beam over the range of 25 pA to 1 nA.

The Ti:sapphire laser used in this study⁵³ produces 790 nm central wavelength, 2 mJ, linearly polarized pulses with a duration of ~ 23 fs fwhm (full width at half-maximum in intensity) at a rate of 10 kHz. The pulse duration is measured with second-harmonic generation frequency-resolved optical gating (SHG FROG).⁵⁴ The laser beam is focused onto the ion beam target by an $f = 203$ mm, 90° off-axis parabolic mirror. Typical intensities utilized in our experiment fall in the range of 10^{13} – 10^{15} W/cm². The peak intensity is determined from the measured power, pulse duration, and focal profile image.⁵⁵ The intensity is controlled by translating the position of the laser focus relative to the ion beam center as well as by attenuating the laser power.^{55,56}

Our setup is schematically illustrated in Figure 1. Signals from a microchannel plate delay-line detector provide arrival

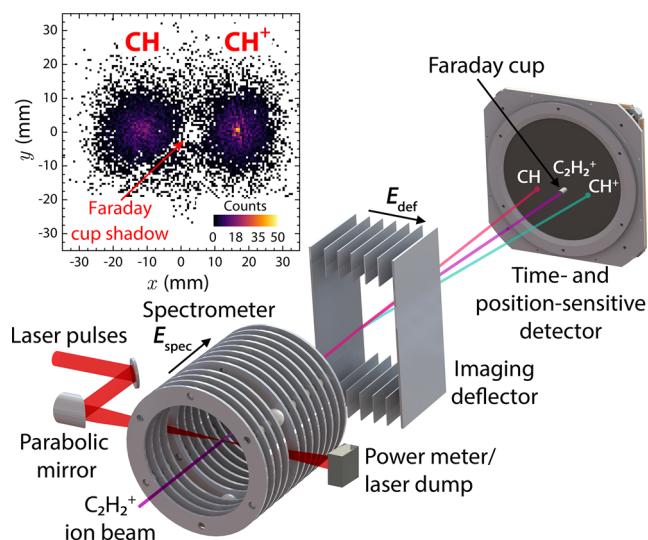


Figure 1. Schematic of our coincidence three-dimensional momentum imaging setup showing the $\text{C}_2\text{H}_2^+ \rightarrow \text{CH}^+ + \text{CH}$ breakup channel. Intense, femtosecond laser pulses interact with the ion beam inside a spectrometer that provides a longitudinal field, E_{spec} , to separate fragments in time. An imaging deflector supplies a transverse field, E_{def} , which separates the fragments in position. The inset shows a sample position spectrum.

time and position information and hence enable calculation of the velocity for each fragment on an event-by-event basis. Each fragmentation channel appears as a diagonal stripe in the coincidence time-of-flight (CTOF) spectrum due to momentum conservation, like those shown in Figure 2a. The typical event rate depends strongly upon the particular ion beam, with event rates for our measurements being at most on the order of 0.1 Hz (for C_2H_2^+ from C_2H_2) to 0.001 Hz (for C_2H_2^+ from $\text{C}_2\text{H}_2\text{F}_2$). Further details regarding our experimental technique can be found in our previous publications.^{57,58}

Momentum conservation allows the discernment of true events from background signals and random coincidences. Random coincidences originate from the interaction of a single laser pulse with multiple molecules, whereby a fragment from one molecule is accidentally paired with another from a different molecule.^{57,59–62} To eliminate these events, we pair fragments from different laser shots to generate a spectrum mimicking the random coincidences, scale this generated spectrum to a purely random feature like those marked in Figure 2a, and then subtract it from the measured data.

Recall that in fragmentation studies, a commonly used signature of molecules being in the acetylene configuration is measurement of the $\text{CH}^{q_1} + \text{CH}^{q_2}$ channel, which we denote as A. Similarly, a signature of molecules in the vinylidene configuration is measurement of the $\text{C}^{q_1} + \text{CH}_2^{q_2}$ channel, which we denote as V. We adopt this notation from this point forward.

For the case of C_2H_2^+ ions being produced from C_2H_2 , these ions predominantly exist in the same configuration as the parent C_2H_2 , HCCH. This assertion is based on electron impact ionization studies.^{63–66} The results for this target are similar to those of the aforementioned photofragmentation studies starting from a neutral C_2H_2 target.^{22–47} Like those studies, we observe both acetylene-like and vinylidene-like two-body breakup channels, visible in the CTOF spectrum in Figure 2a.

Given our ability to measure neutral fragments directly, we focus on the two-body dissociation channels, $\text{CH}^+ + \text{CH}$ (A), $\text{C}^+ + \text{CH}_2$ (V_1), and $\text{C} + \text{CH}_2^+$ (V_2), i.e., those with no ionization of the target, and compute their branching ratios after imposing momentum conservation. Note that the presence of two V channels is merely a consequence of the fact that either the C or the CH_2 fragment is positively charged. The CTOF spectrum for these channels after analysis is shown in Figure 2c. The branching ratios are summarized in Table 1, where the sum of the vinylidene-like channel yields, $V_1 + V_2$, is denoted as V. These results suggest a substantial amount of isomerization of the initially acetylene (linear HCCH configuration) target, as the $V/(A + V)$ branching ratio is $\sim 60\%$.

While the C_2H_2^+ target produced from C_2H_2 likely has a linear HCCH geometry, the C_2H_2^- target possesses the other isomeric configuration of C_2H_2 , vinylidene. This beam is produced by electron attachment in a duoplasmatron ion source loaded with a mixture of 10% ethylene and 90% argon. C_2H_2^- is known to be long-lived in its electronic ground state, $\tilde{X}^2\text{B}_2$. It remains in the vinylidene configuration for >100 s,^{12,67,68} decaying by adiabatic electron loss and subsequent isomerization of the resulting neutral vinylidene to acetylene.⁶⁹ The linear HCCH[–] isomer, on the other hand, has a large negative electron affinity (EA) of -2.6 eV, and it autodetaches within a few femtoseconds.^{69–71} Because the ion beam travel time from the ion source to the laser interaction region in our setup is approximately $20\ \mu\text{s}$, the C_2H_2^- target survives purely in the vinylidene configuration.

In stark contrast to the C_2H_2^+ results, we find that the laser-induced fragmentation of C_2H_2^- is conspicuously reminiscent of the initial target configuration. Predominantly vinylidene-like breakup is observed for all $\text{C}_2\text{H}_2^{q+}$ daughter ions, as evident from the CTOF spectrum in Figure 2b. Figure 2d contains the CTOF spectrum of events fulfilling momentum conservation for two-body dissociation of the C_2H_2^+ daughter

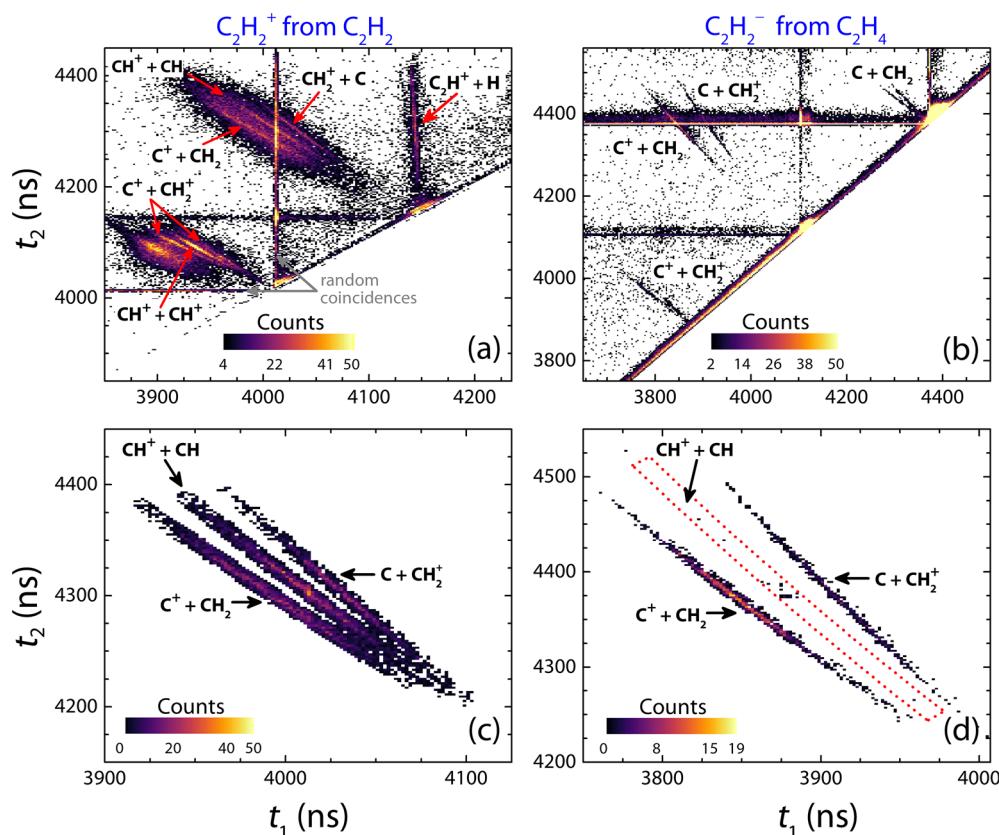


Figure 2. Coincidence time-of-flight (CTOF) spectra for fragmentation of C_2H_2^+ generated from C_2H_2 (left) and C_2H_2^- (right). (a and b) Raw data. (c and d) Fragmentation channels of C_2H_2^+ after momentum conservation. The dashed red box in panel d indicates the expected position of the $\text{CH}^+ + \text{CH}$ channel, i.e., the channel associated with the HCCH final configuration. Gray arrows in panel a mark examples of random coincidence features (see the text). Note that there are other channels not discussed here that will be covered in a future publication(s). For instance, the rate of $\text{CH}^+ + \text{CH}^+$ in panel a is ~ 0.2 Hz.

Table 1. Branching Ratios for Different Initial C_2H_2^q Configurations ($q = -1, 0$, or 1) Exposed to Laser Pulses with an Intensity of $\sim 3 \times 10^{15} \text{ W/cm}^2$

target	source gas	current (nA)	initial configuration	A	$A/(A + V)$ (%)
C_2H_2^+	C_2H_2	1.0	HCCH	$\text{CH}^+ + \text{CH}$	37.6 ± 3.8
C_2H_2^+	$\text{C}_2\text{H}_2\text{F}_2$	0.025	H_2CC	$\text{CH}^+ + \text{CH}$	13.9 ± 2.4
C_2H_2^+	$\text{C}_2\text{H}_2\text{Br}_2$	0.1	<i>cis/trans</i>	$\text{CH}^+ + \text{CH}$	1.7 ± 0.3
C_2H_2^-	C_2H_4	0.2	H_2CC	$\text{CH}^+ + \text{CH}$	0.1 ± 0.5
C_2H_2^-	C_2H_4	0.2	H_2CC	$\text{CH} + \text{CH}$	3.9 ± 2.5
C_2H_2^b	C_2H_4	0.2	H_2CC	$\text{CH}^+ + \text{CH}$	0.1 ± 0.1
C_2H_2^b	C_2H_4	0.2	H_2CC	$\text{CH} + \text{CH}$	3.6 ± 1.2

^aThe labels A and V refer to the $\text{CH}^{q_1} + \text{CH}^{q_2}$ and $\text{C}^{q_1} + \text{CH}_2^{q_2}$ channels, respectively (note $q_1 + q_2 = q_1' + q_2'$ for each case). Also, in some cases, V includes two channels, as either the C or the CH_2 fragment can be positively charged. ^bThe neutral target is produced from the C_2H_2^- beam by a pump pulse 300 fs before interaction with the probe pulse (see the text).

of C_2H_2^- , specifically the $\text{CH}^+ + \text{CH}$ (A), $\text{C}^+ + \text{CH}_2$ (V_1), and $\text{C} + \text{CH}_2^+$ (V_2) channels. The acetylene branching ratio for this process, displayed in Table 1, is consistent with zero. Moreover, this trend of negligible acetylene-like breakup persists for the range of peak laser intensities studied (5×10^{14} to $3 \times 10^{15} \text{ W/cm}^2$).

The contrast between the C_2H_2^+ and C_2H_2^- outcomes discussed so far is surprising. These observations suggest that while C_2H_2^+ generated from C_2H_2 , an HCCH target, isomerizes efficiently, the reverse process does not occur for C_2H_2^- , a target of the H_2CC initial configuration. More precisely, isomerization of the vinylidene target leading to dissociation is not observed. One should note that this does not preclude the possibility of isomerization of molecules that

remain bound. Also, while it is likely that the contrasting C_2H_2^+ and C_2H_2^- results are due to very different excitation pathways, it is not easy to resolve the participating electronic states in the present strong field study, as many electronic states could be populated. For isolating dissociation pathways, further work employing single-photon excitation of these systems^{32,72,73} could be informative.

Some theoretical studies have suggested that neutral vinylidene that results from $\text{C}_2\text{H}_2^- \rightarrow \text{C}_2\text{H}_2 + \text{e}^-$ photo-detachment not only isomerizes to acetylene but then undergoes multiple acetylene \leftrightarrow vinylidene recrossings.^{50,74} This process happens due to the fact that acetylene is highly excited after crossing the vinylidene \rightarrow acetylene barrier. The period of these recrossings is calculated to be on the order of

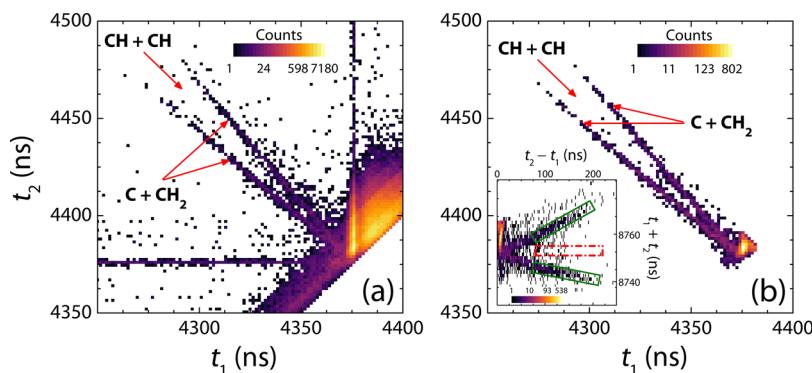


Figure 3. CTOF spectra for fragmentation of C_2H_2^- by a two-pulse scheme zoomed in on the neutral–neutral fragmentation channels. (a) Raw data. (b) After imposing momentum conservation. The inset shows a rotated spectrum. The dashed red and solid green boxes indicate where yields of $\text{CH} + \text{CH}$ and $\text{C} + \text{CH}_2$ were evaluated, respectively.

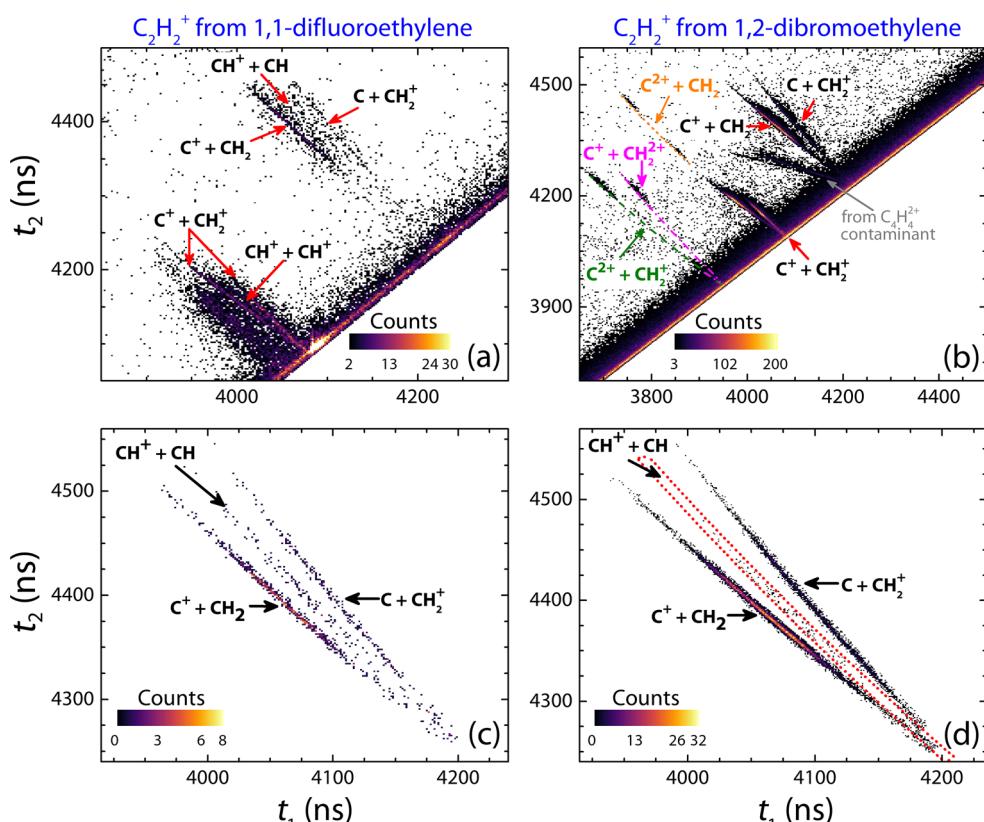


Figure 4. CTOF spectra for laser-induced fragmentation of C_2H_2^+ generated from 1,1-difluoroethylene (left) and 1,2-dibromoethylene (right). (a and b) Raw data. (c and d) The three fragmentation channels of C_2H_2^+ after momentum conservation. The dashed red box in panel d highlights the expected position of the $\text{CH}^+ + \text{CH}$ channel.

100 fs, but the exact period depends heavily upon the initial conditions.⁵⁰

In light of this idea, we performed a two-pulse measurement, where the first pulse causes photodetachment and possibly initiates isomerization and the second pulse probes any ensuing recrossing dynamics. The pump pulse that removes the weakly bound electron of C_2H_2^- (EA = 0.49 eV¹²) has a peak intensity of $2 \times 10^{14} \text{ W/cm}^2$ and efficiently neutralizes the C_2H_2^- molecules within the laser focus, according to our estimates. To arrive at this conclusion, we use our estimated ion beam density, $\sim 6 \times 10^3$ molecules/ cm^3 (i.e., ~ 0.03 molecule/pulse), and the measured rate of neutral C_2H_2 molecules (~ 285 Hz) for a single $5 \times 10^{14} \text{ W/cm}^2$ pulse,

which results in negligible fragmentation. The probe pulse is delayed by 300 fs with respect to the pump pulse, close to the typical recrossing period, and has a much higher peak intensity of $2 \times 10^{15} \text{ W/cm}^2$.

Here too, the vinylidene-like breakup channel, $\text{C} + \text{CH}_2$, heavily dominates, as clearly shown in Figure 3. In contrast, the acetylene-like fragmentation channel, $\text{CH} + \text{CH}$, is hardly visible. Careful analysis including momentum conservation and random coincidence subtraction sets an upper bound for the acetylene branching ratio at $3.6 \pm 1.2\%$. This result is consistent with that of a single pulse similar to the probe pulse, for which the acetylene branching ratio is $3.9 \pm 2.5\%$. Moreover, in the pump–probe measurement, the branching

ratio of the acetylene-like fragmentation channel for which two electrons are removed, $\text{CH}^+ + \text{CH}$, is $0.1 \pm 0.1\%$ (see Table 1), also in accord with the single-pulse results. Theory more closely replicating the initial conditions would aid in guiding future two-pulse experiments sampling recrossing dynamics.

To further explore this strong dependence of isomerization on the initial configuration, we look at another vinylidene-like target, C_2H_2^+ generated from 1,1-difluoroethylene. Charge inversion mass spectrometry measurements indicate that a beam of C_2H_2^+ produced from 1,1-dichloroethylene is long-lived in the vinylidene configuration, with a lifetime of $>8.5 \mu\text{s}$.^{75,76} Given the chemical similarity between 1,1-dichloroethylene and 1,1-difluoroethylene, it is most likely that a C_2H_2^+ beam generated from the latter is also of the vinylidene configuration. Panels a and c of Figure 4 contain the measured CTOF spectra for this target. One can see that both vinylidene-like and acetylene-like breakup occurs, with a vast majority of vinylidene-like breakup.

While a beam of C_2H_2^+ generated from 1,1-difluoroethylene has the vinylidene configuration like C_2H_2^- , there also are compelling reasons to examine the former target. Here, we observe dissociation and thus can avoid any complication ionization might add. In contrast, in the C_2H_2^- case, it is possible that fragmentation of C_2H_2^+ daughter ions involves sequential ionization. For example, the weakly bound electron could be removed early in the laser pulse, followed by propagation on the C_2H_2 surface(s) and eventual removal of the second electron. This process is clearly different from dissociation of a C_2H_2^+ target initially in the vinylidene configuration, such as that produced from 1,1-difluoroethylene.

As a final example, we generate a *cis/trans* C_2H_2^+ beam from 1,2-dibromoethylene. The dibromoethylene sample is composed of a mixture of the *cis* and *trans* isomers with a *cis:trans* ratio of approximately 1:3. This sample and information regarding its *trans:cis* isomer ratio were obtained from Sigma-Aldrich. Intuitively, removal of the bromine atoms will result in C_2H_2^+ far from equilibrium, with excitation likely in the symmetric (*trans*) and antisymmetric (*cis*) C–C–H bending vibrational modes.⁷⁷ What impact will this excitation have on the branching ratio between acetylene-like and vinylidene-like channels?

The CTOF spectra in panels b and d of Figure 4 and the branching ratio in Table 1 show that, similar to the C_2H_2^- results, the predominant breakup is vinylidene-like. This behavior continues over the range of intensities explored (5×10^{13} to $3 \times 10^{15} \text{ W/cm}^2$). For this *cis/trans* target, it should be noted that it is unclear if the initial bending excitation of the ion beam facilitates unimolecular $\text{HCCH} \rightarrow \text{H}_2\text{CC}$ isomerization in flight to the laser interaction or if the isomerization is laser-induced. While in assessing the geometry of this target we have assumed a logic similar to that applied for the H_2CC^+ beam produced from 1,1-difluoroethylene,^{75,76} a future study of the lifetimes of *cis/trans* beams would shed further light on this issue. The key point, however, is that isomerization to vinylidene is much more likely in the *cis/trans* case than in the case of C_2H_2^+ produced from C_2H_2 . Furthermore, the branching ratios measured with the *cis/trans* target are significantly different from those measured with the other targets.

The *cis/trans* C_2H_2^+ measurement suggests that the initial bending excitation of the ions enhances isomerization toward the vinylidene configuration. This observation is in agreement with previous work highlighting the influence of bending

motion on the isomerization process, both for acetylene becoming vinylidene and vice versa. The previously mentioned DF and SEP spectroscopy measurements^{17–21} exploit a high degree of vibrational excitation to examine the coupling of acetylene and vinylidene states resulting in isomerization of HCCH . In these spectroscopic measurements, the minimum energy path for hydrogen migration falls along the C–C–H bending coordinate. Several theoretical studies have also emphasized that bending excitation in acetylene molecules has a decisive impact on the isomerization dynamics.^{50,78–81} Moreover, photoelectron spectroscopy experiments on the vinylidene anion C_2H_2^- by the Neumark group and collaborators have shown that excitation of the in-plane rocking mode of vinylidene readily leads to isomerization, as this state connects to the C–C–H bending vibrational states of acetylene.¹⁶

In summary, our measurements of the laser-induced fragmentation of C_2H_2^q ion beams with assorted unique configurations demonstrate the strong effect of the initial configuration on the isomerization dynamics. While C_2H_2^+ generated from C_2H_2 , a linear HCCH molecule, exhibits acetylene-like and vinylidene-like fragmentation at similar levels, the vinylidene and *cis/trans* configuration C_2H_2^+ molecules that we examined mainly fragment in a vinylidene-like manner, a markedly different and surprising outcome. These intriguing results, which raise the question of the reversibility of the isomerization process, call for future theoretical endeavors to explain them. Measurements of *cis/trans* targets appear to be congruent with experimental and theoretical work pinpointing the pivotal influence of bending excitation in isomerization, but similarly, further work is needed for a detailed understanding and to disentangle the possible unimolecular and laser-driven dynamics.

Finally, we anticipate that our approach for studying C_2H_2^q ($q = -1, 0, \text{ or } 1$) isomerization is not limited to this particular problem. Rather, ion beam targets offer an avenue for examining photoinduced dynamics in a variety of chemical systems in initial configurations that are not easily accessible, a unique vantage point that is complementary to mainstream approaches for studying molecular dynamics.

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Notes

The authors declare no competing financial interest.

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