

Minimal Mass Blowup Solutions for the Patlak-Keller-Segel Equation

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Abstract

We consider the parabolic-elliptic Patlak-Keller-Segel (PKS) model of chemotactic aggregation in two space dimensions which describes the aggregation of bacteria under chemotaxis. When the mass is equal to 8π and the second moment is finite (the doubly critical case), we give a precise description of the dynamic as time goes to infinity and extract the limiting profile and speed. The proof shows that this dynamic is stable under perturbations. © 2018 Wiley Periodicals, Inc.

1 Introduction

We consider the two-dimensional Patlak-Keller-Segel system,

$$(1.1) \quad \begin{cases} \partial_t u(x, t) - \Delta u(x, t) = -\nabla \cdot (u(x, t) \nabla c(x, t)), \\ -\Delta c(x, t) = u(x, t), \\ u(x, t = 0) = u_0 \geq 0, \end{cases}$$

where $x \in \mathbb{R}^2$ and $t > 0$. This system is generally considered the fundamental mathematical model for the study of aggregation by chemotaxis of certain microorganisms [14, 15, 17, 29]. From now on we will refer to (1.1) as Patlak-Keller-Segel (PKS). The first equation describes the motion of the microorganism (u represents the density of cells) as a random walk with drift up the gradient of the *chemoattractant* c . The second equation describes the production and (instantaneous) diffusion of the chemoattractant. PKS and related variants have received considerable mathematical attention over the years, for example, see the review [15] or some of the following representative works [1–3, 5, 7, 8, 11, 13, 16, 27, 33, 35]. The equation is $L^{d/2}$ critical in dimension d ; hence it is mass or L^1 critical in dimension 2, namely: if $u(t, x)$ is a solution to (1.1), then for all $a \in (0, \infty)$, so is

$$u_a(t, x) = \frac{1}{a^2} u\left(\frac{t}{a^2}, \frac{x}{a}\right).$$

Moreover, if $M = 8\pi$, then the second moment is also conserved. This justifies the terminology doubly critical if $M = 8\pi$ and yields very interesting dynamical properties. It has been known for some time that (1.1) possesses a critical mass: if $\|u_0\|_1 \leq 8\pi$, then classical solutions exist for all time (see, e.g., [1–3, 5, 35]), and if $\|u_0\|_1 > 8\pi$, then all classical solutions with finite second moment blow up in finite time [5, 16, 27] and are known to concentrate at least 8π mass into a single point at blowup [35] (see also [13, 33]). Another important property of (1.1) that plays a decisive role in our work is the existence (and uniqueness) of self-similar spreading solutions for all mass $M \in (0, 8\pi)$. These are known to be global attractors for the dynamics if the total mass is less than 8π [5] and for the purposes of our analysis. The mass is conserved in (1.1):

$$(1.2) \quad \frac{d}{dt} \int_{\mathbb{R}^2} u(x, t) dx = 0.$$

Hence, we set

$$M_0 = \int_{\mathbb{R}^2} u_0(x) dx.$$

Notice that the center of mass is also conserved

$$\frac{d}{dt} \int_{\mathbb{R}^2} x_i u(x, t) dx = 0 \quad \text{for } i \in \{1, 2\}.$$

Since the solution of the Poisson equation $-\Delta c = u$ is given by

$$c = \phi_u = -\frac{1}{2\pi} \log |\cdot| \star u,$$

the system becomes

$$(1.3) \quad \begin{cases} u_t = \nabla \cdot (\nabla u - u \nabla \phi_u), \\ \phi_u = -\frac{1}{2\pi} \log |\cdot| \star u & \text{in } [0, T] \times \mathbb{R}^2, \\ u(0) = u_0 \geq 0. \end{cases}$$

In addition, we have that the flow dissipates the free energy

$$(1.4) \quad \frac{d}{dt} \mathcal{F}(u) = \frac{d}{dt} \left(\int_{\mathbb{R}^2} u \log u \, dx - \frac{1}{2} \int_{\mathbb{R}^2} u \phi_u \, dx \right) \leq 0,$$

where \mathcal{F} is the sum of the entropy and the potential energy. The problem is locally well-posed in the finite Borel measures space [6], but the question of global existence requires more conditions. Indeed, we can formally compute a virial identity for a solution of (1.3):

$$\begin{aligned} & \frac{d}{dt} \int_{\mathbb{R}^2} u(x, t) |x|^2 \, dx \\ &= - \int_{\mathbb{R}^2} 2x \cdot \nabla u \, dx + \frac{1}{2\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} \frac{2x \cdot (y - x)}{|x - y|^2} u(x, t) u(y, t) \, dy \, dx = \end{aligned}$$

$$\begin{aligned}
&= 4M_0 + \frac{1}{2\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} \frac{(x-y) \cdot (y-x)}{|x-y|^2} u(x,t)u(y,t) dy dx \\
&= 4M_0 - \frac{M_0^2}{2\pi}.
\end{aligned}$$

So that all solutions with finite second moment and mass bigger than 8π cannot be global in time, we also have an estimation on the maximal existence time T_{\max} if the second moment of the initial data u_0 is finite,

$$T_{\max} \leq \frac{2\pi}{M_0(M_0 - 8\pi)} \int_{\mathbb{R}^2} |x|^2 u_0(x) dx.$$

Thus, this virial computation is a heuristic explanation of the following trichotomy:

- If $\int_{\mathbb{R}^2} u_0(x) dx < 8\pi$, then the solutions are global in time. In addition, it has been proved that the solutions are spreading [5], and the density converges to a self-similar profile in rescaled variables.
- If $\int_{\mathbb{R}^2} u_0(x) dx > 8\pi$, then solutions blow up in finite time [5, 16].
- If $\int_{\mathbb{R}^2} u_0(x) dx = 8\pi$ and $\int_{\mathbb{R}^2} |x|^2 u_0(x) dx < \infty$, solutions are global in time but concentrate at the origin in infinite time [3]. In contrast, if $\int_{\mathbb{R}^2} u_0(x) dx = 8\pi$ and $\int_{\mathbb{R}^2} |x|^2 u_0(x) dx = \infty$ what could happen is not clear except if the solution initially is sufficiently close to a rescaling Q_a of the stationary solution Q [2].

Indeed, the stationary solution of (1.3) is

$$Q(x) = \frac{8}{(1 + |x|^2)^2},$$

its mass is 8π , its second moment is infinite, and Q is the unique minimizer up to symmetries of the free energy $\mathcal{F}(u)$ with $\int_{\mathbb{R}^2} u(x, t) dx = 8\pi$.

When $M_0 < 8\pi$, the solutions are spreading and they converge to a self-similar profile in rescaled variables with a rate $\frac{1}{t}$ in [4, 8, 9], whereas if $M_0 > 8\pi$ the solutions blow up in finite time. A result in [13] describes the dynamic of the blowup by using matching asymptotics, where the authors conjectured that the dynamic should be stable. Recently in [31], the authors proved the stability of the blowup and also provided a description of the dynamic, but for radial solutions. More precisely, in [31] they proved the following:

- Universality of the blowup profile: for all $t \in [0, T)$,

$$u(x, t) = \frac{1}{\lambda(t)^2} (Q + \varepsilon) \left(t, \frac{x}{\lambda(t)} \right)$$

with $\|\varepsilon\|_{H_Q^2} \rightarrow 0$ as $t \rightarrow T$, where

$$\lambda(t) = \sqrt{T - t} e^{-\frac{\sqrt{|\log(T-t)|}}{2}} + O(1).$$

- Stability of the dynamic under small perturbation in $H_Q^2(\mathbb{R}^2) \cap L^1(\mathbb{R}^2)$.

Their proof is based on the modulation theory, which is a strong method for critical problems. For example, it allows people to prove the stability and to describe the dynamic of the flow for the nonlinear Schrödinger equation when the mass is critical [23, 24], for the Schrödinger map when the energy is critical [25], and also for the 1-corotational harmonic heat flow map when the energy is critical [30].

In this paper we are interested in the critical mass case, namely $M_0 = 8\pi$. The problem when $M_0 = 8\pi$ is also energy critical

$$\mathcal{F}(u_\lambda) = \mathcal{F}(u),$$

and note that we have in this case one more conservation law

$$(1.5) \quad \frac{d}{dt} \int_{\mathbb{R}^2} u(x, t) |x|^2 dx = 0.$$

On one hand, if $M_0 = 8\pi$, $\int_{\mathbb{R}^2} u_0(x) |x|^2 dx = +\infty$ and the initial data u_0 is sufficiently close to a rescaling Q_a of Q , then the solution converges to Q_a when $t \rightarrow +\infty$ [2]. On the other hand, if $M_0 = 8\pi$ and $\int_{\mathbb{R}^2} u_0(x) |x|^2 dx < +\infty$, the solution exists globally and concentrates in infinite time [3]. The proof is based on a contradiction argument and a virial identity. When the domain is bounded in [17], the authors proved that radial solutions u collapse in 0 at the rate

$$u(0, t) = 8e^{\frac{5}{2} + 2\sqrt{2}t} (1 + O(t^{-1/2} \log(4t)))$$

as $t \rightarrow +\infty$. The authors in [17] used the partial mass equation to remove the nonlocal difficulty and an argument of subsolutions and supersolutions to bound the solution from above and below.

When the domain is unbounded the infinite speed of propagation of the heat semigroup will send some mass to infinity almost instantly, the concentration will occur only when all the mass 8π is present again in the center of mass. While all the mass concentrates at the origin when $t \rightarrow +\infty$, the second moment is ejected to infinity in space. Recently Senba [34] proved the existence of a family of radial solutions that blow up in infinite time with the speed $u(t, 0) \sim \log(t + 1)^2$. We want to emphasize that the speed found by Senba is different from the one predicted by the physicists in [10, sec. 3.3.3]. His proof is based on a technique of matching asymptotics developed by Herrero and Velazquez for the supercritical mass $M_0 > 8\pi$ in [13]. The goal of the paper is to find a speed similar to the one found by the physicists and describe the dynamic in \mathbb{R}^2 associated to that speed when the mass is critical $M_0 = 8\pi$ and to prove the stability of that dynamic under perturbations of the initial data.

THEOREM 1.1. *For all $A_0 > 0$, there exists $\mu_0^* > 0$ such that for all $\mu_0 < \mu_0^*$ and u_0 initial data of the form*

$$(1.6) \quad u_0 = \frac{1}{\mu_0} \left(Q \left(\frac{x}{\sqrt{\mu_0}} \right) e^{-\frac{|x|^2}{2}} + \tilde{\varepsilon}_0 \left(\frac{x}{\sqrt{\mu_0}} \right) \right)$$

where $\int_{\mathbb{R}^2} \frac{|\tilde{\varepsilon}_0(y)|^2}{Q} e^{\frac{|y|^2}{2}} dy < A_0 \mu_0$

and

$$(1.7) \quad \int_{\mathbb{R}^2} u_0(x) dx = 8\pi, \quad \int_{\mathbb{R}^2} u_0(x) x_i dx = 0 \quad \text{for all } i \in \{1, 2\},$$

$$I := \int_{\mathbb{R}^2} u_0(x) |x|^2 dx,$$

the corresponding solution u blows up in infinite time and satisfies:

- *Universality and stability of the profile: There exists $C > 0$ and $\lambda \in C^1(\mathbb{R}_+)$ such that for t large enough, we have*

$$(1.8) \quad u(x, t) = \frac{1}{\lambda(t)^2} \left(Q \left(\frac{x}{\lambda(t)} \right) e^{-\frac{|x|^2}{2t}} + \tilde{\varepsilon} \left(t, \frac{x}{\lambda(t)} \right) \right),$$

with

$$\int_{\mathbb{R}^2} \frac{|\tilde{\varepsilon}(t, y)|^2}{Q(y)} e^{\frac{|y|^2}{2t \log(2t+1)}} dy \leq \frac{C}{t \log(t)}$$

$$\text{where } \lambda(t) = \frac{\sqrt{I}}{\sqrt{\log(2t+1) + O(\log(\log(t)))}}.$$

REMARK 1.2. The proof requires a spectral gap bound (see Corollary 3.5), which follows from the study of the spectrum of a linearized operator [9]. As of now the proof of Corollary 3.5 in the nonradial case requires some eigenvalue bounds that were only done numerically in [9]. We hope to address this issue rigorously elsewhere.

REMARK 1.3. Notice that if we select any rescaling of u_0 initially like $(u_0)_a(x) = \frac{1}{a^2} u_0\left(\frac{x}{a}\right)$ for some $a > 0$, then the solution

$$u_a(t, x) = \frac{1}{\lambda(t/a^2)^2 a^2} \left(Q \left(\frac{x}{a \lambda(t/a^2)} \right) e^{-\frac{|x|^2}{2t}} + \tilde{\varepsilon} \left(\frac{t}{a^2}, \frac{x}{a \lambda(t/a^2)} \right) \right)$$

induced by $(u_0)_a$ still behaves asymptotically as in Theorem 1.1.

In the radial case, we can remove the smallness assumption on the initial data, using the maximum principle on the partial mass equation, namely:

COROLLARY 1.4. *Let u_0 be a radial initial data satisfying (1.7) and*

$$(1.9) \quad \int_{\mathbb{R}^2} \frac{|u_0|^2}{Q} e^{\frac{|y|^2}{2}} dy < +\infty.$$

Then there exist $C_1 > 0$, $C_2 > 0$, $\lambda_1(t)$, and $\lambda_2(t)$ such that the partial mass m_u of the corresponding solution u (see (8.1)) satisfies for t large enough and $r = |x|$

$$(1.10) \quad -\frac{C_1 \lambda_1^2 e^{-\frac{C_1 r^2}{2t}}}{\lambda_1^2 + r^2} - \frac{C_1 \lambda_1^2}{(\lambda_1^2 + r^2)t |\log t|} \leq 8\pi - m_u(r, t) \leq \frac{C_2 \lambda_2^2 e^{-\frac{C_2 r^2}{2t}}}{\lambda_2^2 + r^2} + \frac{C_2 \lambda_2^2}{(\lambda_2^2 + r^2)t |\log t|},$$

with $\lambda_i(t) = \frac{I_i}{\sqrt{\log(2t+1) + O(\log(\log(t)))}}$ and $I_i > 0$ for $i \in \{1, 2\}$.

The proof of the corollary is given in Section 8.

1.1 Strategy of the Proof

Let's explain the main ideas of the proof that we think could be applicable to many other critical problems. The main idea of the proof is to pass by the subcritical case to reach the critical one. Indeed, the blowup profile of the solution u of (1.3) is the unique stationary solution $Q = \frac{8}{(1+|x|^2)^2}$, but we cannot use the linearization around that stationary solution because we need the second moment of u to be finite and the mass to be 8π . Since Q has an infinite second moment and 8π mass, it seems natural to multiply Q by a cutoff function χ . However, χQ doesn't satisfy a good equation and has an uncontrollable error. The idea is to notice that the stationary solution of the subcritical case after rescaling converges uniformly to Q . Indeed, if we rescale (1.3) with the self-similar variables when the mass M is subcritical, there exists for that M a unique stationary solution n_∞ that depends on M and decays exponentially at infinity. Actually, after another rescaling we prove that n_∞ converges uniformly to Q when $M \rightarrow 8\pi$.

Actually, we prove that n_∞ after a certain rescaling behaves as Q multiplied by a cutoff function, but the advantage is that n_∞ is a stationary solution of some related equation (2.2). The main difficulty now is that everything depends on the mass, i.e., the linearized operator and its spectrum, the energy, and the norms that we are controlling. And since the mass is not fixed, we need that all the inequalities be uniform with respect to the mass. It appears that the energy that we are controlling is coercive but not uniformly with respect to the mass. To overcome this difficulty, we found another bound on the energy coming from a property of the nonlinear term (see Proposition 6.1). Classically, to describe the dynamic of the blowup we need to rescale the solution, and the best rescaling is the blowup speed λ . Hence, we renormalize the previous equation with the following change of variables:

$$u(t, x) = \frac{1}{\lambda^2(t)} v(s, y), \quad \frac{ds}{dt} = \frac{1}{\lambda^2(t)}, \quad \text{and} \quad y = \frac{x}{\lambda(t)},$$

where λ is an unknown to be determined later. Then we obtain

$$(1.11) \quad v_s - \frac{\lambda_s}{\lambda} \Lambda v = \nabla \cdot (\nabla v - v \nabla \phi_v) \quad \text{where} \quad \Lambda_y v = \nabla \cdot (y v).$$

Actually, we don't rescale (1.3) by λ directly, but we arrive at (1.11) in two steps. In the first step we rescale by the forward self-similar variables to make the stationary solution of the subcritical case appear. In the second step we rescale by a certain parameter μ that will be defined later such that the combination of these two rescalings leads to (1.11).

We first start in the second section to set the equations and the bootstrap argument. In the third section we prove the spectral gap estimate. In the fourth section we construct an approximate solution, and in the fifth section we do the derivation of the blowup speed λ . In the sixth section we derive a uniform bound on the potential, and in the seventh section we prove the rigorous control of the error ε .

REMARK 1.5. The idea of solving the critical problem by passing to the limit in the subcritical problem was first used by Matano and Merle in [20–22], where they construct Type II blowup solutions in the energy supercritical heat equation by passing to the limit in Type I blowup solutions of the energy subcritical heat equation. The same idea has also been used by Merle, Raphael, and Szeftel in [26], where the authors constructed scattering bubbles with subcritical mass arbitrarily close to the minimal mass blowup solution. Our result is true for a nonradial solution under a spectral gap assumption, but in the radial case the spectral gap has been proved in [9]. Moreover, the speed of concentration that we derive rigorously from the conservation of the second moment is the one found formally by the physicist [10, sec. 3.3.3].

REMARK 1.6. In this paper we are describing the dynamic of the minimal mass blowup solutions ($M = 8\pi$) in the whole plane \mathbb{R}^2 . Actually, the question of describing the dynamic of minimal mass blowup solutions for partial differential equations has become better understood over the past ten years. Especially on the gKDV equation, Martel, Merle, and Raphael in [19] proved the existence and uniqueness of a minimal mass blowup solution and gave a sharp description of the corresponding blowup. Other works in this direction have been done for NLS; for example, Raphael and Szeftel in [32] proved the existence and uniqueness of a minimal mass blowup solution for an inhomogeneous version of NLS.

REMARK 1.7. The energy estimates that we are using to close the bootstrap argument is based on spectral gaps in time-weighted norms. Actually, this has already been used by Hadzic and Raphael in [12].

1.2 Notation and Conventions

We will use the following conventions:

- C denotes a constant that may change from one line to the next and is uniform with respect to μ .
- For functions f and g , we denote $f = O(g)$ if there exists C uniform with respect to μ such that $f \leq Cg$.
- For quantities A and B , we denote $A \lesssim B$ if there exists C uniform with respect to μ such that $A \leq CB$.

- $\langle r \rangle$ denotes $\sqrt{1 + r^2}$.
- For radial functions such as Q, ϕ_Q , we keep the same notation for $Q(x)$ and $Q(r)$ where $r = |x|$ and write $\Delta\phi = \partial_{rr}\phi + \frac{1}{r}\partial_r\phi$.

2 Bootstrap Argument

2.1 The Subcritical Problem

We expect u to behave like Q at the blowup time. Linearizing u around Q (i.e., $u = Q + \varepsilon$) is not helpful, since the second moment of Q is infinite, and we are looking for solutions with finite second moment. Another alternative would be to multiply the ground state Q by a cutoff function χ to have a finite second moment. However, the error of our solution $Q\chi$ becomes too large and we cannot close our estimates. Nevertheless, if we consider the stationary solution of the subcritical mass problem n_∞ and we rescale it by a certain parameter μ , then the n_∞ rescaled by μ will converge to Q . This sequence Q_μ has much better properties as we will see later. To make n_∞ appear, we rescale (1.3) by the self-similar variables, using $z = \frac{x}{R(t)}$ and $\tau = \log R(t)$ with $R(t) = \sqrt{1 + 2t}$ such that $u(x, t) = \frac{1}{R(t)^2} w(\frac{x}{R}, \log R)$:

$$(2.1) \quad \begin{cases} \partial_\tau w - \Delta w - \Lambda w = -\nabla \cdot (w \nabla \phi_w), \\ -\Delta \phi_w = w, \end{cases}$$

where $\Lambda w = \nabla \cdot (zw)$ and

$$\int_{\mathbb{R}^2} w(z, \tau) dz = \int_{\mathbb{R}^2} u(x, t) dx = 8\pi = \int_{\mathbb{R}^2} u_0(x) dx.$$

It is well-known that (2.1) has stationary solution n_∞ of mass M for all $M < 8\pi$. Actually n_∞ is the solution of the following equation:

$$(2.2) \quad \begin{cases} \Delta n_\infty + \nabla \cdot (zn_\infty - n_\infty \nabla \phi_{n_\infty}) = 0, \\ \phi_{n_\infty} = -\frac{1}{2\pi} \log |\cdot| \star n_\infty \text{ in } \mathbb{R}^2. \end{cases}$$

The equation can also be written as

$$(2.3) \quad -\Delta \phi_{n_\infty} = n_\infty = M_{n_\infty} \frac{e^{\phi_{n_\infty} - (|z|^2/2)}}{\int_{\mathbb{R}^2} e^{\phi_{n_\infty} - (|z|^2/2)} dz}.$$

The linearized operator around n_∞ is also well-known:

$$\begin{aligned} \mathcal{L}^z w &= \Delta w + \nabla \cdot (zw - w \nabla \phi_{n_\infty} - n_\infty \nabla \phi_w) \\ &= \nabla \cdot \left(n_\infty \nabla \left(\frac{w}{n_\infty} - \phi_w \right) \right) \\ &= \nabla \cdot (n_\infty \nabla (\mathcal{M}^z w)), \end{aligned}$$

$$\mathcal{M}^z w = \frac{w}{n_\infty} - \phi_w.$$

Notice that n_∞ can be parametrized by its mass or by its maximum at $z = 0$. To emphasize the dependence of n_∞ in M , we will denote n_∞ as n_∞^M . Throughout the paper M will denote the mass of n_∞ . We will parametrize n_∞ by its maximum and set $n_\infty(0) = \frac{8}{\mu}$; then if μ goes to 0, the mass $M = M(\mu)$ will go to 8π . Hence, if we rescale n_∞ correctly with μ , the maximum of the rescaled n_∞ becomes 8, which is the maximum of Q . This rescaled n_∞ converges uniformly to Q as μ goes to 0. The following proposition summarizes the previous statements:

PROPOSITION 2.1. *Let $\mu = \frac{8}{n_\infty^M(0)}$ and $y = \frac{z}{\sqrt{\mu}}$; then $\mu n_\infty^M(\sqrt{\mu}y) \rightarrow Q(y)$ uniformly as $\mu \rightarrow 0$. Moreover,*

$$Q(y)e^{-\frac{\mu|y|^2}{2}} \leq \mu n_\infty^M(\sqrt{\mu}y) \leq Q(y) \quad \text{for all } y \in \mathbb{R}^2.$$

This is a direct consequence of Proposition 4.1. Since $\mu n_\infty^M(\sqrt{\mu}\cdot)$ converges to Q , it seems natural to rescale (2.1) by using $y = \frac{z}{\sqrt{\mu}}$, $\frac{ds}{d\tau} = \frac{1}{\mu}$ with $w(z, \tau) = \frac{1}{\mu}v(\frac{z}{\sqrt{\mu}}, s)$ where μ will be fixed later:

$$(2.4) \quad \begin{cases} \partial_s v - \Delta v - (\mu + \frac{\partial_s \mu}{2\mu})\Lambda v = -\nabla \cdot (v \nabla \phi_v), \\ -\Delta \phi_v = v. \end{cases}$$

To resume the situation we have three sets of variables (x, t) , (z, τ) , and (y, s) , which are linked in the following way:

$$\begin{aligned} z &= \frac{x}{R(t)}, & y &= \frac{z}{\sqrt{\mu}} = \frac{x}{\sqrt{\mu}R(t)}, \\ \tau(t) &= \log(R(t)), & \frac{ds}{d\tau} &= \frac{1}{\mu}, & \frac{ds}{dt} &= \frac{1}{\mu R(t)^2}, \end{aligned}$$

where $R(t) = \sqrt{1 + 2t}$. We fix $s(0) = e$ to avoid a problem if one divides by s or if one takes the log of s . Now we set

$$Q_\mu(y) = \mu n_\infty^M(\sqrt{\mu}y);$$

hence Q_μ solves

$$(2.5) \quad \Delta Q_\mu - \nabla \cdot (Q_\mu \nabla \phi_{Q_\mu}) = -\mu \nabla \cdot (y Q_\mu).$$

Let $v = Q_\mu + \varepsilon$; we will prove later that $M(\mu) = \int_{\mathbb{R}^2} Q_\mu dy = 8\pi + 2\mu \log(\mu) + O(\mu)$, which implies that the mass of ε will be of order $\mu \log(\mu)$ since the mass of v is 8π . We need ε to be at least of order μ to close our estimates; hence we add a correction to Q_μ . We set $\tilde{Q}_\mu = Q_\mu + T_\mu(y)$ where T_μ is a correction to insure that \tilde{Q}_μ is of mass 8π . It will be constructed in Section 4.

2.2 Modulation

In this subsection, we explain how we fix the parameter μ . We consider initial data of the form

$$v_0 = \tilde{Q}_{\mu_0} + \varepsilon_0,$$

with 8π mass, a finite second moment I such that $\int_{\mathbb{R}^2} \varepsilon_0 dy = 0$, and $\mu_0 > 0$ is sufficiently small. We will show the existence and uniqueness of the following decomposition of our solution:

$$(2.6) \quad v(y, s) = \tilde{Q}_\mu + \varepsilon,$$

where the mass of ε is 0. Set

$$(2.7) \quad L^2_{\tilde{Q}_\mu}(\mathbb{R}^2) = \left\{ f \in L^2(\mathbb{R}^2) \text{ such that } \int_{\mathbb{R}^2} \frac{f^2}{\tilde{Q}_\mu} dx < \infty \right\}$$

and

$$(2.8) \quad L^2_{\tilde{Q}_\mu, 0}(\mathbb{R}^2) = \left\{ f \in L^2_{\tilde{Q}_\mu}(\mathbb{R}^2) \text{ such that } \int_{\mathbb{R}^2} f(x) dx = 0 \right\}.$$

In the following lemma we fix μ such that $(\varepsilon, |\cdot|^2)_{L^2} = 0$.

LEMMA 2.2 (Modulation). *For all $A > 0$, there exists $\bar{\mu} > 0$ such that for all $\mu_* \in (0, \bar{\mu})$, $v \in L^2_{\tilde{Q}_{\mu_*}}$ such that*

$$(2.9) \quad \|v - \tilde{Q}_{\mu_*}\|_{L^2_{\tilde{Q}_{\mu_*}}} \leq \delta(\mu_*) := \sqrt{A\mu_*},$$

there exist a unique $\mu > 0$ and a unique ε such that

$$v(y) = \tilde{Q}_\mu(y) + \varepsilon(y),$$

and

$$(\varepsilon, |\cdot|^2)_{L^2} = 0.$$

In addition,

$$\|\varepsilon\|_{L^2_{\tilde{Q}_{\mu_*}}} \lesssim \delta(\mu_*),$$

and

$$|\mu - \mu_*| \lesssim \mu_*^{3/2} |\log(\mu_*)|^{1/2}.$$

PROOF OF LEMMA 2.2. The proof is based on a careful use of the Intermediate Value Theorem. Given v satisfying (2.9), consider the C^1 function

$$F(\mu) = (v - \tilde{Q}_\mu, |\cdot|^2)_{L^2}.$$

From (5.16) we deduce

$$\partial_\mu F|_{\mu=\mu_*} = ((\partial_\mu \tilde{Q}_\mu)|_{\mu=\mu_*}, |\cdot|^2) = -\frac{C}{\mu_*} + O(|\log \mu_*|^2) \neq 0,$$

and from (5.17)

$$|\partial_{\mu\mu} F| \lesssim \frac{1}{\mu^2}.$$

It follows for $|\mu - \mu_*| \leq \frac{\mu_*}{2}$,

$$F(\mu) = F(\mu_*) + \partial_\mu F(\mu_*)(\mu - \mu_*) + O\left(\frac{1}{\mu_*^2}\right)(\mu - \mu_*)^2,$$

and since

$$F(\mu_*) \leq \|v - \tilde{Q}_{\mu_*}\|_{L^2_{\tilde{Q}_{\mu_*}}} \left(\int_{\mathbb{R}^2} |y|^2 \tilde{Q}_{\mu_*} dy \right)^{\frac{1}{2}} \lesssim \sqrt{\mu_*} |\log \mu_*|^{\frac{1}{2}},$$

we deduce that there exists C^* such that

$$\mu_{\pm} = \mu_* \pm C^* \mu_*^{\frac{3}{2}} |\log \mu_*|^{\frac{1}{2}}$$

satisfy

$$F(\mu_-) < 0 \quad \text{and} \quad F(\mu_+) > 0.$$

Hence, by the Intermediate Value Theorem there exists $\mu \in (\mu_-, \mu_+)$ such that

$$F(\mu) = 0,$$

which concludes the proof. The proof of (5.16) and (5.17) is given in Section 5. \square

2.3 Setting Up the Bootstrap Argument

Let u_0 be as in Theorem 1.1 and let u be the global free energy solution constructed in [3]. Consequently, thanks to Lemma 2.2 the solution admits a unique decomposition on some small time interval $[0, T^*)$,

$$(2.10) \quad u(x, t) = \frac{1}{R^2(t)\mu(t)} \left(\tilde{Q}_{\mu} \left(\frac{x}{R(t)\sqrt{\mu(t)}} \right) + \varepsilon \left(\frac{x}{R(t)\sqrt{\mu(t)}}, t \right) \right),$$

where

$$(\varepsilon(t), 1) = (\varepsilon(t), |\cdot|^2) = 0,$$

and from standard argument (see [18]) $\mu(t)$ satisfies $\mu \in C^1([0, T^*))$. In the rest of this paper, we will abuse notation and consider μ also as a function of s . Using the initial smallness assumption on ε_0 and μ_0 , we prove the following proposition:

PROPOSITION 2.3. *Let u_0 be as in Theorem 1.1 and let u be the global free energy solution constructed in [3]. Assume there exists S^* such that for all $s \in [e, S^*)$, there exist $\eta > 0$ and $\mu := \mu(s)$, a C^1 function from $[e, S^*)$ to \mathbb{R}_+ , $\mu(e) = \mu_0$, and*

$$(2.11) \quad \left| \mu(s) - \frac{1}{2s} \right| < \frac{A_1}{s |\log s|}.$$

Let us define $v(s, y) = \mu R^2 u(t, x) = \tilde{Q}_{\mu}(y) + \varepsilon(s, y)$ where $y = \frac{x}{R\sqrt{\mu}}$ and $\frac{ds}{dt} = \frac{1}{\mu R^2}$. If we assume that for all $s \in [e, S^)$*

$$(2.12) \quad \|\varepsilon\|_{L^2_{\tilde{Q}_{\mu}}}^2 \leq A\mu,$$

then the regime is trapped; in particular when $\frac{S^}{2} < s \leq S^*$ we obtain the same inequalities (2.11) and (2.12) with A_1 and A replaced by $\frac{A_1}{2}$ and $\frac{A}{2}$.*

Hence, if one can increase by continuity the size of the interval $[e, S^*)$, it will induce that we can also increase the size of $[0, T^*)$.

3 The Linearized Operator and Its Spectral Gap

In this section, we describe the properties of the linearized operator \mathcal{L}_μ^y obtained from the linearization around Q_μ .

3.1 Setting Up the Equations

If we plug $v(y, s) = \tilde{Q}_\mu(y) + \varepsilon(y, s)$ into (2.4) we get

$$(3.1) \quad \begin{aligned} \partial_s \varepsilon + \partial_s \tilde{Q}_\mu = & -\nabla \cdot (\varepsilon \nabla \phi_\varepsilon) + \left(\mu + \frac{\mu_s}{2\mu} \right) (\Lambda \varepsilon + \Lambda \tilde{Q}_\mu) + \Delta \tilde{Q}_\mu \\ & + \Delta \varepsilon - \nabla \cdot (\tilde{Q}_\mu \nabla \phi_\varepsilon) - \nabla \cdot (\varepsilon \nabla \phi_{\tilde{Q}_\mu}) - \nabla \cdot (\tilde{Q}_\mu \nabla \phi_{\tilde{Q}_\mu}). \end{aligned}$$

Hence,

$$(3.2) \quad \partial_s \varepsilon = \mathcal{L}_\mu^y \varepsilon + \frac{\mu_s}{2\mu} \Lambda \varepsilon + \Theta_\mu(\varepsilon) + N(\varepsilon) + \tilde{E},$$

where

$$\mathcal{L}_\mu^y \varepsilon = \nabla \cdot (Q_\mu \nabla (\mathcal{M}_\mu^y \varepsilon)) = \Delta \varepsilon + \mu \Lambda \varepsilon - \nabla \cdot (\varepsilon \nabla \phi_{Q_\mu} + Q_\mu \nabla \phi_\varepsilon);$$

the previous equality comes from the fact that $\frac{\nabla Q_\mu}{Q_\mu} = \nabla \phi_{Q_\mu} - \mu y$, which we will prove in Proposition 3.1. In addition,

$$\mathcal{M}_\mu^y \varepsilon = \frac{\varepsilon}{Q_\mu} - \phi_\varepsilon.$$

Note that \mathcal{M}_μ^y corresponds to the linearization around Q_μ of the free energy

$$\mathcal{F}(u) = \int_{\mathbb{R}^2} u \log u \, dx - \frac{1}{2} \int_{\mathbb{R}^2} u \phi_u \, dx.$$

The error term \tilde{E} comes from the fact that \tilde{Q}_μ is not an exact solution of (2.4). It is given by

$$(3.3) \quad \tilde{E} = \Delta \tilde{Q}_\mu + \mu \Lambda \tilde{Q}_\mu - \nabla \cdot (\tilde{Q}_\mu \nabla \phi_{\tilde{Q}_\mu}) + \frac{\mu_s}{2\mu} \Lambda \tilde{Q}_\mu - \partial_s \tilde{Q}_\mu.$$

We split it into two parts:

$$\tilde{E} = E + F$$

with

$$E = \Delta \tilde{Q}_\mu + \mu \Lambda \tilde{Q}_\mu - \nabla \cdot (\tilde{Q}_\mu \nabla \phi_{\tilde{Q}_\mu}) \quad \text{and} \quad F = \frac{\mu_s}{2\mu} \Lambda \tilde{Q}_\mu - \partial_s \tilde{Q}_\mu.$$

The nonlinear term $N(\varepsilon)$ is given by

$$(3.4) \quad N(\varepsilon) = -\nabla \cdot (\varepsilon \nabla \phi_\varepsilon).$$

The linear term Θ_μ measures the error due to the fact that we are linearizing around \tilde{Q}_μ and not around Q_μ . It is given by

$$\Theta_\mu(\varepsilon) = \nabla \cdot [(\tilde{Q}_\mu - Q_\mu) \nabla \phi_\varepsilon + \varepsilon \nabla \phi_{\tilde{Q}_\mu - Q_\mu}].$$

3.2 The Linearized Operator \mathcal{L}_μ^y

We have the following relations between \mathcal{L}_μ^y and \mathcal{L}^z :

$$\mathcal{L}_\mu^y(v(y)) = \mu \mathcal{L}^z\left(\frac{1}{\mu} v\left(\frac{z}{\sqrt{\mu}}\right)\right).$$

We now prove fundamental algebraic relations on Q_μ , \mathcal{M}_μ^y , and \mathcal{L}_μ^y .

PROPOSITION 3.1. *Let Q_μ be the solution of (2.5). Then ϕ_{Q_μ} solves*

$$(3.5) \quad -\Delta \phi_{Q_\mu} = -\phi_{Q_\mu}'' - \frac{1}{r} \phi_{Q_\mu}' = Q_\mu = 8e^{\phi_{Q_\mu} - \mu \frac{|y|^2}{2}};$$

- Algebraic identities for \mathcal{M}_μ^y :

$$\mathcal{M}_\mu^y(\partial_\mu Q_\mu) = -\frac{|y|^2}{2}, \quad \mathcal{M}_\mu^y(\partial_i Q_\mu) = -\mu y_i,$$

$$\mathcal{M}_\mu^y(\Lambda Q_\mu) = 2 - \frac{M}{2\pi} - \mu |y|^2,$$

$$M'(\mu) \mu \partial_M n_\infty(\sqrt{\mu} y) = \partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu,$$

$$\mathcal{M}_\mu^y\left(\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu\right) = -\left(2 - \frac{M}{2\pi}\right).$$

- Algebraic identities for \mathcal{L}_μ^y :

$$\mathcal{L}_\mu^y(\partial_\mu Q_\mu) = -\Lambda(Q_\mu), \quad \mathcal{L}_\mu^y(\partial_i Q_\mu) = -\mu \partial_i Q_\mu,$$

$$\mathcal{L}_\mu^y(\Lambda Q_\mu) = -2\mu \Lambda Q_\mu,$$

$$\mathcal{L}_\mu^y\left(\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu\right) = 0.$$

PROOF.

- Derivation of (3.5).

Let us recall the properties of $n_\infty^M(z)$ for $M < 8\pi$:

$$-\Delta \phi_{n_\infty^M} = n_\infty^M = M \frac{e^{\phi_{n_\infty^M} - (|z|^2/2)}}{\int_{\mathbb{R}^2} e^{\phi_{n_\infty^M} - (|z|^2/2)} dz},$$

with $M = \int_{\mathbb{R}^2} n_\infty^M dz < 8\pi$ and

$$n_\infty^M(0) = M \frac{e^{\phi_{n_\infty^M}(0)}}{\int_{\mathbb{R}^2} e^{\phi_{n_\infty^M} - (|z|^2/2)} dz}$$

or

$$\phi_{n_\infty^M}(0) = \log\left(n_\infty^M(0) \int_{\mathbb{R}^2} e^{\phi_{n_\infty^M} - (|z|^2/2)} dz\right) - \log(M).$$

Hence if we set $\tilde{\phi}_{n_\infty^M} = \alpha + \phi_{n_\infty^M}$ where

$$\alpha = \log(M) - \log\left(\int_{\mathbb{R}^2} e^{\phi_{n_\infty^M} - (|z|^2/2)} dz\right),$$

then we get

$$-\Delta \tilde{\phi}_{n_\infty^M} = n_\infty^M(z) = e^{\tilde{\phi}_{n_\infty^M} - (|z|^2/2)},$$

with $\tilde{\phi}_{n_\infty^M}(0) = \alpha + \phi_{n_\infty^M}(0) = \log(n_\infty^M(0))$. Now we see that if we choose $\mu = \frac{8}{n_\infty^M(0)}$, $y = \frac{z}{\sqrt{\mu}}$, and

$$\phi_{Q_\mu}(y) = \tilde{\phi}_{n_\infty^M}(\sqrt{\mu}y) - \log(n_\infty^M(0)),$$

then ϕ_{Q_μ} solves (3.5) with $\phi_{Q_\mu}(0) = 0$ and $\phi'_{Q_\mu}(0) = 0$.

- Differentiation of (3.5) with respect to y .

Moreover, we have by differentiating (3.6) in y

$$(3.6) \quad \begin{aligned} \log(Q_\mu) &= \log(8) + \phi_{Q_\mu} - \mu \frac{|y|^2}{2}, \\ \frac{\partial_i Q_\mu}{Q_\mu} &= \phi_{\partial_i Q_\mu} - \mu y_i, \end{aligned}$$

then

$$\mathcal{M}_\mu^y(\partial_i Q_\mu) = -\mu y_i.$$

Hence,

$$\mathcal{L}_\mu^y(\partial_i Q_\mu) = -\mu \partial_i Q_\mu.$$

- Differentiation of (3.5) with respect to μ .

We see easily by differentiating (3.6) in μ that

$$(3.7) \quad \frac{\partial_\mu Q_\mu}{Q_\mu} = \phi_{\partial_\mu Q_\mu} - \frac{|y|^2}{2},$$

which means

$$\mathcal{M}_\mu^y(\partial_\mu Q_\mu) = -\frac{|y|^2}{2}.$$

It follows that

$$\mathcal{L}_\mu^y(\partial_\mu Q_\mu) = -\nabla \cdot (y Q_\mu).$$

Actually, we differentiated Q_μ with respect to μ without justifying that Q_μ is smooth with respect to μ . However, this will be done later in Section 4 (see (4.21)).

- Differentiation of (3.5) with respect to λ .

In addition, if we set $Q_\mu^\lambda(y) = \lambda^2 Q_\mu(\lambda y)$ where $\lambda > 0$, then

$$\phi_{Q_\mu}(\lambda y) = -\frac{M(\mu)}{2\pi} \log(\lambda) + \phi_{Q_\mu^\lambda}(y).$$

Using (3.6), we also get

$$\log(Q_\mu^\lambda) = \log(8) + \phi_{Q_\mu^\lambda} + \left(2 - \frac{M}{2\pi}\right) \log(\lambda) - \mu \lambda^2 \frac{|y|^2}{2}.$$

If we differentiate the previous equation in λ and set $\lambda = 1$, we deduce

$$(3.8) \quad \frac{\Lambda Q_\mu}{Q_\mu} - \phi_{\Lambda Q_\mu} = \mathcal{M}_\mu^y(\Lambda Q_\mu) = 2 - \frac{M}{2\pi} - \mu|y|^2.$$

It follows that

$$\mathcal{L}_\mu^y(\Lambda Q_\mu) = -2\mu \Lambda Q_\mu.$$

The relation $M'(\mu)\mu \partial_M n_\infty^M(\sqrt{\mu}y) = \partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu$ is a consequence of an easy computation. From this relation and $\mathcal{L}_\mu^y(\partial_\mu Q_\mu) = -\Lambda(Q_\mu)$, $\mathcal{L}_\mu^y(\Lambda Q_\mu) = -2\mu \Lambda Q_\mu$, we deduce that $\mathcal{L}_\mu^y(\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu) = 0$. One can also recover this from the fact that

$$\mathcal{L}^z(\partial_M n_\infty^M) = 0 \Rightarrow \mathcal{L}_\mu^y\left(\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu\right) = 0. \quad \square$$

LEMMA 3.2. *The operator $Q_\mu \mathcal{M}_\mu^y : L_{Q_\mu}^2 \longrightarrow L_{Q_\mu}^2$, given by*

$$Q_\mu \mathcal{M}_\mu^y u = u - Q_\mu \phi_u,$$

is linear continuous, self-adjoint, and Fredholm. Moreover,

$$\int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y u|^2 \leq K^* \|u\|_{L_{Q_\mu}^2}^2,$$

with K^ uniform with respect to μ ,*

$$\forall (u, v) \in L_{Q_\mu}^2 \times L_{Q_\mu}^2 \quad (\mathcal{M}_\mu^y u, v) = (u, \mathcal{M}_\mu^y v).$$

In addition, $Q_\mu \mathcal{M}_\mu^y : L_{Q_\mu, 0}^2 \longrightarrow L_{Q_\mu}^2$ is positive definite and defines an inner product on $L_{Q_\mu, 0}^2$.

PROOF. We use the following inequality from [31]:

$$\|\phi_u\|_{L^\infty(\{r \leq 1\})} + \left\| \frac{\phi_u}{1 + |\log(r)|} \right\|_{L^\infty(\{r \geq 1\})} \lesssim \|u\|_{L_Q^2}.$$

Hence,

$$\begin{aligned} \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y u|^2 dy &\lesssim \|u\|_{L_{Q_\mu}^2}^2 + \left\| \frac{\phi_u}{1 + |\log|y||} \right\|_{L^\infty}^2 \int_{\mathbb{R}^2} (1 + |\log(|y|)|) Q_\mu dy \\ &\lesssim \|u\|_{L_{Q_\mu}^2}^2. \end{aligned}$$

The self-adjointness follows from the fact that

$$\text{for all } (u, v) \in L_{Q_\mu}^2 \times L_{Q_\mu}^2 \quad (\phi_u, v)_{L^2} = (u, \phi_v)_{L^2}.$$

It is easy to see that the operator is in the form $I - Q_\mu \phi_{(\cdot)}$ with $Q_\mu \phi_{(\cdot)}$ a compact operator, which implies that it is Fredholm. To prove that $(\mathcal{M}_\mu^y v, v) \geq 0$ for any $v \in L_{Q_\mu, 0}^2$, we recall that the free energy $\mathcal{F}(v)$ achieves its minimum at $v = Q_\mu$ according to the logarithmic Hardy-Littlewood-Sobolev inequality and observe that

$$(3.9) \quad (\mathcal{M}_\mu^y v, v) = \lim_{\delta \rightarrow 0} \frac{1}{\delta^2} \mathcal{F}(Q_\mu + \delta v) \geq 0$$

for any smooth function v compactly supported and satisfying $\int_{\mathbb{R}^2} v \, dy = 0$. Hence, by density (3.9) holds for any v in $L_{Q_\mu, 0}^2(\mathbb{R}^2)$. \square

3.3 The Spectrum of \mathcal{L}_μ^y

Let's define the domain of \mathcal{L}_μ^y ,

$$D(\mathcal{L}_\mu^y) = \{f \in L_{Q_\mu}^2 \text{ such that } \mathcal{L}_\mu^y(f) \in L_{Q_\mu}^2\} \subset L_{Q_\mu}^2.$$

PROPOSITION 3.3. \mathcal{L}_μ^y is a self-adjoint operator and has discrete spectrum on $D_0(\mathcal{L}_\mu^y) \subset L_{Q_\mu, 0}^2$ for the inner product $(\cdot, \cdot)_{\mathcal{M}_\mu^y}$, where

$$D_0(\mathcal{L}_\mu^y) = \{f \in L_{Q_\mu, 0}^2 \text{ such that } \mathcal{L}_\mu^y(f) \in L_{Q_\mu, 0}^2\}.$$

Moreover, in $D(\mathcal{L}_\mu^y)$,

$$(3.10) \quad \ker(\mathcal{L}_\mu^y) = \text{span}\left(\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu\right),$$

$$(3.11) \quad \mathcal{L}_\mu^y(\partial_i Q_\mu) = -\mu \partial_i Q_\mu,$$

$$(3.12) \quad \mathcal{L}_\mu^y(\Lambda Q_\mu) = -2\mu \Lambda Q_\mu.$$

Furthermore, Proposition 3.3 was proved in [9] (using the z -coordinate). Campos and Dolbeault obtained the following spectral gap inequality [9]:

THEOREM 3.4. For all $w \in L_{n_\infty^M}^2(\mathbb{R}^2)$, if

$$(3.13) \quad (w, \mathcal{M}^z \partial_M n_\infty^M)_{L^2} = C_1(M)(w, 1) = 0,$$

then

$$(3.14) \quad (\mathcal{M}^z w, w) \leq -(\mathcal{L}^z w, \mathcal{M}^z w).$$

There exists $K_2 > 2$ such that if in addition

$$(w, \mathcal{M}^z(\partial_1 n_\infty^M))_{L^2} = (w, \mathcal{M}^z(\partial_2 n_\infty^M))_{L^2} = (w, \mathcal{M}^z \Lambda n_\infty^M)_{L^2} = 0,$$

then

$$(3.15) \quad K_2(\mathcal{M}^z w, w) \leq -(\mathcal{L}^z w, \mathcal{M}^z w).$$

In the radial case, (3.15) follows from the study of the spectrum of \mathcal{L}^z performed in [9]. Moreover, the numerics performed in [9] (see figure 1 in [9]) show that it is also true in the nonradial case. We hope to address this elsewhere. We rewrite Theorem 3.4 in the y -coordinate with $y = \frac{z}{\sqrt{\mu}}$.

COROLLARY 3.5. *For all $w \in L^2_{Q_\mu}(\mathbb{R}^2)$, if*

$$(3.16) \quad \left(w, \mathcal{M}_\mu^y \left(\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu \right) \right) = C_2(\mu)(w, 1) = 0,$$

then

$$(3.17) \quad \mu(\mathcal{M}_\mu^y w, w) \leq -(\mathcal{L}_\mu^y w, \mathcal{M}_\mu^y w).$$

There exists $K_2 > 2$ such that if in addition

$$(w, \mathcal{M}_\mu^y(\partial_1 Q_\mu))_{L^2} = (w, \mathcal{M}_\mu^y(\partial_2 Q_\mu))_{L^2} = (w, \mathcal{M}_\mu^y \Lambda Q_\mu)_{L^2} = 0,$$

then

$$(3.18) \quad K_2 \mu(\mathcal{M}_\mu^y w, w) \leq -(\mathcal{L}_\mu^y w, \mathcal{M}_\mu^y w).$$

REMARK 3.6. Notice here that the conditions $(w, 1)_{L^2} = (w, |\cdot|^2)_{L^2} = 0$ correspond in the inner product $(\cdot, \cdot)_{\mathcal{M}_\mu^y}$ where \mathcal{L}_μ^y is self-adjoint to the orthogonality on ΛQ_μ and on $\partial_\mu Q_\mu - \frac{1}{2\mu} \Lambda Q_\mu$ the kernel of \mathcal{L}_μ^y . Indeed,

$$(w, \mathcal{M}_\mu^y \Lambda Q_\mu) = \left(2 - \frac{M}{2\pi} \right) (w, 1) - \mu(w, |\cdot|^2) = 0$$

and

$$(w, \mathcal{M}_\mu^y \partial_\mu Q_\mu) = -\frac{1}{2}(w, |\cdot|^2) = 0.$$

In addition, we get for free from the conservation of the center of mass that $(w, y_i)_{L^2} = 0$, which corresponds to the orthogonality on $\partial_i Q_\mu$ in the inner product $(\cdot, \cdot)_{\mathcal{M}_\mu^y}$. Indeed,

$$(w, \mathcal{M}_\mu^y \partial_i Q_\mu) = -\mu \int_{\mathbb{R}^2} w y_i dy = 0.$$

4 Construction of the Approximate Profile

4.1 Properties of Q_μ

In this section, we construct an adequate approximate profile \tilde{Q}_μ that has finite second moment and 8π mass. Recall that μ and M are related by $\mu = \frac{8}{n_\infty^M(0)}$. In the following proposition, we prove that the rescaled n_∞^M by μ converges uniformly to Q the unique stationary solution of (1.1).

PROPOSITION 4.1. *Let n_∞^M be the solution of (2.2), then $\mu n_\infty^M(\sqrt{\mu}y) \rightarrow Q(y)$ uniformly as $\mu \rightarrow 0$. Moreover,*

$$(4.1) \quad \begin{aligned} \phi_Q - \mu \frac{|y|^2}{2} &< \phi_{Q_\mu}(y) - \frac{\mu|y|^2}{2} \\ &< \phi_Q(y) < \phi_{Q_\mu}(y) < 0 \quad \text{for all } y \neq 0, \end{aligned}$$

$$(4.2) \quad \begin{aligned} y \cdot \nabla \phi_{Q_\mu}(y) - \mu|y|^2 &< y \cdot \nabla \phi_Q(y) \\ &< y \cdot \nabla \phi_{Q_\mu}(y) < 0 \quad \text{for all } y \neq 0, \end{aligned}$$

$$(4.3) \quad Q(y)e^{-\frac{\mu|y|^2}{2}} < \mu n_\infty^M(\sqrt{\mu}y) = Q_\mu(y) < Q(y) \quad \text{for all } y \in \mathbb{R}^2.$$

PROOF. We set $r = |y|$ and recall that ϕ_Q is characterized by

$$(4.4) \quad -\phi_Q'' - \frac{1}{r}\phi_Q' = 8e^{\phi_Q} = Q \quad \text{with } \phi_Q(0) = 0 \text{ and } \phi_Q'(0) = 0.$$

The unique solution of (4.4) is given by $\phi_Q(y) = -2\log(1 + |y|^2)$. From Proposition 3.1, we also have

$$-\Delta \phi_{Q_\mu} = 8e^{\phi_{Q_\mu}} e^{-\frac{\mu r^2}{2}} = Q_\mu.$$

We will prove that

$$\phi_Q - \frac{\mu r^2}{2} < \phi_{Q_\mu} - \frac{\mu r^2}{2} < \phi_Q < \phi_{Q_\mu} < 0 \quad \forall r \neq 0.$$

After taking the exponential, we get that

$$Qe^{-\frac{\mu|y|^2}{2}} < Q_\mu(y) < Q(y).$$

Indeed,

$$(4.5) \quad \begin{cases} \Delta \phi_Q + 8e^{\phi_Q} = 0, \\ \Delta \phi_{Q_\mu} + 8e^{\phi_{Q_\mu}} e^{-\frac{\mu r^2}{2}} = 0, \end{cases}$$

with $\phi_Q(0) = \phi_{Q_\mu}(0) = 0 = \phi_Q'(0) = \phi_{Q_\mu}'(0)$. Since $\phi_{Q_\mu}(0) = 0 = \phi_Q(0)$, from

$$(4.6) \quad \begin{cases} \phi_Q(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau e^{\phi_Q(\tau)} d\tau dr', \\ \phi_{Q_\mu}(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau e^{\phi_{Q_\mu}(\tau)} e^{-\frac{\mu \tau^2}{2}} d\tau dr', \end{cases}$$

we deduce easily that $\phi_Q(y) < 0$, $\phi_Q'(y) < 0$, and that $\phi_{Q_\mu}(y) < 0$, $\phi_{Q_\mu}'(y) < 0$. It follows that ϕ_{Q_μ} and ϕ_Q are radially symmetric and decreasing. Set $\psi_\mu(y) = \phi_{Q_\mu} - \frac{\mu|y|^2}{2}$.

PROOF OF $\phi_Q - \frac{\mu r^2}{2} < \phi_{Q_\mu} - \frac{\mu r^2}{2} < \phi_Q < \phi_{Q_\mu} < 0$ FOR r IN A NEIGHBORHOOD OF 0.

We first expand $\phi''_{Q_\mu}, \phi'_{Q_\mu}, \phi_Q, \phi_{Q_\mu}$, and ψ_μ around 0:

$$\begin{aligned}\phi_Q(r) &= -2r^2 + r^4 + o(r^4), \\ \phi''_{Q_\mu}(r) &= \phi''_{Q_\mu}(0) + r\phi_{Q_\mu}^{(3)}(0) + \frac{r^2}{2}\phi_{Q_\mu}^{(4)}(0) + O(r^3), \\ \phi'_{Q_\mu}(r) &= r\phi''_{Q_\mu}(0) + \frac{r^2}{2}\phi_{Q_\mu}^{(3)}(0) + \frac{r^3}{6}\phi_{Q_\mu}^{(4)}(0) + O(r^4), \\ \phi_{Q_\mu}(r) &= \sum_{n=2}^4 \frac{r^n}{n!}\phi_{Q_\mu}^{(n)}(0) + o(r^4),\end{aligned}$$

and

$$\psi_\mu(r) = \sum_{n=0}^4 \frac{r^n}{n!}\psi_\mu^{(n)}(0) + o(r^4).$$

Then, to compare the three functions we just need to find the values of the consecutive derivatives of ϕ_{Q_μ} at 0. To do so, we plug the expansions of $\phi''_{Q_\mu}, \phi'_{Q_\mu}$, and ϕ_{Q_μ} around zero in

$$-\phi''_{Q_\mu} - \frac{1}{r}\phi'_{Q_\mu} = 8e^{\phi_{Q_\mu} - \mu\frac{r^2}{2}};$$

we get

$$\phi''_{Q_\mu}(0) = -4, \quad \phi_{Q_\mu}^{(3)}(0) = 0, \quad \text{and} \quad \phi_{Q_\mu}^{(4)}(0) = 6(4 + \mu).$$

The conclusion follows by substituting the values of the derivatives of ϕ_{Q_μ} in the expansion of ϕ_{Q_μ} and ψ_μ around 0.

PROOF OF $\phi_Q - \frac{\mu r^2}{2} < \phi_{Q_\mu} - \frac{\mu r^2}{2} < \phi_Q < \phi_{Q_\mu} < 0$ FOR r IN A NEIGHBORHOOD OF 0.

We first expand $\phi''_{Q_\mu}, \phi'_{Q_\mu}, \phi_Q, \phi_{Q_\mu}$, and ψ_μ around 0:

$$\begin{aligned}\phi_Q(r) &= -2r^2 + r^4 + o(r^4), \\ \phi''_{Q_\mu}(r) &= \phi''_{Q_\mu}(0) + r\phi_{Q_\mu}^{(3)}(0) + \frac{r^2}{2}\phi_{Q_\mu}^{(4)}(0) + O(r^3), \\ \phi'_{Q_\mu}(r) &= r\phi''_{Q_\mu}(0) + \frac{r^2}{2}\phi_{Q_\mu}^{(3)}(0) + \frac{r^3}{6}\phi_{Q_\mu}^{(4)}(0) + O(r^4), \\ \phi_{Q_\mu}(r) &= \sum_{n=2}^4 \frac{r^n}{n!}\phi_{Q_\mu}^{(n)}(0) + o(r^4),\end{aligned}$$

and

$$\psi_\mu(r) = \sum_{n=0}^4 \frac{r^n}{n!}\psi_\mu^{(n)}(0) + o(r^4).$$

Then, to compare the three functions we just need to find the values of the consecutive derivatives of ϕ_{Q_μ} at 0. To do so, we plug the expansions of ϕ''_{Q_μ} , ϕ'_{Q_μ} , and ϕ_{Q_μ} around zero in

$$-\phi''_{Q_\mu} - \frac{1}{r}\phi'_{Q_\mu} = 8e^{\phi_{Q_\mu} - \mu\frac{r^2}{2}}.$$

We get

$$\phi''_{Q_\mu}(0) = -4, \quad \phi^{(3)}_{Q_\mu}(0) = 0, \quad \text{and} \quad \phi^{(4)}_{Q_\mu}(0) = 6(4 + \mu).$$

The conclusion follows by substituting the values of the derivatives of ϕ_{Q_μ} in the expansion of ϕ_{Q_μ} and ψ_μ around 0.

PROOF OF $\phi_Q - \frac{\mu r^2}{2} < \phi_{Q_\mu} - \frac{\mu r^2}{2} < \phi_Q < \phi_{Q_\mu} < 0$ FOR r IN A NEIGHBORHOOD OF 0.

We first expand ϕ''_{Q_μ} , ϕ'_{Q_μ} , ϕ_Q , ϕ_{Q_μ} , and ψ_μ around 0:

$$\begin{aligned} \phi_Q(r) &= -2r^2 + r^4 + o(r^4), \\ \phi''_{Q_\mu}(r) &= \phi''_{Q_\mu}(0) + r\phi^{(3)}_{Q_\mu}(0) + \frac{r^2}{2}\phi^{(4)}_{Q_\mu}(0) + O(r^3), \\ \phi'_{Q_\mu}(r) &= r\phi''_{Q_\mu}(0) + \frac{r^2}{2}\phi^{(3)}_{Q_\mu}(0) + \frac{r^3}{6}\phi^{(4)}_{Q_\mu}(0) + O(r^4), \\ \phi_{Q_\mu}(r) &= \sum_{n=2}^4 \frac{r^n}{n!}\phi^{(n)}_{Q_\mu}(0) + o(r^4), \end{aligned}$$

and

$$\psi_\mu(r) = \sum_{n=0}^4 \frac{r^n}{n!}\psi_\mu^{(n)}(0) + o(r^4).$$

Then, to compare the three functions we just need to find the values of the consecutive derivatives of ϕ_{Q_μ} at 0. To do so, we plug the expansions of ϕ''_{Q_μ} , ϕ'_{Q_μ} , and ϕ_{Q_μ} around zero in

$$-\phi''_{Q_\mu} - \frac{1}{r}\phi'_{Q_\mu} = 8e^{\phi_{Q_\mu} - \mu\frac{r^2}{2}}.$$

We get

$$\phi''_{Q_\mu}(0) = -4, \quad \phi^{(3)}_{Q_\mu}(0) = 0, \quad \text{and} \quad \phi^{(4)}_{Q_\mu}(0) = 6(4 + \mu).$$

The conclusion follows by substituting the values of the derivatives of ϕ_{Q_μ} in the expansion of ϕ_{Q_μ} and ψ_μ around 0.

PROOF OF $\phi_Q - \frac{\mu r^2}{2} < \phi_{Q_\mu} - \frac{\mu r^2}{2} < \phi_Q < \phi_{Q_\mu} < 0$ FOR ALL $r > 0$.

Let's argue by contradiction. Suppose there exists $r_0 > 0$ such that $\phi_Q(r_0) =$

$\phi_{\mathcal{Q}_\mu}(r_0)$ and $\phi_{\mathcal{Q}_\mu}(r) > \phi_{\mathcal{Q}}(r) > \psi_\mu(r)$ for all $0 < r < r_0$. We use that $\phi_{\mathcal{Q}}$ and $\phi_{\mathcal{Q}_\mu}$ are solutions of

$$(4.7) \quad \begin{cases} (r\phi'_{\mathcal{Q}})' + 8re^{\phi_{\mathcal{Q}}} = 0, \\ (r\phi'_{\mathcal{Q}_\mu})' + 8re^{\phi_{\mathcal{Q}_\mu} - \frac{\mu r^2}{2}} = 0, \end{cases}$$

which can be solved, using the conditions at 0:

$$(4.8) \quad \begin{cases} \phi_{\mathcal{Q}}(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau e^{\phi_{\mathcal{Q}}(\tau)} d\tau dr' \\ \phi_{\mathcal{Q}_\mu}(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau e^{\psi_\mu(\tau)} d\tau dr'. \end{cases}$$

It follows that

$$\phi_{\mathcal{Q}}(r) - \phi_{\mathcal{Q}_\mu}(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau [e^{\phi_{\mathcal{Q}}(\tau)} - e^{\psi_\mu(\tau)}] d\tau dr',$$

and since $e^{\phi_{\mathcal{Q}}} - e^{\psi_\mu} > 0$ for $r \in (0, r_0)$, we get

$$\phi_{\mathcal{Q}}(r_0) - \phi_{\mathcal{Q}_\mu}(r_0) = 0 = -8 \int_0^{r_0} \frac{1}{r'} \int_0^{r'} \tau [e^{\phi_{\mathcal{Q}}(\tau)} - e^{\psi_\mu(\tau)}] d\tau dr' < 0,$$

which is a contradiction. Now suppose there exists r_0 such that $\phi_{\mathcal{Q}}(r_0) = \psi_\mu(r_0)$ and $\phi_{\mathcal{Q}_\mu}(r) > \phi_{\mathcal{Q}}(r) > \psi_\mu(r)$ for all $0 < r < r_0$. Since

$$\phi_{\mathcal{Q}_\mu}(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau e^{\psi_\mu(\tau)} d\tau dr'$$

and $\psi_\mu = \phi_{\mathcal{Q}_\mu} - \mu \frac{r^2}{2}$, it follows that

$$\psi_\mu(r) = -8 \int_0^r \frac{1}{r'} \int_0^{r'} \tau [e^{\psi_\mu(\tau)}] d\tau dr' - \mu \frac{r^2}{2}.$$

Then, using that $\phi_{\mathcal{Q}}(r) < \phi_{\mathcal{Q}_\mu}(r)$, we get

$$\psi_\mu(r_0) < \underbrace{-8 \int_0^{r_0} \frac{1}{r'} \int_0^{r'} \tau [e^{\phi_{\mathcal{Q}}(\tau) - \mu \frac{\tau^2}{2}}] d\tau dr'}_{\tilde{\psi}_\mu(r_0)} - \mu \frac{r_0^2}{2}.$$

Notice that for $\mu = 0$ we have $\tilde{\psi}_0(r_0) = \phi_{\mathcal{Q}}(r_0)$. If we prove that $\partial_\mu \tilde{\psi}_\mu(r_0) < 0$, we deduce that $\psi_\mu(r_0) < \tilde{\psi}_0(r_0) = \phi_{\mathcal{Q}}(r_0)$, which is the desired contradiction. Since

$$\begin{aligned} \partial_r \partial_\mu \tilde{\psi}_\mu(r_0) &= \frac{4}{r_0} \int_0^{r_0} \tau^3 [e^{\phi_{\mathcal{Q}}(\tau) - \mu \frac{\tau^2}{2}}] d\tau - r_0 \\ &< \partial_r \partial_\mu \tilde{\psi}_\mu|_{\mu=0} = \frac{4}{r_0} \int_0^{r_0} \tau^3 e^{\phi_{\mathcal{Q}}} d\tau - r_0, \end{aligned}$$

then it suffices to prove that $\partial_r \partial_\mu \tilde{\psi}_\mu|_{\mu=0} < 0$. Indeed, one can easily calculate

$$\begin{aligned} \partial_r \partial_\mu \tilde{\psi}_\mu|_{\mu=0}(r_0) &= \frac{2}{r_0} \left(-1 + \frac{1}{1+r_0^2} + \log(1+r_0^2) \right) - r_0 \\ &< 0 \quad \text{for all } r_0 > 0, \end{aligned}$$

which concludes the proof.

The proof of (4.2) is similar to (4.1) and is left to the reader. Note that (4.2) implies (4.1). \square

For the rest of our estimations we will need the following proposition. Recall that T_1 was also used in [31].

PROPOSITION 4.2. *Let T_1 be the solution of $\mathcal{L}T_1 = \nabla \cdot (Q \nabla \mathcal{M}T_1) = \Lambda Q$; then we have*

$$\begin{aligned} \phi_{T_1}(y) &= O(r^4 \log(r)) \mathbb{1}_{\{r \leq 1\}} + \left[6(\log r)^2 + O\left(\frac{(\log r)^2}{r^2}\right) \right] \mathbb{1}_{\{r \geq 1\}}, \\ \phi'_{T_1}(r) &= O(r^3 \log(r)) \mathbb{1}_{\{r \leq 1\}} + O\left(\frac{\log r}{r}\right) \mathbb{1}_{\{r \geq 1\}}. \end{aligned}$$

PROOF. First we notice that $\nabla \cdot (Q \nabla \mathcal{M}T_1) = \nabla \cdot (yQ)$, with $T_1(0) = 0$ and $\nabla T_1(0) = 0$, which implies $\mathcal{M}T_1 = \frac{|y|^2}{2}$. And since T_1 is radial and $-\Delta \phi_{T_1} = T_1$, we deduce that ϕ_{T_1} is a solution of

$$L\phi_{T_1} = -\phi''_{T_1} - \frac{1}{r}\phi'_{T_1} - Q\phi_{T_1} = Q\frac{r^2}{2},$$

with $\phi_{T_1}(0) = 0$, $\phi'_{T_1}(0) = 0$. To estimate ϕ_{T_1} we need to invert the operator L . Indeed, the Green's function of L is explicit and the set of radial solutions to the homogeneous problem

$$Lf = 0$$

is spanned by

$$f_0 = 1 - \frac{2}{1+r^2}, \quad f_1 = \frac{r^2 \log(r) - 2 - \log(r)}{1+r^2},$$

with Wronskian

$$W = f'_1 f_0 - f_1 f'_0 = \frac{1}{r}.$$

Hence a solution of

$$Lf = g \quad \text{with } f(0) = f'(0) = 0$$

is given by

$$(4.9) \quad f(r) = -f_0(r) \int_0^r g f_1 \tau \, d\tau + f_1(r) \int_0^r g f_0 \tau \, d\tau.$$

It follows that

$$\phi_{T_1}(r) = -\frac{f_0(r)}{2} \int_0^r Q f_1(\tau) \tau^3 d\tau + \frac{f_1(r)}{2} \int_0^r f_0(\tau) Q \tau^3 d\tau.$$

Hence,

$$\phi_{T_1}(y) = O(r^4 \log \langle r \rangle) \mathbb{1}_{\{r \leq 1\}} + \left[6(\log r)^2 + O\left(\frac{(\log r)^2}{r^2}\right) \right] \mathbb{1}_{\{r \geq 1\}}.$$

Now to estimate ϕ'_{T_1} we use that

$$\phi'_{T_1}(r) = -\frac{f'_0(r)}{2} \int_0^r Q f_1(\tau) \tau^3 d\tau + \frac{f'_1(r)}{2} \int_0^r f_0(\tau) Q \tau^3 d\tau,$$

with

$$f'_1(r) = \frac{4r^2 \log r + r^4 + 4r^2 - 1}{8r} Q \quad \text{and} \quad f'_0(r) = \frac{1}{2} r Q.$$

We finally deduce that

$$\phi'_{T_1}(r) = O(r^3 \log \langle r \rangle) \mathbb{1}_{\{r \leq 1\}} + O\left(\frac{\log r}{r}\right) \mathbb{1}_{\{r \geq 1\}},$$

which concludes the proof. \square

4.2 Expansion of Q_μ

In the previous propositions we just had bounds on Q_μ . Now the following proposition will allow us to do an expansion of Q_μ with respect to μ where Q appears explicitly.

PROPOSITION 4.3. *There exists $\sigma := \sigma(\mu, r)$ such that*

$$(4.10) \quad \phi_{Q_\mu} = \phi_Q + \mu \phi_{T_1} + \sigma$$

where σ satisfies

$$(4.11) \quad |\sigma(\mu, r)| \lesssim \min(\mu^2 r^2 |\log \langle r \rangle|, \mu (\log \langle r \rangle)^2).$$

Moreover,

$$(4.12) \quad \phi'_{Q_\mu} = \phi'_Q + \mu \phi'_{T_1} + \partial_r \sigma,$$

with

$$(4.13) \quad |\partial_r \sigma| \lesssim \min\left(\mu^2 r |\log \langle r \rangle|, \mu \frac{|\log \langle r \rangle|}{r}\right).$$

In addition, we obtain the following expansion of Q_μ :

$$(4.14) \quad Q_\mu = Q e^{\mu \phi_{T_1} - \frac{\mu r^2}{2} + \sigma(\mu, r)}.$$

PROOF. We write (4.1) in a more precise way. We introduce $\psi_1 := \psi_1(r)$ and $\sigma := \sigma(r, \mu)$ such that

$$\phi_{Q_\mu} = \phi_Q + \mu\psi_1(r) + \sigma(\mu, r).$$

Hence, from (3.5) we have

$$-\Delta\phi_Q - \mu\Delta\psi_1 - \Delta\sigma = Qe^{\mu\psi_1 - \frac{\mu r^2}{2} + \sigma(\mu, r)}.$$

By $0 \leq \mu\psi_1 + \sigma(\mu, y) \leq \mu\frac{r^2}{2}$, we get that

$$e^{\mu\psi_1 + \sigma - \frac{\mu r^2}{2}} = 1 + \mu\psi_1 - \mu\frac{r^2}{2} + \sigma(\mu, y) + O\left(\min\left((\mu^2 r^4), \left(\mu\frac{r^2}{2}\right)\right)\right).$$

Hence,

$$\begin{aligned} -\mu\left[(\Delta + Q)\psi_1 - Q\frac{r^2}{2}\right] - (\Delta + Q)\sigma = \\ O\left(\min\left((\mu^2 r^4), \left(\mu\frac{r^2}{2}\right)\right)\right)Q := g(\mu, r). \end{aligned}$$

If we select ψ_1 such that $(\Delta + Q)\psi_1 - Q\frac{r^2}{2} = 0$, then $\psi_1 = \phi_{T_1}$ and we get that

$$-(\Delta + Q)\sigma = L\sigma = g(\mu, r).$$

From (4.9), we obtain that

$$(4.15) \quad \sigma(\mu, r) = -f_0(r) \int_0^r g f_1 \tau \, d\tau + f_1(r) \int_0^r g f_0 \tau \, d\tau.$$

Hence,

$$\sigma(\mu, r) \lesssim \min(\mu^2 r^2 |\log\langle r \rangle|, \mu(\log r)^2).$$

To conclude, (4.13) follows easily by differentiating (4.15) with respect to r . \square

4.3 Bounds on $\partial_\mu^i Q_\mu$

We derive in the following proposition bounds on $\partial_\mu Q_\mu$ and $\phi_{\partial_\mu Q_\mu}$.

PROPOSITION 4.4. *Let Q_μ be the solution of (2.5). Then*

$$(4.16) \quad \phi_{\partial_\mu Q_\mu} = \phi_{T_1} + \partial_\mu \sigma$$

where

$$(4.17) \quad |\partial_\mu \sigma| \lesssim \min(\mu r^2 |\log\langle r \rangle|, |\log\langle r \rangle|^2).$$

Moreover,

$$(4.18) \quad \phi'_{\partial_\mu Q_\mu} = \phi'_{T_1} + \partial_r \partial_\mu \sigma,$$

with

$$|\partial_r \partial_\mu \sigma| \lesssim \min\left(\mu r |\log\langle r \rangle|, \frac{|\log\langle r \rangle|}{r}\right).$$

In addition,

$$(4.19) \quad \partial_\mu Q_\mu = \left(\phi_{T_1} + \partial_\mu \sigma - \frac{r^2}{2} \right) Q_\mu.$$

PROOF. By taking the derivative in μ of (4.15), we get

$$(4.20) \quad \begin{aligned} \partial_\mu \sigma = & -f_0(r) \int_0^r \left(\phi_{T_1} + \partial_\mu \sigma - \frac{r^2}{2} \right) (Q_\mu - Q) f_1 \tau \, d\tau \\ & + f_1(r) \int_0^r \left(\phi_{T_1} + \partial_\mu \sigma - \frac{r^2}{2} \right) (Q_\mu - Q) f_0 \tau \, d\tau. \end{aligned}$$

Since $|f_0| \leq 1$, $|f_1| \leq |\log\langle r \rangle|$, and $|Q_\mu - Q| \lesssim \min(\mu r^2, 1)Q$, it follows that

$$|\partial_\mu \sigma| \lesssim \min(\mu r^2, |\log\langle r \rangle|) |\log\langle r \rangle| + |\log\langle r \rangle| \int_0^r |\partial_\mu \sigma| \min(\mu \tau^2, 1) \tau \, d\tau.$$

Hence, by using the Gronwall inequality, we deduce

$$(4.21) \quad |\partial_\mu \sigma| \lesssim \min(\mu r^2, |\log\langle r \rangle|) |\log\langle r \rangle|.$$

Note that to prove the smoothness in μ rigorously one should do the previous calculation on the finite difference $\frac{\sigma(\mu) - \sigma(\mu_0)}{\mu - \mu_0}$ and then pass to the limit $\mu \rightarrow \mu_0$ in (4.21). We will not detail this here. To prove the bound on $\partial_r \partial_\mu \sigma$, we differentiate (4.20) with respect to r ,

$$(4.22) \quad \begin{aligned} \partial_r \partial_\mu \sigma = & -f'_0(r) \int_0^r \left(\phi_{T_1} + \partial_\mu \sigma - \frac{r^2}{2} \right) (Q_\mu - Q) f_1 \tau \, d\tau \\ & + f'_1(r) \int_0^r \left(\phi_{T_1} + \partial_\mu \sigma - \frac{r^2}{2} \right) (Q_\mu - Q) f_0 \tau \, d\tau. \end{aligned}$$

If we use as before that $|f_0| \leq 1$, $f'_0(r) = \frac{rQ}{2}$, $|f_1| \lesssim |\log\langle r \rangle|$, $|f'_1| \lesssim r^3 Q$, $|Q_\mu - Q| \lesssim \min(\mu r^2, 1)Q$, and $|\partial_\mu \sigma| \lesssim \min(\mu r^2, |\log\langle r \rangle|) |\log\langle r \rangle|$, we deduce

$$|\partial_r \partial_\mu \sigma| \lesssim \min\left(\mu r |\log\langle r \rangle|, \frac{|\log\langle r \rangle|}{r}\right). \quad \square$$

By taking an extra derivative with respect to μ , we can obtain bounds on $\partial_\mu^2 Q_\mu$ and $\phi_{\partial_\mu^2 Q_\mu}$.

PROPOSITION 4.5. *The following bounds hold:*

$$(4.23) \quad |\phi_{\partial_\mu^2 Q_\mu}| \lesssim r^2 |\log\langle r \rangle|$$

and

$$(4.24) \quad \partial_\mu^2 Q_\mu = \left[\phi_{\partial_\mu^2 Q_\mu} + \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right)^2 \right] Q_\mu.$$

PROOF. We differentiate (4.16) with respect to μ ,

$$(4.25) \quad \phi_{\partial_\mu^2} Q_\mu = \partial_\mu^2 \sigma.$$

By differentiating (4.20) in μ and using $\partial_\mu Q_\mu = (\phi_{\partial_\mu} Q_\mu - \frac{r^2}{2}) Q_\mu$ we get

$$\begin{aligned} \partial_\mu^2 \sigma &= -f_0(r) \int_0^r \left[\partial_\mu^2 \sigma (Q_\mu - Q) + Q_\mu \left(\partial_\mu \sigma - \frac{r^2}{2} \right)^2 \right] f_1 \tau \, d\tau \\ &\quad + f_1(r) \int_0^r \left[\partial_\mu^2 \sigma (Q_\mu - Q) + Q_\mu \left(\partial_\mu \sigma - \frac{r^2}{2} \right)^2 \right] f_0 \tau \, d\tau. \end{aligned}$$

Finally, using that $|f_0| \leq 1$, $|f_1| \lesssim |\log \langle r \rangle|$, $|Q_\mu - Q| \lesssim \min(1, \mu r^2) Q$, $|Q_\mu| \lesssim Q$, and $|\partial_\mu \sigma| \lesssim \min(\mu r^2 |\log \langle r \rangle|, |\log \langle r \rangle|^2)$, we deduce that

$$|\partial_\mu^2 \sigma| \lesssim r^2 |\log \langle r \rangle| + |\log \langle r \rangle| \int_0^r |\partial_\mu^2 \sigma| \min(\mu \tau^2, 1) \tau Q \, d\tau.$$

Hence, (4.23) follows from Gronwall inequality. To justify the extra derivative that we took on $\partial_\mu \sigma$ with respect to μ one can do the same calculation above on the finite difference $\frac{\partial_\mu \sigma(\mu) - \partial_\mu \sigma(\mu_0)}{\mu - \mu_0}$ and then pass to the limit $\mu \rightarrow \mu_0$. To conclude, (4.24) follows from differentiating (4.19) with respect to μ . \square

PROPOSITION 4.6. *We have the following bounds:*

$$(4.26) \quad |\phi_{\partial_\mu^3} Q_\mu| \lesssim r^4 |\log \langle r \rangle|,$$

$$(4.27) \quad \begin{aligned} \partial_\mu^3 Q_\mu &= \left[\left(\phi_{\partial_\mu^2} Q_\mu + \left(\phi_{\partial_\mu} Q_\mu - \frac{r^2}{2} \right)^2 + 2\phi_{\partial_\mu^2} Q_\mu \right) \right. \\ &\quad \left. \cdot \left(\phi_{\partial_\mu} Q_\mu - \frac{r^2}{2} \right) + \phi_{\partial_\mu^3} Q_\mu \right] Q_\mu. \end{aligned}$$

PROOF. We differentiate (4.25) with respect to μ . Hence,

$$\phi_{\partial_\mu^3} Q_\mu = \partial_\mu^3 \sigma.$$

By processing as in the proof of Proposition 4.5, one can easily deduce (4.26); the details are left to the reader. To conclude, by differentiating

$$\partial_\mu^2 Q_\mu = \left[\phi_{\partial_\mu^2} Q_\mu + \left(\phi_{\partial_\mu} Q_\mu - \frac{r^2}{2} \right)^2 \right] Q_\mu$$

with respect to μ , (4.27) follows. \square

4.4 Expansions of the Mass M with Respect to μ

In the following lemmas we estimate the mass and the second moment of Q_μ , which are fundamental for either the application of the implicit function theorem in Lemma 2.2 or the derivation of the law of μ in Lemma 5.1.

LEMMA 4.7. *We have the following expansions for the mass of Q_μ :*

$$(4.28) \quad M(\mu) = \int_{\mathbb{R}^2} Q_\mu dy = 8\pi + 2\mu \log(\mu) + O(\mu),$$

$$(4.29) \quad M'(\mu) = \int_{\mathbb{R}^2} \partial_\mu Q_\mu dy = 2 \log(\mu) + O(1),$$

$$(4.30) \quad M''(\mu) = \int_{\mathbb{R}^2} \partial_\mu^2 Q_\mu dy = O\left(\frac{1}{\mu}\right).$$

PROOF. To prove (4.28) we decompose $Q_\mu = Q + Q(e^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} - 1)$ and plug it into $M = \int_{\mathbb{R}^2} Q_\mu dy$. Hence,

$$\begin{aligned} M(\mu) &= \int_{\mathbb{R}^2} Q dy + \int_{\mathbb{R}^2} Q(e^{\phi_{Q_\mu} - \mu\frac{|y|^2}{2}} - 1) dy \\ &= 8\pi + \int_0^{\frac{1}{\sqrt{\mu}}} Q\left(\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2} + O(\mu^2 r^4)\right) r dr \\ &\quad + \int_{\frac{1}{\sqrt{\mu}}}^{+\infty} (Q_\mu - Q) r dr, \end{aligned}$$

Since $|Q_\mu - Q| \lesssim Q$ for $r \geq \frac{1}{\sqrt{\mu}}$ we get that

$$\int_{\frac{1}{\sqrt{\mu}}}^{+\infty} |Q_\mu - Q| r dr \lesssim \int_{\frac{1}{\sqrt{\mu}}}^{+\infty} Q r dr \lesssim \mu.$$

Hence, by using $|\sigma| \lesssim \mu^2 r^2 |\log(r)|$ for $r \leq \frac{1}{\sqrt{\mu}}$ and $|\phi_{T_1}| \lesssim |\log(r)|^2$ we deduce that

$$(4.31) \quad M(\mu) = 8\pi - \frac{\mu}{2} \int_0^{\frac{1}{\sqrt{\mu}}} Q r^3 dr + O(\mu) = 8\pi + 2\mu \log(\mu) + O(\mu).$$

To prove (4.29), we use that $\partial_\mu Q_\mu = (\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2}) Q_\mu$. It follows by using Propositions 4.3 and 4.4 that

$$\begin{aligned} \int_{\mathbb{R}^2} \partial_\mu Q_\mu dy &= \int_0^{\frac{1}{\sqrt{\mu}}} \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right) \\ (4.32) \quad &\quad \cdot Q \left(1 + \mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2} + O(\mu^2 r^4) \right) r dr \\ &\quad + \int_{\frac{1}{\sqrt{\mu}}}^{\infty} \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right) Q_\mu r dr. \end{aligned}$$

Since $|\sigma| \lesssim \min(\mu^2 r^2 |\log\langle r \rangle|, \mu |\log\langle r \rangle|^2)$, $|\phi_{T_1}| \lesssim \min(r^4 |\log\langle r \rangle|, |\log\langle r \rangle|^2)$, and $|\phi_{\partial_\mu \mathcal{Q}_\mu}| \lesssim r^2$, it follows that

$$(4.33) \quad \begin{aligned} \int_{\frac{1}{\sqrt{\mu}}}^{\infty} \left| \phi_{\partial_\mu \mathcal{Q}_\mu} - \frac{r^2}{2} \right| \mathcal{Q}_\mu r \, dr &\lesssim \int_{\frac{1}{\sqrt{\mu}}}^{\infty} r^3 \mathcal{Q}_\mu \, dr \\ &\lesssim \int_{\frac{1}{\sqrt{\mu}}}^{\infty} r^3 \mathcal{Q} e^{\mu(|\log\langle r \rangle|^2 - \frac{r^2}{2})} \, dr \lesssim 1, \end{aligned}$$

and

$$\begin{aligned} \int_0^{\frac{1}{\sqrt{\mu}}} \left| \phi_{\partial_\mu \mathcal{Q}_\mu} - \frac{r^2}{2} \right| \mathcal{Q} \left| \mu \phi_{T_1} + \sigma - \mu \frac{r^2}{2} + O(\mu^2 r^4) \right| r \, dr \\ + \int_0^{\frac{1}{\sqrt{\mu}}} |\phi_{\partial_\mu \mathcal{Q}_\mu}| \mathcal{Q} r \, dr \lesssim 1. \end{aligned}$$

Hence,

$$(4.34) \quad \begin{aligned} M'(\mu) &= \int_{\mathbb{R}^2} \partial_\mu \mathcal{Q}_\mu \, dy = - \int_0^{\frac{1}{\sqrt{\mu}}} \mathcal{Q} \frac{r^3}{2} \, dr + O(1) \\ &= 2 \log(\mu) + O(1). \end{aligned}$$

To prove (4.30) we use Proposition 4.5 to deduce that

$$(4.35) \quad M''(\mu) = \int_{\mathbb{R}^2} \partial_\mu^2 \mathcal{Q}_\mu \, dy = \int_{\mathbb{R}^2} \left[\phi_{\partial_\mu^2 \mathcal{Q}_\mu} + \left(\phi_{\partial_\mu \mathcal{Q}_\mu} - \frac{r^2}{2} \right)^2 \right] \mathcal{Q}_\mu \, dy.$$

From Proposition 4.6 we have for all $r \geq 0$:

$$r^4 - 1 \lesssim \phi_{\partial_\mu^2 \mathcal{Q}_\mu} + \left(\phi_{\partial_\mu \mathcal{Q}_\mu} - \frac{r^2}{2} \right)^2 \lesssim r^4 + 1.$$

Hence,

$$(4.36) \quad \int_0^\infty r^5 \mathcal{Q}_\mu \, dr - 1 \lesssim M''(\mu) \lesssim \int_0^\infty r^5 \mathcal{Q}_\mu(r) \, dr + 1.$$

Since $\mathcal{Q}_\mu = \mathcal{Q} e^{\mu \phi_{T_1} + \sigma - \mu \frac{r^2}{2}}$ we obtain that

$$\begin{aligned} \int_0^\infty r^5 \mathcal{Q}_\mu \, dr &= \int_0^\infty r^5 \mathcal{Q} e^{\mu \phi_{T_1} + \sigma - \mu \frac{r^2}{2}} \, dr \\ &= \frac{1}{\mu} \int_0^\infty \tau^5 \frac{1}{(\mu + \tau^2)^2} e^{\mu \phi_{T_1} \left(\frac{\tau}{\sqrt{\mu}} \right) + \sigma \left(\frac{\tau}{\sqrt{\mu}} \right) - \frac{\tau^2}{2}} \, d\tau, \end{aligned}$$

and since

$$\int_0^\infty \tau^5 \frac{1}{(\mu + \tau^2)^2} e^{\mu \phi_{T_1} \left(\frac{\tau}{\sqrt{\mu}} \right) + \sigma \left(\frac{\tau}{\sqrt{\mu}} \right) - \frac{\tau^2}{2}} \, d\tau$$

is uniformly bounded with respect to μ , we deduce that (4.30) holds, which concludes the proof of Lemma 4.7. \square

4.5 Expansions of the Second Moment with Respect to μ

LEMMA 4.8. *We have the following expression on the second moment of Q_μ and the following bounds on the derivatives of the second moment of Q_μ with respect to μ :*

$$(4.37) \quad \int_{\mathbb{R}^2} Q_\mu(y) |y|^2 dy = \frac{2M(\mu)}{\mu} \left(1 - \frac{M(\mu)}{8\pi} \right).$$

There exists $C > 0$ such that

$$(4.38) \quad \frac{-C}{\mu} \leq \int_{\mathbb{R}^2} \partial_\mu Q_\mu(y) |y|^2 dy \leq \frac{-1}{C\mu},$$

$$(4.39) \quad \frac{1}{C\mu^2} \leq \int_{\mathbb{R}^2} \partial_\mu^2 Q_\mu(y) |y|^2 dy \leq \frac{C}{\mu^2},$$

$$(4.40) \quad \frac{-C}{\mu^3} \leq \int_{\mathbb{R}^2} \partial_\mu^3 Q_\mu(y) |y|^2 dy \leq \frac{-1}{C\mu^3}.$$

PROOF OF (4.37). We start by calculating the second moment of Q_μ . We multiply (2.5) by $|y|^2$ and integrate:

$$\begin{aligned} \int_{\mathbb{R}^2} \Delta Q_\mu |y|^2 dy - \int_{\mathbb{R}^2} \nabla \cdot (Q_\mu \nabla \phi_{Q_\mu}) |y|^2 dy \\ = -\mu \int_{\mathbb{R}^2} \nabla \cdot (y Q_\mu) |y|^2 dy. \end{aligned}$$

Then by integration by parts and using that $\nabla \phi_{Q_\mu} = \frac{1}{|\cdot|} \star Q_\mu$, we get

$$\begin{aligned} 4M + \frac{1}{2\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} \frac{2y \cdot (y-x)}{|x-y|^2} Q_\mu(x) Q_\mu(y) dx dy \\ = 2\mu \int_{\mathbb{R}^2} Q_\mu(y) |y|^2 dy, \\ 4M + \frac{1}{2\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} \frac{(x-y) \cdot (y-x)}{|x-y|^2} Q_\mu(x) Q_\mu(y) dx dy \\ = 2\mu \int_{\mathbb{R}^2} Q_\mu(y) |y|^2 dy \\ 4M - \frac{M^2}{2\pi} = 2\mu \int_{\mathbb{R}^2} Q_\mu(y) |y|^2 dy, \end{aligned}$$

which implies (4.37). □

We now have the following:

PROOF OF (4.38).

$$(4.41) \quad \int_{\mathbb{R}^2} \partial_\mu Q_\mu(y) |y|^2 dy = \int_{\mathbb{R}^2} \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right) Q_\mu(y) |y|^2 dy.$$

Since $Q_\mu = Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}}$ and $\phi_{\partial_\mu Q_\mu} = \phi_{T_1} + \partial_\mu\sigma$, it follows that

$$(4.42) \quad \int_{\mathbb{R}^2} \partial_\mu Q_\mu(y) |y|^2 dy = \int_0^\infty \left(\phi_{T_1} + \partial_\mu\sigma - \frac{r^2}{2} \right) Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} r^3 dr.$$

By using that

$$\begin{aligned} |\phi_{T_1}| &\lesssim \min(r^4 |\log\langle r \rangle|, |\log\langle r \rangle|^2), \\ |\partial_\mu\sigma| &\lesssim \min(\mu r^2 |\log\langle r \rangle|, |\log\langle r \rangle|^2), \end{aligned}$$

we deduce that for all $r \geq 0$

$$-r^2 \lesssim \phi_{T_1} + \partial_\mu\sigma - \frac{r^2}{2} \lesssim -r^2.$$

Hence,

$$\begin{aligned} - \int_0^\infty r^5 Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} dr &\lesssim \int_{\mathbb{R}^2} \partial_\mu Q_\mu(y) |y|^2 dy \\ &\lesssim - \int_0^\infty r^5 Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} dr. \end{aligned}$$

Notice that

$$\int_0^\infty r^5 Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} dr = \frac{1}{\mu} \int_0^\infty \tau^5 \frac{1}{(\mu + \tau^2)^2} e^{\mu\phi_{T_1}(\frac{\tau}{\sqrt{\mu}}) + \sigma(\frac{\tau}{\sqrt{\mu}}) - \frac{\tau^2}{2}} d\tau,$$

and since $\int_0^\infty \tau^5 \frac{1}{(\mu + \tau^2)^2} e^{\mu\phi_{T_1}(\frac{\tau}{\sqrt{\mu}}) + \sigma(\frac{\tau}{\sqrt{\mu}}) - \frac{\tau^2}{2}} d\tau$ is uniformly bounded with respect to μ , we deduce (4.38). \square

PROOF OF (4.39). Notice first that

$$\partial_\mu^2 Q_\mu = \left[\phi_{\partial_\mu^2 Q_\mu} + \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right)^2 \right] Q_\mu.$$

Hence, by using

$$|\phi_{T_1}| \lesssim \min(r^4 |\log\langle r \rangle|, |\log\langle r \rangle|^2), |\partial_\mu\sigma| \lesssim \min(\mu r^2 |\log\langle r \rangle|, |\log\langle r \rangle|^2)$$

and $|\phi_{\partial_\mu^2 Q_\mu}| \lesssim r^2 |\log\langle r \rangle|$, we deduce, for all $r \geq 0$,

$$r^4 - 1 \lesssim \phi_{\partial_\mu^2 Q_\mu} + \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right)^2 \lesssim r^4 + 1.$$

$$\begin{aligned} \int_0^\infty (r^7 - r^3) Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} dr &\lesssim \int_{\mathbb{R}^2} \partial_\mu Q_\mu(y) |y|^2 dy \\ &\lesssim \int_0^\infty (r^7 + r^3) Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} dr. \end{aligned}$$

From (4.31) and (4.33) we get

$$\int_0^\infty r^3 Qe^{\mu\phi_{T_1} + \sigma - \mu\frac{r^2}{2}} dr = O(|\log\mu|).$$

Notice also that

$$\int_0^\infty r^7 Q e^{\mu\phi_{T_1} + \sigma - \mu \frac{r^2}{2}} dr = \frac{1}{\mu^2} \int_0^\infty \tau^7 \frac{1}{(\mu + \tau^2)^2} e^{\mu\phi_{T_1}(\frac{\tau}{\sqrt{\mu}}) + \sigma(\frac{\tau}{\sqrt{\mu}}) - \frac{\tau^2}{2}} d\tau,$$

and since

$$\int_0^\infty \tau^7 \frac{1}{(\mu + \tau^2)^2} e^{\mu\phi_{T_1}(\frac{\tau}{\sqrt{\mu}}) + \sigma(\frac{\tau}{\sqrt{\mu}}) - \frac{\tau^2}{2}} d\tau$$

is uniformly bounded with respect to μ , we deduce (4.39). The last step is to prove (4.40). \square

PROOF OF (4.40). If we use that

$$\partial_\mu^3 Q_\mu = \left[\left(\phi_{\partial_\mu^2 Q_\mu} + \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right)^2 + 2\phi_{\partial_\mu^2 Q_\mu} \right) \left(\phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} \right) + \phi_{\partial_\mu^3 Q_\mu} \right] Q_\mu,$$

then by proceeding as in the proof of the previous inequalities, one can easily deduce (4.40); the rest of the proof is left to the reader. \square

4.6 Correction of the Profile Q_μ and Estimation of the Error E

Now we are ready to decide what the good approximate profile will be. Indeed, we choose the following profile:

$$(4.43) \quad \tilde{Q}_\mu = Q_\mu - \tilde{\mu} \partial_\mu Q_\mu,$$

where

$$(4.44) \quad \tilde{\mu} = \frac{\int_{\mathbb{R}^2} Q_\mu - 8\pi}{\int_{\mathbb{R}^2} \partial_\mu Q_\mu} = \mu + O\left(\frac{\mu}{|\log \mu|}\right),$$

which can be deduced easily from Lemma 4.7. This choice can be justified by two reasons: The first one is that the mass of ε will be zero, and this implies that

$$(\mathcal{M}_\mu^y \varepsilon, \partial_M n_\infty) = (\varepsilon, 1) = 0.$$

The second reason is that if we plug \tilde{Q}_μ into

$$(4.45) \quad \begin{aligned} \Delta \tilde{Q}_\mu + \mu \Lambda \tilde{Q}_\mu - \nabla \cdot (\tilde{Q}_\mu \nabla \phi_{\tilde{Q}_\mu}) &= -\tilde{\mu} \mathcal{L}_\mu^y (\partial_\mu Q_\mu) - \tilde{\mu}^2 \nabla \cdot (\partial_\mu Q_\mu \nabla \phi_{\partial_\mu Q_\mu}) \\ &= \tilde{\mu} \Lambda Q_\mu - \tilde{\mu}^2 \nabla \cdot (\partial_\mu Q_\mu \nabla \phi_{\partial_\mu Q_\mu}), \end{aligned}$$

the identity $-\mathcal{L}_\mu^y (\partial_\mu Q_\mu) = \Lambda Q_\mu$ yields that the error

$$E = \tilde{\mu} \Lambda Q_\mu - \tilde{\mu}^2 \nabla \cdot (\partial_\mu Q_\mu \nabla \phi_{\partial_\mu Q_\mu})$$

is of order μ^2 when it is projected onto $\mathcal{M}_\mu^y \varepsilon$. Indeed, since $(\varepsilon, |\cdot|^2) = (\varepsilon, 1) = 0$, it follows that

$$(\mathcal{M}_\mu^y (\Lambda Q_\mu), \varepsilon) = \left(2 - \frac{M}{2\pi} - \mu |\cdot|^2, \varepsilon \right) = 0,$$

which implies

$$(E, \mathcal{M}_\mu^y \varepsilon)_{L^2} = -\tilde{\mu}^2 (\nabla \cdot (\partial_\mu Q_\mu \nabla \phi_{\partial_\mu Q_\mu}), \mathcal{M}_\mu^y \varepsilon)_{L^2}.$$

This cancellation makes the projection of our error on ε with the inner product $(\cdot, \cdot)_{\mathcal{M}_\mu^\gamma}$ of order μ^2 , which is fundamental for closing the energy estimate.

5 Derivation of the Law of $\mu(s(t))$

Recall that

$$v(y, s) = \tilde{Q}_\mu(y) + \varepsilon(y).$$

The conservation of the second moment in the x -coordinate translates into

$$(5.1) \quad \int_{\mathbb{R}^2} v|y|^2 dy = \frac{I}{\mu R(t)^2},$$

where I is the second moment of the initial data u_0 . We will abuse notation and use $\mu(t)$ for $\mu(s(t))$.

LEMMA 5.1. *Let v be the solution of (2.4) with 8π mass and ε satisfy $(\varepsilon, |\cdot|^2)_{L^2} = 0$. Then,*

$$(5.2) \quad \frac{I}{\mu R(t)^2} = -\frac{M}{2\pi} \log(\mu) + O(1)$$

and

$$(5.3) \quad \mu(t) = \frac{2\pi I}{M(2t+1) \log(2t+1) + O(t \log \log(2t+1))}.$$

If we consider μ as a function of s , we have

$$(5.4) \quad \left| \mu(s) - \frac{1}{2s} \right| \leq \frac{C'}{s \log(s)}.$$

PROOF. If we combine (4.28) and (4.37), then

$$(5.5) \quad \begin{aligned} \int_{\mathbb{R}^2} Q_\mu(y) |y|^2 dy &= \frac{2M}{\mu} \left(\frac{-2\mu \log(\mu) + O(\mu)}{8\pi} \right) \\ &= -\frac{M \log(\mu)}{2\pi} + O(1). \end{aligned}$$

It follows from (5.1) that $(\varepsilon, |\cdot|^2) = 0$ and (5.5) that

$$(5.6) \quad \begin{aligned} \frac{I}{R^2 \mu} &= \int_{\mathbb{R}^2} \tilde{Q}_\mu(y) |y|^2 dy \\ &= -\frac{M}{2\pi} \log(\mu) + O(1) - \tilde{\mu} \int_{\mathbb{R}^2} \partial_\mu Q_\mu |y|^2 dy. \end{aligned}$$

From (4.38) we deduce that

$$(5.7) \quad \int_{\mathbb{R}^2} \partial_\mu Q_\mu(y) |y|^2 dy = O\left(\frac{1}{\mu}\right).$$

Hence, by using $\tilde{\mu} = \mu + O\left(\frac{\mu}{|\log \mu|}\right)$ it follows that

$$(5.8) \quad \frac{I}{R^2 \mu} = \int_{\mathbb{R}^2} \tilde{Q}_\mu(y) |y|^2 dy = -\frac{M}{2\pi} \log(\mu) + O(1),$$

and taking the log of (5.8), we get that

$$\log\left(\frac{1}{\mu}\right) - \log(2t + 1) = \log\log\left(\frac{1}{\mu}\right) + O(1).$$

Consequently,

$$\log\log\left(\frac{1}{\mu}\right) = \log\log(2t + 1) + O(1),$$

which implies (5.3). To find $\mu(s)$ we use that

$$(5.9) \quad \frac{ds}{dt} = \frac{1}{\mu R^2} = \frac{M \log(2t + 1) + O(\log\log(2t + 1))}{2\pi I}.$$

Hence, if we integrate we find

$$s(t) = \frac{2(2t + 1)M \log(2t + 1) + O(t \log\log(2t + 1))}{2\pi I},$$

and (5.4) follows with C' in (5.4) a uniform constant with respect to A_1 , and hence if we choose A_1 sufficiently large, we get $C' < \frac{A_1}{100}$. \square

Now we need to find a bound on μ_s for the energy estimates.

LEMMA 5.2.

$$(5.10) \quad \mu_s = -2\mu^2 + O\left(\frac{\mu^2}{|\log \mu|}\right).$$

PROOF. To prove (5.10), we differentiate $\frac{I}{R^2\mu} = \int_{\mathbb{R}^2} \tilde{Q}_\mu(y)|y|^2 dy$ with respect to s . Hence,

$$(5.11) \quad I \frac{d}{ds} \left(\frac{1}{R^2\mu} \right) = \mu_s \int_{\mathbb{R}^2} \partial_\mu \tilde{Q}_\mu(y)|y|^2 dy.$$

To derive (5.10) we must estimate $\int_{\mathbb{R}^2} \partial_\mu \tilde{Q}_\mu(y)|y|^2 dy$. Indeed, since $\partial_\mu \tilde{Q}_\mu(y) = \partial_\mu Q_\mu - \partial_\mu \tilde{\mu} \partial_\mu Q_\mu - \tilde{\mu} \partial_\mu^2 Q_\mu$, we first have to compute $\partial_\mu \tilde{\mu}$, and from Lemma 4.7 we obtain

$$(5.12) \quad \partial_\mu \tilde{\mu} = 1 + O\left(\frac{1}{|\log \mu|}\right).$$

It follows that

$$(5.13) \quad \partial_\mu \tilde{Q}_\mu = -\mu \partial_\mu^2 Q_\mu + O\left(\frac{\mu \partial_\mu^2 Q_\mu + \partial_\mu Q_\mu}{|\log \mu|}\right).$$

From Lemma 4.7 we compute

$$(5.14) \quad \partial_\mu \tilde{\mu} = 1 + O\left(\frac{1}{|\log \mu|}\right),$$

which implies

$$(5.15) \quad \partial_\mu \tilde{Q}_\mu = -\mu \partial_\mu^2 Q_\mu + O\left(\frac{\mu \partial_\mu^2 Q_\mu + \partial_\mu Q_\mu}{|\log \mu|}\right).$$

From Lemma 4.8 we deduce that

$$\begin{aligned}
 & \int_{\mathbb{R}^2} \partial_\mu \tilde{Q}_\mu |y|^2 dy \\
 &= -\mu \int_{\mathbb{R}^2} \partial_\mu^2 Q_\mu |y|^2 dy + O\left(\int_{\mathbb{R}^2} \frac{\mu \partial_\mu^2 Q_\mu + \partial_\mu Q_\mu}{|\log \mu|} |y|^2 dy\right) \\
 &= -\mu \int_{\mathbb{R}^2} \partial_\mu^2 Q_\mu |y|^2 dy + O\left(\frac{1}{\mu |\log \mu|}\right) \\
 (5.16) \quad &= -\frac{C}{\mu} + O\left(\frac{1}{\mu |\log \mu|}\right),
 \end{aligned}$$

and similarly we obtain from Lemma 4.8,

$$(5.17) \quad \int_{\mathbb{R}^2} \partial_\mu^2 \tilde{Q}_\mu |y|^2 dy = \frac{C}{\mu^2} + O\left(\frac{1}{\mu^2 |\log \mu|}\right),$$

where $C > 0$ is a constant depending on the constant C in (4.39) and (4.40). Finally, since $\frac{dt}{ds} = \mu R^2$ and $R'(t)R(t) = 1$ it follows that

$$\begin{aligned}
 -I\left(\frac{\frac{d}{ds}(\mu R(t)^2)}{\mu^2 R^4}\right) &= -I\left(\frac{\mu_s}{\mu^2 R^2} + 2\frac{R'(t)R}{R^2}\right) = -I\left(\frac{\mu_s}{\mu^2 R^2} + \frac{2}{R^2}\right) \\
 (5.18) \quad &= \mu_s \int_{\mathbb{R}^2} \partial_\mu \tilde{Q}_\mu(y) |y|^2 dy = \mu_s O\left(\frac{1}{\mu}\right),
 \end{aligned}$$

which concludes the proof. \square

6 Bounds on the Potential $\nabla \phi_\varepsilon$

With the orthogonality conditions $(\varepsilon, 1) = (\varepsilon, |\cdot|^2) = 0$, $(\mathcal{M}_\mu^y \varepsilon, \varepsilon) = \|\varepsilon\|_{L_{Q_\mu}^2}^2 - \|\nabla \phi_\varepsilon\|_{L^2}^2$ does not control $\|\varepsilon\|_{L_{Q_\mu}^2}$ uniformly in μ . It turns out that we have a remarkable nonlinear structure that yields a control of $\|\nabla \phi_\varepsilon\|_{L^2}^2$ and hence $\|\varepsilon\|_{L_{Q_\mu}^2}$.

PROPOSITION 6.1.

$$(6.1) \quad \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy = O(\mu).$$

PROOF. We take the L^2 inner product of (3.2) with $|\cdot|^2$:

$$\begin{aligned}
 \frac{d}{ds}(\varepsilon, |\cdot|^2) &= (\mathcal{L}_\mu \varepsilon, |\cdot|^2) - (\nabla \cdot (\varepsilon \nabla \phi_\varepsilon), |\cdot|^2) + \frac{\mu_s}{2\mu} (\Lambda \varepsilon, |\cdot|^2) \\
 &\quad + (\Theta_\mu(\varepsilon), |\cdot|^2) + (F, |\cdot|^2) + (E, |\cdot|^2).
 \end{aligned}$$

Since

$$(\varepsilon, |\cdot|^2) = (\varepsilon, 1) = 0,$$

we have

$$\begin{aligned}
 (\mathcal{L}_\mu \varepsilon, |\cdot|^2) &= -2(\tilde{Q}_\mu \nabla \mathcal{M}_\mu^y \varepsilon, y) = 2(\mathcal{M}_\mu^y \varepsilon, \Lambda \tilde{Q}_\mu) \\
 (6.2) \quad &= 2(\varepsilon, \mathcal{M}_\mu^y \Lambda \tilde{Q}_\mu) = \left(2 - \frac{M}{2\pi}\right)(\varepsilon, 1) - \mu(\varepsilon, |y|^2) = 0
 \end{aligned}$$

and

$$(\Lambda \varepsilon, |\cdot|^2) = -2(\varepsilon, |\cdot|^2) = 0.$$

In addition,

$$(6.3) \quad -(\nabla \cdot (\varepsilon \nabla \phi_\varepsilon), |\cdot|^2) = 2 \int_{\mathbb{R}^2} \varepsilon y \cdot \nabla \phi_\varepsilon \, dy = \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 \, dy.$$

The last equality comes from the fact that $\int_{\mathbb{R}^2} \varepsilon \, dy = 0$, $-\Delta \phi_\varepsilon = \varepsilon$. Since $F = \frac{\mu_s}{2\mu} \Lambda \tilde{Q}_\mu - \mu_s \partial_\mu \tilde{Q}_\mu$, to estimate $(F, |\cdot|^2)$ we first calculate $(\Lambda \tilde{Q}_\mu, |\cdot|^2)$ and $(\partial_\mu \tilde{Q}_\mu, |\cdot|^2)$. Indeed, from (5.6) we get

$$(6.4) \quad (\Lambda \tilde{Q}_\mu, |\cdot|^2) = -2(\tilde{Q}_\mu, |\cdot|^2) = -\frac{2I}{R^2 \mu},$$

and by using (5.16) we deduce

$$(6.5) \quad \mu_s (\partial_\mu \tilde{Q}_\mu, |\cdot|^2) = -I \left(\frac{\mu_s}{\mu^2 R^2} + \frac{2}{R^2} \right).$$

Hence

$$(F, |\cdot|^2) = \frac{2I}{R^2} = -\frac{M\mu \log \mu}{\pi} + O(\mu).$$

Furthermore, since $E = \tilde{\mu} \Lambda Q_\mu - \tilde{\mu}^2 \nabla \cdot (\partial_\mu Q_\mu \nabla \phi_{\partial_\mu Q_\mu})$ and $\tilde{\mu} = \mu + O(\frac{\mu}{|\log \mu|})$, it follows that

$$\begin{aligned}
 (E, |\cdot|^2) &= -2\tilde{\mu}(Q_\mu, |\cdot|^2) + 2\tilde{\mu}^2(\nabla \phi_{\partial_\mu Q_\mu} \cdot y, \partial_\mu Q_\mu) \\
 &= -2\mu(Q_\mu, |\cdot|^2) + 2\tilde{\mu}^2(\nabla \phi_{\partial_\mu Q_\mu} \cdot y, \partial_\mu Q_\mu) + O(\mu) \\
 (6.6) \quad &= \frac{M\mu \log \mu}{\pi} + 2\tilde{\mu}^2(\nabla \phi_{\partial_\mu Q_\mu} \cdot y, \partial_\mu Q_\mu) + O(\mu).
 \end{aligned}$$

From Proposition 4.4 we get that

$$2\tilde{\mu}^2(\nabla \phi_{\partial_\mu Q_\mu} \cdot y, \partial_\mu Q_\mu) \lesssim \mu^2 |\log \mu|^2.$$

Hence,

$$(6.7) \quad (F + E, |\cdot|^2) = \underbrace{-\frac{M\mu \log \mu}{\pi} + \frac{M\mu \log \mu}{\pi}}_{\text{cancellation}} + O(\mu).$$

To estimate $(\Theta_\mu(\varepsilon), |\cdot|^2)$ we use that $\tilde{Q}_\mu - Q_\mu = -\tilde{\mu} \partial_\mu Q_\mu$. Hence,

$$(6.8) \quad (\Theta_\mu(\varepsilon), |\cdot|^2) = -\tilde{\mu}[(\partial_\mu Q_\mu, \nabla \phi_\varepsilon \cdot y) + (\varepsilon, \nabla \phi_{\partial_\mu Q_\mu} \cdot y)].$$

It follows that

$$|(\Theta_\mu(\varepsilon), |\cdot|^2)| \lesssim \mu \left[\|\nabla \phi_\varepsilon\|_{L^2} \left(\int_{\mathbb{R}^2} |\partial_\mu Q_\mu|^2 |y|^2 dy \right)^{\frac{1}{2}} + \|\varepsilon\|_{L^2_{Q_\mu}} \left(\int_{\mathbb{R}^2} \frac{|\nabla \phi_{\partial_\mu Q_\mu}|^2 |y|^2}{Q_\mu} dy \right)^{\frac{1}{2}} \right].$$

Thanks to $|\nabla \phi_{\partial_\mu Q_\mu}| \lesssim \min(\mu|y| |\log|y||, \frac{|\log|y||}{|y|})$, $\partial_\mu Q_\mu \lesssim Q_\mu |y|^2$, and the bootstrap assumption, we deduce

$$(6.9) \quad |(\Theta_\mu(\varepsilon), |\cdot|^2)| \lesssim \sqrt{A} \mu^{\frac{3}{2}} \sqrt{|\log \mu|}.$$

Hence, if we assume that μ was chosen small enough, we deduce that

$$(6.10) \quad \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy \leq C\mu,$$

with C uniform in μ , thus concluding the proof. \square

7 Energy Estimates

We want to prove the following energy bound:

PROPOSITION 7.1.

$$(7.1) \quad \frac{1}{2} \frac{d}{ds} (\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \mu \left(K_2 - \frac{1}{2} - \delta \right) (\mathcal{M}_\mu^y \varepsilon, \varepsilon) \leq C(\delta, A) \mu^2,$$

with $C(A, \delta) \lesssim \frac{1}{\delta} (\sqrt{A}(1 + \delta) + 1)$, where $\delta > 0$ is a sufficiently small constant and A is the bootstrap constant.

PROOF. We multiply (3.2) by $\mathcal{M}_\mu^y \varepsilon$ and integrate

$$(7.2) \quad \begin{aligned} & \frac{1}{2} \frac{d}{ds} (\mathcal{M}_\mu^y \varepsilon, \varepsilon) \\ &= - \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy + \underbrace{\frac{\mu_s}{2\mu} (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) - \frac{\mu_s}{2} \left(\frac{\partial_\mu Q_\mu \varepsilon}{Q_\mu^2}, \varepsilon \right)}_I \\ & \quad + \underbrace{(E, \mathcal{M}_\mu^y \varepsilon)}_{II} + \underbrace{(\Theta_\mu, \mathcal{M}_\mu^y \varepsilon)}_{III} + \underbrace{(N(\varepsilon), \mathcal{M}_\mu^y \varepsilon)}_{IV} + \underbrace{(F, \mathcal{M}_\mu^y \varepsilon)}_V. \end{aligned}$$

Now we estimate each term of (7.2).

7.1 Estimation of I (Linear Terms)

Estimation of $(\partial_\mu Q_\mu \varepsilon / Q_\mu^2, \varepsilon)$

We will use the notation

$$\langle y \rangle = \sqrt{1 + |y|^2}.$$

Notice first that by Proposition 4.3

$$\frac{\partial_\mu Q_\mu}{Q_\mu} = \phi_{\partial_\mu Q_\mu} - \frac{r^2}{2} = -\frac{r^2}{2} + \min(O((\log(r))^2), O(\mu r^2 |\log(r)|)).$$

This term will be very helpful since it has the good sign if we can control

$$\int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \frac{\varepsilon^2}{Q_\mu} dy.$$

Indeed, we can prove that

$$\int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \frac{\varepsilon^2}{Q_\mu} dy \lesssim A\sqrt{\mu} + \frac{\sqrt{A}}{\sqrt{\mu}} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}}.$$

We first use that $\frac{\varepsilon}{Q_\mu} = \mathcal{M}_\mu^y \varepsilon + \phi_\varepsilon$, which implies

$$(7.3) \quad \begin{aligned} \int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \frac{\varepsilon^2}{Q_\mu} dy &= \int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \varepsilon \mathcal{M}_\mu^y \varepsilon dy \\ &\quad + \int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \varepsilon \phi_\varepsilon dy. \end{aligned}$$

Then by Proposition B.2 and the bootstrap assumption

$$\int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \varepsilon \phi_\varepsilon dy \leq \|\phi_\varepsilon\|_{L^\infty} \|\varepsilon\|_{L^2_{Q_\mu}} \|\log(\langle y \rangle)|^2 \sqrt{Q_\mu}\|_{L^2} \lesssim A\mu.$$

To bound the other term we use a bootstrap argument, Lemma 3.2, and Lemma A.2:

$$\begin{aligned} &\int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \varepsilon \mathcal{M}_\mu^y \varepsilon dy \\ &\leq \|\varepsilon\|_{L^2_{Q_\mu}} \left(\int_{\mathbb{R}^2} |\log(\langle y \rangle)|^4 Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\ &\lesssim \|\varepsilon\|_{L^2_{Q_\mu}} \left(\int_{\mathbb{R}^2} |y|^2 Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\ &\lesssim \|\varepsilon\|_{L^2_{Q_\mu}} \left(\frac{4}{\mu^2} \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy + \frac{12}{\mu} \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\ &\lesssim \frac{\sqrt{A}}{\sqrt{\mu}} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} + A\sqrt{\mu}, \end{aligned}$$

which yields

$$\int_{\mathbb{R}^2} |\log(\langle y \rangle)|^2 \frac{\varepsilon^2}{Q_\mu} dy \lesssim A\sqrt{\mu} + \frac{\sqrt{A}}{\sqrt{\mu}} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}}.$$

Hence, if we combine the previous inequality with $\mu_s = -2\mu^2 + O(\frac{\mu^2}{|\log \mu|})$, we get

$$(7.4) \quad -\frac{\mu_s}{2} \left(\frac{\partial_\mu Q_\mu \varepsilon}{Q_\mu^2}, \varepsilon \right) = -\frac{\mu^2}{2} \int_{\mathbb{R}^2} |y|^2 \frac{\varepsilon^2}{Q_\mu} dy + O(\mu^{\frac{5}{2}}).$$

Estimation of $(\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon)$

$$(7.5) \quad \begin{aligned} (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) &= 2(\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \int_{\mathbb{R}^2} y \cdot \nabla \varepsilon \mathcal{M}_\mu^y \varepsilon dy \\ &= 2(\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \int_{\mathbb{R}^2} y \cdot \nabla \varepsilon \frac{\varepsilon}{Q_\mu} dy - \int_{\mathbb{R}^2} y \cdot \nabla \varepsilon \phi_\varepsilon dy. \end{aligned}$$

We calculate each term separately:

$$(7.6) \quad \begin{aligned} \int_{\mathbb{R}^2} y \cdot \nabla \varepsilon \frac{\varepsilon}{Q_\mu} dy &= \frac{1}{2} \int_{\mathbb{R}^2} \frac{y \cdot \nabla (\varepsilon^2)}{Q_\mu} dy \\ &= - \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy + \frac{1}{2} \int_{\mathbb{R}^2} \varepsilon^2 \frac{y \cdot \nabla Q_\mu}{Q_\mu^2} dy. \end{aligned}$$

In addition,

$$- \int_{\mathbb{R}^2} y \cdot \nabla \varepsilon \phi_\varepsilon dy = 2 \int_{\mathbb{R}^2} \varepsilon \phi_\varepsilon dy + \int_{\mathbb{R}^2} \varepsilon y \cdot \nabla \phi_\varepsilon dy.$$

Since for $\int_{\mathbb{R}^2} \varepsilon dy = 0$, we have that

$$\int_{\mathbb{R}^2} \varepsilon \phi_\varepsilon dy = \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy,$$

by using $-\Delta \phi_\varepsilon = \varepsilon$ and integration by parts we deduce that

$$- \int_{\mathbb{R}^2} y \cdot \nabla \varepsilon \phi_\varepsilon dy = 2 \int_{\mathbb{R}^2} \varepsilon \phi_\varepsilon dy + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy = \frac{5}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy.$$

Hence,

$$(7.7) \quad \begin{aligned} (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) &= 2(\mathcal{M}_\mu^y \varepsilon, \varepsilon) - \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy + \frac{1}{2} \int_{\mathbb{R}^2} \varepsilon^2 \frac{y \cdot \nabla Q_\mu}{Q_\mu^2} dy \\ &\quad + \frac{5}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy \\ &= \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy + \frac{1}{2} \int_{\mathbb{R}^2} \varepsilon^2 \frac{y \cdot \nabla Q_\mu}{Q_\mu^2} dy + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy. \end{aligned}$$

Since $\frac{y \cdot \nabla Q_\mu}{Q_\mu} = y \cdot \nabla \phi_{Q_\mu} - \mu|y|^2$ and $y \cdot \nabla \phi_{Q_\mu} > y \cdot \nabla \phi_Q > -4$, we get that

$$\begin{aligned}
 (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) &= \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy - \frac{\mu}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy \\
 &\quad + \frac{1}{2} \int_{\mathbb{R}^2} \frac{\varepsilon^2 y \cdot \nabla \phi_{Q_\mu}}{Q_\mu} dy \geq \\
 &\geq \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy - \frac{\mu}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy \\
 &\quad - 2 \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy \\
 (7.8) \quad &\geq - \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy - \frac{\mu}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy
 \end{aligned}$$

By using $\mu_s = -2\mu^2 + O(\frac{\mu^2}{|\log \mu|})$ we obtain that

$$\begin{aligned}
 \frac{\mu_s}{2\mu} (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) &\leq \mu \int_{\mathbb{R}^2} \frac{\varepsilon^2}{Q_\mu} dy - \frac{\mu}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy + \frac{\mu^2}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy \\
 (7.9) \quad &\leq \mu (\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \frac{\mu}{2} \int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy + \frac{\mu^2}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy.
 \end{aligned}$$

Since $\int_{\mathbb{R}^2} |\nabla \phi_\varepsilon|^2 dy \leq C\mu$ it holds

$$(7.10) \quad \frac{\mu_s}{2\mu} (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) \leq \mu (\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \frac{\mu^2}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy + \frac{C\mu^2}{2},$$

with $\delta > 0$ a sufficiently small constant independent of μ . The previous inequality and (7.4) imply the following bound on I,

$$\begin{aligned}
 \text{I} &= \frac{\mu_s}{2\mu} (\Lambda \varepsilon, \mathcal{M}_\mu^y \varepsilon) - \frac{\mu_s}{2} \left(\frac{\partial_\mu Q_\mu \varepsilon}{Q_\mu^2}, \varepsilon \right) \\
 &\leq \mu (\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \frac{\mu^2}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy \\
 &\quad - \frac{\mu^2}{2} \int_{\mathbb{R}^2} \frac{|y|^2 \varepsilon^2}{Q_\mu} dy + \frac{C\mu^2}{2} + O(\mu^{\frac{5}{2}}) \\
 (7.11) \quad &\leq \mu (\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \frac{\mu^2}{\delta},
 \end{aligned}$$

with $\delta > 0$ sufficiently small.

7.2 Estimation of the Error: II

REMARK 7.2. First we notice that $E = \tilde{\mu} \Lambda Q_\mu - \tilde{\mu}^2 \nabla \cdot (\partial_\mu Q_\mu \nabla \phi_{\partial_\mu Q_\mu})$ and

$$(\Lambda Q_\mu, \mathcal{M}_\mu^y \varepsilon) = (\mathcal{M}_\mu^y (\Lambda Q_\mu), \varepsilon) = \left(2 - \frac{M}{2\pi} \right) (1, \varepsilon) - \mu(| \cdot |^2, \varepsilon) = 0.$$

Thanks to this cancellation, we are able to close our energy estimates.

Hence, by using

$$|\nabla \phi_{\partial_\mu \mathcal{Q}_\mu}| \lesssim \min\left(\mu|y||\log|y||, \frac{|\log|y||}{|y|}\right) \quad \text{and} \quad |\partial_\mu \mathcal{Q}_\mu| \lesssim |y|^2 \mathcal{Q}_\mu$$

(these inequalities are from Proposition 4.4)

$$\begin{aligned} (E, \mathcal{M}_\mu^y \varepsilon) &= -\tilde{\mu}^2 (\nabla \cdot (\partial_\mu \mathcal{Q}_\mu \nabla \phi_{\partial_\mu \mathcal{Q}_\mu}), \mathcal{M}_\mu^y \varepsilon) \\ &\lesssim \mu^2 \left(\int_{\mathbb{R}^2} \frac{|\partial_\mu \mathcal{Q}_\mu|^2 |\nabla \phi_{\partial_\mu \mathcal{Q}_\mu}|^2}{\mathcal{Q}_\mu} dy \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^2} \mathcal{Q}_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\ &\lesssim \mu^2 |\log \mu|^{\frac{3}{2}} \left(\int_{\mathbb{R}^2} \mathcal{Q}_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\ (7.12) \quad &\lesssim \mu \int_{\mathbb{R}^2} \mathcal{Q}_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy + \mu^3 |\log \mu|^3. \end{aligned}$$

7.3 Estimation of III (Linear Terms)

Since

$$\Theta_\mu(\varepsilon) = -\tilde{\mu} \nabla \cdot [\partial_\mu \mathcal{Q}_\mu \nabla \phi_\varepsilon + \varepsilon \nabla \phi_{\partial_\mu \mathcal{Q}_\mu}],$$

and by

$$|\nabla \phi_{\partial_\mu \mathcal{Q}_\mu}| \lesssim \min\left(\mu|y||\log|y||, \frac{|\log|y||}{|y|}\right) \quad \text{and} \quad |\partial_\mu \mathcal{Q}_\mu| \lesssim |y|^2 \mathcal{Q}_\mu,$$

we have

$$\begin{aligned} &(\Theta_\mu(\varepsilon), \mathcal{M}_\mu^y \varepsilon) \\ &= \tilde{\mu} \left[\int_{\mathbb{R}^2} \partial_\mu \mathcal{Q}_\mu \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon dy + \int_{\mathbb{R}^2} \varepsilon \nabla \phi_{\partial_\mu \mathcal{Q}_\mu} \cdot \nabla \mathcal{M}_\mu^y \varepsilon dy \right] \\ &\lesssim \mu \left(\int_{\mathbb{R}^2} \mathcal{Q}_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\ &\quad \times \left[\left(\int_{\mathbb{R}^2} \frac{|\partial_\mu \mathcal{Q}_\mu|^2 |\nabla \phi_\varepsilon|^2}{\mathcal{Q}_\mu} dy \right)^{\frac{1}{2}} + \left(\int_{\mathbb{R}^2} \frac{|\nabla \phi_{\partial_\mu \mathcal{Q}_\mu}|^2 |\varepsilon|^2}{\mathcal{Q}_\mu} dy \right)^{\frac{1}{2}} \right] \\ (7.13) \quad &\lesssim \mu [\|\nabla \phi_\varepsilon\|_{L^2} + \|\varepsilon\|_{L^2_{\mathcal{Q}_\mu}}] \left(\int_{\mathbb{R}^2} \mathcal{Q}_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}}. \end{aligned}$$

Hence, by the bootstrap assumption $\|\varepsilon\|_{L^2_{\mathcal{Q}_\mu}} \leq \sqrt{A\mu}$ and $\|\nabla \phi_\varepsilon\|_{L^2} = O(\sqrt{\mu})$ we get

$$\text{III} = (\Theta_\mu(\varepsilon), \mathcal{M}_\mu^y \varepsilon) \lesssim \mu^{\frac{5}{2}} + \sqrt{\mu} \int_{\mathbb{R}^2} \mathcal{Q}_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy.$$

7.4 Estimation of the Nonlinear Term: IV

For the nonlinear term we will use a Sobolev inequality not on ε but on $\mathcal{M}_\mu^y \varepsilon$ because we control only the gradient of $\mathcal{M}_\mu^y \varepsilon$. To do so we use the identity $\varepsilon = Q_\mu \mathcal{M}_\mu^y \varepsilon + \phi_\varepsilon Q_\mu$:

$$\begin{aligned}
 (N(\varepsilon), \mathcal{M}_\mu^y \varepsilon) &= - \int_{\mathbb{R}^2} \varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy \\
 (7.14) \quad &= - \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy - \int_{\mathbb{R}^2} Q_\mu \phi_\varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy.
 \end{aligned}$$

To estimate $\int_{\mathbb{R}^2} Q_\mu \phi_\varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy$ we use $\sup_{y \in \mathbb{R}^2} Q_\mu = 8$ and Proposition B.2:

$$\begin{aligned}
 (7.15) \quad \int_{\mathbb{R}^2} Q_\mu \phi_\varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy &\leq \\
 &\|\phi_\varepsilon\|_{L^\infty} \|\sqrt{Q_\mu} \nabla \phi_\varepsilon\|_{L^2} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 \, dy \right)^{\frac{1}{2}}.
 \end{aligned}$$

By the Morrey inequality, $-\Delta \phi_\varepsilon = \varepsilon$, and interpolation we get

$$\begin{aligned}
 \|\phi_\varepsilon\|_{L^\infty} &\lesssim \|\phi_\varepsilon\|_{L^4} + \|\nabla \phi_\varepsilon\|_{L^4} \lesssim \|\nabla \phi_\varepsilon\|_{L^2}^{1/2} (\|\Delta \phi_\varepsilon\|_{L^2}^{1/2} + \|\phi_\varepsilon\|_{L^2}^{1/2}) \\
 (7.16) \quad &\lesssim \|\nabla \phi_\varepsilon\|_{L^2}^{1/2} (\|\varepsilon\|_{L_{Q_\mu}^2}^{1/2} + \|\phi_\varepsilon\|_{L^2}^{1/2}).
 \end{aligned}$$

To control $\|\phi_\varepsilon\|_{L^2}$ we use Lemma C.1:

$$(7.17) \quad \|\phi_\varepsilon\|_{L^2}^{1/2} \lesssim \|\varepsilon\|_{L_{Q_\mu}^2}^{1/2}.$$

Hence, by using the bootstrap assumption, we obtain $\|\varepsilon\|_{L_{Q_\mu}^2} \leq \sqrt{A\mu}$

$$\begin{aligned}
 \int_{\mathbb{R}^2} Q_\mu \phi_\varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy &\lesssim \|\nabla \phi_\varepsilon\|_{L^2}^{3/2} \|\varepsilon\|_{L_{Q_\mu}^2}^{1/2} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 \, dy \right)^{\frac{1}{2}} \\
 &\lesssim A^{\frac{1}{4}} \mu \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 \, dy \right)^{\frac{1}{2}} \\
 (7.18) \quad &\lesssim \frac{\sqrt{A}}{\delta} \mu^2 + \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 \, dy,
 \end{aligned}$$

with δ a sufficiently small constant. Since the size of $\int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon \, dy$ is critical, we will take out the biggest component of ε that is in the direction of ΛQ_μ . To do so, we introduce the following decomposition of ε :

$$\varepsilon = \alpha_\mu \Lambda Q_\mu + \hat{\varepsilon},$$

and we fix α_μ such that

$$(7.19) \quad \int_{\mathbb{R}^2} \phi_{\Lambda Q_\mu} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} dy = 0.$$

This specific choice of the orthogonality condition is crucial here. Indeed, this choice will ensure that α_μ has a weak dependence in A (the bootstrap constant) and that the constant C in (A.1) is uniform with respect to μ :

$$\int_B Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy \leq C \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy.$$

First let us prove that

$$|\alpha_\mu| \lesssim \sqrt{\mu}.$$

Recall from (3.8) that

$$\phi_{\Lambda Q_\mu} Q_\mu = \Lambda Q_\mu - \left[\left(2 - \frac{M}{2\pi} \right) - \mu |\cdot|^2 \right] Q_\mu.$$

On the one hand, using Proposition 3.1 and Lemma 4.7 we have

$$\begin{aligned} & \int_{\mathbb{R}^2} \phi_{\Lambda Q_\mu} Q_\mu \mathcal{M}_\mu^y \varepsilon dy \\ &= \alpha_\mu \int_{\mathbb{R}^2} \left[\Lambda Q_\mu - \left[\left(2 - \frac{M}{2\pi} \right) - \mu |\cdot|^2 \right] Q_\mu \right] \mathcal{M}_\mu^y \Lambda Q_\mu dy \\ (7.20) \quad &= \alpha_\mu \int_{\mathbb{R}^2} \left[\Lambda Q_\mu - \left[\left(2 - \frac{M}{2\pi} \right) - \mu |\cdot|^2 \right] Q_\mu \right] \left(2 - \frac{M}{2\pi} - \mu |y|^2 \right) dy. \end{aligned}$$

Since $(\Lambda Q_\mu, 1) = 0$ and $(\Lambda Q_\mu, |\cdot|^2) = -2(Q_\mu, |\cdot|^2) = -\frac{M \log \mu}{2\pi} + O(1)$, it follows that

$$(7.21) \quad \int_{\mathbb{R}^2} \phi_{\Lambda Q_\mu} Q_\mu \mathcal{M}_\mu^y \varepsilon dy = \alpha_\mu \left(\left(2 - \frac{M}{2\pi} \right)^2 M + O(\mu |\log \mu|) \right).$$

On the other hand, using that $\int_{\mathbb{R}^2} \varepsilon dy = \int_{\mathbb{R}^2} \varepsilon |y|^2 dy = (\mathcal{M}_\mu^y \varepsilon, \Lambda Q_\mu)_{L^2} = 0$, we get

$$\begin{aligned} & \int_{\mathbb{R}^2} \phi_{\Lambda Q_\mu} Q_\mu \mathcal{M}_\mu^y \varepsilon dy = \int_{\mathbb{R}^2} \left[\Lambda Q_\mu - \left[\left(2 - \frac{M}{2\pi} \right) - \mu |\cdot|^2 \right] Q_\mu \right] \mathcal{M}_\mu^y \varepsilon dy \\ &= - \int_{\mathbb{R}^2} \left[\left(2 - \frac{M}{2\pi} \right) - \mu |\cdot|^2 \right] Q_\mu \mathcal{M}_\mu^y \varepsilon dy \\ (7.22) \quad &= - \int_{\mathbb{R}^2} \left[\left(2 - \frac{M}{2\pi} \right) - \mu |\cdot|^2 \right] Q_\mu \phi_\varepsilon. \end{aligned}$$

Hence, using that ε has average zero and the decay of Q_μ , we get that

$$(7.23) \quad |\alpha_\mu| \lesssim \int_{\mathbb{R}^2} Q_\mu |\phi_\varepsilon| dy \lesssim \|\nabla \phi_\varepsilon\|_{L^2}.$$

Also, we prove in the following proposition that the difference between $(\mathcal{M}_\mu^y \varepsilon, \varepsilon)$ and $(\mathcal{M}_\mu^y \hat{\varepsilon}, \hat{\varepsilon})$ is of order $\mu^2 |\log \mu|$.

PROPOSITION 7.3. *If*

$$(\varepsilon, 1)_{L^2} = (\varepsilon, |\cdot|^2)_{L^2} = 0, \quad \varepsilon = \alpha_\mu \Lambda Q_\mu + \hat{\varepsilon},$$

$$\text{and} \quad \int_{\mathbb{R}^2} \phi_{\Lambda Q_\mu} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} dy = 0,$$

then

$$(\mathcal{L}_\mu^y \varepsilon, \mathcal{M}_\mu^y \varepsilon) = (\mathcal{L}_\mu^y \hat{\varepsilon}, \mathcal{M}_\mu^y \hat{\varepsilon}) + O(\mu^3 |\log \mu|),$$

$$(\mathcal{M}_\mu^y \varepsilon, \varepsilon) = (\mathcal{M}_\mu^y \hat{\varepsilon}, \hat{\varepsilon}) + O(\mu^2 |\log \mu|).$$

The proof is in the Appendix B. We plug the decomposition of $\varepsilon = \alpha_\mu \Lambda Q_\mu + \hat{\varepsilon}$ into

$$\int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon dy.$$

It then follows that

$$\begin{aligned} & \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon dy \\ &= \alpha_\mu \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y (\Lambda Q_\mu) \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \\ & \quad + \alpha_\mu \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y (\Lambda Q_\mu) dy \\ & \quad + \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \\ & \quad + \alpha_\mu^2 \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y (\Lambda Q_\mu) \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y (\Lambda Q_\mu) dy. \end{aligned}$$

Hence, by using $\mathcal{M}_\mu^y (\Lambda Q_\mu) = 2 - \frac{M}{2\pi} - \mu|y|^2$, we deduce

$$\begin{aligned} & \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \varepsilon \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \varepsilon dy \\ &= \underbrace{\alpha_\mu \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y (\Lambda Q_\mu) \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy}_{\text{IV}_1} \\ (7.24) \quad & - \underbrace{2\mu\alpha_\mu \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot y dy}_{\text{IV}_2} + \underbrace{\int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy}_{\text{IV}_3} \\ & - \underbrace{2\mu\alpha_\mu^2 \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y (\Lambda Q_\mu) \nabla \phi_\varepsilon \cdot y dy}_{\text{IV}_4}. \end{aligned}$$

We start by estimating the term IV_1 .

Estimation of IV_1

Since $\mathcal{M}_\mu^y(\Lambda Q_\mu) = 2 - \frac{M}{2\pi} - \mu|y|^2$ and by (7.23) we obtain

$$\begin{aligned}
 \text{IV}_1 &\lesssim |\alpha_\mu| \int_{\mathbb{R}^2} Q_\mu \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy + |\alpha_\mu| \mu \int_{\mathbb{R}^2} Q_\mu |y|^2 \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \\
 (7.25) \quad &\lesssim \sqrt{\mu} \|\nabla \phi_\varepsilon\|_{L^2} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \\
 &\quad + \mu^{\frac{3}{2}} \|\sqrt{Q_\mu} |y|^2 \nabla \phi_\varepsilon\|_{L^2} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}}.
 \end{aligned}$$

Notice that $\sqrt{Q_\mu} |y|^2$ is uniformly bounded and $\|\nabla \phi_\varepsilon\|_{L^2} \lesssim \sqrt{\mu}$, which implies

$$\begin{aligned}
 \text{IV}_1 &\lesssim \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy + \frac{1}{\delta} \mu^2 + \frac{1}{\delta} \mu^4 \\
 &\lesssim \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy + \frac{1}{\delta} \mu^2.
 \end{aligned}$$

If we use Proposition B.1 we obtain

$$(7.26) \quad \text{IV}_1 \lesssim \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy + \frac{1}{\delta} \mu^2 + \mu^3 |\log \mu|.$$

Estimation of IV_2

Now we estimate IV_2 . We remark that $\sqrt{Q_\mu} |y|$ is uniformly bounded, which implies

$$\begin{aligned}
 \text{IV}_2 &\lesssim \mu |\alpha_\mu| \left(\int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \|\sqrt{Q_\mu} |y| \nabla \phi_\varepsilon\|_{L^2} \\
 (7.27) \quad &\lesssim \mu |\alpha_\mu| \left(\int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \|\nabla \phi_\varepsilon\|_{L^2}.
 \end{aligned}$$

Notice that

$$\begin{aligned}
 \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy &\lesssim \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy + \alpha_\mu^2 \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \Lambda Q_\mu|^2 dy \\
 (7.28) \quad &\lesssim \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy + C \alpha_\mu^2,
 \end{aligned}$$

which implies

$$(7.29) \quad \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \lesssim \int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy + \mu \lesssim A \mu.$$

Hence, if we combine all the previous inequalities, it follows that

$$(7.30) \quad \text{IV}_2 \lesssim \mu^2 A^{\frac{1}{2}} \|\nabla \phi_\varepsilon\|_{L^2} \lesssim \mu^{\frac{5}{2}} A^{\frac{1}{2}}.$$

Estimation of IV₄

To estimate IV₄, we use $\mathcal{M}_\mu^y(\Lambda Q_\mu) = 2 - \frac{M}{2\pi} - \mu|y|^2$:

$$\begin{aligned} \text{IV}_4 &\lesssim \mu \alpha_\mu^2 \left[\int_{\mathbb{R}^2} Q_\mu |\nabla \phi_\varepsilon| |y| dy + \mu \int_{\mathbb{R}^2} Q_\mu |\nabla \phi_\varepsilon| |y|^3 dy \right] \\ (7.31) \quad &\lesssim \mu \alpha_\mu^2 \|\nabla \phi_\varepsilon\|_{L^2} [\|Q_\mu |y|\|_{L^2} + \|Q_\mu |y|^3\|_{L^2}] \lesssim \mu^{\frac{5}{2}}. \end{aligned}$$

Estimation of IV₃

To control the term IV₃ we first use the Gagliardo-Nirenberg inequality,

$$\|f\|_{L^4} \lesssim \|f\|_{L^2}^{1/2} \|\nabla f\|_{L^2}^{1/2},$$

with $f = \sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon}$. Hence,

$$\begin{aligned} &\int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \\ &\leq \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 |\nabla \phi_\varepsilon|^2 dy \right)^{\frac{1}{2}} \\ &\leq \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \|\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon}\|_{L^4} \|\nabla \phi_\varepsilon\|_{L^4} \\ (7.32) \quad &\leq \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \|\nabla \phi_\varepsilon\|_{L^4} \|\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon}\|_{L^2}^{1/2} \\ &\quad \cdot \|\nabla(\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon})\|_{L^2}^{1/2}. \end{aligned}$$

We also apply the Gagliardo-Nirenberg inequality to $\|\nabla \phi_\varepsilon\|_{L^4}$ and use $-\Delta \phi_\varepsilon = \varepsilon$, $L_{Q_\mu}^2 \hookrightarrow L^2$:

$$\begin{aligned} &\int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \leq \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} \|\nabla \phi_\varepsilon\|_{L^2}^{1/2} \|\varepsilon\|_{L_{Q_\mu}^2}^{1/2} \\ (7.33) \quad &\quad \times \|\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon}\|_{L^2}^{1/2} \|\nabla(\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon})\|_{L^2}^{1/2}. \end{aligned}$$

Now we estimate $\|\nabla(\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon})\|_{L^2}^2$. Since $\frac{|\nabla Q_\mu|^2}{Q_\mu} \lesssim (\mu^2 |y|^2 + \frac{|y|^2}{(1+|y|^2)^2}) Q_\mu$, we deduce

$$\begin{aligned} &\|\nabla(\sqrt{Q_\mu} \mathcal{M}_\mu^y \hat{\varepsilon})\|_{L^2}^2 \lesssim \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy + \int_{\mathbb{R}^2} \frac{|\nabla Q_\mu|^2}{Q_\mu} |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \\ (7.34) \quad &\lesssim \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy + \mu^2 \int_{\mathbb{R}^2} |y|^2 Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \\ &\quad + \int_{\mathbb{R}^2} \frac{|y|^2}{(1+|y|^2)^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy. \end{aligned}$$

Then to control $\mu^2 \int_{\mathbb{R}^2} |y|^2 Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy$ and $\int_{\mathbb{R}^2} \frac{|y|^2}{(1+|y|^2)^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy$, we use Proposition A.1 and Lemma A.2. Hence,

$$\begin{aligned}
 & \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \\
 (7.35) \quad & \leq \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{2}} A^{\frac{1}{2}} \mu^{\frac{3}{4}} \\
 & \quad \times \left[\left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{4}} + \mu^{\frac{1}{4}} \left(\int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \hat{\varepsilon}|^2 dy \right)^{\frac{1}{4}} \right].
 \end{aligned}$$

From (7.29) and Holder we get

$$\begin{aligned}
 (7.36) \quad & \int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} \nabla \phi_\varepsilon \cdot \nabla \mathcal{M}_\mu^y \hat{\varepsilon} dy \lesssim \\
 & \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy + \frac{\mu^3 A^2}{\delta} + \mu^3 |\log \mu|.
 \end{aligned}$$

Finally,

$$(N(\varepsilon), \mathcal{M}_\mu^y \varepsilon) \lesssim \frac{1 + A^{\frac{1}{2}}}{\delta} \mu^2 + \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy,$$

with $\delta > 0$ a sufficiently small constant.

7.5 Estimation of the Forcing Term: V

We notice first

$$F = \frac{\mu_s}{2\mu} [\Lambda Q_\mu - \tilde{\mu} \Lambda \partial_\mu Q_\mu] - \mu_s [\partial_\mu Q_\mu - \partial_\mu \tilde{\mu} \partial_\mu Q_\mu - \tilde{\mu} \partial_\mu^2 Q_\mu],$$

and since $(\partial_\mu Q_\mu, \mathcal{M}_\mu^y \varepsilon) = (\Lambda Q_\mu, \mathcal{M}_\mu^y \varepsilon) = 0$, it follows that

$$\begin{aligned}
 (F, \mathcal{M}_\mu^y \varepsilon) &= -\tilde{\mu} \frac{\mu_s}{2\mu} (\Lambda \partial_\mu Q_\mu, \mathcal{M}_\mu^y \varepsilon) + \mu_s \tilde{\mu} (\partial_\mu^2 Q_\mu, \mathcal{M}_\mu^y \varepsilon) \\
 &= -\tilde{\mu} \frac{\mu_s}{2\mu} (y \cdot \nabla \partial_\mu Q_\mu, \mathcal{M}_\mu^y \varepsilon) + \mu_s \tilde{\mu} (\partial_\mu^2 Q_\mu, \mathcal{M}_\mu^y \varepsilon) \\
 &= \tilde{\mu} \frac{\mu_s}{2\mu} (\partial_\mu Q_\mu, \nabla \cdot (y \mathcal{M}_\mu^y \varepsilon)) + \mu_s \tilde{\mu} (\partial_\mu^2 Q_\mu, \mathcal{M}_\mu^y \varepsilon) \\
 (7.37) \quad &= \tilde{\mu} \frac{\mu_s}{2\mu} (\partial_\mu Q_\mu, y \cdot \nabla \mathcal{M}_\mu^y \varepsilon) + \mu_s \tilde{\mu} (\partial_\mu^2 Q_\mu, \mathcal{M}_\mu^y \varepsilon).
 \end{aligned}$$

Hence, if we use Proposition 4.4, $|\mu_s| \leq 2\mu^2$, and $\tilde{\mu} \leq \mu$, it holds that

$$\begin{aligned}
 (F, \mathcal{M}_\mu^y \varepsilon) &\lesssim \mu^3 \left(\int_{\mathbb{R}^2} Q_\mu |\mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^2} \frac{|\partial_\mu^2 Q_\mu|^2}{Q_\mu} dy \right)^{\frac{1}{2}} \\
 &\quad + \mu^2 \left(\int_{\mathbb{R}^2} \frac{|\partial_\mu Q_\mu|^2 |y|^2}{Q_\mu} dy \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy \right)^{\frac{1}{2}} \\
 (7.38) \quad &\lesssim \left(\sqrt{A} + \frac{1}{\delta} \right) \mu^2 + \delta \int_{\mathbb{R}^2} Q_\mu |\nabla \mathcal{M}_\mu^y \varepsilon|^2 dy.
 \end{aligned}$$

Finally, combining all the previous estimates, we obtain

$$(7.39) \quad \frac{1}{2} \frac{d}{ds} (\mathcal{M}_\mu^y \varepsilon, \varepsilon) + \mu (K_2 - 1 - \delta) (\mathcal{M}_\mu^y \varepsilon, \varepsilon) \leq C(\delta, A) \mu^2,$$

with $\delta \ll 1$ and $C(A, \delta) \lesssim \frac{1}{\delta} (\sqrt{A}(1 + \delta) + 1) \leq A$ for A sufficiently large, which concludes the energy estimates. \square

Now we prove Proposition 2.3,

PROOF OF PROPOSITION 2.3. As set before, $\mathcal{E}(s) = (\mathcal{M}_\mu^y \varepsilon, \varepsilon)$, and now set $K = K_2 - 1 - \delta > 1$. If we use that $|\mu(s) - \frac{1}{2s}| \leq \frac{C'}{s|\log s|}$, then

$$(7.40) \quad \frac{1}{2} \mathcal{E}'(s) + \frac{K}{2s} \mathcal{E}(s) \leq \frac{C(\delta, A)}{4s^2}.$$

Hence,

$$\frac{d}{ds} (s^K \mathcal{E}(s)) \leq \frac{C(\delta, A) s^{K-2}}{2};$$

if we integrate, it follows that

$$s^K \mathcal{E}(s) \lesssim C(\delta, A) s^{K-1},$$

which concludes the proof, and we easily see that if we pick originally A sufficiently large, the constant in the bootstrap gets smaller after each step. \square

8 Proof of Corollary 1.4

PROOF OF COROLLARY 1.4. First, let us denote

$$(8.1) \quad m_u(r, s) = 2\pi \int_0^r \tau u(\tau, t) d\tau,$$

where $r = |x|$. Hence, we obtain the following local equation on m_u :

$$\partial_t m_u = \partial_{rr} m_u - \frac{\partial_r m_u}{r} + \frac{m_u \partial_r m_u}{r},$$

$$m_u(t = 0) = m_0 = \int_0^r \tau u_0(\tau) d\tau,$$

$$m_u(r, t) \longrightarrow 8\pi \quad \text{as } r \rightarrow +\infty.$$

The proof of the corollary is just based on the fact that for radial data, we can use the equation on the partial mass and the comparison principle for radial solutions (see [28]): If, for all $r \geq 0$, $m^2(t = 0, r) \leq m^1(t = 0, r)$, then for all $t > 0$ and $r \geq 0$,

$$m^2(t, r) \leq m^1(t, r).$$

Let us choose a radially symmetric initial data u_0^1 satisfying the condition of Theorem 1.1, compactly supported and such that $u_0^1(0) > 0$. Consequently, for $a > 0$ large enough we can achieve that for all $r \geq 0$,

$$m_0(r) \leq m_0^1(r),$$

where

$$m_0^1(r) = 2\pi \int_0^r \tau (u_0^1)_{1/a}(\tau) d\tau.$$

In addition, if we choose

$$u_0^2 = \frac{K}{\mu_0} Q\left(\frac{x}{\sqrt{\mu_0}}\right) e^{-\frac{x^2}{2}},$$

where $K > 1$ is a constant picked to insure that the mass of u_0^2 is 8π . Moreover, u_0^2 verifies the condition of Theorem 1.1, then, for a large enough one can deduce easily that

$$m_0^2(r) \leq m_0(r),$$

where

$$m_0^2(r) = 2\pi \int_0^r \tau (u_0^2)_a(\tau) d\tau.$$

Hence by applying the comparison principle it follows that

$$m_{u^2}(t) \leq m_u(t) \leq m_{u^1}(t),$$

where u^2 is the solution with initial data $(u_0^2)_a$ and u^1 is the solution with initial data $(u_0^1)_{1/a}$. In other words, because we can use the comparison principle we can bound the partial mass of any radial solution between two rescaled solutions for which Theorem 1.1 applies. To prove (1.10), we use that u^1 and u^2 follow the dynamic of the solutions of Theorem 1.1. Indeed, we can decompose u^1 and u^2 as in (1.8), but to be more consistent we use the decomposition (2.6):

$$u^i(x, t) = \frac{1}{\lambda_i^2} \tilde{Q}_{\mu_i}\left(\frac{x}{\lambda_i}, t\right) + \frac{1}{\lambda_i^2} \varepsilon_i\left(\frac{x}{\lambda_i}, t\right) \quad \text{for } i \in \{1, 2\},$$

where \tilde{Q}_{μ_i} is the approximate profile constructed in Section 4 and μ_i is a function of t fixed in Section 5. Actually, with this decomposition we have

$$\int_{\mathbb{R}^2} \varepsilon_i dy = 0.$$

Consequently,

$$(8.2) \quad \int_0^{\frac{r}{\lambda_i}} \varepsilon_i(r') r' dr' = - \int_{\frac{r}{\lambda_i}}^{\infty} \varepsilon_i(r') r' dr'.$$

Moreover, we use that

$$\tilde{Q}_{\mu_i} = Q_{\mu_i} - \tilde{\mu}_i \partial_{\mu_i} Q_{\mu_i},$$

where $\tilde{\mu}_i = \mu_i + O(\mu_i / |\log \mu_i|)$.

From Propositions 4.3 and 4.4, we deduce that

$$\begin{aligned} 8\pi - m \tilde{Q}_{\mu_i} &= \int_{\frac{r}{\lambda_i}}^{\infty} Q_{\mu_i}(r') r' dr' - \tilde{\mu}_i \int_{\frac{r}{\lambda_i}}^{\infty} \partial_{\mu} Q_{\mu_i} r' dr' \\ &= \int_{\frac{r}{\lambda_i}}^{\infty} Q e^{\mu_i \phi_{T_1} + \sigma - \mu_i \frac{r'^2}{2}} r' dr' \\ &\quad - \tilde{\mu}_i \int_{\frac{r}{\lambda_i}}^{\infty} \left(\phi_{T_1} + \partial_{\mu_i} \sigma - \frac{r'^2}{2} \right) Q e^{\mu_i \phi_{T_1} + \sigma - \mu_i \frac{r'^2}{2}} r' dr', \end{aligned}$$

and there exists $0 < c < 1$, uniform with respect to μ_i , such that

$$|\mu_i \phi_{T_1} + \sigma| \leq c \mu_i \frac{r^2}{2}.$$

Hence, for all $r > 0$, we deduce that

$$\begin{aligned} &|8\pi - m \tilde{Q}_{\mu_i}| \\ &\leq \int_{\frac{r}{\lambda_i}}^{\infty} Q e^{-\mu_i(1-c)\frac{r'^2}{2}} r' dr' + \mu_i(1+c) \int_{\frac{r}{\lambda_i}}^{\infty} Q e^{-\mu_i(1-c)\frac{r'^2}{2}} r' dr' \\ &\leq \frac{1+c}{1-c} \int_{\frac{r}{\lambda_i}}^{\infty} Q e^{-\mu_i(1-c)\frac{r'^2}{2}} r' dr' \\ &\quad + \mu_i \frac{(1+c)(1-c)}{1-c} \int_{\frac{r}{\lambda_i}}^{\infty} Q e^{-\mu_i(1-c)\frac{r'^2}{2}} r' dr' \\ &\leq \frac{8(1+c)}{1-c} \left[\int_{\frac{r}{\lambda_i}}^{\infty} \frac{1}{r'^3} e^{-\mu_i(1-c)\frac{r'^2}{2}} dr' \right. \\ &\quad \left. + \frac{\mu_i(1-c)}{2} \int_{\frac{r}{\lambda_i}}^{\infty} \frac{1}{r'} e^{-\mu_i(1-c)\frac{r'^2}{2}} dr' \right]. \end{aligned}$$

By integration by parts we obtain that

$$\begin{aligned} \frac{\mu_i(1-c)}{2} \int_{\frac{r}{\lambda_i}}^{\infty} \frac{1}{r'} e^{-\mu_i(1-c)\frac{r'^2}{2}} dr' = \\ \frac{\lambda_i^2 e^{-\mu_i(1-c)\frac{r^2}{2}}}{2r^2} - \int_{\frac{r}{\lambda_i}}^{\infty} \frac{1}{r'^3} e^{-\mu_i(1-c)\frac{r'^2}{2}} dr'. \end{aligned}$$

Consequently, we get for all $r > 0$ that

$$(8.3) \quad |8\pi - m_{\tilde{Q}_{\mu_i}}| \lesssim \frac{\lambda_i^2 e^{-\mu_i(1-c)\frac{r^2}{2\lambda_i}}}{r^2}.$$

From Lemma 5.1 and Theorem 1.1 we get that

$$\mu_i(t) = \frac{\sqrt{I_i}}{t \log(2t+1) + O(t \log \log(2t+1))}$$

and

$$\lambda_i(t) = R(t)\mu_i(t) \quad \text{with } R(t) = \sqrt{2t+1}.$$

Hence,

$$|8\pi - m_{\tilde{Q}_{\mu_i}}| \lesssim \frac{\lambda_i^2 e^{-C_1 \frac{r^2}{t}}}{r^2}.$$

Since, $\int_{\mathbb{R}^2} \varepsilon dy = 0$, we deduce that

$$\begin{aligned} \int_{\frac{r}{\lambda_i}}^{\infty} \varepsilon_i(r') r' dr' &\lesssim \left(\int_{\frac{r}{\lambda_i}}^{\infty} Q_{\mu_i}(r') dr' \right)^{\frac{1}{2}} \|\varepsilon_i\|_{L^2_{Q_{\mu_i}}} \\ &\lesssim \left(\int_{\frac{r}{\lambda_i}}^{\infty} Q(r') dr' \right)^{\frac{1}{2}} \|\varepsilon_i\|_{L^2_{Q_{\mu_i}}} \\ &\lesssim \frac{\lambda_i^2}{\lambda_i^2 + r^2} \|\varepsilon_i\|_{L^2_{Q_{\mu}}} \lesssim \frac{\lambda_i^2}{(\lambda_i^2 + r^2)t |\log t|}. \end{aligned}$$

Consequently,

$$\begin{aligned} 8\pi - m_{u^2}(r, t) = 8\pi - m_{\tilde{Q}_{\mu_2}} - m_{\varepsilon_2} &\lesssim \frac{\lambda_2^2 e^{-C_2 \frac{r^2}{2}}}{r^2} + m_{\varepsilon_2} \left(\frac{r}{\lambda_2}, t \right) \\ &\lesssim \frac{\lambda_2^2 e^{-C_2 \frac{r^2}{2}}}{r^2} + \frac{\lambda_2^2}{(\lambda_2^2 + r^2)t |\log t|} \end{aligned}$$

and

$$8\pi - m_{u^1}(r, t) \gtrsim -\frac{\lambda_1^2 e^{-C_1 \frac{r^2}{2}}}{r^2} - \frac{\lambda_1^2}{(\lambda_1^2 + r^2)t |\log t|},$$

with $\lambda_i(t) = \frac{\sqrt{I_i}}{\sqrt{\log(2t+1) + O(\log(\log(2t+1)))}}$, where I_i depends on a and the second moment of u . This concludes the proof. \square

Appendix A Hardy Inequalities

PROPOSITION A.1. *Let $f \in H^1(Q_\mu dy)$ be such that $(f, \phi_{\Lambda Q_\mu} Q_\mu)_{L^2} = 0$; then there exists C uniform with respect to μ such that*

$$(A.1) \quad \int_{\mathbb{R}^2} |f|^2 \frac{|y|^2}{(1+|y|^2)^2} Q_\mu dy \leq C \int_{\mathbb{R}^2} Q_\mu |\nabla f|^2 dy.$$

PROOF. We first prove the inequality

$$(A.2) \quad \int_{\mathbb{R}^2} |\nabla f|^2 Q_\mu dy + \int_B |f|^2 Q_\mu dy \geq \frac{3}{4} \int_{\mathbb{R}^2} |f|^2 \frac{|y|^2}{(1+|y|^2)^2} Q_\mu dy.$$

Indeed, let γ be a positive constant that we will fix later,

$$(A.3) \quad \begin{aligned} \int_{\mathbb{R}^2} \left| \nabla f - \gamma \frac{yf}{1+|y|^2} \right|^2 Q_\mu dy &= \int_{\mathbb{R}^2} |\nabla f|^2 Q_\mu dy - 2\gamma \int_{\mathbb{R}^2} \frac{fy \cdot \nabla f}{1+|y|^2} Q_\mu dy \\ &+ \gamma^2 \int_{\mathbb{R}^2} \frac{|f|^2 |y|^2}{(1+|y|^2)^2} Q_\mu dy \geq 0. \end{aligned}$$

By integration by parts, we deduce

$$(A.4) \quad \begin{aligned} -2\gamma \int_{\mathbb{R}^2} \frac{fy \cdot \nabla f}{1+|y|^2} Q_\mu dy &= \gamma \int_{\mathbb{R}^2} \frac{|f|^2 \nabla Q_\mu \cdot y}{1+|y|^2} dy + 2\gamma \int_{\mathbb{R}^2} \frac{|f|^2}{1+|y|^2} Q_\mu dy \\ &- 2\gamma \int_{\mathbb{R}^2} \frac{|f|^2 |y|^2}{(1+|y|^2)^2} Q_\mu dy. \end{aligned}$$

Since

$$\begin{aligned} \nabla Q_\mu \cdot y &= (y \cdot \nabla \phi_Q + \mu y \cdot \nabla \phi_{T_1} + y \cdot \nabla \sigma - \mu |y|^2) Q_\mu \\ &= \left(-\frac{4|y|^2}{1+|y|^2} + \mu y \cdot \nabla \phi_{T_1} + y \cdot \nabla \sigma - \mu |y|^2 \right) Q_\mu, \end{aligned}$$

$$\phi'_{T_1}(r) = O(r^3 \log r) \mathbb{1}_{\{r \leq 1\}} + O\left(\frac{\log r}{r}\right) \mathbb{1}_{\{r \geq 1\}},$$

and

$$\partial_r \sigma \lesssim \min\left(\mu^2 r |\log \langle r \rangle|, \mu \frac{|\log \langle r \rangle|}{r}\right),$$

it follows that

$$\begin{aligned}
 & \gamma \int_{\mathbb{R}^2} \frac{|f|^2 \nabla Q_\mu \cdot y}{1 + |y|^2} dy + 2\gamma \int_{\mathbb{R}^2} \frac{|f|^2}{1 + |y|^2} Q_\mu dy \\
 &= \gamma \int_{\mathbb{R}^2} \frac{|f|^2}{1 + |y|^2} \left[\frac{2(1 - |y|^2)}{1 + |y|^2} \right. \\
 &\quad \left. - \underbrace{\mu \left(|y|^2 - y \cdot \nabla \phi_{T_1} - \frac{y \cdot \nabla \sigma}{\mu} \right)}_{>0} \right] Q_\mu dy \\
 (A.5) \quad & \leq 2\gamma \int_B Q_\mu |f|^2 dy.
 \end{aligned}$$

Hence, if we combine all the previous inequalities we obtain

$$\int_{\mathbb{R}^2} |\nabla f|^2 Q_\mu dy + 2\gamma \int_B Q_\mu |f|^2 dy \geq (2\gamma - \gamma^2) \int_{\mathbb{R}^2} \frac{|f|^2 |y|^2}{(1 + |y|^2)^2} Q_\mu dy;$$

if we select $\gamma = \frac{1}{2}$, then (A.2) follows. Now to show (A.1) we prove that if $(f, \phi_{\Lambda Q_\mu} Q_\mu)_{L^2} = 0$, then

$$(A.6) \quad \int_B Q_\mu |f|^2 \leq C \int_{\mathbb{R}^2} Q_\mu |\nabla f|^2 dy,$$

with C uniform with respect to μ . We first by contradiction prove (A.6) for fixed μ . Assume that (A.6) is false, then there exists $f_n \in H^1(\mathbb{R}^2, Q_\mu dy)$ such that

$$\int_{\mathbb{R}^2} Q_\mu |\nabla f_n|^2 dy \leq \frac{1}{n}, \quad \int_B Q_\mu |f_n|^2 dy = 1, \quad \text{and} \quad (f_n, \phi_{\Lambda Q_\mu} Q_\mu)_{L^2} = 0.$$

Hence, up to a subsequence $f_n \rightarrow f_\infty$ in L^2_{loc} where f_∞ is a constant. In addition, by (A.2) we deduce

$$\int_{\mathbb{R}^2} \frac{|f_n|^2 |y|^2}{(1 + |y|^2)^2} Q_{\mu_n} dy \leq 1.$$

Since $\phi_{\Lambda Q_\mu}(y) = O(1/(1 + |y|^2))$ we can pass to the limit in

$$(A.7) \quad (f_n, \phi_{\Lambda Q_\mu} Q_\mu)_{L^2} = 0,$$

and deduce that

$$(f_\infty, \phi_{\Lambda Q_\mu} Q_\mu)_{L^2} = 0.$$

Since $(1, \phi_{\Lambda Q_\mu} Q_\mu)_{L^2} \neq 0$, we get $f_\infty = 0$. However,

$$\int_B |f_\infty|^2 Q_\mu dy = 1,$$

which contradicts $f_\infty = 0$. To prove the uniformity of C in (A.6), we argue by contradiction. Suppose that C is not uniform with respect to μ , then there exist $C_n > 0$, $\mu_n > 0$, and $f_n \in H^1(\mathbb{R}^2, Q_{\mu_n} dy)$ such that

$$C_n \rightarrow +\infty, \quad \mu_n \rightarrow \mu^*,$$

with $\mu^* \geq 0$.

$$\int_{\mathbb{R}^2} Q_{\mu_n} |\nabla f_n|^2 dy = \frac{1}{C_n}, \quad \int_B Q_{\mu_n} |f_n|^2 dy = 1,$$

and $(f_n, \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n})_{L^2} = 0$.

Actually, $\mu^* = 0$ is the most difficult case, since when $\mu \rightarrow 0$, Q_μ loses its exponential decay at infinity. However, thanks to the fact that $\phi_{\Lambda Q} = O\left(\frac{1}{1+|y|^2}\right)$ has extra decay, one can still pass to the limit in $(f_n, \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n})_{L^2} = 0$. Indeed, for $R > 0$, we have

$$0 = (f_n, \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n})_{L^2} = \int_{B_R} f_n \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n} dy + \int_{\mathbb{R}^2 \setminus B_R} f_n \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n} dy$$

and

$$\int_{B_R} f_n \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n} dy \longrightarrow \int_{B_R} f_\infty \phi_{\Lambda Q} Q dy.$$

Moreover, thanks to (A.2) we get that

$$\begin{aligned} \int_{\mathbb{R}^2 \setminus B_R} f_n \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n} dy &\lesssim \left(\int_{\mathbb{R}^2 \setminus B_R} \frac{|f_n|^2 |y|^2}{(1+|y|^2)^2} Q_{\mu_n} dy \right)^{\frac{1}{2}} \\ &\quad \times \left(\int_{\mathbb{R}^2 \setminus B_R} |\phi_{\Lambda Q_{\mu_n}}|^2 |y|^2 Q_{\mu_n} dy \right)^{\frac{1}{2}} \\ (A.8) \quad &\lesssim \frac{1}{R^2} \left(\int_{\mathbb{R}^2 \setminus B_R} \frac{|f_n|^2 |y|^2}{(1+|y|^2)^2} Q_{\mu_n} dy \right)^{\frac{1}{2}}. \end{aligned}$$

Since

$$\left(\int_{\mathbb{R}^2 \setminus B_R} \frac{|f_n|^2 |y|^2}{(1+|y|^2)^2} Q_{\mu_n} dy \right)^{\frac{1}{2}}$$

is uniformly bounded with respect to n , we deduce that

$$\int_{\mathbb{R}^2 \setminus B_R} f_n \phi_{\Lambda Q_{\mu_n}} Q_{\mu_n} dy \longrightarrow 0,$$

when $R \rightarrow +\infty$ uniformly in n . For the other terms, one can pass to the limit as it has been done before and reach the same contradiction, which concludes the proof. \square

LEMMA A.2. *Let $f \in H^1(\mathbb{R}^2, Q_\mu dy)$ then there exists $c' > 0$ uniform with respect to μ such that*

$$\int_{\mathbb{R}^2} Q_\mu |y|^2 |f|^2 dy \leq \frac{c'}{\mu^2} \int_{\mathbb{R}^2} Q_\mu |\nabla f|^2 dy + \frac{c'}{\mu} \int_{B_1} Q_\mu |f|^2 dy,$$

where B_1 is the unit ball in \mathbb{R}^2 .

PROOF. The proof is based on the following identity with γ a positive constant that will be fixed at the end of the proof:

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla f - \gamma y f|^2 Q_\mu dy &= \int_{\mathbb{R}^2} |\nabla f|^2 Q_\mu dy - 2\gamma \int_{\mathbb{R}^2} f y \cdot \nabla f Q_\mu dy \\ &\quad + \gamma^2 \int_{\mathbb{R}^2} |f|^2 |y|^2 Q_\mu dy \geq 0. \end{aligned}$$

By integrating by parts we get

$$\begin{aligned} -2\gamma \int_{\mathbb{R}^2} f y \cdot \nabla f Q_\mu dy &= \gamma \int_{\mathbb{R}^2} |f|^2 \nabla \cdot (y Q_\mu) dy \\ &= 2\gamma \int_{\mathbb{R}^2} |f|^2 Q_\mu dy + \gamma \int_{\mathbb{R}^2} |f|^2 y \cdot \nabla Q_\mu dy. \end{aligned}$$

Hence,

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla f|^2 Q_\mu dy &\geq -\gamma^2 \int_{\mathbb{R}^2} |f|^2 |y|^2 Q_\mu dy - 2\gamma \int_{\mathbb{R}^2} |f|^2 Q_\mu dy \\ &\quad - \gamma \int_{\mathbb{R}^2} |f|^2 y \cdot \nabla Q_\mu dy. \end{aligned}$$

From Proposition 4.4 there exists $0 < c < 1$ uniform with respect to μ such that

$$y \cdot \nabla \phi_{T_1} + \frac{y \cdot \nabla \sigma}{\mu} \leq c |y|^2.$$

Since

$$y \cdot \nabla Q_\mu = y \cdot \nabla \phi_{Q_\mu} - \mu |y|^2 = -\frac{4|y|^2}{1+|y|^2} - \mu \left(|y|^2 - y \cdot \nabla \phi_{T_1} - \frac{y \cdot \nabla \sigma}{\mu} \right),$$

we get

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla f|^2 Q_\mu dy &\geq \\ &\quad -\gamma(\gamma - (1-c)\mu) \int_{\mathbb{R}^2} |f|^2 |y|^2 Q_\mu dy - 2\gamma \int_{\mathbb{R}^2} \frac{1-|y|^2}{1+|y|^2} |f|^2 Q_\mu dy. \end{aligned}$$

If we choose $\gamma = \frac{(1-c)\mu}{2}$ the conclusion follows. \square

Appendix B Identities between ε and $\hat{\varepsilon}$

PROPOSITION B.1. If $(\varepsilon, 1)_{L^2} = (\varepsilon, |\cdot|^2)_{L^2} = 0$, $\varepsilon = \alpha_\mu \Lambda Q_\mu + \hat{\varepsilon}$, and $\int_{\mathbb{R}^2} Q_\mu \mathcal{M}_\mu^y \hat{\varepsilon} dy = 0$, then

$$\begin{aligned} (\mathcal{L}_\mu^y \varepsilon, \mathcal{M}_\mu^y \varepsilon) &= (\mathcal{L}_\mu^y \hat{\varepsilon}, \mathcal{M}_\mu^y \hat{\varepsilon}) + O(\mu^3 |\log \mu|), \\ (\mathcal{M}_\mu^y \varepsilon, \varepsilon) &= (\mathcal{M}_\mu^y \hat{\varepsilon}, \hat{\varepsilon}) + O(\mu^2 |\log \mu|). \end{aligned}$$

PROOF. The proof is just a consequence of the identities $\mathcal{M}_\mu^y(\Lambda Q_\mu) = 2 - \frac{M}{2\pi} - \mu|y|^2$ and $\mathcal{L}_\mu^y(\Lambda Q_\mu) = -2\mu\Lambda Q_\mu$, and the orthogonality conditions on ε . Indeed, if we replace ε in $(\mathcal{M}_\mu^y\varepsilon, \varepsilon)$ by $\alpha_\mu\Lambda Q_\mu + \hat{\varepsilon}$, it follows that

$$\begin{aligned} (\mathcal{M}_\mu^y\varepsilon, \varepsilon) &= (\mathcal{M}_\mu^y\hat{\varepsilon}, \varepsilon) + \alpha_\mu(\mathcal{M}_\mu^y\Lambda Q_\mu, \varepsilon) \\ &= (\mathcal{M}_\mu^y\hat{\varepsilon}, \varepsilon) = (\mathcal{M}_\mu^y\hat{\varepsilon}, \hat{\varepsilon}) + \alpha_\mu(\mathcal{M}_\mu^y\hat{\varepsilon}, \Lambda Q_\mu) \\ &= (\mathcal{M}_\mu^y\hat{\varepsilon}, \hat{\varepsilon}) + \alpha_\mu\left(\hat{\varepsilon}, 2 - \frac{M}{2\pi} - \mu|\cdot|^2\right). \end{aligned}$$

The fact that $\int_{\mathbb{R}^2} \varepsilon dy = \int_{\mathbb{R}^2} \Lambda Q_\mu dy = 0$ implies

$$\int_{\mathbb{R}^2} \hat{\varepsilon} dy = 0.$$

Hence,

$$(\mathcal{M}_\mu^y\varepsilon, \varepsilon) = (\mathcal{M}_\mu^y\hat{\varepsilon}, \hat{\varepsilon}) - \alpha_\mu\mu(\hat{\varepsilon}, |\cdot|^2).$$

Moreover, since $\int_{\mathbb{R}^2} \varepsilon|y|^2 dy = 0$, it follows

$$\int_{\mathbb{R}^2} \hat{\varepsilon}|y|^2 dy = -\alpha_\mu(\Lambda Q_\mu, |\cdot|^2) = 2\alpha_\mu(Q_\mu, |\cdot|^2) = \alpha_\mu\frac{M}{\pi}|\log \mu| + O(1).$$

Then using that $\alpha_\mu \lesssim \sqrt{\mu}$, we deduce

$$(\mathcal{M}_\mu^y\varepsilon, \varepsilon) = (\mathcal{M}_\mu^y\hat{\varepsilon}, \hat{\varepsilon}) + O(\mu^2|\log \mu|).$$

Now we prove the second identity and we replace ε in $(\mathcal{L}_\mu^y\varepsilon, \mathcal{M}_\mu^y\varepsilon)$ by $\hat{\varepsilon} + \alpha_\mu\Lambda Q_\mu$:

$$\begin{aligned} (\mathcal{L}_\mu^y\varepsilon, \mathcal{M}_\mu^y\varepsilon) &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\varepsilon) + \alpha_\mu(\mathcal{L}_\mu^y\Lambda Q_\mu, \mathcal{M}_\mu^y\varepsilon) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\varepsilon) - 2\mu\alpha_\mu(\Lambda Q_\mu, \mathcal{M}_\mu^y\varepsilon) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\varepsilon) - 2\mu\alpha_\mu\left(2 - \frac{M}{2\pi} - \mu|\cdot|^2, \varepsilon\right) = (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\varepsilon) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\hat{\varepsilon}) + \alpha_\mu(\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\Lambda Q_\mu) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\hat{\varepsilon}) + \alpha_\mu(\mathcal{M}_\mu^y\hat{\varepsilon}, \mathcal{L}_\mu^y\Lambda Q_\mu) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\hat{\varepsilon}) - 2\mu\alpha_\mu(\mathcal{M}_\mu^y\hat{\varepsilon}, \Lambda Q_\mu) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\hat{\varepsilon}) + \alpha_\mu(\mathcal{M}_\mu^y\hat{\varepsilon}, \mathcal{L}_\mu^y\Lambda Q_\mu) \\ &= (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\hat{\varepsilon}) - 2\mu\alpha_\mu\left(\hat{\varepsilon}, 2 - \frac{M}{2\pi} - \mu|\cdot|^2\right). \end{aligned}$$

As before we use that $\int_{\mathbb{R}^2} \hat{\varepsilon} dy = 0$, $\int_{\mathbb{R}^2} \hat{\varepsilon}|y|^2 dy = \alpha_\mu\frac{M}{\pi}|\log \mu| + O(1)$, and $\alpha_\mu \lesssim \sqrt{\mu}$, which imply

$$(B.1) \quad (\mathcal{L}_\mu^y\varepsilon, \mathcal{M}_\mu^y\varepsilon) = (\mathcal{L}_\mu^y\hat{\varepsilon}, \mathcal{M}_\mu^y\hat{\varepsilon}) + O(\mu^3|\log \mu|).$$

This concludes the proof. \square

The following estimations have been proved already in [31] for $\varepsilon \in L^2_Q(\mathbb{R}^2)$, but since $Q_\mu < Q$ we get that $L^2_{Q_\mu}(\mathbb{R}^2) \hookrightarrow L^2_Q(\mathbb{R}^2)$, and we deduce easily from their proposition the following:

PROPOSITION B.2. *Let $\varepsilon \in L^2_{Q_\mu}(\mathbb{R}^2)$, then*

$$\|\nabla \phi_\varepsilon\|_{L^4} \leq C \|\varepsilon\|_{L^2_{Q_\mu}}.$$

If, in addition, $\int_{\mathbb{R}^2} \varepsilon dy = 0$, then

$$\|\phi_\varepsilon\|_{L^\infty} \leq C \|\varepsilon\|_{L^2_{Q_\mu}}, \quad \|\nabla \phi_\varepsilon\|_{L^2} \leq C \|\varepsilon\|_{L^2_{Q_\mu}}.$$

Appendix C Bounds on the Poisson Field

LEMMA C.1. *Let $\varepsilon \in L^2_{Q_\mu}(\mathbb{R}^2)$ such that*

$$\int_{\mathbb{R}^2} \varepsilon(y) dy = \int_{\mathbb{R}^2} y_i \varepsilon(y) dy = \int_{\mathbb{R}^2} |y|^2 \varepsilon(y) dy = 0,$$

then

$$\|\phi_\varepsilon\|_{L^2} \leq C \|\varepsilon\|_{L^2_{Q_\mu}},$$

with C uniform with respect to μ .

PROOF. For $|x| \leq 1$ we have

$$|\phi_\varepsilon(x)| \leq \int_{\mathbb{R}^2} |\log(|x-y|)| |\varepsilon(y)| dy \lesssim \|\varepsilon\|_{L^2_{Q_\mu}},$$

which implies

$$\int_{|x| \leq 1} |\phi_\varepsilon(x)|^2 dx \lesssim \|\varepsilon\|_{L^2_{Q_\mu}}^2.$$

For $|x| \geq 1$ we have

$$(C.1) \quad \left| \phi_\varepsilon(x) - \frac{\log(|x|)}{2\pi} \int_{\mathbb{R}^2} \varepsilon(y) dy \right| = \left| \int_{\mathbb{R}^2} \log\left(\frac{|x-y|}{|x|}\right) \varepsilon(y) dy \right|.$$

We separate the integral into two pieces. Indeed, we look first at $\{|y| \leq |x|\}$:

$$\begin{aligned} & \left| \int_{\{|y| \leq |x|\}} \log\left(\frac{|x-y|}{|x|}\right) \varepsilon(y) dy \right| \\ &= \frac{1}{2} \left| \int_{\{|y| \leq |x|\}} \log\left(\frac{|x-y|^2}{|x|^2}\right) \varepsilon(y) dy \right| \\ &= \left| \int_{\{|y| \leq |x|\}} \log\left(1 + \frac{|y|^2}{|x|^2} - \frac{2x \cdot y}{|x|^2}\right) \varepsilon(y) dy \right| \\ &= \left| \int_{\{|y| \leq |x|\}} \left(\frac{|y|^2}{|x|^2} - \frac{2x \cdot y}{|x|^2} + O\left(\frac{|y|^{2-\eta}}{|x|^{2-\eta}}\right) \right) \varepsilon(y) dy \right|, \end{aligned}$$

with η small. Hence, by using $\int_{\mathbb{R}^2} y_i \varepsilon(y) dy = \int_{\mathbb{R}^2} |y|^2 \varepsilon(y) dy = 0$ we deduce

$$\left| \int_{\{|y| \leq |x|\}} \log\left(\frac{|x-y|}{|x|}\right) \varepsilon(y) dy \right| \lesssim \int_{\{|y| \leq |x|\}} \frac{|y|^{2-\eta}}{|x|^{2-\eta}} |\varepsilon(y)| dy \lesssim \frac{\|\varepsilon\|_{L^2_{\mathcal{Q}\mu}}}{|x|^{2-\eta}}.$$

Now we look at the other part of the integral,

$$\left| \int_{\{|y| \geq |x|\}} \log\left(\frac{|x-y|}{|x|}\right) \varepsilon(y) dy \right|.$$

Since $|x-y|^2 \leq 2(|x|^2 + |y|^2)$ and $\log(1+t) \leq t^{1-\eta}$ with $\eta > 0$ sufficiently small, it follows

$$\begin{aligned} \left| \int_{\{|y| \geq |x|\}} \log\left(\frac{|x-y|}{|x|}\right) \varepsilon(y) dy \right| &= \frac{1}{2} \left| \int_{\{|y| \geq |x|\}} \log\left(\frac{|x-y|^2}{|x|^2}\right) \varepsilon(y) dy \right| \\ &\leq \frac{1}{2} \left| \int_{\{|y| \geq |x|\}} \log\left(2 + 2\frac{|y|^2}{|x|^2}\right) \varepsilon(y) dy \right| \\ (C.2) \quad &\leq \frac{1}{2} \left| \int_{\{|y| \geq |x|\}} \frac{|y|^{2-2\eta}}{|x|^{2-2\eta}} \varepsilon(y) dy \right|. \end{aligned}$$

Hence, we deduce

$$(C.3) \quad \int_{|x| \geq 1} |\phi_\varepsilon(x)|^2 dx \lesssim \|\varepsilon\|_{L^2_{\mathcal{Q}\mu}}^2,$$

which concludes the proof. \square

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