

Broadband Emission in Hybrid Organic-Inorganic Halides of Group 12 Metals

Rachel Roccanova,¹ Matthew Houck,¹ Aymen Yangui,¹ Dan Han,^{2,3,4} Hongliang Shi,⁵ Yuntao Wu,^{6,7} Daniel T. Glatzhofer,¹ Douglas R. Powell,¹ Shiyu Chen,³ Houcem Fourati,⁸ Alain Lussan,⁸ Kamel Boukheddaden,⁸ Mao-Hua Du,^{4} Bayrammurad Saparov^{1*}*

¹Department of Chemistry and Biochemistry, University of Oklahoma, 101 Stephenson Parkway, Norman, OK 73019, USA

²Key Laboratory of Polar Materials and Devices (Ministry of Education), East China Normal University, Shanghai 200241, China

³Department of Physics, East China Normal University, Shanghai 200241, China

⁴Materials Science and Technology Division, Oak Ridge National Laboratory, Oak Ridge, TN 37831, USA

⁵Key Laboratory of Micro-Nano Measurement-Manipulation and Physics (Ministry of Education), Department of Physics, Beihang University, Beijing 100191, China

⁶Scintillation Materials Research Center, University of Tennessee, Knoxville, TN 37996, USA

⁷Department of Materials Science and Engineering, University of Tennessee, Knoxville, TN 37996, USA

⁸Groupe d'Etudes de la Matière Condensée, UMR CNRS 8653-Université de Versailles Saint Quentin En Yvelines, Université Paris-Saclay, 45 Avenue des États-Unis, 78035 Versailles, France

***Corresponding Authors.** mhdu@ornl.gov , saparov@ou.edu

ABSTRACT. We report syntheses, crystal and electronic structures and characterization of three new hybrid organic-inorganic halides (R)ZnBr₃(DMSO), (R)₂CdBr₄·DMSO and (R)CdI₃(DMSO) (where (R) = C₆(CH₃)₅CH₂N(CH₃)₃, and DMSO = dimethyl sulfoxide). The compounds can be conveniently prepared as single crystals and bulk polycrystalline powders using a DMSO-methanol solvent system. Based on the single crystal X-ray diffraction results carried out at room temperature and 100 K, all compounds have zero-dimensional (0D) crystal structures featuring alternating layers of bulky organic cations and molecular inorganic anions based on a tetrahedral coordination around group 12 metal cations. The presence of discrete molecular building blocks in the 0D structures result in localized charges and tunable room temperature light emission including white-light for (R)ZnBr₃(DMSO), bluish-white light for (R)₂CdBr₄·DMSO, and green for (R)CdI₃(DMSO). The highest photoluminescence quantum yield (PLQY) value of 3.07% was measured for (R)ZnBr₃(DMSO), which emits cold white-light based on the calculated correlated color temperature (CCT) of 11044 K. All compounds exhibit fast photoluminescence lifetimes on the timescale of tens of ns, consistent with the fast luminescence decay observed in π -conjugated organic molecules. Temperature dependence photoluminescence study showed the appearance of additional peaks around 550 nm, resulting from the organic salt emission. Density functional theory calculations show that the incorporation of both the low-gap aromatic molecule R and the relatively electropositive Zn and Cd metals can lead to exciton localization at the aromatic molecular cations, which act as luminescence centers.

I. INTRODUCTION

Solution-processable hybrid organic-inorganic halides have been at the forefront of optical and electronic materials research in the past decade owing to their rich crystal chemistry and broadly tunable physical properties.¹⁻⁶ Among these, Pb- and Sn-based hybrid halide perovskites have been the focus of global attention due to their excellent performance in solar cells, which now have record power conversion efficiencies exceeding 22%.⁷ In addition, excellent light emission properties of halide perovskites have also been reported including efficient blue emission⁸⁻⁹ and broadband white-light emission.¹⁰⁻¹⁵ Inspired by these findings, further experimental studies show that non-perovskite hybrid organic-inorganic halides can also have remarkable light emission properties, such as in the case of $(\text{H}_2\text{DABCO})\text{Pb}_2\text{Cl}_6$ and $(\text{H}_3\text{O})(\text{Et}_2\text{-DABCO})_8\text{Pb}_{21}\text{Cl}_{59}$ compounds that demonstrate white-light emission with high color rendering index (CRI) values of 96 and 88, respectively.¹⁶

The light-emission properties of hybrid organic-inorganic perovskites are strongly dependent on their structural dimensionality with 3D perovskites behaving as typical all-inorganic multinary halides with small exciton binding energies (typically below 50 meV).^{1, 17-18} Conversely, low-dimensional hybrid halides demonstrate much stronger charge localization and exciton binding energies in excess of 500 meV, values that are comparable to that of tightly-bound Frenkel excitons in organic materials.^{1, 19} Such high exciton binding energies can accordingly result in room temperature stable excitons and intense excitonic emission. Therefore, preparation and characterization of low-dimensional organic-inorganic hybrid halides for light emission applications has been the focus of several studies published in recent years.²⁰⁻²²

Highest charge localization and exciton binding energies are often observed for zero-dimensional (0D) compounds,⁶ which feature molecular inorganic anions and organic cations in

their crystal structures. Indeed, numerous Cu-, Sn- and Sb-based 0D non-perovskite hybrid halides have recently been reported to exhibit remarkable light emission properties including photoluminescence quantum yields (PLQY) approaching unity²²⁻²⁵ and large Stokes shifts above 300 nm.²⁶ The large Stokes shifts are primarily attributed to the pronounced excited state structural distortions,^{23, 26} and therefore, are rarely observed for materials in the solid state due to the rigidity of solid structures which restricts significant structural reorganization. However, 0D hybrid halides featuring molecular organic and inorganic units are more amenable to structural reorganization, and hence, exhibit broadband photoluminescence (PL) spectra with large Stokes shifts.

Unlike the Pb-based and Sn-based hybrid halides, the optical and electronic properties of the hybrid halides of group 12 metals remain largely unexplored. Notwithstanding a few recent publications focusing on the light emission properties of group 12 metal hybrid halides,^{19-20, 27} most published articles focus on their intriguing structural properties including structural dimensionalities, impact of hydrogen bonding, molecular orientations of organic cations, etc.²⁸⁻³⁰ These group 12 metal-based hybrid halides show a remarkable structural flexibility in their tunable dimensionality^{19, 27} and incorporation of functional organic cations, such as chromophores and polymers,³⁰⁻³¹ in 2D perovskite structures both of which can be utilized to prepare functional materials for practical applications including solid state lighting applications. This fact is further evidenced by a recent discovery of broadband white-light emission in the 2D layered perovskite compound $(C_6H_{11}NH_3)_2CdBr_4$ originating from both organic and inorganic layers.²⁰

In this work, we report the use of the previously reported organic trimethyl(2,3,4,5,6-pentamethylbenzyl)ammonium (denoted as (R) hereafter) halides³² to prepare luminescent 0D hybrid (R)-M-X halides with M = Zn, Cd and X = Br, I. Single crystals of three new compounds within the studied systems were obtained using low-temperature solution reactions. Based on our

X-ray crystallography work, all compounds form 0D crystal structures featuring alternating layers of organic cations and inorganic anions that are based on tetrahedral coordination around metal cations. In (R)ZnBr₃(DMSO) and (R)CdI₃(DMSO), DMSO molecules directly coordinate to metal cations through M-O bonds yielding MX₃DMSO⁻ units, whereas (R)₂CdBr₄·DMSO features CdBr₄²⁻ anions and DMSO molecules that act as spacers in between organic and inorganic ions. No structural phase transitions are observed between room temperature and 100 K. The optical properties were investigated using steady state PL as a function of temperature, photoluminescence excitation (PLE), time-resolved photoluminescence (TRPL), PLQY, as well as density functional theory (DFT). The compounds are shown to demonstrate rare light emission from the organic component of the hybrid organic-inorganic halides including deep blue emission and broadband white-light emission with PLQY values up to 3.07%.

Table 1. Selected room temperature single crystal data collection and refinement parameters for the compounds prepared in this work.

Formula	(R)ZnBr ₃ (DMSO)	(R) ₂ CdBr ₄ ·DMSO	(R)CdI ₃ (DMSO)
Formula weight (g/mol)	603.59	950.90	791.59
Temperature (K)		298 (2)	
Radiation, wavelength (Å)		Mo K α , 0.71073	
Crystal system	Orthorhombic	Triclinic	Orthorhombic
Space group, Z	<i>P</i> 2 ₁ 2 ₁ 2, 4	<i>P</i> -1, 2	<i>P</i> 2 ₁ 2 ₁ 2, 4
<i>a</i> (Å)	9.116(2)	8.9894(6)	9.3959(6)
<i>b</i> (Å)	28.767(7)	14.3500(9)	29.5322(19)

c (Å)	8.892(2)	15.9925(10)	9.3829(6)
α , °	90	72.1342(11)	90
β , °	90	89.0598(12)	90
γ , °	90	89.2993(11)	90
Volume (Å ³)	2331.8(10)	1963.1(2)	2603.6(3)
Density (ρ_{calc}) (g/cm ³)	1.719	1.609	2.019
Absorption coefficient (μ) (mm ⁻¹)	6.289	4.706	4.481
$\theta_{\text{min}} - \theta_{\text{max}}$ (°)	2.29 – 26.41	1.34 – 25.70	1.38 – 29.67
Reflections collected	34552	44472	66620
Independent reflections	4544	7464	7362
R^a indices ($I > 2\sigma(I)$)	$R_1 = 0.0398$ $wR_2 = 0.0892$	$R_1 = 0.0324$ $wR_2 = 0.0690$	$R_I = 0.0370$ $wR_2 = 0.0977$
Goodness-of-fit on F^2	0.918	1.008	1.013
Largest diff. peak and hole (e ⁻ /Å ³)	0.626 and -0.739	0.816 and -0.626	1.320 and -0.926

^a $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$; $wR_2 = [\sum |w(F_o^2 - F_c^2)|^2 / \sum |w(F_o^2)|^2]^{1/2}$, where $w = 1/[\sigma^2 F_o^2 + (AP)^2 + BP]$, with $P = (F_o^2 + 2F_c^2)/3$ and weight coefficients A and B .

II. RESULTS AND DISCUSSION

Crystal structures. Room temperature and low temperature single crystal X-ray diffraction (SXRD) experiments for hybrid bromide and iodides of Zn and Cd metals are summarized in Figure 1, and Tables 1 and S1-S3. Based on the obtained results, there are no structural transitions

between room temperature and 100 K for all reported compounds; instead, normal contraction of the unit cell volume is observed upon cooling (Tables 1 and S1). As expected, the use of bulky non-linear organic cation results in 0D crystal structures for (R)ZnBr₃(DMSO), (R)₂CdBr₄·DMSO and (R)CdI₃(DMSO) featuring alternating layers of isolated inorganic anions and organic cations. In all cases, DMSO solvent molecules are incorporated into the crystal structures of the resultant compounds. In the case of (R)₂CdBr₄·DMSO, the coordination environment around Cd²⁺ only consists of bromide anions giving CdBr₄²⁻ anions and DMSO molecules acting as spacers in between organic and inorganic ions. However, in (R)ZnBr₃(DMSO) and (R)CdI₃(DMSO), DMSO molecules and halide anions form the coordination environment around metal cations through direct M-O bonds yielding MX₃DMSO⁻ units (Figure 2). Notice that the formula for (R)₂CdBr₄·DMSO is written differently compared to (R)ZnBr₃(DMSO) and (R)CdI₃(DMSO) to emphasize direct coordination of DMSO molecules to the metals in the latter. The measured M-O bond distances of $d_{\text{Zn-O}} = 2.024(5)$ Å and $d_{\text{Cd-O}} = 2.247(7)$ Å in (R)ZnBr₃(DMSO) and (R)CdI₃(DMSO) (Table S2), respectively, follow the expected trend based on the sizes of Zn²⁺ and Cd²⁺ ions.³³ These values are slightly longer than the literature-reported Zn-O and Cd-O bond distances in oxyhalides such as Zn_{0.902}SbBrO₂ ($d_{\text{Zn-O}} = 1.897$ Å) and BiCdIO₂ ($d_{\text{Cd/Bi-O}} = 2.239$ Å).³⁴ On the other hand, the metal halide bond distances are in the $d_{\text{Zn-Br}} = 2.38\text{-}2.40$ Å range for (R)ZnBr₃(DMSO), $d_{\text{Cd-Br}} = 2.58\text{-}2.61$ Å for (R)₂CdBr₄·DMSO, and $d_{\text{Cd-I}} = 2.72\text{-}2.73$ Å for (R)CdI₃(DMSO). These values compare well with that reported for other hybrid organic-inorganic halides of Zn and Cd, including [4,4'-H₂bipy][ZnBr₄] ($d_{\text{Zn-Br}} = 2.39\text{-}2.43$ Å),²⁹ (CH₃NH₃)₂CdBr₄ ($d_{\text{Cd-Br}} = 2.57\text{-}2.59$ Å), and (CH₃NH₃)₂CdI₄ ($d_{\text{Cd-I}} = 2.72\text{-}2.75$ Å).¹⁹ In all compounds, the tetrahedral coordination around the metal cations are distorted with bond angles ranging from 98.97(19)° to 118.37(3)° in (R)CdI₃(DMSO) (Table S2). It should be noted that the presence of

DMSO in the coordination sphere of Cd^{2+} undoubtedly contributes to the observed distortion in $(\text{R})\text{CdI}_3(\text{DMSO})$, however, tetrahedral distortions are also observed for cadmium halides with no coordinating solvent molecules including $(\text{R})_2\text{CdBr}_4 \cdot \text{DMSO}$, which features CdBr_4^{2-} units, and in literature, for $(\text{CH}_3\text{NH}_3)_2\text{CdBr}_4$ and $(\text{CH}_3\text{NH}_3)_2\text{CdI}_4$.¹⁹

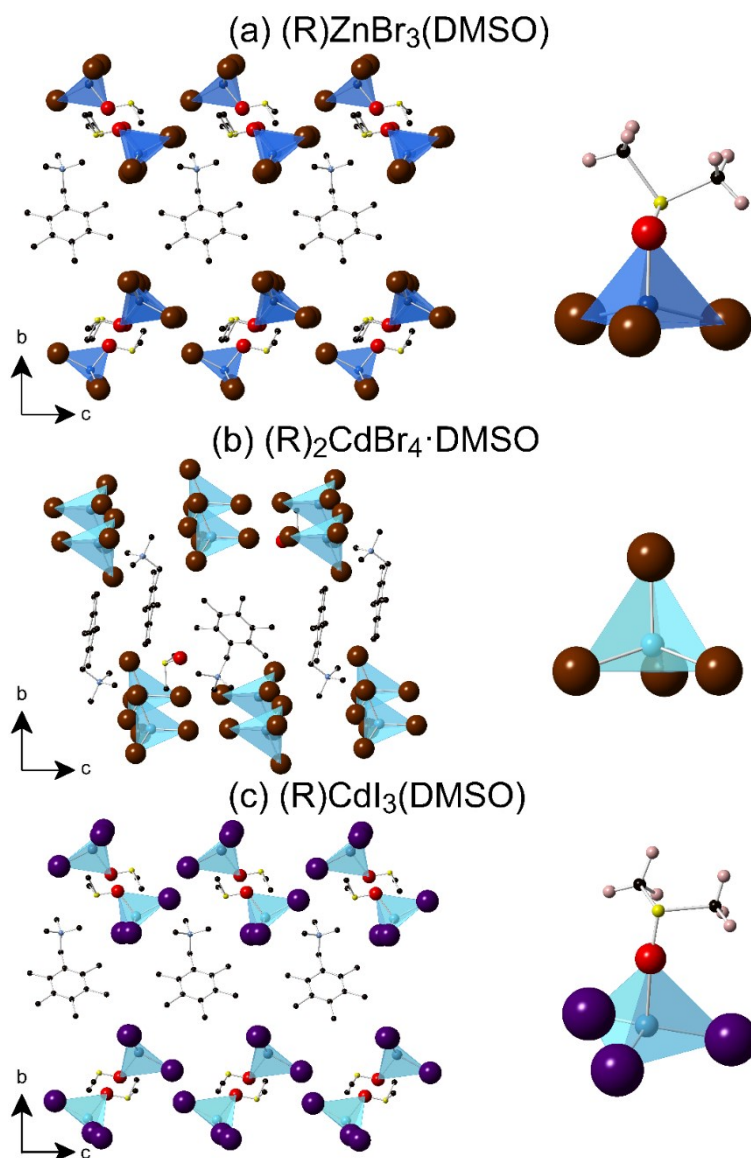


Figure 1. Crystal structures and closeup views of the coordination polyhedra in (a) $(\text{R})\text{ZnBr}_3(\text{DMSO})$, (b) $(\text{R})_2\text{CdBr}_4 \cdot \text{DMSO}$ and (c) $(\text{R})\text{CdI}_3(\text{DMSO})$. Blue and cyan tetrahedra represent coordination environments around Zn and Cd, respectively. Burgundy, purple, red,

yellow, black, and light blue spheres represent Br, I, O, S, C, and N, respectively. For clarity, hydrogen atoms were omitted and only a fraction of organic cations and solvent molecules are shown.

Stability studies. Based on the results of our PXRD measurements, all compounds form phase pure products (Figure 2). The PXRD patterns in all cases have several characteristic low angle peaks indicative of large unit cell parameters stemming from the use of the bulky organic cation. Poor stability in air is an often-cited deficiency of some of the most well-known hybrid organic-inorganic halides including the state-of-the-art photovoltaic material $\text{CH}_3\text{NH}_3\text{PbI}_3$.³⁵⁻³⁶ To test air stability of our compounds, powdered samples were allowed to sit undisturbed under ambient conditions over a period of one month. Periodic PXRD scans taken on the samples exposed to ambient air (Figure S5) indicate that all samples are stable for at least several days. Noticeable changes are recorded for $(\text{R})\text{CdI}_3(\text{DMSO})$, which shows lowered peak intensities and peak broadening after 2 weeks (Figure S5), indicative of reduced crystallinity. These results are consistent with the trend we recently observed for the $(\text{CH}_3\text{NH}_3)_2\text{CdX}_4$ ($\text{X} = \text{Cl}, \text{Br}, \text{I}$) series, in which air stability is markedly worse for the iodide compared to the chloride analog.¹⁹ Furthermore, extensive studies on halide perovskites also suggest reduced air stabilities going from chlorides to iodides.³⁶

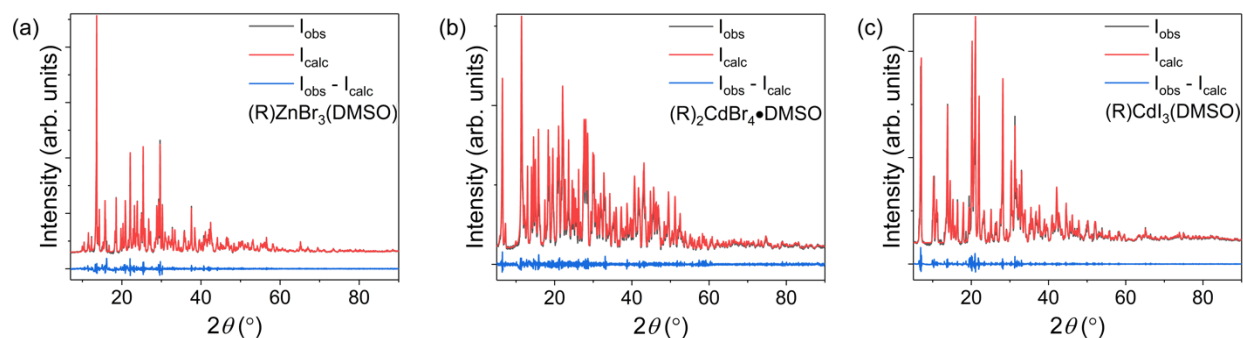


Figure 2. Powder X-ray diffraction (PXRD) patterns (black lines) for (a) (R)ZnBr₃(DMSO), (b) (R)₂CdBr₄•DMSO and (c) (R)CdI₃(DMSO). The Pawley fits and difference plots are shown as red and blue curves, respectively.

Optical properties. The highly tunable light emission properties in this series studied through room temperature optical absorption, PL and PLE measurements are summarized in Figure 4 and Table 2. Optical absorption spectra shows a sharp absorption peak at ~ 290 nm accompanied with a shoulder at 290 nm, 320 nm and 360 nm for (R)ZnBr₃(DMSO), (R)₂CdBr₄•DMSO and (R)CdI₃(DMSO), respectively. These optical absorption bands were attributed to the organic molecule. Moreover, PLE spectra shows the presence of a sharp peak at 386 nm, 399 nm, and 445 nm, for (R)ZnBr₃(DMSO), (R)₂CdBr₄•DMSO and (R)CdI₃(DMSO), respectively. These peaks were assigned to the excitonic absorption.³⁷ Similar to the earlier reports on Cd-based and Pb-based hybrid halide perovskites,^{12, 19-21, 38-40} all compounds studied in this work exhibit broadband emission spectra with full width at half maximum (FWHM) values ranging from 162 nm for (R)CdI₃(DMSO) to 188 nm for (R)ZnBr₃(DMSO). Importantly, bright broadband white-light emission is observed for (R)ZnBr₃(DMSO). The corresponding Commission Internationale de l'Eclairage (CIE) Color Coordinates (x, y) of (0.26, 0.30) for (R)ZnBr₃(DMSO) are close to the

white point (0.33, 0.33) coordinates. $(R)_2CdBr_4 \cdot DMSO$ emits bluish-white light with CIE coordinates of (0.24, 0.35), whereas $(R)CdI_3(DMSO)$ emits green light (0.29, 0.57). For the compounds with emission on the Planckian locus, the calculated correlated color temperatures (CCT) using the approach proposed by Hernández-Andrés et al.⁴¹ are 11044 K for $(R)ZnBr_3(DMSO)$ and 10384 K for $(R)_2CdBr_4 \cdot DMSO$.⁴¹ Such high values suggest that these compounds emit “cool” bluish-white and white colors similar to the cold white-light emitting Pb-based halide perovskites.³⁹

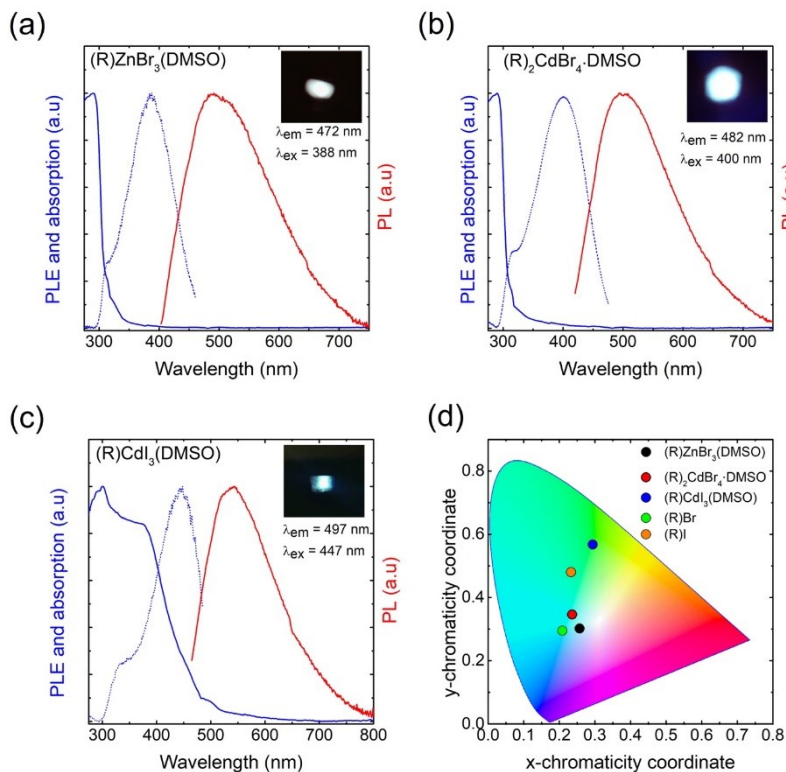


Figure 3. (a-c) Room temperature optical absorption (blue solid lines) PLE (blue dotted lines) and PL (red solid lines) spectra for the compounds studied in this work. The insets show the bright emission from the powder samples under irradiation with their respective maximum excitation

wavelengths. (d) CIE 1931 chromaticity diagram showing the emission colors of the compounds and the organic salts.

Table 2. Summary of the optical properties of the organic-inorganic hybrid compounds (R)ZnBr₃(DMSO), (R)₂CdBr₄·DMSO and (R)CdI₃(DMSO), and the corresponding organic salts (R)Br and (R)I.

Compound	PLQY	FWHM	PLE _{max}	PL _{max}	Stokes	CIE	
	(%)	(nm)	(nm)	(nm)	shift (nm)	coordinates	
(R)ZnBr ₃ (DMSO)	3.07	188	386	491	105	0.26	0.30
(R) ₂ CdBr ₄ ·DMSO	0.32	164	399	501	102	0.24	0.35
(R)CdI ₃ (DMSO)	0.27	162	445	515	70	0.29	0.57
(R)Br	2.44	153	378	462	84	0.21	0.29
(R)I	0.42	168	416	514	98	0.23	0.48

The room-temperature PL lifetime profiles of the compounds prepared in this work are plotted in Figure 4. All the profiles monitoring emission at the wavelength of maximum intensity can be well fitted by two-exponential functions. The first lifetime component of 1-2 ns for all compounds is due to the instrumental response because of the duration of the light pulse (2 ns). The measured PL lifetimes of these compounds range from 9 to 21.1 ns (Figure 4). In the literature, a spread in PL lifetimes of Pb-based hybrid halides are reported ranging from subnanosecond up to 54.1 ns,^{16, 21, 39} attributed to self-trapped excitonic states (STEs), and for some 0D Pb-free hybrid halides phosphorescent lifetimes in microseconds are reported.^{23, 26, 42}

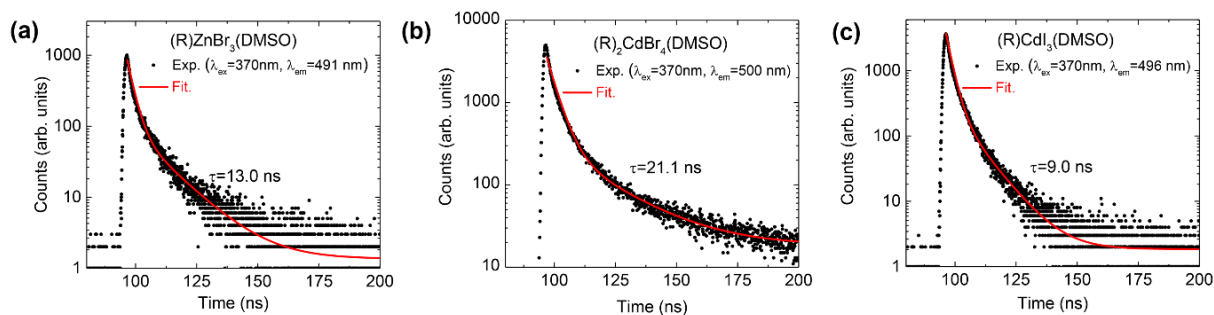


Figure 4. Photoluminescence lifetime profiles (black solid circles) and fitting curves (red lines) for (a) (R)ZnBr₃(DMSO), (b) (R)₂CdBr₄(DMSO), and (c) (R)CdI₃(DMSO). The insets show the excitation and emission wavelengths, and the derived lifetime constants.

In hybrid metal halide perovskites, tunable broadband emission has been attributed to structural distortions, in particular to the distortions of the MX₆ octahedra.^{12, 20, 43} According to these studies, structural distortions directly impact band gaps and the spread in the structural distortions, therefore, broadening the emission spectra. Although the compounds studied in this work are not perovskites and are not based on octahedral inorganic frameworks, tetrahedral distortions around metal cations are clearly evident for each compound (Table S2). However, in the present case, we cannot establish a clear correlation between tetrahedral distortions and the measured FWHM values due to the differing crystal structures and coordination environments around the metal cations in this family. Nevertheless, the structural distortions observed in our compounds could, in principle, aid in the generation of more transient photo-excited STEs similar to those in Pb-halide perovskites.^{39, 44-45} For 0D hybrid halides of Sn and Sb, it has been shown that the significant structural distortions of the excited state also leads to large Stokes shifts of 200-350 nm.^{23, 26} Although we observe smaller Stokes shifts ranging from 70 nm for (R)CdI₃(DMSO) to 105 nm for (R)ZnBr₃(DMSO), these values are still much higher than that for higher dimensional (2D and 3D)

hybrid halide perovskites.⁴⁶ The impact of distortions of the inorganic anions would be greatest if light emission originates from the inorganic molecules, i.e., if the excitons are localized on the inorganic units. Unlike the heavily studied Sn and Pb halide systems, in the present case we employed group 12 metals (Zn and Cd), which have lower electronegativities,⁴⁷ and hence higher lying conduction bands for the inorganic anions. Consequently, excitons in (R)ZnBr₃(DMSO), (R)₂CdBr₄·DMSO, and (R)CdI₃(DMSO) could be localized on the organic molecules, and the relatively smaller Stokes shifts in our compounds (between 70 and 105 nm) could be a direct consequence of this fact. Recently, broadband white-light emission in the 2D layered perovskite compound (C₆H₁₁NH₃)₂CdBr₄ was reported and attributed to both organic and inorganic layers emissions.²⁰ Another recent report of yellowish white-light emission in the 1D hybrid perovskite namely (C₉H₁₀N₂)PbCl₄ attributed its emission properties to a resonant energy transfer mechanism from the inorganic PbCl₄ to the organic molecules.⁴⁰ To further study the origin of the broadband light emission in our compounds, we have also carried out photoluminescence measurements on the precursor organic salts (R)Br and (R)I, which are provided in Figure 6. The corresponding binary inorganic halides do not luminesce at room temperature due to thermal quenching effects.^{19, 42, 48-51} As shown in Table 2, (R)Br emits light blue with CIE coordinates of (0.21, 0.29), whereas (R)I emits light green with CIE coordinates of (0.23, 0.48). Importantly, the PL and PLE spectra of (R)ZnBr₃(DMSO), (R)₂CdBr₄·DMSO, and (R)CdI₃(DMSO) are largely similar to that of (R)Br and (R)I (Figure 6), suggesting that the emission in our hybrid halides originates from the organic molecules. The red shift of the PL peaks of the hybrid halides compared to the organic salts, can be attributed to the difference in the crystal structures of the precursor organic salts and hybrid halides, and the resultant higher charge localization in hybrid halides.

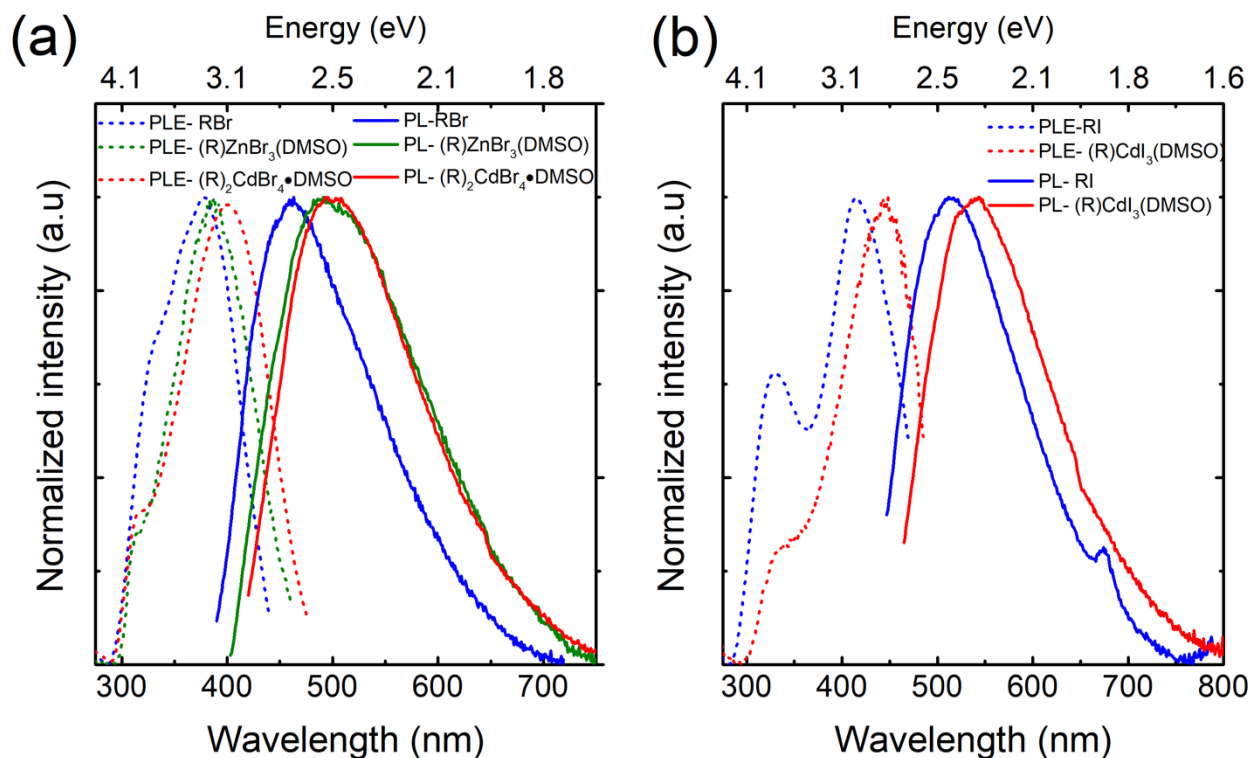


Figure 5. Room temperature PLE (dashed lines) and PL spectra (solid lines) of the (a) (R)Br organic salt and (R)Br-based compounds, and (b) (R)I organic salt and (R)I-based compound. Spectra are normalized for clarity.

The measured room temperature PLQY values range from 0.27% for (R)CdI₃(DMSO) to 3.07% for (R)ZnBr₃(DMSO), which are comparable or higher than most of the reported broadband light emitters based on lead halide perovskites.^{10, 21, 52-53} The precursor salts (R)Br and (R)I demonstrate PLQY values of 2.44 and 0.42%, respectively. Therefore, the emission efficiency is increased upon incorporation of the organic molecule into (R)ZnBr₃(DMSO) but is decreased in (R)₂CdBr₄•DMSO and (R)CdI₃(DMSO). A direct comparison with PLQY values for other hybrid halides of group 12 metals cannot be made due to the apparent lack of in-depth investigations of their light emission properties. However, our recent study on (CH₃NH₃)₂CdX₄ (X = Cl, Br, I)¹⁹

reports that these compounds do not luminesce at room temperature, which seems to confirm the validity of the strategy of using bulky luminescent organic cations to create 0D structures with molecular organic and inorganic ions in order to prepare room-temperature luminescent hybrid halides as employed in this work.

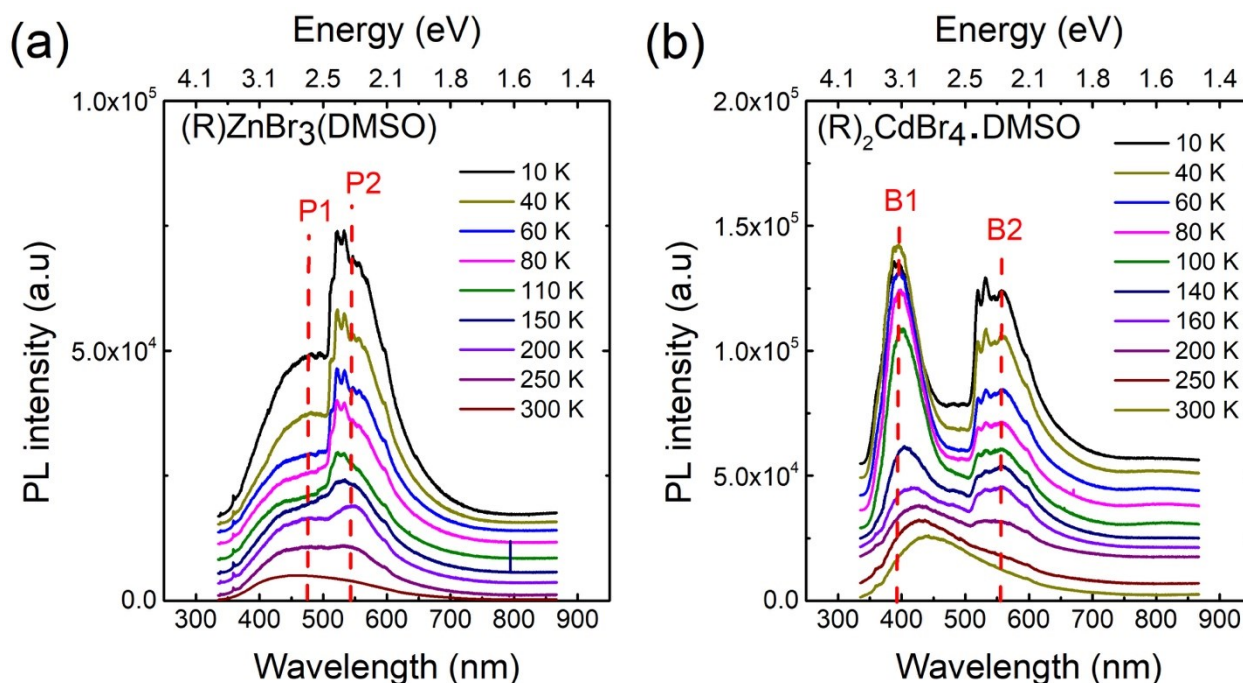


Figure 6. Temperature dependence of PL spectra measured for (a) $(R)ZnBr_3(DMSO)$ and (b) $(R)_2CdBr_4 \cdot DMSO$ under 325 nm irradiation.

To get further information about the thermally-activated processes and the origin of the broadband emission, we measured the temperature dependence PL spectra for two of the more stable compounds in this family, $(R)ZnBr_3(DMSO)$ and $(R)_2CdBr_4 \cdot DMSO$, under 325 nm laser excitation (Figures 6a and 6b). Significant temperature-dependent changes were revealed. For both

Zn-based and Cd-based compounds, the room temperature PL spectrum shows the presence of only one broad band, denoted as P1 and B1 in Figures 6a and 6b for (R)ZnBr₃(DMSO) and (R)₂CdBr₄·DMSO, respectively. Upon cooling, additional PL peaks (denoted as P2 for (R)ZnBr₃(DMSO) and B2 for (R)₂CdBr₄·DMSO emerge at around 550 nm. These peaks split into multiple sub-bands below 200 K. The splitting of PL peaks is often observed in organic-inorganic hybrid materials and may originate from the presence of free- and bi-excitons,⁵⁴ free-excitons and excitons-phonon interactions,⁵⁵ free and bound excitons,⁵ or could even result from the presence of structural phase transitions often observed in these hybrid materials. The latter possibility is ruled out in the present case by our room temperature and 100 K single crystal X-ray diffraction measurements suggesting the absence of structural phase transitions in the 100-300 K range.⁵⁶⁻⁵⁸ In the present case, very similar thermal behavior was observed for (R)ZnBr₃(DMSO) and for (R)₂CdBr₄·DMSO, which suggests that the appearance of additional low temperature PL peaks could result from the emission of the organic cation. Based on the temperature dependence PL data shown in Figure 6b, we plotted in Figure 8 the thermal dependence of the integrated intensity, the FWHM, and the position of B1 PL band of (R)₂CdBr₄·DMSO. The quenching of the B1 peak could be described by the following Arrhenius-type model,⁵⁵

$$I_{\text{PL}} = \frac{I_0}{1 + a \exp\left(\frac{-E_a}{k_B T}\right)}, \quad (1)$$

where I_0 is the low-temperature PL intensity, k_B is the Boltzmann constant, T is the temperature, a is the ratio between the radiative and the nonradiative decay rates, and E_a is the activation energy. The best fit of the experimental data presented in Figure 6a, gives $I_0 = 9.2 \times 10^4 \pm 830$, $a = 76 \pm 8$ and $E_a = 48 \pm 4$ meV. This activation energy presents a good agreement with those reported in similar organic-inorganic hybrid materials such as (C₆H₁₁NH₃)₂PbBr₄¹² and (C₆H₁₁NH₃)₂CdBr₄.²⁰ Moreover, the remarkable broadening of B1 PL

peak as a function of temperature (Figure 7b) can be attributed to exciton-phonons interaction, described within the following law,⁵⁹

$$\Gamma(T) = \Gamma_0 + \Gamma_{AC} \times T + \frac{\Gamma_{LO}}{\exp\left(\frac{E_{LO}}{k_B T}\right) - 1}, \quad (2)$$

where the first term represents the natural line width at 0 K, the second term is the broadening induced by acoustic phonons, and the third term corresponds to the contribution of optical phonons to the peak broadening. The best fit parameters yield $\Gamma_0 = 440 \pm 38$ meV, $\Gamma_{AC} = 0.76 \pm 0.03$ meV.K⁻¹, $\Gamma_{LO} = 22 \pm 4$ meV, and $E_{LO} = 16 \pm 2$ meV. Once again, these values are in good agreement with those reported in organic-inorganic hybrid materials such as (C₆H₁₁NH₃)₂CdBr₄²⁰ as those of the inorganic perovskite CsCdBr₃,⁶⁰ and CdBr₂ crystals.⁶¹ On the other hand, the B1 PL peak of (R)₂CdBr₄·DMSO shows a continuous red shift as a function of temperature (Figure 7c), which is usually the case for excitonic PL peaks accompanied with a clear regime change around 150 K. This is attributed to the impact of the appearance of the B2 PL peak at low temperature, which may cause a competition in the recombination process between B1 and B2 PL bands.

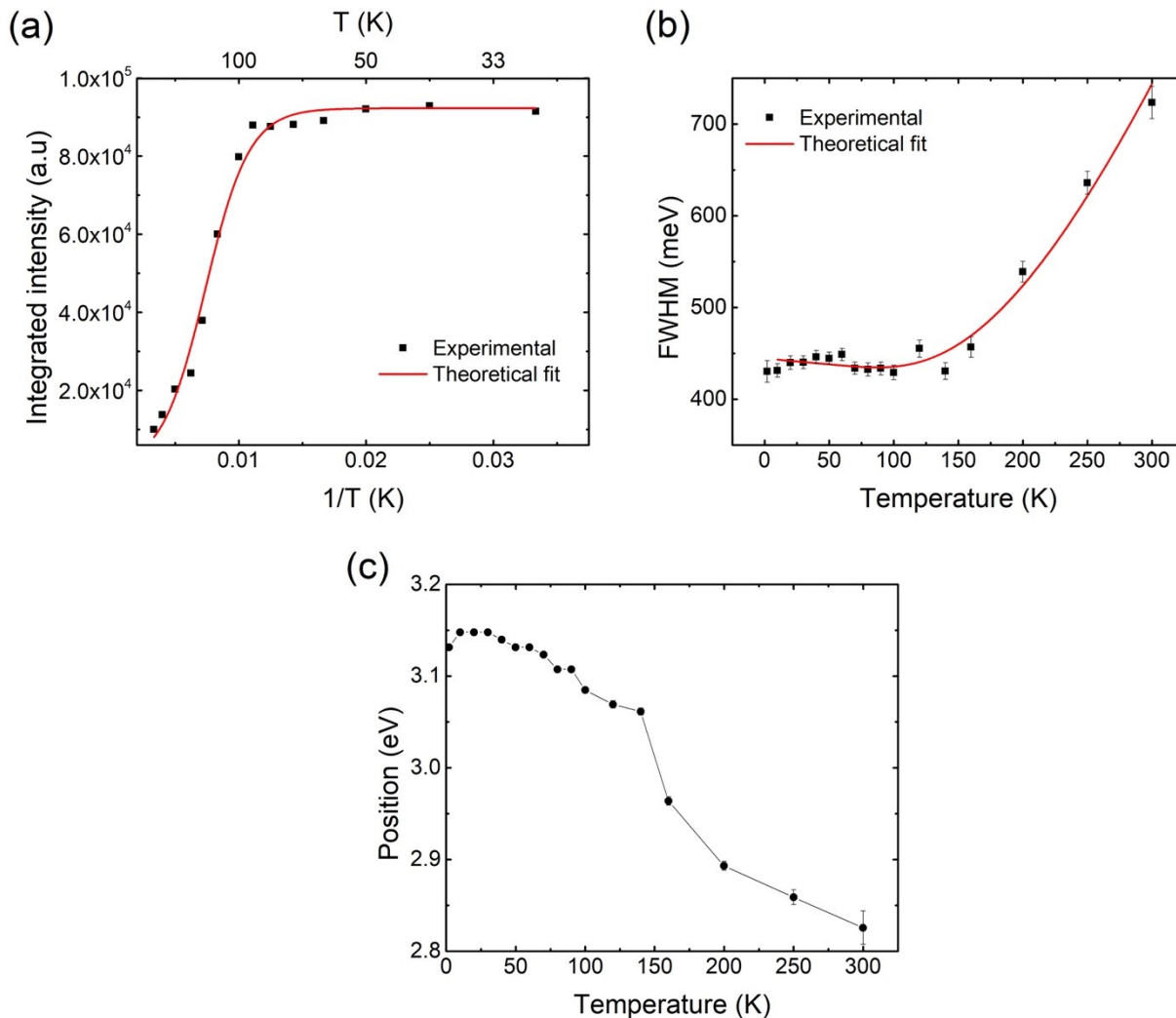


Figure 7. Temperature dependence of (a) the integrated intensity, (b) FWHM, and (c) the position of B1 peak of $(R)_2CdBr_4 \cdot DMSO$ deduced from the PL spectra showed in Figure 5. The red lines (in (a) and (b)) are the best fit of the experimental data, according to eqs (1) and (2), respectively.

First-principle calculations. Electronic band structure calculations were performed in order to further investigate the optical properties of the compounds prepared in this work. Figure 8 displays the electronic band structures and DOS plots for $(R)ZnBr_3(DMSO)$, $(R)_2CdBr_4 \cdot DMSO$ and $(R)CdI_3(DMSO)$, which show the PBE-calculated band gaps of 3.82 eV, 3.61 eV and 3.39 eV,

respectively. Although the PBE band gaps are typically underestimated and do not take into account the exciton binding energies, the trend of the calculated band gaps is consistent with that of the exciton excitation energies, which are 3.21 eV for (R)ZnBr₃(DMSO), 3.11 eV for (R)₂CdBr₄·DMSO, and 2.79 eV for (R)CdI₃(DMSO). All three compounds show small dispersion for the electronic bands near the band gap, indicating that these bands are made up of highly localized electronic states. The band gaps of these compounds are slightly indirect. The VBM and the CBM are located at the Γ and the Y points, respectively, for (R)ZnBr₃(DMSO), at the N and the Γ points for (R)₂CdBr₄·DMSO, and at the Z and the Γ points for (R)CdI₃(DMSO). In luminescent 0D hybrid Sn and Pb based halides, the inorganic Sn/Pb-halide clusters are usually the luminescent centers due to their relatively small energy gaps of the inorganic substructures compared to those of organic molecules.^{23, 26, 62} The low energy gaps of the inorganic substructures in such cases are ensured by the relatively small electronegativity difference between Sn/Pb and halogen elements, which results in valence bands dominated by halogen orbitals and conduction bands primarily made of metal orbitals. In principle, the band alignment in hybrid halides can be altered by (1) combining pairs of metal and halogen elements with large electronegativity difference to accommodate organic molecules' frontier orbitals, or (2) by utilizing low-gap aromatic molecular cations. The hybrid metal halides studied in this work feature electropositive metals Zn and Cd³⁶ and an aromatic organic cation, yielding unusual band alignments at the organic-inorganic interface.¹ The DOS plots in Figure 8 show that (R)ZnBr₃(DMSO) exhibits the type-I band alignment with its valence band maximum (VBM) and conduction band minimum (CBM) derived from the bonding and the antibonding π orbitals of the aromatic molecule cations, while (R)₂CdBr₄·DMSO and (R)CdI₃(DMSO) display the type-II band alignment with the VBM and the CBM containing π -orbitals of the organic cations and the metal-s orbitals of the inorganic

anions, respectively. It is, thus, expected that the optical emission from (R)ZnBr₃(DMSO) is due to the excitons localized on the organic molecular cations. Although the ground-state band structures of (R)₂CdBr₄·DMSO and (R)CdI₃(DMSO) show the type-II band alignment, the lowest unoccupied molecular orbital of the molecular cation is close to the CBM as shown in Figure 8d and 8f. The strong Coulomb binding could localize the exciton at the organic cation, in analogy to the case of (Ph₄P)₂SbCl₅, in which the exciton is localized at the SbCl₅ cluster despite that the ground-state band structure shows the type-II band alignment at the organic-inorganic interface.²³ The fast decay of the luminescence in the three hybrid metal halides as shown in Figure 4 (9 ns - 21 ns) are consistent with the fast emission from the π -conjugated organic molecules (decay times ranging from a few ns to a few tens of ns);⁶³ thus, supporting the attribution of the exciton emission in these hybrid halides to the organic molecules. In comparison, an exciton that is localized at the inorganic metal halide cluster in hybrid metal halides usually has the lifetime on the order of μ s.^{23, 26, 64} Figure 9 shows the hole and the electron wavefunctions of an exciton in (R)ZnBr₃(DMSO), which is optimized by the PBE0 hybrid functional calculation. It can be seen that the exciton is indeed localized on the π orbitals of the C₆ ring in a molecular cation, which, thus, should act as a luminescent center.

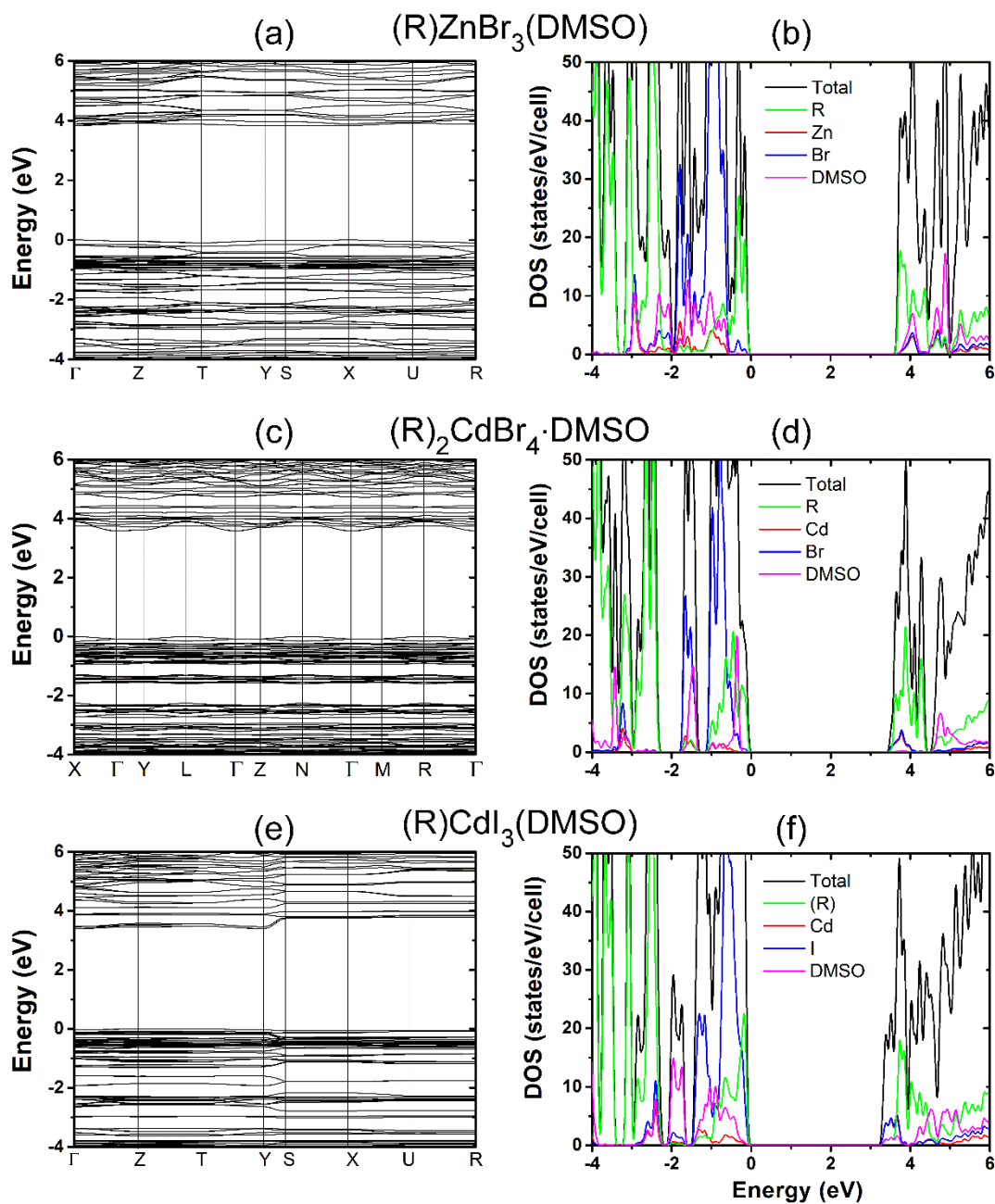


Figure 9. Band structures and density of states (DOS) plots for (a)-(b) $(R)ZnBr_3(DMSO)$, (c)-(d) $(R)_2CdBr_4 \cdot DMSO$ and (e)-(f) $(R)CdI_3(DMSO)$.

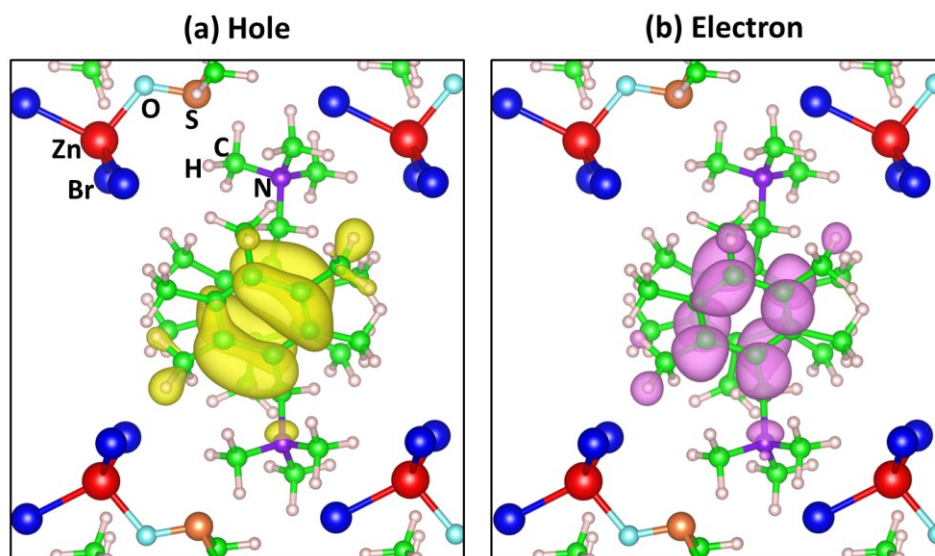


Figure 9. Partial density contours of the hole (a) and the electron (b) in a relaxed exciton in (R)ZnBr₃(DMSO).

III. CONCLUSION

In summary, we report the preparation, crystal and electronic structures, and optical properties of three new hybrid organic-inorganic halides of group 12 metals. Due to the presence of the bulky organic (R)⁺ cation, all compounds form 0D crystal structures featuring alternating layers of isolated inorganic anions based on metal tetrahedra and organic cations. The separation of the crystal structure into distinct inorganic and organic molecular units leads to flat bands near the band gap indicative of the highly localized electronic states. Furthermore, the 0D molecular structures allow for the localization of the excitons on either anionic or cationic molecular units depending on their chemistries. Thus, based on our computational work, the excitons are localized on the organic molecular cation in (R)ZnBr₃(DMSO), suggesting a rare example of light emission from the organic component of hybrid organic-inorganic materials.^{6, 12, 21, 65} The computational results are supported by the PL measurements including the measured fast decay of the

luminescence (9 ns - 21 ns), which is consistent with the fast emission from the π -conjugated organic molecules.

The compounds exhibit broadband emission ranging from bluish-white to white to green light depending on the choice of the halide and metal cations. The luminescence properties derive from the intricacies of their crystal and electronic structures (e.g., packing of the organic cations). Temperature dependence of PL measurements showed the emergence of additional peaks around 550 nm at lower temperatures, attributed to the organic salt emission. For $(R)_2CdBr_4 \cdot DMSO$ an excitonic activation energy $E_a \sim 50$ meV was estimated from the thermal dependence of the intensity of B1 PL peak showing a good agreement with Cd-based organic-inorganic hybrid materials.

The 0D crystal structures designed in this work demonstrate a luminescence efficiency with PLQY values up to 3.07% for the white-light emitting $(R)ZnBr_3(DMSO)$, which also happens to be the most stable (thermally and in air) member of the family. The combinations of the organic cation and inorganic components used in this work, therefore, are responsible for the preparation of room-temperature luminescent materials that have comparable or better PLQY values to other Pb-based and Cd-based hybrid halides reported in literature. Importantly, this work paves the way to band alignment engineering of organic and inorganic components in hybrid organic-inorganic halides to controllably create type I and type II band alignments, and to fabricate materials where the emission originates from organic, inorganic or both components. The fast emission from the hydrogen-rich organic cations, observed for the compounds prepared in this work, may enable the development of fast scintillators that are capable of detecting fast neutrons via proton *recoil*.

IV. EXPERIMENTAL METHODS

Starting materials. Chemicals utilized in this study were either used as purchased or synthesized from the starting materials listed: (i) zinc bromide, 99.999%, Sigma; (ii) zinc iodide, 99.995%, Alfa Aesar; (iii) cadmium bromide, 98% Alfa Aesar; (iv) cadmium iodide, 99.999%, Alfa Aesar; (v) dimethyl sulfoxide, ACS grade, Fisher; (vi) methanol, ACS reagent grade, Pharmco-AAPER. (vii) pentamethylbenzene, Aldrich; (viii) paraformaldehyde, 95%, Aldrich; (ix) hydrobromic acid, 48%, Sigma Aldrich; (x) cetyltrimethyl ammonium bromide, Sigma; (xi) acetic acid, ACS reagent grade, 99.8%, Sigma Aldrich; (xii) o-phosphoric acid, ACS Certified, 85%, Fisher; (xiii) methyl iodide, Aldrich; (xiv) trimethylamine hydrochloride, 98% Aldrich.

Syntheses of hybrid organic inorganic compounds. For the syntheses of hybrid organic-inorganic halides of group 12 metals, various solvents including N,N-dimethylformamide (DMF), dimethyl sulfoxide (DMSO), and methanol were tried. Among these, DMSO was selected as the most suitable solvent due to the highest solubility of the precursor organic salts trimethyl(2,3,4,5,6-pentamethylbenzyl)ammonium halides (R)X (X = Br, I). Details of the syntheses and characterizations of the organic salts are described in the Supplementary Information (see Figures S1-S4). It was determined that single crystals of target compounds can be grown in 20 mL vials at 65 °C by separately solvating (R)X (X = Br, I) in 2 mL of DMSO and MX₂ (M = Zn, Cd) in 1 mL of methanol, then mixing the stoichiometric 1:1 molar ratio solutions together.

Bromide-based compounds: (R)ZnBr₃(DMSO) crystals were grown through slow evaporation of their stoichiometric solutions as described above. (R)₂CdBr₄·DMSO was first obtained by reacting (R)Br and CdBr₂ in a nonstoichiometric 1:1 molar ratio. A subsequent stoichiometric reaction with 2:1 molar ratio of (R)Br:CdBr₂ also afforded (R)₂CdBr₄·DMSO. In all cases, colorless block crystals measuring up to 3 mm formed over a 3 week crystallization period.

Iodide-based compound: (R)CdI₃(DMSO) crystals were grown through slow evaporation of their stoichiometric solutions as light-yellow blocks measuring up to 2 mm in size over a 3 week crystallization period. Syntheses of Zn analogues, (R)ZnI₃ and (R)₂ZnI₄, were tried in different loading ratios of DMF, DMSO, and methanol solutions, but proved unsuccessful.

Single Crystal X-ray Diffraction. The X-ray intensity data for all compounds were collected on a Bruker Apex CCD diffractometer with a graphite monochromated Mo-K α ($\lambda = 0.71073\text{\AA}$) radiation. For room temperature measurements (298(2) K), crystals with good clarity were selected from their mother liquor, cut to size, covered with Super Glue, affixed to a goniometer head, and allowed to dry. For low temperature measurements, crystals were selected from their mother liquor, placed in Dow Chemical vacuum grease, cut to size, and then cooled to 100(2) K. The crystal structures for all compounds were determined from a nonlinear least-squares fit. The data were corrected for absorption by the semi-empirical method based on equivalents and the structures were solved by direct methods using the SHELXTL program and refined by full matrix least-squares on F^2 by use of all reflections.⁶⁶⁻⁶⁷ All non-hydrogen atoms were refined with anisotropic displacement parameters, all occupancies were refined within two standard deviations, and all hydrogen atom positions were determined by geometry. Single crystal data collection and refinement parameters are summarized in Tables 1 and S1. Atomic coordinates, equivalent isotropic displacement parameters and selected interatomic distances and bond angles can be found in Tables S2 – S3. Full details of the crystal structures are summarized in the CIF (Crystallographic Information File) files have been deposited in the Cambridge Crystallographic Data Centre (CCDC) database and can be found under deposition numbers 1840612, 1840614-1840618.

Powder X-ray Diffraction. For powder X-ray diffraction (PXRD) measurements, solution-processed samples were dried overnight under vacuum and then ground into powdered forms.

PXRD measurements were carried out on a Rigaku MiniFlex600 system equipped with a D/tex detector using Ni-filtered Cu-K α radiation source. Typical PXRD scans were performed in the 5-90° (2 θ) range, with a step size of 0.02° to determine phase identity and purity. Data analysis was performed using Rigaku's PDXL2 software package. The baseline originating from the glass slides used to collect data was corrected using the embedded tool in the PDXL2 software, and the data were fitted using the decomposition method (also known as Pawley fitting) embedded in the PDXL2 software package.

Energy Dispersive X-ray (EDX) Spectroscopy Measurements. EDX measurements were conducted to confirm the heavy element compositions of the samples (metal:halogen ratios) using a Zeiss Neon EsB equipped with an Oxford Instruments EDX system.

Optical Measurements. Room temperature PLE, PL and PLQY measurements were performed on a HORIBA Jobin Yvon Fluorolog-3 spectrofluorometer using a Xenon lamp and Quanta- ϕ integrating sphere using the two-curve method in a varied range from 280-860 nm. UV-vis diffuse reflectance were measured to study the optical absorption properties of the compounds. Measurements were performed on polycrystalline powder samples using a PerkinElmer Lambda 750 UV/VIS/NIR Spectrometer equipped with a 100mm Spectralon InGaAs Integrating Sphere attachment over a 250-1100 nm range. TRPL was measured on a HORIBA Jobin Yvon Fluorolog-3 spectrofluorometer using a time-correlated-single-photon counting module. HORIBA Jobin Yvon NanoLEDs (pulsed light-emitting diodes) were used as the excitation source. The duration of the light pulse was shorter than 2 ns. Temperature dependence PL spectra were measured on single crystal samples using a double monochromator U1000 equipped with a photomultiplier. The excitation wavelength was the 325 nm (3.815 eV) line of a Spectra-Physics beamlock 2085 argon laser. The samples were placed in a helium bath cryostat and the measurements were performed

between 2 and 300 K. Room temperature solid-state mid-IR measurements were performed on a Bruker Tensor27 IR over the range of 400 – 4000 cm^{-1} on solid samples ground and then pressed into a matrix of KBr. The IR measurements show the presence of well-defined vibration mode peaks similar to that of their organic salts as shown in Figure S6, which confirms the good quality of compounds.

First Principles Calculations. All calculations were based on DFT implemented in the VASP code.⁶⁸ The interaction between ions and electrons was described by the projector augmented wave method.⁶⁹ The kinetic energy cutoff of 400 eV for the plane-wave basis was used for all calculations. Experimental lattice parameters were used while the atomic positions were fully relaxed until the residual forces were less than 0.02 eV/ Å. Electronic band structures and density of states (DOS) were calculated using the Perdew–Burke–Ernzerhof (PBE) exchange-correlation functional⁷⁰ while excitons were treated by using the more advanced hybrid PBE0 functional,⁶² which has 25% non-local Fock exchange. Previous PBE0 calculations have provided accurate results on exciton properties in hybrid organic-inorganic halide perovskites.^{23, 26, 71}

ASSOCIATED CONTENT

Supporting Information

Syntheses of the organic precursor salts, their NMR spectra, results of stability studies, infrared spectra, the results of low temperature crystallography experiments, tables with atomic coordinates, interatomic distances and bond angles.

The Supporting Information is available free of charge on the ACS Publications website.

Notes

The authors declare no competing financial interests.

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AUTHOR INFORMATION

Rachel Roccanova	rvow@ou.edu
Matthew Houck	houck308@gmail.com
Aymen Yangui	aymen.yangui@ou.edu
Dan Han	handan19912006@126.com
Hongliang Shi	hlshi@buaa.edu.cn
Yuntao Wu	ywu52@utk.edu
Daniel T. Glatzhofer	dtglatzhofer@ou.edu
Douglas R. Powell	d-powell@ou.edu
Shiyu Chen	chensy@ee.ecnu.edu.cn
Houcem Fourati	houcem.fourati@uvsq.fr
Alain Lusson	alain.lusson@uvsq.fr

Kamel Boukheddaden

kamel.boukheddaden@uvsq.fr

Mao-Hua Du* (corresponding author)

mhdu@ornl.gov

Bayrammurad Saparov* (corresponding author)

saparov@ou.edu

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