

1 **MONOTONICITY AND GLOBAL DYNAMICS OF A NONLOCAL
2 TWO-SPECIES PHYTOPLANKTON MODEL***

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4 **Abstract.** We investigate a nonlocal reaction-diffusion-advection system modeling the popula-
5 tion dynamics of two competing phytoplankton species in a eutrophic environment, where nutrients
6 are in abundance and the species are limited by light only for their metabolism. We first demonstrate
7 that the system does not preserve the competitive order in the pointwise sense. Then we introduce a
8 special cone \mathcal{K} involving the cumulative distributions of the population densities, and a generalized
9 notion of super- and subsolutions of the nonlocal competition system where the differential inequal-
10 ties hold in the sense of the cone \mathcal{K} . A comparison principle is then established for such super- and
11 subolutions, which implies the monotonicity of the underlying semiflow with respect to the cone \mathcal{K}
12 (Theorem 2.1). As application, we study the global dynamics of the single species system and the
13 competition system. The latter has implications for the evolution of movement for phytoplankton
14 species.

15 **Key words.** Phytoplankton; competition for light; nonlocal reaction-diffusion equations; mono-
16 tone dynamical system.

17 **AMS subject classifications.** 35B51, 35K57, 47H07, 92D25

18 **1. Introduction.** Phytoplankton are microscopic plant-like photosynthetic or-
19 ganisms that drift in the water columns of lakes and oceans. They grow abundantly
20 around the globe and are the foundation of the marine food chain. Since they trans-
21 port significant amounts of atmospheric carbon dioxide into the deep oceans, they
22 play a crucial role in climate dynamics. Nutrients and light are the essential resources
23 for the growth of phytoplankton. There are three possible ways for phytoplankton
24 to compete for nutrients and light. At one extreme, in oligotrophic ecosystems with
25 an ample supply of light, species compete for limiting nutrients [22, 27]. At the other
26 extreme, in eutrophic ecosystems with ample nutrient supply, species compete for
27 light [8, 16, 17, 33]. In some ecosystems of intermediate conditions, they compete for
28 both nutrients and light [3, 4, 18, 21, 36]. In the water column, phytoplankton diffuse
29 by water turbulence, and also sink or buoy, depending on whether they are heavier
30 than water or not [8].

31 In this paper, we study the two-species nonlocal reaction-diffusion-advection sys-
32 tem proposed by Huisman et al. [16, 18]. The system models the growth of phyto-
33 plankton species in a eutrophic vertical water column, where the species is limited by
34 light only for their metabolism. Consider a water column with unit cross-sectional
35 area and with two phytoplankton species. Let x denote the depth within the water
36 column where x varies from 0 (the top) to L (the bottom), and let $u(x, t), v(x, t)$ stand
37 for the population densities of two phytoplankton species at the location x and time

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38 t , respectively. The following system of reaction-diffusion-advection equations was
 39 proposed in [16] to describe the population dynamics of two phytoplankton species:

40 (1.1)
$$\begin{cases} u_t = D_1 u_{xx} - \alpha_1 u_x + [g_1(I(x, t)) - d_1]u, & 0 < x < L, t > 0, \\ v_t = D_2 v_{xx} - \alpha_2 v_x + [g_2(I(x, t)) - d_2]v, & 0 < x < L, t > 0, \end{cases}$$

41 with no-flux boundary conditions

42 (1.2)
$$\begin{cases} D_1 u_x(x, t) - \alpha_1 u(x, t) = 0, & x = 0, L, t > 0, \\ D_2 v_x(x, t) - \alpha_2 v(x, t) = 0, & x = 0, L, t > 0, \end{cases}$$

43 and initial conditions

44 (1.3)
$$u(x, 0) = u_0(x) \geq, \not\equiv 0, \quad v(x, 0) = v_0(x) \geq, \not\equiv 0, \quad 0 \leq x \leq L,$$

45 where for $i = 1, 2$, $D_i > 0$ is the diffusion coefficient, $\alpha_i \in \mathbb{R}$ is the sinking ($\alpha_i > 0$)
 46 or buoyant ($\alpha_i < 0$) velocity, $d_i > 0$ is the death rate, $g_i(I)$ represents the specific
 47 growth rate of phytoplankton species as a function of light intensity $I(x, t)$.

48 Light intensity is decreasing with depth due to light absorption via phytoplankton
 49 and water. By the Lambert-Beer law [23], the light intensity $I(x, t)$ is given by

50 (1.4)
$$I(x, t) = I_0 \exp \left(-k_0 x - \int_0^x [k_1 u(s, t) + k_2 v(s, t)] ds \right),$$

51 where $I_0 > 0$ is the incident light intensity, $k_0 > 0$ is the background turbidity
 52 that summarizes light absorption by all non-phytoplankton components, and k_i is
 53 the absorption coefficient of the corresponding phytoplankton species. In this model
 54 ample nutrient supply is assumed so that the phytoplankton growth is only limited
 55 by the light availability. We assume that $g_i(I)$ is a smooth function satisfying

56 (1.5)
$$g_i(0) = 0 \quad \text{and} \quad g'_i(I) > 0 \quad \text{for } I \geq 0.$$

A typical example of $g_i(I)$ takes the Michaelis-Menten form

$$g_i(I) = \frac{m_i I}{a_i + I},$$

57 where $m_i > 0$ is the maximal growth rate and $a_i > 0$ is the half saturation constant.

58 Most existing mathematical literatures on phytoplankton are focused on a single
 59 species. The single species model was considered in [33] for the self-shading case (i.e.
 60 $k_0 = 0$) and infinite long water column ($L = \infty$). The existence, uniqueness and
 61 global stability of the steady state are established in [20, 33]. It is shown in [24] that
 62 the self-shading model with any finite water column depth has a stable positive steady
 63 state, which means that the self-shading model has no critical water column depth
 64 beyond which the phytoplankton cannot persist.

65 For the case $k_0 > 0$, it is illustrated in [8] that the condition for phytoplankton
 66 bloom development can be characterized by critical water column depth and some
 67 critical values of the vertical turbulent diffusion coefficient. Du and Hsu [5] studied
 68 both single and two species competing for light with no advection. For the single
 69 species model, the existence, uniqueness, and global attractivity of a positive equilib-
 70 rium was established. Hsu and Lou [13] analyzed the critical death rate, critical water
 71 column depth, critical sinking or buoyant coefficient and critical turbulent diffusion

72 rate. Du and Mei [7] investigated the global dynamics of the single species model for
 73 the case $D = D(x)$, $\alpha = \alpha(x)$ and the asymptotic profiles of the positive steady states
 74 for small or large diffusion and deep water column when D, α are constants. Peng and
 75 Zhao [31,32] considered the effect of time-periodic light intensity I_0 at the surface, due
 76 to diurnal light cycle and seasonal changes. Ma and Ou [28] further studied the model
 77 in [31,32] and assume that $D(t), \alpha(t)$ are time periodic functions. They obtained the
 78 uniqueness and the global attractivity of the positive periodic solution of the single
 79 species model, when it exists.

80 Du et al. [6] studied the effect of photoinhibition on the single phytoplankton
 81 species, and they found that, in contrast to the case of no photoinhibition, where at
 82 most one positive steady state can exist, the model with photoinhibition possesses at
 83 least two positive steady states in certain parameter ranges. Hsu et al. [14] exam-
 84 ined the dynamics of a single species under the assumption that the amount of light
 85 absorbed by individuals is proportional to cell size, which varies for populations that
 86 reproduced by simple cell division into two equal-sized daughter cells.

87 Although many mathematical theories have been developed for single species
 88 phytoplankton model, there are very few results for two or more phytoplankton species
 89 competing for light. The existence of positive steady state and uniform persistence for
 90 two-species model were proved in [5], where there is no sinking or buoyancy. In [29],
 91 Mei and Zhang studied a nonlocal reaction-diffusion-advection system modeling the
 92 growth of multiple competitive phytoplankton species and they found that when the
 93 diffusion of the system is large, there are no positive steady states, and when the
 94 diffusion is not large, there exists at least one positive steady state under proper
 95 conditions.

96 Unlike two-species Lotka-Volterra competition model with diffusion, one main
 97 difficulty for system (1.1)-(1.4) is the lack of comparison principle, i.e.

$$98 \quad u_1(x, 0) \leq u_2(x, 0), \quad v_1(x, 0) \geq v_2(x, 0) \quad \forall x \in [0, L] \\ 99 \quad \not\Rightarrow \quad u_1(x, t) \leq u_2(x, t), \quad v_1(x, t) \geq v_2(x, t) \quad \forall (x, t) \in [0, L] \times (0, \infty),$$

101 due to the nonlocal nature of the nonlinearity. See Remark 3.10.

102 For order-preserving properties in the single species model, Shigesada and Okubo
 103 [33] observed that the cumulative distribution function $U(x, t) := \int_0^x u(s, t) ds$ satisfies
 104 a single reaction-diffusion equation without nonlocal terms. Subsequently, Ishii and
 105 Takagi [20] showed that the flow retains the natural order in U . For a related model
 106 with a water column of infinite depth, they made use of this fact to obtain a complete
 107 classification of the long-time behavior of the population. This fact was used again
 108 in Du and Hsu [5] to determine the long-time dynamics for a single species model
 109 with finite water depth. More recently, Ma and Ou [28] established the comparison
 110 principle for U in the single species model.

111 For the competition model, we will show, by adapting arguments due to Du and
 112 Hsu [5] and Ma and Ou [28], that the cumulative distribution functions

$$(U(x, t), V(x, t)) = \left(\int_0^x u(s, t) ds, \int_0^x v(s, t) ds \right)$$

113 satisfy a nonlocal, strongly coupled system, with non-standard boundary condition
 114 (see (3.3)), and that the resulting system has the strong order-preserving property.

115 Our main result (Theorem 2.1) says that system (1.1)-(1.4) forms a strongly
 116 monotone dynamical system with respect to the order induced by the special cone

115 $\mathcal{K} = \mathcal{K}_1 \times (-\mathcal{K}_1)$, where

116 (1.6)
$$\mathcal{K}_1 = \left\{ \phi \in C([0, L], \mathbb{R}) : \int_0^x \phi(s) ds \geq 0 \text{ for } x \in (0, L] \right\}.$$

117 The new features of this paper can be described as follows: First, Theorem 2.1
 118 is the first monotonicity result for the nonlocal competition system involving two
 119 phytoplankton species. Second, the definition of the relevant cone \mathcal{K} facilitates the
 120 connection with general theory of monotone dynamical systems. Third, generalized
 121 notion of super- and subsolutions (see Definition 3.2), which is new even for the case of
 122 single species, are given. They can potentially be used to obtain qualitative properties
 123 of solutions for the nonlocal system (1.1)-(1.4).

124 The rest of the paper is organized as follows: In Section 2, we state our main
 125 results. In Section 3, we first introduce the notion of super- and subsolutions of
 126 (1.1)-(1.4) with respect to the cone \mathcal{K} , and establish the comparison principle for
 127 the super- and subsolutions. Then we apply the monotonicity result to establish the
 128 global dynamics of the single species model in a general setting. Section 4 is devoted to
 129 the spectral analysis of semi-trivial steady states, and the global dynamics of system
 130 (1.1)-(1.4) are established for three different biological scenarios. In Section 5, we
 131 present some numerical results and discussion.

132 **2. Main Results.** Let \mathbf{X} be a Banach space over \mathbb{R} . We call a subset $K \subset \mathbf{X}$
 133 a *cone* if (i) K is convex, (ii) $\mu K \subset K$ for all $\mu \geq 0$, and (iii) $K \cap (-K) = \{0\}$. A
 134 cone K is said to be solid if it has nonempty interior. Furthermore, for $x, y \in \mathbf{X}$, we
 135 write $x \leqslant_K y$, $x <_K y$ and $x \ll_K y$ if $y - x \in K$, $y - x \in K \setminus \{0\}$ and $y - x \in \text{Int } K$
 136 respectively.

Let \mathcal{K}_1 be given by (1.6). It is straightforward to verify that \mathcal{K}_1 is a solid cone in
 the Banach space $C([0, L]; \mathbb{R})$ with interior

$$\text{Int } \mathcal{K}_1 = \left\{ \phi \in C([0, L]; \mathbb{R}) : \phi(0) > 0, \int_0^x \phi(s) ds > 0 \text{ for } x \in (0, L] \right\}.$$

137 Let $\mathcal{K} = \mathcal{K}_1 \times (-\mathcal{K}_1)$. Then \mathcal{K} is likewise a solid cone in the Banach space $C([0, L]; \mathbb{R}^2)$
 138 with interior given by $\text{Int } \mathcal{K} = \text{Int } \mathcal{K}_1 \times (-\text{Int } \mathcal{K}_1)$. The cone \mathcal{K} induces the partial
 139 order relations $\leqslant_{\mathcal{K}}$, $<_{\mathcal{K}}$ and $\ll_{\mathcal{K}}$ in the usual way.

140 We shall prove that (1.1)-(1.4) is a strongly monotone dynamical system with
 141 respect to the order induced by the cone \mathcal{K} .

142 **THEOREM 2.1.** *Suppose $\{(u_i, v_i)\}_{i=1,2}$ are non-negative solutions of (1.1)-(1.4)
 143 such that $u_2(\cdot, 0) \geq, \not\equiv 0$ and $v_1(\cdot, 0) \geq, \not\equiv 0$ and*

144
$$(u_1(\cdot, 0), v_1(\cdot, 0)) <_{\mathcal{K}} (u_2(\cdot, 0), v_2(\cdot, 0)).$$

145 Then $(u_1(\cdot, t), v_1(\cdot, t)) \ll_{\mathcal{K}} (u_2(\cdot, t), v_2(\cdot, t))$ for all $t > 0$.

146 By Theorem 2.1, system (1.1)-(1.4) is a strongly monotone dynamical system on
 147 $C([0, L]; \mathbb{R}_+^2)$ with respect to the order generated by \mathcal{K} , which together with the theory
 148 of strongly monotone dynamical systems [2, 12, 15, 25, 34, 37], provides a useful tool to
 149 investigate the global dynamics of two-species system (1.1)-(1.4). As a by-product of
 150 our monotonicity result, we also generalize the existing results for single species (see
 151 Subsection 3.2) and give a simple proof based on monotonicity arguments and the
 152 concept of subhomogeneous mappings.

153 As application, we turn our attention to the effects of diffusion and advection on
 154 the global dynamics of (1.1)-(1.4).

155 THEOREM 2.2. *If $D_1 = D_2$, $\alpha_1 < \alpha_2$, $g_1 = g_2$, $d_1 = d_2$, and that both semi-trivial
156 steady states exist, then the first species u drives the second species v to extinction,
157 regardless of initial condition.*

158 Theorem 2.2 shows that the competitor with smaller advection rate has compet-
159 itive advantages, i.e., *smaller advection rate is selected*. By the Lambert-Beer law,
160 the deeper the water column, the weaker the light intensity. Therefore, it is more
161 advantageous for phytoplankton species to move up.

162 THEOREM 2.3. *If $D_1 < D_2$, $\alpha_1 = \alpha_2 \geq [g(1) - d]L$, $g_1 = g_2$, $d_1 = d_2$, and that
163 both semi-trivial steady states exist, then the faster diffuser v drives the slower diffuser
164 u to extinction, regardless of initial condition.*

165 Theorem 2.3 implies that if sinking rate is large, competitor with faster diffusion
166 will always displace the slower one, i.e., *faster diffuser wins*. Intuitively, when both
167 species are sinking with equal and large velocity, faster diffusion can counter balance
168 the tendency to sink and provide individuals with better access to light.

169 THEOREM 2.4. *If $D_1 < D_2$, $\alpha_1 = \alpha_2 \leq 0$, $g_1 = g_2$, $d_1 = d_2$, and that both
170 semi-trivial steady states exist, then the slower diffuser u drives faster diffuser v to
171 extinction, regardless of initial condition.*

172 Theorem 2.4 suggests that if the phytoplankton species are buoyant, the competi-
173 tor with slower diffusion rate will always displace the faster one, i.e., *slower diffusion
174 rate will be selected*. This is in sharp contrast to Theorem 2.3. The reason for this
175 result is that when the phytoplankton are buoyant, turbulent diffusion actually dis-
176 places individuals from the top of the water column, where the light intensity is the
177 strongest.

178 **3. A General Model with Spatio-Temporally Varying Coefficients.** We
179 shall study a generalized version of system (1.1)-(1.4), which allows coefficients to vary
180 explicitly with both space and time. We formulate the nonlocal reaction-diffusion-
181 advection model as follows:

$$(3.1) \quad \begin{cases} u_t = (D_1 u_x - \alpha_1 u)_x + f_1(x, t, \int_0^x u(s, t) ds, \int_0^x v(s, t) ds)u, & 0 < x < L, t > 0, \\ v_t = (D_2 v_x - \alpha_2 v)_x + f_2(x, t, \int_0^x u(s, t) ds, \int_0^x v(s, t) ds)v, & 0 < x < L, t > 0, \\ D_1 u_x - \alpha_1 u = 0, & x = 0, L, t > 0, \\ D_2 v_x - \alpha_2 v = 0, & x = 0, L, t > 0, \\ u(x, 0) = u_0(x) \geq, \not\equiv 0, \quad v(x, 0) = v_0(x) \geq, \not\equiv 0, & 0 \leq x \leq L, \end{cases}$$

183 where, for $i = 1, 2$, $D_i = D_i(x, t) > 0$, $\alpha_i = \alpha_i(x, t)$, and the functions $f_i(x, t, p, q)$ are
184 smooth and satisfy

$$185 \quad (\mathbf{H}) \quad \frac{\partial f_i}{\partial p} < 0, \quad \frac{\partial f_i}{\partial q} < 0 \quad \text{and} \quad \frac{\partial f_i}{\partial x} \leq 0 \quad \text{for all } x \in [0, L] \text{ and } t, p, q \geq 0.$$

186 The assumption holds, e.g. when $f_i(x, t, p, q) = g_i(I_0 \exp(-k_0 x - k_1 p - k_2 q)) -$
187 $d_i(x, t)$ such that g_i is non-decreasing, and d_i is non-decreasing in x . In particular, it
188 includes (1.1)-(1.4), and the previous works [5, 29] as particular cases.

189 **3.1. Strong Monotonicity of (3.1).** This subsection is devoted to proving the
190 monotonicity of system (3.1) with respect to the order induced by cone \mathcal{K} under the
191 assumption (H). First, we state the following standard result (see, e.g. [10, Ch. 3]).

PROPOSITION 3.1. *For continuous, non-negative initial data $(u_0(x), v_0(x))$, sys-
tem (3.1) has a unique solution*

$$(u, v) \in C([0, \infty); C([0, L]; \mathbb{R}_+^2)) \cap C^1((0, \infty); C^\infty([0, L]; \mathbb{R}_+^2)),$$

192 which depends continuously on initial data. Moreover, if $u_0(x) \not\equiv 0$, (resp. $v_0(x) \not\equiv 0$),
 193 then $u(x, t) > 0$ (resp. $v(x, t) > 0$) for $(x, t) \in [0, L] \times (0, \infty)$.

194 Next, we define the following super- and subsolution concepts for (3.1). Note that
 195 the differential inequalities appearing below are to be understood in the sense of cone
 196 \mathcal{K} for each time t . These inequalities hold, in particular, if the differential inequalities
 197 hold in the pointwise sense everywhere.

DEFINITION 3.2. *We say that*

$$(\bar{u}, \underline{v}), (\underline{u}, \bar{v}) \in C([0, T]; C([0, L]; \mathbb{R}_+^2)) \cap C^1((0, T]; C^\infty([0, L]; \mathbb{R}_+^2))$$

198 form a pair of super- and subsolutions of (3.1) in the interval $[0, T]$, if

$$199 \quad (3.2) \quad \begin{cases} \bar{u}_t \geq_{\mathcal{K}_1} (D_1 \bar{u}_x - \alpha_1 \bar{u})_x + f_1(x, t, \int_0^x \bar{u}(s, t) ds, \int_0^x v(s, t) ds) \bar{u}, & 0 < t \leq T, \\ \underline{v}_t \leq_{\mathcal{K}_1} (D_2 \underline{v}_x - \alpha_2 \underline{v})_x + f_2(x, t, \int_0^x \bar{u}(s, t) ds, \int_0^x \underline{v}(s, t) ds) \underline{v}, & 0 < t \leq T, \\ \underline{u}_t \leq_{\mathcal{K}_1} (D_1 \underline{u}_x - \alpha_1 \underline{u})_x + f_1(x, t, \int_0^x \underline{u}(s, t) ds, \int_0^x \bar{v}(s, t) ds) \underline{u}, & 0 < t \leq T, \\ \bar{v}_t \geq_{\mathcal{K}_1} (D_2 \bar{v}_x - \alpha_2 \bar{v})_x + f_2(x, t, \int_0^x \underline{u}(s, t) ds, \int_0^x \bar{v}(s, t) ds) \bar{v}, & 0 < t \leq T, \\ D_1 \bar{u}_x - \alpha_1 \bar{u} \leq 0 \leq D_1 \underline{u}_x - \alpha_1 \underline{u}, & x = 0, 0 < t \leq T, \\ D_1 \bar{u}_x - \alpha_1 \bar{u} \geq 0 \geq D_1 \underline{u}_x - \alpha_1 \underline{u}, & x = L, 0 < t \leq T, \\ D_2 \bar{v}_x - \alpha_2 \bar{v} \leq 0 \leq D_2 \underline{v}_x - \alpha_2 \underline{v}, & x = 0, 0 < t \leq T, \\ D_2 \bar{v}_x - \alpha_2 \bar{v} \geq 0 \geq D_2 \underline{v}_x - \alpha_2 \underline{v}, & x = L, 0 < t \leq T, \\ (\bar{u}(\cdot, 0), \underline{v}(\cdot, 0)) \geq_{\mathcal{K}} (\underline{u}(\cdot, 0), \bar{v}(\cdot, 0)). \end{cases}$$

200 The main result of this section is

THEOREM 3.3. *Assume that f_1, f_2 satisfy (H). Let (\bar{u}, \underline{v}) and (\underline{u}, \bar{v}) be a pair of super- and subsolutions of (3.1) in the interval $[0, T]$. If $\underline{u} > 0$ and $\underline{v} > 0$ in $[0, L] \times [0, T]$, then*

$$(\bar{u}(\cdot, t), \underline{v}(\cdot, t)) \geq_{\mathcal{K}} (\underline{u}(\cdot, t), \bar{v}(\cdot, t)) \quad \text{for } 0 \leq t \leq T.$$

Moreover, if there exists $t_0 \in (0, T]$ such that $\bar{u} > 0$ and $\bar{v} > 0$ in $[0, L] \times (0, t_0]$, and

$$(\bar{u}(\cdot, t_0) - \underline{u}(\cdot, t_0), \underline{v}(\cdot, t_0) - \bar{v}(\cdot, t_0)) \notin \text{Int } \mathcal{K},$$

201 then $(\bar{u}(x, t), \underline{v}(x, t)) \equiv (\underline{u}(x, t), \bar{v}(x, t))$ for $x \in [0, L]$ and $0 \leq t \leq t_0$.

202 A direct consequence of Theorem 3.3 is the strong monotonicity of the continuous
 203 semiflow generated by (3.1). It includes Theorem 2.1 as a particular case.

COROLLARY 3.4. *Assume that f_1, f_2 satisfy (H). Suppose $\{(u_i, v_i)\}_{i=1,2}$ are two
 non-negative solutions of (3.1), such that $u_1(\cdot, 0) \geq, \not\equiv 0$, $v_2(\cdot, 0) \geq, \not\equiv 0$, and*

$$(u_1(\cdot, 0), v_1(\cdot, 0)) >_{\mathcal{K}} (u_2(\cdot, 0), v_2(\cdot, 0)).$$

204 Then $(u_1(\cdot, t), v_1(\cdot, t)) \gg_{\mathcal{K}} (u_2(\cdot, t), v_2(\cdot, t))$ for all $t > 0$.

205 The proof is postponed to later in the section.

206 To show Theorem 3.3, we consider the cumulative distribution functions

$$207 \quad U(x, t) = \int_0^x u(s, t) ds, \quad V(x, t) = \int_0^x v(s, t) ds.$$

208 Then $U(0, t) \equiv 0$, $V(0, t) \equiv 0$ for $t \geq 0$, and $U_x(x, t) = u(x, t)$, $V_x(x, t) = v(x, t)$. In
 209 this way, (3.1) is transformed into the following *strongly coupled, non-local* system of

210 (U, V) (see also [28] for the single species case):

211 (3.3)
$$\begin{cases} U_t = D_1 U_{xx} - \alpha_1 U_x + G_1[U, V, U_x, V_x], & 0 < x < L, t > 0, \\ V_t = D_2 V_{xx} - \alpha_2 V_x + G_2[U, V, U_x, V_x], & 0 < x < L, t > 0, \\ U(0, t) = 0, \quad D_1 U_{xx}(L, t) - \alpha_1 U_x(L, t) = 0, & t > 0, \\ V(0, t) = 0, \quad D_2 V_{xx}(L, t) - \alpha_2 V_x(L, t) = 0, & t > 0, \\ U(x, 0) = \int_0^x u_0(s) ds = U_0(x), & 0 \leq x \leq L, \\ V(x, 0) = \int_0^x v_0(s) ds = V_0(x), & 0 \leq x \leq L, \end{cases}$$

212 where, letting $F_1(x, t, U, V) = \int_0^U f_1(x, t, z, V) dz$, $F_2(x, t, U, V) = \int_0^V f_2(x, t, U, z) dz$,

213
$$G_1[U, V, U_x, V_x](x, t)$$

214 $= \int_0^x f_1\left(s, t, \int_0^s u(y, t) dy, \int_0^s v(y, t) dy\right) u(s, t) ds$
215 $= \int_0^x f_1\left(s, t, U(s, t), V(s, t)\right) U_x(s, t) ds$
216 $= \int_0^x \left\{ \frac{d}{ds} [F_1(s, t, U(s, t), V(s, t))] - \frac{\partial F_1}{\partial x}\left(s, t, U(s, t), V(s, t)\right) \right.$
217 $\quad \left. - \frac{\partial F_1}{\partial V}\left(s, t, U(s, t), V(s, t)\right) V_x(s, t) \right\} ds$
218 $= F_1(x, t, U(x, t), V(x, t)) - \int_0^x \frac{\partial F_1}{\partial x}\left(s, t, U(s, t), V(s, t)\right) ds$
219 (3.4) $\quad - \int_0^x \frac{\partial F_1}{\partial V}\left(s, t, U(s, t), V(s, t)\right) V_x(s, t) ds$

220 and

221
$$G_2[U, V, U_x, V_x](x, t)$$

222 $= \int_0^x f_2\left(s, t, \int_0^s u(y, t) dy, \int_0^s v(y, t) dy\right) v(s, t) ds$
223 $= \int_0^x f_2\left(s, t, U(s, t), V(s, t)\right) V_x(s, t) ds$
224 $= F_2(x, t, U(x, t), V(x, t)) - \int_0^x \frac{\partial F_2}{\partial x}\left(s, t, U(s, t), V(s, t)\right) ds$
225 (3.5) $\quad - \int_0^x \frac{\partial F_2}{\partial U}\left(s, t, U(s, t), V(s, t)\right) U_x(s, t) ds.$

For (3.3), we define the Banach space

$$X_1 = \{\phi \in C^1([0, L], \mathbb{R}) : \phi(0) = 0\}$$

with the usual C^1 norm. The usual cone P_1 in X_1 is

$$P_1 = \{\phi \in X_1 : \phi(x) \geq 0 \text{ for } x \in [0, L]\},$$

with interior

$$\text{Int } P_1 = \{\phi \in X_1 : \phi'(0) > 0, \phi(x) > 0 \text{ for } x \in (0, L]\}.$$

226 Let $X = X_1 \times X_1$, and $P = P_1 \times (-P_1)$. Then P is a cone in X with interior given by
 227 $\text{Int } P = \text{Int } P_1 \times (-\text{Int } P_1)$. The cone P generates the partial order relations $\leq_P, <_P$
 228 and \ll_P on X .

By construction, the solutions (U, V) of (3.3) live in the convex set $E = E_1 \times E_1$, where

$$E_1 = \{\phi \in C^1([0, L]) : \phi(0) = 0, \text{ and } \phi'(x) \geq 0 \text{ for } x \in [0, L]\}.$$

229 From now on we assume the initial data of (3.3) to be in E . Under this assumption,
 230 the existence and uniqueness of the solution $(U(x, t), V(x, t))$ can be derived from
 231 those of $(u(x, t), v(x, t))$.

DEFINITION 3.5. *We say that*

$$(\bar{U}, \underline{V}), (\underline{U}, \bar{V}) \in C([0, T]; E) \cap C^1((0, T]; C^\infty([0, L]; \mathbb{R}_+^2))$$

232 *form a pair of super- and subsolutions of (3.3) in the interval $[0, T]$, if the derivatives
 233 $(\bar{u}, \underline{v}) = (\bar{U}_x, \underline{V}_x)$ and $(\underline{u}, \bar{v}) = (\underline{U}_x, \bar{V}_x)$ form a pair of super- and subsolutions of
 234 (3.1) in the interval $[0, T]$, in the sense of Definition 3.2.*

235 We now prove a strong maximum principle for the system (3.3), which is the key
 236 to proving the strong monotonicity of (3.3).

237 LEMMA 3.6. *Assume that f_1, f_2 satisfy (H). Let (\bar{U}, \underline{V}) and (\underline{U}, \bar{V}) be a pair of
 238 super- and subsolutions of (3.3) in the interval $[0, t^*]$ for some $t^* > 0$, so that*

$$239 \quad (3.6) \quad \bar{U}_x(x, t) > 0 \quad \text{and} \quad \bar{V}_x(x, t) > 0 \quad \text{for } 0 \leq x \leq L, \text{ and } 0 < t \leq t^*,$$

and

$$\underline{U}(x, t) \leq \bar{U}(x, t), \quad \bar{V}(x, t) \geq \underline{V}(x, t) \quad \text{for } 0 \leq x \leq L, \text{ and } 0 \leq t \leq t^*.$$

240 *If one of the following holds:*

241 **(a)** $\underline{U}(x^*, t^*) = \bar{U}(x^*, t^*)$ or $\underline{V}(x^*, t^*) = \bar{V}(x^*, t^*)$ for some $x^* \in (0, L]$;
 242 **(b)** $(\bar{U} - \underline{U})_x(0, t^*) = 0$ or $(\bar{V} - \underline{V})_x(0, t^*) = 0$,
 243 then

$$244 \quad (3.7) \quad (\underline{U}(x, t), \bar{V}(x, t)) \equiv (\bar{U}(x, t), \underline{V}(x, t)) \quad \text{for } 0 \leq x \leq L, 0 \leq t \leq t^*.$$

Proof. In the following we improve upon the arguments of [28] to prove the strong maximum principle for (3.3). We first consider the case when (a) holds. For definiteness assume that $\underline{U}(x^*, t^*) = \bar{U}(x^*, t^*)$ for some $x^* \in (0, L]$. Denote

$$W(x, t) = \bar{U}(x, t) - \underline{U}(x, t).$$

Then by (3.2),

$$(\bar{u} - \underline{u})_t - [D_1(\bar{u} - \underline{u})_x + \alpha_1(\bar{u} - \underline{u})]_x \geq \kappa_1 f_1(x, t, \bar{U}(x, t), \underline{V}(x, t)) - f_1(x, t, \underline{U}(x, t), \bar{V}(x, t))$$

245 Fixing t , and integrating the above from 0 to x , we have, in terms of W ,

$$\begin{aligned} 246 \quad & W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x \\ 247 \quad & \geq \int_0^x f_1(s, t, \bar{U}(s, t), \underline{V}(s, t))\bar{U}_x(s, t) ds - \int_0^x f_1(s, t, \underline{U}(s, t), \bar{V}(s, t))\underline{U}_x(s, t) ds \\ 248 \quad & \geq \int_0^x f_1(s, t, \bar{U}(s, t), \bar{V}(s, t))\bar{U}_x(s, t) ds - \int_0^x f_1(s, t, \underline{U}(s, t), \bar{V}(s, t))\underline{U}_x(s, t) ds \end{aligned} \quad (3.8)$$

250 where we used $\bar{V}(x, t) \geq \underline{V}(x, t)$ for $(x, t) \in [0, L] \times [0, t^*]$. Integrating by parts as in
 251 (3.4), we have

$$\begin{aligned}
 252 \quad & W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x \\
 253 \quad & \geq F_1\left(x, t, \bar{U}(x, t), \bar{V}(x, t)\right) - F_1\left(x, t, \underline{U}(x, t), \bar{V}(x, t)\right) \\
 254 \quad & + \int_0^x \left[\frac{\partial F_1}{\partial \bar{V}}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] \bar{V}_x(s, t) ds \\
 255 \quad & + \int_0^x \left[\frac{\partial F_1}{\partial x}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial x}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] ds \\
 256 \quad & \geq h(x, t)W + \int_0^x \left[\frac{\partial F_1}{\partial \bar{V}}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] \bar{V}_x(s, t) ds, \\
 257 \quad & \quad (3.9)
 \end{aligned}$$

for $x \in [0, L]$, $t \in (0, t^*]$, where

$$h(x, t) = \int_0^1 f_1\left(x, t, \xi \bar{U}(s, t) + (1 - \xi) \underline{U}(s, t), \bar{V}(s, t)\right) d\xi \in L_{loc}^\infty([0, L] \times \mathbb{R}_+).$$

258 Note that we have used $\frac{\partial}{\partial \bar{U}}\left(\frac{\partial F_1}{\partial x}\right) = \frac{\partial f_1}{\partial x} \leq 0$, i.e. $\frac{\partial F_1}{\partial x}$ is non-increasing in \bar{U} in the
 259 last inequality of (3.9). Summarizing, we have

$$\begin{aligned}
 260 \quad & W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x - h(x, t)W \\
 261 \quad & (3.10) \geq \int_0^x \left[\frac{\partial F_1}{\partial \bar{V}}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] \bar{V}_x(s, t) ds.
 \end{aligned}$$

262 Since $\frac{\partial}{\partial \bar{U}}\left(\frac{\partial F_1}{\partial \bar{V}}\right) = \frac{\partial f_1}{\partial \bar{V}} < 0$, i.e. $\frac{\partial F_1}{\partial \bar{V}}$ is non-increasing in \bar{U} , $\bar{U} \geq \underline{U}$, and $\bar{V}_x > 0$,
 263 the last integral is non-negative. Thus $W = \bar{U} - \underline{U}$ satisfies the following linear
 264 differential inequality:

$$265 \quad (3.11) \quad W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x - h(x, t)W \geq 0, \quad \text{for } x \in (0, L], t \in (0, t^*].$$

We claim that $W \equiv 0$ in $[0, L] \times [0, t^*]$. If not, then the parabolic strong maximum principle applied to (3.11) implies that $W(x, t^*) > 0$ for $x \in (0, L)$. Therefore, if there exists some $x^* \in (0, L]$ such that $W(x^*, t^*) = 0$, then $x^* = L$, i.e., $W(L, t^*) = 0$, and hence $W_t(L, t^*) \leq 0$. By the boundary conditions at $(x, t) = (L, t^*)$,

$$D_1 \bar{U}_{xx} - \alpha_1 \bar{U}_x \geq 0 \geq D_1 \underline{U}_{xx} - \alpha_1 \underline{U}_x,$$

266 we have $D_1(L, t^*)W_{xx}(L, t^*) - \alpha_1(L, t^*)W_x(L, t^*) \geq 0$. Then by (3.10) we have

$$\begin{aligned}
 267 \quad & 0 \geq W_t(L, t^*) \\
 268 \quad & \geq \int_0^L \left[\frac{\partial F_1}{\partial \bar{V}}\left(s, t^*, \underline{U}(s, t^*), \bar{V}(s, t^*)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t^*, \bar{U}(s, t^*), \bar{V}(s, t^*)\right) \right] \bar{V}_x(s, t^*) ds.
 \end{aligned}$$

Since $\underline{U}(x, t^*) \leq \bar{U}(x, t^*)$ in $[0, L]$, and $\bar{V}_x > 0$, we deduce that the above inequality holds only if $\underline{U}(x, t^*) \equiv \bar{U}(x, t^*)$ for all $x \in [0, L]$, i.e., $W(x, t^*) \equiv 0$ for all $x \in [0, L]$. This is a contradiction and thus $W = \bar{U} - \underline{U} \equiv 0$ in $[0, L] \times [0, t^*]$. It follows that equality holds everywhere in (3.8) and (3.9), in particular,

$$\int_0^x f_1\left(s, t, \bar{U}(s, t), \underline{U}(s, t)\right) \bar{U}_x(s, t) ds \equiv \int_0^x f_1\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \bar{U}_x(s, t) ds,$$

270 for all $x \in [0, L]$ and $0 < t \leq t^*$. Since $\bar{U}_x > 0$ and $\frac{\partial f_1}{\partial V} < 0$, we deduce that
 271 $\bar{V}(x, t) \equiv \underline{V}(x, t)$ in $[0, L] \times (0, t^*]$ and, by continuity, in $[0, L] \times [0, t^*]$.

272 The remaining case $\bar{V}(x^*, t^*) = \underline{V}(x^*, t^*)$ for some $x^* \in (0, L]$ can be handled
 273 similarly. This completes the proof in case (a) holds.

274 Next, assume (b) holds. We claim that necessarily there is a sequence of $t_j \nearrow t^*$
 275 such that alternative (a) holds, so that we can deduce similarly that $(\bar{U}, \underline{V}) \equiv (\underline{U}, \bar{V})$ in
 276 $[0, L] \times [0, t_j]$ for all j , whence (3.7) holds as well upon letting $t_j \nearrow t^*$. To see the claim,
 277 assume for contradiction that $\bar{U} > \underline{U}$ and $\bar{V} > \underline{V}$ for $(x, t) \in (0, L) \times [t^* - \delta', t^*]$ for some
 278 δ' . Then, observe that the boundary condition ensures $W(0, t^*) = \bar{U}(0, t^*) - \underline{U}(0, t^*) =$
 279 0. Since W satisfies the differential inequality (3.11), we may apply Hopf's Lemma [26,
 280 Lemma 2.8] to deduce that $(\bar{U} - \underline{U})_x(0, t^*) > 0$. Similarly, we can deduce that
 281 $(\bar{V} - \underline{V})_x(0, t^*) > 0$ as well, i.e. alternative (b) does not hold in this case. This
 282 establishes the claim and finishes the proof. \square

283 Theorem 3.3 is a consequence of Lemma 3.6 and the following result:

284 LEMMA 3.7. *Assume that f_1, f_2 satisfy (H). Let (\bar{U}, \underline{V}) and (\underline{U}, \bar{V}) be a pair of
 285 super- and subsolutions of (3.3) in the time interval $[0, T]$. If*

286 (3.12) $\underline{U}_x(x, t) > 0$, and $\underline{V}_x(x, t) > 0$ for $(x, t) \in [0, L] \times [0, T]$,

287 then

288 (3.13) $\bar{U}(x, t) \geq \underline{U}(x, t)$ and $\underline{V}(x, t) \leq \bar{V}(x, t)$ for $0 \leq x \leq L$, $0 \leq t \leq T$.

289 *Proof.* It is enough to prove the result for arbitrary but finite $T > 0$. Given a pair
 290 of super- and subsolutions (\bar{U}, \underline{V}) and (\underline{U}, \bar{V}) in a bounded interval $[0, T]$, we show
 291 (3.13) in two steps.

Step 1. For each small $\delta > 0$, define

$$(\bar{U}^\delta, \underline{V}^\delta) = (\bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2), \quad \text{and} \quad (\underline{U}^\delta, \bar{V}^\delta) = (\underline{U} - \delta\rho_1, \bar{V} + \delta\rho_2),$$

292 where $\rho_i(x, t) := \int_0^x \exp \left(Mt + \int_0^y \frac{\alpha_i(s, t)}{D_i(s, t)} ds \right) dy$ for $i = 1, 2$. By (3.12), there exists
 293 $\delta_0 > 0$ such that for each $\delta \in (0, \delta_0]$,

$$294 \quad (3.14) \quad \begin{cases} (\bar{U}^\delta, \underline{V}^\delta), (\underline{U}^\delta, \bar{V}^\delta) \in E & \text{for } t \in [0, T], \\ \bar{U}_x^\delta > 0, \underline{V}_x^\delta > 0, \underline{U}_x^\delta > 0, \bar{V}_x^\delta > 0 & \text{for } (x, t) \in [0, L] \times [0, T], \\ (\bar{U}^\delta(\cdot, 0), \underline{V}^\delta(\cdot, 0)) \gg_P (\underline{U}^\delta(\cdot, 0), \bar{V}^\delta(\cdot, 0)). & \end{cases}$$

295 It is also clear that there is $C_0 > 0$ (independent of δ) such that

$$296 \quad (3.15) \quad \max_{i=1,2} \|\rho_i\|_{C([0, L] \times [0, T])} \leq C_0 \min_{i=1,2} \inf_{[0, L] \times [0, T]} (\rho_i)_x(x, t).$$

297 We claim that $(\bar{U}^\delta, \underline{V}^\delta)$ and $(\underline{U}^\delta, \bar{V}^\delta)$ forms a pair of super- and subsolutions for
 298 (3.3) in the interval $[0, T]$, in the sense of Definition 3.5. It remains to show the
 299 differential inequalities (3.2) for δ small, as the initial and boundary conditions are
 300 clearly satisfied. A sufficient condition for the first one to hold is

$$301 \quad (3.16) \quad \delta(\rho_1)_{x,t} \geq \kappa_1 [f_1(x, t, \bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2) - f_1(x, t, \bar{U}, \underline{V})] \bar{U}_x + \delta\rho_{1,x} f_1(x, t, \bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2).$$

302 The inequality (3.16) holds since the following holds pointwisely in $[0, L] \times [0, T]$:

$$\begin{aligned} 303 \quad & \delta(\rho_1)_{x,t} - [f_1(x, t, \bar{U} + \delta\rho_1, \bar{V} - \delta\rho_2) - f_1(x, t, \bar{U}, \bar{V})] \bar{U}_x - \delta\rho_{1,x} f_1(x, t, \bar{U} + \delta\rho_1, \bar{V} - \delta\rho_2) \\ 304 \quad & \geq \delta \left(\rho_{1,x} \left[M + \int_0^x \left(\frac{\alpha_1(s, t)}{D_1(s, t)} \right)_t ds - \|f_1\|_\infty \right] - \|Df_1\|_\infty (\rho_1 + \rho_2) \|\bar{U}_x\|_\infty \right), \\ 305 \end{aligned}$$

306 (note that $\bar{U}_x, \bar{V}_x \in C([0, L] \times [0, T])$ by definition of super- and subsolutions) and, by
307 (3.15), the term in the square bracket is non-negative provided the positive parameter
308 $M = M(C_0, \|f\|_{C^1})$ is chosen large enough (but uniformly for $\delta \in (0, \delta_0]$). In the same
309 way, one can show the rest of the differential inequalities. In summary, there is $M > 0$
310 so that for all $\delta \in (0, \delta_0]$, $(\bar{U}^\delta, \bar{V}^\delta)$ and $(\underline{U}^\delta, \underline{V}^\delta)$ form a pair of super- and subsolutions
311 for (3.3) in the interval $[0, T]$. This proves our first claim.

312 **Step 2.** Next, we claim that for all $\delta > 0$,

$$313 \quad (3.17) \quad \bar{U}^\delta(x, t) > \underline{U}^\delta(x, t) \quad \text{and} \quad \underline{V}^\delta(x, t) < \bar{V}^\delta(x, t) \quad \text{for } (x, t) \in (0, L] \times [0, T].$$

314 Suppose not, then it follows from (3.14) that there exists a positive maximal time
315 denoted by $t^* \in (0, T]$ such that $\underline{U}^\delta(x, t) < \bar{U}^\delta(x, t)$, $\bar{V}^\delta(x, t) > \underline{V}^\delta(x, t)$ hold for $0 <$
316 $x \leq L$ and $0 \leq t < t^*$, and $\underline{U}^\delta(x^*, t^*) = \bar{U}^\delta(x^*, t^*)$ or $\bar{V}^\delta(x^*, t^*) = \underline{V}^\delta(x^*, t^*)$ for some
317 $x^* \in (0, L]$. It follows from Lemma 3.6 that $\underline{U}^\delta(x, t) \equiv \bar{U}^\delta(x, t)$ and $\bar{V}^\delta(x, t) \equiv \underline{V}^\delta(x, t)$
318 for all $0 \leq x \leq L$ and $0 \leq t \leq t^*$, which is a contradiction to (3.14). This shows (3.17).
319 Letting $\delta \rightarrow 0$ in (3.17), we deduce that (3.13) holds for $(x, t) \in [0, L] \times [0, T]$. \square

320 Now we prove Corollary 3.4, which includes Theorem 2.1 as a special case.

321 *Proof of Corollary 3.4.* For $i = 1, 2$, let

$$322 \quad (3.18) \quad (U_i(x, t), V_i(x, t)) = \left(\int_0^x u_i(s, t) ds, \int_0^x v_i(s, t) ds \right).$$

323 If we assume in addition that

$$324 \quad (3.19) \quad u_2(x, 0) = U_{2,x}(x, 0) > 0 \quad \text{and} \quad v_1(x, 0) = V_{1,x}(x, 0) > 0 \quad \text{in } [0, L],$$

then by applying the strong maximum principle to the first and second equations of
(3.1) separately, we deduce that

$$u_2 = U_{2,x} > 0 \quad \text{and} \quad v_1 = V_{1,x} > 0 \quad \text{in } [0, L] \times [0, T].$$

325 Therefore, applying Lemma 3.7, we see that if $(U_1(\cdot, 0), V_1(\cdot, 0)) \geq_P (U_2(\cdot, 0), V_2(\cdot, 0))$
326 and (3.19) holds, then

$$327 \quad (3.20) \quad (U_1(\cdot, t), V_1(\cdot, t)) \geq_P (U_2(\cdot, t), V_2(\cdot, t)) \quad \text{for all } t > 0.$$

328 By the fact that initial data satisfying (3.19) is dense in E , we can show that for
329 general initial data in E , if $(U_1(\cdot, 0), V_1(\cdot, 0)) >_P (U_2(\cdot, 0), V_2(\cdot, 0))$, then (3.20) holds.

It remains to show that if $(U_1(\cdot, 0), V_1(\cdot, 0)) >_P (U_2(\cdot, 0), V_2(\cdot, 0))$ and that both
 $U_{1,x}, V_{2,x}$ are non-negative and non-trivial, then

$$(U_1(\cdot, t), V_1(\cdot, t)) \gg_P (U_2(\cdot, t), V_2(\cdot, t)) \quad \text{for all } t > 0.$$

This follows from Lemma 3.6, provided it can be verified that

$$u_1(x, t) = (U_1)_x(x, t) > 0, \quad v_2(x, t) = (V_2)_x(x, t) > 0 \quad \text{for } 0 \leq x \leq L, 0 < t \leq T.$$

330 But this is an immediate consequence of the strong maximum principle applied to the
331 equations of u_1 and v_2 separately. \square

332 **3.2. Global Dynamics of the Single Species Model.** In this section, we gen-
 333 eralize some known results about the following single species model, which is obtained
 334 by setting $v = 0$ in (3.1):

335 (3.21)
$$\begin{cases} \theta_t = (D_1\theta_x - \alpha_1\theta)_x + f_1(x, t, \int_0^x \theta(s, t) ds, 0)\theta, & 0 < x < L, t > 0, \\ D_1\theta_x - \alpha_1\theta = 0, & x = 0, L, t > 0, \\ \theta(x, 0) = \theta_0(x) \geq, \not\equiv 0, & 0 \leq x \leq L, \end{cases}$$

336 where $D_1 = D_1(x, t) > 0$, $\alpha_1 = \alpha_1(x, t)$, and f_1 are smooth and **(H)** holds.

337 The equation (3.21) generates a continuous semiflow in $C([0, L]; \mathbb{R}_+)$ (see, e.g.
 338 [10]). Furthermore, by regarding the nonlocal term $f_1(x, t, \int_0^x \theta(s, t) ds, 0)$ as a given
 339 coefficient, we can view (3.21) as a linear non-autonomous parabolic equation. It
 340 follows from the classical maximum principle that $\theta(x, t) > 0$ for $x \in [0, L]$ and $t > 0$.

341 Define $\bar{\theta} \in C([0, \infty); C([0, L]; \mathbb{R}_+) \cap C^1((0, \infty); C^\infty([0, L]; \mathbb{R}_+))$ to be a superso-
 342 lution of (3.21) if

343 (3.22)
$$\begin{cases} \bar{\theta}_t \geq_{\mathcal{K}_1} (D_1\bar{\theta}_x - \alpha_1\bar{\theta})_x + f_1(x, t, \int_0^x \bar{\theta}(s, t) ds, 0)\bar{\theta}, & t > 0, \\ D_1\bar{\theta}_x - \alpha_1\bar{\theta} = 0, & x = 0, L, t > 0. \end{cases}$$

344 And define $\underline{\theta}$ to be a subsolution of (3.21) if it satisfies the reverse inequality. As a
 345 by-product of the proofs of Lemmas 3.6 and 3.7, we can similarly show that the single
 346 species model is strongly monotone with respect to the order generated by cone \mathcal{K}_1 .

COROLLARY 3.8. *Assume that f_1 satisfies **(H)**. Let $\bar{\theta}$ and $\underline{\theta}$ be super- and subso-
 lution of (3.21) such that*

$$\bar{\theta}(x, t) > 0, \quad \underline{\theta}(x, t) > 0, \quad \text{in } [0, L] \times [0, T], \quad \text{and} \quad \bar{\theta}(\cdot, 0) \geq_{\mathcal{K}_1} \underline{\theta}(\cdot, 0).$$

347 Then $\bar{\theta}(\cdot, t) \geq_{\mathcal{K}_1} \underline{\theta}(\cdot, t)$ for all $t > 0$. Furthermore, if for some $t_0 > 0$ we have
 348 $\bar{\theta}(\cdot, t_0) - \underline{\theta}(\cdot, t_0) \not\in \text{Int } \mathcal{K}_1$, then $\bar{\theta}(\cdot, t) \equiv \underline{\theta}(\cdot, t)$ for $t \in [0, t_0]$.

349 In particular, the continuous semiflow generated by (3.21) is strongly monotone with
 350 respect to the order induced by the cone \mathcal{K}_1 .

351 In contrast to Corollary 3.8, we show here that the pointwise competitive order
 352 is not preserved by (3.21).

PROPOSITION 3.9. *For $i = 1, 2$, let θ_i be a solution of (3.21), with initial condi-
 tions $\theta_{i,0} \in \{\psi \in C^2([0, L]) : D_1\psi_x = \alpha_1\psi \text{ for } x = 0, L\}$. If*

$$\theta_{1,0} \leq, \not\equiv \theta_{2,0} \quad \text{in } [0, L], \quad \text{and} \quad \theta_{1,0} \equiv \theta_{2,0} \quad \text{in } [L - \delta, L] \text{ for some } \delta > 0,$$

353 then $\theta_1(L, t) > \theta_2(L, t)$ for all $0 < t \ll 1$.

354 *Proof.* Since the initial conditions are C^2 and consistent with the boundary con-
 355 dition, the solutions θ_i are of class $C_{x,t}^{2,1}$ in $[0, L] \times [0, \infty)$. Hence, it is enough to show
 356 that $(\theta_1)_t(L, 0) > (\theta_2)_t(L, 0)$. Precisely, at $(x, t) = (L, 0)$,

357
$$\begin{aligned} (\theta_1)_t &= [D_1(\theta_1)_x - \alpha_1\theta_1]_x + f_1(L, 0, \int_0^L \theta_1(s, 0) ds, 0)\theta_1 \\ 358 &> [D_1(\theta_1)_x - \alpha_1\theta_1]_x + f_1(L, 0, \int_0^L \theta_2(s, 0) ds, 0)\theta_1 \\ 359 &= [D_1(\theta_2)_x - \alpha_1\theta_2]_x + f_1(L, 0, \int_0^L \theta_2(s, 0) ds, 0)\theta_2 = (\theta_2)_t. \end{aligned} \quad \square$$

To illustrate Proposition 3.9, we choose initial conditions $\{\theta_{i,0}\}_{i=1,2}$ so that

$$\theta_{1,0} \leq_{P_1} \theta_{2,0} \quad \text{and} \quad \theta_{1,0} \leq_{\mathcal{K}_1} \theta_{2,0},$$

but only the order with respect to \mathcal{K}_1 is preserved by the semiflow; see Figure 1.

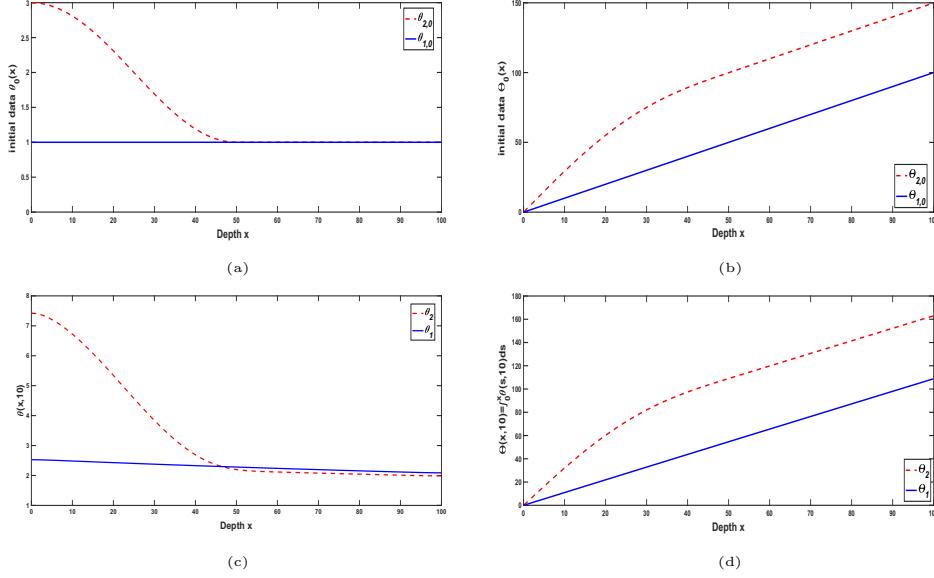


FIG. 1. Numerical solution of (3.21), with $D_1 = 1$, $\alpha_1 = 0$, $L = 100$, $f_1(x, t, \Theta, 0) = g(\exp(-k_0 x - k_1 \Theta))$, where $g(I) = \frac{I}{10+I}$ and $k_0 = k_1 = d = 0.001$, and initial condition $\theta_{1,0} = \chi_{[0, L/2]}(\cos(2\pi x/L) + 1) + 1$ and $\theta_{2,0} = 1$. Panels (a) and (c) are the population densities $\theta_i(x, t)$ ($i = 1, 2$) at times $t = 0$ and $t = 10$ resp.; Panels (b) and (d) are the initial cumulative distribution functions of population densities $\Theta_i(x, t) = \int_0^x \theta_i(s, t) ds$ ($i = 1, 2$) at times $t = 0$ and $t = 10$. The first (resp. second) species is represented by the red/dotted line (resp. blue/solid line).

361 **Remark 3.10.** By choosing $u_i(\cdot, 0) = \theta_{i,0}$ for $i = 1, 2$, and $v_1(\cdot, 0) \equiv v_2(\cdot, 0) \equiv \epsilon$,
362 then $(u_1(\cdot, 0), v_1(\cdot, 0)) \leq_P (u_2(\cdot, 0), v_2(\cdot, 0))$. However, it follows from the above result
363 and continuous dependence on initial data that $(u_1(\cdot, t), v_1(\cdot, t)) \not\leq_P (u_2(\cdot, t), v_2(\cdot, t))$
364 for some $t > 0$.

366 As a consequence of monotone dynamical systems theory, one can show the
367 uniqueness and global asymptotic stability of positive equilibria (in the case of au-
368 tonomous semiflow) or positive periodic solution (in the case of time-periodic semi-
369 flow). We will show the latter here, as the former follows as an easy consequence.

370 The following eigenvalue problem will be useful for our later purposes:

$$371 \quad (3.23) \quad \begin{cases} \varphi_t = (D_1 \varphi_x - \alpha_1 \varphi)_x + f_1(x, t, 0, 0) \varphi + \mu \varphi, & 0 < x < L, 0 < t < T, \\ D_1 \varphi_x - \alpha_1 \varphi = 0, & x = 0, L, 0 < t < T, \\ \varphi(x, 0) = \varphi(x, T), & 0 \leq x \leq L, \\ \varphi(x, t) > 0, & 0 \leq x \leq L, 0 \leq t \leq T. \end{cases}$$

372 It is well known (see, e.g., [11]) that (3.23) has a principal eigenvalue, denoted by μ_1 ,
373 with the corresponding positive eigenfunction.

374 **PROPOSITION 3.11.** *Assume that f_1 satisfies (H), and let D_1, α_1, f_1 be T -periodic
375 in t , and there exists $M_1 > 0$ such that*

$$376 \quad (3.24) \quad \sup_{[0, L] \times [0, T]} f_1(x, t, M_1, 0) < 0 \quad \text{and} \quad \|f_1(\cdot, \cdot, \cdot, 0)\|_{L^\infty([0, L] \times [0, T] \times [0, \infty))} \leq M_1.$$

377 Let μ_1 be the principal eigenvalue of (3.23).

378 (a) If $\mu_1 \geq 0$, then every solution of (3.21) converges to zero;

379 (b) If $\mu_1 < 0$, then (3.21) has a unique positive T -periodic solution. Furthermore,
380 it attracts all non-negative, non-trivial solutions of (3.21).

381 In case $f_1(x, t, p, 0) = g(I_0 \exp(-k_0 x - k_1 p)) - d(x, t)$ where $g(\cdot)$ satisfies (1.5), the
382 condition (3.24) is clearly satisfied, and the above result generalizes all previous re-
383 sults [5, 7, 28, 31, 32]. Our main contribution is a short proof of the boundedness of
384 trajectories, which has not been proven when all coefficients vary periodically with
385 time. This allows the use of the concept of subhomogeneity to show the existence,
386 uniqueness and global stability of positive steady state simultaneously.

387 *Proof of Proposition 3.11.* We will apply [37, Theorem 2.3.4] to prove this propo-
388 sition. Let \tilde{Q}_T be the Poincaré map of time T , generated by the T -periodic equation
389 (3.21). It is obvious that the Poincaré map \tilde{Q}_T is monotone by Corollary 3.8, and
390 compact in $C([0, L])$ by parabolic estimate. Therefore, we need only to verify that ev-
391 ery positive orbit of \tilde{Q}_T in $C([0, L]; \mathbb{R}_+)$ is bounded, \tilde{Q}_T is strongly subhomogeneous,
392 and the Fréchet derivative $\mathcal{D}\tilde{Q}_T(0)$ is compact and strongly positive.

Claim 1. The semiflow is point dissipative, i.e. there exists $M > 0$, independent of
initial data, such that

$$\limsup_{t \rightarrow \infty} \|\theta(\cdot, t)\|_{C([0, L])} \leq M.$$

By the fact that $f_1(x, t, p, 0)$ is uniformly bounded in L^∞ , Harnack inequality [19,
Theorem 2.5] applies, so that there is a uniform positive constant $C' > 0$ such that

$$\sup_{0 < x < L} \theta(x, t) \leq C' \inf_{0 < x < L} \theta(x, t) \quad \text{for all } t \geq 1.$$

393 By (3.24), it is possible to choose a small constant $\delta_2 > 0$ such that

$$394 \quad (3.25) \quad C' \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx < 0 \quad \text{for } 0 \leq t \leq T.$$

395 It suffices to show that $\limsup_{t \rightarrow \infty} \int_0^L \theta dx \leq \max\{M_1, C'LM_1/\delta_2\}$. To this end, it is
396 enough to show the following claim.

CLAIM 3.12. *The differential inequality*

$$\frac{d}{dt} \int_0^L \theta(x, t) dx \leq -\delta_3 \int_0^L \theta(x, t) dx$$

397 holds whenever $\int_0^L \theta(x, t) dx > \max\{M_1, C'LM_1/\delta_2\}$.

Now, denote $\theta_*(t) = \inf_x \theta(x, t)$ and $\theta^*(t) = \sup_x \theta(x, t)$, then

$$M_1 < \frac{\delta_2}{C'L} \int_0^L \theta(x, t) dx \leq \frac{\delta_2}{C'} \theta^*(t) \leq \delta_2 \theta_*(t).$$

398 Integrating the equation of θ over $(0, L)$, we obtain

$$\begin{aligned}
 399 \quad & \frac{d}{dt} \int_0^L \theta(x, t) dx \\
 400 \quad &= \int_0^L f_1 \left(x, t, \int_0^x \theta(s, t) ds, 0 \right) \theta(x, t) dx \\
 401 \quad &\leq \int_0^L f_1(x, t, x\theta_*(t), 0) \theta(x, t) dx \\
 402 \quad &\leq \int_0^{\delta_2} f_1(x, t, 0, 0) \theta(x, t) dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) \theta(x, t) dx \\
 403 \quad &\leq \theta^*(t) \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx \theta_*(t) \\
 404 \quad &\leq \left(C' \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx \right) \theta_*(t) \\
 405 \quad &\leq \left(C' \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx \right) \frac{1}{C'L} \int_0^L \theta(x, t) dx.
 \end{aligned}$$

407 This proves the point dissipativity.

408 **Claim 2.** The Poincaré map is strongly subhomogeneous.

409 We will show that \tilde{Q}_T is strongly subhomogeneous, i.e.

$$410 \quad (3.26) \quad \tilde{Q}_T(\lambda\theta_0) \gg_{\kappa_1} \lambda\tilde{Q}_T(\theta_0) \quad \text{for all } \theta_0 >_{P_1} 0 \text{ and } \lambda \in (0, 1).$$

411 Let $\theta(x, t)$ be solution to (3.21) with initial condition θ_0 . For $(x, t) \in (0, L) \times [0, T]$,

$$\begin{aligned}
 412 \quad & (\lambda\theta)_t = (D_1(\lambda\theta)_x - \alpha_1(\lambda\theta))_x + f_1(x, t, \int_0^x \theta(s, t) ds, 0)(\lambda\theta) \\
 413 \quad & < (D_1(\lambda\theta)_x - \alpha_1(\lambda\theta))_x + f_1(x, t, \int_0^x \lambda\theta(s, t) ds, 0)(\lambda\theta).
 \end{aligned}$$

415 i.e. $\lambda\theta$ is a subsolution to the (3.21) with initial condition $\lambda\theta_0$. Since the above
 416 inequality is strict, $\lambda\theta$ is not identically equal to the solution of (3.21) with initial
 417 condition $\lambda\theta_0$. By Corollary 3.8 and evaluate at time $t = T$, we deduce (3.26).

418 **Claim 3.** The Fréchet derivative $\mathcal{D}\tilde{Q}_T(0)$ is compact and strongly positive.

419 This follows directly from the fact that $\mathcal{D}\tilde{Q}_T(0) = Z(T)$, where $Z(t)$ is the analytic
 420 semigroup generated by the linearized system of (3.21) at $\theta = 0$:

$$421 \quad (3.27) \quad \begin{cases} \theta_t = (D_1\theta_x - \alpha_1\theta)_x + f_1(x, t, 0, 0)\theta, & 0 < x < L, t > 0, \\ D_1\theta_x - \alpha_1\theta = 0, & x = 0, L, t > 0, \\ \theta(x, 0) = \theta_0 \geq, \not\equiv 0, & 0 < x < L. \end{cases}$$

422 That $Z(T)$ is strongly positive follows from standard parabolic maximum principle.
 423 Moreover, by standard parabolic L^p estimate, $Z(T)$ is a bounded map from $C([0, L])$
 424 to $C^2([0, L])$. The map $Z(T)$ is thus compact, by the Arzelà-Ascoli Theorem.

425 If $\mu_1 \geq 0$, then $r(\mathcal{D}\tilde{Q}_T(0)) = \exp(-\mu_1 T) \leq 1$. By [37, Theorem 2.3.4(a)], every
 426 solution of (3.21) converges to zero. If $\mu_1 < 0$, then $r(\mathcal{D}\tilde{Q}_T(0)) = \exp(-\mu_1 T) > 1$.
 427 By [37, Theorem 2.3.4(b)], the map \tilde{Q}_T has a unique positive fixed point $\tilde{\vartheta}$ such that
 428 every positive orbit with non-negative, non-trivial, continuous initial data converges

429 to $\tilde{\vartheta}$. This means that system (3.21) has a unique positive T -periodic solution $\tilde{\theta}$,
430 determined by $\tilde{\theta}(\cdot, 0) = \tilde{\theta}(\cdot, T) = \tilde{\vartheta}$, which attracts all non-negative and non-trivial
431 solutions of (3.21). \square

432 *Remark 3.13.* Within the context of a single species, we improved previous results
433 in [28] by showing a strong maximum principle (which implies strong monotonicity of
434 the semiflow) for super- and subsolutions (which satisfies only differential inequalities),
435 and by allowing the coefficients to be space-time heterogeneous.

436 **4. Global Dynamics for the Nonlocal Two-species Model.** It is well
437 known that diffusion and advection rates have significant effects on the outcome of
438 competition. In this section, we apply Theorem 4.1 to analyze the global dynamics
439 of two-species competition system. To obtain qualitative results, we restrict ourselves
440 for the remainder of the paper to consider the autonomous case (1.1) - (1.3), when
441 D_i, α_i, d_i are constants. In the introduction, the light intensity $I(x, t)$ is given by
442 (1.4), where the shading coefficients of the two species are given by k_1, k_2 . However,
443 by transforming $(\tilde{u}, \tilde{v}) = (k_1 u, k_2 v)$ and $\tilde{g}_i(I_0 \cdot) = g_i(\cdot)$, and by observing that k_1, k_2
444 do not affect the dynamics qualitatively, we may assume $k_1 = k_2 = 1$ and $I_0 = 1$
445 without loss of generality, so that the light intensity (1.4) can be simplified to

$$446 \quad (4.1) \quad I(x, t) = \exp \left(-k_0 x - \int_0^x [u(s, t) + v(s, t)] ds \right).$$

447 We focus on the following three different cases:

- 448 (i) $D_1 = D_2, \alpha_1 < \alpha_2$;
- 449 (ii) $D_1 < D_2, \alpha_1 = \alpha_2 \geq [g(1) - d]L > 0$;
- 450 (iii) $D_1 < D_2, \alpha_1 = \alpha_2 \leq 0$.

451 Due to the strongly monotonicity proved in Theorem 2.1, to a large extent, its
452 dynamics can be determined by the stability/instability of the semi-trivial solution
453 of the stationary problem [2, 12, 15, 25, 34, 37]. For the convenience of the readers, we
454 state the precise abstract theorem here.

455 **THEOREM 4.1** ([15, Theorem B] and [25, Theorem 1.3]). *If the system (1.1)-
456 (1.4) has no positive steady states, and the semi-trivial steady state $(0, \tilde{v})$ (resp. $(\tilde{u}, 0)$)
457 is linearly unstable, then $(\tilde{u}, 0)$ (resp. $(0, \tilde{v})$) is globally asymptotically stable among
458 all non-negative, non-trivial solutions.*

459 *Remark 4.2.* Our setting is slightly more general than that outlined in [15]. In
460 particular, the semiflow Q_t generated by (3.1) is defined in $Y^+ = Y_1^+ \times Y_1^+$, where
461 $Y_1^+ = C([0, L]; \mathbb{R}_+)$, but the semiflow only preserve the order generated by the weaker
462 cone $\mathcal{K} = \mathcal{K}_1 \times (-\mathcal{K}_1)$, with $Y_1^+ \subsetneq \mathcal{K}_1$. However, it is straightforward to observe
463 that [15, Propositions 2.1 and 2.4] are independent of the above assumption, and
464 that the proofs of [15, Theorem B] and [25, Theorem 1.3] both stand in our setting.
465 Therefore, we omit the proof of Theorem 4.1 here.

466 In preparation to apply Theorem 4.1, we will demonstrate that the equation

$$467 \quad (4.2) \quad \begin{cases} \theta_t = D\theta_{xx} - \alpha\theta_x + [g(e^{-k_0 x - \int_0^x \theta(s, t) ds}) - d]\theta = 0, & 0 < x < L, \\ D\theta_x - \alpha\theta = 0, & x = 0, L, \\ \theta(x, 0) = \theta_0(x) \geq 0, & 0 \leq x \leq L, \end{cases}$$

468 has a unique positive steady state $\tilde{\theta}$, which is always linearly stable, and then char-
469 acterize the stability of the two semi-trivial steady states in terms of two standard
470 principal eigenvalue problems.

471 **4.1. An Eigenvalue Problem for the Single Species Model.** For constants
 472 $D > 0$, $\alpha \in \mathbb{R}$ and $h \in C([0, L])$, consider the following standard eigenvalue problem:

473 (4.3)
$$\begin{cases} D\phi_{xx} - \alpha\phi_x + h(x)\phi + \lambda\phi = 0, & 0 < x < L, \\ D\phi_x - \alpha\phi = 0, & x = 0, L. \end{cases}$$

474 By setting $\psi = e^{-(\alpha/D)x}\phi$, the problem (4.3) can be transformed into a self-adjoint
 475 problem

476 (4.4)
$$\begin{cases} -D(e^{(\alpha/D)x}\psi_x)_x - h(x)e^{(\alpha/D)x}\psi = \lambda e^{(\alpha/D)x}\psi, & 0 < x < L, \\ \psi_x(0) = \psi_x(L) = 0. \end{cases}$$

Therefore, all eigenvalues of (4.4) (and thus (4.3)) are real, and we can denote the
 smallest eigenvalue by $\lambda_1(D, \alpha, h)$. Define

$$d_* = -\lambda_1(D, \alpha, -g(e^{-k_0 x})).$$

477 It is easy to show that d_* is positive. In fact, d_* is the critical death rate.

478 THEOREM 4.3 ([5, Theorem 2.1], [13, Theorem 3.1]). *If $0 < d < d_*$, then (4.2)
 479 has a unique positive steady state, denoted by $\tilde{\theta}(x)$. If $d \geq d_*$, then zero is the only
 480 nonnegative steady state of (4.2).*

481 We linearize (4.2) at $\tilde{\theta}$ to obtain the following eigenvalue problem:

482 (4.5)
$$\begin{cases} D\phi_{xx} - \alpha\phi_x + [g(\sigma) - d]\phi - \tilde{\theta}\sigma g'(\sigma) \int_0^x \phi(s) ds + \mu\phi = 0, & 0 < x < L, \\ D\phi_x - \alpha\phi = 0, & x = 0, L, \end{cases}$$

483 where $\sigma = e^{-k_0 x - \int_0^x \tilde{\theta}(s) ds}$.

484 Our result says that $\tilde{\theta}$ is linearly stable. In fact, there is a real eigenvalue of (4.5)
 485 which is strictly less than the real part of all other eigenvalues of (4.5).

486 THEOREM 4.4. *Let $\tilde{\theta}$ be the unique positive steady state of (4.2). The eigenvalue
 487 problem (4.5) admits a real, simple eigenvalue μ_1 and an eigenfunction $\phi \gg_{\mathcal{K}_1} 0$, such
 488 that $\mu_1 < \operatorname{Re} \mu$ for all eigenvalues $\mu \neq \mu_1$. It is characterized as the unique eigenvalue
 489 of (4.5) with the eigenfunction $\phi \gg_{\mathcal{K}_1} 0$. Furthermore, $\mu_1 > 0$.*

490 *Proof.* Assume $\tilde{\theta}(x)$ is a positive steady state of (4.2), and let $\theta_0 \in C([0, L]; \mathbb{R})$.
 491 Then $\theta(\cdot, t) = \hat{\Phi}_t(\theta_0)$, where $\hat{\Phi}_t$ denotes the continuous semiflow generated by (4.2).
 492 Then $z(x, t) = \mathcal{D}\hat{\Phi}_t(\tilde{\theta})[\theta_0](x)$ satisfies the linear equation

493 (4.6)
$$z_t + \mathcal{L}z = 0, \quad z(0) = \theta_0.$$

where the unbounded operator

$$\mathcal{L} = -D\partial_{xx} + \alpha\partial_x - [g(\sigma) - d] + \tilde{\theta}\sigma g'(\sigma) \left(\int_0^x \cdot \right)$$

is defined on the domain

$$\operatorname{Dom}(\mathcal{L}) = \{z \in C^2((0, L)) \cap C^1([0, L]) : \mathcal{L}z \in C([0, L]), Dz_x - \alpha z|_{x=0, L} = 0\}.$$

494 According to [30, Proposition 3.1.4], the linear equation (4.6) generates an analytic
 495 semigroup $e^{-\mathcal{L}t}$ on $C([0, L])$. Thus $\mathcal{D}\hat{\Phi}_t(\tilde{\theta}) = e^{-\mathcal{L}t}$.

496 For $\theta_0 \in \mathcal{K}_1$, $\epsilon > 0$, the monotonicity of $\hat{\Phi}_t$ with respect to cone \mathcal{K}_1 implies

$$497 \quad \frac{\theta(\cdot, t; \tilde{\theta} + \epsilon\theta_0) - \theta(\cdot, t; \tilde{\theta})}{\epsilon} = \frac{\hat{\Phi}_t(\tilde{\theta} + \epsilon\theta_0) - \hat{\Phi}_t(\tilde{\theta})}{\epsilon} \geq_{\mathcal{K}_1} 0.$$

498 Upon taking the limit as $\epsilon \rightarrow 0$, we get $\mathcal{D}\hat{\Phi}_t(\tilde{\theta})[\theta_0] \geq_{\mathcal{K}_1} 0$. In other words, $e^{-\mathcal{L}t} =$
499 $\mathcal{D}\hat{\Phi}_t(\tilde{\theta})$ is a positive operator with respect to the order generated by cone \mathcal{K}_1 in the
500 sense that $\mathcal{D}\hat{\Phi}_t(\tilde{\theta})\mathcal{K}_1 \subset \mathcal{K}_1$ holds for $t \geq 0$.

501 Next, we show that the analytic semigroup $e^{-\mathcal{L}t} = \mathcal{D}\hat{\Phi}_t(\tilde{\theta})$ is strongly positive
502 with respect to the order generated by \mathcal{K}_1 . To prove this, we only need to show that
503 $\int_0^x z(s, t) ds > 0$ and $z(0, t) > 0$ for all $t > 0$. Since $e^{-\mathcal{L}t} = \mathcal{D}\hat{\Phi}_t(\tilde{\theta})$ is a positive opera-
504 tor with respect to cone \mathcal{K}_1 , then $\int_0^x z(s, t) ds \geq 0$. Therefore, if $\int_0^x z(s, t) ds > 0$ does
505 not hold, then there exists some $(x_0, t_0) \in (0, L] \times (0, \infty)$ such that $\int_0^{x_0} z(s, t_0) ds = 0$.

Let $\int_0^x z(s, t) ds = Z(x, t)$. Using the relation

$$[g(\sigma) - d]z - \tilde{\theta}\sigma g'(\sigma)Z = [(g(\sigma) - d)Z]' + k_0 g'(\sigma)\sigma Z,$$

506 we may integrate (4.6) over $(0, x)$ to obtain the differential inequality

$$507 \quad (4.7) \quad Z_t - DZ_{xx} + \alpha Z_x - [g(\sigma) - d]Z = k_0 \int_0^x g'(\sigma)\sigma Z ds \geq 0.$$

508 Since $\theta_0 \not\equiv 0$ and $Z(\cdot, 0) \not\equiv 0$, then the strong maximum principle implies $Z(x, t) >$
509 0 for all $x \in (0, L)$ and $t > 0$, i.e., $x_0 = L$ and $Z(L, t_0) = 0$. Then $Z_t(L, t_0) \leq 0$,
510 and by the boundary condition, we deduce $DZ_{xx}(L, t_0) - \alpha Z_x(L, t_0) = Dz_x(L, t_0) -$
511 $\alpha z(L, t_0) = 0$. It follows from (4.7) that

$$512 \quad (4.8) \quad 0 \geq Z_t(L, t_0) = k_0 \int_0^L g'(\sigma)\sigma Z ds.$$

513 Since $k_0 > 0$, $\sigma > 0$, $g'(\sigma) > 0$, then $Z(x, t_0) \equiv 0$ for all $x \in [0, L]$. Contradiction.

514 Hence, $Z(x, t) = \int_0^x z(s, t) ds > 0$ for all $t > 0$ and $x \in (0, L]$. Since $Z(0, t) \equiv 0$
515 and $Z(x, t)$ satisfies (4.7) for all $t > 0$, then $z(0, t) = Z_x(0, t) > 0$ for all $t > 0$ by the
516 Hopf boundary lemma.

517 Therefore, for each $t > 0$, the operator $e^{-\mathcal{L}t}$ is compact and strongly positive on
518 $C([0, L])$ with respect to the order generated by \mathcal{K}_1 . It follows by standard arguments
519 in [34, Ch. 7] that the elliptic eigenvalue problem (4.5) has a principal eigenvalue
520 $\mu_1 \in \mathbb{R}$ with all the stated properties, except for $\mu_1 > 0$.

To show $\mu_1 > 0$, we suppose to the contrary that $\mu_1 \leq 0$ and use $\phi_1 \gg_{\mathcal{K}_1} 0$ to get

$$\tilde{\theta}\sigma g'(\sigma) \int_0^x \phi_1(s) ds > 0 \quad \text{for } x \in (0, L].$$

Then (4.5) yields that

$$D\phi_{1,xx} - \alpha\phi_{1,x} + [g(\sigma) - d]\phi_1 + \mu_1\phi_1 > 0 \quad \text{for } 0 < x < L.$$

Next, we use the facts $\int_0^x \phi_1(s) ds > 0$ and $\tilde{\theta} > 0$ for $x \in [0, L]$, to obtain the constant
 $c > 0$ such that $\min_{[0, L]}(c\tilde{\theta} - \phi_1) = 0$. Then $\varphi = c\tilde{\theta} - \phi_1$ satisfies

$$\begin{cases} D\varphi_{xx} - \alpha\varphi_x + [g(\sigma) - d]\varphi + \mu_1\varphi < \mu_1 c\tilde{\theta} \leq 0 & \text{for } 0 < x < L, \\ D\varphi_x = \alpha\varphi & \text{for } x = 0, L, \\ \min_{[0, L]} \varphi = 0. \end{cases}$$

521 By the strict differential inequality and non-negativity of φ we must have $\varphi > 0$ in
 522 $(0, L)$ and that $\varphi(x_0) = 0$ for some $x_0 \in \{0, L\}$. But the Hopf boundary lemma says
 523 $\varphi_x(x_0) \neq 0$, which contradicts the boundary condition $\varphi_x(x_0) = \frac{\alpha}{D}\varphi(x_0) = 0$. \square

524 **4.2. Eigenvalue Problems for the Two-species Model.** In this subsection,
 525 we study the linear eigenvalue problem of the two-species model associated with the
 526 stability of semi-trivial steady states.

527 We assume the parameters are chosen so that system (1.1)-(1.4) has two semi-
 528 trivial steady states $(\tilde{u}, 0)$ and $(0, \tilde{v})$ (e.g. if the death rates d_i are not too large). The
 529 associated linearized eigenvalue problem at $(\tilde{u}, 0)$ is

$$(4.9) \quad \begin{cases} D_1\phi_{xx} - \alpha_1\phi_x + [g_1(\sigma_1) - d_1]\phi - \tilde{u}\sigma_1g'_1(\sigma_1)[\int_0^x \phi(s) ds + \int_0^x \varphi(s) ds] + \Lambda\phi = 0, & 0 < x < L, \\ D_2\varphi_{xx} - \alpha_2\varphi_x + [g_2(\sigma_1) - d_2]\varphi + \Lambda\varphi = 0, & 0 < x < L, \\ D_1\phi_x - \alpha_1\phi = D_2\varphi_x - \alpha_2\varphi = 0, & x = 0, L, \end{cases}$$

531 where $\sigma_1(x) = e^{-k_0x - \int_0^x \tilde{u}(s) ds}$.

532 We shall exploit the fact that the second equation is decoupled from the first.

533 Consider the following eigenvalue problem:

$$(4.10) \quad \begin{cases} D_2\varphi_{xx} - \alpha_2\varphi_x + [g_2(\sigma_1) - d_2]\varphi + \lambda\varphi = 0, & 0 < x < L, \\ D_2\varphi_x - \alpha_2\varphi = 0, & x = 0, L. \end{cases}$$

535 As already discussed, (4.10) admits a real principal eigenvalue, denoted by $\lambda_u =$
 536 $\lambda_1(D_2, \alpha_2, g_2(\sigma_1) - d_2)$, which is simple, and its corresponding eigenfunction φ_1 does
 537 not change sign, and $\lambda_u < \lambda$ for all $\lambda \neq \lambda_u$. The stability properties of $(\tilde{u}, 0)$ are
 538 determined by the sign of λ_u , as the next result shows.

539 **PROPOSITION 4.5.** *The problem (4.9) has a principal eigenvalue $\Lambda_1 \in \mathbb{R}$, in the
 540 sense that $\Lambda_1 \leq \text{Re } \Lambda$ for all eigenvalues Λ of (4.9) and that the corresponding eigen-
 541 function can be chosen in $\mathcal{K} \setminus \{(0, 0)\}$. Furthermore, (denote $Y_1^+ = C([0, L]; \mathbb{R}_+)$)*

- 542 (a) *If the principal eigenvalue λ_u of (4.10) is positive, then $\Lambda_1 > 0$.
 543 (b) *If the principal eigenvalue λ_u of (4.10) is non-positive, then $\Lambda_1 = \lambda_u \leq 0$
 544 and the corresponding eigenfunction can be chosen in $\text{Int } \mathcal{K}_1 \times (-\text{Int } Y_1^+)$.**

545 *Proof.* By Theorem 2.1, the semiflow $\{Q_t\}_{t \geq 0}$, generated by the system (1.1)-
 546 (1.4) is strongly monotone with respect to the cone \mathcal{K} . As a result, the linear problem
 547 at any steady state generates a semigroup that is monotone with respect to the cone
 548 \mathcal{K} . Therefore, by standard arguments in [34, Ch. 7], we deduce that the elliptic
 549 problem (4.9), obtained by linearizing (1.1)-(1.4) at the steady state $(\tilde{u}, 0)$, has a
 550 principal eigenvalue Λ_1 with the stated properties. In particular, we can choose the
 551 eigenfunction corresponding to Λ_1 from within $\mathcal{K} \setminus \{(0, 0)\}$.

552 Now, consider the case when the principal eigenvalue λ_u of (4.10) is positive. Let
 553 $\Lambda_1 \in \mathbb{R}$ be the principal eigenvalue of (4.9) with eigenfunction $(\phi_1, \varphi_1) \in \mathcal{K} \setminus \{(0, 0)\}$.
 554 We claim that $\Lambda_1 > 0$. There are two cases to consider: (i) $\varphi_1 \neq 0$; (ii) $\varphi_1 = 0$.

555 In Case (i), (Λ_1, φ_1) furnishes an eigenpair of problem (4.10), the latter of which
 556 has smallest eigenvalue $\lambda_u > 0$. Thus, $\Lambda_1 \geq \lambda_u > 0$.

557 In Case (ii), (Λ_1, ϕ_1) furnishes an eigenpair of

$$(4.11) \quad \begin{cases} D_1\phi_{xx} - \alpha_1\phi_x + [g_1(\sigma_1) - d_1]\phi - \tilde{u}\sigma_1g'_1(\sigma_1)\int_0^x \phi(s) ds + \Lambda\phi = 0, & 0 < x < L, \\ D_1\phi_x - \alpha_1\phi = 0, & x = 0, L. \end{cases}$$

559 By Theorem 4.4, (4.11) has a positive principal eigenvalue μ_1 , and μ_1 is always positive.
 560 Hence, we must have $\Lambda_1 \geq \mu_1 > 0$. This finishes the proof in case $\lambda_u > 0$.

Next, let $\lambda_u \leq 0$ and let $\varphi_1 \in (-\text{Int } Y_1^+) \subset (-\text{Int } \mathcal{K}_1)$ be the corresponding principal eigenfunction of (4.10). It remains to construct $\phi_1 \in \text{Int } \mathcal{K}_1$ such that λ_u is an eigenvalue of (4.9) with eigenfunction $(\phi_1, \varphi_1) \in \text{Int } \mathcal{K}_1 \times (-\text{Int } Y_1^+)$. To that end, define the operator $\mathcal{L}_1 = -D_1 \partial_{xx} + \alpha_1 \partial_x - [g_1(\sigma_1) - d_1] + \tilde{u} \sigma_1 g_1'(\sigma_1) (\int_0^x \cdot)$. By Theorem 4.4, the spectrum $\sigma(\mathcal{L}_1) \subset \{z \in \mathbb{C} : \text{Re } z > 0\}$. And hence for $\lambda_u \leq 0$, 0 is not an eigenvalue of $\mathcal{L}_1 - \lambda_u \mathcal{I}$, and the problem

$$\begin{cases} \mathcal{L}_1 \phi - \lambda_u \phi = -\tilde{u} \sigma_1 g_1'(\sigma_1) \int_0^x \varphi_1(s) ds, & 0 < x < L, \\ D_1 \phi_x - \alpha_1 \phi = 0, & x = 0, L, \end{cases}$$

has a unique solution ϕ_1 . In fact, let $f = -\tilde{u} \sigma_1 g_1'(\sigma_1) \int_0^x \varphi_1(s) ds$, then $f > 0$ and

$$\phi_1 = (\mathcal{L}_1 - \lambda_u)^{-1} f = \int_0^\infty e^{\lambda_u t} S_t f dt,$$

where $S_t = e^{-\mathcal{L}_1 t}$ is the analytic semigroup generated by \mathcal{L}_1 (see, e.g. [9, Theorem 3, Sect. 7.4]). From the proof of Theorem 4.4, S_t is strongly positive with respect to the order generated by cone \mathcal{K}_1 . Therefore, $S_t f \gg_{\mathcal{K}_1} 0$ for all $t > 0$, and

$$\phi_1 \geq_{\mathcal{K}_1} \int_1^\infty e^{\lambda_u t} S_t f dt \geq_{\mathcal{K}_1} 0.$$

561 By construction, we conclude that $\lambda_u \leq 0$ is an eigenvalue of (4.9) with eigenfunction
562 $(\phi_1, \varphi_1) \in \text{Int } \mathcal{K}_1 \times (-\text{Int } Y_1^+)$. Hence $\Lambda_1 \leq \lambda_u \leq 0$. On the other hand, let $(\tilde{\phi}, \tilde{\varphi})$
563 be the eigenfunction of Λ_1 , then $\tilde{\varphi} \not\equiv 0$, since otherwise $(\Lambda, \tilde{\phi})$ is an eigenpair of
564 (4.11), whence $\Lambda \geq \mu_1 > 0$, contradictions. Therefore, $\tilde{\varphi} \not\equiv 0$ and $(\Lambda_1, \tilde{\varphi})$ furnishes an
565 eigenpair of (4.10). Thus $\Lambda_1 \geq \lambda_u$ as well. This completes the proof. \square

566 The linearized eigenvalue problem at semi-trivial steady state $(0, \tilde{v})$ is

$$(4.12) \quad \begin{cases} D_1 \phi_{xx} - \alpha_1 \phi_x + [g_1(\sigma_2) - d_1] \phi + \tilde{\Lambda} \phi = 0, & 0 < x < L, \\ D_2 \varphi_{xx} - \alpha_2 \varphi_x + [g_2(\sigma_2) - d_2] \varphi + \tilde{\Lambda} \varphi = \tilde{v} \sigma_2 g_2'(\sigma_2) [\int_0^x \phi(s) ds + \int_0^x \varphi(s) ds], & 0 < x < L, \\ D_1 \phi_x - \alpha_1 \phi = 0, & x = 0, L, \\ D_2 \varphi_x - \alpha_2 \varphi = 0, & x = 0, L, \end{cases}$$

568 where $\sigma_2(x) = e^{-k_0 x - \int_0^x \tilde{v}(s) ds}$. Let $\lambda_v = \lambda_1(D_1, \alpha_1, g_1(\sigma_2) - d_1)$ denote the principal
569 eigenvalue of the eigenvalue problem

$$(4.13) \quad \begin{cases} D_1 \phi_{xx} - \alpha_1 \phi_x + [g_1(\sigma_2) - d_1] \phi + \lambda \phi = 0, & 0 < x < L, \\ D_1 \phi_x - \alpha_1 \phi = 0, & x = 0, L. \end{cases}$$

571 It follows analogously that the stability properties of $(0, \tilde{v})$ are determined by λ_v .

572 PROPOSITION 4.6. *The problem (4.12) has a principal eigenvalue $\tilde{\Lambda}_1 \in \mathbb{R}$, in
573 the sense that $\tilde{\Lambda}_1 \leq \text{Re } \tilde{\Lambda}$ for all eigenvalues $\tilde{\Lambda}$ of (4.12) and that the corresponding
574 eigenfunction can be chosen in $\mathcal{K} \setminus \{(0, 0)\}$. Furthermore, (denote $Y_1^+ = C([0, L]; \mathbb{R}_+)$)
575 (a) If the principal eigenvalue λ_v of (4.13) is positive, then $\tilde{\Lambda}_1 > 0$.
576 (b) If the principal eigenvalue λ_v of (4.13) is non-positive, then $\tilde{\Lambda}_1 = \lambda_v \leq 0$ and
577 the corresponding eigenfunction can be chosen in $\text{Int } Y_1^+ \times (-\text{Int } \mathcal{K}_1)$.*

578 **4.3. Auxilliary Eigenvalue Lemmas.** In this subsection, we prove several
579 useful lemmas concerning the principal eigenvalue $\lambda_1(D, \alpha, h)$ of (4.3) with positive

580 eigenfunction ϕ_1 . It can be shown that λ_1 and ϕ_1 are smooth functions of α and D
 581 (see, e.g., [1, Lemma 1.2]).

582 We will assume additionally the following:

583 (A) $h(x) \in C^1([0, L])$ such that $h'(x) < 0$ in $[0, L]$.

584 Set $\psi_1 = e^{-(\alpha/D)x} \phi_1$. Then ψ_1 satisfies

$$585 \quad (4.14) \quad \begin{cases} D\psi_{1,xx} + \alpha\psi_{1,x} + h(x)\psi_1 + \lambda_1\psi_1 = 0, & 0 < x < L, \\ \psi_{1,x}(0) = \psi_{1,x}(L) = 0. \end{cases}$$

586 LEMMA 4.7. *If $h(x)$ satisfies (A), then $\psi_{1,x} < 0$ in $(0, L)$.*

587 *Proof.* Multiplying (4.14) by $e^{(\alpha/D)x}$, we rewrite the resulting equation as

$$588 \quad (4.15) \quad \begin{cases} D(e^{(\alpha/D)x}\psi_{1,x})_x + e^{(\alpha/D)x}\psi_1[h(x) + \lambda_1] = 0, & 0 < x < L, \\ \psi_{1,x}(0) = \psi_{1,x}(L) = 0. \end{cases}$$

589 Integrating (4.15) over $(0, L)$, we have

$$590 \quad \int_0^L e^{(\alpha/D)x}\psi_1[h(x) + \lambda_1] dx = 0,$$

591 which implies that $h(x) + \lambda_1$ changes sign in $(0, L)$. Since $h(x)$ is strictly decreasing
 592 in $(0, L)$, then there exists a unique $x_0 \in (0, L)$ such that $h(x) + \lambda_1 > 0$ for $0 < x < x_0$
 593 and $h(x) + \lambda_1 < 0$ for $x_0 < x < L$. Hence, by (4.15) we see that $(e^{(\alpha/D)x}\psi_{1,x})_x < 0$
 594 for $0 < x < x_0$ and $(e^{(\alpha/D)x}\psi_{1,x})_x > 0$ for $x_0 < x < L$. That is, $e^{(\alpha/D)x}\psi_{1,x}$ is strictly
 595 decreasing in $(0, x_0)$, and strictly increasing in (x_0, L) . Since $\psi_{1,x}(0) = \psi_{1,x}(L) = 0$,
 596 we have $\psi_{1,x} < 0$ in $(0, L)$. \square

LEMMA 4.8. *If $h(x)$ satisfies (A), then*

$$\frac{\partial\lambda_1}{\partial\alpha}(D, \alpha, h) > 0 \quad \text{for any } D > 0 \text{ and } \alpha \in \mathbb{R}.$$

597 The proof of Lemma 4.8 is similar to [13, Lemma 5.2], and we omit it here. The
 598 proof of the following Lemma 4.9 is similar to [13, Lemma 7.1] with some modifica-
 599 tions. For the sake of completeness, we give the proof here in detail.

600 LEMMA 4.9. *If $h(x)$ satisfies (A), then the following hold:*

601 (a) $\frac{\partial\lambda_1}{\partial D}(D, \alpha, h) > 0$ for $D > 0$ and $\alpha \leq 0$.

602 (b) *If $\alpha \geq h(0)L$ and $\lambda_1(D^*, \alpha, h) = 0$ for some $D^* > 0$, then $\frac{\partial\lambda_1}{\partial D}(D^*, \alpha, h) < 0$.*

603 *Proof.* Recall that λ_1 and ψ_1 are smooth functions of D . For simplicity of nota-
 604 tion, we denote $\frac{\partial\psi_1}{\partial D}$ by ψ'_1 , etc., where ψ_1 satisfies (4.14). Differentiating (4.14) with
 605 respect to D , we have

$$606 \quad (4.16) \quad \begin{cases} D\psi'_{1,xx} + \psi_{1,xx} + \alpha\psi'_{1,x} + h(x)\psi'_1 + \lambda'_1\psi_1 + \lambda_1\psi'_1 = 0, & 0 < x < L, \\ \psi'_{1,x}(0) = \psi'_{1,x}(L) = 0. \end{cases}$$

607 Multiplying (4.16) by $e^{(\alpha/D)x}\psi_1$ and integrating the resulting equation in $(0, L)$, we
 608 have

$$609 \quad -D \int_0^L e^{(\alpha/D)x}\psi'_{1,x}\psi_{1,x} dx + \int_0^L e^{(\alpha/D)x}\psi_{1,xx}\psi_1 dx + \int_0^L e^{(\alpha/D)x}h(x)\psi'_1\psi_1 dx \\ 610 \quad + \lambda'_1 \int_0^L e^{(\alpha/D)x}\psi_1^2 dx + \lambda_1 \int_0^L e^{(\alpha/D)x}\psi'_1\psi_1 dx = 0. \\ 611 \quad (4.17)$$

612 Similarly, multiplying (4.14) by $e^{(\alpha/D)x}\psi'_1$ and integrating it in $(0, L)$, we have
 613 (4.18)
$$-D \int_0^L e^{(\alpha/D)x}\psi'_{1,x}\psi_{1,x} dx + \int_0^L e^{(\alpha/D)x}h(x)\psi'_1\psi_1 dx + \lambda_1 \int_0^L e^{(\alpha/D)x}\psi'_1\psi_1 dx = 0.$$

614 It follows from (4.17) and (4.18) that

615 (4.19)
$$\lambda'_1 = \frac{-\int_0^L e^{(\alpha/D)x}\psi_{1,xx}\psi_1 dx}{\int_0^L e^{(\alpha/D)x}\psi_1^2 dx}.$$

616 By Lemma 4.7, we have $\psi_{1,x} < 0$ in $(0, L)$. Hence, if $\alpha \leq 0$, then

617 (4.20)
$$\begin{aligned} \int_0^L e^{(\alpha/D)x}\psi_{1,xx}\psi_1 dx &= - \int_0^L \psi_{1,x}(e^{(\alpha/D)x}\psi_1)_x dx \\ &= - \int_0^L e^{(\alpha/D)x}\psi_{1,x}[\psi_{1,x} + (\alpha/D)\psi_1] dx < 0. \end{aligned}$$

619 Thus $\lambda'_1 > 0$ for any $\alpha \leq 0$ and $D > 0$. This proves (a).

620 On the other hand, if $\lambda_1(D^*, \alpha, h) = 0$ for some $D^* > 0$, then the corresponding
 621 eigenfunction ψ_1 satisfies

622 (4.21)
$$\begin{cases} D^*\psi_{1,xx} + \alpha\psi_{1,x} + h(x)\psi_1 = 0, & 0 < x < L, \\ \psi_{1,x}(0) = \psi_{1,x}(L) = 0. \end{cases}$$

Multiplying (4.21) by $e^{(\alpha/D^*)x}$, and integrating over $(0, L)$, we have

$$\int_0^L h(x)\psi_1(x)e^{(\alpha/D^*)x} dx = 0.$$

623 Thus the decreasing function h must change sign, i.e. $h'(x) < 0, h(0) > 0$. Combining
 624 with $\psi_{1,x} < 0$, we have

625
$$\int_0^x h(s)\psi_1(s) ds < \int_0^x h(0)\psi_1(s) ds < h(0) \int_0^x \psi_1(0) ds < h(0)\psi_1(0)L.$$

626 Next, we integrate (4.21) in $(0, x)$, to get

627
$$D^*\psi_{1,x}(x) + \alpha\psi_1(x) = \alpha\psi_1(0) - \int_0^x h(s)\psi_1(s) ds > [\alpha - h(0)L]\psi_1(0) \geq 0,$$

628 provided that $\alpha \geq h(0)L$. By virtue of (4.20), we obtain

629
$$\int_0^L e^{(\alpha/D^*)x}\psi_{1,xx}\psi_1 dx = -\frac{1}{D^*} \int_0^L e^{(\alpha/D^*)x}\psi_{1,x}(D^*\psi_{1,x} + \alpha\psi_1) dx > 0.$$

630 It follows then from (4.19) that $\frac{\partial\lambda_1}{\partial D}(D^*, \alpha, h) < 0$. This proves (b). \square

631 **4.4. The Case $D_1 = D_2, \alpha_1 < \alpha_2$.** To investigate whether stronger or weaker
 632 advection is more beneficial for species to win the competition in the two-species
 633 phytoplankton model, we assume the only phenotypic difference between them is the
 634 advection rate. To be more precise, we assume $D_1 = D_2 \equiv D > 0, \alpha_1 < \alpha_2$. For the
 635 rest of this paper, we assume two phytoplankton species have the same growth rates
 636 and death rates, i.e., $g_1(\cdot) = g_2(\cdot) \equiv g(\cdot)$ and $d_1 = d_2 \equiv d$.

637 *Proof of Theorem 2.2.* By Theorem 4.1, it suffices to establish, for system (1.1)-(1.4) (and that $k_1 = k_2 = I_0 = 1$), the linear instability of $(0, \tilde{v})$, and the non-existence 638 of positive steady states.

640 **Step 1.** $(0, \tilde{v})$ is linearly unstable.

641 Recall that \tilde{v} is the unique positive solution to

$$642 \quad \begin{cases} D\tilde{v}_{xx} - \alpha_2\tilde{v}_x + [g(\sigma_2) - d]\tilde{v} = 0, & 0 < x < L, \\ D\tilde{v}_x - \alpha_2\tilde{v} = 0, & x = 0, L, \end{cases}$$

643 where $\sigma_2(x) = e^{-k_0x - \int_0^x \tilde{v}(s) ds}$. Since $\tilde{v} > 0$ is a positive eigenfunction of (4.3) with 644 $\alpha = \alpha_2$ and $h(x) = g(\sigma_2) - d$, we have $\lambda_1(D, \alpha_2, g(\sigma_2) - d) = 0$.

It follows from Proposition 4.6 that the stability of $(0, \tilde{v})$ is determined by the sign of the principal eigenvalue $\lambda_1(D, \alpha_1, g(\sigma_2) - d)$. Since $\alpha_1 < \alpha_2$, we may apply Lemma 4.8 to yield

$$\lambda_1(D, \alpha_1, g(\sigma_2) - d) < \lambda_1(D, \alpha_2, g(\sigma_2) - d) = 0.$$

645 Thus $(0, \tilde{v})$ is linearly unstable.

646 **Step 2.** The system (1.1)-(1.4) has no co-existence steady states.

647 Suppose to the contrary that (u^*, v^*) be a co-existence steady state of (1.1)-(1.4), 648 then we have

$$649 \quad \begin{cases} Du_{xx}^* - \alpha_1 u_x^* + [g(\sigma^*(x)) - d]u^* = 0, & 0 < x < L, \\ Dv_{xx}^* - \alpha_2 v_x^* + [g(\sigma^*(x)) - d]v^* = 0, & 0 < x < L, \\ Du_x^* - \alpha_1 u^* = 0, \text{ and } Dv_x^* - \alpha_2 v^* = 0, & x = 0, L, \end{cases}$$

where $\sigma^*(x) = \exp(-k_0x - \int_0^x [u^*(s) + v^*(s)] ds)$. Let $h(x) = g(\sigma^*(x)) - d$ so that $h'(x) < 0$. Since $u^*(x) > 0, v^*(x) > 0$, then

$$\lambda_1(D, \alpha_1, h) = \lambda_1(D, \alpha_2, h) = 0.$$

650 This is in contradiction with Lemma 4.8, which says that λ_1 is strictly monotone 651 increasing in α . Therefore, the system (1.1)-(1.4) has no co-existence steady state. \square

652 **4.5. The Case $D_1 < D_2$, $\alpha_1 = \alpha_2 \geq [g(1) - d]L$.** In this and the next subsection, 653 we explore the effect of diffusion on the outcome of competition. According to Lemma 654 4.9, the monotonicity of the principal eigenvalue $\lambda_1(D, \alpha, h)$ with respect to D also 655 depends on the advection rate α . Here, we first consider that both species have large 656 sinking rates, i.e., $\alpha_1 = \alpha_2 \equiv \alpha \geq [g(1) - d]L > 0$. (Note that $g(\cdot)$ satisfies (1.5) and 657 as we assume that the semi-trivial steady states exist, so we always have $g(1) - d > 0$.)

658 *Proof of Theorem 2.3.* By Theorem 4.1, it suffices to establish, for system (1.1)-(1.4), the linear instability of $(\tilde{u}, 0)$, and the non-existence of positive steady states.

660 **Step 1.** $(\tilde{u}, 0)$ is linearly unstable.

661 First, we observe as before from the equation satisfied by \tilde{u} that $\lambda_1(D_1, \alpha, g(\sigma_1) - d) = 0$, where $\sigma_1(x) = e^{-k_0x - \int_0^x \tilde{u}(s) ds}$.

662 Since $D_1 < D_2$ and $\alpha \geq [g(1) - d]L$, we may apply Lemma 4.9(b) to yield

$$\lambda_1(D_2, \alpha, g(\sigma_1) - d) < \lambda_1(D_1, \alpha, g(\sigma_1) - d) = 0.$$

663 It follows from Proposition 4.5 that $(\tilde{u}, 0)$ is linearly unstable.

664 **Step 2.** The system (1.1)-(1.4) has no co-existence steady states.

Suppose to the contrary that (u^*, v^*) is a co-existence steady state of (1.1)-(1.4), then we deduce as before,

$$\lambda_1(D_1, \alpha, g(\sigma^*) - d) = \lambda_1(D_2, \alpha, g(\sigma^*) - d) = 0,$$

665 where $\sigma^*(x) = \exp(-k_0 x - \int_0^x [u^*(s) + v^*(s)] ds)$. But this is in contradiction with
666 Lemma 4.9(b), which says that $D \mapsto \lambda_1(D, \alpha, g(\sigma^*) - d)$ has at most one positive
667 root. Therefore, the system (1.1)-(1.4) has no co-existence steady state. \square

668 **4.6. The Case $D_1 < D_2$, $\alpha_1 = \alpha_2 \leq 0$.** This subsection is devoted to studying
669 whether stronger or weaker diffusion is more beneficial when both species have buoyant
670 rates. Precisely, we assume that $D_1 < D_2$, $\alpha_1 = \alpha_2 \equiv \alpha \leq 0$.

671 *Proof of Theorem 2.4.* By Theorem 4.1, it suffices to establish, for system (1.1)-
672 (1.4), the linear instability of $(0, \tilde{v})$, and the non-existence of positive steady states.

673 **Step 1.** $(0, \tilde{v})$ is linearly unstable.

674 First, we observe as before from the equation satisfied by \tilde{v} that $\lambda_1(D_2, \alpha, g(\sigma_2) - d) = 0$, where $\sigma_2(x) = e^{-k_0 x - \int_0^x \tilde{v}(s) ds}$.

675 Since $D_1 < D_2$ and $\alpha \leq 0$, we may apply Lemma 4.9(a) to yield

$$\lambda_1(D_1, \alpha, g(\sigma_2) - d) < \lambda_1(D_2, \alpha, g(\sigma_2) - d) = 0.$$

676 It follows from Proposition 4.6 that $(0, \tilde{v})$ is linearly unstable.

677 **Step 2.** The system (1.1)-(1.4) has no co-existence steady states.

678 We omit the details here as this is similar to Step 2 of the proofs of Theorems 2.3,
679 where we use part (a) of Lemma 4.9 instead of part (b). This completes the proof. \square

680 **5. Discussion and Numerical Results.** We investigate a nonlocal reaction-
681 diffusion-advection system modeling the growth of two competing phytoplankton
682 species in a eutrophic environment, where nutrients are in abundance and the species
683 are limited by light only for their metabolism. We first demonstrate that the system
684 does not preserve the competitive order in the pointwise sense (Remark 3.10). We
685 introduce a special cone \mathcal{K} involving cumulative distributions of the population den-
686 sities, and a generalized notion of super- and subsolutions of (1.1)-(1.4), where the
687 differential inequalities hold in the sense of the cone \mathcal{K} . A comparison principle is
688 then established for the super- and subsolutions, which implies the monotonicity of
689 the semiflow of (1.1)-(1.4) with respect to the cone \mathcal{K} (Theorem 2.1). From a theo-
690 retical point of view, this paper introduces a new class of reaction-diffusion models
691 with order-preserving property, which may be of independent interest [35].

692 A first application of the monotonicity result yields a simple proof of the existence
693 and global attractivity of the unique positive steady state (or time-periodic solution)
694 to the single species problem (Proposition 3.11). A second application concerns the
695 dynamics of two competing phytoplankton species, as modeled by (1.1)-(1.4), in which
696 sufficient conditions for local (Propositions 4.5 and 4.6) and global stability of semi-
697 trivial steady states (Theorems 2.2-2.4) are obtained.

698 Consider system (1.1)-(1.4) and fix $D_1 < D_2$ and $\alpha_1 = \alpha_2 \equiv \alpha$. Theorems
699 2.3 and 2.4 say that $(\tilde{u}, 0)$ is globally asymptotically stable for $\alpha = 0$, and $(0, \tilde{v})$ is
700 globally asymptotically stable for $\alpha = [g(1) - d]L$, which means there is an exchange
701 of stability between the semi-trivial steady states as α varies from 0 to $[g(1) - d]L$.
702 In this particular case, we conjecture that there exist two constants α_{\min} and α_{\max}
703 such that the following statements hold.

- 704 • When $\alpha \leq \alpha_{\min}$, the semi-trivial steady state $(\tilde{u}, 0)$ is globally asymptotically
705 stable.

- 706 • When $\alpha_{min} < \alpha < \alpha_{max}$, there exists a unique positive steady state (u^*, v^*)
 707 which is globally asymptotically stable.
 708 • When $\alpha \geq \alpha_{max}$, the semi-trivial steady state $(0, \tilde{v})$ is globally asymptotically
 709 stable.

710 In the following, we present some numerical result in support of this conjecture. See
 711 Figure 2.

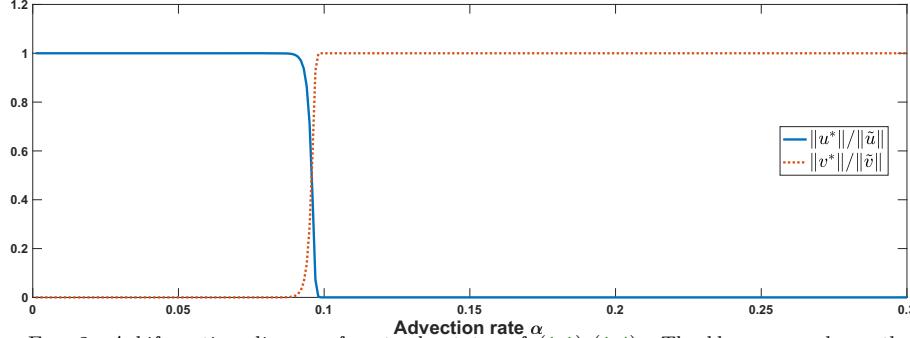


FIG. 2. A bifurcation diagram for steady states of (1.1)-(1.4). The blue curve shows the ratio $\|u^*\|_{L^1}/\|\tilde{u}\|_{L^1}$, and the red curve shows the ratio $\|v^*\|_{L^1}/\|\tilde{v}\|_{L^1}$ as α varies from 0 to 0.3, where (u^*, v^*) is the stable steady state, and $(\tilde{u}, 0)$ and $(0, \tilde{v})$ are the two semi-trivial steady states. The parameters are chosen as $D_1 = 1$, $D_2 = 5$, $d_1 = d_2 = 0.001$, $g_1(I) = g_2(I) = \frac{mI}{a+I}$, $m = 1$, $a = 10$, $I_0 = 1$, $k_0 = k_1 = k_2 = 0.001$, $L = 100$.

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