

# MONOTONICITY AND GLOBAL DYNAMICS OF A NONLOCAL TWO-SPECIES PHYTOPLANKTON MODEL\*

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**Abstract.** We investigate a nonlocal reaction-diffusion-advection system modeling the population dynamics of two competing phytoplankton species in a eutrophic environment, where nutrients are in abundance and the species are limited by light only for their metabolism. We first demonstrate that the system does not preserve the competitive order in the pointwise sense. Then we introduce a special cone  $\mathcal{K}$  involving the cumulative distributions of the population densities, and a generalized notion of super- and subsolutions of the nonlocal competition system where the differential inequalities hold in the sense of the cone  $\mathcal{K}$ . A comparison principle is then established for such super- and subsolutions, which implies the monotonicity of the underlying semiflow with respect to the cone  $\mathcal{K}$  (Theorem 2.1). As application, we study the global dynamics of the single species system and the competition system. The latter has implications for the evolution of movement for phytoplankton species.

**Key words.** Phytoplankton; competition for light; nonlocal reaction-diffusion equations; monotone dynamical system.

**AMS subject classifications.** 35B51, 35K57, 47H07, 92D25

**1. Introduction.** Phytoplankton are microscopic plant-like photosynthetic organisms that drift in the water columns of lakes and oceans. They grow abundantly around the globe and are the foundation of the marine food chain. Since they transport significant amounts of atmospheric carbon dioxide into the deep oceans, they play a crucial role in climate dynamics. Nutrients and light are the essential resources for the growth of phytoplankton. There are three possible ways for phytoplankton to compete for nutrients and light. At one extreme, in oligotrophic ecosystems with an ample supply of light, species compete for limiting nutrients [22, 27]. At the other extreme, in eutrophic ecosystems with ample nutrient supply, species compete for light [8, 16, 17, 33]. In some ecosystems of intermediate conditions, they compete for both nutrients and light [3, 4, 18, 21, 36]. In the water column, phytoplankton diffuse by water turbulence, and also sink or buoy, depending on whether they are heavier than water or not [8].

In this paper, we study the two-species nonlocal reaction-diffusion-advection system proposed by Huisman et al. [16, 18]. The system models the growth of phytoplankton species in a eutrophic vertical water column, where the species is limited by light only for their metabolism. Consider a water column with unit cross-sectional area and with two phytoplankton species. Let  $x$  denote the depth within the water column where  $x$  varies from 0 (the top) to  $L$  (the bottom), and let  $u(x, t), v(x, t)$  stand for the population densities of two phytoplankton species at the location  $x$  and time

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38  $t$ , respectively. The following system of reaction-diffusion-advection equations was  
 39 proposed in [16] to describe the population dynamics of two phytoplankton species:

$$40 \quad (1.1) \quad \begin{cases} u_t = D_1 u_{xx} - \alpha_1 u_x + [g_1(I(x, t)) - d_1]u, & 0 < x < L, \ t > 0, \\ v_t = D_2 v_{xx} - \alpha_2 v_x + [g_2(I(x, t)) - d_2]v, & 0 < x < L, \ t > 0, \end{cases}$$

41 with no-flux boundary conditions

$$42 \quad (1.2) \quad \begin{cases} D_1 u_x(x, t) - \alpha_1 u(x, t) = 0, & x = 0, L, \ t > 0, \\ D_2 v_x(x, t) - \alpha_2 v(x, t) = 0, & x = 0, L, \ t > 0, \end{cases}$$

43 and initial conditions

$$44 \quad (1.3) \quad u(x, 0) = u_0(x) \geq, \neq 0, \quad v(x, 0) = v_0(x) \geq, \neq 0, \quad 0 \leq x \leq L,$$

45 where for  $i = 1, 2$ ,  $D_i > 0$  is the diffusion coefficient,  $\alpha_i \in \mathbb{R}$  is the sinking ( $\alpha_i > 0$ )  
 46 or buoyant ( $\alpha_i < 0$ ) velocity,  $d_i > 0$  is the death rate,  $g_i(I)$  represents the specific  
 47 growth rate of phytoplankton species as a function of light intensity  $I(x, t)$ .

48 Light intensity is decreasing with depth due to light absorption via phytoplankton  
 49 and water. By the Lambert-Beer law [23], the light intensity  $I(x, t)$  is given by

$$50 \quad (1.4) \quad I(x, t) = I_0 \exp \left( -k_0 x - \int_0^x [k_1 u(s, t) + k_2 v(s, t)] ds \right),$$

51 where  $I_0 > 0$  is the incident light intensity,  $k_0 > 0$  is the background turbidity  
 52 that summarizes light absorption by all non-phytoplankton components, and  $k_i$  is  
 53 the absorption coefficient of the corresponding phytoplankton species. In this model  
 54 ample nutrient supply is assumed so that the phytoplankton growth is only limited  
 55 by the light availability. We assume that  $g_i(I)$  is a smooth function satisfying

$$56 \quad (1.5) \quad g_i(0) = 0 \quad \text{and} \quad g'_i(I) > 0 \quad \text{for} \quad I \geq 0.$$

A typical example of  $g_i(I)$  takes the Michaelis-Menten form

$$g_i(I) = \frac{m_i I}{a_i + I},$$

57 where  $m_i > 0$  is the maximal growth rate and  $a_i > 0$  is the half saturation constant.

58 Most existing mathematical literatures on phytoplankton are focused on a single  
 59 species. The single species model was considered in [33] for the self-shading case (i.e.  
 60  $k_0 = 0$ ) and infinite long water column ( $L = \infty$ ). The existence, uniqueness and  
 61 global stability of the steady state are established in [20, 33]. It is shown in [24] that  
 62 the self-shading model with any finite water column depth has a stable positive steady  
 63 state, which means that the self-shading model has no critical water column depth  
 64 beyond which the phytoplankton cannot persist.

65 For the case  $k_0 > 0$ , it is illustrated in [8] that the condition for phytoplankton  
 66 bloom development can be characterized by critical water column depth and some  
 67 critical values of the vertical turbulent diffusion coefficient. Du and Hsu [5] studied  
 68 both single and two species competing for light with no advection. For the single  
 69 species model, the existence, uniqueness, and global attractivity of a positive equilib-  
 70 rium was established. Hsu and Lou [13] analyzed the critical death rate, critical water  
 71 column depth, critical sinking or buoyant coefficient and critical turbulent diffusion

rate. Du and Mei [7] investigated the global dynamics of the single species model for the case  $D = D(x)$ ,  $\alpha = \alpha(x)$  and the asymptotic profiles of the positive steady states for small or large diffusion and deep water column when  $D, \alpha$  are constants. Peng and Zhao [31, 32] considered the effect of time-periodic light intensity  $I_0$  at the surface, due to diurnal light cycle and seasonal changes. Ma and Ou [28] further studied the model in [31, 32] and assume that  $D(t), \alpha(t)$  are time periodic functions. They obtained the uniqueness and the global attractivity of the positive periodic solution of the single species model, when it exists.

Du et al. [6] studied the effect of photoinhibition on the single phytoplankton species, and they found that, in contrast to the case of no photoinhibition, where at most one positive steady state can exist, the model with photoinhibition possesses at least two positive steady states in certain parameter ranges. Hsu et al. [14] examined the dynamics of a single species under the assumption that the amount of light absorbed by individuals is proportional to cell size, which varies for populations that reproduced by simple cell division into two equal-sized daughter cells.

Although many mathematical theories have been developed for single species phytoplankton model, there are very few results for two or more phytoplankton species competing for light. The existence of positive steady state and uniform persistence for two-species model were proved in [5], where there is no sinking or buoyancy. In [29], Mei and Zhang studied a nonlocal reaction-diffusion-advection system modeling the growth of multiple competitive phytoplankton species and they found that when the diffusion of the system is large, there are no positive steady states, and when the diffusion is not large, there exists at least one positive steady state under proper conditions.

Unlike two-species Lotka-Volterra competition model with diffusion, one main difficulty for system (1.1)-(1.4) is the lack of comparison principle, i.e.

$$\begin{aligned} u_1(x, 0) \leq u_2(x, 0), \quad v_1(x, 0) \geq v_2(x, 0) \quad \forall x \in [0, L] \\ \not\Rightarrow u_1(x, t) \leq u_2(x, t), \quad v_1(x, t) \geq v_2(x, t) \quad \forall (x, t) \in [0, L] \times (0, \infty), \end{aligned}$$

due to the nonlocal nature of the nonlinearity. See Remark 3.10.

For order-preserving properties in the single species model, Shigesada and Okubo [33] observed that the cumulative distribution function  $U(x, t) := \int_0^x u(s, t) ds$  satisfies a single reaction-diffusion equation without nonlocal terms. Subsequently, Ishii and Takagi [20] showed that the flow retains the natural order in  $U$ . For a related model with a water column of infinite depth, they made use of this fact to obtain a complete classification of the long-time behavior of the population. This fact was used again in Du and Hsu [5] to determine the long-time dynamics for a single species model with finite water depth. More recently, Ma and Ou [28] established the comparison principle for  $U$  in the single species model.

For the competition model, we will show, by adapting arguments due to Du and Hsu [5] and Ma and Ou [28], that the cumulative distribution functions

$$(U(x, t), V(x, t)) = \left( \int_0^x u(s, t) ds, \int_0^x v(s, t) ds \right)$$

satisfy a nonlocal, strongly coupled system, with non-standard boundary condition (see (3.3)), and that the resulting system has the strong order-preserving property.

Our main result (Theorem 2.1) says that system (1.1)-(1.4) forms a strongly monotone dynamical system with respect to the order induced by the special cone

$\mathcal{K} = \mathcal{K}_1 \times (-\mathcal{K}_1)$ , where

$$(1.6) \quad \mathcal{K}_1 = \left\{ \phi \in C([0, L], \mathbb{R}) : \int_0^x \phi(s) ds \geq 0 \text{ for } x \in (0, L] \right\}.$$

The new features of this paper can be described as follows: First, Theorem 2.1 is the first monotonicity result for the nonlocal competition system involving two phytoplankton species. Second, the definition of the relevant cone  $\mathcal{K}$  facilitates the connection with general theory of monotone dynamical systems. Third, generalized notion of super- and subsolutions (see Definition 3.2), which is new even for the case of single species, are given. They can potentially be used to obtain qualitative properties of solutions for the nonlocal system (1.1)-(1.4).

The rest of the paper is organized as follows: In Section 2, we state our main results. In Section 3, we first introduce the notion of super- and subsolutions of (1.1)-(1.4) with respect to the cone  $\mathcal{K}$ , and establish the comparison principle for the super- and subsolutions. Then we apply the monotonicity result to establish the global dynamics of the single species model in a general setting. Section 4 is devoted to the spectral analysis of semi-trivial steady states, and the global dynamics of system (1.1)-(1.4) are established for three different biological scenarios. In Section 5, we present some numerical results and discussion.

**2. Main Results.** Let  $\mathbf{X}$  be a Banach space over  $\mathbb{R}$ . We call a subset  $K \subset \mathbf{X}$  a *cone* if (i)  $K$  is convex, (ii)  $\mu K \subset K$  for all  $\mu \geq 0$ , and (iii)  $K \cap (-K) = \{0\}$ . A cone  $K$  is said to be *solid* if it has nonempty interior. Furthermore, for  $x, y \in \mathbf{X}$ , we write  $x \leq_K y$ ,  $x <_K y$  and  $x \ll_K y$  if  $y - x \in K$ ,  $y - x \in K \setminus \{0\}$  and  $y - x \in \text{Int } K$  respectively.

Let  $\mathcal{K}_1$  be given by (1.6). It is straightforward to verify that  $\mathcal{K}_1$  is a solid cone in the Banach space  $C([0, L]; \mathbb{R})$  with interior

$$\text{Int } \mathcal{K}_1 = \left\{ \phi \in C([0, L]; \mathbb{R}) : \phi(0) > 0, \int_0^x \phi(s) ds > 0 \text{ for } x \in (0, L] \right\}.$$

Let  $\mathcal{K} = \mathcal{K}_1 \times (-\mathcal{K}_1)$ . Then  $\mathcal{K}$  is likewise a solid cone in the Banach space  $C([0, L]; \mathbb{R}^2)$  with interior given by  $\text{Int } \mathcal{K} = \text{Int } \mathcal{K}_1 \times (-\text{Int } \mathcal{K}_1)$ . The cone  $\mathcal{K}$  induces the partial order relations  $\leq_{\mathcal{K}}$ ,  $<_{\mathcal{K}}$  and  $\ll_{\mathcal{K}}$  in the usual way.

We shall prove that (1.1)-(1.4) is a strongly monotone dynamical system with respect to the order induced by the cone  $\mathcal{K}$ .

**THEOREM 2.1.** *Suppose  $\{(u_i, v_i)\}_{i=1,2}$  are non-negative solutions of (1.1)-(1.4) such that  $u_2(\cdot, 0) \geq, \neq 0$  and  $v_1(\cdot, 0) \geq, \neq 0$  and*

$$(u_1(\cdot, 0), v_1(\cdot, 0)) <_{\mathcal{K}} (u_2(\cdot, 0), v_2(\cdot, 0)).$$

*Then  $(u_1(\cdot, t), v_1(\cdot, t)) \ll_{\mathcal{K}} (u_2(\cdot, t), v_2(\cdot, t))$  for all  $t > 0$ .*

By Theorem 2.1, system (1.1)-(1.4) is a strongly monotone dynamical system on  $C([0, L]; \mathbb{R}_+^2)$  with respect to the order generated by  $\mathcal{K}$ , which together with the theory of strongly monotone dynamical systems [2, 12, 15, 25, 34, 37], provides a useful tool to investigate the global dynamics of two-species system (1.1)-(1.4). As a by-product of our monotonicity result, we also generalize the existing results for single species (see Subsection 3.2) and give a simple proof based on monotonicity arguments and the concept of subhomogeneous mappings.

As application, we turn our attention to the effects of diffusion and advection on the global dynamics of (1.1)-(1.4).

THEOREM 2.2. *If  $D_1 = D_2$ ,  $\alpha_1 < \alpha_2$ ,  $g_1 = g_2$ ,  $d_1 = d_2$ , and that both semi-trivial steady states exist, then the first species  $u$  drives the second species  $v$  to extinction, regardless of initial condition.*

Theorem 2.2 shows that the competitor with smaller advection rate has competitive advantages, i.e., *smaller advection rate is selected*. By the Lambert-Beer law, the deeper the water column, the weaker the light intensity. Therefore, it is more advantageous for phytoplankton species to move up.

THEOREM 2.3. *If  $D_1 < D_2$ ,  $\alpha_1 = \alpha_2 \geq [g(1) - d]L$ ,  $g_1 = g_2$ ,  $d_1 = d_2$ , and that both semi-trivial steady states exist, then the faster diffuser  $v$  drives the slower diffuser  $u$  to extinction, regardless of initial condition.*

Theorem 2.3 implies that if sinking rate is large, competitor with faster diffusion will always displace the slower one, i.e., *faster diffuser wins*. Intuitively, when both species are sinking with equal and large velocity, faster diffusion can counter balance the tendency to sink and provide individuals with better access to light.

THEOREM 2.4. *If  $D_1 < D_2$ ,  $\alpha_1 = \alpha_2 \leq 0$ ,  $g_1 = g_2$ ,  $d_1 = d_2$ , and that both semi-trivial steady states exist, then the slower diffuser  $u$  drives faster diffuser  $v$  to extinction, regardless of initial condition.*

Theorem 2.4 suggests that if the phytoplankton species are buoyant, the competitor with slower diffusion rate will always displace the faster one, i.e., *slower diffusion rate will be selected*. This is in sharp contrast to Theorem 2.3. The reason for this result is that when the phytoplankton are buoyant, turbulent diffusion actually displaces individuals from the top of the water column, where the light intensity is the strongest.

**3. A General Model with Spatio-Temporally Varying Coefficients.** We shall study a generalized version of system (1.1)-(1.4), which allows coefficients to vary explicitly with both space and time. We formulate the nonlocal reaction-diffusion-advection model as follows:

$$(3.1) \quad \begin{cases} u_t = (D_1 u_x - \alpha_1 u)_x + f_1(x, t, \int_0^x u(s, t) ds, \int_0^x v(s, t) ds)u, & 0 < x < L, t > 0, \\ v_t = (D_2 v_x - \alpha_2 v)_x + f_2(x, t, \int_0^x u(s, t) ds, \int_0^x v(s, t) ds)v, & 0 < x < L, t > 0, \\ D_1 u_x - \alpha_1 u = 0, & x = 0, L, t > 0, \\ D_2 v_x - \alpha_2 v = 0, & x = 0, L, t > 0, \\ u(x, 0) = u_0(x) \geq, \neq 0, v(x, 0) = v_0(x) \geq, \neq 0, & 0 \leq x \leq L, \end{cases}$$

where, for  $i = 1, 2$ ,  $D_i = D_i(x, t) > 0$ ,  $\alpha_i = \alpha_i(x, t)$ , and the functions  $f_i(x, t, p, q)$  are smooth and satisfy

$$(H) \quad \frac{\partial f_i}{\partial p} < 0, \frac{\partial f_i}{\partial q} < 0 \text{ and } \frac{\partial f_i}{\partial x} \leq 0 \quad \text{for all } x \in [0, L] \text{ and } t, p, q \geq 0.$$

The assumption holds, e.g. when  $f_i(x, t, p, q) = g_i(I_0 \exp(-k_0 x - k_1 p - k_2 q)) - d_i(x, t)$  such that  $g_i$  is non-decreasing, and  $d_i$  is non-decreasing in  $x$ . In particular, it includes (1.1)-(1.4), and the previous works [5, 29] as particular cases.

**3.1. Strong Monotonicity of (3.1).** This subsection is devoted to proving the monotonicity of system (3.1) with respect to the order induced by cone  $\mathcal{K}$  under the assumption (H). First, we state the following standard result (see, e.g. [10, Ch. 3]).

PROPOSITION 3.1. *For continuous, non-negative initial data  $(u_0(x), v_0(x))$ , system (3.1) has a unique solution*

$$(u, v) \in C([0, \infty); C([0, L]; \mathbb{R}_+^2)) \cap C^1((0, \infty); C^\infty([0, L]; \mathbb{R}_+^2)),$$

192 which depends continuously on initial data. Moreover, if  $u_0(x) \not\equiv 0$ , (resp.  $v_0(x) \not\equiv 0$ ),  
 193 then  $u(x, t) > 0$  (resp.  $v(x, t) > 0$ ) for  $(x, t) \in [0, L] \times (0, \infty)$ .

194 Next, we define the following super- and subsolution concepts for (3.1). Note that  
 195 the differential inequalities appearing below are to be understood in the sense of cone  
 196  $\mathcal{K}$  for each time  $t$ . These inequalities hold, in particular, if the differential inequalities  
 197 hold in the pointwise sense everywhere.

DEFINITION 3.2. We say that

$$(\bar{u}, \bar{v}), (\underline{u}, \underline{v}) \in C([0, T]; C([0, L]; \mathbb{R}_+^2)) \cap C^1((0, T]; C^\infty([0, L]; \mathbb{R}_+^2))$$

198 form a pair of super- and subsolutions of (3.1) in the interval  $[0, T]$ , if

$$199 \quad (3.2) \quad \begin{cases} \bar{u}_t \geq_{\mathcal{K}_1} (D_1 \bar{u}_x - \alpha_1 \bar{u})_x + f_1(x, t, \int_0^x \bar{u}(s, t) ds, \int_0^x \bar{v}(s, t) ds) \bar{u}, & 0 < t \leq T, \\ \bar{v}_t \leq_{\mathcal{K}_1} (D_2 \bar{v}_x - \alpha_2 \bar{v})_x + f_2(x, t, \int_0^x \bar{u}(s, t) ds, \int_0^x \bar{v}(s, t) ds) \bar{v}, & 0 < t \leq T, \\ \underline{u}_t \leq_{\mathcal{K}_1} (D_1 \underline{u}_x - \alpha_1 \underline{u})_x + f_1(x, t, \int_0^x \underline{u}(s, t) ds, \int_0^x \underline{v}(s, t) ds) \underline{u}, & 0 < t \leq T, \\ \underline{v}_t \geq_{\mathcal{K}_1} (D_2 \underline{v}_x - \alpha_2 \underline{v})_x + f_2(x, t, \int_0^x \underline{u}(s, t) ds, \int_0^x \underline{v}(s, t) ds) \underline{v}, & 0 < t \leq T, \\ D_1 \bar{u}_x - \alpha_1 \bar{u} \leq 0 \leq D_1 \underline{u}_x - \alpha_1 \underline{u}, & x = 0, 0 < t \leq T, \\ D_1 \bar{u}_x - \alpha_1 \bar{u} \geq 0 \geq D_1 \underline{u}_x - \alpha_1 \underline{u}, & x = L, 0 < t \leq T, \\ D_2 \bar{v}_x - \alpha_2 \bar{v} \leq 0 \leq D_2 \underline{v}_x - \alpha_2 \underline{v}, & x = 0, 0 < t \leq T, \\ D_2 \bar{v}_x - \alpha_2 \bar{v} \geq 0 \geq D_2 \underline{v}_x - \alpha_2 \underline{v}, & x = L, 0 < t \leq T, \\ (\bar{u}(\cdot, 0), \bar{v}(\cdot, 0)) \geq_{\mathcal{K}} (\underline{u}(\cdot, 0), \underline{v}(\cdot, 0)). \end{cases}$$

200 The main result of this section is

THEOREM 3.3. Assume that  $f_1, f_2$  satisfy **(H)**. Let  $(\bar{u}, \bar{v})$  and  $(\underline{u}, \underline{v})$  be a pair of super- and subsolutions of (3.1) in the interval  $[0, T]$ . If  $\underline{u} > 0$  and  $\underline{v} > 0$  in  $[0, L] \times [0, T]$ , then

$$(\bar{u}(\cdot, t), \bar{v}(\cdot, t)) \geq_{\mathcal{K}} (\underline{u}(\cdot, t), \underline{v}(\cdot, t)) \quad \text{for } 0 \leq t \leq T.$$

Moreover, if there exists  $t_0 \in (0, T]$  such that  $\bar{u} > 0$  and  $\bar{v} > 0$  in  $[0, L] \times (0, t_0]$ , and

$$(\bar{u}(\cdot, t_0) - \underline{u}(\cdot, t_0), \bar{v}(\cdot, t_0) - \underline{v}(\cdot, t_0)) \notin \text{Int } \mathcal{K},$$

201 then  $(\bar{u}(x, t), \bar{v}(x, t)) \equiv (\underline{u}(x, t), \underline{v}(x, t))$  for  $x \in [0, L]$  and  $0 \leq t \leq t_0$ .

202 A direct consequence of Theorem 3.3 is the strong monotonicity of the continuous  
 203 semiflow generated by (3.1). It includes Theorem 2.1 as a particular case.

COROLLARY 3.4. Assume that  $f_1, f_2$  satisfy **(H)**. Suppose  $\{(u_i, v_i)\}_{i=1,2}$  are two non-negative solutions of (3.1), such that  $u_1(\cdot, 0) \geq, \neq 0, v_2(\cdot, 0) \geq, \neq 0$ , and

$$(u_1(\cdot, 0), v_1(\cdot, 0)) >_{\mathcal{K}} (u_2(\cdot, 0), v_2(\cdot, 0)).$$

204 Then  $(u_1(\cdot, t), v_1(\cdot, t)) \gg_{\mathcal{K}} (u_2(\cdot, t), v_2(\cdot, t))$  for all  $t > 0$ .

205 The proof is postponed to later in the section.

206 To show Theorem 3.3, we consider the cumulative distribution functions

$$207 \quad U(x, t) = \int_0^x u(s, t) ds, \quad V(x, t) = \int_0^x v(s, t) ds.$$

208 Then  $U(0, t) \equiv 0, V(0, t) \equiv 0$  for  $t \geq 0$ , and  $U_x(x, t) = u(x, t), V_x(x, t) = v(x, t)$ . In  
 209 this way, (3.1) is transformed into the following *strongly coupled, non-local* system of

210  $(U, V)$  (see also [28] for the single species case):

$$211 \quad (3.3) \quad \begin{cases} U_t = D_1 U_{xx} - \alpha_1 U_x + G_1[U, V, U_x, V_x], & 0 < x < L, t > 0, \\ V_t = D_2 V_{xx} - \alpha_2 V_x + G_2[U, V, U_x, V_x], & 0 < x < L, t > 0, \\ U(0, t) = 0, \quad D_1 U_{xx}(L, t) - \alpha_1 U_x(L, t) = 0, & t > 0, \\ V(0, t) = 0, \quad D_2 V_{xx}(L, t) - \alpha_2 V_x(L, t) = 0, & t > 0, \\ U(x, 0) = \int_0^x u_0(s) ds = U_0(x), & 0 \leq x \leq L, \\ V(x, 0) = \int_0^x v_0(s) ds = V_0(x), & 0 \leq x \leq L, \end{cases}$$

212 where, letting  $F_1(x, t, U, V) = \int_0^U f_1(x, t, z, V) dz$ ,  $F_2(x, t, U, V) = \int_0^V f_2(x, t, U, z) dz$ ,

$$\begin{aligned} 213 & \quad G_1[U, V, U_x, V_x](x, t) \\ 214 & = \int_0^x f_1\left(s, t, \int_0^s u(y, t) dy, \int_0^s v(y, t) dy\right) u(s, t) ds \\ 215 & = \int_0^x f_1\left(s, t, U(s, t), V(s, t)\right) U_x(s, t) ds \\ 216 & = \int_0^x \left\{ \frac{d}{ds} [F_1(s, t, U(s, t), V(s, t))] - \frac{\partial F_1}{\partial x}(s, t, U(s, t), V(s, t)) \right. \\ 217 & \quad \left. - \frac{\partial F_1}{\partial V}(s, t, U(s, t), V(s, t)) V_x(s, t) \right\} ds \\ 218 & = F_1(x, t, U(x, t), V(x, t)) - \int_0^x \frac{\partial F_1}{\partial x}(s, t, U(s, t), V(s, t)) ds \\ 219 \quad (3.4) & \quad - \int_0^x \frac{\partial F_1}{\partial V}(s, t, U(s, t), V(s, t)) V_x(s, t) ds \end{aligned}$$

220 and

$$\begin{aligned} 221 & \quad G_2[U, V, U_x, V_x](x, t) \\ 222 & = \int_0^x f_2\left(s, t, \int_0^s u(y, t) dy, \int_0^s v(y, t) dy\right) v(s, t) ds \\ 223 & = \int_0^x f_2\left(s, t, U(s, t), V(s, t)\right) V_x(s, t) ds \\ 224 & = F_2(x, t, U(x, t), V(x, t)) - \int_0^x \frac{\partial F_2}{\partial x}(s, t, U(s, t), V(s, t)) ds \\ 225 \quad (3.5) & \quad - \int_0^x \frac{\partial F_2}{\partial U}(s, t, U(s, t), V(s, t)) U_x(s, t) ds. \end{aligned}$$

For (3.3), we define the Banach space

$$X_1 = \{\phi \in C^1([0, L], \mathbb{R}) : \phi(0) = 0\}$$

with the usual  $C^1$  norm. The usual cone  $P_1$  in  $X_1$  is

$$P_1 = \{\phi \in X_1 : \phi(x) \geq 0 \text{ for } x \in [0, L]\},$$

with interior

$$\text{Int } P_1 = \{\phi \in X_1 : \phi'(0) > 0, \phi(x) > 0 \text{ for } x \in (0, L]\}.$$



Let  $X = X_1 \times X_1$ , and  $P = P_1 \times (-P_1)$ . Then  $P$  is a cone in  $X$  with interior given by  $\text{Int } P = \text{Int } P_1 \times (-\text{Int } P_1)$ . The cone  $P$  generates the partial order relations  $\leq_P, <_P$  and  $\ll_P$  on  $X$ .

By construction, the solutions  $(U, V)$  of (3.3) live in the convex set  $E = E_1 \times E_1$ , where

$$E_1 = \{\phi \in C^1([0, L]) : \phi(0) = 0, \text{ and } \phi'(x) \geq 0 \text{ for } x \in [0, L]\}.$$

From now on we assume the initial data of (3.3) to be in  $E$ . Under this assumption, the existence and uniqueness of the solution  $(U(x, t), V(x, t))$  can be derived from those of  $(u(x, t), v(x, t))$ .

DEFINITION 3.5. *We say that*

$$(\bar{U}, \underline{V}), (\underline{U}, \bar{V}) \in C([0, T]; E) \cap C^1((0, T]; C^\infty([0, L]; \mathbb{R}_+^2))$$

form a pair of super- and subsolutions of (3.3) in the interval  $[0, T]$ , if the derivatives  $(\bar{u}, \underline{v}) = (\bar{U}_x, \underline{V}_x)$  and  $(\underline{u}, \bar{v}) = (\underline{U}_x, \bar{V}_x)$  form a pair of super- and subsolutions of (3.1) in the interval  $[0, T]$ , in the sense of Definition 3.2.

We now prove a strong maximum principle for the system (3.3), which is the key to proving the strong monotonicity of (3.3).

LEMMA 3.6. *Assume that  $f_1, f_2$  satisfy (H). Let  $(\bar{U}, \underline{V})$  and  $(\underline{U}, \bar{V})$  be a pair of super- and subsolutions of (3.3) in the interval  $[0, t^*]$  for some  $t^* > 0$ , so that*

$$(3.6) \quad \bar{U}_x(x, t) > 0 \quad \text{and} \quad \bar{V}_x(x, t) > 0 \quad \text{for } 0 \leq x \leq L, \text{ and } 0 < t \leq t^*,$$

and

$$\underline{U}(x, t) \leq \bar{U}(x, t), \quad \bar{V}(x, t) \geq \underline{V}(x, t) \quad \text{for } 0 \leq x \leq L, \text{ and } 0 \leq t \leq t^*.$$

If one of the following holds:

(a)  $\underline{U}(x^*, t^*) = \bar{U}(x^*, t^*)$  or  $\underline{V}(x^*, t^*) = \bar{V}(x^*, t^*)$  for some  $x^* \in (0, L]$ ;

(b)  $(\bar{U} - \underline{U})_x(0, t^*) = 0$  or  $(\bar{V} - \underline{V})_x(0, t^*) = 0$ ,

then

$$(3.7) \quad (\underline{U}(x, t), \bar{V}(x, t)) \equiv (\bar{U}(x, t), \underline{V}(x, t)) \quad \text{for } 0 \leq x \leq L, \text{ } 0 \leq t \leq t^*.$$

*Proof.* In the following we improve upon the arguments of [28] to prove the strong maximum principle for (3.3). We first consider the case when (a) holds. For definiteness assume that  $\underline{U}(x^*, t^*) = \bar{U}(x^*, t^*)$  for some  $x^* \in (0, L]$ . Denote

$$W(x, t) = \bar{U}(x, t) - \underline{U}(x, t).$$

Then by (3.2),

$$(\bar{u} - \underline{u})_t - [D_1(\bar{u} - \underline{u})_x + \alpha_1(\bar{u} - \underline{u})]_x \geq \kappa_1 f_1(x, t, \bar{U}(x, t), \underline{V}(x, t)) - f_1(x, t, \underline{U}(x, t), \bar{V}(x, t))$$

Fixing  $t$ , and integrating the above from 0 to  $x$ , we have, in terms of  $W$ ,

$$\begin{aligned} & W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x \\ & \geq \int_0^x f_1(s, t, \bar{U}(s, t), \underline{V}(s, t)) \bar{U}_x(s, t) ds - \int_0^x f_1(s, t, \underline{U}(s, t), \bar{V}(s, t)) \underline{U}_x(s, t) ds \\ & \geq \int_0^x f_1(s, t, \bar{U}(s, t), \bar{V}(s, t)) \bar{U}_x(s, t) ds - \int_0^x f_1(s, t, \underline{U}(s, t), \bar{V}(s, t)) \underline{U}_x(s, t) ds \end{aligned} \quad (3.8)$$



where we used  $\bar{V}(x, t) \geq \underline{V}(x, t)$  for  $(x, t) \in [0, L] \times [0, t^*]$ . Integrating by parts as in (3.4), we have

$$\begin{aligned}
 & W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x \\
 & \geq F_1\left(x, t, \bar{U}(x, t), \bar{V}(x, t)\right) - F_1\left(x, t, \underline{U}(x, t), \bar{V}(x, t)\right) \\
 & + \int_0^x \left[ \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] \bar{V}_x(s, t) ds \\
 & + \int_0^x \left[ \frac{\partial F_1}{\partial x}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial x}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] ds \\
 & \geq h(x, t)W + \int_0^x \left[ \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] \bar{V}_x(s, t) ds,
 \end{aligned}
 \tag{3.9}$$

for  $x \in [0, L]$ ,  $t \in (0, t^*]$ , where

$$h(x, t) = \int_0^1 f_1\left(x, t, \xi \bar{U}(s, t) + (1 - \xi)\underline{U}(s, t), \bar{V}(s, t)\right) d\xi \in L_{loc}^\infty([0, L] \times \mathbb{R}_+).$$

Note that we have used  $\frac{\partial}{\partial \bar{U}}\left(\frac{\partial F_1}{\partial x}\right) = \frac{\partial f_1}{\partial x} \leq 0$ , i.e.  $\frac{\partial F_1}{\partial x}$  is non-increasing in  $U$  in the last inequality of (3.9). Summarizing, we have

$$\begin{aligned}
 & W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x - h(x, t)W \\
 & \geq \int_0^x \left[ \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \underline{U}(s, t), \bar{V}(s, t)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \right] \bar{V}_x(s, t) ds.
 \end{aligned}
 \tag{3.10}$$

Since  $\frac{\partial}{\partial \bar{U}}\left(\frac{\partial F_1}{\partial x}\right) = \frac{\partial f_1}{\partial x} < 0$ , i.e.  $\frac{\partial F_1}{\partial x}$  is non-increasing in  $U$ ,  $\bar{U} \geq \underline{U}$ , and  $\bar{V}_x > 0$ , the last integral is non-negative. Thus  $W = \bar{U} - \underline{U}$  satisfies the following linear differential inequality:

$$(3.11) \quad W_t - D_1(x, t)W_{xx} + \alpha_1(x, t)W_x - h(x, t)W \geq 0, \quad \text{for } x \in (0, L], \quad t \in (0, t^*].$$

We claim that  $W \equiv 0$  in  $[0, L] \times [0, t^*]$ . If not, then the parabolic strong maximum principle applied to (3.11) implies that  $W(x, t^*) > 0$  for  $x \in (0, L)$ . Therefore, if there exists some  $x^* \in (0, L]$  such that  $W(x^*, t^*) = 0$ , then  $x^* = L$ , i.e.,  $W(L, t^*) = 0$ , and hence  $W_t(L, t^*) \leq 0$ . By the boundary conditions at  $(x, t) = (L, t^*)$ ,

$$D_1 \bar{U}_{xx} - \alpha_1 \bar{U}_x \geq 0 \geq D_1 \underline{U}_{xx} - \alpha_1 \underline{U}_x,$$

we have  $D_1(L, t^*)W_{xx}(L, t^*) - \alpha_1(L, t^*)W_x(L, t^*) \geq 0$ . Then by (3.10) we have

$$\begin{aligned}
 & 0 \geq W_t(L, t^*) \\
 & \geq \int_0^L \left[ \frac{\partial F_1}{\partial \bar{V}}\left(s, t^*, \underline{U}(s, t^*), \bar{V}(s, t^*)\right) - \frac{\partial F_1}{\partial \bar{V}}\left(s, t^*, \bar{U}(s, t^*), \bar{V}(s, t^*)\right) \right] \bar{V}_x(s, t^*) ds.
 \end{aligned}$$

Since  $\underline{U}(x, t^*) \leq \bar{U}(x, t^*)$  in  $[0, L]$ , and  $\bar{V}_x > 0$ , we deduce that the above inequality holds only if  $\underline{U}(x, t^*) \equiv \bar{U}(x, t^*)$  for all  $x \in [0, L]$ , i.e.,  $W(x, t^*) \equiv 0$  for all  $x \in [0, L]$ . This is a contradiction and thus  $W = \bar{U} - \underline{U} \equiv 0$  in  $[0, L] \times [0, t^*]$ . It follows that equality holds everywhere in (3.8) and (3.9), in particular,

$$\int_0^x f_1\left(s, t, \bar{U}(s, t), \underline{V}(s, t)\right) \bar{U}_x(s, t) ds \equiv \int_0^x f_1\left(s, t, \bar{U}(s, t), \bar{V}(s, t)\right) \bar{U}_x(s, t) ds,$$

for all  $x \in [0, L]$  and  $0 < t \leq t^*$ . Since  $\bar{U}_x > 0$  and  $\frac{\partial f_1}{\partial V} < 0$ , we deduce that  $\bar{V}(x, t) \equiv \underline{V}(x, t)$  in  $[0, L] \times (0, t^*]$  and, by continuity, in  $[0, L] \times [0, t^*]$ .

The remaining case  $\bar{V}(x^*, t^*) = \underline{V}(x^*, t^*)$  for some  $x^* \in (0, L]$  can be handled similarly. This completes the proof in case (a) holds.

Next, assume (b) holds. We claim that necessarily there is a sequence of  $t_j \nearrow t^*$  such that alternative (a) holds, so that we can deduce similarly that  $(\bar{U}, \underline{V}) \equiv (\underline{U}, \bar{V})$  in  $[0, L] \times [0, t_j]$  for all  $j$ , whence (3.7) holds as well upon letting  $t_j \nearrow t^*$ . To see the claim, assume for contradiction that  $\bar{U} > \underline{U}$  and  $\bar{V} > \underline{V}$  for  $(x, t) \in (0, L] \times [t^* - \delta', t^*]$  for some  $\delta'$ . Then, observe that the boundary condition ensures  $W(0, t^*) = \bar{U}(0, t^*) - \underline{U}(0, t^*) = 0$ . Since  $W$  satisfies the differential inequality (3.11), we may apply Hopf's Lemma [26, Lemma 2.8] to deduce that  $(\bar{U} - \underline{U})_x(0, t^*) > 0$ . Similarly, we can deduce that  $(\bar{V} - \underline{V})_x(0, t^*) > 0$  as well, i.e. alternative (b) does not hold in this case. This establishes the claim and finishes the proof.  $\square$

Theorem 3.3 is a consequence of Lemma 3.6 and the following result:

LEMMA 3.7. Assume that  $f_1, f_2$  satisfy (H). Let  $(\bar{U}, \underline{V})$  and  $(\underline{U}, \bar{V})$  be a pair of super- and subsolutions of (3.3) in the time interval  $[0, T]$ . If

$$(3.12) \quad \underline{U}_x(x, t) > 0, \quad \text{and} \quad \underline{V}_x(x, t) > 0 \quad \text{for } (x, t) \in [0, L] \times [0, T],$$

then

$$(3.13) \quad \bar{U}(x, t) \geq \underline{U}(x, t) \quad \text{and} \quad \underline{V}(x, t) \leq \bar{V}(x, t) \quad \text{for } 0 \leq x \leq L, \quad 0 \leq t \leq T.$$

*Proof.* It is enough to prove the result for arbitrary but finite  $T > 0$ . Given a pair of super- and subsolutions  $(\bar{U}, \underline{V})$  and  $(\underline{U}, \bar{V})$  in a bounded interval  $[0, T]$ , we show (3.13) in two steps.

**Step 1.** For each small  $\delta > 0$ , define

$$(\bar{U}^\delta, \underline{V}^\delta) = (\bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2), \quad \text{and} \quad (\underline{U}^\delta, \bar{V}^\delta) = (\underline{U} - \delta\rho_1, \bar{V} + \delta\rho_2),$$

where  $\rho_i(x, t) := \int_0^x \exp\left(Mt + \int_0^y \frac{\alpha_i(s, t)}{D_i(s, t)} ds\right) dy$  for  $i = 1, 2$ . By (3.12), there exists  $\delta_0 > 0$  such that for each  $\delta \in (0, \delta_0]$ ,

$$(3.14) \quad \begin{cases} (\bar{U}^\delta, \underline{V}^\delta), (\underline{U}^\delta, \bar{V}^\delta) \in E & \text{for } t \in [0, T], \\ \bar{U}_x^\delta > 0, \underline{V}_x^\delta > 0, \underline{U}_x^\delta > 0, \bar{V}_x^\delta > 0 & \text{for } (x, t) \in [0, L] \times [0, T], \\ (\bar{U}^\delta(\cdot, 0), \underline{V}^\delta(\cdot, 0)) \gg_P (\underline{U}^\delta(\cdot, 0), \bar{V}^\delta(\cdot, 0)). \end{cases}$$

It is also clear that there is  $C_0 > 0$  (independent of  $\delta$ ) such that

$$(3.15) \quad \max_{i=1,2} \|\rho_i\|_{C([0,L] \times [0,T])} \leq C_0 \min_{i=1,2} \inf_{[0,L] \times [0,T]} (\rho_i)_x(x, t).$$

We claim that  $(\bar{U}^\delta, \underline{V}^\delta)$  and  $(\underline{U}^\delta, \bar{V}^\delta)$  forms a pair of super- and subsolutions for (3.3) in the interval  $[0, T]$ , in the sense of Definition 3.5. It remains to show the differential inequalities (3.2) for  $\delta$  small, as the initial and boundary conditions are clearly satisfied. A sufficient condition for the first one to hold is

$$(3.16) \quad \delta(\rho_1)_{x,t} \geq \kappa_1 [f_1(x, t, \bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2) - f_1(x, t, \bar{U}, \underline{V})] \bar{U}_x + \delta\rho_{1,x} f_1(x, t, \bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2).$$

The inequality (3.16) holds since the following holds pointwisely in  $[0, L] \times [0, T]$ :

$$\begin{aligned} & \delta(\rho_1)_{x,t} - [f_1(x, t, \bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2) - f_1(x, t, \bar{U}, \underline{V})] \bar{U}_x - \delta\rho_{1,x} f_1(x, t, \bar{U} + \delta\rho_1, \underline{V} - \delta\rho_2) \\ & \geq \delta \left( \rho_{1,x} \left[ M + \int_0^x \left( \frac{\alpha_1(s, t)}{D_1(s, t)} \right) ds - \|f_1\|_\infty \right] - \|Df_1\|_\infty (\rho_1 + \rho_2) \|\bar{U}_x\|_\infty \right), \end{aligned}$$

(note that  $\bar{U}_x, \bar{V}_x \in C([0, L] \times [0, T])$  by definition of super- and subsolutions) and, by (3.15), the term in the square bracket is non-negative provided the positive parameter  $M = M(C_0, \|f\|_{C^1})$  is chosen large enough (but uniformly for  $\delta \in (0, \delta_0]$ ). In the same way, one can show the rest of the differential inequalities. In summary, there is  $M > 0$  so that for all  $\delta \in (0, \delta_0]$ ,  $(\bar{U}^\delta, \underline{V}^\delta)$  and  $(\underline{U}^\delta, \bar{V}^\delta)$  form a pair of super- and subsolutions for (3.3) in the interval  $[0, T]$ . This proves our first claim.

**Step 2.** Next, we claim that for all  $\delta > 0$ ,

$$(3.17) \quad \bar{U}^\delta(x, t) > \underline{U}^\delta(x, t) \quad \text{and} \quad \underline{V}^\delta(x, t) < \bar{V}^\delta(x, t) \quad \text{for } (x, t) \in (0, L] \times [0, T].$$

Suppose not, then it follows from (3.14) that there exists a positive maximal time denoted by  $t^* \in (0, T]$  such that  $\underline{U}^\delta(x, t) < \bar{U}^\delta(x, t)$ ,  $\bar{V}^\delta(x, t) > \underline{V}^\delta(x, t)$  hold for  $0 < x \leq L$  and  $0 \leq t < t^*$ , and  $\underline{U}^\delta(x^*, t^*) = \bar{U}^\delta(x^*, t^*)$  or  $\bar{V}^\delta(x^*, t^*) = \underline{V}^\delta(x^*, t^*)$  for some  $x^* \in (0, L]$ . It follows from Lemma 3.6 that  $\underline{U}^\delta(x, t) \equiv \bar{U}^\delta(x, t)$  and  $\bar{V}^\delta(x, t) \equiv \underline{V}^\delta(x, t)$  for all  $0 \leq x \leq L$  and  $0 \leq t \leq t^*$ , which is a contradiction to (3.14). This shows (3.17). Letting  $\delta \rightarrow 0$  in (3.17), we deduce that (3.13) holds for  $(x, t) \in [0, L] \times [0, T]$ .  $\square$

Now we prove Corollary 3.4, which includes Theorem 2.1 as a special case.

*Proof of Corollary 3.4.* For  $i = 1, 2$ , let

$$(3.18) \quad (U_i(x, t), V_i(x, t)) = \left( \int_0^x u_i(s, t) ds, \int_0^x v_i(s, t) ds \right).$$

If we assume in addition that

$$(3.19) \quad u_2(x, 0) = U_{2,x}(x, 0) > 0 \quad \text{and} \quad v_1(x, 0) = V_{1,x}(x, 0) > 0 \quad \text{in } [0, L],$$

then by applying the strong maximum principle to the first and second equations of (3.1) separately, we deduce that

$$u_2 = U_{2,x} > 0 \quad \text{and} \quad v_1 = V_{1,x} > 0 \quad \text{in } [0, L] \times [0, T].$$

Therefore, applying Lemma 3.7, we see that if  $(U_1(\cdot, 0), V_1(\cdot, 0)) \geq_P (U_2(\cdot, 0), V_2(\cdot, 0))$  and (3.19) holds, then

$$(3.20) \quad (U_1(\cdot, t), V_1(\cdot, t)) \geq_P (U_2(\cdot, t), V_2(\cdot, t)) \quad \text{for all } t > 0.$$

By the fact that initial data satisfying (3.19) is dense in  $E$ , we can show that for general initial data in  $E$ , if  $(U_1(\cdot, 0), V_1(\cdot, 0)) >_P (U_2(\cdot, 0), V_2(\cdot, 0))$ , then (3.20) holds.

It remains to show that if  $(U_1(\cdot, 0), V_1(\cdot, 0)) >_P (U_2(\cdot, 0), V_2(\cdot, 0))$  and that both  $U_{1,x}, V_{2,x}$  are non-negative and non-trivial, then

$$(U_1(\cdot, t), V_1(\cdot, t)) \gg_P (U_2(\cdot, t), V_2(\cdot, t)) \quad \text{for all } t > 0.$$

This follows from Lemma 3.6, provided it can be verified that

$$u_1(x, t) = (U_1)_x(x, t) > 0, \quad v_2(x, t) = (V_2)_x(x, t) > 0 \quad \text{for } 0 \leq x \leq L, \quad 0 < t \leq T.$$

But this is an immediate consequence of the strong maximum principle applied to the equations of  $u_1$  and  $v_2$  separately.  $\square$

**3.2. Global Dynamics of the Single Species Model.** In this section, we generalize some known results about the following single species model, which is obtained by setting  $v = 0$  in (3.1):

$$(3.21) \quad \begin{cases} \theta_t = (D_1\theta_x - \alpha_1\theta)_x + f_1(x, t, \int_0^x \theta(s, t) ds, 0)\theta, & 0 < x < L, \ t > 0, \\ D_1\theta_x - \alpha_1\theta = 0, & x = 0, L, \ t > 0, \\ \theta(x, 0) = \theta_0(x) \geq, \neq 0, & 0 \leq x \leq L, \end{cases}$$

where  $D_1 = D_1(x, t) > 0$ ,  $\alpha_1 = \alpha_1(x, t)$ , and  $f_1$  are smooth and **(H)** holds.

The equation (3.21) generates a continuous semiflow in  $C([0, L]; \mathbb{R}_+)$  (see, e.g. [10]). Furthermore, by regarding the nonlocal term  $f_1(x, t, \int_0^x \theta(s, t) ds, 0)$  as a given coefficient, we can view (3.21) as a linear non-autonomous parabolic equation. It follows from the classical maximum principle that  $\theta(x, t) > 0$  for  $x \in [0, L]$  and  $t > 0$ .

Define  $\bar{\theta} \in C([0, \infty); C([0, L]; \mathbb{R}_+)) \cap C^1((0, \infty); C^\infty([0, L]; \mathbb{R}_+))$  to be a supersolution of (3.21) if

$$(3.22) \quad \begin{cases} \bar{\theta}_t \geq_{\mathcal{K}_1} (D_1\bar{\theta}_x - \alpha_1\bar{\theta})_x + f_1(x, t, \int_0^x \bar{\theta}(s, t) ds, 0)\bar{\theta}, & t > 0, \\ D_1\bar{\theta}_x - \alpha_1\bar{\theta} = 0, & x = 0, L, \ t > 0. \end{cases}$$

And define  $\underline{\theta}$  to be a subsolution of (3.21) if it satisfies the reverse inequality. As a by-product of the proofs of Lemmas 3.6 and 3.7, we can similarly show that the single species model is strongly monotone with respect to the order generated by cone  $\mathcal{K}_1$ .

**COROLLARY 3.8.** *Assume that  $f_1$  satisfies **(H)**. Let  $\bar{\theta}$  and  $\underline{\theta}$  be super- and subsolution of (3.21) such that*

$$\bar{\theta}(x, t) > 0, \quad \underline{\theta}(x, t) > 0, \quad \text{in } [0, L] \times [0, T], \quad \text{and} \quad \bar{\theta}(\cdot, 0) \geq_{\mathcal{K}_1} \underline{\theta}(\cdot, 0).$$

*Then  $\bar{\theta}(\cdot, t) \geq_{\mathcal{K}_1} \underline{\theta}(\cdot, t)$  for all  $t > 0$ . Furthermore, if for some  $t_0 > 0$  we have  $\bar{\theta}(\cdot, t_0) - \underline{\theta}(\cdot, t_0) \notin \text{Int } \mathcal{K}_1$ , then  $\bar{\theta}(\cdot, t) \equiv \underline{\theta}(\cdot, t)$  for  $t \in [0, t_0]$ .*

In particular, the continuous semiflow generated by (3.21) is strongly monotone with respect to the order induced by the cone  $\mathcal{K}_1$ .

In contrast to Corollary 3.8, we show here that the pointwise competitive order is not preserved by (3.21).

**PROPOSITION 3.9.** *For  $i = 1, 2$ , let  $\theta_i$  be a solution of (3.21), with initial conditions  $\theta_{i,0} \in \{\psi \in C^2([0, L]) : D_1\psi_x = \alpha_1\psi \text{ for } x = 0, L\}$ . If*

$$\theta_{1,0} \leq, \neq \theta_{2,0} \quad \text{in } [0, L], \quad \text{and} \quad \theta_{1,0} \equiv \theta_{2,0} \quad \text{in } [L - \delta, L] \text{ for some } \delta > 0,$$

*then  $\theta_1(L, t) > \theta_2(L, t)$  for all  $0 < t \ll 1$ .*

*Proof.* Since the initial conditions are  $C^2$  and consistent with the boundary condition, the solutions  $\theta_i$  are of class  $C_{x,t}^{2,1}$  in  $[0, L] \times [0, \infty)$ . Hence, it is enough to show that  $(\theta_1)_t(L, 0) > (\theta_2)_t(L, 0)$ . Precisely, at  $(x, t) = (L, 0)$ ,

$$\begin{aligned} (\theta_1)_t &= [D_1(\theta_1)_x - \alpha_1\theta_1]_x + f_1(L, 0, \int_0^L \theta_1(s, 0) ds, 0)\theta_1 \\ &> [D_1(\theta_1)_x - \alpha_1\theta_1]_x + f_1(L, 0, \int_0^L \theta_2(s, 0) ds, 0)\theta_1 \\ &= [D_1(\theta_2)_x - \alpha_1\theta_2]_x + f_1(L, 0, \int_0^L \theta_2(s, 0) ds, 0)\theta_2 = (\theta_2)_t. \quad \square \end{aligned}$$

To illustrate Proposition 3.9, we choose initial conditions  $\{\theta_{i,0}\}_{i=1,2}$  so that

$$\theta_{1,0} \leq_{P_1} \theta_{2,0} \quad \text{and} \quad \theta_{1,0} \leq_{\mathcal{K}_1} \theta_{2,0},$$

but only the order with respect to  $\mathcal{K}_1$  is preserved by the semiflow; see Figure 1.

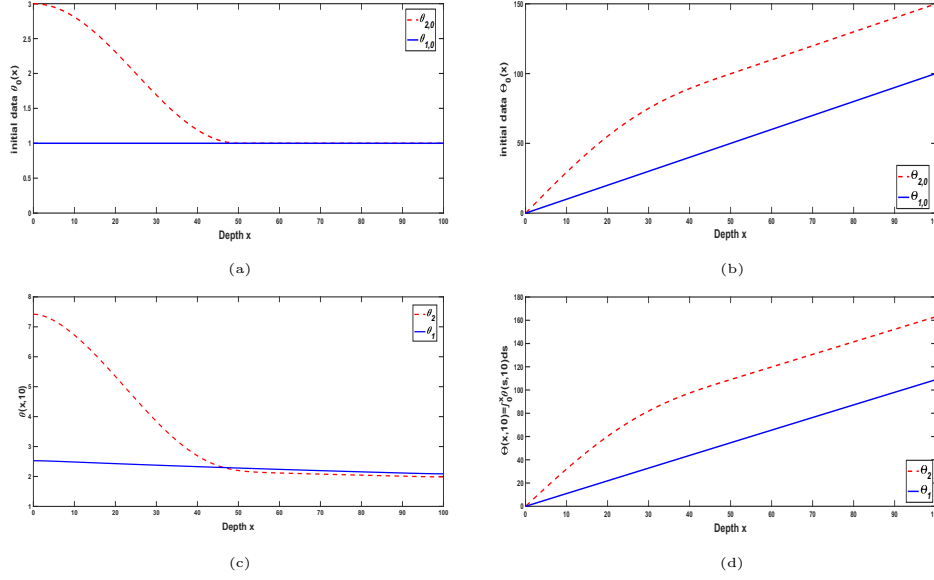


FIG. 1. Numerical solution of (3.21), with  $D_1 = 1$ ,  $\alpha_1 = 0$ ,  $L = 100$ ,  $f_1(x, t, \Theta, 0) = g(\exp(-k_0x - k_1\Theta))$ , where  $g(I) = \frac{I}{10+I}$  and  $k_0 = k_1 = d = 0.001$ , and initial condition  $\theta_{1,0} = \chi_{[0,L/2]}(\cos(2\pi x/L) + 1) + 1$  and  $\theta_{2,0} = 1$ . Panels (a) and (c) are the population densities  $\theta_i(x, t)$  ( $i = 1, 2$ ) at times  $t = 0$  and  $t = 10$  resp.; Panels (b) and (d) are the initial cumulative distribution functions of population densities  $\Theta_i(x, t) = \int_0^x \theta_i(s, t) ds$  ( $i = 1, 2$ ) at times  $t = 0$  and  $t = 10$ . The first (resp. second) species is represented by the red/dotted line (resp. blue/solid line).

*Remark 3.10.* By choosing  $u_i(\cdot, 0) = \theta_{i,0}$  for  $i = 1, 2$ , and  $v_1(\cdot, 0) \equiv v_2(\cdot, 0) \equiv \epsilon$ , then  $(u_1(\cdot, 0), v_1(\cdot, 0)) \leq_P (u_2(\cdot, 0), v_2(\cdot, 0))$ . However, it follows from the above result and continuous dependence on initial data that  $(u_1(\cdot, t), v_1(\cdot, t)) \not\leq_P (u_2(\cdot, t), v_2(\cdot, t))$  for some  $t > 0$ .

As a consequence of monotone dynamical systems theory, one can show the uniqueness and global asymptotic stability of positive equilibria (in the case of autonomous semiflow) or positive periodic solution (in the case of time-periodic semiflow). We will show the latter here, as the former follows as an easy consequence.

The following eigenvalue problem will be useful for our later purposes:

$$(3.23) \quad \begin{cases} \varphi_t = (D_1\varphi_x - \alpha_1\varphi)_x + f_1(x, t, 0, 0)\varphi + \mu\varphi, & 0 < x < L, \ 0 < t < T, \\ D_1\varphi_x - \alpha_1\varphi = 0, & x = 0, L, \ 0 < t < T, \\ \varphi(x, 0) = \varphi(x, T), & 0 \leq x \leq L, \\ \varphi(x, t) > 0, & 0 \leq x \leq L, \ 0 \leq t \leq T. \end{cases}$$

It is well known (see, e.g., [11]) that (3.23) has a principal eigenvalue, denoted by  $\mu_1$ , with the corresponding positive eigenfunction.

**PROPOSITION 3.11.** Assume that  $f_1$  satisfies **(H)**, and let  $D_1, \alpha_1, f_1$  be  $T$ -periodic in  $t$ , and there exists  $M_1 > 0$  such that

$$(3.24) \quad \sup_{[0,L] \times [0,T]} f_1(x, t, M_1, 0) < 0 \quad \text{and} \quad \|f_1(\cdot, \cdot, \cdot, 0)\|_{L^\infty([0,L] \times [0,T] \times [0,\infty))} \leq M_1.$$

Let  $\mu_1$  be the principal eigenvalue of (3.23).

(a) If  $\mu_1 \geq 0$ , then every solution of (3.21) converges to zero;

(b) If  $\mu_1 < 0$ , then (3.21) has a unique positive  $T$ -periodic solution. Furthermore, it attracts all non-negative, non-trivial solutions of (3.21).

In case  $f_1(x, t, p, 0) = g(I_0 \exp(-k_0 x - k_1 p)) - d(x, t)$  where  $g(\cdot)$  satisfies (1.5), the condition (3.24) is clearly satisfied, and the above result generalizes all previous results [5, 7, 28, 31, 32]. Our main contribution is a short proof of the boundedness of trajectories, which has not been proven when all coefficients vary periodically with time. This allows the use of the concept of subhomogeneity to show the existence, uniqueness and global stability of positive steady state simultaneously.

*Proof of Proposition 3.11.* We will apply [37, Theorem 2.3.4] to prove this proposition. Let  $\tilde{Q}_T$  be the Poincaré map of time  $T$ , generated by the  $T$ -periodic equation (3.21). It is obvious that the Poincaré map  $\tilde{Q}_T$  is monotone by Corollary 3.8, and compact in  $C([0, L])$  by parabolic estimate. Therefore, we need only to verify that every positive orbit of  $\tilde{Q}_T$  in  $C([0, L]; \mathbb{R}_+)$  is bounded,  $\tilde{Q}_T$  is strongly subhomogeneous, and the Fréchet derivative  $\mathcal{D}\tilde{Q}_T(0)$  is compact and strongly positive.

**Claim 1.** The semiflow is point dissipative, i.e. there exists  $M > 0$ , independent of initial data, such that

$$\limsup_{t \rightarrow \infty} \|\theta(\cdot, t)\|_{C([0, L])} \leq M.$$

By the fact that  $f_1(x, t, p, 0)$  is uniformly bounded in  $L^\infty$ , Harnack inequality [19, Theorem 2.5] applies, so that there is a uniform positive constant  $C' > 0$  such that

$$\sup_{0 < x < L} \theta(x, t) \leq C' \inf_{0 < x < L} \theta(x, t) \quad \text{for all } t \geq 1.$$

By (3.24), it is possible to choose a small constant  $\delta_2 > 0$  such that

$$(3.25) \quad C' \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx < 0 \quad \text{for } 0 \leq t \leq T.$$

It suffices to show that  $\limsup_{t \rightarrow \infty} \int_0^L \theta dx \leq \max\{M_1, C' L M_1 / \delta_2\}$ . To this end, it is enough to show the following claim.

**CLAIM 3.12.** *The differential inequality*

$$\frac{d}{dt} \int_0^L \theta(x, t) dx \leq -\delta_3 \int_0^L \theta(x, t) dx$$

holds whenever  $\int_0^L \theta(x, t) dx > \max\{M_1, C' L M_1 / \delta_2\}$ .

Now, denote  $\theta_*(t) = \inf_x \theta(x, t)$  and  $\theta^*(t) = \sup_x \theta(x, t)$ , then

$$M_1 < \frac{\delta_2}{C' L} \int_0^L \theta(x, t) dx \leq \frac{\delta_2}{C'} \theta^*(t) \leq \delta_2 \theta_*(t).$$

Integrating the equation of  $\theta$  over  $(0, L)$ , we obtain

$$\begin{aligned}
& \frac{d}{dt} \int_0^L \theta(x, t) dx \\
&= \int_0^L f_1 \left( x, t, \int_0^x \theta(s, t) ds, 0 \right) \theta(x, t) dx \\
&\leq \int_0^L f_1(x, t, x\theta_*(t), 0) \theta(x, t) dx \\
&\leq \int_0^{\delta_2} f_1(x, t, 0, 0) \theta(x, t) dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) \theta(x, t) dx \\
&\leq \theta^*(t) \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx \theta_*(t) \\
&\leq \left( C' \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx \right) \theta_*(t) \\
&\leq \left( C' \int_0^{\delta_2} \max\{f_1(x, t, 0, 0), 0\} dx + \int_{\delta_2}^L f_1(x, t, M_1, 0) dx \right) \frac{1}{C'L} \int_0^L \theta(x, t) dx.
\end{aligned}$$

This proves the point dissipativity.

**Claim 2.** The Poincaré map is strongly subhomogeneous.

We will show that  $\tilde{Q}_T$  is strongly subhomogeneous, i.e.

$$(3.26) \quad \tilde{Q}_T(\lambda\theta_0) \gg_{\kappa_1} \lambda\tilde{Q}_T(\theta_0) \quad \text{for all } \theta_0 >_{P_1} 0 \text{ and } \lambda \in (0, 1).$$

Let  $\theta(x, t)$  be solution to (3.21) with initial condition  $\theta_0$ . For  $(x, t) \in (0, L) \times [0, T]$ ,

$$\begin{aligned}
(\lambda\theta)_t &= (D_1(\lambda\theta)_x - \alpha_1(\lambda\theta))_x + f_1(x, t, \int_0^x \theta(s, t) ds, 0)(\lambda\theta) \\
&< (D_1(\lambda\theta)_x - \alpha_1(\lambda\theta))_x + f_1(x, t, \int_0^x \lambda\theta(s, t) ds, 0)(\lambda\theta).
\end{aligned}$$

i.e.  $\lambda\theta$  is a subsolution to the (3.21) with initial condition  $\lambda\theta_0$ . Since the above inequality is strict,  $\lambda\theta$  is not identically equal to the solution of (3.21) with initial condition  $\lambda\theta_0$ . By Corollary 3.8 and evaluate at time  $t = T$ , we deduce (3.26).

**Claim 3.** The Fréchet derivative  $\mathcal{D}\tilde{Q}_T(0)$  is compact and strongly positive.

This follows directly from the fact that  $\mathcal{D}\tilde{Q}_T(0) = Z(T)$ , where  $Z(t)$  is the analytic semigroup generated by the linearized system of (3.21) at  $\theta = 0$ :

$$(3.27) \quad \begin{cases} \theta_t = (D_1\theta_x - \alpha_1\theta)_x + f_1(x, t, 0, 0)\theta, & 0 < x < L, \ t > 0, \\ D_1\theta_x - \alpha_1\theta = 0, & x = 0, L, \ t > 0, \\ \theta(x, 0) = \theta_0 \geq, \neq 0, & 0 < x < L. \end{cases}$$

That  $Z(T)$  is strongly positive follows from standard parabolic maximum principle. Moreover, by standard parabolic  $L^p$  estimate,  $Z(T)$  is a bounded map from  $C([0, L])$  to  $C^2([0, L])$ . The map  $Z(T)$  is thus compact, by the Arzelà-Ascoli Theorem.

If  $\mu_1 \geq 0$ , then  $r(\mathcal{D}\tilde{Q}_T(0)) = \exp(-\mu_1 T) \leq 1$ . By [37, Theorem 2.3.4(a)], every solution of (3.21) converges to zero. If  $\mu_1 < 0$ , then  $r(\mathcal{D}\tilde{Q}_T(0)) = \exp(-\mu_1 T) > 1$ . By [37, Theorem 2.3.4(b)], the map  $\tilde{Q}_T$  has a unique positive fixed point  $\vartheta$  such that every positive orbit with non-negative, non-trivial, continuous initial data converges



to  $\tilde{\theta}$ . This means that system (3.21) has a unique positive  $T$ -periodic solution  $\tilde{\theta}$ , determined by  $\tilde{\theta}(\cdot, 0) = \tilde{\theta}(\cdot, T) = \tilde{\theta}$ , which attracts all non-negative and non-trivial solutions of (3.21).  $\square$

*Remark 3.13.* Within the context of a single species, we improved previous results in [28] by showing a strong maximum principle (which implies strong monotonicity of the semiflow) for super- and subsolutions (which satisfies only differential inequalities), and by allowing the coefficients to be space-time heterogeneous.

**4. Global Dynamics for the Nonlocal Two-species Model.** It is well known that diffusion and advection rates have significant effects on the outcome of competition. In this section, we apply Theorem 4.1 to analyze the global dynamics of two-species competition system. To obtain qualitative results, we restrict ourselves for the remainder of the paper to consider the autonomous case (1.1) - (1.3), when  $D_i, \alpha_i, d_i$  are constants. In the introduction, the light intensity  $I(x, t)$  is given by (1.4), where the shading coefficients of the two species are given by  $k_1, k_2$ . However, by transforming  $(\tilde{u}, \tilde{v}) = (k_1 u, k_2 v)$  and  $\tilde{g}_i(I_0 \cdot) = g_i(\cdot)$ , and by observing that  $k_1, k_2$  do not affect the dynamics qualitatively, we may assume  $k_1 = k_2 = 1$  and  $I_0 = 1$  without loss of generality, so that the light intensity (1.4) can be simplified to

$$(4.1) \quad I(x, t) = \exp \left( -k_0 x - \int_0^x [u(s, t) + v(s, t)] ds \right).$$

We focus on the following three different cases:

- (i)  $D_1 = D_2, \alpha_1 < \alpha_2$ ;
- (ii)  $D_1 < D_2, \alpha_1 = \alpha_2 \geq [g(1) - d]L > 0$ ;
- (iii)  $D_1 < D_2, \alpha_1 = \alpha_2 \leq 0$ .

Due to the strongly monotonicity proved in Theorem 2.1, to a large extent, its dynamics can be determined by the stability/instability of the semi-trivial solution of the stationary problem [2, 12, 15, 25, 34, 37]. For the convenience of the readers, we state the precise abstract theorem here.

**THEOREM 4.1** ([15, Theorem B] and [25, Theorem 1.3]). *If the system (1.1)-(1.4) has no positive steady states, and the semi-trivial steady state  $(0, \tilde{v})$  (resp.  $(\tilde{u}, 0)$ ) is linearly unstable, then  $(\tilde{u}, 0)$  (resp.  $(0, \tilde{v})$ ) is globally asymptotically stable among all non-negative, non-trivial solutions.*

*Remark 4.2.* Our setting is slightly more general than that outlined in [15]. In particular, the semiflow  $Q_t$  generated by (3.1) is defined in  $Y^+ = Y_1^+ \times Y_1^+$ , where  $Y_1^+ = C([0, L]; \mathbb{R}_+)$ , but the semiflow only preserve the order generated by the weaker cone  $\mathcal{K} = \mathcal{K}_1 \times (-\mathcal{K}_1)$ , with  $Y_1^+ \subsetneq \mathcal{K}_1$ . However, it is straightforward to observe that [15, Propositions 2.1 and 2.4] are independent of the above assumption, and that the proofs of [15, Theorem B] and [25, Theorem 1.3] both stand in our setting. Therefore, we omit the proof of Theorem 4.1 here.

In preparation to apply Theorem 4.1, we will demonstrate that the equation

$$(4.2) \quad \begin{cases} \theta_t = D\theta_{xx} - \alpha\theta_x + [g(e^{-k_0 x - \int_0^x \theta(s, t) ds}) - d]\theta = 0, & 0 < x < L, \\ D\theta_x - \alpha\theta = 0, & x = 0, L, \\ \theta(x, 0) = \theta_0(x) \geq 0, & 0 \leq x \leq L, \end{cases}$$

has a unique positive steady state  $\tilde{\theta}$ , which is always linearly stable, and then characterize the stability of the two semi-trivial steady states in terms of two standard principal eigenvalue problems.

**4.1. An Eigenvalue Problem for the Single Species Model.** For constants

$D > 0$ ,  $\alpha \in \mathbb{R}$  and  $h \in C([0, L])$ , consider the following standard eigenvalue problem:

$$(4.3) \quad \begin{cases} D\phi_{xx} - \alpha\phi_x + h(x)\phi + \lambda\phi = 0, & 0 < x < L, \\ D\phi_x - \alpha\phi = 0, & x = 0, L. \end{cases}$$

By setting  $\psi = e^{-(\alpha/D)x}\phi$ , the problem (4.3) can be transformed into a self-adjoint problem

$$(4.4) \quad \begin{cases} -D(e^{(\alpha/D)x}\psi_x)_x - h(x)e^{(\alpha/D)x}\psi = \lambda e^{(\alpha/D)x}\psi, & 0 < x < L, \\ \psi_x(0) = \psi_x(L) = 0. \end{cases}$$

Therefore, all eigenvalues of (4.4) (and thus (4.3)) are real, and we can denote the smallest eigenvalue by  $\lambda_1(D, \alpha, h)$ . Define

$$d_* = -\lambda_1(D, \alpha, -g(e^{-k_0x})).$$

It is easy to show that  $d_*$  is positive. In fact,  $d_*$  is the critical death rate.

**THEOREM 4.3** ([5, Theorem 2.1], [13, Theorem 3.1]). *If  $0 < d < d_*$ , then (4.2) has a unique positive steady state, denoted by  $\tilde{\theta}(x)$ . If  $d \geq d_*$ , then zero is the only nonnegative steady state of (4.2).*

We linearize (4.2) at  $\tilde{\theta}$  to obtain the following eigenvalue problem:

$$(4.5) \quad \begin{cases} D\phi_{xx} - \alpha\phi_x + [g(\sigma) - d]\phi - \tilde{\theta}\sigma g'(\sigma) \int_0^x \phi(s) ds + \mu\phi = 0, & 0 < x < L, \\ D\phi_x - \alpha\phi = 0, & x = 0, L, \end{cases}$$

where  $\sigma = e^{-k_0x - \int_0^x \tilde{\theta}(s) ds}$ .

Our result says that  $\tilde{\theta}$  is linearly stable. In fact, there is a real eigenvalue of (4.5) which is strictly less than the real part of all other eigenvalues of (4.5).

**THEOREM 4.4.** *Let  $\tilde{\theta}$  be the unique positive steady state of (4.2). The eigenvalue problem (4.5) admits a real, simple eigenvalue  $\mu_1$  and an eigenfunction  $\phi \gg_{\kappa_1} 0$ , such that  $\mu_1 < \text{Re } \mu$  for all eigenvalues  $\mu \neq \mu_1$ . It is characterized as the unique eigenvalue of (4.5) with the eigenfunction  $\phi >_{\kappa_1} 0$ . Furthermore,  $\mu_1 > 0$ .*

*Proof.* Assume  $\tilde{\theta}(x)$  is a positive steady state of (4.2), and let  $\theta_0 \in C([0, L]; \mathbb{R})$ . Then  $\theta(\cdot, t) = \hat{\Phi}_t(\theta_0)$ , where  $\hat{\Phi}_t$  denotes the continuous semiflow generated by (4.2). Then  $z(x, t) = \mathcal{D}\hat{\Phi}_t(\theta)[\theta_0](x)$  satisfies the linear equation

$$(4.6) \quad z_t + \mathcal{L}z = 0, \quad z(0) = \theta_0.$$

where the unbounded operator

$$\mathcal{L} = -D\partial_{xx} + \alpha\partial_x - [g(\sigma) - d] + \tilde{\theta}\sigma g'(\sigma) \left( \int_0^x \cdot \right)$$

is defined on the domain

$$\text{Dom}(\mathcal{L}) = \{z \in C^2((0, L)) \cap C^1([0, L]) : \mathcal{L}z \in C([0, L]), Dz_x - \alpha z|_{x=0, L} = 0\}.$$

According to [30, Proposition 3.1.4], the linear equation (4.6) generates an analytic semigroup  $e^{-\mathcal{L}t}$  on  $C([0, L])$ . Thus  $\mathcal{D}\hat{\Phi}_t(\theta) = e^{-\mathcal{L}t}$ .

For  $\theta_0 \in \mathcal{K}_1$ ,  $\epsilon > 0$ , the monotonicity of  $\hat{\Phi}_t$  with respect to cone  $\mathcal{K}_1$  implies

$$\frac{\theta(\cdot, t; \tilde{\theta} + \epsilon\theta_0) - \theta(\cdot, t; \tilde{\theta})}{\epsilon} = \frac{\hat{\Phi}_t(\tilde{\theta} + \epsilon\theta_0) - \hat{\Phi}_t(\tilde{\theta})}{\epsilon} \geq_{\mathcal{K}_1} 0.$$

Upon taking the limit as  $\epsilon \rightarrow 0$ , we get  $\mathcal{D}\hat{\Phi}_t(\tilde{\theta})[\theta_0] \geq_{\mathcal{K}_1} 0$ . In other words,  $e^{-\mathcal{L}t} = \mathcal{D}\hat{\Phi}_t(\tilde{\theta})$  is a positive operator with respect to the order generated by cone  $\mathcal{K}_1$  in the sense that  $\mathcal{D}\hat{\Phi}_t(\tilde{\theta})\mathcal{K}_1 \subset \mathcal{K}_1$  holds for  $t \geq 0$ .

Next, we show that the analytic semigroup  $e^{-\mathcal{L}t} = \mathcal{D}\hat{\Phi}_t(\tilde{\theta})$  is strongly positive with respect to the order generated by  $\mathcal{K}_1$ . To prove this, we only need to show that  $\int_0^x z(s, t) ds > 0$  and  $z(0, t) > 0$  for all  $t > 0$ . Since  $e^{-\mathcal{L}t} = \mathcal{D}\hat{\Phi}_t(\tilde{\theta})$  is a positive operator with respect to cone  $\mathcal{K}_1$ , then  $\int_0^x z(s, t) ds \geq 0$ . Therefore, if  $\int_0^x z(s, t) ds > 0$  does not hold, then there exists some  $(x_0, t_0) \in (0, L] \times (0, \infty)$  such that  $\int_0^{x_0} z(s, t_0) ds = 0$ .

Let  $\int_0^x z(s, t) ds = Z(x, t)$ . Using the relation

$$[g(\sigma) - d]z - \tilde{\theta}\sigma g'(\sigma)Z = [(g(\sigma) - d)Z]' + k_0 g'(\sigma)\sigma Z,$$

we may integrate (4.6) over  $(0, x)$  to obtain the differential inequality

$$(4.7) \quad Z_t - DZ_{xx} + \alpha Z_x - [g(\sigma) - d]Z = k_0 \int_0^x g'(\sigma)\sigma Z ds \geq 0.$$

Since  $\theta_0 \neq 0$  and  $Z(\cdot, 0) \neq 0$ , then the strong maximum principle implies  $Z(x, t) > 0$  for all  $x \in (0, L)$  and  $t > 0$ , i.e.,  $x_0 = L$  and  $Z(L, t_0) = 0$ . Then  $Z_t(L, t_0) \leq 0$ , and by the boundary condition, we deduce  $DZ_{xx}(L, t_0) - \alpha Z_x(L, t_0) = DZ_x(L, t_0) - \alpha z(L, t_0) = 0$ . It follows from (4.7) that

$$(4.8) \quad 0 \geq Z_t(L, t_0) = k_0 \int_0^L g'(\sigma)\sigma Z ds.$$

Since  $k_0 > 0$ ,  $\sigma > 0$ ,  $g'(\sigma) > 0$ , then  $Z(x, t_0) \equiv 0$  for all  $x \in [0, L]$ . Contradiction.

Hence,  $Z(x, t) = \int_0^x z(s, t) ds > 0$  for all  $t > 0$  and  $x \in (0, L]$ . Since  $Z(0, t) \equiv 0$  and  $Z(x, t)$  satisfies (4.7) for all  $t > 0$ , then  $z(0, t) = Z_x(0, t) > 0$  for all  $t > 0$  by the Hopf boundary lemma.

Therefore, for each  $t > 0$ , the operator  $e^{-\mathcal{L}t}$  is compact and strongly positive on  $C([0, L])$  with respect to the order generated by  $\mathcal{K}_1$ . It follows by standard arguments in [34, Ch. 7] that the elliptic eigenvalue problem (4.5) has a principal eigenvalue  $\mu_1 \in \mathbb{R}$  with all the stated properties, except for  $\mu_1 > 0$ .

To show  $\mu_1 > 0$ , we suppose to the contrary that  $\mu_1 \leq 0$  and use  $\phi_1 \gg_{\mathcal{K}_1} 0$  to get

$$\tilde{\theta}\sigma g'(\sigma) \int_0^x \phi_1(s) ds > 0 \quad \text{for } x \in (0, L].$$

Then (4.5) yields that

$$D\phi_{1,xx} - \alpha\phi_{1,x} + [g(\sigma) - d]\phi_1 + \mu_1\phi_1 > 0 \quad \text{for } 0 < x < L.$$

Next, we use the facts  $\int_0^x \phi_1(s) ds > 0$  and  $\tilde{\theta} > 0$  for  $x \in [0, L]$ , to obtain the constant  $c > 0$  such that  $\min_{[0, L]}(c\tilde{\theta} - \phi_1) = 0$ . Then  $\varphi = c\tilde{\theta} - \phi_1$  satisfies

$$\begin{cases} D\varphi_{xx} - \alpha\varphi_x + [g(\sigma) - d]\varphi + \mu_1\varphi < \mu_1 c\tilde{\theta} \leq 0 & \text{for } 0 < x < L, \\ D\varphi_x = \alpha\varphi & \text{for } x = 0, L, \\ \min_{[0, L]} \varphi = 0. \end{cases}$$

By the strict differential inequality and non-negativity of  $\varphi$  we must have  $\varphi > 0$  in  $(0, L)$  and that  $\varphi(x_0) = 0$  for some  $x_0 \in \{0, L\}$ . But the Hopf boundary lemma says  $\varphi_x(x_0) \neq 0$ , which contradicts the boundary condition  $\varphi_x(x_0) = \frac{\alpha}{D}\varphi(x_0) = 0$ .  $\square$

**4.2. Eigenvalue Problems for the Two-species Model.** In this subsection, we study the linear eigenvalue problem of the two-species model associated with the stability of semi-trivial steady states.

We assume the parameters are chosen so that system (1.1)-(1.4) has two semi-trivial steady states  $(\tilde{u}, 0)$  and  $(0, \tilde{v})$  (e.g. if the death rates  $d_i$  are not too large). The associated linearized eigenvalue problem at  $(\tilde{u}, 0)$  is

$$(4.9) \quad \begin{cases} D_1\phi_{xx} - \alpha_1\phi_x + [g_1(\sigma_1) - d_1]\phi - \tilde{u}\sigma_1 g'_1(\sigma_1) \left[ \int_0^x \phi(s) ds + \int_0^x \varphi(s) ds \right] + \Lambda\phi = 0, & 0 < x < L, \\ D_2\varphi_{xx} - \alpha_2\varphi_x + [g_2(\sigma_1) - d_2]\varphi + \Lambda\varphi = 0, & 0 < x < L, \\ D_1\phi_x - \alpha_1\phi = D_2\varphi_x - \alpha_2\varphi = 0, & x = 0, L, \end{cases}$$

where  $\sigma_1(x) = e^{-k_0 x - \int_0^x \tilde{u}(s) ds}$ .

We shall exploit the fact that the second equation is decoupled from the first. Consider the following eigenvalue problem:

$$(4.10) \quad \begin{cases} D_2\varphi_{xx} - \alpha_2\varphi_x + [g_2(\sigma_1) - d_2]\varphi + \lambda\varphi = 0, & 0 < x < L, \\ D_2\varphi_x - \alpha_2\varphi = 0, & x = 0, L. \end{cases}$$

As already discussed, (4.10) admits a real principal eigenvalue, denoted by  $\lambda_u = \lambda_1(D_2, \alpha_2, g_2(\sigma_1) - d_2)$ , which is simple, and its corresponding eigenfunction  $\varphi_1$  does not change sign, and  $\lambda_u < \lambda$  for all  $\lambda \neq \lambda_u$ . The stability properties of  $(\tilde{u}, 0)$  are determined by the sign of  $\lambda_u$ , as the next result shows.

**PROPOSITION 4.5.** *The problem (4.9) has a principal eigenvalue  $\Lambda_1 \in \mathbb{R}$ , in the sense that  $\Lambda_1 \leq \operatorname{Re} \Lambda$  for all eigenvalues  $\Lambda$  of (4.9) and that the corresponding eigenfunction can be chosen in  $\mathcal{K} \setminus \{(0, 0)\}$ . Furthermore, (denote  $Y_1^+ = C([0, L]; \mathbb{R}_+)$ )*

(a) *If the principal eigenvalue  $\lambda_u$  of (4.10) is positive, then  $\Lambda_1 > 0$ .*

(b) *If the principal eigenvalue  $\lambda_u$  of (4.10) is non-positive, then  $\Lambda_1 = \lambda_u \leq 0$  and the corresponding eigenfunction can be chosen in  $\operatorname{Int} \mathcal{K}_1 \times (-\operatorname{Int} Y_1^+)$ .*

*Proof.* By Theorem 2.1, the semiflow  $\{Q_t\}_{t \geq 0}$ , generated by the system (1.1)-(1.4) is strongly monotone with respect to the cone  $\mathcal{K}$ . As a result, the linear problem at any steady state generates a semigroup that is monotone with respect to the cone  $\mathcal{K}$ . Therefore, by standard arguments in [34, Ch. 7], we deduce that the elliptic problem (4.9), obtained by linearizing (1.1)-(1.4) at the steady state  $(\tilde{u}, 0)$ , has a principal eigenvalue  $\Lambda_1$  with the stated properties. In particular, we can choose the eigenfunction corresponding to  $\Lambda_1$  from within  $\mathcal{K} \setminus \{(0, 0)\}$ .

Now, consider the case when the principal eigenvalue  $\lambda_u$  of (4.10) is positive. Let  $\Lambda_1 \in \mathbb{R}$  be the principal eigenvalue of (4.9) with eigenfunction  $(\phi_1, \varphi_1) \in \mathcal{K} \setminus \{(0, 0)\}$ . We claim that  $\Lambda_1 > 0$ . There are two cases to consider: (i)  $\varphi_1 \neq 0$ ; (ii)  $\varphi_1 = 0$ .

In Case (i),  $(\Lambda_1, \varphi_1)$  furnishes an eigenpair of problem (4.10), the latter of which as smallest eigenvalue  $\lambda_u > 0$ . Thus,  $\Lambda_1 \geq \lambda_u > 0$ .

In Case (ii),  $(\Lambda_1, \phi_1)$  furnishes an eigenpair of

$$(4.11) \quad \begin{cases} D_1\phi_{xx} - \alpha_1\phi_x + [g_1(\sigma_1) - d_1]\phi - \tilde{u}\sigma_1 g'_1(\sigma_1) \int_0^x \phi(s) ds + \Lambda\phi = 0, & 0 < x < L, \\ D_1\phi_x - \alpha_1\phi = 0, & x = 0, L. \end{cases}$$

By Theorem 4.4, (4.11) has a positive principal eigenvalue  $\mu_1$ , and  $\mu_1$  is always positive. Hence, we must have  $\Lambda_1 \geq \mu_1 > 0$ . This finishes the proof in case  $\lambda_u > 0$ .

Next, let  $\lambda_u \leq 0$  and let  $\varphi_1 \in (-\text{Int } Y_1^+) \subset (-\text{Int } \mathcal{K}_1)$  be the corresponding principal eigenfunction of (4.10). It remains to construct  $\phi_1 \in \text{Int } \mathcal{K}_1$  such that  $\lambda_u$  is an eigenvalue of (4.9) with eigenfunction  $(\phi_1, \varphi_1) \in \text{Int } \mathcal{K}_1 \times (-\text{Int } Y_1^+)$ . To that end, define the operator  $\mathcal{L}_1 = -D_1 \partial_{xx} + \alpha_1 \partial_x - [g_1(\sigma_1) - d_1] + \tilde{u} \sigma_1 g'_1(\sigma_1) (\int_0^x \cdot)$ . By Theorem 4.4, the spectrum  $\sigma(\mathcal{L}_1) \subset \{z \in \mathbb{C} : \text{Re } z > 0\}$ . And hence for  $\lambda_u \leq 0$ , 0 is not an eigenvalue of  $\mathcal{L}_1 - \lambda_u \mathcal{I}$ , and the problem

$$\begin{cases} \mathcal{L}_1 \phi - \lambda_u \phi = -\tilde{u} \sigma_1 g'_1(\sigma_1) \int_0^x \varphi_1(s) ds, & 0 < x < L, \\ D_1 \phi_x - \alpha_1 \phi = 0, & x = 0, L, \end{cases}$$

has a unique solution  $\phi_1$ . In fact, let  $f = -\tilde{u} \sigma_1 g'_1(\sigma_1) \int_0^x \varphi_1(s) ds$ , then  $f > 0$  and

$$\phi_1 = (\mathcal{L}_1 - \lambda_u)^{-1} f = \int_0^\infty e^{\lambda_u t} S_t f dt,$$

where  $S_t = e^{-\mathcal{L}_1 t}$  is the analytic semigroup generated by  $\mathcal{L}_1$  (see, e.g. [9, Theorem 3, Sect. 7.4]). From the proof of Theorem 4.4,  $S_t$  is strongly positive with respect to the order generated by cone  $\mathcal{K}_1$ . Therefore,  $S_t f \gg_{\mathcal{K}_1} 0$  for all  $t > 0$ , and

$$\phi_1 \gg_{\mathcal{K}_1} \int_0^\infty e^{\lambda_u t} S_t f dt \gg_{\mathcal{K}_1} 0.$$

By construction, we conclude that  $\lambda_u \leq 0$  is an eigenvalue of (4.9) with eigenfunction  $(\phi_1, \varphi_1) \in \text{Int } \mathcal{K}_1 \times (-\text{Int } Y_1^+)$ . Hence  $\Lambda_1 \leq \lambda_u \leq 0$ . On the other hand, let  $(\tilde{\phi}, \tilde{\varphi})$  be the eigenfunction of  $\Lambda_1$ , then  $\tilde{\varphi} \not\equiv 0$ , since otherwise  $(\Lambda, \tilde{\phi})$  is an eigenpair of (4.11), whence  $\Lambda \geq \mu_1 > 0$ , contradictions. Therefore,  $\tilde{\varphi} \not\equiv 0$  and  $(\Lambda_1, \tilde{\varphi})$  furnishes an eigenpair of (4.10). Thus  $\Lambda_1 \geq \lambda_u$  as well. This completes the proof.  $\square$

The linearized eigenvalue problem at semi-trivial steady state  $(0, \tilde{v})$  is (4.12)

$$\begin{cases} D_1 \phi_{xx} - \alpha_1 \phi_x + [g_1(\sigma_2) - d_1] \phi + \tilde{\Lambda} \phi = 0, & 0 < x < L, \\ D_2 \varphi_{xx} - \alpha_2 \varphi_x + [g_2(\sigma_2) - d_2] \varphi + \tilde{\Lambda} \varphi = \tilde{v} \sigma_2 g'_2(\sigma_2) [\int_0^x \phi(s) ds + \int_0^x \varphi(s) ds], & 0 < x < L, \\ D_1 \phi_x - \alpha_1 \phi = 0, & x = 0, L, \\ D_2 \varphi_x - \alpha_2 \varphi = 0, & x = 0, L, \end{cases}$$

where  $\sigma_2(x) = e^{-k_0 x - \int_0^x \tilde{v}(s) ds}$ . Let  $\lambda_v = \lambda_1(D_1, \alpha_1, g_1(\sigma_2) - d_1)$  denote the principal eigenvalue of the eigenvalue problem

$$(4.13) \quad \begin{cases} D_1 \phi_{xx} - \alpha_1 \phi_x + [g_1(\sigma_2) - d_1] \phi + \lambda \phi = 0, & 0 < x < L, \\ D_1 \phi_x - \alpha_1 \phi = 0, & x = 0, L. \end{cases}$$

It follows analogously that the stability properties of  $(0, \tilde{v})$  are determined by  $\lambda_v$ .

**PROPOSITION 4.6.** *The problem (4.12) has a principal eigenvalue  $\tilde{\Lambda}_1 \in \mathbb{R}$ , in the sense that  $\tilde{\Lambda}_1 \leq \text{Re } \tilde{\Lambda}$  for all eigenvalues  $\tilde{\Lambda}$  of (4.12) and that the corresponding eigenfunction can be chosen in  $\mathcal{K} \setminus \{(0, 0)\}$ . Furthermore, (denote  $Y_1^+ = C([0, L]; \mathbb{R}_+)$ )*

- (a) *If the principal eigenvalue  $\lambda_v$  of (4.13) is positive, then  $\tilde{\Lambda}_1 > 0$ .*
- (b) *If the principal eigenvalue  $\lambda_v$  of (4.13) is non-positive, then  $\tilde{\Lambda}_1 = \lambda_v \leq 0$  and the corresponding eigenfunction can be chosen in  $\text{Int } Y_1^+ \times (-\text{Int } \mathcal{K}_1)$ .*

**4.3. Auxilliary Eigenvalue Lemmas.** In this subsection, we prove several useful lemmas concerning the principal eigenvalue  $\lambda_1(D, \alpha, h)$  of (4.3) with positive

eigenfunction  $\phi_1$ . It can be shown that  $\lambda_1$  and  $\phi_1$  are smooth functions of  $\alpha$  and  $D$  (see, e.g., [1, Lemma 1.2]).

We will assume additionally the following:

(A)  $h(x) \in C^1([0, L])$  such that  $h'(x) < 0$  in  $[0, L]$ .

Set  $\psi_1 = e^{-(\alpha/D)x} \phi_1$ . Then  $\psi_1$  satisfies

$$(4.14) \quad \begin{cases} D\psi_{1,xx} + \alpha\psi_{1,x} + h(x)\psi_1 + \lambda_1\psi_1 = 0, & 0 < x < L, \\ \psi_{1,x}(0) = \psi_{1,x}(L) = 0. \end{cases}$$

LEMMA 4.7. *If  $h(x)$  satisfies (A), then  $\psi_{1,x} < 0$  in  $(0, L)$ .*

*Proof.* Multiplying (4.14) by  $e^{(\alpha/D)x}$ , we rewrite the resulting equation as

$$(4.15) \quad \begin{cases} D(e^{(\alpha/D)x}\psi_{1,x})_x + e^{(\alpha/D)x}\psi_1[h(x) + \lambda_1] = 0, & 0 < x < L, \\ \psi_{1,x}(0) = \psi_{1,x}(L) = 0. \end{cases}$$

Integrating (4.15) over  $(0, L)$ , we have

$$\int_0^L e^{(\alpha/D)x} \psi_1 [h(x) + \lambda_1] dx = 0,$$

which implies that  $h(x) + \lambda_1$  changes sign in  $(0, L)$ . Since  $h(x)$  is strictly decreasing in  $(0, L)$ , then there exists a unique  $x_0 \in (0, L)$  such that  $h(x) + \lambda_1 > 0$  for  $0 < x < x_0$  and  $h(x) + \lambda_1 < 0$  for  $x_0 < x < L$ . Hence, by (4.15) we see that  $(e^{(\alpha/D)x}\psi_{1,x})_x < 0$  for  $0 < x < x_0$  and  $(e^{(\alpha/D)x}\psi_{1,x})_x > 0$  for  $x_0 < x < L$ . That is,  $e^{(\alpha/D)x}\psi_{1,x}$  is strictly decreasing in  $(0, x_0)$ , and strictly increasing in  $(x_0, L)$ . Since  $\psi_{1,x}(0) = \psi_{1,x}(L) = 0$ , we have  $\psi_{1,x} < 0$  in  $(0, L)$ .  $\square$

LEMMA 4.8. *If  $h(x)$  satisfies (A), then*

$$\frac{\partial \lambda_1}{\partial \alpha}(D, \alpha, h) > 0 \quad \text{for any } D > 0 \text{ and } \alpha \in \mathbb{R}.$$

The proof of Lemma 4.8 is similar to [13, Lemma 5.2], and we omit it here. The proof of the following Lemma 4.9 is similar to [13, Lemma 7.1] with some modifications. For the sake of completeness, we give the proof here in detail.

LEMMA 4.9. *If  $h(x)$  satisfies (A), then the following hold:*

(a)  $\frac{\partial \lambda_1}{\partial D}(D, \alpha, h) > 0$  for  $D > 0$  and  $\alpha \leq 0$ .

(b) If  $\alpha \geq h(0)L$  and  $\lambda_1(D^*, \alpha, h) = 0$  for some  $D^* > 0$ , then  $\frac{\partial \lambda_1}{\partial D}(D^*, \alpha, h) < 0$ .

*Proof.* Recall that  $\lambda_1$  and  $\psi_1$  are smooth functions of  $D$ . For simplicity of notation, we denote  $\frac{\partial \psi_1}{\partial D}$  by  $\psi'_1$ , etc., where  $\psi_1$  satisfies (4.14). Differentiating (4.14) with respect to  $D$ , we have

$$(4.16) \quad \begin{cases} D\psi'_{1,xx} + \psi_{1,xx} + \alpha\psi'_{1,x} + h(x)\psi'_1 + \lambda'_1\psi_1 + \lambda_1\psi'_1 = 0, & 0 < x < L, \\ \psi'_{1,x}(0) = \psi'_{1,x}(L) = 0. \end{cases}$$

Multiplying (4.16) by  $e^{(\alpha/D)x}\psi_1$  and integrating the resulting equation in  $(0, L)$ , we have

$$(4.17) \quad \begin{aligned} & -D \int_0^L e^{(\alpha/D)x} \psi'_{1,x} \psi_{1,x} dx + \int_0^L e^{(\alpha/D)x} \psi_{1,xx} \psi_1 dx + \int_0^L e^{(\alpha/D)x} h(x) \psi'_1 \psi_1 dx \\ & + \lambda'_1 \int_0^L e^{(\alpha/D)x} \psi_1^2 dx + \lambda_1 \int_0^L e^{(\alpha/D)x} \psi'_1 \psi_1 dx = 0. \end{aligned}$$

Similarly, multiplying (4.14) by  $e^{(\alpha/D)x}\psi_1'$  and integrating it in  $(0, L)$ , we have

$$(4.18) \quad -D \int_0^L e^{(\alpha/D)x} \psi_{1,x}' \psi_{1,x} dx + \int_0^L e^{(\alpha/D)x} h(x) \psi_1' \psi_1 dx + \lambda_1 \int_0^L e^{(\alpha/D)x} \psi_1' \psi_1 dx = 0.$$

It follows from (4.17) and (4.18) that

$$(4.19) \quad \lambda_1' = \frac{-\int_0^L e^{(\alpha/D)x} \psi_{1,xx} \psi_1 dx}{\int_0^L e^{(\alpha/D)x} \psi_1^2 dx}.$$

By Lemma 4.7, we have  $\psi_{1,x} < 0$  in  $(0, L)$ . Hence, if  $\alpha \leq 0$ , then

$$(4.20) \quad \begin{aligned} \int_0^L e^{(\alpha/D)x} \psi_{1,xx} \psi_1 dx &= - \int_0^L \psi_{1,x} (e^{(\alpha/D)x} \psi_1)_x dx \\ &= - \int_0^L e^{(\alpha/D)x} \psi_{1,x} [\psi_{1,x} + (\alpha/D) \psi_1] dx < 0. \end{aligned}$$

Thus  $\lambda_1' > 0$  for any  $\alpha \leq 0$  and  $D > 0$ . This proves (a).

On the other hand, if  $\lambda_1(D^*, \alpha, h) = 0$  for some  $D^* > 0$ , then the corresponding eigenfunction  $\psi_1$  satisfies

$$(4.21) \quad \begin{cases} D^* \psi_{1,xx} + \alpha \psi_{1,x} + h(x) \psi_1 = 0, & 0 < x < L, \\ \psi_{1,x}(0) = \psi_{1,x}(L) = 0. \end{cases}$$

Multiplying (4.21) by  $e^{(\alpha/D^*)x}$ , and integrating over  $(0, L)$ , we have

$$\int_0^L h(x) \psi_1(x) e^{(\alpha/D^*)x} dx = 0.$$

Thus the decreasing function  $h$  must change sign, i.e.  $h'(x) < 0, h(0) > 0$ . Combining with  $\psi_{1,x} < 0$ , we have

$$\int_0^x h(s) \psi_1(s) ds < \int_0^x h(0) \psi_1(s) ds < h(0) \int_0^x \psi_1(0) ds < h(0) \psi_1(0) L.$$

Next, we integrate (4.21) in  $(0, x)$ , to get

$$D^* \psi_{1,x}(x) + \alpha \psi_1(x) = \alpha \psi_1(0) - \int_0^x h(s) \psi_1(s) ds > [\alpha - h(0)L] \psi_1(0) \geq 0,$$

provided that  $\alpha \geq h(0)L$ . By virtue of (4.20), we obtain

$$\int_0^L e^{(\alpha/D^*)x} \psi_{1,xx} \psi_1 dx = -\frac{1}{D^*} \int_0^L e^{(\alpha/D^*)x} \psi_{1,x} (D^* \psi_{1,x} + \alpha \psi_1) dx > 0.$$

It follows then from (4.19) that  $\frac{\partial \lambda_1}{\partial D}(D^*, \alpha, h) < 0$ . This proves (b).  $\square$

**4.4. The Case  $D_1 = D_2, \alpha_1 < \alpha_2$ .** To investigate whether stronger or weaker advection is more beneficial for species to win the competition in the two-species phytoplankton model, we assume the only phenotypic difference between them is the advection rate. To be more precise, we assume  $D_1 = D_2 \equiv D > 0, \alpha_1 < \alpha_2$ . For the rest of this paper, we assume two phytoplankton species have the same growth rates and death rates, i.e.,  $g_1(\cdot) = g_2(\cdot) \equiv g(\cdot)$  and  $d_1 = d_2 \equiv d$ .



*Proof of Theorem 2.2.* By Theorem 4.1, it suffices to establish, for system (1.1)-(1.4) (and that  $k_1 = k_2 = I_0 = 1$ ), the linear instability of  $(0, \tilde{v})$ , and the non-existence of postive steady states.

**Step 1.**  $(0, \tilde{v})$  is linearly unstable.

Recall that  $\tilde{v}$  is the unique positive solution to

$$\begin{cases} D\tilde{v}_{xx} - \alpha_2\tilde{v}_x + [g(\sigma_2) - d]\tilde{v} = 0, & 0 < x < L, \\ D\tilde{v}_x - \alpha_2\tilde{v} = 0, & x = 0, L, \end{cases}$$

where  $\sigma_2(x) = e^{-k_0x - \int_0^x \tilde{v}(s) ds}$ . Since  $\tilde{v} > 0$  is a positive eigenfunction of (4.3) with  $\alpha = \alpha_2$  and  $h(x) = g(\sigma_2) - d$ , we have  $\lambda_1(D, \alpha_2, g(\sigma_2) - d) = 0$ .

It follows from Proposition 4.6 that the stability of  $(0, \tilde{v})$  is determined by the sign of the principal eigenvalue  $\lambda_1(D, \alpha_1, g(\sigma_2) - d)$ . Since  $\alpha_1 < \alpha_2$ , we may apply Lemma 4.8 to yield

$$\lambda_1(D, \alpha_1, g(\sigma_2) - d) < \lambda_1(D, \alpha_2, g(\sigma_2) - d) = 0.$$

Thus  $(0, \tilde{v})$  is linearly unstable.

**Step 2.** The system (1.1)-(1.4) has no co-existence steady states.

Suppose to the contrary that  $(u^*, v^*)$  be a co-existence steady state of (1.1)-(1.4), then we have

$$\begin{cases} Du_{xx}^* - \alpha_1 u_x^* + [g(\sigma^*(x)) - d]u^* = 0, & 0 < x < L, \\ Dv_{xx}^* - \alpha_2 v_x^* + [g(\sigma^*(x)) - d]v^* = 0, & 0 < x < L, \\ Du_x^* - \alpha_1 u^* = 0, \text{ and } Dv_x^* - \alpha_2 v^* = 0, & x = 0, L, \end{cases}$$

where  $\sigma^*(x) = \exp(-k_0x - \int_0^x [u^*(s) + v^*(s)] ds)$ . Let  $h(x) = g(\sigma^*(x)) - d$  so that  $h'(x) < 0$ . Since  $u^*(x) > 0, v^*(x) > 0$ , then

$$\lambda_1(D, \alpha_1, h) = \lambda_1(D, \alpha_2, h) = 0.$$

This is in contradiction with Lemma 4.8, which says that  $\lambda_1$  is strictly monotone increasing in  $\alpha$ . Therefore, the system (1.1)-(1.4) has no co-existence steady state.  $\square$

**4.5. The Case  $D_1 < D_2, \alpha_1 = \alpha_2 \geq [g(1) - d]L$ .** In this and the next subsection, we explore the effect of diffusion on the outcome of competition. According to Lemma 4.9, the monotonicity of the principal eigenvalue  $\lambda_1(D, \alpha, h)$  with respect to  $D$  also depends on the advection rate  $\alpha$ . Here, we first consider that both species have large sinking rates, i.e.,  $\alpha_1 = \alpha_2 \equiv \alpha \geq [g(1) - d]L > 0$ . (Note that  $g(\cdot)$  satisfies (1.5) and as we assume that the semi-trivial steady states exist, so we always have  $g(1) - d > 0$ .)

*Proof of Theorem 2.3.* By Theorem 4.1, it suffices to establish, for system (1.1)-(1.4), the linear instability of  $(\tilde{u}, 0)$ , and the non-existence of positive steady states.

**Step 1.**  $(\tilde{u}, 0)$  is linearly unstable.

First, we observe as before from the equation satisfied by  $\tilde{u}$  that  $\lambda_1(D_1, \alpha, g(\sigma_1) - d) = 0$ , where  $\sigma_1(x) = e^{-k_0x - \int_0^x \tilde{u}(s) ds}$ .

Since  $D_1 < D_2$  and  $\alpha \geq [g(1) - d]L$ , we may apply Lemma 4.9(b) to yield

$$\lambda_1(D_2, \alpha, g(\sigma_1) - d) < \lambda_1(D_1, \alpha, g(\sigma_1) - d) = 0.$$

It follows from Proposition 4.5 that  $(\tilde{u}, 0)$  is linearly unstable.

**Step 2.** The system (1.1)-(1.4) has no co-existence steady states.

Suppose to the contrary that  $(u^*, v^*)$  is a co-existence steady state of (1.1)-(1.4), then we deduce as before,

$$\lambda_1(D_1, \alpha, g(\sigma^*) - d) = \lambda_1(D_2, \alpha, g(\sigma^*) - d) = 0,$$

where  $\sigma^*(x) = \exp(-k_0 x - \int_0^x [u^*(s) + v^*(s)] ds)$ . But this is in contradiction with Lemma 4.9(b), which says that  $D \mapsto \lambda_1(D, \alpha, g(\sigma^*) - d)$  has at most one positive root. Therefore, the system (1.1)-(1.4) has no co-existence steady state.  $\square$

**4.6. The Case  $D_1 < D_2$ ,  $\alpha_1 = \alpha_2 \leq 0$ .** This subsection is devoted to studying whether stronger or weaker diffusion is more beneficial when both species have buoyant rates. Precisely, we assume that  $D_1 < D_2$ ,  $\alpha_1 = \alpha_2 \equiv \alpha \leq 0$ .

*Proof of Theorem 2.4.* By Theorem 4.1, it suffices to establish, for system (1.1)-(1.4), the linear instability of  $(0, \tilde{v})$ , and the non-existence of positive steady states.

**Step 1.**  $(0, \tilde{v})$  is linearly unstable.

First, we observe as before from the equation satisfied by  $\tilde{v}$  that  $\lambda_1(D_2, \alpha, g(\sigma_2) - d) = 0$ , where  $\sigma_2(x) = e^{-k_0 x - \int_0^x \tilde{v}(s) ds}$ .

Since  $D_1 < D_2$  and  $\alpha \leq 0$ , we may apply Lemma 4.9(a) to yield

$$\lambda_1(D_1, \alpha, g(\sigma_2) - d) < \lambda_1(D_2, \alpha, g(\sigma_2) - d) = 0.$$

It follows from Proposition 4.6 that  $(0, \tilde{v})$  is linearly unstable.

**Step 2.** The system (1.1)-(1.4) has no co-existence steady states.

We omit the details here as this is similar to Step 2 of the proofs of Theorems 2.3, where we use part (a) of Lemma 4.9 instead of part (b). This completes the proof.  $\square$

**5. Discussion and Numerical Results.** We investigate a nonlocal reaction-diffusion-advection system modeling the growth of two competing phytoplankton species in a eutrophic environment, where nutrients are in abundance and the species are limited by light only for their metabolism. We first demonstrate that the system does not preserve the competitive order in the pointwise sense (Remark 3.10). We introduce a special cone  $\mathcal{K}$  involving cumulative distributions of the population densities, and a generalized notion of super- and subsolutions of (1.1)-(1.4), where the differential inequalities hold in the sense of the cone  $\mathcal{K}$ . A comparison principle is then established for the super- and subsolutions, which implies the monotonicity of the semiflow of (1.1)-(1.4) with respect to the cone  $\mathcal{K}$  (Theorem 2.1). From a theoretical point of view, this paper introduces a new class of reaction-diffusion models with order-preserving property, which may be of independent interest [35].

A first application of the monotonicity result yields a simple proof of the existence and global attractivity of the unique positive steady state (or time-periodic solution) to the single species problem (Proposition 3.11). A second application concerns the dynamics of two competing phytoplankton species, as modeled by (1.1)-(1.4), in which sufficient conditions for local (Propositions 4.5 and 4.6) and global stability of semi-trivial steady states (Theorems 2.2-2.4) are obtained.

Consider system (1.1)-(1.4) and fix  $D_1 < D_2$  and  $\alpha_1 = \alpha_2 \equiv \alpha$ . Theorems 2.3 and 2.4 say that  $(\tilde{u}, 0)$  is globally asymptotically stable for  $\alpha = 0$ , and  $(0, \tilde{v})$  is globally asymptotically stable for  $\alpha = [g(1) - d]L$ , which means there is an exchange of stability between the semi-trivial steady states as  $\alpha$  varies from 0 to  $[g(1) - d]L$ . In this particular case, we conjecture that there exist two constants  $\alpha_{min}$  and  $\alpha_{max}$  such that the following statements hold.

- When  $\alpha \leq \alpha_{min}$ , the semi-trivial steady state  $(\tilde{u}, 0)$  is globally asymptotically stable.

- When  $\alpha_{min} < \alpha < \alpha_{max}$ , there exists a unique positive steady state  $(u^*, v^*)$  which is globally asymptotically stable.
- When  $\alpha \geq \alpha_{max}$ , the semi-trivial steady state  $(0, \tilde{v})$  is globally asymptotically stable.

In the following, we present some numerical result in support of this conjecture. See Figure 2.

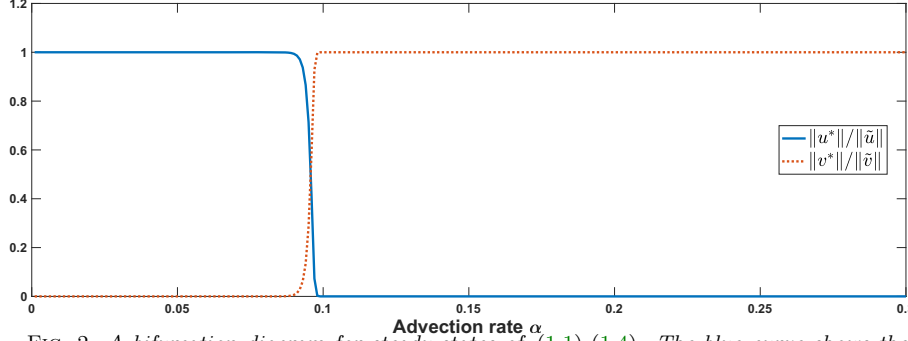


FIG. 2. A bifurcation diagram for steady states of (1.1)-(1.4). The blue curve shows the ratio  $\|u^*\|_{L^1}/\|\tilde{u}\|_{L^1}$ , and the red curve shows the ratio  $\|v^*\|_{L^1}/\|\tilde{v}\|_{L^1}$  as  $\alpha$  varies from 0 to 0.3, where  $(u^*, v^*)$  is the stable steady state, and  $(\tilde{u}, 0)$  and  $(0, \tilde{v})$  are the two semi-trivial steady states. The parameters are chosen as  $D_1 = 1, D_2 = 5, d_1 = d_2 = 0.001, g_1(I) = g_2(I) = \frac{mI}{a+I}, m = 1, a = 10, I_0 = 1, k_0 = k_1 = k_2 = 0.001, L = 100$ .

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## REFERENCES

- [1] R. S. Cantrell and C. Cosner, On the steady-state problem for the Volterra-Lotka competition model with diffusion. *Houston J. Math.* 13 (1987), no. 3, 337-352.
- [2] R. S. Cantrell and C. Cosner, *Spatial Ecology via Reaction-Diffusion Equations*, Wiley Ser. Math. Comput. Biol., John Wiley and Sons, Chichester, UK, 2003.
- [3] Y. Du and S.-B. Hsu, Concentration phenomena in a nonlocal quasi-linear problem modeling phytoplankton I: existence, *SIAM J. Math. Anal.* 40 (2008) 1419-1440.
- [4] Y. Du and S.-B. Hsu, Concentration phenomena in a nonlocal quasi-linear problem modeling phytoplankton II: limiting profile, *SIAM J. Math. Anal.* 40 (2008) 1441-1470.
- [5] Y. Du and S.-B. Hsu, On a nonlocal reaction-diffusion problem arising from the modeling of phytoplankton growth, *SIAM J. Math. Anal.* 42 (2010) 1305-333.
- [6] Y. Du, S.-B. Hsu and Y. Lou, Multiple steady-states in phytoplankton population induced by photoinhibition, *J. Differential Equations* 258 (2015) 2408-2434.
- [7] Y. Du and L. Mei, On a nonlocal reaction-diffusion-advection equation modeling phytoplankton dynamics, *Nonlinearity* 24 (2011) 319-349.
- [8] U. Ebert, M. Arrayas, N. Temme, B. Sommeijer and J. Huisman, Critical condition for phytoplankton blooms, *Bull. Math. Biol.* 63 (2001) 1095-1124.
- [9] L.C. Evans, *Partial differential equations*. Graduate Studies in Mathematics, 19. American Mathematical Society, Providence, RI, 1998.

- [10] D. Henry, "Geometric Theory of Semilinear Parabolic Equations," Lecture Notes in Math., Vol. 840, Springer-Verlag, Berlin, 1981.
- [11] P. Hess, Periodic-Parabolic Boundary Value Problems and Positivity, Pitman Res. Notes Math. Ser., vol. 247, Longman Scientific Technical, Harlow, UK, 1991.
- [12] P. Hess and A. C. Lazer, On an abstract competition model and applications, Nonlinear Analysis T. M. A. 16 (1991), 917-940.
- [13] S.-B. Hsu and Y. Lou, Single species growth with light and advection in a water column, SIAM J. Appl. Math. 70 (2010) 2942-2974.
- [14] S.-B. Hsu, L. Mei and F.-B. Wang, On a nonlocal reaction-diffusion-advection system modeling the growth of phytoplankton with cell quota structure, J. Differential Equations 259 (2015) 5353-5378.
- [15] S.-B. Hsu, H. Smith and P. Waltman, Competitive exclusion and coexistence for competitive systems on ordered Banach spaces. Trans. Am. Math. Soc. 348 (1996) 4083-4094.
- [16] J. Huisman, P. van Oostveen and F. J. Wessing, Species dynamics in phytoplankton blooms: Incomplete mixing and competition for light, Amer. Naturalist, 154 (1999) 46-67.
- [17] J. Huisman, P. van Oostveen and F.J. Wessing, Critical depth and critical turbulence: two different mechanisms for the development of phytoplankton blooms, Limnol. Oceanogr. 44 (1999) 1781-1787.
- [18] J. Huisman, N.N. Pham Thi, D.K. Karl and B. Sommeijer, Reduced mixing generates oscillations and chaos in oceanic deep chlorophyll, Nature 439 (2006) 322-325.
- [19] J. Húska, Harnack inequality and exponential separation for oblique derivative problems on Lipschitz domains, J. Differential Equations 226 (2006) 541-557.
- [20] H. Ishii and I. Takagi, Global stability of stationary solutions to a nonlinear diffusion equation in phytoplankton dynamics, J. Math. Biol. 16 (1982) 1-24.
- [21] C. A. Klausmeier and E. Litchman, Algal games: the vertical distribution of phytoplankton in poorly mixed water columns, Limnol. Oceanogr. 46 (2001) 1998-2007.
- [22] C. A. Klausmeier and E. Litchman, S. A. Levin, Phytoplankton growth and stoichiometry under multiple nutrient limitation, Limnol. Oceanogr. 49 (2004) 1463-1470.
- [23] J. T. Kirk, Light and photosynthesis in aquatic ecosystems. Cambridge university press, 1994.
- [24] T. Kolokolnikov, C. H. Ou and Y. Yuan, Phytoplankton depth profiles and their transitions near the critical sinking velocity, J. Math. Biol. 59 (2009) 105-122.
- [25] K.-Y. Lam and D. Munther, A remark on the global dynamics of competitive systems on ordered Banach spaces. Proc. Am. Math. Soc. 144 (2016) 1153-1159.
- [26] G. M. Lieberman, Second Order Parabolic Differential Equations, World Scientific Publishing, 2003.
- [27] E. Litchman, C. A. Klausmeier, J. R. Miller, O. M. Schofield and P. G. Falkowski, Multi-nutrient, multi-group model of present and future oceanic phytoplankton communities, Biogeosci. 3 (2006) 585-606.
- [28] M. J. Ma and C. H. Ou, Existence, uniqueness, stability and bifurcation of periodic patterns for a seasonal single phytoplankton model with self-shading effect, J. Differential Equations 263 (2017) 5630-5655.
- [29] L. Mei and X. Y. Zhang, Existence and nonexistence of positive steady states in multi-species phytoplankton dynamics, J. Differential Equations 253 (2012) 2025-2063.
- [30] A. Pazy, Semigroups of Linear Operators and Applications to Partial Differential Equations, Springer-Verlag, New York, 1983.
- [31] R. Peng and X.-Q. Zhao, Effects of diffusion and advection on the principal eigenvalue of a periodic parabolic problem with applications, Calc. Var. 54 (2015) 1611-1642.
- [32] R. Peng and X.-Q. Zhao, A nonlocal and periodic reaction-diffusion-advection model of a single phytoplankton species, J. Math. Biol. 72 (3) (2016) 755-791.
- [33] N. Shigesada and A. Okubo, Analysis of the self-shading effect on algal vertical distribution in natural waters, J. Math. Biol. 12 (1981) 311-326.
- [34] H. L. Smith, "Monotone Dynamical Systems, An Introduction to the Theory of Competitive and Cooperative Systems," Mathematical Surveys and Monographs, Vol. 41, Amer. Math. Soc., Providence, RI, 1995.
- [35] H. L. Smith, Monotone dynamical systems: reflections on new advances & applications, Discrete Contin. Dyn. Syst. 37 (2017) 485-504.
- [36] K. Yoshiyama, J. P. Mellard, E. Litchman and C. A. Klausmeier, Phytoplankton competition for nutrients and light in a stratified water column Am. Naturalist 174 (2009) 190-203.
- [37] X.-Q. Zhao, Dynamical Systems in Population Biology. Springer, New York, 2003.