# Magnetic field enhanced coercivity of Fe nanoparticles embedded in antiferromagnetic MnN films 

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#### Abstract

The exchange coupling effect in nanocomposite samples with ferromagnetic (FM) body-centred tetragonal (bct) Fe nanoparticles (NPs) ( $\sim 20 \mathrm{~nm}$ ) embedded in an MnN antiferromagnetic (AFM) matrix is investigated. Both the bct Fe NPs control sample and the $\mathrm{MnN}(\times \mathrm{nm}) \backslash \mathrm{bct} \mathrm{Fe} \mathrm{NPs} \backslash \mathrm{MnN}(20-35 \mathrm{~nm}) \backslash \mathrm{Ta}(5 \mathrm{~nm})$ nanocomposite samples are synthesized by a gas-phase condensation method. Both the coercivity and the remanence ratio of the nanocomposite samples are significantly enhanced compared with the bct Fe NPs sample. The coercivity and the remanence ratio increase with MnN thickness up to 30 nm and then decrease. Additionally, the exchange coupling strength between the bct Fe NPs and AFM MnN matrix is enhanced by magnetic field training. The magnetic field training leads to a higher coercivity and remanence ratio after one cycle of field-cooled hysteresis loop measurement. The coercivity of the composite sample increases by $80 \%$, while the remanence ratio increases by around $30 \%$ compared with the bct Fe NPs sample. This work may provide an alternative approach to the design and manufacturing of FM-AFM exchange-coupled permanent magnets.


Keywords: magnetic nanoparticles, nanocomposite, bct Fe nanoparticles, antiferromagnetic MnN , permanent magnet

S Supplementary material for this article is available online
(Some figures may appear in colour only in the online journal)

## 1. Introduction

Embedding ferromagnetic nanoparticles (NPs) in a matrix could diversify their applications due to the interaction between the NPs and the matrix material [1-6]. In particular, the exchange coupling between antiferromagnets (AFM) and ferromagnets (FM) plays an essential role in many magnetoelectronic devices, such as the magnetic reader for hard disk drives [7], magnetoresistive random access memory [8], and magnetic biosensors [9]. In these applications, AFMs are used to magnetically pin the adjacent FM layer to induce a
unidirectional anisotropy. For AFM materials with high magnetic anisotropy, a hysteresis loop shift is produced [10]. While for AFM materials with low magnetic anisotropy, the AFM spins could be dragged by FM spins during magnetization reversal. In this process, extra energy is required to flip the magnetizations in both the positive and negative branches of the hysteresis loop. The hysteresis loop, as a result, shows an enhanced coercivity and remanence ratio [11]. Experimentally, the enhancement of exchange bias and coercivity could be observed simultaneously due to the distribution of grain size, defects, and anisotropies of the AFM and FM materials. More
details of exchange interactions between FM and AFM could be found in the review papers of this topic [12-14].

Although the interface exchange interaction between AFMs and FMs was first demonstrated from $\mathrm{Co} @ \mathrm{CoO}$ coreshell NPs in 1956 [15], there are limitations in the variety of this core-shell system, where the AFM shells are usually composed from native oxidation, with nitrides or sulfides for the FM core metal [16-18]. In the following decades, there have been few research works on the NPs exchange bias systems, while the main research focus has shifted towards thin films [19] as there are more choices of AFM materials for thin films than in the core-shell structures.

Recently, the study of the exchange effects on NP systems is attracting renewed interest, especially for the systems of FM NPs embedded in an AFM matrix due to significant improvements of fabrication techniques, such as the gas-phase condensation method [4, 20, 21]. These new techniques provide various AFM matrices for FM NPs which had hindered the research on FM/AFM core-shell systems previously. New application perspectives such as thermal stability enhancement of FM NPs, provide approaches to fabricate ultra-high areal density magnetic recording media [2], and to enhance the hard magnetic properties for magnets [1, 22].

In order to prepare nanocomposite samples with NPs embedded in a matrix, a system that could fabricate both NPs and matrix is required. The gas-phase condensation (GPC) system is an ideal candidate for preparing samples with NPs embedded in a matrix [4-6, 23, 24]. The GPC method is a well-controlled high vacuum technique in which NPs produced by a gas aggregation source are co-deposited with a matrix material using a dc sputtering source. The schematic drawing of the GPC system we used for preparing the samples in this paper is shown in figure 1 and can also be found in previous reports [25, 26]. The interfaces between FM NPs and the AFM matrix are free of contamination, which is critical for the FM/AFM exchange coupling investigation. Up to now, only a few groups have reported using this method to prepare FM/AFM nanocomposites (Ni NPs in an IrMn matrix, Co NPs in a Cr matrix, Fe NPs in a Cr matrix, etc) [6, 23, 27].

In this paper, body-centered tetragonal (bct) Fe NPs are the FM material and MnN is used for the AFM matrix for the nanocomposite samples. Fe is abundant on earth and its alloys are typically soft magnetic materials when they are in a body-centred-cubic (bcc) phase and they have been widely used in motors and transformers, whereas bct phase Fe shows high magnetocrystalline anisotropy [28, 29]. We have reported that the magnetic anisotropy of bct Fe NPs prepared by the GPC method was seven times higher than that of bcc Fe [30]. Antiferromagnetic MnN has a Néel temperature ( $T_{\mathrm{N}}$ ) of up to 600 K and is orders of magnitude cheaper than PtMn and IrMn . PtMn and IrMn have been commonly used for thin film exchange stacks [31]. Furthermore, MnN shows better corrosion resistance than FeMn and NiMn [32]. The high electrical resistivity of MnN makes it a good candidate for current in-plane magnetoresistance devices. The $\mathrm{MnN}(\times \mathrm{nm})$ \} bet Fe NPs $\backslash \mathrm{MnN}(\times \mathrm{nm}) \backslash \mathrm{Ta}$ sandwich structure, named as $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite, was prepared by the GPC method with x ranging from 20 to 35 nm . Both the coercivity and the


Figure 1. Schematic drawing of the GPC system synthesizing the nanocomposite sample with nanoparticles embedded in matrix.
remanence ratio of the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite sample are enhanced when compared with the bct Fe NPs sample without the MnN matrix. Additionally, it is found that the exchange bias still exists at 10 K in the nanocomposite samples (see supplementary information (SI, available online at (stacks.iop. org/JPhysD/53/035003/mmedia)). We also report that a high magnetic field (up to 50 kOe ) could enhance the exchange coupling between the FM bct Fe NPs and the AFM MnN matrix. It is observed that after measuring the first cycle of field-cooled hysteresis loops, the coercivities of the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples increased. This increment depends on the thickness of the MnN matrix.

## 2. Methods

The details of the sample preparation and characterization are as follows. A dc sputtering based gas-phase condensation source was used to prepare the bct Fe NPs and a dc magnetron sputtering source was used to prepare the MnN thin films, as shown in figure 1. The working principle of the system can be found in [26, 30]. Argon gas was used as the sputtering gas for the bet Fe NPs. Argon and nitrogen gases were used for fabricating the MnN matrix films. A sandwich structure of the nanocomposite samples was made on an Si substrate\MnN $(\times \mathrm{nm})$ ) bct Fe NPs $\backslash \mathrm{MnN}(\times \mathrm{nm})$ \Ta. The thicknesses of the two MnN layers were identical with $x$ ranging from 20 to 35 nm . The deposition time of the bct Fe NPs for the nanocomposite samples was the same in order to produce consistent experimental results. A bct Fe NP sample with a stack of Silbct Fe NPs $\backslash$ Ta was also prepared as the control sample. Here, $\mathrm{Ta}(\sim 5 \mathrm{~nm})$ served as the capping layer to prevent oxidation.

A magnetic field of up to 6 kOe was applied parallel to the Si substrate during the sample fabrication process to induce a preferred magnetic orientation of the bct Fe NPs [11]. X-ray diffraction (XRD) patterns were taken with a Bruker 2D discover micro-diffractometer. The morphologies of the samples were measured by an FEI T12 ( 120 kV ) transmission electron microscope (TEM). Magnetic characterization was done using a physical property measurement system (PPMS) integrated with a vibrating sample magnetometer (VSM) under varying temperatures from 10 K to 300 K .


Figure 2. Hysteresis loops of bct Fe NPs measured at 300 K (black line) and 10 K (red line), respectively. The inset shows the TEM image of the bct Fe NPs.

## 3. Results and discussion

### 3.1. Magnetic properties and morphology of bct Fe NPs

The morphology and magnetic properties of the bct Fe NPs were investigated. The average size of the bct Fe NPs was around 20 nm as shown in the TEM image from the inset of figure 2. These NPs showed a spherical shape due to the high quenching rate during the formation of the bct Fe NPs. More details of the bct Fe NPs preparation and characterizations can be found in a previous report [30]. Hysteresis loops of the bct Fe NP sample were measured at 300 K and 10 K , respectively. The coercivity $\left(H_{\mathrm{c}}\right)$ was 1.6 kOe at 300 K and 3.2 kOe at 10 K , as shown in figure 2 . Kinks were observed from the 10 K hysteresis loop since there were also some Fe NPs with body-centered cubic (bcc) structures in the bct Fe NP sample. The presence of bcc Fe NPs was due to the inhomogeneous magnetic field distribution along the etching track of the Fe target. The inhomogeneous field distribution could cause different quenching rates along the etching track of the Fe target. As a result, this led to bct Fe NPs at higher quenching rate regions and bcc Fe at lower quenching rate regions. These two phases were also confirmed by the XRD patterns, which can also be found in our previous report [30].

### 3.2. Crystal structure of MnN matrix

MnN thin films were prepared using the dc sputtering gun (matrix gun) as shown in figure 1. A silicon wafer was used as the substrate for the MnN films. The crystal structure of the MnN thin films was characterized by XRD, as shown in figure 3 in which $\mathrm{MnN}(002)$ and $\mathrm{Si}(400)$ peaks were observed. The pattern was consistent with the MnN XRD pattern reported previously [32], where MnN is the face-centered tetragonal structure and antiferromagneric phase (AFM). The exchange bias and the enhancement of coercivity of the $\mathrm{Fe}-$ MnN nanocomposite samples was observed, further demonstrating the nature of MnN as an antiferromagnetic material.


Figure 3. XRD pattern of the MnN thin film deposited by the dc sputtering gun.

### 3.3. Magnetic properties of $\mathrm{Fe}-\mathrm{MnN}$ nanocomposites

3.3.1. Effects of MnN thickness. The magnetic properties of the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples with different thicknesses of MnN layers were also characterized and compared with the bct Fe NPs. Hysteresis loops were measured for all nanocomposite samples at 300 K and 10 K , respectively. The results shown in figure 4 were from the hysteresis loops measured at the second cycle. Figure 4(a) showed the typical hysteresis loops of the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite with 30 nm MnN . An enhanced $H_{\mathrm{c}}$ and remanence ratio ( $M_{\mathrm{r}} / M_{\mathrm{s}}$ ) was obtained compared with the bct Fe NP sample (figure 2). The $H_{\mathrm{c}}$ has increased by $25 \%$ at 300 K and $82 \%$ at 10 K and the $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio has increased by $21 \%$ at 300 K and $30 \%$ at 10 K . Figure 4(b) summarizes the MnN thickness dependence of the $H_{\mathrm{c}}$ and the $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio of the nanocomposite samples. The results of the bct Fe NP sample are shown as the zero nm MnN in figure 4(b). The $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio increased with MnN thickness up to 30 nm and then decreased significantly. This was in good agreement with the polycrystalline bilayer structure with MnN critical thickness around 30 nm [32].

The differences in $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio from the $\mathrm{Fe}-\mathrm{MnN}$ samples might be due to the change of the domain state magnetization of the AFM MnN layer. A model was proposed by Ali et al [33] to understand the thickness dependence of the coercivity and the exchange bias field in an AFM-FM exchange coupled system. There was a critical thickness for the domain wall stability of AFM materials. For an AFM layer below this critical thickness, the domain state magnetization of the AFM was reversible. In this case, additional coercivity with no exchange bias was obtained. When the AFM layer was at critical thickness, some of the domain state magnetizations became irreversible and exchange bias was observed along with reduced coercivity. Beyond this critical thickness, it became more difficult to form domain walls in the AFM layer, thus, larger domains may develop, which could decrease the bias field and coercivity. As discussed above, both exchange bias and coercivity enhancements could be observed experimentally due to the distribution of grain size,


Figure 4. (a) Hysteresis loops of MnN nanocomposite when MnN is 30 nm thick. (b) $H_{\mathrm{c}}$ of MnN nanocomposites with different MnN thickness $(0,20,25,30$, and 35 nm$)$ and the inset is $M_{\mathrm{r}} / M_{\mathrm{s}}$ versus MnN thickness. MnN ( 0 nm ) stands for bet Fe NPs, shown in figure 2.
defects, and anisotropies of the AFM and FM materials. Some $\mathrm{Fe}-\mathrm{N}$ surfaces might also form at the surface of the bct Fe NPs since N atoms were introduced during the growth of the MnN thin films. The $\mathrm{Fe}-\mathrm{N}$ surface may also contribute to the magnetic properties of the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposites.

The $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples may follow this model. The enhancements of $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio were due to the magnetically reversible or partially reversible MnN layers. The exchange bias was also observed from the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples at 10 K . The exchange bias field was around 150 Oe for nanocomposite samples with 20,25 , and 30 nm MnN layers and decreased significantly for the nanocomposite sample with a 35 nm MnN layer, as shown in the SI figure S 1 . The $M_{\mathrm{s}}$ of $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples reduced slightly $(2 \%-14 \%)$ compared with the $M_{\mathrm{s}}$ measured in the first cycle hysteresis loops. However, both the remanence and remanence ratio were improved in the second cycle hysteresis loops, as shown in SI figure S 2 . The critical thickness of MnN in the nanocomposite samples should be around $30-35 \mathrm{~nm}$ based on the model. The enhanced $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio makes the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposites potential candidates for rare-earth-free magnets. The enhancement of $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ in AFM-FM samples is suitable for achieving higher magnetic energy products $(\mathrm{BH})_{\max }$. However, the $M_{\mathrm{s}}$ of the AFM-FM nanocomposite samples could be reduced compared to the FM phase alone since there is no contribution of $M_{\mathrm{s}}$ from the AFM


Figure 5. (a) The coercivity increase ratio of the first and second hysteresis loops versus MnN thickness measured at 300 K and 10 K , respectively. (b) The $H_{\mathrm{c}}$ enhancement at varying MnN thicknesses.
region. The decreased $M_{\mathrm{s}}$ in the AFM-FM nanocomposite samples is expected to lower the $(\mathrm{BH})_{\max }$. Modelling research is under way to quantify the trade-off between the enhanced $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio and the reduced saturation magnetization after introducing AFM matrix in a magnet design [34].
3.3.2. Magnetic field training effect. The magnetic field training effect, which could help enhance the $H_{\mathrm{c}}$ and $M_{\mathrm{r}} / M_{\mathrm{s}}$ ratio for the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples, was investigated. The training effect usually refers to the gradual degradation of the exchange bias field and coercivity during subsequent hysteresis loop measurements. This effect is due to the deviation of the spin configuration of AFM from the equilibrium structure. The decrease of the exchange bias fields and coercivities with increasing hysteresis loop cycles is due to the rearrangement of the spin configuration to an equilibrium state [35]. For the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples, the magnetic field ( $\sim 6 \mathrm{kOe}$ ) applied during deposition assists the alignment of FM bct Fe NPs. However, it is not large enough to align the AFM MnN matrix. Thus, a higher magnetic field of up to 50 kOe was applied to the nanocomposite sample during the hysteresis loop measurement. The field enhances the exchange coupling between the FM bct Fe NPs and the AFM MnN matrix. Magnetic field cooling ( 50 kOe field) was performed on all the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples from 300 K to 10 K . The hysteresis loops were measured with a magnetic field of up to 40 kOe along the same direction as the magnetic field applied during the growth of the nanocomposite samples. Several cycles of field-cooled ( 50 kOe ) hysteresis loops were measured at both 300 K and 10 K to study the magnetic training effect which is usually observed in AFMFM systems. The coercivities and remanence ratios of the $\mathrm{Fe}-$ MnN nanocomposite samples increased in the second cycle of hysteresis loops, as shown in figure 5(a). The magnetic field around $\sim 6 \mathrm{kOe}$ during sample growth could induce the easy axis alignment for the bct Fe NPs. For the as-deposited FeMnN nanocomposite sample, it might not be at its lowest free energy state. After the 50 kOe magnetic field training, these samples are at their lowest energy state with stable magnetic performances. Thus, the coercivities and remanence ratios of
the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples remained constant for the subsequent cycles, indicating the AFM spin configuration was already in the equilibrium state. The enhancement of $H_{c}$ was due to a better magnetic alignment between the FM bct Fe NPs and the AFM MnN matrix after measurements of the first hysteresis loops at 300 K and 10 K .

A high magnetic field, in principle, could help rearrange the spin configurations of the MnN thin film matrix. The effective anisotropy constant of MnN was reported to be around $2.4 \times 10^{5} \mathrm{erg} \mathrm{cm}^{-3}\left(24 \mathrm{~kJ} \mathrm{~m}^{-3}\right)$ in a thin film sample [36], which is relatively small compared with other AFM materials such as IrMn. The applied magnetic field $\sim 50 \mathrm{kOe}$ helped to at least partially enhance the exchange coupling between the AFM MnN matrix and the FM bet Fe NPs [37, 38]. Thus, the enhancement of coercivity and remanence ratio was obtained for the hysteresis loops measured at the second cycle. For the hysteresis loops measured in the subsequent cycles, the alignment between ferromagnetic nanoparticles and the antiferromagnetic matrix was already built, and there was no further increase of $H_{c}$.

The $H_{c}$ enhancement ratio after field training also depended on the MnN thickness. Figure 5(a) shows the $H_{c}$ enhancement of the nanocomposite samples after the second cycle measurement of hysteresis loops with different MnN thickness at 300 K and 10 K , respectively. The enhancement was observed in all the nanocomposite samples, as shown in figure 5(b) with a minimum of around 30 nm MnN . Therefore, the critical AFM MnN thickness for $\mathrm{Fe}-\mathrm{MN}$ nanocomposite samples should be around $30-35 \mathrm{~nm}$ since the coercivity started decreasing dramatically in that range. When the MnN thickness was 35 nm , the coercivity enhancement ratio increased significantly compared to other nanocomposite samples with thinner MnN layers. However, the coercivity of the sample decreased significantly. For 35 nm MnN , some irreversible and large domains may develop in the AFM MnN layers. In this case, the coercivity was low in the first cycle, as revealed in figure 5(a). After the first cycle of hysteresis loops, some MnN grains were exchange coupled to the bct Fe NPs, leading to an enhanced coercivity in the second cycle. Although the coercivity of the $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite sample with 35 nm MnN was small, the coercivity enhancement ratio was high due to the rearrangements of the AFM MnN spin structures.

## 4. Conclusions

In summary, $\mathrm{Fe}-\mathrm{MnN}$ nanocomposite samples were synthesized using a gas-phase condensation method. The prepared samples demonstrated higher coercivity (with an increase of up to $82 \%$ ) and remanence ratio (up to $30 \%$ higher) than that of bct Fe NPs due to exchange coupling between the FM bct Fe NPs and the AFM MnN matrix. Nanocomposite samples with different MnN thicknesses from 20 nm to 35 nm were investigated. The enhancements of coercivity and remanence ratio were dependent on the MnN thickness, which was due to the reversible part of the AFM MnN matrix. Magnetic field training could enhance the exchange coupling effect between the AFM MnN matrix and the FM bet Fe NPs. Thus, enhanced
coercivity and remanence ratio were obtained after one cycle of hysteresis loop measurements. Such an enhancement could enable $\mathrm{Fe}-\mathrm{MnN}$ nanocomposites to be a potential candidate for rare-earth-free magnets.

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