

1                   **Computer Modelling of Close-to-Ground Tornado Wind-Fields for**  
2                                   **Different Tornado Widths**

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## Abstract

Tangential velocity ( $V_t$ ) of tornadoes is the major parameter that causes building damage. In-field tornado measurements are less reliable at less than 20m above ground level (AGL). Laboratory tornado simulators suggest that swirl ratio (S) and radius ( $r_o$ ) are the major tornado parameters that influence the  $V_t$ . However, due to scaling problems, the laboratory simulators also report the  $V_t$  at greater than 20m AGL. Well-refined computational fluid dynamics (CFD) models can evaluate the  $V_t$  at less than 10m AGL. However, the CFD models are limited to  $r_o=1.0$ km, and the effect of  $r_o$  on the  $V_t$  is not investigated. The aim of this study is to investigate the maximum  $V_t$  for different  $r_o$  close to ground. Simulation results show that increasing  $r_o$  decreases the maximum  $V_t$  with respect to  $V_{r_o}$ . Moreover, by increasing  $r_o$ , the corresponding elevation of occurrence of maximum  $V_t$  ( $z_{max}$ ) will increase. However, for all tornado radii, the  $z_{max}$  is between 20m to 64m AGL. In addition, results show that for all  $r_o$ , the radial  $V_t$  profile has two peaks at  $z<10$ m AGL due to strong shear force close to the ground and at higher elevation the profile transit to Rankine Combined Vortex Model (RCVM).

Keywords: Tornado wind field; CFD, Swirl ratio; tornado simulator; axisymmetric flow

## 1. Introduction

The tangential velocity ( $V_t$ ) profile in the field can be obtained from Doppler radar measurements. Doppler radars have been used to collect the data of over 200 individual tornadoes as reported in Wurman et al. (2013). Wurman et al. (2007) asserted that due to the beam limits, the radar measurements are limited to about 20m above the ground. On the other hand, the engineers are interested at the elevations less than 10m above ground level (AGL), where the typical buildings are located. Mathematical technique of Ground-Based Velocity Track Display (GBVTD) uses data of the Doppler radar measurements to find the  $V_t$  close to the ground. Kosiba and Wurman (2010, 2013) and Refan et al. (2017) used this technique to find the tornado features of actual tornadoes. However, they reported the vertical location of the maximum  $V_t$  ( $z_{max}$ ) occurs between 30m to 200m AGL. In addition, Nolan (2013) claimed that the close to ground  $V_t$  profile of the GBVTDs is affected by debris and thus close to ground,  $V_t$  measurements by the GBVTD are biased. To better understand the tornadic flows, the laboratory simulators or tornado vortex chambers (TVCs) are employed. In these simulators,  $V_t$  is influenced by the following parameters as reported

58 by Davies-Jones (1973): Reynolds number (Re), the aspect ratio (AR), and swirl ratio (S), as  
59 defined below:

$$60 \text{ Re} = V_{ro} H_o / \nu \quad (1)$$

61 Where,  $H_o$  is the inlet height of the chamber and the reference length as shown in Figure 1,  $V_{ro}$  is  
62 the radial velocity at  $H_o$  and  $\nu$  is the kinematic viscosity of air. Using  $\text{Re} \geq 4.5 \times 10^4$  in the TVC  
63 models makes the tornado simulations independent of the  $\text{Re}$  as reported by Refan et al. (2017). In  
64 addition, aspect ratio (AR) is defined as:

$$65 \text{ AR} = H_o / r_o \quad (2)$$

66 Where,  $r_o$  is the radius of the tornado or tornado simulator and is equal to half of its width. Also,

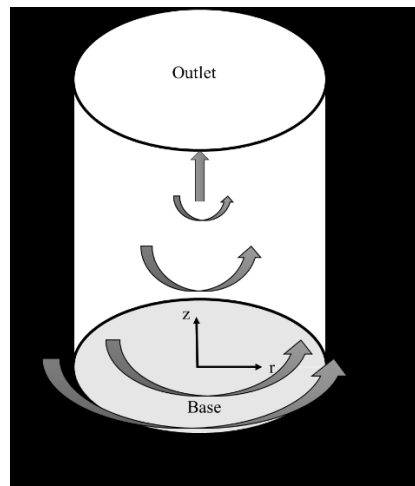
$$67 \text{ Swirl Ratio } S = V_{to} / (2 * \text{AR} * V_{ro}) = V_{to} / (2 * \text{AR} * V_{ro}) \quad (3)$$

68 Here  $V_{to}$  is the tangential velocity at the inlet height  $H_o$ . Equation (3) implies that S, AR,  $V_{ro}$  and  
69  $r_o$  influence the  $V_{to}$ . That is:

$$70 V_{to} = 2SH_o / (r_o V_{ro})$$

71 In our entire write up, variables with a \* like ( $Vt^*$ ) is the non-dimensionalized variable and  
72 without \* is dimensionalized variable.

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Figure 1. Schematic of a simulator and its parameters

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## 78 **1.1. Objective of Current Work**

79 The major objective of this paper is to know the tornadic wind field around 10m from the ground.  
80 This will help to design low rise buildings much better and with lower susceptibility to tornadic  
81 wind hazard. The relatively small size of the laboratory simulators results in large geometric  
82 scaling ratios (Refan et al., 2017) and those simulators cannot evaluate close-to-ground  $V_t$ . In  
83 addition, the scale ratios reported by different researchers are based on either length scale or  
84 velocity scale. The length scale is calculated either using core radius  $r_c$  or location of the maximum  
85 tangential velocity  $z_{max}$  and the velocity scale is based on the maximum tangential velocity. In this  
86 work, none of the scale ratios is introduced. The detailed study conducted by Refan (2014) reports  
87 wind speed from 20 m to 80 m from the ground from both field measurements and experimental  
88 tornado simulator. Hence, it is difficult to collect wind speed around 10m from the ground using  
89 the existing data.

90 Well-refined computational fluid dynamics (CFD) models can compute the  $V_t$  at less than 10m  
91 AGL. Dominguez and Selvam (2017) proposed an axisymmetric CFD model to simulate a tornado  
92 chamber of 1.0km x 2.0km, where  $r_o=1.0$ km,  $H_o=1.0$ km and total height (h) = $2H_o=2.0$ km. They  
93 used a minimum grid spacing (MGS) of  $0.001H_o$  in the vertical axis which amounts to 1.0m from  
94 the ground for  $H_o=1.0$ km. They reported the maximum  $V_t$  occurring at less than 10m AGL.  
95 However, their study was limited to  $r_o=1.0$ km, whereas in actual tornadoes the  $r_o$  may vary. From  
96 observations of different tornadoes by National Weather Service (NWS), it can be inferred that the  
97 significant tornadoes have  $r_o$  in the range of 0.7km to 2.3km (Kashefzadeh, 2018). Therefore, the  
98 specific objectives of this research are:

- 99 1. To vary the  $r_o$  and study its influence on the maximum  $V_t$  with respect to  $V_{r_o}$  and its location.  
100 Hangan and Kim (2008) and Refan (2014) showed that  $V_t$  is dependent on the S parameter and  
101 thus in order to investigate the effect of  $r_o$  on the maximum  $V_t$ , it is necessary to investigate effect  
102 of variation of S on the maximum  $V_t$ .
- 103 2. To investigate effect of  $r_o$  on less than 10m-AGL velocity profile. Typical buildings are located  
104 at elevation of  $z=3.3$ m. Therefore, the maximum  $V_t$  will be investigated at  $z=3.3$ m. Results

105 centered on these objectives will be highly valuable to develop recommendations for safer design  
106 of buildings.

## 107 **2. Numerical Setup**

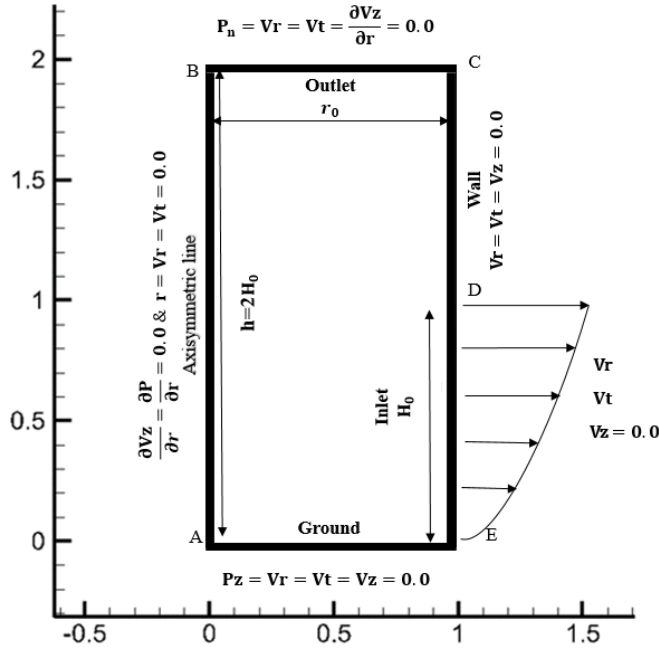
108 *Governing equations:* In this study, non-dimensional Navier-Stokes equation in cylindrical  
109 coordinate system is employed using Large Eddy Simulation (LES) for an axisymmetric model.  
110 This reduces the 3D problem to 2D problem and thus reduces the computational time. Details of  
111 the equations are reported in Kashefzadeh (2018). The governing equations are non-  
112 dimensionalized using  $V_{ro}$  and  $H_o$  as the reference values. The reference value for  $H_o$  and  $V_{ro}$  are  
113 considered to be 1km and 60m/s, respectively. For these reference values, the Re will be greater  
114 than  $1 \times 10^8$ .

115 *Computational domain:* The computational domain in this study is similar to the computational  
116 domain of Dominguez and Selvam (2017). Their non-dimensional computational domain is  $1 \times 2$   
117 ( $r_o = H_o$  &  $h = 2H_o$ ). In this study since  $r_o$  is varied from 0.7km to 2.3 km, the non-dimensional  $r_o^*$   
118 varies from 0.7 to 2.3. The increment of  $r_o^*$  is 0.1, which means that  $r_o^*$  will be 0.7, 0.8, 0.9 and  
119 so on.

120 *Boundary conditions:* The boundary conditions of the axisymmetric model are similar to study of  
121 Wilson and Rotunno (1986) as shown in Figure 2. For cells close to the ground, law of the wall is  
122 used as proposed by Neale et al. (2006).

### 123 **2.1 Mesh of the Computational Domain:**

124 Dominguez and Selvam (2017) used  $MGS = 0.001H_o$  alongside the r- and z- axes. The present study  
125 also uses the same MGS along the r- and z-axis in the vicinity of the axisymmetric line (z-axis).  
126 Then the grid is exponentially increased by a factor of 1.1 and the maximum spacing is considered  
127 to be  $0.1H_o$ . Figure 3 shows the computational domains for non-dimensional  $r_o^*$  of 0.8 and 2. The  
128 grid sizes ranged from  $46 \times 60$  to  $63 \times 60$  in the r and z direction respectively.



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131 Figure 2. Axisymmetric computational domain  
 132 and the boundary conditions  
 133

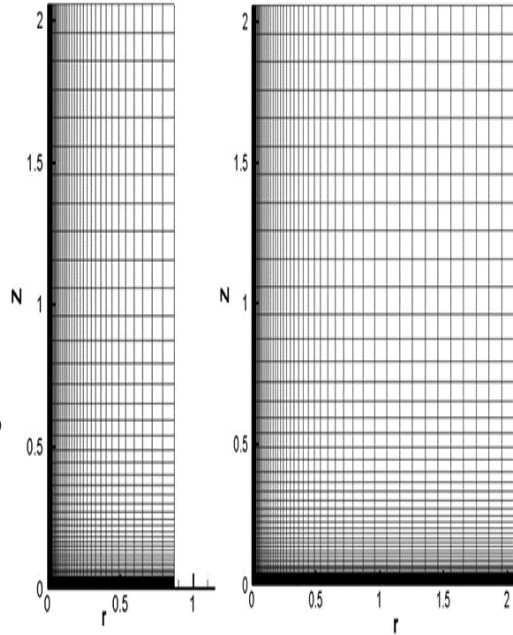


Figure 3. Computational domains for non-  
 dimensional  $r_o$  of (a) 0.8; (b) 2.0

134 **2.2 Radial and Tangential Velocity Components**

135  $V_r^*$  is assumed to vary logarithmically from the ground and the equation for  $V_r^*$  is as follows:

136 
$$V_r^*(z^*) = C_1^* \ln [(z^* + z_o^*) / z_o^*] = C_1 \ln(1 + z^* / z_o^*) \quad (4)$$

137 For open country or Exposure C taking  $z_o = 0.035\text{m}$ , the non-dimensional  $z_o^*$  will be  
 138  $0.035/1000 = 3.5 \times 10^{-5}$ . Keeping the maximum reference  $V_{r_o}^* = 1.0$  the corresponding  $C_1$  becomes:

139 
$$C_1^* = V_{r_o}^* / \ln(1 + H_o^* / z_o^*) = 0.0975 \quad (5)$$

140 Knowing  $V_r^*(z)$ ,  $V_t^*(z)$  is obtained at the inlet by rearranging Equation (3) at height  $z^*$  as follows:

141 
$$V_t^*(z^*) = [2SH_o^* V_r^*(z^*)] / r_o^* \quad (6)$$

142 In Equation 6, the  $V_r^*(z^*)$  and  $H_o^*$  are constant in this work, and the two parameters  $S$  and  $r_o^*$  will  
 143 be varied to determine the  $V_t^*(z^*)$ .

144 **2.3 Solution Scheme**

145 The CFD model uses SOLA-Yaqui type algorithm to solve the equations (Hirt et al, 1975). In this  
146 method, a staggered grid is used where velocities are stored at the nodes and the pressure at the  
147 middle of the cell. In the momentum equation, the diffusion and convection terms are respectively  
148 implicit and explicit. All terms other than convection in the NS equations are approximated using  
149 second order finite volume method (FVM). The QUICK scheme is used for convection term. At  
150 this time, the pressure is solved using SOLA type pressure correction. The advantage of using the  
151 Yaqui-type configuration is to avoid the problem of pressure-velocity decoupling (Harlow and  
152 Welch, 1965; Selvam, 1992). The computer model is run for 5 or 10 time units with a time step of  
153 0.1 to satisfy the CFL condition.

### 154 **3. Results and Findings**

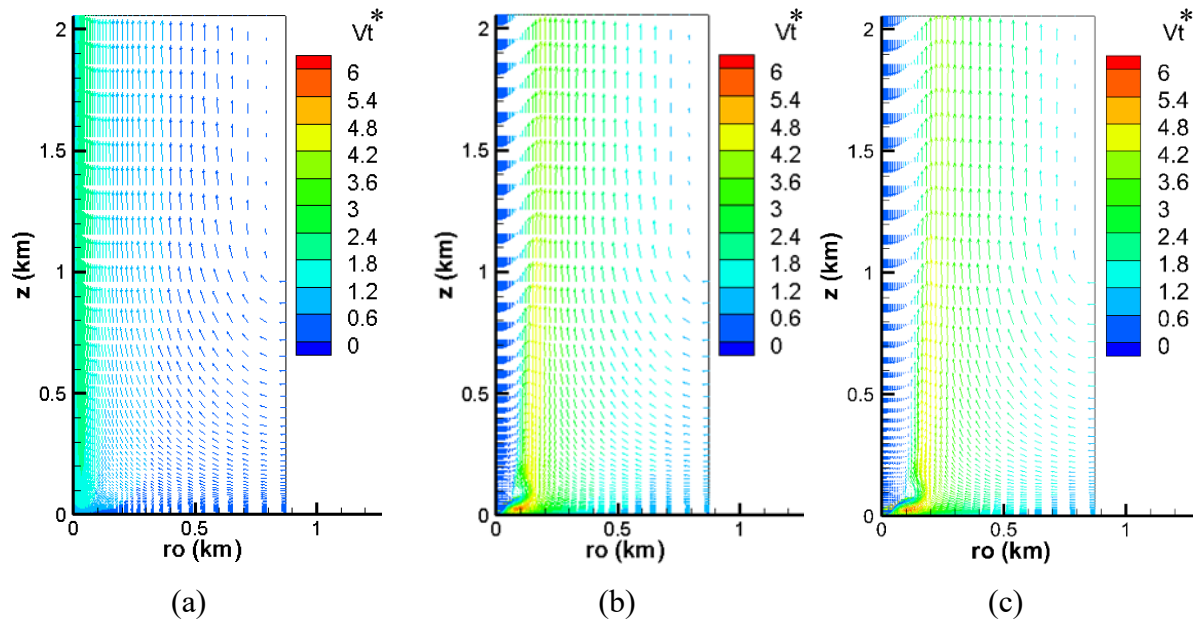
#### 155 **3.1 Swirl ratios for Tornado Touchdown and Maximum $V_t$**

156 For each  $r_o$  in the range of 0.7km to 2.3km, various S parameters in the range of 0.2 to 1.5 are used  
157 and their tornado wind-fields are investigated to determine the swirl ratio that produces the  
158 maximum  $V_t$ .

159 The swirl ratios affect the structure of tornadoes. Hangan and Kim (2008) and Tari et al. (2010)  
160 showed before the touchdown, S is small and tornado has a single-cell structure as shown in  
161 Figures 4 (a) and 6 (a). Then the flow slowly changes from simple jet like flow to touchdown  
162 condition. During the touchdown a vortex breakdown occurs aloft as shown in Figures 4(b) and 6  
163 (b), and afterward the maximum  $V_t$  occurs in transition to a double-cell structure as shown in  
164 Figures 4(c) and 6 (c). To see the flow features for different  $r_o$ ,  $r_o$  varying from 0.8km to 2 km are  
165 considered. Schematics of these stages are also given in Refan (2014). To see clearly, the three  
166 stages of before touchdown, at the touchdown, and double-cell structure, corresponding close up  
167 views are also shown in Figures 5 and 7. The same pattern is observed for all other  $r_o$  but not shown  
168 here.

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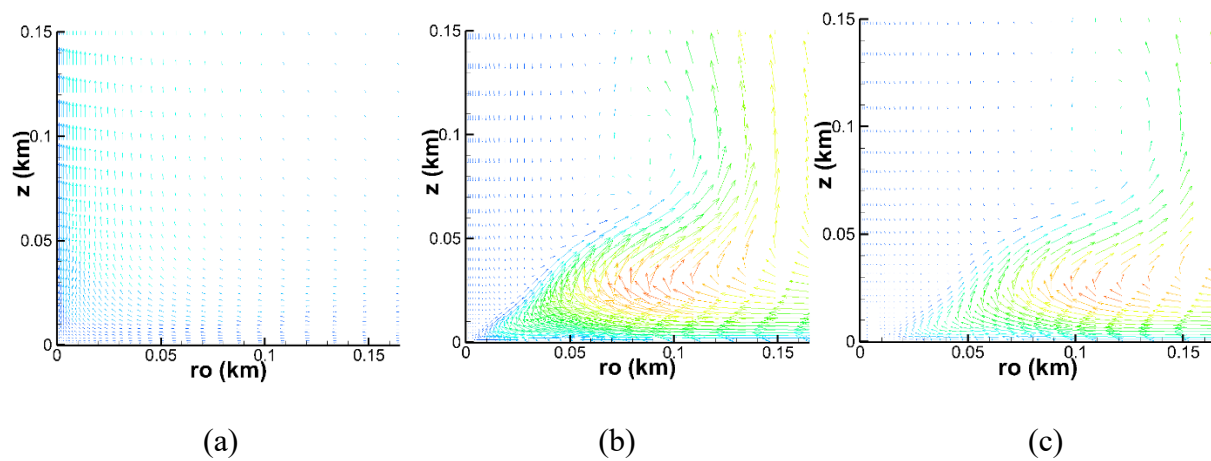


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173 Figure 4. Tornado wind field for  $r_o=0.8\text{km}$  (a)  $S=0.3$ , jet-like and single-cell structure, (b)  $S=0.5$ ,

174 vortex breakdown aloft at touchdown (c)  $S=0.6$ , beyond touchdown, double-cell structure



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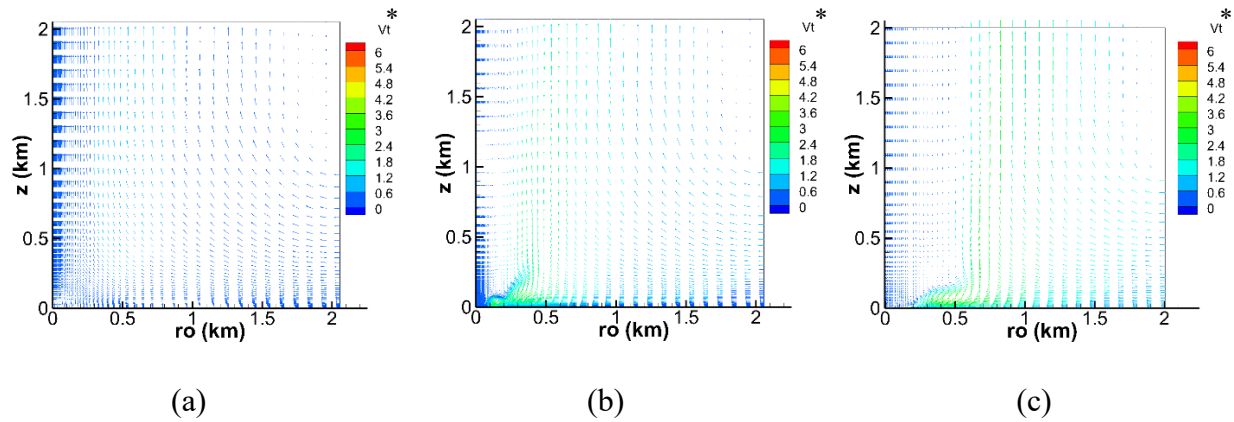
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177 Figure 5. Close up view of Figure 4

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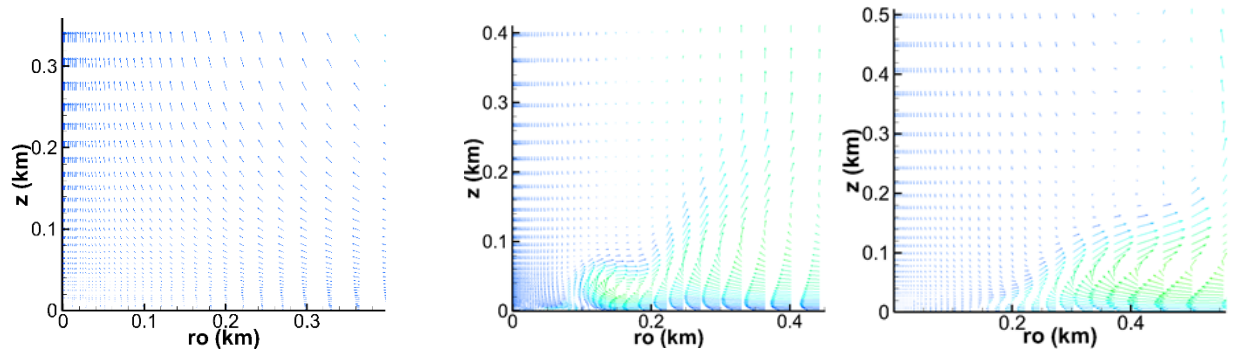




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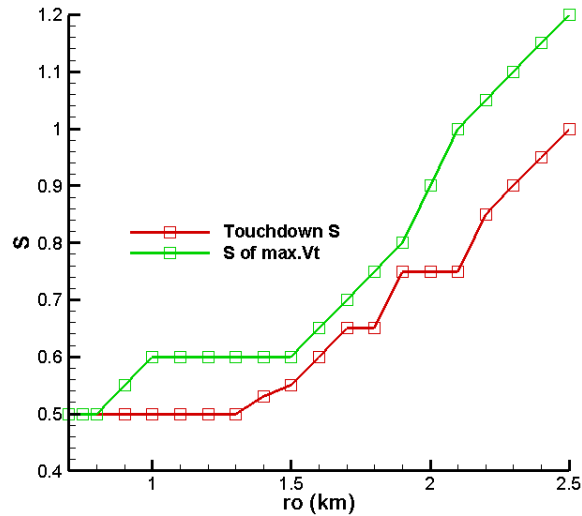
182 Figure 6. Tornado wind field for  $r_o=2.0\text{km}$  (a)  $S=0.3$ , jet-like and single-cell structure, (b)  
 183  $S=0.75$ , vortex breakdown aloft at touchdown and (c)  $S=1.3$ , beyond touchdown, double-cell  
 184 structure



185

186 Figure 7. Close up view of Figure 6

187 The  $S$  for maximum  $V_t$  for each  $r_o$  is determined and plotted with  $S$  for touchdown in Figure 8. It  
 188 can be seen in Figure 8 that the touchdown  $S$  increases by increasing  $r_o$ . The touchdown  $S$  is in the  
 189 range of  $0.40 \leq S \leq 0.9$  for  $0.7\text{km} \leq r_o \leq 2.3\text{km}$ . Similarly, Figure 8 shows that the swirl ratio  $S$  of the  
 190 maximum  $V_t$  increases by increasing  $r_o$  and is in the range of  $0.50 \leq S \leq 1.2$ . This finding is in  
 191 agreement with the previous studies where Lewellen et al (1997) suggested that by increase of  $r_o$ ,  
 192 the  $S$  producing maximum  $V_t$  is likely to increase. Moreover, it can be seen that the  $S$  value  
 193 corresponding to that of the maximum  $V_t$  is always greater than the  $S$  value corresponding to  
 194 touchdown  $S$ , which implies that the maximum  $V_t$  occurs beyond the touchdown. Therefore, in the  
 195 investigation, only swirl ratios that produce tornadoes beyond touchdown are considered because  
 196 these are the ones, which may affect the buildings close to the ground. Therefore, it can be  
 197 concluded that for all radii, the maximum  $V_t$  occurs beyond the touchdown stage.



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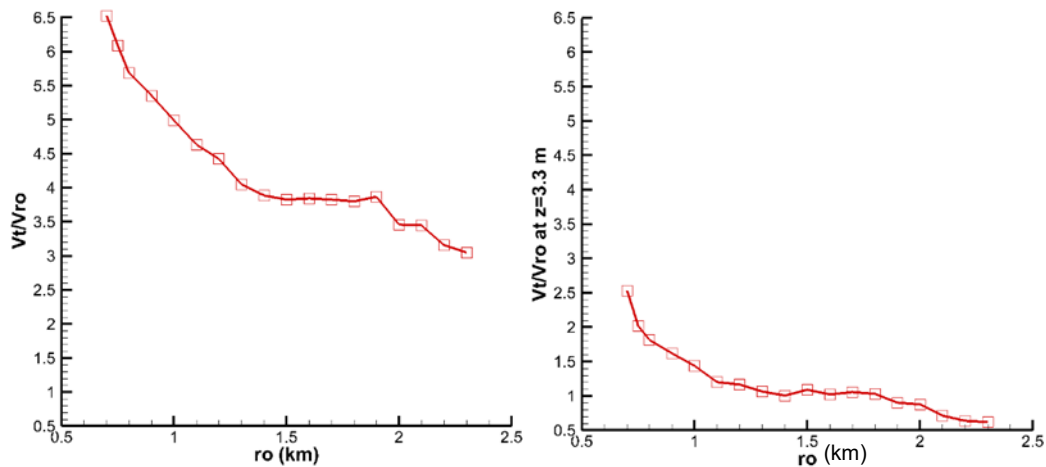
199 Figure 8. Swirl ratios corresponding to the touchdown and maximum  $V_t$  for  $0.7\text{km} \leq r_o \leq 2.3\text{km}$

200 **3.2 Effect of  $r_o$  on Maximum  $V_t$ , Core Radius ( $r_c$ ) and  $z_{\text{max}}$**

201 Figure 9(a) presents the absolute maximum  $V_t$  for  $0.7\text{km} \leq r_o \leq 2.3\text{km}$ . Here on, we will call absolute  
 202 maximum  $V_t$  as  $V_{t_{\text{max}}}$ . It can be seen in this figure that by increasing  $r_o$  from 0.7km to 2.3km, the  
 203  $V_{t_{\text{max}}}$  gradually reduces from  $6.5V_{r_o}$  to almost  $3.5V_{r_o}$ . Likewise, Figure 9(b) shows the maximum  
 204  $V_t$  for various tornado radii at  $z=3.3\text{m}$ , which is the height of a typical low rise building and the  
 205 maximum  $V_t$  gradually reduces from  $2.5V_{r_o}$  to almost  $0.6V_{r_o}$  for  $r_o$  from 0.7km to 2.3km. Similarly,  
 206 Figure 9(c) shows that minimum  $z_{\text{max}}$  is 21m AGL for  $r_o=0.7\text{km}$ , and by increasing  $r_o$ , the  $z_{\text{max}}$   
 207 will also increase. However, for  $r_o \geq 2.0\text{km}$ , the  $z_{\text{max}}$  is constant at 64m. These simulations show  
 208 that the  $z_{\text{max}}$  is in the range of 21m to 64m, whereas radar measurements report  $z_{\text{max}}$  in the range  
 209 of 30m to 200m. Figure 9(d) presents the  $r_c$  for different  $r_o$  where  $r_c$ , is the radial distance of the  
 210 location of the maximum  $V_t$  from the tornado center. It can be seen that  $r_c$  is in the range of  
 211  $100\text{m} \leq r_c \leq 460\text{m}$  for  $0.7\text{km} \leq r_o \leq 2.3\text{km}$ . Table 1 presents a summary of the results.

212 Table 1 shows that for  $r_o=1.0\text{km}$ , the highest peak is  $V_t=4.99V_{r_o}$  at  $S=0.60$  and  $z_{\text{max}}=28.0\text{m}$ .  
 213 Wilson and Rotunno (1986) reported the maximum  $V_t=5.0V_{r_o}$  for  $r_o=1.0\text{km}$ . The reported value is  
 214 for a single study of  $S=0.28$  and  $z_{\text{max}}=1.016\text{km}$ . Lewellen et al. (1997) reported that for  $r_o=1.0\text{km}$   
 215 and  $S=0.94$ , the maximum  $V_t$  is  $6.6V_{r_o}$  at  $z_{\text{max}}=27\text{m}$  AGL. Tari et al (2010) used a laboratory  
 216 simulator and suggested that for  $r_o=1.0\text{km}$ ,  $S=0.68$  produces the maximum  $V_t$  at a height of  $0.34r_o$ .  
 217 The difference of the results from the present study to that of Tari et al. (2010) can be due to  
 218 differences in the geometry of the simulator chamber. The tornado simulator in this work is a based

219 on Ward type, whereas the Tari et al (2010) simulator is similar to Iowa State University. The  
 220 difference in the tornado chamber to touch down condition and other issues needs to be  
 221 investigated further. Likewise, increase in core radius  $r_c$  with increase in chamber radius  $r_o$  is in  
 222 agreement with studies of Ward (1972), Davies-Jones (1973), Jischke and Parang (1974), Church  
 223 et al (1979), Church and Snow (1993), Baker and Church (1979), Tari et al. (2010), Refan (2014),  
 224 and Refan et al. (2017).

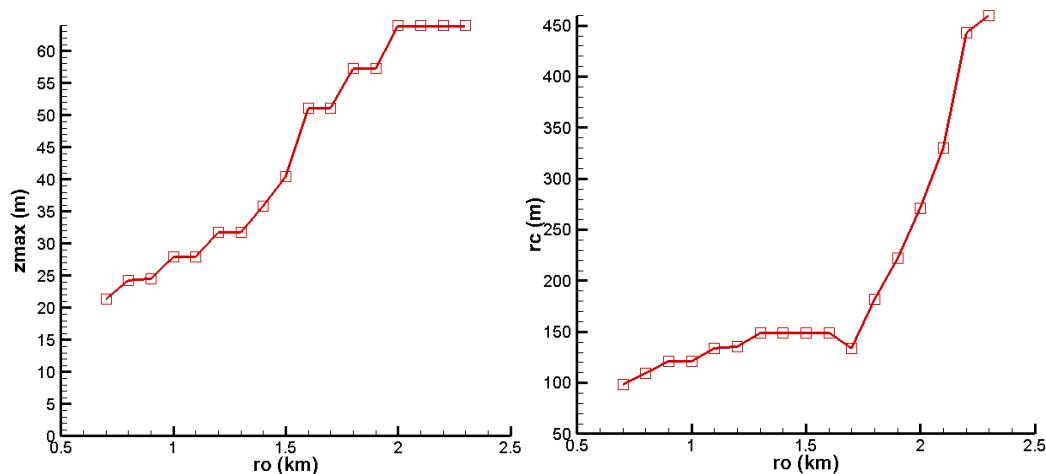


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(a)

(b)



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228

(c)

(d)

229 Figure 9. (a) Absolute maximum  $V_t/V_{r_o}$  for different  $r_o$ , (b) Maximum  $V_t/V_{r_o}$  for different  $r_o$  at  
 230  $z = 3.3$  m, (c)  $z_{max}$  of various  $r_o$  (m); and (d) Core radius of different  $r_o$ .

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232

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Table 1. Summary of the findings for different radii

$r_o(\text{km})$	Touchdown S	S for $V_{t\max}$	$V_{t\max}$ $/V_{r_o}$	Max. $V_t$ $/V_{r_o}$ at $z=3.3$	$z_{\max}(\text{m})$	$r_c(\text{m})$
0.7	0.5	0.5	6.53	2.53	21.4	98.3
0.8	0.5	0.5	5.69	2.02	24.2	109.2
0.9	0.5	0.55	5.35	1.82	24.5	121.1
1	0.5	0.6	4.99	1.62	28	121.1
1.1	0.5	0.6	4.63	1.44	28	134.2
1.2	0.5	0.6	4.43	1.2	31.8	135.2
1.3	0.5	0.6	4.05	1.17	31.8	148.6
1.4	0.53	0.6	3.89	1.07	35.9	148.6
1.5	0.55	0.6	3.83	1.01	40.5	148.6
1.6	0.6	0.65	3.85	1.09	51.1	148.6
1.7	0.65	0.7	3.83	1.02	51.1	134.2
1.8	0.65	0.75	3.87	1.06	57.3	181.9
1.9	0.75	0.8	3.8	1.03	57.3	222.2
2	0.75	0.9	3.46	0.9	64	271
2.1	0.75	1	3.36	0.85	64	330
2.2	0.85	1.1	3.16	0.8	64	443
2.3	0.9	1.2	3.05	0.78	64	460

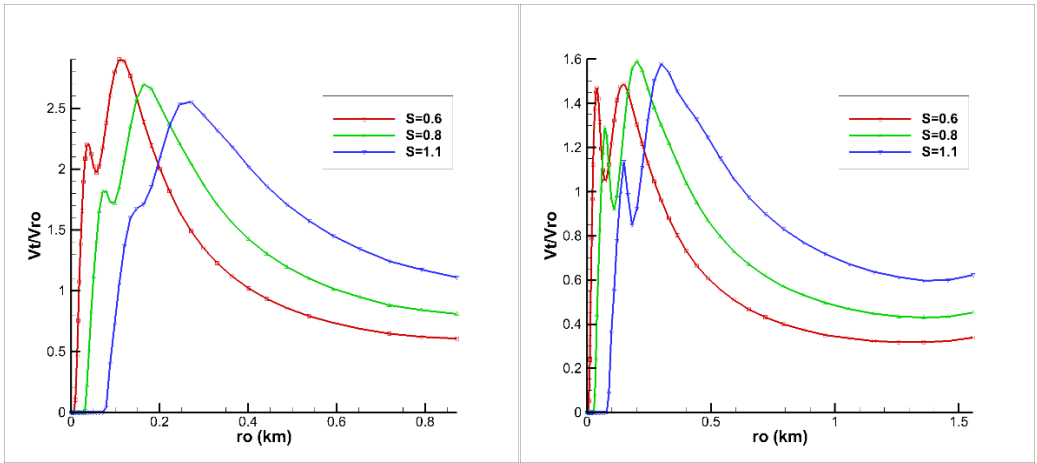
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### 235 3.3 Effect of variation of the swirl ratio on radial $V_t$ profiles

236 In this section, the effect of changing the swirl ratio on radial  $V_t$  profile is investigated. Figures 10  
237 through 13 show the radial  $V_t$  profiles for different  $r_o$  at heights  $z=4.5\text{m}$ ,  $z=9.5\text{m}$ ,  $z=18.5\text{m}$  and  
238  $z=51\text{m}$ . Figure 10 shows the radial  $V_t$  profiles at  $z=4.5\text{m}$  for  $r_o=0.8\text{km}$ ,  $1.5\text{km}$ ,  $1.7\text{km}$ , and  
239  $r_o=2.0\text{km}$ . This figure shows that for all radii at  $z=4.5\text{m}$ , the radial  $V_t$  profile has two peaks and  
240 does not resemble the Rankine Combined Vortex Model (RCVM) profile. Also, one can see that  
241 the double curvature slowly decreases as  $r_o$  increases. Figure 11 shows the radial  $V_t$  profiles for  
242  $r_o=0.8\text{km}$ ,  $1.5\text{km}$ ,  $1.7\text{km}$ , and  $2.0\text{km}$  at  $z=9.5\text{m}$ . Here, for  $r_o$  greater than  $0.8\text{km}$ , double peaks in  
243 the radial profile is distinctly observed. For  $r_o=0.8\text{km}$ , there is a slight kink close to the center. For  
244  $z=18.5\text{m}$  and  $51\text{m}$ , radial velocity profiles are also plotted in Figures 12 and 13. In Figure 12, slight  
245 kink is observed for higher  $r_o$  and in Figure 13 there is no double curvature at all for all radius.

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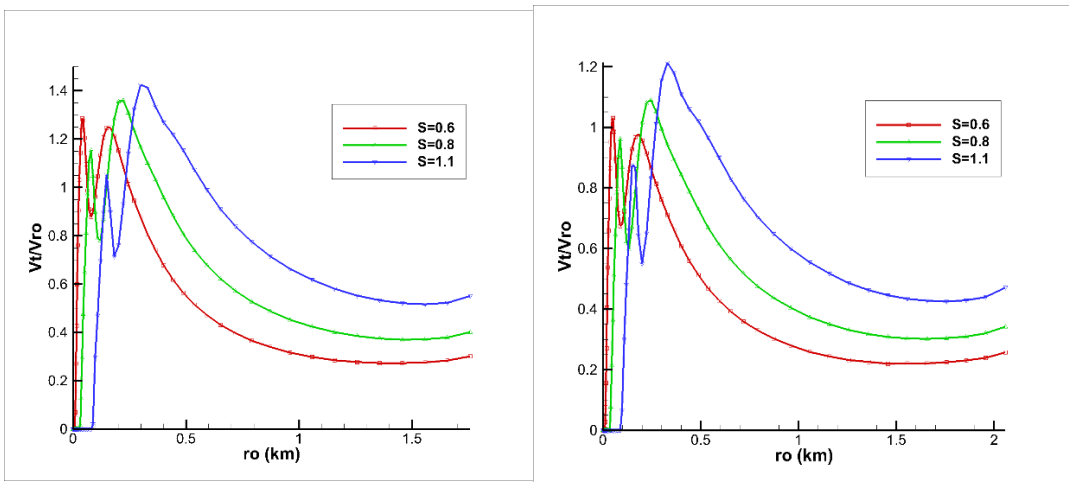


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a)  $r_o=0.8$  km

b)  $r_o=1.5$  km

249



c)  $r_o=1.7$  km

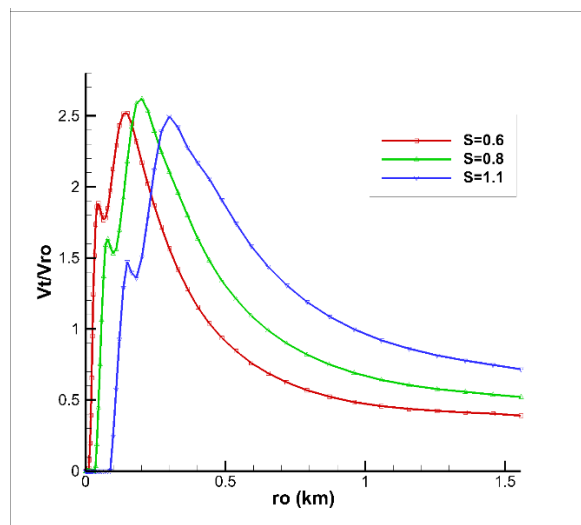
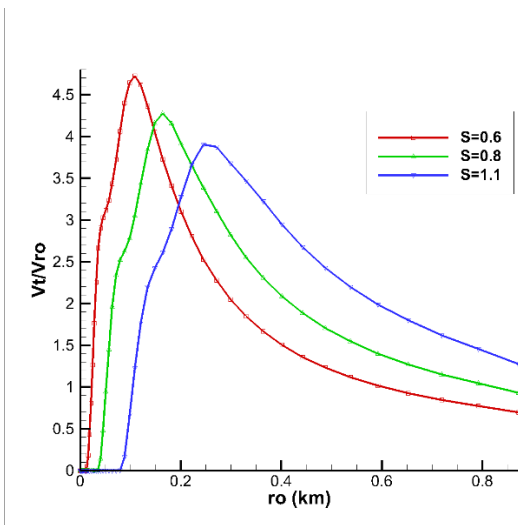
d)  $r_o=2.0$  km

250

251 Figure 10. Radial  $V_t$  profile at  $z=4.5$  m AGL for different tornado radii

252 Refan (2014) also showed the radial  $V_t$  profile having two peaks close to the ground.  
 253 However, Refan (2014) did not make any observation. Also the peaks appeared close to the ground  
 254 and away from the center in their case. These differences may be due to the way vortex chamber  
 255 is built and further detailed studies are warranted. Similarly, Church et al. (1979) showed  
 256 occurrence of two peaks on the velocity profile, but did not report the elevation of occurrence of  
 257 the double-peak. Church et al. (1979) stated that occurrence of the secondary peak on the profile  
 258 is due to the strong shear force close to the ground. It is an important observation which implies  
 259 increased intensity of tornadoes close to ground. Several conclusions are made from this section:

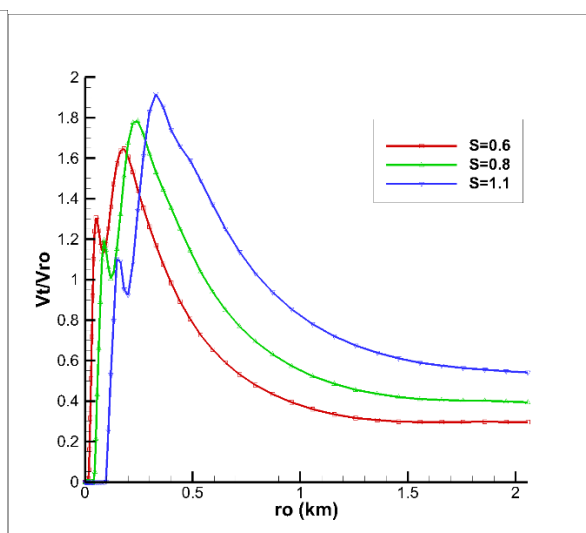
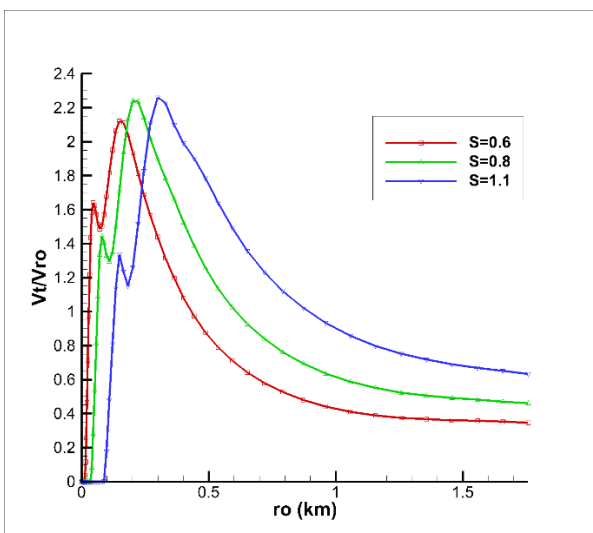
- 260 1. For lower elevation, there are double peaks observed close to the ground for all radius  $r_o$
- 261 considered in this work. When the elevation increases, the double peaks slowly disappear
- 262 from smaller  $r_o$ . Therefore, wider tornadoes have higher intensity due to strong shear
- 263 forces.
- 264 2. Alternatively, these observations imply that RCVM model applies for higher elevation and
- 265 lower  $r_o$ .
- 266 3. For all elevations, it is noted that when the  $r_o$  decreases the maximum  $V_t$  increases or when
- 267  $r_o$  increases the maximum  $V_t$  decreases.



268

269 a)  $r_o=0.8$  km

b)  $r_o=1.5$  km



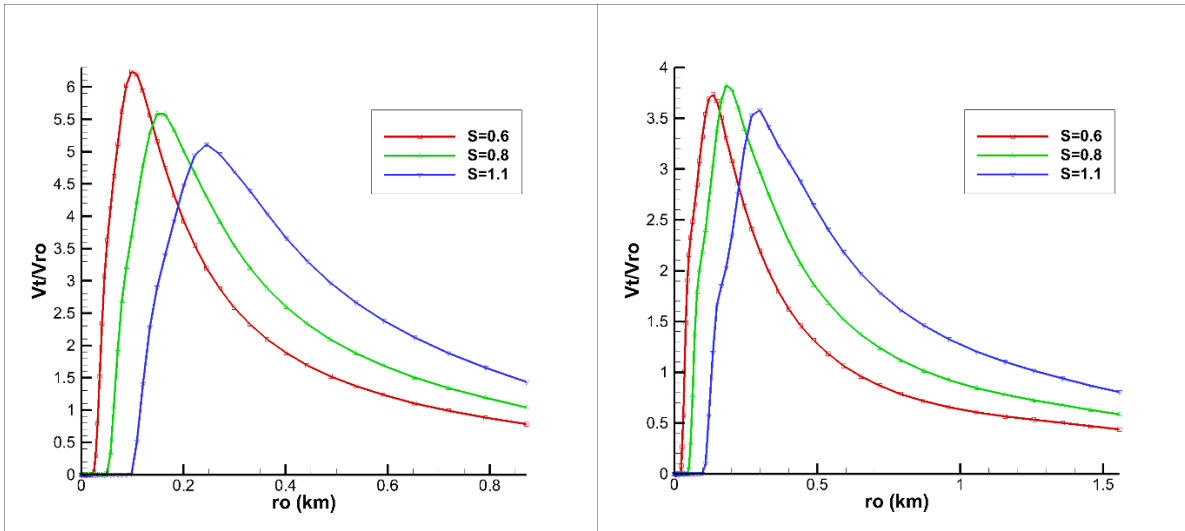
270

271 c)  $r_o=1.7$  km

d)  $r_o=2.0$  km

272 Figure 11. Radial  $V_t$  profile at  $z=9.5\text{m}$  AGL for different tornado radii

273

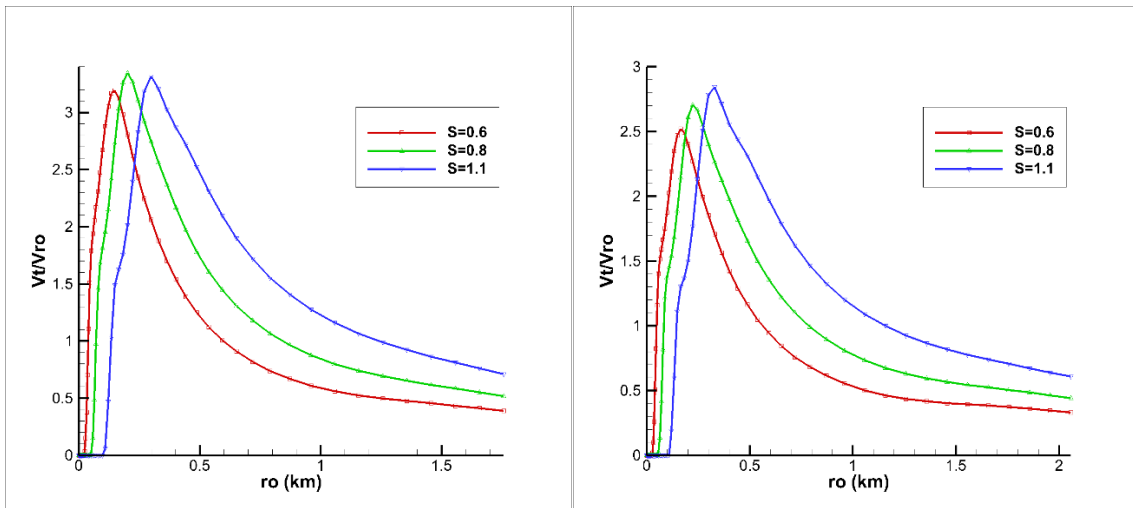


274

a)  $r_o=0.8$  km

b)  $r_o=1.5$  km

275



276

c)  $r_o=1.7$  km

d)  $r_o=2.0$  km

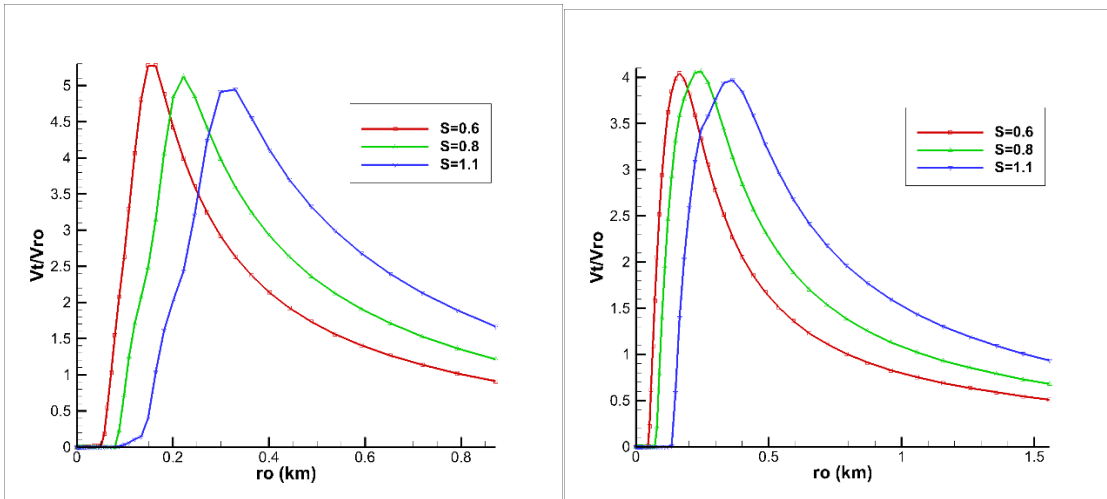
277

278 Figure 12. Radial  $V_t$  profile at  $z=18.5\text{m}$  AGL for different tornado radii

279

280

281

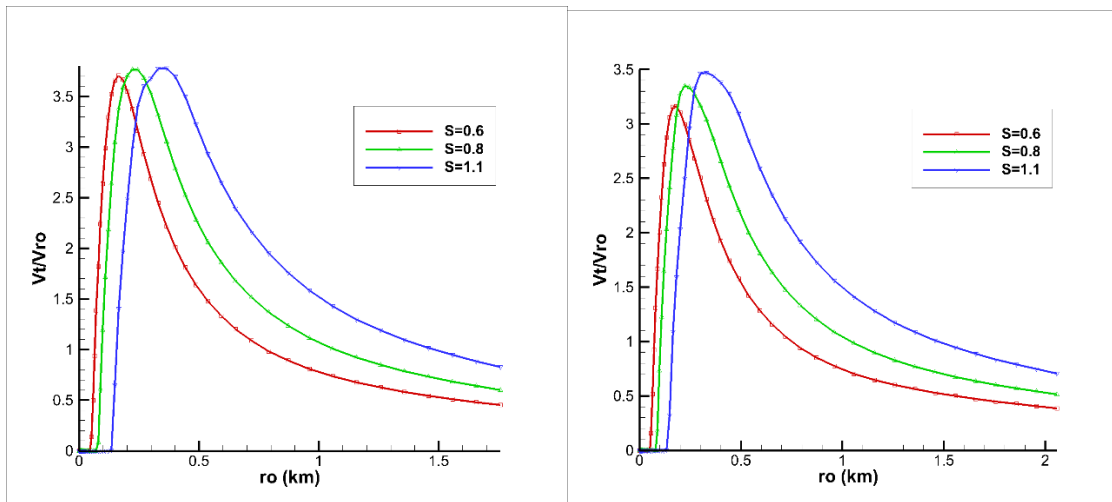


282

283

a)  $r_o=0.8$  km

b)  $r_o=1.5$  km



284

285

c)  $r_o=1.7$  km

d)  $r_o=2.0$  km

286 Figure 13. Radial  $V_t$  profile at  $z=51$  m AGL for different tornado radii

### 287 3.4 Comparison of $V_t$ , $r_c$ and $z_{max}$ against the Actual Tornadoes

288 In this section, the simulation results will be compared against the radar measurements of actual  
 289 tornadoes. For this purpose, the  $r_o$  of actual tornadoes, taken from radar measurements, are used in  
 290 the simulation; the resulting tornado structure,  $r_c$ , and  $z_{max}$  are then compared to the data collected  
 291 from actual tornadoes. This comparison is done for 6 tornadoes as shown in Table 2. Comparison  
 292 of the structure of the tornadoes shows that for all 6 cases, the computational values are in the  
 293 range with radar measurements. Also, comparison of  $r_c$  shows that the radar measurements report



294 a fairly higher value than the simulations. This discrepancy is due to the debris effect in the radar  
295 measurements (Kosiba and Wurman, 2010) which causes the radars measure higher values for the  
296  $r_c$ . The computed  $r_c$  values have error varying from 9% to 37% with respect to field observation.  
297 The error is far more for higher  $r_c$  compared to lower ones. Comparison of the  $z_{max}$  of the radar  
298 measurements to the simulations is possible for three actual tornadoes of Spencer, Manchester,  
299 and Goshen Wyoming tornadoes. Table 2 shows that for these three tornadoes, the  $z_{max}$  of  
300 simulations comply well with the actual tornadoes. Refan et al. (2017) stated that if two scaling  
301 criteria match in comparison of the simulations to the radar measurements, then the simulations  
302 are reliable. Therefore, simulation results in the present study are in close range with field  
303 measurements.

#### 304 4. Conclusions

305 A numerical tornado simulator was proposed in order to investigate effect of the tornado radius  $r_o$   
306 on the maximum tangential velocity  $V_t$  of tornadoes. The following conclusions are made from  
307 the simulations:

- 308 1. Increasing  $r_o$  increases the touchdown swirl ratio and the swirl ratio for maximum  $V_t$  in the range  
309 of  $r_o$  considered for simulation.
- 310 2. Increasing  $r_o$  increases  $z_{max}$ . For  $0.7\text{km} \leq r_o \leq 2.3\text{km}$ ,  $z_{max}$  occurs in the range of  
311  $20\text{m} < z_{max} < 64\text{m}$ , whereas the radar measurements reported  $z_{max}$  in the range of  
312  $30\text{m} < z_{max} < 200\text{m}$ .
- 313 3. Investigating the maximum  $V_t$  at different elevations above and below 10m shows that an  
314 increase of  $r_o$  causes the maximum  $V_t$  to decrease with respect to  $V_{r_o}$ .
- 315 4. For all  $r_o$ , at  $z < 10\text{m}$  AGL, the radial  $V_t$  profile has two peaks. For higher  $z$ , the double peaks in  
316 the radial profile occurs for larger  $r_o$ . In addition, these peaks appear close to the center of the  
317 chamber. This radial profile is different from RCVM flow and the detailed CFD study helped to  
318 visualize this phenomenon. However, the effect of this on force exerted on buildings is yet to be  
319 investigated. Similar double peaks were also observed by Refan (2014) but the double peaks  
320 appear away from the center and this may be due to different type of vortex chamber. Church et  
321 al. (1979) stated that occurrence of the secondary peak on the profile is due to the strong shear

322 force close the ground. More detailed study on the effect of different vortex chamber on double  
323 peak occurrence needs to be conducted.

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329 Table 2. Comparison of the vertical structure,  $r_c$  and  $z_{max}$  of radar measurements to simulations

Tornado	$r_0$ (km)	Technique	Structure	$r_c$ (m)	$z_{max}$ (m)
Spencer (1998)	0.8	Doppler radar (Wurman and Alexander 2005)	Double- celled	120	20
		Simulation results	Double- celled	109	21.4
Manchester (2003)	0.8	Doppler radar (Kuai et al., 2008)	Double- celled	130	20
		Simulation results	Double- celled	109	21.4
Goshen, Wyoming (2009)	1.0	Doppler radar (Wurman et al.2013)	Double- celled	140	30
		Simulation results	Double- celled	121	27.5
Dimmit, Texas (1999)	1.0	Doppler radar (Wurman and Gill, 2000)	Double- celled	150	NA
		Simulation results	Double- celled	121	27.5
El Reno (2013)	2.3	Doppler radar (Bluestein et al, 2015; Wakimoto et al. 2016)	Double- celled	650	NA
		Simulation results	Double- celled	500	65
Bridge Creek Moore(1999)	0.8	Doppler radar (Burgess et al., 2002)	Double- celled	175	NA
		Simulation results	Double- celled	110	21.4

331 **References**

- 332 Baker, G.L., Church, C.R. (1979). Measurements of core radii and peak velocities in modelled  
333 atmospheric vortices. *Journal of Atmospheric Sciences*, 36, 2413-2424.
- 334 Bluestein, H.B., Snyder, J.C., Houser, J.B. (2015). A Multiscale Overview of the El Reno,  
335 Oklahoma, Tornadoic Supercell of 31 May 2013. *Weather and Forecasting*, 30, 525–552,  
336 <https://doi.org/10.1175/WAF-D-14-00152.1>.
- 337 Burgess, D.W., Magsig, M.A., Wurman, J., Dowell, D.C., Richardson, Y. (2002) Radar  
338 observations of the 3 May 1999 Oklahoma City tornado, *Weather and Forecasting*, 17, 456-  
339 471.
- 340 Church, C.R., Snow, J. T., Baker, G.L., Agee, E.M. (1979). Characteristics of tornado-like  
341 vortices as a function of swirl ratio: a laboratory investigation. *Journal of Atmospheric*  
342 *Sciences*, 36, 1755-1776.
- 343 Church, C.R., Snow, J. T. (1993). *The tornado: its structure, dynamics, prediction, and hazards*,  
344 American Geophysical Union, Washington, DC.
- 345 Davies-Jones, R.P. (1973). The Dependence of Core Radius on Swirl Ratio in a Tornado  
346 Simulator. *Journal of Atmospheric Sciences*, 30, 1427–1430.
- 347 Dominguez, D., and R.P. Selvam (2017), Tornado width for computer modeling from Google  
348 Earth data and period of the vortex, *Proceedings: AEI Conference 2017*, ASCE, Oklahoma  
349 City, OK, April 11-13, 470-483.
- 350 Hangan, H., Kim, J.D. (2008). Swirl ratio effects on tornado vortices in relation to the Fujita  
351 scale. *Wind and Structures*, 11, 291-302.
- 352 Harlow, F.H., Welch, J.E. (1965). Numerical calculation for time-dependent viscous  
353 incompressible flow of fluid with free surface. *Physics of Fluids*, 8, 2182-2189.
- 354 Hirt C.W., B.D. Nichols, and N.C. Romero, (1975). SOLA-A Numerical Solution Algorithm for  
355 Transient Fluid Flows, Los Alamos Scientific Laboratory report LA-5852.
- 356 Jischke, M. C., M. Parang, (1979). Properties of simulated tornado-like vortices. *Journal of*  
357 *atmospheric sciences*, 31, 506-512
- 358 Kashefzadeh, M.H. (2018). Computer Modeling of close-to-ground tornado wind-fields for  
359 different tornado widths. Master's thesis, Department of Civil Engineering, University of  
360 Arkansas, USA.
- 361 Karstens, C.D., Gallus, W.A., Lee. B.D. (2013) Analysis of Tornado-Induced Tree Fall Using  
362 Aerial Photography from the Joplin, Missouri, and Tuscaloosa–Birmingham, Alabama,  
363 Tornadoes of 2011. *Journal of Applied Meteorology and Climatology*, 52, 1049-1068.
- 364 Kosiba, K., Wurman, J. (2010). Notes and correspondence, the Three-Dimensional  
365 Axisymmetric Wind Field Structure of the Spencer, South Dakota, 1998 Tornado. *Journal of*  
366 *Atmospheric Sciences*, 67, 3074-3084.
- 367 Kosiba, K., J. Wurman (2013). The Three-Dimensional Structure and Evolution of a Tornado  
368 Boundary Layer, *Journal of Weather and Forecasting*. 28, 1552-1561.
- 369 Kuai, L., Haan, F.L., Gallus, W.A., Sarkar, P. (2008). CFD simulations of the flow field of a  
370 laboratory-simulated tornado for parameter sensitivity studies and comparison with field  
371 measurements. *Wind and Structures*, 11, 75-96.
- 372 Lewellen, W.S., Lewellen, D.C., Sykes, R.I. (1997). Large-eddy simulation of a tornado's  
373 interaction with the surface. *Journal of the Atmospheric Sciences*, 54, 581-605.

374 Neale, A., Derome, D., Blocken, B., Carmeliet, J. (2006). CFD calculation of convective heat  
375 transfer coefficients and validation—Part 2: Turbulent flow. From the web.

376 Nolan, D.S. (2013). On the Use of Doppler Radar–Derived Wind Fields to Diagnose the  
377 Secondary Circulations of Tornadoes. *J. Atmos. Sci.*, 70, 1160–1171

378 Refan, M. (2014). Physical Simulation of Tornado-Like Vortices. PhD thesis, Department of  
379 Mechanical and Materials Engineering, The University of Western Ontario.

380 Refan, M., Hangan, H., Wurman, J., Kosiba, K. (2017). Doppler radar-derived wind field of five  
381 tornado events with application to engineering simulations, *Engineering Structures*, 148,  
382 509-521.

383 Selvam, R.P. (1992), Computation of Pressures on Texas Tech Building, *Journal of Wind*  
384 *Engineering and Industrial Aerodynamics*, 43, 1619-1627.

385 Tari, P.H., Gurka, R., Hangan, H. (2010). Experimental investigation of tornado-like vortex  
386 dynamics with swirl ratio: The mean and turbulent flow fields. *Journal of Wind Eng. Ind.*  
387 *Aerodyn.* 98, 936-944.

388 Wakimoto, R.M., N.T. Atkins, K.M. Butler, H.B. Bluestein, K. Thiem, J.C. Synder, J. Houser,  
389 K. Kosiba, J. Wurman (2016). Aerial Damage Survey of the 2013 El Reno Tornado  
390 Combined with Mobile Radar Data. *American Meteorological Society*, 144, 1749-1776.

391 Ward, N.B., (1972). The exploration of certain features of tornado dynamics using a laboratory  
392 model. *Journal of the Atmospheric Sciences*, 29, 1194-1204.

393 Wilson, T., Rotunno, R. (1986). Numerical simulation of a laminar end-wall vortex and  
394 boundary layer. *Physics of Fluids*, 29, 3993-4005.

395 Wurman, J., Gill, S. (2000). Finescale Radar Observations of the Dimmitt, Texas (2 June 1995),  
396 Tornado. *Monthly Weather Review*, 128, 2135-2164.

397 Wurman, J. and Alexander, C.R. (2005), The 30 May 1998 Spencer, South Dakota, storm,  
398 *Monthly Weather Review*, 133, 97-119.

399 Wurman, J., C. Alexander, P. Robinson, Y. Richardson (2007). Low-Level Winds in Tornadoes  
400 and Potential Catastrophic Tornado Impacts in Urban Areas. *American Meteorological*  
401 *Society*, 89, 87-90.

402 Wurman, J., Kosiba, K., Robinson, P. (2013). In situ, Doppler radar and video observations of  
403 the interior structure of a tornado and the wind-damage relationship. *Bulletin of the American*  
404 *Meteorological Society*, 94, 835-846.

405