

## Scintillation, luminescence and optical properties of Ce-Doped borosilicate glasses

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### ABSTRACT

The structure, ultraviolet (UV) transparency, luminescence, and scintillation of CeO<sub>2</sub>-doped (0.1 mol%) borosilicate glasses containing Sb<sub>2</sub>O<sub>3</sub>, SnO<sub>2</sub> or P<sub>2</sub>O<sub>5</sub> were investigated. Glasses were prepared by melt quenching in air and characterized by several spectroscopy techniques: Raman scattering, optical absorption, photoluminescence (PL), fluorescence lifetime, thermoluminescence, and radioluminescence. Raman scattering results presented three major groups of vibrational modes, within 230–850 cm<sup>-1</sup>, 850–1250 cm<sup>-1</sup>, and 1250–1600 cm<sup>-1</sup>. Luminescence was dominated by a broad band centered at around 360–380 nm, depending on the chemical composition of the glass. Photoluminescence presented a dominant (96–98%) fast decay time of 36–40 ns that was ascribed to Ce<sup>3+</sup>. The P<sub>2</sub>O<sub>5</sub>-containing glass showed the highest scintillation output under X-ray excitation.

### 1. Introduction

Scintillators are materials capable of emitting luminescence due to exposure to ionizing radiation, and find application in several fields, e.g., medical imaging, homeland security, high energy physics, and oil exploration [1–7]. Scintillators are fabricated in many forms, including single crystals, polycrystalline ceramics, and glasses. Glass scintillators are particularly attractive due to advantages over single crystals and transparent polycrystalline ceramics such as low-cost, fast, and large area/volume production, and nearly limitless options for chemical composition. Rare earth ions are normally adopted as luminescence activators in scintillating glasses, with Ce<sup>3+</sup> being the most common because of its fast allowed 5d → 4f transition. However, the luminescence of Ce<sup>3+</sup> is strongly influenced by the chemical composition and atomic arrangement of the nearest neighbors [8–12]. Therefore, investigation of the luminescence of Ce<sup>3+</sup> in different glasses is an active research topic [13–20]. In this work, an exploratory investigation on the effects of the incorporation of Sb<sub>2</sub>O<sub>3</sub>, SnO<sub>2</sub> and P<sub>2</sub>O<sub>5</sub> in Ce-doped borosilicate glasses was executed, focusing on the structure, optical properties, luminescence, and scintillation.

### 1. Experimental procedures

The glass compositions were based on the SiO<sub>2</sub>-B<sub>2</sub>O<sub>3</sub>-Na<sub>2</sub>O-MgO-Ce<sub>2</sub>O<sub>3</sub>-Sb<sub>2</sub>O<sub>3</sub>/SnO<sub>2</sub>/P<sub>2</sub>O<sub>5</sub> system and designed using the proprietary software Reformix 2.0 (LaMaV-UFSCar, Brazil) and the Sciglass® database (ITC Inc.). The following analytical grade oxides and carbonates were used as precursors: high-purity quartz powder (99.99%, SiO<sub>2</sub>, Vitróvita, Brazil), B<sub>2</sub>O<sub>3</sub> (ESPI Metals, USA), Na<sub>2</sub>CO<sub>3</sub> (99.5%, Vetec, Brazil), MgO (96%, Synth, Brazil), CeO<sub>2</sub> (99%, Vetec, Brazil), Sb<sub>2</sub>O<sub>3</sub> (Fluka, Germany), SnO<sub>2</sub> (99.9% Riedel-de Haën, Germany), P<sub>2</sub>O<sub>5</sub> (98%, Sigma Aldrich, UK). The nominal composition of the glasses is given in Table 1. The proper mixtures of these reagents were homogenized in a planetary ball mill (Fritsch Pulverisette 5) at 300 rpm for 12 min using a silicon nitride jar and eight 20 mm-diameter balls. The glasses were melted in an electrical furnace at atmospheric pressure at 1140 °C for 4 h in a platinum crucible. The melts were crushed, ball-milled using the same experimental conditions as before, and re-melted. They were remelted over a total of two re-melting procedures to ensure homogeneity. Finally, the melts were poured on a stainless steel plate, pressed with a stainless steel plaque and annealed at 520 °C for 2 h. Since glass

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fabrication was executed without atmospheric control, partial conversion of  $\text{Ce}^{3+}$  to  $\text{Ce}^{4+}$  was expected to occur (see, e.g., Refs. [21,22]). The surface of both sides of glass slabs were polished using a diamond solution for optical absorption measurements.

Specific density measurements were executed through the application of Archimedes' principle. The weight of each sample was determined both immersed in air ( $W_{\text{air}}$ ) and in water ( $W_{\text{water}}$ ), and the specific density calculated according to:  $\rho = W_{\text{air}}/(W_{\text{air}} - W_{\text{water}})$  taking the density of water to be  $1 \text{ g/cm}^3$ . Three sets of measurements were done for each sample yielding a typical variation of the mass values with  $1\%$  standard deviation. The index of refraction was measured using a Carl Zeiss Jena refractometer Pulfrich-refraktometer PR2 at  $546.1 \text{ nm}$  after thorough polishing and cleaning of the glasses. The index of refraction was also estimated using the Gladstone-Dale relation [23,24] with the experimentally-determined density values as input.

Raman scattering spectroscopy was executed with a Horiba LabRAM HR Evolution Raman microscope equipped with an  $800 \text{ mm}$  focal length spectrograph and a cooled ( $-60 \text{ }^\circ\text{C}$ ) back-illuminated deep-depleted  $1024 \times 256$  pixels CCD detector using a  $100 \text{ mW}$   $532 \text{ nm}$  laser with no attenuation, a  $5\times$  magnification objective, and  $1800 \text{ grooves/mm}$  diffraction gratings. Spectra corresponded to the accumulation of five  $15 \text{ s}$  long scans. Each spectrum was corrected by the pre-recorded instrument-specific response to a calibrated white light source, namely the intensity correction system (ICS).

Optical absorption spectra were collected using a PerkinElmer Lambda 900 UV/Vis/NIR spectrometer.

Photoluminescence emission (PL) and excitation (PLE) measurements were executed using an Edinburgh Instruments FLS-1000 spectrophotometer equipped with a  $450 \text{ W}$  Xe lamp and dual monochromators for emission and excitation. Excitation spectra were obtained within  $240\text{--}360 \text{ nm}$  and emission within  $280\text{--}340\text{--}650 \text{ nm}$ , depending on the sample, with excitation and emission bandwidths fixed at  $0.4 \text{ nm}$ . The number of detected photons was automatically corrected by the instrument's software for the background and spectral differences in excitation intensity and detection sensitivity.

Fluorescence lifetime was executed with an Edinburgh Instruments FLS-1000 spectrophotometer using picosecond pulsed LEDs EPLED-310 emitting at  $317.0 \text{ nm}$  ( $9.6 \text{ nm}$  FWHM bandwidth,  $915.6 \text{ ps}$  pulse width at  $10 \text{ MHz}$ ; glass #3) and EPLED-330 emitting at  $333.8 \text{ nm}$  ( $9.3 \text{ nm}$  FWHM bandwidth,  $789.3 \text{ ps}$  pulse width at  $10 \text{ MHz}$ ; glasses # 1 and 2), depending on the position of the excitation band of the glass. A  $2 \text{ nm}$  bandwidth was used for detection with the monitoring wavelength being within  $368\text{--}382 \text{ nm}$ , while EPLED repetition rate was fixed at  $0.5 \text{ MHz}$ .

Radioluminescence (RL) measurements were executed using a customer-designed Freiberg Instruments Lexsyg Research spectrophotometer equipped with a Varian Medical Systems VF-50J X-ray tube with a tungsten target coupled with an ionization chamber for irradiation dose monitoring. The light emitted by the sample was collected by a lens and converged into an optical fiber connected to an Andor Technology Shamrock 163 spectrograph coupled to an Andor Technology DU920P-BU Newton CCD camera. RL measurements were executed under continuous X-ray irradiation ( $40 \text{ kV}$ ,  $1 \text{ mA}$ ) at room temperature (RT). Powders filled *ca.*  $8 \text{ mm}$  diameter  $1 \text{ mm}$  deep cups thus allowing for relative RL intensity comparison between different samples. The reproducibility of RL integrated intensities was within  $8\%$ . Results were not corrected for the spectral sensitivity of the system.

Thermoluminescence (TL) measurements were executed with a Freiberg Instruments Lexsyg Research spectrophotometer equipped with a Hamamatsu H7360-2 photomultiplier tube using the following procedure: 1) TL readout up to  $450 \text{ }^\circ\text{C}$  with a  $1 \text{ }^\circ\text{C/s}$  heating rate and a  $1 \text{ s}$  integration time followed by cooling down to RT, 2) heating up to  $450 \text{ }^\circ\text{C}$  and keeping the sample at  $450 \text{ }^\circ\text{C}$  for  $300 \text{ s}$  to clean all traps, followed to cooling down to RT, 3) TL readout like in step 1, 4) X-ray irradiation ( $40 \text{ kV}$ ,  $1 \text{ mA}$ ) for  $100 \text{ s}$  at RT, and 5) TL readout like in step 1. TL results presented in this work correspond to the last TL readout normalized by the total mass, *i.e.*, powder and cup, that was within  $401.9$  and  $407.5 \text{ mg}$  (for reference, the average of the mass of ten different cups was  $375.5 \text{ mg}$ ).

## 2. Results and discussion

The Raman spectra of the glasses were similar and presented three groups of bands within about  $230\text{--}850 \text{ cm}^{-1}$ ,  $850\text{--}1250 \text{ cm}^{-1}$ , and  $1250\text{--}1600 \text{ cm}^{-1}$  (Fig. 1). These bands are originated from the vibrations of the silicate and borate structural units. The small amount of  $\text{P}_2\text{O}_5$  in glass #3 is not expected to yield active vibrational modes. It is known that the incorporation of alkali oxides in borate glasses alters the atomic network by disrupting boroxol rings and causing the appearance of a Raman peak at  $765\text{--}775 \text{ cm}^{-1}$  at the expense of the  $805\text{--}810 \text{ cm}^{-1}$  peak [25]. In fact, in both binary alkali borate and ternary alkali borosilicate glasses, the  $805\text{--}810 \text{ cm}^{-1}$  band was reported to disappear for high alkali

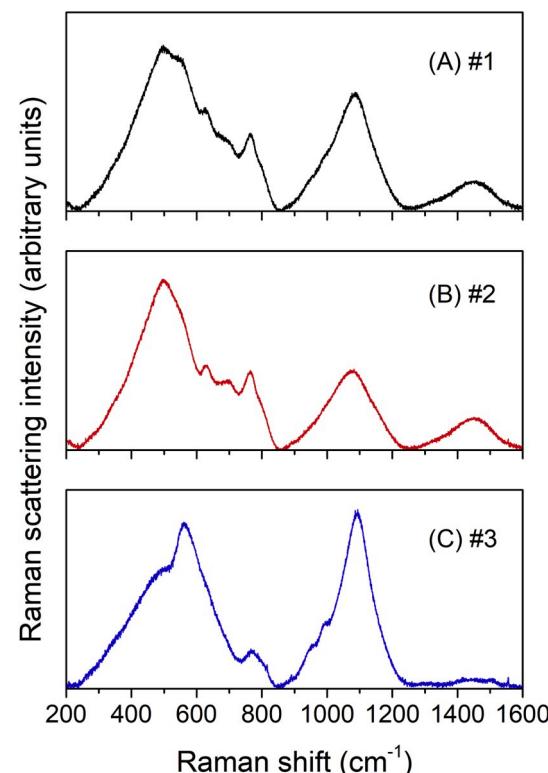


Fig. 1. Raman spectra of: A) glass #1, B) glass #2, and C) glass #3.

Table 1

Glass label, chemical composition, density, experimental index of refraction  $n$ , and calculated index of refraction using the Gladstone-Dale relation,  $n_{\text{GD}}$ .

Glass #	Molar fraction (%)							Density ( $\text{g/cm}^3$ )	$n$	$n_{\text{GD}}$
	$\text{SiO}_2$	$\text{B}_2\text{O}_3$	$\text{Na}_2\text{O}$	$\text{MgO}$	$\text{CeO}_2$	$\text{Sb}_2\text{O}_3$	$\text{SnO}_2$			
1	47	28	14	10.8	0.1	0.1	-	-	2.6	1.528
2	47	28	14	10.8	0.1	-	0.1	-	2.4	1.525
3	47	28	14	10.8	0.1	-	-	0.1	2.4	1.526

concentrations [26–28]. In binary alkali borate glasses, it was shown that the incorporation of low concentrations of an alkali metal occurs by inducing the conversion of the coordination of boron from three to four, *i.e.*, the formation of  $\text{BO}_4^-$  units with an adjacent alkali ion for charge compensation. For higher alkali concentrations, charge compensation is achieved through the formation of non-bridging oxygen (NBO) adjacent to the alkali ion [27]. In ternary sodium borosilicate glasses with  $\text{Na}_2\text{O}/\text{B}_2\text{O}_3$  ratios less than 1 (in this work,  $\text{Na}_2\text{O}/\text{B}_2\text{O}_3 = 0.5$ ), it was shown that the majority of  $\text{Na}^+$  ions are bonded to borate units, *i.e.*, they cause conversion of boron from coordination three to four while causing relatively few  $\text{Si-O}^-$  non-bridging units. In ternary sodium borosilicate glasses, the Raman band within  $450\text{--}550\text{ cm}^{-1}$  was ascribed to bending or rocking of  $\text{Si-O-Si}$  and  $\text{Si-O-B}$  bonds composing the three-dimensional network, *i.e.*, the incorporation of borate groups,  $\text{BO}_4$  and  $\text{BO}_3$  units into the silica network, the band within  $620\text{--}635\text{ cm}^{-1}$  to a vibration of the metaborate ring, the band within  $765\text{--}775\text{ cm}^{-1}$  to vibrations of six-membered borate rings with one or two  $\text{BO}_4$  units, the band within  $805\text{--}810\text{ cm}^{-1}$  to a symmetric vibration of the boroxol ring, the band within  $950\text{--}975\text{ cm}^{-1}$  to the symmetric stretching of the non-bridging bonds of  $\text{SiO}_3$  chain units (*i.e.*,  $\text{Q}^2$  units), the band within  $1050\text{--}1100\text{ cm}^{-1}$  to the stretching of the non-bridging bond of  $\text{SiO}_{5/2}$  sheet units (*i.e.*,  $\text{Q}^3$  units), and the band around  $1460\text{--}1470\text{ cm}^{-1}$  to the stretching mode of  $\text{B-O}$  bonds [26 and references therein]. Also, Raman investigation of  $\text{xMgO}\text{-yNa}_2\text{O}\text{-B}_2\text{O}_3$  glasses with  $\text{x} + \text{y} = 1$  ascribed the band at  $570\text{ cm}^{-1}$  to “loose” diborate units, at  $760\text{ cm}^{-1}$  to six-membered borate rings with two  $\text{BO}_4$  units, at  $840\text{ cm}^{-1}$  to pyroborate units, and at  $940\text{ cm}^{-1}$  to orthoborate units [29]. Accordingly, in this work, the band at about  $497\text{ cm}^{-1}$  was ascribed to bending or rocking of  $\text{Si-O-Si}$  and  $\text{Si-O-B}$  bonds, at around  $550\text{ cm}^{-1}$  to “loose” diborate units, at around  $630\text{ cm}^{-1}$  to a vibration of the metaborate ring, at around  $766\text{ cm}^{-1}$  to vibrations of six-membered borate rings with one or two  $\text{BO}_4$  units, and the shoulder around  $800\text{ cm}^{-1}$  to a symmetric vibration of the boroxol ring. The band at  $975\text{ cm}^{-1}$  was ascribed to  $\text{Q}^2$  units, the band at about  $1083\text{ cm}^{-1}$  to  $\text{Q}^3$  units, and the band at about  $1452\text{ cm}^{-1}$  to the stretching mode of  $\text{B-O}$  bonds. The nature of a small band around  $700\text{ cm}^{-1}$  could not be identified. The Raman spectra of glasses #1 and 2 were similar, while bands at  $630$  and  $700\text{ cm}^{-1}$  were missing in the spectrum of glass #3. Also, the relative contribution of the band at about  $1083\text{ cm}^{-1}$  in the spectrum of glass #3 was much higher than in the other glasses. These differences corresponded to the absence of metaborate rings and a higher content of  $\text{Q}^3$  units in glass #3.

The density of the glasses was determined to be between 2.4 and 2.6  $\text{g/cm}^3$  (cf. Table 1) in agreement with the density values of other similar borosilicate glasses [15,30], and matching the density of commercial Ce-doped glass scintillators [31].

Table 1 shows the experimental and calculated values of the index of refraction. They were within a narrow range of values due to the small variation of density and chemical composition of the glasses, and in agreement with the measured values of the index of refraction of similar glasses [15,32]. While the Gladstone-Dale equation cannot account for dispersion, these results attested for its accuracy in estimating the index of refraction within the visible spectral range, where the relative error to the experimental values was inferior to 2%.

The optical absorption results are shown in Fig. 2 showing the glasses to be transparent up to  $\sim 315\text{ nm}$ , with glass #3 being the most ultraviolet transparent. The spectra presented a strong absorption band centered around  $\sim 302\text{ nm}$  ascribed to the  $4f \rightarrow 5d$  transition of  $\text{Ce}^{3+}$  ions, in agreement with PLE results. The PLE and PL spectra are presented in Fig. 3. The PLE spectra were peaked around  $320\text{ nm}$ , while displaying an asymmetric tail towards shorter wavelengths that was suggestive of the presence of additional excitation bands. The PLE spectrum of glass #3 presented another strong band centered at  $270\text{ nm}$ . All PL spectra presented a single band centered within  $360\text{--}370\text{ nm}$ , and excitation of glass #3 at  $270$  and  $317\text{ nm}$  yielded the same emission spectrum. Glass #3 showed a considerably less self-absorption than the other glasses, as can be inferred by the small superposition between the

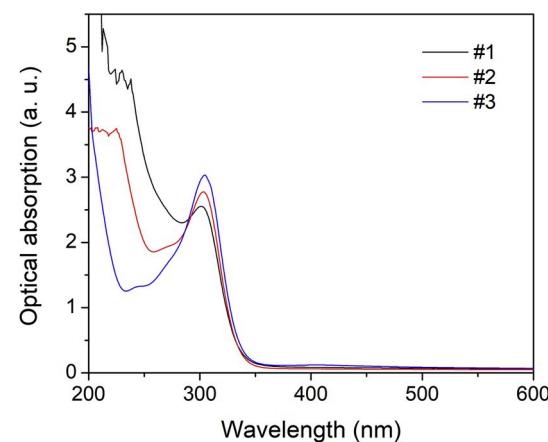


Fig. 2. Optical absorption of Ce-doped borosilicate glasses.

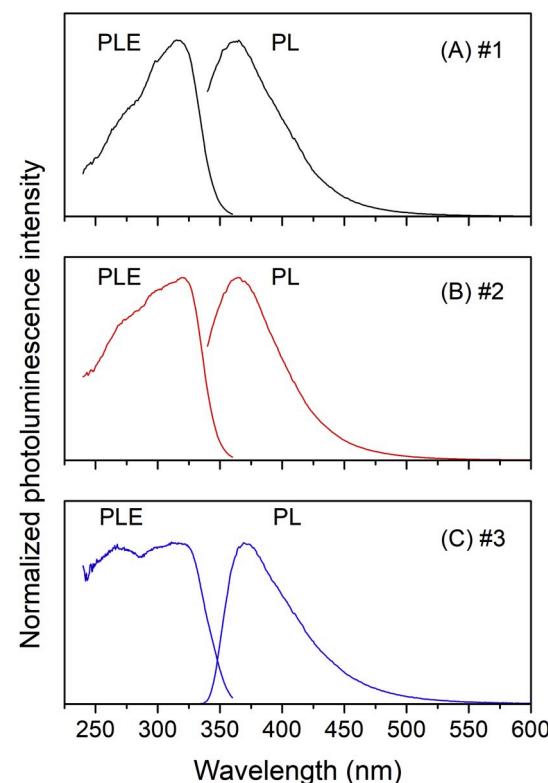


Fig. 3. Normalized PLE and PL spectra of A) glass #1: PLE monitored at  $380\text{ nm}$ , and PL excited at  $317\text{ nm}$ , B) glass #2: PLE monitored at  $382\text{ nm}$ , and PL excited at  $320\text{ nm}$ , and C) glass #3: PLE monitored at  $368\text{ nm}$ , and PL excited at  $262\text{ nm}$ .

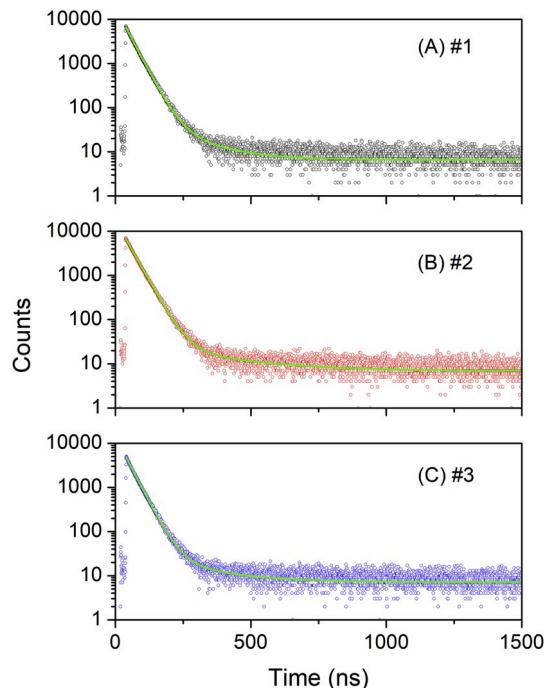
excitation and emission spectra. The emission of the glasses reported in this work was in agreement with that of Ce-doped crystalline borosilicates  $\text{CaB}_2\text{O}_5(\text{Si}_2\text{O}_7)$  at  $343\text{ nm}$ , and  $\text{LaBO}(\text{SiO}_4)$  at  $377\text{ nm}$  [8 and references therein]. Moreover, Table 2 summarizes a literature review on the luminescence of Ce-doped borosilicate glasses where it is possible to check that the peak position of the excitation and emission bands found in this work are in good agreement with the reported results regardless the fabrication conditions [15,16,33–36].

Fluorescence lifetime results (circles) are presented in Fig. 4, together with best-fitting using two exponential functions (continuous lines). The decays were dominated (96–98%) by the fast lifetime component that ranged from 36 to 40 ns, in addition to the contribution of a longer lifetime component amounting to 145 ns for glass #1, 259 ns

**Table 2**

Summary of literature review on Ce-doped borosilicate glasses: glass composition, fabrication details, excitation and emission peak position, lifetime, and respective reference. The glass compositions presented here are indicative of the compositional range used in the original work. See references for details. RT = room temperature; NA = not available.

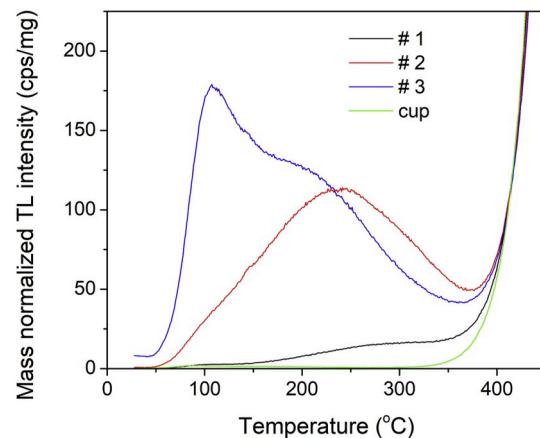
Glass composition	Fabrication details	Excitation Peak (nm)	Emission Peak (nm)	Lifetime (ns)	Ref.
70B <sub>2</sub> O <sub>3</sub> ·(8–14.5)SiO <sub>2</sub> ·(0–20.8)Li <sub>2</sub> O·1Al <sub>2</sub> O <sub>3</sub> ·(0–20)Na <sub>2</sub> O·(0.2–5)Ce <sub>2</sub> O <sub>3</sub>	Melting at 1000 °C in air Poured on graphite at RT	NA	341–353 372–432	NA	15
20SiO <sub>2</sub> ·(28–38)B <sub>2</sub> O <sub>3</sub> ·15BaF <sub>2</sub> ·(10–20)Lu <sub>2</sub> O <sub>3</sub> ·(10–20)Gd <sub>2</sub> O <sub>3</sub> ·2CeF <sub>3</sub>	Melting at 1400 °C in 95/5 N <sub>2</sub> :H <sub>2</sub> Poured on heated stainless steel	ca. 260, 320	385	33–41	16
40Li <sub>2</sub> O·yB <sub>2</sub> O <sub>3</sub> ·(60–y)SiO <sub>2</sub> (0 ≤ y ≤ 40) + xCe in excess (0.1 ≤ x ≤ 2)	Melting at 1000 °C in Ar Poured on 200 °C stainless steel	ca. 323–345	ca. 363–394	37–45	33
15CaO·55B <sub>2</sub> O <sub>3</sub> ·30SiO <sub>2</sub> + 0.04 mol% CeO <sub>2</sub>	Melting at 1500 °C in air Poured on 400 °C stainless steel	319	392	41	34
15SiO <sub>2</sub> ·30B <sub>2</sub> O <sub>3</sub> ·25Al <sub>2</sub> O <sub>3</sub> ·30Gd <sub>2</sub> O <sub>3</sub> + 1 wt% CeO <sub>2</sub>	Melting at 1580 °C in CO Poured on heated stainless steel	ca. 350	410	40 (90%) + 136 (10%)	35
30SiO <sub>2</sub> ·25B <sub>2</sub> O <sub>3</sub> ·30Gd <sub>2</sub> O <sub>3</sub> ·15BaO + 5 wt% CeO <sub>2</sub>	Melting at 1350 °C in N <sub>2</sub> Poured at 1200 °C	360	450	21–31 + 64–74	36



**Fig. 4.** Fluorescence decay of Ce-doped borosilicate glasses (circles) together with best-fit using two exponential functions (continuous lines) of: A) glass #1, B) glass #2, and C) glass #3.

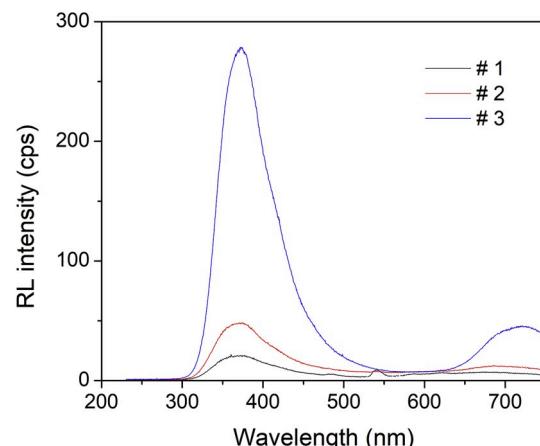
for glass #2, and 181 ns for glass #3. The fast decay was ascribed to Ce<sup>3+</sup> ions, and while the origin of the long lifetime was not known, it was believed to be originated from electrons that are thermally released from shallow traps and radiatively recombine at Ce<sup>3+</sup> sites, in agreement with a small contribution of a few % to the temporal decay and strong TL signal at relatively low temperatures, as discussed below. These results were consistent with previous reports on Ce-doped phosphate and borosilicate glasses [33,35,37].

TL results of the glass powders are presented in Fig. 5 together with the TL response of an empty cup. Overall, the TL response of the glasses corresponded to broad glow peaks in agreement with the disordered nature of glass structures, with a significant contribution at relatively low temperatures. The steep increase for temperatures above about 350 °C corresponded to the blackbody response of the system, as demonstrated by the measurement of the empty cup. These data were analyzed in terms of the integral TL signal, with glass #3 presenting an integrated TL response 12x and glass #2 8x that of glass #1.



**Fig. 5.** TL glow curves of Ce-doped borosilicate glasses, together with the TL response of an empty cup.

RL results are presented in Fig. 6. The spectra are dominated by a broad band centered at 372 nm for glass #1, 369 nm for glass #2, and 373 nm for glass #3, depending on the glass composition, in agreement with PL results. The weaker bands around 700 nm and within ca. 470–640 nm were ascribed to Fe<sup>3+</sup> impurity, in accord with Fe<sup>3+</sup> emission band in quartz at 709 nm [38], and to Eu<sup>3+</sup>, respectively. These impurities could not be detected by energy dispersive X-ray spectroscopy (EDX) measurements, showing them to be trace contamination.



**Fig. 6.** RL spectra of Ce-doped borosilicate glasses.

Their small concentration was not expected to cause any significant effect on the luminescence of  $\text{Ce}^{3+}$ . The dependence of  $\text{Ce}^{3+}$  luminescence on the host is due to the lack of shielding of the  $5d$  orbital to outer electrons making the  $5d \rightarrow 4f$  transition sensitive to the physical arrangement and chemical nature of the surrounding atoms [8–12]. The RL spectra were integrated within 275–470 nm, yielding the relative luminosity of the glasses. Glass #3 was 5.2x brighter than glass #2 and 11.6x brighter than glass #1, while glass #2 was 2.2x brighter than glass #1, in agreement with the optical absorption results and reduced self-absorption presented by glass #3.

### 3. Conclusions

We investigated Ce-doped borosilicate glasses which yielded insights into the effects of the addition of  $\text{Sb}_2\text{O}_3$ ,  $\text{SnO}_2$  and  $\text{P}_2\text{O}_5$  on their optical and luminescent properties. The incorporation of  $\text{P}_2\text{O}_5$  was advantageous over  $\text{Sb}_2\text{O}_3$  and  $\text{SnO}_2$ , leading to superior ultraviolet transparency, reduced self-absorption of  $\text{Ce}^{3+}$  emission, and brighter scintillation under X-ray irradiation.

### Author agreement

This document serves to confirm that the work described in the manuscript has not been published before, it is not under consideration for publication anywhere else, and its publication has been approved by all co-authors and the responsible authorities at the institutions where the work was carried out.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

### CRediT authorship contribution statement

**Linyu Pan:** Formal analysis, Investigation, Writing - review & editing. **Juliana K.M.F. Daguano:** Investigation, Resources, Writing - review & editing. **Neilo M. Trindade:** Investigation, Writing - review & editing. **Mayara Cerruti:** Investigation. **Edgar D. Zanotto:** Conceptualization, Methodology, Resources, Writing - review & editing, Project administration, Funding acquisition. **Luiz G. Jacobsohn:** Conceptualization, Methodology, Formal analysis, Investigation, Resources, Writing - original draft, Writing - review & editing, Project administration, Funding acquisition.

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