

1 **Kinematic Structure and Dynamics of the Denmark Strait Overflow**
2 **from Ship-based Observations**

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May 2020

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Submitted to *Journal of Physical Oceanography*

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20 **Abstract**

21 The dense outflow through Denmark Strait is the largest contributor to the lower limb of the
22 Atlantic Meridional Overturning Circulation, yet a description of the full velocity field across the
23 strait remains incomplete. Here we analyze a set of 22 shipboard hydrographic/velocity sections
24 occupied along the Látrabjarg transect at the Denmark Strait sill, obtained over the time period
25 1993–2018. The sections provide the first complete view of the kinematic components at the sill:
26 the shelfbreak East Greenland Current (EGC), the combined flow of the separated EGC and the
27 North Icelandic Jet (NIJ), and the northward flowing North Icelandic Irminger Current (NIIC).
28 The total mean transport of overflow water is 3.54 ± 0.29 Sv, comparable to previous estimates.
29 The dense overflow is partitioned in terms of water mass constituents and flow components. The
30 mean transports of the two types of overflow water – Atlantic-origin Overflow Water and Arctic-
31 origin Overflow Water – are comparable in Denmark Strait, while the merged NIJ/separated EGC
32 transports 55% more water than the shelfbreak EGC. A significant degree of water mass exchange
33 takes place between the branches as they converge in Denmark Strait. There are two dominant
34 time-varying configurations of the flow that are characterized as a cyclonic state and a non-
35 cyclonic state. These appear to be wind-driven. A potential vorticity analysis indicates that the
36 flow through Denmark Strait is subject to symmetric instability. This occurs at the top of the
37 overflow layer, implying that the mixing/entrainment process that modifies the overflow water
38 begins at the sill.

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43 1. Introduction

44 The dense water formed in the Nordic Seas is the main source of lower North Atlantic Deep Water
45 that plays an essential role in the Atlantic Meridional Overturning Circulation (AMOC) (Dickson
46 and Brown, 1994). Studies have now demonstrated that the dominant contribution to the AMOC
47 is associated with the warm-to-cold transformation that occurs in the Nordics Seas as opposed to
48 that which takes place in the Labrador Sea (Pickart and Spall, 2007; Holte and Straneo, 2017;
49 Lozier et al., 2019). Denmark Strait is one of the key passages through which the dense water from
50 the Nordic domain enters the North Atlantic Ocean. The so-called Denmark Strait Overflow Water
51 (DSOW) accounts for roughly half of the total dense water flowing over the Greenland-Scotland
52 Ridge. The mean transport of DSOW at the sill, which is typically defined as water denser than
53 27.8 kg m^{-3} , is estimated to be 3.2–3.5 Sv (Harden et al., 2016; Jochumsen et al., 2017).

54 There are three different pathways that advect the dense water into Denmark Strait from the north,
55 supplying the overflow water (Fig. 1): the shelfbreak East Greenland Current (EGC); the separated
56 EGC; and the North Icelandic Jet (NIJ). The EGC emanates from Fram Strait and is a surface-
57 intensified flow transporting Atlantic-origin Overflow Water at depth. At these latitudes it is
58 comprised of a shelfbreak branch and an offshore slope branch (Håvik et al., 2017a). Together
59 they advect a combination of warm, salty water that has been modified along the rim current
60 system of the Nordic Seas (Mauritzen, 1996) and also in the high Arctic (Rudels et al., 2005). The
61 shelfbreak EGC transport decreases as it progresses southward, while the slope branch appears to
62 be diverted eastward into the interior north of the Iceland Sea (Håvik et al., 2017a). When the
63 shelfbreak EGC reaches the northern end of the Blosseville Basin it bifurcates to form the
64 separated EGC (Fig. 1). Våge et al. (2013) attribute the bifurcation to local wind stress curl and
65 topography, as well as baroclinic instability of the shelfbreak current.

66 The NIJ is a middepth-intensified current on the north Iceland slope that transports Arctic-origin
67 Overflow Water equatorward. This water is colder, fresher, and denser than the Atlantic-origin
68 Overflow Water. It was hypothesized by Våge et al. (2011) that the NIJ is part of a local
69 overturning loop in the Iceland Sea whereby the subtropical-origin water transported northward
70 by the North Icelandic Irminger Current (NIIC, Fig. 1) is fluxed into the interior of the basin and
71 converted to overflow water by wintertime air-sea heat loss. The dense water then progresses back
72 to the Iceland slope where it sinks and feeds the NIJ. However, it has since been demonstrated that
73 the bulk of the Arctic-origin water must originate from farther north where the wintertime mixed
74 layers are denser (Våge et al., 2015). Recent analysis of the historical data suggests that the water
75 stems from the Greenland Sea (Huang et al., submitted). Based on a large collection of shipboard
76 transects occupied over 15 years, Semper et al. (2019) documented that the NIJ steadily increases
77 in transport as it flows towards Denmark Strait, accounting for a sizable fraction of the dense water
78 that overflows the sill.

79 Using data from a year-long mooring array across the Blossville Basin (roughly 200 km north of
80 the sill), Harden et al. (2016) calculated mean transports for the three individual pathways: $1.50 \pm$
81 0.16 Sv for the shelfbreak EGC, 1.04 ± 0.15 Sv for the separated EGC, and 1.00 ± 0.17 Sv for the
82 NIJ. There was very little seasonal variation, in line with the weak seasonality observed at the sill
83 (Jochumsen et al., 2012). However, Harden et al. (2016) revealed that the transports of the three
84 branches vary on intra-seasonal timescales, and that they tend to compensate each other such that
85 the total overflow transport remains fairly steady. This is suggestive of hydraulic control, which is
86 argued to be active in Denmark Strait (Whitehead, 1989; Nikolopoulos et al., 2003). Harden et al.
87 (2016) argued that wind stress curl forcing causes the compensation between the NIJ and the two
88 EGC branches.

89 On synoptic timescales, the flow of DSOW is highly energetic (Smith, 1976; Bruce, 1995; Rudels
90 et al., 1999; Girton and Sanford, 2003; Käse et al., 2003; von Appen et al., 2017). Using the
91 mooring data from the afore-mentioned Blossville Basin array, Huang et al. (2019) demonstrated
92 that high-frequency variability is driven by mean-to-eddy baroclinic conversion at the shoreward
93 edge of the NIJ. Using a year-long mooring array in Denmark Strait, Moritz et al. (2019) resolved
94 the passage of eddies, finding more anticyclones in the deepest part of the strait and more cyclones
95 west of this. Satellite altimetry data have revealed enhanced levels of surface eddy kinetic energy
96 in the vicinity of the strait (Høyer and Quadfasel, 2001; Håvik et al., 2017b).

97 A series of recent papers have further characterized the high-frequency variability of the DSOW
98 at the sill. Two dominant features have been identified, referred to as boluses and pulses. The
99 former corresponds to the passage of a large lens of overflow water and is associated with cyclonic
100 circulation (von Appen et al., 2017). Mastropole et al. (2017) identified boluses in 46 out of 111
101 transects across the strait occupied since 1990. These features export the very densest DSOW.
102 Pulses correspond to a thinning and acceleration of the DSOW layer, and are associated with anti-
103 cyclonic circulation (von Appen et al., 2017). The two types of features have been identified in a
104 high-resolution numerical model, with characteristics similar to the observations (Almansi et al.,
105 2017). Both the boluses and pulses result in increased transport of DSOW over a period of several
106 days (von Appen et al., 2017). Almansi et al. (2020) have shown that the surges in transport result
107 in the generation of cyclones downstream of the sill. These are the well-known “DSOW cyclones”
108 that emanate from the strait and propagate southward along the East Greenland continental slope
109 (Bruce, 1995; Spall and Price, 1998; von Appen et al., 2017).

110 The numerical study of Spall et al. (2019) determined that boluses and pulses are part of a single
111 dynamical process, associated with baroclinic instability of the hydrographic front in Denmark
112 Strait. This front divides the southward-flowing water emanating from the Nordic Seas and the
113 northward-flowing NIIC. The instability process results in frontal meanders that propagate
114 southwestward through the strait. Meander troughs are associated with boluses, whereby the NIIC
115 shifts towards Iceland and more overflow water is present in the center of the strait. Meander crests
116 are associated with pulses, when the NIIC moves farther into the strait towards Greenland. Spall
117 et al. (2019) demonstrated that this process is dictated by the interplay between the confluent mean
118 flow in the strait that tends to sharpen the front, and the baroclinic instability which works to relax
119 the front. These results show that the dynamics of the DSOW are closely tied to those of the NIIC.

120 Based on the large number of shipboard occupations of the Látrabjarg transect across Denmark
121 Strait (Fig. 1), we now have a good characterization of the two-dimensional hydrographic structure
122 across the strait. However, we lack an analogous view of the kinematic structure. Over the years,
123 moorings have been deployed in the deepest part of the sill, which is referred to as the trough (see
124 Fig. 4). These timeseries have provided information on the vertical structure and transport of the
125 overflow water (e.g. Jochumsen et al., 2017; von Appen et al., 2017; Spall et al., 2019). Recently
126 a five-mooring array was deployed on the western flank of the trough. Using empirical orthogonal
127 function analysis on the velocity timeseries, Jochumsen et al. (2017) found that the first mode
128 reflects a barotropic flow that pulses in time, the second mode represents lateral shifts of the flow,
129 and the third mode corresponds to the mesoscale eddy features noted above, investigated by Moritz
130 et al. (2019). While these measurements have enhanced our understanding of the flow components
131 in Denmark Strait, they are limited in cross-strait coverage and only have near-bottom temperature
132 and salinity information.

133 In this study we analyze the updated collection of shipboard occupations of the Látrabjarg line
134 (Fig. 1). The number of occupations is now 122, and, importantly, 22 of them contain direct
135 velocity measurements. This provides the first-ever robust view of the two-dimensional velocity
136 structure across the strait. It enables us to determine the fate of the three above-mentioned
137 pathways of overflow water into Denmark Strait, including the water masses they advect and their
138 relationship to the NIIC. We are also able to investigate dynamical aspects of the overflow. The
139 paper is organized as follows. We begin with a presentation of the data. This is followed by a
140 description of the mean hydrographic and velocity structure and the partitioning of the overflow
141 transport by water masses and currents. We then characterize the dominant mode of variability and
142 its relationship to local wind forcing. Finally, we address the hydraulic criticality of the overflow,
143 including the occurrence of symmetric instability and implications for mixing.

144 **2. Data and Methods**

145 **2.1. Látrabjarg sections**

146 We use 122 occupations of the Látrabjarg hydrographic transect across Denmark Strait taken
147 between 1990 and 2018 (Fig. 2). This is an updated version of the dataset used by Mastropole et
148 al. (2017), who analyzed 111 of the sections (1990–2012; see Table 1 in Mastropole et al., 2017).
149 A detailed description of the conductivity-temperature-depth (CTD) data and the processing steps
150 are described in Mastropole et al. (2017). Briefly, each occupation is projected onto the standard
151 Látrabjarg line (black line in Fig. 2), and gridded vertical sections of the hydrographic variables
152 are constructed. We followed the same procedure for the 11 additional occupations, which are
153 listed in Table 1. Here we also use direct velocity information obtained on 22 of the sections (Table
154 1; Fig. 3). This consisted of vessel-mounted acoustic Doppler current profiler (ADCP) data (15 of

155 the occupations) and lowered ADCP data (7 of the occupations). Our study focuses primarily on
156 the 22 occupations with velocity data, except for Section 5 where the full historical hydrographic
157 dataset is used.

158 Absolute geostrophic velocity sections were computed using the gridded hydrographic sections in
159 conjunction with gridded sections of the cross-track ADCP velocities, following the same
160 procedure as in Pickart et al. (2016). Errors in the volume transport estimates are associated with
161 instrument uncertainty, the gridding process, and the inability to measure the flow in the “bottom
162 triangles” (the area beneath the deepest common level of adjacent stations). Because of the
163 generally small station spacing of the sections, the latter effect is taken to be negligible. The
164 uncertainties of both the vessel-mounted ADCP and lowered ADCP are taken as 0.02 m s^{-1} (Pickart
165 et al., 2016, 2017). The gridding error was obtained by calculating the differences between the
166 vertically-averaged velocity measurement at each station versus the same quantity determined
167 using the gridded values closest to the station (Nikolopoulos et al., 2009). The final error is taken
168 to be the root of the sum of the squares of the instrument and gridding errors, and is applied over
169 the area of the section where the transport is being calculated. Since this does not assume that the
170 errors are uncorrelated, it represents a conservative estimate.

171 **2.2. Reanalysis Wind Data**

172 Wind fields from the ERA5 reanalysis were obtained from the European Center for Medium-
173 Range Weather Forecasts (ECMWF, <https://www.ecmwf.int/>). This is the 5th generation
174 reanalysis, which uses ECMWF’s Integrated Forecast System (IFS). Previous studies have shown
175 good agreement between IFS products and observations (Harden et al., 2016). The spatial
176 resolution of ERA5 is 0.25° . Here we used the 3-hourly product from 1990 to 2018.

177 **2.3. Satellite Absolute Dynamic Topography**

178 The along-track absolute dynamic topography (ADT) data used in the study were provided by the
179 Copernicus Marine and Environment Monitoring Service (CMEMS,
180 <http://www.marine.copernicus.eu>). The product is processed by the Data Unification and Altimeter
181 Combination System (DUACS) which applies to multi-mission altimeter data. The data are
182 comprised from the TOPEX/POSEDON mission, together with the Jason-1, Jason-2, and Jason-3
183 missions. Since Denmark Strait is close to the northern turning point of the orbits, the along-track
184 data have spatial and temporal resolutions of roughly 12 km and 2 days, respectively. The time
185 period of data coverage used here is 1993 to 2018.

186 **3. Basic Characteristics**

187 **3.1. Mean State**

188 We first present the mean Látrabjarg sections of hydrography and absolute geostrophic velocity
189 using the 22 realizations that include velocity data (Fig. 4). We don't consider the regions on the
190 east and west side of the strait where the number of occupations is less than five. Encouragingly,
191 the mean distributions of potential temperature and salinity are consistent with the analogous
192 means presented in Mastropole et al. (2017) using 111 occupations. It indicates that our mean view
193 using a smaller number of sections is representative. The warm and salty water on the Iceland shelf
194 is the Irminger Water originating from the south (the near-surface fresh water at the eastern end of
195 the section is likely associated with the Iceland coastal current, Logemann et al., 2013). To the
196 west, the vertically varying temperature and salinity reflects several water masses. In the upper
197 layer, the cold and fresh water, referred to as Polar Surface Water, emanates from the Arctic Ocean
198 via Fram Strait (de Steur et al., 2009; Håvik et al., 2017a). Beneath this, the warm water at the

199 western edge of the section, centered near 150 m, is Irminger Water that has recirculated north of
200 the strait (Mastropole et al., 2017; Casanova et al., submitted). Near the bottom is the DSOW,
201 denser than 27.8 kg m^{-3} (this isopycnal is highlighted in Fig. 4). As noted above, this is a
202 combination of Atlantic-origin Overflow Water (AtOW) and Arctic-origin Overflow Water
203 (ArOW), which is banked up on the western side of the trough (see also Våge et al., 2011; Harden
204 et al., 2016; Mastropole et al., 2017). The breakdown between these water masses is addressed in
205 the next section.

206 Figure 4d shows the mean section of absolute geostrophic velocity. This is the first such view of
207 the average, full water column velocity structure across Denmark Strait. The strong poleward flow
208 in the vicinity of the Iceland shelfbreak is the NIIC, which transports Irminger Water into the
209 Iceland Sea. Seaward of the NIIC there are two bands of southward, bottom-intensified flow
210 associated with tilting isopycnals sloping downward from west to east. The stronger band of flow
211 is located on the western side of the deep trough and transports the densest DSOW. The second
212 band is situated near the East Greenland shelfbreak. As noted above, the NIJ, separated EGC, and
213 shelfbreak EGC all advect water into Denmark Strait (Fig. 1). The year-long mooring dataset
214 across the Blosseville Basin used by Harden et al. (2016) revealed that, in the mean, the NIJ and
215 separated EGC were partially merged at that location. Our results demonstrate that, in Denmark
216 Strait, these two currents are fully merged and correspond to the stronger band of flow in Fig. 4d
217 which transports the majority of the DSOW. The weaker band of flow to the west is the shelfbreak
218 EGC. Note, however, that there is only a slight minimum in flow between the shelfbreak EGC and
219 the merged NIJ/separated EGC (Fig. 4d), which indicates that all three branches have combined
220 to some degree in the narrow strait.

221

222 3.2. Partitioning the DSOW Transport

223 The transport of DSOW (denser than 27.8 kg m^{-3}) in the vicinity of Denmark Strait has been
224 estimated in many studies. Harden et al. (2016) reported a year-long mean value of $3.54 \pm 0.16 \text{ Sv}$
225 from the mooring array across the Blosseville Basin in 2011–12. Jochumsen et al. (2012) estimated
226 the value to be $3.40 \pm 0.60 \text{ Sv}$ using one or two moorings in the center of the strait from 1996–
227 2011. This value was later updated by Jochumsen et al. (2017) to be $3.20 \pm 0.50 \text{ Sv}$, accounting
228 for known biases in the near-bottom current measurements and using a new method developed
229 from extended measurements. Based on the mean section of Fig. 4d, we obtain a transport of 3.00
230 $\pm 0.29 \text{ Sv}$ (Table 2). This is lower than the previous estimates because our mean section only
231 extends $\sim 20 \text{ km}$ west of the East Greenland shelfbreak. Transects that extend across the entire
232 Denmark Strait reveal that DSOW is found far onto the Greenland shelf, and the limited velocity
233 information there implies weak mean flow (Brearley et al., 2012; Jochumsen et al., 2012). The
234 Greenland shelf contribution in the model of Macrander (2004) is roughly 0.40 Sv , which has been
235 accounted for in the estimates of Jochumsen et al. (2012). Results from a mooring on the Greenland
236 shelf (30 km west of the trough) implied a similar value of 0.50 Sv (Jochumsen et al., 2017) (the
237 mooring array used by Harden et al. (2016) encompassed the DSOW on the Greenland shelf).
238 Linearly extrapolating our mean section of Fig. 4d to the Greenland coast gives a value of 0.54 Sv
239 for the missing transport, which is in line with the estimates above. Thus, our adjusted total
240 transport of $3.54 \pm 0.29 \text{ Sv}$ is comparable to the previous DSOW transport estimates.

241 Our hydrographic and velocity data provide the opportunity to partition the overflow transport by
242 water masses. To do this, we applied the same water mass end-member technique of Mastropole
243 et al. (2017) to our 22 occupations. Mastropole et al. (2017) defined four endmembers: AtOW,
244 ArOW, Polar Surface Water, and Irminger Water. They devised two mixing triangles, one of which

245 used the first three water masses, and the other using the latter three water masses – the assumption
246 being that AtOW does not mix with Irminger Water. For each station of a given occupation we
247 computed the corresponding end-member percentages, then gridded these to make vertical
248 sections. The mean sections of percentage for the two types of overflow water are shown in Fig.
249 5.

250 The ArOW dominates the deep trough where the merged NIJ/separated EGC is located, accounting
251 for close to 100% of the overflow water near the bottom. By contrast, the AtOW percentage is
252 highest in the vicinity of the Greenland shelfbreak, with large values in the shelfbreak EGC. Note,
253 however, that the percentage of ArOW is comparable to that of the AtOW in this region – this
254 further indicates merging/mixing of the three branches. Using the mean water mass end member
255 percentage sections (Fig. 5) in conjunction with the mean velocity section (Fig. 4d), we get a
256 transport of 1.72 ± 0.15 Sv for ArOW and 0.95 ± 0.10 Sv for AtOW (Table 2). If we assume that
257 the unresolved portion of the flow on the Greenland shelf is predominantly AtOW, this boosts the
258 transport of this water mass to 1.49 ± 0.10 Sv. Hence, we conclude that the mean transports of the
259 two types of overflow water are comparable in Denmark Strait. The remaining transport ($0.33 \pm$
260 0.04 Sv) corresponds to small contributions from the Polar Surface Water and Irminger Water
261 getting mixed into the top of the overflow layer.

262 As noted in the introduction, Harden et al. (2016) partitioned the overflow transport into the three
263 flow branches using data from the upstream Blosseville Basin mooring array. Using four shipboard
264 occupations of the same line, Våge et al. (2013) did the same partitioning with generally consistent
265 results. While the two bands of enhanced southward flow in our mean velocity section reflect the
266 shelfbreak EGC and merged NIJ/separated EGC, respectively (Fig. 4d), the degree to which all

267 three branches have merged/mixed in the strait makes it impossible to do precise partitioning here.
268 It is nonetheless instructive to consider the geographical distribution of the overflow transport.

269 We specify the boundary between the nominal shelfbreak EGC and merged NIJ/separated EGC to
270 be -25 km (i.e. the location of the velocity minimum between the two bands, Fig. 4d). It follows
271 that the shelfbreak EGC transports 0.85 ± 0.14 Sv, while the merged flow accounts for 2.15 ± 0.15
272 Sv (Table 2). By comparison, Harden et al. (2016) calculated 2.04 ± 0.16 Sv for the merged flow
273 and 1.50 ± 0.16 for the shelfbreak EGC. It is safe to assume that the inshore flow on the Greenland
274 shelf at the Látrabjarg line originated from the shelfbreak EGC upstream (recall that Blosseville
275 mooring array captured all of the overflow water on the Greenland side, which was confined to
276 the region of the shelf edge). This increases our shelfbreak EGC transport to 1.39 ± 0.14 Sv, in
277 line with the Blosseville Basin estimate. Hence, our total transport, as well as the geographical
278 distribution of transport across the strait, is consistent with Harden et al.'s (2016) upstream
279 partitioning. With regard to the overflow water masses, our mean sections (Figs. 4d and 5) indicate
280 that the band of flow at the shelfbreak transports comparable amounts of AtOW and ArOW, while
281 the band of flow on the western flank of the trough transports roughly twice as much ArOW as
282 AtOW (Table 2). Again, this attests to the significant degree of exchange between the flow
283 branches as they converge in Denmark Strait.

284 **4. Dominant Variability**

285 We now consider the section to section variability in our 22 occupations, which is a reflection of
286 mesoscale processes. Using a mooring in the center of the Denmark Strait trough, von Appen et
287 al. (2017) showed that the two pronounced mesoscale features, boluses and pulses, are associated
288 with a cyclonic and anticyclonic sense of rotation, respectively. Following the definitions in

289 Mastropole et al. (2017), we identified 8 instances of a bolus and 9 instances of a pulse in our
290 collection of sections, (5 sections could not be classified as either type of feature). We found
291 relatively little difference in the across-strait structure of the alongstream velocity field in these
292 two scenarios. However, inspection of the individual sections revealed 15 cases characterized by
293 a strong cyclonic structure centered in the trough. Figure 6 shows the composite mean of these
294 realizations, compared to the composite of the remaining 7 sections (where again we have only
295 plotted regions with at least five realizations). In the former case, which is referred to as the
296 cyclonic state, both the northward-flowing NIIC near the Iceland shelfbreak and the southward-
297 flowing merged NIJ/separated EGC on the western flank of the trough are intensified, while the
298 shelfbreak EGC is weakened. In the latter case, referred to as the non-cyclonic state, the entire
299 trough contains equatorward flow, but it is weaker and more bottom-trapped. In addition, the NIIC
300 is weaker but there is enhanced poleward flow over much of the Iceland shelf. (The data coverage
301 is insufficient to say anything about the shelfbreak EGC in this state.) The hydrographic structure
302 is not noticeably different in the two states (not shown). The height of the overflow layer (i.e. the
303 height of the 27.8 kg m^{-3} isopycnal) is also similar in both composites, although the stronger flow
304 in the cyclonic composite results in a larger transport of DSOW. It is clear that these two states are
305 not reflective of boluses and pulses, which begs the question: what is the nature of this dominant
306 variability? We argue that it is related to wind forcing.

307 To help demonstrate this, we first characterized the velocity structure in the center of each section
308 by the lateral gradient of the depth-mean velocity across the trough. This is an effective metric that
309 characterizes the degree to which a given section is in the cyclonic state (i.e., the stronger the
310 gradient, the more cyclonic, and vice versa). Using the ERA5 reanalysis wind data, we then created
311 composites of the wind stress curl and wind vectors for the two extremes of the velocity gradient,

312 in particular the five strongest cases and five weakest cases (Fig. 7). The mean wind field during
313 the time of occupation the sections (see Table 1) go into the composites. In the former, the wind
314 in Denmark Strait is strongly out of the northeast and there is pronounced negative wind stress curl
315 over the Blosseville Basin. In the other extreme, the wind is weak and variable, while the wind
316 stress curl is weakly positive over the Blosseville Basin. Våge et al. (2013) showed that negative
317 wind stress curl, together with the closed isobaths of the Blosseville Basin, plays an important role
318 in the bifurcation of the EGC at the northern edge of the basin. This in turn would weaken the
319 shelfbreak EGC. Hence, the wind stress curl pattern in Fig. 7a is conducive for enhancement of
320 the merged NIJ/separated EGC in the trough and decreased flow of the shelfbreak EGC, as seen
321 in the composite of Fig. 6a. In the other extreme the wind stress curl would weaken the merged
322 flow, consistent with the composite of Fig. 6b.

323 The wind stress curl forcing, however, does not explain the variation in the NIIC between the two
324 states. To address the potential role of the along-strait wind, we employed the along-track ADT
325 data (Figure 2 shows the satellite tracks in the region). Using the 25 years of ADT data, we
326 computed the cross-track component of surface geostrophic velocity for each of the satellite
327 crossings and projected these to the Látrabjarg line (see Spall et al. (2019) for details on the
328 methodology). We note that both the NIIC and merged NIJ/separated EGC have a strong surface
329 signature (Fig. 4). Next, we created composites of the surface velocity corresponding to the
330 associated wind conditions in the strait (averaged in a $1^\circ \times 1^\circ$ box around the trough). In particular,
331 we composited all of the satellite crossings for strong northeasterly wind (greater than the mean
332 plus one standard deviation), strong southwesterly wind (same criterion), and for all remaining
333 cases. The results are shown in Figure 8. This demonstrates that when the winds are strongly out
334 of the northeast the NIIC is both stronger and located more seaward, plus the merged NIJ/separated

335 EGC is enhanced as well. This is consistent with the fact that the cyclonic state (Fig. 6a)
336 corresponds to strong northeasterly winds (Fig. 7a).

337 We note that in Fig. 8c that the signatures of the NIIC and merged NIJ/separated EGC are much
338 broader than in the cyclonic velocity composite, plus the NIIC is located on the Iceland shelf and
339 the merged flow is located near the Greenland shelfbreak, versus being situated close to the trough.
340 This is likely due in part to the resolution of the altimeter (12 km) which is not well suited for
341 resolving either flow, plus the compositing process. However, it is also partly due to the fact that
342 not all instances with strong northeasterly wind correspond to a strong NIIC displaced to the west
343 – although this is clearly the case in the mean (Fig. 8c). To assess this, we composited the surface
344 geostrophic velocity for all of the instances when the NIIC was at the edge of the Iceland shelf,
345 regardless of wind conditions. This revealed a significantly narrower, stronger NIIC along with an
346 enhanced southward-flowing merged NIJ/separated EGC. Importantly, the mean wind for these
347 instances was strongly out of the northeast. This, together with Fig. 8c, indicates that the cyclonic
348 state in Denmark Strait is clearly associated with enhanced northeasterly winds through the strait.

349 It remains to be determined what the physical mechanism is behind this change in the NIIC.
350 Upwelling-favorable northeasterly winds should drive southward flow on the Iceland shelf due to
351 Ekman set up, i.e. the opposite of an enhanced NIIC, but the altimeter data are too inaccurate near
352 the coast to confirm this. The strong cyclonic flow offshore, in the vicinity of the trough, is
353 associated with a depression of the sea surface height. Such a signature would arise if there was
354 an increase in wind speed near the center of the strait, due to divergence of the offshore Ekman
355 transport. Unfortunately, the spatial resolution of ERA5 (~30 km) is insufficient to resolve such a
356 change. It should be noted that a two-dimensional view may not be appropriate here because the
357 presence of warm, relatively light water along the south coast of Iceland and cold, dense water

358 along the north slope, will result in an anti-cyclonic propagation of a high sea surface height signal
359 around the west coast of Iceland (Spall et al., 2017). This would act to maintain a high sea surface
360 height over the Iceland shelf to the east of the trough, even in the presence of upwelling-favorable
361 winds. Thus, the enhanced southward flow of the NIJ/separated EGC due to the negative wind
362 stress curl, lowering the sea surface height in the trough, would be concomitant with a stronger
363 northward-flowing NIIC. These ideas warrant further investigation, but are beyond the scope of
364 the present study.

365 **5. Dynamics in the trough**

366 **5.1 Hydraulic criticality**

367 Previous observations have shown that the density structure of the overflow water in Denmark
368 Strait is consistent with that of hydraulic flow over a sill (e.g. Spall and Price, 1998; Nikolopoulos
369 et al., 2003). Despite the fact that the dense water formation in the Nordic Seas is seasonal (e.g.
370 Brakstad et al., 2019), the overflow transport in Denmark Strait shows little to no seasonality
371 (Jochumsen et al., 2012; Harden et al., 2016). Furthermore, the different branches feeding the
372 overflow tend to vary out of phase with each other, such that the total transport remains fairly
373 steady. Together, these results suggest that the overflow through the strait is hydraulically
374 controlled. Using observations and a numerical model, Käse et al. (2003) diagnosed the hydraulic
375 conditions in Denmark Strait using the Froude number $Fr = v/\sqrt{g'D}$, where D is the vertical
376 length scale, $g' = g \Delta\rho/\rho$ is the reduced gravity, g is the gravitational acceleration, and $\Delta\rho$ is the
377 density difference across the interface. Käse et al. (2003) considered different parts of the domain
378 and found that the flow upstream of the sill is subcritical ($Fr < 1$), but, as the flow descends into
379 the Irminger Basin and accelerates, it becomes supercritical ($Fr > 1$). The transition location is

380 roughly 100 km downstream of the sill. This is consistent with theory (Pratt, 1986), observations
381 (Price and Baringer, 1994), and other models (Spall and Price, 1998) regarding the subcritical-to-
382 supercritical transition over a sill.

383 We investigated the Froude number using our 22 sections. In the scenario where the dense water
384 flows beneath a motionless or slowly moving upper layer, the Froude number is the expression
385 given above. In our case, especially for the cyclonic state, there is strong flow throughout the water
386 column. As such, it is more appropriate to use the composite Froude number G for two active
387 layers (Armi, 1986; Kösters, 2004; Pratt, 2008):

$$388 \quad G^2 = Fr_1^2 + Fr_2^2 \quad (1)$$

389 where $Fr_n = v_n/\sqrt{g'D_n}$ is the Froude number in the n^{th} layer. The quantity v_n is the vertically-
390 averaged advective speed in layer n , and the denominator is the internal gravity wave speed where
391 D_n is the layer thickness. A two-layer flow that is laterally uniform is considered supercritical when
392 $G > 1$ and subcritical when $G < 1$. For flows with strong lateral variations in layer thickness and
393 velocity, a local value of $G > 1$ indicates that the flow is locally supercritical, but does not
394 necessarily indicate that the flow as a whole is supercritical. In this case, locally generated
395 disturbances will propagate downstream whereas disturbances that exist over the whole channel
396 width may still propagate upstream (Pratt and Helfrich, 2005). Thus, a flow may be supercritical
397 at certain locations but also subcritical as a whole. We choose the 27.8 kg m⁻³ isopycnal as the
398 interface between the two layers, since this is the top of the dense overflow water and also
399 corresponds to the maximum in stratification (see also von Appen et al., 2017).

400

401 For each occupation we calculated G at the grid points across the section corresponding to the
402 southward flow. Figure 9 shows the results, where we have distinguished between the cyclonic
403 cases (red) and non-cyclonic cases (blue). The individual realizations are plotted as open circles,
404 and the means for the two cases at each cross-stream location are the solid circles. One sees that,
405 for the cyclonic state, the mean G exceeds 1 on the western flank of the trough where the merged
406 NIJ/separated EGC is strongest (Fig. 6a). In all, 11 out of the 15 cyclonic realizations had $G > 1$ in
407 this part of the strait. By contrast, the mean G for the non-cyclonic state is lower than 1 everywhere,
408 although 4 out of the 7 realizations had a value of $G > 1$ somewhere in the domain. As noted above,
409 models and observations indicate that the overflow plume descending from Denmark Strait reaches
410 hydraulic criticality approximately 100 km downstream of the sill. One is tempted to conclude
411 from our measurements that localized hydraulic criticality also occurs intermittently at the sill
412 itself, in the cyclonic configuration when the merged NIJ/separated EGC is intensified on the
413 western flank of the trough. However, the presence of such a confined region where $G > 1$ does
414 not necessarily imply that strait-wide hydraulic control is occurring (Pratt and Helfrich, 2005).
415 Further work is required to shed light on this.

416 **5.2 Mixing and Potential vorticity**

417 Although it remains unclear if the Denmark Strait sill can act as a location of strait-wide hydraulic
418 control akin to what happens farther south, the strong flow at the Látrabjarg line, in conjunction
419 with the weak stratification, result in another important aspect of supercritical flow – that of
420 mixing. This can be assessed by considering the gradient Richardson number, defined as the ratio
421 of the buoyancy frequency to the square of vertical shear in velocity,

422
$$Ri = -\frac{g}{\rho_0} \frac{\partial \rho}{\partial z} \left(\frac{\partial u}{\partial z} \right)^{-2}, \quad (2)$$

423 where ρ is the local density, ρ_0 is the background density (section-wide average), and u is the
 424 along-strait velocity. When Ri is less than the critical value of 0.25 the flow can be subject to
 425 Kelvin-Helmholtz instability, which leads to vertical mixing (in many studies the critical value is
 426 taken to be in the range 0.2–1.0; e.g. Galperin et al., 2007). In Figure 10 we show the vertical
 427 section of Ri for the July 2007 occupation, which is one of the sections where $G > 1$ within the
 428 trough. This reveals a region of $Ri < 0.25$ along the steeply-sloped density front separating the
 429 cold overflow water from the warm Irminger Water. In this case both the weak stratification and
 430 strong velocity shear contribute to the small value of Ri . It is expected that strong vertical mixing
 431 would be occurring in this region.

432 To further investigate the nature and extent of mixing at the Látrabjarg line, we consider the
 433 potential vorticity dynamics of the flow using our 22 occupations. We did this by evaluating the
 434 Ertel potential vorticity (e.g. Spall and Pedlosky, 2008; Lin et al., 2018),

435
$$\Pi = -\frac{f}{\rho_0} \frac{\partial \rho}{\partial z} + \frac{1}{\rho_0} \frac{\partial u}{\partial y} \frac{\partial \rho}{\partial z} - \frac{1}{\rho_0} \frac{\partial u}{\partial z} \frac{\partial \rho}{\partial y}, \quad (3)$$

436 where the y -direction is cross-strait, positive towards Iceland. The Ertel potential vorticity (PV)
 437 has three components: 1) the planetary stretching vorticity term, dictated by the vertical
 438 stratification and earth's rotation; 2) the relative vorticity term, due to the combination of the lateral
 439 gradient of the horizontal velocity and vertical stratification; and 3) the tilting vorticity term, due
 440 to the vertical gradient of horizontal velocity and lateral gradient of density. It is instructive to

441 normalize the second and third terms by the planetary stretching term. For the relative vorticity,
442 this gives

$$443 \quad R_o = -\frac{1}{f} \frac{\partial u}{\partial y}. \quad (4)$$

444 For the tilting vorticity the ratio is

$$445 \quad R_z = \frac{g\alpha}{f^2 \rho_0} \frac{\partial \rho}{\partial y}, \quad (5)$$

446 where α is the isopycnal slope; to derive this, we used the thermal wind relation, $\frac{\partial u}{\partial z} = \frac{g}{f \rho_0} \frac{\partial \rho}{\partial y}$.

447 Using a representative length scale L and velocity scale U , the first ratio (Eq. 4) can be expressed

448 as $R_o = \frac{U}{fL}$, which is the Rossby number. Taking ΔU as the change in velocity over the depth

449 scale, the second ratio (Eq. 5) can be expressed as $R_z = \frac{\Delta U}{fL}$, which has the form of a Rossby

450 number associated with the depth variation in velocity; we refer to this as the shear Rossby number.

451 Note that when the flow is barotropic R_z will be small, even though R_o could be large. When the

452 flow is strongly baroclinic R_z could be large.

453 Returning to the July 2007 occupation, we computed the total Ertel PV and its three components,

454 where the latter two terms have been normalized to show R_o and R_z (Fig. 11). One sees that over

455 most of the section the total PV is qualitatively similar to the stretching term, particularly in the

456 upper layer. However, in the trough the other two terms are significant. The large Rossby number

457 (up to 1.5) changes sign across the merged NIJ/separated EGC, indicating that this flow is highly

458 nonlinear and may be barotropically unstable (Pickart et al., 2005). Furthermore, the lateral

459 gradient of the total PV changes sign with depth in the trough (Fig. 11a), which is a necessary

460 condition for baroclinic instability. This is in line with the model results of Spall et al. (2019) who
461 identified that both the merged NIJ/separated EGC and NIIC are baroclinically unstable. This
462 instability acts to weaken the hydrographic front that is maintained by the convergence of the large
463 scale mean flow.

464 Note also in Fig. 11 that, due to the steeply-sloped isopycnals of the hydrographic front (and
465 corresponding strong velocity shear via thermal wind), the shear Rossby number is less than -1,
466 i.e. the same order as the Rossby number. This results in regions of negative total PV; in particular,
467 note the correspondence between the strong tilting vorticity in Fig. 11d and negative Π in Fig. 11a.
468 The condition of negative total PV can lead to symmetric instability (Haine and Marshall, 1998;
469 D'Asaro et al., 2011), a fast-growing instability that occurs on the order of a few hours (Brearley
470 et al., 2012). At finite amplitude this results in intense, rapid diapycnal mixing (Haine and
471 Marshall, 1998). We now explore further the signature of symmetric instability in our data.

472 **5.3 Symmetric instability**

473 Based on the July 2007 occupation, we seek to elucidate the relationship between the tilting
474 vorticity, or more specifically R_z , and the occurrence of negative PV. Using all the grid points of
475 the 22 realizations, we regressed R_z against Π (Fig. 12). This shows that when $R_z < -1$, 73% of the
476 time this corresponds to negative PV (if the threshold is strengthened to -1.5, the percentage of
477 negative PV is 93%). For the remaining 27% of the data points, the strong positive relative vorticity
478 overcomes the tilting vorticity such that the total PV remains positive. This is seen in Fig. 12,
479 where the value of R_o for each data point is indicated using color. The points in question generally
480 have $0.5 < R_o < 1.5$. Alternatively, the color in Fig. 12 reveals that when negative PV does not
481 correspond to $R_z < -1$ this is due to the flow having a large negative R_o (dark blue symbols in Fig.

482 12). We thus conclude that, outside of extreme instances of large relative vorticity (of either sign),
483 it is generally the case that when the shear Rossby number is less than -1, the total PV is negative
484 – which will result in symmetric instability. This is consistent with the classification of instability
485 in Thomas et al. (2013).

486 Part of our rationale for casting the symmetric instability condition in terms of R_z is that this ratio
487 does not depend on the velocity of the flow, but only on the density structure (see Eq. 5). As such,
488 we can extend the application of the proxy to the complete set of historical hydrographic Látrabjarg
489 sections (we exclude 9 short sections that did not cross the trough). We find that $R_z < -1$ in 60 of
490 the 112 sections, i.e. over 50% of the time (for the more restrictive criterion of $R_z < -1.5$ it is 42%).
491 This suggests that symmetric instability occurs quite frequently in Denmark Strait. Interestingly,
492 the presence of symmetric instability does not seem to be tied to the cyclonic or non-cyclonic
493 velocity states, or to the presence of boluses versus pulses.

494 To determine where in the water column the conditions for symmetric instability occur, we
495 tabulated the occurrences of $R_z < -1$ for all of the sections (Fig. 13). This reveals that the instability
496 occurs mainly in the trough, with a few instances near the surface on the Iceland shelf and near the
497 bottom in the vicinity of the Greenland shelfbreak. However, the majority of cases are clustered
498 into two areas: a deeper region near the western side of the trough, and a shallower region closer
499 to the eastern side of the trough. To shed light on the underlying reasons for this pattern, we
500 constructed a composite hydrographic section for all of the occurrences in the deeper region, then
501 did the same for the shallower region. These are shown in Fig. 14. For the shallower occurrences
502 there is a large amount of cold overflow water filling the trough, while for the deeper occurrences
503 there is only a thin layer of this water banked on the western side of the trough. These two states
504 correspond nearly identically with the hydrographic patterns of boluses and pulses, respectively

505 (von Appen et al., 2017). Note that in both cases the instability takes place in the steep frontal
506 zone, where the tilting vorticity is strongly negative. This result suggests that strong vertical
507 mixing occurs at the top of the overflow layer, regardless of whether there is a large or small
508 amount of dense water present. The implication is that, even though strait-wide hydraulic control
509 may not be achieved until downstream of Denmark Strait, the mixing/entrainment process that
510 modifies the overflow water begins at the sill (also see North et al., 2018).

511 **6. Summary and Discussion**

512 In this study we have used 22 occupations of the Látrabjarg line from 1993–2018, together with
513 reanalysis wind fields and satellite absolute dynamic topography data, to investigate the kinematic
514 structure and dynamics of the Denmark Strait Overflow Water. While the Látrabjarg section has
515 been occupied over 100 times through the years, the unique aspect of the subset considered here
516 is that it includes shipboard velocity data that were used to construct vertical sections of absolute
517 geostrophic velocity. The mean velocity section reveals the presence of the shelfbreak EGC in the
518 vicinity of the Greenland shelf edge, the merged NIJ/separated EGC banked against the western
519 side of the deep trough, and the northward-flowing NIIC near the Iceland shelfbreak.

520 The mean transport of the overflow water (denser than 27.8 kg m^{-3}) is $3.54 \pm 0.29 \text{ Sv}$, which
521 includes an extrapolated estimate of the unresolved component on the Greenland shelf (0.54 Sv).
522 This is close to previously published estimates of the mean overflow transport (Harden et al., 2016;
523 Jochumsen et al., 2017). We partitioned the transport in terms of water masses and current
524 components. For the former we used a hydrographic end-member analysis to distinguish Atlantic-
525 origin Overflow Water (AtOW) from Arctic-origin Overflow Water (ArOW). Assuming that the
526 unresolved overflow transport on the Greenland shelf is AtOW, this gives $1.72 \pm 0.15 \text{ Sv}$ for

527 ArOW and 1.49 ± 0.10 Sv for AtOW, indicating that the mean transports of the two types of
528 overflow water are comparable in Denmark Strait. For the currents, we distinguished the
529 shelfbreak EGC and the merged NIJ/separated EGC using a geographical boundary, and assumed
530 that the unresolved overflow water on the Greenland shelf emanated from the shelfbreak EGC
531 upstream of the strait. This gives 1.39 ± 0.14 Sv for the shelfbreak EGC and 2.15 ± 0.15 Sv for the
532 merged flow, which is in line with similar partitioning done by Harden et al. (2016) upstream in
533 the Blosseville Basin. Notably, both currents transport both types of overflow water, implying a
534 significant degree of exchange between the branches as they converge in Denmark Strait.

535 With regard to temporal variability, there were two dominant configurations of the flow which we
536 refer to as the cyclonic state and the non-cyclonic state. The former is characterized by a strong
537 southward flow of the merged NIJ/separated EGC adjacent to a strong northward flow of the NIIC.
538 This structure was present in 15 of the 22 occupations. In this state the NIIC is located farther to
539 the west and occupies part of the trough. The remaining 7 sections corresponded to weaker
540 southward and northward flows, with the NIIC shifted eastward and the entire trough associated
541 with the merged NIJ/separated EGC. Using the reanalysis wind data, it was demonstrated that the
542 cyclonic state corresponds to negative wind stress curl north of the strait in the Blosseville Basin
543 and strong northeasterly winds within the strait. The former is conducive for an enhanced merged
544 flow as demonstrated previously (Våge et al., 2013). Using the satellite surface geostrophic
545 velocity data, we showed that the NIIC becomes stronger and shifts closer to the shelfbreak under
546 northeasterly winds, although the physical mechanism for this remains unresolved.

547 The hydraulic criticality of the flow was assessed using a composite Froude number that can
548 account for two moving layers – the overflow layer and the lighter water above. This revealed that
549 roughly two thirds of the cyclonic realizations had regions of super-critical flow in the trough, and

550 this condition was present in the mean for the strongest flow in the merged NIJ/separated EGC.
551 This suggests that hydraulic control could be occurring intermittently during the cyclonic state.
552 However, the presence of such a confined region of large Froude number does not necessarily
553 imply that strait-wide hydraulic control is occurring (Pratt and Helfrich, 2005).

554 A potential vorticity (PV) analysis of the 22 occupations indicated that the flow through Denmark
555 Strait is subject to symmetric instability. This occurs when the total PV is negative, which tends
556 to happen when the tilting vorticity becomes strongly negative. We determined that the shear
557 Rossby number (R_z) is a good proxy for determining when symmetric instability is active. In
558 particular, when R_z is less than -1, the total PV is typically negative. This proxy, which does not
559 rely on the flow speed but only the density structure, was then applied to the full set of 122
560 Látrabjarg occupations. This revealed that symmetric instability tends to occur at the top of the
561 overflow layer, regardless of whether there is a large or small amount of dense water in the strait.
562 Symmetric instability is a fast-growing instability that reaches finite amplitude in a matter of hours,
563 leading to intense vertical mixing. This implies that, even though hydraulic criticality may not be
564 achieved until downstream of the strait, the mixing/entrainment process that modifies the overflow
565 water begins at the sill.

566 Previous work has implied that the dominant mesoscale variability in Denmark Strait is due to
567 baroclinic instability of the hydrographic front that separates the overflow water from the
568 subtropical-origin water in the NIIC (Spall et al., 2019). The resulting meanders of the front
569 propagate equatorward through the strait and are associated with the well-known boluses and
570 pulses of overflow water (Mastropole et al., 2017; von Appen et al., 2017). In particular, meander
571 crests are associated with boluses, which correspond to thick layer of overflow water, whereas
572 meander troughs coincide with pulses, which are characterized by a thin layer of overflow water.

573 The results presented here suggest that the dominant variation in alongstream velocity at the sill is
574 wind-driven, rather than being associated with the amount of overflow water present. There are
575 several factors that may help explain this apparent discrepancy.

576 The numerical model results of Almansi et al. (2017) show that, relative to the background state,
577 the biggest difference in the alongstream velocity signature of the boluses and pulses is the bottom
578 intensification associated with the latter. While we don't have enough realizations of the Látrabjarg
579 section with velocity to determine a background state, our composite of pulse realizations shows
580 significantly more bottom-intensification in the trough versus the composite of bolus realizations,
581 in line with Almansi et al.'s (2017) results. Another thing to keep in mind is that the mooring
582 analysis of von Appen et al. (2017) showed that the most conspicuous difference between the
583 passage of boluses versus pulses pertains to the cross-stream velocity signal (cyclonic for boluses,
584 anti-cyclonic for pulses), which we are unable to assess. Both features were associated with an
585 enhancement of the alongstream velocity in the overflow layer. The maximum flow in von Appen
586 et al.'s (2017) bolus composite exceeded 0.40 m s^{-1} , while that for their pulse composite exceeded
587 0.60 m s^{-1} . In our composite vertical sections, the mean near-bottom flow of the pulses is only
588 slightly larger than for the boluses (0.30 m s^{-1} versus 0.24 m s^{-1}), but it must be kept in mind that
589 the mooring composites were based on vastly more data. In any event, both the shear and the
590 magnitude of the alongstream flow – together with the strong hydrographic signals – suggest that
591 we indeed detect these mesoscale features.

592 A final consideration regarding the velocity variability seen in our dataset is the short timescale
593 associated with the passage of the boluses and pulses. The mooring composites of von Appen et
594 al. (2017) indicate that, for both types of features, the strongest signals in alongstream velocity
595 persist for approximately 12 hours. Typical occupations of the Látrabjarg line take a day or more

596 to complete. This means that the timing has to be perfect for a shipboard transect to capture the
597 peak alongstream velocity signature of one these mesoscale features in the trough. On the other
598 hand, the wind-driven flow variability takes place over longer timescales. The ERA5 data indicate
599 that the auto-correlation time for the along-strait winds is 73 hours. Therefore, it is more likely that
600 a given transect will be under the influence of a single wind state. As the collection of Látrabjarg
601 occupations with velocity continues to increase over time, we will be better positioned to elucidate
602 the impacts of external versus internal forcing of the overflow water.

603

604 **Acknowledgements**

605 An inordinate amount of effort was required to obtain, processes, and quality-control the data used
606 in this study. The authors are thankful for the efforts of the many individuals, both at sea and
607 ashore, that made the analysis possible. We are also indebted to M. Spall, T. Haine, and S. Tan for
608 invaluable discussions and insights. Funding for the study was provided by National Science
609 Foundation (NSF) grants OCE-1259618, OCE-1756361, and OCE-1558742. The German research
610 cruises were financially supported through various EU Projects (e.g. THOR, NACLIM) and
611 national projects (most recently TRR 181 "Energy Transfer in Atmosphere and Ocean" funded by
612 the German Research Foundation and RACE II "Regional Atlantic Circulation and Global
613 Change" funded by the German Federal Ministry for Education and Research). GWKM
614 acknowledges the support of the Natural Sciences and Engineering Research Council of Canada.
615 LP was supported by NSF Grant OCE-1657870.

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753 **Tables**

754 Table 1. Occupations of the 22 Látrabjarg transects with velocity measurements used in the study
 755 (VMADCP: vessel-mounted ADCP; LADCP: lowered ADCP). The 11 occupations marked with
 756 an asterisk denote the sections added to the Mastropole et al. (2017) dataset.

Cruise	Ship	Year	date	ADCP
WOCE-93*	R/V <i>Aranda</i>	1993	Aug 30	VMADCP
MSM05-4	R/V <i>Maria S. Merian</i>	2007	Jul 12 – 13	VMADCP
KN194	R/V <i>Knorr</i>	2008	Oct 13 – 14	VMADCP
BS010	R/V <i>Bjarni Saemundsson</i>	2009	Aug 11 – 12	LADCP
M82-1	R/V <i>Meteor</i>	2010	Jul 07	VMADCP
BS001	R/V <i>Bjarni Saemundsson</i>	2011	Feb 10	LADCP
M85-2	R/V <i>Meteor</i>	2011	Aug 19 – 21	VMADCP
KN203	R/V <i>Knorr</i>	2011	Aug 22 – 24	VMADCP
BS002	R/V <i>Bjarni Saemundsson</i>	2012	Feb 08 – 09	LADCP
MSM21-1b	R/V <i>Maria S. Merian</i>	2012	Jun 10 – 11	VMADCP

JR267	R/V <i>James Clark Ross</i>	2012	Jul 28 – 29	LADCP
P437	R/V <i>Poseidon</i>	2012	Aug 10 – 12	VMADCP
BS013*	R/V <i>Bjarni Saemundsson</i>	2013	Feb 06 – 07	LADCP
P471-2*	R/V <i>Poseidon</i>	2014	Jul 06 – 08	VMADCP
P486*	R/V <i>Poseidon</i>	2015	Jun 14 – 25	VMADCP
BS015*	R/V <i>Bjarni Saemundsson</i>	2015	Aug 24	LADCP
P503*	R/V <i>Poseidon</i>	2016	Aug 04 – 06	VMADCP
BS017*	R/V <i>Bjarni Saemundsson</i>	2017	Aug 05 – 06	LADCP
64PE426*	R/V <i>Pelagia</i>	2017	Sep 16 – 18	VMADCP
ALL0118*	NRV <i>Alliance</i>	2018	Mar 20 – 21	VMADCP
MSM76*	R/V <i>Maria S. Merian</i>	2018	Aug 14 – 15	VMADCP
AR306*	R/V <i>Armstrong</i>	2018	Oct 01	VMADCP

757 Table 2. Partitioning of the DSOW transport (Sv) by water masses and current components. The
 758 values in parentheses include the unresolved portion on the Greenland shelf (0.54 Sv).

Water masses \ Currents	Shelfbreak EGC	Merged NIJ/separated EGC	Total
Arctic-origin Water	0.41 ± 0.06	1.31 ± 0.09	1.72 ± 0.15
Atlantic-origin Water	0.33 ± 0.06	0.62 ± 0.04	0.95 ± 0.10 (1.49 ± 0.10)
Other water masses	0.11 ± 0.02	0.22 ± 0.02	0.33 ± 0.04
Total	0.85 ± 0.14 (1.39 ± 0.14)	2.15 ± 0.15	3.00 ± 0.29 (3.54 ± 0.29)

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768 **Figure captions**

769 Figure 1. Schematic circulation in the vicinity of Denmark Strait, including the two branches of
770 the East Greenland Current (EGC) – the shelfbreak EGC and separated EGC – as well as the North
771 Icelandic Jet (NIJ) and the North Icelandic Irminger Current (NIIC). The blue line across Denmark
772 Strait is the Látrabjarg transect from Mastropole et al. (2017). The bathymetry is from ETOPO2v2.
773 Bathymetry contours are in meters.

774 Figure 2. Locations of the hydrographic and satellite measurements used in the study. The red dots
775 are the CTD stations. There are a total of 122 CTD sections (many of them overlap, hence the
776 dense clustering of red dots). The Látrabjarg transect from Mastropole et al. (2017) is the black
777 line. The blue dots are the absolute dynamic topography (ADT) altimeter measurements.

778 Figure 3. Temporal distribution of the 122 Látrabjarg hydrographic sections. Those occupations
779 that include velocity measurements are colored red (for vessel-mounted ADCP data) and magenta
780 (for lowered ADCP). The blue circles correspond to hydrographic measurements only.

781 Figure 4. Mean vertical sections of the 22 occupations of the Látrabjarg transect. (a) data coverage,
782 (b) potential temperature ($^{\circ}\text{C}$), (c) salinity, and (d) absolute geostrophic velocity (m s^{-1}) overlain
783 by potential density (kg m^{-3}) contours. Positive (negative) velocities are equatorward (poleward).
784 The highlighted isopycnal of 27.8 kg m^{-3} is the upper boundary of the overflow water. The Iceland
785 shelf is on the east side of the trough (positive distance), and the Greenland shelf is on the west
786 side (negative distance).

787 Figure 5. Mean vertical sections of the percent presence of the end-member for (a) ArOW and (b)
788 AtOW. The highlighted isopycnal of 27.8 kg m^{-3} is the upper boundary of the overflow water.

789 Figure 6. Composites of absolute geostrophic velocity (m s^{-1} , color) overlain by potential density
790 (kg m^{-3} , contours) for the (a) cyclonic and (b) non-cyclonic cases. Positive velocities are
791 equatorward. The highlighted isopycnal of 27.8 kg m^{-3} is the upper boundary of the overflow
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795 the Denmark Strait trough. The green line denotes the Látrabjarg transect. The trough is marked
796 by the red star.

797 Figure 8. Composites of along-strait surface geostrophic velocity (right column, m s^{-1})
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799 the standard error. (a) Average of all instances where the northeasterly wind in the strait is greater
800 than the mean plus one standard deviation. (b) Instances where the wind is close to the mean. (c)
801 Instances where the southwesterly wind is greater than the mean plus one standard deviation.

802 Figure 9. Composite Froude number G as a function of across-strait distance, for the cyclonic cases
803 (red) and anti-cyclonic cases (blue). The individual values for the 22 occupations are open circles,
804 and the mean values at each location are the filled circles. The critical value of $G = 1$ is indicated
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806 Figure 10. Vertical sections of (a) gradient Richardson number (color) and (b) absolute geostrophic
807 velocity (m s^{-1} ; color) overlain by potential density (kg m^{-3} ; contours) for the Látrabjarg occupation
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822 colored by number of realizations.

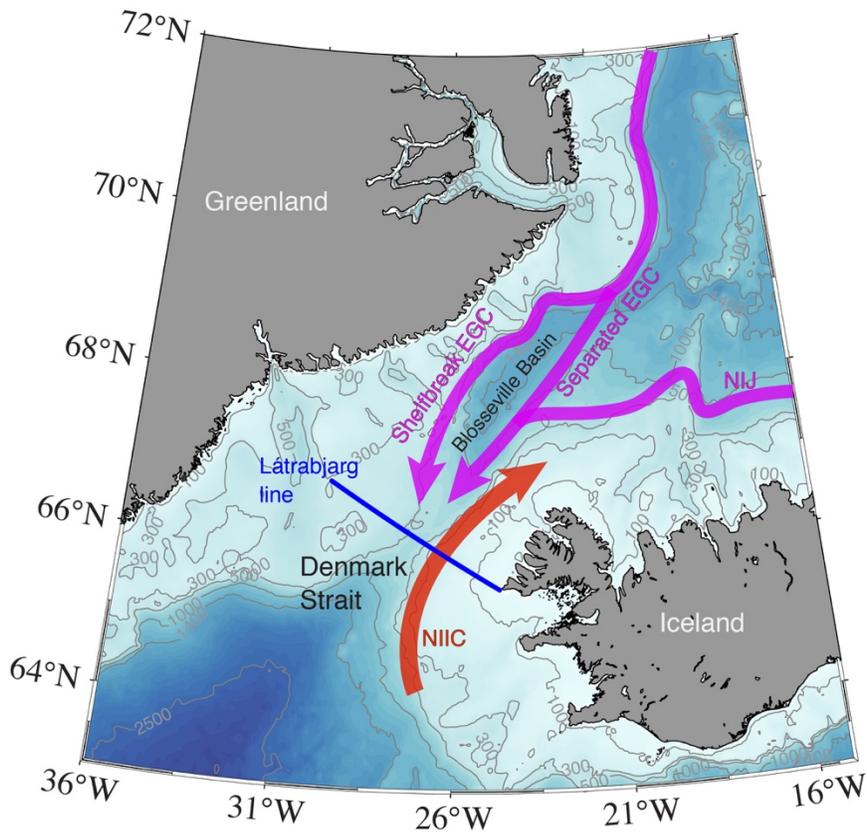
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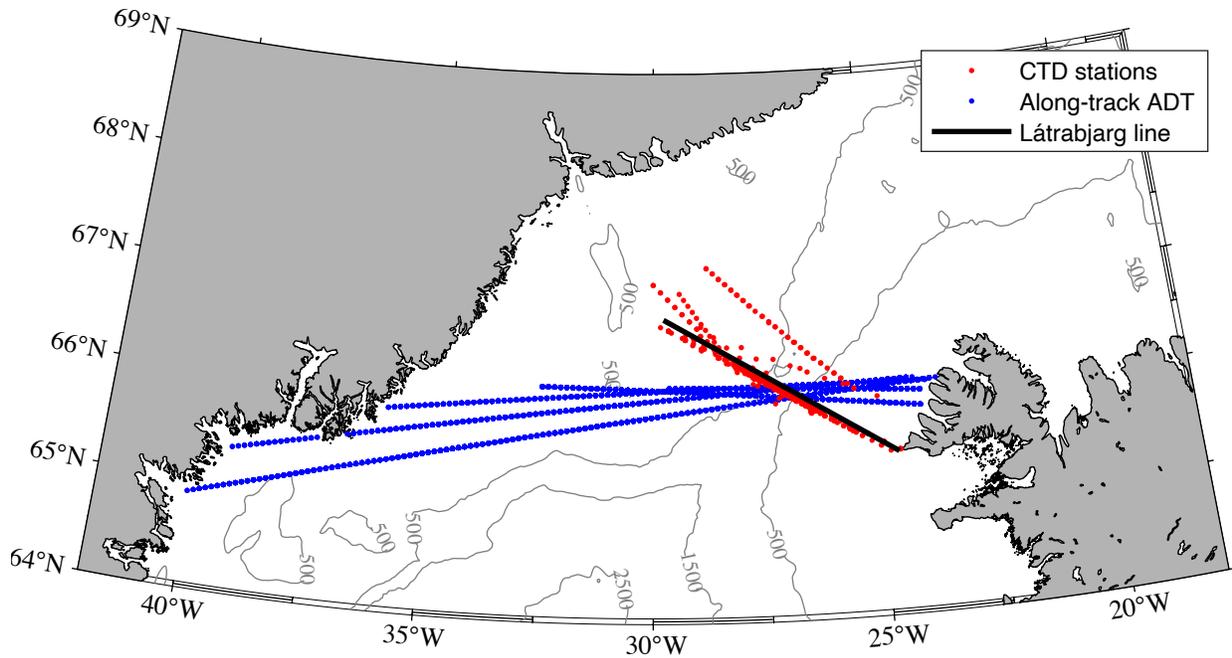
830 **Figures**



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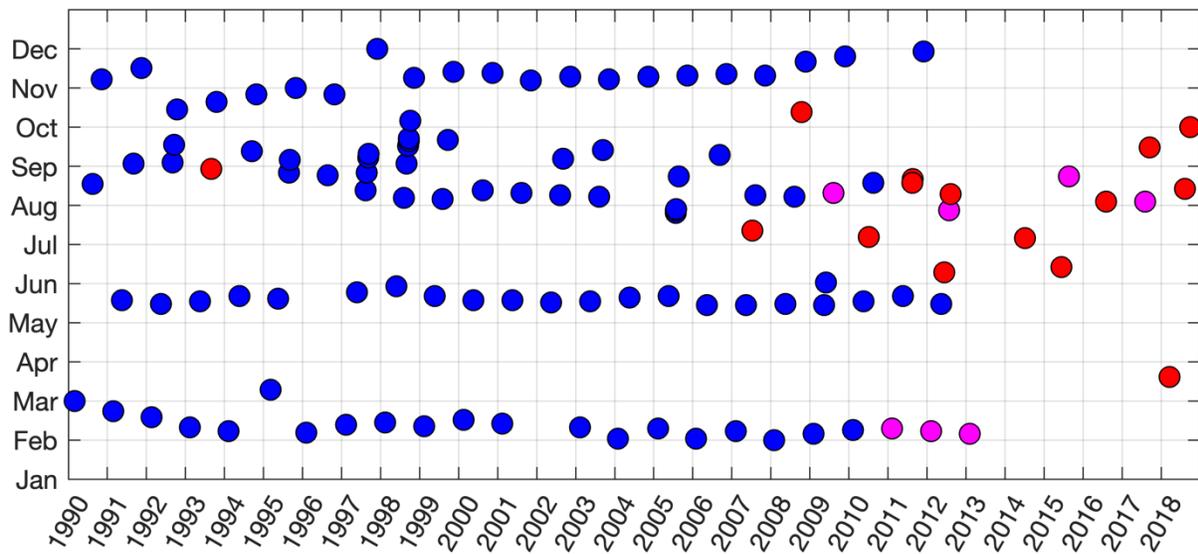
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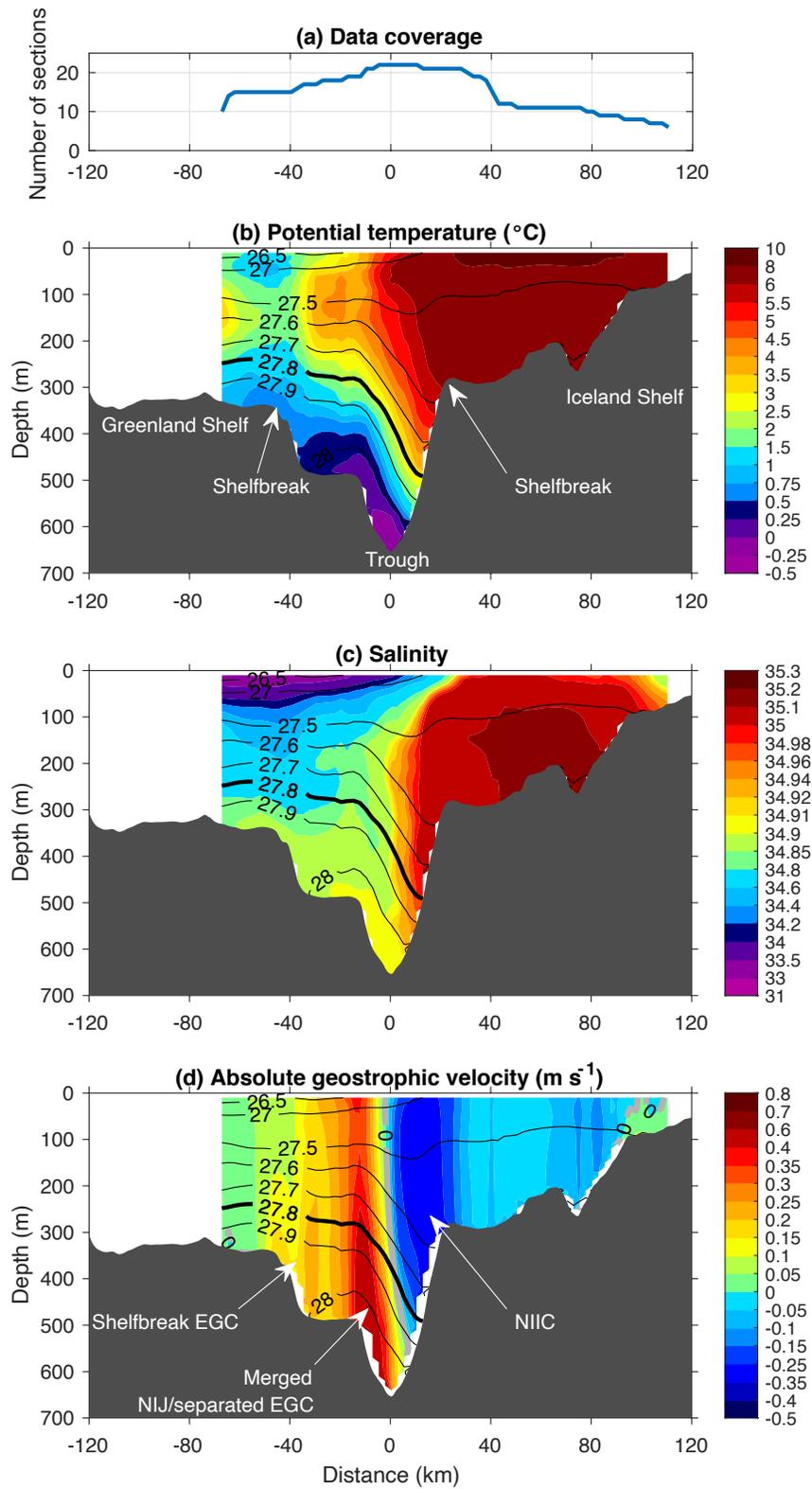
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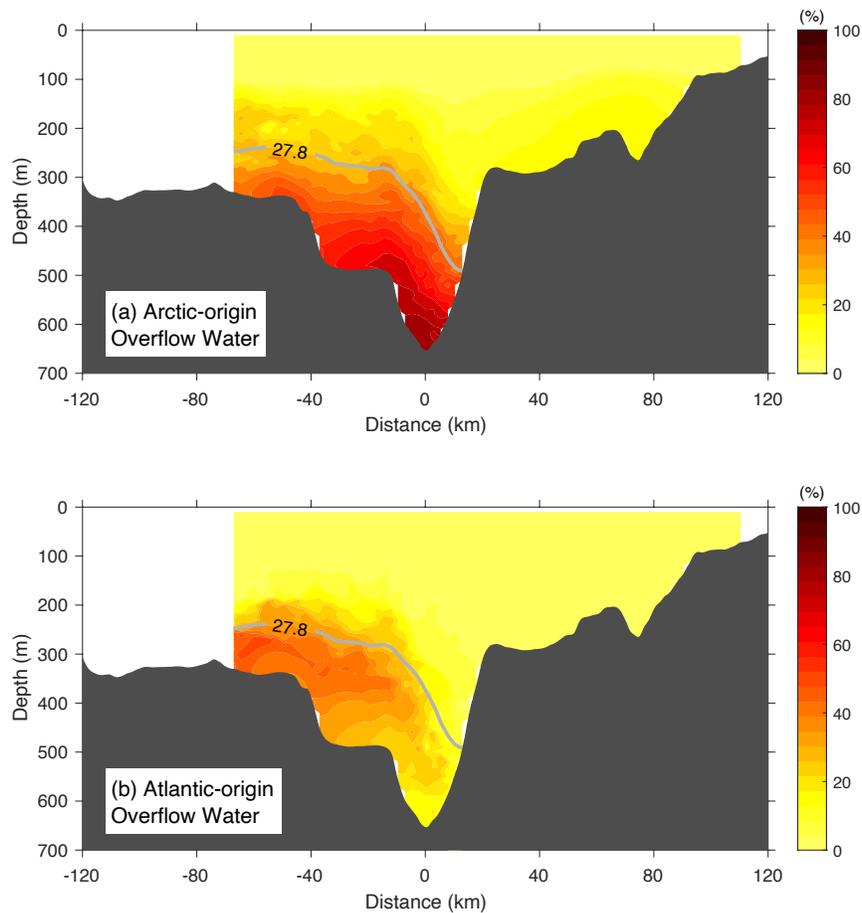
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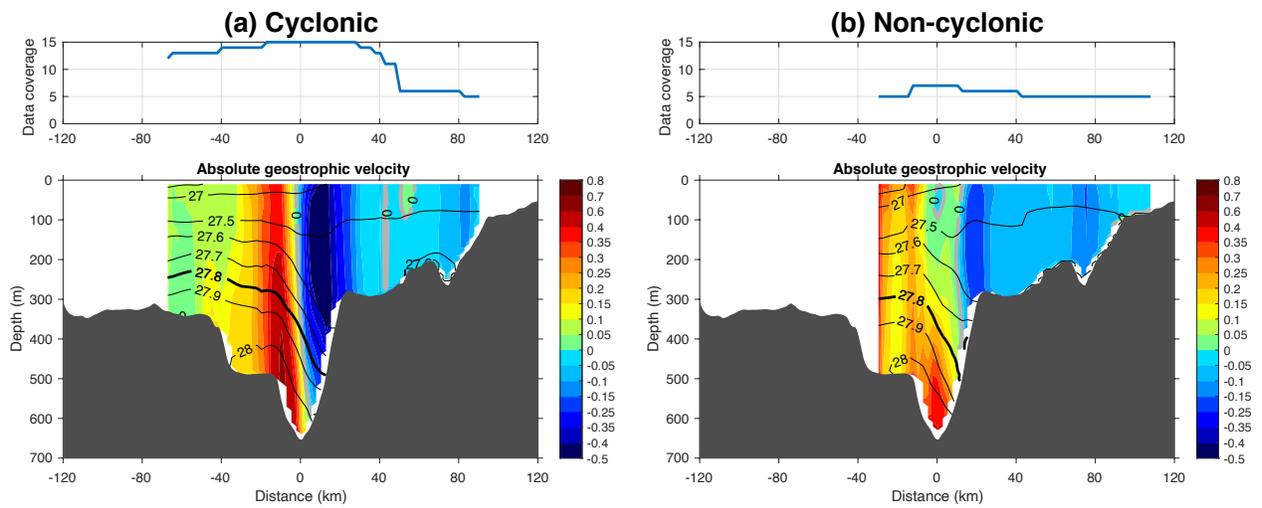
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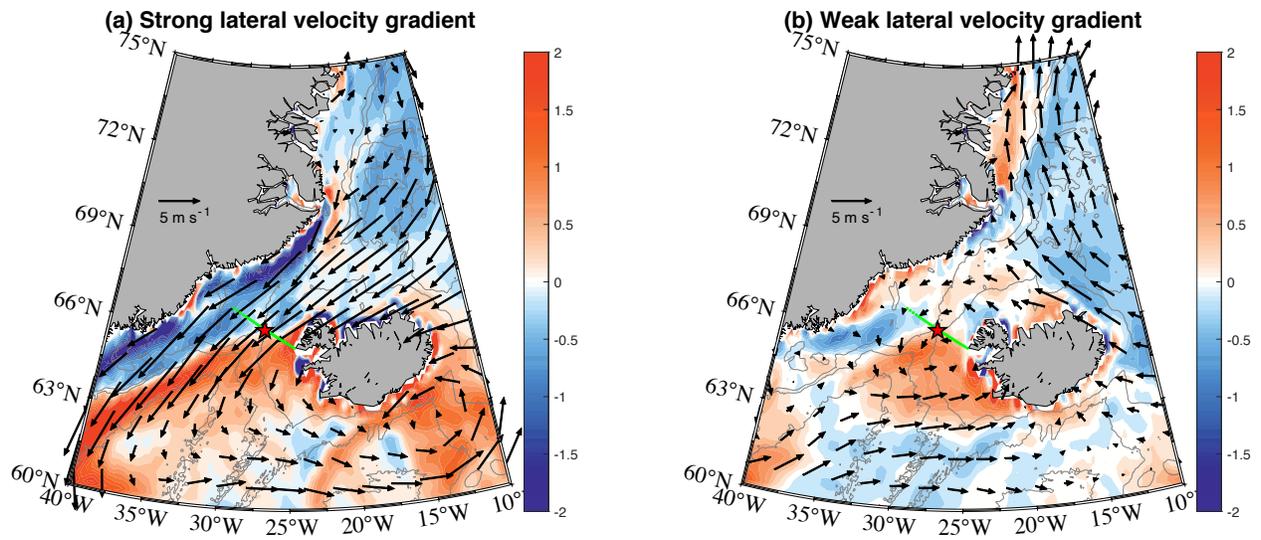
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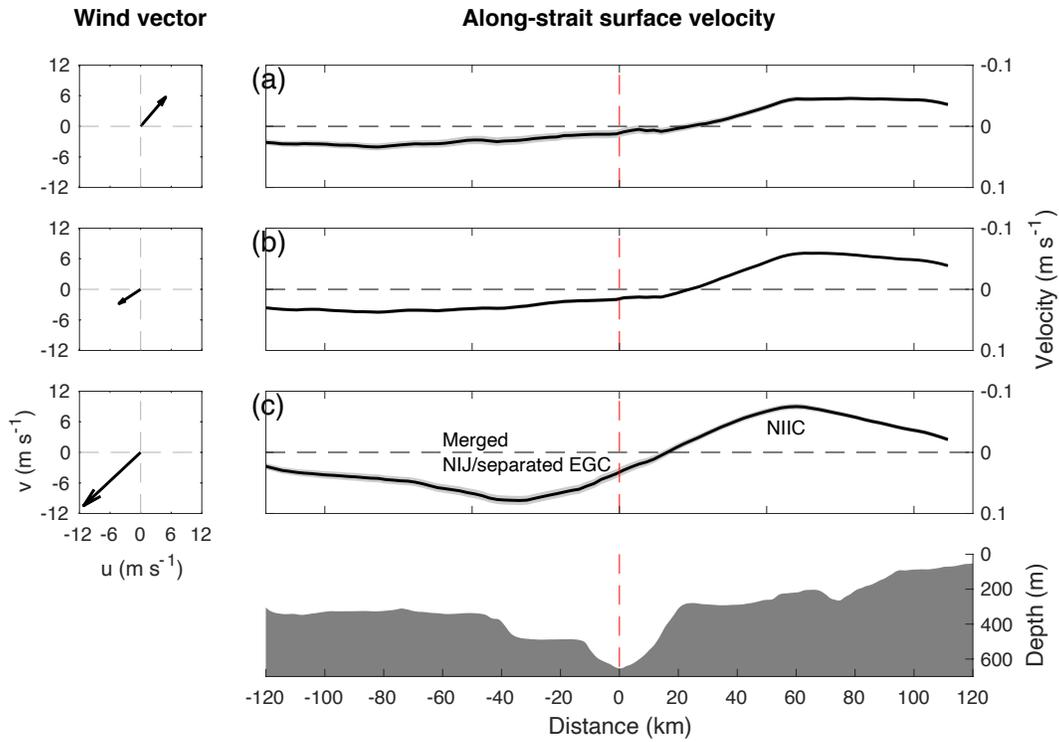


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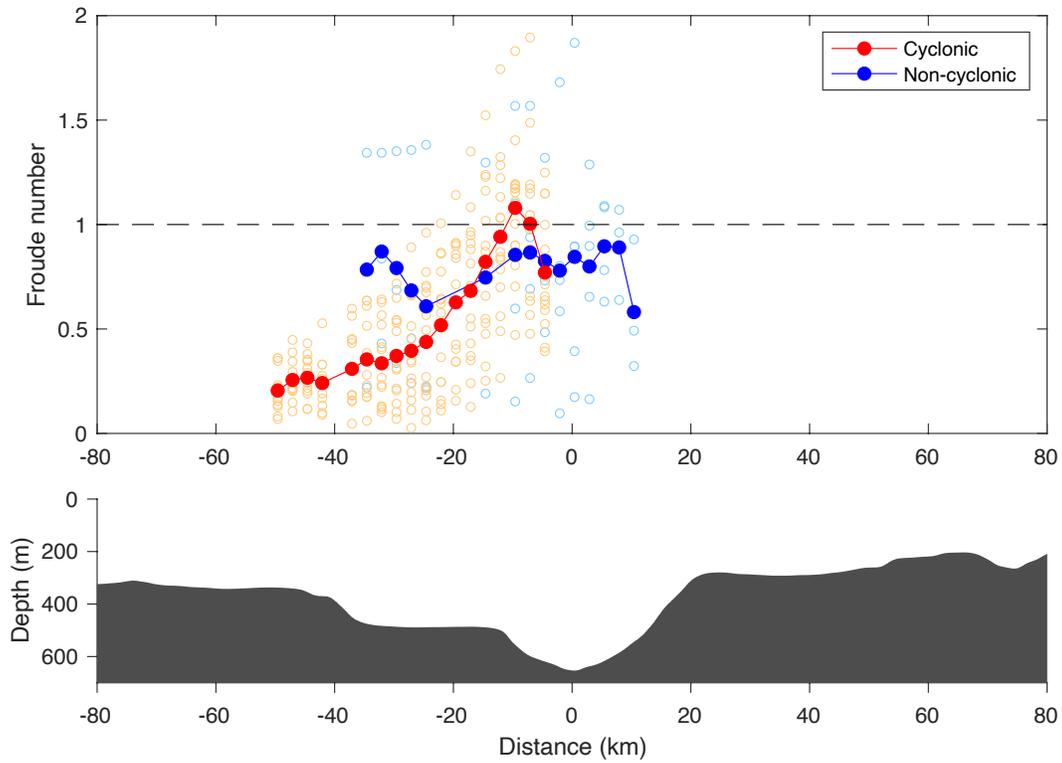
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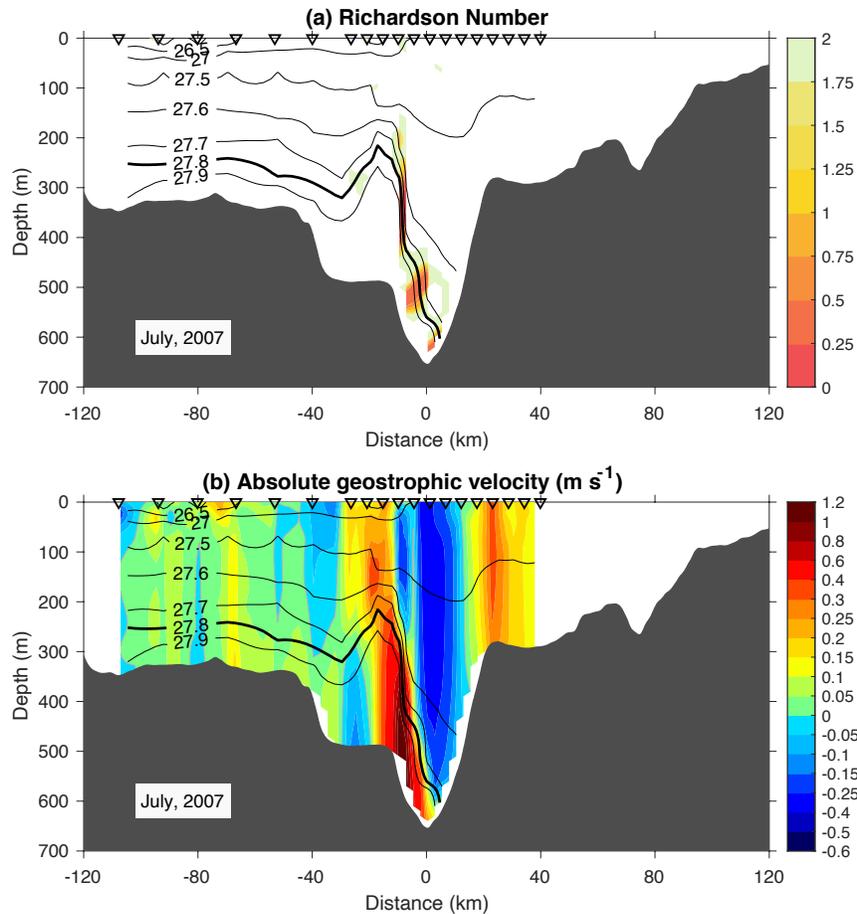
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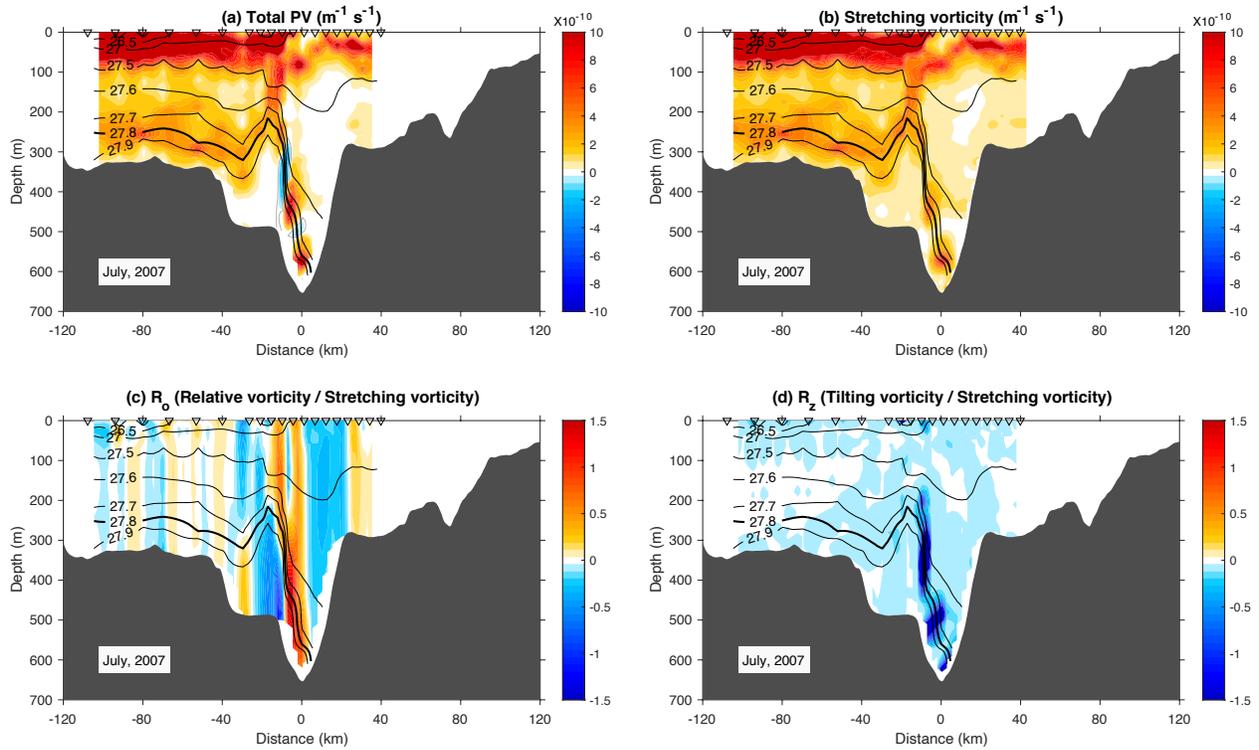
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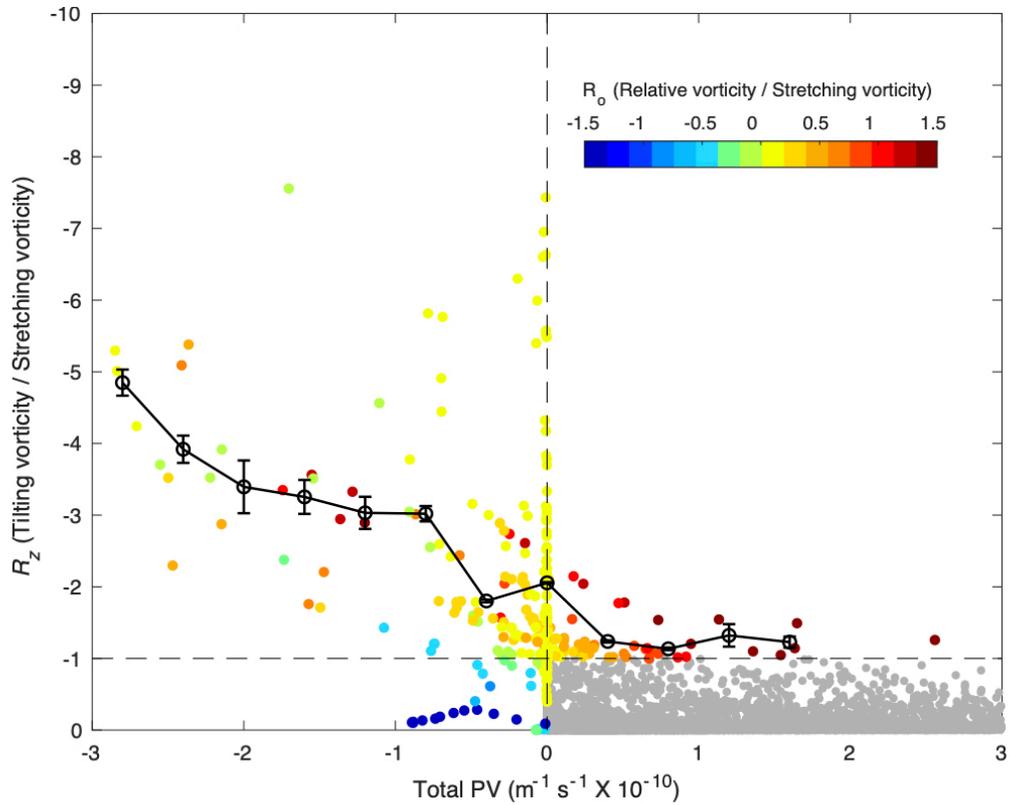
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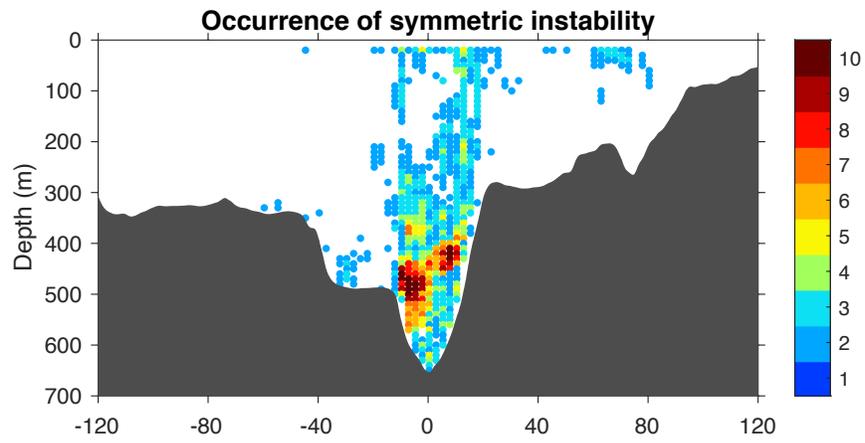
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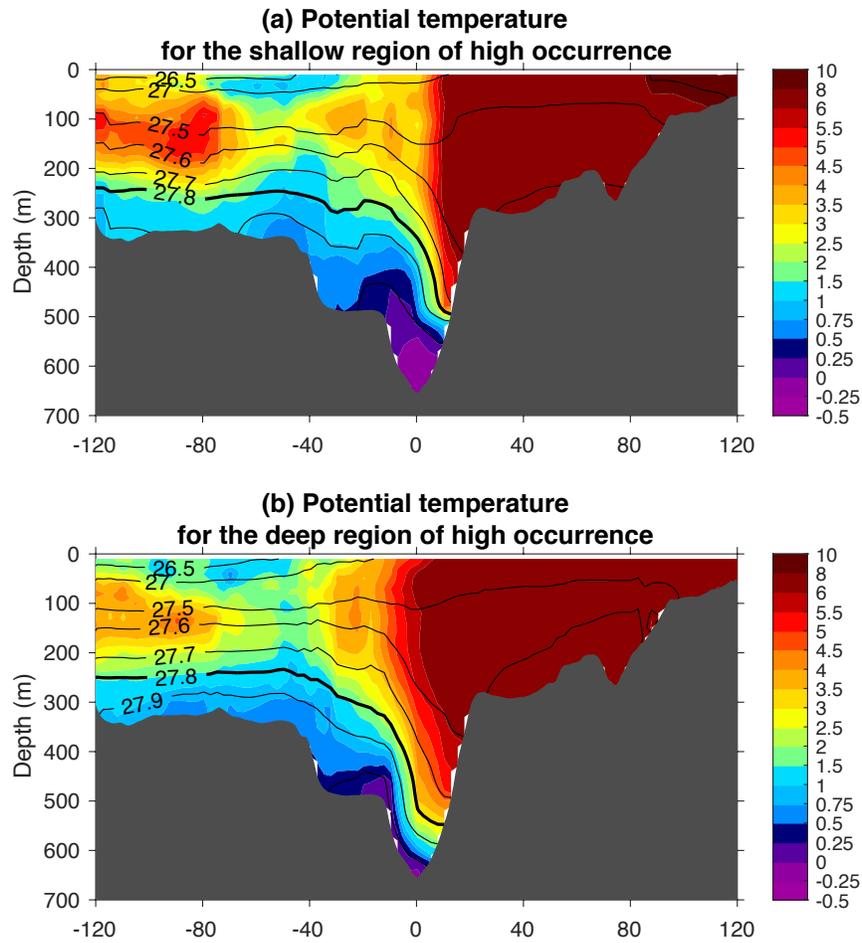


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