# Gromov-Hausdorff Limits of Kähler Manifolds with Ricci Curvature Bounded Below II 

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#### Abstract

We study noncollapsed Gromov-Hausdorff limits of Kähler manifolds with Ricci curvature bounded below. Our main result is that each tangent cone is homeomorphic to a normal affine variety. This extends a result of Donaldson-Sun, who considered noncollapsed limits of polarized Kähler manifolds with two-sided Ricci curvature bounds. © 2019 Wiley Periodicals, Inc.


## 1 Introduction

Consider a sequence $\left(M_{i}, \omega_{i}, p_{i}\right)$ of pointed complete Kähler manifolds of dimension $n$, with $\operatorname{Ric}\left(\omega_{i}\right)>-\omega_{i}$ and $\operatorname{Vol}\left(B\left(p_{i}, 1\right)\right)>v>0$. Suppose that the sequence converges to a metric space $(Z, d, p)$ in the pointed Gromov-Hausdorff sense. By the work of Cheeger-Colding [3], and more recently Cheeger-JiangNaber [6] and others, we have a detailed understanding of the structure of $Z$, even if the $M_{i}$ are merely Riemannian. A starting point for this structure theory is Cheeger-Colding's result [2] that the limit space $Z$ admits tangent cones at each point that are metric cones. In this paper we are interested in studying the additional structure of the tangent cones of $Z$ in the Kähler case.

There are few general results that exploit the Kähler condition: by Cheeger-Colding-Tian [5], the tangent cones are of the form $C(Y) \times \mathbf{C}^{k}$, where $C(Y)$ does not split off a factor of $\mathbf{R}$, while the first author [17] showed that each tangent cone admits a one-parameter group of isometries. Under two-sided Ricci curvature bounds it follows from Anderson [1] that the regular set in $Z$ is a Kähler manifold. In our previous work [18] we showed that under Ricci lower bounds, for sufficiently small $\epsilon>0$, the $\epsilon$-regular set $\mathcal{R}_{\epsilon} \subset Z$ has the structure of a complex manifold, although the metric may not be smooth. Here $x \in \mathcal{R}_{\epsilon}$ if $V_{2 n}-\lim _{r \rightarrow 0} r^{-2 n} \operatorname{vol}(B(x, r))<\epsilon$, where $V_{2 n}$ is the volume of the Euclidean unit ball. An important problem is whether the complex structure can be extended across the singular set of $Z$ to give it the structure of an analytic space. When the $\left(M_{i}, \omega_{i}\right)$ are polarized, i.e., the $\omega_{i}$ are curvature forms of line bundles, then

Donaldson-Sun [11] (under two-sided Ricci bounds) and the authors [18] (under just lower Ricci bounds) showed that this is the case. Without polarizations the question is still open; however, we have the following;

THEOREM 1.1. Every tangent cone of $Z$ is homeomorphic to a normal affine algebraic variety such that under a suitable embedding into $\mathbf{C}^{N}$, the homothetic action on the tangent cone extends to a linear torus action.

Theorem 1.1 was shown previously by Donaldson-Sun [12] under the assumptions that the $\omega_{i}$ are curvature forms of line bundles $L_{i} \rightarrow M_{i}$, and $\left|\operatorname{Ric}\left(\omega_{i}\right)\right|<1$. An important application of their result is that in their setting the holomorphic spectrum of the tangent cones is rigid, which in turn they used to show the uniqueness of tangent cones. While we are not able to show uniqueness, our result does imply the rigidity of the holomorphic spectrum under two-sided Ricci curvature bounds, even when the $\left(M_{i}, \omega_{i}\right)$ are not polarized.

More precisely, recall that for a Kähler cone $C(Y)$, possibly with singularities, the holomorphic spectrum is defined by

$$
\mathcal{S}=\{\operatorname{deg}(f): f \text { is homogeneous and holomorphic on } C(Y)\} \subset \mathbf{R} .
$$

We then have:

Corollary 1.2. Suppose that we have two-sided bounds $\left|\operatorname{Ric}\left(\omega_{i}\right)\right|<1$ along the sequence above. Then for any $q \in Z$ the holomorphic spectrum of every tangent cone at $q$ is the same. In addition, the volume ratio $V_{2 n}^{-1} \operatorname{Vol}(B(o, 1))$ is an algebraic number for every tangent cone $\left(Z_{q}, o\right)$.

As in [12], the rigidity of the holomorphic spectrum follows from the fact that the space of tangent cones at each point is connected, and the holomorphic spectrum consists of algebraic numbers. Note that these results hold in particular for tangent cones at infinity of Calabi-Yau manifolds with Euclidean volume growth.

The method of proof follows the overall strategy of Donaldson-Sun [11, 12] for constructing holomorphic functions on limit spaces. A crucial difference is that in our case the holomorphic functions on a tangent cone $Z_{q}$ are not obtained as limits of holomorphic functions on smooth manifolds. Instead we prove a version of Hörmander's $L^{2}$-estimate on the tangent cone; see Proposition 3.1. In the setting of two-sided Ricci curvature bounds there are substantial simplifications using that the singular set has codimension 4 , and the $L^{2}$-estimate can be proven directly on the tangent cone. We give this argument separately in Section 2, since it may be of independent interest. The case of lower Ricci bounds is treated in Section 3 by proving approximate versions of the basic estimate on smooth spaces converging to the tangent cone. Using this, in Section 4 we follow the strategy of DonaldsonSun [11, 12] to prove Theorem 1.1.

## 2 Two-Sided Ricci Bounds

In this section, we show that when our sequence $\left(M_{i}, \omega_{i}\right)$ has two-sided Ricci curvature bounds $\left|\operatorname{Ric}\left(\omega_{i}\right)\right|<1$, then we can directly prove the Hörmander $L^{2}$ estimate on the tangent cone. While the result is subsumed by the more general setting treated in Section 3, the proof may be of independent interest.

Let $X$ be a tangent cone of the noncollapsed Gromov-Hausdorff limit of a sequence of complete Kähler manifolds with two-sided Ricci curvature bounds. Let $S \subset X$ denote the singular set (see Cheeger-Colding [3] for details on the structure of $X$ ). The regular set $\mathcal{R}=X \backslash S$ admits a Ricci flat Kähler metric $\omega=\frac{1}{2} \sqrt{-1} \partial \bar{\partial} r^{2}$, where $r$ is the distance from the vertex of the cone $X$. Suppose that we have a function $\varphi$ on $\mathcal{R}$ satisfying $c \omega \leq \sqrt{-1} \partial \bar{\partial} \varphi \leq C \omega$ on $\mathcal{R}$ for a nonnegative function $c$ and constant $C>0$. More generally, $C$ could be a locally bounded function. For instance, we could use $\varphi=r^{2}$ or $\varphi=\log \left(1+r^{2}\right)$.

PROPOSITION 2.1. Suppose that $\alpha$ is a smooth ( 0,1 )-form compactly supported in $\mathcal{R}$ such that $\bar{\partial} \alpha=0$. Then there exists a function $f$ on $\mathcal{R}$ satisfying $\bar{\partial} f=\alpha$, and

$$
\int_{\mathcal{R}}|f|^{2} e^{-\varphi} \omega^{n} \leq \int_{\mathcal{R}} c^{-1}|\alpha|^{2} e^{-\varphi} \omega^{n}
$$

Proof. We follow the argument from Demailly [9, theorem 5.1] to prove Hörmander's $L^{2}$-estimate [14], the difference being that in our case we do not know that $\mathcal{R}$ admits a complete Kähler metric, and so we need a more careful argument for approximating $L^{2}$ forms with smooth forms of compact support.

As in Demailly's proof, our goal is to prove the inequality

$$
\begin{equation*}
\left|\int_{\mathcal{R}}\langle\alpha, v\rangle e^{-\varphi} \omega^{n}\right|^{2} \leq\left(\int_{\mathcal{R}} c^{-1}|\alpha|^{2} e^{-\varphi} \omega^{n}\right)\left\|\bar{\partial}_{\varphi}^{*} v\right\|^{2} \tag{2.1}
\end{equation*}
$$

for all smooth ( 0,1 )-forms $v$ with compact support in $\mathcal{R}$. To define $\bar{\partial}_{\varphi}^{*} v$ we are viewing $v$ as a $(0,1)$-form valued in the trivial bundle with metric $e^{-\varphi}$. The existence of the required function $f$ then follows from the Riesz representation theorem.

Given such a smooth ( 0,1 )-form $v$ compactly supported away from $S$, we can decompose $v=v_{1}+v_{2}$ under the $L^{2}$-orthogonal decomposition

$$
L^{2}=\operatorname{ker} \bar{\partial} \oplus(\operatorname{ker} \bar{\partial})^{\perp}
$$

Note that $\bar{\partial}_{\varphi}^{*} v_{2}=0$, and also $\bar{\partial}_{\varphi}^{*} v=0$ near $S$, and therefore $\bar{\partial}_{\varphi}^{*} v_{1}=0$ near $S$. Since also by definition $\bar{\partial} v_{1}=0$, it follows that $v_{1}$ is a harmonic ( 0,1 )-form near $S$. If $u$ is any harmonic $(0,1)$-form valued in a line bundle $L$ with curvature $F_{j \bar{k}}$,
in an orthonormal frame we have

$$
\begin{aligned}
0=\Delta_{\bar{\partial}} u & =\left(\bar{\partial}^{*} \bar{\partial}+\overline{\partial \partial}^{*}\right) u \\
& =\sum_{j}-\nabla_{j} \nabla_{\bar{j}} u_{\bar{k}}+\nabla_{j} \nabla_{\bar{k}} u_{\bar{j}}-\nabla_{\bar{k}} \nabla_{j} u_{\bar{j}} \\
& =\sum_{j}-\nabla_{j} \nabla_{\bar{j}} u_{\bar{k}}+\sum_{j, p} R_{\bar{j} j \bar{k}}^{\bar{p}} u_{\bar{p}}+\sum_{j} F_{j \bar{k}} u_{\bar{j}},
\end{aligned}
$$

and so

$$
\sum_{j} \nabla_{j} \nabla_{\bar{j}} u_{\bar{k}}=\sum_{j, p} R_{\bar{j}_{j} \bar{k}}^{\bar{p}_{\bar{p}}} u_{\bar{j}} F_{j \bar{k}} u_{\bar{j}} .
$$

We also have

$$
\begin{aligned}
\sum_{j} \nabla_{\bar{j}} \nabla_{j} u_{\bar{k}} & =\sum_{j} \nabla_{j} \nabla_{\bar{j}} u_{\bar{k}}-\sum_{j, p} R_{\bar{k} j \bar{j}}^{\bar{p}} u_{\bar{p}}-\sum_{j} F_{j \bar{j}} u_{\bar{k}} \\
& =\sum_{j} F_{j \bar{k}} u_{\bar{j}}-F_{j \bar{j}} u_{\bar{k}}
\end{aligned}
$$

It follows that

$$
\begin{align*}
\Delta|u|^{2}= & \sum_{j} \nabla_{j} \nabla_{\bar{j}}|u|^{2} \\
= & \sum_{j, k}\left(\nabla_{j} \nabla_{\bar{j}} u \bar{k}\right) \overline{u_{\bar{k}}}+\left|\nabla^{1,0} u\right|^{2}+\left|\nabla^{0,1} u\right|^{2}+\sum_{j, k} u_{\bar{k}} \overline{\nabla_{\bar{j}} \nabla_{j} u_{\bar{k}}} \\
= & |\nabla u|^{2}+\sum_{k, p} R_{p \bar{k}} u_{\bar{p}} \overline{u_{\bar{k}}}+\sum_{j, k} F_{j \bar{k}} u \bar{j} \overline{u_{\bar{k}}}+\sum_{j, k} \overline{F_{j \bar{k}}} u_{\bar{k}} \overline{u_{\bar{j}}}  \tag{2.2}\\
& -\sum_{j, k} \overline{F_{j \bar{j}}} u \bar{k} \overline{u_{\bar{k}}} \\
= & |\nabla u|^{2}+\sum_{p, k} R_{p \bar{k}} u \bar{p} \overline{u_{\bar{k}}}+\sum_{j, k} 2 F_{j \bar{k}} u \overline{\bar{j}} \overline{u_{\bar{k}}}-\sum_{j} F_{j \bar{j}}|u|^{2} .
\end{align*}
$$

In our setting the Ricci curvature vanishes, and $F_{j \bar{k}}=\partial_{j} \partial_{\bar{k}} \varphi$. In particular, in a neighborhood of $S$ we have

$$
\Delta\left|v_{1}\right|^{2} \geq\left|\nabla v_{1}\right|^{2}-C n\left|v_{1}\right|^{2}
$$

It follows that $\Delta\left|v_{1}\right| \geq-C\left|v_{1}\right|$ near $S$ for a locally bounded function $C$. Since $\left|v_{1}\right| \in L^{2}(\mathcal{R})$, we can apply Lemma 2.2 below to obtain that $\left|v_{1}\right| \in L_{\text {loc }}^{\infty}$ on a neighborhood of $S$ in $X$ (i.e., $\left|v_{1}\right|$ cannot blow up as we approach the singular set).

Let us now define cutoff functions $\eta_{R}$ and $\chi_{\epsilon}$ as follows. The function $\eta_{R}$ equals 1 in $B_{R}(0)$, vanishes outside of $B_{R+1}(0)$, and satisfies $\left|\nabla \eta_{R}\right|<2$. The function $\chi_{\epsilon}$ equals 1 outside of the $2 \epsilon$-neighborhood of $S$ and vanishes in the $\epsilon$-neighborhood of $S$. In addition, since $S$ has codimension 4 (see [5]), we can arrange that on
compact sets $K$ we have $\left\|\nabla \chi_{\epsilon}\right\|_{L^{2}(K)} \rightarrow 0$ as $\epsilon \rightarrow 0$. Since $\chi_{\epsilon} \eta_{R} v_{1}$ has compact support in $\mathcal{R}$, by the Bochner-Kodaira inequality,

$$
\int_{\mathcal{R}} c\left|\chi_{\epsilon} \eta_{R} v_{1}\right|^{2} e^{-\varphi} \omega^{n} \leq\left\|\bar{\partial}\left(\chi_{\epsilon} \eta_{R} v_{1}\right)\right\|^{2}+\left\|\bar{\partial}_{\varphi}^{*}\left(\chi_{\epsilon} \eta_{R} v_{1}\right)\right\|^{2},
$$

and it follows using the Cauchy-Schwarz inequality that

$$
\begin{align*}
& \left|\int_{\mathcal{R}}\left\langle\alpha, \chi_{\epsilon} \eta_{R} v_{1}\right\rangle e^{-\varphi} \omega^{n}\right|^{2} \\
& \quad \leq \int_{\mathcal{R}} c^{-1}|\alpha|^{2} e^{-\varphi} \omega^{n} \int_{\mathcal{R}} c\left|\chi_{\epsilon} \eta_{R} v_{1}\right|^{2} e^{-\varphi} \omega^{n}  \tag{2.3}\\
& \quad \leq\left(\int_{\mathcal{R}} c^{-1}|\alpha|^{2} e^{-\varphi} \omega^{n}\right)\left(\left\|\bar{\partial}\left(\chi_{\epsilon} \eta_{R} v_{1}\right)\right\|^{2}+\left\|\bar{\partial}_{\varphi}^{*}\left(\chi_{\epsilon} \eta_{R} v_{1}\right)\right\|^{2}\right) .
\end{align*}
$$

Now recall that $v_{1}$ is locally bounded near $S$, and $\left\|\nabla \chi_{\epsilon}\right\|_{L^{2}}$ can be made arbitrarily small. It follows that for fixed $R$ if we let $\epsilon \rightarrow 0$, we have $\chi_{\epsilon} \eta_{R} v_{1} \rightarrow \eta_{R} v_{1}$, $\bar{\partial}\left(\chi_{\epsilon} \eta_{R} v_{1}\right) \rightarrow \bar{\partial}\left(\eta_{R} v_{1}\right)$, and $\bar{\partial}_{\varphi}^{*}\left(\chi_{\epsilon} \eta_{R} v_{1}\right) \rightarrow \bar{\partial}_{\varphi}^{*}\left(\eta_{R} v_{1}\right)$ in $L^{2}$. In addition, using that $\left|\nabla \eta_{R}\right|<2$ and that the supports of $\eta_{R}$ exhaust $\mathcal{R}$ as $R \rightarrow \infty$, we have that $\eta_{R} v_{1} \rightarrow v_{1}, \bar{\partial}\left(\eta_{R} v_{1}\right) \rightarrow \bar{\partial} v_{1}$, and $\bar{\partial}^{*}\left(\eta_{R} v_{1}\right) \rightarrow \bar{\partial}^{*} v_{1}$ in $L^{2}$ as $R \rightarrow \infty$. It follows from (2.3) that

$$
\left|\int_{\mathcal{R}}\left\langle\alpha, v_{1}\right\rangle e^{-\varphi} \omega^{n}\right|^{2} \leq\left(\int_{\mathcal{R}} c^{-1}|\alpha|^{2} e^{-\varphi} \omega^{n}\right)\left(\left\|\bar{\partial} v_{1}\right\|^{2}+\left\|\bar{\partial}_{\varphi}^{*} v_{1}\right\|^{2}\right) .
$$

The required inequality (2.1) now follows since the assumption $\bar{\partial} \alpha=0$ implies that $\alpha$ is orthogonal to $v_{2}$, and at the same time $\bar{\partial} v_{1}=0$ and $\bar{\partial}_{\varphi}^{*} v_{1}=\bar{\partial}_{\varphi}^{*} v$.

We used the following lemma in the proof above. Note that this estimate fails when the singular set has codimension 2.

Lemma 2.2. Let $B$ be a unit ball in $X$, and suppose that $u \in L^{2}(B)$ is such that on $B \backslash S$ the function $u$ is smooth, nonnegative, and $\Delta u \geq-A u$ for a constant $A>0$. Then we have

$$
\sup _{\frac{1}{2} B \backslash S} u \leq C\|u\|_{L^{2}(B)}
$$

for a constant $C$ depending on $A$, the dimension, and asymptotic volume ratio of $X$.

Proof. The function $\tilde{u}(s, x)=e^{\sqrt{A} s} u(x)$ on the product $\mathbf{R} \times B$ satisfies $\Delta_{\mathbf{R} \times X} \tilde{u}=A e^{\sqrt{A} s} u+e^{\sqrt{A} s} \Delta_{X} u \geq 0$, and the $L^{2}$-norm of $\tilde{u}$ on $[-1,1] \times B$ can be bounded in terms of the $L^{2}$-norm of $u$ on $B$ and the constant $A$. Using this we can reduce to the case $A=0$.

Given $y \in \frac{1}{2} B \backslash S$, let $H_{t}(x, y)$ be the heat kernel on $X$ satisfying $\partial_{t} H_{t}=$ $\Delta_{x} H_{t}$ (see Ding [10] for details on the heat kernel on tangent cones). In addition,
let $\eta$ be a cutoff function such that $\eta=1$ on $\frac{2}{3} B$ and $\eta=0$ outside of $\frac{3}{4} B$. We have

$$
\begin{align*}
& \partial_{t} \int_{B} u(x) \eta(x) H_{t}(x, y) d x \\
&=\int_{B} u \eta \Delta H_{t} \\
&=\int_{B} u \Delta\left(\eta H_{t}\right)-\int_{B}\left[u H_{t} \Delta \eta+2 u \nabla \eta \cdot \nabla H_{t}\right]  \tag{2.4}\\
& \quad \geq \int_{B} u \Delta\left(\eta H_{t}\right)-C\|u\|_{L^{2}}
\end{align*}
$$

Here we used that $\Delta \eta \in L^{2}$, and $\nabla H_{t}$ is bounded independently of $t$ on the support of $\nabla \eta$.

We claim that the first term is nonnegative for any $t>0$. For this let $\chi_{\epsilon}$ be cutoff functions such that $1-\chi_{\epsilon}$ is supported in the $\epsilon$-neighborhood of $S$, and in addition $\nabla \chi_{\epsilon}, \Delta \chi_{\epsilon} \in L^{2}$. Such a choice is possible because of remark 1.15 on page 6 of [15]. For a fixed $t>0$ let us write $\psi=\eta H_{t}$. Then

$$
\int_{B} \chi_{\epsilon} u \Delta \psi=\int_{B} u \Delta\left(\chi_{\epsilon} \psi\right)-\int_{B}\left[u \psi \Delta \chi_{\epsilon}+2 u \nabla \psi \cdot \nabla \chi_{\epsilon}\right]
$$

Since now $\chi_{\epsilon} \psi \geq 0$ is compactly supported in $B \backslash S$, the first term is nonnegative, while at the same time the last two terms tend to 0 as $\epsilon \rightarrow 0$. It follows that

$$
\int_{B} u \Delta\left(\eta H_{t}\right) \geq 0
$$

and so (2.4) implies

$$
\int_{B} u \eta H_{1} \geq \lim _{t \rightarrow 0} \int_{B} u \eta H_{t}-C\|u\|_{L^{2}}=u(y)-C\|u\|_{L^{2}}
$$

Since $H_{1}$ is bounded above, we obtain the required estimate for $u(y)$.

## 3 The $L^{\mathbf{2}}$ Estimate on Tangent Cones

Suppose now that $(X, o)$ is a tangent cone at a point of a noncollapsed limit space of $n$-dimensional Kähler manifolds with lower bounds on the Ricci curvature. In particular, for all $R>0$, the ball $B(o, R) \subset X$ is the Gromov-Hausdorff limit of a sequence of balls $B\left(p_{i}, R\right)$ in Kähler manifolds with Ric $>-i^{-1}$, and $\operatorname{vol}\left(B\left(p_{i}, R\right)\right)>\nu R^{2 n}$ for the noncollapsing constant $v>0$, with $p_{i} \rightarrow o$. For sufficiently small $\epsilon$ we know from [18] that the $\epsilon$-regular set $\mathcal{R}_{\epsilon}$ is a complex manifold, and from Cheeger-Jiang-Naber [6] that $X \backslash \mathcal{R}_{\epsilon}$ is $(2 n-2)$-rectifiable with locally finite $(2 n-2)$-dimensional Minkowski content. From now on we choose an $\epsilon=\epsilon(n)$ sufficiently small, but fixed.

For $t \geq 0$, let $w(t)$ satisfy

$$
\begin{equation*}
w^{\prime}>0, \quad w^{\prime \prime} \leq 0, \quad \text { and } \quad w^{\prime}+t w^{\prime \prime}>0 \tag{3.1}
\end{equation*}
$$

For instance, we can take $w(t)=C \log (t+1)$ for any constant $C>0$. In the $L^{2}$-estimate we will use a weight function of the form $\varphi=w\left(r^{2}\right)$, where $r$ is the distance function from the vertex of the cone $X$. We denote by $\mu$ the natural measure on $X$. It will also be useful to denote by $\mu_{0}$ a smooth volume measure on $\mathcal{R}_{\epsilon}$ with respect to the holomorphic charts (up to bounded factors). It follows that on compact sets we have a lower bound $d \mu>C^{-1} d \mu_{0}$.
PROPOSITION 3.1. Suppose that $\alpha=\bar{\partial} h$, where $h$ is a smooth function compactly supported in $\mathcal{R}_{\epsilon}$. Then there exists a function $f$ on $\mathcal{R}_{\epsilon}$ satisfying $\bar{\partial} f=\alpha$, and

$$
\int_{\mathcal{R}_{\epsilon}}|f|^{2} e^{-\varphi} d \mu \leq \int_{\mathcal{R}_{\epsilon}} \frac{|\alpha|^{2}}{w^{\prime}\left(r^{2}\right)+r^{2} w^{\prime \prime}\left(r^{2}\right)} e^{-\varphi} d \mu
$$

Remark 3.2. It will be clear from the proof that we can add weight functions with logarithmic poles on $\mathcal{R}_{\epsilon}$. This will be useful to separate tangents in Proposition 4.1.

Remark 3.3. We must restrict ourselves to exact forms $\alpha$ for technical reasons, but this will not matter in our application. We will define the norm $|\alpha|$ precisely below by taking a limit of the corresponding norms on smooth spaces.

As in the proof of Proposition 2.1, we follow the approach in Demailly [9] to Hörmander's $L^{2}$-estimate, but we need to take care with working on the space $\mathcal{R}_{\epsilon}$ since the metric we are using is not smooth. We construct $f$ satisfying the equation $\bar{\partial} f=\alpha$ in a weak sense, i.e., satisfying

$$
\begin{equation*}
\int_{\mathcal{R}_{\epsilon}} f \bar{\partial} \eta=\int_{\mathcal{R}_{\epsilon}} h \bar{\partial} \eta \tag{3.2}
\end{equation*}
$$

for all smooth ( $n, n-1$ )-forms $\eta$ with compact support in $\mathcal{R}_{\epsilon}$. In order to construct $f$ as an element of $L^{2}\left(e^{-\varphi} d \mu\right)$, we need to define the function $* \bar{\partial} \eta \in$ $L^{2}\left(e^{-\varphi} d \mu\right)$. For this, as well as for the definition of the norm $|\alpha|$, we observe that in any holomorphic chart on $\mathcal{R}_{\epsilon}$ we have a well-defined tensor $g^{j \bar{k}}$ with bounded measurable components, corresponding to the inverse of the metric. More precisely, we have the following:

Lemma 3.4. Suppose that $q_{i} \in B\left(p_{i}, R\right)$ satisfy $q_{i} \rightarrow q \in \mathcal{R}_{\epsilon}$, and $z_{i j}$ are holomorphic charts on small balls $B\left(q_{i}, \rho\right)$ converging to a holomorphic chart $z_{j}$ on $B(q, \rho)$. We use the charts to identify functions on $B\left(q_{i}, \rho\right)$ with functions on $B(q, \rho)$. Then the inverses $g_{i}^{j \bar{k}}$ of the metric tensors on $B\left(q_{i}, \rho\right)$ converge in $L^{p}(d \mu)$ for all $p$ to a tensor $g^{j \bar{k}}$ on $B(q, \rho)$ with bounded components.

Proof. We show that we have

$$
\int_{B(q, \rho)}\left|g_{a}^{j \bar{k}}-g_{b}^{j \bar{k}}\right| d \mu \rightarrow 0
$$

as $a, b \rightarrow \infty$. Note first that we have a uniform upper bound $\left|g_{a}^{j \bar{k}}\right|<C$. Using that the regular points in $\mathcal{R}_{\epsilon}$ have full measure, together with a covering argument, it is enough to show that for all regular $x \in B(q, \rho)$ we have

$$
\begin{equation*}
\lim _{\kappa \rightarrow 0} \lim _{a, b \rightarrow \infty} f_{B(x, \kappa)}\left|g_{a}^{j \bar{k}}-g_{b}^{j \bar{k}}\right| d \mu=0 . \tag{3.3}
\end{equation*}
$$

Consider a ball $B(x, \kappa)$ and $x_{i} \in B\left(p_{i}, R\right)$ such that $x_{i} \rightarrow x$. We can find holomorphic coordinates $w_{i p}$ on $B\left(x_{i}, \kappa\right)$ converging to $w_{p}$ on $B(x, \kappa)$, which give Gromov-Hausdorff approximations to the corresponding Euclidean balls. Let us rescale distances by $\kappa^{-1}$. The Cheeger-Colding estimate [3] shows that the corresponding metric components $g_{i}^{p \bar{q}}$ satisfy

$$
\begin{equation*}
f_{\kappa^{-1} B\left(x_{i}, \kappa\right)}\left|g_{i}^{p \bar{q}}-\delta^{p \bar{q}}\right| d \mu_{i}<\Psi\left(\kappa, i^{-1} \mid x\right) \tag{3.4}
\end{equation*}
$$

Here, and below, $\Psi\left(\delta_{1}, \ldots, \delta_{k} \mid a_{1}, \ldots, a_{l}\right)$ denotes a function converging to 0 as $\delta_{i} \rightarrow 0$ while the $a_{j}$ are fixed.

In addition, we have $g_{i, p \bar{q}}>\left(1-\Psi\left(\kappa, i^{-1} \mid x\right)\right) \delta_{p \bar{q}}$. It follows from this and Colding's volume convergence [8] that in (3.4) we can replace the measure $d \mu_{i}$ with $d \mu$. The gradient estimate for the coordinates $z_{j}$ implies that the Jacobian matrix $d z_{j} / d w_{p}$ is uniformly bounded. It follows that

$$
f_{\kappa^{-1} B\left(x_{i}, \kappa\right)}\left|g_{i}^{j \bar{k}}-\left(d z_{j} / d w_{p}\right) \overline{\left(d z_{k} / d w_{q}\right)} \delta^{\bar{q}}\right| d \mu<\Psi\left(\kappa, i^{-1} \mid x\right)
$$

which implies (3.3). The sequence $g_{a}^{j \bar{k}}$ therefore converges to a limit $g^{j \bar{k}}$ in $L^{1}(d \mu)$, and since these components are uniformly bounded, the convergence is in $L^{p}$ for all $p$ as well.

If we have another set of charts $z_{i j}^{\prime}$ converging to $z_{j}^{\prime}$, then the corresponding components $g^{\prime j} \bar{k}$ are related to $g^{j \bar{k}}$ in the usual way. This follows from the fact that the transition functions $z_{i j} \circ z_{i j}^{\prime-1}$ converge to $z_{j} \circ z_{j}^{\prime-1}$ as $i \rightarrow \infty$.

In terms of the metric components $g^{j \bar{k}}$ we can now define $|\alpha|^{2}=g^{j \bar{k}} \alpha_{\bar{k}} \overline{\bar{j}}$ in the usual way. Note that if we have local coordinates $z_{i j}$ converging to $z_{j}$ on $B(q, \rho)$ as above such that $\varphi_{i}, \rho_{i}$ converge uniformly to $\varphi, r^{2}$, then we have

$$
\lim _{i \rightarrow \infty} \int_{B\left(q_{i}, \rho\right)} \frac{|\alpha|_{g_{i}}^{2} e^{-\varphi_{i}}}{w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)} d \mu_{i}=\int_{B(q, \rho)} \frac{|\alpha|^{2} e^{-\varphi}}{w^{\prime}\left(r^{2}\right)+r^{2} w^{\prime \prime}\left(r^{2}\right)} d \mu .
$$

This follows from the convergence of the metric components $g_{i}^{j \bar{k}}$ to $g^{j \bar{k}}$, together with the convergence [8] of the measures $\mu_{i}$ to $\mu$.

Similarly, we can define $* \bar{\partial} \eta$ in local coordinates by letting

$$
* \overline{\bar{\partial} \eta}=e^{\varphi}(\bar{\partial} \eta)_{1 \overline{1} \cdots n \bar{n}} \operatorname{det}\left(g^{j \bar{k}}\right)
$$

In terms of coordinates $z_{i j}$ converging to $z_{j}$ as above, we then have that $*_{\varphi_{i}} \bar{\partial} \eta$ converges in $L^{p}$ to $* \bar{\partial} \eta$, and in addition we have

$$
\int_{\mathcal{R}_{\epsilon}} f \overline{* \bar{\partial} \eta} e^{-\varphi} d \mu=\int_{\mathcal{R}_{\epsilon}} f \bar{\partial} \eta
$$

for any $f \in L^{2}\left(e^{-\varphi} d \mu\right)$. Note that such an $f$ is also locally in $L^{2}$ with respect to $d \mu_{0}$ since $\varphi$ is locally bounded, and so the right hand side is well defined.

PROOF OF PROPOSITION 3.1. Given a smooth ( $n, n-1$ )-form $\eta$ with compact support in $\mathcal{R}_{\epsilon}$, our goal is to prove the inequality

$$
\begin{align*}
\left(\int_{\mathcal{R}_{\epsilon}} h * \overline{\bar{\partial} \eta} e^{-\varphi} d \mu\right)^{2} & =\left(\int_{\mathcal{R}_{\epsilon}} h \bar{\partial} \eta\right)^{2}  \tag{3.5}\\
& \leq \int_{\mathcal{R}_{\epsilon}} \frac{|\alpha|^{2} e^{-\varphi}}{w^{\prime}\left(r^{2}\right)+r^{2} w^{\prime \prime}\left(r^{2}\right)} d \mu \int_{\mathcal{R}_{\epsilon}}|* \bar{\partial} \eta|^{2} e^{-\varphi} d \mu
\end{align*}
$$

The existence of the required function $f$ then follows in the standard way using the Hahn-Banach and Riesz representation theorems. Indeed, if we denote by $E \subset$ $L^{2}\left(e^{-\varphi} d \mu\right)$ the closure of the subspace of functions of the form $* \bar{\partial} \eta$ with compact support, then (3.5) implies that $h$ defines a linear functional $\lambda: E \rightarrow \mathbf{C}$ with norm

$$
\|\lambda\| \leq\left(\int_{\mathcal{R}_{\epsilon}} \frac{|\alpha|^{2} e^{-\varphi}}{w^{\prime}\left(r^{2}\right)+r^{2} w^{\prime \prime}\left(r^{2}\right)} d \mu\right)^{1 / 2}
$$

We then deduce the existence of $f \in E \subset L^{2}\left(e^{-\varphi} d \mu\right)$ with the same bound on its norm such that (3.2) holds for all smooth $\eta$ with compact support. It then follows that $\bar{\partial} f=\bar{\partial} h$.

We prove the inequality (3.5) using approximations by smooth spaces. Let $\mathcal{U} \subset$ $\mathcal{R}_{\epsilon}$ be an open relatively compact subset containing the closure of the support of $\eta$. Let us fix a large radius $R>0$ such that $\mathcal{U} \subset B(o, R / 2)$. Recall that $B(o, R)$ is a noncollapsed Gromov-Hausdorff limit of a sequence of balls $B\left(p_{i}, R\right)$ in Kähler manifolds with Ric $>-i^{-1}$. For every $q \in B(o, R) \cap \mathcal{U}$, by definition we have a radius $\delta$ such that $B(q, \delta)$ is $\Psi(\epsilon)$-Gromov-Hausdorff close to the Euclidean ball.

Suppose that $q_{i} \in B\left(p_{i}, R\right)$ are such that $q_{i} \rightarrow q$. By theorem 2.1 in [18], we have $\delta^{\prime}$ such that on each $B\left(q_{i}, \delta^{\prime}\right)$ there is a holomorphic chart giving a GromovHausdorff approximation to the Euclidean ball, and a Kähler potential close to $\frac{1}{2} d\left(q_{i}, \cdot\right)^{2}$. These holomorphic charts converge, as $i \rightarrow \infty$, to a holomorphic chart on $B\left(q, \delta^{\prime}\right)$ defining the holomorphic structure on $\mathcal{U}$.

We first construct suitable weight functions $\varphi_{i}$ on $B\left(p_{i}, R\right)$. Since $B\left(p_{i}, R\right) \rightarrow$ $B(o, R)$ in the Gromov-Hausdorff sense, by Cheeger-Colding [3] we have functions $b_{i}$ on $B\left(p_{i}, R\right)$ such that $\Delta b_{i}^{2} / 2=n,\left|\nabla b_{i}\right|<1+\Psi\left(i^{-1}\right)$ on $B\left(p_{i}, R-1\right)$,

$$
\begin{equation*}
\int_{B\left(p_{i}, R\right)}\left|\sqrt{-1} \partial \bar{\partial} b_{i}^{2} / 2-g_{i}\right|^{2}+\left|\left|\nabla b_{i}\right|-1\right|^{2}<\Psi(1 / i) \tag{3.6}
\end{equation*}
$$

and $\left|b_{i}-r_{i}\right|<\Psi(1 / i)$, in terms of the distance $r_{i}$ from $p_{i}$. In other words, $b_{i}^{2} / 2$ is plurisubharmonic in an $L^{2}$ sense, but this is not good enough to get an inequality of the form (3.5). We will therefore modify $b_{i}^{2} / 2$ on the $\operatorname{set} \mathcal{U}$ to make it plurisubharmonic in a pointwise sense there.

Let us cover $\mathcal{U} \subset B(o, R / 2)$ using charts $U_{j}$ such that the smaller charts $\frac{1}{10} U_{j}$ still cover. Under the Gromov-Hausdorff convergence these define charts $U_{i j}$ centered at points $q_{i j} \in B\left(p_{i}, R\right)$. On each $U_{i j}$ there is a Kähler potential $\psi_{i j}$, close to the function $\frac{1}{2} d\left(q_{i j}, \cdot\right)^{2}$, the closeness determined by how close to the Euclidean ball our chart is in the Gromov-Hausdorff sense (this is controlled by our choice of $\epsilon(n)$ above).

At the same time, on each $U_{i j}$ we also have Kähler potentials $\varphi_{i j}$, which are close to $\frac{1}{2} r_{i}^{2}$ (where $r_{i}$ is the distance from $p_{i}$ in the ball $B\left(p_{i}, R\right)$ ) in the sense that

$$
\left|\varphi_{i j}-r_{i}^{2} / 2\right|<\Psi(1 / i)
$$

where $r_{i}=d\left(p_{i}, \cdot\right)$ as above. This is because if we let $\psi_{\infty j}$ be the limit of the potentials $\psi_{i j}$ as $i \rightarrow \infty$ along a subsequence, then $\partial \bar{\partial}\left(\psi_{\infty j}-r^{2} / 2\right)=0$ on $U_{j}$ (see [18, claim 3.1] for a similar result). We can then define $\varphi_{i j}=\psi_{i j}+\left(r^{2} / 2-\right.$ $\psi_{\infty j}$ ) using the local coordinates to identify $U_{i j}$ with $U_{j}$.

It is clear that on the overlap $U_{i j} \cap U_{i k},\left|\varphi_{i j}-\varphi_{i k}\right|<\Psi(1 / i)$. We can ensure that on a smaller compact set of the intersection, $\left|\nabla\left(\varphi_{i j}-\varphi_{i k}\right)\right|$ is small. Using a partition of unity on $\bigcup_{j} \frac{1}{2} U_{i j}$, we can glue the $\varphi_{i j}$ to obtain a function $\tilde{\rho}_{i}$ that satisfies

$$
\begin{array}{r}
\left|\tilde{\rho}_{i}-\frac{r_{i}^{2}}{2}\right|+\left|\partial \bar{\partial} \tilde{\rho}_{i}-g_{i}\right|<\Psi(1 / i)  \tag{3.7}\\
\left|\nabla \tilde{\rho}_{i}\right| \leq r_{i}+\Psi(1 / i),\left|\nabla\left(\widetilde{\rho}_{i}-b_{i}^{2} / 2\right)\right|<\Psi(1 / i)
\end{array}
$$

on $\bigcup_{j} \frac{1}{3} U_{i j}$. Note that since the partition of unity is defined using the charts, the functions giving the partition of unity have uniformly bounded derivatives.

Similarly, we can define a cutoff function $\mu_{i}$, supported on $\bigcup_{j} \frac{1}{3} U_{i j}$, so that $\mu_{i}=1$ on $\bigcup_{j} \frac{1}{6} U_{i j}$. Furthermore,

$$
\begin{equation*}
\left|\nabla \mu_{i}\right|+\left|\partial \bar{\partial} \mu_{i}\right| \leq C \tag{3.8}
\end{equation*}
$$

where $C$ is independent of $i$. Now we define

$$
\begin{equation*}
\rho_{i}=\mu_{i} \tilde{\rho}_{i}+\left(1-\mu_{i}\right) b_{i}+\epsilon_{i} \tag{3.9}
\end{equation*}
$$

where $\epsilon_{i}$ is a positive sequence converging to 0 so that $\rho_{i} \geq 0$. Set

$$
\begin{equation*}
\varphi_{i}=w\left(\rho_{i}\right) \tag{3.10}
\end{equation*}
$$

From (3.7), (3.8), (3.9), (3.10), and standard computation, we find that on $\bigcup_{j} \frac{1}{6} U_{i j}$

$$
\begin{equation*}
\sqrt{-1} \partial \bar{\partial} \varphi_{i} \geq\left(w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)-\Psi(1 / i)\right) g_{i} \tag{3.11}
\end{equation*}
$$

and we also have

$$
\begin{equation*}
\int_{B\left(p_{i}, R-1\right)}\left|\sqrt{-1} \partial \bar{\partial} \varphi_{i}-\left(w^{\prime}\left(\rho_{i}\right) g_{i}+\sqrt{-1} w^{\prime \prime}\left(\rho_{i}\right) \partial \rho_{i} \wedge \bar{\partial} \rho_{i}\right)\right|^{2} \leq \Psi(1 / i) \tag{3.12}
\end{equation*}
$$

Let us consider now the ( $n, n-1$ )-form $\eta$. Let $\gamma_{j}$ be a partition of unity subordinate to the cover $U_{j}$ of $\mathcal{U}$. We can view each $\gamma_{j} \eta$ as an $(n, n-1)$ form on $U_{i j} \subset B\left(p_{i}, R\right)$, and define the form $\eta_{i}$ as their sum. Define $v_{i}=*_{\varphi_{i}} \eta_{i}$, where we are viewing $\eta_{i}$ as an $(n, n-1)$-form with values in the trivial bundle $L^{*}$ with metric $e^{\varphi_{i}}$ (i.e., the dual of the trivial bundle $L$ with metric $e^{-\varphi_{i}}$ ), and $v_{i}$ as a $(0,1)$ form with values in $L$. On $B\left(p_{i}, R\right)$, under the $L^{2}$-product with weight $e^{-\varphi_{i}}$ we decompose $v_{i}=v_{i}^{(1)}+v_{i}^{(2)}$, where $\bar{\partial} v_{i}^{(1)}=0$, and $v_{i}^{(2)} \perp \operatorname{ker} \bar{\partial}$. It follows that $\bar{\partial}_{\varphi_{i}}^{*} v_{i}^{(2)}=0$, and so also $\bar{\partial}_{\varphi_{i}}^{*} v_{i}^{(1)}=0$ outside of the support of $v_{i}$. By Lemma 3.5 below we have that

$$
\begin{equation*}
\left|v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}}<C\left\|v_{i}\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2} \tag{3.13}
\end{equation*}
$$

on the set $B\left(p_{i}, R-1\right) \backslash \bigcup \frac{1}{6} U_{i j}$, for a constant $C$ independent of $i$.
Let $\chi_{R}$ be a cutoff function, supported in $B\left(p_{i}, R-1\right)$, and equal to 1 on $B\left(p_{i}, R / 2\right)$ such that $\left|\nabla \chi_{R}\right|<C^{\prime} / R$ for a constant $C^{\prime}$ independent of $R$. We can regard the $(0,1)$-form $\chi_{R} v_{i}^{(1)}$ as an ( $n, 1$ )-form-valued in the anticanonical line bundle. From the Bochner-Kodaira formula (see (4.7) of [9]), we have

$$
\begin{align*}
& \int_{B\left(p_{i}, R\right)}\left(\sqrt{-1} \partial_{j} \partial_{\bar{k}} \varphi_{i}+\operatorname{Ric}_{j \bar{k}}^{g_{i}}\right)\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{j}} \overline{\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{k}}} e^{-\varphi_{i}}  \tag{3.14}\\
& \quad \leq\left\|\bar{\partial}\left(\chi_{R} v_{i}^{(1)}\right)\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2}+\left\|\bar{\partial}_{\varphi_{i}}^{*}\left(\chi_{R} v_{i}^{(1)}\right)\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2} .
\end{align*}
$$

To estimate the left-hand side, note that writing $\mathcal{V}_{i}=\bigcup \frac{1}{6} U_{i j}$, (3.11) implies that we have

$$
\begin{aligned}
& \int_{\mathcal{V}_{i}}\left(\sqrt{-1} \partial_{j} \partial_{\bar{k}} \varphi_{i}+\operatorname{Ric}_{j \bar{k}}^{g_{i}}\right)\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{j}} \overline{\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{k}}} e^{-\varphi_{i}} \\
& \quad \geq \int_{\mathcal{V}_{i}}\left(w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)-\Psi(1 / i)\right)\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}},
\end{aligned}
$$

while also using (3.1), (3.6), (3.12), and (3.13) we have

$$
\begin{aligned}
& \int_{B\left(p_{i}, R\right) \backslash \mathcal{V}_{i}}\left(\sqrt{-1} \partial_{j} \partial_{\bar{k}} \varphi_{i}+\mathrm{Ric}_{j \bar{k}}^{g_{i}}\right)\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{j}} \overline{\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{k}}} e^{-\varphi_{i}} \\
& \quad \geq-\Psi(1 / i) \int_{B\left(p_{i}, R\right) \backslash \mathcal{v}_{i}}\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}} \\
& \quad+\int_{B\left(p_{i}, R\right) \backslash \nu_{i}}\left(w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)\right)\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}}-
\end{aligned}
$$

$$
\begin{aligned}
& \quad-\int_{B\left(p_{i}, R\right) \backslash \mathcal{V}_{i}}\left|\sqrt{-1} \partial \bar{\partial} \varphi_{i}-\left(w^{\prime}\left(\rho_{i}\right) g_{i}+\sqrt{-1} w^{\prime \prime}\left(\rho_{i}\right) \partial \rho_{i} \wedge \bar{\partial} \rho_{i}\right)\right| \\
& \cdot\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}} \\
& -\int_{B\left(p_{i}, R\right) \backslash \mathcal{V}_{i}} w^{\prime \prime}\left(\rho_{i}\right)\left(\rho_{i}\left(g_{i}\right)_{j \bar{k}}-\sqrt{-1} \partial_{j} \rho_{i} \wedge \partial_{\bar{k}} \rho_{i}\right) \\
& \cdot\left(\chi_{R} v_{i}^{(1)}\right)_{\bar{j}} \overline{\left(\chi_{R} v_{i}^{(1)}\right)} \bar{k}_{\bar{k}} e^{-\varphi_{i}} \\
& \geq \\
& \int_{B\left(p_{i}, R\right) \backslash \mathcal{V}_{i}}\left(w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)\right)\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}}-\Psi\left(i^{-1} \mid R\right) .
\end{aligned}
$$

Notice that we used the assumption $w^{\prime \prime} \leq 0$ above. As for the right-hand side of (3.14), we use that $\bar{\partial} v_{i}^{(1)}=0$ and $\bar{\partial}_{\varphi_{i}}^{*} v_{i}^{(1)}=\bar{\partial}_{\varphi_{i}}^{*} v_{i}$, as well as the bound $\left|\nabla \chi_{R}\right|<$ $C^{\prime} / R$, to get

$$
\begin{aligned}
& \left\|\bar{\partial}\left(\chi_{R} v_{i}^{(1)}\right)\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2}+\left\|\bar{\partial}_{\varphi_{i}}^{*}\left(\chi_{R} v_{i}^{(1)}\right)\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2} \\
& \quad \leq\left\|\bar{\partial}_{\varphi_{i}}^{*} v_{i}\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2}+\frac{C^{\prime}}{R}\left\|v_{i}\right\|_{L^{2}\left(e^{-\varphi_{i}}\right)}^{2}
\end{aligned}
$$

Note that by construction the $v_{i}$ are bounded independently of $i$, and so their $L^{2}$ norms are uniformly bounded. It follows that for sufficiently large $i$ we have

$$
\int_{B\left(p_{i}, R\right)}\left(w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)\right)\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}} \leq\left\|\bar{\partial}_{\varphi_{i}}^{*} v_{i}\right\|_{L^{2}\left(e^{\left.-\varphi_{i}\right)}\right.}^{2}+\frac{C^{\prime}}{R}
$$

We now use the assumption that $\alpha=\bar{\partial} h$ is exact. Using the cutoff functions $\gamma_{j}$ from before, we define smooth functions $h_{i}$ on $B\left(p_{i}, R\right)$ analogously to the way we defined $\eta_{i}$. We then let $\alpha_{i}=\bar{\partial} h_{i}$. By taking $R$ large, we can assume that the supports of $h_{i}, \alpha_{i}$ are in $B\left(p_{i}, R / 2\right)$.

Since $\bar{\partial}_{\varphi_{i}}^{*} v_{i}^{(2)}=0$, we have

$$
\begin{align*}
& \left(\int_{B\left(p_{i}, R\right)} h_{i} \overline{\bar{\partial} \varphi_{\varphi_{i}}^{*} v_{i}} e^{-\varphi_{i}}\right)^{2} \\
& \quad=\left(\int_{B\left(p_{i}, R / 2\right)}\left\langle\bar{\partial} h_{i}, \chi_{R} v_{i}^{(1)}\right\rangle e^{-\varphi_{i}}\right)^{2} \\
& \quad \leq \int_{B\left(p_{i}, R / 2\right)} \frac{\left|\alpha_{i}\right|^{2} e^{-\varphi_{i}}}{w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)}  \tag{3.15}\\
& \quad \cdot \int_{B\left(p_{i} R / 2\right)}\left(w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)\right)\left|\chi_{R} v_{i}^{(1)}\right|^{2} e^{-\varphi_{i}} \\
& \quad \leq \int_{B\left(p_{i}, R / 2\right)} \frac{\left|\alpha_{i}\right|^{2} e^{-\varphi_{i}}}{w^{\prime}\left(\rho_{i}\right)+\rho_{i} w^{\prime \prime}\left(\rho_{i}\right)}\left(\left\|\bar{\partial}_{\varphi_{i}}^{*} v_{i}\right\|_{L^{2}\left(e^{\left.-\varphi_{i}\right)}\right.}^{2}+\frac{C^{\prime}}{R}\right)
\end{align*}
$$

for $i$ large enough. Note that in the first equality we used that the support of $h_{i}$ is contained in $B\left(p_{i}, \frac{R}{2}\right)$ where $\chi_{R}=1$.

Next we want to take the limit as $i \rightarrow \infty$. For this note first that

$$
\int_{B\left(p_{i}, R\right)} h_{i} \overline{\bar{\partial}_{\varphi_{i}}^{*} v_{i}} e^{-\varphi_{i}}=\int_{B\left(p_{i}, R\right)} h_{i} \overline{*_{\varphi_{i}} \bar{\partial} \eta_{i}} e^{-\varphi_{i}} d \mu_{i} .
$$

Let us consider one of our charts $U_{j}$, with coordinates $z_{i k}$ on $U_{i j}$ converging to $z_{k}$ on $U_{j}$. Note that in terms of these coordinates $\eta_{i}$ is not identified exactly with $\eta$ because of the way we defined $\eta_{i}$ in terms of the cutoff functions $\gamma_{j}$. However, if $U_{j}^{\prime}$ is another chart, with corresponding coordinates $z_{i k}^{\prime}$ converging to $z_{k}^{\prime}$, then the transition functions $z_{i k}^{\prime} \circ z_{i k}^{-1}$ converge smoothly to $z_{k}^{\prime} \circ z_{k}^{-1}$. It follows that if we use our coordinates to identify $U_{i j}$ with $U_{j}$, then $\eta_{i}$ converges smoothly to $\eta$ as $i \rightarrow \infty$. Similarly $h_{i}$ converges smoothly to $h$ and in addition $\varphi_{i}$ converges uniformly to $\varphi$. It then follows from Lemma 3.4 that

$$
\lim _{i \rightarrow \infty} \int_{B\left(p_{i}, R\right)} h_{i} \overline{* \varphi_{i} \bar{\partial} \eta_{i}} e^{-\varphi_{i}} d \mu_{i}=\int_{B(p, R)} h \overline{\bar{\partial} \eta} e^{-\varphi} d \mu
$$

Similarly we can take the limit as $i \rightarrow \infty$ of the right-hand side of (3.15), noting that $\bar{\partial}_{\varphi_{i}}^{*} v_{i}=*_{\varphi_{i}} \bar{\partial} \eta_{i}$ and using Lemma 3.4. We therefore have

$$
\begin{aligned}
& \left(\int_{B(p, R)} h \overline{\bar{\partial} \eta} e^{-\varphi} d \mu\right)^{2} \\
& \quad \leq \int_{B(p, R / 2)} \frac{|\alpha|^{2} e^{-\varphi}}{w^{\prime}\left(r^{2}\right)+r^{2} w^{\prime \prime}\left(r^{2}\right)} d \mu\left(\|* \bar{\partial} \eta\|_{L^{2}\left(e^{-\varphi}\right)}^{2}+\frac{C^{\prime}}{R}\right),
\end{aligned}
$$

Finally, we obtain the required inequality (3.5) by letting $R \rightarrow \infty$.
We used the following estimate in the proof above.
Lemma 3.5. Let $B(p, 2)$ be a relatively compact ball in a Kähler manifold with Ric $>-1$ and $\operatorname{vol}(B(p, 2))>v>0$. Let us denote by $L$ the trivial bundle with metric $e^{-\psi}$, and suppose that $|\nabla \psi|$ is bounded. Suppose that $u$ is a harmonic $(0,1)$-form valued in $L$. Then

$$
\sup _{B(p, 1)}|u| \leq C\left(\int_{B(p, 2)}|u|^{2}\right)^{1 / 2}
$$

where $C$ depends on $\nu$, the dimension, and $\sup _{B(p, 2)}|\nabla \psi|$.

Proof. The proof is by Moser iteration. Note first that by (2.2) the norm of $u$ satisfies the differential inequality

$$
\begin{align*}
\Delta|u|^{2}= & |\nabla u|^{2}+R^{j \bar{k}^{u}} \bar{k}_{\overline{u_{j}}}+2 g^{j \bar{k}_{g}} g^{p \bar{q}^{\prime}}\left(\nabla_{j} \nabla_{\bar{q}} \varphi\right) u_{\bar{k}} \overline{u_{\bar{p}}} \\
& -g^{j \bar{k}_{g}{ }^{p \bar{q}}\left(\nabla_{j} \nabla_{\bar{k}} \varphi\right) u_{\bar{q}} \overline{u_{\bar{p}}}} \begin{aligned}
\geq & |\nabla u|^{2}-|u|^{2}+2 g^{j \bar{k}} g^{p \bar{q}}\left(\nabla_{j} \nabla_{\bar{q}} \varphi\right) u_{\bar{k}} \overline{u_{\bar{p}}} \\
& -g^{j \bar{k}} g^{p \bar{q}}\left(\nabla_{j} \nabla_{\bar{k}} \varphi\right) u_{\bar{q}} \overline{u_{\bar{p}}} .
\end{aligned} \tag{3.16}
\end{align*}
$$

Let us write $B_{r}=B(p, r)$ for any $r \leq 2$. Let $\chi$ be a cutoff function supported in $B_{2}$, with $\chi=1$ on $B_{3 / 2}$ and $|\nabla \chi|+|\Delta \chi|<C$. Then we have (each integral being on $B_{2}$ )

$$
\begin{align*}
\int- & \chi^{2} \Delta|u|^{2} \\
\leq & -\int \chi^{2}|\nabla u|^{2}+\int \chi^{2}|u|^{2}-2 \int \chi^{2} g^{j \bar{k}_{g}} g^{p \bar{q}}\left(\nabla_{j} \nabla_{\bar{q}} \varphi\right) u_{\bar{k}} \overline{u_{\bar{p}}} \\
& +\int \chi^{2} g^{j \bar{k}} g^{p \bar{q}}\left(\nabla_{j} \nabla_{\bar{k}} \varphi\right) u_{\bar{q}} \overline{u_{\bar{p}}} \\
= & -\int \chi^{2}|\nabla u|^{2}+\int \chi^{2}|u|^{2}+2 \int 2 \chi g^{j \bar{k}_{g} p \bar{q}} \nabla_{j} \chi \nabla_{\bar{q}} \varphi u_{\bar{k}} \overline{u_{\bar{p}}}  \tag{3.17}\\
& +2 \int \chi^{2} g^{j \bar{k}_{g}{ }^{p \bar{q}} \nabla_{\bar{q}} \varphi\left(\bar{u}_{\bar{p}} \nabla_{j} u_{\bar{k}}+u_{\bar{k}} \overline{\nabla_{\bar{j}} u_{\bar{p}}}\right)} \\
& -\int 2 \chi g^{j \bar{k}_{g}}{ }^{p \bar{q}} \nabla_{j} \chi \nabla_{\bar{k}} \varphi u_{\bar{q}} \overline{u_{\bar{p}}} \\
& -\int \chi^{2} g^{j \bar{k}_{g}}{ }^{p \bar{q}} \nabla_{\bar{k}} \varphi\left(\overline{u_{\bar{p}}} \nabla_{j} u_{\bar{q}}+u_{\bar{q}} \overline{\nabla_{\bar{j}} u \bar{p}}\right) \\
\leq & -\frac{1}{2} \int \chi^{2}|\nabla u|^{2}+C \int|u|^{2},
\end{align*}
$$

for a constant $C$ depending on the dimension, $v$, and $\sup _{B(p, 2)}|\nabla \varphi|$. At the same time

$$
\int \chi^{2} \Delta|u|^{2}=\int \Delta\left(\chi^{2}\right)|u|^{2} \leq C \int|u|^{2}
$$

and combining this with our previous inequality, we get

$$
\begin{equation*}
\int \chi^{2}|\nabla u|^{2}<C \int|u|^{2} \tag{3.18}
\end{equation*}
$$

Under our assumptions we have a bound for the Sobolev constant, and so

$$
\begin{aligned}
\left(\int|\chi| u\left|\left.\right|^{\frac{2 n}{n-2}}\right)^{\frac{n-2}{n}}\right. & \leq \int|\nabla(\chi|u|)|^{2} \\
& =\int|\nabla \chi|^{2}|u|^{2}+\int \chi^{2}|\nabla| u| |^{2}+2 \chi|u| \nabla \chi \cdot \nabla|u| \\
& \leq C \int|u|^{2}
\end{aligned}
$$

using the Cauchy-Schwarz inequality and also (3.18). It follows that

$$
\begin{equation*}
\|u\|_{L^{\frac{2 n}{n-2}\left(B_{3 / 2}\right)}} \leq C\|u\|_{L^{2}\left(B_{2}\right)} \tag{3.19}
\end{equation*}
$$

To estimate higher $L^{p}$ norms, let $p \geq \frac{2 n}{n-2}$, and let $\chi$ be a smooth function compactly supported in $B_{2}$. We can compute, using the Sobolev inequality,

$$
\begin{align*}
& \left(\left.\left.\int|\chi| u\right|^{\frac{p}{2}}\right|^{\frac{2 n}{n-2}}\right)^{\frac{n-2}{n}} \\
& \quad \leq C \int\left|\nabla\left(\chi|u|^{\frac{p}{2}}\right)\right|^{2}  \tag{3.20}\\
& \quad=C \int|\nabla \chi|^{2}|u|^{p}+\left.\left.\chi^{2}|\nabla| u\right|^{\frac{p}{2}}\right|^{2}+2 \chi|u|^{\frac{p}{2}} \nabla \chi \cdot \nabla|u|^{\frac{p}{2}}
\end{align*}
$$

We have

$$
\begin{aligned}
\int & \left.\left.\chi^{2}|\nabla| u\right|^{\frac{p}{2}}\right|^{2} \\
& =\int \frac{p^{2}}{8(p-2)} \chi^{2} \nabla|u|^{p-2} \cdot \nabla|u|^{2} \\
& =-\int \frac{p^{2}}{4(p-2)} \chi|u|^{p-2} \nabla \chi . \nabla|u|^{2}-\int \frac{p^{2}}{8(p-2)} \chi^{2}|u|^{p-2} \Delta|u|^{2}
\end{aligned}
$$

and

$$
\begin{aligned}
\int & 2 \chi|u|^{\frac{p}{2}} \nabla \chi \cdot \nabla|u|^{\frac{p}{2}} \\
& =\int \frac{p}{4}|u|^{p-2} \nabla \chi^{2} \cdot \nabla|u|^{2} \\
& =-\int \frac{p(p-2)}{8} \chi^{2}|u|^{p-4} \nabla|u|^{2} \cdot \nabla|u|^{2}-\int \frac{p}{4} \chi^{2}|u|^{p-2} \Delta|u|^{2} \\
& \leq-\int \frac{p}{4} \chi^{2}|u|^{p-2} \Delta|u|^{2}
\end{aligned}
$$

Combining these with (3.20) we get

$$
\begin{align*}
\left(\left.\left.\int|\chi| u\right|^{\frac{p}{2}}\right|^{\frac{2 n}{n-2}}\right)^{\frac{n-2}{n}} \leq C \int & |\nabla \chi|^{2}|u|^{p}-\frac{p^{2}}{2(p-2)} \chi|u|^{p-1} \nabla \chi \cdot \nabla|u|  \tag{3.21}\\
& -\left(\frac{p^{2}}{8(p-2)}+\frac{p}{4}\right) \chi^{2}|u|^{p-2} \Delta|u|^{2} .
\end{align*}
$$

Using (3.16) we have, similarly to (3.17),

$$
\begin{align*}
& -\int \chi^{2}|u|^{p-2} \Delta|u|^{2}  \tag{3.22}\\
& \quad \leq-\frac{1}{2} \int \chi^{2}|u|^{p-2}|\nabla u|^{2}+C p^{2} \int\left(\chi^{2}+|\nabla \chi|^{2}\right)|u|^{p}
\end{align*}
$$

where $C$ depends on $\sup |\nabla \varphi|$. Since $p \geq \frac{2 n}{n-2}$, we have some constant $c_{n}>0$ such that

$$
c_{n} p<\frac{p^{2}}{2(p-2)}<c_{n}^{-1} p, \quad c_{n} p<\left(\frac{p^{2}}{8(p-2)}+\frac{p}{4}\right)<c_{n}^{-1} p .
$$

It then follows, combining (3.21) and (3.22), and using the Cauchy-Schwarz inequality, that

$$
\begin{equation*}
\left(\left.\left.\int|\chi| u\right|^{\frac{p}{2}}\right|^{\frac{2 n}{n-2}}\right)^{\frac{n-2}{n}} \leq C p^{3} \int\left(\chi^{2}+|\nabla \chi|^{2}\right)|u|^{p} . \tag{3.23}
\end{equation*}
$$

We now choose a sequence of cutoff functions $\chi_{i}$ such that $\chi_{i}$ is supported in $B_{1+2^{-i}}$, and $\chi_{i}=1$ on $B_{1+2^{-i-1}}$. We can arrange that $\left|\nabla \chi_{i}\right|<C 2^{i}$. We also define $\gamma=\frac{n}{n-2}$. Equation (3.23), setting $\chi=\chi_{i}$ and $p=2 \gamma^{i}$, implies that

$$
\left(\int \chi_{i}^{2 \gamma}|u|^{2 \gamma^{i+1}}\right)^{\frac{1}{2 \gamma^{i+1}}} \leq\left(8 C \gamma^{3 i} 2^{2 i}\right)^{\frac{1}{2 \gamma^{i}}}\left(\int_{B_{1+2^{-i}}}|u|^{2 \gamma^{i}}\right)^{\frac{1}{2 \gamma^{i}}} .
$$

From this we have

$$
\|u\|_{L^{2 \gamma^{i+1}}\left(B_{1+2^{-i-1}}\right)} \leq C^{\frac{3 i}{\nu^{i}}}\|u\|_{L^{2 \gamma^{i}}\left(B_{1+2^{-i}}\right)} .
$$

Iterating this, we find that

$$
\sup _{B_{1}}|u| \leq C\|u\|_{L^{2 \gamma}\left(B_{3 / 2}\right)} \leq C^{\prime}\|u\|_{L^{2}\left(B_{1}\right)},
$$

using (3.19).
We will use Proposition 3.1 to construct holomorphic functions on the tangent cone $X$. We will need the following result, which implies that the functions that we construct are actually harmonic across the singular set, and in particular satisfy local $L^{\infty}$ and Lipschitz estimates. Note that in the setting of Section 2 with twosided Ricci bounds, this step is much more straightforward since the singular set has codimension 4.

Proposition 3.6. Let $f: X \rightarrow \mathbf{C}$ be a function such that $\bar{\partial} f=0$ on $\mathcal{R}_{\epsilon}$, and $f$ has polynomial growth in the sense that for some $k, D>0$ we have

$$
\begin{equation*}
\int_{B(o, R)}|f|^{2}<D R^{k} \tag{3.24}
\end{equation*}
$$

for all $R>0$. Then $f$ is harmonic on $X$. In particular, we have local estimates

$$
\sup _{y \in B(x, 1 / 2)}|f(y)|+\sup _{y, z \in B(x, 1 / 2)} \frac{|f(y)-f(z)|}{d(y, z)} \leq C\left(\int_{B(x, 1)}|f|^{2}\right)^{1 / 2}
$$

for a constant $C$ depending only on the dimension and the noncollapsing constant $v$. In addition, $f$ has polynomial growth in the pointwise sense that $|f(x)| \leq$ $C(1+d(o, x))^{d}$ for some $C, d>0$ depending on $k, D$.

Remark 3.7. In the statement we require the global growth condition (3.24) since we use the heat flow to smooth out $f$. The argument could be localized by working with Green's functions on balls instead, but we will not need this.

Proof. Let $H(x, y, t)$ be the heat kernel on $X$. We let $f_{t}$ be the evolution of $f$ under the heat flow:

$$
f_{t}(x)=\int f(y) H(x, y, t) d \mu(y)
$$

Note that the polynomial growth assumption on $f$ together with the exponential decay of $H(x, y, t)$ ensures that this is well-defined. Our goal is to show that $f_{t}=f$ for all $t>0$, which implies that $f$ is harmonic. We will use results about the convergence of heat kernels under Gromov-Hausdorff convergence due to Ding [10], and argue by approximating $f_{t}$ with corresponding flows on smooth spaces.

We fix $T \in(0,1)$ and $q \in \mathcal{R}_{\epsilon} \cap B(o, 1)$. We will show that $f_{T}(q)=f_{1}(q)$. Let $R>0$ be large, and set $\delta=e^{-R}$. The balls $B\left(p_{i}, R\right)$ converge to $B(o, R) \subset X$ in the Gromov-Hausdorff sense. By the results of Ding, the heat kernels $H_{i}(x, y, t)$ on $B\left(p_{i}, R\right)$ (with Dirichlet boundary conditions) satisfy

$$
\left|H_{i}(x, y, t)-H(x, y, t)\right|<\Psi\left(i^{-1}, R^{-1}\right)
$$

for $t=T, 1$ and $x, y \in B(o, R)$, where we use Gromov-Hausdorff approximations to identify points in $B\left(p_{i}, R\right)$ and $B(o, R)$.

Using that $X \backslash \mathcal{R}_{\epsilon}$ has locally finite codimension 2 Minkowski measure, we can find a cutoff function $\eta$ with compact support in $\mathcal{R}_{\epsilon}$ such that $\eta(q)=1$ and $\int_{B(o, R)}|\nabla \eta|^{2}<\delta$. In addition, we can ensure that $\int_{B(o, R)}(1-\eta)^{2}<\delta$. We have corresponding cutoff functions $\eta_{i}$ on $B\left(p_{i}, R\right)$ converging to $\eta$ and satisfying the same estimates.

Using the local holomorphic charts, we can find functions $f_{i}$ on $B\left(p_{i}, R\right)$, converging uniformly to $f$ on the support of $\eta$, such that $\left|\bar{\partial} f_{i}\right|<\Psi\left(i^{-1}\right)$ on $\operatorname{supp}\left(\eta_{i}\right)$.

In particular, this implies that

$$
\int_{B\left(p_{i}, R\right)}\left|\bar{\partial}\left(\eta_{i} f_{i}\right)\right|<\Psi\left(i^{-1}, R^{-1}\right) .
$$

Let $\chi_{i}$ denote cutoff functions supported in $B\left(p_{i}, R\right)$ and equal to 1 on $B\left(p_{i}, R-1\right)$ such that $\left|\nabla \chi_{i}\right|<C$ for a uniform constant $C$. Consider the heat flow for the functions $\chi_{i} \eta_{i} f_{i}$. We let

$$
f_{i, R, t}(q)=\int_{B\left(p_{i}, R\right)} \chi_{i}(y) \eta_{i}(y) f_{i}(y) H_{i}(q, y, t) d \mu_{i}(y)
$$

We have

$$
\begin{aligned}
\frac{d}{d t} f_{i, R, t}(q)= & \int_{B\left(p_{i}, R\right)} \chi_{i} \eta_{i} f_{i} g_{i}^{a \bar{b}} \bar{\partial}_{b} \partial_{a} H_{i}(q, y, t) d \mu_{i}(y) \\
= & -\int_{B\left(p_{i}, R\right)}\left(\bar{\partial}_{b} \chi_{i}\right) \eta_{i} f_{i} g_{i}^{a \bar{b}} \partial_{a} H_{i}(q, y, t) d \mu_{i}(y) \\
& -\int_{B\left(p_{i}, R\right)} \chi_{i} \bar{\partial}_{b}\left(\eta_{i} f_{i}\right) g_{i}^{a \bar{b}} \partial_{a} H_{i}(q, y, t) d \mu_{i}(y)
\end{aligned}
$$

On the support of $\nabla \chi_{i}$ we have the estimate $\left|\nabla_{y} H_{i}(q, y, t)\right|<C_{T} e^{-c R^{2}}$ for $t \in$ [ $T, 1$ ] for some $c>0$ and $C_{T}>0$ depending on $T$, while on $B\left(p_{i}, R\right)$ we have $\left|\nabla_{y} H(q, y, t)\right|<C_{T}$ for $t \in[T, 1]$. It follows that

$$
\left|\frac{d}{d t} f_{i, R, t}(q)\right|<C_{T} \Psi\left(i^{-1} R^{-1}\right)
$$

for $t \in[T, 1]$ and so

$$
\begin{equation*}
\left|f_{i, R, 1}(q)-f_{i, R, T}(q)\right|<C_{T} \Psi\left(i^{-1}, R^{-1}\right) \tag{3.25}
\end{equation*}
$$

The convergence of $\eta_{i} f_{i}$ to $\eta f$, the convergence of the heat kernels as $i \rightarrow \infty$, and their exponential decay imply that

$$
\begin{equation*}
\left|f_{i, R, t}(q)-f_{t}(q)\right|<\Psi\left(i^{-1} \mid R\right) \tag{3.26}
\end{equation*}
$$

for $t=T, 1$. Choosing $R$ sufficiently large (depending on $T$ ) and $i$ large (depending on $T, R$ ), from (3.25) and (3.26) we see that $\left|f_{1}(q)-f_{T}(q)\right|$ can be made arbitrarily small, and therefore $f_{1}(q)=f_{T}(q)$.

## 4 Tangent Cones Are Affine Varieties

As before, we suppose that $X$ is a tangent cone at a point of a noncollapsed limit of Kähler manifolds with Ricci curvature bounded below. Our main goal in this section is to prove that $X$ is homeomorphic to an affine variety. Given the $L^{2}-$ existence result Proposition 3.1 or Proposition 2.1 when we have two-sided Ricci bounds, we can more or less follow the argument in Donaldson-Sun [11, 12] with suitable modifications as in our previous work [18], where only a lower bound for the Ricci curvature is assumed.

Recall that we have a subset $\mathcal{R}_{\epsilon} \subset X$ which has the structure of a complex manifold, and whose complement has locally finite codimension-2 Minkowski content. Our first goal is to show that $X$ is locally homeomorphic to a complex variety by constructing holomorphic functions that embed a neighborhood of the vertex $o \in X$ into $\mathbf{C}^{N}$. The basic ingredient for this is the following.

PROPOSITION 4.1. We can use holomorphic functions of polynomial growth on $X$ to separate points and also to separate tangents at points in $\mathcal{R}_{\epsilon}$. More precisely:
(1) Let $x_{1}, x_{2} \in X$. There exists a holomorphic function $f$ of polynomial growth on $X$ such that $f\left(x_{1}\right) \neq f\left(x_{2}\right)$.
(2) Let $x \in \mathcal{R}_{\epsilon}$. There are holomorphic functions $f_{1}, \ldots, f_{n}$ of polynomial growth on $X$ that define an embedding of a neighborhood of $x$ into $\mathbf{C}^{n}$.
Proof. The proof of this closely follows arguments in Donaldson-Sun [11] as well as the authors' work [18], given the $L^{2}$-existence result Proposition 3.1 together with Proposition 3.6. We can construct holomorphic functions of polynomial growth by using a logarithmic weight $C \log \left(1+r^{2}\right)$; however, in this case the curvature of the metric $e^{-\varphi}$ is not proportional to the metric on the cone. We can, however, use a slightly different weight function $\varphi=w\left(r^{2}\right)$, where $w(t)=t+1$ for $t$ close to 0 , and $w(t)=C \log (1+t)$ for $t$ large, still satisfying the inequalities (3.1). Then in a neighborhood of the vertex the curvature of $e^{-\varphi}$ will be the metric on $X$. Then we can proceed as in [11] or [18] to finish the proof of (1) at least for $x_{1}, x_{2}$ close to the vertex. By scaling this can be extended to all of $X$. Part (2) can be proven in a standard way using weight functions with logarithmic poles (see Remark 3.2).

As in Chen-Donaldson-Sun [7, sec. 2.5] or [18, prop. 3.2], we can show that if at a point $x \in X$ there is a tangent cone that splits off $\mathbf{R}^{2 n-2}$, then polynomial growth holomorphic functions can be used to embed a neighborhood of $x$ into $\mathbf{C}^{n}$.
LEMMA 4.2. A small neighborhood of $o \in X$ is homeomorphic to a normal analytic space.

Proof. The argument is very similar to [12]. By using the separation of points, we can find $N$ polynomial growth holomorphic functions $f_{1}, \ldots, f_{N}$ which all vanish at $o$ so that

$$
\sum_{j=1}^{N}\left|f_{j}\right|^{2} \geq c>0 \quad \text { on } \partial B(o, 1)
$$

One can verify, by using cutoff functions along the singular set, that the image of $\left(f_{1}, \ldots, f_{N}\right)$ near the origin of $\mathbb{C}^{N}$ defines a $d$-closed positive locally rectifiable current of type $(N-n, N-n)$. According to the main theorem in [16] (or equivalently, theorem $1.3\left(a^{\prime}\right)$ of [13]), the image of $\left(f_{1}, \ldots, f_{N}\right)$ defines a complex analytic variety of dimension $n$ near the origin in $\mathbb{C}^{N}$. By adding more holomorphic functions, we can ensure as in [12, props. 2.3, 2.4] that this map is a homeomorphism locally, and the image is a normal variety.

Remark 4.3. Note that unlike in [11, 12] or [18], the holomorphic functions that we construct on $X$ do not necessarily arise as limits of sequences of holomorphic functions on the approximating manifolds.

In order to show that $X$ is homeomorphic to a normal affine variety, we can use the method in Donaldson-Sun [12], which in turn is based on an idea of Van Coevering [20] in the setting of Kähler cones with smooth links. The main point is to decompose polynomial growth holomorphic functions into sums of homogeneous holomorphic functions under the homothetic vector field $r \partial_{r}$. Recall that by [17, theorem 2] there is a one parameter group $\sigma_{t}$ of isometries acting on the tangent cone $X$, preserving the distance function from the vertex $o \in X$. Note that since it acts by isometries, the action of $\sigma_{t}$ preserves $\mathcal{R}_{\epsilon}$. We need the following.

Proposition 4.4. The group $\sigma_{t}$ acts by biholomorphisms on $\mathcal{R}_{\epsilon}$. In addition, the vector field $v$ on $\mathcal{R}_{\epsilon}$ generating this action satisfies $v=J r \partial_{r}$.

Proof. Let us briefly recall the construction of the isometric action from [17]. We consider the geodesic annuli $A_{i}=B\left(p_{i}, 10\right) \backslash B\left(p_{i}, 1 / 3\right)$. From CheegerColding [3] there are smooth functions $\rho_{i}$ on $A_{i}$ such that

$$
\begin{gather*}
\int_{A_{i}}\left|\nabla^{2} \rho_{i}-g_{i}\right|^{2}+\left|\nabla \rho_{i}-\nabla r_{i}^{2} / 2\right|^{2}<\Psi\left(i^{-1}\right),  \tag{4.1}\\
\left|\rho_{i}-r_{i}^{2} / 2\right|<\Psi\left(i^{-1}\right), \quad\left|\nabla \rho_{i}\right| \leq C,
\end{gather*}
$$

where $r_{i}$ is the distance function from $p_{i}$. Define vector fields $v_{i}=J \nabla \rho_{i}$, and let $\sigma_{i, t}$ be the diffeomorphisms generated by $v_{i}$. It is shown in [17] that for small $t$ we can extract a limit $\sigma_{t}$ as $i \rightarrow \infty$, which gives rise to a one-parameter group of isometries on $B(o, 6) \backslash B(o, 5)$ in $X$. Since the action commutes with the homothetic transformations of $X$, it can be extended to all of $X$.

We will show that for small $t$ the limit $\sigma_{t}$ is a biholomorphism. Let $q \in \mathcal{R}_{\epsilon}$ such that $q \in B(o, 6) \backslash B(o, 5)$. Let $q_{i} \in B\left(p_{i}, 10\right)$ such that $q_{i} \rightarrow q$. We have a small $\delta>0$ and holomorphic coordinates $z_{i j}$ on $B\left(q_{i}, \delta\right)$ converging to holomorphic coordinates $z_{j}$ on $B(q, \delta)$.

Abusing notation, we will also denote by $v_{i}$ the $(1,0)$-part of the real vector field defined above. The estimate (4.1) implies that

$$
\begin{equation*}
\int_{A_{i}}\left|\bar{\partial} v_{i}\right|^{2}<\Psi\left(i^{-1}\right) . \tag{4.2}
\end{equation*}
$$

Let us write $v_{i}=v_{i}^{p} \partial_{z_{i p}}$ on $B\left(q_{i}, \delta\right)$ in terms of the holomorphic coordinates. Since in these coordinates we have a lower bound $g_{i, a \bar{b}}>C^{-1} \delta_{a \bar{b}}$, it follows from (4.2) that the components $v_{i}^{p}$ satisfy

$$
\int_{B\left(q_{i}, \delta\right)}\left|\bar{\partial} v_{i}^{p}\right|^{2}<\Psi\left(i^{-1}\right)
$$

On the ball $B\left(q_{i}, \delta\right)$ we can apply the $L^{2}$-estimate to obtain holomorphic functions $\widetilde{v}_{i}^{p}$ with

$$
\begin{equation*}
\int_{B\left(q_{i}, \delta / 2\right)}\left|v_{i}^{p}-\tilde{v}_{i}^{p}\right|^{2}<\Psi\left(i^{-1}\right) \tag{4.3}
\end{equation*}
$$

At the same time, using (4.1), we have uniform bounds $\left|v_{i}^{p}\right|,\left|\widetilde{v}_{i}^{p}\right|<C$. Let us denote by $\tilde{\sigma}_{i, t}$ the flow of (the real part of) $\tilde{v}_{i}=\tilde{v}_{i}^{p} \partial_{z_{i p}}$. As long as $t<T$ for sufficiently small $T>0$, this is well-defined on $B\left(q_{i}, \delta / 3\right)$, say. Up to choosing a subsequence we can assume that the $\tilde{v}_{i}$ converge to a holomorphic vector field $\tilde{v}$ on $B(q, \delta / 3)$, generating a one-parameter group of biholomorphisms $\tilde{\sigma}_{t}$.

We claim that the diffeomorphisms $\sigma_{i, t}$ are approximated by $\tilde{\sigma}_{i, t}$ for $t<T$, and therefore the limiting flow $\sigma_{t}$ coincides with the flow $\tilde{\sigma}_{t}$ by biholomorphisms. More precisely, define the function

$$
u(x, t)=\left|\sigma_{i, t}(x)-\widetilde{\sigma}_{i, t}(x)\right|^{2}
$$

for $x \in B\left(q_{i}, \delta / 3\right)$ and $t<T$. We will show that we have $u(x, t)<\Psi\left(i^{-1}\right)$ for $x$ outside of a set of measure at most $\Psi\left(i^{-1}\right)$. Note that here we are viewing $B\left(q_{i}, \delta\right)$ as a subset of $\mathbf{C}^{n}$ using our coordinates. Define

$$
F(t)=\int_{B(q, \delta / 3)} u(x, t)
$$

Note that

$$
\begin{aligned}
\frac{d}{d t} u(x, t)= & 2\left(\sigma_{i, t}(x)-\widetilde{\sigma}_{i, t}(x)\right) \cdot\left(v_{i}\left(\sigma_{i, t}(x)\right)-\tilde{v}_{i}\left(\widetilde{\sigma}_{i, t}(x)\right)\right) \\
= & 2\left(\sigma_{i, t}(x)-\widetilde{\sigma}_{i, t}(x)\right) \\
& \cdot\left(v_{i}\left(\sigma_{i, t}(x)\right)-\tilde{v}_{i}\left(\sigma_{i, t}(x)\right)+\tilde{v}_{i}\left(\sigma_{i, t}(x)\right)-\tilde{v}_{i}\left(\tilde{\sigma}_{i, t}(x)\right)\right)
\end{aligned}
$$

Using that we have a gradient bound for the holomorphic vector field $\tilde{v}_{i}$, this implies

$$
\begin{aligned}
\frac{d}{d t} u(x, t) & \leq 2 \sqrt{u(x, t)}\left|v_{i}-\tilde{v}_{i}\right|\left(\sigma_{i, t}(x)\right)+C u(x, t) \\
& \leq C u(x, t)+\left|v_{i}-\tilde{v}_{i}\right|_{\mathrm{Euc}}^{2}\left(\sigma_{i, t}(x)\right)
\end{aligned}
$$

where we emphasize that the difference $\left|v_{i}-\tilde{v}_{i}\right|$ is measured using the Euclidean metric given by our coordinates. We can now use the fact that $\sigma_{i, t}$ distorts volumes by at most $\Psi\left(i^{-1}\right)$ (see the proof of theorem 2 in [17]), together with the estimate (4.3) and the uniform bounds for $v_{i}^{p}, \widetilde{v}_{i}^{p}$ to see that

$$
\frac{d}{d t} F(t) \leq C F(t)+\Psi\left(i^{-1}\right)
$$

Since $F(0)=0$, we get $F(t)<\Psi\left(i^{-1}\right)$ for $t<T$, and from this it follows in turn that $u(x, t)<\Psi\left(i^{-1}\right)$ outside of a set of measure at most $\Psi\left(i^{-1}\right)$, as required.

Note that if in terms of the notation above we let $w_{i}=\nabla \rho_{i}$, and we let $\varphi_{i, t}$ denote the one-parameter group of diffeomorphisms generated by $w_{i}$, then $\varphi_{i, t}$
converges to the homothetic expansion map on $X$, generated by $r \partial_{r}$ on $\mathcal{R}_{\epsilon}$. Since $J w_{i}=v_{i}$, it is not hard to see that in the limit we have $J r \partial_{r}=v$.

The action of $\sigma_{t}$ defines a subgroup isomorphic to $\mathbf{R}$ in the isometry group of $X$ preserving the vertex $o$. The latter group is a Lie group by [4]. Taking the closure of this $\mathbf{R}$-subgroup defines a torus $T$ acting isometrically on $X$, which also acts by biholomorphisms on $\mathcal{R}_{\epsilon}$. Fixing $d>0$, let $\mathcal{H}_{d}$ denote the vector space of polynomial growth holomorphic functions $f$ on $X$ of degree at most $d$, i.e., satisfying $|f(x)|<C(1+d(o, x))^{d}$ for a constant $C>0$. Since such $f$ are harmonic, $\mathcal{H}_{d}$ is finite dimensional, and decomposes into weight spaces under the $T$-action. It follows that any $f \in \mathcal{H}_{d}$ can be decomposed into a sum of eigenfunctions $f=f_{\alpha_{1}}+\cdots+f_{\alpha_{m}}$ for $\alpha_{k} \in \operatorname{Lie}(T)^{*}$, where $e^{\sqrt{-1} t} \cdot f_{\alpha_{k}}=e^{\sqrt{-1}\left\langle\alpha_{k}, t\right\rangle} f_{\alpha_{k}}$ for $t \in \operatorname{Lie}(T)$. The one-parameter group of isometries $\sigma_{t}$ is generated by a vector $\xi \in \operatorname{Lie}(T)$. The relation $v=J r \partial_{r}$ then implies that $r \partial_{r} f_{\alpha_{k}}=\left\langle\xi, \alpha_{k}\right\rangle f_{\alpha_{k}}$, i.e., each $f_{\alpha_{k}}$ is homogeneous. We can therefore embed a neighborhood of $o \in X$ into $\mathbf{C}^{N}$ using homogeneous holomorphic functions, and this embedding naturally extends to an embedding of all of $X$. As in [12, lemma 2.19] it follows that the image of $X$ is an affine variety and by construction the torus $T$ acts linearly. This completes the proof of Theorem 1.1.

The proof of Corollary 1.2 follows exactly as proposition 2.21 and lemma 4.2 of Donaldson-Sun [12], based on the work of Martelli-Sparks-Yau [19].

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## Bibliography

[1] Anderson, M. T. Convergence and rigidity of manifolds under Ricci curvature bounds. Invent. Math. 102 (1990), no. 2, 429-445. doi:10.1007/BF01233434
[2] Cheeger, J.; Colding, T. H. Lower bounds on Ricci curvature and the almost rigidity of warped products. Ann. of Math. (2) 144 (1996), no. 1, 189-237. doi:10.2307/2118589
[3] Cheeger, J.; Colding, T. H. On the structure of spaces with Ricci curvature bounded below. I. J. Differential Geom. 46 (1997), no. 3, 406-480.
[4] Cheeger, J.; Colding, T. H. On the structure of spaces with Ricci curvature bounded below. II. J. Differential Geom. 54 (2000), no. 1, 13-35.
[5] Cheeger, J.; Colding, T. H.; Tian, G. On the singularities of spaces with bounded Ricci curvature. Geom. Funct. Anal. 12 (2002), no. 5, 873-914. doi:10.1007/PL00012649
[6] Cheeger, J.; Jiang, W.; Naber, A. Rectifiability of singular sets in spaces with Ricci curvature bounded below. Preprint, 2018. arXiv:1805.07988 [math.DG]
[7] Chen, X.; Donaldson, S.; Sun, S. Kähler-Einstein metrics on Fano manifolds. II: Limits with cone angle less than $2 \pi$. J. Amer. Math. Soc. 28 (2015), no. 1, 199-234. doi:10.1090/S0894-0347-2014-00800-6
[8] Colding, T. Ricci curvature and volume convergence. Ann. of Math. (2) 145 (1997), no. 3, 477501. doi:10.2307/2951841
[9] Demailly, J.-P. Analytic methods in algebraic geometry. Surveys of Modern Mathematics, 1. International Press, Somerville, Mass.; Higher Education Press, Beijing, 2012.
[10] Ding, Y. Heat kernels and Green's functions on limit spaces. Comm. Anal. Geom. 10 (2002), no. 3, 475-514. doi:10.4310/CAG.2002.v10.n3.a3
[11] Donaldson, S.; Sun, S. Gromov-Hausdorff limits of Kähler manifolds and algebraic geometry. Acta Math. 213 (2014), no. 1, 63-106. doi:10.1007/s11511-014-0116-3
[12] Donaldson, S.; Sun, S. Gromov-Hausdorff limits of Kähler manifolds and algebraic geometry, II. J. Differential Geom. 107 (2017), no. 2, 327-371. doi:10.4310/jdg/1506650422
[13] Harvey, R.; Polking, J. Extending analytic objects. Comm. Pure Appl. Math. 28 (1975), no. 6, 701-727. doi:10.1002/cpa. 3160280603
[14] Hörmander, L. $L^{2}$ estimates and existence theorems for the $\bar{\partial}$-operator. Acta Math. 113 (1965), 89-152. doi:10.1007/BF02391775
[15] Jiang, W.; Naber, A. $L^{2}$ curvature bounds on manifolds with bounded Ricci curvature. Preprint, 2016. arXiv:1605.05583 [math.DG]
[16] King, J. R. The currents defined by analytic varieties. Acta Math. 127 (1971), no. 3-4, 185-220. doi:10.1007/BF02392053
[17] Liu, G. On the tangent cone of Kähler manifolds with Ricci curvature lower bound. Math. Ann. 370 (2018), no. 1-2, 649-667. doi:10.1007/s00208-017-1536-0
[18] Liu, G.; Székelyhidi, G. Gromov-Hausdorff limits of Kähler manifolds with Ricci curvature bounded below. Preprint, 2018. arXiv: 1804.08567 [math.DG]
[19] Martelli, D.; Sparks, J.; Yau, S.-T. Sasaki-Einstein manifolds and volume minimisation. Comm. Math. Phys. 280 (2008), no. 3, 611-673. doi:10.1007/s00220-008-0479-4
[20] Van Coevering, C. Examples of asymptotically conical Ricci-flat Kähler manifolds. Math. Z. 267 (2011), no. 1-2, 465-496. doi:10.1007/s00209-009-0631-7

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