

The kinetic study of excited singlet oxygen atom  $O(^1D)$  reactions with acetylene

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Understanding the multi-channel dynamics of  $O(^1D)$  reactions with unsaturated hydrocarbon molecules in low temperature reaction kinetics is critically important in stratospheric chemistry, plasma chemistry, and plasma assisted fuel reforming, materials synthesis, and combustion. A photolysis flow reactor coupled with highly selective mid-infrared Faraday Rotation Spectroscopy (FRS) and direct UV-IR absorption spectroscopy (DAS) techniques was developed for the first time to study the multi-channel dynamics of excited singlet oxygen atom  $O(^1D)$  reactions with  $C_2H_2$  and the kinetics of subsequent reactions. Time-resolved species concentrations of  $O(^1D)$ , OH,  $HO_2$  and  $H_2O$  were obtained and used to develop a validated kinetic model of  $O(^1D)$  reactions with  $C_2H_2$ . The branching ratios of  $O(^1D)$  reaction with  $C_2H_2$  and subsequent  $HO_2$  kinetics were also quantified. It is found that, contrary to  $O(^1D)$  reactions with saturated alkanes, OH formation via direct H abstraction by  $O(^1D)$  is negligible. The results revealed that two chain-branching and propagation reactions via direct  $O(^1D)$  insertion are the major pathways for radical production. Moreover, a new reaction pathway for low temperature  $HO_2$  production via  $C_2H_2OH + O_2$  is identified and its reaction rate is determined. The present study clearly demonstrated the advantage of radical detection and kinetic studies using FRS in the effective suppression of absorption interference from non-paramagnetic hydrocarbons.

## TOC GRAPHICS

**KEYWORDS:** Singlet oxygen atom; Atmospheric Kinetics; Photolysis Herriott cell; Faraday rotation spectroscopy; Balanced Detection; Plasma assisted combustion.

## 1. Introduction

The chemical behavior of electronically excited singlet oxygen atom,  $O(^1D)$ , is of considerable interest and remains experimentally attractive throughout the past few decades due to its substantial roles in establishing chemical composition of atmospheric chemistry, plasma fuel reforming, plasma assisted materials synthesis and combustion.<sup>1-2</sup> Numerous studies<sup>3-5</sup> have shown that  $O(^1D)$  atoms contribute significantly to enhance ignition, reform fuel, and reduce emissions. In particular,  $O(^1D)$  plays an important role in accelerating low temperature chain branching, leading to the formation of cool flames and warm flames.<sup>6-7</sup> Moreover,  $O(^1D)$  also contributes to the oxidation of leaked fuels and unburned hydrocarbons as well as the  $O_3$  cycle in the atmosphere.<sup>8-9</sup> As such, understanding the multi-channel dynamics of the  $O(^1D)$  reaction with fuel molecules and the low temperature reaction kinetics is critical. However,  $O(^1D)$  reactions with fuels are complicated by its multi-channel behavior due to direct oxygen atom insertion process.

Early studies<sup>10-15</sup> of the reaction between  $O(^1D)$  with saturated hydrocarbons have been conducted to understand the role of  $O(^1D)$  initiated chain-branching reactions. In reaction dynamic studies, the products and branching ratios of the reaction channels are quite different because of the competition among different channels after oxygen atom insertion. In these reactions, the dominant reaction mechanism is the oxygen atom insertion into C-H bond, while direct H abstraction in the OH formation channel and oxygen atom insertion into the C-C bond are also important in some cases for larger hydrocarbons.

Unfortunately, for even a simple unsaturated hydrocarbon such as  $C_2H_2$ , the multi-channel dynamics of its reaction with  $O(^1D)$  and the kinetics of subsequent reactions are not well understood. Limited relevant studies are available: Shaun A. C. experimentally studied the reaction between  $O(^1D)$  and  $C_2H_2$  and reported the rate constant which was adopted in the current model.<sup>16</sup>

Yvan and Patrick theoretically studied the primary addition reactions of O(<sup>1</sup>D) with C<sub>2</sub>H<sub>2</sub> without considering secondary bond cleavage reactions, and the secondary products were estimated to be H, C<sub>2</sub>OH, OH, C<sub>2</sub>H and HCCO.<sup>17</sup> Based on this estimation, four potential reaction channels are considered in the current reaction mechanism:



In addition, the radicals formed during these secondary chain-branching reactions initialize further chain-propagation and chain-branching reactions. Formation of radicals (i.e. OH and HO<sub>2</sub>) due to the subsequent hydrocarbon oxidation reactions are therefore extremely important. Understanding the multi-channel dynamics and subsequent reaction kinetics of O(<sup>1</sup>D) with C<sub>2</sub>H<sub>2</sub> is critical to understanding O(<sup>1</sup>D) kinetics with larger unsaturated fuels and unburned hydrocarbons typically containing alkenes and aldehydes.

Time-resolved diagnostics of radical species is crucial in kinetic studies due to the short-lived nature of free radicals. Many spectroscopic techniques have been adopted for radical species detection, the most notable examples include ultraviolet (UV) and mid-infrared (IR) tunable laser absorption spectroscopy (TDLAS), laser-induced fluorescent spectroscopy (LIF), and cavity ring-down spectroscopy (CRDS).<sup>18-24</sup> For LIF, HO<sub>2</sub> is measured indirectly through chemical conversion to OH, and in many cases uncontrolled reactions involving the peroxy radicals

introduce significant experimental uncertainty. Moreover, the sensitivity is limited by fluorescence quenching and uncertainty in OH detection.<sup>21-22</sup>

For CRDS, the acquisition interval is limited by the time required to have sufficient intensity build-up in the cavity and radical quenching on the sampling nozzle.<sup>23</sup> TDLAS in principle can provide much better time resolution; however, sensitive measurement of radicals (e.g. RO<sub>2</sub> and HO<sub>2</sub>) in both CRDS and TDLAS is challenging due to spectral contaminations from H<sub>2</sub>O<sub>2</sub>, H<sub>2</sub>O, and other hydrocarbons in the mid-IR and UV regions<sup>25-26</sup>. Typically, the efficacy of TDLAS depends on identifying a spectral region free from unwanted absorption features, which becomes challenging in higher pressure experiments, and is compounded when using larger fuels with broadband spectral features.<sup>18, 23, 27</sup>

To this end, Brumfield et al. first applied Faraday rotation spectroscopy (FRS) for HO<sub>2</sub> detection at the exit of a laminar flow reactor<sup>27</sup>. A modulated magnetic field was used to improve the detection limit and remove spectral interference. However, no time dependent measurements were attempted due to the slow field modulation.

In this work, we aim to study the multi-channel dynamics of excited singlet oxygen atom O(<sup>1</sup>D) reactions with C<sub>2</sub>H<sub>2</sub> and the kinetics of subsequent reactions via selective and time resolved detection of HO<sub>2</sub> and OH radicals by using FRS in a balanced detection scheme<sup>28</sup> to effectively cancel common-mode spectral interference.

## 2. Experimental Section

In this study, we developed a new photolysis flow reactor to study the reactions between O(<sup>1</sup>D) and C<sub>2</sub>H<sub>2</sub> by combining mid-infrared FRS with direct UV-IR TDLAS. (Figure 1).

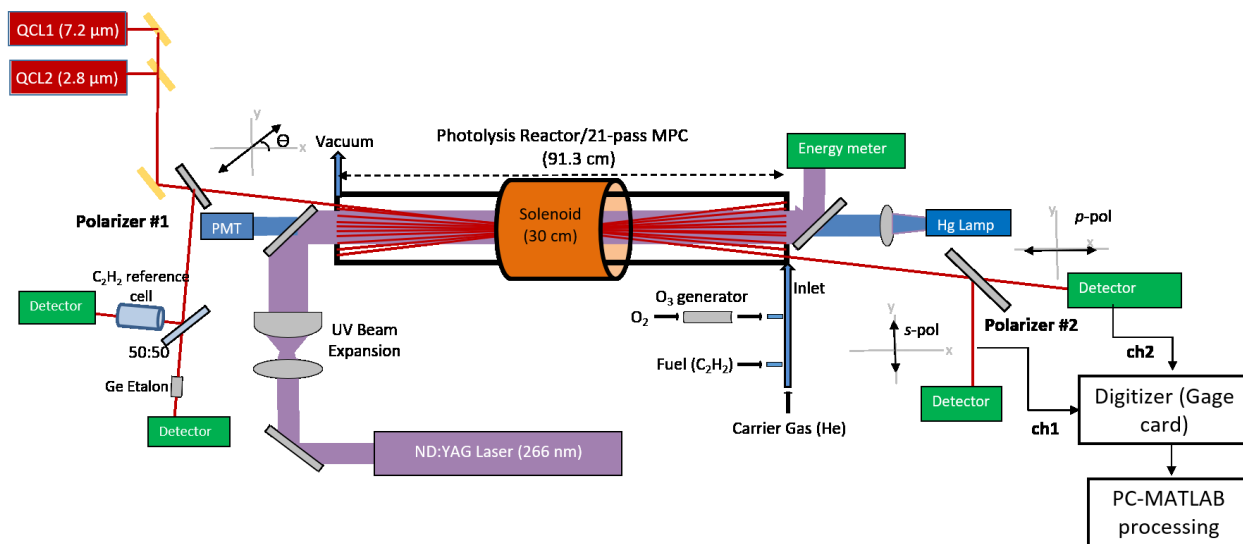


Figure 1. Experimental setup schematic.

The reactor is made of quartz tubing with an inner diameter of 56 mm. A pair of spherical mirrors with 250 mm focal length are installed at the both ends of the reactor 913 mm apart to form the multi-pass configuration (MPC) for the QCLs (7.2  $\mu\text{m}$ , Thorlabs and 2.8  $\mu\text{m}$ , Nanoplus) beams. A protective gold coating extends 8 mm from the edge of the mirror, leaving a 20° section transparent for the IR probe beam to couple into the photolysis cell. The uncoated central part of the spherical mirrors, 40 mm in diameter, allows for the UV photolysis beam to pass through the cell. Two dichroic mirrors (OMEGA OPTICAL) mounted at 45° are placed at the entrance and exit of the flow reactor to allow coupling of both UV photolysis beam and O<sub>3</sub> monitor beam (HG-1 Ocean Optics) in a single-pass configuration. A beam expander consisting of a UV coated plano-concave lens ( $f = -30$  mm) and a UV coated plano-convex lens ( $f = 200$  mm) were used to improve the homogeneity of the photolysis beam. The measured UV beam profile indicates that the uniformity across the reactor cross-section is  $\pm 8.4$  % from the mean value.

Kinetic model is developed based on HP-Mech (C0 – C2)<sup>29-30</sup> with electronically excited species sub-mechanism. The updated HP-Mech was attached in the supplementary material, in which the

O(<sup>1</sup>D) and electronically excited oxygen molecule O<sub>2</sub>(<sup>1</sup>Σ) quenching reactions and reactions with O<sub>3</sub>, H<sub>2</sub>, H<sub>2</sub>O<sub>2</sub>, CH<sub>4</sub>, H<sub>2</sub>O, CH<sub>2</sub>O, H, H<sub>2</sub>O, CH<sub>3</sub> and C<sub>2</sub>H<sub>2</sub> were considered. The half-life of diffusion time for major radicals and atoms out of photolysis beam (50 ms – 100 ms) are more than five times larger than HO<sub>2</sub> radicals decay time (ca. 10 ms) and almost two orders longer compared with OH radical decay time (ca. 1 ms). Thus, the photolysis reactor could be simulated by using closed homogeneous batch reactor (0-D) in Chemkin.

The absolute concentration of O(<sup>1</sup>D) was quantified by measuring the photon flux inside the reactor<sup>1</sup> using the absorption cross-section and the photolysis reaction pathway of O<sub>3</sub> at 266 nm at 60 Torr and 296 K. Ozone concentration is measured by a photomultiplier tube (Hamamatsu R7154) with amplifier (Hamamatsu C6271) using DAS signals from the monochromator (Acton 2500i). To avoid interferences from residual photolysis light or fluorescence from any surface, a 254 nm mercury line filter (Semrock) is installed before the monochromator. The major source of uncertainty in determining absolute concentration of O(<sup>1</sup>D) is quantifying the photolysis laser intensity inside the reactor. Therefore, we used *in-situ* laser light actinometry based on ozone depletion, which is relatively stable at these conditions, with detailed characterization of the cross-section at the monitoring wavelength of 253.65 nm (Mercury line).<sup>31</sup> Ozone depletion in this measurement occurs in two stages, as seen in Figure 2. The first stage is direct photolysis depletion: O<sub>3</sub> + hν → O(<sup>1</sup>D) + O<sub>2</sub> and O<sub>3</sub> + hν → O(<sup>3</sup>P) + O<sub>2</sub>. The second stage is chemical reaction consumption: a decay in ozone concentration due to excited oxygen atoms O(<sup>1</sup>D) and excited oxygen molecules O<sub>2</sub>(<sup>1</sup>Σ) interaction with O<sub>3</sub>. The main reactions for stage 2 are: O(<sup>1</sup>D) + O<sub>2</sub> → O(<sup>3</sup>P) + O<sub>2</sub>(<sup>1</sup>Σ), O<sub>2</sub>(<sup>1</sup>Σ) + O<sub>3</sub> → O(<sup>3</sup>P) + 2O<sub>2</sub> and O(<sup>1</sup>D) + O<sub>3</sub> → 2O<sub>2</sub>. The photon fluence F (photons cm<sup>-2</sup>) can be calculated based on the absorbance (Abs):

$$F = \frac{Abs}{\sigma_{O3,253.65} * L * [O_3] * \sigma_{O3,266}} \quad (5)$$

where the ozone absorption cross-section at 253.65 nm and 266 nm are

$$\sigma_{O3,253.65} = (1.136 \pm 0.012) \times 10^{-17} \text{ cm}^2 \text{ molecule}^{-1} \quad (6)$$

$$\sigma_{O3,266} = (9.04 \pm 0.64) \times 10^{-18} \text{ cm}^2 \text{ molecule}^{-1} \quad (7)$$

$L = 913 \text{ mm}$  is the length of the reactor, and  $[O_3]$  is the concentration. Actinometry measurements were performed before spectroscopy experiments to provide the absolute concentration of  $O(^1D)$  atoms.

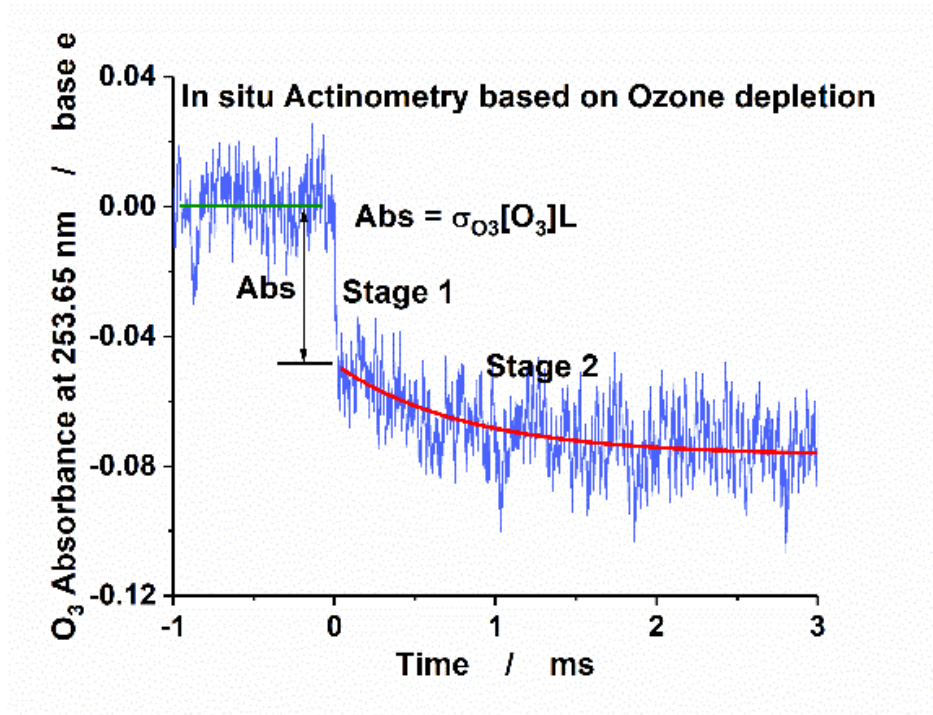


Figure 1. *In-situ* actinometry based on ozone depletion. Ozone is detected at the monitoring wavelength of 253.65 nm (Mercury line). Ozone depletion in this measurement occurs in two stages. The first stage is direct photolysis depletion:  $O_3 + h\nu \rightarrow O(^1D) + O_2$  and  $O_3 + h\nu \rightarrow O(^3P) + O_2$ . The second stage is chemical reaction consumption: a decay in ozone concentration due to excited oxygen atoms  $O(^1D)$  and excited



oxygen molecules  $O_2(^1\Sigma)$  interaction with  $O_3$ . Green line is linear fitting before UV photolysis. Red line is fitted by HP-Mech which photon fluence is fitting parameter.

In the presence of a paramagnetic species such as  $HO_2$ , magnetically induced circular birefringence (MCB) causes rotation of the polarization plane of linearly polarized light, which leads to the FRS signal. As Figure 1 shows, a solenoid is used to generate an axial magnetic field (370 Gauss) along the direction of laser propagation within the flow reactor. A pair of wire-grid polarizers are used: one defines the incident linear polarization, and the second serves as an analyzer to convert laser polarization rotation into intensity changes that can be detected using photodetectors. In this case, the analyzer (Polarizer #2) axis is rotated at an angle of  $45^\circ$  with respect to the initial polarizer (Polarizer #1), hence the exit beam is split into the  $s$  and  $p$  orthogonal polarizations in the transmitted and reflected beam. In this balanced detection configuration, the FRS signal is obtained from the differential measurement between the  $s$  and  $p$  polarizations while other signals, including non-paramagnetic absorptions, appear as common mode signal.<sup>32-33</sup>

The advantage of this detection scheme is demonstrated in Figure 3, where we compare the TDLAS absorption spectra and FRS spectra observed at  $7.2\ \mu\text{m}$  under the same experimental conditions. Figure 3(a) shows the saw-tooth laser frequency scanning across the  $HO_2$  transitions at a rate of 10 kHz, leading to an acquisition time of 100  $\mu\text{s}$  for each spectrum. A total of 5 consecutive spectra are shown here, with ozone photolysis initiated after the second ramp. The TDLAS spectra is shown in panel (b), where 4  $C_2H_2$  lines are clearly visible and labeled in the figure. After the photolysis reactions, the  $HO_2$  absorption signal appears near a dominant  $C_2H_2$  transition as also labeled in panel (b). At 60 Torr, the  $HO_2$  signal is still distinguishable from  $C_2H_2$  but is already affected by the  $C_2H_2$  lines. This spectral interference is effectively suppressed in the

FRS signal as shown in panel (c), where HO<sub>2</sub> formation is clearly visible after photolysis starts. HO<sub>2</sub> concentration is determined by least mean square fitting to the FRS spectrum using HITRAN line parameters. Due to the use of helium as the bath gas, collisional line-broadening parameters deviate from the air-broadening coefficients provided by HITRAN. Based on a well-calibrated frequency axis using a Germanium etalon, the actual broadening coefficient is determined during the nonlinear fitting procedure based on the best fit to the measured spectrum.

It should be noted that at higher pressures, collisional line broadening will result in more severe overlap between this interference and the HO<sub>2</sub> absorption signal, therefore the immunity to spectral interference from non-paramagnetic species makes FRS a uniquely suitable technique for *in-situ* measurements of radical species over a wide range of experimental conditions (e.g. with varying pressure and different fuels). In addition, given the clean signal spectrum achievable using FRS, spectral line-fitting is not required to isolate the HO<sub>2</sub> signal, which would be critical in the case shown in panel (b). Hence line-locked harmonic detection can be performed using fast laser wavelength modulation (e.g. at 100 kHz) to further improve sensitivity and time resolution of FRS detection.

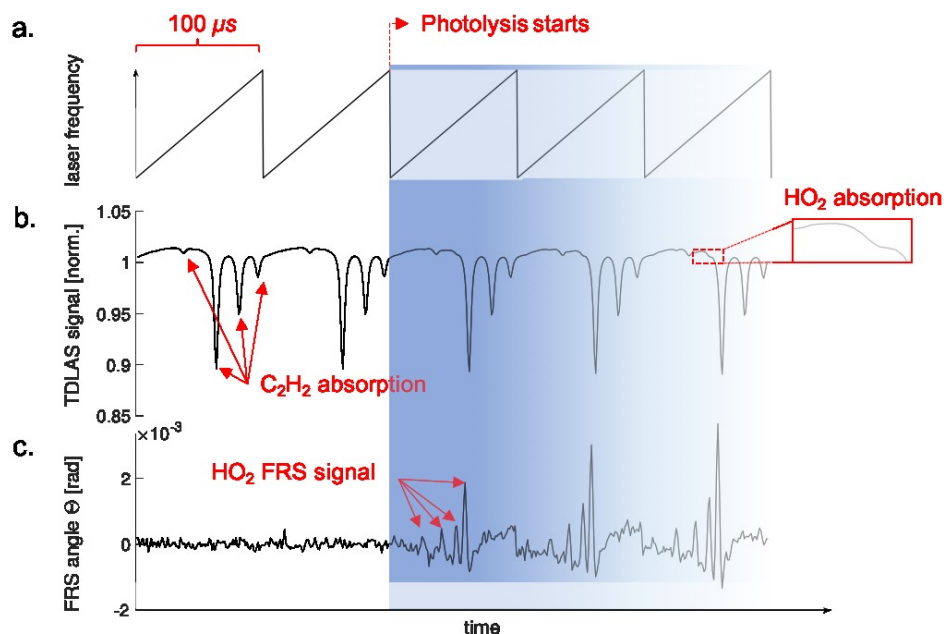


Figure 3. Illustration of spectral interference removal using FRS. (a) Continuous linear laser frequency scanning at a rate of 10 kHz from  $1396.8\text{ cm}^{-1}$  to  $1396.95\text{ cm}^{-1}$ . A total of 5 consecutive spectra are recorded with UV photolysis initiated after the second ramp. (b) TDLAS spectra containing both  $\text{C}_2\text{H}_2$  and  $\text{HO}_2$  absorption signal with prominent  $\text{C}_2\text{H}_2$  features in the same spectral region. A zoomed view of  $\text{HO}_2$  absorption is shown to emphasize the magnitude of spectral interference compared to the  $\text{HO}_2$  signal using TDLAS. c) FRS spectra measured in the same experimental condition as (b). Note the absence of non-paramagnetic  $\text{C}_2\text{H}_2$  absorptions.

In addition to  $\text{HO}_2$ , time-resolved concentration measurements of OH and  $\text{H}_2\text{O}$  are also required for modeling the chemical reaction pathways. In the current setup, another QCL at  $2.8\text{ }\mu\text{m}$  is used for measuring OH and  $\text{H}_2\text{O}$  generated during the reactions following ozone photolysis. Since negligible spectral interference between OH and nearby water lines is observed at the current pressure conditions (60 Torr), TDLAS is used for detecting both OH and  $\text{H}_2\text{O}$ . QCL laser frequency fluctuations up to 0.003 nm is observed during the OH measurements. Since this

fluctuation is 4 times smaller than the full width at half maximum (FWHM) of OH absorption profiles at 60 Torr and 296 K according to HITRAN simulations, insignificant influence to concentration retrieval is assumed. The absorption cross-section of the OH radical at 3568.52 cm<sup>-1</sup> is calibrated using well studied chemical reactions  $O_3 + h\nu \rightarrow O(^1D) + O_2$  and  $O(^1D) + H_2O \rightarrow 2 OH$ . The obtained absorption cross section is:

$$\sigma_{OH,3568.52cm^{-1}} = (3.1 + 0.5) \times 10^{-18} \text{ cm}^2 \text{ molecule}^{-1} \quad (8)$$

It should be noted, however, that with minor modifications to the current system, OH detection using FRS can be implemented to allow measurements at higher reactor pressures where interference becomes significant. The absorption cross section of H<sub>2</sub>O at 3568.29 cm<sup>-1</sup> is:

$$\sigma_{H_2O,3568.29cm^{-1}} = 2.45 \times 10^{-19} \text{ cm}^2 \text{ molecule}^{-1} \quad (9)$$

at 60 Torr, which is taken from the HITRAN database.<sup>34</sup>

### 3. Results and discussion

In Figure 4(a), the OH radical absorption was measured at 3568.52 cm<sup>-1</sup> at 60 Torr and ambient temperature. A large discrepancy has been found when comparing the experimental data with the original HP-Mech mechanism. However, by including the reaction  $C_2H_2 + OH + M$  (R2), which is a dominant channel in low temperature fuel reaction, good agreement between the experimental results and model prediction is observed.

It is interesting to note in Figure 4(a) that the peak of OH production is observed 200 μs after photolysis, which indicates that the channel  $O(^1D) + C_2H_2 \rightarrow C_2H + OH$  (R1d) via H abstraction reaction channel forming OH can be neglected because these channels would result in much faster OH rise (ca. 1 μs). This result is different from the previous studies of  $O(^1D)$  reaction with saturated hydrocarbons,<sup>10, 15, 35</sup> suggesting the uniqueness of  $O(^1D)$  reaction with unsaturated

hydrocarbons. According to the path flux analysis based on the reaction mechanism, channels 1b and 1c, as mentioned in the introduction, have similar sensitivities of OH formation, while channel 1a also displays significant sensitivity but in the opposite direction, to consume OH. As a result, combining a rough adjustment of the ratio of  $k_{1a}/(k_{1b}+k_{1c})$  and a fine adjustment of the ratio of  $k_{1b}/k_{1c}$ , the agreement between simulation and experimental data was greatly improved. The derived branching ratios of channels 1a, 1b and 1c are 56%, 22% and 22%, respectively.

Figure 4(b) shows the  $H_2O$  formation measured using TDLAS at  $3568.29\text{ cm}^{-1}$  to validate the updated kinetic mechanism (HP-Mech) which including the branching ratios of  $O(^1D)$  reaction with  $C_2H_2$  and key elementary reactions shown in table 1. The simulation agrees well with experimental measurement, especially during the initial  $H_2O$  rise but slightly overpredicts the  $H_2O$  concentration at around 1 ms. Signal fluctuation is observed from the minor drifting of the laser wavelength which can not subtracted by the background measurements.

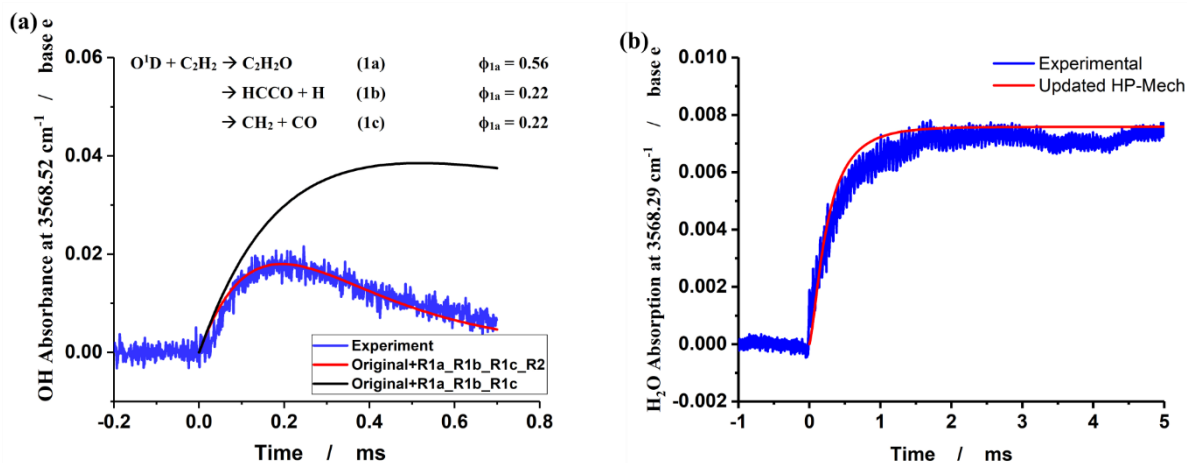


Figure 4. Comparison of the experimental data and simulation based on updated HP-Mech. (a) The time-resolved measurement of OH radicals. The critical reaction  $C_2H_2 + OH + M$  (R2) was missing in the HP-Mech which is the dominant pathway for OH consumption. In this experimental condition, the recommendation channel is dominant (branching ratio  $\phi_{1a} = 56\%$ ), while channel  $O(^1D) +$

$\text{C}_2\text{H}_2 \rightarrow \text{C}_2\text{H} + \text{OH}$  (R1d) can be neglected. (b) The time-resolved measurement of  $\text{H}_2\text{O}$  was measured at  $3568.29 \text{ cm}^{-1}$  wavenumber. There is good agreement between the experimental measurement and theoretical simulation based on the updated HP-Mech.

The path flux analysis indicated that the original mechanism constructed using HP-Mech and ozone measurements with additional reactions involving  $\text{O}(^1\text{D})$  has little sensitivity to the production of  $\text{HO}_2$ , suggesting the absence of some critical reaction channels for  $\text{HO}_2$ . Comparing with vinyl radicals,  $\text{C}_2\text{H}_3$ , 2-hydroxyvinyl radicals,  $\text{C}_2\text{H}_2\text{OH}$  has similar structure which is missing in most combustion models. According to the reaction between  $\text{C}_2\text{H}_3$  and  $\text{O}_2$ , the reaction between 2-hydroxyvinyl radicals,  $\text{C}_2\text{H}_2\text{OH}$ , and  $\text{O}_2$ , might produce  $\text{HCO}$  and  $\text{HCOOH}$ :



In fact, Figure. 5 shows dramatically improved agreement between  $\text{HO}_2$  measurements and the model prediction after including reaction R3. The reaction R3 is directly sensitive towards the  $\text{HCO}$  radical production, which is one of the most important intermediates in  $\text{HO}_2$  formation. Using the  $\text{HO}_2$  temporal profile, the corresponding rate constant of R3 was determined to be  $k_{\text{R3}} = (5 \pm 2.5) \times 10^{11} \text{ cm}^3 \text{ mole}^{-1} \text{ s}^{-1}$  at 296 K and 60 Torr.

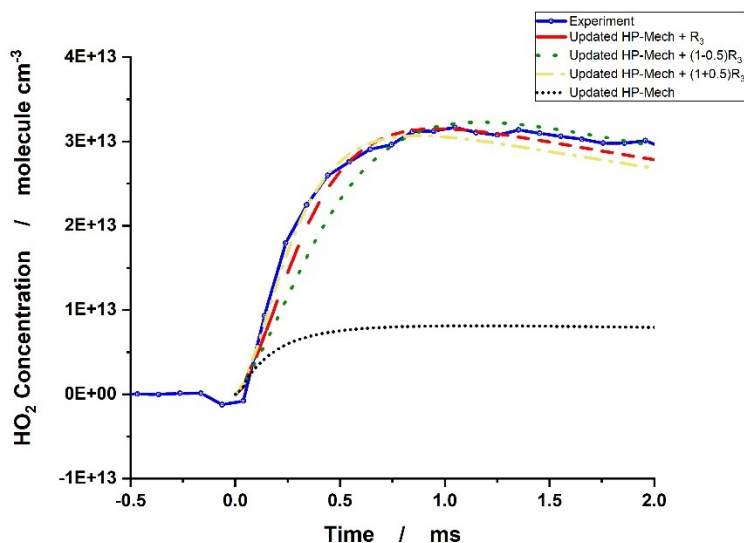


Figure 5. The temporal profile of HO<sub>2</sub> concentration with a gas mixture of C<sub>2</sub>H<sub>2</sub>/O<sub>2</sub>/O<sub>3</sub>/He. A key elementary reaction R3 (C<sub>2</sub>H<sub>2</sub>OH+O<sub>2</sub>) dominates the HO<sub>2</sub> formation. The rate constant is determined by comparing the experimental profile with the updated HP-Mech simulation.

The branching ratios and rate constants of key reactions examined in this study are summarized in Table 1.

Table 1: Rate constants of key elementary reactions used in the simulation. The detailed HP-Mech can be found in the supporting materials. Units are in moles, cubic centimeters, seconds and cal/mole.

<i>Number</i>	<i>Reaction</i>	<i>A</i>	<i>n</i>	<i>Ea</i>	<i>Comments</i>
<i>R1a</i>	O( <sup>1</sup> D)+C <sub>2</sub> H <sub>2</sub> →HCCOH	1.0e14	0	0	56% of k <sub>1total</sub>
<i>R1b</i>	O( <sup>1</sup> D)+C <sub>2</sub> H <sub>2</sub> →HCCO+H	4.0e13	0	0	22% of k <sub>1total</sub>
<i>R1c</i>	O( <sup>1</sup> D)+C <sub>2</sub> H <sub>2</sub> →CH <sub>2</sub> +CO	4.0e13	0	0	22% of k <sub>1total</sub>
<i>R2</i>	C <sub>2</sub> H <sub>2</sub> +OH+M→C <sub>2</sub> H <sub>2</sub> OH+M	2.28e19	0	1071	Ref. <sup>36-37</sup>
<i>R3</i>	C <sub>2</sub> H <sub>2</sub> OH+O <sub>2</sub> →HCO+HCOOH	5e11	0	0	This study

#### 4. Conclusion

In conclusion, the kinetics of O(<sup>1</sup>D) reactions with C<sub>2</sub>H<sub>2</sub> are studied through quantitative time-resolved measurements of HO<sub>2</sub>, OH, H<sub>2</sub>O, and O(<sup>1</sup>D) in a newly developed photolysis flow reactor. In particular, selective and time-resolved detection of HO<sub>2</sub> was successfully demonstrated through use of Faraday rotation spectroscopy by the suppression of spectral interference from non-paramagnetic hydrocarbon absorption via balanced detection. The branching ratios of O(<sup>1</sup>D) reactions with C<sub>2</sub>H<sub>2</sub> as well as its subsequent reaction kinetics are quantified. It was found that differing from O(<sup>1</sup>D) reaction with saturated hydrocarbons, the OH formation channel via direct H abstraction is negligible in O(<sup>1</sup>D) reactions with unsaturated C<sub>2</sub>H<sub>2</sub>. This finding suggests the

uniqueness of  $O(^1D)$  reaction with unsaturated hydrocarbons. Moreover, the retrieved  $HO_2$  temporal profile allows for identification and determination of the rate constant of a missing reaction channel for  $HO_2$  production involving  $C_2H_2OH$  with  $O_2$ . Furthermore, with the time-dependent OH measurement, another missing reaction channel via  $C_2H_2+OH+M \rightarrow C_2H_2OH+M$  was also identified.

With the adoption of these reaction channels and the measured reaction rates and branching ratios, an updated kinetic model for  $O(^1D)$  reactions with  $O(^1D)$  based on HP-Mech is developed. The updated model significantly improved the prediction of experimental results of OH,  $HO_2$ , and  $H_2O$ . The results also demonstrate that the new experimental apparatus combining TDLAS with FRS in the photolysis reactor has the potential for sensitive measurements of many species. (e.g. OH,  $HO_2$ ,  $H_2O$ ,  $CH_2O$ ,  $C_2H_2$ , and  $O_3$ ) for the study more complicated chemical kinetics and dynamics involving  $O(^1D)$  reaction with large unsaturated hydrocarbons and oxygenated fuels. Furthermore, understanding of  $O(^1D)$  reactions with fuels provides a key sub-mechanism of reaction kinetics for combustion, plasma chemistry, materials synthesis, and atmospheric chemistry.

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## Supporting Information:

Experimental conditions; Updated kinetic model based on the modified HP-Mech.



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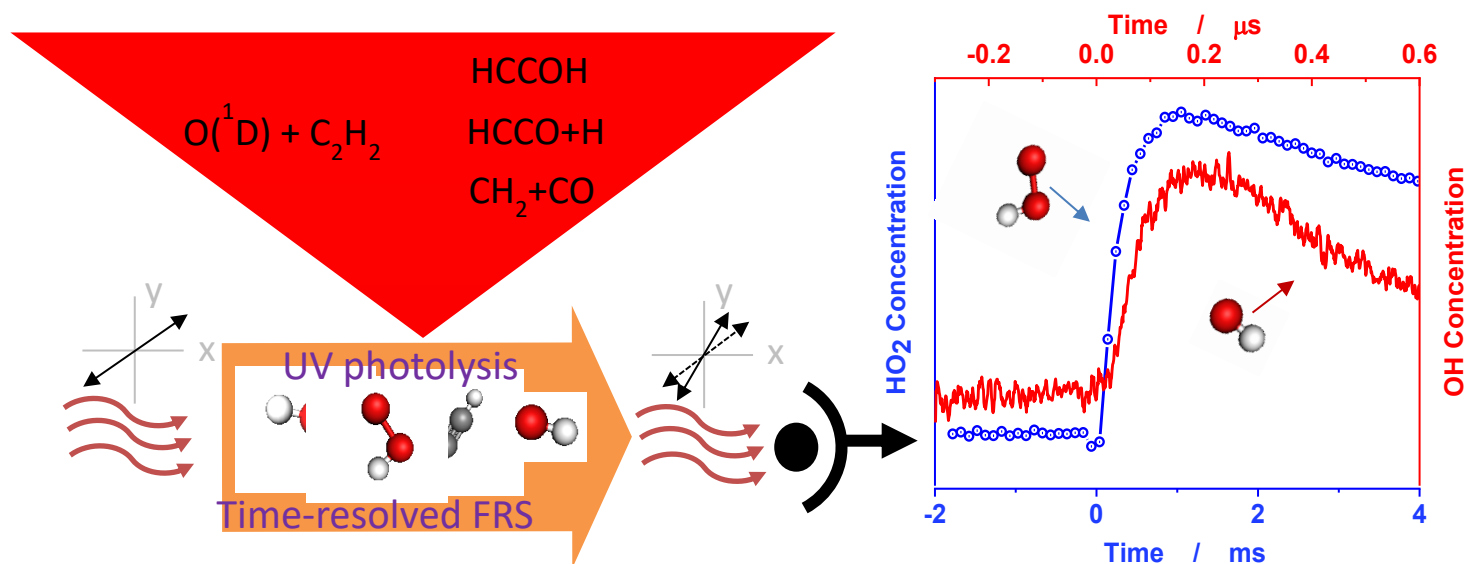
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### **Notes**

The authors declare no competing financial interests.

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