

Subscriber access provided by CARNEGIE MELLON UNIV

Surfaces, Interfaces, and Applications

# **Enhancing Thermal Interface Conductance** to Graphene Using Ni-Pd Alloy Contacts

Dipanjan Saha, Xiaoxiao Yu, Yanhao Du, Zhitao Guo, Feng Xiong, Andrew J Gellman, and Jonathan A Malen ACS Appl. Mater. Interfaces, Just Accepted Manuscript • DOI: 10.1021/acsami.0c06953 • Publication Date (Web): 01 Jul 2020

Downloaded from pubs.acs.org on July 8, 2020

# **Just Accepted**

"Just Accepted" manuscripts have been peer-reviewed and accepted for publication. They are posted online prior to technical editing, formatting for publication and author proofing. The American Chemical Society provides "Just Accepted" as a service to the research community to expedite the dissemination of scientific material as soon as possible after acceptance. "Just Accepted" manuscripts appear in full in PDF format accompanied by an HTML abstract. "Just Accepted" manuscripts have been fully peer reviewed, but should not be considered the official version of record. They are citable by the Digital Object Identifier (DOI®). "Just Accepted" is an optional service offered to authors. Therefore, the "Just Accepted" Web site may not include all articles that will be published in the journal. After a manuscript is technically edited and formatted, it will be removed from the "Just Accepted" Web site and published as an ASAP article. Note that technical editing may introduce minor changes to the manuscript text and/or graphics which could affect content, and all legal disclaimers and ethical guidelines that apply to the journal pertain. ACS cannot be held responsible for errors or consequences arising from the use of information contained in these "Just Accepted" manuscripts.

# **Enhancing Thermal Interface Conductance to Graphene Using Ni-Pd Alloy Contacts**

Dipanjan Saha<sup>1</sup>, Xiaoxiao Yu<sup>2</sup>, Yanhao Du<sup>3</sup>, Zhitao Guo<sup>2</sup>, Feng Xiong<sup>3</sup>, Andrew Gellman<sup>2,4</sup>, Jonathan A Malen<sup>1\*</sup>

<sup>1</sup>Department of Mechanical Engineering, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA

<sup>2</sup>Department of Chemical Engineering, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA

<sup>3</sup>Department of Electrical and Computer Engineering, University of Pittsburgh, Pittsburgh, Pennsylvania 15261, USA

<sup>4</sup>W.E. Scott Institute for Energy Innovation, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA

Keywords: 2D materials, thermoreflectance, alloy thin films, miscibility gap, TEM images

### **Abstract**

To identify superior thermal contacts to graphene we implement a high throughput methodology that systematically explores the Ni-Pd alloy composition spectrum and the effect of Cr adhesion layer thickness on the thermal interface conductance with monolayer CVD graphene. Frequency domain thermoreflectance measurements of two independently prepared Ni-Pd/Cr/graphene/SiO<sub>2</sub> samples both identify a maximum in the metal/graphene/SiO<sub>2</sub> junction thermal interface conductance of 114± (39, 25) MW/m<sup>2</sup>K and 113± (33, 22) MW/m<sup>2</sup>K at ~10 atomic percent Pd in Ni—nearly double the highest reported value for pure metals and three times that of pure Ni or Pd. The presence of Cr, at any thickness, suppresses this maximum. Although the origin of the peak is unresolved, we find that it correlates to a region of the Ni-Pd phase diagram that exhibits a miscibility gap. Cross sectional imaging by high resolution transmission electron microscopy identifies striations in the alloy at this particular composition, consistent with separation into multiple phases. Through this work, we draw attention to alloys in the search for better contacts to 2D materials for next generation devices.

### Introduction

Heat dissipation into metal contacts is critical to the high-performance operation of short channel graphene devices. <sup>1-6</sup> Pure metal thermal contacts to graphene can spread heat laterally, but their poor thermal interface conductance (G) is a bottleneck to the removal of heat generated in the graphene channel and at contact interfaces. <sup>7–9</sup> Heat transfer across metal-graphene interfaces is phonon dominated. <sup>10,11</sup> Alignment of the phonon spectra across the interface is one consideration in the search for high G. The Debye temperature characterizes the thermal activation of a lattice's highest energy phonons. Graphene has a high Debye temperature of 1287 K<sup>12</sup> and thus, metals that also possess a high Debye temperature should be favorable thermal contacts. As exhibited in Figure 1 by the spread in G values for pure metal/graphene/SiO<sub>2</sub> junctions (where the junction is composed of graphene and its two interfaces, since an individual interface cannot be resolved) <sup>11,13–15</sup>, there are additional factors beyond phonon spectrum alignment. One major factor is the adhesion strength of the metal to graphene. For example Hopkins et al., were able to double G in their study by functionalizing the surface of monolayer graphene with oxygen atoms for better bonding to the Al metal contact. 15 Zheng et al. found that strongly bound TiN contacts offered high thermal interface conductance with graphene and measured a value of 135 MW/m<sup>2</sup>K for the TiN/graphene/TiN junction (cannot be compared directly with metal/graphene/SiO<sub>2</sub> junctions in Figure 1). Nonetheless, the inherently low thermal conductivity of TiN thin films (7.5 - 11) $W/m/K^{16,17}$ ) may limit their utility as heat spreaders.

While maximizing adhesion strength may seem desirable, too strong an interaction can result in the metal reacting with the underlying graphene to form a carbide, which is detrimental to the electrical performance of the device. Such a strong interaction is classified as chemisorption (where the metal is chemically bonding to the graphene) in contrast to physisorption

(a weaker interaction where the metal is in physical contact with the graphene). Al, Ag, Cu, Au, Pt, Cr are metals recognized as physisorbing to graphene; while Co, Ti, Ni, and TiN are chemisorbing; and Pd is described in literature as very weakly chemisorbing/strongly physisorbing. These distinctions are based on the binding energy and equilibrium bonding distance between the metal and graphene as calculated by first-principle density functional theory methods. <sup>20</sup>

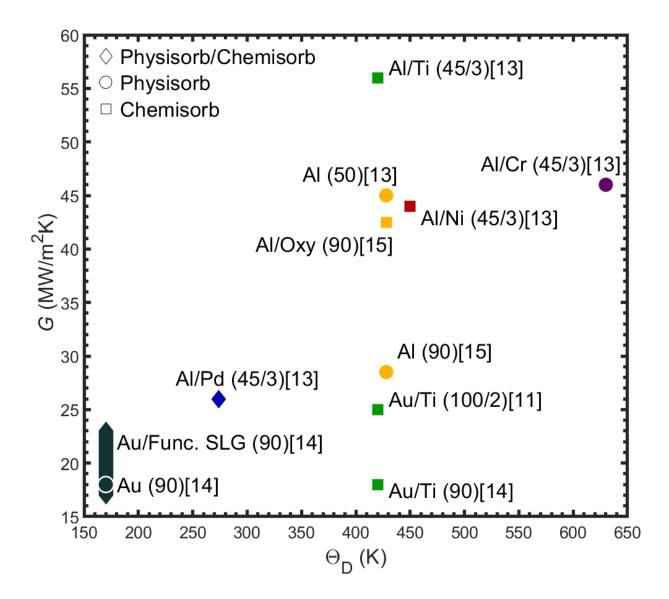


Figure 1. Thermal interface conductance (G) vs Debye temperature ( $\Theta_D$ ) for junctions composed of metal/monolayer graphene/SiO<sub>2</sub>.<sup>11,13–15</sup> The plotted  $\Theta_D$  values correspond to that of the metal in contact with the graphene (grouped by

colors). Al/Oxy refers to an Al thermal contact on oxygen functionalized monolayer graphene.<sup>15</sup> The range of data for Au/functionalized single layer graphene (SLG) encompasses functionalization by oxygen, fluorine, and nitrogen to varying degrees of graphene surface coverage.<sup>14</sup> Parentheses denote thicknesses of metals and adhesion layers in nanometers. Circles denote physisorbing metals, squares chemisorbing, and the diamond marker denotes strongly physisorbing/weakly chemisorbing.<sup>19–25</sup>

The limited choice of pure metals motivates research into whether alloy composition can tune the interaction strength to optimize interfacial heat transport. The ideal interaction would maximize adhesion without damaging the graphene for electrical applications. Unfortunately, optimization studies on alloy contacts to graphene are complicated by the overwhelming nature of the alloy design space.

An additional design consideration is adhesion layer thickness, which has major implications for thermal transport. For example, Jeong et al. reported that *G* between Au and Al<sub>2</sub>O<sub>3</sub> increased from 60±10 MW/m<sup>2</sup>K to 300±50 MW/m<sup>2</sup>K as the thickness of a Cr adhesion layer increased from 0 to 1 nm.<sup>26</sup> For reference, the two most common graphene adhesion layers, Cr and Ti, are used with thicknesses spanning 0.5 to 10 nm according to the electrical graphene contact resistance literature.<sup>6,27–31</sup> Clearly a lack of consensus exists as far as the optimal adhesion layer thickness for graphene contacts.

Utilizing high throughput techniques, we define and implement a new methodology to systematically investigate the relatively unexplored frontier of alloy-graphene contacts, as well as the effect of adhesion layer thickness on G. Specifically, we present measurements of G as a function of alloy composition for a Ni-Pd alloy and thickness for a Cr adhesion layer. Ni-Pd alloys were studied because the spectrum of chemisorption (Ni) to strong physisorption (Pd) is represented. Additionally, Ni and Pd are two of the best electrical graphene contacts.  $^{29}$ 

With this initial study, we find that ~10 at. % Pd in Ni without an adhesion layer yields a maximum thermal interface conductance of  $114 \pm (39, 25)$  MW/m<sup>2</sup>K, double the value of the next highest reported G for a metal/graphene/SiO<sub>2</sub> junction.<sup>13</sup> This major enhancement in G only requires the up-front complexity associated with depositing alloys via evaporation; whether by coevaporation of alloy constituents<sup>32</sup> or evaporation from an alloy target.<sup>33,34</sup> By comparison, other enhancement techniques such as functionalizing the graphene surface achieved 38% of our maximum  $G^{15}$  but can impede electrical transport from the metal to the graphene.<sup>14</sup> While the physics underlying this large G are not yet understood, we observe a clear correlation between the alloy composition where enhancement of G occurs and the existence of a miscibility gap in the alloy phase diagram.

# Results/Discussion

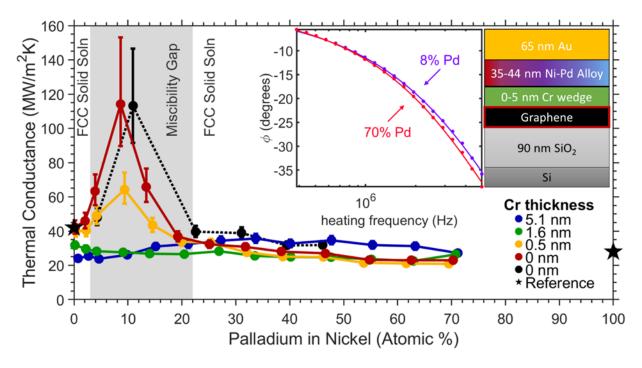


Figure 2. FDTR measurements of thermal interface conductance (G) as a function of Ni-Pd alloy composition and Cr thickness (colored lines). The material stack is shown in the top right inset where G across the metal/graphene/SiO<sub>2</sub>

junction (outlined in red) is reported. A peak in G is observed for zero adhesion layer thickness (red circles), a trend well reproduced by a second sample with no Cr adhesion layer (black circles). A smaller peak is also observable for 0.5 nm of Cr (yellow circles) where the peak value is suppressed relative to 0 nm of Cr. The emergence of the peak clearly correlates with the composition range over which the alloy phase diagram exhibits a miscibility gap (gray shaded region) based on the binary phase diagram calculated for an annealing temperature of  $180^{\circ}$ C.  $^{35,36}$  Data also shows good agreement with the pure metal reference values  $^{13}$  (black stars and shown in Figure 1), and after extrapolating measured data to pure Pd. Vertical error bars represent the  $1^{st}$  and  $3^{rd}$  quartiles in the distribution of G resulting from uncertainty propagation in FDTR fitting parameters (further details in SI). Horizontal error bars estimated as  $\pm 2$  atomic % Pd in Ni, are not shown for legibility. Top-central inset shows a clear difference in the phase lag vs frequency thermoreflectance data at the peak conductance alloy composition as compared to 70% Pd in Ni for 0 nm thickness of Cr.

High throughput thermal interface conductance data acquisition is possible by first depositing a Cr metal wedge adhesion layer, 0-5 nm in thickness, onto graphene on 90 nm SiO<sub>2</sub> on Si. Subsequently, opposing Ni and Pd wedges are co-deposited, to produce a metal contact 35 – 44 nm in thickness with alloy composition varying spatially across the lateral dimension of the sample. Capping the sample with a Au transducer layer enables the sample to be scanned spatially using the laser pump/probe technique, frequency domain thermoreflectance (FDTR), to measure G as a function of adhesion layer thickness and alloy composition (as determined by Energy Dispersive X-ray spectroscopy). Sinusoidal modulation of the pump laser creates a periodic temperature change in the sample with an amplitude and phase, relative to the pump, that depend on the unknown value of G. The temperature response is measured by the probe laser via thermoreflectance, and the phase data are fit over a range of modulation frequencies with an analytical solution to the heat diffusion equation in order to determine G. FDTR phase data and fits are plotted in the Figure 2 inset as a function of frequency for 8 and 70 at. % Pd, illustrating the goodness of fit and its sensitivity to G.

As with the literature values shown in Figure 1, the values of G plotted in Figure 2 represent that of the junction composed of monolayer graphene and its two adjacent interfaces (i.e. the metal/graphene/SiO<sub>2</sub> junction highlighted in red in the inset). In the regions where the Cr adhesion layer is thickest (1.6-5.1 nm), no significant effects of the Ni-Pd composition on G were observed. These Cr thicknesses apparently dampen any alloy-graphene interaction. However, measurements of G corresponding to a Cr thickness of 0.5 nm (2-3 atomic layers of Cr) show an increase in G at ~8 at. % Pd in Ni. Most notably, measurements taken in the region where the Ni-Pd alloy is directly in contact with the graphene (no Cr) exhibit a maximum in G of  $114\pm (39, 25)$  MW/m<sup>2</sup>K. Such a large G is double the highest reported thermal interface conductance for a pure metal/monolayer graphene/SiO<sub>2</sub> junction (see Al/Ti contact of Figure 1).<sup>13</sup> Our measurements at 0 at. % Pd and 70 at. % Pd compare well with reference values to within 10% for pure Ni/graphene/SiO<sub>2</sub> and Pd/graphene/SiO<sub>2</sub> junctions<sup>13</sup>. The highest values of G that we observe for Cr/graphene/SiO<sub>2</sub> are 20% lower than the one previously reported value for Cr/graphene/SiO<sub>2</sub>.<sup>13</sup> As evidenced by the range of G we see as a function of Pd content and Cr thickness, this may occur due to interdiffusion driven by the annealing process (in ref. [13] the Cr is capped with Al and the films are not annealed).38

Due to the surprising appearance of a peak in G at an intermediate Ni-Pd composition, a second sample (Sample 2) without a Cr adhesion layer was fabricated to verify reproducibility (black circles). This second sample does not include a Cr adhesion layer since it only suppresses the peak in conductance. A peak of similar magnitude,  $G = 113\pm (33, 22)$  MW/m<sup>2</sup>K, is observed for the second sample at a similar composition of 11 at. % Pd. While some minor misalignment (as a function of alloy composition) in the peaks of the two datasets is observable, this can be attributed to the positions of the Energy Dispersive X-Ray (EDX) spectroscopy and thermal

measurements not being perfectly co-located. Horizontal error bars on the data, estimated as  $\pm 2$  atomic % Pd in Ni, are not shown in Figure 2 for legibility. Regardless, these two separate samples demonstrate that a peak in G exists for  $\sim 10$  at. % Pd in Ni alloy directly in contact with monolayer graphene on SiO<sub>2</sub>. Although the physical origin of the enhancement in G is unresolved, we find that the peak and its less-enhanced neighboring points (red circles) occur at compositions of Ni-Pd where there exists a miscibility gap in the binary phase diagram at our film annealing temperature of 180°C (Figure S1 in the SI depicts the Ni-Pd binary alloy phase diagram<sup>35,36</sup>). The same phenomena is also observed for 0.5 nm of Cr as an adhesion layer (yellow circles).

The existence of a miscibility gap in the Ni-Pd system itself has been historically contentious due to contradicting experimental observations.  $^{39-42}$  For this reason, we present cross-sectional high-resolution transmission electron microscopy (HRTEM) images of the contacts at ~8 at. % Pd in Ni, the composition of the maximum G. For comparison, HRTEM cross-sectional images were also taken at ~60 at. % Pd. These HRTEM images are shown in Figure 3. In Figure 3a striated areas circled in yellow identify the alloy as phase separating, and compare well with the images of phase separated alloys (characterized as spinodal decompositions) presented in the work of Androulakis et al.  $^{43}$  The physical characteristics of the phase separated regions, such as the striation wavelength, depend on the thermal treatment conditions of the alloy.  $^{44}$  Similar striations are not apparent in the ~60 at. % Pd image shown in Figure 3b. If the miscibility gap is responsible for the enhancement of G, then the observed enhancement would not be expected at higher operating temperatures where the alloy phase is a solid solution (i.e., greater than 550 K).

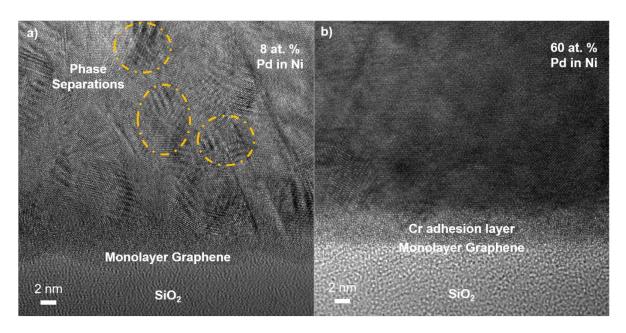


Figure 3. a) Cross sectional HRTEM image of 8 at. % Pd in Ni on monolayer graphene on  $SiO_2$ . The graphene itself is not visible but its relative location in conformance with the  $SiO_2$  morphology is shown. Circled in yellow are three of the multiple regions with visible striations associated with phase separations of the Ni-Pd alloy, demonstrating that the alloy at this composition is in a miscibility gap. b) For comparison, HRTEM image of the sample region where enhancement in G is not observed. Phase separation is not apparent at a composition of 60 at. % Pd in Ni with a 2 nm Cr adhesion layer.

Figure 3 does not visualize the monolayer graphene directly, and thus we cannot rule out the hypothesis that the peak in *G* results from pinholes in the graphene that would allow direct contact between the Ni-Pd alloy and SiO<sub>2</sub>. This hypothesis is motivated by Leong et al., who showed that annealing Ni containing contacts can etch the underlying graphene. Direct measurements of *G* at the interface between Pd and SiO<sub>2</sub> have not been published, but for reference, a recently reported value of *G* of Pt on SiO<sub>2</sub> is 300 MW/m<sup>2</sup>K. Thus parallel heat transfer through 1) pinholes allowing direct alloy/SiO<sub>2</sub> contact and 2) across the targeted alloy/graphene/SiO<sub>2</sub> junction is a possible cause of the observed enhancement. This mechanism would also be expected to produce enhanced *G* for pure Ni, which we do not observe. Cross sectional HRTEM imaging of monolayer graphene has historically been difficult. For example,

Norimatsu and Kusunoki were successful in imaging single layer graphene on SiC with a buffer layer<sup>48</sup> whereas other attempts were not.<sup>49–52</sup> Therefore, the inability to visualize the graphene by our HRTEM imaging should not be prematurely concluded to indicate that it has been compromised relative to its high quality monolayer state prior to metal deposition, as verified by Raman spectroscopy (see Figure S6 in the SI).

### **Conclusion**

Utilizing high throughput techniques, we find that Ni-Pd alloy contacts to graphene exhibit a value of G that is nearly double that of any pure metal contact and over three times that of pure Ni or Pd. The composition of maximum G, ~10 at. % Pd in Ni, clearly correlates with a miscibility gap in the Ni-Pd binary phase diagram. In comparison, the phases of compositions outside of the miscibility gap where G is not enhanced are face centered cubic (FCC) solid solutions. $^{35,36}$ HRTEM images of the peak G alloy composition identify striations that confirm the phase separation of the alloy in agreement with the miscibility gap region of the binary alloy phase diagram, the existence of which was previously debated. While the origin of enhanced G is unknown, this work demonstrates the potential for engineered alloys to form better thermal contacts than pure metals and evokes hope for the relatively unexplored frontier of alloy contacts on 2D materials. If the observed peak in G results from better adhesion to graphene or the creation of pinholes, both of which would be expected to enhance electrical transport<sup>53</sup>, then depositing ~10 at. % Pd in Ni as the metal contacts is a simple method to advance high performance graphene electronics. Further research is required to understand the origin of the enhanced G and whether alloy composition represents a general tool for engineering other interfacial transport properties, such as electrical contact conductance.

#### Methods

# Graphene Acquisition

The samples consist of CVD grown monolayer graphene transferred to 90 nm of SiO<sub>2</sub> on Si, acquired from the commercial supplier Graphene Supermarket. Prior to metal deposition the samples are first annealed at 250°C for four hours in 5% H<sub>2</sub> in Ar by volume (25 sccm H<sub>2</sub> in 475 sccm of Ar in a 2" diameter quartz tube) at atmospheric pressure. Annealing in H<sub>2</sub>/Ar is a common approach to remove organic contaminants from the graphene surface.<sup>54,55</sup> While some polymer residue can still remain<sup>56</sup>, more aggressive surface cleaning techniques present significant risk to damaging the graphene.<sup>27</sup> Raman spectroscopy following the H<sub>2</sub>/Ar annealing procedure validates the graphene as being of sufficiently high quality and monolayer (see Figure S6 in the SI).

# Compositionally Spread Alloy Film (CSAF) Deposition

The metal deposition process onto the annealed graphene involves three steps diagrammed in Figure 4. First, a Cr wedge is deposited ranging in thickness from 0 to 5 nm. Next, in the perpendicular direction of the Cr thickness gradient, Ni and Pd wedges are co-deposited to create an alloy film with composition varying as a function of space – known as a Compositionally Spread Alloy Film (CSAF). Thus with one sample, a high throughput study probing an entire spectrum of alloy compositions can be conducted instead of numerous samples with discrete alloy compositions. Lastly, the sample is coated in 65 nm of Au, a necessary transducer layer for FDTR. The fabrication of a second sample to ensure reproducibility follows a similar procedure without the Cr wedge deposition step. The deposition pressures of the two samples are  $10^{-9}$  and  $10^{-8}$  Torr, respectively, at a rate of 0.2 nm/minute. In order to establish equilibrium phases of the metals, the samples are annealed at  $180^{\circ}$ C for one hour after each deposition step.

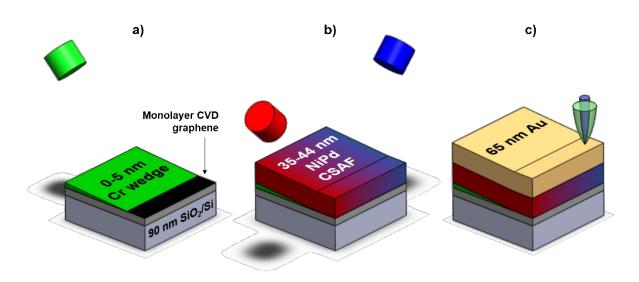


Figure 4. Sample fabrication sequence. a) First a 0-5 nm Cr wedge is deposited onto monolayer CVD graphene/90 nm  $SiO_2/Si$ . b) In the perpendicular direction Ni and Pd wedges are co-deposited to form a 35-44 nm CSAF. c) Final capping with a 65 nm Au film needed for FDTR measurements where co-aligned blue and green lasers scan across the sample to measure G as a function of alloy composition and Cr adhesion layer thickness.

#### Material Characterization

The thicknesses of the metals and alloy composition as a function of space are determined by EDX spectroscopy. A thickness map of Ni, as an example, and the alloy composition map are available in the SI. The Ni wedge shows a relatively linear thickness gradient primarily in one direction. Due to a small spatial offset between the positioning of the graphene sample and the positions of the wedges, the alloy composition ranges from 0 to 70 at. % Pd in Ni.

The EDX system uses a Tescan scanning electron microscope equipped with an Oxford Instruments X-max 80-mm<sup>2</sup> detector. The thicknesses are mapped across a 9 mm x 9 mm area centered on the graphene with a 1 mm grid spacing. The 0-10 keV EDX spectra are measured by rastering a 20 keV electron beam across a 50 x 50 um area at each point. Thicknesses are extracted through spectra fitting done by INCA ThinFilmID software. The fitting procedure models the

layers as Au, Ni-Pd CSAF, Cr (excluded for the second sample), and graphene on an SiO<sub>2</sub>/Si substrate.

### Thermal Characterization

The optical pump-probe technique, FDTR<sup>59</sup>, is used to measure *G* as a function of alloy composition and Cr thickness in contact with graphene. FDTR measurements are collected at positions on the sample corresponding with EDX measurement locations, with linearly interpolated locations also being measured in order to report higher resolution thermal data than composition data. Linear interpolation is appropriate because the spatial thickness gradient of the deposited metal wedges is also linear.

FDTR phase lag data between the reflected pump and probe beams are collected for 20 logarithmically spaced pump modulation frequencies between 300 kHz to 5 MHz. The data is then fit to a solution of the heat diffusion equation for a layered structure<sup>37</sup> composed of Au, Ni-Pd alloy, Cr (excluded for the second sample), the junction interface of interest, 90 nm of SiO<sub>2</sub>, and Si. The junction of interest represents the monolayer graphene and its two adjacent interfaces.

The vertical error bars of our G data represent the 1<sup>st</sup> and 3<sup>rd</sup> quartile in the distribution of G for each alloy composition, resulting from the Monte Carlo method of randomly generating sets of fitting parameter values falling within the normal distribution of each. Further details of the FDTR fitting and example histograms of the Monte Carlo method used for uncertainty analysis can be found in the SI.

Supporting Information: Ni-Pd binary phase diagram, thermal conductivity of Ni-Pd alloy thin

films, FDTR uncertainty analysis, Energy Dispersive X-Ray mapping of alloy thin films, and

Raman spectroscopy of graphene

#### **Contributions to the Work**

Dipanjan Saha: writing, metal deposition, EDX, thermal characterization, electrical device fabrication, electrical characterization, Raman spectroscopy, TEM analysis, 4-point probe device fabrication/measurements

Xiaoxiao Yu: metal deposition, EDX Zhitao Guo: metal deposition, EDX

Yanhao Du: electrical device design/fabrication, electrical characterization, Raman spectroscopy

### Acknowledgements

We would like to acknowledge the contributions of Mohamed Darwish in his help with Raman spectroscopy, designing our 4-point probe devices and Yunus Kesim in his help fabricating/measuring these devices. We would also like to acknowledge the contributions of Noel Tom Nuhfer in the Material Characterization Facility at Carnegie Mellon University, towards the cross-sectional HRTEM images presented in this publication.

D.S. and J.A.M. acknowledge support from the National Science Foundation ENG ECCS 1901972 and Army Research Office Grant W911NF-17-1-0397. Y. D. and F. X. acknowledge support from NSF ENG ECCS 1901864 and MCSI Seed Fund from University of Pittsburgh. X.Y. and A.G. acknowledge support from NSF-CHE1566228 This work was performed, in part, at the Nanoscale Fabrication and Characterization Facility, a laboratory of the Gertrude E. and John M. Petersen Institute of NanoScience and Engineering, housed at the University of Pittsburgh.

#### References

- (1) Geim, A. K.; Novoselov, K. S. The Rise of Graphene. *Nat. Mater.* **2007**, *6*, 183–191. https://doi.org/10.1103/PhysRevB.74.075404.
- (2) Novoselov, K. S.; Fal'Ko, V. I.; Colombo, L.; Gellert, P. R.; Schwab, M. G.; Kim, K. A Roadmap for Graphene. *Nature* **2012**, *490* (7419), 192–200. https://doi.org/10.1038/nature11458.
- (3) Chen, Z.; Jang, W.; Bao, W.; Lau, C. N.; Dames, C. Thermal Contact Resistance between Graphene and Silicon Dioxide. *Appl. Phys. Lett.* **2009**, *95* (16), 93–96. https://doi.org/10.1063/1.3245315.
- (4) Hsu, A.; Wang, H.; Kim, K. K.; Kong, J.; Palacios, T. Impact of Graphene Interface Quality on Contact Resistance and RF Device Performance. *IEEE Electron Device Lett.* **2011**, *32* (8), 1008–1010. https://doi.org/10.1109/LED.2011.2155024.
- (5) Watanabe, E.; Conwill, A.; Tsuya, D.; Koide, Y. Low Contact Resistance Metals for Graphene Based Devices. *Diam. Relat. Mater.* **2012**, *24*, 171–174. https://doi.org/10.1016/j.diamond.2012.01.019.
- (6) Nagashio, K.; Nishimura, T.; Kita, K.; Toriumi, A. Metal/Graphene Contact as a

- Performance Killer of Ultra-High Mobility Graphene Analysis of Intrinsic Mobility and Contact Resistance. *Tech. Dig. Int. Electron Devices Meet. IEDM* **2009**, 1–4. https://doi.org/10.1109/IEDM.2009.5424297.
- (7) Bae, M. H.; Ong, Z. Y.; Estrada, D.; Pop, E. Imaging, Simulation, and Electrostatic Control of Power Dissipation in Graphene Devices. *Nano Lett.* **2010**, *10* (12), 4787–4793. https://doi.org/10.1021/nl1011596.
- (8) Yu, Y. J.; Han, M. Y.; Berciaud, S.; Georgescu, A. B.; Heinz, T. F.; Brus, L. E.; Kim, K. S.; Kim, P. High-Resolution Spatial Mapping of the Temperature Distribution of a Joule Self-Heated Graphene Nanoribbon. *Appl. Phys. Lett.* **2011**, *99* (18), 2009–2012. https://doi.org/10.1063/1.3657515.
- (9) Freitag, M.; Chiu, H. Y.; Steiner, M.; Perebeinos, V.; Avouris, P. Thermal Infrared Emission from Biased Graphene. *Nat. Nanotechnol.* **2010**, *5* (7), 497–501. https://doi.org/10.1038/nnano.2010.90.
- (10) Balandin, A. A.; Ghosh, S.; Bao, W.; Calizo, I.; Teweldebrhan, D.; Miao, F.; Lau, C. N. Superior Thermal Conductivity of Graphene. *Nano Lett.* **2008**, *8* (3), 902–907.
- (11) Koh, Y. K.; Bae, M. H.; Cahill, D. G.; Pop, E. Heat Conduction across Monolayer and Few-Layer Graphenes. *Nano Lett.* **2010**, *10* (11), 4363–4368. https://doi.org/10.1021/nl101790k.
- (12) Tewary, V. K.; Yang, B. Singular Behavior of the Debye-Waller Factor of Graphene. *Phys. Rev. B Condens. Matter Mater. Phys.* **2009**, *79* (12), 1–5. https://doi.org/10.1103/PhysRevB.79.125416.
- (13) Vasquez Guzman, P. A.; Sood, A.; Mleczko, M. J.; Wang, B.; Wong Philip, H. S.; Nishi, Y.; Asheghi, M.; Goodson, K. E. Cross Plane Thermal Conductance of Graphene-Metal Interfaces. *Thermomechanical Phenom. Electron. Syst. -Proceedings Intersoc. Conf.* **2014**, 1385–1389. https://doi.org/10.1109/ITHERM.2014.6892441.
- (14) Foley, B. M.; Hernández, S. C.; Duda, J. C.; Robinson, J. T.; Walton, S. G.; Hopkins, P. E. Modifying Surface Energy of Graphene via Plasma-Based Chemical Functionalization to Tune Thermal and Electrical Transport at Metal Interfaces. *Nano Lett.* 2015, *15* (8), 4876–4882. https://doi.org/10.1021/acs.nanolett.5b00381.
- (15) Hopkins, P. E.; Baraket, M.; Barnat, E. V.; Beechem, T. E.; Kearney, S. P.; Duda, J. C.; Robinson, J. T.; Walton, S. G. Manipulating Thermal Conductance at Metal-Graphene Contacts via Chemical Functionalization. *Nano Lett.* 2012, 12 (2), 590–595. https://doi.org/10.1021/nl203060j.
- (16) Zheng, W.; Huang, B.; Li, H.; Koh, Y. K. Achieving Huge Thermal Conductance of Metallic Nitride on Graphene Through Enhanced Elastic and Inelastic Phonon Transmission. ACS Appl. Mater. Interfaces 2018, 10 (41), 35487–35494. https://doi.org/10.1021/acsami.8b12480.
- (17) Samani, M. K.; Ding, X. Z.; Khosravian, N.; Amin-Ahmadi, B.; Yi, Y.; Chen, G.; Neyts, E. C.; Bogaerts, A.; Tay, B. K. Thermal Conductivity of Titanium Nitride/Titanium Aluminum Nitride Multilayer Coatings Deposited by Lateral Rotating Cathode Arc. *Thin Solid Films* **2015**, *578* (2015), 133–138. https://doi.org/10.1016/j.tsf.2015.02.032.
- (18) Nagareddy, V. K.; Nikitina, I. P.; Gaskill, D. K.; Tedesco, J. L.; Myers-Ward, R. L.; Eddy, C. R.; Goss, J. P.; Wright, N. G.; Horsfall, A. B. High Temperature Measurements of Metal Contacts on Epitaxial Graphene. *Appl. Phys. Lett.* **2011**, *99* (7), 97–100. https://doi.org/10.1063/1.3627167.
- (19) Khomyakov, P. A.; Giovannetti, G.; Rusu, P. C.; Brocks, G.; Van Den Brink, J.; Kelly, P.

- J. First-Principles Study of the Interaction and Charge Transfer between Graphene and Metals. *Phys. Rev. B Condens. Matter Mater. Phys.* **2009**, *79* (19), 1–12. https://doi.org/10.1103/PhysRevB.79.195425.
- (20) Giovannetti, G.; Khomyakov, P. A.; Brocks, G.; Karpan, V. M.; Van Den Brink, J.; Kelly, P. J. Doping Graphene with Metal Contacts. *Phys. Rev. Lett.* **2008**, *101* (2), 4–7. https://doi.org/10.1103/PhysRevLett.101.026803.
- (21) Quiroga, M. A. O.; Cabeza, G. F. Role of Van Der Waals Forces in Graphene Adsorption over Pd, Pt, and Ni. *Brazilian J. Phys.* **2013**, *43* (3), 126–129. https://doi.org/10.1007/s13538-013-0124-3.
- (22) Hu, L.; Hu, X.; Wu, X.; Du, C.; Dai, Y.; Deng, J. Density Functional Calculation of Transition Metal Adatom Adsorption on Graphene. *Phys. B Condens. Matter* **2010**, *405* (16), 3337–3341. https://doi.org/10.1016/j.physb.2010.05.001.
- (23) Kozlov, S. M.; Viñes, F.; Görling, A. Bonding Mechanisms of Graphene on Metal Surfaces. *J. Phys. Chem. C* **2012**, *116* (13), 7360–7366. https://doi.org/10.1021/jp210667f.
- (24) Mao, R.; Kong, B. D. B.; Gong, C.; Xu, S.; Jayasekera, T.; Cho, K.; Kim, K. W. First-Principles Calculation of Thermal Transport in Metal/Graphene Systems. *Phys. Rev. B* **2013**, 87 (16), 165410. https://doi.org/10.1103/PhysRevB.87.165410.
- (25) Valencia, H.; Gil, A.; Frapper, G. Trends in the Adsorption of 3d Transition Metal Atoms onto Graphene and Nanotube Surfaces: A DFT Study and Molecular Orbital Analysis. *J. Phys. Chem. C* **2010**, *114* (33), 14141–14153.
- (26) Jeong, M.; Freedman, J. P.; Liang, H. J.; Chow, C. M.; Sokalski, V. M.; Bain, J. A.; Malen, J. A. Enhancement of Thermal Conductance at Metal-Dielectric Interfaces Using Subnanometer Metal Adhesion Layers. *Phys. Rev. Appl.* **2016**, *5* (1), 1–7. https://doi.org/10.1103/PhysRevApplied.5.014009.
- (27) Robinson, J. A.; Labella, M.; Zhu, M.; Hollander, M.; Kasarda, R.; Hughes, Z.; Trumbull, K.; Cavalero, R.; Snyder, D. Contacting Graphene. *Appl. Phys. Lett.* **2011**, *98* (5), 96–99. https://doi.org/10.1063/1.3549183.
- (28) Huang, B. C.; Zhang, M.; Wang, Y.; Woo, J. Contact Resistance in Top-Gated Graphene Field-Effect Transistors. *Appl. Phys. Lett.* **2011**, *99* (3), 97–100. https://doi.org/10.1063/1.3614474.
- (29) Song, S. M.; Park, J. K.; Sul, O. J.; Cho, B. J. Determination of Work Function of Graphene under a Metal Electrode and Its Role in Contact Resistance. *Nano Lett.* **2012**, *12* (8), 3887–3892. https://doi.org/10.1021/nl300266p.
- (30) Blake, P.; Yang, R.; Morozov, S. V.; Schedin, F.; Ponomarenko, L. A.; Zhukov, A. A.; Nair, R. R.; Grigorieva, I. V.; Novoselov, K. S.; Geim, A. K. Influence of Metal Contacts and Charge Inhomogeneity on Transport Properties of Graphene near the Neutrality Point. *Solid State Commun.* **2009**, *149* (27–28), 1068–1071. https://doi.org/10.1016/j.ssc.2009.02.039.
- (31) Franklin, A. D.; Han, S. J.; Bol, A. A.; Perebeinos, V. Double Contacts for Improved Performance of Graphene Transistors. *IEEE Electron Device Lett.* **2012**, *33* (1), 17–19. https://doi.org/10.1109/LED.2011.2173154.
- (32) Lin, C. J.; Gorman, G. L. Evaporated CoPt Alloy Films with Strong Perpendicular Magnetic Anisotropy. *Appl. Phys. Lett.* **1992**, *61* (13), 1600–1602. https://doi.org/10.1063/1.107509.
- (33) Santala, T.; Adams, C. M. Kinetics and Thermodynamics in Continuous Electron-Beam Evaporation of Binary Alloys. *J. Vac. Sci. Technol.* **1970**, *7* (6), S22–S29.

- https://doi.org/10.1116/1.1315912.
- (34) Powell, A.; Van Den Avyle, J.; Damkroger, B.; Szekely, J.; Pal, U. Analysis of Multicomponent Evaporation in Electron Beam Melting and Refining of Titanium Alloys. *Metall. Mater. Trans. B Process Metall. Mater. Process. Sci.* **1997**, *28* (6), 1227–1239. https://doi.org/10.1007/s11663-997-0078-3.
- (35) Http://www.crct.polymtl.ca/fact/phase\_diagram.php?file=Ni-Pd.jpg&dir=SGsold. SGTE Solder Alloy Phase Diagrams http://www.crct.polymtl.ca/fact/phase\_diagram.php?file=Ni-Pd.jpg&dir=SGsold.
- (36) Bale, C. W.; Bélisle, E.; Chartrand, P.; Decterov, S. A.; Eriksson, G.; Gheribi, A. E.; Hack, K.; Jung, I.-H.; Kang, Y.-B.; Melançon, J.; Pelton, A. D.; Petersen, S.; Robelin, C.; Sangster, J.; Van Ende, M.-A. FactSage Thermochemical Software and Databases, 2010-2016 www.factsage.com.
- (37) Cahill, D. G. Analysis of Heat Flow in Layered Structures for Time-Domain Thermoreflectance. *Rev. Sci. Instrum.* **2004**, *75* (12), 5119–5122. https://doi.org/10.1063/1.1819431.
- (38) Saha, D.; Yu, X.; Jeong, M.; Darwish, M.; Weldon, J.; Gellman, A. J.; Malen, J. A. Impact of Metal Adhesion Layer Diffusion on Thermal Interface Conductance. *Phys. Rev. B* **2019**, *99* (11), 1–9. https://doi.org/10.1103/PhysRevB.99.115418.
- (39) Nash, A.; Nash, P. The Ni-Pd (Nickel-Palladium) System. *Bull. Alloy Phase Diagrams* **1984**, *5* (5), 446–450.
- (40) Fraenkel, W.; Stern, A. Des System Palladium-Nickel (In German). *Z. Anorg. Chem.* **1927**, *166*, 164–165.
- (41) Lin, W.; Sprueill, J. E. The Structure of Nickel-Palladium Solid Solutions. *Acta Metall.* **1971**, *19*, 451–461.
- (42) Hultgren, R.; Zapffe, C. A. An X-Ray Study of the Iron-Palladium and Nickel-Palladium Systems. *Trans. AIME* **1939**, *133*, 58–68.
- (43) Androulakis, J.; Lin, C.; Kong, H.; Uher, C.; Wu, C.; Hogan, T.; Cook, B. A.; Caillat, T.; Paraskevopoulos, K. M.; Kanatzidis, M. G. Spinodal Decomposition and Nucleation and Growth as a Means to Bulk Nanostructured Thermoelectrics: Enhanced Performance in Pb 1 x Sn x Te PbS. *J. Am. Chem. Soc.* **2007**, *129* (31), 9780–9788. https://doi.org/10.1021/ja071875h.
- (44) Huston, E. L.; Cahn, J. W.; Hilliard, J. E. Spinodal Decomposition during Continuous Cooling. *Acta Metall.* **1966**, *14* (9), 1053–1062. https://doi.org/10.1016/0001-6160(66)90193-3.
- (45) Leong, W. S.; Nai, C. T.; Thong, J. T. L. What Does Annealing Do to Metal-Graphene Contacts? *Nano Lett.* **2014**, *14* (7), 3840–3847. https://doi.org/10.1021/nl500999r.
- (46) Leong, W. S.; Gong, H.; Thong, J. T. L. Low-Contact-Resistance Graphene Devices with Nickel-Etched-Graphene Contacts. *ACS Nano* **2014**, *8* (1), 994–1001. https://doi.org/10.1021/nn405834b.
- (47) Kimling, J.; Philippi-Kobs, A.; Jacobsohn, J.; Oepen, H. P.; Cahill, D. G. Thermal Conductance of Interfaces with Amorphous SiO2 Measured by Time-Resolved Magneto-Optic Kerr-Effect Thermometry. *Phys. Rev. B* **2017**, *95* (18), 1–10. https://doi.org/10.1103/PhysRevB.95.184305.
- (48) Norimatsu, W.; Kusunoki, M. Epitaxial Graphene on SiC{0001}: Advances and Perspectives. *Phys. Chem. Chem. Phys.* **2014**, *16* (8), 3501–3511. https://doi.org/10.1039/c3cp54523g.

- (49) Yu, W. J.; Li, Z.; Zhou, H.; Chen, Y.; Wang, Y.; Huang, Y.; Duan, X. Vertically Stacked Multi-Heterostructures of Layered Materials for Logic Transistors and Complementary Inverters. *Nat. Mater.* **2013**, *12* (3), 246–252. https://doi.org/10.1038/nmat3518.
- (50) Wang, K.; Tai, G.; Wong, K. H.; Lau, S. P.; Guo, W. Ni Induced Few-Layer Graphene Growth at Low Temperature by Pulsed Laser Deposition. *AIP Adv.* **2011**, *1* (2), 1–9. https://doi.org/10.1063/1.3602855.
- (51) Sato, M.; Inukai, M.; Ikenaga, E.; Muro, T.; Ogawa, S.; Takakuwa, Y.; Nakano, H.; Kawabata, A.; Nihei, M.; Yokoyama, N. Fabrication of Graphene Directly on SiO 2 without Transfer Processes by Annealing Sputtered Amorphous Carbon. *Jpn. J. Appl. Phys.* **2012**, *51* (4 PART 2). https://doi.org/10.1143/JJAP.51.04DB01.
- (52) Kato, T.; Hatakeyama, R. Site- and Alignment-Controlled Growth of Graphene Nanoribbons from Nickel Nanobars. *Nat. Nanotechnol.* **2012**, 7 (10), 651–656. https://doi.org/10.1038/nnano.2012.145.
- (53) Anzi, L.; Mansouri, A.; Pedrinazzi, P.; Guerriero, E.; Fiocco, M.; Pesquera, A.; Centeno, A.; Zurutuza, A.; Behnam, A.; Carrion, E. A.; Pop, E.; Sordan, R. Ultra-Low Contact Resistance in Graphene Devices at the Dirac Point. *2D Mater.* **2018**, *5* (2). https://doi.org/10.1088/2053-1583/aaab96.
- (54) Kumar, K.; Kim, Y. S.; Yang, E. H. The Influence of Thermal Annealing to Remove Polymeric Residue on the Electronic Doping and Morphological Characteristics of Graphene. *Carbon N. Y.* **2013**, *65*, 35–45. https://doi.org/10.1016/j.carbon.2013.07.088.
- (55) Choi, W.; Seo, Y. S.; Park, J. Y.; Kim, K. B.; Jung, J.; Lee, N.; Seo, Y.; Hong, S. Effect of Annealing in Ar/H2 Environment on Chemical Vapor Deposition-Grown Graphene Transferred with Poly (Methyl Methacrylate). *IEEE Trans. Nanotechnol.* **2015**, *14* (1), 70–74. https://doi.org/10.1109/TNANO.2014.2365208.
- (56) Lin, Y. C.; Lu, C. C.; Yeh, C. H.; Jin, C.; Suenaga, K.; Chiu, P. W. Graphene Annealing: How Clean Can It Be? *Nano Lett.* **2012**, *12* (1), 414–419. https://doi.org/10.1021/nl203733r.
- (57) Fleutot, B.; Miller, J. B.; Gellman, A. J. Apparatus for Deposition of Composition Spread Alloy Films: The Rotatable Shadow Mask. *J. Vac. Sci. Technol. A* **2012**, *30* (6), 061511. https://doi.org/10.1116/1.4766194.
- (58) Yu, X.; Gellman, A. J. Suppression of B2 Phase in PdzCu1-z Alloy Thin Films. *Thin Solid Films* **2018**, *668* (June), 50–55. https://doi.org/10.1016/j.tsf.2018.10.018.
- (59) Schmidt, A. J.; Cheaito, R.; Chiesa, M. A Frequency-Domain Thermoreflectance Method for the Characterization of Thermal Properties. *Rev. Sci. Instrum.* **2009**, *80* (9). https://doi.org/10.1063/1.3212673.
- (60) Bougher, T. L.; Yates, L.; Lo, C. F.; Johnson, W.; Graham, S.; Cola, B. A. Thermal Boundary Resistance in GaN Films Measured by Time Domain Thermoreflectance with Robust Monte Carlo Uncertainty Estimation. *Nanoscale Microscale Thermophys. Eng.* **2016**, *20* (1), 22–32. https://doi.org/10.1080/15567265.2016.1154630.

