Luminescence of Lanthanide Complexes with Perfluorinated Alkoxide Ligands

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ABSTRACT

Four groups of rare earth complexes, comprising eleven new compounds, with fluorinated O-donor ligands ([K(THF)₆][Ln(OC₄F₉)₄(THF)₂], **1-Ln** (Ln = Ce, Nd), [K][Ln(OC₄F₉)₄], **2-Ln** (Ln = Eu, Gd, Dy), [K(THF)₂][Ln(pin^F)₂(THF)₃], **3-Ln** (Ln = Ce, Nd), and [K(THF)₂][Ln(pin^F)₂(THF)₂], **4-Ln** (Ln = Eu, Gd, Dy, Y) have been synthesized and characterized. Single-crystal X-ray diffraction data were collected for all compounds except **2-Ln**. Species **1-Ln**, **3-Ln**, and **4-Ln** are uncommon examples of six- (Eu, Gd, Dy, Y) and seven- (Ce, Nd) coordinate Ln^{III} centers in all-oxygen-donor environments. Species **1-Ln**, **2-Ln**, **3-Ln**, and **4-Ln** are all luminescent (except where Ln = Gd, Y) with the solid-state emission of **1-Ce** being exceptionally blue-shifted for a cerium complex. The emission spectra of the six Nd, Eu, and Dy complexes do not show large differences based on ligand and are generally consistent with the well-known free-ion spectra. Time-dependent DFT results show that **1-Ce** and **3-Ce** undergo allowed $5f \rightarrow 4d$ excitations, consistent with luminescence lifetime measurements in the nanosecond range. Europium-containing **2-Eu** and **4-Eu**, however, were found to have luminescence lifetimes in the millisecond range, indicating phosphorescence rather than fluorescence. The performance of a pair of multi-reference models for the prediction of the Ln = Nd, Eu, Dy absorption spectra was assessed. It was found that

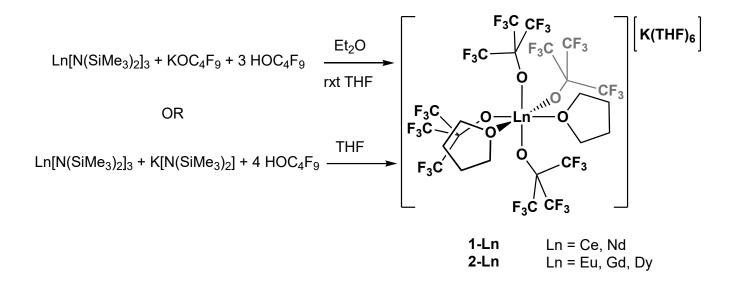
spectroscopy-oriented configuration interaction as applied to a simplified model in which the free ion lanthanide was embedded in ligand-centered Löwdin point charges performed as well (Nd) or better (Eu, Dy) than canonical NEVPT2 calculations where the ligand orbitals were included in the treatment.

KEYWORDS: lanthanides, rare earth, cerium, europium, luminescence, alkoxides, fluorine, oxygen donors

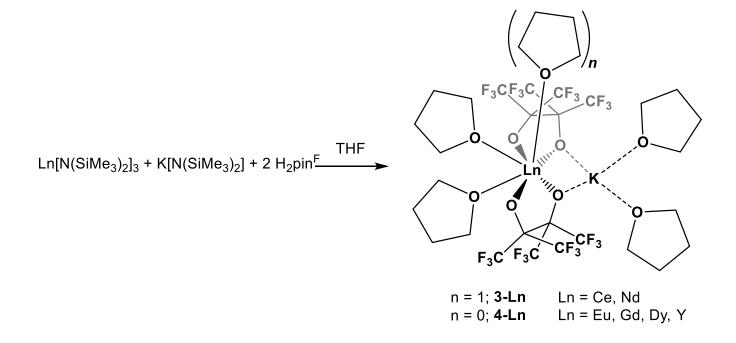
INTRODUCTION

The lanthanide (Ln^{III}) ions are well known for their magnetism and photophysical properties, and are used in lasers, consumer electronics, medical contrast agents, and many other applications.¹⁻⁷ Lanthanide compounds in which structure-property relationships are elucidated are important precursors to these materials. Like other strongly Lewis-acidic rare earth metals, cerium has a strong affinity for hard Lewis bases, such that oxygen donors like β -diketonates⁸⁻¹³ and alkoxides^{5, 14-18} have been extensively investigated. Early Ln^{III} compounds with fluorinated ligands were designed for the deposition of mixed metal oxides, as fluorination improves the volatility and thermal stability of precursors for the manufacture of thin films via chemical vapor deposition (CVD).¹⁹⁻²⁴ In addition, replacing ligand C–H bonds with C–F bonds has been shown to reduce vibrational energy loss that can quench emissions in the NIR range,²⁵ resulting in improved emission quantum yield. To date, however, there are fewer lanthanide complexes with fluorinated ligands than without. Therefore, new Ln-containing complexes with fluorinated ligands have the potential to be useful and informative for their differences from the corresponding proteo analogs.

Our group has previously used fluorinated alkoxides and aryloxides to produce a family²⁶ of complexes including an unusual trivalent Cu²⁷ species, three-coordinate transition metal complexes in exclusively O-donor environments,²⁸ and rare high-spin square-planar metal centers.^{26, 29} To date, five 3*d* metals $(Fe,^{28-30} \text{ Co},^{28, 31} \text{ Ni},^{32} \text{ Cu},^{27-28, 31, 33-35} \text{ Zn}^{28, 36})$ have been prepared as homonuclear, homoleptic $[M(\text{pin}^{F})_2]^{2^-}$ and $[M(\text{OC}_4\text{F9})_n]^{m^-}$ complexes, heteroleptic $[(R_3\text{P})_2\text{M}(\text{OR}^F)_2]$ complexes,³⁷ as well as the main group species $\text{Tl}(\text{OC}_4\text{F9})^{31}$. Thus, with the goal of preparing luminescent Ln^{III} complexes with fluorinated alkoxide ligands and homoleptic O-donor environments, we undertook the synthesis and characterization of a series of Ln^{III} complexes, whose results are shown in Scheme 1, with $[\text{OC}_4\text{F9}]^-$ ligands, and in Scheme 2 for $[\text{pin}^F]^{2^-}$ (perfluoropinacolate) ligands. The compounds may be grouped in four categories based on ligand and coordination number: $([K(\text{THF})_6][\text{Ln}(\text{OC}_4\text{F9})_4(\text{THF})_2]$, **1-Ln** (Ln = Ce, Nd), $[K][\text{Ln}(\text{OC}_4\text{F9})_4]$, **2-Ln** (Ln = Eu, Gd, Dy), $[K(\text{THF})_2][\text{Ln}(\text{pin}^F)_2(\text{THF})_3]$, **3-Ln** (Ln = Ce, Nd), and $[K(\text{THF})_2][\text{Ln}(\text{pin}^F)_2(\text{THF})_2]$, **4-Ln** (Ln = Eu, Gd, Dy, Y).



Scheme 1. Syntheses of the 1-Ln and 2-Ln families with the OC₄F₉ ligand.



Scheme 2. Syntheses of the 3-Ln and 4-Ln families with the pin^F ligand.

EXPERIMENTAL

General procedures.

All syntheses described below were conducted in an inert and anhydrous atmosphere using standard Schlenk line and glovebox techniques at room temperature. The anhydrous solvents THF, Et₂O, and hexanes were dried in an alumina-based solvent purification system (SPS) under Ar, piped directly into a N₂-filled MBraun drybox, and stored over molecular sieves. Potassium hydride (KH) was obtained as a mineral oil dispersion (30 wt %) and purified by washing with hexanes and drying in vacuo prior to storage in a glovebox. The alcohols H₂pin^F and HOC₄F₉ were purchased from Oakwood Chemical, dried over molecular sieves, degassed by three freeze-pump-thaw cycles, and vacuum-transferred prior to use. Toluene and dimethoxyethane (DME) were dried by refluxing over Na/benzophenone under an N₂ atmosphere and distilled. Acetonitrile was distilled from CaH₂ under N₂. Celite (Aldrich) was heated to 125°C under vacuum overnight and stored under N2. KN(SiMe3)2 was purchased from Aldrich, dissolved in toluene, filtered through Celite, and dried under reduced pressure before use. Anhydrous LnCl₃ (Ln = Ce, Nd, Eu, Dy, Y) salts were purchased from Strem and used as received. Anhydrous Ln[N(SiMe₃)₂]₃ (Ln = Ce, Nd, Eu, Dy, Y) were prepared according to the literature.³⁸ UV-vis data were collected with a Shimadzu UV-3600 spectrometer. Elemental analyses were performed by Atlantic Microlabs, Inc. (Norcross, Georgia). Mass Spectrometry data were collected with a Voyager DE STR (MALDI), using a nitrogen laser with DCTB as the matrix.

Synthetic procedures.

[K(THF)₆]**[Ce(OC**₄**F**₉)₄(**THF**)₂], 1-Ce. A solution of HOC₄F₉ (61 µL, 0.44 mmol) in Et₂O (3 mL) was added to a stirred yellow solution of Ce[N(SiMe₃)₂]₃ (91 mg, 0.15 mmol) in Et₂O (3 mL), causing a color change to colorless. A solution of KOC₄F₉ (41 mg, 0.15 mmol) in Et₂O (3 mL) was added to the stirred solution. After 1 hour, the solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing crystals over 7 days (64 mg, 26%). Colorless X-ray quality crystals were grown over 2 days in minimal THF layered with 1:2 THF:hexane at -35 °C. ESI-MS: m/z calcd for [Ce(OC₄F₉)₄]⁻ 1080.23, found 1080. Solution-state maxima: absorption, 298 nm; emission, 406 nm. Solid-state maxima: excitation, 275 nm; emission, 357 nm. Solid-state maxima: excitation, 256 nm; emission, 356 nm.

[K(THF)₆][Nd(OC₄F₉)₄(THF)₂], 1-Nd. HOC₄F₉ (425 μ L, 3.05 mmol) was added dropwise to a stirred blue solution of Nd[N(SiMe₃)₂]₃ (625 mg, 0.999 mmol) in Et₂O (5 mL), causing a color change to light blue. KOC₄F₉ (266 mg, 0.962 mmol) was added and was stirred overnight. The solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, layered with hexane and stored at -35 °C. Dichroic blue/pink X-ray quality crystals were grown over 2 days at -35 °C (500 mg, 31 %). ESI-MS: m/z calcd for [Nd(OC₄F₉)₄]⁻ 1084.36, found 1084. Solid-state maxima: excitation, 350 nm, 525 nm, 580 nm.

[K][**Eu**(**OC**₄**F**₉)4], 2-Eu. A solution of KOC₄F₉ (43 mg, 0.16 mmol) in THF (3 mL) was added to a stirred orange solution of Eu[N(SiMe₃)₂]₃ (98 mg, 0.16 mmol) in THF (3 mL), causing a color change to yellow. HOC₄F₉ (65 μ L, 0.47 mmol) was added in 10 μ L aliquots using a 10 μ L graduated glass syringe causing a color change to colorless and was stirred overnight. The solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing colorless crystals over 4 days (157 mg, 59 %). ESI-MS: m/z calcd for [¹⁵¹Eu(OC₄F₉)₄]⁻ 1090.84, found 1091; calcd for [¹⁵³Eu(OC₄F₉)₄]⁻ 1092.84, found 1093. Solid-state maxima: excitation, 263 nm; emission, 611 nm.

[K][Gd(OC₄F₉)4], 2-Gd. A solution of KN(SiMe₃)₂ (39 mg, 0.20 mmol) in THF (2 mL) was added to a stirred colorless solution of Gd[N(SiMe₃)₂]₃ (125 mg, 0.196 mmol) in THF (4 mL). HOC₄F₉ (110 μ L, 0.788 mmol) was added using a 25 μ L graduated glass syringe. After 10 minutes, the solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for two hours. The solution was layered with room temperature hexane and stored at -35 °C producing colorless crystals over 3 days, which were recrystallized from a chilled THF solution layered with room temperature hexane at -35 °C over 5 days. ESI-MS: m/z calcd for [¹⁵⁸Gd(OC₄F₉)₄]⁻ 1097.85, found 1098; [¹⁵⁵Gd(OC₄F₉)₄]⁻ 1094.85, found 1095.

[K][**Dy**(**OC**₄**F**₉)**4**], **2-Dy.** A solution of KOC₄F₉ (43 mg, 0.16 mmol) in Et₂O (5 mL) was added to a stirred colorless solution of Dy[N(SiMe₃)₂]₃ (101 mg, 0.157 mmol) in Et₂O (5 mL), causing no color change. HOC₄F₉ (66 μ L, 0.47 mmol) in Et₂O (5 mL) was added dropwise and was allowed to stir at room temperature. After 4 days, the solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing colorless crystals over 6 days (105 mg, 39 %). ESI-MS: m/z calcd for [¹⁶⁰Dy(OC₄F₉)₄]⁻ 1099.85, found 1100; [¹⁶¹Dy(OC₄F₉)₄]⁻ 1100.85, found 1101;

 $[^{162}$ Dy(OC₄F₉)₄]⁻ 1101.85, found 1102; $[^{163}$ Dy(OC₄F₉)₄]⁻ 1102.85, found 1103; $[^{164}$ Dy(OC₄F₉)₄]⁻ 1103.85, found 1104. Solid-state maxima: excitation, 298 nm, 324 nm, 350 nm, 364 nm, 378 – 404nm.

[K(THF)₂]**[Ce(pin^F)**₂(**THF)**₃], **3-Ce.** A solution of KN(SiMe₃)₂ (160 mg, 0.802 mmol) and H₂pin^F (287 μ L, 1.61 mmol) in THF (5 mL) was added to a stirred yellow solution of Ce[N(SiMe₃)₂]₃ (500 mg, 0.805 mmol) in THF (5 mL), causing a color change to amber, and was stirred overnight. The solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, layered with hexane and stored at -35 °C. Pale orange X-ray quality crystals were grown over 5 days at -35 °C (459 mg, 47 %). ESI-MS: m/z calcd for {[K][Ce(pin^F)₂]₂]⁻} 1647.51, found 1647. Solution-state maxima: absorption, 248 nm, 322 nm; emission, 452 nm. Solid-state maxima: excitation, 281 nm; emission, 405 nm. Solid-state maxima: excitation, 265 nm; emission, 405 nm.

[K(THF)₂][Nd(pin^F)₂(THF)₃], 3-Nd. A solution of KN(SiMe₃)₂ (160 mg, 0.802 mmol) and H₂pin^F (290 μ L, 1.62 mmol) in THF (5 mL) was added to a stirred blue solution of Nd[N(SiMe₃)₂]₃ (499.7 mg, 0.7995 mmol) in THF (5 mL), causing a color change to yellow, and was stirred overnight. The solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, layered with hexane and stored at -35 °C. Colorless crystals were grown over 4 days at -35 °C. X-ray quality crystals were grown over 2 days in minimal THF layered with hexane at -35 °C (428 mg, 44 %). ESI-MS: m/z calcd for {[K][Nd(pin^F)₂]₂}⁻} 1655.76, found 1655. Solid-state maxima: excitation, 350 nm, 525 nm, 580 nm.

[K(THF)₂][Eu(pin^F)₂(THF)₂], 4-Eu. A solution of KN(SiMe₃)₂ (70 mg, 0.35 mmol) in THF (5 mL) was added to a stirred orange solution of Eu[N(SiMe₃)₂]₃ (216 mg, 0.341 mmol) in THF (10 mL). The solution was stirred for 5 min, during which time it turned deep green. A solution of H₂pin^F (121 µL, 0.680 mmol) in THF (5 mL) was added to the reaction mixture, causing an immediate color change to clear and colorless with a slight yellow tint, and was stirred overnight. The solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing colorless crystals overnight (149 mg, 38 %). X-ray quality crystals were grown by vapor diffusion in concentrated THF against hexane at room temperature over 6 days. ESI-MS: m/z calcd for $\{[K]_3[^{151}Eu(pin^F)_2]\}^+$ 1746.61, found 1747; $\{[K]_3[^{151}Eu(pin^F)_2][^{153}Eu(pin^F)_2]^+\}$ 1748.61, found 1749; $\{[K]_3[^{153}Eu(pin^F)_2]\}^+$ $\{[K][^{151}Eu(pin^{F})_{2}]_{2}[^{153}Eu]^{2+}\}$ 1750.62, found 1751; 1821.61: found 1822; $\{[K]^{153}Eu(pin^{F})_{2}]^{151}Eu(pin^{F})_{2}]^{153}Eu]^{2+}\}$ 1823.61, found 1824; $\{[K]^{153}Eu(pin^{F})_{2}]_{2}^{153}Eu]^{2+}\}$ 1825.61, found 1826. Solid-state maxima: excitation, 464 nm; emission 618 nm.

[K(THF)2][Gd(pin^F)2(THF)2], 4-Gd. A solution of KN(SiMe₃)₂ (80 mg, 0.40 mmol) in THF (3 mL) was added to a stirred colorless solution of Gd[N(SiMe₃)₂]₃ (248 mg, 0.388 mmol) in THF (5 mL). H₂pin^F (140 μ L, 0.784 mmol) was added using a 25 μ L graduated glass syringe. After 60 minutes, the solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing tinted crystals overnight, which were recrystallized in THF layered with hexane and stored at -35 °C over 3 days to produce colorless crystals (226 mg, 51%). X-ray quality crystals were grown by vapor diffusion in a concentrated THF solution against hexane vapor. ESI-MS: m/z calcd for {[K][¹⁶⁰Gd(pin^F)2][¹⁶⁰Gd(pin^F)(O)₂]³⁻} 1386.7, found 1387; {[K][¹⁵⁷Gd(pin^F)2]₂⁻} 1680.7, found 1681; {[K][¹⁵⁷Gd(pin^F)2][¹⁵⁸Gd(pin^F)₂]⁻} 1682.7, found 1683; {[K][¹⁵⁷Gd(pin^F)₂][¹⁶⁰Gd(pin^F)₂]⁻} 1682.7, found 1683; {[K][¹⁵⁷Gd(pin^F)₂][¹⁶⁰Gd(pin^F)₂]⁻} 1684.7, found 1684; {[K][¹⁵⁸Gd(pin^F)₂][¹⁶⁰Gd(pin^F)₂]⁻} 1684.7, found 1685.

[K(THF)₂**][Dy(pin^F)**₂(**THF)**₂**], 4-Dy.** A solution of KN(SiMe₃)₂ (35 mg, 0.18 mmol) in THF (5 mL) was added to a stirred colorless solution of Dy[N(SiMe₃)₂]₃ (112 mg, 0.174 mmol) in THF (5 mL). H₂pin^F (62 μL, 0.35 mmol) was added in 10 μL aliquots using a 10 μL graduated glass syringe, causing a slight yellow tint, and was stirred overnight. The solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing colorless crystals over 3 days (126 mg, 63 %). X-ray quality crystals were grown by vapor diffusion in concentrated THF against hexane at room temperature over 3 days. ESI-MS: m/z calcd for {[K]₃[¹⁶²Dy(pin^F)₂]₂⁺} 1768.63, found 1769; {[K]₃[¹⁶²Dy(pin^F)₂][¹⁶³Dy(pin^F)₂]⁺} 1769.63, found 1770; {[K]₃[¹⁶²Dy(pin^F)₂][¹⁶⁴Dy(pin^F)₂]₂⁺} 1770.63, found 1771; {[K]₃[¹⁶³Dy(pin^F)₂][¹⁶⁴Dy(pin^F)₂]⁺} 1771.63, found 1772; {[K]₃[¹⁶⁴Dy(pin^F)₂]₂⁺} 1772.63, found 1773. Solid-state maxima: excitation, 298 nm, 324 nm, 350 nm, 364 nm, 378 – 404nm.

[K(THF)2][Y(pin^F)2(THF)2], 4-Y. A solution of KN(SiMe₃)₂ (77 mg, 0.39 mmol) in THF (2 mL) was added to a stirred colorless solution of Y[N(SiMe₃)₂]₃ (229 mg, 0.388 mmol) in THF (5 mL). H₂pin^F (140 μ L, 0.784 mmol) was added in 10 μ L aliquots using a 10 μ L graduated glass syringe, creating a slight yellow tint. After 5 minutes, the solvent was removed under vacuum and the resulting solids were dissolved in minimal THF, filtered into a vial, and chilled at -35 °C for an hour. The solution was layered with room temperature hexane and stored at -35 °C producing colorless crystals over 7 days, which were recrystallized from minimal THF layered with hexane at -35 °C over 4 days (223 mg, 56 %). X-ray

quality crystals were grown over 9 days in minimal THF layered with hexane at -35 °C. ESI-MS: m/z calcd for {[K][Y(pin^F)₂]₂⁻} 1544.66, found 1545; {[K₂pin^F]₃[Y(pin^F)₂]⁻} 1982.54, found 1983.

Spectroscopy.

Electronic absorption spectra for 1-Ce, 1-Nd, 3-Ce, and 3-Nd were collected in THF with a Shimadzu UV-3600 spectrometer and quartz cuvettes. Solid-state photoluminescence data for 2-Eu, 2-Dy, 4-Eu, and 4-Dy were collected at 298 K, on a Fluorolog spectrofluorimeter (Horiba Jobin Yvon FL3-22-iHR550) with an excitation monochromator with 1200 grooves/mm gratings blazed at 330 nm and an emission monochromator with 1200 grooves/mm gratings blazed at 500 nm (UV-Vis) or an emission monochromator with 600 grooves/mm gratings blazed at 1000 nm (NIR). An ozone-free xenon lamp of 450 W (Ushio) was used as the radiation source. The excitation spectra, corrected for instrumental function, were measured between 250 and 600 nm. The emission spectra were measured in the range 320-1500 nm in front face mode at 22.5° for the solid-state samples. All emission spectra were corrected for instrumental function. For 1-Ce and 3-Ce, the emission decay curves were obtained using a TCSPC system and a Horiba NanoLED model N-265 (peak wavelength = 265 ± 10 nm, 1 - 2 pJ/pulse) as excitation source. For 2-Eu and 4-Eu, the emission decay curves were obtained using a TCSPC system and a pulsed Xe lamp (pulse width ~5 μ s) as excitation source. Before all decay curve measurements, the spectrum of a blank, using Ludox® solution, was obtained.

Theoretical Calculations

Single-point calculations and gas-phase geometry optimizations of the ground states of all compounds (except 4-Y) were performed at the BP86 level of theory with Gaussian16, Revision A.03.³⁹ The 6-31G* basis set was used for H, C, O, F, and K atoms, and all lanthanides were treated by the Stuttgart RSC ANO/ECP basis set⁴⁰ provided by Basis Set Exchange.⁴¹⁻⁴² 4-Y was treated with the Stuttgart RSC 1997 ECP⁴³ and a double- ζ valence basis set, also from the BSE. Starting coordinates for 1-Ln, 3-Ln, and 4-Ln (except Ln = Y) were obtained from crystal structures. Outer-sphere counter cations were removed from OC₄F₉–containing species, such that calculations were performed on gas-phase [Ln(OC₄F₉)₄(THF)₂]⁻ ions. Optimized coordinates for 1-Nd were modified to be used as starting coordinates for 2-Eu and 2-Dy, which are assumed to be of the same six-coordinates from optimized geometries are available in the supporting information. Time-dependent density functional theory (TD-

DFT) calculations at the B3LYP level of theory were performed on **1-Ce** and **3-Ce** using optimized gasphase structures with Gaussian16. A frozen core of the inner noble gas electrons was used for cerium to reduce computational time. Additional electronic structure analyses were performed using GaussView 6,⁴⁴ ChemCraft 1.8,⁴⁵ and the NBO 6.0⁴⁶⁻⁴⁷ package as provided in the 2016 release of the Amsterdam Density Functional (ADF) program suite.⁴⁸⁻⁴⁹

The absorption spectra of 1-Nd, 2-Eu, 2-Dy, 3-Nd, 4-Eu, and 4-Dy were modeled using the ORCA package, version 4.0.⁵⁰⁻⁵² The recently released SARC2-DKH-QZVP⁵³ basis set was used for all Ln atoms and the relativistically recontracted DKH-DEF2-SV(P)⁵⁴ basis set was used for all other atoms. Scalar relativistic effects were accounted for with the second-order Douglas-Kroll-Hess approximation (DKH2)⁵⁵. The absorption spectra were modeled with quasi-degenerate perturbation theory on top of strongly contracted n-electron valence perturbation theory applied to state-averaged CASSCF wavefunctions (QDPT+NEVPT2/SA-CASSCF),⁵⁶ which has been successfully applied to Ln^{III} systems by Neese et al.^{53, 57} The active space for the SA-CASSCF calculations included 3/6/9 electrons correlated with 7 f-orbitals for the Nd/Eu/Dy species. For the Nd species (1-Nd, 3-Nd), the active space was averaged over the full set of states derived from the free ion terms (35 quartet, 112 doublet [QDPT+SA(35,112)-NEVPT2(3,7)]) for a complete intermediate coupling treatment. For the Eu species (2-Eu, 4-Eu), the active space was averaged over the states derived from the full set of septet and quintet free ion states, and triplet states derived from the ³P, ³O, ³M, ³K, ³H, ³F, and ³I free ion states (7 septet, 140 quintet, 91 triplet [QDPT+SA(7,140,91)-NEVPT2(6,7)]). We have also averaged over the ⁷F, ⁵D, ⁵L, ⁵G, ⁵H, ⁵F, ⁵I, and ³P free ion terms (7 septet, 62 quintet, 3 triplet [QDPT+SA(7,62,3)-NEVPT2(6,7)]) as suggested by Freidzon et al.,⁵⁸ for benchmarking purposes. For the Dy species (**2-Dy**, **4-Dy**), the active space was averaged over the states derived from ⁶H, ⁶F, ⁶P, ⁴F(1), ⁴G(1), ⁴I(1), ⁴M, ⁴K, ⁴L, ⁴P, ⁴D(1), ⁴I(2), ⁴D(2), ⁴H, ⁴G(2), ⁴F(2), ²L, ²K, ²P, ²N, ²F, ²M, ²H, and ²D free ion terms (21 sextet, 133 quartet, 98 doublet [QDPT+SA(21,133,98)-NEVPT2(9,7)]). We have also averaged over the ⁶H, ⁶F, ⁶P, ⁴F, ⁴G, ⁴I, ⁴M, ⁴K, and ⁴L free ion terms as suggested by Freidzon et al.,⁵⁸ (QDPT+SA(21,80)-NEVPT2(9,7)). Spin-orbit coupling was included with quasi-degenerate perturbation theory (QDPT) using the SA-CASSCF wavefunction and strongly contracted NEVPT2⁵⁹⁻⁶¹ diagonal energies. Relativistic picture change effects were included for the spin-orbit mean field operator.

To both estimate the effect of ligand-metal orbital interactions on the energy levels of **1-Nd**, **3-Nd**, **2-Eu**,, **4-Eu**, **2-Dy**, and **4-Dy** and to perhaps improve on the NEVPT2 results described above, trivalent Nd/Eu/Dy ions were embedded in a field of atom-centered Löwdin point charges obtained from PBE0/DEF2-SVP/DKH calculations performed on the tetraanionic (four OC₄F₉ ligands in **1-Nd**, **2-Eu**, **2-Dy**) or trianionic ligand (two pin^F and one K⁺ in **3-Nd**, **4-Eu**, **4-Dy**) combinations from the optimized geometries. The embedded free ions were then modeled with quasi-degenerate perturbation theory on top of spectroscopy-oriented configuration interaction⁶² based on state-averaged CASSCF references (QDPT+SORCI^{emb}/SA-CASSCF^{emb}). Finally, we again note here that the theoretical models listed above were used to predict absorption spectra, not excitation spectra. While we would not expect the differences between the absorption and emission spectra to be significant at higher wavelengths, there might be a greater intensity at shorter wavelengths in the absorption spectra due to non-radiative relaxation processes upon emission.

NBO calculations were based on B3LYP/DKH/SARC2-DKH-QZVP[Ln]/DKH-DEF2-SV(P) [H,C,O,F, K]/cc-pVTZ-DK⁶³(Y) electron densitites.

RESULTS AND DISCUSSION

Synthesis

The $[Ln(OC_4F_9)_4]^-$ compounds, Scheme 1, are prepared from $Ln[N(SiMe_3)_2]_3$ with three equivalents of HOC_4F_9 in the presence of a base, which can be either KN(SiMe_3)_2 or KOC_4F_9, in either Et_2O or THF. X-ray quality single crystals of $[Ln(OC_4F_9)_4]^-$ (1-Ce, 1-Nd) are readily isolated from a concentrated, cold solution of THF layered with hexane or a hexane/THF blend. Crystalline material can be obtained for the analogous compounds with smaller rare earth metals, $[Ln(OC_4F_9)_4]^-$ (2-Eu, 2-Gd, 2-Dy), although robust X-ray quality single crystals could not be obtained. The $[Ln(pin^F)_2]^-$ compounds (3-Ln, 4-Ln), Scheme 2, are all prepared from $Ln[N(SiMe_3)_2]_3$ in the presence of KN(SiMe_3)_2 in THF. These compounds are initially purified by recrystallization from a concentrated, cold solution of THF layered with hexane. X-ray quality single crystals can be obtained either from this method (3-Ce, 3-Nd, 4-Y), or at room temperature by vapor diffusion into concentrated THF of a hexane countersolvent (4-Eu, 4-Gd, 4-Dy).

Solution Behavior

Rare earth complexes of the fluorinated ligands $[OC_4F_9]^-$ and $[pin^F]^{2-}$ are water-sensitive (pKa₁ values of 10.7 and 6.05 respectively). Although transition metal $[M(pin^F)_2]^{2-}$ complexes can be air stable, and even soluble in water, all the complexes reported herein are sensitive to water, perhaps due to the lability of THF and ease of its replacement with H₂O, which could lead to protonolysis and Ln₂O₃ formation. Ce^{III}- containing **1-Ce** and **3-Ce** are also air-sensitive and susceptible to oxidation due to the ease of forming CeO₂.

Structural data for the monodentate $[OC_4F_9]^{1-}$ -containing complexes $[Ln(OC_4F_9)_4(THF)_2]^-$, **1-Ln** (Ln = Ce, Nd) are compared in Table 1, and data collection parameters are given in Table S1. Compounds **1-Ce** and **1-Nd** are monoanionic and are accompanied by a THF-coordinated potassium countercation, $[K(THF)_6]^+$, in the outer sphere (Figure 1). Compound **1-Ce** is six-coordinate at Ce and the four Ce– $O(OC_4F_9)$ bonds have two distinct lengths: those *trans* to the THFs, Ce(1)–O(2), have bond distances of 2.275(4) Å, while those *trans* to each other, Ce(1)–O(1), are slightly longer at 2.296(3) Å (Scheme 3, Table 1). These values are consistent with the literature in which Ce–O distances range from 2.208 Å⁶⁴ to 2.652 Å,¹⁷ with a mean of 2.3(1) Å.^{17, 64-70} The THF ligands are *cis* to each other and identical by symmetry, with a Ce(1)–O(3) bond length of 2.587(4) Å. The same features are seen in **1-Nd**, which maintains the same six-coordinate geometry as shown in Figure S1. The Nd(1)–O(1) bond distances, 2.272(5) Å, are longer than Nd(1)–O(2), 2.242(5) Å, while Nd(1)–O(3) has the longest distance, 2.558(6) Å.

Table 1. Selected distances (Å) and angles (°) in crystal structures of 1-Ce and 1-Nd.

=

	1-Ce	1-Nd
$Ln(1)-O(1)^{a}$	2.296(3)	2.272(5)
$Ln(1)-O(2)^{a}$	2.275(4)	2.242(5)
$Ln(1)-O(3)^{b}$	2.587(4)	2.558(6)
O(1)–Ln(1)–O(2)	93.73(17)	93.9(2)
O(1)–Ln(1)–O(2)	97.96(17)	98.1(2)
O(1)–Ln(1)–O(3)	81.81(14)	81.2(2)
O(1)–Ln(1)–O(3)	82.94(14)	83.3(2)
O(2)–Ln(1)–O(3)	87.60(15)	88.4(2)
O(2)–Ln(1)–O(3)	165.40(15)	165.9(2)

 $^{\it a}O(1)$ and O(2) are oxygen atoms from $[OC_4F_9]^-$

^cO(3) are oxygen atoms from THF

Although no X-ray diffraction quality crystals were obtained for the $[OC_4F_9]^-$ complexes 2-Ln (Ln = Eu, Gd, Dy), ESI-MS data are consistent with the assigned formulae for all. Based upon photoluminescence data (*vide infra*) and the scarcity of non-polymeric four-coordinate complexes of Eu,⁷¹⁻⁸⁹ Gd,^{82, 90-96} and Dy^{82, 89, 97-108} in the literature, it is proposed that 2-Ln also bear two THF ligands as seen in 1-Ln such that the solution structures are six-coordinate. In the absence of crystallographic data for 2-Ln, it is unknown whether the symmetry observed in 1-Ln, which results in only three unique Ln–O distances per

complex, is maintained. The mass spectrometry data also support the structural assignments of 1-Ln, 3-Ln, and 4-Ln that were determined by X-ray crystallography.

Structural data for the bidentate $[pin^{F}]^{2^{-}}$ -containing complexes $[Ln(pin^{F})_{2}(THF)_{3}]^{-}$, **3-Ln**, and $[Ln(pin^{F})_{2}(THF)_{2}]^{-}$, **4-Ln**, are compared in Table 2 and Table S5. An example of the seven coordinate $[K(THF)_{2}][Ln(pin^{F})_{2}(THF)_{3}]$, **3-Ce** is shown in Figure 2 and **3-Nd** is in Figure S2. The complexes are similarly monoanionic but possess an inner-sphere K⁺ countercation. The K⁺ center exhibits long bonds with oxygen atoms of $[pin^{F}]^{2^{-}}$, O(3) and O(4), as well as the fluorine atoms of $[pin^{F}]^{2^{-}}$ and the oxygen atoms of THF (Figure 2). These types of interactions, in which fluorine demonstrates an affinity for oxophilic metal centers, are commonly seen in fluorinated lanthanide complexes.^{19, 109-110} These interactions in the solid state hold the $[pin^{F}]^{2^{-}}$ ligands in a rigid, asymmetric fashion, leaving an open site on the opposite side of the metal center where THF ligands can coordinate.

	3-Ce	3-Nd ^{<i>a</i>}	4-Eu	4-Gd	4-Dy	4- Y
$Ln(1)-O(3)^{b}$	2.404(3)	2.369(3)	2.293(4)	2.257(8)	2.260(2)	2.234(2)
$Ln(1)-O(4)^{b}$	2.399(3)	2.370(3)	2.267(4)	2.284(8)	2.231(2)	2.254(2)
$Ln(1) - O(5)^{c}$	2.337(3)	2.314(3)	2.222(4)	2.226(8)	2.184(2)	2.185(2)
$Ln(1) - O(6)^{c}$	2.349(3)	2.309(3)	2.235(4)	2.209(8)	2.190(3)	2.176(2)
$Ln(1) - O(7)^d$	2.518(4)	2.495(3)	2.413(4)	2.399(8)	2.380(2)	2.352(2)
$Ln(1) - O(8)^d$	2.576(3)	2.516(4)	2.412(4)	2.381(9)	2.385(2)	2.360(2)
$Ln(1)-O(9)^{d}$	2.573(3)	2.536(3)	е	е	е	е
K(1)–O(3)	2.637(3)	2.646(4)	2.653(4)	2.633(8)	2.682(3)	2.646(2)
K(1)–O(4)	2.632(3)	2.629(4)	2.631(4)	2.656(9)	2.625(3)	2.680(2)
O(3)–Ln(1)–O(4)	74.41(11)	74.27(11)	80.18(13)	80.4(3)	80.02(9)	81.06(8)
O(3)–Ln(1)–O(5)	66.45(11)	67.04(12)	69.89(13)	69.8(3)	70.72(9)	71.15(8)
O(4)–Ln(1)–O(6)	66.10(11)	66.92(12)	69.91(13)	69.9(3)	70.34(10)	71.23(7)
Pln1–Pln2 ^f	5.73	4.48	19.23	19.07	17.37	18.84

Table 2. Selected distances (Å) and angles (°) in crystal structures of 3-Ln and 4-Ln

^{*a*}Ln(1) is used to represent atom Nd(2) for **3-Nd**

 ${}^{b}O(3)$ and O(4) are oxygen atoms of $[pin^{F}]^{2-}$ that are proximal to and share an interaction with K⁺

 $^{c}O(5)$ and O(6) are oxygen atoms of $[pin^{F}]^{2-}$ that are distal to K⁺

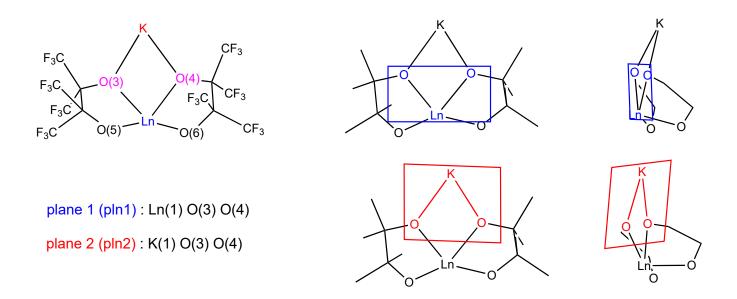
 $^{d}O(7)-O(9)$ are oxygen atoms of THF

^eO(9) of THF is not present for 6-coordinate species **4-Ln**

⁷Pln1 and Pln2 are the Planes of (Ln(1), O(3), O(4)) and (K(1), O(3), O(4)), respectively

respectively

The presence of an inner-sphere K⁺ ion results in a diamond core at the center of the complex marked by atoms Ln(1), O(3), O(4), K(1). For larger rare earth metals (**3-Ln**, Ln = Ce, Nd) this diamond core is close to planar, but for the smaller rare earths (**4-Ln**, Ln = Eu, Gd, Dy, Y) the core is more bent. In the **4-Ln** series, each Ln atom is six-coordinate with one less THF molecule bound than in the **3-Ln** complexes. Figure 3 shows a slightly more open environment at **4-Eu**, and **4-Gd**, **4-Dy**, and **4-Y** are shown in Figures S3, S4, and S5 respectively. Scheme 3 shows a generic **4-Ln** complex and how the two planes, pln1 and pln2 are defined. The angle between the plane of Ln(1), O(3), O(4) (pln1) and the plane of K(1), O(3), O(4) (pln2) generally increases with a decrease in rare earth radial size resulting in a more bent structure for the smaller metals: 5.73° (**3-Ce**); 4.48° (**3-Nd**); 19.23° (**4-Eu**); 19.07° (**4-Gd**); 17.37° (**4-Dy**). Complex **4-Y** is an outlier in this trend: 18.84° (**4-Y**). On the convex side of the angle between pln1 and pln2, there exist long K^{...}F interactions between the inner-sphere K⁺ ion and F atoms of $[pin^{F}]^{2^{-}}$. Interestingly, for the smaller rare earths (**4-Ln**, Ln = Eu, Gd, Dy, Y) where there is a greater pln1–pln2 angle, there are only three K^{...}F interactions, whereas in the complexes of larger rare earth metals (**3-Ln**, Ln = Ce, Nd) there are four, suggesting that the angle of the bent core is independent of K^{...}F interactions.



Scheme 3. Definitions of plane 1 and plane 2, used to calculate angles in Table 2. Sketches have been simplified for clarity with sequential loss of (left) THF ligands, (center) fluorine atoms, and (right) methyl groups.

The two bond distances between the K⁺ and O atoms of $[pin^F]^{2-}$, K(1)–O(3) and K(1)–O(4), are the same for the larger rare earth metals (**3-Ln**, Ln = Ce, Nd). For the smaller rare earths (**4-Ln**, Ln = Eu, Gd, Dy, Y), there is a lack of a consistent trend, where K(1)–O(3) and K(1)–O(4) are the same for the **4-Eu** and **4-Gd**, while an apparent asymmetry exists for **4-Dy** and **4-Y**. An asymmetry in $[pin^F]^{2-}$ coordination resulting from K(1)–O(3) and K(1)–O(4) bonds that is common for all **3-Ln** and **4-Ln** complexes can be seen in how the ligands are bent around the rare earth metal center, toward the K⁺ ion. Bond angles O(3)– Ln(1)–O(4) are significantly shorter than the O(5)–Ln(1)–O(6) angles, and generally increase as radial size decreases. For each complex, angles between the rare earth metal and the O atoms of a given $[pin^F]^{2-}$ ligand, O(3)–Ln(1)–O(4) and O(3)–Ln(1)–O(5), are the same as one another.

The bond lengths between the rare earth centers and the oxygen atoms of $[pin^F]^{2-}$ that share an interaction with K⁺, Ln(1)–O(3) and Ln(1)–O(4), are generally the same length as one another and shorter than the Ln(1)–O(5) and Ln(1)–O(6) distances to the two THF molecules. Exceptions include in **4-Eu** where Eu(1)–O(3), 2.293(4) Å, is longer than Eu(1)–O(4), 2.267(4) Å, and in **4-Gd** where Gd(1)–O(3), 2.257(8) Å, and Gd(1)–O(5), 2.226(8) Å, are the same within error. For the smaller rare earth metals that are six-coordinate, there is no difference in bond length between the metal centers and the two coordinated THF molecules, Ln(1)–O(7) and Ln(1)–O(8). However, for the larger rare earths that are seven-coordinate, two of the THF molecules are at a further distance than the third. In **3-Ce**, Ce(1)–O(8) and Ce(1)–O(9) are similar with bond distances of 2.576(3) Å and 2.573(3) Å, respectively, while the third, Ce(1)–O(7), is demonstrably shorter at 2.518(4) Å. The same pattern is seen in **3-Nd** where Nd(1)–O(8) and Nd–O(9) (2.516(4) Å and 2.536(3) Å) are longer than Nd(1)–O(7) (2.495(3) Å). The longer Ln(1)–O(7) bond may result from repulsion with the closest CF₃ groups of [pin^F]²⁻, which is less present on the other side of the molecule where K···F interactions are found.

Within $[pin^{F}]^{2-}$ the O–C bond distances are similar and unexceptional, as is evident in **3-Ce**: O(3)–C16, 1.364(5) Å; O(4)–C10, 1.360(5) Å; O(5)–C19, 1.347(6) Å; O(6)–C13, 1.350(5) Å. The C10–C13 and C16–C19 distances of the $[pin^{F}]^{2-}$ backbone are similar, 1.658(7) Å and 1.668(7) Å, respectively, but are significantly longer than that of a typical Csp^{3} – Csp^{3} bond of 1.54 Å.¹¹¹ Such elongated bonds are characteristic of the pin^F dianion when chelating a metal center.²⁶ The C–F distances within $[pin^{F}]^{2-}$ vary from 1.326(7) Å (C18–F19) to 1.355(6) Å (C17–C10), but are all the same within error, indicating that the interactions between K⁺ and the F atom of $[pin^{F}]^{2-}$ do not result in a change in C–F bond lengths.

Concomitant with the decrease in coordination number is a change in Ln–O(pin^F) bond lengths.¹¹²⁻¹¹³ Subtracting the ionic radius¹¹⁴ of each Ln^{III} ion from its average Ln–O(pin^F) length normalizes the distances from the metal center to the ligand in a method inverse to that of using a constant ionic radius of O^{2^-} (oxide) to study the ionic radii of the lanthanides.¹¹⁵⁻¹¹⁶ The results of this analysis are presented in Table 3 and show that there is a distinct difference between the six- (**4-Ln**) and seven- (**3-Ln**) coordinate complexes, beyond the change in coordination number: the average Ln–O(pin^F) distances are longer for the smaller ions, despite the steric relief provided by removal of one THF ligand. The difference in Ln–O(pin^F) distances **4-Ln** and **3-Ln** is reduced when the Ln–O(THF) distances are included in the averages (Table S3), indicating that it is primarily the $[pin^F]^{2^-}$ ligands that experience increased strain. This difference is consistent with the rigidity of the molecular structure, as the bidentate nature of $[pin^F]^{2^-}$, coupled with the steric bulk of four CF₃ groups and the interaction of K⁺, restricts the ligands flexibility in the tight coordination environment.

Table 3. Normalization of Ln(1)–O(O(3)-O(6)) distances (Å) in $[pin^F]^{2-}$ complexes.

	3-Ce	3-Nd	4-Eu	4-Gd	4-Dy	4- Y
Coordination Number	7	7	6	6	6	6
Ln ^{III} Ionic Radius ¹¹⁷	1.07	а	0.947	0.938	0.912	0.9
Avg. Ln–O Distance ^b	2.372	2.341	2.255	2.244	2.216	2.212
Normalized Ln–O Distance	1.302		1.308	1.306	1.304	1.312

^anot available for 7-coordinate Nd³⁺

^{*b*}average bond distance of Ln(1)–O(O(3)-O(6))

There are few other crystallographically characterized examples of mononuclear Ce^{III} complexes in which all seven coordinating atoms are oxygen donors. Only five entries in the CSD¹² ([Li₃(THF)₄]](BINOLate)₃Ce(THF)]¹¹⁸, [Li(THF)₂][Li(THF)][Li(OPPh₃)][(BINOLate)₃Ce(OPPh₃)]⁶⁶, [Li(DME)₃][(BINOLate)₃Ce(OPPh₃)]⁶⁶, Ce(Ph₂P(O)NP(O)Ph₂)₃(THF)¹¹⁹, [Ce(OTf)₂(OPPh₃)₄][OTf]¹²⁰) meet these criteria. Ce–O distances range from 2.337 Å to 2.650 Å,¹²⁰ with a mean of 2.44 Å (Table S2).^{65-66, 118-120} Thus, **1-Ce** and **3-Ce** are new contributions to a very small family of complexes. Mononuclear six- and seven-coordinate Nd complexes in all-oxygen donor environments are also rare, with only six^{67,} ^{114, 121-124} and eight^{25, 125-131} entries in the CSD, respectively.

All these fluorinated alkoxide compounds are new members of a rather small family of compounds. There are approximately 2800 structures of non-polymeric lanthanide complexes containing fluorine anywhere in their crystal structures,¹² of which about 500 have only triflate as a fluorine source. The majority feature fluorinated β -diketonates, such as hexafluoroacetylacetonate (hfac), as ligands ancillary to non-fluorinated ligands under principal study. There are only 65 unique species that have neither a β -

diketonate nor a triflate ligand (details of this survey are shown in Scheme S1). One other crystallographically characterized monometallic Ln species¹³² contains two pin^F-derived ligands, but unlike **3-Ce**, it has two phenanthroline co-ligands, namely $[Ce^{III}(pin^F)(Hpin^F)(phen)_2]$ (Scheme S2). Its luminescence behavior is also markedly different (*vide infra*), with an absorption wavelength of 310 nm and primary emission at 380 nm. The rigid phenanthroline moieties are partly responsible for the smaller Stokes shift, as very little vibrational energy is lost through C–H bonds.

Having established the structures of these compounds, it was naturally of interest to compare the absorption and emission spectra of **1-Ln** and **2-Ln** to those of **3-Ln** and **4-Ln**. Calculated spectra were then compared to those obtained experimentally to understand the electronic structure of these complexes and to determine the level of contemporary theory appropriate to treat such complexes.

Spectroscopy.

Normalized solution-state absorption and emission spectra of Ce complexes **1-Ce** and **3-Ce** are shown in Figure 4. Both absorb in the UV range and emit in the blue, with the $[OC_4F_9]^{1-}$ -ligated species **1-Ce** displaying higher energy transitions and a smaller Stokes shift. Although the molecular structure of **1-Ce** is less rigid than **3-Ce**, it is hypothesized that the presence of fewer THF ligands results in fewer C–H bonds through which to lose excitation energy via vibrations.

The solid-state excitation and emission spectra of the Ce^{III} complexes are shown in Figure 5. The broad emission bands, attributed to $d \rightarrow f$ transitions,¹³³⁻¹³⁴ are located in the blue region with maxima at ~357 and ~405 nm for the complexes **1-Ce** and **3-Ce**, respectively. The different coordination environments around the Ce^{III} are reflected in the shift, of ~50 nm, between the maxima of the emission bands. Short luminescence lifetimes in the nanosecond range (Table 4, Figures S6 and S7) for the Ce complexes are characteristic of Laporte-allowed *d-f* transitions.¹³⁵

Table 4: Luminescence lifetimes of Ce and Eu complexes.

	Ce	Eu
Ligand		
[pin ^F] ^{2–}	$30.6 \pm 0.1 \text{ ns}$	$0.935 \pm 0.001 \text{ ms}$
$[OC_4F_9]^{2-}$	$5.0\pm0.1\ ns$	$1.230\pm0.045\ ms$

Compared to the literature (Table S4), the solid-state emission of **1-Ce** is the third most blue-shifted of 49 monometallic Ce species that were reported along with their luminescence properties. Only cerium triflate in THF¹³⁵ and hexachlorocerate in acetonitrile¹³⁶ show emissions at lower wavelengths. A full comparison is available in Table S4. In the solution phase, however, the luminescence of **1-Ce** and **3-Ce** are within a more commonly observed range.

The solid-state excitation and emission spectra of the Nd^{III} complexes **1-Nd** and **3-Nd** are shown in Figure 6. The excitation spectra of the complexes are composed of characteristic Nd^{III} *f-f* intraconfigurational transitions. All the expected ${}^{4}F_{3/2} \rightarrow {}^{4}I_{J}$ (J = 9/2, 11/2 and 13/2) transitions are observed, and the ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ transition is the most intense one. Compared to the spectra of **1-Ce** and **3-Ce** above, those of **1-Nd** and **3-Nd** do not change appreciably with the substitution of the ligands, which is consistent with expected lanthanide emission behavior and further highlights the difference seen in the cerium complexes.

The solid-state excitation and emission spectra of the Eu^{III} complexes 2-Eu and 4-Eu are shown in Figure 7, and the emission decay curves are in Figures S8 and S9, respectively. As for all the lanthanides studied here, except Ce, the excitation spectra are dominated by the f-f transitions, indicating that the emission occurs through direct metal-centered excitation, instead of ligand-mediated sensitization. The Eu-containing complexes have a longer lifetime than Ce, in the millisecond range, characteristic of forbidden nature of the *f*-*f* transitions (Table 4). All the expected characteristic narrow ${}^{5}D_{0} \rightarrow {}^{7}F_{J}$ (J = 0-4) transitions are observed, and the ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ transition is the most intense one. In the emission spectra the differences in fine splitting seen for each ${}^{5}D_{0} \rightarrow {}^{7}F_{J}$ (J = 0 - 4) transition of each complex highlight the change in symmetry around the Eu^{III} ions in both complexes. For example, the ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$ transition is present and fairly intense for 4-Eu, while mostly absent for 2-Eu. The symmetry around Eu in 4-Eu $[K(THF)_2][Eu(pin^F)_2(THF)_3]$ (~C_{2v}) is lower than that in the presumed structure of **2-Eu**, [K(THF)₆][Eu(OC₄F₉)₂(THF)₂] (~O_h),^{137, 140} consistent with the difference in intensity for 0-0 peak. The ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$ transition of the two compounds also differ, with more fine structure for 4-Eu. This behavior is further mirrored in the ${}^5D_0 \rightarrow {}^7F_2$ transition, which has a completely different profile for both complexes, with more intense fine structure components for the complex with lower symmetry, as expected. The differences in symmetry are reflected as well in the calculated excitation spectra (Figures 13 and 14). While the experimental spectra are not sharp enough at room temperature to enable extensive symmetry discussions, the calculated spectra are distinctively different as discussed below.

The excitation and emission spectra of the Dy^{III} complexes **2-Dy** and **4-Dy** are shown in Figure 8. The excitation spectra are composed of only characteristic Dy^{III} *f*-*f* intraconfigurational transitions (${}^{4}H_{13/2} \leftarrow {}^{6}H_{15/2} \sim 298 \text{ nm}$, (${}^{6}P_{3/2} + {}^{4}K_{15/2}$) $\leftarrow {}^{6}H_{15/2} \sim 324 \text{ nm}$, (${}^{4}M_{15/2} + {}^{6}P_{7/2}$) $\leftarrow {}^{6}H_{15/2} \sim 350 \text{ nm}$, (${}^{6}P_{5/2, 3/2}$) $\leftarrow {}^{6}H_{15/2}$

~364 nm, $({}^{4}M_{21/2, 19/2} + {}^{4}K_{17/2} + {}^{4}F_{7/2} + {}^{4}I_{13/2}) \leftarrow {}^{6}H_{15/2} \sim 378 - 404$ nm, ${}^{4}G_{11/2} \leftarrow {}^{6}H_{15/2} \sim 423$ nm).¹³⁸ The ${}^{4}F_{9/2} \rightarrow {}^{6}H_{13/2}$ transition is the most intense.

All the d^0 Y- and f^7 Gd-containing complexes (4-Y, 2-Gd, and 4-Gd, respectively) are expected to be non-emissive under illumination and were not studied in their pure form. Lifetime measurements of 2-Dy and 4-Dy were complicated due to apparent concentration quenching of the excited state.¹³⁹⁻¹⁴¹ The synthesized Gd^{III} complexes were mixed with their corresponding Dy^{III} analogs, in an effort to dilute the Dy^{III} centers in the solid state and minimize the quenching, but the attempts were not successful. Five mixed Gd/Dy samples each for 2-Gd/Dy and 4-Gd/Dy (% Gd = 90, 75, 50, 25, 10; balance Dy) were prepared and attempts were made to record their luminescence lifetimes, but apparent quenching persisted. Attempts were also made with 0.1% and 0.05% Dy in predominantly Gd samples, but no differences were observed versus the pure Dy samples.

Due to the absence of sensitized emission, further in depth photophysical characterization of the complexes was not pursued. We focussed instead on modelling the spectroscopic behaviour in the solid state.

Computational.

Geometry-optimized gas-phase structures of compounds **1-Ce** and **3-Ce** show Ce–O distances all within 0.1 Å of the corresponding values in the crystal structures (Table S5). The symmetry observed in **1-Ce** is preserved in the gas phase, with three distinct Ce–O distances rather than six. The same level of accuracy was obtained for **1-Nd** and **3-Nd**, but in the case of **1-Nd**, the symmetry was lost upon optimization. The average Ln–O(pin^F) distances for **3-Nd**, **4-Eu**, **4-Dy**, and **4-Y** were all within 0.04 Å of the calculated distances. Gas-phase geometries were not able to be obtained for **2-Gd** or **4-Gd**.

Ln–O bonds in the gas-phase optimized structures of **1-Ln**, **2-Ln**, **3-Ln**, **4-Ln**, and **4-Y** were evaluated using both canonical orbital and NBO analyses and two measures of bond order: Mayer (MBO) and Wiberg (WBO). The WBOs are small for Ln–THF bonds, below 0.14 for all complexes studied, and MBOs are similarly reduced relative to Ln–O(pin^F) bond orders. The Ln–O(C4F9) bonds have higher bond orders, consistent with the difference in bond lengths seen in Table 2. There is asymmetry in the $[pin^F]^{2-1}$ ligand environments of **3-Ce**, **3-Nd**, **4-Eu**, **4-Dy**, and **4-Y**, with the bond order of Ln(1)–O(3) and Ln(1)–O(4) being slightly but consistently lower than the others (Table 5). This difference in bond order is

accompanied by longer bond lengths. The $K(1)\cdots O(3)$ and $K(1)\cdots O(4)$ bonds are ionic and very little sharing of electron density is observed in the calculations.

It can also be seen in Table 5 that the Y–O(pin^F) bonds are similar to the other Ln–O(pin^F) bonds. It is commonly observed that Y^{III} is able have stronger orbital bonds with hard ligands, but in this case the WBO suggests that the $[pin^F]^{2-}$ ligands interact predominately through electrostatic means. This idea is supported by the natural electron configuration of the yttrium center that contains only 0.82 non-core d-electrons. This result is similar to the occupation of the d-orbitals in the **3-Ln** and **4-Ln** complexes, which range between 0.66 for **3-Ce** to 0.80 for **4-Dy**.

	Bond	Distance (Å)	MBO	WBO
3-Ce	Ce(1)–O(3)	2.373	0.50	0.21
	Ce(1)–O(4)	2.372	0.50	0.21
	Ce(1)–O(5)	2.356	0.59	0.23
	Ce(1)–O(6)	2.357	0.59	0.23
3-Nd ^{<i>a</i>}	Nd(1)–O(3)	2.371	0.47	0.20
	Nd(1)–O(4)	2.366	0.49	0.20
	Nd(1)–O(5)	2.321	0.58	0.22
	Nd(1)–O(6)	2.309	0.63	0.23
4-Eu	Eu(1)–O(3)	2.327	0.55	0.24
	Eu(1)–O(4)	2.329	0.54	0.23
	Eu(1)–O(5)	2.256	0.68	0.28
	Eu(1)–O(6)	2.255	0.67	0.27
4-Dy	Dy(1)–O(3)	2.241	0.54	0.23
	Dy(1)–O(4)	2.246	0.54	0.23
	Dy(1)–O(5)	2.176	0.70	0.28
	Dy(1)–O(6)	2.182	0.69	0.28
4- Y	Y(1)–O(3)	2.262	0.68	0.24
	Y(1)–O(4)	2.257	0.69	0.25
	Y(1)–O(5)	2.190	0.82	0.28
	Y(1)–O(6)	2.190	0.82	0.28

Table 5: Comparisons of Ln–O distances and bond orders in gas-phase optimized structures of $[pin^F]^{2-}$ species **3-Ln** and **4-Ln**.

^{*a*}Nd(1) is used to represent atom Nd(2) for **3-Nd**

An overlay of the solution-state, solid-state, and simulated gas-phase absorption spectra for Ce-containing **1-Ce** is presented in Figure 9. Whereas vibronic coupling is observed in the solution-phase data, the same features are not found in either the solid state or calculated spectra. In the gas phase, three main transitions comprise the single broad peak, and the MOs contributing the most to each excited state are visualized in Figure 10. Each of the excited states is best characterized as an excitation from the Ce 4*f* SOMO to a 5*d* orbital (Table 6). Good agreement with experiment was found for the calculated absorption spectrum of **3-Ce** (Figure 11) and the calculated absorptions transitions with the main contributing MOs for **3-Ce** (Figure 12).

	Excited State	Acceptor MO	CI Coefficient	% Ce d in MO	% Cefin MO
1-Ce	7	290 A	0.92747	72.99	-
	8	291 A	0.96241	70.07	-
	9	292 A	0.94035	73.52	-
3-Ce	7	292 A	0.43868	40.88	-
		296 A	0.40906	56.63	12.11
		297 A	0.51404	27.97	-
	8	298 A	0.90483	63.75	-

Table 6. Atomic orbital contributions to excited states in **1-Ce** and **3-Ce**, determined from the compositions and CI expansion coefficients of the dominant positive transitions in each case.

The absorption spectra of the Nd^{III} complexes **1-Nd** and **3-Nd** were modeled using QDPT+NEVPT2/SA-CASSCF and QDPT-SORCI^{emb}/SA-CASSCF^{emb} calculations (herein referred to as NEVPT2 and SORCI^{emb}). The results are shown in Figures S10, S11, S12 and S13, with the energy levels and peak assignments shown in Table S6-S7. It can be seen that both the NEVPT2 and SORCI^{emb} calculations adequately reproduce the relative positions of the spectral features for both **1-Nd** and **3-Nd**, while the absolute positions of the excitations are blue-shifted by approximately 40 nm (2778 cm⁻¹).

As was hinted at above, there is very little difference between the excitation spectra of **1-Nd** and **3-Nd**. Nonetheless, we did attempt to validate our theoretical models by assessing their ability to reproduce slight discrepancies in peak positions/intensities/splitting. If the theoretical models are unable to reproduce these discrepancies, then their utility in validating structural assignments based on spectroscopic data is less powerful. We first note that the fine structure of the ${}^{4}G_{5/2} + {}^{4}G_{7/2} \leftarrow {}^{4}I_{9/2}$ peaks at ~590 nm is completely lost when using the QDPT-SORCI^{emb}/SA-CASSCF^{emb} model. This result immediately suggests that this fine structure is a result of orbital interactions between the metal center and the ligands. As will be shown below, the other Eu and Dy species that were modeled did not display fine structure that was not captured by the SORCI^{emb} model. This difference may be due the greater number of empty *f*-orbitals in the Nd^{III} species that are able to act as acceptors for the ligands.

There are two other main discrepancies between the excitation spectra of **1-Nd** and **3-Nd**. The first is the intensity of the ${}^{4}G_{11/2}+{}^{2}K_{15/2}+{}^{2}D_{3/2} \leftarrow {}^{4}I_{9/2}$ peaks at centered at ~470 nm. It can be seen that the peak at the high energy end of this group of peaks is more intense for **3-Nd** and both the NEVPT2 and SORCI^{emb} calculations reflect this feature. The second discrepancy is the splitting of the group of the peak(s) at ~355

nm. This peak is split in the spectrum for **3-Nd** but appears as a single peak for **1-Nd**. The NEVPT2 calculation is able to reproduce this feature, but not the SORCI^{emb} model.

Moving to the Eu^{III}-containing species, the SORCI^{emb} calculations greatly outperform the NEVPT2 calculations as can be seen in Figures S14, S15, S16 and S17. In particular, the NEVPT2 model overstabilizes the ${}^{4}L_{J}$ peaks relative to the other peaks present in the spectra. That being said, the SORCI^{emb} model was not entirely successful, and the high density of both relativistic and non-relativistic states made robust peak assignments unattainable. It is believed that the use of full intermediate couplings would help to overcome some of the shortcomings observed in the models we employed, but currently this calculation is not feasible as this would involve 7 septet roots, 140 quintet roots, 588 triplet roots, and 490 singlet roots.

There are two main visual differences between the excitation spectra of **2-Eu** and **4-Eu**. The first is the peak at ~450 nm in the spectrum of **2-Eu** (Figure 13) that is not present in the excitation spectrum of **4-Eu** (Figure 14) and was not reproduced by either the NEVPT2 or SORCI^{emb} models. The second discrepancy is the intensity of the hypersensitive ${}^{5}D_{2} \leftarrow {}^{7}F_{0}$ peak at ~465 nm. The SORCI^{emb} model was able to reproduce the enhanced intensity of this peak in the excitation spectrum of **4-Eu**.

Both the SORCI^{emb} and NEVPT2 models were more successful in the prediction of the absorption spectra for the Dy^{III} complexes **2-Dy** (Figure S18) and **4-Dy** (Figure S19) as compared with their performance for **2-Eu** and **4-Eu**, with the SORCI^{emb} model once again outperforming the NEVPT2 calculations (Figures S20 and S21). The SORCI^{emb} model was able to reproduce the enhanced intensity of the ${}^{4}G_{11/2} \leftarrow {}^{6}H_{15/2}$ (~425 nm) relative to the ${}^{4}F_{9/2} + {}^{4}I_{15/2} \leftarrow {}^{6}H_{15/2}$ (~450 nm) observed in **4-Eu**. In addition, the SORCI^{emb} and NEVPT2 models both reduced the intensity of the peak at ~325 nm relative to the peaks centered at ~387 nm when moving from **2-Dy** and **4-Dy**.

SUMMARY

The eleven new species reported herein, grouped by ligand and coordination number, $([K(THF)_6][Ln(OC_4F_9)_4(THF)_2]$, **1-Ln** (Ln = Ce, Nd), $[K][Ln(OC_4F_9)_4]$, **2-Ln** (Ln = Eu, Gd, Dy), $[K(THF)_2][Ln(pin^F)_2(THF)_3]$, **3-Ln** (Ln = Ce, Nd), and $[K(THF)_2][Ln(pin^F)_2(THF)_2]$ **4-Ln** (Ln = Eu, Gd, Dy, Y), and are members of the quite small family of mononuclear lanthanide complexes bearing fluorinated alkoxide donors. The fluorescence properties of the cerium complexes **1-Ce** and **3-Ce** depend upon the *d*-orbital splitting effected by these ligands, and further tuning of those ligands (different substituents on pinacolate or the tertiary alkoxide, fluorinated ligands with greater denticity) may potentially lead to highly blue-shifted luminescence. Lifetime data indicated that **2-Eu** and **4-Eu** emissions are phosphorescent. The emission spectra of the six Nd, Eu, and Dy complexes do not show large differences based on ligand and are generally consistent with the well-known free-ion spectra. Time-dependent DFT results show that **1-Ce** and **3-Ce** undergo allowed $5f \rightarrow 4d$ excitations, consistent with luminescence lifetime measurements in the nanosecond range. The ability of a pair of multi-reference models to predict the excitation spectra was evaluated. It was found that for the Nd species, the NEVPT2 model that includes the ligand orbitals was able to reproduce some fine structure that was absent in the SORCI^{emb} model. For the Eu and Dy species, however, the SORCI^{emb} model outperformed the NEVPT2 model, suggesting that the lack of orbital interactions in this model is more than made up for by the more robust treatment of correlation provided by SORCI.

SUPPORTING INFORMATION

The supporting information is available free of charge at ..

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Notes

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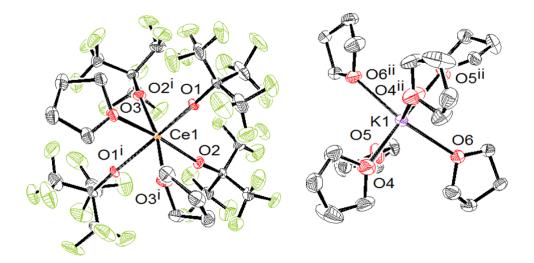


Figure 1. ORTEP of $[K(THF)_6][Ce(OC_4F_9)_4(THF)_2]$ (1-Ce). Ellipsoids are drawn at the 50% level. Hydrogen atoms have been removed for clarity.

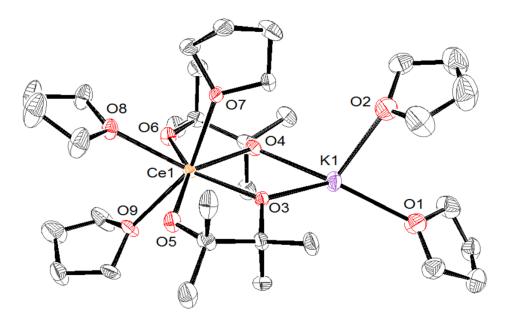


Figure 2. ORTEP of $[K(THF)_2][Ce(pin^F)_2(THF)_3]$ (**3-Ce**). Ellipsoids are drawn at the 50% level. Hydrogen and fluorine atoms have been removed for clarity.

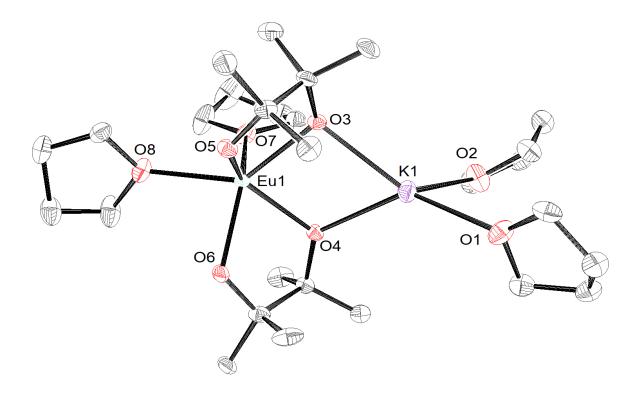


Figure 3. ORTEP of $[K(THF)_2][Eu(pin^F)_2(THF)_2]$ (**4-Eu**). Ellipsoids are drawn at the 50% level. Hydrogen and fluorine atoms have been removed for clarity.

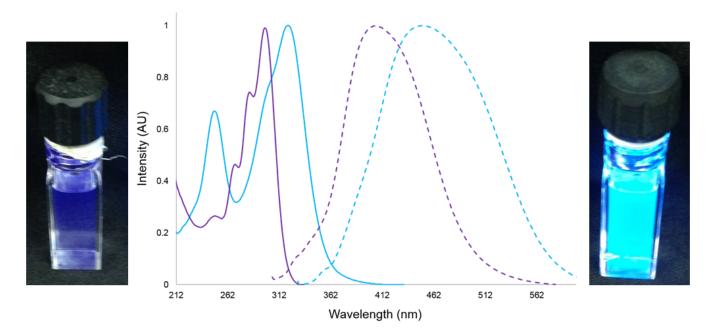


Figure 4. Normalized solution-state excitation and emission spectra of **1-Ce** (purple) and **3-Ce** (blue). Absorption maxima: 298 nm (**1-Ce**); 248, 322 nm (**3-Ce**). Emission maxima: 406 nm (**1-Ce**); 452 nm (**3-Ce**). Photographs of ~0.15 mM **1-Ce** (left) and **3-Ce** (right) in THF were taken under irradiation by 254 nm light.

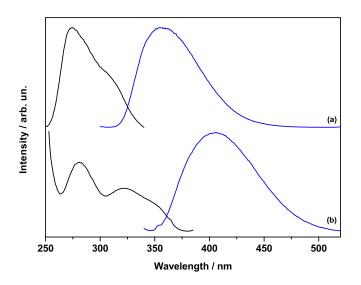


Figure 5. Solid-state excitation and emission spectra of **1-Ce** (top) and **3-Ce** (bottom). Excitation maxima (black traces): 275 nm (**1-Ce**); 281 nm (**3-Ce**). Emission maxima (blue traces): 357 nm (**1-Ce**); 405 nm (**3-Ce**).

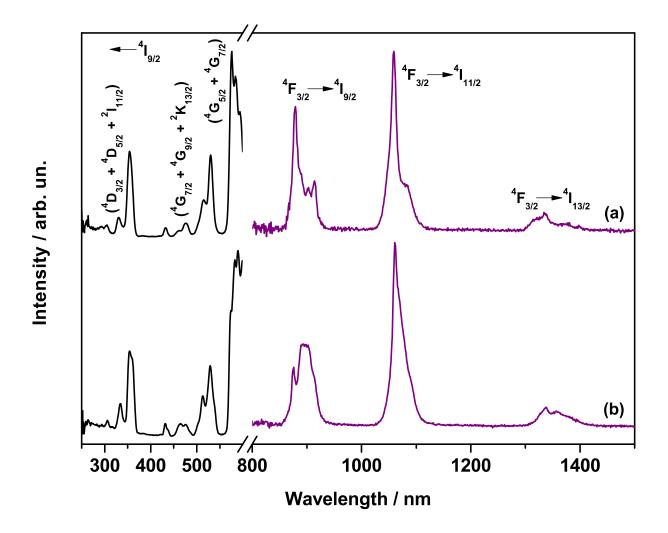


Figure 6. Solid-state excitation and emission spectra of 1-Nd (a) and 3-Nd (b) ($\lambda_{exc} = 350$ nm).

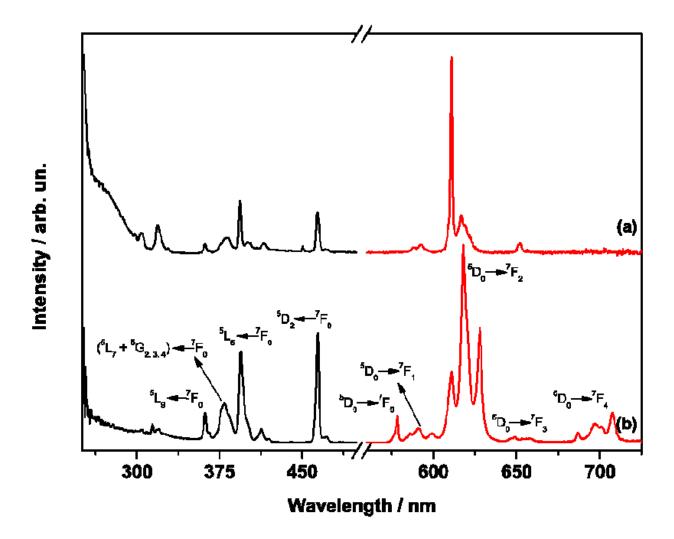


Figure 7. Solid-state excitation and emission spectra of 2-Eu (a) and 4-Eu (b) ($\lambda_{exc} = 464 \text{ nm}$).

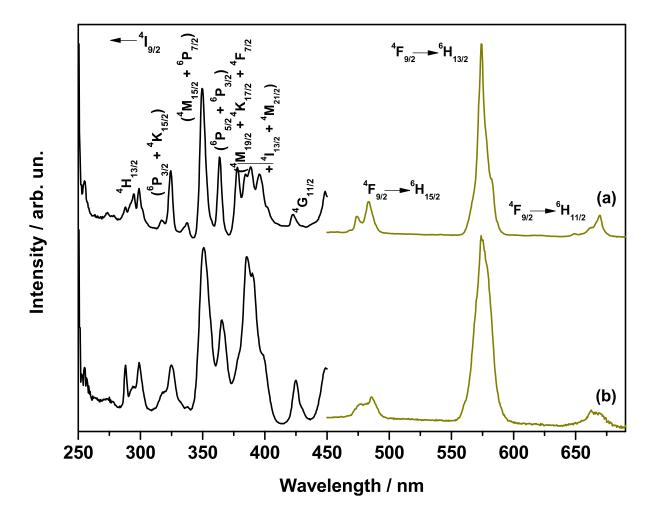


Figure 8. Solid-state excitation and emission spectra of 2-Dy (a) and 4-Dy (b) ($\lambda_{exc} = 350$ nm).

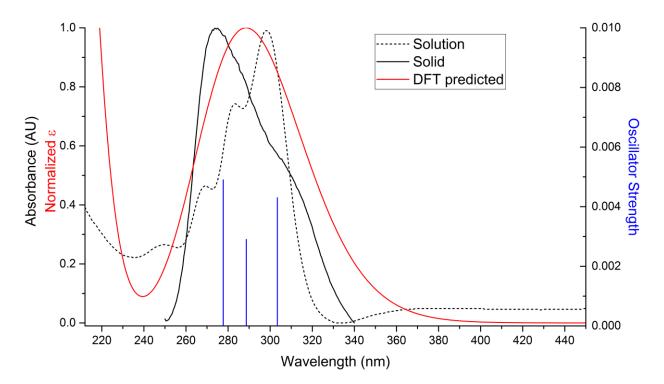


Figure 9. Overlay of solution, solid, and simulated absorption spectra for **1-Ce**. Calculated absorption at 289 nm.

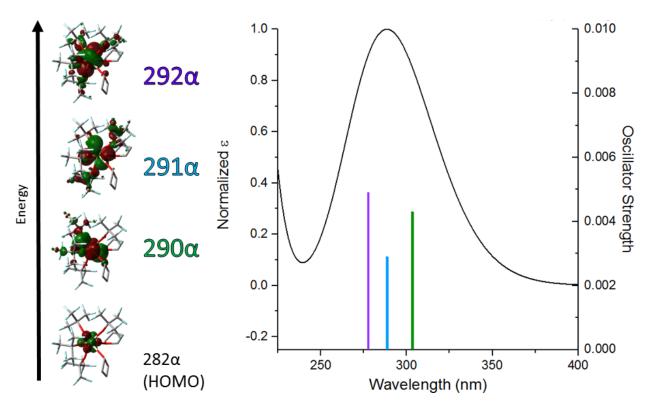


Figure 10. Absorption transitions for 1-Ce and corresponding donor/acceptor orbitals.

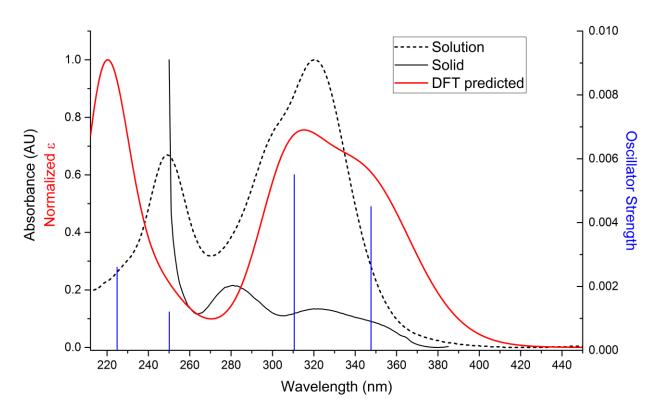


Figure 11. Overlay of solution, solid, and simulated absorption spectra for 3-Ce. Absorption maxima: 250 and 320 nm (solution); 281 and 320 nm (solid); 221 and 315 nm (calculated).

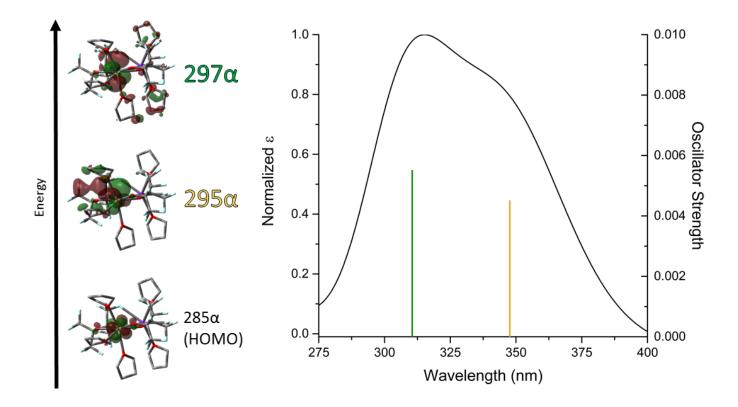


Figure 12. Absorption transitions for 3-Ce and corresponding donor/acceptor orbitals.

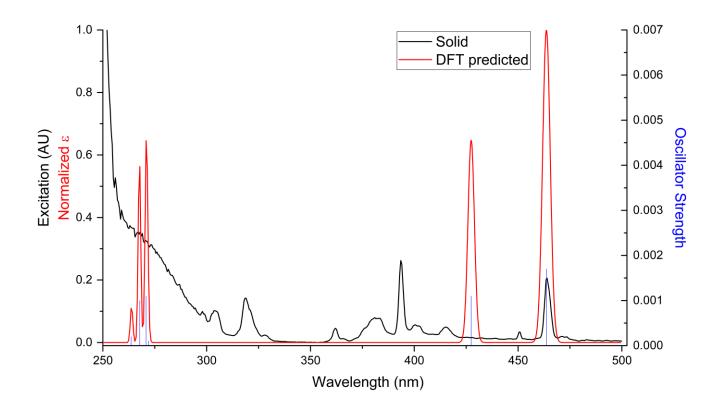


Figure 13. Solid-state excitation (black) and TD-DFT simulated absorption (red) spectra for **2-Eu**, with individual excitations shown in blue. Key excitations in the solid state appear at 395 and 464 nm, with corresponding simulated transitions appearing at 441 and 463 nm.

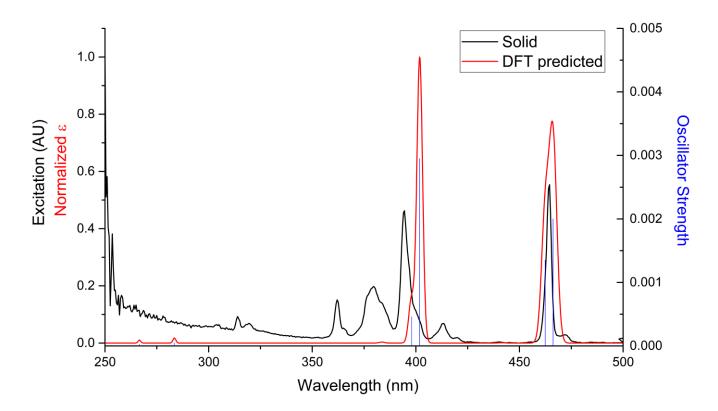


Figure 14. Solid-state excitation (black) and TD-DFT simulated absorption (red) spectra for **4-Eu**, with individual excitations shown in blue. Key excitations in the solid state appear at 395 and 464 nm, with corresponding simulated transitions appearing at 401 and 466 nm.

For Table of Contents Only

Graphic

<mark>Synopsis</mark>

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