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## Highly entangled tensors

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#### **ABSTRACT**

A geometric measure for the entanglement of a unit length tensor  $T \in (\mathbb{C}^n)^{\otimes k}$  is given by  $-2\log_2 \|T\|_{\sigma}$ , where  $\|\cdot\|_{\sigma}$  denotes the spectral norm. A simple induction gives an upper bound of  $(k-1) \log_2(n)$ for the entanglement. We show the existence of tensors with entanglement larger than  $k \log_2(n) - \log_2(k) - o(\log_2(k))$ . Friedland and Kemp (Proc. AMS, 2018) have similar results in the case of symmetric tensors. Our techniques give improvements in this case.

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#### 1. Introduction

We consider the tensor product  $V = V_1 \otimes V_2 \otimes \cdots \otimes V_k$  of finite dimensional Hilbert spaces  $V_1, V_2, \dots, V_k$ . A tensor of the form  $T = v_1 \otimes v_2 \otimes \dots \otimes v_k \in V$  is called a *simple* tensor. A general tensor  $T \in V$  is far from simple. Each vector space  $V_i$  is equipped with positive definite hermitian form  $\langle \cdot, \cdot \rangle$ . The space V is again a Hilbert space, and its positive definite hermitian form  $\langle \cdot, \cdot \rangle$  has the property  $\langle v_1 \otimes v_2 \otimes \cdots \otimes v_k, w_1 \otimes w_2 \otimes \cdots \otimes w_k \rangle =$  $\prod_{i=1}^k \langle v_i, w_i \rangle$ . The euclidean norm on a Hilbert space is defined by  $||v|| = \sqrt{\langle v, v \rangle}$ . We call a tensor  $T \in V$  a unit length tensor if ||T|| = 1.

In quantum physics, a unit length tensor  $T \in V$  corresponds to a pure multipartite state. Such a pure multipartite state  $T \in V$  is called *separable* if T is a simple tensor, and *entangled* if T is not a simple tensor. In the theory of computation, there seems to be a significant upgrade in speed offered by quantum algorithms in comparison with classical ones. A large part of this can be attributed to entanglement.

Various measures for entanglement have been studied owing to different motivations and perspectives. We refer the interested reader to the excellent surveys [1-3]. Using the geometric measure of entanglement (considered previously [4-6]), Gross, Flammia and Eisert [7] argue that most quantum states are too entangled to be useful. On the other hand, observable states for any quantum computer based on photon interactions are symmetric quantum states. For symmetric tensors, the entanglement is much smaller [8]. Friedland and Kemp [8] show that most symmetric quantum states are close to being maximally entangled. In this paper, we generalize and improve some of the results in Refs. [7,8].<sup>1</sup>

We consider here two entanglement measures that are related to (non-euclidean) norms on V. Specifically we consider nuclear and spectral norms. These norms have applications to tensor completion and low rank tensor approximation, see, e.g. [9–11]. In general, the computation of these norms is NP-hard [12,13]. In Nie [14], computation of nuclear norms for symmetric tensors is addressed.

**Definition 1.1:** The *spectral norm*  $||T||_{\sigma}$  of a tensor  $T \in V$  is defined as the maximal value of  $|\langle T, u \rangle|$ , where u ranges over all simple tensors with ||u|| = 1.

**Definition 1.2:** The *nuclear norm*  $||T||_{\star}$  of a tensor  $T \in V$  is the minimal value of  $||u_1|| + ||u_2|| + \cdots + ||u_r||$  over all decompositions  $T = u_1 + u_2 + \cdots + u_r$  where r is a nonnegative integer and  $u_1, u_2, \ldots, u_r$  are simple tensors.

The maximum in Definition 1.1 exists because the unit sphere is compact, and one can show that the minimum in Definition 1.2 exists as well. The spectral and nuclear norms are dual to each other [15]. In particular, we have

$$|\langle T, S \rangle| \le ||T||_{\sigma} ||S||_{\star}. \tag{1}$$

The following measures of entanglement were considered in Ref. [15].

**Definition 1.3:** For a unit length tensor *T* we define two entanglement measures by

$$E(T) = -2\log_2(||T||_{\sigma})$$
 and  $F(T) = 2\log_2(||T||_{\star})$ .

From (1), it follows that  $||T||_{\sigma}||T||_{\star} \geq \langle T, T \rangle = 1$  and  $F(T) \geq E(T)$ . From the Cauchy–Schwarz inequality follows  $||T||_{\sigma} \leq ||T|| = 1$  and  $E(T) \geq 0$ . It is also known that the maximum possible entanglement in both measures is the same [15, Corollary 6.1]. The measure E(T) is the geometric measure of entanglement that we mentioned before.

### 1.1. General tensors

We consider the setting when  $V_1 = V_2 = \cdots = V_k = \mathbb{C}^n$ , and  $V = \mathbb{C}^n \otimes \mathbb{C}^n \otimes \cdots \otimes \mathbb{C}^n = (\mathbb{C}^n)^{\otimes k}$ . Upper bounds on entanglement can be obtained quite easily.

**Proposition 1.4 (Qi et al. [16]):** For all unit length tensors  $T \in (\mathbb{C}^n)^{\otimes k}$ , we have

$$E(T) \le F(T) \le (k-1)\log_2(n).$$

Qi et al. [16] obtained the upper bound for E(T). The same upper bound for F(T) follows easily because the maximum possible entanglement in both the measures is the same [15, Corollary 6.1].

We adapt a technique often found in coding theory to show the existence of tensors whose entanglement is very close to the upper bounds given above.



**Theorem 1.5:** Let  $n, k \geq 2$ , and let

$$C = \min\{\|T\|_{\sigma} \mid T \in (\mathbb{C}^n)^{\otimes k} \text{with} \|T\| = 1\}.$$

Then for any  $0 < \varepsilon < 1$ , we have

$$C^2 \le \frac{k(n-1)\ln\left(\frac{4}{\varepsilon}\right)}{(n^k-1)(1-\varepsilon)^k}.$$

**Corollary 1.6:** Let  $E_{\text{max}} := \max\{E(T) \mid T \in (\mathbb{C}^n)^{\otimes k} \text{with} ||T|| = 1\}$  denote the maximum entanglement of a unit length tensor in  $(\mathbb{C}^n)^{\otimes k}$ . We show the following lower bounds for  $E_{\text{max}}$ .

(1) Let n > 2 be fixed. Then,

$$E_{\text{max}} \ge (k-1)\log_2(n) - \log_2(k) - 1 - \log_2(\ln(k) + 3).$$

Further, if we take k > 21, then,

$$E_{\text{max}} \ge (k-1)\log_2(n) - \log_2(k) - \log_2(\ln(k)) - 2.$$

(2) On the other hand, suppose  $k \ge 2$  is fixed. Then,

$$E_{\text{max}} \ge (k-1)\log_2(n) + O(1).$$

The result in Ref. [7] is a concentration of measure result in the case of n = 2. With a slight modification of the argument we can also recover such a result. Note that the set of all unit length tensors in  $(\mathbb{C}^2)^{\otimes k}$  is just the (real) sphere  $S^{2(2^k)-1} \subset (\mathbb{C}^2)^{\otimes k}$ . In particular, the set of unit length tensors is a sphere which is compact, and has a standard measure on it.

**Corollary 1.7:** Let n = 2, and  $k \ge 4$ . The fraction of unit length tensors T in  $(\mathbb{C}^n)^{\otimes k}$  such that

$$E(T) < k - \log_2(k) - \log_2(\ln(k)) - 3$$

is at most  $e^{-k}$ .

This improves significantly the bound of  $k-2\log_2(k)-3$  of Gross, Flammia and Eisert [7, Theorem 2]. Similar results can be shown for higher n but we omit the details. We remark here that while we improve the bound, the fraction of tensors for which our bound holds is larger. In this sense, both results are incomparable.

So, most tensors have a high entanglement. Yet, finding an explicit<sup>2</sup> tensor with a high entanglement seems quite difficult - this phenomenon is very familiar to researchers in computational complexity. For many complexity measures, one can show that most instances (or a random instance) has high complexity. However, constructing 'explicit' instances with high complexity is extremely challenging, and in many cases still unresolved. Example where explicit instances with high complexity are still far beyond current techniques include the celebrated P vs NP problem (and its algebraic analog VP vs VNP), tensor rank and Waring rank.

Loosely speaking, the geometric measure of entanglement is a measure of complexity. It is an interesting problem to construct explicit examples with a large entanglement.

Problem 1.8: Construct explicit tensors with large entanglement.

We give a construction based on the determinants that constructs explicit tensors whose entanglement is quite large, but still falls short of the bounds in Corollaries 1.6 and 1.7. We define the determinant tensor

$$\det_{n} = \sum_{\pi \in S_{n}} e_{\pi(1)} \otimes e_{\pi(2)} \otimes \cdots \otimes e_{\pi(n)} \in (\mathbb{C}^{n})^{\otimes n}.$$

Consider the tensor  $\det_{n^p} \in (\mathbb{C}^{n^p})^{\otimes n^p}$ . Identifying  $\mathbb{C}^{n^p}$  with  $(\mathbb{C}^n)^{\otimes p}$ , we can think of  $\det_{n^p}$  as a tensor in  $(\mathbb{C}^n)^{\otimes pn^p}$ . We define

$$T_{n,p} = \frac{1}{\sqrt{n^p!}} \det_{n^p} \in (\mathbb{C}^n)^{\otimes pn^p}.$$

Since  $\|\det_{n^p}\| = \sqrt{n^p!}$ , we have  $\|T_{n,p}\| = 1$ .

**Proposition 1.9:** Let n be fixed. For the unit length tensor  $T_{p,n} \in (\mathbb{C}^n)^{\otimes k}$ , with  $k = pn^p$ , we have

$$E(T_{n,p}) \ge k \log_2(n) - o(k).$$

## 1.2. Symmetric tensors

There is a natural action of  $\Sigma_m$ , the symmetric group on m letters on  $(\mathbb{C}^n)^{\otimes m}$  by permuting the tensor factors. A tensor  $T \in (\mathbb{C}^n)^{\otimes m}$  is called a symmetric tensor if it is invariant under the action of  $\Sigma_m$ . We consider the space of symmetric tensors  $S^m(\mathbb{C}^n) \subseteq (\mathbb{C}^n)^{\otimes m}$ . One can identify in a standard way  $S^m(\mathbb{C}^n)$  with polynomials of degree m in n variables.

Symmetric tensors cannot have very large entanglement [8]. We restrict our attention to symmetric tensors of length 1, which we denote by  $S_1^m(\mathbb{C}^n)$ .

**Theorem 1.10 (Friedland and Kemp [8]):** For all  $T \in S_1^m(\mathbb{C}^n)$ , we have

$$E(T) \leq \log_2(d_{n,m}),$$

where  $d_{n,m} = {m+n-1 \choose m}$ .

For fixed n, the existence of tensors  $T \in S_1^m(\mathbb{C}^n)$  with  $E(T) \ge \log_2(d_{n,m}) - \log_2(\log_2 d_{n,m}) - O(1)$  for m sufficiently large is also shown in [8]. Our techniques can be extended to this case and yield improvements. Let  $d_{n,m} = \binom{m+n-1}{m}$ .

**Theorem 1.11:** Let  $C_s = \min\{||T||_{\sigma} \mid T \in S_1^m(\mathbb{C}^n)\}$ . Then for any  $0 < \varepsilon < 1$ , we have

$$C_s^2 \le m^2 \varepsilon^2 + \frac{(n-1)\ln(4/\varepsilon)}{d_{n,m}-1}.$$

Let us define

$$E_{\max}^s = \max\{E(T) \mid T \in S_1^m(\mathbb{C}^n)\}.$$

**Corollary 1.12:** Let  $n \ge 2$  be fixed. Then for m sufficiently large, we have

$$E_{\text{max}}^{s} \ge \log_2(d_{n,m}) - \log_2(\ln(d_{n,m})) - \log_2(n).$$

For n = 2, we obtain a slightly stronger result.

**Corollary 1.13:** Let n = 2. Then for all m, we have

$$E_{\text{max}}^{s} \ge \log_2(m) - \log_2(1 + \ln(4m\sqrt{m})).$$

The strategy in Ref. [8] closely resembles that of Ref. [7]. One significant difficulty they must overcome is to construct ' $\epsilon$ -nets'. We note that our methods bypass the need for this construction. Just like the results in Ref. [7], the results in Ref. [8] are stronger than just showing the existence of tensors with large entanglement. They show a (quantitative) concentration of measure result that says that most symmetric tensors are maximally entangled, see [8, Theorem 1.2]. With our techniques, we can improve upon these results as well by pursuing a strategy similar to Corollary 1.7, but we omit the details.

### 2. Preliminaries

Let  $S^n \subset \mathbb{R}^{n+1}$  denote the *n*-sphere. Observe that the set of unit vectors in  $\mathbb{C}^n = \mathbb{R}^{2n}$  is simply  $S^{2n-1}$ . So, the set of unit length tensors in  $(\mathbb{C}^n)^{\otimes k} = \mathbb{C}^{n^k}$  is just the sphere  $S^{2n^k-1}$ .

**Definition 2.1:** For  $v \in S^{2n-1}$ , i.e. v of unit length in  $\mathbb{C}^n$ , we define

$$B(\nu,\varepsilon) = \{ w \in S^{2n-1} \, | \, |\langle \nu, w \rangle|^2 \geq 1 - \varepsilon \} \subseteq S^{2n-1}.$$

We call  $B(v, \varepsilon)$  the  $\varepsilon$ -ball around v.

**Remark 2.2:** Note that the  $\varepsilon$ -ball as defined above is not a ball in the usual metric. However, it is better suited for our purposes since the spectral norm is defined in terms of inner products.

**Proposition 2.3:** 

$$\operatorname{vol}(B(\nu,\varepsilon)) = \frac{2\pi^n \varepsilon^{n-1}}{(n-1)!}.$$

Before proving the proposition, let us observe that B(v, 1) is the entire sphere  $S^{2n-1}$ . The volume of the sphere is usually calculated by a recursion. As a sanity check, when we plug in  $\varepsilon = 1$  in the above formula, we do recover the (well known) volume of the sphere, i.e.

$$vol(S^{2n-1}) = vol(B(\nu, 1)) = \frac{2\pi^n}{(n-1)!}$$
 (2)

## Corollary 2.4: We have

$$\frac{\operatorname{vol}(B(\nu,\varepsilon))}{\operatorname{vol}(S^{2n-1})} = \varepsilon^{n-1}.$$

**Proof of Proposition 2.3:** Without loss of generality, let us assume  $v = (0, 0, \dots, 0, 1)$  $\in \mathbb{C}^n$ . Then  $B(v, \varepsilon) = \{w \mid |w_n|^2 \ge 1 - \varepsilon\} \subseteq S^{2n-1}$ . In real coordinates, this is the surface described by

(A) 
$$x_1^2 + y_1^2 + x_2^2 + \dots + x_n^2 + y_n^2 = 1$$
 and

(A) 
$$x_1^2 + y_1^2 + x_2^2 + \cdots + x_n^2 + y_n^2 = 1$$
 and  
(B)  $x_n^2 + y_n^2 \ge 1 - \varepsilon$  (or equivalently,  $x_1^2 + y_1^2 + \cdots + x_{n-1}^2 + y_{n-1}^2 \le \varepsilon$ ).

If we restrict to  $y_n \ge 0$ , we get precisely half the surface, and we compute the volume of this half of the surface by writing it as a parametrized surface. We write

$$y_n = f(x_1, y_1, \dots, x_{n-1}, y_{n-1}, x_n) = \sqrt{1 - x_n^2 - R^2} = \sqrt{1 - x_1^2 - y_1^2 - \dots - x_n^2},$$

where  $R = \sqrt{x_1^2 + y_1^2 + \dots + x_{n-1}^2 + y_{n-1}^2}$ . Now, the domain for the parametrized surface is  $\{(x_1, y_1, \dots, x_n, y_n) \mid R^2 \le \varepsilon, -\sqrt{1 - R^2} \le x_n \le \sqrt{1 - R^2}\}$ . The volume of this surface is given by the following formula:

$$\int_{D(0,\sqrt{\varepsilon})} \int_{x_n = -\sqrt{1 - R^2}}^{\sqrt{1 - R^2}} \left( \sqrt{1 + \left(\frac{\partial f}{\partial x_1}\right)^2 + \left(\frac{\partial f}{\partial y_1}\right)^2 + \dots + \left(\frac{\partial f}{\partial y_{n-1}}\right)^2 + \left(\frac{\partial f}{\partial x_n}\right)^2} \right)$$

$$\times dx_n dx_1 dy_1 \dots dy_{n-1},$$

where  $D(0,\sqrt{\varepsilon})=\{(x_1,x_2,\ldots,x_{n-1},y_{n-1})\mid x_1^2+y_1^2+\cdots+y_{n-1}^2\leq \varepsilon\}$ . We compute  $(\partial f/\partial y_i)=-y_i/f$ , and  $(\partial f/\partial x_i)=-x_i/f$ , and so we get that the integrand is simply  $1/f=1/\sqrt{1-x_n^2-R^2}$ . In other words, we have

$$\int_{D(0,\sqrt{\varepsilon})} \int_{x_n=-\sqrt{1-R^2}}^{\sqrt{1-R^2}} \frac{1}{\sqrt{1-R^2-x_n^2}} dx_n dx_1 dy_1 \dots dy_{n-1},$$

But now applying the formula  $\int_{-a}^{a} (1/\sqrt{a^2-x^2}) dx = [\sin^{-1}(x/a)]_{x=-a}^{x=a} = \pi$ , for  $a = \sqrt{1 - R^2}$ , the integral simplifies to

$$\int_{D(0,\sqrt{\varepsilon})} \pi \, \mathrm{d}x_1 \, \mathrm{d}y_1 \dots \mathrm{d}y_{n-1} = \pi \cdot \mathrm{vol}(D(0,\sqrt{\varepsilon})) = \pi \left(\frac{\pi^{n-1}\varepsilon^{n-1}}{(n-1)!}\right) = \frac{\pi^n \varepsilon^{n-1}}{(n-1)!}.$$

This concludes the computation for precisely half the surface, and so we multiply by two to get the required result.

**Proposition 2.5:** Let  $v, z, w \in S^{2n-1} \subset \mathbb{C}^n$  such that  $|\langle v, z \rangle|^2 \geq 1 - \varepsilon$  and  $|\langle z, w \rangle|^2 \geq 1 - \varepsilon$  $1-\varepsilon$ . Then  $|\langle v,w\rangle|^2 > 1-4\varepsilon$ .

**Proof:** By symmetry, we may assume  $|\langle z, w \rangle|^2 \ge |\langle v, z \rangle|^2 = 1 - \delta$  for some  $\delta \le \varepsilon$ . Without loss of generality, let us assume  $\nu = (0, \dots, 0, 1) \in \mathbb{C}^n$ . Then we have  $\sqrt{1 - \delta} = |\langle \nu, z \rangle|$ 



 $|z_n|$ . For  $x=(x_1,\ldots,x_{n-1},x_n)\in\mathbb{C}^n$ , define  $\widetilde{x}=(x_1,\ldots,x_{n-1})\in\mathbb{C}^{n-1}$ . Hence we have  $z = (\widetilde{z}, z_n)$  and  $w = (\widetilde{w}, w_n)$ .

Now, we have

$$\sqrt{1-\delta} \le |\langle z, w \rangle| 
= |\langle \widetilde{z}, \widetilde{w} \rangle + z_n \overline{w_n}| 
\le |\langle \widetilde{z}, \widetilde{w} \rangle| + |z_n \overline{w_n}| 
\le ||\widetilde{z}|| \cdot ||\widetilde{w}|| + (\sqrt{1-\delta})|w_n| 
= \sqrt{\delta} \sqrt{1 - |w_n|^2} + (\sqrt{1-\delta})|w_n|$$

The last equality follows from the fact that for  $x \in \mathbb{C}^n$ , we have  $||x||^2 = ||\widetilde{x}||^2 + |x_n|^2$ . In particular, we can rewrite the inequality as

$$\sqrt{1-\delta}(1-|w_n|) \le \sqrt{\delta}\sqrt{1-|w_n|^2}$$

On squaring the terms, we get

$$(1 - \delta)(1 - |w_n|)^2 \le \delta(1 - |w_n|^2).$$

Now, dividing both sides by  $1 - |w_n|$ , we get

$$(1 - \delta)(1 - |w_n|) \le \delta(1 + |w_n|).$$

Rearranging the terms, we get  $|w_n| \ge 1 - 2\delta$ . We get the required conclusion since we have  $|\langle v, w \rangle|^2 = |w_n|^2 > (1 - 2\delta)^2 > 1 - 4\delta > 1 - 4\varepsilon.$ 

In the next proposition, we give an upper bound on the number of  $\varepsilon$ -balls needed to cover all of  $S^{2n-1}$ . In order to do this, we borrow a standard technique from coding theory and modify it appropriately.

**Proposition 2.6:** There exists  $v_1, v_2, \ldots, v_N \in S^{2n-1}$  such that  $\bigcup_{i=1}^N B(v_i, \varepsilon)$  is all of  $S^{2n-1}$ for some  $N \leq (4/\varepsilon)^{n-1}$ 

**Proof:** Let  $B(v_1, \varepsilon/4), B(v_2, \varepsilon/4), \ldots, B(v_N, \varepsilon/4)$  be a maximal collection of noninteresecting  $\varepsilon/4$ -balls on  $S^{2n-1}$ . Observe that  $N \leq \operatorname{vol}(S^{2n-1})/\operatorname{vol}(B(v,\varepsilon 4)) = (4/\varepsilon)^{n-1}$ . Now for any  $w \in S^{2n-1}$ , we have that  $B(w, \varepsilon/4) \cap B(v_i, \varepsilon/4)$  is non-empty for some i. Hence  $\exists z \in S^{2n-1}$  such that  $|\langle z, w \rangle|^2 > 1 - \varepsilon/4$  and  $|\langle v_i, z \rangle|^2 > 1 - \varepsilon/4$ , and hence by Proposition 2.5, we have that  $w \in B(v_i, \varepsilon)$ . Hence  $\bigcup_i B(v_i, \varepsilon)$  is all of  $S^{2n-1}$ .

#### 3. Main results

## 3.1. Lower bounds on entanglement

In this section, we will give proofs of our main lower bound results on entanglement for general tensors. For this section, we will fix  $v_1, v_2, \dots, v_N \in S^{2n-1} \subset \mathbb{C}^n$  such that  $\bigcup_{i=1}^{N} B(v_i, \varepsilon)$  is all of  $S^{2n-1}$ . The following proposition is the key technical result to derive Theorem 1.5.

**Proposition 3.1:** Let  $C = \min\{||T||_{\sigma} \mid T \in S^{2n^k-1} \subset (\mathbb{C}^n)^{\otimes k}\}$ . Then, for all  $T \in S^{2n^k-1}$ , we have  $|\langle T, v_{i_1} \otimes v_{i_2} \otimes \cdots \otimes v_{i_k} \rangle|^2 \geq C^2(1-\varepsilon)^k$  for some  $1 \leq i_1, i_2, \ldots, i_k \leq N$ .

**Proof:** By definition, for any  $T \in (\mathbb{C}^n)^{\otimes k}$ , we have  $|\langle T, a_1 \otimes a_2 \otimes \cdots \otimes a_k \rangle| \geq C$  for some simple tensor  $a_1 \otimes a_2 \otimes \cdots \otimes a_k$ .

The map  $v \mapsto \langle T, v \otimes a_2 \otimes \cdots \otimes a_k \rangle$  defines a element of  $(\mathbb{C}^n)^*$ , which we denote by f. Using the linear isomorphism  $\mathbb{C}^n$  with  $(\mathbb{C}^n)^*$  given by the inner product, we get a vector  $v_f$  such that  $f(w) = \langle v_f, w \rangle$ . Since  $\bigcup_{i=1}^N B(v_i, \varepsilon)$  is all of  $S^{2n-1}$ , there exists  $v_{i_1}$  such that

$$|\langle v_f, v_{i_1} \rangle| \ge (\sqrt{1-\varepsilon}) \|v_f\| \ge (\sqrt{1-\varepsilon}) |\langle v_f, a_1 \rangle| \ge C\sqrt{1-\varepsilon}.$$

Hence

$$|\langle T, \nu_{i_1} \otimes a_2 \otimes \cdots \otimes a_k \rangle| = |\langle \nu_f, \nu_{i_1} \rangle| \ge C\sqrt{1 - \varepsilon}.$$

Repeating the argument for all the tensor factors, we get a tensor  $v_{i_1} \otimes v_{i_2} \otimes \cdots \otimes v_{i_k}$  such that

$$|\langle T, \nu_{i_1} \otimes \nu_{i_2} \otimes \cdots \otimes \nu_{i_k} \rangle| \geq C(\sqrt{1-\varepsilon})^k$$
.

We write  $[N] := \{1, 2, ..., N\}$  in the following corollary.

**Corollary 3.2:** We have  $\bigcup_{i_1,i_2,...,i_k\in[N]^k} B(v_{i_1}\otimes v_{i_2}\otimes \cdots \otimes v_{i_k}, 1-C^2(1-\varepsilon)^k)$  is all of  $S^{2n^k-1}$ .

Using Corollary 2.4, we get

Corollary 3.3: We have

$$N^k \left( (1 - C^2 (1 - \varepsilon)^k)^{n^k - 1} \right) \ge 1.$$

We have all the tools required to prove Theorem 1.5.

**Proof of Theorem 1.5:** We can use the above inequality to get an upper bound for *C*. We have

$$1 \le N^{k} (1 - C^{2} (1 - \varepsilon)^{k})^{n^{k} - 1}$$
$$\le \left(\frac{4}{\varepsilon}\right)^{k(n-1)} (1 - C^{2} (1 - \varepsilon)^{k})^{n^{k} - 1}.$$

Taking logarithms we get

$$0 \le k(n-1)\ln\left(\frac{4}{\varepsilon}\right) + (n^k - 1)\ln(1 - C^2(1-\varepsilon)^k)$$
  
$$\le k(n-1)\ln\left(\frac{4}{\varepsilon}\right) + (n^k - 1)(-C^2(1-\varepsilon)^k).$$

This gives us

$$C^2 \le \frac{k(n-1)\ln\left(\frac{4}{\varepsilon}\right)}{(n^k-1)(1-\varepsilon)^k},$$

as required.

With the proof of Theorem 1.5, we now proceed to deduce Corollaries 1.6 and 1.7. These are a bit computational, and we will try to keep the computations transparent and succinct.

**Proof of Corollary 1.6:** We first prove part (1). From Theorem 1.5, we get

$$E_{\max} \geq \log_2\left(\frac{n^k-1}{n-1}\right) - \log_2(k) + \log_2(1-\varepsilon)^k - \log_2\left(\ln\left(\frac{4}{\varepsilon}\right)\right).$$

by taking a logarithm on both sides.

First, observe that  $\log_2((n^k - 1)/(n - 1)) \ge \log_2(n^{k-1}) = (k - 1)\log_2(n)$ .

Next, set  $\varepsilon = \delta/k$ . So,  $(1 - \varepsilon)^k \ge 1 - k\varepsilon = 1 - \delta$ . Also, note that  $\log_2(\ln(4/\varepsilon))$ =  $\log_2(\ln(k) + \ln(4/\delta))$ . Now, set  $\delta = 4/e^3$ . Then

$$\log_2(1 - \varepsilon)^k \ge \log_2(1 - \delta) = \log_2(1 - 4/e^3) \ge -1,$$

and

$$\log_2(\ln(k) + \ln(4/\delta)) = \log_2(\ln(k) + 3) \le \log_2(2\ln(k)) = \log_2(\ln(k)) + 1.$$

The inequality above follows because we assume  $k \ge 21$  and consequently  $\ln(k) \ge 3$ . Thus, we get

$$E_{\text{max}} \ge (k-1)\log_2(n) - \log_2(k) - \log_2(\ln(k)) - 2$$

To prove part (2), we just observe that if k is fixed, then  $k \ln(4\varepsilon)/(1-\varepsilon)^k$  is just a constant for fixed  $\varepsilon$ . We get the required result by taking logarithms in the conclusion of Theorem 1.5

We can now prove Corollary 1.7

**Proof of Corollary 1.7:** Let D be the infimum of all numbers d such that the fraction of unit length tensors T having  $||T||_{\sigma} > d$  is at most  $e^{-k}$ . We want a lower bound on  $-2\log_2(D)$ . Similar to Corollary 3.3, we can deduce

$$N^k\left((1-D^2(1-\varepsilon)^k)^{n^k-1}\right) \ge e^{-k}.$$

Again, using that  $N \leq (4/\varepsilon)$ , and following the steps in the proof of Theorem 1.5, we get

$$D^{2} \leq \frac{k \ln\left(\frac{4}{\varepsilon}\right) + k}{(2^{k} - 1)(1 - \varepsilon)^{k}}$$

Taking natural logarithms on both sides, we get

$$-2\log_2(D) \geq \log_2(2^k - 1) + \log_2(1 - \varepsilon)^k - \log(k) - \log_2(\ln(4/\varepsilon) + 1).$$

First, note that for  $k \ge 4$ , we have  $\log_2(2^k - 1) \ge k - \frac{1}{4}$ . Now, set  $\epsilon = 1/4k$ . Then we get  $\log_2(1-\varepsilon)^k \ge \log_2(1-k\varepsilon) = \log_2(3/4) > -3/4$ . We also get  $\log_2(\ln(4/\varepsilon) + 1)$  $= \log_2(\ln(16k) + 1).$ 

Now, for  $k \ge 4$ , one can check that  $\ln(16k) + 1 = \ln(k) + \ln(16) + 1 \le 4 \ln(k)$ . Thus  $\log_2(\ln(16k) + 1) \le \log_2(\ln(k)) + 2$ . Thus, we have:

$$-2\log_2(D) \ge (k - 1/4) - \log(k) - (3/4) - \log_2(\ln(k)) - 2.$$

Thus, we get

$$-2\log_2(D) \ge k - \log(k) - \log_2(\ln(k)) - 3$$

as required.

## 3.2. An explicit tensor with high entanglement

In this section, we prove Proposition 1.9, i.e. we give a lower bound on the entanglement of an explicit tensor ( $T_{n,p}$  as defined in Proposition 1.9). Recall that

$$\det_{n} = \sum_{\pi \in S_{n}} e_{\pi(1)} \otimes e_{\pi(2)} \otimes \cdots \otimes e_{\pi(n)} \in (\mathbb{C}^{n})^{\otimes n}.$$

We know that  $\|\det_n\|_{\sigma}=1$ , by [17]. Consider the tensor  $\det_{n^p}\in(\mathbb{C}^{n^p})^{\otimes n^p}$ . By identifying  $\mathbb{C}^{n^p}$  with  $(\mathbb{C}^n)^{\otimes p}$ , we identify  $(\mathbb{C}^{n^p})^{\otimes n^p}$  with  $(\mathbb{C}^n)^{\otimes pn^p}$ . Thus we can consider  $\det_{n^p}$  as a tensor in  $(\mathbb{C}^n)^{\otimes pn^p}$ . Note that considering  $\det_{n^p}$  as a tensor in  $(\mathbb{C}^n)^{\otimes pn^p}$  cannot increase the spectral norm since the set of simple tensors of unit length in  $(\mathbb{C}^n)^{\otimes pn^p}$  is a subset of simple tensors of unit length in  $(\mathbb{C}^n)^{\otimes pn^p}$ .

On the other hand, if  $e_1, e_2, \ldots, e_n$  is a standard basis for  $\mathbb{C}^n$ , we define  $e_I = e_{i_1} \otimes \cdots \otimes e_{i_p}$  for  $I = (i_1, \ldots, i_p) \in [n]^p$  where  $[n] = \{1, 2, \ldots, n\}$ . Fix any bijection  $h : [n^p] \to [n]^p$ . We define

$$u = \bigotimes_{i=1}^{n^p} e_{h(i)}.$$

Clearly, u is a simple tensor of unit length in  $(\mathbb{C}^n)^{\otimes pn^p}$ , and  $|\langle \det_{n^p}, u \rangle| = 1$ . Hence,  $\|\det_{n^p}\|_{\sigma} = 1$ , when considered as a tensor in  $(\mathbb{C}^n)^{\otimes pn^p}$ . Now, consider the unit length tensor  $T_{n,p}$ .

**Proof of Proposition 1.9:** By the above discussion, we have  $||T||_{\sigma} = 1/\sqrt{n^p!}$ . Hence, we get  $E(T) = 2\log_2 \sqrt{n^p!} = \log_2(n^p!)$ . Using Stirling's formula, we compute

$$\log_2(n^p!) = n^p \log_2(n^p) - \log_2(e)(n^p) + O(\log_2(n))$$
  
=  $pn^p \log_2(n) - o(pn^p)$ .

The proposition follows by replacing  $pn^p$  with k.

## 4. Symmetric tensors

This section is devoted to the study of entanglement for symmetric tensors. The overall strategy parallels the case of general tensors. However, there is one additional ingredient that we will need, and that is Banach's theorem below.



Analogous to the case of general tensors, we define

$$C_s := \min\{\|T\|_{\sigma}^2 \mid T \in S_1^m(\mathbb{C}^n)\}.$$

Our goal is lower bounds on  $C_s$ . The following is the aforementioned result of Banach [18] (see also [8, Theorem 2.7]).

**Theorem 4.1 (Banach):** For  $T \in S^m(\mathbb{C}^n)$ , we have  $||T||_{\sigma} = \max\{|\langle T, v^{\otimes m} \rangle| : ||v|| = 1\}$ .

Observe that by Proposition 2.6, we can take  $v_1, \ldots, v_N$ , with  $N \leq (4/\varepsilon)^{n-1}$  such that  $\bigcup_{\alpha=1}^N B(\nu_\alpha, \varepsilon)$  is all of  $\hat{S}^{n-1}$ . For this entire section  $\nu_1, \ldots, \nu_N$  will denote such a choice. Using Banach's theorem above, we deduce the main technical lemma for the purposes of this section.

**Lemma 4.2:** For all  $T \in S_1^m(\mathbb{C}^n)$ , we have  $|\langle T, v_I^{\otimes m} \rangle|^2 \geq C^2 - m^2 \varepsilon^2$  for some  $I \in \{1, 2, \dots, N\}.$ 

**Proof:** It follows from Theorem 4.1 that for any  $T \in S_1^m(\mathbb{C}^n)$ ,  $\exists v \in \mathbb{C}^n$  with ||v|| = 1 such that  $|\langle T, v^{\otimes m} \rangle| \geq C_s$ . By the choice of  $v_1, \ldots, v_N$ , we have  $|\langle v, v_I \rangle|^2 \geq 1 - \varepsilon$  for some  $I \in \{1, 2, \dots, N\}.$ 

Let  $e_1, \ldots, e_n$  denote an orthonormal basis for  $\mathbb{C}^n$ . Without loss of generality, we can assume  $\nu_I=e_n$ . The vector space  $(\mathbb{C}^n)^{\otimes m}$  has an orthonormal basis  $\{e_{i_1i_2...i_m}:=e_{i_1}\otimes e_{i_1}\}$  $\cdots \otimes e_{i_m} \mid 1 \leq i_1, i_2, \ldots, i_m \leq n$ . We can write

$$(\mathbb{C}^n)^{\otimes m} = \mathbb{C}e_{nn-n} \oplus V,$$

where  $V = \bigoplus_{i_1...i_m \neq nn...n} \mathbb{C} e_{i_1...i_m}$ 

Given  $T \in (\mathbb{C}^n)^{\otimes m}$ , we can write  $T = T_{nn...n} + \widetilde{T}$  uniquely where  $T_{nn...n} \in \mathbb{C}e_{nn...n}$  and  $\widetilde{T} \in V$ .

Now, let  $v = (v_1, \dots, v_n)$  be the coordinates of v in the basis  $e_1, \dots, e_n$ . Then, we have

$$1 - \varepsilon \le |\langle v, v_I \rangle|^2 = |\langle v, e_n \rangle|^2 = |v_n|^2.$$

Thus, we have  $|(v^{\otimes m})_{nn...n}|^2 = |v_n^m|^2 = (|v_n|^2)^m \ge (1-\varepsilon)^m \ge 1-m\epsilon$ . In particular, this implies that  $\|v^{\otimes m}\| \leq m\varepsilon$ .

Now, we have

$$\begin{split} C_s^2 &\leq |\langle T, v^{\otimes m} \rangle|^2 \leq |\langle \widetilde{T}, \widetilde{v^{\otimes m}} \rangle + T_{nn...n} \overline{(v^{\otimes m})_{nn...n}}|^2 \\ &\leq |\langle T, \widetilde{v^{\otimes m}} \rangle|^2 + |T_{nn...n}|^2 |(v^{\otimes m})_{nn...n}|^2 \\ &\leq m^2 \varepsilon^2 + |\langle T, v_I^{\otimes m} \rangle|^2 |(v^{\otimes m})_{nn...n}|^2 \\ &\leq m^2 \varepsilon^2 + |\langle T, v_I^{\otimes m} \rangle|^2. \end{split}$$

The first and second inequalities are clear. The third inequality follows from ||T|| = 1,  $\|\widetilde{v^{\otimes m}}\| \leq m\varepsilon$  and the fact that  $T_{nn...n} = \langle T, e_n^{\otimes m} \rangle = \langle T, v_I^{\otimes m} \rangle$  (since  $v_I = e_n$ ). The last inequality follows from  $|v_{nn...n}^{\otimes m}|^2 = |v_n|^{2m} \le 1$ .

Hence, we get

$$|\langle T, v_I^{\otimes m} \rangle|^2 \ge C_s^2 - m^2 \varepsilon^2.$$

Friedland and Kemp [8] show that we have an isometry between  $S_1^m(\mathbb{C}^n)$  and  $S^{2d_{n,m}-1}$ . Along with the above lemma, we obtain:

**Corollary 4.3:** 
$$\bigcup_{\alpha=1}^{N} B(v_{\alpha}^{\otimes m}, 1 - (C_{s}^{2} - m^{2} \varepsilon^{2})) = S_{1}^{m}(\mathbb{C}^{n}) = S^{2d_{n,m}-1}.$$

Using Corollary 2.4, one obtains the following inequality.

**Corollary 4.4:** We have 
$$(4/\varepsilon)^{n-1}(1-(C_s^2-m^2\varepsilon^2))^{d_{n,m}-1} \ge 1$$
.

**Proof of Theorem 1.11:** Applying the natural logarithm to Corollary 4.4, we get

$$0 \le (n-1)\ln(4/\varepsilon) + (d_{n,m}-1)\ln(1 - (C_s^2 - m^2\varepsilon^2))$$
  
$$\le (n-1)\ln(4/\varepsilon) + (d_{n,m}-1)(-(C_s^2 - m^2\varepsilon^2))$$

Hence, we get

$$(d_{n,m}-1)(C_s^2-m^2\varepsilon^2)\leq (n-1)\ln(4/\varepsilon).$$

Equivalently, we get

$$C_s^2 \le m^2 \varepsilon^2 + \frac{(n-1)\ln(4/\varepsilon)}{d_{n,m}-1}$$

as required.

Once again, by taking a natural logarithm on both sides in Theorem 1.11, we obtain:

$$-2\log_2 C_s \ge -\log_2 \left( m^2 \varepsilon^2 + \frac{(n-1)\ln(4/\varepsilon)}{d_{n,m} - 1} \right)$$

Replacing  $\varepsilon$  by  $\delta/m$ , we get

$$E_{\text{max}}^{s} = -2\log_{2} C_{s} \ge -\log_{2} \left(\delta^{2} + \frac{(n-1)\ln(4m/\delta)}{d_{n,m} - 1}\right)$$
(3)

**Proof of Corollary 1.12:** Substituting  $\delta = (d_{n,m} - 1)^{-1/2}$  in the above expression, we get

$$E_{\text{max}}^{s} \ge \log_2 \frac{d_{n,m} - 1}{1 + (n-1)\ln(4m(d_{n,m} - 1)^{1/2})}$$

For *m* sufficiently large, we claim that

$$\frac{d_{n,m} - 1}{1 + (n-1)\ln(4m(d_{n,m} - 1)^{1/2})} \ge \frac{d_{n,m}}{n\ln(d_{n,m})} \tag{4}$$

from which we can conclude that

$$E_{\text{max}}^{s} \ge \log_2(d_{n,m}) - \log_2(\ln(d_{n,m})) - \log_2(n).$$

So, it remains to prove (4) above. For  $m \gg 0$ , the following is easy to verify because  $d_{n,m} = O(m^{n-1}).$ 

$$1 + (n-1)\ln(4m(d_{n,m}-1)^{1/2}) \le 1 + (n-1)\ln(4m(d_{n,m})^{1/2}) \le 0.99n\ln(d_{n,m}).$$

This means that

$$\frac{1 + (n-1)\ln(4m(d_{n,m}-1)^{1/2})}{n\ln(d_{n,m})} \le 0.99 \le 1 - \frac{1}{d_{n,m}} = \frac{d_{n,m}-1}{d_{n,m}},$$

which rearranged gives us (4) as required.

**Proof of Corollary 1.13:** Setting n = 2, (3) simplifies to

$$E_{\max}^{s} \ge -\log_2\left(\delta^2 + \frac{\ln(4m/\delta)}{m}\right) = \log_2(m) - \log_2(m\delta^2 + \ln(4m/\delta)).$$

If we take  $\delta = 1/\sqrt{m}$ , we get

$$E_{\text{max}}^{s} > \log_2(m) - \log_2(1 + \ln(4m\sqrt{m})),$$

**Notes** 

- 1. These results use a concentration of measure inequality ([7, Equation (2)] and [8, Lemma 3.7]) which needs a correction. While this does not disrupt their results significantly, their results must be corrected.
- 2. This can be given a formal definition, i.e. the usual one in complexity.

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#### References

- [1] Amico L, Fazio R, Osterloh A, et al. Entanglement in many-body systems. Rev Mod Phys. 2008;80:517-576.
- [2] Horodecki R, Horodecki P, Horodecki M, et al. Quantum entanglement. Rev Mod Phys. 2009;81:865-942.
- [3] Plenio MB, Virmani SS. An introduction to entanglement theory. In: Andersson E, Ohberg P, editors. Quantum information and coherence. Scottish Graduate Series. Cham: Springer, 2014. p. 173-209.
- [4] Barnum H, Linden N. Monotones and invariants for multi-particle quantum states. J Phys A. 2001;34(35):6787–6805. Quantum information and computation.
- [5] Shimony A. Degree of entanglement. Fundamental problems in quantum theory (Baltimore, MD, 1994). New York: New York Acad. Sci.; 1995. p. 675–679. (Ann. New York acad. sci.; 755).
- [6] Wei T, Goldbart PM. Geometric measure of entanglement and applications to bipartite and multipartite quantum states. Phys Rev A. 2003;68:042307.
- [7] Gross D, Flammia ST, Eisert J. Most quantum states are too entangled to be useful as computational resources. Phys Rev Lett. 2009;102:190501.
- [8] Friedland S, Kemp T. Most Boson quantum states are almost maximally entangled. Proc Amer Math Soc. 2018;146(12):5035-5049.
- [9] Friedland S, Ottaviani G. The number of singular vector tuples and uniqueness of best rank one approximation of tensors. Found Comput Math. 2014;14(6):1209–1242.
- [10] Nie J, Wang L. Semidefinite relaxations for best rank-1 tensor approximations. SIAM J Matrix Anal Appl. 2014;35:1155-1179.
- [11] Yuan M, Zhang C-H. On tensor completion via nuclear norm minimization. 2014. arXiv:1405.1773 [stat.ML].
- [12] Friedland S, Lim LH. The computational complexity of duality. SIAM J Optim. 2016;26(4): 2378-2393.
- [13] Friedland S, Lim LH. Nuclear norm of higher-order tensors. Math Comput. 2018;87(311): 1255-1281.
- [14] Nie J. Symmetric tensor nuclear norms. SIAM J Appl Algebra Geom. 2017;1(1):599–625.
- [15] Derksen H, Friedland S, Lim LH, et al. Theoretical and computational aspects of entanglement. 2017. arXiv:1705.07160 [quant-ph].
- [16] Qi L, Zhang G, Ni G. How entangled can a multi-party system possibly be? Phys Lett A. 2018;382(22):1465-1471.
- [17] Derksen H. On the nuclear norm and the singular value decomposition of tensors. Found Comput Math. 2016;16(3):779-811.
- [18] Banach S. Uber homogene polynome in (L2). Studia Math. 1938;7:36–44.