

Flat-band magnetism and helical magnetic order in Ni-doped SrCo_2As_2

Yu Li^{1,2,*}, Zhonghao Liu^{3,†}, Zhuang Xu⁴, Yu Song¹, Yaobo Huang⁵, Dawei Shen³, Ni Ma³, Ang Li³, Songxue Chi⁶, Matthias Frontzek⁶, Huibo Cao⁶, Qingzhen Huang⁷, Weiye Wang¹, Yaofeng Xie¹, Rui Zhang¹, Yan Rong⁴, William A. Shelton⁸, David P. Young², J. F. DiTusa^{2,‡} and Pengcheng Dai^{1,4,§}

¹Department of Physics and Astronomy, Rice University, Houston, Texas 77005, USA

²Department of Physics and Astronomy, Louisiana State University, Baton Rouge, Louisiana 70803, USA

³State Key Laboratory of Functional Materials for Informatics and Center for Excellence in Superconducting Electronics, SIMIT, Chinese Academy of Sciences, Shanghai 200050, China

⁴Center for Advanced Quantum Studies and Department of Physics, Beijing Normal University, Beijing 100875, China

⁵Shanghai Synchrotron Radiation Facility, Shanghai Institute of Applied Physics, Chinese Academy of Sciences, Shanghai 201204, China

⁶Neutron Scattering Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA

⁷NIST Center for Neutron Research, National Institute of Standards and Technology, Gaithersburg, Maryland 20899, USA

⁸Department of Chemical Engineering, Louisiana State University, Baton Rouge, Louisiana 70803, USA



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A series of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ single crystals was synthesized, allowing a comprehensive phase diagram with respect to field, temperature, and chemical substitution to be established. Our neutron diffraction experiments revealed a helimagnetic order with magnetic moments ferromagnetically aligned in the ab plane and a helimagnetic wave vector of $q = (0, 0, 0.56)$ for $x = 0.1$. The combination of neutron diffraction and angle-resolved photoemission spectroscopy measurements shows that the tuning of a flat band with $d_{x^2-y^2}$ orbital character drives the helimagnetism and indicates the possibility of a quantum order-by-disorder mechanism.

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I. INTRODUCTION

The exploration of materials in proximity to quantum phase transitions is a fruitful area for discovering exotic states of matter due to the strong influence of quantum fluctuations. This idea can be traced to Villain, who coined the name order by disorder when investigating frustrated insulating magnets [1]. The central idea refers to the fact that quantum fluctuations, akin to classical entropic effects, can lift the large degeneracy of ground states and favor a particular order, hence the term order by disorder [2]. This is considered to be the quantum equivalent to the theme used to explain a variety of phenomena, including surface tension and DNA folding, and has been employed to explain unconventional superconductivity [3] and even emergent gravity [4]. There is great recent theoretical interest in itinerant ferromagnets in which quantum fluctuations dominate and drive the system into a variety of novel magnetic states in the vicinity of a quantum critical point [5–7], such as a helical state driven by ferromagnetic (FM) spin fluctuations. Meanwhile, as an origin of ferromagnetism complementary to that described by Nagaoka [8], flat-band physics [9] has garnered attention since it provides fertile ground for diverse interaction-driven quantum phases, including FM order, a Mott insulating phase [10], and superconductivity [11]. Therefore, it is interesting

to search for helimagnetism associated with a flat band lying close to the chemical potential in a material close to a quantum critical point.

In the AFe_2As_2 ($A = \text{Ca, Sr, Ba}$) family of iron pnictides, it is well known that the superconducting electron pairing is attributed to antiferromagnetic (AF) spin fluctuations [12,13]. However, FM spin fluctuations may also be important to electron pairing [14,15] and were recently observed in Co-substituted compounds by nuclear magnetic resonance [16] and neutron scattering experiments [17,18]. Furthermore, electronic structure calculations [19,20] and angle-resolved photoemission spectroscopy (ARPES) experiments indicate that ACo_2As_2 [17,21,22] is in proximity to a FM instability due to the existence of a flat band near the Fermi level, although these materials remain paramagnetic down to 2 K with AF low-energy spin fluctuations [17]. Since chemical substitution can shift the Fermi level relative to the band structure [23], substitution of Co with Ni will drive the system toward a Van Hove singularity associated with the flat band and efficiently promote the FM instability. Therefore, it is interesting to investigate how AF and FM spin fluctuations evolve in $\text{A}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ and to explore the relevant emergent phenomena.

We successfully synthesized single crystals [Fig. 1(a)] of a series of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ with a range of Ni concentrations between 0 and 0.6 and carried out systematic magnetic susceptibility χ and magnetization M measurements. $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ has a ThCr_2Si_2 -type body-centered tetragonal crystal structure with no indication of a structural transition [Fig. 1(b)] [24]. Considering both SrCo_2As_2 [17,19] and SrNi_2As_2 [25] are paramagnetic without magnetic order

*yuli1@lsu.edu

†lzh17@mail.sim.ac.cn

‡ditusa@phys.lsu.edu

§pdai@rice.edu

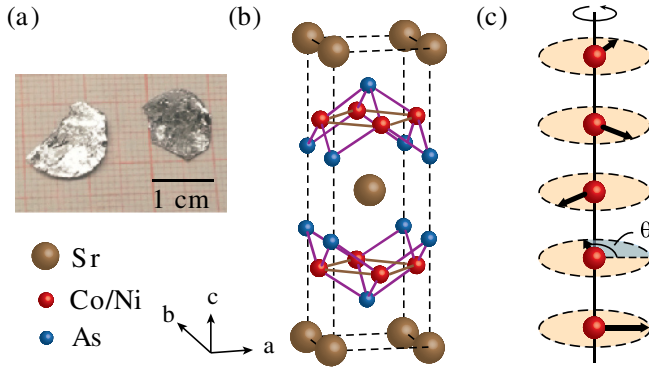


FIG. 1. (a) Single crystals of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$. (b) Tetragonal structure of SrCo_2As_2 . (c) Helical magnetic structure along the c axis with a rotation angle θ between neighboring layers.

at low temperature, it is surprising that we discovered a helical magnetic order in $\text{Sr}(\text{Co}_{0.9}\text{Ni}_{0.1})_2\text{As}_2$ with a magnetic wave vector $\mathbf{q} = (0, 0, 0.56)$ and transition temperature $T_c = 28$ K [Figs. 1(c) and 2(c)]. We find that the transition temperature varies systematically with substitution and has a maximum value at $x = 0.1$ from susceptibility measurements [Fig. 2(c)]. Our results suggest that the ground state is helimagnetic (HM) with moments lying parallel to the ab easy plane while rotating with respect to the c axis in adjacent layers [Fig. 1(c)]. Measurements of $M(H)$ with an in-plane magnetic field at low temperatures ($2 \text{ K} \leq T \leq 4 \text{ K}$) demonstrate a two-step transition from HM order first into a partially polarized magnetic (PPM) phase and ultimately into a fully polarized magnetic (FPM) state. Under a field parallel to the c axis, the magnetization is linearly dependent on H prior to

saturation at the higher field. Combining neutron diffraction and ARPES experiments, we established a close connection between the magnetic moments and a flat band of the $d_{x^2-y^2}$ orbital character. Since it was suggested that BaCo_2As_2 is a quantum paramagnet near a FM critical point [22,26], the helimagnetic order observed in Ni-substituted SrCo_2As_2 may be induced by quantum fluctuations in the vicinity of the putative quantum critical point [2].

II. EXPERIMENTAL DETAILS

A. Sample growth and magnetization measurements

Single crystals of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ were grown from solution using a self-flux method with the ratio $\text{Sr}:(\text{Co}_{1-x}\text{Ni}_x):\text{As} = 1:5:5$. The elements were placed in an aluminum oxide crucible and sealed in an evacuated quartz tube. After heating slowly below 830°C , the mixture was cooked at 1200°C for 20 h and then slowly cooled down to 1050°C at a rate of 3°C/h and then down to 800°C at 10°C/h . Single crystals were obtained by cleaning off the flux. The typical crystals were about 1 cm^2 in size and a few millimeters in thickness [Fig. 1(a)].

Temperature- and field-dependent dc magnetization measurements were performed on a Quantum Design physical property measurement system at Beijing Normal University and a magnetic property measurement system (MPMS) at Louisiana State University.

B. Neutron diffraction

Neutron powder diffraction data for the compound $\text{Sr}(\text{Co}_{0.9}\text{Ni}_{0.1})_2\text{As}_2$ were collected at the NIST Center for Neutron Research on a high-resolution powder neutron diffractometer (BT-1) with neutrons of wavelength $1.5403 \text{ \AA}$ (at 295 K) and $2.0775 \text{ \AA}$ (at 30 and 5 K) produced by using $\text{Cu}(311)$ and $\text{Ge}(311)$ monochromators, respectively. Collimators with horizontal divergences of 15 , 20 , and 7 arc min were used before and after the monochromator and after the sample, respectively. Data were collected in the 2θ range of 3° – 168° with a step size of 0.05° . The structural parameters were refined by Rietveld refinement using the GSAS and EXPGUI programs [27,28]. The atomic neutron scattering factors used in the refinements for Sr, Co, Ni, and As were 0.702 , 0.253 , 1.03 , and $0.658 \times 10^{-12} \text{ cm}$, respectively. About 5 g of $\text{Sr}(\text{Co}_{0.9}\text{Ni}_{0.1})_2\text{As}_2$ single crystals were ground into powder for the neutron powder diffraction measurements. Single-crystal neutron diffraction measurements were carried out on two instruments at the High Flux Isotope Reactor at Oak Ridge National Laboratory, the single-crystal diffractometer HB3A, and the recently upgraded wide-angle neutron diffractometer (WAND2) at the HB-2C beam port. A monochromatic incident beam with a wavelength of $1.546 \text{ \AA}$ was used on the HB3A diffractometer, which is provided by a $\text{Si}(220)$ monochromator [29]. The detector on HB3A is a two-dimensional (2D) scintillation Anger camera with 1-mm spatial resolution. WAND2 uses a monochromatic beam with a wavelength of $1.48 \text{ \AA}$ provided by the $\text{Ge}(311)$ monochromator. It is equipped with a 2D position-sensitive detector (PSD), which covers the scattering angles of 120° and 7.5° in the horizontal and vertical directions, respectively. An

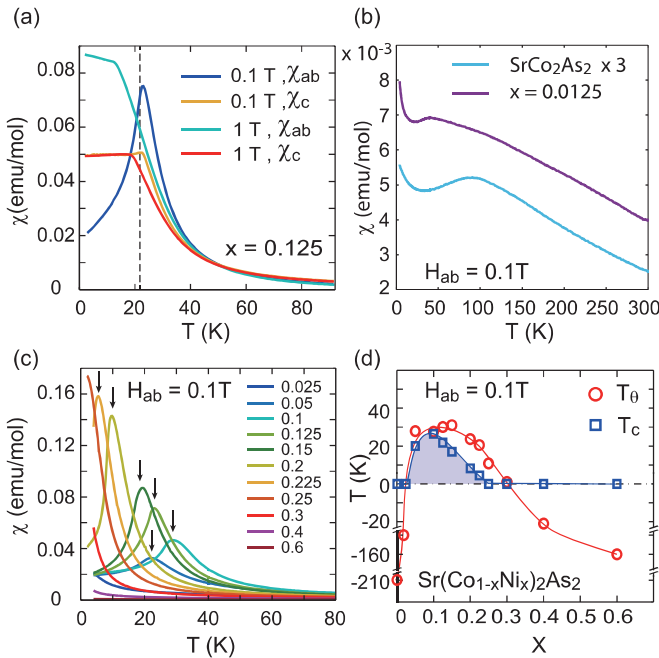


FIG. 2. (a) Magnetic susceptibility χ in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ with field $H = 0.1$ and 1 T parallel and perpendicular to the c axis for $x = 0.125$. (b) and (c) $\chi(T)$ for crystals with $0 \leq x \leq 0.6$. (d) Critical temperature T_c and Weiss temperature T_θ vs x in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$.

oscillating radial collimator is used before the PSD [30]. On both diffractometers the sample was inserted into a closed-cycle refrigerator that provided a temperature range from 4 to 300 K.

For single-crystal diffraction measurements on $\text{Sr}(\text{Co}_{0.9}\text{Ni}_{0.1})_2\text{As}_2$, we define the momentum transfer Q in three-dimensional reciprocal space (in $\text{\AA}^{-1}$) as $Q = H\mathbf{a}^* + K\mathbf{b}^* + L\mathbf{c}^*$, where H , K , and L are Miller indices and $\mathbf{a}^* = 2\pi/\mathbf{a}$, $\mathbf{b}^* = 2\pi/\mathbf{b}$, $\mathbf{c}^* = 2\pi/\mathbf{c}$, with $a = b \approx 3.95 \text{ \AA}$, and $c = 11.62 \text{ \AA}$.

C. Angle-resolved photoemission spectroscopy

ARPES measurements were performed at the Dreamline beamline of the Shanghai Synchrotron Radiation Facility using a Scienta DA80 analyzer and at beamline 13U of the National Synchrotron Radiation Laboratory (Hefei) equipped with a Scienta R4000 analyzer. The energy and angular resolutions were set to 15 meV and 0.2° , respectively. Samples were cleaved in situ, yielding flat, mirrorlike (001) surfaces. During the measurements, the temperature was kept at 20 K, and the pressure was maintained better than 5×10^{-11} Torr.

III. RESULTS

A. Magnetization and magnetic susceptibility data

The magnetic susceptibility $\chi(T)$ measured for $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, with $x = 0.125$, is shown in Fig. 2(a). We find a clear magnetic phase transition at $T_c = 22 \text{ K}$, characterized by a maximum in $d(\chi T)/dT$ [31] for $H = 1 \text{ kOe}$ and H both parallel and perpendicular to the c axis. For $H = 0.1 \text{ T}$, $\chi(T = T_c)$ displays a strong dependence on the direction of H relative to the crystalline lattice, exhibiting a behavior typical of AF order. In contrast to the A-type AF order in CaCo_2As_2 where spins are ordered along the c axis, the anisotropy of $\chi(T)$ for $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ clearly indicates that while the system is isotropic at high temperature, the spins are aligned in the ab plane below T_c [24], exhibiting the typical behavior of an AF magnet with an easy plane. By fitting the high-temperature susceptibility (80–200 K) with the Curie-Weiss law [$\chi(T) \propto (T - T_\theta)^{-1}$], we find a positive T_θ ($\approx 28 \text{ K}$ for H_{ab} and 25 K for H_c) [24], suggesting that the system is dominated by FM interactions.

In classical molecular-field theory [32], the parameter $f \equiv T_\theta/T_c$ is usually found to be smaller than 1, with $-\infty < f < 1$ for antiferromagnets and $f = 1$ for ferromagnets. This is because the real transition (either FM or AF) has to arise before the divergence of $\chi(q)$, which describes the instability of the system towards a particular magnetic state. Therefore, the divergence of Curie-Weiss susceptibility at $T_\theta = T_c$ ($f = 1$) reflects that the system is near a FM instability. However, suppression of the uniform susceptibility [$\chi(q = 0)$] above T_c indicates the influence of low dimensionality since a 2D isotropic Heisenberg system cannot exhibit magnetic order above $T = 0$. This strongly suggests quantum fluctuations may play a significant role. In addition, the large enhancement of $\chi_{ab}(T = T_c)$ with in-plane field at $H = 1 \text{ T}$ suggests the existence of a field-induced metamagnetic transition.

In order to understand the evolution of this magnetic state with x , we systematically explore $\chi(T)$ for $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$

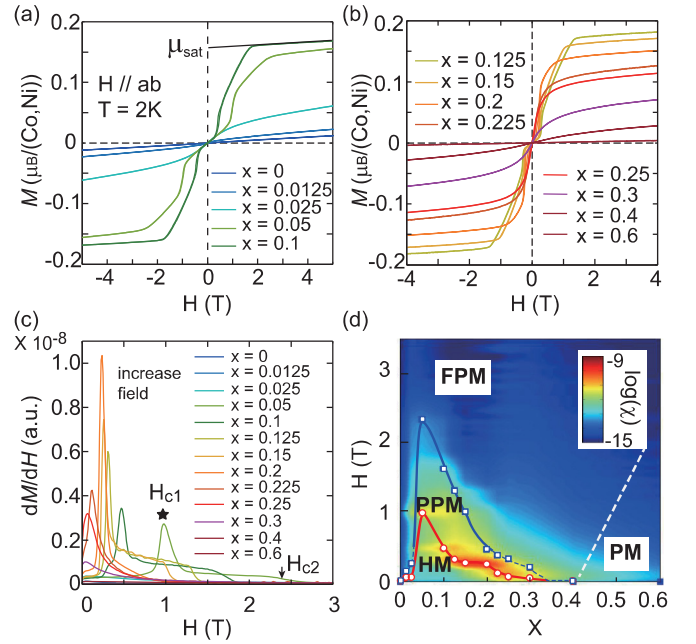


FIG. 3. (a) and (b) $M(H)$ for crystals with $0 \leq x \leq 0.6$. (c) dM/dH vs H for $0 \leq x \leq 0.6$. (d) Magnetic phase diagram with H and x . H_{c1} and H_{c2} are defined in (c). H_{c2} is represented by the blue dashed curve where it is poorly determined by the data. The white dashed line is a speculative separation (or crossover) between paramagnetic (PM) and field-induced FM states. PPM refers to the partially polarized magnetic state, and FPM means the fully polarized magnetic state.

single crystals. In Fig. 2(b), we show magnetic susceptibility of stoichiometric SrCo_2As_2 and slightly Ni doped $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, with $x = 0.0125$. There is a broad maximum arising around 100 K for $x = 0$ and 50 K for $x = 0.0125$, similar to the previous report on SrCo_2As_2 [19] in which the maximum in $\chi(T)$ was attributed to a crossover from a coherent to an incoherent Fermi liquid state with increasing T . However, we note that this broad maximum also exists in other itinerant magnetic materials in the proximity to ferromagnetism, such as CrGe [33]. The gradual suppression with Ni substitution might suggest its close relationship with the novel magnetic phase observed in compounds with larger Ni substitution. With larger x , as shown in Fig. 2(c), we see that the magnetic transition temperature depends on x in a systematic way, creating a dome-like feature in the x - T phase diagram with a maximum T_c at $x \sim 0.1$ [Fig. 2(d)]. In Fig. 2(d), we summarize T_θ by fitting the Curie-Weiss form to $\chi(T)$ for $0 \leq x \leq 0.6$ [24]. Curiously, for Ni concentrations within the dome of magnetic ordering, the dominant interaction, as assessed by T_θ , is ferromagnetic ($T_\theta > 0$) with $f < 1$, while an AF interaction ($T_\theta < 0$) dominates outside this dome. The concurrence of the HM phase transition and the positive T_θ strongly indicates that the magnetic order is driven by FM interactions.

In Figs. 3(a) and 3(b), we applied a magnetic field along the ab plane and measured $M(H)$ for a series of Ni concentrations. Field-induced metamagnetic transitions associated with the reorientation of magnetic moments are clearly observed in the magnetically ordered compounds. In Fig. 3(a), for example, at

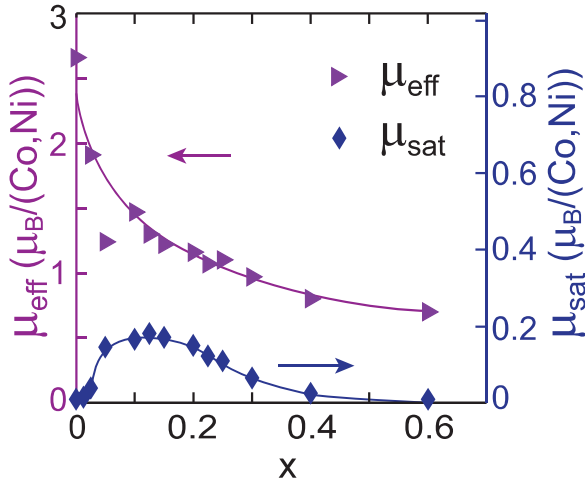


FIG. 4. Effective magnetic moment μ_{eff} by fits of the C-W law to the magnetic susceptibility above T_c and saturated magnetic moment μ_{sat} vs x . The large ratio of $\mu_{\text{eff}}/\mu_{\text{sat}}$ reflects the itinerant nature of magnetism.

$x = 0.1$, the system is in a helical magnetic ground state at low temperatures. The magnetization displays a linear dependence on the in-plane magnetic field at low field [Fig. 3(c)]. As the field increases, a PPM intermediate state arises at $H_{c1} \sim 0.5$ T. This transition is clearly observed in $\chi(H) \equiv dM/dH$ as a sharp peak which defines the critical field as H_{c1} [Fig. 3(c)]. The magnetization around H_{c1} exhibits hysteresis behavior [24] with magnetic field, implying a first-order transition. In the intermediate state, the susceptibility displays a plateau between H_{c1} and H_{c2} beyond which the system evolves into a FPM state. Furthermore, as the Ni concentration changes, the plateau in $\chi(H)$ shrinks and disappears for $x = 0.05$ or $x \geq 0.2$, such that the determination of H_{c2} in these regions is somewhat ambiguous. Therefore, we employ a blue dashed line in the x - H phase diagram to represent the poorly defined H_{c2} [Fig. 3(d)]. We have included the magnetic susceptibility ($\chi \equiv dM/dH$) as a color plot (log scale) in Fig. 3(d) to emphasize the clear delineation of magnetic phases.

In Fig. 4, the effective magnetic moment μ_{eff} , estimated from the Curie constant, and the saturated magnetic moment μ_{sat} , according to the high-field magnetization, are presented. The magnitude of the effective moment is similar to that for an $S = 1/2$ system, while μ_{sat} is about one order of magnitude smaller for $x = 0.05$. This renders a large Rhodes-Wholfarth ratio $\mu_{\text{eff}}/\mu_{\text{sat}}$, suggesting that the magnetism in this system is itinerant in character.

In Figs. 5(a) and 5(b), we investigate the temperature dependence of $M(H)$ and $\chi(H)$ for $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, with $x = 0.1$. The plateau in $\chi(H)$ was observed at all temperatures below T_c [Fig. 5(b)]. We note that even in the paramagnetic state, the magnitude at $T = 30$ K is dramatically enhanced as the system approaches the magnetic transition temperature, in comparison with that at a higher temperature of $T = 100$ K. Based on these measurements, a H - T phase diagram for $x = 0.1$ is presented in Fig. 5(c), and a schematic three-dimensional (H - T - x) phase diagram is shown in Fig. 5(d). In Fig. 6, we plot the magnetization of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ over a range of x with field parallel to the c axis. The magnetization

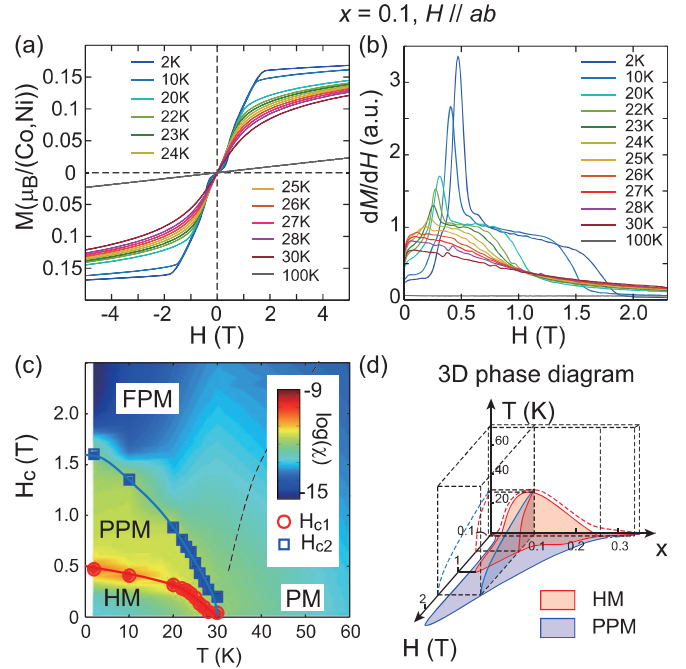


FIG. 5. (a) and (b) M_{ab} and dM_{ab}/dH vs H at temperatures identified in the plots. (c) Phase diagram determined from data in (a) and (b). The dashed line is a speculative crossover separating the fully polarized magnetic (FPM) and paramagnetic (PM) phases. PPM denotes the partially polarized magnetic phase. (d) Three-dimensional phase diagram combining Figs. 2(d), 3(d), and 7(c).

at low fields has a linear dependence on field before saturating at higher fields. In summary, our magnetization and susceptibility measurements for $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ are consistent with the magnetic behavior of a system with helical spin order. However, further neutron diffraction measurements are required to determine the magnetic structure of the ground state.

B. Neutron diffraction

In order to understand the changes that occur with x in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, we probed the magnetic structure at zero field and $T = 4$ K for $x = 0.1$ via neutron diffraction

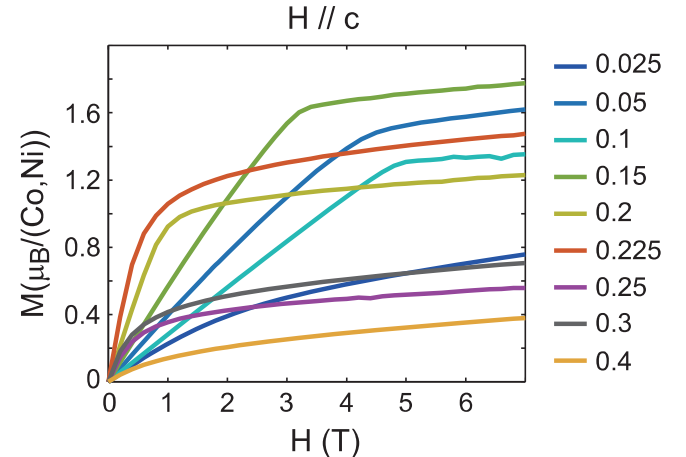


FIG. 6. Magnetization M vs magnetic field H with H parallel to the c axis at a variety of Ni concentrations.

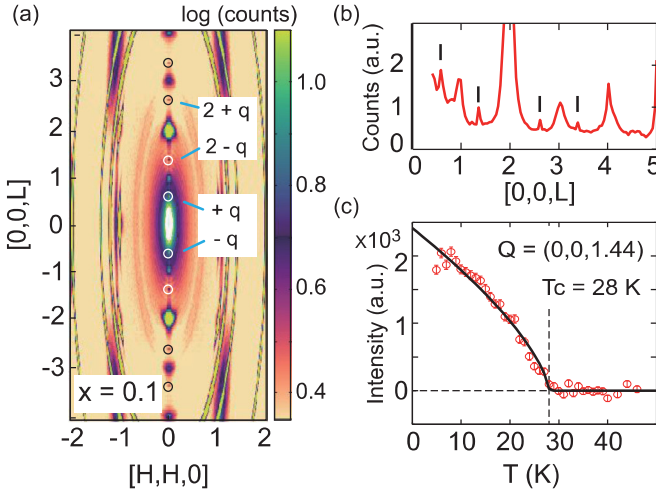


FIG. 7. (a) Neutron diffraction measurements on $\text{Sr}(\text{Co}_{0.9}\text{Ni}_{0.1})_2\text{As}_2$ in the $[\text{H}, \text{H}, \text{L}]$ plane [24]. (b) One-dimensional cut along the $[0, 0, \text{L}]$ direction in (a). The intensities at $(0, 0, \text{L})$ with L being an odd number are due to double scattering [34]. (c) Temperature dependence of the magnetic peak at $Q = (0, 0, 1.44)$.

experiments. We chose to explore a $\text{Sr}(\text{Co}_{0.9}\text{Ni}_{0.1})_2\text{As}_2$ single crystal as these samples displayed the highest T_c with one of the largest sat for this series. Figure 7(a) displays the neutron diffraction pattern in the $[\text{H}, \text{H}, \text{L}]$ scattering plane. Incommensurate peaks at $(0, 0, n - q)$, with n being an even integer and $q = (0, 0, 0.56)$, are clearly observed in Fig 7(b), manifesting either a planar spiral or sinusoidal magnetic structure. However, a sinusoidal magnetic structure is unlikely in this material since it requires a uniaxial spin anisotropy along either the a or b axis and therefore breaks the tetragonal symmetry of the underlying crystalline lattice [24]. This is substantially different from the collinear AF order in iron pnictides [13] or the commensurate A-type AF order in CaCo_2As_2 [35]. In EuCo_2As_2 [36], a similar helical order was observed. However, the magnetic moments are localized on Eu sites with $S = 7/2$, in contrast to the itinerant magnetism in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, implying the underlying physics is significantly different. In addition, the dissimilar spin anisotropy between CaCo_2As_2 [37] and $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ suggests a change from an easy-axis to an easy-plane spin anisotropy [38]. Furthermore, the value of q is close but not equal to $(0, 0, 0.5)$, indicating a short pitch for the helix, about four Co layers or two unit cells, resulting in spins on adjacent Co-layers being nearly perpendicular. This severely constrains the likely exchange interaction responsible for the magnetism. Figure 7(c) demonstrates the T dependence of the scattering intensity at $Q = (0, 0, 1.44)$. The T_c of 28 K agrees well with that determined from $\chi(T)$ in Fig. 2(c). A longitudinal scan along the $[0, 0, \text{L}]$ direction at 4 K (Fig. S4 in Ref. [24]) indicates that the magnetic ordering is long range within the current instrumental resolution.

In Fig. 8, we show the magnetic intensity from neutron diffraction measurements in comparison with the calculated structure factor for helical order. We employed a $d_{x^2-y^2}$ magnetic form factor for Co^{2+} ions [24,39], and the results (solid lines) are in good agreement with our neutron diffraction

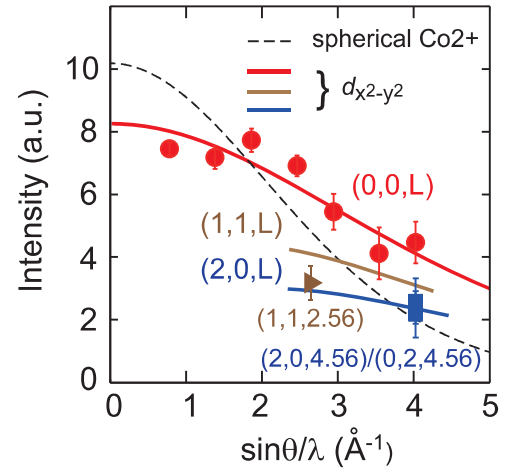


FIG. 8. Integrated intensity of magnetic Bragg peaks. The dashed curve is the magnetic form factor of Co^{2+} using a spherical approximation. The solid curves are magnetic form factors for the $d_{x^2-y^2}$ orbital multiplied by a polarization factor along different measured directions for a helical order. Red, brown, and blue are along the $[0, 0, \text{L}]$, $[1, 1, \text{L}]$, and $[2, 0, \text{L}]/[0, 2, \text{L}]$ directions, respectively.

measurements. This agreement suggests that the magnetic moments in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ are closely associated with electrons having $d_{x^2-y^2}$ orbital character. The spherical Co^{2+} magnetic form factor (dashed curve) is plotted for comparison and is a much poorer representation of our data.

C. Electronic band structure

In nominally pure SrCo_2As_2 , a flat band exists at an energy just above the Fermi level near the M point of the Brillouin zone [17,20]. Substitution of Ni is expected to raise E_F so that this flat band is partially occupied, resulting in E_F being closer to a Van Hove singularity. To establish the connection between the helical magnetic phase and this flat band, we carried out ARPES experiments on a series of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ samples ($x = 0, 0.1, 0.2, 0.4$, and 0.6) to map out their electronic structures. We show in Fig. 9 the Fermi surfaces of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, with $x = 0.1, 0.2, 0.4, 0.6$, as measured by ARPES. A diamond-shaped electronlike Fermi surface (dashed curve) is observed at the Γ point in the sample with $x = 0.1$, similar to that observed in nominally pure $\text{Sr}(\text{Ba})\text{Co}_2\text{As}_2$ [17,21]. This band is associated with a $d_{x^2-y^2}$ orbital and has flat dispersion. It is likely responsible for in-plane ferromagnetism and the observed helimagnetic order. We note that as the Ni concentration increases, this Fermi surface opens along the Γ -M direction and is further separated at $x = 0.6$, consistent with the fact that the flat band shifts down below the Fermi level.

In Fig. 10, we display the cuts of ARPES spectra along the Γ -M and Z-A directions for these samples. While there are barely changes in the $x = 0.1, 0.2$, and 0.4 samples, the whole band structure suddenly shifts down for $x = 0.6$. This shows that the flat band along the Γ -M direction sinks below E_F and thus is only observed in ARPES spectra of Fig. 10(d). In order to further prove the existence of the flat band, we made two cuts along both the horizontal and vertical directions, shown as red lines in Fig. 9(d) (for $x = 0.6$), and display these in

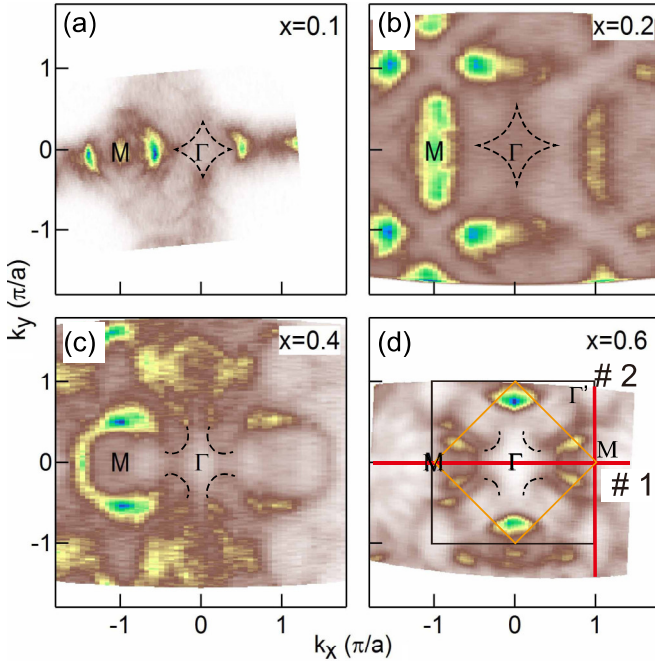


FIG. 9. Plot of the ARPES intensity at E_F of $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ with $x = 0.1, 0.2, 0.4$, and 0.6 . A complex structure of Fermi surfaces is observed, and the topology changes dramatically from $x = 0.1$ to $x = 0.6$ due to the existence of flat-band dispersion. The dashed lines are a guide to the eye.

Fig. 11. Figures 11(a) and 11(c) are the intensity plots of band dispersion along the cut 1 and cut 2 directions. A nearly flat band dispersion around 0.4 eV with an electronlike band minimum at the Γ point is realized in both cuts. These flat-band spectra typically have low intensity in ARPES [40] but

are clearly resolved in the second-derivative plot [Fig. 11(b)] and the energy distribution curves (EDCs) [Fig. 11(d)].

To investigate the x dependence of this flat band, we show the EDC cuts at the M point for samples with $0 \leq x \leq 0.6$ in Fig. 12(a). One peak close to E_F associated with the flat band is identified by the arrows and is degenerate with another electronic band at the M point [17,20]. The Fermi-Dirac function cuts off ARPES spectra at E_F , affecting the shape of these peaks. To avoid this effect, we use the symmetrized EDCs to determine whether the band is located below or above the Fermi energy. Similar methods were applied to determine small superconducting gaps in iron-based superconductors [41] [Fig. 12(b)]. A central peak at E_F manifests that the band is either at or above the Fermi energy and thus gives the lower limit of the band location. We use the value estimated from dynamic mean-field theory (DMFT) calculations on pure SrCo_2As_2 in previous work [17,20] as the upper limit, and the results are summarized in Fig. 12(c). The shaded area roughly establishes the evolution of the energy location of the flat band with Ni substitution. It is clear that the band shifts with x at $x = 0.4$ much more slowly than at larger x . This is a manifestation of the large density of states (DOS) of the flat band for electrons filling as it shifts across E_F ($0 \leq x \leq 0.4$). This flat band has $d_{x^2-y^2}$ orbital character and makes a significant contribution to the DOS near the Fermi energy [17]. Our ARPES results for $0 \leq x \leq 0.6$ confirm that this flat band is partially occupied for the range of x with helimagnetism and suggest that the helical magnetic phase is closely associated with this band.

IV. DISCUSSION AND CONCLUSION

In ACo_2As_2 systems, density functional theory plus DMFT calculations suggest that FM instabilities are due to a flat band

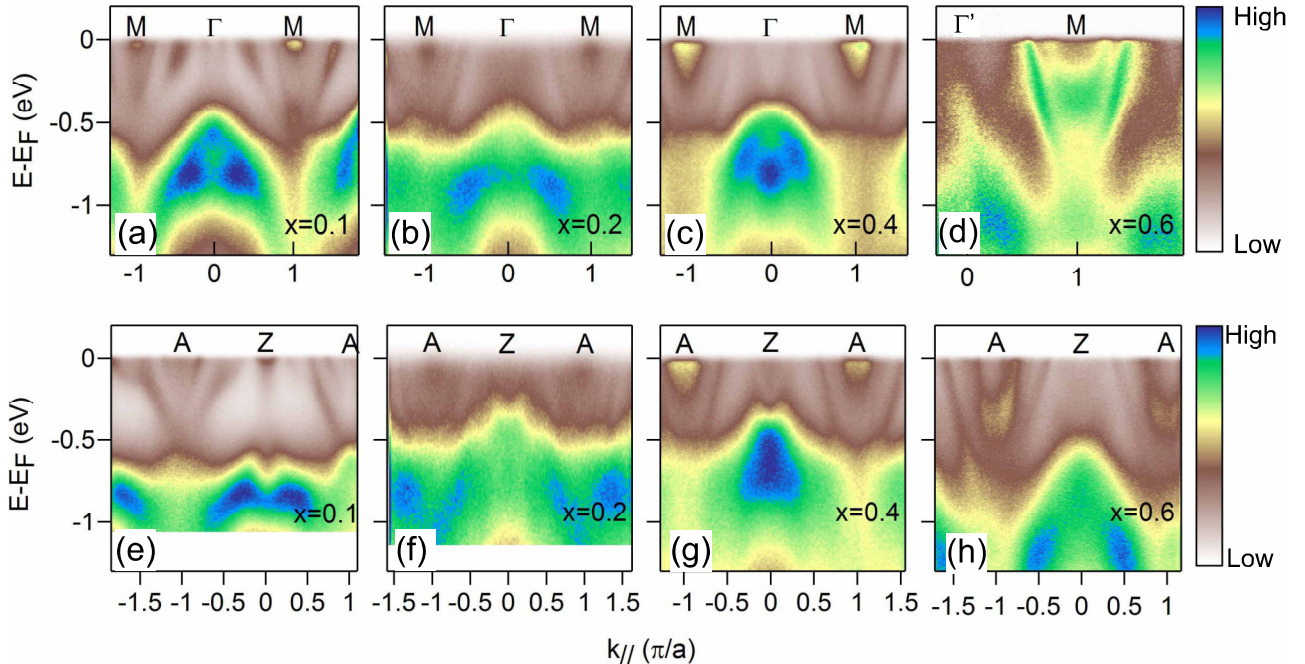


FIG. 10. Intensity plots of the band dispersion along (a)–(d) the Γ -M and (e)–(h) Z-A directions from ARPES measurements for $x = 0.1, 0.2, 0.4$, and 0.6 .

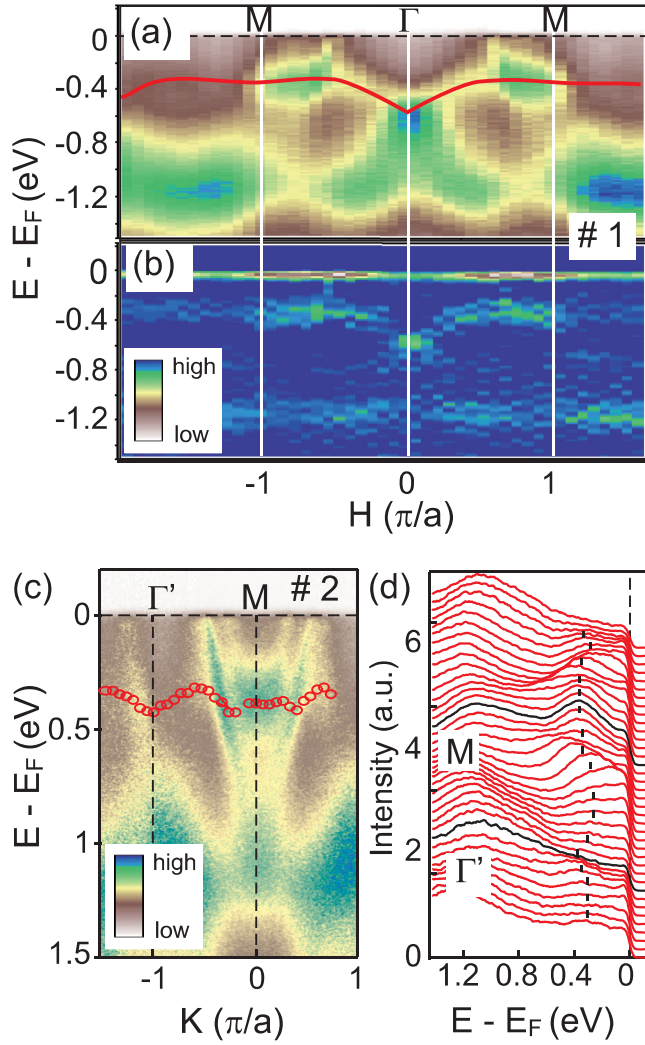


FIG. 11. (a) and (b) ARPES intensity plot and its second derivative along the cut 1 direction as shown in Fig. 9(d). (c) and (d) ARPES spectra along the cut 2 direction and its corresponding energy distribution curve.

residing near E_F . However, experiments show an evolution from AF order in CaCo_2As_2 [18,35] to paramagnetism in BaCo_2As_2 [22] (or $\text{Ba}(\text{Fe}_{1/3}\text{Co}_{1/3}\text{Ni}_{1/3})_2\text{As}_2$ [26]) with evidence of a nearby FM critical point [22,26]. This indicates that the magnetism can be tuned by chemical substitution. In SrCo_2As_2 , previous inelastic neutron scattering experiments [17] revealed the coexistence of FM and AF spin fluctuations. It is likely that the fine tuning of these fluctuations by chemical substitution induces the observed magnetic order in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$.

In many helical magnetic systems such as MnSi [42–44] and $\text{Cr}_{1/3}\text{NbSe}_2$ [45,46], the helical magnetic order is induced by the Dzyaloshinskii-Moriya (DM) interaction [47] arising from the lack of crystalline inversion symmetry. Since $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ has a centrosymmetric crystal structure, the DM interaction is absent. In isostructural EuCo_2As_2 , where a large magnetic moment exists on the Eu site [32], the mechanism for helimagnetism is likely distinct from that in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$. Here, a classical Heisenberg model with a frustrated nearest-neighboring (NN) J_1 and next-nearest-

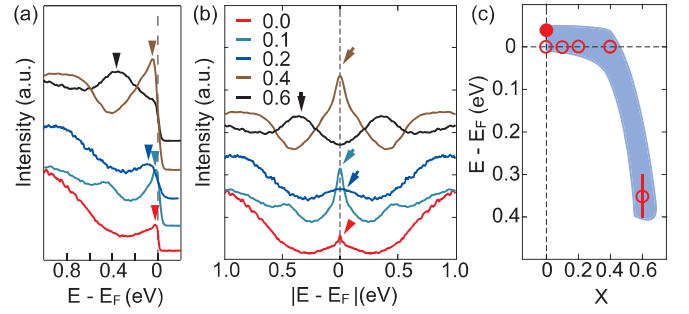


FIG. 12. (a) Energy distribution curves (EDCs) at the M point for $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ with $x = 0, 0.1, 0.2, 0.4, 0.6$. Small peaks close to E_F were identified as residual spectra from the band above the Fermi level truncated by the Fermi-Dirac function. (b) Symmetrized EDCs. The peaks at E_F in symmetrized EDCs set the lower limit of the energy location of the flat band within the instrumental resolution. (c) Energy location of the flat band determined by ARPES and that from the DMFT calculation [20] for $x = 0$ which we used as the upper limit. The shaded area is a guide to the eye. The evolution of the energy of the flat band with respect to the Fermi level is demonstrated.

neighboring J_2 between basal planes can produce helical order [48,49]. Such a mechanism of helical order for centrosymmetric materials was proposed many years ago by incorporating exchange frustration [50]. However, there are difficulties in applying this model to $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$. First, the effective spin S derived from $\chi_{\text{eff}} = g^2 S(S+1)$ is only about $1/2$, making the validity of the classical Heisenberg model ($S \rightarrow \infty$) questionable. Second, according to the theory, the helical arrangement has a q given by $\cos(cq) = J_1/4J_2$, in which c is the interlayer distance. Therefore, the nearly 90° angle between magnetic moments in adjacent layers with $q \sim \frac{\pi}{2c}$ leads to $\cos cq \sim \cos(\pi/2) = 0$ and thus an almost vanishing NN exchange coupling J_1 which appears to be unrealistic. Higher-order exchange couplings such as J_3 may help stabilize the helical state, but the phase space [51] for a stable helix with a rotation of an $\sim \pi/2$ angle between adjacent spins is still limited. Third, the Ruderman-Kittel-Kasuya-Yosida (RKKY) interaction may be responsible for the evolution of magnetism from CaCo_2As_2 to BaCo_2As_2 since the different sizes of intercalated ions can tune the interlayer Co-Co distance and thus the magnitude of exchange interactions. However, the origin of the helimagnetism along the c axis is still ambiguous because an unrealistically large k_F ($\sim \frac{\pi}{c}$) is required to make J_1 nearly zero via the spatial oscillation of the RKKY interaction.

On the other hand, quantum fluctuations in a paramagnet near a FM instability can make significant contributions to the free energy and stabilize unusual magnetic order [2,5,6]. In $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$, the $d_{x^2-y^2}$ band is flat along both the in-plane Γ -M and out-of-plane Γ -Z directions. The massive quantum particle-hole excitations can strongly affect the flat-band dispersion at the mean-field level, and as a consequence, a helical magnetic order with certain momenta can be induced. However, the presence of a flat band requires physics beyond the conventional itinerant spin density wave theory with Fermi surface nesting but may be a realization of the quantum order-by-disorder mechanism.

In conclusion, we discovered a helical magnetic order in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ and established the phase diagram with systematic studies of the magnetic susceptibility. By combining neutron diffraction and ARPES experiments, we have established a close connection between the $d_{x^2-y^2}$ flat band and the helical magnetic order. Based on these results, we suggest that the helical magnetic order in $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$ is driven by quantum fluctuations and the system is likely in the vicinity of at least one quantum critical point.

Note added. Recently, we learned of papers studying [52,53] the same material as the present work. The main experimental results of these studies are consistent with our results.

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