PROCEEDINGS OF SCIENCE

LHC and DIS experimental data in the CT18(Z) global QCD analysis

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We discuss implementation of the LHC experimental data sets in the new CT18 global analysis of quantum chromodynamics (QCD) at the next-to-next-leading order of the QCD coupling strength. New methodological developments in the fitting methodology are discussed. Behavior of the CT18 NNLO PDFs for the conventional and "saturation-inspired" factorization scales in deep-inelastic scattering is reviewed. Four new families of (N)NLO CTEQ-TEA PDFs are presented: CT18, A, X, and Z.

MSUHEP-19-019, SMU-HEP-19-13

XXVII International Workshop on Deep-Inelastic Scattering and Related Subjects - DIS2019
8-12 April, 2019
Torino, Italy

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PoS(DIS2019)021

In the companion contribution [1], we presented the new CT18 global QCD analysis of parton distribution functions (PDFs). The CT18 analysis updates widely used CT14 PDF sets [2] by applying NNLO and NLO global fits to an expanded set of experimental measurements that include high-luminosity data from the ep collider HERA and the Large Hadron Collider. The CT18 experimental data set includes high-statistics measurements from ATLAS, CMS, and LHCb on production of inclusive jets, W/Z bosons, and top quark pairs, while it retaining crucial *legacy* data, such as measurements from the Tevatron and the HERA Run I and Run II combined data. In this contribution, we review implementation of the new data sets in the CT18 global fit and the associated physics issues that affect the resulting PDFs and a wide class of QCD predictions based on them.

By 2018, the LHC collaborations published about three dozen experimental data sets that can potentially constrain the CT PDFs. In light of the unprecedented precision reached in some measurements, the latest LHC data must be analyzed using next-to-next-to-leading order (NNLO) theoretical predictions in perturbative QCD. The final PDFs depend on numerous systematic factors in the experimental data; and the scope of numerical computations needs to be expanded, too. A systematic examination of these effects is essential for trustworthy estimates of PDF uncertainties.

Combined HERA I+II DIS data and an x-dependent factorization scale. Even in the LHC era, the DIS data from ep collider HERA provides the dominant constraints on the CT18 PDFs. This dominance can be established using the `ePump` and `PDFSense` statistical techniques reviewed below. CT18 implements the final ("combined") data set from DIS at HERA-I and II [3] that supersedes the HERA-I only data set [4] used in CT14 [2]. A transitional PDF set, CT14HERA2, was released based on fitting the final HERA data [5]. We found fair overall agreement of the HERA I+II data with both CT14 and CT14HERA2 PDFs, and that both PDF ensembles describe equally well the non-HERA data included in our global analysis. At the same time, we observed some disagreement ("statistical tension") between the e^+p and e^-p DIS cross sections of the HERA I+II data set. We determined that, at the moment, no plausible explanation conclusively explains the full pattern of these tensions, as they are distributed across the whole accessible range of Bjorken x and lepton-proton momentum transfer Q at HERA.

It has been argued that resummation of logarithms $\ln^p(1/x)$ at $x \ll 1$ improves agreement withe HERA Run I+II data by several tens of units of χ^2 [6, 7]. In our analysis, we observe that, by evaluating the DIS cross sections at NNLO with an x-dependent factorization scale, such as a tuned scale $\mu^2_{F,x} = 0.8^2 \left(Q^2 + 0.3 \text{ GeV}^2/x^{0.3}\right)$, instead of the conventional choice $\mu^2_F = Q^2$, we achieve a comparable quality of improvement in the description of the HERA DIS data set by the *fixed-order* NNLO theoretical prediction as the inclusion of the low-x resummation in [6, 7]. Namely, the χ^2(HERAI+II) reduces by > 50 units in the kinematical region $Q > 2$ GeV, $x > 10^{-5}$ of the DIS data included in the CT18 global fit. The parametric form of the x-dependent scale $\mu^2_{F,x}$ is inspired by qualitative saturation arguments (see, e.g., [8]), and the numerical coefficients in $\mu^2_{F,x}$ are chosen to minimize χ^2 for the HERA DIS data.

Fig. 1(left) shows the changes in the candidate CT18 PDFs obtained by fitting DIS with the x-dependent factorization scale, as compared to the CT18 PDFs with the nominal scale. With the scale $\mu^2_{F,x}$, we observe reduced u and d (anti-)quark PDFs and increased gluon and strangeness PDFs at $x < 10^{-2}$ as compared to the nominal CT18 fit, with some compensating changes occuring in the same PDFs in the unconstrained region $x > 0.5$ in order to satisfy the valence and momentum

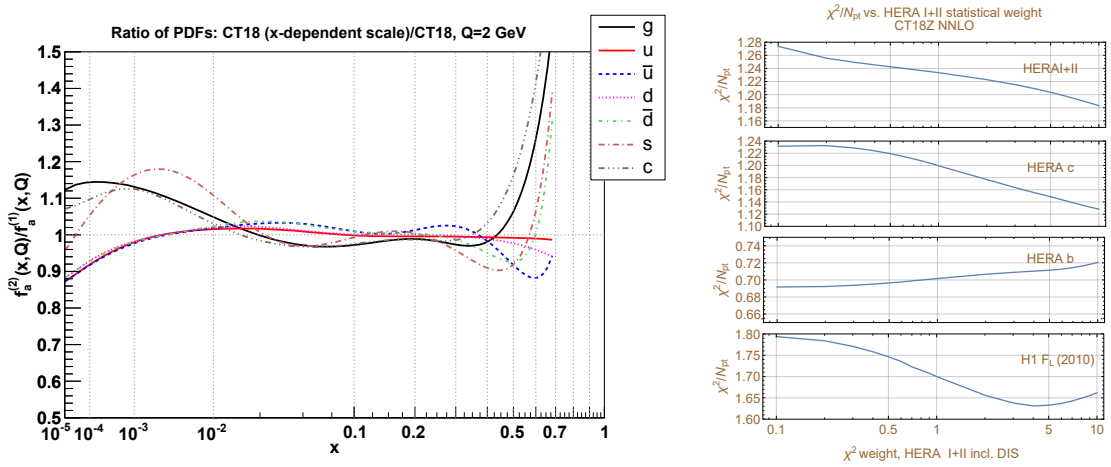

Figure 1: Left: The ratios of the candidate CT18 NNLO PDFs obtained with the x-dependent and standard factorization scales in DIS data sets. Right: The χ^2/N_{pt} values for four HERA data sets in the CT18Z fits with the x-dependent DIS factorization scale and varied statistical weight of the HERA I +II inclusive DIS data set.

sum rules. The right Fig. 1 shows the χ^2/N_{pt} values (divided by the number N_{pt} of experimental data points) for four HERA data sets (inclusive NC+CC DIS [3], reduced charm, bottom production cross sections, and H1 longitudinal function $F_L(x,Q^2)$ [9]) in the fits with the varied statistical weight w of the HERA I+II inclusive DIS data set [3]. The default CT18 fits correspond to $w=1$; with $w=10$, the CT18 fit increasingly behaves as a HERA-only fit. We see that, with the scale $\mu^2_{F,x}$ and $w=10$, χ^2/N_{pt} for the inclusive DIS data set improves almost to the levels observed in the "resummed" HERA-only fits without intrinsic charm [6, 7]. The quality of the fit to the charm SIDIS cross section and H1 F_L also improves.

Selection of new LHC experiments. When selecting the most promising LHC experiments for the CT18 fit, we had to address a recurrent challenge, the presence of statistical tensions among various (sub)sets of the latest experimental data from HERA, LHC, and the Tevatron. The quickly improving precision of the collider data reveals previously irrelevant anomalies either in the experiment or theory. These anomalies are revealed by applying strong goodness-of-fit tests [10]. Figure 2 illustrates the degree of tensions using a representation based on the effective Gaussian variables $S_E \equiv \sqrt{2\chi^2_E} - \sqrt{2N_E - 1}$ [11] constructed from the χ^2 values and numbers of data points for individual data sets E. In a high-quality fit, the probability distribution for S_E must be approximately a standard normal distribution (with a unit half-width). In CTEQ-TEA and global fits from either CTEQ or other groups, we in fact observe wider S_E distributions, cf. Fig. 2, with some most comprehensive and precise data sets (notably, HERA I+II inclusive DIS [3] and ATLAS 7 TeV Z/W production [12]) having S_E values as high as five units or more. The question, then, is how to select the clean and accurate experiments for the global analysis from the list that grows day-by-day, while maximally preserving the consistency of the selected experiments.

For example, there are many LHC experimental data sets [13] that are potentially sensitive to the PDFs, including novel measurements in production of high-p_T Z bosons, $t\bar{t}$ pairs, heavy quarks, and $W+c$ pairs. Including all such candidate experiments into the full global fit is impractical: CPU

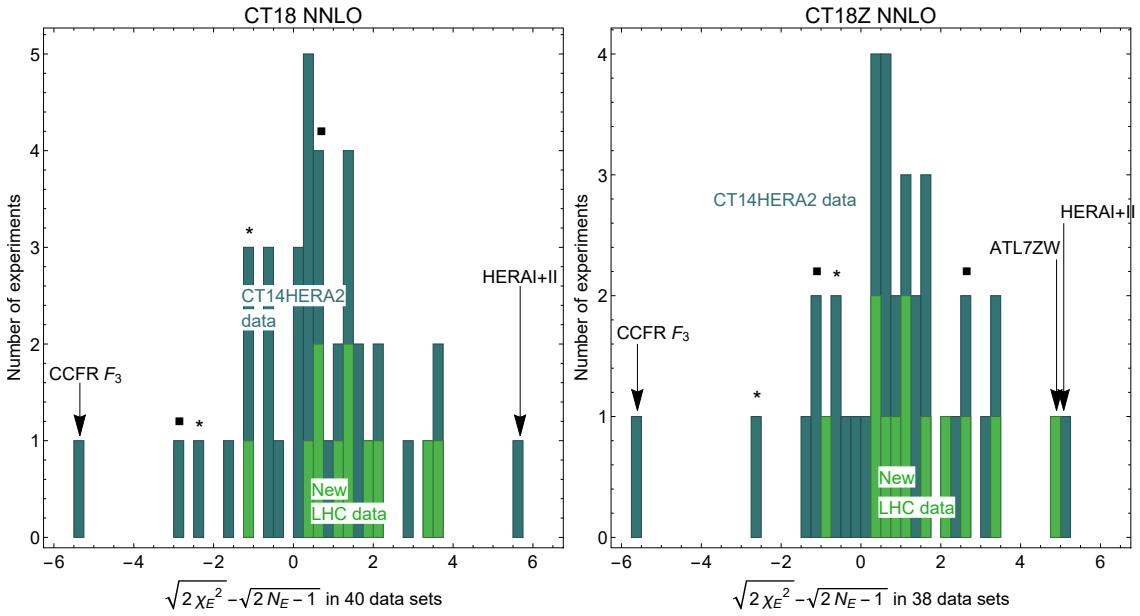

Figure 2: The effective Gaussian variable (S_n) distribution of all (a) CT18 data sets, and (b) CT18Z data sets. Two squares and two stars indicate the S_n values for the NuTeV dimuon and CCFR dimuon data, respectively.

costs grow quickly with the number of experimental data sets at NNLO. Poorly fitted experiments would increase, not decrease, the final PDF uncertainty. The generation of one error PDF set took several days of CPU time in the CT14 fit to 33 experiments in a single-thread mode. Adding 20-30 additional experiments with this setup was thus impossible.

Advancements in fitting methodology. The CTEQ-TEA group resolved these challenges through a multi-prone effort. We developed two programs for fast preliminary analysis to identify the eligible experimental data sets for the global fit. The `PDFSense` program [14] was developed at SMU to predict quantitatively, and before doing the fit, which data sets will have an impact on the global PDF fit. The `ePump` program [15] developed at MSU applies PDF reweighting to quickly estimate the impact of data on the PDFs prior to the global fit. These programs provide helpful guidelines for the selection of the most valuable experiments based entirely on the previously published Hessian error PDFs.

The CTEQ fitting code was parallelized to allow faster turnaround time (one fit within few hours instead of many days) on high-performance computing clusters. For as much relevant LHC data as possible, we computed our own tables for APPLGrid/fastNLO fast interfaces [16, 17] for NLO cross sections (to be multiplied by tabulated point-by-point NNLO/NLO corrections) for various new LHC processes: production of high-p_T Z bosons, jets, $t\bar{t}$ pairs. The APPLgrid tables were cross validated against similar tables from other groups (available in the public domain) and optimized for speed and accuracy.

The resulting family of new PDFs consists of four NNLO PDF ensembles, and the corresponding NLO ones: the default CT18 ensemble and three alternative ensemble, designated as CT18A, X, and Z. Based on the `PDFSense` and `ePump` studies, eleven new LHC data sets have

been included in all four PDF fits, notably, data at 7 and 8 TeV on lepton pair, jet, and $t\bar{t}$ production. Significant effort was spent on understanding the sources of PDF uncertainties. Theoretical uncertainties associated with the scale choice were investigated for the affected processes such as DIS and high-p_T Z production. Other considered theoretical uncertainties were due to the differences among the NNLO/resummation codes (FEWZ, ResBos, MCFM, NNLOJet++,...) and Monte-Carlo integration. The important parametrization uncertainty was investigated by repeating the fits for 90+ trial functional forms of the PDFs. [Our post-CT10 fits parametrize PDFs using Bernstein polynomials, which simplify trying a wide range of parametrization forms to quantify/eliminate potential biases.] In addition to the default CT18 PDF ensemble, the other three sets were obtained under alternative assumptions. (a) The CT18A and CT18Z analyses include high-luminosity ATLAS 7 TeV W/Z rapidity distributions [12] that show some tension with DIS experiments and prefer a larger strangeness PDF than the DIS experiments. Inclusion of the ATLAS 7 TeV W/Z data leads to worse χ^2_E values (higher S_E values) for dimuon SIDIS production data sensitive to the strangeness PDF. This can be seen in the comparison of S_E distributions in Fig. 2, where the S_E values for CCFR and NuTeV dimuon data sets are elevated in the CT18Z fit on the right, as compared to the CT18 fit on the left, as a consequence of inclusion of the ATLAS W/Z data in the CT18Z fit. (b) The CT18X and CT18Z fits use an x-dependent factorization scale in NNLO DIS cross sections to mimic enhanced higher-order logarithms at small Bjorken x and small Q. This choice results in the enhanced gluon PDF at small x and reduced gluon at $x \sim 0.01$, as discussed above. Furthermore, the CDHSW data for DIS on heavy nuclei prefer a somewhat harder gluon PDF at $x > 0.1$ than other data sets. In the CT18Z fit, we have removed the CDHSW data. The combination of these choices in the CT18Z results in the NNLO Higgs production cross section via gluon fusion that is reduced by about 1% compared to the corresponding CT14 and CT18 predictions. Thus, the various choices made during the generation of four CT18(A,X,Z) data sets allow us to more faithfully explore the full range of the PDF behavior at NNLO that is consistent with the available hadronic data, with implications for electroweak precision physics measurements and new physics searches at the LHC.

Acknowledgments

The work of J. Gao was sponsored by the National Natural Science Foundation of China under the Grant No. 11875189 and No.11835005. The work at SMU is supported by the U.S. Department of Energy under Grant No. DE-SC0010129. T. J. Hobbs acknowledges support from an EIC Center Fellowship. The work of M. Guzzi is supported by the National Science Foundation under Grant No. PHY1820818. The work of C.-P. Yuan was supported by the U.S. National Science Foundation under Grant No. PHY-1719914, and he is also grateful for the support from the Wu-Ki Tung endowed chair in particle physics.

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