

# Au-MoS<sub>2</sub> Hybrids as Hydrogen Evolution Electrocatalysts

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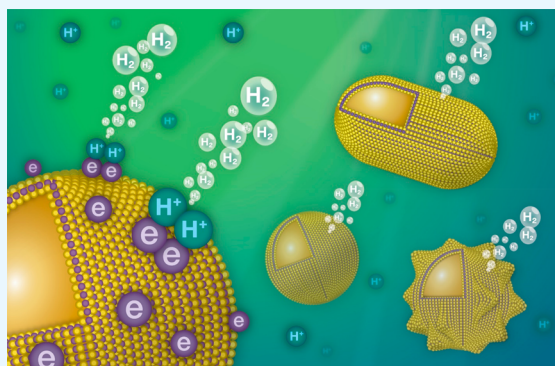
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## Supporting Information

**ABSTRACT:** Core-shell nanoparticles provide a unique morphology to exploit electronic interactions between dissimilar materials, conferring upon them new or improved functionalities. MoS<sub>2</sub> is a layered transition-metal disulfide that has been studied extensively for the hydrogen evolution reaction (HER) but still suffers from low electrocatalytic activity due to its poor electronic conductivity. To understand the fundamental aspects of the MoS<sub>2</sub>–Au hybrids with regard to their electrocatalytic activity, a single to a few layers of MoS<sub>2</sub> were deposited over Au nanoparticles via a versatile procedure that allows for complete encapsulation of Au nanoparticles of arbitrary geometries. High-resolution transmission electron microscopy of the Au@MoS<sub>2</sub> nanoparticles provides direct evidence for the core-shell morphology and also reveals the presence of morphological defects and irregularities in the MoS<sub>2</sub> shell that are known to be more active for HER than the pristine MoS<sub>2</sub> basal plane. Electrochemical measurements show a significant improvement in the HER activity of Au@MoS<sub>2</sub> nanoparticles relative to freestanding MoS<sub>2</sub> or Au-decorated MoS<sub>2</sub>. The best electrochemical performance was demonstrated by the Au nanostars—the largest Au core employed here—encapsulated in a MoS<sub>2</sub> shell. Density-functional theory calculations show that charge transfer occurs from the Au to the MoS<sub>2</sub> layers, producing a more conductive catalyst layer and a better electrode for electrochemical HER. The strategies to further improve the catalytic properties of such hybrid nanoparticles are discussed.

**KEYWORDS:** core-shell, density-functional theory, charge transfer, nanostars, nanorods, electrocatalysis, transition-metal dichalcogenides



## 1. INTRODUCTION

Layered transition-metal dichalcogenides (TMDs), in particular of Group VIB TMDs such as MoS<sub>2</sub>, WS<sub>2</sub>, and MoSe<sub>2</sub>, used traditionally in petrochemistry, have begun to attract interest for key sustainable-energy reactions such as hydrogen evolution reaction (HER) and CO<sub>2</sub> reduction.<sup>1–5</sup> In their ground state (2H phase), Mo and W TMDs are semiconductors with thickness-dependent band-gaps in the range of 1–2 eV<sup>6</sup> and, consequently, are not particularly effective as electrocatalysts. Another limitation of these materials for electrocatalysis, as shown by both calculations and experiments, is that the active sites for HER are located along the (limited) exposed edges while the basal planes are relatively

inert. It is against this background that efforts are currently being dedicated to further modify the MoS<sub>2</sub> structures with the aim of improving performance by increasing the density of their catalytic sites by creating defects,<sup>7–9</sup> modifying the morphology, or by doping with other transition metals.<sup>2,3,10,11</sup> Another approach is to combine MoS<sub>2</sub> with other materials and produce hybrids with improved electrical conductivity,<sup>12</sup> more active basal planes,<sup>13</sup> or improved photocatalytic activity.<sup>14,15</sup> Recently, a study exploring the thickness depend-

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ence (number of layers) of MoS<sub>2</sub> electrocatalysts showed a significant improvement in the catalytic activity for HER with decreasing number of layers.<sup>16</sup> The layer-dependent electrocatalysis was correlated to interlayer hopping of electrons between layers of MoS<sub>2</sub>. Therefore, increasing the hopping efficiency is a crucial factor for the design of MoS<sub>2</sub> with optimal catalytic activity. Although many materials have been tested and researched as heterogeneous catalysts, the use of hybrid materials, such as core–shell structures, is a noteworthy approach for tuning the catalytic properties of layered nanomaterials. In particular, increasing evidence shows that the core is not merely an inert support for the shell and complex interactions (electronic and mechanical) at the core–shell interface can alter catalytic behavior at the surface of few-layer shells.<sup>14,17,18</sup> However, there is lesser knowledge about the influence of such core–shell interactions on the mechanisms of electrocatalytic reactions, which is the focus of this work.

Hybrid nanoparticles made of coinage-metal cores coated with MoS<sub>2</sub> layers have been reported recently by several groups.<sup>12,14,15,17,19,20</sup> The choice of Au as a *model* core material, albeit expensive as a practical catalyst, is quite straightforward in this context: Au is inert toward oxidation; Au interacts strongly with MoS<sub>2</sub> via the S atoms; and the localized surface plasmon frequency of nanostructured Au is strong and in the visible range. The optical characteristics of transition-metal core–MoS<sub>2</sub> shell nanoparticles present a new hybridized state between the surface plasmons of the metal core and the excitonic transitions of the MoS<sub>2</sub> shell.<sup>14,15,17</sup> These hybridized states can be optically tuned by varying the shape of the nanoparticles.<sup>21</sup> Further modulation of this transition by coupling to the excitonic modes of MoS<sub>2</sub> has been demonstrated,<sup>14,15,17</sup> offering myriad potential applications.<sup>22</sup> It is therefore appealing to study hybrid Au–MoS<sub>2</sub> core–shell structures as model photocatalysts<sup>23–25</sup> and electrocatalysts.<sup>26</sup> However, while photocatalysis using plasmonic nanostructures is a complex process, wherein a correct understanding of the process may be hindered by strong thermal effects from the metal nanoparticles,<sup>27</sup> electrocatalysis is a more straightforward route toward understanding the basic mechanisms of reactions at Au–MoS<sub>2</sub> surfaces.

In this work, various MoS<sub>2</sub> and Au structures were produced: three geometries of Au cores coated by a conformal MoS<sub>2</sub> shell, as well as a complementary structure consisting of freestanding MoS<sub>2</sub> nanoflowers decorated with exposed Au nanoparticles. All materials were processed to produce drop-casted electrodes from a suspension of the catalysts in a Nafion-based ink, and their electrocatalytic activity toward hydrogen evolution was studied. Electrochemical measurements reveal that the structures with the largest Au core (Au nanostars) are superior HER electrocatalysts with lower onset potentials and Tafel slopes relative to Au-decorated freestanding MoS<sub>2</sub>. Density-functional theory (DFT) calculations suggest that the dominant mechanism for the improved catalytic activity is most likely one wherein the Au support promotes the formation of basal-plane S vacancies in MoS<sub>2</sub> and further stabilizes H adsorption at these catalytically active sites. These effects are driven by electron transfer from the Au core to the MoS<sub>2</sub> shell (also seen in the XPS measurements), and such charge transfer could also potentially improve the conductivity of the semiconducting 2H phase, thereby benefiting overall electrocatalytic performance.

## 2. EXPERIMENTAL SECTION

**2.1. Experimental Procedures.** All Au cores were preformed and then encapsulated within MoS<sub>2</sub> by overnight stirring with ammonium tetrathiomolybdate ((NH<sub>4</sub>)<sub>2</sub>MoS<sub>4</sub>) to bind MoS<sub>4</sub><sup>2−</sup> anions to the Au cores, using procedures similar to those reported in our previous publication.<sup>17</sup> Detailed procedures for the synthesis of the Au cores, the encapsulation in MoS<sub>2</sub> and the experimental procedures for the preparation of electrodes for the electrocatalytic measurements are available in the [Supporting Information](#). In addition, two reference samples were produced—freestanding MoS<sub>2</sub> and freestanding MoS<sub>2</sub> decorated with Au nanoparticles; detailed synthesis protocols and characterization are provided in the [Supporting Information](#).

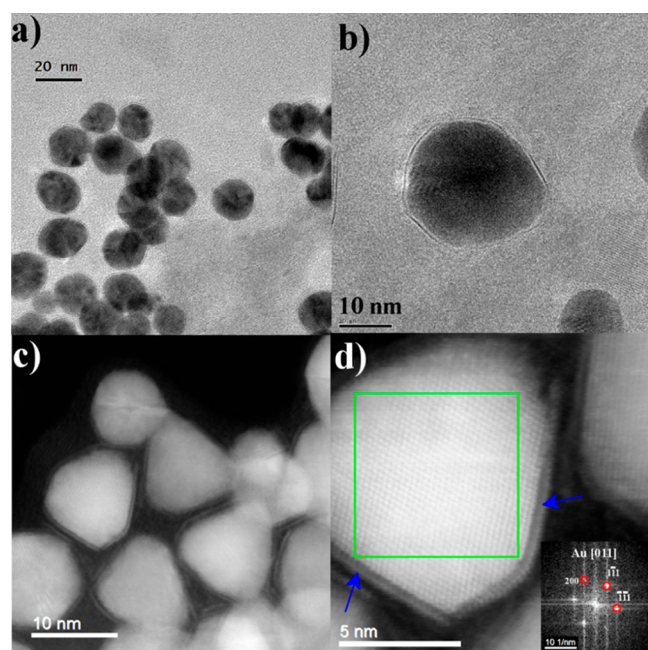
**2.2. Computational Methods.** DFT calculations were performed using Vienna ab Initio Simulation Package (VASP; version 5.4.1).<sup>28,29</sup> The projector-augmented wave (PAW)<sup>30,31</sup> method was used to describe core and valence electrons along with the Perdew–Burke–Ernzerhof (PBE) form<sup>32</sup> of the generalized-gradient approximation to describe electron exchange and correlation. The plane wave kinetic energy cutoff was set to 400 eV, and Gaussian smearing of 0.05 eV was used for Brillouin-zone integrations. Structural optimization was performed using the conjugate-gradient method with a force tolerance of 0.01 eV/Å. The relaxed lattice parameters for bulk Au (FCC) and single-layer MoS<sub>2</sub> (2H phase) are 4.15 and 3.18 Å, respectively. All calculations of isolated MoS<sub>2</sub> were performed using a 4 × 4 supercell; for Au-supported cases, supercells and atomic positions are provided in the [Supporting Information](#). During relaxation of the MoS<sub>2</sub>/Au slab, the bottom two layers of Au were fixed at their bulk positions. To avoid spurious interactions with the periodic images, at least 10 Å of vacuum was inserted normal to the slabs. The Brillouin zone was sampled using a 5 × 5 × 1  $\Gamma$ -centered *k*-point mesh in all calculations. H-adsorption energies were calculated by inserting one H atom on the surface of the slab, and dipole corrections<sup>33,34</sup> were applied along the normal direction. Electronic density-of-states were calculated using 3000 energy-grid points. Similar to prior studies,<sup>35,36</sup> solvent effects have been neglected here, as a first approximation, and are not expected to change the overall conclusions.

## 3. RESULTS AND DISCUSSION

The comparative study presented here is based on various morphologies of Au—spherical nanoparticles, elongated nanorods, and spiky Au nanostars—encapsulated within a few MoS<sub>2</sub> layers. In addition, a sample of Au nanoparticles deposited on nanostructured MoS<sub>2</sub> was produced, to represent samples where the Au is directly exposed to the surrounding medium and can itself potentially catalyze the HER reaction.

**3.1. Materials Analysis and Characterization with Electron Microscopy.** **3.1.1. Gold Particles Embedded within MoS<sub>2</sub> (P-Au@MoS<sub>2</sub>).** Similar to the prior work,<sup>17</sup> the product consisted of nanoparticles with sizes between 10 and 20 nm embedded within almost amorphous material. [Figure 1a](#) shows a TEM image of an assortment of P-Au@MoS<sub>2</sub> nanostructures showing that the Au cores are sheathed conformably by MoS<sub>2</sub> monolayers (blue arrows in [Figure 1d](#)); often, this core–shell morphology is seen more clearly at higher magnification in [Figure 1b](#). [Figure S1](#) shows the core level EELS spectrum of one such nanoparticle, showing Au, Mo, and S signals. [Figure 1c](#) shows a STEM HAADF image of a group of P-Au@MoS<sub>2</sub> nanoparticles while [Figure 1d](#) displays a single nanoparticle at higher magnification using the same technique. Automatic indexing of the FFT performed using the JEMS software<sup>37</sup> revealed that the core of the nanoparticle corresponds to the Au face-centered cubic structure seen along the [011] zone axis. The Au core forms a multiply twinned particle, typical of Au nanostructures, with a truncated octahedron morphology.

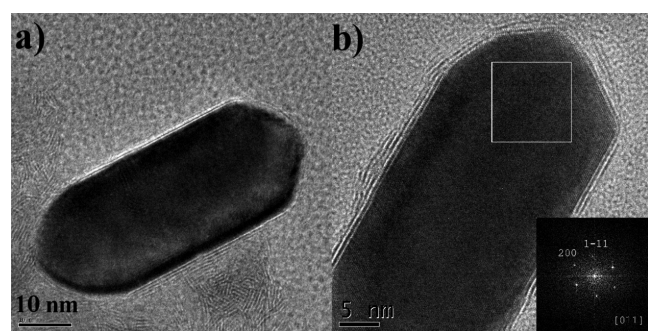




**Figure 1.** Electron microscopy images of P-Au@MoS<sub>2</sub>. (a) TEM image of a few particles. (b) TEM image of a single particle. (c) STEM HAADF image of a group of particles with truncated-tetrahedron Au core structure. (d) HR-STEM HAADF of a single particle. The inset shows the FFT analysis performed at the center of the nanoparticle in the area highlighted by the green square.

### 3.1.2. Gold Rods Embedded within MoS<sub>2</sub> (R-Au@MoS<sub>2</sub>).

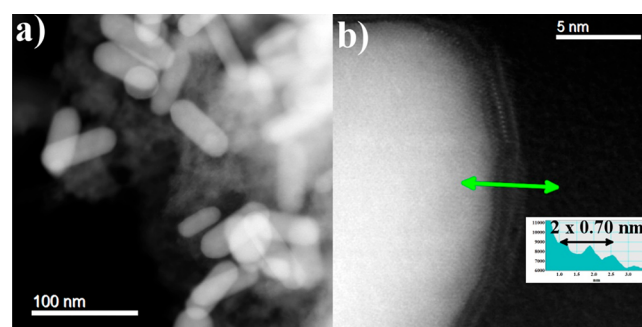
The Au nanorods (aspect ratio of  $\sim 3$ ), were coated with MoS<sub>2</sub> forming a single layer (Figure 2a) or a few-layers thick coating



**Figure 2.** (a, b) TEM image of an individual gold nanorod within MoS<sub>2</sub> (R-Au@MoS<sub>2</sub>). The inset in panel b shows the FFT analysis performed within the area highlighted by the gray square.

(Figure 2b), embedded within MoS<sub>2</sub> amorphous material. The coating of the MoS<sub>2</sub> layer was not as perfect as that for the spherical nanoparticles,<sup>17</sup> frequently consisting of incomplete shells. The Fourier transform (FFT) within the inset of Figure 2b shows the reflections from the (200) and (1 $\bar{1}$ 1) planes of Au. Note that the Au lattice is not coincident with the growth axis in this case. Figure S2 shows another Au nanorod coated with a monolayer of MoS<sub>2</sub>: the perfect contour of the MoS<sub>2</sub> layer and the lattice image of the Au nanorod can be appreciated. The (001) plane of the Au lattice is parallel to the growth axis of the nanorod, which is not always the case (vide infra). It is nonetheless not clear at this point if the lattice of the Au nanorod has any influence on the growth mode of the top MoS<sub>2</sub> layer. Figure S3 shows a tilt series of one such

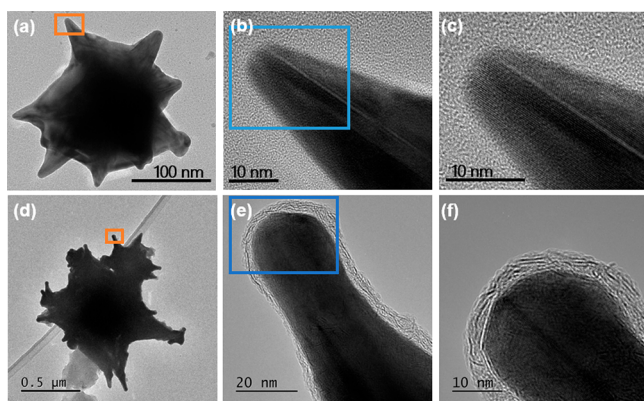
nanorod in three angles ( $0, \pm 25^\circ$ ). The tilt series shows, that, notwithstanding the imperfections in the MoS<sub>2</sub> layer, it encapsulates the entire Au nanorod. Note that the hybrid nanoparticle was placed on a lacy grid with its upper part placed on the amorphous carbon frame and its lower part suspended in vacuum. Therefore, the contrast and image quality of the MoS<sub>2</sub> layer placed on the amorphous carbon substrate is not as good as the one in vacuum. In general, the wrapping of the MoS<sub>2</sub> layer on the tip of the Au nanorod is less continuous than along its facets. Figure 3a shows a STEM-



**Figure 3.** (a) STEM HAADF of an assortment of R-Au@MoS<sub>2</sub>. (b) HR-STEM HAADF micrograph of the tip of a single R-Au@MoS<sub>2</sub> nanoparticle. The green arrow highlights the area used to extract the intensity profile, displayed in the inset, of the STEM HAADF image. The black arrow in the inset marks the distance between three peaks of the intensity profile, showing a 0.70 nm spacing between adjacent layers.

HAADF image of a group of R-Au@MoS<sub>2</sub>. The average value of the aspect ratio is 2.5, with variation between 1.6 and 3.3. Figure 3b shows the tip of one such hybrid nanoparticle in greater detail. The individual MoS<sub>2</sub> layers are clearly discernible on the surface of the nanorod. The MoS<sub>2</sub> interlayer distance has been extracted from this image and is equal to 0.70 nm. This value is slightly higher than the one reported for the interlayer spacing of few-layer MoS<sub>2</sub> (0.62 nm).<sup>38,39</sup> The tip of the Au nanorod is coated with three MoS<sub>2</sub> layers, while two MoS<sub>2</sub> layers are seen at the facet of the nanorod. It is quite common to observe that the coating at the tip of the nanorod contains more MoS<sub>2</sub> layers than at the facets of the nanorods (Figure S4a). Note that the centers of the facets also tend to display a local increase of the number of MoS<sub>2</sub> layers (Figure S4b). Raman measurements of P-Au@MoS<sub>2</sub> and R-Au@MoS<sub>2</sub> showed that the separation of the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> peaks is 22–23 cm<sup>−1</sup> and there is an observed shift of the E<sub>2g</sub><sup>1</sup> to lower frequencies (Figure S5). While a direct comparison with the literature is complex because of the somewhat different morphology of the samples, it is commonly accepted that a peak separation of less than 25 cm<sup>−1</sup> corresponds to two to three molecular layers of MoS<sub>2</sub>,<sup>38</sup> and the E<sub>2g</sub><sup>1</sup> shift is a fingerprint of S vacancies.<sup>40–42</sup>

**3.1.3. Gold Nanostars Embedded within MoS<sub>2</sub> (S-Au@MoS<sub>2</sub>).** Large spiky nanostar Au particles were produced and coated with MoS<sub>2</sub>. TEM images depict the nanostars before (Figure 4a–c) and after the coverage with MoS<sub>2</sub> (Figure 4d–f). The arm length varies between 60 and 360 nm, and the diameter of the entire nanostructure is in the range of 115–475 nm. The coverage on the Au nanostars was more amorphous (Figure 4f), since a thermal annealing stage resulted in the reconstruction of the Au particles into featureless structures and the loss of the spiky nature.

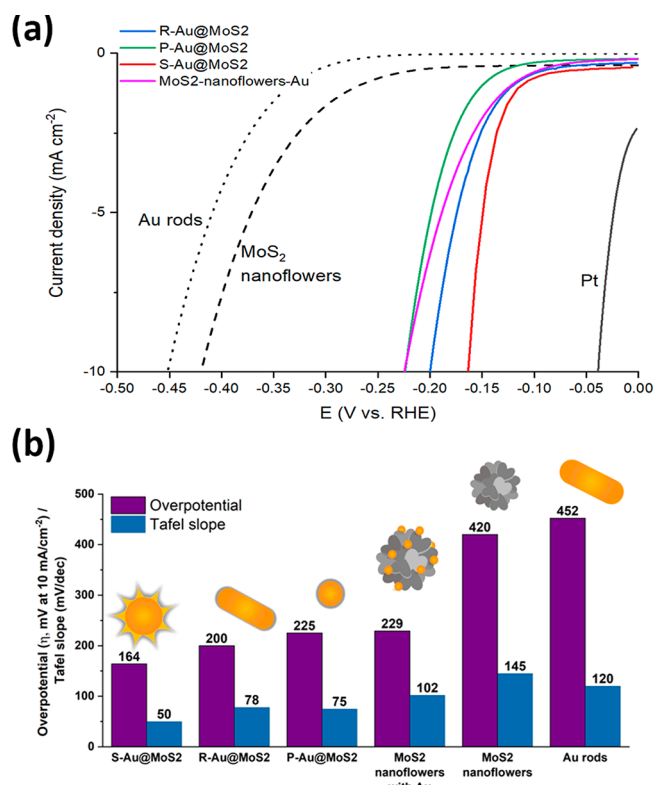


**Figure 4.** TEM images of Au nanostars. (a) Low-magnification image. (b) Higher magnification of the area in the orange frame marked in panel a. (c) Higher magnification of the area in the blue frame marked in panel b. TEM images of Au nanostars coated with MoS<sub>2</sub>: S-Au@MoS<sub>2</sub>. (d) Low-magnification image. (e) Higher magnification of the area in the orange frame marked in panel d. (f) Higher magnification of the area in the blue frame marked in panel e.

Additional data, including the change in absorption spectrum upon MoS<sub>2</sub> deposition, the powder XRD diffraction of the samples, and EDS chemical maps of the S-Au@MoS<sub>2</sub> structures are presented in Figures S6 and S7 in the [Supporting Information](#).

**3.2. Electrocatalytic Measurements.** The electrocatalytic activity toward HER under acidic conditions (0.5 M H<sub>2</sub>SO<sub>4</sub>) was evaluated using a standard, three-electrode setup with the working electrode prepared by drop-casting the catalysts onto a polished glassy carbon (GC) electrode. The polarization curves are presented in [Figure 5a](#) from which we determine the overpotentials at 10 mA·cm<sup>-2</sup> and the Tafel slopes ([Figure 5b](#)). The overpotentials of the hybrid materials are consistently lower than bare MoS<sub>2</sub> or Au. As reference samples, MoS<sub>2</sub> nanoflowers and Au-deposited nanoflowers were also prepared and measured for the catalytic activity toward HER (see the [Supporting Information](#) for detailed synthesis protocol and TEM images in [Figure S8](#)). The hybrids exhibit an improvement of more than 200 mV of the overpotential (at 10 mA·cm<sup>-2</sup>) relative to edge-oriented MoS<sub>2</sub> nanoflowers, regardless of whether the Au particles are embedded within MoS<sub>2</sub> or exposed on its surface. The larger the Au structures, the lower is the overpotential at 10 mA·cm<sup>-2</sup>, reaching 164 mV for S-Au@MoS<sub>2</sub>. Interestingly, the structures with similar sizes of Au particles (P-Au@MoS<sub>2</sub> and nanoflowers deposited with Au particles) showed close results. The HER kinetics is also estimated using the corresponding Tafel plots ([Figure 5b](#)) from which we observe that MoS<sub>2</sub> nanoflowers or Au alone show inferior catalytic properties compared with the hybrid nanoparticles considering both the overpotential and the Tafel slope. The hybrid Au core–MoS<sub>2</sub> shell also shows substantially lower Tafel slopes, reaching 50 mV/dec for S-MoS<sub>2</sub>, indicating a faster rate of the HER for the hybrids. The exchange current density, *j*<sub>0</sub>, is presented in [Table S1](#) and shows moderately higher values for the hybrids.

Additional insights into the improved catalytic activity of the hybrid Au@MoS<sub>2</sub> can be obtained via first-principles DFT calculations. In particular, we seek to understand, via DFT calculations, the relative importance of mechanical strains and electronic interactions that arise at the Au–MoS<sub>2</sub> core–shell interface, and the implications of these effects for the catalytic

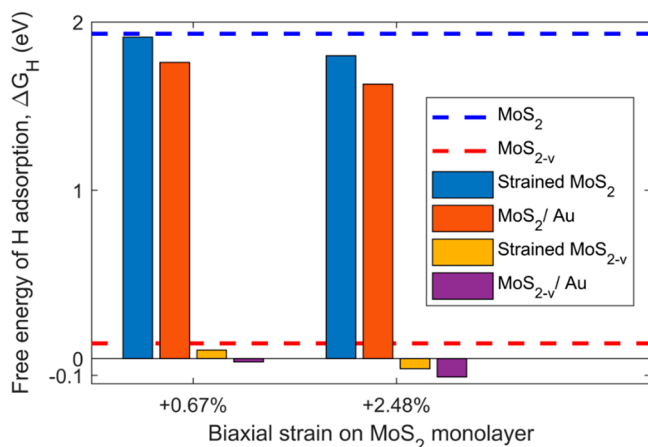


**Figure 5.** (a) HER polarization curves. The polarization curves are based on geometrical current density and not IR corrected. (b) Corresponding Tafel slopes obtained for the hybrid Au/MoS<sub>2</sub> materials, Au (bare rods), MoS<sub>2</sub> nanoflowers, and commercial 10% Pt/C in 0.5 M H<sub>2</sub>SO<sub>4</sub> (pH ~ 0.25). The scan rate of the polarization curves was 10 mV·s<sup>-1</sup>.

properties of the MoS<sub>2</sub> shell. We modeled the Au@MoS<sub>2</sub> nanoparticle as a monolayer of MoS<sub>2</sub> (1L-MoS<sub>2</sub>) supported on an Au(111) slab. As the calculated lattice mismatch between the MoS<sub>2</sub> monolayer and the Au surface is 7.72%, it is unrealistic to enforce epitaxial matching on the MoS<sub>2</sub> layer; the alternate approach of compressing the Au slab<sup>34</sup> leads to artifacts in the electronic structure of the support. While it is reasonable to expect some degree of tensile strain in the curved MoS<sub>2</sub> shell, the precise values for our samples are difficult to extract experimentally as some disorder of the MoS<sub>2</sub> shell is also expected, based on previous data.<sup>43</sup> Hence, we chose to model the supported hybrid Au/MoS<sub>2</sub> system by rotating the MoS<sub>2</sub> layer to produce two coincident-site lattices, one with a small biaxial tensile strain,  $\epsilon = +0.67\%$ , and the second with a larger tensile strain of  $\epsilon = +2.48\%$  (see the [Supporting Information](#) for more details). While it is not possible to account fully for the geometric complexity of the Au cores, the slab models employed here nevertheless allow us to explore separately the strain and support effects that influence the catalytic activity of the MoS<sub>2</sub> shell, thus allowing for broader conclusions beyond the specific set of Au core structures reported here.

[Figure 6](#) displays the free energy of hydrogen adsorption ( $\Delta G_H$ ), a widely used descriptor of HER activity in acidic media,<sup>44–46</sup> for unsupported and Au-supported 1L-MoS<sub>2</sub> at standard conditions [300 K and a potential of 0 V vs reversible hydrogen electrode (RHE)].  $\Delta G_H$  was calculated as



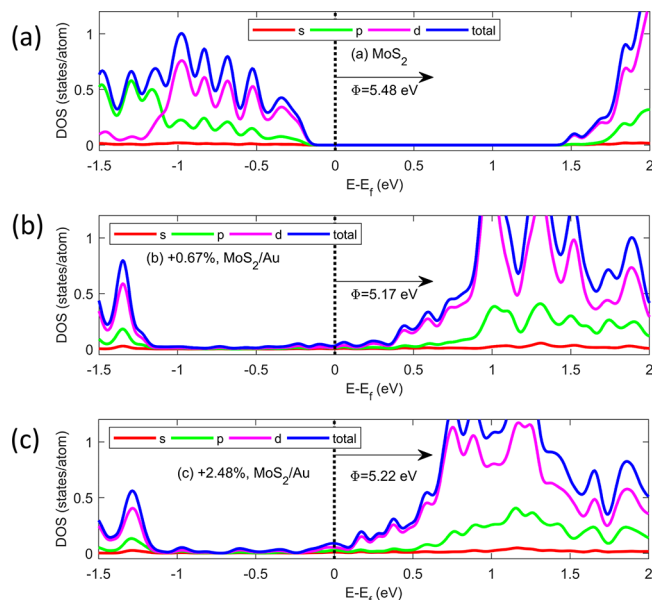


**Figure 6.** Free energy of H adsorption ( $\Delta G_H$ ) at S top sites and S-vacancy sites on the basal plane of freestanding  $\text{MoS}_2$  monolayer and  $\text{Au}(111)$ -supported  $\text{MoS}_2$  monolayer at biaxial tensile strains of +0.67% and +2.48%.  $\text{MoS}_2$  and  $\text{MoS}_{2-v}$  represent the defect-free and defective single layers, respectively. Blue and red dashed lines are the reference values of  $\Delta G_H$  at a S top site and a vacancy site, respectively, in the basal plane of an unstrained, freestanding  $\text{MoS}_2$  monolayer.

$$\Delta G_H = E_{\text{slab}+\text{H}} - E_{\text{slab}} - \frac{E_{\text{H}_2}}{2} + \Delta E_{\text{ZPE}} - T\Delta S \quad (1)$$

where  $E_{\text{slab}+\text{H}}$ ,  $E_{\text{slab}}$ , and  $E_{\text{H}_2}$  are the total energies of either 1L- $\text{MoS}_2$  or the composite 1L- $\text{MoS}_2/\text{Au}(111)$  slab with an adsorbed H atom, the clean slab (no H), and an  $\text{H}_2$  molecule in the gas phase, respectively;  $\Delta E_{\text{ZPE}}$  and  $\Delta S$  are the differences in zero-point energies and entropies, respectively, of an adsorbed H atom and its reference state in the  $\text{H}_2$  molecule. As per Sabatier's principle,<sup>47</sup> a surface with  $\Delta G_H = 0$  is an ideal HER catalyst<sup>46</sup> with negative/positive values of  $\Delta G_H$  leading to over-/underbinding of H. In agreement with previous reports,<sup>36,48–50</sup> we find that the basal plane of the unstrained  $\text{MoS}_2$  monolayer is inert toward HER due to its highly unfavorable  $\Delta G_H \sim 1.9$  eV at a S top site. At +0.67% and +2.48% strain, the adsorption energy decreases by 0.02 and 0.13 eV, respectively; the addition of the  $\text{Au}(111)$  support further lowers  $\Delta G_H$  by 0.15 eV ( $\epsilon = +0.67\%$ ) and 0.16 eV ( $\epsilon = +2.48\%$ ). While the electronic effect of the Au support is more significant than the strain effect, the resulting adsorption energies at the S top sites are still too thermodynamically unfavorable for HER. However, it is well-known that basal-plane S vacancies in  $\text{MoS}_2$  are highly active for HER,<sup>36,48</sup> and indeed, we find that  $\Delta G_H$  drops to  $\sim 0.1$  eV at a S-vacancy site in unstrained  $\text{MoS}_2$  monolayer. At +0.67% and +2.48% strain,  $\Delta G_H$  decreases slightly to 0.05 and  $-0.06$  eV, respectively, and the addition of the Au support further reduces these adsorption energies to  $-0.02$  and  $-0.11$  eV. Thus, at low to moderate tensile strains, the support and strain effects act in synergy and render S-vacancy sites more favorable for HER.

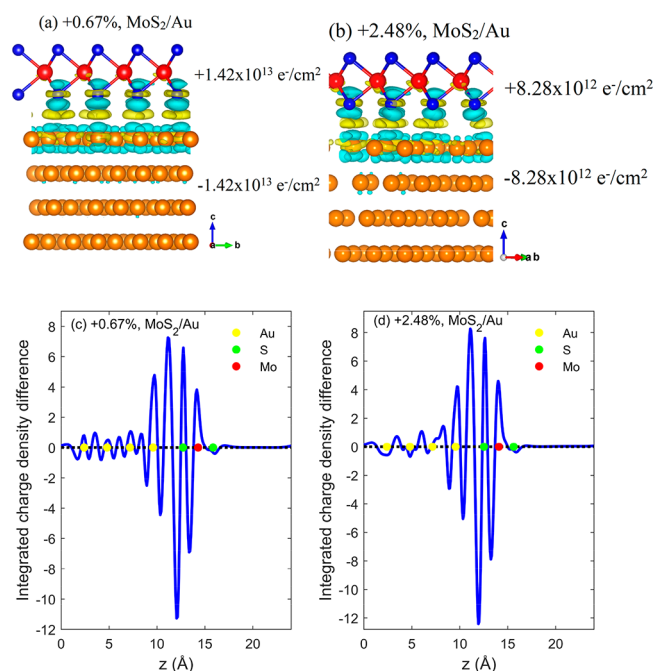
Electronic structure calculations offer insights, beyond thermodynamic calculations, into the microscopic interactions between the  $\text{MoS}_2$  monolayer and the  $\text{Au}(111)$  support. Figure 7 displays the total and angular-momentum-projected DOS of a pristine, unsupported  $\text{MoS}_2$  monolayer and the Au-supported counterparts at +0.67% and +2.48% biaxial strains. The absence of electronic states near the Fermi level for unsupported  $\text{MoS}_2$  monolayer renders H adsorption unfavorable, which is well-documented.<sup>50</sup> However, when an  $\text{MoS}_2$



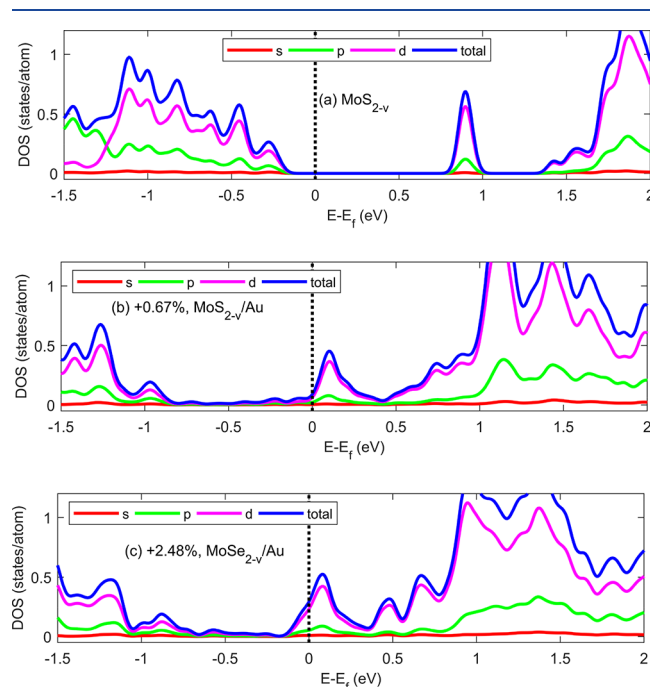
**Figure 7.** Total and angular-momentum-projected density-of-states (DOS) of (a) unsupported and unstrained 1L- $\text{MoS}_2$ , (b) 1L- $\text{MoS}_2$  on  $\text{Au}(111)$  with +0.67% strain, and (c) 1L- $\text{MoS}_2$  on  $\text{Au}(111)$  with +2.48% strain. Only the states projected on to the  $\text{MoS}_2$  layer are displayed for clarity, and the Au states are excluded. The Fermi level is set to zero, and the work function,  $\Phi$ , of the various systems are also indicated in each panel.

monolayer is deposited on an Au support, we notice the gradual emergence of gap states near the Fermi level, which leads to slightly stronger—albeit, still thermodynamically unfavorable—adsorption of H. The density of the gap states increases with increasing tensile strain, and the overall band-gap of  $\text{MoS}_2$  also decreases as known from previous studies.<sup>51–55</sup> A Bader analysis<sup>56–58</sup> reveals that electrons are transferred from the Au slab to the  $\text{MoS}_2$  monolayer ( $1.4 \times 10^{13}$  e $^-/\text{cm}^2$  for  $\epsilon = +0.67\%$  and  $8.28 \times 10^{12}$  e $^-/\text{cm}^2$  for  $\epsilon = +2.48\%$ ). Charge-density-difference plots (Figure 8) are consistent with the Bader analysis and reveal that charge redistribution occurs largely at the  $\text{MoS}_2/\text{Au}$  interface with smaller effects on the upper S basal plane.

For unstrained  $\text{MoS}_2$  monolayer with vacancy defects ( $\text{MoS}_{2-v}$ ; Figure 9), we find that dangling Mo bonds at the S vacancy appear as antibonding gap states,  $\sim 1$  eV above the Fermi level, consistent with previous reports<sup>36,50,59</sup> and these defect states are responsible for enhancing H adsorption. The addition of the Au support injects charge into the  $\text{MoS}_2$  layer and shifts the Fermi level such that the vacancy defect states are much closer to the Fermi level, further stabilizing H adsorption at S-vacancy sites. Thus, our electronic structure studies confirm that, for a monolayer, the support effect, arising from charge transfer from Au to  $\text{MoS}_2$ , is dominant over the strain effect and that the new gap states introduced during this charge-transfer process are responsible for promoting H adsorption on the sulfur basal plane. Furthermore, charge transfer to the  $\text{MoS}_2$  shell can also improve the conductivity of the semiconducting 2H phase, which is, in turn, beneficial for electrocatalysis. XPS measurements on P-Au@ $\text{MoS}_2$  show a shift in the binding energies of the Mo toward lower binding energies (compared with pristine  $\text{MoS}_2$ ) and a shift of the Au binding energies toward the higher energies, usually attributed



**Figure 8.** Charge-density-difference plots for MoS<sub>2</sub> on Au, at (a) +0.67% and (b) +2.48% strain. Isosurfaces are set to  $4.4 \times 10^{-3}$  and  $5.7 \times 10^{-3} \text{ e}^-/\text{\AA}^3$  in panels a and b, respectively, and yellow/cyan colors indicate charge accumulation/depletion. The charge transferred from the Au slab to MoS<sub>2</sub> is indicated in each case. Blue, red, and orange spheres represent the S, Mo, and Au atoms. (c, d) Integrated (over *xy* plane) charge-density differences corresponding to panels a and b as a function of distance normal to the slab (*z*); yellow, green and red dots show the *z*-coordinates of Au, S, and Mo layers in the MoS<sub>2</sub>/Au slab. Charge transfer is strongly localized at the MoS<sub>2</sub>/Au interface.



**Figure 9.** Total and angular-momentum-projected density-of-states of 1L-MoS<sub>2</sub> with a basal-plane S vacancy (MoS<sub>2-v</sub>) for (a) an unsupported, unstrained single layer, (b) MoS<sub>2-v</sub> on Au(111) at +0.67% strain, and (c) MoS<sub>2-v</sub> on Au(111) at +2.48% strain.

to charge transfer from the Au to the Mo (see Figure S10 and Table S2 for the XPS data and analysis).

We also performed a more detailed analysis (see Supporting Information) of the competing thermodynamics of HER and electrochemical desulfurization, the latter having been shown to be a promising route for activation of MoS<sub>2</sub> via generation of sulfur vacancies.<sup>45,60,61</sup> In essence, the calculations show that, at negative potentials, the Au support not only stabilizes H adsorption at the catalytically active sites but also promotes the formation of basal-plane S vacancies in MoS<sub>2</sub>. Tensile strains are slightly detrimental to the formation of vacancies, but once vacancies are formed, strain plays a smaller role in stabilizing H adsorption at these defects (see the Supporting Information). The insights gained from these trends may serve to explain the dominant factors in this system as well as to indicate further possible activation routes for these materials. In a nutshell, the MoS<sub>2</sub> shells are likely defect-rich, in order to encapsulate the Au nanoparticle, producing a combined effect whereby S vacancies, strain, and charge injection from the substrate all make the MoS<sub>2</sub> monolayer a better HER catalyst. Finally, while we have not modeled multilayer MoS<sub>2</sub>@Au structures due to the computational cost, we expect that the improvement in HER activity will be less significant as compared to the monolayer case since the dominant charge-transfer effect is fairly localized to the interface between the first MoS<sub>2</sub> layer and the Au slab; the strain effect is expected to be of nearly similar magnitude between single and multiple layers and, as shown above, is of secondary importance to the charge-transfer effect.

#### 4. CONCLUSIONS

In conclusion, we have presented a method to synthesize core-shell Au@MoS<sub>2</sub> nanoparticles and applied this to realize various MoS<sub>2</sub>-coated Au nanoparticles. Electron microscopy showed that the MoS<sub>2</sub> layers completely encapsulate the Au nanoparticle even though these layers contain defects and seams. Such defects are more common near regions of high curvature such as the tips of the nanorods and nanostars. The hybrid Au@MoS<sub>2</sub> core-shell nanoparticles are superior electrocatalysts relative to freestanding MoS<sub>2</sub>. Theoretical modeling suggests that the improved catalytic activity is likely due to charge transfer from the Au cores to the MoS<sub>2</sub> shells, which improves the electronic conductivity of the semi-conducting MoS<sub>2</sub> shells and also reduces the energetic cost for generation of active chalcogen vacancy sites.

#### ■ ASSOCIATED CONTENT

##### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsam.9b01147.

Extended experimental section (including Materials and Methods); EELS measurements; additional TEM images; UV-vis absorption; EDS maps; XPS and Raman characterization; calculation of the exchange current density; theoretical modeling of the thermodynamics of HER versus electrochemical desulfurization; structural models of Au-supported MoS<sub>2</sub> monolayers (PDF)

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## Notes

The authors declare no competing financial interest.

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