

1 **Characterization of Zirconium Oxides Part II: New Insights on the Growth of Zirconia**
2 **Revealed Through Complementary High-Resolution Mapping Techniques**

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13

14 **Abstract**

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16 Raman mapping, scanning Kelvin probe force microscopy (SKPFM), and scanning electron
17 microscopy with energy dispersive X-ray spectroscopy (SEM/EDS) were combined to investigate
18 oxidized zirconium alloys. Raman provided spatially resolved phase composition and relative
19 stress state. When coupled with SKPFM, phase composition was correlated to Volta potentials
20 differences. The potential of tetragonal zirconia was lower than the metal zirconium, making the
21 tetragonal phase favorable for reaction with diffusing species, thus hindering further oxidation of
22 the relatively cathodic metal. This provides new insight to the theory of the tetragonal phase being
23 an oxidation barrier.

24
25 **Keywords:** Zirconium alloys, Cladding, Oxidation, Raman mapping, SKPFM, SEM/EDS

26 **1. Introduction**

27
28 The adverse environment created from irradiation and variable conditions within a nuclear
29 reactor core induces complex degradation processes of cladding. Zirconium alloys are a viable
30 material choice for fuel cladding, due to its low neutron absorption cross-section and strong
31 corrosion resistance while in a nuclear reactor environment. Dynamic material effects and the
32 extreme environment make it difficult to predict or monitor cladding condition in-core and present
33 significant obstacles for achieving a comprehensive and unified understanding of cladding
34 degradation mechanisms.[1-4]

35
36 Currently, chemical evolution over the lifetime of cladding can only be inferred using post-
37 irradiation examination (PIE),[5] or mimicking the coolant-side conditions with *in-situ*
38 characterization.[6-10] Progress is being made to establish new sensing techniques; for instance,
39 electrochemical impedance spectroscopy (EIS) is becoming established as a cladding degradation
40 sensing technique.[5-7] However, in order to develop accurate EIS models and advance current
41 sensor technology, high-resolution characterization techniques must be used to translate
42 degradation processes along the cladding pathway to equivalent circuit models. This work utilizes
43 scanning Kelvin probe force microscopy (SKPFM), Raman spectroscopy, and scanning electron
44 microscopy (SEM) with energy-dispersive X-ray spectroscopy (EDS) to provide high-resolution
45 co-localized characterization of zirconium oxide and metal/oxide interface.

46
47 Raman spectroscopy provides qualitative to semi-quantitative information on ceramic
48 material composition. This technique has been extensively used for both *in-situ* and post-exposure
49 sectioned analysis of zirconium alloys.[8-13] Raman provided evidence of a bilayer oxide structure
50 grown on zirconium after extended thermal oxidation.[14, 15] The inner layer is a metastable
51 tetragonal phase that has been considered on both ends of the spectrum with regards to being either
52 a protective barrier or a non-participant in the corrosion mechanism of the cladding.[4] The
53 metastable tetragonal phase is believed to be stabilized by a combined effect of interfacial
54 compressive stress, oxygen sub-stoichiometry, and grain size.[8, 9] Via a martensitic phase
55 transformation, tetragonal zirconia transitions to a columnar grained and more porous monoclinic
56 zirconia outer layer. A region of mixed, transient tetragonal-monoclinic zirconia exists between
57 these layers. Raman spectra also provide insight into the stress distribution within the oxide, where
58 high compressive stress is seen in the small, equiaxed tetragonal grains near the metal/oxide
59 interface. With support by oxide porosity, this high compressive stress is believed to reach a critical
value that is followed by a rapid transition to a less stressed state. The stress relief, termed the

60 breakaway phenomenon, is allowed by vertical fracturing in the oxide, providing a pathway for
61 increased oxidation rate.[4, 13]

62 SKPFM is a non-destructive atomic force microscopy (AFM) variant that resolves
63 nanoscale features on the material surface. The traditional Kelvin probe utilizes the capacitive
64 nature between conductors in electrical contact to determine the contact potential difference
65 (CPD), or Volta potential difference (VPD). When applied to a nanoscale and mobile probe,
66 SKPFM spatially resolves local relative Volta potentials that correspond to microstructural
67 heterogeneities on the material's surface. The VPD between two metals in electrical contact
68 directly relates to the difference in their electron work functions.[Bagotsky] While in an inert
69 environment with minimum surface adsorbates, the Kelvin probe can theoretically acquire the
70 work function difference between probe and sample surface. It is a useful surface characterization
71 technique for specifying phase nobility (i.e., microgalvanic coupling),[16-21] hydrogen
72 precipitation detection,[22-26] characterization of semiconductors,[27, 28] and co-localization
73 with elemental makeup from SEM/EDS.[29] However, SKPFM has yet to be utilized for
74 characterizing nuclear cladding zirconium alloys. The thermally grown zirconia is an n-type
75 semiconductor, where oxide growth at the cladding's metal/oxide interface is limited by oxygen
76 anion diffusion.[4] With the existence of an electric field driving spontaneous high-temperature
77 corrosion, the flux of oxygen anions to the metal is balanced by a flux of electrons to the
78 oxide/coolant boundary via hopping mechanism, thus netting zero current.[3, 4] The ionic
79 conductivity of zirconia is ultimately driven by crystal lattice defects, such as Frenkel defects,
80 Schottky disordering, dislocations, and grain boundaries.[3, 30] A region of oxide bordering the
81 metal substrate, stabilizes as tetragonal zirconia due to high compressive stress and oxygen sub-
82 stoichiometry [Refs 8-10 Kurpaska papers, Polatidis 2013], making it a viable electron acceptor
83 (i.e., p-type semiconductor). This thin tetragonal phase is theorized to be a barrier layer to
84 increased oxide growth. [Motta] Additionally, distribution of other species in zirconia produce
85 spatial heterogeneities in electronic properties, and thus provide a driving force for microgalvanic
86 reactions to occur.[Kratochvilova] Given these notable variabilities in the oxide structure, the
87 nanoscale resolution capabilities of SKPFM can provide spatial mapping of ionic/electronic
88 pathways that regulate the corrosion mechanism. SKPFM can provide new insight into the
89 electronic properties of the zirconia layering, oxygen stoichiometry, and metal/oxide interface.
90 Additionally, SKPFM could be utilized to detect heterogeneities such as secondary-phase
91 precipitates, hydrides, and nitrides.

92 The electronic structure of both oxide and metal has great effect on the catalytic activity at
93 the metal/oxide interface. Also, the inclusion of alloying elements plays a role in catalytic activity
94 and selectivity of metals.[31] An ensemble of carefully selected complementary characterization
95 techniques enables new correlations between electronic properties, microstructural
96 heterogeneities, and elemental composition. These correlations provide an insight of zirconia
97 growth mechanisms yet to be observed with individual techniques. With an improved knowledge
98 of cladding degradation, accurate EIS equivalent circuit models can be used for in-core monitoring
99 of cladding degradation.

100
101 **2. Experimental Methods**

102
103 Zirconium (Goodfellow) and Zr-2.65Nb (ATI Metals) were chosen for the current work
104 (Table 1). Plate samples were isothermally oxidized, with thermogravimetric analysis (TGA) used
105 to monitor mass gain rate, as described elsewhere.[32] Samples were exposed to 80% N₂, 20% O₂

106 environment at 700 °C. After oxidation, samples were sectioned and mounted in epoxy. They were
107 then ground with SiC up to 1200 grit, followed by polishing with 1 and 0.05 µm alumina slurries.
108 Following polishing, samples were cleaned with heated Alconox solution on a soft pad, rinsed
109 with ultrapure water, and air dried. Samples were then immediately transferred to the glovebox
110 AFM for SKPFM analysis.

111
112 *Table 1. Zirconium and Zr-2.65Nb compositions.*

	Fe	Sn	Cr	Nb	C	Hf	O	N	H	Zr
Zr (ppm)	200	-	200	-	250	2500	1000	100	10	bal.
Zr-2.65Nb (wt. %)	0.061	-	-	2.62	-	-	0.106	-	-	bal.

113
114 SKPFM was done with a Dimension Icon AFM (Bruker) in an argon filled glovebox
115 (MBraun, <0.1 ppm O₂ and H₂O). PFQNE-AL probes were used to acquire results, operating under
116 a dual-pass method called FM PF-KPFM.[33] In addition to topography, quantitative
117 nanomechanical (QNM) properties were simultaneously acquired from probe/material
118 interactions. Volta potential differences were acquired when the probe was lifted to a user-defined
119 lift height of 100 nm above the surface, creating a nanoscale scenario synonymous to the traditional
120 Kelvin probe. Prior to SKPFM acquisition of the cladding samples, probes were calibrated by
121 imaging a Bruker PFKPFM-SMPL, consisting of an n-doped silicon substrate with patterned
122 islands of aluminum surrounded by a gold interconnect. This sample was utilized to ensure relative
123 consistency of probes by providing a step-wise VPD map from aluminum to silicon to gold.[33]
124 Image processing and analysis were conducted using NanoScope Analysis V1.8 (Bruker).
125 Topography maps underwent a flattening process to remove sample tilt.

126 Raman spectroscopy was accomplished using a Horiba LabRAM HR Evolution (Horiba
127 Scientific) with a monochromatic 532 nm doubled Nd:YAG laser with 50 mW power and ~0.3 µm
128 spectral resolution. Samples were mounted on a motorized stage with ±1 µm X-Y repeatability
129 and accuracy. Spatial resolution depended upon objective lens magnification, ranging from 721
130 nm to 1.18 µm, and thus Raman maps were acquired with 1 µm spacing between collected spectra.
131 Spectral range of 150-700 cm⁻¹ was used to examine peaks of interest. Spectra were processed and
132 analyzed with LabSpec V6.3.x (Horiba). Spectral arrays underwent a baseline correction to remove
133 background noise. Convolution of Gaussian and Lorentzian peak fitting was collected for spectral
134 arrays. Peak position and amplitude were collected for different peaks in each spectrum and formed
135 into X-Y maps. Distinction of zirconia phase, particularly monoclinic and tetragonal, has been
136 well established.[11] Tetragonal phase, which is thermodynamically stable above 1205 °C,[34] is
137 stabilized at lower temperatures in the grown oxide. The stability of tetragonal phase near the
138 metal/oxide interface is driven by high compressive stress and oxygen sub-stoichiometry.[8] Phase
139 content (i.e., volume fraction) of tetragonal and monoclinic zirconia is estimated to calculate the
140 named percent tetragonality. The first tetragonal peak (T₁) versus neighboring monoclinic peaks
141 (M₃ and M₄) was decidedly used to calculate percent tetragonality (Equation 1). Also analyzed
142 from Raman spectra was relative residual stress, where shifts in monoclinic peak positions directly
143 correlates to relative stress.[13] The M₂ peak was chosen to observe relative residual stress, where
144 a decrease in this peak's position correlates with increasingly relative compressive stress. A
145 calibration is commonly done in order to quantify stress – however, a calibration was not
146 conducted for this work. Rather, the relative stress state was observed to show trends as a function
147 of percent tetragonality and distance from the metal/oxide interface.
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$$\%T_{ZrO_2} = \frac{I(T_1)}{I(M_3) + I(T_1) + I(M_4)} \quad \text{Equation 1}$$

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151 SEM/EDS with a Hitachi S-3400N-II (Oxford Instruments Energy+) at 10-15 keV and 10
 152 mm working distance to provide elemental distribution of each sample. Due to the inherent sample
 153 damage that occurs during Raman (laser beam damage [35]) and SEM/EDS (carbon pyrolysis and
 154 electron beam irradiation [20, 21, 29]), these techniques were done in corresponding order after
 155 SKPFM. In addition, samples were carbon-coated prior to SEM/EDS to provide a conductive layer
 156 over the epoxy mount and inhibit charge build-up.

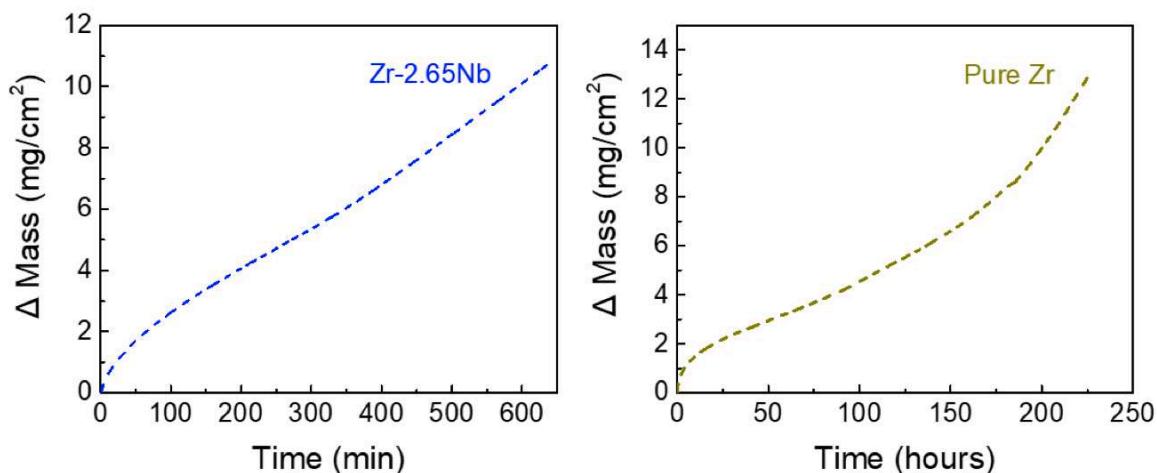
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158 3. Results and Discussion

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160 TGA provided observable mass gain rate for the zirconium samples (Figure 1). Samples
 161 were held to points either before or after breakaway was observed. Breakaway was identified by a
 162 transition of mass gain rate from parabolic to linear kinetics. Faster kinetics and breakaway are
 163 induced much earlier in air than seen in oxygen or steam environments, due to the inclusion of
 164 nitrogen in the mechanism. Nitride formation and subsequent oxidation of those nitrides leads to
 165 increased porosity of the oxide.[32] The post-breakaway pure Zr sample was oxidized for ~226
 166 hours, while the post-breakaway Zr-2.65Nb was oxidized for ~10.5 hours. An additional pure Zr
 167 sample was oxidized for 20 hours, being removed prior to breakaway occurring. The vast
 168 superiority of Zr over Zr-2.65Nb in resisting corrosion and breakaway is opposite of water-side
 169 corrosion performance, where inclusion of niobium improves corrosion and breakaway resistance,
 170 as well as mechanical properties when compared to pure zirconium.[4] In the case of this study,
 171 exposure to high temperature mixed nitrogen/oxygen environment presents different possible
 172 scenarios mimicking an air-cooled core (such as with Transient Reactor Test Facility (TREAT) at
 173 Idaho National Lab) or air ingress and rapid increase in temperature during a nuclear reactor loss-
 174 of-coolant accident (LOCA). Corrosion mechanisms in this particular environment are described
 175 elsewhere.[32]

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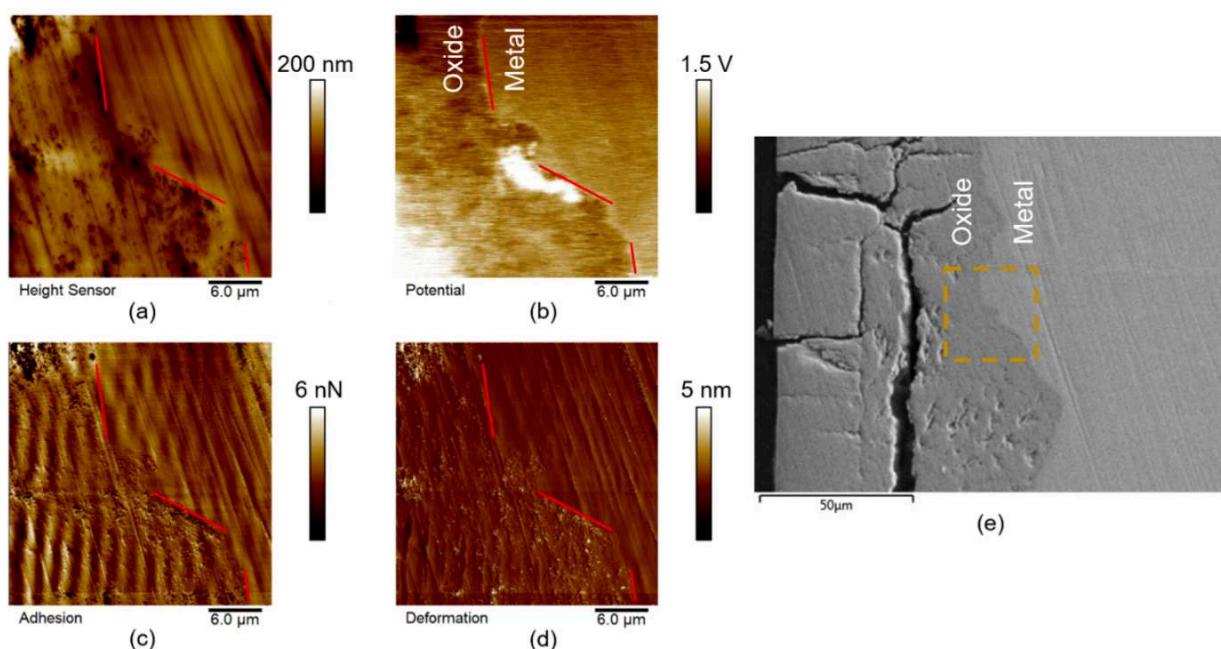
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178 *Figure 1. Normalized mass gain during isothermal oxidation for Zr-2.65Nb (left) and Zr (right).*

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180 Post-oxidation analysis was done on cross sectioned samples. Being the least destructive
181 technique, SKPFM in an argon filled glovebox was done first. Co-localization of SKPFM with the
182 other techniques was accomplished by establishing fiducial marks, such as distinct crack
183 morphology in the oxide or variation in oxide/metal profile.

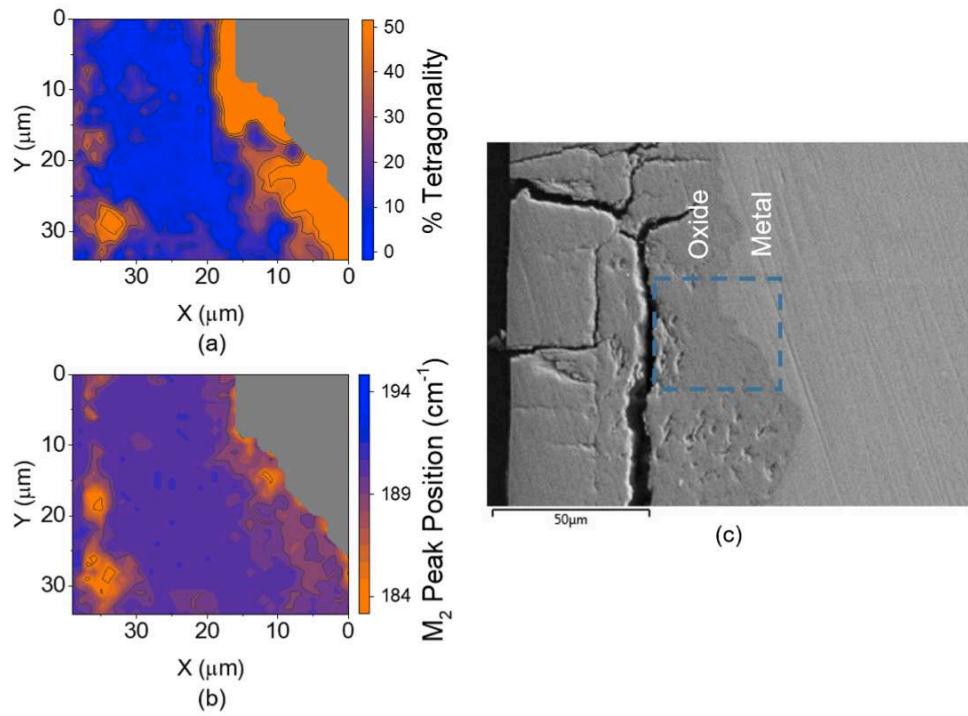
184 For SKPFM, relative Volta potentials (i.e., Volta potential difference, VPD), as well as
185 QNM properties were mapped and used to identify the exact metal/oxide boundary. Figure 2
186 presents height, VPD, adhesion, and deformation acquired for the oxidized Zr sample. The
187 metal/oxide interface is noted with red lines. A clear boundary is seen between the metal and oxide,
188 where the average metal Volta potential is \sim 160 mV higher than the average oxide Volta potential
189 (Figure 2b). Differences in surface morphology (Figure 2a), as well as changes in adhesion (Figure
190 2c) and deformation (Figure 2d) support that the transition point between metal and oxide
191 correspond with the boundary in the Volta potential channel. In addition, a large particle is seen at
192 the metal/oxide interface, displaying higher Volta potential than the surrounding oxide and metal.
193 Based upon the size and location of the particle, as well as the exposing environment, it is likely
194 either a secondary-phase precipitate (SPP) or a nitride. During breakaway in a mixed
195 nitrogen/oxygen environment, nitride particles form at the metal/oxide interface.[32] The
196 difference in lattice density between Zr, ZrO_2 , and ZrN causes lattice mismatch, driving cracking
197 and formation of porous oxide following the oxidation of the nitride particle, leading to pathways
198 for increased cladding degradation.[36-38] On the other hand, elements such as chromium and
199 iron can form SPPs that tend to oxidize at a slower rate than the surrounding zirconium matrix.
200 These slow oxidizing SPPs may support microgalvanic corrosion of the surrounding zirconium, as
201 well as provide pathways via lattice mismatch to cause cracking of the oxide.[4] At higher
202 temperatures, the mobility of iron and chromium should increase, allowing possible formation of
203 SPPs, even when low in concentration. The higher measured Volta potential of the particle seen in
204 Figure 2 supports the theory that it acts as a local cathode relative to the surrounding oxide.
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207 *Figure 2. AFM results for sectioned sample of oxidized Zr. (a) Height, (b) relative Volta potential, (c) adhesion, and*
208 *(d) deformation images with red lines separating metal and oxide. (e) SEM image with area where SKPFM was*
209 *performed.*

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Raman mapping of the same area provided insight into the separation of oxide phases. Figure 3 shows maps of both percent tetragonality (i.e., phase content) and monoclinic peak position (i.e., relative stress state). At the metal/oxide interface (metal is gray in color), tetragonal-rich phase is evident, as are shifts to lower monoclinic peak position that corresponds to increasing compressive stress. As the distance from the interface increases, relative stress decreases, as does the concentration of tetragonal phase. There are notable regions in the bulk oxide with higher percent tetragonality, which may correspond to the “relaxed-tetragonal” phase. It has been shown that this other tetragonal phase stabilizes in the bulk of the oxide via sub-stoichiometry, absent of compressive stress supporting its formation. However, in the case of a few of the tetragonal-rich regions in the bulk of the oxide, a notable correlation between high tetragonality and compressive stress is seen. Looking at the SEM image showing the collected Raman map (Figure 3c), these regions appear to be near a large horizontal crack, along with highly porous oxide.



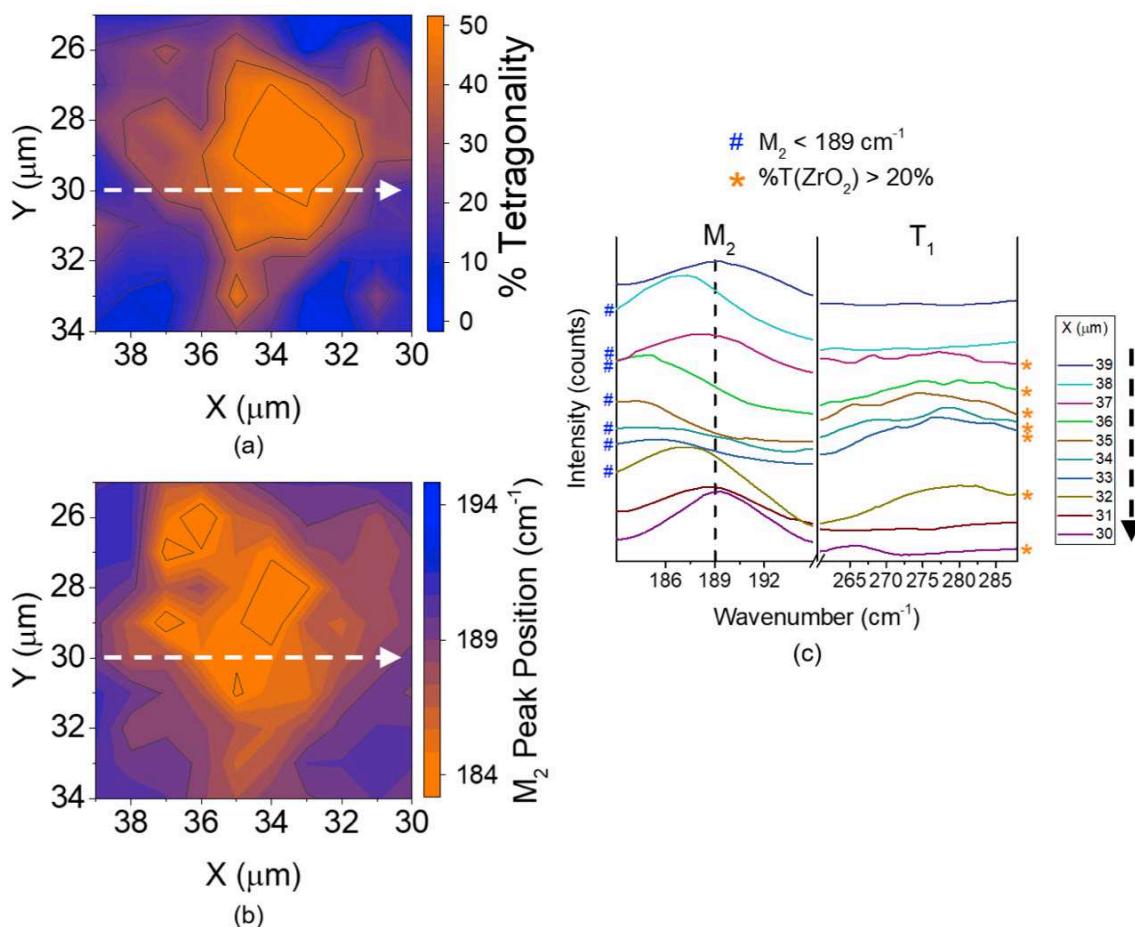
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Figure 3. Raman mapping results for sectioned sample of oxidized Zr. (a) Percent tetragonality and (b) monoclinic (M_2) peak position maps. (c) SEM image with area where Raman mapping was performed.

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A particular region in the bottom left of the Raman map in Figure 3 was further investigated to determine the reason for high compressive stress in the bulk of the oxide. Considering the spectra along the line in Figure 4, a few observations can be made. The spectra that have lower monoclinic peak position and tetragonal-rich peaks are noted in Figure 4c. A majority of spectra with high compressive stress (i.e., low monoclinic peak position) correlate with the presence of tetragonal-rich phase. Two spectra contradict this correlation at $X = 30 \mu\text{m}$ and $38 \mu\text{m}$. For the former, the tetragonal peak is near 265 cm^{-1} , which is different than tetragonal peak positions seen in the other spectra ($275\text{--}285 \text{ cm}^{-1}$). Additionally, there is a lessened presence of compressive stress, as the monoclinic peak is positioned at a higher wavenumber. These factors correspond

237 with the relaxed-tetragonal phase, where relaxed-tetragonal peak position is lower than peak
 238 position of tetragonal phase stabilized by compressive stress.[8] For the $X = 38 \mu\text{m}$ spectra, it
 239 appears that high compressive stress is present, but no tetragonal peak is observed. This appears
 240 to be the transition point from the tetragonal-rich region to monoclinic-rich, even though there is
 241 still high compressive stress. These observations support the theory that oxygen sub-stoichiometry
 242 must be present with notable compressive stress to stabilize tetragonal phase.[8] For all other
 243 spectra with noted tetragonal phase ($X = 32 - 37 \mu\text{m}$), the location of the tetragonal peaks are at
 244 $275-285 \text{ cm}^{-1}$, corresponding with what is typically seen for stress-stabilized tetragonal phase.
 245 Additionally, the M_2 peak position for these spectra supports presence of a relative compressive
 246 stress with lower peak positions. Therefore, the stabilization of this tetragonal zirconia is driven
 247 by both oxygen sub-stoichiometry and compressive stress, much like the tetragonal phase present
 248 near the metal/oxide interface. The presence of the large crack close to this zirconia may support
 249 the change in stress state and thus stabilization of tetragonal phase.
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 252 *Figure 4. (a) Percent tetragonality and (b) monoclinic zirconia (M_2) peak position maps of region within Raman*
 253 *mapped area in Figure 3. (c) Raman spectra across dotted line in (a-b), focusing on M_2 and tetragonal zirconia (T_1)*
 254 *peak positions. Spectra labeled with a blue '#' present M_2 peaks corresponding to compressive stress (i.e., notably*
 255 *lower peak position). Spectra labeled with an orange '*' present tetragonal peaks warranting tetragonal-rich spectra*
 256 *via Equation 1.*

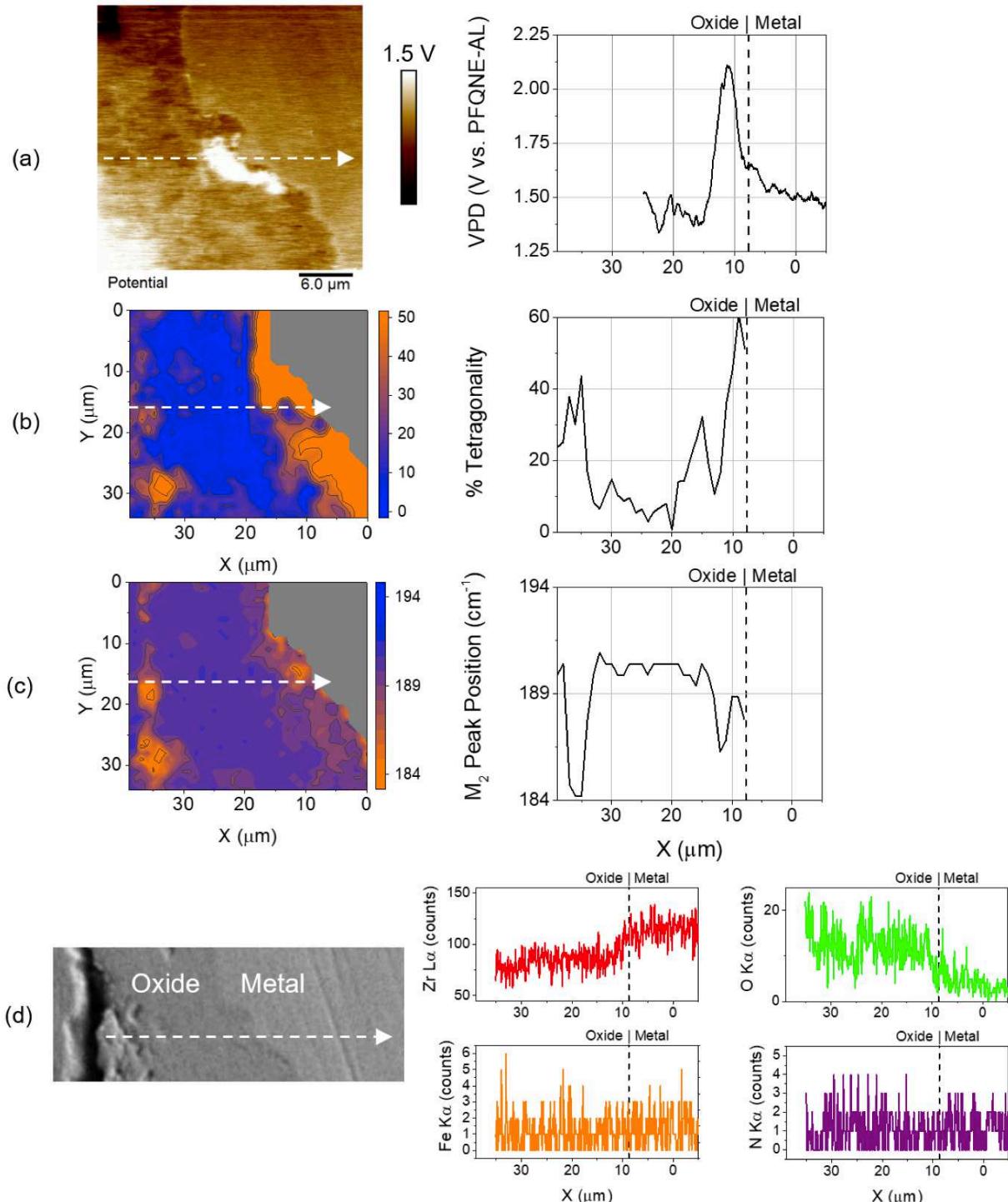
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258 Correlations between the SKPFM and Raman maps were also made. Within the Raman
259 tetragonality map (Figure 3a), at the metal/oxide interface there is an area of lower percent
260 tetragonality with similar shape and size to the cathodic particle seen with SKPFM (Figure 2b). A
261 closer look at this particle is seen in Figure 5. The data acquired via SKPFM (Figure 5a), Raman
262 mapping (Figure 5b-c), and SEM/EDS, (Figure 5d) follow a line that crosses the cathodic particle.
263 Co-localization of X-Y position for each technique was done by aligning the distance from the
264 metal/oxide interface to other notable features. The Volta potential for the particle is ~600 mV
265 greater than the neighboring metal, displaying relative cathodic nature (Figure 5a). This particle's
266 Volta potential maximum is at the same point as the largest decrease in M₂ peak position, and thus
267 largest relative compressive stress (Figure 5c). Additionally, this occurs when percent tetragonality
268 is increasing, but is still below its maximum value (Figure 5b). The end of the oxide, where percent
269 tetragonality and M₂ peak position lines abruptly end, is where the particle ends.

270 Further evidence of either nitride or iron-enrichment for this particle can be seen in Figure
271 6, showing individual Raman spectra across a portion of the noted line scan in Figure 5. Spectra
272 nearest the metal/oxide interface (i.e., near X = 8 μm) reveal a high relative tetragonal peak around
273 280 cm^{-1} . As the distance from the metal/oxide interface increases, evidence of a broad peak
274 around the 200-250 cm^{-1} range is seen. The broad peak in this region can help explain why the
275 percent tetragonality map revealed a particle of lower tetragonality with similar shape to the high
276 Volta potential particle. It is likely due to the M₃ monoclinic peak being used to support calculation
277 of percent tetragonality (Equation 1). Since the broad peak arises in this spectral range, it is
278 recognized as the 'M₃ peak', when in actuality a different compositional feature is revealed.
279 Further away from the metal/oxide interface, this broad peak is slowly lost (Figure 6), and the
280 sharp M₃ peak is revealed (i.e., X = 17 – 19 μm). The broad peak shows similar spectral features
281 as those seen in ZrN,[11] as well as hematite (Fe₂O₃).[39] For ZrN, a sharp peak is seen in the 230-
282 235 cm^{-1} range, with a broad shoulder to the left to ~170 cm^{-1} . For hematite, the A_{1g} mode at 225
283 cm^{-1} and its shouldering E_g mode at 247 cm^{-1} are similar to the spectral features of the broad peak.
284 Due to the shape of this broad peaks seen in Figure 6, it is reasonable to conclude that the particle
285 is an iron-rich SPP. The shouldering that is visible to the right of the broad peak likely correlates
286 to the E_g mode of hematite, while ZrN has a shouldering to the left of its peak. In addition, the Fe-
287 rich SPP is likely highly metallic, thus producing Raman spectra with high amounts of noise. The
288 oxidation of cathodic SPPs are slow relative to Zr, thus leaving a heavily metallic particle in an
289 oxidized matrix.[Motta]

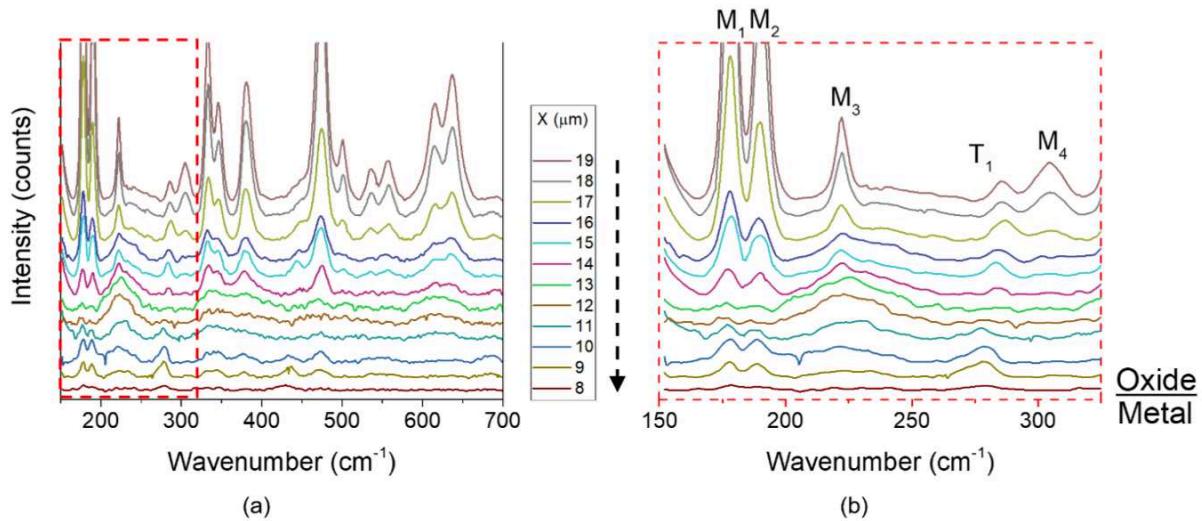
290 For compositional correlation, a line scan of EDS measurements is included (Figure 5d).
291 There is no discernable evidence of concentrated iron or nitrogen counts near the metal/oxide
292 interface. A possible reason for this lack of secondary element response may be due to the damage
293 inflicted upon the sample from the 50 mW Raman laser. When laser power is increased, there is a
294 tradeoff between maximizing signal and increasing surface damage. Damage to the sample surface
295 from the Raman laser makes it difficult to resolve particles even a few microns in size with
296 SEM/EDS. Further evidence of this effect can be seen in Figure 7, where a Zr pre-breakaway
297 sample correlates high Volta potential particles and iron-rich SPPs. Raman mapping was excluded
298 in the analysis of this sample, thus excluding the possible effect of Raman laser damage prior to
299 SEM/EDS. This provides evidence of SPP formation in the Zr sample, regardless of the low
300 concentration of iron in pure Zr (Table 1). Therefore, the absence of iron in the EDS line scan in
301 Figure 5d does not restrict iron from being the contributing element to the cathodic behavior of the
302 particle.

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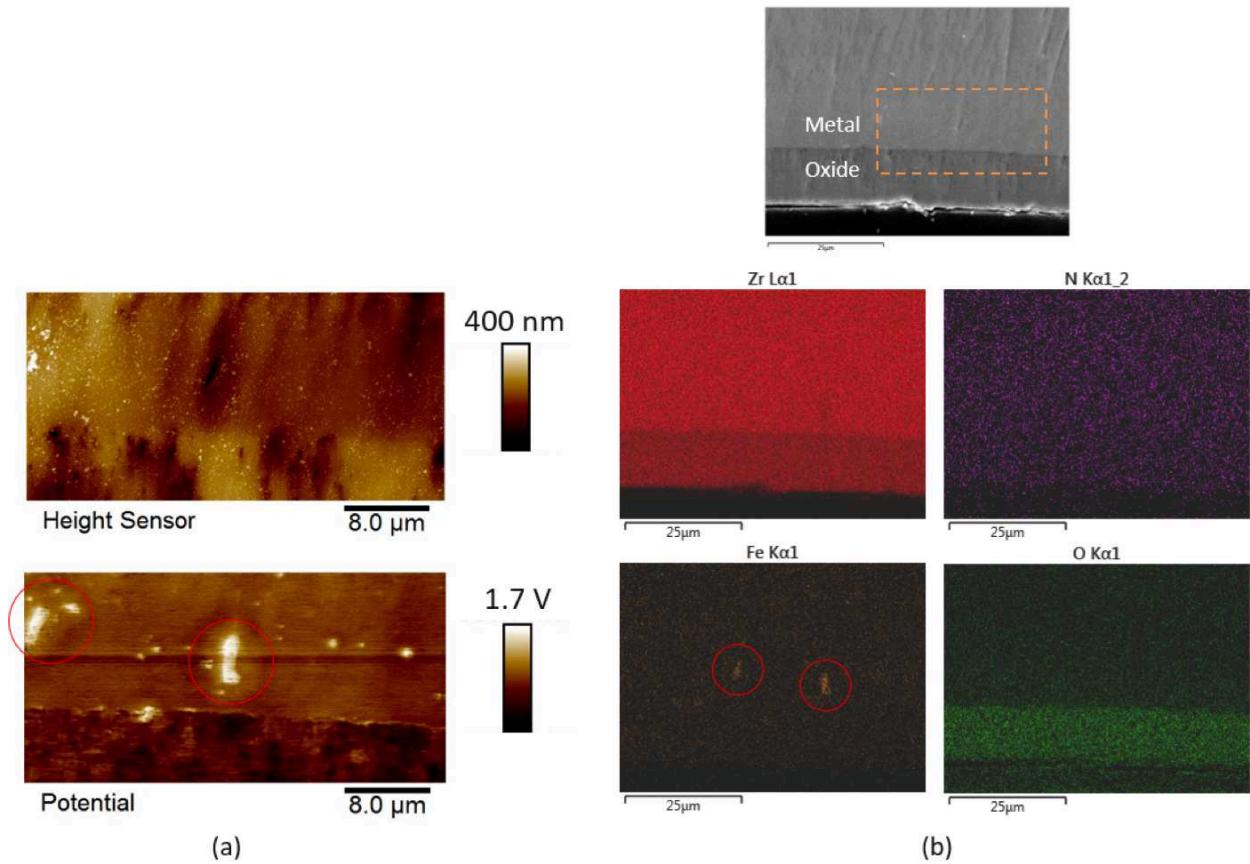
Figure 5. Co-localization between (a) SKPFM Volta potential map with VPD line scan of 1 V range, (b) percent tetragonality and (c) M_2 peak position maps with line scans determined via Raman mapping, and (d) SEM image with EDS elemental line scans for a sectioned sample of oxidized Zr.



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309 *Figure 6. (a) Raman spectra for partial distance across the line scan in Figure 5. (b) Inset of spectra with monoclinic*
 310 *(M) and tetragonal (T) zirconia peaks noted.*

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Figure 7. Co-localization between (a) height and Volta potential via SKPFM and (b) elemental analysis via SEM/EDS for a sectioned sample of oxidized pure Zr (pre-breakaway). Area where SKPFM was performed is included in SEM image. Red circles in Volta potential and EDS maps present correlations between high VPD regions and Fe-rich particles.

318

319 To observe the metal/oxide interface and similarities between features, Figure 8 shows co-
320 localized maps with corresponding line scan data plots for each technique. In this case, a direct
321 transition between metal and oxide is seen at $X \approx 8 \mu\text{m}$. This occurs where the Raman spectra ends,
322 as well as where a sharp transition in Volta potential is seen. The tetragonal-rich region can be
323 distinguished in the $X = 8 - 20 \mu\text{m}$ range (Figure 8b). A drop in Volta potential is visible in this
324 range, providing evidence that the tetragonal-rich region nearest the metal/oxide interface is
325 relatively anodic versus the metal. In fact, the only region within the captured SKPFM image that
326 shows a higher relative Volta potential is in the bottom left area of the image (Figure 8a). This
327 region likely correlates to monoclinic-rich zirconia, while the rest of the zirconia mapped in the
328 SKPFM image is tetragonal-rich, and thus lower in relative Volta potential.

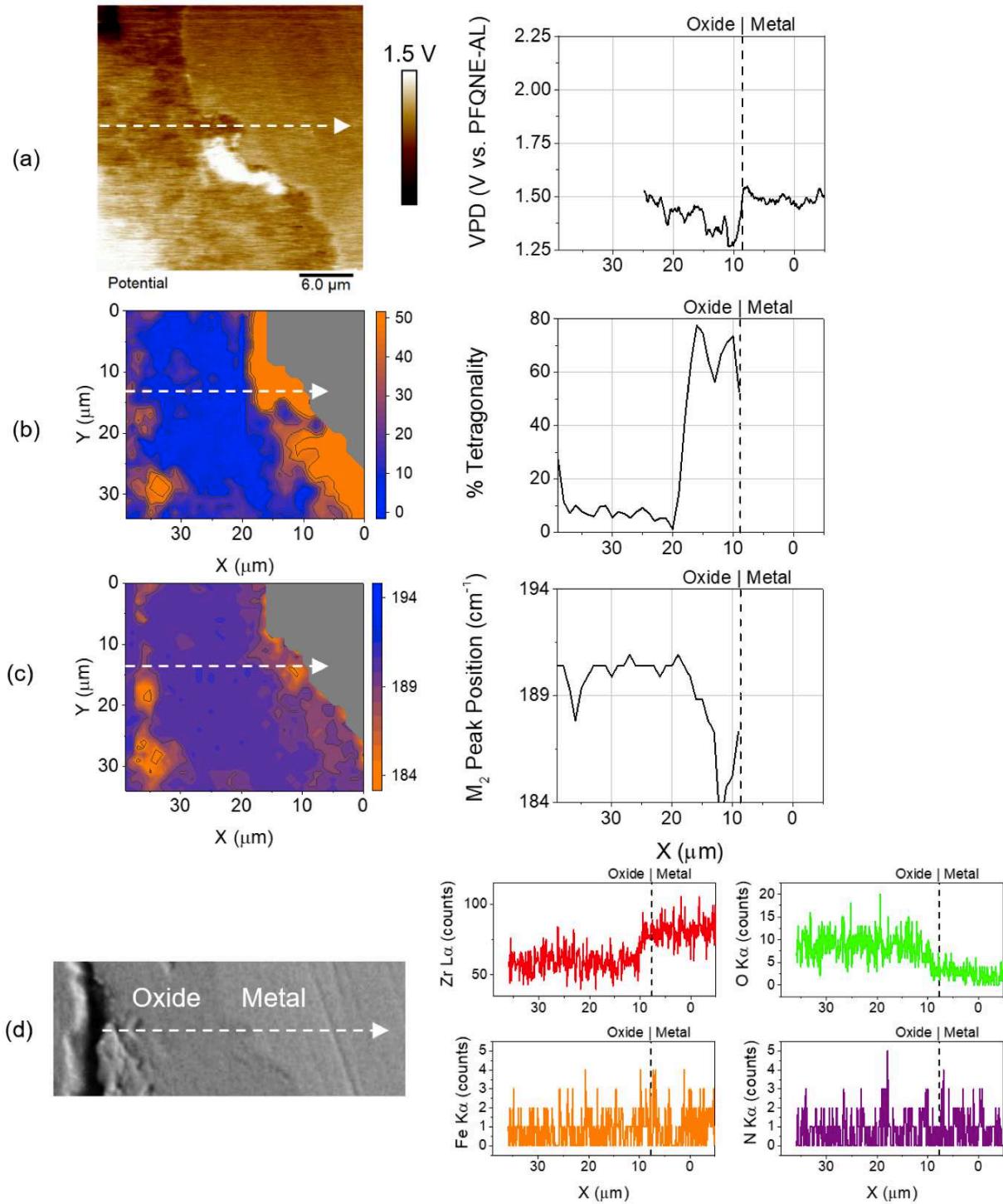
329 Co-localized SKPFM, Raman mapping, and SEM/EDS was also performed on the
330 sectioned Zr-2.65Nb sample (Figure 9). Contrary to Zr, the average Volta potential of the oxide is
331 $\sim 685 \text{ mV}$ greater than that of the metal, revealing relative cathodic behavior for the oxide and
332 relative anodic behavior for the metal (Figure 9a). The Raman map reveals a tetragonal-rich region
333 at the metal/oxide interface, while the bulk of the oxide is monoclinic-rich (Figure 9b). Similarly,
334 compressive stress is seen in the form of lower M_2 peak position for zirconia nearest the
335 metal/oxide interface, while the bulk of the oxide is relatively close to the expected M_2 peak
336 position of 189 cm^{-1} (Figure 9c). Looking at the line scans across the metal/oxide interface,
337 correlations between phase content and VPD can be made. A rapid decline in Volta potential is
338 seen at the same position as an increase in percent tetragonality and decrease in M_2 peak position.
339 In the VPD line scan, there is a thin trench that corresponds to the tetragonal-rich zirconia seen in
340 Figure 9b. After this trench, the VPD slightly increases, coordinating to the Zr-2.65Nb metal. This
341 supports the evidence seen from Zr, where tetragonal-rich zirconia has a lower Volta potential than
342 both the monoclinic-rich zirconia and the metal. A clear metal/oxide boundary is seen in the EDS
343 maps. Due to the substitutional mechanism of Nb in zirconium, no SPPs were seen with EDS.
344 Similarly, no discernible particles were seen in the SKPFM or Raman maps.

345 The tetragonal phase zirconia at the metal/oxide interface has been recognized as a
346 protective barrier, providing a dense layer that limits the oxidation mechanism by diffusion of
347 oxygen anions. The porous monoclinic zirconia layer is theorized to provide an easy pathway for
348 coolant media diffusion through the pores and columnar grain boundaries. Once the corrosive
349 media reaches the tetragonal-rich layer, the small, equiaxed tetragonal grains create an oxidation
350 barrier, forcing transport of oxygen anions via vacancy mechanism to reach the metal and form
351 new oxide. A correlation between tetragonal phase stability and oxygen sub-stoichiometry has
352 already been established.[10] Diffusion resistance of the barrier layer correlates with resistance to
353 charge transfer. Moreover, the lower Volta potential for the tetragonal phase implies that it is in
354 fact more active than the surrounding monoclinic zirconia. Here oxidation is supported by the
355 increased oxygen vacancies near the metal/oxide interface. Oxygen concentration gradients have
356 been seen at the metal/oxide interface prior to the martensitic transition to monoclinic phase. This
357 includes a $\text{Zr}(\text{O})_{\text{sat}}$ region of a few hundred nanometers at the metal/oxide interface, followed by a
358 shallow sub-stoichiometric region (ZrO_{1-x} to ZrO_{1+x}) prior to the stable ZrO_2 zirconia.[4] This
359 coordinates with the stability of tetragonal phase, where both compressive stress and oxygen sub-
360 stoichiometry support its formation. The stable monoclinic bulk zirconia has a balanced
361 stoichiometry (ZrO_2), whereas the tetragonal zirconia is sub-stoichiometric (ZrO_{1-x} to ZrO_{1+x}).
362 This establishes a p-n junction between the n-type monoclinic zirconia and the p-type sub-
363 stoichiometric tetragonal zirconia. This junction provides an additional barrier for anion transfer

364 to the metal substrate.[Ferreira 2002] In addition, the high concentration of holes in this p-type,
365 tetragonal-rich, sub-stoichiometric region causes a reduction in Volta potential relative to the
366 neighboring metal. This provides further support to the barrier layer theory, where diffusing
367 oxygen anions will need to overcome the p-n junction barrier to progress to the oxide/metal
368 interface.

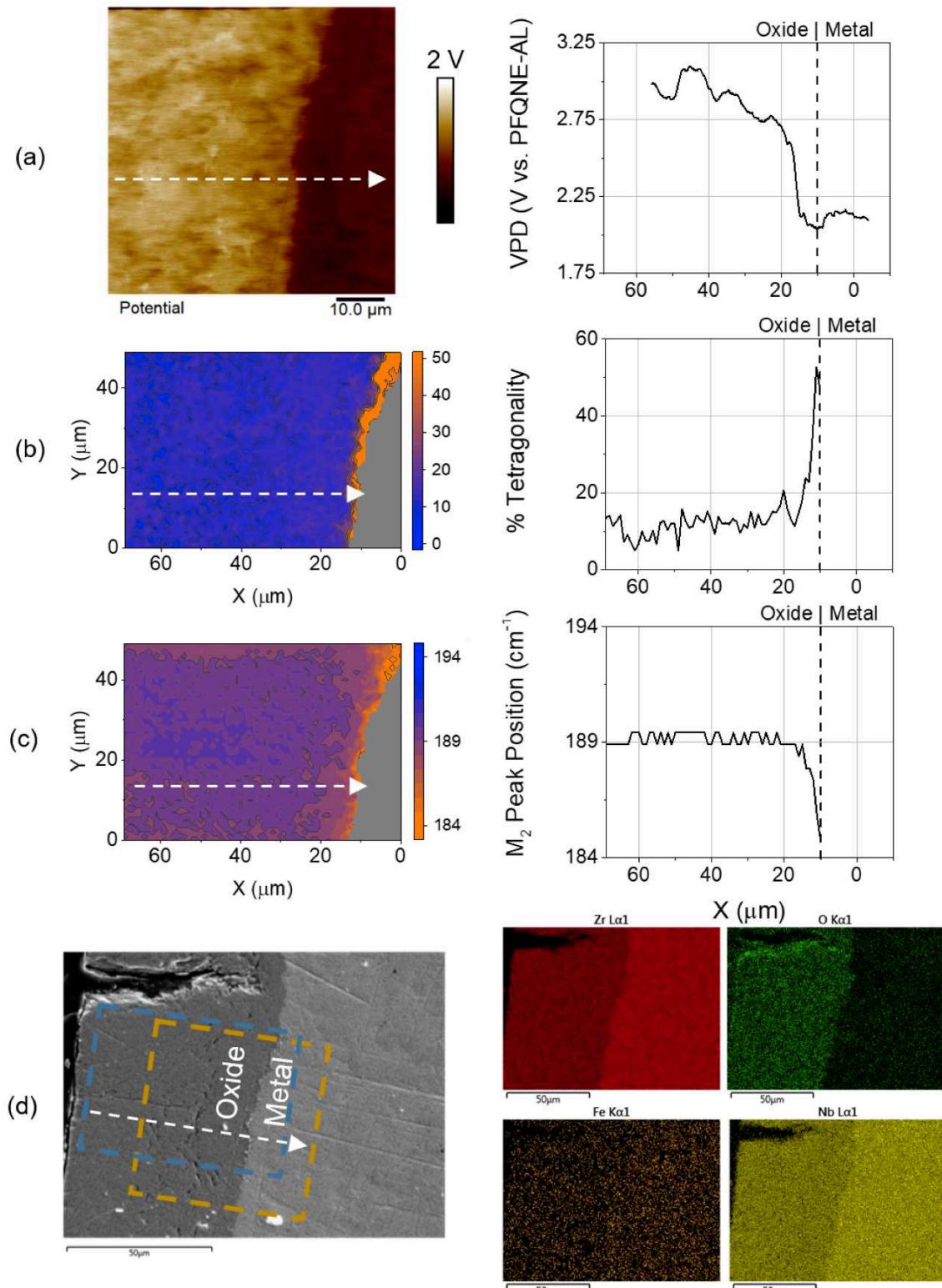
369 All metal/oxide parameters discussed are summarized schematically in Figure 10. The
370 percent tetragonality, stress, Volta potential, and oxygen concentration are shown for different
371 phases of zirconia and zirconium metal. For the metal, the crystal structure is hexagonal closed
372 packed (HCP), under tensile stress, and has minimum oxygen concentration.[3] Crossing the
373 metal/oxide interface, the p-type tetragonal-rich region of zirconia is shown with small, equiaxed
374 grains, high in percent-tetragonality, under high compressive stress, and sub-stoichiometric of
375 oxygen. Further from the metal interface the oxide then undergoes the martensitic phase
376 transformation to an n-type monoclinic zirconia phase, where percent tetragonality is at a
377 minimum, compressive stress is lessened, and the lattice is stoichiometric. Additionally, two small
378 grains of relaxed-tetragonal phase are presented in the bulk of the monoclinic oxide. Here, the
379 percent tetragonality slightly increases and oxygen concentration decreases; however, the stress
380 does not deviate from the rest of the bulk oxide. In this region the oxygen sub-stoichiometry
381 supports stabilization of isolated relaxed-tetragonal phase regions, void of compressive stress. A
382 theoretical representation of a decrease in Volta potential of relaxed-tetragonal phase is included.
383 Based upon the measured anodic behavior of the interface-tetragonal phase, the relaxed-tetragonal
384 phase should also hold relative anodic coupling to the neighboring monoclinic phase with a higher
385 Volta potential. Future studies are needed to further understand the electronic properties and
386 micro-galvanic interactions between oxide phases, secondary particles, and the metal of degraded
387 zirconium cladding.

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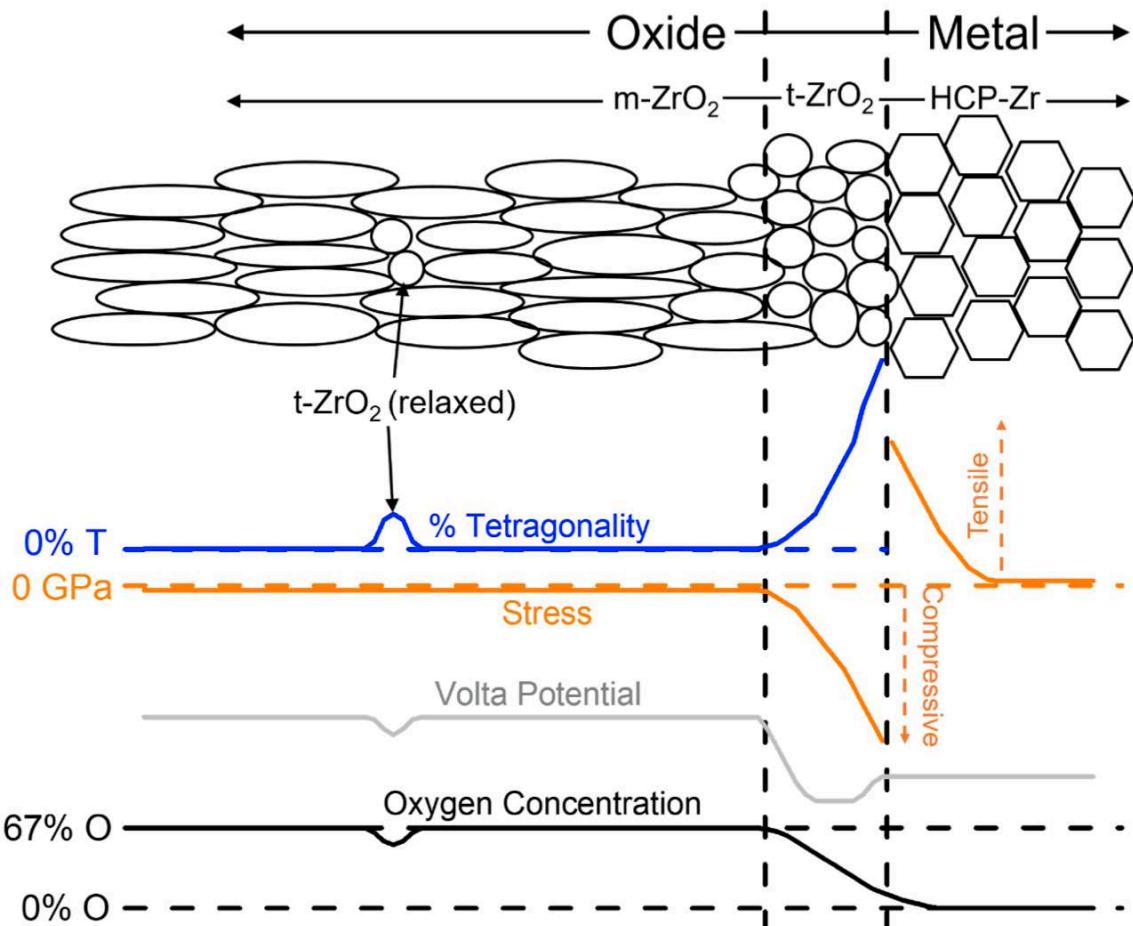
Figure 8. Co-localization between (a) SKPFM Volta potential map with VPD line scan of 1 V range, (b) percent tetragonality and (c) M_2 peak position maps with line scans determined via Raman mapping, and (d) SEM image with EDS elemental line scans for a sectioned sample of oxidized Zr.



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Figure 9. Co-localization between (a) SKPFM Volta potential map with VPD line scan of 1.5 V range, (b) percent tetragonality and (c) M_2 peak position maps with line scans determined via Raman mapping, and (d) SEM image with EDS maps for a sectioned sample of oxidized Zr-2.65Nb.

397



398
399 *Figure 10. Schematic summarizing the different parameters for each zirconia and zirconium phase.*

400
401 **4. Conclusions**
402

403 SKPFM was used to characterize sectioned Zr and Zr-2.65Nb metal/oxide interfaces.
404 SKPFM provided electronic properties for microstructural heterogeneities across these interfaces.
405 For Zr, the average VPD of the oxide near the metal/oxide interface was lower than the average
406 VPD of the metal. A cathodic particle was seen at the metal/oxide interface in the post-breakaway
407 Zr sample. For all samples (pre- and post-breakaway Zr, as well as post-breakaway Zr-2.65Nb),
408 the oxide nearest the metal was lower in VPD than the metal; for the post-breakaway Zr-2.65Nb,
409 oxide further from this interface was much higher in VPD.

410 Raman mapping provided phase and stress distributions for post-breakaway Zr and Zr-
411 2.65Nb. Tetragonal-rich phase was seen near the metal/oxide interface for both samples.
412 Tetragonal phase near the interface is stabilized from compressive stress and oxygen sub-
413 stoichiometry. The relaxed-tetragonal phase was observed in the bulk of the Zr oxide, while it was
414 not observable in high concentrations in the bulk of the Zr-2.65Nb oxide. The relaxed-tetragonal
415 phase is stabilized by oxygen sub-stoichiometry, voided of compressive stress to support stability.
416 Some of the tetragonal phase seen in the bulk of the Zr oxide had a corresponding compressive
417 stress seen by monoclinic zirconia (M_2) peak shift to lower wavenumbers. Therefore, tetragonal
418 phase stabilized with support from compressive stress was seen in the bulk of the Zr oxide. This

419 stress-induced tetragonal phase seen in the bulk Zr oxide was likely supported by nearby crack
420 propagation, creating new stress tensors in surrounding oxide.

421 Co-localization of SKPFM, Raman mapping, and SEM/EDS allowed correlation of
422 different oxide and metal characteristics. For Zr, a broad Raman peak in the 200-250 cm⁻¹ range
423 spatially correlated to the cathodic particle seen with SKPFM. This particle is likely iron-rich, due
424 to the shouldering direction of the characteristic peak and its correlation to the peaks seen in the
425 Raman spectrum of hematite (Fe₂O₃). Raman laser damage removed the opportunity to observe N
426 and Fe EDS counts to confirm the elemental makeup of the cathodic particle. However, when
427 Raman mapping and related laser damage are omitted, correlations between VPD and elemental
428 makeup of particles can be observed, as confirmed by observation of the pre-breakaway Zr sample.
429 Here, the cathodic secondary particles seen with SKPFM were confirmed as iron-rich with EDS.

430 For SKPFM/Raman mapping correlations, tetragonal-rich zirconia at the metal/oxide
431 interface correlated to lower VPD than the neighboring zirconium metal and monoclinic zirconia.
432 Raman spectra revealed a thick tetragonal-rich region near the metal/oxide interface, that also had
433 a lower VPD as seen with SKPFM. A small region of higher VPD was seen in the Zr oxide,
434 correlating to the start of the bulk monoclinic-rich zirconia. This was confirmed with co-localized
435 SKPFM and percent tetragonality mapping of the Zr-2.65Nb sample, where a slight drop in VPD
436 seen at the metal/oxide interface correlates to the thin tetragonal-rich zirconia seen with Raman
437 mapping. Therefore, in addition to the barriers produced by the coolant/oxide and oxide/metal
438 interfaces, the establishment of a p-n junction between the tetragonal and monoclinic zirconia
439 provides another required step in the oxidation mechanism of the metal substrate. This supports
440 the theory that the tetragonal-rich layer at the metal/oxide interface is a protective barrier to further
441 oxidation.

442 SKPFM is a useful characterization technique to support the understanding of cladding
443 oxidation mechanisms. This high-resolution, non-destructive technique can be used in the future
444 for hydride detection and growth, oxide stoichiometry and phase distribution, secondary phases'
445 role in cladding oxidation mechanisms, and irradiation effect of the cladding's electronic
446 properties. When combined with complementary characterization techniques, the methods
447 reported establish novel experimental advancements that can provide new insight into cladding
448 degradation processes.

449

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451

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463 References

464

465 1. B. Cox, Pellet Clad Interaction (Pci) Failures of Zirconium Alloy Fuel Cladding - a
466 Review, *J Nucl Mater* **172** (3), 249-292 (1990) doi: 10.1016/0022-3115(90)90282-R.

467 2. C.G. Yan, R.S. Wang, Y.L. Wang, X.T. Wang and G.H. Bai, Effects of Ion Irradiation on
468 Microstructure and Properties of Zirconium Alloys-a Review, *Nucl Eng Technol* **47** (3), 323-331
469 (2015) doi: 10.1016/j.net.2014.12.015.

470 3. B. Cox, Some thoughts on the mechanisms of in-reactor corrosion of zirconium alloys, *J*
471 *Nucl Mater* **336** (2-3), 331-368 (2005) doi: 10.1016/j.jnucmat.2004.09.029.

472 4. A.T. Motta, A. Couet and R.J. Comstock, Corrosion of Zirconium Alloys Used for
473 Nuclear Fuel Cladding, *Annu Rev Mater Res* **45** 311-343 (2015) doi: 10.1146/annurev-matsci-
474 070214-020951.

475 5. H.K. Jenssen, R. Szoke, N.O. Solum and A. Palencsar, PIE results and mechanistic
476 interpretation on on-line EIS data from the cladding corrosion test IFA-731, HWR-1156, OECD
477 Halden Reactor Project (2017).

478 6. A. Couet, A.T. Motta, A. Ambard and D. Livigni, In-situ electrochemical impedance
479 spectroscopy measurements of zirconium alloy oxide conductivity: Relationship to hydrogen
480 pickup, *Corros Sci* **119** 1-13 (2017) doi: 10.1016/j.corsci.2016.12.008.

481 7. V. Renciukova, J. Macak, P. Sajdl, R. Novotny and A. Krausova, Corrosion of zirconium
482 alloys demonstrated by using impedance spectroscopy, *J Nucl Mater* **510** 312-321 (2018) doi:
483 10.1016/j.jnucmat.2018.08.005.

484 8. L. Kurpaska, J. Favergeon, L. Lahoche, M. El-Marssi, J.L.G. Poussard, G. Moulin and
485 J.M. Roelandt, Raman spectroscopy analysis of air grown oxide scale developed on pure
486 zirconium substrate, *J Nucl Mater* **466** 460-467 (2015) doi: 10.1016/j.jnucmat.2015.06.005.

487 9. L. Kurpaska, J. Favergeon, J.L. Grosseau-Poussard, L. Lahoche and G. Moulin, In-situ
488 stress analysis of the Zr/ZrO₂ system as studied by Raman spectroscopy and deflection test in
489 monofacial oxidation techniques, *Appl Surf Sci* **385** 106-112 (2016) doi:
490 10.1016/j.apsusc.2016.05.074.

491 10. L. Kurpaska, M. Lesniak, R. Jadach, M. Sitarz, J.J. Jasinski and J.L. Grosseau-Poussard,
492 Shift in low-frequency vibrational spectra measured in-situ at 600 degrees C by Raman
493 spectroscopy of zirconia developed on pure zirconium and Zr-1%Nb alloy, *J Mol Struct* **1126**
494 186-191 (2016) doi: 10.1016/j.molstruc.2016.03.001.

495 11. I. Idarraga, M. Mermoux, C. Duriez, A. Crisci and J.P. Mardon, Raman investigation of
496 pre- and post-breakaway oxide scales formed on Zircaloy-4 and M5 (R) in air at high
497 temperature, *J Nucl Mater* **421** (1-3), 160-171 (2012) doi: 10.1016/j.jnucmat.2011.11.071.

498 12. I. Idarraga, M. Mermoux, C. Duriez, A. Crisci and J.P. Mardon, Potentialities of Raman
499 Imaging for the Analysis of Oxide Scales Formed on Zircaloy-4 and M5 (R) in Air at High
500 Temperature, *Oxid Met* **79** (3-4), 289-302 (2013) doi: 10.1007/s11085-012-9331-5.

501 13. K.B. Chong and M.E. Fitzpatrick, Evolution of stress fields and phase content in
502 corroded zirconium cladding materials, *Surf Coat Tech* **324** 140-145 (2017) doi:
503 10.1016/j.surfcoat.2017.05.072.

504 14. P. Barberis, T. Merlemejean and P. Quintard, On Raman spectroscopy of zirconium
505 oxide films, *J Nucl Mater* **246** (2-3), 232-243 (1997) doi: 10.1016/S0022-3115(97)00038-X.

506 15. J. Godlewski, P. Bouvier, G. Lucaleau and L. Fayette, Stress Distribution Measured by
507 Raman Spectroscopy in Zirconia Films Formed by Oxidation of Zr-Based Alloys, in *Zircnoium*
508 *in the Nuclear Industry: Twelfth International Symposium*, West Conshohocken, PA, **1354**, 877-
509 900 (2000).

510 16. F.N. Afshar, J.H.W. De Wit, H. Terryn and J.M.C. Mol, Scanning Kelvin probe force
511 microscopy as a means of predicting the electrochemical characteristics of the surface of a
512 modified AA4xxx/AA3xxx (Al alloys) brazing sheet, *Electrochim Acta* **88** 330-339 (2013) doi:
513 10.1016/j.electacta.2012.10.051.

514 17. J.H.W. De Wit, Local potential measurements with the SKPFM on aluminium alloys,
515 *Electrochim Acta* **49** (17-18), 2841-2850 (2004) doi: 10.1016/j.electacta.2004.01.045.

516 18. N. Sathirachinda, R. Pettersson and J.S. Pan, Depletion effects at phase boundaries in
517 2205 duplex stainless steel characterized with SKPFM and TEM/EDS, *Corros Sci* **51** (8), 1850-
518 1860 (2009) doi: 10.1016/j.corsci.2009.05.012.

519 19. N. Sathirachinda, R. Pettersson, S. Wessman, U. Kivisakk and J.S. Pan, Scanning Kelvin
520 probe force microscopy study of chromium nitrides in 2507 super duplex stainless steel-
521 Implications and limitations, *Electrochim Acta* **56** (4), 1792-1798 (2011) doi:
522 10.1016/j.electacta.2010.08.038.

523 20. C.F. Mallinson, A. Harvey and J.F. Watts, Characterization of Carbide Particles in S-65
524 Beryllium by Scanning Kelvin Probe Force Microscopy, *J Electrochem Soc* **164** (7), C342-C348
525 (2017) doi: 10.1149/2.0271707jes.

526 21. C.F. Mallinson and J.F. Watts, Communication-The Effect of Hydrocarbon
527 Contamination on the Volta Potential of Second Phase Particles in Beryllium, *J Electrochem Soc*
528 **163** (8), C420-C422 (2016) doi: 10.1149/2.0471608jes.

529 22. C. Senoz, S. Evers, M. Stratmann and M. Rohwerder, Scanning Kelvin Probe as a highly
530 sensitive tool for detecting hydrogen permeation with high local resolution, *Electrochem
531 Commun* **13** (12), 1542-1545 (2011) doi: 10.1016/j.elecom.2011.10.014.

532 23. S. Evers, C. Senoz and M. Rohwerder, Hydrogen detection in metals: a review and
533 introduction of a Kelvin probe approach, *Sci Technol Adv Mat* **14** (1), (2013) doi: 10.1088/1468-
534 6996/14/1/014201.

535 24. M.C. Lafouresse, M.L. De Bonfils-Lahovary, C. Charvillat, L. Oger, L. Laffont and C.
536 Blanc, A Kelvin probe force microscopy study of hydrogen insertion and desorption into 2024
537 aluminum alloy, *J Alloy Compd* **722** 760-766 (2017) doi: 10.1016/j.jallcom.2017.06.143.

538 25. L. Oger, M.C. Lafouresse, G. Odemer, L. Peguet and C. Blanc, Hydrogen diffusion and
539 trapping in a low copper 7xxx aluminium alloy investigated by Scanning Kelvin Probe Force
540 Microscopy, *Mat Sci Eng a-Struct* **706** 126-135 (2017) doi: 10.1016/j.msea.2017.08.119.

541 26. Z.L. Hua, S.Y. Zhu, J. Shang, G.X. Cheng, Y.C. Yao and J.Y. Zheng, Scanning Kelvin
542 probe force microscopy study on hydrogen distribution in austenitic stainless steel after
543 martensitic transformation, *Mater Lett* **245** 41-44 (2019) doi: 10.1016/j.matlet.2019.02.089.

544 27. S. Sadewasser, T. Glatzel, M. Rusu, A. Jager-Walda and M.C. Lux-Steiner, High-
545 resolution work function imaging of single grains of semiconductor surfaces, *Appl Phys Lett* **80**
546 (16), 2979-2981 (2002) doi: 10.1063/1.1471375.

547 28. A. Kikukawa, S. Hosaka and R. Imura, Silicon Pn Junction Imaging and
548 Characterizations Using Sensitivity Enhanced Kelvin Probe Force Microscopy, *Appl Phys Lett*
549 **66** (25), 3510-3512 (1995) doi: 10.1063/1.113780.

550 29. M.F. Hurley, C.M. Efaw, P.H. Davis, J.R. Croteau, E. Graugnard and N. Birbilis, Volta
551 Potentials Measured by Scanning Kelvin Probe Force Microscopy as Relevant to Corrosion of
552 Magnesium Alloys, *Corrosion-US* **71** (2), 160-170 (2015) doi: 10.5006/1432.

553 30. N.N. Novik, V.G. Konakov and I.Y. Archakov, Zirconia and Ceria Based Ceramics and
554 Nanoceramics - a Review on Electrochemical and Mechanical Properties, *Rev Adv Mater Sci* **40**
555 (2), 188-207 (2015) doi:

556 31. F.J. Williams and C.M. Aldao, On the origin of metal film work function changes under
557 electrochemical modification, *Surf Sci* **425** (1), L387-L392 (1999) doi: 10.1016/S0039-
558 6028(99)00253-8.

559 32. J.L. Vandegrift, P.M. Price, I.J. Van Rooyen, S. Morrell, D.P. Butt and B.J. Jaques,
560 Oxidation Behavior of Zirconium, Zircaloy-3, Zircaloy-4, Zr-1Nb, and Zr-2.65Nb in Air and
561 Oxygen, *Nuclear Materials and Energy* (2019) doi: 10.1016/j.nme.2019.100692.

562 33. C.M. Efaw, T. Da Silva, P.H. Davis, L. Li, E. Graugnard and M.F. Hurley, Toward
563 Improving Ambient Volta Potential Measurements with SKPFM for Corrosion Studies, *J
564 Electrochem Soc* **166** (11), C3018-C3027 (2019) doi: 10.1149/2.0041911jes.

565 34. O-Zr Phase Diagram, ASM Int., Diagram No. 101191 (2009).

566 35. C. Cardell and I. Guerra, An overview of emerging hyphenated SEM-EDX and Raman
567 spectroscopy systems: Applications in life, environmental and materials sciences, *Trac-Trend
568 Anal Chem* **77** 156-166 (2016) doi: 10.1016/j.trac.2015.12.001.

569 36. C. Duriez, T. Dupont, B. Schmet and F. Enoch, Zircaloy-4 and M5 (R) high temperature
570 oxidation and nitriding in air, *J Nucl Mater* **380** (1-3), 30-45 (2008) doi:
571 10.1016/j.jnucmat.2008.07.002.

572 37. C. Duriez, D. Drouan and G. Pouzadoux, Reaction in air and in nitrogen of pre-oxidised
573 Zircaloy-4 and M5 (TM) claddings, *J Nucl Mater* **441** (1-3), 84-95 (2013) doi:
574 10.1016/j.jnucmat.2013.04.095.

575 38. M. Steinbruck and M. Bottcher, Air oxidation of Zircaloy-4, M5 (R) and ZIRLO (TM)
576 cladding alloys at high temperatures, *J Nucl Mater* **414** (2), 276-285 (2011) doi:
577 10.1016/j.jnucmat.2011.04.012.

578 39. D.L.A. Defaria, S.V. Silva and M.T. Deoliveira, Raman microspectroscopy of some iron
579 oxides and oxyhydroxides, *J Raman Spectrosc* **28** (11), 873-878 (1997) doi: 10.1002/(Sici)1097-
580 4555(199711)28:11<873::Aid-Jrs177>3.3.Co;2-2.

581

582