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# Bankfull Transport Capacity and the Threshold of Motion in Coarse-Grained Rivers

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1       **Bankfull transport capacity and the threshold of motion in coarse-grained rivers**

2

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8

9       **Key points:**

10       • In natural rivers critical Shields stress co-varies with bankfull Shields stress in a manner  
11       predicted by theory

12       • Empirical relations between slope and critical Shields stress are based on a partial sample  
13       of the known parameter space for gravel rivers

14       • Predictions of changes in bankfull transport capacity based on the correlation between  
15       slope and the threshold of motion can be spurious

17 **Abstract**

18 The threshold stress for bed sediment transport exerts a primary control on the geometry and  
19 stability of coarse-grained rivers (diameter  $\geq 5$  mm). Understanding how river bed mobility  
20 couples to channel form is a key mechanistic link for predicting river response to external  
21 perturbations such as land use practices and changing climate. Unfortunately, determination of a  
22 representative threshold stress is notoriously difficult in the field. Empirical studies have  
23 observed that the critical dimensionless shear (Shields) stress ( $\tau^*_c$ ) is correlated with channel  
24 slope, a property that is substantially easier to estimate. Mechanistic models have been  
25 developed to explain the observed correlation; however, limited field data precludes the  
26 widespread application of these models. For practical reasons, the empirical regressions between  
27 slope and  $\tau^*_c$  are utilized as predictive models. Through a large compilation of field data, we  
28 demonstrate that there are two significant problems with using the empirical regressions: (1) they  
29 are based on a partial sampling of the observed parameter space of coarse-grained rivers; and (2)  
30 they do not capture the covariation between the bankfull Shields stress ( $\tau^*_{bf}$ ) and  $\tau^*_c$ . These  
31 regressions provide spurious predictions for the bankfull transport capacity ( $\tau^*_{bf}/\tau^*_c$ ) of gravel-  
32 bed rivers. When site-specific empirical measurements of  $\tau^*_c$  are made, coarse-grained rivers  
33 exhibit a remarkably constant transport capacity that is in close agreement with equilibrium-  
34 channel theory ( $\tau^*_{bf}=1.2\tau^*_c$ ). From these data we advocate that, in the absence of measurements,  
35  $\tau^*_c$  can be reasonably estimated from the  $\tau^*_{bf}$  using equilibrium-channel theory.

36

37 **1. Introduction**

38 A longstanding interest in the science and engineering of rivers is in understanding which flows  
39 are responsible for transporting sediment and how these flows both organize and are shaped by  
40 the channels that convey them (Glover & Florey, 1951; Henderson, 1963; Leopold & Maddock,  
41 1953; Leopold & Wolman, 1957; Wolman & Miller, 1960). This understanding is increasingly  
42 important if we are to predict a river's response under future landscape and climate scenarios  
43 (Hempel, 2018; Phillips et al., 2018; Schmidt & Wilcock, 2008; Slater et al., 2015; Slater &  
44 Singer, 2013). River channel stability is, to first order, controlled by the transport threshold of  
45 material in the bed or channel banks, whichever is harder to entrain (Dunne & Jerolmack, 2018;  
46 Schumm, 1960). In coarse-grained rivers, here considered to be generally gravel-bedded (median  
47 diameter,  $D_{50} \geq 5$  mm), the bed sediment is most often harder to entrain than the banks and

48 hence sets the threshold for altering channel geometry. Single threaded gravel-bed river channel  
49 geometry has been analytically linked to bed-load sediment transport through hydraulics (Parker,  
50 1978, 1979): the summary result is that channel geometry is adjusted so that fluid stress is at the  
51 threshold of motion at the toe of the river banks, and modestly above threshold in the channel  
52 center. This theory has been broadly validated in natural channels and laboratory experiments  
53 (Dade & Friend, 1998; Dunne & Jerolmack, 2018; Métivier et al., 2017; Parker et al., 2007;  
54 Phillips & Jerolmack, 2016; Pitlick et al., 2013; Pitlick & Cress, 2002; Reitz et al., 2014;  
55 Seizilles et al., 2014).

56

57 Yet, the threshold of motion in natural channels remains very challenging to accurately measure,  
58 and data to evaluate hypotheses remain sparse (Buffington & Montgomery, 1997; King et al.,  
59 2004; Mueller et al., 2005) when compared to the available databases of river hydraulic  
60 geometry (see Church & Rood, 1983; Li et al., 2014; Trampush et al., 2014). Standard practice is  
61 to assess mobility through the use of the Shields curve, which is perhaps most valid under  
62 idealized conditions such as normal flow and unimodal bed sediments (Lamb et al., 2008;  
63 Shields, 1936; Wiberg & Smith, 1987). However, even in idealized laboratory conditions the  
64 threshold of motion varies with particle protrusion, bed texture, grain size distribution, and the  
65 structure of the granular bed (Houssais et al., 2015; Kirchner et al., 1990; Masteller & Finnegan,  
66 2017; Pender et al., 2007; Prancevic & Lamb, 2015b; Shvidchenko & Pender, 2000; Wilcock,  
67 1998; Zimmermann et al., 2010). For natural flows where conditions are decidedly less uniform  
68 or steady, the threshold of motion is commonly treated as the measurable motion of a surface  
69 layer (Parker, 1990) and varies both spatially within a reach, temporally within a flood, from  
70 flood to flood, with sediment supply and availability, and with the method of measurement  
71 (Buffington & Montgomery, 1997; Johnson, 2016; Lisle et al., 2000; Marquis & Roy, 2012;  
72 Masteller et al., 2019; Pfeiffer et al., 2017; Prancevic & Lamb, 2015b; Turowski et al., 2011;  
73 Yager et al., 2012, 2018). Within a reach, spatial heterogeneity in grain size organization can  
74 result in unequal mobility especially at shear stresses near the threshold, where smaller particles  
75 are susceptible to turbulent bursts and bed-load flux measurements may strongly reflect partial  
76 transport and not the entire bed (Paola & Seal, 1995; Recking, 2013; Wilcock & McArdell,  
77 1997). These phenomena may be described by some distribution of the threshold for a given  
78 reach; however, additional lines of evidence indicate that even this distribution may be non-

79 stationary and may drift through time due to changes in the river bed state (Charru et al., 2004;  
80 Houssais et al., 2015; Johnson, 2016; Masteller et al., 2019).

81  
82 In natural streams and rivers, a large variety of methods and techniques have been used to  
83 measure bed-load flux and the threshold of motion. These methods largely fall into two  
84 categories, active or passive monitoring. Active methods involve direct measurement of the flux  
85 via physical samplers or registering of impacts (for examples see Bunte et al., 2013; Gray et al.,  
86 2010; King et al., 2004; Reid et al., 1985; Rickenmann et al., 2012) while passive techniques  
87 involve visual determination, tracer particles, competence or largest mobilized particle,  
88 acoustics, and seismometers (see Barton, 2006; Hsu et al., 2011; Phillips et al., 2013; Phillips &  
89 Jerolmack, 2014; Rickenmann et al., 2012; Roth et al., 2016; Wilcock, 1992). However, these  
90 methods remain resource intensive and are not guaranteed to produce a robust estimate of a  
91 threshold of motion unless sampled over larger space and timescales (Monsalve et al., 2016;  
92 Recking, 2013). In some cases, acquiring sufficient bed-load flux measurements could take years  
93 due to the recurrence intervals of floods that are capable of transporting sediment. Despite the  
94 threshold's importance, there is currently a lack of low-cost reliable methodologies to rapidly  
95 assess this variable in the field.

96  
97 Natural and laboratory estimates of the threshold of motion have been observed to possess a  
98 positive correlation with channel slope (Lamb et al., 2008; Mueller et al., 2005). Though initially  
99 counter intuitive, that sediment particles become harder to move at higher gradients, mechanistic  
100 grain-scale models incorporating relative roughness, reduced turbulent intensity, partitioning of  
101 the total shear stress, particle friction angles, and particle lift and drag forces have explained this  
102 correlation under a variety of experimental conditions (Ferguson, 2012; Lamb et al., 2008,  
103 2017a, 2017b; Prancevic et al., 2014; Prancevic & Lamb, 2015a, 2015b; Recking, 2009). Many  
104 of the aforementioned processes necessarily covary with slope and relative roughness, however  
105 additional experiments have isolated reduced turbulent intensity and the lift force as primary  
106 causes for the increase of the threshold of motion with slope (Lamb et al., 2017a). Despite these  
107 advances, the difficulty of measuring the threshold of motion in the field, combined with the  
108 challenge in measuring the necessary parameters to apply the mechanistic models, has led to an  
109 over-reliance on empirical regressions between slope and threshold for natural rivers. In this

110 contribution, we demonstrate that there are two significant problems with using slope-based  
111 empirical regressions as they: (1) are based on a partial sampling of the parameter space of bed-  
112 load rivers, and cannot be extrapolated outside of that range; and (2) do not capture the observed  
113 covariation between the bankfull ( $\tau^{*bf}$ ) and critical ( $\tau^{*c}$ ) Shields stresses and thus can provide  
114 spurious predictions of the bankfull transport capacity ( $\tau^{*bf}/\tau^{*c}$ ). An important point to note here  
115 is that none of the authors of the original studies exploring the correlation between slope and the  
116 threshold of motion (Ferguson, 2012; Lamb et al., 2008; Mueller et al., 2005; Recking, 2009)  
117 suggested that the empirical regressions be used in a predictive fashion, or in lieu of actual  
118 measurements. The use of the empirically based regressions, due to the aforementioned  
119 problems, can lead one to conclude that bankfull transport capacity varies among gravel-bed  
120 rivers, while direct measurements of the threshold indicate that  $\tau^{*bf}/\tau^{*c}$  is remarkably constant  
121 and in quantitative agreement with theoretical predictions (Parker, 1978, 1979). Due to the close  
122 match with theory, we suggest that the bankfull Shields stress can be used to estimate the  
123 threshold of motion as an alternative to the slope-based regressions. Additionally, we illustrate  
124 how using an empirical slope-based regression can lead one to draw potentially incorrect  
125 conclusions, through various sampling strategies of the compiled gravel-bed rivers.  
126

## 127 **2. Data and Methods**

128 Two types of field sites are used within this study: (1) field sites where the threshold of motion  
129 can be reliably estimated from bed-load flux estimates (n=68), and (2) a compilation (n=739) of  
130 coarse-grained river hydraulic geometry to place the hydraulic geometry of the sites where the  
131 threshold of motion was measured into a broader context. For clarity, we refer to the first set as  
132 the ‘threshold’ data and the second as the ‘compilation’ data throughout the remainder of the  
133 manuscript. The compilation and threshold sites both represent samples of the global population  
134 of gravel rivers, however the compilation provides a more complete picture of the variability  
135 within gravel rivers, while the threshold sites represent a smaller partial sample of the of the  
136 compilation’s hydraulic geometry parameter space. The majority of the threshold field sites  
137 come from the study of Mueller et al. (2005) with additional data compiled from the works of  
138 Recking (2010, 2013), King et al. (2004), and others (Andrews, 1994, 2000; Andrews & Erman,  
139 1986; Bunte, 1998; Erwin et al., 2011; Ferguson & Church, 2009; Hinton et al., 2017; Jones &  
140 Seitz, 1980; May et al., 2009; May & Pryor, 2013; McLean et al., 1999; Milhous, 1973; Mueller

141 & Pitlick, 2014; Parker et al., 1982; Phillips & Jerolmack, 2014; Rankl & Smalley, 1992; Ryan  
 142 et al., 2005; Ryan & Emmett, 2002; Smalley et al., 1994; Whitaker & Potts, 2007; Wilcock et al.,  
 143 1996). The compilation data represent field measurements at bankfull conditions of slope ( $S$ ),  
 144 width ( $W$ , m), discharge ( $Q$ ,  $\text{m}^3/\text{s}$ ), depth ( $H$ , m), and the median bed surface grain size ( $D_{50}$ , m).  
 145 These data represent the combination of coarse-grained single thread gravel bed rivers compiled  
 146 by Li et al. (2015), Trampush et al. (2014), Church and Rood (1983), and Phillips and Jerolmack  
 147 (2016). All of the above parameters are not available for each site as the compilation reflects the  
 148 state of the data as collected by the original authors with duplicates removed.

149

## 150 **2.1 Calculating the Threshold of Motion**

151 For the threshold sites we follow the methodology of Mueller et al. (2005) for determining the  
 152 threshold of motion. Their thorough compilation and analysis form the core dataset from which  
 153 the correlations between slope and the threshold of motion have been drawn. We summarize the  
 154 methodology here as it pertains to understanding the current contribution. The primary data  
 155 represent measurements and estimates of the reach-scale fluid driving stress (nondimensionalized  
 156 as the Shields stress,  $\tau^*$ ) and sediment flux per stream width ( $q_s$ ,  $\text{kg}/\text{m}/\text{s}$ ). From these data we  
 157 define the threshold of motion as the dimensionless critical Shields stress ( $\tau_{*c}$ ). The reach and  
 158 cross section scale Shields stress is estimated using hydraulic geometry variables as:

$$159 \quad (1) \quad \tau_* = \frac{\tau}{(\rho_s - \rho)gD_{50}}$$

160 where  $\tau$  is the shear stress ( $\text{Pa}$ ),  $\rho_s$  is the density of sediment (taken here as  $2650 \text{ kg}/\text{m}^3$ ),  $\rho$  is the  
 161 density of water ( $1000 \text{ kg}/\text{m}^3$ ), and  $g$  is the acceleration due to gravity ( $9.81 \text{ m}/\text{s}^2$ ). We  
 162 approximate the shear stress via the depth-slope product as  $\tau = \rho g h S$ , where  $h$  is the flow depth.  
 163 To standardize the measurements across field sites, we nondimensionalize the transport rate as:

$$164 \quad (2) \quad W_* = \frac{Rgq_s}{\rho_s(\tau/\rho)^{1.5}}$$

165 where  $R=1.65$  is the submerged specific density of the sediment (see Parker et al., 1982; Parker,  
 166 1990). The threshold of motion is determined as the median value of Shields stress at which the  
 167 dimensionless transport rate intersects a reference transport rate of  $W^*=0.002$  (Parker, 1990;  
 168 Mueller et al., 2005). The use of a reference transport rate to determine a threshold of motion  
 169 means that, strictly speaking, we have determined the reference Shields stress ( $\tau_{*r}$ ), which is  
 170 close but not necessarily equal to  $\tau_{*c}$ . Throughout the rest of the manuscript references to the

171 threshold of motion and data analysis specifically refer to  $\tau_{*r}$  as a proxy for  $\tau_{*c}$ . For field sites  
172 where the bulk of the sediment flux measurements are either above or below the reference  
173 transport rate, we estimate the point of intersection by fitting a surface-based transport relation  
174  $W* = 0.002(\tau/\tau_r)^{14.2}$  (Parker, 1990). All told we have examined sediment transport data for 132  
175 sites and retained 68 of them for the following analysis. Of the 68 sites, one site (the Mameyes  
176 River) used tracer particles to estimate the threshold of motion (see Phillips and Jerolmack,  
177 2014). The 64 excluded sites were not retained for a variety of reasons, the most common being  
178 missing data, insufficient sample size, and conflicting parameter estimates. In some cases, field  
179 sites were excluded if we could not estimate a reliable reference threshold due to the absence of a  
180 trend within the flux measurements. This occurs in several channels with large sand fractions and  
181 for sites where particles approaching the size of the stream bed  $D_{50}$  were never mobilized. Both  
182 of these issues result in no relation between Shields stress and flux for the range of stresses  
183 reported. For a broader discussion of potential problems with these and similar data see Recking  
184 (2010; 2013).

185

## 186 **2.2 Estimating the Channel Bankfull Transport Capacity**

187 The bankfull transport capacity ( $\tau_{*bf}/\tau_{*r}$ ) is estimated for the threshold data using equation (1) for  
188  $\tau_{*bf}$ , while  $\tau_{*r}$  is determined from bed-load flux measurements. In other words,  $\tau_{*bf}$  is determined  
189 from channel morphology, and hence independently from  $\tau_{*r}$ . In a few cases, our estimates of the  
190 variables required to compute  $\tau_{*bf}$  and  $\tau_{*r}$  based on the raw data differ from the values reported in  
191 the original studies. These differences are generally small, and we use the reported values in  
192 deference to the previous authors. We have only used sites that characterize the bankfull depth as  
193 a morphologic break between the active channel and a flood plain or non-channelized area (refer  
194 to Williams (1978) for additional common metrics of bankfull depth), as expansive flood plains  
195 are not always evident in mountain channels. For several field sites with irregular cross sections  
196 or large gravel bars we estimated the bankfull depth as the average depth across the active  
197 channel, which excludes the banks and large gravel bars from the height estimates. For these  
198 sites, the hydraulic radius becomes increasingly skewed by the bar and no longer reasonably  
199 approximates the stress imposed on the stream bank adjacent to the flow. Overall this results in a  
200 slightly higher estimate of the bankfull depth for some reaches, but matches the hydraulic radius  
201 calculations for more uniform or rectangular cross sections. We have excluded field sites where

202 we cannot reliably estimate both  $\tau_{*r}$  and  $\tau_{*bf}$ ; difficulties in calculating the latter are due to  
 203 insufficient data to calculate  $H$ ,  $S$ , and the  $D_{50}$ . In some cases, we have excluded field sites where  
 204 there are no raw data available and we are not able to acquire reliable estimates of these  
 205 parameters from other sources. Additionally, we have excluded several field sites where multiple  
 206 sources report conflicting estimates of the same hydraulic geometry parameters and we cannot  
 207 determine which ones are representative of the channel in question due to the absence of the raw  
 208 data.

209

210 **3. Results**211 **3.1 Correlation and Covariation between the Reference Shields Stress and Slope**

212 The slope of a mountain river reach is, all things considered, one of the more reliable parameters  
 213 that can be measured. It is one of the few parameters that can be directly measured from aerial  
 214 lidar and satellite derived topographic digital elevation models, whereas  $\tau_{*r}$  is not an easy  
 215 parameter to measure. The correlation between  $\tau_{*r}$  and  $S$  (Figure 1a) previously observed within  
 216 a compilation of mountain rivers (Mueller et al., 2005; Lamb et al., 2008) represented a  
 217 pragmatic path towards estimating a key parameter, even though it lacks the physical basis laid  
 218 out in the mechanistic models (see Prancevic and Lamb (2015b) for a field application of a  
 219 mechanistic model). The correlation between  $S$  and  $\tau_{*r}$  has been reported as following both linear  
 220 and non-linear relations (see Lamb et al., 2008; Mueller et al., 2005; Pitlick et al., 2008; Recking,  
 221 2009). We find that the best fit, in a least-squares sense, is a non-linear relation of the form

$$222 \quad (3) \quad \tau_{*r} = kS^\alpha,$$

223 where  $k=0.27$  and  $\alpha=0.38$  (standard error of 0.045) represent the best fit coefficient and exponent  
 224 ( $R^2=0.53$ ) for these data, respectively (Figure 1a). This fitted equation is similar to that  
 225 determined by Pitlick et al. (2008) and possesses a slightly steeper slope than that of Lamb et al.  
 226 (2008). The differences among the empirical regressions are minor, and arise due to (i) the  
 227 addition of new field sites present here and (ii) the exclusion of laboratory data from the fit in  
 228 equation (3). Throughout the rest of the manuscript we will utilize equation (3) when providing  
 229 examples of  $\tau_{*r}$  as estimated via an empirical regression, though the issues raised in the following  
 230 sections are inherent to all such regressions on these or subsets of these data. The non-linear  
 231 relation provides a better fit to the data, however a linear fit is also reasonable. We also observe  
 232 that  $\tau_{*bf}$  possesses a similar non-linear correlation with slope (also observed by Mueller et al.,

233 2005 and Pitlick et al., 2008). The best-fit regression line between  $\tau_{bf}$  and  $S$  is vertically offset of  
 234 that between  $\tau_r$  and  $S$ , indicating that  $\tau_{bf}$  is slightly larger than  $\tau_r$  in a consistent manner  
 235 ( $\tau_{bf}=0.29S^{0.35}$ ,  $R^2=0.48$  with a standard error of 0.046 for  $\alpha$ , Figure 1a). The compilation data  
 236 represent a larger range of  $\tau_{bf}$  and a steeper trend with  $S$  ( $\tau_{bf}=0.55S^{0.45}$ ,  $R^2=0.28$  with a standard  
 237 error of 0.025 for  $\alpha$ ).

238

239 To understand how the ratio of  $\tau_{bf}/\tau_r$  computed with measured values compares with the same  
 240 ratio where  $\tau_r$  is calculated from the empirical regression we combine equations (1) and (3):

$$241 \quad (4) \quad \frac{\tau_{bf}}{\tau_r} \sim \frac{\tau_{bf}}{kS^\alpha} = \frac{hs}{kS^\alpha RD_{50}} = \frac{hs^{1-\alpha}}{kRD_{50}}.$$

242 The primary effect of estimating  $\tau_{bf}/\tau_r$  with equation (4) is to increase the contribution of slope  
 243 and minimize the role of the median particle size (for  $k<1$ ) within the Shields stress. Recall that  
 244 since the slope is always less than one, as the quantity  $(1-\alpha)$  approaches zero the slope term  
 245 approaches a value of one. Using equation (4) the computed average transport capacity is  $<\tau_{bf}/kS^\alpha>=1.26$  (arithmetic mean of 1.4), which is close to the observed value of  $<\tau_{bf}/\tau_r>=1.27$  ( $<\mathbf{>}$  denotes a geometric mean of the bracketed quantity). However, the shape and variance of the  
 247 distributions are poor matches (Figure 1b), even though equation (4) was fit to these data. Note  
 248 that equation (4) produces transport capacities that extend well below one, which is a non-  
 249 physical result as transport was observed in these rivers for sub bankfull flow. The same exercise  
 250 can be performed for the compilation data using equation (4), which produces an even larger  
 251 variance (gray PDF in Figure 1b). The larger variance is a result of not capturing the covariation  
 252 present in  $\tau_r$  and  $\tau_{bf}$ , as each set of observations are pairs. Where  $\tau_r$  for a given site is  
 253 larger/smaller relative to equation (3),  $\tau_{bf}$  is also larger/smaller for the same site. The covariation  
 254 between  $\tau_{bf}$  and measured  $\tau_r$  is so strong that no trend is apparent when we take the ratio of  
 255  $\tau_{bf}/\tau_r$  and compare it to  $S$ ,  $H/D_{50}$ , and even  $\tau_{bf}$  (Figure 1c). As the required variables to compute  
 256  $\tau_{bf}$  are commonly available or easier to measure, equations (3) and (4) – or variants thereof – are  
 257 used frequently to estimate  $\tau_r$ , despite the fact that the covariation between  $\tau_{bf}$  and  $\tau_r$  produces  
 258 a spurious result for the transport capacity. Through the use of equation (4) the estimated  
 259 bankfull transport capacity for both datasets become positively correlated with relative  
 260 submergence and  $\tau_{bf}$ , even though the actual value of  $\tau_{bf}/\tau_r$  is constant (Figure 1c). These  
 261 correlations are spurious due to the variables' underlying correlation with slope (Figure 1c). The

263 same degree of correlation is not observed between slope and equation (4), because the relation  
264 between  $H/D_{50}$  almost precisely cancels out the effect of plotting slope against itself (Figure 1c).  
265

### 266 **3.2 Parameter Space of the Threshold Field Sites**

267 The second significant issue when computing  $\tau_{*r}$  with equation (3) results from the partial range  
268 of parameter space covered by the field sites from which the regression is based. The broader  
269 compilation of field sites represents a more diverse set of gravel rivers and provides context for  
270 the range and frequency of the different hydraulic variables ( $H$ ,  $S$ ,  $D_{50}$ ,  $W$ , and  $Q$ ). In general,  
271 their probability density histograms do not appear to follow a normal distribution and are better  
272 represented by their natural logarithm (Harman et al., 2008). Thus, we natural-log transform the  
273 hydraulic variables for the threshold and compilation sites prior to computing the histograms.  
274 Compared to the compilation dataset, the threshold site channels are generally shallower,  
275 narrower and steeper, with lower bankfull discharges and coarser beds (Figure 2a-f). In other  
276 words, sites used for the threshold regressions are a non-representative subset of the larger  
277 compilation of gravel rivers. In terms of  $\tau_{*bf}$ , however, these sites match the central tendency of  
278 the larger compilation well (Figure 2f); but they under sample both high and low values of  $\tau_{*bf}$   
279 relative to the larger compilation. The  $\tau_{*bf}$  of the compilation appears to be well described by a  
280 log-normal distribution yielding a geometric average of  $\langle \tau_{*bf} \rangle = 0.054$  for single thread coarse-  
281 grained rivers (Figure 2f). A two-tailed Kolmogorov-Smirnov test for a log-normal distribution  
282 ( $K-S_{\text{stat}}=0.045$  with a  $P_{\text{value}}=0.097$ ) suggests that we cannot reject the null hypothesis that  $\tau_{*bf}$  is  
283 log-normally distributed (i.e. we accept that  $\tau_{*bf}$  does not differ from a log-normal distribution at  
284 a  $P_{\text{value}}$  of  $\sim 0.1$ ). Additionally, the threshold distribution of  $\tau_{*bf}$  also follows a log-normal  
285 distribution ( $K-S_{\text{stat}}=0.095$  with a  $P_{\text{value}}=0.55$ ), though not the same distribution as the  
286 compilation (Two sample KS-test,  $K-S_{\text{stat}}=0.17$  with  $P_{\text{value}}=0.042$ ). It is perhaps not surprising  
287 that the threshold sites are non-representative of the larger compilation, as these sites are by-and-  
288 large geographically biased to mountain rivers primarily within the Rocky Mountains and the  
289 states of Colorado, Wyoming, and Idaho whereas the compilation samples a broader geographic  
290 range (continental United States and Canada).

291  
292 The consequences of the geographic bias in measured  $\tau_{*r}$  and equation (3) can be further  
293 explored by removing the self-correlation (slope occurs in both axis) present within Figure (1).

294 We can examine the relationships between both datasets by exploring the parameter space  
295 between  $H/D_{50}$  and  $S$ , the two free dimensionless variables within Shields stress while holding  
296 the quantity  $\rho/(\rho_s-\rho)$  constant (Figure 3). When viewed this way the compilation data form a  
297 scattered cloud in which bankfull relative submergence ( $H/D_{50}$ ) trends inversely with slope.  
298 These data show that for the same value of  $\tau_{bf}^*$  there exist low gradient rivers with high relative  
299 submergence, and steep gradient rivers with low relative submergence. Particularly, these data  
300 show that there is quite a range in  $\tau_{bf}^*$ , and that high Shields stresses are not solely the domain of  
301 steep rivers. Within this parameter space the threshold sites tend to be overly representative of  
302 the steepest rivers (Figure 3). There are few if any threshold field sites occupying the region of  
303 the parameter space characterized by low slope, high relative submergence rivers. Within these  
304 data a pattern with  $\tau_r^*$  emerges showing an additional dependence on  $H/D$  (also highlighted by:  
305 Mueller et al., 2005; Recking 2009). Equation (3) runs askew to the primary trend of the  
306 compilation data, and doesn't capture the overall pattern of  $\tau_r^*$  within this parameter space.  
307 Interestingly, the pattern in  $\tau_r^*$  follows isolines of increasing Shields stress, in that steep rivers  
308 with low relative submergence appear to have the same value of  $\tau_r^*$  as low gradient rivers with  
309 high relative submergence (Figure 3). For example, using equation (3) to estimate  $\tau_r^*$  at any  
310 value of  $S$  would indicate that as  $H/D_{50}$  increases so too does the transport capacity of the river.  
311 However, the measured values of  $\tau_{bf}^*/\tau_r^*$  do not vary systematically with  $H/D_{50}$  or  $S$ . From these  
312 data it becomes apparent that a third variable, which combines both  $H/D$  and  $S$ , may be a better  
313 predictor of the observed  $\tau_r^*$  pattern. However, it is not necessary to fit such a regression as the  
314 third variable is the bankfull Shields stress (equation 1).

315

### 316 **3.3 Relation between the Bankfull and Reference Shields Stresses**

317 The observed correlation between  $\tau_{bf}^*$  and  $\tau_r^*$  is a strong linear trend (Figure 4). It is important to  
318 note that the methodology that calculates  $\tau_r^*$  is independent of that used to determine  $\tau_{bf}^*$ , as  $\tau_r^*$  is  
319 determined from a range of flow and flux measurements while  $\tau_{bf}^*$  is determined from channel  
320 geometry. The correlation between  $\tau_{bf}^*$  and  $\tau_r^*$ , and by extension  $\tau_c^*$ , was previously shown  
321 (Mueller et al., 2005) to be linear with a subset of the threshold data used here. Mueller et al.  
322 (2005) concluded that this trend indicated that gravel-bedded streams were adjusted to have a  
323 constant bankfull excess Shields stress ( $\tau_{bf}^* - \tau_c^*$ ). The analytic model for the equilibrium channel  
324 geometry of gravel-bedded rivers developed by Parker (1978) provides an explanation for the

325 observed correlation and an expected functional form of  $\tau_{bf}^* = (1+e)\tau_c^*$ , where ‘ $e$ ’ is a small  
 326 positive value. The prediction for a channel with cohesionless unimodal sediment provides that  
 327  $e=0.2$  yielding a predicted relation of  $\tau_{bf}^* = 1.2\tau_c^*$  for a specified value of  $\tau_c^*$ . The best fit relation  
 328 of this form from the data is  $\tau_{bf}^* = 1.19\tau_r^*$  ( $R^2=0.96$ ), which is remarkably close to the analytical  
 329 prediction (Figure 4). Combining these ideas together provides an avenue to predict  $\tau_r^*$  from  $\tau_{bf}^*$   
 330 by rearranging the confirmed analytical equation to yield  $\tau_r^* = 0.83\tau_{bf}^*$ . From this relation we can  
 331 estimate the residuals and compare estimates with equation (3). Histograms of the residuals show  
 332 a positively skewed distribution from equation (3) and a mostly symmetric and narrower  
 333 distribution for the analytical prediction (Figure 4 inset).

334

### 335 **3.4 Illustration of Perceived Differences in Transport Capacity via Subsampling**

336 To illustrate problems with subsampling the parameter space of gravel bed rivers while using  
 337 equation (4) to compute the transport capacity, we created a set of contrived subsamples from the  
 338 larger compilation dataset based on  $S$ ,  $H/D_{50}$ ,  $\tau_{bf}^*$ ,  $Q/W$ , and  $f$  (flow resistance). These  
 339 subsamples are similar to how a researcher might collect field data or values from the literature  
 340 to compute  $\tau_{bf}^*$ , and through equation (4) estimate the bankfull transport capacity to compare  
 341 different regions or catchments. It is important to note that selecting field sites within a particular  
 342 geographic region is something that is commonly done and acceptable practice, but yields a  
 343 selection of rivers with a limited range in values of  $S$ ,  $H/D_{50}$ ,  $\tau_{bf}^*$ ,  $Q/W$ , and  $f$ . Throughout the  
 344 following exercise we show how application of the slope-based regressions (equation 3) to such  
 345 data can yield erroneous conclusions. The expected distribution based on the measured  $\tau_r^*$  is  
 346 narrowly distributed around the theoretical prediction (Figure 1b & reproduced in Figure 5a). To  
 347 facilitate a more direct comparison with the threshold field sites ( $n=68$ ), we selected 70 random  
 348 field sites from the full parameter space of the larger compilation (Figure 5a). We chose 70  
 349 random samples for the contrived sampling scheme as this number is close to the number of  
 350 threshold sites, represents a low relative standard error (< 5%), and produces a reasonable  
 351 distribution. The error in computing the mean is relatively low even for a small set of random  
 352 samples (Figure 5a), because the natural log-transformed distribution of  $\tau_{bf}^*$  is a normal  
 353 distribution. The contrived sampling schemes were created by randomly selecting two sets of 70  
 354 sites from the compilation dataset from above and below the geometric mean for  $S$ ,  $H/D_{50}$ ,  $\tau_{bf}^*$ ,  
 355 and  $Q/W$  (Figure 5b), while sampling criteria for  $f$  was based on the arithmetic mean value. Flow

356 resistance was computed using the Variable Power Equation  $(8/f)^{1/2} = a_1 a_2 (H/D) / [a_1^2 +$   
357  $a_2^2 (H/D)^{5/3}]^{1/2}$  with coefficients  $a_1=7.3$  and  $a_2=2.3$  (Ferguson, 2007), as this equation was  
358 previously demonstrated by Ferguson (2007) to match a large compilation of field data well. The  
359 final category is a combination of  $H/D$  and  $S$  that samples from opposite corners of the  
360 compilation parameter space in Figure 3. The differences in transport capacity through the use of  
361 equation (4) between the two sets of subsamples for each criterion are illustrated in Figure 5b  
362 (see Figure 5c to see the selected sites within the relative submergence-slope parameter space).  
363 Some of the contrived subsamples have similar median values for transport capacity to the  
364 measured threshold sites; however, all subsamples have substantially larger inner quartile ranges  
365 and standard deviations (Figure 5b). Whether a subsample differs from its partner sample is  
366 completely dependent on how the selected sites relate to where equation (3) crosses the  
367 parameter space (Figure 5c). Subsamples showing little difference from each other are those  
368 based on  $S$  and  $Q/W$ , while the rest of the subsamples ( $H/D_{50}$ ,  $\tau_{bf}$ ,  $f$ ,  $H/D_{50}$  &  $S$ ) would indicate  
369 that transport capacity differs for these sets of gravel rivers. These differences, while statistically  
370 significant, are artefacts of the bias that arises by sampling a limited range of the parameter space  
371 relative to equation (3).

372

#### 373 **4. Discussion**

374 Here we start by discussing the quality and bias issues of the threshold data set, as several types  
375 of errors are potentially present. The majority of these issues are likely a consequence of the  
376 complicated nature of measuring sediment transport and channel hydraulic parameters, which  
377 represent snapshots of a dynamic system. For  $\tau_{bf}$  the potential sources of error are in determining  
378  $H$ ,  $D_{50}$ , and  $S$  at each field site. The largest source of error for this study is related to defining the  
379 bankfull depth (Williams, 1978), because errors in both  $S$  and  $D_{50}$  are less likely to affect the  
380 transport capacity as both of these parameters are part of the calculations necessary to compute  $\tau^*$   
381 and  $\tau_{*r}$  through equation (2). The error in measuring the bankfull depth is relatively low and  
382 decreases with the number of cross sections (Harman et al., 2008), though it remains an open  
383 question as to the minimum number of cross sections required to achieve a representative  
384 average bankfull depth, and exactly how to treat the bankfull depth (hydraulic radius or average  
385 active channel depth) for irregular or complicated cross sections. Our intent is to understand the  
386 ratio  $\tau_{bf}/\tau_{*r}$  as it relates to channel stability, therefore we have chosen to calculate  $H$  using the

387 average active channel depth for irregular cross sections and the hydraulic radius where the two  
388 metrics closely agree. A full accounting of this problem is not possible given the current datasets,  
389 as most field sites have no more than three cross sections from which to compute  $H$ ; however,  
390 the difference between both methods is small, at least for the field sites we have examined here.  
391 Defining  $H$  as a morphologic break in the channel cross section, we were able to independently  
392 reproduce the bankfull depths reported by the original authors of the studies from which the  
393 threshold data are compiled. Sources of error for  $\tau_{*r}$  are potentially more numerous as bed-load  
394 transport measurements are notoriously noisy data; for a thorough analysis and discussion of the  
395 potential sources of error see Recking (2013). The largest areas of error for the threshold dataset  
396 are related to sampling bed-load transport at low transport rates and the choice of sampler used to  
397 collect the samples. Mobile samplers (e.g. Hellely-Smith) measure higher flux rates for low  
398 transport conditions, compared to pit and trap type samplers (Bunte et al., 2008). This  
399 oversampling can result in flat (trendless) relations between flux and stress at low transport rates  
400 and may have resulted in the exclusion of several field sites where higher transport rates were not  
401 available to distinguish a trend. A larger concern with these data is in how representative a single  
402 measurement of  $\tau_{*r}$  is of the threshold, as both the spatial and temporal variability of  $\tau_{*r}$  remains  
403 uncertain. The temporal variability, however, may be less worrisome long term as it appears to  
404 be normally distributed where it has been measured (Masteller et al., 2019). Suffice to say,  
405 understanding the dynamics of the threshold of motion remains an area in need of additional  
406 research. Therefore, we caution the reader from focusing on a single field site or exact numerical  
407 values, and instead recommend that the overall trends are more robust.

408  
409 In terms of data bias of the compilation parameter space coverage, we can only speculate given  
410 the available data as to how representative some of these parameters are when compared to the  
411 timescales of channel adjustment. In a sense, the reach-scale channel geometry integrates over  
412 some yet unknown number of flood events, or may even alternate between different states of  
413 adjustment (Pizzuto, 1994; Slater & Singer, 2013; Wolman & Gerson, 1978; Yu & Wolman,  
414 1987). It is not currently definitively known how much the bed composition changes over time  
415 and thus how representative a single grain size measurement is, or how sensitive natural channel  
416 geometry is to changes in bed composition (e.g. MacKenzie & Eaton, 2017). Similarly, it  
417 remains an open question if the sampled bed grain size distribution is reflective of the current

418 measured channel geometry. These are questions though that cannot necessarily be addressed  
419 with the current data compilations, but are worth keeping in mind when considering rivers as  
420 dynamic systems. In addition, it cannot be definitively concluded that the compilation dataset  
421 fully represents the spectrum of  $\tau_{bf}^*$  in coarse-grained rivers, as the compiled field sites  
422 necessarily represent the site selection criteria within the original studies. For example, the  
423 sampling of field sites is strongly biased geographically towards North America. In particular,  
424 field measurements in dryland, arctic or periglacial, and tropical environments are notably  
425 lacking. However, given the close fit to a log-normal distribution (Figure 2f) it is not clear to the  
426 authors that more globally representative sampling would not simply make any fit better.

427

428 Regardless of how representative the compilation dataset is of global rivers, a significant pitfall  
429 of using equation (3) to predict  $\tau_r^*$  is that the field sites on which equation (3) is based do not  
430 sample the full parameter space of the compilation of coarse-grained rivers used in this study  
431 (Figure 2 and 3). We note that this is the largest compilation of coarse-grained rivers to date.  
432 This uneven sampling is strongly biased towards moderate to high slopes and low relative  
433 submergence rivers. Noticeably under sampled are rivers with higher relative submergence for  
434 all slopes ( $H/D > 20$ ), and rivers with lower slopes ( $S < 0.002$ ) in general. Caution should be  
435 exercised when attempting to extrapolate predictions for  $\tau_r^*$  to regions of the parameter space that  
436 are not sampled, especially for sites with higher values of  $H/D$  relative to equation (3) (Figure 3).  
437 Capturing the covariation between  $\tau_r^*$  and  $\tau_{bf}^*$  is especially important for estimating bed load  
438 transport due to flux equations' non-linear dependence on transport stage ( $\tau^*/\tau_r^*$ ) and/or excess  
439 shields stress ( $\tau^* - \tau_r^*$ ) (Mueller et al., 2005). For example, the difference between  $\tau_{bf}^*/\tau_r^* = 1.2$  and  
440  $\tau_{bf}^*/\tau_r^* = 2$ , seems small given the variation in the data; yet when viewed through a common bed-  
441 load transport equation (see Wilcock & Crowe, 2003) this becomes a factor of ~20 in terms of  
442 flux and grows non-linearly with increasing values of  $\tau^*/\tau_r^*$ . As there are very few estimates of  
443  $\tau_r^*$  in high Shields stress ( $\tau^* > 0.12$ ) regions of the parameter space, and especially low slope and  
444 high relative submergence sites, further research is still required to determine the range of  $\tau^*/\tau_r^*$   
445 within these regions. The available data do not, however, support the use of equations (3) and (4)  
446 to determine  $\tau_r^*$  in these regions. These equations would predict that bankfull transport capacity,  
447 and hence bed-load flux, increases with  $\tau_{bf}^*$  – despite the available measured data indicating that  
448  $\tau_{bf}^*/\tau_r^* \sim \text{constant}$  (Figure 1 inset and Figure 4). Interestingly, while the correlation between slope

449 and high values of  $\tau^*_c$  is becoming increasingly well understood (see Lamb et al., 2008; Recking,  
450 2009; Prancevic and Lamb, 2015a), to date the explanation for the increase of  $\tau^*_r$  with  $H/D$  for  
451 regions of low slope remains uncertain to the authors.

452

453 The second issue with the approach laid out in equation (4) is that it does not capture the co-  
454 variation between  $\tau^*_b$  and  $\tau^*_r$  at each site. These two parameters are ‘paired’ in a sense, and using  
455 either fitted regression relation (see figure 1) to estimate  $\tau^*_{bf}$  or  $\tau^*_r$  from slope alone will result in  
456 an incorrect prediction for the ratio  $\tau^*_{bf}/\tau^*_r$ . This pairing of the data is evident when considering  
457 the relation between  $\tau^*_{bf}$  and  $\tau^*_r$  (Figure 4), which closely matches theory (Parker, 1978)  
458 indicating support for a causative relation. While there is some deviation from this trend within  
459 the data (Figure 4) the residuals possess no meaningful correlation with the available hydraulic  
460 variables ( $H$ ,  $W$ ,  $Q$ ,  $D50$ ,  $H/D$ ,  $f$ , and  $S$ ). It remains unclear to the authors if the degree of scatter  
461 in  $\tau^*_{bf}/\tau^*_r$  reflects actual ranges of channel behavior, or represents a combination of error or bias  
462 in the measurements and under-sampling of the various hydraulic parameters or random noise in  
463 a dynamic system. Of the explored relations and correlations, the simple linear relation  $\tau^*_{bf}=1.2\tau^*_r$   
464 remains the best predictor to date. In a sense, this model can serve as a null hypothesis. Absent  
465 independent measurements of  $\tau^*_c$ , this null model states that  $\tau^*_{bf}=1.2\tau^*_c$  or  $\tau^*_c=0.83\tau^*_{bf}$ . With this  
466 in mind and the observation that  $\tau^*_{bf}$  is approximately log-normally distributed ( $\langle\tau^*_{bf}\rangle=0.054$ ),  
467 then  $\tau^*_c$  would also be log-normally distributed and we can estimate its mean as  $\langle\tau^*_c\rangle=0.045$   
468 (close to the prediction of the Shields curve). The value in this approximation is that it places  
469 statistical bounds on the extent of both  $\tau^*_{bf}$  and  $\tau^*_c$  in natural channels. We have demonstrated  
470 several potential outcomes of this sampling bias through a variety of field site selection criteria  
471 (Figure 5b). Though the demonstrated selection criteria were strictly related to simple statistical  
472 splits of hydraulic variables, the discussed bias applies equally to samples drawn based on  
473 geography or regional climate. The upshot here is that an empirical slope-based predictor is not  
474 broadly reliable for estimating the threshold of motion or predicting the transport capacity,  
475 because the data it is based on represent a biased sampling of alluvial gravel rivers and does not  
476 account for the covariation between  $\tau^*_{bf}$  and  $\tau^*_r$ . We advocate that  $\tau^*_r$  can be estimated from  $\tau^*_{bf}$   
477 due to the remarkable consistency observed in channel geometry and the close connection to  
478 theory which provides a physical basis for the prediction. Using channel geometry to predict the  
479 threshold of motion provides an implicitly time and space averaged reach-scale value for  $\tau^*_c$ ,

480 however the extent of the averaging will depend on the number of flows responsible for shaping  
481 the current channel geometry.

482

483 Lacking an alternative easily implementable approach to predicting  $\tau^*_{*c}$ , we recommend using the  
484 null model ( $\tau^*_{bf}/\tau^*_{*c}=1.2$ ) in theory, numerical, and analytical based approaches, while for strictly  
485 empirical approaches error can be incorporated through the observed distribution and standard  
486 deviation ( $\langle\tau^*_{bf}/\tau^*_{*r}\rangle=1.27 \times / \div 1.16$ ). The null model provides a closure for studies in gravel-bed  
487 rivers where  $\tau^*_{*c}$  needs to be estimated. For example, predicting spatial patterns of grain size and  
488 morphology for in stream management and habitat suitability (Phillips & Scatena, 2013; Snyder  
489 et al., 2013), and management of river corridors below major river modifications (Minear, 2010;  
490 Schmidt & Wilcock, 2008). For field sites where the identification of the bankfull depth is  
491 difficult to assess we recommend checking the data against the compilation dataset parameter  
492 space (Figure 3) for a variety of flow depths to assess a likely range. This approach is also  
493 insightful to assess potential bias for geographically based sampling. For steeper river channels  
494 ( $S > 0.01$  and  $H/D < 10$ ) where one absolutely requires an approximation of  $\tau^*_{*r}$  based on metrics  
495 extracted from topography alone, the slope-based regression in equation (3) and other published  
496 variants may be suitable with the former caveats in mind. We recommend the field based  
497 empirical regressions over those that incorporate laboratory measurements, because the field  
498 based regressions capture processes (and errors) inherent to the field that are absent in the lab  
499 such as the difference between measured particle size and mass in steep channels (see Miller et  
500 al., 2014). We do not have a recommendation for bedrock rivers, as the concept of bankfull does  
501 not always exist and a bankfull analogue with morphological significance has yet to be fully  
502 established. Research in this area is still developing, however field and flume experiments  
503 (Johnson et al., 2009; Johnson & Whipple, 2010) have demonstrated that bedrock rivers appear  
504 to adjust towards a condition of steady state to pass the sediment flux and water discharge  
505 supplied. This is similar to the statistical concept of the ‘effective flood’ in alluvial rivers which  
506 is a bankfull analogue and represents the average stress above the threshold of motion (Phillips  
507 & Jerolmack, 2016; Torizzo & Pitlick, 2004; Wolman & Miller, 1960).

508

509 **5. Conclusion**

510 Through the use of an expanded dataset, we demonstrate that empirical regressions based on the  
511 correlation between slope and the threshold of motion can easily result in erroneous conclusions  
512 when considering the channel's transport capacity. This occurs because the empirical regressions  
513 are based on a limited sampling of the parameter space of bed-load rivers and do not capture the  
514 covariation between  $\tau^{*bf}$  and  $\tau^{*r}$ . Predicting the threshold of motion in natural channels remains a  
515 considerable challenge, and a critical knowledge gap for understanding rivers' roles within their  
516 catchments. We recommend site-specific empirical determination of the threshold from  
517 independent measurements of bed-load transport; while this is challenging, the rapid uptake of  
518 seismic, acoustic, and other methods makes this prospect increasingly more feasible. Barring  
519 direct measurement, the threshold of motion's strong correlation with the bankfull Shields stress  
520 remains the most accurate predictor. The covariation of threshold and bankfull Shields stresses is  
521 a consequence of the organization of bed-load rivers to be close to the threshold of motion. The  
522 average bankfull transport capacity determined from available data is in remarkable agreement  
523 with the prevailing theory for gravel-bed river geometry (Parker, 1978). Observed deviation from  
524 this theory when using these empirical regressions is an artefact of the sampling bias inherent in  
525 their construction; future studies will need to address this bias rigorously with independent  
526 observations of the threshold in order to disprove the null hypothesis.

527

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536

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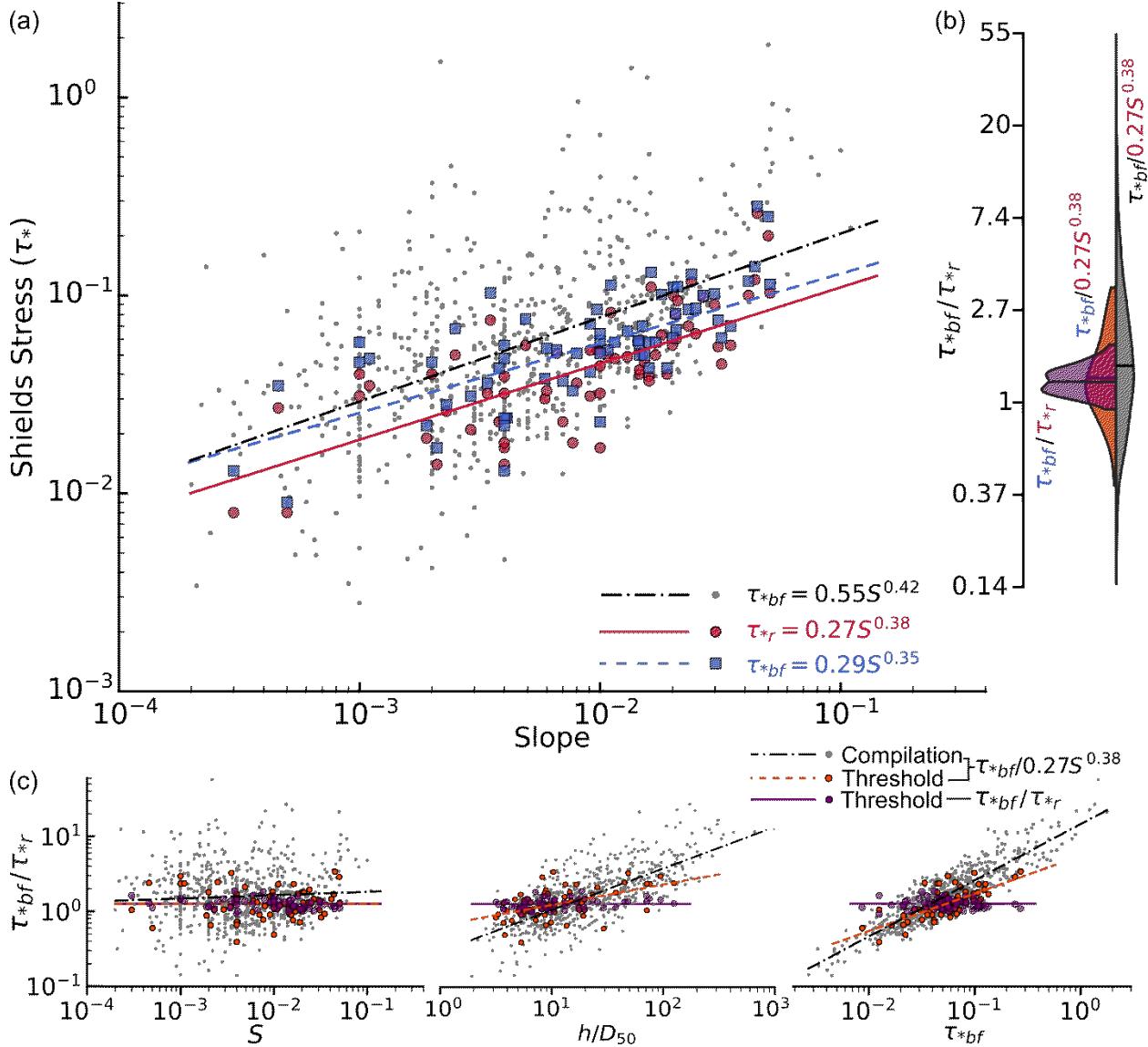
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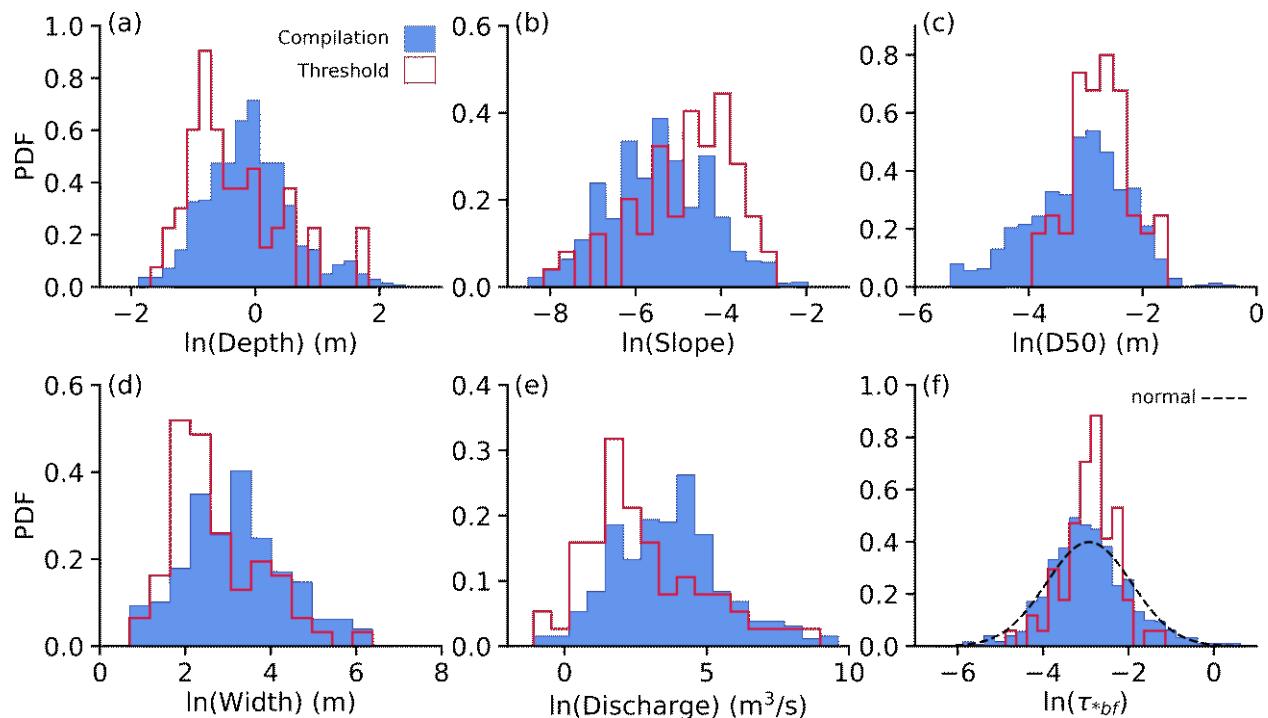
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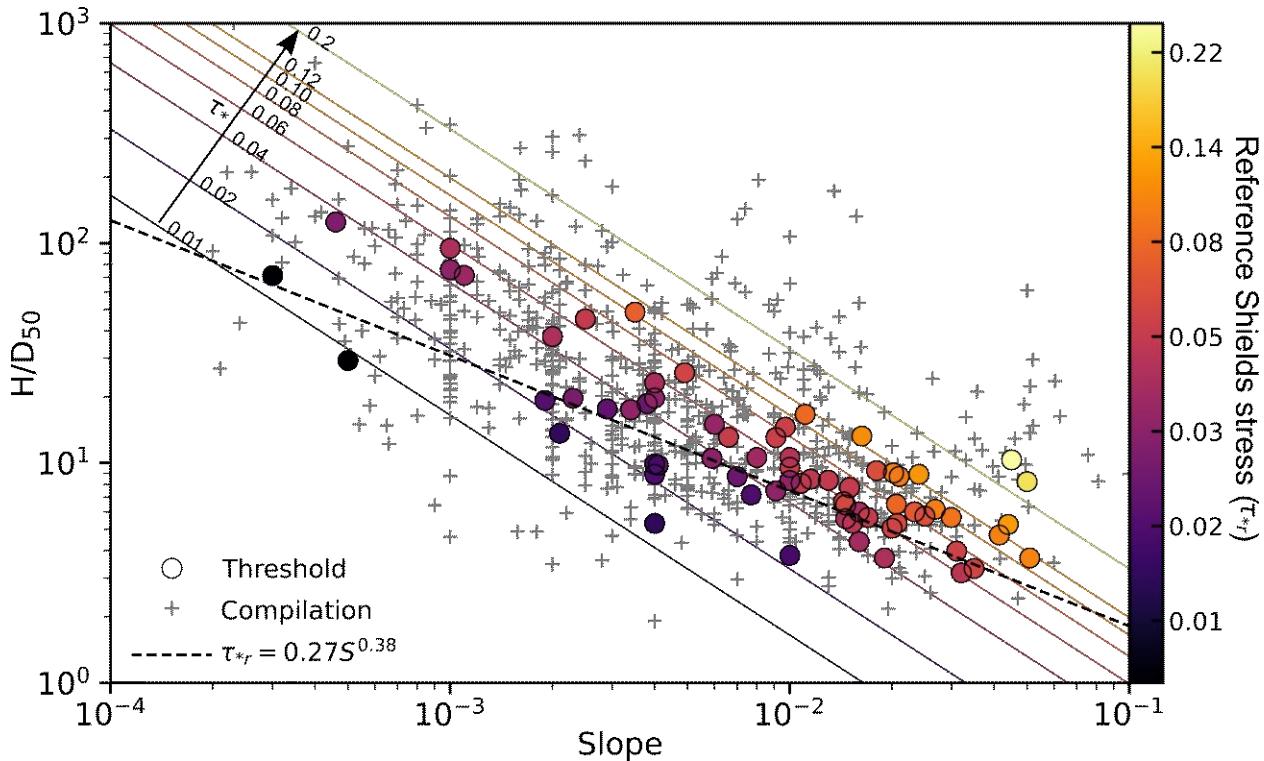
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 852 Figure 1. Correlation between slope, bankfull Shields stress, and the threshold of motion. (a)  
 853 Observed correlation for the threshold data between reach-scale slope, and the reference ( $\tau_{*r}$ , red  
 854 circles) and bankfull Shields stresses ( $\tau_{*bf}$ , blue squares). Red (solid) and blue (dashed) lines  
 855 represent loglog least-squares regressions excluding two outliers ( $\tau_{*r}>0.2$ ). Gray points represent  
 856  $\tau_{*bf}$  for the compilation data with a least-squares regression line (black dash-dot line). (b) Split  
 857 violin plot of the distributions of the bankfull transport capacity ( $\tau_{*bf}/\tau_{*r}$ ) where  $\tau_{*r}$  is estimated  
 858 from flux measurements (purple), and the slope-based regression ( $\tau_{*r}=0.27S^{0.38}$ ) for the threshold  
 859 (orange) and compilation data (gray). The solid line within the distribution represents the median  
 860 and the upper and lower edges of the distribution are clipped at the extents of the data. (c)  
 861 Relations between slope, bankfull relative submergence ( $H/D_{50}$ ), and  $\tau_{*bf}$  with  $\tau_{*bf}/\tau_{*r}$  for the

862 threshold (purple and orange circles) and compilation data (gray points). Horizontal lines  
863 represent the geometric mean ( $\langle \tau^{*bf}/\tau^{*r} \rangle = 1.27$ ) where there is no correlation between the data.  
864 Note that the trends observed for  $H/D_{50}$  and  $\tau^{*bf}$  result from spurious correlation with slope  
865 within equation (4).  
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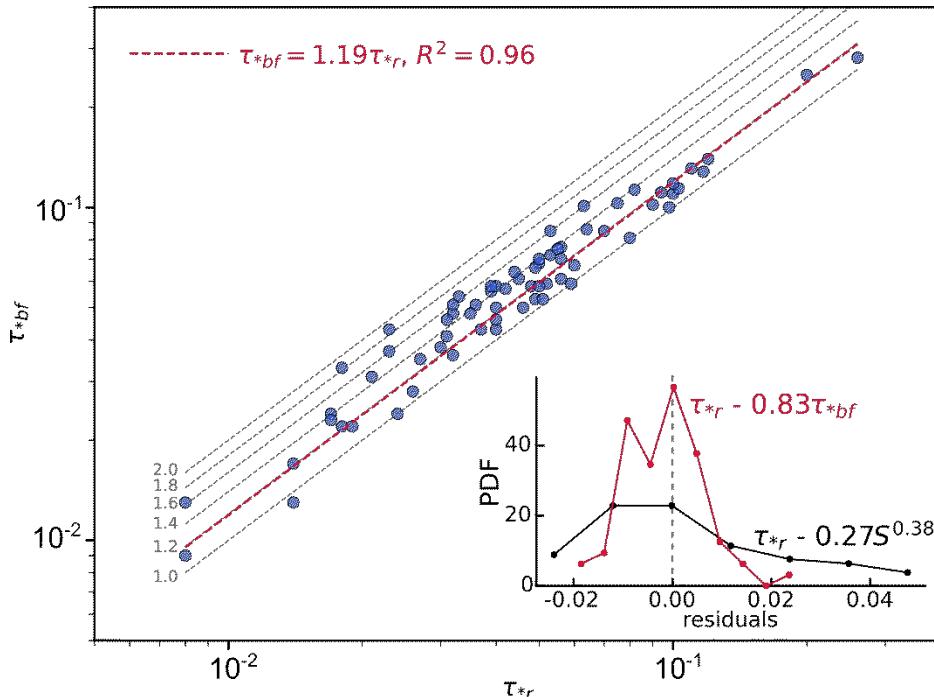
Figure 2. Probability density histograms for sites where  $\tau_{*r}$  was determined from bed-load flux measurements (red line,  $n=68$ ) compared with a larger data compilation (shaded blue,  $n=739$ ) of gravel-bedded rivers. All data were natural-log transformed prior to computing the histograms and bin width for both datasets used the Freedman-Diaconis rule (Freedman & Diaconis, 1981) based on the larger compilation. The variables are (a) bankfull depth ( $n=725$ ), (b) slope ( $n=739$ ), (c)  $D_{50}$  ( $n=739$ ), (d) bankfull width ( $n=272$ ), (e) bankfull discharge ( $n=418$ ), and (f) bankfull Shields stress ( $n=725$ ). Sample sizes vary according to data availability. The bankfull Shields stress (f) is well described by a normal distribution (black dashed line) in natural log space.



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878 Figure 3. Parameter space of bankfull relative submergence ( $H/D_{50}$ ) and slope for sites with  
 879 measured  $\tau_{*r}$  (shaded circles) and the larger river compilation (gray '+' symbol). The shaded  
 880 color and colorbar denote the measured reference Shields stress. The black dashed line represents  
 881 the best fit regression between slope and  $\tau_{*r}$  (equation 3). The multicolored diagonal lines are, by  
 882 definition, the Shields stress. Note that the shaded color pattern is parallel to the Shields stress  
 883 isolines (i.e. light orange points follow the orange isolines and the purple points follow the  
 884 purple lines) and not the regression line.

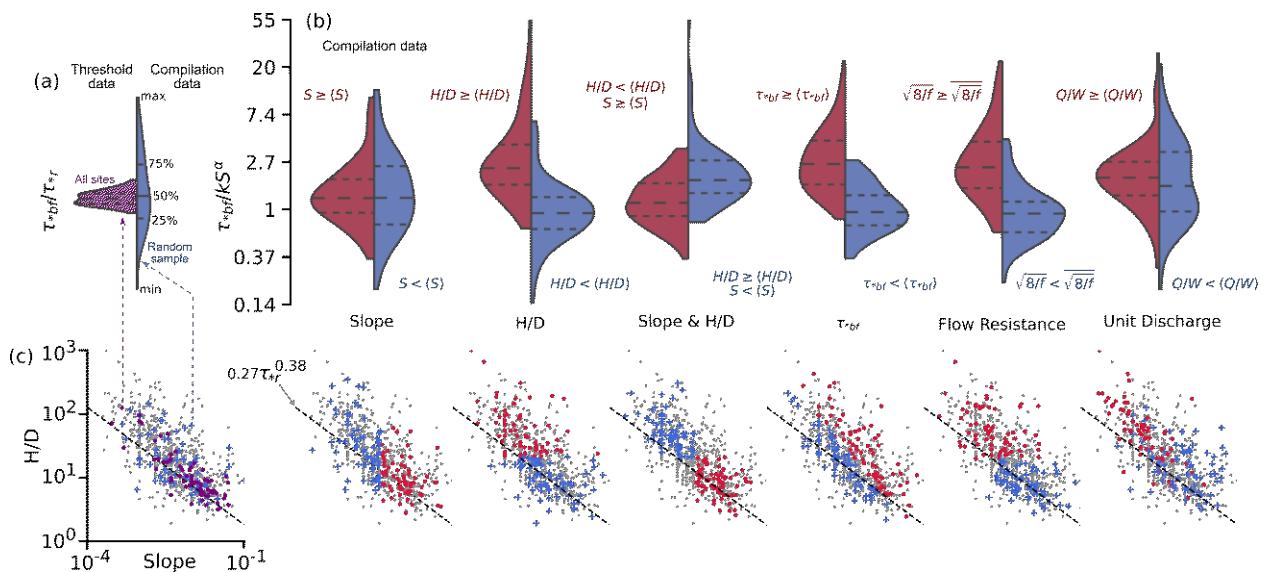
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Figure 4. Relation between the reference and the bankfull Shields stresses. The gray dashed lines represent increasing values for the ratio  $\tau_{*bf}/\tau_{*r}$  in increments of 0.2 for reference, while the red dashed line represents the best fit function of the form  $\tau_{*bf} = (1+e)\tau_{*r}$ . (inset) Residuals for estimating  $\tau_{*r}$  using the relation with  $\tau_{*bf}$  (red line) and equation (3) (black line).

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Figure 5. Illustration of bias in estimating transport capacity for various sampling strategies using equation (4). (a) Split violin showing the measured transport capacity for the threshold field sites (purple,  $n=68$ , identical to Figure 1b) and a random sample (blue,  $n=70$ ) from the compilation dataset where the transport capacity is calculated via equation (4). (a) and (b) share the same vertical axis. (b) Transport capacity calculated via equation (4) for various data sampling strategies from the compilation data. Each column represents sampling the larger compilation based on the variable listed below and each half represents 70 randomly selected field sites for the adjacently labeled condition. All data except flow resistance are natural log transformed prior to computing the distributions. (c) Illustration of the random samples used to compute the distributions in (a) and (b) from the larger compilation (small gray dots). All columns in (c) share the same axes. Blue crosses correspond to the right half and red dots represent the left half of the split violins in (b) directly above each data cloud. The black dashed line is equation (3). Note the distributions illustrate how one can observe a potential difference in transport capacity between gravel-bedded rivers based on how the samples relate to equation (3). The observed difference is spurious due to the selection variable's underlying correlation with slope.