

# Light-dependent Impedance Spectra and Transient Photoconductivity in a Ruddlesden–Popper 2D Lead-halide Perovskite Revealed by Electrical Scanned Probe Microscopy and Accompanying Theory

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## Abstract

Electric force microscopy was used to record the light-dependent impedance spectrum and the probe transient photoconductivity of a film of butylammonium lead iodide,  $\text{BA}_2\text{PbI}_4$ , a 2D Ruddlesden–Popper perovskite semiconductor. The impedance spectrum of  $\text{BA}_2\text{PbI}_4$  showed modest changes as the illumination intensity was varied up to  $1400 \text{ mW cm}^{-2}$ , in contrast with the comparatively dramatic changes seen for 3D lead-halide perovskites under similar conditions.  $\text{BA}_2\text{PbI}_4$ ’s light-induced conductivity had a rise time and decay time of  $\sim 100 \mu\text{s}$ ,  $10^4$  slower than expected from direct electron-hole recombination and yet  $10^5$  faster than the conductivity-recovery times recently observed in 3D lead-halide perovskites and attributed to the relaxation of photogenerated vacancies. What sample properties are probed by electric force microscope measurements remains an open question. A Lagrangian-mechanics treatment of the electric force microscope experiment was recently introduced by Dwyer, Harrell, and Marohn which enabled the calculation of steady-state electric force microscope signals in terms of a complex sample impedance. Here this impedance treatment of the tip-sample interaction is extended, through the introduction of a time-dependent transfer function, to include time-resolved electrical scanned probe measurements. It is shown that the signal in a phase-kick electric force microscope experiment, and therefore also the signal in a time-resolved electrostatic force microscope experiment, can be written explicitly in terms of the sample’s time-dependent resistance (*i.e.*, conductivity).

## 1 Introduction

Recent interest in using 2D Ruddlesden–Popper perovskite for photovoltaic applications has increased due to unique properties that distinguish them from the widely studied 3D perovskite. These properties include higher environmental stability due to the hydrophobic nature of the organic spacer and higher formation energy,<sup>1</sup> increased exciton binding energy due to alternating organic and inorganic layer with disparate dielectric constants,<sup>2</sup> surprisingly reduced in-plane ion motion both in the dark and under illumination presumably due

1 to higher formation energy of vacancies,<sup>3</sup> and the existence of edge states with low energy  
2 that are thought to help in exciton dissociation or lead to unusual charge carrier densities.<sup>4,5</sup>  
3 DC galvanostatic polarization measurements and AC impedance spectroscopy have revealed  
4  $(C_6H_{10}N_2)PbI_4$ , a commonly studied 2D perovskite, is a mixed ionic electronic conductors with  
5 low dark ionic and electronic conductivities on the order of  $1 \times 10^{-10} \text{ S cm}^{-1}$ .<sup>6</sup> Despite the pres-  
6 ence of the insulating organic layer and exciton binding energy of 100's of milli electron volts,  
7 free carrier generation in the quasi 2D perovskites with more than one layer of the inorganic  
8 octahedron ( $n > 1$  where  $n$  is the number of layers of the inorganic octahedron) seems to be  
9 efficient enough such that power conversion efficiency greater than 18% have been reported  
10 for solar cells made from a quasi-2D perovskite.<sup>7</sup> The working principle of such cells for higher  
11  $n$  is thought to involve charge transfer from layered 2D regions to a 3D perovskite network.<sup>8</sup>  
12 An alternative proposal for the efficient generation of free carriers from bound exciton involves  
13 edge states.<sup>4</sup> Anisotropy between in-plane and out-of-plane charge carrier mobility has been  
14 reported; the charge carriers motion is not inhibited completely in the out-of-plane direction,  
15 with reported mobility on the order of  $1 \times 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ <sup>8</sup> for an  $n = 3$  system. Here we study  
16 a 2D Ruddlesden–Popper phase of perovskite butylammonium lead iodide  $((C_4H_9NH_3)_2PbI_4$   
17 or  $BA_2PbI_4$ , herein referred BAPI). BAPI, with  $n = 1$ , where higher dimension/bulk perovskite is  
18 absent, can act as a model system to understand the true nature of charge motion in a quasi 2D  
19 system.

20 We have recently revealed persistent photoinduced conductivity in a variety of 3D per-  
21 ovskite samples using time- and frequency-resolved electrical SPM techniques.<sup>9–11</sup> The con-  
22 ductivity dynamics in these 3D systems, while substrate dependent, generally exhibited a slow,  
23 activated recovery in the dark around room temperature which became faster at low temper-  
24 ature. We hypothesized that light-induced vacancy generation could explain the long lived  
25 changes in conductivity.<sup>9,11</sup> We wanted to test if the same surprisingly persistent light-induced  
26 conductivity was present in the 2D system.

27 Time resolved electrical scanning probe microscopy (SPM) techniques have been used to

1 probe charge dynamics in photovoltaic and ion conducting materials.<sup>12</sup> These techniques com-  
2 bine the very high charge sensitivity of microcantilevers with carefully designed measurement  
3 protocols to push the time resolution of measured sample parameters (such as capacitance,  
4 non-contact friction, dissipation, and surface (photo) voltage) beyond the natural time res-  
5 olution limit of the cantilever period. To harness the true power of these microscopic and  
6 localized measurements, and compare results with bulk electrical characterization techniques,  
7 it is imperative that we establish which sample properties these measurements probe.<sup>13</sup> By  
8 extending a recently developed theoretical framework<sup>14</sup> for quantitatively interpreting elec-  
9 trical scanning probe microscopy data, we show here that changes in sample resistance can be  
10 followed by combination of techniques that measure sample response either in the frequency  
11 and/or in the time domain.

12 Ginger and coworkers have recently performed one of the first electrical SPM studies of the  
13 2D Ruddlesden–Popper phase of perovskite BAPI to characterize light-induced charge carrier  
14 dynamics. They used two electrical scanning probe techniques, namely fast-free time-resolved  
15 electrostatic force microscopy (FF-tr-EFM) and general-acquisition-mode Kelvin probe force  
16 microscopy (G-mode KPFM, or G-KPFM). They measured charging times and time-resolved  
17 surface photovoltage measurements, observing 100 to 1000 $\mu$ s dynamics in these material  
18 in response to light irradiation over a large area.<sup>15</sup> They attributed these dynamics to trap-  
19 mediated electron or ion motion. The two techniques measured time scales that were dif-  
20 ferent by several hundred microseconds. Dielectric response in 3D perovskites is known to be  
21 highly frequency dependent where the low frequency response ( $\leq 10$  kHz) is dominated by ion  
22 motion.<sup>16</sup> Therefore, we expect that measurements that sample different parts of the frequency  
23 response may have a different temporal dependence in response to an external stimuli.

24 Here we use time-resolved EFM (tr-EFM),<sup>17</sup> phase-kick EFM (pk-EFM),<sup>18</sup> broadband local  
25 dielectric spectroscopy (BLDS),<sup>9,19</sup> and dissipation microscopy<sup>9,11,20</sup> to probe changes in the  
26 sample in response to optical irradiation. We show that the transfer function representation of  
27 SPM used to describe broadband local dielectric spectroscopy and dissipation measurements

1 can be extended to understand the origin of the temporal response measured in time resolved  
2 tr-EFM and pk-EFM measurements.<sup>14</sup> We show that these measurements probe changes in  
3 sample resistance instead of the commonly assumed capacitance changes in 2D BAPI. We find  
4 that light-induced conductivity in the 2D BAPI system recovers promptly, in marked contrast to  
5 the slow, activated recovery observed in the many 3D perovskite systems studied to date.

## 6 **2 Experimental Section**

### 7 **2.1 Scanning probe microscopy**

8 All experiments were performed under vacuum ( $5 \times 10^{-6}$  mbar) in a custom-built scanning  
9 Kelvin probe microscope described in detail elsewhere.<sup>18</sup> The cantilever used was a Mikro-  
10 Masch HQ:NSC18/Pt conductive probe. The resonance frequency and quality factor were  
11 obtained from ringdown measurements and found to be  $\omega_c/2\pi = f_c = 60.490\text{kHz}$  and  $Q =$   
12 28000 respectively at room temperature. The manufacturer's specified resonance frequency  
13 and spring constant were  $f_c = 60$  to  $75\text{kHz}$  and  $k = 3.5\text{Nm}^{-1}$ . Cantilever motion was detected  
14 using a fiber interferometer operating at  $1490\text{nm}$  (Corning model SMF-28 fiber). The sample  
15 was illuminated from above with a variable-intensity  $405\text{nm}$  diode laser. More experimental  
16 details regarding the implementation of broadband local dielectric spectroscopy, tr-EFM, and  
17 pk-EFM can be found in the Supporting Information.

### 18 **2.2 Sample preparation**

19 Thin film of  $\text{BA}_2\text{PbI}_4$  was synthesized on ITO adapted from the procedure reported in Ref. 15.  
20 The film was approximately 500 nm thick. Based on the absorption coefficient, we expect the  
21 absorption length to be  $\approx 150\text{nm}$  which is much shorter than the film thickness. Briefly, pre-  
22 cursor solution was prepared by dissolving 1.8 M butylammonium iodide (BAI) and 0.9 M lead  
23 (II) iodide ( $\text{PbI}_2$ ) in 1 mL of dimethylformamide (DMF). Perovskite solution was spin coated on  
24 the ITO substrates at 4000 rpm for 40 s (with high acceleration), followed by thermal annealing

<sup>1</sup> for 10 min at 100 °C. A representative AFM, absorption spectra, and XRD spectra is included in  
<sup>2</sup> the supporting information.

### <sup>3</sup> 3 Results

#### <sup>4</sup> 3.1 Theoretical background

<sup>5</sup> In this section, we extend our previously developed impedance model for SPM experiments<sup>9,14</sup>  
<sup>6</sup> to the situation where the sample properties are time-dependent. Our new model allows us to  
<sup>7</sup> interpret all of our SPM data in a common framework. We explain the key features of this model  
<sup>8</sup> below; a detailed derivation is included in the Supporting Information (S1).

<sup>9</sup> We investigate the sample's time-dependent properties in the dark and under illumination  
<sup>10</sup> using SPM experiments. To interpret these experiments, we need to connect the observed  
<sup>11</sup> cantilever amplitude, frequency, and phase back to the sample properties. These observable  
<sup>12</sup> depend on the tip-sample force  $F_{ts} = C'_{tip}(x)q_t^2/(2C_{tip}(x)^2)$  where  $C_{tip}$  is the tip capacitance,  
<sup>13</sup>  $C'_{tip} = dC_{tip}/dx$  is the derivative of the tip capacitance with respect to the vertical direction,  
<sup>14</sup>  $x$  is the cantilever displacement, and  $q_t$  is the tip charge. (In keeping with the conventional  
<sup>15</sup> notation for one-dimensional harmonic oscillators and to be consistent with the notation used  
<sup>16</sup> in Ref. 14, we represent cantilever displacement with  $x$  rather than  $z$ , with increasing  $x$  cor-  
<sup>17</sup> responding to motion of the cantilever tip away from the sample surface.) The component  
<sup>18</sup> of the tip-sample force evolving in phase with tip motion shifts the cantilever frequency,  $\Delta f_c$ ,  
<sup>19</sup> while the out-of-phase force component causes dissipation,  $\Gamma_s$ , which in turn changes the can-  
<sup>20</sup> tilever amplitude. Tip charge  $q_t$  depends on tip capacitance and sample impedance (Eq. S8 and  
<sup>21</sup> Eq. S9). In Refs. 9 and 14, we used an impedance model and perturbation theory to develop  
<sup>22</sup> an accurate steady-state approximation for  $q_t$  and the cantilever observables. In this model,  
<sup>23</sup> the central quantity that links sample properties to cantilever observables is the transfer func-  
<sup>24</sup> tion  $H(\omega) = V_t/V_{ts}$ . Here  $V_t$  is the voltage dropped across the tip-sample gap;  $V_{ts}$  is the applied  
<sup>25</sup> external voltage, dropped across both the tip-sample gap and the sample; and  $\omega$  is a modu-

1 lation frequency. The tip capacitance  $C_{\text{tip}}$  and sample impedance  $Z(\omega)$  together determine  
 2 the transfer function,  $H(\omega) = (1 + j\omega C_{\text{tip}} Z(\omega))^{-1}$ . The central results of this model are that  
 3  $\Delta f \propto \text{Re}[H(\omega_0)](V_{\text{ts}} - \phi)^2$  and  $\Gamma_s \propto \text{Im}[H(\omega_0)](V_{\text{ts}} - \phi)^2$  with  $\omega_0 = 2\pi f_c$  the cantilever frequency  
 4 and  $\phi$  the surface potential (Eqs. 7 and 8 in Ref. 14). The signal  $\alpha$  in the broadband dielectric  
 5 spectroscopy (BLDS) measurement described below can likewise be expressed in terms of the  
 6 transfer function:  $\alpha \propto |H(\omega_m)|^2(V_{\text{ts}} - \phi)^2$ , with  $\omega_m$  the voltage-modulation frequency (Eq. 10  
 7 in Ref. 14). These expressions straightforwardly connect sample resistance and capacitance,  
 8 through  $Z$  and  $H$ , to changes in cantilever frequency, dissipation, and BLDS signal.

9 The above description is only appropriate when the sample properties are time-  
 10 independent and the experiment is conducted at steady-state. Yet we are also interested in  
 11 analyzing time-dependent changes in sample properties in response to illumination. For this  
 12 reason, we extend the theoretical treatment to the case where the sample parameters may be  
 13 time-dependent. Assuming the applied voltage or light does not induce oscillations of the tip  
 14 charge at the cantilever frequency, our results for time-independent systems can be extended to  
 15 time-varying systems by replacing the time-independent transfer function  $H(\omega)$  with the time-  
 16 varying transfer function  $H(\omega, t)$  defined below.<sup>21</sup> Using these assumptions, we derive a general  
 17 result for the frequency shift in an electric or Kelvin probe probe microscopy experiment

$$\Delta f(t) = -\frac{f_c}{4k_c} \left( C_q'' + \Delta C'' \text{Re}[H(\omega_c, t)] \right) (V_t(t) - \phi)^2 \quad (1)$$

18 where  $f_c$  is the cantilever resonance frequency,  $k_c$  is the cantilever spring constant,  $C_q''$  and  $\Delta C''$   
 19 are related to the second derivative of the tip capacitance,  $\phi$  is the surface potential that rep-  
 20 resents internal voltage source, and  $V_t(t)$  is the tip voltage drop calculated under the assump-  
 21 tion that the tip displacement is fixed at zero. Since  $\phi < V_t$  (Fig. S7), we ignore  $\phi$  to simplify  
 22 our analysis. This result allows pk-EFM and tr-EFM, experiments that inherently involve time-  
 23 dependent dynamics, to be analyzed using the same framework as the steady state measure-  
 24 ments considered previously with this theory. The applicability of this approximation to the

1 experiments described here is illustrated by the close agreement between numerical simula-  
 2 tions of the cantilever frequency and phase and the analytical approximation developed here  
 3 (Fig. S2 and Fig. S3).

4 **3.1.1 Transfer function**

5 In this paper, we model the sample using a parallel resistor  $R_s$  and capacitor  $C_s$ . If  $R_s$  and  $C_s$  are  
 6 time-independent, the transfer function  $H$  is

$$H(\omega) = \frac{1 + j\omega R_s C_s}{1 + j\omega R_s (C_s + C_{\text{tip}})} = \frac{1 + jg^{-1}\omega\omega_{\text{fast}}^{-1}}{1 + j\omega\omega_{\text{fast}}^{-1}} \quad (2)$$

7 where  $\omega_{\text{fast}}(t) = (R_s C_{\text{tot}})^{-1}$  is the time-independent frequency and  $g = C_{\text{tot}}/C_s$  where  $C_{\text{tot}} =$   
 8  $C_{\text{tip}} + C_s$  is the total capacitance.  $C_{\text{tip}}$  is the capacitance between tip and a hypothetical  
 9 ground plane located at the sample surface and  $C_s$  is the sample capacitance when the sample  
 10 impedance is modeled as a parallel resistor and a capacitor. The complex-valued transfer func-  
 11 tion in Eq. 2 has a real part which determines the in-phase forces and an imaginary part which  
 12 determines the out-of-phase forces acting on the cantilever. An increase in cantilever dissipa-  
 13 tion is caused by out-of-phase forces acting on the cantilever arising from changes in the value  
 14 of the imaginary part of the transfer function at the cantilever frequency  $\omega_c$ .<sup>9</sup>

The situation is considerably more complicated if  $R_s$  and/or  $C_s$  vary with time, as they do  
 when the light intensity is changed. In this case, the time-varying response function is

$$H(\omega, t) = \frac{C_{\text{tip}}}{C_{\text{tot}}} \int_0^\infty \exp \left[ - \int_{t-\tau}^t \omega_{\text{fast}}(t') dt' \right] \times \omega_{\text{fast}}(t-\tau) e^{-j\omega\tau} d\tau + \frac{C_s}{C_{\text{tot}}}, \quad (3)$$

15 where  $C_{\text{tot}} = C_{\text{tip}} + C_s$  is the total capacitance, and  $\omega_{\text{fast}}(t) = (R_s C_{\text{tot}})^{-1}$  is the time-dependent  
 16 frequency at which charge responds.

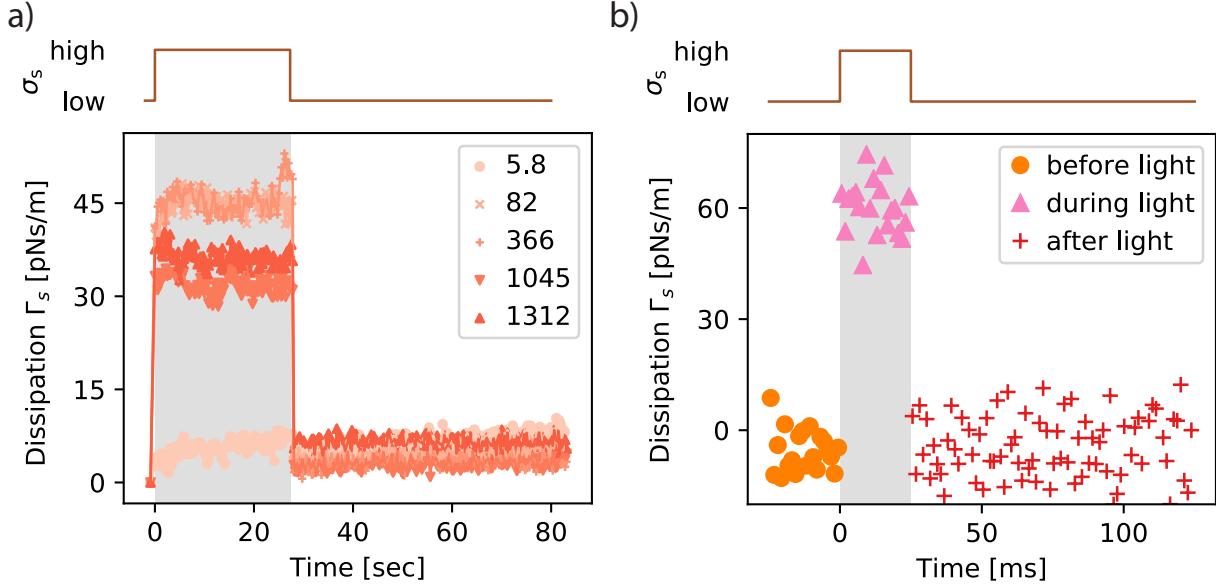


Figure 1: a) Bottom: Sample-induced cantilever dissipation over a BAPI film increases promptly when light is applied (shaded region) and recovers in less than 100ms when light is removed. Top: The corresponding sample conductivity. The dissipation is quantified through ring-down measurements at indicated light intensities (in  $\text{mW cm}^{-2}$ ). Experimental parameters:  $V_{ts} = -4 \text{ V}$ ,  $h = 125 \text{ nm}$ ,  $\lambda = 405 \text{ nm}$ . b) Light-induced sample dissipation (i.e. sample conductivity) recovers in less than 2ms. Data obtained through an implementation of tr-EFM with a constant DC bias applied, as described in the text. Experimental parameters:  $I_{hv} = 120 \text{ mW cm}^{-2}$ ,  $V_{ts} = -6 \text{ V}$ ,  $h = 200 \text{ nm}$ ,  $\lambda = 405 \text{ nm}$ .

### 1 3.2 Dissipation measurements

2 We start our investigation by measuring light induced cantilever dissipation. Dissipation tracks  
 3 changes in the sample resistance  $R_s$  and sample capacitance  $C_s$  when the associated time con-  
 4 stant  $R_s(C_s + C_{\text{tip}})$  is near the cantilever period of  $\omega_c^{-1}$  which is the case here. In our previous  
 5 report, we showed that changes in dissipation caused by light can take tens to hundreds of  
 6 seconds to recover in 3D perovskite samples.<sup>9,11</sup> We perform this measurement by recording  
 7 changes in cantilever ring-down time.<sup>9,11</sup> A DC voltage of  $-4 \text{ V}$  is applied before the start of the  
 8 measurement to ensure that any slow process in response to the tip electric field in the dark  
 9 will not interfere with the measurement. Here we see an increased dissipation when the light  
 10 is turned on that goes away within the time resolution of the measurement,  $\leq 100 \text{ ms}$ , at least  
 11 three orders of magnitude faster than the recovery time seen in 3D  $\text{CsPbBr}_3$  and 3D FAMACs  
 12 films.

1 To better resolve the associated time constant, we turn to an implementation of the tr-EFM  
2 measurement. The cantilever is allowed to freely ring down and the cantilever oscillation is  
3 demodulated to obtain the instantaneous cantilever amplitude, phase, and frequency. The  
4 measurement, shown in Fig. 1b, proceeds as follows. A DC voltage of  $-6\text{ V}$  is applied throughout  
5 the measurement. This voltage ensures that any slow tip charge equilibration on application of  
6 the tip electric field does not affect the measurement. The cantilever is allowed to ring down for  
7 125 ms starting at time  $t = -25\text{ ms}$ . At time  $t = 0\text{ s}$ , a 25 ms light pulse is initiated. This process  
8 is repeated for 600 acquisitions and the averaged cantilever amplitude data is divided into 2 ms  
9 bins. The known cantilever frequency and initial amplitude is used to extract cantilever  $Q$  for  
10 each 2 ms bin giving a time resolution of 2 ms. This time resolution is  $\sim 50$  times better than the  
11 time resolution in Fig. 1a. It is clear from these measurements that the changes in the transfer  
12 function  $H(\omega_c)$  at the cantilever frequency happen in less than 2 ms. Analysis performed with  
13 1 ms bins gave a similar result, albeit with worse SNR, confirming the sub-millisecond time scale  
14 of the underlying charge dynamics that cause dissipation.

### 15 **3.3 Broadband local dielectric spectroscopy**

16 Broadband local dielectric spectroscopy (BLDS) can be used to monitor changes in the  
17 impedance of a semiconducting sample over a wide frequency range.<sup>9,11</sup> In a BLDS spectrum,  
18 we plot a voltage normalized frequency shift alpha ( $\alpha$ ) with units of  $\text{Hz V}^{-2}$  versus the modula-  
19 tion frequency of the applied tip-sample voltage ( $2\pi f_m$ ). We turn to two different implemen-  
20 tations of broadband local dielectric spectroscopy, *i.e.* 1) average frequency-shift BLDS and 2)  
21 amplitude-modulation BLDS, to map out the transfer function over a wide frequency range. In  
22 the first implementation, described in detail elsewhere,<sup>9</sup> we record the average frequency shift  
23 as a function of applied voltage modulation frequency, with a 1:1 “modulation on” to “modu-  
24 lation off” duty cycle employed to allow lock-in detection. While the SNR of this measurement  
25 is relatively low compared to the amplitude-modulation BLDS shown in the proceeding figure,  
26 this measurement allows us to access the low frequency part of the transfer function below

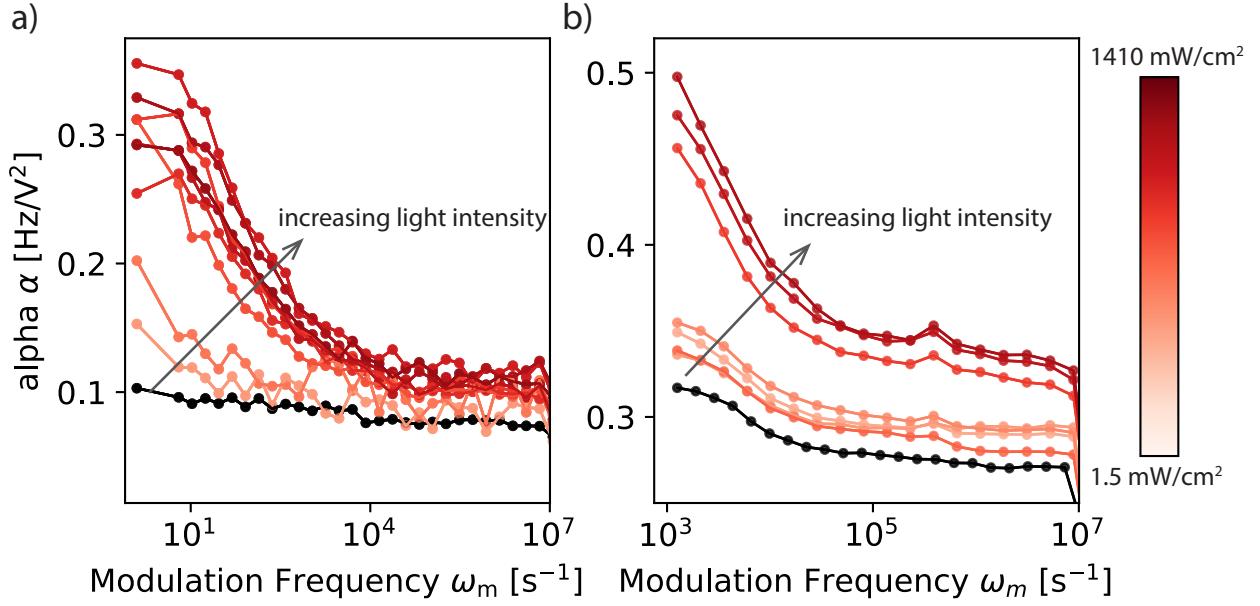


Figure 2: a) Frequency-shift broadband dielectric spectra reveals a large increase in the sample's low frequency ( $< 1000\text{s}^{-1}$ ) response when the light intensity is increased. Experimental parameters:  $V_{\text{pp}} = 6\text{V}$ ,  $h = 125\text{ nm}$ . b) High frequency response at indicated illumination intensities resolved by amplitude modulation broadband local dielectric spectroscopy. Experimental parameters:  $V_{\text{pp}} = 6\text{V}$ ,  $h = 125\text{ nm}$ .

1 1000 s<sup>-1</sup>. In Fig. 2, we can see that the transfer function changes in the low frequency region  
 2 are much more obvious than the small changes in the near-cantilever-frequency region. This  
 3 observation is in line with the relatively small dissipation changes measured in Fig. 1b. We note  
 4 that in the dark, the roll off of the transfer function is below the lowest probed frequency here  
 5 (1 s<sup>-1</sup>).

6 To better resolve changes in the high frequency region, we turn to amplitude-modulated  
 7 BLDS, Fig. 2b. Different numerical values of  $\alpha$  in Fig. 2a and Fig. 2b are attributed to different  
 8 regions in the sample. This measurement's lock-in detection scheme allows us to clearly resolve  
 9 changes in the dielectric spectra in the frequency range of kHz to MHz that were not obvious  
 10 in Fig. 2a. In Fig. 2b, we can see that increasing light intensity increases the value of  $\alpha$ . To  
 11 describe changes in BLDS spectra, we have previously used an *RC* electrical model for the  
 12 sample with one resistor and one capacitor. While a simple *RC* is often not enough to cap-  
 13 ture the full impedance changes in a mixed ionic-electronic conductor,<sup>11</sup> it is nevertheless a  
 14 useful starting point to understand physical process causing the changes in the impedance of

1 the sample. In our model, the measured  $\alpha$  is proportional to  $H(\omega)$  which is a function of two  
2 parameters

$$\omega_{\text{fast}}^{-1} = R_s(C_s + C_{\text{tip}}) \text{ and } g = C_{\text{tot}}/C_s. \quad (4)$$

3 with  $\omega_{\text{fast}}^{-1}$  a time constant describing transfer function roll-off. From Eq. 2, the high-frequency  
4 limit of  $H(\omega)$  is  $1/g$ , Eq. 4. The high-frequency data in Fig. 2 give  $g \sim 2$  to 3, indicating that the  
5 tip capacitance and sample capacitance are of similar size.

6 Light-induced changes in the BLDS spectra can be attributed to resistance and/or capac-  
7 itance. The BLDS spectra show  $g$  decreasing slightly with increasing light intensity. Taking  
8 derivatives,  $\delta g = -C_{\text{tip}}C_s^{-2}\delta C_s$ . If  $C_{\text{tip}}$  is independent of light intensity, the observed decrease  
9 in  $g$  thus indicates a slight increase in  $C_s$  with increasing light intensity. The BLDS spectra also  
10 show a time constant  $\omega_{\text{fast}}^{-1}$  decreasing from seconds to milliseconds as light intensity increases  
11 from 1.5 to  $1410 \text{ mW cm}^{-2}$  (and pinned at a near constant value at high light intensity). This  
12 decrease in time constant requires a decrease in the  $R_s(C_s + C_{\text{tip}})$  product, which can then only  
13 be explained by a large decrease in  $R_s$  with increasing light intensity.

14 These conclusions can be visualized with the help of the simulated spectra shown in Fig. 3a.  
15 As explained in Refs. 9, 11, and 14, the BLDS signal  $\alpha$  is proportional to  $|H(\omega_m)|^2$ , which is  
16 dominated by  $\text{Re}[H(\omega_m)]$  at high frequency. In Fig. 3a we plot representative  $\text{Re}[H(\omega_m)]$  versus  
17  $\omega_m$  spectra in the dark and under three light-on scenarios: increased  $C_s$ , decreased  $R_s$ , and both  
18 changing together. Only the third scenario is qualitatively consistent with the Fig. 2 spectra.

19 These conclusions about photoinduced changes in sample capacitance and resistance are  
20 corroborated by the dissipation data of Fig. 1. The dissipation signal is proportional to <sup>9,11,14</sup>  
21  $\text{Im}[H(\omega_c)]$ , the imaginary part of the transfer function  $H$  evaluated at the cantilever frequency  
22  $\omega_c$ . In Fig. 3b we plot representative  $\text{Im}[H]$  spectra for the four scenarios, with the cantilever  
23 frequency  $\omega_c = 2\pi f_c$  indicated by a dotted vertical line and the value of the transfer function  
24 near  $\omega_c$  indicated by gray shading. Here either the second or third light-on scenario are con-  
25 sistent with the observed increase in dissipation with increasing light intensity. The observed  
26 dissipation unconditionally requires a decrease in  $R_s$  with increasing light intensity.

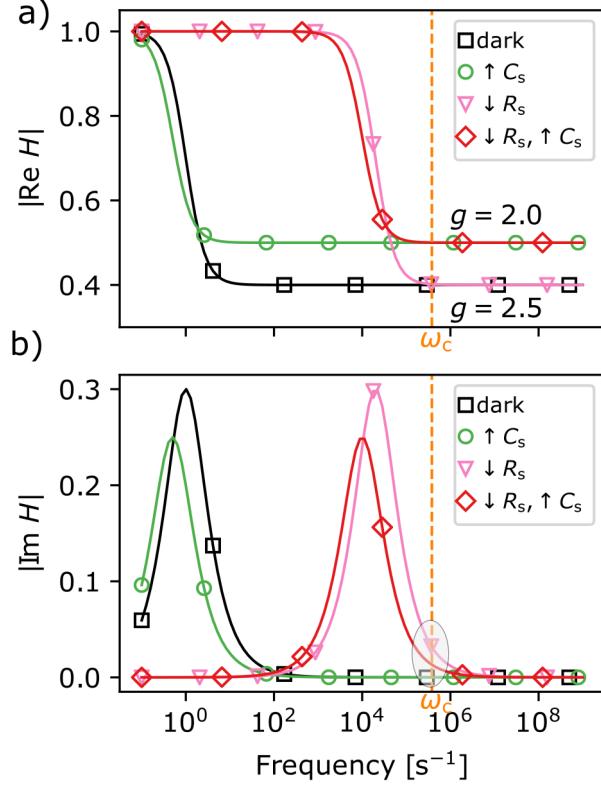


Figure 3: (a) Real and (b) imaginary part of the Eq. 2 transfer function plotted for the following cases: (i) sample in the dark and under illumination with (ii) increased sample capacitance ( $\uparrow C_s$ ), (iii) decreased sample resistance ( $\downarrow R_s$ ), and (iv) increased sample capacitance and decreased sample resistance ( $\downarrow R_s, \uparrow C_s$ ). To mimic an increase in  $C_s$ , the value of  $g$  in Eq. 2 was decreased from 2.5 to 2.0. The cantilever frequency  $\omega_c$  is indicated as a dashed vertical line. The sample-induced frequency shift is proportional to  $\text{Re } H(\omega_c)$ , (a), while sample-induced dissipation is proportional to  $\text{Im } H(\omega_c)$ , (b). The gray zone highlights the region in (b) that determines  $\text{Im } H(\omega_c)$ .

1 We note that the numerical value of the time constant  $\omega_{\text{fast}}^{-1}$  deduced from the BLDS spectra  
2 will depend on both sample capacitance and tip capacitance and this value may not represent  
3 the “life time” of carrier generation or recombination in the sample. Nevertheless, changes in  
4 this time constant can be directly related to changes in sample conductivity.

## 5 4 tr-EFM and pk-EFM

6 Frequency shift (e.g. tr-EFM and FF-tr-EFM)<sup>17,22,23</sup> and phase shift methods (pk-EFM)<sup>14,18</sup> have  
7 been the recent focus of high temporal resolution SPM measurements and have shown the

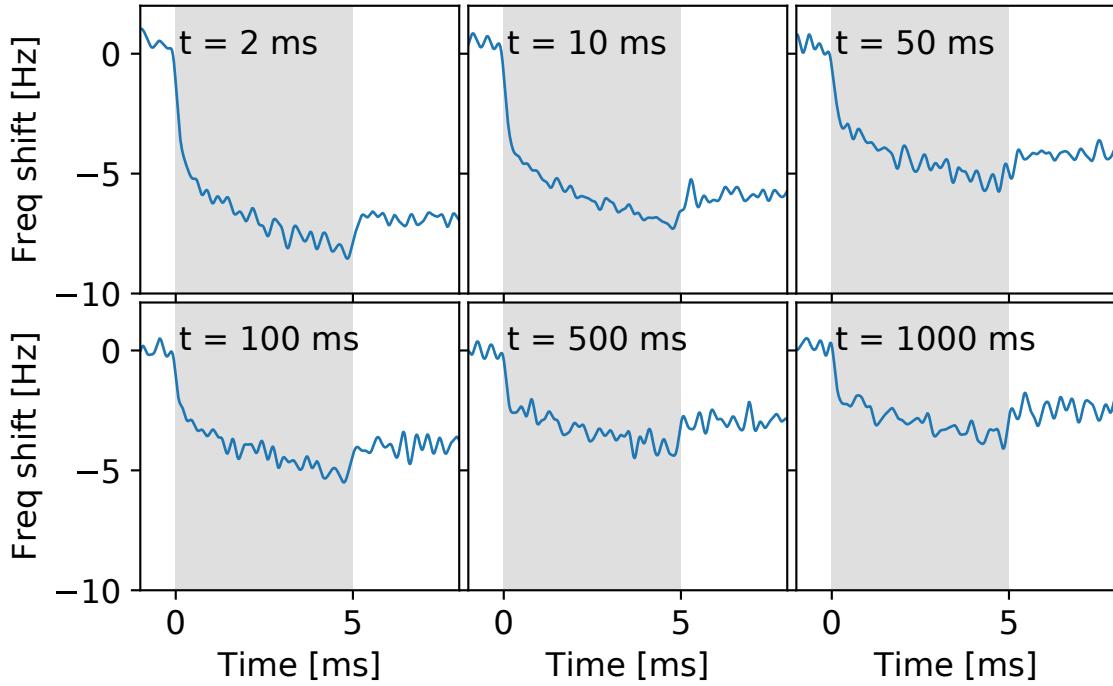


Figure 4: The light-induced frequency shift is determined by the sample bias time  $t$  in the dark. tr-EFM traces with different bias time before the illumination period (shaded region). See Fig. S4 for the associated experimental timing diagram. Experimental parameters:  $I_{hv} = 120 \text{ mW cm}^{-2}$ ,  $V_{ts} = 6 \text{ V}$ ,  $h = 200 \text{ nm}$ , and  $\lambda = 405 \text{ nm}$ .

ability to achieve time resolution down to 1% of the cantilever cycle.<sup>18</sup> To date, the origin of the observed frequency shift and the accompanying phase shift has been attributed solely to changes in sample capacitance.<sup>18,23</sup> Here we show that light-induced changes in sample conductivity better explain the data in samples having an appreciable impedance, such as our lead halide perovskites.

We performed tr-EFM experiments in which the sample was biased for a time  $t$  in the dark before light was applied. In Fig. 4 we show tr-EFM traces of cantilever frequency versus time for different values of the delay time ( $t$ ) between the start of the voltage pulse and the start of the light pulse. For a short delay  $t$  such as 2 ms, we see an 8 Hz light-induced frequency shift with a fast component happening on the  $\approx 100 \mu\text{s}$  time-scale and a slow component evolving on the millisecond timescale. With increasing delay  $t$ , both the fast and slow components decrease in magnitude so that for  $t = 1000 \text{ ms}$ , the light-induced frequency shift is 4 Hz (observation 1).

1 When the light is turned off at 5 ms, most of the light-induced frequency shift is retained (observation 2). In Fig. S5, we show the full trace of the tr-EFM measurement for  $t = 1$  s. These two key observations cannot be easily explained in terms of a photocapacitance, since photocapacitance should not depend on the delay  $t$  and should revert back to the original dark capacitance when the light is turned off.

6 What is the physical explanation for the finding that sample bias in the dark strongly affects  
7 the subsequent light-induced frequency shift? The key insight from Eq. 1 is that the frequency  
8 shift depends on both  $\text{Re}[H(\omega_c, t)]$ , the time-varying transfer function evaluated at the can-  
9 tilever resonance frequency, and the time-dependent tip charge  $q_t$  or tip voltage  $V_t = q_t/C_{\text{tip}}$ .

10 From Fig. 2, we know  $\text{Re}[H(\omega_c)]$  increases with light because the sample conductivity increases.  
11 This increase in sample conductivity explains the two key observations from Fig. 4 once we  
12 properly recognize the role of the tip voltage drop  $V_t$ .

13 In the Fig. 4 tr-EFM experiment, a tip-sample voltage on the order of 5 to 10V was applied  
14 before the light was turned on. The circuit schematics in Fig. 5(a,b) show the modeled tip  
15 voltage  $V_t(t)$  when  $C_s = C_{\text{tip}}$  ( $g = 2$ ). When the tip-sample bias is applied, equal charges build up  
16 on both capacitors and the voltage drop across the tip capacitor is  $V_t = V_{\text{ts}}C_s/(C_s + C_{\text{tip}})$  (Fig. 5a).  
17 The sample capacitor discharges with an  $RC$  time constant  $R_s(C_{\text{tip}} + C_s)$  until eventually all of  
18 the applied tip-sample voltage drops across the tip capacitor:  $V_t = V_{\text{ts}}$  (Fig. 5b). In the dark the  
19 sample is very resistive so it can take seconds for this equilibrium condition to be reached. If the  
20 light is turned on before the tip is fully charged, the increased sample conductivity (decrease in  
21 sample resistance  $R_s$ ) causes the tip to charge more quickly, increasing  $V_t$  and therefore causing  
22 a larger light-induced frequency shift.

23 Figure 5(c,d) shows a simulation to illustrate how this tip charging explains the data of Fig. 4.  
24 It was assumed the sample resistance  $R_s$  decreases when the light is turned on. The evolution of  
25 tip position  $x$  and tip and sample charge are described by a set of coupled differential equations  
26 (see SI Section S3 and Fig. S1 for details). Fig. 5(c,d) shows the frequency shift calculated numer-  
27 ically from  $x(t)$  (points) and analytically from Eq. 1 (curves). The four traces show increasing

1 initial tip voltages that correspond to the increasing bias times in Fig. 4. With reasonable values  
2 of sample resistance in the dark, sample resistance in the light, sample capacitance, and tip  
3 capacitance, the numerical simulations and Eq. 1 do an excellent job qualitatively reproducing  
4 observation 1; the closer the initial tip voltage  $V_t$  is to its steady-state value  $V_t = V_{ts} = 10V$ , the  
5 smaller the magnitude of the observed frequency shift. The slow component of the frequency  
6 shift in Fig. 4 can be explained by the millisecond-scale  $RC$  time apparent in the BLDS spectra  
7 under steady-state illumination. To corroborate this hypothesis, we tracked changes in  $\alpha$  at  
8  $\omega_m = 1257 \text{ s}^{-1}$  ( $1/\omega_m = 0.8 \text{ ms}$ ) in real time when the light was turned on (Fig. S6; temporal res-  
9 olution is limited to 1s by the finite amplitude-modulation frequency and associated lock-in  
10 measurement). The dielectric response changes in  $\leq 1 \text{ s}$ , consistent with the timescales seen in  
11 Fig. 4.

12 The retention of light-induced frequency shift in the dark (observation 2) can also be  
13 explained by carefully considering the evolution of the tip voltage  $V_t$ . When the light is turned  
14 off, the tip does not discharge, since  $V_t = V_{ts}$  is the equilibrium state in both the light and dark.  
15 Therefore any frequency shift caused by an increase in  $V_t$  is retained even after the light is turned  
16 off. This explains why the change in  $\Delta f$  when the light is turned off in Fig. 4 is relatively small  
17 no matter how long the initial pre-bias time was. Any decrease in frequency shift when the light  
18 is turned off is attributed to changes in  $\text{Re } H(\omega_c, t)$  (Eq. 1), since only changes in  $\text{Re } H(\omega_c, t)$  are  
19 reversed when the light turns off.

20 The SI Section S3 analysis reveals that the frequency-shift dynamics are greatly simplified  
21 if the tip charge, *i.e.*  $V_t$ , has reached steady state before light is applied. We can achieve the  
22 steady-state condition by applying a DC voltage for an extended period of time before begin-  
23 ning the measurement (Fig. 6). Pre-biasing the sample with constant applied voltage for 5 min-  
24 utes before the measurement and keeping the bias always on during the measurement keeps  
25 the tip charged ( $V_t = V_{ts}$ ); therefore, all observed frequency shifts reflect changes in the transfer  
26 function at the cantilever frequency  $\text{Re } [H(\omega_c, t)]$ .

27 For  $V_{ts} = 5V$  and  $-5V$ , the light induced frequency changes are essentially identical, with

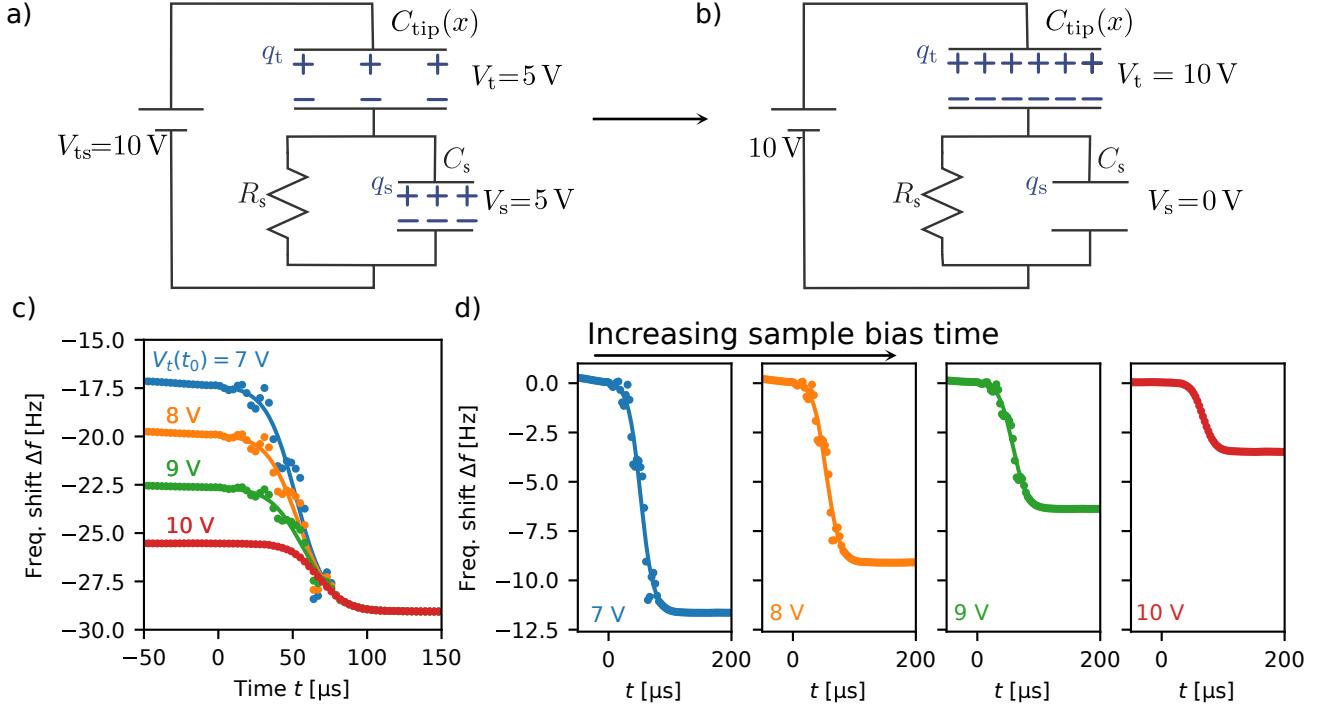


Figure 5: The effect of the initial tip voltage on frequency-shift transients. (a) The model circuit used in this simulation. After the tip-sample voltage  $V_{ts}$  is applied, charge builds up on the tip capacitor and sample capacitor. (b) Eventually, the sample capacitor discharges and  $V_t = V_{ts}$ . (c) Frequency shift versus time calculated by numerical simulation (points) and the Eq. 1 analytical approximation (lines) for different initial tip voltages  $V_t(t_0) = q_t^{(0)}(t_0)/C_{tip} = 7, 8, 9, 10$  V. (d) The data of (c), shifted so  $\Delta f(t=0)=0$  for each trace to more easily compare with Fig. 4. All traces were smoothed by averaging the frequency shift over a single period twice. Simulation parameters:  $R_{dark} = 1 \times 10^7$  MΩ;  $R_{light} = 1 \times 10^4$  MΩ;  $\tau_L = 10$  μs;  $V_t(t_0) = 7, 8, 9$  and 10 V (blue, orange, green, red);  $C_s = 0.05$  fF;  $C_{tip} = 0.1$  fF;  $C'_{tip} = -28$  fF nm<sup>-1</sup>;  $f_c = 62.5$  kHz;  $k_c = 3.5$  N m<sup>-1</sup>;  $Q = 26000$ ,  $A_0 = 50$  nm;  $A_0 = 50$  nm;  $t_0 = -200$  μs;  $\phi_0 = 0$  (see SI Section S3).

1 similar rise and fall time on the order of  $\approx 100$  μs. The insensitivity of the frequency shift, rise  
 2 time, and fall time to the sign of the applied tip-sample voltage rules out any effects on the  
 3 measurement due to surface potential changes. When the light is turned off at  $t = 25$  ms, the  
 4 frequency shift reverts to the dark value, unlike in Fig. 4. This observation is consistent with  
 5 Eq. 1 when  $V_t$  is constant—the light-induced changes in  $\text{Re}[H(\omega_c)]$  are fully reversible. Con-  
 6 sidered together, the data of Fig. 1b and Fig. 6d allow us to conclude that the transfer function  
 7 at  $\omega_c$ , determining the light-induced frequency shift and dissipation, is evolving on the 100 μs  
 8 timescale.

9 To better resolve this time constant, we carried out a pk-EFM experiment in which the

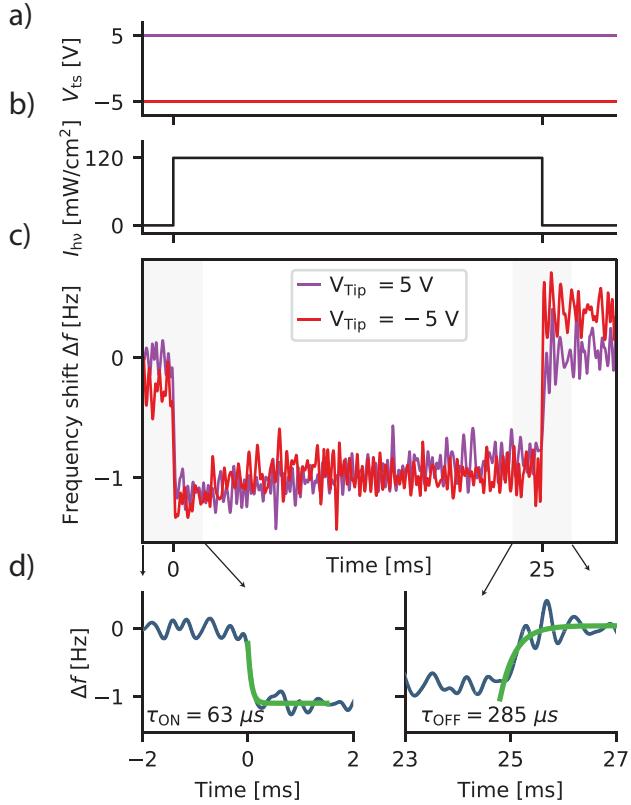


Figure 6: Light-induced changes in cantilever frequency likewise occur promptly. tr-EFM frequency shift signals with a DC bias of  $\pm 5\text{ V}$ . (a) The tip-sample voltage is applied continuously. (b) Light is switched on at time  $t = 0\text{ ms}$  and switched off at  $t = 25\text{ ms}$ . (c) Light induced frequency shift measured for  $V_{\text{ts}} = \pm 5\text{ V}$ . (d) Zoom-in of the  $+5\text{ V}$  trace. Solid (green) lines are a fit to a single exponential, with light-on and light-off time constants indicated. Experimental parameters:  $I_{h\nu} = 120\text{ mW cm}^{-2}$ ,  $h = 200\text{ nm}$ ,  $\lambda = 405\text{ nm}$ .

1 applied tip-sample voltage is applied continuously except during a  $\sim 2\text{ ms}$  time period before the  
 2 end of light pulse. This short time without applied bias is necessary to measure cantilever phase  
 3 changes during the pulse time  $t_p$  arising from time-dependent sample conductivity. From a plot  
 4 of the measured phase shift versus pulse time, we deduce a time-constant of  $(70 \pm 16)\text{ }\mu\text{s}$ .

## 5 Discussion

6 The steady-state impedance spectra and photoconductivity dynamics of the 2D lead-halid per-  
 7 ovskite  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  observed using electrical scanned-probe measurements are qualita-  
 8 tively different from those found in the 3D lead-halide perovskite  $\text{CsBrBr}_3$ <sup>9</sup> and the state-of-

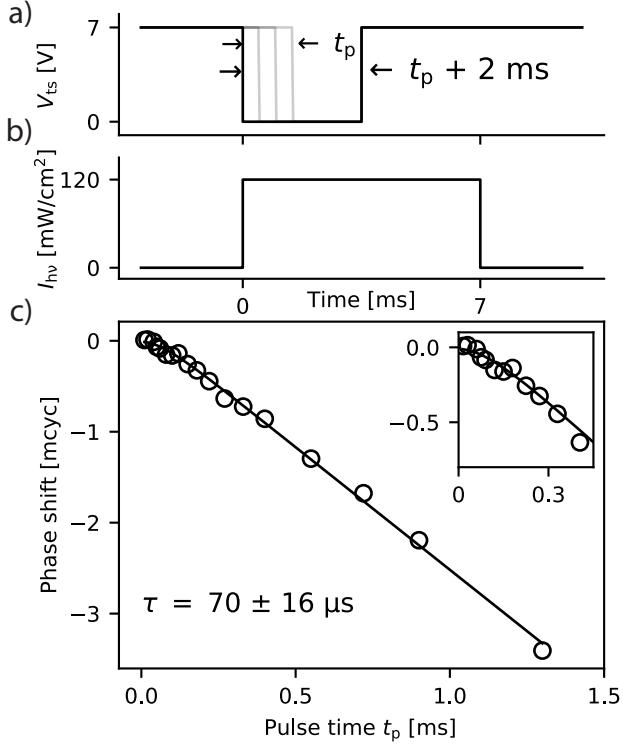


Figure 7: pK-EFM measures a photoconductivity rise time of  $\tau_{\text{on}} = 70 \mu\text{s}$ . Timing diagram of applied (a) tip voltage and (b) light, with  $t_p$  the duration of voltage- and light-pulse overlap. (c) Measured cantilever phase shift  $\Delta\phi$  versus pulse time  $t_p$ . Inset: Short  $t_p$  data. The solid line is a fit to  $\Delta\phi = \Delta\omega_{\text{ss}}(t_p - \tau_{\text{on}} + \tau_{\text{on}}e^{-t_p/\tau_{\text{on}}})$ , Eq. 19 in Ref. 18, with  $\Delta\omega_{\text{ss}}$  the steady-state light-induced cantilever frequency shift and  $\tau_{\text{on}} = 70 \pm 16 \mu\text{s}$  the photoconductivity rise time. Experimental parameters:  $I_{h\nu} = 120 \text{ mW cm}^{-2}$ ,  $V_{\text{ts}} = 7 \text{ V}$ ,  $h = 200 \text{ nm}$ ,  $\lambda = 405 \text{ nm}$ .

<sup>1</sup> the-art FAMACs alloy.<sup>11</sup>

<sup>2</sup> In Ref. 11 we used Maier's transmission-line treatment of impedance spectroscopy<sup>24,25</sup> to  
<sup>3</sup> show that in a mixed electronic-ionic conductor, for reasonable assumptions about size of the  
<sup>4</sup> electronic and ionic conductivities, the scanned-probe broadband local dielectric spectroscopy  
<sup>5</sup> (BLDS) experiment of Fig. 2 measures the *total* sample conductivity. Modeling the sample as a  
<sup>6</sup> parallel resistor-capacitor circuit and assuming  $C_{\text{tip}} \geq C_s$ , which is approximately valid here, the  
<sup>7</sup> associated rolloff frequency depends on sample resistance and is largely independent of sample  
<sup>8</sup> capacitance. In this limit, the rolloff frequency is given by  $\omega_{\text{fast}} \approx A\sigma_s/LC_{\text{tip}}$  with  $A$  the sample  
<sup>9</sup> area under the tip,  $L$  the sample thickness,  $C_{\text{tip}}$  the tip-sample capacitance, and  $\sigma_s = \sigma_{\text{eon}} + \sigma_{\text{ion}}$   
<sup>10</sup> the electronic plus ionic conductivity of the sample. Since values of  $A$ ,  $L$ , and  $C_{\text{tip}}$  used in the  
<sup>11</sup> Fig. 2 BLDS experiment on  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  are similar to values used in related experiments on

1  $\text{CsPbBr}_3$ , <sup>9</sup>  $\sigma_s \propto \omega_{\text{fast}}$ , and we can draw qualitative conclusions about the samples' conductivity  
2 by comparing the measured BLDS spectra and rolloff frequency in the two samples.

3 In  $\text{CsPbBr}_3$ , the impedance spectrum showed two distinct rolloffs<sup>9</sup> — one below 1 Hz,  
4 attributed in the dielectric spectroscopy literature to ionic motion, and a second rolloff that  
5 varied from  $1 \times 10^1$  Hz in the dark to above  $1 \times 10^5$  Hz at  $1000 \text{ mW cm}^{-2}$ . The second rolloff  
6 frequency, and hence the associated  $\sigma_s$ , had a nearly square root dependence on illumina-  
7 tion intensity, as one would expect for free carriers, and yet the conductivity exhibited a slow,  
8 activated recovery consistent with ion or vacancy motion. We subsequently proposed<sup>11</sup> that  
9 these contradictory observations could be rationalized by considering that photoexcitation  
10 can create interstitial iodine atoms and mobile iodine vacancies:<sup>26</sup> in Kröger-Vink notation,  
11  $\text{I}_\text{i}^x \xrightarrow{h\nu} \text{I}_\text{i}^x + \text{h}^\bullet + \text{e}' \rightleftharpoons \text{I}_\text{i}^x + \text{V}_\text{i}^\bullet + \text{e}'$ . The scanned-probe BLDS experiment measures the total con-  
12 ductivity and would therefore be sensitive to both  $\text{V}_\text{i}^\bullet$  and  $\text{e}'$ . In this view, the *circa* 10 s to 100 s  
13 conductivity recovery observed in  $\text{CsPbBr}_3$ <sup>9</sup> and the FAMACs alloy<sup>11</sup> is measuring the slow, acti-  
14 vated return to equilibrium of photogenerated halide vacancies and their geminate electrons.

15 In  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$ , in contrast, only one rolloff was apparent in the BLDS spectrum and  
16 the associated cutoff frequency varied from  $\leq 1$  Hz in the dark to  $10^3$  Hz at  $1000 \text{ mW cm}^{-2}$ . Com-  
17 pared to the  $\text{CsBrBr}_3$  film, the total conductivity of  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  is 100-fold smaller under  
18 illumination and saturates at high illumination intensity. We can make a rough estimate of the  
19 sample's dark conductivity using  $\sigma_s = LC_{\text{tip}}\omega_{\text{fast}}/A$  with  $\omega_{\text{fast}} = 1$  Hz. The sample thickness is  
20  $L = 500 \text{ nm}$  and we estimate  $C_{\text{tip}} = 1 \times 10^{-16} \text{ F}$  and  $A = \pi(h/2)^2$ , with  $h = 150 \text{ nm}$  the tip-sample  
21 separation. Using these parameters, we obtain  $\sigma_s = 4 \times 10^{-13} \text{ S cm}^{-1}$ , in reasonable agreement  
22 with the dark conductivity of  $1 \times 10^{-13} \text{ S cm}^{-1}$  measured for a single crystal of the 2D perovskite  
23  $(\text{PEA})_2\text{PbI}_4$ .<sup>27</sup>

24 In  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  the photoconductivity rise time  $\tau_{\text{on}}$  (Figs. 6 and 7) and decay time  $\tau_{\text{off}}$   
25 (Fig. 6) are  $70 \mu\text{s}$  and  $285 \mu\text{s}$ , respectively. The  $\tau_{\text{on}}$  and  $\tau_{\text{off}}$  times measured here are  $\sim 10^5$  faster  
26 than the conductivity decay times seen in multiple 3D lead-halide perovskites<sup>9,11</sup> but  $\sim 10^4$   
27 slower than the 10 ns electron-hole recombination times observed for  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  using

1 time-resolved microwave conductivity.<sup>28</sup> The  $\tau_{\text{on}} = (70 \pm 16)\mu\text{s}$  measured here using pk-EFM  
2 is in good agreement with the cantilever-frequency equilibration time of  $85\mu\text{s}$  measured near  
3 grain centers in  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  by Giridharagopal, Ginger, and coworkers using FF-tr-EFM.<sup>15</sup>  
4 In Ref. 15, spatial variations in signal were discussed in terms of ionic and electronic carrier  
5 transport, yet in the usual description of tr-EFM the frequency-shift signal, expressed in terms  
6 of a time-dependent tip-sample capacitance, has no overt connection to such transport pro-  
7 cesses. Building on prior work, here we provide that missing theoretical connection. In Ref. 14  
8 we showed, for a sample described by a parallel resistor-capacitor circuit, that the frequency  
9 shift and dissipation probe the real and imaginary parts, respectively, of the Eq. 2 transfer func-  
10 tion  $H(\omega)$ . The BLDS spectra of Fig. 2 and the qualitative changes in  $H(\omega)$ , frequency shift,  
11 and dissipation sketched in Fig. 3 reveal that the frequency shift observed by FF-tr-EFM and  
12 pk-EFM in  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$  is dominated by sample resistance, not capacitance. Motivated by  
13 this observation, here we extended our Ref. 14 treatment of transient electric force microscope  
14 experiments to include a time-dependent sample resistance (*i.e.*, conductivity). This finding  
15 shows that pk-EFM data can in principle be directly compared to time-dependent microwave  
16 conductivity measurements. Since pk-EFM (and FF-tr-EFM) is sensitive to *both* time depen-  
17 dent capacitance and resistance, future time-resolved EFM studies should include measure-  
18 ments of BLDS spectra versus light intensity in order to establish whether  $C_s$  or  $R_s$  is dominating  
19 the sample's response. Because  $\omega_{\text{fast}} \ll \omega_c$  here, the sample's changing conductivity has only  
20 a small effect on the cantilever frequency and dissipation here; imaging spatial variations in  
21  $(\text{C}_4\text{H}_9\text{NH}_3)_2\text{PbI}_4$ 's steady-state photoconductivity would be better accomplished using a BLDS  
22 signal ( $\omega_m \sim 100\text{Hz}$  to  $1000\text{Hz}$ ) than a frequency or dissipation signal.

23 While the electronic and ionic conductivity of  $(\text{C}_6\text{H}_{10}\text{N}_2)_2\text{PbI}_4$  in the dark has been measured  
24 previously,<sup>6</sup> there is little precedent beyond Ref. 15 for observing a *circa*  $100\mu\text{s}$  conductivity  
25 rise time and decay time following optical irradiation for a 2D lead-halide perovskite. There  
26 are two general explanations for such a conductivity-relaxation time. For photogenerated elec-  
27 tronic carriers to equilibrate this slowly would require trapping; hole trapping at acceptor-type

1 iodine vacancies is plausible,<sup>27</sup> as is localization at iodine edge states.<sup>29</sup> In the case of ionic  
2 conductivity, the rise and decay could be due to a dependence of ionic conductivity on the con-  
3 centration of (decaying, trapped) electrons and holes. Alternatively, we could be observing the  
4 formation and relaxation of photogenerated vacancies; in this case we would have to explain  
5 why the formation and relaxation rate is  $10^5$  times faster than in 3D perovskites. In support of  
6 this conjecture, Wang *et al.* found the light-on and light-off transient time for a single-crystal  
7 photodetector of BAPI to be  $10^3$  larger than for a 3D perovskite.<sup>30</sup> The fast rise time is puzzling,  
8 because we would expect the vacancy formation energy to be high in our 2D perovskite based  
9 on the comparatively low dark conductivity seen in related systems.<sup>27</sup> Fast relaxation is more  
10 feasible because of the reduced out-of-plane ion motion expected from the insulating organic  
11 layers; the associate barrier to diffusion would keep the product interstitial halide physically  
12 close to the vacancy and ready for the back reaction. Considering all these observations, we  
13 see little clear evidence of light-induced vacancy formation in BAPI, in contrast with the 3D  
14 perovskites.<sup>9,11</sup>

## 15 **6 Conclusion**

16 We have shown that the transfer function representation of SPM used to describe broadband  
17 local dielectric spectroscopy and dissipation measurements can be extended to write the tem-  
18 poral response measured in tr-EFM and pk-EFM in terms of sample resistance  $R_s$ .<sup>14</sup> We have  
19 established experimentally using a range of steady-state and time-resolved studies that elec-  
20 trical scanned probe measurements in BAPI are primarily observing changes in sample resis-  
21 tance and not capacitance. We find, surprisingly, that light-induced conductivity in the 2D BAPI  
22 system recovers  $10^5$  faster than in many 3D perovskite systems studied to date.<sup>9,11</sup>

23 **SUPPORTING INFORMATION AVAILABLE** The Supporting Information contains:  
24 theoretical derivation for Eq. 1; numerical simulations for frequency and phase shift in tr-  
25 EFM and pk-EFM for variety of experimental conditions;<sup>31</sup> timing diagram for Fig. 4; dielectric

<sup>1</sup> response at  $\omega_m = 1257\text{ s}^{-1}$ ; absorbance spectra; XRD; AFM topography image; surface potential  
<sup>2</sup> at selected light intensities.

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#### <sup>10</sup> Notes

<sup>11</sup> The authors declare no competing financial interest.

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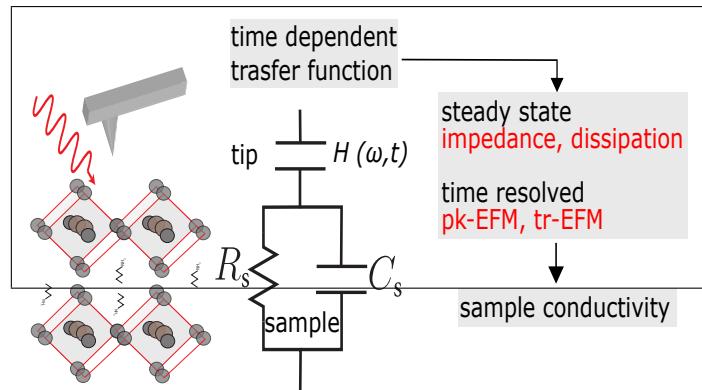
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# 1 Graphical TOC Entry

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**Supporting Information: Light-dependent Impedance Spectra and  
Transient Photoconductivity in a Ruddlesden-Popper 2D Lead-halide  
Perovskite Revealed by Electrical Scanned Probe Microscopy and  
Accompanying Theory**

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## S1. THEORY: DERIVATION OF EQ. 1

To derive Eq. 1, we start from the equation for the phase shift  $\Delta\phi$  induced by the tip-sample force  $F_{\text{ts}}$  [1],

$$\Delta\phi = -\frac{f_c}{k_c A_0^2} \int_0^t F_{\text{ts}}(t') x(t') dt', \quad (\text{S1})$$

where  $f_c$  is the cantilever resonance frequency,  $k_c$  is the cantilever spring constant,  $A_0$  is the cantilever's initial amplitude at  $t = 0$ , and  $x(t)$  is the cantilever's displacement versus time. Following the theory introduced in Ref. 2, the tip-sample force, tip-sample charge, and cantilever displacement are approximated using perturbation theory. We use the notation  $x^{(0)}$  to represent the zeroth-order approximation for  $x(t)$ . From Ref. 2, the zeroth order tip-sample force is

$$F_{\text{ts}}^{(0)}(t) = \frac{1}{2} C'_{\text{tip}} \frac{\left[ q_{\text{t}}^{(0)}(t) \right]^2}{C_{\text{tip}}^2} = \frac{1}{2} C'_{\text{tip}} V_{\text{t}}(t)^2$$

where  $q^{(0)}(t)$  is the zeroth order tip voltage and  $V_{\text{t}}$  is the zeroth order tip voltage

$$V_{\text{t}} = \frac{q_{\text{t}}^{(0)}(t)}{C_{\text{tip}}}. \quad (\text{S2})$$

The zeroth order tip-sample force only causes a phase shift if  $V_{\text{t}}^2$  contains significant content at the cantilever frequency  $\omega_c$ . For now, we assume that  $V_{\text{t}}^2$  varies slowly enough that it does not cause a phase shift.

The first-order tip-sample force  $F_{\text{ts}}^{(1)}$  describes the oscillating forces caused by the oscillating tip. These oscillating forces cause the frequency shift in most KPFM experiments. In the perturbation theory approximation, the first-order tip-sample force is

$$F_{\text{ts}}^{(1)} = \frac{C''_q q_{\text{t}}^{(0)} q_{\text{t}}^{(0)} x^{(0)}}{2C_{\text{tip}}^2} + \frac{C'_{\text{tip}} q_{\text{t}}^{(0)} q_{\text{t}}^{(1)}}{C_{\text{tip}}^2} \quad (\text{S3})$$

where the capacitance derivative  $C''_q$  is

$$C''_q = C''_{\text{tip}} - 2(C'_{\text{tip}})^2/C_{\text{tip}}. \quad (\text{S4})$$

The two terms of Eq. S3 describe the two possible causes of the oscillating tip-sample force. In the first term, the cause is oscillations in the tip-sample energy arising from the oscillating tip displacement (at constant charge). In the second term, the cause is the oscillating charge that

flows in response to the oscillating tip displacement. The resulting phase shift is obtained by substituting back into Eq. S1

$$\Delta\phi = -\frac{f_c}{k_c A_0^2} \int_0^t \left[ \frac{C_q'' q_t^{(0)} q_t^{(0)} x^{(0)}}{2C_{\text{tip}}^2} + \frac{C'_{\text{tip}} q_t^{(0)} q_t^{(1)}}{C_{\text{tip}}^2} \right] x(t') dt'.$$

We approximate  $x(t)$  by its zeroth-order approximation  $x^{(0)}(t)$ . In these experiments, the cantilever is excited at its resonance frequency so  $x^{(0)}(t) = A_0 \cos(\omega_c t + \phi_0)$ . For simplicity, we take the initial cantilever phase to be  $\phi_0 = 0$ ; none of the conclusions below depend on this assumption. After plugging in for  $x(t)$ , the phase shift simplifies to

$$\Delta\phi = \underbrace{-\frac{f_c}{k_c} \int_0^t \frac{C_q'' q_t^{(0)} q_t^{(0)}}{2C_{\text{tip}}^2} \cos^2(\omega_c t') dt'}_{\text{constant charge term}} - \underbrace{\frac{f_c}{k_c A_0} \int_0^t \frac{C'_{\text{tip}} q_t^{(0)} q_t^{(1)}}{C_{\text{tip}}^2} \cos(\omega_c t') dt'}_{\text{oscillating charge term}}. \quad (\text{S5})$$

To determine the cantilever phase shift during these experiments, we consider the two terms in Eq. S5 individually. The first term ( $\Delta\phi_{\text{const}}$ ) describes the phase shift caused by the tip oscillating at constant charge. The second term ( $\Delta\phi_{\text{osc}}$ ) describes an additional phase shift caused by the tip charge oscillating along with the oscillating tip displacement. Using the definition of  $V_t$  (Eq. S2) and the trigonometric identity for  $\cos^2 \theta$ , the first term simplifies to

$$\Delta\phi_{\text{const}} = -\frac{f_c}{4k_c} C_q'' \int_0^t V_t(t')^2 + \underbrace{V_t(t')^2 \cos(2\omega_c t)}_{\text{under-braced term}} dt'. \quad (\text{S6})$$

If  $V_t(t)$  varies slowly, the under-braced term integrates to zero over each cantilever oscillation period. In this case, the overall frequency shift  $\Delta f = d\Delta\phi/dt$  caused by the first term is

$$\Delta f_{\text{const}}(t) = -\frac{f_c}{4k_c} C_q'' V_t(t)^2.$$

For the second term, we take the same basic approach. The second term simplifies to

$$\Delta\phi_{\text{osc}} = -\frac{f_c C'_{\text{tip}}}{k_c A_0 C_{\text{tip}}} \int_0^t V_t q_t^{(1)} \cos(\omega_c t') dt' \quad (\text{S7})$$

We need the first-order tip-charge  $q_t^{(1)}$ . In the time-independent case, this was given by

$$q_t^{(1)} = C_{\text{tip}} h * V_x = C_{\text{tip}} \int_{t_0}^t h(t-t') V_x(t') dt', \quad (\text{S8})$$

where  $h$  is the impulse response function between the applied tip-sample voltage  $V_{\text{ts}}(t)$  and the voltage that drops across the tip capacitor  $V_t = q_t^{(0)}/C_{\text{tip}}$ ,  $*$  denotes convolution, and  $V_x$  is

the oscillating voltage induced by the oscillating tip displacement,  $V_x = C'_{\text{tip}} q_t^{(0)} x^{(0)} / C_{\text{tip}}^2$ . In the time-dependent case, the convolution integral is replaced by the generalized convolution

$$q_t^{(1)} = C_{\text{tip}} \int_{t_0}^t h(t, t') V_x(t') dt'.$$

The time-varying impulse response function  $h(t, t')$  describes the response at time  $t$  to an impulse applied at a time  $t'$ . Substituting in for  $V_x$ , we obtain

$$q_t^{(1)} = \int_{t_0}^t h(t, t') \frac{q_t^{(0)}(t')}{C_{\text{tip}}} C'_{\text{tip}} x^{(0)}(t') dt'.$$

Substituting in for  $x^{(0)}$  and simplifying gives

$$q_t^{(1)} = A_0 C'_{\text{tip}} \int_{t_0}^t h(t, t') V_t(t') \cos(\omega_c t') dt'. \quad (\text{S9})$$

In the experiments we consider, the spectral content of  $V_t$  will be concentrated at low frequencies. Looking back to Eq. S7, the integral will only have non-zero values over a cantilever period when the response is in phase with the oscillating position. Therefore, to a good approximation, we need the component of the oscillating tip charge at the cantilever frequency  $\omega_c$ . Since the spectral content of  $V_t$  is concentrated at low frequencies, this will be given by the *time-varying* response function  $H(\omega_c, t)$

$$q_t^{(1)} \approx A_0 C'_{\text{tip}} V_t(t) [\text{Re}[H(\omega_c, t)] \cos(\omega_c t) - \text{Im}[H(\omega_c, t)] \sin(\omega_c t)] \quad (\text{S10})$$

Substituting this equation for  $q_t^{(1)}$  into Eq. S7, we obtain

$$\Delta\phi_{\text{osc}} = -\frac{f_c (C'_{\text{tip}})^2}{k_c C_{\text{tip}}} \int_0^t V_t(t')^2 [\cos^2(\omega_c t') \text{Re}[H(\omega_c, t')] + \sin(\omega_c t') \cos(\omega_c t') \text{Im}[H(\omega_c, t')]] dt'$$

Simplifying using the double angle formulas and the definition  $\Delta C'' = 2(C'_{\text{tip}})^2 / C_{\text{tip}}$ ,

$$\Delta\phi_{\text{osc}} = -\frac{f_c \Delta C''}{4k_c} \int_0^t V_t(t')^2 \left[ \underbrace{\text{Re}[H(\omega_c, t')]}_{+ \cos(2\omega_c t') \text{Re}[H(\omega_c, t')]} + \sin(2\omega_c t') \text{Im}[H(\omega_c, t')] \right] dt'. \quad (\text{S11})$$

If  $V_t^2$  varies slowly and  $H(\omega_c, t)$  is linear within each cantilever period, then only the underbraced term contributes to the integral over a cantilever oscillation period. In this case, the frequency shift  $\Delta f = d\Delta\phi/dt$  is

$$\Delta f_{\text{osc}}(t) = -\frac{f_c \Delta C''}{4k_c} V_t(t)^2 \text{Re}[H(\omega_c, t)].$$

Adding the two terms, we obtain the overall frequency shift given in Eq. 1:

$$\Delta f(t) = -\frac{f_c}{4k_c} \left( C_q'' + \Delta C'' \operatorname{Re}[H(\omega_c, t)] \right) V_t(t)^2.$$

The zeroth order tip voltage  $V_t(t)$  is

$$V_t(t) = \int_{t_0}^t h(t, t') V_{ts}(t') dt'. \quad (\text{S12})$$

The two key approximations made in this derivation were (1)  $V_t$  and  $V_t^2$  do not have significant spectral content at the cantilever resonance frequency  $\omega_c$  and (2) the real and imaginary components of  $H(\omega_c, t)$  are linear over each cantilever period.

## S2. DETERMINING THE TIME-VARYING RESPONSE FUNCTION

The time-varying response function is the Fourier transform of the modified time-varying impulse response  $\bar{h}(t, \tau)$ , where  $\tau = t - t'$  is the delay between the time  $t$  at which the response is measured and the time  $t'$  at which the impulse was applied. The modified time-varying impulse response is

$$\bar{h}(t, \tau) = h(t, t - \tau) \quad (\text{S13})$$

and the Fourier transform is

$$H(\omega, t) = \int_{-\infty}^{\infty} \bar{h}(t, \tau) e^{-j\omega\tau} d\tau$$

Next we describe how  $H(\omega_c, t)$  can be calculated for the single parallel sample resistance and capacitance model used in the text ( $R_s \parallel C_s$ , circuit shown in Fig. S1a).

In this case, the differential equations describing the evolution of the tip charge  $q_t$  can be expressed in terms of  $q_R$ , where  $\dot{q}_R$  is the current through the sample resistance  $R_s$ . The state variable  $q_R$  is described by the differential equation

$$\dot{q}_R = -\omega_{\text{fast}}(t) q_R + C_{\text{tip}} \omega_{\text{fast}}(t) V_{ts}(t) \quad (\text{S14})$$

where  $\omega_{\text{fast}}(t) = (R_s(t)(C_s + C_{\text{tip}}))^{-1}$ . To zeroth order, the dependence of the tip capacitance on distance is negligible and the system is linear. The system is time-varying through the time-dependence of the resistance  $R_s$ . The tip charge and tip voltage are

$$q_t = \left( \frac{C_{\text{tip}}}{C_{\text{tip}} + C_s} \right) q_R + \left( \frac{C_{\text{tip}} C_s}{C_{\text{tip}} + C_s} \right) V_{ts}(t) \quad (\text{S15})$$

$$V_t = \left( \frac{1}{C_{\text{tip}} + C_s} \right) q_R + \left( \frac{C_s}{C_{\text{tip}} + C_s} \right) V_{ts}(t) \quad (\text{S16})$$

The propagator (or state transition matrix)  $\Phi$  describes the evolution of the state variable, with  $q_R(t) = \Phi(t, t_0)q_R(t_0)$  in the absence of a voltage input. For this system, the propagator is

$$\Phi(t, t') = e^{-\int_{t'}^t \omega_{\text{fast}}(\theta) d\theta}. \quad (\text{S17})$$

The time-varying impulse response<sup>1</sup> gives the response of the tip voltage at a time  $t$  to a voltage input applied at time  $t'$

$$h(t, t') = \left( \frac{C_{\text{tip}}}{C_s + C_{\text{tip}}} \right) \omega_{\text{fast}}(t') \Phi(t, t') u(t - t') + \left( \frac{C_s}{C_s + C_{\text{tip}}} \right) \delta(t - t'),$$

where  $u(t)$  is the Heaviside step function ( $u(t) = 0$  for  $t < 0$ ,  $u(t) = 1$  for  $t > 0$ ) and  $\delta(t)$  is the Dirac delta function. The modified time varying impulse response is given by defining the delay  $\tau = t - t'$ ,

$$\bar{h}(t, \tau) = \left( \frac{C_{\text{tip}}}{C_s + C_{\text{tip}}} \right) \omega_{\text{fast}}(t - \tau) \Phi(t, t - \tau) u(\tau) + \left( \frac{C_s}{C_s + C_{\text{tip}}} \right) \delta(\tau).$$

The sought-after time-varying frequency response is the Fourier transform of  $\bar{h}$  with respect to  $\tau$ <sup>2</sup>

$$\begin{aligned} H(\omega, t) &= \int_{-\infty}^{\infty} \left[ \left( \frac{C_{\text{tip}}}{C_s + C_{\text{tip}}} \right) \omega_{\text{fast}}(t - \tau) \Phi(t, t - \tau) u(\tau) + \left( \frac{C_s}{C_s + C_{\text{tip}}} \right) \delta(\tau) \right] e^{-j\omega\tau} d\tau, \\ &= \int_0^{\infty} \left( \frac{C_{\text{tip}}}{C_s + C_{\text{tip}}} \right) \omega_{\text{fast}}(t - \tau) \Phi(t, t - \tau) e^{-j\omega\tau} d\tau + \int_{-\infty}^{\infty} \left( \frac{C_s}{C_s + C_{\text{tip}}} \right) \delta(\tau) e^{-j\omega\tau} d\tau. \end{aligned}$$

Simplifying the second integral and substituting the expression for  $\Phi$  from Eq. S17, we obtain Eq. 3. In words,  $H(\omega, t)$  gives the response of  $V_t$  to an applied external voltage  $e^{j\omega t}$ . In the limit that  $\omega_{\text{fast}}$  is constant,  $H(\omega, t)$  becomes time-independent and reduces to  $H(\omega)$  given by Eq. 2.

### S3. TR-EFM AND PK-EFM SIMULATIONS

The approximation for  $\Delta f$  given by Eq. 1 was compared against the results of numerical simulations of the cantilever's dynamics for tr-EFM and pk-EFM experiments. To simulate a tr-EFM or pk-EFM experiment in which the light is turned on at  $t = 0$ , the sample resistance was taken to respond to light with a time constant  $\tau_L$ ,

$$R_s(t) = \begin{cases} R_{\text{dark}} & t \leq 0 \\ R_{\text{dark}} + (R_{\text{light}} - R_{\text{dark}})(1 - e^{-t/\tau_L}) & t > 0, \end{cases} \quad (\text{S18})$$

---

<sup>1</sup> From Kailath Ch. 9, Eq. 23 [3],  $h(t, t') = C(t)\Phi(t, t')B(t') + D(t)\delta(t - t')$ , where  $B, C, D$  have their usual definitions for state space representations.

<sup>2</sup> From Shmaliy Eq. 6.31 [4],  $H(\omega, t) = \bar{H}(\omega, t)$ .

where  $R_{\text{dark}}$  is the sample's dark resistance and  $R_{\text{light}}$  is the sample's final resistance after the light has been on for a long time.

To simulate a tr-EFM experiment, the tip-sample voltage was left constant  $V_{\text{ts}}(t) = V$ . To simulate a pk-EFM experiment, the tip-sample voltage  $V_{\text{ts}}$  was stepped back to zero after a time  $t_p$

$$V_{\text{ts}}(t) = \begin{cases} V & t \leq t_p \\ 0 & t > t_p. \end{cases} \quad (\text{S19})$$

In our simulations, the tip-sample voltage was always  $V = 10\text{V}$ .

In Eq. 1, the two factors that affect the frequency shift are the zeroth order tip voltage  $V_t$  and the time-varying transfer function  $H(\omega_c, t)$  at the cantilever frequency  $\omega_c$ . Figure S1 illustrates the effect of each factor for a sample where the light decreases the sample resistance from  $R_{\text{dark}} = 100\text{G}\Omega$  to  $R_{\text{light}} = 10\text{M}\Omega$  with an exponential risetime of  $\tau_L = 10\mu\text{s}$ . These values were chosen to clearly illustrate how each factor affects the frequency shift.

The experiment begins when the applied voltage  $V_{\text{ts}}$  is switched from 0 to 10V. Initially, equal charges build up on both capacitors and the voltage drop across the tip capacitor is  $V_t = V_{\text{ts}}C_s/(C_s + C_{\text{tip}})$  (Fig. S1a). The sample capacitor discharges with an  $RC$  time constant  $R_{\text{dark}}(C_{\text{tip}} + C_s) = 1.1\text{ms}$ , so that eventually, all of the applied voltage drops across the tip capacitor (Fig. S1b). When the light is turned on, the sample resistance  $R_s$  decreases by 6 orders of magnitude, decreasing the  $RC$  time constant from 1.1 ms to 1.1 ns.

Dramatic differences in the frequency shift are observed depending on when the light is turned on. Fig. S1c shows the tip voltage versus time after the light is turned on ( $t = 0$ ) when the tip is only partially charged (blue), and, for comparison, when the tip is fully charged (orange). If the tip is only partially charged, the light-induced decrease in sample resistance speeds up the charging of the tip capacitor—the tip charges fully in 200  $\mu\text{s}$  compared to the 3 ms it would take in the dark. If allowed to charge fully before turning on the light,  $V_t$  remains constant (orange).

Figure S1d shows the resulting frequency shift versus time signal calculated from numerical simulations (points) and approximated using Eq. 1 (lines). According to Eq. 1, the frequency shift  $\Delta f$  is proportional to  $V_t^2$ , so the blue trace and circles show an increase in frequency shift from  $t = 0$  to 200  $\mu\text{s}$  as the tip charges. The messy oscillations in the numerically simulated frequency-shift transient (blue circles) occur because the tip charge and dc displacement are

changing rapidly on the timescale of the cantilever period, so the cantilever frequency shift is poorly defined during these periods. The phase shift is still well-defined, however, and Fig. S1d shows that Eq. 1 accurately captures the numerically simulated phase shift remarkably well.

The increase in frequency shift after 200  $\mu$ s is caused by an increase in the time-varying transfer function at the cantilever frequency ( $H(\omega_c, t)$ ). Physically, as the circuit's  $RC$  time constant drops below the cantilever's inverse resonance frequency  $\omega_c^- 1 = 2.5 \mu$ s, more of the tip charge is able to oscillate on and off the tip during motion, so that the oscillating force in-phase with the tip displacement increases. The predicted frequency and phase shift agrees closely with the numerical simulations.

All of these light and bias-dependent effects can be understood using the time-varying transfer function  $H(\omega, t)$ . Fig. S1f plots the real part of  $H$  versus angular frequency at 40  $\mu$ s intervals after the light is turned on at  $t = 0$ . Over the first 160  $\mu$ s after the light is turned on, the main effect of the reduced sample resistance is an increase in  $H$  at frequencies well below the cantilever frequency  $\omega_c$ . This low-frequency response can affect the frequency shift through the tip voltage  $V_t$ , which is the generalized convolution of the time-varying impulse response function and the applied tip-sample voltage (Eq. S12). Notice that the increase in  $H$  at low-frequencies has no effect if all of the tip-sample voltage is already dropping across the tip, as for the orange traces in Fig. S1(c,d). In this case, light-induced frequency shifts can be attributed to changes in  $\text{Re } H(\omega_c, t)$  alone. The experimental data of Fig. 4 shows this same dependence of the light-induced frequency shift on the initial tip charge.

While the effect of the increase in  $H$  at low-frequencies depends on the initial tip charge, a change in  $\text{Re } H(\omega_c, t)$  directly affects the cantilever frequency shift. Between 160  $\mu$ s and 280  $\mu$ s, the real part of the transfer function  $\text{Re } H(\omega_c, t)$  increases dramatically at the cantilever frequency (Fig. S1a). The cantilever frequency shift (Fig. S1c) increases further as a result.

Simulations were performed for a wide variety of experimental parameters, designed to cover regimes where each factor ( $V_t$  and  $\text{Re } H(\omega_c, t)$ ) influences the frequency shift. For sample tr-EFM experiments, Figure S2 shows that Eq. 1 is a good approximation of the cantilever frequency shift.

To assess pk-EFM experiments, the measured and predicted phase shifts were compared for 1080 simulated experiments. Figure S3 shows that over a wide variety of sample and cantilever parameters, the predicted and measured phase shift in pk-EFM experiments agreed closely. The residuals  $r = \Delta\phi_{\text{simulated}} - \Delta\phi_{\text{predicted}}$  had a mean and standard deviation of 0.027 mcy and

0.049 mcyc respectively. The maximum residual was 0.20 mcyc.

### S3A. Numerical simulations

The numerical simulations used to generate Figs. S2– S3 were performed in Python. The code for these simulations is publicly available [9]. The coupled first-order ordinary differential equations that described the cantilever dynamics are

$$\dot{x} = v \quad (\text{S20a})$$

$$\dot{v} = -\omega_c^2 x - \frac{\omega_c}{Q} v + \frac{C'_{\text{tip}}(x)}{2mC_{\text{tip}}(x)^2} q_t^2 \quad (\text{S20b})$$

$$\dot{q}_R = -\frac{1}{R_s(t)(C_{\text{tip}}(x) + C_s)} q_R + \frac{C_{\text{tip}}(x)}{R_s(t)(C_{\text{tip}}(x) + C_s)} V_{\text{ts}}(t). \quad (\text{S20c})$$

where  $v$  is the cantilever velocity. For convenient comparison to the model developed above, the zeroth order charge was also computed using

$$\dot{q}_R^{(0)} = -\frac{1}{R_s(t)(C + C_s)} q_R^{(0)} + \frac{C}{R_s(t)(C + C_s)} V_{\text{ts}}(t), \quad (\text{S21})$$

where  $C = C_{\text{tip}}(0)$ . The tip capacitance and its derivatives were given by  $C_{\text{tip}}(x) = C + C'x + C''x^2/2$ . In Eq. S20b, the tip charge  $q_t$  was calculated using Eq. S15. The numerical integration was performed in Python using Scipy's odeint function, storing the state vector  $\mathbf{y} = (x \ v \ q_R \ q_R^{(0)})^T$  every 1  $\mu\text{s}$ .

The initial state vector for the cantilever was

$$\mathbf{y}_0 = \begin{pmatrix} x \\ v \\ q_R \\ q_R^{(0)} \end{pmatrix} = \begin{pmatrix} A_0 \cos(\omega_c t_0 + \phi_0) \\ A_0 \omega_c \sin(\omega_c t_0 + \phi_0) \\ \beta CV \\ \beta CV \end{pmatrix},$$

where  $A_0$  is the cantilever's initial amplitude,  $\phi_0$  is the cantilever's initial phase,  $t_0 = -100 \mu\text{s}$  is the initial time the numerical integration was started,  $\beta$  is the initial fraction of the tip sample voltage that drops across the tip capacitor, and the applied tip-sample voltage  $V = 10 \text{ V}$  (Eq. S19).

The output of the numerical integration was used to determine the cantilever's amplitude, phase, and frequency. The cantilever amplitude  $A = |z|$  and phase  $\phi = \arg z$  were calculated using the complex number

$$z = (x - x_{\text{eq}}) - j v / \omega_c. \quad (\text{S22})$$

where  $x_{\text{eq}} = C'_{\text{tip}} V_{\text{ts}}^2 / (2k_c)$ . The amplitude and phase data were filtered by averaging over a single cantilever period (16 data points). The frequency shift was calculated by numerically differentiating the filtered phase shift (using second order central differences via numpy's gradient function).

To compare the impedance theory approximation to the numerical integration, the zeroth order tip voltage  $V_t$  and the time-varying frequency response  $H(\omega, t)$  must be known. The zeroth order tip voltage was calculated using Eq. S16 with  $q_R = q_R^{(0)}$ . To approximate  $H(\omega, t)$ , the double integral of Eq. 3 was broken into pieces depending on the value of  $t$  and  $\omega_{\text{fast}}(t)$ . When  $\omega_{\text{fast}} \ll \omega_c$ , the necessary integral is oscillatory and decays slowly which makes converging a numerical approximation difficult. In this case, the integrand was sampled at points equally spaced through the cantilever cycle (160 points per cycle) and Simpson's rule was used to compute the value of the integral.

#### S4. SCANNING PROBE MICROSCOPY

The scanning probe microscopy set up used to perform different measurements here has been described in our previous reports [5]. Cantilever motion was detected using a fiber interferometer operating at 1490 nm (Corning SMF-28 fiber). The laser diode's (QPhotonics laser diode QFLD1490-1490-5S) dc current was set using a precision current source (ILX Lightwave LDX-3620), and the current was modulated at radio frequencies using the input on the laser diode mount (ILX Lightwave LDM 4984, temperature-controlled with ILX Lightwave LDT-5910B). The interferometer light was detected with a 200-kHz bandwidth photodetector (New Focus model 2011, built-in high-pass filter set to 200 kHz) and digitized at 1 MHz (National Instruments, PCI-6259). The cantilever was driven using a commercial PLL cantilever controller (RHK Technology, PLLPro2 Universal AFM controller) with PLL feedback loop integral gain  $I = 2.5$  Hz, proportional gain  $P = -5$  ° Hz. The sample was illuminated from above with a fiber-coupled 405 nm laser (Thorlabs model LP405-SF10, held at 25 °C with a Thorlabs model TED200C temperature controller). The laser current was controlled using the external modulation input of the laser's current controller (Thorlabs model LDC202, 200 kHz bandwidth). The light was coupled to the sample through a multimode, 50 μm diameter core, 0.22 NA optical fiber (Thorlabs model FG050LGA). The intensity at sample surface was calculated based on an estimated spot size of  $\approx 0.26$  mm<sup>2</sup>

Implementation of broad band local dielectric spectroscopy has been described previously in Refs. 6 and 7. The procedure is reproduced below for reference. For amplitude modulation BLDS (Fig. 2b), we applied a time-dependent voltage to the cantilever tip :

$$V_m(t) = V_{pp} \left( \frac{1}{2} + \frac{1}{2} \cos(2\pi f_{am} t) \right) \cos(2\pi f_m t). \quad (S23)$$

In the experiments reported in the manuscript,  $f_m = 200\text{Hz}$  to  $1.5\text{MHz}$ ,  $f_{am} = 45\text{Hz}$ , and the amplitude was set to  $V_{pp} = 6\text{V}$ . The time-dependent voltage in Equation S23 was generated using a digital signal generator (Keysight 33600). The cantilever frequency shift was measured in real time using a phase-locked loop (PLL; RHK Technology, model PLLPro2 Universal AFM controller), the output of which was fed into a lock-in amplifier (LIA; Stanford Research Systems, model 830). The LIA time constant and filter bandwidth were  $300\text{ms}$  and  $6\text{dB/oct}$ , respectively. At each stepped value of  $f_m$ , a wait time of  $1500\text{ms}$  was employed, after which frequency-shift data were recorded for an integer number of frequency cycles corresponding to  $\approx 2\text{ sec}$  of data acquisition at each  $f_m$ . The measurable  $\Delta f_{BLDS}$  primarily probes the response at  $\omega_m$  (Eq. S24).

$$\Delta f_{BLDS}(\omega_m) = -\frac{f_c V_m^2}{16k} \left[ C_q'' + \Delta C'' \operatorname{Re} \left( \hat{H}(\omega_m + \omega_c) + \hat{H}(\omega_m - \omega_c) \right) \right] |\hat{H}(\omega_m)|^2 \quad (S24)$$

$\Delta f_{BLDS}$  is related to the plotted voltage-normalized frequency shift  $\alpha$  by Eq. S25.

$$\alpha = \frac{\Delta f_{BLDS}(\omega_m)}{V_m^2}. \quad (S25)$$

The  $\Delta f_{BLDS}$  frequency-shift signal was obtained from the LIA outputs as follows. From the (real) in-phase and out-of-phase voltage signals  $V_X$  and  $V_Y$ , respectively, a single (complex signal) in hertz was calculated using the formula

$$Z_{\text{Hz}} = (V_x + j V_y) \frac{S}{10} \times \sqrt{2} \times 20 \frac{\text{Hz}}{\text{V}} \quad (S26)$$

From  $Z_{\text{Hz}}$  we calculate  $\alpha$

$$\alpha = \frac{4|Z_{\text{Hz}}|}{V_{pp}^2} = \frac{8\sqrt{2}S\sqrt{V_x^2 + V_y^2}}{V_{pp}^2}. \quad (S27)$$

In frequency shift BLDS (Fig. 2a) described in Ref. 7, the applied waveform is not amplitude modulated at  $45\text{Hz}$  and instead an equal period ON/OFF amplitude-modulating is applied. The resultant frequency shift is calculated by software demodulation of the cantilever response by subtracting the average frequency shift during the ON period from the OFF period.

#### S4A. tr-EFM and pk-EFM

Implementation and data work up details for tr-EFM and pk-EFM have been extensively described in Ref. 5 and Ref. 8 and is reproduced briefly below for reference. A commercial pulse and delay generator (Berkeley Nucleonics, BNC565) was used to generate tip voltage and light modulation pulses, as well as to turn off the cantilever drive voltage. The PLLPro2 generated the cantilever drive voltage and also output a 1 V, phase-shifted sine wave copy of the cantilever oscillation, generated by an internal lock-in amplifier coupled to the phase-locked loop. A home built gated cantilever clock circuit converted this 1 V phase-shifted sine wave to a square wave, which was used as a clock for timing tip voltage and light pulses. A 5 V digital signal output by the National Instruments PCI-6259 gates the clock, controlling the start of the experiment. The BNC565 was used to trigger all signals relative to the cantilever clock. Cantilever drive was switched off (2 ms to 10 ms) before the start of light pulse. The raw cantilever oscillation data (digitized at 1 MHz) was saved along with counter timings (PCI-6259, 80-MHz counter) indicating the precise starting time of the light pulse (synchronized to the cantilever oscillation), allowing the start of the light pulse to be determined to within 12.5 ns. Along with each pk-EFM phase shift data point, a control data point, identical except without turning on the light, was collected.

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## **FIGURES**

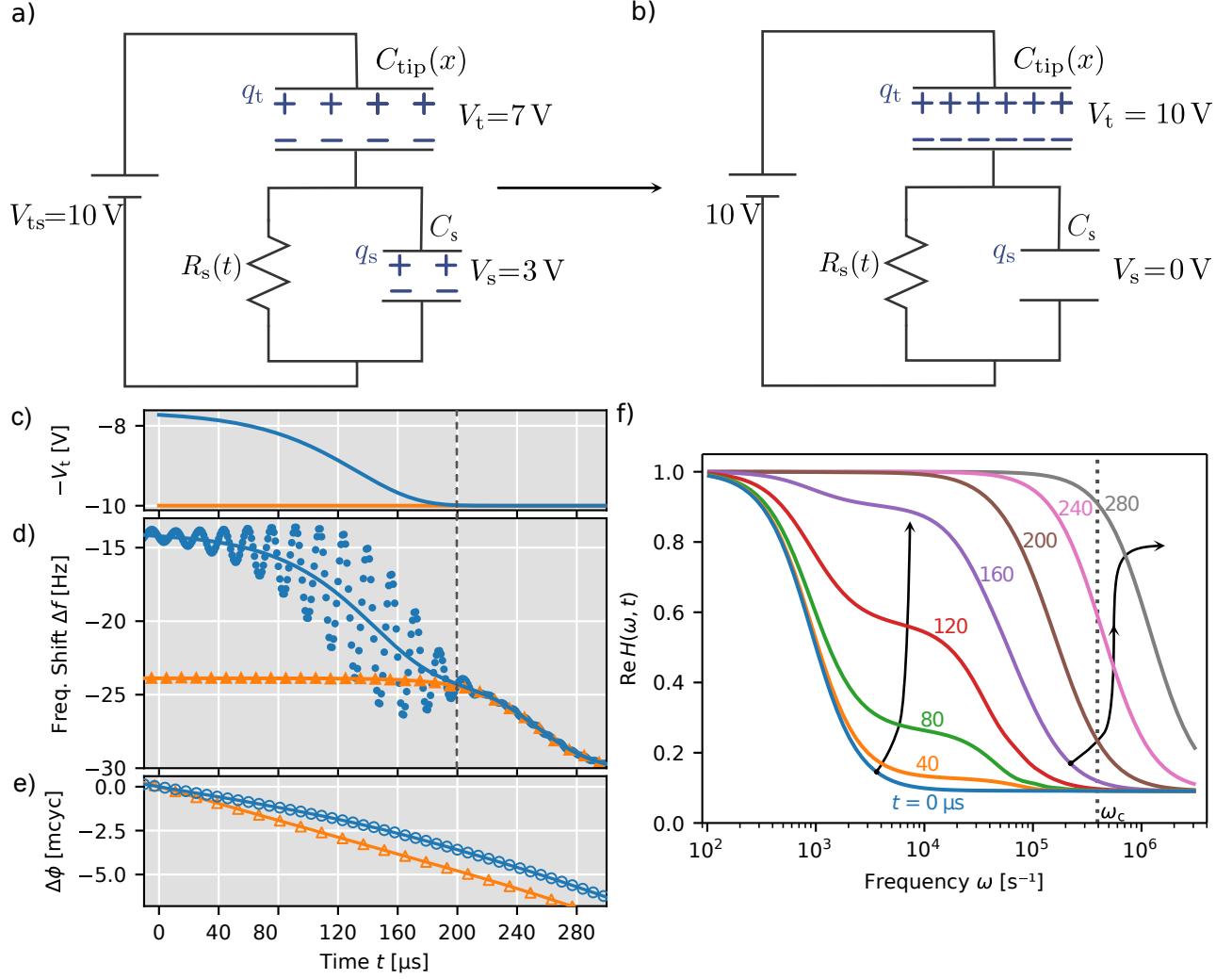


FIG. S1. Frequency shift and time-varying transfer function versus time for a representative pk-EFM experiment. (a) The model circuit used in this simulation. After the tip-sample voltage is applied, charge builds up on the tip and sample capacitors. (b) Eventually, the sample capacitor discharges and  $V_t = V_{ts}$ . The (c) negative zeroth order tip voltage  $-V_t$ , (d) frequency shift, and (e) phase shift calculated using numerical simulations (points) and analytically using Eq. 1 (line). For (c)–(e),  $V_t(t_0) = 7\text{ V}$  (blue circles), 10V (orange triangles). (f) The time-varying response function  $H(\omega, t)$  shown at times  $t = 0, 40, 80, \dots, 280\text{ }\mu\text{s}$ . Constant experimental parameters:  $R_{\text{dark}} = 10\text{ T}\Omega$ ,  $R_{\text{light}} = 10\text{ M}\Omega$ ,  $\tau_L = 40\text{ }\mu\text{s}$ ,  $k_c = 3.5\text{ N m}^{-1}$ ,  $f_c = 62\text{ kHz}$ ,  $Q = 26000$ ,  $C_{\text{tip}} = 1 \times 10^{-4}\text{ pF}$ ,  $C'_{\text{tip}} = -2.8 \times 10^{-5}\text{ pF }\mu\text{m}^{-1}$ ,  $C_s = 1 \times 10^{-5}\text{ pF}$ ,  $C''_{\text{tip}} = 6.77 \times 10^{-5}\text{ pF }\mu\text{m}^{-2}$ ,  $V_{ts} = 10\text{ V}$ ,  $t_0 = -200\text{ }\mu\text{s}$ .

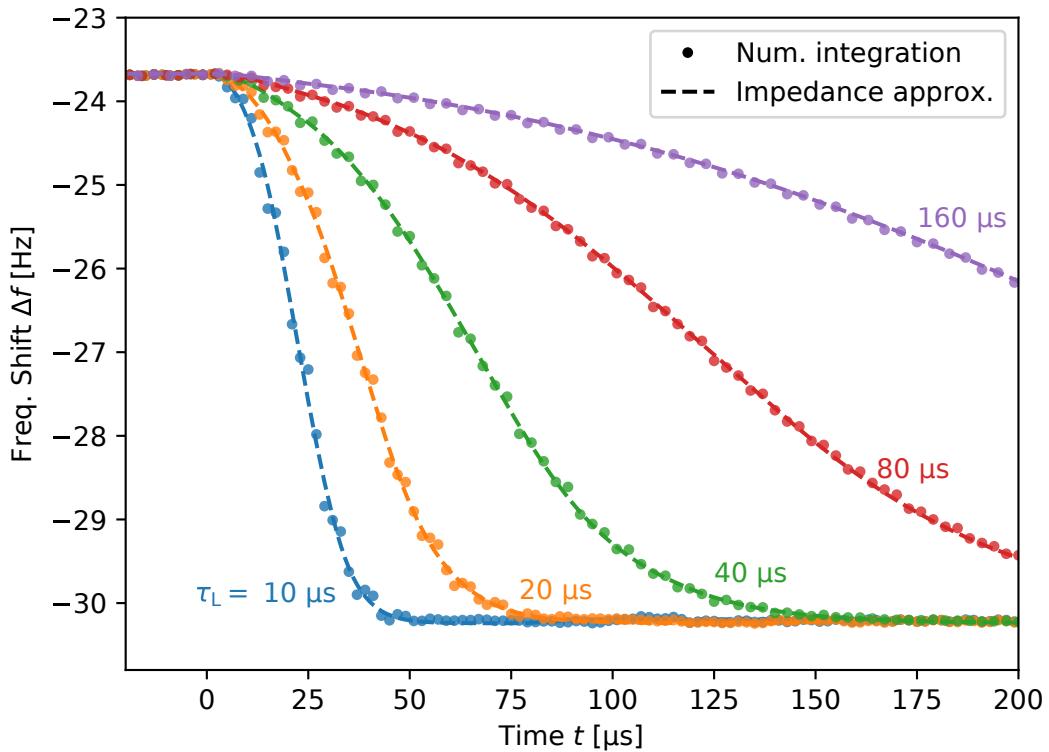


FIG. S2. Frequency shift versus time for representative simulated EFM experiments. For different light-induced time constants  $\tau_L$  (Eq. S18), points show the frequency shift calculated from numerical integration of Eq. S20 and dashed lines show the results calculated using the approximation in Eq. 1. For each trace, the following experimental parameters were held constant:  $f_c = 62.5\text{ kHz}$ ,  $k_c = 3.5\text{ Nm}^{-1}$ ,  $Q = 26000$ ,  $A_0 = 50\text{ nm}$ ,  $\phi_0 = \pi$ ,  $t_0 = -100\mu\text{s}$ ,  $R_{\text{dark}} = 100\text{ G}\Omega$ ,  $R_{\text{light}} = 10\text{ M}\Omega$ ,  $C_s = 0$ ,  $C_{\text{tip}} = 1 \times 10^{-4}\text{ pF}$ ,  $C_q'' = 52.1\text{ pF nm}^{-2}$ ,  $\Delta C'' = 15.7\text{ pF nm}^{-2}$ ,  $V_t(t_0) = 10\text{ V}$ ,  $V_{\text{ts}} = 10\text{ V}$ .

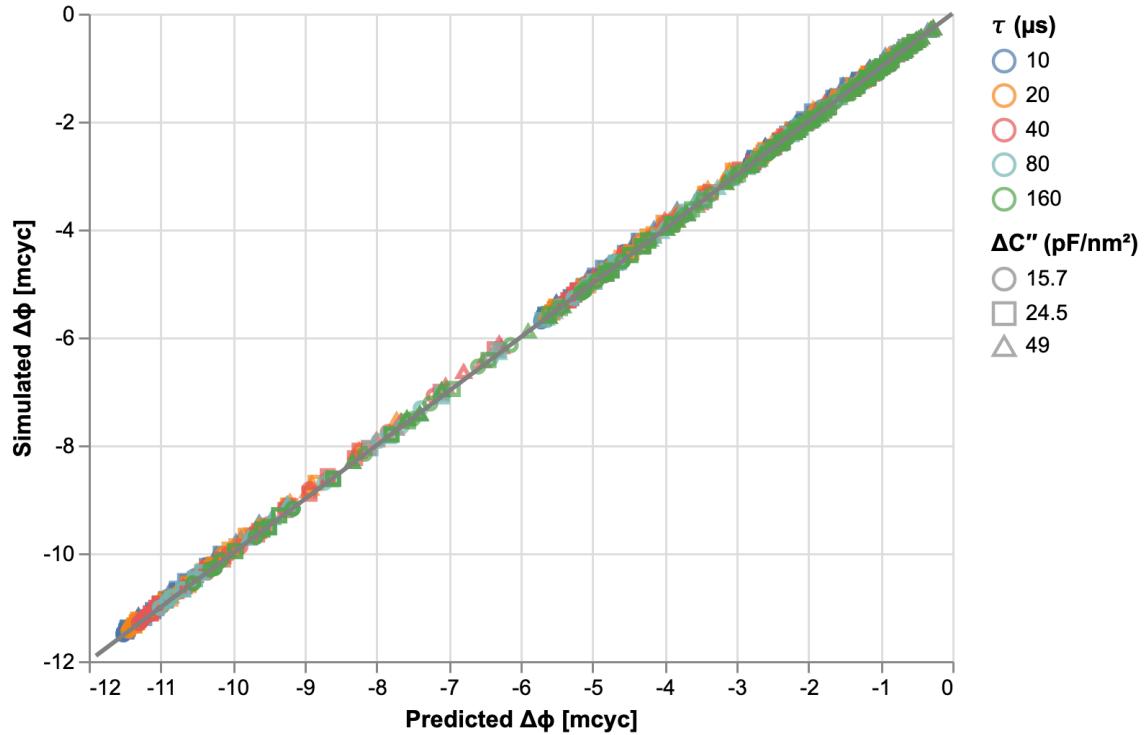


FIG. S3. Simulated phase shift agrees closely with predicted phase shift in pk-EFM experiments. Experimental parameters:  $R_{\text{dark}} = 1 \times 10^5, 1 \times 10^7, 1 \times 10^8 \text{ M}\Omega$ ;  $R_{\text{light}} = 10 \text{ M}\Omega$ ;  $\tau_L = 10, 20, 40, 80, 160 \mu\text{s}$ ;  $V_t(t_0) = 7, 10 \text{ V}$ ;  $t_p = 48, 96, 192, 384 \mu\text{s}$ ;  $C_s = 0.01, 0.04, 0.1 \text{ fF}$ ;  $C'_{\text{tip}} = -28, -35, -49.5 \text{ fF nm}^{-1}$ ;  $C_{\text{tip}} = 0.1 \text{ fF}$ ;  $f_c = 62.5 \text{ kHz}$ ;  $k_c = 3.5 \text{ N m}^{-1}$ ;  $Q = 26000$ ,  $A_0 = 50 \text{ nm}$ ;  $A_0 = 50 \text{ nm}$ ;  $t_0 = -200 \mu\text{s}$ ;  $\phi_0 = 0$ .

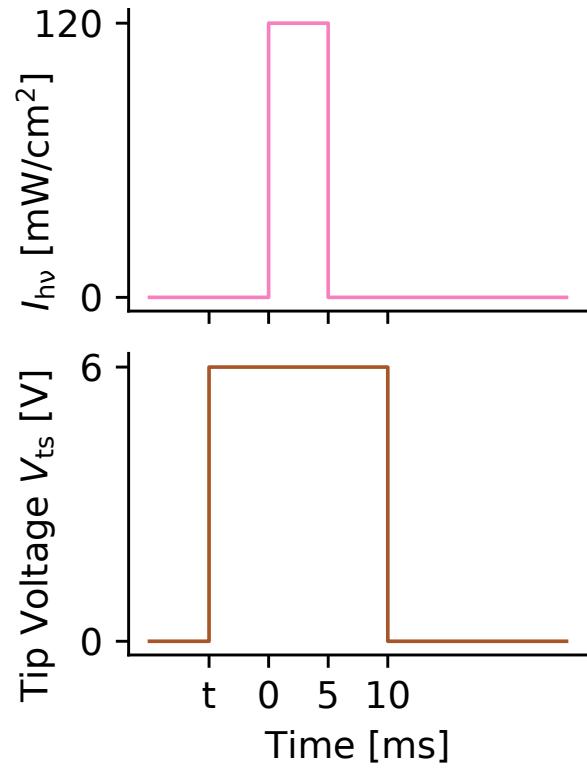


FIG. S4. Timing diagram for voltage and light pulse used in Fig. 4 experiments. Bias time in dark  $t$  is varied from -2 ms to -1000 ms.

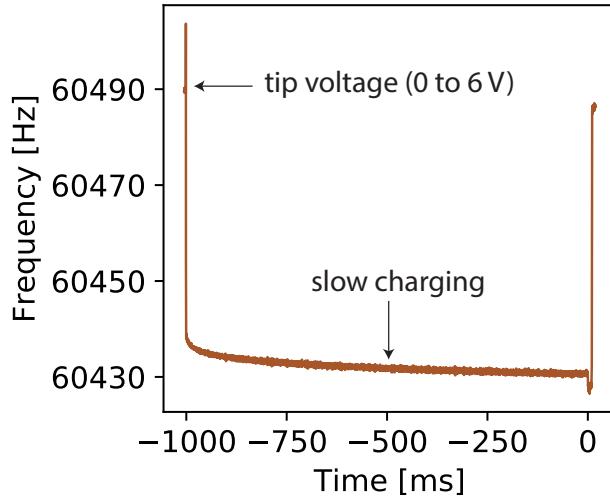


FIG. S5. tr-EFM frequency shift for  $t_p = 1$  s shown in Fig. 4. Additional frequency shift due to slow charging is apparent in the measured frequency shift after the tip voltage  $V_{ts}$  is changed to 6 V even after several hundred millisecond wait.

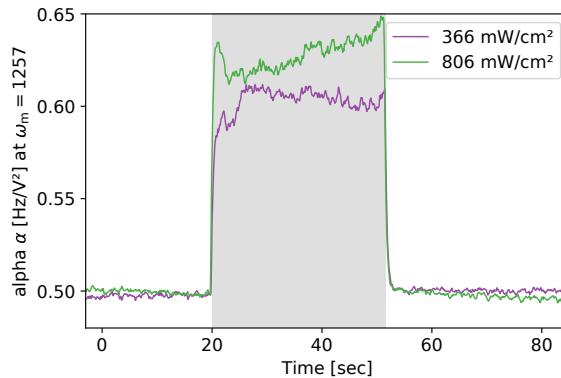


FIG. S6. Dielectric response at  $\omega_m = 1257 \text{ s}^{-1}$  for a period of illumination at two different light intensities. Dielectric response measured by  $\alpha$  increases and decreases within the time resolution of the measurement ( $\approx 1 \text{ s}$ ). Light was turned on at time  $t = 20 \text{ s}$  and turned off at time  $t = 52 \text{ s}$ . Experimental parameters:  $V_{ts} = 6 \text{ V}$ ,  $h = 200 \text{ nm}$ .

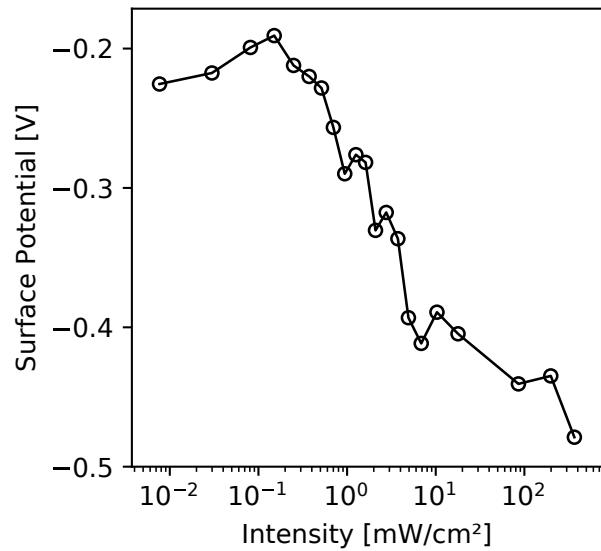


FIG. S7. Surface potential measured through frequency voltage parabolas at selected light intensities.

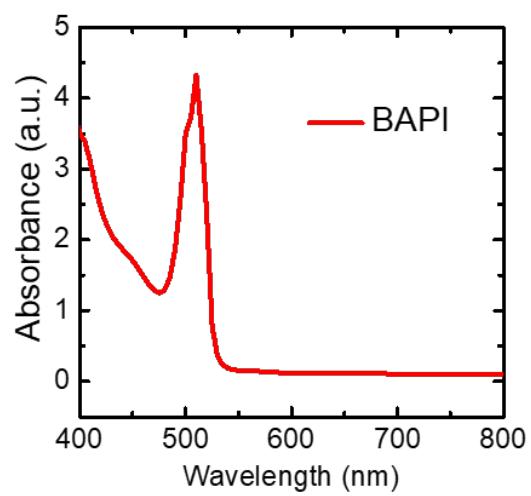


FIG. S8. Absorbance spectra of  $\text{BA}_2\text{PbI}_4$ .

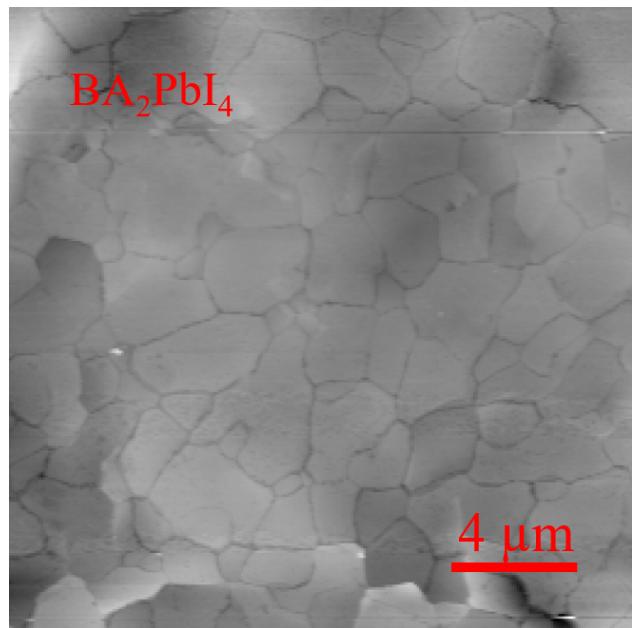


FIG. S9. AFM topography image of the perovskite film.

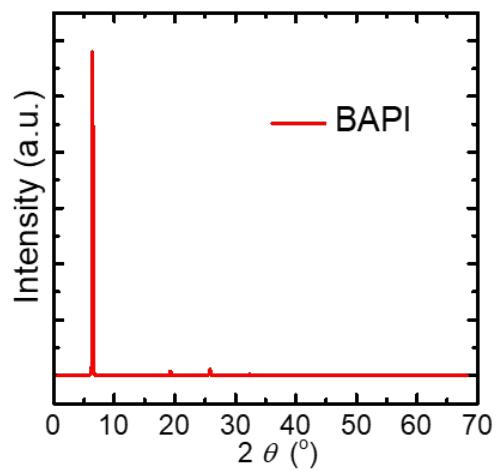


FIG. S10. XRD spectra of the perovskite film.