



# Transformation of trigonal planar B<sub>4</sub> into zigzag B<sub>4</sub> units within the new boride series Ti<sub>2-x</sub>M<sub>1+x-δ</sub>Ir<sub>3+δ</sub>B<sub>3</sub> (x = 0.5 for M = V–Mn, x = 0 for M = Mn–Ni and δ < 0.2)

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## ABSTRACT

In metal-rich borides, numerous boron fragments B<sub>n</sub> (n = 2, 3, 4, 5, 6) have been discovered, the B<sub>4</sub> units being the most versatile with four different shapes (bent, zigzag, trigonal planar and tetrahedral). We report on the new boride series Ti<sub>2-x</sub>M<sub>1+x-δ</sub>Ir<sub>3+δ</sub>B<sub>3</sub> (x = 0.5 for M = V–Mn, x = 0 for M = Mn–Ni and δ < 0.2), in which a structural change occurs by successive substitution of the 3d transition metal M = V, Cr, Mn, Fe, Co and Ni. It is found that the change in structure from the Ti<sub>1+x</sub>Os<sub>2-x</sub>RuB<sub>2</sub>-type structure (*P6̄2m*, no. 189) to the Ti<sub>1+x</sub>Rh<sub>2-x+y</sub>Ir<sub>3-y</sub>B<sub>3</sub>-type (*Pbam*, no. 55) leads to a change of B<sub>4</sub> shape from trigonal planar B<sub>4</sub> (M = V–Mn) to zigzag B<sub>4</sub> fragment (M = Mn–Ni). Even though there is no group–subgroup relationship between the two structures, we present how the Ti<sub>1+x</sub>Os<sub>2-x</sub>RuB<sub>2</sub>-type structure can easily be geometrically derived from the Ti<sub>1+x</sub>Rh<sub>2-x+y</sub>Ir<sub>3-y</sub>B<sub>3</sub>-type.

## 1. Introduction

Metal borides are known for some set of characteristic properties such as high melting points, superior hardness, and chemical inertness that originate from their extraordinary crystal chemistry, in which strong boron–boron and metal–boron bonds play a key role [1–10]. The crystal chemistry of metal-rich borides has been enriched in recent years by the discovery of complex phases in which different boron fragments such as B<sub>2</sub>, B<sub>3</sub>, B<sub>4</sub>, B<sub>5</sub> and B<sub>6</sub> were found [11–15]. In some cases, these compounds exhibit highly exciting properties: Superconductivity was observed in NbRuB (containing B<sub>2</sub> dumbbells) [16], in TaRuB and NbOsB (both containing zigzag B<sub>4</sub> units) [17], while itinerant ferrimagnetism was reported for TiCrIr<sub>2</sub>B<sub>2</sub> (trigonal planar B<sub>4</sub> units) [18], itinerant ferromagnetism in arc-melted Nb<sub>6</sub>Fe<sub>1-x</sub>Ir<sub>6+x</sub>B<sub>8</sub> (B<sub>6</sub> rings) [19] which becomes an itinerant ferrimagnet when annealed at 1600 °C [20]. Most of these boron fragments can be derived from the honeycomb boron layer of the AlB<sub>2</sub>-type structure [3]. Accordingly, the fragments arise from boron-centered trigonal metal prisms sharing a rectangular face, thus leading to short boron–boron distances (typically

1.7 to 1.8 Å). For fragments with more than three boron atoms, different configurations are possible such as for the B<sub>4</sub> (zigzag–B<sub>4</sub> in Mo<sub>2</sub>IrB<sub>2</sub> [21] and Ti<sub>1+x</sub>Rh<sub>2-x+y</sub>Ir<sub>3-y</sub>B<sub>3</sub> types [22], bent–B<sub>4</sub> in β–Cr<sub>2</sub>IrB<sub>2</sub> [13] or trigonal planar B<sub>4</sub> unit in the hexagonal Ti<sub>1+x</sub>Os<sub>2-x</sub>RuB<sub>2</sub> (*P6̄2m*, no. 189) [23]. Recently, weak spin frustration and ferrimagnetic ordering below 275 K was found in TiCrIr<sub>2</sub>B<sub>2</sub> (Ti<sub>1+x</sub>Os<sub>2-x</sub>RuB<sub>2</sub>-type) and attributed to the unique substructure of the magnetic element, namely a chain of equilateral Cr<sub>3</sub> triangles [18] which strongly interacts with trigonal planar B<sub>4</sub> units. The discovery of this phase was the result of our increased efforts of designing new phases by incorporation of magnetic 3d elements into new boride structures. Naturally, after discovering TiCrIr<sub>2</sub>B<sub>2</sub>, the next target was to study the substitution of Cr by other 3d transition metals and its impact on the crystal structure and possibly magnetic properties. In the present work we report on the new boride series Ti<sub>2-x</sub>M<sub>1+x-δ</sub>Ir<sub>3+δ</sub>B<sub>3</sub> (x = 0.5 for M = V–Mn, x = 0 for Mn–Ni and δ < 0.2) in which a structural change occurs by successive substitution of the 3d transition metal M = V, Cr, Mn, Fe, Co and Ni. Finally, we explain how the Ti<sub>1+x</sub>Os<sub>2-x</sub>RuB<sub>2</sub>-type (with trigonal planar B<sub>4</sub> fragment) and the Ti<sub>1+x</sub>Rh<sub>2-x+y</sub>Ir<sub>3-y</sub>B<sub>3</sub>-type (with zigzag B<sub>4</sub> fragment) structures can be

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geometrically derived from each other.

## 2. Experimental procedures

The boride phases were synthesized by a high temperature reaction in an electric arc furnace directly from the elements (more details on the synthesis can be found in the SI). Overall, three sets of samples were prepared with the starting compositions  $2\text{Ti}:\text{M}:\text{3Ir}:\text{3B}$ ,  $1.5\text{Ti}:\text{1.5M}:\text{3Ir}:\text{3B}$  and  $\text{Ti}:\text{2M}:\text{3Ir}:\text{3B}$ , where  $\text{M} = \text{V}, \text{Cr}, \text{Mn}, \text{Fe}, \text{Co}$  or  $\text{Ni}$ . These three sets were prepared in order to investigate if the  $\text{Ti}:\text{M}$  ratio affects the type of boride obtained. Powders of the respective elements are mixed thoroughly in the corresponding ratio and pressed into dense pellets. The pellets are then arc-melted several times in the arc furnace (20 V voltage and 20 A applied current) which is filled with Ar. Small gray beads with metallic luster were obtained. Their weight was compared to the initial weight to keep track of any material lost during the reaction, but the loss turned out to be negligible. All samples were stable in air even after a long exposure time.

The beads were crushed, and single crystals were isolated from the powders under an optical microscope. The samples were then finely ground in an agate mortar and used to collect powder X-ray diffraction (PXRD) data, utilizing  $\text{Cu K}\alpha$  radiation ( $\lambda = 1.54059 \text{ \AA}$ ) of a Stadi P powder diffractometer (STOE & Cie GmbH, Darmstadt, Germany) equipped with a Ge monochromator and an Image-plate detector. Each measurement took 3 h and covered a  $2\theta$  range from  $10^\circ$  to  $90^\circ$ . The powder diffraction data was refined by a full-matrix least-squares refinement (Rietveld) as implemented in the FULLPROF Suite [24,25].

In case suitable single crystals were isolated, they were subsequently used to collect single-crystal X-ray diffraction (SCXRD) data, utilizing  $\text{Mo-K}\alpha_1$  radiation ( $\lambda = 0.71073 \text{ \AA}$ ) of a Smart APEX diffractometer by Bruker equipped with a CCD detector. The structure was solved by direct methods using the SHELXS package and a full-matrix least-squares refinement was performed using SHELXL [26,27]. Additionally, the presence of the metals in the single crystals were confirmed by energy dispersive X-ray spectroscopy (EDX) with a LEO 1450VP electron microscope from Carl ZEISS.

## 3. Results and discussion

### 3.1. Phase analysis

In the powder diffraction data of all V- and Cr-containing samples the  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type phase was identified. However, these samples contained additional compounds, mainly the  $\text{Cu}_3\text{Au}$ -type  $\text{TiIr}_3\text{B}_{1-y}$  phase [28] in all of them and the  $\tau$ -boride  $\text{Cr}_{7.9}\text{Ir}_{4.1}\text{B}_6$  [13] for the two Cr samples with higher Cr-content ( $\text{Ti}:\text{M} = 1$  and  $1/2$ ). For the Mn-containing samples with low Mn-content ( $\text{Ti}:\text{M} = 2$  and  $1$ ), a structure change was observed as the  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type phase [22] was obtained as main product, while for higher Mn-content ( $\text{Ti}:\text{M} = 1/2$ ) only a minority  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type phase [23] is present, the main phase being the  $\tau$ -boride ( $(\text{Fe}_{0.54}\text{Ir}_{0.46})_{20}\text{Fe}_3\text{B}_6$ -type) [29]. The  $(\text{Ti}_{1-x}\text{Fe}_x)_3\text{Ir}_3\text{B}_3$  samples contain a  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type boride phase together with a  $\tau$ -boride ( $(\text{Fe}_{0.54}\text{Ir}_{0.46})_{20}\text{Fe}_3\text{B}_6$ -type) [29] for all  $\text{Ti}:\text{M}$  ratios, and a third phase ( $\text{TiIr}_3\text{B}_{1-y}$ ) [28] is found only for  $\text{Ti}:\text{M} = 2$ . The  $(\text{Ti}_{1-x}\text{M}_x)_3\text{Ir}_3\text{B}_3$  samples with  $\text{M} = \text{Co}$  and  $\text{Ni}$  contain a  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type boride only for  $\text{Ti}:\text{M} = 2$ . In these samples, some unindexed peaks are present, thus the weight fractions given in Tables S1, S2 and S3 (supporting information, SI) are just relative values.

### 3.2. Single-crystal structure determination

In order to refine the proper compositions of the new phases, single-crystal refinements coupled with EDX analyses were carried out. Due to the irregular shapes of the single crystals studied by EDX only qualitative measurements were possible, except for the orthorhombic Mn-based phase which yielded the metal ratio  $2.4(2):0.7(2):2.9(2)$  for  $\text{Ti}:\text{Mn}:\text{Ir}$ ,

**Table 1**

Crystallographic and structure refinement data of  $\text{Ti}_{1-x}\text{V}_x\text{Ir}_2\text{B}_2$  and  $\text{TiMn}_{0.92}\text{(2)Ir}_{2.08(2)}\text{B}_2$  (obtained from the starting composition  $\text{Ti}:\text{2Mn}:\text{3Ir}:\text{3B}$ ).

Formula	Ti <sub>1-x</sub> V <sub>x</sub> Ir <sub>2</sub> B <sub>2</sub>	TiMn <sub>0.92(2)</sub> Ir <sub>2.08(2)</sub> B <sub>2</sub>		
space group, Z	$\overline{P}6_2m$ (no. 189); 3	$\overline{P}6_2m$ (no. 189); 3		
formula weight (g/mol)	504.86	519.38		
F(000)	627	645		
lattice parameters				
a (Å)	8.601(2)	8.610(2)		
b (Å)	8.601(2)	8.610(2)		
c (Å)	3.1870(6)	3.1880(6)		
V (Å <sup>3</sup> )	204.18(6)	204.65(6)		
calc. density (g/cm <sup>3</sup> )	12.318	12.634		
abs. coefficient (mm <sup>-1</sup> )	103.253	107.599		
θ-range (°)	4.74 < θ < 35.69	4.74 < θ < 35.78		
hkl-range	-4 ≤ h ≤ 14 -14 ≤ k ≤ 13 -4 ≤ l ≤ 5	-13 ≤ h ≤ 14 -14 ≤ k ≤ 11 -4 ≤ l ≤ 5		
no. of reflections, R <sub>int</sub>	1809; 0.0881	2166; 0.1373		
indep. reflections	382	385		
no. parameters	17	20		
R <sub>1</sub> ; wR <sub>2</sub> (all I)	0.0639; 0.1037	0.0585; 0.1480		
Goof	1.110	1.114		
diff. peak/hole (e·Å <sup>-3</sup> )	-3.583/4.406	-6.070/7.444		
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indep. reflections	1082	1019	1063	618
no. parameters	46	37	46	41
R <sub>1</sub> ; wR <sub>2</sub> (I > 2sigma (I))	0.0588; 0.0912	0.0518; 0.0952	0.0610; 0.1133	0.0462; 0.1004
R <sub>1</sub> ; wR <sub>2</sub> (all I)	0.1111; 0.1045	0.0772; 0.1064	0.1103; 0.1294	0.0684; 0.1099
Goof	1.067	1.090	1.064	1.042
diff. peak/hole (e·Å <sup>-3</sup> )	-5.427/ 7.830	-4.965/ 5.601	-5.896/ 5.794	-4.455/ 3.404

which is in fair agreement with the results obtained from single-crystal diffraction below. The crystallographic and structure refinement data are given in Tables 1 and 2, while Table 3, S4, S5, S6 and S7 contain the refined atomic positions and displacement parameters and Table 4 contain selected bond lengths (for a representative of each structure

**Table 2**

Crystallographic and single-crystal structure data of  $\text{Ti}_2\text{M}_{1-x}\text{Ir}_{3+x}\text{B}_3$  obtained from the starting composition  $2\text{Ti}:\text{M}:\text{3Ir}:\text{3B}$  with  $\text{M} = \text{Mn}, \text{Fe}$  and  $\text{Co}$ .

Formula	$\text{Ti}_2\text{Mn}_{0.83}\text{(1)Ir}_{3.17(1)}\text{B}_3$	$\text{Ti}_2\text{Fe}_{0.91}\text{(1)Ir}_{3.09(1)}\text{B}_3$	$\text{Ti}_2\text{Co}_{0.88}\text{(1)Ir}_{3.12(1)}\text{B}_3$	$\text{Ti}_2\text{Ni}_{0.88}\text{(1)Ir}_{3.12(1)}\text{B}_3$
space group, Z	$P6_3/m$ (no. 55); 4			
formula weight (g/mol)	783.10	772.27	779.09	778.89
$F(000)$	1295	1281	1291	1295
lattice parameters				
a ( $\text{\AA}$ )	8.639(2)	8.618(1)	8.60(2)	8.598(2)
b ( $\text{\AA}$ )	15.004(4)	14.975(7)	15.00(3)	14.970(4)
c ( $\text{\AA}$ )	3.228(7)	3.221(1)	3.225(5)	3.238(7)
V ( $\text{\AA}^3$ )	418.4(9)	415.7(2)	416.0(2)	416.8(2)
calc. density ( $\text{g/cm}^3$ )	12.432	12.340	12.439	12.413
abs. coefficient ( $\text{mm}^{-1}$ )	106.281	104.932	106.143	106.429
$\theta$ -range ( $^\circ$ )	$4.71 < \theta < 35.93$	$5.46 < \theta < 35.38$	$4.72 < \theta < 35.76$	$4.72 < \theta < 28.52$
$hkl$ -range	$-12 \leq h \leq 14$ $-16 \leq k \leq 24$ $-4 \leq l \leq 5$	$-13 \leq h \leq 14$ $-8 \leq k \leq 24$ $-4 \leq l \leq 5$	$-11 \leq h \leq 14$ $-22 \leq k \leq 24$ $-4 \leq l \leq 5$	$-11 \leq h \leq 7$ $-19 \leq k \leq 20$ $-4 \leq l \leq 4$
no. of reflections, $R_{\text{int}}$	4318; 0.1311	3320; 0.0681	4238; 0.1119	2331; 0.0934
indep. reflections	1082	1019	1063	618
no. parameters	46	37	46	41
$R_1$ ; $wR_2$ ( $I > 2\sigma$ )	0.0588; 0.0912	0.0518;	0.0610;	0.0462;
( $I$ )		0.0952	0.1133	0.1004
$R_1$ ; $wR_2$ (all $I$ )	0.1111; 0.1045	0.0772;	0.1103;	0.0684;
		0.1064	0.1294	0.1099
Goof	1.067	1.090	1.064	1.042
diff. peak/hole ( $\text{e.\AA}^{-3}$ )	$-5.427/7.830$	$-4.965/5.601$	$-5.896/5.794$	$-4.455/3.404$

**Table 3**Atomic coordinates, site occupation factors (SOF) and displacement parameters of  $\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$  and  $\text{Ti}_2\text{Ni}_{0.88(1)}\text{Ir}_{3.12(1)}\text{B}_3$ .  $U_{13}$  and  $U_{23} = 0$ .

$\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$										
Atom	Site	x	Y	z	SOF	$U_{\text{eq}}$	$U_{11}$	$U_{22}$	$U_{33}$	$U_{12}$
Ir1	6k	0.4636(2)	0.2760(2)	0.5	1	0.0096(3)	0.0058(4)	0.0058(5)	0.0190(5)	0.0042(4)
Mn2/Ir2	3f	0.1786(7)	0.1786(7)	0	0.93(2)/0.07(2)	0.0010(2)	0.014 (2)	0.005(2)	0.009(2)	—
Ti3	3f	0.4174(9)	0	0	1	0.005(1)	0.006(2)	0.003(2)	0.011(2)	0.003(2)
B1	1b	0	0	0.5	1	0.028(9)	—	—	—	—
B2	3g	0	0.258(9)	0.5	1	0.028(9)	—	—	—	—
B3	2c	0.6667	0.3333	0	1	0.028(9)	—	—	—	—

$\text{Ti}_2\text{Ni}_{0.88(1)}\text{Ir}_{3.12(1)}\text{B}_3$										
atom	Site	x	y	z	SOF	$U_{\text{eq}}$	$U_{11}$	$U_{22}$	$U_{33}$	$U_{12}$
Ir1	4 h	0.0197(2)	0.09511(9)	0.5	1	0.0117(3)	0.0024(5)	0.0092(7)	0.0236(8)	0.0007(4)
Ir2	4 h	0.7148(2)	0.14228(9)	0.5	1	0.0118(4)	0.0024(6)	0.0095(7)	0.0236(7)	−0.0009(4)
Ir3	4 h	0.9337(2)	0.26982(9)	0.5	1	0.0129(4)	0.0029(5)	0.0097(7)	0.0262(8)	−0.0006(4)
Ir4/Ni4	4g	0.4748(4)	0.0858(2)	0	0.11(2)/0.89(2)	0.010(2)	0.006(2)	0.015 (2)	0.009 (2)	0.003(2)
Ti5	4g	0.1819(6)	0.2069(4)	0	1	0.0050(8)	0.005(2)	0.007(3)	0.002(3)	−0.001(2)
Ti6	4g	0.1985(6)	0.9962(4)	0	1	0.0050(8)	0.005(2)	0.007(3)	0.002(3)	−0.001(2)
B7	4 h	0.290(4)	0.099(3)	0.5	1	0.010(4)	—	—	—	—
B8	4 h	0.605(4)	0.006(3)	0.5	1	0.010(4)	—	—	—	—
B9	4g	0.890(4)	0.165(3)	0	1	0.010(4)	—	—	—	—

**Table 4**Interatomic distances (Å) of  $\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$  and  $\text{Ti}_2\text{Ni}_{0.88(1)}\text{Ir}_{3.12(1)}\text{B}_3$ .

$\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$		
Atom 1	Atom 2	d(atom1 – atom2) (Å)
B1	B2	2.22(6)
	Mn2	2.21(1)
B2	Mn2	2.10(4)
	Ti3	2.53(4)
	Ir1	2.14(2)
B3	Ir1	2.23(1)
	Ti3	2.59(1)
Ir1	Ir1	2.70(1) – 2.80(1)
Mn2	Mn2	2.61(1) – 3.19(1)

$\text{Ti}_2\text{Ni}_{0.88(1)}\text{Ir}_{3.12(1)}\text{B}_3$		
B7	B8	1.81(6)
B8	B8	1.81(4)
B7	Ir1–3	2.28(2) – 2.47(3)
	Ti5, Ti6 & Ir4/Ni4	2.32(3) – 2.33(3)
B8	Ir1–3	2.25(3)
	Ti5, Ti6 & Ir4/Ni4	2.23(2) – 2.34(2)
B9	Ir1–3	2.23(2) – 2.29(2)
	Ti5, Ti6 & Ir4/Ni4	2.53(4) – 2.62(4)
Ir1–3	Ir1–3	2.68(1) – 2.87(1)
Ir4/Ni4	Ir4/Ni4	2.61(1) – 3.22(1)

type).

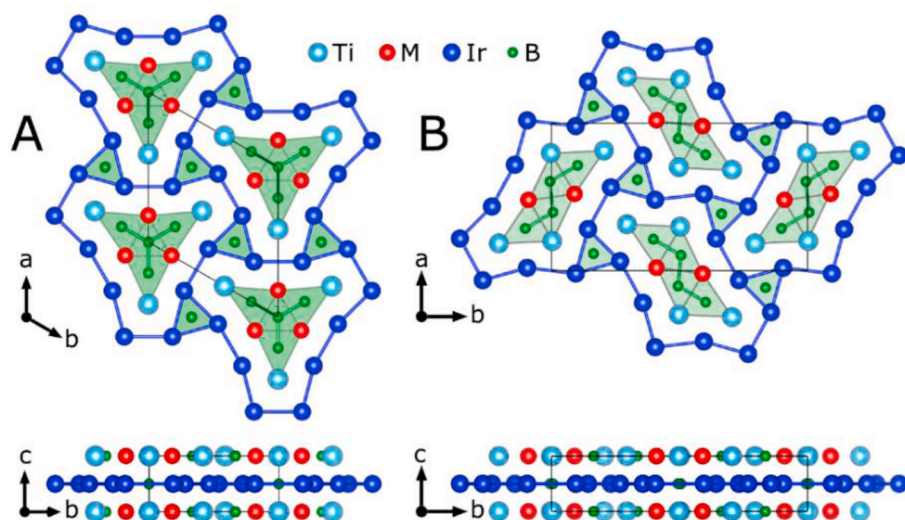
$\text{TiMnIr}_2\text{B}_2$  ( $M = \text{V}, \text{Mn}$ ): The structure of  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$  [23] was used as model for the refinement of single-crystal data of the V- and Mn-based compounds since the PXRD data indicated isotypism of both compounds with  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ . From the  $\text{TiOs}_2\text{RuB}_2$  model, Os was substituted by Ir and Ru by M (V, Mn) leading to  $\text{TiMnIr}_2\text{B}_2$  as the starting model. The refinements converged in both cases smoothly. Since Ti and V cannot be distinguished using XRD (both are 1 electron apart and have similar atomic radius) and because EDX analysis had confirmed the presence of the two elements in the single crystal analyzed, the Ti- and V-sites were refined with a fixed 50:50 (Ti:V) mixed occupancy, thus the composition of this phase should be viewed as  $(\text{Ti}_{1-x}\text{V}_x)_2\text{Ir}_2\text{B}_2$ . In the case of  $M = \text{Mn}$ , the large difference in atomic radius between Ti and Mn was enough to differentiate both elements (similar to the reported  $M = \text{Cr}$ ) [18]. However, disorder on the Mn site with the heavier Ir (8%) was necessary to fully occupy the site, leading to the refined chemical formula  $\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$ . Occupational disorder between a first-row transition metal and a 4d or 5d metal is frequently observed in metal borides such as  $\text{Nb}_6\text{Fe}_{1-x}\text{Ir}_{0.46}\text{B}_8$  or  $(\text{Fe}_{0.54}\text{Ir}_{0.46})_{20}\text{Fe}_3\text{B}_6$  [19,29]. Furthermore, the B–B distance had to be restrained to 1.85 Å (average

B–B distance in  $\text{Ti}_2\text{MnIr}_3\text{B}_3$ , see below) as the refined distance was very large (2.22 Å), exceeding all other distances (1.75–1.90 Å) found for all  $\text{B}_4$  units in this and related structure types.

$\text{Ti}_2\text{MnIr}_3\text{B}_3$  ( $M = \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}$ ): PXRD data established isotypism of these four compounds with  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type structure. Suitable single crystals could be isolated in all cases, and the composition “ $\text{Ti}_{1+x}(\text{Rh}_{1-x})(\text{Rh}_{1+y})\text{Ir}_{3-y}\text{B}_3$ ” was used as the starting model but replacing  $(\text{Rh}_{1-x})$  and  $(\text{Rh}_{1+y})$  by Ti and M, respectively, leading to the ideal starting formula  $\text{Ti}_2\text{MnIr}_3\text{B}_3$ . In all cases, the refinements converged smoothly. However, a mixed occupancy on the M site with the heavier Ir was necessary to fully occupy the M site, leading to the general chemical formula  $\text{Ti}_2\text{M}_{1-y}\text{Ir}_{3+y}\text{B}_2$  ( $y < 0.2$ ). The B–B bond lengths range from 1.81 (5) Å to 1.90(5) Å for the outer B7–B8 bond and from 1.76(6) Å to 1.84 (6) Å for the central B8–B8 bond of the zigzag  $\text{B}_4$  fragment. A list of the bond lengths for all compounds can be found in Tables S8 and S9 in the SI. In general, they fall into the expected range for intermetallic borides. The M–B and Ir–B bonds are roughly between 2.2 Å and 2.3 Å, while the Ti–B distances are mostly 2.4 Å or longer.

### 3.3. Structural descriptions

$(\text{Ti}_{1-x}\text{V}_x)_2\text{Ir}_2\text{B}_2$  and  $\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$ : All atoms are distributed in two layers at  $z = 0$  and  $z = 0.5$  along the  $c$  direction. All Ir atoms are found in the layer at  $z = 0.5$  together with two out of three boron atoms (Fig. 1). The two boron atoms (B1 and B2) build trigonal planar  $\text{B}_4$  units with a B–B distance of 1.78(5) Å in  $(\text{Ti}_{1-x}\text{V}_x)_2\text{Ir}_2\text{B}_2$  and 1.85 Å in  $\text{TiMn}_{0.92(2)}\text{Ir}_{2.08(2)}\text{B}_2$ . The layer at  $z = 0$  contains all other metals (Ti, Mn, Ti/V) and the third boron atom (B3). In this layer the M (Ti/V or Mn) atoms form equilateral triangles with a M–M distance of 2.72(2) Å ( $M = \text{Ti/V}$ ) or 2.61(1) Å ( $M = \text{Mn}$ ). Two  $\text{M}_3$  triangles are stacked on top of each other, forming  $\text{M}_3$  chains with an intrachain distance corresponding to the  $c$  lattice parameter at ca. 3.19 Å. The shortest distance between two Mn-atoms of two adjacent triangles in the same layer is 6.404(7) Å, which is more than twice as large as the distance in the  $c$  direction. These distances and the boron–metal and metal–metal distances are comparable to those found in  $\text{TiCrIr}_2\text{B}_2$  and reported  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type phases such as  $\text{Ti}_{1.6}\text{Os}_{2.4}\text{B}_2$  and  $\text{Ti}_{1-x}\text{Fe}_x\text{Os}_2\text{RhB}_2$  ( $0 < x < 0.5$ ) [30]. All atoms at  $z = 0$  are coordinated by prisms build by atoms in the  $z = 0.5$  layers: trigonal prisms for boron and pentagonal prisms for Ti and M (Fig. 1a). Likewise, the boron atoms at  $z = 0.5$  are coordinated by prisms of metals at  $z = 0$ , whereas the iridium atoms are coordinated by strongly distorted icosahedra of all atom. For clarity reasons, we will use the simplified formula  $\text{TiMnIr}_2\text{B}_2$  in the following discussion to represent these two phases.



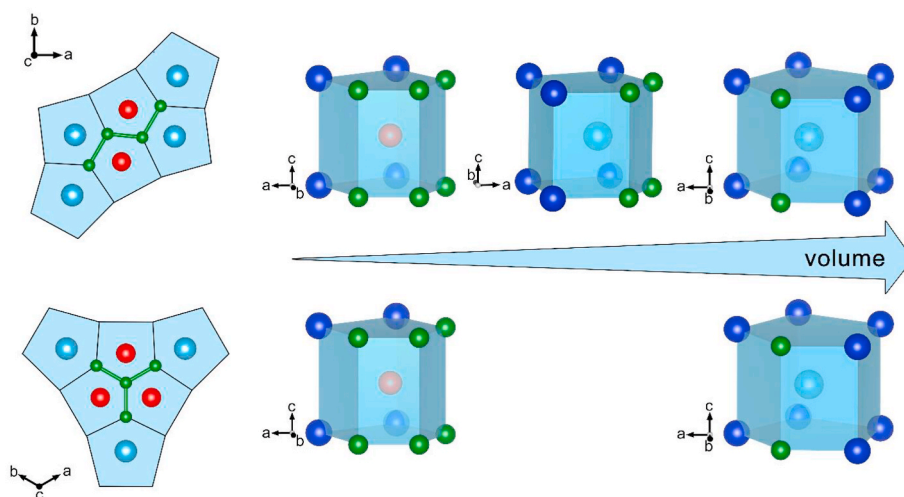
**Fig. 1.** Projections of the structures of (a)  $\text{TiMnIr}_2\text{B}_2$  ( $M = \text{V, Mn}$ ;  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type) and of (b)  $\text{Ti}_2\text{M}_{1-x}\text{Ir}_{3+x}\text{B}_3$  ( $x < 0.2$ ;  $M = \text{Mn-Ni}$ ;  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type) phases along [001] on top and along [100] at the bottom. The highlighted green triangles indicate boron-filled trigonal prisms. Note that for the  $M = \text{V}$  phase, the Ti and M positions are all filled by the mixed Ti/V atoms. (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

The  $\text{TiMnIr}_2\text{B}_2$ -phase contains a magnetically active element (Mn), thus the Mn substructure (chains of  $\text{Mn}_3$  triangles) has the potential to produce interesting magnetic properties. Unfortunately, only 6 wt.-% of this phase has been synthesized until now, so the magnetic properties cannot be studied at this time. Nevertheless, comparing this substructure with other will lead to a better understanding of the potential magnetic properties. For example, the similar Cr substructure has produced the isomorphous  $\text{TiCrIr}_2\text{B}_2$  ferrimagnet. Another boride phase containing a similar substructure of a magnetic element is  $\text{NbFeB}$  ( $\text{ZrNiAl}$ -type structure) [31], also crystallizes in the hexagonal space group ( $P6_2/m$ , no 189, Pearson symbol hP9) and contains trigonal boron-centered  $\text{M}_6\text{B}$  prisms. Thus, both structure types contain triangles of magnetic 3d transition metals. In contrast to the boron capped  $\text{M}_6\text{B}$  ( $M = \text{Cr, Mn}$ ) prisms, the rectangular faces of the  $\text{Fe}_6\text{B}$  prisms in  $\text{NbFeB}$  are capped by Nb-atoms, resulting in isolated B atoms in the latter compound instead of trigonal  $\text{B}_4$  fragments as found in  $\text{TiMnIr}_2\text{B}_2$ . Nevertheless, the  $\text{M}_3$  ( $M = \text{Cr, Mn, Fe}$ ) triangles in these three compounds are very similar, having comparable intra-triangle distances of 2.64 Å (in  $\text{Cr}_3$ ), 2.61 Å (in  $\text{Mn}_3$ ) and 2.67 Å (in  $\text{Fe}_3$ ) as well as similar inter-triangle distances of 3.18 Å ( $\text{Fe}_3 - \text{Fe}_3$ ), 3.19 Å ( $\text{Mn}_3 - \text{Mn}_3$ ) and 3.22 Å ( $\text{Fe}_3 - \text{Fe}_3$ ). Consequently, the distances within the chains of  $\text{M}_3$  triangles in these three compounds are nearly the same, but the distance between these chains is

structure-type dependent: The interchain distance of 8.61 Å in  $\text{TiMnIr}_2\text{B}_2$  is similar to that in  $\text{TiCrIr}_2\text{B}_2$  (8.55 Å), whereas it is only 6.01 Å in  $\text{NbFeB}$  (both center-to-center), hinting at different magnetic properties for these two structure types despite the similar magnetic chains of  $\text{M}_3$  triangles. Such studies of magnetic properties would be highly interesting both experimentally and theoretically, as they would help shed some light on the effects of different interchain distances.

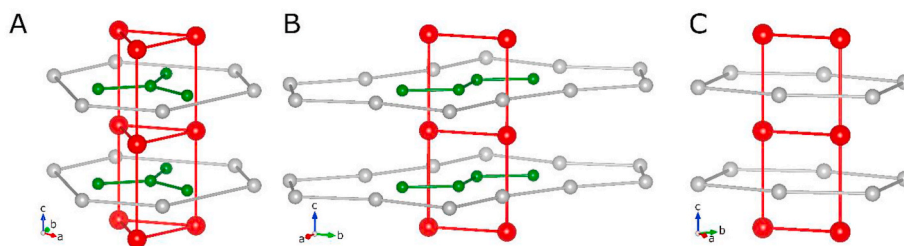
$\text{Ti}_2\text{M}_{1-x}\text{Ir}_{3+x}\text{B}_3$  ( $x < 0.2$ ;  $M = \text{Mn-Ni}$ ): Analogous to the  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type structure discussed above, the structure of  $\text{Ti}_2\text{M}_{1-x}\text{Ir}_{3+x}\text{B}_3$  compounds ( $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type) displays two different layers at  $z = 0$  and  $z = 0.5$ : Ti, M and isolated boron atoms are found at  $z = 0$ , while Ir and other B atoms (building zigzag  $\text{B}_4$  fragments) fill the  $z = 0.5$  layer.

Two  $\text{M}_2$  dumbbells and two Ti atoms form a trigonal prism ( $[\text{Ti}_2\text{M}_4]\text{B}_8$  prism) which is centered by the inner B atom (B8) of the zigzag  $\text{B}_4$  fragment. The outer B atoms (B7) of the zigzag  $\text{B}_4$  fragment resides in the nearby trigonal prisms which have the inverse Ti:M ratio ( $[\text{Ti}_4\text{M}_2]\text{B}_7$  prisms). The zigzag  $\text{B}_4$  fragment arises from the rectangular face sharing of  $[\text{Ti}_2\text{M}_4]\text{B}_8$  and  $[\text{Ti}_4\text{M}_2]\text{B}_7$  prisms (Fig. 1b). The two other rectangular faces of these prisms are all capped by Ir atoms at  $z = 0.5$ . Six of those Ir atoms form another set of trigonal prisms centered by the isolated B9 atoms ( $[\text{Ir}_6]\text{B}_9$  prism). Because the two layers at  $z = 0$  and  $z = 0.5$



**Fig. 2.** Sets of pentagonal prisms centered by M (red) and Ti (cyan) in the orthorhombic  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type structure (top) and the hexagonal  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type structure (bottom). (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)





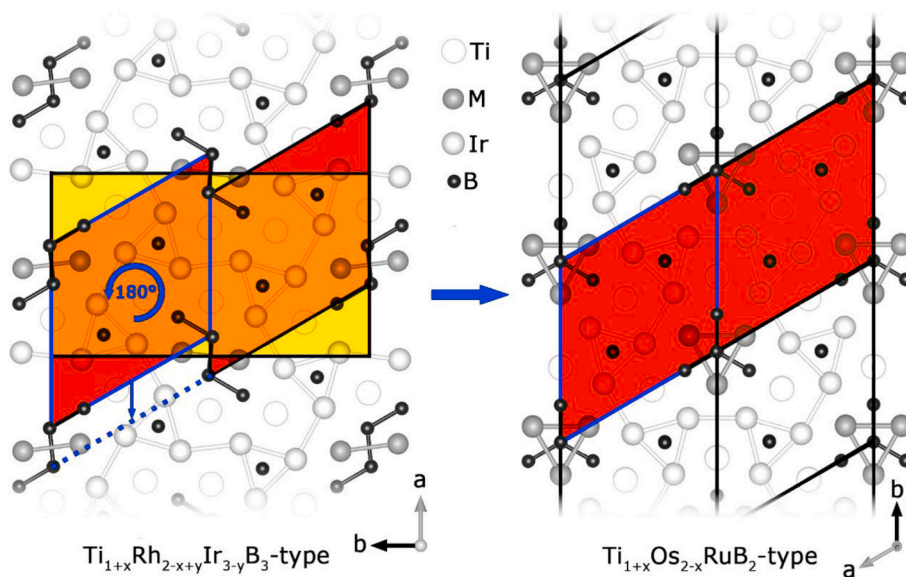
**Fig. 3.** Detailed view of the different M-substructures ( $M_3$  chain and  $M_2$  ladder) and coordination found in (a) the hexagonal  $TiIr_2B_2$  structure ( $Ti_{1+x}Os_{2-x}RuB_2$ -type), (b) in the orthorhombic  $Ti_2M1R_3B_3$  structure ( $Ti_{1+x}Rh_{2-x+y}Ir_{3-y}B_3$ -type, right) and in the tetragonal  $Ti_9M_2Ru_{18}B_8$  ( $M = Fe$ ) structure.

alternate along the  $c$ -axis, all trigonal prisms form infinite columns parallel to  $[001]$  by sharing their trigonal faces. Just as in the previous structure type, all atoms at  $z = 0$  are coordinated by prisms build by atoms in the  $z = 0.5$  layers including pentagonal prisms for Ti and M (Figs. 1 and 2). In contrast, only the boron atoms at  $z = 0.5$  are coordinated by prisms of metals at  $z = 0$ , whereas the iridium atoms are coordinated by strongly distorted icosahedra of all atom types as previously described in  $Ti_{1+x}Rh_{2-x+y}Ir_{3-y}B_3$  [22]. The B–B distance in this series varies from 1.81(4) Å for  $M = Ni$  to 1.90(4) Å for  $M = Mn$ . These phases are the first of this structure type containing a magnetic element (M), thus it is important to describe the substructure build by M: The shortest M–M distance, which is very similar in all phases and varies from 2.61(1) Å for  $M = Mn$  and  $Ni$  to 2.67(1) Å for  $M = Fe$  and  $Co$ , is consistent with  $M_2$  dumbbells. The distance between dumbbells (inter-dumbbell distance) has the same length as the  $c$  lattice parameter (ca. 3.23 Å), thereby building a ladder substructure ( $M_2$  chain). A similar ladder-like arrangement of magnetic Fe atoms (with Fe–Fe distances of 2.49 Å and 2.97 Å) has been observed in  $Ti_9Fe_2Ru_{18}B_8$ , and it is believed to have induced ferromagnetism in this phase [32]. However, the significant differences between these two types of ladder substructures (ca. 0.2 Å larger distances in the new substructure) can affect the magnetic properties. In particular, the proximity of the nearby  $B_4$  fragments which strongly interact with the new ladders (Fig. 3b) in the new phases (absent in  $Ti_9Fe_2Ru_{18}B_8$ , Fig. 3c) may drastically impact the magnetic properties, as strong bonding between M and B will lower the electron density on M and thus reduce the strength of M–M magnetic interactions (a related manuscript on theoretical calculations will be communicated soon).

### 3.4. Structural relationship between $Ti_{1-x}Rh_{2-x+y}Ir_{3-y}B_3$ - and $Ti_{1-x}Os_{2-x}RuB_2$ -type structures

Using different 3d transition metals for M ( $M = V - Ni$ ) in the series  $(Ti_{1-x}M_x)_3Ir_3B_3$  ( $x = 1/3, 1/2$ ) results in the formation of hexagonal  $Ti_{1-x}Os_{2-x}RuB_2$ -type structure for  $M = V, Cr$  and  $Mn$  on the one hand and of orthorhombic  $Ti_{1-x}Rh_{2-x+y}Ir_{3-y}B_3$ -type structure for  $M = Mn, Fe, Co$  and  $Ni$  on the other hand. Moreover, in the case of  $M = Mn$ , both structures occur side by side. Comparing the two structure types, a close relationship between them becomes apparent: Both are layered along the  $c$ -axis, both have a metal-to-boron ratio M:B of 2:1 and both contain isolated boron as well as  $B_4$  fragments. However, the  $B_4$  unit is trigonal planar in the  $Ti_{1-x}Os_{2-x}RuB_2$ -type structure while it is zigzag shaped in the  $Ti_{1-x}Rh_{2-x+y}Ir_{3-y}B_3$ -type structure. Each boron atom of both  $B_4$ -units is located inside a trigonal prism made from M and Ti. The different topology of the  $B_4$ -fragment and ultimately the different space group symmetry of the structures depends only on how those trigonal prisms are connected. While three trigonal prisms in the  $Ti_{1-x}Os_{2-x}RuB_2$ -type structure are connected via one of their rectangular faces to a central trigonal prism, the trigonal prisms in the  $Ti_{1-x}Rh_{2-x+y}Ir_{3-y}B_3$ -type structure are all successively connected by their rectangular faces (Figs. 1 and 2).

The close relationship between the two structures raises the question why one is preferred over the other. By looking at the whole series one finds that for the larger 3d transition metals (V, Cr, Mn) the hexagonal structure is preferred, while the smaller 3d metals (Mn, Fe, Co, Ni) prefer the orthorhombic structure. Analyzing the coordination polyhedra (Fig. 2) around the metal atoms M and Ti, two different polyhedra exist



**Fig. 4.** Structural relation between the orthorhombic  $Ti_{1+x}Rh_{2-x+y}Ir_{3-y}B_3$ -type structure (left, unit cell shaded in yellow) and the hexagonal  $Ti_{1+x}Os_{2-x}RuB_2$ -type structure (right, unit cells shaded in red). (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

for the hexagonal structure while three different polyhedra occur in the orthorhombic structure. Two out of three polyhedra of the orthorhombic structure are identical with the two polyhedra of the hexagonal structure. Accordingly, the orthorhombic structure contains a coordination polyhedron that does not exist in the hexagonal structure; thus, this explains the absence of a group-subgroup relationship between them.

The polyhedra have different volumes (Fig. 2), and the volume increases with decreasing number of boron atoms being part of the coordination polyhedron. The hexagonal structure contains the small polyhedron (six boron atoms participating) and the large polyhedron (two boron atoms involved). In the orthorhombic structure, however, an additional third coordination polyhedron exists and contains four boron atoms; thus, its volume lies in-between the previous two.

Naturally, we expect the large atoms to occupy the larger polyhedron, as it was found in the hexagonal structure of  $\text{TiCrIr}_2\text{B}_2$ , where Ti seats in the larger polyhedron while the smaller Cr atom occupies the smaller polyhedron. The same holds true for the hexagonal compounds with  $M = \text{V}$  and  $\text{Mn}$ .

Due to the presence of three different polyhedra, the situation is slightly more complex in the orthorhombic structures, but the general trend is the same. The Ti atoms occupy the large and the intermediate polyhedra, while the smaller M atoms ( $M = \text{Mn}, \text{Fe}, \text{Co}$  or  $\text{Ni}$ ) are found inside the small polyhedron mixed with a small amount of Ir (the second largest metal atom,  $r_{\text{Ti}} > r_{\text{Ir}} > r_{\text{M}}$ ).

The two different pentagonal prisms of the hexagonal structure occur three times each, while the three different pentagonal prisms of the orthorhombic structure occur twice each. Thus, the total number of pentagonal prisms is the same for both structures. Consequently, the ratio between the number of large sites (preferred by Ti) to the number of small sites (preferred by M) in the hexagonal structure is 1:1. In comparison, the ratio between the number of sites preferred by Ti to the number of sites preferred by M in the orthorhombic structure is 2:1, hinting at larger unit cells for the orthorhombic structure. Indeed, the average cell volume is  $406.5 \text{ \AA}^3$  for the hexagonal compounds and  $415.8 \text{ \AA}^3$  for the orthorhombic compounds.

As mentioned above, a group-subgroup relationship does not exist between the two structure types. Nevertheless, starting from the orthorhombic structure, one can derive the hexagonal structure with the following simple steps:

- 1) Identify two similar hexagonal unit cells (Fig. 4, left, red shadings) within the orthorhombic  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type unit cell (Fig. 4, left, yellow shadings),
- 2) Rotate one hexagonal unit cell (Fig. 4, left, blue outline) by  $180^\circ$  around the  $c$ -axis, which produces the other adjacent hexagonal unit cell, and
- 3) Shift the rotated left hexagonal unit with respect to the other hexagonal cell along  $a$  by ca.  $1.8 \text{ \AA}$  (blue arrow).
- 4) The hexagonal  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$ -type structure is obtained (Fig. 4, right).

No modification in  $c$  direction is required because of the same layered arrangement of both structures. This transformation can be considered as a form of glide-reflection twinning of the hexagonal structure [33].

#### 4. Conclusion

We have synthesized and fully characterized new phases within the series  $\text{Ti}_{2-x}\text{M}_{1+x-\delta}\text{Ir}_{3+\delta}\text{B}_3$  ( $x = 0.5$  for  $M = \text{V-Mn}$ ,  $x = 0$  for  $M = \text{Mn-Ni}$  and  $\delta < 0.2$ ), the main product is isostructural to hexagonal  $\text{Ti}_{1+x}\text{Os}_{2-x}\text{RuB}_2$  for  $M = \text{V-Mn}$ , while for  $M = \text{Fe}, \text{Co}$  or  $\text{Ni}$  the  $\text{Ti}_{1+x}\text{Rh}_{2-x+y}\text{Ir}_{3-y}\text{B}_3$ -type structure is adopted, independent of the Ti:M ratio of the starting composition. For  $M = \text{Mn}$ , both structures were obtained side by side. The different Ti:M ratio (1:1 in the hexagonal and 2:1 in the orthorhombic) is attributed to differently sized polyhedra around these

atoms, thereby directing the formation of each structure type. Moreover, the two structure types can be geometrically derived from each other, even though there is no obvious group-subgroup relationship.

#### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### CRediT authorship contribution statement

**Jan P. Scheifers:** Investigation, Writing - original draft, Writing - review & editing. **Michael Küpers:** Methodology, Investigation. **Yue-mei Zhang:** Supervision, Writing - review & editing. **Laura Lutz-Kappelman:** Visualization, Writing - review & editing. **Gordon J. Miller:** Supervision, Writing - review & editing. **Boniface P.T. Fokwa:** Conceptualization, Supervision, Writing - original draft, Writing - review & editing.

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#### Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.solidstatesciences.2020.106294>.

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