



The Gaia–Kepler Stellar Properties Catalog. I. Homogeneous Fundamental Properties for 186,301 Kepler Stars

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Abstract

An accurate and precise Kepler Stellar Properties Catalog is essential for the interpretation of the Kepler exoplanet survey results. Previous Kepler Stellar Properties Catalogs have focused on reporting the best-available parameters for each star, but this has required combining data from a variety of heterogeneous sources. We present the Gaia–Kepler Stellar Properties Catalog, a set of stellar properties of 186,301 Kepler stars, homogeneously derived from isochrones and broadband photometry, Gaia Data Release 2 parallaxes, and spectroscopic metallicities, where available. Our photometric effective temperatures, derived from g to K_s colors, are calibrated on stars with interferometric angular diameters. Median catalog uncertainties are 112 K for T_{eff} , 0.05 dex for $\log g$, 4% for R_* , 7% for M_* , 13% for ρ_* , 10% for L_* , and 56% for stellar age. These precise constraints on stellar properties for this sample of stars will allow unprecedented investigations into trends in stellar and exoplanet properties as a function of stellar mass and age. In addition, our homogeneous parameter determinations will permit more accurate calculations of planet occurrence and trends with stellar properties.

Unified Astronomy Thesaurus concepts: Catalogs (205); Fundamental parameters of stars (555); Exoplanet systems (484)

Supporting material: machine-readable tables

1. Introduction

The Kepler mission, officially retired in 2018, has left an unprecedented legacy data set for stellar astrophysics and exoplanet science. Due to the long-baseline, high-precision observations and subsequent follow-up efforts, the Kepler target stars have become one of the best-characterized samples of stars (Huber et al. 2014; Mathur et al. 2017; Berger et al. 2018b).

The original Kepler Input Catalog (KIC; Brown et al. 2011) compiled data for the purpose of target selection. It included optical photometry ($griz$), T_{eff} , $\log g$, and metallicities. From the ~ 13 million stars within the KIC, $\sim 200,000$ stars were chosen for monitoring based on the KIC stellar properties. The exact selection function is complex, but solar-type stars were prioritized according to more precise determinations of the Kepler sample’s stellar properties (Batalha et al. 2010; Berger et al. 2018b). Overall, these $\sim 200,000$ target stars either had imprecise stellar parameters—0.3–0.4 dex uncertainties in $\log g$ (Brown et al. 2011; Huber et al. 2016) and ≈ 200 K uncertainties in T_{eff} —or lacked parameters altogether, such as masses, ages, radii, densities, and distances.

The first Kepler Stellar Properties Catalog (KSPC; Huber et al. 2014) was published to consolidate all of the follow-up work done for Kepler stars and improve the estimated planetary occurrence rates (e.g., Howard et al. 2012; Petigura et al. 2013; Burke et al. 2015; Fulton et al. 2017). This catalog included follow-up spectroscopy, spectroscopic surveys, and asteroseismic analysis. In addition to the 12,000 Kepler stars with asteroseismic constraints prior to 2014, Huber et al. (2014) analyzed another ~ 3000 oscillating stars, providing a total of $\approx 15,500$ stars with asteroseismic radii and masses. However,

radii and masses for most stars remained imprecise due to the vast majority of stars having only photometric constraints.

In the years following the first KSPC, the number of Kepler stars with spectroscopic constraints increased considerably due to two large-scale spectroscopic surveys: (1) the Apache Point Observatory for Galactic Evolution Experiment (APOGEE; Majewski et al. 2017) and (2) the Large Sky Area Multi-Object Fiber Spectroscopic Telescope survey (LAMOST; Zhao et al. 2012; Luo et al. 2015). Mathur et al. (2017) implemented this spectroscopy in addition to $\log g$ constraints from the stellar granulation-driven flicker method (Bastien et al. 2016) to produce the Data Release 25 (DR25) KSPC. In combination with improved methodology, these additional data led to typical uncertainties of $\approx 27\%$ in radius, $\approx 17\%$ in mass, and $\approx 51\%$ in density. The large median catalog uncertainties on radius and density remained due to a lack of additional data (e.g., spectroscopy and parallaxes) for the majority of stars.

Fortunately, Gaia DR2 (Arenou et al. 2018; Gaia Collaboration et al. 2018; Lindegren et al. 2018) recently provided parallaxes to 1.3 billion stars, including $\leq 20\%$ parallaxes for $\sim 180,000$ Kepler stars (Berger et al. 2018b). Combining Gaia DR2 parallaxes with T_{eff} and K_s -band magnitudes from the DR25 KSPC (Mathur et al. 2017), Berger et al. (2018b) recomputed stellar radii and luminosities for 177,911 Kepler stars, updating our census of the Kepler targets with median radius precisions of 8% and allowing us to determine the fraction of main-sequence (67%), subgiant (21%), and giant (12%) stars in the Kepler target list. However, this work did not provide masses, $\log g$, or densities because isochrones were not used. Isochrones are required to derive physical parameters such as mass, $\log g$, and density from bulk observables such as parallaxes and photometry.

In this paper, we utilize Gaia DR2 parallaxes, homogeneous stellar g and K_s photometry, and spectroscopic metallicities,

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where available, to improve on previous analyses and present the most accurate, homogeneous, and precise analysis of stars in the Kepler field. We re-derive stellar T_{eff} , $\log g$, radii, masses, densities, luminosities, and ages for 186,301 Kepler targets, and investigate the stellar properties of a number of noteworthy Kepler exoplanet-hosting stars.

2. Methodology

2.1. Sample Selection

To identify our sample, we use the same Gaia–Kepler cross-match detailed in Berger et al. (2018b), which included 195,710 stars. We removed stars lacking “AAA” 2MASS photometry (Skrutskie et al. 2006) and any stars lacking measured parallaxes in Gaia DR2. Requiring “AAA” 2MASS photometry means we removed the brightest stars due to saturation and the faintest stars due to photon noise. In addition, after cross-matching our sample with the binaries of Kraus et al. (2016), we found that 248 out of the 263 matched primaries had “AAA” 2MASS photometry, while seven of those without “AAA” photometry had low-contrast ($\Delta m \lesssim 2$ mag), moderately resolved ($2' \leq \theta_{\text{sep}} \leq 4''$) companions. Stars lacking parallaxes are typically moderately close (~ 200 – 400 mas) equal brightness binaries, whereas other binaries at least have parallaxes, even if the errors are larger (A. Kraus et al. 2020, in preparation). These cuts reduced the sample to 190,213 stars and then to 186,672 stars, successively. Requiring g -band photometry from either the KIC or the Kepler-INT Survey (KIS; Greiss et al. 2012) reduced the catalog to 186,548 stars.

2.2. Input Photometry

To create a homogeneous catalog for the entire Kepler target sample, we mainly used Sloan g and 2MASS K_s photometry. We chose these two passbands to maximize both our T_{eff} sensitivity and the number of stars included in our final catalog. We avoided using additional SDSS or other ground-based bandpasses which would further reduce our sample. However, we did use J magnitudes to derive estimated K_s magnitudes for binary secondaries, although they are not used directly as an input to our isochrone fitting process. While Gaia G , B_p , and R_p are available for the vast majority of Kepler stars, there remains ongoing work to sufficiently characterize their transmission profiles for synthetic photometry.

We took our K_s photometry solely from 2MASS, which has an effective angular resolution of $4''$, where any fainter sources between $1''.5$ and $4''$ were omitted (Skrutskie et al. 2006). We then used Gaia DR2 photometry, which has a resolution of $\lesssim 1''$, to identify contaminating sources within $4''$. To do this, we first cross-matched Gaia DR2 sources within $4''$ of our catalog of Kepler stars. Next, we cross-matched Gaia-detected secondary sources with the United Kingdom Infrared Telescope (UKIRT) J -band observations using the WFCAM Science Archive and the WSERV4v20101019 database. We chose a matching radius of $0''.5$ from the Gaia secondaries based on the minimum in the distribution of angular separation of matches. We do not use these J magnitudes in our isochrone fitting procedure, as K_s magnitudes are less affected by extinction than UKIRT J and maximize our T_{eff} sensitivity.

Without taking any additional steps, we might wrongly confuse some UKIRT J -band magnitudes as belonging to the

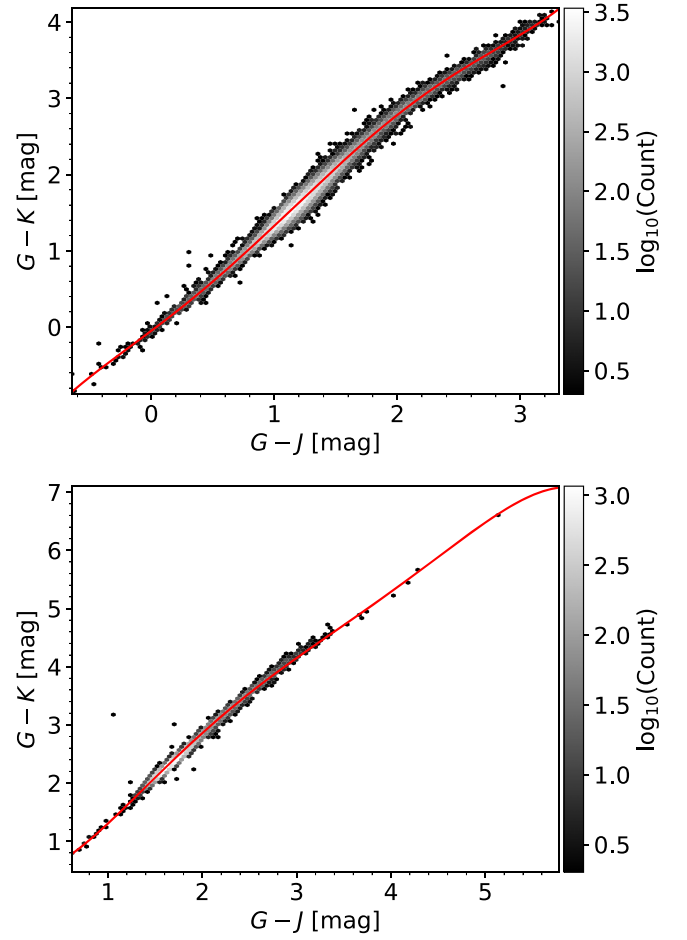


Figure 1. $G - K_s$ vs. $G - J$ of all stars without Gaia companions within $4''$ for dwarfs (top panel) and giants (bottom panel). Color-coding represents logarithmic number density. The red line displays the best-fit fifth-order polynomial to the locus of points.

secondaries when they belong to the primaries. To find these false matches, we plotted the distribution of $J_{2\text{MASS}} - J_{\text{UKIRT}}$ and found two peaks: (1) one narrow peak occurring at 0.0 ± 0.2 mag and (2) a broader peak occurring at 3 ± 1 mag. The first peak corresponds to false secondary-as-primary identifications while the second peak represents true secondaries. Therefore, we excluded all secondaries that have both $|J_{2\text{MASS}} - J_{\text{UKIRT}}| \leq 0.2$ mag and angular separations $\leq 1''.5$.

Figure 1 displays our computed fifth-order polynomial fits to the $G - K_s$ versus $G - J_{2\text{MASS}}$ curves of the non-binary Kepler dwarfs (top) and giants (bottom); stars are designated by their evolutionary state (Table 1, Berger et al. 2018b). We removed all stars with binary flags ≥ 0 in Table 1 of Berger et al. (2018b) as well as those with Gaia-detected companions to avoid contaminated secondary K_s magnitudes. We computed secondary K_s magnitudes using the difference of the Gaia G magnitude and a $G - K_s$ color. The $G - K_s$ color was computed from the best-fit polynomial evaluated at $G - J_{\text{UKIRT}}$, where J_{UKIRT} is the “jAperMag3” UKIRT J photometry. For secondaries without a UKIRT J magnitude, we did not compute a secondary K_s magnitude.

To compute amended K_s magnitudes for those primary stars with Gaia DR2 companions within $4''$, we used the following

Table 1
Gaia–Kepler Stellar Input Parameters

KIC ID	g (mag)	σ_g (mag)	K_s (mag)	σ_K (mag)	π (mas)	σ_π (mas)	[Fe/H]	$\sigma_{[\text{Fe}/\text{H}]}$	RUWE	# Companions	K_s Prov	Ev. State
757076	12.351	0.020	9.559	0.017	1.524	0.048			0.947			
757099	13.704	0.020	11.094	0.018	2.708	0.027			2.173			
757137	10.052	0.028	6.722	0.017	1.753	0.025			0.913			RGB
757280	12.133	0.020	10.627	0.018	1.214	0.022			0.870			
757450	15.895	0.020	13.118	0.029	1.199	0.026	0.229	0.15	1.030			
891901	13.631	0.020	11.928	0.020	0.857	0.116			7.356			
891916	15.354	0.020	13.076	0.026	0.453	0.215			8.712			
892010	12.617	0.021	9.041	0.017	0.541	0.024			1.014			
892107	13.131	0.023	10.163	0.017	1.064	0.023			0.940			clump
892195	14.340	0.020	11.814	0.019	2.080	0.017			1.121			
892203	14.033	0.020	11.950	0.020	1.802	0.016			1.245			
892376	15.521	0.020	10.721	0.015	3.038	0.360			16.182			
892667	13.424	0.020	11.818	0.020	0.851	0.015			0.867			
892675	13.860	0.020	11.940	0.020	1.711	0.014			1.065			
892678	12.536	0.021	10.580	0.018	1.024	0.026			0.985	1.000	BinaryCorr	
892713	12.466	0.021	10.509	0.017	0.960	0.023			0.987	1.000	BinaryCorr	

Note. KIC ID, g -mag, K_s -mag, parallax, metallicity, Renormalized Unit-Weight Error (RUWE) flag, number of companions within 4'' detected by Gaia, K_s -mag flag indicating potential corrections compared to 2MASS K_s (empty rows indicate no correction), and giant branch evolutionary state flag from Vrard et al. (2016) and Hon et al. (2018) as input parameters for our sample of 186,548 Kepler stars. A subset of our input parameters is provided here to illustrate the form and format. The full table, in machine-readable format, can be found online.

(This table is available in its entirety in machine-readable form.)

expression:

$$K_{\text{prim}} = K_{\text{sec}} - 2.5 \log(10^{-0.4*(K_{2\text{MASS}} - K_{\text{sec}})} - f(\theta)) \quad (1)$$

$$\text{where } f(\theta) = \begin{cases} 1.0 & \theta \leq 1''.25 \\ -\frac{4}{9}\theta + \frac{14}{9} & 1.25 < \theta < 3''.5 \\ 0.0 & \theta \geq 3''.5 \end{cases}$$

and θ is the angular separation between the primary and the secondary in arcseconds, K_{prim} is the corrected primary K_s magnitude, K_{sec} is the secondary K_s magnitude computed as described above, and $K_{2\text{MASS}}$ is the original magnitude provided by 2MASS. The expression for f above represents the fraction of flux of fainter sources contained within the 2MASS aperture. This expression was derived from a comparison of the Kraus & Hillenbrand (2007) point-spread function fit magnitudes to the 2MASS catalog (Skrutskie et al. 2006) magnitudes. The comparison reveals that $\leq 1''.25$ and $\geq 3''.5$ are the places where 100% and 0% of the flux appear to be captured, and we use a linear relationship both for simplicity and because it is consistent with the binary fitting results of Kraus & Hillenbrand (2007).

We provide the status of these corrections in Table 1, flagging Kepler stars with one Gaia-resolved companion but no correction as “BinDetNoCorr” and those with one resolved companion with a correction as “BinaryCorr.” Stars with multiple resolved companions are flagged as “TerDetNoCorr,” “TerDetBinCorr,” and “TertiaryCorr” depending on whether we computed K_s -magnitude corrections for zero, one, or multiple companions, respectively. We also include the number of companions as an additional column in the table. Finally, we adopted the photometric errors reported by 2MASS.

Unlike the K_s photometry, both the KIC and KIS g photometry require calibration. To convert the KIC g photometry to the SDSS g of the MESA Isochrones and Stellar Tracks (MIST v1.2, Choi et al. 2016; Dotter 2016;

Paxton et al. 2011, 2013, 2015) grid, we used Equation (1) in Pinsonneault et al. (2012). We solved the SDSS–Vega system equations provided in Section 4 of González-Solares et al. (2011) for g_{SDSS} to convert the Vega system KIS photometry back to the SDSS AB system.

Next, we used KIS photometry for all sources where it was both available and the corresponding photometry flag indicated neither saturation nor bad pixels. We used the calibrated KIC photometry for all other stars. Our final input catalog utilized KIS g photometry for 148,410 stars, and KIC photometry for the remaining 38,138 stars. We did not amend the g -band magnitudes for contamination from secondaries like we did for K_s magnitudes, due to the $\approx 1''.5$ effective resolutions of the KIC and KIS catalogs.

Neither Brown et al. (2011) nor Greiss et al. (2012) report uncertainties for individual sources; hence, we computed the errors of our g magnitudes by utilizing the photometric scatter relations provided in Brown et al. and Greiss et al.:

$$\sigma_{g_{\text{KIS}}} = \sqrt{0.02^2 + (1.1 * e^{0.6456 * g_{\text{KIS}} - 16.1181})^2} \quad (2)$$

$$\sigma_{g_{\text{KIC}}} = \sqrt{0.02^2 + (0.01 * (g_{\text{KIC}} - 12))^2}. \quad (3)$$

We report our g and K_s photometry and their errors in Table 1.

2.3. Isochrone Fitting

Isochrone fitting allows the straightforward determination of stellar parameters, such as T_{eff} , masses, and ages, from a set of input observables. We used the most recent MIST models (v1.2 with rotation, Paxton et al. 2011, 2013, 2015; Choi et al. 2016; Dotter 2016), which we interpolated from the grids provided on the MIST website. Our final grid contained ~ 7 million models with 117 ages from 0.1 Gyr to 3.68 Gyr in 52 logarithmic steps and 3.68 Gyr to ≈ 20 Gyr in linear steps of 0.25 Gyr to sufficiently sample both young and old stellar models and avoid preferential “snapping” to sparse bands of the model grid. We chose 20 Gyr as our maximum age because it is the largest age in the MIST grid and it minimizes posterior

truncation. Posterior truncation produces an under/overestimation of derived stellar parameters and an underestimation of their errors; this deleterious effect is inevitable with any finite grid, but we minimized its magnitude by including models up to the grid’s maximum age. We also did not use any pre-main sequence models. The grid had initial metallicities ranging between -2.0 and 0.5 dex with 0.05 dex steps. The grid also accounted for element diffusion, which affects the surface abundances.

Isochrone grids frequently struggle to reproduce empirical constraints from M-dwarfs due in part to the presence of starspots and strong magnetic fields (Boyajian et al. 2012a; Feiden & Chaboyer 2012; Mann et al. 2019). Therefore, we implemented the empirical Mann et al. (2015, 2019) $r - J - T_{\text{eff}}$, M_{K_s} -radius, and M_{K_s} -mass and metallicity relations to compute T_{eff} , stellar radius, and stellar mass, respectively. These relations are mainly calibrated on absolutely flux-calibrated spectra (T_{eff} , radius) and the binary orbital parameters (mass) of nearby M-dwarfs. We do not extrapolate the Mann et al. (2015, 2019) relations and only change models that are within the empirical fits: $1.9 < r - J < 5.5$ mag, $M_{K_s} \geq 4.0$ mag, and $[\text{Fe}/\text{H}] \geq -0.6$ dex. Given these photometric and metallicity constraints, we applied the Mann empirical relations to stars with masses below $\approx 0.75 M_{\odot}$; hence, we revised $\approx 178,000$ models. Redder M-dwarf models ($r - J \geq 5.5$ or $T_{\text{eff}} \lesssim 2800$ K) that would require extrapolation are dropped from our grid altogether.

Utilizing solely the individual g and K_s photometric errors added in quadrature yielded T_{eff} fractional errors of $\lesssim 1\%$. These errors are too small given that there are systematic errors in interferometric angular diameters which, in turn, set the fundamental limit on T_{eff} errors: $\approx 2\%$. Therefore, we computed the best-fitting 12th-order polynomial to the relationship of T_{eff} to the $g - K_s$ color of all models within our model grid using numpy’s polyfit routine (Figure 2). We chose a 12th-order polynomial because all lower-order polynomials do not accurately trace the center of the T_{eff} -color curve, while higher-order polynomials minimally improve the resulting correlation coefficient. For our eventual isochrone fitting, we adopted the maximum $g - K_s$ error: either (1) the g and K_s errors added in quadrature or (2) the 2% T_{eff} curve-computed value of the input $g - K_s$ color in the bottom panel of Figure 2. We also found that we underestimate our stellar mass errors for M-dwarfs because of the tight Mann et al. (2019) M_{K_s} -mass relation. Hence, we inflate our M-dwarf ($g - K_s \geq 4$, $M_{K_s} \geq 4$) M_{K_s} uncertainties by adding an error term in quadrature corresponding to a 2.1% mass error, identical to the scatter in the empirical relation (Mann et al. 2019).

We employed `isoclassify` (Huber et al. 2017) and the Green et al. (2019) reddening map to derive stellar parameters from our input observables: (1) SDSS g and 2MASS K_s photometry, (2) Gaia DR2 parallaxes (Gaia Collaboration et al. 2016, 2018; Arenou et al. 2018; Lindegren et al. 2018), (3) red giant evolutionary state flags (RGB versus the Red Clump; Vrad et al. 2016; Hon et al. 2018), and (4) spectroscopic metallicities from the Kepler Stellar Properties Catalog (Mathur et al. 2017), the California-Kepler Survey (CKS; Petigura et al. 2017), APOGEE DR14 (Majewski et al. 2017; Abolfathi et al. 2018), and LAMOST DR5 (Zhao et al. 2012; Ren et al. 2018), where available. We computed the 16th, 50th, and 84th percentile values for T_{eff} , $\log g$, $[\text{Fe}/\text{H}]$, radius, mass, density,

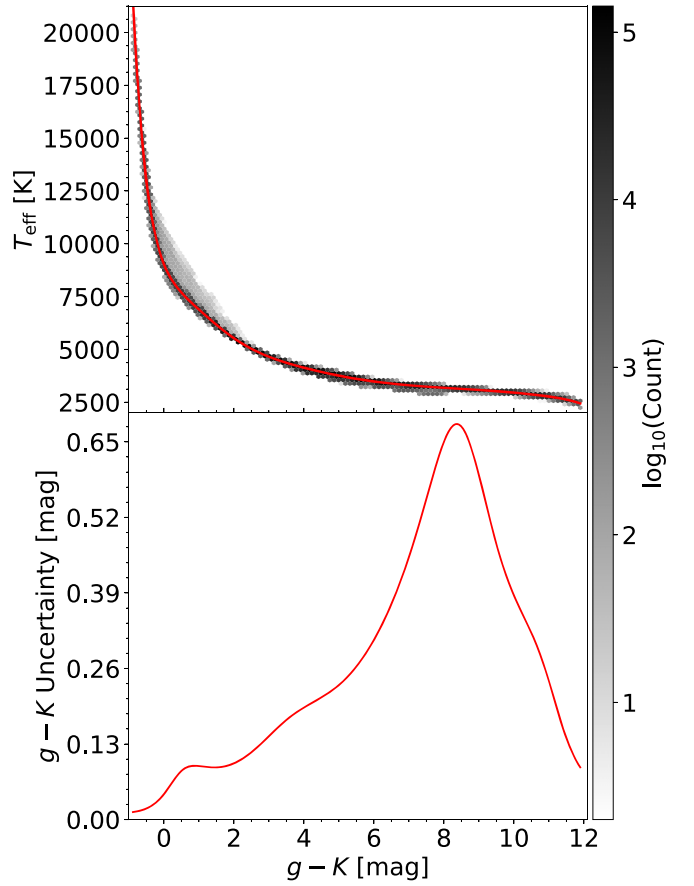


Figure 2. Corrections to the initial $g - K_s$ magnitude errors utilizing our interpolated MIST model grid. Top: T_{eff} vs. $g - K_s$ color for our model grid, with the logarithmic density of points illustrated by the two-dimensional greyscale histogram with the corresponding colorbar. We plot the best-fit 12th-order polynomial in red. Bottom: the red curve represents the required $g - K_s$ minimum uncertainty to reach a 2% T_{eff} error for all stars, dependent on their $g - K_s$ color.

age, distance, and V-band extinction from the marginalized posteriors.

We used conservative 0.15 dex metallicity errors instead of the quoted errors for all stars with spectroscopic metallicity constraints because the pipeline-to-pipeline uncertainty in metallicities is $\gtrsim 0.1$ dex (Furlan et al. 2018). For stars that do not have spectroscopic metallicity constraints ($\sim 120,000$), we used a prior centered on solar metallicity with a standard deviation of ~ 0.20 dex, which is appropriate for the Kepler field (Dong et al. 2014). As demonstrated in Howes et al. (2019), a lack of a metallicity constraint typically results in an age-metallicity degeneracy depending on the stellar properties/evolutionary state.

2.4. Accounting for Binaries

As we described in Section 2.2, we addressed Gaia-detected stellar companions within $4''$ that contaminate the 2MASS K_s photometry. In addition, Gaia DR2 astrometric flags appear to be another useful tool for identifying binaries that are not resolved by the satellite. For instance, Evans (2018), Rizzuto et al. (2018), and Ziegler et al. (2018) have already demonstrated that these astrometric flags are useful for identifying smaller separation binaries that are not spatially resolved.

To identify these smaller separation binaries, we computed Gaia DR2’s Renormalized Unit-Weight Error (RUWE) by interpolating the tabular data detailed in Lindegren (2018). RUWE is the magnitude- and color-independent re-normalization of the astrometric χ^2 of Gaia DR2 (unit-weight error or UWE). The RUWE values are reported for all Kepler stars in the “RUWE” column of Table 1. Any stars with $\text{RUWE} \gtrsim 1.2$ are likely to be binaries (A. Kraus et al. 2020, in preparation). Of the 186,548 stars remaining, 170,845 had $\text{RUWE} \leq 1.4$ and 164,736 had $\text{RUWE} \leq 1.2$; we provide RUWE values for every star, but analyze both those with high RUWE and low RUWE similarly. Any derived stellar parameters for these high-RUWE stars should be treated with extra caution, along with those that have “NoCorr” in their K_s magnitude flags also in Table 1. We also did not amend the input magnitudes of Kepler planet host stars with adaptive optics-detected stellar companions (Furlan et al. 2017) in order to preserve the homogeneity of our catalog.

3. Validating the Output Stellar Parameters

3.1. Accuracy of Derived Effective Temperatures

To ensure our grid-computed stellar effective temperatures are accurate, we compared them to interferometric T_{eff} measurements for a sample of 108 stars from Boyajian et al. (2013) and Huang et al. (2015) with Tycho B and V photometry as well as 2MASS K_s photometry. Although these stars do have g -band photometry from the American Association of Variable Star Observers Photometric All-Sky Survey (Henden et al. 2018), it is saturated for these stars. Therefore, we had to convert the Tycho B and V photometry into SDSS g using the following procedure. First, we converted Tycho B (B_T) and V (V_T) photometry into Johnson B and V photometry using Table 2 from Bessell (2000). Then we converted our B and V magnitudes into g magnitudes with the transformation given in the bottom portion of Table 1 of Jester et al. (2005). We found color-dependent systematics in our comparison of g and V_T magnitudes, which we eliminated utilizing the Tycho $B_T - V_T$ colors. Hence, we computed our interferometric sample g magnitudes as follows:

$$g(B, V, B_T, V_T) = V - 0.03 + 0.60(B - V) - 0.11(B_T - V_T) + 0.09(B_T - V_T)^2. \quad (4)$$

We present the comparison to the interferometric stars in Figure 3. Based on the reported median shifts and median absolute deviations in T_{eff} for the dwarfs and giants, our derived Kepler stellar T_{eff} appear to be accurate within our 2% T_{eff} errors for the T_{eff} range covered here. We found that the residuals of stars with K_s -band errors ≥ 0.25 mag were particularly discrepant in their residuals, and ignore them here. For solar T_{eff} and late-G and early-K-dwarfs, we estimate hotter T_{eff} than interferometric determinations by 50–60 K, while we underestimate the T_{eff} of F-dwarfs by ≈ 100 K, also within our reported errors. We caution that the derived effective temperatures for M-dwarfs are systematically offset by ≈ 75 K, and the giants appear to demonstrate a strong trend with interferometric T_{eff} . The trend in the giants is likely due to insufficient color transformations and/or saturated photometry, as all interferometric stars are close to 2MASS saturation due to their proximity. Upcoming work on M-dwarfs (E. Gaidos et al. 2020, in preparation) and previously published work on the

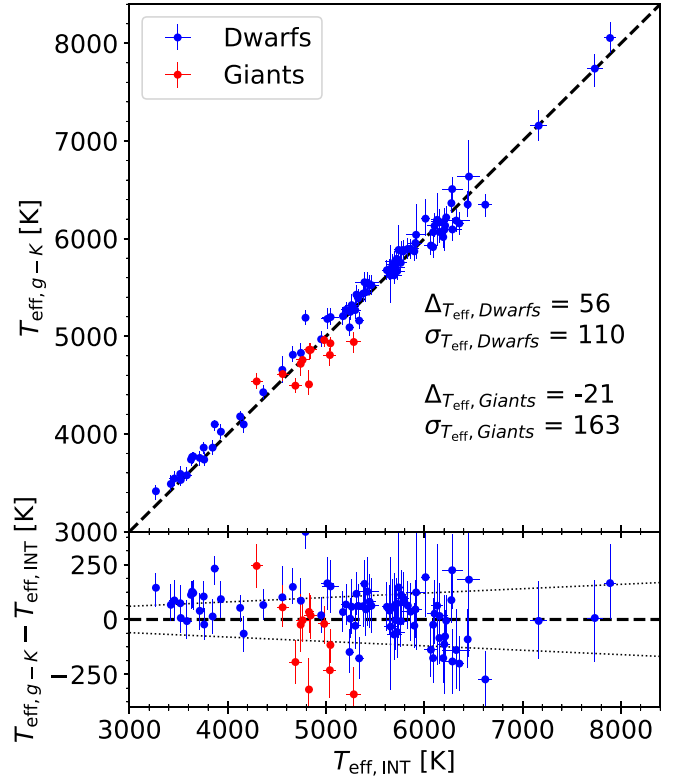


Figure 3. Top: grid-modeled T_{eff} vs. interferometrically derived T_{eff} (Boyajian et al. 2013; Huang et al. 2015). We plot the dwarfs as blue points and giants as red points, as well as their respective uncertainties. The black dashed line is the 1:1 line. The text in the plot indicates the median shifts (Δ) and the median absolute deviations (σ) for both the dwarfs and giants, as labeled. Bottom: residuals as a function of interferometric T_{eff} . The black dotted lines represent 2% fractional uncertainties above and below equality.

APOGEE Kepler Giants with asteroseismic data (APOKASC; Serenelli et al. 2017; Pinsonneault et al. 2018) are more specialized and hence better alternatives to the data we derive here for these specific samples of stars.

3.2. Binary Effects on Stellar Properties

Regardless of how much we account for binaries in our above analysis, there will inevitably be systems that still have unresolved, unidentified stellar companions. These companions will be at angular separations less than both the Gaia photometric resolution and RUWE effective resolution, and they will affect our estimates of the stellar properties. To quantify these effects, we took a set of models from our MIST grid of isochrones (Section 2.3) at a particular age and initial metallicity and combined their photometry with that from all stellar models less than or equal to the mass of the primary in that same isochrone. We then ran the photometry for each of these modified models through *isoclassify*. We used typical Kepler field Gaia DR2 parallaxes and errors and primary surface metallicities, assuming zero reddening for simplicity. We chose primary stars of 0.85 , 1.15 , and $1.45 M_{\odot}$, at the mode of ages (2.59 Gyr) and metallicities (0.0 dex) determined for the Kepler target sample.

In Figure 4, we plot the effect of binarity on the age, radius, and T_{eff} estimates for our representative set of age, metallicity, and stellar masses for Kepler stars. In red, yellow, and blue are the binary fitting results for the $0.85 M_{\odot}$, $1.15 M_{\odot}$, and $1.45 M_{\odot}$ primaries, respectively.

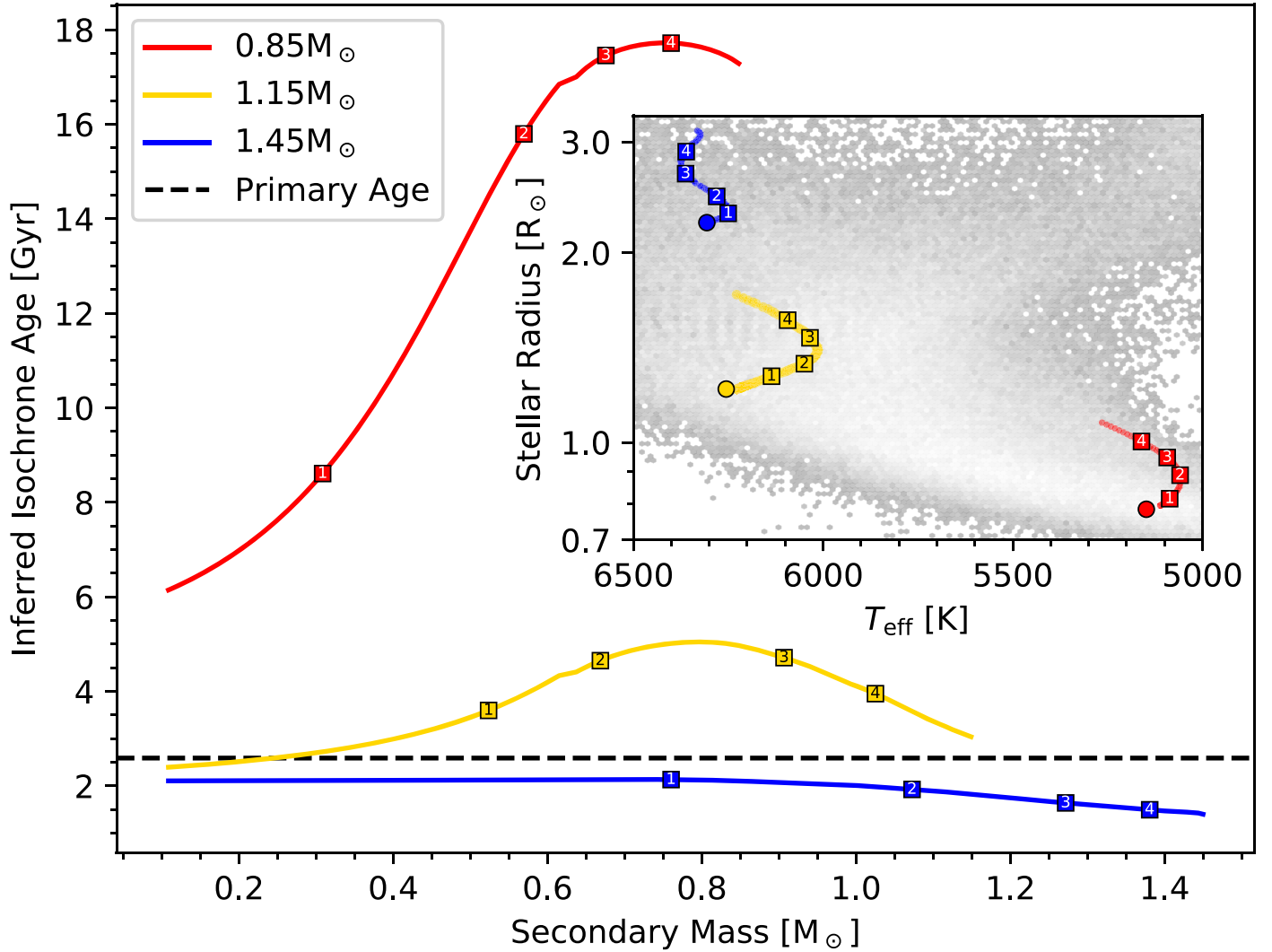


Figure 4. *isoclassify*-derived age vs. the secondary’s mass for simulated binary systems with primary masses $0.85 M_{\odot}$ (red), $1.15 M_{\odot}$ (yellow), and $1.45 M_{\odot}$ (blue). Both the primary and the secondaries are 2.59 Gyr old (black dashed line) and have initial surface metallicities of 0.0 dex. The inset shows where the sole primary and primary + secondary combinations occur on the Hertzsprung–Russell diagram. The large, color-matched circles represent the primary’s properties and the colored curves are the primary + secondary composite parameters. The color- and number-coordinated squares represent the same solutions on both diagrams. To guide the eye, we plot a two-dimensional histogram for the entire Kepler catalog underneath, where darker grays represent areas with lower logarithmic number density.

We see that the age of the $0.85 M_{\odot}$ star is severely affected by the addition of a secondary: even with the least massive companion, the age is overestimated by ≈ 3.5 Gyr. This is expected: isochrones have very little predictive power for stellar ages for such low-mass stars, as they do not evolve significantly on timescales similar to the age of the universe. Hence, the age posteriors for even the lowest-mass secondaries approach our flat prior and herd age estimates toward the center of our age distribution (≈ 10 Gyr). The rest of the age–secondary mass curve behaves mostly like we expect for the $0.85 M_{\odot}$ star. As we add more massive secondaries, the binary system will mimic the photometry of an older, more evolved star. The slight turnover to younger ages at the end of the curve means that the system mimics the photometry of a slightly younger but more massive primary. For all unresolved, equal-mass binary systems, we would expect to determine the same effective temperature as a single star, but a $\sqrt{2}$ times larger radius. However, we do not see this in the inset Hertzsprung–Russell (H-R) diagram for the $0.85 M_{\odot}$ star because of the

grid-edge behavior of the red “backwards-c” curve. Due to the finite age limit of the MIST models, the only single-star models that are consistent with the higher luminosity of the binary are hotter, not purely larger, older models. The “backward-c” curve also occurs because secondaries of different colors and magnitudes affect the resultant color of the system: the least massive, coolest secondaries hardly contribute to the overall photometry and the most massive secondaries are very similar in color and magnitude to the primary. Any unresolved companion $\geq 0.3 M_{\odot}$ will significantly affect the derived age.

The age bias is reduced for the $1.15 M_{\odot}$ system (yellow), which only becomes significantly affected once the secondary mass reaches half the mass of the primary. Eventually, the curve turns over as more massive secondaries push the system to larger luminosities and hence more evolved versions of higher-mass stars at similar ages. In the H-R diagram, this system appears as a “backward-c,” mimicking the behavior of the $0.85 M_{\odot}$ primary system. However, the $1.15 M_{\odot}$ primary star system produces larger radii and similar T_{eff} , and is

unaffected by grid-edge effects when combined with similar-mass secondaries. Typical age uncertainties (≈ 2 Gyr) for these stars mean that the majority of the inferred age curve is within 1σ of the primary star’s true age except for the broad bump occurring for $0.6\text{--}1.0 M_{\odot}$.

Finally, the blue $1.45 M_{\odot}$ binary age–secondary mass curve in Figure 4 shows a slight underestimation of the ages of high-mass stars. The curve increases only slightly until the secondary masses exceed $0.8 M_{\odot}$, at which point it begins to predict even smaller ages. This slight increase in the derived system age occurs because the system’s photometry appears to move slowly along its evolutionary track for companions $\leq 0.8 M_{\odot}$. For larger secondary masses, we determine smaller ages because the combined photometry places the system on higher-mass tracks where stars are both younger and at similar T_{eff} . The behavior of this system in the H-R diagram is qualitatively different from the other systems, as the $1.45 M_{\odot}$ star is past the main-sequence turn-off at 2.59 Gyr. The top of the blue curve represents an equal-mass binary, while evolved, smaller-mass secondaries make up the top half of the “S” pattern (i.e., squares labeled 3 and 4). The remaining lower-mass secondaries produce the bottom half of the “S” curve (i.e., squares labeled 1 and 2), similar to the “backward-c” seen in the 0.85 and $1.15 M_{\odot}$ curves. Typical age uncertainties for these stars (≈ 0.5 Gyr) place the curve within uncertainties of the primary star’s input age except for the most massive ($\gtrsim 1 M_{\odot}$) secondaries.

Ultimately, we find that while binaries will affect our derived stellar ages, the magnitude of the effect scales with the typical error bars for each type of star. Low-mass star ages are biased significantly by as many as 10 Gyr by the presence of a mass ratio ≥ 0.5 secondary companion, but the age uncertainties we derive for these stars are on the order of 6–7 Gyr. For higher-mass stars, we find our ages are biased by only a few Gyr, where the error bars are typically on the order of 1 Gyr.

3.3. Accuracy of Derived Stellar Ages

3.3.1. Cluster Ages

To independently confirm our derived stellar ages, we used the 1 Gyr open cluster NGC 6811 (Meibom et al. 2011). We did not use NGC 6791 because its main-sequence turn-off was too faint for our photometry.

We utilized the Gaia DR2-KIC matches provided by D. Godoy-Rivera et al. (2020, in preparation) as our sample, resulting in 287 matches. We used the same output parameters as those provided in Table 2. This is because the cluster’s metallicity $[\text{Fe}/\text{H}] = 0.05$ dex (Molenda-Zakowicz et al. 2014) is well within the Gaussian prior centered at solar metallicity (0.0 dex) with the standard deviation of 0.2 dex that we assume for Kepler field stars without spectroscopic metallicities. To ensure that our NGC 6811 ages are reliable, we removed all stars with terminal-age main sequence (TAMS) ≥ 14 Gyr and all giant stars using an ad hoc cut in stellar radius- T_{eff} space, similar to that of Equation (1) in Fulton et al. (2017): $\frac{R_{\star}}{R_{\odot}} < 10^{0.00035(T_{\text{eff}} - 4500) + 0.15}$. In addition, we removed all stars with $\text{RUWE} \geq 1.2$ to minimize potential age-contaminating binaries (see Section 3.2). This left us with 146 matches in NGC 6811.

NGC 6811 shows good agreement between the literature cluster age and the isochrone-dependent ages we derive here

(Figure 5). There are a few stars in the high-age tail of the distribution, but these are stars that populate the coolest part of the remaining main sequence (TAMS $\lesssim 14$ Gyr) and have the largest radii. Although we already removed all stars with high RUWE, these targets might be close or unresolved binaries with a hotter and cooler component, which biases the resulting photometry to indicate a cooler star of larger radius and hence older age. This effect can be large, as discussed in Section 3.2 and demonstrated in Figure 4. However, only a few stars lie in such areas of parameter space.

3.3.2. Asteroseismic Ages

We also compared our derived ages to those of Kepler stars that have asteroseismic ages. We utilized the “boutique” frequency-modeled ages from the Kepler legacy sample detailed in Silva Aguirre et al. (2017), which includes results from a number of analysis pipelines. In Figure 6, the ages that we derive are in reasonable agreement with those provided by a variety of asteroseismic pipelines. In the top panel of Figure 6, we plot each asteroseismic pipeline as a different color, so horizontal rows of colored points indicate results for the same star. The horizontal scatter of the colored points are typically larger than their reported errors, which indicates that systematic pipeline differences dominate. Any deviations from the 1:1 dashed line are sufficiently accounted for by a combination of the typical error bars (bottom right, top panel) and any systematic scatter depending on the asteroseismic pipeline one chooses. In addition, the asteroseismic ages do not fall above the age of the universe (likely due to a model grid age-cutoff), hence we see the largest differences at older ages. However, even the largest discrepancies are within 1σ of the 1:1 line.

Ultimately, we report a median offset and scatter of 5% and 29%, respectively. We conclude that our isochrone-derived ages are consistent with ages determined through more precise methods within the uncertainties that we report.

3.3.3. Kinematic Ages

In Figure 7 we display a comparison of our isochrone-derived ages for Kepler exoplanet host stars with spectroscopic metallicities, $\text{RUWE} \leq 1.2$, $\text{TAMS} \leq 20$ Gyr, and $\text{iso_gof} \geq 0.99$ with their total space (UVW) velocities relative to the local standard of rest (LSR). We computed the latter from Gaia DR2 proper motions and parallaxes and CKS (Petigura et al. 2017) radial velocities, following the method outlined in Newton et al. (2016): we used Equation (1) from Johnson & Soderblom (1987) and the transformation matrix from Perryman et al. (1997) with the LSR defined by Schönrich et al. (2010). We computed uncertainties from Equation (2) in Johnson & Soderblom (1987), where we used 0.1 km s^{-1} radial velocity uncertainties (Petigura et al. 2017) and the formal Gaia DR2 uncertainties on proper motions and parallaxes (Gaia Collaboration et al. 2018; Lindegren et al. 2018). Most total space velocity uncertainties are smaller than the markers. Figure 7 reveals that both the mean total space velocities and their dispersion increase at higher isochrone ages, which matches the expectations that old stars have had more time to be perturbed by gravitational interactions with other stars (Soderblom 2010; Newton et al. 2016).

Table 2
Gaia–Kepler Stellar Output Parameters

KIC ID	T_{eff} (K)	$\log g$ (dex)	[Fe/H]	M_* (M_{\odot})	R_* (R_{\odot})	$\log \rho_*$ (ρ_{\odot})	$\log L_*$ (L_{\odot})	Age (Gyr)	Distance (pc)	A_V (mag)	GOF	TAMS (Gyr)
757076	5052 $^{+103}_{-86}$	3.37 $^{+0.07}_{-0.08}$	-0.14 $^{+0.16}_{-0.19}$	1.40 $^{+0.18}_{-0.22}$	4.00 $^{+0.14}_{-0.15}$	-1.67 $^{+2.38}_{-2.40}$	0.98 $^{+0.04}_{-0.10}$	2.5 $^{+1.8}_{-0.7}$	651 $^{+22}_{-21}$	0.37	1.0000	2.7
757099	5364 $^{+102}_{-84}$	4.32 $^{+0.04}_{-0.03}$	0.08 $^{+0.14}_{-0.13}$	0.87 $^{+0.05}_{-0.04}$	1.07 $^{+0.02}_{-0.02}$	-0.15 $^{+1.13}_{-1.25}$	-0.07 $^{+1.23}_{-1.29}$	15.2 $^{+3.0}_{-4.0}$	367 $^{+7}_{-6}$	0.34	1.0000	17.2
757137	4628 $^{+84}_{-76}$	2.39 $^{+0.08}_{-0.09}$	-0.11 $^{+0.15}_{-0.17}$	1.67 $^{+0.31}_{-0.30}$	13.59 $^{+0.32}_{-0.33}$	-3.19 $^{+3.83}_{-3.93}$	1.88 $^{+0.74}_{-0.64}$	1.5 $^{+1.1}_{-0.6}$	568 $^{+12}_{-11}$	0.34	1.0000	1.7
757280	6856 $^{+144}_{-139}$	3.83 $^{+0.03}_{-0.03}$	-0.03 $^{+0.21}_{-0.11}$	1.71 $^{+0.09}_{-0.09}$	2.61 $^{+0.07}_{-0.07}$	-1.02 $^{+2.01}_{-2.06}$	1.13 $^{+0.05}_{-0.00}$	1.2 $^{+0.2}_{-0.2}$	822 $^{+19}_{-21}$	0.50	1.0000	1.6
757450	5301 $^{+111}_{-103}$	4.43 $^{+0.05}_{-0.04}$	0.24 $^{+0.13}_{-0.13}$	0.91 $^{+0.06}_{-0.06}$	0.96 $^{+0.03}_{-0.03}$	0.01 $^{+0.85}_{-0.91}$	-0.18 $^{+1.24}_{-1.28}$	9.5 $^{+5.4}_{-5.1}$	829 $^{+24}_{-23}$	0.46	1.0000	16.1
891901	6350 $^{+130}_{-131}$	3.96 $^{+0.10}_{-0.09}$	0.02 $^{+0.15}_{-0.14}$	1.41 $^{+0.12}_{-0.12}$	2.02 $^{+0.28}_{-0.26}$	-0.80 $^{+1.16}_{-1.35}$	0.78 $^{+0.30}_{-0.17}$	2.2 $^{+0.7}_{-0.5}$	1122 $^{+156}_{-146}$	0.34	1.0000	2.9
891916	5650 $^{+131}_{-137}$	4.13 $^{+0.22}_{-0.25}$	0.02 $^{+0.15}_{-0.17}$	1.01 $^{+0.16}_{-0.11}$	1.35 $^{+0.59}_{-0.36}$	-0.46 $^{+0.37}_{-0.68}$	0.24 $^{+0.25}_{-0.08}$	7.6 $^{+3.7}_{-3.3}$	1193 $^{+515}_{-322}$	0.36	1.0000	9.7
892010	4555 $^{+141}_{-92}$	2.30 $^{+0.16}_{-0.12}$	-0.02 $^{+0.16}_{-0.20}$	1.71 $^{+0.71}_{-0.40}$	15.19 $^{+0.77}_{-0.76}$	-3.32 $^{+3.65}_{-3.92}$	1.96 $^{+1.07}_{-0.99}$	1.4 $^{+1.7}_{-0.8}$	1832 $^{+87}_{-87}$	0.37	1.0000	1.6
892107	4894 $^{+83}_{-85}$	3.24 $^{+0.08}_{-0.10}$	-0.05 $^{+0.14}_{-0.15}$	1.24 $^{+0.21}_{-0.24}$	4.41 $^{+0.13}_{-0.12}$	-1.84 $^{+2.53}_{-2.51}$	1.01 $^{+0.11}_{-0.15}$	4.1 $^{+4.5}_{-1.7}$	936 $^{+25}_{-24}$	0.28	1.0000	4.3
892195	5333 $^{+101}_{-84}$	4.37 $^{+0.04}_{-0.03}$	0.07 $^{+0.14}_{-0.13}$	0.86 $^{+0.06}_{-0.04}$	1.00 $^{+0.02}_{-0.02}$	-0.07 $^{+1.02}_{-1.17}$	-0.14 $^{+1.31}_{-1.36}$	14.3 $^{+3.5}_{-4.6}$	479 $^{+8}_{-8}$	0.25	1.0000	18.0
892203	5712 $^{+108}_{-105}$	4.39 $^{+0.04}_{-0.04}$	0.01 $^{+0.15}_{-0.15}$	0.97 $^{+0.07}_{-0.07}$	1.04 $^{+0.02}_{-0.02}$	-0.07 $^{+0.99}_{-1.02}$	0.02 $^{+1.14}_{-1.19}$	6.5 $^{+4.1}_{-3.5}$	553 $^{+10}_{-10}$	0.22	1.0000	11.3
892667	6704 $^{+148}_{-128}$	3.95 $^{+0.03}_{-0.04}$	0.01 $^{+0.16}_{-0.17}$	1.55 $^{+0.08}_{-0.09}$	2.17 $^{+0.06}_{-0.06}$	-0.83 $^{+1.81}_{-1.84}$	0.94 $^{+0.15}_{-0.19}$	1.6 $^{+0.3}_{-0.3}$	1171 $^{+29}_{-29}$	0.50	1.0000	2.2
892675	5929 $^{+108}_{-108}$	4.39 $^{+0.04}_{-0.04}$	-0.02 $^{+0.14}_{-0.16}$	1.04 $^{+0.07}_{-0.08}$	1.08 $^{+0.02}_{-0.02}$	-0.09 $^{+1.09}_{-1.06}$	0.12 $^{+1.05}_{-1.08}$	3.6 $^{+3.2}_{-2.2}$	583 $^{+11}_{-10}$	0.25	1.0000	8.5
892678	5890 $^{+121}_{-114}$	3.57 $^{+0.03}_{-0.03}$	0.02 $^{+0.19}_{-0.22}$	1.58 $^{+0.07}_{-0.06}$	3.38 $^{+0.12}_{-0.10}$	-1.40 $^{+2.37}_{-2.41}$	1.10 $^{+0.11}_{-0.01}$	1.8 $^{+0.1}_{-0.1}$	967 $^{+34}_{-17}$	0.25	1.0000	2.1
892713	6238 $^{+123}_{-129}$	3.55 $^{+0.07}_{-0.04}$	0.08 $^{+0.21}_{-0.20}$	1.73 $^{+0.22}_{-0.09}$	3.64 $^{+0.12}_{-0.11}$	-1.45 $^{+2.20}_{-2.40}$	1.25 $^{+0.25}_{-0.15}$	1.4 $^{+0.1}_{-0.3}$	1033 $^{+32}_{-29}$	0.50	1.0000	1.7
892718	5000 $^{+97}_{-90}$	4.57 $^{+0.03}_{-0.04}$	-0.08 $^{+0.14}_{-0.13}$	0.78 $^{+0.04}_{-0.05}$	0.76 $^{+0.03}_{-0.03}$	0.25 $^{+0.77}_{-0.74}$	-0.49 $^{+1.46}_{-1.50}$	6.3 $^{+7.1*}_{-4.5*}$	874 $^{+32}_{-31}$	0.31	1.0000	23.1

Note. KIC ID, effective temperature, surface gravity, surface metallicity, stellar mass, stellar radius, density, luminosity, age, distance, V -magnitude extinction, combined likelihood goodness-of-fit (GOF), and terminal-age main-sequence (TAMS) parameters and their errors for 186,301 Kepler stars, output from our isochrone placement routine detailed in Section 2. Ages with asterisks are either those with uninformative posteriors (TAMS ≥ 20 Gyr) or unreliable ages (GOF ≤ 0.99). Stars within Table 1 and not in this table have fewer than 10 models within 4σ of the input observables. A subset of our output parameters is provided here to illustrate the form and format. The full table, in machine-readable format, can be found online.

(This table is available in its entirety in machine-readable form.)

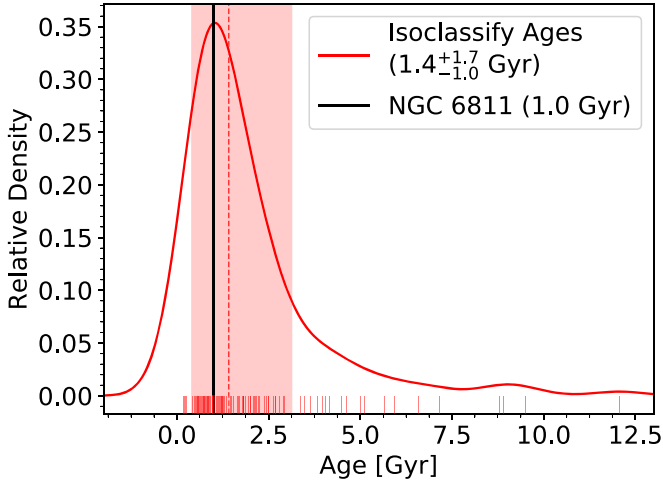


Figure 5. Age comparison for open cluster NGC 6811 at 1.0 Gyr (Meibom et al. 2011), located within the Kepler field. The solid red distribution represents the Gaussian Kernel Density Estimate of the ages of individual stars within each cluster as derived in this work, with the median and the 1σ confidence interval represented by the vertical red dashed line and shaded region, respectively. We use Scott’s rule (Scott 1992) bandwidths to produce the overall distribution. Translucent vertical red lines represent the inferred ages for each star within the sample. The black, solid vertical line represents the cluster ages from the literature in each panel. We only include non-giant stellar constituents with TAMS ≤ 14 Gyr.

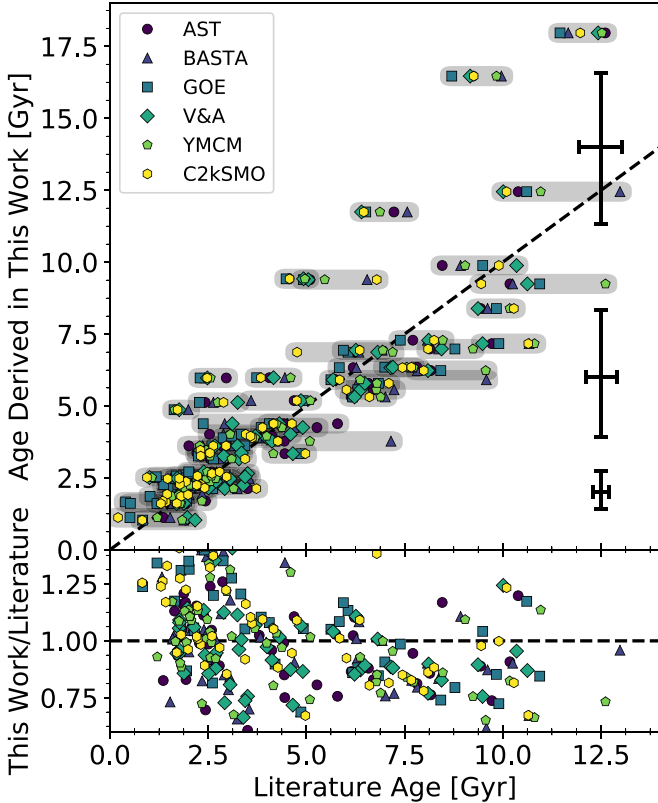


Figure 6. Ages derived in this work vs. those with frequency-modeled asteroseismic ages from a variety of pipelines (Silva Aguirre et al. 2017). The black dashed line represents agreement. The various colors/shapes represent the different pipelines. The translucent gray rounded rectangles represent the ranges of age estimates from different asteroseismic pipelines for each system, which includes not only seismic differences but also differences between model grids. The bottom panel is the ratio of the two age determinations. We have also plotted median error bars in the right-hand portion of the top panel, where, from bottom to top, the error bars represent the median uncertainties of stars with isochrone ages between 0–4 Gyr, 4–8 Gyr, and ≥ 8 Gyr, respectively.

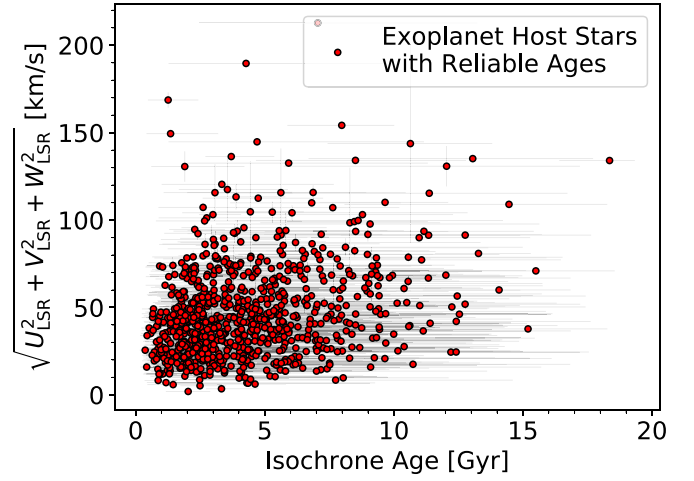


Figure 7. Total space (UVW) velocities relative to the local standard of rest derived from Gaia DR2 proper motions and parallaxes and CKS (Petigura et al. 2017) radial velocities versus isochrone ages computed in this work for Kepler exoplanet host stars with reliable ages. In this case, host stars with reliable ages are dwarfs with spectroscopic metallicities, RUWE ≤ 1.2 , TAMS ≤ 20 Gyr, and $\text{iso_gof} \geq 0.99$. We plot uncertainties as gray, translucent error bars.

4. Revised Properties of Kepler Stars

4.1. Catalog Description

Here, we investigate the properties of all 186,301 Kepler stars. We tabulate stellar T_{eff} , $\log g$, metallicity, mass, radius, luminosity, mean stellar density, age, distance, and extinction in Table 2. Some stars in Table 1 are not included in Table 2. These stars have fewer than 10 models within 4σ of their observational uncertainties, so their under-sampled posteriors are insufficiently constrained. Stars without a solution frequently appear in unphysical areas of parameter space for a single star, such as above the lower main sequence. In addition, we flagged the ages (and corresponding uncertainties) of stars which we deem unreliable or uninformative with asterisks. Unreliable ages are flagged according to the goodness-of-fit parameter (GOF), which is computed using the overall likelihood of the closest model grid point to the set of input observables found in Table 1. We provide the GOF parameter in Table 2; we recommend treating any GOF ≤ 0.99 stars with extra caution. Stars with GOF ≤ 0.99 are outliers in stellar radius– T_{eff} space and they have extremely small fractional error bars compared to typical stars within our catalog. We chose 0.99 as our GOF cut because it represents a compromise between keeping too many outliers (GOF ≤ 0.9) and removing too many stars with reasonable solutions (GOF ≤ 0.999) based on the density of our computed grid.

We flag uninformative ages based on the TAMS for that star. We compute the TAMS by performing 2D interpolation on MIST evolutionary tracks of stars of similar mass to each of our derived stellar masses. If the TAMS of the star is greater than the maximum age of our grid, 20 Gyr, we do not expect to derive any informative age information from that star, given the observational uncertainties and the limitations of isochrone placement. We choose to use 20 Gyr as our age cutoff rather than the age of the universe because we still determine informative, non-truncated age posteriors for stars older than ≈ 14 Gyr from our 20 Gyr maximum age grid. About 14% of stars within the Kepler field have TAMS ≥ 20 Gyr. Therefore, most K- and all M-dwarfs have uninformative ages, as these

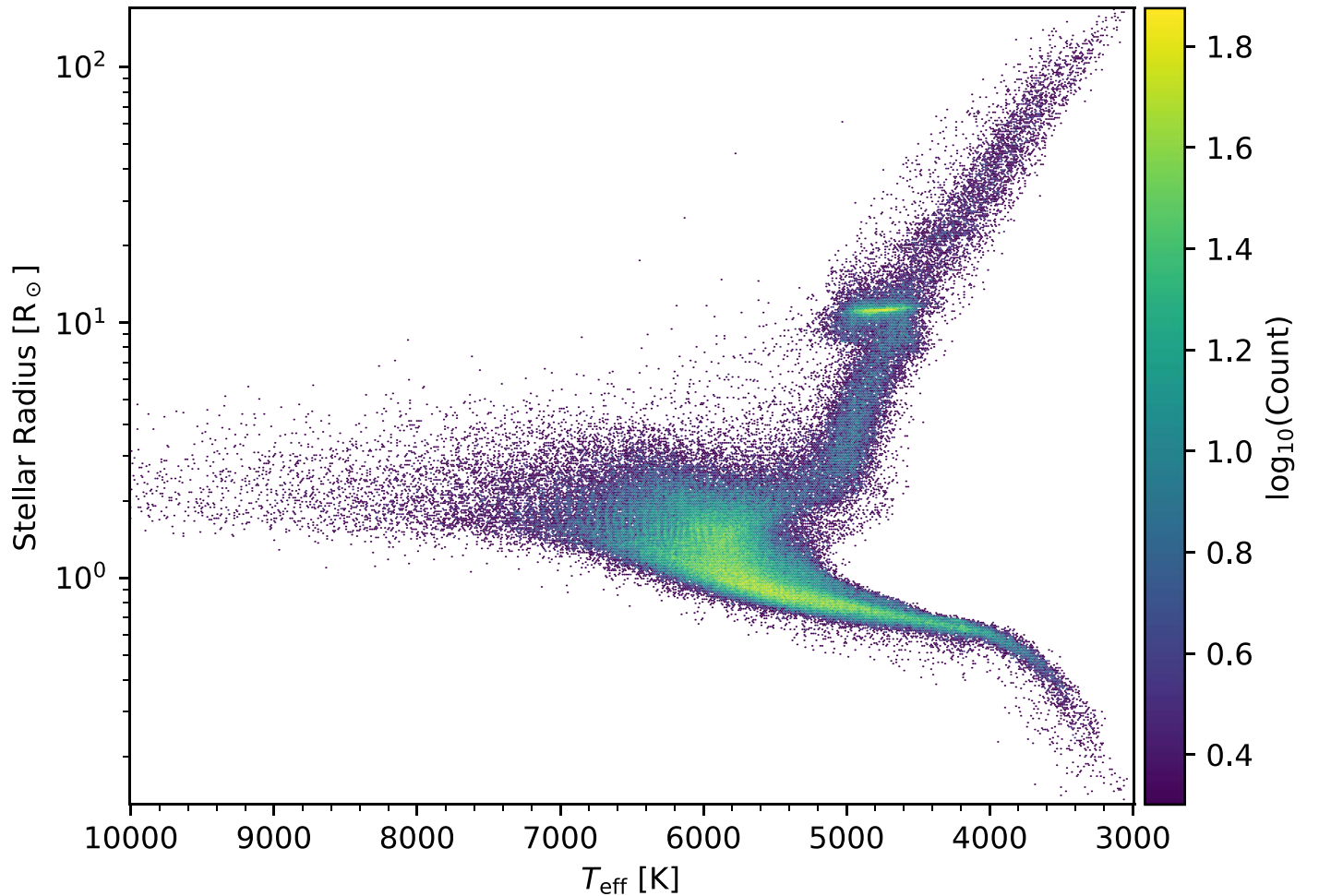


Figure 8. Radius vs. effective temperature for $\sim 186,000$ Kepler stars with radii and T_{eff} derived based on Gaia DR2 parallaxes and $g - K$ photometry presented above. Color-coding represents logarithmic number density.

stars have not had enough time within the age of the universe to evolve substantially in the H-R diagram.

4.2. The Grid-Modeled H-R Diagram of Kepler Stars

Figure 8 shows stellar radius versus effective temperature for the Kepler stars with grid-modeled radii and T_{eff} determined by this work. We see a clear main sequence, from M-dwarfs at $T_{\text{eff}} = 3000$ K and $R_* \approx 0.2 R_{\odot}$, through A-stars at $T_{\text{eff}} \lesssim 9000$ K and $R_* \approx 2 R_{\odot}$. The main-sequence turn-off at $T_{\text{eff}} \approx 6000$ K and $R_* \approx 2 R_{\odot}$ is visible, along with the giant branch. We identify the “red clump” as the concentration of stars surrounding $T_{\text{eff}} \approx 4900$ K and $R_* \approx 11 R_{\odot}$. As expected, the Kepler catalog is dominated by F- and G-type stars as a result of the selection bias for solar-type stars to detect Earth-like transiting exoplanets (Batalha et al. 2010).

Unlike all previous KSPCs (Huber et al. 2014; Mathur et al. 2017; Berger et al. 2018b), the T_{eff} gap around 4000 K is no longer present. We expect this given that we are deriving T_{eff} from our continuous set of input $g - K_s$ colors and the Mann et al. (2015) color- T_{eff} relation overlaps at 4200 K, as we discussed in Sections 2.2 and 3.1. Stellar input observables and their uncertainties must include at least 10 MIST models to produce a solution, so we do not find any stellar solutions outside of our model grid. Therefore, we do not report output

parameters for 247 stars in Table 2. While some stars fall below the nominal main sequence, the discrepancies are not as large as those reported by Berger et al. (2018b). A number of stars (~ 500) below the main sequence that are inferred to be subdwarfs ($T_{\text{eff}} = 3600\text{--}5400$ K and $R_* \leq 0.6 R_{\odot}$) or in other extreme parameter regimes could have erroneous T_{eff} values or excess noise in the astrometry according to Gaia DR2. In addition, if we ignore all stars with $\text{RUWE} \geq 1.4$ (reducing 186,301 to $\approx 170,000$ stars), the putative subdwarfs disappear, as well as a number of other stars in sparsely populated areas of the H-R diagram. Therefore, most of the inferred subdwarfs have high RUWE values and thus potentially erroneous parallaxes.

The binary main sequence identified in Berger et al. (2018b) is not prominent here. This has two reasons: (1) we have corrected photometry for stars with secondaries between 1 and 4'' (Section 2.4) and (2) lower main-sequence grid models, even out to ages of 20 Gyr, are still not luminous (and hence large) enough to emulate the absolute magnitude of a multiple-star system; thus, such results are not allowed by our analysis.

The striping pattern ranging from 5800 to 7000 K and $\approx 1\text{--}2 R_{\odot}$ is an artifact of our model grid. The brighter colored stripes are stellar isochrones at solar metallicity, where individual stellar solutions preferentially “snap” to the isochrone grid. Increasing the age resolution of our model grid or only using

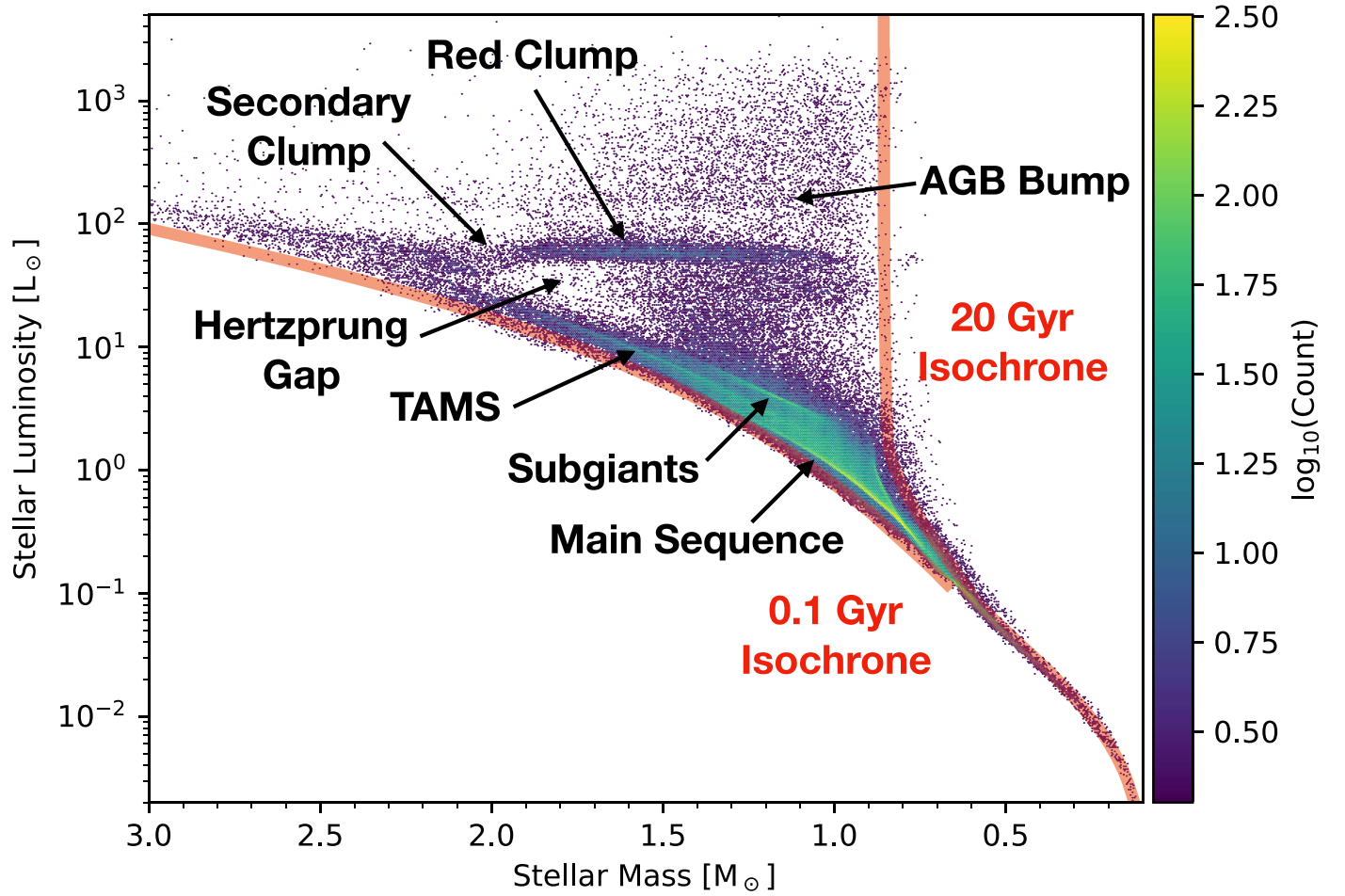


Figure 9. Luminosity vs. mass for $\sim 186,000$ Kepler stars. Color-coding represents logarithmic number density. The red, translucent curves represent the 0.1 (left) and 20 Gyr (right), $[\text{Fe}/\text{H}] = 0.0$ dex isochrones. We have labeled all features in the distribution accordingly.

stars with spectroscopic metallicity constraints would significantly reduce the contrast of the stripes, but computational constraints and the lack of spectroscopic metallicities for 2/3 of the Kepler sample prevent us from doing so here.

4.3. The Grid-modeled Mass–Luminosity Diagram of Kepler Stars

Another benefit of isochrone fitting over the work in Berger et al. (2018b) is the determination of stellar masses. In Figure 9, we plot our grid-modeled luminosities versus our grid-modeled masses. This diagram shows a variety of features resulting from the processes of stellar evolution, similar to the radius- T_{eff} H-R diagram in Figure 8. It is even easier to see the stellar radius evolution in this plot, given the ability to choose a mass on the x -axis and follow the change in density of points as the stellar age (and luminosity) increases.

For masses below $0.8 M_{\odot}$, the main sequence grows thinner due to the lack of luminosity evolution within the age of the universe. We see that for stellar masses between 0.6 and $0.8 M_{\odot}$, there is some scatter around the zero-age main sequence (ZAMS). This is close to where we replaced the MIST model parameters with the M-dwarf empirical relations from Mann et al. (2015, 2019) (Section 2.3). From there, the

luminosities drop off as well as the apparent scatter as the masses approach the hydrogen-burning limit.

At masses $\geq 0.8 M_{\odot}$, we see the distribution expands vertically. In addition, the highest density of points occurs between 0.8 and $1.0 M_{\odot}$, representing the large fraction of solar-type stars within the Kepler sample. For masses $\gtrsim 0.9 M_{\odot}$, luminosities begin to span from the main sequence up the giant branch. The smooth curves tracing the outermost models on both the left and right represent the minimum and maximum age solar-metallicity isochrones that we used in our analysis. The lower-left edge is a 100 Myr isochrone, while the upper-right edge is a 20 Gyr isochrone.

There are a few features that are prominent as a function of mass and luminosity at masses $\gtrsim 1.0 M_{\odot}$. Starting from the bottom of the distribution, we see that there is an over-density of points arcing from the yellow, highest densities up and to the left to subsequently higher masses and luminosities. This branch represents some of the youngest stars in our sample, less than half their TAMS. Just above this main-sequence curve is the higher-mass TAMS, the long, arcing over-density of points from $M_{\star} \approx 1.1\text{--}3.0 M_{\odot}$ to $L_{\star} \approx 2.5\text{--}140 L_{\odot}$. Next, we see an over-density of points ranging from $M_{\star} \approx 1\text{--}1.7 M_{\odot}$ and $L_{\star} \approx 3\text{--}20 L_{\odot}$. These stars are all subgiants, where luminosities stay roughly constant at a particular stellar mass as they move toward the red giant branch.

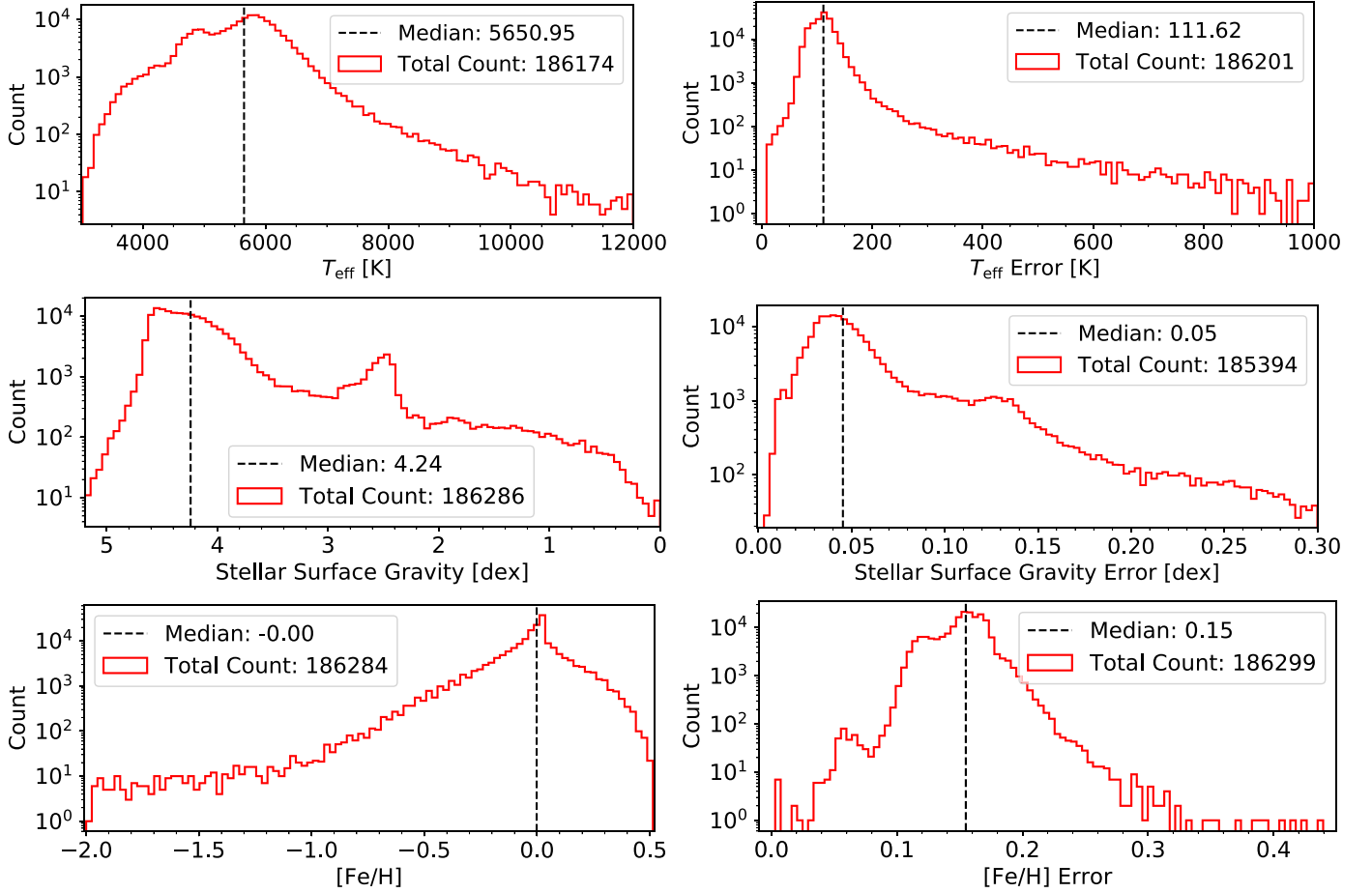


Figure 10. T_{eff} , $\log g$, and metallicity parameter and uncertainty distributions from our catalog. The black dashed vertical lines illustrate the median value for each parameter, the value of which is given in the legend belonging to each plot. In addition, the total number of stars displayed in each histogram is provided in the legend. These numbers vary due to choices in parameter cutoffs, and they are usually smaller than the total number of stars presented here, 186,301. Some histograms have logarithmic scaling on the x and/or y-axes. Top row: stellar effective temperatures and their absolute uncertainties. Middle row: stellar surface gravities and their absolute uncertainties. Bottom row: stellar metallicities and their absolute uncertainties.

As we increase in stellar luminosity, we see a lack of stars at $M_{\star} \approx 1.6\text{--}2.0 M_{\odot}$ and $L_{\star} \approx 15\text{--}50 L_{\odot}$, an illustration of the Hertzsprung Gap. This under-density occurs because these massive stars evolve so quickly during their subgiant and giant phases that they reach the red clump almost instantaneously. The red clump is the swath of points from just above the Hertzsprung Gap ($M_{\star} \approx 2.0 M_{\odot}$ and $L_{\star} \approx 60 L_{\odot}$) directly to the right and lower masses at the same luminosity ($M_{\star} \approx 1.0 M_{\odot}$ and $L_{\star} \approx 60 L_{\odot}$). The secondary clump (for massive stars) is the clustering of points below and to the left of the main clump for $M_{\star} \approx 2.0 M_{\odot}$. Finally, the slight over-density of points at $M_{\star} \approx 1.1\text{--}2.0 M_{\odot}$ and $L_{\star} \approx 140 L_{\odot}$ represents the asymptotic giant branch (AGB) bump, where a few stars have reached the AGB phase. The AGB bump is analogous to the AGB as the RGB bump is to the RGB, where the luminosity decreases after the He-core and H-shell burning stops for a short time, which then causes contraction and subsequent reignition. The star then continues to grow more luminous as it moves up the AGB. This temporary deceleration of evolution on the AGB produces the AGB bump (Gallart 1998).

We cannot glean additional structure from luminosities $\geq 140 L_{\odot}$, but we do see a wide range of masses at these high luminosities as expected. Moreover, we find that the masses of

Kepler stars range between 0.1 and $5.1 M_{\odot}$, and only ~ 400 Kepler stars have masses exceeding $3.0 M_{\odot}$. This is consistent with the inferred lack of very young stars in the Kepler field.

4.4. Parameter and Uncertainty Distributions

4.4.1. Temperatures, Surface Gravities, and Metallicities

In Figure 10, we plot the histograms of T_{eff} , $\log g$, and metallicity for the Kepler parent sample. Unsurprisingly, the T_{eff} histogram peaks close to the solar T_{eff} , and the giants cause the slight bump around 4800 K. The errors in T_{eff} peak strongly around the median of 112 K by design (see Section 2.3 for details), as this represents a fractional error of $\sim 2\%$ for the median star in our sample.

The peak in the stellar surface gravity ($\log g$) histogram occurs at 4.24 dex in cgs units. This is consistent with the larger-than-expected percentage of Kepler subgiant targets (21%, Berger et al. 2018b). Typical errors are on the order of 0.05 dex, which is a dramatic improvement over the ≈ 0.2 dex $\log g$ median error provided in Mathur et al. (2017) due to the strong radius constraints from Gaia DR2 parallaxes.

Because the vast majority of Kepler stars do not have spectroscopic metallicities ($\approx 120,000$ out of $\approx 186,000$), the

metallicity distribution and its uncertainties are not particularly informative. The left plot shows a peak at solar metallicity with a sharp drop-off to either higher or lower metallicities. This is unsurprising, given that our priors are centered on solar metallicity. The uncertainty distribution has one major peak around the median uncertainty of 0.15 dex; this peak represents a convolution of metallicities derived from a Gaussian prior centered at solar metallicity with a ≈ 0.2 dex width and the remaining objects that have spectroscopic metallicities with fixed 0.15 dex uncertainties. Almost all of the ≈ 400 stars with metallicity uncertainties ≤ 0.08 dex have $\text{GOF} \leq 0.99$ and are therefore unreliable. The remaining low-metallicity uncertainty stars have observables which place them in sparse areas of the model grid.

4.4.2. Radii, Masses, Densities, Luminosities, and Ages

Figure 11 contains the remaining important parameters we derived for all Kepler stars. Because many of our output posteriors are asymmetric, we compute our fractional errors by taking the maximum of the upper and lower uncertainties and then divide that maximum by the computed value for that parameter.

The first row of Figure 11 contains stellar masses. We see that the mass distribution peaks at $1 M_{\odot}$, with a median that is slightly super-solar at $1.06 M_{\odot}$. The fractional uncertainty distribution in mass peaks at $\approx 7\%$, which is half that reported in Mathur et al. (2017). The peak of fractional mass errors close to 2% corresponds to the Mann et al. (2019) empirical M_K -mass uncertainties by design (Section 2.3).

The second, third, and fourth rows of Figure 11 display the distributions of stellar radii, mean stellar densities (ρ_*), and luminosities, respectively and their uncertainties. Each histogram peaks near solar values and plateaus toward larger radii and luminosities and smaller densities due to subgiant contamination. The more narrow peaks that occur around $11 R_{\odot}$, $10^{-3} \rho_{\odot}$, and $60 L_{\odot}$ represent the red clump. The uncertainty distributions for stellar radius, density, and luminosity peak at $\approx 3\%$, $\approx 10\%$, and $\approx 8\%$ and have a median of $\approx 4\%$, 13% , and 10% , respectively. Each has a broad tail to larger fractional uncertainties, which is dependent mostly on the precision of the parallax from Gaia DR2. T_{eff} errors are held fixed to $\approx 2\%$ as described in Section 2.3. The 13% median fractional error in ρ_* represents a factor of ≈ 4 improvement over the previous KSPC’s fractional uncertainties (Mathur et al. 2017). These precise ρ_* values will be a critical input for refitting Kepler transits.

The fifth row of Figure 11 contains stellar ages for the entire sample of Kepler stars. The median value of 4.58 Gyr is close to solar. The distribution peaks around 2.5 Gyr and gradually falls off to larger ages. There is a bump at 10 Gyr, half the age of the grid, where most of the M-dwarfs fall. This occurs because H-R diagram constraints are essentially uninformative for M-dwarfs given that even the most massive M-dwarfs have main-sequence lifetimes ≥ 50 Gyr, over twice the maximum age of our grid. Encouragingly, the distribution also qualitatively matches the red giant asteroseismology-derived age distributions in Silva Aguirre et al. (2018) and Pinsonneault et al. (2018), as well as the rotation-based ages in Claytor et al. (2020) and the Galactic Archaeology with HERMES-Gaia ages in Buder et al. (2019).

The right histogram displaying the fractional age errors has a median fractional age uncertainty of 56%. The peaks in the histogram represent various areas of parameter space where maximum fractional age errors are common. The first peak occurring at slightly less than 0.25 fractional age errors is one that corresponds to ≈ 0.9 – $1.3 M_{\odot}$ subgiants and ≈ 1.4 – $2.0 M_{\odot}$ TAMS Kepler stars. The second, largest peak at ≈ 0.35 corresponds to the age uncertainties of (1) highest-mass stars ($\gtrsim 1.3 M_{\odot}$) on the main sequence, (2) intermediate-mass stars (≈ 1.3 – $1.7 M_{\odot}$) at the TAMS, and (3) low-mass stars (≈ 0.7 – $1.3 M_{\odot}$) on the subgiant branch, TAMS, and on the upper edge of the grid. Highlighting these stars in the H-R diagram outlines the main-sequence turn-off “hook.” The third peak, which occurs just below 0.5 includes (1) high-mass stars ($\gtrsim 1.2 M_{\odot}$) on the main sequence and (2) low-mass stars (≈ 0.7 – $1.4 M_{\odot}$) at the TAMS and at the maximum ages within our grid. Finally, the fourth peak occurs at 0.65 fractional age errors because of the M-dwarfs and their uninformative ages. M-dwarfs do not evolve at all in 20 Gyr, and hence have flat age posteriors with medians at ≈ 10 Gyr and 1σ uncertainties between 6 and 7 Gyr. From there, the distribution smoothly decreases until 1.5 fractional age errors. Solar-type stars at the ZAMS do not produce fractional age errors larger than 1.5 due to grid edge effects and the typical observational uncertainties, resulting in the sudden dip in the age distribution. Following this dip, the number of stars with larger and larger fractional age errors declines gradually.

4.5. T_{eff} Comparison to the DR25 KSPC

Figure 12 shows a comparison of stellar T_{eff} in the DR25 stellar properties catalog (Mathur et al. 2017) to those derived in this paper. The distribution approximately tracks the 1:1 line. Our T_{eff} are offset by -3% for giant stars, $+5\%$ for M-dwarfs, $\leq 1\%$ for late-K-dwarfs, -3% for early K-dwarfs, -3% for solar-type stars, $+1\%$ for F-dwarfs, and $+10\%$ for A-stars. For stars with $T_{\text{eff}} \geq 10,000$ K, our T_{eff} are $\geq 10\%$ larger. The majority of Mathur et al. (2017) T_{eff} come from Pinsonneault et al. (2012), which used KIC extinction values based on a simple extinction model (Brown et al. 2011). These extinction values were later shown to be overestimated (Rodrigues et al. 2014). In addition, the extinction used by Pinsonneault et al. (2012) does not account for each star’s distance, where stars that are farther away will experience more extinction.

Due to reddening’s increasing effect over longer distances, the Green et al. (2019) reddening map can account for the major differences that we see from the most distant, hottest stars down to the closer F-dwarfs. The solar-type stars and early K-dwarfs experience slightly less extinction than predicted (Rodrigues et al. 2014). Similarly, Huber et al. (2017) found that after accounting for the underestimated metallicity and overestimated extinction values used by Pinsonneault et al. (2012), the T_{eff} scales should be cooler by -20 to -65 K. This cooler T_{eff} scale brings us closer to the spectroscopic T_{eff} (see Figure 13 in Pinsonneault et al. 2012). However, the M-dwarfs are still too hot given their 2% fractional errors while the giants are too cool. Much of this is likely due to the systematic issues displayed in Figure 3, where the M-dwarfs are too hot and the giants exhibit a strong trend, likely created by systematic errors in color transformations.

The T_{eff} gap at ≈ 4200 K is visible in the Mathur et al. (2017) data. In addition, we observe banding structures which are

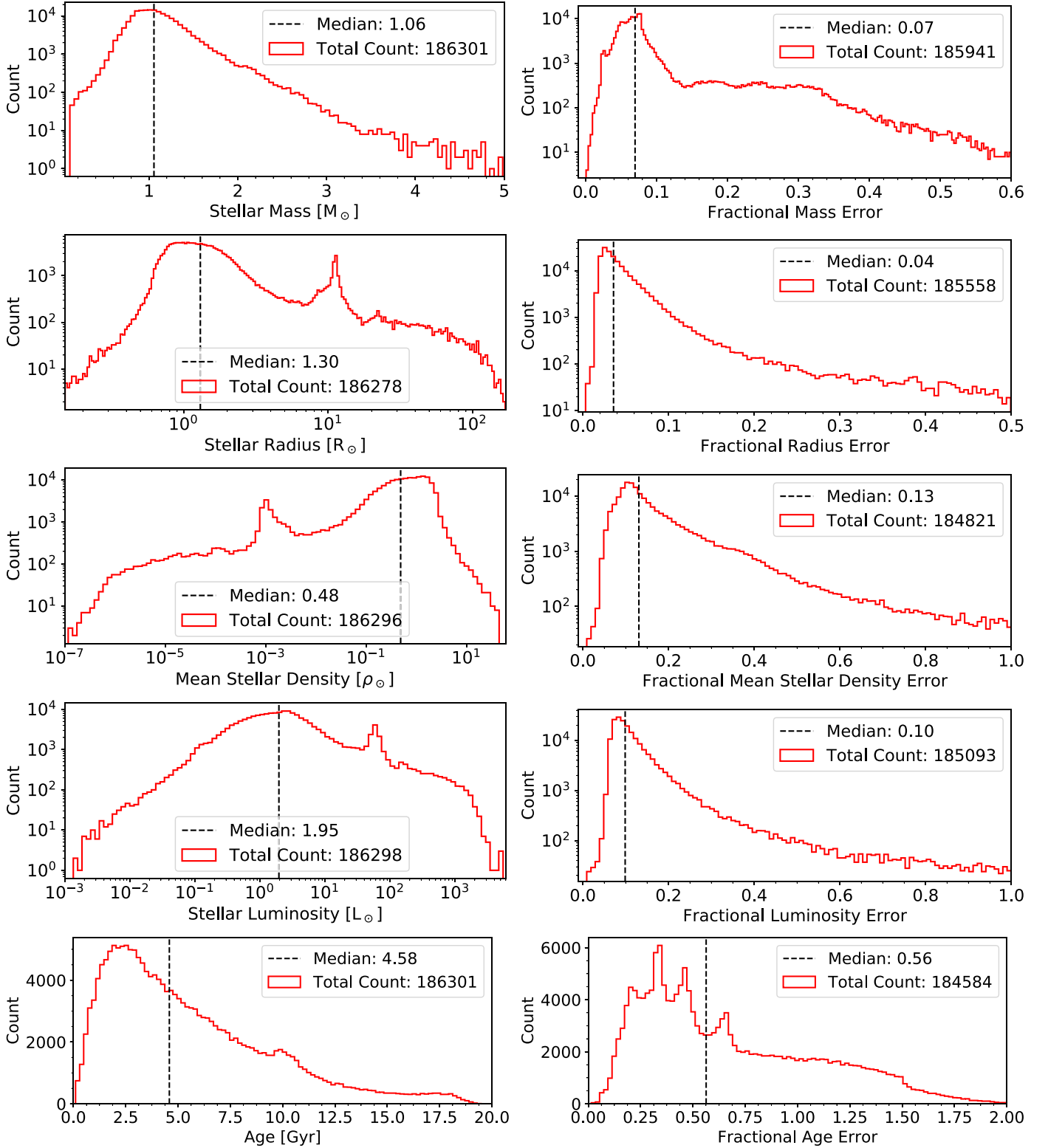


Figure 11. Stellar parameter and uncertainty distributions from our catalog. The black dashed vertical lines illustrate the median value for each parameter, the value of which is given in the legend belonging to each plot. In addition, the total number of stars in displayed in each histogram is provided in the legend. These numbers vary due to choices in parameter cutoffs, and they are usually smaller than the total number of stars presented here, 186,301. Some histograms have logarithmic scaling on the x - and/or y -axes.

visible as horizontal lines in the top plot and diagonal lines in the bottom residuals plot. This structure is an artifact in the input effective temperatures in Mathur et al. (2017). It is

unclear where exactly this banding comes from, but it appears to be dependent on a set of models, as the peaks are evenly spaced every ≈ 100 K.

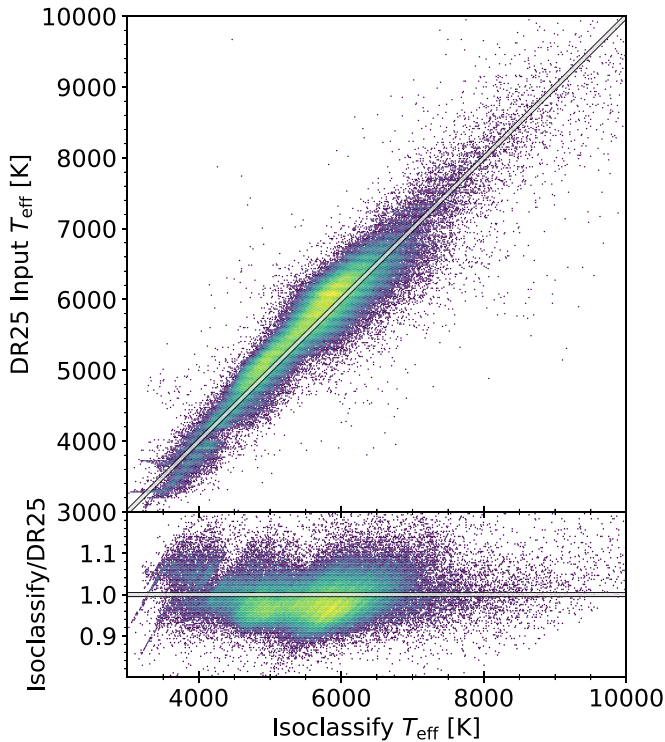


Figure 12. Comparison of T_{eff} of the inputs to the DR25 Kepler Stellar Properties Catalog (Mathur et al. 2017) and the T_{eff} derived in this paper. The colors represent the logarithmic density of points. The white and black line is the 1:1 comparison between DR25 T_{eff} and our derived T_{eff} . The bottom panel shows the ratio between DR25 stellar T_{eff} and our stellar T_{eff} .

5. Guidelines for Catalog Use

Our catalog includes multi-parameter solutions to 186,301 Kepler stars. In this section, we provide important guidelines, caveats, and limitations for the reader to implement/consider when utilizing this catalog.

1. Our input catalog contains 186,548 stars, while our output catalog contains 186,301 stars. This is because the input parameters of 247 stars were too far removed from our grid of MIST models. Another 1543 stars have GOF parameters less than our threshold ($\text{GOF} \leq 0.99$ in Table 2). These stars should be used with caution.
2. We only use a single model grid. Thus, differences due to input physics in model grids are not captured in the reported uncertainties.
3. Our output metallicities ($[\text{Fe}/\text{H}]$ in Table 2) for the 120,000 stars constrained by the Kepler field’s 0.2 dex dispersion solar metallicity prior should be treated with caution. The remaining 66,000 stars have spectroscopic metallicity constraints with 0.15 dex uncertainty. Both sets of stars have large metallicity uncertainties.
4. We do not treat (likely) binaries differently in our isochrone fitting analysis. We amend the 2MASS K_s -band magnitudes where possible, but do not modify the input observables or output stellar parameters of stars with large RUWE ($\text{RUWE} \geq 1.2$ in Table 1). These large RUWE stars and other likely binaries should be removed or treated with caution.
5. We systematically overestimate M-dwarf T_{eff} by $\approx 2\%$, while our giant T_{eff} exhibit a strong trend compared to interferometric determinations. This will affect our

estimates of masses and radii for both M-dwarfs and giants, and giant masses, in particular, are extremely sensitive to metallicity and T_{eff} , both of which are not constrained well in our catalog. E. Gaidos et al. (2020, in preparation) and APOKASC catalogs (Serenelli et al. 2017; Pinsonneault et al. 2018) will provide better and more reliable parameters for Kepler M-dwarfs and giants, respectively.

6. We flag stellar ages which we deem unreliable ($\text{GOF} \leq 0.99$ in Table 2) or uninformative ($\text{TAMS} \geq 20$ Gyr in Table 2) with asterisks, resulting in 14% of catalog stars with suspect ages. We still provide the median and 1σ confidence intervals for posterity, but these are stars whose ages cannot be constrained by our analysis. Due to degeneracies between stellar age and stellar metallicity, our most reliable stellar ages are dwarfs with spectroscopic metallicities ($[\text{Fe}/\text{H}]$ constrained in Table 1).
7. We also caution against the use of our giant ages, given their strong dependence on stellar mass, which is strongly dependent on T_{eff} and metallicity. For more reliable ages for many of the giants included in this catalog, see Serenelli et al. (2017) and Pinsonneault et al. (2018).

6. Stellar Parameter Comparisons for Noteworthy Kepler Systems

To demonstrate the effectiveness of our catalog, we take a look at a few Kepler systems that had stellar radius and mass estimates that were in tension before Gaia DR2. Figure 13 plots the stellar radius and mass measurements from a variety of sources for these systems. A more thorough investigation of updated planet radii of the Kepler sample will be presented in a companion paper (T. A. Berger et al. 2020, in preparation).

6.1. Kepler-11

Kepler-11 hosts six low-density planets and was one of the first multiplanet systems discovered by Kepler (Lissauer et al. 2011). The host star was analyzed most recently in Bedell et al. (2017), where it was classified as a solar twin. Lissauer et al. (2013) also investigated the host’s stellar properties by using transit timing variations (TTVs) to determine the star’s density. Bedell et al. (2017) performed two analyses on Kepler-11: (1) a spectroscopic determination of stellar parameters, and (2) a photodynamical light curve analysis. The spectroscopic analysis led to an estimation of $R_* \approx 1.02 \pm 0.03 R_\odot$ and $M_* \approx 1.04 \pm 0.01 M_\odot$ based on both Yonsei-Yale and Dartmouth isochrones. Keeping the stellar mass fixed at $1.04 M_\odot$, the photodynamical analysis yielded a stellar radius of $1.07^{+0.04}_{-0.01} R_\odot$. Bedell et al. (2017) then computed the mean stellar density for each of the methods, finding that they were at tension. The photodynamical analysis produced a mean stellar density in agreement with Lissauer et al. (2013), while the spectroscopic analysis did not agree.

Our results (purple, center) appear to be in better agreement with the photodynamical analysis of the Kepler-11 lightcurve, and hence also with the prediction of the stellar density computed in Lissauer et al. (2013). While our reported stellar mass is greater than the mass derived in both the lightcurve and spectroscopic analysis in Bedell et al. (2017), the 1σ uncertainty includes the Bedell et al. $1.04 M_\odot$ estimates. Gaia DR2 parallaxes provide the strongest constraints on the stellar

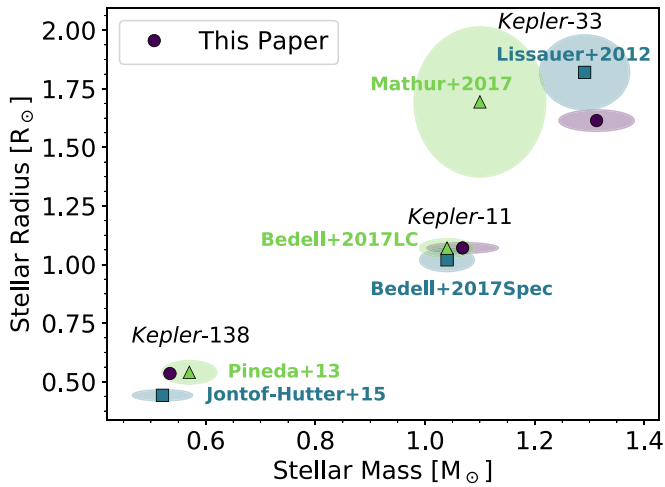


Figure 13. Stellar mass and radius comparisons for three particular Kepler systems with stellar parameters at tension in the past. Points are colored and marked according to the source of their information, and the individual systems are labeled accordingly. Plum points and the purple 1σ error ellipses are values and uncertainties determined from the analysis discussed above, while teal squares and green triangles and their respective error ellipses are taken from the literature. Kepler-11 was investigated in Bedell et al. (2017), Kepler-33 in Lissauer et al. (2012) and Mathur et al. (2017), and Kepler-138 in Pineda et al. (2013) and Jontof-Hutter et al. (2015).

radius, which are in good agreement with the photodynamical and TTV analyses of the Kepler-11 lightcurve, providing further evidence for the larger radius predicted by these methods.

6.2. Kepler-33

Kepler-33, investigated in Lissauer et al. (2012), has five planetary companions, all between $1.5\text{--}6.0 R_{\oplus}$. In Lissauer et al. (2012), both the mass and age posteriors are bimodal due to the star’s location on the H-R diagram, close to the main-sequence turn-off and subgiant branch. Lissauer et al. (2012) reports a mass of $1.29 \pm 0.08 M_{\odot}$ and a radius of $1.82 \pm 0.16 R_{\odot}$ based on the Yonsei-Yale (Yi et al. 2001, 2003; Kim et al. 2002; Demarque et al. 2004) isochrone placement of the spectroscopically derived parameters (teal square). However, while Mathur et al. (2017) used the spectroscopic parameters of Lissauer et al. (2012), they used Dartmouth Stellar Evolution Database (Dotter et al. 2008) models to derive $R_{*}=1.69 \pm 0.32 R_{\odot}$ and $M_{*}=1.10 \pm 0.12 M_{\odot}$ (green triangle). Therefore, the stellar mass tension is result of different stellar model grids.

Our result is more than 1σ discrepant from the Mathur et al. (2017) mass and the Lissauer et al. (2012) radius. Because the stellar masses derived by Lissauer et al. and Mathur et al. are model grid-dependent, we focus on the discrepancy in stellar radius. The stellar radius in Lissauer et al. is mostly constrained by the spectroscopic estimate of $\log g$. However, spectroscopic $\log g$ values are often degenerate with T_{eff} and $[\text{Fe}/\text{H}]$ (Torres et al. 2012). Alternatively, parallaxes, which constrain distances and hence radii, provide a more straightforward approach to determining stellar radii. Hence, we are confident in our parallax-constrained properties of Kepler-33.

6.3. Kepler-138

Investigated in both Pineda et al. (2013) and Jontof-Hutter et al. (2015), Kepler-138 has three planetary companions, one of which was determined to be a Mars-sized planet by

Jontof-Hutter et al. The two results are at tension, largely because different methods were used to determine the parameters. Pineda et al. utilized an empirical relation that determined absolute K_s magnitudes from the equivalent widths of molecular lines (TiO, VO) and photometric colors. These absolute magnitudes were then converted to masses using an absolute magnitude–mass relation from Delfosse et al. (2000) and then to masses using the mass–radius relation of Boyajian et al. (2012b). Alternatively, Jontof-Hutter et al. computed the stellar parameters of Kepler-138 by fitting the light curve constraint of ρ_{*} as well as the spectroscopic constraints of T_{eff} and $[\text{Fe}/\text{H}]$ of Muirhead et al. (2012) to Dartmouth Stellar Evolution models.

According to our analysis, Kepler-138 has a mass of $0.53 M_{\odot}$ and a radius of $0.54 R_{\odot}$. Because this particular star is an M-dwarf, we caution that our mass is suspect given the systematically overestimated T_{eff} of M-dwarfs demonstrated in Figure 3 above. For instance, if we corrected Kepler-138’s T_{eff} by the 70 K systematic offset seen for interferometric stars, we would compute $\approx 0.5 M_{\odot}$ and $\approx 0.5 R_{\odot}$ for the star’s mass and radius, which would agree with neither Pineda et al. (2013) nor Jontof-Hutter et al. (2015). We note that both the radius and the mass error bars are smaller than the marker size.

Our radius and mass estimates for Kepler-138 agree with the reported masses and radii in Pineda et al. (2013), while they disagree with those reported in Jontof-Hutter et al. (2015). Although Jontof-Hutter et al. cite a few potential inaccuracies of the Pineda et al. results, it appears that our solution breaks the tension in favor of the Pineda et al. result. Possible reasons could be inaccuracies in the light curve modeling or photodynamical modeling for the determination of ρ_{*} . Any inaccuracies in the mean stellar density will scale as R_{*}^{-3} while only linearly in mass, which explains the large discrepancy in radius as compared to the one in mass. However, due to the systematic offset in M-dwarf T_{eff} seen in Figure 3, we caution against drawing any strong conclusions for the properties of Kepler-138.

7. Summary and Conclusions

We presented a re-classification of stellar parameters— T_{eff} , masses, radii, luminosities, densities, surface gravities, ages, and metallicities for 186,301 stars observed by the Kepler mission by combining Gaia DR2 parallaxes and spectroscopic metallicities with calibrated KIC (Brown et al. 2011) and KIS g (Greiss et al. 2012) and visual-binary-decontaminated 2MASS K_s (Skrutskie et al. 2006) photometry. We utilized a custom-interpolated set of MIST models (Choi et al. 2016) and *isoclassify* (Huber et al. 2017) to derive stellar parameters. Our main results are as follows.

1. We determine parameters for 186,301 stars. The median (fractional) precisions of our T_{eff} , $\log g$, radii, masses, mean stellar densities, luminosities, and ages are 112 K, 0.05 dex, 4%, 7%, 13%, 10%, and 56%, respectively.
2. We provide the first Kepler Stellar Properties Catalog (KSPC) with a homogeneous T_{eff} scale. M-dwarf T_{eff} may be 75 K hotter than similar stars with interferometry, and our FGK-dwarf T_{eff} are cooler than those of Mathur et al. (2017) by ≈ 110 K.
3. We derive a median Kepler target age of ≈ 4.6 Gyr. Our ages are in good agreement with cluster and asteroseismic ages, where we find that our median age is 3% larger than

the asteroseismic estimate with a scatter of 29%. We caution that 14% of our ages are uninformative, due to the constraints of isochrone fitting for low-mass stars. Our ages are most reliable for both the most massive stars and those on the subgiant branch.

4. We provide the first KSPC that attempts to account for the binarity of all Kepler stars when performing isochrone fits to absolute K_s -band photometry. In addition, we find that at least 12% of the Kepler sample is affected by binary companions. When holding age and metallicity constant, we observe that age biases of companions are functions of both primary and secondary mass; we find binary companions will bias the ages of lower-main-sequence stars by as much as 10 Gyr (where age uncertainties are $\gtrsim 6$ Gyr), and higher-mass stars by a few Gyr (where age uncertainties are $\lesssim 1$ Gyr).
5. We derive accurate and precise stellar masses and radii for three Kepler systems with tension in their reported parameters based on previous analyses. Our results typically break the tension and favor one result over another, although we suggest the reader carefully consider the methods used in each analysis before drawing any strong conclusions.

All of the homogeneous parameters reported here will prove useful for future Kepler exoplanet occurrence rate computations, as homogeneous treatment for all stars ensures that both the host star and field star parameters are considered similarly. For instance, Bryson et al. (2020) utilized the parameters presented in this work to investigate the DR25 catalog's reliability and completeness.

In addition, the masses and ages presented here provide important constraints for Kepler exoplanet host stars. The masses determined here will allow us to constrain the stellar mass dependence of the planet radius gap (Fulton & Petigura 2018; Gupta & Schlichting 2020; Wu 2019). Stellar ages also have interesting implications for exoplanets. Previous analyses have hinted at age-dependent effects on the radii of small exoplanets, particularly those at or near the gap, for subsamples of the Kepler exoplanets (Fulton et al. 2017; Mann et al. 2017; Berger et al. 2018a). With our stellar mass and age constraints for the entire Kepler exoplanet host sample, we will investigate stellar mass and age-dependent exoplanet trends in our companion paper (T. A. Berger et al. 2020, in preparation). Ultimately, we look forward to future investigations which will discover both new features about the Kepler sample and confirm previous results, continuing the legacy of the Kepler telescope well beyond its final observation.

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Software: astropy (Astropy Collaboration et al. 2013), dustmaps (Green et al. 2018), GNU Parallel (Tange 2018), isoclassify (Huber et al. 2017), Matplotlib (Hunter 2007), mw dust (Bovy et al. 2016), Pandas (McKinney 2010), SciPy (Virtanen et al. 2020).

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References

- Abolfathi, B., Aguado, D. S., Aguilar, G., et al. 2018, *ApJS*, **235**, 42
 Arenou, F., Luri, X., Babusiaux, C., et al. 2018, *A&A*, **616**, A17
 Astropy Collaboration, Robitaille, T. P., Tollerud, E. J., et al. 2013, *A&A*, **558**, A33
 Bastien, F. A., Stassun, K. G., Basri, G., & Pepper, J. 2016, *ApJ*, **818**, 43
 Batalha, N. M., Borucki, W. J., Koch, D. G., et al. 2010, *ApJL*, **713**, L109

- Bedell, M., Bean, J. L., Meléndez, J., et al. 2017, *ApJ*, **839**, 94
- Berger, T. A., Howard, A. W., & Boesgaard, A. M. 2018a, *ApJ*, **855**, 115
- Berger, T. A., Huber, D., Gaidos, E., & van Saders, J. L. 2018b, *ApJ*, **866**, 99
- Bessell, M. S. 2000, *PASP*, **112**, 961
- Bovy, J., Rix, H.-W., Green, G. M., Schlafly, E. F., & Finkbeiner, D. P. 2016, *ApJ*, **818**, 130
- Boyajian, T. S., McAlister, H. A., van Belle, G., et al. 2012a, *ApJ*, **746**, 101
- Boyajian, T. S., von Braun, K., van Belle, G., et al. 2012b, *ApJ*, **757**, 112
- Boyajian, T. S., von Braun, K., van Belle, G., et al. 2013, *ApJ*, **771**, 40
- Brown, T. M., Latham, D. W., Everett, M. E., & Esquerdo, G. A. 2011, *AJ*, **142**, 112
- Bryson, S., Coughlin, J., Batalha, N. M., et al. 2020, *AJ*, **159**, 279
- Buder, S., Lind, K., Ness, M. K., et al. 2019, *A&A*, **624**, A19
- Burke, C. J., Christiansen, J. L., Mullally, F., et al. 2015, *ApJ*, **809**, 8
- Choi, J., Dotter, A., Conroy, C., et al. 2016, *ApJ*, **823**, 102
- Clayton, Z. R., van Saders, J. L., Santos, A. R. G., et al. 2020, *ApJ*, **888**, 43
- Delfosse, X., Forveille, T., Ségransan, D., et al. 2000, *A&A*, **364**, 217
- Demarque, P., Woo, J.-H., Kim, Y.-C., & Yi, S. K. 2004, *ApJS*, **155**, 667
- Dong, S., Zheng, Z., Zhu, Z., et al. 2014, *ApJL*, **789**, L3
- Dotter, A. 2016, *ApJS*, **222**, 8
- Dotter, A., Chaboyer, B., Jevremović, D., et al. 2008, *ApJS*, **178**, 89
- Evans, D. F. 2018, *RNAAS*, **2**, 20
- Feiden, G. A., & Chaboyer, B. 2012, *ApJ*, **757**, 42
- Fulton, B. J., & Petigura, E. A. 2018, *AJ*, **156**, 264
- Fulton, B. J., Petigura, E. A., Howard, A. W., et al. 2017, *AJ*, **154**, 109
- Furlan, E., Ciardi, D. R., Cochran, W. D., et al. 2018, *ApJ*, **861**, 149
- Furlan, E., Ciardi, D. R., Everett, M. E., et al. 2017, *AJ*, **153**, 71
- Gaia Collaboration, Brown, A. G. A., Vallenari, A., et al. 2018, *A&A*, **616**, A1
- Gaia Collaboration, Prusti, T., de Bruijne, J. H. J., et al. 2016, *A&A*, **595**, A1
- Gallart, C. 1998, *ApJL*, **495**, L43
- González-Solares, E. A., Irwin, M., McMahon, R. G., et al. 2011, *MNRAS*, **416**, 927
- Green, G. M., Schlafly, E. F., Finkbeiner, D., et al. 2018, *MNRAS*, **478**, 651
- Green, G. M., Schlafly, E. F., Zucker, C., Speagle, J. S., & Finkbeiner, D. P. 2019, *ApJ*, **887**, 93
- Greiss, S., Steeghs, D., Gänsicke, B. T., et al. 2012, *AJ*, **144**, 24
- Gupta, A., & Schlichting, H. E. 2020, *MNRAS*, **493**, 792
- Hadden, S., & Lithwick, Y. 2016, *ApJ*, **828**, 44
- Henden, A. A., Levine, S., Terrell, D., et al. 2018, AAS Meeting, **232**, 223
- Hon, M., Stello, D., & Yu, J. 2018, *MNRAS*, **476**, 3233
- Howard, A. W., Marcy, G. W., Bryson, S. T., et al. 2012, *ApJS*, **201**, 15
- Howes, L. M., Lindegren, L., Feltzing, S., Church, R. P., & Bensby, T. 2019, *A&A*, **622**, A27
- Huang, Y., Liu, X.-W., Yuan, H.-B., et al. 2015, *MNRAS*, **454**, 2863
- Huber, D., Bryson, S. T., Haas, M. R., et al. 2016, *ApJS*, **224**, 2
- Huber, D., Silva Aguirre, V., Matthews, J. M., et al. 2014, *ApJS*, **211**, 2
- Huber, D., Zinn, J., Bojesen-Hansen, M., et al. 2017, *ApJ*, **844**, 102
- Hunter, J. D. 2007, *CSE*, **9**, 90
- Jester, S., Schneider, D. P., Richards, G. T., et al. 2005, *AJ*, **130**, 873
- Johnson, D. R. H., & Soderblom, D. R. 1987, *AJ*, **93**, 864
- Jontof-Hutter, D., Rowe, J. F., Lissauer, J. J., Fabrycky, D. C., & Ford, E. B. 2015, *Natur*, **522**, 321
- Kim, Y.-C., Demarque, P., Yi, S. K., & Alexander, D. R. 2002, *ApJS*, **143**, 499
- Kraus, A. L., & Hillenbrand, L. A. 2007, *ApJ*, **662**, 413
- Kraus, A. L., Ireland, M. J., Huber, D., Mann, A. W., & Dupuy, T. J. 2016, *AJ*, **152**, 8
- Lindegren, L. 2018, Re-normalising the astrometric chi-square in Gaia DR2, http://www.rssd.esa.int/doc_fetch.php?id=3757412
- Lindegren, L., Hernandez, J., Bombrun, A., et al. 2018, *A&A*, **616**, A2
- Lissauer, J. J., Fabrycky, D. C., Ford, E. B., et al. 2011, *Natur*, **470**, 53
- Lissauer, J. J., Jontof-Hutter, D., Rowe, J. F., et al. 2013, *ApJ*, **770**, 131
- Lissauer, J. J., Marcy, G. W., Rowe, J. F., et al. 2012, *ApJ*, **750**, 112
- Luo, A.-L., Zhao, Y.-H., Zhao, G., et al. 2015, *RAA*, **15**, 1095
- Majewski, S. R., Schiavon, R. P., Frinchaboy, P. M., et al. 2017, *AJ*, **154**, 94
- Mann, A. W., Dupuy, T., Kraus, A. L., et al. 2019, *ApJ*, **871**, 63
- Mann, A. W., Feiden, G. A., Gaidos, E., Boyajian, T., & von Braun, K. 2015, *ApJ*, **804**, 64
- Mann, A. W., Gaidos, E., Vanderburg, A., et al. 2017, *AJ*, **153**, 64
- Mathur, S., Huber, D., Batalha, N. M., et al. 2017, *ApJS*, **229**, 30
- McKinney, W. 2010, in Proc. 9th Python in Science Conf., ed. S. van der Walt & J. Millman (Austin, TX: SciPy), 51
- Meibom, S., Barnes, S. A., Latham, D. W., et al. 2011, *ApJL*, **733**, L9
- Molenda-Žakowicz, J., Brogaard, K., Niemczura, E., et al. 2014, *MNRAS*, **445**, 2446
- Muirhead, P. S., Johnson, J. A., Apps, K., et al. 2012, *ApJ*, **747**, 144
- Newton, E. R., Irwin, J., Charbonneau, D., et al. 2016, *ApJ*, **821**, 93
- Paxton, B., Bildsten, L., Dotter, A., et al. 2011, *ApJS*, **192**, 3
- Paxton, B., Cantiello, M., Arras, P., et al. 2013, *ApJS*, **208**, 4
- Paxton, B., Marchant, P., Schwab, J., et al. 2015, *ApJS*, **220**, 15
- Perryman, M. A. C., Lindegren, L., Kovalevsky, J., et al. 1997, *A&A*, **500**, 501
- Petigura, E. A., Howard, A. W., Marcy, G. W., et al. 2017, *AJ*, **154**, 107
- Petigura, E. A., Marcy, G. W., & Howard, A. W. 2013, *ApJ*, **770**, 69
- Pineda, J. S., Bottom, M., & Johnson, J. A. 2013, *ApJ*, **767**, 28
- Pinsonneault, M. H., An, D., Molenda-Žakowicz, J., et al. 2012, *ApJS*, **199**, 30
- Pinsonneault, M. H., Elsworth, Y. P., Tayar, J., et al. 2018, *ApJS*, **239**, 32
- Ren, J. J., Rebassa-Mansergas, A., Parsons, S. G., et al. 2018, *MNRAS*, **477**, 4641
- Rizzuto, A. C., Vanderburg, A., Mann, A. W., et al. 2018, *AJ*, **156**, 195
- Rodrigues, T. S., Girardi, L., Miglio, A., et al. 2014, *MNRAS*, **445**, 2758
- Schönrich, R., Binney, J., & Dehnen, W. 2010, *MNRAS*, **403**, 1829
- Scott, D. W. 1992, *Multivariate Density Estimation* (2nd ed.; Hoboken, NJ: John Wiley and Sons)
- Serenelli, A., Johnson, J., Huber, D., et al. 2017, *ApJS*, **233**, 23
- Silva Aguirre, V., Bojesen-Hansen, M., Slumstrup, D., et al. 2018, *MNRAS*, **475**, 5487
- Silva Aguirre, V., Lund, M. N., Antia, H. M., et al. 2017, *ApJ*, **835**, 173
- Skrutskie, M. F., Cutri, R. M., Stiening, R., et al. 2006, *AJ*, **131**, 1163
- Soderblom, D. R. 2010, *ARA&A*, **48**, 581
- Tange, O. 2018, GNU Parallel 2018, Zenodo, doi:10.5281/zenodo.1146014
- Torres, G., Fischer, D. A., Sozzetti, A., et al. 2012, *ApJ*, **757**, 161
- Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, *Nat. Methods*, **17**, 261
- Vrard, M., Mosser, B., & Samadi, R. 2016, *A&A*, **588**, A87
- Wu, Y. 2019, *ApJ*, **874**, 91
- Yi, S., Demarque, P., Kim, Y.-C., et al. 2001, *ApJS*, **136**, 417
- Yi, S. K., Kim, Y.-C., & Demarque, P. 2003, *ApJS*, **144**, 259
- Zhao, G., Zhao, Y.-H., Chu, Y.-Q., Jing, Y.-P., & Deng, L.-C. 2012, *RAA*, **12**, 723
- Ziegler, C., Law, N. M., Baranec, C., et al. 2018, *AJ*, **155**, 161