

Extracting the spin polarizabilities of the proton by measurement of Compton double-polarization observables

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The Compton double-polarization observable Σ_{2z} has been measured for the first time in the $\Delta(1232)$ resonance region using a circularly polarized photon beam incident on a longitudinally polarized proton target at the Mainz Microtron. This paper reports these results, together with the model-dependent extraction of four proton spin polarizabilities from fits to additional asymmetry data using either a dispersion relation calculation or a baryon chiral perturbation theory calculations with the weighted average of these two fits resulting in $\gamma_{E1E1} = -2.87 \pm 0.52$, $\gamma_{M1M1} = 2.70 \pm 0.43$, $\gamma_{E1M2} = -0.85 \pm 0.72$, and $\gamma_{M1E2} = 2.04 \pm 0.43$, in units of 10^{-4} fm^4 .

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I. INTRODUCTION

The electromagnetic interaction of a photon with a nucleon can be studied through Compton scattering experiments. It is best described using an effective Hamiltonian expanded in terms of the incident photon energy. Structure observables of these composite systems are experimentally accessible by elastically scattering real photons from the nucleon in real Compton scattering (RCS). Over decades, RCS has been established as a benchmark for understanding the ground-state properties of the nucleon, such as the magnetic moment. However, the leading-order properties that are sensitive to the internal quark dynamics of the nucleon are still poorly understood experimentally. This paper uses RCS in the $\Delta(1232)$ resonance region as a probe to understand some internal structure observables of a nucleon, the nucleon polarizabilities. These are fundamental properties that describe how its internal structure deforms under an applied electromagnetic field [1,2].

The electromagnetic field of the photon induces transitions of certain definite multipolarities while attempting to deform the nucleon. The effective Hamiltonian at second order in incident photon energy, E_γ , depends on the electric and magnetic scalar polarizabilities, α_{E1} and β_{M1} , and at third order depends on the spin polarizabilities (SPs).

The third-order effective Hamiltonian term in the spin-dependent interaction is

$$H_{\text{eff}}^{(3)} = -4\pi \left[\frac{1}{2} \gamma_{E1E1} \vec{\sigma} (\vec{E} \times \dot{\vec{E}}) + \frac{1}{2} \gamma_{M1M1} \vec{\sigma} (\vec{H} \times \dot{\vec{H}}) - \gamma_{M1E2} E_{ij} \sigma_i H_j + \gamma_{E1M2} H_{ij} \sigma_i E_j \right], \quad (1)$$

where $\dot{\vec{E}}$, $\dot{\vec{H}}$, E_{ij} , and H_{ij} are the partial derivatives with respect to time and space defined as $\dot{\vec{E}} = \partial_t \vec{E}$, $\dot{\vec{H}} = \partial_t \vec{H}$, $E_{ij} = \frac{1}{2}(\partial_i E_j + \partial_j E_i)$, and $H_{ij} = \frac{1}{2}(\partial_i H_j + \partial_j H_i)$, and γ_{E1E1} , γ_{M1M1} , γ_{M1E2} , and γ_{E1M2} are the four SPs. The physics behind these leading-order SPs involves the excitation of the spin- $\frac{1}{2}$ target nucleon to some intermediate state (Δ or N^*) via an electric or magnetic ($E1$ or $M1$) dipole transition and a successive de-excitation back to a spin- $\frac{1}{2}$ nucleon final state via an electric or magnetic dipole ($E1$ or $M1$) or quadrupole ($E2$ or $M2$) transition. These internal structure constants are manifestations of the spin structure of the nucleon, which parametrize the “stiffness” of the nucleon’s spin against the electromagnetically induced deformations relative to the spin axis.

Measurements of two linear combinations of these four SPs—the forward spin polarizability, γ_0 [3,4], and the backward spin polarizability, γ_π [5]—have been reported for the proton by several experiments. An extraction of the individual proton SPs was recently published via measurement of the double-polarization Compton asymmetry— Σ_{2z} —using a transversely polarized proton target at the Mainz Microtron (MAMI) [6] in conjunction with the γ_0 and γ_π results and measurement of the beam-polarization Compton asymmetry Σ_3 performed at the LEGS facility [7]. This paper describes an improvement to the extraction of these proton SPs from the measurement of the double-polarization asymmetry Σ_{2z} using a longitudinally polarized proton target at MAMI. Σ_{2z}

is defined as

$$\Sigma_{2z} = \frac{1}{P_{\text{circ}}^\gamma P_z^t} \left[\frac{(N_{+z}^R + N_{-z}^L) - (N_{+z}^L + N_{-z}^R)}{(N_{+z}^R + N_{-z}^L) + (N_{+z}^L + N_{-z}^R)} \right], \quad (2)$$

where $N_{\pm z}^R$ and $N_{\pm z}^L$ are the normalized yield for right-handed and left-handed helicity states of the beam with the target polarized in the $\pm z$ direction, and P_{circ}^γ and P_z^t are the degrees of the photon beam circular polarization and target polarization, respectively.

II. EXPERIMENTAL SETUP

The experiment was performed in the A2 hall at MAMI [8,9], a facility composed of a cascade of three race track microtrons (RTM) that can provide both unpolarized and longitudinally polarized electron beams with energies up to 1.6 GeV [8]. The longitudinally polarized electron beam was produced by irradiating a strained GaAsP III-V semiconductor with circularly polarized laser light [10]. A 180° polarization flip was provided by reversing the helicity of the laser light with a Pockels cell at a rate of approximately 1 Hz. A standard Mott polarimeter [11], installed near the MAMI accelerator cascade, was used for polarization measurements. The average beam polarization was $86.8 \pm 0.1\%$ [12]. For this measurement, a 450 MeV polarized electron beam passed through an alloy radiator of cobalt and iron, producing circularly polarized Bremsstrahlung photons. The photon polarization, P_γ , was determined by the helicity transfer relationship

$$P_\gamma = P_e \frac{4E_\gamma E_e - E_\gamma^2}{4E_e^2 - 4E_\gamma E_e + 3E_\gamma^2}, \quad (3)$$

where P_e is the electron beam polarization, E_e is the electron beam energy, and E_γ is the energy of the radiated photon. E_γ was determined by detecting the Bremsstrahlung electrons in the tagged photon spectrometer [13], and only photons in the energy range $E_\gamma = 265\text{--}305$ MeV were used for this analysis. The previously mentioned flip of the electron beam polarization direction results in a flip of the photon beam helicity, which, given the relatively fast rate of 1 Hz, provides the Σ_{2z} asymmetry relatively free of systematic effects. The photon beam was passed through a 2.5-mm-diameter lead collimator, resulting in a beam spot size of 9 mm on the longitudinally polarized frozen spin target (FST) [14] located in the center of the Crystal Ball spectrometer (CB) [15].

The FST used dynamic nuclear polarization, and its polarization was measured with a nuclear magnetic resonance coil; both are described in detail in Ref. [14]. Polarization of up to 80% and relaxation times of nearly 1000 h were achieved [13,16], and the direction of proton polarization was reversed approximately once per week. While flipping the photon helicity is enough to produce the Σ_{2z} asymmetry, additionally reversing the target polarization is useful to further remove systematic effects. Polarization measurements were completed at the start and end of each data taking period for different polarization orientations. Corrections to the target polarization due to ice buildup on the NMR coil [17] were determined with π^0 asymmetries as well as comparisons of

unpolarized and polarized total inclusive and π^0 cross sections [18]. To reflect inconsistencies between these methods, a liberal systematic error of 10% for the target polarization was utilized.

Data were collected during two beamtimes in 2014 and 2015 using the nearly 4π CB-TAPS detector system [19]: the CB as a central calorimeter, and TAPS as a forward calorimeter. The CB consists of 672 optically isolated NaI(Tl) crystals with a truncated triangular pyramid shape arranged in two hemispheres. It covers about 94% of 4π steradians and an angular range of $21^\circ \leq \theta \leq 159^\circ$ [7]. TAPS consists of 366 hexagonal BaF₂ crystals and two inner rings totaling 72 PbWO₄ crystals and covers an angular range of $2^\circ \leq \theta \leq 20^\circ$ [20]. Charged particles were identified using energy deposition in the particle identification detector and tracked by a pair of multi-wire proportional chambers or TAPS-veto detectors and their corresponding calorimetric detector. Although the CB-TAPS system covers the angular range of 2 – 159° , there are regions near the entrance and exit through the detectors that are less efficient. These regions are: (i) the forward hole in the TAPS detector, 2 – 6° , and (ii) the backward hole in the CB, 150 – 159° . Fiducial cuts were applied to remove all the data from these angular regions of reduced detection efficiency.

III. DATA ANALYSIS

The Compton scattering channel, $\gamma p \rightarrow \gamma p$, has a simple final state, but it is very important to correctly identify background from competing reactions because its cross section is only about 1% of the cross section for the dominant π^0 photoproduction process. In addition, under certain conditions, π^0 photoproduction can mimic the Compton scattering signature if one of the π^0 decay photons escapes the detector, or if the electromagnetic showers from the two photons overlap due to finite angular resolution. The Compton channel was identified by selecting events having a total deposited energy above 40 MeV, where a single neutral and a single charged track are reconstructed, with the former in coincidence with a hit in the tagger. In order to remove uncorrelated events between CB-TAPS and the photon tagger, the timing differences between the neutral track and hits in the tagger were checked against a 20 ns wide prompt (coincidence) window and a 910 ns wide random window split in two with one on either side of the prompt peak. The random sample was normalized by the relative window widths and subtracted from the prompt timing signal.

To eliminate competing backgrounds from coherent and incoherent Compton scattering and π^0 photoproduction off of nonhydrogen nuclei in the FST from the windows and shells of the cryostat material (mainly $^3\text{He}/^4\text{He}$, ^{12}C , and ^{16}O), separate data were taken by inserting a carbon foam target with density 0.55 g/cm³ into the same cryostat and the normalized yield was subtracted. A base scaling factor was determined by the ratio of live-time corrected tagger scalars for the butanol and carbon data sets. Comparison of π^0 photoproduction simulations with the data showed that a correction to this ratio of approximately 10% was necessary to account for a higher contribution from the helium in the target

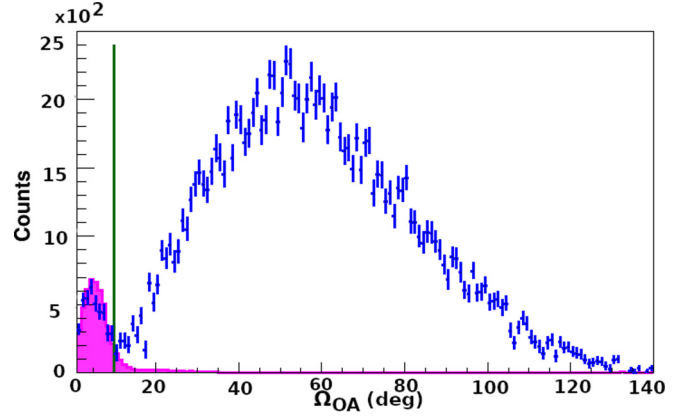


FIG. 1. Opening angle distribution for simulated Compton scattering events (magenta) compared with the carbon-subtracted data (blue) at $E_\gamma = 285$ – 305 MeV and over all Compton angles. A cut of 10° on the opening angle is indicated by the vertical line (green).

to this background. The carbon target density was chosen such that the number of nucleons equals the number of nonhydrogen nucleons from the $^3\text{He}/^4\text{He}$, ^{12}C , and ^{16}O in the target. To remove background from π^0 photoproduction off of the proton, the coincidence of a recoil charged track in addition to the neutral track was required, as mentioned above. However, since protons suffer a significant amount of energy loss when they travel from the event vertex through the target material, a $^3\text{He}/^4\text{He}$ refrigeration bath, various cryostat shells, and a longitudinal holding coil on their way to a detector crystal, the analysis was limited to an incident photon energy range of $E_\gamma = 265$ – 305 MeV. Further details on the background cuts, subtractions, and normalization factors can be found in Ref. [12,18].

To identify events of interest, four-momentum conservation was used to constrain the observed reaction kinematics. As the background varies significantly across both energy and angle, their dependencies were studied. The tagged photon energy bins below $\gamma p \rightarrow \pi^0 \pi^0 p$ threshold were divided into five θ bins, and were analyzed separately. The opening angle (Ω_{OA}), defined as the angle between the detected proton, \vec{p}_{recoil} , and where the proton was expected assuming RCS kinematics, $\vec{p}_{\text{miss}} = \vec{p}_{\gamma_i} - \vec{p}_{\gamma_f}$, $\cos(\Omega_{\text{OA}}) = \frac{\vec{p}_{\text{miss}} \cdot \vec{p}_{\text{recoil}}}{|\vec{p}_{\text{miss}}| \times |\vec{p}_{\text{recoil}}|}$, was used for a two-body reaction selection. The Monte Carlo simulated opening angle results show a sharp peak around 5° , which is in good agreement with the data. The large background, as seen in Fig. 1, is mainly due to the π^0 photoproduction process from the proton. This can be suppressed by applying a 10° opening angle cut, as indicated by the green vertical line. The Compton coplanarity angle, defined as the difference in the azimuthal angles of a scattered photon and a recoil proton, $\Delta\phi = |\phi_{\gamma_f} - \phi_p|$, was used to suppress additional background. A cut on the fixed coplanarity angle, $\Delta\phi = 180 \pm 15^\circ$, as indicated by the two vertical green lines in Fig. 2, was applied to the reconstructed events. For those events with a single neutral and a single charged track, the missing mass is calculated with

$$M_{\text{miss}}^2 = (E_{\gamma_i} + m_p c^2 - E_{\gamma_f})^2 - (\vec{p}_{\gamma_i} - \vec{p}_{\gamma_f})^2 c^2, \quad (4)$$

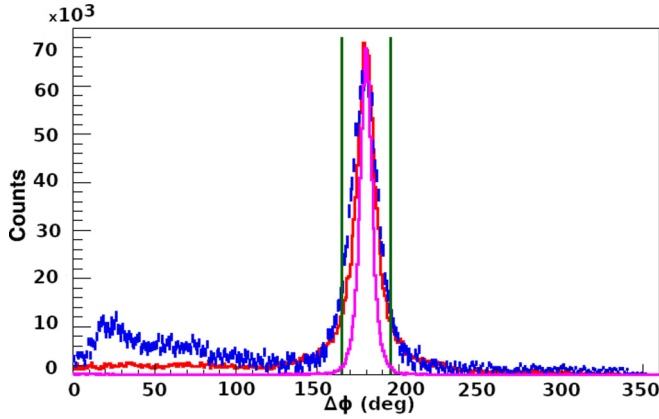


FIG. 2. Coplanarity distribution for simulated Compton scattering events (magenta), and simulated π^0 events that were analyzed as if they were a Compton photon (red), compared with the carbon-subtracted data (blue) at $E_\gamma = 285$ –305 MeV and over all Compton angles (Ω_{OA} cut from Fig. 1 is applied).

where $(E_{\gamma_i}, \vec{p}_{\gamma_i}c)$ and $(E_{\gamma_f}, \vec{p}_{\gamma_f}c)$ are the four vectors of the incident and scattered photon, respectively, and m_p is the proton mass.

The carbon-subtracted M_{miss} spectrum using the corrected carbon target scaling factors [12,18] is shown in Fig. 3. Simulations of both Compton scattering and π^0 production were passed through the same analysis chain, with the same cuts applied. The distributions from these two reactions were added together (according to their known cross section at a given energy and angle). From these spectra, there is clearly good agreement of the data with the expected distribution up to $M_{\text{miss}} \approx 980$ MeV/ c^2 . It is observed that π^0 photoproduction is the major source of background above a M_{miss} of approximately 940 MeV/ c^2 , and hence it is necessary to set a clear upper M_{miss} limit that coincides with the turn-on point of this background.

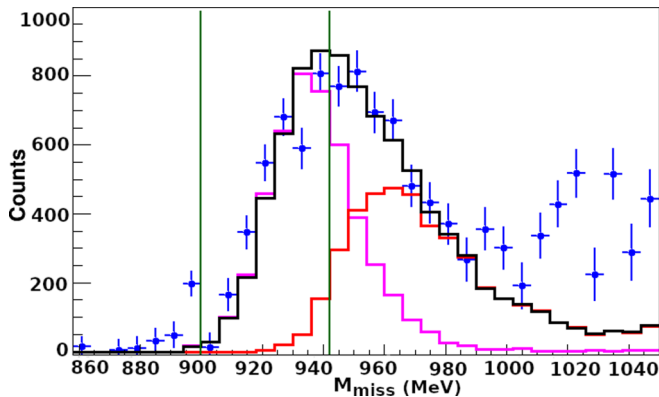


FIG. 3. Missing mass spectrum for carbon-subtracted data (blue), Monte Carlo simulated results from Compton scattering (magenta), and π^0 photoproduction (red) satisfying Compton cuts, and the sum of the two simulated contributions to show an expected distribution (black), all at $\theta_\gamma = 125$ –140° and $E_\gamma = 285$ –305 MeV. Two vertical lines (green) represent the missing mass integration limit.

TABLE I. Summary of results and uncertainties for the Compton Σ_{2z} asymmetry.

E_γ (MeV)	θ_γ	Σ_{2z}	Rand.	Syst.
265–285	87.5°	0.193	± 0.056	± 0.024
	102.5°	0.290	± 0.040	± 0.035
	117.5°	0.402	± 0.037	± 0.048
	132.5°	0.672	± 0.036	± 0.077
	147.5°	0.672	± 0.042	± 0.081
285–305	87.5°	0.121	± 0.040	± 0.016
	102.5°	0.279	± 0.034	± 0.033
	117.5°	0.428	± 0.038	± 0.048
	132.5°	0.591	± 0.029	± 0.066
	147.5°	0.751	± 0.046	± 0.085

Though M_{miss} spectra can be integrated up to the most conservative limit of 938 MeV/ c^2 (proton mass), the following steps were taken to maximize the integrated yield. First, the lower M_{miss} limit was fixed at 900 MeV/ c^2 for each energy and angle bin. Second, the asymmetry determined using the conservative M_{miss} upper limit of 938 MeV/ c^2 was taken as a reference. Finally, the asymmetry was allowed to vary a maximum of 5% by moving the M_{miss} upper limit to higher values compared to the reference. This $\pm 5\%$ is based on the systematic uncertainties from the choice of carbon target length and the ratio of π^0 photoproduction background to Compton scattering determined from simulation. As the asymmetry shifts either up or down for different bins, there is no concern about introducing a systematic shift from the ‘correct’ asymmetry. As an additional check the central value on the spin polarizabilities, as extracted by the method described below, were compared between the reference and final asymmetries, which indicated only small effects on γ_{E1E1} and γ_{M1E2} of approximately 20% of their errors and negligible effects on the other two. The resulting final M_{miss} upper limits are between 940–948 MeV/ c^2 , and further details on this work can be found in Ref. [12,18].

IV. RESULTS

The Σ_{2z} asymmetries for $E_\gamma = 265$ –285 MeV and $E_\gamma = 285$ –305 MeV, obtained by combining the results from the two beam times via their weighted average, are tabulated in Table I and shown in Fig. 4 along with determinations at 0° through dispersive sum rules [21,22]. While the absolute statistical errors only vary between 0.029–0.056, the relative errors vary between 5–33 % due to the small asymmetry at 90°. The systematic errors from the three different sources: target polarization (10%), beam polarization (2.7%), and carbon subtraction (3–6 %), were added in quadrature and their average between the 2014 and 2015 beam times for each Compton angle is listed in the table and shown as a separate block above the horizontal axis in the figures. These total systematic errors vary between 0.016–0.085 absolute, or 11–13 % relative. To study the sensitivity of the Σ_{2z} results on the SPs, a fixed- t dispersion relation code (HDPV) [2,23,24] was used to generate predicted asymmetries at fixed laboratory energies for

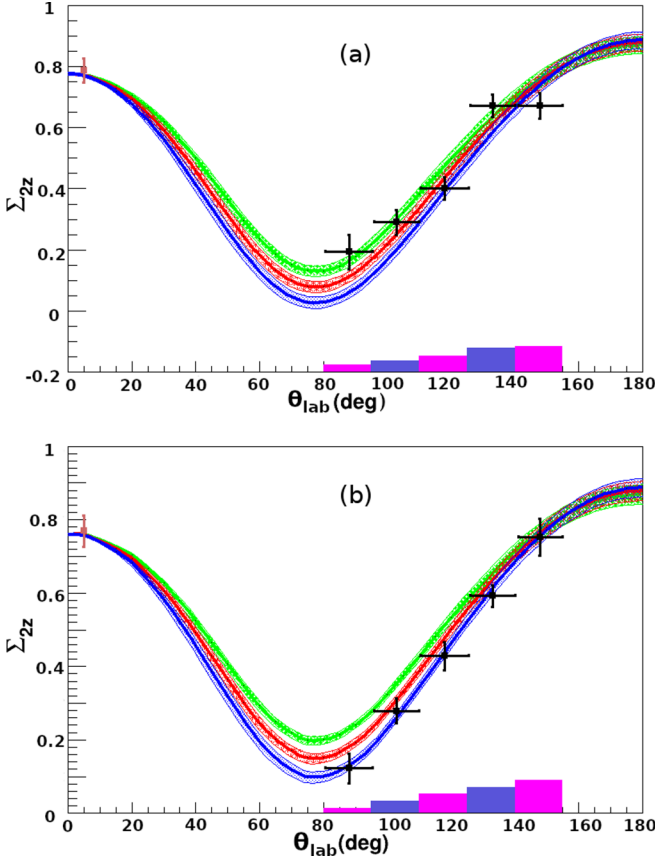


FIG. 4. Compton Σ_{2z} for $E_\gamma = 265\text{--}285$ MeV (a) and $E_\gamma = 285\text{--}305$ MeV (b). The red point is the value for Σ_{2z} at 0° , plotted at 5° for readability, as determined by dispersive sum rules [21,22]. The curves are from the HDPV dispersion theory calculation of Pasquini *et al.*, [2,23,24], where γ_{E-} [27] is fixed at $-3.5 \times 10^{-4} \text{ fm}^4$ and γ_{M-} [27] is set at $-0.5, 1.5$, or $3.5 \times 10^{-4} \text{ fm}^4$, in the green, red, or blue bands, respectively. The width of each band represents the other parameters, $\gamma_0, \gamma_\pi, \alpha_{E1} + \beta_{M1}$, and $\alpha_{E1} - \beta_{M1}$ varying within their experimental errors. The error bars shown are point-to-point statistical plus random systematic errors added in quadrature. The correlated systematic uncertainties are shown as a separate block above the horizontal axis for each Compton angle.

various values of the scalar and spin polarizabilities. Predictions within baryon chiral perturbation theory (B χ PT) [25] and heavy baryon chiral perturbation theory (HB χ PT) [26,27] are also available, but are not shown in Fig. 4 to preserve readability. The code used nominal values for the scalar polarizabilities of $\alpha_{E1} + \beta_{M1} = 13.8 \pm 0.4$ (Baldin sum rule) [28] and $\alpha_{E1} - \beta_{M1} = 8.7 \pm 0.7$ (in units of 10^{-4} fm^3) [29], and for the SPs of $\gamma_0 = -0.929 \pm 0.105$ [21,22] and $\gamma_\pi = 8 \pm 1.8$ (in units of 10^{-4} fm^4) [5]. It should be noted that the value for $\alpha_{E1} - \beta_{M1}$ was chosen as the current PDG numbers [29], despite the debate regarding them [30,31], as the focus of this study is on the spin polarizabilities. It should also be noted that this value for γ_π does not include the π^0 -pole component, set as $-46.7 \times 10^{-4} \text{ fm}^4$ [31] in all of these studies.

Though γ_0 and γ_π can form a basis of the SPs with γ_{E1E1} and γ_{M1M1} , they can alternatively form an orthogonal basis with $\gamma_{E-} = \gamma_{E1E1} - \gamma_{E1M2}$ and $\gamma_{M-} = \gamma_{M1M1} - \gamma_{M1E2}$ [27].

TABLE II. Polarizabilities in 10^{-4} fm^4 from fitting Σ_{2z} , Σ_{2x} , and Σ_3^{LEGS} asymmetries using either a HDPV [2,23,24] or a B χ PT [25] calculation, and weighted average of the SPs.

	Σ_{2z}, Σ_{2x} , and Σ_3^{LEGS} data fits		
	HDPV	B χ PT	Weighted average
γ_{E1E1}	-3.18 ± 0.52	-2.65 ± 0.43	-2.87 ± 0.52
γ_{M1M1}	2.98 ± 0.43	2.43 ± 0.42	2.70 ± 0.43
γ_{E1M2}	-0.44 ± 0.67	-1.32 ± 0.72	-0.85 ± 0.72
γ_{M1E2}	1.58 ± 0.43	2.47 ± 0.42	2.04 ± 0.43
χ^2/dof	1.14	1.36	

In Fig. 4, γ_{E-} was fixed at $-3.5 \times 10^{-4} \text{ fm}^4$ and γ_{M-} was set at $-0.5, 1.5$, or 3.5 in the same units. The various bands represent the different values for γ_{M-} , while the spread of each band is a result of allowing $\gamma_0, \gamma_\pi, \alpha_{E1} + \beta_{M1}$, and $\alpha_{E1} - \beta_{M1}$, to vary by their experimental errors. It is clear

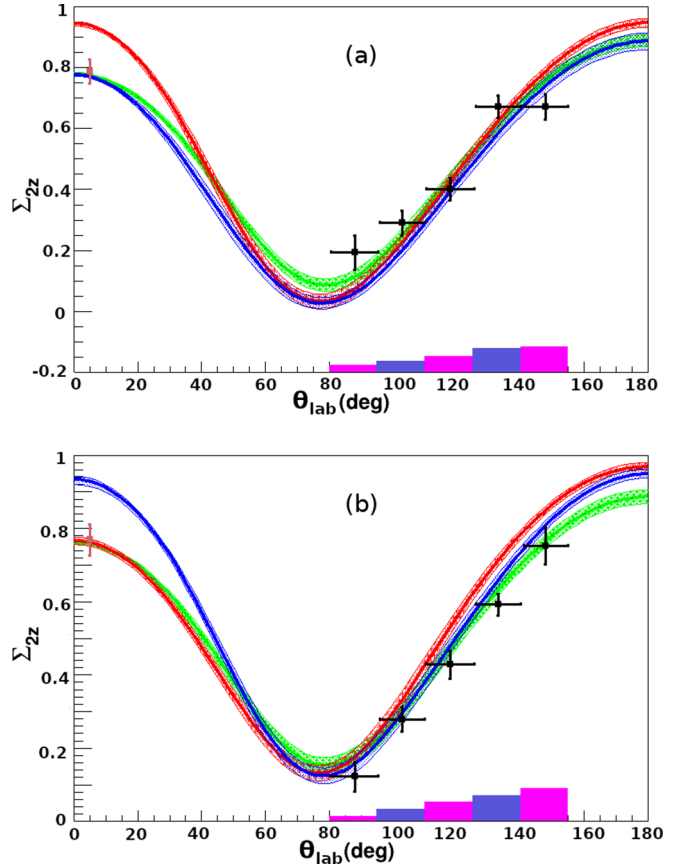


FIG. 5. Compton Σ_{2z} for $E_\gamma = 265\text{--}285$ MeV (a) and $E_\gamma = 285\text{--}305$ MeV (b). The red point is the value for Σ_{2z} at 0° , plotted at 5° for readability, as determined by dispersive sum rules [21,22]. The green, red, and blue curves are from HDPV [2,23,24], B χ PT [25], and HB χ PT [26,27] calculations, respectively. For each, the central curve uses the weighted average values from Table II, and the width of each band represents the parameters varying within the errors quoted in the same table. The error bars shown are point-to-point statistical plus random systematic errors added in quadrature. The correlated systematic uncertainties are shown as a separate block above the horizontal axis for each Compton angle.

from Fig. 4 that the Σ_{2z} data in this energy range indicate a sensitivity to γ_{M-} of approximately ± 2 in the standard units. Alternatively, γ_{M-} can be fixed at $1.5 \times 10^{-4} \text{ fm}^4$ and γ_{E-} set at -5.5 , -3.5 , or -1.5 in the same units. Unlike the previous case, Σ_{2z} in this energy range showed a weak sensitivity to γ_{E-} .

V. GLOBAL FITS

A global analysis of Σ_{2z} data from this measurement, along with the published Σ_{2x} and Σ_3^{LEGS} results, and the prior values of γ_0 and γ_π , was performed to study the model dependence and extract the SPs. This was done by fitting the asymmetry data using the HDPV calculation [2,23,24] and a $B\chi$ PT calculation [25]. The extracted SPs determined using each model are summarized in Table II. The fit with HDPV results in $\gamma_{E-} = -2.74 \times 10^{-4} \text{ fm}^4$ and $\gamma_{M-} = 1.4$, in the same units, similar to the values used for the theoretical bands in Fig. 4. The values from the two models are fairly consistent, and the best estimate of a central value is given by the weighted average in the last column of Table II. The errors for the weighted average values were conservatively taken as the larger of the two fits. These errors were chosen in favor of the weighted error, because the weighted errors assume the uncertainties in the theoretical calculations are uncorrelated, for which this paper makes no statement. The data are again shown in Fig. 5, now with theoretical calculations for HDPV [2,23,24], $B\chi$ PT [25], and $HB\chi$ PT [26,27], using the weighted average values for the SPs.

VI. SUMMARY

In summary, model-dependent extractions of the SPs from a combined data fit of double- and single-polarized Compton scattering asymmetry results in the $\Delta(1232)$ resonance

region are presented. These extracted SPs are also in good agreement with dispersion relation [2,23,24], baryon chiral perturbation theory [32], heavy baryon chiral perturbation theory [26,27], K -matrix theory [33], and chiral Lagrangian [34] predictions. Although the uncertainties in the SPs are significantly improved compared to previously reported results [6], forthcoming Σ_3 results from MAMI experiments [35] are expected to provide further improvements in the determination of these fundamental nuclear structure terms.

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- [1] F. Hagelstein, R. Miskimen, and V. Pascalutsa, *Prog. Part. Nucl. Phys.* **88**, 29 (2016).
 - [2] B. Pasquini, D. Drechsel, and M. Vanderhaeghen, *Phys. Rev. C* **76**, 015203 (2007).
 - [3] J. Ahrens *et al.*, *Phys. Rev. Lett.* **87**, 022003 (2001).
 - [4] H. Dutz *et al.*, *Phys. Rev. Lett.* **91**, 192001 (2003).
 - [5] M. Camen *et al.*, *Phys. Rev. C* **65**, 032202(R) (2002).
 - [6] P. P. Martel *et al.*, *Phys. Rev. Lett.* **114**, 112501 (2015).
 - [7] G. Blanpied *et al.*, *Phys. Rev. C* **64**, 025203 (2001).
 - [8] K. H. Kaiser *et al.*, *Nucl. Instrum. Methods Phys. Res. A* **593**, 159 (2008).
 - [9] A. Jankowiak *et al.*, *Eur. Phys. J. A* **28**, 149 (2006).
 - [10] K. Aulenbacher *et al.*, *Nucl. Instrum. Methods Phys. Res. A* **391**, 498 (1997).
 - [11] V. Tioukine, K. Aulenbacher, and E. Riehn, *Rev. Sci. Instrum.* **82**, 033303 (2011).
 - [12] D. Paudyal, Ph.D. thesis, University of Regina (2017), <http://www2.kph.uni-mainz.de/images/publications/phd/Paudyal-PhD-thesis.pdf>.
 - [13] J. C. McGeorge *et al.*, *Eur. Phys. J. A* **37**, 129 (2008).
 - [14] A. Thomas, *Eur. Phys. J. Spec. Top.* **198**, 171 (2011).
 - [15] A. Starostin *et al.*, *Phys. Rev. C* **64**, 055205 (2001).
 - [16] C. M. Tarbert *et al.*, *Phys. Rev. Lett.* **100**, 132301 (2008).
 - [17] F. Afzal, Ph.D. thesis, University of Bonn (2019), <http://hss.ulb.uni-bonn.de/2019/5551/5551.htm>.
 - [18] D. Paudyal, G. Huber, P. P. Martel, and D. Hornidge, Technical Report, http://lichen.phys.uregina.ca/cbtaps/techreport_pi02019.pdf.
 - [19] M. Unverzagt *et al.*, *Eur. Phys. J. A* **39**, 169 (2009).
 - [20] R. Novotny, *IEEE Trans. Nucl. Sci.* **38**, 379 (1991).
 - [21] O. Gryniuk, F. Hagelstein, and V. Pascalutsa, *Phys. Rev. D* **92**, 074031 (2015).
 - [22] O. Gryniuk, F. Hagelstein, and V. Pascalutsa, *Phys. Rev. D* **94**, 034043 (2016).
 - [23] B. R. Holstein, D. Drechsel, B. Pasquini, and M. Vanderhaeghen, *Phys. Rev. C* **61**, 034316 (2000).
 - [24] D. Drechsel, B. Pasquini, and M. Vanderhaeghen, *Phys. Rep. C* **378**, 99 (2003).
 - [25] V. Lensky and V. Pascalutsa, *Eur. Phys. J. C* **65**, 207 (2010).
 - [26] H. W. Griebhammer, J. A. McGovern, and D. R. Phillips, *Eur. Phys. J. A* **52**, 139 (2016).
 - [27] H. W. Griebhammer, J. A. McGovern, and D. R. Phillips, *Eur. Phys. J. A* **54**, 37 (2018).

- [28] V. Olmos de León *et al.*, [Eur. Phys. J. A **10**, 207 \(2001\)](#).
- [29] M. Tanabashi *et al.*, [Phys. Rev. D **98**, 030001 \(2018\)](#).
- [30] H. W. Griebhammer, J. A. McGovern, D. R. Phillips, and G. Feldman, [Prog. Part. Nucl. Phys. **67**, 841 \(2012\)](#).
- [31] B. Pasquini, P. Pedroni, and S. Sconfiatti, [J. Phys. G: Nucl. Part. Phys. **46**, 104001 \(2019\)](#).
- [32] V. Lensky, J. A. McGovern, and V. Pascalutsa, [Eur. Phys. J. C **75**, 604 \(2015\)](#).
- [33] S. Kondratyuk and O. Scholten, [Phys. Rev. C **64**, 024005 \(2001\)](#).
- [34] A. M. Gasparyan, M. F. M. Lutz, and B. Pasquini, [Nucl. Phys. A **866**, 79 \(2011\)](#).
- [35] C. Collicott *et al.* (in preparation, 2020).