

Electrical Breakdown Characteristics of Supercritical Trifluoroiodomethane-Carbon Dioxide (CF₃I-CO₂) Mixtures

Jia Wei^{1*}, Alfonso Cruz¹, Farhina Haque², Chanyeop Park², and Lukas Graber¹

1. School of Electrical and Computer Engineering, Georgia Institute of Technology, Atlanta, Georgia, USA

2. Department of Electrical and Computer Engineering, Mississippi State University, Mississippi State, MS, USA
jia.wei@gatech.edu

Abstract- Supercritical fluids (SCFs), as a new type of dielectric medium, combine highly desirable properties of gaseous and liquid dielectrics, most notably: low viscosity, high heat transfer capability, and high dielectric strength. This paper, for the first time, reports the results of breakdown experiments on the mixture of trifluoroiodomethane (CF₃I) and carbon dioxide (CO₂) under the supercritical conditions. Experiments were conducted in a 0.1 mm gap under a uniform electric field. The state of substance during experiments was obtained by the calculation from the equation of state (EoS), and by evaluating the optical appearance. The experiment result shows a dielectric strength of the supercritical CF₃I-CO₂ mixture at an average of 278 kV/mm with a 10% CF₃I mass fraction. Our experiments also suggest that the dielectric behavior of this supercritical CF₃I-CO₂ mixture manifests a discontinuity of the dielectric strength near the critical point. This paper indicates the suitability of using supercritical CF₃I-CO₂ mixture for high power density applications.

I. INTRODUCTION

Supercritical fluids (SCFs), a state achieved when temperature and pressure are above the critical point, as shown in Fig. 1, have been widely studied and applied in the field of chemistry, including polymer processing, drying, and extraction of products from natural materials [1]. Recently, SCFs have drawn attention in the area of electrical discharge: not only for chemical synthesis and nanomaterials fabrication [2] but also on utilizing SCFs as dielectric media for power applications [3, 4]. Specifically, Zhang *et al.* demonstrated a supercritical nitrogen (N₂) plasma switch and showed that supercritical N₂ exhibits excellent dielectric properties and recovery behavior from breakdown [5, 6]. Kiyani *et al.* conducted research on the pre-breakdown and breakdown in supercritical carbon dioxide (CO₂) with different electrode geometries and revealed a polarity effect in supercritical CO₂ [7, 8]. Besides investigations on pure substances such as supercritical N₂ and supercritical CO₂, our previous study has focused on the breakdown characteristics of supercritical carbon dioxide-ethane (CO₂-C₂H₆) mixtures. The study found that mixtures of SCFs can yield a useful combination of properties such as the dielectric strength and the critical point, thus allowing for a broader range of applications [9]. A theoretical model was developed to describe the mechanism of electrical breakdown in supercritical CO₂ with molecular clusters formation. By analyzing the cross-section data of

clusters, it has been confirmed that the cross-section data of clusters decrease from those of gases [10-12].

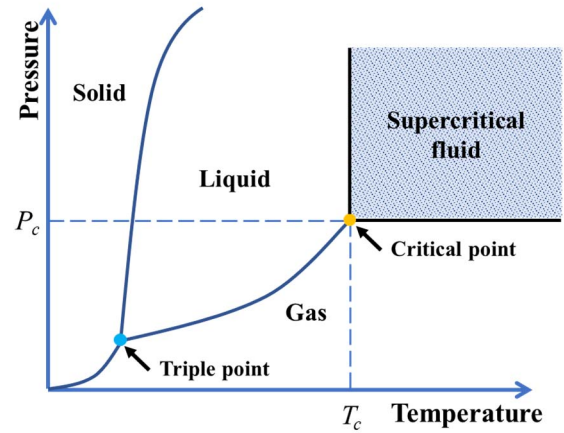


Fig. 1 A typical phase diagram. It shows four distinct thermodynamic phases, including solid, liquid, gaseous, and supercritical fluid.

Trifluoroiodomethane (CF₃I), as one of the candidates of sulfur hexafluoride (SF₆) alternatives, has several beneficial physical properties. It is nonflammable, colorless, and chemically stable. In addition, CF₃I has a low global warming potential (GWP) of 0.4 against approximately 22,800 for SF₆. Also, the overall atmospheric lifetime of CF₃I is very short due to the weak chemical bond C-I. A comparison of some important physical properties is given in Table 1 [13-15].

Table 1. Physical properties of SF₆, CO₂, and CF₃I.

	Molecular formula	T_c (K)	P_c (MPa)	GWP	Atmospheric lifetime (Years)
Carbon dioxide	CO ₂	304.2	7.39	1	200
Trifluoroiodomethane	CF ₃ I	395.2	4.04	0.4	30 (days)
Sulfur Hexafluoride	SF ₆	318.7	3.76	22,800	3,200

Although extensive research has been carried out on the electric discharge generated in supercritical CO₂ and N₂, there is no consideration of many other highly promising supercritical substances or mixtures. Especially, for power applications, there has not been any investigation of the SF₆

alternatives in the supercritical state. Similarly, although numerous research efforts have focused on the insulation performance of CF_3I and its mixtures with CO_2 and N_2 over the past few years [16-18], these studies were all conducted in the gaseous phase, which suffers under limited heat transfer and dielectric strength compared to SCFs. In this paper, we demonstrate, to the best of our knowledge, the first experimental measurement on the breakdown characteristics of a supercritical CF_3I - CO_2 mixture.

II. EXPERIMENTAL SETUP

A. Breakdown test setup

A complete setup with the high pressure chamber (A), a temperature-controlled water bath (B), pressure gauges (C), a thermal control system (D and E), a high precision scale, and control valves, is assembled and built as shown in Fig. 2. The electrode on the high voltage side is attached to the brass conductor, and the grounded electrode is clamped by a stainless steel plate that is supported and electrically grounded by three stainless steel rods (F). The copper electrodes (G) have a spherical geometry with a diameter of 20 mm.

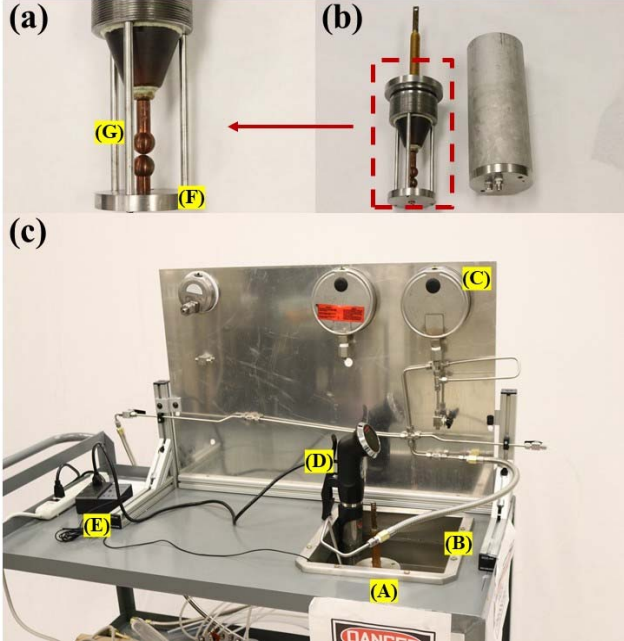


Fig. 2 The electrical breakdown experimental setup. (a) the internal configuration of the chamber, (b) the external view of the high pressure chamber, (c) the testbed for the SCF breakdown experiment, which contains a temperature-controlled bath, pressure gauges, and a thermal control system.

The experiments were conducted under DC voltage (Spellman SL60PN1200, rated 60 kV and 20 mA). The voltage was ramped up at 0.5 kV/s, and the gap between the two electrodes was set to be 0.1 mm. The gap is measured by a stainless steel feeler gauge with accuracy within $\pm 5\%$. The breakdown voltage was then measured and recorded fifteen times at each condition. The accuracy of the voltage measurement is estimated to be within $\pm 3\%$. The average breakdown voltage and the corresponding scattering data were

calculated based on these fifteen measurements. The time interval between two successive breakdown events was about two minutes to reduce the effect of the medium state variation caused by the arc discharge and allow discharge products to move away from the zone of the highest electric field.

B. Experimental conditions and isothermal properties

Sometimes, SCF can be represented as an intermediate state between gas and liquid states; it combines the properties of both phases such as low viscosity like a gas, and effective heat transfer capabilities like a liquid. In addition to these physical properties, which are promising for a wide range of applications, SCFs also offer high dielectric strength due to its high density. Therefore, more research projects have been conducted in utilizing SCFs as dielectric and insulating media in power applications.

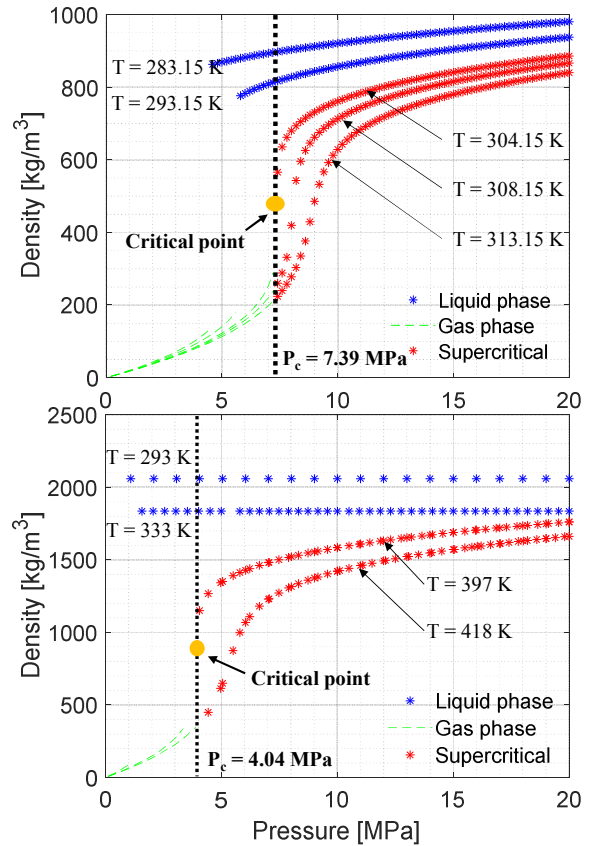


Fig. 3. Density – pressure (ρ - P) diagrams of CO_2 (above) and CF_3I (below) with isothermal lines. The critical points for each substance are represented by the yellow dot at $T_c = 304.2$ K, $P_c = 7.39$ MPa (CO_2) and $T_c = 395.2$ K, $P_c = 4.04$ MPa (CF_3I) [19-21].

Fig. 3 shows density – pressure (ρ - P) diagrams of both CO_2 (above) and CF_3I (below) with isothermal lines. Green dashed lines represent the fluid in the gaseous phase. Blue and red asterisks represent the fluid in the liquid phase and the supercritical state, respectively. The critical points are represented by the yellow dots. Near the critical point, small changes in the medium pressure result in significant changes in the density. Critical anomalies can be observed in the

vicinity of the critical point, such as maximum values of thermal conductivity and heat capacity.

The critical point of the $\text{CF}_3\text{I}-\text{CO}_2$ mixture at 10% CF_3I mass fraction is acquired from the equation of state (EoS) [19, 20]. Specifically, the calculation of the critical point (T_c , P_c) of a supercritical mixture is based on solving equations including the mathematical definition of critical point [22] and the EoS of corresponding substances:

$$\left(\frac{\partial P}{\partial V}\right)_{T_c} = \left(\frac{\partial^2 P}{\partial V^2}\right)_{T_c} = 0 \quad (1)$$

and two determinants of multicomponent mixture provided by Gibbs [23]:

$$D_1 = \det \left| \frac{\partial^2 G}{\partial x_i \partial x_j} \right| = 0 \quad (2)$$

$$D_2 = \det \left| (1 - \delta_{ki}) \frac{\partial^2 G}{\partial x_i \partial x_j} + \delta_{ki} \frac{\partial D_1}{\partial x_j} \right| = 0 \quad (3)$$

where V is the volume, P is the pressure, T_c represents the critical temperature, G is Gibbs energy, x_i is composition of component i , k is a constant.

The calculated critical point of this mixture ($T_c = 315.0$ K, $P_c = 7.20$ MPa) is in good agreement with the optically determined critical point using a high pressure optical vessel.

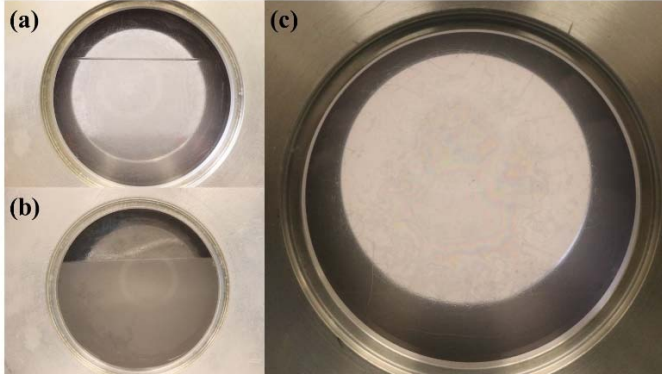


Fig 4. Determination of the critical point of a sample of $\text{CF}_3\text{I}-\text{CO}_2$ at 10% CF_3I mass fraction by observing the critical opalescence. (a) the medium at a temperature below its critical point ($T < T_c$), (b) the medium at a temperature near its critical point ($T \approx T_c$), (c) the medium at a temperature above its critical point ($T > T_c$).

III. RESULTS & DISCUSSIONS

To identify the state of the $\text{CF}_3\text{I}-\text{CO}_2$ mixture inside the high pressure chamber, the critical point needs to be determined. Besides the critical values calculated by the EoS, a static equilibrium apparatus, as shown in Fig. 4, was also built to replicate the experimental condition and to confirm the critical points visually. Specifically, the critical opalescence is observed to determine the critical state of such a binary mixture. Fig. 4 shows a sample of this mixture with gravity directed downward. In Fig. 4(a) the temperature is below its critical point: two distinct phases exist. The boundary between

gas and liquid is visible. As the temperature rises, the gas-liquid phase boundary starts to vanish, as shown in Fig. 4(b). Fig. 4(c) shows the sample significantly above its critical point, the phase boundary disappears. As a result, the critical point observed from the optical chamber is $T_c = 317.5$ K, $P_c = 7.25$ MPa, which is in good agreement with the critical point calculated from the EoS ($T_c = 315.0$ K, $P_c = 7.20$ MPa).

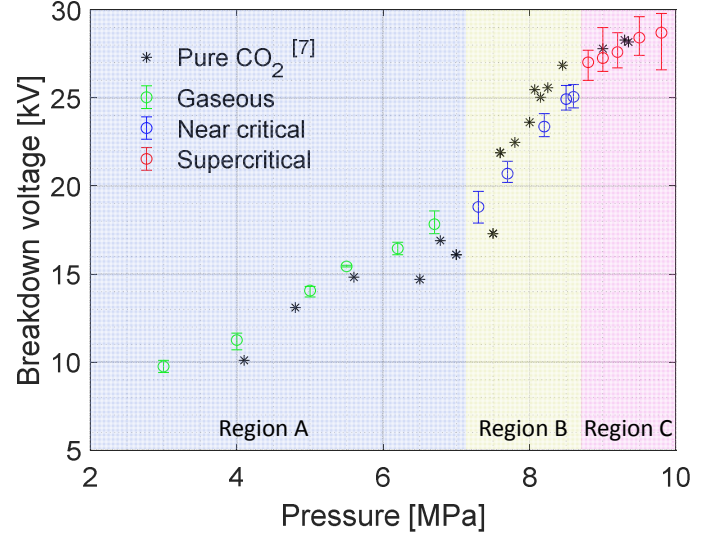


Fig 5. Experimental results for the breakdown voltage of supercritical $\text{CF}_3\text{I}-\text{CO}_2$ mixture up to 10 MPa in a uniform electric field at a 0.1 mm gap at a temperature of 323 K. The mass fraction of CF_3I is 10%. The black asterisks indicate the average of the measurement data of pure CO_2 [9]. Open circles represent the average breakdown voltage, and vertical error bars represent the corresponding experimental scattering data.

The breakdown voltage measurements were carried out at a temperature of 323 K (50°C) from the gaseous phase to the supercritical condition. Fig. 5 shows the measured breakdown voltages as a function of the pressure. The average breakdown voltage of fifteen measurements and their scattering data at one experimental condition are represented by an open circle and a vertical error bar. The result also shows a comparison of the dielectric strength measurement of this $\text{CF}_3\text{I}-\text{CO}_2$ mixture with experimental results of pure CO_2 [9]. The result indicates that the breakdown strength of the $\text{CF}_3\text{I}-\text{CO}_2$ mixture in the gaseous phase is, in general, higher than those of pure CO_2 , even for only a 10% CF_3I mass fraction (region A). At the situation near the critical point, the $\text{CF}_3\text{I}-\text{CO}_2$ mixture exhibits a discontinuity of the slope where the substance undergoes a phase change (region B). Such a discontinuity can be explained from the microscopic point of view by the molecular clustering, as well as the macroscopic point of view by the density fluctuation. It can also be estimated that the density fluctuation of this binary mixture reaches a maximum near 8.8 MPa, which also represents a boundary surface that separates the supercritical region into gas-like phases and liquid-like phases. In the supercritical state, the breakdown voltage of both $\text{CF}_3\text{I}-\text{CO}_2$ mixture and pure CO_2 have a trend to saturate at higher densities (region C). However, no

significant difference is observed on the dielectric strength of the pure CO₂ and CF₃I-CO₂ (10% mass fraction of CF₃I) mixture at the supercritical state. This is not very surprising since a 10% of CF₃I mass fraction in this binary mixture fluid is only equivalent to a 2.4% of mole fraction. Although it is easier to use the mass fraction for experiments, the mole fraction is a better measure to understand the impact onto the dielectric strength from an electronic point of view since it correlates with the probability of electron-molecule collisions with one over the other constituents. In this case, a 2.4% mole fraction of CF₃I indicates that electrons only rarely collide with the strongly electronegative CF₃I molecules, i.e. the overall effect onto the dielectric strength is limited.

IV. CONCLUSIONS

In this study, we conducted breakdown measurements of supercritical CF₃I-CO₂ mixture with a CF₃I mass fraction of 10% at 323 K. We also compared the result with our previous breakdown strength measurement of pure supercritical CO₂. Experimental conditions such as electrode geometry and gap distance were maintained, except for the experimental temperature (two dielectric fluids have different critical temperatures). Moreover, we showed the breakdown strength of CF₃I-CO₂ at the supercritical state, which is higher than that in the gaseous phase. Meanwhile, since the critical temperature of the CF₃I-CO₂ mixture at 10% CF₃I mass fraction is only about 10 K above the critical temperature of pure CO₂, no additional temperature withstanding capability is necessary to implement this new dielectric medium. Experimental investigations on the dielectric strength of this mixture at higher CF₃I mass fraction will be conducted in the near future. This study also indicates the suitability of using SCF mixtures as dielectric and insulating media for applications that require efficient heat dissipation, high-speed motion, and strong electric fields.

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