



Quivers with subadditive labelings: classification and integrability

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Abstract

Strictly subadditive, subadditive and weakly subadditive labelings of quivers were introduced by the second author, generalizing Vinberg’s definition for undirected graphs. In our previous work we have shown that quivers with strictly subadditive labelings are exactly the quivers exhibiting Zamolodchikov periodicity. In this paper, we classify all quivers with subadditive labelings. We conjecture them to exhibit a certain form of integrability, namely, as the T -system dynamics proceeds, the values at each vertex satisfy a linear recurrence. Conversely, we show that every quiver integrable in this sense is necessarily one of the 19 items in our classification. For the quivers of type $\hat{A} \otimes A$ we express the coefficients of the recurrences in terms of the partition functions for domino tilings of a cylinder, called *Goncharov–Kenyon Hamiltonians*. We also consider tropical T -systems of type $\hat{A} \otimes A$ and explain how affine slices exhibit solitonic behavior, i.e. soliton resolution and speed conservation. Throughout, we conjecture how the results in the paper are expected to generalize from $\hat{A} \otimes A$ to all other quivers in our classification.

Keywords Cluster algebras · Zamolodchikov integrability · Discrete solitons · Domino tilings · Linear recurrence

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Introduction

A *quiver* Q is a directed graph without 1-cycles (i.e. loops) and directed 2-cycles. For a vertex v of a quiver, one can define a certain operation called a *mutation*, which produces a new quiver denoted $\mu_v(Q)$ (see Definition 1.1.5). We say that a quiver is bipartite if its underlying graph is bipartite, in which case we say that a map $\epsilon: \text{Vert}(Q) \rightarrow \{0, 1\}$, $v \mapsto \epsilon_v$ is a *bipartition* if for every edge $u \rightarrow v$ of Q we have $\epsilon_u \neq \epsilon_v$. Here $\text{Vert}(Q)$ is the set of vertices of Q .

It is clear from Definition 1.1.5 that μ_u and μ_v commute if u, v are not connected by an edge in Q . Therefore, we can define

$$\mu_\circ = \prod_{u: \epsilon_u=0} \mu_u; \quad \mu_\bullet = \prod_{v: \epsilon_v=1} \mu_v.$$

We say that Q is *recurrent* if $\mu_\circ(Q) = \mu_\bullet(Q) = Q^{\text{op}}$ where Q^{op} is the same quiver as Q but with all the arrows reversed.

Let Q be a bipartite recurrent quiver. Denote $\mathbf{x} := \{x_v\}_{v \in \text{Vert}(Q)}$ to be the set of indeterminates, one for each vertex of Q , and let $\mathbb{Q}(\mathbf{x})$ be the field of rational functions in these variables. The T -system associated with Q is a family $T_v(t)$ of elements of $\mathbb{Q}(\mathbf{x})$ satisfying the following relations for all $v \in \text{Vert}(Q)$ and all $t \in \mathbb{Z}$:

$$T_v(t+1)T_v(t-1) = \prod_{u \rightarrow v} T_u(t) + \prod_{v \rightarrow w} T_w(t).$$

Here the products are taken over all *arrows* connecting the two vertices.

It is clear that the parity of $t + \epsilon_v$ in all of the terms is the same, so the T -system associated with Q splits into two completely independent ones. Without loss of generality we may consider only one of them. From now on we assume that the T -system is defined only for $t \in \mathbb{Z}$ and $v \in \text{Vert}(Q)$ satisfying

$$t + \epsilon_v \equiv 0 \pmod{2}.$$

The T -system is set to the following initial conditions:

$$T_v(\epsilon_v) = x_v$$

for all $v \in \text{Vert}(Q)$.

Let us say that the T -system associated with a recurrent quiver Q is *integrable* if for every vertex $v \in \text{Vert}(Q)$, there exists an integer N and elements $J_0, J_1, \dots, J_N \in \mathbb{Q}(\mathbf{x})$ satisfying $J_0, J_N \neq 0$ and

$$\sum_{j=0}^N J_j T_v(t + 2j) = 0$$

for all $t \in \mathbb{Z}$ with $t + \epsilon_v$ even. We also refer to a recurrent quiver Q as *Zamolodchikov integrable* if the associated T -system is integrable. If the recurrence has the form $T_v(t + 2N) = T_v(t)$ for all $t \in \mathbb{Z}$ and $v \in \text{Vert}(Q)$, then we call Q *Zamolodchikov periodic*.

Just as in the periodic case, we call a quiver Zamolodchikov integrable when the *bipartite* T -system is integrable. More general notions of T -systems can be found in [28].

Zamolodchikov periodicity for the case when Q is a tensor product of two finite *ADE* Dynkin diagrams has been studied extensively (see [10,12,14,24,25,31,37,41,42]) and was proven in full generality in [20] and later in [17,18], where tropical Y -systems played a major role. By analyzing tropical T -systems (see Part 4), we have classified Zamolodchikov periodic quivers in [13], where we showed that these are exactly the quivers admitting a strictly subadditive labeling (Definition 1.1.4).

Besides thermodynamic Bethe ansatz [42], T -systems and Y -systems arise naturally in a lot of different contexts in physics and representation theory, e.g. [11,22,23,25,27,29,32], see [26] for a survey.

Assem, Reutenauer and Smith [1] showed that the affine Dynkin diagrams of types $\hat{A}$ and $\hat{D}$ are Zamolodchikov integrable, and later Keller and Scherotzke [21] extended this result to all affine Dynkin diagrams. Conversely, it was shown in [1] that if every vertex of a Zamolodchikov integrable quiver Q is either a source or a sink then Q is necessarily an affine Dynkin diagram.

In Sects. 1.2 and 4.2 we prove (Theorem 1.2.5) that if a bipartite recurrent quiver is Zamolodchikov integrable then it admits a subadditive labeling (see Definition 1.1.4).

We then classify (Theorem 2.2.5) quivers that admit subadditive labelings in Part 2. We conjecture all of them to be Zamolodchikov integrable (Conjecture 3.6.1).

When Q is a tensor product (see Definition 1.1.1) of type $\hat{A} \otimes A$, it was shown in [30] that Q is Zamolodchikov integrable. In Sects. 3.1–3.4, we express the recurrence coefficients $J_1, \dots, J_N$ for the vertices of such Q in terms of partition functions of domino tilings on the cylinder, called *Goncharov–Kenyon Hamiltonians*. See Theorem 3.2.5 and Corollary 3.4.2.

In Sect. 3.5 we show that the above Goncharov–Kenyon Hamiltonians belong to the upper cluster algebra.

In Part 4, we analyze the tropical T -system associated with quivers admitting a subadditive labeling. We show (Corollary 4.2.3) that when $t \gg 0$ or $t \ll 0$, every affine slice of the tropical T -system moves with some constant speed. We explain how this can be seen as *soliton resolution* in Sect. 4.2, and then we proceed to show *speed conservation* in Sect. 4.3: for the quiver of type $\hat{A}_{2n-1} \otimes A_m$, the speeds of the solitons at $t \gg 0$ are equal to the speeds of solitons at $t \ll 0$ after one applies a diagram automorphism to A_m . See Example 4.3.2 for an illustration of these solitonic phenomena.

Finally, in Sects. 3.6 and 4.4 we conjecture most of our results for all other quivers in our classification.

Part 1. Zamolodchikov integrable quivers.

1.1 Preliminaries

1.1.1 Bigraphs

In [36] Stembridge studies admissible W -graphs for the case when $W = I(p) \times I(q)$ is a direct product of two dihedral groups. These W -graphs encode the structure of representations of Iwahori-Hecke algebras, and were first introduced by Kazhdan and Lusztig in [19]. The following definitions are adapted from [36] with slight modifications. A *bigraph* is an ordered pair of undirected loopless graphs (Γ, Δ) which share a common set of vertices $V := \text{Vert}(\Gamma) = \text{Vert}(\Delta)$ and do not share edges. A bigraph is called *bipartite* if there is a map $\epsilon : V \rightarrow \{0, 1\}$ such that for every edge (u, v) of Γ or of Δ we have $\epsilon_u \neq \epsilon_v$.

There is a simple one-to-one correspondence between bipartite quivers and bipartite bigraphs. Namely, to each bipartite quiver Q with a bipartition $\epsilon : \text{Vert}(Q) \rightarrow \{0, 1\}$ we associate a bigraph $G(Q) = (\Gamma(Q), \Delta(Q))$ on the same set of vertices defined as follows:

- $\Gamma(Q)$ contains an (undirected) edge (u, v) if and only if Q contains a directed edge $u \rightarrow v$ with $\epsilon_u = 0, \epsilon_v = 1$;
- $\Delta(Q)$ contains an (undirected) edge (u, v) if and only if Q contains a directed edge $u \rightarrow v$ with $\epsilon_u = 1, \epsilon_v = 0$.

Similarly, we can direct the edges of any given bipartite bigraph G to get a bipartite quiver $Q(G)$.

It is convenient to think of (Γ, Δ) as of a single graph with edges of two colors: red for the edges of Γ and blue for the edges of Δ .

Definition 1.1.1 Let S and T be two bipartite undirected graphs. Then their *tensor product* $S \otimes T$ is a bipartite bigraph $G = (\Gamma, \Delta)$ with vertex set $\text{Vert}(S) \times \text{Vert}(T)$ and the following edge sets:

- for each edge $\{u, u'\} \in S$ and each vertex $v \in T$ there is an edge between (u, v) and (u', v) in Γ ;
- for each vertex $u \in S$ and each edge $\{v, v'\} \in T$ there is an edge between (u, v) and (u, v') in Δ ;

An example of a tensor product is given in Fig. 1.

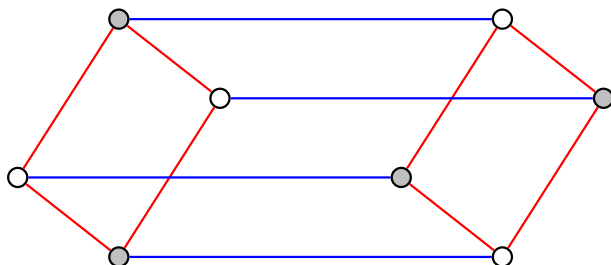


Fig. 1 A tensor product of a square (type $\hat{A}_3$) and a single edge (type A_2)

1.1.1.1 Reformulation of the dynamics in terms of bigraphs

Let $G = (\Gamma, \Delta)$ be a bipartite bigraph with a vertex set V . Then the associated T -system for G is defined as follows:

$$T_v(t+1)T_v(t-1) = \prod_{(u,v) \in \Gamma} T_u(t) + \prod_{(v,w) \in \Delta} T_w(t);$$

$$T_v(\epsilon_v) = x_v.$$

It is easy to see that this system is equivalent to the corresponding system defined for $Q(G)$ in the Introduction.

1.1.2 Finite and affine ADE Dynkin diagrams and their Coxeter numbers

By a *finite ADE Dynkin diagram* we mean a Dynkin diagram of type A_n , D_n , E_6 , E_7 , or E_8 . An *affine ADE Dynkin diagram* is a Dynkin diagram of type $\hat{A}_n$, $\hat{D}_n$, $\hat{E}_6$, $\hat{E}_7$, or $\hat{E}_8$, see Fig. 2.

The following characterization of finite and affine ADE Dynkin diagrams is due to Vinberg [40]:

Theorem 1.1.2 *Let $G = (V, E)$ be an undirected graph with possibly multiple edges. Then:*

- *G is a finite ADE Dynkin diagram if and only if there exists a map $v : V \rightarrow \mathbb{R}_{>0}$ such that for all $v \in V$,*

$$2v(v) > \sum_{(u,v) \in E} v(u).$$

- *G is an affine ADE Dynkin diagram if and only if there exists a map $v : V \rightarrow \mathbb{R}_{>0}$ such that for all $v \in V$,*

$$2v(v) = \sum_{(u,v) \in E} v(u). \quad (1.1.1)$$

The values of v satisfying (1.1.1) are given in Fig. 2, normalized so that they are integers with greatest common divisor equal to 1.

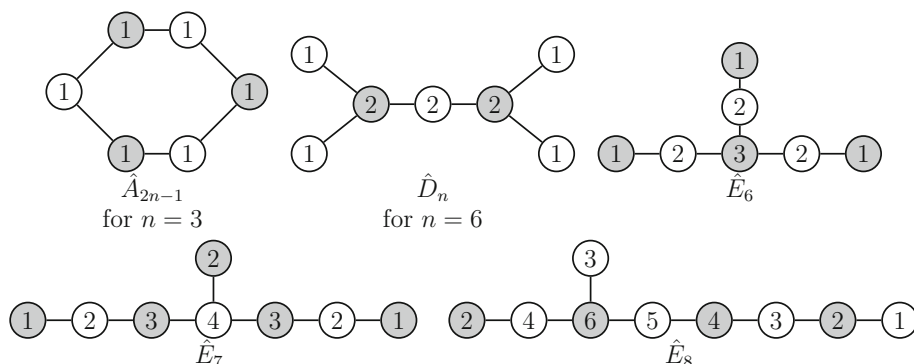


Fig. 2 Affine ADE Dynkin diagrams and their additive labelings

Table 1 Coxeter numbers of finite ADE Dynkin diagrams

Λ	A_n	D_m	E_6	E_7	E_8
$h(\Lambda)$	$n + 1$	$2m - 2$	12	18	30

For each finite ADE Dynkin diagram Λ there is an associated integer $h(\Lambda)$ called *Coxeter number*. We list Coxeter numbers of finite ADE Dynkin diagrams in Table 1. If $\hat{\Lambda}$ is an affine Dynkin diagram, we set $h(\hat{\Lambda}) = \infty$.

It is well known that the Coxeter number has a nice interpretation in terms of eigenvalues of the adjacency matrix:

Proposition 1.1.3 ([3, Exercise V.6.4])

- If Λ is a finite ADE Dynkin diagram then the dominant eigenvalue of its adjacency matrix equals $2 \cos(\pi/h(\Lambda))$;
- if $\hat{\Lambda}$ is an affine ADE Dynkin diagram then the dominant eigenvalue of its adjacency matrix equals 2.

□

In particular, the second claim justifies setting $h(\hat{\Lambda}) := \infty$.

1.1.3 Subadditive labelings

Let $G = (\Gamma, \Delta)$ be a bipartite bigraph on vertex set V . A *labeling* of its vertices is a function $\nu : V \rightarrow \mathbb{R}_{>0}$, which assigns to each vertex v of G a positive real label $\nu(v)$.

Definition 1.1.4 A labeling $\nu : V \rightarrow \mathbb{R}_{>0}$ is called

- *strictly subadditive* if for any vertex $v \in V$,

$$2\nu(v) > \sum_{(u,v) \in \Gamma} \nu(u), \text{ and } 2\nu(v) > \sum_{(v,w) \in \Delta} \nu(w).$$

- *subadditive* if for any vertex $v \in V$,

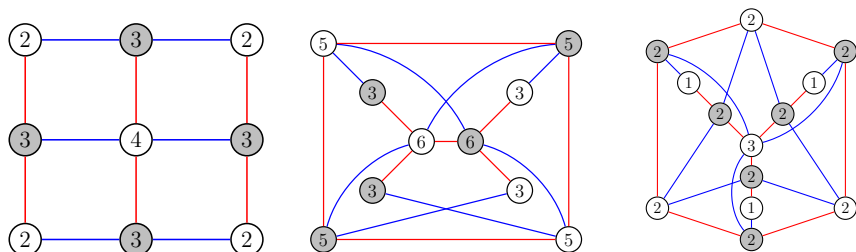
$$2\nu(v) \geq \sum_{(u,v) \in \Gamma} \nu(u), \text{ and } 2\nu(v) > \sum_{(v,w) \in \Delta} \nu(w).$$

- *weakly subadditive* if for any vertex $v \in V$,

$$2\nu(v) \geq \sum_{(u,v) \in \Gamma} \nu(u), \text{ and } 2\nu(v) \geq \sum_{(v,w) \in \Delta} \nu(w).$$

Examples of each type can be found in Fig. 3.

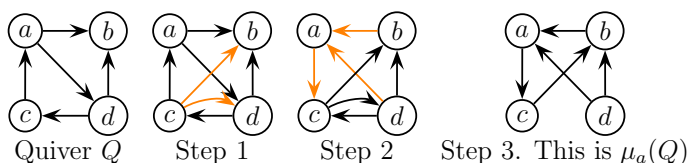
Strictly subadditive, subadditive and weakly subadditive labelings of quivers have been introduced in [30]. The terminology is motivated by Vinberg's *subadditive labelings* [40] for undirected graphs (see Theorem 1.1.2).



a strictly subadditive labeling

a subadditive labeling

a weakly subadditive labeling

Fig. 3 Different kinds of labelings

Fig. 4 Mutating a quiver at vertex a

1.1.4 Quivers

Definition 1.1.5 For a vertex v of Q one can define the *quiver mutation* μ_v at v as follows:

- (1) for each pair of edges $u \rightarrow v$ and $v \rightarrow w$ create an edge $u \rightarrow w$;
- (2) reverse the direction of all edges adjacent to v ;
- (3) if some directed 2-cycle is present, remove both of its edges; repeat until there are no more directed 2-cycles.

Let us denote the resulting quiver $\mu_v(Q)$. See Fig. 4 for an example of each step.

Now, let Q be a bipartite quiver. Recall that $\mu_\circ$ (resp., $\mu_\bullet$) is the simultaneous mutation at all white (resp., all black) vertices of Q , and that Q is *recurrent* if $\mu_\circ(Q) = \mu_\bullet(Q) = Q^{\text{op}}$.

As we have observed in [13], this property translates nicely into the language of bigraphs:

Corollary 1.1.6 A bipartite quiver Q is recurrent if and only if the associated bipartite bigraph $G(Q)$ has commuting adjacency matrices A_Γ, A_Δ .

We define three variations of Stembridge's admissible ADE bigraphs (see [36]):

Definition 1.1.7 Let $G = (\Gamma, \Delta)$ be a bipartite bigraph, and assume that the adjacency $|V| \times |V|$ matrices A_Γ and A_Δ of Γ and Δ commute. In this case we encode the three definitions in Table 2. For instance, G is an affine $\boxtimes$ finite ADE bigraph if each connected component of Γ is an affine ADE Dynkin diagram and each connected component of Δ is a finite ADE Dynkin diagram. We similarly define the notions of *admissible* and affine $\boxtimes$ affine ADE bigraphs.

The following fact is an easy consequence of [36, Lemma 4.3]:

Lemma 1.1.8 Let $G = (\Gamma, \Delta)$ be a bipartite bigraph and assume that the adjacency matrices A_Γ, A_Δ commute. Then the dominant eigenvalues of all components of Γ are equal to the

Table 2 Three types of bigraphs

Notion	All components of Γ are	All components of Δ are
Admissible ADE bigraph	Finite ADE Dynkin diagrams	Finite ADE Dynkin diagrams
Affine $\boxtimes$ finite ADE bigraph	Affine ADE Dynkin diagrams	Finite ADE Dynkin diagrams
Affine $\boxtimes$ affine ADE bigraph	Affine ADE Dynkin diagrams	Affine ADE Dynkin diagrams

same value λ_Γ , and the dominant eigenvalues of all components of Δ are equal to the same value λ_Δ . Matrices A_Γ and A_Δ have a common dominant eigenvector $\mathbf{v}$ such that

$$A_\Gamma \mathbf{v} = \lambda_\Gamma \mathbf{v}; \quad A_\Delta \mathbf{v} = \lambda_\Delta \mathbf{v}.$$

Corollary 1.1.9 *Let $G = (\Gamma, \Delta)$ be a bipartite bigraph and assume that the adjacency matrices A_Γ, A_Δ commute, and assume that all connected components of Γ and of Δ are either affine or finite ADE Dynkin diagrams. Then all connected components of Γ have the same Coxeter number denoted $h(\Gamma)$, and all connected components of Δ have the same Coxeter number denoted $h(\Delta)$.*

Combining Lemma 1.1.8, Definitions 1.1.7, 1.1.4, Vinberg's characterization (Theorem 1.1.2), and Proposition 1.1.3, we get the following proposition, whose part (1) was shown in [13, Proposition 5.1]. The proof for parts (2) and (3) is completely analogous and we refer the reader to [13] for details.

Proposition 1.1.10 *Let Q be a bipartite recurrent quiver Q and $G(Q) = (\Gamma, \Delta)$ be the corresponding bipartite bigraph. Then*

- (1) *Q admits a strictly subadditive labeling if and only if $G(Q)$ is an admissible ADE bigraph;*
- (2) *Q admits a subadditive labeling which is not strictly subadditive if and only if $G(Q)$ is an affine $\boxtimes$ finite ADE bigraph;*
- (3) *Q admits a weakly subadditive labeling which is not subadditive if and only if $G(Q)$ is an affine $\boxtimes$ affine ADE bigraph. $\square$*

1.2 Zamolodchikov integrable quivers admit weakly subadditive labelings

Recall that a bipartite recurrent quiver Q is called *Zamolodchikov integrable* if for every vertex $v \in \text{Vert}(Q)$, there exists an integer N and rational functions $J_0, \dots, J_N \in \mathbb{Q}(\mathbf{x})$ such that $J_0, J_N \neq 0$ and

$$\sum_{j=0}^N J_j T_v(t + 2j) = 0$$

for all $t \in \mathbb{Z}$ with $t + \epsilon_v$ even.

The following lemma is the first step towards the proof of Theorem 1.2.5.

Lemma 1.2.1 *If a bipartite recurrent quiver Q is Zamolodchikov integrable then Q admits a weakly subadditive labeling.*

Proof For $v \in \text{Vert}(Q)$, $t \in \mathbb{Z}$, define a positive number $a(v, t) := T_v(2t + \epsilon_v) \mid_{\mathbf{x}:=1}$ to be the value of $T_v(2t + \epsilon_v)$ if one substitutes $x_u := 1$ for all $u \in \text{Vert}(Q)$. By the Laurent Phenomenon (see [9]), the numbers $a(v, t)$ are integers which are positive since $T_v(2t + \epsilon_v)$ can be written as a subtraction-free expression in $\mathbf{x}$. Note that, unlike $T_v(t)$, the numbers $a(v, t)$ are defined for all v, t , regardless of parity.

Since $a(v, t)$ is always a positive integer, it is easy to see that the sequences $a(v, *) := (a(v, t))_{t \in \mathbb{Z}}$ are either simultaneously bounded or simultaneously unbounded (for all v). Assume for the sake of contradiction that for some vertex v , the sequence $a(v, *)$ is unbounded, but there is another vertex u for which the sequence $a(u, *)$ is bounded, say, $a(u, t) < C$ for all $t \in \mathbb{Z}$. We may assume that Q is connected and that u and v are neighbors in Q . Let t be such that $a(v, t) > C^2$. Then by the definition of the T -system, we have

$$a(u, t+1) > \frac{a(v, t)}{a(u, t)} > C,$$

where the first inequality uses the fact that all the numbers involved are positive integers, hence each of them is at least 1. This leads to an immediate contradiction.

If all the sequences are simultaneously bounded then they are periodic with the same period. This implies that the T -system associated with Q is periodic for *any* initial data, see [13, Remark 7.2]. In particular, such Q admits a strictly subadditive labeling by [13, Theorem 1.10]. Thus the only case left for us to consider is when the sequence $a(v, *)$ is unbounded for every v .

We need to show that if Q is Zamolodchikov integrable then Q admits a weakly subadditive labeling. The way to find such a labeling is going to be very similar to the proof of [1, Theorem 1].

The fact that Q is Zamolodchikov integrable implies that for each v , the sequences $a(v, *)$ satisfy a linear recurrence. Knowing that each of them is unbounded suggests using [1, Lemma 1] that describes the asymptotic behavior of sequences $a(v, *)$. Before we state it, let us denote $A(k) \approx B(k)$ for two functions of k if their ratio tends to a positive constant as $k \rightarrow \infty$.

Lemma 1.2.2 (see [1, Lemma 1]) *Let $a(v, *)$ be an unbounded sequence of positive integers satisfying a linear recurrence for each $v \in \text{Vert}(Q)$. Then there exist:*

- an integer $p \geq 1$;
- real numbers $\lambda(v, l) \geq 1$ for each $v \in \text{Vert}(Q)$, $l = 0, \dots, p$;
- integers $d(v, l) \geq 0$ for each $v \in \text{Vert}(Q)$, $l = 0, \dots, p$;
- a strictly increasing sequence $(n_k)_{k \in \mathbb{Z}_{\geq 0}}$ of nonnegative integers

such that the following things hold:

- (1) for every $v \in \text{Vert}(Q)$ and every $l = 0, \dots, p$, $a(v, pn_k + l) \approx \lambda(v, l) n_k^{d(v, l)}$;
- (2) for every $v \in \text{Vert}(Q)$ there exists $l = 0, \dots, p$ such that $\lambda(v, l) > 1$ or $d(v, l) \geq 1$;
- (3) for every $v \in \text{Vert}(Q)$ we have $\lambda(v, 0) = \lambda(v, p)$ and $d(v, 0) = d(v, p)$.

□

Clearly, the sequences $a(v, *)$ satisfy all the requirements of Lemma 1.2.2. For each $v \in \text{Vert}(Q)$, define

$$\lambda(v) := \prod_{l=0}^{p-1} \lambda(v, l) \in \mathbb{R}_{\geq 1}; \quad d(v) := \sum_{l=0}^{p-1} d(v, l) \in \mathbb{Z}_{\geq 0}.$$

For all $v \in \text{Vert}(Q)$ and $t \in \mathbb{Z}$, define

$$b(v, t) := \prod_{l=0}^{p-1} a(v, t + l).$$

Applying Lemma 1.2.2 yields

$$b(v, pn_k) \approx \lambda(v)^{n_k} n_k^{d(v)}. \quad (1.2.1)$$

By property (3) of Lemma 1.2.2, we have $a(v, pn_k) \approx a(v, pn_k + p)$ for every $v \in \text{Vert}(Q)$ and thus we can write

$$\begin{aligned} b(v, pn_k)^2 &\approx \prod_{l=0}^{p-1} a(v, pn_k + l) a(v, pn_k + l + 1) \\ &= \prod_{l=0}^{p-1} T_v(2(pn_k + l) + \epsilon_v) T_v(2(pn_k + l + 1) + \epsilon_v) \\ &= \prod_{l=0}^{p-1} \left(\prod_{u \rightarrow v} T_u(2(pn_k + l) + \epsilon_v + 1) + \prod_{v \rightarrow w} T_w(2(pn_k + l) + \epsilon_v + 1) \right) \\ &\geq \left(\prod_{u \rightarrow v} \prod_{l=0}^{p-1} T_u(2(pn_k + l) + \epsilon_v + 1) \right) \\ &\quad + \left(\prod_{v \rightarrow w} \prod_{l=0}^{p-1} T_w(2(pn_k + l) + \epsilon_v + 1) \right) \\ &\approx \left(\prod_{u \rightarrow v} b(u, pn_k) \right) + \left(\prod_{v \rightarrow w} b(w, pn_k) \right). \end{aligned}$$

The last equality is justified as follows: if $\epsilon_v = 0$ then $T_u(2(pn_k + l) + \epsilon_v + 1)$ is indeed equal to $a(u, pn_k + l)$. If $\epsilon_v = 1$ then $T_u(2(pn_k + l) + \epsilon_v + 1) = a(u, pn_k + l + 1)$ and then we again use $a(u, pn_k) \approx a(u, pn_k + p)$ in order to get to the last line.

By analyzing the asymptotics (1.2.1) of $b(v, pn_k)$, we see that for all $v \in \text{Vert}(Q)$,

$$\lambda(v)^2 \geq \max \left(\prod_{u \rightarrow v} \lambda(u), \prod_{v \rightarrow w} \lambda(w) \right); \quad d(v) \geq \max \left(\sum_{u \rightarrow v} d(u), \sum_{v \rightarrow w} d(w) \right). \quad (1.2.2)$$

Note that $\log \lambda(v) \geq 0$ and $d(v) \geq 0$ for all $v \in \text{Vert}(Q)$. Define $v(v) := \log \lambda(v) + d(v)$. By property (2) of Lemma 1.2.2, $v(v) > 0$ for all $v \in \text{Vert}(Q)$. By (1.2.2), v is a weakly subadditive labeling of $\text{Vert}(Q)$. $\square$

Applying Proposition 1.1.10, we get

Corollary 1.2.3 *If a bipartite recurrent quiver Q is Zamolodchikov integrable then $G(Q)$ is either*

- (1) *an admissible ADE bigraph, or*
- (2) *an affine $\boxtimes$ finite ADE bigraph, or*
- (3) *an affine $\boxtimes$ affine ADE bigraph.*

$\square$

Remark 1.2.4 We have shown in [13] that case (1) of the above corollary holds if and only if Q is *Zamolodchikov periodic*, that is, the T -system associated with Q is periodic. Obviously, this is a special case of Zamolodchikov integrability.

In fact, only cases (1) and (2) of Corollary 1.2.3 are possible when Q is Zamolodchikov integrable:

Theorem 1.2.5 *If a bipartite recurrent quiver Q is Zamolodchikov integrable then $G(Q)$ is either*

- (1) *an admissible ADE bigraph, or*
- (2) *an affine $\boxtimes$ finite ADE bigraph.*

We postpone the proof of this theorem until Sect. 4.2.

Part 2. The classification of affine $\boxtimes$ finite ADE bigraphs

Each affine ADE Dynkin diagram $\hat{\Lambda}$ has the associated dominant eigenvector $\mathbf{v}_{\hat{\Lambda}} : \text{Vert}(\hat{\Lambda}) \rightarrow \mathbb{R}$ corresponding to the eigenvalue 2. In other words, for every $v \in \hat{\Lambda}$ we have

$$2\mathbf{v}_{\hat{\Lambda}}(v) = \sum_{(v,w) \in \text{Edges}(\hat{\Lambda})} \mathbf{v}_{\hat{\Lambda}}(w).$$

We normalize $\mathbf{v}_{\hat{\Lambda}}$ so that its entries are positive integers with the smallest entry equal to 1. The values of $\mathbf{v}_{\hat{\Lambda}}$ are given in Fig. 2.

2.1 Self and double bindings

In this section, we classify all the bipartite affine $\boxtimes$ finite ADE bigraphs $G = (\Gamma, \Delta)$ such that Γ has either one or two connected components. If Γ has just one connected component then G is called a *self binding*, and if Γ has two connected components then G is called a *double binding*. We start with self bindings.

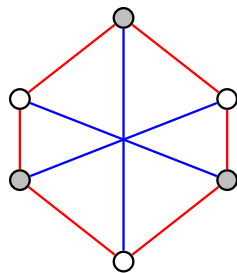
Throughout this section we assume that $h(\Delta) > 2$, i.e. that Δ has at least one edge (because if $h(\Delta) = 2$ then all connected components of Δ are of type A_1).

2.1.1 Self bindings

Lemma 2.1.1 *If $G = (\Gamma, \Delta)$ is a self binding then all the connected components of Δ are of type A_2 .*

Proof Let $\mathbf{v} : \text{Vert}(G) \rightarrow \mathbb{R}$ be the common eigenvector for A_{Γ} and A_{Δ} from Lemma 1.1.8. Thus $A_{\Gamma}\mathbf{v} = 2\mathbf{v}$. Since Γ has just one connected component, we may rescale $\mathbf{v}$ so that it is equal to $\mathbf{v}_{\Gamma}$. Now, let $\lambda_{\Delta} := 2 \cos(\pi/h(\Delta))$ be the dominant eigenvalue for A_{Δ} . We have that for every $v \in \text{Vert}(G)$,

$$\sum_{(v,w) \in \Delta} \mathbf{v}(w) = \lambda_{\Delta} \mathbf{v}(v).$$

Fig. 5 Self binding $\mathcal{S}_5$ 

Since there exists a vertex v for which $\mathbf{v}(v) = 1$, it follows that λ_Δ is an integer. This can only happen when $h(\Delta) = 3$, that is, when all the connected components of Δ have Coxeter number 3. The only finite ADE Dynkin diagram with Coxeter number 3 is A_2 . $\square$

Proposition 2.1.2 • For every $n \geq 1$, there is a self binding $\mathcal{S}_{4n+1} = (\Gamma_n, \Delta_n)$ where Γ_n is an affine ADE Dynkin diagram of type $\hat{A}_{4n+1}$, that is, a single cycle with $4n+2$ vertices, and two vertices of Γ_n are connected by an edge of Δ_n iff they are the opposite vertices of that cycle (see Fig. 5);

- There are no other self bindings.

Proof Let $G = (\Gamma, \Delta)$ be a self binding. By Lemma 2.1.1, all the components of Δ are just isolated single edges. Let us define an involution $\mathbf{i} : \text{Vert}(G) \rightarrow \text{Vert}(G)$ such that v and $\mathbf{i}(v)$ are exactly the vertices connected by the edges of Δ . This is a fixed point free involution, otherwise Δ would have a connected component of type A_1 . Moreover, since G is bipartite, $\mathbf{i}$ should reverse the colors of vertices. Finally, if $(u, v) \in \Gamma$ then one must also have $(\mathbf{i}(u), \mathbf{i}(v)) \in \Gamma$ because otherwise the adjacency matrices A_Γ and A_Δ would not commute. Thus $\mathbf{i}$ is a color-reversing involutive automorphism of G without fixed points. The only affine ADE Dynkin diagram admitting such an automorphism is $\hat{A}_{4n+1}$ for $n \geq 1$, where the automorphism is just a rotation by 180° . $\square$

2.1.2 Double bindings: scaling factor

The classification of double bindings is going to be much richer than that of self bindings. Throughout the rest of this section, we assume that $G = (\Gamma, \Delta)$ is a double binding, and that $\text{Vert}(G) = X \sqcup Y$, where X and Y are the two connected components of Γ , and recall that they are affine ADE Dynkin diagrams. A *parallel binding* is a bigraph of type $\hat{\Lambda} \otimes A_2$ and, following [36], is denoted $\hat{\Lambda} \equiv \hat{\Lambda}$.

Definition 2.1.3 The *scaling factor* of G (denoted $\text{scf}(G)$) is the number λ_Δ^2 where $\lambda_\Delta = 2 \cos(\pi/h(\Delta))$ is the dominant eigenvalue for A_Δ .

Proposition 2.1.4 The scaling factor $\text{scf}(G)$ is an integer equal to either 1, 2, or 3. Moreover,

- (1) if $\text{scf}(G) = 1$ then all connected components of Δ are of type A_2 ;
- (2) if $\text{scf}(G) = 2$ then all connected components of Δ are of type A_3 ;
- (3) if $\text{scf}(G) = 3$ then all connected components of Δ are either of type A_5 or of type D_4 .

Proof We view maps $\tau : \text{Vert}(G) \rightarrow \mathbb{R}$ as pairs $\begin{pmatrix} \tau_X \\ \tau_Y \end{pmatrix}$ where $\tau_X : \text{Vert}(X) \rightarrow \mathbb{R}$ and $\tau_Y : \text{Vert}(Y) \rightarrow \mathbb{R}$ are restrictions of τ to the corresponding subsets. Let $\tau = \begin{pmatrix} \tau_X \\ \tau_Y \end{pmatrix}$ be the

common dominant eigenvector for A_Γ and A_Δ from Lemma 1.1.8. We may rescale it so that $\tau_X = \alpha \mathbf{v}_X$ and $\tau_Y = \mathbf{v}_Y$ for some $\alpha \in \mathbb{R}$. Since the entries of the dominant eigenvector are positive, we may assume $\alpha > 0$. Since $A_\Delta \tau = \lambda_\Delta \tau$, we have

$$\sum_{(v,w) \in \Delta} \mathbf{v}_Y(w) = \lambda_\Delta \alpha \mathbf{v}_X(v), \quad \forall v \in X; \quad (2.1.1)$$

$$\sum_{(v,w) \in \Delta} \alpha \mathbf{v}_X(v) = \lambda_\Delta \mathbf{v}_Y(w), \quad \forall w \in Y. \quad (2.1.2)$$

If we substitute $v \in X$ such that $\mathbf{v}_X(v) = 1$ in (2.1.1), we will get that $\lambda_\Delta \alpha \in \mathbb{Z}_{>0}$. Similarly, if we substitute $w \in Y$ such that $\mathbf{v}_Y(w) = 1$ in (2.1.2), we will get that $\lambda_\Delta / \alpha \in \mathbb{Z}_{>0}$. Therefore their product λ_Δ^2 belongs to $\mathbb{Z}_{>0}$ as well. A straightforward case analysis shows that this can only happen when $h(\Delta) = 3, 4$, or 6 , and the result follows. $\square$

A simple consequence of the proof is the following observation:

Corollary 2.1.5 *Up to switching X and Y , we have:*

$$\sum_{(v,w) \in \Delta} \mathbf{v}_Y(w) = \text{scf}(G) \mathbf{v}_X(v), \quad \forall v \in X; \quad (2.1.3)$$

$$\sum_{(v,w) \in \Delta} \mathbf{v}_X(v) = \mathbf{v}_Y(w), \quad \forall w \in Y. \quad (2.1.4)$$

Proof We know that $\lambda_\Delta^2 \in \{1, 2, 3\}$ and thus $\lambda_\Delta \in \{1, \sqrt{2}, \sqrt{3}\}$. Thus the only $\alpha \in \mathbb{R}$ satisfying $\lambda_\Delta / \alpha \in \mathbb{Z}_{>0}$ and $\lambda_\Delta \alpha \in \mathbb{Z}_{>0}$ is either $\alpha = \lambda_\Delta$ or $\alpha = 1/\lambda_\Delta$. $\square$

By the same reasoning as in the proof of Proposition 2.1.2, if $\text{scf}(G) = 1$ then G is a parallel binding. It remains to classify double bindings with scaling factor 2 and 3. We say that a double binding is *nontrivial* if it is not a parallel binding, i.e. if the scaling factor is 2 or 3.

Definition 2.1.6 When X is an affine ADE Dynkin diagram of type $\hat{\Lambda}$ and Y is an affine ADE Dynkin diagram of type $\hat{\Lambda}'$ then we say that G is a double binding of type $\hat{\Lambda} * \hat{\Lambda}'$.

Note that Corollary 2.1.5 is not symmetric in X and Y , so if G is a double binding of type $\hat{\Lambda} * \hat{\Lambda}'$ then necessarily X has type $\hat{\Lambda}$, Y has type $\hat{\Lambda}'$ and (2.1.3) and (2.1.4) hold. In other words, we treat double bindings of types $\hat{\Lambda} * \hat{\Lambda}'$ and $\hat{\Lambda}' * \hat{\Lambda}$ differently.

A simple consequence of (2.1.4) is

Corollary 2.1.7 *For any double binding G , the maximal value of $\mathbf{v}_X$ is less than or equal to the maximal value of $\mathbf{v}_Y$.*

Proof Let $\mathbf{v}_X(u)$ be the maximal value of $\mathbf{v}_X$, then clearly $\mathbf{v}_Y(w) \geq \mathbf{v}_X(u)$ for any $(u, w) \in \Delta$ by (2.1.4). $\square$

For $p \in \mathbb{Z}$, denote by $\mathbf{v}_Y^{-1}(p)$ the set of vertices u of Y with $\mathbf{v}_Y(u) = p$.

Proposition 2.1.8 *There are no non-trivial double bindings of type $\hat{\Lambda} * \hat{\Lambda}$ (i.e. when X and Y have the same type).*

Proof Let M be the maximal value of $\mathbf{v}_X$ and $\mathbf{v}_Y$, and let $W = \mathbf{v}_X^{-1}(M)$, $U = \mathbf{v}_Y^{-1}(M)$ be the sets of vertices where $\mathbf{v}_X$ (resp., $\mathbf{v}_Y$) takes the maximal value. It is clear from (2.1.4) that every vertex from U is Δ -connected to at most one vertex from W . By the same reason, every vertex from $\text{Vert}(Y) \setminus U$ is not Δ -connected to any vertex from W . Thus every vertex from W is allowed to be Δ -connected only to vertices from U , and by (2.1.3), each of them should be connected to at least two vertices in U . We get a contradiction since the sizes of W and U are supposed to be the same. $\square$

2.1.3 Double bindings involving type $\hat{E}$

We say that Y is *one-two-bipartite* if for every $u, w \in \text{Vert}(Y)$ with $\mathbf{v}_Y(u) = 1$ and $\mathbf{v}_Y(w) = 2$, we have $\epsilon_u \neq \epsilon_w$ (that is, all ones in $\mathbf{v}_Y$ are white and all twos in $\mathbf{v}_Y$ are black, or vice versa). Note that if Y is of type $\hat{E}_6$ or $\hat{E}_8$ then Y is one-two-bipartite, see Fig. 2.

Lemma 2.1.9 *Let G be a double binding, and assume that Y is one-two-bipartite. Then $\text{scf}(G)$ divides $\#\mathbf{v}_Y^{-1}(1)$.*

Proof Let $w \in \mathbf{v}_Y^{-1}(1)$. By (2.1.4), there is exactly one vertex $v \in X$ with $(w, v) \in \Delta$, and moreover, $\mathbf{v}_X(v) = 1$. By (2.1.3),

$$\sum_{(v,u) \in \Delta} \mathbf{v}_Y(u) = \text{scf}(G).$$

Since $\text{scf}(G) \leq 3$ and $\mathbf{v}_Y(w) = 1$ is one of the terms in the left hand side, all the other terms in the left hand side are equal to either 1 or 2. But all vertices u with $(v, u) \in \Delta$ must be of the same color, since the graph is bipartite. The set of Δ -neighbors of v consists of exactly $\text{scf}(G)$ vertices u with $\mathbf{v}_Y(u) = 1$. By (2.1.4), v is the only Δ -neighbor of each such u . Therefore, the set $\mathbf{v}_Y^{-1}(1)$ is partitioned into classes, and each class has $\text{scf}(G)$ members that have the same Δ -neighbor. $\square$

Corollary 2.1.10 *If G is a non-trivial double binding of type $\hat{\Lambda} * \hat{E}_6$ then $\text{scf}(G) = 3$.*

Proposition 2.1.11 (1) *There are no non-trivial double bindings of type $\hat{\Lambda} * \hat{E}_8$;*
 (2) *the only non-trivial double binding of type $\hat{E}_n * \hat{\Lambda}$ is the double binding $\hat{E}_6 * \hat{E}_7$ depicted in Fig. 6.*

Proof To prove (1), just observe that if Y is of type $\hat{E}_8$ then $\#\mathbf{v}_Y^{-1}(1) = 1$ and apply Lemma 2.1.9.

To prove (2), we can first eliminate all the cases except for $\hat{E}_6 * \hat{E}_7$:

- by (1), there are no bindings of type $\hat{E}_n * \hat{E}_8$;
- by Proposition 2.1.8, there are no bindings of types $\hat{E}_6 * \hat{E}_6$ or $\hat{E}_7 * \hat{E}_7$;
- by Corollary 2.1.7, there are no bindings of types $\hat{E}_7 * \hat{E}_6$, $E_n * A_m$, or $E_n * D_m$.

Now we need to prove that there is only one double binding of type $\hat{E}_6 * \hat{E}_7$. Let $\{w_1, w_2, w_3\}$ be all the vertices of X (which is of type $\hat{E}_6$) with $\mathbf{v}_X(w_i) = 2$ for $i = 1, 2, 3$. Since Y is of type $\hat{E}_7$, it has 5, say, white vertices and 3 black vertices. Let $\{u_1, u_2, u_3\}$ be these three black vertices. Since w_1, w_2, w_3 are all of the same color, it is clear from (2.1.4) that they are white (because if the left hand side of (2.1.4) is even then the right hand side should be also even), and thus the other 4 vertices of X are black. To sum up, the edges of Δ connect the vertices u_1, u_2, u_3 to the vertices w_1, w_2, w_3 , and we have

$$\mathbf{v}_X(w_1) = \mathbf{v}_X(w_2) = \mathbf{v}_X(w_3) = \mathbf{v}_Y(u_1) = \mathbf{v}_Y(u_3) = 2, \quad \mathbf{v}_Y(u_2) = 4.$$

A simple case analysis shows that u_2 is Δ -connected to two vertices, say, to w_1 and w_2 while u_1 and u_3 are then both connected to w_3 . Now, using the fact that the adjacency matrices A_Γ and A_Δ commute, there is only one way to recover the rest of the double binding, and we get exactly $\hat{E}_6 * \hat{E}_7$ from Fig. 6. $\square$

We conclude the analysis of double bindings for which one of the components is of type $\hat{E}_n$ with the following proposition:

- Proposition 2.1.12** (1) *There are no non-trivial double bindings of type $\hat{A}_m * \hat{E}_n$;*
 (2) *there is exactly one non-trivial double binding of type $\hat{D}_m * \hat{E}_6$, namely, the binding $\hat{D}_4 * \hat{E}_6$ depicted in Fig. 7;*
 (3) *there is exactly one non-trivial double binding of type $\hat{D}_m * \hat{E}_7$, namely, the binding $\hat{D}_6 * \hat{E}_7$ depicted in Fig. 7.*

Proof First, we show (1). If X is of type $\hat{A}_m$ then $\mathbf{v}_X(w) = 1$ for all $w \in X$. If Y has type $\hat{E}_7$ then there is a vertex $u \in Y$ with $\mathbf{v}_Y(u) = 4$ which is impossible since u has at most three neighbors, so by (2.1.4), $\mathbf{v}_Y(u) \leq 3$. By Proposition 2.1.11, Y cannot be of type $\hat{E}_8$, so assume now that Y is of type $\hat{E}_6$. Let $\mathbf{v}_Y^{-1}(1) = \{u_1, u_2, u_3\}$. Then there is a vertex $w_1 \in X$ connected by Δ to each of u_1, u_2, u_3 . Let w_2 be such that $(w_2, w_1) \in \Gamma$. Since w_2 has different color than w_1 , it can only be Δ -connected to vertices $u \in Y$ with $\mathbf{v}_Y(u) = 2$. But the sum $\sum_{(u, w_2) \in \Delta} \mathbf{v}_Y(u)$ should be equal to 3 which is impossible because it is even. Thus (1) follows.

Next, we prove (2), so assume X has type $\hat{D}_m$ and Y has type $\hat{E}_6$. By Corollary 2.1.10, the scaling factor in this case equals to 3. Let $\mathbf{v}_Y^{-1}(1) = \{u_1, u_2, u_3\}$. Then each of u_1, u_2, u_3 is connected to some vertex $w_1 \in X$ with $\mathbf{v}_X(w_1) = 1$. Therefore w_1 has a unique Γ -neighbor $w_2 \in X$, and $\mathbf{v}_X(w_2) = 2$. Since the adjacency matrices A_Γ and A_Δ commute, w_2 should be connected to all three vertices u_4, u_5, u_6 of Y satisfying $\mathbf{v}_Y(u_i) = 2$ for $i = 4, 5, 6$. Since X has three more vertices w_3, w_4, w_5 with $\mathbf{v}_X(w_i) = 1$ for $i = 3, 4, 5$, each of them has to be connected to the remaining vertex u_7 of Y with $\mathbf{v}_Y(u_7) = 3$. It follows that there are no more vertices in X , so we are done with (2).

Finally, we show (3), so let X have type $\hat{D}_m$ and let Y have type $\hat{E}_7$. Assume first that the scaling factor is 2, and let $u \in \text{Vert}(Y)$ be a vertex with $\mathbf{v}_Y(u) = 3$. Then by (2.1.3), if $(u, w) \in \Delta$ for some $w \in \text{Vert}(X)$, then $\mathbf{v}_X(w) \geq 2$, but since X is of type $\hat{D}_m$, $\mathbf{v}_X(w)$ must be equal to 2. Since $\mathbf{v}_Y(u)$ is odd, this contradicts (2.1.4).

Thus the scaling factor has to be equal to 3. Because Y is of type $\hat{E}_7$, Y has 3, say, black vertices u_1, u_2, u_3 , and 5 white vertices, and we have $\mathbf{v}_Y(u_1) = \mathbf{v}_Y(u_3) = 2$, $\mathbf{v}_Y(u_2) = 4$. It follows now that:

- X has exactly 2 white vertices w_1 and w_2 ;
- one of the components of Δ has type A_5 and connects the vertices $u_1 - w_1 - u_2 - w_2 - u_3$.

Again, using commuting adjacency matrices, one can reconstruct the rest of the double binding and see that it is in fact $\hat{D}_6 * \hat{E}_7$ in Fig. 7. $\square$

2.1.4 Double bindings involving type $\hat{A}$

One can identify the vertices of the cycle A_{2m-1} with $\mathbb{Z}_m := \mathbb{Z}/m\mathbb{Z}$. We define double and triple coverings to be the following double bindings: in a *double covering* $\hat{A}_{2n-1} * \hat{A}_{4n-1}$, a vertex $j \in \mathbb{Z}_{4m}$ of Y is connected by a blue edge to a vertex $i \in \mathbb{Z}_{2m}$ of X iff $i \equiv j \pmod{2m}$. Similarly, in a *triple covering* $\hat{A}_{2n-1} * \hat{A}_{6n-1}$, a vertex $j \in \mathbb{Z}_{6m}$ of Y is connected

by a blue edge to a vertex $i \in \mathbb{Z}_{2m}$ of X iff $i \equiv j \pmod{2m}$. These are obviously affine $\boxtimes$ finite ADE bigraphs.

Proposition 2.1.13 *The only possible double bindings of type $\hat{A}_{m-1} * \hat{A}_{k-1}$ are:*

- (1) *parallel bindings $\hat{A}_{2n-1} \equiv \hat{A}_{2n-1}$;*
- (2) *double coverings $\hat{A}_{2n-1} * \hat{A}_{4n-1}$;*
- (3) *triple coverings $\hat{A}_{2n-1} * \hat{A}_{6n-1}$.*

Proof By (2.1.4), each vertex of Y has exactly one blue neighbor, and each vertex of X has exactly $\text{scf}(G)$ blue neighbors. Let $(w_i)_{i \in \mathbb{Z}_k}$ be the vertices of Y listed in cyclic order, and let $(v_i)_{i \in \mathbb{Z}_m}$ be the vertices of X in cyclic order. Let $f : \mathbb{Z}_k \rightarrow \mathbb{Z}_m$ be the map such that $v_{f(i)}$ is the unique blue neighbor of w_i for all $i \in \mathbb{Z}_k$. Since the adjacency matrices have to commute, we get that $\{f(i+1), f(i-1)\} = \{f(i)+1, f(i)-1\}$ which immediately yields the result of the proposition. $\square$

By Corollary 2.1.7, there are no double bindings of type $\hat{D}_m * \hat{A}_n$ so the only case left in this section is $\hat{A}_n * \hat{D}_m$.

Proposition 2.1.14 *The only possible double bindings of type $\hat{A}_n * \hat{D}_m$ are the double bindings of type $\hat{A}_{2n-1} * \hat{D}_{n+2}$ in Fig. 6 and the exceptional double binding of type $\hat{A}_3 * \hat{D}_5$ in Fig. 7.*

Proof We have two options: either $\text{scf}(G) = 2$ or $\text{scf}(G) = 3$. If $\text{scf}(G) = 2$ then we know that each non-leaf vertex of Y is connected to exactly two vertices of X , and is the only blue neighbor of each of them. On the other hand, there are two more vertices v_1, v_2 in X , and each of them has two blue neighbors which are leaves in Y . Now using commuting adjacency matrices condition one can easily recover that G is the double binding of type $\hat{A}_{2n-1} * \hat{D}_{n+2}$ from Fig. 6.

Now assume that $\text{scf}(G) = 3$. This means that each vertex of X is connected to an odd number of leaves of Y . Since Y has exactly four leaves, it follows that X has either two or four vertices. If X has two vertices then the sum of values of $\mathbf{v}_Y$ is six so Y has type $\hat{D}_4$ but then all the leaves of Y have the same color so one of the vertices of X is not going to be connected to any of them. We are left with the case when X has four vertices and each of them is connected to a leaf of Y and to a non-leaf of Y . Therefore Y has type $\hat{D}_5$ from which one can quickly see that G is the unique double binding of type $\hat{A}_3 * \hat{D}_5$ from Fig. 7. $\square$

2.1.5 Double bindings of type $\hat{D}_{m+1} * \hat{D}_{k+1}$

Proposition 2.1.15 *The only possible double bindings of type $\hat{D}_{m+1} * \hat{D}_{k+1}$ are the double bindings of type $\hat{D}_n * \hat{D}_{2n-2}$ and the double bindings of type $\hat{D}_{n+1} * \hat{D}_{3n-1}$ constructed in the proof of this proposition and depicted for small n in Fig. 7.*

Proof Let $v_1^+, v_1^-, v_2, \dots, v_{m-1}, v_m^+, v_m^-$ be the vertices of X , the component of type $\hat{D}_{m+1}$, and $w_1^+, w_1^-, w_2, \dots, w_{k-1}, w_k^+, w_k^-$ be the vertices of Y which has type $\hat{D}_{k+1}$. Here we assume that v_1^+, v_1^- are the leaves attached to v_2 and so on. By (2.1.4), each of $w_1^+, w_1^-, w_k^+, w_k^-$ is connected to exactly one leaf of X . Without loss of generality assume that w_1^+ is connected to v_1^+ by a blue edge. Since the adjacency matrices commute, w_2 has to be connected to v_2 by blue edges. We claim that w_1^- cannot be connected to v_1^+ . Indeed, otherwise there would be at least two blue–red paths and at most one red–blue path from v_1^+ to w_2 , so the matrices would not commute. On the other hand, w_1^- is connected to a leaf, and

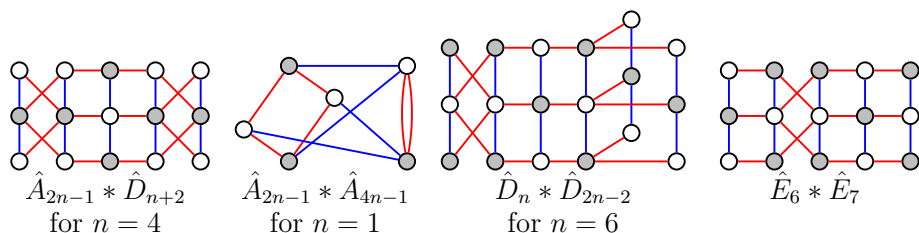


Fig. 6 Three infinite and one exceptional family of double bindings with scaling factor 2. All blue components have type A_3 (color figure online)

this leaf has to be a neighbor of v_2 . So without loss of generality we may assume that w_1^- is connected to v_1^- (we only make a choice here when Y has type $\hat{D}_4$ in which case all the four leaves of Y are connected to w_2).

By Proposition 2.1.8, we have $m \neq k$ and by (2.1.3–2.1.4) we actually have $m < k$. We claim that for each $i = 2, \dots, m-1$, w_i is connected to v_i , and thus to nothing else by (2.1.4). We show it by induction on i , where the base $i = 2$ has already been shown. Assume that w_i is connected to v_i . Then there is a red-blue path from v_{i+1} to w_i , and v_{i+1} is not connected to w_{i-1} so it has to be connected to w_{i+1} , and the claim follows for $i = 2, \dots, m-1$. Now there is a red-blue path from v_m^+ to w_{m-1} so by the same reasoning v_m^+ and v_m^- are connected to w_m . Now there is a red-blue path from w_{m+1} to v_m^+ and to v_m^- so w_{m+1} is connected to v_{m-1} . Again using induction we can show that for $i = 1, 2, \dots, m-2$, w_{m+i} is connected to v_{m-i} . This includes the fact that $2m-2 < k$ which is true since before we stop we need to add another blue edge to v_1^+ in order to satisfy (2.1.3).

If $\text{sfc}(G) = 2$ then (2.1.3) is satisfied for all vertices of X except for v_1^+ and v_1^- so we complete the construction of the graph by joining v_1^+ to w_k^+ and v_1^- to w_k^- , where necessarily $k = 2m-1$. This can be considered to be the definition of $\hat{D}_n * \hat{D}_{2n-2}$.

If $\text{sfc}(G) = 3$ then (2.1.3) is not satisfied for v_m^+ yet so we note that $2m-1 < k$ and thus have to connect both v_1^+ and v_1^- to w_{2m-1} . But then there is a red-blue path from w_{2m} to v_1^+ so w_{2m} has to be connected to v_2 . Now for $i = 1, 2, \dots, m-2$ it follows that w_{2m+i-1} is connected to v_{i+1} . After that (2.1.3) fails only for v_m^+ and v_m^- which we connect to w_k^+ and w_k^- respectively. Here k is necessarily equal to $3m-1$ yielding the double binding of type $\hat{D}_{n+1} * \hat{D}_{3n-1}$. $\square$

2.1.6 The classification of self and double bindings

We summarize the results of Sects. 2.1.1–2.1.5 in the following theorem:

Theorem 2.1.16 • *The only possible self bindings are S_{4n+1} for $n \geq 1$.*

- *all the double bindings with scaling factor 2 are listed in Fig. 6;*
- *all the double bindings with scaling factor 3 are listed in Fig. 7;*
- *the only other double bindings are parallel bindings $\hat{\Lambda} \equiv \hat{\Lambda}$.*

$\square$

Remark 2.1.17 In [13, Section 9.1], we introduced duality of symmetric bigraphs (not to be confused with Stembridge's dual bigraphs in [36]). Here we briefly list some pairs of dual symmetric bigraphs for certain choices of the auxiliary data¹ which we omit:

¹ i.e. $\mathbf{i}$, V_+ , V_0 , V_- , X in the notation of [13].

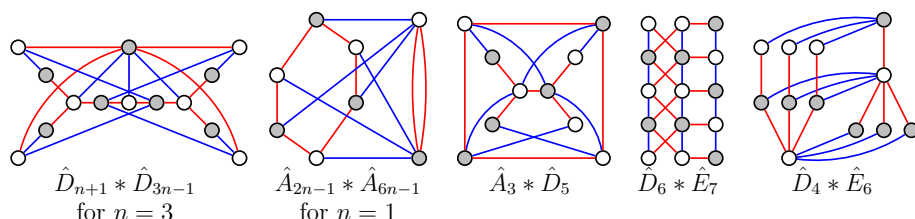


Fig. 7 Two infinite and three exceptional families of double bindings with scaling factor 3. All blue components have types A_5 or D_4 (color figure online)

- $\hat{A}_{2n-1} * \hat{D}_{n+2}$ is dual to $\hat{A}_{2n-1} \otimes A_3$;
- $\hat{A}_{2n-1} * \hat{A}_{4n-1}$ is dual to $\hat{A}_{4n-1} \otimes A_3$;
- $\hat{D}_n * \hat{D}_{2n-2}$ is dual to $\hat{A}_{2n-1} * \hat{A}_{4n-1}$;
- $\hat{E}_6 * \hat{E}_7$ is dual to itself;
- $\hat{D}_{n+1} * \hat{D}_{3n-1}$ is dual to $\hat{A}_{2n-1} * \hat{A}_{6n-1}$;
- $\hat{D}_6 * \hat{E}_7$ is dual to $\hat{D}_4 * \hat{E}_6$;
- $\hat{A}_3 * \hat{D}_5$ is dual to the triple covering $\hat{A}_1 * \hat{A}_5$.

2.2 The classification

To classify affine $\boxtimes$ finite ADE bigraphs, we mostly follow the strategy of [36]: we are going to show that the *component graph* $\mathcal{C}$ of Γ defined below is a path with either at most one loop (in case there is a self binding) or at most one non-parallel double binding.

Definition 2.2.1 Let $G = (\Gamma, \Delta)$ be a bigraph. Let $C_1, C_2, \dots, C_m$ be the connected components of Γ . Define the graph $\mathcal{C} = \mathcal{C}(G)$ with vertex set $[m] := \{1, 2, \dots, m\}$ such that (i, j) is an edge of $\mathcal{C}$ iff there is a blue edge $(u, v) \in \Delta$ with $u \in C_i$ and $v \in C_j$.

Let G be an affine $\boxtimes$ finite ADE bigraph. We define its reduced version $\tilde{G}$ to be the same as G but with all the blue edges removed from each self binding in G . Clearly, $\mathcal{C}(\tilde{G})$ is $\mathcal{C}(G)$ with all the loops removed. It is also clear that $\tilde{G}$ is going to be an affine $\boxtimes$ finite ADE bigraph as well.

Several properties of affine $\boxtimes$ finite ADE bigraphs have literally the same statements as their analogs for admissible ADE bigraphs of [36], so we list them with the corresponding references to the parts of [36] where they are proved:

Lemma 2.2.2 *Let G be an affine $\boxtimes$ finite ADE bigraph. Then:*

- (1) *the component graph $\mathcal{C}(\tilde{G})$ is acyclic (see [36, proof of Lemma 2.5(b)]);*
- (2) *in fact, $\mathcal{C}(\tilde{G})$ is a path (see [36, Proof of (ii) in Section 5]);*
- (3) *G contains at most one non-parallel double binding (see [36, Proof of (ii) in Section 5]).*

□

These properties allow us to describe every affine $\boxtimes$ finite ADE bigraph by a string of symbols $\hat{A}_n, \hat{D}_n, \hat{E}_n, \hat{S}_n$ with symbols $*, \equiv$ inserted between them, for example, $\hat{A}_1 \equiv \hat{A}_1 * \hat{A}_2$ has three red connected components (i.e. $m = 3$) and C_1 and C_2 form a parallel binding while C_2 and C_3 form a double binding of type $\hat{A}_1 * \hat{A}_2$.

Lemma 2.2.3 *Assume that G is an affine $\boxtimes$ finite ADE bigraph containing a self binding. Then it contains exactly one self binding and all the double bindings in G are parallel.*

Proof Assume for the sake of contradiction that G has at least two self bindings. We may remove everything else so that they occur at the ends of $\mathcal{C}(\tilde{G})$ (which is a path on $[m]$). After some relabeling, the edges of $\mathcal{C}(\tilde{G})$ become exactly $\{(i, i+1)\}_{i \in [m-1]}$. We are going to construct a blue cycle in G as follows: let $v_1^1 \in C_1$ be any vertex, then there is a blue path $v_1^1, v_2^1, \dots, v_m^1$ with $v_i^1 \in C_i$. Since C_m is a self binding, v_m^1 is connected by a blue edge to some other vertex $v_m^2 \in C_m$, from which we can construct a blue path $v_m^2, v_{m-1}^2, \dots, v_1^2$ with $v_i^2 \in C_i$ again. But now v_2^1 is connected by a blue edge to some other vertex $v_3^1 \in C_1$, so we may continue our path until it crosses itself yielding a blue cycle in G which is a contradiction since all the finite ADE Dynkin diagrams are acyclic.

Assume now that there is a self binding and a non-parallel double binding in G . Again, we may assume that the self binding occurs in C_1 and the double binding occurs between C_{m-1} and C_m with $\mathcal{C}(\tilde{G})$ being a path on $[m]$. Take the maximal blue path P in G . Since all the vertices in $C_1, \dots, C_{m-1}$ have blue degree at least 2, both endpoints of P belong to C_m and have blue degree 1. But since the blue components of the double binding $C_{m-1} * C_m$ are either A_3, A_5 , or D_4 (see Proposition 2.1.4), the vertices of P adjacent to the endpoints have blue degrees at least 3. Therefore they coincide because every finite ADE Dynkin diagram contains at most one vertex of degree 3. So P has length at most 3, and therefore $m = 2$. It is clear that adding a self binding to any of the double bindings involving type $\hat{A}$ yields either a cycle or a blue component with at least two vertices of degree 3. $\square$

Proposition 2.2.4 *The only affine $\boxtimes$ finite bigraphs involving self-bindings are*

$$(\hat{A}_{4n+1})^m := \mathcal{S}_{4n+1} \equiv \hat{A}_{4n+1} \equiv \dots \equiv \hat{A}_{4n+1} \quad (m \text{ factors}, m \geq 1, n \geq 1).$$

Proof If G is an affine $\boxtimes$ finite bigraph with a self binding then we know that $\mathcal{C}(\tilde{G})$ is a path by Lemma 2.2.2, so let $C_1, C_2, \dots, C_m$ be its connected components with the self binding happening in C_l for some $l \in [m]$. If $l \neq 1, m$ then we immediately get two vertices of degree 3 in every blue component, so we may assume that $l = 1$. By Lemma 2.2.3, all the double bindings are parallel and the result follows. $\square$

Theorem 2.2.5 *Let G be an affine $\boxtimes$ finite ADE bigraph. Then G is isomorphic to exactly one of the following bigraphs:*

- (1) $\hat{\Lambda} \otimes \Lambda'$ where $\hat{\Lambda}$ and Λ' are an affine and a finite ADE Dynkin diagram respectively;
- (2) $(\hat{A}_{4n+1})^m$ ($m \geq 1, n \geq 1$);
- (3) $(\hat{A} \hat{D}^{m-1})_n := \hat{A}_{2n-1} * \hat{D}_{n+2} \equiv \dots \equiv \hat{D}_{n+2}$ (m factors, $m \geq 2, n \geq 2$);
- (4) $(\hat{A}^{m-1} \hat{D})_n := \hat{A}_{2n-1} \equiv \dots \equiv \hat{A}_{2n-1} * \hat{D}_{n+2}$ (m factors, $m \geq 3, n \geq 2$);
- (5) $(\hat{A} \hat{A}^{m-1})_n := \hat{A}_{2n-1} * \hat{A}_{4n-1} \equiv \dots \equiv \hat{A}_{4n-1}$ (m factors, $m \geq 2, n \geq 1$);
- (6) $(\hat{A}^{m-1} \hat{A})_n := \hat{A}_{2n-1} \equiv \dots \equiv \hat{A}_{2n-1} * \hat{A}_{4n-1}$ (m factors, $m \geq 3, n \geq 1$);
- (7) $(\hat{D} \hat{D}^{m-1})_n := \hat{D}_n * \hat{D}_{2n-2} \equiv \dots \equiv \hat{D}_{2n-2}$ (m factors, $m \geq 2, n \geq 4$);
- (8) $(\hat{D}^{m-1} \hat{D})_n := \hat{D}_n \equiv \dots \equiv \hat{D}_n * \hat{D}_{2n-2}$ (m factors, $m \geq 3, n \geq 4$);
- (9) $\hat{E} \hat{E}^{m-1} := \hat{E}_6 * \hat{E}_7 \equiv \dots \equiv \hat{E}_7$ (m factors, $m \geq 2$);
- (10) $\hat{E}^{m-1} \hat{E} := \hat{E}_6 \equiv \dots \equiv \hat{E}_6 * \hat{E}_7$ (m factors, $m \geq 3$);
- (11) $\hat{A}_{2n-1} \equiv \hat{A}_{2n-1} * \hat{D}_{n+2} \equiv \hat{D}_{n+2}$ ($n \geq 2$);
- (12) $\hat{A}_{2n-1} \equiv \hat{A}_{2n-1} * \hat{A}_{4n-1} \equiv \hat{A}_{4n-1}$ ($n \geq 1$);
- (13) $\hat{D}_n \equiv \hat{D}_n * \hat{D}_{2n-2} \equiv \hat{D}_{2n-2}$ ($n \geq 4$);
- (14) $\hat{A}_{2n-1} * \hat{A}_{6n-1}$;

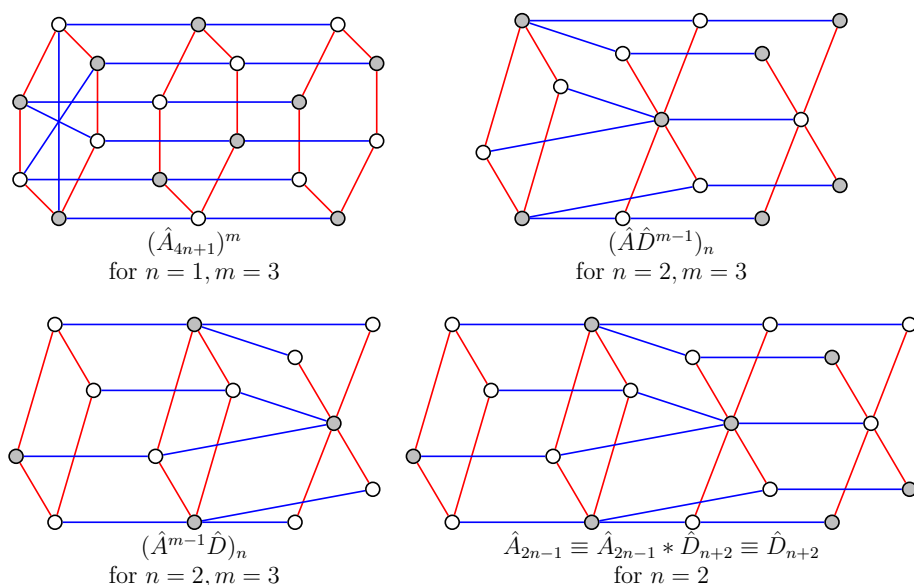


Fig. 8 Items (2), (3), (4), and (11) of our classification

- (15) $\hat{D}_{n+1} * \hat{D}_{3n-1}$;
- (16) $\hat{E}_6 \equiv \hat{E}_6 * \hat{E}_7 \equiv \hat{E}_7$;
- (17) $\hat{A}_3 * \hat{D}_5$;
- (18) $\hat{D}_6 * \hat{E}_7$;
- (19) $\hat{D}_4 * \hat{E}_6$.

Note that the infinite families are (1)–(15), so there are 15 infinite families and 4 exceptional bigraphs. Please see Fig. 8 for examples.

Proof By Proposition 2.2.4, we may assume that G has no self bindings. If all the double bindings in G are parallel then G is a tensor product. Otherwise consider the unique double binding $C_l * C_{l+1}$ of G . If it has scaling factor 3 then all of its components are of type either A_5 or D_4 by Proposition 2.1.4, so it is clear that adding an edge to all vertices of the same color in A_5 or in D_4 does not produce a finite ADE Dynkin diagram (in fact, it always produces an affine ADE Dynkin diagram). Therefore if the scaling factor is 3 then $m = 2$ and G is just the double binding itself. If the scaling factor is 2 then all the blue components are of type A_3 so obviously either $l = 1$ or $l = m - 1$ and m is arbitrary, or $l = 2$ and $m = 4$, and the theorem follows. $\square$

Part 3. T -systems of type $A \otimes \hat{A}$.

3.1 Variables via domino tilings

3.1.1 Speyer's formula

In [34] Speyer gives the following formula for variables of the octahedron recurrence. One can think of it as the T -system associated with the tensor product of type $A_\infty \otimes A_\infty$, where

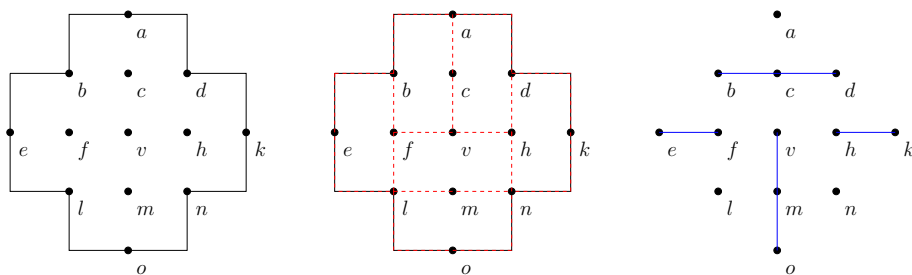


Fig. 9 An Aztec diamond $\mathcal{Z}_v(2)$, its domino tiling $\mathcal{D}$, and the associated graph $G_{\mathcal{D}}$

A_{∞} is the infinite path graph with vertex set $\mathbb{Z}$ and edge set $\{(i, i + 1) \mid i \in \mathbb{Z}\}$. The tensor product is understood in the sense of Definition 1.1.1, which works the same way for infinite graphs.

Let $\mathcal{Z}_v(t)$ be the Aztec diamond of radius t centered at vertex v . By abuse of notation, denote $\mathcal{Z}_v(t)$ also the set of vertices of $\mathcal{Z}_v(t)$ that are not its outer corners. Let $\mathcal{D}$ be a domino tiling of $\mathcal{Z}_v(t)$. Each domino has a *cut edge* that separates its two halves. Let $G_{\mathcal{D}}$ be the graph obtained by taking $\mathcal{Z}_v(t)$ as the set of vertices, and the set of all cut edges in $\mathcal{D}$ as the set of edges. For a vertex $u \in \mathcal{Z}_v(t)$ let $d_{\mathcal{D}}(u)$ be the degree of vertex u in graph $G_{\mathcal{D}}$. It is easy to see that each $d_{\mathcal{D}}(u)$ can take values 0, 1, 2 only.

Theorem 3.1.1 [34] *The formula for the variable $T_v(t)$ in an $A_{\infty} \otimes A_{\infty}$ T -system is as follows:*

$$T_v(2t + 1) = T_v(2t + 2) = \sum_{\mathcal{D}} \prod_{u \in \mathcal{Z}_v(2t+1)} u^{1-d_{\mathcal{D}}(u)},$$

$$T_v(-2t) = T_v(-2t - 1) = \sum_{\mathcal{D}} \prod_{u \in \mathcal{Z}_v(2t)} u^{1-d_{\mathcal{D}}(u)},$$

where $t \geq 0$ and the sum is taken over all domino tilings $\mathcal{D}$ of $\mathcal{Z}_v(2t + 1)$ and of $\mathcal{Z}_v(2t)$, respectively.

Example 3.1.2 In Fig. 9 we see an example of an Aztec diamond $\mathcal{Z}_v(2)$, its domino tiling $\mathcal{D}$, and the associated graph $G_{\mathcal{D}}$. The Laurent monomial this tiling contributes is $\frac{aln}{cm}$, which is easily seen to be one of the monomials in

$$T_v(-2) = \frac{\frac{ev+bl}{f} \frac{vk+dn}{h} + \frac{av+bd}{c} \frac{ln+vo}{m}}{v}$$

$$= \frac{evk}{fh} + \frac{edn}{fh} + \frac{blk}{fh} + \frac{bldn}{fhv} + \frac{aln}{cm} + \frac{avo}{cm} + \frac{bdo}{cm} + \frac{bdln}{cmv}.$$

Remark 3.1.3 Alternative approaches to giving explicit formulas for the octahedron recurrence can be found in the works of Di Francesco and Kedem [5–8] and Henriques [16]. We shall use Speyer’s language as the most convenient for our purposes.

3.1.2 Formula with cylindric boundary conditions

Consider now the case of T -system of type $A_m \otimes \hat{A}_{2n-1}$. The quiver is naturally embedded on a cylinder. Consider the lifting of the quiver to the universal cover of the cylinder, where

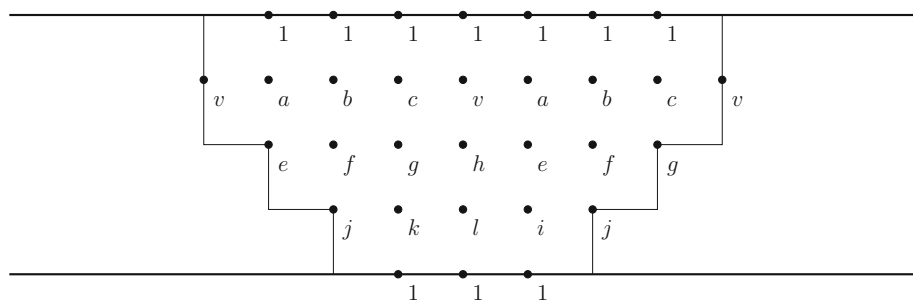


Fig. 10 An example of region $\mathcal{Z}_v(4)$ on the universal cover of a cylinder with $n = 2$ and $m = 3$

the vertex variables are periodic. We claim that the following variation of Speyer's theorem holds.

Let $\mathcal{Z}_v(t)$ now be the intersection of the Aztec diamond of radius t centered at vertex v with the universal cover of the cylinder, where we include two layers of frozen variables with values 1 on both boundaries. An example for $n = 2$ and $m = 3$ is shown in Fig. 10. For each domino tiling $\mathcal{D}$ of $\mathcal{Z}_v(t)$ define $G_{\mathcal{D}}$ and $d_{\mathcal{D}}(u)$ as before, but now using the periodicity of variables on the universal cover.

Theorem 3.1.4 *The formula for the variable $T_v(t)$ in an $A_m \otimes \hat{A}_{2n-1}$ T -system is as follows:*

$$T_v(2t+1) = T_v(2t+2) = \sum_{\mathcal{D}} \prod_{u \in \mathcal{Z}_v(2t+1)} u^{1-d_{\mathcal{D}}(u)},$$

$$T_v(-2t) = T_v(-2t-1) = \sum_{\mathcal{D}} \prod_{u \in \mathcal{Z}_v(2t)} u^{1-d_{\mathcal{D}}(u)},$$

where $t \geq 0$ and the sum is taken over all domino tilings $\mathcal{D}$ of $\mathcal{Z}_v(2t+1)$ and of $\mathcal{Z}_v(2t)$, respectively.

Proof We are going to apply Speyer's theorem to the $A_{\infty} \otimes A_{\infty}$ case with variables as shown in Fig. 11.

There are three logical steps to the proof. First, we claim that as we run the T -system dynamics, the Laurent monomials with the minimal power of ϵ remain the same in the vertices which carry variables $1, \epsilon, \epsilon^2, \dots$ at the beginning, while at the same time the minimal degree of ϵ in Laurent monomials in the rest of the vertices (i.e. the ones in the middle of the universal cover) is 0. Indeed, let us argue this by induction. Applying a mutation at a vertex with minimal Laurent monomial 1 (i.e. with value $1 + O(\epsilon)$), we see that the new value is

$$\frac{(1 + O(\epsilon))(1 + O(\epsilon)) + (\epsilon + O(\epsilon^2))(O(1) + O(\epsilon))}{1 + O(\epsilon)},$$

where $O(\epsilon^k)$ denotes terms with ϵ -degree at least k . It is clear then that specializing at $\epsilon = 0$ we get 1, which must be then the Laurent monomial with the smallest degree of ϵ in the result. A similar argument applies in other locations carrying a power of ϵ at the beginning.

Next, we claim that plugging in $\epsilon = 0$ into the formulas for the T -system of type $A_{\infty} \otimes A_{\infty}$ constructed as above returns exactly the formulas for T -system of type $A_m \otimes \hat{A}_{2n-1}$. Again, we can argue this by induction. At the very beginning the claim is obvious. The step is also easy to see from the first claim above. This is because by induction assumption the exchange

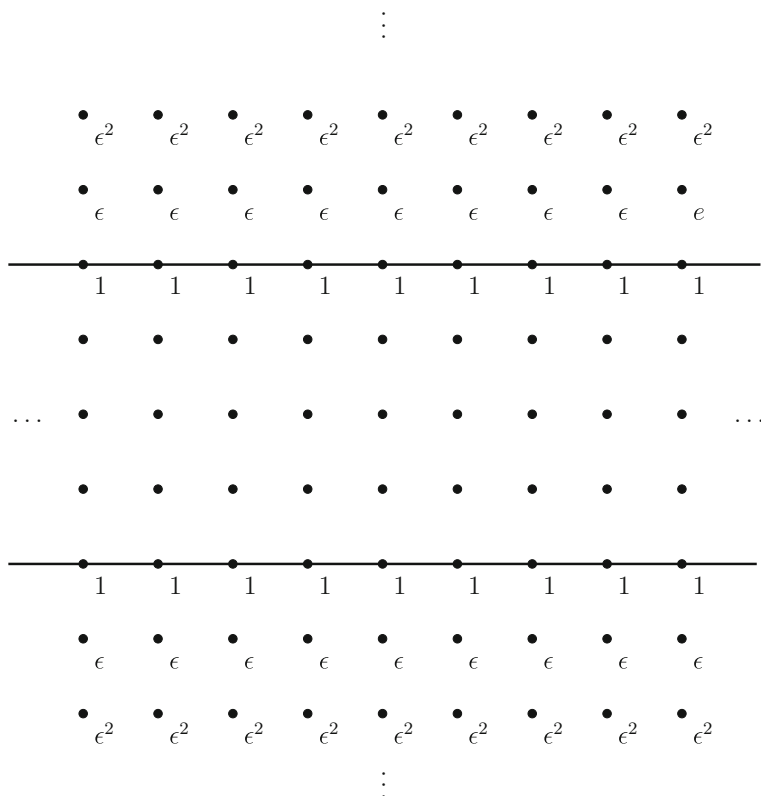


Fig. 11 An assignment of variables outside of the universal cover

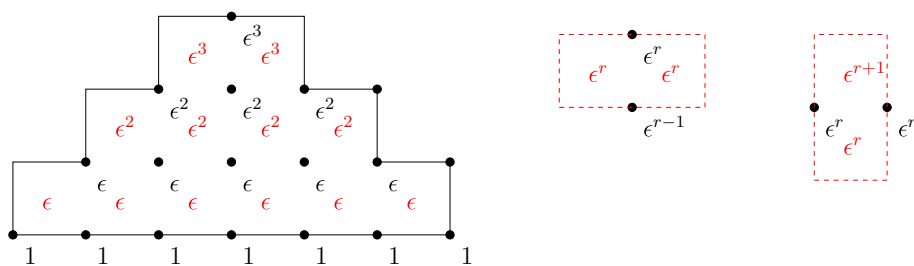


Fig. 12 An assignment of variables outside of the universal cover

relations for $A_\infty \otimes A_\infty$ T -system specialize to exchange relations for $A_m \otimes \hat{A}_{2n-1}$ T -system, and all the powers of ϵ involved are non-negative.

Finally, we want to argue that Speyer's formula applied to the above $A_\infty \otimes A_\infty$ case and specialized at $\epsilon = 0$ indeed returns the formula stated in the theorem. For that, we claim that in order for a domino tiling $\mathcal{D}$ to contribute a term with degree of ϵ equal 0 (i.e. a term which will not die after specializing) the chunks of Aztec diamonds $\mathcal{Z}_v(t)$ that are outside of the universal cover need to be tiled with horizontal tiles only. Such $\mathcal{D}$ -s are then in bijection with the tilings of the part of $\mathcal{Z}_v(t)$ that is inside the universal cover strip, as desired.

Let us look at a chunk of $\mathcal{Z}_v(t)$ that falls outside of the universal cover. Give each potential domino square weight ϵ^r equal to the larger weight a vertex adjacent to this square has. Considering both ways a domino can be positioned, it is clear that the weight picked up by the corresponding edge in $G_{\mathcal{D}}$ is equal to the weight of its squares minus one (see Fig. 12):

$$r + (r - 1) = r + r - 1 \text{ and } r + r = (r + 1) + r - 1.$$

From this it is easy to see that the dominos lying in this chunk can pick up maximal weight of at most the weight of all squares minus potential number of dominos, which is

$$2 \cdot R + 4 \cdot (R - 1) + \cdots + 2R \cdot 1 - R(R + 1)/2,$$

where ϵ^R is the maximal power of ϵ in the chunk. On the other hand, the total weight to burn in the chunk is

$$1 \cdot R + 3 \cdot (R - 1) + \cdots + (2R - 1) \cdot 1,$$

which is easily seen to be the same. Thus, in order for the ϵ to not enter the resulting overall weight picked up by $G_{\mathcal{D}}$ inside the chunk, we need the equality to hold, which happens only if every square in the chunk is covered by a domino that lies in this chunk. This happens only when the chunk is tiled by the horizontal dominoes. $\square$

3.2 Boundary affine slices and Goncharov–Kenyon Hamiltonians

Let us refer to copies of $\hat{A}_{2n-1}$ in $A_m \otimes \hat{A}_{2n-1}$ as *affine slices*. We will distinguish *boundary affine slices* which correspond to the two boundary vertices of the Dynkin diagram A_m , and *internal affine slices* which correspond to the internal vertices of the Dynkin diagram A_m . In this section, we identify the recurrence coefficients of boundary affine slices as Goncharov–Kenyon Hamiltonians introduced in [15]. We shall see in Sect. 3.4 that the recurrence coefficients of the internal affine slices can be expressed through the Goncharov–Kenyon Hamiltonians using plethysm of symmetric functions. While we leave the question of an explicit formula for internal affine slices coefficients open, we will be able to deduce some of their properties in Sects. 3.4 and 3.5.

3.2.1 Thurston height

Recall from [39] the following definition of *Thurston height* function associated to a domino tiling. Consider a cylinder $\mathfrak{C}_{m,2n}$ which we can think of as $(m + 1) \times 2n$ rectangle with sides of length $m + 1$ glued. We can identify the $m \times 2n$ non-boundary nodes with vertices of the quiver $A_m \otimes \hat{A}_{2n-1}$. Fix a chessboard coloring of the cylinder, and fix a node O at its bottom boundary such that the square to the right of O is colored black. Let $\mathcal{D}$ be a domino tiling of $\mathfrak{C}_{m,2n}$. Define the function

$$h : \text{nodes of } \mathfrak{C}_{m,2n} \longrightarrow \mathbb{Z}$$

as follows:

- $h(O) = 0$;
- if $a \rightarrow b$ is a directed edge of a domino in $\mathcal{D}$ and the cell to the right of it is black, then $h(b) - h(a) = 1$;
- if $a \rightarrow b$ is a directed edge of a domino in $\mathcal{D}$ and the cell to the right of it is white, then $h(b) - h(a) = -1$.

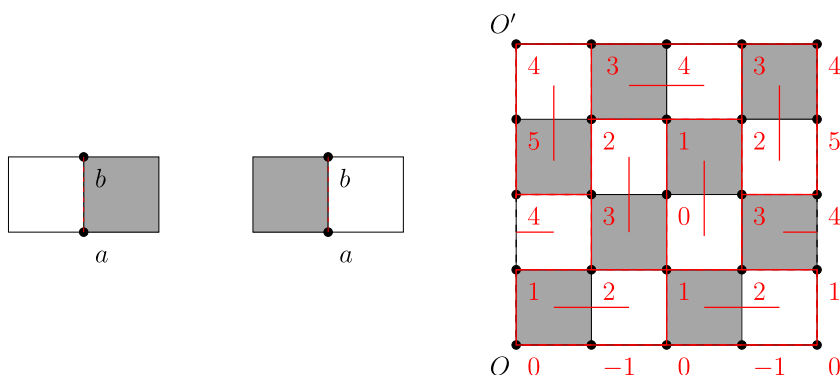


Fig. 13 Two cases to consider when constructing h and an example on $\mathfrak{C}_{3,4}$

If there is no cell to the right, we can still decide between the two options by looking at the cell to the left and assuming the cell to the right has an opposite color.

Theorem 3.2.1 *Thurston height h is a well-defined function on the nodes of $\mathfrak{C}_{m,2n}$.*

An example of a domino tiling of $\mathfrak{C}_{m,2n}$ and the associated Thurston height function can be seen in Fig. 13.

Proof It is known [39] that Thurston height function is well-defined for regions in the plane without holes. Thus, it is well-defined on the infinite periodic tiling obtained by lifting $\mathcal{D}$ to the universal cover of $\mathfrak{C}_{m,2n}$. It remains to argue that this height function is also periodic, and thus can be folded back onto the cylinder. Assume it is not periodic, then it must steadily grow or steadily decline as we circle around the cylinder. However, then it would reach arbitrary high or arbitrary low values, which is impossible since any node is within distance m from the lower boundary, which is filled with 0-s and -1 -s. The contradiction implies the desired property. $\square$

We can now define the height of a tiling $\mathcal{D}$ as

$$h(\mathcal{D}) = h(O') - h(O),$$

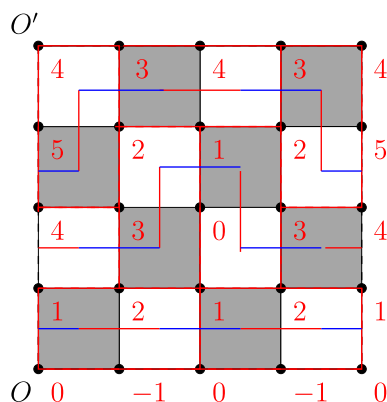
where O' is the node on the top boundary component opposite of O .

Proposition 3.2.2 *The function $h(\mathcal{D})$ takes values $4k$, $k = -(m+1)/2, \dots, (m+1)/2$ if m is odd, and takes values $4k+1$, $k = -m/2-1, \dots, m/2$ if m is even. There is a unique tiling having minimal height and a unique tiling having maximal height.*

Proof As we walk from O to O' straight up, at each step the height changes either by ± 1 or ± 3 , depending on whether the step cuts a domino and what the colors on the sides are. The claims of the proposition then easily follow. $\square$

Let us refer to the tiling with the minimal height as *the sea*, and denote it $\mathfrak{S}$. One can give an alternative definition of the height of a tiling $h(\mathcal{D})$ as follows. For any tiling $\mathcal{D}$, put $\mathfrak{S}$ and $\mathcal{D}$ on the same picture. What we get is a *double dimer model*, where all dominos will split into closed cycles. An example of such superposition for the tiling in Fig. 13 is given in Fig. 14. The dominos of the sea $\mathfrak{S}$ are shown in blue.

Fig. 14 The superposition of $\mathfrak{S}$ and $\mathcal{D}$



The cycles created in the process may include contractible cycles and non-contractible cycles. Let us refer to the latter as *hula hoops*. Note that the contractible cycles may be just double edges, if $\mathfrak{S}$ and $\mathcal{D}$ share dominos. The example in Fig. 14 has zero contractible cycles and three hula hoops. Denote $h(\mathcal{D})$ the number of hula hoops created by superposing $\mathcal{D}$ and $\mathfrak{S}$.

Proposition 3.2.3 *We have*

$$h(\mathcal{D}) = h(\mathfrak{S}) + 4h(\mathcal{D}).$$

Proof One can always walk from O to O' so that the only steps that cross dominos, rather than follow their boundaries, are the ones crossing the hula hoops. It is easy to see that each such crossing is responsible for a difference of 4 between the accumulated parts of $h(\mathfrak{S})$ and $h(\mathcal{D})$. Furthermore, it is easy to see that since $\mathfrak{S}$ is the tiling with minimal height, each such crossing must make $h(\mathcal{D})$ larger by 4 than $h(\mathfrak{S})$, as opposed to smaller. Otherwise we could change $\mathfrak{S}$ by using the dominos of $\mathcal{D}$ from the hula hoop, and decrease its height even further, which is impossible. The proposition claim follows. $\square$

3.2.2 The recurrence

Define *Goncharov–Kenyon Hamiltonians* to be the sums

$$H_r = \sum_{h(\mathcal{D})=r} \prod_{u \in \mathfrak{C}_{m,2n}} u^{1-d_{\mathcal{D}}(u)},$$

where $d_{\mathcal{D}}(u)$ is as before the degree of u in the associated graph $G_{\mathcal{D}}$ on the cylinder, and the sum is taken over all tilings $\mathcal{D}$ of height $4r + h(\mathfrak{S})$. Here on the boundary we always have $u = 1$, which makes H_r -s into functions of variables at the vertices of the quiver $A_m \otimes \hat{A}_{2n-1}$.

Example 3.2.4 Take $m = 3$ and $n = 1$. We have six variables a, b, c, d, e, f at the vertices of the quiver $A_3 \otimes \hat{A}_1$. Figure 15 shows the domino tilings contributing to H_1 and the monomials they contribute. As a result, we find

$$H_1 = \frac{ab}{de} + \frac{a}{be} + \frac{b}{ad} + \frac{c}{f} + \frac{d}{a} + \frac{ef}{bc} + \frac{e}{cf}.$$

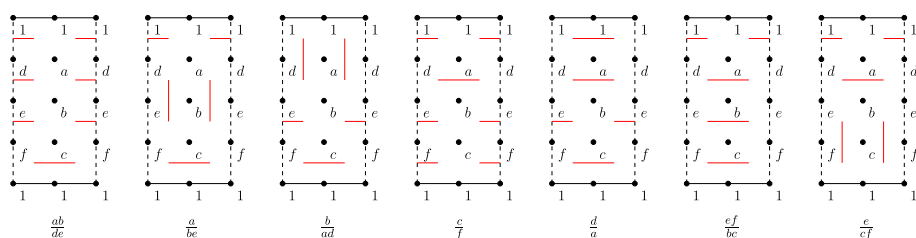


Fig. 15 The tilings $\mathcal{D}$ with height $h(\mathcal{D}) = -4$

Similarly we find

$$\begin{aligned}
 H_2 &= \frac{abc}{def} + \frac{ab}{dcf} + \frac{bc}{adf} + \frac{ac}{be} + \frac{be}{acdf} + \frac{ef}{ad} \\
 &\quad + \frac{af}{cd} + \frac{e}{b} + \frac{cd}{af} + \frac{de}{acf} + \frac{b}{e} + \frac{df}{be} + \frac{def}{abc}, \text{ and} \\
 H_3 &= \frac{bc}{ef} + \frac{c}{be} + \frac{b}{cf} + \frac{a}{d} + \frac{f}{c} + \frac{de}{ab} + \frac{e}{ad}.
 \end{aligned}$$

The only tiling contributing to H_0 is the sea $\mathfrak{S}$, and it is easy to see that

$$H_0 = H_4 = 1.$$

Let v' be the vertex diametrically opposite to v on the same affine slice. Let

$$v(j) = \begin{cases} v & \text{if } j \text{ is even;} \\ v' & \text{if } j \text{ is odd.} \end{cases}$$

We are ready to state the main theorem of the section.

Theorem 3.2.5 *For any vertex v on the top boundary affine slice of the quiver $A_m \otimes \hat{A}_{2n-1}$ the T -system satisfies for any t the following recursion*

$$\begin{aligned}
 T_{v(0)}(t + (m+1)n) - H_1 T_{v(1)}(t + mn) + \cdots + (-1)^m H_m T_{v(m)}(t + n) \\
 + (-1)^{m+1} T_{v(m+1)}(t) = 0.
 \end{aligned}$$

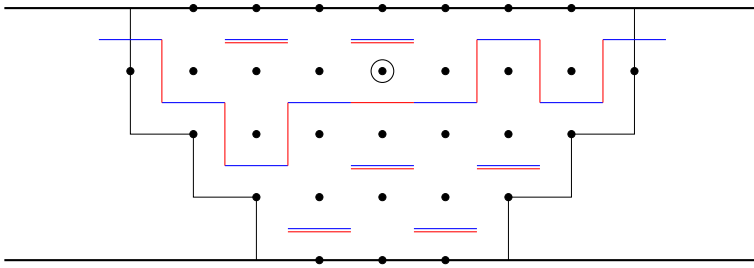
Similarly, for any vertex v on the bottom boundary affine slice and any t we have

$$\begin{aligned}
 T_{v(0)}(t + (m+1)n) - H_m T_{v(1)}(t + mn) + \cdots + (-1)^m H_1 T_{v(m)}(t + n) \\
 + (-1)^{m+1} T_{v(m+1)}(t) = 0.
 \end{aligned}$$

3.3 Proof of the recurrence

In this section we prove Theorem 3.2.5. We consider the case of v lying in the top affine slice of the quiver $A_m \otimes \hat{A}_{2n-1}$. The case of the bottom affine slice is similar.

As we have seen, the Laurent monomials entering both the H_i -s and the T_v -s have an interpretation in terms of weights of domino tilings. We are going to construct an involution which associates each Laurent monomial in the expansion by linearity of $H_i T_{v(i)}(t + (m+1-n))$ to an equal Laurent monomial in either $H_{i-1} T_{v(i-1)}(t + (m+2-i)n)$ or $H_{i+1} T_{v(i+1)}(t +$



$(m - i)n$). This implies that all of the terms cancel out as desired, since thus created pairs of Laurent monomials are equal but have opposite signs.

Let $\mathcal{Z}_{v(i)}(t + (m + 1 - i)n)$ be the fragment of an Aztec diamond lying inside the universal cover, as defined above. Let $\mathcal{D}_{\mathcal{Z}}$ be a domino tiling of $\mathcal{Z}_{v(i)}(t + (m + 1 - i)n)$, contributing a term into $T_{v(i)}(t + (m + 1 - i)n)$. Superpose $\mathcal{D}_{\mathcal{Z}}$ with the universal cover of the sea $\mathfrak{S}$, which is a tiling of the universal cover of the cylinder. Consider the part of the result that intersects $\mathcal{Z}_{v(i)}(t + (m + 1 - i)n)$.

Example 3.3.2 In Fig. 16 an example is presented of a superposition of $\mathcal{D}_{\mathbb{Z}}$, shown in red, with the universal cover of $\mathfrak{S}$, shown in blue. Here $m = 3, n = 2$, the vertex $v(2) = v$ is circled and $\mathcal{Z}_v(4)$ is shown. One can clearly see the hose, while the rest of the dominos form 2-cycles.

Now we are ready to define the involution. Assume we are given a pair $(\mathcal{D}_{\mathfrak{E}}, \mathcal{D}_{\mathfrak{Z}})$ with corresponding Laurent monomials contributing to the product $H_i T_{v(i)}(t + (m + 1 - i)n)$. Take the hose associated with $\mathcal{D}_{\mathfrak{Z}}$ and start following its edges *on the cylinder* $\mathfrak{E}_{m,2n}$. One of the two events is going to occur:

- In the first case, take the first such self-intersection, and *extract* from it the corresponding hula hoop. By this we mean cut out from the hose the dominos of the part between endpoints of self-intersection, and add the corresponding red dominos to $\mathcal{D}_{\mathcal{C}}$ instead of the blue ones it is currently using. In the second case, take the first such intersection with H and *insert* H

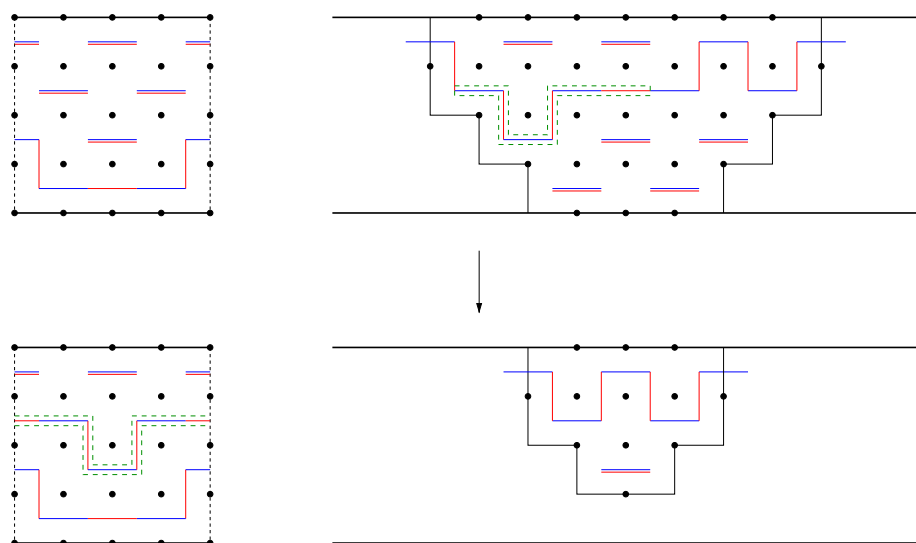


Fig. 17 The effect of the first event on a pair $(\mathcal{D}_{\mathcal{C}}, \mathcal{D}_{\mathcal{Z}})$

to extend the hose, by pasting it at this first point of intersection. We then remove the red edges of H from $\mathcal{D}_{\mathcal{C}}$, substituting the blue sea edges instead. In either case we get a new pair $(\mathcal{D}'_{\mathcal{C}}, \mathcal{D}'_{\mathcal{Z}})$.

Lemma 3.3.3 *The resulting pair is a well-defined pair of domino tilings that contributes either to $H_{i-1}T_{v(i-1)}(t + (m + 2 - i)n)$ or to $H_{i+1}T_{v(i+1)}(t + (m - i)n)$, depending on which of the two events occurred.*

Example 3.3.4 An example of a pair $(\mathcal{D}_{\mathcal{C}}, \mathcal{D}_{\mathcal{Z}})$ contributing to $H_1T_{v(1)}(4)$ for which the first event occurs is shown in Fig. 17. The new pair $(\mathcal{D}'_{\mathcal{C}}, \mathcal{D}'_{\mathcal{Z}})$ in this case contributes to $H_{i+1}T_{v(i+1)}(t + (m - i)n) = H_2T_{v(2)}(2)$.

An example of a pair $(\mathcal{D}_{\mathcal{C}}, \mathcal{D}_{\mathcal{Z}})$ contributing to $H_1T_{v(1)}(4)$ for which the second event occurs is shown in Fig. 18. The new pair $(\mathcal{D}'_{\mathcal{C}}, \mathcal{D}'_{\mathcal{Z}})$ in this case contributes to $H_{i-1}T_{v(i-1)}(t + (m + 2 - i)n) = H_0T_{v(0)}(6)$.

In both cases, the fragments that get either extracted or inserted are circled by a green dashed line.

Proof The only somewhat non-trivial part of the claim is why after a hula hoop is extracted from a hose, what remains is still a proper hose. The reason is that all blue dominos in the hose flow East, which means that the red dominos must flow North, East or South, but not West. This means that the red and the blue dominos that we need to connect after the extraction are compatible. $\square$

The final claim we need to conclude the theorem is the following.

Lemma 3.3.5 *This map is a weight-preserving involution on Laurent monomials.*

Proof If the first event occurred in $(\mathcal{D}_{\mathcal{C}}, \mathcal{D}_{\mathcal{Z}})$ and $(\mathcal{D}'_{\mathcal{C}}, \mathcal{D}'_{\mathcal{Z}})$ was created, then the second event occurs in $(\mathcal{D}'_{\mathcal{C}}, \mathcal{D}'_{\mathcal{Z}})$ at exactly the same place, and $(\mathcal{D}_{\mathcal{C}}, \mathcal{D}_{\mathcal{Z}})$ is created. Same holds vice versa. Thus, the map is an involution. The fact that it is weight preserving is easy to see from the way we assign weights to domino tilings. $\square$

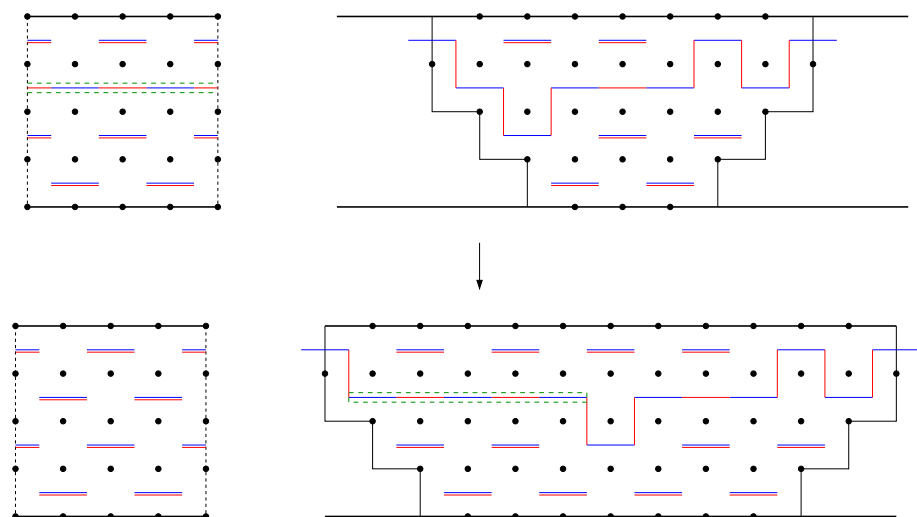


Fig. 18 The effect of the second event on a pair $(\mathcal{D}_{\mathcal{C}}, \mathcal{D}_{\mathcal{Z}})$

3.4 Affine slices and plethysm

In this section, we explain how to express the recurrence coefficients of the affine slices in $A_m \otimes \hat{A}_{2n-1}$ through the Goncharov–Kenyon Hamiltonians H_i . Note that Theorem 3.2.5 did not quite answer that yet for the boundary slices, since it involved variables at *two* vertices v and v' , rather than a single v . We rely here on results of [30], as well as on the language of tensors introduced there.

Recall that in [30] the T -system variables are interpreted as certain polynomial SL_{m+1} -invariants of a collection of $2n$ vectors in $\mathbb{C}^{m+1}$ and one matrix $A \in SL_{m+1}$. The key theorem is the following strengthening of [30, Theorem 1.11].

Theorem 3.4.1 *The variables on the r -th slice of $A_m \otimes \hat{A}_{2n-1}$, $r = 1, \dots, m$ satisfy the same recurrence as the exterior powers $\wedge^r(\hat{A}^q)$, $q \in \mathbb{Z}$, where $\hat{A} = A^2$.*

In particular, according to the Cayley–Hamilton theorem, the recurrence for $r = 1$ is given by the characteristic polynomial of $\hat{A}$.

Proof As q grows, we keep repeating the Dehn twists, which inserts $\hat{A} \otimes \dots \otimes \hat{A}$ into the tensor. Thus, we obtain the tensor $\hat{A}^q \otimes \dots \otimes \hat{A}^q$ in the middle. Furthermore, this tensor is attached to the anti-symmetrizing Levi-Cevita tensor, which results in the anti-symmetrization of $\otimes^r(\hat{A}^q)$, which is $\wedge^r(\hat{A}^q)$. $\square$

Corollary 3.4.2 *The recurrence coefficients of the affine slices $r = 1, \dots, m$ are expressed in terms of the Goncharov–Kenyon Hamiltonians H_i as plethysms of elementary symmetric functions and the power sum symmetric function $e_j[e_r[p_2]]$ are expressed through the original elementary symmetric functions $e_i = H_i$.*

Proof Theorem 3.2.5 tells us the recurrence satisfied by the sequence of T_v -s and $T_{v'}$ -s. To obtain the recurrence satisfied by T_v -s only, we need to take every second term of the sequence. In terms of the recurrence, this means we just need to square the roots of the

recurrence polynomial. This means that if $H_i = e_i(\lambda)$, then the coefficients on boundary levels are just the plethysms $e_i[p_2]$, which of course can be expressed as polynomials in the H_i -s.

Since in the construction of the ring of invariants in [30] the dimension count forces the vectors and the matrix A to be generic, the $r = 1$ affine slice cannot satisfy any linear recurrence of length shorter than $2n(m+1)$. This means that any two such linear recurrences must coincide, and thus the plethysms $e_i[p_2]$ of Goncharov–Kenyon Hamiltonians H_i are the coefficients of the characteristic polynomial of $\hat{A}$.

Now, if $\lambda_i, i = 1, \dots, m+1$ are the eigenvalues of $\hat{A}$, then products

$$\lambda_{i_1} \cdots \lambda_{i_r}, \quad 1 \leq i_1 < \cdots < i_r \leq m+1$$

are the eigenvalues of $\wedge^r(\hat{A})$. Then the coefficients of the corresponding characteristic polynomial of $\wedge^r(\hat{A}^q)$ are exactly the plethysms $e_j[e_r[p_2(\lambda)]]$. $\square$

Corollary 3.4.3 *The Goncharov–Kenyon Hamiltonians H_i are conserved quantities of the T -system.*

Proof As in the previous proof, the minimal recurrence satisfied by the boundary affine slice is unique, and thus its coefficients are the same no matter which moment we pick as $t = 0$. $\square$

Corollary 3.4.4 *The recurrence for the r th affine slice has the form*

$$T_v\left(t + 2n\binom{m+1}{r}\right) - \cdots \pm T_v(t) = 0$$

with exactly $\binom{m+1}{r} + 1$ terms on the left.

Proof This is clear since the size of $\wedge^r(\hat{A})$ is $\binom{m+1}{r}$. $\square$

Corollary 3.4.5 *The recurrence coefficients of r th and $(m+1-r)$ th affine slices are the same up to the reversal of the order. If m is odd, then the coefficients of the slice $r = (m+1)/2$ are palindromic.*

Proof Since $\hat{A} \in SL_{m+1}$, we know that the constant term of the characteristic polynomial is 1. Alternatively, we have already seen that $H_{m+1} = 1$. Either way, we see that $\prod_{i=1}^{m+1} \lambda_i = 1$. This means that the eigenvalues of $\wedge^r(\hat{A}^q)$ and of $\wedge^{m+1-r}(\hat{A}^q)$ are inverses of each other, and the claim follows. $\square$

Example 3.4.6 Consider the case $m = 3$. In this case we have 3 affine slices, two boundary and one internal. The recurrence relations satisfied by the T -system are as follows.

- If v lies on the $r = 1$ affine slice,

$$\begin{aligned} T_v(t + 8n) - (H_1^2 - 2H_2)T_v(t + 6n) + (H_2^2 - 2H_1H_3 + 2)T_v(t + 4n) \\ - (H_3^2 - 2H_2)T_v(t + 2n) + T_v(t) = 0. \end{aligned}$$

- If v lies on the $r = 2$ affine slice,

$$\begin{aligned} T_v(t + 12n) - (H_2^2 - 2H_1H_3 + 2)T_v(t + 10n) \\ + ((H_1^2 - 2H_2)(H_3^2 - 2H_2) - 1)T_v(t + 8n) \\ - ((H_1^2 - 2H_2)^2 + (H_3^2 - 2H_2)^2 - 2(H_2^2 - 2H_1H_3 + 2))T_v(t + 6n) \\ + ((H_1^2 - 2H_2)(H_3^2 - 2H_2) - 1)T_v(t + 4n) \\ - (H_2^2 - 2H_1H_3 + 2)T_v(t + 2n) + T_v(t) = 0. \end{aligned}$$

- If v lies on the $r = 3$ affine slice,

$$T_v(t + 8n) - (H_3^2 - 2H_2)T_v(t + 6n) + (H_2^2 - 2H_1H_3 + 2)T_v(t + 4n) - (H_1^2 - 2H_2)T_v(t + 2n) + T_v(t) = 0.$$

Here for example $(H_1^2 - 2H_2)^2 + (H_3^2 - 2H_2)^2 - 2(H_2^2 - 2H_1H_3 + 2)$ is determined by the plethysm

$$\begin{aligned} e_3[e_2(\lambda_1, \lambda_2, \lambda_3, \lambda_4)] &= e_3(\lambda_1\lambda_2, \lambda_1\lambda_3, \lambda_1\lambda_4, \lambda_2\lambda_3, \lambda_2\lambda_4, \lambda_3\lambda_4) \\ &= (\lambda_1^3\lambda_2\lambda_3\lambda_4 + \cdots) + (\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \cdots) \\ &\quad + 2(\lambda_1^2\lambda_2^2\lambda_3\lambda_4 + \cdots) = e_4e_1^2 + e_3^2 - 2e_2e_4 = e_1^2 + e_3^2 - 2e_2, \end{aligned}$$

followed by the plethysm $e_1[p_2] = e_1^2 - 2e_2$, $e_2[p_2] = e_2^2 - 2e_1e_3 + 2$, $e_3[p_2] = e_3^2 - 2e_2$. Throughout we use $e_4 = 1$.

3.5 Laurent property and positivity

Recall that the *upper cluster algebra* $\mathfrak{U}_A$ associated with a cluster algebra A is the algebra of all elements of the fraction field of A that can be expressed as Laurent polynomials in *any* cluster of A . Due to Laurent property of cluster algebras [9] we know that $A \subseteq \mathfrak{U}_A$. The equality holds in some cases, while in other cases $\mathfrak{U}_A$ is strictly larger. We refer the reader to [2] for a rigorous definition and properties of upper cluster algebras.

Theorem 3.5.1 *Goncharov–Kenyon Hamiltonians H_i are elements of the upper cluster algebra associated with the quiver $A_m \otimes \hat{A}_{2n-1}$.*

Of course, the H_i -s are Laurent expressions in terms of the initial cluster of this T -system by definition. Since we know they are conserved quantities, the same holds for any cluster in the T -system. However, the claim of the theorem is much stronger, since the T -system represents only one way to mutate the quiver, while the Laurentness is true for *any* such way.

Corollary 3.5.2 *The coefficients of recurrence polynomials of all vertices of $A_m \otimes \hat{A}_{2n-1}$ lie in the upper cluster algebra.*

Proof Since those coefficients are polynomials in the H_i -s by Corollary 3.4.2, the statement follows. $\square$

Sherman and Zelevinsky [33] have defined a *positive cone* inside the upper cluster algebra $\mathfrak{U}_A$ to be the subset of all elements of $\mathfrak{U}_A$ that are expressible as *positive* Laurent expression in any cluster of A .

Conjecture 3.5.3 *Goncharov–Kenyon Hamiltonians H_i are elements of the positive cone of the corresponding upper cluster algebras.*

Again, by definition the H_i -s are positive in terms of the clusters along time evolution of the T -system, but the claim of the conjecture is much stronger.

3.5.1 Proof of Theorem 3.5.1

We are going to use the standard proof technique, see, e.g., [2, Theorem 1.5]. Specifically, to know that a certain H_i lies in the upper cluster algebra, it suffices to check the Laurent

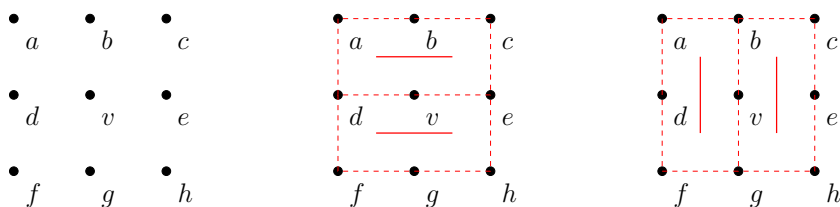


Fig. 19 The two possible ways to get v in the denominator of the corresponding monomial

condition with respect to some seed together with all the seeds obtained from it by a single mutation. The fact that the H_i -s are positive in the initial seed of the T -system is true by definition. Thus, it remains to check positivity in all seeds obtained by mutating just a single variable in the initial seed.

Let v be the variable that is mutated, and assume the surrounding variables are as in Fig. 19. Note that some of the variables may be equal to 1 if v is close to the boundary.

When we mutate at v , we make a substitution

$$v \leftarrow \frac{bg + de}{v'}.$$

Let us consider the effect of this substitution on the Laurent monomials entering H_i , which as we know correspond to domino tilings $\mathcal{D}$:

$$H_i = \sum_{\mathfrak{h}(\mathcal{D})=i} \prod_{u \in \mathfrak{C}_{m,2n}} u^{1-d_{\mathcal{D}}(u)}.$$

For each tiling $\mathcal{D}$ where v does not appear at all in the monomial, i.e. where $d_{\mathcal{D}}(u) = 1$, or where v appears in the numerator, i.e. $d_{\mathcal{D}}(u) = 0$, the Laurentness is not violated by the substitution $v \leftarrow \frac{bg+de}{v'}$. Thus, it remains to consider the terms where v appears in the denominator, i.e. $d_{\mathcal{D}}(u) = 2$.

The key observation is that such tilings $\mathcal{D}$ come in pairs. This is because locally around vertex v they need to look in one of the two ways shown in Fig. 19. Furthermore, the local move swapping between those two ways to tile the surrounding 2×2 square does not change the height of the tiling. Thus, all tilings $\mathcal{D}$ contributing to the terms of H_i with v in the denominator indeed come in pairs, differing by the application of this local 2×2 square swap. Let $\mathcal{D}$ and $\mathcal{D}'$ be such a pair. Then

$$\prod_{u \in \mathfrak{C}_{m,2n}} u^{1-d_{\mathcal{D}}(u)} + \prod_{u \in \mathfrak{C}_{m,2n}} u^{1-d_{\mathcal{D}'}(u)} = \left(\frac{bg}{v} + \frac{de}{v} \right) b^{-d_{\mathcal{D}}(b)} g^{-d_{\mathcal{D}}(g)} \prod_{u \in \mathfrak{C}_{m,2n}, u \neq b, v, g} u^{1-d_{\mathcal{D}}(u)}.$$

We see that after the substitution this becomes

$$\begin{aligned} & \left[\left(\frac{bg}{v} + \frac{de}{v} \right) b^{-d_{\mathcal{D}}(b)} g^{-d_{\mathcal{D}}(g)} \prod_{u \in \mathfrak{C}_{m,2n}, u \neq b, v, g} u^{1-d_{\mathcal{D}}(u)} \right]_{v \leftarrow \frac{bg+de}{v'}} \\ &= v' b^{-d_{\mathcal{D}}(b)} g^{-d_{\mathcal{D}}(g)} \prod_{u \in \mathfrak{C}_{m,2n}, u \neq b, v, g} u^{1-d_{\mathcal{D}}(u)}, \end{aligned}$$

which is a Laurent expression. The statement follows. $\square$

Remark 3.5.4 Alternatively, both Theorem 3.5.1 and Corollary 3.4.3 can be deduced directly from Urban Renewal Theorem, see, for example, [34, Section 5.2].

3.6 Conjectures

Let v be any vertex of any of the quivers in the classification of Theorem 2.2.5. We formulate here conjectures generalizing the results of this paper from $A_m \otimes \hat{A}_{2n-1}$, which is a special case of the first family, to *all* families from our classification.

The following conjecture generalizes Corollaries 3.4.2 and 3.4.4. In light of Theorem 2.2.5 this conjecture is a stronger version of [30, Conjecture 1.9].

Conjecture 3.6.1 *For every vertex v , there exist numbers i and N and rational functions $J_0 = 1, J_1, \dots, J_N, J_{N+1} = 1$ in $\mathbb{Q}(\mathbf{x})$ such that*

- *The J_k -s are the conserved quantities of the T -system;*
- *For any t we have*

$$J_0 T_v(t + i(N + 1)) - J_1 T_v(t + iN) + \dots \pm J_{N+1} T_v(t) = 0.$$

Note that among the linear recurrences satisfied by the sequences there is a *minimal one*. This is because if two recurrences are satisfied, then so is one given by the greatest common divisor of their characteristic polynomials. Let us from now on assume that the choices of i, N and J_k -s are made so that the resulting recurrence is minimal.

The following conjecture generalizes Theorem 3.5.1 and Corollary 3.5.2.

Conjecture 3.6.2 *The J_k -s belong to the upper cluster algebra of the cluster algebra associated with Q . In particular, they are Laurent polynomials in the variables at any time t .*

The following conjecture generalizes Conjecture 3.5.3.

Conjecture 3.6.3 *The J_k -s are positive Laurent expressions in terms of any cluster of the cluster algebra. In other words, they are elements of the positive cone inside the upper cluster algebra, as defined by Sherman and Zelevinsky [33].*

Our next conjecture is open even in Type $A_m \otimes \hat{A}_{2n-1}$. Let Q be any affine $\boxtimes$ finite quiver and let the J_k -s be as above. Consider an infinite Toeplitz matrix $\mathcal{H} = \mathcal{H}(m, n)$ where the entries are defined as follows:

$$\mathcal{H}_{i,j} = \begin{cases} J_{j-i} & \text{if } 0 \leq j - i \leq m + 1; \\ 0 & \text{otherwise.} \end{cases}$$

Conjecture 3.6.4 *All minors of $\mathcal{H}$ are either identically 0 or positive Laurent polynomials in $\mathbf{x}$.*

In other words, we conjecture that the J_i -s form a *totally positive sequence*, or *Pólya frequency sequence*, see [4] for the background. We also state the following weaker version of Conjecture 3.6.4:

Conjecture 3.6.5 *The roots of the recurrence polynomial*

$$J(z) = z^{N+1} - J_1 z^N + \dots \pm J_N z \mp 1$$

are positive real numbers.

Each of the types A_m , D_m and E_6 has a canonical involution on the Dynkin diagram, sending the diagram to itself. Denote this involution η . Assume our T -system is of the tensor product type, and more specifically of the form $\Lambda' \otimes \hat{\Lambda}$, where Λ' is a finite type Dynkin diagram of type A_m , D_{2m+1} or E_6 , and $\hat{\Lambda}$ is an arbitrary extended Dynkin diagram. Let v' be the vertex of $\Lambda' \otimes \hat{\Lambda}$ having the same $\hat{\Lambda}$ coordinate, but whose Λ' coordinate is obtained from that of v via involution η . The following conjecture generalizes Corollary 3.4.5.

Conjecture 3.6.6 *The recurrence polynomials of v and v' have the same coefficients but in the opposite order. In particular, if $v = v'$, then the recurrence polynomial $J(z)$ is palindromic.*

Assume now we are in *any other case*, i.e. either our T -system belongs to a different family of the classification, or it is a tensor product but Λ' is not of type A_m , D_{2m+1} or E_6 . The following conjecture again generalizes Corollary 3.4.5.

Conjecture 3.6.7 *The recurrence polynomial $J(z)$ of v is palindromic.*

Part 4. Tropical T -systems.

4.1 The behavior of tropical T -systems

4.1.1 Tropical T -systems: definition

Each bipartite recurrent quiver Q has the corresponding T -system which we will call *the geometric T -system* associated with Q in order to distinguish it from another system which we introduce in this section. We refer the reader to Example 4.3.2 for an illustration of most of the statements that we prove in Sects. 4.1.1–4.3.

Definition 4.1.1 Let Q be a bipartite recurrent quiver, and let $\lambda : \text{Vert}(Q) \rightarrow \mathbb{Z}$ be any map. Then the *tropical T -system* associated with Q is a family of integers $t_v^\lambda(t) \in \mathbb{Z}$ for every $v \in \text{Vert}(Q)$, $t \in \mathbb{Z}$ with $t + \epsilon_v$ even satisfying the following relations:

$$t_v^\lambda(t+1) + t_v^\lambda(t-1) = \max \left(\sum_{u \rightarrow v} t_u^\lambda(t), \sum_{v \rightarrow w} t_w^\lambda(t) \right);$$

$$t_v^\lambda(\epsilon_v) = \lambda(v).$$

It is apparent from the definition that $t_v^\lambda(t)$ is the *tropicalization* of $T_v(t)$. One can define a tropical T -system with values in $\mathbb{Q}$ or $\mathbb{R}$, but for our purposes it is sufficient to consider only the integer-valued version (see also Remark 4.2.2). The defining recurrence relation can be translated into the language of bigraphs as follows: if $G(\Gamma, \Delta)$ is a bipartite bigraph then the relation becomes

$$t_v^\lambda(t+1) + t_v^\lambda(t-1) = \max \left(\sum_{(u,v) \in \Gamma} t_u^\lambda(t), \sum_{(v,w) \in \Delta} t_w^\lambda(t) \right).$$

If $P(\mathbf{x}) \in \mathbb{Z}[\mathbf{x}^{\pm 1}]$ is a multivariate Laurent polynomial in variables $(x_v)_{v \in \text{Vert}(Q)}$ then define $P|_{\mathbf{x}=q^\lambda} \in \mathbb{Z}[q^{\pm 1}]$ to be the univariate Laurent polynomial in q obtained from P by substituting $x_v = q^{\lambda(v)}$ for all $v \in \text{Vert}(Q)$. Further, define $\deg_{\max}(q, P|_{\mathbf{x}=q^\lambda})$ to be the maximal degree of q in $P|_{\mathbf{x}=q^\lambda}$. The following claim gives a connection between the geometric and tropical T -systems:

Proposition 4.1.2 (see [13, Lemma 6.3]) For every $v \in \text{Vert}(Q)$, $t \in \mathbb{Z}$ with $t + \epsilon_v$ even and any $\lambda : \text{Vert}(Q) \rightarrow \mathbb{Z}$, we have

$$t_v^\lambda(t) = \deg_{\max}(q, T_v(t) \mid_{\mathbf{x}=q^\lambda}).$$

4.1.2 Linear algebraic properties of the affine Coxeter transformation

Let $\hat{\Lambda}$ be a bipartite affine ADE Dynkin diagram, and let w and b be the numbers of white and black vertices in $\hat{\Lambda}$ respectively. One can view a map $\mathbf{u} : \text{Vert}(\hat{\Lambda}) \rightarrow \mathbb{Z}$ as a vector $\begin{pmatrix} \mathbf{u}_W \\ \mathbf{u}_B \end{pmatrix} \in \mathbb{Z}^{w+b}$. Then the adjacency matrix $A_{\hat{\Lambda}}$ of $\hat{\Lambda}$ has the form

$$A_{\hat{\Lambda}} = \begin{pmatrix} 0 & A \\ A^t & 0 \end{pmatrix}$$

where A is a $w \times b$ matrix and t denotes matrix transpose. Define the *mutation matrices*

$$\omega_W := \begin{pmatrix} -I_w & A \\ 0 & I_b \end{pmatrix}; \quad \omega_B := \begin{pmatrix} I_w & 0 \\ A^t & -I_b \end{pmatrix}.$$

Here I_k is the identity $k \times k$ matrix. Finally, the *Coxeter transformation* for $\hat{\Lambda}$ is defined as a product $\mathbf{C} = \omega_B \omega_W$. By Lemma 1.1.8, the matrix $A_{\hat{\Lambda}}$ has a dominant eigenvector $\mathbf{v} = \begin{pmatrix} \mathbf{v}_W \\ \mathbf{v}_B \end{pmatrix}$ corresponding to the eigenvalue 2. This means

$$A\mathbf{v}_B = 2\mathbf{v}_W; \quad A^t\mathbf{v}_W = 2\mathbf{v}_B.$$

Just as in Part 2, all the coordinates of $\mathbf{v}$ are assumed to be positive integers with greatest common divisor equal to 1. For a vector $\mathbf{u} = \begin{pmatrix} \mathbf{u}_W \\ \mathbf{u}_B \end{pmatrix}$, we define three linear functionals as follows:

$$\begin{aligned} \text{SPEED}_W(\mathbf{u}) &:= \langle \mathbf{v}_B, \mathbf{u}_B \rangle - \langle \mathbf{v}_W, \mathbf{u}_W \rangle; & \text{SPEED}_B(\mathbf{u}) &:= -\text{SPEED}_W(\mathbf{u}); \\ \text{SUM}(\mathbf{u}) &:= \langle \mathbf{v}_W, \mathbf{u}_W \rangle + \langle \mathbf{v}_B, \mathbf{u}_B \rangle. \end{aligned}$$

Here $\langle \cdot, \cdot \rangle$ denotes the standard inner product in $\mathbb{R}^w$ and in $\mathbb{R}^b$.

Proposition 4.1.3 For any vector $\mathbf{u} = \begin{pmatrix} \mathbf{u}_W \\ \mathbf{u}_B \end{pmatrix}$, the following holds:

$$\begin{aligned} \text{SPEED}_B(\omega_W(\mathbf{u})) &= \text{SPEED}_W(\mathbf{u}); \\ \text{SPEED}_W(\omega_B(\mathbf{u})) &= \text{SPEED}_B(\mathbf{u}); \\ \text{SUM}(\omega_W(\mathbf{u})) &= \text{SUM}(\mathbf{u}) + 2 \text{SPEED}_W(\mathbf{u}); \\ \text{SUM}(\omega_B(\mathbf{u})) &= \text{SUM}(\mathbf{u}) + 2 \text{SPEED}_B(\mathbf{u}). \end{aligned}$$

Proof We will only prove the equalities for ω_W , and the argument is a pretty straightforward calculation:

$$\begin{aligned} \text{SPEED}_B(\omega_W(\mathbf{u})) &= \langle \mathbf{v}_W, A\mathbf{u}_B - \mathbf{u}_W \rangle - \langle \mathbf{v}_B, \mathbf{u}_B \rangle = \langle A^t\mathbf{v}_W, \mathbf{u}_B \rangle - \langle \mathbf{v}_W, \mathbf{u}_W \rangle - \langle \mathbf{v}_B, \mathbf{u}_B \rangle \\ &= \langle \mathbf{v}_B, \mathbf{u}_B \rangle - \langle \mathbf{v}_W, \mathbf{u}_W \rangle = \text{SPEED}_W(\mathbf{u}); \end{aligned}$$

Table 3 (see [35, Table 4.1])
Affine Coxeter numbers of affine
ADE Dynkin diagrams

$\hat{\Lambda}$	$\hat{A}_{2n-1}$	$\hat{D}_{2n}$	$\hat{D}_{2n+1}$	$\hat{E}_6$	$\hat{E}_7$	$\hat{E}_8$
$h_a(\hat{\Lambda})$	n	$2n - 2$	$2(2n - 1)$	6	12	30
$\theta(\hat{\Lambda})$	2	1	2	1	1	1

$$\begin{aligned} \text{SUM}(\omega_W(\mathbf{u})) &= \langle \mathbf{v}_W, A\mathbf{u}_B - \mathbf{u}_W \rangle + \langle \mathbf{v}_B, \mathbf{u}_B \rangle = \langle A^t \mathbf{v}_W, \mathbf{u}_B \rangle - \langle \mathbf{v}_W, \mathbf{u}_W \rangle + \langle \mathbf{v}_B, \mathbf{u}_B \rangle \\ &= \langle A^t \mathbf{v}_W, \mathbf{u}_B \rangle - \langle \mathbf{v}_W, \mathbf{u}_W \rangle + \langle \mathbf{v}_B, \mathbf{u}_B \rangle = \text{SUM}(\mathbf{u}) + 2 \text{SPEED}_W(\mathbf{u}). \end{aligned}$$

□

Proposition 4.1.3 says that SPEED is preserved while SUM grows linearly as we mutate. It turns out that up to a shift by $\mathbf{v}$, the mutation action is periodic:

Proposition 4.1.4 *For any affine ADE Dynkin diagram $\hat{\Lambda}$, there exists an integer $h_a(\hat{\Lambda})$ called the affine Coxeter number and an integer $\theta(\hat{\Lambda})$ such that for any vector $\mathbf{u} = \begin{pmatrix} \mathbf{u}_W \\ \mathbf{u}_B \end{pmatrix}$, we have*

$$\mathbf{C}^{h_a(\hat{\Lambda})} \mathbf{u} = \mathbf{u} + \theta(\hat{\Lambda}) \text{SPEED}_W(\mathbf{u}) \mathbf{v}. \quad (4.1.1)$$

Moreover,

$$\theta(\hat{\Lambda}) = \frac{4h_a(\hat{\Lambda})}{\langle \mathbf{v}, \mathbf{v} \rangle},$$

and the values of $h_a(\hat{\Lambda})$ and $\theta(\hat{\Lambda})$ are given in Table 3.

Proof Stekolshchik [35, Remark 4.3] gives complete information on the Jordan normal form of $\mathbf{C}$: all eigenvalues of $\mathbf{C}$ are roots of unity and the greatest common divisor of their periods is $h_a(\hat{\Lambda})$. Moreover, all of them have multiplicity one except for one of them ($\lambda = 1$) which has multiplicity 2. In our notation, the eigenvector attached to eigenvalue 1 is precisely $\mathbf{v}$ and the adjoint vector is $\mathbf{v}' := \frac{1}{4} \begin{pmatrix} \mathbf{v}_W \\ -\mathbf{v}_B \end{pmatrix}$ (see [35, Proposition 3.10]). We have $\mathbf{C}\mathbf{v}' = \mathbf{v} + \mathbf{v}'$, and they are orthogonal to each other and to all other eigenvectors. The result follows. □

Example 4.1.5 Let $\hat{\Lambda} = \hat{D}_4$. Then $h_a(\hat{\Lambda}) = 2$ and the dominant eigenvector is given by Fig. 2:

$$\mathbf{v} = \begin{pmatrix} 1 & 1 \\ 2 & 1 \end{pmatrix}.$$

Thus

$$\theta(\hat{\Lambda}) = \frac{4 \cdot 2}{1^2 + 1^2 + 1^2 + 1^2 + 2^2} = 1.$$

Consider the following vector of initial values:

$$\mathbf{u} = \begin{pmatrix} 1 & 1 \\ 2 & 2 \end{pmatrix}.$$

Let us assume that the vertex in the middle is white. Then $\text{SPEED}_W(\mathbf{u}) = 1 + 1 + 1 + 2 - 2 \cdot 2 = 1$. Since $h_a(\hat{\Lambda}) = 2$, we need to calculate

$$\mathbf{C}^2(\mathbf{u}) = \omega_B \omega_W \omega_B \omega_W \mathbf{u}.$$

The sequence of vectors that we will get is:

$$\mathbf{u} = \begin{smallmatrix} 1 & 1 \\ 1 & 2 \end{smallmatrix} \xrightarrow{\omega_W} \begin{smallmatrix} 1 & 1 \\ 1 & 3 \end{smallmatrix} \xrightarrow{\omega_B} \begin{smallmatrix} 2 & 2 \\ 2 & 3 \end{smallmatrix} \xrightarrow{\omega_W} \begin{smallmatrix} 2 & 2 \\ 2 & 4 \end{smallmatrix} \xrightarrow{\omega_B} \begin{smallmatrix} 2 & 2 \\ 2 & 3 \end{smallmatrix} = \mathbf{u} + \mathbf{v}.$$

We indeed see that

$$\mathbf{C}^{h_a(\hat{\Lambda})}(\mathbf{u}) = \mathbf{C}^2(\mathbf{u}) = \mathbf{u} + \mathbf{v} = \mathbf{u} + \theta(\hat{\Lambda}) \text{SPEED}(\mathbf{u})\mathbf{v}.$$

We would like to apply these observations to the tropical T -system $t_v^\lambda(t)$ defined above. Let $G = (\Gamma, \Delta)$ be a bigraph, and let $\hat{\Lambda}$ be a connected component of Γ isomorphic to an affine ADE Dynkin diagram. We define $t_{\hat{\Lambda}}^\lambda(t)$ to be the vector in $\mathbb{R}^w$ for t even and in $\mathbb{R}^b$ for t odd

which sends $v \in \text{Vert}(\hat{\Lambda})$ to $t_v^\lambda(t)$ when $t + \epsilon_v$ is even. In particular, the vectors $\begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix}$ and $\begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t+2) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix}$ belong to $\mathbb{R}^{w+b}$. Moreover, they satisfy the following inequality:

$$\begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t+2) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix} \geq \omega_W \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix}.$$

Here $\geq$ means that each coordinate of the vector on the left hand side is at least the corresponding coordinate of the vector on the right hand side. This inequality holds trivially by the definition of the tropical T -system. Moreover, it is an equality if and only if for every white vertex v of $\hat{\Lambda}$, we have

$$\sum_{(u,v) \in \Gamma} t_u^\lambda(2t+1) \geq \sum_{(v,w) \in \Delta} t_w^\lambda(2t+1). \quad (4.1.2)$$

Now, using the positivity of the coordinates of $\mathbf{v}$ and Proposition 4.1.3, we get that

$$\begin{aligned} \text{SPEED}_B \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t+2) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix} &\geq \text{SPEED}_W \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix}; \\ \text{SUM} \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t+2) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix} &\geq \text{SUM} \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix} + 2 \text{SPEED}_W \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(2t) \\ t_{\hat{\Lambda}}^\lambda(2t+1) \end{pmatrix}. \end{aligned}$$

And again, each inequality becomes an equality if and only if (4.1.2) holds for every white vertex v of $\hat{\Lambda}$.

Define the following functions of t :

$$\begin{aligned} \text{SPEED}_{\hat{\Lambda}}(t) &:= \begin{cases} \text{SPEED}_W \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(t) \\ t_{\hat{\Lambda}}^\lambda(t+1) \end{pmatrix}, & \text{if } t \text{ is even;} \\ \text{SPEED}_B \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(t+1) \\ t_{\hat{\Lambda}}^\lambda(t) \end{pmatrix}, & \text{if } t \text{ is odd;} \end{cases} \\ \text{SUM}_{\hat{\Lambda}}(t) &:= \begin{cases} \text{SUM} \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(t) \\ t_{\hat{\Lambda}}^\lambda(t+1) \end{pmatrix}, & \text{if } t \text{ is even;} \\ \text{SUM} \begin{pmatrix} t_{\hat{\Lambda}}^\lambda(t+1) \\ t_{\hat{\Lambda}}^\lambda(t) \end{pmatrix}, & \text{if } t \text{ is odd.} \end{cases} \end{aligned}$$

We have thus shown the following:

Proposition 4.1.6 *Let $G = (\Gamma, \Delta)$ be a bipartite bigraph, and let $\hat{\Lambda}$ be a connected component of Γ isomorphic to an affine ADE Dynkin diagram. Then for every $t \in \mathbb{Z}$ we have*

$$\text{SPEED}_{\hat{\Lambda}}(t+1) \geq \text{SPEED}_{\hat{\Lambda}}(t); \quad \text{SUM}_{\hat{\Lambda}}(t+1) \geq \text{SUM}_{\hat{\Lambda}}(t) + 2 \text{SPEED}_{\hat{\Lambda}}(t).$$

Moreover, either (4.1.2) holds or both inequalities are strict. $\square$

4.2 Solitonic behavior: soliton resolution

It turns out that for Zamolodchikov integrable quivers, the tropical T -system behaves linearly for all but finitely many moments of time. Namely, let Q be a bipartite recurrent quiver and assume Q is Zamolodchikov integrable but not Zamolodchikov periodic. Let $G(Q) = (\Gamma, \Delta)$ be the corresponding bipartite bigraph. Then Corollary 1.2.3 together with Remark 1.2.4 imply that all connected components of Γ are affine ADE Dynkin diagrams.

The following proposition will be later illustrated by Example 4.3.2.

Proposition 4.2.1 *Assume that Q is Zamolodchikov integrable and all connected components of Γ are affine ADE Dynkin diagrams as above. Then for every map $\lambda : \text{Vert}(Q) \rightarrow \mathbb{Z}$ there exists an integer t_0 such that for every $|t| > t_0$ and for every $v \in \text{Vert}(Q)$ with $t + \epsilon_v$ even we have*

$$\sum_{(u,v) \in \Gamma} t_u^\lambda(t+1) \geq \sum_{(v,w) \in \Delta} t_w^\lambda(t+1).$$

In other words, for any initial data λ , the inequality (4.1.2) is violated only finitely many times.

Proof If the inequality (4.1.2) is violated infinitely many times, then there exists a connected component $\hat{\Lambda}$ of Γ such that $\text{SPEED}_{\hat{\Lambda}}(t) \rightarrow +\infty$ as $t \rightarrow +\infty$, because each time (4.1.2) is violated, $\text{SPEED}_{\hat{\Lambda}}(t)$ increases by at least 1 (see Proposition 4.1.6). In this case, again, by Proposition 4.1.6, $\text{SUM}_{\hat{\Lambda}}(t)$ grows superlinearly. By Proposition 4.1.2, $\text{SUM}_{\hat{\Lambda}}(t)$ is just a linear combination of $\deg_{\max}(q, T_v(t) |_{\mathbf{x}=q^\lambda})$ for $v \in \hat{\Lambda}$, and thus there is a vertex $v \in \hat{\Lambda}$ for which $\deg_{\max}(q, T_v(t) |_{\mathbf{x}=q^\lambda})$ grows superlinearly. But the values of $T_v(t)$ satisfy a linear recurrence, and thus $\deg_{\max}(q, T_v(t) |_{\mathbf{x}=q^\lambda})$ cannot grow faster than linearly. $\square$

Remark 4.2.2 This proof works exactly the same way if the values of t^λ are assumed to lie in $\mathbb{Q}$ instead of $\mathbb{Z}$. We do not know whether the result of Proposition 4.2.1 holds when the values of t^λ belong to $\mathbb{R}$.

Now we are finally able to deduce Theorem 1.2.5:

Proof of Theorem 1.2.5 By Corollary 1.2.3, Remark 1.2.4 and Proposition 1.1.10, we need to show that if all components of Γ and of Δ are affine ADE Dynkin diagrams then Q cannot be recurrent. By Proposition 4.2.1, the inequality

$$\sum_{(u,v) \in \Gamma} t_u^\lambda(t+1) \geq \sum_{(v,w) \in \Delta} t_w^\lambda(t+1)$$

is violated finitely many times. By symmetry between Γ and Δ , the reverse inequality

$$\sum_{(u,v) \in \Gamma} t_u^\lambda(t+1) \leq \sum_{(v,w) \in \Delta} t_w^\lambda(t+1)$$

is also violated finitely many times. Therefore after finitely many steps we will have

$$\sum_{(u,v) \in \Gamma} t_u^\lambda(t+1) = \sum_{(v,w) \in \Delta} t_w^\lambda(t+1) \quad (4.2.1)$$

for all $v \in \text{Vert}(Q)$. To see that this is impossible, consider the following integers $y_v^\lambda(t)$ defined for $t + \epsilon_v$ even:

$$y_v^\lambda(t) = \sum_{(u,v) \in \Gamma} t_u^\lambda(t+1) - \sum_{(v,w) \in \Delta} t_w^\lambda(t+1).$$

It is well known that the numbers $y_v^\lambda(t)$ give (up to a sign) a solution to the *tropical Y-system* associated with Q , see, for example, [17]. Since the mutations for the tropical Y -system are involutions as well, they are invertible, so we get a contradiction with (4.2.1) because it states that for all initial data λ , the tropical Y -system $y_v^\lambda(t)$ eventually becomes zero. $\square$

Combining Proposition 4.2.1 with Proposition 4.1.4, we get the following corollary which we call “soliton resolution”:

Corollary 4.2.3 *Let Q be a Zamolodchikov integrable quiver and let $\hat{\Lambda}$ be a component of Γ isomorphic to an affine ADE Dynkin diagram. Then for every map $\lambda : \text{Vert}(Q) \rightarrow \mathbb{Z}$, there exist integers $\text{SPEED}_{\hat{\Lambda}}^+(\lambda)$ and $\text{SPEED}_{\hat{\Lambda}}^-(\lambda)$ such that*

- for all $t \gg 0$ and all $v \in \hat{\Lambda}$ we have

$$t_v^\lambda(t + 2h_a(\hat{\Lambda})) = t_v^\lambda(t) + \theta(\hat{\Lambda}) \text{SPEED}_{\hat{\Lambda}}^+(\lambda) \mathbf{v}(v); \quad (4.2.2)$$

- for all $t \ll 0$ and all $v \in \hat{\Lambda}$ we have

$$t_v^\lambda(t - 2h_a(\hat{\Lambda})) = t_v^\lambda(t) + \theta(\hat{\Lambda}) \text{SPEED}_{\hat{\Lambda}}^-(\lambda) \mathbf{v}(v); \quad (4.2.3)$$

In other words, the values of t_v^λ grow linearly for $|t| \gg 0$.

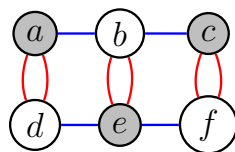
For instance, the integers $\text{SPEED}_{\hat{\Lambda}}^+(\lambda)$ and $\text{SPEED}_{\hat{\Lambda}}^-(\lambda)$ are calculated in example 4.3.2.

Let us explain the soliton terminology. Assume Q is an affine $\boxtimes$ finite ADE bigraph and consider the associated tropical T -system t^λ . Its restriction to each affine slice $\hat{\Lambda}$ behaves independently of other slices when $|t| \gg 0$. We treat it as a particle (a *1-soliton*). Then what happens is that when t grows from $-\infty$, the particles move independently with constant speeds given by (4.2.3). Then for small values of t they start interacting with each other and eventually they again start moving independently with constant speeds given by (4.2.2). Such a phenomenon is commonly called *soliton resolution*, see [38].

Corollary 4.2.4 *Let Q be a Zamolodchikov integrable quiver and let $\hat{\Lambda}$ be a component of Γ isomorphic to an affine ADE Dynkin diagram. Then for every map $\lambda : \text{Vert}(Q) \rightarrow \mathbb{Z}$, the following are equivalent:*

- (1) $\text{SPEED}_{\hat{\Lambda}}^+(\lambda) = 0$;
- (2) $\text{SPEED}_{\hat{\Lambda}}^-(\lambda) = 0$;
- (3) t_v^λ is a periodic sequence for every $v \in \text{Vert}(Q)$.

Proof It is obvious that (3) implies (1) and (2). The fact that each of them implies (3) follows from Corollary 4.2.3: if $\text{SPEED}_{\hat{\Lambda}}^+(\lambda) = 0$ then t_v^λ is periodic for $t \gg 0$, but then it is periodic for all t . $\square$

Fig. 20 The bigraph $A_3 \otimes \hat{A}_1$


4.3 Solitonic behavior: speed conservation

In this section we show that the speeds with which affine slices move get preserved after the scattering process is over, in the sense of Corollary 4.2.3. This can be viewed as a tropical version of Corollary 3.4.5.

Specifically, let $\hat{\Lambda}_1, \dots, \hat{\Lambda}_m$ be the m affine slices of $A_m \otimes \hat{A}_{2n-1}$, and for $r = 1, 2, \dots, m$ denote

$$\text{SPEED}_r^+ := \text{SPEED}_{\hat{\Lambda}_r}^+(\lambda); \quad \text{SPEED}_r^- := \text{SPEED}_{\hat{\Lambda}_r}^-,$$

where $\lambda : \text{Vert}(Q) \rightarrow \mathbb{Z}$ is fixed throughout this section.

Theorem 4.3.1 *For any $1 \leq r \leq m$ we have*

$$\text{SPEED}_r^+ = \text{SPEED}_{m+1-r}^-.$$

Example 4.3.2 Let us give an example of the kind of phenomenon in Theorem 4.3.1. Let us say that $m = 3$ and $n = 2$, so our quiver Q is $A_3 \otimes \hat{A}_1$ depicted in Fig. 20.

We will compactly draw this quiver as $\begin{smallmatrix} a & b & c \\ d & e & f \end{smallmatrix}$. Let $\hat{\Lambda}_1, \hat{\Lambda}_2, \hat{\Lambda}_3$ be the three red connected components, and assume we start our mutation sequence with black vertices. Then we have

$$\text{SPEED}_{\hat{\Lambda}_1}(t) = d - a; \quad \text{SPEED}_{\hat{\Lambda}_2}(t) = b - e; \quad \text{SPEED}_{\hat{\Lambda}_3}(t) = f - c.$$

We denote $\mathbf{S}(t) = (\text{SPEED}_{\hat{\Lambda}_1}(t), \text{SPEED}_{\hat{\Lambda}_2}(t), \text{SPEED}_{\hat{\Lambda}_3}(t))$. Now, for $t \ll 0$, $\text{SPEED}_{\hat{\Lambda}_r}(t) = -\text{SPEED}_r^-$ and for $t \gg 0$, $\text{SPEED}_{\hat{\Lambda}_r}(t) = \text{SPEED}_r^+$ for $r = 1, 2, 3$. The mutations and speeds for initial values $\begin{smallmatrix} 6 & 6 & 7 \\ 3 & 10 & 5 \end{smallmatrix}$ are given in Table 4. It is clear from the table that

$$\text{SPEED}_1^- = \text{SPEED}_3^+ = 3, \quad \text{SPEED}_2^- = \text{SPEED}_2^+ = 4, \quad \text{SPEED}_3^- = \text{SPEED}_1^+ = 2.$$

This agrees with the statement of Theorem 4.3.1. Next, it is also apparent from the table that the entries of $\mathbf{S}(t)$ weakly increase, and each of them changes if and only if for at least one vertex in the corresponding connected component, the sum of blue neighbors is strictly larger than the sum of red neighbors. This is precisely the statement of Proposition 4.1.6. Finally, observe that for every vertex $v \in \hat{\Lambda}_r$, we have

$$t_v^\lambda(4) = t_v^\lambda(2) + 2 \text{SPEED}_r^+,$$

which is an application of Corollary 4.2.3.

Let

$$H_r^\oplus = \max_{\mathfrak{h}(\mathcal{D})=r} \sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}}(u)) \lambda(u),$$

be the tropicalizations of Goncharov–Kenyon Hamiltonians H_r . Here $d_{\mathcal{D}}(u)$ is as before the degree of u in the associated graph $G_{\mathcal{D}}$ on the cylinder, and the sum is taken over all tilings $\mathcal{D}$ of height $4r + \mathfrak{h}(\mathfrak{S})$. Here on the boundary we always have $u = 0$.

Table 4 The evolution of the tropical T -system of type $A_3 \otimes \hat{A}_1$. The blue boldface numbers are the ones for which the sum of blue neighbors was strictly larger than the sum of red neighbors

t	-5	-4	-3	-2	-1	0	1	2	3	4
$a\ b\ c$	6 6 7	0 6 3	0 -2 3	-2 -2 -1	-2 -1 -1	1 -1 0	1 2 0	5 2 6	5 10 6	9 10 12
$d\ e\ f$	3 10 5	3 2 5	-1 2 1	-1 -2 1	-1 -2 -3	-1 0 -3	3 0 3	3 6 3	7 6 9	7 14 9
$\mathbf{S}(t)$	-3 -4 -2	-3 -4 -2	-1 -4 -2	-1 0 -2	1 1 -2	2 1 3	2 2 3	2 4 3	2 4 3	2 4 3

Lemma 4.3.3 *The $H_r^\oplus$ are conserved quantities of the tropical T -system of type $A_m \otimes \hat{A}_{2n-1}$.*

Proof Follows from Proposition 4.1.2 combined with the fact that H_r 's themselves are conserved quantities of the corresponding geometric T -system, see Corollary 3.4.3. $\square$

Our strategy consists of proving the following proposition.

Proposition 4.3.4 *Both SPEED_r^+ and SPEED_{m+1-r}^- are equal to $H_r^\oplus$ for respectively $t \gg 0$ and $t \ll 0$.*

Example 4.3.5 Let us continue Example 4.3.2. Recall the formula for H_1 calculated in Example 3.2.4:

$$H_1 = \frac{ab}{de} + \frac{a}{be} + \frac{b}{ad} + \frac{c}{f} + \frac{d}{a} + \frac{ef}{bc} + \frac{e}{cf}.$$

Thus,

$$H_1^\oplus = \max(a + b - d - e, a - b - e, b - a - d, c - f, d - a, e + f - b - c, e - c - f).$$

For instance, at $t = -4$ we have

$$H_1^\oplus = \max(2, -11, -2, -3, 2, -1, -9) = 2.$$

Or we can take $t = 0$ instead and get

$$H_1^\oplus = \max(0, 1, 2, 2, 1, -3, 2) = 2.$$

We encourage the reader to check that for other moments of time, $H_1^\oplus$ is always equal to 2, which is a statement of Lemma 4.3.3. In agreement with Proposition 4.3.4, we have

$$H_1^\oplus = \text{SPEED}_1^+ = \text{SPEED}_3^-.$$

Theorem 4.3.1 follows trivially from Proposition 4.3.4 and Lemma 4.3.3, since as a conserved quantity $H_r^\oplus$ is the same at any point in time, including $t \gg 0$ and $t \ll 0$. Let us prove Proposition 4.3.4. We are going to prove the $\text{SPEED}_r^+ = H_r^\oplus$ part, the other part is essentially verbatim. Let us formulate several key lemmas.

Consider the time $t \gg 0$ large enough for all speeds to have stabilized. Adopt the convention $\text{SPEED}_0^+ = \text{SPEED}_{m+1}^- = 0$.

Lemma 4.3.6 *The speeds form a weakly subadditive sequence, i.e. for any $1 < r < m$ we have*

$$2 \text{SPEED}_r^+ \geq \text{SPEED}_{r-1}^+ + \text{SPEED}_{r+1}^+.$$

Informally, our strategy is to show that the maximum in the definition of $H_r^\oplus$ is achieved on the term equal to SPEED_r^+ . Note that we do not claim that this is the only term where the maximum is achieved, just that it is one of such terms. The following lemma is a major step. We postpone its proof, and first show how to use it to imply Proposition 4.3.4.

Lemma 4.3.7 *The maximum in the expression*

$$\max_{\mathfrak{h}(\mathcal{D})=r} \sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}}(u)) \lambda(u)$$

is achieved at one of the tilings $\mathcal{D}$ consisting entirely of horizontal dominos.

Recall that we can compute $h(\mathcal{D})$ by walking up from vertex O to vertex O' on the cylinder, collecting a contribution of ± 1 or ± 3 on each step. Let $\epsilon_i = +1$ if the i th step along this path contributes a positive value and let $\epsilon_i = -1$ if it contributes a negative value. It is easy to see that

$$h(\mathcal{D}) = \frac{m+1 + \sum_{i=1}^{m+1} \epsilon_i}{2}.$$

If all dominos of $\mathcal{D}$ are horizontal, each layer of the cylinder $\mathfrak{C}_{m,2n}$ has exactly two ways to be tiled, one contributing $\epsilon_i = +1$ and the other contributing $\epsilon_i = -1$.

Lemma 4.3.8 *If all dominos in $\mathcal{D}$ are horizontal, we have*

$$\sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}}(u))\lambda(u) = \frac{1}{2} \sum_{i=1}^{m+1} \epsilon_i \cdot (\text{SPEED}_i^+ - \text{SPEED}_{i-1}^+).$$

Proof Depending on which of the two ways to tile the i th layer of $\mathfrak{C}_{m,2n}$ is used, this part of the tiling contributes into degrees $d_{\mathcal{D}}(u)$ for exactly half of vertices on each of the affine slices i and $i-1$. Specifically, it either contributes to degrees of white u -s on the i th affine slice and degrees of black u -s for the $(i-1)$ -st affine slice, or the other way around. The statement of the lemma is the numerical expression of this observation. $\square$

Now we are ready to prove Proposition 4.3.4.

Proof By Lemma 4.3.7 we know that the maximum in the definition of $H_r^{\oplus}$ is achieved at one of the tilings $\mathcal{D}$ with all dominos horizontal. There are 2^{m+1} such tilings, corresponding to 2^{m+1} choices one can make for each ϵ_i : either $\epsilon_i = 1$ or $\epsilon_i = -1$. Furthermore, the maximum is taken over $\mathcal{D}$ -s with $h(\mathcal{D}) = r$, which means exactly r among ϵ -s are $+1$.

According to Lemma 4.3.6 we know that

$$\begin{aligned} \text{SPEED}_1^+ &= \text{SPEED}_1^+ - \text{SPEED}_0^+ \geq \text{SPEED}_2^+ - \text{SPEED}_1^+ \geq \dots \\ &\geq \text{SPEED}_{m+1}^+ - \text{SPEED}_m^+ = -\text{SPEED}_m^+. \end{aligned}$$

Thus the maximum is obviously achieved when *the first* r among ϵ -s are equal to $+1$, and the rest of them are equal to -1 . The terms cancel out resulting in

$$\frac{1}{2} \sum_{i=1}^{m+1} \epsilon_i \cdot (\text{SPEED}_i^+ - \text{SPEED}_{i-1}^+) = \text{SPEED}_r^+.$$

Thus, the maximal term in the expression for $H_r^{\oplus}$ is equal to SPEED_r^+ , as desired. $\square$

4.3.1 Proof of Lemma 4.3.7

Consider a domino tiling $\mathcal{D}$ of the cylinder $\mathfrak{C}_{m,2n}$ which has vertical dominos. Our strategy will be to construct a different tiling $\mathcal{D}'$, which has strictly less vertical dominos than $\mathcal{D}$, and such that

$$\sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}}(u))\lambda(u) \leq \sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}'}(u))\lambda(u).$$

Recall that there is a distinguished tiling $\mathfrak{S}$ which we call the sea, such that superposition of $\mathfrak{S}$ with any other tiling does not contain contractible closed cycles, except for possibly

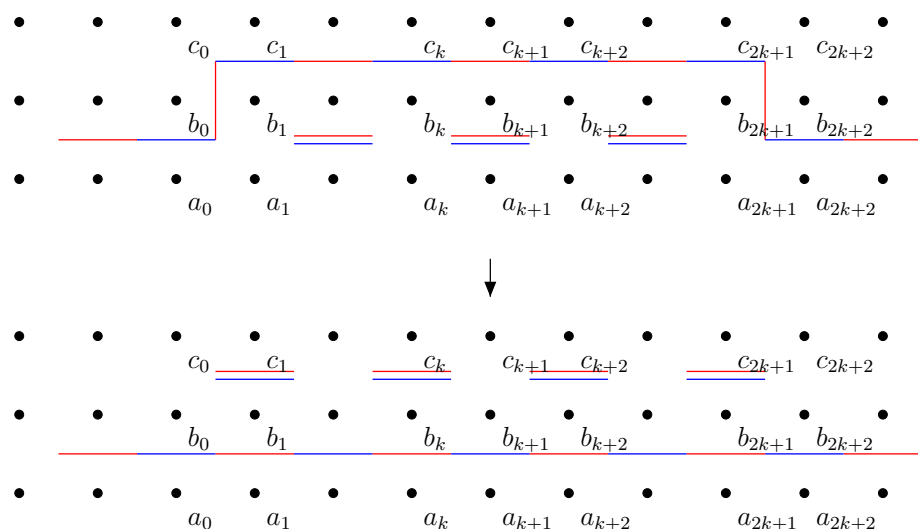


Fig. 21 The superpositions $\mathfrak{S} \cup \mathcal{D}$ (top) and $\mathfrak{S} \cup \mathcal{D}'$ (bottom) at the local part where hula hoop H is being straightened

double dominos. Consider the double dimer $\mathcal{D} \cup \mathfrak{S}$ obtained by taking superposition of our $\mathcal{D}$ and $\mathfrak{S}$. Several hula hoops are formed. In fact, if $\mathcal{D}$ is contributing to $H_r^\oplus$, then $\mathfrak{h}(\mathcal{D}) = r$ and exactly r hula hoops are formed.

Consider the lowest of the hula hoops H which has vertical dominos. It exists because we assume $\mathcal{D}$ has some vertical dominos, and the part of $\mathcal{D} \cup \mathfrak{S}$ not covered by hula hoops consists of horizontal double dominos. Choose one of the highest points in H and consider the horizontal part of H that contains this point together with two vertical dominos on its ends, see Fig. 21. Note that H may have several such parts, we just pick one of them. We create $\mathcal{D}'$ by straightening H in this local spot, as shown at the bottom of Fig. 21. It is clear that $\mathcal{D}'$ has strictly less vertical dominos than $\mathcal{D}$ does.

Proposition 4.3.9 *We have $\sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}}(u)) \lambda(u) \leq \sum_{u \in \mathfrak{C}_{m,2n}} (1 - d_{\mathcal{D}'}(u)) \lambda(u)$.*

Denote the vertices surrounding this part of H at time t by a_i -s, b_i -s, and c_i -s, $0 \leq i \leq 2k + 2$ as shown in Fig. 21. Let us put $\mu(v) := t_v^\lambda(t)$ for all $v \in \text{Vert}(Q)$. We assume $t \gg 0$ is sufficiently large for the claim of Corollary 4.2.3 to hold. Then the time evolution of a -s depends only on the values of μ at a -s, etc. More formally, the following lemma holds, describing the values of the tropical T -system at time $t + k$ for all $0 \leq k \leq n - 1$:

Lemma 4.3.10 *Define ϵ_k to be 0 if k is even and 1 if k is odd. Then we have*

$$\begin{aligned} t_{a_{k+1}}^\lambda(t+k) &= \sum_{i=1}^{k+1} \mu(a_{2i-1}) - \sum_{i=1}^k \mu(a_{2i}), \\ t_{b_k}^\lambda(t+k) &= \sum_{i=0}^k \mu(b_{2i}) - \sum_{i=1}^k \mu(b_{2i-1}), \\ t_{b_{k+2}}^\lambda(t+k) &= \sum_{i=1}^{k+1} \mu(b_{2i}) - \sum_{i=1}^k \mu(b_{2i+1}), \end{aligned}$$

$$t_{c_{k+1}}^\lambda(t+k) = \sum_{i=1}^{k+1} \mu(c_{2i-1}) - \sum_{i=1}^k \mu(c_{2i}).$$

Proof A straightforward application of recurrences

$$t_{a_i}^\lambda(t+j+1) = t_{a_{i-1}}^\lambda(t+j) + t_{a_{i+1}}^\lambda(t+j) - t_{a_i}^\lambda(t+j-1), \text{ etc.}$$

which hold due to Proposition 4.2.1 and our choice of large enough initial time t . $\square$

Now we are ready to prove Proposition 4.3.9.

Proof Each edge of $G(\mathcal{D})$ subtracts from the corresponding term of $H_r^\oplus$ two variables on its ends. Thus, we will compare those contributions for $G(\mathcal{D})$ and $G(\mathcal{D})$. We want to show that one is bigger than the other, which translates into

$$\begin{aligned} \mu(b_0) + \mu(b_1) + \sum_{i=1}^k (\mu(a_{2i}) + 2\mu(b_{2i}) + \mu(c_{2i})) + \mu(b_{2k+1}) + \mu(b_{2k+2}) \\ \geq \sum_{i=1}^k (\mu(a_{2i-1}) + 2\mu(b_{2i-1}) + \mu(c_{2i-1})). \end{aligned}$$

By Lemma 4.3.10 this is easily seen to be equivalent to

$$t_{b_k}^\lambda(t+k) + t_{b_{k+2}}^\lambda(t+k) \geq t_{a_{k+1}}^\lambda(t+k) + t_{c_{k+1}}^\lambda(t+k),$$

which holds by Proposition 4.2.1 and our choice of large enough t . $\square$

4.4 Conjectures

We conjecture that both soliton resolution and speed conservation properties hold for all families of our classification in Theorem 2.2.5.

For soliton resolution, we make the following conjecture, generalizing Proposition 4.2.1. It can also be viewed as a tropical analog of Conjecture 3.6.5.

Conjecture 4.4.1 *For any quiver Q in our affine $\boxtimes$ finite classification and any initial conditions either over $\mathbb{Z}$, or more generally over $\mathbb{R}$, there exists t_0 such that for $|t| > t_0$ the edges of finite component graph Δ do not affect the dynamics, i.e. for any vertex $v \in Q$ we have*

$$\sum_{(u,v) \in \Gamma} t_u^\lambda(t) \geq \sum_{(v,w) \in \Delta} t_w^\lambda(t).$$

In other words, for large enough time in both directions the affine slices of Q evolve as separate particles.

For speed conservation, we need to consider two cases, just as we did in Conjectures 3.6.6 and 3.6.7.

Each of the types A_m , D_m and E_6 has a canonical involution on the Dynkin diagram, sending the diagram to itself. As before, denote this involution η . Assume our tropical T -system is of the tensor product type, and more specifically of the form $\Lambda' \otimes \hat{\Lambda}$, where Λ' is a finite type Dynkin diagram of type A_m , D_{2m+1} or E_6 , and $\hat{\Lambda}$ is an arbitrary extended Dynkin diagram. Let $\hat{\Lambda}_v$ and $\hat{\Lambda}_{\eta(v)}$ be two affine slices of $\Lambda' \otimes \hat{\Lambda}$ such that their Λ' coordinates are related by η . Let $\text{SPEED}_v^\pm$ and $\text{SPEED}_{\eta(v)}^\pm$ be the corresponding speeds for $t \gg 0$ and

$t \ll 0$. Here we assume that Conjecture 4.4.1 holds and thus the speeds are well-defined. The following conjecture generalizes Theorem 4.3.1. It can also be viewed as a tropical analog of Conjecture 3.6.6.

Conjecture 4.4.2 *We have*

$$\text{SPEED}_v^+ = \text{SPEED}_{\eta(v)}^-.$$

Assume now we are in *any other case*, i.e. either our tropical T -system belongs to a different family of the classification, or it is a tensor product but Λ' is not of types A_m , D_{2m+1} or E_6 . Let $\hat{\Lambda}$ be any affine slice of the quiver, and let $\text{SPEED}_{\hat{\Lambda}}^\pm$ be the corresponding speeds as $t \gg 0$ and $t \ll 0$. The following conjecture again generalizes Corollary 3.4.5.

Conjecture 4.4.3 *We have*

$$\text{SPEED}_{\hat{\Lambda}}^+ = \text{SPEED}_{\hat{\Lambda}}^-.$$

References

1. Assem, I., Reutenauer, C., Smith, D.: Friezes. Adv. Math. **225**(6), 3134–3165 (2010)
2. Berenstein, A., Fomin, S., Zelevinsky, A.: Cluster algebras. III. Upper bounds and double Bruhat cells. Duke Math. J. **126**(1), 1–52 (2005)
3. Bourbaki, N.: Lie groups and Lie algebras. Chapters 4–6. Elements of Mathematics (Berlin). Springer, Berlin, Translated from the 1968 French original by Andrew Pressley (2002)
4. Brenti, F.: Unimodal, log-concave and Pólya frequency sequences in combinatorics. Mem. Am. Math. Soc. **81**(413), viii+106 (1989)
5. Di Francesco, P.: The solution of the A_r T -system for arbitrary boundary. Electron. J. Combin. **17**(1), R89 (2010). (43)
6. Di Francesco, P.: T -systems, networks and dimers. Comm. Math. Phys. **331**(3), 1237–1270 (2014)
7. Di Francesco, P., Kedem, R.: Positivity of the T -system cluster algebra. Electron. J. Combin. **16**(1), R140 (2009). (39)
8. Di Francesco, P., Kedem, R.: T -systems with boundaries from network solutions. Electron. J. Combin. **20**(1), P3 (2013). (62)
9. Fomin, S., Zelevinsky, A.: Cluster algebras. I. Foundations. J. Am. Math. Soc. **15**(2), 497–529 (2002). (electronic)
10. Fomin, S., Zelevinsky, A.: Y -systems and generalized associahedra. Ann. Math. (2) **158**(3), 977–1018 (2003)
11. Frenkel, E., Reshetikhin, N.: The q -characters of representations of quantum affine algebras and deformations of W -algebras. In Recent developments in quantum affine algebras and related topics (Raleigh, NC, 1998), volume 248 of Contemp. Math., pp 163–205. American Mathematical Society, Providence, RI, (1999)
12. Frenkel, E., Szenes, A.: Thermodynamic Bethe ansatz and dilogarithm identities. I. Math. Res. Lett. **2**(6), 677–693 (1995)
13. Galashin, P., Pylyavskyy, P.: The classification of Zamolodchikov periodic quivers. Am. J. Math. **141**(2), 447–484 (2019)
14. Gliozzi, F., Tateo, R.: Thermodynamic Bethe ansatz and three-fold triangulations. Int. J. Modern Phys. A **11**(22), 4051–4064 (1996)
15. Goncharov, A.B., Kenyon, R.: Dimers and cluster integrable systems. Ann. Sci. Éc. Norm. Supér. (4) **46**(5), 747–813 (2013)
16. Henriques, A.: A periodicity theorem for the octahedron recurrence. J. Algebraic Combin. **26**(1), 1–26 (2007)
17. Inoue, R., Iyama, O., Keller, B., Kuniba, A., Nakanishi, T.: Periodicities of T -systems and Y -systems, dilogarithm identities, and cluster algebras I: type B_r . Publ. Res. Inst. Math. Sci. **49**(1), 1–42 (2013)
18. Inoue, R., Iyama, O., Keller, B., Kuniba, A., Nakanishi, T.: Periodicities of T -systems and Y -systems, dilogarithm identities, and cluster algebras II: types C_r , F_4 , and G_2 . Publ. Res. Inst. Math. Sci. **49**(1), 43–85 (2013)

19. Kazhdan, D., Lusztig, G.: Representations of Coxeter groups and Hecke algebras. *Invent. Math.* **53**(2), 165–184 (1979)
20. Keller, B.: The periodicity conjecture for pairs of Dynkin diagrams. *Ann. Math. (2)* **177**(1), 111–170 (2013)
21. Keller, B., Scherotzke, S.: Linear recurrence relations for cluster variables of affine quivers. *Adv. Math.* **228**(3), 1842–1862 (2011)
22. Kirillov, A.N., Reshetikhin, N.Yu.: Exact solution of the XXZ Heisenberg model of spin S . *Zap. Nauchn. Sem. Leningrad. Otdel. Mat. Inst. Steklov. (LOMI)*, 145(Voprosy Kvant. Teor. Polya i Statist. Fiz. 5):109–133, 191, 195, (1985)
23. Knight, H.: Spectra of tensor products of finite-dimensional representations of Yangians. *J. Algebra* **174**(1), 187–196 (1995)
24. Kuniba, A., Nakanishi, T.: Spectra in conformal field theories from the Rogers dilogarithm. *Modern Phys. Lett. A* **7**(37), 3487–3494 (1992)
25. Kuniba, A., Nakanishi, T., Suzuki, J.: Functional relations in solvable lattice models. I. Functional relations and representation theory. *Int. J. Modern Phys. A* **9**(30), 5215–5266 (1994)
26. Kuniba, A., Nakanishi, T., Suzuki, J.: T -systems and Y -systems in integrable systems. *J. Phys. A* **44**(10), 103001, 146 (2011)
27. Nakajima, H.: t -analogs of q -characters of Kirillov-Reshetikhin modules of quantum affine algebras. *Represent. Theory* **7**, 259–274 (2003). (electronic)
28. Nakanishi, T.: Periodicities in cluster algebras and dilogarithm identities. In *Representations of Algebras and Related Topics*, EMS Ser. Congr. Rep., pp 407–443. Eur. Math. Soc., Zürich, (2011)
29. Ogievetsky, E., Wiegmann, P.: Factorized S -matrix and the Bethe ansatz for simple Lie groups. *Phys. Lett. B* **168**(4), 360–366 (1986)
30. Pylyavskyy, P.: Zamolodchikov integrability via rings of invariants. [arXiv:1506.05378](https://arxiv.org/abs/1506.05378), (2015)
31. Ravanini, F., Valleriani, A., Tateo, R.: Dynkin TBAs. *Int. J. Modern Phys. A* **8**(10), 1707–1727 (1993)
32. Reshetikhin, N.Y.: The spectrum of the transfer matrices connected with Kac-Moody algebras. *Lett. Math. Phys.* **14**(3), 235–246 (1987)
33. Sherman, P., Zelevinsky, A.: Positivity and canonical bases in rank 2 cluster algebras of finite and affine types. *Mosc. Math. J.* **4**(4), 947–974 (2004). 982
34. Speyer, D.E.: Perfect matchings and the octahedron recurrence. *J. Algebraic Combin.* **25**(3), 309–348 (2007)
35. Stekolshchik, R.: Notes on Coxeter Transformations and the McKay Correspondence. Springer Monographs in Mathematics. Springer, Berlin (2008)
36. Stembridge, J.R.: Admissible W -graphs and commuting Cartan matrices. *Adv. Appl. Math.* **44**(3), 203–224 (2010)
37. Szenes, A.: Periodicity of Y -systems and flat connections. *Lett. Math. Phys.* **89**(3), 217–230 (2009)
38. Tao, T.: Why are solitons stable? *Bull. Am. Math. Soc. (N.S.)* **46**(1), 1–33 (2009)
39. Thurston, W.P.: Conway’s tiling groups. *Am. Math. Monthly* **97**(8), 757–773 (1990)
40. Vinberg, È.B.: Discrete linear groups that are generated by reflections. *Izv. Akad. Nauk SSSR Ser. Mat.* **35**, 1072–1112 (1971)
41. Volkov, A.Y.: On the periodicity conjecture for Y -systems. *Commun. Math. Phys.* **276**(2), 509–517 (2007)
42. Zamolodchikov, A.B.: On the thermodynamic Bethe ansatz equations for reflectionless ADE scattering theories. *Phys. Lett. B* **253**(3–4), 391–394 (1991)