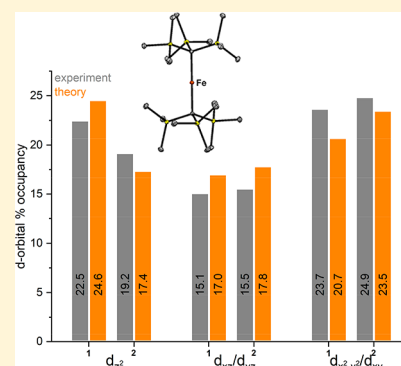


## Insights into Single-Molecule-Magnet Behavior from the Experimental Electron Density of Linear Two-Coordinate Iron Complexes

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## Supporting Information

**ABSTRACT:** A breakthrough in the study of single-molecule magnets occurred with the discovery of zero-field slow magnetic relaxation and hysteresis for the linear iron(I) complex  $[\text{Fe}(\text{C}(\text{SiMe}_3)_2)_2]^-$  (**1**), which has one of the largest spin-reversal barriers among mononuclear transition-metal single-molecule magnets. Theoretical studies have suggested that the magnetic anisotropy in **1** is made possible by pronounced stabilization of the iron  $d_z^2$  orbital due to  $3d_z^2-4s$  mixing, an effect which is predicted to be less pronounced in the neutral iron(II) complex  $\text{Fe}(\text{C}(\text{SiMe}_3)_2)_2$  (**2**). However, experimental support for this interpretation has remained lacking. Here, we use high-resolution single-crystal X-ray diffraction data to generate multipole models of the electron density in these two complexes, which clearly show that the iron  $d_z^2$  orbital is more populated in **1** than in **2**. This result can be interpreted as arising from greater stabilization of the  $d_z^2$  orbital in **1**, thus offering an unprecedented experimental rationale for the origin of magnetic anisotropy in **1**.



## INTRODUCTION

Since the discovery of magnetic blocking in a dodecanuclear manganese cluster more than two decades ago,<sup>1</sup> the search for single-molecule magnets with improved operating temperatures has burgeoned beyond the study of transition metal clusters to encompass complexes of the lanthanides<sup>2</sup> and actinides<sup>3</sup> as well as mononuclear transition-metal compounds.<sup>4–9</sup> Arguably, the original impetus for this expansion arose from both the discovery of slow magnetic relaxation and magnetic hysteresis in mononuclear lanthanide complexes<sup>10</sup> and the realization that magnetic anisotropy is the key factor in dictating slow magnetic relaxation for transition-metal compounds.<sup>11,12</sup> Indeed, contemporary efforts in the design of single-molecule magnets have focused largely on various strategies to exploit and enhance magnetic anisotropy toward the realization of ever larger energy barriers to spin relaxation and high-temperature magnetic hysteresis.<sup>13–15</sup> Magnetic anisotropy is a phenomenon that derives ultimately from the

presence of nonzero orbital angular momentum in the electronic ground state of a magnetic ion, which itself results either from the presence of (quasi) degenerate electronic configurations (a first-order effect)<sup>16</sup> or via mixing with excited electronic states possessing nonzero orbital angular momentum (a second-order effect).<sup>17</sup> Coupling of the spin and orbital angular momentum in the former case<sup>16</sup> removes the degeneracy of the ground state and creates the observed anisotropy, which can also be enhanced by the ligand field.

The magnitude of the spin–orbit coupling increases with atomic number and can reach maximal values with heavy elements such as the lanthanides and actinides. Many trivalent lanthanide ions in particular exhibit both a large spin and unquenched orbital angular momentum that arise from the near-degeneracy of the well-shielded 4f orbitals, and thus

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perhaps unsurprisingly the pursuit of lanthanide-based single-molecule magnets represents one of the most promising avenues to date. Indeed, lanthanide compounds have set recent records for both magnetic relaxation barrier and hysteresis temperature.<sup>13,15,18</sup> Various levels of theory have been developed to predict the magnitude and orientation of the magnetic anisotropy in lanthanide systems and are in some cases quite successful,<sup>19</sup> while experimentally the anisotropy can be quantified using neutron diffraction.<sup>20</sup>

For the transition-metal ions, first-order orbital angular momentum is typically quenched by the ligand field, but a second-order effect can arise when the absolute value of the d orbital splitting is optimized such that excited states exhibiting orbital angular momentum are close in energy to the ground state.<sup>17</sup> Alternatively, a strategy to achieve first-order orbital angular momentum in mononuclear transition-metal complexes is to target linear coordination environments, which ensure the presence of degenerate orbitals that will be unaffected by Jahn–Teller distortions. This approach was first realized in the development of a series of linear, two-coordinate high-spin iron(II) single-molecule magnets, including the complex  $\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2$  (**2**).<sup>8,21–26</sup> However, an applied magnetic field was required in order to observe slow magnetic relaxation for these complexes, due at least in part to pronounced ground-state quantum tunneling of the magnetization associated with the non-Kramers integer-spin iron(II) ion.<sup>8</sup> Notably, successful chemical reduction of **2** was later achieved to generate the linear, two-coordinate compound  $[\text{K}(\text{crypt-222})][\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]$  (**1**),<sup>27</sup> with a half-integer-spin iron(I) center. Compound **1** behaves as a single-molecule magnet under zero dc field and exhibits a relaxation barrier<sup>28,29</sup> of  $246(3) \text{ cm}^{-1}$  as well as magnetic hysteresis to 6 K.<sup>27</sup>

These results can be successfully interpreted within a ligand field model. In a strictly cylindrically symmetric coordination environment (i.e., strictly axial symmetry, with the *z* axis oriented along the axial direction), the d orbitals split into ( $d_{x^2-y^2}$ ,  $d_{xy}$ ), ( $d_{xz}$ ,  $d_{yz}$ ), and ( $d_z^2$ ) sets. In the limit of purely electrostatic crystal field theory, the d orbitals are split depending on the projection of the angular momentum on the *z* axis ( $m_l$ ). Larger absolute values of  $m_l$  are associated with less repulsion due to the axial ligands, and thus these orbitals are lowest in energy. However, symmetry-allowed  $3d_z^2-4s$  orbital mixing may lower the energy of the  $d_z^2$  molecular orbital below that of the  $m_l = \pm 1$  and  $m_l = \pm 2$  d orbital sets. Assuming electron configurations of  $3d^7$  (not  $3d^6 4s^1$ ) and  $3d^6$  for the iron(I) and iron(II) centers, respectively,<sup>30</sup> the d orbital sets for **1** and **2** will necessarily require different ordering to achieve ground states with nonzero orbital angular momentum. In particular, a nonzero orbital angular momentum requires an odd electron count in one of the degenerate sets of d orbitals—i.e., in ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) or ( $d_{xz}$ ,  $d_{yz}$ ). For this scenario to occur, the  $d_z^2$  orbital in **1** must be lowest in energy, while in **2** it must either be the highest in energy or between the ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) and ( $d_{xz}$ ,  $d_{yz}$ ) sets.

Ab initio ligand field theory<sup>31</sup> calculations have predicted that the  $d_z^2$  orbital is indeed lowest in energy for **1**,<sup>27</sup> whereas the ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) set is predicted to have the lowest energy for **2**,<sup>24</sup> providing rationale for the reported magnetic relaxation behavior of both compounds.<sup>27,8</sup> However, direct experimental evidence for this unexpected energy ordering of the d orbitals has remained lacking. Spectroscopic approaches only provide information about the relative energies of the different states but do not reveal the exact nature of the involved states. In

contrast, X-ray diffraction probes the electronic spatial distribution, and it is in principle possible to use this technique to reconstruct a full and complete description of the electron density. Recently, some of us reported the first application of this technique to a single-molecule magnet.<sup>32</sup> Here, we introduce the use of the experimental electron density ( $\rho$ )—derived from accurate high-resolution, low-temperature single-crystal X-ray diffraction data<sup>33</sup>—to explain the magnetic anisotropy in **1** and **2**. In this work, we have chosen to use a multipole-based reconstruction of the experimental electron density, which allows for unambiguous and direct estimation of d orbital occupancies.<sup>34</sup>

## EXPERIMENTAL AND COMPUTATIONAL METHODS

**General Methods.** The syntheses of **1** and **2** were performed as reported in the literature.<sup>25,27</sup>

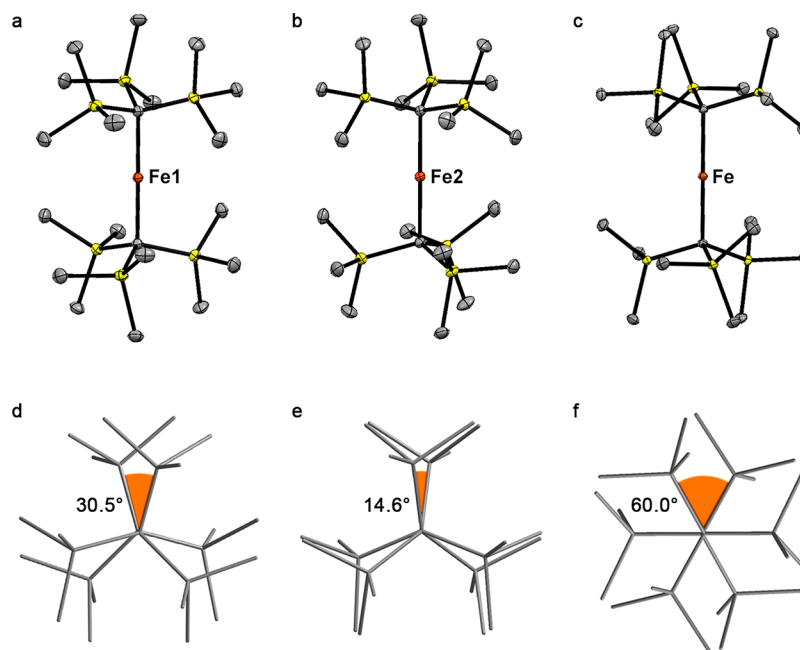
**Single-Crystal X-ray Diffraction.** Single-crystal X-ray diffraction data for **1** were collected using an Oxford Diffraction Supernova instrument equipped with a Mo X-ray source at Aarhus University, an Atlas charge-coupled device detector, and a four-circle goniometer. The crystal was cooled to 100(1) K using an Oxford Cryosystems liquid nitrogen Cryostream device. The data collection consisted of 26  $\omega$  scans. The unit cell parameters were determined, and the Bragg intensities were integrated using CrysAlisPRO.<sup>35</sup> The intensities were empirically corrected for absorption using SCALE3 ABSPACK implemented in CrysAlisPRO.<sup>35</sup> Single-crystal X-ray diffraction data for **2** were measured at the beamline BL02B1 at the Spring-8 synchrotron in Japan. The diffractometer at BL02B1 is equipped with a large curved imaging plate detector with automatic readout mounted on a three-circle goniometer. The crystal was kept in a He gas stream ( $\sim 50 \text{ K}$ ), and a set of 30 frames was collected. The unit cell parameters were determined and the Bragg intensities were integrated using the software RAPID-AUTO.<sup>36</sup> An empirical absorption correction was applied. Equivalent reflections were averaged using SORTAV<sup>37,38</sup> for both **1** and **2**. Further details regarding the data collection and reduction are included in the Supporting Information. The structures of **1** and **2** were solved with direct methods using SHELXT<sup>39</sup> via the Olex2 interface,<sup>40</sup> and the refinement was carried out by the least-squares method using SHELXL-97<sup>41</sup> in Olex2.<sup>40</sup>

**Electron Density Analysis.** The multipole models of the electron density distributions in **1** and **2** were refined in XD2006/XD2015<sup>42,43</sup> using the atom-centered Hansen–Coppens formalism.<sup>44</sup> This model uses density-normalized spherical harmonics,  $d_{lm\pm}$ , to describe the angular density redistribution, as shown in eq 1. The population parameters,  $P_v$  and  $P_{lm\pm}$ , and the radial expansion/contraction parameters,  $\kappa$  and  $\kappa'_l$ , are refined to obtain the final experimental electron density.

$$\rho_{\text{atom}}(\mathbf{r}) = P_c \rho_c(r) + P_v \rho_v \kappa^3(\kappa r) + \sum_{l=0}^{l_{\text{max}}} \kappa_l'^3 R_l(\kappa'_l, r) \sum_{m=0}^l P_{lm\pm} d_{lm\pm}(\theta, \varphi) \quad (1)$$

The intimate relationship between the spherical harmonic functions used in the multipole model and the hydrogenic orbitals means that it is straightforward to derive experimental d orbital occupancies by a linear transformation of the multipole population parameters. Further details regarding the multipole modeling for **1** and **2** are included in the Supporting Information. A topological analysis<sup>45</sup> of the multipole model electron density in **1** and **2** was carried out with the XDPROP module in XD.<sup>42,43</sup> This module was also used to calculate experimental d orbital occupancy estimates for the iron atoms in **1** and **2** from the multipole populations, according to the procedure from Holladay et al.<sup>34</sup>

**Ab Initio Calculations.** The electronic energy levels as well as relativistic and nonrelativistic electron densities of sites Fe1 and Fe2 in **1** and the single Fe site in **2** were calculated on the basis of the



**Figure 1.** Structural data for compounds **1** and **2**. (a–c) Structures of the unique  $[\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]^-$  complexes in the asymmetric unit of **1** (a, b) and the  $\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2$  molecular unit (twice the asymmetric unit) of **2** (c), as determined from the experimental multipole models. Thermal ellipsoids are rendered at the 50% probability level. Gray, yellow, and orange ellipsoids represent C, Si, and Fe atoms, respectively; H atoms have been omitted for clarity. (d–f) Wireframe structures of the two independent  $[\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]^-$  moieties of **1** (d, e) and the  $\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2$  molecular unit of **2** (f), illustrating the dihedral angles between Si–C–C\* and C–C\*–Si\* planes as described in the text (average values of 30.5, 14.6, and  $\sim 60.0^\circ$ , respectively, are illustrated in d–f; carbon atoms used to define the planes are those bonded to iron). The Fe–C bonds in these structures are oriented roughly orthogonal to the plane of the text, and hydrogen atoms are omitted for clarity.

complex geometries from the crystallographic CIF files. State average complete active space self-consistent field (CASSCF) calculations were performed, and dynamic correlation was included using the *N*-electron valence perturbation theory (NEVPT2) method. All 10 spin-free quartet states of **1** (active space consisting of 7 electrons distributed over the 5 3d molecular orbitals, CAS(7,5)) and the 5 quintet states of **2** (active space consisting of 6 electrons distributed over the 5 3d molecular orbitals) were taken into account (see the [Supporting Information](#) for input files and selected output results). Inclusion of states of lesser multiplicity (specifically, 40 doublets for **1** and the 45 triplets and 50 singlets for **2**) did not essentially change the results. Nonrelativistic and relativistic (including scalar and spin-orbit effects) densities used in the analysis were read from the correspondingly adjusted “\*.gbw” files (adjusted for spin-orbit sublevels of the  $^4\text{E}$  (**1**) and  $^5\text{E}$  (**2**) states). The relativistic and nonrelativistic densities lead to essentially the same results, supporting the observation that magnetic sublevels for each complex share essentially the same spatial part of the many-electron wave functions. All calculations were done with a development version of the ORCA program package<sup>46,47</sup> that was appropriately adjusted for the purpose of the present project.

## RESULTS

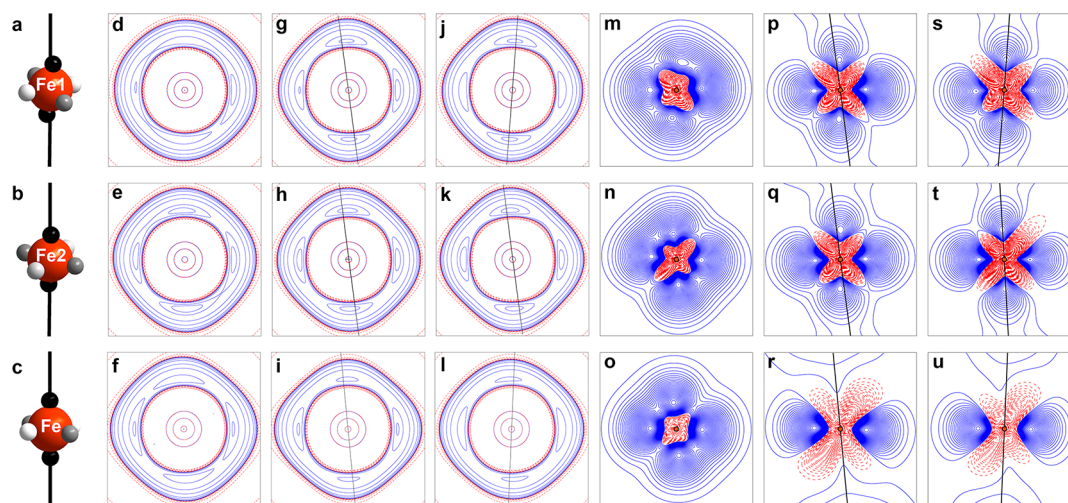
**Structural Results from Multipole Models.** The complex anion of **1** and the molecular unit of **2** derived from the experimental multipole models are shown in [Figure 1a–c](#). Surprisingly, the unit cell found for **1** is twice as large as that reported previously<sup>27</sup>—although the space group symmetry remains triclinic  $P\bar{1}$ , the number of independent molecules is doubled such that there are two geometrically distinct  $[\text{K}(\text{crypt-222})][\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]$  units in the asymmetric unit. In each unique iron complex of **1**, the  $[\text{C}(\text{SiMe}_3)_3]^-$  ligands are nearly eclipsed, and their relative orientations can be described by the dihedral angles between Si–C–C\* and C–C\*–Si\* planes (the carbon atoms are those

bonded to iron), which range from  $30.1$  to  $30.9^\circ$  in  $[\text{Fe1}(\text{C}(\text{SiMe}_3)_3)_2]^-$  and  $14.4$  to  $14.8^\circ$  in  $[\text{Fe2}(\text{C}(\text{SiMe}_3)_3)_2]^-$  ([Figure 1d,e](#)). The average of these ranges is notably close to the range determined for the previously reported structure ( $22.6^\circ$  versus  $21.8$ – $22.2^\circ$ , respectively).<sup>27</sup> In compound **2**, the  $[\text{C}(\text{SiMe}_3)_3]^-$  ligands are in a staggered conformation with much larger dihedral angles in the range from  $57.8$  to  $61.3^\circ$  ([Figure 1f](#)). While the C–Fe–C angle in **2** is strictly linear due to symmetry (space group  $C2/c$ ), the C–Fe–C angles in **1** are nearly linear at  $179.09(3)$  and  $179.38(3)^\circ$ .

The bond distances in **1** and **2** are also distinct, highlighting the different iron formal oxidation states in each complex—in  $[\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]^-$ , the iron formal oxidation state is +1 and the Fe–C distances range from  $2.0610(7)$  to  $2.0653(7)$  Å, whereas the iron formal oxidation state in  $\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2$  is +2 and the Fe–C distance is  $2.0535(4)$  Å. Integration of the experimental multipole model electron densities within atomic basins derived from Bader’s quantum theory of atoms in molecules<sup>45</sup> afforded iron atomic charges of approximately +0.6 and +1.4 in **1** and **2**, respectively. These values confirm that iron is more positive in **2** than in **1**, as expected, although they are smaller than the iron formal oxidation states in each complex.

**Laplacian and Deformation Electron Density Distributions.** The Laplacian of the electron density,  $\nabla^2\rho$ , is the sum of the diagonal elements of the Hessian matrix—obtained by taking the second derivative of  $\rho$  in the *x*, *y*, and *z* directions—and is used to highlight local concentrations ( $\nabla^2\rho < 0$ ) or depletions ( $\nabla^2\rho > 0$ ) of electron density. Three-dimensional Laplacian distributions for **1** and **2**, obtained from the experimental multipole models, revealed six significantly negative minima, located  $\sim 0.32$  Å from each unique iron atom (i.e., in the iron M shell<sup>48</sup>). These minima roughly form an





**Figure 2.** Electron density based distributions in **1** and **2** from the experimental multipole models. (a–c) Locations of the Laplacian minima (charge concentrations) determined for the iron atoms in **1** (a,b) and **2** (c) marked as white, gray, and black spheres. (d–l) Laplacian contour maps with a width of 1 Å for iron atoms in **1** and **2**. (m–u) Static model deformation electron density maps with a width of 2.2 Å for iron atoms in **1** and **2**. The planes in these maps are defined from the positions of the minima in a–c. The maps in the first column (d–f, m–o) include white and gray spheres, the maps in the second column (g–i, p–r) include the black and white spheres, and the maps in the third column (j–l, s–u) include the black and gray spheres. Solid blue lines represent positive contours of  $-\nabla^2\rho$  (local charge concentration) and the static model deformation electron density. Dashed red lines represent negative contours of  $-\nabla^2\rho$  (local charge depletion) and the static model deformation electron density. Laplacian contours are drawn at  $\pm 2 \times 10^n$ ,  $4 \times 10^n$ , and  $8 \times 10^n$  eÅ $^{-5}$ , where  $n$  is an integer ranging from  $-3$  to  $+3$ . An additional contour is drawn at  $-\nabla^2\rho = 1000$  e/Å $^3$  to emphasize the different numerical values of the Laplacian minima. The static model deformation electron density contours in m–u are drawn in increments of 0.05 e/Å $^3$ .

octahedron around each iron atom (Figure 2a–c), with axial corners located along the C–Fe–C direction (i.e., close to the local  $z$  axis on iron) and equatorial corners positioned near the plane orthogonal to this direction (i.e., close to the local iron  $xy$  plane).<sup>49</sup> Relative to the C–Fe–C direction, the octahedra are tilted slightly by 6–8° (as seen in Figure 2a–c). The Laplacian minima for each complex are quite distinct, as illustrated by two-dimensional Laplacian contour maps corresponding to selected planes containing an iron and four of the six minima (Figure 2d–l). For compound **1**, the four minima positioned close to the local iron  $xy$  plane vary between  $-1084$  and  $-965$  e/Å $^5$  for Fe1 and between  $-1128$  and  $-1083$  e/Å $^5$  for Fe2. The two minima near the C–Fe–C direction are  $-1045$  and  $-1040$  e/Å $^5$  for Fe1 and  $-1067$  e/Å $^5$  for Fe2. Thus, the four “equatorial” Laplacian minima are quite similar to the two “axial” minima in **1**, differing by only 4–8% and indicating that the concentrations of electron density are comparable along all three directions. In contrast, for **2** the four “equatorial” minima (values of  $-956$  and  $-1021$  e/Å $^5$ ) differ from the two “axial” minima ( $-851$  e/Å $^5$ ) by 12–20%.

Contour maps of the static model deformation electron density surrounding the iron atoms in **1** and **2** are also depicted in Figure 2m–u, in the same planes as those presented for the Laplacian maps. The deformation density maps illustrate the difference between the multipole-expressed electron densities and hypothetical, independent atom model densities, wherein the atoms are represented by spherical isolated atom densities. Thus, these maps highlight regions where redistribution due to bonding has taken place. Within the framework of the ligand field model used here, the Laplacian and deformation density maps indicate a depletion of electron density in the  $d_{xz}$  and the  $d_{yz}$  orbitals for **1** and instead a preferential and similar filling of the  $d_{x^2-y^2}$ ,  $d_{xy}$ , and  $d_z^2$  orbitals. For **2**, the maps indicate a preferential filling of the  $d_{x^2-y^2}$  and  $d_{xy}$  orbitals.

**Experimental d Orbital Occupancies.** Experimental d orbital occupancies can be estimated from the multipole populations in the multipole model of the electron density, according to Holladay et al.<sup>34</sup> (see **Experimental and Computational Methods** for more details). It is important to emphasize that the angular component of the multipoles—and thus the subsequently derived d orbital occupancies—relate to user-defined local atomic coordinate systems and are therefore unrelated to the unit cell axes. Here, we define the local  $z$  axis to point along the respective Fe–C directions for all iron atoms; however, for the  $x$  and  $y$  directions within the  $xy$  plane, there is no clear optimal choice for these linear compounds.<sup>50</sup> The extraction of d orbital occupancies assumes that the electron density near the Fe site is primarily due to the atomic orbitals of Fe: i.e., that there is vanishing Fe–C electronic overlap. Although this approximation is somewhat simplistic, it is justified in this scenario where the ligand field is weak for both **1**<sup>27</sup> and **2**<sup>8</sup> and the bonding between the ligands and the central Fe ion is primarily ionic in nature. The experimental d orbital occupancies resulting from using this procedure are given in Table 1 for the three unique Fe atoms in **1** and **2**. The sum of the occupancies of the nearly degenerate  $d_{x^2-y^2}$  and  $d_{xy}$  orbitals is invariant with respect to rotation of the  $x$  and  $y$  axes

**Table 1.** d Orbital Occupancies Obtained from the Experimental Multipole Populations of **1** and **2**<sup>a</sup>

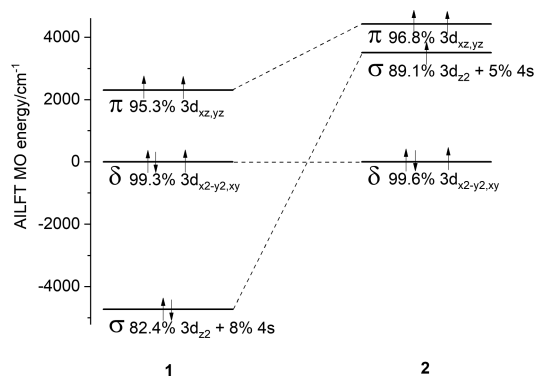
		( $d_{xz}$ , $d_{yz}$ )	( $d_{x^2-y^2}$ , $d_{xy}$ )	$d_z^2$
<b>1</b>	Fe1	1.05(2) 15.5(2)	1.58(2) 23.3(2)	1.52(3) 22.4(4)
	Fe2	1.00(2) 14.8(2)	1.62(2) 23.9(2)	1.53(3) 22.6(4)
<b>2</b>	Fe	0.98(1) 15.5(1)	1.57(1) 24.9(1)	1.21(1) 19.2(2)

<sup>a</sup>Occupancy averages are reported for nearly degenerate ( $d_{xz}$ ,  $d_{yz}$ ) and ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) pairs of orbitals. Absolute occupancies and percent occupancies of the valence (“d”) electron count are given (latter in italics).

around the  $z$  axis, but the relative occupancies do depend on the orientation of these axes in the  $xy$  plane. The same observation is true for the  $d_{xz}$  and  $d_{yz}$  orbitals. Thus, only the average occupancies are reported for these pairs in Table 1.

For **1**, the average occupancy of the ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) orbitals and the occupancy of the  $d_z^2$  orbital are quite similar (22.4(4)–23.9(2)%) and are significantly higher than the average occupancy of the ( $d_{xz}$ ,  $d_{yz}$ ) orbitals (14.8(2)–15.5(2)%). For **2**, the average occupancy of the ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) orbitals is the highest (24.9(1)%) followed by the intermediate occupancy of the  $d_z^2$  orbital (19.2(2)%) and the lowest average occupancy of the ( $d_{xz}$ ,  $d_{yz}$ ) orbitals (15.5(1)%). Interestingly, although the two  $[\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]^-$  moieties in **1** differ with respect to the relative orientation of their  $[\text{C}(\text{SiMe}_3)_3]^-$  ligands, their  $d$  orbital occupancies are identical within a few standard deviations. The  $d$  orbital occupancy estimates from the multipole populations do not provide direct information about the quantitative ordering of the  $d$  orbital energies. However, assuming the Aufbau principle is fulfilled, the experimentally derived  $d$  orbital occupancies obtained for **1** clearly support a model that requires that the  $d_z^2$  orbital is significantly stabilized relative to the ( $d_{xz}$ ,  $d_{yz}$ ) orbitals, in contrast to their anticipated relative energies based on crystal field theory. For **2**, the experimentally derived occupancies suggest less stabilization of the  $d_z^2$  orbital than in **1**.

**Theoretical Electronic Structure Analysis.** Energies and populations of the 3d-like molecular orbitals derived from an ab initio ligand field theory (AILFT) analysis of the CASSCF/NEVPT2 results are reported in Figure 3. These results



**Figure 3.** Energies and populations of the  $d$  orbitals in **1** and **2**. The results are obtained from a ligand field analysis of the ab initio calculations. Selected 3d and 4s atomic orbital contributions to the 3d-like molecular orbitals are given in percent.

confirm the previous finding<sup>24,27</sup> that the existence of magnetic blocking in **1** and **2** can be explained by the large differences in their relative  $d$  orbital energies (i.e., that the  $d_z^2$  orbital is significantly more stabilized in **1** than in **2** and that only the population of the  $d_z^2$  orbital is expected to change between the two compounds).<sup>51</sup> Selected atomic orbital contributions are included in Figure 3. The large percentage of 3d character in the molecular orbitals illustrates that the free ion approximation is well suited for the analysis. The percentage of 4s character in the  $\sigma$  ( $d_z^2$ -like) molecular orbital is 8% for **1** and less (5%) for **2**, i.e., the larger stabilization of the  $d_z^2$  molecular orbital in **1** than in **2** is associated with more  $3d_z^2$ –4s mixing in **1** than in **2**.

**Theoretical  $d$  Orbital Occupancies.** The  $d$ -orbital populations derived using a ligand field analysis of ab initio

results are not necessarily quantitatively comparable to the experimental  $d$  orbital occupancies derived in this work from a multipole-based reconstruction of the total electron density. To enable such a comparison, we used ab initio calculations to compute structure factors representing the molecular electron density and used these calculated values to refine multipole population parameters—again defining the local  $z$  axis along the Fe–C directions for all iron atoms. In this manner, we obtained theoretical  $d$  orbital occupancies (Table 2) that are

**Table 2.**  $d$  Orbital Occupancies Obtained from the Theoretical Multipole Populations of **1** and **2**<sup>a</sup>

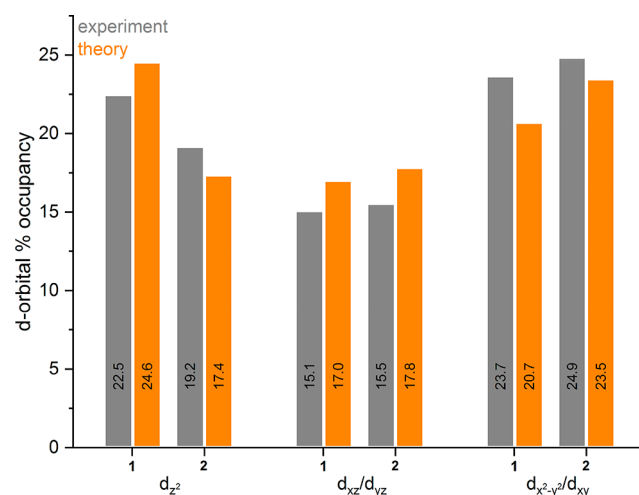
	( $d_{xz}$ , $d_{yz}$ )	( $d_{x^2-y^2}$ , $d_{xy}$ )	$d_z^2$
Fe1 ( <b>1</b> )	1.09 17.0	1.33 20.7	1.57 24.6
Fe ( <b>2</b> )	1.14 17.8	1.50 23.5	1.11 17.4

<sup>a</sup>Occupancy averages are reported for the nearly degenerate ( $d_{xz}$ ,  $d_{yz}$ ) and ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) pairs of orbitals. Absolute occupancies and percent occupancies of the valence (“ $d$ ”) electron count are given (latter in italics). Theoretical structure factors were calculated for only one of the two independent iron moieties in **1**.

directly comparable to the experimental occupancies given in Table 1. The details of the structure factor calculations and subsequent multipole modeling from ab initio calculations are given in the Supporting Information. The theoretical  $d$  orbital occupancies in Table 2 obtained from the theoretical multipole model populations clearly reflect the  $d$  orbital energy ordering and populations in Figure 3 resulting from AILFT approaches. This suggests that the procedure used here to extract  $d$  orbital occupancies from a multipole expansion of the electron density is highly reliable and is a viable alternative to AILFT analysis.

## DISCUSSION

A comparison of the experimental and theoretical  $d$  orbital occupancies determined from multipole parametrization of the electron density is given in Figure 4. Overall, the two methods are in reasonably good agreement and differ within 1–3% for both compounds. For **1**, theory predicts a larger accumulation



**Figure 4.**  $d$  orbital occupancies in **1** and **2** from experiment and theory. Occupancies (in percentages of the valence (“ $d$ ”) electron count) are derived from multipole population parameters using experimental (gray) and theoretical (orange) structure factors. Occupancy averages are reported for the nearly degenerate ( $d_{xz}$ ,  $d_{yz}$ ) and ( $d_{x^2-y^2}$ ,  $d_{xy}$ ) pairs of orbitals.

of electron density along the local  $z$  axis than in the  $xy$  plane, whereas experimentally these regions are found to have very similar densities. For **2**, theory predicts similar occupancies of the  $d_{z^2}$  and  $(d_{xz}, d_{yz})$  orbitals, whereas experimentally the  $d_{z^2}$  orbital has a higher occupancy than the  $(d_{xz}, d_{yz})$  orbitals. Importantly, the  $d_{z^2}$  orbital occupancy is found to be 3–7% higher in **1** than in **2** using both approaches, and thus the experimental results support the present and previous finding from theory<sup>24,27</sup> that the  $d_{z^2}$  orbital is substantially more stabilized in **1** than in **2**.

## CONCLUSION

In summary, we have used single-crystal X-ray diffraction data to determine the experimental electron density of the two linear iron compounds  $[\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]^-$  and  $\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2$  and analyzed the data to reveal the existence of fundamentally different  $d$  orbital occupancies in the two compounds. Spatial partitioning of the density also enabled demonstration of the different oxidation states and  $d$  electron counts in both compounds. The combined data for each compound importantly provides experimental evidence for the origin of the large magnetic anisotropy in  $[\text{Fe}(\text{C}(\text{SiMe}_3)_3)_2]^-$  and the corresponding exceptional slow magnetic relaxation behavior. An obvious extension of this method is to the study of electronic density distribution and magnetic anisotropy in lanthanide-based single-molecule magnets. Although the origin of magnetic anisotropy in these systems is different from that of transition metals, the asymmetry in  $4f$  orbital populations arising from the ligand field will directly reveal the ground  $M_J$  state. In addition, we will be able to quantify directly the degree of oblateness or prolateness of the valence electrons, which plays a particular role in the design of single-molecule magnets of the lanthanide ions.<sup>52</sup> Work along these lines is underway.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.inorgchem.8b03301.

Details about data collection, data reduction, multipole modeling, theoretical calculations, and theoretical structure factors (PDF)

### Accession Codes

CCDC 1882141–1882142 contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via [www.ccdc.cam.ac.uk/data\\_request/cif](http://www.ccdc.cam.ac.uk/data_request/cif), or by emailing [data\\_request@ccdc.cam.ac.uk](mailto:data_request@ccdc.cam.ac.uk), or by contacting The Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: +44 1223 336033.

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## Notes

The authors declare no competing financial interest.

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(49) In the case of exact cylindrical symmetry, there would obviously not be any extrema in the Laplacian (or the density) in the  $xy$  plane. The extrema that we observe result from the influence of deviations from the exact symmetry.

(50) It is obvious to choose the minima in the Laplacian distribution in the  $xy$  plane as defining the directions, but this does not minimize the cross terms which otherwise are good indicators for a good choice of axes. In addition, the sum of the populations in  $d_{xy}$  and  $d_{x^2-y^2}$  and in  $d_{xz}$  and  $d_{yz}$  are constant upon rotation of the  $xy$  axes around the  $z$  direction. We will therefore in the following use the sum of these, instead of individual  $d$  orbital populations.

(51) The quantitative differences between the energies reported in Figure 3 and those previously reported (refs 24 and 27) are mainly assigned to the changes in the CASSCF and NEVPT2 routines and the implementation of the ALLFT method in the updated version of the ORCA program package (refs 46 and 47) utilized in the present study. The geometrical differences between structures of compounds

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