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Effect of polar side chains on neutral and p-doped polythiophene

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Incorporation of polar side chains on organic semiconducting materials have been used recently in thermoelectric materials to increase dopant:semiconductor miscibility and stability to further increase the performance and durability of devices. However investigations into how polar side chains can affect the structure and energetics of polythiophenes compared to non-polar alkyl side chains are usually carried out using materials with no common morphological structure. Within this work we systematically investigate the increase in polar side chain content on poly(3-hexylthiophene) (P3HT) and how the optical, electrochemical, and structural properties are affected. We find a decreasing degree of aggregation with increasing polar side chain content leading to lower charge carrier mobilities. Upon doping with 2,3,5,6-tetrafluoro-7,7,8,8-tetracyanoquinodimethane (F4TCNQ), we find that the electrical conductivity is reduced when incorporating the polar side chain and no stabilising effect is demonstrated when annealing the doped thin films at raised temperatures. This study emphasises that polar functionalities do not always increase dopant:semiconductor interactions and can harm desirable structural and electrical characteristics, and therefore should be incorporated into organic semiconductors with caution.

1. Introduction

Poly(3-hexylthiophene) (P3HT), shown in Figure 1, is a widely studied organic semiconductor in the field of organic electronics due to its ease of synthesis, commercial availability, solution processability, structural organisation and good charge transport properties. P3HT, a hole transporting semiconducting polymer with a highly regioregular structure, has been used in research for many applications including: organic bioelectronics, organic photovoltaics, organic field effect transistors (OFETs) and organic thermoelectric generators.¹⁻⁶ The performance of organic semiconductors in electronic devices is often improved by molecular doping which introduces additional positive (p-type doping) or negative (n-type doping) charge carriers.⁷⁻¹¹ A well-researched system for molecular doping is P3HT in combination with the oxidant 2,3,5,6-tetrafluoro-7,7,8,8-tetracyanoquinodimethane (F4TCNQ) depicted in Figure 1. F4TCNQ has a deep lowest unoccupied molecular orbital (LUMO) which makes it energetically favourable to remove an electron from the shallower highest occupied molecular orbital (HOMO) of P3HT

thereby introducing additional positive charge carriers in the system. High conductivities have been achieved in this system by different processing techniques, including highly oriented polymer films, vapour doping and sequential doping from solution.¹²⁻¹⁵ Nevertheless, there is still an issue with dopant stability, since F4TCNQ can diffuse in the solid state, in particular at elevated temperatures, and drift due to applied electric fields, potentially reducing device lifetimes.^{16,17} Recently, molecular engineering of semiconductors and dopants has been employed as a tool to control dopant diffusion and polymer:dopant interactions. One well-explored strategy is to introduce ethylene glycol motifs into the polymer side chain.¹⁸⁻²⁴ These glycol groups increase the local dielectric constant, stabilizing charged species formed upon doping, and therefore enhancing dopant stability at elevated temperatures. In n-type materials this strategy has been particularly effective, yielding dramatic improvements in conductivity and thermoelectric performance, presumably due to the favourable interactions between glycol groups and cations.^{20,21}

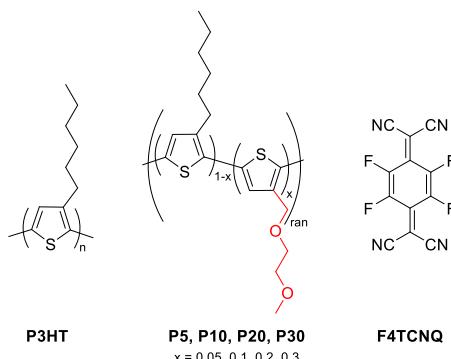


Figure 1. Structure of P3HT (left), P5-P30 (middle) and F4TCNQ (right).

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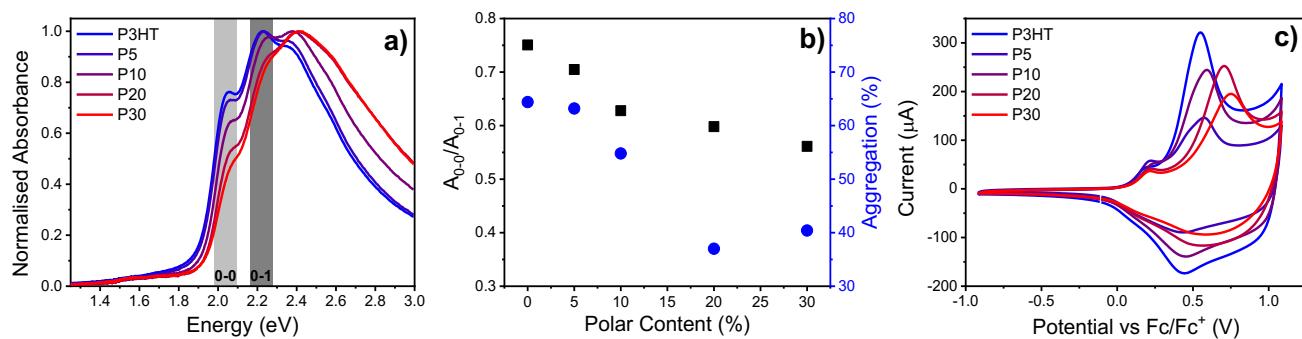


Figure 2. a) Normalised thin film UV-Vis spectra of polymer series spin-cast from 10 mg ml^{-1} ODCB on quartz slides at 2000 rpm for 90 seconds. Greyed areas indicate the 0-0 and 0-1 vibronic transitions. b) Plot showing degree of crystallinity and $A_{0,0}/A_{0,1}$ calculated from the fittings of the normalised thin film UV-Vis spectra versus polar content of co-polymers. c) Thin film cyclic voltammograms of drop-cast polymer series vs Fc/Fc^+ (0.1 M TBAPF_6 as supporting electrolyte in acetonitrile) at a scan rate of 50 mV s^{-1} .

In the work presented herein, to investigate the role of polar side chains in the context of molecular doping of p-type organic semiconductors, we have synthesised and characterised a series of P3HT derivatives with increasing content of ethylene glycol based polar side chains as shown in Figure 1. P3HT is chosen as the reference system due to its well-studied properties both in its neutral and doped states. The polar side chain contains one ethylene glycol unit separated from the backbone by one carbon atom in order to maintain a six-atom long side chain, and simultaneously diminish resonance effects from the oxygen atoms that will otherwise change frontier molecular orbital energies and doping behaviour drastically.^{25,26} The random co-polymers were synthesised with 5, 10, 20 and 30 % polar side chain content, thus affording **P5**, **P10**, **P20** and **P30** respectively, to systematically cover polar side chain content matching typical doping concentrations used for organic semiconductors. The polymers were subsequently doped with F4TCNQ via sequential doping methods to investigate the electrical properties and thermal stability.¹³

2. Results and Discussion

2.1. Polymer Synthesis

P3HT and the random co-polymers were synthesised using the Grignard metathesis method, a well-known polymerisation reaction to produce regioregular P3HT.^{27,28} The full synthetic method is shown in the supporting information. The general synthesis of the random co-polymers includes trans-metallation of the two 2,5-dibrominated thiophene monomers using $i\text{PrMgCl} \cdot \text{LiCl}$ in separate flasks, then after combining, using $\text{Ni}(\text{dppp})\text{Cl}_2$ as the catalyst to afford the polymers in 56-77% yield. Nuclear magnetic resonance spectroscopy confirmed that the monomer feed ratio matches the monomer ratio in the co-polymers as shown in the supporting information (Table S1). As shown in Table 1 and supporting information (Fig S11), the molecular weights of the polymers determined by gel permeation chromatography in chloroform are similar across the series. Thermogravimetric analysis reveals two thermal

degradation events around 300 and 450 °C. The lower temperature event increases in weight percentage loss with increasing polar side chain content indicating that the alkyl side chain is more thermally stable than the ethylene glycol based side chain, as shown in the supporting information (Fig S12 and Table S2). Differential scanning calorimetry of P3HT reveals melting and crystallisation transitions at 226 and 192 °C respectively, however no strong transitions are observed upon incorporating polar side chains (Fig S13, S14 and Table S2).

2.2. Optical, Electrochemical and Structural Properties

The UV-Vis spectra of the polymers dissolved in o-dichlorobenzene (ODCB) show a gradual blue-shift in the absorption maximum (λ_{max}) of ~9nm with increasing polar content, as summarised in Table 1 and the supporting information (Fig S15). UV-Vis spectra of the polymer thin films spin-cast from ODCB reveal the degree of structural order in the solid state, as shown in Figure 2a. P3HT shows the distinct $\pi-\pi^*$ transition with a λ_{max} at 2.23 eV (557 nm) and a shoulder at 2.06 eV (602 nm) which are assigned the 0-1 and 0-0 vibronic transitions respectively.²⁹ Addition of polar sidechains causes a gradual blue-shift and change in absorption of the 0-1 and 0-0 transitions. Spano calculated that these transitions can be related to J- and H-like aggregates, where the 0-0 absorption feature is ascribed to planarity of the polymer backbone within aggregates.²⁹⁻³¹ Therefore a decrease in the ratio between the absorption height of the first two vibronic peaks ($A_{0,0}/A_{0,1}$) correlates to increased disorder within the aggregates. When incorporating the polar side chain, a gradual blue-shift is observed, along with a significant decrease in $A_{0,0}/A_{0,1}$, suggesting an increase in disorder within the polymer aggregates as shown in Figure 2b.

The solid-state UV-Vis spectra were fitted using four Gaussian curves representing the aggregate regions and the experimental solution spectrum representing the amorphous regions as detailed in the supporting information (Fig S16). From these fittings, a relative degree of aggregation of 64 % was

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Table 1. Optical and electrochemical properties and molecular weight data.

Polymer	M_w/M_n [kg mol ⁻¹] ^a (PDI)	Solution λ_{max} [nm] ^b	Thin film λ_{max} [nm] ^c	E_g^{opt} [eV] ^d	$A_{0-0} \lambda_{max}$ [eV] ^e	$E_{1/2}^{(ox)}$ [V] ^f vs Fc/Fc ⁺
P3HT	48.8/38.0 (1.28)	462	557	1.90	2.05	0.50
P5	46.2/33.2 (1.39)	460	555	1.91	2.05	0.51
P10	53.2/39.1 (1.36)	456	521	1.92	2.06	0.53
P20	58.2/40.5 (1.44)	453	515	1.92	2.07	0.64
P30	47.3/33.6 (1.41)	453	516	1.93	2.08	0.67

^aMeasured using gel permeation chromatography in chloroform at 40 °C vs. polystyrene standards. ^bMeasured from 10 µg ml⁻¹ ODCB solutions. ^cThin films were spin-cast from 10 mg ml⁻¹ ODCB solutions onto quartz slides at 2000 rpm for 90 seconds. ^dCalculated from thin film λ_{onset} using: E [eV] = 1240/ λ [nm]. ^eCalculated from fittings of normalised thin films UV-Vis spectra. ^fCV recorded vs Ag/Ag⁺ (Ag/Ag⁺ - 0.115 V = Fc/Fc⁺) of polymer drop-cast onto carbon electrode (0.1 mg ml⁻¹ CHCl₃ solution) using 0.1 M TBAPF₆ in acetonitrile as supporting electrolyte at 50 mV s⁻¹. $E_{1/2}^{(ox)}$ = (E_{pc} + E_{pa})/2.

estimated for P3HT, with slightly lower values of 63 and 55 % for **P5** and **P10** respectively, dropping further to around 40 % for **P20** and **P30** as depicted in Figure 2b. Analysis of the Gaussian fitting curves representing the 0-0 and 0-1 transitions reveals that A_{0-0}/A_{0-1} decreases with increasing polar side chain content, suggesting that incorporation of polar side chains decreases the planarity of the polymer backbone, as discussed earlier. Also shown in Table 1, the onset of absorption (E_g^{opt}) and A_{0-0} transition energy both blue shift 0.03 eV upon incorporating 30 % polar content indicating a slight increase in band gap energy.

The thin film electrochemical properties examined with cyclic voltammetry show a clear and reversible oxidative process for all polymers in the series, as depicted in Figure 2c and Table 1. The half-wave potential increases by 0.17 V going from P3HT to **P30**, which we ascribe to the decrease in aggregation with increasing polar content. The electron withdrawing inductive effect from the oxygen atom removed one carbon from the polymer backbone is also expected to contribute to the shift in half-wave potential observed with increasing polar content.²⁵ The major oxidative peak is for all polymers preceded by a minor oxidative event around 0.1 – 0.2 V which we ascribe to the oxidation of highly ordered crystalline regions; the gradual disappearance of this minor peak with increasing scan rate as depicted in Figure 2c and S18 is in agreement with relatively slow ion diffusion into crystalline regions, as would be expected. Interestingly the ascribed crystalline oxidation peak remains at similar potentials across the series, however the major oxidative peak shifts to higher potential, indicating that the polar side chain does not affect the microstructure in a uniform manner.

Density functional theory calculations of gas phase decamers reveal a larger dihedral angle between adjacent thiophene units upon incorporation of polar side chains compared to the corresponding decamer with only alkyl chains,

as illustrated in Figure 3 and the supporting information (Fig S19-S21). This suggests that the increased disorder with increasing polar content is not only a solid-state effect due to poor packing, but an intrinsic attribute of the mixed alkyl/glycol polymer system, in agreement with the blue-shifted solution UV-Vis data.

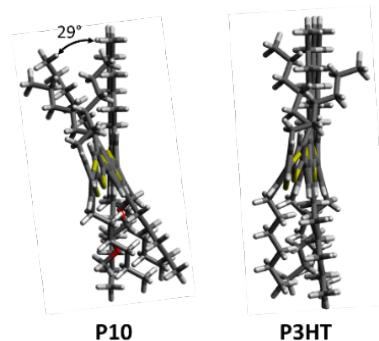


Figure 3. Along the chain image of energy-minimised decamers of P3HT and **P10** (polar side chain in the middle) in the gas phase using B3LYP/6-31g(d) basis sets. The SCCS dihedral angle between a polar and alkyl side chain unit is indicated by the arrow.

The morphological structure of drop-cast films onto silicon substrates were investigated using X-ray diffraction, revealing a d(100) peak before and after incorporation of polar side chains, as shown in the supporting information (Fig S22 and Table S4). Interestingly, heating and cooling cycles from differential scanning calorimetry of the polar copolymers show no crystalline melt transitions, however X-ray diffraction reveals the existence of crystalline domains. The structure and orientation of the crystalline polymer domains in spin-cast thin films were probed with grazing incidence wide-angle X-ray scattering (GIWAXS) depicted in

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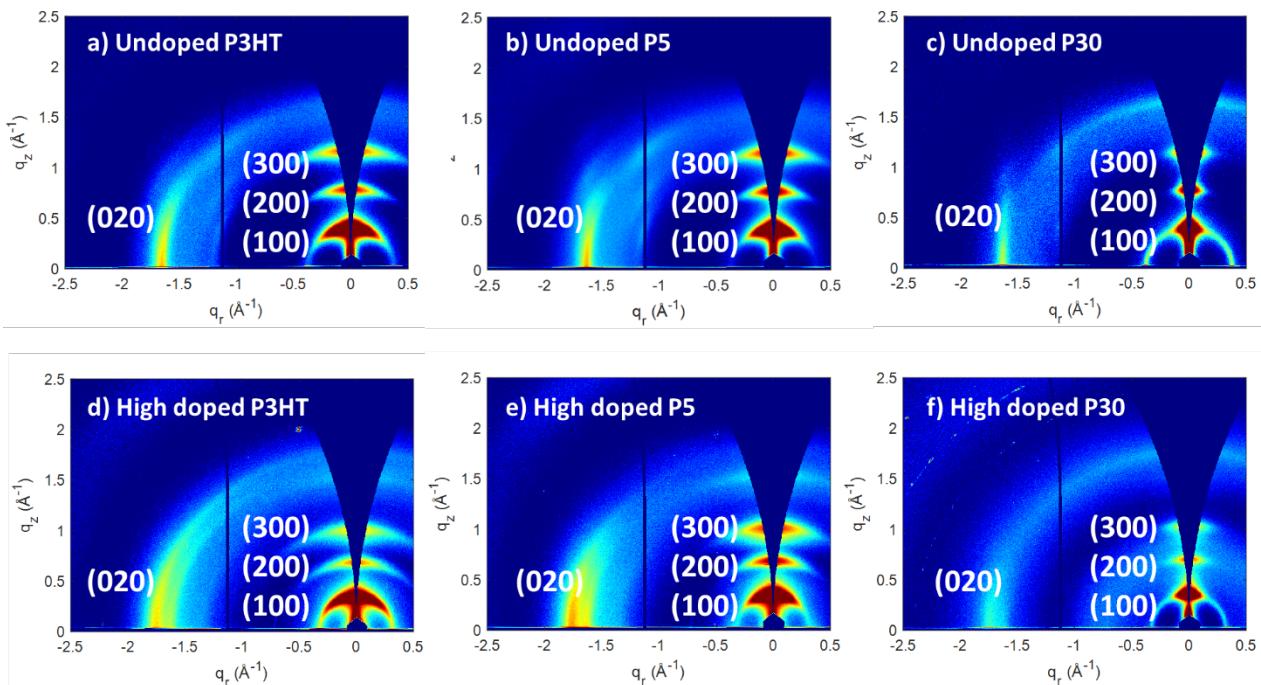


Figure 4. a-c) GIWAXS patterns of P3HT, **P5** and **P30** spin- cast onto silicon substrates at 2000 rpm for 90 s from a 10 mg ml⁻¹ ODCB solution, respectively. d-f) GIWAXS patterns of P3HT, P5 and P30 thin films doped with F4TCNQ using high doping conditions respectively (1 mg ml⁻¹ F4TCNQ in MeCN then 1 mg ml⁻¹ F4TCNQ in ODCB solutions with penetration times of 60 seconds each).

Figure 4 and the supporting information (Fig S23 and S24). Pristine P3HT thin film structures determined by GIWAXS have previously been reported.³²⁻³⁴ The scattering pattern, depicted in Figure 4a, shows that P3HT crystallites orient edge-on relative to the substrate with a π - π stack scattering peak (020) in-plane and multiple orders of strong lamellar scattering (h00) out-of-plane, consistent with previous results.³³ The structure for newly developed material, poly(3-((2-(2-methoxyethoxy)ethoxy)methyl)thiophene) (P3MEEMT), has been studied as well and has shown to manifest a mix of face-on and edge-on crystallites, with π - π stack scattering peak (020) and lamellar scattering (100) occurring both in- and out-of-plane.^{35,36}

Incorporation of glycol chains in the polar copolymers presented herein narrows the azimuthal angle distribution of the out-of-plane lamellar scattering (h00) and the near in-plane (020) scattering appears increasingly less arc-like and more rod-like with increasing polar side chain content depicted in Figure 4 and the supporting information (Fig S23). These both indicate an improved orientation of the edge-on crystallites. At the same time, the increasing glycol compositions lead to the strengthening of an isotropic lamellar ring shown in Figure 4b

and c, implying the growth of a sub-population of randomly oriented crystallites. In- and out-of-plane line cuts from the 2D plots were fit to extract d-spacings from peak centres and all polymers in the series show similar lamellar d-spacings (16.1 – 16.2 Å) and π - π stack d-spacings (3.80 – 3.84 Å) presented in the supporting in formation (Fig S23, S24 and Table S5, S6).

From peak widths, the coherence lengths in the direction of lamellar stacking were extracted from the fitted peaks (Figure S25). For the oriented crystallites, the coherence length increased from 160 Å (~10 repeat units) to 240 Å (~15 repeat units) with increasing polar content up to 20%, however at 30% a decrease is observed. For isotropic crystallites, the coherence length continually increased with increasing polar side chain content from a minimum of 48 Å (~3 repeat units) for **P5**, to 170 Å (~10 repeat units) for **P30**, mirroring the increase in isotropic scattering qualitatively apparent in 2-D scattering plots presented in the supporting information (Fig S23). Lacking a film thickness normalized absolute measure of incident and scattered photon flux, the GIWAXS data cannot be quantitatively correlated with the 0-1 absorption feature and spectroscopically determined degree of aggregation. However, qualitatively the in-plane π -stack scattering intensity appears to

decrease somewhat with increased glycol content in intensity versus the scattering background.

Atomic force microscopy of spin-cast thin films, shown in the supporting information (Fig S29), reveal surface morphologies of the polymer series in good agreement with reduced aggregation with increasing polar content. Although AFM images cannot unambiguously reveal the crystallinity of films, we observe a steady decrease in RMS roughness from 3.6 nm for P3HT to \sim 1 nm for **P20** and **P30**, shown in the supporting information (Fig S30), consistent with a reduction in aggregation with increasing polar content, as observed by UV-Vis spectroscopy.³⁷

2.3. Electrical Properties

The charge transport properties were investigated by fabrication of top gate bottom contact organic field-effect transistors (OFETs) with a CYTOP dielectric layer. Output and transfer characteristics are shown in the supporting information (Fig S33 and Table S7). The saturated hole mobility decreases gradually by approximately one order of magnitude across the series from $3.8 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for P3HT to $0.5 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for **P30**. The charge carrier mobility can be related to the free exciton bandwidth (W) which is also related to A_{0-0}/A_{0-1} of the 0-0 and 0-1 optical transitions as depicted in Equation 1 according to work by Spano. et. al.²⁹

$$\frac{A_{0-0}}{A_{0-1}} \approx \left(\frac{1 - 0.24W/E_p}{1 + 0.073W/E_p} \right)^2 \quad (1)$$

Where E_p is the intramolecular C=C vibration ($E_p = 0.18 \text{ eV}$). As illustrated in Figure 5a, the increase in exciton bandwidth with increasing polar content, originating from the decreasing backbone planarity and degree of aggregation, correlates reasonably with a gradual decrease in hole mobility across the series.

The electrical conductivity of the polymer series, using 2,3,5,6-tetrafluorotetracyanoquinodimethane (F4TCNQ) as a molecular p-type dopant, was investigated using a series of different sequential doping conditions. Sequential doping was chosen to afford high doping efficiencies and prevent solution aggregation during thin film fabrication.^{13,38} The polymers were

spin-cast from ODCB and subsequently doped with F4TCNQ in acetonitrile or ODCB at low and high dopant concentrations to emulate the weak and strong doping regimes, respectively, observed for P3HT.^{13,39,40} A gradual drop in conductivity of an order of magnitude is seen when going from P3HT to **P30** for both doping levels as illustrated in Figure 5b. This is in agreement with the decreasing trend in charge carrier mobilities of the neutral polymer films as discussed below.

2.4. Polymer:Dopant Interactions

To further examine the doped polymers, we investigated the structural and optical properties using GIWAXS, UV-Vis and Fourier-transform infrared (FTIR) spectroscopy. F4TCNQ doped P3HT thin film structures determined by GIWAXS have been widely reported.^{33,39,40} The 2-D GIWAXS scattering patterns and line cuts are shown in Figure 4d-f and in the supporting information (Figure S24, S26-S27). After doping with F4TCNQ, the backbone peak disappears due to loss of registry.³³ For the oriented crystallites, upon low doping the lamellar spacings increase from \sim 16.1 to \sim 18.0 Å, while at high doping level, the lamellar spacing for P3HT, **P5**, and **P10** increase a further \sim 0.3 Å, as shown in Figure 5c and Table S5. On the other hand, for the isotropic crystallites, upon low doping the lamellar spacing for all copolymers increase by \sim 1.7 Å compared to the undoped films, and then increase a further \sim 1.0 Å upon high doping. (Table S5 and Fig S28) These increases are consistent with penetration of F4TCNQ into the crystalline lamellar stacks.³³ For both oriented and isotropic crystallites, the lamellar coherence lengths of all polymers first increase with low F4TCNQ doping, then decrease for high doped samples (Figure S25). Upon low doping, the π -stacking scattering shows the appearance of a shoulder and clearly split into two peaks upon high doping, which is similar to previous reports.^{33,39,40} On the whole, the GIWAXS data indicate that the structural evolution with doping is similar across the polymer series. However, the in-plane π -stack scattering intensity versus the scattering background saw a bigger decrease with increasing polar content indicating loss in structural order. This result can be correlated with reduction of conductivity.

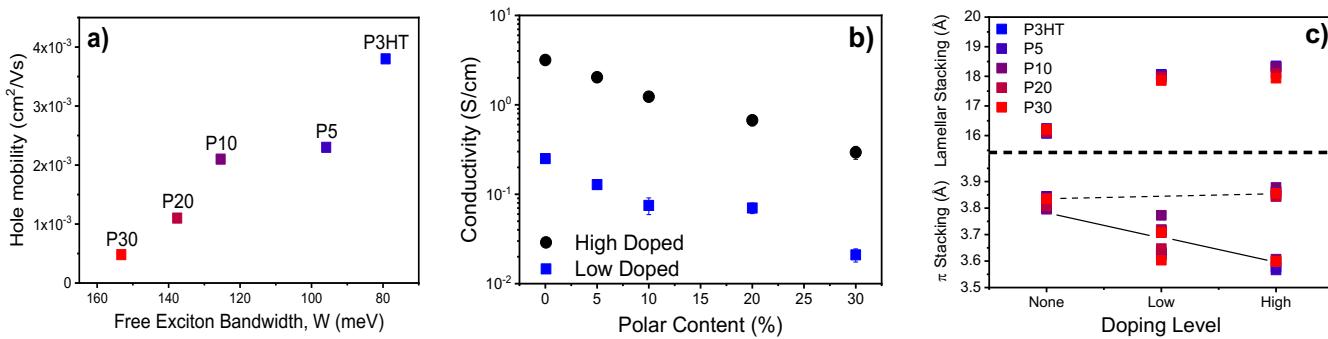


Figure 5. a) OFET hole mobility plotted vs free exciton bandwidth from Equation 1. b) Electrical conductivity measured on polymer thin films spin-coated (from 10 mg ml⁻¹ ODCB polymer solutions) onto glass slides with 50 nm Gold electrodes in Van de Pauw geometry. Thin Film thickness were measured to be between 27 – 30 nm for P3HT and co-polymers. Co-polymers were doped at high (1 mg ml⁻¹ F4TCNQ in MeCN then 1 mg ml⁻¹ F4TCNQ in ODCB solutions with penetration times of 60 seconds each) and low (0.1 mg ml⁻¹ F4TCNQ MeCN with a penetration time of 10 seconds) doping conditions. c) π -stacking distance ($d(020)$) and lamellar stacking distance ($d(h00)$) of oriented crystallites from GIWAXS measurements of

as-cast (from 10 mg ml⁻¹ ODCB), low doped and high doped polymer thin films. The solid line and the dash line for π -stacking spacing are used for tracking the contracted π -stacking polymorph and the expanded π -stack polymorph, respectively.

The surface morphology of doped polymer films was probed with AFM as shown in the supporting information (Fig S31 and S32). After doping, the same trend as seen for undoped films is observed where increasing polar side chain content leads to decreased roughness.

Thin film UV-Vis spectroscopy of the high and low doped polymers confirm integer charge transfer between polymer and dopant, in all cases evidenced by bleaching of the neutral polymer π - π^* transition and appearance of the distinctive four peaks from the F4TCNQ radical anion at 2.95, 1.78, 1.61, and 1.45 eV.¹³ The relative intensity of the F4TCNQ radical anion peaks at high doping (Figure 6a) appear to show a slight reduction with increasing polar group content. However, fits of the high and low doped thin film UV-Vis spectra with KF4TCNQ solution spectrum as the F4TCNQ radical anion peaks shown in the supporting information (Fig S37-S38) reveal that the F4TCNQ anion content in each sample is similar. In the high doping samples, the fits reveal a slight (~10%) increase in dopant density for **P5** and **P10** vs. P3HT, while for **P20** and **P30** the values are within error of P3HT. In the low doping samples, there is a slight (10-20%) increase in dopant density for all polymers compared to P3HT. These carrier density values, along with the measured charge carrier mobilities were then used to estimate an effective doping efficiency (Fig S39 and Table S9). From these calculations we conclude that the doping efficiency is only slightly reduced when introducing polar side chains, and that the drop in conductivity seen in Figure 5b almost exclusively results from a decrease in mobility (Figure 5a).

Since the P2 sub-gap optical transition of the P3HT radical cation at ~1.5 eV overlaps with the F4TCNQ radical anion peaks in the UV-Vis spectrum, FTIR was used to analyse the lower-energy P1 sub-gap transition at ~0.5 eV.³³ The position of this broad peak correlates with the degree of localisation of the

polaron, with a shift to lower energies indicative of increased delocalisation.⁴¹⁻⁴³ For both the high and low doped polymer series, a blue-shift of ~0.13 eV was observed for the P1 transition when going from P3HT (0.45 eV) to **P30** (0.55 eV), indicative of increasing localisation of the polaron with increasing polar content as shown in Figure 6b and supporting information (Fig S34). This localization signature is reflected in the reduced carrier mobility observed in FET measurements.

In the C≡N stretching region of the FTIR spectra of high and low doped polymer thin films we see three peaks at 2194, 2187 and 2169 cm⁻¹ which are labelled A, B and C, respectively as illustrated in Figure 6c (high doping) and Fig S36 (low doping). The A and B bands are related to the b_{1u} and a_g C≡N stretching modes in the F4TCNQ radical anion, respectively, while the C band is attributed to the b_{2u} and b_{3g} C≡N stretching modes.⁴⁴⁻⁴⁶ It is hypothesised that the a_g (B) and b_{3g} (C) modes, which are not normally IR active, are observed due to a loss of planarity of F4TCNQ when in the polymer film.⁴⁷ There is a linear relationship between the degree of charge transfer and the change in the IR frequency of the b_{1u} stretching mode (A) from neutral to reduced F4TCNQ.^{48,49} Since all the polymers herein display a 33 cm⁻¹ shift, we conclude that complete transfer between polymer and dopant takes place for both P3HT and the copolymers with polar side chains.

For high doped polymer films, the relative intensities of the B and C bands decrease with increasing polar content, concurrent with a slight shift to higher wavenumbers. We ascribe this behaviour to the charge transfer complex between the oxidised P3HT and reduced F4TCNQ to be more localised in the case of the polythiophenes with increased polar content. This is in agreement with the blue-shifted P1 transition discussed above and previous work by Pingel et al., where they altered the conjugation length of P3HT through random co-polymerisation with different amounts of the conjugation breaker tetrafluorobenzene (TFB).⁴⁶ They found that with increasing amounts of TFB, the peak at 2188 cm⁻¹, labelled B herein, decreases in intensity, and they ascribe that to the F4TCNQ charge transfer becoming more localised. A decrease in the intensity of the peaks is also indicative of decreased concentration in integer charge transfer formation.

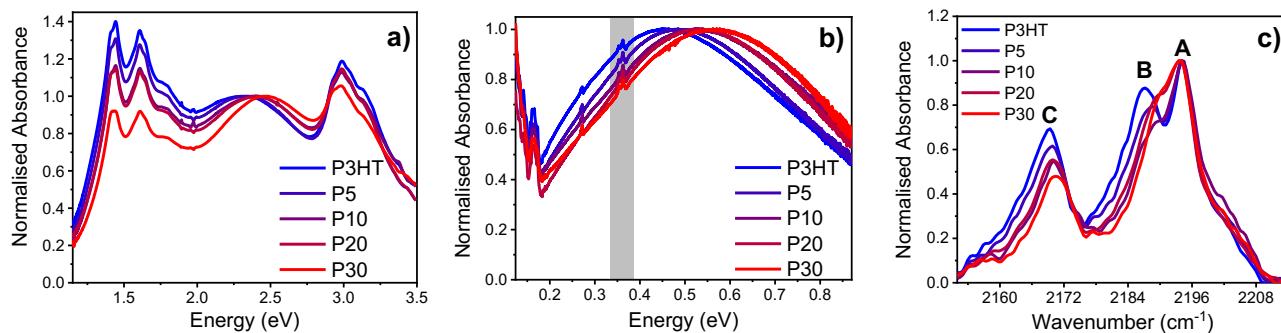


Figure 6. a) UV-Vis spectra of high doped polymer thin films normalised to the π - π^* transition. b) Normalised FTIR spectra of high doped polymer thin films (grey area highlighting the region of the F4TCNQ C≡N stretches). The spectra are normalised to λ_{max} of the P1 polaron. c) IR region showing the C≡N stretches of F4TCNQ where the A, B and C labelled peaks are at 2194, 2187 and 2169 cm⁻¹ respectively. The spectra are normalised to peak A.

The FTIR data strongly suggests that the gradual loss of backbone planarity with increasing polar content observed in the neutral polymer films also exists in the doped polymer films.

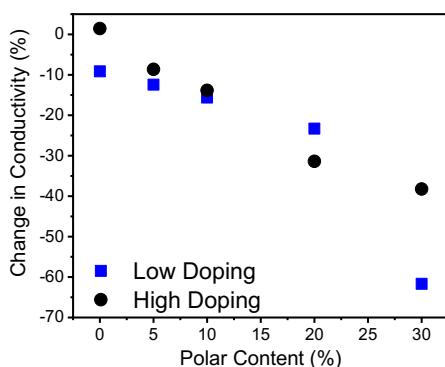


Figure 7. Relative change in electrical conductivity of F4TCNQ doped polymer thin films after annealing at 80 °C for 20 mins.

We subsequently investigated the thermal stability of the doped polymers by annealing the thin films under a nitrogen atmosphere for 20 mins at 80 °C. We chose this temperature to observe the effect of diffusion within the film, as it is below the temperature at which F4TCNQ starts to sublimate out of P3HT films.^{18,50} Conductivity measurements before and after annealing reveal a drop in conductivity across the series for high and low doping conditions as illustrated in Figure 7. At high doping levels, the conductivity of P3HT increases slightly upon thermal annealing, whereas the relative drop in conductivity upon annealing increases with increasing polar content for the copolymers. Under low doping conditions, a similar trend is observed, however a larger decrease in conductivity around 60% is seen for P30 upon annealing. Under these conditions, there is thus no indication of a stabilising interaction between the dopant and the polar side chain. The reduced thermal stability with increased polar group content may be explained by recent work by Watts *et. al.* who observe thermally induced formation of HF4TCNQ anions in F4TCNQ doped P3HT via proton abstraction of the alpha proton on the side chain.⁵¹ Within our system, the carbon spacer with an oxygen adjacent could promote this proton abstraction, leading to a larger decrease in conductivity with increasing polar content.

3. Conclusions

Polar side chain functionalisation of semiconducting polymers, where an ethylene glycol based motif is directly attached via a C-O bond to the polymer backbone, will have a pronounced effect on frontier energy levels due to resonance effects, and on polymer backbone planarity if attractive intramolecular interactions such as non-covalent S-O interactions can be established. This makes direct comparison with the fully alkylated analogue inaccurate. In this study, we sought to overcome these obstacles by introducing a methylene spacer group between the polymer backbone and the ethylene glycol motif. We find that in this situation, ethylene glycol based side chains have a significant negative impact on polymer microstructure and energetics. This modification is in turn detrimental to properties important in many organic electronic devices, such as field-effect transistors and thermoelectric generators.

We have studied in detail the effect of ethylene glycol based polar side chains incorporated in a random fashion into the well-known semiconducting polymer P3HT. Compared to P3HT, introduction of 5 and 10 % polar side chain content has a relatively small impact on optical, electrochemical and structural properties, whereas polar content above 10% drastically impairs the polymer crystallinity and shifts the main oxidative event to significantly higher electrochemical potentials. OFET hole mobilities decrease gradually with increasing polar content, initially ascribed to loss of backbone planarity for low polar content, and subsequently further impacted by weakened interchain interactions for higher polar content polymers. Conductivity measurements after doping with F4TCNQ reveal a similar gradual decrease in electrical conductivity with increasing polar content. Changes in dopant concentration and doping efficiency are small across all polymers, indicating that the reduction in carrier mobility dominates the drop in conductivity.

Our study highlights that the incorporation of polar side chains is not a universal solution to e.g. increasing a material's dielectric constant, polymer:dopant miscibility or favourable interactions with ionic species while simultaneously maintaining other desirable properties such as ordered solid-state structural organisation and effective charge carrier transport. We moreover find no evidence in our system that the incorporation of polar side chains enhances the doping stability at elevated temperatures. Upon thermal annealing of F4TCNQ-doped polymer films, conductivity values in fact decrease more with increasing polar content.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 B. Lüssem, C. M. Keum, D. Kasemann, B. Naab, Z. Bao and K. Leo, *Chem. Rev.*, 2016, **116**, 13714–13751.
- 2 A. Marocchi, D. Lanari, A. Facchetti and L. Vaccaro, *Energy*

ARTICLE

Journal Name

3 *Environ. Sci.*, 2012, **5**, 8457–8474.

4 H. Sirringhaus, *Adv. Mater.*, 2014, **26**, 1319–1335.

5 O. Bubnova and X. Crispin, *Energy Environ. Sci.*, 2012, **5**, 9345–9362.

6 S. Holliday, R. S. Ashraf, A. Wadsworth, D. Baran, S. A. Yousaf, C. B. Nielsen, C. H. Tan, S. D. Dimitrov, Z. Shang, N. Gasparini, M. Alamoudi, F. Laquai, C. J. Brabec, A. Salleo, J. R. Durrant and I. McCulloch, *Nat. Commun.*, DOI:10.1038/ncomms11585.

7 J. Borges-González, C. J. Kousseff and C. B. Nielsen, *J. Mater. Chem. C*, 2019, **7**, 1111–1130.

8 Y. Lin, Y. Firdaus, M. I. Nugraha, F. Liu, S. Karuthedath, A. H. Emwas, W. Zhang, A. Seitkhan, M. Neophytou, H. Faber, E. Yengel, I. McCulloch, L. Tsetseris, F. Laquai and T. D. Anthopoulos, *Adv. Sci.*, 2020, **7**, 1–9.

9 K. Kang, S. Watanabe, K. Broch, A. Sepe, A. Brown, I. Nasrallah, M. Nikolka, Z. Fei, M. Heeney, D. Matsumoto, K. Marumoto, H. Tanaka, S. Kuroda and H. Sirringhaus, *Nat. Mater.*, 2016, 1–8.

10 V. Vijayakumar, Y. Zhong, V. Untilova, M. Bahri, L. Herrmann, L. Biniek, N. Leclerc and M. Brinkmann, *Adv. Energy Mater.*, 2019, **1900266**, 1900266.

11 I. E. Jacobs and A. J. Moulé, *Adv. Mater.*, 2017, **29**, 1–39.

12 A. F. Paterson, A. Savva, S. Wustoni, L. Tsetseris, B. D. Paulsen, H. Faber, A. H. Emwas, X. Chen, G. Nikiforidis, T. C. Hidalgo, M. Moser, I. P. Maria, J. Rivnay, I. McCulloch, T. D. Anthopoulos and S. Inal, *Nat. Commun.*, 2020, 1–11.

13 E. Lim, K. A. Peterson, G. M. Su and M. L. Chabinyc, *Chem. Mater.*, 2018, **30**, 998–1010.

14 I. E. Jacobs, E. W. Aasen, J. L. Oliveira, T. N. Fonseca, J. D. Roehling, J. Li, G. Zhang, M. P. Augustine, M. Mascal and A. J. Moulé, *J. Mater. Chem. C*, 2016, **4**, 3454–3466.

15 J. Hynynen, D. Kiefer, L. Yu, R. Kroon, R. Munir, A. Amassian, M. Kemerink and C. Müller, *Macromolecules*, 2017, **50**, 8140–8148.

16 A. Hamidi-Sakr, L. Biniek, J. L. Bantignies, D. Maurin, L. Herrmann, N. Leclerc, P. Lévéque, V. Vijayakumar, N. Zimmermann and M. Brinkmann, *Adv. Funct. Mater.*, 2017, **27**, 1–13.

17 J. Li, C. Koshnick, S. O. Diallo, S. Ackling, D. M. Huang, I. E. Jacobs, T. F. Harrelson, K. Hong, G. Zhang, J. Beckett, M. Mascal and A. J. Moulé, *Macromolecules*, 2017, **50**, 5476–5489.

18 L. Müller, S. Y. Rhim, V. Sivanesan, D. Wang, S. Hietzschold, P. Reiser, E. Mankel, S. Beck, S. Barlow, S. R. Marder, A. Pucci, W. Kowalsky and R. Lovrincic, *Adv. Mater.*, 2017, **29**, 1–7.

19 R. Kroon, D. Kiefer, D. Stegerer, L. Yu, M. Sommer and C. Müller, *Adv. Mater.*, 2017, **1700930**, 1700930.

20 J. Li, C. W. Rochester, I. E. Jacobs, E. W. Aasen, S. Friedrich, P. Stroeve and A. J. Moulé, *Org. Electron.*, 2016, **33**, 23–31.

21 D. Kiefer, A. Giovannitti, H. Sun, T. Biskup, A. Hofmann, M. Koopmans, C. Cendra, S. Weber, L. J. Anton Koster, E. Olsson, J. Rivnay, S. Fabiano, I. McCulloch and C. Müller, *ACS Energy Lett.*, 2018, **3**, 278–285.

22 J. Liu, M. P. Garman, J. Dong, B. Van Der Zee, L. Qiu, G. Portale, J. C. Hummelen and L. J. A. Koster, *ACS Appl. Energy Mater.*, 2019, **2**, 6664–6671.

23 A. I. Hofmann, R. Kroon, L. Yu and C. Müller, *J. Mater. Chem. C*, 2018, **6**, 6905–6910.

24 J. Li, G. Zhang, D. M. Holm, I. E. Jacobs, B. Yin, P. Stroeve, M. Mascal and A. J. Moulé, *Chem. Mater.*, 2015, **27**, 5765–5774.

25 D. Kiefer, R. Kroon, A. I. Hofmann, H. Sun, X. Liu, A. Giovannitti, D. Stegerer, A. Cano, J. Hynynen, L. Yu, Y. Zhang, D. Nai, T. F. Harrelson, M. Sommer, A. J. Moulé, M. Kemerink, S. R. Marder, I. McCulloch, M. Fahlman, S. Fabiano and C. Müller, *Nat. Mater.*, 2019, **18**, 149–155.

26 C. B. Nielsen, A. Giovannitti, D.-T. T. Sbircea, E. Bandiello, M. R. Niazi, D. A. Hanifi, M. Sessolo, A. Amassian, G. G. Malliaras, J. Rivnay and I. McCulloch, *J. Am. Chem. Soc.*, 2016, **138**, 10252–10259.

27 R. S. Loewe, P. C. Ewbank, J. Liu, L. Zhai and R. D. McCullough, *Macromolecules*, 2001, **34**, 4324–4333.

28 R. D. McCullough and R. D. Lowe, *J. Chem. Soc. Chem. Commun.*, 1992, 70–72.

29 F. C. Spano, *Acc. Chem. Res.*, 2009, **43**, 429–439.

30 J. Clark, J. F. Chang, F. C. Spano, R. H. Friend and C. Silva, *Appl. Phys. Lett.*, 2009, **94**, 2007–2010.

31 F. C. Spano and C. Silva, *Annu. Rev. Phys. Chem.*, 2014, **65**, 477–500.

32 R. J. Kline, M. D. McGehee and M. F. Toney, *Nat. Mater.*, 2006, **5**, 222–228.

33 D. T. Scholes, P. Y. Yee, J. R. Lindemuth, H. Kang, J. Onorato, R. Ghosh, C. K. Luscombe, F. C. Spano, S. H. Tolbert and B. J. Schwartz, *Adv. Funct. Mater.*, 2017, **27**, 1–13.

34 N. Kayunkid, S. Uttiya and M. Brinkmann, *Macromolecules*, 2010, **43**, 4961–4967.

35 C. G. Bischak, L. Q. Flagg, K. Yan, T. Rehman, D. W. Davies, R. J. Quezada, J. W. Onorato, C. K. Luscombe, Y. Diao, C. Z. Li and D. S. Ginger, *J. Am. Chem. Soc.*, 2020, **142**, 7434–7442.

36 B. X. Dong, C. Nowak, J. W. Onorato, J. Strzalka, F. A. Escobedo, C. K. Luscombe, P. F. Nealey and S. N. Patel, *Chem. Mater.*, 2019, **31**, 1418–1429.

37 B. Meng, J. Liu and L. Wang, *Polym. Chem.*, 2020, **11**, 1261–1270.

38 D. T. Scholes, S. A. Hawks, P. Y. Yee, H. Wu, J. R. Lindemuth, S. H. Tolbert and B. J. Schwartz, *J. Phys. Chem. Lett.*, 2015, **6**, 4786–4793.

39 B. Neelamraju, K. E. Watts, J. E. Pemberton and E. L. Ratcliff, *J. Phys. Chem. Lett.*, 2018, **9**, 6871–6877.

40 D. T. Duong, C. Wang, E. Antono, M. F. Toney and A. Salleo, *Org. Electron. physics, Mater. Appl.*, 2013, **14**, 1330–1336.

41 R. Ghosh, C. K. Luscombe, M. Hambach, S. C. B. Mannsfeld, A. Salleo and F. C. Spano, *Chem. Mater.*, DOI:10.1021/acs.chemmater.9b01704.

42 R. Ghosh, A. R. Chew, J. Onorato, V. Pakhnyuk, C. K. Luscombe, A. Salleo and F. C. Spano, *J. Phys. Chem. C*, 2018, **122**, 18048–18060.

43 R. Ghosh, C. M. Pochas and F. C. Spano, *J. Phys. Chem. C*,
2016, **120**, 11394–11406.

44 N. L. Haworth, J. Lu, N. Vo, T. H. Le, C. D. Thompson, A. M.
Bond and L. L. Martin, *Chempluschem*, 2014, **79**, 962–972.

45 H. Hase, K. O'Neill, J. Frisch, A. Opitz, N. Koch and I.
Salzmann, *J. Phys. Chem. C*, 2018, **122**, 25893–25899.

46 P. Pingel, L. Zhu, K. S. Park, J. O. Vogel, S. Janietz, E. G. Kim,
J. P. Rabe, J. L. Brédas and N. Koch, *J. Phys. Chem. Lett.*,
2010, **1**, 2037–2041.

47 E. F. Aziz, A. Vollmer, S. Eisebitt, W. Eberhardt, P. Pingel, D.
Neher and N. Koch, *Adv. Mater.*, 2007, **19**, 3257–3260.

48 H. Méndez, G. Heimel, S. Winkler, J. Frisch, A. Opitz, K.
Sauer, B. Wegner, M. Oehzelt, C. Röthel, S. Duhm, D.
Többens, N. Koch and I. Salzmann, *Nat. Commun.*, 2015, **6**,
8560.

49 J. S. Chappell, W. A. Bryden, M. Maxfield, D. O. Cowan, A.
N. Bloch and T. O. Poehler, *J. Am. Chem. Soc.*, 1981, **103**,
2442–2443.

50 K. Kang, S. Schott, D. Venkateshvaran, K. Broch, G.
Schweicher, D. Harkin, C. Jellett, C. B. Nielsen, I. McCulloch
and H. Sirringhaus, *Mater. Today Phys.*, 2019, **8**, 112–122.

51 K. E. Watts, B. Neelamraju, M. Moser, I. McCulloch, E. L.
Ratcliff and J. E. Pemberton, *J. Phys. Chem. Lett.*, 2020, **11**,
6586–6592.