

# 1 Suppression of CO<sub>2</sub> outgassing by gas bubbles under a hurricane

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## 18 Key Points:

- 19 • Gas bubbles reduce CO<sub>2</sub> efflux from supersaturated water by a factor of 50 under  
20 hurricane winds for the simulated conditions.
- 21 • Bubble-induced supersaturation should be included in gas flux parameterization for CO<sub>2</sub>  
22 in high wind conditions
- 23 • Parameterizations derived using O<sub>2</sub>, N<sub>2</sub> and inert gases may not be accurate for CO<sub>2</sub> that  
24 has a very different solubility.

## 26 Plain Language Summary

27 Carbon dioxide (CO<sub>2</sub>) is the primary anthropogenic greenhouse gases in the atmosphere and is  
28 the gas most responsible for global warming. The ocean is an important sink of anthropogenic  
29 atmospheric CO<sub>2</sub>, yet the exchange of CO<sub>2</sub> between the ocean and the atmosphere is not fully  
30 understood. This study re-examines the exchange of CO<sub>2</sub> under hurricanes over the low-latitude  
31 ocean, using hurricane Frances (2004) as an example. Previous studies show that hurricanes  
32 significantly facilitate the outgassing of CO<sub>2</sub> due to the extreme wind. Those studies, however,  
33 do not explicitly consider gas bubbles. Gas bubbles, entrained into the ocean when ocean wave  
34 breaks, are ubiquitous under hurricanes also due to the extreme wind. While in the ocean, gas  
35 bubbles not only move around under the influence of the chaotic wind-driven currents and their  
36 own buoyancy, they also exchange gases with the water. Our study, using state-of-the-art  
37 computer models that concurrently simulate the chaotic ocean currents, gas bubbles, and  
38 dissolved gases, demonstrates that hurricane plays a significantly smaller role in the ocean-  
39 atmosphere transfer of CO<sub>2</sub> than previously estimated.

## 40 Abstract

41 The role of gas bubbles on the air-sea CO<sub>2</sub> flux during Hurricane Frances (2004) is studied using  
42 a large-eddy simulation model that couples ocean surface boundary layer turbulence, gas bubbles  
43 and dissolved gases. In the subtropical surface ocean where gases are slightly supersaturated,  
44 gases in bubbles can still dissolve due to hydrostatic pressure and surface tension exerted on  
45 bubbles. Under the simulated conditions, the CO<sub>2</sub> efflux with an explicit bubble effect is less  
46 than 2% of that calculated using a gas flux formula without explicit inclusion of bubble effect.  
47 The use of a gas flux parameterization without bubble-induced supersaturation contributes to  
48 uncertainty in the global carbon budget. The results highlight the importance of bubbles under  
49 high winds even for soluble gases such as CO<sub>2</sub> and demonstrate that gas flux parameterization  
50 derived from gases of certain solubility may not be accurate for gases of very different solubility.

## 51 1 Introduction

52 The transfer of carbon dioxide (CO<sub>2</sub>) through the ocean-atmosphere interface modulates the  
53 cycling of CO<sub>2</sub> in the earth system and influences the amount of anthropogenic CO<sub>2</sub>, an  
54 important greenhouse gas, absorbed in the ocean. It has now been qualitatively understood that  
55 the ocean takes in CO<sub>2</sub> at the high latitudes where the surface water is cold and is undersaturated

56 in CO<sub>2</sub>, and releases CO<sub>2</sub> at low latitudes where the water is warm and is supersaturated in CO<sub>2</sub>.  
57 The quantitative estimates of ocean-atmosphere CO<sub>2</sub> flux, however, still contain strong  
58 uncertainties [e.g., *Wanninkhof et al.* 2009; *Woolf et al.* 2019], associated with errors in the data  
59 used to calculate CO<sub>2</sub> flux including gas concentrations and wind speed, and with the functional  
60 form of gas flux parameterization due to insufficient understanding of gas transfer processes  
61 particularly when the ocean is rough and gas bubbles entrained during wave breaking play a key  
62 role.

63 Gas bubbles alter air-sea gas exchange in two ways: First, the gas transfer rate is enhanced as  
64 gases go in and out of the ocean through bubble surfaces in addition to the sea surface. Secondly,  
65 gases still dissolve through bubbles at saturated and supersaturated water as gases in bubbles are  
66 squeezed by water pressure and surface tension; as a result, the ocean is supersaturated when the  
67 total air-sea gas flux is zero [e.g., *Woolf* 1997]. Although the enhanced gas transfer rate is  
68 already (mostly implicitly) in popular gas flux parameterizations, the bubble-induced  
69 supersaturation is thought to be unimportant for highly soluble gases such as CO<sub>2</sub> and is never  
70 considered in any realistic CO<sub>2</sub> flux calculations. A recent observational study [*Leighton et al.*  
71 2018], however, concluded that bubble-induced supersaturation is more important than  
72 previously thought, implying that existing CO<sub>2</sub> calculations neglecting bubble-induced  
73 supersaturation may overestimate (underestimate) efflux (influx). On the other hand, the  
74 importance of bubble-induced supersaturation for less soluble gases such as O<sub>2</sub> [e.g., *Vagle et al.*  
75 2010; *Bushinsky et al.* 2016; *Wang et al.* 2019], inert gases [*Stanley et al.* 2009] and their ratios  
76 [e.g., *Hamme and Emerson* 2006] has already been recognized. It was also established in those  
77 studies that accurate quantification of the total air-sea gas flux requires the separation of the total  
78 gas flux into three components, viz., the gas flux through the ocean surface, that through bubbles  
79 that completely dissolved, and that through bubbles that eventually burst at the ocean surface  
80 (partially dissolved).

81 Since gas transfer rate increases with wind speed, it is plausible that gas transfer under hurricanes  
82 makes a significant contribution to global CO<sub>2</sub> flux. Based on observational data, *Bates et al.*  
83 [1998] concluded that hurricanes contribute up to half of the global CO<sub>2</sub> efflux. *Levy et al.*  
84 [2012] noted that the gross effect of hurricanes on CO<sub>2</sub> efflux is much less and is about 10% of  
85 the global CO<sub>2</sub> efflux because a considerable portion of the ocean under a hurricane path is  
86 undersaturated in CO<sub>2</sub> and vertical mixing is anomalously weak after the passage of hurricanes.

87 Several other studies [e.g., *Bate* 2007; *Huang and Imberger* 2010] also discuss the CO<sub>2</sub> efflux  
88 under hurricane conditions. While those studies debate the quantitative integral effect of  
89 hurricanes on CO<sub>2</sub> efflux, they all agree that strong CO<sub>2</sub> outgassing occurs over surface water  
90 supersaturated in CO<sub>2</sub> during hurricane passage due to the enhanced transfer rate associated with  
91 strong winds. Gas flux estimates in most studies implicitly include bubble-enhanced transfer rate  
92 but not bubble-induced supersaturation [e.g., *Wanninkhof* 1992; *Fairall et al.* 2011]. While  
93 parameterizations are based on gas flux measurements, including CO<sub>2</sub>, the data for high wind are  
94 collected at high latitudes where the wind is strong but CO<sub>2</sub> is strongly undersaturated [e.g., *Ho*  
95 *et al.* 2011; *Bell et al.* 2017]. At strongly undersaturated conditions, both the gas flux through the  
96 ocean surface and that through bubbles are from the atmosphere to the ocean. It remains  
97 unknown if the gas transfer formulas derived from undersaturated water, with CO<sub>2</sub> influx, apply  
98 to supersaturated water, with CO<sub>2</sub> efflux.

99 The purpose of this study is to examine the effect of gas bubbles on CO<sub>2</sub> outgassing in  
100 supersaturated waters under Hurricane Frances (2004). The study is conducted by synthesizing in  
101 situ observations including gases and numerical solutions from a process model that  
102 simultaneously models turbulent flows, gas bubbles, and dissolved bubbles. Section 2 describes  
103 the data and the model; Section 3 presents and discusses the results; and Section 4 is a summary.

## 104 **2 Data and Model Description**

### 105 **2.1 Data**

106 Hurricane Frances (2004) was a category-4 hurricane in the northwestern Atlantic Ocean causing  
107 significant damage to the Bahamas and southeastern states of the U.S. A large campaign was  
108 conducted starting August 31<sup>st</sup>, 2004 to observe both sides of the air-sea interface in the path of  
109 the hurricane. Details of the field campaign, instrumentation, and available data are in *Black et*  
110 *al.* [2007]. For this study, the concentrations of dissolved O<sub>2</sub> and N<sub>2</sub> derived from an O<sub>2</sub> sensor  
111 and a gas tension device on a Lagrangian float that constantly transited across the upper 40-m or  
112 so of the ocean under the maximum hurricane winds are used (Fig. 1a) [*D'Asaro and McNeil*  
113 2007]. Temperature and salinity from a nearby autonomous profiling Electromagnetic  
114 Autonomous Profiling Explorer (EM-APEX) float [*Sanford et al.* 2011] are also used.

### 115 **2.2 Model**

#### 116 **2.2.1 Model description**

117 The computer solutions are obtained using the coupled-ocean-bubble-gas model [Liang *et al.*  
118 2017]. Turbulent currents under the hurricane are simulated with the National Center for  
119 Atmospheric Research Large Eddy Simulation (NCAR-LES) model [e.g., *Sullivan and*  
120 *McWilliams* 2010]. The evolution of gas bubbles is tracked as Lagrangian particles and the  
121 concentration of dissolved gases are simulated using advection-diffusion equations. Model  
122 formulation and implementation are repeated in the supporting information. The next two  
123 paragraphs highlight the improvements from Liang *et al.* [2017] that are important for bubble  
124 simulations.

125 Bubbles are entrained into the ocean by breaking waves. In our past studies [e.g., Liang *et al.*  
126 2017], breaking wave number density is an exponential function of breaking wave speed ( $c$ ) (see  
127 equation 3.15 in *Sullivan et al.* 2007). The functional form is based on visual observation by  
128 *Melville and Matusov* [2002]. In this study, we implemented the breaking wave front distribution  
129 function ( $\Lambda(c)$ ) from more recent infrared observations [Zappa *et al.* 2012; *Sutherland and*  
130 *Melville* 2013 and 2015; *Romero* 2019] that are power-law functions of  $c$  (see supporting  
131 information for details). There are more large breaking waves and less small breaking waves  
132 using the current power-law distribution function than the exponential distribution function in  
133 our previous studies (see Figure S1 in supporting information).

134 In addition to  $O_2$  and  $N_2$  that were simulated in Liang *et al.* [2017],  $CO_2$  is also explicitly  
135 simulated in the same way as in Liang *et al.* [2011]., In the model, dissolved inorganic carbon  
136 ( $[DIC]=[CO_3^{2-}]+[HCO_3^{2-}]+[H_2CO_3]$ ) and alkalinity are simulated. The partial pressure of  $CO_2$  is  
137 diagnosed by assuming equilibrium chemistry [e.g., *Sarmiento and Gruber* 2006; *Emerson and*  
138 *Hedges* 2008]. Since the carbonate reaction is assumed instantaneous, our model does not have  
139 the chemical enhancement effect of  $CO_2$  exchange that is negligible at normal ocean turbulence  
140 conditions [e.g., *Wanninkhof and Knox* 1996]. The use of equilibrium assumption in LES studies  
141 of upper ocean  $CO_2$  is also confirmed by Smith *et al.* [2018], who compared LES simulations  
142 resolving carbonate reactions and those assuming carbonate system equilibrium and concluded  
143 that explicit consideration of carbonate reaction has minimal effect on  $CO_2$  flux.

#### 144 **2.2.2 Model configuration**

145 The model was configured in a rectangular domain of  $400\times400\times200$ m with  $256\times256\times192$  grids.  
146 The simulation starts on day 245.15 ( $\sim 3:36$  AM Sept. 1<sup>st</sup> GMT) of 2004 when the dissolved gas

147 data are available and lasts for about 1.16 days within which the wind strengthened to its  
148 maximum and weakened to about the pre-storm strength. Wind speeds at 10-m above the ocean  
149 surface ( $U_{10}$ ) and sea level pressure (Fig. 1b) were interpolated from the National Oceanic and  
150 Atmospheric Administration/Hurricane Research Division (NOAA/HRD) real-time wind  
151 analysis (H\*WIND) product [Powell *et al.* 1988]. The float is located on the right flank of the  
152 hurricane and the largest  $U_{10}$  over the float is greater than 52 m/s at day 0.6. Surface heat flux  
153 and freshwater flux were interpolated from the hourly NCEP Climate Forecast System  
154 Reanalysis (CFSR) product at a spatial resolution of  $0.312^\circ \times 0.312^\circ$  [Saha *et al.* 2010]. The  
155 strong wind and precipitation also lead to strong cooling flux and freshwater flux to the ocean  
156 (Fig. 1c). Wave parameters, including Stokes drift, peak wave period, and wave energy to the  
157 ocean, are based on the solutions of wave spectrum computed using the wind-wave model  
158 WAVEWATCH III® (WWIII) [Tolman *et al.* 2009]. The WWIII model was also driven by the  
159 H\*WIND product. Details about the configuration of the WWIII model were described in Fan *et*  
160 *al.* [2009]. The calculation of wave energy flux to the ocean follows *Fan and Hwang* [2017]. The  
161 waves in the front-right quadrant of the hurricane track are higher and longer due to the  
162 resonance effect caused by the movement of the storm, while those in the rear-left quadrant are  
163 lower and shorter. The atmospheric fraction of gases is constant through the simulation and is  
164 20.9%, 78.1% and 374 ppm for O<sub>2</sub>, N<sub>2</sub> and CO<sub>2</sub>, respectively. The value for atmospheric CO<sub>2</sub>  
165 concentration was inferred from the updated observation-based global monthly gridded sea  
166 surface CO<sub>2</sub> and air-sea CO<sub>2</sub> flux product [Landschützer *et al.* 2017].

167 Initial profiles of temperature, salinity, dissolved O<sub>2</sub> and dissolved N<sub>2</sub> were set based on the  
168 measurements on the floats (Figs. 2a and 2b). The initial mixed layer depth is less than 30 m.  
169 Temperature is higher and salinity is lower in the mixed layer than below, as is typical of the  
170 region [Sanford *et al.* 2011]. Dissolved N<sub>2</sub> is lowest in the mixed layer and increases with depth  
171 below the mixed layer. It is about 1% supersaturated in the mixed layer. Dissolved O<sub>2</sub> is more  
172 than 2% supersaturated in the mixed layer. In the studied region with strong surface heating and  
173 sunlight, the euphotic zone is deeper than the mixed layer, leading to higher dissolved O<sub>2</sub> in the  
174 thermocline than in the mixed layer. Dissolved O<sub>2</sub> is the largest at around 60 m, below which  
175 dissolved O<sub>2</sub> decreases with depth due to remineralization. Initial profiles of dissolved inorganic  
176 carbon and alkalinity (Fig. 2c) are the same as those used by *Huang and Imberger* [2010] for the  
177 same storm. The profiles are based on measurement by the R/V Knorr cruise at (22.21°N,

178 66.00°W) in 1997, but are increased by 9.59 mmol/kg and 4.27 mmol/kg for DIC and alkalinity,  
179 respectively, based on the observed increase in mean DIC and alkalinity in the same region and  
180 between 1997 and 2004 by *Bates* [2007].

181 To understand the effect of gas bubbles, two simulations were carried out, one without gas  
182 bubbles (run NB) and the other one with bubbles (run B). In the simulation without gas bubbles,  
183 the total gas flux is calculated using the parameterization by *Wanninkhof* [1992] that is  
184 commonly used in earth system models. In the parameterization, the enhanced gas transfer rate  
185 due to bubbles is implicitly included while the bubble-induced supersaturation is neglected. In  
186 the simulation with gas bubbles, gas bubbles are represented by 8 million Lagrangian particles.  
187 The total gas flux is the sum of the gas flux through bubbles and the gas flux through the ocean  
188 surface. The gas flux through bubbles is explicitly calculated from the bubble fields and the  
189 dissolved gas fields, and the gas transfer rate through the ocean surface is calculated using the  
190 formula proposed by *Goddijn-Murphy et al.* [2012] for *in situ* wind, i.e.,  $k_{s,660}=2.6U_{10}-5.7$ .

191 **3 Results**

192 During the passage of the hurricane, the ocean surface cools by about two degrees. The cooling  
193 is primarily driven by the turbulent entrainment of thermocline water. Only 10% of the total  
194 cooling is caused by surface heat flux. Detailed analysis of the cold wake of the hurricane has  
195 been carried out by *D'Asaro et al.* [2007]. It plays an important role in the local atmosphere-  
196 ocean coupling [e.g., *D'Asaro et al.* 2007], and in the heat exchange between the surface and the  
197 interior oceans [*Mei et al.* 2013]. Although there is strong freshwater flux associated with  
198 precipitation during the passage of the hurricane, sea surface salinity increases (Fig 3b) due to  
199 the dominant effect of entrainment of saltier thermocline water below the mixed layer. The  
200 model captures the observed upper ocean response well, particularly during the strongest winds.

201 Mixed-layer dissolved O<sub>2</sub>, dissolved N<sub>2</sub> and DIC all increase after the passage of the hurricane in  
202 simulations with and without gas bubbles (red lines in Figs. 3c to 3e). Similar to heat and salt,  
203 dissolved gas concentrations are controlled by air-sea surface gas flux and by entrainment of  
204 thermocline water which has different dissolved gas concentrations. Throughout the simulated  
205 period, all three gases are supersaturated with respect to their atmospheric pressure (see the  
206 comparison between the solid lines and the dashed lines in Figs. 3c, 3d and 3f), so that the  
207 surface flux is from the ocean to the atmosphere. The supersaturated condition is a combined

208 consequence of the initial supersaturation condition, the low atmospheric pressure during the  
209 storm, the entrainment of the thermocline water immediately below the mixed layer, and the  
210 mixing of waters of different temperatures. For the three gases, the effect of entrainment  
211 dominates that of surface outgassing, and the mixed layer concentrations of gases increase.

212 In the simulation with bubbles, the mixed layer concentrations of all three gases increase, and are  
213 larger than in the simulation without bubbles throughout the simulated period. Given that the  
214 effect of bottom entrainment is the same for both simulations, the difference between the two  
215 simulations is from the explicit inclusion of gas bubbles. The difference between the two  
216 simulations is much more evident for O<sub>2</sub> and N<sub>2</sub> than for DIC and pCO<sub>2</sub>. This is consistent with  
217 the conclusion in *Koch et al. [2009]* that air-sea gas flux has a minimal influence on the change  
218 of mixed-layer DIC under a hurricane and the decrease in pCO<sub>2</sub> is due to changes in mixed layer  
219 temperature and salinity.

220 To better understand the role of bubbles in air-sea gas flux, especially for CO<sub>2</sub> flux, the total gas  
221 flux between the ocean and the atmosphere, and the respective contribution from the two types of  
222 bubbles and the ocean surface are presented in Figs. 4a to 4c. In the simulation without bubbles,  
223 the total gas flux (black dashed lines in Figs. 4a to 4c) is from the ocean to the atmosphere for all  
224 three gases due to the supersaturation of the gases. In the simulation with bubbles, the total gas  
225 flux for both O<sub>2</sub> and N<sub>2</sub> (black solid lines in Figs. 4a and 4b) is from the atmosphere to the ocean,  
226 in the opposite direction from the surface flux (red lines) and the total flux without bubbles  
227 (black dashed lines). The gas flux through bubbles dominates the surface flux for both O<sub>2</sub> and  
228 N<sub>2</sub>. The fractional contribution by completely dissolved bubbles (blue lines) is larger for N<sub>2</sub> than  
229 for O<sub>2</sub> because the solubility of N<sub>2</sub> is smaller than that of O<sub>2</sub>.

230 With bubbles, the total gas flux of CO<sub>2</sub> is from the ocean to the atmosphere aside from a brief  
231 period between around day 0.58 to day 0.69 when the winds are strongest (Fig. 4c). Under the  
232 simulated conditions, the integrated CO<sub>2</sub> flux during the simulated period is still from the ocean  
233 to the atmosphere but is less than 2% of the flux estimated using a parameterization without  
234 explicit consideration of bubble-induced supersaturation. Except for the first 0.32 days when  
235 CO<sub>2</sub> is more than 8% supersaturated and U<sub>10</sub><26m/s, bubbles contribute to CO<sub>2</sub> dissolution  
236 although the gas is at least 3% supersaturated. In traditional gas flux parameterization without  
237 bubble-induced supersaturation, bubbles will always contribute to gas efflux in supersaturated

238 conditions. This implies that previous studies using a traditional gas flux parameterization  
239 significantly overestimates CO<sub>2</sub> efflux during the passage of a hurricane. Since CO<sub>2</sub> is the most  
240 soluble among the three simulated gases, the relative contribution by bubbles to total gas flux is  
241 the smallest. The effect of gas bubbles becomes evident (the black and the red lines deviate)  
242 when U<sub>10</sub>>30m/s. The contribution of completely dissolved bubbles is much smaller than that  
243 from partially dissolved bubbles for CO<sub>2</sub> that is highly soluble. This is because the fraction of  
244 gas that dissolves through partially dissolved (large) bubbles increases with increasing solubility  
245 while the fraction of a gas through completely dissolved bubbles is determined only by the  
246 atmospheric fraction of the gas.

247 While parameterizations for bubble-mediated gas flux including bubble-induced supersaturation  
248 have been examined for weakly soluble gases such as inert gases [Stanley *et al.* 2009], oxygen  
249 [e.g., Atamanchuk *et al.* 2020] and nitrogen [e.g., Emerson and Bushinsky 2016], they have  
250 never been tested with CO<sub>2</sub>. That is likely because all CO<sub>2</sub> gas flux studies are conducted at high  
251 latitudes where the water is undersaturated. We compare CO<sub>2</sub> flux through bubbles from the LES  
252 solutions with that calculated using three parameterizations, including Nicholson *et al.* [2011],  
253 Liang *et al.* [2013]; and Nicholson *et al.* [2016] (denoted as N11, L13, and N16 hereafter). The  
254 three parameterizations explicitly include the effect of bubble-induced supersaturation that is  
255 required to predict an influx or a reduced efflux under supersaturated conditions. The  
256 parameterizations also compare well with observed concentrations of a few weakly soluble gases  
257 [e.g., Manning *et al.* 2016; Nicholson *et al.* 2016; Liang *et al.* 2017]. The three parameterizations  
258 show different prediction skill for total gas fluxes (Figs. 4d to 4f) and the gas fluxes through  
259 bubbles when compared with the LES models (Figs. 4g to 4i). Note that the three  
260 parameterizations neglect wave condition as an additional parameter. Wave condition regulates  
261 the distribution of breaking waves and boundary layer turbulence, thereby modulating both  
262 bubble entrainment [Deike *et al.* 2017] and bubble penetration [Liang *et al.* 2012]. It was shown  
263 to be important for determining gas flux [e.g., Liang *et al.* 2017] and gas transfer rate [e.g.,  
264 Brumer *et al.* 2017; Deike *et al.* 2018; Esters *et al.* 2018; Reichl and Deike 2020].  
265 Parameterization L13 predicts a CO<sub>2</sub> efflux until day 0.5 and then switches to CO<sub>2</sub> influx  
266 afterward. The change from efflux to influx, also in the LES solutions, is associated with the  
267 drastic decrease in pCO<sub>2</sub> due to the entrainment of subsurface water into the mixed layer. In  
268 parameterizations N11 and N16, the surface efflux dominates the influx through bubbles

269 throughout the studied period. The predicted total CO<sub>2</sub> fluxes by the three parameterizations are  
270 all much larger in magnitude than the results from the LES model (Fig. 4f). It is likely because  
271 the equilibration time is much shorter for CO<sub>2</sub> than for N<sub>2</sub>, O<sub>2</sub> and inert gases (see section 4b of  
272 *Woolf and Thorpe [1991]*), As a result, in many of the partially dissolved bubbles, equilibration  
273 is reached for CO<sub>2</sub>, but not for less soluble gases such as O<sub>2</sub> and N<sub>2</sub>. When generalized for CO<sub>2</sub>,  
274 parameterizations derived from weakly soluble gases assume that equilibration is not reached for  
275 CO<sub>2</sub>, therefore overestimating the flux through bubbles for CO<sub>2</sub> and other highly soluble gases.

## 276 4 Summary

277 This study shows that gas bubbles have a substantial effect on the transport of a soluble gas,  
278 CO<sub>2</sub>. The outgassing of CO<sub>2</sub> over supersaturated water during the passage of a hurricane is  
279 smaller than previous estimates based on traditional gas flux parameterizations. The small efflux  
280 is due to gas bubbles that transfer gases into the ocean even under supersaturated conditions. At  
281 the right flank of Hurricane Frances (2004) close to the location of maximum winds, CO<sub>2</sub> efflux  
282 including the effect of gas bubbles is less than 2% of that calculated using a popular gas flux  
283 parameterization without bubble-induced supersaturation under the simulated conditions. These  
284 results underscore the significance of previously overlooked bubble-induced supersaturation  
285 conditions for CO<sub>2</sub>, a highly soluble gas. They also demonstrate that parameterizations derived  
286 using weakly soluble gases such as O<sub>2</sub>, N<sub>2</sub> and inert gases are not accurate for highly soluble  
287 gases such as CO<sub>2</sub>, and vice versa.

288 This study focuses on changes in mixed layer dissolved gas concentration and air-sea gas flux  
289 during the passage of a hurricane. Although the CO<sub>2</sub> outgassing flux is substantially smaller than  
290 previous estimates, hurricanes likely still have considerable impacts on CO<sub>2</sub> outgassing well after  
291 their passage. As shown in Fig. 3, the mixed-layer DIC concentration after the passage of the  
292 hurricane is substantially higher than its pre-storm values due to interior mixing. Although pCO<sub>2</sub>  
293 is not significantly different right after the passage of the hurricane because of cooling from  
294 entrainment of thermocline water, anomalous warming over the cold wake after the hurricane  
295 [e.g., *Price et al. 2008*] increases pCO<sub>2</sub>, leading to anomalous outgassing in the hurricane wake.  
296 On the other hand, it is also possible that the anomalously high DIC is consumed during the  
297 enhanced plankton bloom after the passage of hurricanes and is exported to the deep ocean  
298 through the sinking of organic particles. Massive ocean phytoplankton blooms after the passage

299 of a hurricane are commonly observed [Lin *et al.* 2003; Walker *et al.* 2005; Liu *et al.* 2019]. The  
300 study of CO<sub>2</sub> flux over a complete hurricane wake would require a regional or a global model  
301 together with a parameterization explicitly including bubble-induced supersaturation. Our  
302 ongoing efforts are to develop a parameterization including bubble-induced supersaturation for  
303 CO<sub>2</sub>, with both wind and wave as parameters. Future field campaign could focus on measuring  
304 CO<sub>2</sub> flux over supersaturated water, those measurements should include air-side fluxes and  
305 concentrations at the two sides of the ocean surface.

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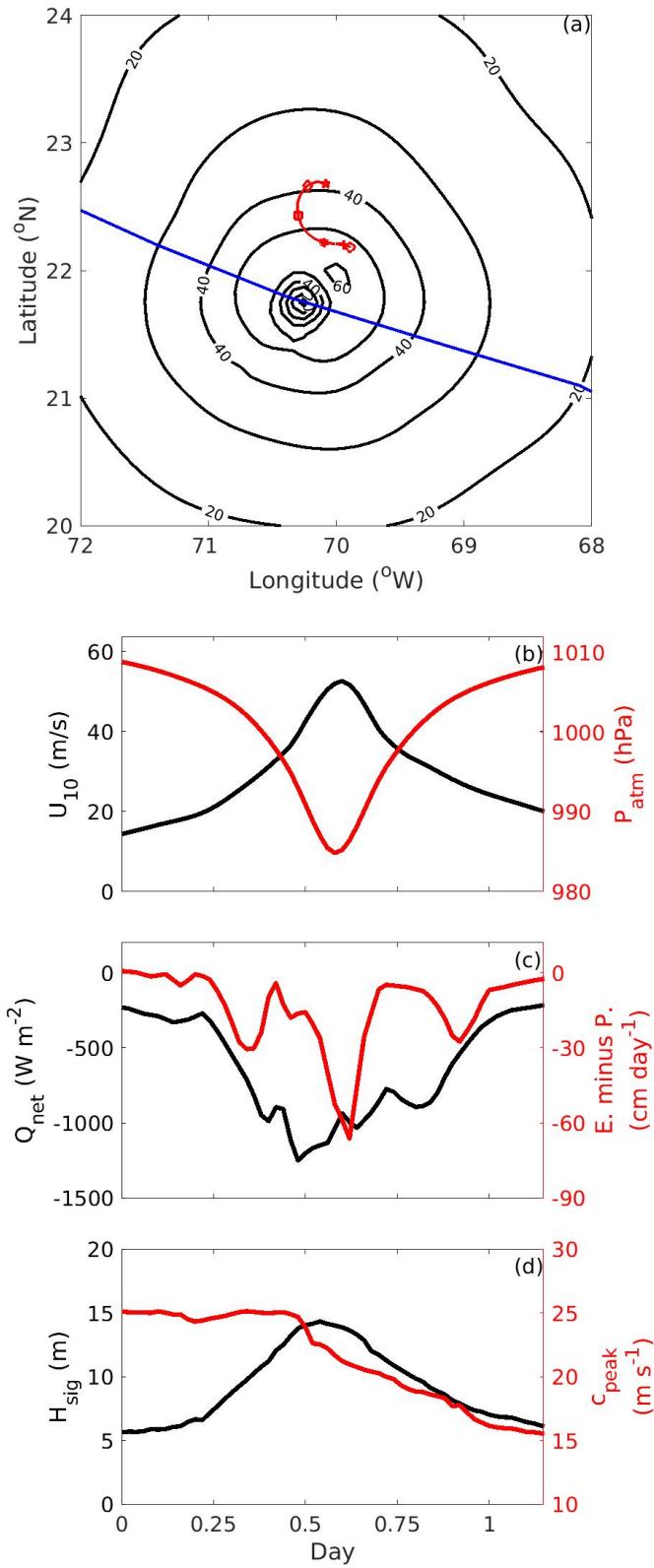
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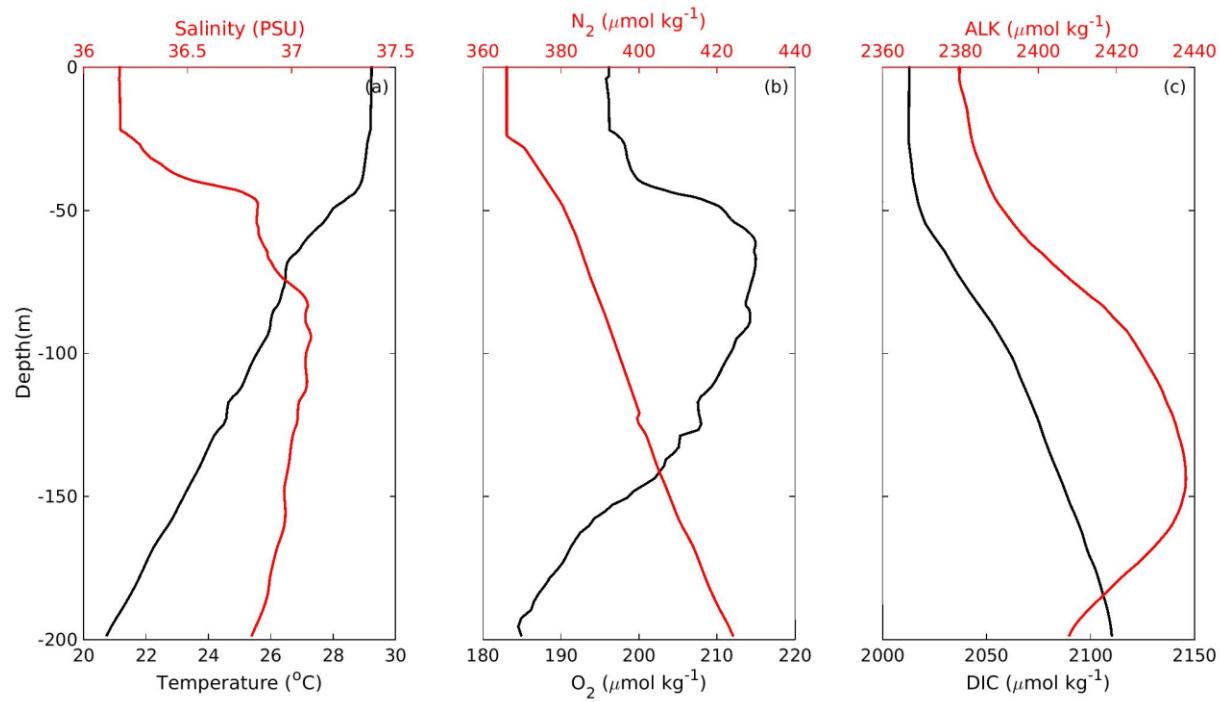




518 Figure 1. (a) Contours of wind speed [m/s] for Hurricane Frances (2004) at day 245.75. The blue  
 519 line is the track of the hurricane. The red line is the path of the float where measurement is taken,  
 520 with symbols indicating a quarter-day increment. At the location of the float, (b) 10-m wind  
 521 speed ( $U_{10}$ ) and sea level pressure ( $P_{atm}$ ). (c) Surface heat flux ( $Q_{net}$ ) and evaporation minus  
 522 precipitation (E.minusP.). (d) Significant wave height ( $H_{sig}$ ) and peak wave speed ( $c_{peak}$ ). Time 0  
 523 corresponds to day 245.15 of year 2004.

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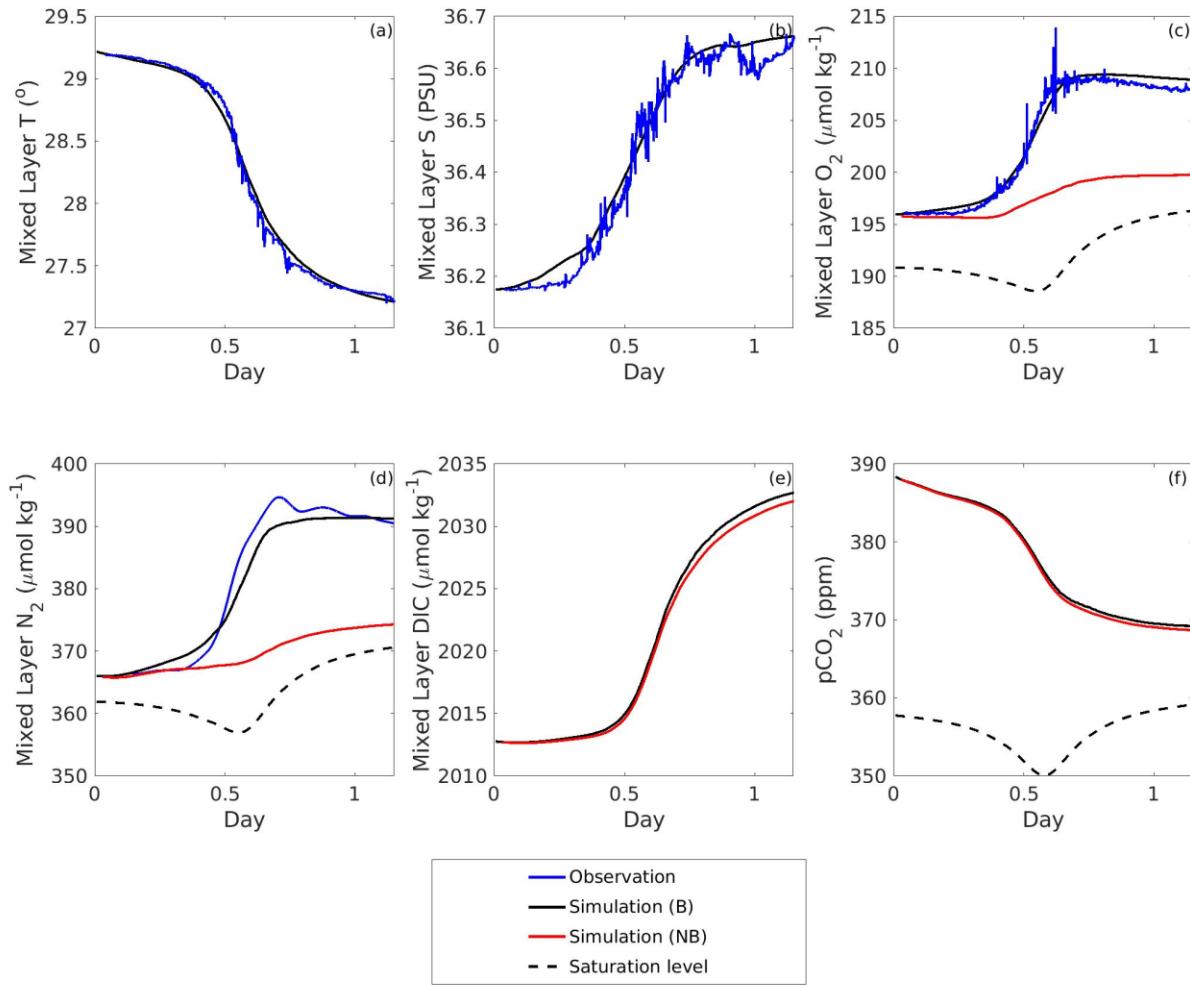


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527 Figure 2. Initial profiles of (a) temperature and salinity; (b) dissolved O<sub>2</sub> and N<sub>2</sub>; and (c) DIC and  
 528 alkalinity.

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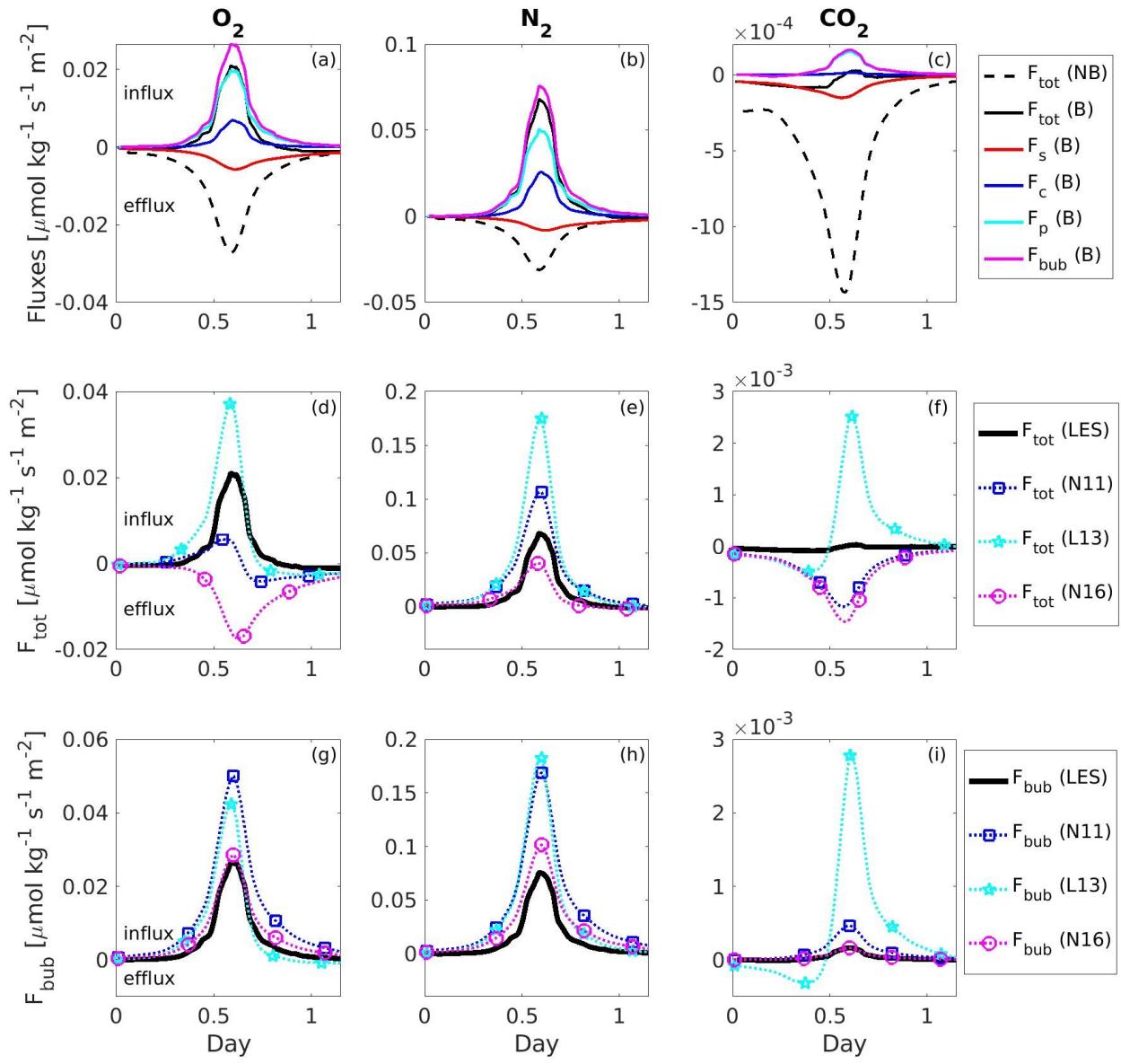
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Figure 3. The evolution of simulated and observed mixed-layer (a) temperature; (b) salinity; (c) dissolved O<sub>2</sub> concentration; (d) dissolved N<sub>2</sub> concentration; (e) DIC; and (f) pCO<sub>2</sub>. Blue lines are from observation; black lines are from simulation with bubbles; red lines are from simulation without bubbles; and dashed lines are the saturation levels. The black line and the blue line overlaps in panels (a) and (b).

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545 Figure 4. Upper row: The evolution of gas fluxes for total gas flux ( $F_{\text{tot}}$ ), surface gas flux ( $F_s$ ),  
 546 gas flux through completely dissolved bubbles ( $F_c$ ), gas flux through partially dissolved bubbles  
 547 ( $F_p$ ), and  $F_{\text{bub}} = F_c + F_p$  in simulation without bubbles (Run NB) and with bubbles (Run B). Middle  
 548 row: The comparison of  $F_{\text{tot}}$  between three parameterizations (N11: *Nicholson et al. 2011*; L13:  
 549 *Liang et al. 2013*; N16: *Nicholson et al. 2016*) and LES solutions. Lower row: The comparison  
 550 of  $F_{\text{bub}}$  between the three parameterizations and LES solutions. The left, middle and right  
 551 columns are for  $\text{O}_2$ ,  $\text{N}_2$ , and  $\text{CO}_2$ , respectively. Positive means influx and negative indicates  
 552 efflux.