CrossMark

# A Survey of Active Galaxies at TeV Photon Energies with the HAWC Gamma-Ray Observatory 

A. Albert ${ }^{1}{ }^{(1)}$, C. Alvarez ${ }^{2}$, J. R. Angeles Camacho ${ }^{3}$, J. C. Arteaga-Velázquez ${ }^{4}$, K. P. Arunbabu ${ }^{5}$ (D) D. Avila Rojas ${ }^{3}$, H. A. Ayala Solares ${ }^{6}$ (1) , V. Baghmanyan ${ }^{7}$ (D) E. Belmont-Moreno ${ }^{3}$ © , S. Y. BenZvi ${ }^{8}$ (i) , C. Brisbois ${ }^{9}$ (iD) K. S. Caballero-Mora ${ }^{2}$ (iD),  B. L. Dingus ${ }^{1}\left(\mathbb{D}\right.$, M. A. DuVernois ${ }^{13}$ (1) , M. Durocher ${ }^{1}$, J. C. Díaz-Vélez ${ }^{12}$ © ${ }^{(1)}$ K. Engel ${ }^{9}$ (1) C. Espinoza ${ }^{3}$ (1), K. L. Fan ${ }^{9}$, M. Fernández Alonso ${ }^{6}$, H. Fleischhack ${ }^{14}$ © , N. Fraija ${ }^{15}$ (1), A. Galván-Gámez ${ }^{15}$, D. García ${ }^{3}$, J. A. García-González ${ }^{3}$ (D), F. Garfias ${ }^{15}$ (1), M. M. González ${ }^{15}$ (1), J. A. Goodman ${ }^{9}$ (1) , J. P. Harding ${ }^{1}$ (1), S. Hernández ${ }^{3}$ (D) B. Hona ${ }^{14}$, D. Huang ${ }^{14}$ (1) , F. Hueyotl-Zahuantitla ${ }^{2}$ (1) , P. Hüntemeyer ${ }^{14}$, A. Iriarte ${ }^{15}{ }^{(1)}$, A. Jardin-Blicq ${ }^{16,17,18}{ }^{(1)}$, V. Joshi ${ }^{19}{ }^{(1)}$, D. Kieda ${ }^{20}$ (D) G. J. Kunde ${ }^{1}$, <br>K. Malone ${ }^{1}{ }^{(1)}$, O. Martínez ${ }^{11}{ }^{(1 D}$, I. Martinez-Castellanos ${ }^{9}$ (1) , J. Martínez-Castro ${ }^{23}$ (1) , J. A. Matthews ${ }^{24}$ (1) ,  M. U. Nisa ${ }^{21}$ (1), R. Noriega-Papaqui ${ }^{25}{ }^{(1)}$, L. Olivera-Nieto ${ }^{16}$, A. Peisker ${ }^{21}$, E. G. Pérez-Pérez ${ }^{22}$ (1) C. D. Rho ${ }^{27}$ (D) ,  H. Schoorlemmer ${ }^{16}$ (10) A. J. Smith ${ }^{9}$ (10) R. W. Springer ${ }^{20}$ © ${ }^{(1)}$, K. Tollefson ${ }^{21}$ © I I. Torres ${ }^{10}$ © , R. Torres-Escobedo ${ }^{12}$, F. Ureña-Mena ${ }^{10}$, L. Villaseñor ${ }^{11}$ (10), T. Weisgarber ${ }^{29}$, E. Willox ${ }^{9}$, A. Zepeda ${ }^{30}{ }^{(1)}$, H. Zhou ${ }^{31}$ (iD), and C. de León ${ }^{4}$ (iD The HAWC collaboration<br>${ }^{1}$ Physics Division, Los Alamos National Laboratory, Los Alamos, NM, USA<br>${ }^{2}$ Universidad Autónoma de Chiapas, Tuxtla Gutiérrez, Chiapas, México<br>${ }^{3}$ Instituto de Física, Universidad Nacional Autónoma de México, Ciudad de México, México<br>${ }^{4}$ Universidad Michoacana de San Nicolás de Hidalgo, Morelia, México<br>${ }^{5}$ Instituto de Geofísica, Universidad Nacional Autónoma de México, Ciudad de México, México<br>${ }^{6}$ Department of Physics, Pennsylvania State University, University Park, PA, USA<br>${ }^{7}$ Institute of Nuclear Physics Polish Academy of Sciences, PL-31342 IFJ-PAN, Krakow, Poland<br>${ }^{8}$ Department of Physics \& Astronomy, University of Rochester, Rochester, NY, USA<br>${ }^{9}$ Department of Physics, University of Maryland, College Park, MD, USA<br>${ }^{10}$ Instituto Nacional de Astrofísica, Óptica y Electrónica, Tonantzintla, Puebla, México; alberto@inaoep.mx, sara@inaoep.mx, annalia@inaoep.mx, danrosa@inaoep.mx<br>${ }^{11}$ Facultad de Ciencias Físico Matemáticas, Benemérita Universidad Autónoma de Puebla, Puebla, México<br>${ }^{12}$ Departamento de Física, Centro Universitario de Ciencias Exactas e Ingenierías, Universidad de Guadalajara, Guadalajara, México<br>${ }^{13}$ Department of Physics, University of Wisconsin-Madison, Madison, WI, USA<br>${ }^{14}$ Department of Physics, Michigan Technological University, Houghton, MI, USA<br>15 Instituto de Astronomía, Universidad Nacional Autónoma de México, Ciudad de México, México<br>${ }^{16}$ Max-Planck Institute for Nuclear Physics, 69117 Heidelberg, Germany<br>${ }^{17}$ Department of Physics, Faculty of Science, Chulalongkorn University, 254 Phayathai Road, Pathumwan, Bangkok 10330, Thailand<br>${ }^{18}$ National Astronomical Research Institute of Thailand (Public Organization), Don Kaeo, MaeRim, Chiang Mai 50180, Thailand<br>${ }^{19}$ Erlangen Centre for Astroparticle Physics, Friedrich-Alexander-Universität Erlangen-Nürnberg, Erlangen, Germany<br>${ }^{20}$ Department of Physics and Astronomy, University of Utah, Salt Lake City, UT, USA<br>${ }^{21}$ Department of Physics and Astronomy, Michigan State University, East Lansing, MI, USA<br>${ }^{22}$ Universidad Politécnica de Pachuca, Pachuca, Hidalgo, México<br>${ }^{23}$ Centro de Investigación en Computación, Instituto Politécnico Nacional, Ciudad de México, México<br>${ }^{24}$ Dept of Physics and Astronomy, University of New Mexico, Albuquerque, NM, USA<br>${ }^{25}$ Universidad Autónoma del Estado de Hidalgo, Pachuca, México<br>${ }^{26}$ Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Ciudad de México, México<br>${ }^{27}$ University of Seoul, Seoul, Republic of Korea<br>${ }^{28}$ Instituto de Física Corpuscular, CSIC, Universidad de València, E-46980, Paterna, València, Spain<br>${ }^{29}$ Department of Chemistry and Physics, California University of Pennsylvania, California, Pennsylvania, USA<br>${ }^{30}$ Departamento de Física, Centro de Investigación y de Estudios Avanzados del IPN, Ciudad de México, México<br>${ }^{31}$ Tsung-Dao Lee Institute \& School of Physics and Astronomy, Shanghai Jiao Tong University, Shanghai, People's Republic of China Received 2020 September 18; revised 2020 November 9; accepted 2020 November 12; published 2021 January 29


#### Abstract

The High Altitude Water Cherenkov (HAWC) Gamma-Ray Observatory continuously detects TeV photons and particles within its large field of view, accumulating every day a deeper exposure of two-thirds of the sky. We analyzed 1523 days of HAWC live data acquired over four and a half years, in a follow-up analysis of 138 nearby $(z<0.3)$ active galactic nuclei from the Third Catalog of Hard Fermi-LAT sources culminating within $40^{\circ}$ of the zenith at Sierra Negra, the HAWC site. This search for persistent TeV emission used a maximum-likelihood analysis assuming intrinsic power-law spectra attenuated by pair production of gamma-ray photons with the extragalactic background light. HAWC clearly detects persistent emission from Mkn 421 and Mkn 501, the two brightest blazars in the TeV sky, at $65 \sigma$ and $17 \sigma$ level, respectively. Marginal evidence, just above the $3 \sigma$ level, was found for three other known very high-energy emitters: the radio galaxy M87 and the BL Lac objects VER J0521 +211 and 1ES $1215+303$, the latter two at $z \sim 0.1$. We find a $4.2 \sigma$ evidence for collective emission


from the set of 30 previously reported very high-energy sources, with Mkn 421 and Mkn 501 excluded. Upper limits are presented for the sample under the power-law assumption and in the predefined (0.5-2.0), (2.0-8.0), and (8.0-32.0) TeV energy intervals.

Unified Astronomy Thesaurus concepts: Active galactic nuclei (16); Blazars (164); Gamma rays (637); Gamma-ray sources (633); Sky surveys (1464); Radio galaxies (1343)

## 1. Introduction

The nuclei of active galaxies are remarkable in possessing observed bolometric luminosities up to $10^{48}-10^{49} \mathrm{erg} \mathrm{s}^{-1}$, surpassing the energy output of their host galaxies, while being unresolved down to the smallest physical scales observable. The stringent Requirements of extreme energy outputs in restricted volumes, together with rapid variability and the presence of powerful relativistic jets, led early on to modeling active galactic nuclei (AGNs) as accreting supermassive black holes of masses up to $10^{9} M_{\odot}$ and accretion rates exceeding $1 M_{\odot} \mathrm{yr}^{-1}$ (Hoyle \& Fowler 1963; Kembhavi \& Narlikar 1999). The geometry of an inner accretion disk and an outer dusty torus surrounding the black hole allows for a qualitative view of the different types of AGNs in terms of the orientation of the line of sight relative to the disk and torus axes aligned with the black hole rotation axis (Antonucci 1993). Anisotropic emission causes overestimates of the AGN luminosity for privileged lines of sight. In standard AGN scenarios, jets can be powered either by the inner regions of the radiation-dominated accretion disk (Blandford \& Payne 1982; Hawley et al. 2015), or by the rapid rotation of the black hole (Blandford \& Znajek 1977; Tchekhovskoy et al. 2011).

AGNs constitute the most common type of $\mathrm{GeV} \gamma$-ray source in the sky. Most of the objects detected with the Large Area Telescope (LAT) on board Fermi are extragalactic, and the vast majority of them are blazars, either BL Lac objects or flatspectrum radio quasars (Acero et al. 2015; Ajello et al. 2017). Relativistic models considering the alignment of AGN jets with our line of sight can account, in most cases, for the high energetics and rapid variability observed in $\gamma$-ray emission up to TeV energies (Dermer \& Schlickeiser 1993; Liao 2018). The detection of radio galaxies and Seyfert galaxies with the Fermi-LAT and Imaging Atmospheric Cherenkov Telescopes (IACTs) provides further insight into the jet-powered view of AGNs under a unified scheme, as probes of off-axis $\gamma$-ray emission (Rieger \& Levinson 2018).

For more than three decades, AGNs have been suspected sources of ultra-high-energy cosmic rays: acceleration may occur in the central engine, the relativistic jet, or distant radio lobes (Hillas 1984; Kotera \& Olinto 2011). The coincidence of the 290 TeV neutrino event IceCube-170922A with the $\gamma$-ray emitting BL Lac TXS 0506+056 provided fresh observational support (IceCube Collaboration et al. 2018). Observations at the highest photon energies are important to characterize the extreme energetics of AGN. Although space-borne instruments conduct deep and wide field-of-view observations leading to all-sky surveys in the high-energy range ( $\mathrm{HE} ; 0.1-100 \mathrm{GeV}$ ), most of our knowledge in the very high-energy (VHE; $\gtrsim 100 \mathrm{GeV}$ ) regime comes from pointed observations with IACTs, nowadays able to detect individual sources with fluxes greater than $1 / 1000$ that of the Crab Nebula. IACT observatories rely on performing deep, but sparse, follow-up observations of active objects; adequate selections of targets; and surveying regions of particular interest. They are limited by their relatively small fields of view, with the most extensive
survey performed with these instruments to date, the dedicated H.E.S.S. Galactic Plane Survey, covering $\sim 0.3 \mathrm{sr}$ in the course of a decade (H.E.S.S. Collaboration et al. 2018).

Compensating their lower instantaneous sensitivity with steradian fields of view, high-altitude air shower arrays are now able to perform unbiased continuous monitoring of known AGNs, in particular Markarian 421 and Markarian 501 (Abeysekara et al. 2017a). This precludes their potential to conduct large surveys with sufficient depth to detect extragalactic TeV sources. The main hurdle for reaching the extragalactic sky is the access to the lowest-photon energies, because our view of the TeV sky is impaired by extragalactic background light (EBL). Photon-photon pair production of $\mathrm{TeV} \gamma$-rays with EBL photons sets a physical limit on how far in distance and in spectral range extragalactic sources can be observed (Salamon \& Stecker 1998). Pair production proceeds efficiently just above its kinematic threshold, making TeV $\gamma$-rays prone to interact with infrared radiation.

This paper presents an AGN follow-up survey performed with four and a half years of full-operations data from the High Altitude Water Cherenkov (HAWC) Gamma-Ray Observatory, building on the preliminary release of this work with a somewhat smaller HAWC data set (Carramiñana et al. 2019). The HAWC time-integrated data cover $\sim 60 \%$ of the sky, extending the Fermi-LAT all-sky survey to a search for persistent $\mathrm{TeV} \gamma$-ray emission. The HAWC survey encompasses all AGNs in the Third Catalog of Hard Fermi-LAT sources (3FHL) accessible from the HAWC site with a redshift $z \leqslant 0.3$, of which about $20 \%$ have been previously reported in the VHE range through IACT-pointed observations in different states of activity. The paper is organized as follows: Section 2 describes the HAWC observatory, its data, and standard analysis; Section 3 provides a brief review of the EBL and its effect on TeV spectra; Section 4 presents the sample of AGNs drawn from the 3FHL catalog, with due considerations of ground-based IACT observations; and Section 5 discusses the dedicated follow-up study of the sample, leading to the summary and conclusions in Section 6.

## 2. The HAWC Gamma-Ray Observatory

HAWC is a wide field-of-view $\mathrm{TeV} \gamma$-ray observatory optimized for surveying cosmic high-energy sources. It is located inside the Parque Nacional Pico de Orizaba, in the Mexican state of Puebla. The HAWC array occupies a relatively flat area of the Volcán Sierra Negra mountain, at an altitude of 4100 m , centered at geographical latitude $18.995^{\circ}$ N and longitude $97.308^{\circ} \mathrm{W}$. HAWC has achieved a $\gtrsim 95 \%$ duty cycle, allowing it to survey two-thirds of the sky every sidereal day with sufficient depth to detect the Crab Nebula at the $5 \sigma$ level (Abeysekara et al. 2017b). This study improves on the 2HWC survey performed with the first year and a half of data, which allowed the detection of 39 sources of $\mathrm{TeV} \gamma$ rays (Abeysekara et al. 2017c). The most recent HAWC allsky survey, $3 H W C$, uses the same data as here. It contains $65 \mathrm{TeV} \gamma$-ray sources, most of them along the Galactic Plane,
including the $189 \sigma$ detection of the Crab Nebula (Albert et al. 2020).

### 2.1. The HAWC Detector

HAWC is an extensive air shower (EAS) array, sampling in detail secondary particles produced by primary cosmic rays in the upper atmosphere. HAWC data analysis can distinguish between hadronic and $\gamma$-ray-induced cascades through their different charge distributions at the ground. The observatory consists of a dense array of 300 large water Cherenkov detectors (WCDs) covering collectively a physical area larger than $22,000 \mathrm{~m}^{2}$. Each WCD is a cylindrical tank of 7.3 m diameter and 5 m height, filled with $180 \mathrm{~m}^{3}$ of water and instrumented with four upward-facing photomultiplier tubes (PMTs) at its base. The signals from the 1200 PMT channels are brought to the data acquisition system, located near the center of the array, to be processed in real time. HAWC has been in full operation since its inauguration on 2015 March 20, after two years of gathering data with a partial array configuration (Abeysekara et al. 2016). Further details about the observatory can be found in Abeysekara et al. (2017b).

### 2.2. HAWC Data and Standard Analysis

The HAWC array records about 25,000 events per second, the vast majority caused by hadronic cosmic rays. Each data record contains the particle arrival timing and deposited charge on each of the 1200 PMTs, which are used to locate the event in the sky and to perform photon/hadron discrimination. The analysis presented here follows the validation observation of the Crab Nebula, both in the $\gamma$-hadron cuts used and in partitioning the data in nine bins according to the fraction of channels hit, as indicated in Table 2 of Abeysekara et al. (2017b). The bin number provides a coarse measure of the primary energy, with an important overlap in the energy distributions of different bins due to the fluctuations inherent in the development of particle cascades. Lower bins relate to lower energies, and the spatial resolution improves with increasing bin number, with the detailed detector response depending on the spectrum of the source and its declination. The lowest bin used in this analysis, $\mathcal{B}=1$, has peak sensitivity around 0.5 TeV for a source with a power-law spectrum of index 2.63 culminating at the zenith, as the Crab Nebula.

The HAWC sky surveys have a mean photon energy of about 7 TeV and one-year sensitivity between 50 and 100 mCrab (Abeysekara et al. 2017c). HAWC data analysis is based on computing the likelihood ratio of a source +background to a background-only model, given by the test statistic,

$$
\begin{equation*}
\mathrm{TS}=2 \ln \left\{\frac{\mathcal{L}(S+B)}{\mathcal{L}(B)}\right\} \tag{1}
\end{equation*}
$$

where $\mathcal{L}(A)$ is the likelihood of model $A$, given by the product of the probability density function computed at each point of the region of interest. The test statistics (Equation (1)) refers to the comparison of a background-only model $(B)$ and a background+source model $(S+B)$. Given a TS value, its statistical significance can be approximated by $s= \pm \sqrt{\mathrm{TS}}$, with the sign indicating an excess or deficit relative to the background. For all-sky surveys, like 2 HWC , the analysis is
performed by optimizing TS on every pixel of a $N_{\text {side }}=1024$ HEALPix grid model of the observable sky (Górski et al. 2005). The source model generally consists of either a pointsource or an extended-source hypothesis following a simple power-law spectrum of fixed spectral index, with free normalization. Joint normalization and spectral index optimizations are then performed to further characterize detected sources. The analysis presented here is performed similarly to that of the 2 HWC , although on predefined sky locations and including the attenuation of TeV photons caused by their interaction with extragalactic background light. We use 1523 days of live data acquired between 2014 November 26 and 2019 June 3. The live data comprises $92.3 \%$ of the total time span. The data deficit is due mostly to quality cuts and run losses during bad weather.

The comparison of HAWC and Fermi-LAT data requires the consideration of the respective systematic uncertainties of each experiment. HAWC fluxes presented here have an estimated $15 \%$ systematic uncertainty (Section 2.2 ); those in the 3 FHL catalog are quoted to have uncertainties of $9 \%$ in the $150-500 \mathrm{GeV}$ band, and $15 \%$ in $0.5-2.0 \mathrm{TeV}$ (Ajello et al. 2017).

## 3. Photon-Photon Attenuation by Extragalactic Background Light

The astrophysical relevance of the $\gamma \gamma \rightarrow \mathrm{e}^{+} \mathrm{e}^{-}$process as an absorption mechanism for distant sources was pointed out by Gould \& Schréder (1966, 1967a, 1967b), first in consideration of the cosmic microwave background, and later for more generic backgrounds. Photon-photon pair production is described by the cross section $\sigma_{\gamma \gamma}=\pi r_{e}^{2} \psi(\omega)$, with $r_{e}$ the classical electron radius, and $\psi$ an analytical function of $\omega=\sqrt{E_{1} E_{2}(1-\mu) / 2}$, the energy of each photon in the center of momentum frame, a relativistic invariant given by the product of the energies of the two photons in an arbitrary frame, $E_{1}=E_{\gamma}, E_{2}=h \nu$, and $\mu=\cos \theta$, where $\theta$ is their interaction angle. Pair production requires $\omega \geqslant$ $m_{e} c^{2}$, with the cross section maximized at $\omega \approx 1.4 m_{e} c^{2} \Rightarrow$ $E_{\gamma} h \nu \approx 2\left(m_{e} c^{2}\right)^{2} \approx 0.5 \mathrm{TeV} \cdot \mathrm{eV}$. Hence, $1 \mathrm{TeV} \gamma$-rays are prone to interact with near-infrared photons of $0.5 \mathrm{eV}(\lambda \simeq$ $2.5 \mu \mathrm{~m}$ ), while a 100 TeV photon is to interact with far-infrared extragalactic light, $\lambda \simeq 250 \mu \mathrm{~m}$.
The absorption of high-energy photons from a distant source of redshift $z$ traversing through intervening radiation is governed by the optical depth,

$$
\begin{equation*}
\tau\left(E_{\gamma}, z\right)=\frac{1}{2} \int_{0}^{d(z)} \int_{0}^{\infty} \int_{-1}^{+1} \sigma_{\gamma \gamma}(\omega) n_{\nu}(\ell) d \mu d \nu d \ell \tag{2}
\end{equation*}
$$

where the photon density $n_{\nu}$ may describe local and/or cosmic intervening radiation fields. ${ }^{32}$ In the case of an evolving cosmic field, a dependence on redshift $n_{\nu}=n_{\nu}(z)$ may be introduced. The photon path is integrated using the cosmological lighttravel distance $d \ell=c d z /(1+z) H(z)$. The probability that the $\gamma$-ray survives the journey is $\exp (-\tau)$. Given the usual functional form of $\tau\left(E_{\gamma}, z\right)$, the survival probability behaves close to a cutoff once $\tau=1$ is reached. In principle, the opacity of the universe to VHE $\gamma$-rays is calculated given $n_{\nu}(z)$. In practice, measurements of the light backgrounds are difficult to

[^0]perform, particularly in the infrared and far-infrared, and observations of distant $\gamma$-ray sources become relevant for constraining the spectral shape of the EBL, both in the local universe and as a function of redshift (Acciari et al. 2019). Here we use the EBL model of Domínguez et al. (2011), which fits well observations by Fermi-LAT and IACTs.

We note that two of the sources in our sample, NGC 1068 and M87, are located below the lower bound of the redshift range of most EBL models, which start at $z=0.01$. For these sources we assume that the photon density has not changed between now and then, to approximate,

$$
\begin{equation*}
\tau\left(E_{\gamma}, z\right) \approx \tau\left(E_{\gamma}, 0.01\right)(z / 0.01) \tag{3}
\end{equation*}
$$

The analysis here assumes intrinsic power-law spectra for the sources. We can describe the overall effect of the EBL, estimating the observed integrated photon flux ( $N_{\text {obs }}$; photons $\mathrm{cm}^{-2} \mathrm{~s}^{-1}$ ) for an intrinsic spectrum of index $\alpha$, introducing the relation,

$$
\begin{align*}
N_{\mathrm{obs}}\left(\geqslant E_{0}\right) & =\int_{E_{0}}^{\infty}\left(\frac{d N}{d E}\right)_{\mathrm{intr}} e^{-\tau(E, z)} d E  \tag{4}\\
& =e^{-z / z_{h}} N_{\mathrm{intr}}\left(\geqslant E_{0}\right),
\end{align*}
$$

where $N_{\text {intr }}\left(\geqslant E_{0}\right) \propto E_{0}^{-\alpha+1} /(\alpha-1)$ is the integral of the differential intrinsic spectrum. The resulting horizon scale $z_{h}$ depends strongly on $E_{0}$ and weakly on the power-law index $\alpha$. Using the EBL model of Domínguez et al. (2011), we get $z_{h}=0.106$ for $\alpha=2.5$ and $E_{0}=0.5 \mathrm{TeV}$, justifying the bound $z \leqslant 0.3$ considered for this study. The value of $z_{h}$ at 0.5 TeV ranges from 0.096 for $\alpha=2.0$ to 0.113 for $\alpha=3.0$; on the other hand, the dependence of $z_{h}$ with $E_{0}$ is exponential, going from $z_{h}=0.728$ at 0.1 TeV to $z_{h}=0.068$ at $E_{0}=$ 1.0 TeV , for $\alpha=2.5$, as presented in Carramiñana et al. (2019).

We compared three main EBL models for the case $\alpha=2.5$, $E_{0}=0.5 \mathrm{TeV}$. When considering the upper and lower uncertainties of Domínguez et al. (2011), we get that $z_{h}$ is in the interval $(0.098,0.121)$, which is consistent with $z_{h}=0.102$ obtained with Gilmore et al. (2012) and $z_{h}=0.099$ with Franceschini \& Rodighiero (2017). These three models coincide within a $\sim 10 \%$ systematic uncertainty.

## 4. Active Galaxies above $10 \mathrm{GeV}: 3 \mathrm{FHL}$ and TeV Pointed Observations

### 4.1. A Sample of Active Galaxies from the 3FHL Catalog

The 3FHL catalog contains 1556 objects detected at photon energies between 10 GeV and 2 TeV in the first seven years of Fermi operations, from 2008 August 4 to 2015 August 2 (Ajello et al. 2017). Of the 3FHL entries, 79\% are identified or associated with extragalactic objects, mostly BL Lac objects and flat-spectrum radio quasars ( $48 \%$ and $11 \%$ of the 3FHL, respectively). As defined in the different LAT catalogs, an association refers to the positional coincidence of the HE $\gamma$-ray source with an object having suitable properties, while an identification requires measuring correlated variability between the $\gamma$-ray source and its associated counterpart. These criteria result in $9 \%$ of the sources in 3FHL listed as identified and 78\% as associated, while the remaining $13 \%$ are unassociated or unclassified. It is customary in Fermi catalogs to distinguish between identifications using uppercase letters, such as RDG

Table 1
Classes of AGN Selected

| Source Class | Number of Sources |  |  |
| :--- | :---: | :---: | :---: |
|  | Identified | Associated | Total |
| BL Lac objects (BLL + bll) | 6 | 111 | 117 |
| Blazars candidates of uncertain <br> type (bcu) <br> Radio galaxies (RDG + rdg) <br> Flat-spectrum radio quasars (FSRQ <br> $\quad$fsrq) <br> Starburst galaxies (SBG + sbg) <br> Total number of sources in sample | 1 | 8 | 8 |

for identified radio galaxies, and associations using lowercase letters, such as rdg for an association with a radio galaxy.

The 3FHL catalog partitions its nominal wide $10 \mathrm{GeV}-$ 2 TeV spectral interval into five bands. The majority of AGNs are detected at TS $>10$ in the two lower-energy bands, 10-20 and $20-50 \mathrm{GeV}$. Spectral cutoffs at energies $\lesssim 50 \mathrm{GeV}$ are common, as can be noticed in our sample: of the 138 AGNs, 50 (14) are detected above $5 \sigma$ in the LAT $50-150 \mathrm{GeV}$ ( $150-500 \mathrm{GeV}$ ) intermediate band(s). Furthermore, Mkn 421 and Mkn 501 are the only two AGNs studied here with $\mathrm{TS}>25$ in the $0.5-2.0 \mathrm{TeV}$ LAT band, the spectral intersection with HAWC.

The 3FHL catalog assigns a flag to sources as follows: $\mathrm{TeV}=\mathrm{P}$, when reported at VHE energies; $\mathrm{TeV}=\mathrm{C}$ for candidates for TeV detection; and $\mathrm{TeV}=\mathrm{N}$ for nonreported and not candidates. TeV candidates are by definition sources undetected with VHE ground-based instruments whose LAT data satisfy three conditions: (i) significance $s>3$ above 50 GeV ; (ii) spectral index $<3$; (iii) integrated photon flux $N(>50 \mathrm{GeV})>10^{-11} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$. In addition, the 3FHL catalog assesses source variability through the $V_{\text {bayes }}$ parameter, the number of Bayes blocks needed to model the light curve. A source with $V_{\text {bayes }}=1$ is consistent with a constant flux (Ajello et al. 2017).

We select 3FHL catalog sources identified or associated with AGNs with redshifts $z \leqslant 0.3$ that culminate within $40^{\circ}$ of the zenith as viewed from the HAWC site. For the selection of the follow-up sample we used the current version of the 3FHL catalog available at the Fermi Science Support Center. ${ }^{33}$ The sample of 138 objects is summarized in Table 1. It contains 32 objects flagged $\mathrm{TeV}=\mathrm{P}$ (positive VHE detections), and $32 \mathrm{TeV}=\mathrm{C}$ (candidates). The sample is grouped in five source classes, defined in the 3FHL, of distinct properties (Table 1):

1. Starburst galaxies are the nearest and apparently less luminous AGNs in Fermi-LAT. While prone to host active nuclei, the prevailing $\gamma$-ray emission is dominated by cosmic rays produced in star formation processes. NGC 1068 is the only 3FHL starburst inside the decl. range of our selection, and the lowest redshift AGN in our sample (Figure 1). Intriguingly, it has recently been associated with one of the hot spots in the neutrino sky, as observed by the IceCube observatory (Aartsen et al. 2020).

[^1]

Figure 1. Luminosity as a function of redshift for the 3FHL sample studied with HAWC. Note the clear separation in redshift and luminosities between the single starburst (yellow point, the nearest object), radio galaxies (red), intermediate distance bcus (in blue), and the more distant group of FSRQs (orange). BL Lac objects, in gray, span most of the redshift interval and have the highest observed luminosities.
2. Radio galaxies (RDG) are the nearest extragalactic GeV sources dominated by an active nucleus. They appear up to three orders of magnitude more luminous than starbursts (Figure 1). Radio galaxies are attractive targets for HAWC because the relatively close distance translates in a reduced photon-photon attenuation, potentially allowing sampling of the far-infrared portion of the EBL through observations above $\sim 10-30 \mathrm{TeV}$. Six 3FHL catalog radio galaxies transit through the field of view of HAWC, with redshifts between $z=0.0042$, for M87, and $z=0.029$, for NGC 1218. Four of them have been claimed as VHE sources by IACT collaborations ${ }^{34}$ (Rieger \& Levinson 2018).
3. BL Lacertae objects (BLL) constitute the majority of known GeV and VHE $\gamma$-ray emitters. They dominate the 3FHL catalog, in particular for redshifts $z \lesssim 0.7$. As expected, BL Lac objects completely dominate our sample with 6 identifications and 111 associations. They span most of the $z \leqslant 0.3$ range, led by Mkn 421 at $z=0.031$, the nearest and brightest 3FHL AGN.
4. Flat-spectrum radio quasars (FSRQ) are the most distant and seemingly luminous blazars. Our sample includes six such sources, with redshifts ranging from $z=0.158$ to $z=0.222$. Of these, only PKS $0736+017$ has been reported as a TeV source, the nearest FSRQ claimed in the VHE range so far (H.E.S.S. Collaboration et al. 2020).
5. Blazars candidates of uncertain type (bcu) are AGNs poorly characterized across the electromagnetic spectrum. The 3FHL catalog contains 290 bcus, a good fraction of them in the Southern sky, all with radio-loud associations, and $90 \%$ of them lacking redshift measurements. Only eight such objects satisfy our selection criteria. They are relatively near objects, between $z=0.036$ and $z=0.128$ (Figure 1). None is a VHE source, but three of them are flagged as TeV candidates.

[^2]Figure 1 shows the luminosity per solid angle unit, $d_{L}(z) f_{e}$, with $f_{e}$ being the ( $10 \mathrm{GeV}-1 \mathrm{TeV}$ ) energy flux from the 3 FHL catalog, and $d_{L}(z)$ the luminosity distance. The actual luminosity of each source depends on the unknown solid angle of emission.

Prior to the analysis, we identified five objects in our sample close in the sky to bright 2HWC sources that could affect their analysis. We set a conservative distance threshold of $5^{\circ}$, equivalent to five times the $68 \%$ containment radius of $\mathcal{B}=1$, the bin where contamination by a bright source nearby is more likely. As our sample excludes by construction low Galactic latitudes, with all selected sources located at $|b|>5^{\circ}$, the only concerns were for:

1. 3FHL J0521.7+2112, located at 3.07 degrees from the Crab Nebula;
2. 3FHL J1041.7+3900, 3FHL J1100.3 + 4020, and 3FHL $\mathrm{J} 1105.8+3944$, located respectively at $4.51,2.28$, and 1.57 degrees from Markarian 421;
3. 3FHL J1652.7+4024, located at 0.68 degrees from Markarian 501.

Only 3FHL J0521.7 +2112 is associated with a well-known VHE source, the BL Lac object VER J0521+211. The other four sources are flagged as $\mathrm{TeV}=\mathrm{N}$. These five objects were analyzed and tested later for contamination from the bright neighbor source, as detailed for VER J0521+211 in Section 5.3.2. From the respective tests we decided to exclude 3FHL J1105.8 +3944 and 3FHL J1652.7 +4024 from our sample.
We also revised the redshift measurements listed in the 3FHL catalog. Our sample is dominated by BL Lac objects, often with questionable redshifts. Redshifts from the 3FHL catalog were systematically collated with the SIMBAD and NED databases, with the Sloan Digital Sky Survey (SDSS) database consulted in some particular cases. We used as references the redshift surveys of Shaw et al. (2013) and Healey et al. (2008), and in particular the dedicated survey of TeV sources performed with the 10.4 m Gran Telescopio Canarias (GTC) by Paiano et al. (2017). While we decided to systematically use the redshifts listed in the 3FHL catalog, we list in Table 2 sources with disputed values, for future reference.

### 4.2. VHE Ground-based Observations Related to Our Sample

AGNs have been extensively studied with IACTs, with evidence of VHE $\gamma$-ray emission in almost a hundred of them (Wakely \& Horan 2008; Madejski \& Sikora 2016). These observations have shown that TeV flaring is an intrinsic characteristic of blazars and radio galaxies. IACT deep observations, able to reach down a few percent of the flux of the Crab Nebula in a single run, have identified high, medium, and low states of activity in several sources. However, continuous long-term coverage of the AGNs population cannot be performed with IACTs. The definition of the base level of AGN VHE emission is a pending task. Despite their lower instantaneous sensitivity, EAS arrays with efficient $\gamma /$ hadron discrimination can quantify better time-averaged fluxes, integrated over long periods of time, while monitoring for flaring activity, as they continuously drift over large portions of the sky.

The first significant AGN detections at TeV energies were those of Mkn 421 and Mkn 501 (Punch et al. 1992;

Table 2
Sources with Redshifts Under Question

| 3FHL Name <br> (1) | Associated Source <br> (2) | Redshift (3) | Other Redshifts <br> (4) | Notes (5) | References <br> (6) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 3FHL J0112.1+2245 | S2 0109+22 | 0.265 | $z>0.35, z<0.67$ | Detected up to 200 GeV . | (1), (2), (3) |
| 3FHL J0521.7+2112 | TXS 0518+211 | 0.108 | $>0.18$ | Emission up $\gtrsim 1 \mathrm{TeV}$. | (4), (5), (6) |
| 3FHL J0650.7+2503 | 1ES 0647+250 | 0.208 | $>0.29$ | $z \simeq 0.41$ host properties. | (4), (5), (7) |
| 3FHL J1230.2+2517 | ON 246 | 0.135 | $>0.10$ | ... | (5), (8) |
| 3FHL J2323.8+4210 | BZB J2323+4210 | 0.059 | $\geqslant 0.267$ | TeV candidate | (5) |

Note. Column (2) shows the object associated with the 3FHL source, whose redshift is under question; column (3) display the entry for the redshift in the 3 FHL catalog; column (4) shows other redshift estimates, or bounds; column (6) shows the reference for the redshift value/bound and for the comments shown in column (5).

References. (1) Healey et al. (2008); (2) Paiano et al. (2016); (3) MAGIC Collaboration et al. (2018); (4) Shaw et al. (2013); (5) Paiano et al. (2017); (6) Prokoph et al. (2015); (7) Kotilainen et al. (2011); (8) Nass et al. (1996).

Quinn et al. 1996). These two Markarian galaxies have been extensively studied in the TeV range for more than two decades. They are highly variable but remain bright enough over long periods of time to have been detected by EAS arrays like the Tibet Air Shower Array and MILAGRO, both providing first unbiased views of their TeV emission on timescales of years (Bartoli et al. 2011; Abdo et al. 2014). In the last five years, HAWC has been performing an increasingly deeper monitoring of these two BL Lac objects (Abeysekara et al. 2017a). Measurements of their long-term averaged emission are presented in this paper.
IACTs have been able to go deeper and beyond these two well-known blazars by implementing sophisticated technologies to achieve lower energy thresholds, around or below 100 GeV , with improved sensitivities. This has permitted us to peer through the EBL horizon up to redshifts $z \lesssim 0.9$ (Abeysekara et al. 2015; Ahnen et al. 2015), while sampling different types of extragalactic sources, such as:

1. The starbursts galaxies M82 and NGC 253, detected by VERITAS and H.E.S.S. respectively, with fluxes below $1 \%$ of the Crab (Ohm 2016).
2. Four Faranoff-Riley type I radio galaxies: Centaurus A (too Southern for HAWC); NGC 1275, the most prominent galaxy of the massive Perseus cluster; 3C 264, newly reported as a VHE source; and M87, the massive central elliptical galaxy in the nearby Virgo cluster. We note the ambiguous classification of IC 310 and PKS 0625-35, referred as RDG in the 3FHL, and as unknown type of AGN in TeVCat (Wakely \& Horan 2008) and by Rieger \& Levinson (2018).
3. At least 60 BL Lacertae objects as VHE $\gamma$-ray sources (Wakely \& Horan 2008; Ajello et al. 2017), mostly high-frequency peaked BL Lac objects (HBL; 50 sources), plus a few intermediate-frequency peaked BLL (IBL; 8 sources, including BL Lac, W Comae, and VER J0521 +211 , in our sample); and only two lowfrequency peak BL Lacs (LBL). While most of the IACT observations have been reported with sub- TeV thresholds, we note the observations of H $1426+428$ by VERITAS and HEGRA at energies between 0.25 and 2.5 TeV , reporting fluxes from $3 \%$ to $10 \%$ of the Crab (Aharonian et al. 2002; Petry et al. 2002).
4. Seven FSRQs, of which PKS $0736+017$ is the only FSRQ detected at VHE with $z<0.3$ and entering the field of view of HAWC. This object was found in a flaring state at a flux level of 100 mCrab between 100 and


Figure 2. Photon flux extrapolations to $E>0.5 \mathrm{TeV}$ using the spectral models of the 3FHL catalog and $\gamma \gamma$ attenuation by the EBL. Colors in the dots indicate the TeV flag assigned in the 3 FHL catalog. The horizontal red dotted line represents a flux equivalent to 30 mCrab , which we take as a benchmark of potential HAWC detectability. We identify the VHE sources VER J0521+211 and 1ES $1215+303$ with respective green rings below the red line.

300 GeV , and showing a steep spectrum (H.E.S.S. Collaboration et al. 2020).

### 4.3. Photon Flux Extrapolations

We calculated the expected integral photon flux above 0.5 TeV by extrapolating the spectral models and parameters listed in the 3FHL. Although in HAWC a distinction is made between the intrinsic spectrum of a source (as emitted) and the related observable parameters (attenuated by the EBL), FermiLAT spectral models do not need to make this distinction. The difference is minor in most of the LAT energy regime, with only a moderate increase in spectral indices for sources with $z \gtrsim 1$ (Ajello et al. 2017). Therefore, we added the attenuation effect of the EBL to the LAT spectral models. The extrapolated fluxes as a function of redshift are shown in Figure 2. Note the uncertainties in the flux extrapolations due to the uncertainties in the fit parameters, specially as they are propagated more than one order of magnitude above the 3FHL pivot energies. The horizontal red line corresponds to a photon flux of $3 \%$ of the Crab Nebula, indicative of potential HAWC detectability within the current data. Mkn 421 and Mkn 501 stand clearly at about an order of magnitude above the red line, followed by $1 \mathrm{ES} 2344+514$ and I Zw 187 (1ES1727+502), two of the
nearest BL Lac objects known. The radio galaxies M87 and IC 301 are also above the flux level of interest, although with fairly large uncertainties in the extrapolation. Of the sources potentially detectable, only TXS $0210+515$ is tagged as undetected in the VHE range in the 3FHL. However, we note its recent detection with MAGIC (Acciari et al. 2020).
We also computed photon flux extrapolations of VHE sources within TeVCat, using simple power laws and approximating the EBL attenuation by a direct exponential in redshift (Section 3). The predictions from this second set are also rather uncertain, because IACTs observations are by nature short and sparse, reflecting different activity states, and spectra are not always fitted above 1 TeV . Five candidates stood out with extrapolated fluxes $N(>0.5 \mathrm{TeV}) \geqslant 20 \mathrm{mCrab}$, four of them in common with the 3FHL extrapolation: M87, Mkn 421, Mkn 501, H 1426+428, and 1ES 2344+514. Results for the sources named in this section are shown and discussed in Section 5.

## 5. HAWC Follow-up Survey of 3FHL AGNs

The HAWC follow-up survey consists of a systematic search for $\mathrm{TeV} \gamma$-ray emission from each of the 138 AGNs selected from the Fermi-LAT 3FHL catalog. We performed a maximum-likelihood test assuming a point source at the location of the presumed 3FHL counterpart with an intrinsic power-law differential flux spectra of index $\alpha$ attenuated by the EBL,

$$
\begin{equation*}
\left(\frac{d N}{d E}\right)_{\mathrm{obs}}=K\left(E / E_{0}\right)^{-\alpha} e^{-\tau(E, z)} \tag{5}
\end{equation*}
$$

A first analysis run was performed with a fixed spectral index, $\alpha=2.5$, fitting only the normalization $K$ for a pivot energy $E_{0}=1 \mathrm{TeV}$. From this we obtained:

1. The compilation of TS values and 1 TeV normalizations $K$ for the revised sample of 136 AGN; ${ }^{35}$
2. Statistics of significances ( $s \equiv \pm \sqrt{\mathrm{TS}}$ ) for source classes and TeV flags. Following Wilk's theorem, under the null hypothesis the behavior of $s$ tends to a Gaussian distribution, of mean $\mu(s) \rightarrow 0$ and standard deviation $\sigma(s) \rightarrow 1 / \sqrt{N}$, for $N$ points;
3. The corresponding list of $2 \sigma$ ( $\sim 95 \%$ confidence level) upper limits on the normalization, $K_{2 \sigma}$, computed under the Feldman \& Cousins (1998) approach, which allows only non-negative fluxes;
4. The comparison of flux values and upper limits with the extrapolations performed in Section 4.3;
5. The characterization of upper limits as function of declination and redshift.

For those sources with test statistic TS $>9$ we then computed optimized spectral fits, allowing both $K$ and $\alpha$ to vary. The discussion of the sources complying with this criterion is presented in Sections 5.2 and 5.3. Finally, we also computed quasi-differential limits for all the sources in the sample in three energy intervals, (0.5-2.0), (2.0-8.0), and (8.0-32) TeV, following the procedure implemented in IceCube Collaboration et al. (2017). The first energy interval, the HAWC overlap with 3FHL, is where we expect AGN emission to be brighter, while the second interval includes the peak of HAWC sensitivity. The

[^3](8.0-32.0) TeV range put bounds where the EBL attenuation is more severe. These quasi-differential limits are computed assuming $\alpha=2.0$ without consideration of EBL attenuation, as detailed in Section 4.3.3 of IceCube Collaboration et al. (2017).

### 5.1. The $\alpha=2.5$ AGN Search

The default spectral model for the AGN search was an index $\alpha=2.5$, as used in the 3HWC (Albert et al. 2020). However, note that here the index refers to the intrinsic spectrum and that the observed spectra will be softer. While we know from previous analyses that this index can represent fairly well the data from both Mkn 421 and Mkn 501 (Coutiño de Leon et al. 2019), photon indices in extragalactic 3FHL sources peak in the interval between 2.0 and 2.5 (Ajello et al. 2017).

### 5.1.1. Overall Results

The results for the complete AGN sample are shown in Table 3. The left-hand side of Figure 3 shows the histogram of significances for the sample, also displayed as the first entry of Table 4. Mkn 421 is detected with a significance $\sqrt{\mathrm{TS}}=+64.6$, and Mkn 501 with $\sqrt{\mathrm{TS}}=+16.6$. These two high-significance detections drive up the statistics of the complete sample to an overall $9 \sigma$ deviation from the null hypothesis. When removing Mkn 421 and Mkn 501, the joint significance drops to a $p$-value of $2.2 \%$ (Figure 3, righthand side).

Three more objects showed marginally significant test statistic $\mathrm{TS}>9$ : M87 (TS = 12.9); 1ES 1215+303 (TS = 11.4); and VER J0521+211 ( $\mathrm{TS}=9.5$ ). We note that the random probability of having three TS $>9$ values in 134 trials, once Mkn 421 and Mkn 501 are excluded, is $8.5 \times 10^{-4}$. These three tentative sources have favorable declinations for HAWC, $\left|\delta-19^{\circ}\right| \lesssim 10^{\circ}$, allowing for transits of more than 6 hr . Spectral fits for the five TS $>9$ sources are presented in Section 5.3.

### 5.1.2. Upper Limits and Sensitivity

The $K_{2 \sigma}$ upper limits span a large range of values: from $2.4 \times 10^{-13} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ (3FHL J1543.6+0452 at $z=$ 0.040 ) to $10^{-10} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ (3FHL J0325.6-1646 at $z=$ 0.291 ). Upper limits are sensitive to both the declination and redshift of the source, as shown in Figure 4. We performed a fit to the set of $K_{2 \sigma}$ values that assumes a Gaussian dependence in decl. and exponential in redshift. Three parameters were computed: the normalization limit at the reference point $\delta=19^{\circ}$ and $z=0$, the angular width in decl., and the redshift exponential scale. The fit obtained is

$$
\begin{align*}
\log _{10} K_{2 \sigma}= & -12.33+\frac{1}{2 \ln 10}\left(\frac{\delta-19^{\circ}}{18^{\circ} .93}\right)^{2}  \tag{6}\\
& +\frac{1}{\ln 10}\left(\frac{z}{0.089}\right)
\end{align*}
$$

This fit excludes the five $\mathrm{TS}>9$ sources. Its correlation coefficient is $r=+0.870$ and the dispersion with the data is 0.241 dex. The value for the fit of $K_{2 \sigma}$ at $z=0, \delta=19^{\circ}$ corresponds to 14 mCrab , a factor of two lower than the predefined depth of the survey. This would apply to near and optimally located sources only. The sensitivity of the survey degrades with a Gaussian angle of 18.9 degrees, leading to a

Table 3
HAWC Power-law Fits and Significances for the Revised Sample of 136 AGNs from the 3FHL Catalog

| 3FHL Source <br> (1) | Counterpart <br> (2) | Class <br> (3) | Redshift (4) | $\begin{aligned} & \text { TS } \\ & \text { (5) } \end{aligned}$ | $\begin{gathered} \pm \sqrt{\mathrm{TS}} \\ (6) \end{gathered}$ | $K \pm \Delta K$ <br> (7) | $\begin{gathered} K_{2 \sigma} \\ (8) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| J0007.9+4711 | RX J0007.9+4711 | bll | 0.2800 | -4.49 | -2.12 | $-31.9 \pm 15.1$ | 9.16 |
| J0013.8-1855 | RBS 0030 | bll | 0.0949 | -0.10 | -0.32 | $-2.20 \pm 6.81$ | 11.5 |
| J0018.6+2946 | RBS 0042 | bll | 0.1000 | +0.01 | +0.11 | $+0.10 \pm 0.93$ | 1.97 |
| J0037.8+1239 | NVSS J003750+123818 | bll | 0.0890 | +0.47 | +0.69 | $+0.52 \pm 0.76$ | 2.03 |
| J0047.9+3947 | B3 0045+395 | bll | 0.2520 | -0.18 | -0.42 | $-2.68 \pm 6.38$ | 10.3 |
| J0056.3-0936 | TXS 0053-098 | bll | 0.1031 | -0.21 | -0.46 | $-1.33 \pm 2.88$ | 4.53 |
| J0059.3-0152 | 1RXS J005916.3-015030 | bll | 0.1440 | +0.95 | +0.98 | $+2.61 \pm 2.67$ | 7.93 |
| J0112.1+2245 | S2 $0109+22$ | BLL | 0.2650 | +0.33 | +0.57 | $+1.99 \pm 3.47$ | 8.88 |
| J0123.0+3422 | 1ES 0120+340 | bll | 0.2720 | +6.53 | +2.56 | $+13.1 \pm 5.13$ | 23.4 |
| J0131.1+5546 | TXS 0128+554 | bcu | 0.0365 | -0.85 | -0.92 | $-1.61 \pm 1.75$ | 2.10 |
| J0152.6+0147 | PMN J0152+0146 | bll | 0.0800 | -0.89 | -0.94 | $-0.89 \pm 0.94$ | 1.10 |
| J0159.5+1047 | RX J0159.5+1047 | bll | 0.1950 | +1.09 | +1.04 | $+2.52 \pm 2.41$ | 7.31 |
| J0211.2+1051 | MG1 J021114+1051 | BLL | 0.2000 | -0.58 | -0.76 | $-1.92 \pm 2.51$ | 3.26 |
| J0214.5+5145 | TXS $0210+515$ | bll | 0.0490 | +2.63 | +1.62 | $+2.43 \pm 1.51$ | 5.43 |
| J0216.4+2315 | RBS 0298 | bll | 0.2880 | -0.08 | -0.29 | $-1.14 \pm 3.97$ | 6.74 |
| J0217.1+0836 | ZS 0214+083 | bll | 0.0850 | +4.58 | +2.14 | $+1.70 \pm 0.80$ | 3.31 |
| J0219.1-1723 | 1RXS J021905.8-172503 | bll | 0.1287 | +0.01 | +0.10 | $+1.0 \pm 9.9$ | 20.9 |
| J0232.8+2017 | 1ES $0229+200$ | bll | 0.1400 | -0.28 | -0.53 | $-0.70 \pm 1.32$ | 1.98 |
| J0242.7-0002 | NGC 1068 | sbg | 0.0038 | +1.45 | +1.21 | $+0.24 \pm 0.20$ | 0.65 |
| J0308.4+0408 | NGC 1218 | rdg | 0.0288 | +0.24 | +0.49 | $+0.17 \pm 0.35$ | 0.87 |
| J0312.8+3614 | V Zw 326 | bll | 0.0710 | +0.48 | +0.69 | $+0.53 \pm 0.77$ | 2.07 |
| J0316.6+4120 | IC 310 | RDG | 0.0189 | +0.74 | +0.86 | $+0.31 \pm 0.36$ | 1.01 |
| J0319.8+1845 | RBS 0413 | bll | 0.1900 | +0.05 | +0.22 | $+0.45 \pm 2.08$ | 4.64 |
| J0319.8+4130 | NGC 1275 | RDG | 0.0176 | +0.36 | +0.60 | $+0.21 \pm 0.35$ | 0.91 |
| J0325.6-1646 | RBS 0421 | bll | 0.2910 | +0.07 | +0.27 | $+12.6 \pm 46.4$ | 107 |
| J0326.3+0226 | 1H 0323+022 | bll | 0.1470 | -0.22 | -0.47 | $-1.02 \pm 2.16$ | 3.35 |
| J0334.3+3920 | $4 \mathrm{C}+39.12$ | rdg | 0.0203 | -0.91 | -0.96 | $-0.33 \pm 0.34$ | 0.40 |
| J0336.4-0348 | 1RXS J033623.3-034727 | bll | 0.1618 | +0.19 | +0.44 | $+1.64 \pm 3.75$ | 9.10 |
| J0339.2-1736 | PKS 0336-177 | bcu | 0.0656 | +0.27 | +0.52 | $+1.83 \pm 3.55$ | 8.97 |
| J0349.3-1159 | 1ES 0347-121 | bll | 0.1850 | +0.74 | +0.86 | $+9.00 \pm 10.5$ | 30.0 |
| J0416.8+0105 | 1ES 0414+009 | bll | 0.2870 | +1.71 | +1.31 | $+9.16 \pm 7.00$ | 22.9 |
| J0424.7+0036 | PKS 0422+00 | bll | 0.2680 | +1.60 | +1.27 | $+8.17 \pm 6.46$ | 21.1 |
| J0521.7+2112 | TXS 0518+211 | bll | 0.1080 | +9.49 | +3.08 | $+2.85 \pm 0.93$ | 4.72 |
| J0602.0+5316 | GB6 J0601+5315 | bcu | 0.0520 | -0.17 | -0.42 | $-0.79 \pm 1.88$ | 3.01 |
| J0617.6-1715 | TXS 0615-172 | bll | 0.0980 | +1.51 | +1.23 | $+7.46 \pm 6.07$ | 19.7 |
| J0648.7+1517 | RX J0648.7+1516 | bll | 0.1790 | +0.26 | +0.51 | $+1.00 \pm 1.97$ | 4.92 |
| J0650.7+2503 | 1ES 0647+250 | bll | 0.2030 | +0.01 | +0.10 | $+0.24 \pm 2.43$ | 5.12 |
| J0656.2+4235 | $4 \mathrm{C}+42.22$ | bll | 0.0590 | -0.65 | -0.81 | $-0.71 \pm 0.88$ | 1.13 |
| J0725.8-0056 | PKS 0723-008 | bcu | 0.1270 | -2.76 | -1.66 | $-3.48 \pm 2.09$ | 1.58 |
| J0730.4+3307 | 1RXS J073026.0+330727 | bll | 0.1120 | +0.96 | +0.98 | $+1.21 \pm 1.23$ | 3.66 |
| J0739.3+0137 | PKS 0736+01 | fsrq | 0.1910 | -0.96 | -0.98 | $-3.40 \pm 3.47$ | 3.93 |
| J0753.1+5354 | $4 \mathrm{C}+54.15$ | bll | 0.2000 | +0.78 | +0.88 | $+15.2 \pm 17.3$ | 49.6 |
| J0757.1+0957 | PKS 0754+100 | bll | 0.2660 | -0.12 | -0.34 | $-1.39 \pm 4.04$ | 6.80 |
| J0809.7+3457 | B2 0806+35 | bll | 0.0830 | +0.09 | +0.30 | $+0.26 \pm 0.87$ | 2.03 |
| J0809.8+5218 | 1ES 0806+524 | bll | 0.1371 | +0.27 | +0.52 | $+3.59 \pm 6.9$ | 17.3 |
| J0816.4-1311 | PMN J0816-1311 | bll | 0.0460 | -4.67 | -2.16 | $-3.23 \pm 1.49$ | 0.89 |
| J0816.4+5739 | SBS 0812+578 | bll | 0.2940 | -0.55 | -0.74 | $-2.32 \pm 3.14$ | 4.16 |
| J0816.9+2050 | SDSS J081649.78+205106.4 | bll | 0.0583 | -2.72 | -1.65 | $-0.77 \pm 0.47$ | 0.36 |
| J0828.3+4153 | GB6 B0824+4203 | bll | 0.2262 | +0.84 | +0.92 | $+5.80 \pm 6.33$ | 18.4 |
| J0831.8+0429 | PKS 0829+046 | bll | 0.1738 | +2.00 | +1.41 | $+3.62 \pm 2.56$ | 8.72 |
| J0847.2+1134 | RX J0847.1+1133 | bll | 0.1982 | +0.09 | +0.29 | $+0.72 \pm 2.45$ | 5.65 |
| J0850.6+3454 | RX J0850.5+3455 | bll | 0.1450 | +1.43 | +1.20 | $+2.33 \pm 1.95$ | 6.20 |
| J0908.9+2311 | RX J0908.9+2311 | bll | 0.2230 | -2.15 | -1.47 | $-3.98 \pm 2.71$ | 2.32 |
| J0912.4+1555 | SDSS J091230.61+155528.0 | bll | 0.2120 | -0.47 | -0.69 | $-1.73 \pm 2.52$ | 3.54 |
| J0930.4+4952 | 1ES 0927+500 | bll | 0.1867 | -4.58 | -2.14 | $-20.4 \pm 9.53$ | 5.60 |
| J1015.0+4926 | 1H 1013+498 | bll | 0.2120 | +0.03 | +0.19 | $+2.14 \pm 11.5$ | 25.1 |
| J1027.0-1749 | 1RXS J102658.5-174905 | bll | 0.2670 | +0.10 | +0.32 | $+14.6 \pm 45.5$ | 105 |
| J1041.7+3900 | B3 1038+392 | bll | 0.2084 | -0.51 | -0.71 | $-3.15 \pm 4.43$ | 6.01 |
| J1053.6+4930 | GB6 J1053+4930 | bll | 0.1404 | +0.31 | +0.56 | $+2.97 \pm 5.33$ | 13.8 |
| J1058.6+5628 | TXS 1055+567 | BLL | 0.1433 | -1.62 | -1.27 | $-15.5 \pm 12.2$ | 11.7 |
| J1100.3+4020 | RX J1100.3+4019 | bll | 0.2250 | +2.67 | +1.63 | $+9.08 \pm 5.56$ | 20.3 |
| J1104.4+3812 | Mkn 421 | BLL | 0.0310 | +4166.97 | +64.55 | $+29.5 \pm 0.5$ | 30.6 |
| J1117.0+2014 | RBS 0958 | bll | 0.1380 | +5.28 | +2.30 | $+3.00 \pm 1.31$ | 5.63 |

Table 3
(Continued)

| 3FHL Source <br> (1) | Counterpart <br> (2) | Class <br> (3) | Redshift <br> (4) | $\begin{aligned} & \text { TS } \\ & (5) \end{aligned}$ | $\begin{gathered} \pm \sqrt{\mathrm{TS}} \\ (6) \end{gathered}$ | $K \pm \Delta K$ <br> (7) | $\begin{gathered} K_{2 \sigma} \\ (8) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| J1120.8+4212 | RBS 0970 | bll | 0.1240 | +1.42 | +1.19 | $+2.73 \pm 2.29$ | 7.29 |
| J1125.9-0743 | 1RXS J112551.6-074219 | bll | 0.2790 | -2.74 | -1.65 | $-23.2 \pm 14.0$ | 10.8 |
| J1136.8+2549 | RX J1136.8+2551 | bll | 0.1560 | -1.13 | -1.06 | $-1.74 \pm 1.64$ | 1.75 |
| J1140.5+1528 | NVSS J114023+152808 | bll | 0.2443 | -0.59 | -0.77 | $-2.41 \pm 3.13$ | 4.00 |
| J1142.0+1546 | MG1 J114208+1547 | bll | 0.2990 | -0.93 | -0.96 | $-4.08 \pm 4.24$ | 4.94 |
| J1145.0+1935 | 3C 264 | rdg | 0.0216 | +3.61 | +1.90 | $+0.44 \pm 0.24$ | 0.92 |
| J1150.3+2418 | OM 280 | bll | 0.2000 | +7.15 | +2.67 | $+6.24 \pm 2.33$ | 10.9 |
| J1154.1-0010 | 1RXS J115404.9-001008 | bll | 0.2535 | -0.28 | -0.53 | $-3.27 \pm 6.19$ | 9.46 |
| J1204.2-0709 | 1RXS J120417.0-070959 | bll | 0.1850 | -2.30 | -1.52 | $-9.58 \pm 6.31$ | 5.09 |
| J1217.9+3006 | 1ES 1215+303 | bll | 0.1300 | +11.36 | +3.37 | $+4.64 \pm 1.38$ | 7.39 |
| J1219.7-0312 | 1RXS J121946.0-031419 | bll | 0.2988 | -0.94 | -0.97 | $-10.0 \pm 10.3$ | 12.4 |
| J1221.3+3010 | PG 1218+304 | bll | 0.1837 | +5.02 | +2.24 | $+5.23 \pm 2.34$ | 9.92 |
| J1221.5+2813 | W Comae | bll | 0.1029 | +6.03 | +2.45 | $+2.33 \pm 0.95$ | 4.24 |
| J1224.4+2436 | MS 1221.8+2452 | bll | 0.2187 | $+0.55$ | +0.74 | $+1.99 \pm 2.68$ | 7.32 |
| J1229.2+0201 | 3C 273 | FSRQ | 0.1583 | -1.97 | -1.40 | $-3.48 \pm 2.48$ | 2.15 |
| J1230.2+2517 | ON 246 | bll | 0.1350 | +0.43 | +0.66 | $+0.86 \pm 1.31$ | 3.46 |
| J1230.8+1223 | M87 | rdg | 0.0042 | +12.93 | +3.60 | $+0.56 \pm 0.16$ | 0.88 |
| J1231.4+1422 | GB6 J1231+1421 | bll | 0.2559 | -0.09 | -0.30 | $-1.01 \pm 3.40$ | 5.82 |
| J1231.7+2847 | B2 1229+29 | bll | 0.2360 | $+0.00$ | +0.05 | $+0.18 \pm 3.30$ | 6.75 |
| J1253.7+0328 | MG1 J125348+0326 | bll | 0.0657 | -3.80 | -1.95 | $-1.36 \pm 0.70$ | 0.46 |
| J1256.2-1146 | PMN J1256-1146 | bcu | 0.0579 | +0.26 | +0.51 | $+0.84 \pm 1.63$ | 4.11 |
| J1310.3-1158 | TXS 1307-117 | bll | 0.1400 | -3.50 | -1.87 | $-11.4 \pm 6.10$ | 4.11 |
| J1341.2+3959 | RBS 1302 | bll | 0.1715 | -0.97 | -0.99 | $-3.35 \pm 3.40$ | 3.92 |
| J1402.6+1559 | MC 1400+162 | bll | 0.2440 | +1.78 | +1.33 | $+4.14 \pm 3.10$ | 10.3 |
| J1411.8+5249 | SBS 1410+530 | bcu | 0.0765 | +0.24 | +0.49 | $+1.40 \pm 2.84$ | 7.12 |
| J1418.0+2543 | 1E 1415.6+2557 | bll | 0.2363 | +0.27 | +0.52 | $+1.60 \pm 3.07$ | 7.78 |
| J1419.4+0444 | SDSS J141927.49+044513.7 | bll | 0.1430 | +0.16 | +0.40 | $+0.74 \pm 1.85$ | 4.50 |
| J1419.7+5423 | OQ 530 | bll | 0.1525 | $+0.30$ | +0.55 | $+5.94 \pm 10.8$ | 27.6 |
| J1428.5+4240 | H $1426+428$ | bll | 0.1292 | +1.58 | +1.26 | $+3.18 \pm 2.54$ | 8.18 |
| J1436.9+5639 | RBS 1409 | bll | 0.1500 | +2.38 | +1.54 | $+21.0 \pm 13.6$ | 48.1 |
| J1442.8+1200 | 1ES 1440+122 | bll | 0.1631 | +3.53 | +1.88 | $+3.37 \pm 1.80$ | 6.95 |
| J1449.5+2745 | B2.2 1447+27 | bll | 0.2272 | +0.24 | +0.49 | $+1.49 \pm 3.03$ | 7.46 |
| J1500.9+2238 | MS 1458.8+2249 | bll | 0.2350 | +3.06 | +1.75 | $+5.10 \pm 2.91$ | 11.0 |
| J1508.7+2708 | RBS 1467 | bll | 0.2700 | +0.15 | +0.38 | $+1.48 \pm 3.89$ | 9.29 |
| J1512.2+0203 | PKS 1509+022 | fsrq | 0.2195 | -1.84 | -1.36 | $-5.73 \pm 4.23$ | 3.81 |
| J1518.5+4044 | GB6 J1518+4045 | bll | 0.0652 | +0.91 | +0.96 | $+0.83 \pm 0.87$ | 2.57 |
| J1531.9+3016 | RX J1531.9+3016 | bll | 0.0653 | +0.03 | +0.18 | $+0.10 \pm 0.57$ | 1.25 |
| J1543.6+0452 | CGCG 050-083 | bcu | 0.0400 | -5.82 | -2.41 | $-1.04 \pm 0.43$ | 0.24 |
| J1554.2+2010 | 1ES 1552+203 | bll | 0.2223 | +1.76 | +1.33 | $+3.50 \pm 2.64$ | 8.81 |
| J1603.8+1103 | MG1 J160340 +1106 | bll | 0.1430 | -0.17 | -0.41 | $-0.61 \pm 1.50$ | 2.41 |
| J1615.4+4711 | TXS 1614+473 | fsrq | 0.1987 | -1.08 | -1.04 | $-8.29 \pm 7.99$ | 8.82 |
| J1643.5-0646 | NVSS J164328-064619 | bcu | 0.0820 | -0.00 | -0.01 | $-0.02 \pm 1.61$ | 3.23 |
| J1647.6+4950 | SBS 1646+499 | bll | 0.0475 | +0.10 | +0.32 | $+0.38 \pm 1.21$ | 2.82 |
| J1653.8+3945 | Mkn 501 | BLL | 0.0330 | +276.97 | +16.64 | $+7.74 \pm 0.49$ | 8.72 |
| J1719.2+1745 | PKS 1717+17 | bll | 0.1370 | -0.32 | -0.56 | $-0.72 \pm 1.28$ | 1.90 |
| J1725.4+5851 | 7C 1724+5854 | bll | 0.2970 | -1.94 | -1.39 | $-99.0 \pm 71.0$ | 63.5 |
| J1728.3+5013 | I Zw 187 | bll | 0.0550 | -0.05 | -0.22 | $-0.32 \pm 1.44$ | 2.57 |
| J1730.8+3715 | GB6 J1730+3714 | bll | 0.2040 | -0.21 | -0.45 | $-1.71 \pm 3.78$ | 5.86 |
| J1744.0+1935 | S3 1741+19 | bll | 0.0830 | +1.62 | +1.27 | $+0.83 \pm 0.66$ | 2.17 |
| J1745.6+3950 | B2 1743+39C | bll | 0.2670 | +0.04 | +0.21 | $+1.47 \pm 7.11$ | 15.5 |
| J1813.5+3144 | B2 1811+31 | bll | 0.1170 | -1.68 | -1.30 | $-1.59 \pm 1.23$ | 1.14 |
| J1917.7-1921 | 1H 1914-194 | bll | 0.1370 | +0.01 | +0.12 | $+1.64 \pm 13.6$ | 28.9 |
| J2000.4-1327 | NVSS J200042-132532 | fsrq | 0.2220 | -0.76 | -0.87 | $-15.2 \pm 17.4$ | 21.1 |
| J2014.4-0047 | PMN J2014-0047 | bll | 0.2310 | +1.23 | +1.11 | $+6.12 \pm 5.51$ | 17.1 |
| J2039.4+5219 | 1ES 2037+521 | bll | 0.0540 | -0.94 | -0.97 | $-1.70 \pm 1.75$ | 2.04 |
| J2042.0+2428 | MG2 J204208+2426 | bll | 0.1040 | $+0.58$ | +0.76 | $+0.67 \pm 0.89$ | 2.44 |
| J2055.0+0014 | RGB J2054+002 | bll | 0.1508 | -0.26 | -0.51 | $-1.27 \pm 2.51$ | 3.78 |
| J2108.8-0251 | TXS 2106-030 | bll | 0.1490 | -0.55 | -0.74 | $-2.23 \pm 3.01$ | 4.03 |
| J2143.5+1742 | OX 169 | fsrq | 0.2110 | +0.18 | +0.42 | $+1.03 \pm 2.43$ | 5.82 |
| J2145.8+0718 | MS 2143.4+0704 | bll | 0.2350 | -1.02 | -1.01 | $-3.65 \pm 3.61$ | 4.14 |
| J2150.2-1412 | TXS 2147-144 | bll | 0.2290 | +0.01 | +0.10 | $+2.02 \pm 20.2$ | 42.7 |
| J2202.7+4216 | BL Lacertae | BLL | 0.0690 | -0.25 | -0.50 | $-0.50 \pm 1.0$ | 1.53 |
| J2250.0+3825 | B3 2247+381 | bll | 0.1190 | -0.27 | -0.52 | $-0.86 \pm 1.67$ | 2.54 |

Table 3
(Continued)

| 3FHL Source | Counterpart | Class | Redshift | TS | (4) <br> $(1)$ | $(2)$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |

Note. Column (3) uses the identification (uppercase) or association (lowercase) class as given in the 3FHL (described in Section 4.1). Column (4) shows the redshift entry in the 3FHL, used in our analysis. Columns (5) and (6) show the test statistics (TS) and significances ( $\pm \sqrt{\text { TS }}$ ) estimated, allowing fluxes to be positive or negative. Column (7) shows that $K$ is the power-law normalization at 1 TeV in units of $10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, and column (8), $K_{2 \sigma}$, its corresponding $2 \sigma$ upper limit in the same units, computed following Feldman \& Cousins (1998).
$11 \%$ response at $40^{\circ}$ from zenith. The dependence of the upperlimit normalizations with redshift gives $z_{h}=0.089$, matching that of an observed power-law spectrum integrated from $E_{0}=0.63 \mathrm{TeV}$, close to the initial assumption $E_{0}=0.5 \mathrm{TeV}$ (Section 3).

### 5.1.3. Comparison with LAT Extrapolations

In Figure 5 we compare photon fluxes, $N(>0.5 \mathrm{TeV})$, computed from the power-law fit (Equation (5)) with the extrapolations of LAT spectra. HAWC photon fluxes relate to the normalizations $K$ through the integration of the differential spectra (Equations (4) and (5)),

$$
\begin{equation*}
N_{\mathrm{obs}}(>0.5 \mathrm{TeV})=\frac{4 \sqrt{2}}{3} K \cdot 1 \mathrm{TeV} \cdot e^{-z / z_{h}} \tag{7}
\end{equation*}
$$

The green $(\mathrm{TeV}=\mathrm{P})$ dots shown in Figure 5 are for the objects with $\mathrm{TS}>9$, while the rest are shown with the respective HAWC upper limit, using $K_{2 \sigma}$ in Equation (7). Thirteen objects have HAWC limits below the LAT extrapolations, although compatible within the uncertainties. These include IC 310, I Zw 187 (1ES 1727+502), B3 2247+381, and 1ES $2344+514$, identified in Figure 2.

### 5.1.4. Source Classes

Table 4 presents average significances for distinct groups of sources. While GeV emission is known to occur in these sources, $80 \%$ of the objects in our sample have not been detected in the VHE range. Furthermore, HAWC is testing long-term average emission, in contrast to the relatively brief IACTs observations. The group statistics, reinforced by the HAWC measurements as a function of redshift for different source classes (Figures 6 and 7), are briefly summarized as follows:

1. Except for the one starburst, radio galaxies constitute the nearest class of sources in our sample (Figure 6). All six are nearer than $z=0.03$, and happen to be at favorable declinations, from $\delta=+4^{\circ}$ (NGC 1278) to $\delta=+41^{\circ}$ (IC 310 and NGC 1275). Four of them are known VHE sources, including 3C 264 , reported after the publication of the 3FHL, which shows a $+1.9 \sigma$ excess in the HAWC
data. As shown in Figure 6, their bounds are at the $N(>0.5 \mathrm{TeV}) \sim 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ level. The mean significance for this group has a $p$-value of $0.4 \%$.
2. The blazars candidates of uncertain type studied here are at redshifts intermediate between those of radio galaxies and FSRQs. None of the bcus studied here has been reported in the VHE regime. The HAWC upper limits shown in Figure 6 are not particularly constraining, owing to the unfavorable declinations of these sources, seven out of eight outside $+0^{\circ} \leqslant \delta \leqslant+50^{\circ}$.
3. Flat-spectrum radio quasars constitute the most distant class of source in our sample. Five out of six have not been detected as VHE sources, and do not comply with the LAT requirement for the $\mathrm{TeV}=\mathrm{C}$ flag. They appear mostly as underfluctuations in the HAWC data. Figure 6 shows the upper limits on photon fluxes, most of them below the 30 mCrab level.
4. BL Lacertae objects constitute the clear majority of our sample, and of VHE sources. Still, 89 of the 117 BL Lac objects studied here have not been reported in the VHE regime. Figure 7 shows the HAWC fluxes for Mkn 421, Mkn 501, VER J0521+211, and 1ES 1215+303, at about (2-3) $\times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ for the two weaker cases.
The $\mathrm{TeV}=\mathrm{P}$ flagged sources appear with a $p$-value at the $10^{-5}$ level, excluding Mkn 421 and Mkn 501, as expected for a subthreshold persistent TeV emission in these known sub- TeV emitters. The $\mathrm{TeV}=\mathrm{C}$ candidates and $\mathrm{TeV}=\mathrm{N}$ groups do not provide any collective hint of emission. The $p$-values quoted do not account for the number of trials used in testing different groups; they are quoted as indicative of the potential presence of persistent TeV emission at levels $\lesssim 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$.

### 5.2. Spectral Fits for High-significance Sources

We computed optimized spectra for the TS $>9$ sources, fitting together normalizations and spectral indices. The fits are summarized in Figure 8, together with the systematic uncertainties. These have been quantified as $15 \%$ in $K$, the 1 TeV normalization, and $5 \%$ in $\alpha$, the spectral index. We discuss here the two high-significance sources, Mkn 421 and Mkn 501.



Figure 3. Histogram of significances for the sample whole (left) and for the clean sample, excluding Mkn 421 and Mkn 501 (right). The dotted lines correspond to normal distributions of mean zero and standard deviation one for the given number of objects. The point with $\sqrt{\mathrm{TS}}=65$ is Mkn 421 , while that at 16 is Mkn 501 . The three bins at $s>3$ on the right-hand-side histogram contain VER J0521+211, 1ES 1215+303, and M87, each in one bin.

Table 4
Average Significances for Different Groups of Sources

| Source | Number <br> of Objects | Significances |  | $p$ Mean |
| :--- | :---: | :---: | :---: | :---: |
| Group |  | Std. dev |  |  |
|  | $(2)$ | $\mu(s)$ | $\sigma_{s}$ | $P[\mu(s)>x]$ |
| $(1)$ | 136 | +0.769 | 5.762 | $5.58 \times 10^{-19}$ |
| All | 134 | +0.175 | 1.209 | 0.022 |
| All-Mkn | 1 | +1.21 | $\ldots$ | 0.113 |
| Starburst | 6 | +1.082 | 1.403 | $4.03 \times 10^{-3}$ |
| Radio galaxies | 113 | +0.220 | 1.179 | $9.75 \times 10^{-3}$ |
| BL Lacs-Mkn | 6 | -0.872 | 0.609 | 0.984 |
| FSRQ + fsrq | 8 | -0.487 | 1.028 | 0.916 |
| bcu | 30 | +0.768 | 1.228 | $1.29 \times 10^{-5}$ |
| $\mathrm{TeV}=\mathrm{P}$ (clean) | 32 | +0.124 | 1.152 | 0.241 |
| $\mathrm{TeV}=\mathrm{C}$ (clean) | 72 | -0.050 | 1.142 | 0.665 |
| $\mathrm{TeV}=\mathrm{N}$ (clean) |  |  |  |  |

### 5.2.1. Markarian 421

Markarian 421 is the brightest persistent extragalactic object in the TeV sky. For the default $\alpha=2.5$ search, we computed $\mathrm{TS}=4167$ for $K=\left(29.5 \pm 0.5_{\text {stat }} \pm 4.4_{\text {syst }}\right) \times$ $10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$. The optimized power-law fit for the intrinsic spectrum resulted in

$$
\begin{align*}
& \frac{d N}{d E}=\left(33.0 \pm 0.6_{\text {stat }} \pm 4.9_{\text {syst }}\right) \times 10^{-12} \\
& \quad \times\left(\frac{E}{1 \mathrm{TeV}}\right)^{-2.63 \pm 0.02_{\text {stat }} \pm 0.13_{\text {syst }}} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1} \tag{8}
\end{align*}
$$

with $\mathrm{TS}=4193$, i.e., an increase in the test statistic $\Delta \mathrm{TS}=26$ relative to $\alpha$ fixed at 2.5. The integration to the observed photon flux, using $z_{h}=0.116$ for $\alpha=2.63$ in Equation (4), gives $N_{\text {obs }}=(48 \pm 8) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, somewhat larger than the $\alpha=2.5$ estimate shown in Figure 5. The intrinsic energy flux, $f_{E}=(132 \pm 22) \times 10^{-12} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1}$, translates into a luminosity per solid angle of $f_{E} d_{L}(z)^{2}=L(>0.5 \mathrm{TeV}) / \Delta \Omega=(2.1 \pm$ $0.4) \times 10^{43} \mathrm{erg} \mathrm{s}^{-1} \mathrm{sr}^{-1}$, about one-third ( $32 \%$ ) of the $(10 \mathrm{GeV}-$ 1 TeV ) luminosity per steradian inferred from the Fermi-LAT


Figure 4. Fit of $K_{2 \sigma}$, the upper-limit normalizations at 1 TeV , as a function of declination and redshift for the "clean" sample. The curved dashed lines are the fits given by Equation (6) to the upper limits at the redshifts indicated in the plot. Blue arrows are upper limits for AGNs with $z \leqslant 0.1$, green for $0.1<z \leqslant 0.2$, and red for $0.2<z \leqslant 0.3$. Measured values of $K$ for the TS $>9$ sources are indicated by stars. They are at declinations $12^{\circ}$ (M87, in blue), $21^{\circ}$ (VER J0521+211, in green); $30^{\circ}$ (1ES 1215+303, in green); and close to $40^{\circ}$ (Mkn 421 as the blue star above the red curve and Mkn 501 as the blue star intersecting the green curve). Statistical errors bars are shown, thought smaller than the markers for the Markarian sources.
measurements (Figure 1). The HAWC spectrum is consistent with IACT observations made by MAGIC between 2007 and 2009 (Ahnen et al. 2016). Using an energy threshold of 400 GeV , these authors found photon fluxes varying from $N_{\text {min }}(>0.5 \mathrm{TeV})=9.33 \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ to $N_{\max }(>0.5 \mathrm{TeV})=$ $2.22 \times 10^{-10} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, weakly dependent on their assumed differential index of 2.5 .

The 3FHL catalog and HAWC spectra for Mkn 421 are shown together in Figure 9. There is a fairly good match between the two fits, with the curves intersecting at about 0.66 TeV . The local LAT spectral index at $1 \mathrm{TeV}, 2.5 \pm 0.2$, is consistent within uncertainties with the HAWC spectral index. IACT observations between 100 GeV and 5 TeV , performed around 2005 with MAGIC, resulted in a spectral index of ( $2.20 \pm 0.08$ ), with indications of a cutoff in the intrinsic


Figure 5. Flux extrapolations of 3 FHL spectra vs. HAWC measurements. Color code indicates the TeV flag: green $\mathrm{TeV}=\mathrm{P}$; orange $\mathrm{TeV}=\mathrm{C}$; red $\mathrm{TeV}=\mathrm{N}$. The green points are the measured fluxes of the five objects with TS $>9$, the rest of the sample indicated by upper limits. Errors shown are statistical only.


Figure 6. HAWC photon flux upper limits for radio galaxies (RDG), blazars candidates of uncertain type (bcus), and flat-spectrum radio quasars (FSRQ), naturally segregated in redshift integrals. The green point at the lowest redshift corresponds to M87. Only three RDGs and one FSRQ have a $\mathrm{TeV}=\mathrm{P}$ flag. 3C 264 was reported as a relatively faint VHE source after the release of the 3FHL catalog. The red dotted line is the 30 mCrab reference, while the vertical dotted lines separate the types of sources.
spectrum (Albert et al. 2007a). A detailed spectral analysis of Mkn 421 and Mkn 501 with HAWC data will be presented in a separate publication. Preliminary results can be found in Coutiño de Leon et al. (2019).

### 5.2.2. Markarian 501

The TeV emission from Markarian 501 observed by HAWC is not as steady as that of Mkn 421 (Abeysekara et al. 2017a). In fact, the statistical significance of the time-averaged TeV emission of Mkn 501 has decreased with increased HAWC exposure. The $\alpha=2.5$ search resulted in a test statistic $\mathrm{TS}=276.97$ for $K=\left(7.74 \pm 0.49_{\text {stat }} \pm 1.16_{\text {syst }}\right) \times 10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$. The


Figure 7. HAWC photon flux upper limits for BL Lac objects and the starburst galaxy NGC 1068. Only BL Lac objects with LAT extrapolation above $2 \times 10^{-13} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ are shown, as in Figure 2. The redshift interval is restricted to $z<0.23$, as in Figure 6. No BL Lac is extrapolated above that flux beyond that distance. Four flux points are shown: Mkn 421 and Mkn 501 above $10^{-11} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, and VER J0521+211 and 1ES $1215+303$ slightly below and above the red line at redshifts 0.11 and 0.13 , respectively.


Figure 8. Optimized spectral fits for five selected sources. Mkn 421 and Mkn 501, the two highest points, present the best statistics. VER J5021+211 and 1ES $1215+303$ have the extreme spectral indices, although with a larger uncertainty. M87 appears with the smallest normalization and a spectral index consistent with $\alpha=2.5$. The black and gray oval in the upper left represents systematic uncertainties of $15 \%$ in the 1 TeV normalization $(K)$ and $5 \%$ on spectral index $(\alpha)$.
optimized power-law fit is

$$
\begin{align*}
& \frac{d N}{d E}=\left(6.21 \pm 0.69_{\text {stat }} \pm 0.93_{\text {syst }}\right) \times 10^{-12} \\
& \quad \times\left(\frac{E}{1 \mathrm{TeV}}\right)^{-2.31 \pm 0.08_{\text {stat }} \pm 0.12_{\text {syst }}} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1} \tag{9}
\end{align*}
$$

with $\mathrm{TS}=280.28$, representing a moderate increase in the test statistic $\Delta \mathrm{TS}=3.33$. The integrated observed photon flux is $N_{\text {obs }}=(8.5 \pm 2.6) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, accounting for EBL attenuation. The intrinsic energy flux, $f_{E}=(40 \pm 16) \times$ $10^{-12} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1}$, translates into a luminosity per solid angle of $d_{L}(z)^{2} f_{E}=L(>0.5 \mathrm{TeV}) / \Delta \Omega=(7.3 \pm 2.8) \times 10^{42} \mathrm{erg} \mathrm{s}^{-1} \mathrm{sr}^{-1}$, which is about $25 \%$ of the $(10 \mathrm{GeV}-1 \mathrm{TeV})$ luminosity per


Figure 9. Spectra of Mkn 421 measured by LAT, including EBL attenuation (black and gray, on the left) and by HAWC (blue, on the right). The full-dotted line represents the intrinsic spectrum, with statistical errors, and the full line represents the fit with EBL attenuation. The two attenuated fits match well at their (0.5-2.0) TeV intersection.
steradian measured by Fermi-LAT (Figure 1). This fraction is similar to that observed for Mkn 421.

The 3FHL catalog and HAWC spectra for Mkn 501 are shown together in Figure 10. The agreement is not as good as for Mkn 421: the Mkn 501 spectrum is harder and lies below the measurement by the LAT. The local LAT spectral index at 1 TeV is $2.58 \pm 0.35$, just consistent with HAWC when accounting for the propagation of the uncertainty in the curvature parameter $\beta$ of the 3FHL fit. We note that Mkn 501 has a variability index $V_{\text {bayes }}=4$ in the 3FHL catalog.

The literature on the VHE characteristics of Mkn 501 is extensive. Its high TeV variability was noted shortly after its 1996 discovery, when the HEGRA group reported flux variations of an order of magnitude observed in mid1997 (Quinn et al. 1996; Aharonian et al. 1997). Contemporaneous 10 m Whipple data confirmed the high state, adding indications of a curved spectrum favored over a simple power law (Samuelson et al. 1998). Further monitoring showed strong variations in flux, although with stable spectra best-fitted by a power law plus an intrinsic cutoff at around 6 TeV (Aharonian et al. 1999). Observations in 1998-1999 showed lower activity, with evidence for spectral curvature and steepening (Aharonian et al. 2001). Mkn 501 observations renewed in 2005, when the MAGIC telescope measured strong and very fast variability, with spectral indices ranging from $\sim 2.0$ to 2.7 (Albert et al. 2007b). Data taken the following year (2006) by MAGIC permitted the characterization of a low-activity state, with fluxes similar to those reported by VERITAS and MAGIC from data taken in the 2009 joint observations with the FermiLAT (Anderhub et al. 2009; Acciari et al. 2011). The contemporaneous LAT data showed spectral variability also present in the GeV range, with index variations $\Delta \alpha \sim 1$ during the first 480 days of Fermi observations (Abdo et al. 2011).
In the last decade, EAS arrays have been able to reach the sensitivities needed to perform long-term monitoring of Mkn 501: the ARGO collaboration reported variations of a factor of six in flux observed between 2011 October and 2012 April (Bartoli et al. 2012). The HAWC collaboration presented light curves for the Crab, Mkn 421, and Mkn 501 from its first 17 months of observations (Abeysekara et al. 2017a). The


Figure 10. Spectrum of Mkn 501 measured by LAT and EBL attenuated (black and gray) and by HAWC (blue). The full-dotted line represents the HAWC intrinsic spectrum, with statistical errors, and the full line represents the fit with EBL attenuation.

HAWC light curve of Mkn 501 showed a low flux baseline with a handful of very short strong flares. The simple powerlaw fit (of index 2.84) was disfavored against a power law with exponential cutoff.

The variability of Mkn 501 has prevented a baseline characterization of this object. In Coutiño de Leon et al. (2019) we presented the HAWC spectrum of Mkn 501 using data acquired between 2015 June and 2017 December. Even though spectra with exponential cuts were tested, the Mkn 501 data proved consistent with a pure power law of index $2.40 \pm 0.06$. While there might be a slight decrease in the spectral index ( $2.31 \pm 0.08 \pm 0.20$ ), the flux reported here is about half that reported in Coutiño de Leon et al. (2019). The four year average TeV flux observed by HAWC is a factor of two above the lowest activity observed so far (Anderhub et al. 2009; Acciari et al. 2011).

### 5.3. Marginally Detected Sources

### 5.3.1. M87

M87 is the central galaxy of the Virgo cluster, a giant elliptical at a distance of just $(16.4 \pm 0.5) \mathrm{Mpc}$, as measured independently of its redshift, $z=0.0042$ (Bird et al. 2010). With an optical magnitude $V=8.6$ and an angular diameter of about $8^{\prime}$, M87 has been imaged in detail for more than a century, showing a distinct bright active nucleus and a single optical jet (Curtis 1918; Tsvetanov et al. 1998). The mass of the central black hole has been measured to be (6.5 $\pm 0.7) \times 10^{9} M_{\odot}$, through its imaging with the Event Horizon Telescope (Event Horizon Telescope Collaboration et al. 2019).

Also known as Virgo A, it is a bright object all throughout the electromagnetic spectrum. As 3FHL J1230.8+1223, it is a $12.1 \sigma$ detection above 10 GeV , as reported in the 3 FHL catalog, with $4.8 \sigma$ in $150-500 \mathrm{GeV}$, and a $2 \sigma$ upper limit in the top ( $0.5-2.0$ ) TeV band. M87 has been observed frequently in the VHE regime since its 2003 discovery by the HEGRA collaboration (Aharonian et al. 2003). The temporal behavior of M87 in the high-energy and VHE bands has been reviewed by Ait Benkhali et al. (2019). Fluxes can vary by a factor of 10 between low, mid, and high states, with indications of the


Figure 11. High-energy spectra of M87 from the LAT catalog and HAWC observations. The data suggest a steepening of the spectrum at energies $\sim 1 \mathrm{TeV}$. The HAWC optimized power-law fit has $\mathrm{TS}=13.2$. The quasidifferential computation for the $(2.0-8.0) \mathrm{TeV}$ interval is at $\mathrm{TS}=13.1$.
spectral index varying from 2.6 in low states to 2.2 in high states. In addition to the low and high states, time variability on single-day timescales during high states has been reported by the H.E.S.S. and MAGIC collaborations (Beilicke et al. 2007; Albert et al. 2008).

The default HAWC search at the M87 location gave a test statistic of $\mathrm{TS}=12.93$ for a normalization $K=(0.56 \pm$ $\left.0.16_{\text {stat }} \pm 0.08_{\text {syst }}\right) \times 10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$. The optimized power-law fit for the intrinsic spectrum is

$$
\begin{align*}
& \frac{d N}{d E}=\left(0.69 \pm 0.22_{\text {stat }} \pm 0.10_{\text {syst }}\right) \times 10^{-12} \\
& \quad \times\left(\frac{E}{1 \mathrm{TeV}}\right)^{-2.63 \pm 0.18_{\text {stat }} \pm 0.13_{\text {syst }}} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}, \tag{10}
\end{align*}
$$

with $\mathrm{TS}=13.19$, i.e., a nonsignificant increase in the test statistic, $\Delta \mathrm{TS}=0.26$, showing $\alpha=2.5$ to be an acceptable solution within the statistics of the optimized fit. The integrated photon flux is $N_{\text {obs }}(>0.5 \mathrm{TeV})=(1.3 \pm 0.9) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, with less than $4 \%$ of attenuation by the EBL. The energy flux, $f_{E}=(2.7 \pm 2.4) \times 10^{-12} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1}$, translates into a luminosity per solid angle of $L(>0.5 \mathrm{TeV}) / \Delta \Omega=(6.9 \pm 6.3) \times$ $10^{39} \mathrm{erg} \mathrm{s}^{-1} \mathrm{sr}^{-1}$, which is about $25 \%$ of that in the $10 \mathrm{GeV}-$ 1 TeV range, from the respective 3 FHL parameters. The lower apparent power with respect to the two nearest BL Lacs is attributed to the off-axis viewing of the jet. Still, the relative power when compared with the LAT regime appears to be similar.

From Ait Benkhali et al. (2019), we computed photon fluxes from M87 observations by IACTs, in different activity states:

1. High state: $N(>0.5 \mathrm{TeV})=\left(5.74_{-1.147}^{+1.14}\right) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$,
2. Mid state: $N(>0.5 \mathrm{TeV})=\left(1.39_{-0.43}^{+0.47}\right) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$,
3. Low state: $N(>0.5 \mathrm{TeV})=\left(2.85_{-1.53}^{+2.27}\right) \times 10^{-13} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$.

The 4.5 yr averaged emission, as indicated by the HAWC data, matches the mid state, with a relatively steep spectral index. The comparison with the 3FHL data, shown in Figure 11, points to a steepening in the spectrum. While the LAT and HAWC data are not contemporaneous, with a variability index $V_{\text {bayes }}=1$, M87 does not stand as a variable in the 3FHL
catalog. Still, an analysis of joint contemporaneous LATHAWC data is desirable.

### 5.3.2. VER J0521 + 211

VER J0521 +211 was discovered as a $\mathrm{TeV} \gamma$-ray source by the VERITAS collaboration, during observations following up its detection above 30 GeV by LAT (Ong 2009). 3FHL $\mathrm{J} 0521.7+2112$ itself is associated with the radio-loud BL Lac object TXS $0518+211$, and the X-ray ROSAT source RX J0521.7+2112. The redshift survey of Shaw et al. (2013) assigned $z=0.108$ to the optical counterpart. This is the value listed in the 3FHL catalog and assumed for this analysis. Archambault et al. (2013) did not confirm this redshift, and neither did Paiano et al. (2017), who only set a lower limit $z>0.18$ using deep spectroscopy with the GTC 10.4 m . The VHE emission from this object has been measured at least up to 1 TeV , with no apparent decline in the spectrum (Prokoph et al. 2015). This makes the high-energy characterization of VER J0521+211 of particular interest.

The $\alpha=2.5$ search resulted in $\mathrm{TS}=9.49$ for $K=$ $\left(2.85 \pm 0.93_{\text {stat }} \pm 0.43_{\text {syst }}\right) \times 10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$. The best power-law fit for the intrinsic spectrum is

$$
\begin{align*}
& \frac{d N}{d E}=\left(2.39 \pm 0.89_{\text {stat }} \pm 0.36_{\text {syst }}\right) \times 10^{-12} \\
& \quad \times\left(\frac{E}{1 \mathrm{TeV}}\right)^{-2.01 \pm 0.38_{\text {stat }} \pm 0.10_{\text {syst }}} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1} \tag{11}
\end{align*}
$$

with $\mathrm{TS}=10.34$, representing a modest increase in the test statistic $\Delta \mathrm{TS}=0.85$ with respect to the default search. This is the hardest AGN spectrum for the $\mathrm{TS}>9$ HAWC sample. The integrated observed photon flux is $N_{\text {obs }}(>0.5 \mathrm{TeV})=$ $(1.5 \pm 1.1) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, attenuated by a factor of $2 / 3$ due to the EBL. The energy flux and luminosity per solid angle cannot be accurately determined with the spectral index so close to 2.0 . If we take $E^{2} d N / d E$ at 1 TeV as indicative, we get an estimated $f_{E} \sim(3.8 \pm 1.4) \times 10^{-12} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1}$ and $d_{I}^{2} f_{E} \sim(8.4 \pm 3.1) \times 10^{42} \mathrm{erg} \mathrm{s}^{-1} \mathrm{sr}^{-1}$. Notwithstanding the uncertain distance, this source appears particularly luminous in the GeV regime, with $L(10 \mathrm{GeV}-1 \mathrm{TeV}) / \Delta \Omega \sim 1.4 \times$ $10^{44} \mathrm{erg} \mathrm{s}^{-1} \mathrm{sr}^{-1}$, standing above the two Markarians in Figure 1.

The 3FHL data, shown together with the HAWC spectral fit in Figure 12, have strong detections up to the $150-500 \mathrm{GeV}$ band, with a highest-energy photon of 370 GeV . The 3FHL fit is a log-parabola, transiting from a hard spectrum at about 30 GeV to a very steep local spectral index of $3.7 \pm 0.7$ at 1 TeV ( 4.3 when attenuating the 3 FHL spectrum). The attenuated HAWC spectrum corresponds to an observed spectral index of $\sim 2.7$, mostly inconsistent with the LAT fit. The data may be reconciled through an intrinsic hardening at about 200 or 300 GeV . The HAWC quasi-differential bound in the $0.5-2.0 \mathrm{TeV}$ band is a factor of 4.6 lower than the 3 FHL limit in the same band. We note that the quasi-differential analysis gave $\mathrm{TS}=9.1$ for the $2.0-8.0 \mathrm{TeV}$ band, optimal in terms of the HAWC instrumental response. With the assumed redshift, $\tau$ is due to range from 1.6 to 3.2 in that energy interval, pointing to an intrinsically hard spectrum. The analysis gave $\mathrm{TS}=0$ for $8.0-32.0 \mathrm{TeV}$, expected to be heavily attenuated for the assumed redshift. Additional HAWC data analysis should allow the further constraint of the shape of the TeV spectrum VER J0521+211, testing the redshift assumption.


Figure 12. Spectral fits of VER J0521.7+2112 from Fermi-LAT (gray) and HAWC (blue) data. The best fit of the HAWC data is for an intrinsic power law of index 2.0, which remains harder than the 3FHL log-parabola fit after considering EBL attenuation at the assumed redshift. The HAWC quasidifferential limit at $(2.0-8.0) \mathrm{TeV}$ is for a $\mathrm{TS}=9.1$.

We also note that VER J0521 +211 has a variability index $V_{\text {bayes }}=4$ in the 3FHL catalog, indicating that a joint analysis of contemporaneous LAT and HAWC data would be relevant.

As mentioned in Section 4.1, VER J0521+211 is located 3.07 from the Crab Nebula, the brightest source in the 2HWC and 3HWC catalogs. This angular distance corresponds to three times the $68 \%$ containment angle $\psi_{68}$ for $\mathcal{B}=1$, and $>6 \psi_{68}$ for $\mathcal{B}>2$. We tested for potential contamination, repeating the maximum-likelihood test with $\alpha=2.5$ at five locations equidistant from the Crab Nebula, forming together with VER J0521+211 a hexagon around the Crab. These provided test statistics TS between -4.46 and +1.54 , in contrast with $\mathrm{TS}=+9.49$ at the location of VER $\mathrm{J} 0521+211$.

### 5.3.3. 1ES $1215+303$

1ES $1215+303$ is one of six $\gamma$-ray emitting BL Lac objects located in the Northern part of the Coma Berenices constellation, five of them known to be VHE sources. The redshift of this HBL is now confirmed to be $z=0.130$, discarding the early measurement $z=0.237$ by Lanzetta et al. (1993). The lower value was confirmed through optical spectroscopy at the GTC and, more recently, through the direct identification of a Ly $\alpha$ emission line (Paiano et al. 2017; Furniss et al. 2019). Cataloged as 3 FHL J1217.9 +3006 , this object is well detected up to 500 GeV , and modeled with a power law of index $\alpha=2.3 \pm 0.1$. It has a variability index $V_{\text {bayes }}=2$.

The default HAWC search gave $\mathrm{TS}=11.36$ for $K=$ $\left(4.64 \pm 1.38_{\text {stat }} \pm 0.70_{\text {syst }}\right) \times 10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ as the intrinsic normalization. The optimized power-law fit for the intrinsic spectrum is rather softer,

$$
\begin{align*}
& \frac{d N}{d E}=\left(3.78 \pm 1.36_{\text {stat }} \pm 0.57_{\text {syst }}\right) \times 10^{-12} \\
& \quad \times\left(\frac{E}{1 \mathrm{TeV}}\right)^{-3.07 \pm 0.37_{\text {stat }} \pm 0.15_{\text {syst }}} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1} \tag{12}
\end{align*}
$$

with $\mathrm{TS}=12.80$, representing a test-statistic increase $\Delta \mathrm{TS}=$ 1.44. The attenuated photon flux is $N_{\text {obs }}(>0.5 \mathrm{TeV})=$ $(2.53 \pm 1.35) \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$. The integrated energy flux results


Figure 13. High-energy spectra of 1ES 1215+303 from the 3FHL catalog and HAWC observations. We note the intersection of both spectra at $\gtrsim 1 \mathrm{TeV}$ and its decay, as it suffers EBL attenuation. This is the most distant source detected by HAWC so far.
in a luminosity per solid angle of $L(>0.5 \mathrm{TeV}) / \Delta \Omega=$ $(3.9 \pm 2.7) \times 10^{43} \mathrm{erg} \mathrm{s}^{-1} \mathrm{sr}^{-1}$, given the luminosity distance of 585 Mpc . This is about $30 \%$ of the corresponding value between 10 GeV and 1 TeV . The 3FHL and HAWC spectra, compatible within the uncertainties, are shown in Figure 13.

1ES $1215+303$ was first detected as a VHE source by MAGIC in 2011 (Lombardi et al. 2011; Aleksić et al. 2012). It is usually observed together with PG 1218+304, located just 0.88 degrees away in the sky, for which we obtained a test statistic TS $=+2.24$. Given that both sources are known VHE emitters, the distance is not large enough to be certain that there is no overlap between both sources. With a redshift $z=0.184$, PG $1218+304$ is more distant and prone to be heavily attenuated above 1 TeV . A dedicated study with improved HAWC analysis tools is now pending.

Long-term monitoring of this source by VERITAS was presented by Aliu et al. (2013), prior to the report of a single short and intense flare seen in 2014 with VERITAS and FermiLAT (Abeysekara et al. 2017d). During this episode the VHE flux of this HBL reached 2.4 times that of the Crab Nebula, with a variability timescale $\lesssim 3.6 \mathrm{hr}$. The MAGIC spectrum of Aleksić et al. (2012), ranging from 70 GeV to 1.8 TeV , had an intrinsic spectral index of 2.96, deattenuated with the model of Domínguez et al. (2011), used here. The MAGIC and HAWC spectral indices are in good agreement. The long-term joint monitoring of this source by LAT and VERITAS, spanning more than 10 years, confirms the spectral index and points to four strong and short flares that occurred during full HAWC operations by Valverde et al. (2020). We checked the dates of these four episodes and did not find evidence for them in the HAWC data. An optimized timedependent analysis of this source with HAWC, beyond the scope of this paper, will be the subject of future work.

### 5.4. Limits on Sources of Interest

The extrapolation of 3 FHL catalog spectra (Figure 2) and of IACT spectral fits allowed us to identify AGN targets for potential HAWC detection (Section 4.3). We presented evidence of persistent TeV emission for three of those targets-Mkn 421, Mkn 501, and at a weaker level, M87-and, also marginally, for two sources below the 30 mCrab reference limit (marked as a red

Table 5
Upper Limits on 3FHL AGN Candidate Targets

| 3FHL Entry | Source | $z$ | $\pm \sqrt{\text { TS }}$ | $K_{2 \sigma}$ | $N_{0.5}^{x}$ |
| :--- | :---: | :---: | :---: | :---: | ---: |
| 3FHL J0214.5+5145 | TXS 0210+515 | 0.049 | +1.62 | 5.43 | 2.40 |
| 3FHL J0316.6+4120 | IC 310 | 0.019 | +0.86 | 4.01 | 4.37 |
| 3FHL J1428.5+4240 | H 1426+428 | 0.129 | +1.26 | 8.18 | 1.21 |
| 3FHL J1728.3+5013 | I Zw 187 | 0.055 | -0.22 | 2.57 | 3.25 |
| 3FHL J2250.0+3825 | B3 2247+381 | 0.119 | -0.27 | 2.54 | 2.37 |
| 3FHL J2347.0+5142 | 1ES 2344+514 | 0.044 | +1.45 | 4.62 | 7.12 |

Note. $K_{2 \sigma}$, the spectrum normalization at 1 TeV , is in units of $10^{-12} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1} . N_{0.5}$ are extrapolated photon fluxes above 0.5 TeV and the corresponding upper limit, in units of $10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$.
line in Figure 2), VER J0521+211 and 1ES $1215+303$. On the other hand, sources like IC 310, 1ES 2344+514, TXS 0210+515, 1ES $1727+502$, B3 $2247+381$, and $\mathrm{H} 1426+428$ were not detected. As shown in Figure 4, the HAWC sensitivity is dependent on source decl. We note that while the sources with $\sqrt{\mathrm{TS}}>9$ are all in the range $+12^{\circ}<\delta<+40^{\circ}$, five of the undetected candidates are either North of $\delta=+50^{\circ}$, or farther than $z=0.1$. Decl. is of particular relevance here, as the response of EAS arrays to low-energy events is compromised at large zenith angles. Still, as shown in Table 5 and in Section 5.1.3, four of the corresponding upper limits are below the extrapolation of the corresponding 3FHL spectrum.

IC 310 is a target that is both nearby and culminates at an adequate zenith angle, but remained undetected in this HAWC analysis. The upper limit set in the photon flux is close to a third of the rather uncertain LAT extrapolation, and the HAWC upper limit in the $0.5-2.0 \mathrm{TeV}$ band is 2.5 lower than the one in the 3 FHL catalog. We note that while the 3 FHL catalog reports a relatively low variability index, $V_{\text {bayes }}=2$, this source is known to display extreme variability in the VHE range, on timescales as low as five minutes, challenging models and severely constraining the emission region to scales smaller than its event horizon (Aleksić et al. 2014).

In addition to the preselected targets, we point here to two additional sources of intrinsic interest: 3C 264 and NGC 1068. 3C 264 is a radio galaxy hosted by the elliptical galaxy NGC 3862, at a distance of about 90 Mpc . A compact radio source powers a relativistic jet imaged in radio and in the optical (Crane et al. 1993; Lara et al. 2004). 3C 264 shows in the LAT data up to a highest-energy photon of 97 GeV . It was listed as a $\mathrm{TeV}=\mathrm{N}$ source in the 3FHL catalog, and considered as such in this analysis. However, it was later detected in the VHE range by the VERITAS Collaboration (Mukherjee 2018). The photon flux measured by VERITAS above 300 GeV indicates that this object should be too faint for HAWC. The HAWC data show a $+1.9 \sigma$ excess, statistically consistent with such a low flux and providing an upper limit $N(>0.5 \mathrm{TeV})<1.42 \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$.

NGC 1068 has become a source of interest due to its coincidence with a hotspot in the IceCube all-sky map (Aartsen et al. 2020). While two other starburst galaxies, NGC 253 and M82, have been detected in the VHE range as faint sources, with fluxes $1 \%$ of the Crab Nebula (Ohm 2016), NGC 1068 remains undetected and an unlikely candidate for HAWC detection. It is a weak $5.3 \sigma$ detection in the 3FHL catalog, with practically no signal above 20 GeV and a rather steep spectral index $\alpha=3.8 \pm 1.0$. The HAWC data have $\mathrm{a}+1.2 \sigma$ excess at the location of NGC 1068, for an upper-limit normalization $K_{2 \sigma}=6.46 \times 10^{-13} \mathrm{TeV}^{-1} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ at 1 TeV . The quasidifferential HAWC limit in the common energy band,
$N(0.5-2.0 \mathrm{TeV}) \leqslant 2.32 \times 10^{-12} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$, is a factor of eight lower than the respective LAT limit.

## 6. Summary

The HAWC Gamma-Ray Observatory has performed an extensive follow-up survey of known $\mathrm{GeV} \gamma$-ray emitting active galaxies at TeV energies. We investigated all AGNs in the 3 FHL catalog with a redshift lower than 0.3 and transiting within $40^{\circ}$ of latitude $19^{\circ} \mathrm{N}$, the HAWC zenith, searching for $\mathrm{TeV} \gamma$-ray emission averaged over a 4.5 yr period. The HAWC data show clear signals from Mkn 421 and Mkn 501, from which we quantified their long-term averaged spectra. In addition, we obtained marginal evidence for TeV emission from the radio galaxy M87 and the BL Lac objects VER J0521 +211 and 1ES $1215+303$. The fluxes estimated for these sources are compatible with values previously reported for mid- or low-activity states. When excluding Mkn 421 and Mkn 501, we find collective evidence for long-term averaged TeV emission from radio galaxies and BL Lac objects with a $p$ value $\sim 1 \%$, and for known VHE emitters ( $\mathrm{TeV}=\mathrm{P}$ ) with $p$ value $\sim 10^{-5}$. No evidence was found for other source classes or for LAT sources not previously reported in the VHE range.
Upper limits were set for the whole sample assuming intrinsic power-law spectra of index 2.5 attenuated by the interaction of $\gamma$ rays with extragalactic background radiation. These limits are dependent on the decl. and redshift of the sources, confirming a redshift attenuation of exponential scale $z_{h} \simeq 0.1$ for HAWC. Bounds on observed photon fluxes in three energy intervals, (0.5-2.0), (2.0-8.0), and (8.0-32.0) TeV, are also given. HAWC measurements were compared with the mostly noncontemporaneous 3 FHL catalog long-term data and with specific IACT observations.

As the exposure of HAWC continues to deepen, the increased sensitivity will allow for deeper searches of extragalactic sources. Long-term variability is an area of opportunity for ground-based EAS $\gamma$-ray observatories not explored in this paper. Analyses of multiyear AGN light curves are under way in order to expand the investigation presented here to the time regime. The continuous and improved operation of HAWC is leading to a better understanding of EAS arrays, and in particular of water Cherenkov detectors. New analysis tools to improve the sub- TeV sensitivity of HAWC have been developed and implemented. These will provide improved energy response and the reduction of noise at low energies, two requirements for a deeper access to the extragalactic sky. These upgrades will allow HAWC to build on the analysis presented here, laying the path to future TeV survey instruments, in both the Northern and Southern Hemispheres.

We acknowledge the support from the US National Science Foundation (NSF); the US Department of Energy Office of High-Energy Physics; the Laboratory Directed Research and Development (LDRD) program of Los Alamos National Laboratory; Consejo Nacional de Ciencia y Tecnología (CONACyT), México, grants 271051, 232656, 260378, 179588, 254964, 258865, 243290, 132197, A1-S-46288, A1-S-22784, Cátedras $873,1563,341,323$, Red HAWC, México; DGAPA-UNAM grants IG100317, IG101320, IN111315, IN111716-3, IN111419, IA102019, IN112218; VIEP-BUAP; PIFI 2012, 2013; PROFOCIE 2014, 2015; the University of Wisconsin Alumni Research Foundation; the Institute of Geophysics, Planetary Physics, and Signatures at Los Alamos National Laboratory; Polish Science Centre grant DEC-2017/ 27/B/ST9/02272; Coordinación de la Investigación Científica de la Universidad Michoacana; the Royal Society-Newton Advanced Fellowship 180385; la Generalitat Valenciana, grant CIDEGENT/2018/034; Chulalongkorn University's CUniverse (CUAASC) grant. Thanks to Scott Delay, Luciano Díaz, and Eduardo Murrieta for technical support.

This work used data from the Fermi Science Support Center and the TeVCat online source catalog (http://tevcat.uchicago. edu). It also made use of the SIMBAD, NED (NASA/IPAC Extragalactic Database), ADS, and SDSS databases.

Facility: The High Altitude Water Cherenkov (HAWC) Gamma-Ray Observatory.

## ORCID iDs

A. Albert © https://orcid.org/0000-0003-0197-5646
K. P. Arunbabu © https:// orcid.org/0000-0002-3032-663X
H. A. Ayala Solares (10 https:// orcid.org/0000-0002-2084-5049
V. Baghmanyan (©) https:// orcid.org/0000-0003-0477-1614
E. Belmont-Moreno (©) https:// orcid.org/0000-0002-

1114-2640
S. Y. BenZvi © https://orcid.org/0000-0001-5537-4710
C. Brisbois © https:// orcid.org/0000-0002-5493-6344
K. S. Caballero-Mora © https:// orcid.org/0000-0002-4042-3855
T. Capistrán © https://orcid.org/0000-0003-2158-2292
A. Carramiñana (1) https://orcid.org/0000-0002-8553-3302
S. Casanova (1) https://orcid.org/0000-0002-6144-9122
U. Cotti © https://orcid.org/0000-0002-7607-9582
J. Cotzomi © https://orcid.org/0000-0002-1132-871X
S. Coutiño de León © https://orcid.org/0000-0002-7747-754X
E. De la Fuente (1) https:// orcid.org/0000-0001-9643-4134
B. L. Dingus © https://orcid.org/0000-0001-8451-7450
M. A. DuVernois (1) https://orcid.org/0000-0002-2987-9691
J. C. Díaz-Vélez © https:// orcid.org/0000-0002-0087-0693
K. Engel © https:// orcid.org/0000-0001-5737-1820
C. Espinoza © https:// orcid.org/0000-0001-7074-1726
H. Fleischhack (10) https://orcid.org/0000-0002-0794-8780
N. Fraija (i) https:// orcid.org/0000-0002-0173-6453
J. A. García-González © https://orcid.org/0000-0002-

4188-5584
F. Garfias © https://orcid.org/0000-0003-1122-4168
M. M. González © https://orcid.org/0000-0002-5209-5641
J. A. Goodman © https://orcid.org/0000-0002-9790-1299
J. P. Harding (1) https:// orcid.org/0000-0001-9844-2648
S. Hernández © https://orcid.org/0000-0002-2565-8365
D. Huang © https://orcid.org/0000-0002-3808-4639
F. Hueyotl-Zahuantitla © https:// orcid.org/0000-0002-5527-7141
A. Iriarte (1) https:// orcid.org/0000-0001-5811-5167
A. Jardin-Blicq © https://orcid.org/0000-0002-6738-9351
V. Joshi © https:// orcid.org/0000-0003-4467-3621
D. Kieda © https:// orcid.org/0000-0003-4785-0101
A. Lara © https://orcid.org/0000-0001-6336-5291
W. H. Lee © https:// orcid.org/0000-0002-2467-5673
H. León Vargas (© https://orcid.org/0000-0001-5516-4975
J. T. Linnemann (©) https:// orcid.org/0000-0003-2696-947X
A. L. Longinotti (1) https://orcid.org/0000-0001-8825-3624
G. Luis-Raya (1) https:// orcid.org/0000-0003-2810-4867
K. Malone © https://orcid.org/0000-0001-8088-400X
O. Martínez © https:// orcid.org/0000-0001-9052-856X
I. Martinez-Castellanos © https://orcid.org/0000-0001-9035-1290
J. Martínez-Castro (©) https:// orcid.org/0000-0002-2824-3544
J. A. Matthews © https:// orcid.org/0000-0002-2610-863X
P. Miranda-Romagnoli (1) https://orcid.org/0000-0002-8390-9011
M. Mostafá © https:// orcid.org/0000-0002-7675-4656
A. Nayerhoda (1) https:// orcid.org/0000-0003-0587-4324
L. Nellen (1) https://orcid.org/0000-0003-1059-8731
M. Newbold (i) https://orcid.org/0000-0001-9428-7572
M. U. Nisa © https:// orcid.org/0000-0002-6859-3944
R. Noriega-Papaqui (©) https:// orcid.org/0000-0001-7099-108X
E. G. Pérez-Pérez © https://orcid.org/0000-0001-5998-4938
C. D. Rho © https://orcid.org/0000-0002-6524-9769
D. Rosa-González (1) https://orcid.org/0000-0003-1327-0838
E. Ruiz-Velasco © https:// orcid.org/0000-0001-6939-7825
F. Salesa Greus © https:// orcid.org/0000-0002-8610-8703
A. Sandoval (1) https:// orcid.org/0000-0001-6079-2722
M. Schneider (1) https://orcid.org/0000-0001-8644-4734
H. Schoorlemmer (©i) https://orcid.org/0000-0002-8999-9249
A. J. Smith © https://orcid.org/0000-0002-1012-0431
R. W. Springer (©) https://orcid.org/0000-0002-1492-0380
K. Tollefson © https://orcid.org/0000-0001-9725-1479
I. Torres (©) https:// orcid.org/0000-0002-1689-3945
L. Villaseñor © https:// orcid.org/0000-0001-6876-2800
A. Zepeda (© https://orcid.org/0000-0001-9976-2387
H. Zhou (1) https://orcid.org/0000-0003-0513-3841
C. de León (1) https://orcid.org/0000-0002-8528-9573

## References

Aartsen, M. G., Ackermann, M., Adams, J., et al. 2020, PhRvL, 124, 051103
Abdo, A. A., Abeysekara, A. U., Allen, B. T., et al. 2014, ApJ, 782, 110
Abdo, A. A., Ackermann, M., Ajello, M., et al. 2011, ApJ, 727, 129
Abeysekara, A. U., Albert, A., Alfaro, R., et al. 2017a, ApJ, 841, 100
Abeysekara, A. U., Albert, A., Alfaro, R., et al. 2017b, ApJ, 843, 39
Abeysekara, A. U., Albert, A., Alfaro, R., et al. 2017c, ApJ, 843, 40 Abeysekara, A. U., Alfaro, R., Alvarez, C., et al. 2016, ApJ, 817, 3 Abeysekara, A. U., Archambault, S., Archer, A., et al. 2015, ApJL, 815, L22 Abeysekara, A. U., Archambault, S., Archer, A., et al. 2017d, ApJ, 836, 205 Acciari, V. A., Ansoldi, S., Antonelli, L. A., et al. 2019, MNRAS, 486, 4233 Acciari, V. A., Ansoldi, S., Antonelli, L. A., et al. 2020, ApJS, 247, 16 Acciari, V. A., Arlen, T., Aune, T., et al. 2011, ApJ, 729, 2
Acero, F., Ackermann, M., Ajello, M., et al. 2015, ApJS, 218, 23 Aharonian, F., Akhperjanian, A., Barrio, J., et al. 2001, ApJ, 546, 898
Aharonian, F., Akhperjanian, A., Barrio, J., et al. 2002, A\&A, 384, L23 Aharonian, F., Akhperjanian, A., Beilicke, M., et al. 2003, A\&A, 403, L1 Aharonian, F., Akhperjanian, A. G., Barrio, J. A., et al. 1997, A\&A, 327, L5 Aharonian, F. A., Akhperjanian, A. G., Barrio, J. A., et al. 1999, A\&A, 349, 11 Ahnen, M. L., Ansoldi, S., Antonelli, L. A., et al. 2015, ApJL, 815, L23 Ahnen, M. L., Ansoldi, S., Antonelli, L. A., et al. 2016, A\&A, 593, A91 Ait Benkhali, F., Chakraborty, N., \& Rieger, F. M. 2019, A\&A, 623, A2 Ajello, M., Atwood, W. B., Baldini, L., et al. 2017, ApJS, 232, 18 Albert, A., Alfaro, R., Alvarez, C., et al. 2020, ApJ, 905, 76

Albert, J., Aliu, E., Anderhub, H., et al. 2007a, ApJ, 663, 125
Albert, J., Aliu, E., Anderhub, H., et al. 2007b, ApJ, 669, 862
Albert, J., Aliu, E., Anderhub, H., et al. 2008, ApJL, 685, L23
Aleksić, J., Alvarez, E. A., Antonelli, L. A., et al. 2012, A\&A, 544, A142
Aleksić, J., Ansoldi, S., Antonelli, L. A., et al. 2014, Sci, 346, 1080
Aliu, E., Archambault, S., Arlen, T., et al. 2013, ApJ, 779, 92
Anderhub, H., Antonelli, L. A., Antoranz, P., et al. 2009, ApJ, 705, 1624
Antonucci, R. 1993, ARA\&A, 31, 473
Archambault, S., Arlen, T., Aune, T., et al. 2013, ApJ, 776, 69
Bartoli, B., Bernardini, P., Bi, X. J., et al. 2011, ApJ, 734, 110
Bartoli, B., Bernardini, P., Bi, X. J., et al. 2012, ApJ, 758, 2
Beilicke, M., Aharonian, F., Benbow, W., et al. 2007, in AIP Conf. Ser. 921, The First GLAST Symposium, ed. S. Ritz, P. Michelson, \& C. A. Meegan (Melville, NY: AIP), 147
Bird, S., Harris, W. E., Blakeslee, J. P., \& Flynn, C. 2010, A\&A, 524, A71
Blandford, R. D., \& Payne, D. G. 1982, MNRAS, 199, 883
Blandford, R. D., \& Znajek, R. L. 1977, MNRAS, 179, 433
Carramiñana, A., Rosa González, D., Coutiño de León, S., \& Longinotti, A. L. 2019, Proc. ICRC (Madison, WI), 36, 644
Coutiño de Leon, S., Alonso, A. C., Rosa-Gonzalez, D., \& Longinotti, A. L. 2019, Proc. ICRC (Madison, WI), 36, 654
Crane, P., Peletier, R., Baxter, D., et al. 1993, ApJL, 402, L37
Curtis, H. D. 1918, PLicO, 13, 9
Dermer, C. D., \& Schlickeiser, R. 1993, ApJ, 416, 458
Domínguez, A., Primack, J. R., Rosario, D. J., et al. 2011, MNRAS, 410, 2556
Event Horizon Telescope Collaboration, Akiyama, K., Alberdi, A., et al. 2019, ApJL, 875, L1
Feldman, G. J., \& Cousins, R. D. 1998, PhRvD, 57, 3873
Franceschini, A., \& Rodighiero, G. 2017, A\&A, 603, A34
Furniss, A., Worseck, G., Fumagalli, M., et al. 2019, AJ, 157, 41
Gilmore, R. C., Somerville, R. S., Primack, J. R., \& Domínguez, A. 2012, MNRAS, 422, 3189
Górski, K. M., Hivon, E., Banday, A. J., et al. 2005, ApJ, 622, 759
Gould, R. J., \& Schréder, G. 1966, PhRvL, 16, 252
Gould, R. J., \& Schréder, G. P. 1967a, PhRv, 155, 1404
Gould, R. J., \& Schréder, G. P. 1967b, PhRv, 155, 1408
Hawley, J. F., Fendt, C., Hardcastle, M., Nokhrina, E., \& Tchekhovskoy, A. 2015, SSRv, 191, 441
Healey, S. E., Romani, R. W., Cotter, G., et al. 2008, ApJS, 175, 97
H.E.S.S. Collaboration, Abdalla, H., Abramowski, A., et al. 2018, A\&A, 612, A1
H.E.S.S. Collaboration, Abdalla, H., Adam, R., et al. 2020, A\&A, 633, A162 Hillas, A. M. 1984, ARA\&A, 22, 425
Hoyle, F., \& Fowler, W. A. 1963, MNRAS, 125, 169
Icecube Collaboration, Aartsen, M. G., Ackermann, M., et al. 2017, A\&A, 607, A115
IceCube Collaboration, Aartsen, M. G., Ackermann, M., et al. 2018, Sci, 361, 147
Kembhavi, A. K., \& Narlikar, J. V. 1999, Quasars and Active Galactic Nuclei: An Introduction (Cambridge: Cambridge Univ. Press)
Kotera, K., \& Olinto, A. V. 2011, ARA\&A, 49, 119
Kotilainen, J. K., Hyvönen, T., \& Falomo, R. 2011, A\&A, 534, L2
Lanzetta, K. M., Turnshek, D. A., \& Sandoval, J. 1993, ApJS, 84, 109
Lara, L., Giovannini, G., Cotton, W. D., Feretti, L., \& Venturi, T. 2004, A\&A, 415, 905
Liao, N.-h. 2018, Galax, 6, 68
Lombardi, S., Lindfors, E., Becerra Gonzalez, J., et al. 2011, arXiv:1110.6786
Madejski, G. G., \& Sikora, M. 2016, ARA\&A, 54, 725
MAGIC Collaboration, Ansoldi, S., Antonelli, L. A., et al. 2018, MNRAS, 480, 879
Mukherjee, R. 2018, ATel, 11436, 1
Nass, P., Bade, N., \& Kollgaard, R. I. 1996, A\&A, 309, 419
Ohm, S. 2016, CRPhy, 17, 585
Ong, R. A. 2009, ATel, 2260, 1
Paiano, S., Landoni, M., Falomo, R., et al. 2017, ApJ, 837, 144
Paiano, S., Landoni, M., Falomo, R., Scarpa, R., \& Treves, A. 2016, MNRAS, 458, 2836
Petry, D., Bond, I. H., Bradbury, S. M., et al. 2002, ApJ, 580, 104
Prokoph, H., Schultz, C., \& Da Vela, P. 2015, Proc. ICRC (The Hague), 34, 864
Punch, M., Akerlof, C. W., Cawley, M. F., et al. 1992, Natur, 358, 477
Quinn, J., Akerlof, C. W., Biller, S., et al. 1996, ApJL, 456, L83
Rieger, F., \& Levinson, A. 2018, Galax, 6, 116
Salamon, M. H., \& Stecker, F. W. 1998, ApJ, 493, 547
Samuelson, F. W., Biller, S. D., Bond, I. H., et al. 1998, ApJL, 501, L17
Shaw, M. S., Romani, R. W., Cotter, G., et al. 2013, ApJ, 764, 135
Tchekhovskoy, A., Narayan, R., \& McKinney, J. C. 2011, MNRAS, 418, L79
Tsvetanov, Z. I., Hartig, G. F., Ford, H. C., et al. 1998, ApJL, 493, L83
Valverde, J., Horan, D., Bernard, D., et al. 2020, ApJ, 891, 170
Wakely, S. P., \& Horan, D. 2008, Proc. ICRC (Mérida, Yucatán, Mexico), 3, 1341


[^0]:    ${ }^{32}$ Equation (2) assumes an isotropic $n_{\nu}$; this may not describe a local radiation field.

[^1]:    $\overline{33}$ Fits file gll_psch_v13.fits, dated 2017 July at https://fermi.gsfc.nasa.gov/ ssc/.

[^2]:    $\overline{34} 3 \mathrm{C} 264$ has a $\mathrm{TeV}=\mathrm{N}$ flag, because its VHE detection occurred after the publication of the 3 FHL catalog.

[^3]:    35 3FHL J1105.8+3944 and 3FHL J1652.7 +4024 excluded.

