

Anomalous normal-state magnetotransport in an electron-doped cuprate

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We report magnetoresistance and Hall angle measurements of the electron-doped cuprate $\text{La}_{2-x}\text{Ce}_x\text{CuO}_4$ over a wide range of dopings from $x = 0.08$ – 0.17 . Above 100 K, we find an unconventional $\sim H^{3/2}$ magnetic field dependence of the magnetoresistance observed in all samples doped within the superconducting dome. Further, the measured magnetoresistance violates Kohler's rule. Given the ubiquity of this anomalous magnetoresistance at high temperatures above the superconducting dome, we speculate that the origin of this behavior is linked to the unusual $\rho \sim T^2$ resistivity observed over the same wide parameter range at high temperatures. We also find a strong doping dependence of the Hall angle with an unconventional temperature dependence of $\cot \theta_H \sim T^4$ ($T^{2.5}$) for samples doped below (above) the Fermi surface reconstruction doping $x_{\text{FSR}} = 0.14$.

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I. INTRODUCTION

The normal state of cuprate high-temperature superconductors has captivated the interest of the condensed matter physics community for the past decade while defying every attempt at theoretical explanation. The transport phenomena observed in these materials is believed to depart from the conventional Landau Fermi liquid theory of metals [1] and such non-Fermi liquid behaviors appear to be a common feature of disparate families of high-temperature superconductors [2]. Consequently, it is reasonable to imagine that the mechanism of high-temperature superconductivity may naturally emerge from this non-Fermi liquid strange metal just as phonon-mediated superconductivity naturally emerges from a conventional Fermi liquid, making an understanding of this strange metallic state a potential stepping stone toward identifying the origin of high-temperature superconductivity.

The typical hallmark of the strange metallic state is the infamous linear-in- T resistivity of hole-doped cuprates, which persists over an anomalously large temperature range, from T_c to 1000 K in some systems [3,4]. However, electron-doped compounds also exhibit a plethora of strange metallic behavior which differ sharply from the conventional properties of a Fermi liquid [5]. Further, these materials display two distinct regimes of strange metallicity with different behaviors at high and low temperatures, and which may or may not be of a common origin.

The high-temperature strange metallic phase of electron-doped cuprates is characterized by a universal quadratic-in- T resistivity, seen in all compounds and for all dopings [6,7] from roughly 100 K to above 600 K [8]. This T^2 behavior is what one might naively expect for a Fermi liquid. However, a more thoughtful consideration of the large magnitude of the resistivity and the high temperature scale at which it is observed lead one to conclude that there is, in fact, an extremely

strange transport behavior [7]. In fact, it is arguably stranger than the high-temperature linear-in- T resistivity of the hole-doped materials, which could potentially be explained by electron-phonon scattering in a low carrier density system [9] and is a generic feature of most conventional metals.

At low temperatures, the nature of the strange metallic ground state is strongly doping dependent and nonuniversal. Generically, the phase diagram of the electron-doped cuprates is dominated by a Fermi surface reconstruction (FSR) which occurs inside the superconducting dome [10–14] and is believed to be driven by short-ranged antiferromagnetic order [15] or the onset of topological order [16]. For the material of interest in this study, $\text{La}_{2-x}\text{Ce}_x\text{CuO}_4$ (LCCO), this FSR occurs at cerium concentration $x = .14$ [17] (for reference, the SC dome extends from $x = .07$ to $x = .175$). In samples doped below the FSR (i.e., $x < .14$ for LCCO), the low-temperature resistivity exhibits an upturn, increasing with decreasing temperature [17,18]. The origin of this upturn, also seen in hole-doped cuprates [19,20], is not well understood but is thought to be associated with the underdoped materials' proximity to an antiferromagnetic insulating phase [21,22].

The low-temperature transport behavior of samples doped beyond the FSR ($x > .14$ in LCCO) has proven to be particularly intriguing. Remarkably, the resistivity in this region of the phase diagram varies linearly with temperature from a doping-dependent crossover temperature of the order of tens of Kelvin down to the lowest measured temperature of 30 mK when superconductivity is suppressed with an external magnetic field [23]. This is in stark contrast to the Fermi liquid expectation of a low-temperature T^2 resistivity, and is perhaps the most compelling evidence available for a non-Fermi-liquid ground state. In addition, it has recently been found that the low-temperature magnetoresistance (MR) of these overdoped samples is linear in magnetic field, in contrast to the conventional H^2 dependence expected for weak fields [24].

In light of this correspondence between the temperature-dependent resistivity and MR in the overdoped strange metallic ground state, in this paper we investigate whether

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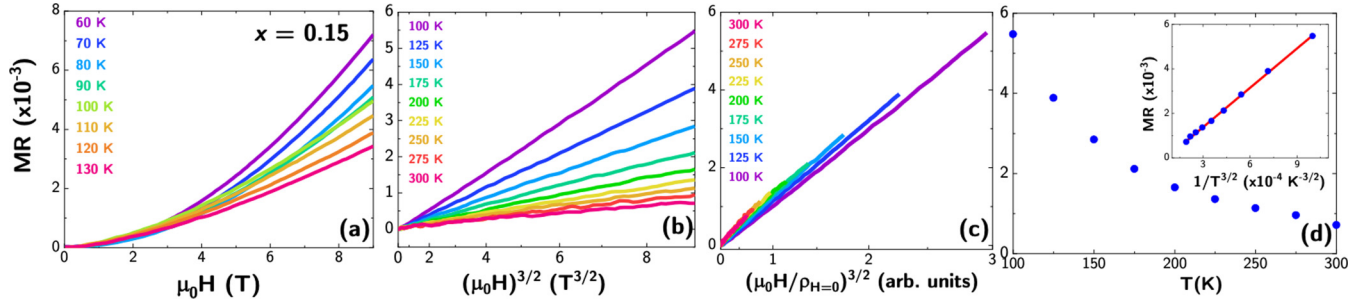


FIG. 1. Magnetoresistance (MR) data for an $x = 0.15$ LCCO sample measured up to 9 T. (a) raw $MR = [\rho(T, H) - \rho(T, 0)]/\rho(T, 0)$ curves for temperatures between 60–130 K. The MR is quadratic in fields below 80 K [24], but the crossing of the curves taken at 80, 90, and 100 K suggests the MR at 80 K grows more rapidly with field than at 100 K. That is, $MR \sim H^n$ with $n < 2$ at 100 K and above. (b) The MR from 100–300 K plotted against $(\mu_0 H)^{3/2}$. The linearity of the curves clearly implies that $MR \sim H^{3/2}$. (c) Kohler's scaling plot of the MR from (b) plotted against $[\mu_0 H/\rho(T, 0)]^{3/2}$. Collapse of all the data onto one curve would suggest Kohler's rule is obeyed. Note that deviations from collapse are visible even at the lowest fields. (d) The magnitude of the MR at 9 T plotted as a function of temperature; inset: the same data plotted against $1/T^{3/2}$. The linearity of the data on this scale suggests that $MR(T) \sim 1/T^{3/2}$ which is incompatible with Kohler's rule.

a similar correspondence exists for the high-temperature metallic state where $\rho \sim T^2$. To this end, we report measurements of the low-field transverse MR ($H \perp ab$ -plane) for LCCO samples spanning the phase diagram, from $x = 0.08$ to $x = 0.17$. All measurements are performed on c -axis oriented epitaxial thin films of LCCO grown via pulsed laser deposition on SrTiO_3 substrates. Details of the sample preparation can be found in the literature [17].

II. RESULTS AND DISCUSSION

In Fig. 1, we show representative $MR = [\rho(T, H) - \rho(T, 0)]/\rho(T, 0)$ data for an $x = 0.15$ LCCO sample. In Fig. 1(a), one sees that the curves taken at 80, 90, and 100 K cross one another, indicating the MR increases more rapidly with field at 80 K than at 100 K. Further, by inspection, the MR at lower temperatures appears to be quadratic in field, while the higher temperature curves seem to have a weaker field dependence. This is confirmed in Fig. 1(b), where we plot the MR from 100 K to room temperature as a function of $(\mu_0 H)^{3/2}$. The linearity of the curves on this scale clearly shows that the MR scales as $MR \sim H^{3/2}$, in contrast to the typical weak-field behavior of $MR \sim H^2$ anticipated for conventional metals.

A useful lens through which to consider the MR behavior of a metallic system is Kohler's rule, the statement that the MR should depend on the ratio of the mean free path to the cyclotron radius in a simple semiclassical picture. More formally, as can be seen from the Boltzmann equation, it is the statement that the MR depends only on the product of the magnetic field and scattering time or, more practically (since $\rho(T, 0) \sim \tau^{-1}$), is a function of only the ratio of the magnetic field to the zero-field resistivity [25], i.e., $MR = F[H/\rho(T, 0)]$ for some function $F(x)$. In Fig. 1(c), we attempt to assess the validity of Kohler's rule in this system by plotting the MR against $(\mu_0 H)^{3/2}$, where the near collapse of the data onto a single curve appears to suggest that Kohler's rule is obeyed. However, one can still see directly in the figure that there are deviations from a precise collapse even at modest fields.

To assess the validity of Kohler's rule in a less qualitative manner, we note that since $MR \sim H^{3/2}$ and $\rho(T, 0) \sim T^2$, a corollary to Kohler's rule is that the temperature dependence of the MR (at fixed field) must be $MR \sim H^{3/2} \sim [H/\rho(T, 0)]^{3/2} \sim 1/T^3$. The temperature-dependent MR at 9 T is plotted in Fig. 1(d), which, as can be seen in the inset, follows a clear $MR \sim 1/T^{3/2}$ behavior, in violation with Kohler scaling. That is, despite what visually appears to be a near collapse of the data in Fig. 1(c), the temperature dependence of the MR unambiguously demonstrates that Kohler's rule is violated. The failure of Kohler's rule in LCCO is hardly surprising, given the extreme simplicity of the theory and the complexity of the material under study.

Moving on from this single doping ($x = 0.15$), we plot the MR against $(\mu_0 H)^{3/2}$ for dopings $0.13 < x < 0.17$ in Fig. 2, and find that the $MR \sim H^{3/2}$ behavior is generic above 100 K. Further, in Fig. 3, we show the temperature dependence of the MR is also $\sim 1/T^{3/2}$ for all dopings studied. Consequently, we find that Kohler's rule is violated across the LCCO phase diagram.

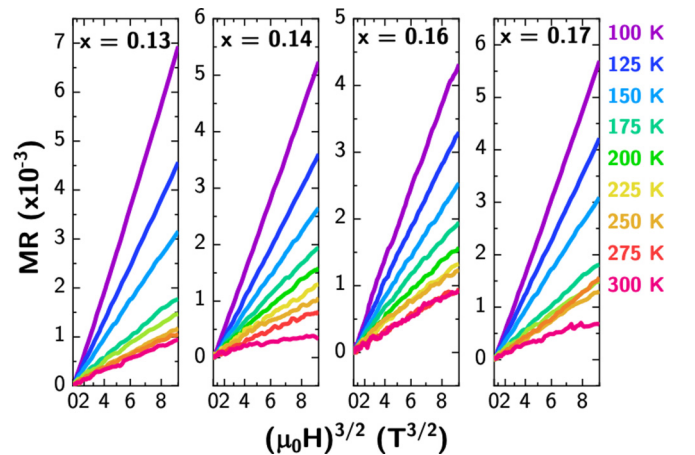


FIG. 2. Plots of the magnetoresistance (MR) versus $(\mu_0 H)^{3/2}$ for dopings $x = 0.13, 0.14, 0.16$, and 0.17 at numerous temperatures between 100–300 K. The $MR \sim H^{3/2}$ behavior is observed for each doping.

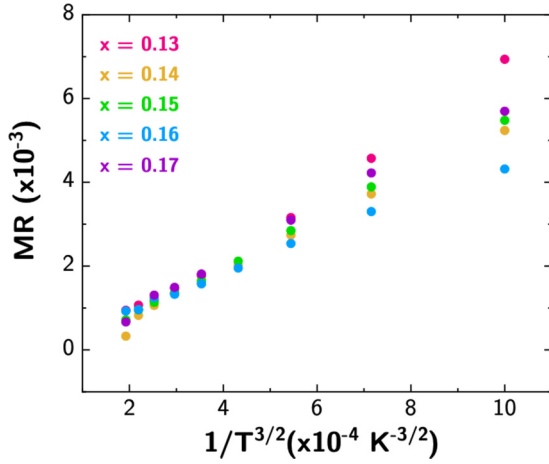


FIG. 3. Temperature-dependent magnetoresistance (MR) at 9 T for dopings $x = 0.13$ – 0.17 plotted against $1/T^{3/2}$.

To further illustrate the ubiquity of the $MR \sim H^{3/2}$ scaling at high temperatures, we plot the MR at 100 K for several dopings well below the FSR doping $x_{\text{FSR}} = 0.14$, where the $MR \sim H^{3/2}$ dependence is clearly seen. The temperature-dependent MR was not studied in detail for these dopings.

Taken together, Figs. 1–4 establish a universal $\sim H^{3/2}$ magnetic field dependence of the MR for $T \gtrsim 100$ K across the LCCO phase diagram from $x = 0.08$ – 0.17 , i.e., for all dopings with a superconducting ground state. That is, the anomalous MR power law occurs over the same wide region of the phase diagram where the anomalous high-temperature $\rho \sim T^2$ dependence of the zero-field resistivity is observed. It is then natural to speculate that these two unconventional transport phenomena are driven by a common resistive scattering mechanism.

To supplement our MR measurements, in the inset of Fig. 5(a) we present measurements of the Hall coefficient for several dopings throughout the superconducting dome. Note that even at room temperature, the Hall coefficient has a non-trivial temperature dependence, departing from single-carrier

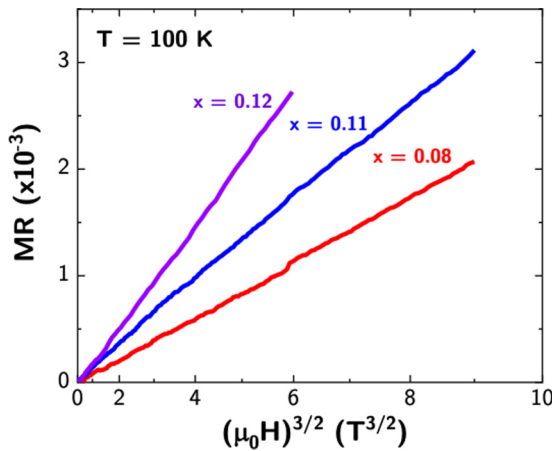


FIG. 4. Magnetoresistance (MR) at 100 K plotted against $(\mu_0 H)^{3/2}$ for several dopings below the Fermi surface reconstruction doping $x_{\text{FSR}} = 0.14$. Again, the $MR \sim H^{3/2}$ behavior is seen for all dopings.

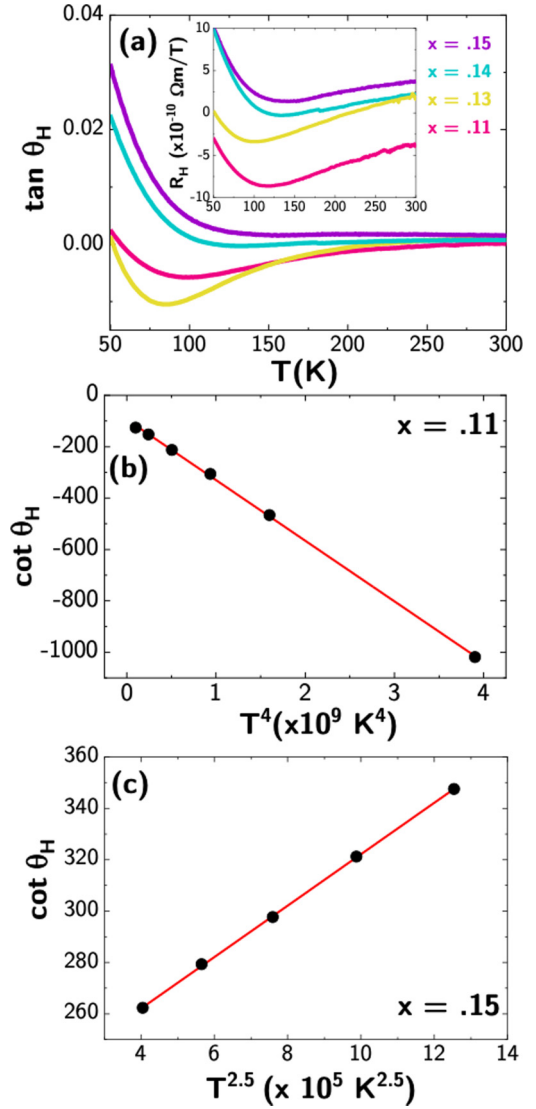


FIG. 5. High-temperature Hall response: (a) $\tan \theta_H = \rho_{xy}/\rho_{xx}$ versus temperature for several dopings of LCCO from 50 to 300 K; inset: Hall number for the same samples from 50 to 300 K. (b) $\cot \theta_H = \rho_{xx}/\rho_{xy}$ for the $x = 0.11$ sample from 150 to 300 K, plotted against T^4 . The red line is a linear fit. (c) $\cot \theta_H$ for the $x = 0.15$ sample from 150 K to 300 K, plotted against $T^{2.5}$. The red line is a linear fit.

Fermi liquid expectations. Moreover, R_H changes signs at high temperatures for dopings near the FSR, which naively suggests that both electron- and holelike carriers may be relevant to the high-temperature transport properties of LCCO within this doping range. However, given the longstanding confusion over the meaning of the Hall coefficient in the cuprates, such a conclusion may very well be premature.

To further our characterization of the high-temperature transport phenomenology, in Fig. 5(b) the tangent of the Hall angle, $\tan \theta_H \equiv \rho_{xy}/\rho_{xx}$ is shown for several dopings up to room temperature. Although the cotangent of the Hall angle, $\cot \theta_H = \rho_{xx}/\rho_{xy}$ is the typical quantity of theoretical interest, the zeros of ρ_{xy} which, as mentioned above, are present in some dopings up to 100–200 K, prevent the evaluation of this ratio for every doping.

For the $x = 0.11$ and $x = 0.15$ samples, ρ_{xy} does not change signs above 100 K, allowing for $\cot \theta_H$ to be analyzed. As shown in Figs. 5(b) and 5(c), $\cot \theta_H$ exhibits a T^4 temperature dependence for the $x = 0.11$ sample and a $T^{2.5}$ temperature dependence for the $x = 0.15$ sample. We emphasize that the in-plane resistivity ρ_{ab} has a continuous $\sim T^2$ temperature dependence over the entire temperature range under study [17], so the changes in the temperature dependence of $\cot \theta_H$ are due entirely to the Hall response.

In contrast, a Fermi liquid is expected to have $\rho_{xx} \sim T^2$ and $\rho_{xy} \sim T^0$, and thus $\cot \theta_H \sim T^2$, which is in fact one of the defining features of a Fermi liquid, and a behavior observed in the hole-doped cuprates [26,27]. Although other temperature dependencies can be realized due to changes in the Fermi surface geometry [28], multiple relaxation times [26], or antiferromagnetic spin fluctuations [29,30], a definitive explanation for the anomalous temperature dependence of the Hall angle reported here is well beyond the scope of the present paper.

A prior study [31] of another electron-doped cuprate, $\text{Pr}_{2-x}\text{Ce}_x\text{CuO}_4$, found $\cot \theta_H \sim T^\alpha$ with $3 < \alpha < 4$ for underdoped samples, which is consistent with our finding of $\cot \theta_H \sim T^4$ for our $x = 0.11$ sample. However, our finding of $\cot \theta_H \sim T^{2.5}$ in an overdoped electron-doped cuprate is of particular interest given the especially exotic transport behaviors observed in this region of the phase diagram [5,23,24]. We further note that the strong doping dependence of $\cot \theta_H$ (that is, the very different power laws exhibited on either side

of the FSR doping) contrasts sharply with the largely doping independent longitudinal response.

III. CONCLUSIONS

Altogether, we have demonstrated an anomalous $\sim H^{3/2}$ field dependence of the MR in LCCO which coexists with the well-known $\rho \sim T^2$ temperature-dependent resistivity at high temperatures for all dopings within the superconducting dome. We hope that this aspect of the high-temperature metallic state of the electron-doped cuprates will help identify the resistive scattering mechanism responsible for these unconventional transport phenomena, and perhaps its relationship to high-temperature superconductivity and/or the low-temperature strange metallic state. In contrast, we find the transverse response to be strongly doping dependent, with drastically different temperature dependencies of the Hall angle for samples doped above or below the FSR doping $x_{\text{FSR}} = 0.14$.

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