

Dynamical Methods for Studying Stability and Noise in Frequency Combs Sources

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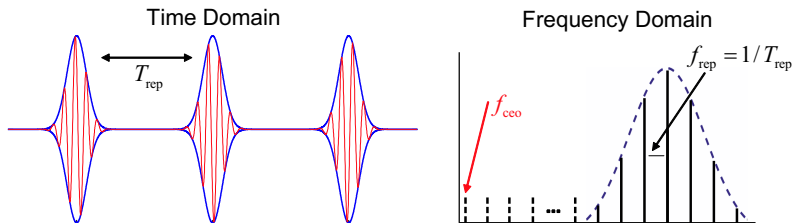
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Background

In 1999–2000: Frequency combs were invented

“The excitement surrounding the rapid evolution in these fields since 1999 gives us a hint of what it must have been like after 1927 when the first ideas of quantum mechanics were introduced. . .”

— J. L. Hall and T. W. Hänsch, 2005 Nobel prize winners in Physics

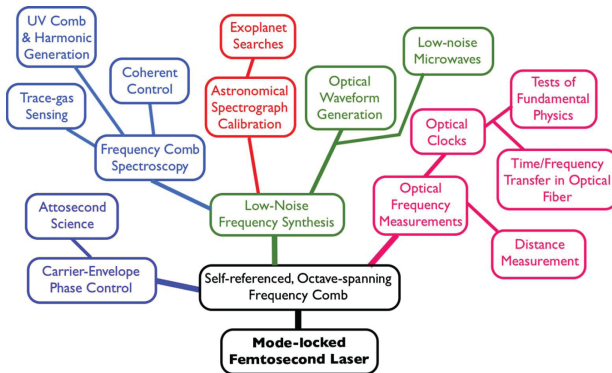


The key advance was electronically locking f_{ceo} and f_{rep} !

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Background

- Frequency and time measurement was revolutionized
- Many new applications opened up¹



¹ S. Diddams, J. Opt. Soc. Am. B **27**, 51 (2010).

Motivation

- Advances have all been through “cut-and-try” experimentation
- Theoretical tools for analyzing and designing frequency combs are primitive
 - ▶ “Brute force” simulations or rough analytical approximations
 - ★ Adequate for post-hoc analysis; inadequate for design
 - ★ yield limited insight into the sources of instability

Key theoretical questions:

- Where in the adjustable parameter space are combs stable?
- What is the noise performance?
- How can we optimize the comb?
 - ▶ high output power and/or large bandwidth and/or low noise

Approach

Combine: 400 years of dynamical systems theory

+

modern computers and algorithms

(linear algebra + root-finding)

to answer these questions

Origins of the Dynamical Approach

These dynamical ideas are very old!

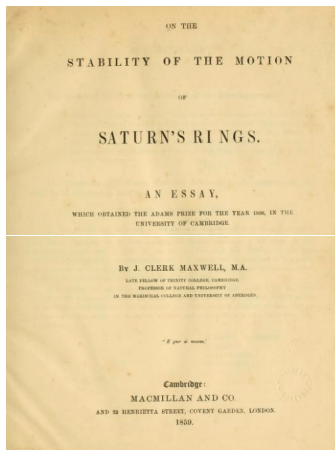
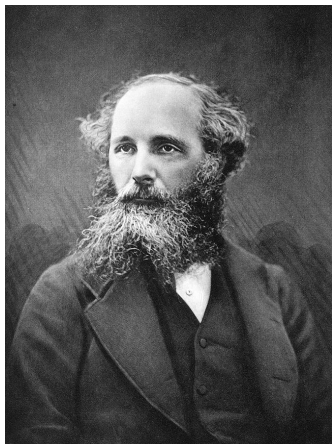
- The pendulum clock:
(Galileo – 1632; Huygens – 1673; Euler – 1736)
- Stability of the solar system:
 - ▶ Two body problem: Newton – 1686
 - ▶ Three body problem:
In general, not solvable, but . . .
stable fixed points found by Lagrange – 1772 (observed 1906)
- Application to continuous systems . . .
(described by partial differential equations)

Origins of the Dynamical Approach

Before Maxwell's Equations (1861) . . .

Before the Maxwell-Boltzmann distribution (1865) . . .

Maxwell explained the stability of Saturn's rings (1859)



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Numerous Applications

- Control of satellite orbits
- Electronic system design
- Plasma systems (tokamaks, ionosphere, . . .)
- Mechanical, chemical, and fluid systems
- Biological systems (heart, animal populations)
- Economic systems
- Lasers/other optical resonators
 - ▶ but mostly in highly simplified, almost analytical approximations!

Frequency Comb Systems

Passively modelocked lasers (slow saturable gain):

$$\frac{\partial u}{\partial T} = \left[-i\phi + t_s \frac{\partial}{\partial t} + \frac{g(|u|)}{2} \left(1 + \frac{1}{2\omega_g^2} \frac{\partial^2}{\partial t^2} \right) - \frac{l}{2} - i\frac{\beta''}{2} \frac{\partial^2}{\partial t^2} + i\gamma|u|^2 - f_{sa}(|u|) \right] u,$$

$$g(|u|) = g_0 [1 + w_0/(P_{\text{sat}} T_R)]^{-1}$$

cubic-quintic model (fast saturable absorber)¹

$$f_{sa}(|u|) = \delta|u|^2 - \sigma|u|^4$$

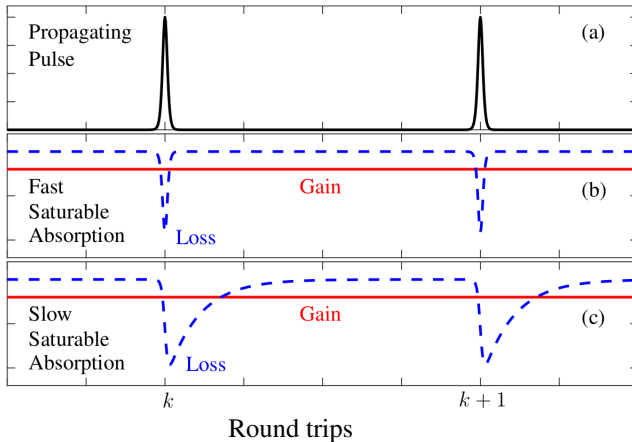
SESAM model (slow saturable absorber)²

$$f_{sa}(|u|) = -\frac{\rho}{2}n(t, T), \quad \frac{\partial n(t, T)}{\partial T} = \frac{1-n}{T_A} - \frac{|u(t, T)|}{w_A}n$$

¹S. Wang et al., J. Opt. Soc. Am. B **31**, 2914 (2014).

²S. Wang et al., Opt. Lett. **42**, 2362 (2017).

Saturable Absorption



The loss is saturated by the incoming pulse and then recovers

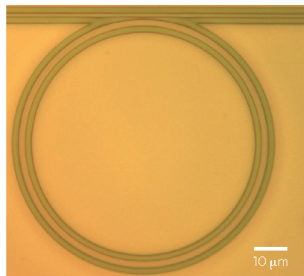
- almost instantaneously with fast absorbers
- slowly compared to the pulse duration with slow absorbers

Frequency Comb Systems

Microresonators

The Lugiato-Lefever Equation:³

$$\frac{\partial \psi}{\partial t} - i \frac{\partial^2 \psi}{\partial x^2} - i |\psi|^2 \psi + (1 + i\alpha) \psi - F = 0$$



System Parameters:

- α : frequency detuning (−5 to 10)
- F : pump amplitude (0 to 4)
- L : microresonator length (25 to 200)

³Z. Qi et al., Conf. Lasers Elect.-Opt. (2018), paper SF2A.6.

References

- Haus modelocking equation:
 - ▶ S. Wang et al., J. Opt. Soc. Am. B **31**, 2914 (2014)
 - ▶ Martinez, Fork, and Gordon, J. Opt. Soc. Am. B **2**, 753 (1985)
- Cubic-quintic modelocking equation:
 - ▶ Moores, Opt. Comm. **96**, 65 (1993)
 - ▶ Soto-Crespo et al., J. Opt. Soc. Am. B **13**, 1439 (1996)
 - ▶ Articles in *Dissipative Solitons*, Akhmediev and Ankiewicz ed. (2005)
- SESAM equation:
 - ▶ Kärtner et al., IEEE J. Sel. Quantum Electron. **2**, 540 (1996)
- Lugiato-Lefever equation:
 - ▶ Lugiato and Lefever, Phys. Rev. Lett. **58**, 2209 (1987)
 - ▶ Haelterman et al., Opt. Comm. **91**, 401 (1992)
 - ▶ Matsko et al., Opt. Lett. **36**, 2845 (2011)
 - ▶ Coen et al., Opt. Lett. **38**, 1790 (2013)
 - ▶ Chembo and Menyuk, Phys. Rev. A **87**, 053852 (2013)

Standard, “brute-force” approach

- Solve the evolution equations for many roundtrips
- Use a noisy initial condition
- Convergence $\iff$ existence + stability
- Change parameters; repeat

Advantages:

- Easy to program
- Intuitive (mimics experiments)

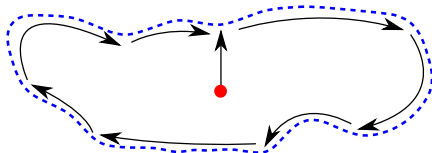
Disadvantages:

- Computationally slow
- Ambiguous near a stability boundary
- Limited insight into sources of instability

This approach is better for analysis than synthesis!

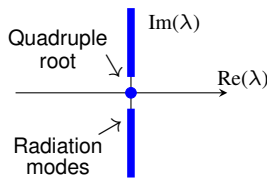
Our Approach

- Solve the evolution equations **once** to find a stationary solution
 - ▶ in a highly stable case
- Determine the stationary solution as parameters vary by solving a **root-finding problem**
- In parallel, find the eigenvalues of the linearized evolution equation
 - ▶ **The dynamical spectrum**
 - ▶ A stable solution has no eigenvalues with positive real parts
- Find parameters where one or more eigenvalues hit the imaginary axis
- Track the stability boundary

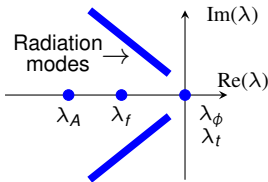


Atlas of Dynamical Spectra

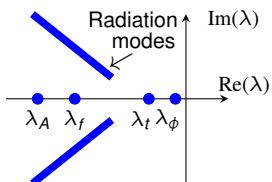
Classical NLS Spectrum



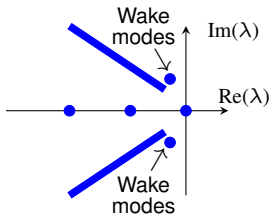
Soliton Laser



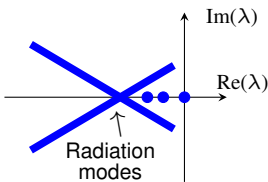
With locking



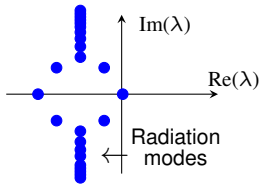
SESAM laser



Normal dispersion lasers



Microresonator



Advantages

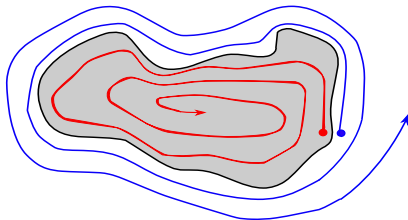
- $10^3 - 10^5$ times faster than brute-force solutions
- Unambiguous determination of stable operating parameter regimes
- Allows rapid mapping and optimization of solution properties
 - ▶ bandwidth, power, noise. . .
- Yields insight into the sources of instability

Disadvantages

- More difficult to program
- **The concepts are unfamiliar to many optical experimentalists**

Two important caveats

- Accessibility vs. stability
 - ▶ Dynamical methods do not tell you how to access stable solutions
 - ★ Example: single solitons are hard to access in microresonators



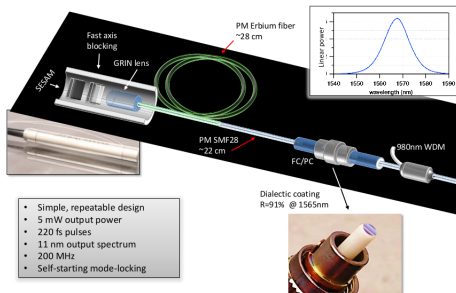
Two important caveats

- Accessibility vs. stability
 - ▶ Dynamical methods do not tell you how to access stable solutions
 - ★ Example: single solitons are hard to access in microresonators
- Unstable system evolution
 - ▶ Dynamical methods do not tell you how an unstable solution evolves
 - ★ Chaos, another stable solution, breathers are all possible

Dynamical and evolutionary methods are complementary!

The SESAM Modelocked Fiber Laser

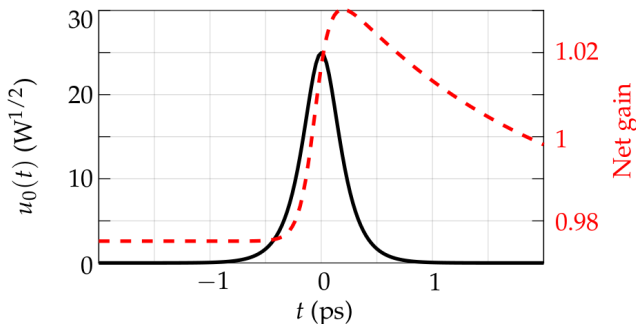
- The system is built using
 - ▶ telecom grade polarization-maintaining (PM) components
 - ▶ highly-doped erbium-doped fiber
 - ▶ highly non-linear PM fibers
 - ▶ a semiconductor saturable absorber mirror (SESAM)
- Output: highly stable 200 MHz combs with $P_{av} = 5$ mW



¹ Sinclair et al., Opt. Express **22**, 6996 (2014).

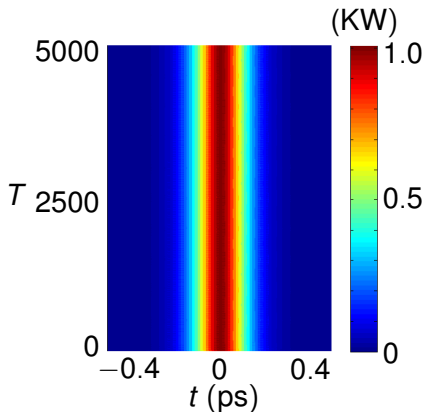
Balance of Energy

- The SESAM and the linear gain open a gain window that allows the pulse to grow
- A soliton wake instability will occur when g_0 becomes sufficiently large or β'' becomes sufficiently small

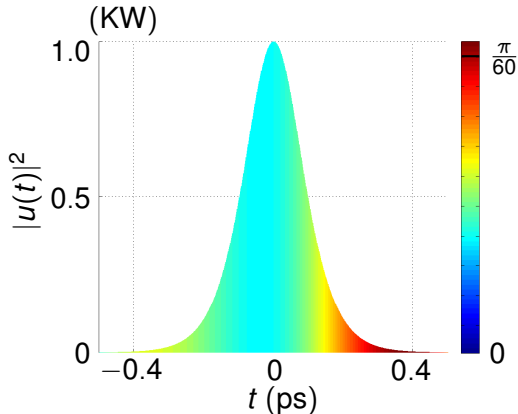


Stable Operation

The stable pulse is close in shape to a sech pulse
(soliton solution of the nonlinear Schrödinger equation)



Color indicates the pulse power

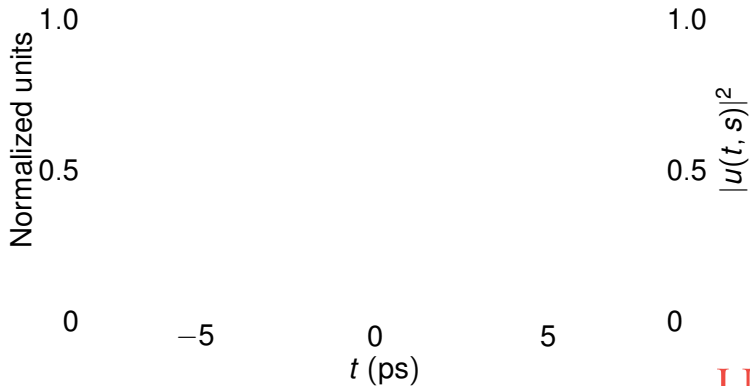


Color indicates the phase of the pulse

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Soliton Wake Instability

Quasi-periodicity is observed in the evolution



Soliton Wake Instability

Quasi-periodic Letter is observed in the evolution

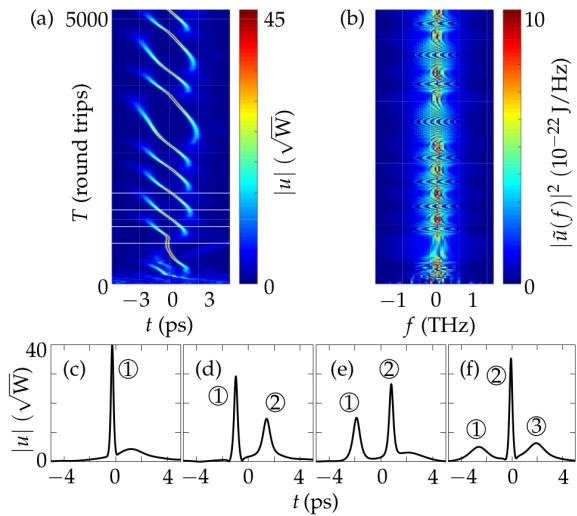
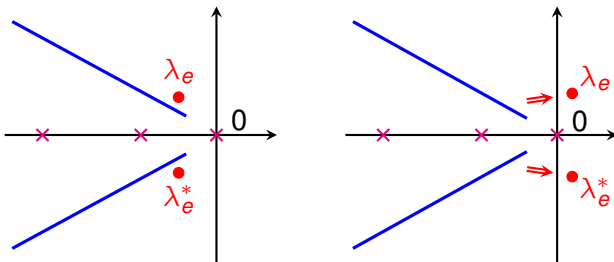


Fig. 4. The saturated growth rate $\lambda_{w\pm} = -7.909 \times 10^{-3}$ indicate how the eigenvalue

Dynamical Spectrum

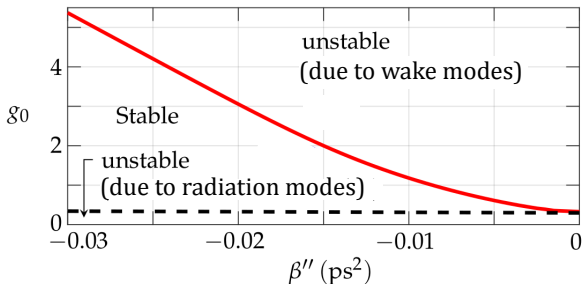
- As g_0 increases:
 - a pair of eigenvalues λ_e, λ_e^* emerge from the radiation modes (edge bifurcation)
- $\text{Re}[\lambda_e], \text{Re}[\lambda_e^*]$ become positive, leading to instability (Hopf bifurcation)¹



¹S. Wang et al., Opt. Lett. **42**, 2362 (2017).

Stable Region

- Continuous modes become unstable when the gain is too low
- The wake mode instability occurs when
 - ▶ the unsaturated gain becomes large
 - ▶ the group delay dispersion becomes small



Once the stable region is known, optimization is possible!

Cnoidal Waves in Microresonators

- Solitons are a special case of a broader class of periodic waveforms: **cnoidal waves**

- Cnoidal waves include:

- ▶ Single solitons
- ▶ Soliton crystals
- ▶ Turing rolls

- Cnoidal wave attributes:

On the one hand, they . . .

- ▶ have clean spectra with evenly spaced comb lines (like single solitons)
- ▶ have analytical solutions that exist in the no-loss limit (like single solitons)¹
- ▶ can have a broad bandwidth (like solitons)

¹Z. Qi et al., J. Opt. Soc. Am. B **34**, 785 (2017).

Cnoidal Waves in Microresonators

- Solitons are a special case of a broader class of periodic waveforms: **cnoidal waves**
- Cnoidal waves include:
 - ▶ Single solitons
 - ▶ Soliton crystals
 - ▶ Turing rolls
- Cnoidal wave attributes:

On the other hand, they . . .

 - ▶ can be easily and deterministically accessed
(unlike single solitons or soliton crystals)
 - ▶ use the pump more efficiently and produce higher power comb lines

BUT

with lines spaced farther apart

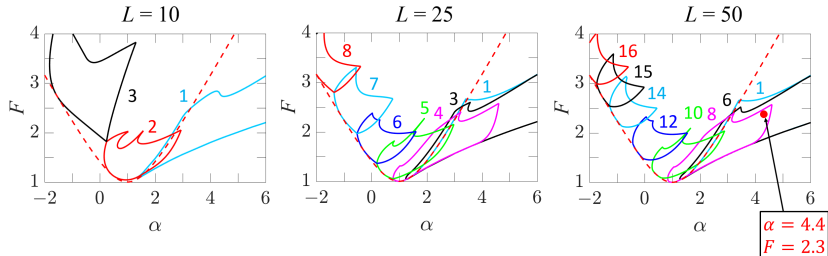
Stable Regions

Stable regions for different periodicities

F = normalized pump power

α = normalized detuning

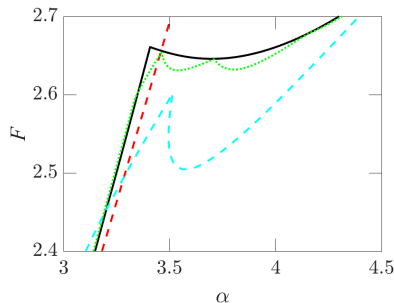
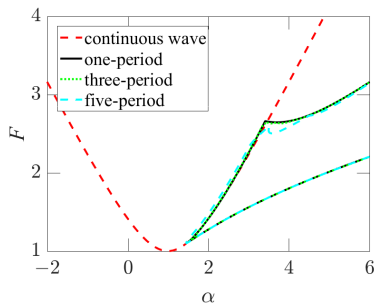
L = normalized diameter



- Continuous waves are stable below the red-dashed line
- Below $\alpha = 41/30 \simeq 1.37$, cnoidal waves can be easily accessed by raising the pump power
- Cnoidal waves are stable in a U-shaped region in α - F space
- Moving along this region, different values of N_{per} can be deterministically accessed

Stable Regions ($L = 50$)

Why are single solitons hard to access

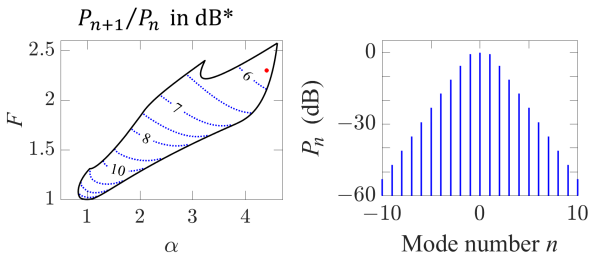


- Substantial overlap with other cnoidal waves
- Almost complete overlap with continuous waves

By contrast the periodicity-8 cnoidal wave can be deterministically accessed!

Periodicity-8 Cnoidal Waves ($L = 50$)

Controllability and bandwidth



- With FSR = 125 GHz; 30 dB down bandwidth = 24 THz
- Uses the pump more efficiently than a single soliton

Large bandwidth can be obtained!

* large n limit ($n \gtrsim 4$)

Conclusions

- Powerful dynamical methods can be combined with modern computer algorithms to:
 - ▶ rapidly determine where in the parameter space stable solutions lie
 - ▶ yield important insights into the sources of instability
 - ▶ rapidly determine the noise performance
 - ▶ optimize the system performance

Conclusions

- We have applied these methods to:
 - ▶ Laser models with slow saturable gain
 - ★ Identified parameters ranges where two stable solutions exist
 - ★ Compare the cubic-quintic model to other models
 - ▶ SESAM laser
 - ★ Characterized the wake mode instability
 - ★ Determined where stable solutions exist
 - ★ Explained the appearance of sidebands
 - ★ Characterized the noise performance
 - ★ Optimized the system parameters
 - ▶ Microresonators
 - ★ Determined where cnoidal waves are stable
 - ★ Explained why single solitons are hard to access
 - ★ Found broadband, easily accessible cnoidal wave solutions
 - ★ Determined the impact of thermal effect
 - ★ Determined the impact of avoided crossings

Summary

These methods are an important complement to widely used evolutionary methods and should be commonly used!

Our software is available at:

<http://www.umbc.edu/photonics/software.html>

See: “Dynamical method to evaluate noise. . .”
 “Boundary tracking algorithms”
 “Stability boundary tracking for cnoidal wave solutions”