

THE SHARP TIME DECAY RATE OF THE ISENTROPIC NAVIER-STOKES SYSTEM IN $\mathbb{R}^3$

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ABSTRACT. We investigate the sharp time decay rates of the solution U for the compressible Navier-Stokes system (1.1) in $\mathbb{R}^3$ to the constant equilibrium $(\bar{\rho} > 0, 0)$ when the initial data is a small smooth perturbation of $(\bar{\rho}, 0)$. Let $\tilde{U}$ be the solution to the corresponding linearized equations with the same initial data. Under a mild non-degenerate condition on initial perturbations, we show that $\|U - \tilde{U}\|_{L^2}$ decays at least at the rate of $(1+t)^{-\frac{5}{4}}$, which is faster than the rate $(1+t)^{-\frac{3}{4}}$ for the $\tilde{U}$ to its equilibrium $(\bar{\rho}, 0)$. Our method is based on a combination of the linear sharp decay rate obtained from the spectral analysis and the energy estimates.

1. Introduction. In this paper, we are concerned with the sharp decay rates of solutions to the Cauchy problem for the isentropic Navier-Stokes equations:

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho u) = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \partial_t(\rho u) + \operatorname{div}(\rho u \otimes u) + \nabla p(\rho) = \operatorname{div} T, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \lim_{|x| \rightarrow \infty} \rho = \bar{\rho}, \quad \lim_{|x| \rightarrow \infty} u = 0, & t \in \mathbb{R}^+, \\ (\rho, u)|_{t=0} = (\rho_0, u_0), & x \in \mathbb{R}^3, \end{cases} \quad (1.1)$$

which governs the motion of a isentropic compressible viscous fluid. The unknown functions ρ and u represent the density and velocity of the fluid respectively. The pressure $p = p(\rho)$ is a smooth function in a neighborhood of a positive constant $\bar{\rho}$ s.t. $p'(\bar{\rho}) > 0$. T is the viscosity stress tensor given by $T = \mu(\nabla u + (\nabla u)^t) + \nu(\operatorname{div} u)I$ with I the identity matrix. We assume that the constant viscosity coefficients $\mu > 0$

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and ν satisfy $\nu + \frac{2}{3}\mu > 0$. Throughout this article, by *optimal time decay rate*, we refer to the best possible decay rate in upper bound as many literatures, and the *sharp time decay rate* includes the best possible upper and lower bounds.

Using the classical spectral method, the optimal time decay rate (upper bound) of the linearized equations of the isentropic Navier-Stokes equations are well known. One may then expect that the small solution of the nonlinear equations (1.1) have the same decay rate as the linear one. Our work is devoted to proving the sharp time decay rate (for both upper and lower bound) for the nonlinear system.

In the case of one space dimension, Zeng [24] and Liu-Zeng [15] offered a detailed analysis of the solution to a class of hyperbolic-parabolic system through point-wise estimate, including the isentropic Navier-Stokes system. For multi-dimensional Navier-Stokes equations (and/or Navier-Stokes-Fourier system), the H^s global existence and time-decay rate of strong solutions with the initial perturbation small in $H^s \cap L^1$ are obtained in whole space first by A. Matsumura and T. Nishida [17], [18]. When the small initial perturbation belongs to H^3 only, using a weighted energy method, A. Matsumura [16] showed the time-decay rate $(1+t)^{-\frac{3}{4}}$ of upper bound in L^∞ -norm. Since then, there are concrete development on the upper bound time-decay estimates: the optimal L^p (with $2 \leq p \leq \infty$) upper bound decay rate was proved by G. Ponce [19], combining the spectral analysis on linearized system and the energy method for small initial perturbation in L^1 . For the isentropic Navier-Stokes equations with artificial viscosity, D. Hoff and K. Zumbrun [6], [7] studied the Green's function and derived the L^p ($1 \leq p \leq \infty$) upper bound time decay rate of diffusive waves for the small initial perturbation belongs to $H^m \cap L^1$ with $m \geq 4$. Liu and Wang [14] studied the point-wise estimates of the Green function of the linearized isentropic Navier-Stokes system in 3D and then analyzed the coupling of nonlinear diffusion waves, obtained the optimal (upper bound) decay rate. These results were further extended to the exterior problem [12], [11], or the half space problem [9], [10], [8]. Recently, Guo and Wang in [5] developed a new general energy method for proving the optimal (upper bound) time decay rates of the solutions to the dissipative equations in the whole space, using a family of scaled energy estimates with minimum derivative counts and interpolations among them without linear decay analysis.

When additional external force is taken into account, the external force does affect the long time behavior of solutions. The upper bound of time decay rates were studied intensively, see for instance [1] and [2] on unbounded domain, [22], [23] on the convergence of the non-stationary flow to the corresponding steady flow when the initial data are small in $H^3 \cap L^{\frac{6}{5}}$, and [4], [3], on the optimal $L^p - L^q$ upper bound decay rates for potential forces.

The main goal of current paper is to establish the *sharp decay rate*, on both upper and lower bounds, to the solutions of (1.1) using relatively simple energy method. We remark that similar results had been pursued by M. Schonbek [20], [21] for incompressible Navier-Stokes equations, and by Li, Matsumura-Zhang [13] for isentropic Navier-Stokes-Poisson system. Although they share the same spirit in obtaining the lower bound decay rates, the feature of the spectrum near zero exhibits quite different behaviors, leading to different analysis. For instance, we explored the elegant structure of the higher order nonlinear terms of Navier-Stokes, when choosing conservative variables: density and momentum. The conservative form of the sharp equations provided a natural derivative structure in these terms,

leading to the possibility of a faster decay rate estimate. We will make a more detailed comparison later in this paper.

Define $n = \rho - \bar{\rho}$, and let $m = \rho u = (n + \bar{\rho})u$ be the momentum. We rewrite (1.1) as

$$\begin{cases} \partial_t n + \operatorname{div} m = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \partial_t m + c^2 \nabla n - \bar{\mu} \Delta m - (\bar{\mu} + \bar{\nu}) \nabla \operatorname{div} m = F, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \lim_{|x| \rightarrow \infty} n = 0, \quad \lim_{|x| \rightarrow \infty} m = 0, & t \in \mathbb{R}^+, \\ (n, m)|_{t=0} = (\rho_0 - \bar{\rho}, \rho_0 u_0), & x \in \mathbb{R}^3, \end{cases} \quad (1.2)$$

where $\bar{\mu} = \frac{\mu}{\bar{\rho}}$, $\bar{\nu} = \frac{\nu}{\bar{\rho}}$, $c = \sqrt{p'(\bar{\rho})} > 0$ is the sound speed, and

$$\begin{aligned} F = & -\operatorname{div} \left\{ \frac{m \otimes m}{n + \bar{\rho}} + \bar{\mu} \nabla \left(\frac{nm}{n + \bar{\rho}} \right) \right\} \\ & - \nabla \left\{ (\bar{\mu} + \bar{\nu}) \operatorname{div} \left(\frac{nm}{n + \bar{\rho}} \right) + (p(n + \bar{\rho}) - p(\bar{\rho}) - c^2 n) \right\}. \end{aligned}$$

It is this structure of F that plays an important role in our analysis.

Our aim is to obtain a clear picture of the large time behavior of $U = (n, m)$ in $L^2(\mathbb{R}^3)$ when $U_0 = (\rho_0 - \bar{\rho}, \rho_0 u_0)$ is sufficiently smooth and small. We introduce the following initial value problem of the linearized Navier-Stokes system corresponding to (1.2):

$$\begin{cases} \partial_t \tilde{n} + \operatorname{div} \tilde{m} = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \partial_t \tilde{m} + c^2 \nabla \tilde{n} - \bar{\mu} \Delta \tilde{m} - (\bar{\mu} + \bar{\nu}) \nabla \operatorname{div} \tilde{m} = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \lim_{|x| \rightarrow \infty} \tilde{n} = 0, \quad \lim_{|x| \rightarrow \infty} \tilde{m} = 0, & t \in \mathbb{R}^+, \\ (\tilde{n}, \tilde{m})|_{t=0} = (\rho_0 - \bar{\rho}, \rho_0 u_0), & x \in \mathbb{R}^3, \end{cases} \quad (1.3)$$

where $\bar{\mu} = \frac{\mu}{\bar{\rho}}$, $\bar{\nu} = \frac{\nu}{\bar{\rho}}$, $c = \sqrt{p'(\bar{\rho})}$. It is known that the L^2 -norm of $\tilde{U} = (\tilde{n}, \tilde{m})$ decays at the optimal upper bound rate $(1+t)^{-\frac{3}{4}}$ for generic small initial data, see for instance [18]. A detailed proof on the optimal lower and upper bound rate will be given in the section 3 of this paper. In section 4, we prove that $\|(U - \tilde{U})(\cdot, t)\|_{L^2}$ decays at a faster rate than $\|\tilde{U}(\cdot, t)\|_{L^2}$, under some reasonable conditions on the initial data. Therefore, $\|U(\cdot, t)\|_{L^2}$ shares the sharp decay rate of $(1+t)^{-\frac{3}{4}}$.

Notation. For $a \lesssim b$, we mean that there is a uniform constant C , which may be different on different lines, such that $a \leq Cb$. And $a \approx b$ stands for $a \lesssim b$ and $b \lesssim a$.

We now state our main result.

Theorem 1.1. *Assume that $(n_0, m_0) \in L^1(\mathbb{R}^3) \cap H^3(\mathbb{R}^3)$, $\delta_0 =: \|(n_0, m_0)\|_{L^1(\mathbb{R}^3) \cap H^3(\mathbb{R}^3)}$ is sufficiently small, and*

$$\int_{\mathbb{R}^3} (n_0, m_0) dx \neq 0, \quad (1.4)$$

then there is a unique global classical solution $\tilde{U} = (\tilde{n}, \tilde{m}) \in \mathcal{C}([0, \infty); H^3(\mathbb{R}^3))$ of the linearized system (1.3) satisfying for some positive constant C

$$\begin{aligned} C^{-1}(1+t)^{-\frac{3}{4}-\frac{k}{2}} & \leq \|\nabla^k \tilde{n}(t)\|_{L^2(\mathbb{R}^3)} \leq C(1+t)^{-\frac{3}{4}-\frac{k}{2}}, \quad k = 0, 1, 2, 3, \\ C^{-1}(1+t)^{-\frac{3}{4}-\frac{k}{2}} & \leq \|\nabla^k \tilde{m}(t)\|_{L^2(\mathbb{R}^3)} \leq C(1+t)^{-\frac{3}{4}-\frac{k}{2}}, \quad k = 0, 1, 2, 3, \end{aligned}$$

and the initial value problem (1.2) has a unique solution $U = (n, m) \in \mathcal{C}([0, \infty); H^3(\mathbb{R}^3))$. Moreover, let $n_h = n - \tilde{n}$ and $m_h = m - \tilde{m}$, then it holds that

$$\begin{aligned} \|\nabla^k(n_h, m_h)(t)\|_{L^2(\mathbb{R}^3)} &\lesssim \delta_0^2(1+t)^{-\frac{5}{4}-\frac{k}{2}}, \quad k = 0, 1, 2, \\ \|\nabla^3 m_h(t)\|_{L^2(\mathbb{R}^3)} &\lesssim \delta_0^2(1+t)^{-\frac{11}{4}}, \quad \|\nabla^3 n_h(t)\|_{L^2(\mathbb{R}^3)} \lesssim \delta_0(1+t)^{-\frac{7}{4}}. \end{aligned}$$

As a consequence, there exists a positive constant C_1 such that

$$\begin{aligned} C_1^{-1}(1+t)^{-\frac{3}{4}-\frac{k}{2}} &\leq \|\nabla^k n(t)\|_{L^2(\mathbb{R}^3)} \leq C_1(1+t)^{-\frac{3}{4}-\frac{k}{2}}, \quad k = 0, 1, 2, \\ C_1^{-1}(1+t)^{-\frac{3}{4}-\frac{k}{2}} &\leq \|\nabla^k m(t)\|_{L^2(\mathbb{R}^3)} \leq C_1(1+t)^{-\frac{3}{4}-\frac{k}{2}}, \quad k = 0, 1, 2, 3. \end{aligned}$$

Remark 1.1. We remark that this theorem is valid under the condition (1.4) which is important in the lower bound estimate to the linearized problem. When (1.4) fails, the decay rate of the linearized system (1.3) depends on the order of the degeneracy of moments. Assume $(n_0, m_0) \in L^1 \cap H^3$ and belong to certain appropriate weighted L^p spaces, similar situation happened also in the incompressible Navier-Stokes equations, c.f. [20], [21]. We also note that our condition (1.4) is weaker than those in most of previous results where the differentiability of Fourier transform of initial disturbance is required in general.

Remark 1.2. In [13], Li, Matsumura-Zhang proved the lower bound decay rate of the linearized isentropic Navier-Stokes-Poisson system, they only require $|\hat{n}_0(\xi)| > c_0 > 0$ for $|\xi| \ll 1$ with c_0 a constant due to the special structure of the spectrum from the help of the Poisson term. This condition is proposed in Fourier space, similar to (1.4) in some sense. In our case, the spectrum is different and the different structure leads to different sharp decay rates.

2. Basic energy estimates. In what follows, we will set $n = \rho - \bar{\rho}$, $u = u - 0$. We rewrite (1.1) in the perturbation form as

$$\begin{cases} \partial_t n + \bar{\rho} \operatorname{div} u = -n \operatorname{div} u - u \cdot \nabla n, \\ \partial_t u + \gamma \bar{\rho} \nabla n - \bar{\mu} \Delta u - (\bar{\mu} + \bar{\nu}) \nabla \operatorname{div} u \\ \quad = -u \cdot \nabla u - \bar{\mu} f(n) \Delta u - (\bar{\mu} + \bar{\nu}) f(n) \nabla \operatorname{div} u - g(n) \nabla n, \\ \lim_{|x| \rightarrow \infty} n = 0, \quad \lim_{|x| \rightarrow \infty} u = 0, \\ (n, u)|_{t=0} = (\rho_0 - \bar{\rho}, u_0), \end{cases} \quad (2.1)$$

where $\bar{\mu} = \frac{\mu}{\bar{\rho}}$, $\bar{\nu} = \frac{\nu}{\bar{\rho}}$, $\gamma = \frac{p'(\bar{\rho})}{\bar{\rho}^2}$, and the nonlinear functions f and g are defined by

$$f(n) := \frac{n}{n + \bar{\rho}}, \quad g(n) := \frac{p'(n + \bar{\rho})}{n + \bar{\rho}} - \frac{p'(\bar{\rho})}{\bar{\rho}}. \quad (2.2)$$

We assume that there exist a time of existence $T > 0$ and sufficiently small $\delta > 0$, such that *a priori* estimate

$$\|n(t)\|_{H^3} + \|u(t)\|_{H^3} \leq \delta, \quad (2.3)$$

holds for any $t \in [0, T]$. First of all, by (2.3) and Sobolev's inequality, we obtain that

$$\frac{\bar{\rho}}{2} \leq n + \bar{\rho} \leq 2\bar{\rho}.$$

Hence, we immediately have

$$|f(n)|, |g(n)| \leq C|n|, \quad |\nabla^k f(n)|, |\nabla^k g(n)| \leq C \quad \forall k \in \mathbb{N}^+, \quad (2.4)$$

where $f(n)$ and $g(n)$ are nonlinear functions of n defined by (2.2).

Next, we begin with the energy estimates including n and u themselves. The following results is essentially due to A. Matsumura and T. Nishida [17], [18].

Theorem 2.1. *Assume that $(n_0, u_0) \in H^3(\mathbb{R}^3)$, then there exists a constant $\delta_0 > 0$ such that if*

$$\|n_0\|_{H^3} + \|u_0\|_{H^3} \leq \delta_0,$$

then the problem (2.1) admits a unique global solution $(n(t), u(t))$ satisfying that for all $t \geq 0$,

$$\|n(t)\|_{H^3}^2 + \|u(t)\|_{H^3}^2 + \int_0^t (\|\nabla n(\tau)\|_{H^2}^2 + \|\nabla u(\tau)\|_{H^3}^2) d\tau \leq C (\|n_0\|_{H^3}^2 + \|u_0\|_{H^3}^2),$$

where C is a positive constant independent of time.

The proof of this theorem is divided into several subsections.

2.1. Energy estimate for $0 \leq k \leq 3$. For $k = 0$, multiplying the first equation in (2.1) by γn and the second equation in (2.1) by u , summing up and then integrating the result over $\mathbb{R}^3$ by parts. By virtue of Hölder's inequality, Sobolev's inequality and the fact (2.4), we obtain that

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^3} (\gamma |n|^2 + |u|^2) dx + \int_{\mathbb{R}^3} (\bar{\mu} |\nabla u|^2 + (\bar{\mu} + \bar{\nu}) |\operatorname{div} u|^2) dx \\ &= \int_{\mathbb{R}^3} \gamma (-n \operatorname{div} u - u \cdot \nabla n) n - (u \cdot \nabla u + \bar{\mu} f(n) \Delta u \\ & \quad + (\bar{\mu} + \bar{\nu}) f(n) \nabla \operatorname{div} u + g(n) \nabla n) \cdot u dx \\ &\lesssim \|n\|_{L^3} \|\nabla u\|_{L^2} \|n\|_{L^6} + (\|u\|_{L^3} \|\nabla u\|_{L^2} + \|n\|_{L^3} \|\nabla n\|_{L^2}) \|u\|_{L^6} \\ & \quad + (\|u\|_{L^\infty} \|\nabla n\|_{L^2} + \|n\|_{L^\infty} \|\nabla u\|_{L^2}) \|\nabla u\|_{L^2} \\ &\lesssim (\|n\|_{L^3} + \|u\|_{L^3} + \|n\|_{L^\infty} + \|u\|_{L^\infty}) (\|\nabla n\|_{L^2}^2 + \|\nabla u\|_{L^2}^2). \end{aligned} \tag{2.5}$$

Now for $1 \leq k \leq 3$, applying ∇^k to (2.1) and then multiplying the first equation by $\gamma \nabla^k n$ and the second equation by $\nabla^k u$, summing up and integrating over $\mathbb{R}^3$. For $k = 1$ we have

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^3} (\gamma |\nabla n|^2 + |\nabla u|^2) dx + \int_{\mathbb{R}^3} (\bar{\mu} |\nabla^2 u|^2 + (\bar{\mu} + \bar{\nu}) |\nabla \operatorname{div} u|^2) dx \\ &\lesssim (\|n\|_{L^\infty} + \|u\|_{L^\infty} + \|\nabla n\|_{L^\infty} + \|\nabla u\|_{L^\infty}) (\|\nabla n\|_{L^2}^2 + \|\nabla u\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2). \end{aligned} \tag{2.6}$$

For $k = 2$ we have

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^3} (\gamma |\nabla^2 n|^2 + |\nabla^2 u|^2) dx + \int_{\mathbb{R}^3} (\bar{\mu} |\nabla^3 u|^2 + (\bar{\mu} + \bar{\nu}) |\nabla^2 \operatorname{div} u|^2) dx \\ &\lesssim (\|n\|_{L^\infty} + \|u\|_{L^\infty} + \|\nabla n\|_{L^\infty} + \|\nabla u\|_{L^\infty}) (\|\nabla^2 n\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2 + \|\nabla^3 u\|_{L^2}^2). \end{aligned} \tag{2.7}$$

For $k = 3$ we have

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^3} (\gamma |\nabla^3 n|^2 + |\nabla^3 u|^2) dx + \int_{\mathbb{R}^3} (\bar{\mu} |\nabla^4 u|^2 + (\bar{\mu} + \bar{\nu}) |\nabla^3 \operatorname{div} u|^2) dx \\
& \lesssim (\|n\|_{L^\infty} + \|u\|_{L^\infty} + \|\nabla n\|_{L^\infty} + \|\nabla u\|_{L^\infty}) (\|\nabla^3 n\|_{L^2}^2 + \|\nabla^3 u\|_{L^2}^2 + \|\nabla^4 u\|_{L^2}^2) \\
& \quad + \|\nabla n\|_{L^3} \|\nabla^4 u\|_{L^2} \|\nabla^2 n\|_{L^6} + \|\nabla u\|_{L^3} \|\nabla^4 u\|_{L^2} \|\nabla^2 u\|_{L^6} \\
& \quad + \|\nabla^2 n\|_{L^3} (\|\nabla^3 n\|_{L^2} + \|\nabla^4 u\|_{L^2}) \|\nabla^2 u\|_{L^6}.
\end{aligned} \tag{2.8}$$

Summing up the above estimates, noting that $\delta > 0$ is small, we obtain that

$$\frac{d}{dt} \sum_{0 \leq k \leq 3} (\gamma \|\nabla^k n\|_{L^2}^2 + \|\nabla^k u\|_{L^2}^2) + C_1 \sum_{1 \leq k \leq 4} \|\nabla^k u\|_{L^2}^2 \leq C_2 \delta \sum_{1 \leq k \leq 3} \|\nabla^k n\|_{L^2}^2. \tag{2.9}$$

2.2. Recover the dissipation estimates for n . For $0 \leq k \leq 2$, applying ∇^k to the second equation in (2.1) and then multiplying by $\nabla^{k+1} n$. The key idea is to integrate by parts in the t -variable and to use the continuity equation. Thus integrating the results by parts for both the t - and x -variables, we obtain for $k = 0$ that

$$\begin{aligned}
& \frac{d}{dt} \int_{\mathbb{R}^3} u \cdot \nabla n dx + \gamma \bar{\rho} \int_{\mathbb{R}^3} |\nabla n|^2 dx \\
& \lesssim \|\nabla u\|_{L^2}^2 + \|\nabla n\|_{L^2} \|\nabla^2 u\|_{L^2} + (\|n\|_{L^\infty} + \|u\|_{L^\infty}) (\|\nabla n\|_{L^2}^2 + \|\nabla u\|_{L^2}^2),
\end{aligned} \tag{2.10}$$

for $k = 1$, we get

$$\begin{aligned}
& \frac{d}{dt} \int_{\mathbb{R}^3} \nabla u \cdot \nabla^2 n dx + \gamma \bar{\rho} \int_{\mathbb{R}^3} |\nabla^2 n|^2 dx \\
& \lesssim \|\nabla^2 u\|_{L^2}^2 + \|\nabla^2 n\|_{L^2} \|\nabla^3 u\|_{L^2} + (\|(n, u)\|_{L^\infty} + \|(\nabla n, \nabla u)\|_{L^\infty}) \\
& \quad \times (\|\nabla n\|_{L^2}^2 + \|\nabla^2 n\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2),
\end{aligned} \tag{2.11}$$

and for $k = 2$ we have

$$\begin{aligned}
& \frac{d}{dt} \int_{\mathbb{R}^3} \nabla^2 u \cdot \nabla^3 n dx + \gamma \bar{\rho} \int_{\mathbb{R}^3} |\nabla^3 n|^2 dx \\
& \lesssim \|\nabla^3 u\|_{L^2}^2 + \|\nabla^3 n\|_{L^2} \|\nabla^4 u\|_{L^2} + (\|(n, u)\|_{L^\infty} + \|(\nabla n, \nabla u)\|_{L^\infty}) \\
& \quad \times (\|\nabla^2 n\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2 + \|\nabla^3 n\|_{L^2}^2 + \|\nabla^3 u\|_{L^2}^2).
\end{aligned} \tag{2.12}$$

Plugging the above estimates, using the smallness of $\delta > 0$, we obtain that

$$\frac{d}{dt} \sum_{0 \leq k \leq 2} \int_{\mathbb{R}^3} \nabla^k u \cdot \nabla^{k+1} n dx + C_3 \sum_{1 \leq k \leq 3} \|\nabla^k n\|_{L^2}^2 \leq C_4 \sum_{1 \leq k \leq 4} \|\nabla^k u\|_{L^2}^2. \tag{2.13}$$

Proof of Theorem 2.1. Multiplying (2.13) by $\frac{2C_2\delta}{C_3}$, adding it with (2.9), with the help of smallness of $\delta > 0$, we deduce that there exists a constant $C_5 > 0$ such that

$$\begin{aligned}
& \frac{d}{dt} \left\{ \sum_{0 \leq k \leq 3} (\gamma \|\nabla^k n\|_{L^2}^2 + \|\nabla^k u\|_{L^2}^2) + \frac{2C_2\delta}{C_3} \sum_{0 \leq k \leq 2} \int_{\mathbb{R}^3} \nabla^k u \cdot \nabla^{k+1} n dx \right\} \\
& \quad + C_5 \left\{ \sum_{1 \leq k \leq 3} \|\nabla^k n\|_{L^2}^2 + \sum_{1 \leq k \leq 4} \|\nabla^k u\|_{L^2}^2 \right\} \leq 0.
\end{aligned} \tag{2.14}$$

Next, we define $\mathcal{E}(t)$ to be C_5^{-1} times the expression under the time derivative in (2.14). Then we may write (2.14) as

$$\frac{d}{dt}\mathcal{E}(t) + \|\nabla n(t)\|_{H^2}^2 + \|\nabla u(t)\|_{H^3}^2 \leq 0. \quad (2.15)$$

Observe that since δ is small, then there exists a constant $C_6 > 0$ such that

$$C_6^{-1} (\|n(t)\|_{H^3}^2 + \|u(t)\|_{H^3}^2) \leq \mathcal{E}(t) \leq C_6 (\|n(t)\|_{H^3}^2 + \|u(t)\|_{H^3}^2).$$

Then integrating (2.15) directly in time, we get

$$\begin{aligned} & \sup_{0 \leq t \leq T} (\|n(t)\|_{H^3}^2 + \|u(t)\|_{H^3}^2) + C_6 \int_0^T (\|\nabla n(\tau)\|_{H^2}^2 + \|\nabla u(\tau)\|_{H^3}^2) d\tau \\ & \leq C_6^2 (\|n_0\|_{H^3}^2 + \|u_0\|_{H^3}^2). \end{aligned}$$

Using a standard continuity argument along with classical local wellposedness theory, this closes the *a priori* assumption (2.3) if we assume $\|n_0\|_{H^3} + \|u_0\|_{H^3} \leq \delta_0$ is sufficiently small. We can then extend the solution globally in time and complete the proof of Theorem 2.1. $\square$

3. Spectral analysis of the semigroup. In this section, we consider the initial value problem for the linearized Navier-Stokes system

$$\begin{cases} \partial_t \tilde{n} + \operatorname{div} \tilde{m} = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \partial_t \tilde{m} + c^2 \nabla \tilde{n} - \bar{\mu} \Delta \tilde{m} - (\bar{\mu} + \bar{\nu}) \nabla \operatorname{div} \tilde{m} = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \lim_{|x| \rightarrow \infty} \tilde{n} = 0, \quad \lim_{|x| \rightarrow \infty} \tilde{m} = 0, & t \in \mathbb{R}^+, \\ (\tilde{n}, \tilde{m})|_{t=0} = (\rho_0 - \bar{\rho}, \rho_0 u_0), & x \in \mathbb{R}^3, \end{cases} \quad (3.1)$$

where $\bar{\mu} = \frac{\mu}{\rho}$, $\bar{\nu} = \frac{\nu}{\rho}$, $c = \sqrt{p'(\bar{\rho})}$.

In terms of the semigroup theory for evolutionary equations, the solution $(\tilde{n}, \tilde{m})$ of the linearized Navier-Stokes problem (3.1) can be expressed for $\tilde{U} = (\tilde{n}, \tilde{m})^t$ as

$$\tilde{U}_t = B\tilde{U}, \quad t \geq 0, \quad \tilde{U}(0) = \tilde{U}_0,$$

which gives rise to

$$\tilde{U}(t) = S(t)\tilde{U}_0 = e^{tB}\tilde{U}_0, \quad t \geq 0,$$

where B is defined as

$$B = \begin{pmatrix} 0 & -\operatorname{div} \\ -c^2 \nabla & \bar{\mu} \Delta + (\bar{\mu} + \bar{\nu}) \nabla \operatorname{div} \end{pmatrix}.$$

What left is to analyze the differential operator B in terms of its Fourier expression $A(\xi)$ and show the long time properties of the semigroup $S(t)$. Applying the Fourier transform to system (3.1), we have

$$\partial_t \widehat{\tilde{U}}(t, \xi) = A(\xi) \widehat{\tilde{U}}(t, \xi), \quad t \geq 0, \quad \widehat{\tilde{U}}(0, \xi) = \widehat{\tilde{U}}_0(\xi),$$

where $\xi = (\xi_1, \xi_2, \xi_3)^t$, and $A(\xi)$ is defined as

$$A(\xi) = \begin{pmatrix} 0 & -i\xi^t \\ -c^2 i\xi & -\bar{\mu}|\xi|^2 I_{3 \times 3} - (\bar{\mu} + \bar{\nu})\xi \otimes \xi \end{pmatrix}.$$

The eigenvalues of the matrix A can be computed by

$$\det(A(\xi) - \lambda I) = -(\lambda + \bar{\mu}|\xi|^2)^2 (\lambda^2 + (2\bar{\mu} + \bar{\nu})|\xi|^2 \lambda + c^2|\xi|^2) = 0,$$

which implies

$$\lambda_0 = -\bar{\mu}|\xi|^2 (\text{double}), \quad \lambda_1 = \lambda_1(|\xi|), \quad \lambda_2 = \lambda_2(|\xi|).$$

The semigroup e^{tA} is expressed as

$$e^{tA} = e^{\lambda_0 t} P_0 + e^{\lambda_1 t} P_1 + e^{\lambda_2 t} P_2,$$

where the project operators P_i can be computed as

$$P_i = \prod_{i \neq j} \frac{A(\xi) - \lambda_j I}{\lambda_i - \lambda_j}.$$

By a direct computation, we can verify the exact expression for the Fourier transform $\widehat{G}(t, \xi)$ of Green's function $G(t, x) = e^{tB}$ as

$$\widehat{G}(t, \xi) = e^{tA} = \begin{pmatrix} \frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} & -\frac{i\xi^t(e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \\ -\frac{c^2 i\xi(e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} & e^{-\lambda_0 t} \left(I - \frac{\xi \otimes \xi}{|\xi|^2} \right) + \frac{\xi \otimes \xi}{|\xi|^2} \frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} \end{pmatrix} = \begin{pmatrix} \widehat{N} \\ \widehat{M} \end{pmatrix}.$$

Indeed, we can make the following decomposition for $(\tilde{n}, \tilde{m}) = G * \tilde{U}_0$ as

$$\widehat{\tilde{n}} = \widehat{N} \cdot \widehat{\tilde{U}}_0 = (\widehat{\mathcal{N}} + \widehat{\mathfrak{N}}) \cdot \widehat{\tilde{U}}_0, \quad \widehat{\tilde{m}} = \widehat{M} \cdot \widehat{\tilde{U}}_0 = (\widehat{\mathcal{M}} + \widehat{\mathfrak{M}}) \cdot \widehat{\tilde{U}}_0,$$

where

$$\widehat{\mathcal{N}} = \begin{pmatrix} \frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} & 0 \end{pmatrix}, \quad \widehat{\mathfrak{N}} = \begin{pmatrix} 0 & -\frac{i\xi^t(e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \end{pmatrix},$$

$$\widehat{\mathcal{M}} = \begin{pmatrix} -\frac{c^2 i\xi(e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} & 0 \end{pmatrix}, \quad \widehat{\mathfrak{M}} = \begin{pmatrix} 0 & e^{-\lambda_0 t} \left(I - \frac{\xi \otimes \xi}{|\xi|^2} \right) + \frac{\xi \otimes \xi}{|\xi|^2} \frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} \end{pmatrix}.$$

We further decompose the Fourier transform $\widehat{N}$, $\widehat{M}$ into low frequency term and high frequency term below.

Define

$$\widehat{\mathcal{N}} = \widehat{\mathcal{N}}_1 + \widehat{\mathcal{N}}_2, \quad \widehat{\mathfrak{N}} = \widehat{\mathfrak{N}}_1 + \widehat{\mathfrak{N}}_2, \quad \widehat{\mathcal{M}} = \widehat{\mathcal{M}}_1 + \widehat{\mathcal{M}}_2, \quad \widehat{\mathfrak{M}} = \widehat{\mathfrak{M}}_1 + \widehat{\mathfrak{M}}_2,$$

where $(\cdot)_1 = \chi(\xi)(\cdot)$, $(\cdot)_2 = (1 - \chi(\xi))(\cdot)$, and $\chi(\xi)$ is a smooth cut off function such that

$$\chi(\xi) = \begin{cases} 1, & |\xi| \leq R, \\ 0, & |\xi| \geq R + 1. \end{cases}$$

Then we have the following decomposition for $(\tilde{n}, \tilde{m}) = G * \tilde{U}_0$ as

$$\begin{aligned} \widehat{\tilde{n}} &= \widehat{N} \cdot \widehat{\tilde{U}}_0 = \widehat{\mathcal{N}}_1 \cdot \widehat{\tilde{U}}_0 + \widehat{\mathcal{N}}_2 \cdot \widehat{\tilde{U}}_0 = (\widehat{\mathcal{N}}_1 + \widehat{\mathfrak{N}}_1) \cdot \widehat{\tilde{U}}_0 + (\widehat{\mathcal{N}}_2 + \widehat{\mathfrak{N}}_2) \cdot \widehat{\tilde{U}}_0, \\ \widehat{\tilde{m}} &= \widehat{M} \cdot \widehat{\tilde{U}}_0 = \widehat{\mathcal{M}}_1 \cdot \widehat{\tilde{U}}_0 + \widehat{\mathcal{M}}_2 \cdot \widehat{\tilde{U}}_0 = (\widehat{\mathcal{M}}_1 + \widehat{\mathfrak{M}}_1) \cdot \widehat{\tilde{U}}_0 + (\widehat{\mathcal{M}}_2 + \widehat{\mathfrak{M}}_2) \cdot \widehat{\tilde{U}}_0. \end{aligned} \tag{3.2}$$

To derive the long time decay rate of solution, we need to use accurate approximation to the Fourier transform $\widehat{G}(t, x)$ of Green's function for both lower frequency and high frequency. In terms of the definition of the eigenvalues, we are able to obtain that it holds for $|\xi| \leq \eta$ for some small positive constant η that

$$\begin{aligned} \lambda_1 &= -\frac{2\bar{\mu} + \bar{\nu}}{2} |\xi|^2 + \frac{i}{2} \sqrt{4c^2 |\xi|^2 - (2\bar{\mu} + \bar{\nu})^2 |\xi|^4} = a + bi, \\ \lambda_2 &= -\frac{2\bar{\mu} + \bar{\nu}}{2} |\xi|^2 - \frac{i}{2} \sqrt{4c^2 |\xi|^2 - (2\bar{\mu} + \bar{\nu})^2 |\xi|^4} = a - bi, \end{aligned} \tag{3.3}$$

and we have

$$\begin{aligned} \frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} &= e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[\cos(bt) + \frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 \frac{\sin(bt)}{b} \right] \\ &\sim O(1) e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t}, \quad |\xi| \leq \eta, \end{aligned}$$

$$\begin{aligned} \frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} &= e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[\cos(bt) - \frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 \frac{\sin(bt)}{b} \right] \\ &\sim O(1) e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t}, \quad |\xi| \leq \eta, \end{aligned}$$

$$\frac{e^{\lambda_1 t} - e^{\lambda_2 t}}{\lambda_1 - \lambda_2} = e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \frac{\sin(bt)}{b} \sim O(1) \frac{1}{|\xi|} e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t}, \quad |\xi| \leq \eta,$$

where

$$b = \frac{1}{2} \sqrt{4c^2|\xi|^2 - (2\bar{\mu} + \bar{\nu})^2|\xi|^4} \sim c|\xi| + O(|\xi|^3), \quad |\xi| \leq \eta.$$

For the high frequency $|\xi| \geq \eta$, we are also able to obtain that it holds for $|\xi| \geq \eta$ that

$$\begin{aligned} \lambda_1 &= -\frac{2\bar{\mu} + \bar{\nu}}{2}|\xi|^2 - \frac{1}{2}\sqrt{(2\bar{\mu} + \bar{\nu})^2|\xi|^4 - 4c^2|\xi|^2} = a - b, \\ \lambda_2 &= -\frac{2\bar{\mu} + \bar{\nu}}{2}|\xi|^2 + \frac{1}{2}\sqrt{(2\bar{\mu} + \bar{\nu})^2|\xi|^4 - 4c^2|\xi|^2} = a + b, \end{aligned} \tag{3.4}$$

and we have

$$\frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} = \frac{1}{2} e^{(a+b)t} [1 + e^{-2bt}] - \frac{a}{2b} e^{(a+b)t} [1 - e^{-2bt}] \sim O(1) e^{-R_0 t}, \quad |\xi| \geq \eta,$$

$$\frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} = \frac{a+b}{2b} e^{(a+b)t} [1 - e^{-2bt}] + e^{(a-b)t} \sim O(1) e^{-R_0 t}, \quad |\xi| \geq \eta,$$

$$\frac{e^{\lambda_1 t} - e^{\lambda_2 t}}{\lambda_1 - \lambda_2} = \frac{1}{2b} e^{(a+b)t} [1 - e^{-2bt}] \sim O(1) \frac{1}{|\xi|^2} e^{-R_0 t}, \quad |\xi| \geq \eta,$$

where

$$b = \frac{1}{2} \sqrt{(2\bar{\mu} + \bar{\nu})^2|\xi|^4 - 4c^2|\xi|^2} \sim \frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 - \frac{2c^2}{2\bar{\mu} + \bar{\nu}} + O(|\xi|^{-2}), \quad |\xi| \geq \eta.$$

Here R_0, η are some fixed positive constants.

4. Time decay rate for linear semigroup. In this section, we apply the spectral analysis to the semigroup for the linearized Navier-Stokes system. We will establish the L^2 and L^p ($2 \leq p \leq \infty$) time decay rate of the global solutions for the linearized Navier-Stokes system.

4.1. L^2 -time decay rate. With the help of the formula for Green's function in Fourier space and the asymptotic analysis on its elements, we are able to establish the L^2 time decay rate. Indeed, we have the L^2 -time decay rate of the global strong solution to the problem for the linearized Navier-Stokes system as follows.

Proposition 4.1. *Let $U_0 = (n_0, m_0) \in L^1(\mathbb{R}^3) \cap H^l(\mathbb{R}^3)$ with $l \geq 3$, then $(\tilde{n}, \tilde{m})$ solves the linearized Navier-Stokes system (3.1) and satisfies for $0 \leq k \leq l$ that*

$$\|\nabla^k(\tilde{n}, \tilde{m})(t)\|_{L^2(\mathbb{R}^3)} \leq C(1+t)^{-\frac{3}{4}-\frac{k}{2}} (\|U_0\|_{L^1(\mathbb{R}^3)} + \|\nabla^k U_0\|_{L^2(\mathbb{R}^3)}),$$

where C is a positive constant independent of time.

Proof. A straightforward computation together with the formula of the Green's function $\widehat{G}(t, \xi)$ gives

$$\begin{aligned}\widehat{n}(t, \xi) &= \frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} \widehat{n}_0 - \frac{i\xi \cdot \widehat{m}_0 (e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \\ &\sim \begin{cases} O(1) e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} (|\widehat{n}_0| + |\widehat{m}_0|), & |\xi| \leq \eta, \\ O(1) e^{-R_0 t} (|\widehat{n}_0| + |\widehat{m}_0|), & |\xi| \geq \eta, \end{cases} \\ \widehat{m}(t, \xi) &= -\frac{e^2 i\xi (e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \widehat{n}_0 + e^{-\lambda_0 t} \widehat{m}_0 + \left(\frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} - e^{-\lambda_0 t} \right) \frac{\xi(\xi \cdot \widehat{m}_0)}{|\xi|^2} \\ &\sim \begin{cases} O(1) e^{-\bar{\mu}|\xi|^2 t} (|\widehat{n}_0| + |\widehat{m}_0|), & |\xi| \leq \eta, \\ O(1) e^{-R_0 t} (|\widehat{n}_0| + |\widehat{m}_0|), & |\xi| \geq \eta, \end{cases}\end{aligned}$$

here and below, R_0, η are some fixed positive constants. Therefore, we have the L^2 -decay rate for $(\widehat{n}, \widehat{m})$ as

$$\begin{aligned}\|(\widehat{n}, \widehat{m})(t)\|_{L^2(\mathbb{R}^3)}^2 &= \int_{|\xi| \leq \eta} |(\widehat{n}, \widehat{m})(t, \xi)|^2 d\xi + \int_{|\xi| \geq \eta} |(\widehat{n}, \widehat{m})(t, \xi)|^2 d\xi \\ &\lesssim \int_{|\xi| \leq \eta} e^{-2\bar{\mu}|\xi|^2 t} (|\widehat{n}_0|^2 + |\widehat{m}_0|^2) d\xi + \int_{|\xi| \geq \eta} e^{-2R_0 t} (|\widehat{n}_0|^2 + |\widehat{m}_0|^2) d\xi \\ &\lesssim (1+t)^{-\frac{3}{2}} \|(n_0, m_0)\|_{L^1(\mathbb{R}^3) \cap L^2(\mathbb{R}^3)}^2.\end{aligned}$$

And the L^2 -decay rate on the derivatives of $(\widehat{n}, \widehat{m})$ as

$$\begin{aligned}\|(\nabla^k \widehat{n}, \nabla^k \widehat{m})(t)\|_{L^2(\mathbb{R}^3)}^2 &= \int_{|\xi| \leq \eta} |\xi|^{2k} |(\widehat{n}, \widehat{m})(t, \xi)|^2 d\xi + \int_{|\xi| \geq \eta} |\xi|^{2k} |(\widehat{n}, \widehat{m})(t, \xi)|^2 d\xi \\ &\lesssim \int_{|\xi| \leq \eta} e^{-2\bar{\mu}|\xi|^2 t} |\xi|^{2k} (|\widehat{n}_0|^2 + |\widehat{m}_0|^2) d\xi + \int_{|\xi| \geq \eta} e^{-2R_0 t} |\xi|^{2k} (|\widehat{n}_0|^2 + |\widehat{m}_0|^2) d\xi \\ &\lesssim (1+t)^{-\frac{3}{2}-k} (\|(n_0, m_0)\|_{L^1(\mathbb{R}^3)}^2 + \|(\nabla^k n_0, \nabla^k m_0)\|_{L^2(\mathbb{R}^3)}^2).\end{aligned}$$

The proof of the Proposition 4.1 is completed. $\square$

It should be noted that the L^2 -time decay rates derived above are optimal.

Proposition 4.2. *Let $U_0 = (n_0, m_0) \in L^1(\mathbb{R}^3) \cap H^l(\mathbb{R}^3)$ with $l \geq 3$, assume that $M_n = \int_{\mathbb{R}^3} n_0(x) dx$ and $M_m = \int_{\mathbb{R}^3} m_0(x) dx$ satisfies that M_n, M_m are at least not all zeros, then the solution $(\widetilde{n}, \widetilde{m})$ of the linearized Navier-Stokes system (3.1) given by Proposition 4.1 satisfies for $0 \leq k \leq l$*

$$\begin{aligned}C^{-1}(1+t)^{-\frac{3}{4}-\frac{k}{2}} &\leq \|\nabla^k \widetilde{n}(t)\|_{L^2(\mathbb{R}^3)} \leq C(1+t)^{-\frac{3}{4}-\frac{k}{2}}, \\ C^{-1}(1+t)^{-\frac{3}{4}-\frac{k}{2}} &\leq \|\nabla^k \widetilde{m}(t)\|_{L^2(\mathbb{R}^3)} \leq C(1+t)^{-\frac{3}{4}-\frac{k}{2}},\end{aligned}$$

where C is a positive constant independent of time.

Proof. We only show the case of $k = 0$ for simplicity, the argument applies to the other orders of derivatives. From the formula of the Green's function $\widehat{G}(t, \xi)$, we

deduce that

$$\begin{aligned}
& \widehat{n}(t, \xi) \\
&= \frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} \widehat{n}_0 - \frac{i\xi \cdot \widehat{m}_0 (e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \\
&= e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[\cos(bt) \widehat{n}_0 - i\xi \cdot \widehat{m}_0 \frac{\sin(bt)}{b} \right] + e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 \frac{\sin(bt)}{b} \widehat{n}_0 \right] \\
&= T_1 + T_2, \quad \text{for } |\xi| \leq \eta,
\end{aligned}$$

$$\begin{aligned}
& \widehat{m}(t, \xi) \\
&= -\frac{c^2 i\xi (e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \widehat{n}_0 + e^{-\lambda_0 t} \widehat{m}_0 + \left(\frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} - e^{-\lambda_0 t} \right) \frac{\xi(\xi \cdot \widehat{m}_0)}{|\xi|^2} \\
&= \left[e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[\cos(bt) \frac{\xi(\xi \cdot \widehat{m}_0)}{|\xi|^2} - c^2 i\xi \frac{\sin(bt)}{b} \widehat{n}_0 \right] + e^{-\bar{\mu}|\xi|^2 t} \left[\widehat{m}_0 - \frac{\xi(\xi \cdot \widehat{m}_0)}{|\xi|^2} \right] \right] \\
&\quad - e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 \frac{\sin(bt)}{b} \frac{\xi(\xi \cdot \widehat{m}_0)}{|\xi|^2} \right] \\
&= S_1 + S_2, \quad \text{for } |\xi| \leq \eta,
\end{aligned}$$

here and below, η is a sufficiently small but fixed constant.

It is easy to check that

$$\begin{aligned}
\|\widehat{n}(t, \xi)\|_{L^2}^2 &= \int_{|\xi| \leq \eta} |\widehat{n}(t, \xi)|^2 d\xi + \int_{|\xi| \geq \eta} |\widehat{n}(t, \xi)|^2 d\xi \\
&\geq \int_{|\xi| \leq \eta} |T_1 + T_2|^2 d\xi \geq \int_{|\xi| \leq \eta} \frac{1}{2} |T_1|^2 - |T_2|^2 d\xi.
\end{aligned} \tag{4.1}$$

We then calculate that

$$\begin{aligned}
\int_{|\xi| \leq \eta} |T_2|^2 d\xi &\lesssim \|\widehat{n}_0\|_{L^\infty}^2 \int_{|\xi| \leq \eta} e^{-(2\bar{\mu} + \bar{\nu})|\xi|^2 t} |\xi|^4 \left(\frac{\sin(bt)}{b} \right)^2 d\xi \\
&\lesssim \|\widehat{n}_0\|_{L^\infty}^2 \int_{|\xi| \leq \eta} e^{-(2\bar{\mu} + \bar{\nu})|\xi|^2 t} |\xi|^2 d\xi \lesssim (1+t)^{-\frac{5}{2}} \|n_0\|_{L^1}^2.
\end{aligned} \tag{4.2}$$

Since $n_0(x) \in L^1$ implies $\widehat{n}_0(\xi) \in C(\mathbb{R}^3)$. If $\widehat{n}_0(0) = \int_{\mathbb{R}^3} n_0(x) dx \neq 0$, we deduce that $\widehat{n}_0(\xi) \neq 0$ for $|\xi| \leq \eta$ when η is sufficiently small. One finds that, when $M_n \neq 0$,

$$|\widehat{n}_0(\xi)|^2 \geq \frac{1}{C} \left| \int_{\mathbb{R}^3} n_0(x) dx \right|^2 \geq \frac{M_n^2}{C}, \quad \text{for } |\xi| \leq \eta.$$

For $\widehat{m}_0$, a similar argument yields that, when $M_m \neq 0$, we have

$$\frac{|\xi \cdot \widehat{m}_0(\xi)|^2}{|\xi|^2} \geq \frac{|\xi \cdot M_m|^2}{C|\xi|^2}, \quad \text{for } |\xi| \leq \eta.$$

When $M_n \neq 0$, $M_m \neq 0$, with the help of the above analysis, using $b \sim c|\xi| + O(|\xi|^3)$ for $|\xi| \leq \eta$, we obtain that

$$\begin{aligned}
& \int_{|\xi| \leq \eta} |T_1|^2 d\xi \\
&\geq \frac{M_n^2}{C} \int_{|\xi| \leq \eta} e^{-(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \cos^2(bt) d\xi + \frac{1}{C} \int_{|\xi| \leq \eta} \frac{|\xi \cdot M_m|^2}{b^2} e^{-(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \sin^2(bt) d\xi
\end{aligned}$$

$$\begin{aligned}
&\geq \frac{\min\{M_n^2, \frac{M_m^2}{3c^2}\}}{C} \int_{|\xi| \leq \eta} e^{-(2\bar{\mu}+\bar{\nu})|\xi|^2 t} (\cos^2(bt) + \sin^2(bt)) d\xi \\
&\geq C_1 \int_{|\xi| \leq \eta} e^{-(2\bar{\mu}+\bar{\nu})|\xi|^2 t} d\xi \\
&\geq C^{-1}(1+t)^{-\frac{3}{2}}.
\end{aligned} \tag{4.3}$$

If $M_n \neq 0$, $M_m = 0$, and by the continuity of $\widehat{m}_0$ near $\xi = 0$, there exists a small enough constant ϵ such that $\epsilon \rightarrow 0$ as $\xi \rightarrow 0$, and

$$|\widehat{m}_0(\xi)|^2 < \epsilon, \quad \text{for } |\xi| \leq \eta.$$

We thus use the help of spherical coordinates and the change of variables $r = |\xi|\sqrt{t}$ to obtain that

$$\begin{aligned}
&\int_{|\xi| \leq \eta} |T_1|^2 d\xi \\
&\geq \frac{M_n^2}{C} \int_{|\xi| \leq \eta} e^{-(2\bar{\mu}+\bar{\nu})|\xi|^2 t} \cos^2(bt) d\xi - \frac{\epsilon}{Cc^2} \int_{|\xi| \leq \eta} e^{-(2\bar{\mu}+\bar{\nu})|\xi|^2 t} \sin^2(bt) d\xi \\
&\geq \frac{M_n^2}{C} t^{-\frac{3}{2}} \int_0^{\eta\sqrt{t}} e^{-(2\bar{\mu}+\bar{\nu})r^2} \cos^2(cr\sqrt{t}) r^2 dr - \frac{\epsilon}{Cc^2} t^{-\frac{3}{2}} \int_0^{\eta\sqrt{t}} e^{-(2\bar{\mu}+\bar{\nu})r^2} \sin^2(cr\sqrt{t}) r^2 dr \\
&\geq \frac{M_n^2}{C} t^{-\frac{3}{2}} \sum_{k=0}^{\lfloor \frac{c\eta t}{\pi} \rfloor - 1} \int_{\frac{k\pi}{c\sqrt{t}}}^{\frac{(k+1)\pi}{c\sqrt{t}}} e^{-(2\bar{\mu}+\bar{\nu})r^2} \cos^2(cr\sqrt{t}) r^2 dr - \frac{\epsilon}{Cc^2} (1+t)^{-\frac{3}{2}} \\
&\geq \frac{M_n^2}{2C} t^{-\frac{3}{2}} \sum_{k=0}^{\lfloor \frac{c\eta t}{\pi} \rfloor - 1} \int_{\frac{k\pi}{c\sqrt{t}}}^{\frac{(k+1)\pi}{c\sqrt{t}}} e^{-(2\bar{\mu}+\bar{\nu})r^2} r^2 dr - \frac{\epsilon}{Cc^2} (1+t)^{-\frac{3}{2}} \\
&\geq C_1^{-1} (1+t)^{-\frac{3}{2}} - C_2^{-1} \epsilon (1+t)^{-\frac{3}{2}} \\
&\geq C^{-1} (1+t)^{-\frac{3}{2}}.
\end{aligned} \tag{4.4}$$

In the case of $M_n = 0$, $M_m \neq 0$, we can use a similar argument to obtain that

$$\begin{aligned}
&\int_{|\xi| \leq \eta} |T_1|^2 d\xi \\
&\geq -\frac{\epsilon}{C} \int_{|\xi| \leq \eta} e^{-(2\bar{\mu}+\bar{\nu})|\xi|^2 t} \cos^2(bt) d\xi + \frac{M_m^2}{3Cc^2} \int_{|\xi| \leq \eta} e^{-(2\bar{\mu}+\bar{\nu})|\xi|^2 t} \sin^2(bt) d\xi \\
&\geq C^{-1} (1+t)^{-\frac{3}{2}}.
\end{aligned} \tag{4.5}$$

Combining the above estimates (4.1), (4.2), (4.3), (4.4) and (4.5), we obtain the lower bound of the time decay rate for $\widetilde{n}(t, x)$ as

$$\|\widetilde{n}(t, x)\|_{L^2}^2 = \|\widehat{\widetilde{n}}(t, \xi)\|_{L^2}^2 \geq C^{-1} (1+t)^{-\frac{3}{2}}.$$

The lower bound of the time decay rate for $\widetilde{m}(t, x)$ can be shown in a similar fashion. It is not difficult to derive that

$$\|\widehat{\widetilde{m}}(t, \xi)\|_{L^2}^2 \geq \int_{|\xi| \leq \eta} \frac{1}{2} |S_1|^2 - |S_2|^2 d\xi, \tag{4.6}$$

then we find that

$$\int_{|\xi| \leq \eta} |S_2|^2 d\xi \lesssim (1+t)^{-\frac{5}{2}} \|m_0\|_{L^1}^2. \tag{4.7}$$

We then calculate that

$$\begin{aligned}
& \int_{|\xi| \leq \eta} |S_1|^2 d\xi \\
& \geq \left\{ \frac{c^4 M_n^2}{C} \int_{|\xi| \leq \eta} \frac{|\xi|^2}{b^2} e^{-(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \sin^2(bt) d\xi \right. \\
& \quad \left. + \frac{1}{C} \int_{|\xi| \leq \eta} \frac{|\xi \cdot M_m|^2}{|\xi|^2} e^{-(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \cos^2(bt) d\xi \right\} \\
& \quad + \left\{ \int_{|\xi| \leq \eta} e^{-\frac{1}{2}(4\bar{\mu} + \bar{\nu})|\xi|^2 t} \cos(bt) \frac{\xi(\xi \cdot \hat{m}_0)}{|\xi|^2} \left(\hat{m}_0 - \frac{\xi(\xi \cdot \hat{m}_0)}{|\xi|^2} \right) d\xi \right\} \\
& = J_1 + J_2.
\end{aligned}$$

A direct computation gives rise to

$$J_1 \geq C^{-1}(1+t)^{-\frac{3}{2}}, \quad J_2 = 0. \quad (4.8)$$

Combining the above estimates (4.6), (4.7) and (4.8), we obtain the lower bound of the time decay rate for $\tilde{m}(t, x)$ as

$$\|\tilde{m}(t, x)\|_{L^2}^2 = \|\hat{\tilde{m}}(t, \xi)\|_{L^2}^2 \geq C^{-1}(1+t)^{-\frac{3}{2}}.$$

Then the proof of Proposition 4.2 is completed. $\square$

4.2. L^p -time decay rate. In this subsection, we establish the following L^p -time decay rate of the global strong solution to the linearized Navier-Stokes system with $p \in [2, +\infty]$.

Proposition 4.3. *Let $U_0 = (n_0, m_0) \in L^1(\mathbb{R}^3) \cap W^{l,p}(\mathbb{R}^3)$ with $l \geq 3$, then $(\tilde{n}, \tilde{m})$ solves the linearized Navier-Stokes system (3.1) and satisfies for $0 \leq k \leq l$ and $p \in [2, +\infty]$ that*

$$\|\nabla^k(\tilde{n}, \tilde{m})(t)\|_{L^p(\mathbb{R}^3)} \leq C(1+t)^{-\frac{3}{2}(1-\frac{1}{p})-\frac{k}{2}} (\|U_0\|_{L^1(\mathbb{R}^3)} + \|\nabla^k U_0\|_{L^p(\mathbb{R}^3)}),$$

where C is a positive constant independent of time.

To prove Proposition 4.3, the following two lemmas in [6] are helpful.

Lemma 4.1. *Let $n \geq 1$ and assume that $\hat{f}(\xi) \in L^\infty \cap C^{n+1}(\mathbb{R}^n \setminus \{0\})$, with*

$$|\nabla_\xi^\alpha \hat{f}(\xi)| \leq C' \begin{cases} |\xi|^{-|\alpha|+\sigma_1}, & |\xi| \leq R, |\alpha| = n, \\ |\xi|^{-|\alpha|-\sigma_2}, & |\xi| \geq R, |\alpha| = n-1, n, n+1, \end{cases}$$

where $\sigma_1, \sigma_2 > 0$ and $n > 2 - 2\sigma_2$. Then $\hat{f}(\xi)$ is continuous at 0 and ∞ , and

$$f = m_1 + m_2 \delta,$$

where $m_1 \in L^1(\mathbb{R}^n)$ satisfies $\|m_1\|_{L^1(\mathbb{R}^n)} \leq C(C')$, m_2 is the constant

$$m_2 = (2\pi)^{-\frac{n}{2}} \lim_{|\xi| \rightarrow \infty} \hat{f}(\xi),$$

and δ is the Dirac distribution. In particular, $\hat{f}(\xi)$ is a strong L^p multiplier, $1 \leq p \leq \infty$, in the sense that, for any $g \in L^p$,

$$\|f * g\|_{L^p} \leq C \|g\|_{L^p}, \quad 1 \leq p \leq \infty,$$

where C depends only on $|m_2| \leq \|\hat{f}\|_{L^\infty}$ and the constant C' above.

Lemma 4.2. *Let $\hat{g}(t, \xi) = \hat{K}(t, \xi)\hat{f}(\xi)$, where $\hat{K}(t, \xi) = e^{-\vartheta|\xi|^2 t}$, $\hat{f}(\xi) \in L^\infty \cap C^{n+1}(\mathbb{R}^n)$, and*

$$|\nabla_\xi^\beta \hat{f}(\xi)| \leq C' |\xi|^{-|\beta|}, \quad |\beta| \leq n+1.$$

Then $\nabla_x^\alpha g(t, \cdot) \in L^p$ for $t > 0$, and for all α , $1 \leq p \leq \infty$, we have

$$\|\nabla_x^\alpha g(t, \cdot)\|_{L^p} \leq C(|\alpha|) t^{-\frac{n}{2}(1-\frac{1}{p})-\frac{|\alpha|}{2}}.$$

In particular, $\widehat{\nabla_x^\alpha g(t, x)} = (i\xi)^\alpha \hat{g}(t, \xi)$ is a strong L^p multiplier, with norm bounded by $C(|\alpha|, \vartheta) C' t^{-\frac{|\alpha|}{2}}$, where the constant $C(|\alpha|, \vartheta)$ depends only on $|\alpha|$ and ϑ .

Now let us turn to the proof of Proposition 4.3.

Proof of Proposition 4.3. We first analyze above higher frequency terms denoted by $(\cdot)_2$. Recall that

$$\lambda_1 = -(2\bar{\mu} + \bar{\nu})|\xi|^2 + \frac{2c^2}{2\bar{\mu} + \bar{\nu}} + O(|\xi|^{-2}), \quad \lambda_2 = -\frac{2c^2}{2\bar{\mu} + \bar{\nu}} + O(|\xi|^{-2}), \quad |\xi| \geq \eta.$$

We shall prove that the higher frequency terms are L^p Fourier multipliers with an exponential time decay coefficient $Ce^{-c_1 t}$ for some constants $c_1 > 0$. For simplicity, we only show that $\widehat{\mathcal{N}}_2$ is an L^p Fourier multiplier at higher frequency as follows. It holds

$$\frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} = e^{\lambda_2 t} + \frac{\lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} - \frac{\lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2}.$$

By a direct computation, it is easy to verify

$$|\nabla_\xi^k \lambda_2| \lesssim |\xi|^{-2-k}, \quad |\xi| \geq \eta,$$

which gives rise to

$$\left| \nabla_\xi^k \left[(1 - \chi(\cdot)) e^{\lambda_2 t} \right] \right|, \left| \nabla_\xi^k \left[(1 - \chi(\cdot)) \frac{\lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} \right] \right| \lesssim \begin{cases} 0, & |\xi| \leq R, \\ e^{-c_1 t} |\xi|^{-2-k}, & |\xi| \geq R, \end{cases}$$

here and below, $R > 0$ is a given constant. Thus, from Lemma 4.1 it follows that the inverse Fourier transform of the term $(1 - \chi(\cdot)) \left(e^{\lambda_2 t} + \frac{\lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2} \right)$ is an L^p multiplier with the coefficient $Ce^{-c_1 t}$. The other part of $\widehat{\mathcal{N}}_2$ at higher frequency can be written as

$$(1 - \chi(\cdot)) \frac{\lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2} \sim e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t} \left[(1 - \chi(\cdot)) \frac{e^{(-\lambda_2 - \frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2)t}}{\lambda_1 - \lambda_2} \right].$$

We can regard $e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t}$ as the function $K(t, \xi)$ of Lemma 4.2, and the rest term satisfies the condition. Thus, the inverse Fourier transform of $(1 - \chi(\cdot)) \frac{\lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2}$ is also an L^p multiplier with the coefficient $Ce^{-c_1 t}$. These facts imply that $\widehat{\mathcal{N}}_2$ at higher frequency is an L^p multiplier with the coefficient $Ce^{-c_1 t}$. Applying the similar analysis to the terms $\widehat{\mathfrak{N}}_2$, $\widehat{\mathcal{M}}_2$, and $\widehat{\mathfrak{M}}_2$, we can show that their inverse Fourier transform are all L^p multiplier with the constant coefficient $Ce^{-c_1 t}$. Then

$$\|(\nabla_x^k (\mathcal{N}_2 * f), \nabla_x^k (\mathfrak{N}_2 * f), \nabla_x^k (\mathcal{M}_2 * f), \nabla_x^k (\mathfrak{M}_2 * f))(t)\|_{L^p} \leq Ce^{-c_1 t} \|\nabla_x^k f\|_{L^p}, \quad (4.9)$$

for all integer $k \geq 0$, and $p \in [2, \infty]$.

We also need to deal with the corresponding lower frequency terms denoted by $\widehat{(\cdot)}_1$. Recall that

$$\frac{\lambda_1 e^{\lambda_2 t} - \lambda_2 e^{\lambda_1 t}}{\lambda_1 - \lambda_2}, \frac{\lambda_1 e^{\lambda_1 t} - \lambda_2 e^{\lambda_2 t}}{\lambda_1 - \lambda_2}, \frac{|\xi|(e^{\lambda_1 t} - e^{\lambda_2 t})}{\lambda_1 - \lambda_2} \sim O(1)e^{-\frac{1}{2}(2\bar{\mu} + \bar{\nu})|\xi|^2 t}, \quad |\xi| \leq \eta,$$

which imply that for $|\xi| \leq \eta$ that

$$\begin{aligned} |\widehat{\mathcal{N}}_1| &\sim O(1)e^{-c_2|\xi|^2 t}, & |\widehat{\mathfrak{N}}_1| &\sim O(1)e^{-c_2|\xi|^2 t}, \\ |\widehat{\mathcal{M}}_1| &\sim O(1)e^{-c_2|\xi|^2 t}, & |\widehat{\mathfrak{M}}_1| &\sim O(1)e^{-c_2|\xi|^2 t}, \end{aligned}$$

for some constants $c_2 > 0$. Thus, by Hausdorff-Young's inequality with $p \in [2, +\infty]$, we can obtain

$$\begin{aligned} \|(\nabla^k \mathcal{N}_1, \nabla^k \mathfrak{N}_1, \nabla^k \mathcal{M}_1, \nabla^k \mathfrak{M}_1)(t)\|_{L^p} &\leq C \left(\int_{|\xi| \leq \eta} |\xi|^k e^{-c_2|\xi|^2 t} d\xi \right)^{\frac{1}{q}} \\ &\leq C(1+t)^{-\frac{3}{2}(1-\frac{1}{p})-\frac{k}{2}}. \end{aligned} \quad (4.10)$$

Combining (4.9) and (4.10), we finally have for $t > 0$ that

$$\begin{aligned} \|(\nabla^k(N * f), \nabla^k(M * f))(t)\|_{L^p} &= \|(\nabla^k((N_1 + N_2) * f), \nabla^k((M_1 + M_2) * f))(t)\|_{L^p} \\ &\leq C(1+t)^{-\frac{3}{2}(1-\frac{1}{p})-\frac{k}{2}} \|f\|_{L^1} + Ce^{-c_1 t} \|\nabla^k f\|_{L^p} \\ &\leq C(1+t)^{-\frac{3}{2}(1-\frac{1}{p})-\frac{k}{2}} (\|f\|_{L^1} + \|\nabla^k f\|_{L^p}). \end{aligned}$$

The proof of Proposition 4.3 is completed. $\square$

5. L^2 -time decay rate for nonlinear system. We are ready to prove Theorem 1.1 on the *sharp time decay rate* of the global solution to the initial value problem for the nonlinear Navier-Stokes system.

In what follows, we will set $n_h = n - \tilde{n}$ and $m_h = m - \tilde{m}$, then we have

$$\begin{cases} \partial_t n_h + \operatorname{div} m_h = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \partial_t m_h + c^2 \nabla n_h - \bar{\mu} \Delta m_h - (\bar{\mu} + \bar{\nu}) \nabla \operatorname{div} m_h = F, & (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \lim_{|x| \rightarrow \infty} n_h = 0, \quad \lim_{|x| \rightarrow \infty} m_h = 0, & t \in \mathbb{R}^+, \\ (n_h, m_h)|_{t=0} = (0, 0), & x \in \mathbb{R}^3, \end{cases} \quad (5.1)$$

where $\bar{\mu} = \frac{\mu}{\bar{\rho}}$, $\bar{\nu} = \frac{\nu}{\bar{\rho}}$, $c = \sqrt{p'(\bar{\rho})}$, and

$$\begin{aligned} F = & -\operatorname{div} \left\{ \frac{(m_h + \tilde{m}) \otimes (m_h + \tilde{m})}{n_h + \tilde{n} + \bar{\rho}} + \bar{\mu} \nabla \left(\frac{(n_h + \tilde{n})(m_h + \tilde{m})}{n_h + \tilde{n} + \bar{\rho}} \right) \right\} \\ & - \nabla \left\{ (\bar{\mu} + \bar{\nu}) \operatorname{div} \left(\frac{(n_h + \tilde{n})(m_h + \tilde{m})}{n_h + \tilde{n} + \bar{\rho}} \right) + (p(n_h + \tilde{n} + \bar{\rho}) - p(\bar{\rho}) - c^2(n_h + \tilde{n})) \right\}. \end{aligned}$$

Denote $U_h = (n_h, m_h)^t$, we have the equivalent form of system (5.1) in vector form

$$\partial_t U_h = B U_h + H, \quad t \geq 0, \quad U_h(0) = 0,$$

where the nonlinear term $H(\tilde{U}, U_h) = (0, F(\tilde{U}, U_h))^t$. Thus, we can represent the solution in term of the semigroup

$$U_h(t) = S(t) * U_h(0) + \int_0^t S(t-\tau) * H(\tilde{U}, U_h)(\tau) d\tau,$$

which (n_h, m_h) can be decomposed as

$$n_h = N * U_h(0) + \int_0^t \mathfrak{N}(t - \tau) * H(\tau) d\tau, \quad (5.2)$$

$$m_h = M * U_h(0) + \int_0^t \mathfrak{M}(t - \tau) * H(\tau) d\tau. \quad (5.3)$$

Furthermore, in view of the above definition for $\widehat{\mathfrak{N}}(\xi)$ and $\widehat{\mathfrak{M}}(\xi)$, it is easy to verify for some constants $c_3 > 0$, $c_4 > 0$, $R_0 > 0$, we discover that

$$|\widehat{\mathfrak{N}}(\xi)| \sim O(1)e^{-c_3|\xi|^2 t}, \quad |\widehat{\mathfrak{M}}(\xi)| \sim O(1)e^{-c_3|\xi|^2 t}, \quad |\xi| \leq \eta,$$

$$|\widehat{\mathfrak{N}}(\xi)| \sim O(1)\frac{1}{|\xi|}e^{-R_0 t}, \quad |\widehat{\mathfrak{M}}(\xi)| \sim O(1)\frac{1}{|\xi|^2}e^{-R_0 t} + O(1)e^{-c_4|\xi|^2 t}, \quad |\xi| \geq \eta.$$

Thus, applying a similar argument as in the proof of Proposition 4.1, we have

$$\|(\nabla^k \mathfrak{N} * H, \nabla^k \mathfrak{M} * H)(t)\|_{L^2} \leq C(1+t)^{-\frac{3}{2}(\frac{1}{q}-\frac{1}{2})-\frac{1}{2}-\frac{k}{2}}(\|Q\|_{L^q} + \|\nabla^{k+1}Q\|_{L^2}), \quad q = 1, 2, \quad (5.4)$$

$$\|(\nabla^k \mathfrak{N} * H, \nabla^k \mathfrak{M} * H)(t)\|_{L^2} \leq C(1+t)^{-\frac{3}{2}(\frac{1}{q}-\frac{1}{2})-\frac{1}{2}-\frac{k}{2}}(\|Q\|_{L^q} + \|\nabla^k Q\|_{L^2}), \quad q = 1, 2, \quad (5.5)$$

$$\|\nabla^k \mathfrak{M} * H(t)\|_{L^2} \leq C(1+t)^{-\frac{3}{2}(\frac{1}{q}-\frac{1}{2})-\frac{1}{2}-\frac{k}{2}}(\|Q\|_{L^q} + \|\nabla^{k-1}Q\|_{L^2}), \quad q = 1, 2, \quad (5.6)$$

for any non-negative integer k and

$$\begin{aligned} Q = & \left| \frac{(m_h + \tilde{m}) \otimes (m_h + \tilde{m})}{n_h + \tilde{n} + \bar{\rho}} + \bar{\mu} \nabla \left(\frac{(n_h + \tilde{n})(m_h + \tilde{m})}{n_h + \tilde{n} + \bar{\rho}} \right) \right| \\ & + \left| (\bar{\mu} + \bar{\nu}) \operatorname{div} \left(\frac{(n_h + \tilde{n})(m_h + \tilde{m})}{n_h + \tilde{n} + \bar{\rho}} \right) + (p(n_h + \tilde{n} + \bar{\rho}) - p(\bar{\rho}) - c^2(n_h + \tilde{n})) \right|. \end{aligned} \quad (5.7)$$

For readers' convenience, we show how to estimate $\|\nabla^k \mathfrak{M} * H(t)\|_{L^2}$ as an example. The other two estimates can be obtained by the similar argument. Indeed,

$$\begin{aligned} & \|\nabla^k \mathfrak{M} * H(t)\|_{L^2}^2 \\ & \lesssim \int_{|\xi| \leq \eta} e^{-2c_3|\xi|^2 t} |\xi|^{2k} |\widehat{H}|^2 d\xi + \int_{|\xi| \geq \eta} e^{-2R_0 t} |\xi|^{2k-4} |\widehat{H}|^2 d\xi \\ & \quad + \int_{|\xi| \geq \eta} e^{-2c_4|\xi|^2 t} |\xi|^{2k} |\widehat{H}|^2 d\xi \\ & \lesssim \int_{|\xi| \leq \eta} e^{-2c_3|\xi|^2 t} |\xi|^{2k+2} |\widehat{Q}|^2 d\xi + \int_{|\xi| \geq \eta} e^{-2R_0 t} |\xi|^{2k-2} |\widehat{Q}|^2 d\xi \\ & \quad + \int_{|\xi| \geq \eta} e^{-2c_4|\xi|^2 t} |\xi|^{2k+2} |\widehat{Q}|^2 d\xi \\ & \lesssim (1+t)^{-3(\frac{1}{q}-\frac{1}{2})-1-k} (\|Q\|_{L^q(\mathbb{R}^3)}^2 + \|\nabla^{\tilde{k}} Q\|_{L^2(\mathbb{R}^3)}^2), \quad q = 1, 2, \quad k-1 \leq \tilde{k} \in \mathbb{N}^+. \end{aligned}$$

5.1. L^2 -time decay rate. In this subsection, we establish the faster decay rate for (n_h, m_h) . We will start with an *a priori* assumption on a carefully chosen quantity $\Lambda(t)$ defined in (5.8), and then later prove a better estimate with the help of the smallness of initial data.

We begin with following Lemma.

Lemma 5.1. *Let $r_1, r_2 > 0$ be real, one has*

$$\begin{aligned} \int_0^{\frac{t}{2}} (1+t-\tau)^{-r_1} (1+\tau)^{-r_2} d\tau &= \int_0^{\frac{t}{2}} (1+\frac{t}{2}+\tau)^{-r_1} (1+\frac{t}{2}-\tau)^{-r_2} d\tau \\ &\lesssim \begin{cases} (1+t)^{-r_1}, & \text{for } r_2 > 1, \\ (1+t)^{-(r_1-\epsilon)}, & \text{for } r_2 = 1, \\ (1+t)^{-(r_1+r_2-1)}, & \text{for } r_2 < 1, \end{cases} \end{aligned}$$

and

$$\begin{aligned} \int_{\frac{t}{2}}^t (1+t-\tau)^{-r_1} (1+\tau)^{-r_2} d\tau &= \int_0^{\frac{t}{2}} (1+t-\tau)^{-r_2} (1+\tau)^{-r_1} d\tau \\ &\lesssim \begin{cases} (1+t)^{-r_2}, & \text{for } r_1 > 1, \\ (1+t)^{-(r_2-\epsilon)}, & \text{for } r_1 = 1, \\ (1+t)^{-(r_1+r_2-1)}, & \text{for } r_1 < 1, \end{cases} \end{aligned}$$

where $\epsilon > 0$ is a small but fixed constant.

Proposition 5.1. *Under the assumptions of Theorem 1.1, the solution (n_h, m_h) of the nonlinear system (5.1) satisfies for $k = 0, 1, 2$ that*

$$\begin{aligned} \|(\nabla^k n_h, \nabla^k m_h)\|_{L^2} &\leq C\delta_0^2 (1+t)^{-\frac{5}{4}-\frac{k}{2}}, \\ \|\nabla^3 m_h\|_{L^2} &\leq C\delta_0^2 (1+t)^{-\frac{11}{4}}, \quad \|\nabla^3 n_h\|_{L^2} \leq C\delta_0 (1+t)^{-\frac{7}{4}}, \end{aligned}$$

where C is a positive constant independent of time.

From (5.7), we deduce

$$Q(\tilde{U}, U_h) = Q_1 + Q_2 + Q_3 + Q_4,$$

which implies for a smooth solution (n, m) satisfying $\|(n, m)\|_{H^3} < \infty$ that

$$\begin{aligned} Q_1 &= Q_1(\tilde{U}, U_h) \sim O(1) (n_h^2 + m_h \otimes m_h + \tilde{n}^2 + \tilde{m} \otimes \tilde{m}), \\ Q_2 &= Q_2(\tilde{U}, U_h) \sim O(1) (\tilde{n}n_h + \tilde{m} \otimes m_h), \\ Q_3 &= Q_3(\tilde{U}, U_h) \sim O(1) (\nabla(n_h \cdot m_h) + \nabla(\tilde{n} \cdot \tilde{m})), \\ Q_4 &= Q_4(\tilde{U}, U_h) \sim O(1) (\nabla(\tilde{n} \cdot m_h) + \nabla(n_h \cdot \tilde{m})). \end{aligned}$$

Define

$$\begin{aligned} \Lambda(t) =: \sup_{0 \leq s \leq t} \left\{ \sum_{k=0}^2 (1+s)^{\frac{5}{4}+\frac{k}{2}} \delta_0^{-\frac{3}{4}} \|(\nabla^k n_h, \nabla^k m_h)(s)\|_{L^2} \right. \\ \left. + (1+s)^{\frac{7}{4}} \|(\nabla^3 n_h, \nabla^3 m_h)(s)\|_{L^2} \right\}. \end{aligned} \quad (5.8)$$

Proposition 5.2. *Under the assumptions of Theorem 1.1, if for some $T > 0$, $\Lambda(t) \leq \delta_0^{\frac{1}{2}}$ for any $t \in [0, T]$, then it holds that*

$$\Lambda(t) \leq C\delta_0^{\frac{3}{4}}, \quad t \in [0, T],$$

where C is a positive constant independent of time.

The proof of this Proposition 5.2 consists of following three steps.

5.2. Step 1: The basic energy estimates. Starting with (5.4), (5.5), (5.6) and (5.8), we have after a complicate but straightforward computation that

$$\begin{aligned}
\|(n_h, m_h)\|_{L^2} &\lesssim \int_0^t \|(\mathfrak{N}(t-\tau) * H(\tau), \mathfrak{M}(t-\tau) * H(\tau))\|_{L^2} d\tau \\
&\lesssim \int_0^t (1+t-\tau)^{-\frac{5}{4}} (\|Q(\tau)\|_{L^1} + \|Q(\tau)\|_{L^2}) d\tau \\
&\lesssim \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t)\right) \int_0^t (1+t-\tau)^{-\frac{5}{4}} (1+\tau)^{-\frac{3}{2}} d\tau \\
&\lesssim (1+t)^{-\frac{5}{4}} \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t)\right).
\end{aligned} \tag{5.9}$$

It is easy to verify that

$$\begin{aligned}
\|Q(t)\|_{L^1} &\lesssim \|Q_1\|_{L^1} + \|Q_2\|_{L^1} + \|Q_3\|_{L^1} + \|Q_4\|_{L^1} \\
&\lesssim \|(\tilde{n}, \tilde{m})\|_{L^2}^2 + \|(n_h, m_h)\|_{L^2}^2 + \|(n_h, m_h)\|_{L^2} (\|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^2} \\
&\quad + \|(\nabla n_h, \nabla m_h)\|_{L^2}) + \|(\tilde{n}, \tilde{m})\|_{L^2} (\|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^2} + \|(\nabla n_h, \nabla m_h)\|_{L^2}) \\
&\lesssim (1+t)^{-\frac{3}{2}} \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t)\right).
\end{aligned}$$

Indeed, by virtue of Hölder's inequality and Gagliardo-Nirenberg's inequality, we obtain that

$$\|u\|_{L^\infty} \lesssim \|\nabla u\|_{L^2}^{\frac{1}{2}} \|\nabla^2 u\|_{L^2}^{\frac{1}{2}},$$

which implies that

$$\begin{aligned}
&\|Q(t)\|_{L^2} \\
&\lesssim \|(\tilde{n}, \tilde{m})\|_{L^\infty} (\|(\tilde{n}, \tilde{m})\|_{L^2} + \|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^2} + \|(n_h, m_h)\|_{L^2} \\
&\quad + \|(\nabla n_h, \nabla m_h)\|_{L^2}) + \|(n_h, m_h)\|_{L^\infty} (\|(n_h, m_h)\|_{L^2} + \|(\nabla n_h, \nabla m_h)\|_{L^2}) \\
&\quad + \|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^\infty} \|(n_h, m_h)\|_{L^2} \\
&\lesssim (1+t)^{-\frac{9}{4}} \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t)\right).
\end{aligned}$$

Furthermore, exactly as in the estimate of the high order derivatives, we have

$$\begin{aligned}
&\|(\nabla n_h, \nabla m_h)\|_{L^2} \\
&\lesssim \int_0^{\frac{t}{2}} \|(\nabla \mathfrak{N}, \nabla \mathfrak{M})(t-\tau) * H(\tau)\|_{L^2} d\tau + \int_{\frac{t}{2}}^t \|(\mathfrak{N}, \mathfrak{M})(t-\tau) * \nabla H(\tau)\|_{L^2} d\tau \\
&\lesssim \int_0^{\frac{t}{2}} (1+t-\tau)^{-\frac{7}{4}} (\|Q(\tau)\|_{L^1} + \|\nabla Q(\tau)\|_{L^2}) d\tau + \int_{\frac{t}{2}}^t (1+t-\tau)^{-\frac{1}{2}} \|\nabla Q(\tau)\|_{L^2} d\tau \\
&\lesssim \left(\delta_0^2 + \delta_0^{\frac{9}{8}} \Lambda^2(t)\right) \left(\int_0^{\frac{t}{2}} (1+t-\tau)^{-\frac{7}{4}} (1+\tau)^{-\frac{3}{2}} d\tau + \int_{\frac{t}{2}}^t (1+t-\tau)^{-\frac{1}{2}} (1+\tau)^{-\frac{11}{4}} d\tau\right) \\
&\lesssim (1+t)^{-\frac{7}{4}} \left(\delta_0^2 + \delta_0^{\frac{9}{8}} \Lambda^2(t)\right),
\end{aligned} \tag{5.10}$$

Similarly, it holds that

$$\begin{aligned}
&\|\nabla Q(t)\|_{L^2} \\
&\lesssim \|(\tilde{n}, \tilde{m})\|_{L^\infty} (\|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^2} + \|(\nabla^2 \tilde{n}, \nabla^2 \tilde{m})\|_{L^2} + \|(\nabla n_h, \nabla m_h)\|_{L^2}
\end{aligned}$$

$$\begin{aligned}
& + \|(\nabla^2 n_h, \nabla^2 m_h)\|_{L^2} + \|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^\infty} (\|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^2} + \|(n_h, m_h)\|_{L^2} \\
& + \|(\nabla n_h, \nabla m_h)\|_{L^2}) + \|(n_h, m_h)\|_{L^\infty} (\|(\nabla^2 \tilde{n}, \nabla^2 \tilde{m})\|_{L^2} + \|(\nabla n_h, \nabla m_h)\|_{L^2} \\
& + \|(\nabla^2 n_h, \nabla^2 m_h)\|_{L^2}) + \|(\nabla n_h, \nabla m_h)\|_{L^\infty} \|(\nabla n_h, \nabla m_h)\|_{L^2} \\
& \lesssim (1+t)^{-\frac{11}{4}} \left(\delta_0^2 + \delta_0^{\frac{9}{8}} \Lambda^2(t) \right).
\end{aligned}$$

Thus, we also get that

$$\begin{aligned}
& \|(\nabla^2 n_h, \nabla^2 m_h)(t)\|_{L^2} \\
& \lesssim \int_0^{\frac{t}{2}} \|(\nabla^2 \mathfrak{N}, \nabla^2 \mathfrak{M})(t-\tau) * H(\tau)\|_{L^2} d\tau \\
& \quad + \int_{\frac{t}{2}}^t \|(\mathfrak{N}, \mathfrak{M})(t-\tau) * \nabla^2 H(\tau)\|_{L^2} d\tau \\
& \lesssim \int_0^{\frac{t}{2}} (1+t-\tau)^{-\frac{9}{4}} (\|Q(\tau)\|_{L^1} + \|\nabla^2 Q(\tau)\|_{L^2}) d\tau \\
& \quad + \int_{\frac{t}{2}}^t (1+t-\tau)^{-\frac{1}{2}} \|\nabla^2 Q(\tau)\|_{L^2} d\tau \\
& \lesssim \left(\delta_0^2 + \delta_0 \Lambda(t) + \delta_0^{\frac{3}{4}} \Lambda^2(t) \right) \left(\int_0^{\frac{t}{2}} (1+t-\tau)^{-\frac{9}{4}} (1+\tau)^{-\frac{3}{2}} d\tau \right. \\
& \quad \left. + \int_{\frac{t}{2}}^t (1+t-\tau)^{-\frac{1}{2}} (1+\tau)^{-\frac{13}{4}} d\tau \right) \\
& \lesssim (1+t)^{-\frac{9}{4}} \left(\delta_0^2 + \delta_0 \Lambda(t) + \delta_0^{\frac{3}{4}} \Lambda^2(t) \right). \tag{5.11}
\end{aligned}$$

Finally, we have

$$\begin{aligned}
& \|\nabla^2 Q(t)\|_{L^2} \\
& \lesssim (\|(\tilde{n}, \tilde{m})\|_{L^\infty} + \|(n_h, m_h)\|_{L^\infty}) (\|(\nabla^3 \tilde{n}, \nabla^3 \tilde{m})\|_{L^2} + \|(\nabla^3 n_h, \nabla^3 m_h)\|_{L^2}) \\
& \quad + (\|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^\infty} + \|(\nabla n_h, \nabla m_h)\|_{L^\infty}) (\|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^2} + \|(\nabla n_h, \nabla m_h)\|_{L^2}) \\
& \quad + (\|(\tilde{n}, \tilde{m})\|_{L^\infty} + \|(n_h, m_h)\|_{L^\infty} + \|(\nabla \tilde{n}, \nabla \tilde{m})\|_{L^\infty} + \|(\nabla n_h, \nabla m_h)\|_{L^\infty}) \\
& \quad \times (\|(\nabla^2 \tilde{n}, \nabla^2 \tilde{m})\|_{L^2} + \|(\nabla^2 n_h, \nabla^2 m_h)\|_{L^2}) \\
& \lesssim (1+t)^{-\frac{13}{4}} \left(\delta_0^2 + \delta_0 \Lambda(t) + \delta_0^{\frac{3}{4}} \Lambda^2(t) \right).
\end{aligned}$$

5.3. Step 2: The higher order energy estimates. In this subsection, we will close the *a priori* estimates and complete the proof of Proposition 5.2. For this purpose, we need to derive the time decay rate of higher order derivatives of (n_h, m_h) . We will establish the following lemma.

Lemma 5.2. *Under the assumption of Theorem 1.1, one has*

$$\|\nabla^2 n(t)\|_{H^1} + \|\nabla^2 u(t)\|_{H^1} \lesssim (1+t)^{-\frac{7}{4}} \left(\delta_0 + \delta_0^{\frac{3}{4}} \Lambda(t) \right).$$

In particular, it holds that

$$\|\nabla^3(n_h, m_h)(t)\|_{L^2} \lesssim (1+t)^{-\frac{7}{4}} \left(\delta_0 + \delta_0^{\frac{3}{4}} \Lambda(t) \right).$$

Proof. First of all, in view of (2.12), recovering the dissipation estimate for n , we see that

$$\begin{aligned} & \frac{d}{dt} \int_{\mathbb{R}^3} \nabla^2 u \cdot \nabla^3 n dx + C_1 \|\nabla^3 n\|_{L^2}^2 dx \\ & \leq C_2 (\|\nabla^3 u\|_{L^2}^2 + \|\nabla^4 u\|_{L^2}^2) + C(1+t)^{-\frac{3}{2}} \left(\delta_0 + \delta_0^{\frac{3}{8}} \Lambda(t) \right) \\ & \quad \times (\|\nabla^2 n\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2 + \|\nabla^3 u\|_{L^2}^2). \end{aligned} \quad (5.12)$$

Summing up (2.7) and (2.8) in the energy estimate for (n, u) , we can directly derive

$$\begin{aligned} & \frac{d}{dt} \int_{\mathbb{R}^3} (\gamma |\nabla^2 n|^2 + |\nabla^2 u|^2 + \gamma |\nabla^3 n|^2 + |\nabla^3 u|^2) dx + C_3 (\|\nabla^3 u\|_{L^2}^2 + \|\nabla^4 u\|_{L^2}^2) \\ & \leq C(1+t)^{-\frac{3}{2}} \left(\delta_0 + \delta_0^{\frac{3}{8}} \Lambda(t) \right) (\|\nabla^2 n\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2 + \|\nabla^3 n\|_{L^2}^2). \end{aligned} \quad (5.13)$$

Multiplying (5.12) by $\epsilon_1 \frac{C_3}{C_2}$ with $\epsilon_1 > 0$ a small but fixed constant, adding it with (5.13), we deduce that there exists a constant $C_4 > 0$ such that

$$\begin{aligned} & \frac{d}{dt} \left\{ \sum_{2 \leq k \leq 3} (\gamma \|\nabla^k n\|_{L^2}^2 + \|\nabla^k u\|_{L^2}^2) + \epsilon_1 \frac{C_3}{C_2} \int_{\mathbb{R}^3} \nabla^2 u \cdot \nabla^3 n dx \right\} \\ & \quad + C_4 \left(\|\nabla^3 n\|_{L^2}^2 + \sum_{3 \leq k \leq 4} \|\nabla^k u\|_{L^2}^2 \right) \\ & \leq C(1+t)^{-\frac{3}{2}} \left(\delta_0 + \delta_0^{\frac{3}{8}} \Lambda(t) \right) (\|\nabla^2 n\|_{L^2}^2 + \|\nabla^2 u\|_{L^2}^2). \end{aligned}$$

Next, we define

$$\mathcal{E}_1(t) = \left\{ \sum_{2 \leq k \leq 3} (\gamma \|\nabla^k n\|_{L^2}^2 + \|\nabla^k u\|_{L^2}^2) + \epsilon_1 \frac{C_3}{C_2} \int_{\mathbb{R}^3} \nabla^2 u \cdot \nabla^3 n dx \right\}.$$

Observe that since $\epsilon_1 \frac{C_3}{C_2}$ is small, then there exists a constant $C_5 > 0$ such that

$$C_5^{-1} (\|\nabla^2 n(t)\|_{H^1}^2 + \|\nabla^2 u(t)\|_{H^1}^2) \leq \mathcal{E}_1(t) \leq C_5 (\|\nabla^2 n(t)\|_{H^1}^2 + \|\nabla^2 u(t)\|_{H^1}^2).$$

Then we arrive at

$$\frac{d}{dt} \mathcal{E}_1(t) + C_4 (\|\nabla^3 n(t)\|_{L^2}^2 + \|\nabla^3 u(t)\|_{H^1}^2) \leq C(1+t)^{-5} \left(\delta_0 + \delta_0^{\frac{3}{8}} \Lambda(t) \right) \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t) \right).$$

Denote $S(t) = \left\{ \xi \mid |\xi| \leq \sqrt{\frac{3(1+\gamma)}{C_4}} (1+t)^{-\frac{1}{2}} \right\}$ the time-dependent n -dimensional sphere. This decomposition allows us to estimate L^2 time decay depend on $(\hat{n}, \hat{u})$ for frequency values $\xi \in S(t)$, then we obtain that

$$\begin{aligned} & \frac{C_4}{3} \|\nabla^3(n, u)(x)\|_{L^2}^2 \geq \frac{C_4}{3} \int_{S(t)^c} |\xi|^6 |(\hat{n}, \hat{u})(\xi)|^2 d\xi \\ & \geq (1+\gamma)(1+t)^{-1} \int_{\mathbb{R}^3} |\xi|^4 |(\hat{n}, \hat{u})(\xi)|^2 d\xi - (1+\gamma)(1+t)^{-1} \int_{S(t)} |\xi|^4 |(\hat{n}, \hat{u})(\xi)|^2 d\xi. \end{aligned}$$

Hence we have

$$\begin{aligned} & \frac{d}{dt} \mathcal{E}_1(t) + (1+t)^{-1} \mathcal{E}_1(t) + \|\nabla^3 n\|_{L^2}^2 + \|\nabla^3 u\|_{H^1}^2 \\ & \lesssim (1+t)^{-5} \left(\delta_0 + \delta_0^{\frac{3}{2}} \Lambda(t) \right) \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t) \right) + (1+t)^{-1} \int_{S(t)} |\xi|^4 |(\widehat{n}, \widehat{u})(\xi)|^2 d\xi \\ & \quad + (1+t)^{-1} \int_{\mathbb{R}^3} \nabla^2 u \cdot \nabla^3 n dx. \end{aligned}$$

Multiplying the above equation by $(1+t)^5$, we obtain that

$$\frac{d}{dt} \left\{ (1+t)^5 \mathcal{E}_1(t) \right\} + (1+t)^5 \left(\|\nabla^3 n\|_{L^2}^2 + \|\nabla^3 u\|_{H^1}^2 \right) \lesssim (1+t)^{\frac{1}{2}} \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t) \right).$$

Integrating it with respect to time from 0 to T , then we have

$$\begin{aligned} & (1+t)^5 \mathcal{E}_1(t) + \int_0^T (1+t)^5 \left(\|\nabla^3 n\|_{L^2}^2 + \|\nabla^3 u\|_{H^1}^2 \right) dt \\ & \lesssim \mathcal{E}_1(0) + (1+t)^{\frac{3}{2}} \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t) \right), \end{aligned}$$

which implies that

$$\|\nabla^3 n\|_{L^2}^2 + \|\nabla^3 u\|_{L^2}^2 \lesssim \mathcal{E}_1(t) \lesssim (1+t)^{-5} \delta_0^2 + (1+t)^{-\frac{7}{2}} \left(\delta_0^2 + \delta_0^{\frac{3}{2}} \Lambda^2(t) \right).$$

Finally, we have

$$\|\nabla^3 n_h\|_{L^2} + \|\nabla^3 m_h\|_{L^2} \lesssim (1+t)^{-\frac{7}{4}} \left(\delta_0 + \delta_0^{\frac{3}{4}} \Lambda(t) \right).$$

This completes the proof of this Lemma. $\square$

5.4. Step 3: Close the estimates. In this subsection, we first combine the above a priori estimates of (5.8), (5.9), (5.10), (5.11) and Lemma 5.2 together to give the proof of the Proposition 5.2. In deed, for any $t \in [0, T]$, we have shown that

$$\Lambda(t) \leq C \left(\delta_0 + \delta_0^{\frac{1}{4}} \Lambda(t) + \Lambda^2(t) \right) \leq C \delta_0^{\frac{3}{4}}. \quad (5.14)$$

With the help of standard continuity argument, Proposition 5.2 and the smallness of $\delta_0 > 0$, implies that $\Lambda(t) \leq C \delta_0^{\frac{3}{4}}$ for any $t > 0$. Moreover, we deduce the time decay estimate for (n_h, m_h) from (5.9), (5.10), (5.11), Lemma 5.2 and (5.14) that

$$\begin{aligned} & \|(\nabla^k n_h, \nabla^k m_h)\|_{L^2} \lesssim \delta_0^2 (1+t)^{-\frac{5}{4}-\frac{k}{2}}, \quad k = 0, 1, \\ & \|\nabla^2(n_h, m_h)\|_{L^2} \lesssim \delta_0^{\frac{7}{4}} (1+t)^{-\frac{9}{4}}, \quad \|\nabla^3(n_h, m_h)\|_{L^2} \lesssim \delta_0 (1+t)^{-\frac{7}{4}}. \end{aligned}$$

Consequently, for any $t \in [0, T]$ we have

$$\Lambda(t) \leq C \delta_0. \quad (5.15)$$

From (5.11) and (5.15), thus we also get that

$$\|\nabla^2(n_h, m_h)\|_{L^2} \lesssim \delta_0^2 (1+t)^{-\frac{9}{4}}.$$

For $\nabla^3 m_h$, in view of the (5.6), we see that

$$\begin{aligned}
& \|\nabla^3 m_h(t)\|_{L^2} \\
& \lesssim \int_0^{\frac{t}{2}} (1+t-\tau)^{-\frac{11}{4}} (\|Q(\tau)\|_{L^1} + \|\nabla^2 Q(\tau)\|_{L^2}) d\tau \\
& \quad + \int_{\frac{t}{2}}^t (1+t-\tau)^{-\frac{1}{2}} \|\nabla^2 Q(\tau)\|_{L^2} d\tau \\
& \lesssim \delta_0^2 \left(\int_0^{\frac{t}{2}} (1+t-\tau)^{-\frac{11}{4}} (1+\tau)^{-\frac{3}{2}} d\tau + \int_{\frac{t}{2}}^t (1+t-\tau)^{-\frac{1}{2}} (1+\tau)^{-\frac{13}{4}} d\tau \right) \\
& \lesssim \delta_0^2 (1+t)^{-\frac{11}{4}}.
\end{aligned}$$

Hence, we finish the proof of the Proposition 5.1. Theorem 1.1 follows.

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REFERENCES

- [1] K. Deckelnick, [Decay estimates for the compressible Navier-Stokes equations in unbounded domains](#), *Math. Z.*, **209** (1992), 115–130.
- [2] K. Deckelnick, [\$L^2\$ decay for the compressible Navier-Stokes equations in unbounded domains](#), *Comm. Partial Differential Equations*, **18** (1993), 1445–1476.
- [3] R. Duan, H. Liu, S. Ukai and T. Yang, [Optimal \$L^p\$ - \$L^q\$ convergence rates for the compressible Navier-Stokes equations with potential force](#), *J. Differential Equations*, **238** (2007), 220–233.
- [4] R. Duan, S. Ukai and T. Yang and H. Zhao, [Optimal convergence rates for the compressible Navier-Stokes equations with potential forces](#), *Math. Models Methods Appl. Sci.*, **17** (2007), 737–758.
- [5] Y. Guo and Y. Wang, [Decay of dissipative equations and negative Sobolev spaces](#), *Comm. Partial Differential Equations*, **37** (2012), 2165–2208.
- [6] D. Hoff and K. Zumbrun, [Multi-dimensional diffusion waves for the Navier-Stokes equations of compressible flow](#), *Indiana Univ. Math. J.*, **44** (1995), 603–676.
- [7] D. Hoff and K. Zumbrun, [Pointwise decay estimates for multidimensional Navier-Stokes diffusion waves](#), *Z. Angew. Math. Phys.*, **48** (1997), 597–614.
- [8] Y. Kagei and S. Kawashima, [Stability of planar stationary solutions to the compressible Navier-Stokes equation on the half space](#), *Comm. Math. Phys.*, **266** (2006), 401–430.
- [9] Y. Kagei and T. Kobayashi, [On large-time behavior of solutions to the compressible Navier-Stokes equations in the half space in \$\mathbb{R}^3\$](#) , *Arch. Ration. Mech. Anal.*, **165** (2002), 89–159.
- [10] Y. Kagei and T. Kobayashi, [Asymptotic behavior of solutions of the compressible Navier-Stokes equations on the half space](#), *Arch. Ration. Mech. Anal.*, **177** (2005), 231–330.
- [11] T. Kobayashi, [Some estimates of solutions for the equations of motion of compressible viscous fluid in the three-dimensional exterior domain](#), *J. Differential Equations*, **184** (2002), 587–619.
- [12] T. Kobayashi and Y. Shibata, [Decay estimates of solutions for the equations of motion of compressible viscous and heat-conductive gases in an exterior domain in \$\mathbb{R}^3\$](#) , *Comm. Math. Phys.*, **200** (1999), 621–659.
- [13] H.-L. Li, A. Matsumura and G. Zhang, [Optimal decay rate of the compressible Navier-Stokes-Poisson system in \$\mathbb{R}^3\$](#) , *Arch. Ration. Mech. Anal.*, **196** (2010), 681–713.

- [14] T.-P. Liu and W. Wang, [The pointwise estimates of diffusion wave for the Navier-Stokes systems in odd multi-dimensions](#), *Comm. Math. Phys.*, **196** (1998), 145–173.
- [15] T.-P. Liu and Y. Zeng, [Large time behavior of solutions for general quasilinear hyperbolic-parabolic systems of conservation laws](#), *Mem. Amer. Math. Soc.*, **125** (1997), 120 pp.
- [16] A. Matsumura, An energy method for the equations of motion of compressible viscous and heat-conductive fluids, *MRC-Technical Summary Report*, **2194** (1981), 1–16.
- [17] A. Matsumura and T. Nishida, [The initial value problem for the equations of motion of compressible viscous and heat-conductive fluids](#), *Proc. Japan Acad. Ser. A Math. Sci.*, **55** (1979), 337–342.
- [18] A. Matsumura and T. Nishida, [The initial value problem for the equations of motion of viscous and heat-conductive gases](#), *J. Math. Kyoto Univ.*, **20** (1980), 67–104.
- [19] G. Ponce, [Global existence of small solutions to a class of nonlinear evolution equations](#), *Nonlinear Anal.*, **9** (1985), 399–418.
- [20] M. E. Schonbek, [Large time behaviour of solutions to the Navier-Stokes equations](#), *Comm. Partial Differential Equations*, **11** (1986), 733–763.
- [21] M. E. Schonbek, [Lower bounds of rates of decay for solutions to the Navier-Stokes equations](#), *J. Amer. Math. Soc.*, **4** (1991), 423–449.
- [22] Y. Shibata and K. Tanaka, [On the steady flow of compressible viscous fluid and its stability with respect to initial disturbance](#), *J. Math. Soc. Japan*, **55** (2003), 797–826.
- [23] Y. Shibata and K. Tanaka, [Rate of convergence of non-stationary flow to the steady flow of compressible viscous fluid](#), *Comput. Math. Appl.*, **53** (2007), 605–623.
- [24] Y. Zeng, [\$L^1\$ asymptotic behavior of compressible, isentropic, viscous 1-D flow](#), *Comm. Pure Appl. Math.*, **47** (1994), 1053–1082.

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