



The Nature of Ionized Gas in the Milky Way Galactic Fountain

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Abstract

We address the spatial scale, ionization structure, mass, and metal content of gas at the Milky Way disk–halo interface detected as absorption in the foreground of seven closely spaced, high-latitude halo blue horizontal branch stars with heights $z = 3\text{--}14$ kpc. We detect transitions that trace multiple ionization states (e.g., Ca II, Fe II, Si IV, C IV) with column densities that remain constant with height from the disk, indicating that the gas most likely lies within $z < 3.4$ kpc. The intermediate ionization state gas traced by C IV and Si IV is strongly correlated over the full range of transverse separations probed by our sight lines, indicating large, coherent structures greater than 1 kpc in size. The low ionization state material traced by Ca II and Fe II does not exhibit a correlation with either N_{H_1} or transverse separation, implying cloudlets or clumpiness on scales less than 10 pc. We find that the observed ratio $\log(N_{\text{Si IV}}/N_{\text{C IV}})$, with a median value of -0.69 ± 0.04 , is sensitive to the total carbon content of the ionized gas under the assumption of either photoionization or collisional ionization. The only self-consistent solution for photoionized gas requires that Si be depleted onto dust by 0.35 dex relative to the solar Si/C ratio, similar to the level of Si depletion in DLAs and in the Milky Way interstellar medium. The allowed range of values for the areal mass infall rate of warm, ionized gas at the disk–halo interface is $0.0003 < dM_{\text{gas}}/dtdA [M_{\odot} \text{ kpc}^{-2} \text{ yr}^{-1}] < 0.006$. Our data support a physical scenario in which the Milky Way is fed by complex, multiphase processes at its disk–halo interface that involve kiloparsec-scale ionized envelopes or streams containing parsec-scale, cool clumps.

Unified Astronomy Thesaurus concepts: Ultraviolet astronomy (1736); Photoionization (2060); Warm ionized medium (1788); Intergalactic dust clouds (810); Circumgalactic medium (1879); Galaxy evolution (594); Halo stars (699); High resolution spectroscopy (2096)

1. Introduction

The gaseous halo of the Milky Way (MW) supports a dynamic baryon cycle that provides fuel for future star formation and collects the by-products of stellar evolution (Putman et al. 2012). This cycle is often envisioned as a Galactic fountain, in which hot, overpressurized gas is ejected from the disk into the halo and subsequently cools, recombines, and rains back down onto the disk (Shapiro & Field 1976; Fraternali & Binney 2008). However, the physical mechanisms that inject energy and material into the fountain and the detailed gas physics that follow are likely considerably more complex (e.g., Kim & Ostriker 2018). Our ability to detect the gas flowing into and out of the Galaxy at small distances—just as it emerges from the disk, or accretes to fuel the interstellar medium (ISM)—presents a unique opportunity to constrain the physics that drive the Galactic fountain.

Despite a rich observational history of absorption and emission studies, following the complete gas cycle of star-forming galaxies, from halo to disk, has remained a challenge. Since Münch & Zirin (1961) first observed absorption from Ca II H + K lines in the

spectra of high-latitude early-type stars, it has been hypothesized that the MW is enveloped in a diffuse, multiphase halo (Spitzer 1956). More recently, absorption-line experiments using bright background sources, including halo stars and quasars, have been used to probe the diffuse circumgalactic gas surrounding our own Galaxy and others at projected distances >10 kpc (e.g., Spitzer 1965; Bergeron 1986; Savage & Sembach 1994; Lanzetta et al. 1995; Sembach et al. 2003; Fox et al. 2004; Lehner et al. 2012; Prochaska et al. 2011). H I 21 cm emission maps of our own Galaxy's halo reveal giant arches, shells, and complexes of neutral clouds moving with $25 \text{ km s}^{-1} < |\nu_{\text{LSR}}| < 250 \text{ km s}^{-1}$ (e.g., Wakker & van Woerden 1997). The Wisconsin H α Mapper (WHAM; Haffner et al. 1999) has found extended, diffuse Galactic H α emission above the Galactic plane, not associated with H II regions. Studies of the pulsar dispersion measure (DM) in the Galactic halo constrain the electron density of a warm ionized medium (WIM) with a characteristic height of 2 kpc above the disk midplane (Gaensler et al. 2008). The physical interpretation of these rich data sets is hampered by limited distance constraints, sparsely sampled background sight lines, and degeneracies between local “cloud” sizes and the unknown phase structure of the gas.

Absorption-line studies of the MW halo itself have used both distant QSOs and bright halo stars to constrain the nature and



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Table 1
BHB Star Properties

BHB Name, ID	R.A. (deg)	Decl. (deg)	l (deg)	b (deg)	m_{FUV}	m_g	HRV (km s $^{-1}$)	d (kpc)	x (kpc)	y (kpc)	z (kpc)
J1527+4027 BHB1	231.9586	40.4580	65.6478	55.2262	16.30	13.78	-105.6 ± 0.9	4.1	7.0	-2.1	3.4
J1534+5015 BHB2	233.5468	50.2657	80.9667	51.4485	16.95	14.20	-129.4 ± 1.8	5.2	7.5	-3.2	4.1
J1344+1842 BHB3	206.0184	18.7164	0.8905	75.2642	17.27	14.46	-74.7 ± 1.4	5.9	6.5	0.0	5.7
J1325+2232 BHB4	201.4767	22.5473	4.4384	80.9410	17.83	14.90	-100.3 ± 1.8	7.4	6.8	-0.1	7.3
J1415+3716 BHB5	213.9881	37.2828	67.9148	69.5097	17.97	15.57	-134.0 ± 1.2	9.9	6.7	-3.2	9.2
J1341+2823 BHB6	205.3327	28.3996	42.3913	78.8949	18.07	16.07	-124.8 ± 2.2	11.3	6.4	-1.5	11.1
J1459+4128 BHB7	224.9749	41.4809	69.9352	60.0843	18.47	16.72	-183.1 ± 4.4	15.3	5.4	-7.2	13.3

Note. Properties of target BHB stars based on SDSS DR8 SEGUE data (see Xue et al. 2011, for details); all coordinates are in epoch J2000. *GALEX* FUV magnitudes are determined by crossmatching the closest point-source detection to each BHB target.

amount of diffuse, high-velocity ($|\nu_{\text{LSR}}| > 90 \text{ km s}^{-1}$) material beyond the disk. The Galactic corona exhibits a range of ionization states traced by numerous ionized metal line transitions. From observations with the *Far Ultraviolet Spectroscopic Explorer* (FUSE) and, later, the Cosmic Origins Spectrograph (COS), we see widespread, yet patchy, absorption from Si IV, C IV, and O VI that indicates $T \approx 10^5 \text{ K}$ gas at scale heights of 2–5 kpc (e.g., Savage & Wakker 2009). Furthermore, ionized, diffuse, high-velocity gas is distributed over large spatial scales, with a covering fraction of $\sim 60\%$ (Savage et al. 2003; Sembach et al. 2003). Based on conservative ionization model assumptions, this gas provides enough potential material to sustain star formation in the MW for many gigayears (Lehner & Howk 2011).

We present observations that directly constrain the kinematics, physical coherence scale, mass, and metal content of high-latitude ionized gas at the MW disk–halo interface. The disk–halo interface region of the Galaxy includes extraplanar, ionized material, sometimes referred to as the WIM, thick disk, and corona, and roughly extends to $\pm 10 \text{ kpc}$ above and below the disk. It additionally contains well-known structures mapped in H I 21 cm emission commonly referred to as intermediate-velocity clouds (IVCs) and high-velocity clouds (HVCs). Using a bracketing method that requires spectroscopy toward foreground and background halo stars with known distances, the heights to IVCs are constrained to lie within $\lesssim 3 \text{ kpc}$, and those of HVCs are mostly constrained to lie within heights $\lesssim 15 \text{ kpc}$ (e.g., Danly et al. 1993; Wakker et al. 1996, 2001, 2007; Thom et al. 2008; Wakker et al. 2008).

In this study, we map the absorption signatures toward a sample of seven halo blue horizontal branch (BHB) stars at different heights out to 14 kpc in the direction of the north Galactic pole. To select our targets, we used a catalog of > 4500 BHB stars with well-calibrated distances and velocities (Xue et al. 2011). The BHB stars in our sample lie at high Galactic latitudes, $b > 60^\circ$. At these latitudes, we minimize ambiguities in tracking inflowing and outflowing material via blueshifts and redshifts along the lines of sight. We also minimize the absorption due to the MW ISM along each line of sight. With a small sample of background probes at various scale heights, our goal is to place physical constraints on extraplanar ionized and neutral material in the MW.

Throughout this manuscript we will report gas radial velocities, ν_{rad} , along given lines of sight from our vantage point at the location of the Sun in the MW. As we have no information about the transverse velocity of the gas perpendicular to the line of sight, it is important to distinguish these 1D

velocities from true spatial velocities. Furthermore, we report all velocities in the frame of the local standard of rest (LSR), which accounts for the peculiar velocity of the Sun with respect to the rotation of the Galaxy.

In Section 2 we describe the data and observations used in this study. Section 3 presents our line profile analysis of the absorption-line data. In Section 4, we present our results based on the column density and velocity centroid measurements. Section 5 examines how the ratio of Si IV/C IV constrains both photoionization and collisional ionization models of the high-latitude, intermediate-velocity (IV) gas we observe. In Section 6, we summarize our key results and discuss the implied spatial scales of the observed clouds and/or interfaces in the context of physical models of the Galactic fountain, and finally we discuss the implications of possibly dusty, ionized disk–halo interface material.

2. Data and Sample: Target Blue Horizontal Branch Stars

We targeted BHB stars in the halo of the MW toward the north Galactic pole extending 13 kpc above the disk into the halo. Our catalog of BHB stars comes from Xue et al. (2011), who selected BHB stars from Sloan Digital Sky Survey (SDSS) DR8 stars with fiber spectra. Halo BHB stars are good beacons for four reasons: (1) they are blue point sources and therefore have bright continua at UV and optical wavelengths; (2) their UV spectra are characterized by a fairly featureless bright continuum with broad Lyman series lines (Brown et al. 2012) because they are quite metal-poor; (3) they are approximate “standard candles,” since their distances can be estimated from multiband photometry with an accuracy of 10% (Deason et al. 2011); and (4) many of them are moving with $|\nu_{\text{rad}}| > 100 \text{ km s}^{-1}$, which allows us to cleanly separate any absorption features intrinsic to the star from the disk–halo interface gas (the BHB intrinsic stellar spectra do not exhibit any detectable Si IV or C IV but do seem to show a modest amount of C I, Si II, Fe II, and Ca II; see Section 2.3).

The seven BHB stars in our COS sample have far-UV (FUV) magnitudes $m_{\text{FUV}} < 18.5$, bright enough to be observed with COS in less than eight orbits per star. All lie at a high Galactic latitude ($l > 60^\circ$), which in principle allows us to distinguish inflows from outflows above the MW disk. Table 1 presents the properties of our sample of seven BHB stars, with coordinates, distances, and stellar velocities. Reported BHB distances serve as an upper limit on the distance to any gas detected along the line of sight. Figure 1 shows our seven sight lines on an all-sky map of the MW, along with the approximate locations of well-studied IV extraplanar H I structures for reference.

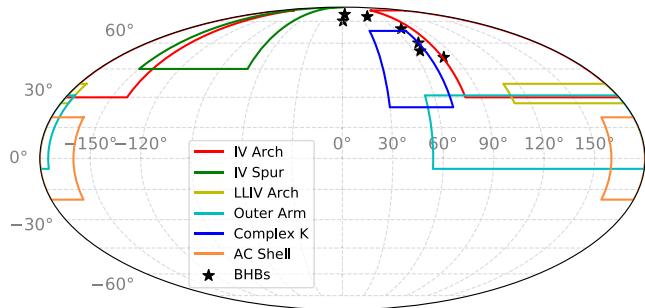


Figure 1. All-sky map in Galactic coordinates, with Galactic center at $l = 0^\circ$, $b = 0^\circ$, showing the positions of the seven BHB sight lines as black stars. The colored shapes mark the approximate locations of well-studied IV complexes identified in H I 21 cm maps, with coordinates based on those given in van Woerden et al. (2005). The positions of the BHB stars skirt the very edge of the large, high-latitude structure known as the IV arch. For reference, Complex K has very patchy coverage in the area shown; velocities centered on $v_{\text{LSR}} \approx -80 \text{ km s}^{-1}$, which is higher than the absorption velocities we observe; and much lower overall integrated H I column densities than the IV arch (Wakker 2001).

In Table 1, we additionally provide 3D Cartesian coordinates (given as xyz), commonly used in dynamical studies of the MW halo. This coordinate system is particularly useful for our purposes in examining the physical structure of gas in the frame of the MW. For reference, in this coordinate system, Galactic center is at the origin and the Sun lies at $(x, y, z) = (8, 0, 0)$ kpc. The x -axis extends from Galactic center in the direction of the Sun, y is perpendicular to x in the plane of the disk, and z is perpendicular to the disk such that the system is right-handed, with positive z toward the north Galactic pole. We will refer to z as height for the remainder of the paper.

2.1. HST/COS Spectroscopy

FUV spectra were obtained with COS (Froning & Green 2009; Green et al. 2012) on the *Hubble Space Telescope* (Program ID: 14140). We observed our BHB targets with the COS G160M grating with a central wavelength of 1577 Å with exposure times adjusted to achieve a signal-to-noise ratio (S/N) ≈ 15 per resolution element over $\lambda_{\text{obs}} \approx 1385\text{--}1760$ Å. Table 2 contains a summary of the COS observations.

Our COS data probe absorption due to the following metal ions: Si II ($\lambda 1526$), Fe II ($\lambda 1608$), Al II ($\lambda 1670$), Si IV ($\lambda\lambda 1393, 1402$), and C IV ($\lambda\lambda 1548, 1550$). We refer to singly ionized species as low ions. These low ions have ionization potential energies ranging from ≈ 16 to 18 eV and are most likely arising in cool gas photoionized by stellar radiation from the MW, with some contribution from the extragalactic UV background (EUVB). We refer to Si IV and C IV as intermediate ions. The ionization potential energies of Si IV and C IV are 45 and 64 eV, respectively. The intermediate ions may trace warm photoionized material and/or may be produced in part by collisional ionization.

We combine the CALCOS-generated x1D files using v3.1.1 of the COADD_X1D routine provided by the COS-GTO team (Danforth et al. 2016), which properly treats the error arrays of the input files using Poisson statistics. The code aligns the different exposures by determining a constant offset determined by cross-correlating strong ISM lines in a 10 Å wide region of the spectrum. With COS spectra, misalignments with magnitudes up to $\pm 30 \text{ km s}^{-1}$ between spectral lines in the same target have been described, which leads to smearing and

unphysical velocity offsets between different absorption lines. We discuss these wavelength calibration issues in detail in Section 2.3.

The COS line-spread function (LSF) is well described by a Gaussian convolved with a power law that extends to many tens of pixels beyond the line center (Ghavamian et al. 2009). These broad wings both affect the precision of our equivalent width measurements and complicate assessments of line saturation. We mediate these effects when we fit absorption lines (described in Section 3) by convolving with the real LSF. Each COS resolution element at $R \sim 18,000$ covers 16 km s^{-1} and is sampled by six raw pixels. We performed our analysis on the data binned by three native spectral pixels to a dispersion of $\Delta\lambda \approx 0.0367$ Å. The resulting science-grade spectra are characterized by an FWHM $\approx 18 \text{ km s}^{-1}$.

2.2. Keck/HIRES Spectroscopy

As a complement to our *HST*/COS data, in 2015A we began obtaining Keck HIRES spectra covering Ca II H + K ($\lambda\lambda 3934, 3969$) and Na I ($\lambda\lambda 5892, 5897$) along the BHB lines of sight through several programs. The Keck HIRES data are fully described in Bish et al. (2019) and are summarized here. Exposure times varied with the magnitude of the star and are given in Table 2. Generally, we aimed to achieve an S/N of $\gtrsim 20$ over the range of 3900–5900 Å. Our final wavelength ranges varied slightly with different cross-disperser and echelle angles for the BHB stars on different nights, depending on the program for which it was observed. We obtained HIRES spectra with the C5 decker with the image rotator in vertical-angle mode, and the kv370 order blocking filter with 2×1 binning. The C5 decker with a 1.148 slit provides spectroscopic resolution of $1.3 \text{ km s}^{-1} \text{ pixel}^{-1}$ and $R = 36,000$, which corresponds to a velocity FWHM of 8.3 km s^{-1} . The C1 decker with a 0.861 slit provides spectroscopic resolution of $1.0 \text{ km s}^{-1} \text{ pixel}^{-1}$ and $R = 48,000$, which corresponds to a velocity FWHM of 6.2 km s^{-1} . Our final velocity resolution of $5\text{--}10 \text{ km s}^{-1}$ allows us to distinguish material from the disk of the MW and the star and resolve different absorption components from each other.

The ionization potential energies of Ca II and Na I are 11.9 and 5.1 eV, respectively. Because they lie below 1 ryd, these ions are likely associated with cool, predominantly neutral material. However, Na I is known to show up sporadically along MW ISM lines of sight and exhibit little correlation with gas-phase metallicity or H I column density (Jenkins 2009). Thus, the interpretation of the Na I line is complex, whereas Ca II is a typical ion tracing neutral and low ionization state gas.

2.3. Wavelength Calibration Errors in COS

Here we describe our analysis to ensure that the COS wavelength solution of the co-added, aligned spectra produced by the Danforth routine (Danforth et al. 2016) is accurate. Our strategy is to check our solution against Ca II stellar absorption lines detected in Keck HIRES spectra and SDSS SEGUE stellar radial velocities of the same targets. We fit unsaturated, unblended stellar absorption lines of Ca II in the Keck spectra and C I, Si II, and Fe II in the COS spectra to check for potential systematic errors in the COS wavelength calibration that have been reported and discussed in detail by Wakker et al. (2015). We fit stellar absorption profiles along with gas absorption

Table 2
Summary of Keck/HIRES and *HST*/COS G160M Observations

BHB Name ID	HIRES: Date-Obs	Exptime (s)	S/N	decker, λ Range (Å)	COS: Date-Obs	N_{orbit}	S/N
J1527+4027 BHB1	2015 May 14	400	48	C1: 3824–6880	2015 Dec 29	1	14.1
J1534+5015 BHB2	2015 Apr 12	420	31	C1: 3678–6576	2015 Dec 24	2	17.0
J1344+1842 BHB3	2016 Mar 30	600	30	C5: 3830–8350	2016 May 15	3	16.3
J1325+2232 BHB4	2015 May 14	800	25	C1: 3824–6880	2016 May 24	4	15.2
J1415+3716 BHB5	2016 Mar 30	900	24	C5: 3830–8350	2016 Feb 21, 27	6	16.0
J1341+2823 BHB6	2015 May 14	1600	28	C1: 3824–6880	2016 May 26, Jun 25	6	15.7
J1459+4128 BHB7	2017 Feb 20	1800	18	C5: 3830–8350	2015 Dec 17, 21	8	15.3

profiles as described in Section 3. We compare our fitted stellar radial velocities from HIRES and COS spectra with the SDSS DR8 radial velocities for the BHB stars in the v_{LSR} rest frame.

Existing methods to test and correct for systematic errors in the COS wavelength calibration assume alignment in MW interstellar absorption lines of various species, sometimes with a range of ionization potential energies (e.g., Wakker et al. 2015; Danforth et al. 2016). The primary benefit of these methods is that they are broadly applicable to a variety of data sets, including QSO spectra. Instead, our method assumes that stellar absorption lines of different ionic species observed with different instruments should give a consistent value for the stellar radial velocity of each BHB star. Notably, our method will allow for velocity offsets between different ionic species owing to foreground gas absorption at $z \approx 0$ along each BHB line of sight. Thus, any differences between gas kinematics of different ionic species at the disk–halo interface of the MW may correspond to physical velocity offsets between different gas phases.

Figure 2 shows velocity offsets of stellar absorption lines between those measured in SDSS spectra (Xue et al. 2011) and those measured in our Keck or COS spectra. We show the offsets (Δv_{SDSS}) for the different ionic species that consistently exhibit stellar absorption: Ca II $\lambda 3934$ (Keck), C I $\lambda 1656$, Fe II $\lambda 1608$, and Si II $\lambda 1526$ (COS). We use a rest wavelength for C I $\lambda 1656$ of 1656.9283, which is the wavelength of the 3P_0 ground-state transition of C I, and do not include any contribution from weaker excited fine-structure levels of the C I multiplet that span $\pm 335 \text{ km s}^{-1}$ since they do not appear to be present in our data. We note that Al II $\lambda 1670$ often shows absorption from the BHB star; however, we exclude it here because it is invariably saturated and often blended with the MW interstellar absorption, which makes its fitted velocity centroid more uncertain.

Figure 2 shows that for most BHB stars the stellar radial velocities measured in Keck and COS spectra are consistent with SDSS radial velocities within approximately $\pm 10 \text{ km s}^{-1}$. The velocity centroids of Ca II $\lambda 3934$ as measured in the Keck data (small, red squares) are generally in better agreement with the SDSS radial velocities, with a median $|\Delta v_{\text{SDSS}}$ of 5.5 km s^{-1} . Next, we note that there are no apparent trends between Δv_{SDSS} and wavelength for the stellar absorption lines detected in the COS G160M spectra. The gray shaded regions indicate the approximate mean rms of the dispersion solutions for our Keck HIRES (dark gray; $\sim 2 \text{ km s}^{-1}$) and COS G160M (light gray; $\sim 7.5 \text{ km s}^{-1}$) spectra. Reported errors in the SDSS heliocentric radial velocities for each BHB range between 1 and 3 km s^{-1} .

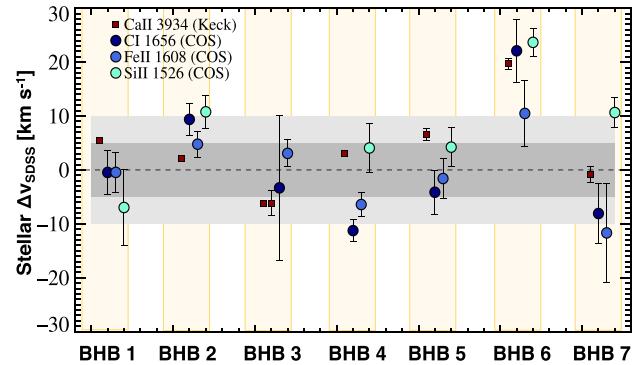


Figure 2. Difference between SEGUE BHB stellar radial velocities as measured in SDSS DR8 spectra and those we measure in our Keck/HIRES and *HST*/COS spectra for different ionic transitions. The data are organized along the x -axis by BHB ID number, as given in Table 1, and by ionic transition as given in the key on the upper left of the panel. Stellar absorption is apparent in Ca II $\lambda 3934$ (Keck), C I $\lambda 1656$, Fe II $\lambda 1608$, and Si II $\lambda 1526$ (COS) for most of our lines of sight. The gray shaded regions indicate the approximate mean rms of the dispersion solutions for our Keck HIRES (dark gray; $\sim 2 \text{ km s}^{-1}$) and COS G160M (light gray; $\sim 7.5 \text{ km s}^{-1}$) spectra.

In general, we find that SDSS radial velocities are in agreement with those measured from the Keck and COS spectra. One exception is the stellar absorption we detect along the line of sight to BHB6 (J1341+2823). The fits to the stellar velocities in both Keck HIRES and COS G160M spectra show $> 10 \text{ km s}^{-1}$ offset from the SDSS-reported heliocentric radial velocity (HRV). Despite this lack of agreement with SDSS, the HIRES and COS spectra give self-consistent results for BHB6. The remainder of our analysis does not use SDSS spectra, so if there is an error in the SDSS spectra, it has no effect on our analysis. Overall, Figure 2 indicates that stellar radial velocities inferred from our HIRES and COS spectra are largely consistent within typical wavelength calibration errors for each instrument. Hidden systematic errors in the COS G160M wavelength solution that have been reported elsewhere (e.g., Wakker et al. 2012) do not appear to significantly impact our BHB star COS spectra.

2.4. EBHIS H I 21 cm Profiles

H I 21 cm emission can provide constraints on the distribution of neutral hydrogen in the general vicinity of our BHB sight lines. However, direct comparisons of measured column densities in absorption and emission data are ill-advised for two reasons. First, the spatial resolution of the H I data is $10''/8$, orders of magnitude larger than the much narrower $\sim 1''$ HIRES slit or the $2''/5$ circular COS aperture. This disparity implies that

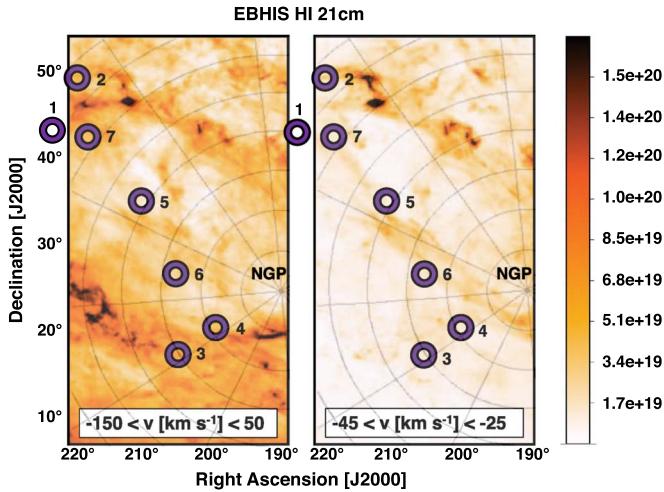


Figure 3. Integrated H I 21 cm column density maps from EBHIS (beam FWHM = 10'8; Winkel et al. 2016) near the north Galactic pole (NGP) in the vicinity of the seven BHB targets for two velocity ranges: a broad range, $-150 \text{ km s}^{-1} < v_{\text{LSR}} < 50 \text{ km s}^{-1}$ (left), and a narrower range, centered on the bulk of the IV gas, $-45 \text{ km s}^{-1} < v_{\text{LSR}} < -25 \text{ km s}^{-1}$ (right). The NGP is marked on each image, and the grid curves show rings of constant Galactic latitude (b) in 5° increments, while the diagonal lines show constant Galactic longitude (l) in 30° increments (BHB stars 3 and 4 lie near the $l \approx 0^{\circ}$ line). Additionally, the approximate R.A. and decl. in degrees are given on the x - and y -axes, respectively. The location of each BHB is marked with a purple circle (unfortunately, BHB1 lies just off the edge of the processed images). The high-latitude BHB stars lie on the very low l edge ($0^{\circ} < l < 80^{\circ}$) of the well-known IV structure known as the “IV arch” covering a large angular area on the sky, constrained along several lines of sight to have $0.4 \text{ kpc} < d < 1.8 \text{ kpc}$ (e.g., Wakker et al. 2001).

emission and absorption measurements do not trace material on the same spatial scales. Second, there are only radial velocity and no distance constraints for material detected in 21 cm emission. Nonetheless, these H I data are useful because of their all-sky coverage and their excellent constraints on the large-scale structure of the IV material we wish to examine.

We use data from the Effelsberg-Bonn H I Survey (EBHIS; Winkel et al. 2016), which images the full sky with a bandwidth of 100 MHz. The effective beam FWHM is 10'8, and the effective velocity resolution is 1.44 km s^{-1} . We show in Figure 3 integrated H I 21 cm column density maps from EBHIS over two velocity ranges (left: a broad velocity range $-150 \text{ km s}^{-1} < v_{\text{LSR}} < 50 \text{ km s}^{-1}$; right: the velocity range over which we detect the bulk of our IV absorption, $-45 \text{ km s}^{-1} < v_{\text{LSR}} < -25 \text{ km s}^{-1}$) in the approximate location of our sample of BHB stars. Figure 3 shows that none of our BHB lines of sight intersect the dense neutral clumps visible in this image, and instead seem to trace the wispy structures that extend across tens of degrees.

At the position of each BHB, we download the H I spectra that are generated using a weighted interpolation with a Gaussian kernel from the H I Survey server of the Argelander-Institut für Astronomie.¹¹ We show these EBHIS H I 21 cm brightness temperature profiles as a function of v_{LSR} on the right column of Figure 4. IV neutral gas with $v_{\text{LSR}} < -25 \text{ km s}^{-1}$ is apparent along each of our seven lines of sight. In order to determine the H I column density in the IV gas along each line of sight, we perform simple, dual Gaussian fits to these double-peaked, asymmetric profiles seen in Figure 4. In the optically thin limit, the H I column density is directly

proportional to the integrated 21 cm emission line brightness (T_b), which we calculate using the standard relation (Draine 2011):

$$\frac{N_{\text{H}_1}}{\text{cm}^{-2}} \approx 1.82 \times 10^{18} \int \left[\frac{T_b(v)}{K} \right] d\left(\frac{v}{\text{km s}^{-1}} \right). \quad (1)$$

We use the FWHM, velocity centroid, and peak T_b from the simple Gaussian fits to the IV components to determine their column densities within the 10'8 beam. The relatively large beam size gives the average N_{H_1} on scales of 9.4 pc at a distance of $z = 3 \text{ kpc}$ [or $l \approx (3.1 \text{ pc}) \times (z/1 \text{ kpc})$]. Smaller-scale structures will be averaged out, including any unresolved structures probed by the absorption-line data along the BHB sight lines. Thus, values of N_{H_1} determined from EBHIS data cannot be used as either upper or lower limits to the neutral gas column density associated with the absorption we see in Keck/HIRES or HST/COS spectra. However, the variation in N_{H_1} over the area traced by our sight lines will prove useful for our analysis.

3. Line Profile Analysis

Lines of sight toward high Galactic latitude halo stars pass through MW disk gas; extraplanar, ionized “thick-disk” gas at the disk–halo interface; and the inner MW circumgalactic medium (CGM; Savage et al. 2003; Savage & Wakker 2009; Zheng et al. 2019). Figures 4 and 5 show sections of the Keck/HIRES and HST/COS spectra that cover the ionized and neutral transitions we analyze in this work. These continuum-normalized spectra are given as a function of v_{LSR} , and the region we define as tracing the MW ISM is highlighted in light green, $-20 < v_{\text{LSR}} (\text{km s}^{-1}) < 20$. The data themselves are shown as a black histogram, while the line profile fits are shown as colored curves and are fully described later in this section. We adopt terminology such that absorption features with LSR absolute velocities within $\sim 20\text{--}90 \text{ km s}^{-1}$ of $v_{\text{LSR}} = 0 \text{ (km s}^{-1}\text{)}$ are referred to as “IV” clouds. Features with absolute radial velocities $v_{\text{LSR}} > 90 \text{ km s}^{-1}$ are generally referred to as HVCs (Wakker et al. 2003); however, we do not detect any HVCs in our data set.

Thus, our study focuses on extraplanar IV material, which must be understood within a rich context of previous work. Figure 1 shows our sight lines, along with the locations of well-known IV structures identified via H I 21 cm maps. At a distance of $< 3.5 \text{ kpc}$ (Ryans et al. 1997), from $l = 80^{\circ}\text{--}220^{\circ}$ and $b = 30^{\circ}\text{--}80^{\circ}$, lies the $\sim 10^5 M_{\odot}$ IV arch, moving toward us at $40\text{--}90 \text{ km s}^{-1}$. This is the closest IV structure to our lines of sight, which lie just on its edge. Other less massive IV structures include the IV spur and the IV arch extension (LLIV), both of which exhibit scale heights $< 2 \text{ kpc}$ (Wakker 2001, 2004) and appear to be accreting at a mean velocity of $\approx -50 \text{ km s}^{-1}$. For reference, lower and upper limits on the distances to most HVCs and IVCs typically place them in the inner ($1 \text{ kpc} < d < 10 \text{ kpc}$) MW halo (Putman et al. 2012).

Along each halo star sight line, there is absorption from the MW ISM immediately visible in low ionization state and neutral gas at approximately the LSR velocity ($v_{\text{LSR}} = 0 \text{ km s}^{-1}$) in its direction. Additionally, we often detect absorption lines of neutral and low-ion metals roughly coincident with the BHB star’s radial velocity reported by SEGUE. The extent to which there are shifts between the centers of the beige bands and the line centroids of the stellar absorption in COS and Keck data highlights the uncertainties in the wavelength calibration that

¹¹ https://www.astro.uni-bonn.de/hisurvey/AllSky_profiles/

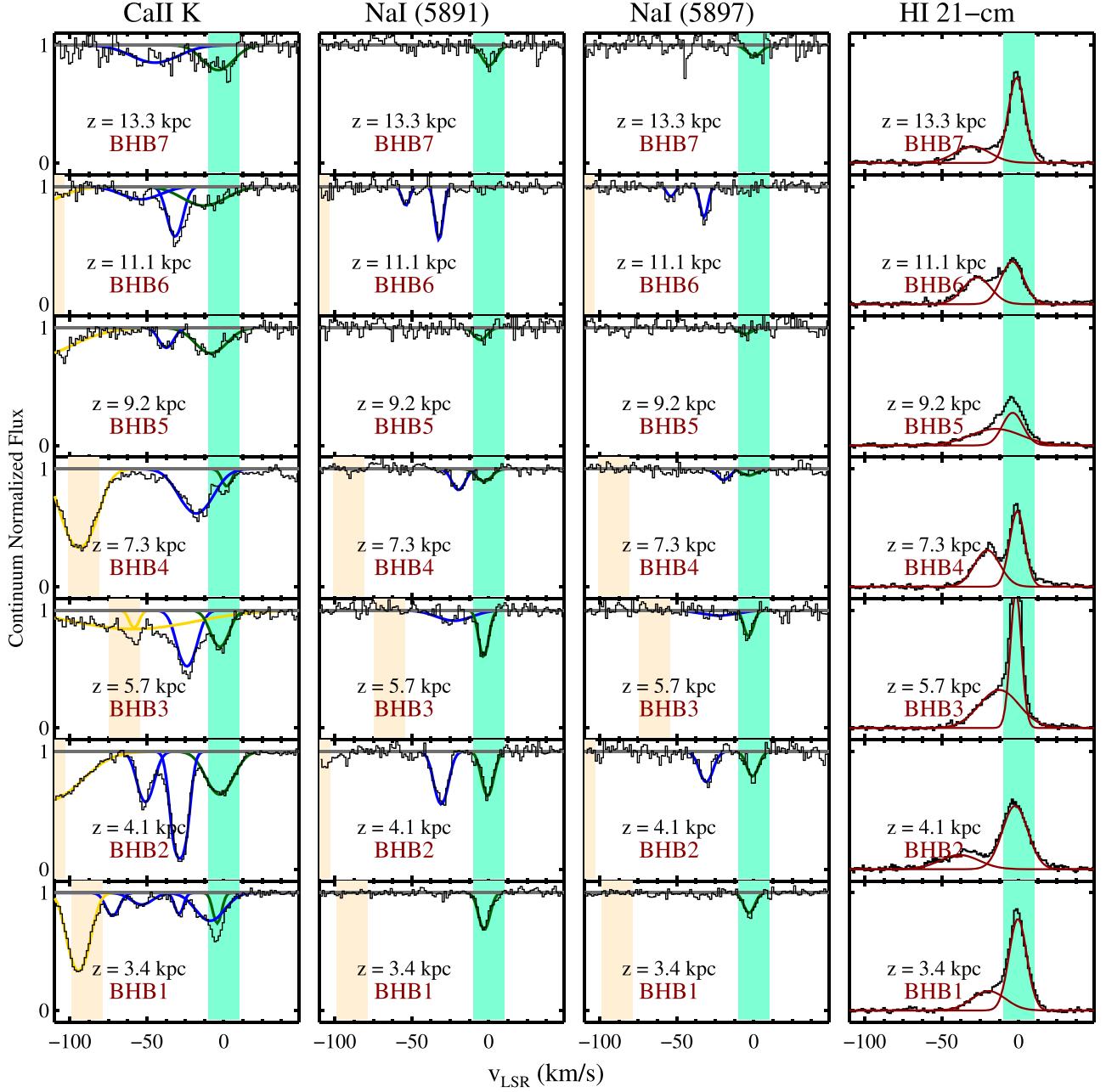


Figure 4. The first three columns show absorption-line stacks of the Na I D doublet and Ca II K drawn from the Keck/HIRES spectra. The velocities are relative to v_{LSR} , and the spectra have been continuum normalized. The rightmost column shows the EBHIS H I 21 cm brightness temperature as a function of v_{LSR} . The peak brightness temperature, T_b , on the y-axis of the rightmost column is ≈ 5.3 K. The 25 km s^{-1} wide, green highlighted band centered on Galactic emission at $v_{\text{LSR}} = 0 \text{ km s}^{-1}$ matches the strongest H I feature in every case. In the Keck spectra, the beige highlighted band marks the radial velocity of the BHB as v_{LSR} from SEGUE, for reference. We detect broad stellar absorption features in Ca II, but not in Na I. The blueshifted features relative to the MW ISM represent gas along the line of sight to the BHB, which has its height from the disk given in the lower left corner of each panel. Our Voigt profile fits are shown in different colors in each panel. Features identified as IV absorption are shown in blue, while components we associate with the MW ISM are shown in dark green. Two-component Gaussian fits to the H I brightness temperature plots are shown in deep red in the rightmost panels.

manifest as differences with the SEGUE stellar radial velocities in the LSR frame, shown in Figure 2. Together, the MW ISM and stellar absorption do not account for the total observed absorption along each BHB line of sight. In the Keck/HIRES spectra, there are kinematically distinct absorption lines owing to an additional IV gas component(s). In the *HST/COS* spectra, the spectral resolution is too low to completely separate IV gas from MW ISM, although we may assume that both gaseous components are present, as there is absorption over the full velocity range covered by the multiple components seen in the

Keck/HIRES spectra. In this section, we describe our approach to a Voigt profile fitting analysis with the goal of examining the extraplanar “IV” gas properties in both Keck/HIRES and *HST/COS* spectra.

Before proceeding, we note two general caveats important for interpreting our results. First, in principle, distinguishing extraplanar and/or halo gas clouds from a coherent MW ISM is accomplished by imposing a velocity cut in the LSR frame on absorption seen along the line of sight as described above. However, gas with velocities inconsistent with MW differential

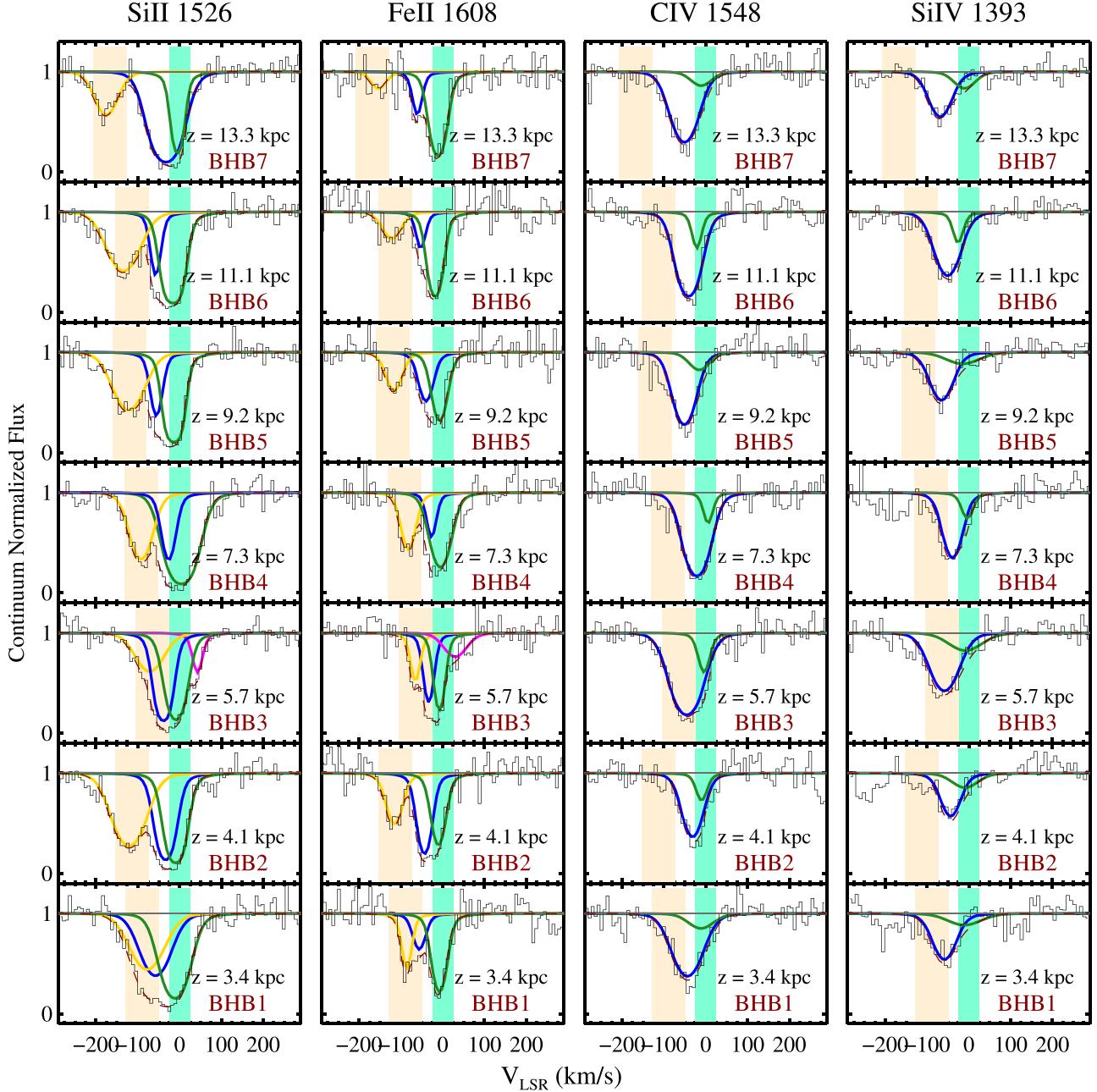


Figure 5. Ionic species drawn from the *HST/COS* absorption-line spectra of the seven target BHB stars, centered on the transitions relevant to this study: Si II, Fe II, C IV, and Si IV. The Voigt profile fits for individual components are shown with thick colored lines, while the total normalized line profile fit is shown with a dashed line over the detected features. Blue lines highlight the IV absorption, gold lines highlight stellar absorption, and green lines show the MW components. We show the SEGUE radial velocity of the BHB star in v_{LSR} , for reference, as a beige shaded area. Stellar absorption is present and fitted in the low-ionization transitions of Si II and Fe II, while it is absent from the intermediate ions. The central green shaded region is meant to guide the eye to $v_{\text{LSR}} = 0 \text{ km s}^{-1}$. The low ions in BHB3 distinctly show a weak redshifted component, highlighted in pink, although we do not discuss redshifted material in this work.

rotation technically need not be spatially distinct from the MW ISM. Second, extraplanar material detected in absorption at coincident line-of-sight velocities may not itself be entirely cospatial in a 3D sense. This last caveat is one to which we will return when comparing the column densities of intermediate- and low-ionization transitions.

Voigt profile fitting is an indispensable tool used to model and interpret absorption-line data and is the best method for “deblending” stellar absorption, MW ISM, and IV gas. Generally, these profiles are generated using rules-based algorithms that make assumptions about the number of absorption components to fit, and they sometimes demand correspondence between

component structures in different ionic species. Each of these assumptions introduces a prior that can constrain parameter estimates in ways that are statistically valid but do not capture the full physical uncertainty of the models. In particular, with COS, one has to worry about unresolved saturation effects, which are notorious for introducing specious column densities. Furthermore, higher spectral resolution tends to resolve more and more component structure, which is apparent in visually comparing the higher-resolution HIRES with the lower-resolution COS absorption profiles. The spectral resolution and S/N of the data therefore set clear limits as to how many components one can fit and to what extent one can trust the derived fits.

The standard approach to Voigt profile fitting is to fit as few components as required by the data for a reasonably good fit, with good being defined in a number of possible ways. In Section 3.2, we detail all of the assumptions that go into our line profile analysis. By adopting a consistent approach to Voigt profile fitting of absorption along each line of sight for every ion observed in our COS spectra, we may understand the ionization structure of the material in the MW’s disk–halo interface. Specifically, we constrain the properties of the IV gas detected in the *HST/COS* spectra given a set of kinematic priors that come from higher spectral resolution Keck/HIRES spectra.

3.1. Qualitative Assessment and Overview

The FWHM of a single resolution element in our COS G160M spectra is more than two times broader than that of our HIRES spectra (18 km s^{-1} vs. 7 km s^{-1}). As a result, the MW ISM is blended with absorption owing to the IV extraplanar features we aim to examine in the COS data. The HIRES data allow a cleaner separation of MW local ISM and additional absorption features, seen in Figure 4. The Ca II K absorption at $v_{\text{LSR}} \approx 0 \text{ km s}^{-1}$ lines up with the narrow, strong emission at $v_{\text{LSR}} \approx 0 \text{ km s}^{-1}$ visible in the EBHIS H I 21 cm data. Consistent with previous studies of Na I that report its patchy detection in the MW, the $v_{\text{LSR}} \approx 0 \text{ km s}^{-1}$ component of Na I absorption appears in only four of seven cases (e.g., Meyer & Lauroesch 1999). In addition, absorption consistent with the BHB stellar radial velocity is visible in Ca II K, C I, Si II, Al II, and Fe II (Ca II H is catastrophically blended with broad stellar He absorption). We do not see stellar absorption in Na I, Si IV, and C IV in our HIRES or COS data for the targeted BHB stars.

For all BHB stellar sight lines, a simple visual inspection of the line profiles for the “intermediate ions,” Si IV and C IV, compared to those of the “low ions,” i.e., the neutral and singly ionized species, can instruct our profile fitting. In the COS spectra, the blended intermediate ions consistently exhibit broader line widths and more highly blueshifted velocity centroids ($\gtrsim 25 \text{ km s}^{-1}$) than the blended low ions. This apparent velocity offset most likely arises because the contribution of the predominantly neutral, higher-density, $v_{\text{LSR}} \approx 0 \text{ km s}^{-1}$, MW ISM to the line profile is much greater in low ions than it is for intermediate ions. Adopting this assumption leads directly to the inference that extraplanar IV gas must exhibit a higher ionization parameter than the MW ISM. The precise quantitative difference in ionization parameter (and density, by extension) between the MW and IV gas depends strongly on the total H I column of each component, the gas-phase metallicity of each component, and the background ionizing spectrum affecting each component. We note that with our current data set we cannot independently constrain any of these quantities.

3.2. A Fitting Procedure to Isolate Intermediate-velocity Gas from Milky Way Absorption

Our iterative fitting program for absorption lines detected with both Keck/HIRES and COS uses the MPFIT software¹² (Markwardt 2009). The line profiles we derive from Voigt profile fitting of the COS data are convolved with the COS LSF at LP3 as given at the nearest observed wavelength grid point

in the compilation by Ghavamian et al. (2009). Different transitions of the same ionic species (e.g., Si IV $\lambda 1393$ and Si IV $\lambda 1402$) are required to have the same component structure and are therefore fit simultaneously to give a single solution. Each fit gives a value of the component column density N , Doppler b , and velocity centroid v in the LSR frame. The best fits emerge after a χ^2 minimization process implemented in our code and depend strongly on the number of components assumed to be present. Reported errors in fitted parameters and overall goodness of fit via reduced χ^2 are not independent of the assumptions in the analysis or the limitations of the data (e.g., due to saturation).

For the Keck/HIRES line profile analysis, we adopt a minimum number of components that results in profile fits to the data such that the reduced $\chi^2_{\nu} < 1.1$. Not including stellar absorption, we fit an average of two to three Ca II components per sight line in the Keck data. Stellar absorption is seen in Ca II (yellow fits shown in Figure 4) and is generally well separated from the MW and IV gas. Each sight line has a distinctly separable MW component visible in Ca II, sometimes detected in Na I. When Na I is detected in either IV or MW gas, its components’ velocity alignment with the Ca II components is excellent.

In our analysis of the COS spectra, we proceed under the assumption that both MW ISM and IV components are present with varying contributions in the low and intermediate ions. The main reason for “forcing” a two-component fit is to be able to compare low and intermediate ionic column densities directly without worrying about different levels of contamination from an MW component. This assumption is physically motivated and supported by the Keck/HIRES spectra, which show a distinct MW component in Ca II and H I along all sight lines. To ensure that two components are present in the COS data, our fitting procedure imposes lower limits on the column density of the MW components (details below). If these limits are not imposed, the iterative Voigt profile fits sometimes converge on a solution such that a single component dominates the result (i.e., the minimized second component contributes a column $< 10^9 \text{ cm}^{-2}$). In general, the precise values of our imposed lower limits, as long as they are not unreasonably large so as to be completely inconsistent with the data, have very little impact on the final fit parameters. Finally, we note that our results, which focus primarily on IV Si IV and C IV, depend very little on the fitting assumptions described in this section.

We present the component fits for each line of sight for select transitions in both Keck and COS spectra in Figures 4 and 5. We tabulate our component fitting results for the column density N , Doppler parameter b , and velocity offset v from v_{LSR} , along with associated errors, in Tables 3 and 4 for the MW and IV components, respectively. We also show the results from the dual Gaussian fits to the H I 21 cm T_b profiles; in this case the b column is in fact a Gaussian σ and not a Doppler b parameter. The Voigt profile fits give typical values of χ^2 ranging from 200 to 400, corresponding to values of reduced $\chi^2_{\nu} < 1.4$. Our best constraints on the MW low-ion column densities in the COS data are from Fe II, which is the least saturated low ion. Fits for MW C IV and Si IV are poorly constrained by the data, and thus their best-fit parameters are often given as limits. Below, we describe our COS absorption-line fitting procedure and all associated assumptions in detail.

¹² <http://cow.physics.wisc.edu/~craigm/idl/fitting.html>

Table 3
BHB Milky Way ISM Component Fits

BHB ID	Ion	$\log N_{\text{ion}}$	v_{LSR}	b
BHB1	Ca II	11.48 \pm 0.06	-4.24 \pm 0.20	3.84 \pm 0.36
	Fe II	14.57 \pm 0.08	-10.30 \pm 2.27	16.36 \pm 3.50
	Si II	14.50 \pm 0.16	<-14.24	35.43 \pm 3.78
	Al II	13.33 \pm 0.11	-12.41 \pm 3.24	22.30 \pm 4.72
	Si IV	<12.80	<-14.24	>50.00
	C IV	<13.20	<-14.24	37.06 \pm 21.44
	H I	19.97 \pm 0.01	-0.49 \pm 0.10	8.16 \pm 0.14
BHB2	Ca II	12.14 \pm 0.01	-2.65 \pm 0.19	11.64 \pm 0.28
	Fe II	14.44 \pm 0.07	<-12.65	16.48 \pm 2.73
	Si II	14.65 \pm 0.30	<-12.65	22.33 \pm 3.86
	Al II	13.81 \pm 1.75	<-12.65	12.83 \pm 11.73
	Si IV	<12.80	<-12.65	37.25 \pm 20.33
	C IV	<13.20	<-12.65	13.05 \pm 6.84
	H I	19.97 \pm 0.01	-2.30 \pm 0.14	11.72 \pm 0.20
BHB3	Ca II	11.89 \pm 0.00	-2.60 \pm 0.37	8.00 \pm 0.42
	Fe II	14.64 \pm 0.38	-10.39 \pm 3.57	<10.00
	Si II	<14.50	<-12.60	23.72 \pm 5.05
	Al II	13.27 \pm 0.37	<-12.60	15.19 \pm 8.21
	Si IV	<12.80	<-12.60	>50.00
	C IV	13.42 \pm 0.34	-6.80 \pm 4.39	13.67 \pm 9.68
	H I	19.94 \pm 0.01	-1.78 \pm 0.04	5.09 \pm 0.08
BHB4	Ca II	11.25 \pm 0.08	1.92 \pm 0.43	4.38 \pm 0.67
	Fe II	14.57 \pm 0.06	<-8.08	26.22 \pm 3.75
	Si II	14.68 \pm 0.22	0.85 \pm 13.57	38.13 \pm 7.83
	Al II	<13.50	<-8.08	28.61 \pm 2.47
	Si IV	<12.80	-7.06 \pm 3.81	12.88 \pm 7.16
	C IV	<13.20	2.91 \pm 5.31	<10.00
	H I	19.82 \pm 0.01	-0.90 \pm 0.11	7.01 \pm 0.14
BHB5	Ca II	11.94 \pm 0.03	-8.31 \pm 0.81	13.77 \pm 1.19
	Fe II	<14.80	-9.96 \pm 2.22	18.52 \pm 4.09
	Si II	15.17 \pm 0.83	-16.57 \pm 3.84	15.67 \pm 6.07
	Al II	14.46 \pm 4.24	<-18.31	13.76 \pm 18.28
	Si IV	<12.80	<-18.31	>50.00
	C IV	<13.20	<-18.31	29.05 \pm 17.77
	H I	19.59 \pm 0.02	-3.91 \pm 0.19	9.55 \pm 0.35
BHB6	Ca II	11.95 \pm 0.08	-12.74 \pm 2.89	19.12 \pm 3.41
	Fe II	<14.80	-20.31 \pm 2.35	15.40 \pm 2.36
	Si II	14.85 \pm 0.24	-18.64 \pm 3.14	21.21 \pm 4.44
	Al II	13.42 \pm 0.14	-22.30 \pm 7.58	25.99 \pm 6.21
	Si IV	12.87 \pm 0.19	<-22.74	<10.00
	C IV	13.27 \pm 0.37	-22.68 \pm 7.35	<10.00
	H I	19.74 \pm 0.01	-4.15 \pm 0.28	10.35 \pm 0.34
BHB7	Ca II	11.92 \pm 0.04	-3.63 \pm 1.09	13.33 \pm 1.54
	Fe II	14.73 \pm 0.12	<-13.63	18.58 \pm 3.22
	Si II	14.85 \pm 1.16	-7.46 \pm 6.21	<10.00
	Al II	15.50 \pm 0.46	<-13.63	10.44 \pm 1.86
	Si IV	<12.80	<-13.63	>50.00
	C IV	<13.20	<-13.63	>50.00
	H I	19.92 \pm 0.01	-1.26 \pm 0.09	7.78 \pm 0.13

Note. Voigt profile fit parameters for the MW ISM components for a variety of ionic species observed along each BHB sight line. The parameters are column density as $\log N$ (cm^{-2}), velocity centroid v_{LSR} (km s^{-1}), and Doppler parameter b , which quantifies the line width in km s^{-1} . We also tabulate the results from the Gaussian fits to the H I 21 cm T_b profiles; in this case the b column is given as the Gaussian $\sigma \times \sqrt{2}$.

1. *Identifying the stellar absorption by SEGUE radial velocity.* As described in Section 2.3, the low-ion transitions in COS show absorption roughly consistent

Table 4
BHB Intermediate-velocity Component Fits

BHB ID	Ion	$\log N_{\text{ion}}$	v_{LSR}	b
BHB1	Ca II	11.95 \pm 0.02	-8.85 \pm 0.55	12.58 \pm 0.66
	Ca II	11.30 \pm 0.04	-29.11 \pm 0.24	3.87 \pm 0.36
	Ca II	11.43 \pm 0.04	-53.27 \pm 0.76	9.84 \pm 1.19
	Ca II	11.52 \pm 0.03	-72.30 \pm 0.30	5.90 \pm 0.44
	Fe II	13.89 \pm 0.27	-58.63 \pm 7.61	16.20 $^{+17.42}_{-15.20}$
	Si II	<14.22	>-60.59	>50.00
	Al II	12.98 \pm 0.33	-59.41 \pm 2.89	10.05 \pm 7.20
BHB2	Si IV	13.35 \pm 0.04	-61.80 \pm 2.75	31.39 \pm 4.18
	C IV	14.04 \pm 0.03	-46.75 \pm 2.71	41.39 \pm 3.03
	H I	19.57 \pm 0.04	-19.45 \pm 1.11	15.34 \pm 1.30
	Ca II	12.56 \pm 0.01	-28.60 \pm 0.05	5.86 \pm 0.06
	Ca II	12.02 \pm 0.01	-50.83 \pm 0.13	7.14 \pm 0.17
	Fe II	14.62 \pm 0.09	-46.07 \pm 2.42	15.49 \pm 3.34
	Si II	<14.50	>-36.40	22.30 \pm 21.77
BHB3	Al II	<13.34	>-38.84	25.39 \pm 24.84
	Si IV	13.26 \pm 0.04	-46.91 \pm 3.67	26.27 \pm 4.37
	C IV	13.92 \pm 0.03	-34.19 \pm 1.60	27.07 \pm 2.21
	H I	19.51 \pm 0.03	-38.40 \pm 0.75	18.12 \pm 1.16
	Ca II	12.14 \pm 0.02	-23.96 \pm 0.23	8.20 \pm 0.37
	Fe II	14.41 \pm 0.26	-35.92 \pm 3.51	<10.00
	Si II	14.65 \pm 0.53	-41.45 \pm 10.73	17.91 \pm 14.57
BHB4	Al II	<14.06	>-42.09	<10.00
	Si IV	13.58 \pm 0.02	-61.84 \pm 2.09	39.11 \pm 3.07
	C IV	14.30 \pm 0.05	-48.63 \pm 4.44	40.57 \pm 3.82
	H I	19.94 \pm 0.01	-12.62 \pm 0.47	18.34 \pm 0.42
	Ca II	12.25 \pm 0.01	-17.97 \pm 0.33	14.36 \pm 0.49
	Fe II	13.95 \pm 0.32	-29.32 \pm 6.49	<10.00
	Si II	14.05 \pm 0.64	>-28.92	15.98 \pm 13.65
BHB5	Al II	<13.41	>-22.05	16.09 \pm 24.48
	Si IV	13.54 \pm 0.03	-42.39 \pm 1.77	23.98 \pm 2.90
	C IV	14.27 \pm 0.03	-24.32 \pm 1.80	32.37 \pm 2.05
	H I	19.73 \pm 0.02	-20.40 \pm 0.31	11.71 \pm 0.47
	Ca II	11.51 \pm 0.06	-37.40 \pm 0.73	6.57 \pm 1.04
	Fe II	14.12 \pm 0.12	-42.46 \pm 4.24	18.75 \pm 6.56
	Si II	13.97 \pm 0.28	-58.31 \pm 7.22	12.88 \pm 8.44
BHB6	Al II	13.46 \pm 0.41	-35.50 \pm 13.82	28.80 \pm 6.80
	Si IV	13.38 \pm 0.03	-68.77 \pm 2.10	31.21 \pm 3.18
	C IV	14.07 \pm 0.02	-53.98 \pm 1.91	29.28 \pm 2.40
	H I	19.69 \pm 0.01	-14.83 \pm 1.02	23.52 \pm 0.76
	Ca II	11.98 \pm 0.06	-31.69 \pm 0.32	6.59 \pm 0.55
	Ca II	11.72 \pm 0.09	-53.88 \pm 2.75	17.50 \pm 3.78
	Fe II	13.81 \pm 0.19	-54.75 \pm 5.47	<10.00
BHB7	Si II	13.98 \pm 0.22	-59.95 \pm 3.23	<10.00
	Al II	12.97 \pm 0.69	-61.79 \pm 5.87	<10.00
	Si IV	13.61 \pm 0.04	-54.30 \pm 2.46	34.11 \pm 2.08
	C IV	14.27 \pm 0.04	-44.00 \pm 1.92	29.75 \pm 1.85
	H I	19.63 \pm 0.02	-26.95 \pm 0.52	12.73 \pm 0.68
	Ca II	11.94 \pm 0.04	-44.74 \pm 2.12	21.32 \pm 0.97
	Fe II	13.93 \pm 0.11	-63.15 \pm 3.01	10.27 \pm 5.44
BHB7	Si II	14.66 \pm 0.06	-35.81 \pm 2.57	39.64 \pm 2.23
	Al II	12.99 \pm 0.07	-64.94 \pm 3.74	19.82 \pm 5.69
	Si IV	13.34 \pm 0.03	-71.76 \pm 2.08	31.10 \pm 3.21
	C IV	14.10 \pm 0.02	-53.76 \pm 1.56	36.50 \pm 2.26
	H I	19.52 \pm 0.03	-30.00 \pm 0.67	16.38 \pm 1.06

Note. Voigt profile fit parameters for the IV components for a variety of ionic species observed along each BHB sight line. The parameters are column density as $\log N$ (cm^{-2}), velocity centroid v_{LSR} (km s^{-1}), and Doppler parameter b , which quantifies the line width in km s^{-1} . We also tabulate the results from the Gaussian fits to the IV component of the H I 21 cm T_b profiles; in this case the b column is given as the Gaussian $\sigma \times \sqrt{2}$.

with SEGUE-reported stellar radial velocities. When present, we fit this stellar component allowing N , b , and v_{cen} to vary freely, and the results are consistent with expectations from HIRES spectra in which stellar absorption in Ca II is cleanly separated from the MW and IV components. In Figures 4 and 5, the stellar component is colored in gold. In one case, BHB3, absorption centered on the expected stellar velocity in Ca II demands two components, one very broad and one weak, narrow component. Physically, both may not be associated with the star, but we proceed by simply noting that the IV column density for Ca II along this sight line may be an underestimate.

2. *Defining the MW component in velocity space.* For every BHB line of sight observed with COS, we fit one MW ISM component within $\pm 10 \text{ km s}^{-1}$ of the LSR velocity for the Ca II K MW line fit as defined by the HIRES spectra ($v_{\text{cen,Ca II}}$). We force this MW component to have a minimum column density as defined below. In this analysis, we do not use the Na I line directly, since an MW component is not consistently detected in Na I owing to its rather dramatic patchiness in the ISM (e.g., Münch & Zirin 1961). However, we note that when Na I is detected in the HIRES spectra at $v_{\text{LSR}} \approx 0 \text{ km s}^{-1}$, its velocity centroid is consistent with Ca II MW absorption within a few kilometers per second. In Figures 4 and 5, the MW component is colored green. Table 3 shows that in approximately 50% of the cases our MW Voigt profile fits do not fully converge because they are constrained not to venture outside of the $\pm 10 \text{ km s}^{-1}$ band. In these cases, the values for the velocity centroid of the MW component are reported as limits: e.g., $v_{\text{cen}} < v_{\text{cen,Ca II}} - 10 \text{ km s}^{-1}$.
3. *Lower limit on the Low-ion MW component column densities.* The combined IV and MW components of the strong low-ion transitions detected in the COS spectra (Si II, Al II, Fe II) show signs of saturation along every line of sight. Due to its non-Gaussian LSF, COS is known to saturate at normalized fluxes of ~ 0.2 . In our fitting procedure, we assume that the MW ISM dominates the column density contribution to the observed *saturated* absorption features near $v_{\text{LSR}} \approx 0 \text{ km s}^{-1}$. This assumption corresponds to a minimum column density contribution from each transition, which we can compute using its oscillator strength and rest wavelength, and assuming that the lines are unresolved, with $b \approx 20 \text{ km s}^{-1}$. To compute this lower limit on the MW column density, we use the XIDL routine “ism_pltlin.pro.”¹³ Lower-limit $\log N [\text{cm}^{-2}]$ values for Si II, Fe II, and Al II, respectively, are determined to be 14.5, 14.8, and 13.3. In practice, this limit serves to demand a dominant MW component for the neutral, low ionization state ISM. Our minimization procedure converges to this MW column density lower limit for the low ions in only three cases seen in Table 3. Because the Voigt profile fits would prefer a lower column density than we allow, we report $\log N_{\text{ion}}$ in these three cases as conservative upper limits.
4. *Setting limits on the intermediate-ion MW component column densities.* Our primary assumption is that there are two *detected* components along the BHB lines of

sight corresponding to MW ISM and IV gas in the halo. For the intermediate ions that do not have a dominant MW component visible in the data, we use the XIDL routine “ism_pltlin.pro” to determine the column density floors for Si IV and C IV. The median equivalent width 2σ detection limits computed on each COS spectrum are 36 and 51 mÅ for Si IV and C IV, respectively. Corresponding $\log N [\text{cm}^{-2}]$ values for Si IV and C IV are determined to be 12.8 and 13.2. We artificially impose the requirement that the column densities not go below these upper limits when we fit for two velocity components. In practice, this limit serves to demand a detected but weak MW component for the ISM in intermediate ionization states. The fits for the majority of the sight lines hit this lower column density limit in the intermediate-ionization species because there are simply not easily separable MW components visible in the data. Indeed, most of the Si IV and C IV prefers a single component fit that is very similar to the reported IV fits in Table 4, with a slightly higher column density. Nonetheless, to present a consistent analysis, we demand an MW component in Si IV and C IV with the above minimum columns. Thus, in the tables, the reported fits are given as *upper* limits, since the Voigt profiles would prefer to converge on an arbitrarily low column density for the MW intermediate ions.

The above steps define the primary assumptions that set our best two-component IV + MW fits in the COS spectra. In addition, we assume that each component in the COS spectra cannot have a Doppler b parameter $< 10 \text{ km s}^{-1}$ or $> 50 \text{ km s}^{-1}$. These limits ensure that the iterative fitting procedure does not simply minimize a second component by giving it a very small or large line width. We define a lower limit of 10 km s^{-1} to be the narrowest resolvable line given the COS line FWHM of 18 km s^{-1} . For reference, the relation between the two quantities is given by $\text{FWHM} \approx 1.665b$. The precise values of the imposed Doppler b parameter limits make very little difference to the IV and MW column density values, and the cases in which our profile fits hit these limits are mostly for the poorly constrained MW components. In the following sections, we primarily focus on the properties of the IV component. In general, this procedure produces well-constrained column densities and velocities for the IV components tabulated in Table 4, which are the focus of this study.

4. Results

4.1. Lack of Trends with Height from the Disk

The estimated distance and position of each BHB in our sample allow us to place an upper limit on the distance and height to the observed IV gas along the line of sight. Examining how the measured absorption properties (N , v , Doppler b) vary with both distance upper limits and height upper limits may provide insight into the 3D distribution of the gas. For example, an increase in column density with z would signify gas located at the full range of scale heights probed by our BHB stars ($3 \text{ kpc} < z < 14 \text{ kpc}$). However, we find no significant correlation between any of these quantities and either distance or height. We note, however, that our data are not sensitive to variations in column density below 3.4 kpc. Below, we limit our discussion to the Galactic-latitude-corrected “vertical” projections (e.g., $\log N \sin b$, $v \sin b$) and height.

¹³ https://github.com/profxj/xidl/blob/master/ISM/ism_pltlin.pro

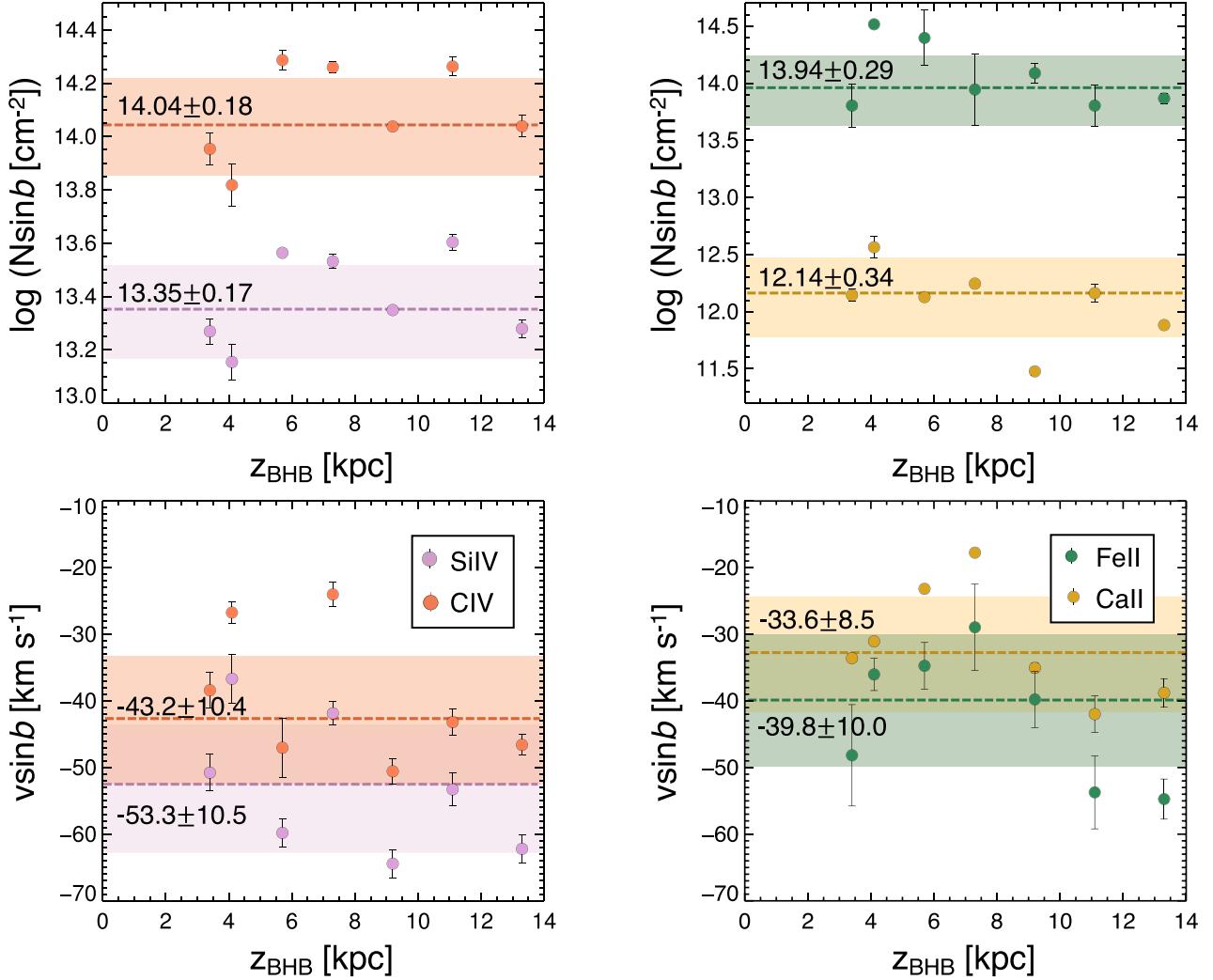


Figure 6. Lack of trends with maximum height, given as z , from the MW disk for derived properties of IV gas traced by both intermediate ionization state species (Si IV: plum; C IV: coral; left panels) and low ionization state species (Fe II: green; Ca II: gold; right panels). The top two panels examine vertical column density as a function of z , with the correction factor for Galactic latitude applied. The bottom panels show variations in the velocity centroids of the IV components, again corrected for Galactic latitude, b . Median values, shown as dashed lines, with standard deviations, delineated by the shaded regions, of $\log N \sin b$ and $v \sin b$ are shown in the corresponding colors on each panel. There are no significant trends with height for these quantities (regardless of the $\sin b$ scaling), implying that the majority of the gas in absorption traced by these observations lies at $z \lesssim 3.4$ kpc.

The top panels of Figure 6 show that the vertical component of the IV gas column density in every ionization state, from Ca II to C IV, remains roughly constant over the range of heights probed by our observations. Each panel shows a median of the value for $\log N \sin b$ or $v \sin b$ for each ion in the IV component only. For the low ions, we show only Ca II and Fe II IV absorption properties since these two ions have the best-constrained IV Voigt profile fits. There is only small scatter of ± 0.2 dex for the intermediate ions (top left) and a more moderate scatter of ± 0.3 – 0.4 dex for the low ions (top right). This narrow span of low- and intermediate-ion column densities indicates that the majority of the gas traced by absorption lies at heights $\lesssim 3.4$ kpc.

The lack of any discernible increase in column density with z above 3.4 kpc is in agreement with the analysis of Savage & Wakker (2009). In their study of numerous halo star and QSO sight lines, they determined an exponential scale height of a given ion by comparing column densities projected along the

z -axis to a simple “plane-parallel” model of column densities with an exponential decline in the z -direction. In this plane-parallel framework, they find that Si IV and C IV have similar exponential scale heights of 3.2 and 3.6 kpc, with typical errors on the order of ± 1 kpc. Low ions tend to exhibit much smaller exponential scale heights < 1 kpc.

The median values of $\log N \sin b$ of both Si IV and C IV along our BHB sight lines agree with the median $\log N \sin b$ reported by Savage & Wakker (2009) for extragalactic lines of sight with $b > 45^\circ$. Their median values of $\log N \sin b$ are 13.63 ± 0.21 and 14.12 ± 0.25 for Si IV and C IV, respectively, while our corresponding values of $\log N \sin b$ are 13.35 ± 0.17 and 14.04 ± 0.18 . We note that the apparent optical depth method column densities they report include any contribution from the MW itself, which is consistent with our median values of $\log N \sin b$ being slightly lower than theirs. Adding back in the MW contribution to our values for $\log N \sin b$, we get medians of 13.48 ± 0.22 and 14.10 ± 0.24

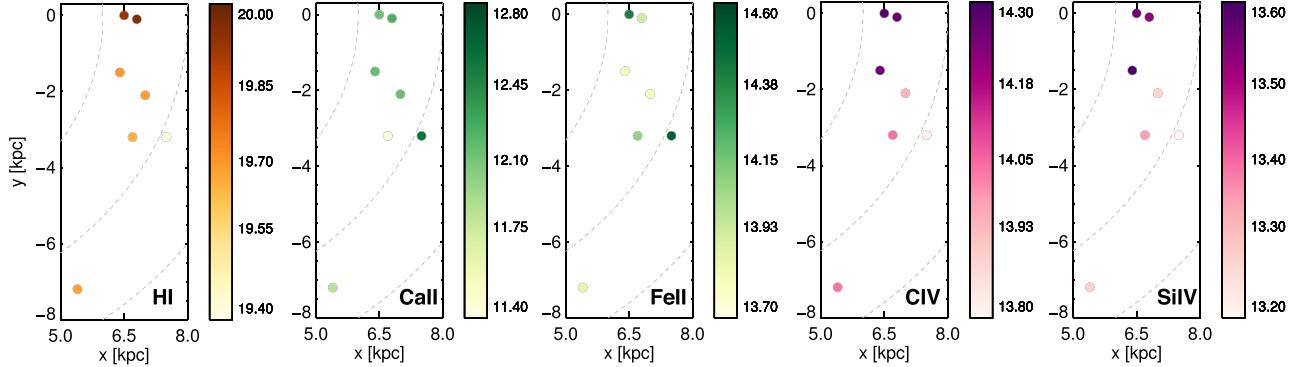


Figure 7. Maps in the x - y plane of IV vertical ionic column densities ($\log N \sin b$), including the EBHIS 21 cm $\log N \sin b$, at the position of each BHB. In this reference frame, the MW center is at the origin (0, 0) off the upper left side of the figure. The physical scale on each axis is given in kpc assuming a distance of 3 kpc for the IV gas. The associated color bar on the right of each panel gives the column density range. The low ions (Ca II, Fe II) show a larger scatter in column densities and little coherence. In the intermediate ions (Si IV, C IV) and H I, the structure is more coherent, with a decrease in vertical column density of ~ 0.05 dex kpc^{-1} from top to bottom. For reference, gray dashed curves show circles of constant galactocentric radius centered on $(x, y) = (0, 0)$ plotted every 2 kpc.

for Si IV and C IV, respectively. This additional contribution brings our median vertical column densities into even better agreement with those of Savage & Wakker (2009).

There is minor tension, at the $\sim 1\sigma$ level, between the best-fit vertical velocity values for the intermediate (Si IV) and low ions (Ca II) seen in the bottom two panels of Figure 6. This tension increases when we consider IV H I 21 cm emission radial velocities (in the range of $-40 < v_{\text{LSR}} [\text{km s}^{-1}] < -15$). The IV gas traced by intermediate-ion absorption appears to be moving toward the disk of the MW approximately 10 km s^{-1} faster than the material traced by the low-ion absorption. This offset could be due to systematic uncertainty introduced by our Voigt profile fitting assumptions not accounted for by the errors propagated during the iterative fitting procedure. It may also arise if we have underestimated the wavelength calibration uncertainty (but see Figure 2). While there may be physical implications of a velocity offset between different gas phases, we do not comment on it further given its low significance in our data.

4.2. Column Density Variation in the x - y Plane

Having established that our data show no significant trends with z , here we focus on a comparison between column densities of the fitted IV components as a function of their transverse separation from one another. Column densities of Si IV and C IV, with ionization potential energies of 45.1 and 64.5 eV, respectively, exhibit the strongest correlation with each other from sight line to sight line. For $\log N_{\text{Si IV}}$ versus $\log N_{\text{C IV}}$, the Spearman's rho rank correlation coefficient (a value of 1 = perfectly increasing monotonic correlation) is 0.86 and the linear Pearson correlation coefficient (a value of 1 = perfectly increasing linear correlation) is 0.97. Figure 7 clearly shows that the seven IV column density measurements of Si IV and C IV are strongly correlated with each other. The low ions (e.g., Ca II, Fe II) with ionization potential energies < 20 eV exhibit weak to moderate correlations with each other (Spearman's rho = 0.2–0.3; Pearson's coefficient = 0.3–0.5), the larger scatter likely due to the large uncertainties in the fitted column densities of the weak IV component.

Next, we examine whether the column densities of low ionization state gas traced by Si II, Fe II, and Ca II correlate with the columns of intermediate ionization state gas traced by Si IV and C IV. The Spearman's rho rank correlation coefficient

between $\log N_{\text{lowion}}$ and $\log N_{\text{highion}}$ ranges between -0.1 and 0.1 for the various species, strongly indicating no monotonic correlation of any kind between the species of different ionization states. As we will explore in Section 6, this result is consistent with pictures put forth by many previous studies of the MW halo and Magellanic Stream (Putman et al. 2003; Fox et al. 2004; Richter et al. 2017), in which larger envelopes of ionized gas contain patchy clumps of cooler, more neutral gas.

Finally, we search for trends between ionic column densities and the IV $\log N_{\text{H I}}$ in order to assess whether the gas seen in absorption toward the BHB stars is tracing the same material as that seen in emission in Figure 3. First, it is important to recall that the H I column densities based on the EBHIS 21 cm emission data are averaged over a much larger angular scale than the absorption-line column densities. Furthermore, the precise value of $N_{\text{H I}}$ associated with N_{ion} is impossible to determine using the 21 cm data alone. Any dense clumps or “cloudlets” along each line of sight with angular sizes $< 10.8'$ would be washed out by the larger EBHIS beam (Wakker et al. 2001; Lockman et al. 2002; Saul et al. 2012). Any differential depletion would add additional scatter. For these reasons, it is not surprising that gas column densities of the low ions traced by Ca II, Fe II, and Si II in Keck and COS data do not correlate at all with $N_{\text{H I}}$. We show this lack of correlation in the right panel of Figure 8 for Ca II, which offers the best constraints on the absorption properties of the low ions.

There is, however, a moderate to strong correlation between both Si IV and C IV and H I, seen in the left panel of Figure 8. For both intermediate ions, the linear least-squares fit is shown as a gray line, indicating positive slopes of 0.59 ± 0.076 (Si IV) and 0.75 ± 0.096 (C IV). The Spearman's rho rank coefficient (such that a value of 1 = a perfectly monotonic positive correlation) in each case indicates a moderately strong correlation, with values of 0.78 and 0.75 for Si IV and C IV, respectively. What this likely means is that the H I column densities in the $10.8'$ beam on average are tracing the larger-scale, diffuse structures rather than the smaller-scale, dense “cloudlets” probed by the low ions.

Directly addressing the angular scale of the absorbing and emitting structures, Figure 9 shows the difference in column density between each unique IV component “pair” for the seven lines of sight (21 pairs total) versus the pair angular separation. The C IV and Si IV are tracing a coherent, large-

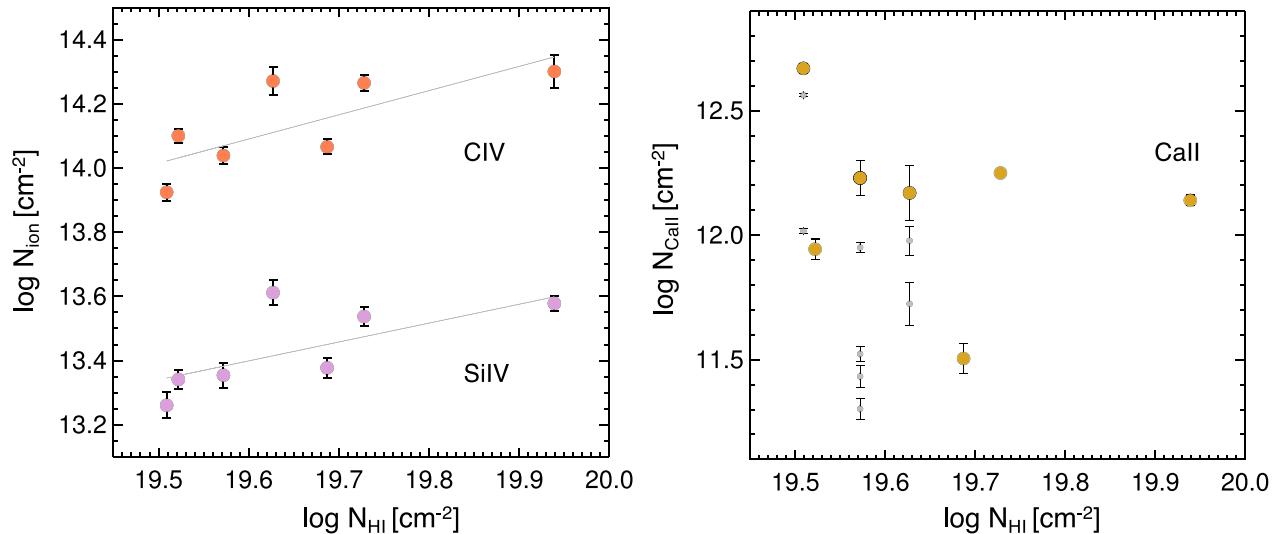


Figure 8. Left: correlation between column densities of intermediate ions C IV (coral) and Si IV (plum) and the EBHIS 21 cm $\log N_{\text{H}\alpha}$ in the IV gas components. The linear least-squares fits are shown as gray lines, giving positive slopes of 0.59 ± 0.076 (Si IV) and 0.75 ± 0.096 (C IV). Right: column densities of all individual Ca II IV components (in gray) and the sum of the Ca II IV component column densities along BHB lines of sight (in gold). There is no correlation between Ca II IV component column densities and the IV EBHIS 21 cm $\log N_{\text{H}\alpha}$. Although we do not show it here, there is a similar lack of a trend and large scatter in other low ion column densities measured in the COS data (e.g., Fe II, Si II), albeit with considerably larger errors in the fitted IV component columns.

scale structure, as evidenced by the strong correlation between $\Delta \log N$ and the separation over 40° . Both intermediate-ion correlations give a Spearman’s rho coefficient of 0.82, with a two-sided significance of its deviation from 0 of <0.000004 . The scatter in this correlation is only ~ 0.15 dex, and within 10° the median $\Delta \log N$ is only 0.05 dex, indicating a relatively smooth structure. Furthermore, the IV H I 21 cm column densities (shown for reference as gray diamonds) seem to follow a similar trend, but with a slightly larger scatter. At a distance of 3 kpc, roughly the upper limit we can set with our data, $10^\circ \approx 0.5$ kpc. At the lower distance limit to the IV arch of 0.9 kpc (Wakker 2001), the structure of which the BHB sight lines skirt the edge, $10^\circ \approx 0.15$ kpc.

In the right panel of Figure 9, we show the same pairwise $\Delta \log N$ versus angular separation for two low ionization state transitions, Ca II (gold) and Fe II (green). While the Ca II IV component column densities are well determined from the Keck HIRES data, the $\log N_{\text{Fe II}}$ of the IV component has a much larger error owing to the many assumptions and limits imposed on the lower-resolution COS spectrum. Nonetheless, we can rule out any correlation between $\Delta \log N_{\text{lowion}}$ and angular separation with high confidence. Furthermore, the scatter of the data over the 40° angular scales is >0.4 dex. This large scatter and lack of a correlation are consistent with patchy, small-scale clumps of material. Considering the monotonic increase with small scatter to 40° of the intermediate ions and the large scatter and lack of trend seen in the low ions (right panel of Figure 9), our data directly support the physical picture of large ionized envelopes or streams containing cooler condensations or clumps.

In principle, additional archival UV absorption-line data could support this picture of \sim kiloparsec-scale ionized structures with even larger area coverage. As an initial step, we harvested Si IV column densities from the COS-GAL QSO sight-line database (Zheng et al. 2019) to explore $\Delta \log N$ variations over a similar area in the northern Galactic sky to that of our BHB stars. In total, we selected 11 COS-GAL sight lines against which to compare our BHB Si IV IV component

column densities. The results can be seen in Figure 10, where the gray points track the pairwise Si IV column density differences for the extragalactic COS-GAL sight lines. This archival sample shows only a weak to moderate increasing correlation between pair column density difference and angular separation, with a Spearman’s rho correlation coefficient of 0.35 and a two-sided significance of its deviation from 0 of ≈ 0.04 .

The most likely reason for the weaker correlation in Figure 10 is that there is a contribution to the absorption along the QSO sight lines from the more extended, ionized CGM up to much larger heights, ~ 200 kpc (Zheng et al. 2019). Each QSO sight line probes the full MW halo virial radius and likely pierces multiple kiloparsec-scale structures in Si IV. The effect would be to wash out the individual structure coherence seen for our stellar sight lines. Furthermore, the apparent optical depth method column densities reported by COS-GAL include an MW component, whereas we explicitly do not. The MW component’s contribution to the total column density is minimal for Si IV and C IV but still exists at the 20% level and likely adds additional scatter to the relation. Thus, we consider Figure 10 to offer tentative support for our physical picture (explored in detail in the discussion). Further support would come from analyzing all stellar sight lines with distance constraints probed by COS and STIS, including a full Voigt profile analysis similar to that which we present in this work. Such an analysis over large angular area, while beyond the scope of our study, would potentially be sensitive to the *maximum* size scale of these large, ionized structures.

5. Ionization Models

As intermediate-ionization species, Si IV and C IV may be a part of the WIM (10^4 – 10^5 K) or may lie in a separate, tenuous “transitional” phase at $\sim 10^5$ K. The former would indicate photoionization, while the latter would point toward collisional ionization, or gas cooling out of the ambient hot halo, as the dominant process giving rise to the Si IV and C IV. In this

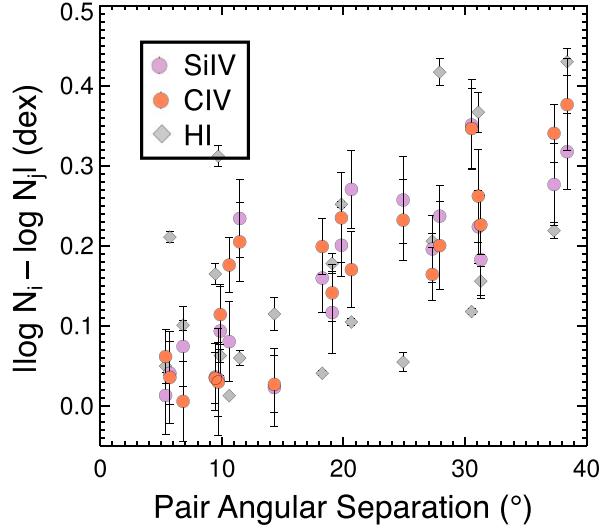


Figure 9. Pairwise ionic column density differences for the IV components along all seven sight lines (21 unique pairs) vs. their angular separation on the sky in degrees. At a distance of 3.0 kpc, 10° corresponds to ~ 0.5 kpc. There is a highly significant monotonic, increasing correlation between pair separation and column density differences in the intermediate ions (Si IV in plum, C IV in coral; Spearman’s rho ≈ 0.82 for both). H I 21 cm emission (in gray diamonds; Spearman’s rho ≈ 0.46) shows only a weak to moderate correlation with much larger scatter. In contrast, in the right panel, there is no correlation between low-ion column density differences between pairs and their angular separation on the sky, and much larger column density differences overall (Fe II in green, Ca II in gold; Spearman’s rho $\lesssim 0.05$).

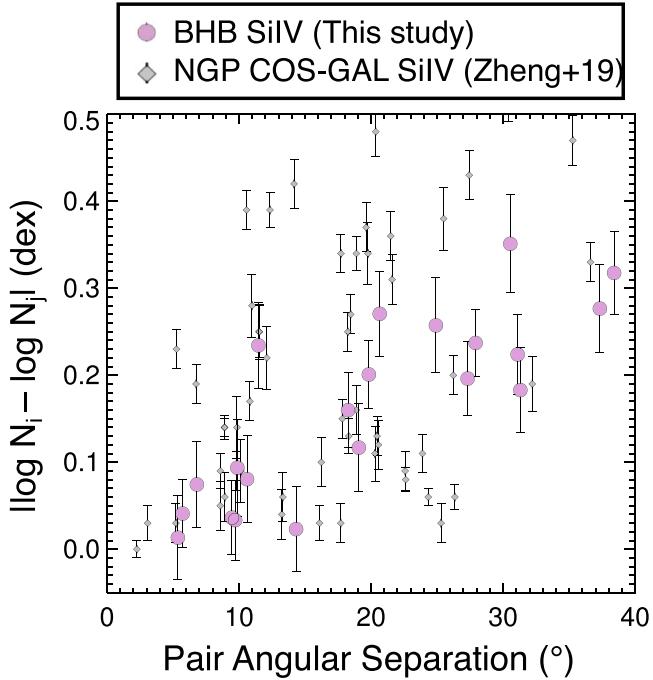
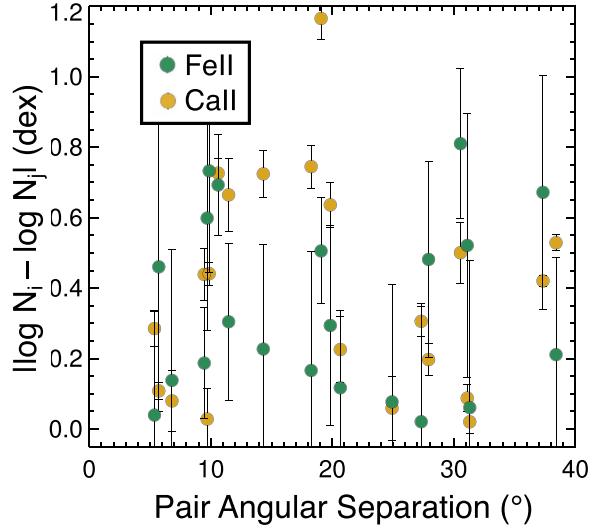


Figure 10. Pairwise Si IV column density differences for the IV components along all seven sight lines (plum circles; 21 unique pairs) vs. their angular separation on the sky in degrees, compared to a sample of 11 sight lines from the COS-GAL database (gray diamonds; Zheng et al. 2019) selected to be in the same region of the northern Galactic sky as the BHB stars. The scatter in the COS-GAL QSO sight-line database is larger for reasons outlined in Section 4.2, but there remains a weak to moderate increasing correlation between pair column density difference and angular separation, with a Spearman’s rho correlation coefficient of 0.35.

section, we examine the constraints imposed by our observations for both cases. Ultimately, this analysis drives at the origin of the material and possibly at the nature of its



surroundings, informing our understanding of relevant Galactic fountain astrophysics.

5.1. Constraints from Photoionization Models

Photoionization models, such as those constructed with the spectral synthesis code CLOUDY (Ferland et al. 2017), perform detailed radiative transfer through a parcel of gas characterized by some neutral column density, $N_{\text{H I}}$, metallicity, [M/H], and ionization parameter, U ($U \equiv (n_{\text{H}} \times c)/\Phi_{\text{tot}}$, where Φ_{tot} is the integrated ionizing photon flux generated by the assumed photoionizing spectrum). One must make several assumptions to run such a CLOUDY model: (1) thermal and ionization equilibrium, (2) an ionizing background radiation field, and (3) a density structure for the parcel of gas, often approximated as a uniform density slab (but see Stern et al. 2016). Another common assumption of these models is relative solar abundances of the elements, although these element ratios can be allowed to vary to account for depletion onto dust (Jenkins 2009). Comparing the outputs of a grid of ionization models with the data (e.g., systematically varying $N_{\text{H I}}$, [M/H], and/or U) ultimately allows one to find the model or set of models that are consistent with the constraints set by the ionic column densities determined from the observations.

We construct such a grid of CLOUDY models to constrain the physical state of the IV gas probed by our observations of Si IV and C IV *only*, assuming that they are predominantly photoionized (we address collisional ionization in and out of equilibrium in Section 5.2). We assume solar ratios of the elements and exclude the low ions from this analysis. We are justified in doing so, as our previous analysis has shown that they are almost certainly not cospatial with the intermediate ionization state material. We note that if both the low ions and the intermediate ions are part of the photoionized WIM, then the WIM material itself cannot be characterized by a single density “slab.” However, our CLOUDY models of the single-phase intermediate-ionization IV gas assume a uniform density layer irradiated by the extragalactic UV background and

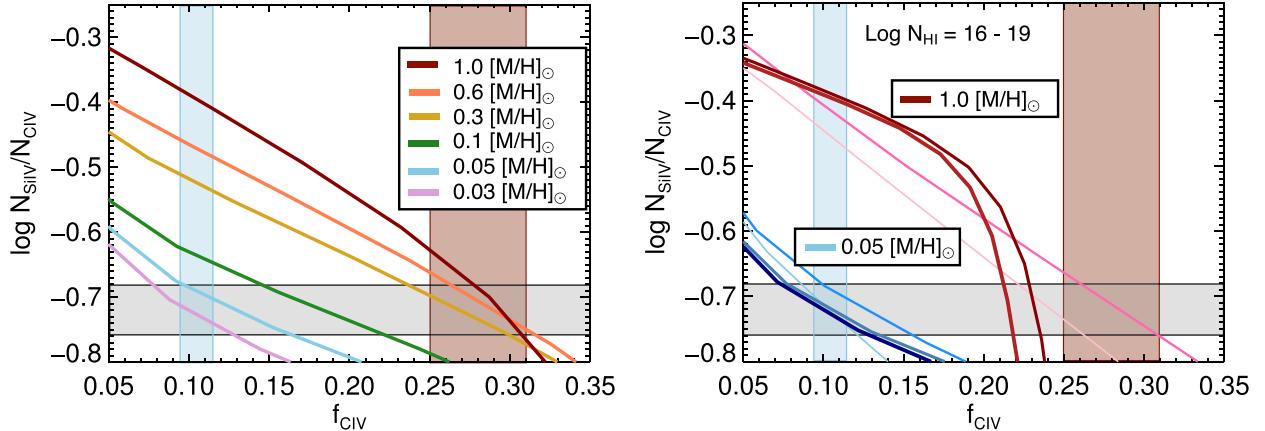


Figure 11. Fraction of carbon seen as C IV, $f_{\text{C IV}}$, vs. the column density ratio $\log N_{\text{Si IV}}/N_{\text{C IV}}$ for six CLOUDY photoionization models of varying metallicity $[\text{M}/\text{H}] = 0.03\text{--}1.0$ (left) and varying $\log N_{\text{H I}}$ from 16 to 19 in 1 dex steps (right). The range of observed $\log N_{\text{Si IV}}/N_{\text{C IV}}$ is shown in the gray horizontal regions in each panel, with well-constrained values ranging from ≈ -0.76 to -0.68 . In the right panel, we show increasing $\log N_{\text{H I}} [\text{cm}^{-2}]$ with line thickness and darkness from 16 to 19. We show these sets of curves at two values of metallicity, $[\text{M}/\text{H}] = 0.05$ (blue curves) and $[\text{M}/\text{H}] = 1.0$ (brownish-red curves). We highlight two potential “solutions” for gas-phase metallicity, $f_{\text{C IV}} \approx 11\%$ in vertical bands of light blue, and $f_{\text{C IV}} \approx 25\text{--}31\%$ in brownish-red. We discuss these two divergent solutions in Section 5.1.

escaping radiation from the MW disk, as presented by Fox et al. (2005). These models represent a simplified picture that nonetheless may prove informative.

Φ_{tot} (the number of ionizing photons emitted per second by the assumed radiation field) at the MW disk–halo interface is highly uncertain and depends strongly on (1) the absorbing feature’s distance from the MW and (2) an assumed escape fraction of the ionizing photons from the disk. In practice, adjusting Φ_{tot} affects primarily the calculated hydrogen number densities, $n_{\text{H}} = \Phi_{\text{tot}}/(U \times c)$. To address this uncertainty, we ran several CLOUDY models, assuming the Fox et al. (2005) ionizing spectrum at different heights above the disk. For reference, we find that at 10, 2, and 0 kpc above the MW disk, $\Phi_{\text{tot}} = 340,000$, 5,500,000, and 37,000,000 s^{-1} , respectively. We note that the ionization parameters and ionized fractions of hydrogen and metals are consistent among these models owing to identical *spectral shapes*. However, it is apparent that such a range in disk heights produces ionizing photon fluxes, and thus gas volume densities, that vary over two orders of magnitude. Specifically, the assumed hydrogen number densities shift upward by two orders of magnitude with a corresponding increase in Φ_{tot} . Thus, constraining the volume density of this IV material is not possible without additional constraints on key parameters (namely, N_{H} and $[\text{M}/\text{H}]$, aka metallicity). In the following analysis, we report metallicities relative to solar abundances, $X[\text{M}/\text{H}]/[\text{M}/\text{H}]_{\odot}$, using the compact notation $[\text{M}/\text{H}] = X$, where X is some fraction of solar.

Under the above assumptions, the column density ratio $\log N_{\text{Si IV}}/N_{\text{C IV}}$ is a sensitive probe of the fraction of total carbon seen as C IV $\equiv f_{\text{C IV}}$. This dependence is shown in Figure 11. The median $\log N_{\text{Si IV}}/N_{\text{C IV}} = -0.69 \pm 0.04$ and shows little variation from sight line to sight line. The precise value of $f_{\text{C IV}}$ inferred by the measured range of $-0.76 < \log N_{\text{Si IV}}/N_{\text{C IV}} < -0.68$ depends strongly on the assumed metallicity of the gas. This metallicity dependence is a result of the shape of the cooling curve and the strong dependence of the high and intermediate ionic fractions on the metal line cooling. The left panel shows that these models are very sensitive to the gas-phase metallicity. We chose an $N_{\text{H I}} = 17.5$ at which to plot these CLOUDY models on the left, and we note that the

models are only sensitive to the choice of $N_{\text{H I}}$ at $[\text{M}/\text{H}] \gtrsim 1.0$ (at metallicities \lesssim solar, there is no sharp turnover for optically thick models, and the spread in $f_{\text{C IV}}$ looks similar to the blue curves in the left panel). In the right panel, we see that these curves shift slightly with variations in assumed $N_{\text{H I}}$ over a range of $16.0 < \log N_{\text{H I}} < 19.0$, with greater differences between optically thin and thick gas at high metallicity. The fraction of carbon in C IV does not depend at all on the assumed Φ_{tot} (although it does depend on the slope of the assumed ionizing spectrum). For the range of $\log N_{\text{Si IV}}/N_{\text{C IV}}$ in the observed IV gas along our BHB sight lines, these models indicate that anywhere between 7% and 32% of the carbon is in C IV. Thus, for our median $\log N_{\text{C IV}} \sin b$ of 14.04, this implies $14.5 < \log N_{\text{C, total}} \sin b < 15.2$. In principle, if we had some way of independently constraining either the metallicity of the material or the total hydrogen column density along the line of sight, N_{H} , we would fully constrain the metal content of the IV gas along our sight lines.

5.1.1. Independent Constraints on $[\text{M}/\text{H}]$ and N_{H}

Addressing the metallicity of the gas, Wakker (2001) and Wakker et al. (2001) present a thorough and exhaustive spectroscopic analysis to determine the metallicity of the nearby IV arch and other IVC and HVC halo structures. They use reliable values of $N_{\text{H I}}$ and measurements of N I, O I, and S II along four lines of sight, all low ions that are presumed to trace predominantly neutral material. Wakker (2001) finds that the neutral phase of the IV arch has a metallicity of 0.8–1.2 solar. In general, IVCs present metallicities typically between 0.5 and 1.5 solar, indicating an origin in the disk of the MW (Richter et al. 2017).

Assuming that the IV neutral and ionized gas phases are “well mixed,” they would have the same metallicity. There are at least a few lines of evidence to suggest that this is not an unreasonable approximation (e.g., Richter et al. 2017). Adopting solar metallicity for the intermediate ionization state material traced by our observations puts us on the brownish-red curve in the left panel of Figure 11. This curve then leads us to adopt the maximum value of $f_{\text{C IV}}$, $\sim 30\%$, where the constraints of our measurements on $\log N_{\text{Si IV}}/N_{\text{C IV}}$ indicated by the

horizontal gray shaded region intersect with the curves shown by the vertical rectangular brownish-red shaded region. Thus, we may adopt the lower bound on the total carbon content of the IV gas of $\log N_{\text{C, total}} \sin b \approx 14.5$. For self-consistency with the assumed solar metallicity, considering the solar carbon abundance of $12 + \log (\text{C/H}) = 8.43$, this estimate implies an $N_{\text{H}} \sin b$ of ≈ 18.1 .

However, we may additionally constrain $N_{\text{H}} \sin b$ using an entirely different method. To do so, we turn to pulsar DM estimates for the MW halo (e.g., Cordes & Lazio 2002; Howk et al. 2006; Gaensler et al. 2008). As the light from a pulsar located in the MW halo interacts with free electrons along its line of sight, pulses at lower frequencies experience a delay relative to the higher-frequency pulses (e.g., Ferrière 2001). Sensitive pulsar timing experiments toward thousands of pulsars in the MW halo allow constraints on the DM of radio pulsars that translates directly into a free-electron column, N_{e} (e.g., Taylor & Cordes 1993). Several works have combined these measurements into a physical model of the MW WIM by assuming some large-scale structure (e.g., exponential disk). These models can then be used to estimate N_{e} at any given set of galactic coordinates for any assumed height (e.g., Cordes & Lazio 2002; Berkhuijsen et al. 2006; Gaensler et al. 2008; Yao et al. 2017).

We use the model of Gaensler et al. (2008), with the updates from Savage & Wakker (2009) that account for inhomogeneities in the gas, to estimate $N_{\text{e}} \sin b$ for high-latitude lines of sight in the region of the halo probed by our BHB sight lines. Integrating the N_{e} of the model to the maximum height of our BHB sample results in $\log N_{\text{e}} \sin b = 19.83 \text{ cm}^{-2}$. This value is similar to that of Howk et al. (2006), who use a value of $\log N_{\text{e}}$ toward their high-latitude pulsars of 19.91 cm^{-2} , although these lie at distances of $\sim 10 \text{ kpc}$. Accounting for the distance constraints on our IV gas clouds, determined to predominantly lie at $z < 3.4 \text{ kpc}$, we reduce our adopted value of $\log N_{\text{e}} \sin b$ to 19.79 (see Equation (6) of Gaensler et al. 2008, with $z = 3400 \text{ pc}$, $H_n = 1410 \text{ pc}$, and $n_0 = 0.0155 \text{ cm}^{-3}$ as in Savage & Wakker 2009; here n_0 is the density of gas at the midplane, and H_n is the corresponding scale height).

Additional assumptions, such as the fraction of free electrons due to helium and the fraction of free electrons due to hot, coronal gas, allow us to convert N_{e} to N_{H} . For these additional assumptions, we rely on the pioneering work of Howk et al. (2006), who used sight lines toward globular clusters containing radio pulsars to independently constrain N_{H} and derive accurate gas-phase abundances. We refer readers to that paper for details on the method we follow and the limits we adopt. For the helium correction, Howk et al. (2006) adopt a minimal case such that the fraction of doubly ionized helium, $x(\text{He}^{+2}) = 0$, and $\eta_{\text{WIM,min}}$, a numerical correction factor to the conversion between N_{e} and N_{H} , is equal to 0.98. The maximum correction factor for helium, $\eta_{\text{WIM,max}}$, is assumed to be 0.81. Following a two-component prescription for the hot ionized medium (HIM) contribution presented in Equation (3) of Howk et al. (2006), we assume $(N_{\text{e}} \sin b)_{\text{HIM}} = 19.10$. Thus, our minimum and maximum values of $\log N_{\text{H II,WIM}} \sin b$ are 19.59 and 19.67, respectively.

Converting $N_{\text{H II}} \sin b$ to $N_{\text{H}} \sin b$ requires knowing the neutral fraction, which requires that we return to our ionization model. We note that for the full range of models shown in Figure 11, over the allowed values of $\log (N_{\text{Si IV}}/N_{\text{C IV}})$, the neutral gas fraction ranges from $\sim 0.1\%$ to 0.0001% . Thus, we

assume that $N_{\text{H II}} \sin b \approx N_{\text{H}} \sin b$ for the intermediate ionization state material. Immediately we can see there is a mismatch with the solar-metallicity model discussed earlier in this section, on the order of 1.5 dex ($\log N_{\text{H}} \sin b$ of ≈ 18.1 vs. 19.6)! Another way to examine this discrepancy is to find the set of self-consistent models such that the derived abundance from $N_{\text{C}}/N_{\text{H}}$ is equal to the value one would derive from the [M/H] curve in Figure 11 and the associated $f_{\text{C IV}}$. The only self-consistent models in this framework lie at $[\text{M/H}] \approx 0.05$. Tuning the assumptions within the allowed ranges does not bring the two models into agreement. For example, adopting a metallicity at the low end of the neutral IVC metallicity range, 0.5 solar, and choosing a maximal value of $(N_{\text{e}} \sin b)_{\text{HIM}} = 19.22$ (Howk et al. 2006), the values continue to be highly discrepant, on the order of ~ 1 dex.

We note that these limits are robust to uncertainties in the photoionization modeling barring a very significant alternation to the assumed slope of the ionizing spectrum from 1 to 5 ryd. One way to resolve the discrepancy with ionization modeling is to arbitrarily assign Si IV and C IV to distinct phases themselves. Doing so would not be justified by the data, especially given the extent to which Si IV and C IV column densities, velocity centroids, and line widths are well matched along every single line of sight in our sample. Another possibility, which we explore in detail in Section 5.2, is that Si IV and C IV represent neither WIM nor HIM material, but rather are the manifestations of a tenuous, transitional temperature, $T \sim 10^5 \text{ K}$ plasma cooling out of the ambient $\sim 10^6 \text{ K}$ coronal medium.

In summary, the two unacceptable photoionization model options presented above are that either the intermediate ionization state IV gas is metal-poor and not well mixed with the neutral IV gas or $(N_{\text{e}} \sin b)_{\text{HIM}} \approx (N_{\text{e}} \sin b)_{\text{WIM}}$. That is, for the models to converge on a metal-rich solution such that $[\text{M/H}]_{\text{IV}} \approx 0.5$ or above, the HIM or some other ionized gas phase would have to contribute $\gtrsim 95\%$ of the free electrons along the lines of sight to pulsars. Such a large value of $\log N_{\text{H II,HIM}}$ on the order of 19.7 is inconsistent with observed constraints on $n_{\text{H,HIM}} \lesssim 8 \times 10^{-4} \text{ cm}^{-3}$ (Bregman & Lloyd-Davies 2007; Grcevich & Putman 2009). Inefficient mixing of gas phases has been invoked in the context of Galactic winds (Oppenheimer et al. 2012) but is not a particularly elegant solution in a Galactic fountain scenario in which this material originated in the disk of the MW.

5.1.2. A Self-consistent Photoionization Solution Requires an Si-depleted Ionized Medium

Can we resolve this tension in the models, or must we reject photoionization of Si IV and C IV altogether? There is a solution we have not yet considered. That is, the photoionization models are perfectly self-consistent if we assume that there is a substantial fraction of silicon depleted onto dust in the warm ionized disk-halo interface of the MW. Turning again to Figure 11, we see that the orange and brownish-red curves (metallicities of 0.5–1.0 solar) are brought into consistency with the pulsar dispersion constraints (light-blue shaded region) when we allow for silicon depletion relative to carbon between -0.25 and -0.35 dex. That is, the undepleted ratio of $\log N_{\text{Si IV}}/N_{\text{C IV}}$ should be closer to ≈ -0.4 if the gas is metal enriched at a level of 0.5–1.0 solar.

Addressing the photoionization model's self-consistency, we note that the length scale implied by this fiducial model

($\ell = N_{\text{H}}/n_{\text{H}}$) is

$$\ell \approx 3.2 \text{ kpc} \left(\frac{N_{\text{H}}}{10^{19.7} \text{ cm}^{-2}} \right) \left(\frac{10^{-2.3} \text{ cm}^{-3}}{n_{\text{H}}} \right). \quad (2)$$

This value of $n_{\text{H}} = 0.0055 \text{ cm}^{-3}$ occurs at $\Phi_{\text{tot}} = 5,500,000 \text{ s}^{-1}$ (the value from a Fox et al. 2005 spectrum 2 kpc from the disk) precisely at an $f_{\text{C IV}}$ of 11%. While there is a slight mismatch (2 kpc vs. 3.2 kpc), the length scales are consistent enough within the uncertainty of Φ_{tot} not to raise any red flags. Additionally, a ≈ 3 kpc length scale is consistent with the scale heights of Si IV and C IV given by Savage & Wakker (2009). Finally, we do note that this gas density is more than an order of magnitude lower than the volume density of the WIM of the MW disk inferred from measurements of singly and doubly ionized species (Reynolds 1993; Sembach et al. 2000). If IV Si IV and C IV are an extension of the MW’s photoionized WIM, they would represent a lower-density extension (e.g., Stern et al. 2016).

5.2. The Case in Which Si IV and C IV Are Collisionally Ionized in Transition Temperature Gas

Collisional ionization is the dominant process in low-density plasma in the absence of a significant ionizing radiation field. In collisional ionization equilibrium (CIE), the ion fractions produced by this process are a function of only gas temperature and are determined by detailed balance between electron impact collisional ionization, electron–ion recombination, and charge transfer reactions (Gnat & Sternberg 2007). Outside of equilibrium, the ionization state of the plasma depends on whether the gas is cooling at a constant density (isochoric) or at a constant pressure (isobaric) and also on the gas-phase metallicity (e.g., Gnat & Sternberg 2009; Wiersma et al. 2009). Furthermore, complex nonequilibrium collisional ionization processes have been calculated for a number of physical processes, including supersonic shocks, conductive interfaces, and turbulent mixing layers between cool clouds and a hot ambient medium (Allen et al. 2008; Gnat et al. 2010; Kwak & Shelton 2010; Wakker et al. 2012). Here we explore a physical scenario in which Si IV and C IV trace so-called “transitional temperature” 10^5 K gas and are produced predominantly via collisionally driven cooling from a hot, ionized coronal plasma (e.g., Savage & Wakker 2009).

The physical conditions that would give rise to a scenario in which, e.g., Si II, Si III, C II, and C III are photoionized, but Si IV and C IV are collisionally ionized, would involve an ionizing spectrum that does not produce sufficient photons with energies $>45 \text{ eV}$. In their study of HVCs exposed to a radiation field owing to Galactic ionizing photons and the extragalactic background, Fox et al. (2005) find that it is difficult to put both C IV and O VI into the same photoionized phase as the low ions. While this may result from C IV and O VI simply arising in lower-density material that is nonetheless still photoionized, it may instead be indicative of collisional processes at work. Figure 8 of Fox et al. (2005) shows several ionization “edges” in their assumed spectrum and a particularly steep drop at 54 eV caused by the He II edge in hot stars. There are still photons with energies $>54 \text{ eV}$, but the uncertainty in their overall flux is large enough to make the collisional ionization scenario for Si IV and C IV plausible. In our analysis that follows, we ignore photoionization and require the majority of

the Si IV and C IV to be collisionally ionized in a $\sim 10^5 \text{ K}$ medium.

Figure 12 shows predictions of gas column densities and ionization fractions from collisional ionization models both in and out of equilibrium (Gnat & Sternberg 2007; Gnat et al. 2010). In CIE, as shown with the black curve in Figure 12, the fraction of C in C IV at these temperatures is $\sim 15\%$ for the observed range of $\log N_{\text{Si IV}}/N_{\text{C IV}}$. This fraction implies a total vertical carbon column density of $\log N_{\text{C, total}} \sin b \approx 14.9$. Nonequilibrium models, seen in brownish-red and green in Figure 12, suggest an even higher total carbon column density.

We now turn to the implied total hydrogen column in the supposed tenuous gas phase in CIE for a range of gas-phase metallicities. If the gas is metal enriched at the solar level, this implies $\log N_{\text{H}} \sin b = 18.4$, or $\sim 5\%$ of the total $N_{\text{H}} \sin b$ implied by pulsar DM estimates described in the previous section. If this gas is instead at $[\text{M}/\text{H}] = 0.3$, the amount of hydrogen would increase to $\log N_{\text{H}} \sin b = 18.9$, or $\sim 16\%$ of the total $N_{\text{H}} \sin b$ measurement in the MW thick disk. Given well-reasoned estimates by Howk et al. (2006) that the $\log N_{\text{H,HIM}} \sin b \approx 19.0$, the CIE model strongly implies an $[\text{M}/\text{H}] > 0.3$ for the warm and hot components of the MW halo. We further note that the lower $f_{\text{C IV}}$ values implied by nonequilibrium models with $[\text{M}/\text{H}]$ ranging from 0.1 to 1.0 (brownish-red and green curves in Figure 12) would require $N_{\text{H}} \sin b > 19.4$, or that $>50\%$ of the ionized hydrogen column in the halo be in this collisionally ionized 10^5 K phase. This picture is not physical and would directly contradict the observational evidence that the WIM and the HIM are likely dominant contributors to the line-of-sight N_{e} (Reynolds 1993; Howk et al. 2006; Gaensler et al. 2008).

Previous considerations of collisionally ionized 10^5 K gas in the MW halo have envisioned layers and/or interfaces (e.g., Fox et al. 2009; Wakker et al. 2012). In these structures, either conductive interfaces or turbulent mixing layers, predictions of $\log N_{\text{Si IV}}/N_{\text{C IV}}$ lie in our observed range, but the total C IV column densities of each layer are closer to $\sim 10^{12} \text{ cm}^{-2}$, two orders of magnitude lower than we observe (Kwak & Shelton 2010). This inconsistency is alleviated if we demand that our line of sight intersect hundreds of such overlapping layers. However, in this scenario, hundreds of layers cooling out of equilibrium would tend to imply values of $\log N_{\text{H}} \sin b > 19.4$ in this 10^5 K “tenuous” gas phase, just as we described in the previous paragraph. Furthermore, hundreds of interfaces or layers between small cool clouds and the HIM are difficult to reconcile with the roughly kiloparsec-scale spatial correlations of the Si IV and C IV in the x - y plane.

Gas cooling behind supersonic shocks has also been invoked to explain the presence of highly ionized species such as Si IV and C IV in the CGM (Dopita & Sutherland 1996; Werk et al. 2016; McQuinn & Werk 2018). Under the assumption of a steady, 1D flow and a given metallicity, several models predict column densities of high and intermediate ions in the post-shock cooling layers. Here we use the results of Allen et al. (2008), who produce tables for gas at solar metallicity, shock speeds that range from 100 to 1000 km s^{-1} , and pressure exerted by the transverse magnetic field ($B_{\text{t}} n^{-\frac{1}{2}}$) ranging from 10^{-4} to $10 \mu\text{G cm}^{-3/2}$. The magnetic pressure is important because it limits the compression through the shock, and thus higher magnetic parameters produce higher gas column densities. These shock models predict values of $\log N_{\text{Si IV}}/N_{\text{C IV}}$ to range widely from -1.0 to 0 , mostly dependent on the shock speed. However, when we consider the predicted

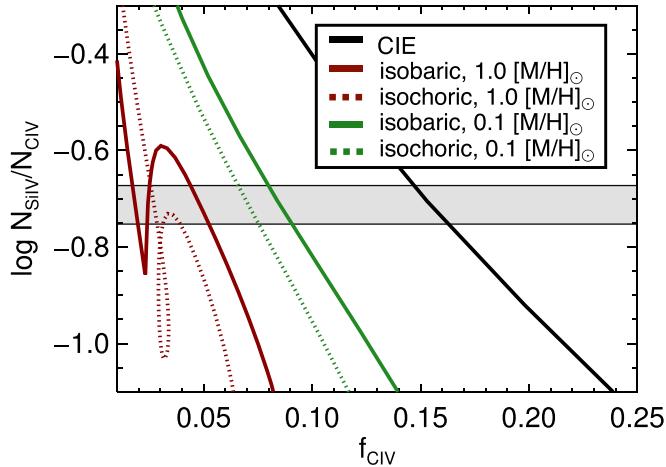


Figure 12. Fraction of carbon seen as C IV, $f_{\text{C IV}}$, vs. the column density ratio $\log N_{\text{Si IV}}/N_{\text{C IV}}$ for several models of collisional ionization based on Gnat & Sternberg (2007). CIE produces a trend shown as the black line and is independent of density and metallicity. Out of equilibrium (brownish-red and green curves) the predicted column densities and ionization fractions depend on the metallicity of the gas and whether its cooling is isobaric or isochoric. The range of our measured $N_{\text{Si IV}}/N_{\text{C IV}}$ is shown by the horizontal gray rectangle.

values of the gas column densities at the observed ratio of -0.7 , they are closer to a few times 10^{12} cm^{-2} and 10^{13} cm^{-2} for Si IV and C IV, respectively, or more than an order of magnitude lower than we observe.

We conclude that the collisional ionization case is most plausible if the intermediate-ionization gas is in CIE, presumably cooling from the hot coronal medium. In this picture, the 10^6 K gas and the 10^5 K gas would have the same metallicity, and the majority of the cooling would be taking place within 3 kpc of the MW disk. The contribution of the HIM to the value of N_e is uncertain, but if we return to the discussion in Section 5.1, we may use $\log N_{\text{H, HIM}} \sin b \approx 19.0$ (Howk et al. 2006). Using the corresponding constraint on $\log N_{\text{C, total}} \sin b \approx 14.9$ (Figure 12) in the collisionally ionized medium (CIM) and requiring that $N_{\text{H, CIM}} < N_{\text{H, HIM}}$, we find that the hot corona of the MW must have $[\text{M}/\text{H}] > 0.3$ in this physical scenario.

5.3. Estimate of Areal Mass Infall Rate

Given the above constraints on the physical conditions of the C IV-bearing gas, we may estimate areal gas accretion rate for the observed region of sky containing our BHB lines of sight, parameterized as

$$\frac{dM_{\text{gas}}}{dtdA} \equiv \frac{M_{\text{gas}} v_{\text{infall}}}{d_{\text{gas}} \pi R_{\text{cloud}}^2} [\text{M}_\odot \text{ kpc}^{-2} \text{ yr}^{-1}]. \quad (3)$$

We determine the total gas mass from column density measurements, using an effective surface area calculation:

$$M_{\text{gas}} = 310 \left(\frac{Z_\odot}{Z} \right) \pi R_{\text{cloud}}^2 \langle N_{\text{C IV}} \sin b \rangle m_C \left(\frac{1}{f_{\text{C IV}}} \right) [\text{M}_\odot]. \quad (4)$$

Equation (4) assumes the solar abundance of carbon, $12 + \log (\text{C}/\text{H})$, to be 8.43 (Asplund et al. 2009). Among the quantities relevant to this calculation, we allow for the following values and ranges of values: (1) the vertical component of the mean C IV column density of the IV gas, $\langle N_{\text{C IV}} \sin b \rangle = 10^{14.04}$

cm^{-2} ; (2) the fraction of carbon in C IV, $0.09 < f_{\text{C IV}} < 0.17$; (3) the gas-phase metallicity, $0.3 < Z/Z_\odot < 1.0$; (4) the physical size scale of the absorbing structure $R_{\text{cloud}} \approx 1 \text{ kpc}$; (5) the vertical infall velocity of the IV component of C IV absorption, $v_{\text{infall}} \equiv -v \sin b \approx 45 \text{ km s}^{-1}$; and (6) the distance to the absorbing structure, $1.0 \text{ kpc} < d_{\text{gas}} < 3 \text{ kpc}$. We note that this areal mass infall rate considers only the material in this intermediate ionization state gas containing C IV and Si IV but captures constraints from both photo- and collisional ionization within the range of allowed $f_{\text{C IV}}$ values and metallicities.

Following Equations (3) and (4), the resulting bounds on the areal mass accretion rate in warm, ionized gas at high Galactic latitudes are

$$0.0003 < \frac{dM_{\text{gas}}}{dtdA} < 0.006 [\text{M}_\odot \text{ kpc}^{-2} \text{ yr}^{-1}]. \quad (5)$$

This range of areal mass accretion rate encompasses the mean star formation rate surface density of the MW disk, $\langle \Sigma_{\text{SFR, MW}} \rangle = 0.0033 \text{ M}_\odot \text{ yr}^{-1} \text{ kpc}^{-2}$ (Kennicutt & Evans 2012). Our implied total mass infall rate of high-latitude WIM, dM_{gas}/dt , is $0.02 \text{ M}_\odot \text{ yr}^{-1}$ at most and at least $0.001 \text{ M}_\odot \text{ yr}^{-1}$. This range is lower than estimates from all-sky studies of the mass infall rates from the ionized components of HVCs and IVCs, excluding the Magellanic Stream, which find ranges of dM_{gas}/dt to be $0.45\text{--}1.40 \text{ M}_\odot \text{ yr}^{-1}$ at $d < 15 \text{ kpc}$ (Lehner & Howk 2011; Fox et al. 2019). It is also considerably lower than the mass infall rate for the entire MW halo, $\gtrsim 5 \text{ M}_\odot \text{ yr}^{-1}$ (Fox et al. 2014; Richter et al. 2017). It is unlikely that the small neutral clumps we have observed in this region provide a substantive boost to this mass infall rate (e.g., Putman et al. 2012). However, it is possible that material from elsewhere in the halo makes it to the disk via other routes followed by radially mixing within the disk (e.g., Sellwood & Binney 2002). In contrast, our observations may instead imply that much of the more distant ionized halo material that is moving toward the MW disk never in fact makes it through the disk–halo interface to eventually form new stars (Joung et al. 2012).

6. Summary and Discussion

6.1. Summary of Key Results

Our data provide novel, independent constraints on the coherence scales of IV gas in low and intermediate ionization states at the disk–halo interface. We find that C IV and Si IV are strongly correlated over the full range of transverse separations probed by our sight lines ($\sim 40^\circ$ or 2 kpc at $z = 3 \text{ kpc}$) and with $N_{\text{H I}}$ from 21 cm emission. In contrast, the low ionization state material traced by Ca II and Fe II exhibits neither a correlation with $N_{\text{H I}}$ nor a transverse separation (Figures 8 and 9). This readily apparent difference in ionic behavior indicates the existence of large, coherent ionized structures $\gtrsim 1 \text{ kpc}$ in size, containing cool, neutral cloudlets or clumps with spatial scales $\lesssim 10 \text{ pc}$ ($10.8' \text{ at } z = 3 \text{ kpc}$). We explore the implications of this result in Section 6.2. Finally, we find that the implied areal mass infall rate of the high-latitude WIM encompasses the present-day mean star formation rate surface density of the MW disk, although it is slightly lower than previous estimates from all-sky data.

The Si IV/C IV ratio of the IV gas along our sight lines, -0.69 ± 0.04 , shows remarkably little variation over $\sim 40^\circ$.

We argue that this column density ratio places two mutually exclusive constraints on the physical conditions of the extraplanar gas under two distinct sets of assumptions: photoionization and collisional ionization. In the photoionization case:

1. We assume that Si IV and C IV are produced by photoionization from a combination of starlight from the MW disk and the extragalactic UV background at $z = 0$ (Fox et al. 2005).
2. We adopt constraints from in-depth work that constrains the cool, low ionization state IV gas to have a metallicity of approximately solar value, requiring that cool and warm gas are well mixed in terms of their metal content (Wakker et al. 2008).
3. Finally, we assume that the total hydrogen column associated with the Si IV and C IV is given by the total electron column measured from pulsar DM experiments at high Galactic latitude, with appropriate corrections for He, the HIM, and the neutral gas fraction.

Proceeding in this framework of a photoionized WIM, we find that $\log N_{\text{Si IV}}/N_{\text{C IV}} \approx -0.7$ requires that Si be depleted relative to C by 0.25–0.35 dex compared to the solar value of Si/C. Under these assumptions, this paper would present observational evidence for a dusty WIM at the MW disk–halo interface. We explore the implications of this result in Section 6.3.

The other constraint on the physical state of the material relies only on the production of Si IV and C IV by collisional ionization in a $10^{5.2}$ K gas-phase cooling out of a hotter, 10^6 K ambient medium with a density $< 8 \times 10^{-4} \text{ cm}^{-3}$ (e.g., Bregman & Lloyd-Davies 2007). When we consider the implied total hydrogen content of this transient, cooling phase, we find that there is significant tension if we allow halo $[\text{M}/\text{H}] < 0.3$. Thus, in the collisional ionization framework, our constraint on the total carbon content implies an MW corona that is enriched at a level of greater than one-third solar. This constraint would represent the first observational constraint on the metallicity of the hot corona of the MW.

Additional observations are required to determine which of these two models of the physical state of the IV gas is accurate. For example, access to doubly ionized species such as Si III, S III, Al III, and Fe III would better constrain the photoionization models and possibly provide further evidence of dust depletion in the ionized extraplanar material (e.g., Howk & Savage 1999). The detection of more highly ionized species such as O VI, O VII, or O VIII in the IV gas would be crucial for a complete understanding of the nature of collisionally ionized material in the halo (Faerman et al. 2017). Finally, our results underscore the importance of ionization modeling that considers multiple gas phases under different ionization conditions, with the additional complicating consideration of significant depletion onto dust in ionized extraplanar material. Thus, to convincingly model the physical state of diffuse gas requires observers to study as many ionic species as possible along any given sight line.

6.2. Cloud Length Scales and Physical Models

We find that IV absorption from C IV and Si IV likely traces a coherent ionized structure at the disk–halo interface spanning at least ~ 1 kpc across. The much larger scatter in the low ions

(e.g., Ca II) on the same spatial scales is consistent with a picture in which cool clumps of gas are considerably more compact (< 10 pc in length scale). All of this gas most likely lies within 3.4 kpc of the MW disk and may contain enough total mass to maintain the current star formation rate of the MW.

6.2.1. Comparison with Complimentary Observational Analyses

Small, compact neutral structures are observed directly via 21 cm emission in the MW halo. In particular, as one improves the spatial resolution with which the neutral gas in the MW halo is observed, small, parsec-scale structures are detected with increasing frequency (Lockman et al. 2002; Saul et al. 2012). Our work finds that both ionized and neutral phases are detected along all seven sight lines, moving at similar velocities, indicating that these small, cool clumps are likely contained within a larger-scale highly ionized structure moving toward the MW disk at a velocity of $30\text{--}50 \text{ km s}^{-1}$. Previous work has highlighted the discrepant covering fractions of the low and intermediate ionization state material in the MW halo (Lehner et al. 2012; Richter et al. 2017) as evidence supporting the multiphase structure of the gas. Our evidence for this multiphase structure of the inner MW halo is both complimentary and more direct. In addition, our physical interpretation of the phase structure is consistent with that of Fox et al. (2004), who find the high-density neutral HVCs to be surrounded by much larger envelopes of diffuse, ionized material in absorption-line experiments.

The physical scenario in which both intermediate and low ionization state gas is photoionized by a combination of the EUVB and star formation in the MW disk (Fox et al. 2005) therefore requires that at least two distinct gas phases characterized by different densities (or length scales) be considered. One straightforward means of modeling multiphase, photoionized gas is to have the multiple phases arise from a smooth gas density gradient in the absorbing material. Stern et al. (2016) find that an outwardly declining density gradient of a factor of 100 kpc^{-1} can successfully reproduce the multiphase ion column densities from Mg II to O VI in L^* halos at $z \sim 0.2$. This density gradient implies length scales of Mg II-traced material to be ~ 50 pc, while Si IV-traced material has a characteristic length scale of 3.9 kpc , a factor of ~ 80 difference. We find a factor of ≈ 100 difference in the length scales of Ca II-traced (10 pc) and C IV (Si IV)-traced gas (1 kpc), suggesting a similarly steep gas density gradient if the multiphase gas in the MW disk–halo interface is predominantly photoionized.

Recently, Zheng et al. (2019) conducted a careful, statistical analysis of the Si IV absorption along 132 sight lines that pierce the MW disk–halo interface and found evidence for two components contributing to the absorption: (1) a flat-slab component similar to that of Savage & Wakker (2009), and (2) a global component analogous to that traced by ionized metals observed along extragalactic lines of sight toward background QSOs within projected distances of $10\text{--}150$ kpc from L^* galaxies (e.g., Werk et al. 2013). It is primarily the lack of a correlation between Galactic latitude and $N_{\text{Si IV}}$ that allows them to disfavor a flat-slab model alone, without some global contribution. Explicitly, the flat-slab model tends to predict larger $\log N$ at lower b because path lengths at lower latitudes are longer; such a trend is distinctly absent from the data.

Additional model parameters such as patchiness tend to confound the interpretations, however.

The two-component model of Zheng et al. (2019) predicts that the disk–halo component would contribute $\log N_{\text{Si IV}} = 12.1$ at high Galactic latitudes, whereas our data imply $\log N_{\text{Si IV}} = 13.35$ toward the north Galactic pole within 3.4 kpc. Our results, then, appear to be at odds with the analysis of Zheng et al. (2019). However, we note that relative column density contributions from the disk–halo interface and global components are rather poorly constrained in the two-component model, especially when one allows for patchiness or clumpiness. Our data consist of only seven sight lines toward high-latitude BHB stars and therefore cannot rule out the existence of a global component as put forth by Zheng et al. (2019), especially if the column density of nearby gas dominates over the global component. However, our observations do seem to imply that the relative contribution of the global component to ionized gas column densities in the MW halo is probably lower than that reported by Zheng et al. (2019). Indeed, because Zheng et al. (2019) include only QSO sight lines in their analysis, they likely underestimate the contribution from nearby gas to the observed column densities. Their model could be improved by adding in nearby stellar sight lines, such as those analyzed in this study, to achieve better control of the parameter fitting.

6.2.2. Comparison with Simulations

Observed limits on length scale, prevalence, velocity, and density of cool, extraplanar clouds with $z < 3.4$ kpc can ultimately constrain detailed feedback models that seek to explain the origin and influence of galaxy-scale winds. For example, small clumps < 10 pc in size may affect the relevant cooling processes and thus evolution of material at the disk–halo interface. Recent progress in hydrodynamical modeling has led to increasingly sophisticated treatments of diffuse halo gas that move well beyond the Galactic fountain as it was originally envisioned, in which spontaneous radiative instabilities cause clouds to cool and fall back down onto the disk ballistically (Shapiro & Field 1976; Bregman 1980). Such idealized, high-resolution simulations of Galactic fountain processes, guided by a range of observational constraints, are well poised to address the uncertain physics at the disk–halo interface (e.g., Marinacci et al. 2011; Armillotta et al. 2017; Fraternali 2017; Fielding et al. 2018; Grønnow 2018; Gronke & Oh 2018; Kim & Ostriker 2018; Schneider et al. 2018; Melso et al. 2019; Sormani & Sobacchi 2019). These models include a wide range of complex physics that impact the gas evolution to varying degrees: magnetic fields, halo rotation, conductive interfaces, and stratified turbulent layers.

Here we highlight a few case studies of disk–halo interface or fountain gas to show the promise of high-resolution idealized simulations in untangling the complex gas physics at work in this regime. In the “Three-Phase ISM in Galaxies Resolving Evolution with Star formation and SN feedback” (TIGRESS) simulation suite, Kim & Ostriker (2018) present a comprehensive, self-consistent model of the multiphase star-forming ISM in which a fountain flow arises naturally in warm ionized material energized by local supernovae. The TIGRESS simulation allows for a phase-by-phase analysis of gas flows in a 3D box $1 \text{ kpc} \times 1 \text{ kpc}$ in the horizontal direction and $-4.5 \text{ kpc} < z < 4.5 \text{ kpc}$ in the vertical direction, centered on a region of a galaxy disk with parameters similar to those in the

solar neighborhood. The TIGRESS setup is therefore directly analogous to our MW BHB experiment.

Most saliently, the self-consistent treatment of star formation, FUV emission from massive stars, and supernovae in TIGRESS generates outflows in a hot phase ($T > 5 \times 10^5 \text{ K}$) and periods of upwelling and subsequent fallback (a fountain flow) in the cool, warm, and ionized phases with $T < 5 \times 10^5 \text{ K}$. The temporal variation in the SFRs leads to ~ 50 Myr periods in which the warm and ionized gas phases are dominated by successive periods of outflow and inflow. Thus, in this rendering of the Galactic fountain, our observations would indicate that the WIM above the solar neighborhood is currently experiencing a temporary inflow period. Furthermore, the gas inflow velocities in the simulations of $\sim 60 \text{ km s}^{-1}$ are similar to our values of $v \sin b$, which range between 40 and 50 km s^{-1} for the warm ionized gas traced by Si IV and C IV.

Additionally, TIGRESS provides a natural explanation for why we do not observe substantial outflowing material in conjunction with inflowing material in our area. First, $> 70\%$ of the outflowing gas by mass is in a hot phase inaccessible by our COS observations. Second, although inflows and outflows are contemporaneous in the simulations in every phase, warm gas, when outflowing, is pressure confined by the hot gas into small clodlets, which would have a comparably small covering fraction. Once an episode of star formation ceases, the warm ionized material falls back and expands to fill nearly an order of magnitude more volume than when it is in the outflowing period.

Finally, we note the areal mass infall rate in the warm ionized fountain flows in the TIGRESS simulations range between 0.001 and $0.01 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$ between 1 and 3 kpc above the disk, which is similar to the time-averaged star formation rate of their simulated disk of $\sim 0.006 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$. Our areal mass infall rate estimated in Section 5.3 is similar, $0.0003\text{--}0.006 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$ within 3 kpc. We note that several facets of the simulations impair such direct comparisons with observations. TIGRESS does not include photoionization, so it is not clear how the gas column densities we measure (e.g., $N_{\text{C IV}}$) translate to a warm, ionized phase in the model. In addition, TIGRESS does not include a global geometry for the simulated galaxy, which may matter for any larger-scale gas flows, especially those that may originate in the center of the galaxy, where star formation is more concentrated.

Other high-resolution idealized simulations inform cloud survival times and mass-loading rates of the gas flows in the Galactic fountain. Armillotta et al. (2017) find from extremely high resolution 2D simulations of the evolution of cool clouds passing through an ambient hot medium that the characteristic length scale of cool gas clumps will influence their survival times in a halo environment, such that cool clouds > 250 pc in size are able to retain their cool gas mass (although they are subject to fragmentation owing to thermal instabilities) over > 250 Myr timescales. However, the recent analysis of 3D simulations of cold cloud evolution by Gronke & Oh (2018) concluded that the survival time of these structures is in fact governed by the cooling time of “mixed,” intermediate-temperature material surrounding the cloud, rather than by the formal cloud-crushing time. This enables the survival of much smaller cold structures with lengths $\gtrsim 2$ pc and could potentially explain the existence of such small structures consistent with our observations. In the Fielding et al. (2018) study of clustered supernova-driven winds in idealized 3D hydrodynamical simulations of sections of galaxy disks, the

ability of superbubbles to vent into the halo is regulated in part by the disk scale height. Once breakout occurs, the wind mass loading is determined by the mixing and cooling of material stripped off of entrained cold clouds (Fielding et al. 2018). Finally, the CGOLS suite of high-resolution isolated galaxy models likewise demonstrates that interactions between fast-moving outflow material and small, cool, dense clouds can lead to increased mass loading of the hot phase, which has dramatic implications for the impact of feedback in galactic halos (Schneider et al. 2018).

We note that length scales <100 pc are not currently resolvable in state-of-the-art cosmological hydrodynamical models that aim to probe diffuse halo gas. Such simulations now reach spatial resolutions in the CGM on the order of hundreds of parsecs (Hummels et al. 2019; Peebles et al. 2019; Suresh et al. 2019; van de Voort et al. 2019) and find that even a factor of ~ 10 increase in spatial resolution can produce dramatic effects on the physical processes and phase structure of the gas. In these simulations, there is a substantial neutral hydrogen abundance in clouds with sizes <1 kpc (and such high-density clouds are predicted to occur on even smaller scales as resolution is increased yet further). Our results support the conclusions of these studies, namely, that higher-resolution simulations in a cosmological context are required to trace the relevant physical processes and phase structure of the CGM and disk–halo interface. Recent progress, however, indicates that a full understanding of the Galactic fountain, including its timescales, extent, and mass and energy loading factors, is on the horizon.

6.3. A Dusty WIM

Our photoionization modeling of extraplanar Si IV and C IV column densities suggests the presence of Si-bearing dust in the high-latitude WIM. Silicon depletions on the order of a few tenths of a dex are consistent with levels of depletion in DLAs of $\delta_{\text{Si}} = -0.27 \pm 0.16$ dex (measured with Si/Zn) assuming that carbon is not depleted itself (Vladilo et al. 2011). Si/C ratios of ~ -0.7 appear to occur in the neutral MW ISM at a value of F^* of 0.3–0.5, where F^* is the line-of-sight depletion strength factor (Jenkins 2009) that increases with increasing depletion. Notably, the 1σ sight-line-to-sight-line variation in Si/C we observe is only ± 0.04 dex. Under the assumption of photoionization, this constant ratio would indicate remarkably uniform dust coverage over the kiloparsec-sized spatial scales of the material traced by Si IV and C IV.

We note that depletion onto dust from ionized material is not entirely unexpected. Howk & Savage (1999) discovered the first evidence for dust in the diffuse ionized gas of the Galactic ISM from modeling the observed ratio of Al III/S III, which traces the ratio of a depleted element (aluminum) to a nondepleted element (sulfur) in the ISM. Their result, like this one, is contingent on S III and Al III not having significant contributions from collisional ionization. Further confirmation of dust in an ionized phase came from Lagache et al. (1999, 2000), who first detected far-IR emission from high-latitude dust using *COBE* and H I data and refined this measurement with data from the Wisconsin H α Mapper (WHAM). Additionally, Dobler et al. (2009) find evidence for a spinning dust component of the MW WIM by analyzing the H α -correlated emission spectrum of the 5 yr *WMAP* data. Finally, a study by Gringel et al. (2000) found an absorption component of molecular hydrogen in an ORFEUS-Echelle spectrum at -62 km s $^{-1}$, which they associated

with the MW halo, a sign that there are grains present on which H₂ can form. Both absorption and emission studies support a picture in which dust temperature and emissivity are similar between neutral and ionized gas, and that the processing of dust grains in the WIM is similar to that in the H I gas.

Dust may persist up to several megaparsecs from the disks of galaxies, as traced by a reddening coefficient similar to that of interstellar dust (Ménard et al. 2010; Peek et al. 2015). Furthermore, a substantial fraction of the metals generated within the galaxy may be lurking in a circumgalactic dust phase (Peebles et al. 2014). The coexistence of dust and highly ionized gas in galactic coronae has far-reaching implications for self-shielding, magnetic fields, molecule formation, gas cooling, and grain destruction in rarefied environments. Future studies of gas physics in the disk–halo interface and the CGM should incorporate the physics of dust and allow for nonsolar ratios of the elements to accurately interpret observations with hydrodynamical simulations.

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Facility: *HST*: COS Keck: HIRES.

Software: CALCOS (<https://github.com/spacetelescope/calcos/>); MPFIT1 (<http://cow.physics.wisc.edu/~craigm/idl/fitting.html>); CLOUDY (<https://www.nublado.org/>).

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