



Ultra-high-quality two-dimensional electron systems

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Two-dimensional electrons confined to GaAs quantum wells are hallmark platforms for probing electron-electron interactions. Many key observations have been made in these systems as sample quality has improved over the years. Here, we present a breakthrough in sample quality via source-material purification and innovation in GaAs molecular beam epitaxy vacuum chamber design. Our samples display an ultra-high mobility of $44 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ at an electron density of $2.0 \times 10^{11} \text{ cm}^{-2}$. These results imply only 1 residual impurity for every 10^{10} Ga/As atoms. The impact of such low impurity concentration is manifold. Robust stripe and bubble phases are observed, and several new fractional quantum Hall states emerge. Furthermore, the activation gap (Δ) of the fractional quantum Hall state at the Landau-level filling ($\nu = 5/2$), which is widely believed to be non-Abelian and of potential use for topological quantum computing, reaches $\Delta \approx 820 \text{ mK}$. We expect that our results will stimulate further research on interaction-driven physics in a two-dimensional setting and substantially advance the field.

Single-crystal GaAs thin-film structures grown by molecular beam epitaxy (MBE) are often considered to be among the purest materials that can be made in the laboratory. Being nearly defect free, these structures provide an exceptional platform for exploring a diverse range of physical sciences, with extensive electronic and photonic applications. The highlight of ultra-high-quality GaAs films, however, is their utilization in the investigation of electron-electron interaction phenomena. Typically this is achieved by studying the low-temperature magnetotransport of two-dimensional (2D) electron systems (2DESs) hosted in modulation-doped GaAs quantum wells, where the electrons are spatially separated from the dopants to reduce scattering between the electrons and the intentional impurities. A magnetic field applied perpendicular to the 2DES enhances the relative scale of Coulomb energy in the system by quenching the Fermi energy via Landau quantization. A plethora of many-body phases have materialized in GaAs 2DESs using this framework; the discovery of the odd-denominator¹ and even-denominator² fractional quantum Hall (FQH) effects as well as the observation of Wigner solid^{3–6}, stripe or nematic^{7,8} and bubble phases⁹ are some notable examples (for reviews, see the literature^{10,11}).

Naturally, these fascinating phases only started to emerge in experiments as sample quality increasingly improved. For example, after the first observation of the FQH effect¹, it took about two decades of growth condition improvements to realize stripe or nematic phases in GaAs 2DESs^{7,8}. These developments continue to motivate the community to search for methods to produce ever-better-quality samples, as it is stimulating to anticipate what other physics is yet to be uncovered. Indeed, these efforts have extended to several other material systems with varying electronic properties, and numerous many-body phases have been observed in the magnetotransport/capacitance features of 2DESs hosted in AlAs^{12–17}, Si^{18,19}, graphene^{20–24}, ZnO^{25,26}, Ge^{27,28} and WSe₂ (ref. ²⁹). While each and every one of these systems provides the opportunity to study exquisite interaction-driven physics in different perspectives, GaAs-based 2DESs remain important; see the literature for examples^{30,31}.

A useful metric to quantify the quality of a 2DES is the electron mobility because it is straightforward to measure and is inversely proportional to the average electron scattering rate in the 2DES. High mobility values in 2DESs imply that electrons are less likely to scatter over prolonged trajectories in such samples, making them valuable platforms to study delicate many-body electron phases as well as ballistic or phase-coherent transport. After decades of development, electron mobility values as high as $\mu \approx 35 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ have been observed in GaAs 2DESs with densities near $n \approx 3 \times 10^{11} \text{ cm}^{-2}$ (refs. ^{32–35}). However, despite subsequent efforts, the mobility of modern state-of-the-art GaAs 2DESs has been in a stalemate for more than a decade now. Here we present a breakthrough in the MBE crystal growth of ultra-high-quality GaAs 2DESs that enhances the mobility to $\mu \approx 44 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ at electron densities as low as $n \approx 2 \times 10^{11} \text{ cm}^{-2}$. In the low-density ($n \lesssim 1.5 \times 10^{11} \text{ cm}^{-2}$) regime, where scattering by residual (background) impurities is dominant, our samples show twice the mobility of previous samples^{32–35}. Low-temperature magnetotransport traces taken in a sample of this class show new FQH states, indicating great prospects for future studies of interaction-driven physics in the 2D setting.

Numerical calculations have suggested that residual impurities are the limiting factor for mobility in ultra-high-quality GaAs 2DESs^{36–38}. Efforts to systematically purify the Ga and Al source materials are in line with this understanding, as these metals are the most likely origin of unwanted impurities in a well-maintained, ultra-high-vacuum growth chamber^{35,39,40}. Despite the different history of the source materials loaded into the MBE chambers and the varying purification techniques, ultra-high-quality samples from several groups around the world seem to converge at roughly similar values in electron mobility^{32–35}. These results hint that the impurities that are limiting the mobility in the modern ultra-high-quality GaAs 2DESs come from somewhere other than the source material.

While the environment in ultra-high-vacuum MBE chambers is certainly very sparse in atomic and molecular density, it is not completely void of matter. Even in well-equipped, conventional vacuum chambers, it is common that the mass spectrometer data of the

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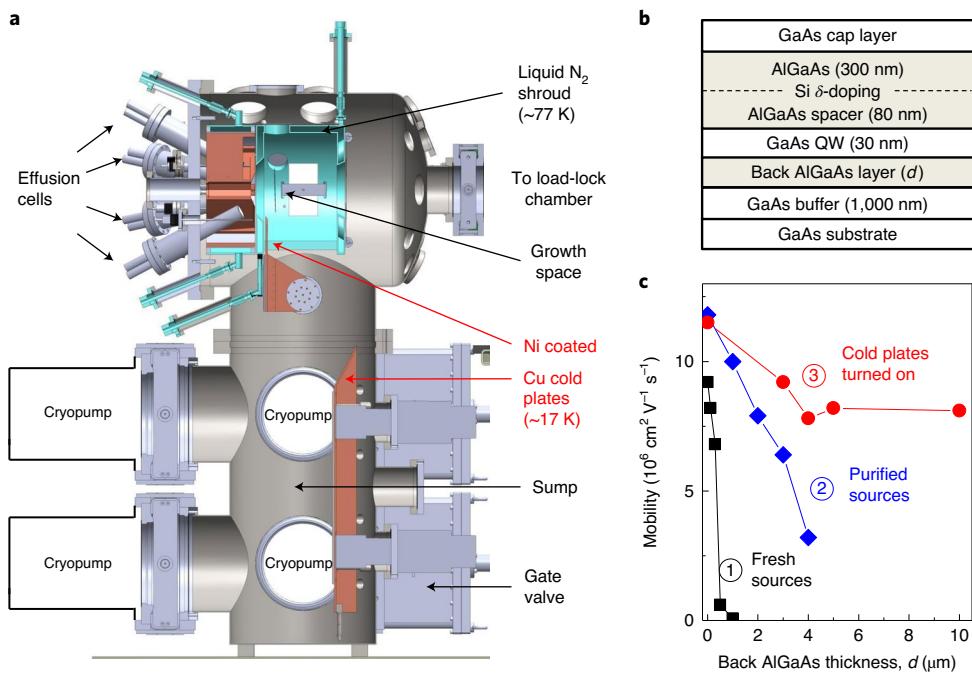


Fig. 1 | Improving vacuum quality and its assessment in a state-of-the-art MBE chamber. **a**, Schematic diagram of the MBE chamber used in this study. In addition to four large ($3,000\text{ s}^{-1}$) cryopumps operating at $\sim 10\text{ K}$, there are three auxiliary cryo-cooled Cu cold plates ($\sim 17\text{ K}$ on their back side) that pump the chamber. One $5 \times 5\text{ inch}^2$ cold plate is located in close proximity to the growth space, while two $12 \times 15\text{ inch}^2$ cold plates are in the sump. **b**, Sample structure used to analyse vacuum quality. A δ -function-like Si-dopant profile is introduced to the top side of the GaAs quantum well (QW) to define a 2DES in the QW. If the source material is pure enough, the surface segregation of impurities during the growth of the back AlGaAs layer can be exploited to evaluate the vacuum quality by analysing the mobility of the GaAs 2DES. **c**, Mobility of samples that have the structure shown in **b**, with varying back AlGaAs thicknesses (d) and different vacuum and source purity conditions. The 2DES density is $n \approx 2.0 \times 10^{11}\text{ cm}^{-2}$. The datasets for different conditions are colour coded, where the data shown in black, blue and red were taken from samples grown in that chronological order; they correspond to samples grown with (1) a fresh batch of sources, (2) purified sources and the Cu cold plates turned on, respectively.

growth space show traces of H_2O , N_2 , O_2 , CH_4 and their derivatives. If incorporated during growth, these species would act as impurities in the structure and cause degradation in the quality of the GaAs 2DES. Once the source materials become so pure that they are no longer the primary supplier of impurities to the growth environment, the background vacuum quality would determine the impurity concentration in the sample.

As shown in Fig. 1a, we built an MBE chamber to verify this hypothesis. Our design includes conventional MBE chamber components such as a load-lock chamber and a liquid N_2 shroud, as well as three auxiliary cryo-cooled ($\sim 17\text{ K}$) cold plates that augment four large ($3,000\text{ s}^{-1}$) cryopumps to achieve extreme levels of vacuum during sample growth. The cold plates are made of Cu to maximize cooling power and are coated with Ni to prevent possible corrosion when exposed to As and Ga, which are ubiquitously present in the growth chamber. It is difficult to quantitatively assess how much the base pressure improved in our growth chamber with the cold plates operating because the ion extractor gauges installed in the chamber cannot reliably measure pressures below $P \sim 2\text{--}3 \times 10^{-12}\text{ torr}$, and the base pressure already reaches this range even when only the four cryopumps are pumping on the chamber. However, when the cold plates are turned on, the mass spectrometer data shows a factor of 10 improvement in the partial pressures of N_2 and O_2 species and a factor of 2.5 improvement in the partial pressures of H_2O -related species. The mass spectrometer data, as well as more details concerning the cold plates, can be found in Section I of the Supplementary Discussion.

As previously demonstrated, it is possible to systematically evaluate the cleanliness of the growth environment during the MBE of GaAs-AlGaAs heterostructures¹⁰. A brief summary of the concept

is to use the mobility of a GaAs 2DES with the sample structure shown in Fig. 1b as a very sensitive probe to gauge impurity accumulation during growth. This method utilizes the fact that impurities surface-segregate on the growth front of the back AlGaAs layer and deposit at the AlGaAs-GaAs interface when the growing layer is changed from AlGaAs to GaAs. For a given back AlGaAs layer thickness, the mobility is lower when the growth environment is worse. The strength of this procedure is that the back AlGaAs layer thickness (d) can be made extremely large to detect even the most minute amounts of impurities incorporated in the crystal during the growth process.

Figure 1c shows the mobility of such structures grown in varying growth environments as a function of d . The enhanced mobility of samples grown after the sources were sufficiently outgassed (blue symbols) compared to those grown with a fresh batch of source materials (black symbols) demonstrates the importance of preparing pure source materials when aiming for the cleanest growth conditions. Under these conditions, the source materials were clean enough that high-mobility GaAs 2DESs made from them displayed mobilities on par with ultra-high-quality samples in the literature. Despite further purification efforts, we could not obtain a striking improvement from the blue dataset under the normal operating vacuum conditions of having four cryopumps operating. This implied that at this point, our source materials had been amply purified so that they were no longer the primary supplier of impurities for our samples. It is then plausible to assume that the background vacuum starts to play a more important role. Consistent with this assumption, samples grown with only one cryopump operating exhibited worse mobilities compared to those grown with all four pumps turned on. Following this test, we investigated a series of

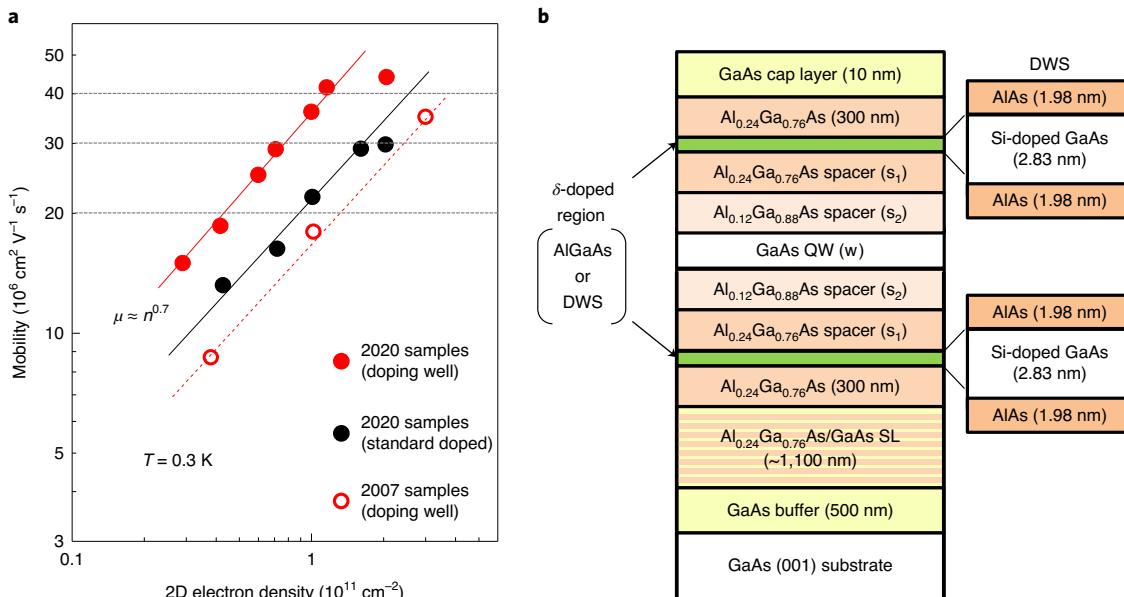


Fig. 2 | Mobility versus 2D electron density for our GaAs 2DESS. a, Data from samples grown under the improved vacuum conditions are shown as solid symbols for the DWS (red) and standard modulation-doped structure (black). The mobility values from our previous ultra-high-quality doping-well samples are shown as red open symbols for comparison. These previous samples have electron mobilities similar to ultra-high-quality samples reported in the literature^{32–35}. The solid and dashed lines through each set of data points serve as guides to the eye, and their slopes are consistent with a power-law relation that is roughly $\mu \approx n^{0.7}$. **b**, Layer structure of the samples used in **a**. The spacer thicknesses (s_1, s_2) and well widths (w) of some representative samples can be found in Section II of the the Supplementary Discussion. A superlattice (SL) structure is implemented in all samples to reduce any potentially detrimental effects from the substrate.

samples grown with all the cryopumps and the cold plates operating during growth (data shown in red in Fig. 1c). The data reveal that there is a notable improvement in GaAs 2DES mobility at all values of d compared to the case when the cold plates are off. Remarkably, the mobility of the samples grown with all cold plates operating sustains the high value of $\mu \approx 8 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ even when $d = 10 \mu\text{m}$.

These results demonstrate that vacuum integrity plays a crucial role in determining the amount of unintentional impurities deposited on the sample during growth once the source material has been extensively purified. Consequently, we grew several GaAs samples with a wide range of 2D electron densities to investigate the impact of having an ultra-clean vacuum environment in the MBE chamber. Figure 2a compares the mobility of these samples with the previous ultra-high-quality GaAs 2DESS^{32–35}. Figure 2b shows the layer structure of the samples used to obtain the data presented in Fig. 2a. We used two types of structure: the standard modulation-doped structure and the doping-well structure (DWS)⁴¹. While the donor energy level, which determines the position of the Fermi level, is tied to the AlGaAs barrier in the standard modulation-doped structure, in the DWS it is tied to the narrow AlAs layers that flank a narrow GaAs doping quantum well⁴¹. The DWS is advantageous in comparison to the standard modulation-doped structure because the electrons confined to the AlAs layer in the doped region provide additional screening for the 2DES from both residual impurities and intentional dopant ions^{38,41}. It is striking that with the improvement in vacuum, even the standard modulation-doped samples (black circles in Fig. 2a) have considerably higher mobility values for all densities when compared to previous state-of-the-art DWSs (red open circles). This is particularly noteworthy considering that DWSs were necessary to achieve the previously reported ultra-high-mobility values^{32–35}. It seems that the reduction in impurities from better vacuum conditions by implementing the cold plates is pronounced enough to overcome the lack of such screening in our standard modulation-doped structures.

Furthermore, when we grow DWSs with the cold plates operating, we see an even larger increase in the mobility (red solid circles in Fig. 2a). These samples display mobility values as high as $\mu \approx 44 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ at the density of only $n \approx 2.0 \times 10^{11} \text{ cm}^{-2}$. This implies a remarkable enhancement in sample quality considering that previous ultra-high-quality GaAs 2DESSs had mobilities of $\mu \approx 35 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ at $n \approx 3.0 \times 10^{11} \text{ cm}^{-2}$ (refs. ^{32–35}). Sample improvement is evident over a wide range of 2D electron densities as shown in Fig. 2a. When $n < 1.5 \times 10^{11} \text{ cm}^{-2}$, our samples have mobility values that are roughly twice that of previous ultra-high-quality samples. For example, the mobility of our $n \approx 1.0 \times 10^{11} \text{ cm}^{-2}$ sample is $\mu \approx 36 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, whereas previous ultra-high-mobility values for GaAs 2DESSs with a similar density are less than $\mu \approx 18 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ (refs. ^{35,42}).

The power-law dependence observed for the mobility versus 2D electron density profiles plotted in Fig. 2a is also noteworthy. Within the same category of samples, we observe a $\mu \propto n^{0.7}$ relation for all cases. A similar power-law dependence was reported in high-quality, low-density GaAs 2DESSs with very large spacer-layer thicknesses, and it is usually interpreted as an indication that the mobility is limited by the residual impurities in the structure^{33–45}; there is also theoretical justification for such an interpretation⁴⁶. This understanding is certainly consistent with our results, as the primary improvement we have made to our samples compared to the previous data is the reduction of residual impurities in the structure. Based on models for the two types of structures used for our samples^{38,46}, we estimate that the residual impurity concentration in our GaAs quantum wells is $\sim 1 \times 10^{13} \text{ cm}^{-3}$. Considering that there are $\sim 1 \times 10^{23} \text{ atoms cm}^{-3}$ in single-crystal GaAs, this means there is roughly one impurity for every 10 billion atoms in these samples.

As mentioned in the introduction, probing intricate many-body phenomena is a core application of ultra-high-quality GaAs 2DESSs. In this context, we also studied the low-temperature ($T \approx 30 \text{ mK}$) magnetotransport of one of our representative samples with

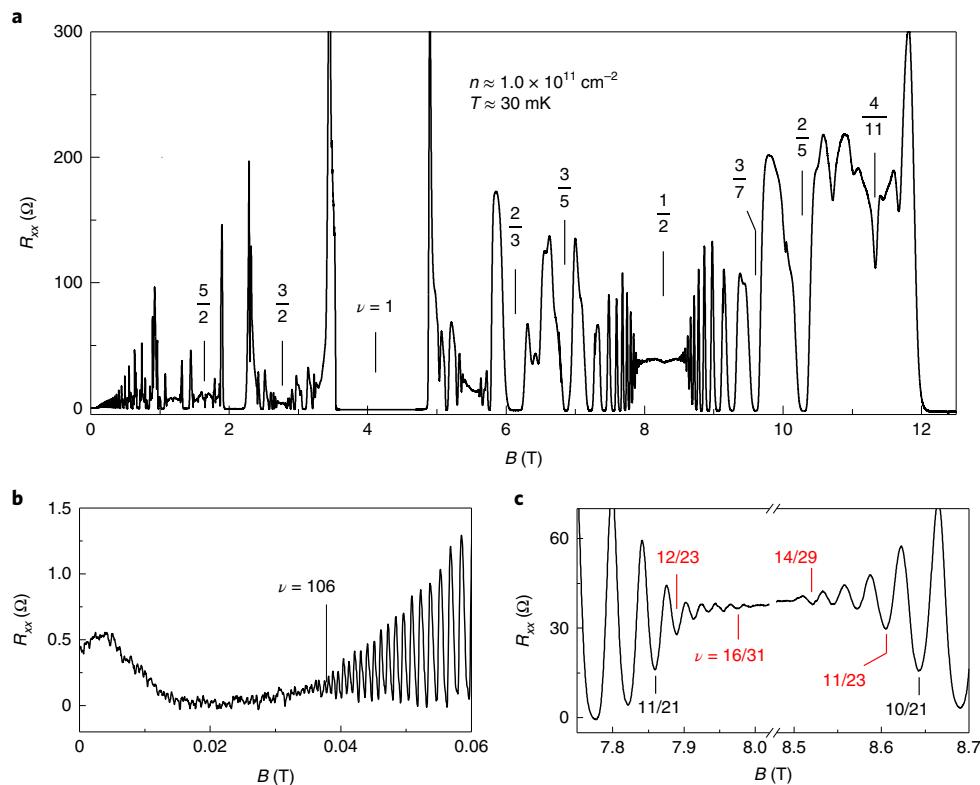


Fig. 3 | Low-temperature ($T \approx 30 \text{ mK}$) magnetoresistance data of a GaAs 2DES with density $n \approx 1.0 \times 10^{11} \text{ cm}^{-2}$. **a**, A full-field R_{xx} versus magnetic field trace. The Landau-level fillings (ν) of several quantum Hall features are marked. **b**, Expanded view of the low-field magnetoresistance near $B=0$. Resistance minima at fillings as high as $\nu=106$ can be seen in the Shubnikov-de Hass oscillations. **c**, Expanded view of the magnetoresistance near $\nu=1/2$. High-order FQH states up to $\nu=11/21$ and $10/21$ have been observed in previous ultra-high-quality samples with similar electron density (marked in black)⁴². New FQH states at $\nu=12/23$, $13/25$, $14/27$, $15/29$ and $16/31$ are seen on the left of $\nu=1/2$, and FQH states at $\nu=11/23$, $12/25$, $13/27$ and $14/29$ on the right flank of $\nu=1/2$. The lowest-order and highest-order new FQH state on each side of $\nu=1/2$ is marked in red.

$n \approx 1.0 \times 10^{11} \text{ cm}^{-2}$. The specific layer structure for this sample is provided in Supplementary Table 1. Figure 3a shows a full-field longitudinal magnetoresistance (R_{xx}) trace of the sample, while Fig. 3b,c focuses on specific regions near zero magnetic field and near $\nu=1/2$, respectively; $\nu=hn/eB$ is the Landau-level filling factor, where h is the Planck constant, e is the fundamental charge and B is the perpendicular magnetic field. It is clear from the data that the sample has very high quality. For example, as seen in Fig. 3b, there are prominent signatures of Shubnikov-de Hass oscillations up to $\nu=106$ at $B < 0.04$ T. This implies that the Landau-level broadening in this sample is smaller than the $B=0.04$ T cyclotron energy gap of $heB/m^* \approx 68 \text{ meV}$ ($m^*=0.067$ is the effective mass of electrons in GaAs in units of the free-electron mass). In addition, the data plotted in Fig. 3c display high-order FQH states up to $\nu=16/31$ and $\nu=14/29$ on the left and right flanks of $\nu=1/2$. We compare these results to those reported previously for epitaxially grown samples and 2D materials with ultra-high quality. In total, compared to previous ultra-high-quality GaAs samples with similar density⁴², nine extra FQH states are observed near $\nu=1/2$ in our sample, whose lowest-order and highest-order Landau-level fillings are marked in red on each side of $\nu=1/2$. By comparison, in ultra-high-quality monolayer graphene samples, high-order FQH states have been observed up to $\nu=8/15$ and $\nu=7/15$ on the left and right flanks of $\nu=1/2$ at similar temperatures but higher magnetic fields (~ 14 T)²⁴; the data presented in Fig. 3 exhibit 15 additional FQH states with respect to these samples.

The quality of our samples in this density range also stands out at higher Landau-level fillings. Figure 4a,b shows expanded R_{xx} traces of the Fig. 3 sample near $\nu=3/2$ and $\nu=5/2$, respectively.

Remarkably, even at this relatively low density, FQH states up to $\nu=20/13$ are observable in the vicinity of $\nu=3/2$. Furthermore, the $\nu=5/2$ and other FQH states in the second orbital Landau level ($N=1$) are extraordinarily strong, considering that they occur at $B < 1.9$ T. In fact, the activation gap we measure for the $\nu=5/2$ FQH state is $^{5/2}\Delta \approx 820 \text{ mK}$ (Supplementary Fig. 4), surpassing the value of $^{5/2}\Delta \approx 625 \text{ mK}$ measured in previous ultra-high-quality samples by a considerable margin⁴⁷. This observation is particularly noteworthy given that the density of this sample is only $\sim 1/3$ of that of the 2DESS used in previous evaluations.

Given the potentially non-Abelian nature of the $\nu=5/2$ FQH state^{48,49}, the data presented here have exciting implications for the realization of fault-tolerant, topological quantum computing devices. Using optimistic estimates for qubit error rate⁴⁸, our larger $^{5/2}\Delta$ value implies a factor of $\sim 10^6$ improvement compared to previous ultra-high-quality samples when operations are performed at $T=5 \text{ mK}$. In principle, if the samples presented here do not suffer from severe detrimental effects caused by lithographic procedures and are amenable to gating, one should be able to perform substantially more robust qubit operations at much lower magnetic fields. Previously, samples with similar structure have been used for interferometry experiments that can be considered a basis for qubit operation in GaAs 2DESS⁴⁹, so we are optimistic that the technical details can be worked out.

Figure 4c shows magnetotransport data of a different sample with a similar density of $n \approx 1.1 \times 10^{11} \text{ cm}^{-2}$ in magnetic field ranges that correspond to higher ($N=2$ and 3) orbital Landau levels. Well-quantized, reentrant integer quantum Hall states, as well as signatures of stripe or nematic phases, are observable at magnetic

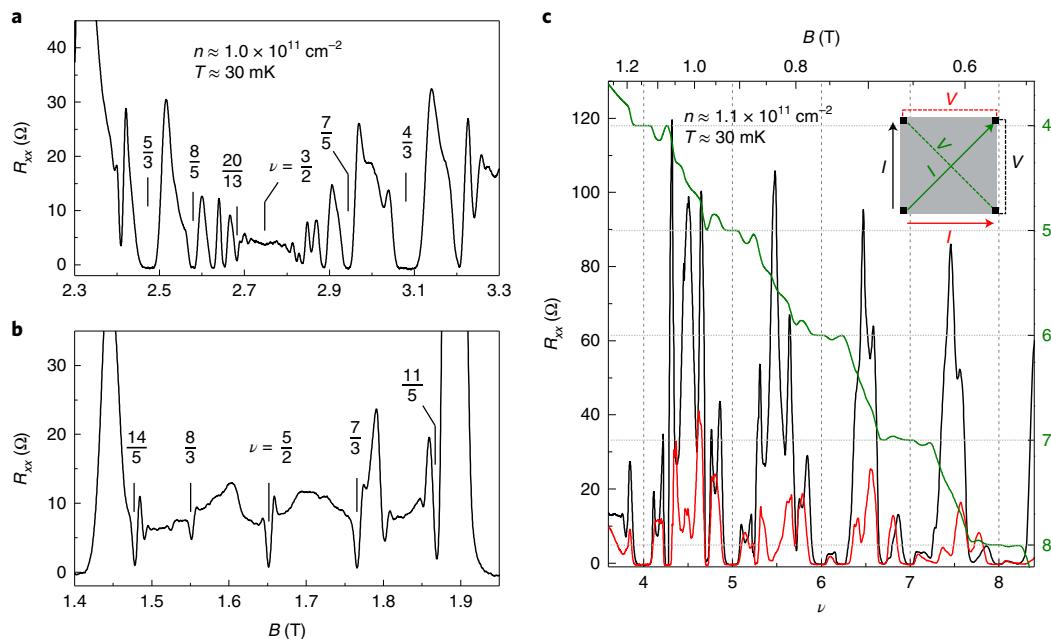


Fig. 4 | Low-temperature ($T \approx 30$ mK) magnetotransport data of $n \approx 1.0 \times 10^{11}$ cm $^{-2}$ samples at higher Landau-level fillings. a, b, Magnetoresistance data near $\nu = 3/2$ ($N = 0$ orbital Landau level, **a**) and $\nu = 5/2$ ($N = 1$, **b**). The number of observed FQH states as well as the strengths of each of the features are remarkable, considering the relatively low density of the sample. **c**, Longitudinal (black and red) and Hall (green) resistance data for the high-index ($N = 2$ and 3) orbital Landau levels. The presence of reentrant integer quantum Hall states and stripe or nematic phases is evident at magnetic fields $B < 1$ T. The inset shows the configuration of current (I) and voltage (V) contacts used for the measurements.

fields $B \lesssim 1.0$ T. These correlated states are known to be fragile, and are typically only observed in much higher density samples at larger magnetic fields and lower temperatures^{7,8,30,50}.

The results presented here suggest a bright future for the investigation of interaction-driven physics in GaAs 2DESSs. With this improvement in sample quality, several new FQH states have emerged, and numerous correlated phases display strong robustness. For example, the $\nu = 5/2$ FQH state exhibits a gap value of $5/2\Delta \approx 820$ mK, and we demonstrate that reentrant integer quantum Hall and stripe or nematic states are visible even at very low magnetic fields and electron densities. Moreover, we have experimentally shown that vacuum integrity limits sample quality in current state-of-the-art MBE-grown GaAs. This gives a clear direction for further improvement in the quality of GaAs 2DESSs. We speculate that further reducing the residual background impurity concentration by a factor of ~ 3 should allow one to obtain mobility values exceeding 10^8 cm 2 V $^{-1}$ s $^{-1}$ as predicted by theory³⁶. Perhaps this can be realized by implementing even more cold plates or pumps in the MBE chamber.

Some mysteries have also developed. When $n \geq 1.5 \times 10^{11}$ cm $^{-2}$, the electron mobilities seem to deviate to lower values than the power-law relation $\mu \propto n^{0.7}$ would predict. We are currently unsure of the origin of this behaviour. It is possible that the remote ionized impurities from the intentional dopant atoms become relevant as a smaller spacer thickness is required to achieve higher 2D electron densities. If this is the case, in the future it may be useful to start from a low-density sample with large spacer thickness and increase the density by applying gate voltages to circumvent this issue. Another option is to vary structural parameters in the doped region of the DWS so that the spacer thickness can be increased while maintaining the same 2DES density in the main quantum well⁴¹.

Online content

Any methods, additional references, Nature Research reporting summaries, source data, extended data, supplementary information,

acknowledgements, peer review information; details of author contributions and competing interests; and statements of data and code availability are available at <https://doi.org/10.1038/s41563-021-00942-3>.

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Methods

Sample preparation. All of our samples are grown on 2-inch-diameter GaAs substrates in the vacuum chamber set-up shown in Fig. 1a. The substrates are outgassed for 30 minutes at $T \approx 610^\circ\text{C}$ in an As beam flux of $P \approx 6.0 \times 10^{-6}$ torr prior to growth, where the temperature of the substrate is evaluated using a factory-calibrated pyrometer (Incon Modline 7). We always confirm clear single-crystalline features in the reflection high-energy electron diffraction patterns of the substrate after this process. The substrate temperature is typically $T \approx 640^\circ\text{C}$ during growth. This growth temperature was chosen based on our past experience for growing ultra-high-quality samples. For such samples, the typical temperature window for optimal growth was approximately $\pm 10^\circ\text{C}$. The deposition rate of GaAs is calibrated to be $\sim 2.83 \text{ \AA s}^{-1}$ for all growth by tuning the temperature of the Ga oven based on reflection high-energy electron diffraction oscillations. We tune the Al growth rate in a similar fashion to obtain the barrier alloy fraction of choice. The barrier alloy fraction is 32% for the samples whose data are shown in Fig. 1c, while the ultra-high-mobility samples use a stepped-barrier structure with alloy fractions 24% and 12%. The specifics of the sample structure of the ultra-high-mobility samples can be found in Supplementary Discussion Section II.

Transport measurements. We performed all electronic measurements in the van der Pauw configuration using low-frequency lock-in amplifiers. Our samples have a square shape and a typical size of $4 \text{ mm} \times 4 \text{ mm}$. The mobility values of the GaAs 2DESs are evaluated in a ^3He cryostat with a base temperature of $T \approx 0.3 \text{ K}$. A simple Drude formula $\mu = 1/\rho ne$ is used to obtain the mobility, where n is the 2DES density, e is the fundamental electron charge and ρ is the resistivity of the 2DES. Quantum Hall features in the magnetoresistance data are used to deduce n . For ρ , we take the average value of the resistance (R_{ave}) measured between all the four-probe contact configurations in the sample, and use the standard van der Pauw geometry expression $\rho = \pi R_{\text{ave}}/\ln(2)$. The low-temperature magnetotransport data presented in the main text are measured in a dilution refrigerator with a base temperature of $T \approx 30 \text{ mK}$. It is well known that the illumination and cooldown procedure can affect the electronic properties of the GaAs 2DES at low temperatures. For the measurements in this work, we illuminate all samples for 5 minutes at $T \approx 10 \text{ K}$ with a red light-emitting diode (LED) before turning the LED off and waiting for an additional 30 minutes at $T \approx 10 \text{ K}$. A current of 6 mA is passed through the LED during illumination. We then cool the samples to the base temperature. A similar procedure is adopted for the dilution refrigerator measurements, but the LED illumination is done at $T \approx 4 \text{ K}$ instead of 10 K because

of the limitations of the apparatus. By repeating this procedure, we saw less than 5% variance in the 2DES density and mobility even when the samples experienced a full thermal cycle to room temperature. For the magnetotransport measurements, a magnetic field sweep rate of 1 T per hour was used, except for the case of Fig. 1b, where a slower sweep rate of 0.1 T per hour was used.

Data availability

Data supporting the results in this paper and the Supplementary Information are available on request to the corresponding author. Source data are provided with this paper.

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Author contributions

Y.J.C. and L.N.P. conceived the work. K.W.B., K.W.W. and L.N.P. designed and built the MBE chamber. Y.J.C., K.W.B., K.W.W. and L.N.P. designed, grew and evaluated the quality of all samples at $T \approx 0.3 \text{ K}$. K.A.V.-R. and P.T.M. performed the dilution refrigerator measurements. Y.J.C. and M.S. wrote the manuscript with input from all co-authors.

Competing interests

The authors declare no competing interests.

Additional information

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