Battery characterization via eddy-current imaging with nitrogen-vacancy centers in diamond

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Sensitive and accurate diagnostic technologies with magnetic sensors are of great importance for identifying and localizing defects of rechargeable solid batteries in a noninvasive detection. We demonstrate a microwave-free AC magnetometry method with negatively charged NV centers in diamond based on a cross-relaxation feature between NV centers and individual substitutional nitrogen (P1) centers occurring at $51.2 \, \text{mT}$. We apply the technique to non-destructive solid-state battery imaging. By detecting the eddy-current-induced magnetic field of the battery, we distinguish a defect on the external electrode and identify structural anomalies inside the battery body. The achieved spatial resolution is $360\,(2)\,\mu\text{m}$ at a $100\,(0.5)\,\mu\text{m}$ distance between the sample and the probe.

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I. INTRODUCTION

Sensitive and accurate diagnostic technologies with magnetic sensors are of great importance for identifying and localizing defects of rechargeable solid batteries in a noninvasive detection. Negatively charged nitrogen-vacancy (NV) centers have been extensively exploited as precise nanoscale probes in applications such as measurement of magnetic fields ¹, temperature ², strain ³, rotation ^{4,5}, electric fields ^{6,7} and so forth. Recently, NV centers have also been utilized for eddy-current imaging. Eddy-current detection was demonstrated with vapor-cell magnetometers ^{8–10}, and later also with NV diamond ¹¹. Compared with other sensors, diamond-based devices can be used over a wide temperature range, can have nanoscale spatial resolution, high sensitivity and wide bandwidth. In ref. ¹¹, a microwave-free eddy-current imaging device based on NV centers in diamond was demonstrated making use of a NV-NV cross-relaxation feature between 0 and 20 mT. In this paper, we report imaging based on a much narrower cross-relaxation feature between the NV centers and P1 centers and apply it to perform non-destructive evaluation experiments on small solid-state batteries.

The imaged sample is an all-ceramic multilayer solid-state battery produced by TDK Corporation. It incorporates inner electrodes, electrolyte and external electrodes. When the battery is placed in an oscillating magnetic field (primary field), eddy currents flow in the electrodes and the electrolyte, which, in turn in our case induce an alternating magnetic field (secondary field).

We demonstrate all-optical AC magnetometry to detect the secondary field as a function of position. The secondary field is anti-parallel to the primary field and exhibits a phase delay. These quantities relate to a number of characteristics in the imaged sample, such as: shape, dimensions, conductivity and susceptibility. As a result, one can use NV centers as probes to distinguish textures and identify structural anomalies in- and outside the battery.

II. EXPERIMENTAL APPARATUS

A schematic of the experimental setup is shown in Fig. 1. Linearly polarized light from a continuous-wave green laser at 532 nm is continuously pumping the NV centers in a diamond sample. The laser power can be adjusted with a polarizer and a half-wave-plate. NV centers in the excited state decay back to the ground state both by radiative and nonradiative transitions, with red photoluminescence (PL) emitted in the former case. The fact that NV centers in the

 $m_s = \pm 1$ magnetic sublevels have significantly higher probability to undergo an intersystem crossing (ISC) ¹², i.e a nonradiative transition from the triplet to the singlet excited state, results in the majority of NV centers being optically pumped into the $m_s = 0$ spin projection. In order to minimize the noise caused by intensity fluctuations in the laser beam, we use a light-power stabilization loop incorporating a photodiode (PD2), a proportional-integral-differential (PID) controller, and an acousto-optic modulator (AOM).

A dichroic mirror and a notch filter are used to filter out red PL from the green light. The diamond used in this setup is a type-Ib, (111)-cut, high-pressure-high-temperature-grown sample with dimension $3.0 \times 3.0 \times 0.4 \,\mathrm{mm}^3$. The initial nitrogen concentration was specified less than 110 ppm. Then nitrogen-rich sample was electron-irradiated and annealed ¹¹. The diamond is glued on the top plane of a parabolic concentrator used to collect fluorescence. A radio-frequency (RF) coil is made from copper wire of 0.1 mm diameter. It has five turns and is wound around the diamond to provide a modulation field to both the diamond and battery. The modulation field is provided by the Lock-in amplifier (LIA) used for the AC magnetometric scheme. The frequency of the field (maximum 4 MHz) is also used as the LIA reference for detection.

A 3D translation stage is used to raster scan the battery in front of the diamond. To reduce the effect of magnetic noise produced by the 3D translation stage. The battery is glued on top of a 24 cm rod made from polyvinyl chloride (PVC) onto a motorized three-dimension-translation-stage controlled by a computer. A variable background field is provided by a custom-made electromagnet (EM). This field is required for the microwave-free magnetic field detection method. The EM consists of approximately 200 turns wound with a rectangular cross-section $(1.4 \times 0.8 \,\mathrm{mm}^2)$ wire around a 5 cm-diameter bore. The coil is wound on a water-cooled copper mount, and produces a background field of 2.9 mT per ampere supplied. Diamond, battery and RF coil are all placed inside the bore of the EM. A LIA detects the amplitude (R) and phase (θ) of PL modulation. The LIA is connected to the computer, and R and θ are recorded along with position of the battery (the 3D translation stage).

III. AC MAGNETOMETRY

The PL of a diamond sample displays various cross relaxation features, extensively discussed in the literature ^{13–15}, as a function of the bias magnetic field, as shown in Fig. 2. In this letter, we exploit a cross-relaxation feature between NV and P1 centers occurring at 51.2 mT [area B

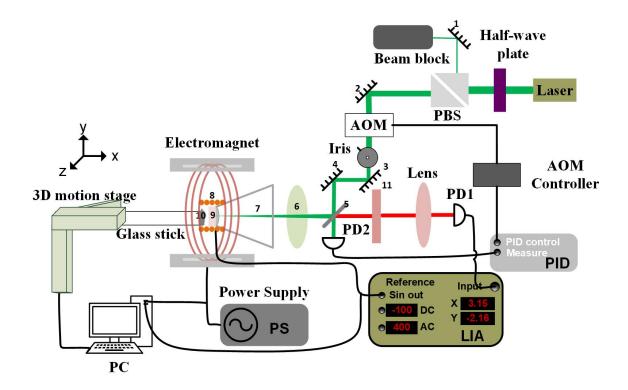


FIG. 1: Experimental setup. 1-4, mirror 5, dichroic mirror 6, lens 7, Parabolic concentrator 8, RF coil 9, diamond 10, battery 11, Notch filter. PBS: polarizing beamsplitter; AOM: acousto-optic modulator; PD1,2: photodiodes; PC: personal computer; PS: power supply.

in Fig. 2 (a)]. This feature is used in the battery measurement because of its improved sensitivity compared to the broad NV-NV cross-relaxation feature near zero field (area A), and of its reduced dependence on magnetic field compared to cross relaxation with NV centers not aligned along the magnetic field at 60 mT (area C) and to the ground state level-anticrossing (GSLAC) feature at $102.4 \, \text{mT}$ (area D). Figure 2 (b) shows the corresponding signal amplitude R detected with the LIA, demodulated at the referenced frequency. To maximize the response to alternating signals, R has to be maximized. The NV features seen in Fig. 2 are expected to display a temperature dependence. For example, the NV-P1 cross-relaxation features are expected to shift by $-1.34 \, \mu \text{T/K}$ at room temperature 16 , which can lead to drifts and be interpreted as magnetic field noise. As noted in the inset (d), we choose to use the features at $52.5 \, \text{mT}$ to detect the AC magnetic signals. The noise sensitivity of the sensor for the parameters used in this experiment leads to an estimated sensitivity of $40 \, \text{nT} / \sqrt{\text{Hz}}$ with a bandwidth of $100 \, \text{kHz}$. It should be pointed out that the bandwidth of this technique can be expanded up to MHz, depending on pump-power intensity and NV-axis alignment with respect to the external magnetic field 11 .

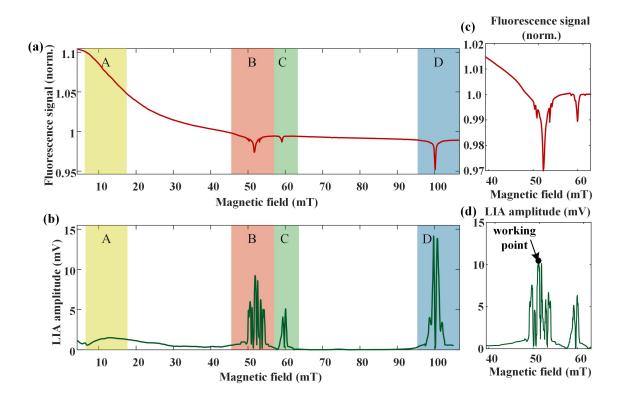


FIG. 2: (a). PL of NV centers as a function of the background magnetic field. The four areas A, B, C, D that can be exploited for magnetometry indicate features due to spin mixing, cross relaxation between NV centers and P1 centers, cross relaxation with NV centers that are not aligned along the magnetic field, and ground state level anticrossing, respectively. (b) Corresponding amplitude of LIA referenced at 100 Hz for those features. (c) Zoomed-in B area of (a), PL recorded with field-modulation depth of 0.48 mT at 100 Hz. (d) Zoomed-in B area of (b), indicating the working point at 52.5 mT.

IV. BATTERY MEASUREMENT

The dimensions of the solid state battery measured in this work is $4.0 \times 3.0 \times 1.0 \,\mathrm{mm^3}$, and are shown in Fig. 3 (a)¹⁷. It comprises Li-based-ceramic-oxide inner electrodes (noted as 2), Li-based-ceramic-oxide electrolyte (denoted as 3), copper charge collector (indicated as 1), and external electrodes made of aluminum. Figure 3 (b) shows a photograph of the solid battery, which has an artificial 1 mm-length defect at the left external electrode. The battery sample is placed at distance of 0.1 mm to the diamond at the perpendicular direction (*x*-direction), and scanned over the transverse plane (*y*-*z* plane). Figure 3 (g)-(j) show the spatially resolved phase signals of the LIA at the corresponding modulation frequencies of 1 kHz, 5 kHz, 10 kHz, 40 kHz, respectively.

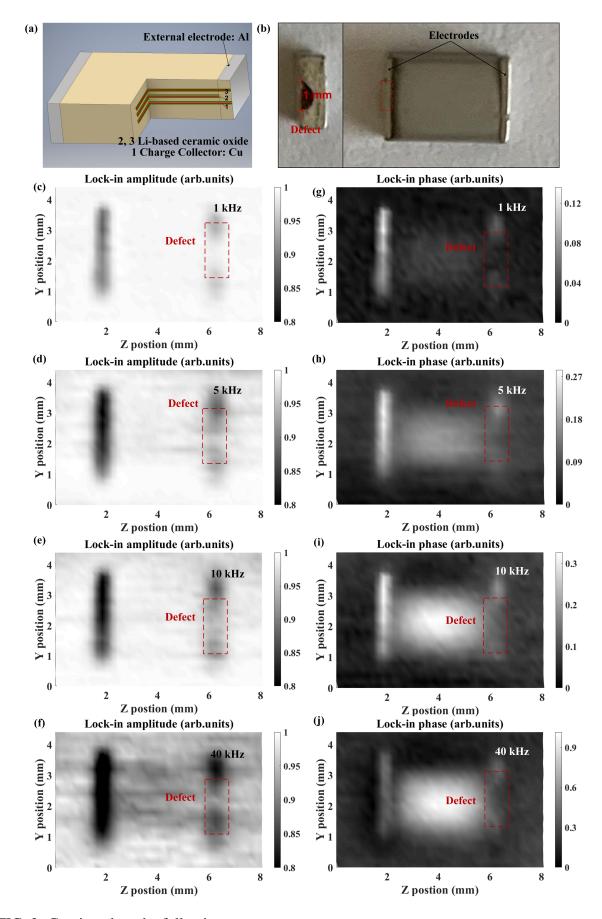


FIG. 3: Continued on the following page.

FIG. 3: (a) Battery schematic. The ceramic solid battery consists of external aluminium electrodes, Li-based ceramic oxide inner electrodes (noted as 2), Li-based ceramic oxide electrolyte (noted as 3), and copper charge collector (indicated as 1). (b) A photograph of the battery with dimension $4 \times 3 \times 1 \text{ mm}^3$. Distinctly visible is the defect on the external electrode. (c)-(f) Maps of the *R* of LIA for modulation frequencies of 1 kHz, 5 kHz, 10 kHz and 40 kHz. (g)-(j) Maps of the θ of LIA for modulation frequencies of 1 kHz, 5 kHz, 10 kHz and 40 kHz, respectively.

The external electrodes including the defect are distinctly visible under the modulation frequency of 1 kHz as shown in Fig. 3 (g). As the modulation frequency is increasing, the phase θ shows an image of the battery structure between the electrodes. There the induced field from different portions of the battery can be differentiated, as shown in Fig. 3 (h)-(j). Figure 3 (c)-(f) show the spatially resolved amplitude signals of LIA at the corresponding modulation frequencies of 1 kHz, $5 \, \text{kHz}$, $10 \, \text{kHz}$, $40 \, \text{kHz}$.

In order to investigate the ability of diamond-based magnetometer to identify structural anomalies inside batteries, we introduced two kinds of artificial anomalies. First, an iron-containing brass cylinder with a height of 1 mm and diameter of 1 mm is placed inside the battery, serving as an impurity. The results are shown in Fig. 4 (a)-(d), displaying the spatially resolved amplitude and phase signals of the LIA for the solid battery at frequencies of 5 kHz and 20 kHz respectively. This impurity is visible at both frequencies in the phase signal, indicated with red dashed circles in Fig 4. The impurity image is absent in the amplitude signal at frequencies lower than 5 kHz, this is because the amplitude variations caused by the impurity, are smaller than the background shifts caused by temperature changes, for the amplitude signal. The other anomaly is made by cutting the battery in half and then attaching the halves back together, leaving no defect visible to the naked eye. The results are shown in Fig 5, where (a)-(d) display spatially resolved amplitude and phase signals of LIA for the solid battery at frequencies of 5 kHz and 20 kHz respectively. The crevice is distinctly visible in the middle of the battery in the phase signals at both frequencies, indicated by the yellow ellipses in (b) and (d). In the amplitude signals the crevice is absent again.

The primary field produced at the center of the coil where the battery is located is given by:

$$B_{\text{primary}} = \frac{\mu_0 \times N \times I}{2r_{\text{coil}}},\tag{1}$$

where $\mu_0 = 4\pi \times 10^{-7} \,\text{T} \times \text{m/A}$ is the permeability of free space, N = 4 is the number of turns of the RF coil, $I = 0.8 \,\text{A}$ is the amplitude of the alternating current through the coil, and $r_{\text{coil}} = 4.2 \,\text{mm}$

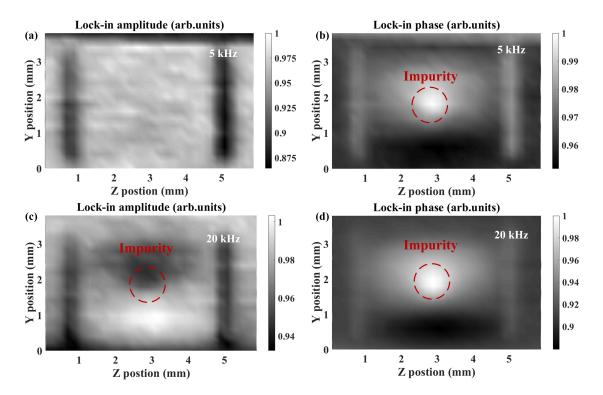


FIG. 4: Imaging pictures of solid batteries with the impurity. (a) and (b), (c) and (d) are amplitudes and phases of LIA for the battery with an impurity at frequencies of 5 kHz and 20 kHz respectively. (e) and (g), (f) and (h) are amplitudes and phases of LIA for the battery with a crevice in the middle at frequencies of 5 kHz and 20 kHz respectively.

is the radius of the coil. To perform the LIA detection scheme the modulated field amplitude is selected to be 0.48 mT, confirmed experimentally by the GSLAC PL scan and theoretically by eq.1.

For AC measurements, the magnetic susceptibility χ_m yields two quantities ^{18,19}: the magnitude and phase of the susceptibility. Alternatively, we can express it in the form of complex susceptibility with the real component χ_r and imaginary component χ_i ,

$$\chi_m = \chi_r + i\chi_i \tag{2}$$

where χ_r is the slope of DC magnetization curve M(H) of the battery, and χ_i indicates dissipative processes in the battery. The secondary field produced by the battery also depends on the skin depth in the conducting material:

$$\delta = \sqrt{\frac{2}{\sigma\omega\mu}},\tag{3}$$

where $\omega = 2\pi f$ denotes the angular frequency, σ is conductivity of the material and μ is mutual

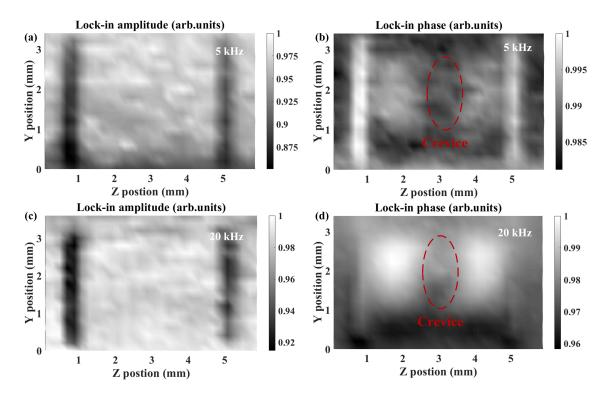


FIG. 5: Imaging pictures of solid batteries with the crevice. (a) and (b), (c) and (d) are amplitudes and phases of LIA for the battery with a crevice at frequencies of 5 kHz and 20 kHz respectively.

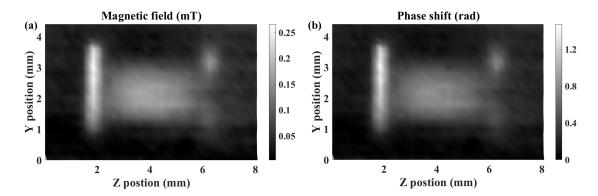


FIG. 6: Magnetic field and phase shifts generated by the battery at frequency of 5 kHz. (a) Calibrated magnetic field generated by the battery. (b) Phase shifts generated by the battery.

permeability of the material. In order to estimate the AC magnetic field generated by the battery, we adopt the calibration formula in ref⁸,

$$\delta\phi = \arctan\left(\frac{R_0\cos(\phi_0) - R_m\cos(\phi_m)}{R_m\sin(\phi_m)}\right), \tag{4}$$

$$\delta R = \sqrt{R_m^2\sin(\phi_m)^2 + (R_0\cos(\phi_0) - R_m\cos(\phi_m))^2}, \tag{5}$$

$$\delta R = \sqrt{R_m^2 \sin(\phi_m)^2 + (R_0 \cos(\phi_0) - R_m \cos(\phi_m))^2},$$
 (5)

where $\delta \phi$ is phase shifts generated by the battery, δR is phase shifts generated by the battery,

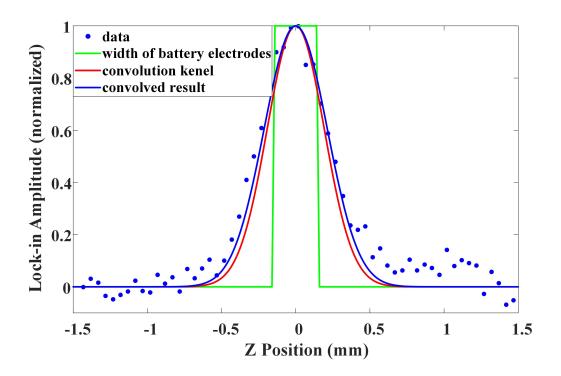


FIG. 7: Determination of spatial resolution. The blues dots signify the average cross section of the 0.3 mm-width electrode imaged by the sensor, the green trace denotes the width of electrode of battery as a square function, and the red curve represents a Gaussian function, as a convolution kernel of the green trace to fit the experimental data, noted as blue trace.

 R_0 and ϕ_0 are the background-corresponding amplitude and phase of the LIA, R_m and ϕ_m are the measured lock-in amplitude and the phase. The calibrated results are shown in Fig 6. We calibrated the magnetic field and phase shifts generated by the battery in Fig 3 at modulation frequency of 5 kHz. The external electrode without the defect induces the maximum magnetic field of 0.25 mT and phase shift of 1.3 rad.

V. SPATIAL RESOLUTION

One of the advantages of diamond-based sensors is the spatial resolution. For determining the spatial resolution of this battery-measurement system, we use data from one of the aluminium-made electrodes of the solid battery. As shown in Fig. 7, the average cross section of 0.3 mm-width electrode is expressed by a square function, noted as the green plot, a Gaussian-function, noted with red is employed as a kernel to convolve the square function to fit the measured data, denoted as blue dots. All traces have been normalized for better comparison. By adjusting the full width

at half-maximum of Gaussian function, we can estimate the spatial resolution of the system. We represent that results with blue. The spatial resolution of the system can be estimated form the full width at half maximum (FWHM) of the kernel and is estimated as $360 \pm 2 \,\mu$ m. It is mostly restrained by the distance between the battery and diamond. Closer proximity could increase measurement contrast, allowing more legible image of small samples.

VI. CONCLUSION AND OUTLOOK

We demonstrate microwave-free eddy-current imaging with NV centers in diamond exploiting the cross-relaxation feature between the NV centers and P1 centers occurring at 51.2 mT. The magnetic sensitivity is $40\,\mathrm{nT}/\sqrt{\mathrm{Hz}}$ with 100 kHz bandwidth, but as previously demonstrated in literature 11 it can be extended to a few MHz. The diamond-based sensor is used to noninvasively image a mm-sized solid battery. The spatially resolved amplitude and phase lock-in signals can be used to determine the dimension and shape of the battery features, including a 1 mm-width artificial defect and 0.3 mm-width electrodes, from which the spatial resolution is calibrated as $360(2)\,\mu\mathrm{m}$ limited by the distance to the battery. Furthermore, diamond-based probe also identifies an anomaly and a crevice inside the battery, demonstrating the ability to detect internal anomalies which are not visible to the naked eye. This paper mainly focuses on testing solid batteries. For future studies, the spatial resolution could be further improved by either using a thinner diamond sample or a diamond with a shallow implanted NV layer to shorten the sensor-sample distance. By tuning the bandwidth, it would be possible to identify structures, at different depth inside the battery including electrodes, electrolytes 20 and even material components, which could be useful for battery assessment and development.

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