Diapiric relamination of the Orocopia Schist during low-angle

2 subduction

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Abstract

The Orocopia and related schists are sediments subducted during the Laramide orogeny and are thought to have been underplated as a laterally extensive layer at the base of the crust in the southwest U.S. Cordillera. This concept is hard to reconcile with the existence of continental mantle lithosphere in southeast California and western Arizona. Analytical solutions and numerical modeling suggest that the Orocopia Schist may have ascended through the mantle lithosphere as sediment diapirs or sub-solidus crustal plumes to become emplaced in the middle to lower crust. Modeled time-temperature cooling paths are consistent with the exhumation history of the Orocopia Schist and explain an initial period of rapid cooling shortly after peak metamorphism. The Orocopia Schist represents a potential example of relaminated sediment observable at the surface.

Introduction

Shallow to flat-slab subduction is one of the primary causes of large-scale sediment subduction during which trench, accretionary complex, and forearc basin material is tectonically eroded and transported beneath the upper plate (Ducea and Chapman, 2018). The fate of this subducted crustal material has global implications for the composition and evolution of the

lithosphere and may be recycled into the mantle, underplated, or relaminated (Clift and Vannucchi, 2004; Behn et al., 2011; Hacker et al., 2011; Kelemen and Behn, 2016). In the southern U.S. Cordillera, shallow subduction during the Laramide orogeny resulted in the sublithospheric emplacement of the Pelona, Rand, Orocopia, San Emigdio, and Sierra de Salinas schists that were subsequently exhumed to the surface during the Late Cenozoic (Chapman, 2017). Current models propose that these schists were emplaced as a laterally extensive sheet beneath the crust (e.g., Jacobson et al., 2007). These models are supported by outcrops near the paleo-plate margin where the schists are demonstrably underplated directly to the base of the crust and contributed to extinguishing the Late Cretaceous Sierra Nevada continental arc (Kidder and Ducea, 2006). Farther inland, however, the emplacement mechanism of these schists is less certain. Recent discoveries of Orocopia Schist in western Arizona (Haxel et al., 2015; Jacobson et al., 2017; Strickland et al., 2018) expanded the landward limit for known schist emplacement (Fig. 1). Models suggesting that a sheet of Orocopia Schist was underplated to the crust imply that mantle lithosphere was removed during low-angle subduction. However, mantle xenolith, isotopic, and geophysical studies indicate the mantle lithosphere was not fully removed during the Laramide orogeny in the southwest U.S. (Livacarri and Perry, 1993; DePaolo and Daley, 2000; Allison et al., 2013), and structural restorations and thermochronologic data indicate that the Orocopia Schist was emplaced in the lower to middle crust shortly after peak metamorphism, but prior to Late Cenozoic exhumation (Jacobson et al., 2007; Spencer et al., 2018; Strickland et al., 2018). These observations beg the question: After the Orocopia Schist was subducted, how did it come to reside in the crust if the mantle lithosphere and possibly lower crust was in its way?

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Emplacement conditions of the Orocopia Schist

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The events associated with the emplacement of the Orocopia and related schists resulted in the dismemberment and destruction of the plate margin, including nearly complete removal of a massive accretionary complex and forearc basin (Jacobson et al., 2011). These schists are estimated to have subducted as a 10 ± 5 km thick layer based on outcrop exposures, structural reconstructions, and low-velocity anomalies in the middle to lower crust (Porter et al., 2011; Xia and Platt, 2017; Spencer et al., 2018). This detail alone distinguishes the Orocopia Schist from most other documented instances of sediment subduction that have subducted layer thicknesses of ≤ 1 km (Clift and Vannucchi, 2004). Layer thickness is a critical length scale that determines how rapidly and at what rheological conditions sediment detaches from the down-going slab. Sediment diapirs are proposed to form for any sediment layer thicker than ~300 m, at temperatures as low as 500 °C (Currie et al., 2007; Behn et al., 2011; Miller and Behn, 2012). Peak metamorphic conditions recorded by the Orocopia Schist are 8-14 kbar and 550-650 °C (Chapman, 2017). New maximum temperature estimates using Raman spectroscopy of carbonaceous material in the Orocopia Schist from recently discovered locations in western Arizona are 599-649 °C at Cemetery Ridge and 546-583 °C at the Plomosa Mountains. (Supplementary File 1). Phase modeling of subducted metasedimentary rocks at these P-T conditions (Behn et al., 2011) indicates a density contrast of ~550 kg/m³ between the Orocopia Schist and the continental mantle lithosphere and a density contrast of up to 200 kg/m³ with the lower crust. Density contrasts as low as 100 kg/m³ have been linked to instability growth (Currie et al., 2007; Behn et al., 2011). Viscosity also influences the formation of density instabilities and is largely dependent on temperature and strain rate. Strain rates as high 10⁻¹³ s⁻¹ have been reported from the Pelona

Schist based on deformed quartz fabrics (Xia and Platt, 2017). Strain rate in the continental mantle lithosphere near the slab interface in the region of flat-slab subduction has been estimated at 10^{-15} to 10^{-16} s⁻¹ based on the results of numerical modeling (Liu and Currie, 2019). At these strain rates and temperatures, the wet quartzite flow law of Hirth et al. (2001) suggests that the Orocopia Schist had an effective viscosity of 10^{19} to 10^{21} Pa·s. The lower end of this range is comparable to models for lower crustal flow (e.g., Beaumont et al., 2001). Post-seismic relaxation and xenolith studies from the area underlain by the Orocopia Schist indicate that modern mantle lithosphere has an effective viscosity of 10¹⁹ to 10²⁰ Pa·s (Behr and Hirth, 2014). Laramide-age, upper mantle xenoliths from the stable Colorado Plateau indicate effective viscosities of 10¹⁹ to 10²³ Pa·s at 550-750°C and an assumed pressure of 2 GPa (Behr and Smith, 2016). Analytical solutions for Rayleigh-Taylor instabilities in two-layer systems (Hess and Parmentier, 1995) applied to the Orocopia Schist and the mantle lithosphere at the conditions described produce characteristic instability growth times of 0.5-5 Myr, diapir diameters of 30-80 km, and maximum ascent velocities of 0.1-2 km/Myr, assuming a 10 km thick layer of schist (Fig. 2). Reconstructed time-temperature histories for the Orocopia Schist indicate that

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Reconstructed time-temperature histories for the Orocopia Schist indicate that subduction, prograde to peak metamorphism, and initial cooling all occurred in <15 Myr (Grove et al., 2003; Jacobson et al., 2007; 2011) (Fig. 3). The period of initial cooling (Fig. 3) is enigmatic and has been variably attributed to erosional exhumation, thrust faulting, extensional faulting, isostacy, and subduction channel return flow (Chapman, 2017). Rapid initial cooling to 300-400 °C supports emplacement of the Orocopia Schist in the lower to middle crust where it resided for 20-50 Myr before being exhumed to the surface during the Miocene (Fig. 3). The most extreme example of crustal emplacement comes from the Plomosa Mountains where the

Orocopia Schist may have been located as shallow as 5 km prior to exhumation (Spencer et al., 2018; Strickland et al., 2018). I propose that the period of rapid initial cooling can be explained by diapiric ascent of the Orocopia Schist from the Farallon slab to the middle crust.

Numerical experiments and model set up

To explore the possibility of relamination, I consider two end-member numerical models involving density instabilities that can place the Orocopia Schist above the mantle lithosphere (Fig. 4). Model 1 is inspired by studies that suggest isotherms are nearly parallel to the plate interface in flat-slab segments (e.g., Liu and Currie, 2019) and supposes that several discrete instabilities formed nearly simultaneously (Chapman, 2017). Model 2 supposes that instability growth is a direct function of temperature/depth and that the Orocopia Schist detached from the Farallon slab at a point-like source in a fixed upper plate reference frame. This model is broadly analogous to the "extrusion" model that Strickland et al. (2018) proposed for the Plomosa Mountains or the more general "crustal plume" model of Currie et al. (2007). Model 2 introduces Orocopia Schist-like material into the bottom left of the model at a rate of 10 cm/yr, consistent with conservative estimates of North America-Farallon plate convergence rate during the Laramide orogeny (Torsvik et al., 2008).

The numerical experiments were conducted using the finite-element software COMSOL Multiphysics. Fig. 4B describes the geometry and initial conditions for a reference model based on the thickness, density, viscosity, temperature, and pressure constraints outlined above. The thickness of the mantle lithosphere is constrained by maximum pressure estimates from metamorphic studies (Chapman, 2017 and references therein). The initial geotherm was scaled to match the thermal structure of numerical models of flat-slab subduction at 50-110 km depth that

suggest a slab-mantle interface of ca. 600 °C (Liu and Currie, 2019), consistent with peak metamorphic temperatures. Internal heat generation, strain heating, and adabatic heating were neglected. The density and viscosity of layers representing the crust and Orocopia Schist were coupled to the thermal model and varied as a function of temperature using the wet quartz flow law of Hirth et al. (2001). Model details are presented in Supp. File 1.

Generation and characteristics of sediment diapirs

The numerical modeling results predict diapiric rise of the Orocopia Schist at the conditions considered. Figures 4C and 4D show model results at intermediate time steps using the starting conditions described in the reference model. Both Model 1 and Model 2 produce large-scale diapir-like features that become emplaced in the lower to middle crust (as shallow as 20 km depth) in < 10 Myr. Model 1 generates classic Rayleigh-Taylor instabilities that reach a maximum diameter of ca. 65 km after emplacement (Fig. 4C). The lower crust is thickened between the instabilities. Model 2 produces a plume-like structure that rises rapidly into the middle crust and then starts to grow horizontally for as long as Orocopia Schist is input into the model. The position of the plume at the end of the model run is shown in Fig. 4D. The "crustal plume" exhibits only minor vertical thickening during its lateral growth. The lower crust is thickened in front of the growing plume.

Tens of model trials were run for Model 1, deviating from the reference model to explore variations in subducted sediment thickness and viscosity, which is affected primarily by temperature and strain rate. Sedimentary diapirs continued to be produced with a layer thickness of 1 km, the smallest thickness considered. Increasing the maximum mantle lithosphere viscosity to 10^{22} Pa·s delayed instability development (ca. 40 Myr after the start of a model run) and

increasing the mantle lithosphere viscosity to 10^{23} Pa·s failed to produce instabilities in a 100 Myr-long model run. Increasing the maximum viscosity of the Orocopia Schist layer to 10^{21} Pa·s delayed instability formation by 2-3 Myr, but otherwise had little effect on diapir generation. This analysis suggests that mantle lithosphere viscosity has the largest effect on the development of sediment diapirs and represents the greatest uncertainty within the model parameters considered.

Thermal models for subduction zones suggest that mantle temperatures initially increase upward from the slab, reaching a peak in the core of the mantle wedge, which may facilitate the formation of diapirs (e.g., Jull and Kelemen, 2001; Miller and Behn, 2012). During flat slab subduction, where no asthenospheric mantle is present between the slab and the lithosphere, temperature inversions within the lithospheric mantle are predicted to be subdued (e.g., Liu and Currie, 2019), but may still enhance diapirism. The addition of a "hot core" to the mantle lithosphere in Model 1 (~150°C above the ambient geotherm) was investigated using a wet olivine flow law for viscosity (Hirth and Kohlstedt, 2003), but did not produce significantly different results compared to model runs without the "hot core." The relatively thin mantle lithosphere in Model 1 (20-30 km) limited viscous decay caused by temperature gradients.

The temperature of the sediment diapirs was tracked to compare the model results with observed time-temperature histories of Orocopia Schist (Fig. 3). Model 1 indicates rapid cooling (ca. 40 °C/Myr) during initial ascent of the instability and then a step-like shift to slow cooling (3-5 °C/Myr) for the remainder of the model run as the diapir slowly equilibrates to the ambient geotherm. The shape of the cooling path and the final temperature of the modeled sediment diapir in Model 1 are consistent with the observed time-temperature paths (Fig. 3). Model 2 also

indicates initial rapid cooling (ca. 35 °C/Myr) during diapir ascent, but the cooling rate gradually slows through time until it reaches a similar cooling rate to Model 1 (3-5 °C/Myr).

Discussion and Conclusions

The Orocopia Schist, outcropping at the surface of the Earth and located up to 450 km from the paleo-trench (present day coordinates), is one of the most unambiguous examples of subducted sediment being re-incorporated into the crust. The thickness, density, and rheology of the Orocopia Schist suggest that it was susceptible to the formation of sediment diapirs if subducted beneath the lithospheric mantle.

However, the possibility of relatively high viscosities ($\geq 10^{23}$ Pa·s) in the upper mantle may limit the applicability of diapir models. The wet olivine flow law of Hirth and Kohlstedt (2003) predicts upper mantle viscosities of 10^{23} to 10^{25} Pa·s at the strain rate and temperatures considered – higher than required by the modeling results ($< 10^{22}$ Pa·s). The modeled values overlap with viscosity estimates determined from Laramide-age mantle xenoliths (10^{19} to 10^{23} Pa·s), but these xenoliths came from the deepest, most viscous, parts of the lithosphere (Behr and Smith, 2016). Lithospheric mantle xenoliths from southern California have diverse mineralogies and exhibit evidence for metasomatic alteration (Luffi et al., 2009; Chin et al., 2104), which suggests that viscosities calculated assuming a homogenous olivine composition are maximum estimates. Another possibility is that pre-existing structural heterogeneities may have focused buoyant ascent of the Orocopia Schist.

Analytical solutions for Rayleigh-Taylor instabilities in two-layer systems (Fig. 2) and numerical models (Fig. 4) suggest that Orocopia Schist would have detached from the Farallon slab as diapirs or sub-solidus crustal plumes and ascended through the mantle lithosphere and

into the lower crust - a form of relamination. Time-temperature cooling paths produced from the numerical models are consistent with the observed exhumation history of the Orocopia Schist and may explain a period of rapid initial cooling (Fig. 3). Diapirism is consistent with mechanical mixing of mantle rocks (peridotite, actinolite/tremolite pods) and schist observed at Cemetery Ridge and the Plomosa Mountains (Haxel et al., 2015; Strickland et al., 2018).

The results suggest that conceptual models for a laterally extensive sheet of Orocopia Schist lying below the base of the crust throughout the southwest U.S. may be unrealistic, particularly where mantle lithosphere is still present (e.g., Allison et al., 2013). The emplacement mechanisms proposed here predict that Orocopia Schist should be non-uniform and localized in the subsurface. The counterpart to rising diapirs of Orocopia Schist are descending "drips" of mantle lithosphere that may have been prone to delamination after the removal of the Farallon plate. This could help to explain a long-standing issue related to the discontinuous and patchwork-like preservation of mantle lithosphere in the southwest U.S. Cordillera (e.g., Livacarri and Perry, 1993). Finally, the modeling predicts crustal thickening roughly equivalent to the thickness of subducted sediment layer. The emplacement of the Orocopia Schist may help explain significant crustal thickening during the Laramide orogeny in southern Arizona, which cannot be explained by horizontal shortening alone (Chapman et al., 2020).

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viscosity, μ_1 . The viscosity of the upper layer, μ_2 (mantle lithosphere) is 10^{21} Pa·s and the density 320 difference, $\Delta \rho$, of the two layers is 550 kg/m³. 321 322 Fig. 3. Comparison of the observed cooling history of the Orocopia Schist (Grove et al., 2003; 323 Jacobson et al., 2007; 2011) with time-temperature paths generated by numerical modeling of 324 sediment diapirism. The models are consistent with a period of rapid initial cooling from peak 325 metamorphic conditions (ca. 600°C) to 300-400°C, which is interpreted to represent lower to 326 middle crustal depths. 327 328 Fig. 4. A) Schematic cross-section showing subduction of Orocopia Schist beneath the mantle 329 lithosphere. B) Reference model geometry showing temperature, density, and viscosity 330 variations with depth. C) Results of Model 1 at t=8 Myr showing diapir formation. D) Results of 331 Model 2 at t=10 Myr, showing a "crustal plume." Heavy dashed white lines show the position of 332 "Orocopia Schist" at subsequent time steps. Thinner, dashed white lines are schematic flow lines. 333