

## Connection between zero chromaticity and long in-plane polarization lifetime in a magnetic storage ring

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In this paper, we demonstrate the connection between a magnetic storage ring with additional sextupole fields set so that the  $x$  and  $y$  chromaticities vanish and the maximizing of the lifetime of in-plane polarization (IPP) for a 0.97-GeV/ $c$  deuteron beam. The IPP magnitude was measured by continuously monitoring the down-up scattering asymmetry (sensitive to sideways polarization) in an in-beam, carbon-target polarimeter and unfolding the precession of the IPP due to the magnetic anomaly of the deuteron. The optimum operating conditions for a long IPP lifetime were made by scanning the field of the storage ring sextupole magnet families while observing the rate of IPP loss during storage of the beam. The beam was bunched and electron cooled. The IPP losses appear to arise from the change of the orbit circumference, and consequently the particle speed and spin tune, due to the transverse betatron oscillations of individual particles in the beam. The effects of these changes are canceled by an appropriate sextupole field setting.

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## I. INTRODUCTION

A recent paper [1] reports the achievement of a 1000-second lifetime for the in-plane polarization (IPP) of a spin-polarized deuteron beam circulating at 0.97-GeV/ $c$  in the COSY storage ring [2]. This lifetime was made possible through a combination of beam bunching, electron cooling, trimming of sextupole field components, and limiting the beam current ( $< 10^9$  deuterons/fill). This demonstration fulfills a requirement [3], albeit at a lower circulating beam current, for the use of a storage ring to search for an intrinsic electric dipole moment (EDM) on the nuclei in a circulating charged-particle beam. This paper describes in more detail the connection noted in that letter between the setting of the COSY sextupole trim fields for long IPP lifetime and the vanishing of the  $x$  and  $y$  chromaticities in the ring. The COSY storage ring was operated at horizontal and vertical betatron tunes near 3.6 (but not equal to each other) and well away from machine resonances [4]. The features of an EDM storage ring experiment have been reviewed elsewhere [1,3,5].

The precession of the beam's IPP in a storage ring is governed by the interaction between the particles' magnetic anomaly and any vertical (perpendicular to the ring plane) magnetic field. The precession rate is given by the product of the beam revolution frequency  $f_{\text{rev}}$  and the spin tune  $\nu_S = G\gamma$ , where  $G$  is the deuteron's magnetic anomaly and  $\gamma$  is the relativistic factor (see Table I where  $C$  is the ring circumference). Any variation in the speeds among the beam particles creates a distribution in  $\gamma$  that leads to a spread in individual particle spin directions and eventual depolarization of the beam. With a coasting beam distributed around the ring, the leading first-order contribution to

$\nu_S$  variations comes from the momentum spread of the beam, which is typically in the vicinity of  $10^{-4}$  of the nominal momentum for the usual sorts of beam injection schemes. This leads to an IPP lifetime of tens of milliseconds for  $f_{\text{rev}} \simeq 10^6 \text{ s}^{-1}$ . This effect can be removed on average by bunching the beam, thus forcing long-term isochronicity on all of the beam particles. Variations in individual particle momenta still exist, but the rf cavity now traps all particles within the bunch where they oscillate in synchrotron orbits of different sizes. Since each particle's momentum now oscillates about the average value, depolarization is significantly reduced.

Second-order contributions to the spin tune spread may come from oscillations of the particles about the nominal orbit through either betatron (transverse) or synchrotron (longitudinal) motion. The transverse motion increases the length of the path that the particles must travel as they go around the ring. In combination with beam bunching, this forces those particles with large oscillation amplitudes to travel with a larger speed, thus increasing  $\gamma$  and generating depolarization.

If the oscillations in the longitudinal direction are of small amplitude compared to the ring circumference and the rf cavity is driven by a sinusoidal voltage, then the motion closely approximates a simple harmonic system. If the effect of the speed increase as the particle moves forward through the beam bunch is balanced by a speed decrease while going in the backward direction, then there is no net effect on average and these speed changes do not contribute to a growing depolarization. Typically, the frequency of this oscillation is about four orders of magnitude smaller than the revolution frequency, while betatron oscillations in a strong focusing system have frequencies that are larger than the revolution frequency by a factor of a few.

For large synchrotron oscillations at higher order, an additional quadratic dependence of  $\gamma$  on speed appears and this term may result in a shift of the average spin tune from the reference orbit value. The size of the spin tune shift depends on the second-order momentum compaction factor [7], a property of the ring lattice. This effect is hard to measure for the COSY lattice at the lower beam momenta relevant for this experiment due to limitations in the acceptance of the ring for momentum changes and the requirement to measure a quadratic term. This and other evidence suggest that these higher-order effects of synchrotron motion are small enough that we can neglect them for the studies described here.

The second-order effects from transverse orbit oscillations may be reduced by applying electron cooling to the beam, which reduces the spread of particle orbits in phase space (see Table I for typical cooled beam properties). Under these conditions, the IPP lifetime may reach seconds instead of tens of milliseconds.

A better way to manage second-order depolarization from transverse oscillations is to cancel the effect entirely

TABLE I. Machine and beam parameters.

Assumed
$c = 299792458 \text{ m/s}$
$m_d = 1875.612793(47) \text{ MeV}/c^2$
$G = -0.14298754(26)$
Measured
$f_{\text{rev}} = 750602.5(5) \text{ Hz}$
$f_{\text{res}} = 871431.29(10) \text{ Hz}$
Calculated
$p = 970.060(24) \text{ MeV}/c$
$C = 183.4827(37) \text{ m}$
$T = 236.007(11) \text{ MeV}$
$\gamma = 1.1258296(59)$
$\beta = 0.4593915(89)$
$\nu_S = G\gamma = -0.16097960(78)$
$f_{\text{rev}}\nu_S = -120831.29(59) \text{ Hz}$
Beam Properties (from Ref. [6])
$\Delta p/p = 4.91 \pm 0.13 \times 10^{-5}$
$\epsilon_X \text{ emittance} = 0.09 \text{ } \mu\text{m}$
$\epsilon_Y \text{ emittance} = 0.42 \text{ } \mu\text{m}$

rather than just reduce the beam's phase space size. Sextupole fields in conjunction with quadrupole focusing fields can shift the center point of the transverse oscillation pattern. The amount of the shift depends quadratically on the size of the transverse oscillation. Such adjustments made in the arcs of the COSY ring can result in shorter orbits when the shift is toward the center of the ring. In combination with the path lengthening created by the oscillations in the first place, a proper sextupole strength can match these two effects, making the modified orbits the same length as unperturbed orbits with no transverse oscillation. Since the isochronicity of the orbit is maintained by the rf cavity, all orbits will have the same average spin tune. The COSY ring was built with sextupole correction magnets located where there are large horizontal and vertical beta functions as well as a large horizontal dispersion. The available sextupole fields appear to be sufficient to create the needed cancellations. So the study of the beam properties undertaken in this paper explores how to measure and then correct the effects of depolarization in the presence of a distribution of transverse oscillation amplitudes.

Section II contains a brief summary of transverse oscillations and their connection to chromaticity. In Sec. III we turn to the explanation of the experiment, beginning with a review of the time-marking system used to measure the IPP [8] and a description of the context provided by the map of chromaticity values for the data-taking plan. The results are shown in Sec. IV along with the constraints on their interpretation. This section also compares the best IPP lifetimes with the chromaticity values. Conclusions are presented in Sec. V.

## II. TRANSVERSE OSCILLATIONS, CHROMATICITY, AND SEXTUPOLE FIELD ADJUSTMENTS

In a strong focusing machine, particles that are undergoing transverse oscillations will see a restoring force toward the reference orbit that is strong enough to generate a few oscillations for one revolution of the beam around the machine. The number of oscillations per revolution is known as the betatron tune,  $\nu$ . The horizontal ( $x$ ) and vertical ( $y$ ) focusing histories are different and not strongly coupled to each other. So it makes sense to define betatron tunes for both directions, either  $\nu_x$  or  $\nu_y$ . In general, these tunes depend on the particle momentum  $p$ . The linear coefficient in this relationship is called the chromaticity,  $\xi_{x,y} = \partial\nu_{x,y}/\partial(p/p)$ . Since an increase in the beam momentum reduces the time that is available for oscillations during one revolution, and since the restoring force from the quadrupoles remains roughly the same for all momenta, chromaticities in lattices without other alterations tend to be negative. This is called the natural chromaticity  $\xi_{\text{nat}}$ . This result may be altered, as we will

see later, by the addition of sextupole fields, some of which may arise in standard components such as the ring dipoles.

The natural chromaticity for one particle may be expressed in terms of the Courant-Snyder parametrization of betatron oscillations in a velocity vs position phase space with the symbols  $\alpha_{cs}$ ,  $\beta_{cs}$  and  $\gamma_{cs}$  (see Chap. 2 of Ref. [9]). In the case of vertical oscillations, for example, this leads to

$$4\pi\xi_{\text{nat}} = \oint K\beta_{cs}ds = \oint (\gamma_{cs} + \alpha'_{cs})ds = \oint \gamma_{cs}ds, \quad (1)$$

where  $K$  is the focusing function and  $ds$  is a differential element of the beam path around the ring. A common substitution,  $K\beta_{cs} = \gamma_{cs} + \alpha'_{cs}$ , is used. The derivative of  $\alpha_{cs}$  is with respect to  $s$  and its integral around the ring vanishes for one complete turn. For horizontal oscillations, Eq. (1) would begin with a minus ( $-$ ) sign.

Betatron oscillations increase the path length or circumference around the ring as

$$C = \oint \sqrt{1 + y'^2}ds \approx C_0 + \frac{1}{2} \oint y'^2ds, \quad (2)$$

where  $C_0$  is the circumference of the reference orbit and  $y'$  is the slope of the particle path with respect to  $s$ . The fractional change in the circumference is given by

$$\frac{\Delta C}{C_0} = \frac{1}{2C_0} \oint y'^2ds = \frac{J}{C_0} \oint \gamma_{cs}ds = 4\pi \frac{J}{C_0} \xi_{\text{nat}}, \quad (3)$$

where  $J$  is the action (see Chap. 2.1.1 of Ref. [10] for the substitution). The last substitution comes from Eq. (1). This result is similar to Ref. [11] for an ensemble of particles.

Minimizing the depolarization means that the path length changes must become small or zero for particles away from the reference orbit. This means that  $\Delta C$  must become zero. This can only be accomplished in Eq. (3) by including an additional contribution to the chromaticity from sextupole fields, as given by

$$\frac{\Delta C}{C_0} = 4\pi \frac{J}{C_0} \xi_{\text{nat}} - \oint \frac{J}{C_0} D\beta \frac{B''\ell}{(B\rho)} \delta_{MXi}(s)ds, \quad (4)$$

where  $D$  is the dispersion,  $B''$  is the curvature of the sextupole field,  $\ell$  is the sextupole magnet length,  $(B\rho)$  is the rigidity of the ring (see Chap. 2.1.1 of Ref. [10]), and the  $\delta$ -function  $\delta_{MXi}(s)$  locates the sextupole magnet belonging to one of the MXS, MXL, or MXG families. Making  $\Delta C = 0$  means that the total chromaticity vanishes.

It has been argued independently along these same lines that the simultaneous appearance of zero chromaticity and a long IPP lifetime should occur [12]. This property was used in the preparation of the electron and positron beams for a comparison of their magnetic moments [13,14].

The appearance of the action  $J$  in both terms of Eq. (4) means that the strength of the sextupole, adjusted by

changing  $B''$ , may be used to make the variations in the circumference zero for all values of  $J$  or oscillation sizes, at least to this order.

This concept was investigated using the particle tracking program COSY Infinity [15] with a ring lattice that was the model for the machine setup described here (no dispersion in the machine straight sections). Various orbits were calculated that involved horizontal or vertical betatron oscillations (in  $x$  and  $y$ ) about the reference orbit. A series of different betatron amplitudes were used. For each orbit the precession of the IPP was calculated for one turn. With no field corrections, the spin tune  $\nu_S$  varied quadratically with the betatron amplitude. Including sextupole fields from the three COSY sextupole families caused the spin tune values to vary for nonvanishing betatron amplitudes. This variation was linear in the sextupole strength. It was possible to obtain values that were above or below the spin tune of the reference orbit. At one particular value of the sextupole field, all of the spin tunes for a series of betatron amplitudes (either  $x$  or  $y$ ) became equal, indicating that a cancellation of this depolarizing effect is possible. Since all sextupole families affect all betatron amplitudes, we can write a more general condition for minimal spin tune spread as follows

$$0 = |A + a_1 I_S + a_2 I_L + a_3 I_G| \langle x'^2 \rangle + |B + b_1 I_S + b_2 I_L + b_3 I_G| \langle y'^2 \rangle + O(\langle \Delta p/p \rangle^2), \quad (5)$$

where  $A$  and  $B$  represent the contributions to the natural chromaticity, or the chromaticity with no adjustable sextupole strength applied. The sextupole currents,  $I_S$ ,  $I_L$ , and  $I_G$  are the currents applied to the three families of sextupoles located in the arcs of the COSY ring. These families are located at places of large Courant-Snyder  $\beta_{x,cs}$ ,  $\beta_{y,cs}$  and dispersion  $D$  respectively, as shown in Fig. 1. The subscripts  $S$ ,  $L$ , and  $G$  refer to designations of the sextupole families (MXS, MXL, and MXG) used in the magnet-naming scheme for COSY. These beam properties define the sensitivity to the sextupole strength described by the  $a_i$  and  $b_i$  coefficients. But the currents are the same in both of the major terms of Eq. (5), and this means that any solution must be found for both terms together. It may also be the case that there is an additional dependence, but through synchrotron oscillations rather than betatron, that adds an additional term that depends on  $\langle (\Delta p/p)^2 \rangle$ . Such a term arises only as the synchrotron motion departs from simple harmonic, as it eventually must for large oscillations in the sinusoidal potential provided by the ring rf bunching cavity.

The requirement as sextupole fields are adjusted is that the IPP lifetime improves. In practical usage, we often tracked the reciprocal of the IPP lifetime and looked for a setting where the reciprocal went to zero (or close to zero). This will come up again in Sec. IV, where lifetime measurements are explained in more detail. In an independent check, model calculations of betatron-induced

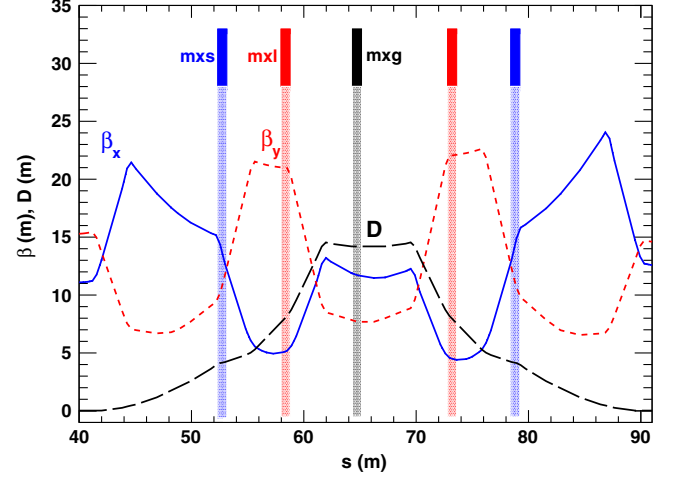


FIG. 1. Variation of the beta functions  $\beta_x$  and  $\beta_y$  along with the dispersion  $D$  as a function of the distance  $s$  along the reference orbit through the first arc of the COSY storage ring. The opposite arc is similar. The locations of the MXS, MXL, and MXG sextupole magnets are indicated with different colors. All of the magnets in each of these three families (in both arcs) are wired in series and controlled with a single setting of the current.

depolarization made using the no-lattice spin transport model of Benati [6] showed that the reciprocal of the polarization lifetime varies linearly with the size of the betatron oscillations as represented by the square of the transverse velocity or particle path slope shown in Eq. (1). Thus the behavior of the reciprocal near zero chromaticity should be linear with sextupole strength.

These arguments provided a framework to search for long IPP lifetimes with the assumption that the best conditions under which to find the optimum lifetime would occur when  $\xi_x = 0$  and  $\xi_y = 0$  at the same time. This may be understood in a simple way by the statement that all particle paths have the same length regardless of the magnitude of any transverse (betatron) oscillation that may be present.

### III. EVENT TIME MARKING, CHROMATICITY MEASUREMENT, AND RUNNING PLAN

A summary of the experimental method was presented by Guidoboni [1]. The beam polarization was continuously monitored using the scattering of beam deuterons from a thick carbon target into a set of polarimeter scintillators employed previously by the EDDA experiment [16–18]. Each polarimeter event was marked with a clock time from which it was possible to unfold the anomalous precession based on the total time that had elapsed since the beginning of data acquisition on a particular beam store and an estimate of the spin tune [8]. From this, we assigned an IPP direction or phase to each event. In each of a series of time bins during the store, the polarimeter events were sorted by direction and the asymmetry computed for each direction



bin from a comparison of the down and up counting rates assigned to that bin. A fit of a sinusoidal curve to these asymmetries then gave the magnitude of the IPP and the direction, or phase, of the polarization at the beginning of the time bin.

In this analysis, the spin tune must be accurate to about one part in  $10^6$  or else the accumulated difference between the calculated and actual phases in even a one-second time bin is enough to reduce the apparent IPP magnitude. To remedy this, the spin tune in the analysis was scanned over a narrow range near the starting value and a search was conducted for the spin tune with the maximum IPP. This allowed each time interval to have its own spin tune, but in practice this led to an unstable series of phase values. In a second pass through the data, the spin tune was fixed at an average value among all the time bins (sometimes with a smooth trend added) and the analysis repeated. With this constraint, the new set of phases had a smooth time dependence so long as the IPP magnitude was not close to zero [8].

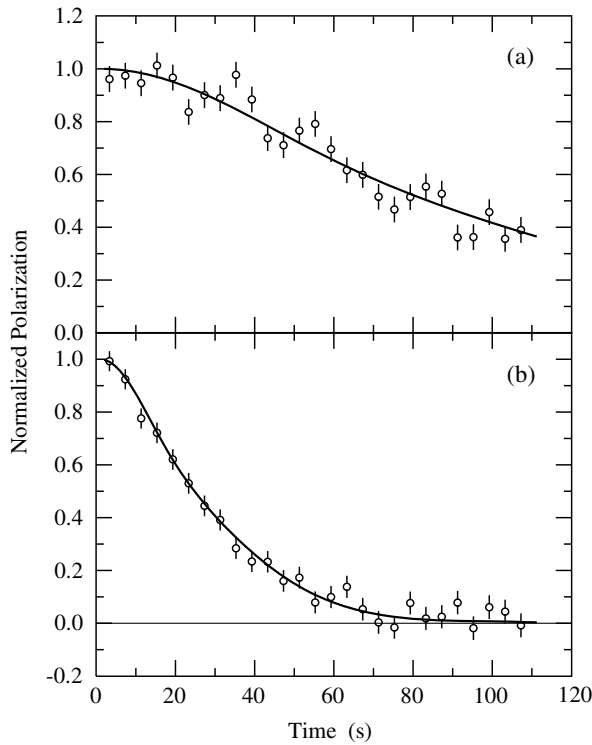


FIG. 2. (from Ref. [1]) Measurements of the IPP as a function of time after the polarization was rotated into the horizontal plane. The scale is set so that the normalized polarization is unity at  $t = 0$  s. The measurements were made with two beam polarization states, separately normalized and then averaged. (Both the spin tune and the phase were constrained in order to improve the analysis accuracy for small IPP values [1].) The two panels represent two different sextupole magnet settings, resulting in polarization lifetimes, defined as the time for the normalized IPP to reach 0.606, or  $64.7 \pm 5.4$  s and  $18.6 \pm 2.6$  s in (a) and (b), respectively.

During a series of runs recorded during a 2015 experiment, the injected beam going into COSY was equal to or less than  $10^9$  deuterons/fill. This led to measurements of the IPP as a function of the time during the store as shown in Fig. 1 of Ref. [1] and reproduced here as Fig. 2. However, the more detailed comparison of sextupole settings for a good IPP lifetime with values of the chromaticity, as shown in Fig. 3 of Ref. [1] and reproduced here as Fig. 3, was made with higher beam currents (up to an order of magnitude) in a longer study completed in 2014. It is this longer study that we describe in this paper.

At the start of the experiment, the chromaticities of the COSY ring were measured for a large grid of sextupole settings. The beam was prepared in coasting mode (no bunching) with electron cooling on, and small velocity changes were made by adjusting the electron beam voltage. The resulting changes in the  $x$  and  $y$  betatron tunes gave the chromaticities. Scatter in the tune measurements was used to estimate the errors in the measurement.

The results for the chromaticities were linear with changes to the  $S$  and  $G$  sextupole current values, as shown in Figs. 4 and 5 for some sections through the  $x$  and  $y$  measurements. Altogether, the chromaticity was measured at 24 locations. More points were taken in single scans for the cases where either the MXS or MXG currents were held constant at zero. All 24 points were used to reproduce the chromaticities with a plane described by

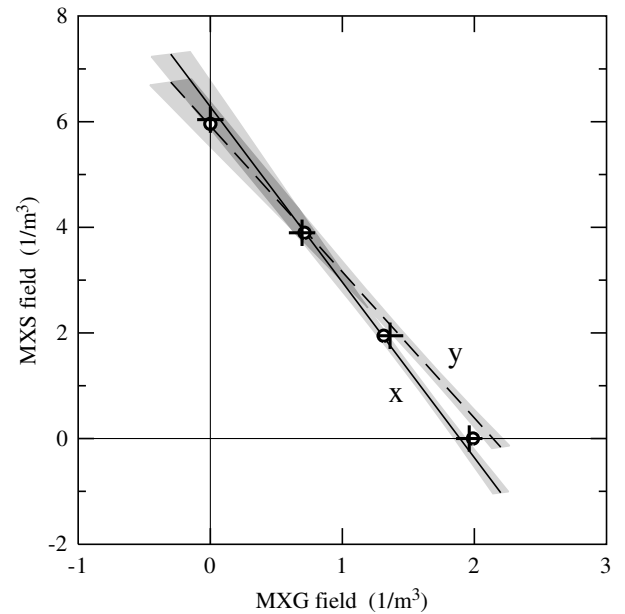


FIG. 3. (from Ref. [1]) This figure represents the  $\text{MXS} \times \text{MXG}$  plane where  $B_{\text{MXL}} = -0.14 \text{ m}^{-3}$ . The gray bands represent the error about the lines of zero  $x$  and  $y$  chromaticity. Scans of the MXS and MXG sextupole family currents found several local points of maximal IPP lifetime. The circles represent a beam setup with horizontal heating; the plus signs represent runs with enhanced longitudinal width. The error bars are smaller than the points.

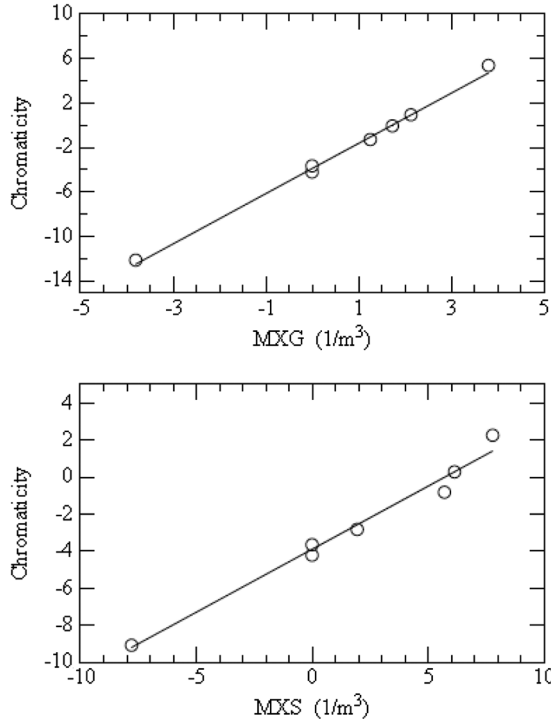


FIG. 4. Measurements of the  $x$  chromaticity as a function of the sextupole setting for either the MXG or MXS magnet families while the other was held at zero. The lines are the slices through the plane corresponding to the scan, which is a fit to all 24 chromaticity points. The field for MXL is  $-0.145 \text{ m}^{-3}$ .

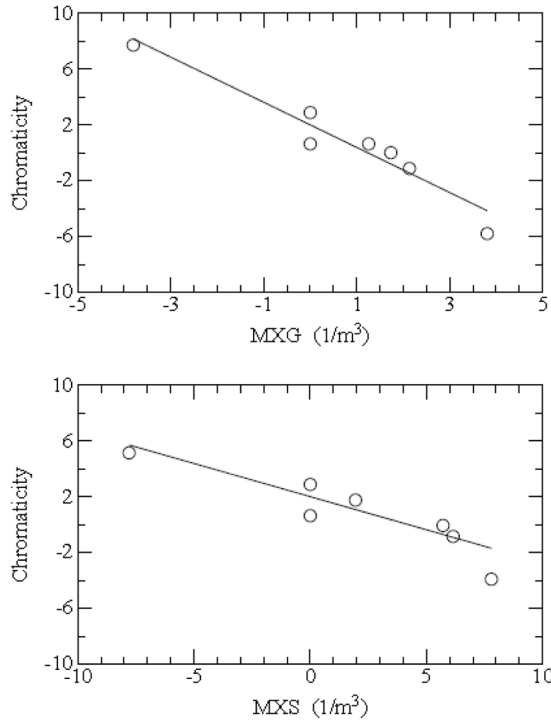


FIG. 5. Measurements of the  $y$  chromaticity in a manner similar to Fig. 1.

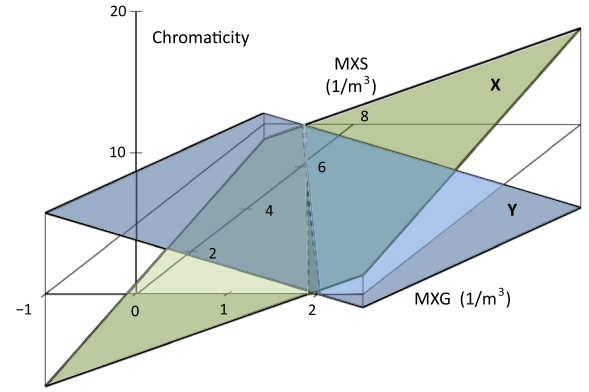


FIG. 6. Values of the  $x$  (green) and  $y$  (blue) chromaticities as a function of the fields in the  $S$  and  $G$  sextupole magnet families. The planes are fits to a set of individual chromaticity measurements. The place where each plane crosses zero chromaticity is indicated by a dashed line.

$$\xi_{x,y} = T + UI_S + WI_G, \quad (6)$$

with  $I_L = -0.145 \text{ m}^{-3}$  and  $T$ ,  $U$ , and  $W$  determined by the fitting program. The panels in Figs. 4 and 5 show the level of agreement between the plane and the  $x$  and  $y$  chromaticity measurements for scans along either axis where  $I_G$  or  $I_S = 0 \text{ m}^{-3}$ . The quality of the agreement may be judged by the scatter of the data points about each curve. The  $x$  and  $y$  chromaticity planes are shown in Fig. 6.

As shown in Fig. 6, the lines of zero chromaticity lie in the  $SG$  plane. The initial IPP lifetime scans showed little sensitivity to  $I_L$  (discussed later), so this sextupole family was set at  $-0.145 \text{ m}^{-3}$  since this value brought the two zero lines (dashed in Fig. 3) as close together as possible. The MXL field was left at that place for the remainder of the study. A typical MXL scan is discussed at the end of Sect. IV.

The chromaticity results shown in Fig. 6 suggest a simpler strategy for adjusting sextupole strengths in a general 3-dimensional search: leave  $I_L$  fixed and scan either  $I_S$  or  $I_G$  so that the scanning path crosses the zero chromaticity line. Then look at the resulting IPP lifetime values for the appearance of a maximum. Or look at the reciprocal of the lifetime for a place where it goes to zero.

In the 2014 study we chose two beam preparations in order to be sensitive in different ways to the origins of the depolarization away from the optimum setting. The first, a “wide” beam, involved electron cooling while bunching as the initial preparation, followed by a short period of horizontal heating using white noise applied to a set of horizontal electric field plates. The white noise frequency spectrum was a band overlapping a harmonic of the betatron tune so that coupling to the beam was enhanced. This resulted in a beam that was physically wider on average than it was tall. At the time when observations of the polarization needed to begin, vertical heating was

applied to the beam so that deuterons began to intersect with the front face of the carbon block target [19]. The target limited the betatron amplitude for vertical motion, and thus limited the contribution of this degree of freedom to the beam depolarization.

The second preparation was intended to increase the spread of the longitudinal phase space in order to determine whether or not this degree of freedom mattered. The beam was first electron cooled without bunching, then, with electron cooling off, the bunching cavity was restarted. The expectation was that this would result in a beam that was physically longer (“long” beam) without having substantially more transverse motion than was there following the cooling process.

In each case, four scans were planned, as shown by bars in Fig. 7. All four scans were completed for the wide beam. A magnet cooling system failure stopped the experiment after the first three scans were completed for the long beam. Consequently our study contains only seven determinations of the optimum IPP polarization along one of these scan paths. After the cooling system was repaired, the fourth scan was completed for the long beam and included in Fig. 2 (from Ref. [1]). It is omitted in the analysis reported here since its consistency with the first seven measurements is not guaranteed.

The beam stores for the wide beam were divided into several time intervals. The first 60 s were devoted to injecting, accelerating, and electron-cooling the beam. The beam was bunched throughout this time. Then cooling was turned off and for another 5 s horizontal heating white noise

was applied to the beam. Beginning at 67 s, a weaker vertical white noise was applied, which drove the edges of the beam into the 17-mm thick carbon block target mounted about 3 mm above the beam center. This lasted for 10 s with the polarization vertical to allow for a measurement of the polarization. At 77 s, the rf solenoid was started, operating at the first harmonic,  $(1 - G\gamma)f_{\text{rev}}$ , of the revolution frequency. The solenoid was on for a little over 2 s, a time adjusted so that the vertical polarization went to zero. Beginning at 80 s, the data acquisition was started. At that time, an adjustment was made to the vertical position of the beam to increase extraction of deuterons onto the carbon target. Then the window for data taking was 90 s long.

The preparation of the long beam was similar except that the horizontal heating time was used instead to start beam bunching. With cooling off at that point, the coasting beam had spread around the ring, so imposing bunching created some large amplitude synchrotron oscillations. The longer beam length was clearly visible on the beam pickups.

Each run consisted of a series of stores. There were three polarization states (positive vector, negative vector, and unpolarized), and each one in turn was used for the injected beam. The stores for each polarization state were separately accumulated and then the results added together during the data analysis. Since the purpose of the study was to produce measurements of the polarization lifetime, no effort was made to quantify the beam polarization value absolutely. What was important was knowing the starting asymmetry measured by the EDDA detector assembly for each polarization state while the polarization was still vertical. At modest beam intensities, this asymmetry was matched at 5 s after the start of the rf solenoid used to precess the polarization into the horizontal plane, indicating no significant polarization loss during rotation. These asymmetries were 0.250 for spin down and 0.229 for spin up. These were both recorded as a positive IPP since it is only the magnitude of the polarization as it rotates in the horizontal plane that is determined. On some occasions, the spin tune for a specific store was far enough away from the average for the set that there was concern the measured IPP would be reduced due to smearing of the phase due to the length of the individual time bins. In these cases, that particular store was dropped from the analysis.

Along with IPP values, detector data rates were also obtained for each time interval. In runs where the initial rate was particularly high, the measured IPP was found to be less than the typical values of 0.250 and 0.229. This led us to consider a correction to the asymmetry for rate. The larger rate came with pile-up in the signal and a background of essentially spin-independent events appearing in the data set. As the largest IPP values typically occurred for the time bin at 5 s, this bin was taken as the standard for comparison. Plotting as a function of rate in the EDDA detector, Fig. 8 shows a clear trend downward with increasing rate. For each time bin, the IPP was divided by the value of the

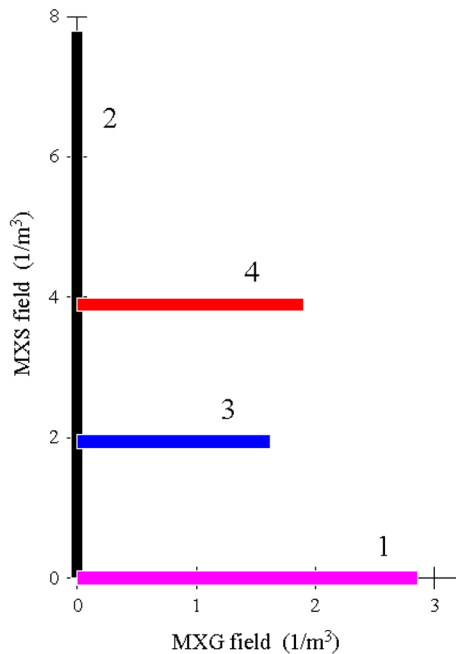


FIG. 7. In this picture of the  $S \times G$  plane, four bars show the paths of four planned scans of the IPP lifetime. They are labeled “1” through “4” for later reference.

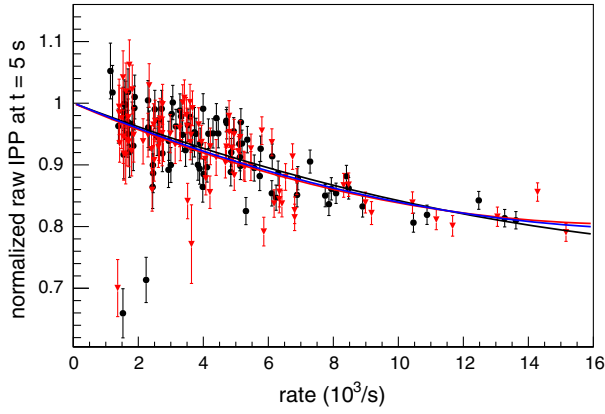


FIG. 8. The normalized value of the measured IPP in time bin 5 as a function of the EDDA detector rate. The curves are a polynomial fit to the measurements. The black curve fits the black points for spin down; the red curve fits the red points for spin up; and the blue curve fits both. The IPP values have been scaled upward by a factor for each polarization state that makes the typical low-rate bin-5 polarization equal to one. In the analysis for all time bins, the measured IPP was adjusted upward by the appropriate factor to compensate for this loss.

average curve in Fig. 8 evaluated for the data rate associated with that bin. This correction is particularly important since we also determined that the most reliable estimate of the IPP lifetime comes from the slope of the tangent curve fit to the data in the early part of each store. A sample of IPP data with a low rate (left panel) and a higher rate (right panel) is shown in Fig. 9.

At the end of the experiment, all the runs were replayed. Cycles with poor information were removed from the data sample. Examples of the results for Scan 3 are shown in Fig. 10. Each one of these panels represents the results of several stores with the same polarization state and sextupole

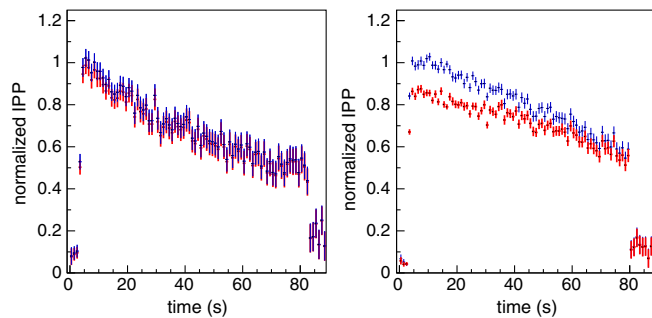


FIG. 9. Measurement of the IPP for two runs, normalized to one at  $t = 5$  s based on the typical asymmetry values described in the text. In the left panel, the detector rate is low and not much change is apparent between red (original measurement) and blue (rate corrected measurement). The right panel contains another example for a higher rate run with a considerable initial slope change.

magnet settings. The average spin tune was  $0.1609710 \pm 0.0000001$  among all of the runs. Rate corrections were applied.

There are various options for defining the IPP lifetime,  $\tau$ , based on the shape of the polarization as a function of time (see Fig. 10). An ensemble of spins with little initial spread often loses polarization slowly at first, giving a time function that is similar to a Gaussian, as seen in Fig. 1 of [1] (see Fig. 1 here) for deuteron beams less than  $10^9/\text{fill}$ . In this case the Gaussian width (half-width at 0.606 of maximum amplitude) makes sense as an IPP lifetime. This feature is present to a lesser extent in the time dependence of the model from Ref. [8], also created for lower beam currents. In other cases, the half-life may be appropriate (see Fig. 4 of [1], reproduced here as Fig. 11) or the width at  $1/e = 0.368$  from the exponential function. The shapes of Fig. 10 present a more complicated situation (associated with beams of about  $10^{10}/\text{fill}$ ) without a single universal shape.

An examination of these shapes showed that the information crucial to a maximization of the IPP lifetime was often contained in the first few seconds of the time function, which could be taken from a polynomial fit to the early part of the data. In Fig. 10, a linear function did well for part of the time in three of the six panels. A quadratic reproduced part of the function in the other three panels, but resulted in two positive and one negative curvature values with no apparent pattern. Other data points were similarly varied. Finally we chose to use either a linear or quadratic polynomial as seemed appropriate for the early part of the time function and to take only the constant and linear terms to determine the IPP lifetime,  $\tau$ . This lifetime was chosen to be the time over which the linear part fell to 0.606 of the initial asymmetry at  $t = 5$  s. In each case a polynomial was fit to the measurements for as long a time as the fit reproduced the data. In some cases, a decline in polarization that began as nearly a straight line deviated within a few seconds to another course. At that point the polynomial reproduction became meaningless. Prior work with such data has indicated that the thick carbon block, which always extracts particles for the polarization measurement from the fringe of the beam, may be sampling first the halo followed by the core of the beam. With a wide beam, these two components may have different polarization distributions, resulting in some runs where there appears to be a break in the loss rate for the IPP. The time function shapes vary run to run, and may be sensitive to orbit distortions or other changes associated with the varying sextupole strengths.

In Fig. 10, the fit is shown by a red line. The red band indicates the width of the error in the line based on the statistics of the fit. The extension of the constant and linear terms in the polynomial is shown by the green line, which is visible wherever it differs from the red band. The time  $\tau$  needed to reach 60.6% of the initial asymmetry was obtained



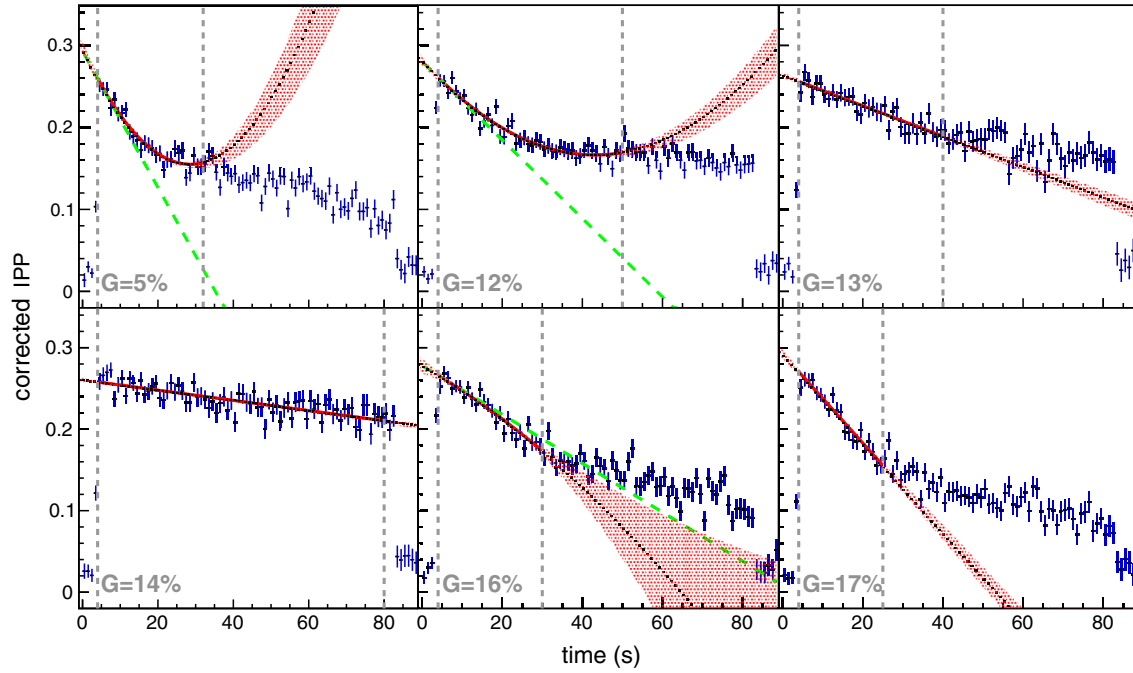


FIG. 10. Panels for representative sextupole values involved in Scan 3 of Fig. 7 for a wide beam and negative polarization. Zero time begins 80 s into the store when the data acquisition was started. The value of  $I_G$  in terms of a percentage of full sextupole current are included in each panel. A region (indicated here between vertical dashed lines) near the beginning of each data set was reproduced by either a linear or a quadratic polynomial (heavy red line). The red region shows the extension of this polynomial, including statistical errors from the fitting process. The green line represents the tangent to this polynomial at 5 s where it differs from a nonlinear polynomial fit.

from the constant and linear polynomial coefficients. Figs. 12 and 13 show panels in which these values are represented logarithmically in a projected representation of the  $S \times G$  plane.

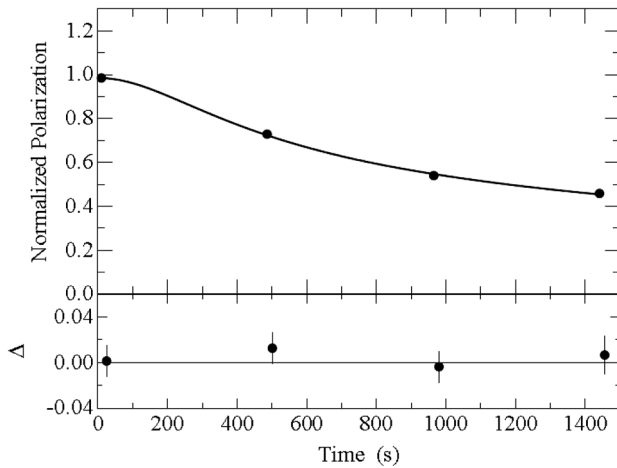


FIG. 11. (from Ref. [1]) Measurements of the normalized polarization. Least squares fits were made to the data acquired after rotation into the horizontal plane during four periods when the beam was extracted onto the polarimeter target. Data points taken with target extraction off are not shown. (The fitting template is based on an unfolded Gaussian beam distribution.) The residuals of the fit are given in the lower panel. The IPP lifetime is  $782 \pm 117$  s.

#### IV. RESULTS AND INTERPRETATION

The curves in Fig. 12 for the wide beam show peaks that indicate the location of the longest IPP lifetimes. These peaks are more narrow than their counterparts in Fig. 13 for the long beam. This is likely a result of the deliberate horizontal heating that was applied to the beam. This spreads the distribution of betatron oscillation amplitudes and spin tunes, thus making it harder to obtain a good match for the various settings of the sextupole fields. So the IPP lifetime drops quickly away from the optimum setting. For the data in Fig. 13, no horizontal heating was applied. Instead, the beam was first cooled as a coasting beam. Afterward bunching was turned on. The resulting beam in general had a larger longitudinal extent in COSY than the wide beam. Again, there are clear peaks in the scans of IPP lifetime, but they are broader. Without the horizontal beam spreading, the lifetime is less sensitive to sextupole setting. This would suggest that the greater longitudinal extent of the beam has not introduced another source of depolarization to replace that caused by the horizontal heating. In fact, for scans 2 (black) and 3 (blue) the best IPP lifetimes are longer than for the wide beam. This suggests that, even at its best, the sextupole cancellation of the horizontal heating was not completely effective. At the same time, the higher order contribution that was expected to arrive for the long beam from  $(\Delta p/p)^2$  terms has not appeared at a similar level of importance. The large differences between the lowest and highest IPP lifetimes point to the importance

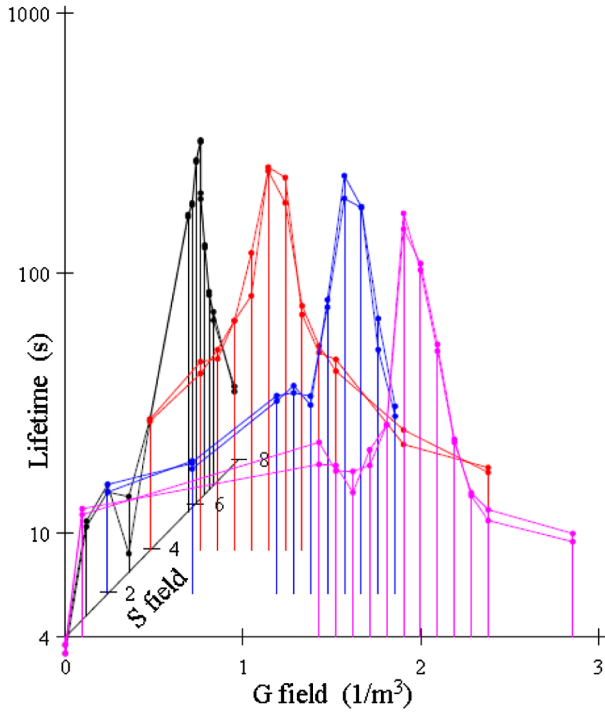


FIG. 12. Representations of the IPP lifetime  $\tau$  on a projected graph as a function of  $S$  and  $G$  field values for the wide beam. The various scans are shown with different colors. Scan 1 is lavender, 2 is black, 3 is blue, and 4 is red. The two polarization states are plotted separately; their values overlap closely and show essentially the same information.

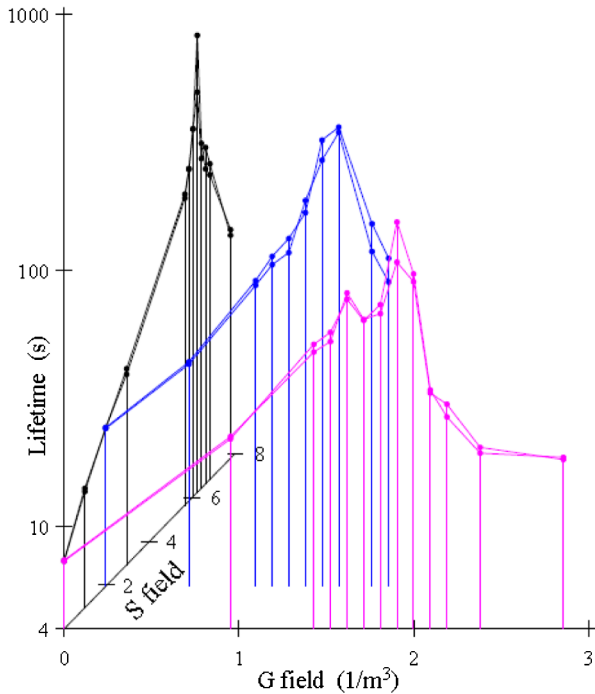


FIG. 13. Same as Fig. 12 except for the long beam. Scan 4 did not have enough data to be considered.

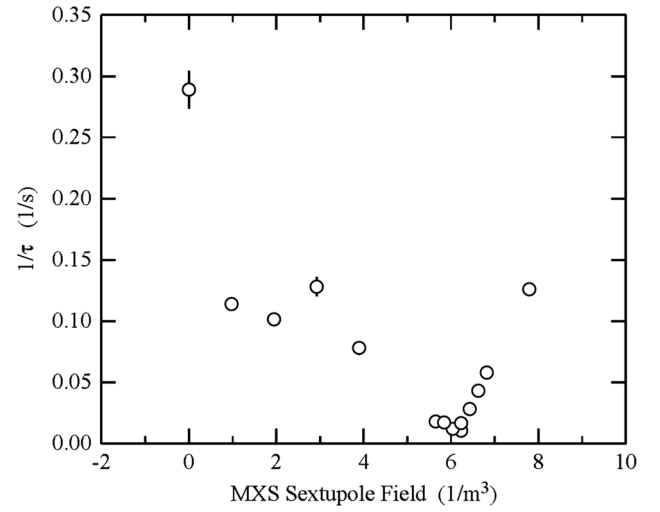


FIG. 14. Measurements of  $1/\tau$  of the IPP for scan 2 and a wide beam. The value of  $MXG$  was zero. Errors shown arise only from the fit of the polynomial to the early time measurements as shown in Fig. 10. Only the points for the spin down polarization state are shown. In general, the two polarization states agree well.

of making sextupole corrections in a ring where the long lifetime is important, as it is for a dedicated EDM storage ring.

The next question to consider is whether the values of the best IPP lifetimes are associated with small or zero values of the chromaticities. These locations may be obtained roughly from the peaks in Figs. 12 and 13. Earlier it was suggested that plots of  $1/\tau$  from Fig. 10 versus sextupole field would have a “V” appearance as suggested by Eq. (5), especially if there was only one important contribution to the depolarization. One example from scan 2 from the wide beam runs is shown in Fig. 14.

There is a feature in the scan just above  $6 \text{ m}^{-3}$  that would suggest the formation of a “V” pattern. The other observation is that the errors, most of which are less than the size of the circles representing the points, are much too small to account for what appear to be fluctuations in  $1/\tau$  from point to point. It is clear that, as the sextupole scans were being created, the setup of the machine itself was slowly changing over a matter of hours in addition to any response to the sextupole changes themselves. One such change might be steering of the orbit. In all cases, the sextupole values for each scan were taken in a random order to avoid the drifts appearing as trends in the measurements. But this may only have caused drifts to appear as random changes from one point to the next. The decision to limit the range of the polynomial fit to only those points in Fig. 10 that could be described with a quadratic curve also served to reduce the size of the statistical errors. This problem will need to be addressed again when we attempt to estimate the errors on the location in sextupole space where the IPP lifetime is the longest. For points that are well away from the “V”, the values of  $1/\tau$  tend to be lower (longer IPP lifetime) than

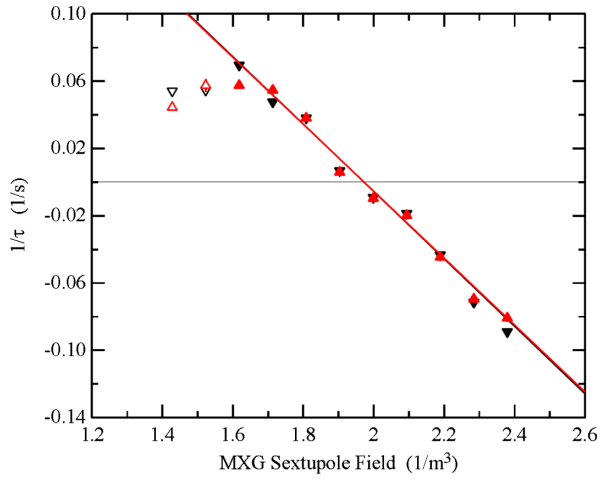


FIG. 15.  $1/\tau$  points for the wide beam scan 1 as a function of the  $G$  sextupole field ( $S$  is zero). The black triangles are for the down polarization state and the red triangles are for the up polarization state. The filled symbols indicate those points included in the straight line fit. Points shown with open symbols were not included. Weights based on the statistical errors were used. This emphasizes points in the fit that are near the zero crossing.

would be expected from a simple extrapolation of a linear fit to the sides of the “V”. This likely indicates the growing importance of higher-order corrections to the simple model of Eq. (5).

In the next series of four figures, we will concentrate only on the region of the “V” for the wide beam. In order to locate more precisely the point with the best IPP lifetime, the “V” pattern may be reproduced with the absolute value of a straight line. Alternatively, if the points that lie to the right of the bottom of the “V” change sign, then the data may be reproduced with a single (two parameter) straight line. As mentioned above, some points begin to deviate from the line when they are away from the point of the “V”. These will be excluded, as well as any points clearly having some large systematic error. The result in the area of the zero crossing for the wide beam scan 1 is shown in Fig. 15. Points outside the range of the plot were ignored in the fit. Points on the plot that were also excluded are shown with an open symbol.

It is important to note that there does not seem to be any excessive scatter in the points near the zero crossing. Points near zero that are systematically shifted away from zero might indicate other contributions to the decoherence. Such contributions must exist as the IPP lifetime cannot reach infinity, as suggested by the straight line. The quality of the fit in this region suggests that these contributions are below our present sensitivity to detect with this point spacing. The overlay of the two straight lines for the two polarization states supports the reproducibility of these results with each given run. Nevertheless, the reduced chi square value ( $\chi^2/\nu$ ) of these fits lies in the hundreds. So when the point where the lines cross zero (the position of the longest IPP lifetime)

TABLE II. Sextupole settings for the best IPP lifetime.

Scan	Sextupole	Spin down	Spin up
Wide beam			
1	MXG	$1.972 \pm 0.010$	$1.973 \pm 0.011$
2	MXS	$6.114 \pm 0.024$	$6.110 \pm 0.027$
3	MXG	$1.365 \pm 0.007$	$1.372 \pm 0.006$
4	MXG	$0.709 \pm 0.006$	$0.697 \pm 0.009$
Long beam			
5	MXG	$1.93 \pm 0.05$	$1.93 \pm 0.05$
6	MXS	$6.23 \pm 0.10$	$6.23 \pm 0.10$
7	MXG	$1.33 \pm 0.05$	$1.33 \pm 0.05$

is calculated, its statistical error is enhanced by the square root of the reduced chi square value in order to compensate for systematic changes point to point. The resulting zero crossing values are given in Table II. The remaining fits for the wide beam are shown in Figs. 16, 17, and 18.

Data for the long beam runs is characterized by smaller values of  $1/\tau$ , or longer IPP lifetimes. At the same time the “V” pattern appears to be shifted positively so that the point of the “V” no longer rests on zero. Such a shift causes a discontinuity in the straight line procedure used for the wide beam scans. A third parameter, such as the size of the shift, may be incorporated into the fitting process. But in practice the decisions about which points to include in the fit and where to locate the point of the “V” (to decide which points are to its left and right) leads to ambiguities in the results. In some cases, these are outside the expected bounds, even considering the systematic deviations point to point in these measurements.

So the measurements with the long beam will be converted into an optimum sextupole setting by simple visual inspection of the data, choosing a place in a location near the lowest values of the slope. The error will be

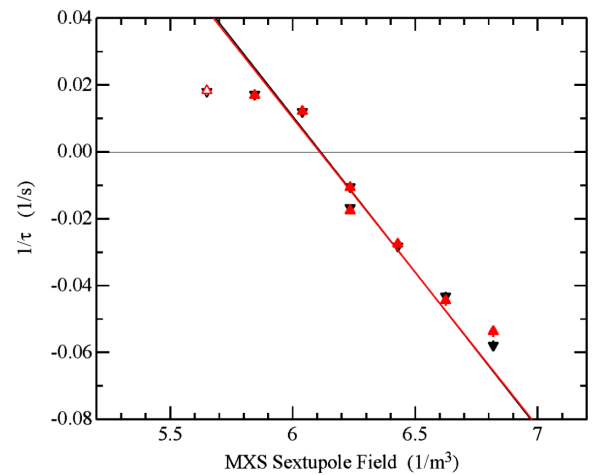


FIG. 16. Same as Fig. 15 but for wide beam scan 2, shown as a function of  $S$  ( $G$  is zero). This is a subset of the data shown in Fig. 14.

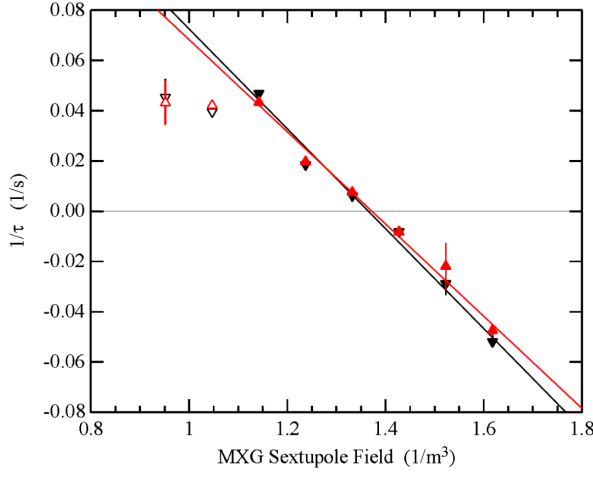


FIG. 17. Same as Fig. 15 but for wide beam scan 3, shown as a function of  $G$  ( $S = 1.95 \text{ m}^{-3}$ ).

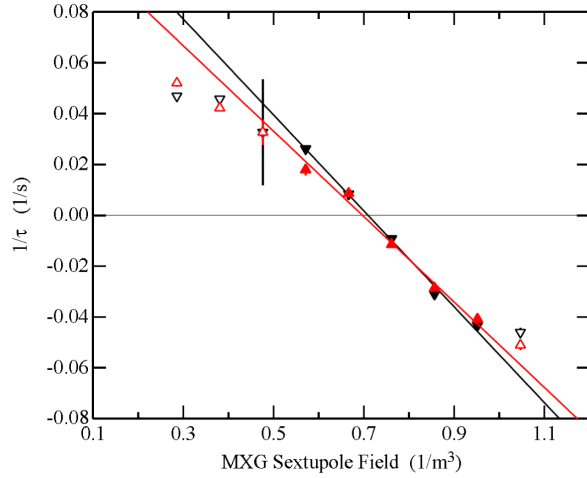


FIG. 18. Same as Fig. 15 but for wide beam scan 4, shown as a function of  $G$  ( $S = 3.90 \text{ m}^{-3}$ ).

$\pm 0.05 \text{ m}^{-3}$  in  $G$  and  $\pm 0.10 \text{ m}^{-3}$  in  $S$ . Also, there do not appear to be any significant differences in the value taken for the two polarization states. These scans are shown in Figs. 19, 20, and 21.

The amount by which the point of the “V” in  $1/\tau$  for the long beam rises above zero is clearly less than  $0.01 \text{ s}^{-1}$ . This could indicate the extent to which the long beam had introduced depolarization effects associated with this beam preparation method and which cannot be effectively compensated with the sextupole field manipulation. Additional tests would be needed to identify this depolarization source.

The values listed in Table II may be placed on the same template that was used for Fig. 3 of Ref. [1] (shown as Fig. 3 here). The result is Fig. 22. The two bands show where the  $x$  and  $y$  chromaticities were found to be zero, as indicated in Fig. 6. In this case, errors based on the internal consistency of the chromaticity measurements were used to

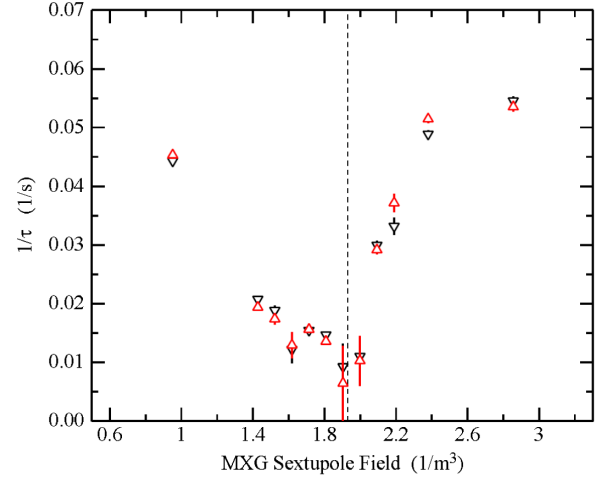


FIG. 19. Measurements made with the long beam of  $1/\tau$  as a function of the  $G$  sextupole field (Scan 1). The  $S$  field is zero. The optimum field value is indicated by a dashed line.

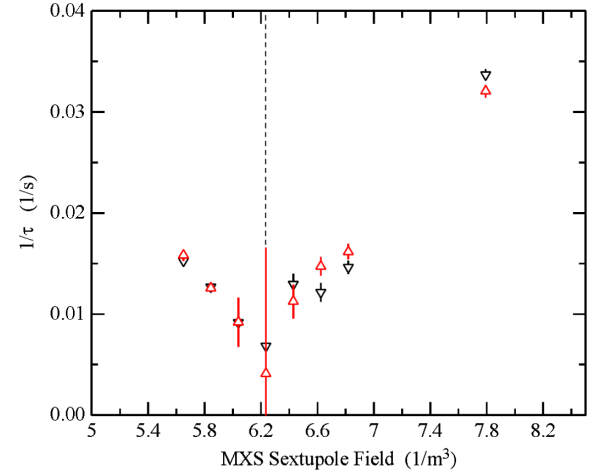


FIG. 20. Measurements made with the long beam of  $1/\tau$  as a function of the  $S$  sextupole field (Scan 2). The  $G$  field is zero. The optimum field value is indicated by a dashed line.

specify bands that illustrate the one standard deviation width in the measurements of the chromaticity. The symbols used to locate the points of maximum IPP lifetime are circles for the wide beam and plus signs for the long beam. In both cases, the errors are smaller than the symbol size. The circles and plus signs are consistent with the error bands on the chromaticity, as would be expected on the basis of the formalism described in Sec. II.

Following the power failure, a check was made of the setting of the MXL sextupole determined originally from the value chosen to make the zero chromaticity lines coincide (see Fig. 6). A scan was conducted with  $S = 6.04 \text{ m}^{-3}$  and  $G = 0 \text{ m}^{-3}$ , a place close to the optimum IPP result for scan 2. Because of the different machine setup following recovery, a new calibration of the rate correction



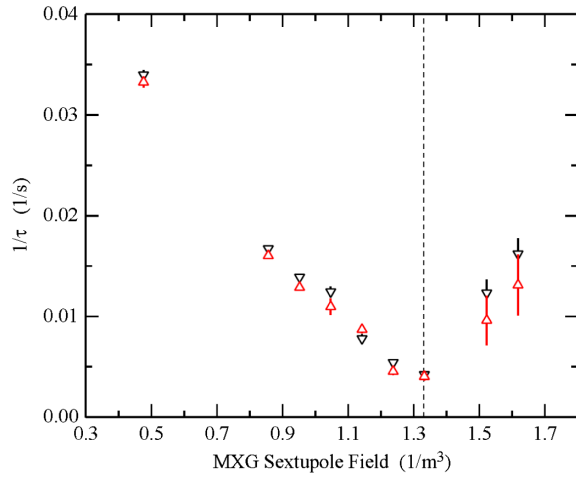


FIG. 21. Measurements made with the long beam of  $1/\tau$  as a function of the  $G$  sextupole field (Scan 3). The  $S$  field is  $1.95 \text{ m}^{-3}$ . The optimum field value is indicated by a dashed line.

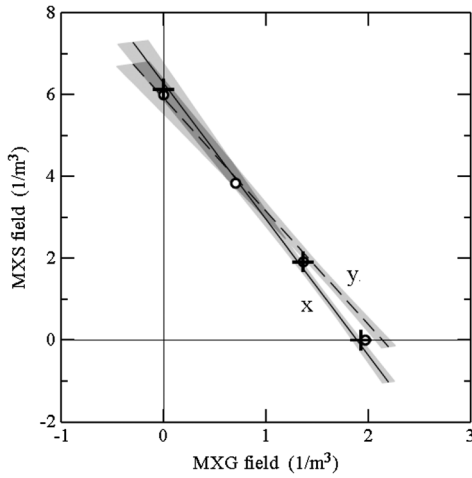


FIG. 22. Two lines with error bands show the places where the  $x$  and  $y$  chromaticities were consistent with zero. The locations of the points of largest IPP lifetime are shown by the circles for the wide beam and the plus signs for the long beam. For the long beam, the point from Scan 4 is missing.

was required and the nominal IPP magnitudes for polarization down and up were changed to 0.230 and 0.219, respectively. Fig. 23 shows the values of  $1/\tau$  for a scan of the MXL field on either side of the optimum point determined from the chromaticity measurements. In this case the dependence does not follow the “V” shape noted for the other scans. A fit using a quadratic polynomial shows a minimum value close to the setting obtained at the beginning but at a slightly more negative value of the magnet field. This supports the general conclusion that vanishing chromaticities are associated with sextupole magnet settings that create simultaneous zeros for the  $x$  and  $y$  chromaticities.

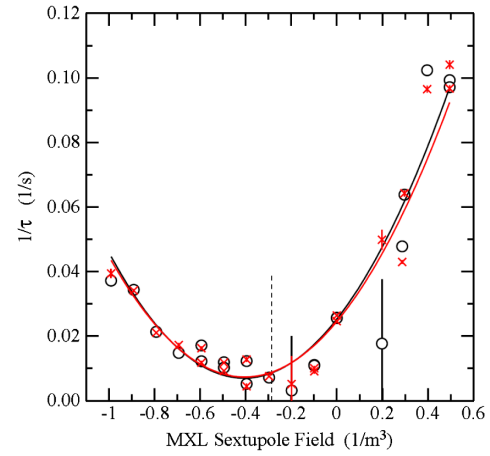


FIG. 23. Measurements of  $1/\tau$  as a function of changing MXL field for  $S = 6.04 \text{ m}^{-3}$  and  $G = 0 \text{ m}^{-3}$ . The black circle points are for polarization negative and the red X points for polarization positive. The dashed line marks the point chosen to bring the  $x$  and  $y$  chromaticities close together. The corresponding parabolic curves are a guide to the eye for each polarization state.

## V. CONCLUSIONS

The search for an EDM using a storage ring requires that we begin with the polarization of the beam parallel to the velocity, a situation that is unstable against momentum spread and depolarizing effects arising from betatron oscillations. In order to have a sufficient time for observation (about 1000 s) in any particular beam store, techniques are needed to suppress the depolarization of the beam once the polarization is rotated into the storage ring plane. Bunching and electron cooling together are not sufficient, so we embarked on a study to determine whether or not further progress could be made using the fields of the sextupole magnet families built into the arcs of the COSY storage ring. It was expected that the quadratic nature of the sextupole field away from the stable beam orbit could be adjusted to match the path lengthening that arises from betatron oscillations, which are still present to some extent even after electron cooling. With the right sign and scale, one would cancel the other and a major source of depolarization would be eliminated.

This proved to be correct, and a search of sextupole space on the COSY storage ring showed that there are places where the IPP lifetimes become very long. These places are not unique, but fall along a line in sextupole space. Independent measurements found that along the same line, measurements reveal almost zero chromaticity. With some sextupole field adjustment prior to running, these lines of  $x$  and  $y$  chromaticity may be brought together, and that combined line becomes the locus of settings with long IPP lifetimes. A simple analysis confirms this association.

Guidoboni [1] has shown that, under optimal conditions, the best IPP lifetimes may exceed 1000 s, as shown in

Fig. 11. This achievement may also be considered from the perspective of an oscillating system of particles whose motion remains coherent for a large number of periods. Such coherence may be expressed as the “quality factor”  $Q$ . With a  $1/e$  lifetime [1] of  $2280 \pm 336$  s in Fig. 11, then  $Q = \omega_0 \tau / 2 = -\pi f_{\text{rev}} G \gamma \tau = 8.6 \times 10^8$ , a very large value.

The question of operating above  $10^9$  deuterons/fill in the COSY experiment has to do with complications that arise at larger currents. In particular, the beam appears to separate into two components, a halo with a particularly short IPP lifetime that responds well to sextupole corrections and an inner core that does not. In addition, collective effects, such as head-to-tail oscillations of the entire beam, appear. Their impact on the polarization lifetime is not understood. These effects are likely to depend on the ring lattice used for the EDM experiment, and need to be reevaluated in that context.

### ACKNOWLEDGMENTS

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