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**New observational constraints on warm rain processes and their climate implications**

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**Main points:**

1. Constrain a microphysics scheme using new cloud and drizzle retrievals
2. Implemented the updated scheme in NCAR CESM to simulate the warm rain globally
3. The updated scheme alleviates a long-lasting problem in most climate models

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32 **Abstract:**

33 Low stratiform clouds have profound impacts on the hydrological cycle and the Earth's radiation  
34 budget. However, realistic simulation of low clouds in climate models presents a major challenge.  
35 Here we employ the newly retrieved cloud and drizzle microphysical properties to improve the  
36 autoconversion and accretion parameterizations in a microphysical scheme. We find that the new  
37 autoconversion (accretion) rate contributes 14% lower (greater) to total drizzle water content than  
38 the original scheme near the cloud top. Compared to satellite results, the simulated cloud liquid  
39 water path (LWP) and shortwave cloud radiative effect using the original scheme in a climate  
40 model agree well on global average but with large regional differences. Simulations using the  
41 updated scheme show a 7.3% decrease in the light rain frequency, and a 10% increase in LWP.  
42 The updated microphysics scheme alleviates the long-lasting problem in most climate models, i.e.  
43 'too frequent and too light precipitation'.

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46 **Plain-language summaries**

47 There has been a growing concern that most climate models predict too frequent and too light  
48 precipitation, which is primarily due to lack of reliable sub-grid variability and vertical variations  
49 of microphysical processes in low clouds. With the newly retrieved cloud and drizzle  
50 microphysical properties from a recent field campaign, we updated the classic warm rain  
51 microphysical scheme which was developed by Khairoutdinov & Kogan (2000) and widely used  
52 in weather and climate models. We examined relative contributions of different microphysical  
53 processes to the rain drop formation and growth processes. The altered scheme reflects the advance  
54 in process-level understanding of warm rain processes and reveals their relative contributions to  
55 the rain drop formation and growth processes at different cloud heights. The altered scheme has  
56 the potential of mitigating the long-lasting problem in most climate models and achieving more  
57 accurate climate assessments. Our findings unambiguously attest the paramount importance of  
58 cloud microphysical parameterizations in climate simulation.

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## 63 1. Introduction

64 Low-level stratiform clouds (hereafter called low clouds) have been a topic of considerable  
65 interest because they strongly reflect incoming shortwave radiation (Stephens et al., 2015) and  
66 exert complex feedbacks on the climate system (Stephens, 2005; Wood, 2012; IPCC, 2013;  
67 L’Ecuyer et al., 2019; Hang et al. 2019). The radiative effect of low clouds contributes to one of  
68 the largest uncertainties in climate modeling (Stephens et al., 2015; IPCC, 2013) and has been well  
69 known to be influenced by aerosols (Penner et al., 2004; Ghan et al., 2016; Seinfeld et al., 2016;  
70 Fan et al., 2016; Li et al., 2020). Drizzle is common in maritime low clouds (Dong et al., 2014  
71 a&b; Wu et al., 2015 & 2017). The formation of drizzle significantly modulates stratocumulus-to-  
72 cumulus transition (Yamaguchi et al., 2017) and plays an important role in determining cloud  
73 lifetime (Albrecht, 1989). Furthermore, they have profound impacts on the hydrological cycle and  
74 the Earth’s radiation budget (Stephens et al., 2015; Wood et al., 2009; Stephens et al., 2010; Suzuki  
75 et al., 2010; Kay et al., 2018), and consequently on the Earth’s climate (Bony et al., 2005; Schmidt  
76 et al. 2006). Despite their importance, it is challenging to simulate low clouds realistically in  
77 climate models where they disagree substantially in the magnitude of cloud feedback for the  
78 regimes of low clouds. As a result, most general circulation models (GCMs) predict too frequent  
79 and too light precipitation (Stephens et al., 2010; Donner et al. 2011; Soden et al., 2011; Lebsock  
80 et al., 2013; Jing et al., 2017 & 2018; Wu et al., 2018; Zhang et al., 2019).

81 Another great challenge in GCMs is how to evaluate the cloud microphysical processes, such  
82 as autoconversion ( $R_{auto}$ ) and accretion ( $R_{accr}$ ) rates in low clouds (Wu et al., 2018; Zhang et al.,  
83 2019) since these processes cannot be directly measured. In fact, these processes are often  
84 parameterized as power law relationships with cloud and drizzle properties in model simulations.  
85 Satellite results have been widely used to evaluate these processes and concluded that GCM  
86 simulations are very sensitive to the choices of threshold cloud droplet radius in simulating the  
87 cloud-to-rain particle conversion and growth processes (Suzuki et al., 2010, 2013 & 2015;  
88 Nakajima et al., 2010). However, satellite retrievals suffer relatively large uncertainties,  
89 originating from their measurement and retrieval errors, as well as their limitations in observing  
90 clouds, especially for drizzling clouds (Suzuki et al., 2010; Ma et al., 2018). Most GCMs predict  
91 too frequent and too light precipitation due to lack of reliable sub-grid variability and vertical  
92 variations of  $R_{auto}$  and  $R_{accr}$  (Jing et al., 2017; Wu et al., 2018; Zhang et al., 2019; Suzuki et al.,  
93 2010; Golaz et al., 2002; Liu et al., 2007; Cheng and Xu, 2009; Wood and Hartmann, 2006).

94 The retrieved cloud and drizzle microphysical properties during the Aerosol and Cloud  
 95 Experiments in the Eastern North Atlantic (ACE-ENA) field campaign (Wu et al., 2020) have 1-  
 96 min temporal and 30-m vertical resolution, which are important for studying warm rain processes.  
 97 In this study, we used the retrievals to recalibrate the  $R_{auto}$  and  $R_{accr}$  parameterizations in  
 98 Khairoutdinov & Kogan (2000) scheme (hereafter called KK) into the new KK scheme (hereafter  
 99 called NKK). The profiles of  $R_{auto}$  and  $R_{accr}$  ( $R_{auto}(Z)$  and  $R_{accr}(Z)$ , where  $Z$  is the in-cloud height)  
 100 can be used to advance the process-level understanding of warm rain process. To further test the  
 101 NKK scheme, we implemented the  $R_{auto}(Z)$  and  $R_{accr}(Z)$  into the National Center for Atmospheric  
 102 Research (NCAR) Community Earth System Model (CESM, Morrison and Gettelman, 2008;  
 103 Hurrell et al., 2013) to simulate the warm rain frequency and intensity globally.

104

## 105 2. Methods

### 106 2.1 Recalibrate the KK warm rain scheme

107 A brief summary about the ground-based retrieval (Wu et al., 2020) is presented in the  
 108 supplementary. The retrieved microphysics include cloud-droplet (drizzle) number concentration  
 109  $N_c$  ( $N_d(Z)$ ) liquid water content  $LWC_c(Z)$  ( $LWC_d(Z)$ ) and mass weighted mean radius  $r_c(Z)$  (=  $\frac{3LWC_c}{4N_c\pi\rho_w}$ )<sup>1/3</sup>) and  $r_d(Z)$ . The units of  $LWC$ ,  $r$ , and  $N$  in this study are  $\text{gm}^{-3}$ ,  $\mu\text{m}$ , and  $\text{cm}^{-3}$ ,  
 110 respectively. Figures S1a-S1d show the retrieved cloud and drizzle microphysical properties. The  
 111 retrieved  $LWC_c(Z)$  and  $r_c(Z)$  increased from the cloud base, peaked just below the cloud top, and  
 112 then decreased toward the cloud top. The retrieved  $r_d(Z)$  and  $LWC_d(Z)$  (Figs. S1c & S1d), opposite  
 113 to their cloud counterparts, increased from the cloud top downward, peaked in the middle or  
 114 bottom of the cloud, and decreased further down.

116 The  $R_{auto}$  and  $R_{accr}$  are usually parameterized as power law relationships with cloud and rain  
 117 water mixing ratios ( $q_c$  and  $q_r$ ) and  $N_c$  (KK; Kessler, 1969; Tripoli et al., 1980; Beheng, 1994; Liu  
 118 and Daum, 2004), and in the forms of

$$119 \quad R_{auto}(Z) = \left(\frac{\partial q_r}{\partial t}\right)_{auto} = A q_c^{a1}(Z) N_c^{a2}, \quad (1)$$

$$120 \quad R_{accr}(Z) = \left(\frac{\partial q_r}{\partial t}\right)_{accr} = B (q_c(Z) q_r(Z))^b, \quad (2)$$

121 where  $A$ ,  $a1$ ,  $a2$ ,  $B$ , and  $b$  are coefficients in different schemes and are usually constants. In this  
 122 study,  $q_c(Z)$  and  $q_r(Z)$  can be calculated from retrieved  $LWC_c(Z)$  and  $LWC_d(Z)$  and dry air density  
 123 ( $\rho_{air}$ ), therefore,  $R_{auto}(Z)$  and  $R_{accr}(Z)$  are a function of height  $Z$  in Eqs. (1) and (2). In addition to

124  $LWC_c(Z)$  and  $LWC_d(Z)$ , the retrieved layer-mean  $N_c$  is also used in Eq. (1). As a proof of concept,  
 125 the KK scheme is used as an example in this study, in which  $A = 1350$ ,  $a1 = 2.47$ ,  $a2 = -1.79$ ,  
 126  $B = 67$ , and  $b = 1.15$ .

127 The summation of  $R_{auto}(Z)$  and  $R_{accr}(Z)$  is the total drizzle water production rate ( $P_r(Z)$ ), which  
 128 can be converted to  $LWC_d(KK)$  ( $= \int \rho_{air} P_r(Z) * dt$ ,  $dt$  is 1 min here) to directly compare with  
 129 retrieved  $LWC_d(Z)$  with 1-min temporal resolution. The retrieved  $LWC_d(Z)$  is then used to scale the  
 130 time interval of  $LWC_d(KK)$  within 1 minute as

$$131 \quad R'_{auto}(Z) = \frac{LWC_d(Z)}{\int \rho_{air} P_r(Z) dt} R_{auto}(Z) = A'(Z) q_c^{2.47}(Z) N_c^{-1.79}, \quad (3)$$

$$132 \quad R'_{accr}(Z) = \frac{LWC_d(Z)}{\int \rho_{air} P_r(Z) dt} R_{accr}(Z) = B'(Z) (q_c(Z) q_r(Z))^{1.15}. \quad (4)$$

133  $A'(Z)$  and  $B'(Z)$ , which are functions of height  $Z$  with unitless, can be calculated from Eq. (3) and  
 134 Eq. (4). To clarify the terms used in this study,  $R_{auto}(Z)$  and  $R_{accr}(Z)$ , and constants  $A$  and  $B$  are  
 135 used in KK scheme, while  $R'_{auto}(Z)$  and  $R'_{accr}(Z)$ , and  $A'(Z)$  and  $B'(Z)$  are used in NKK scheme.

136 A sensitivity study has shown that the coefficients  $A$  and  $B$  in Eqs. (1) and (2) are more or less  
 137 dependent on the retrieved  $r_c$  and  $r_d$  than other coefficients  $a1$ ,  $a2$ , and  $b$ . Also, because of the  
 138 linear constraining in Eq. (3) and Eq. (4), the linear coefficients  $A$  and  $B$  are modified in this study  
 139 and the exponential terms are retained. The relationship between  $q_c$  and  $N_c$  in KK does constrain  
 140  $r_c$ , but this constraint seems too weak in the lower part of cloud and too strong in the upper (Fig.  
 141 S1f).

142 The  $LWC_d(KK)$  profiles peaked in the center and upper part of the cloud (Fig. S1e), which are  
 143 different from the retrieved  $LWC_d(ret)$ . The ratios of  $LWC_d(KK)$  to  $LWC_d(ret)$  (Fig. S1f) show that  
 144  $LWC_d(KK)$  were overestimated in the upper part and underestimated in the lower part of the cloud.  
 145 The higher  $LWC_d$  ratios in the upper part and reduced ratios in the lower part of the cloud suggest  
 146 that it is imperative to recalibrate and constrain the KK scheme using observations. The profiles  
 147 of the  $LWC_d$  ratios in Fig. S1f motivate us to modify the coefficients  $A$  and  $B$  as a function of  
 148 height  $Z$ , such as  $A'(Z)$  and  $B'(Z)$ , not constants with height. Physically,  $A'(Z)$  and  $B'(Z)$  should  
 149 strongly correlate with the profiles of cloud and drizzle microphysical properties, more precisely,  
 150 to  $r_c(Z)$  and  $r_d(Z)$ .

151 Note that KK scheme is used as an example in this study in which we first calculate  $P_r(Z)$  by  
 152 summing  $R_{auto}(Z)$  and  $R_{accr}(Z)$ , and then constrain the original  $R_{auto}(Z)$  and  $R_{accr}(Z)$  with retrieved  
 153  $LWC_d(Z)$  in Eq. (3) and Eq. (4). The approaches can only be applied to schemes having both  $R_{auto}$

154 and  $R_{accr}$  parameterizations, while those having one of them, such as having  $A_{auto}$  in Liu and Daum  
155 (2004), cannot be modified in this study.

156

## 157 **2.2 CESM Simulations**

158 The NSF/DOE Community Earth System Model (CESM) version 1.2.2 (Hurrell et al., 2013) is  
159 used in this study to access the impacts of altered warm rain microphysics on cloud and  
160 precipitation on the regional and global scale. The atmospheric component, Community  
161 Atmosphere Model version 5.3 (CAM5), possesses notable improvements in simulating cloud and  
162 precipitation. In particular, a prognostic two-moment stratiform cloud microphysics scheme  
163 (Morrison and Gettelman, 2008) was implemented in the CAM5 for the first time. The KK scheme  
164 is used for the warm rain process.

165 Six-year equilibrium present-day forcing simulations (2000-2005) are performed for each  
166 microphysical configuration. The first year (2000) is considered as spin-up, and the last five-year  
167 results (2001-2005) are analyzed. To assess the Probability Density Function (PDF) of  
168 precipitation rates, typically high frequency (e.g., hourly) precipitation output is required.  
169 However, it is computationally expensive to store such high-frequency output in a global climate  
170 model. In this study, we adopted an in-situ diagnostic method (Wang et al., 2016) to generate  
171 precipitation probability density functions (PDFs) based on rain rates at each model time step (the  
172 hourly time scale) because this method can also facilitate the comparison between GCM simulated  
173 and satellite retrieved transient precipitation rates (Aumann et al., 2018).

174

## 175 **3. Constrain cloud microphysics scheme using ground-based retrievals**

176 The  $R_{auto}(Z)$  and  $R_{accr}(Z)$  profiles, calculated from ground-based retrievals using Eqs. (1) and  
177 (2), for the case of July 18, 2017 are presented in Fig. S2. As expected,  $R_{auto}(Z)$  increases with  
178 height, basically follows the  $LWC_c(Z)$  profiles. During drizzle drops falling process,  $R_{accr}(Z)$   
179 becomes increasingly important as demonstrated in Fig. S2b where  $R_{accr}(Z)$  is the largest in the  
180 cloud center. To recalibrate  $R_{auto}(Z)$  and  $R_{accr}(Z)$  in the KK scheme, we used the retrievals to derive  
181  $A'(Z)$  and  $B'(Z)$  profiles as demonstrated in Figs. S2c & S2d. Figures 1a and 1b show the  
182 probability density functions (PDFs) and cumulative density functions (CDFs) of  $A'(Z)$  and  $B'(Z)$ .  
183 The prescribed  $A$  and  $B$  values (constants) in the KK scheme fall in the same bins as their mode  
184 values, suggesting that the prescribed values are representative for the most scenarios.

185 To be applicable of  $A'(Z)$  and  $B'(Z)$  in model simulations, we parameterized  $A'(Z)$  and  $B'(Z)$   
 186 as exponential functions of  $r_c(Z)$  and  $r_c(Z)/r_d(Z)$  in Figs. 1c and 1d as:

$$187 \quad A'(Z) = 121683 \exp(-0.528 r_c(Z)) + 364, \quad (5)$$

$$188 \quad B'(Z) = 632 \exp\left(-24.5 \frac{r_c(Z)}{r_d(Z)}\right) + 51. \quad (6)$$

189 The profiles of  $A'(Z)$  and  $B'(Z)$  in Fig. S2 are smaller than their prescribed values in the upper  
 190 part and greater in the lower part of the cloud. The joint PDFs of  $A'(Z)$  and  $r_c(Z)$  with an  
 191 exponential relationship between them (Eq. 5) are illustrated in Fig. 1c. We find that  $A'(Z)$   
 192 decreases with increasing  $r_c(Z)$  and becomes smaller than the prescribed  $A$  when  $r_c(Z)$  is greater  
 193 than  $\sim 9 \mu\text{m}$ . Physically,  $R_{auto}$  should increase with increasing  $r_c(Z)$ , the result here suggests that  
 194 the power law functions for  $q_c(Z)$  and  $N_c$  in Eq. (1) are too strong and the fitted exponential formula  
 195  $A'(Z)$  acts to reduce the power law relationship and bring  $R'_{auto}(Z)$  to more reasonable values to  
 196 correct the overestimated  $LWC_d(KK)$  in the upper part and underestimated  $LWC_d(KK)$  in the lower  
 197 part of the cloud as shown in Fig. S1f. Introducing  $r_c(Z)$  in Eq. (3) adds a more direct constraint  
 198 on the autoconversion process than the relationship originally used in KK because the  
 199 autoconversion process is primarily a conversion process from cloud droplets to drizzle drops near  
 200 the cloud top.

201 Similarly, we fitted an exponential function between  $B'(Z)$  and  $r_c(Z)/r_d(Z)$  (Eq. 6) in Fig. 1d.  
 202  $B'(Z)$  decreases with increasing  $r_c(Z)/r_d(Z)$  until the ratios reach  $\sim 0.2$ . The fitted formula and the  
 203 pattern of the joint PDF, as well as Fig. S1f, reveal that  $B'(Z)$  values should change with height.  
 204 Near the cloud top,  $r_c(Z)$  is the largest while  $r_d(Z)$  is the smallest, resulting in the greatest  $r_c(Z)/r_d(Z)$ ,  
 205 where  $B'(Z)$  from the NKK scheme are the smallest and remain nearly constant ( $\sim 50$ ) for  
 206  $r_c(Z)/r_d(Z) > 0.2$ , even smaller than the prescribed  $B$  ( $B=67$ ). From the cloud top to the cloud base,  
 207  $r_c(Z)$  decreases but  $r_d(Z)$  increases, resulting in the smallest  $r_c(Z)/r_d(Z)$  at the bottom of the cloud.  
 208  $B'(Z)$  increases with decreasing  $r_c(Z)/r_d(Z)$  from the top to the base and reaches the largest value  
 209 at the bottom of the cloud. This change in  $B'(Z)$  will counterbalance the overestimated  $LWC_d(KK)$   
 210 in the upper part and underestimated  $LWC_d(KK)$  in the lower part of the cloud as demonstrated in  
 211 Fig. S1f.

212 Theoretically, the collision efficiency is the highest and reaches nearly unity for  $r_c(Z)/r_d(Z) >$   
 213  $0.2$ , while the collision efficiency decreases significantly with decreasing  $r_c(Z)/r_d(Z)$  (Rogers and  
 214 Yau, 1989). With fixed drizzle drop size, larger cloud droplets have higher collision efficiency and

215 correspondingly larger  $R_{auto}$  values, which typically happens near the cloud top. For smaller cloud  
 216 droplets, their collision efficiencies are much lower because they tend to follow the streamlines  
 217 around a falling drizzle drop. However, the coalescence efficiency is opposite to the collision  
 218 efficiency, that is, smaller cloud droplets more easily stay with drizzle drops and remain joined.  
 219 This argument is further proved from the retrieved  $r_c(Z)$ ,  $r_d(Z)$ ,  $LWC_c(Z)$  and  $LWC_d(Z)$  in Fig. S1  
 220 and  $R_{auto}$  and  $R_{accr}$  in Figs. 1e and 1f.  $R_{auto}$  contribution to drizzle water content increases with  
 221 height, peaking near the cloud top which basically follows the  $r_c(Z)$  and  $LWC_c(Z)$  variations, while  
 222  $R_{accr}$  contributes most near the cloud base which is attributed by  $r_d(Z)$  and  $LWC_d(Z)$ .

223 The fitted exponential formula between  $B'(Z)$  and  $r_c(Z)/r_d(Z)$  in Fig. 1d is opposite to the  
 224 theoretical collision efficiency. This is because the KK scheme tends to overestimate  
 225  $LWC_d(KK)$  near the cloud top and  $B'(Z)$  should be decreased in order to lessen  $R_{accr}$ . Near the  
 226 cloud base,  $B'(Z)$  is usually the largest from the NKK scheme, which acts to enhance  $R_{accr}$  to  
 227 compensate the underestimation of  $LWC_d(KK)$ .

228 For warm rain processes, cloud droplets normally form at the cloud base, grow with height  
 229 through condensation in updrafts into the largest cloud droplets ( $r_c \sim 20 \mu\text{m}$ , Rogers and Yau, 1996;  
 230 Wood, 2005a&b; Wallace and Hobbs, 2006; Takahashi et al., 2017), and become drizzle-sized  
 231 drops through the collision-coalescence near the cloud top in which  $R_{auto}$  becomes important (Wu  
 232 et al., 2015; Cheng and Xu, 2009; Liu and Daum, 2004, Wu et al., 2015). These drizzle drops fall  
 233 from near the cloud top grow by collecting cloud droplets and small drizzle drops. As drizzle drops  
 234 fall,  $R_{accr}$  becomes increasingly important.

235 To quantify the cloud-to-rain particle conversion and growth processes, we normalized the  
 236 individual profiles in cloud height coordinate.  $R_{auto}(Z)$  and  $A_{accr}(Z)$  are calculated from prescribed  
 237  $A$  and  $B$  (constants, white dashed lines in Figs. 1c and 1d) in the KK scheme, while the NKK  
 238 scheme  $A'(Z)$  and  $B'(Z)$  are function of  $r_c(Z)$  and  $r_c(Z)/r_d(Z)$  as shown in Figs. 1c and 1d (solid  
 239 while lines). Figure 1e shows the composite profiles of  $R_{auto}$  and  $R_{accr}$  for all the drizzle cases  
 240 during ACE-ENA. The normalized  $R_{auto}$  increased significantly with height, with a peak at  $z_i \sim$   
 241 0.75, and then decreased toward the cloud top. The smaller  $R_{auto}$  values at the cloud top are mainly  
 242 caused by cloud droplet evaporation associated with cloud-top entrainment as observed by aircraft  
 243 in situ measurements (Wu et al., 2020). The normalized  $R_{accr}$  values are, in general, one order of  
 244 magnitude greater than the  $R_{auto}$  except at the cloud top where they are closer. The  $R'_{auto}(Z)$  values

245 are slightly less than the  $R_{auto}(Z)$  values in the upper part of the cloud and greater in the lower part.  
246 The  $R'_{accr}(Z)$  values, on the other hand, are greater than the  $R_{auto}(Z)$  values at all levels.

247 The relative contributions of  $R_{auto}$  and  $R_{accr}$  to total drizzle water production rate ( $P_r = A_{auto} +$   
248  $A_{accr}$ ) are presented in Fig. 1f.  $R_{auto}(Z)$  and  $R_{accr}(Z)$  contribute  $\sim 45\%$  and  $55\%$  of  $P_r(Z)$  near the  
249 cloud top, respectively. As drizzle drops fall,  $R_{accr}$  becomes increasingly important. For the NKK  
250 scheme, the  $R'_{auto}(Z)$  and  $R'_{accr}(Z)$  contribute  $\sim 31\%$  and  $\sim 69\%$  of  $P_r(Z)$  near the cloud top, which  
251 are 14% less and more, respectively, than the contributions from the KK scheme. The relative  
252 contributions of autoconversion (accretion) gradually decrease (increase) toward the cloud base  
253 and have nearly the same in both schemes below  $z_i = 0.3$ . Near the cloud top, the 14% lower  
254 contribution from  $R'_{auto}(Z)$  corroborates that the NKK scheme has lower precipitation frequency  
255 than the KK scheme. On the other hand, the 14% greater contribution from  $R'_{accr}(Z)$  confirms that  
256 the NKK scheme has higher precipitation intensity than the KK scheme. At the upper part of the  
257 clouds, the less (more) autoconversion (accretion) contributions from the NKK scheme  
258 corroborate the notion that the KK scheme overestimated autoconversion rates and underestimated  
259 accretion rates, which could be a reason that most GCMs predict ‘too frequent and too light  
260 precipitation’. Meanwhile, the NKK scheme has the potential to mitigate the outstanding problem  
261 in GCM precipitation simulations and shed light on future model development.

262 Notice that the focus of this study is on the vertical distributions of  $R_{auto}$  and  $R_{accr}$ , and their  
263 impacts on precipitation simulation. The spatial variations of  $R_{auto}$  and  $R_{accr}$ , especially their  
264 subgrid variabilities, should share the equal importance in precipitation simulation. For example,  
265 Wu et al. (2018) calculated the so-called enhancement factors,  $E_{auto}$  and  $E_{accr}$ , using ARM ENA  
266 ground-based observations and retrievals. They found both enhancement factors increase with the  
267 increase of model grid size. These results are similar to those from Xie & Zhang (2015) and results  
268 from satellite observations in Lebsock et al. (2013) and Zhang et al. (2019). Comparing the  
269 prescribed enhancement factors in Morrioso and Gettelman (2008) to the observed ones, a higher  
270  $E_{auto}(3.2)$  and a lower  $E_{accr}(1.07)$  at small grids were used in Morrioso and Gettelman (2008). In  
271 this study, however, we only investigate the vertical distribution of  $A_{auto}$  and  $A_{accr}$  and their impact  
272 on precipitation with prescribed enhancement factors in CESM simulations.

273

#### 274 **4. Impacts of the updated microphysics scheme in climate simulations**

275 The KK scheme has been widely used in cloud-resolving (Seinfeld et al., 2016) and global  
276 climate models, including the NCAR/DOE CESM (Morrisson and Gettleman, 2008; Hurrell et al.,  
277 2013; Gettleman et al., 2019). In this study, we implemented the updated schemes  $R'_{auto}(Z)$  and  
278  $R'_{accr}(Z)$  in CESM version 1.2 (CESM1) to assess the climatic influence of recalibrated cloud  
279 microphysical processes. We first compare the standard CESM1 simulations with the satellite  
280 products to justify the rationale of updating microphysics scheme. The CERES Edition 4 cloud  
281 liquid water path (LWP) retrievals from the Moderate Resolution Imaging Spectroradiometer  
282 (MODIS) and shortwave cloud radiative forcing (SWCF) from the Clouds and the Earth's Radiant  
283 Energy System (CERES) on board of the Terra and Aqua satellites (Minnis et al., 2020) will serve  
284 as the benchmark. Figure 2 shows the CESM1 simulated spatial distributions of maritime LWPs  
285 between  $60^\circ$  S and  $60^\circ$  N with a mean of  $43.0 \text{ g/m}^2$ , which is close to the satellite retrieval ( $44.1$   
286  $\text{g/m}^2$ ). However, large differences exist over some regions. For example, there are positive biases  
287 of LWP in CESM1 over the Inter Tropical Convergence Zone (ITCZ), whereas over the  
288 stratocumulus-prevailing regions like the Southeast Pacific and Southeast Atlantic, the negative  
289 biases can be up to  $-20 \text{ g/m}^2$  as shown in Fig. 2e., which is consistent with the common problem  
290 of GCM, i.e., too frequent drizzle precipitation for stratus and stratocumulus clouds.

291 The spatial distributions of observed and modeled SWCF values have strong negative  
292 correlations with their corresponding LWPs, that is, larger LWP corresponds to stronger negative  
293 SWCF as illustrated in Figs. 2b and 2d. The spatial distribution of the biases in SWCF (Fig. 2f)  
294 mirrors those in LWP (Fig. 2e), indicating that the SWCF biases are largely contributed by those  
295 in LWP. Over  $60^\circ$  S to  $60^\circ$  N, the oceanic SWCF bias is  $-2.3 \text{ W/m}^2$ . In addition to CESM1  
296 simulations, we also simulate cloud LWP using CESM2 (version 2.1.1) whose microphysical  
297 scheme includes an enhancement factor in the KK scheme. However, the LWP in CESM2 is found  
298 to be overestimated by 66% in comparison with satellite observations (Fig. S3). Therefore, we  
299 choose not to test our observational constraints in CESM2.

300 To reveal the relative importance of the changes in  $R_{auto}$  and  $R_{accr}$  parameterizations, we  
301 conducted two model sensitivity studies by using  $R'_{auto}(Z)$  first, and then using both  $R'_{auto}(Z)$  and  
302  $R'_{accr}(Z)$  in CESM1. Figure 3a shows the differences in maritime stratiform cloud LWP between  
303 the simulations using  $R'_{auto}(Z)$  and  $R_{auto}(Z)$  in which  $R'_{auto}(Z)$  significantly increased cloud LWP.  
304 Such an increase is more evident in the mid-latitude regions than the tropics, which can be  
305 attributed to the fact that the stratiform clouds are more prevalent in the mid-latitudes. The

306 increased  $LWPs$  in the mid-latitudes using NKK greatly counterbalance the negative biases in LWP  
307 using KK scheme (Fig. 3e), bring the modeled  $LWPs$  closer to satellite retrievals. In particular, the  
308 simulated  $LWPs$  using  $R'_{auto}(Z)$  increased  $11.8 \text{ g m}^{-2}$  over  $60^\circ \text{ S}-60^\circ \text{ N}$  oceanic regions (Fig. 3a)  
309 which is more than 20% fractional changes, and  $9.8 \text{ g m}^{-2}$  globally (Fig. S4b). The increases in  
310 mean stratiform cloud fractions (CFs) were only 0.5% and 0.8% for the mid-latitudes and globally  
311 (Fig. S5), but much more for fractional changes, up to 10% over the regions like subtropics and  
312 the Arctic as seen in Fig. S5d.

313 The reduced  $R'_{auto}(Z)$  near the cloud top shown in Fig. 1f corroborates the notion that the  
314 overestimation of  $R_{auto}$  is more important in determining the overall  $R_{auto}$  effect than the  
315 underestimation of  $R_{auto}$  in the bottom of the cloud. Hence,  $R'_{auto}(Z)$  exerts a larger influence on  
316 the height dependency of precipitation processes in the cloud. In contrast,  $R_{accr}$  is generally  
317 underestimated throughout the whole cloud profile (Figs. 1e and 1f). Therefore, a stronger  $R'_{accr}(Z)$   
318 can be expected when implementing it in the model simulations, which can result in more cloud  
319 LWP as evident in Fig. 3a. Taking changes by both  $R'_{auto}(Z)$  and  $R'_{accr}(Z)$  into account together,  
320 the net cloud LWP changes (Fig. 3b) are much less than those simulated with  $R'_{auto}(Z)$  only in Fig.  
321 3a, but are still dominated by the impact of the autoconversion change, with a mean increase of  
322  $4.5 \text{ g m}^{-2}$  over  $60^\circ \text{ S}-60^\circ \text{ N}$  oceanic regions, corresponding to a 10% increase. No significant  
323 changes in stratiform CFs using  $R'_{auto}(Z)$  only or both  $R'_{auto}(Z)$  and  $R'_{accr}(Z)$  are found in this study  
324 (Fig. S5).

325 Cloud-to-rain particle conversion is also crucial for drizzle formation process in clouds. With  
326 suppressed autoconversion rates near the cloud top, the  $R'_{auto}(Z)$  results in significant reductions  
327 in precipitation frequency, particularly in the subtropical regions (Figs. S4c and S4d). Figure S4d  
328 illustrates the decreased precipitation frequency corresponded with increased cloud LWP (Fig. S4b)  
329 and stratiform CF (Fig. S5), although they were imperfectly matched in their spatial distributions.  
330 Similar to the cloud responses, the rain formation process is dominated by the autoconversion  
331 change. This conclusion is further confirmed in Figs. S4c and S4d where the mean absolute  
332 changes in global precipitation frequency are -3.1% with  $R'_{auto}(Z)$  only and -2.6% with both  
333  $R'_{auto}(Z)$  and  $R'_{accr}(Z)$ , with a significant decrease over the tropical regions.

334 To further probe the surface precipitation changes as a function of rain intensity, we employed  
335 an in situ diagnostic method (Wang et al., 2016) to generate precipitation PDFs based on rain rates  
336 on the hourly time scale. During the model integration, at each model time step, the new diagnostic

337 accumulates instantaneous precipitation rates into 30 predefined bins. At the end of each month,  
338 the corresponding percentage for each bin can be calculated to obtain a PDF and output it in the  
339 monthly data. The model sensitivity run shows that for the stratiform clouds, the mean frequency  
340 of drizzle or light precipitation (intensities less than 5 mm/day) was reduced from 25.2% to 22.7%  
341 when  $R'_{accr}(Z)$  was applied (Fig. 3e), corresponding to a 10% fractional decrease. With an elevated  
342 accretion rate in the NKK scheme updating both processes, a 7.3% (fractional) drizzle reduction  
343 still exists. The reduced precipitation frequency with the NKK scheme alleviates a long-lasting  
344 problem related with the precipitation in GCMs (Stephens et al., 2010). It is difficult to obtain  
345 stratiform precipitation from observations, so we do not compare the stratiform precipitation PDF  
346 with observations in this study. Wang et al. (2016) examined the total precipitation PDF in the  
347 CESM1 simulations using the KK scheme, and found that the simulated precipitation frequency  
348 for light precipitation frequency is 5% higher than the Tropical Rainfall Measuring Mission  
349 (TRMM, Lau and Wu, 2011) observations (54%) over the tropical region of 25° S-25° N. Although  
350 it is not the same region as this study (60° S-60° N), this result corroborates that the simulated  
351 precipitation frequency using the NKK scheme is changing towards to observed one.

352

## 353 **5. Conclusions**

354 It is a great challenge to realistically simulate low clouds and associated warm rain in climate  
355 models without reliable vertical variations of microphysical processes. In this work, we use the  
356 newly retrieved cloud and drizzle microphysical properties to constrain the autoconversion and  
357 accretion parameterizations in a widely used microphysical scheme, and then implement the  
358 updated scheme into the NCAR CESM to examine the responses of warm rain frequency and cloud  
359 properties. Climate simulations with the updated cloud microphysical scheme exhibit the reduced  
360 precipitation frequency and increased precipitation intensity, indicating that the new scheme has  
361 the potential of mitigating the outstanding problem in GCM precipitation simulations and  
362 achieving more accurate climate assessments.

363 The findings from this study attest the paramount importance of cloud microphysics  
364 parameterizations in GCM simulations. In particular, we show that it is critical to take the in-cloud  
365 vertical variations of warm rain processes into account when developing cloud microphysical  
366 schemes. We note that the robustness of our findings is subject to the representative of new  
367 parameterizations derived from a field campaign. Therefore, it is imperative to use more ground-

368 based observations from different field campaigns and ARM permanent sites as well as a single  
369 column modeling framework to test if these new parameterizations are valid over other oceans and  
370 land surfaces. Future study will also focus on how altered warm rain processes can influence the  
371 aerosol indirect effect, cloud feedback, and climate sensitivity.

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388 **References**

- 389 Albrecht, B. A. Aerosols, Cloud Microphysics, and Fractional Cloudiness. *Science*. **245** (4923),  
390 1227–1230 (1989).
- 391 Aumann, H. H., Behrangi, A. & Wang, Y. Increased Frequency of Extreme Tropical Deep  
392 Convection: AIRS Observations and Climate Model Predictions. *Geophys. Res. Lett.* **45**,  
393 13,530–13,537 (2018).
- 394 Beheng, K. D. A parameterization of warm cloud microphysical conversion processes. *Atmos.*  
395 *Res.* **33**, 193–206 (1994).
- 396 Bony, S. & Dufresne, J. L. Marine boundary layer clouds at the heart of tropical cloud feedback  
397 uncertainties in climate models. *Geophys. Res. Lett.* **32**, 1–4 (2005).
- 398 Cheng, A. & Xu, K.-M. A PDF-Based Microphysics Parameterization for Simulation of  
399 Drizzling Boundary Layer Clouds. *J. Atmos. Sci.* **66**, 2317–2334 (2009).
- 400 Dong, X., Xi, B., Kennedy, A., Minnis, P. & Wood, R. A 19-Month Record of Marine Aerosol–  
401 Cloud–Radiation Properties Derived from DOE ARM Mobile Facility Deployment at the  
402 Azores. Part I: Cloud Fraction and Single-Layered MBL Cloud Properties. *J. Clim.* **27**,  
403 3665–3682 (2014a).
- 404 Dong, X. *et al.* Investigation of the Diurnal Variation of Marine Boundary Layer Cloud  
405 Microphysical Properties at the Azores. *J. Clim.* **27**, 8827–8835 (2014b).
- 406 Donner, L. J. *et al.* The dynamical core, physical parameterizations, and basic simulation  
407 characteristics of the atmospheric component AM3 of the GFDL global coupled model  
408 CM3. *J. Clim.* **24**, 3484–3519 (2011).
- 409 Fan, J., Y. Wang, D. Rosenfeld, X. Liu: Review of Aerosol-Cloud Interactions: Mechanisms,  
410 Significance and Challenges: *J. Atmo. Sci.* 73(11), 4221–4252 (2016).
- 411 Gettelman, A. *et al.* High climate sensitivity in the Community Earth System Model Version 2  
412 (CESM2), *Geophys. Res. Lett.*, **46**, 8329–8337 (2019).
- 413 Golaz, J.-C., Larson, V. E. & Cotton, W. R. A PDF-Based Model for Boundary Layer Clouds.  
414 Part II: Model Results. *J. Atmos. Sci.* **59**, 3552–3571 (2002).
- 415 Ghan, S. *et al.*: Challenges in constraining anthropogenic aerosol effects on cloud radiative  
416 forcing using present-day spatiotemporal variability. *Proc. Natl Acad. Sci. USA*. 113 (21):  
417 5804–5811 (2016).
- 418 Hang, Y., L’Ecuyer, T. S., Henderson, D. S., Matus, A. V. & Wang, Z. Reassessing the effect of  
419 cloud type on earth’s energy balance in the age of active spaceborne observations. Part II:  
420 Atmospheric heating. *J. Clim.* **32**, 6219–6236 (2019).
- 421 Hurrell, J. W. *et al.* The Community Earth System Model: A Framework for Collaborative  
422 Research. *Bull. Am. Meteorol. Soc.* **94**, 1339–1360 (2013).
- 423 IPCC. *Climate Change 2013: The Physical Science Basis. Contribution of Working Group I to*  
424 *the Fifth Assessment Report of the Intergovernmental Panel on Climate Change [Stocker,*  
425 *T.F., D. Qin, G.-K. Plattner, M. Tignor, S.K. Allen, J. Boschung, A. Nauels, Y. Xia,*  
426 *(Cambridge University Press, Cambridge, United Kingdom and New York, NY, USA,*  
427 *2013). doi:10.1017/CBO9781107415324.*
- 428 Jing, X. *et al.* A Multimodel Study on Warm Precipitation Biases in Global Models Compared to  
429 Satellite Observations. *J. Geophys. Res. Atmos.* **122**, 11,806–11,824 (2017).
- 430 Jing, X. & Suzuki, K. The Impact of Process-Based Warm Rain Constraints on the Aerosol  
431 Indirect Effect. *Geophys. Res. Lett.* **45**, 10,729–10,737 (2018).
- 432 Kay, J. E. *et al.* Scale-Aware and Definition-Aware Evaluation of Modeled Near-Surface  
433 Precipitation Frequency Using CloudSat Observations. *J. Geophys. Res.*, **123**, 4294–4309

434 (2018).

435 Kessler, E. On the Distribution and Continuity of Water Substance in Atmospheric Circulations.  
436 in *On the Distribution and Continuity of Water Substance in Atmospheric Circulations* 1–  
437 84 (American Meteorological Society, 1969). doi:10.1007/978-1-935704-36-2\_1

438 Khairoutdinov, M. & Kogan, Y. A New Cloud Physics Parameterization in a Large-Eddy  
439 Simulation Model of Marine Stratocumulus. *Mon. Weather Rev.* **128**, 229–243 (2000).

440 Lau, K. M., & Wu H. T. Climatology and changes in tropical oceanic rainfall characteristics  
441 inferred from Tropical Rainfall Measuring Mission (TRMM) data (1998-2009). *J.*  
442 *Geophys. Res.* 116 D1711, doi:1029/2011D015827 (2011).

443 Lebsock, M., Morrison, H. & Gettelman, A. Microphysical implications of cloud-precipitation  
444 covariance derived from satellite remote sensing. *J. Geophys. Res.* **118**, 6521–6533  
445 (2013).

446 L’Ecuyer, T. S., Hang, Y., Matus, A. V. & Wang, Z. Reassessing the effect of cloud type on  
447 earth’s energy balance in the age of active spaceborne observations. Part I: Top of  
448 atmosphere and surface. *J. Clim.* **32**, 6197–6217 (2019).

449 Li, Z. et al. East Asian Study of Tropospheric Aerosols and Impact on Regional Cloud,  
450 Precipitation, and Climate (EAST-AIRCPC). *J. Geophys. Res.* **124**, 2019JD030758 (2020).

451 Liu, Y., Daum, P. H., McGraw, R. L., Miller, M. A. & Niu, S. Theoretical expression for the  
452 autoconversion rate of the cloud droplet number concentration. *Geophys. Res. Lett.* **34**,  
453 (2007).

454 Liu, Y. & Daum, P. H. Parameterization of the Autoconversion Process. Part I: Analytical  
455 Formulation of the Kessler-Type Parameterizations. *J. Atmos. Sci.* **61**, 1539–1548 (2004).

456 Ma, X., Jia, H., Yu, F. & Quaas, J. Opposite Aerosol Index-Cloud Droplet Effective Radius  
457 Correlations Over Major Industrial Regions and Their Adjacent Oceans. *Geophys. Res.*  
458 *Lett.* **45**, 5771–5778 (2018).

459 Morrison, H. & Gettelman, A. A New Two-Moment Bulk Stratiform Cloud Microphysics  
460 Scheme in the Community Atmosphere Model, Version 3 (CAM3). Part I: Description  
461 and Numerical Tests. *J. Clim.* **21**, 3642–3659 (2008).

462 Minnis, P., et al. (2020). CERES MODIS cloud product retrievals for Edition 4 – Part I: Algorithm  
463 changes, *IEEE Trans. Geosci. Remote Sens.*, 58, doi:10.1109/TGRS.2020.3008866.

464 Nakajima, T. Y. et al. Droplet Growth in Warm Water Clouds Observed by the A-Train. Part I:  
465 Sensitivity Analysis of the MODIS-Derived Cloud Droplet Sizes. *J. Atmos. Sci.* **67**, 1884–  
466 1896 (2010).

467 Penner, J.E., X. Dong, and Y. Chen: Observational evidence for a change in radiative forcing due  
468 to the indirect aerosol effect. *Nature*, **427**, 231-234 (2004).

469 Rogers, R. R. R. & Yau, M. K. Short Course in Cloud Physics. 295 (1996).

470 Seinfeld, J.H. et al.: Improving our fundamental understanding of the role of aerosol–cloud  
471 interactions in the climate system. *Proc. Natl Acad. Sci. USA*. 113 (21): 5781-5790 (2016).

472 Stephens, G. L. et al. The albedo of earth. *Reviews of Geophysics* **53**, 141–163 (2015).

473 Stephens, G. L. Cloud Feedbacks in the Climate System: A Critical Review. *J. Clim.* **18**, 237–  
474 273 (2005).

475 Stephens, G. L. et al. Dreary state of precipitation in global models. *J. Geophys. Res.* **115**,  
476 (2010).

477 Schmidt, G. A. et al. Present-Day Atmospheric Simulations Using GISS ModelE: Comparison to  
478 In Situ, Satellite, and Reanalysis Data. *J. Clim.* **19**, 153–192 (2006).

479 Soden, B. J. & Vecchi, G. A. The vertical distribution of cloud feedback in coupled ocean-

480 atmosphere models. *Geophys. Res. Lett.* **38**, L12704 (2011).

481 Seinfeld, J.H. et al.: Improving our fundamental understanding of the role of aerosol–  
482 cloud interactions in the climate system. *Proc. Natl Acad. Sci. USA.* **113** (21): 5781–5790  
483 (2016).

484 Suzuki, K., Stephens, G. L. & Lebsock, M. D. Aerosol effect on the warm rain formation  
485 process: Satellite observations and modeling. *J. Geophys. Res.* **118**, 170–184 (2013).

486 Suzuki, K. *et al.* Evaluation of the Warm Rain Formation Process in Global Models with  
487 Satellite Observations. *J. Atmos. Sci.* **72**, 3996–4014 (2015).

488 Suzuki, K., Nakajima, T., Nakajima, T. Y. & Stephens, G. L. Effect of the droplet activation  
489 process on microphysical properties of warm clouds. *Environ. Res. Lett.* **5**, 024012 (2010).

490 Takahashi, H., Lebsock, M., Suzuki, K., Stephens, G. & Wang, M. An investigation of  
491 microphysics and subgrid-scale variability in warm-rain clouds using the A-train  
492 observations and a multiscale modeling framework. *J. Geophys. Res.* **122**, 7493–7504  
493 (2017).

494 Tripoli, G. J., Cotton, W. R., Tripoli, G. J. & Cotton, W. R. A Numerical Investigation of Several  
495 Factors Contributing to the Observed Variable Intensity of Deep Convection over South  
496 Florida. *J. Appl. Meteorol.* **19**, 1037–1063 (1980).

497 Wallace, J. M. & Hobbs, P. V. *Atmospheric Science: An Introductory Survey: Second Edition.*  
498 *Atmospheric Science: An Introductory Survey: Second Edition* (Elsevier Inc., 2006).  
499 doi:10.1016/C2009-0-00034-8

500 Wang, Y., Ma, P.-L., Jiang, J. H., Su, H. & Rasch, P. J. Toward reconciling the influence of  
501 atmospheric aerosols and greenhouse gases on light precipitation changes in Eastern  
502 China. *J. Geophys. Res. Atmos.* **121**, 5878–5887 (2016).

503 Wood, R. Drizzle in Stratiform Boundary Layer Clouds. Part I: Vertical and Horizontal  
504 Structure. *J. Atmos. Sci.* **62**, 3011–3033 (2005a).

505 Wood, R. Drizzle in Stratiform Boundary Layer Clouds. Part II: Microphysical Aspects. *J.*  
506 *Atmos. Sci.* **62**, 3034–3050 (2005b).

507 Wood, R. & Hartmann, D. L. Spatial variability of liquid water path in marine low cloud: The  
508 importance of mesoscale cellular convection. *J. Clim.* **19**, 1748–1764 (2006).

509 Wood, R., Kubar, T. L. & Hartmann, D. L. Understanding the Importance of Microphysics and  
510 Macrophysics for Warm Rain in Marine Low Clouds. Part II: Heuristic Models of Rain  
511 Formation. *J. Atmos. Sci.* **66**, 2973–2990 (2009).

512 Wood, R. Stratocumulus Clouds. *Mon. Weather Rev.* **140**, 2373–2423 (2012).

513 Wu, P., Dong, X. & Xi, B. Marine boundary layer drizzle properties and their impact on cloud  
514 property retrieval. *Atmos. Meas. Tech.* **8**, 3555–3562 (2015).

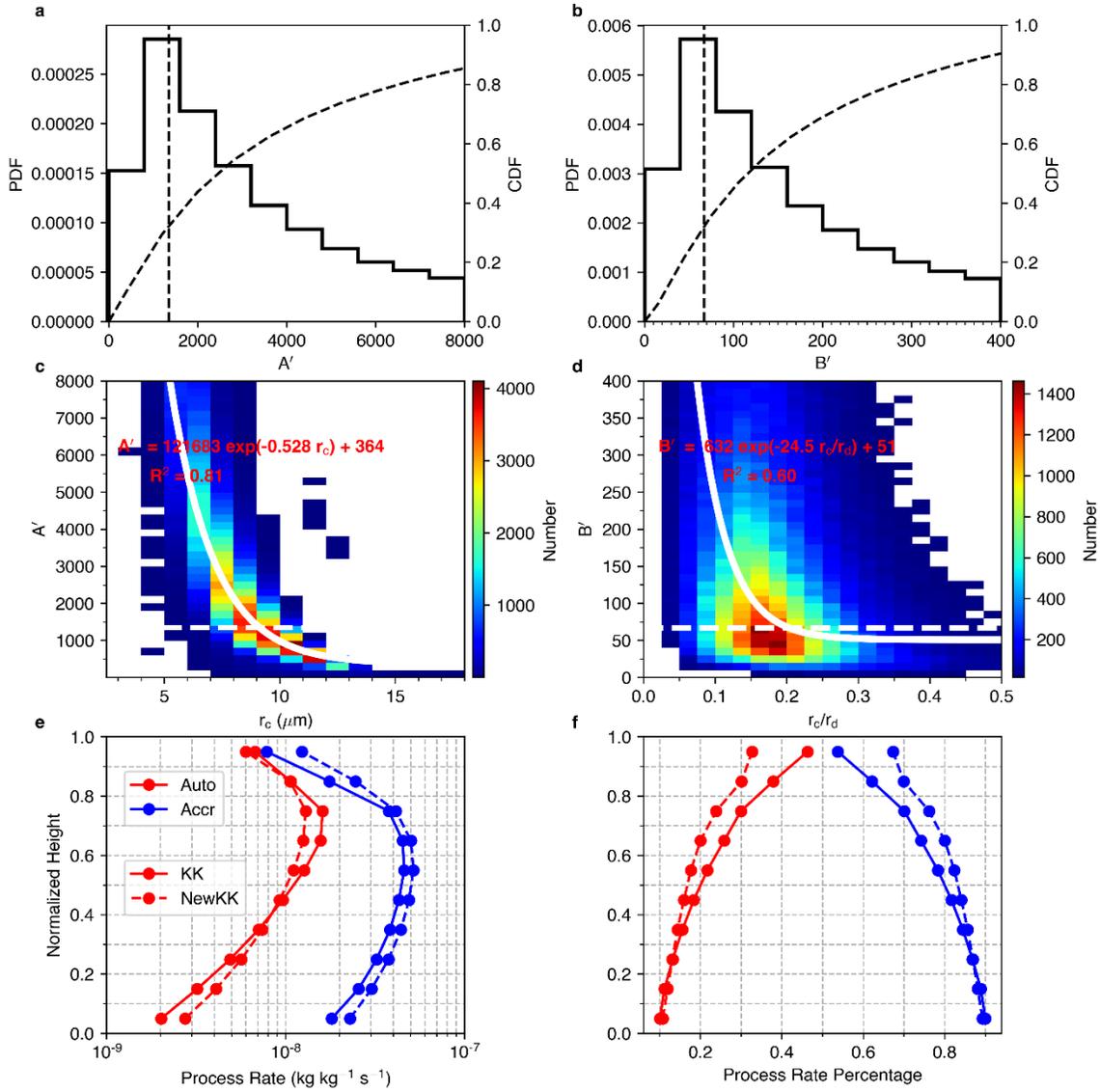
515 Wu, P. *et al.* Effects of environment forcing on marine boundary layer cloud-drizzle processes. *J.*  
516 *Geophys. Res.* **122**, 4463–4478 (2017).

517 Wu, P., Xi, B., Dong, X. & Zhang, Z. Evaluation of autoconversion and accretion enhancement  
518 factors in general circulation model warm-rain parameterizations using ground-based  
519 measurements over the Azores. *Atmos. Chem. Phys.* **18**, 17405–17420 (2018).

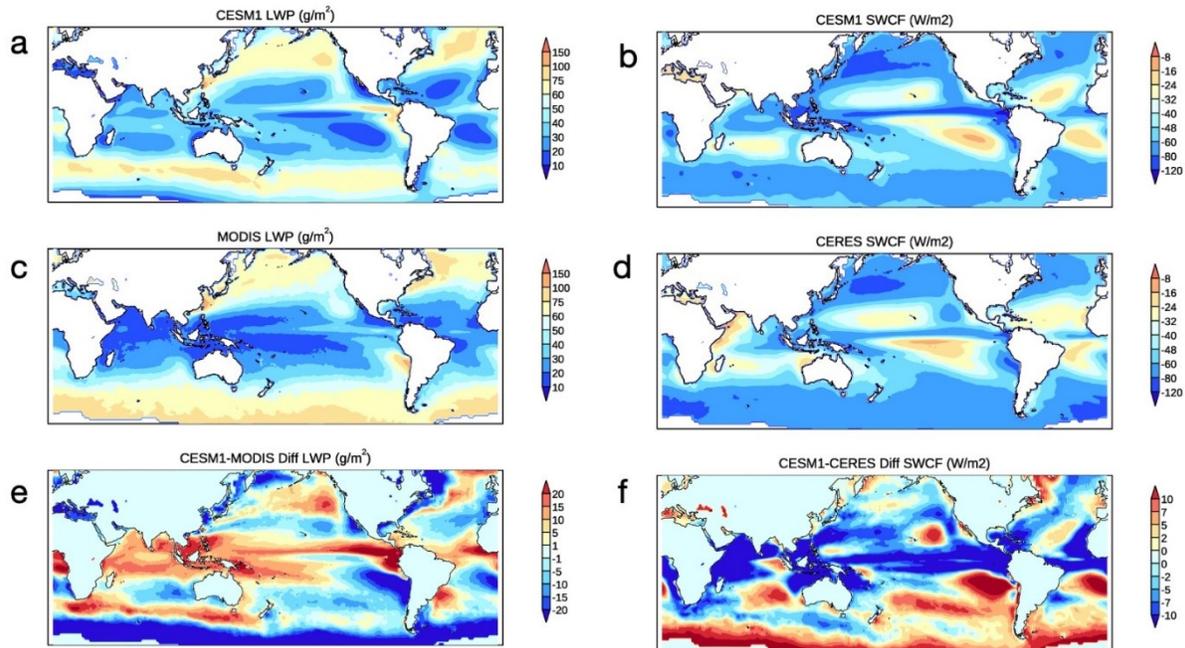
520 Wu, P., Dong, X., Xi, B., Tian, J. & Ward, D. M. Profiles of MBL Cloud and Drizzle  
521 Microphysical Properties Retrieved From Ground-Based Observations and Validated by  
522 Aircraft In Situ Measurements Over the Azores. *J. Geophys. Res.* **125**, (2020).

523 Xie, X., & Zhang, M. Scale-aware parameterization of liquid cloud inhomogeneity and its  
524 impact on simulated climate in CESM. *Journal of Geophysical Research: Atmospheres*,  
525 *120*(16), 8359–8371. <https://doi.org/10.1002/2015JD023565> (2015).

526 Yamaguchi, T., Feingold, G. & Kazil, J. Stratocumulus to Cumulus Transition by Drizzle. *J.*  
527 *Adv. Model. Earth Syst.* **9**, 2333–2349 (2017).  
528 Zhang, Z. *et al.* Subgrid variations of the cloud water and droplet number concentration over the  
529 tropical ocean: Satellite observations and implications for warm rain simulations in  
530 climate models. *Atmos. Chem. Phys.* **19**, 1077–1096 (2019).  
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562 **Fig. 1.** (a,b) Probability density functions (PDFs, solid lines) and cumulative density functions  
563 (CDFs, dashed curves) of coefficients  $A'(Z)$  and  $B'(Z)$  calculated from ground-based retrievals.  
564 Vertical dashed lines mark the constants in the KK scheme. Joint histograms of (c)  $A'(Z)$  with  $r_c(Z)$   
565 and (d)  $B'(Z)$  and the ratio of  $r_c(Z)$  to  $r_d(Z)$ . White solid lines are the exponential fittings. White  
566 dashed lines mark the prescribed  $A$  and  $B$  in the KK scheme. (e) Normalized profiles of  $R_{auto}$  (red  
567 lines) and  $R_{accr}$  (blue lines) for all the drizzle cases during ACE-ENA (a total 9,213 1-min profiles)  
568 by cloud thickness ( $z_i = \frac{z-z_b}{z_t-z_b}$ , where subscripts  $b$  and  $t$  denote cloud base and top, respectively).  
569 The solid and dashed lines represent the profiles from KK and NKK schemes, respectively. (f) The  
570 percentages of total drizzle water production rate ( $R_{auto}+R_{accr}$ ) contributed by  $R_{auto}$  and  $R_{accr}$ .

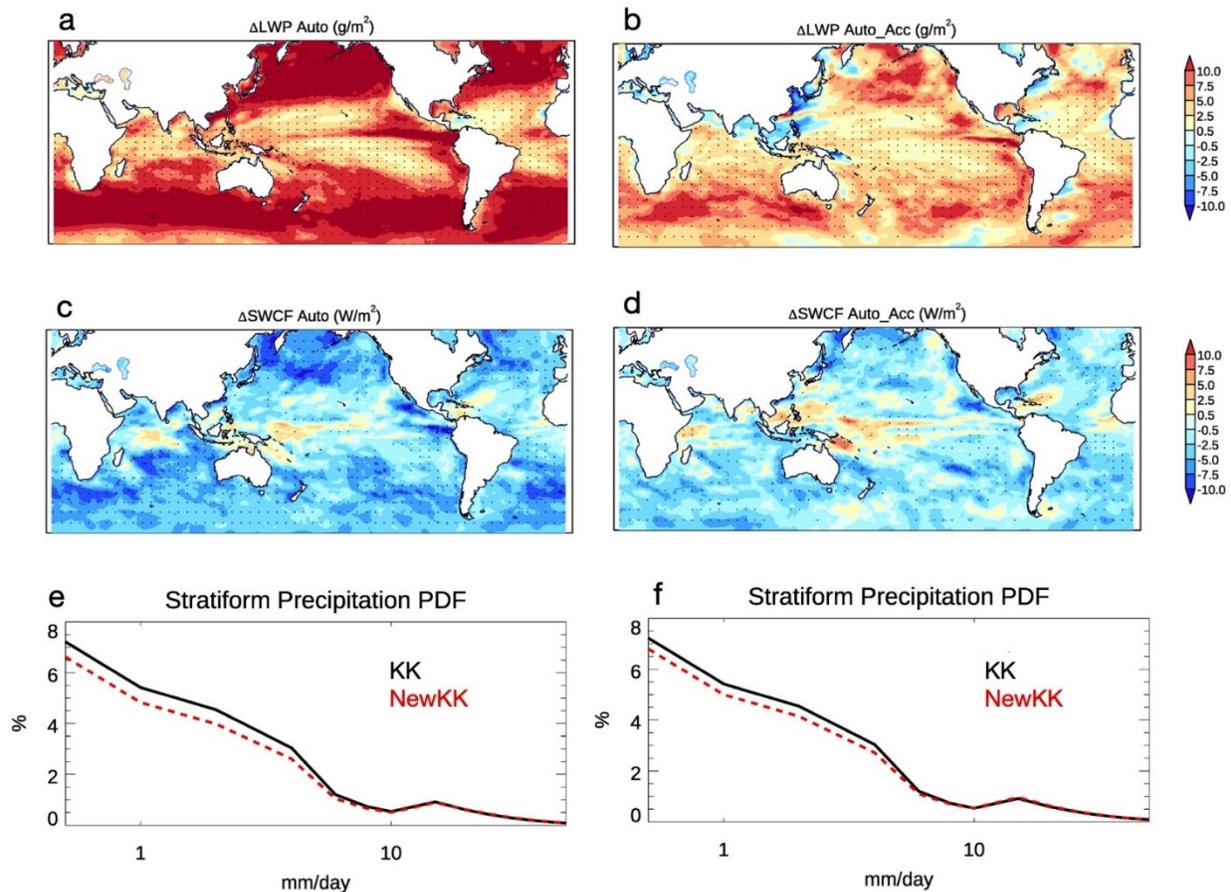


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572 **Fig. 2.** Comparisons of cloud liquid water path (LWP, left column) and shortwave cloud forcing  
 573 (SWCF, right column) between CESM1 present-day scenario simulations (a,b) and CERES-  
 574 MODIS satellite cloud and radiation climatologies (c,d), as well as their differences (e,f). Model  
 575 simulations consist of five ensemble members. Satellite data are averaged over 2001-2019 from  
 576 both Terra and Aqua satellites.

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581 **Fig. 3.** Changes in LWP (a,b), SWCF (c,d), and Probability distribution functions (PDFs) of large-  
 582 scale stratiform precipitation (e,f) in CESM simulations by different warm rain schemes (NKK –  
 583 KK). Left column: autoconversion only. Right column: both autoconversion and accretion. The  
 584 stippling indicates the statistically significant changes that are larger than the model internal  
 585 variability (calculated as the standard deviation among the ensemble members). The precipitation  
 586 PDF are averaged over 60°S to 60°N oceanic regions. The spreads of precipitation frequency in  
 587 each bin are all less than 0.1% among different ensemble members, so they are too small to be  
 588 shown in the panels e and f.

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