1	Noble gas incorporation into silicate glasses:
2	implications for planetary volatile storage
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13	Abstract
14	Incorporation of small molecules in silicate melts may provide an important mechanism for
15	storing noble gases in the deep Earth, yet the means by which chemically inert noble gases enter
16	and are retained in silica-based materials is not understood. High-pressure, room-temperature
17	sound velocity measurements on silica and natural-basalt glasses in different pressure-
18	transmitting media reveal that neon enters the structure of silicate glasses and enhances their
19	elastic strengths, whereas ethanol-methanol mixture does not. Combined with literature data, we
20	found the incorporation of small molecules into silica and basalt glasses is controlled by the void
21	size distribution of the glass and size of the molecules. Pressure primarily reduces the size of
22	noble gases, thereby increasing their solubilities in silicate melts and glasses.
23	

24 Introduction

25

26	Radiogenic heat production generates <sup>40</sup> Ar, <sup>21</sup> Ne and <sup>4</sup> He inside Earth, and the ratios of these
27	isotopes to non-radiogenic isotopes have been used to infer the style of mantle convection and
28	the source of ocean island basalts (e.g. Mukhopadhyay and Parai, 2019). Owing to their changing
29	reactivity and volatility with pressure, noble gases are also useful geochemical tracers for interior
30	processes of planets (e.g. Sanloup et al., 2005). However, how these noble gases are distributed
31	among potential geochemical reservoirs, and how they alter the physical properties of their host
32	with increasing depth (and therefore pressure) is still unclear. The storage of noble gases in
33	quartz, ferropericlase, and bridgmanite at high pressure has been experimentally verified (Rosa et
34	al., 2020; Sanloup et al., 2005). Nevertheless, the partition coefficients of noble gases between
35	minerals and melts are on the order of 10 <sup>-3</sup> (Karato, 2016), implying significant storage of noble
36	gases in silicate melts. This deep storage could also potentially alter atmospheric composition.
37	As silicate glasses and melts share structural similarities (Morard et al., 2020; Williams and
38	Jeanloz, 1988), with glass being the kinetically hindered state of the corresponding melt,
39	understanding the incorporation of noble gases into silicate glasses can shed light on their storage
40	in natural silicate melts.
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42 Noble gases are widely used as pressure-transmitting media in high-pressure diamond-anvil cell 43 experiments. These gases are chemically inactive and display relatively low mechanical strength, 44 and thus minimize pressure gradients and deviatoric stresses in the sample chamber (Angel *et al.*, 45 2007; Klotz *et al.*, 2009). Use of noble gases as pressure-transmitting media presumes minimal 46 interaction with the pressurized sample, yet there have been several reports that helium

penetrates into the structure of silica glass at room temperature. (Sato *et al.*, 2011; Shen *et al.*,
2011; Weigel *et al.*, 2012). These results indicate that helium gets into the structure of silica
glass, enhancing both its incompressibility and rigidity. Another study on basalt and enstatite

- 50 glasses also indicates neon can enter their structure at high pressure (Clark *et al.*, 2016).
- 51

52 Void-space analysis of silica could shed light on the incorporation of noble gases into its 53 structure. Theoretical simulations of the structure of silica glass and void-size analysis have 54 provided statistics on the interstitial space (i.e., the largest spherical site not occupied by Si or O) 55 that could potentially be available for incorporating noble gases (Malavasi et al., 2006; 56 Shackelford and Masaryk, 1978). Assuming the distribution of the interstitial sites follows a log-57 normal distribution, Shackelford and co-workers found the highest density of interstitial site 58 diameters at d = 1.81 Å, with the distribution extending to around d = 4.0 Å. Based on this result, 59 the four smaller-sized noble gases, He (d = 2.551 Å), Ne (d = 2.820 Å), Ar (d = 3.542 Å) and Kr 60 (d = 3.655 Å), can occupy some fraction of the interstitial sites in silica glass, and potentially 61 modify the structure and physical properties. Another theoretical structure simulation based on 62 molecular dynamics found an overall smaller void size, suggesting limited incorporation of Ar 63 and larger-sized molecules. However, systematic study of the high-pressure solubility of these 64 gases in silica glass has been lacking, despite a number of studies on helium (Sato et al., 2011; 65 Shen et al., 2011; Weigel et al., 2012).

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To help clarify the mechanism of noble-gas incorporation into amorphous silica and natural
silicate glasses, we measured high-pressure Brillouin spectra of silica and basalt glasses using
different pressure transmitting media at room temperature. The measured elastic properties of the

70	material provide insight into the structural evolution and molecule incorporation of each glass
71	with compression. Elasticity is a useful monitor of the solution process because in situ
72	measurement of gas solubility is challenging at high pressure. Together with existing literature
73	data, we provide a comprehensive review of noble gases migrating into silica glass and natural
74	basalt glasses under pressure. We find that solubility is controlled by the atomic sizes of the
75	noble gases relative to the size of available interstitial spaces in the silicate glasses. Pressure
76	alters both factors, thereby affecting the solubility of noble gases and other volatile species in
77	glasses and melts. Our results support the potential for increased storage of noble gases in
78	Earth's and other deep planetary interiors.
79	
80	Results
81	Details on sample synthesis, compositions, and data collection can be found in the
82	Supplementary Information. We found that the sound velocities of silica glass depend on the
82 83	Supplementary Information. We found that the sound velocities of silica glass depend on the pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For
83	pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For
83 84	pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For both the compressional and shear velocities of silica glass, we observed a drop in velocity when
83 84 85	pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For both the compressional and shear velocities of silica glass, we observed a drop in velocity when increasing pressure between 1 and 3 GPa, followed by a slightly increasing or nearly unchanged
83 84 85 86	pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For both the compressional and shear velocities of silica glass, we observed a drop in velocity when increasing pressure between 1 and 3 GPa, followed by a slightly increasing or nearly unchanged velocity at higher pressures (Fig. 1). Silicate glasses with natural compositions share a similar
83 84 85 86 87	pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For both the compressional and shear velocities of silica glass, we observed a drop in velocity when increasing pressure between 1 and 3 GPa, followed by a slightly increasing or nearly unchanged velocity at higher pressures (Fig. 1). Silicate glasses with natural compositions share a similar framework structure with silica glass, and the velocity drop upon initial compression was also
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83 84 85 86 87 88 89	pressure-transmitting medium, as measured at pressures ranging from 0-10 GPa (Fig. 1). For both the compressional and shear velocities of silica glass, we observed a drop in velocity when increasing pressure between 1 and 3 GPa, followed by a slightly increasing or nearly unchanged velocity at higher pressures (Fig. 1). Silicate glasses with natural compositions share a similar framework structure with silica glass, and the velocity drop upon initial compression was also documented in other polymerized silicate glasses, such as basalt, jadeite and albite glasses (Liu and Lin, 2014; Sakamaki <i>et al.</i> , 2014). In contrast, depolymerized glasses like diopside or

3 GPa, the SiO<sub>4</sub> tetrahedra in the glass rotate into the void space to form a high-density structure
(Clark *et al.*, 2016). The rotation does not involve substantial compression of the interatomic
bonds, so the elastic moduli of the material remain largely unaltered. Therefore, the velocities,
given by the square root of the ratio of the moduli and density, decreases during this stage.

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98 However, after the void space is filled, tetrahedral rotation is replaced by the interatomic bonds 99 shortening, and the sound velocities then increase under compression (Clark et al., 2016). Depolymerized silicate glasses, which contain larger 'modifier' cations like Mg<sup>2+</sup>, Na<sup>2+</sup> or Ca<sup>2+</sup>, 100 101 have less void space and consequently less flexibility. The densification may also involve some 102 chemical bond shortening and leads to the unchanged velocity profile with increasing pressure. 103 This explanation is supported by molecular dynamics (MD) simulations on silicate glasses and 104 melts over a wide compositional range (Guillot and Sator, 2007; Salmon et al., 2019), with a 105 decrease in Si-O-Si angle documented between 0-5 GPa.

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107 For basalt glasses, sound velocities at high pressure are also influenced by the pressure media 108 (Fig. 2). The BIR-1 sample has higher velocities in a neon medium, as compared with basalt in 109 M-E and M-E-W (methanol: ethanol: water = 16:3:1); it also has an earlier transition pressure at 110 which the velocities start to increase. Comparing the M-E and M-E-W cases, we see that water 111 seems to lower the decreasing slope below 5 GPa, but it does not change the transition point for the change in velocity trends (Fig. 2). In fact, water enhancing glass incompressibility has been 112 reported (Murakami, 2018) and the velocity difference between the two cases here could be due 113 114 to water penetration (Fig. 3).

- 116 **Discussion**
- 117 Sound velocity data for amorphous materials can be very useful to calculate its density at high
- 118 pressure (Zha *et al.*, 1994). However, this method would fail if pressure media penetrates the
- 119 sample (Weigel *et al.*, 2012). We calculated the *P-V* curve from velocities observed in different
- 120 media and used this to examine whether incorporation of pressure media occurred in our
- 121 experiments (See Supplementary Information for details). Neon and helium are able to penetrate
- 122 into the silica structure while water and methanol molecules seem to also be able to penetrate
- 123 into basalt glasses (Fig. S-3).
- 124

125 We compared the molecular size of the pressure media with the size of the interstitial space in 126 the silica structure (Fig. 3 and S-4). It is found that at ambient conditions, the sizes of helium, 127 neon and water molecules are smaller than some voids in the silica structure. At higher pressures, 128 the void size distribution generally shifts to smaller volume, but the peak position only slightly 129 moves and is still larger than 1.75 Å. On the other hand, the sizes of highly compressible gases 130 decrease dramatically with increasing pressure, especially for helium and neon (Fig. S-4). At 5.7 131 GPa and  $\frac{40}{00}$  K, almost half of the interstitial sites are open to helium and around 20% are open 132 to neon. These results indicate that helium and neon elevate the elastic stiffness of silica by 133 supporting the structure in the void space, while molecules larger than argon are too big to be 134 incorporated into silica and did not show this effect. Pressure makes these atoms smaller, 135 enhancing the solubility of neon and helium into silica (Fig. 3). Although argon also becomes 136 smaller at high pressure, it is still larger than most of the voids in silica. Its solubility is limited 137 and does not influence the elastic properties significantly (Fig. 1).

138

139 Our measurements do not provide solubility values, but by comparing the gas and non-gas 140 experiments we can make an estimate of this number (Sato *et al.*, 2011) (Fig. S-3). In Figure 4, 141 the upper limit is constrained by the maximum available space in the silica structure. This space 142 is calculated as the volume difference between normal silica and the ultra-dense six-coordinated 143 silica extropolated to lower pressure. On the other hand, the lower limit is given by the difference 144 between gas and non-gas curves, assuming the 'expansion' should wholy or partly come from 145 the volume of gas in the structure. Since the partial volume of a component in a mixture is 146 smaller than the volume of its own existence (Bajgain et al., 2015), the real solubility should be 147 higher than the lower limit here.

148

149 Basalt glass is compositionally more complex than silica glass with the addition of other cations. 150 These cations can be classified into two categories: the network formers like Ti, Al and network 151 modifiers like Mg, Ca, Na and K. The two sets of cations have distinct effects on gas solubility. 152 Network modifiers tend to form bonds between the bridging SiO<sub>4</sub> tetrahedra and lower the 153 volume of void space. Their negative correlation with noble gas solubility has been 154 experimentally observed (Tournour and Shelby, 2008a; Tournour and Shelby, 2008b). On the 155 other hand, network formers, which reside in the Si site, seem to have less of an influence on gas 156 solubility.

157

## 158 Geochemical Implications

159

160 The geometrical packing and coordination of atoms in silicate melts and glasses are similar at 161 ambient and high pressure conditions based on experimental observations (Morard *et al.*, 2020;

162 Williams and Jeanloz, 1988). Hence, the void space distribution in melt structure is likely to be 163 comparable and our results here support significant solubility of helium and neon in high 164 pressure silicate melts (Fig. 3). Furthermore, the partition coefficients between mineral and melts 165 for noble gases are in the order of  $10^{-3}$  (Karato, 2016). It is expected silicate melts should be an important host for noble gases. He and Ne are the 2<sup>nd</sup> and 5<sup>th</sup> most abundant element in the solar 166 system (Palme et al., 2014). For <sup>3</sup>He/<sup>4</sup>He and <sup>20,21</sup>Ne/<sup>22</sup>Ne isotopic ratios, the discrepancy of 167 168 upper mantle material value from atmospheric value has been a hot topic in geochemistry (e.g. 169 Bekaert et al., 2019; Mukhopadhyay and Parai, 2019). Most answers to this question require a 170 deep primordial reservoir which has unique geochemical features and survive mantle convection 171 for the last 4.5 billion years. It has been noticed that some patches making of partially molten 172 rock might exist at the core mantle boundary (e.g. Wen *et al.*, 2001). These melts may be able to 173 host large amounts of noble gases like helium and neon with primordial and less-radiogenic 174 features. Other than the core-mantle boundary, partial melting may also occur at the top of the 175 lower mantle due to dehydration melting (Fu et al., 2019). These layers might be perturbed by 176 mantle convection more often and host noble gases with more radiogenic features. Therefore, the 177 observed difference in noble gases ratios in OIBs and MORBs could be possibly due to sampling 178 different melt reservoirs for noble gases. The storage of helium or neon discussed here reaches 179 conditions beyond the range of this experiment, and due to the complex coordination 180 environment change of silicon at higher pressures (Wang et al., 2014), and the high temperature 181 conditions in deep Earth, directly applying our results to these conditions may not be suitable. 182 However, the mechanism revealed in this study and previous studies (Clark et al., 2016; Sato et 183 *al.*, 2011; Weigel *et al.*, 2012) (i.e., availability of interstitial void to compressed noble gases) is 184 still valid and future structure simulation and void space analysis of heated silica/basalt glass at

high pressures is needed. Geodynamic simulations are also needed to better estimate the degreeof mixing during these processes.

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188	Noble gas-silicate interaction may also have important implications for the composition of the
189	atmospheres of other planetary bodies like Jupiter. It is found that the abundances of helium and
190	neon in Jupiter's atmosphere are significantly lower than other noble gases, when compared to
191	solar composition (Fortney, 2010). Our results suggest that this discrepancy could be related to
192	interior processes in the planet. Forming a He-Ne-silicate composite at Jupiter's rocky core could
193	be a viable option. Whether such a mechanism could explain the deficit of neon and helium in
194	Jupiter's atmosphere requires further experimental and computational work. Our study
195	demonstrates the controlling factors for noble gas solubility in a silicate melt are the noble gas
196	size compared to void size, indicating that data on the structure of silicate melts with natural
197	compositions at higher <i>P</i> - <i>T</i> is crucially needed in order to estimate the storage capacity of noble
198	gases in deep planetary interiors.
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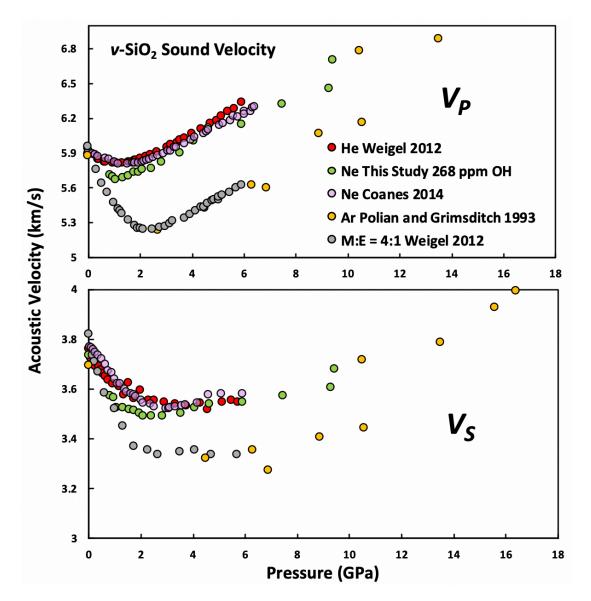
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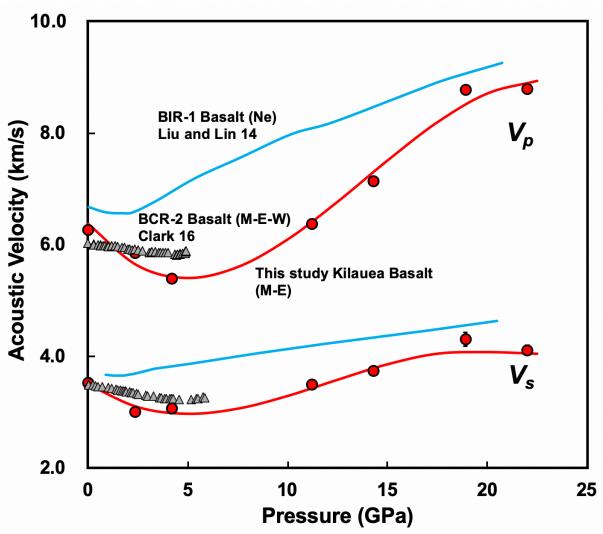


## 287 288

Figure 1. Sound velocities of vitreous silica under high pressure in different pressure

**media.** M-E represents 4:1 Methanol-Ethanol mixture. Errors of the velocities are estimated from statistical uncertainties arising from the peak fitting. Typical errors are less than 1.5 % and smaller than the size of the symbols. For both the  $V_P$  and  $V_S$  of silica in different noble gas media, we found a consistent trend for the acoustic velocities—He>Ne>Ar $\approx$ M-E. The abnormal velocity minimum at around 2-5 GPa can be attributed to the rearrangement of SiO<sub>4</sub> tetrahedra in the vitreous silica structure (Clark et al., 2016).

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- 297





301 Figure 2. Sound velocities of basalt glasses at high pressure. M-E-W: 16:3:1 Methanol-

Ethanol-Water mixture, M-E: 4:1 Methanol-Ethanol mixture. Errors of the velocities are 302 303 estimated from statistical uncertainties arising from the peak fitting. Error bars smaller than the

304 symbols plotted are not shown. Although the three basalt glasses have slightly different

305 composition, their degrees of polymerization (NBO/T= 0.6, 0.9 and 0.8 for BCR-2, BIR-1 and

- 306 KB, respectively) are quite similar (BIR-1, blue colored line, Liu and Lin, 2014; BCR-2,
- 307 triangular points, Clark et al., 2016). The extents of the velocity drops are very different among
- 308 the glasses. We observed that BIR-1 has a 2 % drop for both  $V_P$  and  $V_S$ , BCR-2 has a 2.8 % drop
- 309 for  $V_P$  and a 7.2 % drop for  $V_S$ , while KB in M-E has a 14 % drop for both  $V_P$  and  $V_S$ . We 310 attribute this variation to be mostly due to the different pressure media used. H<sub>2</sub>O and Ne have
- 311 relatively small molecule size that can possibly penetrate into the structure of silicate glass and
- 312 make it stiffer (Fig. 3 and S-4).
- 313

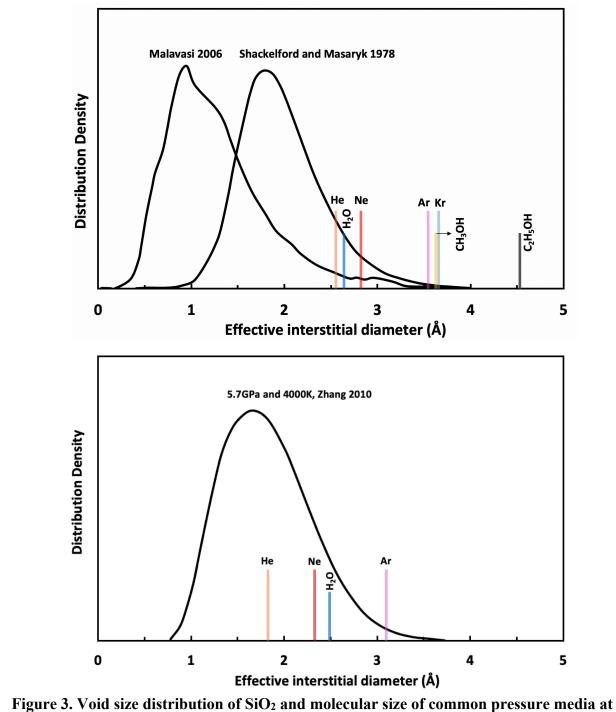
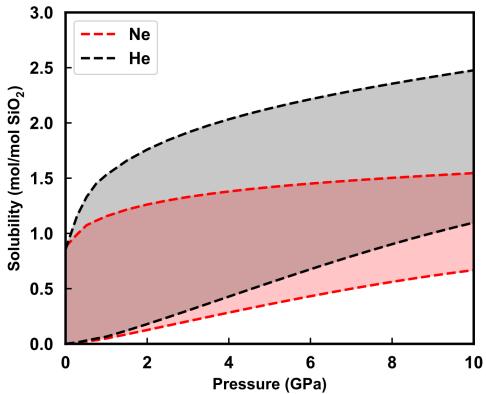




Figure 3. Void size distribution of SiO<sub>2</sub> and molecular size of common pressure media at
ambient condition and high pressure. Top panel: Molecular size data at ambient conditions
were adapted from Reid et al., 1987. Bottom panel: High pressure molecule sizes were calculated

319 using equations of state (He: Loubeyre et al., 1993; Ne: Dewaele et al., 2008; Ar: Ross et al.,

- 320 1986, H<sub>2</sub>O: Yoshimura et al., 2006).
- 321





322 323 Figure 4. Solubility of neon and helium in vitreous silica at high pressure. The two lines for 324 each medium represent upper limit and lower limit, respectively. The lower and upper limits 325 were estimated by (Vrigid-V normal)/Vgas and (Vrigid-Vsixfold)/Vgas, respectively, where Vrigid and V<sub>normal</sub> represent the molar volume of SiO<sub>2</sub> glass in noble gas media and non-gas media 326 conditions, accordingly;  $V_{gas}$  – the molar volume of noble gas;  $V_{sixfold}$  – the molar volume of 327 328 sixfold-coordinated SiO<sub>2</sub> glass (Sato et al., 2011). We do not have an accurate determination of 329 volume of silica in Ne, instead we assume the volume change under pressure is same as the He 330 case, as similar volume curves were suggested by integration method (Fig. S-3).