

# The Influence of Gravity Waves on the Slope of the Kinetic Energy Spectrum in Simulations of Idealized Midlatitude Cyclones

MAXIMO Q. MENCHACA AND DALE R. DURRAN

*Department of Atmospheric Sciences, University of Washington, Seattle, Washington*

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## ABSTRACT

The influence of gravity waves generated by surface stress and by topography on the atmospheric kinetic energy (KE) spectrum is examined using idealized simulations of a cyclone growing in baroclinically unstable shear flow. Even in the absence of topography, surface stress greatly enhances the generation of gravity waves in the vicinity of the cold front, and vertical energy fluxes associated with these waves produce a pronounced shallowing of the KE spectrum at mesoscale wavelengths relative to the corresponding free-slip case. The impact of a single isolated ridge is, however, much more pronounced than that of surface stress. When the mountain waves are well developed, they produce a wavenumber to the  $-5/3$  spectrum in the lower stratosphere over a broad range of mesoscale wavelengths. In the midtroposphere, a smaller range of wavelengths also exhibits a  $-5/3$  spectrum. When the mountain is 500 m high, the waves do not break, and their KE is entirely associated with the divergent component of the velocity field, which is almost constant with height. When the mountain is 2 km high, wave breaking creates potential vorticity, and the rotational component of the KE spectrum is also strongly energized by the waves. Analysis of the spectral KE budgets shows that the actual spectrum is the result of continually shifting balances of direct forcing from vertical energy flux divergence, conservative advective transport, and buoyancy flux. Nevertheless, there is one interesting example where the  $-5/3$ -sloped lower-stratospheric energy spectrum appears to be associated with a gravity-wave-induced *upscale* inertial cascade.

## 1. Introduction

The seminal analysis of data collected on commercial aircraft by [Nastrom and Gage \(1985\)](#) established that, on average, the midlatitude horizontal kinetic energy (KE) spectrum in the upper troposphere and lower stratosphere follows a  $k^{-3}$  slope at large scales with a gradual transition at wavelengths around 400 km to a  $k^{-5/3}$  slope at smaller scales, where  $k$  is the horizontal wavenumber. The  $k^{-3}$  portion of the spectrum is widely accepted as arising from an enstrophy cascade from large to small scales in approximate agreement with two-dimensional turbulence theory ([Kraichnan 1967](#); [Lindborg 1999](#)) or quasigeostrophic turbulence ([Charney 1971](#)). In contrast, there is much less agreement on the processes responsible for the  $k^{-5/3}$  slope in the mesoscale, which encompasses horizontal scales between roughly 4 and 400 km.

As shown by Kolmogorov and Obukhov ([Vallis 2017](#), p. 442) a  $k^{-5/3}$  spectral slope can be produced by an

inertial cascade of energy, and in three-dimensional homogeneous isotropic turbulence, this cascade is from large to small scales. But in the atmosphere, the time-mean horizontal KE spectrum exhibits a  $k^{-5/3}$  slope at scales far too large to be characterized by homogeneous isotropic motions. Several authors have therefore proposed alternative mechanisms through which horizontal KE might cascade through the mesoscale, including from large to small scales via gravity waves ([Dewan 1979](#); [VanZandt 1982](#)), or from small to large scales via stratified turbulence developing in regions such as thunderstorm anvils ([Gage 1979](#); [Lilly 1983](#)).

Nevertheless, the  $k^{-5/3}$  spectral slope does not have to be produced through an inertial cascade, and recent research supports the idea that energy injected into the small scales by gravity waves and deep convection might directly force the spectrum. [Waite and Snyder \(2009\)](#) simulated cyclones developing in a baroclinically unstable channel flow using 25-km horizontal grid spacing and found that gravity waves radiating from regions of geostrophic imbalance created a  $k^{-5/3}$  slope over the short-wavelength end of the spectrum in the lower

*Corresponding author:* Dale R. Durran, drdee@uw.edu

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stratosphere. The spectral energy budget for their simulations showed the stratospheric KE distribution was not produced solely by an inertial cascade, but also by direct forcing from vertical pressure flux divergence over much of the mesoscale. The simulations in [Waite and Snyder \(2009\)](#) were dry and underenergized compared to observations. When these simulations were extended to include moisture and parameterized deep convection, gravity wave activity was enhanced and transitions from  $k^{-3}$  to  $k^{-5/3}$  spectral slopes developed in both the troposphere and the stratosphere ([Waite and Snyder 2013](#)). As in the dry case, the  $k^{-5/3}$  spectral slope was not produced exclusively by an inertial cascade; rather the spectral energy budget showed contributions to the horizontal KE from buoyancy forces over a broad range of scales. [Peng et al. \(2015a\)](#) provided a more detailed analysis of the spectral energy budget for an idealized developing cyclone in a moist baroclinically unstable channel flow similar to that considered in [Waite and Snyder \(2013\)](#). Again using a horizontal grid spacing of 25 km, they found that direct forcing, from vertical pressure fluxes and horizontal KE fluxes, provided important contributions to the lower stratosphere KE spectrum over a broad range of mesoscale wavelengths. [Augier and Lindborg \(2013\)](#) examined 10 days of global simulations with the Atmospheric General Circulation Model for the Earth Simulator (AFES) run at 60-km horizontal resolution (T639) and 25 days of global simulations with the European Centre for Medium-Range Weather Forecasts (ECMWF) Integrated Forecast System (IFS) at roughly 30-km resolution (T1279). In both datasets they found that the stratospheric mesoscales were directly forced by upward energy fluxes associated with gravity waves propagating up from the troposphere, although in the AFES simulations, the vertical energy flux divergence was less important than a strong downscale energy cascade.

Using convection-permitting horizontal grid spacings of 1 km, [Durrán and Weyn \(2016\)](#) showed a  $k^{-5/3}$  horizontal KE spectrum can be generated in an initially quiescent horizontally uniform environment by idealized convective systems. [Sun et al. \(2017\)](#) analyzed the spectral KE budget in simulations similar to those in [Durrán and Weyn \(2016\)](#) and showed that buoyancy forces and vertical energy fluxes play an important role in regulating the KE spectrum across a wide range of scales, again suggesting the  $k^{-5/3}$  slope is not exclusively generated through an inertial cascade.

Mountains have long been identified as both a source of internal gravity waves and as regions in which observations show enhanced horizontal velocity variance ([Nastrom and Gage 1985](#); [Nastrom et al. 1987](#); [Jasperson et al. 1990](#); [Nastrom and Fritts 1992](#)). Numerical simulations

of the global atmospheric circulation can reproduce the transition between the  $k^{-3}$  and  $k^{-5/3}$  slopes in the horizontal KE spectrum ([Koshyk and Hamilton 2001](#); [Hamilton et al. 2014](#); [Skamarock et al. 2014](#)). [Malardel and Wedi \(2016\)](#) used a dry adiabatic version (no surface friction, no surface heat fluxes) of one such model to investigate the influence of orography on the horizontal KE spectrum and found Earth's topographic forcing greatly enhanced the energy in the mesoscale portion of the spectrum, shallowing the spectral slope from  $k^{-3}$  to something only slightly steeper than  $k^{-5/3}$ .

In this paper we again consider a dry baroclinically unstable channel flow, similar to that in [Waite and Snyder \(2009\)](#), to investigate the influence of a single isolated ridge on the domain-averaged horizontal KE spectrum. Our previous research on mountain waves generated in such flows ([Menchaca and Durrán 2017](#), hereafter [MD17](#)) showed that in the absence of surface friction, the surface winds around the developing cyclone become too strong to produce realistic mountain waves. Further evidence for the potential importance of surface friction was obtained by [Malardel and Wedi \(2016\)](#), who found that the horizontal KE on scales shorter than about 400 km in their free-slip dry adiabatic model with terrain exceeded the KE in a full-complexity reference simulation that included moisture and the full suite of standard convective and other physical parameterizations. We therefore include a simple parameterization of boundary layer stress in all of our simulations, except for a no-drag, no-mountain reference case similar to that in [Waite and Snyder \(2009\)](#).

In the following, [section 2](#) provides a brief overview of the numerical model and the simulations. The horizontal and vertical KE spectra for the suite of simulations are presented in [section 3](#). [Section 4](#) gives an analysis of the spectral energy budgets. [Section 5](#) contains the conclusions.

## 2. Numerical model and overview of the simulations

The large-scale flow and the initiation of the cyclone are described in [MD17](#), along with the shape of the isolated ridge, whose approximate  $x$  and  $y$  extents are 80 and 640 km. In addition to simulations with no topography, the same two mountain heights are again considered: 500 m and 2 km. All simulations are conducted with the Advanced Research version of the Weather Research and Forecasting Model (WRF-ARW) using a horizontal domain  $L_x = 16200$  km wide and periodic along the east–west ( $x$ ) axis with symmetric side walls bounding its  $L_y = 9000$ -km north–south ( $y$ ) extent. The domain is 20.5 km deep with the top 6 km devoted to a

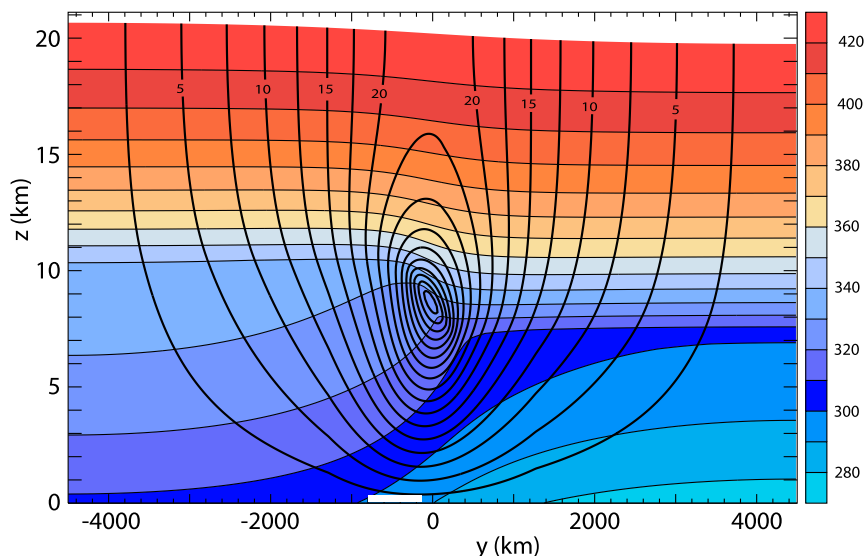


FIG. 1. North-south cross section of background shear flow:  $\theta$  (colors, 10-K intervals) and zonal velocity (black,  $2.5 \text{ m s}^{-1}$  intervals). The white bar extending from  $-800 \leq y \leq -160 \text{ km}$  shows the north-south extent of the ridge. Data are not plotted in the wave-absorbing layer.

Rayleigh-damping layer tuned to absorb vertically propagating gravity waves (as detailed in MD17). The horizontal grid spacing is  $\Delta x = \Delta y = 12 \text{ km}$ ; there are 80 vertical levels spaced at  $30 \text{ m}$  near the surface, with  $\Delta z$  increasing to roughly  $240 \text{ m}$  around the tropopause and  $400 \text{ m}$  near the model top. The time step is  $40 \text{ s}$ . The only physical parameterizations we employ are 2D Smagorinsky mixing in the horizontal and a modified version of the Yonsei University (YSU) boundary layer scheme (Hong et al. 2006) to limit the winds around the cyclone to realistic values through surface friction and vertical mixing. Default values of the parameters are used for the boundary layer scheme except that the heat flux is set to zero,  $z_0$  is a uniform  $0.1 \text{ m}$ , and the boundary layer height is fixed at the sixth model level, approximately  $282 \text{ m}$  above the surface.

Figure 1 is a  $y$ - $z$  cross section showing isotachs of the westerly wind component and isentropes of potential temperature in the baroclinically unstable shear flow. The north-south position of the mountain is shown by the white bar above the  $y$  axis, which is just south of the core of the jet. As discussed in MD17, a cyclone is triggered by an isolated PV perturbation; its evolution between 3.5 and 6.5 days is illustrated by the surface isobars and surface potential temperature fields plotted in Fig. 2 for the simulation with flat terrain. The position occupied by the mountain, in those simulations in which it is present, is indicated in Fig. 2 by the black vertical bar centered at  $(x, y) = (0, -480) \text{ km}$ . Note that the mountain covers roughly  $0.1\%$  of the surface area containing significant surface pressure perturbations at 6.5 days.

In the simulations with topography, mountain waves are weak at 3.5 days, when the cyclone is well upstream of the terrain (Fig. 2a). As the cold front passes the mountain at 4.5 days, waves amplify (Fig. 2b), and strong transient wave activity persists past 5.5 days (Fig. 2c), before slackening toward the end of the simulation as the cyclone propagates zonally past the mountain (Fig. 2d).

Throughout the following, we focus on four simulations; three use the YSU boundary layer scheme and terrain heights of 0, 500, and  $2000 \text{ m}$  and are denoted H0, H500, and H2000, respectively. The fourth simulation, H0FS, has flat terrain and a free-slip lower boundary. As shown by the comparison of the vertical velocity fields in Fig. 3, at 5.5 days many more short-wavelength gravity waves are generated in H0 than in the corresponding free-slip simulation H0FS. These waves originate along the cold front and in the cold air mass southeast of the low center and propagate upward into the stratosphere, where they are quite prominent at  $z = 12 \text{ km}$  (Fig. 3d), although their vertical-velocity amplitudes are not high [ $O(1) \text{ cm s}^{-1}$ ].

Several previous investigations of gravity waves generated during free-slip simulations of the life cycle of baroclinic waves have found a broad spectrum of gravity waves generated through a variety of mechanisms in different sectors around the cyclone (e.g., Plougonven and Snyder 2007; Kim et al. 2016). In comparison to these previous free-slip simulations, the gravity wave packets appearing in the surface-stress simulation H0 are unusually monochromatic and well defined, with an average wavelength of about  $70 \text{ km}$ , or roughly  $6\Delta x$ .

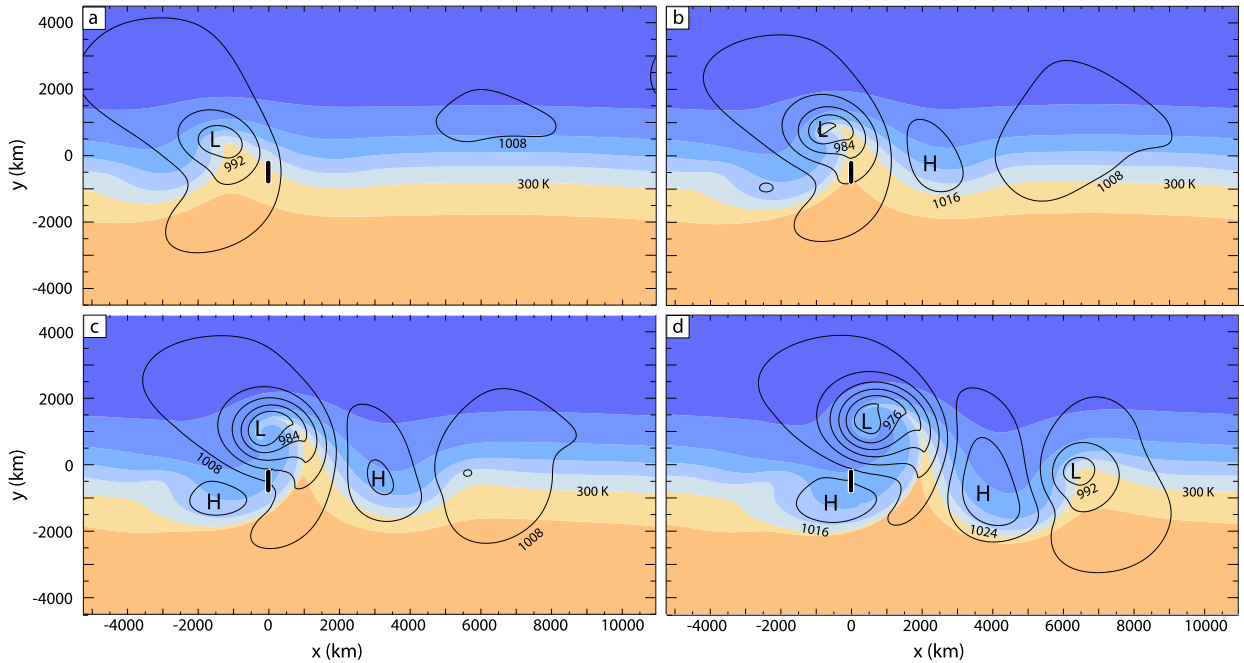


FIG. 2. Surface isobars (black, 8-hPa intervals) and surface  $\theta$  (color, 5-K intervals, 300-K isoline is labeled) for the developing cyclone in the no-mountain simulation at (a) 3.5, (b) 4.5, (c) 5.5, and (d) 6.5 days. The position occupied by the mountain in simulations where it is present is depicted by the black vertical bar at  $x = 0$  km. Lows and highs are labeled by “L” and “H,” respectively.

As such the waves are poorly resolved, and, in light of the findings of Snyder et al. (1993) about spurious gravity waves appearing in studies of 2D frontogenesis, it is possible they are a numerical artifact. To test this, the simulation was repeated with a nested mesh having a horizontal grid spacing of 4 km. As shown in Fig. 4 the vertical velocities in the waves on the finer mesh are stronger by roughly a factor of 3, and shorter wavelengths appear. Nevertheless, the reduction in the wavelength in the region cut by the green line segment is modest, and it is nowhere near as great as the factor-of-3 decrease in the grid spacing. The average wavelength along the 340-km-long line segment is 68 km in the 12-km-mesh simulation and 52 km in the 4-km simulation. The wavelength in the 4-km simulation is, therefore, roughly  $13\Delta x$ , which is adequate numerical resolution. Closer to the surface, at  $z = 1$  km, the well-organized wave pattern in Fig. 3 becomes a more complex superposition of waves on the 4-km mesh. A thorough investigation of the generation and propagation of these waves is left for future study.

### 3. Spectra

Energy spectra are obtained by first linearly interpolating the WRF data to constant- $z$  surfaces at vertical intervals of 250 m. The fields, which are already periodic

in  $x$ , are made doubly periodic using the rigid-wall boundary conditions to extend the data in  $y$ : the meridional velocity  $v$  is odd-extended, while the remaining variables are even-extended (Waite and Snyder 2009). The resulting data, which now fill a domain of size  $L_x$  by  $2L_y$  are Fourier transformed from  $(x, y)$  physical space to wavenumber space, where  $\mathbf{k} = k_x\mathbf{i} + k_y\mathbf{j}$  is the wave-number vector. The horizontal and vertical KE spectra are calculated as

$$E_h(\mathbf{k}) \equiv \frac{1}{2} \bar{\rho} [\hat{u}^*(\mathbf{k})\hat{u}(\mathbf{k}) + \hat{v}^*(\mathbf{k})\hat{v}(\mathbf{k})] = \frac{1}{2} \bar{\rho} \hat{\mathbf{u}}^*(\mathbf{k}) \cdot \hat{\mathbf{u}}(\mathbf{k}), \quad (1)$$

$$E_v(\mathbf{k}) \equiv \frac{1}{2} \bar{\rho} \hat{w}^*(\mathbf{k})\hat{w}(\mathbf{k}), \quad (2)$$

where  $\mathbf{u} = (u, v)$  is the horizontal velocity vector,  $w$  is the vertical velocity, Fourier-transformed variables are indicated by hats, complex conjugates are denoted by the star superscript, and  $\bar{\rho}(z)$  is a horizontally uniform background density profile. These 2D spectra are binned by the magnitude of the horizontal wavenumber  $k = |\mathbf{k}|$ , with bin size  $\Delta k = 2\pi/L_y$ , such that for an arbitrary field  $a$ ,

$$\hat{a}(k) = \sum_{\mathbf{k} \in R(k)} \hat{a}(\mathbf{k}), \quad (3)$$

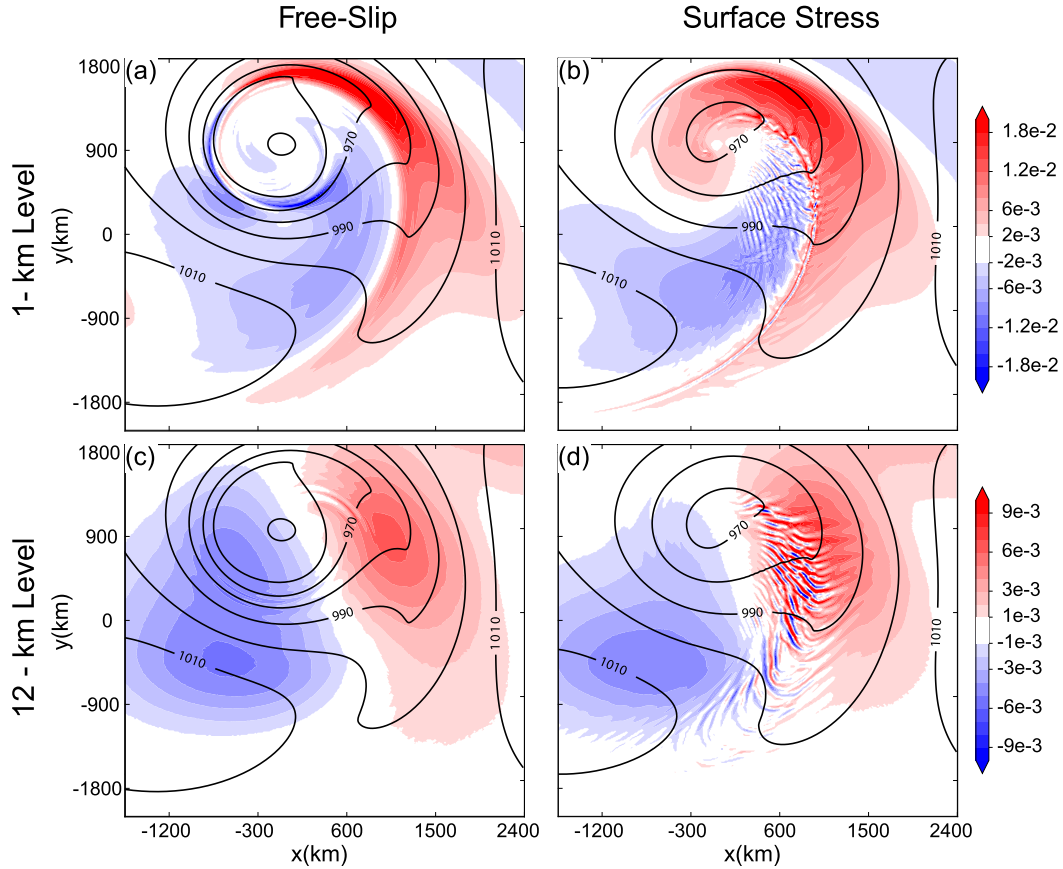


FIG. 3. The vertical velocity  $w$  (color,  $\text{m s}^{-1}$ ) at 5.5 days at (a), (b)  $z = 1$  and (c), (d)  $z = 12$  km from the (a), (c) free-slip simulation (H0FS) and (b), (d) surface stress simulation (H0). Also plotted for reference are isobars on the 50 m AGL surface, contoured every 10 hPa.

where  $R(k)$  is the set of  $\mathbf{k}$  satisfying

$$k - \frac{\Delta k}{2} \leq |\mathbf{k}| < k + \frac{\Delta k}{2}. \quad (4)$$

Data are plotted as a function of the nondimensional discrete horizontal wavenumber  $\tilde{k} \equiv (L_y/2\pi)k$ , and smoothed to reduce the noise introduced by 2D binning using the technique in Durrán et al. (2017).

#### a. Horizontal KE spectra

Spectra of the kinetic energy of the horizontal wind  $E_h(\tilde{k})$  for all four simulations are plotted in Fig. 5. Spectra from the midtroposphere, vertically averaged between 3.5 and 5.5 km, appear in the left column; spectra from the lower stratosphere, averaged between 11.5 and 13.5 km are in the right column. The 11.5–13.5-km layer is within the 9–14-km range over which data collected by commercial aircraft showed that, on average, the slopes of the atmospheric horizontal kinetic energy spectrum are  $k^{-3}$  and  $k^{-5/3}$  on the large-scale and mesoscale, respectively (Nastrom and Gage 1985). Reference lines

with slopes of  $\tilde{k}^{-3}$  and  $\tilde{k}^{-5/3}$  are therefore also plotted in each panel.

We first examine the influence of surface stress by comparing the two no-mountain simulations, whose spectra are plotted as solid (H0) and dashed (H0FS) red curves. At both 4.5 (Fig. 5a) and 6.5 days (Fig. 5c), the lower-tropospheric spectrum for H0 clearly has a flatter slope than that for H0FS for mesoscale wavelengths  $\lambda$  between roughly 50 and 110 km. By day 6.5, the slope in this wavelength band for H0 has become very close to  $\tilde{k}^{-5/3}$ , while that for H0FS is much steeper than  $\tilde{k}^{-3}$ . On the larger synoptic scale, for wavelengths between 450 and 4500 km, the spectral slopes for both H0 and H0FS are close to  $\tilde{k}^{-3}$  by day 6.5. The situation is qualitatively similar in the lower stratosphere (Figs. 5b,d), although at this higher level, the synoptic-scale slopes for both simulations remain steeper than  $\tilde{k}^{-3}$ , while the mesoscale slopes for H0 are flatter than  $\tilde{k}^{-5/3}$ . The mesoscale slopes for the free-slip case H0FS are significantly shallower than those at the corresponding times in the midtroposphere, but still steeper than the slopes for H0.



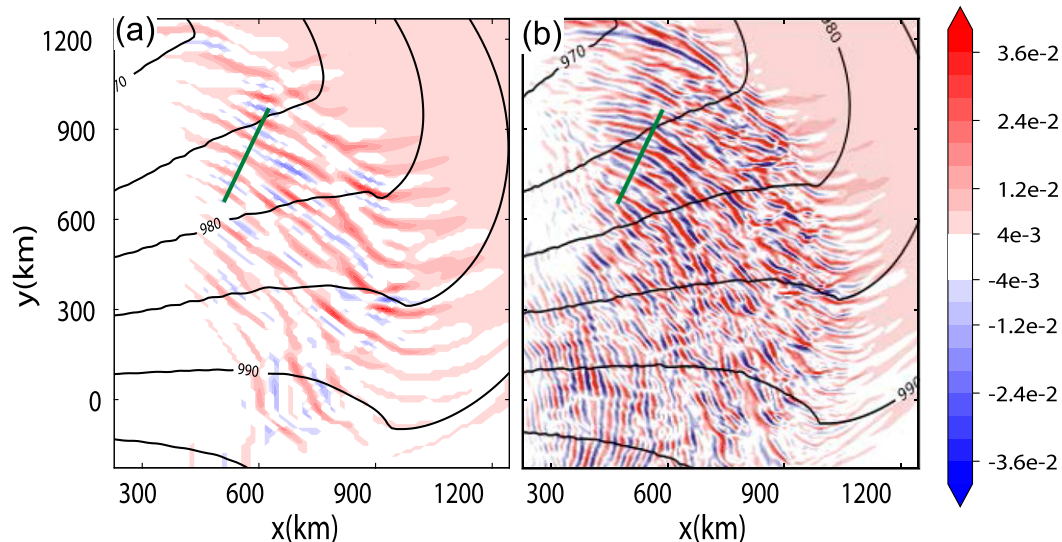


FIG. 4. The vertical velocity  $w$  (color,  $\text{m s}^{-1}$ ) at  $z = 12$  km in (a) a zoom of Fig. 3d into the region southeast of the cyclone center and (b) the same region from a nested simulation of the same case with 4-km horizontal grid spacing. Also plotted for reference are isobars on the 50 m AGL surface, contoured every 10 hPa and a 340-km-long reference line roughly perpendicular to the wave fronts (green).

The spectral slopes steepen due to numerical dissipation at wavelengths shorter than 50–60 km ( $4-5\Delta x$ ). As will be discussed later in more detail, the mesoscale shallowing of the spectrum of H0 relative to H0FS is associated with short-wavelength gravity waves generated near the surface that are not present in H0FS, as apparent in the vertical velocity fields in Figs. 3 and 4.

Much stronger gravity waves are generated in the simulations with the mountain ridge, and as evident in Fig. 5, there is substantially more  $E_h(\bar{k})$  at mesoscale wavelengths in these simulations and a broader range of wavelengths over which the spectral slopes are flatter than in the H0 simulation. At 4.5 days in the mid-troposphere (Fig. 5a), both H500 (green) and H2000 (blue) are clearly more energetic at scales below  $\lambda = 300$  km, with the spectral slopes equal to or shallower than  $\bar{k}^{-5/3}$  (Fig. 5a). By day 6.5 (Fig. 5c), the H2000 simulation shows increased values of  $E_h(\bar{k})$  beyond those in H0 for all wavelengths shorter than 750 km. Similarly, the H500 simulation has increased values of  $E_h(\bar{k})$  relative to H0 for all wavelengths shorter than 250 km. An approximate  $\bar{k}^{-5/3}$  slope is seen in  $E_h$  for the H500 case between wavelengths of 90 and 180 km, whereas the H2000 simulation exhibits a roughly  $\bar{k}^{-3}$  slope between wavelengths of 900 km and the dissipation scale.

The influence the mountain is even more dramatic in the lower stratosphere (Figs. 5b,d). Unlike in the troposphere, the spectrum of  $E_h$  decays more rapidly than  $\bar{k}^{-3}$  on the synoptic scales. Just slightly before the cold front arrives at the mountain, flattened, and even positive, spectral slopes are present on the mesoscale for

both H500 and H2000 (Fig. 5b) and  $E_h$  is enhanced relative to the H0 case through longer wavelengths than those in the troposphere. The horizontal KE spectra exhibit relative minima around wavelengths of 360 km for H500 and 180 km for H2000, suggesting that at this early time, the terrain's influence is strongest at small scales. By 6.5 days, the spectra for both H500 and H2000 exhibit clear  $\bar{k}^{-5/3}$  slopes between wavelengths of roughly 1000 km and the dissipation scale, and the  $E_h$  values for the 2-km-high-mountain case are about an order of magnitude higher than those for the 500-m-mountain case. These results suggest that even isolated mountains of modest height can at least temporarily generate the observed  $\bar{k}^{-5/3}$  slope of the mesoscale kinetic energy spectrum.

#### b. Rotational and divergent components of the horizontal KE spectrum

The horizontal kinetic energy at wavenumber  $\mathbf{k}$  is the sum of the rotational  $E_r(\mathbf{k})$  and divergent  $E_d(\mathbf{k})$  components of the horizontal wind,<sup>1</sup>

$$E_r(\mathbf{k}) = \frac{1}{2} \bar{\rho} \frac{|\hat{\zeta}(\mathbf{k})|^2}{k^2} \quad \text{and} \quad E_d(\mathbf{k}) = \frac{1}{2} \bar{\rho} \frac{|\hat{\delta}(\mathbf{k})|^2}{k^2}, \quad (5)$$

where  $\zeta = \partial_x v - \partial_y u$  is the vertical component of the vorticity, and  $\delta = \partial_x u + \partial_y v$  is the horizontal divergence.

<sup>1</sup> Because our domain is periodic in  $x$  and has symmetric northern and southern boundaries, there is no nontrivial contribution from a nondivergent irrotational velocity field.

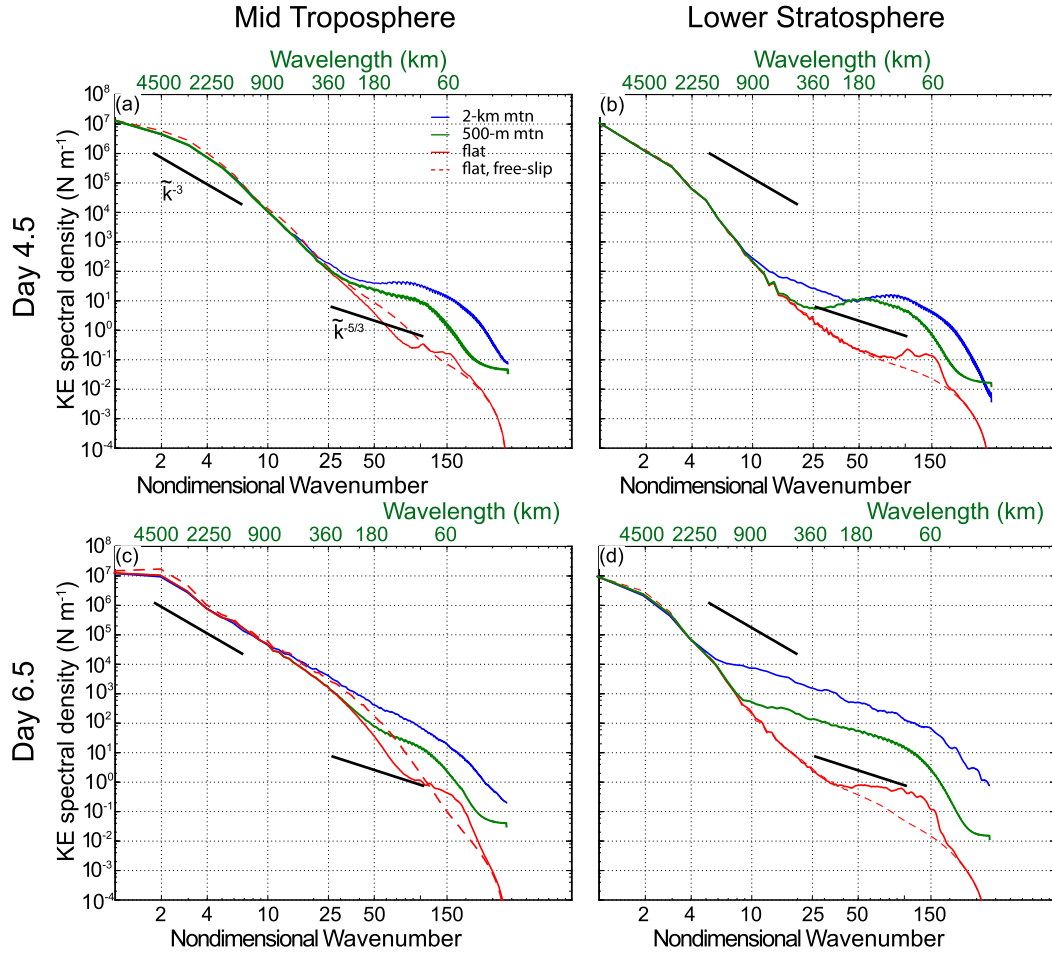


FIG. 5. Horizontal KE spectra  $E_h$  ( $\text{N m}^{-1}$ ) as a function of nondimensional wavenumber  $\tilde{k}$  and wavelength (km), averaged over (a),(c) the midtroposphere (3.5–5.5 km) and (b),(d) the lower stratosphere (11.5–13.5 km) at (a),(b) 4.5 and (c),(d) 6.5 days. Each panel shows the four simulations—H0FS (red dashed), H0 (red solid), H500 (green solid), and H2000 (blue solid)—along with reference lines having slopes of  $\tilde{k}^{-3}$  and  $\tilde{k}^{-5/3}$ .

Although they did not examine the influence of mountains, previous investigators have demonstrated that the divergent component is largely responsible for the shallowing of the kinetic energy spectrum at mesoscale wavelengths (Waite and Snyder 2009, 2013; Peng et al. 2015a). Figure 6 compares spectra for  $E_h$ ,  $E_r$ , and  $E_d$  in the lower stratosphere at 6.5 days for all four simulations; the curves for  $E_h$  (blue lines) are therefore identical to those plotted in Fig. 5d.

In the free-slip case, Fig. 6a,  $E_r$  (green line, mostly underneath the blue  $E_h$  curve) exceeds  $E_d$  (red line) at all wavelengths, and the difference is very large except in the mesoscale between roughly 100 and 200 km. The  $E_r$  and  $E_d$  spectra have similar slopes on the large scales that are steeper than  $\tilde{k}^{-3}$ , but these slopes are reduced on the mesoscale with the slope of  $E_d$  becoming shallower than both the slope of  $E_r$  and  $\tilde{k}^{-5/3}$ . These results are roughly similar to those obtained by Waite and

Snyder (2009, Fig. 7b) for  $E_r$  and  $E_d$  spectra in a free-slip baroclinic wave simulation except they found  $E_d$  to be sufficiently energetic to exceed  $E_r$  on the mesoscale. The values of  $E_d$  in Waite and Snyder (2009) are probably higher relative to  $E_r$  because they are from a later time in the development of the cyclone when scale contraction at the fronts and other inviscid processes have produced stronger gravity waves.

Surface stresses in the boundary layer hardly change  $E_r$  from the values for the free-slip simulation, but a substantial mesoscale enhancement of  $E_d$  in H0 relative to that in H0FS is immediately apparent in Fig. 6b. In H0,  $E_r$  and  $E_d$  again have similar steep slopes on the large scales, but  $E_d$  flattens out and even obtains a positive slope as it crosses the  $E_r$  curve at a wavelength of 180 km. The overall difference in horizontal kinetic energy spectrum between H0FS and H0 can be attributed to the difference in  $E_d$  produced by the stronger gravity waves in H0.

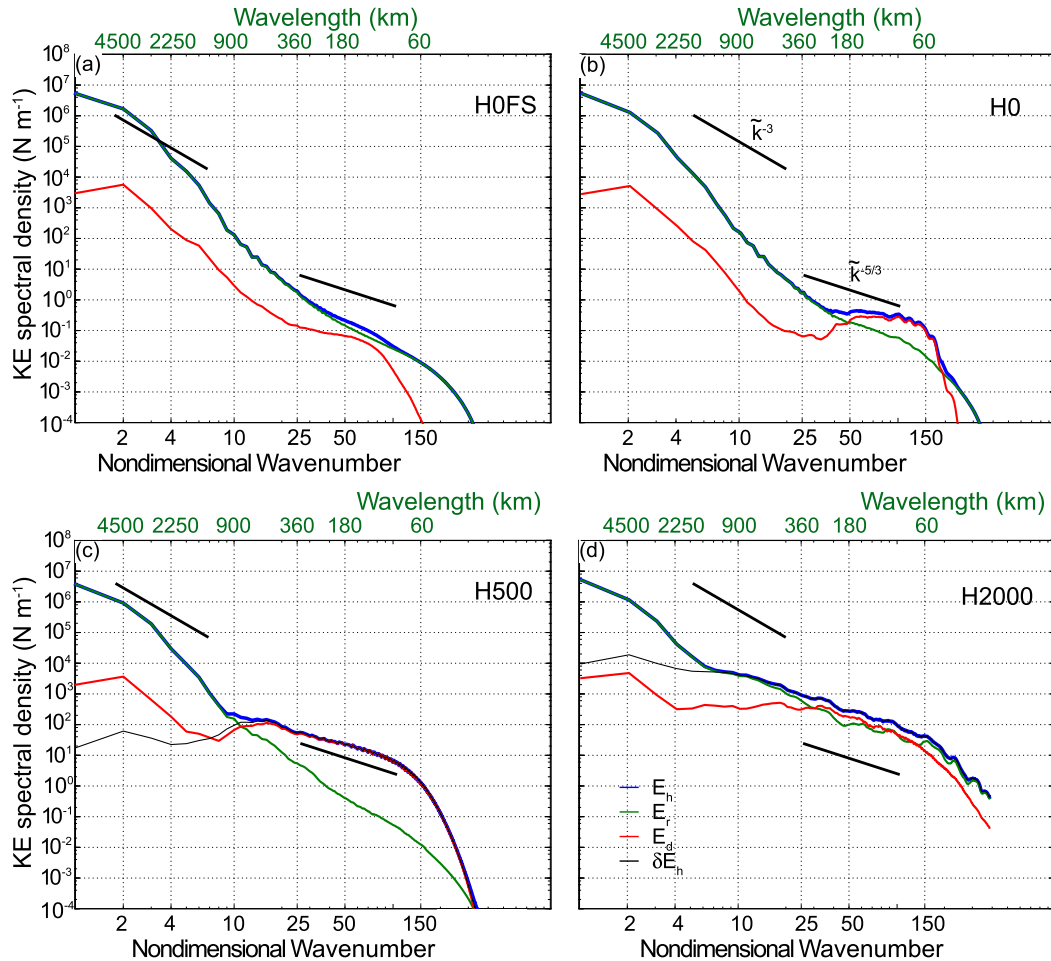


FIG. 6. Horizontal KE spectra as a function of nondimensional wavenumber  $\tilde{k}$  and wavelength (km) averaged over the lower stratosphere at 6.5 days for simulations (a) H0FS, (b) H0, (c) H500, and (d) H2000. Depicted are the total KE  $E_h$  (blue), the rotational component  $E_r$  (green), and the divergent component  $E_d$  (red), along with reference lines having slopes of  $\tilde{k}^{-3}$  and  $\tilde{k}^{-5/3}$ . In addition (c) and (d) show the mountain–no-mountain perturbation KE  $\delta E_h$  (black).

The influence of an isolated 500-m-high mountain is more dramatic than the domain-wide imposition of surface stress. As shown in Fig. 6c,  $E_r$  is only slightly more energetic in H500 than in H0, but  $E_d$  is much stronger and its spectrum follows a  $\tilde{k}^{-5/3}$  slope over a broad range of mesoscale wavelengths. While  $E_r$  continues to dominate  $E_d$  on the largest scales,  $E_d$  becomes dominant on the mesoscale, with the cross-over point occurring at a wavelength of 900 km. The mountain–no-mountain perturbation kinetic energy spectrum  $\delta E_h(\tilde{k})$  is also plotted in Fig. 6c, where  $\delta E_h(\tilde{k})$  is computed using the same formula as for  $E_h(\tilde{k})$  except that the horizontal velocities  $\tilde{\mathbf{u}}$  in (1) are replaced by the difference in the horizontal velocities between simulations with and without the mountain (in this case H500 and H0). At wavelengths less than 1200 km,  $\delta E_h$  is a very good approximation to  $E_d$ , suggesting that the difference in divergent kinetic energy in the two simulations is produced by vertically propagating

mountain waves. Moreover, just as in H0, the transition in slope of the horizontal kinetic energy spectrum  $E_h$  develops because the gravity wave generated  $E_d$  dominates  $E_r$  on the mesoscale.

The vertically propagating mountain waves generated by the 500-m-high mountain do not break (MD17, Figs. 6–9a,b), but there is substantial breaking and wave–mean flow interaction when the mountain is 2 km high (MD17, Figs. 11, 12).<sup>2</sup> In the 2-km-mountain

<sup>2</sup> The figures referenced in MD17 show results on a 5-km mesh, nested inside a 15-km outer domain; those referenced in Menchaca and Durran (2018) show results on a 4-km mesh, nested inside a 12-km outer domain. The spectra presented here were computed on a periodic 12-km unnested domain. Nevertheless, the breaking or nonbreaking character of the waves and general structure of the wake generated by the 2-km-high mountain remains the same in all simulations.



simulation, the perturbations induced by wave-breaking generate potential vorticity anomalies, some of which grow in scale as they propagate downstream (Menchaca and Durran 2018, Fig. 12), thereby providing forcing for the rotational component of the horizontal KE at intermediate or large scales. As a consequence, both  $E_d$  and  $E_r$  are significantly energized in H2000 relative to H0, and the mountain–no-mountain perturbation kinetic energy  $\delta E_h$  exceeds  $E_d$  at all wavelengths (Fig. 6d). As expected,  $E_r$  dominates  $E_d$  at large scales; their amplitudes become similar at wavelengths around 300 km, but unlike the spectra in H0 and H500,  $E_d$  does not dominate  $E_r$  at shorter wavelengths. Instead  $E_r$  and  $E_d$  exhibit similar amplitudes down to the dissipation scale of roughly 60 km ( $5\Delta x$ ). The small-scale  $E_r$  appears to be generated directly by intense small-scale vorticity perturbations within and downstream of the wave breaking regions (Menchaca and Durran 2018, Fig. 12). Similar tendencies for localized patches of turbulent flow to generate equal-amplitude divergent and rotational KE spectra were present at the anvil outflow level in simulations of isolated convective systems (Weyn and Durran 2017, Figs. 14, 15).

Several previous studies have suggested that the transition between steep spectral decay at large scales (nominally  $k^{-3}$ ) and the flatter  $k^{-5/3}$  slope on the mesoscale occurs at longer wavelengths in the stratosphere than in the troposphere because the spectral power in  $E_d$  is nearly constant with height, while that in  $E_r$  is distinctly lower in the stratosphere (Waite and Snyder 2009; Burgess et al. 2013; Skamarock et al. 2014). This is quite clearly the case in the H500 simulation, as illustrated by the  $E_r$  and  $E_d$  spectra plotted in Fig. 7 using data from day 6.5. Curves are plotted together for the midtroposphere (3.5–5.5 km), upper troposphere (6.5–8.5 km), and lower stratosphere (11.5–13.5 km). At wavelengths below 750 km,  $E_d$  (dashed lines) is virtually identical at all three levels. In contrast,  $E_r$  (solid lines) is very similar in both tropospheric layers, but one or two orders of magnitude weaker in the stratosphere.

### c. Vertical kinetic energy spectra

The vertical kinetic energy spectra  $E_v$  are plotted in Fig. 8. First consider the results for the midtropospheric layer between elevations of 3.5 and 5.5 km, shown in the left column. Similar large-scale signatures from the baroclinic wave are present in all simulations, where the slopes of  $E_v$  evolve toward  $\tilde{k}^{-3}$  at 6.5 days (Fig. 8c). This slope is also seen over the large scales in the  $E_v$  spectra for the dry baroclinic wave simulation of Peng et al. (2015a). At smaller scales,  $E_v$  decreases monotonically with decreasing wavelength in the H0FS simulation, whereas the corresponding spectra for H0 show a clear

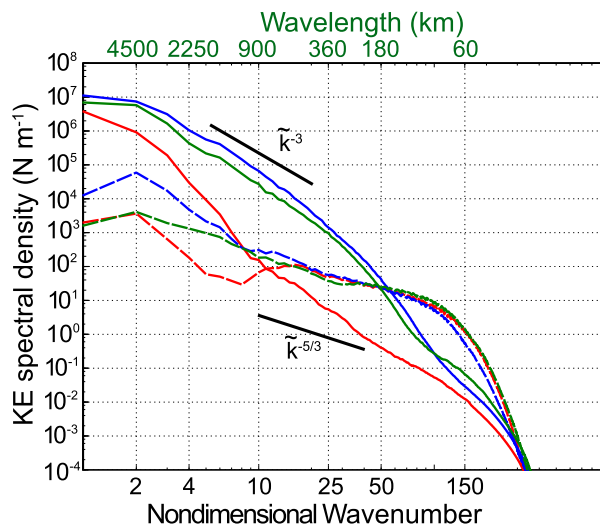


FIG. 7. Rotational and divergent components of the horizontal KE spectra as a function of nondimensional wavenumber  $\tilde{k}$  and wavelength (km) at 6.5 days for the H500 simulation. Parameters  $E_r$  (solid) and  $E_d$  (dashed) are averaged over the midtroposphere (3.5–5.5 km, green), the upper troposphere (6.5–8.5 km, blue), and the lower stratosphere (11.5–13.5 km, red). Reference lines having slopes of  $\tilde{k}^{-3}$  and  $\tilde{k}^{-5/3}$  are also shown.

separation between the values of  $E_v$  associated with the cyclone at large scales and a gravity-wave-produced secondary maximum in  $E_v$  at wavelengths between 90 km and those shorter scales at which numerical viscosity reduces the energy. As was the case with the horizontal KE, the waves triggered by the isolated ridge in H500 and H2000 produce much more mesoscale  $E_v$  than does surface stress acting alone in H0. At day 6.5, the  $E_v$  spectrum in both H500 and H2000 exhibits a modest minimum around a wavelength of 900 km and then gradually intensifies at shorter wavelengths. The H500 spectrum at this time is similar to that obtained (Skamarock et al. 2014, Fig. 10) in a global model with 15-km resolution that included both terrain and moist convection. Similar, relatively flat  $E_v$  spectra were generated by convection in simulations of a moist baroclinic wave (Peng et al. 2015a) and in 2.8-km-resolution ensemble simulations of warm-season high-precipitation events over Europe (Bierdel et al. 2012). In comparison to day 6.5, at day 4.5 the mesoscale-wavelength values of  $E_v$  are weaker and the local minimum between the large and small scales is much more distinct (Fig. 8a).

Because it lies above the strongest baroclinically unstable circulations, the large-scale  $E_v$  in the lower-stratospheric layer between 11.5 and 13.5 km (Figs. 8b,d) is over an order of magnitude smaller than in the midtroposphere. In contrast, the gravity waves generated by the terrain or surface friction yield mesoscale  $E_v$  whose amplitudes in the lower stratosphere are closer those in

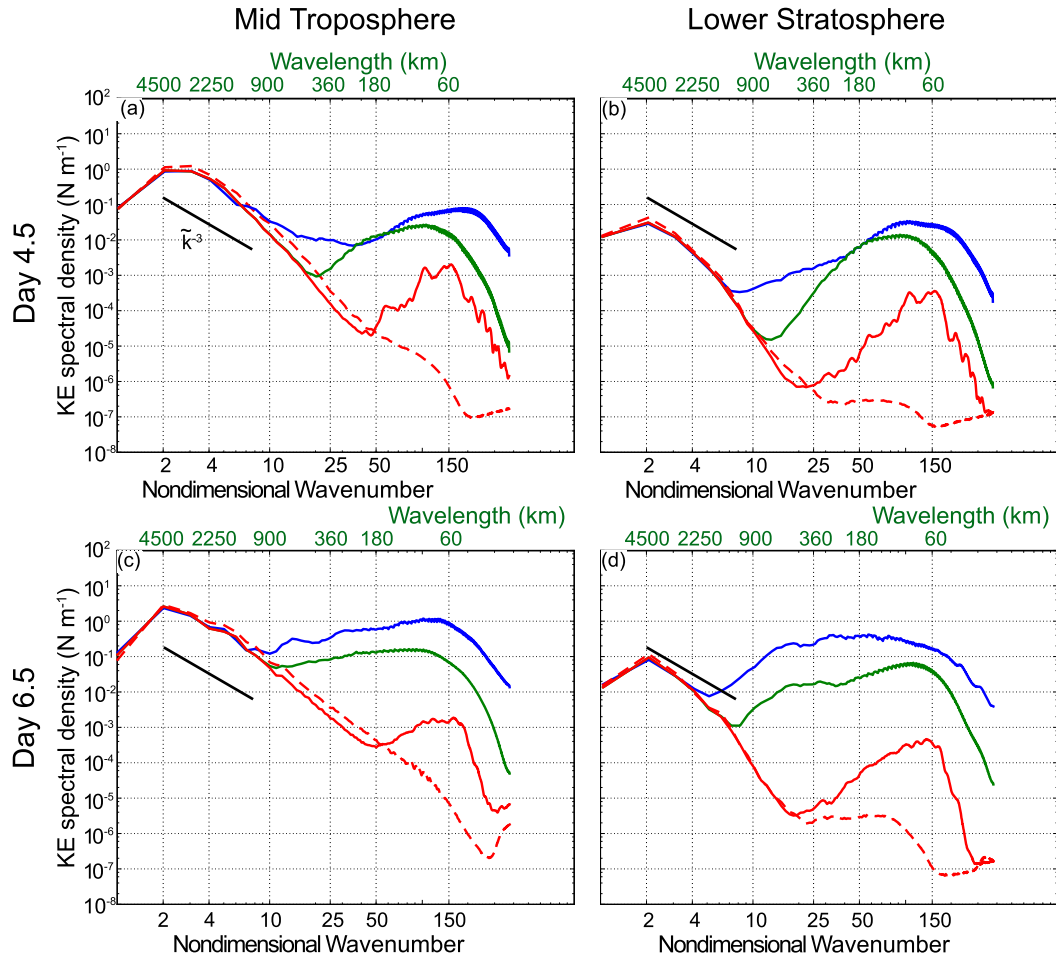


FIG. 8. As in Fig. 5, but the vertical KE spectrum  $E_v$  with  $\tilde{k}^{-3}$  reference lines.

the midtroposphere. At intermediate wavelengths, the relative minimum in  $E_v$  between the large scale and the mesoscale becomes more pronounced in the stratosphere than in the midtroposphere. At all times and levels, the terrain and the surface stress are more pronounced sources of small-scale power for the  $E_v$  spectra than for  $E_h$  (cf. Figs. 5 and 8).

#### 4. Spectral energy budgets

As demonstrated in the previous section, terrain-induced gravity waves have a strong influence on the KE spectra, energizing the spectrum at mesoscale wavelengths and tending to produce a  $k^{-5/3}$  slope on the mesoscale. In the following we analyze the spectral energy budget for  $E_h(\mathbf{k})$  to help identify the processes governing the evolution of the horizontal KE spectra, particularly at mesoscale wavelengths. As detailed in the appendix, the spectral horizontal energy budget equation may be expressed (Peng et al. 2014, 2015a; Sun et al. 2017)

$$\frac{\partial}{\partial t} E_h(\mathbf{k}) = T(\mathbf{k}) + B(\mathbf{k}) + V(\mathbf{k}) + D(\mathbf{k}), \quad (6)$$

where  $T$  is the contribution from advection written in the form (A8) that is conservative in the sense that it integrates to zero over all  $\mathbf{k}$ ,  $B$  is the buoyancy forcing, which represents the conversion from available potential energy (APE) to horizontal kinetic energy,  $V$  is the divergence of the vertical KE flux, and  $D$  is dissipation.

We are particularly interested in examining the extent to which KE is cascading between scales, and for such purposes it is convenient to plot the contributions from the forcing terms in (6) as the cumulative sum over all 2D wavenumbers larger than a given  $k$  (Augier and Lindborg 2013; Malardel and Wedi 2016). For an arbitrary variable  $X(k)$  we can define

$$\Pi_X(k) = \sum_{k \leq l \leq N} X(l), \quad (7)$$

where  $N$  is the maximum 2D wavenumber represented on the grid. If  $X(k)$  is one of the source terms in (6) and if

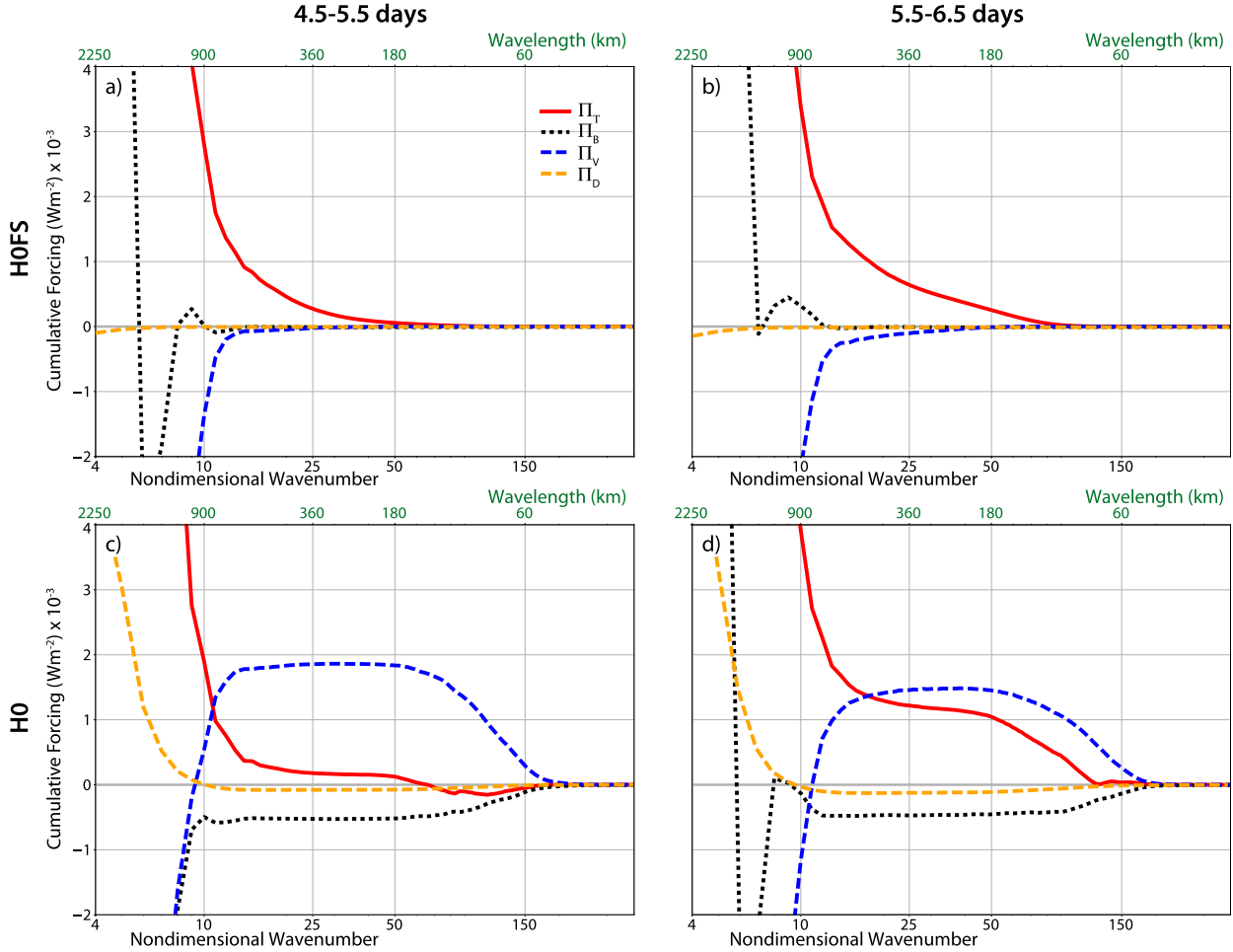


FIG. 9. Cumulative spectral budget terms averaged over the lower stratosphere (11.5–13.5 km) for (a),(b) H0FS and (c),(d) H0 averaged over (a),(c) 4.5–5.5 and (b),(d) 5.5–6.5 days. Plotted are  $\Pi_T$  (red solid),  $\Pi_B$  (black dotted),  $\Pi_V$  (blue dashed), and  $\Pi_D$  (orange dashed).

$\Pi_X(k)$  is plotted as a function of  $k$ , then the points in this plot where  $\Pi_X$  has a negative slope ( $\partial\Pi_X/\partial k < 0$ ) correspond to wavenumbers at which  $X(k)$  is tending to increase  $E_h(k)$  (Malardel and Wedi 2016, Fig. 3). Similarly, points with a positive slope correspond to wavenumbers at which  $X(k)$  acts to decrease  $E_h(k)$ . Because  $T(k)$  describes conservative energy transfer between horizontal wavenumbers, the sign of  $\Pi_T$  indicates whether it is acting to transfer  $E_h$  upscale or downscale. As an example, if  $\Pi_T(k) > 0$ ,  $T(k)$  is acting to produce a net transfer of  $E_h$  to wavenumbers greater than  $k$ ; which would be a downscale transfer to shorter wavelengths. Similarly, negative values of  $\Pi_T$  represent upscale transfer. As before, the values of  $\Pi_X$  will be plotted as a function of the nondimensional discrete horizontal wavenumber  $\tilde{k} \equiv (L_y/2\pi)k$ .

#### a. Lower stratosphere

The terms in the cumulative spectral KE budget, averaged over the lower stratospheric layer  $11.5 \leq z \leq 13.5$  km,

are plotted for the simulations without topography in Fig. 9. The longest wavelengths are omitted in these plots for better visibility, because 1) our focus is not on baroclinic instability, but rather on mesoscale processes, 2) the magnitudes of the terms are typically off-scale at the longest wavelengths, and 3) because the resolution in wavenumber space is very coarse at the longest wavelengths.<sup>3</sup>

In the free-slip case H0FS, downscale advective transfers are the primary source of KE at scales less than about 750 km ( $\Pi_T > 0$ ,  $\partial\Pi_T/\partial\tilde{k} < 0$ , Figs. 9a,b). In the 5.5–6.5-day average, there is also a tendency for the vertical energy flux divergence to remove KE from the layer ( $\partial\Pi_V/\partial\tilde{k} > 0$ ). The contribution from dissipation  $\Pi_D$  is negligible.

<sup>3</sup> Nevertheless, as expected,  $\Pi_T(0) = 0$ .

The budget for H0 is quite different: vertical energy flux convergence acts to increase KE over wavelengths of  $50 \leq \lambda \leq 180$  km ( $\partial \Pi_V / \partial \tilde{k} < 0$ , Figs. 9c,d). This forcing is opposed by buoyancy forces converting KE to PE over wavelengths of  $50 \leq \lambda \leq 150$  km, which is roughly the range over which there are significant differences in the lower stratospheric KE spectra of H0FS and H0 at 4.5 days in Fig. 5b. The 4.5–5.5-day forcing indicated by  $\Pi_T$  is weak, and the advective energy transfer switches from downscale to upscale for wavelengths shorter than about 130 km. Even at wavelengths longer than 130 km, the presence of boundary layer dissipation does not increase the contribution from  $T(k)$  relative to the free-slip case, although there is a region of  $200 < \lambda < 400$  km in Fig. 9c ( $\partial \Pi_T / \partial \tilde{k} \approx 0$ ) through which KE is transferred downscale without deposition. In contrast, in the 5.5–6.5-day averaged H0 budget there is pronounced downscale advective transfer of KE at all wavelengths and a strong tendency for  $T(k)$  to increase KE in the range  $70 \leq \lambda \leq 180$  km. The wavelength range in Fig. 5d where H0 deviates from H0FS is roughly the same range where there is significant contributions from  $\Pi_V$ ,  $\Pi_T$ , or  $\Pi_B$  in the 5.5–6.5-day cumulative budget. The contribution from  $\Pi_D$  is very small compared to the other terms, but surprisingly, it is greatest at large scales, and it acts to increase the KE at those scales. The tendency to increase KE comes entirely from the vertical component of the eddy diffusivity [see (A14)], and represents the convergence of vertical diffusive fluxes.<sup>4</sup>

The horizontal KE in the H500 case is strongly forced by vertical energy flux convergence (Figs. 10a,b), which is to be expected based on linear theory for steady-state mountain waves.<sup>5</sup> As derived by Eliassen and Palm (1960), the momentum flux carried by steady linear 2D mountain waves in the  $x$ – $z$  plane is constant with height,

$$\frac{d}{dz} \left( \bar{p} \int u' w' dx \right) = 0, \quad (8)$$

and the ratio of the vertical energy flux to the momentum flux is determined by the basic state cross-mountain wind speed  $U(z)$  such that

$$\int p' w' dx = -U \bar{p} \int u' w' dx, \quad (9)$$

<sup>4</sup> It is not clear why the dissipation so strongly forces the longest wavelengths, but it is worth repeating that, at these long wavelengths, the magnitude of  $\Pi_D$  is still far smaller than that of the other three terms.

<sup>5</sup> Wave action, not perturbation KE, is conserved as gravity waves propagate through a shear flow.

where in the preceding, the primes denote perturbations,  $p$  is pressure, and the integrals are either taken over a wavelength or the full periodic domain. The relation (9) also holds in observed mountain waves (Smith et al. 2008, 2016). Since the momentum flux associated with the mountain waves in the H500 simulation is almost constant with height (Durran and Menchaca 2019, Fig. 4), (9) implies that the energy flux associated with such waves would decrease with height at levels above the jet, where the environmental wind speed decreases with height. Such vertical energy flux convergence, which acts to increase KE, is strongest in the wavelength range  $60 < \lambda < 180$  km, particularly at 5.5–6.5 days when  $\partial \Pi_V / \partial \tilde{k}$  is strongly negative.

Over both the 4.5–5.5- and 5.5–6.5-day averaging periods,  $T(k)$  is transferring energy upscale across wavelengths in the range  $90 < \lambda < 1500$  km. This is roughly similar to the  $90 < \lambda < 1100$ -km range of wavelengths over which the KE spectrum for H500 exhibits a clean  $\tilde{k}^{-5/3}$  slope in Fig. 5d. KE is removed by advective transfer from almost the same small scales where it is forced by vertical energy flux divergence,  $70 < \lambda < 150$  km. Conversely,  $T(k)$  deposits KE at wavelengths greater than about 1100 km, where  $V(k)$  is acting to reduce KE. Downscale advective transfer is also a source of KE at small wavelengths less than 70 km. Vertically propagating gravity waves have zero buoyancy flux, so not surprising that  $\Pi_B$  is not an important term in the budget. The dissipation  $\Pi_D$  is also unimportant.

The overall scenario suggested by the lower-stratospheric KE budget from the H500 simulation is that vertical energy flux divergence associated with gravity waves propagating from below deposits energy at small scales, which then cascades to larger scales via conservative advection. Since, at wavelengths shorter than about 800 km, the KE is entirely divergent and equal to difference  $\delta KE_h$  between the KE in the mountain and no-mountain simulations (Fig. 6c), the inertial cascade is likely associated with gravity waves. The upscale cascade is surprising because gravity waves are thought to cascade energy *downscale* to create the  $k^{-5/3}$  KE spectrum in the ocean (Garrett and Munk 1979). Other recently computed lower stratospheric spectral KE budgets suggest a variety of alternative behaviors when there is forcing from parameterized moist convection. Augier and Lindborg (2013) found a downscale advective cascade across the mesoscale in AFES simulations, but very little advective transfer in the ECMWF IFS model. On the other hand, Peng et al. (2015a) found a weak upscale energy transfer by  $\Pi_T$  at wavelengths longer than 360 km in a channel-model simulation of moist baroclinic instability. Even increasing the mountain height to 2 km in these simulations

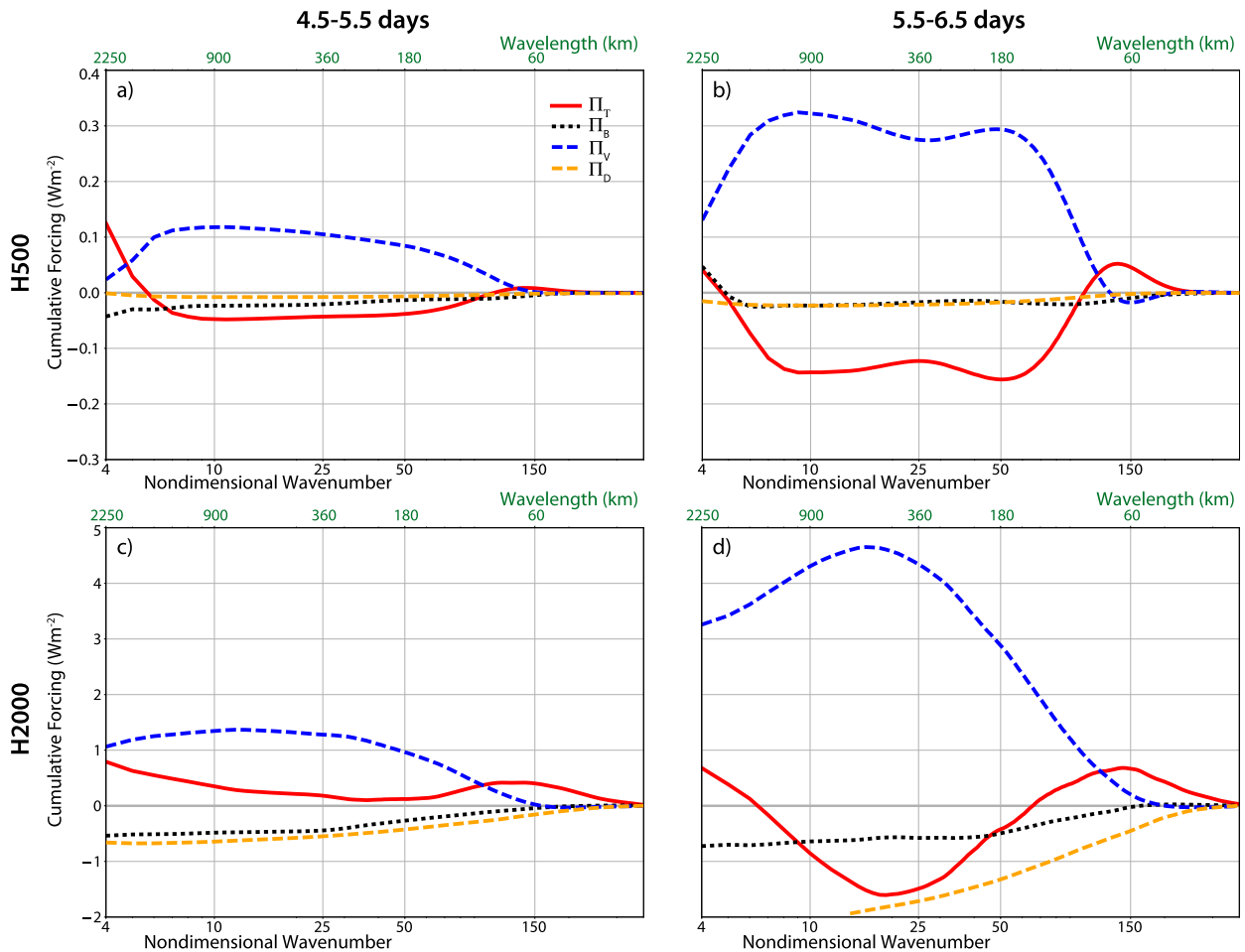


FIG. 10. As in Fig. 9, but for (a),(b) H500 and (c),(d) H2000. Note the increases in the scale on the vertical axis as the mountain height is increased.

changes the direction of the KE transfer associated with  $\Pi_T$  at 4.5–5.5 days to weakly downscale.

In contrast to the H500 simulation, extensive wave breaking, turbulent mixing, and associated secondary circulations occur in the presence of the 2-km-high mountain, although there is no wave breaking in the stratosphere until after day 4.5 (MD17). As a consequence, the cumulative buoyancy forcing in H2000 increases by more than an order of magnitude from that found in H500 (Figs. 10c,d, note the increase the vertical scale from Figs. 10a,b):  $B(k)$  tends to convert KE to available potential energy at essentially all wavelengths longer than 60 km, but the effect is strongest up through wavelengths of 360 km in the 4.5–5.5-day average and up through wavelengths of 180 km in the 5.5–6.5-day average. At 4.5–5.5 days, the distribution of  $\Pi_V$  is roughly similar to that for H500, but at 5.5–6.5 days, there is a strong tendency for the vertical energy flux divergence to increase the KE over the larger range of wavelengths

( $50 < \lambda < 500$  km), while acting to decrease the KE at longer wavelengths. Large differences from the H500 case are also apparent in  $\Pi_T$ . At 4.5–5.5 days, upscale energy transfer has been replaced by downscale transfer and  $T(k)$  also tends to increase the KE at all wavelengths longer than 300 km.

Finally, in contrast to the other cases,  $\Pi_D$  is nontrivial at 4.5–5.5 days in H2000, where it acts as a sink of KE at a rate similar to  $\Pi_B$  over all wavelengths. At 5.5–6.5 days,  $\Pi_D$  acts as an even stronger sink on the KE (Fig. 10d). One might suppose that the dissipative forcing would be focused at the highest wavenumbers, and the contribution to  $\partial\Pi_D/\partial k$  from horizontal eddy diffusion is indeed limited to wavelengths shorter than 180 km. The contribution from vertical eddy diffusion, on the other hand, acts across a much larger range of scales. One reason why forcing from  $\Pi_D$  appears across a broad range of scales is because the regions of vigorous subgrid-scale mixing produced by wave breaking are



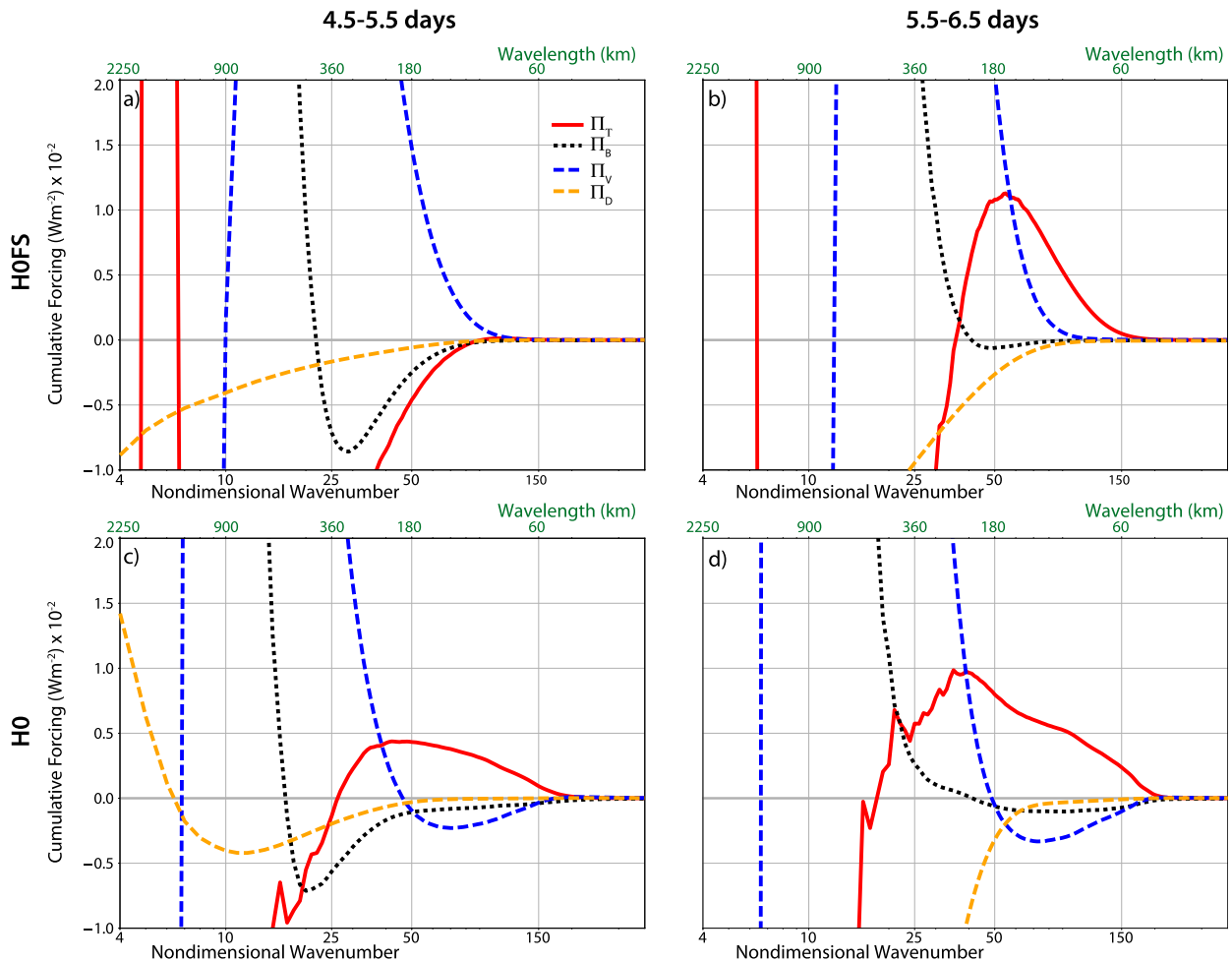


FIG. 11. As in Fig. 9, but averaged over the midtroposphere (3.5–5.5 km).

localized in physical space to a small region above the mountain, and as such, the distribution of their dissipative forcing is very broad in wavenumber space. Peng et al. (2015a) found a roughly similar wavenumber dependence for the dissipation in the lower stratosphere in their moist (RH60) baroclinic wave simulation.

As with the other terms in the cumulative budget, the differences in  $\Pi_T$  between H500 and H2000 are greatest in the 5.5–6.5-day average (cf. Figs. 10b and 10d). Instead of a broad region where  $\partial\Pi_T/\partial\tilde{k}$  is near zero, the conservative advective flux largely counteracts the forcing from  $V(k)$ , tending to decrease KE in the wavelength range  $50 < \lambda < 500$  km while acting to increase KE at longer and shorter wavelengths. The cumulative advective flux is upscale for  $150 < \lambda < 1500$  km and downscale at longer and shorter wavelengths. Given the broad range of wavelengths over which direct forcing is implied by significant variations in  $\Pi_V$ ,  $\Pi_T$ , and  $\Pi_D$ , it is apparent that a nearly nondissipative advective cascade of KE, such as that found for H500, is not required to produce

the  $\tilde{k}^{-5/3}$  spectral slope shown for H2000 in Fig. 5d over wavelengths in the range  $60 < \lambda < 1500$  km.

#### b. Midtroposphere

We now consider the cumulative spectral horizontal KE budget for the 3.5–5.5-km layer in the middle troposphere. In both H0FS and H0, a negative  $\Pi_B$  slope at synoptic scales (Fig. 11,  $\Pi_B$  itself is largely off scale) is associated with the conversion of APE to KE as expected in a baroclinically unstable flow (Augier and Lindborg 2013; Malardel and Wedi 2016; Peng et al. 2015b). In the mesoscale, the 4.5–5.5-day-averaged budget for H0FS (Fig. 11a) shows horizontal KE being created by vertical energy flux divergence  $V(k)$ , while it is removed by both  $B(k)$  and  $T(k)$  between roughly  $90 < \lambda < 360$  km, and transferred upscale. In comparison, at 5.5–6.5 days the tendency of  $B(k)$  to diminish KE is much smaller, while  $T(k)$  acts more strongly to reduce KE at scales greater than 150 km, and to increase KE at shorter scales (Fig. 11b). The advective energy transfer

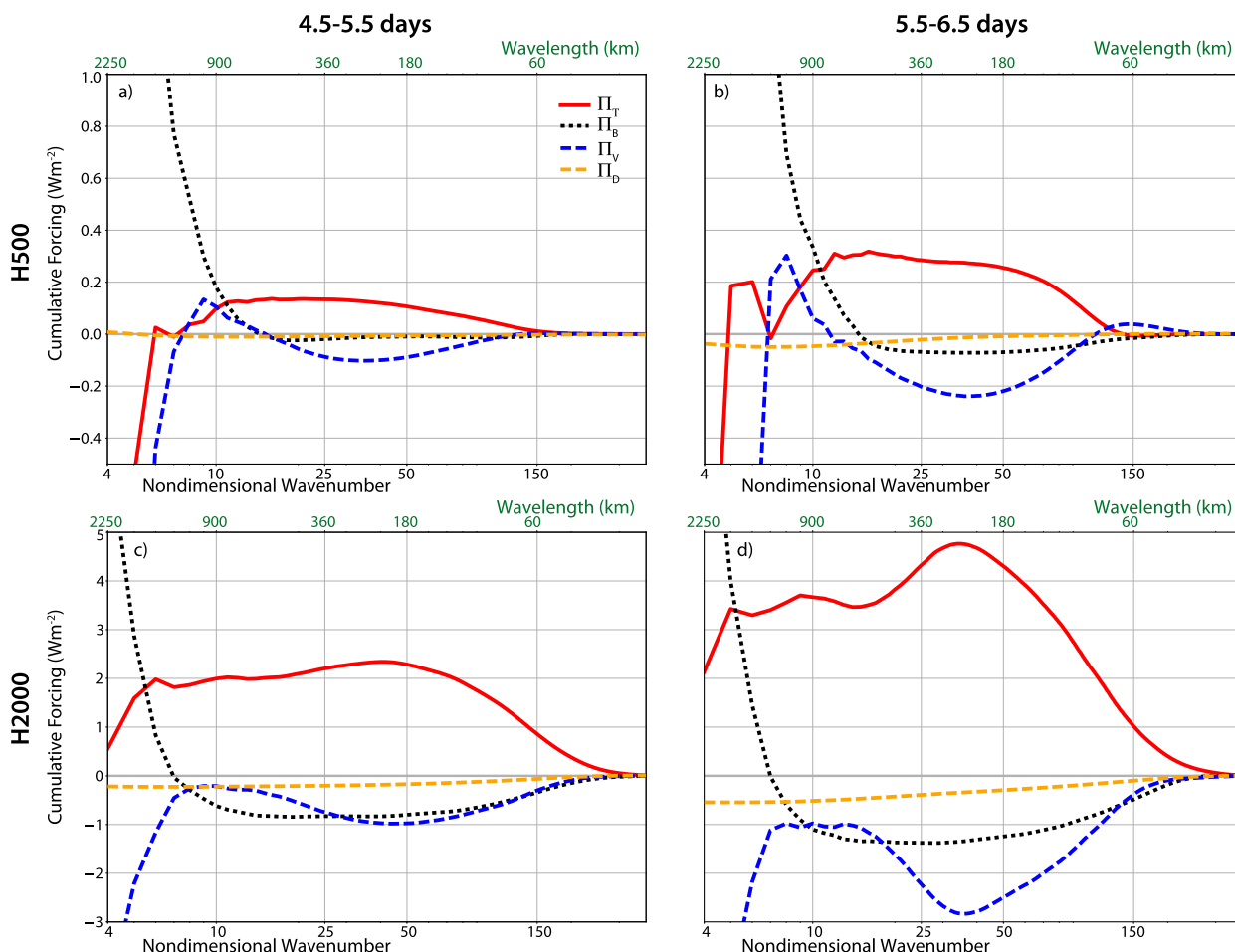


FIG. 12. As in Fig. 10, but averaged over the midtroposphere (3.5–5.5 km). Note the further increases in the scale on the vertical axis.

is downscale for  $\lambda < 250$  km and upscale for longer mesoscale wavelengths. The vertical scale is sufficiently exaggerated (to show the mesoscale region) that  $\Pi_D$  looks potentially nontrivial, but in fact it is not significant in comparison to the dominant budget terms.

Under the presence of surface stress,  $V(k)$  acts to remove KE over wavelengths between  $50 < \lambda < 150$  km, which are approximately the same wavelengths at which vertical energy flux divergence in the same H0 simulation supplies KE to the lower stratosphere (cf. Figs. 9c,d and 11c,d). This negative KE tendency is largely balanced by downscale advective transport of KE beginning around wavelengths of 220 km. The region with a spectral slope approximating  $k^{-5/3}$  in the midtroposphere at day 6.5 in Fig. 5c also occupies most of this same wavenumber band. The contribution of  $\Pi_D$  to the spectral KE budget for H0 is again small compared to the dominant terms in the budget, although as was the case in the lower stratosphere, vertical subgrid scale diffusive fluxes make a positive contribution to the KE at very long wavelengths.

Turning to the influence of topography, in the midtroposphere the contributions of  $V(k)$  and  $T(k)$  in H500 are roughly opposite those in the lower stratosphere at wavelengths shorter than 250 km. Since the budget is computed below the height of the jet maximum, (9) suggests steady nonbreaking mountain waves will produce energy flux divergence. As shown in Figs. 5a and 5c, over approximately the same  $90 < \lambda < 180$  km interval, where H500 exhibits a  $k^{-5/3}$  spectrum the vertical energy flux divergence acts to reduce KE, and the conservative advective transfer is downscale and tending to increase the KE (Figs. 12a,b). Nevertheless, at scales larger than the mountain waves ( $250 < \lambda < 1100$  km)  $V(k)$  exhibits a weak tendency to increase KE, which is similar to the behavior in the lower stratosphere. The buoyancy forcing  $B(k)$  plays about the same modest role in both the midtroposphere and the lower stratosphere for wavelengths less than 450 km (cf. Figs. 10a,b and 12a,b), but at longer wavelengths there is substantial conversion of APE to KE in

the growing baroclinically unstable wave. The contributions from dissipation are negligible.

In H2000 the midtropospheric  $\Pi_V$  shows a tendency for  $V(k)$  to reduce KE at mesoscale wavelengths shorter than 230 km, which is again opposite its behavior in the lower stratosphere (Figs. 12c,d). The advective KE transfer is downscale at all wavelengths and tends to increase KE for  $\lambda < 230$  km. At slightly larger scales, the slopes of  $\Pi_V$  and  $\Pi_T$  switch signs; particularly in the 5.5–6.5-day average, there is a tendency for  $V(k)$  to increase, and  $T(k)$  to decrease, KE in the wavelength band of  $230 < \lambda < 750$  km. At wavelengths less than 180 km,  $B(k)$  tends to more strongly decrease KE than in any of the other cases. As in all the midtropospheric budgets, at synoptic-scale wavelengths  $B(k)$  tends to increase KE in the baroclinically unstable flow. The net effect of all this forcing at mesoscale wavelengths is to increase the 6.5-day horizontal KE relative to that in the H0 simulation at all wavelengths shorter than 900 km (Fig. 5c), but instead of a  $k^{-5/3}$  slope over these wavelengths, the spectrum follows an approximate  $k^{-3}$  slope. Because there is almost no wavebreaking in the midtroposphere, the dissipative contribution of dissipation  $\Pi_D$  to the spectral KE budget is much weaker, than it is in the lower stratosphere, and this difference is particularly pronounced at 5.5–6.5 days (cf. Figs. 10d and 12d).

Also of note is the systematic increase in the scale at which the conversion of APE to KE ceases as the topographic forcing increases. In the 5.5–6.5-day average,  $\Pi_B$  exhibits strong negative slopes down through wavelengths of about 360 km in H0, 750 km in H500, and 900 km in H2000 (Figs. 11d, 12b,d). The mountains appear to be inhibiting the cyclogenetic intensification of KE at small synoptic scales.

## 5. Conclusions

We have examined the influence of gravity waves on the KE spectrum associated with a developing cyclone in a dry baroclinically unstable shear flow. Although deviations from geostrophic balance eventually become strong enough to trigger gravity wave emission in free-slip simulations of such cyclones (Waite and Snyder 2009), we have focused on somewhat earlier stages of the cyclone life cycle when the gravity waves triggered in the free-slip simulation are weak. In comparison to the free-slip case, the gravity waves in the vicinity of the cold front are significantly enhanced by surface stress. The precise generation mechanism for these waves, and the extent to which they are accurately captured in the numerical simulations, are left for future analysis. The surface-stress-induced gravity waves energize the divergent component  $E_d$  of the horizontal KE spectrum at

wavelengths less than about 200 km, but they have almost no influence on the rotational component  $E_r$ . The total horizontal KE spectrum in both the troposphere and the stratosphere develops a pronounced region of shallow slope as a result of these frictionally generated gravity waves. The simulations with topography also incorporated the same surface boundary layer parameterization to avoid unrealistically large surface winds and high-amplitude mountain waves. Although the terrain only covers about 0.1% of the surface area filled by significant large-scale disturbances at 6.5 days, mountain waves generated by flow over the single isolated ridge have a major impact on the domain-integrated KE spectrum, and this impact is much greater than that generated by surface stress alone.

The mountain waves generated by the 500-m-high ridge do not break and their contribution to the horizontal KE spectrum  $E_h$  is entirely captured by the divergent component  $E_d$ . When the waves are mature, at 6.5 days, the slope of the  $E_d$  spectrum closely matches  $k^{-5/3}$ , and the amplitude of  $E_d$  is constant with height between the midtroposphere and lower stratosphere at wavelengths shorter than about 750 km. The transition from a  $k^{-3}$  slope at large scales to a  $k^{-5/3}$  slope at small scales occurs at longer wavelengths in the stratosphere than in the troposphere because of the lower level of background  $E_r$  in the stratosphere, as previously suggested by Waite and Snyder (2009), Burgess et al. (2013), and Skamarock et al. (2014).

The mountain waves generated by the 2-km-high ridge break down in both the troposphere and the stratosphere (MD17), creating patches of vertical vorticity, and thereby enhancing the rotational as well as the divergent components of  $E_h$  at essentially all scales. In the stratosphere at day 6.5,  $E_r$  and  $E_d$  have roughly equal amplitude over wavelengths  $60 < \lambda < 360$  km, and their sum  $E_h$  has developed an approximate  $k^{-5/3}$  slope over the very broad spectral band of  $60 < \lambda < 1800$  km. In the troposphere, on the other hand, where  $E_r$  is largely energized by baroclinic instability, the flattening of the  $E_h$  spectral slope is much less pronounced. The creation of patches of vertical vorticity in wave breaking regions is at least qualitatively comparable to circulations generated in simulations of deep convective systems in horizontally homogeneous environments [compare Fig. 13a of Menchaca and Durran (2018) with Figs. 17d–f of Weyn and Durran (2017)], and the rotational and divergent components of the horizontal KE spectrum were also found to have roughly equal amplitude in the upper troposphere and lower stratosphere in those convective systems.

Spectral budgets for the horizontal KE were evaluated in the midtroposphere and the lower stratosphere.

The most surprising results were obtained for the 500-m-mountain case in the lower stratosphere, where the 5.5–6.5-day averaged budget shows energy input in the mesoscale from vertical energy flux convergence, energy removal at the same scale by conservative advection, and a near inertial advective cascade *upscale* from wavelengths of about 180 km to wavelengths of about 1100 km (Fig. 10b). Nevertheless, as apparent in the corresponding KE budget for the simulation with the 2-km-high mountain, the horizontal KE spectrum can also develop an approximate  $k^{-5/3}$  slope over an even broader range of wavelengths without any inertial cascade whatsoever. Instead, the  $k^{-5/3}$  slope is produced by nontrivial contributions from all four terms in the spectral KE budget: vertical energy flux convergence, conservative advection, buoyancy forcing, and dissipation (Fig. 10d). Moreover, the forcing from the two largest terms, vertical energy flux convergence and conservative horizontal advection, flips sign around a wavelength of about 500 km without producing any obvious deviations from the  $k^{-5/3}$  slope. Given the previously discussed similarities between the local perturbations produced in wave breaking regions and those generated aloft by deep convection, our finding that direct forcing across a wide range of scales is responsible for the  $k^{-5/3}$  spectrum in the 2-km-mountain simulation might have been anticipated from Sun et al. (2017), who obtained a similar result analyzing simulations of deep convection.

The contributions from vertical energy flux divergence and conservative advection in the midtroposphere KE budgets are opposite in sign relative to those in the lower stratospheric budgets. Vertical energy fluxes are removing KE from the midtroposphere at mesoscale wavelengths, transporting it upward, and depositing it in the lower stratosphere. The conservative advective transport of KE also switches from downscale in the midtroposphere to upscale in the lower stratosphere, although in the 5.5–6.5-day averaged budgets, the only near inertial cascade occurs in the lower stratosphere for the 500-m-high mountain.

The slope of the mesoscale KE spectrum varies significantly over time and does not seem to be produced by any single universal dynamical mechanism. At 6.5 days, for example, the spectral slopes at mesoscale wavelengths in the lower stratosphere are close to  $k^{-5/3}$ , but are far more variable at 4.5 days (Figs. 5b,d). Substantial variability in the slope of the mesoscale energy spectrum has also been documented in previous studies. Selz et al. (2019) found a strong correlation between precipitation rate and variability in the spectral slope at short wavelengths in three years of analysis data over Europe. Case-study ensemble simulations of weakly forced diurnal convection have demonstrated that the

slope can approach  $k^{-5/3}$  during hours of active convection and then drop back to steeper values when the convection dies down (Weyn and Durran 2019, Fig. 5). It may be that the  $k^{-5/3}$  mesoscale KE spectrum in the atmosphere is not as universal as sometimes supposed, and that its slope should not be interpreted as providing evidence for the action of any particular dynamical mechanism (Stumpf and Porter 2012).

Among the different ways that mountain waves produce a  $k^{-5/3}$  spectrum, the H500 case is worthy of further comment. The mesoscale perturbations in this case are essentially all gravity waves, and the spectral energy budget suggests an inertial cascade is occurring between the large and small scales, yet in contrast to the commonly understood action of gravity waves in creating the  $k^{-5/3}$  KE spectrum in the ocean (Garrett and Munk 1979), the energy transfer is *upscale* (Figs. 10a,b). We hypothesize this upscale transfer is produced as the waves propagate through a background horizontal flow with strong horizontal variations. The WKB ray tracing equations for the horizontal wavenumbers ( $k_x, k_y$ ) in quasi-linear gravity waves propagating upward and westward relative to a background velocity [ $U(x, y), V(x, y)$ ] are

$$\frac{D_g k_x}{Dt} = -k_x \frac{\partial U}{\partial x} - k_y \frac{\partial V}{\partial x}, \quad (10)$$

$$\frac{D_g k_y}{Dt} = -k_x \frac{\partial U}{\partial y} - k_y \frac{\partial V}{\partial y}. \quad (11)$$

Chen et al. (2005) demonstrated that the changes in wavenumber encountered by mountain waves propagating through a translating barotropic jet were sufficient to change the magnitude of the horizontally averaged momentum flux. In future work, it would be interesting to use ray tracing to follow the mountain waves as they are modified while propagating through the complex large-scale flow to test whether such WKB effects are responsible for filling out the  $k^{-5/3}$  KE spectrum.

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## APPENDIX

## Terms in the Spectral KE Budget

Following Peng et al. (2015a), for two scalar fields  $a$  and  $b$  define  $(a, b)_k = \hat{a}^*(\mathbf{k})\hat{b}(\mathbf{k})$ , where  $\hat{a}$  is the Fourier transform of  $a$  and the asterisk denotes the complex conjugate. For the vector fields  $\mathbf{a}$  and  $\mathbf{b}$ , define  $(\mathbf{a}, \mathbf{b})_k = \hat{\mathbf{a}}^*(\mathbf{k}) \cdot \hat{\mathbf{b}}(\mathbf{k})$ . Using this notation and (1),

$$\partial_t E_h(\mathbf{k}) = \frac{\bar{\rho}}{2}[(\mathbf{u}, \partial_t \mathbf{u})_k] + \text{c.c.}, \quad (\text{A1})$$

where “c.c.” denotes the complex conjugate of the preceding expression. Letting  $\nabla$  be the 2D horizontal gradient operator, the horizontal momentum equation may be well approximated as

$$\partial_t \mathbf{u} = -\mathbf{u} \cdot \nabla \mathbf{u} - w \partial_z \mathbf{u} - c_p \bar{\theta} \nabla \pi' + \mathbf{F}_h, \quad (\text{A2})$$

where  $\mathbf{u}$  is the horizontal velocity vector,  $\mathbf{F}_h$  is the turbulent subgrid-scale diffusivity,  $\bar{\theta}(z)$  is a vertically varying approximation to the full  $\theta(x, y, z, t)$ , and  $\pi$  is the Exner function, which is split between a horizontally uniform component  $\bar{\pi}(z)$  in hydrostatic balance with  $\bar{\theta}(z)$  and the remainder  $\pi'(x, y, z, t)$ . Substituting (A2) into (A1) gives

$$\partial_t E_h(\mathbf{k}) = A(\mathbf{k}) + P(\mathbf{k}) + D(\mathbf{k}). \quad (\text{A3})$$

Here  $A$  is the tendency arising from advection,

$$A(\mathbf{k}) = -\frac{\bar{\rho}}{2}(\mathbf{u}, \mathbf{u} \cdot \nabla \mathbf{u} + w \partial_z \mathbf{u})_k + \text{c.c.}, \quad (\text{A4})$$

$P$  is from the pressure gradient force,

$$P(\mathbf{k}) = -\frac{\bar{\rho}}{2}c_p \bar{\theta}(\mathbf{u}, \nabla \pi')_k + \text{c.c.}, \quad (\text{A5})$$

and  $D$  is from dissipation,

$$D(\mathbf{k}) = \frac{\bar{\rho}}{2}(\mathbf{u}, \mathbf{F}_h)_k + \text{c.c.} \quad (\text{A6})$$

As detailed in Peng et al. (2014), Peng et al. (2015a), and Sun et al. (2017), the advection term can be decomposed as

$$A(\mathbf{k}) = T(\mathbf{k}) + V_a(\mathbf{k}) + \varepsilon_1(\mathbf{k}), \quad (\text{A7})$$

where

$$T(\mathbf{k}) = -\bar{\rho} \left\{ \left[ \mathbf{u}, \mathbf{u} \cdot \nabla \mathbf{u} + \frac{1}{2} \mathbf{u} (\nabla \cdot \mathbf{u}) \right]_k - \frac{1}{2} (\partial_z \mathbf{u}, w \mathbf{u})_k + \frac{1}{2} (\mathbf{u}, w \partial_z \mathbf{u})_k \right\} + \text{c.c.}, \quad (\text{A8})$$

represents that component of the transfer of energy between wavenumbers that is conservative in the sense that its sum over all  $\mathbf{k}$  is zero. All derivatives are computed using a sixth-order compact difference [Lele 1992; Durran 2010, Eq. (3.50), p. 115]; horizontal derivatives are computed on constant  $z$  surfaces. Our numerical formulas are, therefore, not algebraically identical to the corresponding terms in the WRF Model.

In (A7),

$$V_a(\mathbf{k}) = -\frac{1}{2} \partial_z [\bar{\rho}(\mathbf{u}, w \mathbf{u})_k] + \text{c.c.}, \quad (\text{A9})$$

is the divergence of the vertical advective energy flux, and

$$\varepsilon_1(\mathbf{k}) = \frac{\bar{\rho}}{2} \{ [\mathbf{u}, \mathbf{u} (\nabla \cdot \mathbf{u} + \partial_z w)]_k + \partial_z \ln(\bar{\rho})(\mathbf{u}, w \mathbf{u})_k \} + \text{c.c.} \quad (\text{A10})$$

If the anelastic continuity equation,  $\bar{\rho} \nabla \cdot \mathbf{u} + \partial_z (\bar{\rho} w) = 0$ , is satisfied,  $\varepsilon_1(\mathbf{k})$  is zero. Although we are not using an anelastic model,  $\varepsilon_1(\mathbf{k})$  is quite small in our simulations and will be neglected.

Peng et al. (2014) used the pseudo-incompressible approximation (Durran 1989) to write the contribution from the pressure gradient force as

$$P(\mathbf{k}) = c_p \bar{\rho} \bar{\theta}(w, \partial_z \pi')_k - c_p \partial_z [\bar{\rho} \bar{\theta}(w, \pi')_k] + \text{c.c.} \quad (\text{A11})$$

Under the hydrostatic approximation,  $\partial_z \pi'$  is proportional to  $\theta'$  and the first term in (A11) represents buoyancy forcing, which will be subsequently written as

$$B(\mathbf{k}) = c_p \bar{\rho} \bar{\theta}(w, \partial_z \pi')_k + \text{c.c.} \quad (\text{A12})$$

The second term in (A11) is the divergence of the vertical flux of energy due to pressure working and will be combined with  $V_a(\mathbf{k})$  to give the total vertical energy flux divergence

$$V(\mathbf{k}) = -\frac{1}{2} \partial_z [\bar{\rho}(\mathbf{u}, w \mathbf{u})_k] - c_p \partial_z [\bar{\rho} \bar{\theta}(w, \pi')_k] + \text{c.c.} \quad (\text{A13})$$

The expression for  $\mathbf{F}_h$  in the dissipation term, (A6), is evaluated as

$$K_h \nabla^2 \mathbf{u} + K_v \frac{\partial^2 \mathbf{u}}{\partial z^2}, \quad (\text{A14})$$

where  $K_h$  is the horizontal eddy viscosity computed by the WRF model from a function of the horizontal deformation and  $K_v$  is vertical eddy viscosity computed



from the boundary layer scheme. Because in most cases the dissipation is insignificant, we used the easy-to-compute approximate expression (A14) rather than the full expression for diffusion in physical space used in the actual WRF simulations. The correct full expression would compute the divergence of the deviatoric stress tensor and would move the eddy viscosities inside that divergence operator.

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