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3 **Modeling the Load of SARS-CoV-2 Virus in Human Expelled Particles**  
4 **during Coughing and Speaking**  
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6 Yang Wang,<sup>1\*</sup> Guang Xu,<sup>2</sup> Yue-Wern Huang<sup>3</sup>  
7

8 <sup>1</sup>Department of Civil, Architectural and Environmental Engineering  
9 Missouri University of Science and Technology, Rolla, MO, 65401  
10

11 <sup>2</sup>Department of Mining and Nuclear Engineering,  
12 Missouri University of Science and Technology, Rolla, MO, 65401  
13

14 <sup>3</sup>Department of Biological Sciences  
15 Missouri University of Science and Technology, Rolla, MO, 65401  
16

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21 \* To whom correspondence should be addressed:

22 Tel: +1-573-341-4597

23 E-mail address: [yangwang@mst.edu](mailto:yangwang@mst.edu)

24 **Abstract**

25 Particle size is an essential factor when considering the fate and transport of virus-containing  
26 droplets expelled by human, because it determines the deposition pattern in the human  
27 respiratory system and the evolution of droplets by evaporation and gravitational settling.  
28 However, the evolution of virus-containing droplets and the size-dependent viral load have not  
29 been studied in detail. The lack of this information leads to uncertainties in understanding the  
30 airborne transmission of respiratory diseases, such as the COVID-19. In this study, through a set  
31 of differential equations describing the evolution of respiratory droplets and by using the SARS-  
32 CoV-2 virus as an example, we investigated the distribution of airborne virus in human expelled  
33 particles from coughing and speaking. More specifically, by calculating the vertical distances  
34 traveled by the respiratory droplets, we examined the number of viruses that can remain airborne  
35 and the size of particles carrying these airborne viruses after different elapsed times. From a  
36 single cough, a person with a high viral load in respiratory fluid ( $2.35 \times 10^9$  copies per ml) may  
37 generate as many as  $1.23 \times 10^5$  copies of viruses that can remain airborne after 10 seconds,  
38 compared to 386 copies of a normal patient ( $7.00 \times 10^6$  copies per ml). Masking, however, can  
39 effectively block around 94% of the viruses that may otherwise remain airborne after 10 seconds.  
40 Our study found that no clear size boundary exists between particles that can settle and can  
41 remain airborne. The results from this study challenge the conventional understanding of disease  
42 transmission routes through airborne and droplet mechanisms. We suggest that a complete  
43 understanding of the respiratory droplet evolution is essential and needed to identify the  
44 transmission mechanisms of respiratory diseases.  
45  
46 Keywords: COVID-19, aerosol, model, airborne transmission, coughing, speaking

47 **Introduction**

48 The ongoing pandemic of COVID-19 highlights the urgent need to understand the transport and  
49 evolution of pathogen-containing aerosols and droplets, because there are contradictory evidence  
50 and conclusions on the potential transmission route of SARS-CoV-2 [1-7]. At the very beginning  
51 of the disease outbreak, the World Health Organization (WHO) [8] and Centers for Disease  
52 Control and Prevention (CDC) [9] stated that the transmission of SARS-CoV-2 through the  
53 airborne route, which is by inhaling virus-containing aerosols, is unlikely. Instead, droplet  
54 transmission, which is through exposure to respiratory droplets, and contact transmission, which  
55 is the infection through direct or indirect contact with an infected person, are believed to be the  
56 major transmission routes. The traditional distinction between a “droplet” and an “aerosol (or  
57 droplet nuclei)” is based on size, where droplets are suspended particles above 5  $\mu\text{m}$  in diameter,  
58 and aerosols are those below 5  $\mu\text{m}$  [10]. To avoid confusion, in this study, we will use  
59 “particles” to refer to a summation of “aerosols” and “droplets.” It is thought that droplets can  
60 settle to ground in a few seconds, but aerosols can remain airborne for an extended period of  
61 time. Although there is no such definition in atmospheric studies, this traditional distinction  
62 between droplets and aerosols has been useful for setting clinical guidelines on the use of  
63 personal protective equipment for healthcare workers [11]. However, the conventional  
64 distinction between aerosols and droplets has led to a “false dichotomy” [12] in understanding  
65 airborne pathogens, because whether a respiratory particle can remain airborne depends on many  
66 factors.

67

68 Existing studies show that human activities such as coughing, sneezing, and speaking generate  
69 particles, with more than 90% of the total particle numbers less than 5  $\mu\text{m}$  after evaporation [13-  
14].

70 17]. Evaporation can significantly extend the dispersion lifetime of particles before they settle,  
71 enhancing the infection risk of airborne viruses. For example, the sizes of the largest droplets  
72 that would totally evaporate before settling 2 m are between 60 and 100  $\mu\text{m}$ , and these expelled  
73 large droplets are carried more than 6 m away by exhaled air at a velocity of  $50 \text{ m s}^{-1}$  (sneezing),  
74 more than 2 m away at a velocity of  $10 \text{ m s}^{-1}$  (coughing) and less than 1 m away at a velocity of  
75  $1 \text{ m s}^{-1}$  (breathing) [14]. Many of these existing studies, including a recent one [18] investigated  
76 the droplet lifetime influenced by the ambient temperature and humidity using the evaporating  
77 drop mathematical model, but the virus contained in the particles, and the associated viral load as  
78 a function of particle size were not included in the model. This particle size-dependent viral load  
79 is crucial to our understanding of the relative importance of airborne and droplet transmission  
80 because if a significant number of viruses remain in airborne, appropriate precautions should be  
81 taken, such as universal masking, stronger indoor ventilation rate, and air disinfection. Until  
82 now, more evidence is also showing that similar to other pathogens such as influenza viruses and  
83 *Mycobacterium tuberculosis* [19], SARS-CoV-2 can be carried by aerosols [20-25].

84

85 Theoretically, coughing, sneezing, and speaking generate particles by aerosolizing the  
86 respiratory fluid, and the number of viruses in a particle is determined by the viral concentration  
87 in the respiratory fluid and the volume of the particle. Therefore, the number of viruses in a  
88 single particle should scale with the cube of the particle diameter. Based on the typical  
89 concentration of the SARS-CoV-2 viruses in respiratory fluid [26], one can calculate that a  
90 considerable number of human expelled particles do not contain viruses due to their small  
91 volume. During the evolution of the respiratory droplets, evaporation complicates the size-  
92 dependent viral load in aerosols and droplets, as the size of the particles changes with time.

93 Gravitational settling will remove larger droplets that contain more viruses. Collectively, they  
94 ensure the necessity to examine the load of viruses in human expelled particles of different sizes.

95

96 Using the most recent SARS-CoV-2 data, this study used the Monte-Carlo method to simulate  
97 the particles generated from coughing and speaking and used a Poisson distribution function to  
98 determine the virus load in the particles. The particle size-dependent viral load and its variation  
99 as a function of time during evaporation and gravitational settling are modeled using mass and  
100 heat transfer equations and the momentum equation. The detailed modeling methods are  
101 elaborated in the Methods section. In the Results and Discussion section, we show that most of  
102 the virus-containing particles can remain airborne for an extended period of time longer than 10  
103 seconds. We analyzed how the elapsed time and viral load in the respiratory fluid affect the  
104 transport of the virus-containing particles, and examined the particle emission from coughing  
105 and speaking. Finally, we discuss the uncertainties associated with this analysis.

106

## 107 **Methods**

### 108 **Size distributions of human expelled particles**

109 Accurate size distributions of human expelled droplets are required to estimate the particle size-  
110 dependent viral load. Existing studies commonly used an Aerodynamic Particle Sizer (APS, TSI  
111 Inc.) to measure the size distributions of human-emitted droplets [16, 27-29]. However, droplets  
112 will evaporate during their transport in the measurement setup, leading to uncertainties in  
113 measuring the original droplet sizes. The size distributions of directly emitted droplets can be  
114 more accurately measured by in-situ light scattering experiments conducted near the human  
115 mouth [30, 31]. In this study, we adopted such droplet size distributions measured by Chao et al.

116 [30], where speaking generates particles with a geometric mean diameter ( $D_{d,g}$ ) of 16.0  $\mu\text{m}$  and a  
 117 geometric standard deviation ( $\sigma_{d,g}$ ) of 0.55, and coughing generates particles with a  $D_{d,g}$  of 13.5  
 118  $\mu\text{m}$  and a  $\sigma_{d,g}$  of 0.50. We further assume that speaking and coughing generate a total number  
 119 ( $N_d$ ) of 50 per second and 3000 per cough, respectively [16, 30]. The droplet size ( $D_d$ ) follows a  
 120 lognormal size distribution, where

$$n_d(D_d) = \frac{N_d}{\sqrt{2\pi}\ln(\sigma_{d,g})} \exp\left[-\frac{(\ln(D_d)-\ln(D_{d,g}))^2}{2(\ln(\sigma_{d,g}))^2}\right]. \quad (1)$$

121 We adopted a Monte-Carlo method to randomly generate  $N_d$  number of droplets following the  
 122 lognormal size distribution. The number of viruses in a droplet with a size of  $D_d$  can be  
 123 calculated by

$$VL(D_d) = \frac{\pi}{6} D_d^3 C_V, \quad (2)$$

124 where  $C_V$  is the viral load of SARS-CoV-2 in the respiratory fluid. Existing studies show that  $C_V$   
 125 has an average value of  $7.00 \times 10^6$  copies per ml, with a maximum of  $2.35 \times 10^9$  copies per ml,  
 126 which is largely dependent on the number of days after onset of symptoms [32]. We further  
 127 assume that the liquid content of the respiratory fluid is composed of 0.9% NaCl-water solution.  
 128 Therefore, after evaporation, the expelled particles can leave a solid core containing salt and  
 129 viruses, which is a more realistic model of respiratory particles.

130

131 We should note that the number of viruses calculated by Eq. (2) is hardly an integer.  $VL(D_d)$   
 132 reflects the expected number of viruses in a droplet, but the actual number will take integer  
 133 values above or below  $VL(D_d)$ . To reflect the randomness of this process, we assume that the  
 134 actual number of viruses enclosed in a droplet follows the Poisson distribution [33]. We have

$$f(x) = \frac{(VL(D_d))^x}{x!} \exp[-VL(D_d)]. \quad (3)$$

135 In this equation,  $f(x)$  is the probability the droplet with a size  $D_d$  containing exactly  $x$  ( $x =$   
136 0, 1, 2, ... ) number of viruses.

137

## 138 **Evaporation and gravitational settling**

139 After being emitted, a droplet undergoes evaporation and gravitational settling. The size of the  
140 droplet is determined by the following mass and heat transfer equations:

$$\dot{m}_d = \rho_d \frac{d}{dt} \left( \frac{\pi}{6} D_d^3 \right) = -A_d h_m (p_{v,s} - p_{v,\infty}), \text{ and} \quad (4)$$

$$m_d C_{pd} \frac{dT_d}{dt} = A_d h (T_\infty - T_d) + L \dot{m}_d. \quad (5)$$

141 The droplet evaporation rate  $\dot{m}_d$  is driven by the difference between the vapor pressure in the  
142 surrounding air  $p_{v,\infty}$  and the vapor pressure at the droplet surface  $p_{v,s}$ .  $p_{v,s}$  is assumed as  
143 saturated vapor pressure at droplet temperature  $T_d$ , considering the Kelvin and Raoult effects.  $A_d$   
144 is the droplet surface area,  $L$  is the latent heat of vaporization, and  $C_{pd}$  is the heat capacity of the  
145 droplet. The mass transfer coefficient  $h_m$  and the heat transfer coefficient  $h$  can be solved using  
146 the Ranz-Marshall correlations for the Sherwood and Nusselt numbers [34]. The ultimate droplet  
147 size is determined by the solid components in the droplet. Previous studies on respiratory droplet  
148 evaporation commonly ignored the influence of microorganisms enclosed in the droplet, leading  
149 to an underestimate of the final particle size and overestimate of the particle lifetime. In this  
150 model simulation, we further considered the influence of SARS-CoV-2 on the physical size of  
151 the evaporated droplet, by assuming that the enclosed SARS-CoV-2 virus has a spherical shape  
152 and diameter of 100 nm (65 to 125 nm according to Astuti et al. [35]) and a density of 1.35 g cm<sup>-</sup>  
153 <sup>3</sup>, similar to common protein [36].

154

155 The gravitational settling of the human expelled particles can be solved by the momentum  
156 balance equation, where

$$m_d \frac{d^2z}{dt^2} = \frac{1}{2} \rho_g V_z^2 A_d C_D. \quad (6)$$

157 In Eq. (6),  $z$  is the droplet settling distance,  $\rho_g$  is ambient air density,  $V_z$  is droplet velocity in the  
158 vertical direction,  $A_d$  is the cross section area of the droplet ( $A_d = \frac{\pi}{4} D_d^2$ ), and  $C_D$  is the drag  
159 coefficient, which is dependent on the Reynolds number of the particle motion [37]. In this  
160 study, we focus on the vertical movement of the particles in order to estimate whether the  
161 particles can remain airborne after different elapsed time. The horizontal movement of the  
162 particles will largely depend on the activity that generates the particles, and they will be  
163 examined briefly at the end of the analyses.

164

165 The differential equations in Eqs. (4-6) can be solved simultaneously, where the droplet  
166 diameter, droplet surface temperature, and droplet settling distance can be derived as a function  
167 of time. Assuming that these human expelled droplets are generated at the height of 1.7 m with  
168 no initial vertical velocity, we can further calculate the lifetime of a droplet, which is the time  
169 corresponding to  $z = 1.7$  m. For all the calculations, we assume an indoor environmental  
170 condition, where the temperature is 23 °C and the relative humidity is 50%. Conceivably,  
171 temperature and relative humidity can affect the droplet evolution through evaporation, as shown  
172 in Chen 2020a [18]. Moreover, they will likely influence the viability of viruses and, thereby the  
173 infection risk [38], which is discussed at the end of the following section. However, this study  
174 focuses on modeling the number of viruses that can remain airborne after being emitted by the  
175 patient. We should note that there are other modeling methods available to understand the  
176 dynamics associated with biological and physical systems [39-41]. In this work, we used the

177 relatively simplified differential equations to understand the transport of the virus-containing  
178 aerosols and estimate the load of viruses in human expelled particles.

179

## 180 **Results and Discussion**

181 In the following analysis, we demonstrate how the airborne viral load depends on the size of the  
182 human expelled particles and its variation as a function of time. We first analyze the load of the  
183 airborne virus on particles generated from a single cough, and then examine its dependence on  
184 elapsed time and the viral load in the respiratory fluid. We also compare the airborne viral load  
185 associated with speaking against that of coughing.

186

### 187 **Droplet properties at the point of emission**

188 Fig 1 shows an example solution demonstrating the evolution of droplets generated during a  
189 single cough. Fig 1a displays the size distribution of 3000 coughing droplets randomly generated  
190 following the lognormal distribution in Eq. (1). At a viral load of  $7.00 \times 10^6$  copies per ml in the  
191 respiratory fluid, viruses are mostly contained in droplets larger than 10  $\mu\text{m}$ , because the product  
192 of the droplet volume and the viral concentration in smaller droplets is far below 1. Among the  
193 3000 droplets generated by a single cough, approximately  $390 \pm 16$  droplets contain viruses, and  
194 the total number of viruses in these virus-containing droplets is  $9.8 \times 10^3 \pm 6.4 \times 10^3$  copies  
195 (Table 1). This large standard deviation is a result of a few giant droplets, which contain a  
196 substantial number of viruses. However, these giant droplets are also subject to rapid removal by  
197 gravitational settling as time progresses.

198

199 **Fig 1. Evolution of droplets emitted by a cough over an elapsed time of ten seconds at**  
200 **respiratory viral loads of (a – c)  $7.00 \times 10^6$  and (d – f)  $2.35 \times 10^9$  copies per ml.** (a) and (d)  
201 Size distribution of droplets and virus-containing droplets at point of emission. (b) and (e) Size  
202 distribution of non-virus-containing (airborne), virus-containing (airborne), and settled particles  
203 at an elapsed time of ten seconds. (c) and (f) Distribution of vertical distances traveled by the  
204 virus-containing particles at an elapsed time of ten seconds. The inset figure in panel (c) shows a  
205 schematic of the modeled system.

206

207 **Table 1.** Number of virus-containing particles and number of viral copies remain suspended in  
208 the air after different elapsed times in a cough.

	Viral load in respiratory fluid (copies per ml)	
	$7.00 \times 10^6$	$2.35 \times 10^9$
Virus-containing droplets after 0 s	$390 \pm 16$	$2021.6 \pm 22.4$
Viral copies after 0 s	$9.8 \times 10^3 \pm 6.4 \times 10^3$	$2.6 \times 10^6 \pm 1.7 \times 10^6$
Virus-containing particles after 1 s	$380 \pm 6$	$2017 \pm 25$
Viral copies after 1 s	$4.4 \times 10^3 \pm 0.7 \times 10^3$	$1.33 \times 10^6 \pm 0.11 \times 10^6$
Virus-containing particles after 3 s	$349 \pm 16$	$1990 \pm 23$
Viral copies after 3 s	$1.2 \times 10^3 \pm 0.1 \times 10^3$	$4.15 \times 10^5 \pm 0.11 \times 10^5$
Virus-containing particles after 10 s	$250 \pm 7$	$1855 \pm 13$
Viral copies after 10 s	$386 \pm 7$	$1.23 \times 10^5 \pm 0.05 \times 10^5$
Virus-containing particles after 30 s	$232 \pm 14$	$1871 \pm 7$
Viral copies after 30 s	$333 \pm 12$	$1.13 \times 10^5 \pm 0.01 \times 10^5$

209

210

## 211 **Effect of elapsed time**

212 After ten seconds of evaporation and gravitational settling, the peak size of the expelled particles  
213 shifted to around 2.2  $\mu\text{m}$  (Fig 1b). Due to the salt and viruses in the droplet, the virus-containing  
214 particles now have a size above 2  $\mu\text{m}$ . Approximately 5.1% of virus-containing particles are  
215 below 5  $\mu\text{m}$ , which traditionally would be categorized as "aerosols." The number of viruses  
216 contained in these sub-5  $\mu\text{m}$  particles is  $20 \pm 2$  copies. However, 59.5% of virus-containing  
217 particles remain airborne (settle less than 1.7 m), and the number of viruses contained in the  
218 evaporated droplets is  $386 \pm 7$  copies. This result shows that one cannot simply use a specific  
219 size to determine whether a respiratory particle settle or remain airborne. Droplet evaporation  
220 and heat transfer over time need to be incorporated to be more accurately depict the respiratory  
221 particle behavior. Fig 1c also shows the vertical distance traveled by the virus-containing  
222 particles and the number of viruses contained in the particles after ten seconds of droplet  
223 emission. It demonstrates that around 80% of the virus-containing particles settle with a vertical  
224 distance within 0.5 m, meaning that these suspended particles can linger in the inhalation zone of  
225 people surrounding the patient.

226

227 The number of viral copies contained in the particles decreases rapidly with the elapsed time,  
228 from  $9.8 \times 10^3$  at the point of emission to  $333 \pm 12$  at an elapsed time of 30 s. It is because larger  
229 particles that enclose more viral copies settle faster (Fig 1b). On the other hand, the number of  
230 virus-containing particles that remain airborne is relatively insensitive to elapsed time, from 390  
231  $\pm 16$  at the point of emission to  $232 \pm 14$  at 30 s. This insensitivity is caused by the fact that most

232 of the virus-containing droplets shrink to sizes that cannot be effectively settled by gravity.

233 Therefore, these particles will have a longer lifetime and pose a higher infection risk.

234

## 235 **Effect of viral load in respiratory fluid**

236 The viral load in the respiratory fluid drastically affects the evolution of human expelled virus-  
237 containing particles (Figs. 1d-1f). At a viral load of  $2.35 \times 10^9$  copies per ml, droplets as small as  
238 4  $\mu\text{m}$  start to contain viruses (Fig 1d), and around 67.4% of droplets contain viruses. The fraction  
239 of virus-containing particles remaining airborne after an elapsed time of ten seconds is also high  
240 (Fig 1e), reaching 61.8%. Again, it is not realistic to use a cut-off size of 5  $\mu\text{m}$  to differentiate  
241 “aerosols” from “droplets.” Due to the high viral load in the respiratory fluid ( $2.35 \times 10^9$  copies  
242 per ml), the number of viral copies in the evaporated particles ( $1.23 \times 10^5$ ) is orders of magnitude  
243 higher compared to the average condition (386 under  $7.00 \times 10^6$  copies per ml). The vertical  
244 distribution of the virus-containing particles and the copies of viruses in Fig 2f show  
245 considerably higher values in shorter vertical distances (0 to 0.5 m), meaning that a patient with a  
246 higher viral load in the respiratory fluid would pose a significantly higher infection risk to the  
247 surrounding people.

248

## 249 **Airborne viral load during speaking**

250 Compared to coughing, speaking is a process that continuously generated respiratory droplets.

251 Therefore, when examining the evolution of droplets as a function of time, we need to consider

252 the droplets emitted at different times of speaking cumulatively. Fig 2a shows the properties of

253 droplets during one second of speaking at the point of emission for a patient with a viral load of

254  $2.35 \times 10^9$  copies per ml in the respiratory fluid. Due to the few numbers of droplets generated,

255 the droplet size distribution is subject to high uncertainty. Fig 2b shows the size distribution of  
256 speaking-generated particles ten seconds after a one-minute speech. The size distribution is not  
257 significantly different from that of coughing, as shown in Fig 1e. However, due to the longer  
258 elapsed time of particles emitted at the beginning of the speaking period (up to 70 seconds),  
259 particles of 20  $\mu\text{m}$  can settle down to the ground, compared to 40  $\mu\text{m}$  for coughing. However, the  
260 vertical distribution of the numbers of virus-containing droplets and viral copies still show higher  
261 numbers in shorter vertical distances (0 to 0.5 m), meaning that a considerable fraction of  
262 speaking-generated droplets can remain airborne due to evaporation.

263

264 **Fig 2. Evolution of droplets emitted by one-minute of speaking after an elapsed time of ten**  
265 **seconds at a respiratory viral load of  $2.35 \times 10^9$  copies per ml.** (a) Size distribution of droplets  
266 and virus-containing droplets at point of emission during one-second of speaking. (b) Size  
267 distribution of non-virus-containing (airborne), virus-containing (airborne), and settled particles  
268 at an elapsed time of ten seconds. (c) Distribution of vertical distances traveled by the virus-  
269 containing particles at an elapsed time of ten seconds. (d) Size-dependent filtration efficiency  
270 curves for a surgical mask (earloop) extracted from Chen et al. (1992) and Hao et al. (2020). (e)  
271 Size distribution of non-virus-containing (airborne), virus-containing (airborne), and settled  
272 particles at an elapsed time of ten seconds with mask-wearing. (f) Distribution of vertical  
273 distances traveled by the virus-containing particles at an elapsed time of ten seconds with mask-  
274 wearing.

275

276 **Effect of mask-wearing**

277 Using the proposed model, we could also evaluate the effectiveness of face masks in preventing  
278 the spread of viruses. Fig 2d shows the size-dependent filtration efficiency of aerosols from 0.03  
279 to 10  $\mu\text{m}$  for common surgical mask materials [42, 43]. Due to the combined mechanisms of  
280 inertial impaction, interception, Brownian diffusion, and electrostatic interaction, the filtration  
281 efficiency curves generally show an “escape window” where particles with hundreds of  
282 nanometers can penetrate through the filter, resulting in lower efficiencies. Existing literature  
283 also uses the term “most penetrating particle size (MPPS)” to describe the reduced filtration  
284 efficiency in this size range [44]. Unlike medical respirators, face masks have the issue of flow  
285 leakage between the mask and the wearer [45]. Here, we assume a flow leakage of 5%, and  
286 calculated the evolution of droplets generated from speaking using the average filtration  
287 efficiency in Fig 2d. The numbers of both the non-virus-containing and virus-containing droplets  
288 reduced significantly (Fig 2e) compared to the unmasked speaking (Fig 2b), with the total  
289 number of airborne virus-containing droplets decreased by 94.9% (from  $2122 \pm 17$  to  $108 \pm 5$ ),  
290 and with the total number of viral copies decreased by 95.6% (from  $1.4 \times 10^5 \pm 0.1 \times 10^5$  to  $6.2 \times$   
291  $10^3 \pm 0.2 \times 10^3$ ). Although the number of virus-containing particles is still the highest near the  
292 point of emission (within the vertical distance of 0.5 m, Fig 2f), the number of viral copies  
293 decreased significantly within this distance. Due to the effective removal of virus-containing  
294 particles, the vertical distribution of the number of viral copies becomes more random, and the  
295 two peaks in the distance between 1 and 1.7 m in Fig 2f are caused by a few large droplets that  
296 escaped from the mask. Compared to the unmasked condition (Fig 2b), the number fraction of  
297 evaporated particles below 1  $\mu\text{m}$  becomes higher under the masked condition (Fig 2e), mainly  
298 due to the lower filtration efficiencies of the masks for particles between 0.1 and 1  $\mu\text{m}$ .

299

300 **Uncertainties associated with the analysis**

301 The above analysis shows that a significant fraction of respiratory droplets can remain airborne  
302 after they are emitted. Note that the horizontal movement of the droplets is not shown in this  
303 study, because the horizontal velocity of respiratory droplets depends strongly on human activity,  
304 age, and ambient environment [46-48]. The trajectory of the exhaled respiratory droplets is  
305 affected by both the expired air flows profile and surrounding air flow patterns. Existing studies  
306 treated the exhaled air as a turbulent round jet [49, 50], and the turbulent flow will enhance the  
307 heat and mass transfer between the droplet and the surrounding air. Therefore, respiratory  
308 droplets will likely evaporate faster than the simulated results in this study, and a larger fraction  
309 of respiratory droplets and viruses may remain airborne for a longer period of time. Here, we  
310 adopt a simplified flow field derived from a previous experimental study [51], where the  
311 horizontal velocity of air expelled from coughing follows the equation

$$V_x = 0.875/(l_x + 0.333)^2. \quad (7)$$

312 In Eq. (7),  $V_x$  is the velocity of the respiratory droplet in the horizontal direction in  $\text{m s}^{-1}$  when  
313 there is no ambient air flow and  $l_x$  is the horizontal distance from the point of emission in m.  
314 According to this relationship, the distance traveled by the respiratory droplets as a function of  
315 time can be derived as:

$$l_x = \sqrt[3]{(2.625t + 0.0369)} - 0.333. \quad (8)$$

316 According to this simplified solution, airborne droplets can travel a horizontal distance of 2.64 m  
317 after 10 s, and 3.95 m after 30 s. Considering that virus-containing particles can remain airborne  
318 after 30 seconds (Table 1), the “six-feet (or 2 m) rule” is not sufficient in preventing disease  
319 transmission. Nonetheless, universal masking may be a better option for disease transmission, as

320 it can capture the respiratory droplets effectively through impaction and interception at the  
321 source of generation [43, 52].

322

323 In this study, we did not consider the viability of viruses in particles with different sizes. Since  
324 pathogen viability is dependent on the surface properties of materials [53], the viability of  
325 viruses in droplets may also change as a function of time, because evaporation continuously  
326 increases the droplet surface tension and expose the components of the droplet to the surface of  
327 the droplet. For example, virus deactivation may occur after exposure to the air-water interface,  
328 where irreversible rearrangement and folding of the viruses' protein take place [54, 55].  
329 Moreover, the distribution of viruses in droplets of different sizes may not be uniform. For  
330 example, studies on airborne virus sampling show that viable viruses tend to be sampled in  
331 particles below 5  $\mu\text{m}$  [56, 57]. One possible explanation is that droplets of different sizes may  
332 originate from different regions of the respiratory system, where smaller droplets are formed  
333 from regions of a higher viral load. The measurement of virus-laden aerosols in outbreaks in  
334 farms also indicated that certain viruses tend to be associated with particles below 0.4  $\mu\text{m}$  [58],  
335 which may be due to the mechanism of aerosol generation. Therefore, future work can further  
336 study how the expired air flows and size-dependent viability of the viruses affect the  
337 concentration of the airborne viruses generated from coughing and speaking.

338

## 339 Conclusion

340 In this work, we investigated the dependence of airborne viral load on the size distributions of  
341 the human expelled particles. We found that differentiating "aerosols" and "droplets" using a  
342 specific size, e.g., 5  $\mu\text{m}$ , does not reflect the actual evolution of virus-containing particles over

343 time and space, because a large number of particles above 5  $\mu\text{m}$  can remain airborne after an  
344 extended period of time. Our simulation result showed that after ten seconds of a cough,  
345 although most evaporated particles are larger than 5  $\mu\text{m}$ , 59.5% of the original virus-containing  
346 particles are still able to remain airborne. Although the numbers of airborne viral copies and  
347 virus-containing particles decrease with elapsed time, this dependence becomes weaker at long  
348 elapsed times due to the significantly longer residence time of the smaller particles. We further  
349 show that a high viral load in the respiratory fluid will lead to a significantly higher infection risk  
350 due to the large number of virus-containing aerosols that remain airborne after an extended  
351 elapsed time. Our simulation also shows that wearing a mask can effectively reduce the spread of  
352 the viruses. The simulation results challenge the false dichotomy of using aerosols and droplets  
353 to separate the modes of disease transmission.

354

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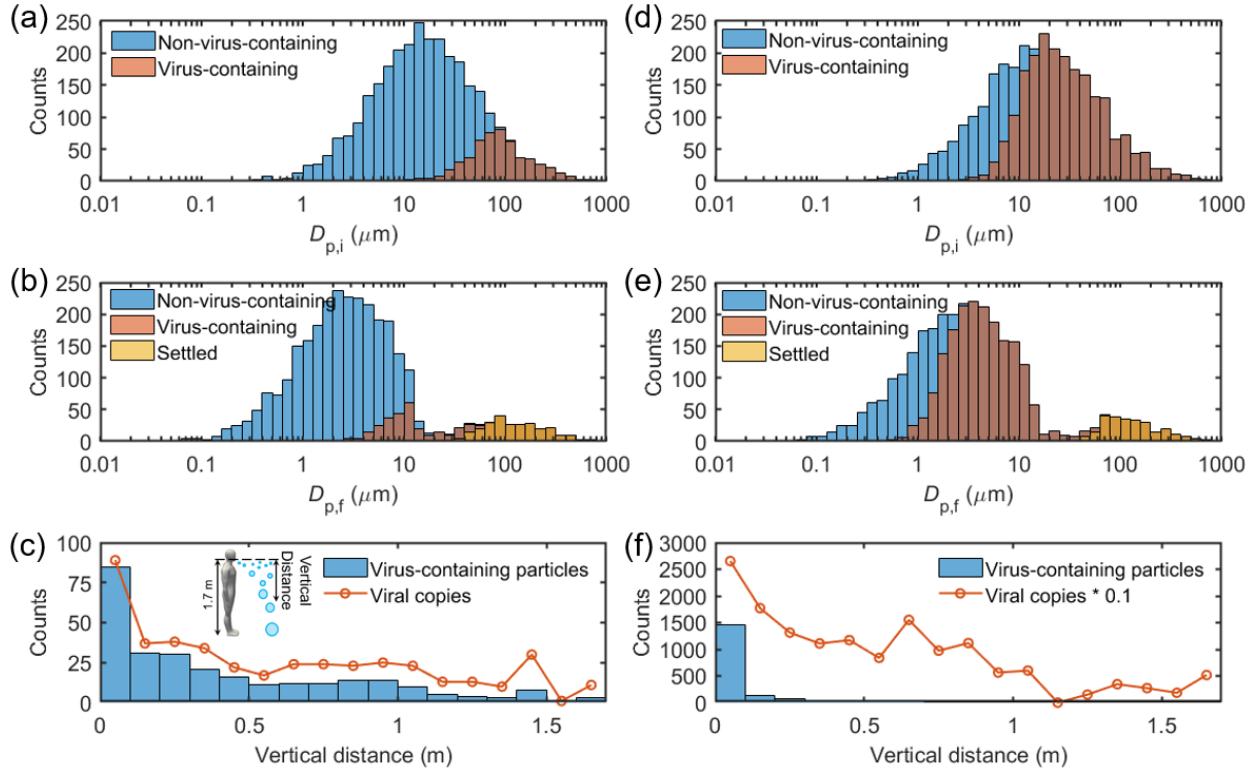
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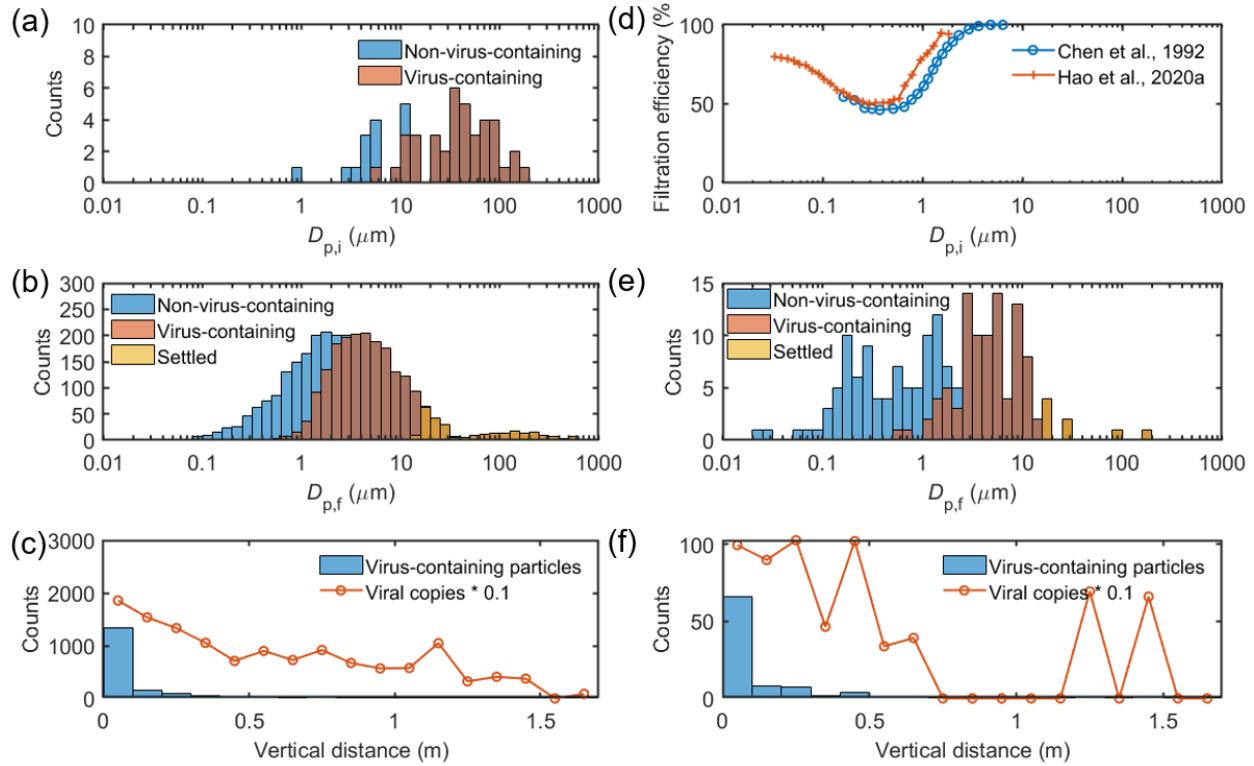
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513 **Fig. 1.** Evolution of droplets emitted by a cough over an elapsed time of ten seconds at  
 514 respiratory viral loads of (a – c)  $7.00 \times 10^6$  and (d – f)  $2.35 \times 10^9$  copies per ml. (a) and (d) Size  
 515 distribution of droplets and virus-containing droplets at point of emission. (b) and (e) Size  
 516 distribution of non-virus-containing (airborne), virus-containing (airborne), and settled particles  
 517 at an elapsed time of ten seconds. (c) and (f) Distribution of vertical distances traveled by the  
 518 virus-containing particles at an elapsed time of ten seconds. The inset figure in panel (c) shows a  
 519 schematic of the modeled system.



520

521 **Fig. 2.** Evolution of droplets emitted by one-minute of speaking after an elapsed time of ten  
 522 seconds at a respiratory viral load of  $2.35 \times 10^9$  copies per ml. (a) Size distribution of droplets  
 523 and virus-containing droplets at point of emission during one-second of speaking. (b) Size  
 524 distribution of non-virus-containing (airborne), virus-containing (airborne), and settled particles  
 525 at an elapsed time of ten seconds. (c) Distribution of vertical distances traveled by the virus-  
 526 containing particles at an elapsed time of ten seconds. (d) Size-dependent filtration efficiency  
 527 curves for a surgical mask (earloop) extracted from Chen et al. (1992) and Hao et al. (2020). (e)  
 528 Size distribution of non-virus-containing (airborne), virus-containing (airborne), and settled  
 529 particles at an elapsed time of ten seconds with mask-wearing. (f) Distribution of vertical  
 530 distances traveled by the virus-containing particles at an elapsed time of ten seconds with mask-  
 531 wearing.

532