

Irreducibility of the Fermi Surface for Planar Periodic Graph Operators

WEI LI AND STEPHEN P. SHIPMAN¹

Department of Mathematics, Louisiana State University, Baton Rouge

Abstract. We prove that the Fermi surface of a connected doubly periodic self-adjoint discrete graph operator is irreducible at all but finitely many energies provided that the graph (1) can be drawn in the plane without crossing edges (2) has positive coupling coefficients (3) has two vertices per period. If “positive” is relaxed to “complex”, the only cases of reducible Fermi surface occur for the graph of the tetrakis square tiling, and these can be explicitly parameterized when the coupling coefficients are real. The irreducibility result applies to weighted graph Laplacians with positive weights.

Key words: Graph operator, Fermi surface, Floquet surface, reducible algebraic variety, planar graph, graph Laplacian

MSC: 47A75, 47B39, 39A70, 39A14, 39A12

1 Introduction

The Fermi surface (or Fermi curve) of a doubly periodic operator at an energy E is the analytic set of complex wavevectors (k_1, k_2) admissible by the operator at that energy. Whether or not it is irreducible is important for the spectral theory of the operator because reducibility is required for the construction of embedded eigenvalues induced by a local defect [10, 11, 14] (except, as for graph operators, when an eigenfunction has compact support). (Ir)reducibility of the Fermi surface has been established in special situations. It is irreducible for the discrete Laplacian plus a periodic potential in all dimensions [13] (see previous proofs for two dimensions [1], [7, Ch. 4] and three dimensions [2, Theorem 2]) and for the continuous Laplacian plus a potential that is separable in a specific way in three dimensions [3, Sec. 2]. Reducibility of the Fermi surface is attained for multilayer graph operators constructed by appropriately coupling discrete graph operators [14] or metric graph (quantum graph) operators [14, 15, 6].

The multilayer graphs in [14, 15] are inherently non-planar by their very construction—they cannot be rendered in the plane without crossing edges. This led us to ask whether planarity prohibits reducibility of the Fermi surface. In this work, we address this problem by direct computation, and we are able to give a complete answer for discrete graph operators with real coefficients and two vertices per period. It turns out that planarity together with positivity of the coefficients prohibits reducibility. Our results apply to discrete weighted graph Laplacians (Theorem 2), where the positivity of the weights are a discrete version of coercivity of second-order elliptic operators.

Additionally, we find that irreducibility occurs even for complex operator coefficients, including magnetic Laplacians, except for the planar graph whose faces are the tetrakis square tiling (Fig. 2). For that graph, reducibility occurs for all energies for certain non-positive choices of the coefficients. All the reducible cases can be explicitly parameterized when the coupling coefficients are real. Our results are stated in section 2.3.

Our study of the Fermi surface is part of a more general effort to understand fine properties of the dispersion function $D(k_1, k_2, E)$ of wavevector and energy, particularly using computational methods to attack problems that hitherto evade theoretical methods. Recently, computational techniques from algebraic geometry [5] have been used to analyze the genericity of the extrema of the zero set of D , also focusing on doubly periodic discrete graph operators with two vertices per period, as in our present work.

¹Corresponding author; shipman@lsu.edu

2 Periodic graph operators and their Fermi surfaces

Let Γ be a graph with vertex set $\mathcal{V}(\Gamma)$ and edge set $\mathcal{E}(\Gamma)$ drawn in the plane \mathbb{R}^2 , such that Γ is invariant with respect to shifts along two linearly independent vectors ξ_1 and ξ_2 . For any $n = (n_1, n_2) \in \mathbb{Z}^2$, denote the shift of a vertex $v \in \mathcal{V}(\Gamma)$ or an edge $e \in \mathcal{E}(\Gamma)$ along $n_1\xi_1 + n_2\xi_2$ by nv or ne . This realizes the shift group \mathbb{Z}^2 as a free group of isomorphisms of Γ . We assume from the beginning that $\mathcal{V}(\Gamma)$ has exactly two orbits; a fundamental domain contains one vertex from each orbit, and the set of these two vertices is denoted by $\dot{\mathcal{V}} = \{v_1, v_2\}$. Denote the \mathbb{Z}^2 -orbit of v_i by \mathcal{V}_i , so that $\mathcal{V}(\Gamma) = \mathcal{V}_1 \dot{\cup} \mathcal{V}_2$.

2.1 Periodic operators

This section sets the background and notation for periodic operators on the space $\ell^2(\mathcal{V}(\Gamma), \mathbb{C})$ of square-integrable functions from the vertex set of Γ to the complex field \mathbb{C} . Given an operator on $\ell^2(\mathcal{V}(\Gamma), \mathbb{C})$, the edges of the graph Γ correspond to nonzero matrix elements of this operator.

Any function $g : \mathcal{V} \rightarrow \mathbb{C}$ is identified with the vector $(g_1, g_2) \in \mathbb{C}^2$ with $g_j = g(v_j)$; and any function $\tilde{f} \in \ell^2(\mathcal{V}(\Gamma), \mathbb{C})$ is identified with a function $f \in \ell^2(\mathbb{Z}^2, \mathbb{C}^2)$ by $f(n)_j = \tilde{f}(nv_j)$. Denote the standard elementary basis vectors in \mathbb{C}^2 by $\{\epsilon_j\}_{j=1}^2$. A periodic, or shift-invariant, operator A on $\ell^2(\mathcal{V}(\Gamma), \mathbb{C}) \cong \ell^2(\mathbb{Z}^2, \mathbb{C}^2)$ is a convolution operator. Precisely, for each $n \in \mathbb{Z}^2$, let $A(n)$ be a 2×2 matrix. For $f \in \ell^2(\mathbb{Z}^2, \mathbb{C}^2)$, define

$$(Af)(n) = \sum_{m \in \mathbb{Z}^2} A(m)f(n-m). \quad (2.1)$$

The operator A is said to be of *finite extent* if there is a number M such that $A(n) = 0$ for $|n| > M$. A is self-adjoint if $A(-n) = A(n)^*$ for all $n \in \mathbb{Z}^2$. This class of operators includes the standard and magnetic discrete Laplacians (see [8, §2] and [9, eqn. (1.4)]).

The graph Γ is called the *graph associated to the operator* A provided each entry of A , as a matrix indexed by $\mathcal{V}(\Gamma)$, is nonzero if and only if the corresponding pair of vertices is connected by an edge in $\mathcal{E}(\Gamma)$. More precisely, for all $i, j \in \{1, 2\}$ and all $n \in \mathbb{Z}^2$,

$$\{v_i, nv_j\} \in \mathcal{E}(\Gamma) \iff (\epsilon_i, A(n)\epsilon_j) \neq 0. \quad (2.2)$$

Double edges do not make sense for discrete graph operators, but loops do.

2.2 The Fermi surface

The z -transform (or Floquet transform $\hat{f}(n, z)$ at $n = 0$) of a function $f \in \ell^2(\mathbb{Z}^2, \mathbb{C}^2)$ is the formal Laurent series

$$\hat{f}(z) = \sum_{m \in \mathbb{Z}^2} f(m)z^{-m}, \quad z = (z_1, z_2). \quad (2.3)$$

The notation is $z^m = z_1^{m_1} z_2^{m_2}$, for $m = (m_1, m_2)$. The Floquet transform of the finite-extent operator A is a Laurent polynomial in $z = (z_1, z_2)$ with 2×2 matrix coefficients,

$$\hat{A}(z) = \sum_{m \in \mathbb{Z}^2} A(m)z^{-m}. \quad (2.4)$$

Under the Floquet transform, the operator A becomes a matrix multiplication operator,

$$(Af)^\wedge(z) = \hat{A}(z)\hat{f}(z). \quad (2.5)$$

The *dispersion function* for A is the determinant

$$D(z_1, z_2, E) = \det(\hat{A}(z_1, z_2) - E). \quad (2.6)$$

It is a Laurent polynomial in z_1 and z_2 and a polynomial in E . The *Floquet surface* Φ_E for A at energy E is the algebraic set

$$\Phi_E = \{(z_1, z_2) \in (\mathbb{C}^*)^2 : D(z_1, z_2, E) = 0\}, \quad (2.7)$$

where $\mathbb{C}^* = \mathbb{C} \setminus \{0\}$. A pair (z_1, z_2) lies on Φ_E if and only if A admits a Floquet mode u at (z_1, z_2, E) , which is a function $u : \mathcal{V}(\Gamma) \rightarrow \mathbb{C}$ (not in ℓ^2) such that, for all $v \in \mathcal{V}(\Gamma)$ and $n = (n_1, n_2) \in \mathbb{Z}^2$,

$$Au = Eu \quad \text{and} \quad u(nv) = z_1^{n_1} z_2^{n_2} u(v), \quad (2.8)$$

that is, u is a simultaneous eigenfunction of A and of the \mathbb{Z}^2 action. When considered as a relation between the wavenumbers $(k_1, k_2) \in \mathbb{C}^2$ (for fixed E), where $z_1 = e^{ik_1}$ and $z_2 = e^{ik_2}$, the Floquet surface is called the *Fermi surface*.

The Floquet surface (or the Fermi surface) Φ_E at energy E is *reducible* whenever $D(z_1, z_2, E)$ can be factored nontrivially, that is,

$$D(z_1, z_2, E) = D_1(z_1, z_2, E) D_2(z_1, z_2, E), \quad (2.9)$$

in which D_1 and D_2 are Laurent polynomials in z_1 and z_2 , neither of which is a monomial. That is to say, reducibility of Φ_E is the same as nontrivial factorization of $D(z_1, z_2, E)$.

Figure 1 shows an example of (one period of) a periodic (nonplanar) graph with two vertices and ten edges per period, including two loops. The matrix elements of a self-adjoint periodic operator A with this associated graph are labeled. The matrix $\hat{A}(z) - E$ is

$$\begin{aligned} \hat{A}(z) - E = & \begin{bmatrix} a_0 - E & b_0 \\ \bar{b}_0 & c_0 - E \end{bmatrix} + z_1 \begin{bmatrix} a_1 & 0 \\ d_1 & c_1 \end{bmatrix} + z_2 \begin{bmatrix} a_2 & 0 \\ d_2 & 0 \end{bmatrix} + z_1 z_2 \begin{bmatrix} a_3 & 0 \\ d_3 & 0 \end{bmatrix} + \\ & + z_1^{-1} \begin{bmatrix} \bar{a}_1 & \bar{d}_1 \\ 0 & \bar{c}_1 \end{bmatrix} + z_2^{-1} \begin{bmatrix} \bar{a}_2 & \bar{d}_2 \\ 0 & 0 \end{bmatrix} + z_1^{-1} z_2^{-1} \begin{bmatrix} \bar{a}_3 & \bar{d}_3 \\ 0 & 0 \end{bmatrix}. \end{aligned} \quad (2.10)$$

In section 3, we show that, by collecting all planar connected periodic graphs into isomorphism classes, it is sufficient to consider only the edges shown in Fig. 1.

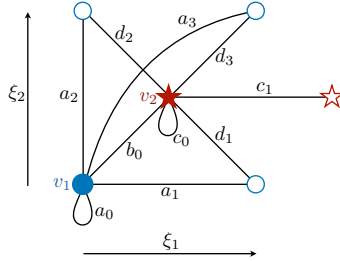


Figure 1: All possible edges in a fundamental domain for the sixteen equivalence classes discussed in Section 3.

2.3 Theorems

The main result of this work is Theorem 1, and Theorem 2 is an application of it to weighted graph Laplacians with positive coefficients. Subsequent sections develop their proofs.

Theorem 1 (Irreducible Fermi surface for planar graphs). *Let A be a doubly periodic self-adjoint discrete graph operator (with complex graph coefficients) whose associated graph Γ is connected and planar with two vertices per fundamental domain. The Fermi surface for A is either (i) irreducible for all but a finite number of energies E or (ii) reducible for all energies E .*

(a) *If the matrix elements of A corresponding to the edges of Γ are real and those corresponding to non-loops are positive, then (i) holds.*

(b) *If (ii) holds, then Γ minus its loops is the graph of the tetrakis square tiling. If a factorization of type $D(z_1, z_2) = D_1(z_1)D_2(z_2)$ is admitted, it must be of the form*

$$D(z_1, z_2) = C(\alpha_1 z_1 + \bar{\alpha}_1 z_1^{-1} + (\beta_1 + \gamma_1 E))(\alpha_2 z_2 + \bar{\alpha}_2 z_2^{-1} + (\beta_2 + \gamma_2 E)), \quad (2.11)$$

in which C is a constant, α_1 and α_2 are complex numbers, and β_i and γ_i are real numbers. If the graph coefficients are real, then (2.11) is the only possible factorization and α_i are real. The set of all operators with real coefficients and reducible Fermi surface is parameterized explicitly by four real variables.

The parameterization of the set of graph operators with real coefficients associated to the tetrakis tiling will be given in section 4.2 (equations (4.75)).

Part (a) of this theorem applies to weighted discrete graph Laplace operators. Such an operator A acts on a function $\tilde{f} \in \ell^2(\mathcal{V}(\Gamma), \mathbb{C})$ by

$$(A\tilde{f})(v) = \sum_{e=(w,v) \in \mathcal{E}(\Gamma)} a_e(\tilde{f}(w) - \tilde{f}(v)). \quad (2.12)$$

By identifying \tilde{f} with $f \in \ell^2(\mathbb{Z}^2, \mathbb{C}^2)$ as described in section 2.1, A can be written in the convolution form (2.1). The matrices $A(m) = \{a_{ij}^m\}_{i,j \in \{1,2\}}$ and the off-diagonal elements of $A(0)$ can be chosen independently through the choice of the weights a_e . Then the diagonal entries of $A(0)$ are determined by the following relations for $i = 1, 2$:

$$\sum_{m \in \mathbb{Z}^2} a_{i1}^m + a_{i2}^m = 0. \quad (2.13)$$

The graph associated with this operator (as described in section 2.1) typically has loops ($a_{ii}^0 \neq 0$) due to the “self-interactions” coming from the $-\tilde{f}(v)$ terms in (2.12). If all the numbers a_e in (2.12) are positive, then only the loops contribute non-positive (but real) coefficients to A . Thus part (a) of Theorem 1 yields a corollary.

Theorem 2. *Let A be a doubly periodic discrete graph Laplace operator with positive weights, whose associated graph Γ is connected and planar with two vertices per fundamental domain. Then the Fermi surface for A is irreducible for all but a finite number of energies E .*

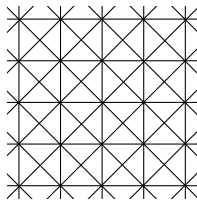


Figure 2: The periodic graph corresponding to the tetrakis square tiling of the plane.

3 Reduction to sixteen equivalence classes

This section describes the reduction of all graphs associated to the operators we are considering to sixteen cases of equivalent graphs. Each of these cases is represented by a distinguished graph drawn in the plane whose fundamental domain contains a subset of the edges drawn in Fig. 1.

Let us call a graph *admissible* if it is periodic, connected, and planar and has two vertices per fundamental domain and no loops. We argue in this section that each admissible graph is isomorphic to one of sixteen graphs. A loop may be added to any of these graphs without losing planarity. The strategy is to show that each admissible graph has one of three special periodic tilings of the plane as a subgraph, where the boundaries of the tiles contain certain vertices and their shifts. These special graphs, denoted by Γ_A , Γ_B , and Γ_C , are drawn in Figs. 3 and 4. New edges are then added to the tiles to produce all admissible graphs, which are depicted in Figs. 5, 6, and 8.

Recall the set decomposition $\mathcal{V}(\Gamma) = \mathcal{V}_1 \dot{\cup} \mathcal{V}_2$ of the vertices of Γ into the two disjoint sets of equivalent vertices (the two orbits of the \mathbb{Z}^2 action). For $i \in \{1, 2\}$, let Γ_i denote the subgraph induced by \mathcal{V}_i , that is, the graph with vertex set \mathcal{V}_i and edge set containing those edges in $\mathcal{E}(\Gamma)$ both of whose vertices are in \mathcal{V}_i .

There are three disjoint cases to consider.

Case A. Γ is bipartite with respect to the decomposition $\mathcal{V}(\Gamma) = \mathcal{V}_1 \dot{\cup} \mathcal{V}_2$, that is, both Γ_1 and Γ_2 have no edges.

Case B. Either Γ_1 or Γ_2 has (or both have) one edge per fundamental domain, and neither has more than one edge per fundamental domain.

Case C. Either Γ_1 or Γ_2 has at least two edges per fundamental domain.

In each of these cases, we determine a minimal subgraph that Γ must contain, which will be Γ_A , Γ_B , or Γ_C . Afterwards, we will find all possible admissible graphs containing that subgraph.

3.1 Minimal subgraphs

For each Case $X \in \{A, B, C\}$, we determine a minimal admissible graph Γ_X . The graph Γ_X satisfies the conditions of Case X and has the property that any graph Γ satisfying the conditions of Case X must contain Γ_X up to interchanging \mathcal{V}_1 and \mathcal{V}_2 and transformation by a matrix in $GL(2, \mathbb{Z})$.

Case A. This construction is illustrated in Fig. 3. Γ must contain an edge e_0 connecting a vertex v_1 in \mathcal{V}_1 to a vertex v_2 in \mathcal{V}_2 . Set $W = \{v_1, v_2\}$. Since Γ is connected, it has an edge e_1 connecting a vertex in W with a vertex in a translate gW of W for some nonzero $g \in \mathbb{Z}^2$. The translates of the vertices v_1 and v_2 and the edges e_0 and e_1 by elements $ng \in \mathbb{Z}^2$, with $n \in \mathbb{Z}$, form an infinite periodic subchain C_1 of Γ , with vertices alternating between \mathcal{V}_1 and \mathcal{V}_2 . The \mathbb{Z}^2 orbits of v_1, v_2, e_0 and e_1 form an array of identical, disjoint copies of C_1 that contains all vertices of Γ . Since Γ is connected, there is an edge e_2 connecting a vertex in C_1 with a vertex in a translate of C_1 (that is not equal to C_1 itself); and since Γ is planar, this connection must be between adjacent chains. The \mathbb{Z}^2 orbits of v_1, v_2, e_0, e_1 and e_2 is a hexagonal lattice, and contains all vertices of Γ .

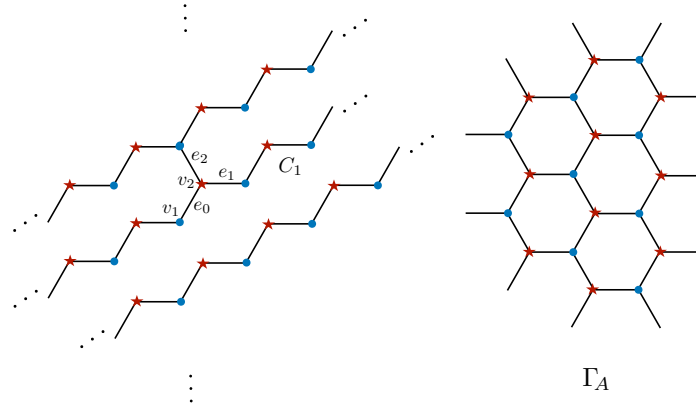


Figure 3: **Case A.** Every graph Γ must have Γ_A as a subgraph. The chain C_1 is described in the text for the construction of Γ_A . Vertices in \mathcal{V}_1 are depicted by blue dots, and vertices in \mathcal{V}_2 are depicted by red stars.

Case B. This construction is illustrated in Fig. 4 (left). Let Γ have an edge e_1 between a vertex $v_1 \in \mathcal{V}_1$ and gv_1 for some nonzero $g \in \mathbb{Z}^2$. The translates of this edge and its vertices by elements $ng \in \mathbb{Z}^2$, with $n \in \mathbb{Z}$, form a chain C_1 connecting a one-dimensional array of vertices in \mathcal{V}_1 . All the \mathbb{Z}^2 translates produce an array of disconnected translates of C_1 , which is just the graph Γ_1 induced by \mathcal{V}_1 since Case B disallows more than one edge per fundamental domain of Γ_1 . Since Γ is connected and planar, and no edge can join a vertex in C_1 with one of its disjoint translates, there is an edge e_2 connecting v_1 to a vertex $v_2 \in \mathcal{V}_2$ and another edge e_3 connecting v_2 to a vertex $hv_1 \in \mathcal{V}_1$ in an adjacent translate of C_1 . The \mathbb{Z}^2 orbit of the vertices v_1 and v_2 and the edges e_1, e_2 , and e_3 , form a connected periodic graph Γ_B , as depicted in Fig. 4. As in Case A, since Γ is planar, Γ_B contains all the vertices of Γ .

Case C. This construction is illustrated in Fig. 4 (right). Let Γ have an edge e_1 connecting two vertices v_1 and gv_1 in \mathcal{V}_1 . The \mathbb{Z}^2 orbits of v_1, gv_1 and e_1 form an array of disjoint chains, which contain all vertices

in \mathcal{V}_1 . Since Γ is connected and planar, and has an edge from v_1 to a vertex in \mathcal{V}_1 that is different from v_1 , gv_1 and $g^{-1}v_1$, v_1 has to be connected to a vertex in \mathcal{V}_1 on an adjacent chain by an edge e_2 . The \mathbb{Z}^2 orbits of v_1 , e_1 and e_2 form a square lattice and contain all vertices in \mathcal{V}_1 . When Γ is drawn in the plane, this grid can be depicted as a square grid with each square containing exactly one vertex from \mathcal{V}_2 in its interior. Since Γ is connected, a vertex of \mathcal{V}_1 (and therefore every vertex of \mathcal{V}_1 by periodicity) must be connected by an edge to a vertex of \mathcal{V}_2 . Let e_0 be an edge connecting v_1 to a vertex $v_2 \in \mathcal{V}_2$, such that e_0 does not cross any edge of the square grid. Define Γ_C to be this grid together with all the vertices in \mathcal{V}_2 and the \mathbb{Z}^2 translates of e_0 . This subgraph is depicted in Fig. 4.

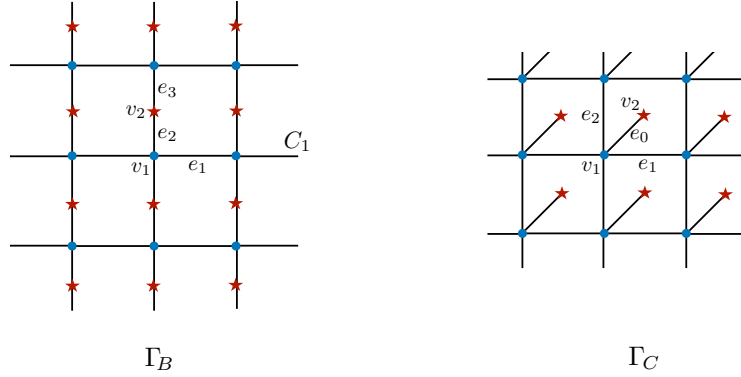


Figure 4: **Case B.** Every graph Γ must have Γ_B as a subgraph. Its construction is described in the text in the construction of Γ_B . **Case C.** Every graph Γ must have Γ_C as a subgraph. Vertices in \mathcal{V}_1 are depicted by blue dots, and vertices in \mathcal{V}_2 are depicted by red stars.

3.2 Equivalent graphs for each case

For each of the three cases, we determine up to graph isomorphism all of the planar periodic graphs, with two vertices per period, that contain the given subgraph.

Case A. The minimal graph Γ_A (Fig. 3) has as a fundamental domain two vertices and three edges—one edge of each of the three orientations on the boundary of a hexagon. Besides Γ_A itself, the only way to produce an admissible graph satisfying Case A and containing Γ_A is to connect two opposite vertices within the hexagon in Fig. 5. All three ways of doing this lead to periodic graphs that are isomorphic by rotation. Thus, the two subcases A1 and A2 in Fig. 5 represent the only equivalence classes for Case A. Each of these subcases can be identified with the graph in Fig. 1, in which only certain edges are retained.

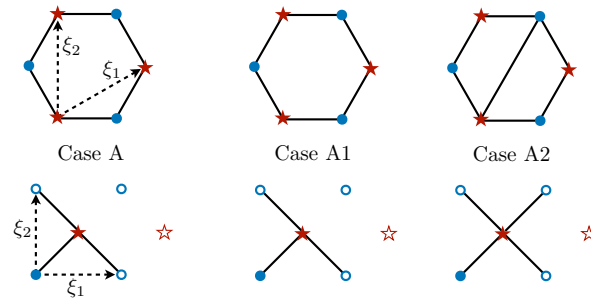


Figure 5: The hexagon is one face of the minimal graph Γ_A shown in Fig. 3 for Case A. In each subcase, the edges in one fundamental domain are identified with edges in the standard representation in Fig. 1.

Case B. The minimal graph Γ_B (Fig. 4) has a fundamental domain consisting of two vertices and three edges on the boundary of a rectangular face, as illustrated in Fig. 6. All other graphs of Case B are obtained

by joining vertices on the boundary of the face by edges passing through the face provided that no edges cross and no two vertices in \mathcal{V}_1 (blue circles) are connected. We argue that there are six equivalence classes of isomorphic graphs, represented by the diagrams in Fig. 6. First consider the graphs for which the vertices in \mathcal{V}_2 (stars) are not connected. There are four ways in which a single edge connects a circle to a star, and all of these are isomorphic by reflections (Case B2). There are four ways to have two edges connecting circle to star (Case B3), and all of them are isomorphic. The one illustrated in Fig. 6 yields only one of the other four upon reflection. To obtain the other two (both interior edges sloped the same way), one has to shear by the matrix $\begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} \in GL(2, \mathbb{Z})$ while keeping the two points in any shift of \mathcal{V} (circle and star) rigid; this is illustrated in Fig. 7. The other three equivalence classes (Cases B4–B6) are obtained by adding an edge connecting the two stars to each of Cases B1–B3.

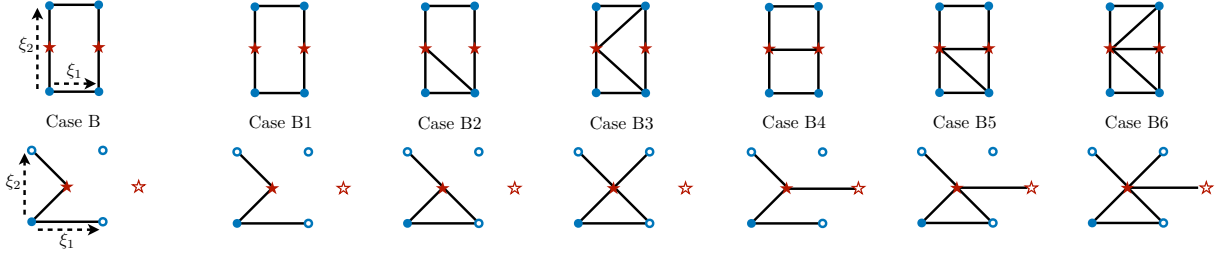


Figure 6: The rectangle is one face of the minimal graph Γ_B shown in Fig. 4 for Case B. In each of the six subcases, the edges in one fundamental domain are identified with edges in the standard representation in Fig. 1.

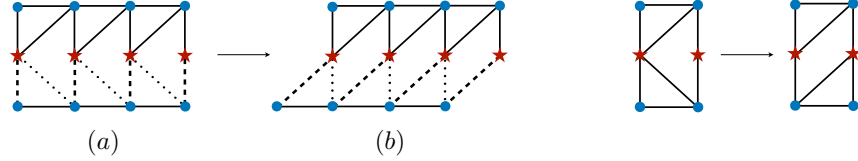


Figure 7: In the description of Case B in section 3.2, it is mentioned how Case B3 in Fig. 6 can be represented by either of the two diagrams on the right through an application of a shear matrix to Γ . This is illustrated by the periodic shift that transforms the drawing (a) to the drawing (b) of the same graph Γ .

Case C. The minimal graph Γ_C (Fig. 4) has a fundamental domain consisting of two edges and one vertex of the square, the vertex at the center, and the edge connecting to it, as depicted in Case C1 of Fig. 8. We argue that each admissible graph containing Γ_C as a subgraph is isomorphic to one of the eight cases shown in Fig. 8. Without adding any diagonal edge, one obtains Cases C1–C4, taking into account equivalence through rotation and reflection, and Case C8. The other cases are obtained by adding a diagonal edge to Cases C1–C4. This results in only three new Cases C5–C7 because in fact Case C4 with an extra diagonal is isomorphic to Case C6. This is seen through applying a shear and then a reflection, similarly to the argument for Case B. This is illustrated in Fig. 9.

4 (Ir)reducibility of the Fermi surface

Let the graph variables not corresponding to loops be denoted by the multi-variable $S = (a_1, a_2, a_3, b_0, c_1, d_1, d_2, d_3)$, and let \bar{S} denote the complex conjugates as independent variables. A graph operator of any of the sixteen types is determined by a choice of values of the variables a_0, c_0, S .

We argue that, without sacrificing the generality of our analysis, we may take $c_0 = 0$. Indeed, with the alternatives (i) and (ii) in Theorem 1 in mind, let S be fixed at some arbitrary values. For any $\alpha \in \mathbb{C}$, factorability of $D(z_1, z_2, E)$ for all energies E with $a_0 = \alpha$ and $c_0 = 0$ is equivalent to factorability for all energies E for any choice of a_0 and c_0 with $a_0 - c_0 = \alpha$; this is evident from the definition of D in (2.6)

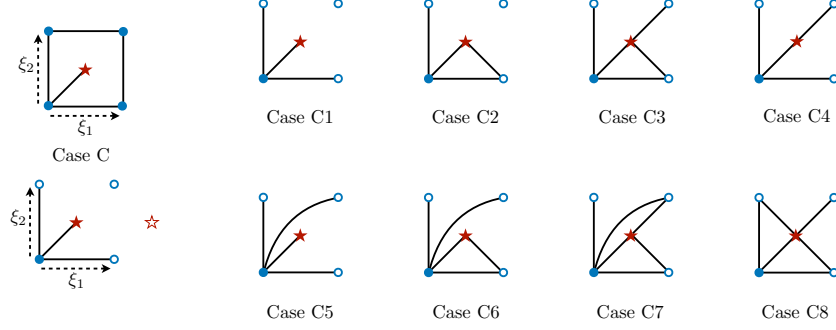


Figure 8: The repeating unit of the minimal graph Γ_C is a face of the subgraph Γ_1 (grid on circle vertices) plus an extra edge connecting a star vertex. The edges of each of the eight subcases are drawn according to the standard representation in Fig. 1.

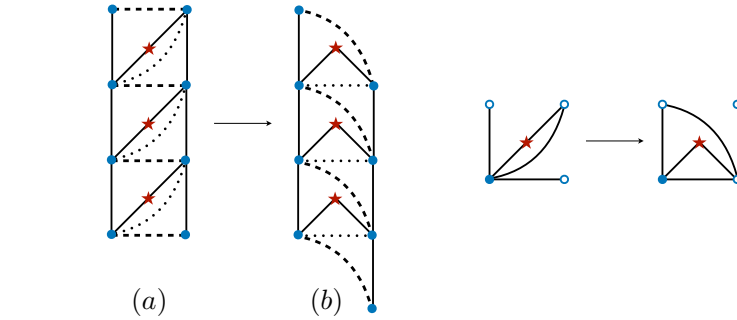


Figure 9: In the description of Case C in section 3.2, it is mentioned how Case C4 with an extra diagonal is isomorphic to Case C6. This is illustrated, up to reflections, by the periodic shift that transforms the drawing (a) to the drawing (b) of the same graph Γ .

and (2.10). Similarly, non-factorability of $D(z_1, z_2, E)$ at all but finitely many energies E with $a_0 = \alpha$ and $c_0 = 0$ is equivalent to non-factorability at all but finitely many energies E for any choice of a_0 and c_0 with $a_0 - c_0 = \alpha$. Therefore, we will henceforth take

$$c_0 = 0. \quad (4.14)$$

4.1 Alternatives (i) and (ii) in Theorem 1

In this subsection, we justify that either (i) or (ii) in Theorem 1 holds.

The dispersion function $D(z_1, z_2, E)$ is a Laurent polynomial in (z_1, z_2) with coefficients that are polynomials in (E, a_0, S, \bar{S}) . There is a nontrivial factorization $D(z_1, z_2, E) = D_1(z_1, z_2, E)D_2(z_1, z_2, E)$ at given values of (E, a_0, S) if and only if a certain finite set P_1 of polynomials (P_1 is $GB_C(I)$ as defined precisely in Step 7 of Section 4.2) in the variables (E, a_0, S, \bar{S}) vanishes at those values, with the values \bar{S} being numerically conjugate to the values S .

Let P_2 be the set of all coefficients of the polynomials in P_1 , viewed as polynomials in E (P_2 is Coeff_E as defined precisely in Step 8 of in Section 4.2). The elements of P_2 are polynomials in (a_0, S, \bar{S}) . First, let (a_0, S, \bar{S}) take on values, with S and \bar{S} being numerical conjugates, at which all elements of P_2 vanish. Then all elements of P_1 vanish at all energies E , which implies that $D(z_1, z_2, E)$ admits a nontrivial factorization. Alternatively, let (a_0, S, \bar{S}) take on values at which some element of P_2 does not vanish. Thus some element of P_1 is a nonzero polynomial in E . For any value of E that is not a root of this polynomial, there is then an element of P_1 that does not vanish at the values (E, a_0, S) . This implies that $D(z_1, z_2, E)$ is not factorable at such (E, a_0, S) . Thus, in the second alternative, $D(z_1, z_2, E)$ is not factorable at each value of E that is not a root of a certain polynomial.

4.2 Algorithm

We present an algorithm for determining reducibility of the Fermi surface, which leads to a proof of Theorem 1. We implement it in later sections, with a combination of Mathematica® and by-hand computations.

Loop I: do the following steps for each of the sixteen graphs shown in Figs. 5, 6, and 8.

1. Let Γ denote one of the sixteen cases shown in Figs. 5, 6, and 8. Let $S_\Gamma \subset S$ be the set of variables corresponding to edges of Γ . That is, the values of the variables in S_Γ are nonzero but are otherwise undetermined and the values of the variables in $S \setminus S_\Gamma$ are zero. Set $K_\Gamma = |S_\Gamma|$ and let \bar{S}_Γ denote the set of the complex conjugates of the variables in S_Γ , viewed as independent from S_Γ . Compute $D(z_1, z_2, E) := \det(\hat{A}(z) - E)$, where $\hat{A}(z)$ is defined in (2.10) with the variables in $S \setminus S_\Gamma$ and their complex conjugates set equal to zero. $D(z_1, z_2, E)$ is a Laurent polynomial in z_1 and z_2 with coefficients that are polynomials in the variables $(E, a_0, S_\Gamma, \bar{S}_\Gamma)$. It possesses the symmetry

$$D(z_1, z_2) = \bar{D}(z_1^{-1}, z_2^{-1}), \quad (4.15)$$

in which \bar{D} indicates conjugation of the coefficients $(S_\Gamma, \bar{S}_\Gamma)$ and not conjugation of the variables (z_1, z_2) .

2. Compute $P_c := z_1^{m_0} z_2^{n_0} D$, where m_0 and n_0 are the highest powers of z_1 and z_2 in the monomials of D , which, by (4.15) are also minus the lowest powers of z_1 and z_2 in the monomials of D . In all cases we compute below, it turns out that, for all choices of $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$, P_c is a polynomial in z_1 and z_2 with neither z_1 nor z_2 as a common factor for all but finitely many values of E . Let M denote the degree of P_c . Factoring D nontrivially (neither factor is a monomial) into Laurent polynomials in z_1 and z_2 is equivalent to factoring P_c into nonconstant polynomials in z_1 and z_2 .

3. The polynomial P_c is of the form $P_c = \sum_{(m,n) \in \mathcal{I}_c} c_{m,n} z_1^m z_2^n$, where $\mathcal{I}_c = \{(m,n) \in U : m+n \leq M\}$ and $U = \mathbb{Z}_{\geq 0} \times \mathbb{Z}_{\geq 0}$. Compute $c_{m,n}$ as a polynomial in the variables $(E, a_0, S_\Gamma, \bar{S}_\Gamma)$. Identify the set C_Z of “always zero” coefficients $c_{m,n}$ and the set C_N of “almost never zero” coefficients $c_{m,n}$. More precisely, $(m,n) \in C_Z$ if and only if $c_{m,n} = 0$ for all real E and a_0 and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$, that is, $c_{m,n}$ is identically zero as a polynomial in these variables; and $(m,n) \in C_N$ if and only if for any real a_0 and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$, $c_{m,n} = 0$ for at most finitely many real E , that is, $c_{m,n}$ is a nonzero polynomial in E for each choice of real a_0 and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$.

4. The polynomial P_c factors if and only if there exist polynomials P_a and P_b of the form $P_a = \sum_{(m,n) \in \mathcal{I}_a} a_{m,n} z_1^m z_2^n$ and $P_b = \sum_{(m,n) \in \mathcal{I}_b} b_{m,n} z_1^m z_2^n$ such that

$$c_{m,n} = \sum_{(i,j)+(k,l)=(m,n), (i,j) \in \mathcal{I}_a, (k,l) \in \mathcal{I}_b} a_{i,j} b_{k,l} \quad \forall (m,n) \in U, \quad (4.16)$$

in which $\mathcal{I}_a = \mathcal{I}_b = \{(m,n) \in U : m+n \leq M-1\}$. To decrease the computational cost, the index sets \mathcal{I}_a and \mathcal{I}_b can be reduced by exploiting the relations (4.16). Consider all $(m,n) \in C_Z$ such that the right-hand side of (4.16) is a monomial in the a and b coefficients; each of these relations gives a product of coefficients equal to zero. For each of these (m,n) , one of the coefficients (say $a_{i,j}$) in the monomial must be set to zero and its index removed from the corresponding index set (say, (i,j) removed from \mathcal{I}_a). There are different ways in which this can be done, and each way yields a reduced pair of index sets $(\mathcal{I}_a, \mathcal{I}_b)$ with a corresponding factorization $P_c = P_a P_b$. Each of these new pairs of index sets leads to new relations obtained by setting the appropriate coefficients in (4.16) to zero. The new relations may again contain monomials in the right-hand side, and one can further reduce the index set in different ways. This process can be continued to obtain a tree of index pairs $(\mathcal{I}_a, \mathcal{I}_b)$ with terminal nodes being those for which the relations (4.16) have no monomial right-hand sides for all $(m,n) \in C_Z$. Some of these index pairs can be ruled out, namely those for which the right-hand side of (4.16) vanishes identically for some $(m,n) \in C_N$. In practice, we terminate any part of the tree whenever this occurs.

Let N_Γ be the number of reduced index pairs. (We find that, for each of the sixteen cases there are zero to six reduced index pairs.) Obviously, for any fixed real a_0 and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$, P_c factors at energy E if and only if P_c factors into P_a and P_b with one of the reduced index pairs at energy E .

Loop II: do Steps 5–8 for each of the N_Γ reduced index pairs $(\mathcal{I}_a, \mathcal{I}_b)$ of index sets obtained in Step 4.

5. Construct the set

$$I := \left\{ c_{m,n} - \sum_{(i,j)+(k,l)=(m,n), (i,j) \in \mathcal{I}_a, (k,l) \in \mathcal{I}_b} a_{i,j} b_{k,l}, \quad \forall (m,n) \in U \right\}. \quad (4.17)$$

I is a finite set. Recall that for given $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$, P_c factors into $P_a = \sum_{(m,n) \in \mathcal{I}_a} a_{m,n} z_1^m z_2^n$ and $P_b = \sum_{(m,n) \in \mathcal{I}_b} b_{m,n} z_1^m z_2^n$ if and only if there exist $\{a_{m,n} : (m,n) \in \mathcal{I}_a\}$ and $\{b_{m,n} : (m,n) \in \mathcal{I}_b\}$ such that every element of I vanishes.

6. Let $\ell = |\mathcal{I}_a| + |\mathcal{I}_b|$, and fix an ordering O of $\{a_{i,j}, b_{k,l}, c_{m,n} : (i,j) \in \mathcal{I}_a, (k,l) \in \mathcal{I}_b, (m,n) \in \mathcal{I}_c\}$, where the variables $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$ are placed at the end. Compute the Groebner basis of the $(\ell+1)$ th elimination ideal of I with the ordering O , and denote it by $GB_C(I)$. By standard nonlinear elimination theory [4], the Groebner basis $GB_C(I)$ is a set of polynomials in $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$, whose zero set is the projection of the zero set of I onto $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$. That is, P_c factors into some polynomials P_a and P_b with reduced index sets \mathcal{I}_a and \mathcal{I}_b if and only if $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$ is a root of all elements of $GB_C(I)$. Note that the $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$ being a root of $GB_C(I)$ is independent of the choice of the ordering O , as long as $\{c_{m,n} : (m,n) \in \mathcal{I}_c\}$ are placed at the end.

7. Compute the set GB of polynomials in $(E, a_0, S_\Gamma, \bar{S}_\Gamma)$, which is obtained from $GB_C(I)$ by replacing each $c_{m,n} : (m,n) \in \mathcal{I}_c$ by its expression as a polynomial in $(E, a_0, S_\Gamma, \bar{S}_\Gamma)$, according to Step 3. Compute Coeff_E , which is defined to be the totality of all coefficients of the elements of GB when treated as polynomials in the variable E . It can be seen that Coeff_E is a set of polynomials in the variables $(a_0, S_\Gamma, \bar{S}_\Gamma)$.

8. Determine whether all elements of Coeff_E vanish for some allowed a_0 and S_Γ , where S_Γ is an element in $(\mathbb{R}^+)^{K_\Gamma}$ or $(\mathbb{R}^*)^{K_\Gamma}$ or $(\mathbb{C}^*)^{K_\Gamma}$, with S_Γ and \bar{S}_Γ related through conjugation. When a choice of values of (a_0, S_Γ) is a root of all elements of Coeff_E , then P_c factors into P_a and P_b of types \mathcal{I}_a and \mathcal{I}_b for all real E . When a choice of values of (a_0, S_Γ) is not a root of all elements of Coeff_E then P_c factors into P_a and P_b of types \mathcal{I}_a and \mathcal{I}_b for at most finitely many values of E . This is because, for such (a_0, S_Γ) , GB contains a nonzero polynomial q in the variable E , whose value will be nonzero for all E except for the finite set of roots of q . (Note that if Coeff_E never identically vanishes for any real a_0 and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$, then it never identically vanishes for any real a_0 and $S_\Gamma \in (\mathbb{R}^*)^{K_\Gamma}$ or $S_\Gamma \in (\mathbb{R}^+)^{K_\Gamma}$.)

Here is a summary of the sets of polynomials involved:

I	contains the polynomials in $(c_{m,n}, a_{i,j}, b_{k,l})$ whose vanishing is equivalent to a factorization $P_c = P_a P_b$.
$GB_C(I)$	is a Groebner basis for I that eliminates $(a_{i,j}, b_{k,l})$ and retains $(c_{m,n})$.
GB	results from substituting $c_{m,n}$ in $GB_C(I)$ by its expression in the variables $(E, a_0, S_\Gamma, \bar{S}_\Gamma)$.
Coeff_E	contains the coefficients of the elements of GB as polynomials in E ; they are polynomials in $(a_0, S_\Gamma, \bar{S}_\Gamma)$.

4.3 Example for the algorithm

We apply the algorithm to Case C5 as an example.

Loop I.

1. Case C5 has edges $S_\Gamma = \{a_1, a_2, a_3, b_0\} \in (\mathbb{C}^*)^4$, and $K_\Gamma = |S_\Gamma| = 4$. Thus

$$\hat{A}(z_1, z_2) - E = \begin{bmatrix} -E + a_0 + a_1 z_1 + \bar{a}_1 z_1^{-1} + a_2 z_2 + \bar{a}_2 z_2^{-1} + a_3 z_1 z_2 + \bar{a}_3 z_1^{-1} z_2^{-1} & b_0 \\ \bar{b}_0 & -E \end{bmatrix}. \quad (4.18)$$

The highest powers of z_1 and z_2 in $D(z_1, z_2, E)$, the determinant of (4.18), are 1 and 1.

2. We obtain $P_c = z_1 z_2 D$,

$$P_c = -a_3 z_1^2 z_2^2 E - a_1 z_1^2 z_2 E - a_2 z_1 z_2^2 E + z_1 z_2 (-|b_0|^2 - a_0 E + E^2) - \bar{a}_2 z_1 E - \bar{a}_1 z_2 E - \bar{a}_3 E. \quad (4.19)$$

3. The degree of P_c is $M = 4$. The coefficients $c_{m,n}$, as polynomials in $(E, a_0, S_\Gamma, \bar{S}_\Gamma)$, are

$$\begin{aligned} c_{2,2} &= -a_3 E, & c_{2,1} &= -a_1 E, & c_{1,2} &= -a_2 E, & c_{1,1} &= -b_0 \bar{b}_0 - a_0 E + E^2, \\ c_{1,0} &= -\bar{a}_2 E, & c_{0,1} &= -\bar{a}_1 E, & c_{0,0} &= -\bar{a}_3 E. \end{aligned} \quad (4.20)$$

The index sets of zero and non-zero coefficients are

$$C_Z = \{(4, 0), (3, 1), (1, 3), (0, 4), (3, 0), (0, 3), (2, 0), (0, 2)\} \cup \{(m, n) \in U : m + n > 4\}, \quad (4.21)$$

$$C_N = U \setminus C_Z. \quad (4.22)$$

4. Using the constraints in C_Z and $\tilde{C}_N = \{(2, 2), (0, 0)\} \subset C_N$, we obtain $N_\Gamma = 3$ pairs of reduced index sets \mathcal{I}_a and \mathcal{I}_b . The corresponding P_a and P_b are of the forms

$$P_a = a_{1,1}z_1z_2 + a_{0,0}, \quad P_b = b_{1,1}z_1z_2 + b_{1,0}z_1 + b_{0,1}z_2 + b_{0,0}; \quad (4.23)$$

$$P_a = a_{1,1}z_1z_2 + a_{0,1}z_2 + a_{0,0}, \quad P_b = b_{1,1}z_1z_2 + b_{1,0}z_1 + b_{0,0}; \quad (4.24)$$

$$P_a = a_{0,1}z_2 + a_{0,0}, \quad P_b = b_{2,1}z_1^2z_2 + b_{1,1}z_1z_2 + b_{1,0}z_1 + b_{0,0}. \quad (4.25)$$

For Loop II, we take Pair (4.23) as an example.

5. The set of coefficients of $P_c - P_aP_b$ in terms of $\{a_{i,j}, b_{k,l}, c_{m,n} : (i, j) \in \mathcal{I}_a, (k, l) \in \mathcal{I}_b, (m, n) \in \mathcal{I}_c\}$ is

$$I = \{c_{0,0} - b_{0,0}, c_{0,1} - b_{0,1}, c_{1,0} - b_{1,0}, c_{1,1} - a_{1,1}b_{0,0}, c_{1,2} - a_{1,1}b_{0,1}, c_{2,1} - a_{1,1}b_{1,0}, c_{2,2} - a_{1,1}b_{1,1}\}. \quad (4.26)$$

Here, $a_{0,0}$ is set equal to 1 without loss of generality.

6. Using the ordering $O = \{a_{1,1}, a_{0,0}, b_{1,1}, b_{1,0}, b_{0,1}, b_{0,0}, c_{2,2}, c_{2,1}, c_{1,2}, c_{1,1}, c_{1,0}, c_{0,1}, c_{0,0}\}$, we compute $GB_C(I)$ in Mathematica®:

$$GB_C(I) = \{-c_{1,0}c_{1,2} + c_{2,1}c_{0,1}, c_{0,0}c_{1,2}^2 - c_{1,1}c_{1,2}c_{0,1} + c_{2,2}c_{0,1}^2, \\ c_{0,0}c_{2,1}c_{1,2} - c_{1,0}c_{1,1}c_{1,2} + c_{1,0}c_{2,2}c_{0,1}, c_{0,0}c_{2,1}^2 + c_{1,0}^2c_{2,2} - c_{1,0}c_{2,1}c_{1,1}\}. \quad (4.27)$$

7. Substitute the relations in (4.20) for $c_{m,n}$ in (4.27) to obtain GB . The set of coefficients of GB as polynomials in E is

$$\text{Coeff}_E = \{a_1\bar{a}_1 + a_2\bar{a}_2, -\bar{a}_1a_2b_0\bar{b}_0, -a_0\bar{a}_1a_2 - \bar{a}_1^2a_3 - a_2^2\bar{a}_3, \bar{a}_1a_2, -a_2\bar{a}_2b_0\bar{b}_0, \\ -a_0a_2\bar{a}_2 - \bar{a}_1\bar{a}_2a_3 + a_1a_2\bar{a}_3, a_2\bar{a}_2, a_1\bar{a}_2b_0\bar{b}_0, a_0a_1\bar{a}_2 - \bar{a}_2^2a_3 - a_1^2\bar{a}_3, -a_1\bar{a}_2\}. \quad (4.28)$$

When seeking common solutions of the polynomials in Coeff_E , we keep in mind that the values of S_Γ and \bar{S}_Γ are complex conjugates of each other.

8. The second term in Coeff_E is $-\bar{a}_1a_2|b_0|^2$, which can not be zero when $S_\Gamma \in (\mathbb{C}^*)^4$. This means that, for all choices of $S_\Gamma \in (\mathbb{C}^*)^4$, one of the elements of GB is a nonzero polynomial in E . So if E is not in the finite set of roots of this polynomial, one of the elements of GB does not vanish at (E, S_Γ) . We conclude that, for all real a_0 and $S_\Gamma \in (\mathbb{C}^*)^4$, Case C5 factors into Pair (4.23) for at most finitely many values of E .

Repeating Steps 5–8 for the pair (4.24) reveals that Coeff_E contains the element $a_1a_2a_3|b_0|^2$. Thus for all a_0 and $S_\Gamma \in (\mathbb{C}^*)^4$, Case C5 factors into the pair (4.24) for at most finitely many values of E . Similarly, for the pair (4.25), Coeff_E contains the element $a_1a_3|b_0|^2$, and thus for all a_0 and $S_\Gamma \in (\mathbb{C}^*)^4$, Case C5 factors into the pair (4.25) for at most finitely many E values.

The conclusion in this example is that, for all a_0 and $S_\Gamma \in (\mathbb{C}^*)^4$, Case C5 factors for at most finitely many values of E .

4.4 Proof of Theorem 1(a): Irreducibility

Now we present the proof of Theorem 1 by showing the key results when the Algorithm is performed for each of the sixteen cases.

Case C7. $S_\Gamma = \{a_1, a_2, a_3, b_0, d_1, d_3\}$. Thus $K_\Gamma = 6$ and

$$P_c = (-b_0d_3 - a_3E)z_1^2z_2^2 + (-b_0d_1 - a_1E)z_1^2z_2 + (-\bar{d}_1d_3 - a_2E)z_1z_2^2 \\ + (-|b_0|^2 - |d_1|^2 - |d_3|^2 - a_0E + E^2)z_1z_2 + (-d_1\bar{d}_3 - \bar{a}_2E)z_1 + (-\bar{b}_0\bar{d}_1 - \bar{a}_1E)z_2 - \bar{b}_0\bar{d}_3 - \bar{a}_3E. \quad (4.29)$$

Case C6. $S_\Gamma = \{a_1, a_2, a_3, b_0, d_1\}$ and P_c is equal to (4.29) with $d_3 = 0$.

Case C4. $S_\Gamma = \{a_1, a_2, b_0, d_3\}$ and P_c is equal to (4.29) with $a_3 = d_1 = 0$.

Case C3. $S_\Gamma = \{a_1, a_2, b_0, d_1, d_3\}$ and P_c is equal to (4.29) with $a_3 = 0$.

Case B2. We use the isomorphic graph shown in Fig. 10. Then $S_\Gamma = \{a_1, b_0, d_1, d_3\}$ and P_c is equal to (4.29) with $a_2 = a_3 = 0$.

Case A1. We use the isomorphic graph shown in Fig. 10. Then $S_\Gamma = \{b_0, d_1, d_3\}$ and P_c is equal to (4.29) with $a_1 = a_2 = a_3 = 0$.

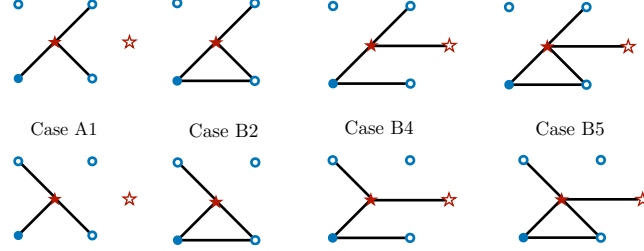


Figure 10: These figures of a fundamental domain depict isomorphisms of the corresponding periodic graph, which are used for convenience in the analysis of the dispersion function $D(z_1, z_2; E)$.

In each of these cases, the coefficients in (4.29) do not identically vanish, and thus they all have a common index set C_Z ,

$$C_Z = \{(4, 0), (3, 1), (1, 3), (0, 4), (3, 0), (0, 3), (2, 0), (0, 2)\} \cup \{(m, n) \in U : m + n > 4\}. \quad (4.30)$$

And since $|b_0 d_3|^2 + |a_3|^2 \neq 0$ for all these cases, they have a common subset \tilde{C}_N of C_N ,

$$\tilde{C}_N = \{(2, 2), (0, 0)\}. \quad (4.31)$$

The constraints from (4.30) and (4.31) allow three pairs of reduced index sets, whose corresponding P_a and P_b are of the forms (4.23), (4.24) and (4.25). Thus, we can compute Coeff_E for Case C7, and obtain Coeff_E for the other cases by setting the appropriate variables equal to zero. Table 1 displays the nonvanishing entries in Coeff_E for each pair in each of these cases.

For Case C4, Coeff_E for pair (4.24) does not directly contain an element that is nonzero for $a_0 \in \mathbb{R}$ and $\{a_1, a_2, b_0, d_3\} \in (\mathbb{C}^*)^4$. We show that two particular elements of Coeff_E cannot be simultaneously zero for $a_0 \in \mathbb{R}$ and $\{a_1, a_2, b_0, d_3\} \in (\mathbb{C}^*)^4$. These elements are

$$\{a_1^2 a_2^2 + a_1 a_2 b_0 d_3, -a_1 a_2 b_0^2 \bar{b}_0 d_3 + a_1 \bar{a}_1 b_0^2 d_3^2 + a_2 \bar{a}_2 b_0^2 d_3^2 - a_1 a_2 b_0 d_3^2 \bar{d}_3\}. \quad (4.32)$$

The vanishing of the first element in (4.32) requires that $a_1 = -b_0 d_3 / a_2$ and $\bar{a}_1 = -\bar{b}_0 \bar{d}_3 / \bar{a}_2$. With these constraints, the second element in (4.32) becomes $b_0^2 d_3^2 (|a_2|^2 + |b_0|^2) (|a_2|^2 + |d_3|^2) / |a_2|^2$, which is nonvanishing for $a_0 \in \mathbb{R}$ and $\{a_1, a_2, b_0, d_3\} \in (\mathbb{C}^*)^4$.

For Case C3, Coeff_E for pair (4.24) does not directly contain an element that is nonzero for $a_0 \in \mathbb{R}$ and $\{a_1, a_2, b_0, d_1, d_3\} \in (\mathbb{C}^*)^5$. We show that three particular elements of Coeff_E cannot be simultaneously zero for $a_0 \in \mathbb{R}$ and $\{a_1, a_2, b_0, d_1, d_3\} \in (\mathbb{C}^*)^5$. These elements are

$$\begin{aligned} & \{\bar{a}_1^2 \bar{a}_2^2 + \bar{a}_1 \bar{a}_2 \bar{b}_0 \bar{d}_3, \\ & 2\bar{a}_1 \bar{a}_2^2 \bar{b}_0 \bar{d}_1 - a_0 \bar{a}_1 \bar{a}_2 \bar{b}_0 \bar{d}_3 + 2\bar{a}_1^2 \bar{a}_2 d_1 \bar{d}_3 + \bar{a}_2 \bar{b}_0^2 \bar{d}_1 \bar{d}_3 + \bar{a}_1 \bar{b}_0 d_1 \bar{d}_3^2, \\ & a_2^2 b_0^2 d_1^2 - a_1 a_2 b_0^2 \bar{b}_0 d_3 - a_0 a_2 b_0^2 d_1 d_3 + 3a_1 a_2 b_0 d_1 \bar{d}_1 d_3 + a_1 \bar{a}_1 b_0^2 d_3^2 + a_2 \bar{a}_2 b_0^2 d_3^2 - a_0 a_1 b_0 \bar{d}_1 d_3^2 \\ & + b_0^2 d_1 \bar{d}_1 d_3^2 + a_1^2 \bar{d}_1^2 d_3^2 - a_1 a_2 b_0 d_3^2 \bar{d}_3\}. \end{aligned} \quad (4.33)$$

The vanishing of the first two elements in (4.33) requires that $b_0 = -a_1 a_2 / d_3$, $a_0 = (a_2^2 d_1 - \bar{d}_1 d_3^2) / (a_2 d_3)$. With these constraints, the third element in (4.33) becomes $a_1^2 a_2^2 (|a_1|^2 + |d_3|^2) (|a_2|^2 + |d_3|^2) / (d_3 \bar{d}_3)$, which is nonvanishing for all $a_0 \in \mathbb{R}$ and $\{a_1, a_2, b_0, d_1, d_3\} \in (\mathbb{C}^*)^5$.

Class	Pair	non-vanishing element in Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$
C7	(4.23)	$b_0 d_1 \bar{d}_1^3 d_3$
	(4.24)	$\bar{a}_1 \bar{a}_2 \bar{a}_3$
	(4.25)	$b_0^3 \bar{b}_0 d_3$
C6	(4.23)	$b_0 b_0 d_1 d_1$
	(4.24)	$\bar{a}_1 \bar{a}_2 \bar{a}_3$
	(4.25)	$-\bar{a}_1 \bar{a}_3$
C4	(4.23)	$a_1 \bar{a}_2$
	(4.24)	explained in text
	(4.25)	$-a_1 \bar{a}_1$
C3	(4.23)	$-a_1 \bar{a}_2$
	(4.24)	explained in text
	(4.25)	$\bar{d}_1^3 \bar{b}_0 d_1 d_3$
B2	(4.23)	$b_0 d_1 \bar{d}_1^3 d_3$
	(4.24)	$\bar{a}_1 \bar{b}_0 d_1 \bar{d}_3^2$
	(4.25)	$b_0^3 \bar{b}_0 d_1 d_3$
A1	(4.23)	$b_0 \bar{d}_1^4 d_3$
	(4.24)	$b_0^2 d_1 \bar{d}_1 d_3^2$
	(4.25)	$b_0^3 \bar{b}_0 d_1 d_3$

Table 1: Nonvanishing entries in Coeff_E for the pairs (4.23), (4.24) and (4.25) for Cases C7, C6, C4, C3, B2, and A1.

In Case C8, $S_\Gamma = \{a_1, a_2, b_0, d_1, d_2, d_3\}$. Thus $K_\Gamma = 6$ and

$$\begin{aligned}
P_c = & (-b_0 d_3) z_1^2 z_2^2 + (-b_0 d_1 - \bar{d}_2 d_3 - a_1 E) z_1^2 z_2 + (-b_0 d_2 - \bar{d}_1 d_3 - a_2 E) z_1 z_2^2 \\
& + (-d_1 \bar{d}_2) z_1^2 + (-|b_0|^2 - |d_1|^2 - |d_2|^2 - |d_3|^2 - a_0 E + E^2) z_1 z_2 + (-\bar{d}_1 d_2) z_2^2 \\
& + (-\bar{b}_0 \bar{d}_2 - d_1 \bar{d}_3 - \bar{a}_2 E) z_1 + (-\bar{b}_0 \bar{d}_1 - d_2 \bar{d}_3 - \bar{a}_1 E) z_2 - \bar{b}_0 \bar{d}_3
\end{aligned} \tag{4.34}$$

For case B3, $S_\Gamma = \{a_1, b_0, d_1, d_2, d_3\}$ and P_c is equal to (4.34) with $a_2 = 0$.

For case A2, $S_\Gamma = \{b_0, d_1, d_2, d_3\}$ and P_c is equal to (4.34) with $a_1 = a_2 = 0$.

In each of these three cases, the coefficients in (4.34) do not identically vanish, and thus they all have a common index set C_Z ,

$$C_Z = \{(4, 0), (3, 1), (1, 3), (0, 4), (3, 0), (0, 3)\} \cup \{(m, n) \in U : m + n > 4\}, \tag{4.35}$$

and since $b_0 d_3 \neq 0$ in these three cases, they have a common subset \tilde{C}_N of C_N ,

$$\tilde{C}_N = \{(2, 2), (0, 0)\}. \tag{4.36}$$

The constraints from (4.36) allow three pairs of reduced index sets, corresponding to factorizations $P_a P_b$ of the forms

$$(a_{1,1} z_1 z_2 + a_{1,0} z_1 + a_{0,1} z_2 + a_{0,0}) (b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2 + b_{0,0}), \tag{4.37}$$

$$(a_{0,2} z_2^2 + a_{0,1} z_2 + a_{0,0}) (b_{2,0} z_1^2 + b_{1,0} z_1 + b_{0,0}), \tag{4.38}$$

$$(a_{0,1} z_2 + a_{0,0}) (b_{2,1} z_1^2 z_2 + b_{2,0} z_1^2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2 + b_{0,0}). \tag{4.39}$$

Therefore, we can compute Coeff_E for Case C8 and obtain Coeff_E for the other two cases by setting the corresponding variables equal to zero. Table 2 displays the nonvanishing entries in Coeff_E for each pair in each case. The determinant D in Case C8 may factor for all E for some $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$; in fact it also may factor for all E for some $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{R}^*)^{K_\Gamma}$, but it only factors for at most finitely many E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{R}^+)^{K_\Gamma}$. This is discussed in more detail in Sec. 4.5. The items listed for Case C8 in Table 2 are items that do not vanish for any $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{R}^+)^{K_\Gamma}$.

Class	Pair	nonvanishing element in Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{R}^+)^{K_\Gamma}$
C8	(4.37)	$-a_1^2 a_2^3 - a_1 a_2^2 d_1 d_2 - a_1 a_2^2 b_0 d_3 - a_2 b_0 d_1 d_2 d_3$
	(4.38)	$a_1 a_2 + d_1 d_2$
	(4.39)	$a_1^2 a_2^2 + a_1 a_2 d_1 d_2 + a_1 a_2 b_0 d_3 + b_0 d_1 d_2 d_3$
Class	Pair	nonvanishing element in Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^{K_\Gamma}$
B3	(4.37)	$\bar{a}_1^2 b_0 d_1 d_2 d_3$
	(4.38)	$-\bar{b}_0 \bar{d}_3$
	(4.39)	$\bar{b}_0 \bar{d}_1 d_2 \bar{d}_3$
A2	(4.37)	$b_0 d_1 d_2 d_3$
	(4.38)	$-\bar{b}_0 \bar{d}_3$
	(4.39)	$\bar{b}_0 \bar{d}_1 d_2 \bar{d}_3$

Table 2: Nonvanishing entries in Coeff_E for Pairs (4.37), (4.38) and (4.39) for Cases C8, B3, and A2.

Case B6: $S_\Gamma = \{a_1, b_0, d_1, d_2, d_3, c_1\}$. Thus $K = 6$ and

$$\begin{aligned}
P_c = & a_1 c_1 z_1^4 z_2 - b_0 d_3 z_1^3 z_2^2 + (a_0 c_1 - b_0 d_1 - \bar{d}_2 d_3 - (a_1 + c_1)E) z_1^3 z_2 + (-b_0 d_2 - \bar{d}_1 d_3) z_1^2 z_2^2 \\
& - d_1 \bar{d}_2 z_1^3 + (-|b_0|^2 + \bar{a}_1 c_1 + a_1 \bar{c}_1 - |d_1|^2 - |d_2|^2 - |d_3|^2 - a_0 E + E^2) z_1^2 z_2 - \bar{d}_1 d_2 z_1 z_2^2 \\
& + (a_0 \bar{c}_1 - \bar{b}_0 \bar{d}_2 - d_1 \bar{d}_3) z_1^2 + (-\bar{b}_0 \bar{d}_1 - d_2 \bar{d}_3 - (\bar{a}_1 + \bar{c}_1)E) z_1 z_2 \\
& - \bar{b}_0 \bar{d}_3 z_1 + \bar{a}_1 \bar{c}_1 z_2.
\end{aligned} \tag{4.40}$$

Since $a_1 c_1 \neq 0$ and $b_0 d_3 \neq 0$, we obtain C_Z and a subset \tilde{C}_N of C_N :

$$\begin{aligned}
C_Z = & \{(5, 0), (2, 3), (1, 4), (0, 5), (4, 0), (1, 3), (0, 4), (0, 3), (0, 2), (0, 0)\} \cup \{(m, n) \in U : m + n > 5\} \\
\tilde{C}_N = & \{(4, 1), (3, 2), (1, 0), (0, 1)\}.
\end{aligned} \tag{4.41}$$

The constraints from (4.41) yield six pairs of reduced index sets, whose corresponding factorizations $P_a P_b$ are of the forms

$$(a_{1,1} z_1 z_2 + a_{0,0}) (b_{3,0} z_1^3 + b_{2,1} z_1^2 z_2 + b_{2,0} z_1^2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2), \tag{4.42}$$

$$(a_{2,0} z_1^2 + a_{1,0} z_1 + a_{0,0}) (b_{2,1} z_1^2 z_2 + b_{1,2} z_1 z_2^2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2), \tag{4.43}$$

$$(a_{2,0} z_1^2 + a_{1,1} z_1 z_2 + a_{1,0} z_1 + a_{0,0}) (b_{2,1} z_1^2 z_2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2), \tag{4.44}$$

$$(a_{2,0} z_1^2 + a_{1,1} z_1 z_2 + a_{1,0} z_1 + a_{0,1} z_2) (b_{2,1} z_1^2 z_2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,0}), \tag{4.45}$$

$$(a_{1,0} z_1 + a_{0,0}) (b_{3,1} z_1^3 z_2 + b_{2,2} z_1^2 z_2^2 + b_{2,1} z_1^2 z_2 + b_{1,2} z_1 z_2^2 + b_{2,0} z_1^2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2), \tag{4.46}$$

$$(a_{1,0} z_1 + a_{0,1} z_2) (b_{3,1} z_1^3 z_2 + b_{2,1} z_1^2 z_2 + b_{2,0} z_1^2 + b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,0}). \tag{4.47}$$

Nonvanishing elements of Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^6$ are, for these six cases respectively, $a_1 c_1 d_1 \bar{d}_2$, $\bar{b}_0^2 d_1 \bar{d}_2 \bar{d}_3^2$, $a_1 b_0^2 c_1 \bar{d}_1 d_2 d_3^2$, $\bar{b}_0^2 d_1^2 \bar{d}_2^2 \bar{d}_3^2$, $b_0^2 d_1^2 d_2^2 d_3^2$, and $\bar{a}_1 \bar{b}_0 \bar{c}_1 \bar{d}_3$.

For Cases B4 and B5, using the isomorphic graphs shown in Fig. 10, we obtain $S_\Gamma = \{a_1, b_0, d_3, c_1\}$ and $S_\Gamma = \{a_1, b_0, d_1, d_3, c_1\}$, respectively. By setting $d_1 = d_2 = 0$ or $d_1 = 0$ in (4.40), we find that these two cases possess a common C_Z and a common subset \tilde{C}_N of C_N ,

$$\begin{aligned}
C_Z = & \{(5, 0), (2, 3), (1, 4), (0, 5), (4, 0), (1, 3), (0, 4), (3, 0), (1, 2), (0, 3), (0, 2), (0, 0)\} \cup \{(m, n) \in U : m + n > 5\} \\
\tilde{C}_N = & \{(4, 1), (3, 2), (1, 0), (0, 1)\}.
\end{aligned} \tag{4.48}$$

The constraints in (4.48) do not allow any reduced index pairs, and so these cases are done.

For Case C1, $S_\Gamma = \{a_1, a_2, b_0\}$ and

$$P_c = -\bar{a}_2 z_1 E - \bar{a}_1 z_2 E - a_1 z_1^2 z_2 E - a_2 z_1 z_2^2 E + z_1 z_2 (-b_0 \bar{b}_0 - a_0 E + E^2). \tag{4.49}$$

For Case C2, $S_\Gamma = \{a_1, a_2, b_0, d_1\}$ and

$$P_c = -\bar{a}_2 z_1 E - a_2 z_1 z_2^2 E + z_1^2 z_2 (-b_0 d_1 - a_1 E) + z_2 (-\bar{b}_0 \bar{d}_1 - \bar{a}_1 E) + z_1 z_2 (-b_0 \bar{b}_0 - d_1 \bar{d}_1 - a_0 E + E^2). \quad (4.50)$$

For Case B1, $S_\Gamma = \{a_1, b_0, d_2\}$ and

$$P_c = -\bar{b}_0 \bar{d}_2 z_1 - b_0 d_2 z_1 z_2^2 - \bar{a}_1 z_2 E - a_1 z_1^2 z_2 E + z_1 z_2 (-b_0 \bar{b}_0 - d_2 \bar{d}_2 - a_0 E + E^2). \quad (4.51)$$

In each of these cases P_c is of degree 3, and

$$\begin{aligned} C_Z &= \{(3, 0), (0, 3), (2, 0), (0, 2), (0, 0)\} \cup \{(m, n) \in U : m + n > 3\}, \\ \tilde{C}_N &= \{(2, 1), (1, 2)\} \subset C_N. \end{aligned} \quad (4.52)$$

The constraints in (4.52) do not allow any reduced index pairs, so we have taken care of these three cases.

4.5 Proof of Theorem 1(b): Reducibility for the tetrakis graph

We examine the factorizations in Case C8, which is the tetrakis graph. We saw in the previous section that this is the only case in which the Fermi surface can be reducible. Recall that reducibility can result from three possible factorizations (4.37, 4.38, 4.39). To determine the graph coefficients that realize these factorizations, we seek common roots $(a_0, a_1, a_2, b_0, d_1, d_2, d_3)$ of all the elements of Coeff_E for each of these three pairs of factors. For pairs (4.38) and (4.39), we find all roots with a_0 real and $(a_1, a_2, b_0, d_1, d_2, d_3)$ complex, and for pair (4.37), we find all real roots $(a_0, a_1, a_2, b_0, d_1, d_2, d_3)$. Each of these 7-tuples of coefficients $(a_0, a_1, a_2, b_0, d_1, d_2, d_3)$ corresponds to a discrete graph whose Fermi surface is reducible. The difficulty in finding all complex solutions to pair (4.37) is discussed at the end of this section.

4.5.1 Complex roots of Coeff_E for factorization (4.38)

A subset of Coeff_E for the factorization (4.38) in Case C8 is

$$\begin{aligned} \mathcal{J} = \{ & -\bar{a}_1 \bar{a}_2 - \bar{b}_0 \bar{d}_3, \bar{a}_1 a_2 + \bar{d}_1 d_2, -\bar{a}_2 \bar{b}_0 \bar{d}_1 - \bar{a}_1 \bar{b}_0 \bar{d}_2 + a_0 \bar{b}_0 \bar{d}_3 - \bar{a}_1 d_1 \bar{d}_3 - \bar{a}_2 d_2 \bar{d}_3, \\ & -\bar{b}_0^2 \bar{d}_1 \bar{d}_2 + b_0 \bar{b}_0^2 \bar{d}_3 - d_1 d_2 \bar{d}_3^2 + \bar{b}_0 d_3 \bar{d}_3^2, -b_0^2 d_1 d_2 + b_0^2 \bar{b}_0 d_3 - \bar{d}_1 \bar{d}_2 d_3^2 + b_0 d_3^2 \bar{d}_3 \}. \end{aligned} \quad (4.53)$$

A common root of Coeff_E has to be a common root of \mathcal{J} . We will find all roots of (4.53) and check that they are in fact roots of Coeff_E . From the first two entries in \mathcal{J} , $a_1 = -b_0 d_3 / a_2$ and $d_1 = \bar{a}_2 b_0 d_3 / (a_2 \bar{d}_2)$. Note that the condition for a_1 is not achievable if all non-loop edge weights are positive. By substituting these relations and their complex conjugates into (4.53), the five elements of \mathcal{J} become (in the same order)

$$\begin{aligned} \mathcal{J} = \{ & 0, 0, \bar{d}_3 (a_0 \bar{b}_0 - a_2 \bar{b}_0^2 / d_2 - \bar{a}_2 d_2 + \bar{b}_0^2 \bar{d}_2 / \bar{a}_2 + b_0 \bar{b}_0 d_3 \bar{d}_3 / (a_2 \bar{d}_2)), \\ & -(-\bar{a}_2 b_0 d_2 + a_2 \bar{b}_0 \bar{d}_2) \bar{d}_3 (a_2 \bar{b}_0^2 \bar{d}_2 - \bar{a}_2 d_2 d_3 \bar{d}_3) / (a_2 \bar{a}_2 d_2 \bar{d}_2), \\ & -(-\bar{a}_2 b_0 d_2 + a_2 \bar{b}_0 \bar{d}_2) d_3 (-\bar{a}_2 b_0^2 d_2 + a_2 \bar{d}_2 d_3 \bar{d}_3) / (a_2 \bar{a}_2 d_2 \bar{d}_2) \}. \end{aligned} \quad (4.54)$$

Vanishing of the third entry in (4.54) requires that $a_0 = a_2 \bar{b}_0 / d_2 + \bar{a}_2 d_2 / \bar{b}_0 - \bar{b}_0 \bar{d}_2 / \bar{a}_2 - b_0 d_3 \bar{d}_3 / (a_2 \bar{d}_2)$ and that this expression be real (since a_0 must be real). By making this substitution for a_0 in (4.54), \mathcal{J} is updated to

$$\begin{aligned} \mathcal{J} = \{ & 0, 0, 0, -(-\bar{a}_2 b_0 d_2 + a_2 \bar{b}_0 \bar{d}_2) \bar{d}_3 (a_2 \bar{b}_0^2 \bar{d}_2 - \bar{a}_2 d_2 d_3 \bar{d}_3) / (a_2 \bar{a}_2 d_2 \bar{d}_2), \\ & -(-\bar{a}_2 b_0 d_2 + a_2 \bar{b}_0 \bar{d}_2) d_3 (-\bar{a}_2 b_0^2 d_2 + a_2 \bar{d}_2 d_3 \bar{d}_3) / (a_2 \bar{a}_2 d_2 \bar{d}_2) \}. \end{aligned} \quad (4.55)$$

For (4.55) to vanish, either $\bar{a}_2 = a_2 \bar{b}_0 \bar{d}_2 / (b_0 d_2)$ or $\bar{a}_2 = a_2 \bar{d}_2 d_3 \bar{d}_3 / (b_0^2 d_2)$. In the former case, (4.55) identically vanishes and $a_0 = a_2 \bar{b}_0 / d_2 + \bar{a}_2 d_2 / \bar{b}_0 - \bar{b}_0 \bar{d}_2 / \bar{a}_2 - b_0 d_3 \bar{d}_3 / (a_2 \bar{d}_2) \in \mathbb{R}$. Note that $\bar{a}_2 = a_2 \bar{b}_0 \bar{d}_2 / (b_0 d_2)$ is a constraint on a_2 and \bar{a}_2 but does not eliminate a_2 and \bar{a}_2 . We obtain that the substitutions

$$\begin{aligned} a_0 &= a_2 \bar{b}_0 / d_2 + \bar{a}_2 d_2 / \bar{b}_0 - \bar{b}_0 \bar{d}_2 / \bar{a}_2 - b_0 d_3 \bar{d}_3 / (a_2 \bar{d}_2), \\ a_1 &= -b_0 d_3 / a_2, \\ \bar{a}_2 &= a_2 \bar{b}_0 \bar{d}_2 / (b_0 d_2), \\ d_1 &= \bar{a}_2 b_0 d_3 / (a_2 \bar{d}_2) \end{aligned} \quad (4.56)$$

cause (4.53) to vanish. In the latter case, by substituting $\bar{a}_2 = a_2 \bar{d}_2 d_3 \bar{d}_3 / (b_0^2 d_2)$ into (4.55), \mathcal{J} is updated into

$$\mathcal{J} = \{0, 0, 0, 0, -(b_0 \bar{b}_0 - d_3 \bar{d}_3)^2 (b_0 \bar{b}_0 + d_3 \bar{d}_3) / (\bar{b}_0 \bar{d}_3)\}. \quad (4.57)$$

Thus $b_0 \bar{b}_0 - d_3 \bar{d}_3 = 0$. In fact, the relations $\bar{a}_2 = a_2 \bar{d}_2 d_3 \bar{d}_3 / (b_0^2 d_2)$ and $b_0 \bar{b}_0 - d_3 \bar{d}_3 = 0$ combine to give $\bar{a}_2 = a_2 \bar{b}_0 \bar{d}_2 / (b_0 d_2)$, which turns the latter case ($\bar{a}_2 = a_2 \bar{d}_2 d_3 \bar{d}_3 / (b_0^2 d_2)$) into a subcase of the former case ($\bar{a}_2 = a_2 \bar{b}_0 \bar{d}_2 / (b_0 d_2)$). We check in Mathematica® that Coeff_E vanishes identically for (4.56). Thus all complex solutions to Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^6$ are represented by (4.56).

We conclude that P_c factors into the form (4.38) if and only if the relations (4.56) are satisfied. These relations provide a parameterization of all factorizations of P_c by the four variables $\{a_2, b_0, d_2, d_3\}$. The factorizations are

$$|d_2|^2 P_c = -(\bar{d}_2 d_3 z_1^2 + d_2 \bar{d}_3 + (|d_3|^2 + |d_2|^2 + a_2 \bar{d}_2 E / b_0) z_1) (b_0 d_2 z_2^2 + \bar{b}_0 \bar{d}_2 + (|d_2|^2 + |b_0|^2 + b_0 d_2 E / a_2) z_2). \quad (4.58)$$

4.5.2 Complex roots of Coeff_E for factorization (4.39)

A subset of Coeff_E is

$$\begin{aligned} \mathcal{K} = & \{(\bar{a}_1 a_2 + \bar{d}_1 d_2)(\bar{a}_1 \bar{a}_2 + \bar{b}_0 \bar{d}_3), (-d_1 d_2 + \bar{b}_0 \bar{d}_3)(\bar{d}_1 \bar{d}_2 - b_0 d_3)(\bar{b}_0 \bar{d}_1 - d_2 \bar{d}_3)^2, \\ & -a_1 \bar{a}_1 d_1 \bar{d}_1 d_2 \bar{d}_2 + \bar{a}_1^2 b_0 d_1 \bar{d}_2 d_3 + a_1^2 \bar{b}_0 \bar{d}_1 d_2 \bar{d}_3 - a_1 \bar{a}_1 b_0 \bar{b}_0 d_3 \bar{d}_3, \\ & \bar{a}_2 b_0^3 d_1^2 d_2 - a_2 b_0 d_1^3 \bar{d}_1 \bar{d}_2 + a_1 b_0^2 \bar{b}_0 d_1 d_2 \bar{d}_2 - a_0 b_0^2 d_1^2 d_2 \bar{d}_2 - a_1 b_0 d_1^2 \bar{d}_1 d_2 \bar{d}_2 + a_2 b_0 d_1^2 d_2 \bar{d}_2^2 - a_1 b_0 d_1 d_2^2 \bar{d}_2^2 \\ & - \bar{a}_2 b_0^3 \bar{b}_0 d_1 d_3 - a_1 b_0^2 \bar{b}_0^2 \bar{d}_2 d_3 + a_0 b_0^2 \bar{b}_0 d_1 \bar{d}_2 d_3 + a_0 b_0 d_1^2 \bar{d}_1 \bar{d}_2 d_3 - a_1 d_1^2 \bar{d}_1^2 \bar{d}_2 d_3 - \bar{a}_2 b_0^2 d_1 d_2 \bar{d}_2 d_3 \\ & - a_2 b_0 \bar{b}_0 d_1 \bar{d}_2^2 d_3 + a_2 d_1^2 \bar{d}_1 \bar{d}_2^2 d_3 + a_1 b_0 \bar{b}_0 d_2 \bar{d}_2^2 d_3 + a_0 b_0 d_1 d_2 \bar{d}_2^2 d_3 - a_1 d_1 \bar{d}_1 d_2 \bar{d}_2^2 d_3 - a_2 d_1 d_2 \bar{d}_2^3 d_3 \\ & + \bar{a}_2 b_0^2 \bar{b}_0 \bar{d}_2 d_3^2 - \bar{a}_2 b_0 d_1 \bar{d}_1 \bar{d}_2 d_3^2 - a_0 b_0 \bar{b}_0 \bar{d}_2^2 d_3^2 + 2a_1 \bar{b}_0 \bar{d}_1 \bar{d}_2^2 d_3^2 - a_0 d_1 \bar{d}_1 \bar{d}_2^2 d_3^2 + a_2 \bar{b}_0 \bar{d}_2^2 d_3^2 + \bar{a}_2 \bar{d}_1 \bar{d}_2^2 d_3^3 \\ & + a_2 b_0^2 d_1^2 \bar{d}_3 + 2a_1 b_0^2 d_1^2 d_2 \bar{d}_3 - a_1 b_0^2 \bar{b}_0 d_1 d_3 \bar{d}_3 - a_0 b_0^2 d_1^2 d_3 \bar{d}_3 + a_1 b_0 d_1^2 \bar{d}_1 d_3 \bar{d}_3 - a_2 b_0 d_1^2 \bar{d}_2 d_3 \bar{d}_3 \\ & + \bar{a}_2 b_0^2 d_1 d_3^2 \bar{d}_3 - a_1 b_0 \bar{b}_0 \bar{d}_2 d_3^2 \bar{d}_3 + a_0 b_0 d_1 \bar{d}_2 d_3^2 \bar{d}_3 + a_1 d_1 \bar{d}_1 \bar{d}_2 d_3^2 \bar{d}_3 - \bar{a}_2 b_0 \bar{d}_2 d_3^2 \bar{d}_3 - a_1 b_0 d_1 d_3^2 \bar{d}_3^2, \\ & a_1 \bar{a}_1 a_2 \bar{a}_2 b_0 \bar{b}_0 + a_0 \bar{a}_1 a_2 \bar{a}_2 b_0 d_1 + a_0 a_1 a_2 \bar{a}_2 \bar{b}_0 \bar{d}_1 + a_1 \bar{a}_1 a_2 \bar{a}_2 d_1 \bar{d}_1 - a_2 \bar{a}_2 b_0 \bar{b}_0 d_1 \bar{d}_1 + a_0 a_1 \bar{a}_1 \bar{a}_2 b_0 d_2 \\ & - \bar{a}_1 \bar{a}_2 b_0^2 d_1 d_2 - a_0^2 a_1 \bar{a}_2 \bar{d}_1 d_2 + a_1 a_2 \bar{a}_2^2 \bar{d}_1 d_2 + a_1 \bar{a}_2 b_0 \bar{b}_0 \bar{d}_1 d_2 + 2a_0 \bar{a}_2 b_0 d_1 \bar{d}_1 d_2 + 2a_1 \bar{a}_2 d_1 \bar{d}_1^2 d_2 \\ & + a_0 a_1 \bar{a}_1 a_2 \bar{b}_0 \bar{d}_2 - a_0^2 \bar{a}_1 a_2 d_1 \bar{d}_2 + \bar{a}_1 a_2^2 \bar{a}_2 d_1 \bar{d}_2 + \bar{a}_1 a_2 b_0 \bar{b}_0 d_1 \bar{d}_2 - a_1 a_2 \bar{b}_0^2 \bar{d}_1 \bar{d}_2 + 2a_0 a_2 \bar{b}_0 d_1 \bar{d}_1 \bar{d}_2 \\ & + 2\bar{a}_1 a_2 d_1^2 \bar{d}_1 \bar{d}_2 + a_1 \bar{a}_1 a_2 \bar{a}_2 d_2 \bar{d}_2 - a_1 \bar{a}_1 b_0 \bar{b}_0 d_2 \bar{d}_2 + 2a_0 \bar{a}_1 b_0 d_1 d_2 \bar{d}_2 + 2a_0 a_1 \bar{b}_0 \bar{d}_1 d_2 \bar{d}_2 - 3a_0^2 d_1 \bar{d}_1 d_2 \bar{d}_2 \\ & + 2a_2 \bar{a}_2 d_1 \bar{d}_1 d_2 \bar{d}_2 + b_0 \bar{b}_0 d_1 \bar{d}_1 d_2 \bar{d}_2 + 3d_1^2 \bar{d}_1^2 d_2 \bar{d}_2 + 2a_1 \bar{a}_2 \bar{d}_1 d_2^2 \bar{d}_2 + 2\bar{a}_1 a_2 d_1 d_2 \bar{d}_2^2 + 3d_1 \bar{d}_1 d_2^2 \bar{d}_2^2 \\ & - \bar{a}_1 a_2 \bar{a}_2^2 b_0 d_3 + a_0 a_1 \bar{a}_1 \bar{a}_2 \bar{d}_1 d_3 - \bar{a}_1 \bar{a}_2 b_0 d_1 \bar{d}_1 d_3 - a_1 \bar{a}_2 \bar{b}_0 \bar{d}_1^2 d_3 - \bar{a}_2^2 b_0 \bar{d}_1 d_2 d_3 + a_0 \bar{a}_1 a_2 \bar{a}_2 \bar{d}_2 d_3 \\ & - a_1 \bar{a}_1 \bar{b}_0 \bar{d}_1 \bar{d}_2 d_3 - a_2 \bar{a}_2 \bar{b}_0 \bar{d}_1 \bar{d}_2 d_3 - 2a_0 \bar{a}_1 d_1 \bar{d}_1 \bar{d}_2 d_3 - \bar{b}_0 d_1 \bar{d}_1^2 \bar{d}_2 d_3 - \bar{a}_1 \bar{a}_2 b_0 d_2 \bar{d}_2 d_3 + 2a_0 \bar{a}_2 \bar{d}_1 d_2 \bar{d}_2 d_3 \\ & - \bar{a}_1 a_2 \bar{b}_0 \bar{d}_2^2 d_3 - \bar{b}_0 \bar{d}_1 d_2 \bar{d}_2^2 d_3 - \bar{a}_1 \bar{a}_2 \bar{d}_1 \bar{d}_2 d_3^2 - a_1 a_2^2 \bar{a}_2 \bar{b}_0 \bar{d}_3 + a_0 a_1 \bar{a}_1 a_2 d_1 \bar{d}_3 - \bar{a}_1 a_2 b_0 d_1^2 \bar{d}_3 \\ & - a_1 a_2 \bar{b}_0 d_1 \bar{d}_1 \bar{d}_3 + a_0 a_1 a_2 \bar{a}_2 d_2 \bar{d}_3 - a_1 \bar{a}_1 b_0 d_1 d_2 \bar{d}_3 - a_2 \bar{a}_2 b_0 d_1 d_2 \bar{d}_3 + 2a_0 a_1 d_1 \bar{d}_1 d_2 \bar{d}_3 - b_0 d_1^2 \bar{d}_1 d_2 \bar{d}_3 \\ & - a_1 \bar{a}_2 b_0 d_2^2 \bar{d}_3 - a_2^2 \bar{b}_0 \bar{d}_1 \bar{d}_2 \bar{d}_3 - a_1 a_2 \bar{b}_0 d_2 \bar{d}_2 \bar{d}_3 + 2a_0 a_2 d_1 d_2 \bar{d}_2 \bar{d}_3 - b_0 d_1 d_2^2 \bar{d}_2 \bar{d}_3 + a_1 \bar{a}_1 a_2 \bar{a}_2 d_3 \bar{d}_3 \\ & - a_1 \bar{a}_1 d_1 \bar{d}_1 d_3 \bar{d}_3 + a_1 \bar{a}_2 \bar{d}_1 d_2 d_3 \bar{d}_3 + \bar{a}_1 a_2 d_1 \bar{d}_2 d_3 \bar{d}_3 - a_2 \bar{a}_2 d_2 \bar{d}_2 d_3 \bar{d}_3 + d_1 \bar{d}_1 d_2 \bar{d}_2 d_3 \bar{d}_3 - a_1 a_2 d_1 d_2 \bar{d}_3^2 \}. \end{aligned} \quad (4.59)$$

A root of Coeff_E has to be a root of \mathcal{K} . From the first entry in \mathcal{K} , $a_1 = -d_1 \bar{d}_2 / \bar{a}_2$ or $a_1 = -b_0 d_3 / a_2$. From the second entry in \mathcal{K} , $d_1 = \bar{b}_0 d_3 / d_2$ or $d_1 = \bar{d}_2 d_3 / b_0$. Next, following the approach in Sec. 4.5.1, it turns out that all roots of \mathcal{K} take four forms; one of them is equivalent to (4.56), and the others are subsets of the set defined by (4.56). A factorization of the form (4.38) implies a factorization of the form (4.39). Thus all roots of Coeff_E for pair (4.39), and therefore all factorizations of the form (4.39), are given exactly by the parameterization (4.56). This implies that each factorization of the form (4.39) is in fact a factorization of the form (4.38).

4.5.3 Real roots of Coeff_E for factorization (4.37)

Observe that, if P_c factors into the form (4.38) for some graph operator, then it also factors into the form (4.37). We now seek all graphs for which P_c factors into the form (4.37). As stated in Theorem 1, we are able to parameterize all such graphs with real coefficients.

Let us briefly summarize the methods of the previous two subsections. For pairs (4.38) and (4.39) of Case C8, the method for obtaining all roots of Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^6$ goes as follows. Coeff_E is generically a set of polynomials in variables $\{x_i : 1 \leq i \leq N\}$. Suppose Coeff_E contains an element that factors into M factors and, for each $k : 1 \leq k \leq M$, when the k th factor is zero, one variable, say x_{i_k} , can be solved in terms of other variables (complex conjugates are treated as independent variables). In this situation, finding roots of Coeff_E is equivalent to considering M cases where each factor is zero. In the k th case, we substitute for x_{i_k} its expression in terms of other variables into all the elements of Coeff_E . When this expression does not involve \bar{x}_{i_k} , we also substitute \bar{x}_{i_k} by the conjugate of the expression for x_{i_k} . After this substitution, Coeff_E is updated and the number of variables decreases. If this process can be repeated until all polynomials in Coeff_E either identically vanish or one of them is identically nonzero (such as $a_1\bar{a}_1 + b_2\bar{b}_2$), then we have obtained all common roots of the elements of Coeff_E and therefore all graphs that factor into the given form $P_c = P_a P_b$.

It turns out that, by this method and symmetry analysis in this section, all admissible graph operators with real coefficients and factorable dispersion function can be determined. But the algebraic equations become prohibitively complicated when the graph coefficients are allowed to be complex, even with a combination of computer and by-hand computations. Thus the only part of this study that remains incomplete is the determination of all graph operators associated to the tetrakis graph that have non-real coefficients and for which the Fermi surface is reducible.

For $E, a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{R}^*)^6$, we can obtain all factorizations of P_c into pair (4.37) by exploiting symmetry. When all these variables are real, the symmetry $D(z_1, z_2; E) = \bar{D}(z_1^{-1}, z_2^{-1}; E)$ becomes

$$D(z_1, z_2; E) = D(z_1^{-1}, z_2^{-1}; E). \quad (4.60)$$

Suppose that $D(z_1, z_2; E)$ factors into two Laurent, non-monomial, polynomials $D(z_1, z_2; E) = A(z_1, z_2)B(z_1, z_2)$. Then

$$D(z_1, z_2; E) = A(z_1, z_2)B(z_1, z_2) = A(z_1^{-1}, z_2^{-1})B(z_1^{-1}, z_2^{-1}). \quad (4.61)$$

Multiplying by $z_1^{m_0} z_2^{n_0}$ as in Step 3 of the algorithm yields

$$P_c(z_1, z_2) = z_1^{m_0} z_2^{n_0} A(z_1, z_2)B(z_1, z_2) = z_1^{m_0} z_2^{n_0} A(z_1^{-1}, z_2^{-1})B(z_1^{-1}, z_2^{-1}). \quad (4.62)$$

If P_c factors into pair (4.37), then (4.62) implies

$$P_c = F_1 F_2 = F_3 F_4, \quad (4.63)$$

in which

$$F_1 = a_{1,1} z_1 z_2 + a_{1,0} z_1 + a_{0,1} z_2 + a_{0,0}, \quad (4.64)$$

$$F_2 = b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2 + b_{0,0}, \quad (4.65)$$

$$F_3 = a_{1,1} + a_{1,0} z_2 + a_{0,1} z_1 + a_{0,0} z_1 z_2, \quad (4.66)$$

$$F_4 = b_{1,1} + b_{1,0} z_2 + b_{0,1} z_1 + b_{0,0} z_1 z_2. \quad (4.67)$$

If one of F_1, F_2, F_3 and F_4 factors, then it is reduced to pair (4.39). All the graphs in this case are parameterized by (4.56) with the values of $\{a_2, b_0, d_2, d_3\}$ being real.

Thus we need only to consider the situation in which F_1, F_2, F_3 and F_4 are all irreducible. In this case, we have $F_3|F_1$ or else $F_3|F_2$, since $\mathbb{C}[x, y]$ is a unique-factorization domain.

Claim 3. $F_3 \nmid F_1$

Proof. The assumption that $F_3|F_1$ implies that $(a_{1,1}, a_{1,0}, a_{0,1}, a_{0,0}) = C(a_{0,0}, a_{0,1}, a_{1,0}, a_{1,1})$. This leads to $C^2 = 1$ and therefore $C = \pm 1$. Thus P_c factors into either

$$\begin{aligned} P_a &= a_{0,0}z_1z_2 + a_{0,1}z_1 + a_{0,1}z_2 + a_{0,0} \\ P_b &= b_{1,1}z_1z_2 + b_{1,0}z_1 + b_{0,1}z_2 + b_{0,0} \end{aligned} \quad (4.68)$$

or else

$$\begin{aligned} P_a &= -a_{0,0}z_1z_2 - a_{0,1}z_1 + a_{0,1}z_2 + a_{0,0} \\ P_b &= b_{1,1}z_1z_2 + b_{1,0}z_1 + b_{0,1}z_2 + b_{0,0}. \end{aligned} \quad (4.69)$$

Performing Steps 5-8 of the algorithm for the pair (4.68) reveals that Coeff_E contains the element -1 ; and performing the algorithm for the pair (4.69) reveals that Coeff_E contains the element 1 . In either case, P_c is not factorable, which is inconsistent with the assumption (4.63). \square

In the case in which F_1, F_2, F_3 and F_4 are all irreducible and $F_3|F_2$, we have $(a_{1,1}, a_{1,0}, a_{0,1}, a_{0,0}) = C(b_{0,0}, b_{0,1}, b_{1,0}, b_{1,1})$ for some complex C . That is, P_c factors into

$$\begin{aligned} P_a &= C(b_{0,0}z_1z_2 + b_{0,1}z_1 + b_{1,0}z_2 + b_{1,1}), \\ P_b &= b_{1,1}z_1z_2 + b_{1,0}z_1 + b_{0,1}z_2 + b_{0,0}. \end{aligned} \quad (4.70)$$

Performing the algorithm for the pair (4.70), we obtain that Coeff_E consists of the following four polynomials in the variables $(a_0, a_1, a_2, b_0, d_1, d_2, d_3)$:

$$\begin{aligned} \text{Coeff}_E &= \{ (a_1a_2 + d_1d_2)(a_1a_2 + b_0d_3), \\ & 2a_1a_2^2b_0d_1 + 2a_1^2a_2b_0d_2 - a_0a_1a_2d_1d_2 + a_2b_0d_1^2d_2 + a_1b_0d_1d_2^2 - a_0a_1a_2b_0d_3 + 2a_1^2a_2d_1d_3 + a_2b_0^2d_1d_3 \\ & + 2a_1a_2^2d_2d_3 + a_1b_0^2d_2d_3 - 2a_0b_0d_1d_2d_3 + a_1d_1^2d_2d_3 + a_2d_1d_2^2d_3 + a_1b_0d_1d_3^2 + a_2b_0d_2d_3^2, \\ & a_2^2b_0^2d_1^2 + 3a_1a_2b_0^2d_1d_2 - a_0a_2b_0^2d_1^2d_2 - a_1a_2d_1^3d_2 + a_1^2b_0^2d_2^2 - a_0a_1b_0d_1d_2^2 + a_1^2d_1^2d_2^2 + a_2^2d_1^2d_2^2 \\ & + b_0^2d_1^2d_2^2 - a_1a_2d_1d_2^3 - a_1a_2b_0^3d_3 - a_0a_2b_0^2d_1d_3 + 3a_1a_2b_0d_1^2d_3 - a_0a_1b_0^2d_2d_3 + a_0^2b_0d_1d_2d_3 - b_0^3d_1d_2d_3 \\ & - a_0a_1d_1^2d_2d_3 - b_0d_1^3d_2d_3 + 3a_1a_2b_0d_2^2d_3 - a_0a_2d_1d_2^2d_3 - b_0d_1d_2^3d_3 + a_1^2b_0^2d_3^2 + a_2^2b_0^2d_3^2 - a_0a_1b_0d_1d_2^3 \\ & + a_1^2d_1^2d_3^2 + b_0^2d_1^2d_3^2 - a_0a_2b_0d_2d_3^2 + 3a_1a_2d_1d_2d_3^2 + a_2^2d_2^2d_3^2 + b_0^2d_2^2d_3^2 + d_1^2d_2^2d_3^2 - a_1a_2b_0d_3^3 - b_0d_1d_2d_3^3, \\ & a_2b_0^3d_1^2d_2 - a_2b_0d_1^4d_2 + a_1b_0^3d_1d_2^2 - a_0b_0^2d_1^2d_2^2 + a_1b_0d_1^3d_2^2 + a_2b_0d_1^2d_2^3 - a_1b_0d_1d_2^4 - a_2b_0^4d_1d_3 + a_2b_0^2d_1^3d_3 \\ & - a_1b_0^4d_2d_3 + a_0b_0^3d_1d_2d_3 - 2a_1b_0^2d_1^2d_2d_3 + a_0b_0d_1^3d_2d_3 - a_1d_1^4d_2d_3 - 2a_2b_0^2d_1d_2^2d_3 + a_2d_1^3d_2^2d_3 + a_1b_0^2d_2^3d_3 \\ & + a_0b_0d_1d_2^3d_3 + a_1d_1^2d_2^3d_3 - a_2d_1d_2^4d_3 + a_1b_0^3d_1d_2^3 - a_0b_0^2d_1^2d_2^3 + a_1b_0d_1^3d_2^3 + a_2b_0^3d_2d_2^3 - 2a_2b_0d_1^2d_2d_2^3 \\ & - a_0b_0^2d_2^2d_2^3 - 2a_1b_0d_1d_2^2d_2^3 - a_0d_1^2d_2^2d_2^3 + a_2b_0d_2^3d_2^3 + a_2b_0^2d_1d_2^3 + a_1b_0^2d_2d_2^3 + a_0b_0d_1d_2d_2^3 + a_1d_1^2d_2d_2^3 \\ & + a_2d_1d_2^2d_2^3 - a_1b_0d_1d_2^4 - a_2b_0d_2d_2^4 \}. \end{aligned} \quad (4.71)$$

We will find all simultaneous real roots $(a_0, a_1, a_2, b_0, d_1, d_2, d_3)$ to the set (4.71). The vanishing of the first element in Coeff_E implies that $a_1 = -d_1d_2/a_2$ or $a_1 = -b_0d_3/a_2$. Updated with $a_1 = -d_1d_2/a_2$, the second element in Coeff_E becomes

$$a_2^{-1}(-d_1d_2 + b_0d_3)(-a_0a_2d_1d_2 - d_1d_2(b_0d_2 + d_1d_3) + a_2^2(b_0d_1 + d_2d_3)). \quad (4.72)$$

Thus $d_1 = b_0d_3/d_2$ or $a_0 = (-d_1d_2(b_0d_2 + d_1d_3) + a_2^2(b_0d_1 + d_2d_3))/(a_2d_1d_2)$. With these two substitutions separately, the third element becomes

$$d_3^2(-a_0a_2b_0d_2 + a_2^2(b_0^2 + d_2^2) - b_0^2(d_2^2 + d_3^2))^2/(a_2^2d_2^2) \quad (4.73)$$

or

$$(a_2^2 + d_1^2)(a_2^2 + d_2^2)(d_1d_2 - b_0d_3)^2/a_2^2, \quad (4.74)$$

respectively. By setting (4.73) to zero and solving for a_0 or setting (4.74) to zero and solving for d_1 , one obtains exactly the same expressions for a_0 and d_1 . Thus in the case $a_1 = -d_1 d_2 / a_2$, we must have

$$\begin{aligned} a_0 &= (a_2^2(b_0^2 + d_2^2) - b_0^2(d_2^2 + d_3^2)) / (a_2 b_0 d_2), \\ a_1 &= -b_0 d_3 / a_2, \\ d_1 &= b_0 d_3 / d_2. \end{aligned} \quad (\text{parametrization for real coefficients}) \quad (4.75)$$

Similarly in the case that $a_1 = -b_0 d_3 / a_2$, we obtain exactly these same expressions for d_1 and a_0 (for this, the fourth element of Coeff_E is used). Thus, (4.75) holds in any case. It can be checked that all elements of Coeff_E vanish when (4.75) holds. Thus all solutions to Coeff_E for $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{R}^*)^6$ are given by (4.75).

Summarizing, we have proved in the case of real graph coefficients that if the polynomial P_c admits a factorization of the form (4.37) and the polynomials F_1 , F_2 , F_3 and F_4 are irreducible, then the graph coefficients must satisfy the relations (4.75). One finds that the set of real tuples $(a_0, a_1, a_2, b_0, d_1, d_2, d_3)$ that satisfy (4.75) coincide exactly with the relations (4.56) when all variables are real. Those relations correspond to the graphs for which one (and therefore all) of the F_i is reducible, and, as mentioned above, such graphs are subsumed by the factorization form (4.39) and therefore also of the form (4.38). This implies that the assumption that all F_i are irreducible is untenable.

This concludes the proof that all graphs with real coefficients and reducible Fermi surface must admit a factorization of the form (4.38) and that all such factorizations are parameterized by four real variables according to (4.75). Explicitly, the factorizations are

$$d_2^2 P_c = -(d_2 d_3 z_1^2 + d_2 d_3 + (d_3^2 + d_2^2 + a_2 d_2 E / b_0) z_1) (b_0 d_2 z_2^2 + b_0 d_2 + (d_2^2 + b_0^2 + b_0 d_2 E / a_2) z_2), \quad (4.76)$$

and the corresponding factorizations of the dispersion function are

$$D(z_1, z_2; E) = - \left(d_3(z_1 + z_1^{-1}) + \frac{d_3^2}{d_2} + d_2 + \frac{a_2}{b_0} E \right) \left(b_0(z_2 + z_2^{-1}) + d_2 + \frac{b_0^2}{d_2} + \frac{b_0}{a_2} E \right). \quad (4.77)$$

We end this section with a short discussion on complex graph coefficients, particularly, explaining why symmetry does not help with the search for all factorizations of P_c . A factorization of the form (4.37) implies eight variables in the polynomials in the set I in Step 5 of the algorithm. When $a_0 \in \mathbb{R}$ and $S_\Gamma \in (\mathbb{C}^*)^6$, this factorization, together with the symmetry $D(z_1, z_2; E) = \bar{D}(z_1^{-1}, z_2^{-1}; E)$ implies $P_c = F_1 F_2 = F_3 F_4$ with

$$F_1 = z_1^{m_1} z_2^{n_1} A(z_1, z_2) = a_{1,1} z_1 z_2 + a_{1,0} z_1 + a_{0,1} z_2 + a_{0,0}, \quad (4.78)$$

$$F_2 = z_1^{m_2} z_2^{n_2} B(z_1, z_2) = b_{1,1} z_1 z_2 + b_{1,0} z_1 + b_{0,1} z_2 + b_{0,0}, \quad (4.79)$$

$$F_3 = z_1^{m_3} z_2^{n_3} A(z_1^{-1}, z_2^{-1}) = \bar{a}_{1,1} + \bar{a}_{1,0} z_2 + \bar{a}_{0,1} z_1 + \bar{a}_{0,0} z_1 z_2, \quad (4.80)$$

$$F_4 = z_1^{m_4} z_2^{n_4} B(z_1^{-1}, z_2^{-1}) = \bar{b}_{1,1} + \bar{b}_{1,0} z_2 + \bar{b}_{0,1} z_1 + \bar{b}_{0,0} z_1 z_2. \quad (4.81)$$

If F_1 , F_2 , F_3 and F_4 are all irreducible, then either $F_3|F_1$ and $F_4|F_2$ or $F_3|F_2$ and $F_4|F_1$. In the case $F_3|F_2$ and $F_4|F_1$, we obtain

$$C(\bar{a}_{1,1}, \bar{a}_{1,0}, \bar{a}_{0,1}, \bar{a}_{0,0}) = (b_{0,0}, b_{0,1}, b_{1,0}, b_{1,1}), \quad (4.82)$$

where C is some complex number. This means that P_c factors into

$$\begin{aligned} P_a &= a_{1,1} z_1 z_2 + a_{1,0} z_1 + a_{0,1} z_2 + a_{0,0} \\ P_b &= C(\bar{a}_{0,0} z_1 z_2 + \bar{a}_{0,1} z_1 + \bar{a}_{1,0} z_2 + \bar{a}_{1,1}). \end{aligned} \quad (4.83)$$

Since complex conjugates must be treated as independent variables in the process of variable elimination using Groebner-basis routines, the symmetry has not reduced the number of variables in the polynomials of the set I in Step 5 of the algorithm. In fact, $GB_C(I)$ has not changed at all. Determining all factorizations of the form (4.37) for the tetrakis graph with complex coefficients remains unsolved.

5 Necessity of positivity and planarity: Reducible Fermi surfaces

The examples in this section demonstrate that the positivity of the coefficients and the planarity conditions for irreducibility of the Fermi surface cannot be dispensed with.

5.1 A non-planar graph with arbitrary coefficients

Non-planar graphs obtained by coupling two identical copies of a given periodic graph provide easy examples of reducible Fermi surfaces, as demonstrated in [14]. Fig. 11 depicts two identical copies of a (planar) square graph coupled by edges between corresponding vertices. The resulting graph is not planar. The matrices $A(n_1, n_2)$ for this example are

$$A(0,0) = \begin{bmatrix} 0 & b \\ \bar{b} & 0 \end{bmatrix}, \quad A(1,0) = \begin{bmatrix} a & 0 \\ 0 & a \end{bmatrix}, \quad A(0,1) = \begin{bmatrix} c & 0 \\ 0 & c \end{bmatrix}, \quad (5.84)$$

plus $A(-1,0) = A(1,0)^*$ and $A(0,-1) = A(0,1)^*$ being determined by self-adjointness and the rest vanishing. The Floquet transform for the operator $A - E$ is

$$\hat{A}(z_1, z_2) - E = \begin{bmatrix} az_1 + \bar{a}z_1^{-1} + cz_2 + \bar{c}z_2^{-1} - E & b \\ \bar{b} & az_1 + \bar{a}z_1^{-1} + cz_2 + \bar{c}z_2^{-1} - E \end{bmatrix}, \quad (5.85)$$

and its determinant is

$$\det(\hat{A}(z_1, z_2) - E) = (az_1 + \bar{a}z_1^{-1} + cz_2 + \bar{c}z_2^{-1} + |b| - E) (az_1 + \bar{a}z_1^{-1} + cz_2 + \bar{c}z_2^{-1} - |b| - E), \quad (5.86)$$

which demonstrates the reducibility of the Fermi surface for all energies E . This type of graph is a case of a more general construction of coupled graph operators with reducible Fermi surface [14].

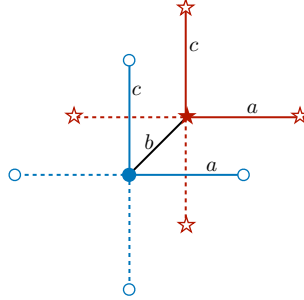


Figure 11: A periodic non-planar graph with reducible Fermi surface. The solid vertices and edges form a fundamental domain.

5.2 Planar graphs with indefinite coefficients

According to Theorem 1(a), the Fermi surface of admissible graphs with positive coefficients is always irreducible. But for the tetrakis graph, factorizations of the form (2.11) described in Theorem 1(b) and parameterized in section 4.5.1 can be obtained with indefinite graph coefficients. Fig. 12 shows some of these graphs with coefficients equal to ± 1 . The associated factorizations of the dispersion function are

$$\begin{aligned} D_{(a)}(E, z_1, z_2) &= (z_1 + z_1^{-1} + 2 - E) (z_2 + z_2^{-1} + 2 + E), \\ D_{(b)}(E, z_1, z_2) &= (z_1 + z_1^{-1} + 2 - E) (z_2 + z_2^{-1} - 2 - E), \\ D_{(c)}(E, z_1, z_2) &= (z_1 + z_1^{-1} - 2 + E) (z_2 + z_2^{-1} - 2 - E). \end{aligned}$$

These three graph operators are unitarily equivalent by gauge transformation since they have identical circuit fluxes, by extending [12, Lemma 2.1] to periodic graph operators. Their Fermi surfaces are transformed by $z_1 \mapsto z_1 e^{i\phi_1}$ and $z_2 \mapsto z_2 e^{i\phi_2}$, where ϕ_j is the phase advance along the j -th period vector. Each of these phases is 0 or $\pi \pmod{2\pi}$, depending on the sign of the coefficient along the horizontal or vertical edge. This implies that the spectra of these three operators are identical; in each case, the first factor of $D_{(j)}$

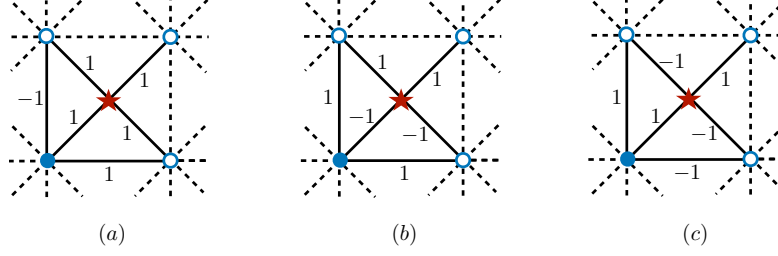


Figure 12: Doubly periodic planar graphs with indefinite coefficients and reducible Fermi surface. The solid vertices and edges form a fundamental domain, and ± 1 are coupling coefficients, as described in section 5.2.

($j \in \{a, b, c\}$) contributes the interval $E \in [0, 4]$, and the second contributes $E \in [-4, 0]$. The dependence of the spectrum of periodic discrete magnetic operators on fluxes is studied in [9].

Acknowledgment. This work was supported by NSF research Grants DMS-1411393 and DMS-1814902.

Note. On behalf of both authors, the corresponding author (SPS) asserts that there exists no conflict of interest regarding this article.

References

- [1] D. Bättig. *A toroidal compactification of the two-dimensional Bloch-manifold*. PhD thesis, ETH-Zürich, 1988.
- [2] D. Bättig. A toroidal compactification of the Fermi surface for the discrete Schrödinger operator. *Comment. Math. Helvetici*, 67:1–16, 1992.
- [3] D. Bättig, H. Knörrer, and E. Trubowitz. A directional compactification of the complex Fermi surface. *Compositio Math*, 79(2):205–229, 1991.
- [4] David Cox, John Little, and Donal O’Shea. *Ideals, Varieties, and Algorithms*. Springer, 2007.
- [5] Ngoc Do, Peter Kuchment, and Frank Sottile. Generic properties of dispersion relations for discrete periodic operators. *arXiv:1910.06472*, 2019.
- [6] Lee Fisher, Wei Li, and Stephen P. Shipman. Reducible fermi surface for multi-layer quantum graphs including stacked graphene. *arXiv:2005.13764 [math-ph]*, 2020.
- [7] D. Gieseke, H. Knörrer, and E. Trubowitz. *The Geometry of Algebraic Fermi Curves*. Academic Press, Boston, 1993.
- [8] P. G. Harper. Single band motion of conduction electrons in a uniform magnetic field. *Proceedings of the Physical Society. Section A*, 68(10):874–878, oct 1955.
- [9] Evgeny Korotyaev and Natalia Saburova. Magnetic Schrödinger operators on periodic discrete graphs. *Journal of Functional Analysis*, 272(4):1625 – 1660, 2017.
- [10] Peter Kuchment and Boris Vainberg. On absence of embedded eigenvalues for Schrödinger operators with perturbed periodic potentials. *Commun. Part. Diff. Equat.*, 25(9–10):1809–1826, 2000.
- [11] Peter Kuchment and Boris Vainberg. On the structure of eigenfunctions corresponding to embedded eigenvalues of locally perturbed periodic graph operators. *Comm. Math. Phys.*, 268(3):673–686, 2006.

- [12] Elliott H. Lieb and Michael Loss. Fluxes, Laplacians, and Kasteleyn's theorem. *Duke Mathematical Journal*, 71(2):337–363, August 1993.
- [13] Wencai Liu. Irreducibility of the Fermi variety for discrete periodic Schrödinger operators and embedded eigenvalues. *arXiv:2006.04733*, 2020.
- [14] Stephen P. Shipman. Eigenfunctions of unbounded support for embedded eigenvalues of locally perturbed periodic graph operators. *Comm. Math. Phys.*, 332(2):605–626, 2014.
- [15] Stephen P. Shipman. Reducible Fermi surfaces for non-symmetric bilayer quantum-graph operators. *J. Spectral Theory*, (in press), 2019.