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Algebraic Calderón-Zygmund theory

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ABSTRACT

Calderón-Zygmund theory has been traditionally developed on metric measure spaces satisfying additional regularity properties. In the lack of good metrics, we introduce a new approach for general measure spaces which admit a Markov semigroup satisfying purely algebraic assumptions. We shall construct an abstract form of ‘Markov metric’ governing the Markov process and the naturally associated BMO spaces, which interpolate with the L_p -scale and admit endpoint inequalities for Calderón-Zygmund operators. Motivated by noncommutative harmonic analysis, this approach gives the first form of Calderón-Zygmund theory for arbitrary von Neumann algebras, but is also valid in classical settings like Riemannian manifolds with nonnegative Ricci curvature or doubling/nondoubling spaces. Other less standard commutative scenarios like fractals or abstract probability spaces are also included. Among our applications in the noncommutative setting, we improve recent results for quantum Euclidean

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spaces and group von Neumann algebras, respectively linked to noncommutative geometry and geometric group theory.

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Introduction

The analysis of linear operators associated to singular kernels is a central topic in harmonic analysis and partial differential equations. A large subfamily of these maps is under the scope of Calderón-Zygmund theory, which exploits the relation between metric and measure in the underlying space to provide sufficient conditions for L_p boundedness. The Calderón-Zygmund decomposition [5] or the Hörmander smoothness condition for the kernel [26] combine the notions of proximity in terms of the metric with that of smallness in terms of the measure. The doubling and/or polynomial growth conditions between metric and measure yield more general forms of the theory [13,46,47,61,62]. To the best of our knowledge, the existence of a metric in the underlying space is always assumed in the literature.

In this paper, we introduce a new form of Calderón-Zygmund theory for general measure spaces admitting a Markov semigroup which only satisfies purely algebraic assumptions. This is especially interesting for measure spaces where the geometric information is poor. It includes abstract probability spaces or fractals like the Sierpinski gasket, where a Dirichlet form is defined. It is also worth mentioning that our approach recovers Calderón-Zygmund theory for classical spaces and provides alternative forms over them. In spite of these promising directions —very little explored here— our main motivation has been to develop a noncommutative form of Calderón-Zygmund theory for noncommutative measure spaces (von Neumann algebras) where the notions of quantum metric [37,55,56] seem inefficient.

A great effort has been done over the last years to produce partial results towards a noncommutative Calderón-Zygmund theory [21,25,29,31,45,49]. The model cases considered so far are all limited to (different) noncommutative forms of Euclidean spaces, described as follows:

A) Tensor products. Let $f = (f_{jk}) : \mathbb{R}^n \rightarrow \mathcal{B}(\ell_2)$ be a matrix-valued function and consider the tensor product extension of a standard Calderón-Zygmund operator acting on f , formally given by

$$Tf(x) = \int_{\mathbb{R}^n} k(x, y) f(y) dy = \left(Tf_{jk}(x) \right)_{jk} \quad \text{for } x \notin \text{supp} f.$$

The L_p -boundedness of this map in the associated (tensor product) von Neumann algebra $\mathcal{M} = L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{B}(\ell_2)$ trivially follows for $p > 1$ from the vector-valued theory,

due to the UMD nature of Schatten p -classes. On the contrary, no endpoint estimate for $p = 1$ is possible using vector-valued methods. The original argument in [49] —also in a recent simpler form [3]— combines noncommutative martingales with a pseudolocalization principle for classical Calderón-Zygmund operators. More precisely, a quantification of how much L_2 -mass of a singular integral is concentrated around the support of the function on which it acts. This inequality has been the key tool in the recent solution of the Nazarov-Peller conjecture [7], a strengthening of the celebrated Krein conjecture [58] on operator Lipschitz functions.

B) Crossed products. New L_p estimates for Fourier multipliers in group von Neumann algebras have recently gained considerable momentum for its connections to geometric group theory. The first Hörmander-Mikhlin type theorem in this direction [31] exploited finite-dimensional cocycles of the given group G to transfer the problem to the cocycle Hilbert space $\mathcal{H} = \mathbb{R}^n$. To find sufficient regularity for L_p -boundedness amounts to study Calderón-Zygmund operators in the crossed products $L_\infty(\mathbb{R}^n) \rtimes G$ induced by the cocycle action. Nonequivariant extensions of CZOs on these von Neumann algebras were investigated in [31], after identifying the right BMO space for the length function determined by the cocycle. These operators have the form

$$\int_G f_g \rtimes \lambda(g) d\mu(g) \mapsto \int_G T_g(f_g) \rtimes \lambda(g) d\mu(g).$$

Here μ, λ respectively denote the Haar measure and left regular representation on the (unimodular) group G , whereas $T_g = \alpha_g T \alpha_{g^{-1}}$ is a twisted form of a classical CZO T on \mathbb{R}^n by the cocycle action α . We refer to [32,50] for further results.

C) Quantum deformations. PDEs in matrix algebras and ‘noncommutative manifolds’ appear naturally in theoretical physics. Pseudodifferential operators were introduced by Connes in 1980 to study a quantum form of the Atiyah-Singer index theorem over these algebras. These techniques have been underexploited over the last 30 years, due to fundamental obstructions to understand singular integral theory in this context. The core of singular integrals and pseudodifferential operator L_p -theory was developed in [21] over the archetypal algebras of noncommutative geometry. It includes quantum tori, Heisenberg-Weyl algebras and other quantum deformations of \mathbb{R}^n of great interest in quantum field theory, string theory and quantum probability. This was the first approach to a ‘fully noncommutative’ Calderón-Zygmund theory for CZOs not acting on copies of \mathbb{R}^n as tensor or crossed product factors, but still related to Euclidean methods.

We introduce in this paper the first form of Calderón-Zygmund theory valid for general (semifinite) von Neumann algebras. As we explained above, the main difficulty arises from the lack of very standard geometric tools, like the existence of a nice underlying metric or the construction of suitable covering lemmas. We shall circumvent it using a very different approach based on algebraic properties of a given Markov process. Our applications cover a wide variety of scenarios which will be discussed, giving especial emphasis to noncommutative forms of Euclidean spaces and locally compact abelian groups,

which are our main classical models. In the first case, we shall weaken/optimize the CZ conditions on quantum Euclidean spaces [21]. In the second case, LCA groups correspond to quantum groups which are both commutative and cocommutative [59]. We shall give CZ conditions for convolution maps over quantum groups. In the cocommutative (non necessarily commutative) context, this includes group von Neumann algebras.

Calderón-Zygmund extrapolation

Based on the behavior of the Hilbert transform in the real line, the main goal of Calderón-Zygmund theory is to establish regularity properties on the kernel of a singular integral operator, so that L_2 -boundedness automatically extrapolates to L_p boundedness for $1 < p < \infty$. A singular integral operator in a Riemannian manifold (X, d, μ) admits the kernel representation

$$T_k f(x) = \int_X k(x, y) f(y) d\mu(y) \quad \text{for } x \notin \text{supp } f.$$

Namely, T_k is only assumed a priori to send test functions into distributions, so that it admits a distributional kernel in $X \times X$ which coincides in turn with a locally integrable function k away from the diagonal $x = y$, where the kernel presents certain singularity. This already justifies the assumption $x \notin \text{supp } f$ in the kernel representation. The work in [5,26] culminated in the following sufficient conditions on a singular integral operator in \mathbb{R}^n for its L_p -boundedness:

i) L_2 -boundedness

$$\|T_k : L_2(\mathbb{R}^n) \rightarrow L_2(\mathbb{R}^n)\| < \infty.$$

ii) Hörmander kernel smoothness

$$\text{ess sup}_{x, y \in \mathbb{R}^n} \int_{|x-z| > 2|x-y|} |k(x, z) - k(y, z)| + |k(z, x) - k(z, y)| dz < \infty.$$

Historically, this was used to prove a weak endpoint inequality in L_1 . The same holds for Riemannian manifolds with nonnegative Ricci curvature [1]. Alternatively it is simpler to use L_2 -boundedness and the kernel smoothness condition to prove $L_\infty \rightarrow \text{BMO}$ boundedness. The result then follows by well-known duality and interpolation arguments. Our strategy resembles this approach:

P1. Identify the appropriate BMO spaces.

P2. Prove the expected interpolation results with L_p spaces.

P3. Provide conditions on CZO's which yield $L_\infty \rightarrow \text{BMO}$ boundedness.

In the classical setting, we typically find H_1/BMO spaces associated to a metric or a martingale filtration. Duong and Yan [17,18] extended this theory replacing some averages over balls in the metric space by semigroups of positive operators, although the existence of a metric was still assumed. This assumption was removed in [30,44] providing a theory of semigroup type BMO spaces with no further assumptions on the given space. In particular, we could say that Problems 1 and 2 were solved in [30], but it has been unclear since then how to provide natural CZ conditions which imply $L_\infty \rightarrow \text{BMO}$ estimates. In this paper we solve P3 by splitting it into:

- P3a. Construct a ‘metric’ governing the Markov process.
- P3b. Define ‘metric BMO’ spaces which still interpolate with the L_p scale.
- P3c. Provide CZ conditions giving $L_\infty \rightarrow \text{BMO}$ boundedness for metric BMO’s.

P3a. Markov metrics. Given a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ on the semifinite von Neumann algebra (\mathcal{M}, τ) —in other words, formed by normal self-adjoint cpu maps S_t — we introduce a *Markov metric* for it as any family

$$\mathcal{Q} = \left\{ (R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+ \right\}$$

composed of normal completely positive unital (cpu) maps $R_{j,t} : \mathcal{M} \rightarrow \mathcal{M}$ and elements $\sigma_{j,t}, \gamma_{j,t}$ of \mathcal{M} with $\gamma_{j,t} \geq \mathbf{1}_{\mathcal{M}}$, such that the following estimates (which show how the Markov metric governs the Markov semigroup in a controlled way) hold:

- i) *Hilbert module majorization*: $\langle \xi, \xi \rangle_{S_t} \leq \sum_{j \geq 1} \sigma_{j,t}^* \langle \xi, \xi \rangle_{R_{j,t}} \sigma_{j,t}$.
- ii) *Metric integrability condition*: $k_{\mathcal{Q}} = \sup_{t > 0} \left\| \sum_{j \geq 1} \sigma_{j,t}^* \gamma_{j,t}^2 \sigma_{j,t} \right\|_{\mathcal{M}}^{\frac{1}{2}} < \infty$.

Here $\langle \cdot, \cdot \rangle_{\Phi}$ is the \mathcal{M} -valued inner product on $\mathcal{M} \bar{\otimes} \mathcal{M}$ for any cpu map Φ , given by $\langle a \otimes b, a' \otimes b' \rangle_{\Phi} = b^* \Phi(a^* a') b'$. Markov metrics are a priori unrelated to Rieffel’s quantum metric spaces [55,56]. They present on the contrary vague similarities with abstract formulations of classical CZ theory in the absence of CZ kernels and/or doubling measures [2,61]. We shall explain what motivates our definition below and we shall also illustrate how Euclidean and other classical metrics fit in.

P3b. Metric type BMO spaces. Let

$$\|f\|_{\text{BMO}_{\mathcal{S}}^c} = \sup_{t \geq 0} \left\| \left(S_t(f^* f) - (S_t f)^* (S_t f) \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}$$

and $\|f\|_{\text{BMO}_{\mathcal{S}}} = \max\{\|f\|_{\text{BMO}_{\mathcal{S}}^c}, \|f^*\|_{\text{BMO}_{\mathcal{S}}^c}\}$. We shall define the semigroup type BMO space $\text{BMO}_{\mathcal{S}}(\mathcal{M})$ as the weak-* closure of \mathcal{M} in certain direct sum of Hilbert modules determined by $\mathcal{S} = (S_t)_{t \geq 0}$. $\text{BMO}_{\mathcal{S}}(\mathcal{M})$ interpolates with $L_p(\mathcal{M})$ under certain notion

of regularity [30] to be recalled below. Given a Markov metric \mathcal{Q} associated to this semigroup, let us define in addition

$$\|f\|_{\text{BMO}_{\mathcal{Q}}} = \max \left\{ \|f\|_{\text{BMO}_{\mathcal{Q}}^c}, \|f^*\|_{\text{BMO}_{\mathcal{Q}}^c} \right\},$$

$$\|f\|_{\text{BMO}_{\mathcal{Q}}^c} = \sup_{t>0} \inf_{\substack{M_t \text{ cpu} \\ M_t: \mathcal{M} \rightarrow \mathcal{M}}} \sup_{j \geq 1} \left\| \left(\gamma_{j,t}^{-1} [R_{j,t}|f|^2 - |R_{j,t}f|^2 + |R_{j,t}f - M_tf|^2] \gamma_{j,t}^{-1} \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}.$$

Theorem A1. *Let (\mathcal{M}, τ) be a semifinite von Neumann algebra equipped with a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$. Let us consider a Markov metric \mathcal{Q} associated to $\mathcal{S} = (S_t)_{t \geq 0}$. Then, we find*

$$\|f\|_{\text{BMO}_{\mathcal{S}}} \lesssim k_{\mathcal{Q}} \|f\|_{\text{BMO}_{\mathcal{Q}}}.$$

If $\mathcal{S} = (S_t)_{t \geq 0}$ is regular, then $\text{BMO}_{\mathcal{Q}}(\mathcal{M}) \subset \text{BMO}_{\mathcal{S}}(\mathcal{M})$ interpolates with $L_p(\mathcal{M})$.

Theorem A1 solves P3b. Its proof is not hard after having defined the right notion of Markov metric and the right BMO norm. Let us note in passing that the term $R_{j,t}f - M_tf$ is there to accommodate nondoubling spaces to our definition in the spirit of Tolsa's RBMO space [61]. As a consequence of Theorem A1, proving $L_{\infty} \rightarrow \text{BMO}$ boundedness for metric BMO's (Problem 3c) implies the same result for semigroup BMO spaces (Problem 3). Of course, one could try to prove such a statement directly, but it seems that the metric/measure relation found with these new notions is crucial for a noncommutative CZ theory.

P3c. Calderón-Zygmund operators. The commutative idea behind the notion of Markov metric (explained in more detail in the body of the paper) is to find pointwise majorants of the integral kernels of our semigroup $\mathcal{S} = (S_t)_{t \geq 0}$, so that we can dominate S_t by certain sum of averaging operators over a distinguished family of measurable sets $\Sigma_{j,t}(x)$. These sets may be considered as the 'balls' in the Markov metric. In the noncommutative setting, this pointwise estimates must be written in terms of the given Hilbert module majorization and the cpu maps $R_{j,t}$ must be averages over certain projections $q_{j,t}$. Making this precise in full generality is one of the challenges of our algebraic approach and too technical to be explained at this point of the paper. A simple model case is given by

$$R_{j,t}f = \left((id \otimes \tau)(q_{j,t}) \right)^{-\frac{1}{2}} (id \otimes \tau)(q_{j,t}(\mathbf{1} \otimes f)q_{j,t}) \left((id \otimes \tau)(q_{j,t}) \right)^{-\frac{1}{2}} \quad (\text{Avg})$$

for certain family of projections $q_{j,t} \in \mathcal{M} \bar{\otimes} \mathcal{M}$. The linear map $\widehat{R}_{j,t}(\mathbf{1} \otimes f) = R_{j,t}f$ trivially amplifies to $\mathcal{M} \bar{\otimes} \mathcal{M}$. We may also consider similar formulas for the cpu maps M_t in the metric BMO norm. (Avg) allows to identify the Markov metric in terms of the 'balls' $q_{j,t}$ instead of the corresponding averaging maps $R_{j,t}$.

As it happens in classical Calderón-Zygmund theory, we need to impose some additional properties in the Markov metric to establish a good relation with the underlying

(noncommutative) measure. We have split these into *algebraic* and *analytic* conditions, further details will be given in the text. Let us just mention that the algebraic ones are inherent to noncommutativity and hold trivially in commutative cases. The analytic ones provide forms of Jensen's inequality and a measure/metric growth condition. Once we know the Markov metric satisfies these conditions, we may introduce Calderón-Zygmund operators. Assume that $T(\mathcal{A}_{\mathcal{M}}) \subset \mathcal{M}$ for a map T acting on a weak-* dense subalgebra $\mathcal{A}_{\mathcal{M}} \subset \mathcal{M}$ such that $\mathcal{A}_{\mathcal{M}}$ has dense intersection with $L_p(\mathcal{M})$ for any $1 \leq p < \infty$. The goal is to establish sufficient *Calderón-Zygmund conditions* on T for $L_\infty \rightarrow \text{BMO}_c$ boundedness. These are noncommutative forms of standard properties. Again, it is unnecessary to introduce them here in full generality, we do it in Section 2. In the model case above, our CZ conditions are:

i) $L_\infty(L_2^\mathcal{S})$ -boundedness

$$\left\| (id \otimes \tau) \left((id \otimes T)(x)^* (id \otimes T)(x) \right)^{\frac{1}{2}} \right\|_{\mathcal{M}} \lesssim \left\| (id \otimes \tau)(x^* x)^{\frac{1}{2}} \right\|_{\mathcal{M}},$$

for $x \in \mathcal{M} \bar{\otimes} \mathcal{M}$ such that $\|(id \otimes \tau)(x^* x)\|_{\mathcal{M}} < \infty$.

ii) *Size 'kernel' conditions*

- $\widehat{M}_t \left(|(id \otimes T)((1 \otimes f)(A_{j,t} - a_t))|^2 \right) \lesssim \gamma_{j,t}^2 \|f\|_{\mathcal{M}}^2,$
- $\widehat{R}_{j,t} \left(|(id \otimes T)((1 \otimes f)(A_{j,t} - a_{j,t}))|^2 \right) \lesssim \gamma_{j,t}^2 \|f\|_{\mathcal{M}}^2,$

for any $f \in \mathcal{A}_{\mathcal{M}}$ and certain family of operators $A_{j,t}, a_{j,t} \in \mathcal{M} \bar{\otimes} \mathcal{M}$ with $A_{j,t} \geq a_{j,t}$.

iii) *Hörmander 'kernel' conditions*

- $\Phi_{j,t} \left(|\delta((id \otimes T)((1 \otimes f)(1 - a_{j,t})))|^2 \right) \lesssim \gamma_{j,t}^2 \|f\|_{\mathcal{M}}^2,$
- $\Psi_{j,t} \left(|\delta((id \otimes T)((1 \otimes f)(1 - A_{j,t})))|^2 \right) \lesssim \gamma_{j,t}^2 \|f\|_{\mathcal{M}}^2,$

for certain family of normal cpu linear maps $\Phi_{j,t}, \Psi_{j,t} : \mathcal{M} \bar{\otimes} \mathcal{M} \rightarrow \mathcal{M}$.

In condition ii), $A_{j,t}$ and $a_{j,t}$ play the role of 'dilated balls' from $q_{j,t}$. In the last condition, δ is the derivation $x \mapsto x \otimes 1 - 1 \otimes x$ acting on the second leg of the tensor product. In the Euclidean case, these conditions reduce to L_2 -boundedness and the classical size/smoothness conditions for the kernel. Our general conditions include many more amplification algebras and derivations, other than $\mathcal{M} \bar{\otimes} \mathcal{M}$ and δ . Any map $T : \mathcal{A}_{\mathcal{M}} \rightarrow \mathcal{M}$ satisfying the above CZ-conditions will be called a *column CZ-operator*.

Theorem A2. *Let (\mathcal{M}, τ) be a semifinite von Neumann algebra equipped with a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ with associated Markov metric \mathcal{Q} fulfilling our algebraic and analytic assumptions. Then, any column CZ-operator T defines a bounded operator*

$$T : \mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{\mathcal{Q}}^c(\mathcal{M}).$$

If \mathcal{S} is regular, interpolation and duality give similar conditions for L_p -boundedness.

A generalized form of Theorem A2 is the main result of this paper. It is easy to recover Euclidean CZ-extrapolation from it. In the Euclidean and many other doubling scenarios, the size kernel condition ii) does not play any role. Our next goal is to explore how the general form of Theorem A2 applies in concrete von Neumann algebras with specific Markov metrics. It is worth mentioning that our regularity assumption holds for all the relevant examples in this paper.

Applications

Algebraic Calderón-Zygmund theory applies in classical and noncommutative measure spaces. In the commutative context, we shall limit ourselves to prove that algebraic and classical theories match in three important cases: Euclidean spaces with both Lebesgue or Gaussian measures and Riemannian manifolds with non-negative Ricci curvature. We shall not explore further implications in new commutative scenarios, like abstract probability spaces or fractals equipped with specific Dirichlet forms. In the noncommutative context, we start by analyzing the model case of matrix-valued functions from a very general viewpoint. We also consider Calderón-Zygmund operators over matrix algebras, generalizing triangular truncations as the archetype of singular integral operator. Most importantly, our abstract theory applies to quantum Euclidean spaces and quantum groups, which constitute our main motivations in this paper.

It will be useful to specify the form that our Calderón-Zygmund operators take when come associated to a concrete kernel. Our applications below include CZ conditions on the kernel. In the basic model case above, we set

$$T_k f = (id \otimes \tau)(k(\mathbf{1} \otimes f)) \quad (\text{Ker 1})$$

for some kernel k affiliated to $\mathcal{M} \bar{\otimes} \mathcal{M}_{\text{op}}$. Recall that the opposite structure $(\mathcal{M}_{\text{op}}$ is the same algebra \mathcal{M} endowed with the reversed product $a \cdot b = ba$) in the second tensor leg of the kernel for this (standard) model was already justified in [21]. It is a feature of CZ theory which can only be witnessed in noncommutative algebras. It will also be useful to generalize a bit our model case before analyzing any concrete application. Consider an auxiliary von Neumann algebra \mathcal{A} equipped with a n.s.f. trace φ , a $*$ -homomorphism $\sigma : \mathcal{M} \rightarrow \mathcal{A} \bar{\otimes} \mathcal{M}$ and the representation

$$S_{\tilde{k}} f = (id \otimes \varphi)(\tilde{k} \text{ flip} \circ \sigma(f)) \quad (\text{Ker 2})$$

for some kernel \tilde{k} affiliated to $\mathcal{M} \bar{\otimes} \mathcal{A}_{\text{op}}$. Of course, when $\mathcal{A} = \mathcal{M}$ and $\sigma(f) = \mathbf{1} \otimes f$ we recover our model case above, with kernel representation (Ker 1). This more general framework requires to redefine $R_{j,t}$ in (Avg) and the CZ conditions, as we shall do in the body of the paper. The advantage is to take \mathcal{A} as an elementary (commutative) algebra, from which we can transfer metric information. One may think of σ as a corepresentation in the terminology of quantum groups. Theorem A2 still holds in this case. We shall refer to *intrinsic* or *transferred* theories when using the model case $\mathcal{A} = \mathcal{M}$ or its generalization respectively.

Quantum Euclidean spaces. As geometrical spaces with noncommuting spatial coordinates, quantum Euclidean spaces have appeared frequently in the literature of mathematical physics, in the contexts of string theory and noncommutative field theory. These algebras play the role of a central and testing example in noncommutative geometry as well. The singular integral operators on quantum Euclidean spaces naturally appear in the recent study of Connes' quantized calculus [40,42,57] and noncommutative harmonic analysis [10,21,22,64]. Let

$$\Theta \in M_n(\mathbb{R})$$

be anti-symmetric. Briefly, the quantum Euclidean space \mathcal{R}_Θ is the von Neumann algebra generated by certain family of unitaries $\{u_j(s) : 1 \leq j \leq n, s \in \mathbb{R}\}$ satisfying

$$\begin{aligned} u_j(s)u_j(t) &= u_j(s+t), \\ u_j(s)u_k(t) &= e^{2\pi i \Theta_{jk}st} u_k(t)u_j(s). \end{aligned}$$

Define $\lambda_\Theta(\xi) = u_1(\xi_1)u_2(\xi_2) \cdots u_n(\xi_n)$ and set

$$f = \int_{\mathbb{R}^n} \check{f}_\Theta(\xi) \lambda_\Theta(\xi) d\xi = \lambda_\Theta(\check{f}_\Theta),$$

for $\check{f}_\Theta \in \mathcal{C}_c(\mathbb{R}^n)$. The trace on \mathcal{R}_Θ is determined by

$$\tau_\Theta(f) = \tau_\Theta \left(\int_{\mathbb{R}^n} \check{f}_\Theta(\xi) \lambda_\Theta(\xi) d\xi \right) = \check{f}_\Theta(0).$$

When $\Theta = 0$, $L_p(\mathcal{R}_\Theta, \tau_\Theta)$ reduces to $L_p(\mathbb{R}^n)$ with the Lebesgue measure. Precise definitions and a theory of singular integrals for \mathcal{R}_Θ appear in [21]. The main result relies on gradient kernel conditions for the intrinsic model (Ker 1). Remarkably, we show in this paper that the transference model (Ker 2)

$$\sigma_\Theta : \mathcal{R}_\Theta \ni \lambda_\Theta(\xi) \mapsto \exp_\xi \otimes \lambda_\Theta(\xi) \in L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta$$

goes further, since it just requires Hörmander type smoothness for the kernel. Here \exp_ξ stands for the ξ -th character $\exp(2\pi i \langle \xi, \cdot \rangle)$ in \mathbb{R}^n . There is a close relation between both models in this case

$$\tilde{\pi}_\Theta(m \otimes \exp_\xi) = m \lambda_\Theta(\xi)^* \otimes \lambda_\Theta(\xi), \quad \forall m \in \mathcal{R}_\Theta \quad \text{and} \quad T_k(f) = S_{\tilde{k}}(f) \quad \text{for } k = \tilde{\pi}_\Theta(\tilde{k}).$$

Another crucial map is the $*$ -homomorphism

$$\pi_\Theta : L_\infty(\mathbb{R}^n) \ni \exp_\xi \mapsto \lambda_\Theta(\xi) \otimes \lambda_\Theta(\xi)^* \in \mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta^{\text{op}}.$$

If B_R denotes the Euclidean R -ball centered at the origin, define the projections $a_R = \pi_\Theta(1_{5B_R})$ and $a_R^\perp = 1 - a_R$. Set $k_\sigma = (\sigma_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})(k) \in L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta^{\text{op}}$ and define the derivation $\delta\varphi(x, y) = \varphi(x) - \varphi(y)$ to set the kernel condition in $L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta^{\text{op}}$

$$\sup_{|x| \leq R, |y| \leq R} \left| \delta \left((id \otimes id \otimes \tau_\Theta) [k_\sigma(1 \otimes 1 \otimes f)(1 \otimes a_R^\perp)] \right) (x, y) \right| \lesssim \|f\|_{\mathcal{R}_\Theta}. \quad (\text{Hör})$$

As we shall justify in the paper, (Hör) is the right form of Hörmander kernel condition in this framework. The column BMO-norm admits in \mathcal{R}_Θ an equivalent form

$$\|f\|_{\text{BMO}_c(\mathcal{R}_\Theta)} \approx \|\sigma_\Theta(f)\|_{\text{BMO}_c(\mathbb{R}^n; \mathcal{R}_\Theta)}$$

for the operator-valued BMO space $\text{BMO}_c(\mathbb{R}^n; \mathcal{R}_\Theta)$ from [43]. These are all the ingredients to obtain Calderón-Zygmund extrapolation over quantum Euclidean spaces. Namely, the general form of Theorem A2 then yields the following theorem.

Theorem B1. T_k is bounded from \mathcal{R}_Θ to $\text{BMO}_c(\mathcal{R}_\Theta)$ provided:

- i) T_k is bounded on $L_2(\mathcal{R}_\Theta)$.
- ii) The kernel condition (Hör) holds.

Interpolation and duality give similar (symmetrized) conditions for L_p -boundedness.

Theorem B1 improves the main CZ extrapolation theorem in [21] by reducing the gradient kernel condition there to the (more flexible) Hörmander integral condition above, as we shall prove along the paper. In fact, the result which we shall finally prove is slightly more general than the statement above.

Quantum groups. Let G be a locally compact group with a left invariant Haar measure μ . When G is abelian, the Fourier transform carries the convolution algebra $L_1(G, \mu)$ into the multiplication algebra $L_\infty(\widehat{G}, \widehat{\mu})$ associated to the dual group with its (normalized) Haar measure. However, when G is not abelian, we can not construct the dual group and the multiplication algebra above becomes the group von Neumann algebra which is generated by the left regular representation of G . These algebras are basic models of (noncommutative, but still cocommutative) quantum groups, over which we shall study singular integrals.

Let \mathbb{G} be a locally compact quantum group of Kac type —precise definitions in the body of the paper— with comultiplication Δ and left-invariant and right-invariant Haar weights ψ, φ . Given a weak- $*$ dense subalgebra \mathcal{A} of $L_\infty(\mathbb{G})$ and a linear map T satisfying $T(\mathcal{A}) \subset L_\infty(\mathbb{G})$, it is a convolution map when

$$(T \otimes id_{\mathbb{G}}) \circ \Delta = \Delta \circ T = (id_{\mathbb{G}} \otimes T) \circ \Delta.$$

To simplify the problem, we shall consider the case where \mathbb{G} admits an α -doubling intrinsic Markov metric. That is, the projections which generate the cpu maps $R_{j,t}$'s satisfy

$$\frac{\psi(q_{\alpha(j),t})}{\psi(q_{j,t})} \leq c_\alpha$$

for a strictly increasing function $\alpha : \mathbb{N} \rightarrow \mathbb{N}$ with $\alpha(j) > j$ and a constant c_α .

Theorem B2. *Let \mathbb{G} be a locally compact quantum group of Kac type and assume it comes equipped with a convolution semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ which admits an α -doubling intrinsic Markov metric. Let $T : \mathcal{A} \rightarrow L_\infty(\mathbb{G})$ be a convolution map defined on a weakly dense $*$ -subalgebra \mathcal{A} of $L_\infty(\mathbb{G})$ such that*

- i) T is bounded on $L_2(\mathbb{G})$.
- ii) $\frac{1}{|\psi(q_{j,t})|^2} (\psi \otimes \psi) \left((q_{j,t} \otimes q_{j,t}) |\delta(T(fq_{\alpha(j),t})^\perp)|^2 \right) \lesssim \|f\|_{L_\infty(\mathbb{G})}^2$.

Then, the linear map T extends to a bounded map $T : L_\infty(\mathbb{G}) \rightarrow \text{BMO}_\mathcal{S}^c(L_\infty(\mathbb{G}))$.

As usual, L_p estimates follow from symmetrized conditions by interpolation and duality when the Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ is regular. In fact, we shall prove a more general statement which incorporates tensor products with an additional algebra (\mathcal{M}, τ) . Theorem B2 is proved one more time from Theorem A2. In fact, it is conceivable to remove the α -doubling restriction and still make the convolution map bounded under an additional size kernel condition.

Noncommutative transference. In a different direction, we shall finish this paper with a section devoted to noncommutative forms of Calderón-Cotlar method of transference [4,12,14]. The basic idea is to transfer L_p estimates of convolution maps on quantum groups to a much wider class of maps which arise by transference. We refer to [6,8,10,48,50,54] for other forms of transference in the context of group von Neumann algebras and quantum tori.

1. Markov metrics

An abstract form of Calderón-Zygmund theory incorporating noncommutative algebras lacks standard geometrical tools. Given a Markov semigroup on a von Neumann algebra—a semigroup of normal cpu self-adjoint maps on the given algebra—we shall construct a ‘metric’ governing the Markov process. Our model case in a commutative measure space (Ω, μ) is a Markov semigroup of linear maps of the form

$$S_t f(x) = \int_{\Omega} s_t(x, y) f(y) d\mu(y).$$

The idea is to find pointwise majorants of the form

$$s_t(x, y) \leq \sum_{j=1}^{\infty} \frac{|\sigma_{j,t}(x)|^2}{\mu(\Sigma_{j,t}(x))} \chi_{\Sigma_{j,t}(x)}(y), \quad (1.1)$$

so that $S_t f(x)$ is dominated by a given combination of averaging operators over certain measurable sets $\Sigma_{j,t}(x)$. These sets will determine some sort of metric on (Ω, μ) under additional integrability properties. Naively, we may think of them as balls or coronas around x in the hidden metric with radii depending on (j, t) . In this section we formalize this idea and construct BMO spaces with respect to the associated ‘Markov metric’ which satisfy the expected interpolation results.

1.1. Hilbert modules

A noncommutative measure space is a pair (\mathcal{M}, τ) formed by a semifinite von Neumann algebra \mathcal{M} and a n.s.f. trace τ . We assume in what follows that the reader is familiar with basic terminology from noncommutative integration theory [36,60]. Non-expert readers may proceed by fixing a measure space (Ω, μ) with $\mathcal{M} = L_{\infty}(\Omega)$ and τ the integral operator associated to μ . Given a cpu map $\Phi : \mathcal{M} \rightarrow \mathcal{M}$, we may construct the $L_p(\mathcal{M})$ -Hilbert module $L_p(\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M})$. Namely consider the seminorm on $\mathcal{M} \otimes \mathcal{M}$

$$\|\xi\|_{L_p(\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M})} = \|\sqrt{\langle \xi, \xi \rangle_{\Phi}}\|_{L_p(\mathcal{M})}$$

determined by the $L_{\frac{p}{2}}(\mathcal{M})$ -valued inner product

$$\left\langle \sum_j a_j \otimes b_j, \sum_k a'_k \otimes b'_k \right\rangle_{\Phi} = \sum_{j,k} b_j^* \Phi(a_j^* a'_k) b'_k.$$

Then $L_p(\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M})$, $1 \leq p < \infty$ stands for the norm completion of the algebraic tensor product, and $\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}$ denotes for the completion in the strong operator topology determined by $\xi_{\alpha} \rightarrow \xi$ when $\tau(\langle \xi - \xi_{\alpha}, \xi - \xi_{\alpha} \rangle_{\Phi} g) \rightarrow 0$ for all $g \in L_1(\mathcal{M})$. When Φ is normal, $\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}$ is the dual space of $L_1(\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M})$ (see [34]) and the abstract characterization of Hilbert modules [51] yields a weak-* continuous right \mathcal{M} -module map $\rho : \mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M} \rightarrow L_{\infty}(\mathcal{M}; \mathcal{H}^c)$, which takes values in the Hilbert space valued noncommutative L_{∞} space $L_{\infty}(\mathcal{M}; \mathcal{H}^c)$ (see e.g. [28]) and satisfies $\langle \xi, \eta \rangle_{\Phi} = \rho(\xi)^* \rho(\eta)$. Let us collect a few properties which will be instrumental along this paper.

Lemma 1.1. *Given a normal cpu map $\Phi : \mathcal{M} \rightarrow \mathcal{M}$:*

- i) $\langle \xi_1 + \xi_2, \xi_1 + \xi_2 \rangle_{\Phi} \leq 2\langle \xi_1, \xi_1 \rangle_{\Phi} + 2\langle \xi_2, \xi_2 \rangle_{\Phi},$
- ii) $\|f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes \Phi f\|_{\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}} = \|\Phi|f|^2 - |\Phi f|^2\|_{\mathcal{M}}^{\frac{1}{2}},$
- iii) $|\Phi f - g|^2 \leq \langle f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes g, f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes g \rangle_{\Phi},$

- iv) $\|f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes \Phi f\|_{\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}} \sim \inf_{g \in \mathcal{M}} \|f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes g\|_{\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}},$
 v) If $\Phi \leq_{cp} \sum_k \beta_k \Psi_k$, then

$$\|\xi\|_{\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}} \leq \left(\sum_k \beta_k \|\xi\|_{\mathcal{M} \bar{\otimes}_{\Psi_k} \mathcal{M}}^2 \right)^{\frac{1}{2}}.$$

Proof. The first inequality follows from hermitianity of the inner product and the identity $\langle \xi, \eta \rangle_{\Phi} = \rho(\xi)^* \rho(\eta)$ explained above. The second one is straightforward from the definition of $\mathcal{M} \bar{\otimes}_{\Phi} \mathcal{M}$. The third inequality follows from Kadison-Schwarz inequality after expanding both sides. The lower estimate in iv) holds trivially with constant 1, while the upper estimate holds with constant 2 since

$$f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes \Phi f = (f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes g) - (\mathbf{1}_{\mathcal{M}} \otimes (\Phi f - g))$$

and the second term on the right hand side is estimated using iii). Finally, for the last inequality let $\xi = \sum_k a_k \otimes b_k$ and define the column matrices $A^* = \sum_k a_k^* \otimes e_{k1}$ and $B = \sum_k b_k \otimes e_{k1}$. Then we find

$$\begin{aligned} \langle \xi, \xi \rangle_{\Phi} &= \sum_{j,k} b_j^* \Phi(a_j^* a_k) b_k = B^* \Phi(A^* A) B \\ &\leq \sum_k \beta_k B^* \Psi_k(A^* A) B = \sum_k \beta_k \langle \xi, \xi \rangle_{\Psi_k}. \quad \square \end{aligned}$$

1.2. Markov semigroups and metrics

Throughout the article we will assume that the underlying noncommutative measure space (\mathcal{M}, τ) comes equipped with a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$. That is, a weak-* continuous semigroup of normal cpu (*completely positive unital*) operators $(S_t)_{t \geq 0}$ on \mathcal{M} with the additional symmetric assumption

$$\tau(S_t(f)g) = \tau(fS_t(g)) \quad \text{for all } f, g \in \mathcal{M} \cap L_1(\mathcal{M}).$$

Note that $\tau(S_t f) = \tau(f)$ for all f , so S_t 's are faithful and are contractive on $L_1(\mathcal{M})$. By interpolation, S_t 's extend to contractions on $L_p(\mathcal{M})$ ($1 \leq p < \infty$) and satisfy $\lim_{t \rightarrow 0} S_t f = f$ in $L_p(\mathcal{M})$ for all $f \in L_p(\mathcal{M})$.

Also recall that \mathcal{S} admits an infinitesimal negative generator

$$Af = \lim_{t \rightarrow 0} t^{-1}(f - S_t(f))$$

defined on $\text{dom}(A) = \bigcup_{1 \leq p \leq \infty} \text{dom}_p(A)$, where $\text{dom}_p(A)$ is given by

$$\text{dom}_p(A) = \left\{ f \in L_p(\mathcal{M}) : \lim_{t \rightarrow 0} t^{-1}(S_t(f) - f) \text{ converges in } L_p(\mathcal{M}) \right\}.$$

It is easy to see that $\frac{1}{s} \int_0^s S_t(f) dt \in \text{dom}_p(A)$ for any $s > 0, f \in L_p(\mathcal{M})$, so $\text{dom}_p(A)$ is dense in $L_p(\mathcal{M})$ for $p < \infty$, and is weak-* dense in \mathcal{M} for $p = \infty$. We will denote by A_p

the restriction of A on $\text{dom}_p(A)$.

A Markov metric associated to (\mathcal{M}, τ) and \mathcal{S} is determined by a family

$$\mathcal{Q} = \left\{ (R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+ \right\}$$

where $R_{j,t} : \mathcal{M} \rightarrow \mathcal{M}$ are normal completely positive unital maps satisfying $R_{j,t}f = f$ for $f \in \ker A_\infty$, the fixed-point subspace of the semigroup \mathcal{S} , and $\sigma_{j,t}, \gamma_{j,t}$ are elements of the von Neumann algebra \mathcal{M} with $\gamma_{j,t} \geq \mathbf{1}_{\mathcal{M}}$, so that the estimates below hold:

- i) *Hilbert module majorization*: $\langle \xi, \xi \rangle_{S_t} \leq \sum_{j \geq 1} \sigma_{j,t}^* \langle \xi, \xi \rangle_{R_{j,t}} \sigma_{j,t}, \forall \xi \in \mathcal{M} \otimes \mathcal{M}$,
- ii) *Metric integrability condition*: $k_{\mathcal{Q}} = \sup_{t > 0} \left\| \sum_{j \geq 1} \sigma_{j,t}^* \gamma_{j,t}^2 \sigma_{j,t} \right\|_{\mathcal{M}}^{\frac{1}{2}} < \infty$.

Our notion of Markov metric is easily understood for our (commutative) model case above. Let $\mathcal{S} = (S_t)_{t \geq 0}$ be a Markov semigroup on (Ω, μ) with associated kernels $s_t(x, y)$ satisfying the pointwise estimate (1.1). Given $\xi : \Omega \times \Omega \rightarrow \mathbb{C}$ essentially bounded, we have

$$\langle \xi, \xi \rangle_{S_t} = \int_{\Omega} s_t(x, y) |\xi(x, y)|^2 d\mu(y) \leq \sum_{j=1}^{\infty} \frac{|\sigma_{j,t}(x)|^2}{\mu(\Sigma_{j,t}(x))} \int_{\Sigma_{j,t}(x)} |\xi(x, y)|^2 d\mu(y). \quad (1.2)$$

This means that $R_{j,t}f(x)$ is the average of $f \in L_\infty(\Omega)$ over the set $\Sigma_{j,t}(x)$. Reciprocally, if we take $\xi_k(x, y) = \phi_k(y_0 - y)$ to be an approximation of identity around y_0 , we recover the pointwise estimates for the kernel $s_t(x, y_0)$. In other more general contexts, the upper bounds for the kernel or even the kernel description of the semigroup might not have the same form. As we shall see, many of these cases can still be handled via Hilbert module majorization. We shall provide along the paper a wide variety of examples which fall into these possible classes.

1.3. Semigroup BMOs

Given a noncommutative measure space (\mathcal{M}, τ) and a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ acting on (\mathcal{M}, τ) , we may define the *semigroup BMO $_{\mathcal{S}}$ -norm* as

$$\|f\|_{\text{BMO}_{\mathcal{S}}} = \max \left\{ \|f\|_{\text{BMO}_{\mathcal{S}}^r}, \|f\|_{\text{BMO}_{\mathcal{S}}^c} \right\},$$

where the row and column BMO norms are given by

$$\begin{aligned} \|f\|_{\text{BMO}_{\mathcal{S}}^r} &= \sup_{t \geq 0} \left\| \left(S_t(f f^*) - (S_t f)(S_t f)^* \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}, \\ \|f\|_{\text{BMO}_{\mathcal{S}}^c} &= \sup_{t \geq 0} \left\| \left(S_t(f^* f) - (S_t f)^*(S_t f) \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}. \end{aligned}$$

This definition makes sense since we know from the Kadison-Schwarz inequality that $|S_t f|^2 \leq S_t |f|^2$. The null space of this seminorm is $\ker A_\infty$, the fixed-point subspace of our semigroup. Indeed, if $\|f\|_{\text{BMO}_S} = 0$ we know from [11] that f belongs to the multiplicative domain of S_t , so that

$$\tau(gf) = \tau(S_{t/2}(gf)) = \tau(S_{t/2}(g)S_{t/2}(f)) = \tau(gS_t(f)).$$

This proves that f is fixed by the semigroup. Reciprocally, $\ker A_\infty$ is a $*$ -subalgebra of \mathcal{M} by [27,35]. [27] only deals with the state case, but the proof there carries over to the semifinite trace case since we have the additional symmetric assumption on S_t . Thus, the seminorm vanishes on $\ker A_\infty$. In particular, we obtain a norm after quotienting out $\ker A_\infty$. Letting $w_t(f) = f \otimes \mathbf{1} - \mathbf{1} \otimes S_t f$, this provides us with a map

$$f \in \mathcal{M} \xrightarrow{w} (w_t(f))_{t \geq 0} \in \bigoplus_{t \geq 0} \mathcal{M} \bar{\otimes}_{S_t} \mathcal{M}$$

which becomes isometric when we equip \mathcal{M} with the norm in BMO_S^c . Given a net $(f_\lambda)_\lambda \in \mathcal{M}$ with uniform bounded BMO_S^c -norm (i.e. $\sup_\lambda \|f_\lambda\|_{\text{BMO}_S^c} < \infty$), we say $(f_\lambda)_\lambda$ \mathcal{S} -weak- $*$ converges if $(w(f_\lambda))_\lambda$ weak- $*$ -converges in $\bigoplus_{t \geq 0} \mathcal{M} \bar{\otimes}_{S_t} \mathcal{M}$. That is to say, for any $t, m > 0$, $a_j, b_j \in L^2(\mathcal{M})$, and $x = \sum_{j=1}^m a_j \otimes b_j \in L_1(\mathcal{M} \bar{\otimes}_{S_t} \mathcal{M})$, the following equivalent inner products

$$\begin{aligned} \tau(\langle w_t(f_\lambda), x \rangle_{S_t}) &= \tau\left(\langle f_\lambda \otimes \mathbf{1} - \mathbf{1} \otimes S_t f_\lambda, \sum_{j=1}^m a_j \otimes b_j \rangle_{S_t}\right) \\ &= \tau\left(f_\lambda^* \left(\sum_{j=1}^m a_j S_t(b_j) - S_t(S_t(a_j)b_j)\right)\right) \end{aligned} \quad (1.3)$$

converges. This allows us to define the atomic h_1^c space. Let Atom_c be the collection of all ξ 's of the form

$$\xi = \sum_{j=1}^m a_j S_t(b_j) - S_t(S_t(a_j)b_j)$$

for some $t, m > 0$, $a_j, b_j \in L_2(\mathcal{M})$ and

$$\left\| \sum_{j=1}^m a_j \otimes b_j \right\|_{L_1(\mathcal{M} \bar{\otimes}_{S_t} \mathcal{M})} = 1.$$

For any $g \in \text{Span}(\text{Atom}_c)$, let

$$\|g\|_{h_{1,at}^c} = \inf \left\{ \sum_k |c_k| : g = \sum_{k=1}^n c_k \xi_k, \xi_k \in \text{Atom}_c \right\}$$

It is easy to see that for any $g \in \text{Span}(\text{Atom}_c)$, $\|g\|_{h_{1,at}^c} \geq \|g\|_{L_1(\mathcal{M})}$. We define $h_{1,at}^c(\mathcal{S})$ to be the $\|\cdot\|_{h_{1,at}^c}$ -closure of $\text{Span}(\text{Atom}_c)$, which is contained in $L_1(\mathcal{M})$. Let $L_\infty^\circ(\mathcal{M}) = L_\infty(\mathcal{M})/\ker A_\infty$. By (1.3), it is clear that

$$(L_\infty^\circ(\mathcal{M}), \|\cdot\|_{\text{BMO}_\mathcal{S}^\circ}) \subset (h_{1,at}^c(\mathcal{S}))^*$$

and a $\|\cdot\|_{\text{BMO}_\mathcal{S}^\circ}$ -bounded net \mathcal{S} -weak-* converges iff it weak-* converges in $(h_{1,at}^c(\mathcal{S}))^*$. It is also clear that $(L_\infty^\circ(\mathcal{M}), \|\cdot\|_{\text{BMO}_\mathcal{S}^\circ})$ separates points in the whole space $h_{1,at}^c(\mathcal{S})$, so it is weak-* dense in $(h_{1,at}^c(\mathcal{S}))^*$. Recall two convergent nets are equivalent if their differences converge to 0. Define $\text{BMO}_\mathcal{S}^\circ$ as the collection of all the equivalence classes of \mathcal{S} -weak-* convergent nets $(f_\lambda)_\lambda$ with $\sup_\lambda \|f_\lambda\|_{\text{BMO}_\mathcal{S}^\circ} < \infty$. So $\text{BMO}_\mathcal{S}^\circ$ is the weak-* closure of $(L_\infty^\circ(\mathcal{M}), \|\cdot\|_{\text{BMO}_\mathcal{S}^\circ})$ in $(h_{1,at}^c(\mathcal{S}))^*$, thus equals $(h_{1,at}^c(\mathcal{S}))^*$.

Similarly, we may define $\text{BMO}_\mathcal{S}$ as the intersection $\text{BMO}_\mathcal{S}^r \cap \text{BMO}_\mathcal{S}^\circ$, where the row BMO follows by taking adjoints above. The natural operator space structure is given by

$$M_m(\text{BMO}_\mathcal{S}(\mathcal{M})) = \text{BMO}_{\widehat{\mathcal{S}}}(M_m(\mathcal{M})) \quad \text{with} \quad \widehat{S}_t = id_{M_m} \otimes S_t.$$

Remark 1.2. Incidentally, we note that $\text{BMO}_\mathcal{S}$ is written as $bmo(\mathcal{S})$ in [30].

1.4. Markov regularity and interpolation

It will be essential for us to provide interpolation results between semigroup type BMO spaces and the corresponding noncommutative L_p spaces [53]. It is a hard problem to identify the minimal regularity on the semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ which suffices for this purpose. We will recall in this subsection an interpolation theory from [30].

We say that a semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ admits a *reversed Markov dilation* if there exists a larger semifinite von Neumann algebra $(\mathcal{N}, \tilde{\tau})$, a decreasing filtration $(\mathcal{N}_{[s]})_{s \geq 0}$ with corresponding conditional expectations $E_{[s]}$, and a family of trace preserving *-homomorphisms $\pi_s : \mathcal{M} \rightarrow \mathcal{N}_{[s]}$ such that

$$E_{[s]}(\pi_t(f)) = \pi_s(S_{s-t}f), \quad \text{for all } t < s, x \in \mathcal{M}.$$

This implies that $m(f) = (m_s(f))_{s \geq 0}$ with

$$m_s(f) = \pi_s(S_s f),$$

is a martingale with respect to the reversed filtration $(\mathcal{N}_{[s]})$. A Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ is called *regular* if it admits a reversed Markov dilation and for each $2 \leq p < \infty$, there exists a weakly dense subset $B_p \subset L_p(\mathcal{N})$ so that, for every $f \in B_p$ the associated martingale $m(f) = (m_s(f))_{s \geq 0}$ has almost uniformly continuous path. In other words, for every $T > 0$ and $\varepsilon > 0$ there exists a projection e such that $\tilde{\tau}(1 - e) < \varepsilon$ and $f_e : [0, T] \ni t \mapsto m_t(f)e \in \mathcal{N}$ is continuous.

Define

$$L_p^\circ(\mathcal{M}) = L_p(\mathcal{M})/\ker A_p = \left\{ f \in L_p(\mathcal{M}) : \lim_{t \rightarrow \infty} S_t f = 0 \right\}.$$

As it was explained in [30], the space $L_p^\circ(\mathcal{M})$ is complemented in $L_p(\mathcal{M})$, $1 \leq p \leq \infty$, and $[L_1^\circ(\mathcal{M}), \text{BMO}_\mathcal{S}]$ form an interpolation couple. The following result will be crucial in what follows, we refer the reader to [30] for a detailed proof.

Theorem 1.3. *If $\mathcal{S} = (S_t)_{t \geq 0}$ is a regular Markov semigroup*

$$[\text{BMO}_\mathcal{S}, L_p^\circ(\mathcal{M})]_{p/q} \simeq_{cb} L_q^\circ(\mathcal{M}) \quad \text{for all } 1 \leq p < q < \infty.$$

Note that interpolation against the full space $L_p(\mathcal{M})$ is meaningless since $\text{BMO}_\mathcal{S}$ does not distinguish the fixed-point space of the semigroup. Very roughly, we shall typically apply the above result to a CZO which is bounded on $L_2(\mathcal{M})$ and sends a weak-* dense subalgebra $\mathcal{A} \subset \mathcal{M}$ to $\text{BMO}_\mathcal{S}$, and \mathcal{A} has dense intersection with $L_p(\mathcal{M})$ for any $1 \leq p < \infty$. Recalling the projection map $J_p : L_p(\mathcal{M}) \rightarrow L_p^\circ(\mathcal{M})$ and letting T denote the CZO, we find by interpolation that

$$J_p T : L_p(\mathcal{M}) \cap \mathcal{A} \subset [\mathcal{A}, L_2(\mathcal{M})]_{2/p} \rightarrow [\text{BMO}_\mathcal{S}, L_2^\circ(\mathcal{M})]_{2/p} = L_p^\circ(\mathcal{M}) \subset L_p(\mathcal{M}).$$

Then the density of $L_p(\mathcal{M}) \cap \mathcal{A}$ give us the boundedness of $J_p T$ from $L_p(\mathcal{M})$ to $L_p^\circ(\mathcal{M})$. To obtain L_p boundedness of T , it suffices to assume that T leaves the fixed-point space invariant and is bounded on it. It should be noticed though, that in many cases the L_p boundedness of the CZO follows automatically. For instance, in \mathbb{R}^n with the Lebesgue measure and the heat semigroup, it turns out that $L_p = L_p^\circ$. On the other hand, the fixed-point space for the Poisson semigroup on the n -torus is just composed of constant functions and the corresponding projection can be estimated apart regarded as a conditional expectation. Moreover, the same applies for Fourier multipliers on arbitrary discrete groups. The L_p boundedness for $1 < p < 2$ will follow by taking adjoints under certain symmetry on the hypotheses.

Remark 1.4. On commutative or semicommutative von Neumann algebras, the existence of a Markov dilation is called the martingale problem and a Markov semigroup is regular if it is diffusion. Noncommutative diffusion Markov semigroups and their regularity were recently studied in [33] for the case of finite von Neumann algebras. At the moment, there is no written argument on the existence of Markov dilations on semifinite von Neumann algebras. However, we should recall that all the concrete examples in this paper are regular. Our commutative models are the heat semigroup, the Ornstein-Uhlenbeck semigroup and the Laplace-Beltrami semigroup and all of them satisfy our regularity assumption. Regular Markov dilations in our noncommutative examples (matrix algebras and quantum Euclidean spaces) can be deduced by transference. Indeed, we shall embed

these algebras into larger semicommutative ones and the Markov process arises from an amplification of the classical heat semigroup.

Remark 1.5. Let us comment on the case of general semifinite von Neumann algebras, where the results of [33] do not formally apply. Given a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0} = (e^{-tA})_{t \geq 0}$, we may first fix a finite projection e , a small $s > 0$ and consider the new generator

$$A_{s,e} = \frac{id_{e\mathcal{M}e} - S_s^e}{s}$$

with $S_s^e(f) = eS_s(efe)e$ being selfadjoint. Then $A_{s,e}$ generates a regular Markov semigroup on $e\mathcal{M}e$. So the interpolation between the semigroup BMOs and L_p and the corresponding semigroup H_p - L_p equivalence which appear in [30] hold for $e^{-tA_{s,e}}$. We expect that an approximation argument is enough to prove the same equivalence for the original semigroup S_t for $2 < p < \infty$, and that a BMO interpolation theory would hold for general Markov semigroups on semifinite von Neumann algebras.

1.5. Markov metric $\text{BMO}_{\mathcal{Q}}$

Let us now introduce a Markov metric type BMO space for von Neumann algebras and relate it with the semigroup type BMO spaces defined above. Given a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ acting on (\mathcal{M}, τ) , consider a Markov metric $\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+\}$ as defined above and define $\|f\|_{\text{BMO}_{\mathcal{Q}}} = \max\{\|f\|_{\text{BMO}_{\mathcal{Q}}^c}, \|f^*\|_{\text{BMO}_{\mathcal{Q}}^c}\}$, where the column BMO-norm is given by

$$\sup_{t>0} \inf_{M_t \text{ cpu}} \sup_{j \geq 1} \left\| \left(\gamma_{j,t}^{-1} [R_{j,t}|f|^2 - |R_{j,t}f|^2 + |R_{j,t}f - M_t f|^2] \gamma_{j,t}^{-1} \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}$$

and the infimum runs over normal cpu maps $M_t : \mathcal{M} \rightarrow \mathcal{M}$ satisfying $M_t(f) = f$ for $f \in \ker A_{\infty}$. The row norm is estimated in the same way. Define, for $g \in L_1(\mathcal{M})$,

$$\|g\|_{h_{1,\mathcal{Q}}^c} = \sup\{|\tau(f^*g)| : f \in L_{\infty}^{\circ}(\mathcal{M}), \|f\|_{\text{BMO}_{\mathcal{Q}}^c} \leq 1\}. \quad (1.4)$$

We will see in the theorem below that, $\|\cdot\|_{\text{BMO}_{\mathcal{S}}^c} \lesssim \|\cdot\|_{\text{BMO}_{\mathcal{Q}}^c}$ on $L_{\infty}^{\circ}(\mathcal{M})$. So $\|f\|_{h_{1,\mathcal{Q}}^c} < \infty$ for any $f \in h_{1,at}^c(\mathcal{S})$ (see its definition in Section 1.3). Let $h_{1,\mathcal{Q}}^c$ be the $\|\cdot\|_{h_{1,\mathcal{Q}}^c}$ -closure of $(h_{1,at}^c(\mathcal{S}), \|\cdot\|_{h_{1,\mathcal{Q}}^c})$, which is contained in $L_1(\mathcal{M})$. It is clear that $(L_{\infty}^{\circ}(\mathcal{M}), \|\cdot\|_{\text{BMO}_{\mathcal{Q}}^c}) \subset (h_{1,\mathcal{Q}}^c)^*$. Let $\text{BMO}_{\mathcal{Q}}^c$ be its weak-* closure in $(h_{1,\mathcal{Q}}^c)^*$. That is the linear space of all equivalence classes of $\|\cdot\|_{\text{BMO}_{\mathcal{Q}}^c}$ -bounded, weak-* convergent nets $(f_{\lambda})_{\lambda} \in L_{\infty}^{\circ}(\mathcal{M})$ in $(h_{1,\mathcal{Q}}^c)^*$. For such a net $(f_{\lambda})_{\lambda}$, define the norm of its equivalence class as

$$\|[(f_{\lambda})_{\lambda}]\|_{\text{BMO}_{\mathcal{Q}}^c} = \sup_{\|g\|_{h_{1,\mathcal{Q}}^c} \leq 1} \lim_{\lambda} |\tau(f_{\lambda}^*g)|.$$

We have $\text{BMO}_{\mathcal{Q}}^c$ is equal to the whole space $(h_{1,\mathcal{Q}}^c)^*$, since it separates points in $h_{1,\mathcal{Q}}^c$. Similarly, we may define $\text{BMO}_{\mathcal{Q}}$ as the intersection $\text{BMO}_{\mathcal{Q}}^r \cap \text{BMO}_{\mathcal{Q}}^c$, where the row BMO follows by taking adjoints above. Since $\gamma_{j,t} \geq \mathbf{1}_{\mathcal{M}}$, the inverses exist and $L_{\infty}(\mathcal{M})$ embeds in $\text{BMO}_{\mathcal{Q}}$. Indeed, using that $R_{j,t}$ and M_t are cpu, the square bracket above is bounded by $2\|f\|_{\infty}^2 \mathbf{1}_{\mathcal{M}}$ and $\gamma_{j,t}^{-2} \leq \mathbf{1}_{\mathcal{M}}$.

Now, recalling the value of the constant $k_{\mathcal{Q}}$ in our definition of Markov metric, we prove that $\text{BMO}_{\mathcal{Q}}$ embeds in $\text{BMO}_{\mathcal{S}}$.

Theorem 1.6. *Let (\mathcal{M}, τ) be a noncommutative measure space equipped with a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$. Let us consider a Markov metric \mathcal{Q} associated to $\mathcal{S} = (S_t)_{t \geq 0}$. Then, we find*

$$\|f\|_{\text{BMO}_{\mathcal{S}}} \lesssim k_{\mathcal{Q}} \|f\|_{\text{BMO}_{\mathcal{Q}}}.$$

In particular, we see that $L_{\infty}(\mathcal{M}) \subset \text{BMO}_{\mathcal{Q}} \subset \text{BMO}_{\mathcal{S}}$ and

$$[\text{BMO}_{\mathcal{Q}}, L_p^{\circ}(\mathcal{M})]_{p/q} \simeq L_q^{\circ}(\mathcal{M}) \quad \text{for all } 1 \leq p < q < \infty$$

for any Markov metric \mathcal{Q} associated to a regular Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$.

Proof. Let us set

$$\begin{aligned} \xi_t &= f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes M_t f \\ &= (f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes R_{j,t} f) + (\mathbf{1}_{\mathcal{M}} \otimes (R_{j,t} f - M_t f)) = \xi_{j,t}^1 + \xi_{j,t}^2. \end{aligned}$$

The assertion follows from Lemma 1.1 and our definition of Markov metric

$$\begin{aligned} \|f\|_{\text{BMO}_{\mathcal{S}}} &= \sup_{t>0} \left\| \left(S_t |f|^2 - |S_t f|^2 \right)^{\frac{1}{2}} \right\|_{\mathcal{M}} \\ &= \sup_{t>0} \|f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes S_t f\|_{\mathcal{M} \bar{\otimes}_{S_t} \mathcal{M}} \\ &\lesssim \sup_{t>0} \|f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes M_t f\|_{\mathcal{M} \bar{\otimes}_{S_t} \mathcal{M}} = \sup_{t>0} \|\langle \xi_t, \xi_t \rangle_{S_t}\|_{\mathcal{M}}^{\frac{1}{2}} \\ &\lesssim \sup_{t>0} \left\| \sum_{j \geq 1} \sigma_{j,t}^* [\langle \xi_{j,t}^1, \xi_{j,t}^1 \rangle_{R_{j,t}} + \langle \xi_{j,t}^2, \xi_{j,t}^2 \rangle_{R_{j,t}}] \sigma_{j,t} \right\|_{\mathcal{M}}^{\frac{1}{2}} \\ &\leq k_{\mathcal{Q}} \|f\|_{\text{BMO}_{\mathcal{Q}}}. \end{aligned}$$

The identities are clear. The first inequality follow from Lemma 1.1 iv), the second one from the Hilbert module majorization associated to the Markov metric and Lemma 1.1 i). To justify the last inequality, note that the square bracket inside the term on the left equals $R_{j,t} |f|^2 - |R_{j,t} f|^2 + |R_{j,t} f - M_t f|^2$. Hence, left multiplication by $\gamma_{j,t} \gamma_{j,t}^{-1}$ and right multiplication by $\gamma_{j,t}^{-1} \gamma_{j,t}$ yields the given inequality with $k_{\mathcal{Q}}$ the metric integrability

constant. From the above inequalities, we see that the null space of the seminorm $\|\cdot\|_{\text{BMO}_{\mathcal{Q}}}$ on \mathcal{M} is again $\ker A_{\infty}$. Combining this with (1.3), we see that $h_{1,\mathcal{Q}}^c$ contains $h_{1,at}^c$ as a dense subset. So $\text{BMO}_{\mathcal{Q}} \subset \text{BMO}_{\mathcal{S}}$. The interpolation result follows from Theorem 1.3 and the embeddings $L_{\infty}(\mathcal{M}) \subset \text{BMO}_{\mathcal{Q}} \subset \text{BMO}_{\mathcal{S}}$. The proof is complete. \square

Remark 1.7. Let $\xi = \sum_j A_j \otimes B_j$, with

$$A_j = \begin{pmatrix} a_{\alpha\beta}^j \end{pmatrix} \quad \text{and} \quad B_j = \begin{pmatrix} b_{\alpha\beta}^j \end{pmatrix}$$

elements of $M_m(\mathcal{M})$. If $\widehat{S}_t = id_{M_m} \otimes S_t$, it turns out that

$$\langle \xi, \xi \rangle_{\widehat{S}_t} = \left(\sum_{\alpha=1}^m \left\langle \underbrace{\sum_{j,\beta} a_{\alpha\beta}^j \otimes b_{\beta\gamma_1}^j}_{\eta_{\alpha,\gamma_1}}, \underbrace{\sum_{j,\beta} a_{\alpha\beta}^j \otimes b_{\beta\gamma_2}^j}_{\eta_{\alpha,\gamma_2}} \right\rangle_{S_t} \right)_{\gamma_1, \gamma_2} \in M_m(\mathcal{M}).$$

This can be used to provide an operator space structure on $\text{BMO}_{\mathcal{Q}}$. Namely, the canonical choice for the matrix norms is $M_m(\text{BMO}_{\mathcal{Q}}(\mathcal{M})) = \text{BMO}_{\widehat{\mathcal{Q}}}(M_m(\mathcal{M}))$ where the Markov metric on $M_m(\mathcal{M})$

$$\widehat{\mathcal{Q}} = \left\{ (id_{M_m} \otimes R_{j,t}, \mathbf{1}_{M_m} \otimes \sigma_{j,t}, \mathbf{1}_{M_m} \otimes \gamma_{j,t}) \right\}$$

is associated to the extended semigroup $(\widehat{S}_t)_{t \geq 0}$. Then, we trivially obtain that $k_{\widehat{\mathcal{Q}}} = k_{\mathcal{Q}} < \infty$. However, according to the identity above for $\langle \xi, \xi \rangle_{\widehat{S}_t}$, the Hilbert module majorization takes the form

$$\left(\sum_{\alpha=1}^m \langle \eta_{\alpha,\gamma_1}, \eta_{\alpha,\gamma_2} \rangle_{S_t} \right)_{\gamma_1, \gamma_2} \leq \sum_{j \geq 1} \left(\sigma_{j,t}^* \sum_{\alpha=1}^m \langle \eta_{\alpha,\gamma_1}, \eta_{\alpha,\gamma_2} \rangle_{R_{j,t}} \sigma_{j,t} \right)_{\gamma_1, \gamma_2}.$$

This gives a matrix-valued generalization of our Hilbert module majorization for $\mathcal{S} = (S_t)_{t \geq 0}$ on \mathcal{M} , to be checked when we use this o.s.s. Theorem 1.6 yields a cb-embedding of $\text{BMO}_{\mathcal{Q}}$ into $\text{BMO}_{\mathcal{S}}$ under this assumption. According to the characterization (1.2), it holds for Markov metrics on commutative spaces (Ω, μ) .

1.6. The Euclidean metric

Before using Markov metrics in our approach to Calderón-Zygmund theory, it is illustrative to recover the Euclidean metric from a suitable Markov semigroup. Let $\mathcal{S} = (H_t)_{t \geq 0}$ denote the (regular) heat semigroup on \mathbb{R}^n , with kernels

$$h_t(x, y) = \frac{1}{(4\pi t)^{\frac{n}{2}}} \exp\left(\frac{-|x - y|^2}{4t}\right).$$

Take $\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+\}$ determined by

- $\sigma_{j,t}^2 \equiv \frac{2e}{\sqrt{\pi}} j^{\frac{n}{2}} e^{-j}$ and $\gamma_{j,t}^2 \equiv j^{\frac{n}{2}} \geq 1$,
- $R_{j,t}f(x) = \frac{1}{|B_{\sqrt{4jt}}(x)|} \int_{B_{\sqrt{4jt}}(x)} f(y) dy$.

Note that $\sigma_{j,t}$ and $\gamma_{j,t}$ are allowed to be essentially bounded functions in \mathbb{R}^n , but in this case it suffices to take constant functions. In the definition of $R_{j,t}$, we write $B_r(x)$ to denote the Euclidean ball in \mathbb{R}^n centered at x with radius r . It is clear that $R_{j,t}$ defines a cpu map on $L_\infty(\mathbb{R}^n)$ and $R_{j,t}f = f$ for any $f \in \ker A_\infty$, since $\ker A_\infty$ is the space of constant functions in this case. To show that \mathcal{Q} defines a Markov metric, we need to check that it provides a Hilbert module majorization of the heat semigroup and the metric integrability condition holds. The latter is straightforward, while the Hilbert module majorization reduces to check that

$$h_t(x, y) \leq \frac{2e}{\sqrt{\pi}} \sum_{j \geq 1} \frac{j^{\frac{n}{2}} e^{-j}}{|B_{\sqrt{4jt}}(x)|} \chi_{B_{\sqrt{4jt}}(x)}(y).$$

This can be justified by determining the unique corona centered at x with radii $\sqrt{4(j-1)t}$ and $\sqrt{4jt}$ where y lives, details are left to the reader. Note that we could have taken $\gamma_{j,t} \equiv 1$ and still obtain a Markov metric. Our choice will be justified below and also in the next section, where we shall need $\gamma_{j,t} \equiv j^{\frac{n}{2}}$ to compare $\text{BMO}_{\mathcal{Q}}$ with other BMO spaces which interpolate. Before that, our only evidences that this is the right Markov metric in the Euclidean case are the fact that the $R_{j,t}$'s are averages over Euclidean balls and the isomorphism

$$\text{BMO}_{\mathcal{Q}} = \text{BMO}_{\mathbb{R}^n},$$

where the latter space is the usual BMO space in \mathbb{R}^n

$$\|f\|_{\text{BMO}_{\mathbb{R}^n}} = \sup_{B \subset \mathbb{R}^n} \left(\frac{1}{|B|} \int_B |f(x) - f_B|^2 dx \right)^{\frac{1}{2}}.$$

Here, the supremum is taken over all Euclidean balls B in \mathbb{R}^n and f_B stands for the average of f over B . Let us justify this isomorphism. If we pick $M_t f(x) = R_{1,t} f(x)$ it follows from a standard calculation that

$$\begin{aligned} & |R_{j,t} f(x) - M_t f(x)|^2 \\ &= \left| \frac{1}{|B_{\sqrt{4t}}(x)|} \int_{B_{\sqrt{4t}}(x)} (f(y) - f_{B_{\sqrt{4t}}(x)}) dy \right|^2 \end{aligned} \tag{1.5}$$

$$\leq \frac{j^{\frac{n}{2}}}{|B_{\sqrt{4jt}}(x)|} \int_{B_{\sqrt{4jt}}(x)} |f(y) - f_{B_{\sqrt{4jt}}(x)}|^2 dy = j^{\frac{n}{2}} (R_{j,t}|f|^2 - |R_{j,t}f|^2)(x).$$

This automatically yields the following inequality

$$\|f\|_{\text{BMO}_{\mathcal{Q}}}^2 \lesssim \sup_{j,t} \text{ess sup}_{x \in \mathbb{R}^n} \frac{1}{|B_{\sqrt{4jt}}(x)|} \int_{B_{\sqrt{4jt}}(x)} |f(y) - f_{B_{\sqrt{4jt}}(x)}|^2 dy \leq \|f\|_{\text{BMO}_{\mathbb{R}^n}}^2.$$

The converse is even simpler, since taking $j = 1$ we obtain

$$\begin{aligned} \|f\|_{\text{BMO}_{\mathbb{R}^n}}^2 &= \sup_{t>0} \text{ess sup}_{x \in \mathbb{R}^n} \frac{1}{|B_{\sqrt{4t}}(x)|} \int_{B_{\sqrt{4t}}(x)} |f(y) - f_{B_{\sqrt{4t}}(x)}|^2 dy \\ &= \sup_{t>0} \left\| \gamma_{1,t}^{-1} [R_{1,t}|f|^2 - |R_{1,t}f|^2] \gamma_{1,t}^{-1} \right\|_{L_{\infty}(\mathbb{R}^n)} \leq \|f\|_{\text{BMO}_{\mathcal{Q}}}^2. \end{aligned}$$

Remark 1.8. The term $|R_{j,t}f - M_tf|$ did not play a significant role at this point. More generally, the above argument also works for any doubling metric space Ω equipped with a Borel measure μ : $\mu(B(x, 2r)) \leq C\mu(B(x, r))$ for every $x \in \Omega$ and $r > 0$, with $B(x, r) = \{y \in \Omega : \text{dist}(x, y) \leq r\}$. As we shall see later, the additional term $|R_{j,t}f - M_tf|$ in the $\text{BMO}_{\mathcal{Q}}$ -norm appears to include Tolsa's RBMO spaces [61] in those measure spaces (Ω, μ) for which we can find an appropriate Dirichlet form which provides us with a Markov semigroup acting on (Ω, μ) .

Remark 1.9. A related semigroup BMO norm is

$$\|f\|_{\text{BMO}_{\mathcal{S}}} = \sup_{t \geq 0} \left\| \left(H_t[|f - H_tf|^2] \right)^{\frac{1}{2}} \right\|_{\infty}.$$

All the norms consider so far are equivalent for the heat semigroup $\mathcal{S} = (H_t)_{t \geq 0}$ on \mathbb{R}^n , generated by the Laplacian $\Delta = \sum_{j=1}^n \partial_{x_j}^2$. In fact, we may also consider by subordination the Poisson semigroup $\mathcal{P} = (P_t)_{t \geq 0}$ on \mathbb{R}^n generated by the square root $\sqrt{-\Delta}$, or even other subordinations [20]. Then, elementary calculations give the following norm equivalences up to dimensional constants

$$\|f\|_{\text{BMO}_{\mathbb{R}^n}} \sim \|f\|_{\text{BMO}_{\mathcal{P}}} \sim \|f\|_{\text{BMO}_{\mathcal{P}}} \sim \|f\|_{\text{BMO}_{\mathcal{S}}} \sim \|f\|_{\text{BMO}_{\mathcal{S}}} \sim \|f\|_{\text{BMO}_{\mathcal{Q}}}.$$

Moreover, let $\mathcal{R} = L_{\infty}(\mathbb{R}^n) \bar{\otimes} \mathcal{M}$ denote the von Neumann algebra tensor product of $L_{\infty}(\mathbb{R}^n)$ with a noncommutative measure space (\mathcal{M}, τ) . Define the norm in $\text{BMO}_{\mathcal{R}}$ as $\|f\|_{\text{BMO}_{\mathcal{R}}} = \max\{\|f\|_{\text{BMO}_{\mathcal{R}}^c}, \|f^*\|_{\text{BMO}_{\mathcal{R}}^c}\}$, where

$$\|f\|_{\text{BMO}_{\mathcal{R}}^c} = \sup_{B \text{ balls}} \left\| \left(\frac{1}{|B|} \int_B |f(x) - f_B|^2 dx \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}.$$

Then, the same norm equivalences hold in the semicommutative case

$$\|f\|_{\text{BMO}_{\mathcal{R}}} \sim \|f\|_{\text{BMO}_{\mathcal{P}_{\otimes}}} \sim \|f\|_{\text{BMO}_{\mathcal{P}_{\otimes}}} \sim \|f\|_{\text{BMO}_{\mathcal{S}_{\otimes}}} \sim \|f\|_{\text{BMO}_{\mathcal{S}_{\otimes}}},$$

where $S_{\otimes,t} = S_t \otimes id_{\mathcal{M}}$ and $P_{\otimes,t} = P_t \otimes id_{\mathcal{M}}$. Moreover, by Remark 1.7, all these norms are in turn equivalent to the norm in $\text{BMO}_{\mathcal{Q}_{\mathcal{R}}}$, with the Markov metric which arises tensorizing the canonical one with the identity/unit of \mathcal{M} .

2. Algebraic CZ theory

In classical Calderón-Zygmund theory, L_p boundedness of CZOs follows from L_2 boundedness under a smoothness condition on the kernel. Our next goal is to identify which are the analogues of these conditions for semifinite von Neumann algebras equipped with a Markov metric, and to show L_p boundedness of CZOs fulfilling them. Our new conditions are certainly surprising. The boundedness for $p = 2$ must be replaced by a certain mixed-norm estimate (which reduces in the classical theory to L_2 boundedness), while Hörmander kernel smoothness will be formulated intrinsically without any reference to the kernel. These abstract assumptions will adopt a more familiar form in the specific cases that we shall consider in the forthcoming sections.

In order to give a Calderón-Zygmund framework for von Neumann algebras we start with some initial considerations, which determine the general form of Markov metrics that we shall work with. Consider a Markov metric \mathcal{Q} associated to a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ acting on (\mathcal{M}, τ) . Then, we shall assume that the cpu maps $R_{j,t}$ from \mathcal{Q} are of the following form

$$\begin{aligned} \mathcal{M} &\xrightarrow{\rho_j} \mathcal{N}_{\rho} \xrightarrow{\mathbb{E}_{\rho}} \rho_1(\mathcal{M}) \simeq \mathcal{M}, \\ R_{j,t}f &= \mathbb{E}_{\rho}(q_{j,t})^{-\frac{1}{2}} \mathbb{E}_{\rho}(q_{j,t} \rho_2(f) q_{j,t}) \mathbb{E}_{\rho}(q_{j,t})^{-\frac{1}{2}}, \end{aligned} \quad (2.1)$$

where $\rho_1, \rho_2 : \mathcal{M} \rightarrow \mathcal{N}_{\rho}$ are injective $*$ -homomorphisms into certain von Neumann algebra \mathcal{N}_{ρ} , the map $\mathbb{E}_{\rho} : \mathcal{N}_{\rho} \rightarrow \rho_1(\mathcal{M})$ is an operator-valued weight which will be defined later and the $q_{j,t}$'s are projections in \mathcal{N}_{ρ} . In particular, we shall assume that our formula for $R_{j,t}f$ makes sense so that $q_{j,t}$ and $q_{j,t} \rho_2(f) q_{j,t}$ belong to the domain of \mathbb{E}_{ρ} , see Section 2.1 for further details. Our model provides a quite general form of Markov metric which includes the Markov metric for the heat semigroup considered before. Indeed, take $\mathcal{N}_{\rho} = L_{\infty}(\mathbb{R}^n \times \mathbb{R}^n)$ with $\rho_1 f(x, y) = f(x)$ and $\rho_2 f(x, y) = f(y)$. Let \mathbb{E}_{ρ} be the integral in \mathbb{R}^n with respect to the variable y and set

$$q_{j,t}(x, y) = \chi_{B_{\sqrt{4jt}}(x)}(y) = \chi_{B_{\sqrt{4jt}}(y)}(x) = \chi_{|x-y| < \sqrt{4jt}}.$$

Then, it is straightforward to check that we recover from (2.1) the $R_{j,t}$'s for the heat semigroup. Note that the $q_{j,t}(x, \cdot)$'s reproduce in this case all the Euclidean balls in \mathbb{R}^n . Morally, this is why we call \mathcal{Q} a Markov metric, since it codifies some sort of underlying

metric in (\mathcal{M}, τ) . According to our definition of $\text{BMO}_{\mathcal{Q}}$, we shall also consider projections q_t in \mathcal{N}_{ρ} and cpu maps

$$M_t f = E_{\rho}(q_t)^{-\frac{1}{2}} E_{\rho}(q_t \rho_2(f) q_t) E_{\rho}(q_t)^{-\frac{1}{2}}. \quad (2.2)$$

In this case, to show that $\text{BMO}_{\mathcal{Q}}$ is equivalent to $\text{BMO}_{\mathbb{R}^n}$, we have seen in Paragraph 1.6, the non trivial part is to prove $\|f\|_{\text{BMO}_{\mathcal{Q}}} \lesssim \|f\|_{\text{BMO}_{\mathbb{R}^n}}$. Note that in the definition of $\text{BMO}_{\mathcal{Q}}$ -norm, the infimum is taken for all cpu maps M_t . So it is enough to show the inequality for some particular M_t .

2.1. Operator-valued weights

In this subsection we briefly review the definition and basic properties of operator-valued weights from [23,24]. A unital, weakly closed $*$ -subalgebra is called a von Neumann subalgebra. A conditional expectation $\mathcal{E}_{\mathcal{M}} : \mathcal{N} \rightarrow \mathcal{M}$ onto a von Neumann subalgebra \mathcal{M} is a positive unital projection satisfying the bimodular property $\mathcal{E}_{\mathcal{M}}(a_1 f a_2) = a_1 \mathcal{E}_{\mathcal{M}}(f) a_2$ for all $a_1, a_2 \in \mathcal{M}$. It is called normal if $\sup_{\alpha} \mathcal{E}_{\mathcal{M}}(f_{\alpha}) = \mathcal{E}_{\mathcal{M}}(\sup_{\alpha} f_{\alpha})$ for bounded increasing nets (f_{α}) in \mathcal{N}_{+} . A normal conditional expectation such that $\tau \circ \mathcal{E}_{\mathcal{M}} = \tau$ exists if and only if the restriction of τ to the von Neumann subalgebra \mathcal{M} remains semifinite [60].

The *extended positive part* $\widehat{\mathcal{M}}_{+}$ of the von Neumann algebra \mathcal{M} is the set of lower semicontinuous maps $m : \mathcal{M}_{*,+} \rightarrow [0, \infty]$ which are linear on the positive cone, $m(\lambda_1 \phi_1 + \lambda_2 \phi_2) = \lambda_1 m(\phi_1) + \lambda_2 m(\phi_2)$ for $\lambda_j \geq 0$ and $\phi_j \in \mathcal{M}_{*,+}$. The extended positive part is closed under addition, increasing limits and is fixed by the map $x \mapsto a^* x a$ for any $a \in \mathcal{M}$. It is clear that \mathcal{M}_{+} sits in the extended positive part. When \mathcal{M} is abelian, we find $\mathcal{M} \simeq L_{\infty}(\Omega, \mu)$ for some measure space (Ω, μ) and the extended positive part corresponds in this case to the set of μ -measurable functions on Ω (module sets of zero measure) with values in $[0, \infty]$. A harder characterization of the extended positive part for arbitrary von Neumann algebras was found by Haagerup in [23]. Assume that \mathcal{M} acts on \mathcal{H} and consider a positive operator A affiliated with \mathcal{M} with spectral resolution $A = \int_{\mathbb{R}_{+}} \lambda d e_{\lambda}$. Then, we may construct an associated element in $\widehat{\mathcal{M}}_{+}$

$$m_A(\phi) = \int_{\mathbb{R}_{+}} \lambda d(\phi(e_{\lambda})).$$

In general, any $m \in \widehat{\mathcal{M}}_{+}$ has a unique spectral resolution

$$m(\phi) = \int_{\mathbb{R}_{+}} \lambda d(\phi(e_{\lambda})) + \infty \phi(p)$$

where the e_{λ} 's form an increasing family of projections in \mathcal{M} and p is the projection $\mathbf{1}_{\mathcal{M}} - \lim_{\lambda} e_{\lambda}$. Moreover, the map $\lambda \mapsto e_{\lambda}$ is strongly continuous from the right and we

find that $e_0 = 0$ iff m does not vanish on $\mathcal{M}_*^+ \setminus \{0\}$, while $p = 0$ iff the family of $\phi \in \mathcal{M}_*^+$ with $m(\phi) < \infty$ is dense in \mathcal{M}_*^+ .

Operator-valued weights appear as “unbounded conditional expectations” and the simplest nontrivial model is perhaps a partial trace $E_{\mathcal{M}} = \text{tr}_{\mathcal{A}} \otimes \text{id}_{\mathcal{M}}$ with $\mathcal{N} = \mathcal{A} \bar{\otimes} \mathcal{M}$ and \mathcal{A} a semifinite non-finite von Neumann algebra. In general, an *operator-valued weight* from \mathcal{N} to \mathcal{M} is just a linear map

$$E_{\mathcal{M}} : \mathcal{N}_+ \rightarrow \widehat{\mathcal{M}}_+ \quad \text{satisfying} \quad E_{\mathcal{M}}(a^*fa) = a^*E_{\mathcal{M}}(f)a$$

for all $a \in \mathcal{M}$. As usual, $E_{\mathcal{M}}$ is called normal when $\sup_{\alpha} E_{\mathcal{M}}(f_{\alpha}) = E_{\mathcal{M}}(\sup_{\alpha} f_{\alpha})$ for bounded increasing nets (f_{α}) in \mathcal{N}_+ . Since $a^*fb = \frac{1}{4} \sum_{k=0}^3 i^{-k} (a + i^k b)^* f (a + i^k b)$ by polarization, we see that bimodularity of conditional expectations is equivalent to $E_{\mathcal{M}}(a^*fa) = a^*E_{\mathcal{M}}(f)a$ for $a \in \mathcal{M}$. In particular, the fundamental properties which operator-valued weights loose with respect to conditional expectations are unitality and the fact that unboundedness is allowed for the image. Additionally, when $\mathcal{M} = \mathbb{C}$ the map $E_{\mathcal{M}}$ becomes an ordinary weight on \mathcal{N} . In analogy with ordinary weights, we take

$$L_{\infty}^c(\mathcal{N}; E_{\mathcal{M}}) = \left\{ f \in \mathcal{N} : \|E_{\mathcal{M}}(f^*f)\|_{\mathcal{M}} < \infty \right\}.$$

Note that when $E_{\mathcal{M}} = \text{tr}_{\mathcal{A}} \otimes \text{id}_{\mathcal{M}}$ with $\mathcal{N} = \mathcal{A} \bar{\otimes} \mathcal{M}$, $L_{\infty}^c(\mathcal{N}; E_{\mathcal{M}})$ are the Hilbert space valued noncommutative L_{∞} spaces defined in [28], which we denote by $L_2^c(\mathcal{A}) \bar{\otimes} \mathcal{M}$. Let $\mathcal{N}_{E_{\mathcal{M}}}$ be the linear span of $f_1^*f_2$ with $f_1, f_2 \in L_{\infty}^c(\mathcal{N}; E_{\mathcal{M}})$. Then we find

- i) $\mathcal{N}_{E_{\mathcal{M}}} = \text{span}\{f \in \mathcal{N}_+ : \|E_{\mathcal{M}}f\| < \infty\}$,
- ii) $L_{\infty}^c(\mathcal{N}; E_{\mathcal{M}})$ and $\mathcal{N}_{E_{\mathcal{M}}}$ are two-sided modules over \mathcal{M} ,
- iii) $E_{\mathcal{M}}$ has a unique linear extension $E_{\mathcal{M}} : \mathcal{N}_{E_{\mathcal{M}}} \rightarrow \mathcal{M}$ satisfying

$$E_{\mathcal{M}}(a_1fa_2) = a_1E_{\mathcal{M}}(f)a_2 \quad \text{with} \quad f \in \mathcal{N}_{E_{\mathcal{M}}} \text{ and } a_1, a_2 \in \mathcal{M}.$$

In particular, if $E_{\mathcal{M}}(\mathbf{1}) = \mathbf{1}$ we recover a conditional expectation onto \mathcal{M} . An operator-valued weight $E_{\mathcal{M}}$ is called faithful if $E_{\mathcal{M}}(f^*f) = 0$ implies $f = 0$ and semifinite when $L_{\infty}^c(\mathcal{N}; E_{\mathcal{M}})$ is σ -weakly dense in \mathcal{N} . It is of interest to determine for which pairs $(\mathcal{N}, \mathcal{M})$ we may construct n.s.f. operator-valued weights. Among other results, Haagerup proved in [24] that this is the case when both von Neumann algebras are semifinite and there exists a unique trace preserving one. Note that conditional expectations do not always exist in this case. He also proved that given $E_{\mathcal{M}_j}$ n.s.f. operator-valued weights in $(\mathcal{N}_j, \mathcal{M}_j)$ for $j = 1, 2$, there exists a unique n.s.f. operator-valued weight $E_{\mathcal{M}_1 \otimes \mathcal{M}_2}$ associated to $(\mathcal{N}_1 \bar{\otimes} \mathcal{N}_2, \mathcal{M}_1 \bar{\otimes} \mathcal{M}_2)$ such that $(\phi_1 \otimes \phi_2) \circ E_{\mathcal{M}_1 \otimes \mathcal{M}_2} = (\phi_1 \circ E_{\mathcal{M}_1}) \otimes (\phi_2 \circ E_{\mathcal{M}_2})$ for any pair (ϕ_1, ϕ_2) of normal semifinite faithful weights on $(\mathcal{M}_1, \mathcal{M}_2)$.

2.2. Algebraic/analytic conditions

The identity

$$Tf(x) = \int_{\Omega} k(x, y)f(y) d\mu(y)$$

is just a vague expression to consider classical Calderón-Zygmund operators. It is well-known that particular realizations as above are only meaningful outside the support of f and understanding k as a distribution which coincides with a locally integrable function on $\mathbb{R}^n \times \mathbb{R}^n \setminus \Delta$. Instead of that, we shall not specify any kernel representation of our CZOs since our conditions below will be formulated in a more intrinsic way. These kernel representations will appear later on in this paper with the concrete examples that we shall consider.

Let T be a densely defined operator on \mathcal{M} , which means that $Tf \in \mathcal{M}$ for all f in a weak-* dense subalgebra $\mathcal{A}_{\mathcal{M}}$ of \mathcal{M} . Our assumption does not necessarily hold for classical Calderón-Zygmund operators defined in abelian von Neumann algebras $(\mathcal{M}, \tau) = L_{\infty}(\Omega, \mu)$, but it is true for the truncated singular integral operators satisfying the standard size condition for the kernel, take for instance $\mathcal{A}_{\mathcal{M}} = \mathcal{M} \cap L_1(\mathcal{M})$. In particular, this is not a crucial restriction since we shall be able to take L_p -limits as far as our estimates below are independent of T . Our aim is to settle conditions on T of CZ type assuring that $T : L_{\infty}(\mathcal{M}) \rightarrow \text{BMO}_{\mathcal{Q}}^c$, provided (\mathcal{M}, τ) comes equipped with a Markov metric \mathcal{Q} . In this paragraph, we establish some preliminary algebraic and analytic conditions on the Markov metric and the CZO. Consider injective *-homomorphisms $\pi_1, \pi_2 : \mathcal{M} \rightarrow \mathcal{N}_{\pi}$ and an operator-valued weight $E_{\pi} : \mathcal{N}_{\pi} \rightarrow \pi_1(\mathcal{M})$ which may or may not coincide with ρ_1, ρ_2 and E_{ρ} from (2.1). Assume there exists a map defined on a weak-* dense subalgebra $\mathcal{A}_{\mathcal{N}_{\pi}}$ of \mathcal{N}_{π}

$$\begin{aligned} \widehat{T} : \mathcal{A}_{\mathcal{N}_{\pi}} \subset \mathcal{N}_{\pi} &\rightarrow \mathcal{N}_{\rho} \\ \text{satisfying } \widehat{T} \circ \pi_2 &= \rho_2 \circ T \quad \text{on } \mathcal{A}_{\mathcal{M}}. \end{aligned} \tag{2.3}$$

Algebraic conditions:

i) \mathcal{Q} -monotonicity of E_{ρ}

$$E_{\rho}(q_{j,t}|\xi|^2 q_{j,t}) \leq E_{\rho}(|\xi|^2)$$

for all $\xi \in \mathcal{N}_{\rho}$ and every projection $q_{j,t}$ determined by \mathcal{Q} via the identity in (2.1). Similarly, we assume the same inequality holds when we replace the $q_{j,t}$'s by the q_t 's appearing in (2.2).

ii) Right \mathcal{B} -modularity of \widehat{T}

$$\widehat{T}(\eta \pi_1 \rho_1^{-1}(b)) = \widehat{T}(\eta)b$$

for all $\eta \in \mathcal{A}_{\mathcal{N}_\pi}$ and all b lying in some von Neumann subalgebra \mathcal{B} of $\rho_1(\mathcal{M})$ which includes $E_\rho(q_t)$, $E_\rho(q_{j,t})$ and $\rho_1(\gamma_{j,t})$ for every projection q_t and $q_{j,t}$ determined by \mathcal{Q} via the identities in (2.1) and (2.2).

As we shall see both conditions trivially hold in the classical theory, where the first condition essentially says that integrating a positive function over a “Markov metric ball” is always smaller than integrating it over the whole space, while the second condition allows to take out x -dependent functions from the y -dependent integral defining T . Our conditions remain true in many other situations, which will be explored below in this paper. Nevertheless, condition i) suggests that certain amount of commutativity might be necessary to work with Markov metrics.

To state our analytic conditions we introduce an additional von Neumann algebra \mathcal{N}_σ containing \mathcal{N}_ρ as a von Neumann subalgebra. Then, we consider derivations $\delta : \mathcal{N}_\rho \rightarrow \mathcal{N}_\sigma$ given by the difference $\delta = \sigma_1 - \sigma_2$ of two $*$ -homomorphisms, so that $\delta(ab) = \sigma_1(a)\sigma_1(b) - \sigma_2(a)\sigma_2(b) = \delta(a)\sigma_1(b) + \sigma_2(a)\delta(b)$ as expected. We also consider the natural amplification maps which are normal cpu maps defined on \mathcal{N}_ρ ,

$$\widehat{R}_{j,t} : \mathcal{N}_\rho \ni \xi \mapsto E_\rho(q_{j,t})^{-\frac{1}{2}} E_\rho(q_{j,t} \xi q_{j,t}) E_\rho(q_{j,t})^{-\frac{1}{2}} \in \rho_1(\mathcal{M}),$$

$$\widehat{M}_t : \mathcal{N}_\rho \ni \xi \mapsto E_\rho(q_t)^{-\frac{1}{2}} E_\rho(q_t \xi q_t) E_\rho(q_t)^{-\frac{1}{2}} \in \rho_1(\mathcal{M}).$$

Analytic conditions:

i) Mean differences conditions

- $\widehat{R}_{j,t}(\xi^* \xi) - \widehat{R}_{j,t}(\xi)^* \widehat{R}_{j,t}(\xi) \leq \Phi_{j,t}(\delta(\xi)^* \delta(\xi)),$
- $[\widehat{R}_{j,t}(\xi) - \widehat{M}_t(\xi)]^* [\widehat{R}_{j,t}(\xi) - \widehat{M}_t(\xi)] \leq \Psi_{j,t}(\delta(\xi)^* \delta(\xi)),$

for some derivation $\delta : \mathcal{N}_\rho \rightarrow \mathcal{N}_\sigma$ and cpu maps $\Phi_{j,t}, \Psi_{j,t} : \mathcal{N}_\sigma \rightarrow \rho_1(\mathcal{M})$.

ii) Metric/measure growth conditions

- $\mathbf{1} \leq \pi_1 \rho_1^{-1} E_\rho(q_t)^{-\frac{1}{2}} E_\pi(a_t^* a_t) \pi_1 \rho_1^{-1} E_\rho(q_t)^{-\frac{1}{2}} \lesssim \pi_1 \rho_1^{-1}(\gamma_{j,t}^2),$
- $\mathbf{1} \leq \pi_1 \rho_1^{-1} E_\rho(q_{j,t})^{-\frac{1}{2}} E_\pi(a_{j,t}^* a_{j,t}) \pi_1 \rho_1^{-1} E_\rho(q_{j,t})^{-\frac{1}{2}} \lesssim \pi_1 \rho_1^{-1}(\gamma_{j,t}^2),$

for some family of operators $a_t, a_{j,t} \in \mathcal{N}_\pi$ to be determined later on.

A complete determination of the operators a_t and $a_{j,t}$ is only possible after imposing additional size and smoothness conditions in our definition of Calderón-Zygmund operator below. Nevertheless, we shall see that these operators will play the role of “dilated Markov balls” as it is the case in classical CZ theory. In fact, in the classical case our last condition trivially holds for doubling measures, and also for measures of polynomial or even exponential growth provided we find a Markov metric with large enough $\gamma_{j,t}$ ’s. Our assertions will be illustrated below. The first condition takes the form in the classical case of a couple of easy consequences of Jensen’s inequality, namely

$$\begin{aligned} \left| \int_{B_1} |f|^2 d\mu - \left| \int_{B_1} f d\mu \right|^2 \right| &\leq \int_{B_1 \times B_1} |f(y) - f(z)|^2 d\mu(y) d\mu(z), \\ \left| \int_{B_1} f d\mu - \int_{B_2} f d\mu \right|^2 &\leq \int_{B_1 \times B_2} |f(y) - f(z)|^2 d\mu(y) d\mu(z). \end{aligned} \quad (2.4)$$

2.3. CZ extrapolation

Now we introduce CZOs in this context. As we already mentioned, we consider a priori densely defined (unbounded) maps $T : \mathcal{A}_{\mathcal{M}} \rightarrow \mathcal{M}$ whose amplified maps are right \mathcal{B} -modules according to our algebraic assumptions above. In addition, we impose three conditions generalizing L_2 boundedness, the size and the smoothness conditions for the kernel.

Calderón-Zygmund type conditions:

i) *Boundedness condition*

$$\widehat{T} : L_{\infty}^c(\mathcal{N}_{\pi}; \mathbf{E}_{\pi}) \rightarrow L_{\infty}^c(\mathcal{N}_{\rho}; \mathbf{E}_{\rho}).$$

ii) *Size “kernel” condition*

$$\begin{aligned} \bullet \widehat{M}_t \left(\left| \widehat{T}(\pi_2(f)(A_{j,t} - a_t)) \right|^2 \right) &\lesssim \gamma_{j,t}^2 \|f\|_{\infty}^2, \\ \bullet \widehat{R}_{j,t} \left(\left| \widehat{T}(\pi_2(f)(A_{j,t} - a_{j,t})) \right|^2 \right) &\lesssim \gamma_{j,t}^2 \|f\|_{\infty}^2, \end{aligned}$$

for for any $f \in \mathcal{A}_{\mathcal{M}}$ and a family of operators $A_{j,t} \in \mathcal{N}_{\pi}$ with $A_{j,t} \geq a_{j,t}, a_t$ to be determined.

iii) *Smoothness “kernel” condition*

$$\begin{aligned} \bullet \Phi_{j,t} \left(\left| \delta(\widehat{T}(\pi_2(f)(\mathbf{1} - a_{j,t}))) \right|^2 \right) &\lesssim \gamma_{j,t}^2 \|f\|_{\infty}^2, \\ \bullet \Psi_{j,t} \left(\left| \delta(\widehat{T}(\pi_2(f)(\mathbf{1} - A_{j,t}))) \right|^2 \right) &\lesssim \gamma_{j,t}^2 \|f\|_{\infty}^2. \end{aligned}$$

Let $T : \mathcal{A}_{\mathcal{M}} \rightarrow \mathcal{M}$ be a densely defined map which admits an amplification \widehat{T} satisfying (2.3). Any such T will be called an *algebraic column CZO* whenever the amplification map is right \mathcal{B} -modular and satisfies the CZ conditions we have given above. At first sight, our boundedness assumption might appear to be unrelated to the classical condition. The reader could have expected the L_2 boundedness of T , but our assumption is formally equivalent to it in the classical case and gives the right condition for more general algebras. On the other hand, our size and smoothness conditions are intrinsic in the sense that the kernel is not specified under this degree of generality. We shall recover classical kernel type estimates from our conditions in our examples below. As explained above, the operators $a_t, a_{j,t}$ and $A_{j,t}$ play the role of dilated Markov balls and our conditions were somehow modeled by Tolsa’s arguments in [61]. Perhaps a significant difference—in contrast to Tolsa’s approach—is that our smoothness condition is analog to a Hörmander type condition, more than the (stronger) Lipschitz regularity assumption.

Theorem 2.1. *Let (\mathcal{M}, τ) be a noncommutative measure space equipped with a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ with associated Markov metric \mathcal{Q} which fulfills our algebraic and analytic assumptions. Then, any algebraic column CZO T defines a bounded operator*

$$T : \mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{\mathcal{Q}}^c.$$

Proof. The first goal is to estimate the norm of

$$A = \gamma_{j,t}^{-1} \left(R_{j,t} |Tf|^2 - |R_{j,t} Tf|^2 \right) \gamma_{j,t}^{-1}.$$

The map $\Pi_{j,t} : \mathcal{M} \bar{\otimes}_{R_{j,t}} \mathcal{M} \ni a \otimes b \mapsto \mathbf{1} \otimes R_{j,t}(a)b \in \mathbf{1} \otimes \mathcal{M}$ extends to a right $(\mathbf{1} \otimes \mathcal{M})$ -module projection, which is well-defined in the sense that $\langle \xi, \xi \rangle_{R_{j,t}} = 0$ implies $\Pi_{j,t}(\xi) = 0$. Now, since

$$A = \gamma_{j,t}^{-1} \left\langle Tf \otimes \mathbf{1} - \mathbf{1} \otimes R_{j,t} Tf, Tf \otimes \mathbf{1} - \mathbf{1} \otimes R_{j,t} Tf \right\rangle_{R_{j,t}} \gamma_{j,t}^{-1},$$

we may use $\Pi_{j,t}$ to deduce the following identity

$$A = \left\langle (id - \Pi_{j,t})(Tf \otimes \gamma_{j,t}^{-1}), (id - \Pi_{j,t})(Tf \otimes \gamma_{j,t}^{-1}) \right\rangle_{R_{j,t}}.$$

Consider the amplification maps $\widehat{R}_{j,t}$ defined on \mathcal{N}_{ρ} and $\widehat{\Pi}_{j,t}$ defined by $\mathcal{N}_{\rho} \otimes \mathcal{M} \ni n \otimes b \mapsto \mathbf{1} \otimes \widehat{R}_{j,t}(n)b \in \mathbf{1} \otimes \mathcal{M}$. It is easy to see that

$$R_{j,t} = \widehat{R}_{j,t} \circ \rho_2 \quad \text{and} \quad \Pi_{j,t} = \widehat{\Pi}_{j,t} \circ (\rho_2 \otimes id).$$

By (2.3), it turns out that $A = \langle \mathbf{a}, \mathbf{a} \rangle_{\widehat{R}_{j,t}}$ where

$$\begin{aligned} \mathbf{a} &= (id - \widehat{\Pi}_{j,t})(\rho_2 Tf \otimes \gamma_{j,t}^{-1}) \\ &= (id - \widehat{\Pi}_{j,t})(\widehat{T} \pi_2 f \otimes \gamma_{j,t}^{-1}) \\ &= (id - \widehat{\Pi}_{j,t})(\widehat{T}(\pi_2(f) a_{j,t}) \otimes \gamma_{j,t}^{-1}) \\ &\quad + (id - \widehat{\Pi}_{j,t})(\widehat{T}(\pi_2(f)(\mathbf{1} - a_{j,t})) \otimes \gamma_{j,t}^{-1}) = \mathbf{a}_1 + \mathbf{a}_2 \end{aligned}$$

According to Lemma 1.1 i), we may estimate A as follows

$$A \lesssim \langle \mathbf{a}_1, \mathbf{a}_1 \rangle_{\widehat{R}_{j,t}} + \langle \mathbf{a}_2, \mathbf{a}_2 \rangle_{\widehat{R}_{j,t}} = A_1 + A_2.$$

Since $\widehat{\Pi}_{j,t}(n \otimes b) = \mathbf{1} \otimes \widehat{R}_{j,t}(n)b$, the Kadison-Schwarz inequality yields

$$\left\langle \widehat{\Pi}_{j,t}(n \otimes b), \widehat{\Pi}_{j,t}(n \otimes b) \right\rangle_{\widehat{R}_{j,t}} \lesssim \langle n \otimes b, n \otimes b \rangle_{\widehat{R}_{j,t}}.$$

In conjunction with Lemma 1.1 i) again, we deduce the following estimate for A_1

$$A_1 \lesssim \left\langle \widehat{T}(\pi_2(f)a_{j,t}) \otimes \gamma_{j,t}^{-1}, \widehat{T}(\pi_2(f)a_{j,t}) \otimes \gamma_{j,t}^{-1} \right\rangle_{\widehat{R}_{j,t}} = \gamma_{j,t}^{-1} \widehat{R}_{j,t} \left(|\widehat{T}(\pi_2(f)a_{j,t})|^2 \right) \gamma_{j,t}^{-1}.$$

In order to bound the term in the right hand side, we apply (2.1) and the properties of the operator-valued weight E_ρ together with our algebraic conditions. Indeed, we first use the \mathcal{Q} -monotonicity of E_ρ ; then the fact that it commutes with the left/right multiplication by elements affiliated to \mathcal{M} (like $\gamma_{j,t}^{-1}$ or $E_\rho(q_{j,t})^{-1/2}$); finally we use the right \mathcal{B} -modularity of the amplification of T :

$$\begin{aligned} & \gamma_{j,t}^{-1} \widehat{R}_{j,t} \left(|\widehat{T}(\pi_2(f)a_{j,t})|^2 \right) \gamma_{j,t}^{-1} \\ & \leq \gamma_{j,t}^{-1} E_\rho(q_{j,t})^{-\frac{1}{2}} E_\rho \left(|\widehat{T}(\pi_2(f)a_{j,t})|^2 \right) E_\rho(q_{j,t})^{-\frac{1}{2}} \gamma_{j,t}^{-1} \\ & = E_\rho \left(\gamma_{j,t}^{-1} E_\rho(q_{j,t})^{-\frac{1}{2}} |\widehat{T}(\pi_2(f)a_{j,t})|^2 E_\rho(q_{j,t})^{-\frac{1}{2}} \gamma_{j,t}^{-1} \right) \\ & = E_\rho \left| \widehat{T} \left(\underbrace{\pi_2(f)a_{j,t} \pi_1 \rho_1^{-1} (E_\rho(q_{j,t})^{-\frac{1}{2}} \gamma_{j,t}^{-1})}_{\xi_1} \right) \right|^2 = E_\rho |\widehat{T}(\xi_1)|^2. \end{aligned}$$

Now, our first CZ condition i) gives the boundedness we need since

$$\begin{aligned} \|A_1\|_{\mathcal{M}} & \leq \|\widehat{T}(\xi_1)\|_{L_\infty^c(\mathcal{N}_\rho; E_\rho)}^2 \lesssim \|\xi_1\|_{L_\infty^c(\mathcal{N}_\pi; E_\pi)}^2 \\ & = \left\| E_\pi \left(|\pi_2(f)a_{j,t} \pi_1 \rho_1^{-1} (E_\rho(q_{j,t})^{-\frac{1}{2}} \gamma_{j,t}^{-1})|^2 \right) \right\|_{\mathcal{M}} \\ & \leq \left\| \pi_1 \rho_1^{-1} (\gamma_{j,t}^{-1} E_\rho(q_{j,t})^{-\frac{1}{2}})^* E_\pi(a_{j,t}^* a_{j,t}) \pi_1 \rho_1^{-1} (E_\rho(q_{j,t})^{-\frac{1}{2}} \gamma_{j,t}^{-1}) \right\|_{\mathcal{M}} \|f\|_\infty^2. \end{aligned}$$

The last term on the right is dominated by $\|f\|_\infty^2$ according to our second analytic condition on metric/measure growth. The estimate for A_2 is simpler. Indeed, if we set $\xi_2 = \widehat{T}(\pi_2(f)(1 - a_{j,t}))$ then

$$\begin{aligned} A_2 & = \left\langle (id - \widehat{\Pi}_{j,t})(\xi_2 \otimes \gamma_{j,t}^{-1}), (id - \widehat{\Pi}_{j,t})(\xi_2 \otimes \gamma_{j,t}^{-1}) \right\rangle_{\widehat{R}_{j,t}} \\ & = \gamma_{j,t}^{-1} \left(\widehat{R}_{j,t} |\xi_2|^2 - |\widehat{R}_{j,t}(\xi_2)|^2 \right) \gamma_{j,t}^{-1} \leq \gamma_{j,t}^{-1} \Phi_{j,t}(|\delta \xi_2|^2) \gamma_{j,t}^{-1} \lesssim \|f\|_\infty^2 \mathbf{1}, \end{aligned}$$

where the first inequality holds for some derivation $\delta : \mathcal{N}_\rho \rightarrow \mathcal{N}_\sigma$ and some normal cpu map $\Phi_{j,t} : \mathcal{N}_\sigma \rightarrow \rho_1(\mathcal{M})$ by our first analytic condition on mean differences. Then our CZ condition iii) on kernel smoothness justifies our last estimate. Our estimates so far prove the desired estimate

$$\sup_{t>0} \sup_{j \geq 1} \left\| \left(\gamma_{j,t}^{-1} [R_{j,t} |Tf|^2 - |R_{j,t} T f|^2] \gamma_{j,t}^{-1} \right)^{\frac{1}{2}} \right\|_{\mathcal{M}} \lesssim \|f\|_\infty.$$

Therefore, it remains to estimate the norm of

$$B = \gamma_{j,t}^{-1} \left(|R_{j,t}Tf - M_tTf|^2 \right) \gamma_{j,t}^{-1}.$$

To do so, we decompose the middle term using (2.3) as follows

$$\begin{aligned} R_{j,t}Tf - M_tTf &= \widehat{R}_{j,t}(\rho_2Tf) - \widehat{M}_t(\rho_2Tf) \\ &= \widehat{R}_{j,t}(\widehat{T}(\pi_2(f)a_{j,t})) - \widehat{M}_t(\widehat{T}(\pi_2(f)a_t)) \\ &\quad + \left[\widehat{R}_{j,t}(\widehat{T}(\pi_2(f)(1 - a_{j,t}))) - \widehat{M}_t(\widehat{T}(\pi_2(f)(1 - a_t))) \right] = \mathbf{b}_1 - \mathbf{b}_2 + \mathbf{b}_3. \end{aligned}$$

Letting $B_j = \gamma_{j,t}^{-1} |\mathbf{b}_j|^2 \gamma_{j,t}^{-1}$ we get $B \lesssim B_1 + B_2 + B_3$. By Kadison-Schwarz we get

$$B_1 \leq \gamma_{j,t}^{-1} \widehat{R}_{j,t} \left(|\widehat{T}(\pi_2(f)a_{j,t})|^2 \right) \gamma_{j,t}^{-1} \lesssim \|f\|_\infty^2,$$

where the last inequality was justified in our estimate of A_1 above. Replacing $q_{j,t}$ by q_t , the same argument serves to control the term B_2 . To estimate B_3 we decompose \mathbf{b}_3 as follows

$$\begin{aligned} \mathbf{b}_3 &= \left[\widehat{R}_{j,t}(\widehat{T}(\pi_2(f)(1 - A_{j,t}))) - \widehat{M}_t(\widehat{T}(\pi_2(f)(1 - A_{j,t}))) \right] \\ &\quad + \widehat{R}_{j,t}(\widehat{T}(\pi_2(f)(A_{j,t} - a_{j,t}))) - \widehat{M}_t(\widehat{T}(\pi_2(f)(A_{j,t} - a_t))) = \mathbf{b}_{31} + \mathbf{b}_{32} - \mathbf{b}_{33}. \end{aligned}$$

Taking $\xi_3 = \widehat{T}(\pi_2(f)(1 - A_{j,t}))$ and applying our analytic condition i) on mean differences together with our CZ condition iii) on kernel smoothness, we obtain that

$$\gamma_{j,t}^{-1} |\mathbf{b}_{31}|^2 \gamma_{j,t}^{-1} \lesssim \gamma_{j,t}^{-1} \Psi_{j,t}(|\delta \xi_3|^2) \gamma_{j,t}^{-1} \lesssim \|f\|_\infty^2.$$

It remains to estimate the terms B_{32} and B_{33} . Applying the Kadison-Schwarz inequality, it is easily checked that these terms are also dominated by $\|f\|_\infty^2$ by means of our CZ size kernel condition ii). Altogether, we have justified that

$$\sup_{t>0} \inf_{M_t \text{cpu}} \sup_{j \geq 1} \left\| \left(\gamma_{j,t}^{-1} [R_{j,t}f - M_t f]^2 \right) \gamma_{j,t}^{-1} \right\|_{\mathcal{M}}^{\frac{1}{2}} \lesssim \|f\|_\infty.$$

Combining our estimates for A and B , we deduce that $T : \mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{\mathcal{Q}}^c$. \square

The $\mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{\mathcal{Q}}^r$ boundedness of the map T is equivalent to the $\mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{\mathcal{Q}}^c$ boundedness of the map $T^\dagger(f) = T(f^*)^*$. According to this, an *algebraic* CZO is any column CZO T for which T^\dagger remains a column CZO. By Theorem 2.1, we know that any algebraic CZO T associated to $(\mathcal{M}, \tau, \mathcal{Q})$ as above is automatically $\mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{\mathcal{Q}}$

bounded. Assuming L_2 boundedness and regularity of the Markov semigroup, we may interpolate via Theorem 1.6. Under the same assumptions for T^* , we may also dualize and obtain the following extrapolation result.

Corollary 2.2. *Let (\mathcal{M}, τ) be a noncommutative measure space equipped with a Markov regular semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ and a Markov metric $\mathcal{Q} = (R_{j,t}, \sigma_{j,t}, \gamma_{j,t})$ fulfilling our algebraic and analytic assumptions. Then, every L_2 -bounded algebraic CZO T satisfies that $J_p T : L_p(\mathcal{M}) \rightarrow L_p^\circ(\mathcal{M})$ for $p > 2$. Applying duality, similar conditions for T^* yield L_p -boundedness of TJ_p for every $1 < p < 2$.*

Remark 2.3. Theorem 2.1 admits a completely bounded version in the category of operator spaces. Since the operator space structure [19,52] of BMO is determined by

$$M_m(\text{BMO}_{\mathcal{S}}(\mathcal{M})) = \text{BMO}_{\widehat{\mathcal{S}}}(M_m(\mathcal{M}))$$

for $\widehat{\mathcal{S}} = (id_{M_m} \otimes S_t)_{t \geq 0}$, we just need to replace \mathcal{M} by $M_m(\mathcal{M})$ everywhere, amplify all the involved maps by tensorizing with id_{M_m} and require that the hypotheses hold with constants independent of m . Then, we obtain the cb-boundedness of T .

Remark 2.4. As noticed in the Introduction, a common scenario is given by the choice $\mathcal{N}_\rho = \mathcal{M} \bar{\otimes} \mathcal{M}$ with $\rho_1(f) = f \otimes \mathbf{1}$ and $\rho_2(f) = \mathbf{1} \otimes f$, together with $E_\rho = id \otimes \tau$ and $\pi_j = \rho_j$ for $j = 1, 2$. In this case, it is clear that the amplification map is given by

$$\widehat{T} = id_{\mathcal{M}} \otimes T \quad \text{so that} \quad \widehat{T}\pi_2 = \rho_2 T.$$

In particular, it turns out that the L_2 boundedness of T in Corollary 2.2 follows automatically from our CZ boundedness condition i). This is the case in classical Calderón-Zygmund theory. It is also true when $\mathcal{N}_\rho = \mathcal{M} \bar{\otimes} \mathcal{A}$ for an auxiliary algebra \mathcal{A} and $\rho_2 = \text{flip} \circ \sigma$, where $\sigma : \mathcal{M} \rightarrow \mathcal{A} \bar{\otimes} \mathcal{M}$ is a $*$ -homomorphism satisfying $E_\rho \circ \rho_2(f) = \tau(f)\mathbf{1}_{\mathcal{M}}$. This leads to another significant family of examples. It is however surprising that in general, the L_2 boundedness and the CZ boundedness assumptions are a priori unrelated. Thus, CZ extrapolation requires in this context to verify two boundedness conditions. It would be quite interesting to explore the corresponding “ $T(1)$ problems” that arise naturally.

2.4. The classical theory revisited

We now illustrate our algebraic approach in the classical context of Euclidean spaces with the Lebesgue measure. This will help us to understand some of our conditions and will show how some others are automatic in a commutative framework. Take $\mathcal{M} = L_\infty(\mathbb{R}^n)$ with the Lebesgue measure and $\mathcal{S} = (H_t)_{t \geq 0}$ the heat semigroup $H_t = \exp(t\Delta)$. In Paragraph 1.6 we introduced the Markov metric \mathcal{Q} given by

- $\sigma_{j,t}^2 \equiv \frac{2e}{\sqrt{\pi}} j^{\frac{n}{2}} e^{-k}$ and $\gamma_{j,t}^2 \equiv j^{\frac{n}{2}} \geq 1$,
- $R_{j,t} f(x) = \frac{1}{|B_{\sqrt{4jt}}(x)|} \int_{B_{\sqrt{4jt}}(x)} f(y) dy$.

Moreover, as we explained at the beginning of this section

$$R_{j,t} f = E_{\rho}(q_{j,t})^{-\frac{1}{2}} E_{\rho}(q_{j,t} \rho_2(f) q_{j,t}) E_{\rho}(q_{j,t})^{-\frac{1}{2}}$$

satisfies our basic assumption (2.1). Here the amplification von Neumann algebra is $\mathcal{N}_{\rho} = L_{\infty}(\mathbb{R}^n \times \mathbb{R}^n)$, the $*$ -homomorphisms $\rho_j f(x_1, x_2) = f(x_j)$, the projections $q_{j,t}(x, y) = \chi_{B_{\sqrt{4jt}}(x)}(y)$, and the operator-valued weight E_{ρ} is the integration map with respect to the second variable. The cpu map M_t which appears in the definition of $BMO_{\mathcal{Q}}$ is still taken by $M_t = R_{1,t}$ as in Paragraph 1.6.

Taking $\mathcal{N}_{\pi} = \mathcal{N}_{\rho}$, $\mathcal{A}_{\mathcal{M}} = L_2(\mathbb{R}^n) \cap L_{\infty}(\mathbb{R}^n)$ and T a standard CZO in \mathbb{R}^n , the algebraic conditions trivially hold in this case. Let $E_{j,k,t}^x = B_{\sqrt{4jt}}(x) \times B_{\sqrt{4kt}}(x)$. Taking $\Phi_{j,t}$ and $\Psi_{j,t}$ to be the averaging maps over $E_{j,j,t}^x$ and $E_{j,1,t}^x$ respectively and the family of dilated balls $(A_{j,t}(x, y), a_{j,t}(x, y)) = (\chi_{\alpha B_{\sqrt{4jt}}(x)}(y), \chi_{5B_{\sqrt{4jt}}(x)}(y))$ with $a \geq 5$, we may recover the conditions as we explained right after stating them. Let us now show how our algebraic CZ conditions hold from the classical ones. The boundedness condition reduces to the classical one, see Remark 4.2 A). Our size conditions can be rewritten as follows:

- $\text{ess sup}_{x \in \mathbb{R}^n} \int_{B_{\sqrt{4t}}(x)} \left| \int_{\alpha B_{\sqrt{4jt}}(x) \setminus 5B_{\sqrt{4t}}(x)} k(y, z) f(z) dz \right|^2 dy \lesssim j^{\frac{n}{2}} \|f\|_{\infty}^2$,
- $\text{ess sup}_{x \in \mathbb{R}^n} \int_{B_{\sqrt{4jt}}(x)} \left| \int_{\alpha B_{\sqrt{4t}}(x) \setminus 5B_{\sqrt{4jt}}(x)} k(y, z) f(z) dz \right|^2 dy \lesssim j^{\frac{n}{2}} \|f\|_{\infty}^2$.

The above conditions follow from the usual size condition

$$|k(y, z)| \lesssim \frac{1}{|z - y|^n}.$$

Next, taking $E_{j,k,t}^x$ as above, our smoothness conditions are:

- $\text{ess sup}_{x \in \mathbb{R}^n} \int_{E_{j,1,t}^x} \left(\int_{(5B_{\sqrt{4jt}}(x))^c} (k(y_1, z) - k(y_2, z)) f(z) dz \right)^2 dy_1 dy_2 \lesssim j^{\frac{n}{2}} \|f\|_{\infty}^2$,
- $\text{ess sup}_{x \in \mathbb{R}^n} \int_{E_{j,j,t}^x} \left(\int_{(\alpha B_{\sqrt{4jt}}(x))^c} (k(y_1, z) - k(y_2, z)) f(z) dz \right)^2 dy_1 dy_2 \lesssim j^{\frac{n}{2}} \|f\|_{\infty}^2$.

The above conditions easily follow from the usual Hörmander condition

$$\operatorname{ess\,sup}_{y_1, y_2 \in \mathbb{R}^n} \int_{|y_1 - z| \geq 2|y_1 - y_2|} |k(y_1, z) - k(y_2, z)| \, dz < \infty.$$

Note that our algebraic CZ conditions are slightly weaker than the classical CZ conditions, but still sufficient for $L_\infty \rightarrow \text{BMO}$ and L_p -boundedness of the CZ map.

Remark 2.5. Our size condition is only used to estimate B in the proof of Theorem 2.1. We saw in Paragraph 1.6 that $B \lesssim A$ for the Euclidean metric. Thus, our size condition is not necessary here, as it also happens in the classical formulation.

3. Applications I — Commutative spaces

In this section we give specific constructions of Markov metrics on two basic commutative spaces: Riemannian manifolds with nonnegative Ricci curvature and Gaussian measure spaces. Beyond the Euclidean-Lebesguean setting considered above, these are the most relevant settings over which Calderón-Zygmund theory has been studied. As a good illustration of our algebraic method, we shall recover the extrapolation results. Noncommutative spaces will be explored later on.

3.1. Riemannian manifolds

Let (Ω, μ) be a measure space equipped with a Markov semigroup, so that we may construct the corresponding semigroup type BMO space. In order to study the $L_\infty \rightarrow \text{BMO}$ boundedness of CZOs in (Ω, μ) it is essential to identify a Markov metric to work with. Now we provide sufficient conditions for a semigroup on a Riemannian manifold to yield a Markov metric satisfying our algebraic/analytic conditions, so that Theorem 2.1 is applicable. Let us consider an n -dimensional complete Riemannian manifold (M, g) equipped with the geodesic distance d determined by the Riemannian metric g . Denote the volume of a geodesic ball centered at x with radius r by $\operatorname{vol}_g(B_r(x))$. Let S_M be a Markov semigroup on M given by

$$S_{M,t}f(x) = \int_M s_t(x, y)f(y) \, dy.$$

Proposition 3.1. *Assume that*

- i) *M has Ricci curvature ≥ 0 .*
- ii) *The kernel admits an upper bound*

$$s_t(x, y) \lesssim \frac{\phi(t)^{n+\varepsilon}}{\operatorname{vol}_g(B_{\phi(t)}(x))(d(x, y) + \phi(t))^{n+\varepsilon}},$$

for some strictly positive function ϕ and some parameter $\varepsilon > 0$.

Then \mathcal{S}_M admits a Markov metric satisfying the algebraic/analytic conditions.

Proof. If $\Sigma_{j,t}(x) = B_{2^j\phi(t)}(x)$, by estimating the upper bound on the corona centered at x with radii $2^{j-1}\phi(t)$ and $2^j\phi(t)$, our assumption gives

$$s_t(x, y) \lesssim \sum_{j=1}^{\infty} \frac{2^{-j(n+\varepsilon)}}{\text{vol}_g(\Sigma_{0,t}(x))} \chi_{\Sigma_{j,t}(x)}(y). \quad (3.1)$$

According to Davies [15, Theorem 5.5.1], non-negative Ricci curvature implies

$$\begin{aligned} \text{vol}_g(B_r(x)) &\leq c_n r^n, \\ \text{vol}_g(B_{\gamma r}(x)) &\leq \gamma^n \text{vol}_g(B_r(x)) \end{aligned}$$

for all $x \in M$, $r > 0$ and $\gamma > 1$. In particular, $\text{vol}_g(\Sigma_{j,t}(x)) \leq 2^{jn} \text{vol}_g(\Sigma_{0,t}(x))$. By (3.1), this implies that $(\sigma_{j,t}, \gamma_{j,t}) = (2^{-j\varepsilon/2}, 1)$ forms a Markov metric for \mathcal{S}_M in conjunction with the averaging maps

$$R_{j,t}f(x) = \int_{\Sigma_{j,t}(x)} f(y) dy \quad \text{for } (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+.$$

By assumption ii), $\ker A_\infty$ the fixed-point subspace of \mathcal{S}_M is the space of constant functions on M , so $R_{j,t}f = f$ for any $f \in \ker A_\infty$.

Our construction for $\mathcal{M} = L_\infty(M)$ and $\mathcal{A}_M = L_\infty(M) \cap L_2(M)$ follows the basic model in the Introduction and the one used above in the Euclidean setting: $\mathcal{N}_\rho = \mathcal{N}_\pi = \mathcal{M} \bar{\otimes} \mathcal{M}$ with ρ_j the canonical inclusion maps and $q_{j,t}(x, y) = \chi_{\Sigma_{j,t}(x)}(y) = \chi_{\Sigma_{j,t}(y)}(x)$. Then, the algebraic conditions for the Markov metric are obviously satisfied. Let us now check the analytic conditions. Taking $\mathcal{N}_\sigma = \mathcal{M} \bar{\otimes} \mathcal{M} \bar{\otimes} \mathcal{M}$, the derivation $\delta : \mathcal{N}_\rho \rightarrow \mathcal{N}_\sigma$ given by $\delta(a \otimes b) = a \otimes (\mathbf{1} \otimes b - b \otimes \mathbf{1})$ and the maps $M_t = R_{1,t}$, it turns out that the mean difference conditions follow from Jensen's inequality on normalized balls of (M, g) as it follows from our comments after the definition of the analytic conditions. It remains to consider the metric/measure growth conditions. By taking $a_{j,t}(x, y) = \chi_{\Sigma_{j+1,t}(x)}(y)$ and $(q_t, a_t) = (q_{1,t}, a_{1,t})$, these conditions reduce to show that

$$\text{vol}_g(B_{2^{j+1}\phi(t)}(x)) \approx \text{vol}_g(B_{2^j\phi(t)}(x)).$$

This follows in turn from the fact that M has a non-negative Ricci curvature. \square

Let (M, g) be a complete Riemannian manifold with non-negative Ricci curvature and let Δ be the Laplace-Beltrami operator. The heat semigroup \mathcal{S}_Δ generated by Δ is regular and admits a kernel on (M, g) satisfying the upper bound estimate mentioned in

the above proposition. We know from Davies [15, Theorem 5.5.11] that the heat kernel satisfies

$$h_t(x, y) \leq \frac{a_\delta}{\text{vol}_g(B_{\sqrt{t}}(x))} \exp\left(-\frac{d(x, y)^2}{4(1+\delta)t}\right) \quad (3.2)$$

for any $\delta > 0$ and certain constant a_δ . This implies that

$$\begin{aligned} h_t(x, y) &\lesssim \frac{a_\delta}{\text{vol}_g(B_{\sqrt{t}}(x))} \frac{(4(1+\delta)t)^{\frac{n+\varepsilon}{2}}}{(d(x, y)^2 + 4(1+\delta)t)^{\frac{n+\varepsilon}{2}}} \\ &\lesssim \frac{(\sqrt{4(1+\delta)t})^{n+\varepsilon}}{\text{vol}_g(B_{\sqrt{4(1+\delta)t}}(x))(d(x, y) + \sqrt{4(1+\delta)t})^{n+\varepsilon}}, \end{aligned}$$

which gives the expected upper bound with $\phi(t) = \sqrt{4(1+\delta)t}$.

Remark 3.2. Once we have confirmed that algebraic and analytic conditions hold for the Markov process generated by the Laplace-Beltrami operator Δ , it should be noticed that our CZ conditions are again implied by the classical ones. Arguing as in Remarks 1.8 and 2.5, we see that the Ricci curvature assumption allows us to ignore our size kernel conditions. Next, it is straightforward to check that the boundedness condition reduces in this case to standard L_2 -boundedness. Finally, our discussion in section 2.4 shows that our smoothness kernel condition is guaranteed under the classical Hörmander kernel condition. Note in addition that our conditions also hold in the row case. In particular, classical CZOs in (M, g) become algebraic CZOs. Moreover, the gaussian upper estimate (3.2) indicates that in (M, g) with the heat semigroup \mathcal{S}_Δ we have $L_p^\circ(M) = L_p(M, g)$ for $1 < p < \infty$.

By the discussion above, we have all the ingredients to apply Theorem 2.1 and Corollary 2.2. Let us illustrate it for the Riesz transforms on (M, g) . Consider the Riemannian gradient $\nabla = (\partial_1, \partial_2, \dots, \partial_n)$ on (M, g) . The Riesz transform on (M, g) is formally defined by

$$R = (R_j) = \nabla(-\Delta)^{-\frac{1}{2}} \quad \text{with} \quad R_j = \partial_j(-\Delta)^{-\frac{1}{2}}.$$

Then we may recover Bakry's theorem [1] using our algebraic approach. Indeed integration by parts gives $\|\nabla f\|_2 = \|\Delta^{\frac{1}{2}} f\|_2$ which implies L_2 -boundedness of Riesz transforms. Moreover, the Hörmander condition follows from [9, 41].

Corollary 3.3. *Let (M, g) be a complete n -dimensional Riemannian manifold with non-negative Ricci curvature. Then for all $1 < p < \infty$, there exists a constant $C_p > 0$ such that*

$$\|R_j f\|_{L_p(M, g)} \leq C_p \|f\|_{L_p(M, g)} \quad \text{for all} \quad 1 \leq j \leq n.$$

3.2. The Gaussian measure

Now we study the Ornstein-Uhlenbeck semigroup on the Euclidean space equipped with its Gaussian measure, which satisfies as well our regularity assumption. We shall first construct a Markov metric for it. Then we shall prove that our algebraic/analytic and Calderón-Zygmund conditions hold for the standard CZOs in this setting. The generator of the Ornstein-Uhlenbeck semigroup $O = (O_t)_{t \geq 0}$ is the operator

$$L = \frac{\Delta}{2} - x \cdot \nabla$$

on (\mathbb{R}^n, μ) with $d\mu(y) = \exp(-|y|^2)dy$. We have

$$\begin{aligned} O_t f(x) &= \frac{1}{(\pi - \pi e^{-2t})^{\frac{n}{2}}} \int_{\mathbb{R}^n} \exp\left(-\frac{|e^{-t}x - y|^2}{1 - e^{-2t}}\right) f(y) dy \\ &= \frac{1}{(\pi - \pi e^{-2t})^{\frac{n}{2}}} \int_{\mathbb{R}^n} \exp\left(|x|^2 - \frac{|e^t x - y|^2}{e^{2t} - 1}\right) f(y) d\mu(y). \end{aligned}$$

First, we establish a useful lemma showing that the local behavior —i.e. for small values of t — of the semigroup type BMO norm for the Ornstein-Uhlenbeck semigroup determines it completely.

Lemma 3.4. *Given $\delta > 0$, there exists $C_\delta > 0$ such that*

$$\sup_{t \geq 0} \|O_t |f|^2 - |O_t f|^2\|_\infty \leq C_\delta \sup_{t < \delta} \|O_t |f|^2 - |O_t f|^2\|_\infty.$$

Proof. It is easy to check that

$$O_t f(x) = H_{v(t)} f(e^{-t}x), \quad (3.3)$$

for $v(t) = \frac{1}{4}(1 - e^{-2t})$ and the heat semigroup $H_t = \exp(t\Delta)$. Given $t > 0$ and $f \in L_\infty(\mathbb{R}^n)$, let $F(s) = H_s |H_{t-s} f|^2$ for $0 \leq s \leq t$. According to the definition of H_t , we obtain the following identity

$$\begin{aligned} \partial_s F &= (\partial_s H_s) |H_{t-s} f|^2 + H_s [(\partial_s H_{t-s} f)^* (H_{t-s} f)] + H_s [(H_{t-s} f)^* (\partial_s H_{t-s} f)] \\ &= \Delta H_s |H_{t-s} f|^2 - H_s [(\Delta H_{t-s} f)^* (H_{t-s} f)] - H_s [(H_{t-s} f)^* (\Delta H_{t-s} f)] \\ &= H_s [\Delta |H_{t-s} f|^2 - (\Delta H_{t-s} f)^* (H_{t-s} f) - (H_{t-s} f)^* (\Delta H_{t-s} f)] \\ &= 2H_s |\nabla H_{t-s} f|^2. \end{aligned}$$

Kadison-Schwarz inequality gives for $0 < u < s$

$$H_u |\nabla H_{t-u} f|^2 = H_u |H_{s-u} \nabla_x H_{t-s} f|^2 \leq H_s |\nabla_x H_{t-s} f|^2$$

which implies that $\partial_s F$ is increasing and F is convex. Rearranging the inequality $F(s) \leq \frac{1}{2}(F(0) + F(2s))$, we get $H_{2t}|f|^2 - |H_{2t}f|^2 \leq 2H_t(H_t|f|^2 - |H_t f|^2)$ for any $t \geq 0$. Then, the L_∞ contractivity of H_t gives

$$\|H_{2^k t}|f|^2 - |H_{2^k t}f|^2\|_\infty \leq 2^k \|H_t|f|^2 - |H_t f|^2\|_\infty. \quad (3.4)$$

On the other hand, choosing k_δ such that $2^{k_\delta} v(\delta) \geq \frac{1}{4}$ and applying (3.3) and (3.4)

$$\begin{aligned} \sup_{t \geq 0} \|O_t|f|^2 - |O_t f|^2\|_\infty &= \sup_{t < \frac{1}{4}} \|H_t|f|^2 - |H_t f|^2\|_\infty \\ &\leq \sup_{t < 2^{k_\delta} v(\delta)} \|H_t|f|^2 - |H_t f|^2\|_\infty \\ &\leq 2^{k_\delta} \sup_{t < \delta} \|O_t|f|^2 - |O_t f|^2\|_\infty. \quad \square \end{aligned}$$

By the lemma above, it suffices to construct a Markov metric for $(O_t)_{t \geq 0}$ with $0 < 2t < \frac{1}{18}$. Let $v = \sqrt{e^{2t} - 1}$ and consider the following family of balls and coronas in the gaussian space for $(j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+$

$$\Sigma_{j,t}(x) = B(e^t x, \sqrt{j}v) \quad \text{and} \quad \Omega_{j,t}(x) = \Sigma_{j,t}(x) \setminus \Sigma_{j-1,t}(x).$$

Let $j_0 = j_0(x, t)$ be the smallest possible integer j satisfying that $0 \in \Sigma_{j,t}(x)$.

The case $n = 1$. If $1 \leq j < j_0$, let

$$\begin{aligned} D_{j,t}^-(x) &= \left\{ y \in \Omega_{j,t}(x) : e^t|x| - \sqrt{j}v \leq |y| \leq e^t|x| - \sqrt{j-1}v \right\}, \\ D_{j,t}^+(x) &= \left\{ y \in \Omega_{j,t}(x) : e^t|x| + \sqrt{j-1}v \leq |y| \leq e^t|x| + \sqrt{j}v \right\}. \end{aligned}$$

Then, $D_{j,t}^-(x) \cup D_{j,t}^+(x) = \Omega_{j,t}(x)$ and we get

$$O_t f(x) \lesssim \frac{1}{v} \left(\sum_{\substack{\varepsilon = \pm \\ 1 \leq j < j_0}} \exp(|x|^2 - j) \int_{D_{j,t}^\varepsilon(x)} f d\mu + \sum_{j \geq j_0} \exp(|x|^2 - j) \int_{\Sigma_{j,t}(x)} f d\mu \right) \quad (3.5)$$

for any positive $f \in L_\infty(\mathbb{R}, \mu)$. The above estimate indicates the natural candidates for the cpu maps $R_{j,t}$ and $\sigma_{j,t} \in L_\infty(\mathbb{R}, \mu)$. When $1 \leq j < j_0$ and $\varepsilon = \pm$, we define

$$R_{j,t,\varepsilon} f(x) = \frac{1}{\mu(D_{j,t}^\varepsilon(x))} \int_{D_{j,t}^\varepsilon(x)} f d\mu \quad \text{and} \quad \sigma_{j,t,\varepsilon}^2(x) = \frac{1}{v} \exp(|x|^2 - j) \mu(D_{j,t}^\varepsilon(x)).$$

Note here we need an extra index for $R_{j,t}$'s when $j < j_0$. This is consistent with the assumptions (i) and (ii) in our definition of Markov metric, since we only need the index-set of $R_{j,t}$'s to be countable. On the other hand, if $j \geq j_0$, we set

$$R_{j,t}f(x) = \frac{1}{\mu(\Sigma_{j,t}(x))} \int_{\Sigma_{j,t}(x)} f d\mu \quad \text{and} \quad \sigma_{j,t}^2(x) = \frac{1}{v} \exp(|x|^2 - j) \mu(\Sigma_{j,t}(x)).$$

Since $\ker A_\infty$ the fixed-point subspace of O is the space of constant functions, so $R_{j,t,\varepsilon}$ and $R_{j,t}$ fix all the functions in $\ker A_\infty$. In order to find $\gamma_{j,t,\varepsilon}$ and $\gamma_{j,t}$ satisfying the metric integrability condition, we need to estimate $\mu(D_{j,t}^\varepsilon(x))$ and $\mu(\Sigma_{j,t}(x))$ respectively. Since the density function μ is monotone on $D_{j,t}^\varepsilon(x)$ we get

$$\mu(D_{j,t}^-(x)) = \int_{D_{j,t}^-(x)} e^{-y^2} dy \leq \exp(-|e^t x| - \sqrt{j}v)^2 \frac{v}{\sqrt{j}}, \quad (3.6)$$

$$\mu(D_{j,t}^+(x)) = \int_{D_{j,t}^+(x)} e^{-y^2} dy \leq \exp(-|e^t x| + \sqrt{j-1}v)^2 \frac{v}{\sqrt{j}}. \quad (3.7)$$

When $j \geq j_0$ we use the trivial estimate

$$\mu(\Sigma_{j,t}(x)) = \int_{\Sigma_{j,t}(x)} e^{-y^2} dy \leq 2\sqrt{j}v.$$

Combining the estimates obtained above, we deduce for $1 \leq j < j_0$

$$\sigma_{j,t,-}^2 \leq \frac{1}{\sqrt{j}} \exp(-|v|x - e^t \sqrt{j}|^2) \quad \text{and} \quad \sigma_{j,t,+}^2 \leq \frac{1}{\sqrt{j}} \exp(-|v|x + e^t \sqrt{j-1}|^2).$$

When $j \geq j_0$, we have $e^t |x| \leq \sqrt{j}v$ and $|x|^2 \leq jv^2 e^{-2t} < j/4$. Therefore

$$\sigma_{j,t}^2 \leq 2\sqrt{j} \exp(|x|^2 - j) < 2\sqrt{j} \exp\left(-\frac{3}{4}j\right).$$

Now we are ready to choose the optimal γ 's for the metric integrability condition in the definition of Markov metric. We respectively define for $1 \leq j < j_0$ and $j \geq j_0$

$$\gamma_{j,t,\varepsilon}^2(x) = \exp\left(\frac{|v|x| + e^t \sqrt{j-1}|^2}{2}\right) \quad \text{and} \quad \gamma_{j,t}^2(x) = \frac{1}{\sqrt{j}} \exp\left(\frac{j}{4}\right).$$

Then it turns out that

$$\begin{aligned} \sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} \sum_{1 \leq j < j_0} \sigma_{j,t,-}^2(x) \gamma_{j,t,-}^2(x) &\leq \sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} \sum_{1 \leq j < j_0} \frac{1}{\sqrt{j}} \exp\left(-\frac{|6v|x| - e^t \sqrt{j}|^2}{4}\right) \\ &\leq \sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} 2 \int_{\mathbb{R}} \exp\left(-\frac{|6v|x| - e^t u|^2}{4}\right) du < \infty. \end{aligned}$$

$$\begin{aligned} \sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} \sum_{1 \leq j < j_0} \sigma_{j,t,+}^2(x) \gamma_{j,t,+}^2(x) &\leq \sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} \sum_{1 \leq j < j_0} \frac{1}{\sqrt{j}} \exp\left(-\frac{|v|x| + e^t \sqrt{j-1}|^2}{2}\right) \\ &\leq 1 + \sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} 2 \int_{\mathbb{R}} \exp\left(-\frac{|v|x| + e^t u|^2}{2}\right) du < \infty. \end{aligned}$$

On the other hand, it is clear that

$$\sup_{\substack{x \in \mathbb{R} \\ 0 < t < \frac{1}{36}}} \sum_{j \geq j_0} \sigma_{j,t}^2(x) \gamma_{j,t}^2(x) \leq \sum_{j \geq j_0} 2 \exp\left(-\frac{j}{2}\right) < \infty.$$

We have constructed a Markov metric for the Ornstein-Uhlenbeck semigroup.

Let us now verify the analytic conditions, since the algebraic conditions are trivial by commutativity. As we mentioned in Subsection 2.2, the first condition is an easy consequence of Jensen's inequality for the Gaussian measure. By the definition of $R_{j,t,\varepsilon}$ and $R_{j,t}$, we get $q_{j,t,\varepsilon}(x, y) = \chi_{D_{j,t,\varepsilon}^\varepsilon(x)}(y)$ and $q_{j,t}(x, y) = \chi_{\Sigma_{j,t}(x)}(y)$. Thus, it remains to find proper $a_{j,t,\varepsilon}$ and $a_{j,t}$ to make sure that

$$\mu(D_{j,t}^\varepsilon(x)) \lesssim \int_{\mathbb{R}} a_{j,t,\varepsilon}^2(x, y) d\mu(y) \lesssim \gamma_{j,t,\varepsilon}^2(x) \mu(D_{j,t}^\varepsilon(x))$$

and similarly for the pairs $(\Sigma_{j,t}(x), a_{j,t}^2(x, y))$. When $j < j_0$ we consider the functions $a_{j,t,\varepsilon}(x, y) = \chi_{2D_{j,t,\varepsilon}^\varepsilon(x)}(y)$, so the lower estimates are trivial. Denote by $c_{j,t,\varepsilon}(x)$ the center of $D_{j,t}^\varepsilon(x)$. Let $\beta = v(\sqrt{j} - \sqrt{j-1})$. Arguing as in (3.6), we get

$$\begin{aligned} \exp\left(-\left|c_{j,t,\varepsilon}(x) + \frac{\beta}{2}\right|^2\right) &\leq \frac{\mu(D_{j,t}^\varepsilon(x))}{\beta} \leq \exp\left(-\left|c_{j,t,\varepsilon}(x) - \frac{\beta}{2}\right|^2\right), \\ \exp\left(-\left|c_{j,t,\varepsilon}(x) + \beta\right|^2\right) &\leq \frac{\mu(2D_{j,t}^\varepsilon(x))}{2\beta} \leq \exp\left(-\left|c_{j,t,\varepsilon}(x) - \beta\right|^2\right). \end{aligned}$$

Since $c_{j,t,\pm}(x) = e^t|x| \pm \frac{1}{2}v(\sqrt{j} + \sqrt{j-1})$, this implies that

$$\frac{\mu(2D_{j,t}^\varepsilon(x))}{\mu(D_{j,t}^\varepsilon(x))} \leq 2 \exp\left(3|c_{j,t,\varepsilon}(x)|\beta - \frac{3}{4}\beta^2\right) \leq 2\gamma_{j,t,\varepsilon}^2(x).$$

The estimate for $j \geq j_0$ is easier. Take $a_{j,t}(x, y) = \chi_{2\Sigma_{j,t}(x)}(y)$. Note

$$\begin{aligned} 2\sqrt{j}v \exp(-|2\sqrt{j}v|^2) &\leq \mu(\Sigma_{j,t}(x)) \leq 2\sqrt{j}v, \\ 4\sqrt{j}v \exp(-|3\sqrt{j}v|^2) &\leq \mu(2\Sigma_{j,t}(x)) \leq 4\sqrt{j}v. \end{aligned}$$

Then, since $0 < 2t < \frac{1}{18}$, we get

$$\frac{\mu(2\Sigma_{j,t}(x))}{\mu(\Sigma_{j,t}(x))} \leq 2 \exp(4jv^2) \leq 2 \exp\left(\frac{j}{4}\right) = 2\gamma_{j,t}^2(x).$$

Our choice for q_t and a_t correspond to the balls $\Sigma_{1,t}(x)$ and $2\Sigma_{1,t}(x)$.

The case $n > 1$. The argument is similar, so we just point out the necessary modifications. Since $v(\sqrt{j} - \sqrt{j-1}) < vj^{-\frac{1}{2}}$, we may pick $c_n j^{n-1}$ balls $D_{j,t}^s(x)$ for $1 \leq s \leq c_n j^{n-1}$ with radius $\frac{v}{2\sqrt{j}}$, centered on the sphere

$$\left\{ y : |y - e^t x| = \frac{v(\sqrt{j} + \sqrt{j-1})}{2} \right\}$$

such that

$$\Omega_{j,t}(x) \subset \bigcup_{s \geq 1} D_{j,t}^s(x)$$

and each $D_{j,t}^{s_0}(x)$ overlaps with at most c'_n other balls $D_{j,t}^s(x)$. Then, if $f \geq 0$

$$O_t f(x) \lesssim \frac{1}{v^n} \left(\sum_{\substack{j < j_0 \\ 1 \leq s \leq c_n j^{n-1}}} \exp(|x|^2 - j) \int_{D_{j,t}^s(x)} f d\mu + \sum_{j \geq j_0} \exp(|x|^2 - j) \int_{\Sigma_{j,t}(x)} f d\mu \right).$$

Then, we may consider the following Markov metric

$$\begin{aligned} & \left(\sigma_{j,t,s}^2(x), \gamma_{j,t,s}^2(x), R_{j,t,s} f(x) \right) \\ &= \left(\frac{\exp(|x|^2 - j)}{v^n} \mu(D_{j,t}^s(x)), j^{-\frac{n-1}{2}} \exp\left(\frac{|v|x| + e^t \sqrt{j-1}|^2}{2}\right), \int_{D_{j,t}^s(x)} f d\mu \right) \end{aligned}$$

for $j < j_0$ and $1 \leq s \leq c_n j^{n-1}$. When $j \geq j_0$, we set

$$\left(\sigma_{j,t}^2(x), \gamma_{j,t}^2(x), R_{j,t} f(x) \right) = \left(\frac{\exp(|x|^2 - j)}{v^n} \mu(\Sigma_{j,t}(x)), j^{-\frac{n}{2}} \exp\left(\frac{j}{4}\right), \int_{\Sigma_{j,t}(x)} f d\mu \right).$$

The analytic conditions hold under the same choices we made for $n = 1$.

Corollary 3.5. *Let $\mathcal{O} = (O_t)_{t \geq 0}$ be the Ornstein-Uhlenbeck semigroup and T be a singular integral operator defined on $L_\infty(\mathbb{R}^n, d\mu)$ with kernel k . More precisely, we have the kernel representation*

$$Tf(x) = \int_{\mathbb{R}^n} k(x, y) f(y) d\mu(y) \quad \text{for } x \notin \text{supp } f.$$

Suppose T is bounded on $L_2(\mathbb{R}^n, \mu)$ and it satisfies

$$\sup_B \sup_{\text{ball}} \sup_{\substack{z \in B \\ j \geq 1}} \int_{2^{j+1}B \setminus 2^j B} |k(z, y)| d\mu(y) < \infty, \quad (3.8)$$

$$\sup_B \sup_{\text{ball}} \sup_{z_1, z_2 \in B} \int_{(5B)^c} |k(z_1, y) - k(z_2, y)| d\mu(y) < \infty. \quad (3.9)$$

Then T is a bounded map from $L_\infty(\mathbb{R}^n, \mu)$ to the semigroup $\text{BMO}_\mathcal{O}$ space.

Proof. It suffices to prove that our CZ conditions hold. The row and column boundedness conditions reduce to L_2 -boundedness. Let M_t be the averaging map in $L_\infty(\mathbb{R}^n, \mu)$ over the ball $\Sigma_{1,t}(x)$. Given $1 \leq j < j_0$, define $A_{j,t,s}(x, \cdot)$ as the characteristic function over the ball $\Sigma_{j+1,t}(x)$. Then $A_{j,t,s} \leq \chi_{2^r D_{j,t}^s(x) \cap \Sigma_{1,t}(x)}$ with $r = [2 \log_2(j+1)] + 3$, where $[\cdot]$ stands for the integer part. Therefore, applying (3.8), we have

$$\begin{aligned} \sup_{z \in \Sigma_{1,t}(x)} \int_{\Sigma_{j+1,t}(x) \setminus 2\Sigma_{1,t}(x)} |k(z, y)| d\mu(y) &\lesssim r \lesssim \gamma_{j,t,s}^2(x), \\ \sup_{z \in D_{j,t}^s(x)} \int_{\Sigma_{j+1,t}(x) \setminus 2D_{j,t}^s(x)} |k(z, y)| d\mu(y) &\lesssim r \lesssim \gamma_{j,t,s}^2(x). \end{aligned}$$

For $j \geq j_0$, let $A_{j,t}(x, \cdot) = \chi_{\Sigma_{4j,t}(x)} \leq \chi_{2^u \Sigma_{1,t}(x)}$ with $u = [\log_2(2\sqrt{j})] + 1$. Applying (3.8) as above, we see that T satisfies our size conditions. Moreover, (3.9) implies our smoothness conditions as in the Euclidean-Lebesgue setting, Section 2.4. \square

Remark 3.6. Since the Gaussian measure is non-doubling, the term $R_{j,t}f - M_t f$ in the Markov metric BMO space $\text{BMO}_\mathcal{O}$ is essential to characterise the changes of the mean values of the function f . This explains the relevance of the size kernel condition in the Calderón-Zygmund theory for the gaussian measure.

4. Applications II — Noncommutative spaces

In this section we apply our algebraic approach to study Calderón-Zygmund operators in flag von Neumann algebras which originally motivated us and include matrix algebras, quantum Euclidean spaces and quantum groups. We start by reconstructing and refining the semicommutative theory, which deals with tensor and crossed products with metric measure spaces.

4.1. Operator-valued theory

Let (Ω, μ) be a doubling metric space —as in Remark 1.8— and consider a Markov semigroup $S_t : L_\infty(\Omega) \rightarrow L_\infty(\Omega)$. Let \mathcal{M} be a semifinite von Neumann algebra with a n.s.f. trace τ . Then we call the semigroup $\mathcal{S} = (S_t \otimes id_\mathcal{M})_{t \geq 0}$ a *semicommutative Markov*

semigroup. Consider the algebra of essentially bounded functions $f : \Omega \rightarrow \mathcal{M}$ equipped with the trace

$$\varphi(f) = \int_{\Omega} \tau(f(y)) d\mu(y).$$

Its weak-* closure $\mathcal{R} = L_{\infty}(\Omega) \bar{\otimes} \mathcal{M}$ is a von Neumann algebra. Assume that there exists a Markov metric $\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j,t) \in \mathbb{Z}_+ \times \mathbb{R}_+\}$ associated to the original Markov semigroup on $L_{\infty}(\Omega)$. Let $q_{j,t}(x, y) = \chi_{\Omega_{j,t}^x}(y)$ stand for the projections determined by \mathcal{Q} via (2.1). We assume in addition that \mathcal{Q} satisfies the metric/measure growth condition

$$\frac{\mu(\Sigma_{j,t}^x)}{\mu(\Omega_{j,t}^x)} \leq \gamma_{j,t}(x) \quad (4.1)$$

by choosing $a_{j,t}(x, y) = \chi_{\Sigma_{j,t}^x}(y)$. The remaining algebraic and analytic conditions trivially hold in this case. Indeed, the algebraic conditions follow by commutativity and analytic conditions just require to pick the right averaging maps according to Jensen's inequality, as explained in (2.4). Note that \mathcal{Q} satisfies an operator-valued generalization of the Hilbert module majorization in the line of Remark 1.7. Thus \mathcal{Q} extends to a Markov metric in \mathcal{R} by tensorizing with $id_{\mathcal{M}}$ and $\mathbf{1}_{\mathcal{M}}$ respectively.

Our goal is to study CZO's formally given by

$$Tf(x) = \int_{\Omega} k(x, y) (f(y)) d\mu(y) \quad \text{with} \quad \begin{cases} f : \Omega \rightarrow \mathcal{M}_1 \text{ and } x \notin \text{supp}_{\Omega} f, \\ k(x, y) \in \mathcal{L}(L_0(\mathcal{M}_1), L_0(\mathcal{M}_2)). \end{cases}$$

That is, $k(x, y)$ is linear from τ_1 -measurable to τ_2 -measurable operators. If we set $\mathcal{R}_j = L_{\infty}(\Omega) \bar{\otimes} \mathcal{M}_j$, we should emphasize that $L_p(\mathcal{R}_j) = L_p(\Omega; L_p(\mathcal{M}_j))$. In particular, this framework does not fall in the vector-valued theory because we take values in different Banach spaces for different values of p , see [49] for further explanations. This class of operators is inspired by two distinguished examples with $\mathcal{M}_1 = \mathcal{M} = \mathcal{M}_2$:

- Operator-valued case

$$Tf(x) = \int_{\Omega} k_{ov}(x, y) \cdot f(y) d\mu(y).$$

- Noncommutative model

$$Tf(x) = \int_{\Omega} (id_{\mathcal{M}} \otimes \tau) [k_{nc}(x, y) \cdot (\mathbf{1}_{\mathcal{M}} \otimes f(y))] d\mu(y).$$

In the first case, the kernel takes values in \mathcal{M} or even in the complex field and acts on $f(y)$ by left multiplication $k(x, y)(f(y)) = k_{ov}(x, y) \cdot f(y)$. It is the canonical map

when $L_p(\mathcal{R})$ is regarded as the Bochner space $L_p(\Omega; L_p(\mathcal{M}))$. On the contrary if we simply think of $L_p(\mathcal{R})$ as a noncommutative L_p space, a *natural* CZO should be an integral map with respect to the full trace $\varphi = \int_{\Omega} \otimes \tau$ and the kernel should be a $\varphi \otimes \varphi$ -measurable operator $k : \Omega \times \Omega \rightarrow \mathcal{M} \bar{\otimes} \mathcal{M}$. The noncommutative model provides the resulting integral formula. Note that this model also falls in our general framework by taking $k(x, y)(f(y)) = (id_{\mathcal{M}} \otimes \tau)[k_{nc}(x, y) \cdot (\mathbf{1}_{\mathcal{M}} \otimes f(y))]$.

Theorem 4.1. *Let $\mathcal{S} = (S_t)_{t \geq 0}$ be a Markov semigroup on (Ω, μ) which admits a Markov metric $\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+\}$ satisfying the above assumptions. Let $a_{j,t}(x, y) = \chi_{\Sigma_{j,t}^x(y)}$ be the projections determined by \mathcal{Q} via (4.1). Consider the CZO formally given by*

$$Tf(x) = \int_{\Omega} k(x, y)(f(y)) d\mu(y).$$

Then, T maps $L_{\infty}(\mathcal{R}_1)$ to $BMO_{\mathcal{S}}^c(\mathcal{R}_2)$ provided the conditions below hold

i) L_2^c -boundedness condition,

$$\left\| \left(\int_{\Omega} |Tf|^2 d\mu \right)^{\frac{1}{2}} \right\|_{\mathcal{M}_2} \lesssim \left\| \left(\int_{\Omega} |f|^2 d\mu \right)^{\frac{1}{2}} \right\|_{\mathcal{M}_1}.$$

ii) Smoothness condition for the kernel,

$$\int_{\Omega \setminus \Sigma_{j,t}^x} \|(k(y_1, z) - k(y_2, z))(f(z))\|_{\mathcal{M}_2} d\mu(z) \lesssim \|f\|_{\mathcal{R}_1}$$

uniformly in $j \geq 1$, $t > 0$, $x \in \Omega$ and $y_1, y_2 \in \Omega_{j,t}^x$.

Proof. The proof follows from Theorem 2.1. Since the underlying space (Ω, μ) is a doubling metric space, the size kernel condition is unnecessary. Thus, it remains to check the L_2^c -boundedness condition and the kernel smoothness condition. Consider $\mathcal{A}_{\mathcal{M}} = L_{\infty}(\mathcal{R}_1) \cap (L_2^c(\Omega) \bar{\otimes} L_{\infty}(\mathcal{M}))$ (see [28] for the definition of Hilbert space valued L_{∞} spaces), $\mathcal{N}_{\pi} = L_{\infty}(\Omega \times \Omega) \bar{\otimes} \mathcal{M}_1$, $\mathcal{N}_{\rho} = L_{\infty}(\Omega \times \Omega) \bar{\otimes} \mathcal{M}_2$, $\omega(\varphi)(x, y) = \varphi(y)$ for $\varphi \in L_{\infty}(\Omega)$ and $(\pi_2, \rho_2) = (\omega \otimes id_{\mathcal{M}_1}, \omega \otimes id_{\mathcal{M}_2})$. Let $\hat{T} = id_{\Omega} \otimes T$, $\Phi_{j,t}$ be the averaging map over $\Omega_{j,t}^x \times \Omega_{j,t}^x$ and $\Delta = \delta \otimes id_{\mathcal{M}_2}$ with $\delta\varphi(x, y) = \varphi(x) - \varphi(y)$. Then condition i) yields the L_2^c -boundedness condition. It is also easy to see that condition ii) implies our kernel smoothness condition. Thus, the result follows from Theorem 2.1. \square

Remark 4.2. We continue with a few comments:

A) When $\mathcal{M}_1 = \mathcal{M}_2 = \mathcal{M}$ and the kernel $k(x, y)(f(y)) = k(x, y) \cdot f(y)$ acts by left multiplication, the boundedness condition i) becomes equivalent to the usual L_2 boundedness. Indeed, using that $\mathcal{M} \subset \mathcal{B}(L_2(\mathcal{M}))$ we obtain

$$\begin{aligned}
& \left\| \left(\int_{\mathbb{R}^n} |Tf(y)|^2 dy \right)^{\frac{1}{2}} \right\|_{\mathcal{M}} \\
&= \sup_{\|h\|_2 \leq 1} \left(\int_{\mathbb{R}^n} \langle h, |Tf(y)|^2 h \rangle dy \right)^{\frac{1}{2}} = \sup_{\|h\| \leq 1} \|(Tf)(\mathbf{1}_{\mathbb{R}^n} \otimes h)\|_{L_2(\mathcal{R})} \\
&= \sup_{\|h\| \leq 1} \|T(f(\mathbf{1}_{\mathbb{R}^n} \otimes h))\|_{L_2(\mathcal{R})} \leq \|T\|_{\mathcal{B}(L_2(\mathcal{R}))} \left\| \left(\int_{\mathbb{R}^n} |f(y)|^2 dy \right)^{\frac{1}{2}} \right\|_{\mathcal{M}}.
\end{aligned}$$

B) We have used so far semigroup type BMO's. When (Ω, μ) comes equipped with a doubling metric, we may replace it by other standard (equivalent) forms of BMO, as pointed in Remark 1.9. By well-known arguments [49], our kernel smoothness condition reduces to

$$\sup_{R>0} \operatorname{ess\,sup}_{y_1, y_2 \in B_R} \left\| \int_{(B_{\lambda R})^c} (k(y_1, z) - k(y_2, z))(f(z)) dz \right\|_{\mathcal{M}} \lesssim \|f\|_{\mathcal{R}}, \quad (\text{Sm}_{\lambda})$$

for $\lambda > 1$. The classical Hörmander condition

$$\operatorname{ess\,sup}_{y_1, y_2 \in \mathbb{R}^n} \int_{d(y_1, z) > \lambda d(y_1, y_2)} \|k(y_1, z) - k(y_2, z)\|_{\mathcal{M}} dz < \infty, \quad (\text{Hr}_{\lambda})$$

satisfies $(\text{Hr}_{\lambda}) \Rightarrow (\text{Sm}_{2\lambda+1})$. In fact, an even weaker condition suffices

$$\sup_{R>0} \left\| \int_{B_R \times B_R} \left| \int_{(B_{\lambda R})^c} (k(y_1, z) - k(y_2, z))f(z) dz \right|^2 dy_1 dy_2 \right\|_{\mathcal{M}} \lesssim \|f\|_{\mathcal{R}}^2.$$

C) We recall that $L_{\infty}(\mathcal{R}) \rightarrow \text{BMO}_{\mathcal{S}}$ boundedness requires that $T^{\dagger}f = T(f^*)^*$ satisfies the same assumptions as T . If $k(x, y) \in \mathcal{M}$ is given by left multiplication the only effect in T^{\dagger} is that $k(x, y)$ is replaced by $k(x, y)^*$ and now operates by right multiplication. This left/right condition was formulated in [49] in terms of \mathcal{M} -bimodular maps. Moreover, a counterexample was constructed to show that the bimodularity is indeed essential. It is also quite interesting to note that in the ‘noncommutative model’ we have

$$\int_{\mathbb{R}^n} (id_{\mathcal{M}} \otimes \tau) [k(x, y) \cdot (\mathbf{1}_{\mathcal{M}} \otimes f(y))] dy = \int_{\mathbb{R}^n} (id_{\mathcal{M}} \otimes \tau) [(\mathbf{1}_{\mathcal{M}} \otimes f(y)) \cdot k(x, y)] dy$$

by traciality and this pathology does not occur. Finally, the L_p boundedness is guaranteed for $2 < p < \infty$ since the classical heat semigroup has a regular Markov metric and $J_p = id_{L_p(\mathbb{R}^n)}$ in this case. As for $1 < p < 2$, it suffices to take adjoints which leads to Hörmander smoothness in the second variable

$$\operatorname{ess\,sup}_{z_1, z_2 \in \mathbb{R}^n} \int_{|y - z_1| > \lambda |z_1 - z_2|} \|k(y, z_1) - k(y, z_2)\|_{\mathcal{M}} dy < \infty.$$

Of course, this is still consistent with the classical CZ theory $\mathcal{M} = \mathbb{C}$.

We now study the $L_\infty \rightarrow \text{BMO}$ boundedness of twisted CZO's on homogeneous spaces. Given a discrete group G with left regular representation $\lambda : G \rightarrow \mathcal{B}(\ell_2(G))$ let $\mathcal{L}(G)$ denote its group von Neumann algebra. Let (\mathcal{M}, τ) with $\mathcal{M} \subset \mathcal{B}(\mathcal{H})$ be a noncommutative probability space and $\alpha : G \rightarrow \text{Aut}(\mathcal{M})$ be a trace preserving action. Consider two $*$ -representations

$$\begin{aligned} \rho : \mathcal{M} \ni f &\mapsto \sum_{h \in G} \alpha_{h^{-1}}(f) \otimes e_{h,h} \in \mathcal{M} \bar{\otimes} \mathcal{B}(\ell_2(G)), \\ \Lambda : G \ni g &\mapsto \sum_{h \in G} \mathbf{1}_{\mathcal{M}} \otimes e_{gh,h} \in \mathcal{M} \bar{\otimes} \mathcal{B}(\ell_2(G)), \end{aligned}$$

where $e_{g,h}$ is the matrix unit for $\mathcal{B}(\ell_2(G))$. Now we define the crossed product algebra $\mathcal{M} \rtimes_\alpha G$ as the weak operator closure in $\mathcal{M} \bar{\otimes} \mathcal{B}(\ell_2(G))$ of the $*$ -algebra generated by $\rho(\mathcal{M})$ and $\Lambda(G)$. A generic element of $\mathcal{M} \rtimes_\alpha G$ can be formally written as $\sum_{g \in G} f_g \rtimes_\alpha \lambda(g)$ with $f_g \in \mathcal{M}$. With this convention, we may embed the crossed product algebra $\mathcal{M} \rtimes_\alpha G$ into $\mathcal{M} \bar{\otimes} \mathcal{B}(\ell_2(G))$ via the map $j = \rho \rtimes \Lambda$. Indeed, we have

$$\begin{aligned} j\left(\sum_{g \in G} f_g \rtimes_\alpha \lambda(g)\right) &= \sum_{g \in G} \rho(f_g) \Lambda(g) \\ &= \sum_{g \in G} \left(\sum_{h, h' \in G} (\alpha_{h^{-1}}(f_g) \otimes e_{h,h}) (\mathbf{1}_{\mathcal{M}} \otimes e_{gh',h'}) \right) \\ &= \sum_{g \in G} \left(\sum_{h \in G} \alpha_{h^{-1}}(f_g) \otimes e_{h,g^{-1}h} \right) \\ &= \sum_{g \in G} \left(\sum_{h \in G} \alpha_{(gh)^{-1}}(f_g) \otimes e_{gh,h} \right). \end{aligned}$$

Since the action α will be fixed, we relax the terminology and write $\sum_{g \in G} f_g \lambda(g)$ instead of $\sum_{g \in G} f_g \rtimes_\alpha \lambda(g)$. We say that a Markov semigroup $\mathcal{S} = (S_t)_{t \geq 0}$ in \mathcal{M} is G -equivariant if

$$\alpha_g S_t = S_t \alpha_g \quad \text{for } (t, g) \in \mathbb{R}_+ \times G.$$

If \mathcal{S} is a G -equivariant Markov semigroup on \mathcal{M} , let $\mathcal{S}_\rtimes = (S_t \rtimes id_G)_{t \geq 0}$ and $\mathcal{S}_\otimes = (S_t \otimes id_{\mathcal{B}(\ell_2(G))})_{t \geq 0}$ denote the crossed/tensor product amplification of our semigroup on $\mathcal{M} \rtimes G$ and $\mathcal{M} \bar{\otimes} \mathcal{B}(\ell_2(G))$ respectively. Note that \mathcal{S}_\rtimes is Markovian due to the G -equivariance of \mathcal{S} . In the following result, our CZO's are of the form

$$Tf(x) = \int_{\Omega} k(x, y)(f(y)) d\mu(y)$$

for all $f \in (\mathcal{R}_1, \varphi_1)$, where $(\mathcal{R}_j, \varphi_j) = L_\infty(\Omega, \mu) \bar{\otimes} (\mathcal{M}_j, \tau_j)$ and $k(x, y) : \mathcal{M}_1 \rightarrow \mathcal{M}_2$. In other words, we keep the same terminology as for Theorem 4.1. We shall also use the notation

$$\widehat{\mathcal{M}}_j = \mathcal{M}_j \bar{\otimes} \mathcal{B}(\ell_2(G)) \quad \text{and} \quad \widehat{\mathcal{R}}_j = \mathcal{R}_j \bar{\otimes} \mathcal{B}(\ell_2(G)).$$

Corollary 4.3. *Let $G \curvearrowright L_\infty(\Omega, \mu)$ be an action α which is implemented by a measure preserving transformation β , so that $\alpha_g f(x) = f(\beta_{g^{-1}}x)$. Let $\mathcal{S} = (S_t)_{t \geq 0}$ be a G -equivariant Markov semigroup on (Ω, μ) which admits a Markov metric $\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+\}$ satisfying the assumptions above. Let us consider a family of CZO's formally given by*

$$T_g f(x) = \int_{\Omega} k_g(x, y)(f(y)) d\mu(y) \quad \text{for } g \in G.$$

Then, $\sum_g f_g \lambda(g) \mapsto \sum_g T_g(f_g) \lambda(g)$ is bounded $\mathcal{R}_1 \rtimes G \rightarrow \text{BMO}_{\mathcal{S}_\times}^c(\mathcal{R}_2 \rtimes G)$ if

i) L_2^c -boundedness condition,

$$\left\| \left(\int_{\Omega} |(T_{gh^{-1}}) \bullet \xi|^2 d\mu \right)^{\frac{1}{2}} \right\|_{\widehat{\mathcal{M}}_2} \lesssim \left\| \left(\int_{\Omega} |\xi|^2 d\mu \right)^{\frac{1}{2}} \right\|_{\widehat{\mathcal{M}}_1},$$

where \bullet stands for the generalized Schur product of matrices. In other words, the CZO $T_{gh^{-1}}$ only acts on the (g, h) -th entry of ξ for each $g, h \in G$.

ii) Smoothness condition for the kernel,

$$\int_{\Omega} \left\| (K(y_1, z) - K(y_2, z)) \bullet (\xi(z)(1 - a_{j,t}(x, z))) \right\|_{\widehat{\mathcal{M}}_2} d\mu(z) \lesssim \|\xi\|_{\widehat{\mathcal{R}}_1},$$

uniformly on $j \geq 1$, $t > 0$, $x \in \Omega$ and $y_1, y_2 \in \Omega_{j,t}^x$. Here, the CZ kernel $K(y, z) = \sum_{g,h} k_{gh^{-1}}(\beta_g y, \beta_g z) \otimes e_{g,h}$ acts once more as a Schur multiplier.

Proof. Letting $\xi = \sum_{g,h} a_{g,h} \otimes e_{g,h} \in \widehat{\mathcal{R}}_1$, we define the map

$$\begin{aligned} \Phi : \widehat{\mathcal{R}}_1 &\rightarrow \text{BMO}_{\mathcal{S}_\otimes}^c(\widehat{\mathcal{R}}_2), \\ \Phi(\xi)(x) &= \sum_{g,h} \alpha_{g^{-1}} \int_{\Omega} k_{gh^{-1}}(x, y)(a_{g,h}(\beta_{g^{-1}}(y))) d\mu(y) \otimes e_{g,h}. \end{aligned}$$

By the definition of j , it is easy to check that

$$j\left(\sum_g T_g(f_g) \lambda(g)\right) = \Phi\left(j\left(\sum_g f_g \lambda(g)\right)\right).$$

Since \mathcal{S} is G -equivariant, according to [31, Lemma 2.1], we have

$$\|g\|_{\text{BMO}_{\mathcal{S} \rtimes \mathcal{G}}^c(\mathcal{R}_2 \rtimes G)} = \sup_{t \geq 0} \left\| \left(S_{\otimes, t} |j(g)|^2 - |S_{\otimes, t} j(g)|^2 \right)^{\frac{1}{2}} \right\|_{\widehat{\mathcal{R}}_2}.$$

Therefore, it suffices to show that Φ is $\widehat{\mathcal{R}}_1 \rightarrow \text{BMO}_{\mathcal{S}_{\otimes}}^c(\widehat{\mathcal{R}}_2)$ bounded. We find

$$\Phi(\xi)(x) = \int_{\Omega} K(x, y)(\xi(y)) d\mu(y).$$

Thus, we may regard Φ as a semicommutative CZO and apply Theorem 4.1 where \mathcal{M}_j is replaced by $\widehat{\mathcal{M}}_j$. Since $\Phi(\xi) = \sum_{g, h} (\alpha_{g^{-1}}) \bullet (T_{gh^{-1}}) \bullet (\alpha_g) \bullet \xi$ and β is measure preserving, we immediately find that the L_2^c -boundedness assumption implies that the map

$$\Phi : L_2^c(\Omega) \bar{\otimes} \widehat{\mathcal{M}}_1 \rightarrow L_2^c(\Omega) \bar{\otimes} \widehat{\mathcal{M}}_2$$

is bounded. Moreover, the smoothness condition matches that of Theorem 4.1. \square

Remark 4.4. Our work so far yields sufficient conditions for the $L_{\infty} \rightarrow \text{BMO}$ boundedness of $T \rtimes id_G$ in more general settings. In particular, if $T_g = T$ and $\alpha_g T = T \alpha_g$ for all $g \in G$, then we find for any T fulfilling the assumption of Theorem 4.1, $T \rtimes id_G : \mathcal{R}_1 \rtimes G \rightarrow \text{BMO}_{\mathcal{S} \rtimes \mathcal{G}}^c(\mathcal{R}_2 \rtimes G)$ is bounded. This gives an example where the hypothesis of Lemma 2.2 in [31] is satisfied.

4.2. Matrix algebras

In this paragraph, we introduce a Markov metric for the matrix algebra $\mathcal{B}(\ell_2)$. The triangular truncation plays the noncommutative form of the Hilbert transform on $\mathcal{B}(\ell_2)$. We shall reprove the L_p -boundedness of the triangular truncation for $1 < p < \infty$ and a new $\text{BMO} \rightarrow \text{BMO}$ estimate by means of this Markov metric and our algebraic approach. Consider the $*$ -homomorphism $u : \mathcal{B}(\ell_2) \rightarrow L_{\infty}(\mathbb{R}) \bar{\otimes} \mathcal{B}(\ell_2)$ determined by

$$u(e_{mk}) = e^{2\pi i(m-k)} \cdot e_{mk}.$$

Given $A = \sum_{m, k} a_{mk} e_{mk}$, define the semigroup

$$S_t(A) = \sum_{m, k} e^{-t|m-k|^2} a_{mk} e_{mk}.$$

It is not difficult to see that it defines a Markov semigroup of convolution type. In fact, u is a corepresentation of $L_{\infty}(\mathbb{R})$ (equipped with its natural comultiplication map $\Delta f(x, y) = f(x + y)$) in $\mathcal{B}(\ell_2)$ and it turns out that $\mathcal{S} = (S_t)_{t \geq 0}$ is the regular semigroup associated to the heat semigroup on \mathbb{R} by transference

$$u \circ S_t = (H_t \otimes id_{\mathcal{B}(\ell_2)}) \circ u.$$

Define the cpu map $R_{j,t}$ on $\mathcal{B}(\ell_2)$ by $u \circ R_{j,t} = (\tilde{R}_{j,t} \otimes id_{\mathcal{B}(\ell_2)}) \circ u$, where $\tilde{R}_{j,t}f(x)$ denotes the average of $f \in L_\infty(\mathbb{R})$ over the interval $B_{\sqrt{4jt}}(x)$. Now, given a matrix $A = \sum_{m,k} a_{mk} e_{mk}$ we find

$$\begin{aligned} u \circ R_{j,t}(A)(x) &= \int_{B_{\sqrt{4jt}}(x)} u(A)(y) dy \\ &= \int_{B_{\sqrt{4jt}}(x)} \sum_{m,k} e^{2\pi i(m-k)y} a_{mk} e_{mk} dy \\ &= \sum_{m,k} \frac{\sin(4\sqrt{jt}\pi(m-k))}{4\sqrt{jt}\pi(m-k)} e^{2\pi i(m-k)x} a_{mk} e_{mk}. \end{aligned}$$

Thus, we find the following identity

$$R_{j,t}(A) = \sum_{m,k} \frac{\sin(4\sqrt{jt}\pi(m-k))}{4\sqrt{jt}\pi(m-k)} a_{mk} e_{mk}.$$

Taking $\sigma_{j,t}^2 = 2e\sqrt{j/\pi}e^{-j}\mathbf{1}_{\mathcal{B}(\ell_2)}$ and $\gamma_{j,t}^2 = \sqrt{j}\mathbf{1}_{\mathcal{B}(\ell_2)}$, we obtain a Markov metric in $\mathcal{B}(\ell_2)$. Indeed, the metric integrability condition holds trivially, as for the Hilbert module majorization it reduces to prove that $B_1 \leq B_2$ with

$$\begin{aligned} B_1 &= u(\langle \xi, \xi \rangle_{S_t}) = \langle u \otimes u(\xi), u \otimes u(\xi) \rangle_{H_t \otimes id_{\mathcal{B}(\ell_2)}}, \\ B_2 &= \sum_j \sigma_{j,t}^2 u(\langle \xi, \xi \rangle_{R_{j,t}}) = \sum_j \sigma_{j,t}^2 \langle u \otimes u(\xi), u \otimes u(\xi) \rangle_{\tilde{R}_{j,t} \otimes id_{\mathcal{B}(\ell_2)}}. \end{aligned}$$

In other words, it suffices to note that the canonical Markov metric in \mathbb{R} —which recovers the Euclidean metric, as proved in Paragraph 1.6— admits a matrix-valued extension, as it was justified in Remark 1.7. Let us now consider the triangular truncation

$$\triangle(A) = \sum_{m>k} a_{mk} e_{mk}.$$

Corollary 4.5. *We have*

$$\|\triangle(A)\|_{\text{BMO}_S} \lesssim \|A\|_{\text{BMO}_S}.$$

In particular, given $1 < p < \infty$ we obtain $\|\triangle(A)\|_{S_p} \lesssim \frac{p^2}{p-1} \|A\|_{S_p}$.

Proof. Recall that

$$u \circ \Delta = (L \otimes id_{\mathcal{B}(\ell_2)}) \circ u,$$

for $L = \frac{1}{2}(id + iH)$ and $\widehat{Hf}(\xi) = -i\operatorname{sgn}(\xi)\widehat{f}(\xi)$, the Hilbert transform in the real line. We may also regard $u : \mathcal{B}(\ell_2) \rightarrow L_\infty(\mathbb{T}) \bar{\otimes} \mathcal{B}(\ell_2)$ as a corepresentation of \mathbb{T} instead of \mathbb{R} and the above identity holds replacing H by the Hilbert transform in the torus. In this case, u becomes a trace preserving $*$ -homomorphism and the well-known S_p inequalities for Δ reduce to the boundedness of the Hilbert transform in $L_p(\mathbb{T}; S_p(\ell_2))$, which is also well-known and follows in passing from the semicommutative theory in the previous paragraph. Alternatively, the second assertion follows from the first one by interpolation and duality. According to Remark 1.9, to prove the first assertion it suffices to show that the map

$$T = i(id_{B(\ell_2)} - 2\Delta)$$

is $\operatorname{BMO} \rightarrow \operatorname{BMO}$ bounded for the semigroup BMO space which is associated to the transferred Poisson semigroup P_t on $B(\ell_2)$ given by

$$P_t : (a_{ij}) \mapsto (e^{-t|i-j|}a_{ij}).$$

Given $A = (a_{jk})_{j,k}$ in $B(\ell_2)$ then

$$\begin{aligned} |A|^2 &= A^*A = \left(\sum_k \overline{a_{ki}}a_{kj} \right)_{i,j} \\ T(A) &= i \left(\operatorname{sgn}(k-j)a_{jk} \right)_{j,k} \quad (TA)^* = i \left(\operatorname{sgn}(k-j)\overline{a_{kj}} \right)_{j,k} \end{aligned}$$

Then $(P_t|A|^2 - |P_tA|^2)_{ij} = \sum_k (e^{-t|i-j|} - e^{-t|k-j|}e^{-t|i-k|})\overline{a_{ki}}a_{kj}$ and

$$\begin{aligned} &(P_t|T(A)|^2 - |P_tT(A)|^2)_{ij} \\ &= \sum_k (e^{-t|i-j|} - e^{-t|k-j|}e^{-t|i-k|})\operatorname{sgn}(k-i)\operatorname{sgn}(k-j)\overline{a_{ki}}a_{kj}. \end{aligned}$$

Since $\operatorname{sgn}(k-i)\operatorname{sgn}(k-j) \neq 1$ iff $e^{-t|i-k|}e^{-t|k-j|} = e^{-t|i-j|}$, we get

$$P_t|A|^2 - |P_tA|^2 = P_t|T(A)|^2 - |P_tT(A)|^2.$$

The last identity implies that T is an isometry on the Poisson BMO space. \square

4.3. Quantum Euclidean spaces

Given an integer $n \geq 1$, fix an anti-symmetric \mathbb{R} -valued $n \times n$ matrix Θ . We define A_Θ as the universal C^* -algebra generated by a family $u_1(s), u_2(s), \dots, u_n(s)$ of one-parameter unitary groups in $s \in \mathbb{R}^n$ which are strongly continuous and satisfy the following Θ -commutation relations

$$u_j(s)u_k(t) = e^{2\pi i \Theta_{jk}st} u_k(t)u_j(s).$$

If $\Theta = 0$, by Stone's theorem we can take $u_j(s) = \exp(2\pi i s \langle e_j, \cdot \rangle)$ and A_Θ is the space of bounded continuous functions on \mathbb{R}^n . In general, given $\xi \in \mathbb{R}^n$, define the unitaries $\lambda_\Theta(\xi) = u_1(\xi_1)u_2(\xi_2) \cdots u_n(\xi_n)$. Let E_Θ be the closure in A_Θ of $\lambda_\Theta(L_1(\mathbb{R}^n))$ with

$$f = \int_{\mathbb{R}^n} \check{f}_\Theta(\xi) \lambda_\Theta(\xi) d\xi.$$

If $\Theta = 0$, $E_\Theta = C_0(\mathbb{R}^n)$. Define

$$\tau_\Theta(f) = \tau_\Theta \left(\int_{\mathbb{R}^n} \check{f}_\Theta(\xi) \lambda_\Theta(\xi) d\xi \right) = \check{f}_\Theta(0)$$

for $\check{f}_\Theta : \mathbb{R}^n \rightarrow \mathbb{C}$ integrable and smooth. τ_Θ extends to a n.s.f. trace on E_Θ . Let $\mathcal{R}_\Theta = A''_\Theta = E''_\Theta$ be the von Neumann algebra generated by E_Θ in the GNS representation of τ_Θ . Note that if $\Theta = 0$, $\mathcal{R}_\Theta = L_\infty(\mathbb{R}^n)$. In general we call \mathcal{R}_Θ a quantum Euclidean space. There are two maps which play important roles while doing analysis over quantum Euclidean spaces. The first one is the corepresentation map $\sigma_\Theta : \mathcal{R}_\Theta \rightarrow L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta$, given by $\lambda_\Theta(\xi) \mapsto \exp_\xi \otimes \lambda_\Theta(\xi)$ where \exp_ξ stands for the Fourier character $\exp(2\pi i \langle \xi, \cdot \rangle)$. Note that σ_Θ is a normal injective $*$ -homomorphism. The second map is $\pi_\Theta : \exp_\xi \mapsto \lambda_\Theta(\xi) \otimes \lambda_\Theta(\xi)^*$, which extends to a normal $*$ -homomorphism from $L_\infty(\mathbb{R}^n)$ to $\mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta^{\text{op}}$, where $\mathcal{R}_\Theta^{\text{op}}$ is the opposite algebra of \mathcal{R}_Θ , which is obtained by preserving the linear and adjoint structures but reversing the product. We refer the readers to [21] for more detailed information of quantum Euclidean spaces and these two maps.

BMO and Markov metric. Our first goal is to construct a natural Markov metric for quantum Euclidean spaces. Let us recall the heat semigroup on \mathbb{R}^n acting on $\varphi : \mathbb{R}^n \rightarrow \mathbb{C}$ admits the following form

$$H_t \varphi(x) = \int_{\mathbb{R}^n} \hat{\varphi}(\xi) e^{-t|\xi|^2} \exp_\xi(x) d\xi.$$

This induces a regular semigroup on \mathcal{R}_Θ determined by

$$\sigma_\Theta \circ S_{\Theta,t} = (H_t \otimes id_{\mathcal{R}_\Theta}) \circ \sigma_\Theta.$$

$S_{\Theta,t}$ gives a Markov semigroup on \mathcal{R}_{Θ} which formally acts as

$$S_{\Theta,t}(f) = \int_{\mathbb{R}^n} \check{f}_{\Theta}(\xi) e^{-t|\xi|^2} \lambda_{\Theta}(\xi) d\xi. \quad (4.2)$$

The corresponding semigroup column BMO norm is given by

$$\begin{aligned} \|f\|_{\text{BMO}_c(\mathcal{R}_{\Theta})} &= \sup_{t>0} \left\| \left(S_{\Theta,t}(|f|^2) - |S_{\Theta,t}(f)|^2 \right)^{\frac{1}{2}} \right\|_{\mathcal{R}_{\Theta}} \\ &\approx \sup_{B \text{ ball in } \mathbb{R}^n} \left\| \left(\int_B |\sigma_{\Theta}(f) - \sigma_{\Theta}(f)_B|^2 d\mu \right)^{\frac{1}{2}} \right\|_{\mathcal{R}_{\Theta}} = \|\sigma_{\Theta}(f)\|_{\text{BMO}_c(\mathbb{R}^n; \mathcal{R}_{\Theta})}. \end{aligned}$$

The null space of the above seminorm is the preimage of the space of constant operator-valued functions under σ_{Θ} . According to Remark 1.7, the semicommutative extension $H_t \otimes id_{\mathcal{R}_{\Theta}}$ of the heat semigroup, together with the extension of the corresponding Markov metric from Paragraph 2.4 still satisfies the Hilbert module majorization

$$\langle \xi, \xi \rangle_{H_t \otimes id_{\mathcal{R}_{\Theta}}} \leq \sum_{j \geq 1} \sigma_{j,t}^* \langle \xi, \xi \rangle_{R_{j,t} \otimes id_{\mathcal{R}_{\Theta}}} \sigma_{j,t} \quad (4.3)$$

as well as the integrability condition, where $\sigma_{j,t}^2 \equiv 2e\sqrt{j^n/\pi}e^{-j}$, $\gamma_{j,t}^2 \equiv j^{\frac{n}{2}}$ and $R_{j,t}f(x)$ is the average of f over $B_{\sqrt{4jt}}(x)$. Then we can easily produce a Markov metric on \mathcal{R}_{Θ} . Let $B_{j,t}$ be the Euclidean ball in \mathbb{R}^n centered at the origin with radius $\sqrt{4jt}$ and consider the projections $q_{j,t} = \chi_{B_{j,t}} \otimes \mathbf{1}_{\mathcal{R}_{\Theta}}$. Define the cpu maps

$$R_{\Theta,j,t}(f) = \frac{1}{|B_{j,t}|} \int_{B_{j,t}} \sigma_{\Theta}(f)(x) dx = \frac{1}{|B_{j,t}|} \int_{\mathbb{R}^n} \hat{\chi}_{B_{j,t}}(\xi) \check{f}_{\Theta}(\xi) \lambda_{\Theta}(\xi) d\xi.$$

It is easy to check that

$$\sigma_{\Theta} \circ R_{\Theta,j,t} = (R_{j,t} \otimes id_{\mathcal{R}_{\Theta}}) \circ \sigma_{\Theta}, \quad (4.4)$$

and $R_{\Theta,j,t}f = f$ for any f in the null space of $\text{BMO}_c(\mathcal{R}_{\Theta})$ -seminorm. The Hilbert module majorization

$$\langle \xi, \xi \rangle_{S_{\Theta,t}} \leq \sum_{j \geq 1} \sigma_{j,t}^* \langle \xi, \xi \rangle_{R_{\Theta,j,t}} \sigma_{j,t}$$

for $\xi \in \mathcal{R}_{\Theta} \bar{\otimes}_{S_{\Theta,t}} \mathcal{R}_{\Theta}$ is equivalent to the same inequality after composing with the $*$ -homomorphism σ_{Θ} , which follows in turn by the intertwining identities (4.2) and (4.4), together with the majorization (4.3). Therefore, we obtain a Markov metric on \mathcal{R}_{Θ} associated to S_{Θ}

Table 1
Algebraic skeleton for \mathcal{R}_Θ .

Generic algebraic objects	Quantum Euclidean spaces
\mathcal{M}	\mathcal{R}_Θ
\mathcal{N}_ρ	$L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta$ $\rho_1 = \mathbf{1} \otimes \cdot, \rho_2 = \sigma_\Theta$ $E_\rho = \text{Lebesgue integral}$
\mathcal{N}_π	$\mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta^{\text{op}}$ $\pi_1 = \mathbf{1} \otimes \cdot, \pi_2 = \cdot \otimes \mathbf{1}$ $E_\pi = \tau_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}}$
\mathcal{N}_σ	$L_\infty(\mathbb{R}^n) \bar{\otimes} L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta$ $\mathcal{N}_\sigma = (\delta \otimes id_{\mathcal{R}_\Theta})(\mathcal{N}_\rho)$ $\delta\varphi(x, y) = \varphi(x) - \varphi(y)$

$$\mathcal{Q}_\Theta = \{(R_{\Theta,j,t}, \sigma_{j,t}, \gamma_{j,t}) \mid (j, t) \in \mathbb{Z}_+ \times \mathbb{R}_+\}.$$

The algebraic structure. We start with the kernel representation of our CZOs over the (fully noncommutative) von Neumann algebra \mathcal{R}_Θ . Given a kernel k affiliated to $\mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta^{\text{op}}$, the linear map associated to it is formally given by

$$T_k f = (id_{\mathcal{R}_\Theta} \otimes \tau_\Theta)(k(\mathbf{1}_{\mathcal{R}_\Theta} \otimes f)) = (id_{\mathcal{R}_\Theta} \otimes \tau_\Theta)((\mathbf{1}_{\mathcal{R}_\Theta} \otimes f)k).$$

The reader is referred to [21] for more details. Our goal is to provide sufficient conditions for the $L_\infty \rightarrow \text{BMO}$ boundedness of T_k . Consider the $*$ -homomorphism $\sigma_\Theta: \mathcal{R}_\Theta \rightarrow L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta$. In the case of quantum Euclidean spaces, we need the full algebraic skeleton introduced in Section 2. In Table 1 there is a little dictionary to identify the main objects. Next, note that

$$\sigma_\Theta \circ T_k(f) = (id_{\mathbb{R}^n} \otimes id_{\mathcal{R}_\Theta} \otimes \tau_\Theta)(k_\sigma(\mathbf{1}_{\mathbb{R}^n} \otimes \mathbf{1}_{\mathcal{R}_\Theta} \otimes f)),$$

where $k_\sigma = (\sigma_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})(k)$. Denote $\sigma_\Theta \circ T_k$ by T_{k_σ} . Define

$$\widehat{T}_k: \mathcal{R}_\Theta \bar{\otimes} \mathcal{R}_\Theta \ni f \otimes a \mapsto T_{k_\sigma}(f)(\mathbf{1}_{\mathbb{R}^n} \otimes a) \in L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta.$$

Then it is clear that the compatibility condition (2.3) holds since $\widehat{T}_k \circ \pi_2 = \sigma_\Theta \circ T_k$.

Lemma 4.6. *If T_k is bounded on $L_2(\mathcal{R}_\Theta)$, then*

$$\|\widehat{T}_k: L_2^c(\mathcal{R}_\Theta) \bar{\otimes} \mathcal{R}_\Theta \rightarrow L_2^c(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta\| \leq \|T_k: L_2(\mathcal{R}_\Theta) \rightarrow L_2(\mathcal{R}_\Theta)\|,$$

the rigorous definition of $L_2^c(\mathcal{R}_\Theta) \bar{\otimes} \mathcal{R}_\Theta$ can be found in [28].

Proof. We need to introduce two maps:

$$j_\Theta: L_2(\mathbb{R}^n) \ni \int_{\mathbb{R}^n} \varphi(\xi) \exp_\xi d\xi \mapsto \int_{\mathbb{R}^n} \varphi(\xi) \lambda_\Theta(\xi) d\xi \in L_2(\mathcal{R}_\Theta),$$

$$W : L_2^c(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta \ni \int_{\mathbb{R}^n} \exp_\xi \otimes a(\xi) d\xi \mapsto \int_{\mathbb{R}^n} \exp_\xi \otimes \lambda_\Theta(\xi) a(\xi) d\xi \in L_2^c(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta.$$

It is straightforward to show that W extends to an isometry. Moreover, j_Θ is also an L_2 -isometry, we refer the reader to [21, Section 1.3.2] for the proof. Observe that

$$\begin{aligned} \sigma_\Theta(f)(1_{\mathbb{R}^n} \otimes a) &= \int_{\mathbb{R}^n} \check{f}_\Theta(\xi) \exp_\xi \otimes \lambda_\Theta(\xi) a d\xi \\ &= W\left(\int_{\mathbb{R}^n} \check{f}_\Theta(\xi) \exp_\xi \otimes a d\xi\right) = W \circ (j_\Theta^* \otimes id_{\mathcal{R}_\Theta})(f \otimes a). \end{aligned}$$

Letting $f = T_k g$, we get

$$\widehat{T}_k(g \otimes a) = W(j_\Theta^* T_k \otimes id_{\mathcal{R}_\Theta})(g \otimes a).$$

The properties of the maps j_Θ and W readily imply the assertion. \square

Now let us introduce a weak-* dense subalgebra of \mathcal{R}_Θ , which is the analogue of the classical Schwartz class. Let $\mathcal{S}(\mathbb{R}^n)$ denote the classical Schwartz class in the Euclidean space \mathbb{R}^n and define

$$\mathcal{S}_\Theta = \{f \in \mathcal{R}_\Theta : \check{f}_\Theta \in \mathcal{S}(\mathbb{R}^n)\}.$$

Using the unitary map

$$j_\Theta : \mathcal{S}(\mathbb{R}^n) \rightarrow \mathcal{S}_\Theta$$

as defined in the proof of Lemma 4.6, \mathcal{S}_Θ carries a natural locally convex topology. Its topological dual \mathcal{S}'_Θ is the quantum space of tempered distributions on \mathcal{R}_Θ . Consider a continuous linear operator $T \in \mathcal{L}(\mathcal{S}_\Theta, \mathcal{S}'_\Theta)$, we have $j_\Theta^* T j_\Theta \in \mathcal{L}(\mathcal{S}(\mathbb{R}^n), \mathcal{S}(\mathbb{R}^n)')$. By a result of Schwartz, there exists a unique kernel $K \in \mathcal{S}'(\mathbb{R}^{2n}) = (\mathcal{S}(\mathbb{R}^n) \otimes_\pi \mathcal{S}(\mathbb{R}^n))'$ such that T admits the kernel $k = (j_\Theta \otimes j_\Theta)(K) \in (\mathcal{S}_\Theta \otimes_\pi \mathcal{S}_\Theta)'$. Actually, the kernel representations T_k satisfying the Calderón-Zygmund type conditions in the following theorem belong to $\mathcal{L}(\mathcal{S}_\Theta, \mathcal{S}'_\Theta)$. It provides sufficient conditions for the $L_\infty(\mathcal{R}_\Theta) \rightarrow \text{BMO}_c(\mathcal{R}_\Theta)$ boundedness of CZO operators associated to kernels in $(\mathcal{S}_\Theta \otimes_\pi \mathcal{S}_\Theta)'$. We shall use the quantum analogue of the bands around the diagonal

$$a_B = \pi_\Theta(\chi_{5B}) = \int_{\mathbb{R}} \widehat{\chi}_{5B}(\xi) \lambda_\Theta(\xi) \otimes \lambda_\Theta(\xi)^* d\xi.$$

Theorem 4.7. *Let $T_k \in \mathcal{L}(\mathcal{S}_\Theta, \mathcal{S}'_\Theta)$ and assume*

i) *Cancellation*

$$\|T_k : L_2(\mathcal{R}_\Theta) \rightarrow L_2(\mathcal{R}_\Theta)\| < \infty.$$

ii) *For any $f \in \mathcal{R}_\Theta$ and any Euclidean ball B centered at the origin*

$$\int_{B \times B} |\Sigma_{\Theta,k,f,B}(y_1) - \Sigma_{\Theta,k,f,B}(y_2)|^2 dy_1 dy_2 \lesssim \|f\|_{\mathcal{R}_\Theta}^2,$$

$$\text{where } \Sigma_{\Theta,k,f,B} = (id_{\mathbb{R}^n} \otimes id_{\mathcal{R}_\Theta} \otimes \tau_\Theta) [k_\sigma(\mathbf{1}_{\mathbb{R}^n} \otimes \mathbf{1}_{\mathcal{R}_\Theta} \otimes f)(\mathbf{1}_{\mathbb{R}^n} \otimes a_B^\perp)].$$

Then, the Calderón-Zygmund operator T_k is bounded from $L_\infty(\mathcal{R}_\Theta)$ to $BMO_c(\mathcal{R}_\Theta)$.

Proof. By Theorem 1.6, it suffices to prove

$$T_k : L_\infty(\mathcal{R}_\Theta) \rightarrow BMO_{\mathcal{Q}_\Theta}^c.$$

Arguing as in Paragraph 1.6, the Markov metric BMO norm takes the simpler form

$$\|f\|_{BMO_{\mathcal{Q}_\Theta}^c} = \sup_{t>0} \sup_{j \geq 1} \left\| \left(\gamma_{j,t}^{-1} [R_{\Theta,j,t}(|f|^2) - |R_{\Theta,j,t}(f)|^2] \gamma_{j,t}^{-1} \right)^{\frac{1}{2}} \right\|_{\mathcal{R}_\Theta}.$$

In other words, the extra term in the definition of BMO is dominated by the above expression as in (1.5). As noticed in Remark 2.5, the size kernel condition is then superfluous. This also reduces the analytic conditions and the smooth kernel conditions to be checked. In summary, according to the proof of Theorem 2.1, the assertion will follow if we can justify:

C0) Initial condition

$$T_k : \mathcal{A}_\Theta \rightarrow \mathcal{R}_\Theta \quad \text{for } \mathcal{A}_\Theta \subset \mathcal{R}_\Theta \text{ weak-* dense.}$$

Al1) \mathcal{Q}_Θ -monotonicity of E_ρ

$$E_\rho(q_{j,t}|\xi|^2 q_{j,t}) \leq E_\rho(|\xi|^2).$$

Al2) Right modularity of \widehat{T}_k

$$\widehat{T}_k(\eta\pi_1(b)) = \widehat{T}_k(\eta)\rho_1(b).$$

An1) Mean differences

$$\widehat{R}_{\Theta,j,t}(\xi^*\xi) - \widehat{R}_{\Theta,j,t}(\xi)^* \widehat{R}_{\Theta,j,t}(\xi) \leq \Phi_{j,t}(\delta(\xi)^* \delta(\xi)) \text{ for some cpu } \Phi_{j,t}.$$

An2) Metric/measure growth

$$\mathbf{1} \leq \pi_1 \rho_1^{-1} \mathbb{E}_\rho(q_{j,t})^{-\frac{1}{2}} \mathbb{E}_\pi(a_{j,t}^* a_{j,t}) \pi_1 \rho_1^{-1} \mathbb{E}_\rho(q_{j,t})^{-\frac{1}{2}} \lesssim \pi_1 \rho_1^{-1} (\gamma_{j,t}^2).$$

CZ1) L_2^c -boundedness condition

$$\widehat{T}_k : L_\infty^c(\mathcal{N}_\pi; \mathbb{E}_\pi) \rightarrow L_\infty^c(\mathcal{N}_\rho; \mathbb{E}_\rho).$$

CZ2) Kernel smoothness condition

$$\Phi_{j,t} \left(\left| \delta(\widehat{T}_k(\pi_2(f)(\mathbf{1} - a_{j,t}))) \right|^2 \right) \lesssim \gamma_{j,t}^2 \|f\|_\infty^2.$$

The initial conditions trivially hold for good kernels $k \in \mathcal{S}_\Theta \otimes_{\text{alg}} \mathcal{S}_\Theta$ for $\mathcal{A}_\Theta = \mathcal{S}_\Theta$. In [21] it was required to extend the main result from this class of kernels to general ones in $\mathcal{S}'_{\Theta \oplus \Theta}$, by reproving certain auxiliary results in the context of distributions. In our case, this is much simpler. Indeed, when dealing with general kernels, we just note that $T_k(f) \in L_2(\mathcal{R}_\Theta)$ for all $f \in \mathcal{S}_\Theta$ by assumption. Given the form of $R_{\Theta,j,t}$, it trivially follows that $R_{\Theta,j,t}(|T_k f|^2)$ and $R_{\Theta,j,t} T_k f$ are well-defined operators in $L_1(\mathcal{R}_\Theta)$ and $L_2(\mathcal{R}_\Theta)$ respectively. In particular, the proof of Theorem 2.1 follows exactly as it was written there under this more flexible assumption. Therefore, the initial condition can be relaxed to the condition

$$T_k : \mathcal{S}_\Theta \rightarrow L_2(\mathcal{R}_\Theta).$$

In fact, according to [21, Proposition 2.17], every algebraic column CZO is normal. Thus, it suffices —as we did in Theorem 2.1— to justify that $T_k : \mathcal{S}_\Theta \rightarrow \text{BMO}_{\mathcal{Q}_\Theta}^c$ is bounded, as we shall do by justifying the remaining conditions.

Al1 holds trivially since $q_{j,t} = \chi_{B_{j,t}} \otimes \mathbf{1}_{\mathcal{R}_\Theta}$ lives in the center of \mathcal{N}_ρ . On the other hand, according to the definition of ρ_1, π_1 from Table 1, the algebraic condition Al2 can be rewritten as follows

$$\widehat{T}_k(\eta(\mathbf{1}_{\mathcal{R}_\Theta} \otimes b)) = \widehat{T}_k(\eta)(\mathbf{1}_{\mathbb{R}^n} \otimes b).$$

This is clear from the definition of \widehat{T}_k . Next, condition An1 reads as

$$\int_{B_{j,t}} |\xi|^2 d\mu - \left| \int_{B_{j,t}} \xi d\mu \right|^2 \leq \int_{B_{j,t} \times B_{j,t}} |\xi(x) - \xi(y)|^2 d\mu(x) d\mu(y)$$

for \mathcal{R}_Θ -valued functions, when $\Phi_{j,t}$ is chosen to be the average over $B_{j,t} \times B_{j,t}$. As in (2.4), this is a consequence of the operator-valued Jensen's inequality. Next recalling that $a_{j,t} = \pi_\Theta(\chi_{5B_{j,t}})$, condition An2 takes the form

$$|B_{j,t}| \mathbf{1}_{\mathcal{R}_\Theta} \leq (\tau_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})(\pi_\Theta(\chi_{5B_{j,t}})) \lesssim j^{\frac{n}{2}} |B_{j,t}| \mathbf{1}_{\mathcal{R}_\Theta}.$$

To verify it we note that

$$(\tau_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})(\pi_\Theta(\varphi)) = (\tau_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})\left(\int_{\mathbb{R}^n} \widehat{\varphi}(\xi) \lambda_\Theta(\xi) \otimes \lambda_\Theta(\xi)^* d\xi\right) = \widehat{\varphi}(0) \mathbf{1}_{\mathcal{R}_\Theta}.$$

Then we get $(\tau_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})(\pi_\Theta(\chi_{5B_{j,t}})) = 5|B_{j,t}| \mathbf{1}_{\mathcal{R}_\Theta}$. Condition CZ1 reduces to our L_2 -boundedness assumption by Lemma 4.6. Finally, the smoothness condition ii) in the statement readily implies condition CZ2 for all values of j, t . \square

The smoothness condition in Theorem 4.7 is of Hörmander type, while the one in the main result of [21] is a gradient condition. As expected, we shall show that our condition in this paper is more flexible than that of [21, Theorem 2.6]. We use \bullet for the product in $\mathcal{M} \bar{\otimes} \mathcal{M}^{\text{op}}$, so that

$$(a \otimes b) \bullet (a' \otimes b') = (aa') \otimes (b'b).$$

The quantum analogue of the metric is defined by

$$d_\Theta = \pi_\Theta(| \cdot |)$$

for the Euclidean norm $| \cdot |$. Moreover, we also introduce the Θ -deformation of the free gradient. Let $\mathcal{L}(\mathbb{F}_n)$ denote the group von Neumann algebra associated to the free group with n generators \mathbb{F}_n . It is well-known from (say) [63] that $\mathcal{L}(\mathbb{F}_n)$ is generated by n semicircular random variables s_1, s_2, \dots, s_n . Note that there exist derivations ∂_Θ^j in \mathcal{S}_Θ which are determined by

$$\partial_\Theta^j(\lambda_\Theta(\xi)) = 2\pi i \xi_j \lambda_\Theta(\xi)$$

for $1 \leq j \leq n$. Define the Θ -deformed free gradient as

$$\nabla_\Theta = \sum_{j=1}^n s_j \otimes \partial_\Theta^j : \mathcal{S}_\Theta \rightarrow \mathcal{L}(\mathbb{F}_n) \bar{\otimes} \mathcal{R}_\Theta.$$

If ∇ denotes the free gradient for $\Theta = 0$, it is easy to check that

$$\begin{aligned} (id_{\mathcal{L}(\mathbb{F}_n)} \otimes \sigma_\Theta) \circ \nabla_\Theta &= \sum_{j=1}^n s_j \otimes (\sigma_\Theta \circ \partial_\Theta^j) \\ &= \sum_{j=1}^n s_j \otimes (\partial_j \circ \sigma_\Theta) = (\nabla \otimes id_{\mathcal{R}_\Theta}) \circ \sigma_\Theta. \end{aligned} \tag{4.5}$$

For the convenience of the reader, we cite Theorem 2.6 from [21] below.

Theorem 4.8. Let $T_k \in \mathcal{L}(\mathcal{S}_\Theta, \mathcal{S}'_\Theta)$ and assume:

i) *Cancellation*

$$\|T_k : L_2(\mathcal{R}_\Theta) \rightarrow L_2(\mathcal{R}_\Theta)\| \leq A_1.$$

ii) *Gradient condition.* There exists

$$\alpha < \frac{n}{2} < \beta < \frac{n}{2} + 1$$

satisfying the gradient conditions below for $\rho = \alpha, \beta$

$$\left| d_\Theta^\rho \bullet (\nabla_\Theta \otimes id_{\mathcal{R}_\Theta^{\text{op}}})(k) \bullet d_\Theta^{n+1-\rho} \right| \leq A_2.$$

Then, we find the following $L_\infty \rightarrow \text{BMO}_c$ estimate

$$\|T_k : L_\infty(\mathcal{R}_\Theta) \rightarrow \text{BMO}_c(\mathcal{R}_\Theta)\| \leq C_n(\alpha, \beta)(A_1 + A_2).$$

To simplify notation, we shall write in what follows Σ for $\Sigma_{\Theta, k, f, B}$. According to the semicommutative Poincaré type inequality introduced in [21, Proposition 1.6] we obtain

$$\left\| \int_{B \times B} |\delta(\Sigma)|^2 d\mu \times \mu \right\|_{\mathcal{R}_\Theta} \leq 16R^2 \left\| (\mathbf{1} \otimes \chi_B \otimes \mathbf{1})(\nabla \otimes id_{\mathcal{R}_\Theta})(\Sigma) \right\|_{\mathcal{L}(\mathbb{F}_n) \bar{\otimes} L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta}^2$$

for R = radius of B . By (4.5), we may rewrite

$$\begin{aligned} & (\mathbf{1} \otimes \chi_B \otimes \mathbf{1})(\nabla \otimes id_{\mathcal{R}_\Theta})(\Sigma) \\ &= (id^{\otimes 3} \otimes \tau_\Theta) \left((\mathbf{1} \otimes \chi_B \otimes \mathbf{1}^{\otimes 2})(\nabla \otimes id^{\otimes 2})(k_\sigma)(\mathbf{1}^{\otimes 3} \otimes f)(\mathbf{1}^{\otimes 2} \otimes a_{j,t}^\perp) \right) \\ &= (id^{\otimes 3} \otimes \tau_\Theta) \left((\mathbf{1} \otimes \chi_B \otimes \mathbf{1}^{\otimes 2})(id \otimes \sigma_\Theta \otimes id)(\nabla_\Theta \otimes id)(k)(\mathbf{1}^{\otimes 3} \otimes f)(\mathbf{1}^{\otimes 2} \otimes a_{j,t}^\perp) \right) \\ &= (id^{\otimes 3} \otimes \tau_\Theta)(\mathbf{K} \bullet (\mathbf{1}^{\otimes 3} \otimes f)) \end{aligned}$$

with

$$\mathbf{K} = (\mathbf{1} \otimes \chi_B \otimes \mathbf{1}^{\otimes 2})(id \otimes \sigma_\Theta \otimes id)(\nabla_\Theta \otimes id)(k) \bullet (\mathbf{1}^{\otimes 2} \otimes a_{j,t}^\perp)$$

in $\mathcal{L}(\mathbb{F}_n) \bar{\otimes} (\mathcal{S}(\mathbb{R}^n) \otimes_\pi \mathcal{S}_\Theta \otimes_\pi \mathcal{S}'_\Theta)'$. Thus, $(\mathbf{1} \otimes \chi_B \otimes \mathbf{1})(\nabla \otimes id_{\mathcal{R}_\Theta})(\Sigma) = T_{\mathbf{K}}(f)$. We turn to the proofs of Theorem 2.6, Proposition 2.15 and Remark 2.16 (as the generalizations of Theorem 2.6) in [21], they show that the condition ii) in Theorem 4.8 implies

$$\|T_{\mathbf{K}}(f)\|_{\mathcal{L}(\mathbb{F}_n) \bar{\otimes} L_\infty(\mathbb{R}^n) \bar{\otimes} \mathcal{R}_\Theta} \leq C_n(\alpha, \beta) \frac{A_2}{R} \|f\|_{\mathcal{R}_\Theta},$$

which is inequality (2.2) in [21]. Combining the calculations above, we deduce that condition ii) in Theorem 4.8 is stronger than condition ii) in Theorem 4.7. In conclusion, the Calderón-Zygmund extrapolation on \mathcal{R}_Θ that we obtain by applying Theorem 2.1 improves the corresponding result in [21].

Remark 4.9. By regularity of \mathcal{S}_Θ , L_p -estimates follows as in Corollary 2.2.

4.4. Quantum Fourier multipliers

We now refine our abstract result for locally compact quantum groups. We shall need some basic notions from the theory of quantum groups, details can be found in Kustermans/Vaes' papers [38,39]. Let us consider a von Neumann algebra \mathcal{N} equipped with a comultiplication map, a normal injective unital $*$ -morphism $\Delta : \mathcal{N} \rightarrow \mathcal{N} \bar{\otimes} \mathcal{N}$ satisfying the coassociativity law

$$(id_{\mathcal{N}} \otimes \Delta)\Delta = (\Delta \otimes id_{\mathcal{N}})\Delta.$$

Assume also the existence of two n.s.f weights ψ and φ on \mathcal{N} such that

$$(id_{\mathcal{N}} \otimes \psi)\Delta(a) = \psi(a)\mathbf{1}_{\mathcal{N}} \quad \text{and} \quad (\varphi \otimes id_{\mathcal{N}})\Delta(a) = \varphi(a)\mathbf{1}_{\mathcal{N}} \quad \text{for } a \in \mathcal{N}_+.$$

We call ψ and φ the *left-invariant Haar weight* and the *right-invariant Haar weight* on \mathcal{N} respectively. Then the quadruple $\mathbb{G} = (\mathcal{N}, \Delta, \psi, \varphi)$ is called a (von Neumann algebraic) *locally compact quantum group* and we write $L_\infty(\mathbb{G})$ for the quantum group von Neumann algebra \mathcal{N} . Using the Haar weights, one can construct an antipode S on \mathcal{N} which is a densely defined anti-automorphism on \mathcal{N} satisfying the identity

$$(id_{\mathcal{N}} \otimes \psi)((\mathbf{1}_{\mathcal{N}} \otimes a^*)\Delta(b)) = S((id_{\mathcal{N}} \otimes \psi)(\Delta(a^*)(\mathbf{1}_{\mathcal{N}} \otimes b))).$$

The comultiplication map Δ determines a multiplication on the predual $L_1(\mathbb{G})$ given by convolution $\varphi_1 \star \varphi_2(a) = (\varphi_1 \otimes \varphi_2)\Delta(a)$. The pair $(L_1(\mathbb{G}), \star)$ forms a Banach algebra. In what follows, if not specified otherwise, the quantum groups \mathbb{G} we shall work with admit a tracial left-invariant Haar weight ψ . The simplest model of noncommutative quantum groups are group von Neumann algebras $\mathcal{L}(G)$ associated to discrete groups. If λ is the left regular representation of G , the comultiplication is determined by $\Delta(\lambda(g)) = \lambda(g) \otimes \lambda(g)$. Its isometric nature follows from Fell's absorption principle and the convolution is abelian. The standard trace on $\mathcal{L}(G)$ is a left and right-invariant Haar weight. Moreover, in this case, the antipode is bounded and $S(\lambda(g)) = \lambda(g^{-1})$.

A convolution semigroup of states is a family $(\phi_t)_{t \geq 0}$ of normal states on $L_\infty(\mathbb{G})$ such that $\phi_{t_1} \star \phi_{t_2} = \phi_{t_1+t_2}$. The corresponding semigroup of completely positive maps is given by

$$S_{\Delta,t}(a) = (\phi_t \otimes id_{\mathbb{G}}) \circ \Delta(a).$$

When $\mathcal{S}_\Delta = (S_{\Delta,t})_{t \geq 0}$ is a Markov semigroup such that $\psi \circ S_{\Delta,t} = \psi$ for all $t \geq 0$, we call it a *convolution semigroup*.

Lemma 4.10. *Let \mathbb{G} be a locally compact quantum group equipped with a convolution semigroup of states $(\phi_t)_{t \geq 0}$. Then, $\mathcal{S}_\Delta = ((\phi_t \otimes id_{\mathbb{G}}) \circ \Delta)_{t \geq 0}$ is a Markov semigroup on $L_\infty(\mathbb{G})$ whenever*

- i) $\phi_t \circ S = \phi_t$ for all $t \geq 0$,
- ii) $S_{\Delta,t}(a) \rightarrow a$ as $t \rightarrow 0^+$ in the weak-* topology of $L_\infty(\mathbb{G})$.

Proof. Let us begin with the self-adjointness

$$\begin{aligned} \psi(a^* S_{\Delta,t}(b)) &= \psi(a^* (\phi_t \otimes id_{\mathbb{G}}) \Delta(b)) \\ &= \phi_t \otimes \psi((\mathbf{1}_{\mathbb{G}} \otimes a^*) \Delta(b)) \\ &= \phi_t \left((id_{\mathbb{G}} \otimes \psi)((\mathbf{1}_{\mathbb{G}} \otimes a^*) \Delta(b)) \right) \\ &= \phi_t \left(S(\underbrace{id_{\mathbb{G}} \otimes \psi(\Delta(a^*)(\mathbf{1}_{\mathbb{G}} \otimes b))}_{\rho}) \right). \end{aligned}$$

This means that $\psi(a^* S_{\Delta,t}(b)) = \phi_t(S(\rho)) = \phi_t(\rho)$ and we get

$$\begin{aligned} \psi(a^* S_{\Delta,t}(b)) &= \phi_t \otimes \psi(\Delta(a^*)(\mathbf{1}_{\mathbb{G}} \otimes b)) \\ &= \psi((\phi_t \otimes id_{\mathbb{G}}) \circ \Delta(a^*)b) = \psi(S_{\Delta,t}(a)^* b). \end{aligned}$$

The remainder properties are straightforward. Indeed, identity $S_{\Delta,t}(\mathbf{1}_{\mathbb{G}}) = \mathbf{1}_{\mathbb{G}}$ is obvious. The weak-* convergence of the $S_{\Delta,t}(a)$'s as $t \rightarrow 0^+$ is assumed and the complete positivity is clear. The normality follows from the weak-* continuity of ϕ_t and Δ . Finally, the semigroup law easily follows from coassociativity. \square

In what follows, we shall assume that the hypotheses of Lemma 4.10 hold. Let us fix a quantum group $\mathbb{G} = (\mathcal{N}, \Delta, \psi, \varphi)$ and consider a convolution semigroup \mathcal{S}_Δ associated to it. A Markov metric

$$\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t}) : j, t \in \mathbb{Z}_+ \times \mathbb{R}_+\}$$

in $L_\infty(\mathbb{G}) = \mathcal{N}$ associated to \mathcal{S}_Δ will be called an *intrinsic Markov metric* when there exists an increasing family of projections $p_{j,t}$ in $L_\infty(\mathbb{G})$ such that the cpu maps take the form

$$R_{j,t}f = \frac{1}{\psi(p_{j,t})}(\psi \otimes id_{\mathbb{G}})((p_{j,t} \otimes \mathbf{1}_{\mathbb{G}})\Delta(f)). \quad (4.6)$$

In other words, we use the algebraic skeleton

$$(\mathcal{N}_\rho = \mathcal{N}_\pi, \rho_1, \rho_2, \mathbf{E}_\rho, q_{j,t}) = (L_\infty(\mathbb{G}) \bar{\otimes} L_\infty(\mathbb{G}), \mathbf{1} \otimes \cdot, \Delta, \psi \otimes id_{\mathbb{G}}, p_{j,t} \otimes \mathbf{1}_{\mathbb{G}}).$$

Remark 4.11. Assume that

$$\gamma_{j,t} \in \mathbb{R}_+ \quad \text{and} \quad \gamma_{j,t}^2 \geq \frac{\psi(p_{j,t})}{\psi(p_{1,t})}.$$

Then, the term $|R_{j,t}f - M_tf|$ in the metric BMO norm satisfies for $M_t = R_{1,t}$ that

$$\begin{aligned} |R_{j,t}f - M_tf|^2 &= \left| \frac{1}{\psi(p_{1,t})} (\psi \otimes id_{\mathbb{G}})((\Delta(f) - \mathbf{1} \otimes R_{j,t}f)(p_{1,t} \otimes \mathbf{1})) \right|^2 \\ &\leq \frac{1}{\psi(p_{1,t})} (\psi \otimes id_{\mathbb{G}}) \left[|(\Delta(f) - \mathbf{1} \otimes R_{j,t}f)(p_{1,t} \otimes \mathbf{1})|^2 \right] \\ &= \frac{\psi(p_{j,t})}{\psi(p_{1,t})} (R_{j,t}|f|^2 - |R_{j,t}f|^2) \leq \gamma_{j,t} (R_{j,t}|f|^2 - |R_{j,t}f|^2) \gamma_{j,t}. \end{aligned}$$

According to Theorem 1.6, this yields

$$\|f\|_{\text{BMO}_{\mathcal{S}_\Delta}^c} \lesssim \|f\|_{\text{BMO}_{\mathbb{Q}}^c} \lesssim \sup_{t>0} \sup_{j \geq 1} \left\| \left(R_{j,t}|f|^2 - |R_{j,t}f|^2 \right)^{\frac{1}{2}} \right\|_{L_\infty(\mathbb{G})}. \quad (4.7)$$

Additionally, we may consider transferred Markov metrics in other von Neumann algebras. Consider a convolution semigroup of states $(\phi_t)_{t \geq 0}$ on a locally compact quantum group $L_\infty(\mathbb{G})$ and a von Neumann algebra equipped with a n.s.f. trace (\mathcal{M}, τ) . A corepresentation $\pi : \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$ is a normal injective $*$ -representation satisfying the identity

$$(id_{\mathbb{G}} \otimes \pi) \circ \pi = (\Delta \otimes id_{\mathcal{M}}) \circ \pi.$$

Every such π yields a *transferred convolution semigroup* $\mathcal{S}_\pi = (S_{\pi,t})_{t \geq 0}$ with

$$\begin{aligned} S_{\pi,t} : \mathcal{M} &\rightarrow \mathcal{M}, \\ S_{\pi,t}f &= (\phi_t \otimes id_{\mathcal{M}}) \circ \pi(f). \end{aligned}$$

Lemma 4.12. Assume that

- $\tau(S_{\pi,t}(f_1)^* f_2) = \tau(f_1^* S_{\pi,t}(f_2))$,
- $S_{\pi,t}f \rightarrow f$ as $t \rightarrow 0^+$ in the weak- $*$ topology of \mathcal{M} .

Then \mathcal{S}_π defines a Markov semigroup on \mathcal{M} such that $\pi \circ S_{\pi,t} = (S_{\Delta,t} \otimes id_{\mathcal{M}}) \circ \pi$.

Proof. It is easy to check that $S_{\pi,t}$ is cpu and the normality follows from the weak-* continuity of ϕ_t and π . Hence, it remains to show the identity $\pi \circ S_{\pi,t} = (S_t \otimes id_{\mathbb{G}}) \circ \pi$ and the semigroup law. We first observe that $\pi(\phi_t \otimes id_{\mathcal{M}}) = (\phi_t \otimes id_{\mathbb{G}} \otimes id_{\mathcal{M}})(id_{\mathbb{G}} \otimes \pi)$ as maps on $L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M}$. Indeed, by weak-* continuity, it suffices to test the identity on elementary tensors $n \otimes m$, for which the identity is trivial. Therefore, we have

$$\begin{aligned} (S_{\Delta,t} \otimes id_{\mathcal{M}})\pi &= (\phi_t \otimes id_{\mathbb{G}} \otimes id_{\mathcal{M}})(\Delta \otimes id_{\mathcal{M}})\pi \\ &= (\phi_t \otimes id_{\mathbb{G}} \otimes id_{\mathcal{M}})(id_{\mathbb{G}} \otimes \pi)\pi \\ &= \pi(\phi_t \otimes id_{\mathcal{M}})\pi = \pi S_{\pi,t}. \end{aligned}$$

For the semigroup law we note that

$$\begin{aligned} S_{\pi,t_1} S_{\pi,t_2} &= (\phi_{t_1} \otimes id_{\mathcal{M}})(\phi_{t_2} \otimes id_{\mathbb{G}} \otimes id_{\mathcal{M}})(id_{\mathbb{G}} \otimes \pi)\pi \\ &= (\phi_{t_1} \otimes id_{\mathcal{M}})(\phi_{t_2} \otimes id_{\mathbb{G}} \otimes id_{\mathcal{M}})(\Delta \otimes id_{\mathcal{M}})\pi \\ &= (\phi_{t_2} \otimes \phi_{t_1} \otimes id_{\mathcal{M}})(\Delta \otimes id_{\mathcal{M}})\pi = (\phi_{t_2} \star \phi_{t_1} \otimes id_{\mathcal{M}})\pi = S_{\pi,t_1+t_2}. \quad \square \end{aligned}$$

In the sequel, we shall assume that the assumptions in Lemma 4.12 hold. Intrinsic Markov metrics on $L_{\infty}(\mathbb{G})$ yield *transferred Markov metrics* on \mathcal{M} associated to the transferred convolution semigroup S_{π} . Indeed, given any intrinsic Markov metric $\mathcal{Q} = \{(R_{j,t}, \sigma_{j,t}, \gamma_{j,t})\}$ in \mathbb{G} with cpu maps $R_{j,t}$ given by (4.6), the transferred cpu maps $R_{\pi,j,t}$ are given by

$$R_{\pi,j,t}f = \frac{1}{\psi(p_{j,t})}(\psi \otimes id_{\mathcal{M}})((p_{j,t} \otimes \mathbf{1})\pi(f)).$$

It is easy to check that $\pi \circ R_{\pi,j,t} = (R_{j,t} \otimes id_{\mathcal{M}}) \circ \pi$. Assume in addition that $\sigma_{j,t} \in \mathbb{R}_+$. Then, arguing as we did before Corollary 4.5 for the corepresentation u of \mathbb{R} in $\mathcal{B}(\ell_2)$, we get a Markov metric in \mathcal{M}

$$\mathcal{Q}_{\pi} = \{(R_{\pi,j,t}, \sigma_{j,t}, \gamma_{j,t}) : j \in \mathbb{Z}_+, t \in \mathbb{R}_+\}.$$

Let $\alpha : \mathbb{N} \rightarrow \mathbb{N}$ be a strictly increasing function with $\alpha(j) > j$. This Markov metric is called α -doubling if there exists some constant c_{α} such that $\psi(q_{\alpha(j),t}) \leq c_{\alpha}\psi(q_{j,t})$.

Remark 4.13. In what follows, we impose our Markov metrics to be α -doubling for some function $\alpha : \mathbb{N} \rightarrow \mathbb{N}$, to satisfy $\sigma_{j,t} \in \mathbb{R}_+$ as well as the condition in Remark 4.11. Altogether, this allows to eliminate the size CZ condition and reduce the number of analytic and smoothness CZ conditions to be checked for both the intrinsic Markov metric and the transferred one.

Observe that the transferred formulation above includes the intrinsic formulation by taking $(\mathcal{M}, \pi) = (\mathbb{G}, \Delta)$. Let us now state the corresponding Calderón-Zygmund theory.

Given $\mathcal{A}_{\mathcal{M}}$ a weakly dense $*$ -subalgebra of \mathcal{M} , let T be a (not necessarily bounded) operator $T : \mathcal{A}_{\mathcal{M}} \rightarrow \mathcal{M}$. We say T is a *transferred map* if there exists an amplification map

$$\widehat{T} : \mathcal{D} \subset L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M} \rightarrow L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M}$$

satisfying the identity

$$\pi \circ T = \widehat{T} \circ \pi|_{\mathcal{A}_{\mathcal{M}}}. \quad (4.8)$$

Again, \mathcal{D} is a weakly dense $*$ -subalgebra for which $\pi(\mathcal{A}_{\mathcal{M}}) \subset \mathcal{D}$. In the case $(\mathcal{M}, \pi) = (\mathbb{G}, \Delta)$, we can always take the amplification $T \otimes id_{\mathcal{M}}$ and condition above just imposes that T is a quantum Fourier multiplier. In the following theorem, we provide sufficient conditions on the amplification map to make a given transferred CZO T bounded from $\mathcal{A}_{\mathcal{M}}$ to $\text{BMO}_{S_{\pi}}^c$.

Theorem 4.14. *Let*

$$\pi : \mathcal{M} \rightarrow L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M}$$

be a corepresentation of a locally compact quantum group \mathbb{G} in a semifinite von Neumann algebra (\mathcal{M}, τ) . Assume that $L_{\infty}(\mathbb{G})$ comes equipped with an α -doubling intrinsic Markov metric \mathcal{Q} determined by an increasing family of projections $p_{j,t}$ as above. Then, a transferred map T (with amplification for which (4.8) holds) will be bounded from $\mathcal{A}_{\mathcal{M}}$ to $\text{BMO}_{S_{\pi}}^c$ provided:

- i) $\widehat{T} : L_2^c(\mathbb{G}) \bar{\otimes} \mathcal{M} \rightarrow L_2^c(\mathbb{G}) \bar{\otimes} \mathcal{M}$ is bounded,
- ii) $\frac{(\psi \otimes \psi \otimes id_{\mathcal{M}})}{\psi(p_{j,t})^2} \left((p_{j,t} \otimes p_{j,t} \otimes \mathbf{1}_{\mathcal{M}}) | \delta_{\mathbb{G}}(\widehat{T}(\pi(f)p_{\alpha(j),t}^{\perp})) |^2 \right) \lesssim \|f\|_{\mathcal{M}}^2$.

Proof. We use the algebraic skeleton

$$(\mathcal{M}, \mathcal{N}_{\rho} = \mathcal{N}_{\pi}, \rho_1, \rho_2, E_{\rho}, q_{j,t}) = (\mathcal{M}, L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M}, \mathbf{1} \otimes \cdot, \pi, \psi \otimes id_{\mathbb{G}}, p_{j,t} \otimes \mathbf{1}_{\mathbb{G}}).$$

Identity (4.8) is the compatibility condition (2.3). Let us justify the algebraic conditions. The second one is trivial since both $E_{\rho}(q_{j,t})$ and $\rho_1(\gamma_{j,t})$ belong to \mathbb{R}_+ in this case. For the first one, consider the product \bullet in $L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M}_{\text{op}}$. Then, we just observe that

$$\begin{aligned} E_{\rho}(q_{j,t} |\xi|^2 q_{j,t}) &= (\psi \otimes id_{\mathcal{M}})((p_{j,t} \otimes \mathbf{1}) \xi^* \xi) \\ &= (\psi \otimes id_{\mathcal{M}})(\xi \bullet (p_{j,t} \otimes \mathbf{1}) \bullet \xi^*) \\ &\leq (\psi \otimes id_{\mathcal{M}})(\xi \bullet \xi^*) = (\psi \otimes id_{\mathcal{M}})(\xi^* \xi) = E_{\rho}(|\xi|^2). \end{aligned}$$

Define the amplifications

$$\widehat{R}_{\pi,j,t} : L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M} \ni \xi \mapsto \frac{1}{\psi(p_{j,t})} (\psi \otimes id_{\mathcal{M}})((p_{j,t} \otimes \mathbf{1}_{\mathcal{M}})\xi) \in \mathcal{M}.$$

Consider also the cpu maps

$$\begin{aligned} \Phi_{j,t} &: L_{\infty}(\mathbb{G}) \bar{\otimes} L_{\infty}(\mathbb{G}) \bar{\otimes} \mathcal{M} \rightarrow \mathcal{M}, \\ \Phi_{j,t}(\eta) &= \frac{1}{\psi(p_{j,t})^2} (\psi \otimes \psi \otimes id_{\mathcal{M}})((p_{j,t} \otimes p_{j,t} \otimes \mathbf{1}_{\mathcal{M}})\eta). \end{aligned}$$

Recalling that $\delta_{\mathbb{G}}(x) = x \otimes \mathbf{1} - \mathbf{1} \otimes x$, the identity

$$\Phi_{j,t}(|\delta_{\mathbb{G}}(\xi)|^2) = 2\widehat{R}_{\pi,j,t}(|\xi|^2) - 2|\widehat{R}_{\pi,j,t}(\xi)|^2,$$

is straightforward. This readily implies the first analytic condition. On the other hand, since the auxiliary Markov metric is α -doubling, the second analytic condition reduces to note that

$$(\psi \otimes id_{\mathcal{M}})(q_{\alpha(j),t} \otimes \mathbf{1}_{\mathcal{M}}) \leq c_{\alpha}(\psi \otimes id_{\mathcal{M}})(q_{j,t} \otimes \mathbf{1}_{\mathcal{M}}).$$

Thus, according to inequality (4.7), the assertion follows from Theorem 2.1. \square

Remark 4.15. As noticed, the main particular case of Theorem 4.14 arises for $(\mathcal{M}, \pi) = (\mathbb{G}, \Delta)$ with amplification $T \otimes id_{\mathbb{G}}$. Condition (4.8) becomes the identity

$$\Delta \circ T = (T \otimes id_{\mathbb{G}}) \circ \Delta.$$

In other words, these are translation invariant CZ operators. We also call them quantum Fourier multipliers in this paper and it can be checked, as expected, that these maps are of convolution type in the sense that there exists a kernel k affiliated to $L_{\infty}(\mathbb{G})$ so that

$$Tf = k \star f = (id_{\mathbb{G}} \otimes \psi)(\Delta(k)(\mathbf{1}_{\mathbb{G}} \otimes Sf)).$$

In this particular case, it is not difficult to prove that our conditions in Theorem 4.14 reduce to those in Theorem B2 from the Introduction. Of course, Theorem 4.14 also applies as well for nonconvolution CZ operators on quantum groups, or even for transferred forms of them to other von Neumann algebras \mathcal{M} .

Remark 4.16. One may consider twisted convolution CZO's on quantum groups applying Theorem 4.14. As an illustration, assume that $\mathbb{G} \curvearrowright L_{\infty}(\mathbb{G})$ by a trace preserving action α and that \mathbb{G} is a quantum group satisfying $(\alpha_g \otimes \alpha_g)\Delta = \Delta\alpha_g$ for all $g \in \mathbb{G}$. This property is quite natural in the commutative case, where quantum groups come from locally compact groups and α is typically implemented by a measure preserving map β . Note that the underlying Haar measure is translation invariant and the condition above just imposes that β is an homomorphism. Let us see what we get for a map

$$\sum_g f_g \lambda(g) \mapsto \sum_g T_g(f_g) \lambda(g),$$

where the T_g 's are normal convolution maps on $L_\infty(\mathbb{G})$. Assume $L_\infty(\mathbb{G})$ comes equipped with a convolution G -equivariant semigroup \mathcal{S}_Δ which admits a η -doubling intrinsic Markov metric. Then, we get a bounded map $L_\infty(\mathbb{G}) \rtimes G \rightarrow \text{BMO}_{\mathcal{S}_\Delta}^c$ when the following conditions hold:

i) We have a bounded map

$$L_2^c(\mathbb{G}) \bar{\otimes} \mathcal{B}(\ell_2(G)) \ni \xi \mapsto (T_{gh^{-1}}) \bullet \xi \in L_2^c(\mathbb{G}) \bar{\otimes} \mathcal{B}(\ell_2(G)),$$

where \bullet stands once more for the generalized Schur product of matrices.

ii) Letting $\mathcal{R} = L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{B}(\ell_2(G))$ and $\Psi(\xi) = \sum_{g,h} (\alpha_{g^{-1}}) \bullet (T_{gh^{-1}}) \bullet (\alpha_g) \bullet \xi$,

$$\frac{(\psi \otimes \psi \otimes id_{\mathcal{B}(\ell_2(G))})}{\psi(p_{j,t})^2} \left((p_{j,t} \otimes p_{j,t} \otimes \mathbf{1}) | \delta_{\mathbb{G}}(\Psi(\xi q_{\eta(j),t}^\perp)) |^2 \right) \lesssim \|\xi\|_{\mathcal{R}}^2.$$

Remark 4.17. All our results in this paragraph impose the additional assumption that our quantum groups admit a tracial Haar weight. We believe however that our results can be extended to the general non-tracial case. We leave this generalization open to the interested reader.

5. Noncommutative transference

Originally motivated by Cotlar's paper [14] and the method of rotations, Calderón developed a circle of ideas [4] which was called the *transference method* after the systematic study of Coifman/Weiss in their monograph [12]. The fundamental work of K. de Leeuw [16] also had a big impact in this line of research. Let us consider an amenable locally compact group G with left Haar measure μ , a σ -finite measure space (Ω, ν) and a uniformly bounded representation $\beta : G \rightarrow \mathcal{B}(L_p(\Omega))$. Roughly, Calderón's transference is a technique which allows to transfer the L_p boundedness of a convolution operator $f \mapsto k \star f$ on $L_p(G)$ to the corresponding transferred operator on $L_p(\Omega)$

$$Vf(w) = \int_G k(g) \beta_{g^{-1}} f(w) d\mu(g),$$

for some compactly supported kernel k in $L_1(G)$. A case by case limiting procedure also allows to consider more general (singular) kernels. In the rest of this section we shall develop a noncommutative form of Calderón-Coifman-Weiss technique.

Our first task is to clarify what we mean by 'representation' and 'amenable' in the context of quantum groups. Using the commutative locally compact quantum group $L_\infty(G)$ as above, a representation $\beta : G \rightarrow \text{Aut}(\mathcal{M})$ induces a $*$ -representation $\pi_\beta : \mathcal{M} \rightarrow L_\infty(G; \mathcal{M})$ by

$$\pi_\beta f(g) = \beta_{g^{-1}} f.$$

Note that we have

$$\begin{aligned} (id_{\mathbb{G}} \otimes \pi_\beta)(\pi_\beta f)(g, h) &= \pi_\beta(\beta_{g^{-1}} f)(h) = \beta_{h^{-1}} \beta_{g^{-1}} f \\ &= \beta_{(gh)^{-1}} f = (\Delta_{\mathbb{G}} \otimes id_{\mathcal{M}})(\pi_\beta f)(g, h). \end{aligned}$$

Given a semifinite von Neumann algebra (\mathcal{M}, τ) and a locally compact quantum group \mathbb{G} , this leads us to consider corepresentations $\pi : \mathcal{M} \mapsto L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$ satisfying $(id_{\mathbb{G}} \otimes \pi) \circ \pi = (\Delta \otimes id_{\mathcal{M}}) \circ \pi$. Note that comultiplication is a corepresentation by coassociativity. To show what we mean by ‘uniformly bounded’, let us go back to our motivating example $\beta : \mathbb{G} \rightarrow \text{Aut}(\mathcal{M})$, where we take $\mathcal{M} = L_\infty(\Omega)$ for some σ -finite measure space (Ω, ν) . In the classical case

$$\|\beta_g f\|_p \sim \|f\|_p \quad \text{for all } g \in \mathbb{G}$$

up to an absolute constant independent of f, g . We say that a corepresentation $\pi : \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$ is *uniformly bounded in $L_p(\mathcal{M})$* if for any $f \in \mathcal{M} \cap L_p(\mathcal{M})$ we have

$$\frac{1}{c_\pi} \|f\|_{L_p(\mathcal{M})}^p \leq (id_{\mathbb{G}} \otimes \tau)(|\pi(f)|^p) \leq c_\pi \|f\|_{L_p(\mathcal{M})}^p$$

for some absolute constant c_π independent of f . Note that our notion again reduces to the classical one on $L_\infty(\mathbb{G})$. Note also that, since $|\pi(f)|^p = \pi(|f|^p)$, our definition reduces to the p -independent condition

$$\frac{1}{c_\pi} \|f\|_{L_1(\mathcal{M})} \leq (id_{\mathbb{G}} \otimes \tau)(\pi(f)) \leq c_\pi \|f\|_{L_1(\mathcal{M})} \quad \text{for all } f \in \mathcal{M}_+ \cap L_1(\mathcal{M}).$$

Now we introduce what we mean by an ‘amenable’ quantum group. We say that \mathbb{G} satisfies *Følner’s condition* if for every projection $q \in L_1(\mathbb{G})$ and every $\varepsilon > 0$, there exists two non-zero projections $q_1, q_2 \in L_1(\mathbb{G})$ such that

$$\Delta(q_1)(q \otimes q_2) = q \otimes q_2 \quad \text{and} \quad \psi(q_1) \leq (1 + \varepsilon) \psi(q_2).$$

In the standard example for a locally compact group G , where $(L_\infty(G), \psi)$ is $L_\infty(G)$ equipped with the left Haar measure μ and Δ is given by $\Delta_G(\xi)(g, h) = \xi(gh)$ the classical comultiplication, it turns out that \mathbb{G} is amenable iff G is an amenable group. Indeed, our notion can be rephrased in this case by saying that for any compact set K in G and any $\varepsilon > 0$, there exists a non-empty neighborhood of the identity W of finite measure such that

$$\mu(KW) \leq (1 + \varepsilon) \mu(W),$$

which corresponds to $(q, q_1, q_2) = (\chi_K, \chi_{KW}, \chi_W)$ in our formulation. This is exactly the classical characterization of amenability, known as Følner's condition, used by Coifman and Weiss in [12]. Given an amenable locally compact group G with left Haar measure μ , it is clear that $L_\infty(G, \mu)$ with its natural quantum group structure is amenable. On the other hand, as expected, any compact quantum group is amenable just by taking $q_1 = q_2 = \mathbf{1}_G$.

Assume that \mathbb{G} admits a corepresentation $\pi : \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$. Given $\mathcal{A}_\mathcal{M}$ a weakly dense $*$ -subalgebra of \mathcal{M} , we say that a linear operator $V : \mathcal{A}_\mathcal{M} \rightarrow \mathcal{M}$ is a *transferred convolution map* if there exists $\Phi : \mathcal{D} \subset L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$, an auxiliary convolution map such that $\pi \circ V = \Phi \circ \pi|_{\mathcal{A}_\mathcal{M}}$. The classical transferred operator

$$V = \int_G k(g) \beta_{g^{-1}} f(w) d\mu(g)$$

comes from

$$\Phi(\xi)(g, w) = \int_G k(h) \xi(hg, w) d\mu(h) = (\varphi \otimes id_G \otimes id_\Omega)(\Delta_G \otimes id_\Omega).$$

If $\pi_\beta f(g) = \beta_{g^{-1}} f$ denotes the corresponding corepresentation, we may then apply the identities in the proof of Lemma 4.12 again to deduce the following identities

$$\Phi \circ \pi_\beta = (\varphi \otimes id_G \otimes id_\Omega)(\Delta_G \otimes id_\Omega) \pi_\beta = (\varphi \otimes id_G \otimes id_\Omega)(id_G \otimes \pi_\beta) \pi_\beta = \pi_\beta(\varphi \otimes id_\Omega) \circ \pi_\beta.$$

By injectivity of π_β , we must have

$$Vf(w) = (\varphi \otimes id_\Omega) \pi_\beta f(w) = \int_G k(g) \beta_{g^{-1}} f(w) d\mu(g)$$

as expected. This shows how we recover the classical construction.

Let us now settle the framework for our transference result. Assume that \mathbb{G} is amenable and consider $\pi : \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$ a uniformly bounded corepresentation in $L_2(\mathcal{M})$. We say that $T : L_p(\mathbb{G}) \rightarrow L_p(\mathbb{G})$ is a convolution map with finitely supported L_1 kernel when the map T has the form $T = (\phi \otimes id_\mathbb{G}) \circ \Delta$ for some functional $\phi = \psi(d \cdot)$, with d an element in $L_1(\mathbb{G})$ whose left support q satisfies $\psi(q) < \infty$. In the commutative case, this is the kind of operators which are transferred. Roughly, the goal is to show how a limit operator $T = \lim_\gamma T_\gamma$ of such maps which is bounded on $L_2(\mathbb{G})$ and $L_\infty(\mathbb{G}) \rightarrow \text{BMO}_S$ can be transferred under suitable conditions to a bounded map on $L_p(\mathcal{M})$.

Remark 5.1. Young's inequality extends to this setting as

$$\|d \star f\|_p = \|(\phi \otimes id_\mathbb{G}) \Delta(f)\|_p \leq 4 \|d\|_1 \|f\|_p,$$

where $\phi = \psi(d \cdot)$ and $1 \leq p \leq \infty$. Indeed, when d and f are positive the inequality holds with constant 1. This can be justified by interpolation. When $p = 1$ we use Fubini and the left-invariance of ψ , while for $p = \infty$ it follows from the fact that $(\phi \otimes id_{\mathbb{G}})\Delta$ is a positive map with $\mathbf{1}_{\mathbb{G}} \mapsto \psi(d)$. In the general case, we split d, f into their positive parts and obtain the constant 4. In fact, the same argument still holds after matrix amplification and we deduce that $(\phi \otimes id_{\mathbb{G}})\Delta$ is completely bounded on $L_p(\mathbb{G})$ with cb-norm $4\|d\|_1$. This is however not enough for transference, since the norms $\|d_\gamma\|_1$ might not be uniformly bounded.

Theorem 5.2. *Let \mathbb{G} be an amenable quantum group and consider a uniformly bounded corepresentation $\pi : \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$ in $L_2(\mathcal{M})$ associated with a semifinite von Neumann algebra (\mathcal{M}, τ) . Let $T : L_2(\mathbb{G}) \rightarrow L_2(\mathbb{G})$ be a bounded map and assume that*

$$(T \otimes id_{\mathcal{M}}) = \text{SOT} - \lim_{\gamma} (T_\gamma \otimes id_{\mathcal{M}})$$

for some net $T_\gamma = (\phi_\gamma \otimes id_{\mathbb{G}}) \circ \Delta$ of convolution maps with finitely supported L_1 kernels and such that $\lim_{\gamma} \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))} \leq \|T\|_{\mathcal{B}(L_2(\mathbb{G}))}$. Then, the net of transferred operators $V_\gamma = (\phi_\gamma \otimes id_{\mathcal{M}}) \circ \pi$ satisfies the inequalities

$$\|V_\gamma\|_{\mathcal{B}(L_2(\mathcal{M}))} \leq c_\pi \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))}.$$

We thus find a WOT-cluster point V satisfying $\|V\|_{\mathcal{B}(L_2(\mathcal{M}))} \leq c_\pi \|T\|_{\mathcal{B}(L_2(\mathbb{G}))}$.

Proof. Note that we have

$$\pi V_\gamma = (\phi_\gamma \otimes id)(id_{\mathbb{G}} \otimes \pi)\pi = (\phi_\gamma \otimes id)(\Delta \otimes id_{\mathcal{M}})\pi = (T_\gamma \otimes id_{\mathcal{M}})\pi.$$

Hence, the uniform boundedness of π yields

$$\frac{1}{c_\pi} \|V_\gamma f\|_2^2 \leq (\rho \otimes \tau)(|\pi V_\gamma(f)|^2) = (\rho \otimes \tau)(|(T_\gamma \otimes id_{\mathcal{M}})\pi(f)|^2)$$

for any state ρ on $L_\infty(\mathbb{G})$. On the other hand, if $\phi_\gamma = \psi(d_\gamma \cdot)$ and q_γ denotes the left support of d_γ , we know from the amenability assumption that for any $\varepsilon > 0$ we may find projections $q_{1\gamma}$ and $q_{2\gamma}$ such that

$$\Delta(q_{1\gamma})(q_\gamma \otimes q_{2\gamma}) = q_\gamma \otimes q_{2\gamma} \quad \text{and} \quad \psi(q_{1\gamma}) \leq (1 + \varepsilon) \psi(q_{2\gamma}).$$

Taking $\rho = \psi(q_{2\gamma} \cdot)/\psi(q_{2\gamma})$, we obtain the inequality

$$\frac{1}{c_\pi} \|V_\gamma f\|_2^2 \leq \frac{(\psi \otimes \tau)}{\psi(q_{2\gamma})} \left(\left| (\phi \otimes id)((\Delta \otimes id_{\mathcal{M}})\pi(f)(\mathbf{1}_{\mathbb{G}} \otimes q_{2\gamma} \otimes \mathbf{1}_{\mathcal{M}})) \right|^2 \right)$$

since ρ is supported by $q_{2\gamma}$. Moreover, $d_\gamma \otimes q_{2\gamma}$ is supported on the left by $q_\gamma \otimes q_{2\gamma}$ and amenability provides $d_\gamma \otimes q_{2\gamma} = \Delta(q_{1\gamma})(d_\gamma \otimes q_{2\gamma})$. Once we have created $\Delta(q_{1\gamma})$, we can eliminate $q_{2\gamma}$. Altogether gives

$$\frac{1}{c_\pi} \|V_\gamma f\|_2^2 \leq \frac{1}{\psi(q_{2\gamma})} \left\| (T_\gamma \otimes id)(\pi(f)(q_{1\gamma} \otimes \mathbf{1}_\mathcal{M})) \right\|_{L_2(L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M})}^2.$$

Now we use the L_2 boundedness of T_γ and uniform boundedness of π to conclude

$$\begin{aligned} \frac{1}{c_\pi} \|V_\gamma f\|_2^2 &\leq \frac{1}{\psi(q_{2\gamma})} \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))}^2 \psi\left((q_{1\gamma} \otimes \mathbf{1}_\mathcal{M})(id_\mathbb{G} \otimes \tau)(|\pi(f)|^2)\right) \\ &\leq \frac{c_\pi}{\psi(q_{2\gamma})} \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))}^2 \psi(q_{1\gamma}) \|f\|_2^2 \leq c_\pi (1 + \varepsilon) \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))}^2 \|f\|_2^2. \end{aligned}$$

Letting $\varepsilon \rightarrow 0$, we prove the inequality

$$\|V_\gamma\|_{\mathcal{B}(L_2(\mathcal{M}))} \leq c_\pi \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))}.$$

Since T is bounded on $L_2(\mathbb{G})$ and $\lim_\gamma \|T_\gamma\|_{\mathcal{B}(L_2(\mathbb{G}))} \leq \|T\|_{\mathcal{B}(L_2(\mathbb{G}))}$, the operators V_γ are eventually in a ball of radius $c_\pi(1 + \delta)\|T\|_{\mathcal{B}(L_2(\mathbb{G}))}$ for any $\delta > 0$. The closure of such ball is weak operator compact and thus we find our cluster point. \square

We now study $L_\infty \rightarrow \text{BMO}$ transference and then interpolate/dualize to obtain L_p -transference. This approach seems to be new even in the classical theory and where our semigroup formulation becomes an essential ingredient.

Corollary 5.3. *Let \mathbb{G} be a compact (hence amenable) quantum group equipped with a $L_2(\mathcal{M})$ -uniformly bounded corepresentation $\pi : \mathcal{M} \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$. Let $(\phi_t)_{t \geq 0}$ be a convolution semigroup of states on $L_\infty(\mathbb{G})$, giving rise to Markov semigroups \mathcal{S}_Δ on (\mathbb{G}, ψ) and \mathcal{S}_π on (\mathcal{M}, τ) . Let $T = \text{SOT-}\lim_\gamma T_\gamma$ be as above and take $\mathcal{A}_\mathcal{M} = \mathcal{M} \cap L_2(\mathcal{M})$. Then, if $T : L_\infty(\mathbb{G}) \rightarrow \text{BMO}_{\mathcal{S}_\Delta}$ is completely bounded, we find that*

$$V = \text{WOT-}\lim_\gamma V_\gamma : \mathcal{A}_\mathcal{M} \rightarrow \text{BMO}_{\mathcal{S}_\pi}$$

is completely bounded. Moreover, if T_π is regular, the complete boundedness of $J_p V : L_p(\mathcal{M}) \rightarrow L_p(\mathcal{M})$ follows for every $2 < p < \infty$ by interpolation. In addition the complete boundedness of $V J_p : L_p(\mathcal{M}) \rightarrow L_p(\mathcal{M})$ for $1 < p < 2$ holds under the same assumptions for T^ .*

Proof. By uniform boundedness of π we have

$$\left\| (id_\mathbb{G} \otimes \tau)(|\pi(f)|^2)^{\frac{1}{2}} \right\|_{L_\infty(\mathbb{G})} \leq c_\pi \|f\|_2,$$

which implies that $\pi : L_2(\mathcal{M}) \rightarrow L_\infty(\mathbb{G}) \bar{\otimes} L_2^c(\mathcal{M})$ is bounded by c_π . According to the finiteness of $L_\infty(\mathbb{G})$, we deduce that in fact $\pi : L_2(\mathcal{M}) \rightarrow L_2(L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M})$ is still bounded with the same norm. This proves that

$$\pi V = \text{WOT} - \lim_{\gamma} \pi V_{\gamma} = \text{WOT} - \lim_{\gamma} (T_{\gamma} \otimes id_{\mathcal{M}}) \pi = (T \otimes id_{\mathcal{M}}) \pi.$$

In particular, $\pi V = (T \otimes id_{\mathcal{M}}) \pi$ over $\mathcal{A}_{\mathcal{M}}$ and identity $\pi S_{\pi,t} = (S_t \otimes id_{\mathcal{M}}) \pi$ yields

$$\begin{aligned} \|Vf\|_{\text{BMO}_{S_{\pi}}^c} &= \sup_{t \geq 0} \left\| S_{\pi,t} |Vf|^2 - |S_{\pi,t} Vf|^2 \right\|_{\mathcal{M}}^{\frac{1}{2}} \\ &= \sup_{t \geq 0} \left\| \pi S_{\pi,t} |Vf|^2 - |\pi S_{\pi,t} Vf|^2 \right\|_{L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}}^{\frac{1}{2}} \\ &= \|(T \otimes id_{\mathcal{M}}) \pi(f)\|_{\text{BMO}_{S_{\pi}}^c} \leq \|T\|_{cb} \|\pi(f)\|_{L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}} = \|T\|_{cb} \|f\|_{\mathcal{M}} \end{aligned}$$

for $f \in \mathcal{A}_{\mathcal{M}}$. Since the same inequality holds after matrix amplification, we deduce that $V : \mathcal{A}_{\mathcal{M}} \rightarrow \text{BMO}_{S_{\pi}}^c$ is completely bounded with $\text{cb-norm} \leq \|T\|_{cb}$. The row case is similar because

$$\pi V^{\dagger} = (\pi V)^{\dagger} = (T\pi)^{\dagger} = T^{\dagger} \pi.$$

The assertions on L_p boundedness follow as usual from Theorem 1.3. \square

Remark 5.4. Under the above assumptions, we see that for $V = \text{WOT} - \lim_{\gamma} V_{\gamma}$ we can find the concrete form of its amplification map Φ defined on $L_\infty(\mathbb{G}) \bar{\otimes} \mathcal{M}$. In this case, by applying Theorems 4.14 to $\Phi = T \otimes id_{\mathcal{M}}$, we get Calderón-Zygmund extrapolation for the transferred convolution map V on \mathcal{M} .

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