

A new neutron lifetime experiment with cold neutron beam decay in superfluid helium-4

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Abstract

The puzzle remains in the large discrepancy between neutron lifetime measured by the two distinct experimental approaches—counts of beta decays in a neutron beam and storage of ultracold neutrons in a potential trap, namely, the beam method versus the bottle method. In this paper, we propose a new experiment to measure the neutron lifetime in a cold neutron (CN) beam with a sensitivity goal of 0.1% or sub-1 s. The neutron beta decays will be counted in a superfluid helium-4 scintillation detector at 0.5 K, and the neutron flux will be simultaneously monitored by the helium-3 captures in the same volume. The CN beam must be of wavelength $\lambda > 16.5 \text{ \AA}$ to eliminate scattering with superfluid helium. A new precise measurement of neutron lifetime with the beam method of unique inherent systematic effects will greatly advance in resolving the puzzle.

Keywords: neutron lifetime, neutron lifetime discrepancy, beam neutron lifetime, neutron beta decay spectrum, superfluid helium-4 scintillator

(Some figures may appear in colour only in the online journal)

1. Introduction

A precise measurement on the neutron lifetime is important to many fundamental questions in particle physics, astrophysics and cosmology, such as Cabibbo–Kobayashi–Maskawa unitarity and primordial helium abundance in big bang nucleosynthesis [1–3]. So far, its values obtained from the two distinct methods significantly differ from each other [4–6], possibly due to unaccounted systematic effects in either or both of the methods; yet otherwise it implies new physics [7–14], many theories of which remain controversial [15–17]. On one side, the measurement is done in a neutron beam by counting the number of neutrons undergoing beta decay

when the neutron flux passes through a defined volume. It is thus called the beam method. The weighted average of the recent two beam lifetime measurements with a proton quasi-Penning trap in a cold neutron (CN) beam is $\tau_n = 888.0 \pm 2.0$ s [18, 19]. On the other side, ultra-cold neutrons (UCN) can be stored in a material box or magneto-gravity trap, and the neutron lifetime is measured by counting the surviving neutrons after a period of storage. It is called the bottle method. The weighted average of several recent bottle lifetime measurements is $\tau_n = 879.4 \pm 0.6$ s [20–26]. The difference is as large as 8.7 ± 2.1 s (4.1σ). Many further experimental efforts are on the way to address the discrepancy. While existing experiments are upgrading to improve their statistics and searching for hidden systematic effects, new experimental strategies with a distinct set of systematic effects are being proposed and carried out. For instance, researchers in J-PARC started a new measurement in the pulsed CN beam by characterization of the electron recoils in the beta decay events and the helium-3 capture events in a time projection chamber filled with gaseous mixture of helium and carbon dioxide [27, 28]. It is a revival of the beam experiment originally proposed by Kossakowski *et al* in 1989 [29]. Researchers at Los Alamos National Laboratory are prototyping a beam/bottle hybrid experiment, named UCNProBe, to measure the number of decays and helium-3 captures via detection of scintillation in a UCN storage box [30].

In this paper, we propose a new experimental method with a different combination of existing technologies, in order to resolve the neutron lifetime enigma [31]. The proposed experiment is essentially a beam lifetime measurement. It counts the decay product—electrons, rather than protons, via detection of electron recoil scintillation in superfluid helium-4 at 0.5 K. In order to eliminate neutron scattering with superfluid helium, the CN beam must be of wavelength $\lambda > 16.5$ Å, where kinematics of scattering can never be satisfied [32]. The neutron flux is monitored by helium-3 captures via nuclear recoil scintillation in the same volume of superfluid helium-4. The decay events may be distinguished from the capture events, as the features of scintillation differ between the electron and nuclear recoils. Meanwhile, a precise beta spectrum of neutrons, in addition to the neutron capture peak, will be constructed in a wide energy window with good resolution. A fit of beta spectrum can separate the overlapping counts of decay events from capture events, and complement the total counts with the missing number of electrons outside the detectable energy window. In the end, an accurate neutron lifetime can be obtained with a good knowledge of the ratio of capture-to-decay event rates and the helium-3 density in superfluid helium-4.

2. Experimental method

Suppose the decay volume is cylindrical with a length L of 75 cm and a diameter D of 7.5 cm, i.e. the length to diameter ratio is $L/D = 10$. The CN beam of 3 cm diameter passes the decay volume along the axis of the cylinder. The detectable neutron decay rate in the volume is given as

$$\dot{N}_\beta = \tau_\beta^{-1} \epsilon_\beta L \int_{A_b} da \int_v dv I(v, \vec{r}) \frac{1}{v} \quad (1)$$

where τ_β is the neutron lifetime, ϵ_β is the detection efficiency of the beta decay in the given geometry, A_b is the cross sectional area of the beam, and $I(v, \vec{r})$ is CN fluence rate with respect to the neutron velocity v and cross sectional distribution of positions \vec{r} .

When the CN beam passes through the decay volume, the ^3He nuclei in superfluid helium-4 capture neutrons via nuclear reaction $n + ^3\text{He} \rightarrow p + t + 764$ keV. The detectable capture rate

is given as

$$\dot{N}_{p+t} = \epsilon_{\text{He3}} \sigma_{\text{He3}}^{\text{th}} v_n^{\text{th}} n_{\text{He3}} L \int_{A_b} da \int_v dv I(v, \vec{r}) \frac{1}{v}, \quad (2)$$

where ϵ_{He3} is the detection efficiency of the capture events, $\sigma_{\text{He3}}^{\text{th}}$ is the absorption cross section of ${}^3\text{He}$ nuclei for thermal neutrons at a velocity $v_n^{\text{th}} = 2200 \text{ m s}^{-1}$, and n_{He3} is the ${}^3\text{He}$ density. The neutron lifetime τ_β can be obtained from the ratio of the observed neutron ${}^3\text{He}$ capture rate to the beta decay rate.

$$\tau_\beta = \frac{\dot{N}_{p+t}}{\dot{N}_\beta} \cdot \frac{\epsilon_\beta}{\epsilon_{\text{He3}}} \cdot \frac{1}{\sigma_{\text{He3}}^{\text{th}} v_n^{\text{th}} n_{\text{He3}}} \quad (3)$$

Equation (3) is the key expression in this experiment. It explicitly shows the measurement of τ_β is independent of the neutron flux as well as the geometry of the decay volume. The overall accuracy relies on that of the observed ratio of event rates $\kappa = \dot{N}_{p+t}/\dot{N}_\beta$, the helium-3 density n_{He3} in superfluid helium-4, and the detection efficiency of scintillation events ϵ_{He3} and ϵ_β . The former two quantities will be experimentally acquired, and the detection efficiencies will be determined through simulations considering the calibration and background discrimination.

Here, we provide an estimate of count rates based on the published performance of the fundamental neutron physics beam line (FnPB) in the spallation neutron source (SNS) at Oak Ridge National Laboratory, as shown in figure 1 [33]. The neutron flux at 17 Å is about $2.4 \times 10^6 \text{ Hz } \text{\AA}^{-1} \text{ cm}^{-2} \text{ MW}^{-1}$. With a time-averaged proton power of 1.8 MW at 60 Hz of double-chopper, the incident rate of 17 Å neutrons with 0.5 Å pulse width is about $1.5 \times 10^7 \text{ Hz}$. In this estimate, neutron lifetime is taken as the Particle Data Group suggested value $\tau_\beta = 880 \text{ s}$ [34]. It takes about 3.2 ms for the 17 Å neutrons to pass 75 cm long decay volume, and the neutron decay probability is 3.66×10^{-6} for a CN beam with a cross section of 3 cm diameter. There are an average of 55.9 Hz of neutron decay events. The natural abundance of ${}^3\text{He}$ in liquid helium is $X_{\text{He3}} = 5 \times 10^{-7}$ in fractional concentration. Near isotopically pure ${}^4\text{He}$ with $X_{\text{He3}} < 2.5 \times 10^{-13}$ has been produced as reported by Hendry and McClintock [35]. Assuming superfluid helium-4 with $X_{\text{He3}} = 2 \times 10^{-10}$ can be prepared, an average of 252.1 Hz neutron capture events will occur simultaneously when the CN beam passes the decay volume.

3. Scintillation signals in liquid helium

The number of beta decays and neutron captures will be counted via scintillation signals in liquid helium. Liquid helium is an ideal scintillator that has been proposed and studied for detection of neutrino [36] and dark matter [37, 38]. Compared to those experiments, the detection volume in the current beam lifetime experiment is compact. High detection efficiency of scintillation photons can be practically achieved, as well as a wide energy window with good resolution and small lower bound.

In a neutron beta decay event, scintillation is generated by the recoiling electron, as the counterpart proton is too slow. About 35% of the total energy in each electron recoil above 1 keV creates He_2^* molecules of excited singlet state $\text{He}_2 (A^1 \Sigma_u^+)$ in liquid helium. The singlets radiatively decay in less than 10 ns and emit about 22 extreme ultraviolet (EUV) photons per keV of electron recoil energy KE_e , with a spectrum spanning from 13 eV to 20 eV and centering at 16 eV [39]. It forms the prompt pulse of scintillation light. There are about 1.7×10^4 photons per decay event at the end point energy of 782 keV in the neutron beta decay spectrum. On the other hand, the neutron capture is purely a nuclear recoil event, the scintillation process of

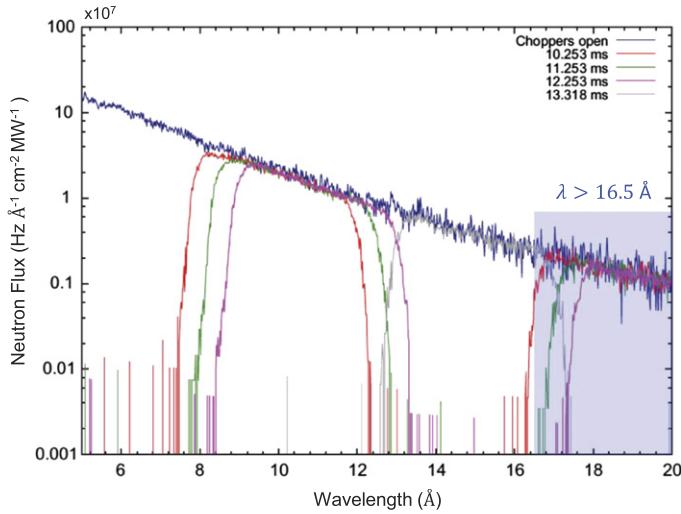


Figure 1. Spectrum of CN beam at the SNS FnPB beamline with choppers (a reprint of figure 5 in Fomin 2015 [33]). The portion of wavelength $\lambda > 16.5 \text{ \AA}$ is highlighted in blue shadows.

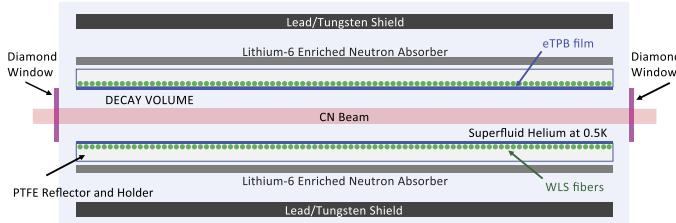


Figure 2. Schematic of the conceptual detector (non-scaled).

which is similar but of different features. About 13% of the recoil energy of 764 keV converts into a prompt light pulse and results in about 6.4×10^3 photons per capture event [40]. The stopping power dE/dx for a recoiling nucleus in liquid helium of a density $\rho = 0.145 \text{ g cm}^{-3}$ is $2 \times 10^4 \text{ eV } \mu\text{m}^{-1}$. The typical stopping range for a 800 keV recoiling nucleus is $40 \mu\text{m}$. By contrast, the stopping power for a 800 keV recoiling electron is only $40 \text{ eV } \mu\text{m}^{-1}$ [41] on average, and its stopping range can reach up to 2 cm. Therefore, a diameter of 7.5 cm is sufficient to prevent almost all of the recoiling electron born in the 3 cm diameter CN beam from touching the inner surface of the decay volume. It nearly guarantees no quenching of the prompt scintillation on the wall. Because of the dramatic difference in track length, the scintillation light of decay events is much more dispersed spatially than that of capture events. The former appears as a line of chained point sources, whereas the latter as a single point source.

In addition, both electron and nuclear recoils also generate a large amount of triplet He_2^* excimers ($a^3\Sigma_u^+$), which has a 13 s lifetime in liquid helium. The radiative decay of the triplet excimers is forbidden as it requires a spin flip; yet it can occur via the bimolecular Penning ionization that converts a portion of the triplet into singlet, most likely in a high density of

triplet excimers along the recoil track. This type of scintillation light appears as a large number of after-pulses of EUV photons, following the prompt pulse, and temporally scattered over tens of micro-seconds. Each of them is much weaker than the prompt pulse, and therefore mostly registered as pulses of single or a few photo-electrons (PE) in the same detector. The occurrence rate of after-pulses decreases as to a combination of two components dependent exponentially and inversely on time, respectively, $g(t) = Ae^{-t/\tau_s} + B/t + C$. It has been experimentally demonstrated that the $1/t$ component of the electron recoils is much weaker than that of the nuclear recoils [42, 43]. This feature offers an important tool to distinguish the decay events from the capture events.

4. Detection of scintillation

A standard method has been well developed to detect the EUV scintillation in liquid helium by many experiments [39, 40, 42–45]. Based on the known technologies, we describe a conceptual design as a baseline for a quantitative analysis. A schematic of the detector is shown in figure 2. The EUV scintillation light is first converted into a blue spectrum near 400 nm by an organic fluor—tetraphenyl butadiene (TPB). A thin layer of evaporated TPB (eTPB) can be coated on an acrylic film and wrapped into a cylinder as the boundary of the decay volume. The eTPB coating faces the inside of the decay volume. Optical fibers can be molded with a structural support as if wound on the outside of the film cylinder to collect light. The fibers cover the full length of the decay volume so as to maximize the light collection. Wavelength shifting (WLS) fibers are a common option to convert the emitted blue light into a green spectrum near 500 nm along with a redistribution of photon phase space. A portion of the shifted light can be trapped inside the fiber by total internal reflection and transmitted to the photon sensors.

The overall light conversion efficiency is estimated as follows. Owing to the large L/D ratio of the decay volume, more than 96% in solid angle of the scintillation light can be converted by TPB for the events occurring in the central region, as shown in figure 3. The conversion efficiency of eTPB has been demonstrated to be greater than unity [46]. Since a thick eTPB coating often appear opaque for visible light due to its surface roughness, the blue photons heading inwards the decay volume might reflect and diffuse on the eTPB coating. It is thus difficult to characterize the distribution of these inwards-going photons that are collected by WLS fibers upon multiple scattering in the eTPB coating. As a moderate estimate, we only take into account the 50% of eTPB re-emitted blue photons that travel outwards to the adjacent WLS fibers. Approximately 90% of them can impinge on the fiber cores with the help of a polytetrafluoroethylene (PTFE) reflector, which is also a structural holder clamped on the outside, and then about 80% is absorbed and shifted into green light. The double cladding WLS fibers made by Kuraray have a trapping efficiency of 5.4% in one direction [47]. When read on both ends, 10.8% of the shifted green light can be conveyed towards the 2 photon sensors. Since the fiber has a bending loss of about 4% per turn on a 7.5 cm diameter curve and an attenuation length longer than 7.5 m, the fiber length must be constrained. In each detector unit, a round WLS fiber of 1 m long and 1 mm diameter is helically wound around the decay volume by 3 turns, and the extra length on each free end is routed to a separate photon sensor. It needs 250 units in a tight packing to cover the whole length of the decay volume and set up an axial resolution. The average transmission efficiency of such a configuration is about 90% along the fiber. With regard to the difficulty of making large amount of superfluid-leak-tight fiber feedthroughs, there must be two optical breaks at the windows of the liquid helium vessel, each of which has a 90% transmission. As for the 500 photon sensors, we may employ silicon photomultipliers, which are compact in size and have a typical quantum efficiency of 34% for the versions with large microcells [48]. The overall conversion efficiency η_{tot} is 0.9%, i.e. an

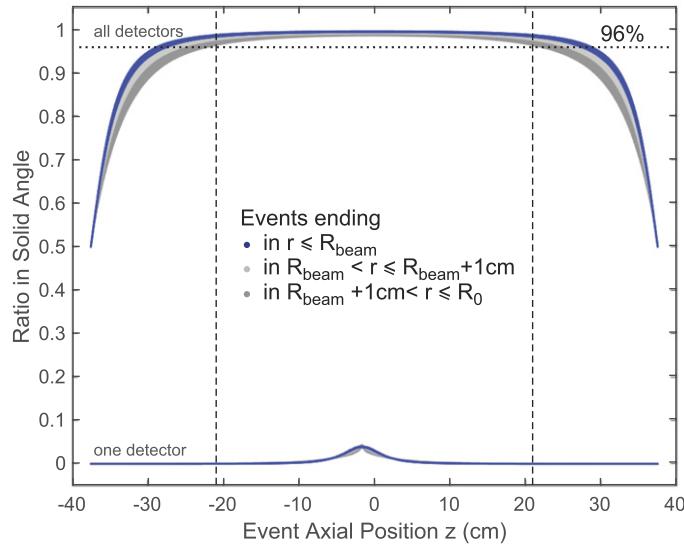


Figure 3. The coverage of solid angle for events ending at different axial positions in the decay volume. The upper plot is the total coverage of solid angle by the sum of all the detectors; and the lower plot is that by one unit detector centered at -1.65 cm. Events ending in the beam pass ($r \leq R_{\text{beam}}$) are marked with blue, within 1 cm away from the beam pass ($R_{\text{beam}} < r \leq R_{\text{beam}} + 1 \text{ cm}$) marked with light grey, and outside the regions above ($R_{\text{beam}} + 1 \text{ cm} < r \leq R_0$) marked with dark grey, where R_0 is the maximum radius that a recoiling electron can reach.

average of 9 PE can be detected per 1×10^3 EUV scintillation photons. The average prompt PE numbers for decay events with the spectrum peak energy at 245 keV and the endpoint energy at 782 keV are $\bar{N}_{\text{PE}}^{(\text{peak})} = 49.8$ and $\bar{N}_{\text{PE}}^{(\text{endpt})} = 158.9$, respectively. Every PE corresponds to about 5 keV of electron recoil energy. On the other hand, the average prompt PE number for the neutron capture events of recoil energy at 764 keV is $\bar{N}_{\text{PE}}^{(\text{p+})} = 57.7$. It coincides with beta events of 283.8 keV, close to the peak of the beta spectrum.

5. Detector response and event reconstruction

We perform a preliminary study on the response of detectors by Monte Carlo simulations. As listed below, several assumptions have been adopted to simplify the model but present the essential physics as a proof of principle. Further modelling with more details is needed.

- Only the prompt scintillation signals are recorded for all the events. It means the decay and capture events cannot be distinguished among the simulated data. Yet in real experiment, they are distinguishable by the difference in the $1/t$ time-dependent occurrence rates of the after-pulses. This additional information will improve data analysis and understanding of systematic effects.
- The scintillation light for the capture events is emitted from a point source as their track length is tens of microns, whereas that for the beta events is from an energy-dependent straight tracks of length up to 2 cm. For electron recoils, more energy deposits in the vicinity of the track end as it slows down. The spatial energy deposition approximately follows $dE/dr \propto r^2$, where r is the geometric distance from the starting point of electron recoil

[49]. Further studies can be performed on simulated scattering tracks with productions of secondary δ -electrons.

- (c) Only the outward-going portion of the eTPB converted light can be collected by the adjacent fibers, but none of the inward-going, as the latter reflects and diffuses on the coating into a broader distribution over all the detectors, yet much weaker in intensity than the former.
- (d) The EUV light converted by eTPB will be collected by the fibers tightly wound against the thin film at the same axial position. It means the solid angle of light from an event projecting on the section of the eTPB film is equivalent to that on the detector lying against the film.
- (e) All the detectors have the same efficiency. In reality, the efficiency of detectors are different and may vary with time. Calibrations are necessary and will be discussed in subsection 5.3.
- (f) There is no timing information in this simulation. We assume all the events are in the coincident time window as the CN beam passes the decay volume. Timing information will greatly improve the position reconstruction and be a practical way of minimizing background.

A sampling on starting position \vec{x}_0 of events as to a uniform distribution function, $\Pr(\vec{x}_0) = \text{const}$, is carried out in the beam-occupied volume¹. 4×10^7 random events are generated and assigned as either capture or decay according to a preset ratio, $\kappa = 4.4975$. The capture events are of a point source in the beam-occupied volume, whereas the decay events are of a straight line source that may extend out. For each event, emission of scintillation light is isotropic, and the portion of solid angle received by each and every of the 250 detectors is simulated as the hit probability of each detector given the source position \vec{x} , $\Pr(n_{\text{det}}|\vec{x})$, where $n_{\text{det}} \in [1, 250]$ is the detector number. Figure 3 plots the accepted portion of solid angle by one unit detector centered at -1.65 cm, and sum of that over all the detectors for each event at various axial positions z , respectively. The events ending in the beam path are highlighted in blue. It shows the total accepted solid angle has a weak dependence on the radial ending position, which correlates with the recoil energy of electrons. With all the detectors functioning, it can cover more than 96% of solid angle for events in a central region spanning 42 cm in the axial direction, as shown by the upper plot in figure 3. The events lying within the axial edges of a unit detector have about 4% of chance to be registered by this detector, as shown in the lower plot in figure 3. The overall detector hit probability $\Pr(n_{\text{det}})$ is derived by the integral over the entire volume V_0 , $\Pr(n_{\text{det}}) = \int_{V_0} d\vec{x} \Pr(n_{\text{det}}|\vec{x}) \Pr(\vec{x})$, and plotted in the figure 4. For detectors in the central region, the overall hit probability is about 0.40% for scintillation light of an event at any position to be registered; while for those close to the ends, the chance is naturally much less.

5.1. Reconstruction of event position and the fiducial cut

With the registered PE numbers $N_{\text{PE}}(n_{\text{det}})$ from an event on a series of detectors n_{det} , $1 \leq n_{\text{det}} \leq 250$, the probability of event position within a given sub-volume $\Delta V(z_k)$ centered at z_k is derived in equation (4) by the Bayes' theorem. It is on the assumption that every detector is independent, i.e. no cross talk. The probability of an event within a certain region of interest, e.g. the 42 cm long central region, can be calculated as the accumulated probability, $\sum_{z_k} \Pr(\Delta V(z_k)|N_{\text{PE}}(n_{\text{det}}))$. Figure 5(a) shows an example of the simulated electron recoil event with 77 observed PE's distributed on several detectors. The entire volume V_0 is divided into

¹ The cross sectional distribution of the beam fluent rate $I(v, \vec{r})$ is independent of \vec{r} in this simulation.

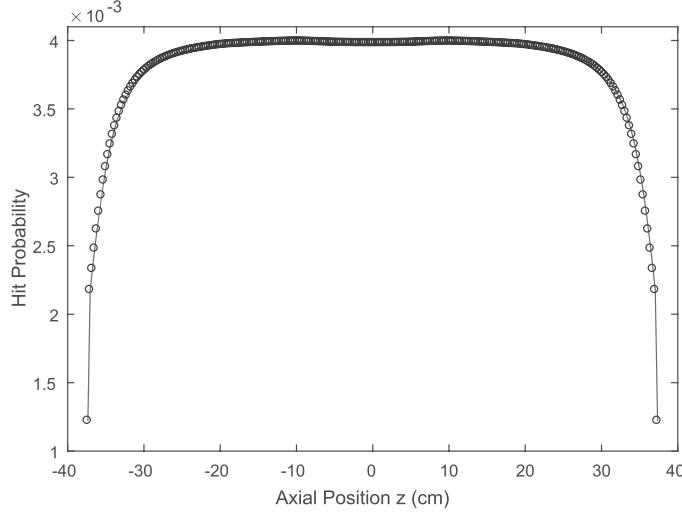


Figure 4. The overall hit probability on each of the 250 detectors.

sub-volumes $\Delta V(z_k)$ as disks of 1 cm thick, and the distribution probability of event position for each sub-volume is calculated and plotted in figure 5(b). The accumulated probability for this event to occur within the 42 cm long central region is 93.7%. In general, the more PE's observed, the more accurate the reconstructed position of the event. However, since the events near the ends of the decay volume lose a significant portion of scintillation light on the end windows, i.e. information is truncated, the reconstructed positions are biased towards the center. Therefore, the accumulated probability of positions inside the central region of 42 cm generally performs better in identifying events thereof than that of reconstructed positions, especially for low PE events. Yet neither is satisfactory in selecting events in a region with relatively identical position distribution and uniform ratio of capture-to-decay event rates. Later, we find a combination of both actually forms a good fiducial cut. It is demonstrated in figure 6. Figures 6(a) and (b) show the events are selected by a combination of the following criteria, (i) reconstructed positions within ± 15 cm, and (ii) accumulated probability of more than 80% inside the central region of 42 cm. The number of selected events is about 42% of the total. Figure 6(c) shows the distributions of the original axial positions between the decay and capture events are almost identical; and figure 6(d) shows the capture-to-decay ratio κ has a flat plateau and relatively sharp edges in the selected region with respect to the original axial positions.

$$\begin{aligned} \Pr(\Delta V(z_k) | N_{\text{PE}}(n_{\text{det}}), 1 \leq n_{\text{det}} \leq 250) \\ = \int_{\Delta V(z_k)} d\vec{x} \frac{\sum_{n_{\text{det}}=1}^{250} N_{\text{PE}}(n_{\text{det}}) \Pr(n_{\text{det}} | \vec{x}) \Pr(\vec{x})}{\sum_{n_{\text{det}}=1}^{250} N_{\text{PE}}(n_{\text{det}}) \Pr(n_{\text{det}})} \quad (4) \end{aligned}$$

5.2. Determination on the ratio of capture-to-decay event rates

The spectrum of the selected events is plotted against the PE number in figure 7. The capture events overlap with the decay events in the mid range and the spectrum cuts off at a lower bound of 4 PE's, equivalent to a recoil energy $KE_e = 20$ keV, on purpose to exclude random backgrounds of few photons. In order to resolve the ratio κ of capture-to-decay event rates from

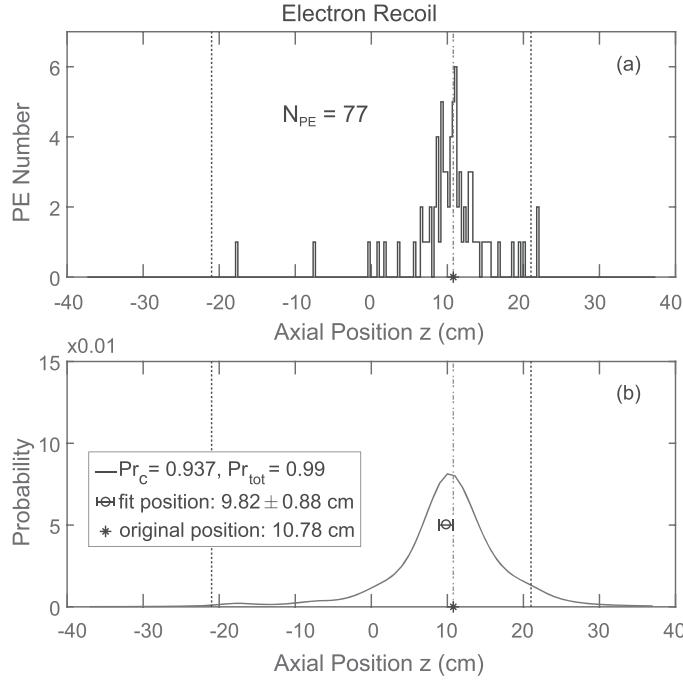


Figure 5. An example on the deduced possibility of event axial position based on the observed distribution of 77 PE's. (a) The spatial distribution of observed PE numbers on different detectors; (b) the spatial distribution of probabilities on different event axial positions. The total probability Pr_{tot} is 0.99 and the probability in the central region Pr_c is 0.937. The fit position based on the observed PE's is 9.82 ± 0.88 cm, and in comparison, the original position fed in the simulation is 10.78 cm.

the acquired spectrum, a theoretical model for fitting is constructed. It consists of three components: the neutron decay spectrum, the single capture peak and the background. In this study, we only simulate signals of the former two, but omit the effect of the background, because it will be poorly defined without the knowledge of the actual system. A discussion on the possible backgrounds will be presented in section 6.

The neutron beta decay spectrum is formulated as

$$\frac{d\Gamma_n}{dE_e} \propto F_n(E_e) p E_e (E_e^{(\text{endpt})} - E_e)^2 \quad (5)$$

where $E_e^{(\text{endpt})}$ is the endpoint energy, E_e is the total electron energy, and $p = (E_e^2 - m_e^2)^{1/2}$ is the momentum of electron. m_e is the electron mass. $F_n(E_e)$ is the Fermi function for neutrons defined as

$$F_n(E_e) = 4 \exp(\pi\alpha E_e/p) \frac{|\Gamma(1 - i\alpha E_e/p)|^2}{\Gamma(3)^2} \quad (6)$$

where $\alpha = 1/137$ is the fine structure constant, and $\Gamma(z)$ is the gamma function. In reality, the spectrum contains a Poisson smearing due to the random process in the light transportation and conversion into PE's. Though the deposited energy from neutron capture events is single-valued at 764 keV, it appears as a much broadened peak. The Poisson probability function,

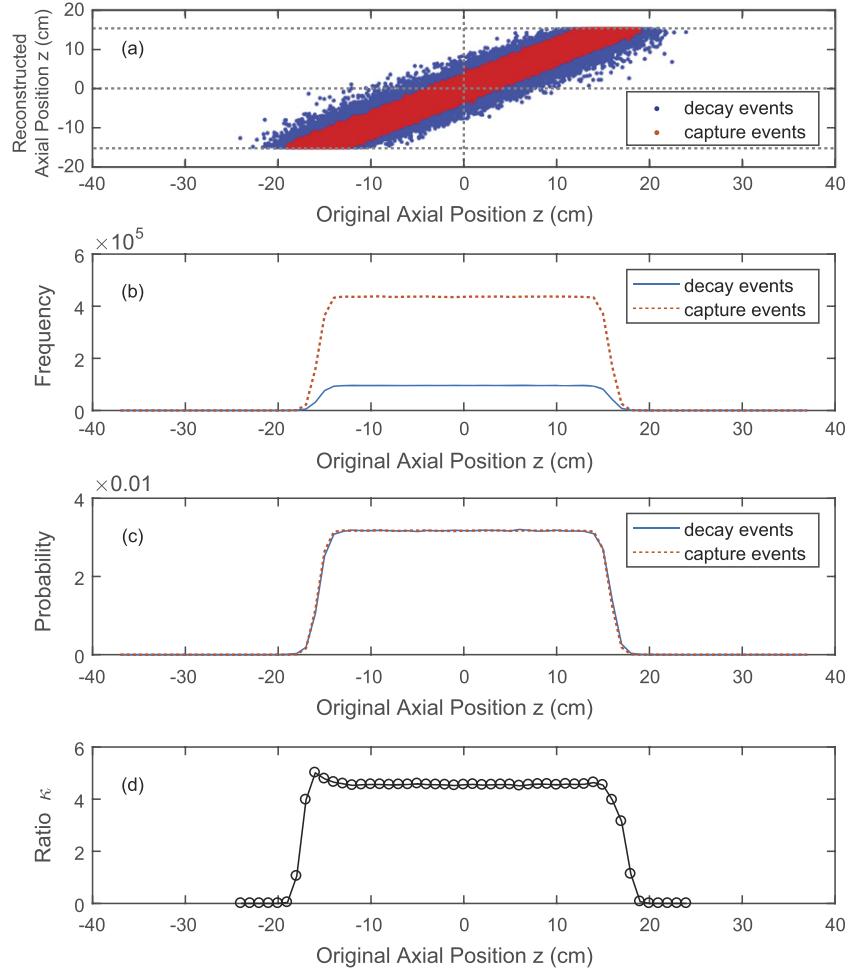


Figure 6. Analysis of the combined fiducial cut on the simulated data. (a) The reconstructed axial position z of both the simulated decay events (blue dots) and capture events (red dots) plotted against their original position; (b) and (c) frequency and probability distribution of the selected decay events (blue solid line) and capture events (red dashed line) by the fiducial cut as to their original positions, respectively; (d) the ratio κ calculated within each bin of original position z .

$\text{Pois}[N_{\text{PE}}, \lambda]$, with the PE number N_{PE} and the mean λ , is thus embedded in the probability function $f(N_{\text{PE}})$ of the observed spectrum as

$$f(N_{\text{PE}}) = \frac{1}{1 + \tilde{\kappa}} \frac{\epsilon_{\text{cut}}(N_{\text{PE}})}{C_e} \sum_{M_{\text{PE}}} \text{Pois}[N_{\text{PE}}, M_{\text{PE}}] \cdot \mathcal{F}_e(M_{\text{PE}}) + \frac{\tilde{\kappa}}{1 + \tilde{\kappa}} \frac{1}{C_{p+t}} \text{Pois}[N_{\text{PE}}, \lambda_{p+t}] + B \quad (7)$$

where $\tilde{\kappa}$ is the observed ratio of capture-to-decay event rates, C_e and C_{p+t} are the normalization factors for the spectrum bins above the lower bound of 4 PE's. The third term B represents

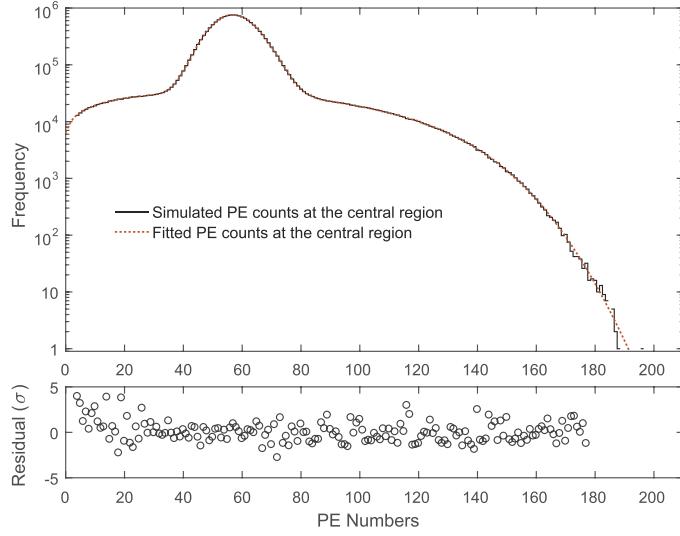


Figure 7. Simulated neutron beta decay spectrum in addition to the neutron capture peak (black solid line) and the result of the maximum likelihood (ML) fit (red dotted line). The residual of fitting with respect to the PE numbers is shown in the lower plot $\chi^2/ndf = 1.4$.

the background. The second term is the broaden peak of the neutron capture events, where λ_{p+t} is the mean PE number for the capture events. The first term is the modified neutron beta decay spectrum. In a Poisson process, events of different recoil energies and their corresponding deterministically-converted PE number M_{PE} all contribute to a given bin of observed PE number N_{PE} . The share of contribution from each M_{PE} is actually the probability of events with a deterministic M_{PE} in the neutron beta decay spectrum that follows equation (5), namely $\mathcal{F}_e(M_{\text{PE}})$, and given by

$$\mathcal{F}_e(M_{\text{PE}}) = \int_{(M_{\text{PE}}-0.5)/\eta_e\Omega(M_{\text{PE}})}^{(M_{\text{PE}}+0.5)/\eta_e\Omega(M_{\text{PE}})} \frac{d\Gamma_n}{dE_e} dE_e \quad (8)$$

where $\eta_e \cdot \Omega(M_{\text{PE}})$ is the KE_e -to-PE's conversion coefficient for electron recoils. As indicated in figure 3, the coverage ratio of solid angle $\Omega(M_{\text{PE}})$ has a weak dependence on the electron recoil energy, i.e., on the PE number M_{PE} in the central region of the decay volume. The particular coefficient $\epsilon_{\text{cut}}(N_{\text{PE}})$, in the first term of equation (7), is the efficiency of detection related to the fiducial cut. It comes from a significant position uncertainty due to few PE's and long tracks for the low and high PE events, respectively. It inevitably introduces a detection inefficiency, $1 - \epsilon_{\text{cut}}(N_{\text{PE}})$, and hence, a distortion of spectrum. Such an effect is extracted from the simulation and fitted as shown in figure 8.

There remain 3 fitting parameters in equation (7): η_e , λ_{p+t} , and most importantly, $\tilde{\kappa}$. Maximum Likelihood (ML) method is employed to acquire the best fit and associated errors of the above 3 parameters. κ in equation (7) is marked as $\tilde{\kappa}$, as the fitting only acquires the observed capture-to-decay ratio within the data bins. $\tilde{\kappa}$ must be corrected for the actual κ with the detection efficiencies $\epsilon_{\text{He3}} = 1$ and $\epsilon_{\beta} = 0.9878$, where $\epsilon_{\beta} < 1$ is due to the missing portion of spectrum less than 4 PE's. The result is listed in table 1, and the plots in figure 7 show the

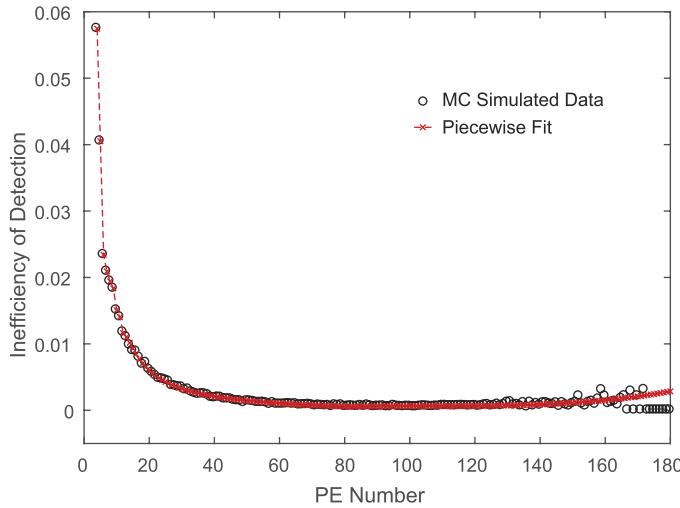


Figure 8. Inefficiency of detection, $1 - \epsilon_{\text{cut}}(N_{\text{PE}})$, related to the combined fiducial cut of reconstructed positions within ± 15 cm and 80% accumulated probability within ± 21 cm. The simulation data is plotted in black circles, and the piecewise fit is plotted in red dashed line. Piecewise fit can better capture the features of inefficiency as a sharp rise in low PE region and a slow tilt in the high PE region.

Table 1. Result of ML fit to the simulated spectrum.

	$\kappa = \Delta N_{p+t} / \Delta N_{\beta}$	η_e	λ_{p+t}
Preset value	4.4975	0.2032	57.254
Monte Carlo	4.4997	0.2032	57.292
ML fit	4.4983 ± 0.0040	0.2037 ± 0.0001	57.291 ± 0.002

result of the ML fit (in red dotted line) on top of the acquired spectrum (in black solid line). It demonstrates the fit can extract κ value at an accuracy well within 0.1%.

5.3. Reconstruction of event energy and detector calibration

Reconstruction of event energy will be performed in calibration of all the detectors as a cross reference. In the previous subsection, the conversion efficiency η_e for electron recoils is fitted through the analysis of the neutron decay spectrum, and the reconstructed energy is obtained as $E_e = N_{\text{PE}} / \eta_e \Omega(N_{\text{PE}})$. In reality, each photon sensor has a different quantum efficiency $\eta_{\text{SiPM}}^{(i)}$, and each fiber has a variation in transmission efficiency $\eta_{\text{fiber}}^{(i)}$. The overall detection coefficient $\eta_{\text{tot}}(x)$ thus varies for different sub-volumes, due to the variation of solid angles upon detectors of different quantum efficiency. It is a common approach to use conversion electron sources, such as ^{109}Cd (63, 84 keV), ^{139}Ce (127, 160 keV), ^{113}Sn (364, 388 keV), ^{207}Bi (481, 975, 1047 keV) for calibrations in between production runs [50]. These sources can be placed in many designated positions to map out the response of different detectors. During the production runs, the calibration can also be done with the neutron capture peak, and additional deposits of α source or lithium neutron capture film on the end windows of the decay volume.

6. Background suppression

In order to achieve a highly accurate measurement on the ratio of capture-to-decay rates, background signals must be properly suppressed, discriminated or subtracted. Cosmic ray muons can be easily identified by coincidence in the veto detectors surrounding the apparatus. The static radioactive backgrounds from materials of the apparatus can be shielded by a thick layer of lead or tungsten as shown in figure 2 and characterized in the background runs. The gamma rays from the CN source can be greatly suppressed by bending the beam direction out of sight with proper neutron optics [33]. The most harmful type of backgrounds are the gamma rays produced by the neutron-induced activation near the decay volume and undergoing Compton scattering on liquid helium inside the decay volume and the surrounding WLS fibers. The Compton electrons are identical to the decay events on the features of scintillation signals. It hence will be pooled in the PE spectrum and modelled as B in equation (7). The spectrum of the Compton electrons produced by gamma rays above 4 MeV is mostly flat in the region of neutron decay spectrum. Some of the delayed gamma rays can be characterized during the intervals between the CN beam pulses, such as the 1.6 MeV gamma rays emitted at a half-life of 11.16 s from the neutron activated fluorine. In this paper, we focus on two types of prompt gamma rays due to the neutron captures by the window material, and by hydrogen, $p + n \rightarrow d + \gamma$ (2.2 MeV), in the plastic components surrounding the decay volume, such as the WLS fibers, etc. They are believed to be the major contributors to the backgrounds.

The first measure to suppress neutron-induced gamma rays is to reduce the capture and scattering of neutrons on the window materials. Polycrystalline CVD diamond is a good option, because carbon has relative small scattering and capture cross sections, and a thin window of 5 cm diameter and 1 mm thickness [51] is commercially available with a good mechanical strength. The capture cross section of carbon for the 17 Å CN's is 0.033 barns and the capture fraction is 5.66×10^{-4} . For a CN flux of 1.5×10^7 CN s⁻¹, about 8.6×10^3 Hz of neutrons are captured with an emission of prompt gamma rays mostly at energies of 1.3, 3.7 and 4.9 MeV. A simulation shows the intensive prompt gamma rays result in more than 600 Hz Compton events in liquid helium inside the decay volume, and more than 100 Hz in the polystyrene WLS fibers. Most events distribute spatially near the windows, and temporally at the moments when the neutron flux passes the windows. Although the number of the window-originated Compton events greatly overwhelms that of the decay events, they can be separated in time if the neutron beam can be chopped into sharp pulses both in time and energy. It requires the decay volume to be set up close to the source. The neutron decays will then appear as scattered single events in time sequence between two intensive bursts of Compton events when neutrons pass the entrance and exit windows, respectively. Considering it takes about 3.2 ms for the 17 Å neutrons to pass the 75 cm long decay volume, the middle 1.6 ms is the time interval when the beam pulse passes the central region. The typical recovery time for the SiPM sensors are hundreds of nano-seconds [48], whereas the rate of after pulses following each prompt signal decays in tens of micro-seconds [42]. Therefore, such a time cut can effectively distinguish the events occurring in the central region, which are crucial in construction of the energy spectrum, and eliminate impacts of the window-originated Compton events. In return, the bright bursts of Compton events can be used as a calibration reference of the beam flux and spectrum.

The scattering of CN's on the windows at 0.5 K is dominated by incoherent scattering, which is an s-wave scattering independent of the incident velocity. The incoherent cross section of carbon for the 17 Å CN's is 0.001 barns, and the scattered fraction is 1.76×10^{-5} . For the same CN flux as above, about 270 Hz of neutrons are scattered isotropically from both the windows into the decay volume and interact with hydrogens on the fibers. The incoherent scattering cross section of hydrogen is 80.3 barns, much larger than the capture cross section of 3.1 barns

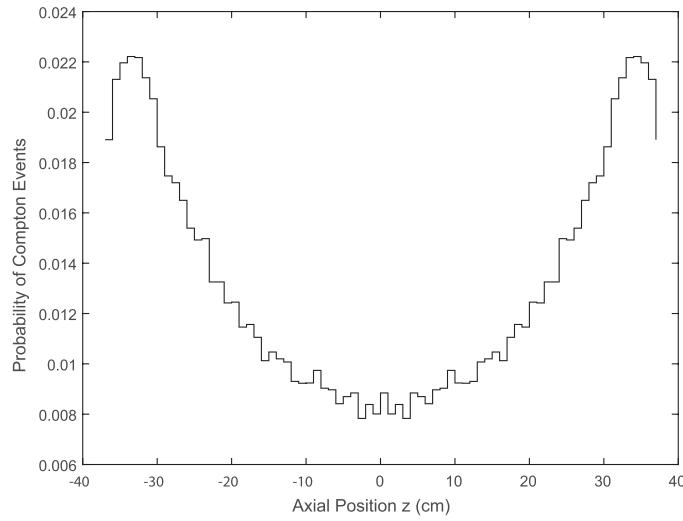


Figure 9. Probability of Compton electron events in the axial axis of the decay volume induced by the prompt gamma rays due to hydrogen-captures of neutrons scattered from the beam inlet and outlet diamond windows.

for 17 Å CN's. The stray neutrons will mostly scatter incoherently in the plastics and has a small chance of being captured by the hydrogen nuclei. In order to minimize the chance of captures, a second measure is to deploy an effective neutron absorber on the outside of the plastics so as to capture all the outgoing stray neutrons. Lithium-6 enriched material is an ideal option, since ${}^6\text{Li}$ has a large neutron absorption cross section of 8887 barns for 17 Å CN's, and there is no associated emission of gamma rays in the reaction, ${}^6\text{Li} + n \rightarrow \alpha + t + 4.78 \text{ MeV}$.

A simulation on neutron scattering and capture is carried out on the geometry of a 1 mm thick polystyrene fibers around the decay volume, a 5 mm thick PTFE holder clamped on the fibers, and a sufficiently thick lithium absorber at the outermost shell that absorbs all the stray neutrons, as illustrated in figure 2. It is found about 4.3% of the scattered neutrons are captured by the fibers and 0.4% by the PTFE holder. The neutron captures on hydrogen in the fibers do not induce any significant scintillation in the polystyrene, as the kinetic energy of deuterium is merely about 1.3 keV. The resultant prompt gamma rays of 2.2 MeV contribute a background of Compton events at 0.28 Hz in liquid helium inside the decay volume, and 0.31 Hz in the polystyrene WLS fibers. As shown in figure 9, the Compton events have a higher chance to occur near the windows. Since the majority of Compton electrons are at energies near the Compton peak of 2.0 MeV, it adds about 0.05% to the total counts of neutron decay events with a fiducial cut of ± 15 cm on the central region. It can be characterized and corrected in data analysis.

7. Accurate measurement of helium-3 density with UCN's

The last quantity crucial to determine the neutron lifetime in equation (3) is the ${}^3\text{He}$ density $n_{\text{He3}} = 2.18 \times 10^{22} X_{\text{He3}}$, where the fractional concentration X_{He3} of about 2×10^{-10} needs to be prepared and characterized in high precision. One possible way of measurement is to employ the atom trap trace analysis, which has been demonstrated to measure the abundance of rare isotope ${}^{39}\text{Ar}$ at the level of 10^{-16} [52]. Similar technology may be developed for detection of the

^3He concentration at a precision well below 0.1%. Then, the uncertainty of the combined term $\sigma_{\text{He}3}^{\text{th}} v_n^{\text{th}} n_{\text{He}3}$ in equation (3) will be dominated by that of the ^3He capture cross section of thermal neutrons, $\sigma_{\text{He}3}^{\text{th}} = 5333 \pm 7 \text{ b}$ [53], equivalent to 0.13%. However, in this paper, we provide another possible method to measure the ^3He density, which takes advantage of the scintillation rates correlated to the neutron captures on ^3He nuclei in the sample liquid helium. These deployed neutrons are not CN's in a beam but rather UCN stored in a material bottle filled with the sample liquid helium. Since UCN's uniformly distribute in the storage volume and scintillations can partially quench on the walls, it is impractical to construct a well-defined spectrum as in subsection 5.2, and therefore, the difference in the $1/t$ responses of after-pulses will be the key tool to distinguish the decay and capture events. We will characterize the purity of ^4He and accurately measure ^3He concentration in the later prepared helium mixture via the time-dependent rates of both decay and capture scintillation events. Furthermore, it will be shown that value of the combined term $\sigma_{\text{He}3}^{\text{th}} v_n^{\text{th}} n_{\text{He}3}$ will be directly obtained in experiment. Therefore, our goal sensitivity in neutron lifetime of below 0.1% can be achieved with enough statistics, and is no longer limited by the uncertainty of 0.13% in $\sigma_{\text{He}3}^{\text{th}}$. Details will be articulated in the following subsections.

7.1. UCN storage in a neutron-friendly volume

Suppose the sample liquid helium fills a storage volume made of UV transmitting acrylic tube of 75 cm long, 7 cm ID and 7.5 cm OD. It is sealed at both ends and coated with deuterated films on the inside so that it is hermetic and friendly to UCN's. It can be installed inside the detector setup as described in section 4, except the PTFE reflector can be as thin as 50 μm because of the following two reasons: (i) the structural support can be loaded to the storage tube; and (ii) PTFE generates a high level of background due to the delayed gamma rays from neutron-activated fluorine as discussed in subsection 7.6.

The UCN storage follows

$$\dot{N}_{\text{UCN}}(t) = -\frac{N_{\text{UCN}}(t)}{\tau_{\text{tot}}}, \quad (9)$$

$$N_{\text{UCN}}(t) = N_0 \exp\left(-\frac{t}{\tau_{\text{tot}}}\right), \quad (10)$$

where N_0 is the initial number of UCN's and τ_{tot} is the storage time constant. Several factors contribute to τ_{tot} of this volume,

$$\tau_{\text{tot}}^{-1} = \tau_{\text{He}3}^{-1} + \tau_{\beta}^{-1} + \tau_{\text{up}}^{-1} + \tau_{\text{loss}}^{-1}. \quad (11)$$

$\tau_{\text{He}3}$ is the neutron ^3He capture time constant of interest as given in equation (12). For $X_{\text{He}3} = 2 \times 10^{-10}$, $\tau_{\text{He}3}$ is about 195.2 s, which dominates the total storage time, compared to the neutron lifetime $\tau_{\beta} \approx 880$ s.

$$\frac{1}{\tau_{\text{He}3}} = n_{\text{He}3} \sigma_{\text{He}3}^{\text{th}} v_{\text{He}3}^{\text{th}} = 2.56 \times 10^7 X_{\text{He}3} (\text{s}^{-1}) \quad (12)$$

UCN's suffer a loss from captures or up-scattering of the wall nuclei. Such an effect can be described by Schrödinger equations with one-dimensional potential and characterized by a loss probability per bounce, $f(E_{\text{UCN}})$. The rate of wall collisions is given by $(vA/4V)$, where v is the UCN velocity, A is the area of the storage chamber, and V is its volume. The contribution

of wall losses to the storage time is then given as

$$\frac{1}{\tau_{\text{wall}}} = f(E_{\text{UCN}}) \left(\frac{vA}{4V} \right). \quad (13)$$

Generally, the hydrogen in the organic materials has a large up-scattering cross section for the UCN's. Therefore the hydrogen in contact with UCN must be replaced with deuterium. The inner wall of the storage volume needs to be coated with a layer of deuterated polystyrene, whose Fermi potential is about 160 neV [54]. The TPB converter coating also needs to be deuterated. These technologies are under development and tests by the SNS nEDM collaboration [55]. With a loss probability per bounce of $f(E_{\text{UCN}}) = 10^{-5}$, the same requirement as the SNS nEDM UCN storage cells [56], the time constant due to wall loss τ_{wall} is about 1672.0 s.

τ_{up} is the loss rate due to upscattering of neutrons by quasi-particles, phonons and rotons, in superfluid helium. It is greatly suppressed by a Boltzmann factor. At $T < 0.6$ K, the dominant process is multi-phonon scattering following

$$\frac{1}{\tau_{\text{up}}} = \frac{T^7}{100} [\text{s}^{-1}]. \quad (14)$$

At $T = 0.5$ K, $\tau_{\text{up}} = 12800$ s. With all the contributions above included in equation (11), the storage time constant τ_{tot} is about 144.2 s. This is merely an estimate. τ_{tot} will be accurately measured in experiments in order to acquire a high-precision determination on the ${}^3\text{He}$ concentration $X_{\text{He}3}$ in subsection 7.5.

7.2. UCN production

The UCN's for this measurement are produced *in situ* in the neutron decay volume by super-thermal process: a 8.9 Å beam of CN's are down-scattered inelastically into UCN's via exciting a single phonon in superfluid helium [57]. The UCN density can build up in the decay volume with the time constant τ_{tot} . The accumulated UCN density in liquid helium exposed to the CN beam is given by

$$\rho_{\text{UCN}}(t_{\text{fill}}) = R\tau_{\text{tot}} \left[1 - \exp \left(-\frac{t_{\text{fill}}}{\tau_{\text{tot}}} \right) \right]. \quad (15)$$

τ_{tot} is the storage time constant as given in equation (11), and the production rate per unit volume R is given by

$$R = 2.2 \times 10^{-8} \left(\frac{d\Phi}{dE} \right) [\text{cm}^{-3} \text{ s}^{-1}] \quad (16)$$

where an incident flux spectrum of $(d\Phi/dE)$ is in units of $(\text{cm}^{-2} \text{ s}^{-1} \text{ Å}^{-1})$, and the production of UCN's is up to the maximum storage Fermi energy of 160 neV [57, 58]. As shown in figure 1, the 8.9 Å CN flux in the SNS FnPB is about $5.9 \times 10^7 \text{ Hz cm}^{-2} \text{ Å}^{-1}$. The production rate per unit volume is deduced to be $R \approx 1.3 \text{ UCN's cm}^{-3} \text{ s}^{-1}$. So the steady state UCN density ρ_{UCN} can reach an average of about 141.4 UCN's cm^{-3} with 200 s of beam filling. i.e. a total number $N_0 \approx 7.5 \times 10^4$ of UCN's can be filled in the storage cell with a beam-occupied volume of 530.1 cm^3 . However, the initial UCN filling number N_0 varies in different runs depending on the CN beam intensity and stability, and is of little use in the data analysis due to the large uncertainty. The estimate above is for the sake of presenting the order of magnitude.

7.3. A shortcut to examine neutron lifetime enigma

The observed capture and decay rates of UCN's in the storage volume as well as their ratio are given by

$$\dot{N}_{\text{UCN}}^{(p+t)}(t) = -\epsilon'_{\text{He3}} \frac{N_0}{\tau_{\text{He3}}} \exp\left(-\frac{t}{\tau_{\text{tot}}}\right), \quad (17)$$

$$\dot{N}_{\text{UCN}}^{(\beta)}(t) = -\epsilon'_{\beta} \frac{N_0}{\tau_{\beta}} \exp\left(-\frac{t}{\tau_{\text{tot}}}\right), \quad (18)$$

$$\tilde{\kappa}_{\text{UCN}} = \frac{\dot{N}_{\text{UCN}}^{(p+t)}(t_1)}{\dot{N}_{\text{UCN}}^{(\beta)}(t_2)} = \frac{\epsilon'_{\text{He3}}}{\epsilon'_{\beta}} \frac{\tau_{\beta}}{\tau_{\text{He3}}} \exp\left(-\frac{t_1 - t_2}{\tau_{\text{tot}}}\right), \quad (19)$$

where ϵ'_{He3} and ϵ'_{β} are the detection efficiencies of UCN capture and decay events that can be obtained via simulation. Since the ${}^3\text{He}$ capture events can be distinguished from the decay events via the difference in the $1/t$ decay rate of after-pulses, $\dot{N}_{\text{UCN}}^{(p+t)}(t)$ and $\dot{N}_{\text{UCN}}^{(\beta)}(t)$ can be directly acquired in the measurement. The ratio $\tilde{\kappa}_{\text{UCN}}$ can then be calculated as in equation (19), by counting both event rates in the identical time bins, i.e. $t_i = t_1 = t_2$,

$$\tilde{\kappa}_{\text{UCN}}(t_i) = \frac{\epsilon'_{\text{He3}}}{\epsilon'_{\beta}} \frac{\tau_{\beta}}{\tau_{\text{He3}}} = \frac{\epsilon'_{\text{He3}}}{\epsilon'_{\beta}} \tau_{\beta} n_{\text{He3}} \sigma_{\text{He3}}^{\text{th}} v_{\text{He3}}^{\text{th}}. \quad (20)$$

$\tilde{\kappa}_{\text{UCN}}(t_i)$ should statistically fluctuate around its true value. In this scenario, the neutron lifetime enigma can be examined by simply comparing $\tau_{\beta} n_{\text{He3}}$ obtained above by equation (20) in the UCN storage volume with that by equation (3) in the beam decay volume. The real value of ${}^3\text{He}$ density n_{He3} is no longer necessary. But there is a caveat: the success of this ‘shortcut’ trick greatly depends on how well one can characterize the background events, especially the Compton events induced by the gamma rays, and separate them from decay events through analysis and modelling.

7.4. Characterization on the purity of helium-4

The production of isotopically pure ${}^4\text{He}$ can be carried out in a purifier similar to that designed by Hendry and McClintock [35]. For a residual ${}^3\text{He}$ concentration, $X_{\text{He3}}^{(0)} < 0.1\% X_{\text{He3}} \approx 2 \times 10^{-13}$, the ${}^3\text{He}$ capture time constant $\tau_{\text{He3}}^{(0)}$ is expected to be more than 2×10^5 s. With the property of the UCN storage volume described in subsection 7.1, the total storage time τ_{tot} is expected to be about 551.7 s, and the UCN density $\rho_{\text{UCN}}^{(0)}$ can reach about $219.2 \text{ UCN's cm}^{-3}$. With 200 s of beam filling, it may achieve a fill of $N_0^{(0)} \approx 1.16 \times 10^5$ UCN's in the storage volume.

Because of the scarcity of the ${}^3\text{He}$ atoms in the isotopically pure ${}^4\text{He}$ liquid, a total number of neutron capture events $\Delta N_{\text{UCN}}^{(p+t,0)}$ may be counted over a period of storage time Δt via identification on the after-pulses of scintillation events. So is the total number of the neutron decay events $\Delta N_{\text{UCN}}^{(\beta,0)}$. Integration on equations (17) and (18), respectively, gives

$$\Delta N_{\text{UCN}}^{(p+t,0)} = -\epsilon'_{\text{He3}} N_0^{(0)} \frac{\tau_{\text{tot}}^{(0)}}{\tau_{\text{He3}}^{(0)}} \left[1 - \exp\left(-\frac{\Delta t}{\tau_{\text{tot}}^{(0)}}\right) \right], \quad (21)$$

$$\Delta N_{\text{UCN}}^{(\beta,0)} = -\epsilon'_{\beta} N_0^{(0)} \frac{\tau_{\text{tot}}^{(0)}}{\tau_{\beta}} \left[1 - \exp\left(-\frac{\Delta t}{\tau_{\text{tot}}^{(0)}}\right) \right]. \quad (22)$$

With the ratio of two equations above, the residual ${}^3\text{He}$ concentration, $X_{\text{He}3}^{(0)}$, is obtained via the following equation,

$$2.56 \times 10^7 X_{\text{He}3}^{(0)} (\text{s}^{-1}) = \tau_{\text{He}3}^{(0)-1} = \frac{1}{\tau_\beta} \frac{\epsilon'_\beta}{\epsilon'_{\text{He}3}} \frac{\Delta N_{\text{UCN}}^{(p+t,0)}}{\Delta N_{\text{UCN}}^{(\beta,0)}}. \quad (23)$$

It is expected the counts of background Compton events might be comparable to or even larger than that of capture events with the residual ${}^3\text{He}$ concentration. It undoubtedly introduces a large uncertainty on $X_{\text{He}3}^{(0)}$, but may not appear as large after propagated to that of the measured $X_{\text{He}3}$ in the later prepared helium mixture. The same applies to the value of τ_β used in equation (23). Despite simply fed with the Particle Data Group value, the error should be negligible after propagated to that of $X_{\text{He}3}$.

Meanwhile, since $\tau_{\text{He}3}^{(0)-1}$ is negligible, the storage time $\tau_{\text{tot}}^{(0)}$ in isotopically pure ${}^4\text{He}$ is expressed as

$$\tau_{\text{tot}}^{(0)-1} = \tau_{\text{up}}^{-1} + \tau_{\text{loss}}^{-1}. \quad (24)$$

$\tau_{\text{tot}}^{(0)}$ can be accurately measured by fitting the time-dependent decline of $\dot{N}_{\text{UCN}}^{(\beta)}(t)$ as in equation (18) with normalization to an arbitrary time zero after the beam stops. The uncertainty of $\tau_{\text{tot}}^{(0)}$ greatly depends on a good understanding of the Compton event rate.

7.5. Accurate determination of helium-3 concentration in the helium mixture

Once the isotopically pure ${}^4\text{He}$ is characterized, the desired ${}^3\text{He}$ concentration of about $X_{\text{He}3} = 2 \times 10^{-10}$ can be prepared by mixing with natural helium of a known ${}^3\text{He}$ abundance as to a preset volume ratio. The ${}^3\text{He}$ atoms may be expelled from the volume via heat flush to fine tune its concentration [35, 59]. It is noteworthy that equations (19) and (23) are equivalent but neither can be used in the accurate determination of $X_{\text{He}3}$, because they proportionally link τ_β to $\tau_{\text{He}3}$. Therefore, $X_{\text{He}3}$ should be extracted from the difference between $\tau_{\text{tot}}^{(0)}$ in the isotopically pure ${}^4\text{He}$ and τ_{tot} in the prepared helium mixture, according to the following expression,

$$2.56 \times 10^7 X_{\text{He}3} [\text{s}^{-1}] = n_{\text{He}3} \sigma_{\text{He}3}^{\text{th}} v_{\text{He}3}^{\text{th}} = \tau_{\text{He}3}^{-1} = \tau_{\text{tot}}^{-1} - \tau_{\text{tot}}^{(0)-1}. \quad (25)$$

In a storage measurement, both the decay and capture event rates, $\dot{N}_{\text{UCN}}^{(p+t)}(t)$ and $\dot{N}_{\text{UCN}}^{(\beta)}(t)$, decline with an identical time dependence on the total storage time constant τ_{tot} of the volume, as in equations (17) and (18). τ_{tot} can be measured via both channels independently, though the decay rate channel might have a higher uncertainty due to the contamination of background Compton events. The detection efficiencies, $\epsilon'_{\text{He}3}$ and ϵ'_β , are irrelevant as long as they remain stable in one measurement cycle. However, it has been experimentally found the UCN loss rate in a material bottle storage is not completely described by an exponential decline as in equations (10), (17) or (18). A dual exponential fit often works much better than the single exponential fit in a storage volume [54]. It is possibly due to the change of UCN spectrum during the storage, especially as UCN's of higher energy have more chances to be up-scattered and absorbed on the wall, or to leak out at some spots of wall materials with lower Fermi potentials. As a result, it may impose systematic uncertainties on the measurement of both τ_{tot} and $\tau_{\text{tot}}^{(0)}$, which is tightly related to the energy spectrum of the UCN's generated in the storage volume and quality of the wall material.

Nevertheless, $X_{\text{He}3}$ given by equation (25) is not a good option to calculate the beam neutron lifetime via equation (3), because the prefactor 2.56×10^7 has inherited an uncertainty of 0.13% from $\sigma_{\text{He}3}^{\text{th}}$. In fact, equation (3) asks for the combined term $n_{\text{He}3} \sigma_{\text{He}3}^{\text{th}} v_{\text{He}3}^{\text{th}}$, which can be

directly acquired once and for all by equation (25). Therefore, $\sigma_{\text{He3}}^{\text{th}}$ by itself is no longer in the path of final calculation of the beam neutron lifetime, and hence, its uncertainty becomes irrelevant. In the end, the sub-0.1% accuracy of the beam neutron lifetime relies on adequate statistics acquired for both τ_{tot} and $\tau_{\text{tot}}^{(0)}$, as well as the ratio of event rates $\kappa = \dot{N}_{p+t}/\dot{N}_{\beta}$ discussed earlier in subsection 5.2.

7.6. Gamma ray backgrounds in the UCN storage measurement

Similarly, gamma rays generated from the neutron captures on the surrounding materials are the major contributor to the background, but the UCN storage measurement is mostly sensitive to the delayed components rather than the prompt. In addition to the capture and scattering by the windows, the 8.9 Å CN beam has a scattering cross section of about 0.025 barns on liquid helium at 0.5 K [60, 61]. 4% of CN's will be inelastically scattered by the phonons in liquid helium. A very small portion is down-converted into UCN's and trapped in the storage volume, whereas majority of the scattered neutrons project at angles of around 84 degrees off the incident direction. About 12% of the scattered neutrons are captured by hydrogen, on top of those captured by the diamond window. The resultant prompt gamma rays are intensive but vanish right after the beam halts. Therefore, it does not affect the measurement of UCN storage time τ_{tot} . More troublesome are the delayed gamma rays from the 0.002% of scattered neutrons captured on the fluorine in the 50 μm thick PTFE reflector. After 200 s of UCN filling with the 8.9 Å CN beam, the neutron-activated fluorine saturates. The resultant delayed gamma rays of 1.6 MeV induce Compton scintillation at 2.5 Hz in liquid helium and 1.7 Hz in the fibers at time-zero when the beam turns off. Since the half-life of the activated fluorine is 11.16 s, the rate of Compton scintillation in liquid helium drops below 1 Hz after 15 s. At the meantime, the rate of decay events is 76.8 Hz, assuming the storage time is 144.2 s and the initial UCN number is $N_0 \approx 7.5 \times 10^4$. The delayed gamma ray background discussed above can be eliminated by replacing the PTFE reflector with a polymer reflector, such as Vikuiti VM2000 [62], but others may still remain due to impurities in the surrounding materials.

8. Conclusion

This proposed experiment has great potential to reach a sensitivity of 0.1% or sub-1 second in neutron lifetime measurement. It offers an entirely different set of systematic uncertainties from the existing beam experiments. As explicitly expressed in equation (3), the calculation of τ_{β} is independent of the neutron flux and the geometry of the decay volume. Most uniquely, it does not require any magnetic field, and may be set up to test the hypothesis of neutron-mirror neutron $n - n'$ oscillations, where the intensity of magnetic field plays an important role [7, 8, 63, 64]. Its apparent disadvantage is the flux of a CN beam at the wavelength $\lambda > 16.5$ Å being much weaker than that of the most commonly in use 4–5 Å CN beams. It will take more beam time to gain adequate statistics. There might be a possibility to optimize the CN beam output for the long wavelength specifically needed in this experiment. Nevertheless, the low event rate allows a thorough characterization of the temporally spaced events with less interference. High precision measurement is required on the three quantities: the ratio of event rates κ in the CN decay volume, as well as the storage time constants τ_{tot} with liquid helium mixed with a proper concentration of helium-3, and $\tau_{\text{tot}}^{(0)}$ with isotopically pure liquid helium-4 in the UCN storage volume. Both the volumes can be connected and share the same batch of prepared liquid helium, where statistics on two measurements can be gained simultaneously. The 8.9 Å CN beam can be extracted from the main beam by a monochromator [33, 65] and sent to the UCN storage volume, whereas the 17 Å CN's can be selected by a set of double or triple

choppers in the main beam and sent to the CN decay volume. Hopefully, a good design of the double or triple choppers can effectively filter the contamination of neutrons in the unwanted spectrum. The atom trap trace analysis can be set up to characterize the helium-3 concentration in samples as a verification of the result offline. Furthermore, while the intrinsic UCN storage time $\tau_{\text{tot}}^{(0)}$ is being measured in the storage volume filled with isotopically pure liquid ^4He , the neutron beta decay spectrum can be simultaneously obtained in the decay volume at a good resolution and accuracy with this detector. It may provide a measurement on the Fierz interference term b in the energy dependent neutron beta decay rate at a potentially high precision compared to the most recent results [66, 67]. This paper only covers the proof of principle for the proposed experiment. Many unknown technical issues will unsurprisingly emerge as the engineering designs and prototyping tests proceed. We hope it can eventually help resolve the neutron lifetime enigma.

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