



Particle filters for data assimilation based on reduced order data models[†]

John Maclean^{a*}, Erik S. Van Vleck^b

^a*School of Mathematical Sciences, University of Adelaide, South Australia.*

<http://www.adelaide.edu.au/directory/john.maclean>

^b*Department of Mathematics, University of Kansas, USA. <http://people.ku.edu/~erikvv/>*

*Correspondence to: John.Maclean@adelaide.edu.au

We introduce a framework for Data Assimilation (DA) in which the data is split into multiple sets corresponding to low-rank projections of the state space. Algorithms are developed that assimilate some or all of the projected data, including an algorithm compatible with any generic DA method. The major application explored here is PROJ-PF, a projected Particle Filter. The PROJ-PF implementation assimilates highly informative but low-dimensional observations. The implementation considered here is based upon using projections corresponding to Assimilation in the Unstable Subspace (AUS). In the context of particle filtering, the projected approach mitigates the collapse of particle ensembles in high dimensional DA problems while preserving as much relevant information as possible, as the unstable and neutral modes correspond to the most uncertain model predictions. In particular we formulate and numerically implement a projected Optimal Proposal Particle Filter (PROJ-OP-PF) and compare to the standard optimal proposal and to the Ensemble Transform Kalman Filter.

Key Words: Data Assimilation, Numerical Analysis, Dimension Reduction

Received...

1. Introduction

Many data assimilation techniques were developed based on extending assumptions of linearity in the state space and data models and under the assumption of Gaussian errors. Several techniques have proven to be successful in weakening these

assumptions, while other techniques have been developed to explicitly overcome these obstacles. Important among these are particle filters (Doucet *et al.* 2000), a key subject of this paper. Particle filters have proven to be successful for low dimensional assimilation problems but tend to have difficulty with higher dimensional problems. Different variants of particle filters have been developed to combat these difficulties, including implicit

[†]JM, ONR, grant: N00014-18-1-2204; ARC, grant DP180100050; EVV, NSF, DMS-1714195 and DMS-1722578.

13 particle filters, proposal density methods, the optimal proposal,
 14 etc. (Chorin *et al.* 2010; Snyder *et al.* 2008; van Leeuwen 2010;
 15 Snyder 2011; Morzfeld *et al.* 2012). Recent work has often
 16 focused on the issue of localization (Farchi and Bocquet (2018),
 17 e.g.), and two localised particle filtering algorithms (Poterjoy
 18 and Anderson 2016; Potthast *et al.* 2019) have been applied in
 19 an operational geophysical framework. In the localised particle
 20 filter of Potthast *et al.* (2019) observations are projected onto the
 21 subspace spanned by the ensemble of model forecasts to reduce
 22 the dimension of the observations.

23 Our contribution in this paper is to develop a framework for
 24 data assimilation schemes in which the data are constrained by
 25 an arbitrary projection to lie in some subspace of observation
 26 or model space. We explicitly obtain a form for the reduction
 27 in data dimension, and an expression that determines how much
 28 the posterior of the Bayesian DA scheme is affected by use
 29 of the projection. While the projection is not specified, the
 30 key idea is that some physically based reduction technique can
 31 then be employed in concert with a DA scheme. In such a
 32 way the assimilation step is performed in a space of very low
 33 dimension.

34 The derivation in this paper was motivated in large part by
 35 assimilation in the unstable subspace (AUS) techniques. These
 36 techniques have largely focused on projection in the tangent space
 37 of the nonlinear model using Lyapunov vectors while employing
 38 the original data or observational model. The techniques and
 39 framework developed in this paper allow for combinations of
 40 (time dependent) projected and unprojected physical and data
 41 models, and their formulation is independent of the source of
 42 the projections. The framework and techniques lead to several
 43 natural applications. In particular, we develop a new particle filter
 44 algorithm that makes use of the original, unprojected physical
 45 and observation models for the particle update together with a
 46 weight update employing the projected observation model and
 47 a resampling scheme that restricts perturbations to the projected
 48 space.

49 We now discuss the historical antecedents of the projections
 50 in this manuscript, and connect them to other recent filtering

51 approaches. The AUS techniques (Carrassi *et al.* 2008a; Trevisan
 52 *et al.* 2010; Palatella *et al.* 2013) to improve speed and reliability
 53 of data assimilation specifically address the partitioning of
 54 the tangent space into stable, neutral and unstable subspaces
 55 corresponding to Lyapunov vectors associated with negative,
 56 zero and positive Lyapunov exponents. In particular, Trevisan,
 57 d'Isidoro & Talagrand propose a modification of 4DVar, so-called
 58 4DVar-AUS, in which corrections are applied only in the unstable
 59 and neutral subspaces (Trevisan *et al.* 2010; Palatella *et al.* 2013).
 60 These techniques are based on updating in the unstable portion of
 61 the tangent space and may be interpreted in terms of projecting
 62 covariance matrices during the assimilation step. Motivated by
 63 these techniques for assimilation in the unstable subspace, in
 64 de Leeuw *et al.* (2018) a new method is developed for data
 65 assimilation that utilizes distinct treatments of the dynamics in
 66 the stable and non-stable directions. The key piece of de Leeuw
 67 *et al.* (2018) related to this work is the following projected model
 68 update. For a smooth discrete time model $u_{n+1} = F_n(u_n)$ and
 69 projection Π_n , and for $\{u_n^{(0)}\}_{n=0}^N$ any reference solution, solve
 70 for $\{d_n\}_{n=0}^N$:

$$u_{n+1}^{(0)} + d_{n+1} = \Pi_{n+1} F_n(u_n^{(0)} + d_n), \quad n = 0, \dots, N-1. \quad (1)$$

71 Unlike most past work related to AUS our primary focus is on
 72 developing a systematic approach to confining the data, not the
 73 model, to the unstable subspace. In some of the initial works on
 74 AUS (Carrassi *et al.* 2007, 2008b), either target observations at
 75 the location where the unstable mode attains its maximum value,
 76 or only the observations falling in the vicinity of the maximum,
 77 were assimilated. Albeit empirical, that choice already signified
 78 using only data projected on an approximation of the unstable
 79 subspace, that was obtained by Breeding on the Data Assimilation
 80 Cycle (BDAS). Furthermore, González-Tokman and Hunt (2013);
 81 Bocquet *et al.* (2017); Grudzien *et al.* (2018b,a); Frank and
 82 Zhuk (2018); Reddy *et al.* (2020); Tranninger *et al.* (2020) are
 83 all at least in part devoted to discussing the necessary and/or
 84 sufficient criteria for filter stability in terms of the projection of
 85 the observations into the unstable/neutral/weakly stable directions

86 and this is directly related to the choice of adaptive observation
 87 operators in Law *et al.* (2016).

88 Another branch of projected DA schemes use the ‘Dynamically
 89 Orthogonal’ (DO) formulation (Sapsis and Lermusiaux 2009;
 90 Sapsis 2010), in which the forecast model is broken into a partial
 91 differential equation governing the mean field and a number
 92 of stochastic differential equations describing the evolution of
 93 components in a time-dependent stochastic subspace of the
 94 original differential equation. The DO approach was used to
 95 assimilate with different DA schemes in the subspace and mean
 96 field space in Sondergaard and Lermusiaux (2013); Majda *et al.*
 97 (2014); Qi and Majda (2015). These techniques use both a
 98 projected and mean field model to make a forecast, similar to
 99 using (1).

100 Projection-based DA schemes have been developed to assimilate
 101 coherent structures (Maclean *et al.* 2017) or features (Morzfeld
 102 *et al.* 2018) in the data. These approaches have used likelihood-
 103 free sequential Monte Carlo methods, or an ad hoc ‘perturbed
 104 observations’ approach, to deal with the difficulty of calculating
 105 the likelihood function for a coherent structure. The derivation
 106 in this paper may lead to an explicit likelihood for data-derived
 107 coherent structures/features obtained via a projection.

108 This paper is organized as follows. Data assimilation is reviewed
 109 in section 2 and projected DA is formulated in section 3.
 110 Algorithms for using the new projected data are introduced
 111 (section 4) and applied in several numerical experiments
 112 (section 5). A discussion (section 7) and bibliography conclude
 113 the paper.

114 2. Data Assimilation

Data assimilation methods combine orbits from a dynamical system model with measurement data to obtain an improved estimate for the posterior probability density function (pdf) of a physical system. In this paper we develop a data assimilation method in the context of the discrete time stochastic model

$$u_{n+1} = F_n(u_n) + \sigma_n, \quad n = 0, 1, \dots \quad (2)$$

where $u_n \in \mathbb{R}^N$ are the state variables at time n and $\sigma_n \sim \mathcal{N}(0, \mathbf{Q})$, i.e., drawn from a normal distribution with mean zero and model error covariance \mathbf{Q} . Let the sequence $\{u_0^t, u_1^t, \dots\}$, be a distinguished orbit of this system, referred to as the *true solution* of the model, and presumed to be unknown. As each time t_n is reached we collect an observation y_n related to u_n^t via

$$y_n = \mathbf{H}u_n^t + \eta_n, \quad y_n \in \mathbb{R}^M \quad (3)$$

where $\mathbf{H} : \mathbb{R}^N \rightarrow \mathbb{R}^M$, $M \leq N$, is the observation operator, and the noise variables η_n are drawn from a normal distribution $\eta_n \sim \mathcal{N}(0, \mathbf{R})$ with zero mean and known observational error covariance matrix \mathbf{R} . In general the observation operator can be nonlinear.

We formulate DA under the ubiquitous Bayesian approach. Consider the assimilation of a single observation, y_n , at time step n . Given a prior estimate $p(u_n)$ of the state, Bayes’ Law gives $p(u_n|y_n) \propto p(y_n|u_n)p(u_n)$. Using (3) the likelihood function is, up to a normalization constant,

$$p(y_n|u_n) \propto \exp \left[-\frac{1}{2} (y_n - \mathbf{H}u_n)^T \mathbf{R}^{-1} (y_n - \mathbf{H}u_n) \right]. \quad (4)$$

This procedure, which we have written for the assimilation of data at a single observation time, readily extends to the sequential assimilation of observations at multiple times under the assumptions that the state is Markovian and the observation errors at different times are conditionally independent (see for example Budhiraja *et al.* (2017)).

In the following we introduce some key DA schemes. Not much detail is given here, but the interested reader is referred in particular to three recent books on DA, (Reich and Cotter 2015; Law *et al.* 2015; Asch *et al.* 2016).

2.1. Kalman Filtering

The Kalman Filter and later extensions are ubiquitous in DA, and are now briefly described. For a linear model, i.e. where (2)

is

$$u_{n+1} = \mathbf{A}_n u_n + \sigma_n, \quad (5)$$

and for the linear observation operator \mathbf{H} , the Kalman Filter calculates the exact posterior $u_n|y_n \sim \mathcal{N}(u_n^a, \mathbf{P}_n^a)$, where the analysis variables are

$$u_n^a = u_n^f + \mathbf{K}_n(y_n - \mathbf{H}u_n^f), \quad (6)$$

$$\mathbf{P}_n^a = (\mathbf{I} - \mathbf{K}_n \mathbf{H}) \mathbf{P}_n^f. \quad (7)$$

The weight matrix \mathbf{K}_n is the Kalman gain matrix

$$\mathbf{K}_n = \mathbf{P}_n^f \mathbf{H}^T \left(\mathbf{H} \mathbf{P}_n^f \mathbf{H}^T + \mathbf{R} \right)^{-1}. \quad (8)$$

133 The superscript f is reserved for *forecast* variables, obtained at
 134 time n by using (5) to update $\{u_{n-1}^a, \mathbf{P}_{n-1}^a\}$,

$$u_n^f = \mathbf{A}_{n-1} u_{n-1}^a + \sigma_{n-1}, \quad (9)$$

$$\mathbf{P}_n^f = \mathbf{A}_{n-1} \mathbf{P}_{n-1}^a \mathbf{A}_{n-1}^T + \mathbf{Q}. \quad (10)$$

136 Two extensions of the Kalman Filter are prevalent in nonlinear
 137 DA, the Extended Kalman Filter (EKF) and Ensemble Kalman
 138 Filter (EnKF). Neither give the exact posterior for a nonlinear
 139 model.

140 2.1.1. Extended Kalman Filter

141 The nonlinear model (2) is used to make the forecast u_n^f , and
 142 then the update of the covariance (10) is applied using the
 143 linearisation

$$\mathbf{A}_n = \left. \frac{\partial F_n}{\partial u} \right|_{u_n^a}. \quad (11)$$

144 If the observation operator is a nonlinear function $h()$, the
 145 linearization

$$\mathbf{H}_n = \left. \frac{dh}{du} \right|_{u_n^f} \quad (12)$$

146 is used everywhere except to compute the *innovation* $y_n - h(u_n^f)$
 147 in the calculation of y_n^a .

148 The EKF is suitable for low dimensional nonlinear filtering, but
 149 the required linearizations are nontrivial for high-dimensional

filtering. The EnKF by contrast is well suited to high
 150 dimensions. 151

152 2.1.2. Ensemble Kalman Filter

153 The Ensemble Kalman Filter is a Monte Carlo approximation of
 154 the Kalman Filter that is well suited to high dimensional filtering
 155 problems, introduced in Evensen (1994); Burgers *et al.* (1998).
 156 An ensemble of forecasts $u_n^{f,i}$ are made at time t_n , i from 1 to L .
 157 Then the forecast covariance \mathbf{P}_n^f is approximated by the sample
 158 covariance of the ensemble, and the analysis ensemble $u_n^{a,i}$ is
 159 obtained in such a way that its mean $\bar{u}_n^a = \frac{1}{L} \sum_i u_n^{a,i}$ satisfies
 160 (6) and its sample covariance satisfies (7). In this paper we will
 161 use analysis updates corresponding to the Ensemble Transform
 162 Kalman Filter (ETKF) (Bishop *et al.* 2001). For more details and
 163 a modern introduction to the Ensemble Kalman Filter, see e.g.
 164 Evensen (2009) and Carrassi *et al.* (2018). 164

165 2.2. The Particle Filter

166 Particle Filters (PF) are a collection of particle based data
 167 assimilation schemes that do not rely on linearization of the
 168 dynamics or Gaussian representations of the posterior; see Doucet
 169 *et al.* (2001) for a comprehensive review. The basic idea is to
 170 represent the prior distribution $p(u_n)$, previously the forecast,
 171 and the posterior distribution $p(u_n|y_n)$, previously the analysis,
 172 by discrete probability measures. Suppose that at time $n-1$ we
 173 have the posterior distribution (u_{n-1}^i, w_{n-1}^i) , supported on points
 174 $u_{n-1}^1, \dots, u_{n-1}^L$ and with weights $w_{n-1}^1, \dots, w_{n-1}^L$. Each $w_{n-1}^i \geq$
 175 0 and $\sum_{i=1}^L w_{n-1}^i = 1$. Here L is the number of particles that
 176 are used to approximate the distribution. The two key steps in the
 177 Particle Filter are as follows:

178 *Prediction step.* Propagate each of the particles $u_{n-1}^i \mapsto u_n^i$. One
 179 simple choice, the bootstrap PF, is to use the state dynamics (2) to
 180 forecast each particle.

181 This gives the forecast probability distribution as a discrete
 182 probability measure concentrated on L points $\{u_n^i\}_{i=1}^L$ with
 183 weights $\{w_n^i\}_{i=1}^L$.

184 *Filtering step.* Update the weights $\{w_{n-1}^i\}_{i=1}^L$ using the
 185 observation y_n . In the bootstrap PF the update is

$$w_n^i = c w_{n-1}^i p(y_n|u_n^i), \quad (13)$$

186 where c is chosen so that $\sum_{i=1}^L w_n^i = 1$.

187 This scheme is easy to implement but suffers from severe
 188 degeneracy, especially for high dimensional observations: even
 189 ten independent observations are sufficient to produce degenerate
 190 weights. That is, after a few time steps all the weight tends to
 191 concentrate on a few particles. A partial remedy is to monitor the
 192 Effective Sample Size (ESS) and resample when the ESS drops
 193 below some threshold in order to refresh the particle cloud; see
 194 e.g. Doucet *et al.* (2001); Budhiraja *et al.* (2017).

195 2.2.1. The Optimal Proposal

196 The optimal proposal particle filter (OP-PF) (Snyder *et al.* 2008;
 197 Doucet *et al.* 2000; Snyder 2011; Van Leeuwen 2012, e.g.)
 198 attempts to address the degeneracy issue in particle filters with the
 199 aim of ensuring that all posterior particles have similar weights.
 200 The ‘proposal’ is the distribution used to update the particles
 201 from one time step to the next. In the prediction step in the
 202 basic particle filter above, the particles are updated using the
 203 model, so the proposal density in that approach is (compare (2))
 204 $u_n^i | u_{n-1}^i \sim \mathcal{N}(F_{n-1}(u_{n-1}^i), \mathbf{Q})$.

205 The optimal proposal density is $p(u_n^i | u_{n-1}^i, y_n)$. Given the
 206 additive noise of the model (2) and a linear observation operator
 207 \mathbf{H} , the optimal proposal update in each particle is Gaussian
 208 with $u_n^i | u_{n-1}^i, y_n \sim \mathcal{N}(m_n^i, \mathbf{Q}_p)$, and we obtain the explicit
 209 update

$$210 u_n^i = m_n^i + \phi, \quad \phi \sim \mathcal{N}(0, \mathbf{Q}_p) \quad (14)$$

where

$$211 m_n^i = F_{n-1}(u_{n-1}^i) + \mathbf{Q} \mathbf{H}^T \left(\mathbf{H} \mathbf{Q} \mathbf{H}^T + \mathbf{R} \right)^{-1} I_n^i, \quad (15)$$

$$212 \mathbf{Q}_p^{-1} = \mathbf{Q}^{-1} + \mathbf{H}^T \mathbf{R}^{-1} \mathbf{H}, \quad (16)$$

213 and $I_n^i := y_n - \mathbf{H} F_{n-1}(u_{n-1}^i)$. The mean of the particle m_n^i is
 214 obtained by a Kalman filter step, albeit with $\mathbf{P}_n^f = \mathbf{Q}$.

Two applications of Bayes’ law (e.g. in Snyder (2011)) show that
 the weight update for the i -th particle drawn from this proposal
 satisfies $w_n^i \propto p(y_n | u_{n-1}^i) w_{n-1}^i$ with a Gaussian likelihood

function,

$$215 w_n^i \propto \exp \left[-\frac{1}{2} (I_n^i)^T \left(\mathbf{H} \mathbf{Q} \mathbf{H}^T + \mathbf{R} \right)^{-1} (I_n^i) \right] w_{n-1}^i. \quad (17)$$

216 As mentioned in the previous section, degeneracy - characterised
 217 by a single particle with weight of approximately 1 - is a common
 218 problem in the PF. In Snyder *et al.* (2015) it is shown that, of all PF
 219 schemes that obtain u_n^i using u_{n-1}^i and y_n , the ‘optimal proposal’
 220 above has the minimum variance in the weights. That is, it suffers
 221 the least from weight degeneracy. In van Leeuwen *et al.* (2019)
 222 this result is extended to any PF scheme that obtains u_n^i using
 223 $u_{n-1}^{1:L}$ and y_n .

224 However, in Snyder (2011) it is shown that the optimal proposal
 225 requires an ensemble size L satisfying $\log L \propto N \times M$ for a
 226 linear model, or will suffer from filter degeneracy. That is, filter
 227 degeneracy is intimately connected to model and observation
 228 dimension, and is a fundamental obstacle to Particle Filtering in
 229 high dimensional problems.

226 3. Projected Data Models

227 We now develop an approach to decompose the observations
 228 using projections defined in state space. A wealth of techniques
 229 from dynamical systems theory can then be used to obtain low-
 230 dimensional data models.

231 Suppose that at time n a dynamically significant rank p orthogonal
 232 projection $\Pi_n \in \mathbb{R}^{N \times N}$ is available, as well as data $y_n \in \mathbb{R}^M$.
 233 We next derive a projected data model consisting of a projected
 234 obeservation operator, projected observation, and associated
 235 observation error covariance.

236 *Step One: lift the data into model space*

237 In order to apply the projection Π_n to data, we first need to find
 238 an equivalent representation of the data in model space.
 239 Assuming \mathbf{H} has full row rank, we define an N -dimensional
 240 vector $\tilde{y}_n = \mathbf{H}^\dagger y_n$ where $\mathbf{H}^\dagger = \mathbf{H}^T (\mathbf{H} \mathbf{H}^T)^{-1}$, the pseudo
 241 inverse of \mathbf{H} . The data model for \tilde{y}_n is

$$242 \tilde{y}_n = \mathbf{H}^\dagger y_n = \Pi_{\mathbf{H}} u_n^t + \mathbf{H}^\dagger \eta_n = \Pi_{\mathbf{H}} u_n^t + \psi_n \quad (18)$$

237 where $\Pi_{\mathbf{H}} = \mathbf{H}^\dagger \mathbf{H}$ is an orthogonal projection, and
 238 $\psi_n \sim \mathcal{N}(0, \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T)$.
 239

240 Using that $\mathbf{H}\mathbf{H}^\dagger = \mathbf{I}$ one readily confirms that $\mathbf{H}\tilde{y}_n = y_n =$
 241 $\mathbf{H}u_n^t + \eta_n$. That is, the observation operator collapses \tilde{y}_n onto
 242 the standard data model. The transformation through \mathbf{H}^\dagger has
 243 not affected the output of a DA scheme, as $p(\tilde{y}_n|x) = p(y_n|x)$;
 244 however \tilde{y}_n is of compatible dimension with Π_n .

245 *Step Two: project the data into a rank p subspace*

246 We now make use of the orthogonal projection Π_n . The idea
 247 is to formulate a new data model, along the lines of $\Pi_n\tilde{y}_n =$
 248 $\Pi_n\Pi_{\mathbf{H}}u_n^t + \Pi_n\mathbf{H}^\dagger\eta_n$, that contains only the components of
 249 the observation that align with the projection. The projected
 250 data models that are developed here may be considered as
 251 generalizations of the construction of observation operators (see
 252 Grudzien *et al.* (2018a) Def. 13 and Law *et al.* (2016)).

Define $y_n^p = \Pi_n\tilde{y}_n = \Pi_n\mathbf{H}^\dagger y_n \in \mathbb{R}^N$, the projected observation.
 The data model is

$$y_n^p = \Pi_n\mathbf{H}^\dagger y_n = \Pi_n\Pi_{\mathbf{H}}u_n^t + \xi_n \quad (19)$$

where $\xi_n \sim \mathcal{N}(0, \Pi_n\mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \Pi_n)$. The data model y_n^p has a singular normal distribution with support in the p -dimensional subspace of model space spanned by the projection Π_n . Some information from the observations is typically lost in this step by applying the projection. At the end of Section 3 we derive a data model for the orthogonal data (that is discarded in (19)), and in Section 4.1 we sketch a result that establishes the difference between assimilating with the full data (3) and with the projected data. The likelihood $y_n^p|u$ of this distribution has an explicit form using the pseudo-inverse (see e.g. Tsukuma and Kubokawa (2015)) as

$$p(y_n^p|u) \propto \exp\left(-\frac{1}{2}(I_n^p)^T \left(\Pi_n\mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \Pi_n\right)^\dagger I_n^p\right) \quad (20)$$

253 where $I_n^p := y_n^p - \Pi_n\Pi_{\mathbf{H}}u$.

Remark 1. The product $\Pi_n\Pi_{\mathbf{H}}$ is not generally an orthogonal projection, and in some circumstances it might be desired to instead identify the projection $\Pi_n^{\mathbf{H}}$ that is the intersection of Π_n and $\Pi_{\mathbf{H}}$. This projection $\Pi_n^{\mathbf{H}}$ may be approximated by Von Neumann's algorithm or Dykstra's projection algorithm. The projection $\Pi_n^{\mathbf{H}}$ should only be used if the transversality condition $p + M - N > 0$ is satisfied; otherwise there is no guarantee of any intersection between Π_n and $\Pi_{\mathbf{H}}$. If the transversality condition is satisfied, e.g., with a high dimensional observation space, then we can replace the product of projections $\Pi_n\Pi_{\mathbf{H}}$ with the projection into the intersection $\Pi_n^{\mathbf{H}}$. Since in most applications this will not be satisfied, throughout the rest of the paper we focus on the use of the product of projections.

268 *Step Three: reduce the projected data to a p -vector*

To make explicit the reduction in the data dimension that has been obtained by y_n^p we introduce a low dimensional data model. Denote by \mathbf{U}_n the matrix with orthonormal columns satisfying $\Pi_n = \mathbf{U}_n \mathbf{U}_n^T$. This matrix may be already known (in the examples in Section 3 \mathbf{U}_n is obtained first, and then Π_n is calculated from $\mathbf{U}_n \mathbf{U}_n^T$), or \mathbf{U}_n may be found via the singular value or Schur decompositions.

Define $y_n^q = \mathbf{U}_n^T y_n^p \equiv \mathbf{U}_n^T \tilde{y}_n \in \mathbb{R}^p$, with the associated data model

$$y_n^q = \mathbf{H}_n^q u_n^t + \gamma_n, \quad (21)$$

where $\mathbf{H}_n^q = \mathbf{U}_n^T \Pi_{\mathbf{H}}$, $\gamma_n \sim \mathcal{N}(0, \mathbf{R}_n^q)$, and $\mathbf{R}_n^q = \mathbf{U}_n^T \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \mathbf{U}_n$.

The transformations between, and dimensions of, the different data variables defined in this section are illustrated in Figure 1.

3.1. Properties of the projected data

Theorem 3.1 (Equivalence of y_n^p and y_n^q). For the data models associated with y_n^p and y_n^q given by (19) and (21), respectively, $p(y_n^q|u) = p(y_n^p|u)$.

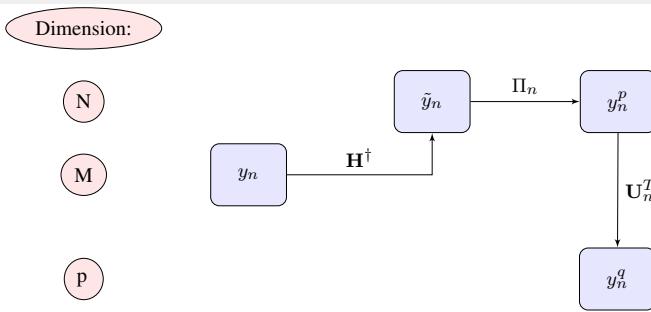


Figure 1. The progression from the original data y_n to low-dimensional, projected data y_n^q . The rectangular boxes contain data, or data-derived constructs. The vertical placement of each box corresponds to the dimension of the data at each step: $N \geq M \geq p$. Note that in practice one does not need to compute \tilde{y}_n or y_n^p .

Proof. The matrix \mathbf{U}_n has orthonormal columns, so $\mathbf{U}_n^\dagger = \mathbf{U}_n^T$ and for any matrix \mathbf{B}

$$(\mathbf{U}_n \mathbf{B})^\dagger = \mathbf{B}^\dagger \mathbf{U}_n^\dagger = \mathbf{B}^\dagger \mathbf{U}_n^T, \quad (22)$$

$$(\mathbf{B} \mathbf{U}_n^T)^\dagger = (\mathbf{U}_n^T)^\dagger \mathbf{B}^\dagger = \mathbf{U}_n \mathbf{B}^\dagger. \quad (23)$$

279 Applying these results to (20), and using that $\Pi_n = \mathbf{U}_n \mathbf{U}_n^T$,
280 $y_n^p = \mathbf{U}_n y_n^q$, $\mathbf{U}_n^T \mathbf{U}_n = \mathbf{I}$, $I_n^p := y_n^p - \Pi_n \Pi_{\mathbf{H}} u$, and $I_n^q := y_n^q -$
281 $\mathbf{U}_n^T \Pi_{\mathbf{H}} u$,

$$\begin{aligned} p(y_n^p | u) &= \frac{1}{c} \exp \left(-\frac{1}{2} (I_n^p)^T (\Pi_n \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \Pi_n)^{-1} I_n^p \right) \\ &= \frac{1}{c} \exp \left(-\frac{1}{2} (\mathbf{U}_n I_n^q)^T (\tilde{\mathbf{R}}_n)^\dagger (\mathbf{U}_n I_n^q) \right) \\ &= \frac{1}{c} \exp \left(-\frac{1}{2} (I_n^q)^T \mathbf{U}_n^T \mathbf{U}_n (\mathbf{R}_n^q)^\dagger \mathbf{U}_n^T \mathbf{U}_n I_n^q \right) \\ &= \frac{1}{c} \exp \left(-\frac{1}{2} (I_n^q)^T (\mathbf{R}_n^q)^\dagger I_n^q \right) \\ &= p(y_n^q | u) \end{aligned} \quad (24)$$

282 where $\tilde{\mathbf{R}}_n := \mathbf{U}_n \mathbf{U}_n^T \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \mathbf{U}_n \mathbf{U}_n^T$ and c is a normalising
283 constant. \square

284 If in addition $p \leq M$ (or $0 < p + M - N \leq M$ for $\Pi_n \equiv \Pi_n^{\mathbf{H}}$),
285 and if $\mathbf{H} \mathbf{U}_n$ is full rank, then the covariance matrix \mathbf{R}_n^q of
286 y_n^q is invertible and y^q has a standard normal distribution.
287 More generally for $(\mathbf{H} \mathbf{H}^T)^{-1} \mathbf{R}(\mathbf{H} \mathbf{H}^T)^{-1} = \mathbf{L}^T \mathbf{L}$, the Cholesky
288 factorization, consider the SVD of $\mathbf{L} \mathbf{H} \mathbf{U}_n = \mathbf{S} \mathbf{\Sigma} \mathbf{V}^T$. The rank of
289 the covariance matrix $\mathbf{R}_n^q = \mathbf{U}_n^T \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \mathbf{U}_n = \mathbf{V} \mathbf{\Sigma}^T \mathbf{V}^T$ is
290 equal to the number of non-zero singular values of $\mathbf{\Sigma}$.

291 Theorem 3.1 provides a blueprint for any DA system with a linear
292 observation operator to be efficiently implemented with projected
293 observations, involving the following changes: the observation y_n
294 is replaced with y_n^q , the observation operator \mathbf{H} is replaced with

\mathbf{H}_n^q , and the assumed measurement covariance \mathbf{R} is replaced with
295 \mathbf{R}_n^q .

3.2. The orthogonal data model

Though the focus of this paper is on the projected data, a data model for the complementary orthogonal projection $\mathbf{I} - \Pi_n$ is easy to write down. Define

$$y_n^{q\perp} = (\mathbf{U}_n^\perp)^T \tilde{y}_n \in \mathbb{R}^{N-p}, \quad (25)$$

where $\mathbf{U}_n^\perp (\mathbf{U}_n^\perp)^T = \mathbf{I} - \Pi_n$. The two projected data models are not independent in general and have joint distribution

$$\begin{bmatrix} y_n^q \\ y_n^{q\perp} \end{bmatrix} \sim \mathcal{N} \left(\begin{bmatrix} \mathbf{H}_n^q u_n^t \\ \mathbf{H}_n^{q\perp} u_n^t \end{bmatrix}, \begin{bmatrix} \mathbf{R}_n^q & \mathbf{R}_{12,n}^q \\ \mathbf{R}_{21,n}^q & \mathbf{R}_n^{q\perp} \end{bmatrix} \right), \quad (26)$$

where $\mathbf{H}^{q\perp} = (\mathbf{U}_n^\perp)^T \Pi_{\mathbf{H}}$, $\mathbf{R}^{q\perp} = (\mathbf{U}_n^\perp)^T \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \mathbf{U}_n^\perp$, and the off-diagonal covariances are $\mathbf{R}_{12,n}^q = \mathbf{U}_n^T \mathbf{H}^\dagger \mathbf{R}(\mathbf{H}^\dagger)^T \mathbf{U}_n^\perp$ 299 and $\mathbf{R}_{21,n}^q = (\mathbf{R}_{12,n}^q)^T$.
300

301 The joint distribution (26) is not used in this manuscript. The
302 cross-covariance term $\mathbf{R}_{12,n}^q$ measures the information about the
303 projected subspace that is lost by not assimilating the orthogonal
304 component $y^{q\perp}$. In ongoing work we are developing approaches
305 to factorise the posterior into two components, via (26), and apply
306 different DA methods to each component, incorporating the cross-
307 covariance terms.

4. Algorithms for Projected DA

308 In this section we discuss how some combination of the
309 standard/projected forecast models (2), (1) and data models (3),
310 (21), (25)–(26) may be used to form a ‘projected DA scheme’.
311

312 A projected data model changes the innovation, the observation
313 operator, and the observation error covariance. A projected
314 physical model changes the prior and model error covariances.
315 We want combinations of physical models, data models, and
316 DA techniques that optimize the assimilation, particularly of the
317 Particle Filtering schemes discussed in Section 2.2.

318 We identify the following approaches to assimilating with
319 projected data using the results of this paper:

320 **Algorithm 1** (Project data only, and discard the orthogonal
 321 component). *Apply a standard DA scheme using the unprojected
 322 forecast model (2), but replace the standard data (3) with the
 323 projected data y_n^q of (21). The observation operator is replaced
 324 by \mathbf{H}_n^q , and the covariance matrix of the observations is replaced
 325 by \mathbf{R}_n^q .*

326 By careful construction, the unique component of ‘Algorithm 1’
 327 is actually a modified Bayesian posterior. This construction of a
 328 modified DA scheme as a modification to the posterior enables
 329 modularity: any DA scheme suitable for nonlinear filtering may
 330 be implemented with projected data models. In particular, this
 331 projected DA algorithm is fully compatible with localized DA
 332 schemes.

PROJ-PF uses the standard forecast model (2) to update the
 particles, but computes the weight update with

$$w_n^i \propto \exp \left[-\frac{1}{2} (y_n^q - \mathbf{H}_n^q u_n)^T (\mathbf{R}_n^q)^{-1} (y_n^q - \mathbf{H}_n^q u_n) \right] w_{n-1}^i. \quad (27)$$

333

334 Another algorithm to be described is a novel, efficient PF
 335 scheme taking advantage of the Optimal Proposal PF described
 336 in section 2.2.1.

Algorithm 2 (PROJ-OP-PF: Blend projected and unprojected data
 in the assimilation step). *This algorithm describes a Particle
 Filter, PROJ-OP-PF, that uses the typical optimal proposal
 equations (14)–(15) for the particle update. The weight update
 for each particle is computed using the projected data model only,
 i.e. using the projected form of (17),*

$$w_n^i \propto \exp \left[-\frac{1}{2} (I_n^q)^T \left(\mathbf{H}_n^q \mathbf{Q} (\mathbf{H}_n^q)^T + \mathbf{R}_n^q \right)^{-1} (I_n^q) \right] w_{n-1}^i. \quad (28)$$

337 where $I_n^q \equiv I_n^q(u_{n-1}^i) := y_n^q - \mathbf{H}_n^q F_{n-1}(u_{n-1}^i)$.

338 Algorithm 2 uses the original observation error covariance and the
 339 original observation operator for the particle update but employs
 340 the projected observation error covariance and the projected
 341 observation operator for the weight update. This strategy will be

342 tested on the chaotic Lorenz-96 system in Section 6. One major
 343 advantage of this approach is that it requires no modification of
 344 the numerical simulation used to obtain the forecast. A second
 345 advantage is its efficiency; the full data are used for the particle
 346 update step, over which the update is straightforward and the
 347 dimension of the data does not lead to filter degeneracy; and
 348 only the projected data are used to avoid filter degeneracy in the
 349 weight update step. The scheme will prove to be more accurate
 350 than either, OP-PF or an Algorithm 1 implementation of OP-PF,
 351 in numerical tests.

352 Particle Filters can benefit from adding noise on resampling,
 353 particularly with deterministic forecast models. The correct way
 354 to do this is to generate noise sampled from a Markov chain that
 355 leaves the target pdf unchanged, see e.g. Doucet *et al.* (2000). For
 356 example, one can implement an accept-reject step for the particle.
 357 For simplicity we either do not add noise, or add noise sampled
 358 from $\mathcal{N}(\mathbf{0}, \omega^2 \mathbf{I})$, where $\omega \in \mathbb{R}$ must be tuned. We consider an
 359 algorithm for projecting the noise on resampling.

Algorithm 3 (PROJ-RESAMP: Resampling in the Unstable
 360 Subspace). *When adding noise ω to particles after resampling,
 361 multiply this random vector by $\alpha \Pi_n + (1 - \alpha) \mathbf{I}$ for some $\alpha \in
 362 [0, 1]$.*

363 When $\alpha = 0$ this algorithm is no different to the normal
 364 resampling approach, but for $\alpha > 0$ some proportion of the
 365 uncertainty in resampling is constrained to lie in the space spanned
 366 by the columns of \mathbf{U}_n . If employed in concert with an accept-
 367 reject step, this may improve the chances of acceptance. For AUS
 368 the resampling scheme should add more noise in the directions
 369 of greatest uncertainty in the forecast model, which provides one
 370 advantage; a second advantage is that the algorithm does not shift
 371 particles as far off the attractor.

4.1. Convergence results for projected algorithms

372 A normal line of inquiry for a new DA algorithm is to quantify
 373 the conditions under which it will well represent the posterior
 374 distribution, which neglecting time subscripts we write as $p(u|y)$.
 375 The projected algorithms above do not generally converge to
 376 $p(u|y)$, and so there are two questions: ‘Does the algorithm
 377

379 converge to a known distribution?", and 'How different is that
 380 distribution to the actual posterior?'.
 381

382 Algorithm 1 clearly implements an approximation of the
 383 distribution $p(u|y^q)$. That is, a Particle Filter implementation
 384 would converge to $p(u|y^q)$ in the limit as the number of
 385 particles approaches infinity. The distribution approximated by
 386 Algorithm 2 is a blending of $p(u|y)$ and $p(u|y^q)$ that is non-trivial
 to obtain in closed form.
 387

We now quantify how the Algorithm 1 distribution $p(u|y^q)$ relates to the standard posterior $p(u|y)$. For this we will employ the Hellinger distance: given two probability measures μ and μ' , with associated probability distributions ρ and ρ' , the Hellinger distance between the two is

$$d_H(\mu, \mu') = \left[\frac{1}{2} \int \left(\sqrt{\rho(u)} - \sqrt{\rho'(u)} \right)^2 du \right]^{1/2}. \quad (29)$$

387 To bound this distance for Algorithm 1 we write $\rho(u) = p(u|y)$
 388 and $\rho'(u) = p(u|y^q)$. The second distribution is written as

$$p(u|y^q) = p(u|y) \frac{p(y|y^q)}{p(y|u, y^q)}, \quad (30)$$

389 obtained via Bayes' law in the form $p(u) = p(u|y) p(y)/p(y|u)$,
 390 conditioning on y^q , and using $p(u|y, y^q) = p(u|y)$. Factorising
 391 both numerator and denominator, we obtain the final form

$$p(u|y^q) = p(u|y) \frac{p(y^{q\perp}|y^q)}{p(y^{q\perp}|u, y^q)}. \quad (31)$$

Substituting into (29) we obtain a bound for the consistency of
 Algorithm 1 with the original posterior $p(u|y)$,

$$\begin{aligned} d_H(\mu, \mu') &= \left[\frac{1}{2} \int \left(1 - \sqrt{\frac{p(y^{q\perp}|y^q)}{p(y^{q\perp}|u, y^q)}} \right)^2 \rho(u) du \right]^{1/2} \\ &= \left[\frac{1}{2} \mathbb{E}^\mu \left(1 - \sqrt{\frac{p(y^{q\perp}|y^q)}{p(y^{q\perp}|u, y^q)}} \right)^2 \right]^{1/2}. \end{aligned} \quad (32)$$

392 An intuitive example of the above bounds in practice is a slow-
 393 fast system with a slow manifold onto which the fast variables
 394 are attracted. Choosing Π_n to identify the slow variables will lead

395 to a small value of $d_H(\mu, \mu')$ for Algorithm 1 data assimilation
 396 schemes, since knowledge of the slow variables is sufficient to
 397 constrain the fast variables. In the case where there are few slow
 398 variables and many fast variables, then, an Algorithm 1 Particle
 399 Filter will be a much less degenerate implementation of the
 400 Particle Filter that converges close to the desired posterior $p(u|y)$.
 401 The bounds (29)–(32) at present help to build intuition rather than
 402 providing a practical tool: they only apply at a single assimilation
 403 time, and it is difficult to infer a bound for the difference between
 404 the Algorithm 1 posterior and the standard posterior over multiple
 405 time steps.
 406

5. Application: Assimilation in the Unstable 407 Subspace

For the remainder of the paper we will study the case where the projection identifies the most unstable modes in the forecast model. To determine these modes we employ the discrete QR algorithm (Dieci and Van Vleck 2007, 2015). For the discrete time model $u_{n+1} = F_n(u_n) + \sigma_n$ with $u_n \in \mathbb{R}^N$, let $\mathbf{U}_0 \in \mathbb{R}^{N \times p}$ ($p \leq N$) denote a random matrix such that $\mathbf{U}_0^T \mathbf{U}_0 = \mathbf{I}$,

$$\mathbf{U}_{n+1} \mathbf{T}_n = F'_n(u_n) \mathbf{U}_n \approx \frac{1}{\epsilon} [F_n(u_n + \epsilon \mathbf{U}_n) - F_n(u_n)], \quad n = 0, 1, \dots \quad (33)$$

where $\mathbf{U}_{n+1}^T \mathbf{U}_{n+1} = \mathbf{I}$ and \mathbf{T}_n is upper triangular with positive
 408 diagonal elements. With a finite difference approximation the cost
 409 is that of an ensemble of size p plus a reduced QR via modified
 410 Gram-Schmidt to re-orthogonalize. Time dependent orthogonal
 411 projections to decompose state space are $\Pi_n = \mathbf{U}_n \mathbf{U}_n^T$ and $\mathbf{I} -$
 412 $\Pi_n = \mathbf{I} - \mathbf{U}_n \mathbf{U}_n^T$. In order to apply (33) to an ensemble DA
 413 method, \mathbf{U}_0 must be specified and we must choose how to obtain
 414 u_n from the ensemble of particles at time t_n . We initialise \mathbf{U}_0
 415 from a modified Gram-Schmidt orthonormalization of a random
 416 $N \times p$ matrix, and choose $u_n := \sum_i w_n^i u_n^i$, the weighted particle
 417 mean.
 418

5.1. Projected approach and classical AUS techniques

This somewhat technical section establishes the relationship
 between existing AUS algorithms and the projected data approach.
 We consider the EKF-AUS (Trevisan and Palatella 2011; Palatella

et al. 2013, e.g.). EKF-AUS is a modified EKF in which the forecast covariance matrix \mathbf{P}_n^f is replaced by the projected matrix $\Pi_n \mathbf{P}_n^f \Pi_n$, leading to the Kalman gain

$$\mathbf{K}_n = \Pi_n \mathbf{P}_n^f \Pi_n \mathbf{H}^T \left[\mathbf{H} \Pi_n \mathbf{P}_n^f \Pi_n \mathbf{H}^T + \mathbf{R} \right]^{-1}, \quad (34)$$

where the EKF forecast covariance matrix \mathbf{P}_n^f and observation operator $\mathbf{H}_n \equiv \mathbf{H}$ are described in Section 2.1.1. It is clear that the EKF-AUS Kalman gain can be written as a combination of the columns of \mathbf{U}_n .

For comparison, we write down the Kalman gain associated with the data model (19),

$$\mathbf{K}_n = \mathbf{P}_n^f \Pi_{\mathbf{H}} \Pi_n \left[\Pi_n \mathbf{H}^\dagger \left(\mathbf{H} \mathbf{P}_n^f \mathbf{H}^T + \mathbf{R} \right) (\mathbf{H}^\dagger)^T \Pi_n \right]^\dagger. \quad (35)$$

We choose this form to most closely resemble EKF-AUS; the arguments of Theorem 3.1 guarantee that (35) is identical to the Algorithm 1 implementation of the EKF.

The difference between the two Kalman gains is essentially that (35) interchanges the position of \mathbf{H} and Π_n , requiring the use of \mathbf{H}^\dagger in order to do so, but manages to project all terms in the covariance-weighting inverse instead of only the forecast covariance matrix. Unlike the classical AUS gain (34), (35) does not restrict the analysis increment to the unstable subspace. The innovation is $y_n - \mathbf{H}u_n^f$ in classical AUS, but with (35) would be $y_n^p - \Pi_n \Pi_{\mathbf{H}} u_n^f$.

That is, classical AUS uses the full data but restricts the assimilation update to the unstable subspace via (34); Algorithm 1 restricts the innovation to the unstable subspace but the assimilation update can distribute this innovation across the whole of model space. The comparison between these algorithms here is pedagogical, not competitive; the advantages of the EKF-AUS algorithm are well established, while Algorithm 1 effects a reduction in data dimension that we will explore for Particle Filters, not the EKF.

Finally we obtain a form of EKF associated with the projected model (1) and unprojected data. This is essentially a re-derivation of EKF-AUS from the projected framework employed in this paper, confirming that the two are compatible.

Consider a linear or linearized physical model $u_{n+1} = \mathbf{A}_n u_n + \sigma_n$. The projected model has the form $v_{n+1} = \mathbf{B}_n v_n + \Pi_{n+1} \sigma_n$ with $\mathbf{B}_n := \Pi_{n+1} \mathbf{A}_n \Pi_n$, $v_n = \Pi_n u_n$, model error covariances $\mathbf{Q}_{n+1} := \Pi_{n+1} \mathbf{Q} \Pi_{n+1}$, and observation operators $\mathbf{H}_n := \mathbf{H} \Pi_n$. The reduced dimensional projected model has the form $w_{n+1} = \mathbf{T}_n w_n + \mathbf{U}_{n+1}^T \sigma_n$ with $\mathbf{T}_n := \mathbf{U}_{n+1}^T \mathbf{A}_n \mathbf{U}_n$, $w_n = \mathbf{U}_n^T u_n$, model error covariances $\mathbf{Q}_{n+1} := \mathbf{U}_{n+1}^T \mathbf{Q} \mathbf{U}_{n+1}$, and observation operators $\mathbf{H}_n := \mathbf{H} \mathbf{U}_n$. Let $\tilde{\mathbf{P}}_n^f$ and $\hat{\mathbf{P}}_n^f$ denote the forecast covariance matrices for the projected and reduced order models, respectively. Initialising with $\tilde{\mathbf{P}}_0^f = \Pi_0 \mathbf{P}_0^f \Pi_0$ and $\hat{\mathbf{P}}_0^f = \mathbf{U}_0^T \mathbf{P}_0^f \mathbf{U}_0$ we see that by projecting and appropriate modification to the observation operators we obtain projected and reduced order EKF-AUS forecast covariance matrices, respectively.

We will now explore the benefits of the projected data algorithms in an AUS framework, using (33) to calculate the projections, and present examples from the Lorenz 96 system.

6. Numerics

Consider the system of ordinary differential equations introduced in Lorenz (1996),

$$\dot{u}_i = (u_{i+1} - u_{i-2}) u_{i-1} - u_i + F, \quad (36)$$

for $i = 1, \dots, 40$ and $F = 8$. This ‘Lorenz-96’ system is chaotic with 14 positive and 1 neutral Lyapunov exponents. We present experiments in which the deterministic part of the model (2) is given by an integration of (36) for a fixed time.

The primary focus of this section is Algorithm 2, PROJ-OP-PF, compared to the OP-PF and ETKF.

6.1. Methods

The following five DA methods are compared.

1. PROJ-OP-PF, adding no noise on resampling. We compute with each p from 1 to 20.
2. PROJ-OP-PF, adding Gaussian noise on resampling and employing Algorithm 3 with $\alpha = 0.99$. We consider each

478 p from 1 to 20, and consider ten values of the noise on
 479 resampling ω , linearly spaced between 10^{-4} and 10^{-1} .

480 3. OP-PF, adding no noise on resampling.

481 4. OP-PF, adding Gaussian noise on resampling. We consider
 482 ten values of the noise on resampling ω , linearly spaced
 483 between 10^{-4} and 10^{-1} .

484 5. ETKF. We consider ten values of multiplicative inflation for
 485 the forecast ensemble between 1 and 1.1, and additionally
 486 employ additive inflation of \mathbf{Q} on the forecast covariance
 487 matrix.

488 The first and third DA methods are benchmark schemes
 489 that, compared to the second and fourth scheme, clarify the
 490 contribution of noise on resampling to the good performance of
 491 OP-PF and PROJ-OP-PF.

492 The goal of these experiments is to improve on the Optimal
 493 Proposal Particle Filter; the ETKF is present as an example of a
 494 very good DA method for this model. As outlined in Section 2,
 495 the OP-PF is the least degenerate of a wide class of Particle
 496 Filters, and further improving on its performance within the
 497 Particle Filtering class of algorithms is a significant achievement.
 498 The major advantage of PROJ-OP-PF over the ETKF is, as with
 499 any Particle Filter, when filtering nonGaussian prior/posterior
 500 distributions.

501 6.2. Experimental setup

502 We perform identical twin experiments: the true system state is
 503 generated by the model update equation (2), and the role of DA is
 504 to correct for errors in the initial condition and the realizations of
 505 the model noise. All experiments fix the following.

- 506 • the number of ensemble members, $L = 20$.
- 507 • the initial conditions for each ensemble sampled from a
 508 Gaussian with spread equal to the model noise, centred on
 509 the true initial condition.
- 510 • model and true system dynamics are simulated with the
 511 fourth order Runge-Kutta scheme with time steps of 0.01,
 512 repeated until an observation time is reached.

- 513 • the methods are spun up by computing and then discarding
 514 1000 analysis steps; performance is then measured over the
 515 following 10,000 analysis steps.
- 516 • each method is repeated twenty times at each of the possible
 517 combinations of the tuning parameters listed above.

518 We present results for six distinct scenarios, or experiments.
 519 The key parameters in each experiment are the model noise \mathbf{Q} ,
 520 the observation covariance \mathbf{R} , the time between observations
 521 $\Delta \equiv t_n - t_{n-1}$, and the proportion of the state that is observed.
 522 The first experiment produces a well-known test regime for the
 523 Ensemble Kalman Filter. The final experiment is chosen according
 524 to recommendations in Majda *et al.* (2014):

525 "Demanding tests for filter performance are the
 526 regimes of spatially sparse, infrequent in time, high-
 527 quality (low observational noise) observations for a
 528 strongly turbulent dynamical system."

529 Our implementation of this test is less extreme than in Majda *et al.*
 530 (2014), but we only employ $L = 20$ particles (compared to 10,000
 531 in the cited paper).

532 Table 1 displays the key experimental parameters.

533 The first two experiments involve small values of \mathbf{Q} . For these the
 534 OP-PF essentially reduces to a bootstrap Particle Filter, which is
 535 not accurate for the high-dimensional Lorenz-96 system (RMSE
 536 around 5). In these two experiments *only*, we additively inflate
 537 the value of \mathbf{Q} used in the OP-PF and PROJ-OP-PF, (14)–(17)
 538 and (28), by $0.3\mathbf{I}_{40}$. This modification retains some of the optimal
 539 proposal update and stabilises both OP-PF and PROJ-OP-PF
 540 (though the best performing PROJ-OP-PF, in experiment 3, does
 541 not use it).

542 6.3. Results

543 On average out of all experiments, PROJ-OP-PF has 13% less
 544 RMSE than OP-PF, and resamples about half as often (49% of the
 545 time). The advantages of PROJ-OP-PF are more pronounced in
 546 Experiments 1–3, but there is some improvement over OP-PF in all
 547 cases. These statistics are computed by comparing the best-tuned
 548 PROJ-OP-PF to the best-tuned OP-PF (each averaged over 10^4

	Q	R	Δ	observation spacing
Exp. 1	$(0.01)^2 \mathbf{I}_{40}$	\mathbf{I}_{40}	0.05	all vars observed
Exp. 2	$(0.01)^2 \mathbf{I}_{40}$	$(0.5)^2 \mathbf{I}_{40}$	0.05	all vars
Exp. 3	$(0.1)^2 \mathbf{I}_{40}$	$(0.5)^2 \mathbf{I}_{40}$	0.05	all vars
Exp. 4	$(0.1)^2 \mathbf{I}_{40}$	$(0.5)^2 \mathbf{I}_{20}$	0.05	every second var
Exp. 5	$(0.1)^2 \mathbf{I}_{40}$	$(0.1)^2 \mathbf{I}_{20}$	0.05	every second var
Exp. 6	$(0.1)^2 \mathbf{I}_{40}$	$(0.1)^2 \mathbf{I}_{20}$	0.1	every second var

Table 1. Key parameters for the Lorenz-96 experiments in Section 6. Each row is a different experiment and the columns show the model noise, observation noise, time interval between observations and proportion of state space observed. The first experiment resembles a classic Lorenz-96 filtering experiment. The final experiment has larger model uncertainty, more accurate observations (harder for Particle Filters to avoid degeneracy), and infrequent, partial observations. Only one parameter changes between each experiment; in the online version of this article, that parameter is colored blue.

	OP-PF:		PROJ-OP-PF:	
	RMSE	Resamp	RMSE	Resamp
Exp. 1	0.71	59%	0.53	8%
Exp. 2	0.42	57%	0.35	5%
Exp. 3	0.42	58%	0.36	5%
Exp. 4	1.78	58%	1.68	49%
Exp. 5	0.81	62%	0.72	53%
Exp. 6	0.57	61%	0.53	57%

Table 2. Summary, for each of the experiments described in Section 6.2, of the RMSE and percentage of assimilation steps that trigger resampling. Results only listed for the best tuned (minimal RMSE) implementations of OP-PF and PROJ-OP-PF. The projected scheme beats OP-PF in all cases, ranging from a modest (7%) improvement in both RMSE and resampling in Experiment 6, to a 25% improvement in RMSE and 85% reduction in the number of resampling steps in Experiment 1.

549 time steps and over 20 repetitions of each experiment, at all tuning
 550 parameters described in section 6.1). Full details for the tuning and
 551 error scaling are presented in detail in the following subsection.
 552 The best-tuned results for each experiment are summarised in
 553 Table 2.

554 6.4. Detailed results

555 The optimally tuned parameters for each experiment are reported
 556 in Table 3. Figures 2–7 show how varying the projection
 557 rank p affects PROJ-OP-PF, compared to optimally tuned
 558 implementations of the OP-PF and ETKF, in each experiment. For
 559 each plot the tuned ETKF shows roughly what is achievable with
 560 an ensemble of size 20, while OP-PF displays the poor filtering
 561 performance of Particle Filters in high dimensions. PROJ-OP-PF
 562 reduces, by up to half, the gap in RMSE between OP-PF and
 563 ETKF.

	OP-PF:	PROJ-OP-PF:
	ω	p
Exp. 1	0	1
Exp. 2	0	1
Exp. 3	0	1
Exp. 4	0.02 (−1%)	5 0.01 (−1%)
Exp. 5	0.01 (−4%)	7 10^{-4} (−5%)
Exp. 6	10^{-4} (−2%)	9 10^{-4} (−3%)

Table 3. Summary, for each of the experiments described in Section 6.2, of the optimal tuning parameters for OP-PF and PROJ-OP-PF. Numbers in parentheses show the percentage reduction of the RMSE by employing noise on resampling, compared to the benchmark OP-PF and PROJ-OP-PF schemes. The optimal inflation for ETKF was 1.09 or 1.1 in all cases.

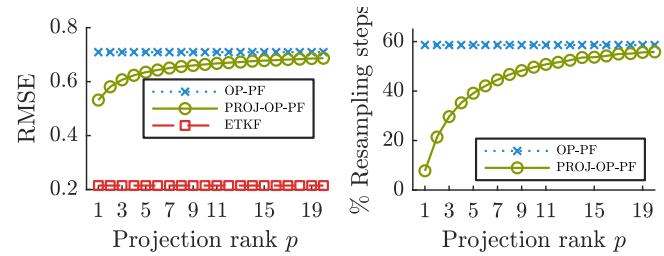


Figure 2. PROJ-OP-PF results scaling p for experiment 1 (compare Tables 1, 2 and 3). The best tuned PROJ-OP-PF covers 36% of the gap from the OP-PF error to the, probably optimal, ETKF error, and reduces the proportion of resampling to almost none.

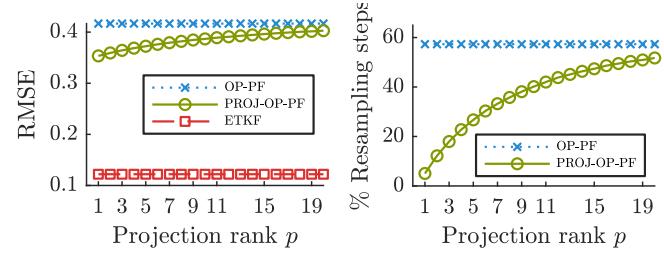


Figure 3. PROJ-OP-PF results scaling p for experiment 2. The best tuned PROJ-OP-PF scheme covers 21% of the gap from the best tuned OP-PF to the ETKF.

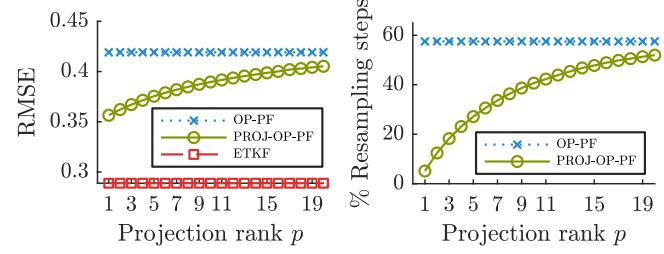


Figure 4. PROJ-OP-PF results scaling p for experiment 3. The best tuned PROJ-OP-PF scheme covers 48% of the gap from the best tuned OP-PF to the ETKF.

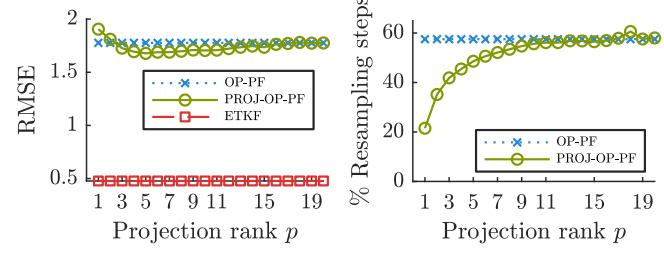


Figure 5. PROJ-OP-PF results scaling p for experiment 4. This experiment is by far the most challenging for PROJ-OP-PF, which at best has RMSE 1.68 (and resamples on 49% of assimilation steps) compared to RMSE 1.78 (58%) for OP-PF. One can instead ($p = 2$) reduce resampling by a factor of about half, corresponding to a much less tight ensemble and better higher order statistics, without greatly affecting the RMSE.

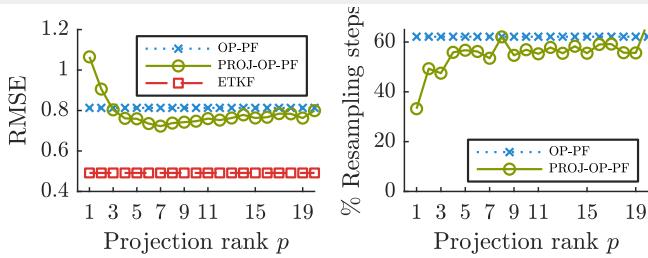


Figure 6. PROJ-OP-PF results scaling p for experiment 5. Compared to OP-PF one can tune PROJ-OP-PF for RMSE, covering 28% of the gap from the best tuned OP-PF to the ETKF, or ($p = 3$) reduce resampling without affecting the RMSE.

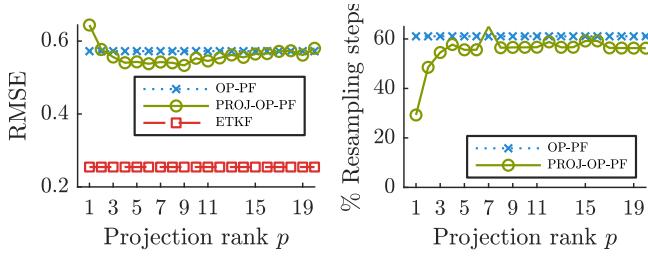


Figure 7. PROJ-OP-PF results scaling p for experiment 6. Compared to OP-PF one can tune PROJ-OP-PF for RMSE, covering 12% of the gap from the best tuned OP-PF to the ETKF, or ($p = 2$) reduce resampling without affecting the RMSE.

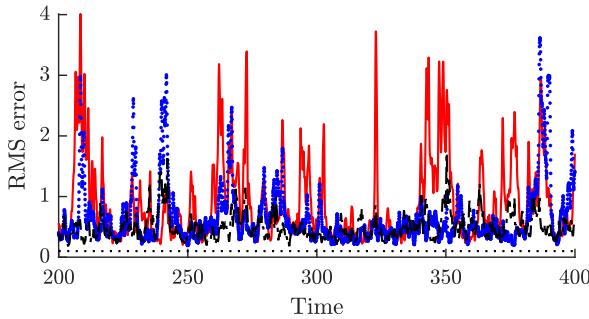


Figure 8. RMSE over time of OP-PF (solid line, red in the online version), PROJ-OP-PF (dots, blue in the online version) and ETKF (dashed, black in the online version). OP-PF frequently spikes in error, while PROJ-OP-PF spikes less frequently. For visualisation we plot the rolling mean RMSE of every five assimilation steps and restrict the plot to a subset of assimilation times.

564 The RMSE over time from all methods from a subset of one run
 565 from experiment 5 is visualised in Figure 8. We see that both
 566 Particle Filter methods have similar *minimum* RMSE to the ETKF,
 567 but also frequently spike in error. The advantage of PROJ-OP-
 568 PF is that the error spikes are much less frequent, and smaller
 569 in amplitude. The more severe degeneracy of OP-PF, compared
 570 to PROJ-OP-PF, is visible in Figure 9, where an ensemble from
 571 experiment 5 is shown by plotting six state variables at the 6000-
 572 th assimilation step.

573 The ETKF is generally stable with low RMSE for multiplicative
 574 inflation equal to or larger than 1.02 (1.04 in experiments 4-
 575 6). However it may be under-dispersive with spread lower than
 576 RMSE, and for this reason the optimal inflation was 1.1 in all
 577 cases. For each experiment in order, the spread in the ETKF was

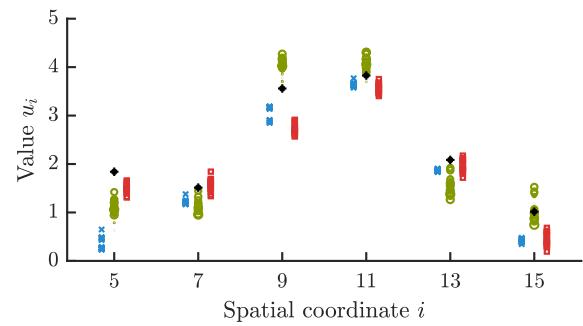


Figure 9. Details of all ensembles after an assimilation step, plotting six spatial coordinates. Coloured splotches show the ensembles, each plotted slightly apart in the horizontal direction for visualisation. From left to right in each coordinate: OP-PF (blue crosses in the online version), which has just resampled, is tightly spread in four of the six plotted coordinates, and has significant error in many; PROJ-OP-PF (green circles) has a larger spread and tends to be closer to the truth (black star). The EnKF (red squares) is closest to the truth in four of six coordinates. Note that none of the variables plotted in this figure were observed.

0.20, 0.10, 0.13, 0.19, 0.06, 0.06. The final three spreads can be
 578 widened to approximately equal the RMSE by taking a much
 579 larger value of the multiplicative inflation, 1.3. The additive
 580 inflation of \mathbf{Q} is needed in experiments 4-6, else the ETKF suffers
 581 large RMSE values (> 3).
 582

6.5. Compete with EnKF by including model error

The previous experiments outline, over diverse scenarios, the
 584 advantage of PROJ-OP-PF over OP-PF. We now adapt Harlim
 585 and Hunt (2007)'s example scenario in which non-Gaussian
 586 filters can compete with the EnKF, and confirm the advantage
 587 of PROJ-OP-PF in that case. The idea is to introduce model
 588 error: all DA methods will obtain forecasts using forcing $F =$
 589 6 in (36), whereas the true system state is generated with
 590 the original forcing $F = 8$. Experiments in this section will
 591 otherwise replicate the prior sections: experimental setup as
 592 given in Section 6.2, and key parameters for each experiment
 593 given in Table 1. The principal result is that PROJ-OP-PF out-
 594 performs the tuned ETKF (which, as before, performs best with
 595 larger multiplicative inflation) in three of six experiments, with
 596 roughly equal performance in a fourth. Given that the posterior
 597 for the Lorenz-96 system is generally almost Gaussian, this is a
 598 remarkable result. Figure 10 shows the relative performance of
 599 all three schemes for Experiment Three. Summary statistics are
 600 provided for all DA schemes in Table 4.

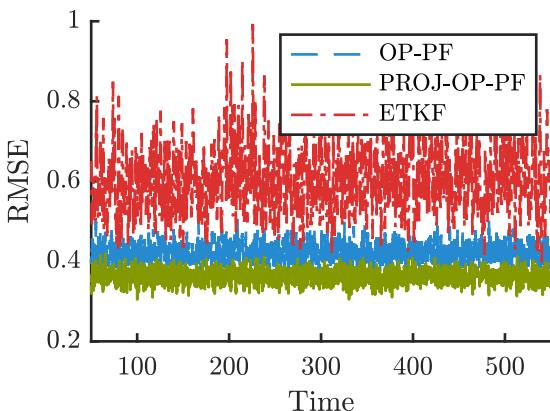


Figure 10. RMSE over time for a single realization of Experiment Three with model error, described in Section 6.5. For visualisation purposes we plot the rolling mean of the RMSE over ten assimilation steps. Our algorithm PROJ-OP-PF now out-performs the ETKF, and maintains its advantage over OP-PF.

	OP-PF:		PROJ-OP-PF:		ETKF:
	RMSE	Res	RMSE	Res	
Exp. 1	0.73	58%	0.60	10%	1.29
Exp. 2	0.42	57%	0.36	5%	1.25
Exp. 3	0.42	57%	0.36	6%	0.61
Exp. 4	1.97	58%	1.88	50%	1.25
Exp. 5	1.40	84%	1.17	100*%	1.16
Exp. 6	1.25	95%	1.15	100*%	0.91

Table 4. Statistics for experiments with model error (compare Table 2 without model error). Tuned parameters are given in Table 3, except for Experiments 5 and 6 in which PROJ-OP-PF uses $\omega = 0.1$. The 100% resampling for PROJ-OP-PF in those experiments can be reduced if ω is smaller (at a commensurate cost to RMSE.) For three experiments, including the standard benchmark for the EnKF, Experiment One, PROJ-OP-PF significantly outperforms both OP-PF and the ETKF, with drastically reduced resampling compared to OP-PF.

6.6. Outlook

We now place the results of our experiments in more context. A key difference between our tuned PROJ-OP-PF and related works is the optimal selection of the subspace dimension: for our Experiments One to Three, the optimal dimension was $p = 1$, and the largest value considered was $p = 9$. This is in stark contrast with AUS: in (Trevisan and Palatella 2011; Palatella *et al.* 2013, e.g.) the optimal subspace dimension is typically the number of positive and stable Lyapunov exponents (14 in our experiments), and Grudzien *et al.* (2018b,a) incorporate model error and conclude the optimal subspace dimension is larger still. Quinn *et al.* (2020) consider reduced rank DA in three coupled Lorenz-63 systems and conclude that the optimal subspace dimension is the Kaplan-Yorke dimension, again larger than the dimension of the unstable-neutral subspace. In some manner, all of the above papers are considering lowering the dimension of the model—whether for reduced order modelling

or to represent some ensemble in model space. By contrast our approach projects the data, and preserves the full rank model forecast.

We consider three reasons for the smaller subspace dimension. First note that the PROJ-OP-PF algorithm uses the data in two ways: the original, non-projected form is used in an EnKF-like step to update the particles, then the reduced order data is used to update the particle weights. So we can understand that our method may still perform well with small p if the *combination* of the full rank particle update with the low rank weight update is sufficient to approximate the posterior distribution.

Secondly, we conjecture that the effect of reduced data may relate to synchronization. Pecora and Carroll (1990, 1991); Pecora *et al.* (1997) show that if one inserts values for certain variables into a dynamical system, and solves the dynamical system with essentially arbitrary initial conditions in the remaining variables, then the solution synchronizes with the true solution. In this work we are forcing with projected data that for each time is a linear combination of the model variables. We conjecture that this may provide a synchronizing effect even if the dimension of the projected data is lower than the number of positive Lyapunov exponents.

Lastly we note that there is a reason to anticipate that the optimal dimensions for projected data and projected models may differ. In the reduced order physical models, the optimal subspace dimension has more to do with sufficiently resolving solutions so that particles provide good approximation of the true solution. If we are projecting the data (only), then the particles can well approximate the true solution and we only need to determine their importance, i.e., their weights. The projected data models may, due to containing less information, update the weights more slowly, but they also avoid the degeneracy issues with large data dimension (see, e.g., Bengtsson *et al.* (2008); Snyder *et al.* (2008); Van Leeuwen (2009)).

This manuscript provides a framework for the use of projections in reduced order data models. The framework empowers future research to investigate optimal projections for various classes of problems. Some contemporary works provide insight: Beeson

657 and Sri Namachchivaya (2020) find best results using projections
 658 calculated from future time intervals, and (Carrassi *et al.* 2020,
 659 Section 5.1) argue that there is limited benefit to the projected
 660 approach if the model is deterministic. The second reference
 661 includes an explicit proof that, for a perfect deterministic model
 662 with uncorrelated observation errors ($\mathbf{R} = \mathbf{I}$), projecting to a
 663 lower dimension does not reduce filter degeneracy; therefore we
 664 learn that the benefit of employing projections is related to model
 665 error, model noise, and/or correlated observation errors. Along
 666 these lines, Albarakati *et al.* (2021) extend the PROJ-OP-PF
 667 scheme to project both the model and data and consider other
 668 reduced order modeling techniques in addition to AUS. Numerical
 669 results with a 38100-dimensional Shallow Water Equation show
 670 that the projected model dimension should be at least the Kaplan-
 671 Yorke dimension, but the projected data dimension can be much
 672 smaller.

673 7. Discussion

674 In this work a new approach to DA has been derived that allows
 675 for dimension reduction of the data using a projection defined
 676 in state space. The chief application has been Particle Filters
 677 Assimilating in the Unstable Subspace, which the classical AUS
 678 approach is unsuitable for because ensemble methods already
 679 project the forecast strongly into the unstable subspace (Bocquet
 680 and Carrassi 2017). By contrast the new approach sharply reduces
 681 filter degeneracy in a predictable fashion, improves filter accuracy
 682 and allows one to construct a sensible resampling scheme that
 683 adds more noise in more uncertain directions. Algorithms resting
 684 on the projected DA approach were tested on the chaotic Lorenz
 685 96 system that provides a challenging scenario for particle filters.
 686 The main algorithm tested was a particle filter, PROJ-OP-PF
 687 that mixes projected and unprojected data based on the optimal
 688 proposal. The discrete QR technique used to find the unstable
 689 subspace in this work is rigorously justified and the additional
 690 cost incurred by it is proportional to employing an ensemble size
 691 of the dimension of the projected subspace. The PROJ-OP-PF
 692 scheme makes only the most basic use of the data in the space
 693 orthogonal to whatever projection is employed. Projecting the data
 694 to lower the dimension of the observations generally reduces both
 695 the RMSE and the resampling as compared to OP-PF. Resampling

696 is essentially always better in our experiments and the RMSE is
 697 only worse for PROJ-OP-PF with extremely low rank projections
 698 in experiments 4 and 5.

699 Interesting extensions to these techniques include the use of
 700 nonlinear observation operators and the application of spatial
 701 localization techniques. Due to the modular approach we have
 702 taken by introducing projected observations, existing localization
 703 techniques for particle filters can be applied as they would with the
 704 original unprojected observations. In addition, the combination
 705 of projected physical model and projected observational model
 706 allow for a combination of localized and unlocalized systems
 707 to be employed much like the way projected and unprojected
 708 data models are used in Proj-OP-PF. In this work we have
 709 considered full rank, linear observation operators. There are some
 710 straightforward extensions to nonlinear observation operators both
 711 when using state space based projections and observation space
 712 based projections. For state space based projections, the full
 713 rank assumption is not necessary, only the existence of \mathbf{H}^\dagger
 714 which can be formed using the SVD. For nonlinear observation
 715 operators, the linearization of \mathbf{H} could be used to form $D\mathbf{H}(x)^\dagger$.
 716 Observation space projections, e.g., using PCA, remove the
 717 need for \mathbf{H}^\dagger or $D\mathbf{H}(x)^\dagger$ and can be applied directly to obtain
 718 projected data models for both linear and nonlinear observation
 719 operators.

720 Future work will generalise the projected DA approach to employ
 721 two assimilation methods, one in the projected and one in the
 722 orthogonal space. Such manipulations are done in Majda *et al.*
 723 (2014); Slivinski *et al.* (2015), for example, and formulated for
 724 model error in AUS in section 3.2 of Grudzien *et al.* (2018b).
 725 The methods formulated in this paper have also been extended
 726 to particle filters that project both model and data; the resulting
 727 scheme was successfully applied to a high-dimensional shallow
 728 water model in Albarakati *et al.* (2021).

729 Acknowledgements

730 JM acknowledges the support of ONR grant N00014-18-1-
 731 2204, NSF grant DMS-1722578, and the Australian Research
 732 Council Discovery Project DP180100050. EVV acknowledges the
 733 support of NSF grants DMS-1714195 and DMS-1722578. The

734 authors are grateful to Alberto Carrassi for helpful feedback on
 735 an early version of this work. This work was supported with
 736 supercomputing resources provided by the Phoenix HPC service
 737 at the University of Adelaide.

738 **References**

- 734 Carrassi A, Trevisan A, Descamps L, Talagrand O, Uboldi F. 2008b. 779
 735 Controlling instabilities along a 3DVar analysis cycle by assimilating in 780
 736 the unstable subspace: a comparison with the EnKF. *Nonlinear Process. 781
 737 Geophys.* **15**: 503–521. 782
- 734 Carrassi A, Trevisan A, Uboldi F. 2007. Adaptive observations and 783
 735 assimilation in the unstable subspace by breeding on the data-assimilation 784
 736 system. *Tellus A: Dynamic Meteorology and Oceanography* **59**(1): 101– 785
 737 113. 786
- 734 Chorin A, Morzfeld M, Tu X. 2010. Implicit particle filters for data 787
 735 assimilation. *Communications in Applied Mathematics and Computational 788
 736 Science* **5**(2): 221–240. 789
- 734 de Leeuw B, Dubinkina S, Frank J, Steyer A, Tu X, Van Vleck E. 2018. 790
 735 Projected shadowing-based data assimilation. *SIAM Appld. Dyn. Sys.* . 791
- 734 Dieci L, Van Vleck ES. 2007. Lyapunov and Sacker-Sell spectral intervals. *J. 792
 735 Dynam. Differential Equations* **19**(2): 265–293. 793
- 734 Dieci L, Van Vleck ES. 2015. Lyapunov exponents: Computation. In: 794
 735 *Encyclopedia of Applied and Computational Mathematics*, Engquist B (ed). 795
 736 Springer-Verlag, pp. 834–838. 796
- 734 Doucet A, De Freitas N, Gordon N. 2001. An introduction to sequential Monte 797
 735 Carlo methods. In: *Sequential Monte Carlo methods in practice*, Doucet A, 798
 736 De Freitas N, Gordon N (eds), Springer, pp. 3–14. 799
- 734 Doucet A, Godsill S, Andrieu C. 2000. On sequential Monte Carlo sampling 800
 735 methods for Bayesian filtering. *Statistics and computing* **10**(3): 197–208. 801
- 734 Evensen G. 1994. Sequential data assimilation with a nonlinear quasi- 802
 735 geostrophic model using Monte Carlo methods to forecast error statistics. 803
Journal of Geophysical Research: Oceans **99**(C5): 10 143–10 162. 804
- 734 Evensen G. 2009. *Data assimilation: the ensemble Kalman filter*. Springer 805
 735 Science & Business Media. 806
- 734 Farchi A, Bocquet M. 2018. Comparison of local particle filters and new 807
 735 implementations. *Nonlinear Processes in Geophysics* **25**(4): 765–807. 808
- 734 Frank J, Zhuk S. 2018. A detectability criterion and data assimilation for 809
 735 nonlinear differential equations. *Nonlinearity* **31**(11): 5235–5257. 810
- 734 González-Tokman C, Hunt BR. 2013. Ensemble data assimilation for 811
 735 hyperbolic systems. *Physica D: Nonlinear Phenomena* **243**(1): 128–142. 812
- 734 Grudzien C, Carrassi A, Bocquet M. 2018a. Asymptotic forecast uncertainty 813
 735 and the unstable subspace in the presence of additive model error. 814
SIAM/ASA J. Uncertain. Quantif. **6**(4): 1335–1363. 815
- 734 Grudzien C, Carrassi A, Bocquet M. 2018b. Chaotic dynamics and the role 816
 735 of covariance inflation for reduced rank kalman filters with model error. 817
Nonlinear Processes in Geophysics **25**(3): 633–648. 818
- 734 Harlim J, Hunt BR. 2007. A non-Gaussian Ensemble Filter for Assimilating 819
 735 Infrequent Noisy Observations. *Tellus A: Dynamic Meteorology and 820
 736 Oceanography* **59**(2): 225–237. 821
- 734 Law K, Stuart A, Konstantinos Z. 2015. *Data assimilation: A mathematical 822
 735 introduction*, vol. 62. Springer Texts in Applied Mathematics, ISBN 978-3- 823
 736 319-20324-9. 824

- 825 Law KJH, Sanz-Alonso D, Shukla A, Stuart AM. 2016. Filter accuracy for the 826 Lorenz 96 model: fixed versus adaptive observation operators. *Phys. D* **325**: 827 1–13.
- 828 Lorenz EN. 1996. Predictability - a problem partly solved. In: *Proceedings of* 829 *seminar on Predictability*, vol. 1, Palmer T, Hagedorn R (eds). ECMWF, 830 Cambridge University Press: Reading, UK, pp. 1–18.
- 831 Maclean J, Santitissadeekorn N, Jones CKRT. 2017. A coherent structure 832 approach for parameter estimation in Lagrangian data assimilation. *Phys.* 833 *D* **360**: 36–45.
- 834 Majda AJ, Qi D, Sapsis TP. 2014. Blended particle filters for large-dimensional 835 chaotic dynamical systems. *Proc. Natl. Acad. Sci. USA* **111**(21): 7511– 836 7516.
- 837 Morzfeld M, Adams J, Lunderman S, Orozco R. 2018. Feature-based data 838 assimilation in geophysics. *Nonlinear Processes in Geophysics* **25**(2): 355– 839 374.
- 840 Morzfeld M, Tu X, Atkins E, Chorin AJ. 2012. A random map implementation 841 of implicit filters. *J. Comput. Phys.* **231**(4): 2049–2066.
- 842 Palatella L, Carrassi A, Trevisan A. 2013. Lyapunov vectors and assimilation 843 in the unstable subspace: theory and applications. *Journal of Physics A: 844 Mathematical and Theoretical* **46**(25): 254 020.
- 845 Pecora LM, Carroll TL. 1990. Synchronization in chaotic systems. *Physical 846 Review Letters* **64**(8): 821–824. Publisher: American Physical Society.
- 847 Pecora LM, Carroll TL. 1991. Driving systems with chaotic signals. *Physical 848 Review A* **44**(4): 2374–2383. Publisher: American Physical Society.
- 849 Pecora LM, Carroll TL, Johnson GA, Mar DJ, Heagy JF. 1997. Fundamentals 850 of synchronization in chaotic systems, concepts, and applications. *Chaos: An 851 Interdisciplinary Journal of Nonlinear Science* **7**(4): 520–543. 852 Publisher: American Institute of Physics.
- 853 Poterjoy J, Anderson JL. 2016. Efficient assimilation of simulated 854 observations in a high-dimensional geophysical system using a localized 855 particle filter. *Monthly Weather Review* **144**(5): 2007–2020.
- 856 Potthast R, Walter A, Rhodin A. 2019. A localized adaptive particle filter 857 within an operational nwp framework. *Monthly Weather Review* **147**(1): 858 345–362.
- 859 Qi D, Majda AJ. 2015. Blended particle methods with adaptive subspaces for 860 filtering turbulent dynamical systems. *Phys. D* **298/299**: 21–41.
- 861 Quinn C, O& TJ, apos, Kane, Kitsios V. 2020. Application of a local 862 attractor dimension to reduced space strongly coupled data assimilation for 863 chaotic multiscale systems. *Nonlinear Processes in Geophysics* **27**(1): 51– 864 51. Publisher: Copernicus GmbH.
- 865 Reddy AS, Apté A, Vadlamani S. 2020. Asymptotic properties of linear filter 866 for deterministic processes. *Systems Control Lett.* **139**: 104 676, 8.
- 867 Reich S, Cotter C. 2015. *Probabilistic forecasting and Bayesian data 868 assimilation*. Cambridge University Press.
- 869 Sapsis T. 2010. Dynamically orthogonal field equations. PhD thesis, 870 Massachusetts Institute of Technology, Department of Mechanical 871 Engineering.
- 872 Sapsis TP, Lermusiaux PFJ. 2009. Dynamically orthogonal field equations for 873 continuous stochastic dynamical systems. *Phys. D* **238**(23–24): 2347–2360. 873
- 874 Slivinski L, Spiller E, Apté A, Sandstede B. 2015. A hybrid particle–ensemble 875 kalman filter for lagrangian data assimilation. *Monthly Weather Review* 876 **143**(1): 195–211.
- 877 Snyder C. 2011. Particle filters, the "optimal" proposal and high-dimensional 878 systems. In: *Proceedings of the ECMWF Seminar on Data Assimilation for 879 atmosphere and ocean*. pp. 1–10.
- 880 Snyder C, Bengtsson T, Bickel P, Anderson J. 2008. Obstacles to high- 881 dimensional particle filtering. *Monthly Weather Review* **136**(12): 4629– 882 4640.
- 883 Snyder C, Bengtsson T, Morzfeld M. 2015. Performance bounds for particle 884 filters using the optimal proposal. *Monthly Weather Review* **143**(11): 4750– 885 4761.
- 886 Sondergaard T, Lermusiaux PF. 2013. Data assimilation with Gaussian mixture 887 models using the dynamically orthogonal field equations. Part I: Theory and 888 scheme. *Monthly Weather Review* **141**(6): 1737–1760.
- 889 Tranninger M, Seeber R, Steinberger M, Horn M. 2020. Uniform detectability 890 of linear time varying systems with exponential dichotomy. *IEEE Control 891 Systems Letters* : 1–1.
- 892 Trevisan A, D'Isidoro M, Talagrand O. 2010. Four-dimensional variational 893 assimilation in the unstable subspace and the optimal subspace dimension. 893
- 894 *Q.J.R. Meteorol. Soc.* **136**: 487–496.
- 895 Trevisan A, Palatella L. 2011. On the Kalman filter error covariance collapse 896 into the unstable subspace. *Nonlinear Processes in Geophysics* **18**(2): 243– 897 250.
- 898 Tsukuma H, Kubokawa T. 2015. Estimation of the mean vector in a singular 899 multivariate normal distribution. *J. Multivariate Anal.* **140**: 245–258.
- 900 Van Leeuwen PJ. 2009. Particle Filtering in Geophysical Systems. *Monthly 901 Weather Review* **137**(12): 4089–4114. Publisher: American Meteorological 902 Society.
- 903 van Leeuwen PJ. 2010. Nonlinear data assimilation in geosciences: 904 an extremely efficient particle filter. *Quarterly Journal of the Royal 905 Meteorological Society* **136**(653): 1991–1999.
- 906 van Leeuwen PJ. 2012. Particle filters for the geosciences. *Advanced Data 907 Assimilation for Geosciences: Lecture Notes of the Les Houches School of 908 Physics: Special Issue* : 291.
- 909 van Leeuwen PJ, Künsch HR, Nerger L, Potthast R, Reich S. 2019. Particle 910 filters for high-dimensional geoscience applications: A review. *Quarterly 910 Journal of the Royal Meteorological Society* **145**(723): 2335–2365.
- 911