

Measurement of branching fractions and search for CP violation in $D^0 \rightarrow \pi^+\pi^-\eta$, $D^0 \rightarrow K^+K^-\eta$, and $D^0 \rightarrow \phi\eta$ at Belle



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ABSTRACT: We measure the branching fractions and CP asymmetries for the singly Cabibbo-suppressed decays $D^0 \rightarrow \pi^+\pi^-\eta$, $D^0 \rightarrow K^+K^-\eta$, and $D^0 \rightarrow \phi\eta$, using 980 fb⁻¹ of data from the Belle experiment at the KEKB e^+e^- collider. We obtain

$$\begin{aligned} \mathcal{B}(D^0 \rightarrow \pi^+\pi^-\eta) &= [1.22 \pm 0.02 (\text{stat}) \pm 0.02 (\text{syst}) \pm 0.03 (\mathcal{B}_{\text{ref}})] \times 10^{-3}, \\ \mathcal{B}(D^0 \rightarrow K^+K^-\eta) &= [1.80^{+0.07}_{-0.06} (\text{stat}) \pm 0.04 (\text{syst}) \pm 0.05 (\mathcal{B}_{\text{ref}})] \times 10^{-4}, \\ \mathcal{B}(D^0 \rightarrow \phi\eta) &= [1.84 \pm 0.09 (\text{stat}) \pm 0.06 (\text{syst}) \pm 0.05 (\mathcal{B}_{\text{ref}})] \times 10^{-4}, \end{aligned}$$

where the third uncertainty (\mathcal{B}_{ref}) is from the uncertainty in the branching fraction of the reference mode $D^0 \rightarrow K^-\pi^+\eta$. The color-suppressed decay $D^0 \rightarrow \phi\eta$ is observed for the first time, with very high significance. The results for the CP asymmetries are

$$\begin{aligned} A_{CP}(D^0 \rightarrow \pi^+\pi^-\eta) &= [0.9 \pm 1.2 (\text{stat}) \pm 0.5 (\text{syst})]\%, \\ A_{CP}(D^0 \rightarrow K^+K^-\eta) &= [-1.4 \pm 3.3 (\text{stat}) \pm 1.1 (\text{syst})]\%, \\ A_{CP}(D^0 \rightarrow \phi\eta) &= [-1.9 \pm 4.4 (\text{stat}) \pm 0.6 (\text{syst})]\%. \end{aligned}$$

The results for $D^0 \rightarrow \pi^+\pi^-\eta$ are a significant improvement over previous results. The branching fraction and A_{CP} results for $D^0 \rightarrow K^+K^-\eta$, and the A_{CP} result for $D^0 \rightarrow \phi\eta$, are the first such measurements. No evidence for CP violation is found in any of these decays.

KEYWORDS: Branching fraction, Charm physics, CP violation, e^+e^- Experiments

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1 Introduction

Singly Cabibbo-suppressed (SCS) decays of charmed mesons provide a promising opportunity to study CP violation in the charm sector. Within the Standard Model, CP violation in charm decays is expected to be of the order of 10^{-3} or smaller [1, 2], and thus challenging to observe. SCS decays are of special interest, as interference that includes a new physics amplitude could lead to large CP violation. The CP asymmetry between $D^0 \rightarrow f$ and $\bar{D}^0 \rightarrow \bar{f}$ decays (A_{CP}) is defined as

$$A_{CP} = \frac{\mathcal{B}(D^0 \rightarrow f) - \mathcal{B}(\bar{D}^0 \rightarrow \bar{f})}{\mathcal{B}(D^0 \rightarrow f) + \mathcal{B}(\bar{D}^0 \rightarrow \bar{f})}. \quad (1.1)$$

The only observation of CP violation in the charm sector to date is from the LHCb experiment, where a difference in A_{CP} between the SCS $D^0 \rightarrow K^+K^-$ and $D^0 \rightarrow \pi^+\pi^-$ decays [3] was observed: $\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$. In this paper, we investigate two analogous SCS decays, $D^0 \rightarrow \pi^+\pi^-\eta$ and $D^0 \rightarrow K^+K^-\eta$. A search for a CP asymmetry in the first decay was performed by the BESIII experiment; the resulting precision was 6% [4]. There have been no results for $D^0 \rightarrow K^+K^-\eta$ decays to date. Theoretically, it is difficult to predict CP asymmetries for three-body decays, while some predictions exist for intermediate two-body processes: $A_{CP}(D^0 \rightarrow \rho^0\eta)$ is predicted to be -0.53×10^{-3} from tree amplitudes

alone, and -0.23×10^{-3} after considering QCD-penguin and weak penguin-annihilation [1]. The asymmetry $A_{CP}(D^0 \rightarrow \phi\eta)$ is predicted to be zero in several theoretical models [1]. A precise measurement of branching fractions (\mathcal{B}) for these three-body decays is an important step towards searching for CP violation in these channels.

In this paper we utilize the full Belle data sample of 980 fb^{-1} to measure \mathcal{B} and A_{CP} for three SCS decays: $D^0 \rightarrow \pi^+\pi^-\eta$, $D^0 \rightarrow K^+K^-\eta$, and $D^0 \rightarrow \phi\eta$. All \mathcal{B} measurements are performed relative to the Cabibbo-favored (CF) decay $D^0 \rightarrow K^-\pi^+\eta$, which has been well-measured (with a fractional uncertainty $\delta\mathcal{B}/\mathcal{B} \sim 3\%$ [5]) by both Belle [6] and BESIII [7]. The current world average for $\mathcal{B}(D^0 \rightarrow \pi^+\pi^-\eta)$ has a fractional uncertainty $\delta\mathcal{B}/\mathcal{B} \sim 6\%$ [5]. The branching fraction for $D^0 \rightarrow \phi\eta$ was previously measured by Belle with 78 fb^{-1} of data [8]; the measurement reported here uses an order of magnitude more data and supersedes that result. BESIII found evidence for $D^0 \rightarrow \phi\eta$ (4.2σ) [9] and observed a non- ϕ $D^0 \rightarrow K^+K^-\eta$ component (5.2σ) [7].

To identify the flavor of the neutral D meson when produced, we reconstruct $D^{*+} \rightarrow D^0\pi_s^+$ and $D^{*-} \rightarrow \bar{D}^0\pi_s^-$ decays; the charge of the daughter π_s^\pm (which has low momentum and is referred to as the “slow” pion) identifies whether the D meson is D^0 or \bar{D}^0 . The raw asymmetry measured (A_{raw}) receives contributions from several sources:

$$A_{\text{raw}} = A_{CP}^{D^0 \rightarrow f} + A_{\text{FB}}^{D^{*+}} + A_{\epsilon}^{\pi_s}, \quad (1.2)$$

where $A_{CP}^{D^0 \rightarrow f}$ is the CP asymmetry for $D^0 \rightarrow f$; $A_{\text{FB}}^{D^{*+}}$ is the forward-backward asymmetry due to γ - Z^0 interference and higher-order QED effects [10] in $e^+e^- \rightarrow c\bar{c}$ collisions; and $A_{\epsilon}^{\pi_s}$ is the asymmetry resulting from a difference in reconstruction efficiencies between π_s^+ and π_s^- . This asymmetry depends on the transverse momentum $p_T(\pi_s)$ and polar angle $\theta(\pi_s)$ of the π_s in the laboratory frame. We correct for this by weighting signal events by a factor $[1 - A_{\epsilon}^{\pi_s}(p_T, \cos\theta)]$ for D^0 decays, and by a factor $[1 + A_{\epsilon}^{\pi_s}(p_T, \cos\theta)]$ for \bar{D}^0 decays. After this weighting, we are left with the π_s -corrected asymmetry

$$A_{\text{corr}}(\cos\theta^*) = A_{CP} + A_{\text{FB}}(\cos\theta^*). \quad (1.3)$$

Since A_{FB} is an odd function of the cosine of the D^{*+} polar angle θ^* in the e^+e^- center-of-mass (CM) frame, and A_{CP} is independent of $\cos\theta^*$, we extract A_{CP} and $A_{\text{FB}}(\cos\theta^*)$ via

$$A_{CP} = \frac{A_{\text{corr}}(\cos\theta^*) + A_{\text{corr}}(-\cos\theta^*)}{2}, \quad (1.4)$$

$$A_{\text{FB}}(\cos\theta^*) = \frac{A_{\text{corr}}(\cos\theta^*) - A_{\text{corr}}(-\cos\theta^*)}{2}. \quad (1.5)$$

Fitting the values of A_{CP} for different $\cos\theta^*$ bins to a constant gives our final measurement of A_{CP} for $D^0 \rightarrow f$.

2 Belle detector and data sets

This measurement is based on the full data set of the Belle experiment, which corresponds to a total integrated luminosity of 980 fb^{-1} [11] collected at or near the $\Upsilon(nS)$ ($n = 1, 2, 3, 4, 5$) resonances. The Belle experiment ran at the KEKB energy-asymmetric collider [12, 13].

The Belle detector is a large-solid-angle magnetic spectrometer consisting of a silicon vertex detector (SVD), a 50-layer central drift chamber (CDC), an array of aerogel threshold Cherenkov counters (ACC), a barrel-like arrangement of time-of-flight scintillation counters (TOF), and an electromagnetic calorimeter comprising CsI(Tl) crystals located inside a superconducting solenoid coil providing a 1.5 T magnetic field. An iron flux-return located outside the coil is instrumented to detect K_L^0 mesons and to identify muons. A detailed description of the detector is given in refs. [11, 14].

We use Monte Carlo (MC) simulated events to optimize selection criteria, study backgrounds, and evaluate the signal reconstruction efficiency. Signal MC events are generated by EVTGEN [15] and propagated through a detector simulation based on GEANT3 [16]. Final-state radiation from charged particles is simulated using the PHOTOS package [17]. Three-body decays are generated according to phase space. An MC sample of “generic” events, corresponding to an integrated luminosity four times that of the data, is used to develop selection criteria. It includes $B\bar{B}$ events and continuum processes $e^+e^- \rightarrow q\bar{q}$, where $q = u, d, s, c$. At the $\Upsilon(5S)$ resonance, the MC includes $B_s^{(*)0}\bar{B}_s^{(*)0}$ events. Selection criteria are optimized by maximizing a figure-of-merit $N_{\text{sig}}/\sqrt{N_{\text{sig}} + N_{\text{bkg}}}$, where N_{sig} and N_{bkg} are the numbers of signal and background events, respectively, expected in a two-dimensional signal region in variables M and Q . The variable M is the invariant mass of the $h^+h^-\eta$ ($h = \pi, K$) combination, and $Q = [M(h^+h^-\eta\pi_s^+) - M(h^+h^-\eta) - m_{\pi_s^+}] \cdot c^2$ is the kinetic energy released in the D^{*+} decay.

3 Event selection and optimization

We reconstruct the signal decays $D^0 \rightarrow \pi^+\pi^-\eta$ and $D^0 \rightarrow K^+K^-\eta$, and the reference decay $D^0 \rightarrow K^-\pi^+\eta$, in which the D^0 originates from $D^{*+} \rightarrow D^0\pi^+$, as follows.¹ Charged tracks are identified as K^\pm or π^\pm candidates using a likelihood ratio $\mathcal{R}_K \equiv \mathcal{L}_K/(\mathcal{L}_K + \mathcal{L}_\pi)$, where \mathcal{L}_K (\mathcal{L}_π) is the likelihood that a track is a K^\pm (π^\pm) based on the photon yield in the ACC, dE/dx information in the CDC, and time-of-flight information from the TOF [18]. Tracks having $\mathcal{R}_K > 0.60$ are identified as K^\pm candidates; otherwise, they are considered as π^\pm candidates. The corresponding efficiencies are approximately 90% for kaons and 95% for pions. Tracks that are highly electron-like ($\mathcal{R}_e > 0.95$) or muon-like ($\mathcal{R}_\mu > 0.95$) are rejected, where the electron and muon likelihood ratios \mathcal{R}_e and \mathcal{R}_μ are determined mainly using information from the ECL and KLM detectors, respectively [19, 20]. Charged tracks are required to have at least two SVD hits in the $+z$ direction (defined as the direction opposite that of the positron beam), and at least two SVD hits in the x - y (transverse) plane. The nearest approach of the π_s^+ track to the e^+e^- interaction point (IP) is required to be less than 1.0 cm in the x - y plane, and less than 3.0 cm along the z axis.

Photon candidates are identified as energy clusters in the ECL that are not associated with any charged track. The photon energy (E_γ) is required to be greater than 50 MeV in the barrel region (covering the polar angle $32^\circ < \theta < 129^\circ$), and greater than 100 MeV in the endcap region ($12^\circ < \theta < 31^\circ$ or $132^\circ < \theta < 157^\circ$). The ratio of the energy deposited

¹Throughout this paper, charge-conjugate modes are implicitly included unless stated otherwise.

in the 3×3 array of crystals centered on the crystal with the highest energy, to the energy deposited in the corresponding 5×5 array of crystals, is required to be greater than 0.80.

Candidate $\eta \rightarrow \gamma\gamma$ decays are reconstructed from photon pairs having an invariant mass satisfying $500 \text{ MeV}/c^2 < M(\gamma\gamma) < 580 \text{ MeV}/c^2$. This range corresponds to about 3σ in $M(\gamma\gamma)$ resolution. The absolute value of the cosine of the $\eta \rightarrow \gamma_1\gamma_2$ decay angle, defined as $\cos\theta_\eta \equiv E(\eta)/p(\eta) \cdot (E_{\gamma_1} - E_{\gamma_2})/(E_{\gamma_1} + E_{\gamma_2})$, is required to be less than 0.85. This retains around 89% of the signal while reducing backgrounds by a factor of two. To further suppress backgrounds, we remove η candidates in which both photon daughters can be combined with other photons in the event to form $\pi^0 \rightarrow \gamma\gamma$ candidate decays satisfying $|M_{\gamma\gamma} - m_{\pi^0}| < 10 \text{ MeV}/c^2$, where m_{π^0} is the nominal π^0 mass [5]. This veto requirement has an efficiency of 95% while reducing backgrounds by a factor of three ($D^0 \rightarrow K^+K^-\eta$) and four ($D^0 \rightarrow \pi^+\pi^-\eta$).

Candidate $D^0 \rightarrow \pi^+\pi^-\eta$, $D^0 \rightarrow K^+K^-\eta$, and $D^0 \rightarrow K^-\pi^+\eta$ decays are reconstructed by combining π^\pm and K^\pm tracks with η candidates. A vertex fit is performed with the two charged tracks to obtain the D^0 decay vertex position; the resulting fit quality is labeled χ_v^2 . To improve the momentum resolution of the η , the γ daughters are subjected to a fit in which the photons are required to originate from the D^0 vertex position, and the invariant mass is constrained to be that of the η meson [5]. The fit quality of this mass constraint (χ_m^2) is required to satisfy $\chi_m^2 < 8$, and the resulting η momentum is required to be greater than $0.70 \text{ GeV}/c$. For $D^0 \rightarrow \pi^+\pi^-\eta$ candidates, we veto events in which $|M(\pi^+\pi^-) - m_{K_s^0}| < 10 \text{ MeV}/c^2$, where $m_{K_s^0}$ is the nominal K_s^0 mass [5], to suppress background from CF $D^0 \rightarrow K_s^0\eta$ decays. This veto range corresponds to about 3σ in resolution. The D^0 invariant mass M is required to satisfy $1.850 \text{ GeV}/c^2 < M < 1.878 \text{ GeV}/c^2$ for $D^0 \rightarrow K^+K^-\eta$ candidates; $1.840 \text{ GeV}/c^2 < M < 1.884 \text{ GeV}/c^2$ for $D^0 \rightarrow \pi^+\pi^-\eta$ candidates; and $1.842 \text{ GeV}/c^2 < M < 1.882 \text{ GeV}/c^2$ for $D^0 \rightarrow K^-\pi^+\eta$ candidates. These ranges correspond to about 2σ in resolution.

Candidate $D^{*+} \rightarrow D^0\pi_s^+$ decays are reconstructed by combining D^0 candidates with π_s^+ tracks. We first fit for D^{*+} decay vertex using the D^0 momentum vector and decay vertex position, and the IP as a constraint (i.e., the D^{*+} nominally originates from the IP). The resulting goodness-of-fit is labeled χ_{IP}^2 . To improve the resolution in Q , another vertex fit is performed: in this case we constrain the π_s^+ daughter to originate from the D^{*+} decay vertex, and the resulting fit quality is labeled χ_s^2 . The sum of the above three fit qualities, $\sum \chi_{\text{vtx}}^2 = \chi_v^2 + \chi_{\text{IP}}^2 + \chi_s^2$, is required to be less than 50; this requirement has a signal efficiency of about 97%. Those D^{*+} candidates satisfying $0 < Q < 15 \text{ MeV}$ are retained for further analysis. To eliminate D^{*+} candidates originating from B decays, and to also suppress combinatorial background, the D^{*+} momentum in the CM frame is required to be greater than $2.70 \text{ GeV}/c$.

After the above selection criteria are applied, about 2.1% of $D^0 \rightarrow \pi^+\pi^-\eta$ events, 1.3% of $D^0 \rightarrow K^-\pi^+\eta$ events, and $< 0.1\%$ of $D^0 \rightarrow K^+K^-\eta$ events have two or more D^{*+} candidates. For such multi-candidate events, we choose a single candidate: that which has the smallest value of the sum $\sum \chi_{\text{vtx}}^2 + \chi_m^2(\eta)$. This criterion, according to MC simulation, identifies the correct candidate 54% of the time.

Region	Component	$D^0 \rightarrow K^- \pi^+ \eta$	$D^0 \rightarrow \pi^+ \pi^- \eta$	$D^0 \rightarrow K^+ K^- \eta$
Fitted region	signal	180369 ± 837	12982 ± 198	1482 ± 60
	background	57752 ± 761	101011 ± 357	5681 ± 88
Signal region	signal	162456 ± 754	12053 ± 184	1343 ± 54
	background	7578 ± 100	11274 ± 40	678 ± 11

Table 1. Yields of signal and background events in the fitted region $0 < Q < 15$ MeV, and in the signal region $|Q - 5.86| < 0.80$ MeV.

4 Measurement of the branching fractions

4.1 Measurement of $\mathcal{B}(D^0 \rightarrow \pi^+ \pi^- \eta)$ and $\mathcal{B}(D^0 \rightarrow K^+ K^- \eta)$

We extract the signal yield via an unbinned maximum-likelihood fit to the Q distribution. The probability density function (PDF) used for signal events is taken to be the sum of a bifurcated Student's t-function (S_{bif}), which is defined in appendix A, and one or two asymmetric Gaussians (G_{asym}), with all having a common mean. The PDF used for $D^0 \rightarrow K^+ K^- \eta$ signal events is simply an S_{bif} function. These PDFs are explicitly

$$\mathcal{P}_{\text{sig}}^{K\pi\eta} = f_1[f_s S_{\text{bif}}(\mu, \sigma_0, \delta_0, n_l, n_h) + (1 - f_s)G_{\text{asym}}(\mu, r_1\sigma_0, \delta_1)] + (1 - f_1)G_{\text{asym}}(\mu, r_2r_1\sigma_0, \delta_2), \quad (4.1)$$

$$\mathcal{P}_{\text{sig}}^{\pi\pi\eta} = f_s S_{\text{bif}}(\mu, \sigma_0, \delta_0, n_l, n_h) + (1 - f_s)G_{\text{asym}}(\mu, r_1\sigma_0, \delta_1), \quad (4.2)$$

$$\mathcal{P}_{\text{sig}}^{KK\eta} = S_{\text{bif}}(\mu, \sigma_0, \delta_0, n_l, n_h). \quad (4.3)$$

In these expressions, δ_i is an asymmetry parameter characterizing the difference between left-side and right-side widths: $\sigma_{R,L} = \sigma(1 \pm \delta)$. Most of these parameters are fixed to values obtained from MC simulation. However, the parameters μ , σ_0 , and, for the higher-statistics $D^0 \rightarrow K^- \pi^+ \eta$ channel, $n_{l,h}$, are floated to account for possible differences in resolution between data and MC. For backgrounds, the PDF is taken to be a threshold function $f(Q) = Q^\alpha e^{-\beta Q}$; for the CF mode $D^0 \rightarrow K^- \pi^+ \eta$, we include an additional symmetric Gaussian to describe a small background component originating from misreconstructed $D^0 \rightarrow K^- \pi^+ \pi^0 \pi^0$ decays. The parameters of this Gaussian are fixed to values obtained from MC simulation, while all other parameters are floated. No other peaking backgrounds, such as misreconstructed D^0 decays or signal decays in which a pion from the D^0 is swapped with that from the D^{*+} decay, are found in the MC simulation.

The results of the fit are shown in figure 1, along with the pull $(N_{\text{data}} - N_{\text{fit}})/\sigma$, where σ is the error on N_{data} . All fit residuals look satisfactory. The signal yields in the fitted region $0 < Q < 15$ MeV, and in the signal region $|Q - 5.86| < 0.80$ MeV, are listed in table 1.

To measure the branching fraction, we must divide these signal yields by their reconstruction efficiencies. However, the reconstruction efficiency for a decay can vary across the Dalitz plot of three-body phase space, and the Dalitz-plot distribution of $D^0 \rightarrow \pi^+ \pi^- \eta$ and $D^0 \rightarrow K^+ K^- \eta$ decays has not been previously measured. Thus, to avoid systematic uncertainty due to the unknown Dalitz distribution (or decay model), we correct our signal yields for reconstruction efficiencies as follows. We divide the Dalitz plot of the data into

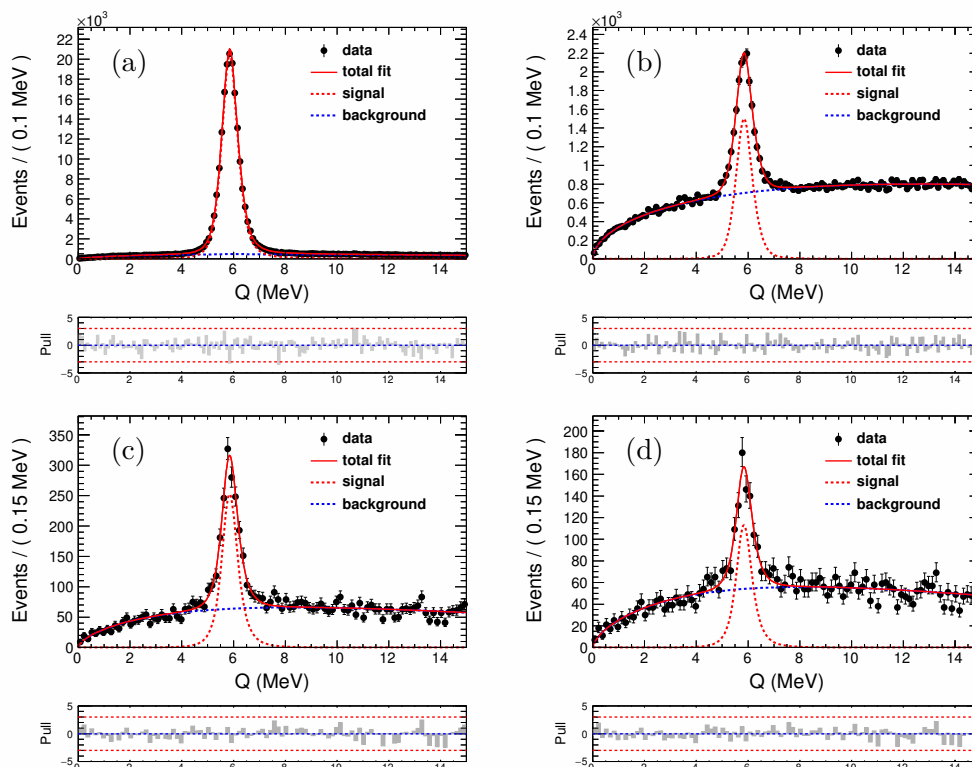


Figure 1. Distributions of the released energy Q in $D^{*+} \rightarrow D^0 \pi_s^+$ decay for (a) $D^0 \rightarrow K^- \pi^+ \eta$, (b) $D^0 \rightarrow \pi^+ \pi^- \eta$, (c) $D^0 \rightarrow K^+ K^- \eta$, and (d) $D^0 \rightarrow K^+ K^- \eta$ with the ϕ -peak excluded by requiring $|M_{KK} - m_\phi| > 20 \text{ MeV}/c^2$. Points with error bars show the data; the dashed red curve shows the signal; the dashed blue curve shows the background; and the solid red curves show the overall fit result. The pull plots underneath the fit results show the residuals divided by the errors in the histogram.

bins of $M^2(h^+ h^-)$ and $M^2(h^- \eta)$, where $h = \pi$ or K , determine the reconstruction efficiency independently for each bin, and calculate the corrected signal yield via the formula

$$N^{\text{cor}} = \sum_i \frac{N_i^{\text{tot}} - N_i^{\text{bkg}} f_i^{\text{bkg}}}{\varepsilon_i}, \quad (4.4)$$

where i runs over all bins. The values of $M^2(h^+ h^-)$ and $M^2(h^- \eta)$ are calculated subject to the constraint $M(h^+ h^- \eta) = m_{D^0}$. The formula (4.4) has the following terms:

- ε_i is the signal reconstruction efficiency for bin i , as determined from a large sample of MC events. The resolutions in Q of the MC samples are adjusted to match those of the data. The efficiencies for $D^0 \rightarrow \pi^+ \pi^- \eta$ are plotted in figure 2(a), and those for $D^0 \rightarrow K^+ K^- \eta$ are plotted in figure 3(a). These efficiencies include a small ($\sim 2\%$) correction for K^\pm and π^\pm particle identification (PID) efficiencies, to account for small differences observed between data and MC simulation. This correction is determined using a sample of $D^{*+} \rightarrow [D^0 \rightarrow K^- \pi^+] \pi_s^+$ decays.

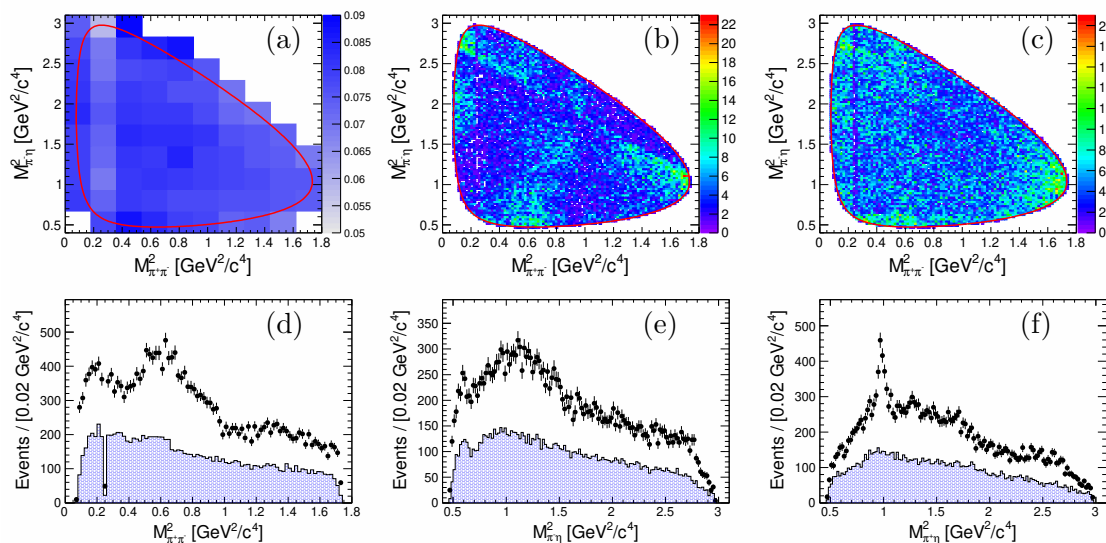


Figure 2. For $D^0 \rightarrow \pi^+\pi^-\eta$: (a) distribution of reconstruction efficiencies over the Dalitz plot, divided into 10×10 bins of $M_{\pi^+\pi^-}^2$ vs $M_{\pi^-\eta}^2$. The red lines indicate the Dalitz plot boundaries. (b) Dalitz plot for events in the Q signal region $|Q - 5.86| < 0.80$ MeV. (c) Dalitz plot for events in the sideband region $2.5 < |Q - 5.86| < 4.90$ MeV, used to estimate the background shape. (d, e, f) Projections of Dalitz variables $M_{\pi^+\pi^-}^2$, $M_{\pi^-\eta}^2$, and $M_{\pi^+\eta}^2$, respectively. Points with error bars show events in the signal region; blue-filled histograms show the estimated background (see text). The dip in $M_{\pi^+\pi^-}^2$ near $0.25 \text{ GeV}^2/c^4$ is due to the K_s^0 veto.

- N_i^{tot} is the number of events in the Q signal region and the i^{th} bin of the Dalitz plot. These yields are plotted in figure 2(b) for $D^0 \rightarrow \pi^+\pi^-\eta$ and in figure 3(b) for $D^0 \rightarrow K^+K^-\eta$.
- N^{bkg} is the total background yield in the Q signal region, as obtained from fitting the Q distribution (see figure 1).
- f_i^{bkg} is the fraction of background in the i^{th} -bin, with $\sum_i f_i = 1$. These fractions are obtained from the Dalitz plot distribution of events in the Q sideband region $2.5 \text{ MeV} < |Q - 5.86| < 4.9 \text{ MeV}$. The distribution of sideband events is shown in figure 2(c) for $D^0 \rightarrow \pi^+\pi^-\eta$ and in figure 3(c) for $D^0 \rightarrow K^+K^-\eta$.

There are $10 \times 10 = 100$ bins in total for $D^0 \rightarrow \pi^+\pi^-\eta$, and $5 \times 5 = 25$ bins total for $D^0 \rightarrow K^+K^-\eta$. The final corrected yields obtained using eq. (4.4) are $N^{\text{cor}} = (1.536_{-0.020}^{+0.021}) \times 10^5$ for $D^0 \rightarrow \pi^+\pi^-\eta$, and $N^{\text{cor}} = (2.263_{-0.077}^{+0.084}) \times 10^4$ for $D^0 \rightarrow K^+K^-\eta$.

The branching fraction of a signal mode relative to that of the normalization mode is determined from the ratio of their respective efficiency-corrected yields:

$$\frac{\mathcal{B}(D^0 \rightarrow h^+h^-\eta)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = \frac{N^{\text{cor}}(D^0 \rightarrow h^+h^-\eta)}{N^{\text{cor}}(D^0 \rightarrow K^-\pi^+\eta)}, \quad (4.5)$$

where $h = K$ or π . The efficiency-corrected yield for the normalization channel $D^0 \rightarrow K^-\pi^+\eta$ is evaluated in a different manner than those of the signal modes. As the Dalitz

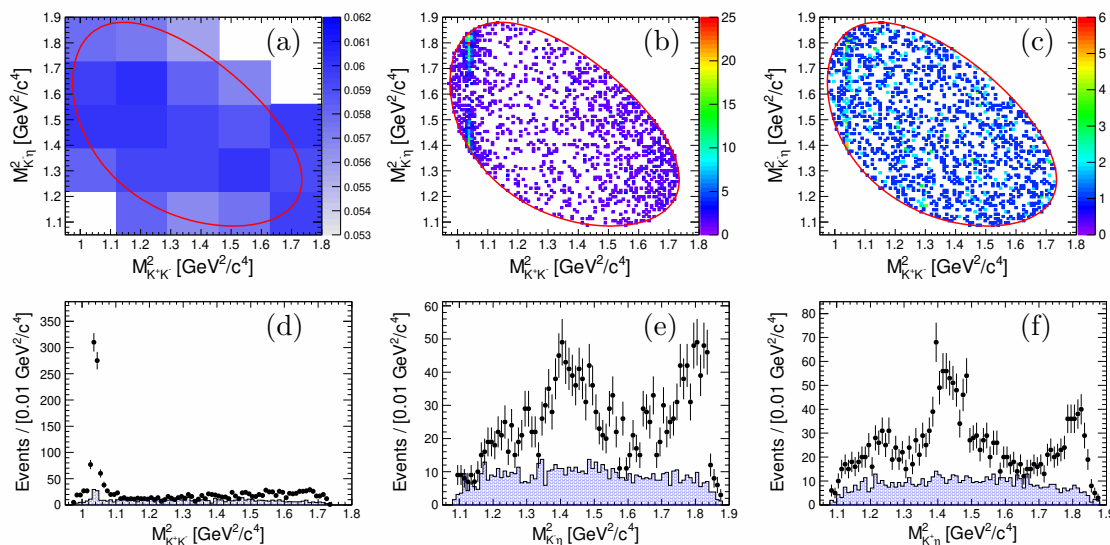


Figure 3. For $D^0 \rightarrow K^+K^-\eta$: (a) distribution of reconstruction efficiencies over the Dalitz plot, divided into 5×5 bins of $M_{K^+K^-}^2$ vs $M_{K^-\eta}^2$. The red lines indicate the Dalitz plot boundaries. (b) Dalitz plot for events in the Q signal region $|Q - 5.86| < 0.80$ MeV. (c) Dalitz plot for events in the sideband region $2.5 < |Q - 5.86| < 4.90$ MeV, used to estimate the background shape. (d, e, f) Projections of Dalitz variables $M_{K^+K^-}^2$, $M_{K^-\eta}^2$, and $M_{K^+\eta}^2$, respectively. Points with error bars show events in the signal region; blue-filled histograms show the estimated background (see text).

plot of $D^0 \rightarrow K^-\pi^+\eta$ decays has been measured with high statistics [6], we use the resulting decay model to generate an MC sample, and use that sample to evaluate the overall reconstruction efficiency. The result, including the small PID efficiency correction, is $\varepsilon_{K\pi\eta} = (6.870 \pm 0.014)\%$. Dividing the fitted yield for $D^0 \rightarrow K^-\pi^+\eta$ (see table 1) by this value gives $N^{\text{cor}}(D^0 \rightarrow K^-\pi^+\eta) = (2.365 \pm 0.011) \times 10^6$.

Inserting all efficiency-corrected yields into eq. (4.5) gives the ratios of branching fractions

$$\frac{\mathcal{B}(D^0 \rightarrow \pi^+\pi^-\eta)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [6.49 \pm 0.09 (\text{stat}) \pm 0.12 (\text{syst})] \times 10^{-2}, \quad (4.6)$$

$$\frac{\mathcal{B}(D^0 \rightarrow K^+K^-\eta)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [9.57^{+0.36}_{-0.33} (\text{stat}) \pm 0.20 (\text{syst})] \times 10^{-3}. \quad (4.7)$$

The second error listed is the systematic uncertainty, which is evaluated below (section 4.3). Multiplying both sides of eqs. (4.6) and (4.7) by the world average value $\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta) = (1.88 \pm 0.05)\%$ [5] gives

$$\mathcal{B}(D^0 \rightarrow \pi^+\pi^-\eta) = [1.22 \pm 0.02 (\text{stat}) \pm 0.02 (\text{syst}) \pm 0.03 (\mathcal{B}_{\text{ref}})] \times 10^{-3}, \quad (4.8)$$

$$\mathcal{B}(D^0 \rightarrow K^+K^-\eta) = [1.80^{+0.07}_{-0.06} (\text{stat}) \pm 0.04 (\text{syst}) \pm 0.05 (\mathcal{B}_{\text{ref}})] \times 10^{-4}, \quad (4.9)$$

where the third uncertainty listed is due to the branching fraction for the reference mode $D^0 \rightarrow K^-\pi^+\eta$. The result (4.8) is consistent with the world average value $(1.17 \pm 0.07) \times 10^{-3}$ [5] but has improved precision. The result (4.9) is the first such measurement.

The Dalitz plots and projections are shown in figure 2 for $D^0 \rightarrow \pi^+\pi^-\eta$ and in figure 3 for $D^0 \rightarrow K^+K^-\eta$. The background plotted is taken from the Q sideband region, with the entries scaled to match the background yield in the signal region obtained from the Q fit (figure 1). Several intermediate structures are clearly visible. For $D^0 \rightarrow \pi^+\pi^-\eta$ events, the $M_{\pi^+\pi^-}^2$ projection in figure 2(d) shows the $D^0 \rightarrow \rho^0(770)\eta$, $\rho(770) \rightarrow \pi^+\pi^-$ decay process to be dominant. The $M^2(\pi^+\eta)$ distribution in figure 2(f) shows a sharp peak near $1.0 \text{ GeV}^2/c^4$, which indicates $D^0 \rightarrow a_0(980)^+\pi^-$, $a_0(980)^+ \rightarrow \pi^+\eta$ decay. In contrast, the $M^2(\pi^-\eta)$ distribution in figure 2(e) shows no indication of $D^0 \rightarrow a_0(980)^-\pi^+$, $a_0(980)^- \rightarrow \pi^-\eta$. This is unexpected, as the branching fraction for $D^0 \rightarrow a_0(980)^-\pi^+$ is predicted to be two orders of magnitude larger than that for $D^0 \rightarrow a_0(980)^+\pi^-$ [21].

For $D^0 \rightarrow K^+K^-\eta$ events, the $M_{K^+K^-}^2$ distribution shows the $D^0 \rightarrow \phi\eta$, $\phi \rightarrow K^+K^-$ decay process to be dominant. However, a non- ϕ contribution is also visible. We thus measure $\mathcal{B}(D^0 \rightarrow K^+K^-\eta)_{\phi\text{-excluded}}$ by requiring $|M_{K^+K^-} - m_\phi| > 20 \text{ MeV}/c^2$. The signal yield is obtained as before by fitting the Q distribution. The result is 599 ± 45 events in the signal region, as shown in figure 1(d). The change in likelihood, with and without including a signal component in such Q fitting, is $\Delta \ln L = 214$. As the number of degrees of freedom for the fit with no signal component is three less than that for the nominal fit (parameters N_{sig} , μ , and σ_0 are dropped), this value of $\Delta \ln L$ corresponds to a statistical significance for the signal of 20σ .

We divide the signal yields obtained for bins of the Dalitz plot by the efficiencies for these bins [see eq. (4.4)] to obtain $N^{\text{cor}} = 12443_{-893}^{+1071}$ (for $D^0 \rightarrow K^+K^-\eta$ with the ϕ excluded). Thus

$$\frac{\mathcal{B}(D^0 \rightarrow K^+K^-\eta)_{\phi\text{-excluded}}}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [5.26_{-0.38}^{+0.45} (\text{stat}) \pm 0.11 (\text{syst})] \times 10^{-3}. \quad (4.10)$$

The second error listed is the systematic uncertainty, which is evaluated below (section 4.3). Multiplying each side of eq. (4.10) by $\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta) = (1.88 \pm 0.05)\%$ [5] gives the branching fraction for $D^0 \rightarrow K^+K^-\eta$ with the ϕ component excluded by requiring $|M_{K^+K^-} - m_\phi| > 20 \text{ MeV}/c^2$:

$$\mathcal{B}(D^0 \rightarrow K^+K^-\eta)_{\phi\text{-excluded}} = [0.99_{-0.07}^{+0.08} (\text{stat}) \pm 0.02 (\text{syst}) \pm 0.03 (\mathcal{B}_{\text{ref}})] \times 10^{-4}. \quad (4.11)$$

This result is somewhat higher (but more precise) than a similar measurement by BESIII, $(0.59 \pm 0.19) \times 10^{-4}$ [7].

4.2 Measurement of $\mathcal{B}(D^0 \rightarrow \phi\eta)$

As shown in figure 3, the decay $D^0 \rightarrow K^+K^-\eta$ is dominated by the Cabibbo- and color-suppressed decay $D^0 \rightarrow \phi\eta$, $\phi \rightarrow K^+K^-$. We thus measure the branching fraction for $D^0 \rightarrow \phi\eta$ by performing a two-dimensional fit to the M_{KK} and Q distributions of $D^0 \rightarrow K^+K^-\eta$ events. The fitted region is $M_{KK} < 1.08 \text{ GeV}/c^2$ and $Q < 15 \text{ MeV}$. In this region, signal decays and background are straightforward to identify: the non- ϕ $D^0 \rightarrow K^+K^-\eta$ component peaks in Q but not in M_{KK} , whereas combinatorial background containing $\phi \rightarrow K^+K^-$ decays peak in M_{KK} but not in Q . The signal PDF for M_{KK} is taken to be

the sum of a Gaussian and two asymmetric Gaussians, with a common mean for the $M(\phi)$ peak. The PDF for Q is taken to be a bifurcated Student's t-function:

$$\begin{aligned} \mathcal{P}_{\text{sig}}(M_{KK}) &= f_2[f_1 G(\mu_m, \sigma_{m0}) + (1 - f_1)G_{\text{asym}}(\mu_m, r_1\sigma_{m0}, \delta_2)] \\ &\quad + (1 - f_2)G_{\text{asym}}(\mu_m, r_2r_1\sigma_{m0}, \delta_3), \end{aligned} \quad (4.12)$$

$$\mathcal{P}_{\text{sig}}(Q) = S_{\text{bif}}(\mu_q, \sigma_0, \delta_0, n_l, n_h), \quad (4.13)$$

$$\mathcal{P}_{\text{sig}}^{D^0 \rightarrow \phi\eta}(M_{KK}, Q) = \mathcal{P}_{\text{sig}}(M_{KK}) \cdot \mathcal{P}_{\text{sig}}(Q). \quad (4.14)$$

Most of the signal shape parameters are fixed to MC values; the parameters μ_m, σ_{m0} for M_{KK} and μ_q, σ_0 for Q are floated to account for possible differences in resolution between data and MC simulation.

The non- ϕ $D^0 \rightarrow K^+K^-\eta$ component includes several processes such as $D^0 \rightarrow a_0(980)\eta$, $D^0 \rightarrow K_0^*(1430)^\pm K^\mp$, and non-resonant $D^0 \rightarrow K^+K^-\eta$ decays. We parameterize this component in M_{KK} by a threshold function, and in Q by the same PDF as that used for the signal:

$$\mathcal{P}_{\text{peak}}(M_{KK}) = \sqrt{M_{KK} - m_0} \cdot e^{-\beta_{\text{peak}}(M_{KK} - m_0)}, \quad (4.15)$$

$$\mathcal{P}_{\text{peak}}(Q) = \mathcal{P}_{\text{sig}}(Q), \quad (4.16)$$

$$\mathcal{P}_{\text{peak}}(M_{KK}, Q) = \mathcal{P}_{\text{peak}}(M_{KK}) \cdot \mathcal{P}_{\text{peak}}(Q), \quad (4.17)$$

where the threshold value $m_0 = 2m_K = 0.987354 \text{ GeV}/c^2$. We consider possible interference between the non- ϕ component and the $D^0 \rightarrow \phi\eta$ signal as a systematic uncertainty.

For combinatorial background, the Q distribution is parameterized with a threshold function. The M_{KK} distribution has two parts: (1) a ϕ -peak, which is taken to be the same as that of signal decay, and (2) a threshold function. These PDFs take the forms

$$\mathcal{P}_{\text{bkg}}(M_{KK}) = f_\phi \mathcal{P}_{\text{sig}}(M_{KK}) + (1 - f_\phi)(M_{KK} - m_0)^{\alpha_m} e^{-\beta_m(M_{KK} - m_0)}, \quad (4.18)$$

$$\mathcal{P}_{\text{bkg}}(Q) = Q^{\alpha_q} e^{-\beta_q Q}, \quad (4.19)$$

$$\mathcal{P}_{\text{bkg}}(M_{KK}, Q) = \mathcal{P}_{\text{bkg}}(M_{KK}) \cdot \mathcal{P}_{\text{bkg}}(Q). \quad (4.20)$$

The relative fraction f_ϕ and the Q threshold parameters are floated; all other parameters are fixed to MC values.

The results of the two-dimensional likelihood fit are shown in figure 4. We obtain a signal yield $N_{\text{sig}} = 728 \pm 36$ in the full fitted region, and $N_{\text{sig}} = 600 \pm 29$ in the signal region $|Q - 5.86| < 0.8 \text{ MeV}$ and $|M_{K^+K^-} - m_\phi| < 10 \text{ MeV}/c^2$. The difference in likelihood, with and without including a signal component, is $\Delta \ln \mathcal{L} = 464.8$. As the number of degrees of freedom for the fit with no signal component is one less than that for the nominal fit (parameter N_{sig} is dropped), this value of $\Delta \ln L$ corresponds to a statistical significance for $D^0 \rightarrow \phi\eta$ of 31σ .

We evaluate the signal reconstruction efficiency using a large MC sample of $D^0 \rightarrow \phi\eta$ decays. We obtain, for events in the M_{KK} - Q signal region, an efficiency $\varepsilon = (5.262 \pm 0.021)\%$. Thus, $N^{\text{cor}}(D^0 \rightarrow \phi\eta, \phi \rightarrow K^+K^-) = (1.140 \pm 0.055) \times 10^4$, and

$$\frac{\mathcal{B}(D^0 \rightarrow \phi\eta, \phi \rightarrow K^+K^-)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [4.82 \pm 0.23 (\text{stat}) \pm 0.16 (\text{syst})] \times 10^{-3}. \quad (4.21)$$

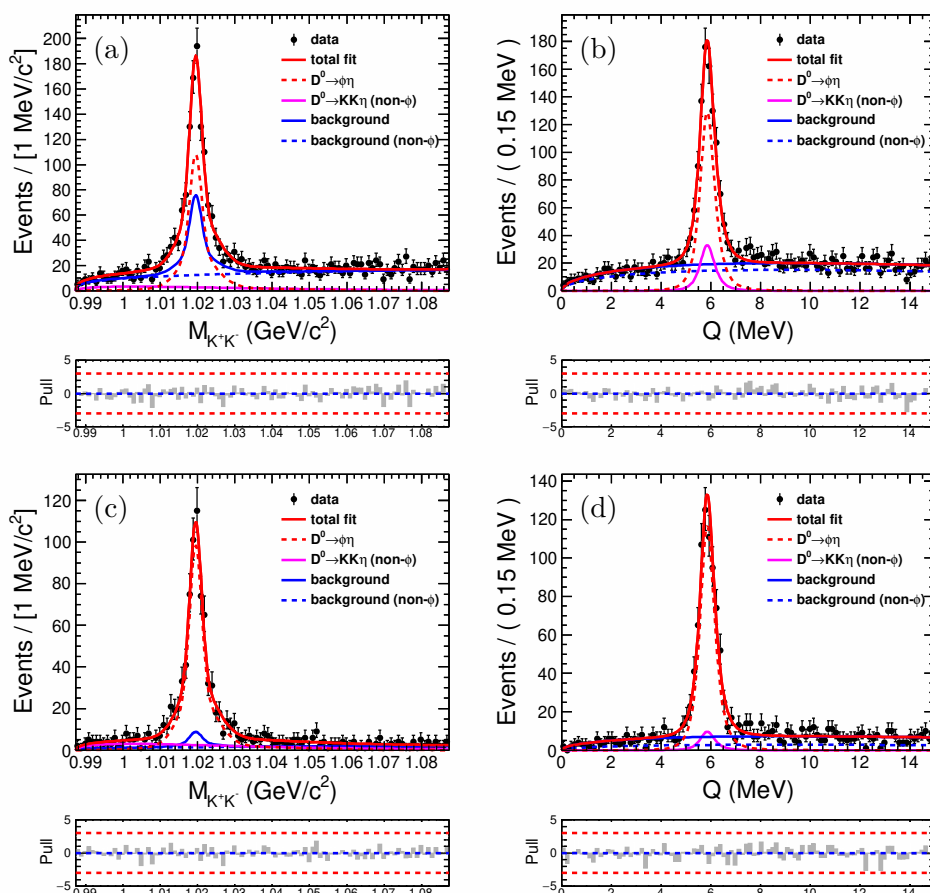


Figure 4. Projections of K^+K^- invariant mass distributions in Q (a) fit region and (c) signal region ($|Q - 5.86| < 0.8 \text{ MeV}$) and Q distributions in $M_{K^+K^-}$ (b) fit region and (d) signal region ($|M_{K^+K^-} - m_\phi| < 10 \text{ MeV}/c^2$) from the $M_{K^+K^-}-Q$ two-dimensional fit. Points with error bars are the data. The red solid line is total fit result. The red dashed curves are signal of $D^0 \rightarrow \phi\eta$, and the magenta solid curves show the Q -peaking background from non- ϕ component in $D^0 \rightarrow K^+K^-\eta$. The blue dash line is the non- ϕ component of total non- Q -peaking background (blue line).

The second error listed is the systematic uncertainty, which is evaluated below (section 4.3). Multiplying each side by the world average value $\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta) = (1.88 \pm 0.05)\%$ [5] and dividing by $\mathcal{B}(\phi \rightarrow K^+K^-) = (49.2 \pm 0.5)\%$ [5], we obtain

$$\mathcal{B}(D^0 \rightarrow \phi\eta) = [1.84 \pm 0.09 (\text{stat}) \pm 0.06 (\text{syst}) \pm 0.05 (\mathcal{B}_{\text{ref}})] \times 10^{-4}, \quad (4.22)$$

where the systematic uncertainty includes the small uncertainty on $\mathcal{B}(\phi \rightarrow K^+K^-)$. This result is consistent with, but notably more precise than, the current world average of $(1.8 \pm 0.5) \times 10^{-4}$ [5]. It is also consistent with theoretical predictions [22, 23]. As a consistency check, we calculate the branching fraction of the non- ϕ $D^0 \rightarrow K^+K^-\eta$ component by subtracting the $D^0 \rightarrow \phi\eta$ branching fraction from the total $D^0 \rightarrow K^+K^-\eta$ result:

$$\mathcal{B}(D^0 \rightarrow K^+K^-\eta) - \mathcal{B}(D^0 \rightarrow \phi\eta, \phi \rightarrow K^+K^-) = (0.90 \pm 0.08) \times 10^{-4}, \quad (4.23)$$

Systematic sources	$\frac{\mathcal{B}(D^0 \rightarrow \pi^+ \pi^- \eta)}{\mathcal{B}(D^0 \rightarrow K^- \pi^+ \eta)}$	$\frac{\mathcal{B}(D^0 \rightarrow K^+ K^- \eta)}{\mathcal{B}(D^0 \rightarrow K^- \pi^+ \eta)}$	$\frac{\mathcal{B}(D^0 \rightarrow (\phi \rightarrow K^+ K^-) \eta)}{\mathcal{B}(D^0 \rightarrow K^- \pi^+ \eta)}$
PID efficiency correction	1.8%	1.9%	1.9%
Signal PDF	0.3%	0.5%	0.9%
Background PDF	0.0%	0.0%	0.1%
Mass resolution calibration	0.1%	0.3%	0.0%
Yield correction with efficiency map	0.3%	0.7%	–
MC statistics	0.3%	0.4%	0.4%
K_S^0 veto	0.1%	–	–
Interference in M_{KK}	–	–	2.5%
Total syst. error	1.9%	2.1%	3.3%

Table 2. Systematic uncertainties (fractional) for the branching ratio measurements.

which is very close to our measurement of $\mathcal{B}(D^0 \rightarrow K^+ K^- \eta)_{\phi\text{-excluded}}$ in eq. (4.11).

4.3 Systematic uncertainties

The sources of systematic uncertainty in measuring the branching fractions are listed in table 2. These uncertainties are evaluated as follows.

- A correction for PID efficiency is applied to K^\pm and π^\pm tracks, to account for a difference in efficiency between data and MC simulation. The correction depends on track momentum and is small; the uncertainty on the correction is even smaller, in the range (0.90–0.97)%. When evaluating this uncertainty for a ratio of branching fractions, we conservatively assume the efficiency corrections for K^+ and π^+ tracks (which appear separately in numerator and denominator, or vice-versa) are anticorrelated.
- The uncertainty due to the parameters fixed in the fit for the signal yield is evaluated as follows. We sample these parameters simultaneously from Gaussian distributions, accounting for their correlations, and re-fit for the signal yield. The procedure is repeated 1000 times and these yields are plotted. The ratio of the root-mean-square (RMS) to the mean value of the resulting distribution of signal yields is taken as the systematic uncertainty due to the fixed parameters. The Gaussian distributions from which the parameters are sampled have mean values equal to the fixed values of the parameters, and widths equal to their respective uncertainties.
- For background PDFs, all parameters are floated in the fits except for those describing the amount of background and its shape for $D^0 \rightarrow K^- \pi^+ \eta$, which are taken from MC simulation. We evaluate this uncertainty due to these fixed parameters as done above, by simultaneously sampling these parameters from Gaussian distributions having mean values equal to the fixed values and widths equal to their respective uncertainties. The RMS of the resulting distribution of $D^0 \rightarrow K^- \pi^+ \eta$ yields is taken as the systematic uncertainty due to the peaking background.

- We correct for differences in mass resolutions (including M , Q and M_{KK}) between data and MC when calculating reconstruction efficiencies (in eq. (4.4) for $D^0 \rightarrow \pi^+\pi^-\eta$ and $D^0 \rightarrow K^+K^-\eta$, as well as for $D^0 \rightarrow \phi\eta$ and the reference mode). We take the systematic uncertainty of this procedure to be the difference in the ratio of efficiency-corrected signal yields to that of the reference mode obtained both with and without this resolution correction.
- The efficiency for $D^0 \rightarrow \pi^+\pi^-\eta$ and $D^0 \rightarrow K^+K^-\eta$ decays is evaluated in bins of the Dalitz plot; see eq. (4.4). This efficiency has uncertainty arising from the number of bins used, from the efficiency values ε_i for the various bins, and from the bin-by-bin background subtraction.
 - For the first uncertainty, we vary the numbers of bins used, and the corresponding change in the efficiency is taken as the systematic uncertainty. For $D^0 \rightarrow \pi^+\pi^-\eta$ decays, our nominal result uses 10×10 bins; thus we also try 8×8 and 12×12 bins. For $D^0 \rightarrow K^+K^-\eta$ decays, our nominal result uses 5×5 bins; thus we also try 3×3 and 7×7 bins. We obtain 0.25% for $D^0 \rightarrow \pi^+\pi^-\eta$, and 0.50% for $D^0 \rightarrow K^+K^-\eta$.
 - to evaluate the effect of uncertainties in ε_i , we sample the ε_i from Gaussian distributions having mean values equal to the nominal values, and widths equal to their uncertainties. For each sampling, we re-calculate the yield N^{cor} using eq. (4.4) and plot the result. The RMS of this distribution is taken as the systematic uncertainty: 0.21% for $D^0 \rightarrow \pi^+\pi^-\eta$, and 0.43% for $D^0 \rightarrow K^+K^-\eta$.
 - the background subtraction procedure depends on the distribution of background over the Dalitz plot. We take this distribution from a data Q sideband region. To evaluate the uncertainty due to this Dalitz distribution, we shift the Q sideband region used by ± 0.4 MeV and repeat the procedure. The change is assigned as the systematic uncertainty: 0.02% for $D^0 \rightarrow \pi^+\pi^-\eta$, and 0.03% for $D^0 \rightarrow K^+K^-\eta$.
- The efficiency for $D^0 \rightarrow K^-\pi^+\eta$ is evaluated from MC simulation using a Dalitz decay model. The uncertainty due to this model is evaluated by varying the model and re-calculating ε . Specifically, our nominal model uses eight intermediate resonances as measured in ref. [6]; as an alternative, we include all thirteen intermediate resonances listed in ref. [6]. The resulting change in our reconstruction efficiency is very small ($\delta\varepsilon/\varepsilon < 0.01\%$). As a cross check, we calculate N^{cor} for $D^0 \rightarrow K^-\pi^+\eta$ using eq. (4.4); the result is $(2.369 \pm 0.007) \times 10^6$, which is almost identical with our nominal result (the difference is much smaller than the uncertainty).
- There are small uncertainties in the reconstruction efficiencies ε , which are evaluated from MC simulation, due to the finite statistics of the MC samples used.
- There is an uncertainty arising from the K_S^0 veto required for $D^0 \rightarrow \pi^+\pi^-\eta$ decays. We evaluate this by changing the veto region from $m_{K_S^0} \pm 10 \text{ MeV}/c^2$ ($\sim 3\sigma$) to

$m_{K_S^0} \pm 15 \text{ MeV}/c^2$ ($\sim 5\sigma$); the resulting change in the $D^0 \rightarrow \pi^+\pi^-\eta$ signal yield is taken as the systematic uncertainty.

- We consider possible interference between the $D^0 \rightarrow \phi\eta$ amplitude and that of non- ϕ $D^0 \rightarrow K^+K^-\eta$. Such interference could alter the M_{KK} distribution used to fit for the $D^0 \rightarrow \phi\eta$ yield. We evaluate this effect by introducing a relative phase θ between the two amplitudes, which modifies the PDF to be

$$\mathcal{P}_{\text{total}}(M_{KK}, Q) = \left| A_\phi(M_{KK}) + r e^{i(\theta+k\cdot\pi)} \sqrt{F_{\text{non-}\phi}(M_{KK})} \right|^2 \times F_{\text{sig}}(Q). \quad (4.24)$$

In this expression, A_ϕ is a relativistic Breit-Wigner function, $F_{\text{non-}\phi}$ is the shape function used in the nominal fit to model non- ϕ decays, and k is a factor that adds a phase shift of π depending on whether the cosine of the K^+ helicity angle (θ_h) is positive or negative, i.e., $k = 0$ for $\cos \theta_h < 0$, and $k = 1$ otherwise [24]. The helicity angle θ_h is defined as the angle between the momentum of the K^+ and the η in the K^+K^- rest frame. After fitting with this PDF, we calculate the $D^0 \rightarrow \phi\eta$ yield as the product of the total yield obtained and the fraction f_ϕ given by

$$f_\phi = \int |A_\phi|^2 dM_{KK} / \int \left| A_\phi + r e^{i(\theta+k\cdot\pi)} \sqrt{F_{\text{non-}\phi}} \right|^2 dM_{KK}. \quad (4.25)$$

The result is that the $D^0 \rightarrow \phi\eta$ yield decreases by 2.5%, and thus we assign this value as the systematic uncertainty due to possible interference.

The total systematic uncertainty is obtained by adding in quadrature all the above contributions. The results are listed in table 2.

5 Measurement of CP asymmetries

5.1 Measurement of $A_{CP}(D^0 \rightarrow \pi^+\pi^-\eta)$ and $A_{CP}(D^0 \rightarrow K^+K^-\eta)$

To measure the CP asymmetries, we divide the sample for each channel into D^0 and \bar{D}^0 decays, where the flavor of the D^0 or \bar{D}^0 is tagged by the charge of the π_s^\pm from the $D^{*+} \rightarrow D^0\pi_s^+$ or $D^{*-} \rightarrow \bar{D}^0\pi_s^-$ decay. To correct for an asymmetry in π_s^\pm reconstruction efficiencies, we weight events according to the π_s^\pm efficiency mapping of ref. [25].

We simultaneously fit the Q distributions of these weighted samples for D^0 and \bar{D}^0 decays with parameters fixed in the same way as done for the branching fraction measurements. The parameters N_{sig} and A_{corr} are fitted, where the D^0 and \bar{D}^0 signal yields are given by $N_{\text{sig}}(D^0, \bar{D}^0) = (N_{\text{sig}}/2) \cdot (1 \pm A_{\text{corr}})$. The fit results are shown in figure 5. We obtain $N_{\text{sig}} = 12975 \pm 198$ and $A_{\text{corr}} = (1.44 \pm 1.24)\%$ for $D^0 \rightarrow \pi^+\pi^-\eta$ decays, and $N_{\text{sig}} = 1482 \pm 60$ and $A_{\text{corr}} = (-0.25 \pm 2.96)\%$ for $D^0 \rightarrow K^+K^-\eta$ decays.

These values for A_{corr} include the forward-backward asymmetry A_{FB} . We correct for A_{FB} by calculating A_{corr} in eight bins of $\cos \theta^*$, where θ^* is the polar angle of the D^{*+} with respect to the $+z$ axis in the e^+e^- CM frame. The bins used are $[-1.0, -0.6]$, $[-0.6, -0.4]$, $[-0.4, -0.2]$, $[-0.2, 0.0]$, $[0.0, 0.2]$, $[0.2, 0.4]$, $[0.4, 0.6]$, and $[0.6, 1.0]$. The

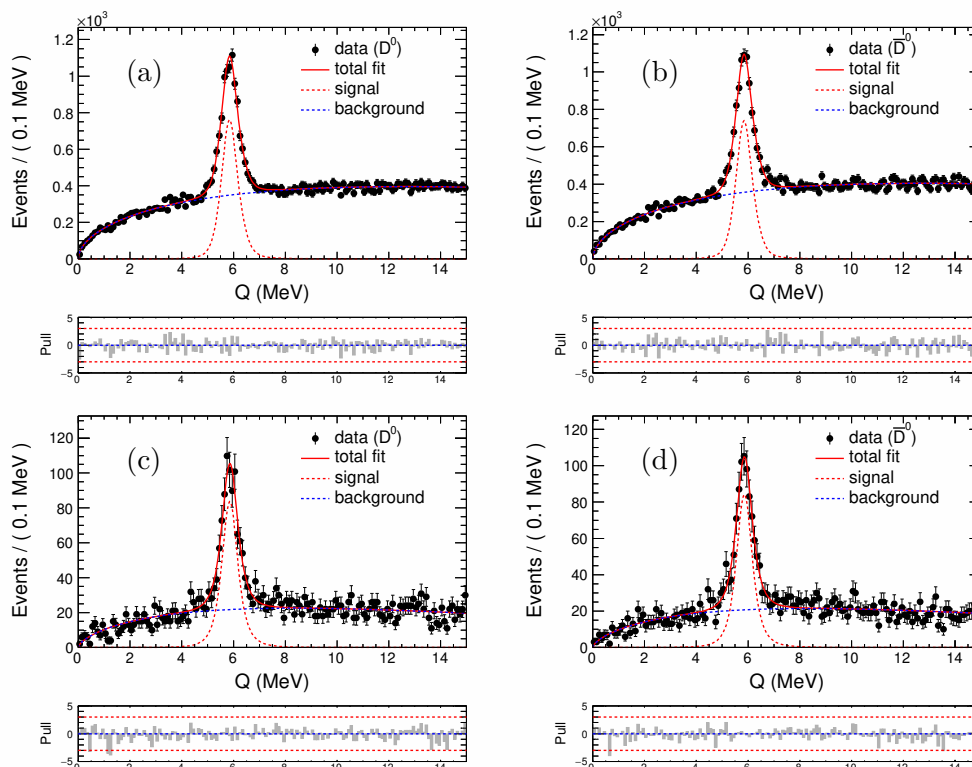


Figure 5. Simultaneous fit for $D^0 \rightarrow \pi^+\pi^-\eta$ (a) and $\bar{D}^0 \rightarrow \pi^+\pi^-\eta$ (b) candidates; and $D^0 \rightarrow K^+K^-\eta$ (c) and $\bar{D}^0 \rightarrow K^+K^-\eta$ (d) candidates.

asymmetries A_{CP} and A_{FB} are then extracted via eqs. (1.4) and (1.5). The resulting four values of A_{CP} and A_{FB} are plotted in figures 7(a, d) for $D^0 \rightarrow \pi^+\pi^-\eta$ and in figures 7(b, e) for $D^0 \rightarrow K^+K^-\eta$. Fitting the A_{CP} values to constants yields the final results

$$A_{CP}(D^0 \rightarrow \pi^+\pi^-\eta) = [0.9 \pm 1.2 (\text{stat}) \pm 0.5 (\text{syst})]\%, \quad (5.1)$$

$$A_{CP}(D^0 \rightarrow K^+K^-\eta) = [-1.4 \pm 3.3 (\text{stat}) \pm 1.1 (\text{syst})]\%. \quad (5.2)$$

The second error listed is the systematic uncertainty, which is evaluated below (section 5.3). The first result is a factor of four more precise than a recent measurement by BESIII [4], while the latter result is the first such measurement. The A_{FB} values plotted in figures 7(d-e) decrease with $\cos\theta^*$ and are consistent with (somewhat lower in $D^0 \rightarrow \pi^+\pi^-\eta$ than) the leading-order prediction [26] at $\sqrt{s} = 10.6 \text{ GeV}$ of $A_{FB}^{c\bar{c}} = -0.029 \cdot \cos\theta^*/(1 + \cos^2\theta^*)$, at the current level of statistics.

5.2 Measurement of $A_{CP}(D^0 \rightarrow \phi\eta)$

We repeat the above procedure to determine A_{CP} for $D^0 \rightarrow \phi\eta$ decays. Here, to determine parameters N_{sig} and A_{corr} , we perform a two-dimensional fit in $[M_{KK}, Q]$ for the D^0 and \bar{D}^0 samples simultaneously. We allow N_{sig} and A_{corr} to float separately for the $\phi\eta$ and non-resonant $K^+K^-\eta$ components. The projections of the fit result are shown

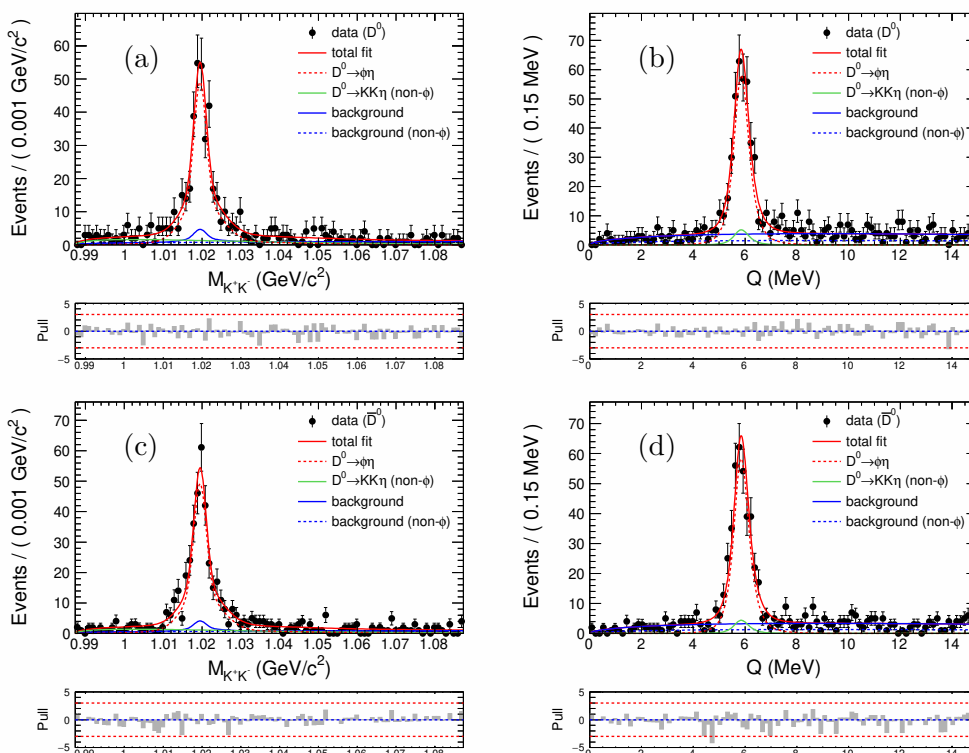


Figure 6. Projections in $M_{K^+K^-}$ (left) and Q (right) of the two-dimensional $(M_{K^+K^-}, Q)$ fit, for $D^0 \rightarrow \phi\eta$ (top) and $\bar{D}^0 \rightarrow \phi\eta$ (bottom). In both cases, the $M_{K^+K^-}$ projection corresponds to the Q signal region ($|Q - 5.86| < 0.8$ MeV), and the Q projection corresponds to the $M_{K^+K^-}$ signal region ($|M_{K^+K^-} - m_\phi| < 0.01$ GeV/ c^2).

in figure 6, and the results are $N_{\text{sig}} = 728 \pm 36$ and $A_{\text{corr}} = (-0.17 \pm 4.44)\%$. We perform this fit separately to obtain the A_{corr} values for the eight bins of $\cos\theta^*$ and use eqs. (1.4) and (1.5) to extract A_{CP} and A_{FB} . The resulting four values of A_{CP} and A_{FB} are plotted in figures 7(c,f). Fitting these A_{CP} values to a constant gives

$$A_{CP}(D^0 \rightarrow \phi\eta) = [-1.9 \pm 4.4(\text{stat}) \pm 0.6(\text{syst})]\%, \quad (5.3)$$

where the second error listed is the systematic uncertainty, evaluated below (section 5.3). This result is consistent with zero, as expected [22].

5.3 Systematic uncertainties

Fortunately, most systematic uncertainties in measuring A_{CP} cancel. The remaining sources of systematic uncertainty are listed in table 3 and are evaluated as follows.

- There is an uncertainty arising from fixed parameters in the fit used to describe signal and background shapes. We evaluate this uncertainty using the sampling method described previously to evaluate uncertainties for the branching fraction measurement. The resulting uncertainties for A_{CP} are small: < 0.001 for both $D^0 \rightarrow \pi^+\pi^-\eta$ and $D^0 \rightarrow K^+K^-\eta$, and 0.002 for $D^0 \rightarrow \phi\eta$. We also consider different possible Q

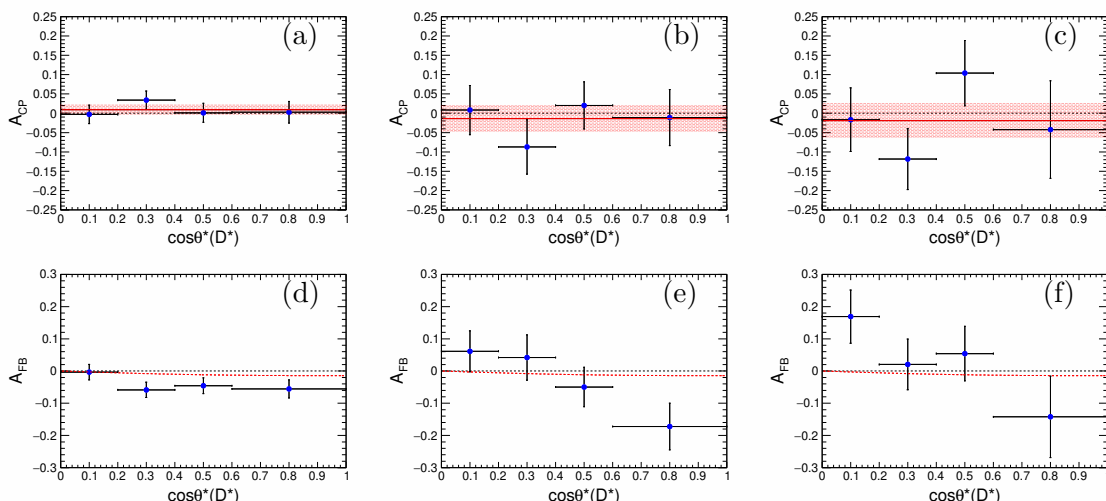


Figure 7. CP -violating asymmetry A_{CP} (top) and forward-backward asymmetry A_{FB} (bottom) values as a function of $\cos\theta^*(D^{*+})$ for (a, d) $D^0 \rightarrow \pi^+\pi^-\eta$, (b, e) $D^0 \rightarrow K^+K^-\eta$, and (c, f) $D^0 \rightarrow \phi\eta$, respectively. The solid red lines with a band region are the averaged values with their uncertainties. The dashed red curves in the A_{FB} plots show the leading-order prediction for $A_{FB}(e^+e^- \rightarrow c\bar{c})$.

Sources	$\sigma_{A_{CP}}(D^0 \rightarrow \pi^+\pi^-\eta)$	$\sigma_{A_{CP}}(D^0 \rightarrow K^+K^-\eta)$	$\sigma_{A_{CP}}(D^0 \rightarrow \phi\eta)$
Signal and bkg	0.004	0.010	0.006
$\cos\theta^*$ binning	0.002	0.004	0.002
$A_\varepsilon(\pi_s)$ map	0.001	0.001	0.001
Total syst. error	0.005	0.011	0.006

Table 3. The absolute systematic uncertainties for A_{CP} measurement in each SCS decay mode.

and M_{KK} resolutions for the D^0 and \bar{D}^0 samples by allowing the σ_0 parameter in eqs. (4.1)–(4.3) and the σ_{m0} and σ_0 parameters in eqs. (4.12), (4.13) to vary between the two samples. The change in A_{CP} is 0.004 for $D^0 \rightarrow \pi^+\pi^-\eta$, 0.010 for $D^0 \rightarrow K^+K^-\eta$, and 0.006 for $D^0 \rightarrow \phi\eta$. Combining these two uncertainties in quadrature gives the values listed in table 3.

- We extract A_{CP} via a binning procedure in $\cos\theta^*$ [see eqs. (1.4)–(1.5)], and there is possible uncertainty arising from the choice of bins used. We thus change the number of bins from eight to six, with bin divisions $(-1.0, -0.55, -0.27, 0.0, 0.27, 0.55, 1.0)$. The resulting change in A_{CP} is taken as the systematic uncertainty due to this source. There is a small uncertainty arising from a difference in the detector acceptance near the boundaries $\cos\theta^* = \pm 1$; we evaluate this by considering only events with $|\cos\theta^*| < 0.90$.
- We correct for a possible asymmetry in π_s^\pm reconstruction efficiencies by weighting events according to a mapping of efficiencies $\varepsilon(\pi_s)$. There are 56 bins in this map, and the efficiencies for each bin has some uncertainty. We thus vary these efficiencies

individually by their uncertainties to create 56 new efficiency maps with $+1\sigma$ shifts and 56 maps with -1σ shifts. We subsequently weight the D^0 and \bar{D}^0 samples by these efficiency maps and repeat the fit for N_{sig} and A_{CP}^{cor} . The resulting deviations from the nominal fit result are summed in quadrature to give the systematic uncertainty arising from this source. We obtain $^{+0.050}_{-0.087}\%$ for $D^0 \rightarrow K^+K^-\eta$, $^{+0.065}_{-0.072}\%$ for $D^0 \rightarrow K^+K^-\eta$, and $^{+0.043}_{-0.100}\%$ for $D^0 \rightarrow \phi\eta$.

The total systematic uncertainty is obtained by adding in quadrature all the above contributions. The results are listed in table 3.

6 Conclusion

In summary, based on a data set corresponding to an integrated luminosity of 980 fb^{-1} recorded by the Belle experiment, we report measurements of the branching fractions of the SCS decays $D^0 \rightarrow \pi^+\pi^-\eta$ and $D^0 \rightarrow K^+K^-\eta$ relative to that for the CF decay $D^0 \rightarrow K^-\pi^+\eta$. We also measure the relative branching fraction for the resonant decay $D^0 \rightarrow \phi\eta$; this measurement uses an order of magnitude more data than used for our previous measurement [8] and supersedes it. Our results are:

$$\frac{\mathcal{B}(D^0 \rightarrow \pi^+\pi^-\eta)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [6.49 \pm 0.09 (\text{stat}) \pm 0.13 (\text{syst})]\%, \quad (6.1)$$

$$\frac{\mathcal{B}(D^0 \rightarrow K^+K^-\eta)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [0.957^{+0.036}_{-0.033} (\text{stat}) \pm 0.021 (\text{syst})]\%, \quad (6.2)$$

$$\frac{\mathcal{B}(D^0 \rightarrow \phi\eta, \phi \rightarrow K^+K^-)}{\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta)} = [4.82 \pm 0.23 (\text{stat}) \pm 0.16 (\text{syst})] \times 10^{-3}. \quad (6.3)$$

The color-suppressed decay $D^0 \rightarrow \phi\eta$ is observed for the first time, with high statistical significance. Multiplying the above results by the world average value $\mathcal{B}(D^0 \rightarrow K^-\pi^+\eta) = (1.88 \pm 0.05)\%$ [5] gives the following absolute branching fractions:

$$\mathcal{B}(D^0 \rightarrow \pi^+\pi^-\eta) = [1.22 \pm 0.02 (\text{stat}) \pm 0.02 (\text{syst}) \pm 0.03 (\mathcal{B}_{\text{ref}})] \times 10^{-3}, \quad (6.4)$$

$$\mathcal{B}(D^0 \rightarrow K^+K^-\eta) = [1.80^{+0.07}_{-0.06} (\text{stat}) \pm 0.04 (\text{syst}) \pm 0.05 (\mathcal{B}_{\text{ref}})] \times 10^{-4}, \quad (6.5)$$

$$\mathcal{B}(D^0 \rightarrow \phi\eta) = [1.84 \pm 0.09 (\text{stat}) \pm 0.06 (\text{syst}) \pm 0.05 (\mathcal{B}_{\text{ref}})] \times 10^{-4}, \quad (6.6)$$

where the third uncertainty is due to the branching fraction for the reference mode $D^0 \rightarrow K^-\pi^+\eta$. These results are the most precise to date.

The time-integrated CP asymmetries are measured to be

$$A_{CP}(D^0 \rightarrow \pi^+\pi^-\eta) = [0.9 \pm 1.2 (\text{stat}) \pm 0.4 (\text{syst})]\%, \quad (6.7)$$

$$A_{CP}(D^0 \rightarrow K^+K^-\eta) = [-1.4 \pm 3.3 (\text{stat}) \pm 1.0 (\text{syst})]\%, \quad (6.8)$$

$$A_{CP}(D^0 \rightarrow \phi\eta) = [-1.9 \pm 4.4 (\text{stat}) \pm 0.6 (\text{syst})]\%. \quad (6.9)$$

The first result represents a significant improvement in precision over the previous result [4]. The latter two are the first such measurements. No evidence for CP violation is found.

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A Bifurcated Student's t-function

The bifurcated Student's t-function is defined as:

$$S_{\text{bif}}(x; \mu, \sigma, \delta, n_l, n_h) = \frac{2P_H P_L}{(P_H + P_L)\sqrt{\pi}} \begin{cases} \left[1 + \frac{1}{n_h} \left(\frac{x-\mu}{\sigma(1+\delta)}\right)^2\right]^{-\frac{n_h+1}{2}}, & \text{for } x \geq \mu; \\ \left[1 + \frac{1}{n_l} \left(\frac{x-\mu}{\sigma(1-\delta)}\right)^2\right]^{-\frac{n_l+1}{2}}, & \text{for others.} \end{cases} \quad (\text{A.1})$$

Here the factors P_H and P_L are calculated as: $P_H = \frac{\Gamma(\frac{n_h+1}{2})}{\sigma \cdot (1+\delta) \Gamma(\frac{n_h}{2})} \frac{1}{\sqrt{n_h}}$ and

$P_L = \frac{\Gamma(\frac{n_l+1}{2})}{\sigma \cdot (1-\delta) \Gamma(\frac{n_l}{2})} \frac{1}{\sqrt{n_l}}$, where Γ is the Gamma function.

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